

IFAE 2008 Bologna, 26-28 marzo 2008 Umberto De Sanctis, Università di Milano & INFN

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#### Outline

- Searches for Supersymmetry at LHC
  - Search strategies for mSUGRA models
  - Commissioning of the detector
  - Measurement and control of backgrounds
- After discovery
  - Measurement of masses and other properties

#### SUPERSYMMETRY REMINDER

#### Adds to each SM fermion (boson) a bosonic (fermionic) partner.

SM Particles	SUSY Particles			
quarks: q	q	squarks: $\tilde{q}$		
leptons: <i>l</i>	l	sleptons: $\tilde{l}$		
gluons: g	g	gluino: $\tilde{g}$		
charged weak boson: $W^\pm$	$W^{\pm}$	Wino: $\widetilde{W}^{\pm}$	~±	
	$H^{\pm}$	charged higgsino: ${\widetilde{H}}^{\pm}$	$\int \chi_{1,2}$ chargino	
niggs. 11	$h^{\circ}, A^{\circ}, H^{\circ}$	neutral higgsino: ${ ilde h}^{\circ}, { ilde A}^{\circ},$	Ĵ	
neutral weak boson: $Z^{0}$	$Z^{o}$	Zino: $\widetilde{Z}^{\circ}$	$\widetilde{\chi}^{0}_{1,2,3,4}$ neutralino	
photon: $\gamma$	γ	photino: $\tilde{\gamma}$		

• **R-parity**  $R = (-1)^{3(B-L)+2S}$  can be conserved (**RPC**) or violated (**RPV**)

- RPC implies:
  - SUSY particles produced in pairs
  - stable and neutral lightest SUSY particle (LSP)
  - no proton decay

LSP is a good candidate for cold Dark Matter

**MSSM** Lagrangian depends on 105 parameters **mSUGRA** requires only 5 parameters

- Also other SUSY models exist: GMSB, AMSB, ...

Par.	Description			
m <sub>0</sub>	Common scalar mass			
m <sub>1/2</sub>	Common gaugino mass			
A <sub>0</sub>	Common trilinear term			
tanβ	Ratio of Higgs vev			
sign(μ)	$\mu$ from Higgs sector			

#### A needle in an hay stack



Only one event (i.e. pp collision) in **one bilion** may contain an Higgs boson or a squark....

#### Need high luminosity

Need an **efficient online selection** (trigger) to select interesting events: cannot register everything electronically for further processing

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#### What do we do when we get the data?

Before we can claim discovery of "New Physics" we have to do some homework...

- Understand and calibrate detector and trigger in situ using well-known physics samples:  $Z/W \rightarrow$  leptons, semileptionic tt
- Understand basic SM physics at 14 TeV: first measurements and publications
  - $\bullet$  jets and  $W\!,Z$  cross-section top mass and cross-section
  - Event features: Min. bias, jet distributions, PDF constraints
- Understand tails of SM processes as backgrounds (tt, W/Z + jets), go for discovery: Z', SUSY, Higgs

But let's have a look at our main SUSY discovery strategy, to understand what we need to understand to get there...

mSUGRA benchmark points

SUSY benchmark points chosen in the  $(m_0, m_{1/2})$  plane for different  $tan\beta$  values:

- ✓ Systematically exploring phenomenological signatures
- Scanning the parameter phase space constrained by latest experimental data

Coannihilation: Light  $\tilde{\tau}_1$  in equilibrium with  $\tilde{\chi}_1^0$ , so annihilate via  $\tilde{\chi}_1^0 \tilde{\tau}_1 \rightarrow \gamma \tau$ .

*Bulk:* bino  $\tilde{\chi}_1^0$ ; light  $\tilde{\ell}_R$  enhances annihilation.

*Funnel:* H,A poles enhance  $\overleftarrow{\boldsymbol{\xi}}^{T}$  annihilation for tan  $\beta \gg 1$ .

Focus point: Small  $\mu^2$ , so Higgsino  $\tilde{\chi}_1^0$  annihilate. Heavy s-fermions, so small FCNC.



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#### SUSY signatures at an hadronic collider

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- Assuming R-parity conservation
- Strongly interacting sparticles (squarks, gluinos) should dominate production unless very heavy.
- Cascade decays to the stable, weakly interacting lightest neutralino follows.
- Event topology:
  - high p<sub>T</sub> jets (from squark/gluino decay)
  - Large E<sub>T</sub><sup>miss</sup> signature (from LSP)
  - High p<sub>T</sub> leptons, b-jets, τ-jets (depending on model parameters).

Several other possibilities exist (even if not mentioned in this talk), but our effort has to be as more "model independent" as possible.

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## Event topologies and baseline selection

Early searches try to cover a broad range of experimental signatures, but they are classified based on the event topology:

Large E <sub>T</sub> miss +	Jet multiplicity	Additional SUSY scenario		Backgrounds
	<i>≥</i> 4	No lepton	mSUGRA, AMSB, split SUSY, heavy squark	QCD, ttbar, W/Z
		One lepton (e, $\mu$ )	One lepton (e, $\mu$ ) mSUGRA, AMSB, split SUSY, heavy squark	
		di-lepton	mSUGRA, AMSB, GMSB	ttbar
		di-tau	GMSB, large tan <i>β</i>	ttbar, W
		γγ	GMSB	free
	~2		light squark	Z

Baseline selection (to be optimized)

- Jet multiplicity  $\ge 4$ ,  $p_T^{1st} > 100 \text{GeV}$ ,  $p_T^{others} > 50 \text{GeV}$
- $E_T^{miss} > max(100 GeV, 0.2 x M_{eff})$
- Transverse sphericity > 0.2
- Additional cuts depending on signature: Transverse mass > 100GeV,  $p_T^{lepton} > 20GeV$  (for one-lepton mode), harder cuts on  $M_{eff}$ ... Bologna, 27/03/2008 IFAE 2008 U. De Sanctis

# ET MISS Commissioning

- Biggest challenge fake E<sup>T</sup><sub>MISS</sub>: Tevatron experience shows can be hard to control
  - Machine backgrounds, calorimeter problems, jet mis-measurement ...
  - Devote much effort to event-cleaning
  - Focus on channels where minimised until understood (e.g. leptons)



- min bias events, dijet  $P_T$  balance,  $m_T(W \rightarrow l\nu)$ ,  $Z \rightarrow ll$ ,  $Z \rightarrow \tau \tau$  etc.
- Beyond this: non-Gaussian tails absolutely crucial
  - Event cleaning/rejection
  - Estimation

After this:

- Validation of MC with control sample  $(Z \rightarrow \ell \ell, tt \rightarrow bqqb\ell v)$ ;
- Estimation of W/Z+jets and tt backgrounds from  $m_T(\ell, E^T_{MISS})$  or Z- $\rightarrow \ell\ell$  control sample;
- Optimization of inclusive search event selection with major backgrounds estimates (MC or data driven) normalized to data;

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Missing ET in MHT30 skim

#### SUSY search strategies: No lepton mode

Most promising search strategy: jets + E<sub>T</sub><sup>miss</sup> + n-leptons

- Real missing energy from SM
  processes with hard neutrino (tt,
  W+jets, Z+jets, bb\*, cc\*)
- \*  $\nu$  from semileptonic B/D decay
- Fake missing energy from detector
- Jet energy resolution (expecially non-gaussian tails) critical A good understanding of both SM physics and detector (missing energy expecially) critical to claim excess over SM predictions







#### Z+ jets background

• Data-driven estimation of  $Z \rightarrow vv$  bkg  $E_T^{miss}$  shape and normalization using  $Z \rightarrow \mu\mu$  sample and replacing  $\mu$  with v



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#### ttbar background

- 1. Top mass is largely uncorrelated with  $E_T^{miss}$ 
  - used as a calibration variable
- 2. Select semi-leptonic top candidates
  - mass window: 140-200 GeV
- 3. Contributions of combinatorial BG to top mass are estimated from the side-band events (200GeV<m<sub>top</sub><260GeV)
- 4. Normalize the  $E_T^{miss}$  distribution in low  $E_T^{miss}$  region where SUSY signal contamination is small.
- 5. Extrapolate it to high  $E_T^{MISS}$  region and estimate the background with SUSY signal selection.

Several other techniques also under investigation



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#### Other strategy: 1-lepton channel

- Removing the lepton-veto: 1 lepton + Jets + E<sup>T</sup><sub>MISS</sub> channel
  - Lepton can usually come from chargino/neutralino decays into LSP
  - Heavily suppression of the QCD background (difficult to estimate from data and also with MC) requiring 1 isolated lepton with  $P_T > 20$  GeV/c.
  - Dominant background are the same as the 0-lepton channel (except QCD) : top seems to be the dominant one, but W + jets is not negligible.



#### Back to SUSY discovery



 $5\sigma$  discovery potential on m1/2-m0 (mgluino-msquark) space for 1 fb-1

1 fb<sup>-1</sup> of ATLAS/CMS data should be enough to discover SUSY if squark/gluino mass lower than 1.5+2 TeV.

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# Di-Lepton Edge mass measurement (1)

- In case of a discovery of SUSY, particle properties can be measured to verify that they are indeed SUSY partners
- Edge(s) of di-lepton invariant mass correlated with slepton and neutralino masses



 $\widetilde{\chi}^{0}_{2} \rightarrow \widetilde{l} \ l \rightarrow \widetilde{\chi}^{0}_{1} \ l^{+} \ l^{-}$ 

- Impossible to reconstruct peaks because  $\chi_1^0$  (LSP) escapes detection, more complicated relations between masses of particles  $M_{II}^{\max} = M(\tilde{\chi}_2^0) \sqrt{1 - \frac{M^2(\tilde{l}_R)}{M^2(\tilde{\chi}_2^0)}} \sqrt{1 - \frac{M^2(\tilde{\chi}_1^0)}{M^2(\tilde{l}_R)}}$ involved.
  - Uncorrelated (SUSY+SM) background (two leptons from independent chains) removed by flavour subtraction:
     e<sup>+</sup>e<sup>-</sup> + β<sup>2</sup> μ<sup>+</sup>μ<sup>-</sup> β (e<sup>+</sup>μ<sup>-</sup>-e<sup>-</sup>μ<sup>+</sup>), β=ε<sub>e</sub>/ε<sub>μ</sub>
  - Leptons can also be combined with jets of the full decay chain to look for other kinematical edges ( $M_{III}$  or  $M_{III}$ )

#### Di-Lepton Edge mass measurement (2)

#### Flavour subtraction at work....





## **Conclusions**

- A brief review of the search strategies for SUSY in ATLAS & CMS has been presented;
  - New discoveries possible with <u>early LHC data (O(100)pb<sup>-1</sup></u>)
- Accurate knowledge of SM physics and of detector performance needed for any new discovery
  - First data taking period devoted to understanding of detector
- Any claim of new physics requires check of trigger refinements and data-driven estimates of syst./background
  - First, focus on less systematic-affected analyses (e.g. striking signatures and resonances)
- Larger statistics needed for full scan over SUSY parameters space and **discrimination** between different models.

# **BACKUP SLIDES**

#### **Detector calibration and alignment**

The jet energy scale affects directly SUSY discovery plots trough the cut on the presence of hard jets.

Also,  $E_T^{miss}$  depends on the correct reconstruction of the energies of jets, photons, electrons, and muons!

- We will start from the knowledge obtained from test-beam data, electronics calibrations, survey measurements during installation of the tracking detectors, and cosmics data.
- We will then use well-known SM processes (standard candles) to improve **Examples: leptonic decays of Z, W mass in semileptonic top events**

E	Expected performance day-1		Physics samples to improve (examples)		
ECAL uniformity e/γ1-2% (~0.5% locally) ~ 2 % ~ 3 %      Y E-scale Fracking alignmentHCAL uniformity Jet E-scale Tracking alignment1-2% (~0.5% locally) ~ 3 %       10-200 μm in R\phi Pixels/SCT ?				Isolated electrons, $Z \rightarrow ee$ Single pions, QCD j $\gamma/Z + 1j$ , $W \rightarrow jj$ Generic tracks, isola	Z→ee ets in tt events ated μ, Ζ → μμ
Process	$\sigma \times BR$	Eff.	Events	selected for 100 $\rm pb^{-1}$	Available statistics
$W \to \ell \nu$	20 nb	$\sim 20\%$		$\sim 400000$	estimates of
$Z \to \mu \mu$	2 nb	$\sim 20\%$		$\sim 40000$	reconstruction
$\bar{t}t$ (semileptonic)	370 pb	$\sim 1.5\%$		< 1000	efficiencies
Pologna 27/02/2009				סר	22

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#### ETMISS Commissioning (2)

Just two examples, several other physics process can be used: minimum bias, Z(II), ttbar, ...

W(lv) sample: Shape of transverse mass distribution depends on  $E_T^{miss}$  scale and resolution.

 $Z(\tau\tau)$  sample: Z mass can be reconstructed with collinear approximation (since the  $\tau$  are boosted, v are along visible  $\tau$  energy). Can be used to calibrate  $E_T^{miss}$  scale.

 $\sigma$  = 14.2 GeV

80

60

80

70

50

30

20

10

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Z and W background

## Estimate $Z \rightarrow vv$ from $Z \rightarrow \ell^+\ell^-$



Similarly we can use the Z distribution to estimate the  $W \rightarrow \ell v$  background

#### Z and W background (O-lepton mode)

 $Z{\rightarrow}\nu\nu$  and  $W{\rightarrow}l\nu$  can be estimated from  $Z{\rightarrow}\ell^+\ell^-$ 

Either **replace** the two leptons with neutrinos correcting for acceptance and efficiency

Or determine the **MC normalization** from Z(ll) and apply it to normalize the MC distribution of Z(vv) and W(lv) (almost same production mechanism)



#### Transverse mass method



- 1. Define control sample with transverse mass <100GeV
- 2. Estimate the  $E_T^{miss}/M_{eff}$  shapes of background processes using control sample
- 3. Determine the normalization of backgrounds with low  $E_T^{miss}$  regions of control and signal samples.

# Can be used for both W and top backgrounds in 0-lepton, 1-lepton and 2 lepton channels (results shown here for 1-lepton)

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#### Transverse mass method

Including SUSY signal (SU3)



• Satisfying performances with the  $M_{\tau}$  discrimination technique.

• However, taking account of SUSY signal contamination in the control sample, this estimate appears to be over the mark (by a factor of 2.5 for SU3). It would not prevent discovery.

#### Other strategy: 2-lepton channel

- Increasing the number of leptons
  - Reduces the signal because of (model dependent) leptonic BRs
  - Heavily suppresses the background
  - Statistical significance is smaller but S/B ratio larger. Top is dominant background
  - The Same Sign channel has the best S/B ratio but limited by signal rate



#### What next ?

"Observation of an excess of events in multijet+MET events in pp collisions at 14 TeV with the ATLAS detector"



Large (>100GeV) Missing ET events: Smoking gun of Supersymmetry



Measurement of the "effective mass" peak correlates with the SUSY mass scale (average squark, gluino mass) Meff = MET+PT,1+PT,2+PT,3+PT,4 15% (40%) precision on M(SUSY) with 10fb<sup>-1</sup> for mSUGRA (MSSM)

#### Measurement of neutralino spin (1)

Important to measure the spin of new particles: it's the fundamental check to ensure that what we have discovered is SUSY!!



The charge asymmetry is **diluted** because:

- Usually it is not possible to discriminate the *near* and *far* leptons: we sum m(ql<sup>far</sup>) and m(ql<sup>near</sup>) invariant masses
- 2. The charge conjugated cascade decay (from the anti-squark) gives the opposite asymmetry. However, cancelation is not exact because at LHC a larger number of squarks than anti-squarks is produced (pp collider)

#### Measurement of neutralino spin (2)



#### **Other SUSY scenarios**

Across the MSSM, there is a rich variation in the SUSY phenomenology. The signatures expected at the LHC can be very different from the "mainstream" scenario discussed so far.

- <u>GMSB</u>: the lightest SUSY particle is the gravitino. The next-to-lightest particle (NLSP) decay only trough gravitational interactions and may live longer than the time-of-flight across the detector.
- <u>Split SUSY:</u> scalars are much heavier than the electroweak scale. The gluino decays trough virtual squarks, and may live longer than the time of flight across the detector.
- <u>R-parity violation</u>: the neutralino decays. Less missing energy and more jets or other particles.
- <u>Light stop models</u>: a scalar top with 120-150 GeV mass is still allowed.

Each scenario is covered by dedicated search strategies. I will discuss the GMSB scenario here...

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#### Gauge Mediated Supersymmetry Breaking. Models for SUSY breaking, alternative to mSUGRA

- SUSY breaking transmitted from Hidden sector to visible sector via gauge interactions ("messengers")
- Why interesting?
  - more natural suppression of FCNC
  - not huge  $\sigma$  but clear signature to claim early discovery or exclusion
    - $\sigma \sim 0.1 \div 1 \text{ pb}$  (model dependent)

LSP is the **Gravitino** (m≤keV)

- light, stable and weakly interacting
- possible candidate for Dark Matter

Present limits: Tevatron,  $\Lambda > 80$  TeV, m(neutralino,chargino) > 108, 195 GeV

Par.	Description		
Λ	SUSY breaking scale		
M <sub>m</sub>	Messenger mass scale		
tanβ	Ratio of Higgs vev		
N <sub>m</sub>	Number of SU(5) messenger multiplets		
sign(μ)	$\mu$ from Higgs sector		
C <sub>grav</sub>	Sets NLSP lifetime		

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#### GMSB Model (2)

Phenomenology depends on nature and lifetime of the second lightest state (NLSP):

	${\widetilde {\cal T}}$ or ${\widetilde l}$ is NLSP	$\widetilde{\chi}_1$ is NLSP		
cτ >> L	Like an heavy μ	Like mSUGRA		
cτ ≈ L	NLSP decays in the detector, life	etimes measurements.		
cτ << L	Decays into 2 τ	Decays into 2 $\gamma$		

- **τ trigger and reconstruction** in early data not trivial
- **Decay into 2**γ **promising** (good ECAL performance early enough?)
- Lifetime measurements: need to understand vertexing in early data
  - For longer lifetimes, need to understand **background**:
    - Hard radiation from high-p<sub>T</sub> cosmic muons
    - Delayed hadronic showers  $(K_L^0 and neutrons)$

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## GMSB Model: Performances (2)

- Heavy slow "stable" leptons can be tagged with Time-Of-Flight measurements in muon drift tubes.
- Large calorimetric  $E_{MISS}^{T}$  due to quasi-stable leptons, like in mSUGRA.
- Timing/triggering issues most critical (association to the correct BCID problematic if  $\beta < 0.7$ , recoverable with MDT but specific algorithm for long-lived heavy particles will be useful).



#### Light stop scenario

- Direct limits allows the scalar top to be lighter than top
- There are models which explain baryogenesis (the generation of matterantimatter asymmetry of the Universe) and Dark Matter at once using SUSY
  - Give up SUGRA-like unification of SUSY masses
  - Require a very light stop
  - ... and of course CP violation
- Consider direct production stop pairs:



- Looks a lot like top pair production
  - Cross section is comparable (400 pb for a 140 GeV stop)
  - Same final state, but "wrong" invariant mass combination (no W, top peak)
  - Still two unobserved neutralinos: no mass peak!
  - Softer leptons, jets and missing energy than in ttbar
  - Biggest problem is ttbar background Bologna, 27/03/2008 IFAE 2008 U. De Sanctis



The signal is "visible" on top of the SM background <u>if we assume</u> we know (from Montecarlo predictions) how many SM events (and the shape of distribution) pass event selection on average.

Since we may not trust the MC prediction to this level of accuracy, we developed a technique to estimate the shape of the SM contribution to the distribution.

Once we know the shape, we can fix the normalization of the background using the events at large invariant mass, where no SUSY contribution is expected.

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Search for light stop (II)

#### We developed a technique to estimate the ttbar bckg from data:

**Control sample 1:** tight extra cuts on hadronic side [M(jjb) = m(top), M(jj)=M(W)] select ttbar - Used to measure shape of M(bl) in ttbar events

**Control sample 2:** tight cuts on leptonic side [M(lb,xEt) = M(top)] to select ttbar -Used to measure shape of M(bjj) in ttbar events

# Signal visible after background subtraction with $\sim 1 \text{ fb}^{-1}$





**Solid line:** SUSY events among those passing event selection

**Points:** Measured distribution, after subtracting the SM contribution estimated with control

#### mSUGRA models

- A random choice of the 105 MSSM parameters violates limits from B/D/K physics, electric dipole moments, FCNC, ...
- Need some assumption on the structure of SUSY breaking lagrangian. In mSUGRA (5 free parameters, most studied by ATLAS and CMS):
  - Conserved R-parity
  - **Common mass**  $m_0$  for susy scalars,  $m_{1/2}$  for fermions (at GUT scale).
  - Common value A<sub>0</sub> for the trilinear coupling of the s-fermions with the 2 Higgs doublets.

Then 5 free parameters:  $m_0$ ,  $m_{1/2}$ ,  $A_0$ ,  $\tan \beta$ , sgn  $\mu$ 

Further constraints if it is required that the Big Bang has produced the right amount of stable neutralinos to explain observed Dark Matter density May be too constrained. Experiments at colliders are interested mostly in ider Mingreight ANAPS to develop and stExEy29@arch strategies 40 U. De Sanctis

#### Supersymmetry: what is?

# Supersymmetry (SUSY) in a nutshellStandard particlesSuperpartnersQuarks, leptons, neutrinos (spin 1/2)Squarks, sleptons, sneutrinos (spin-0)W, Z, gluino (spin-1)Wino, zino, gluino (spin 1/2)Higgs (spin-0)Higgsino (spin 1/2)At least two Higgs doublets are needed → five Higgs bosons

Wino, Zino, Higgsino mix  $\rightarrow$  4 charged (chargino) and 4 neutral (neutralino) states

SUSY particles not observed yet  $\rightarrow$  must be heavy  $\rightarrow$  symmetry is broken

It is possible to put directly SUSY mass terms in the lagrangian. This gives about **100 free parameters** with the minimal field content above (MSSM model)

**Constrained models** (with assumptions on the structure of SUSY breaking) have only a few parameters – but assumptions may be wrong.

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#### Kinematical structures



Same-Flavour Other-sign Lepton combinations.

ee + μμ - eμ Entries / 2 GeV



The invariant mass of the two leptons has a kinematical endpoint which measures:

$$M_{ll}^{\max} = M(\tilde{\chi}_{2}^{0}) \sqrt{1 - \frac{M^{2}(\tilde{l}_{R})}{M^{2}(\tilde{\chi}_{2}^{0})}} \sqrt{1 - \frac{M^{2}(\tilde{\chi}_{1}^{0})}{M^{2}(\tilde{l}_{R})}}$$

It may be possible to observe two edges, if both decays are open:

$$\begin{array}{ccc} \chi_{2}^{0} \rightarrow l \widetilde{l}_{L} \rightarrow l l \chi_{1}^{0} \\ \chi_{2}^{0} \rightarrow l \widetilde{l}_{R} \rightarrow l l \chi_{1}^{0} \end{array}$$

The SM and SUSY combinatorial backgrounds have two leptons from independent decay chains. The background cancel in the flavour subtraction FAE 2008 42 U. De Sanctis

#### Getting mass peaks

• The 4-momentum of the  $\chi^0_2$  can be reconstructed from the approximate relation

 $p(\chi_2^0) = (1-m(\chi_1^0)/m(II)) p_{II}$ valid when m(II) near the edge.

• The  $\chi_2^0$  can be combined with b-jets to reconstruct the gluino and sbottom mass peaks from  $\tilde{g} \rightarrow b \tilde{b} \rightarrow b b \chi_2^0$ 

SPS1a, 300 fb<sup>-1</sup>, stat. errors only:  $m(\tilde{g})-0.99m(\chi^0_1) = (500.0 \pm 6.4) \text{ GeV}$   $m(\tilde{g})-m(\tilde{b}_1) = (103.3 \pm 1.8) \text{ GeV}$  $m(\tilde{g})-m(\tilde{b}_2) = (70.6 \pm 2.6) \text{ GeV}$ 





 $m(\chi bb)$ - $m(\chi b) (GeV)$ 

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#### Other mass measurements



#### Supersymmetry: why?

Supersymmetry can solve several problems of the Standard Model at once

#### Hierarchy problem:

- Fermions and bosons contribute with opposite sign to the Higgs mass
- $\delta m_{\rm H} \sim m_{\rm SUSY}$  [SUSY mass scale]
- Hierarchy ok if SUSY masses near the Higgs scale (accessible to a TeV-scale collider) True also for other SM extensions addressing hierarchy. The TeV-scale new physics and the Higgs are the main motivations for the Large Hadron Collider

#### **Dark Matter**

- Need a conserved quantum number to avoid
- proton decay: R = +1 for SM particles, R = -1
- for SUSY particles. Consequences:
- SUSY particles are produced in pairs
- The **lightest SUSY particle is stable**. If weakly interacting, it's a good candidate for Dark Matter

#### **Unification of forces:**

Better convergence of interaction strength as a function of energy





An other variable which has small correlation with MET is

#### HT2 = $\sum$ (pt jets 2,3,4) + $\sum$ (pt e, $\mu$ )

- leading jet is not included in order to avoid correlation with MET
- use MET significance rather than MET to reduce correlation

#### one lepton mode



#### also works for OS di-lepton mode

#### **Getting SUSY particle masses**

- Combine measurements from edges of different jet/lepton combinations to obtain 'model-independent' mass measurements.
- LSP mass uncertainty large, all other masses strongly correlated with it. A future Linear Collider measurement of  $\chi^0_1$  mass would improve the precision on all masses.



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masses (GeV)		LHCC5	SPS1a			
$m(\tilde{\chi_1^0})$		122	96			
$m(\tilde{l}_R)$		157	143			
$m(\tilde{\chi_2^0})$		233	177			
$m(\widetilde{q_L})$		687-690	537-543			
Spartic	le Expe	Expected precision (100 fb <sup>-1</sup> )				
q̃∟		± 3%				
$\widetilde{\gamma}^0$		± 6%				

+ 9%

± 12%

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 $\widetilde{\chi}^{0}_{1}$ 

#### From masses to model parameters

From a given set of measurements one scans the **parameter space** and finds the points compatible with data. These points are fed to relic density calculators to get constraints on **neutralino dark matter abundance** 

			$\Omega_{\gamma}h$		
Variable	Value (GeV)	Stat. (GeV)	Scale (GeV)	Total	~
$m_{\ell\ell}^{max}$	77.07	0.03	0.08	0.08	$\log_{10}($
$m_{\ell\ell\sigma}^{max}$	428.5	1.4	4.3	4.5	n 1200
$m_{lg}^{low}$	309.3	0.9	3.0	3.1	
$m_{\ell\sigma}^{high}$	378.0	1.0	3.8	3.9	ti da la
min	201.9	1.6	2.0	2.6	පි 1000 -
$m_{\ell\ell b}^{min}$	183.1	3.6	1.8	4.1	Š.
$m(\ell_L)-m( ilde{\chi}^0_1)$	106.1	1.6	0.1	1.6	
$m_{\ell\ell}^{max}( ilde{\chi}_4^0)$	280.9	2.3	0.3	2.3	800 -
$m_{\tau\tau}^{max}$	80.6	5.0	0.8	5.1	-
$m(\tilde{g}) = 0.99 \times m(\tilde{\chi}_1^0)$	500.0	2.3	6.0	6.4	-
$m( ilde{q}_{R})-m( ilde{\chi}_{1}^{0})$	424.2	10.0	4.2	10.9	600 -
$m( ilde{g}) - m( ilde{b}_1)$	103.3	1.5	1.0	1.8	-
$m( ilde{g}) - m(b_2)$	70.6	2.5	0.7	2.6	-
			<b>400</b> –		
Parameter	Expecte	d precisio	on (300 fb <sup>-</sup>	<sup>1</sup> )	-
mo		± 2%			200
m <sub>4/2</sub>		± 0.6			
tan(B)		+ 9%			
Λ		± 160/	0.170		
<b>A</b> 0		± 10%			•
Bologna, 27/03	/2008		IFAE	2008	



 $\Omega_{\chi}h^2 = 0.1921 \pm 0.0053$ 



#### How much data will we need?

Statistical reach with 100 pb<sup>-1</sup> is in the TeV region, well beyond Tevatron limits (~400 GeV) BUT

- only in a few cases SUSY has distinctive kinematical features

- main selection tool at both trigger and analysis level is to select event with large missing Et, difficult to muster experimentally

More luminosity (for control samples) and/or time may be needed to understand backgrounds

Let's go back to detector commissioning and SM background studies...