

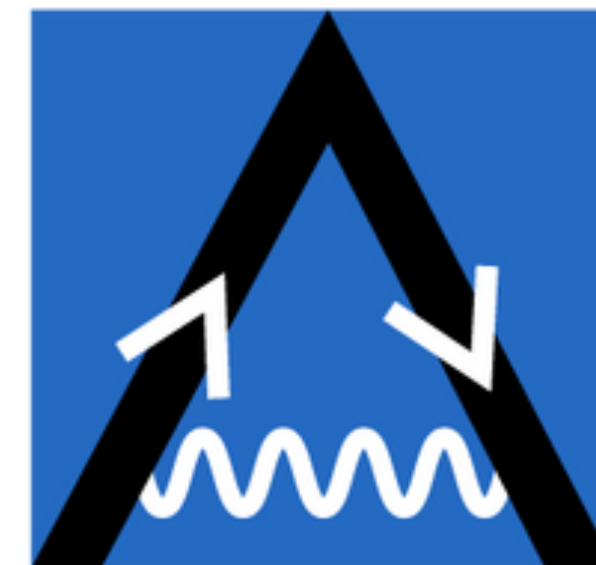
Study of the simulations of the readout for the future LHCb Electromagnetic Calorimeter and characterization of photomultipliers

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2nd examiner: Dr. Stefano Perazzini



International Master
Advanced Methods
in Particle Physics

Outline

- Introduction: the LHCb ECAL Upgrade II
- Simulation studies of a Pb-Polystyrene module
- Analysis of testbeam data from the CERN SPS
- Ageing campaign of the PMT Hamamatsu R14755U-100
- Conclusions

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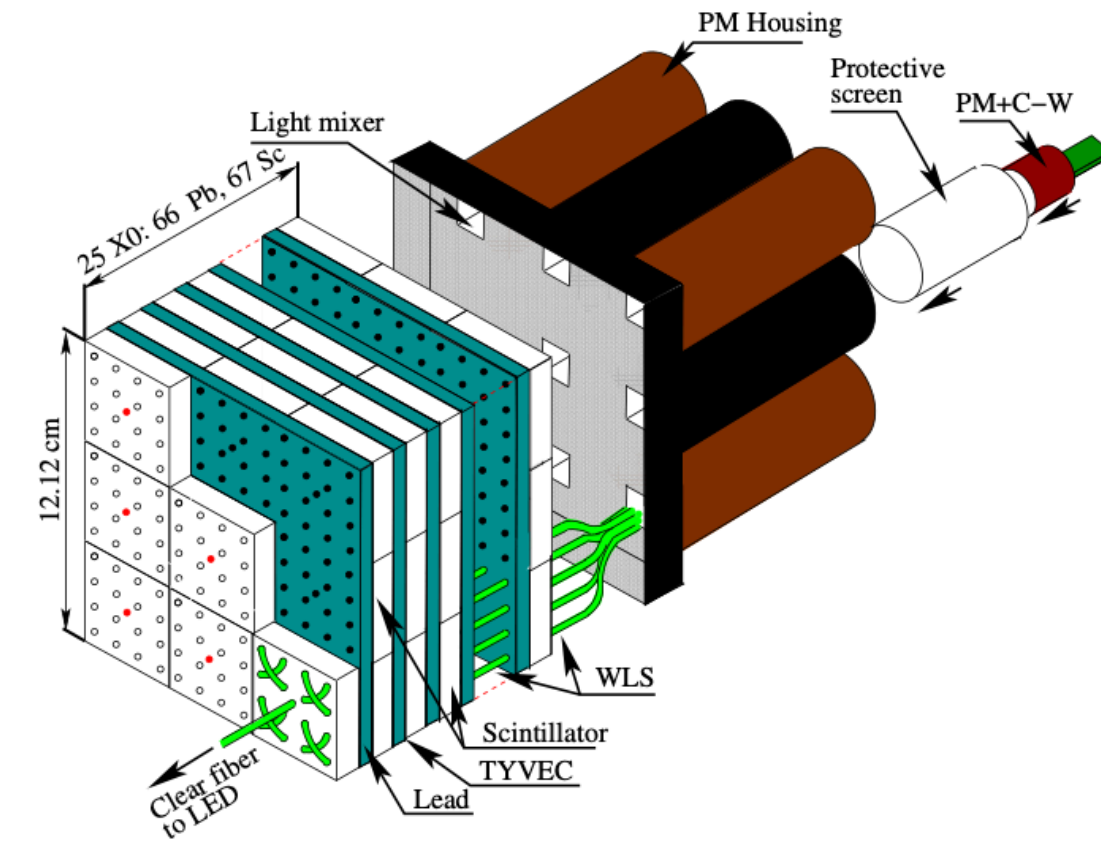
Working principle of a sampling ECAL

(Scintillator-based)

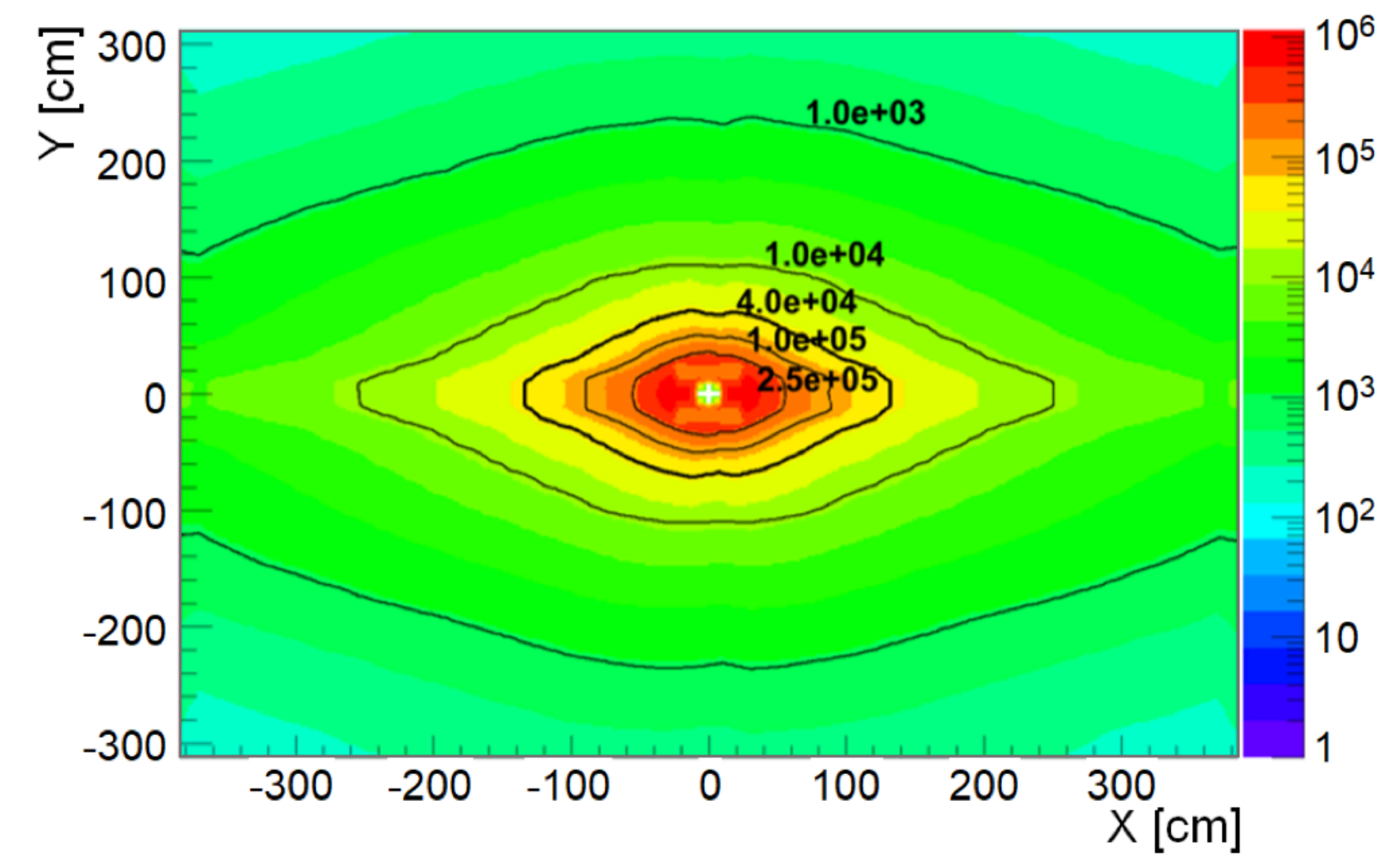
- Electromagnetic Calorimeters (**ECALs**) measure the incident particles' energy (e^\pm, π^0) providing position information, too
- Two main materials involved:
 - **Scintillator** (active): converts the incident particles' energy into **light**
 - **Absorber** (passive): generates the **electromagnetic shower** and absorbs the energy
- The produced light has to be propagated to the photodetectors
- Finally, the photodetectors (e.g. **PMTs**) create the signals

LHCb ECAL Upgrade II

- Currently: sampling ECAL composed of Shashlik modules
- Radiation doses ~ 1 MGy foreseen for Run 5 and Run 6
- Luminosity up to $1.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- The high luminosity environment will require:
 - ▶ **Time resolution \sim few tens of picoseconds**
 - ▶ **Radiation hardness**
 - ▶ Energy resolution at the level of the current one (10% sampling term, 1% constant term)



Scheme of a currently-used Shashlik module (Irina Machikhiliyan and LHCb calorimeter group. <https://iopscience.iop.org/article/10.1088/1742-6596/160/1/012047>)



Expected radiation dose for the High Luminosity phase, in Gy (“Framework TDR for the LHCb Upgrade II: Opportunities in flavour physics, and beyond, in the HL-LHC era.” <https://inspirehep.net/literature/2707810>)

LHCb ECAL Upgrade II

- Future: **Spaghetti Calorimeter (SpaCal)**
- Scintillating fibres inserted into a dense passive absorber
 - Fibres: **polystyrene** / **garnet crystal**
 - Absorber: **lead (Pb)** / **tungsten (W)**

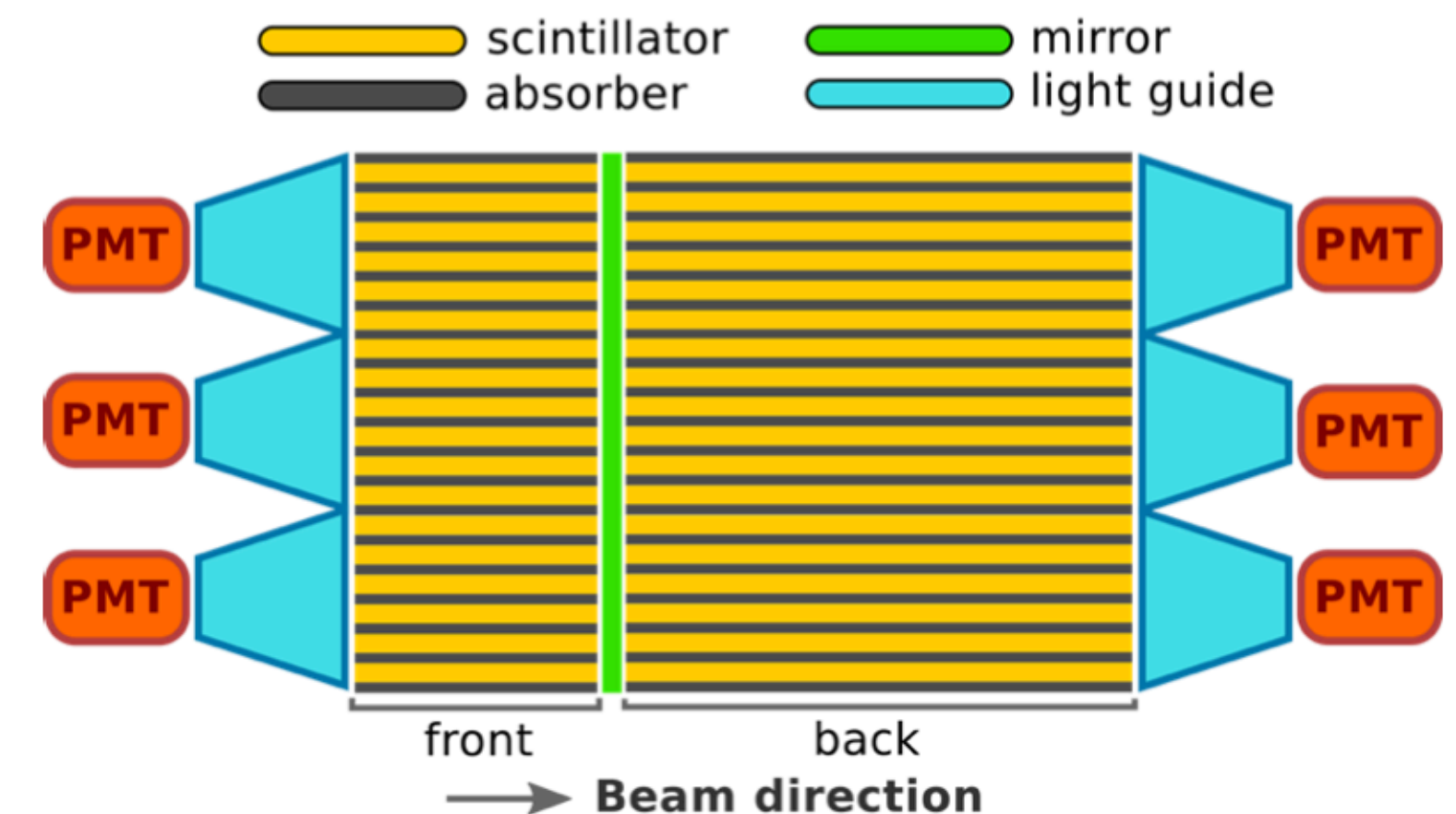
➔ Run 4:

- **W-Poly**
- Single-side readout

➔ Run 5 & 6:

- **W-Crystal** and **Pb-Poly**
- Double-side readout

Picture of a Pb-polystyrene prototype in a test-beam setup

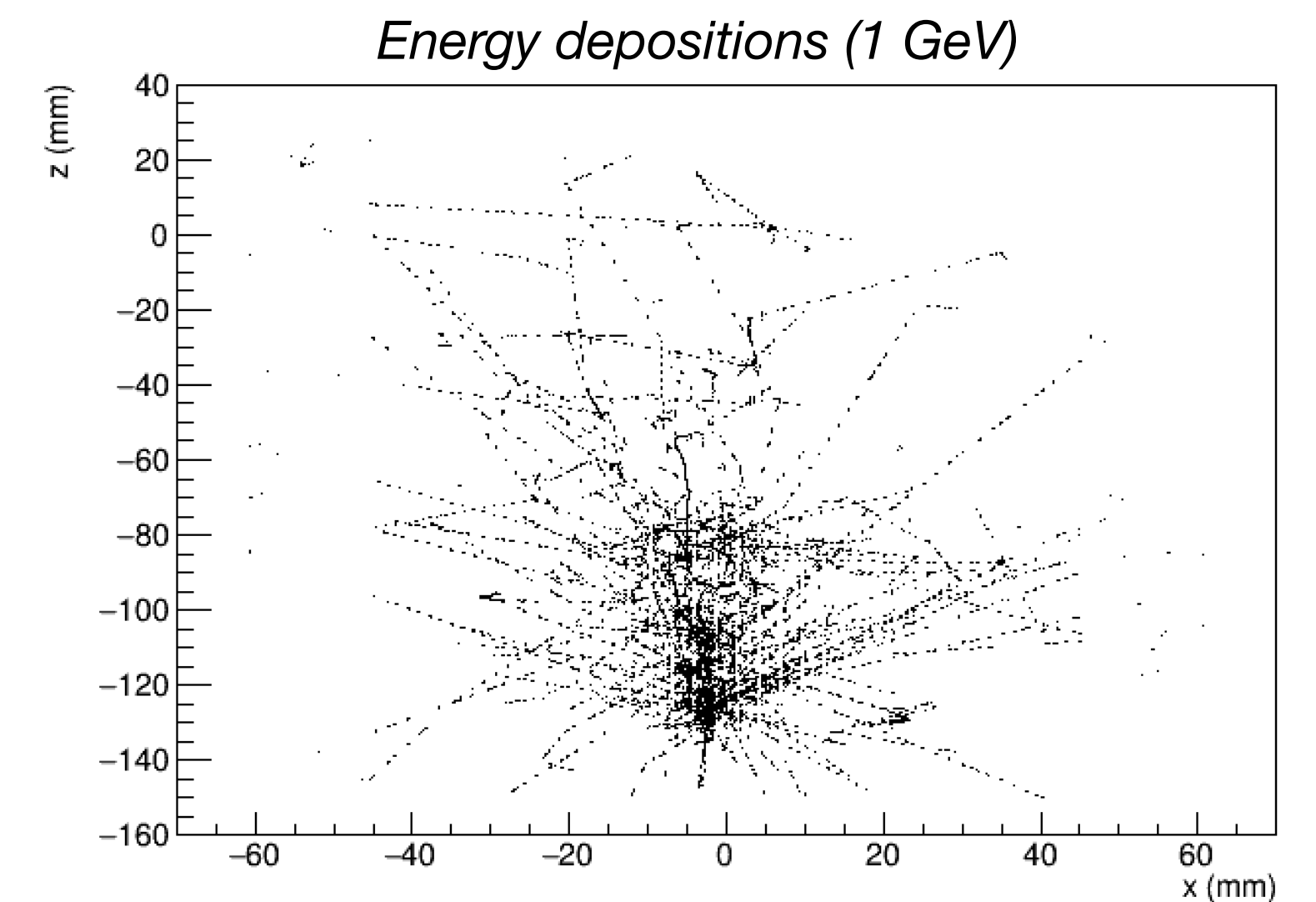
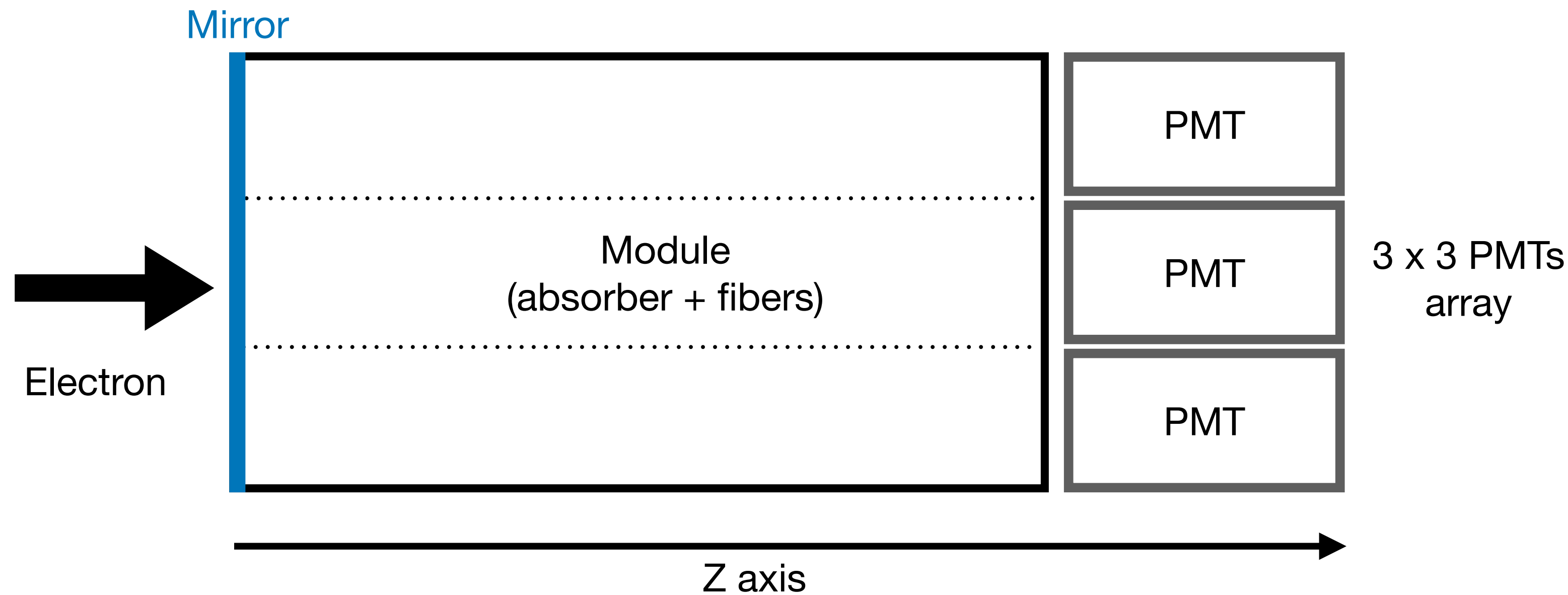
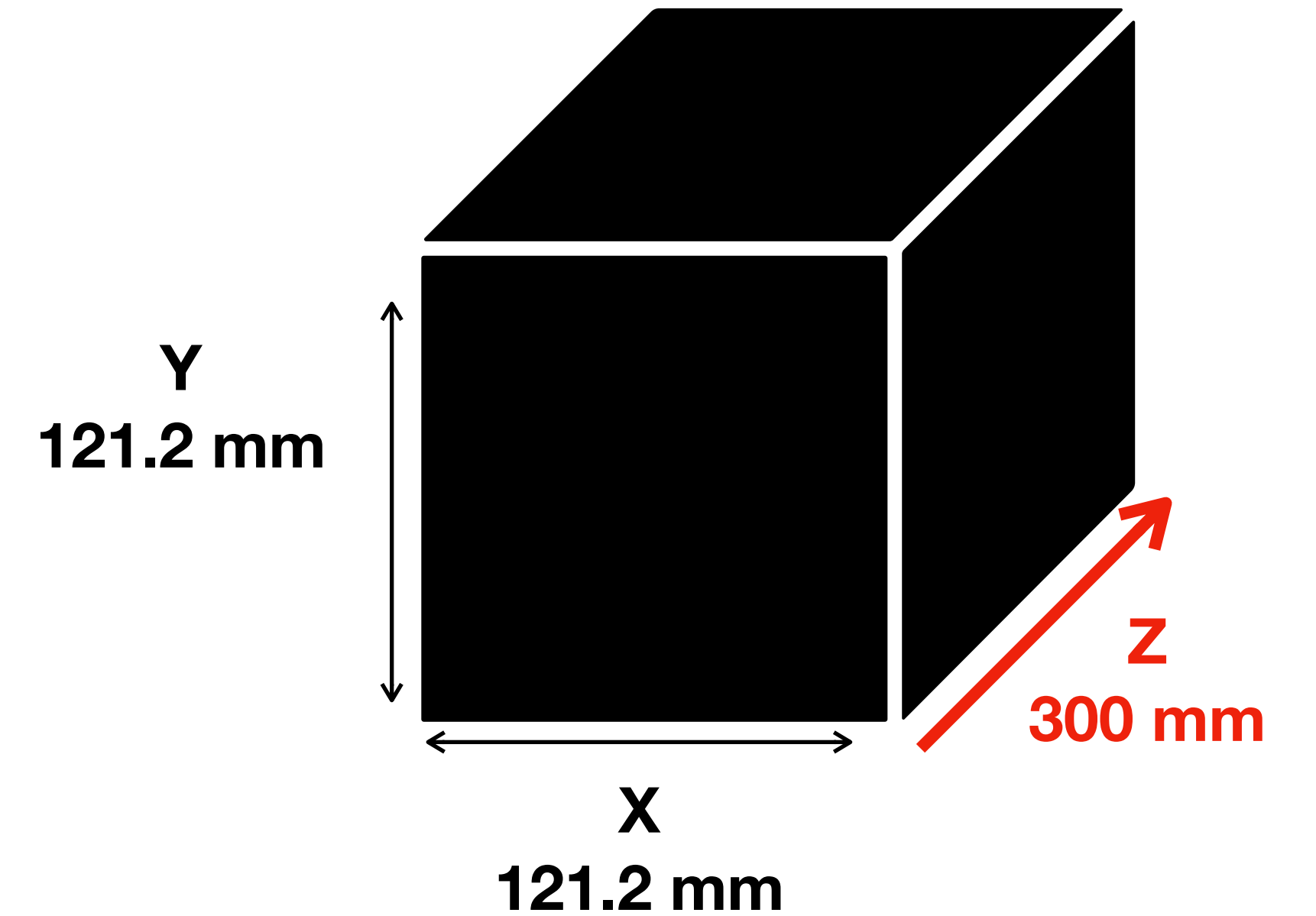


Scheme of a SPACAL module featuring a double-side readout
(“Framework TDR for the LHCb Upgrade II: Opportunities in flavour physics, and beyond, in the HL-LHC era.” <https://inspirehep.net/literature/2707810>)

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- **Goal:** study the **time resolution** of a simulated module using GEANT4
- 1000 incident e^- at 1 GeV and 10 GeV
- $3^\circ/3^\circ$ angles (w.r.t. x and y axes)
- Module under study: **Pb + Polystyrene**
- **Single side readout** (back)
- 3 x 3 cells, each one read out by a PMT



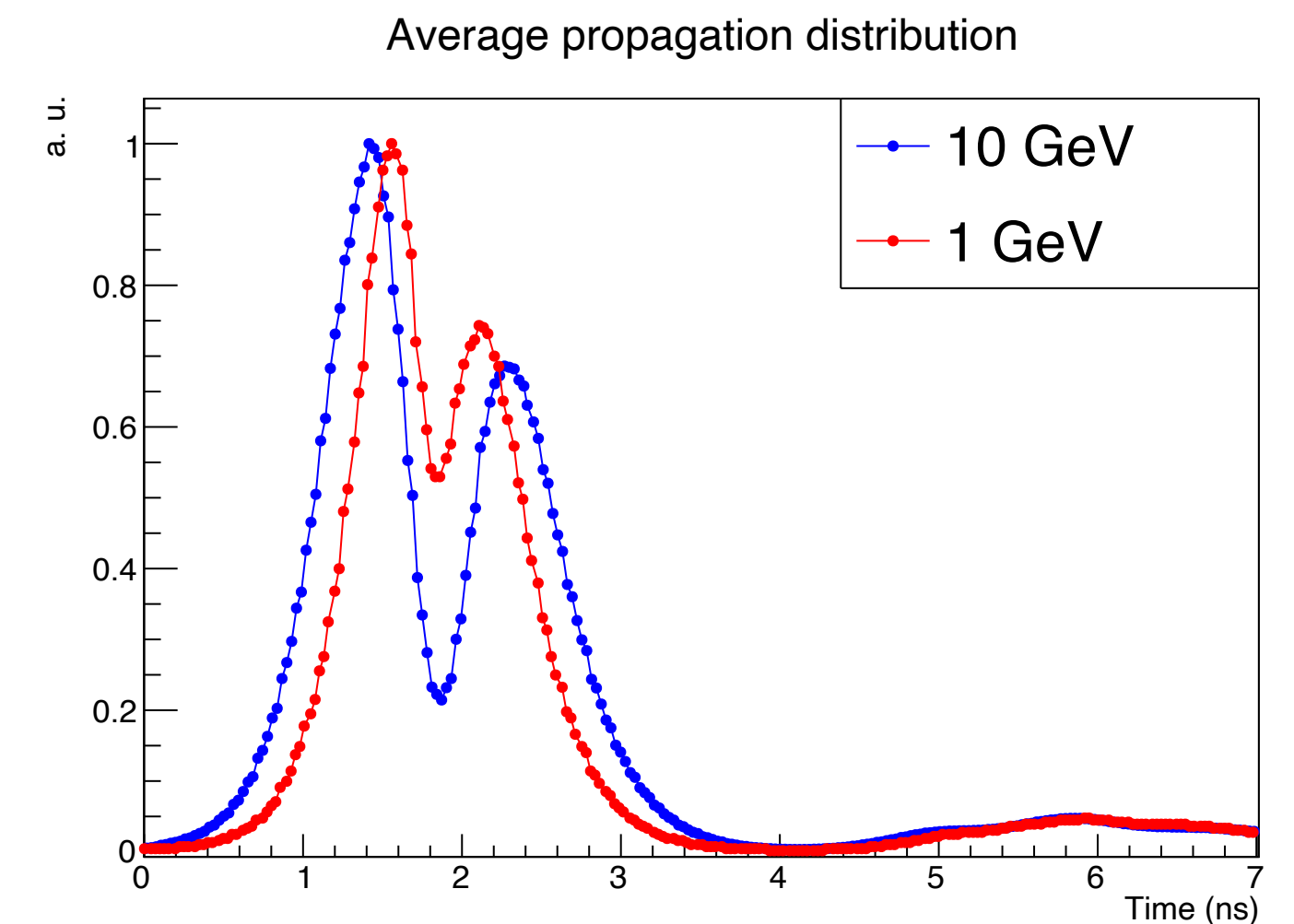
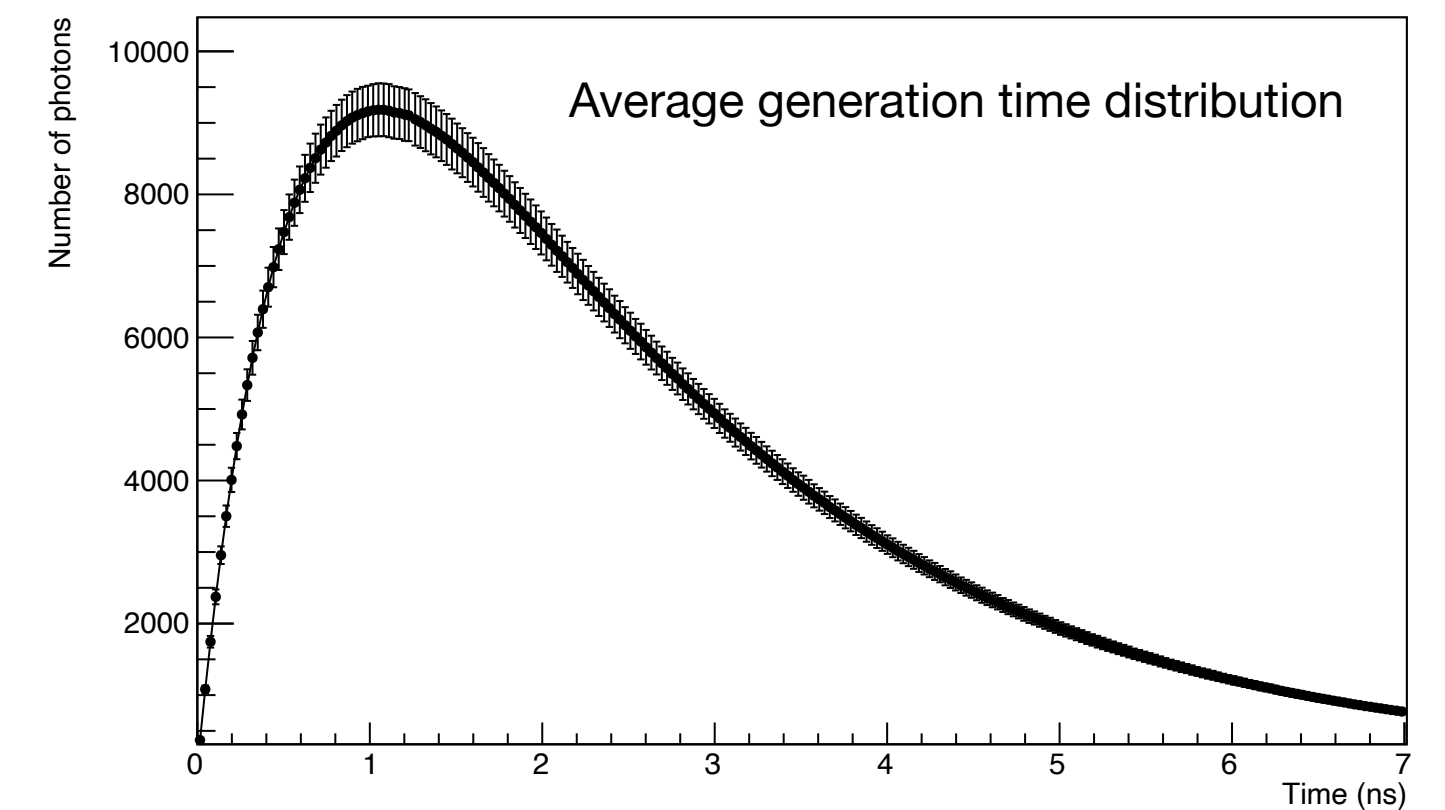
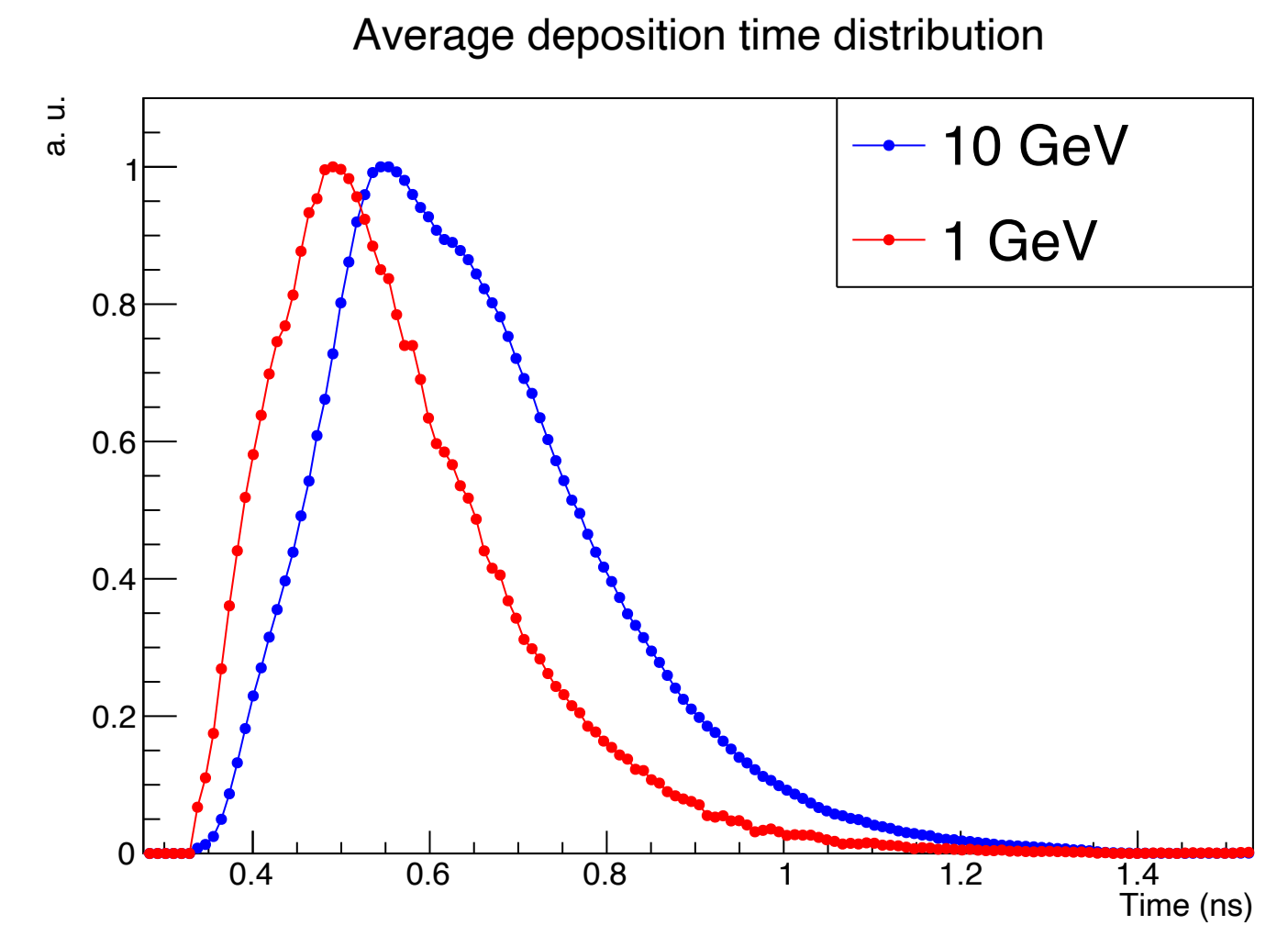
Simulation of the photons

Each optical photon has:

- **Deposition time** = time stamp of the single energy deposition which triggers the scintillation
- **Generation time** = time required by the scintillation process to generate the photon
- **Propagation time** = time to reach the PMT window

$$t_{total} = t_{deposition} + t_{generation} + t_{propagation}$$

(where the incident electron is created at $t = 0$)



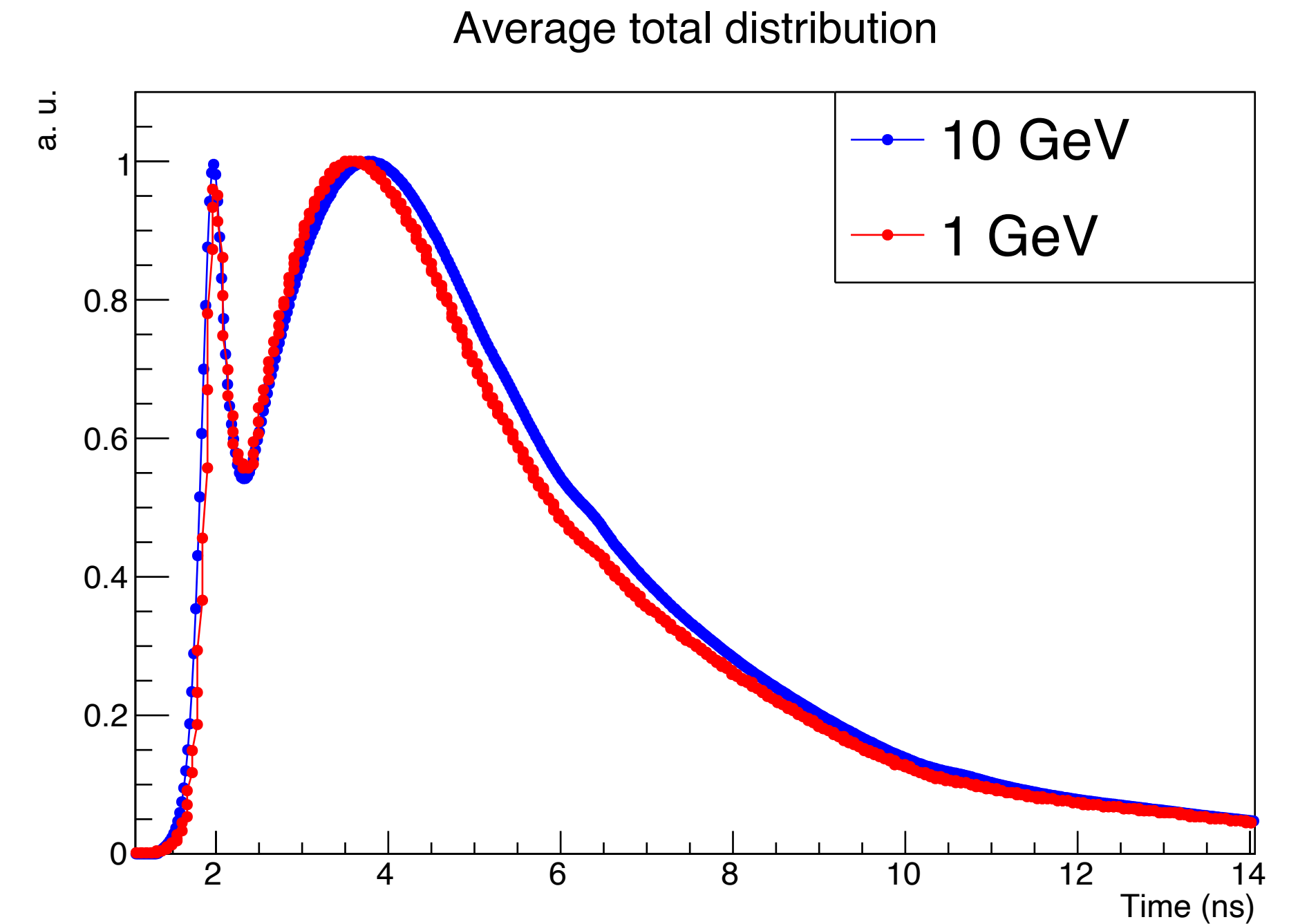
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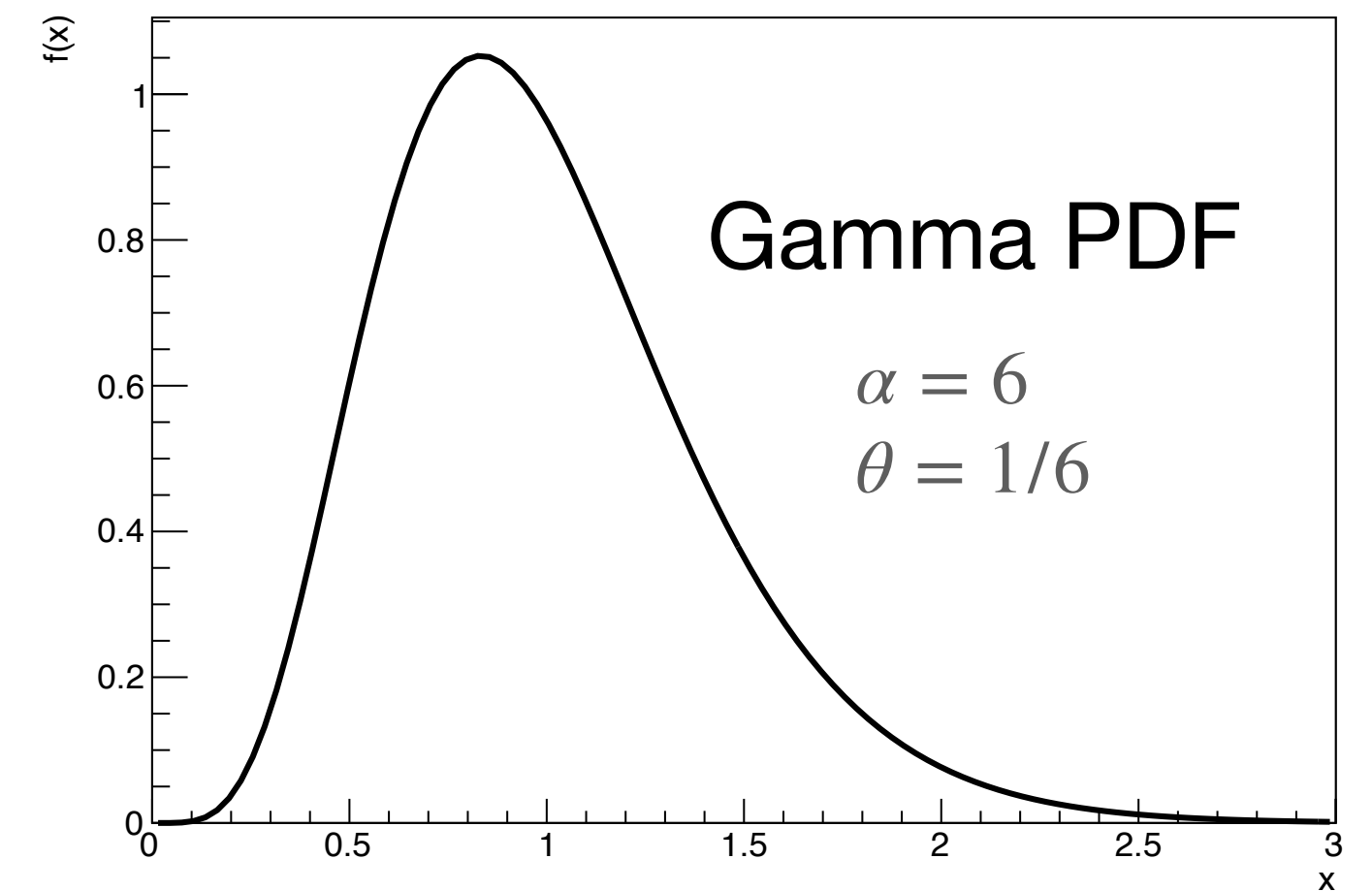
$$t_{total} = t_{deposition} + t_{generation} + t_{propagation}$$

(where the incident electron is created at $t = 0$)



PMTs simulations

- 1024 samples in 204.8 ns time window $\longrightarrow \Delta t = 200 \text{ ps}$
- SPTR (or TTS, Transit Time Spread)
- Amplitude fluctuations following a Gamma pdf: $p(x) = \frac{1}{\Gamma(\alpha) \cdot \theta^\alpha} x^{\alpha-1} e^{-x/\theta}$



Single photoelectron pulse:

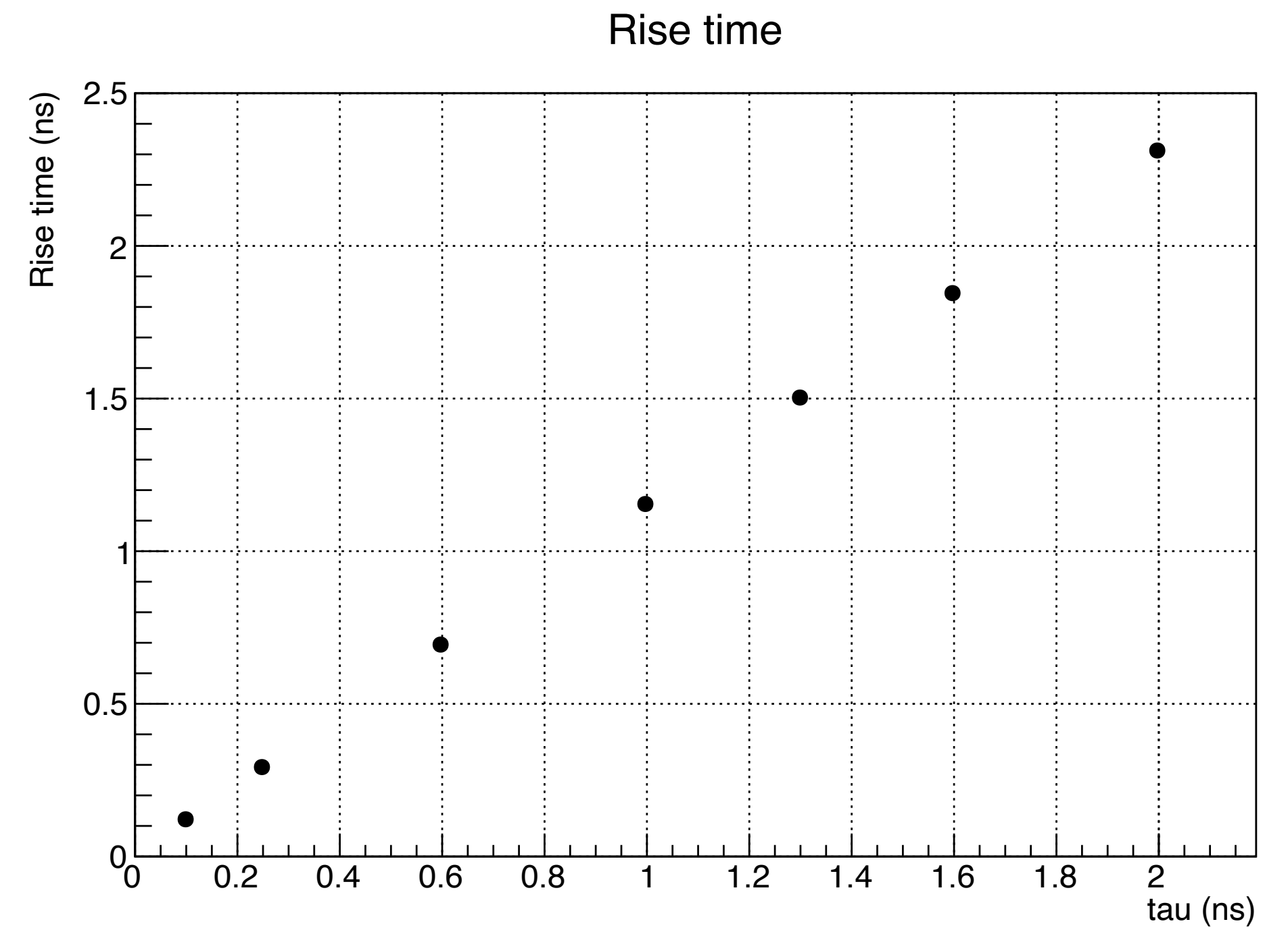
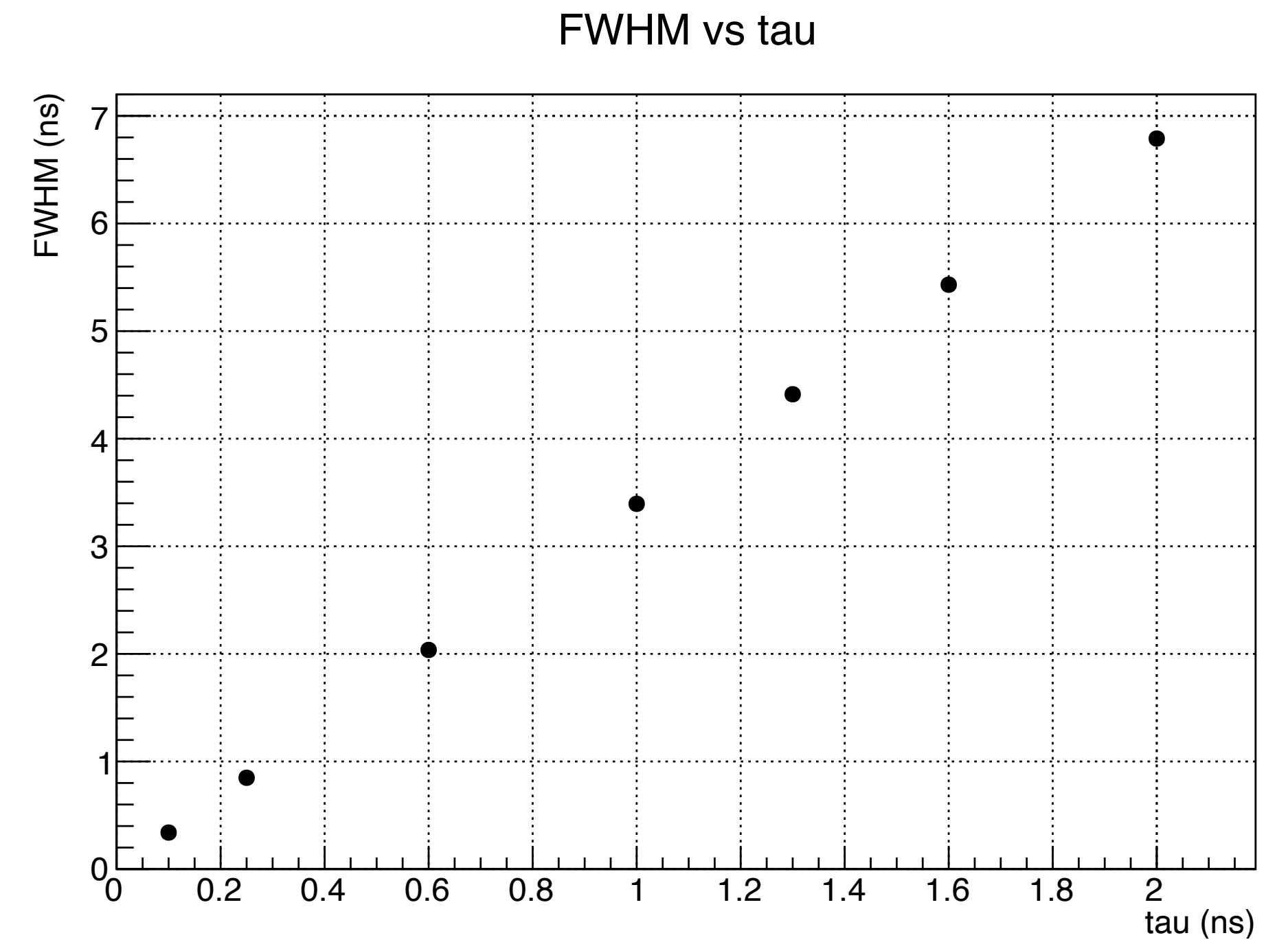
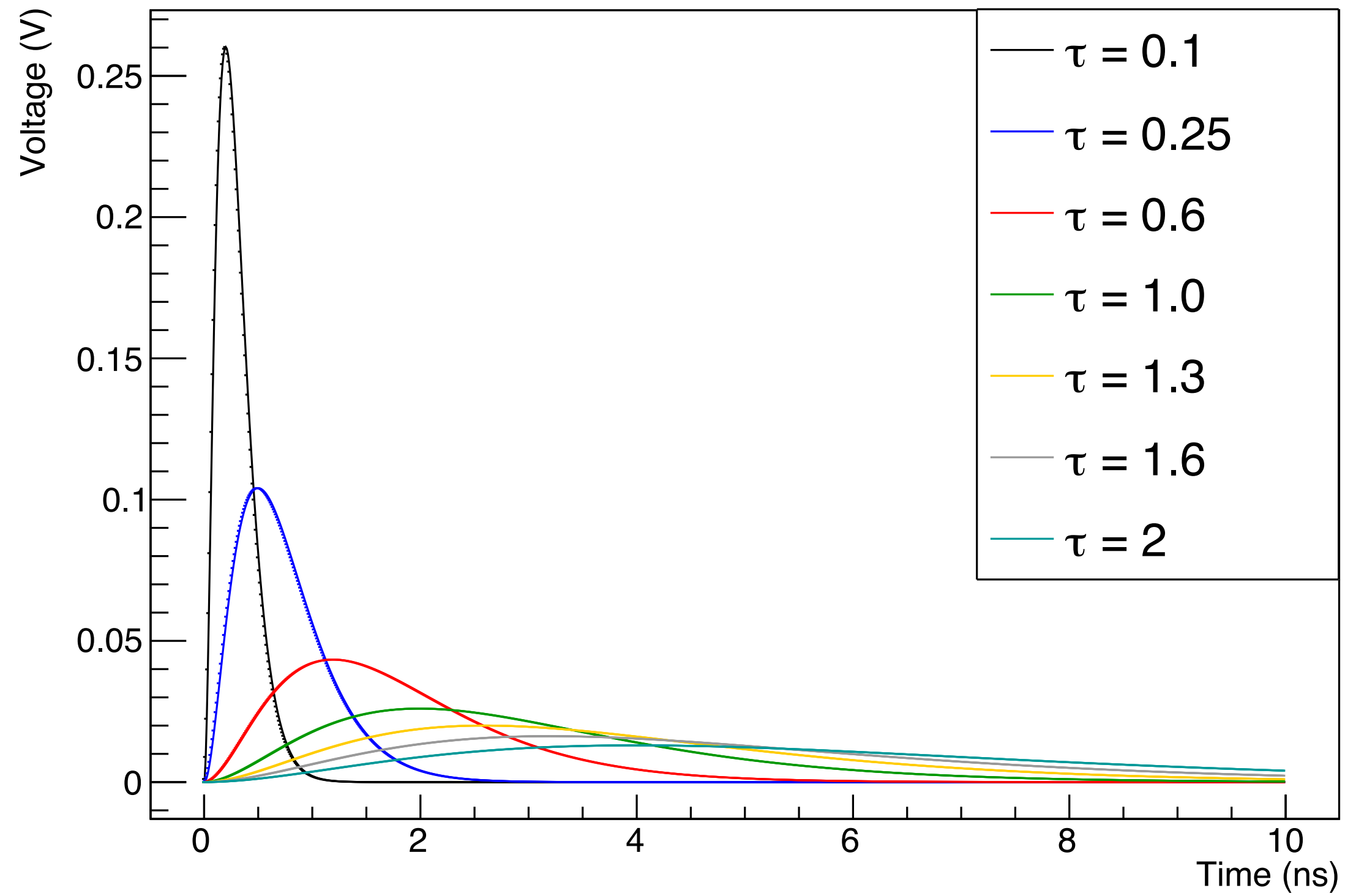
$$f(t) = A \cdot t^2 \cdot e^{-t/\tau} \quad \left| \quad A = \frac{R \cdot \text{gain} \cdot q_e}{\tau^3} \cdot 10^9$$

(FWHM proportional to tau)

Single p.h.e. pulses

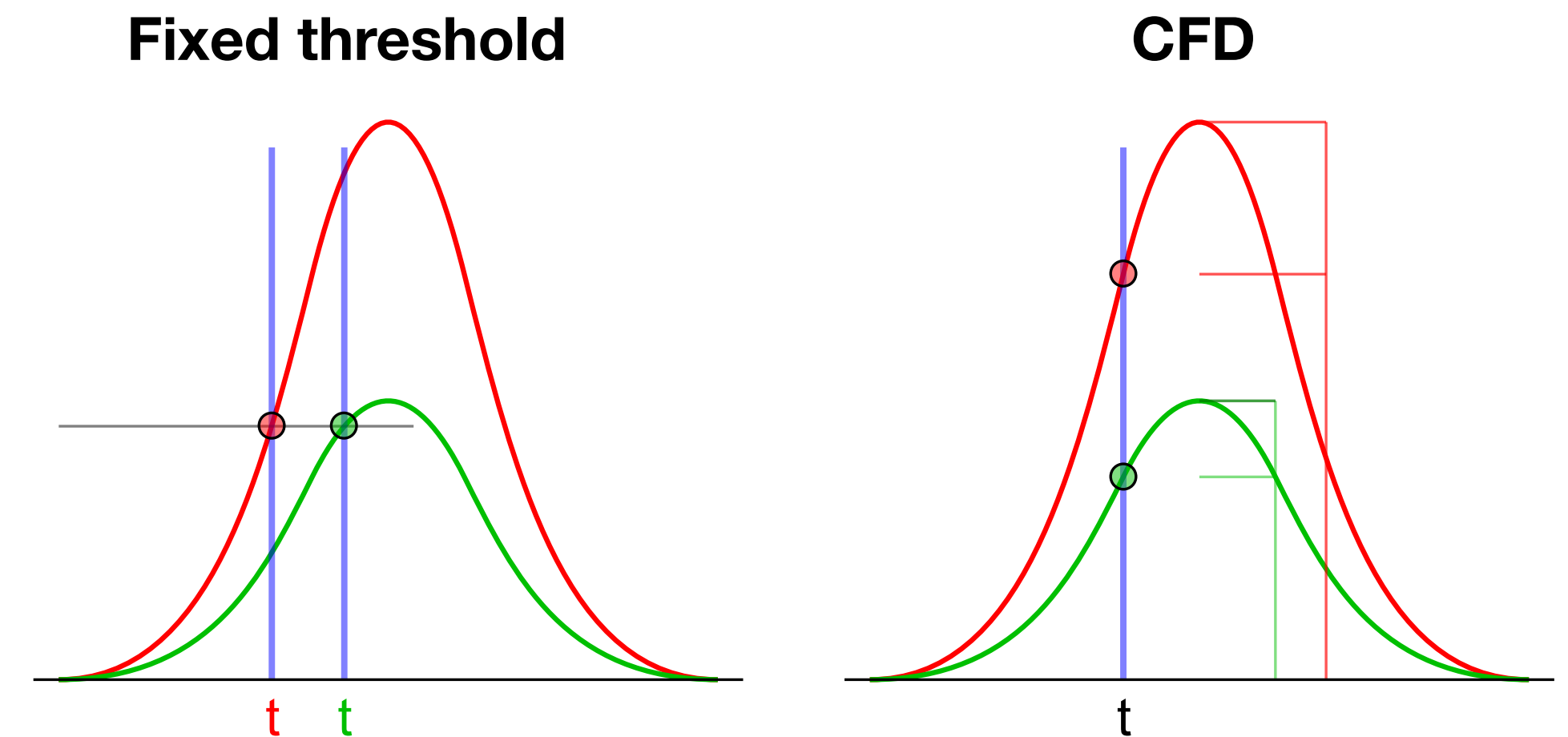
FWHM, rise time and increase linearly with tau

Single phe pulses



What is the “time stamp” of a signal?

- Time stamp computed with the “**C**onstant **F**raction **D**iscriminator” (**CFD**) algorithm
- **Time stamp** = time at which the signal exceeds a defined fraction of the pulse’s amplitude
- The “best” fraction must be properly chosen in order to optimize the time resolution
- **Time resolution** = std. dev. of the time stamps sample



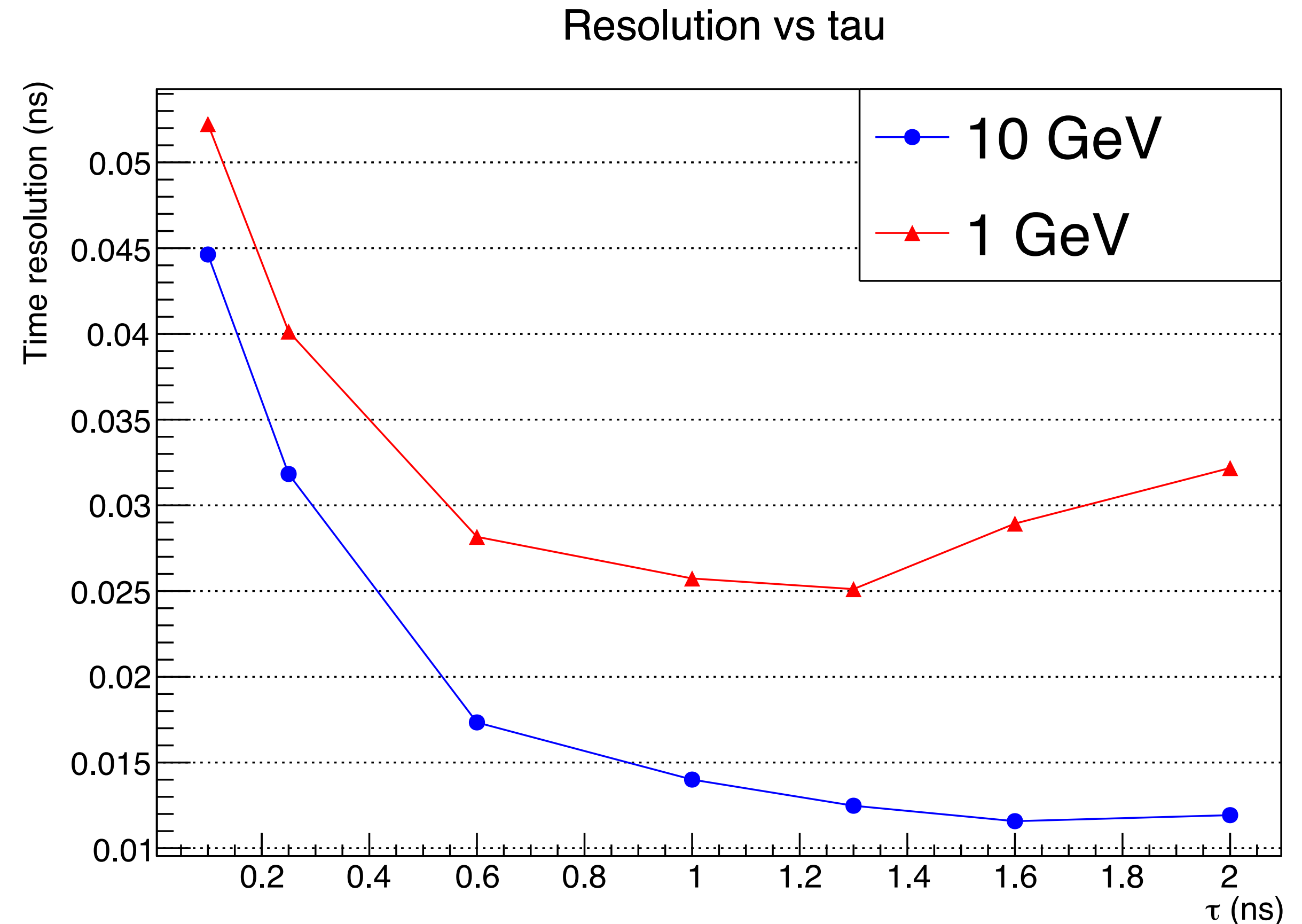
https://en.wikipedia.org/wiki/Constant_fraction_discriminator

First results

- As expected, better resolution at higher energies (photostatistics contribution)
- Slow PMTs perform better

Why?

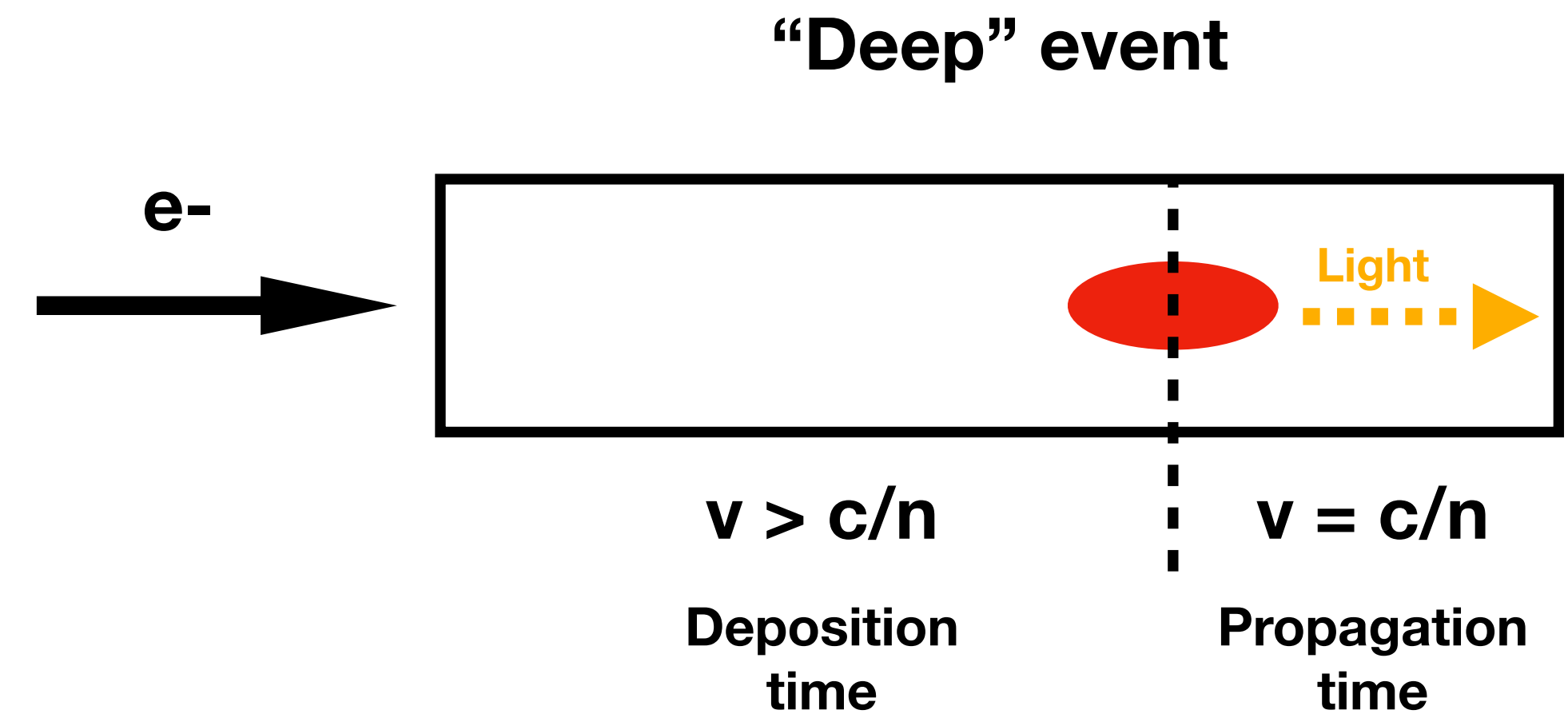
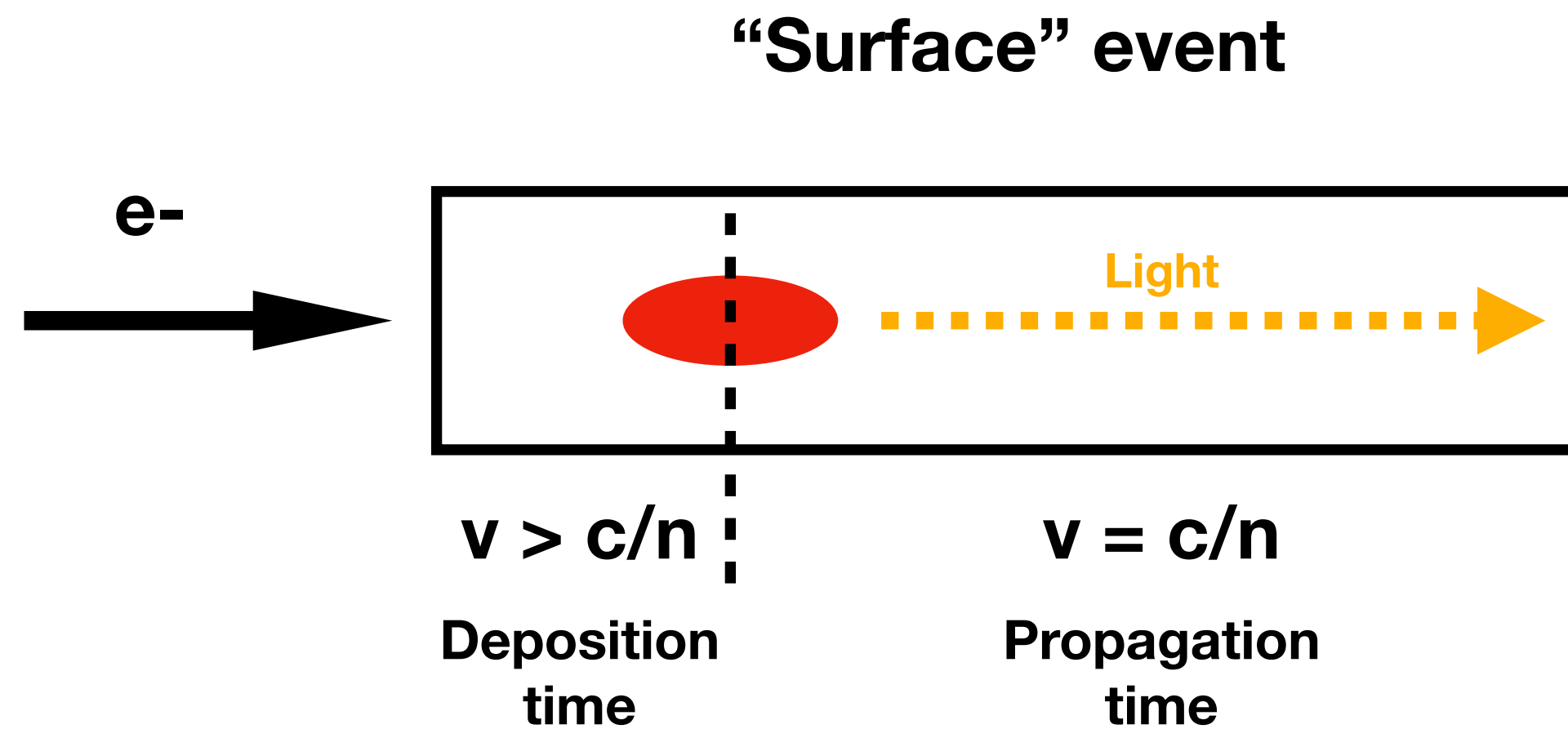
➔ Slow PMTs are less affected by the longitudinal fluctuations of the showers



Why are slower PMTs better ?

- Shower depth and time stamp are correlated. For **deeper showers**:
 - **Direct photons arrive earlier** to the PMT → Negative correlation

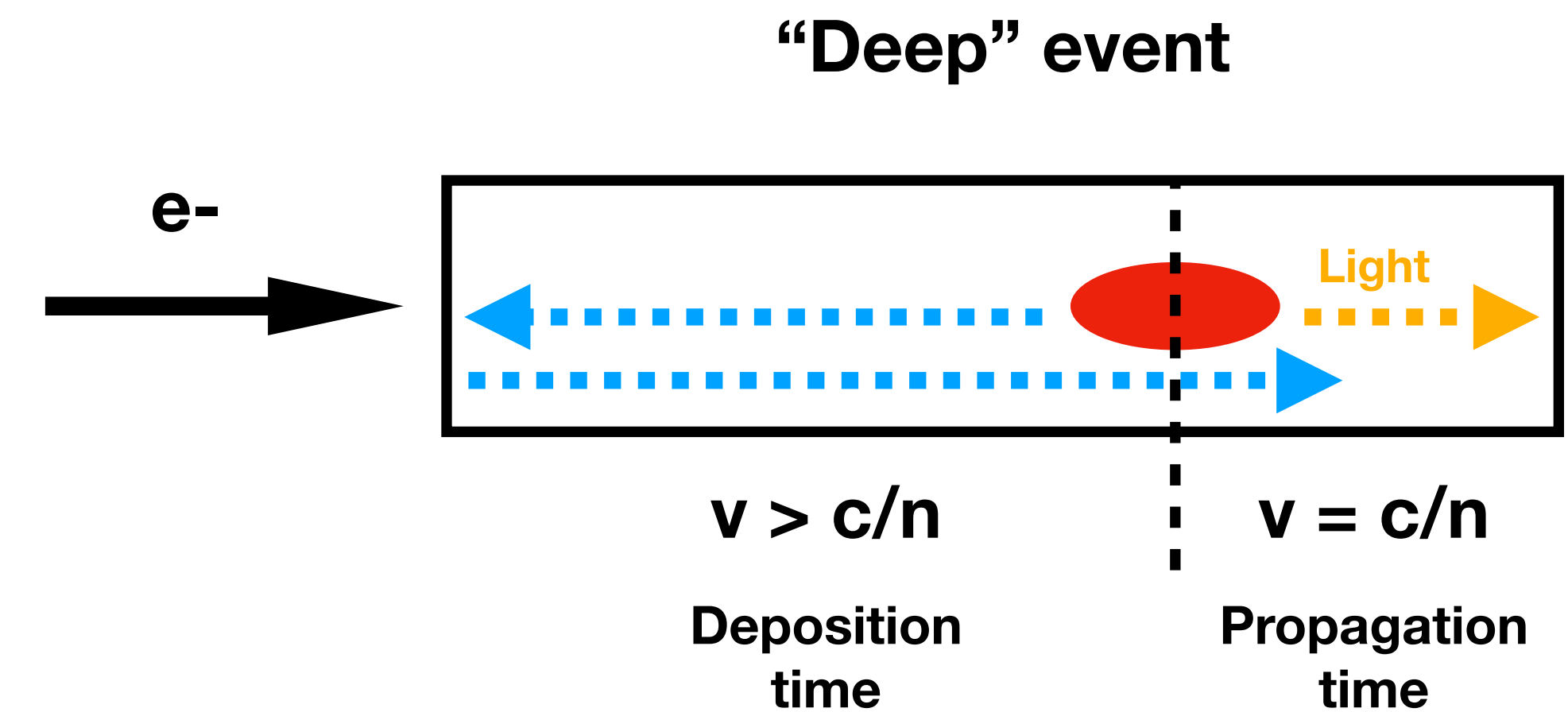
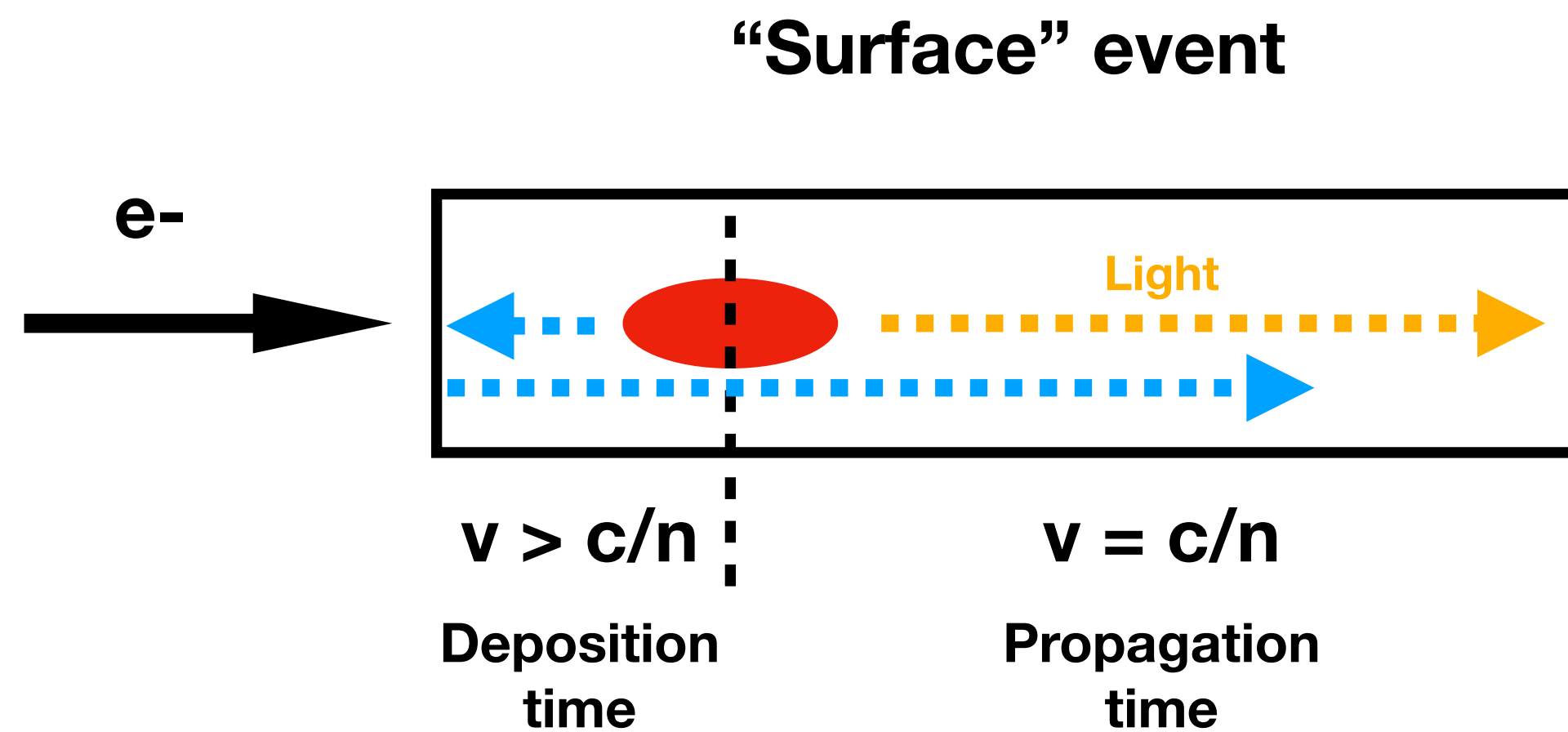
- *Barycenter of the energy depositions*
- *Direct photons*



Why are slower PMTs better ?

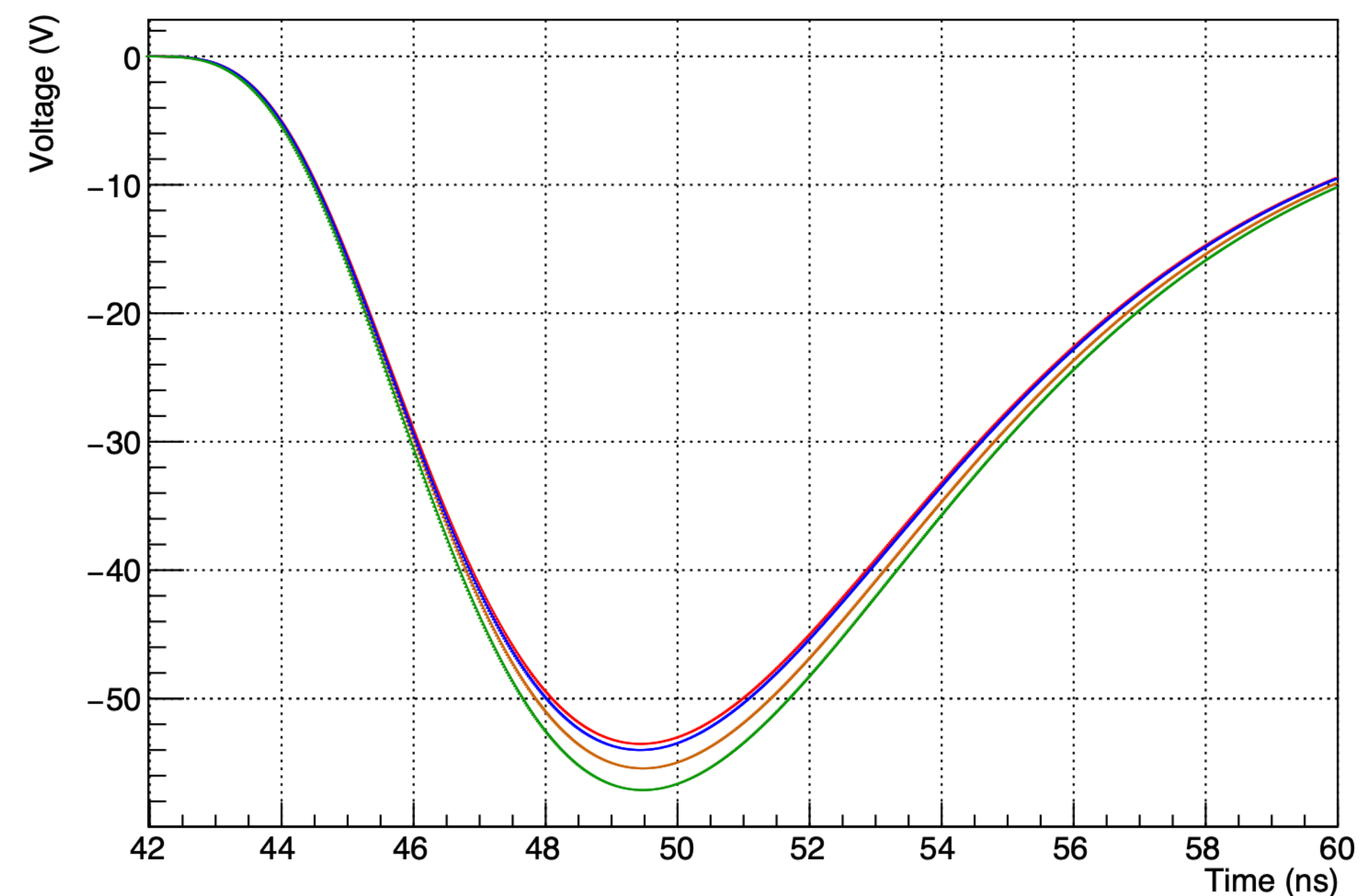
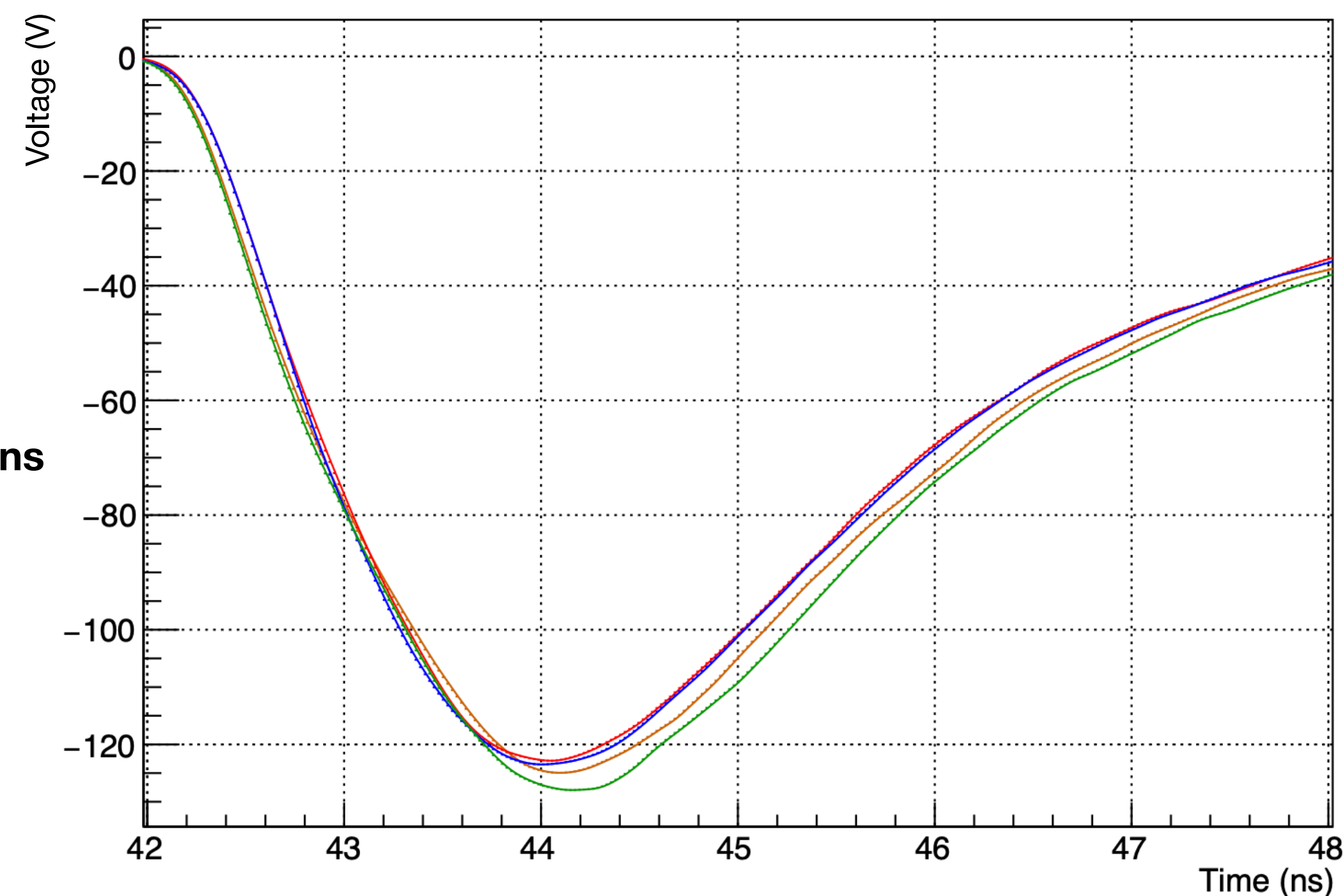
- Shower depth and time stamp are correlated. For **deeper showers**:
 - ▶ **Direct photons arrive earlier** to the PMT → Negative correlation
 - ▶ **Reflected photons arrive later** → Positive correlation
- ➔ The CFD time stamp is biased by the shower depth
- ➔ This bias worsens the time resolution

- *Barycenter of the energy depositions*
- *Reflected photons*
- *Direct photons*



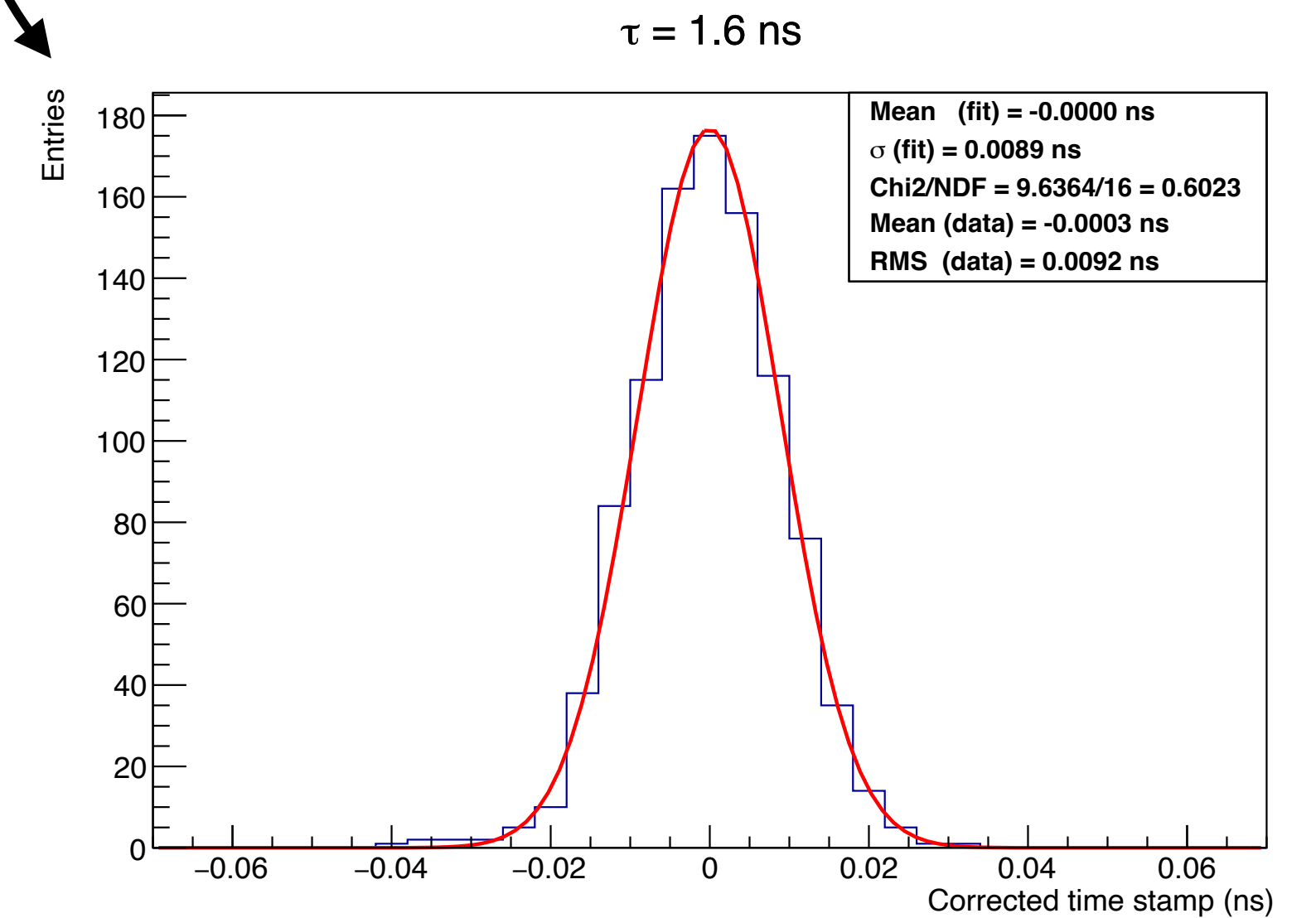
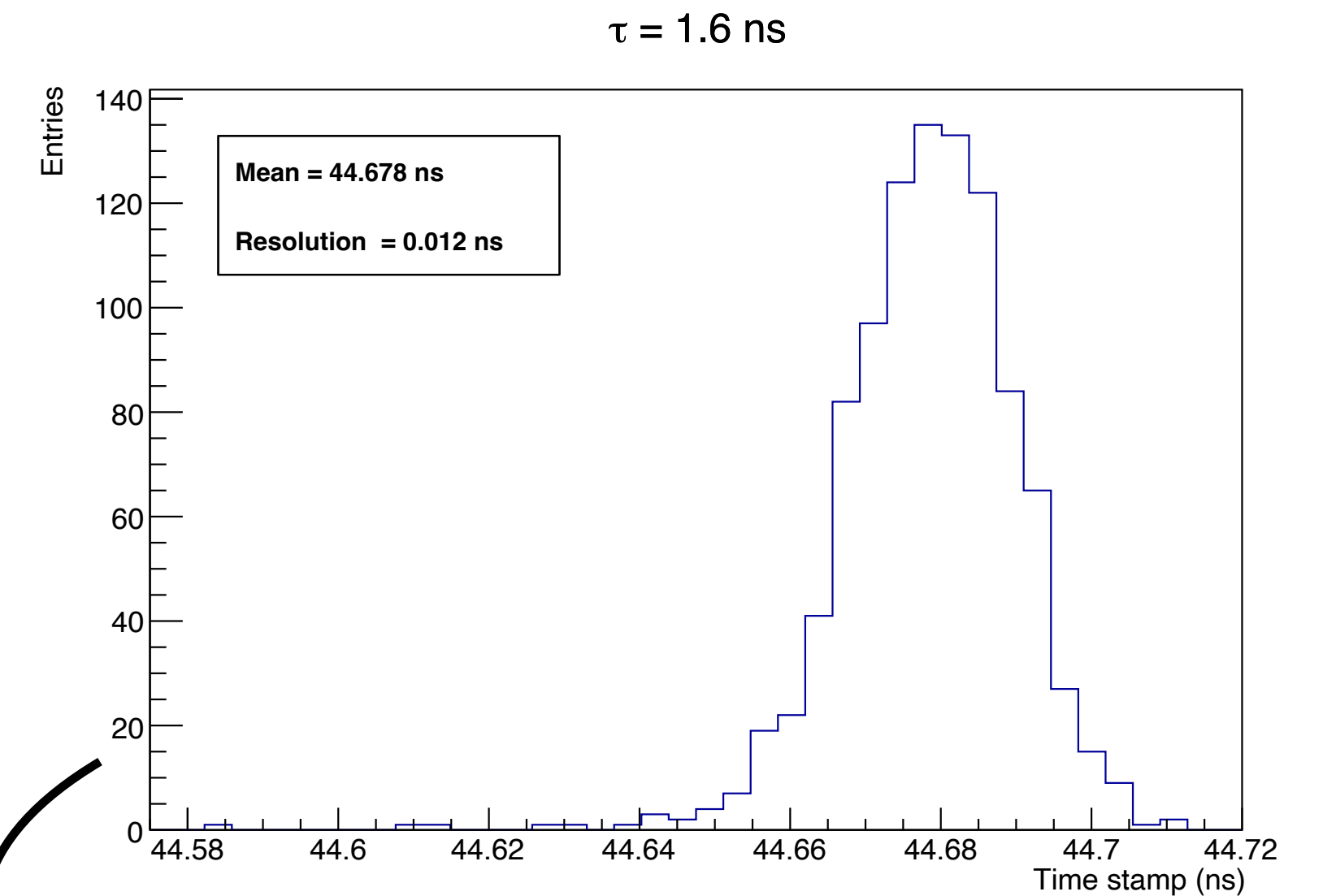
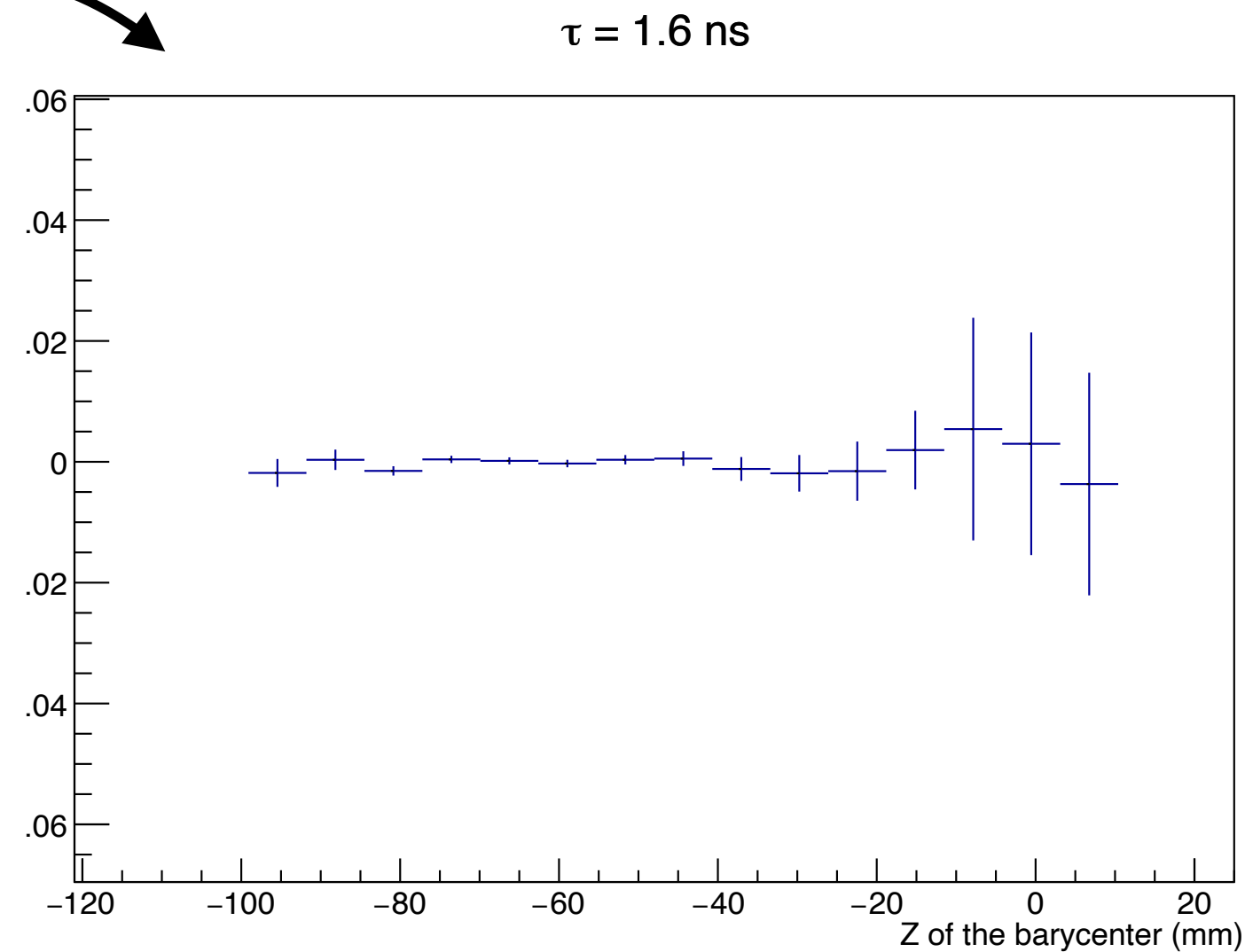
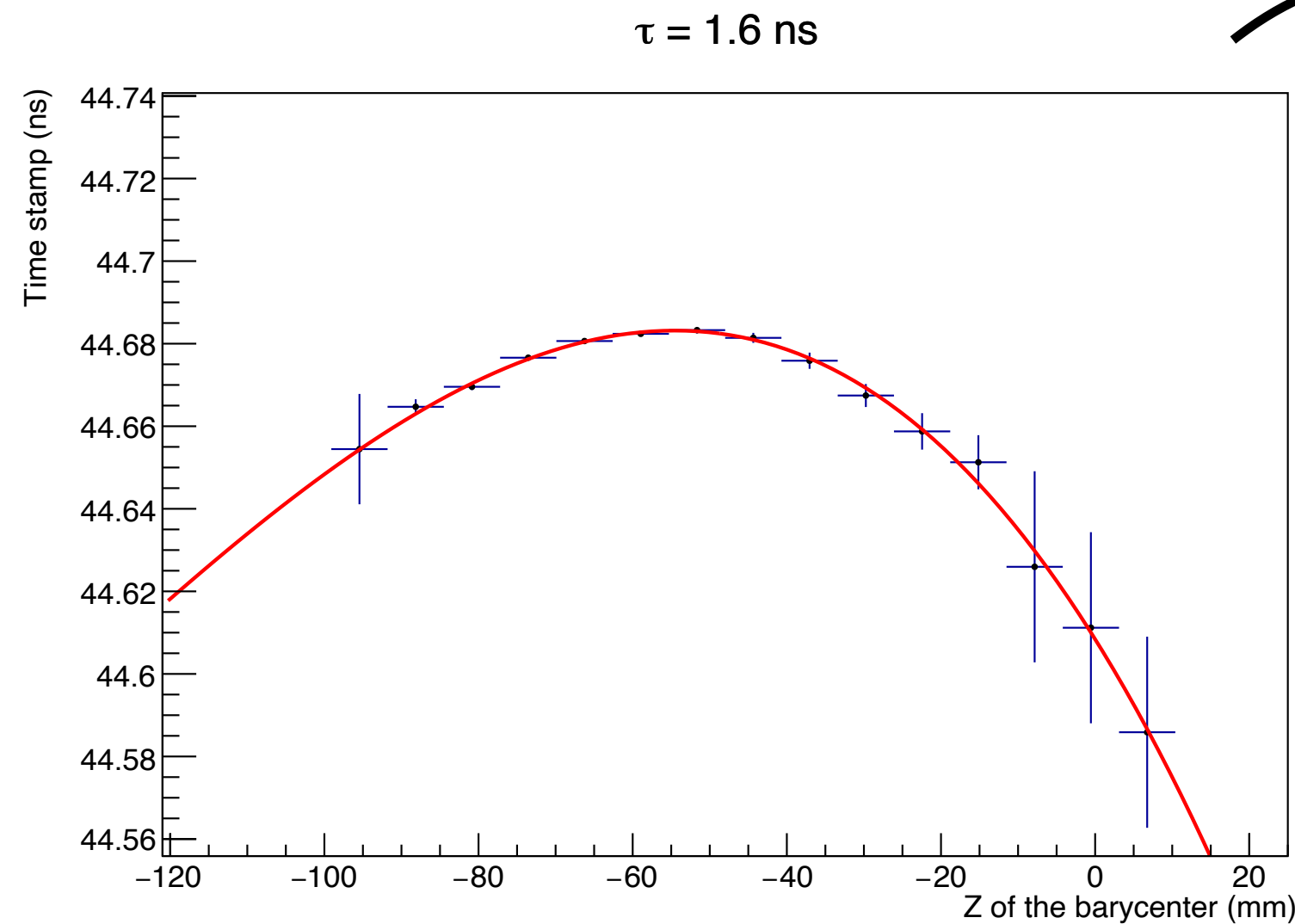
Why are slower PMTs better ?

- This effect is **more relevant for fast PMTs** (they better distinguish between direct and reflected photons)
- It **affects the shape of the PMTs signals** → The CFD method can't take it into account
- It **depends on the CFD threshold**
 - Low thresholds mostly detect direct photons
 - For some thresholds the two correlations partially cancel out each other, removing the overall bias



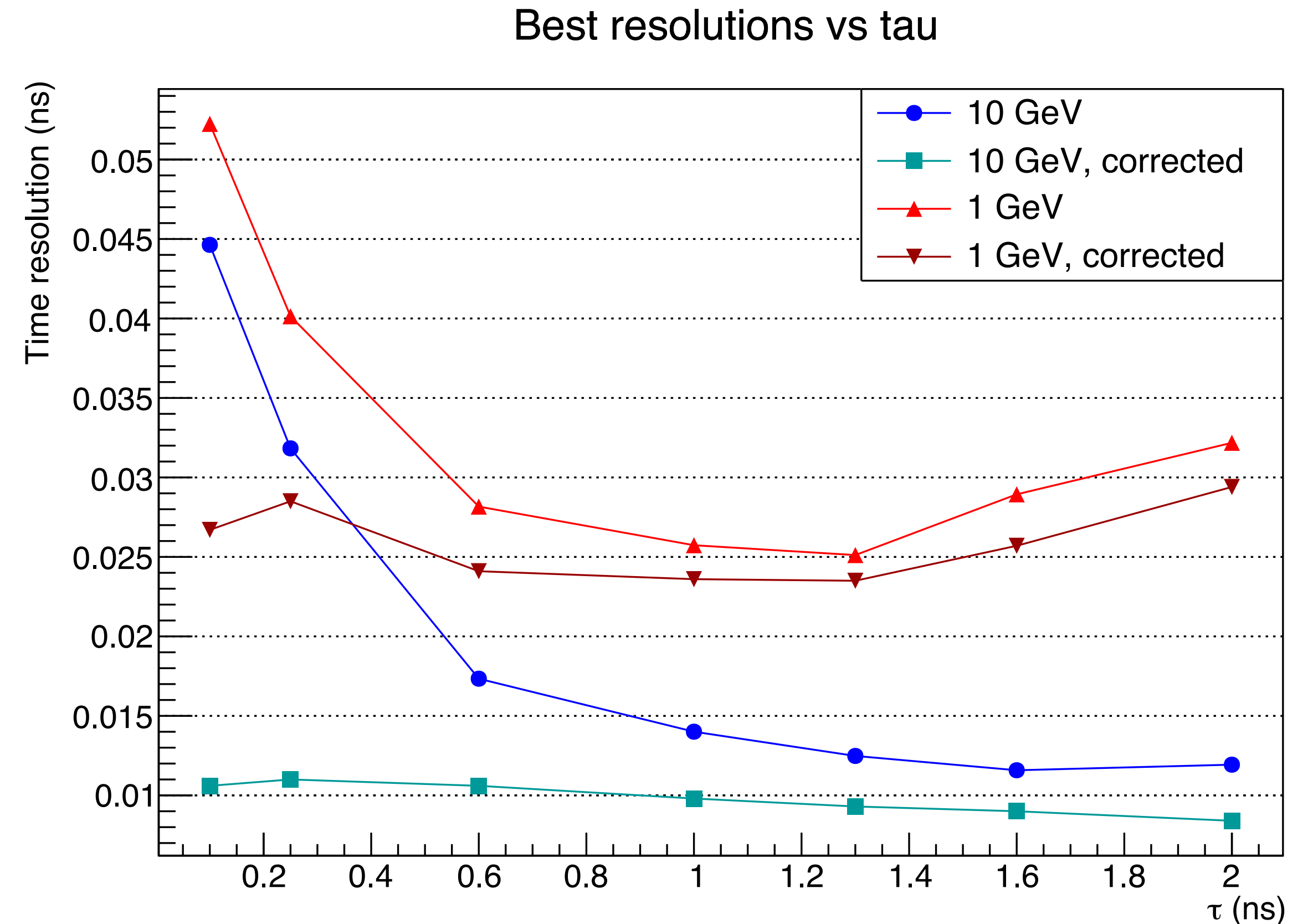
Correction procedure

- **Polynomial fit** to the profiled scatter plot of time stamp t vs shower depth of each event
 - ➔ Find the **correction curve** f
- **Corrected time stamp** for the j^{th} event defined as: $\hat{t}_j = t_j - f_j$
- The best CFD threshold after the bias correction may be different



Results

- **Corrected resolution** = std. dev. of the unbiased time stamps \hat{t}
- **Faster PMTs (lower τ) undergo wider corrections**
- The best CFD threshold after the correction is **always ~ 10% or ~ 90%**
 - At these levels: correlation between time stamp and shower depth is maximum
 - Highest corrections



Here, the best CFD threshold is the optimal one for each point (different before/after the correction)

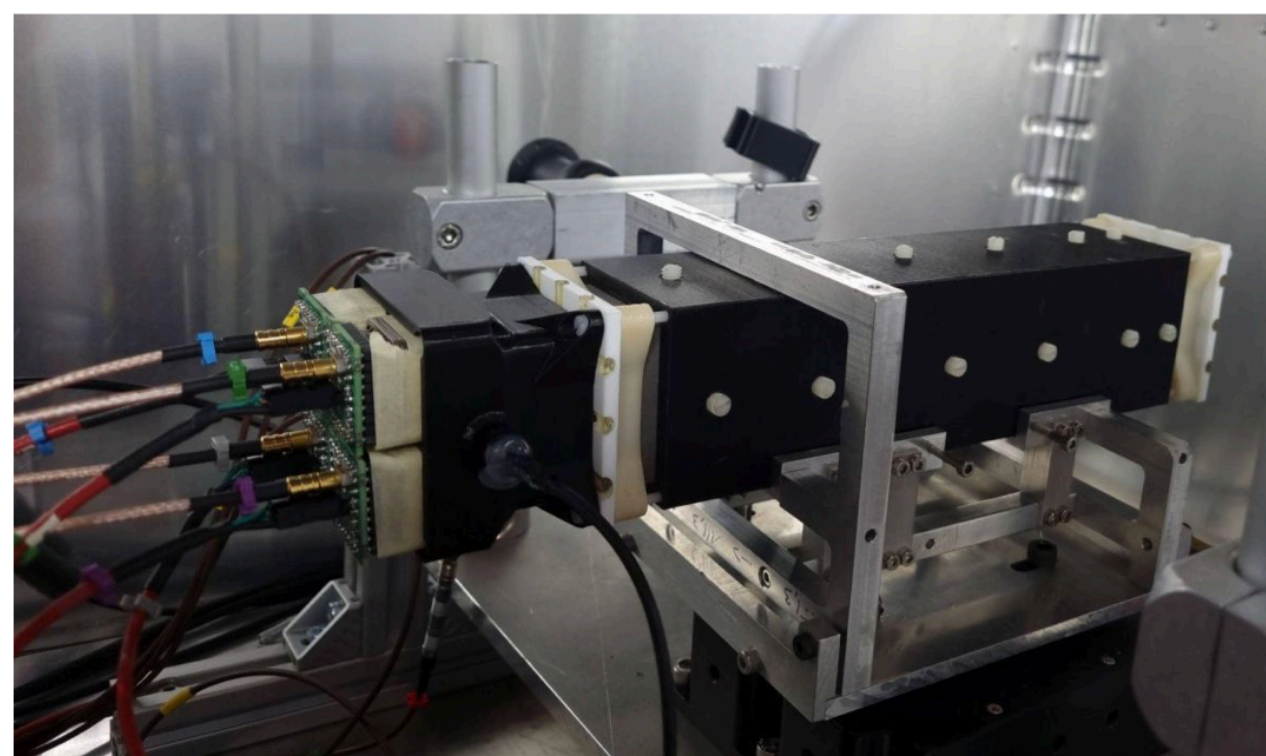
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Prototypes

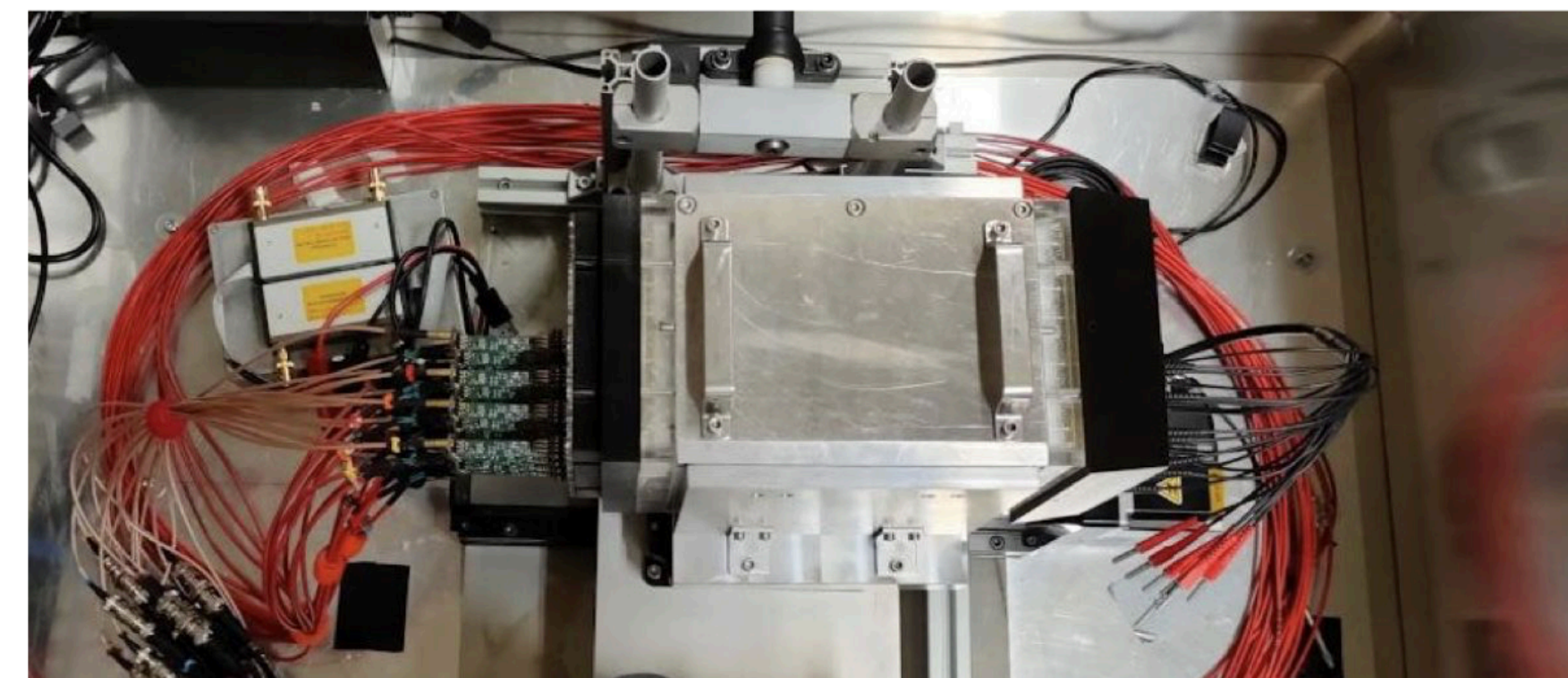
“Small module”

- SpaCal **W-Polystyrene**
- 4 cells only (4.5 x 4.5 cm²)
- 5 cm - square to octagon light guides
- Kuraray **SCSF-78 (blue)** or **3HF (green)** fibres
- Readout with 4 different PMTs:
 - **R7600U**, **R9880U**, **R14755U**, R11187 (a.k.a. Tilecal)



“Module 0”

- SpaCal **W-Polystyrene**
- Full-size module: 16 cells (12.1 x 12.1 cm²)
- 10 cm - square to octagon light guides
- Kuraray **SCSF-78 (blue)** fibres
- Readout with **R9880U** PMT



Time resolution model

- The time resolution as a function of the number of photons impinging the PMTs is well described by

$$\sigma_T(N_{ph}) = \frac{a}{N_{ph}} \oplus \frac{b}{\sqrt{N_{ph}}} \oplus c$$

- Noise term
- Sampling term
- Constant term

- Assuming linearity: $N_{ph} \propto E$ (energy of the incident e^-)

$$\sigma_T(E) = \frac{a'}{E} \oplus \frac{b'}{\sqrt{E}} \oplus c'$$

Time resolution model

- **Noise term:** caused by the electronic noise fluctuations
- Faster PMTs (quicker rise time) lead to smaller noise terms
- When exploiting the CFD algorithm, it can be estimated as

$$\sigma_{T_{noise}} = \sqrt{\frac{2}{3}} \frac{\sigma_n}{dA/dt}$$

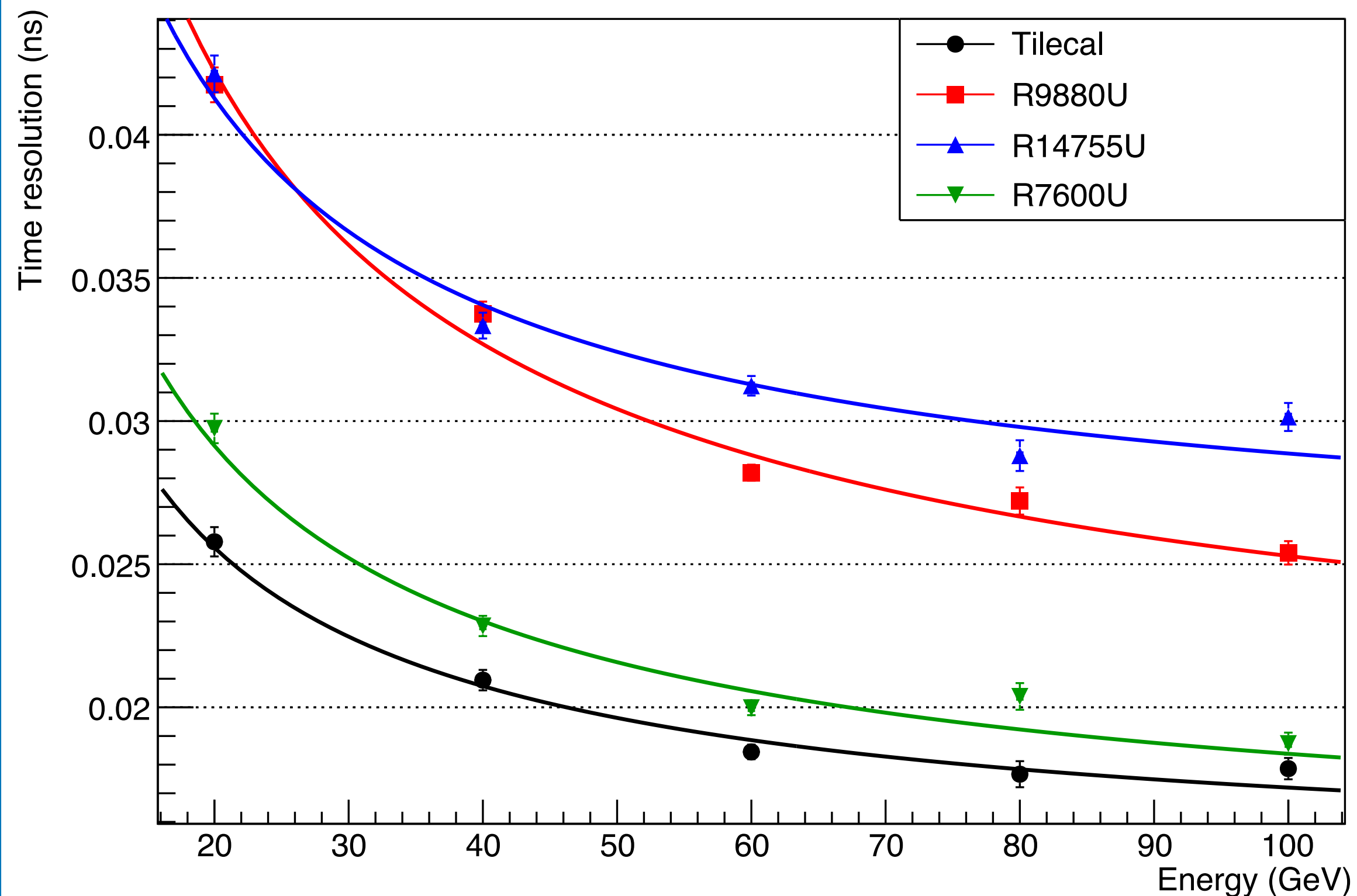
- σ_n = std. dev. of the electronic noise
- A = pulse's amplitude

- If it is subtracted in quadrature, the resolution as a function of the energy becomes

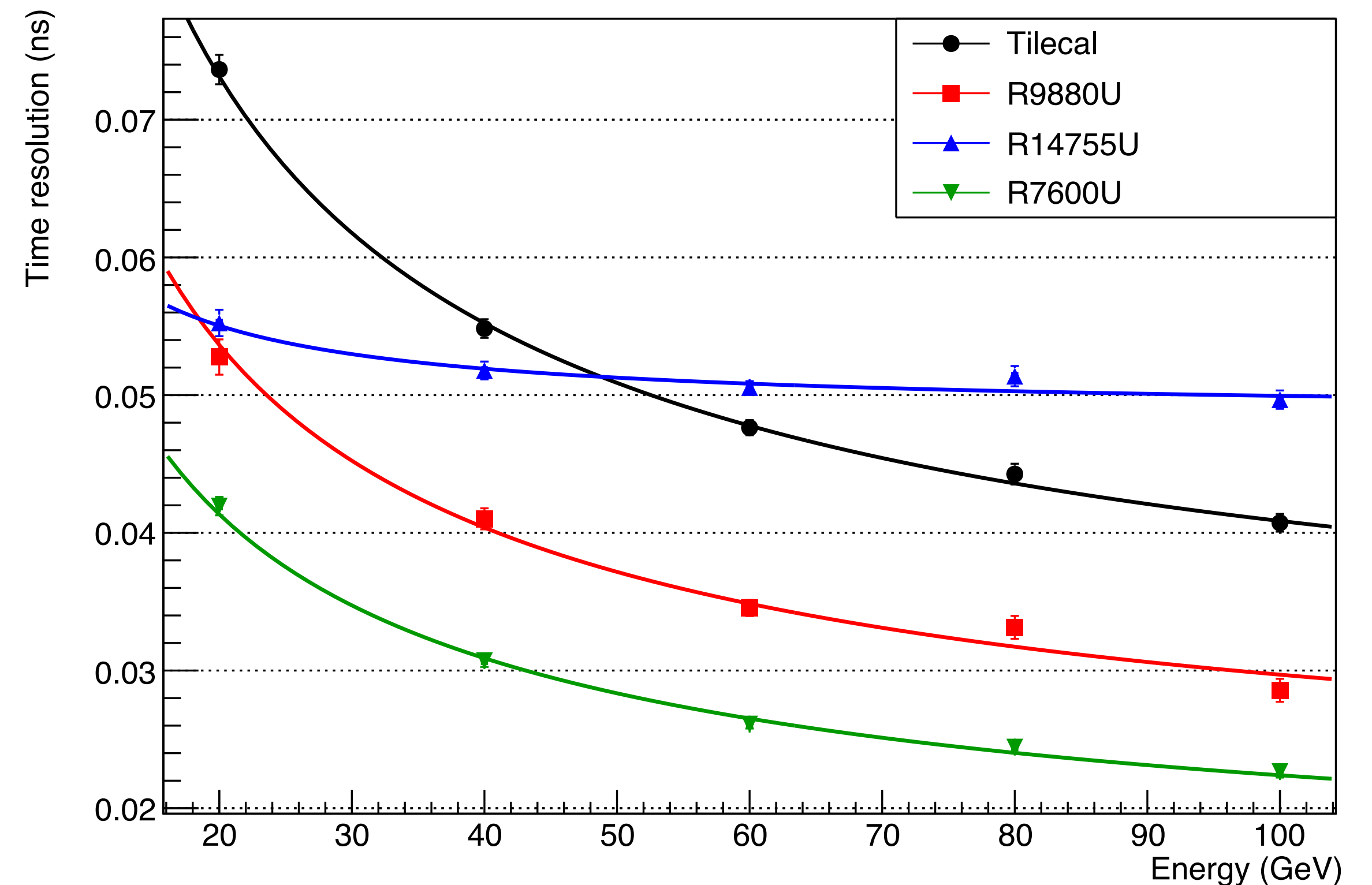
$$\sigma_T(E) = \frac{b'}{\sqrt{E}} \oplus c'$$

Time resolution results - Small module

Time resolution vs energy - SCSF-78 (blue) fibres



Time resolution vs energy - 3HF (green) fibres



- **SCSF-78** results are systematically better due to faster decay time of the fibres
- **SCSF-78** fibres: slow PMTs (**R7600U** and **Tilecal**) perform better → Less biased by shower depth
- **3HF** fibres: best results for PMTs with **Extended Red Multi Alkali (ERMA)** photocathode

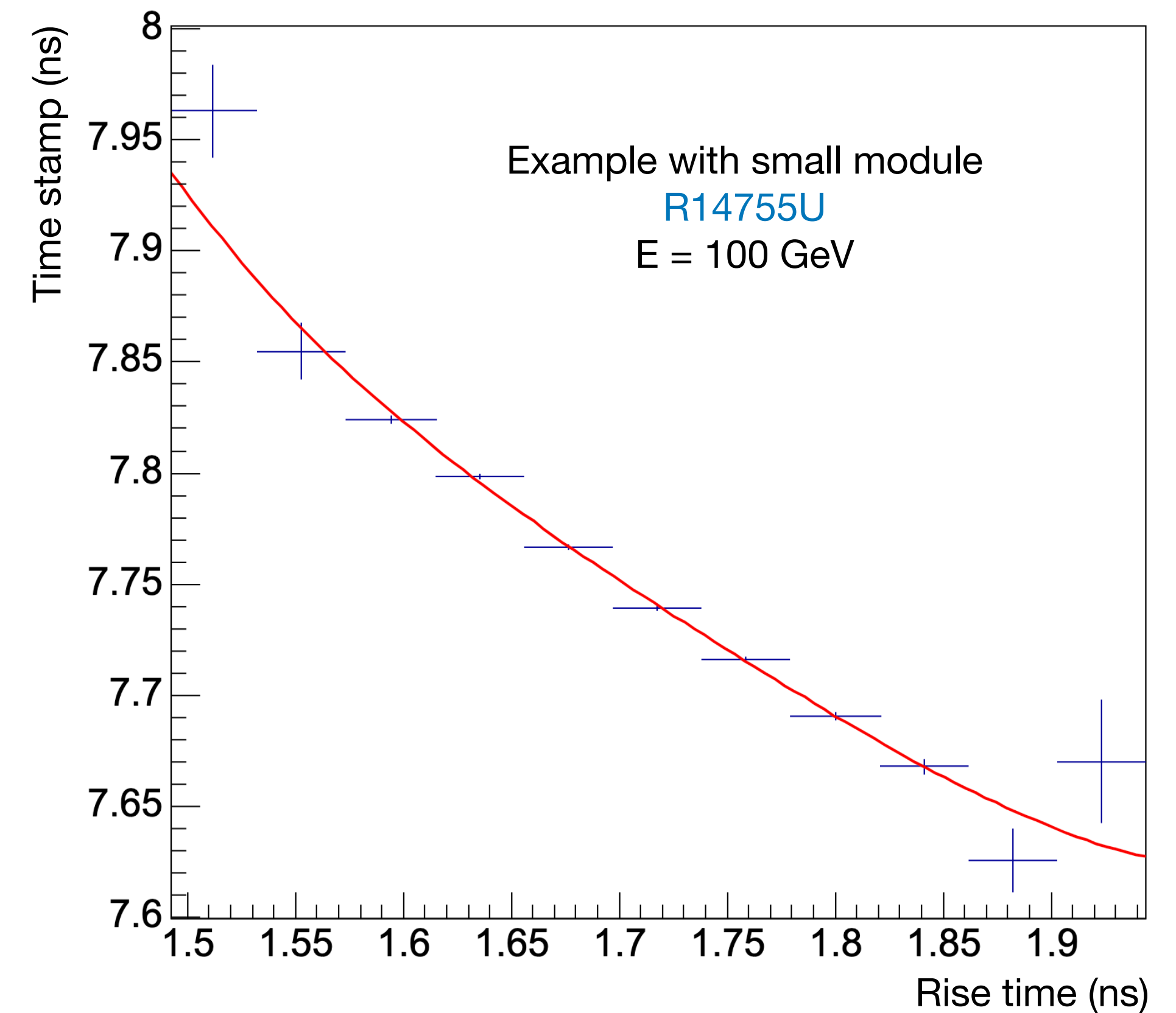
Correction to the time stamp

More info in backup slides

Simulations show that the rise time is highly correlated to the shower depth

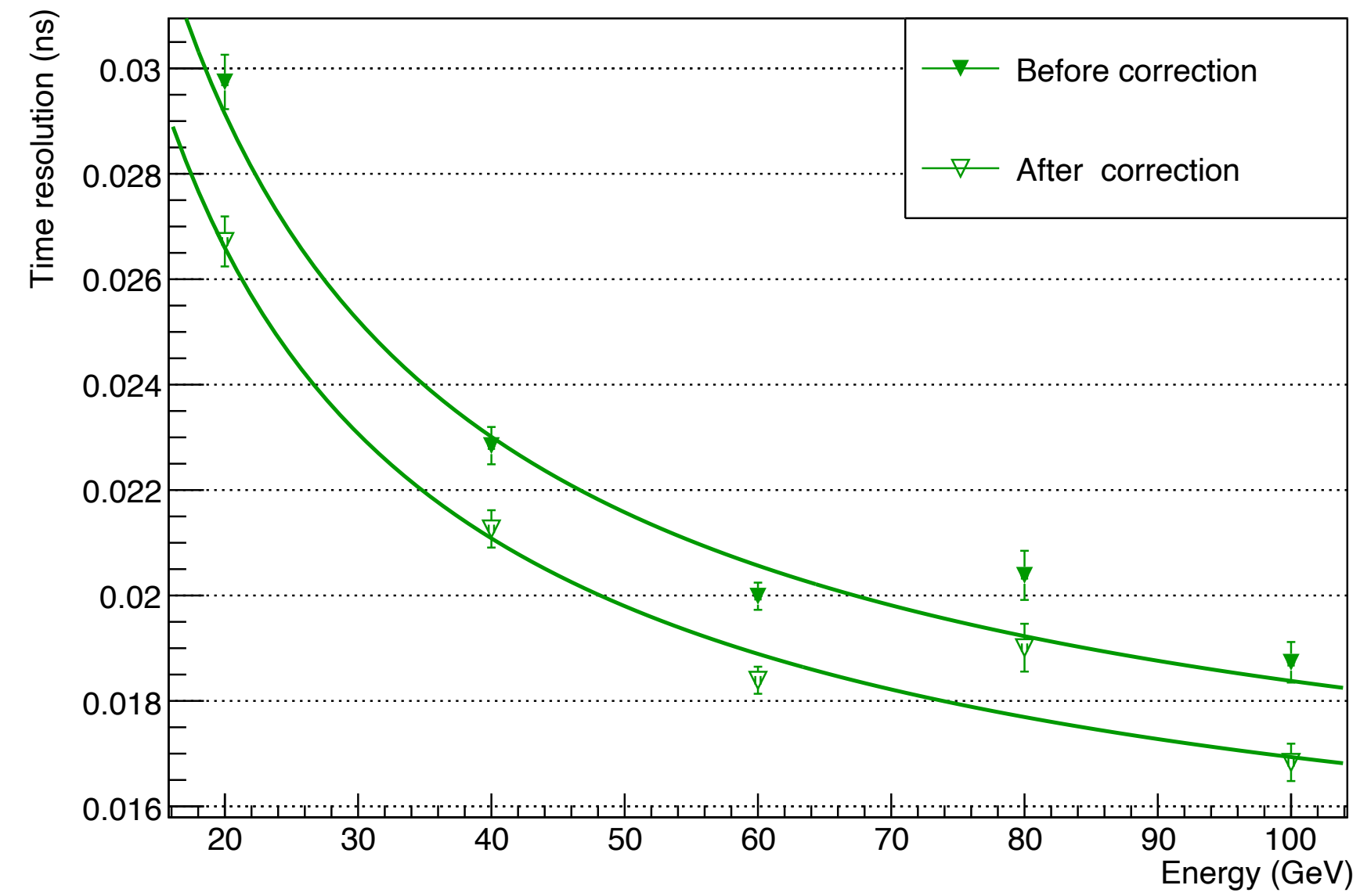
➔ **Idea: exploit the rise time to remove the bias**

- Polynomial fit to the profiled scatter plot of time stamp t vs rise time of each signal —> Find the **correction curve** f
- **Corrected time stamp** for the j^{th} event defined as: $\hat{t}_j = t_j - f_j$
- The corrected time resolution is the standard deviation of \hat{t}

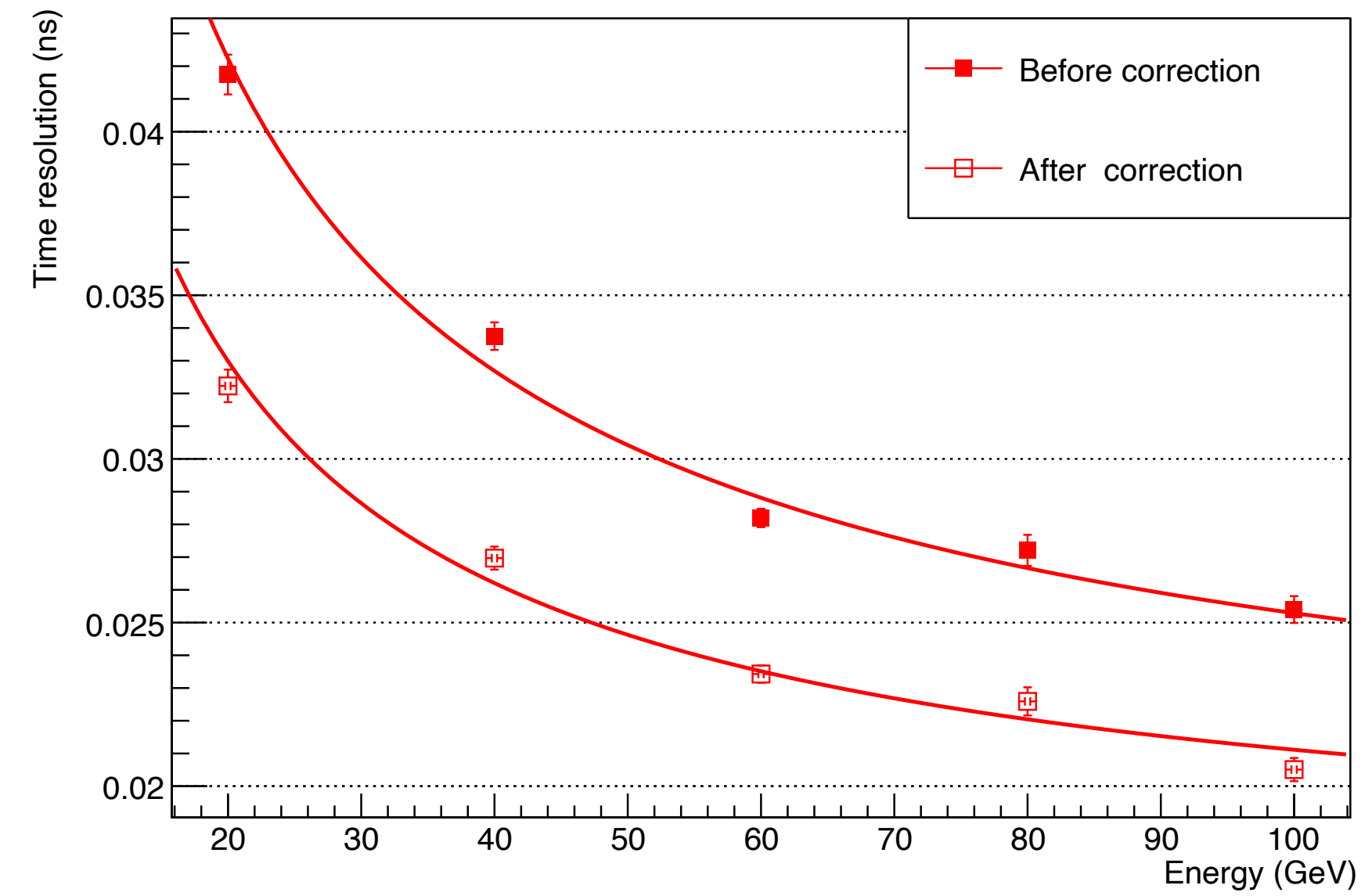


Corrected resolution - Small module with [SCSF-78 fibres](#)

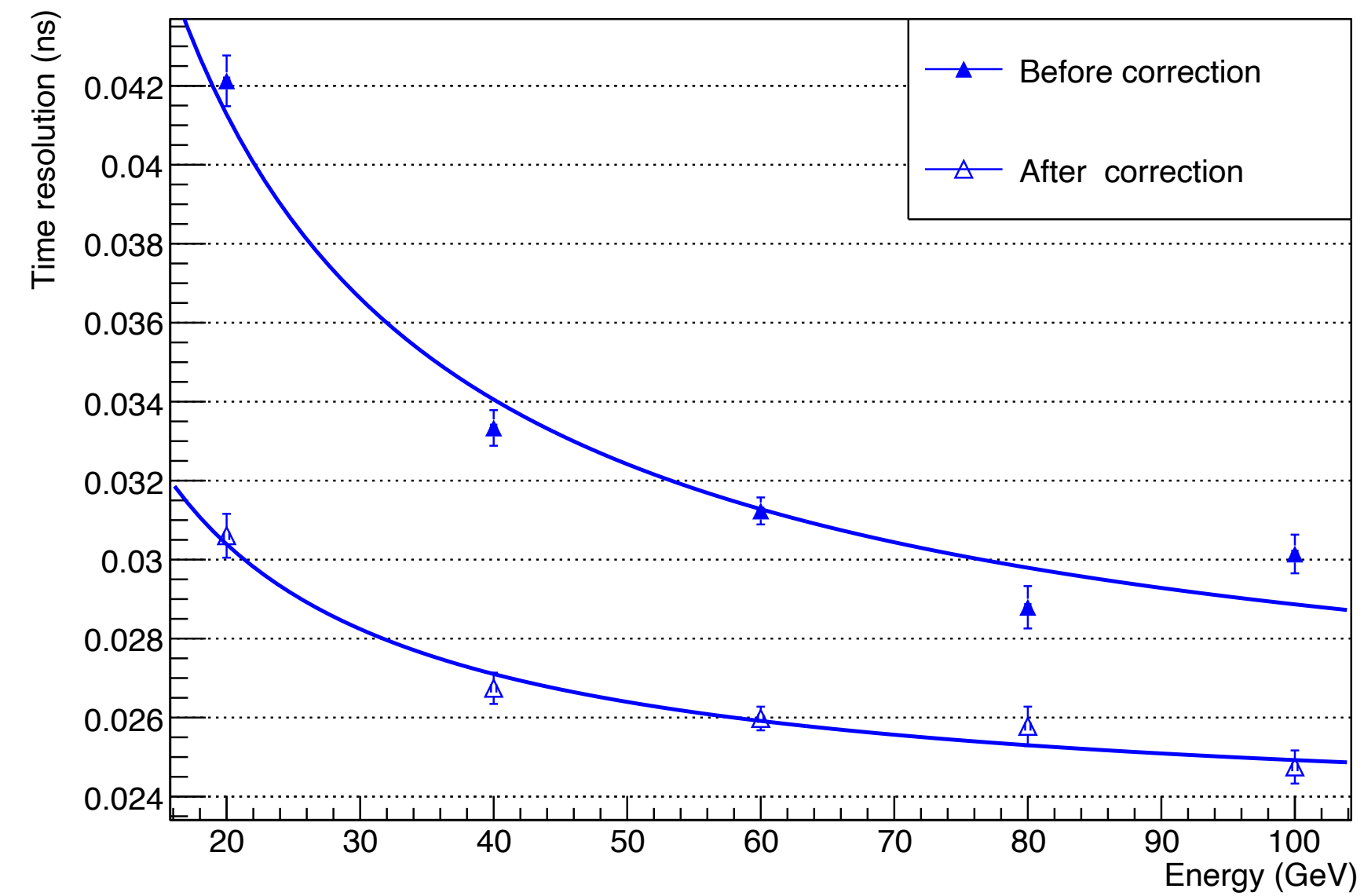
R7600U - SCSF-78 (blue) fibres



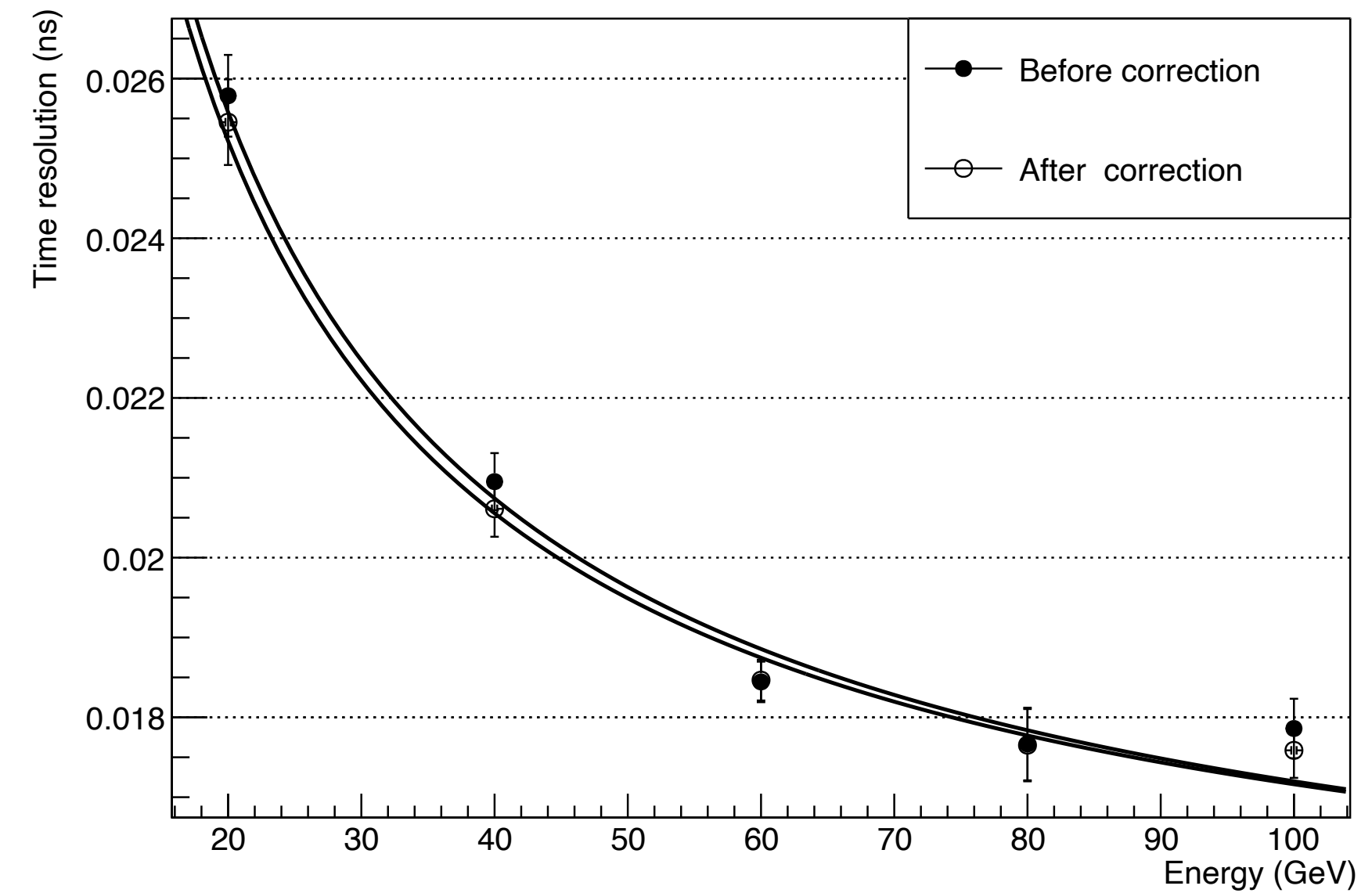
R9880U - SCSF (blue) fibres



R14755U - SCSF (blue) fibres

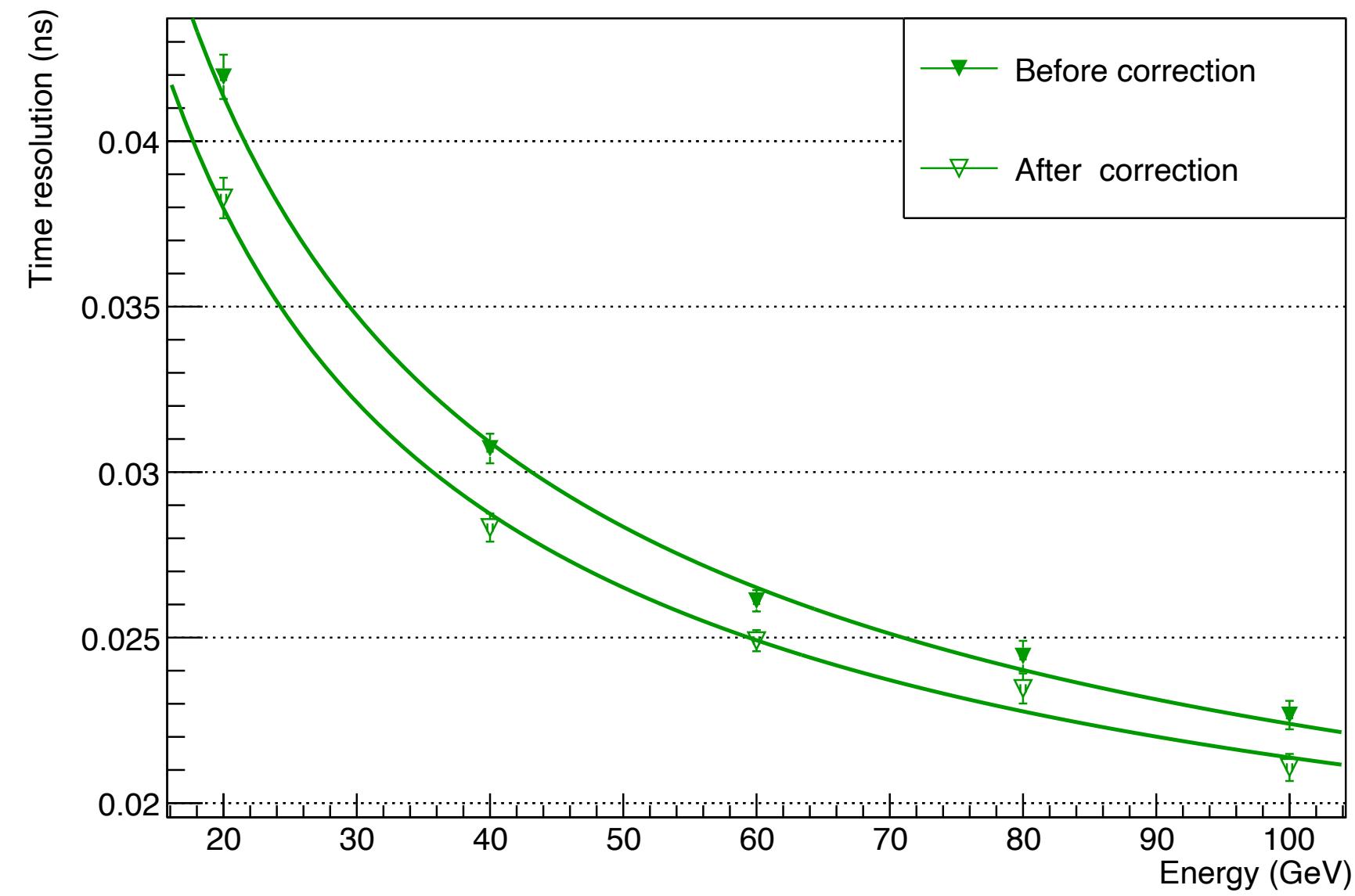


Tilecal - SCSF-78 (blue) fibres

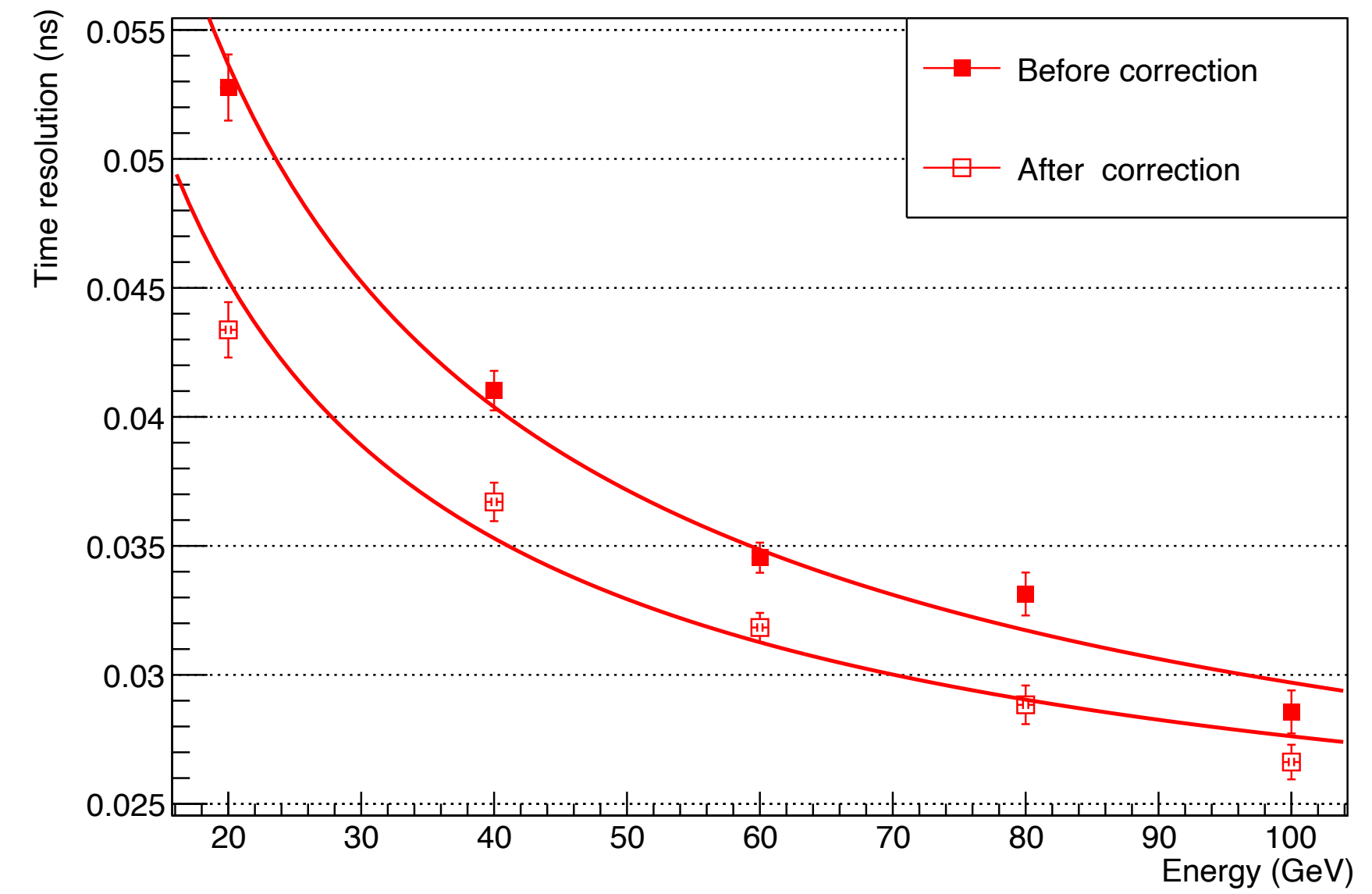


Corrected resolution - Small module with 3HF fibres

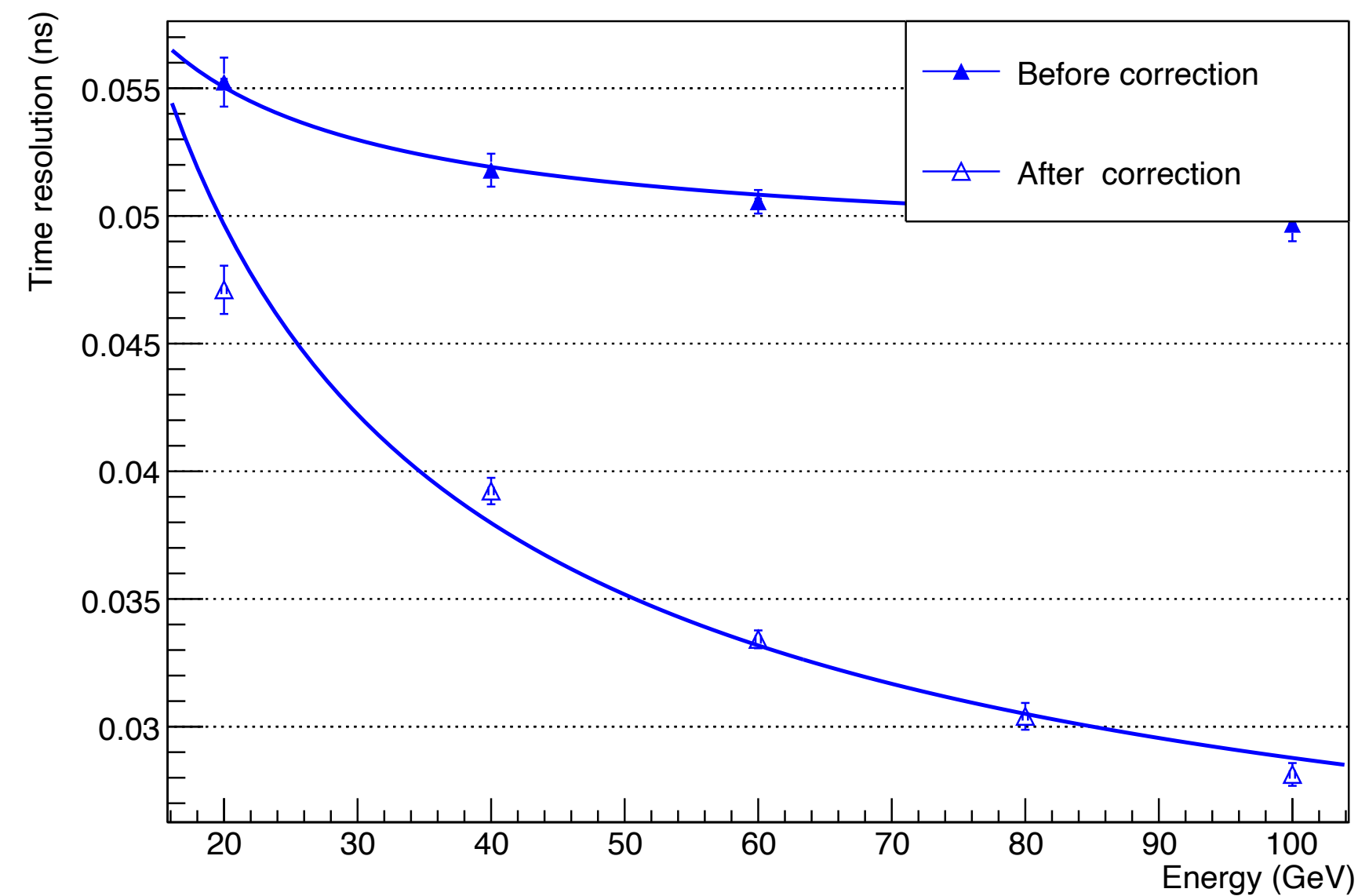
R7600U - 3HF (green) fibres



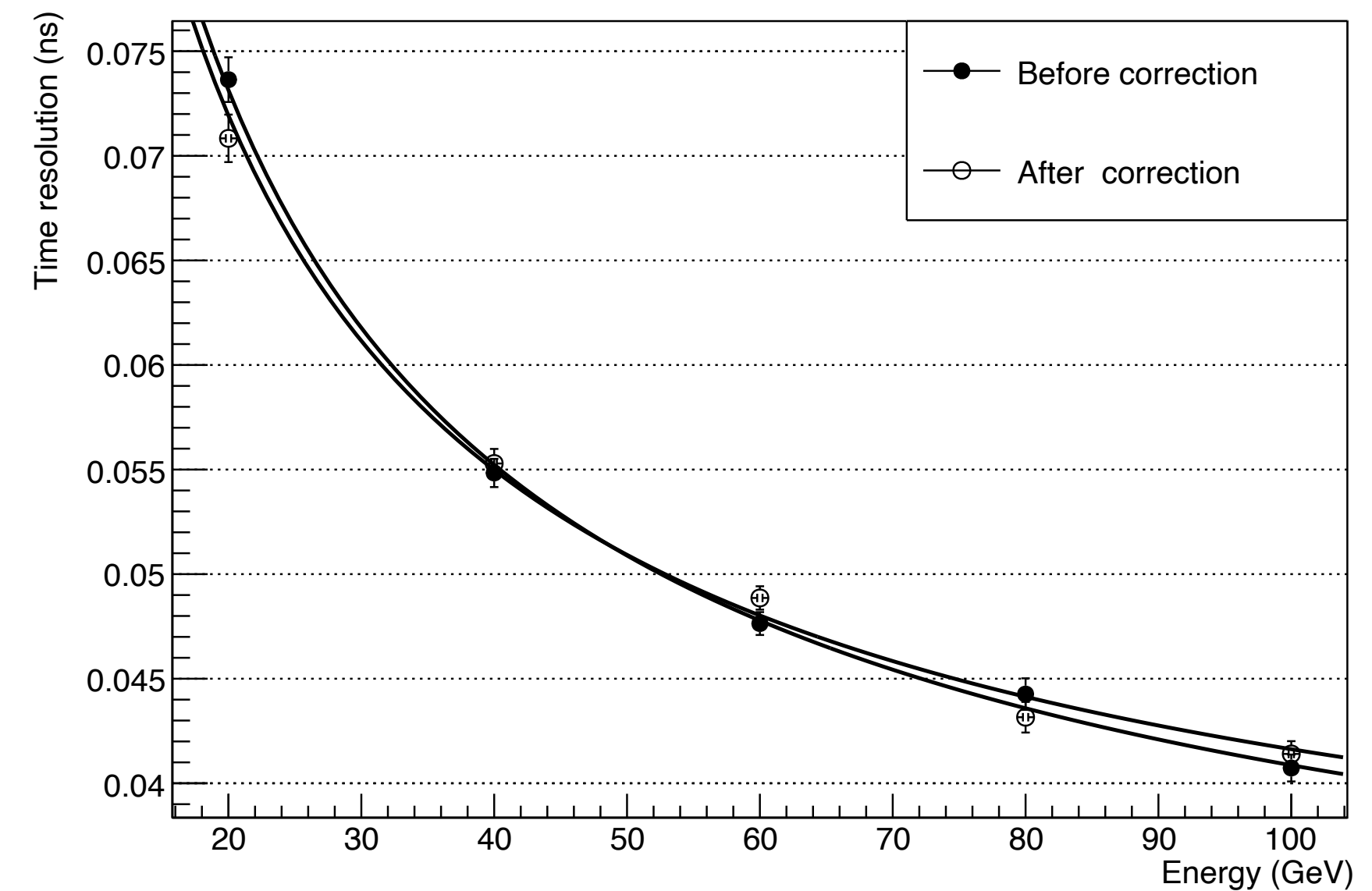
R9880U - 3HF (green) fibres



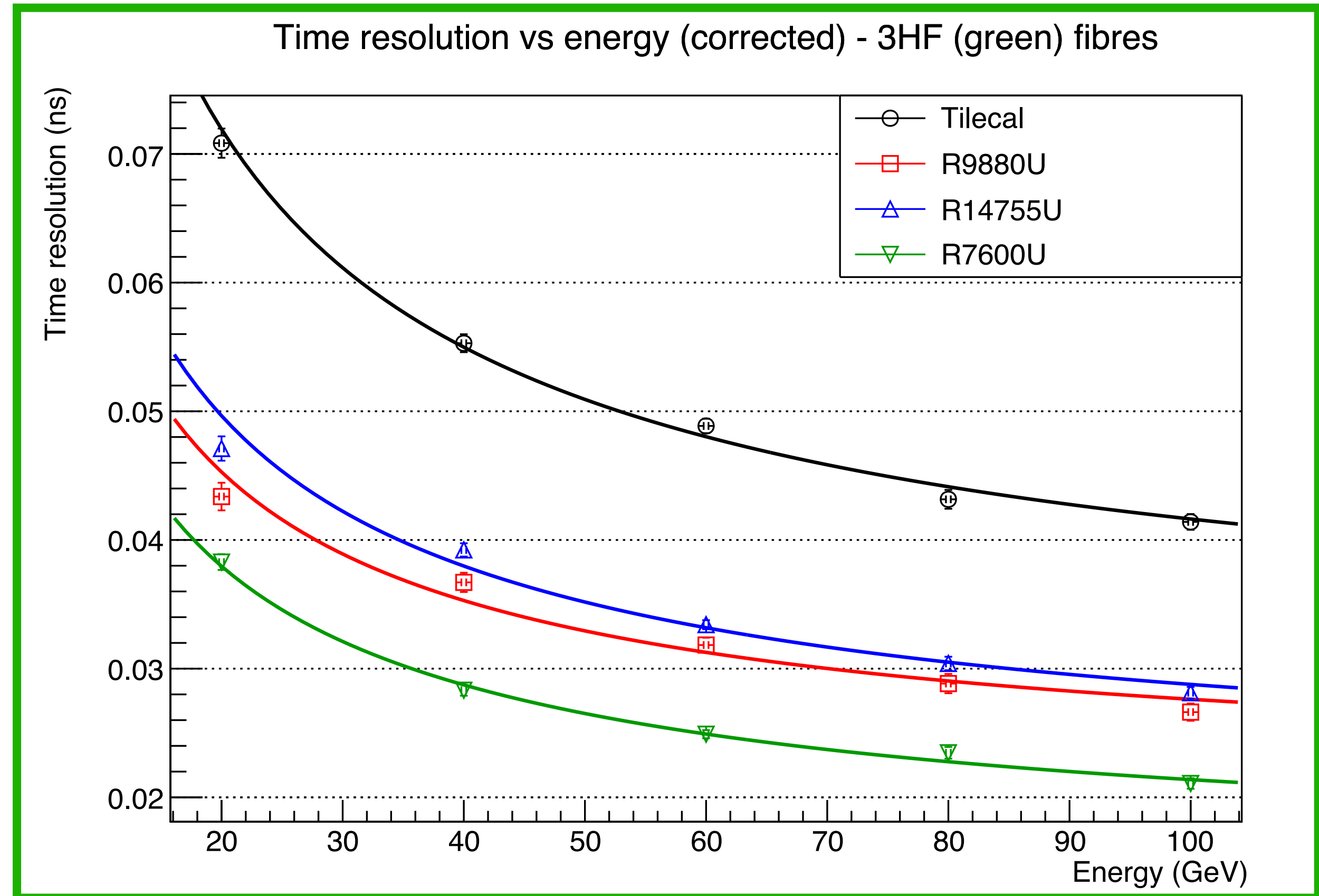
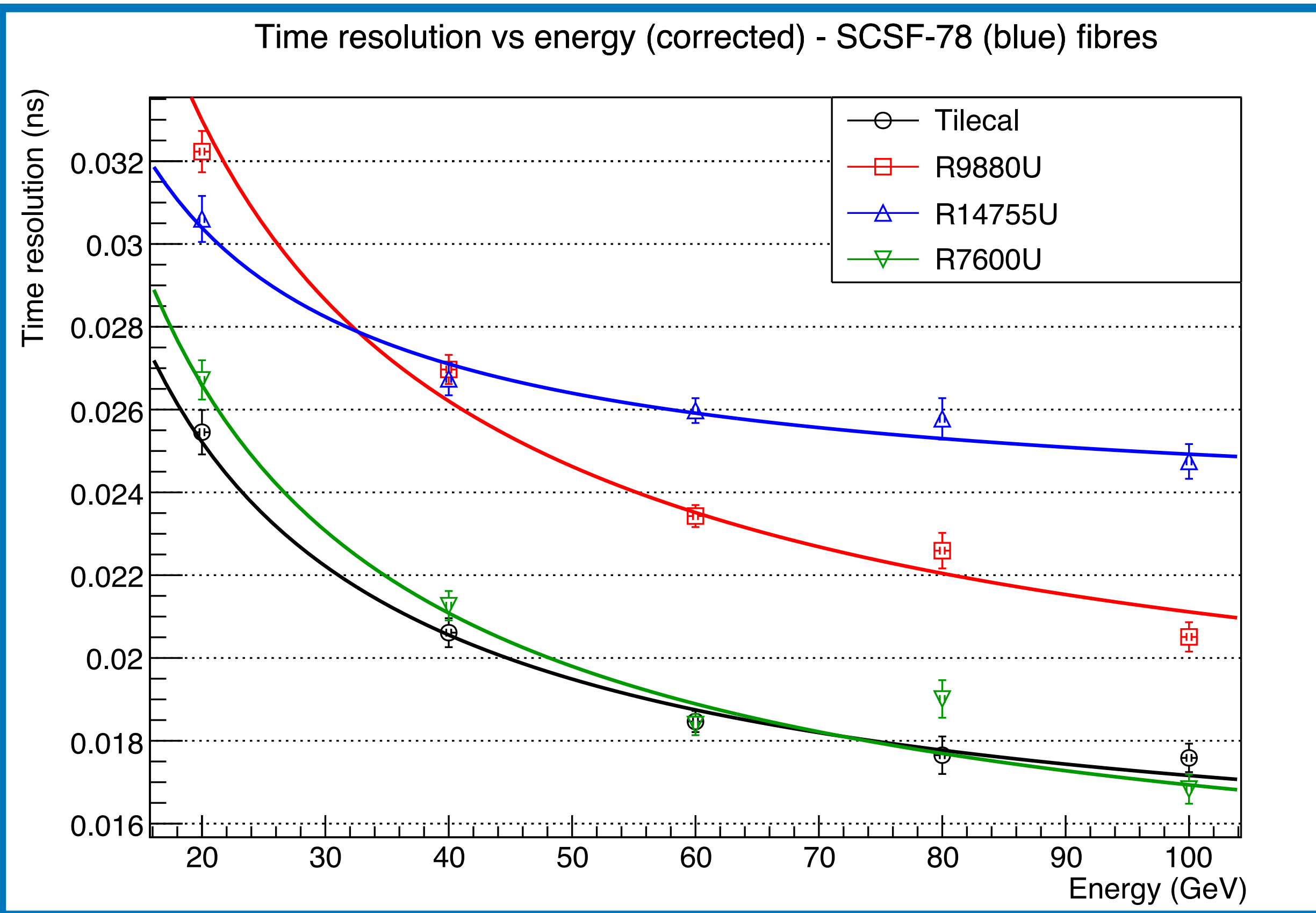
R14755U - 3HF (green) fibres



Tilecal - 3HF (green) fibres



Corrected resolution - Small module

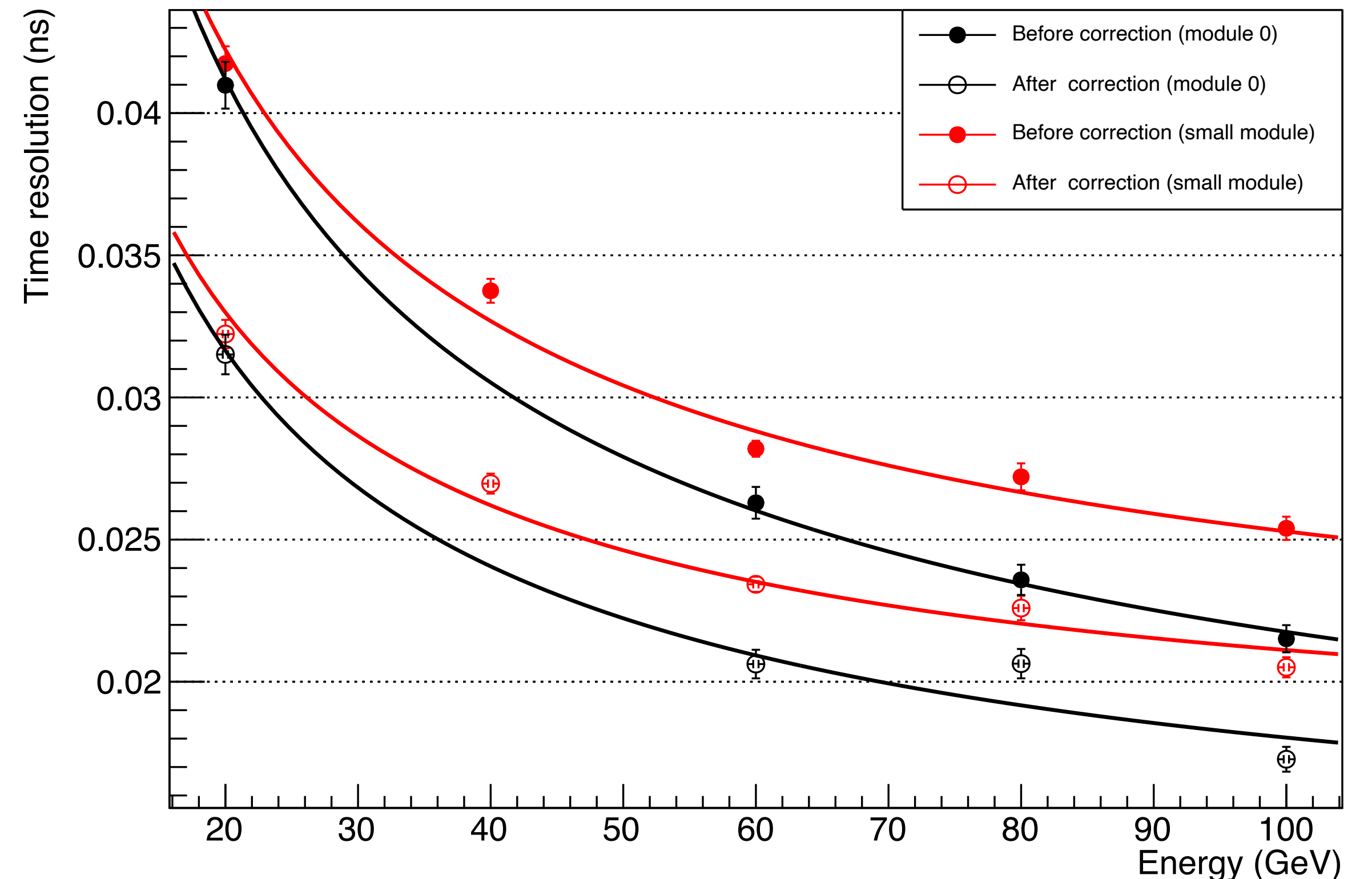


- As expected, **fast PMTs (R9880U and R14755U) undergo wider corrections**
- Still not enough for the fast PMTs to do better than **R7600U**
- The best threshold is always ~ 10% or 90%

“Module 0” vs “Small module”

- Both equipped with **SCSF-78** fibres and **R9880U** PMT
- **Module 0 performs better**
 - Down to 17 ps at 100 GeV
- Differences probably caused by **different light guide length**
 - 10 cm (module 0) vs 5 cm (small module)
 - Different photons loss

W-Poly small module vs module 0



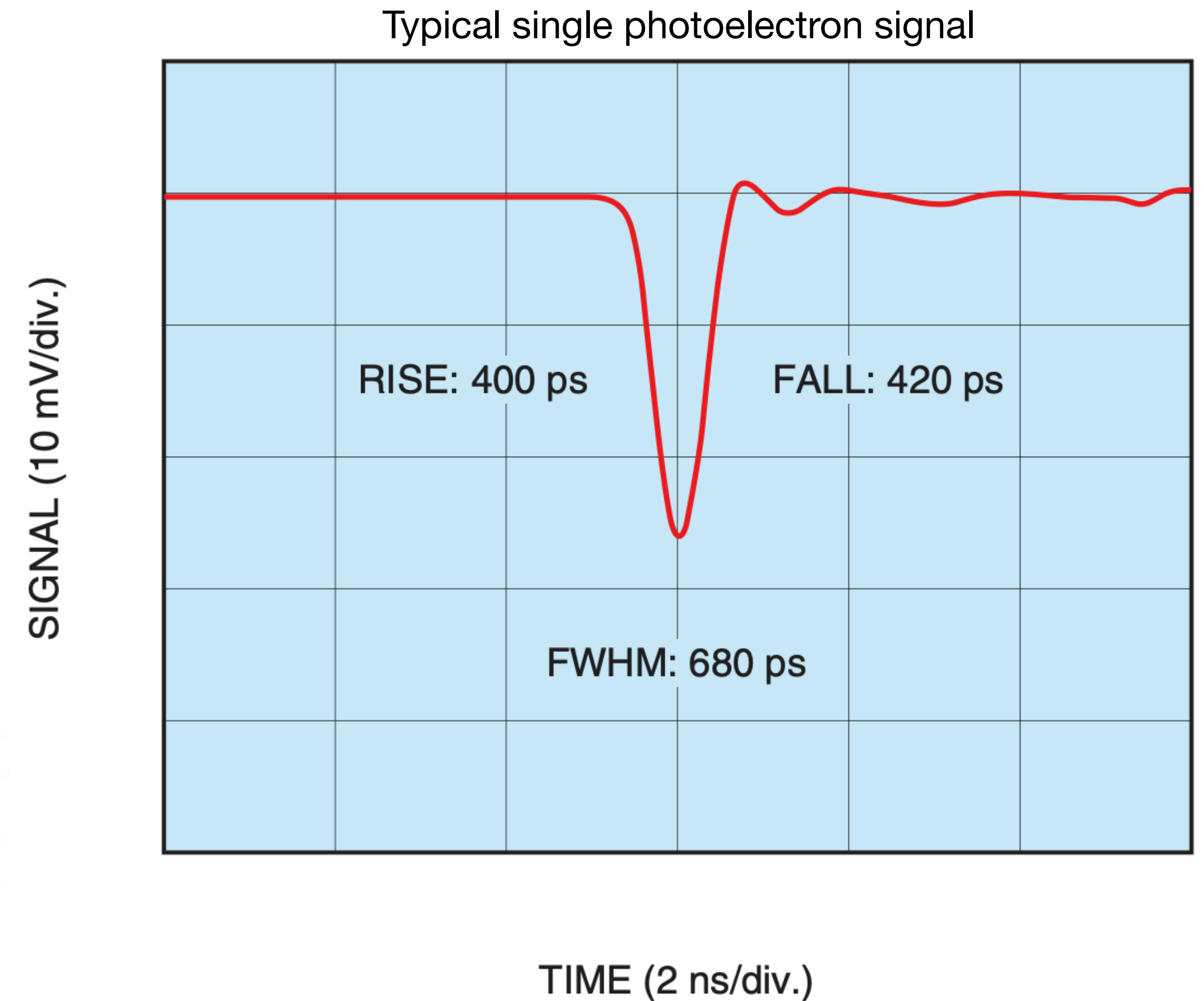
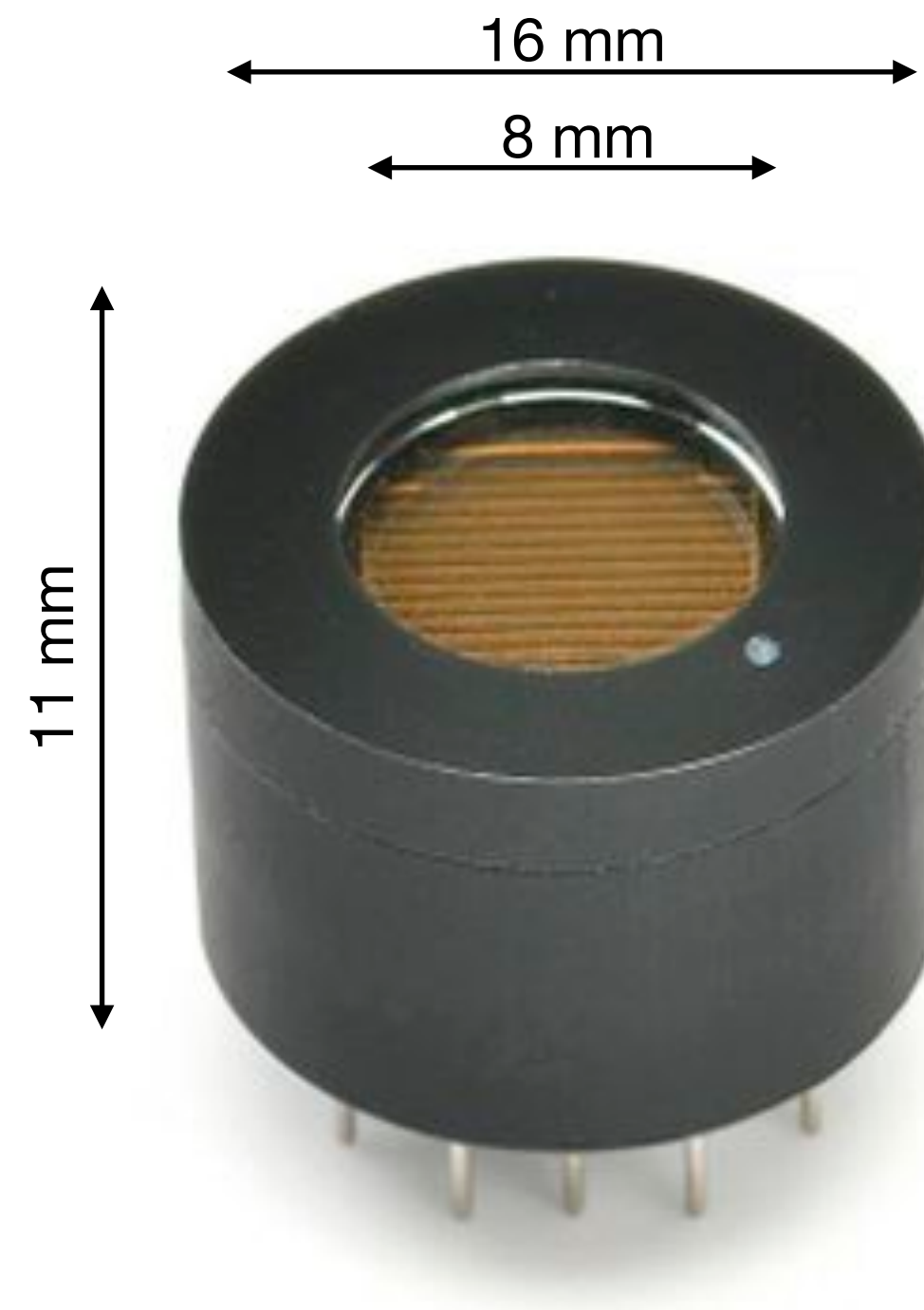
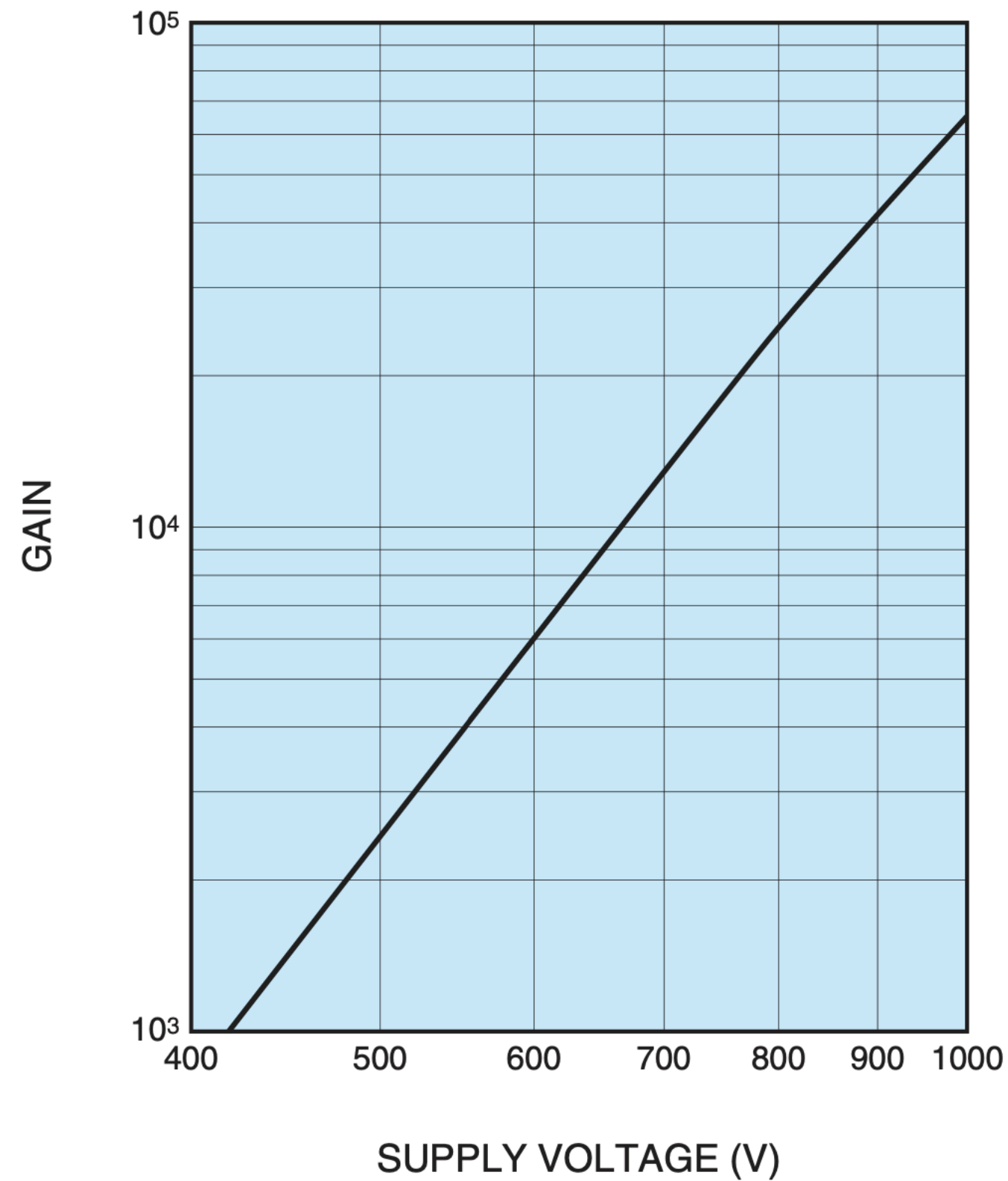
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The Hamamatsu R14755U-100 PMT

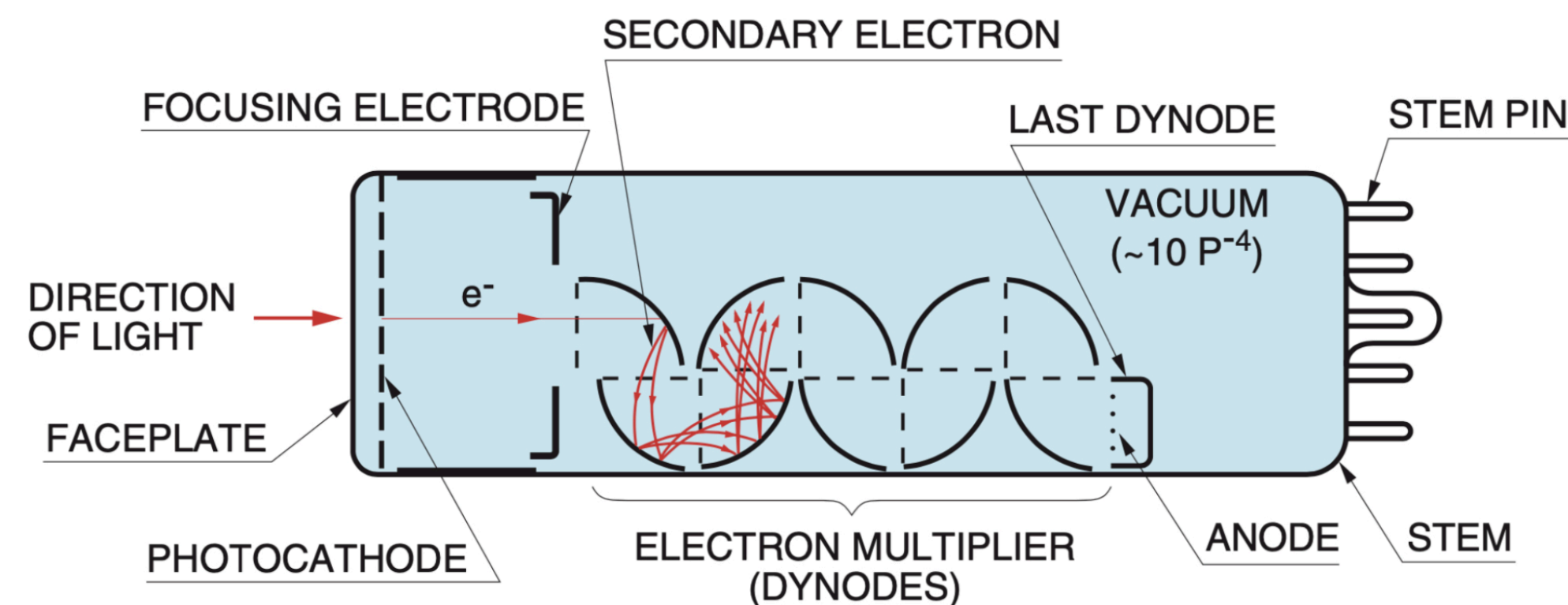
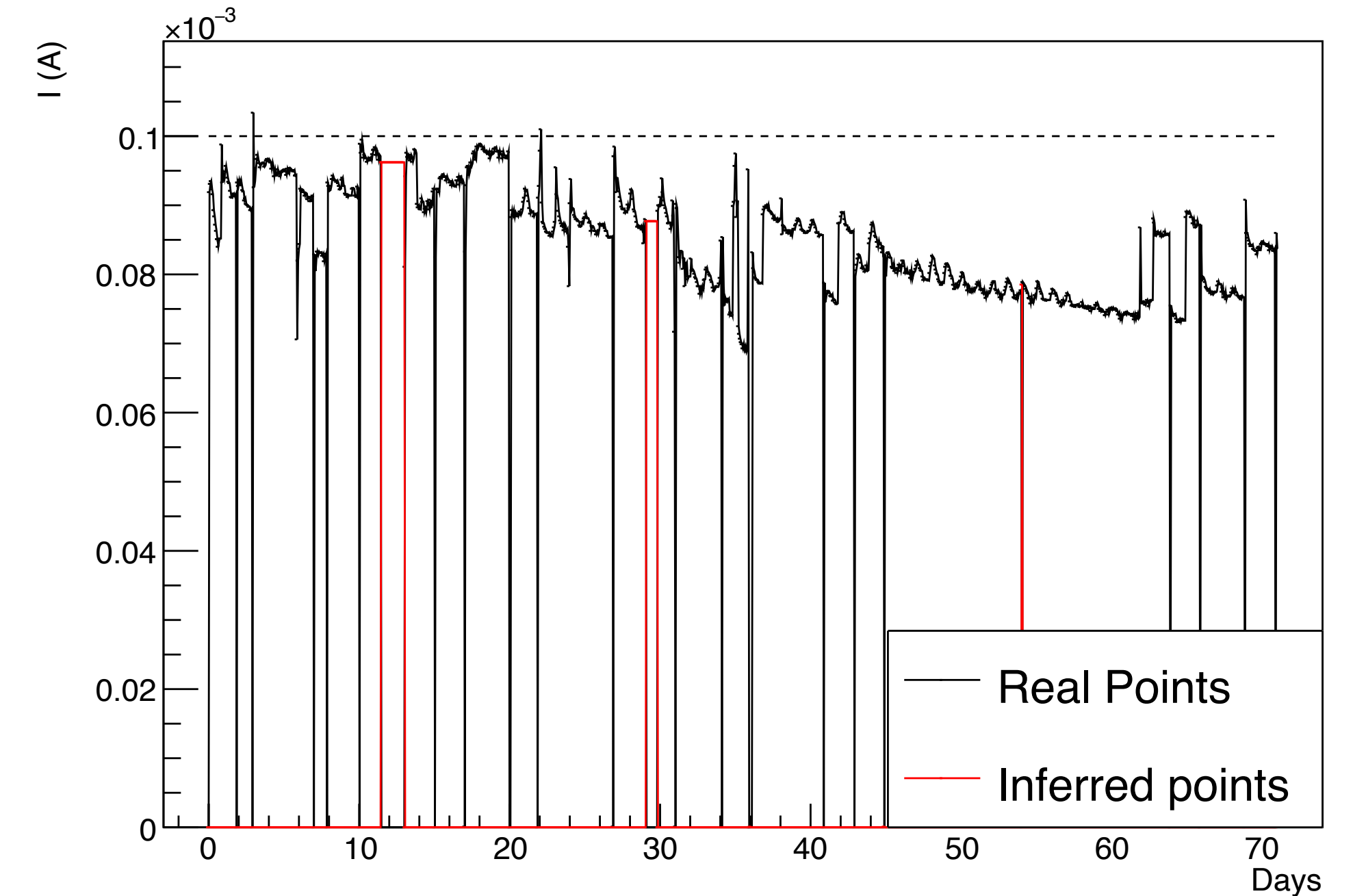
- Candidate PMT for the future LHCb ECAL thanks to its fast response
- Low gain, 6 dynode stages

Figures from the Hamamatsu data sheet
(https://www.hamamatsu.com/content/dam/hamamatsu-photonics/sites/documents/99_SALES_LIBRARY/etd/R14755U-100_TPMH1380E.pdf)



Ageing campaign & experimental setup

- Ageing affects the PMT performances:
 - Wear of the dynode system
 - Increase of the dark current
 - Decrease of the photocathode efficiency
- **Goal:** study the variation of the PMT gain up to 500 C of integrated charge
- Order of 10'000 C expected for new ECAL readout PMTs



$$\text{Gain} \doteq \frac{I_{anode}}{I_{cathode}}$$

Experimental setup & method

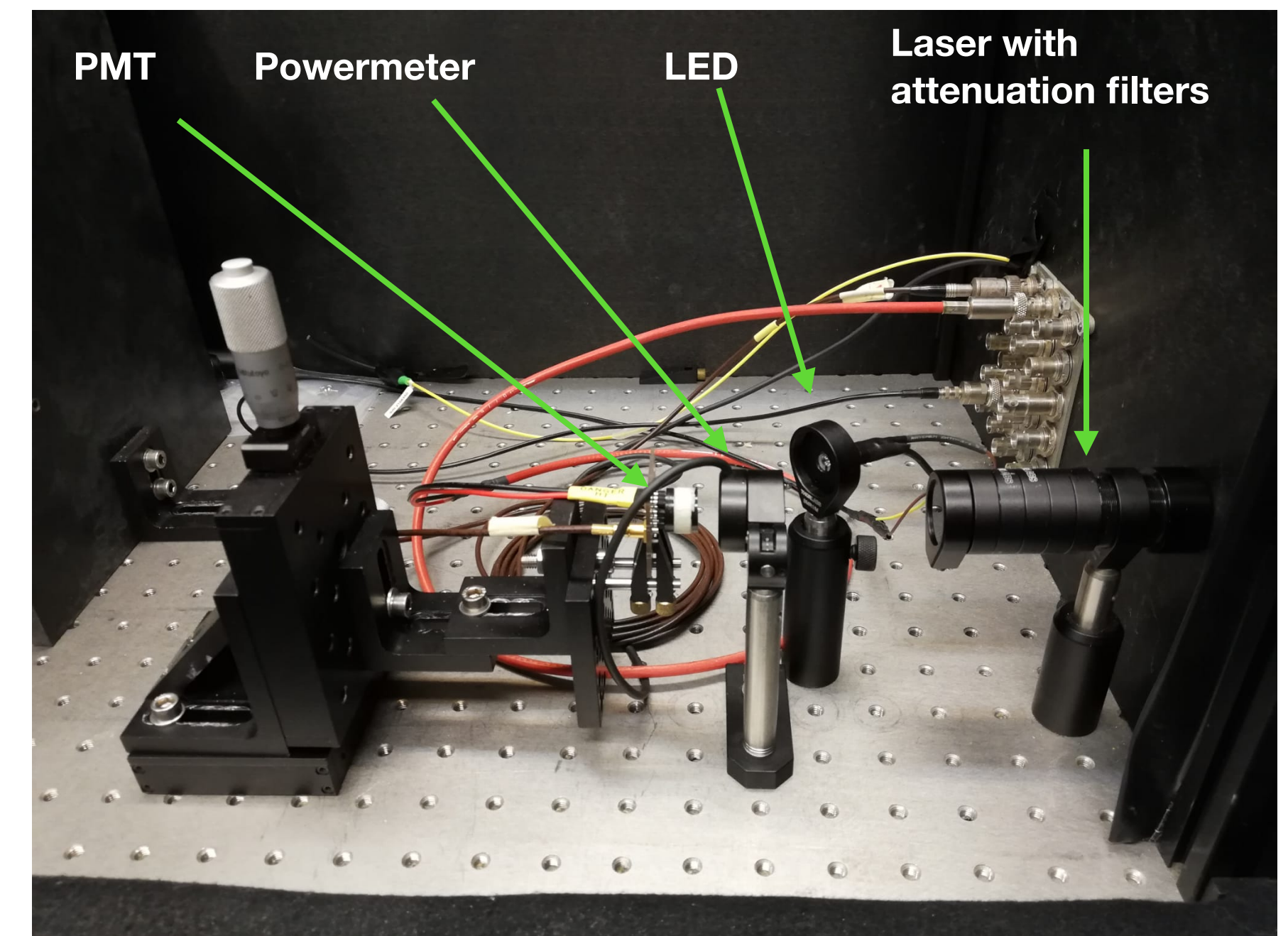
- **Ageing** performed with a white LED light continuously impinging the PMT
- **Gain measurements:** using a laser light ($\lambda = 405 \text{ nm}$), temporarily switching off the LED
- Measure the power (P) of the laser
 - ➔ get the number of photons per second impinging the PMT window:

$$P = n_{\gamma} \cdot \frac{hc}{\lambda}$$

- By measuring the anodic and dark current, we get the gain:

$$G = \frac{n_e}{Q_E \cdot n_{\gamma}} = \frac{hc}{Q_E \cdot P \cdot \lambda} \cdot \frac{(I_{anode} - I_{dark})}{q_{e^-}}$$

- Procedure validated by measuring the gain of another PMT also with another method (compatibility within 12%)



Gain and HV

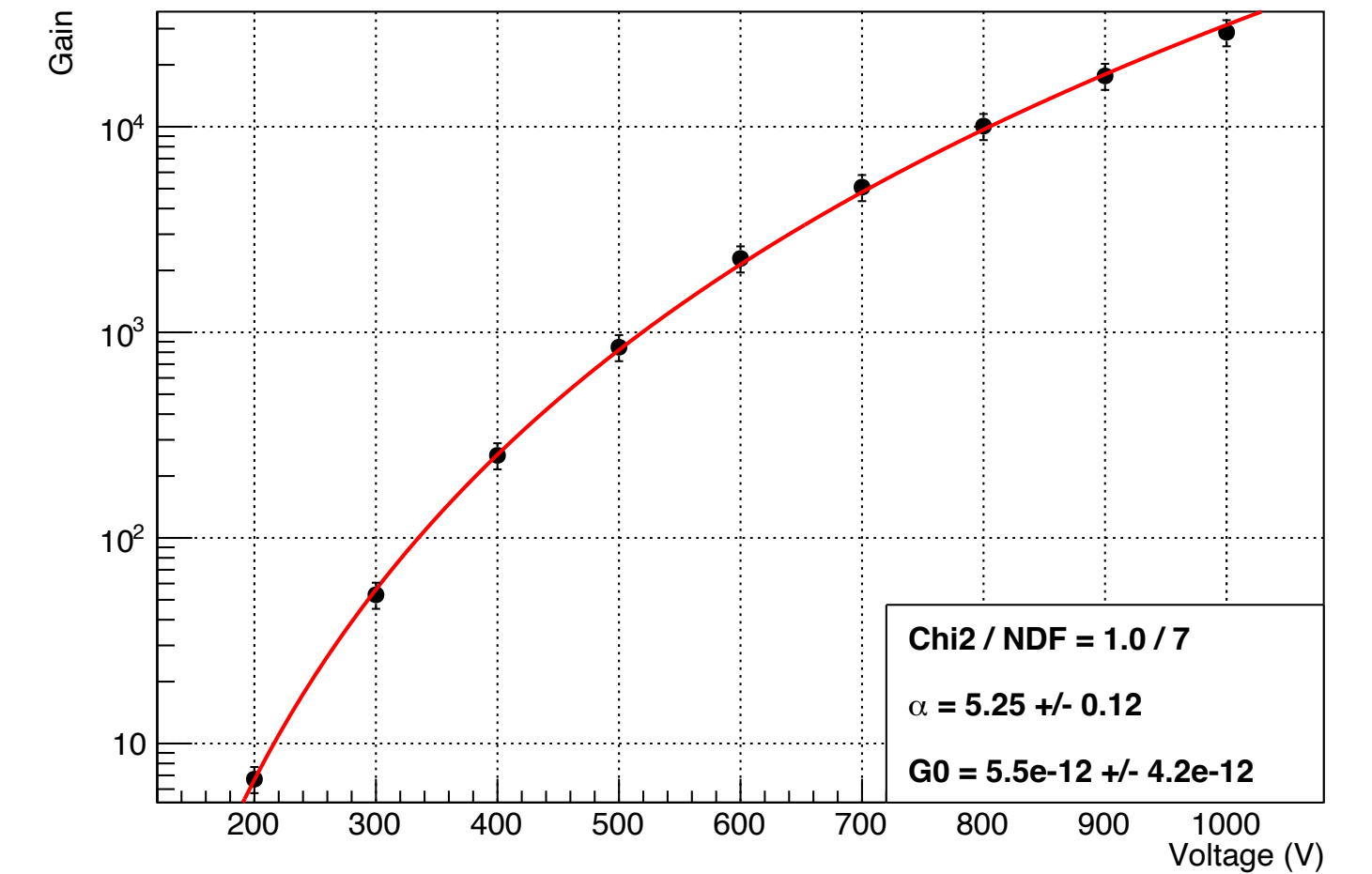
- The PMT gain grows with the applied voltage following:

$$G = G_0 \cdot V^\alpha$$

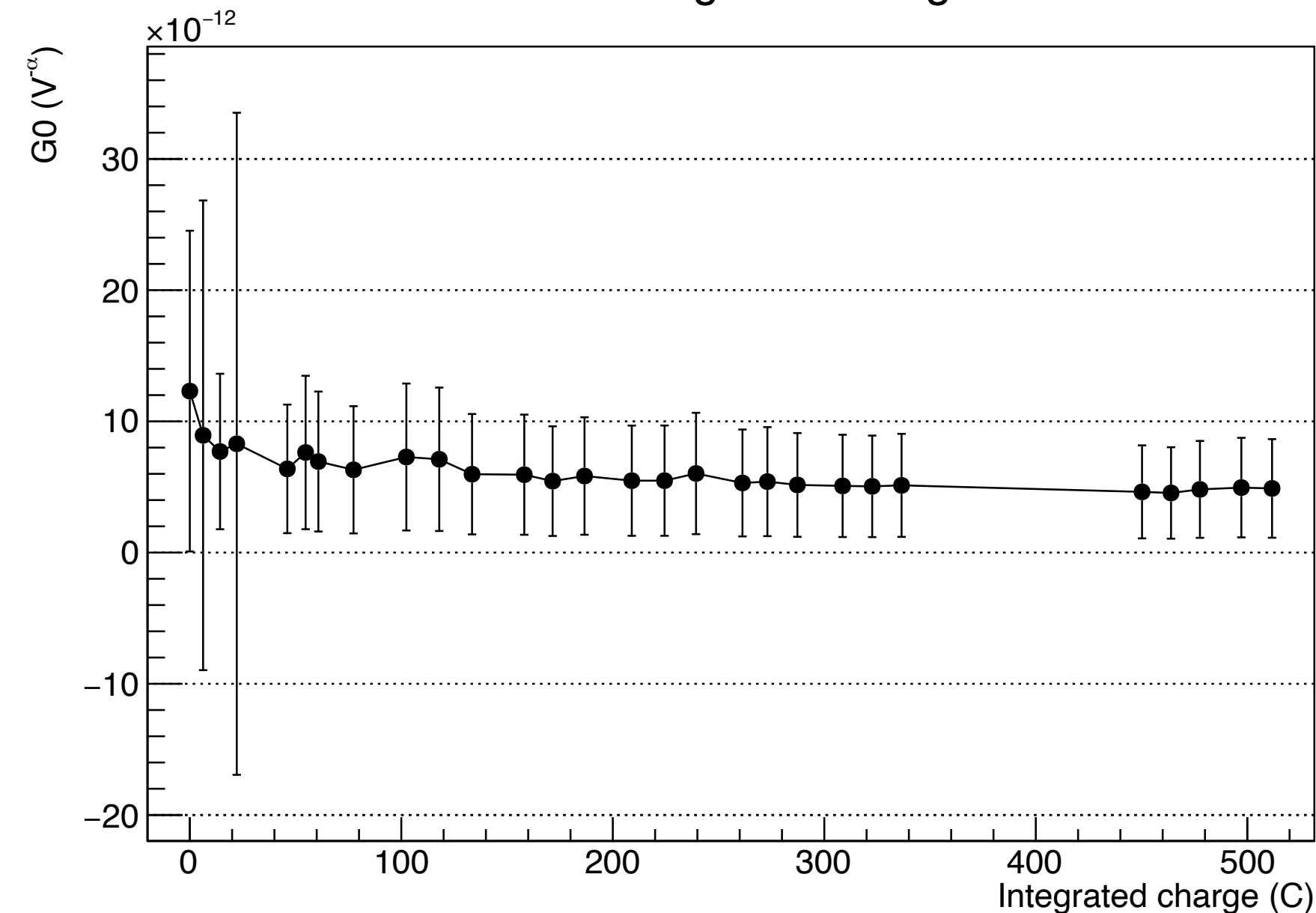
- Each time, the gain has been measured from 200 V to 1000 V (with 100 V steps)
- Power-law fit** to the data to get α and G_0

Example of a power-law fit

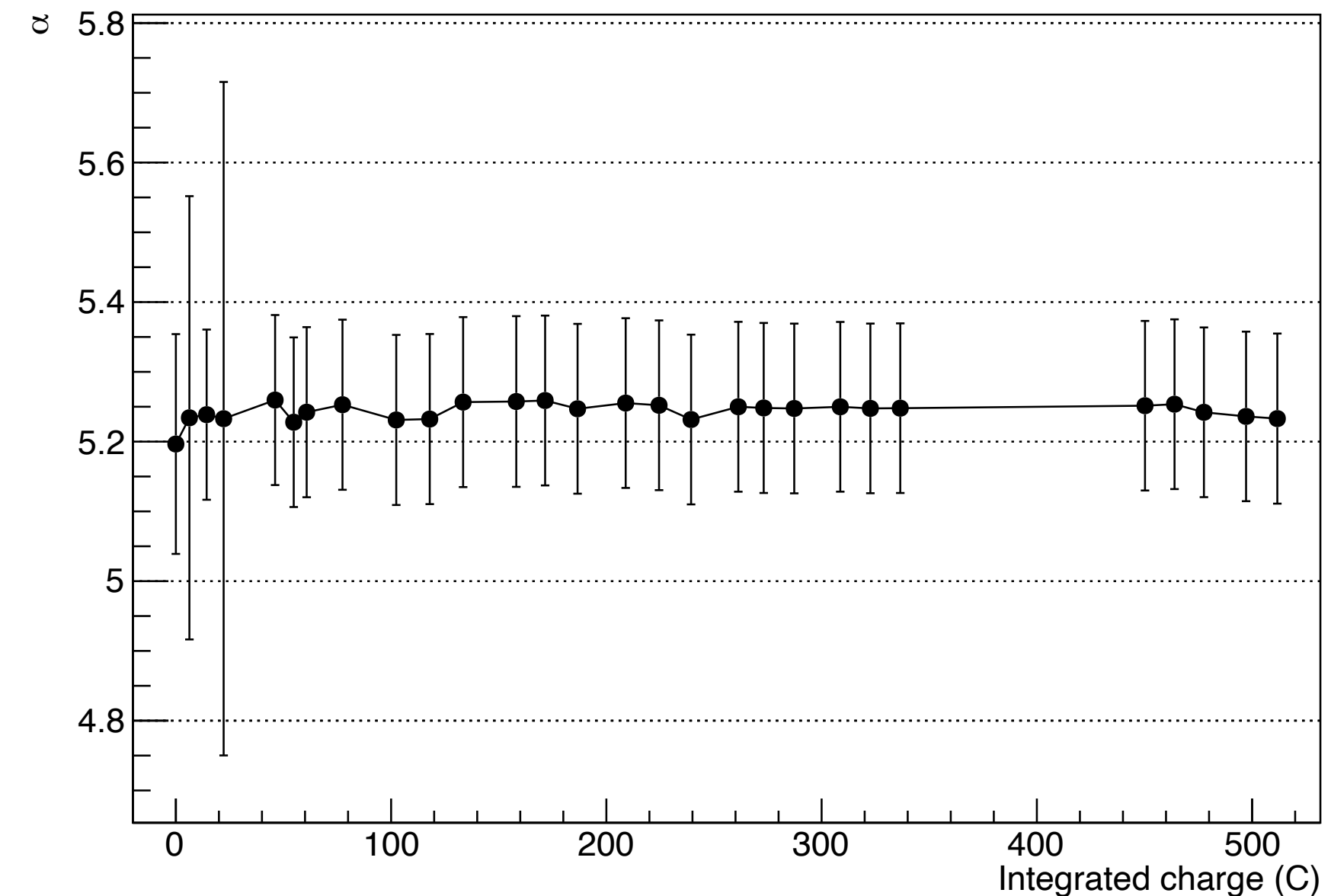
224 C



G0 vs integrated charge



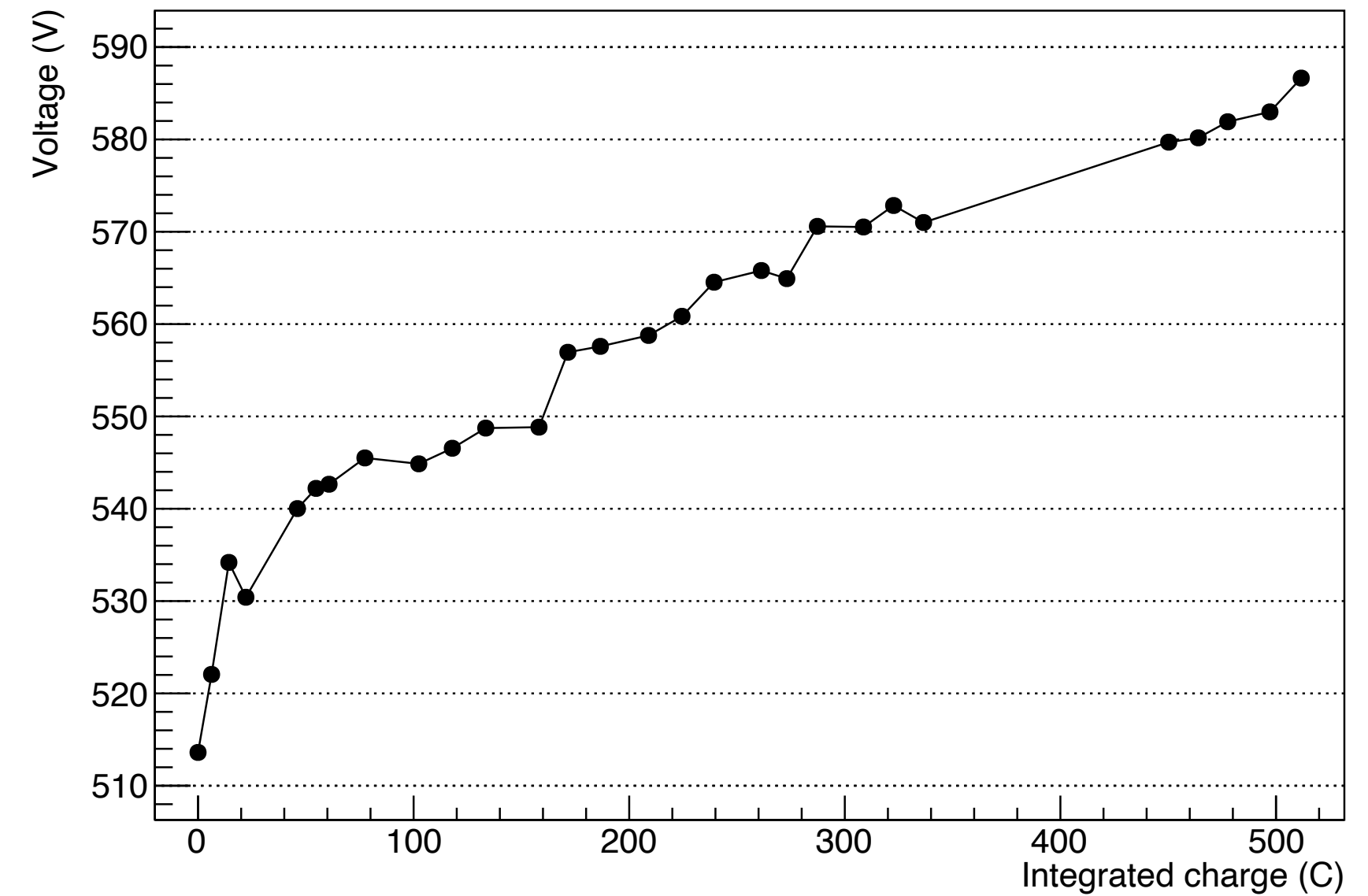
alpha vs integrated charge



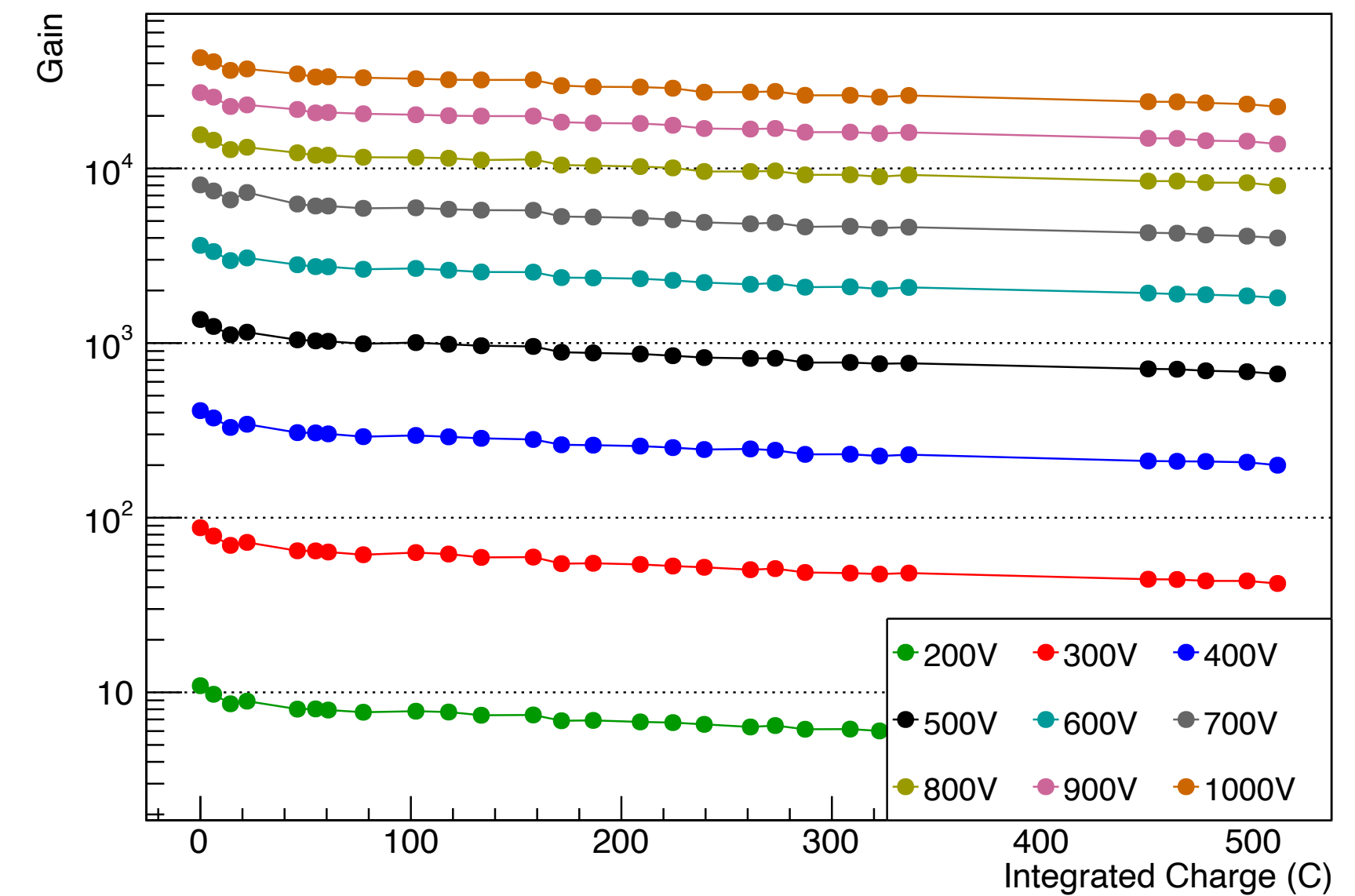
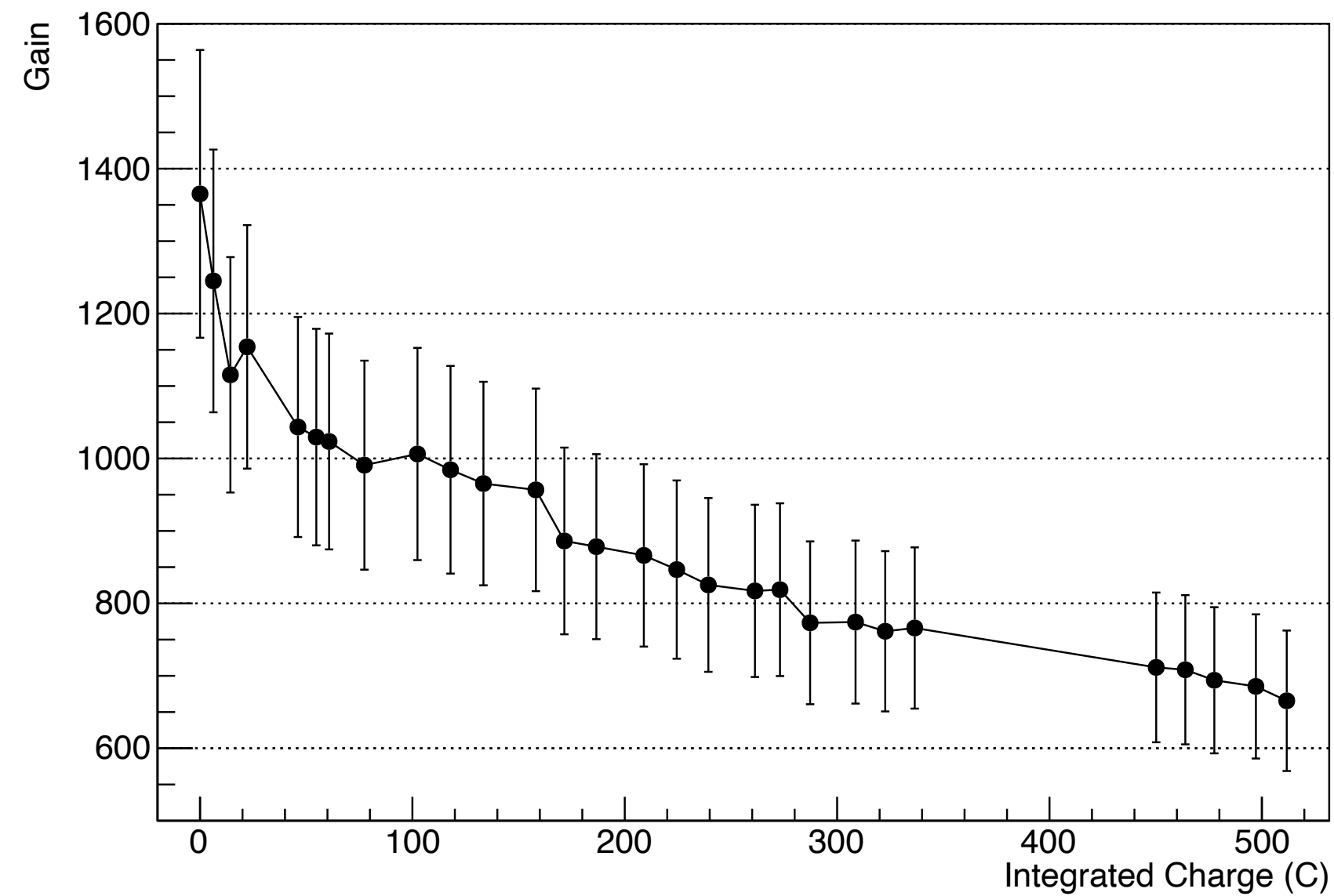
Results

- Quick drop in the first 50 C
- Then, stable and slow decrease up to 500 C
- Overall decrease by a factor 2
- Same behaviour at all the applied voltages

Voltage to keep gain at 1500



500V



Outline

- Introduction: the LHCb ECAL Upgrade II
- Simulation studies of a Pb-Polystyrene module
- Analysis of testbeam data from the CERN SPS
- Ageing campaign of the PMT Hamamatsu R14755U-100
- **Conclusions**

Conclusions - Time resolution

Simulation studies

- The signal formation inside the single-side readout SpaCal modules has been studied by means of GEANT4 simulations
- The time resolution is worsened by the longitudinal fluctuations of the showers affecting the pulses' shape
 - ➔ **The CFD algorithm can't take this into account**
- A procedure aiming at removing the shower depth bias has been developed and applied to testbeam data, exploiting the signals' rise time

Conclusions - Time resolution

Simulation studies

- The signal formation inside the single-side readout SpaCal modules has been studied by means of GEANT4 simulations
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SPS Testbeam data

- A comparison among 4 PMT models and 2 types of fibres has been performed for two SpaCal W-Polystyrene modules
- Resolutions below 20 ps obtained at high energies
 - ➔ **Good timing capabilities of the SpaCal even in single-side readout mode**

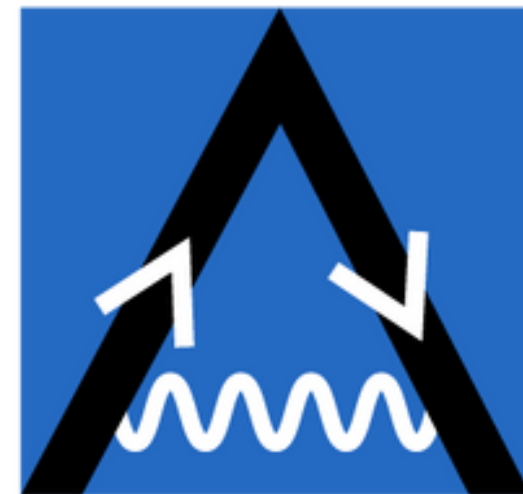
Open questions & Outlook

- What is the physical origin of the sampling and constant terms of the time resolution?
- Why do fast PMTs still do worse than R7600U after the correction procedure?
 - Corrections are never ideal
 - CFD + polynomial fit may not be the best approach
- **The CFD algorithm may not be the best one**
 - ➔ **Is there a better way to define the time stamps?**
 - A Machine Learning method will be tested in the near future

Conclusions - Ageing campaign

- An ageing campaign of the Hamamatsu R14755U-100 PMT has been carried out up to 500 C
- The gain has been measured at applied voltages between 200 V and 1000 V
 - **Stable decrease** (up to 500 C)
 - Same behaviour at all the applied voltages
- The PMT showed a very good (slow) aging behaviour
 - ➔ **Good candidate for the future LHCb ECAL in terms of ageing**

Thank you for your attention



International Master
Advanced Methods
in Particle Physics

Backup slides

To get an idea:



Hamamatsu R7600U-00-M4

FWHM ~ 2.1 ns

Tau ~ 0.6 ns



Hamamatsu R14755U-100

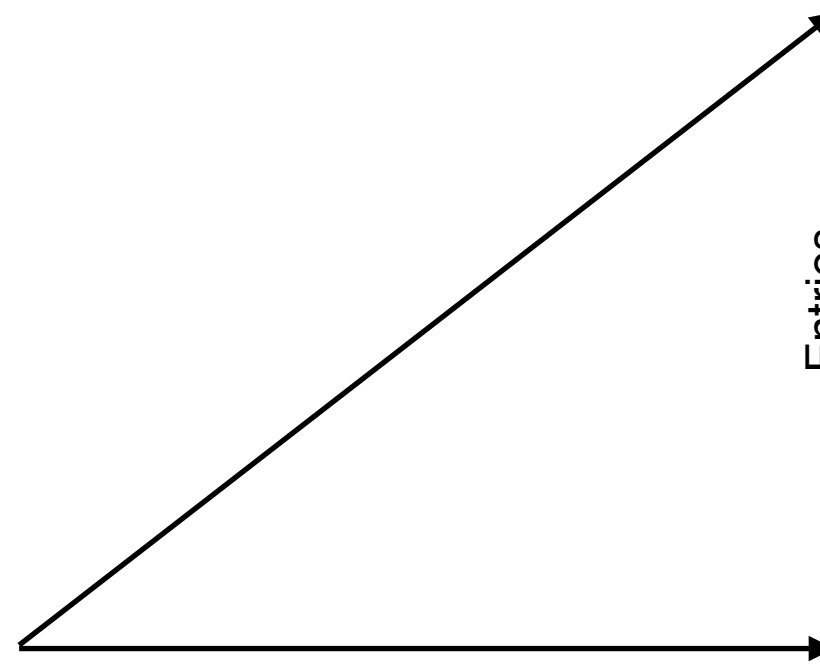
FWHM ~ 0.68 ns

Tau ~ 0.2 ns

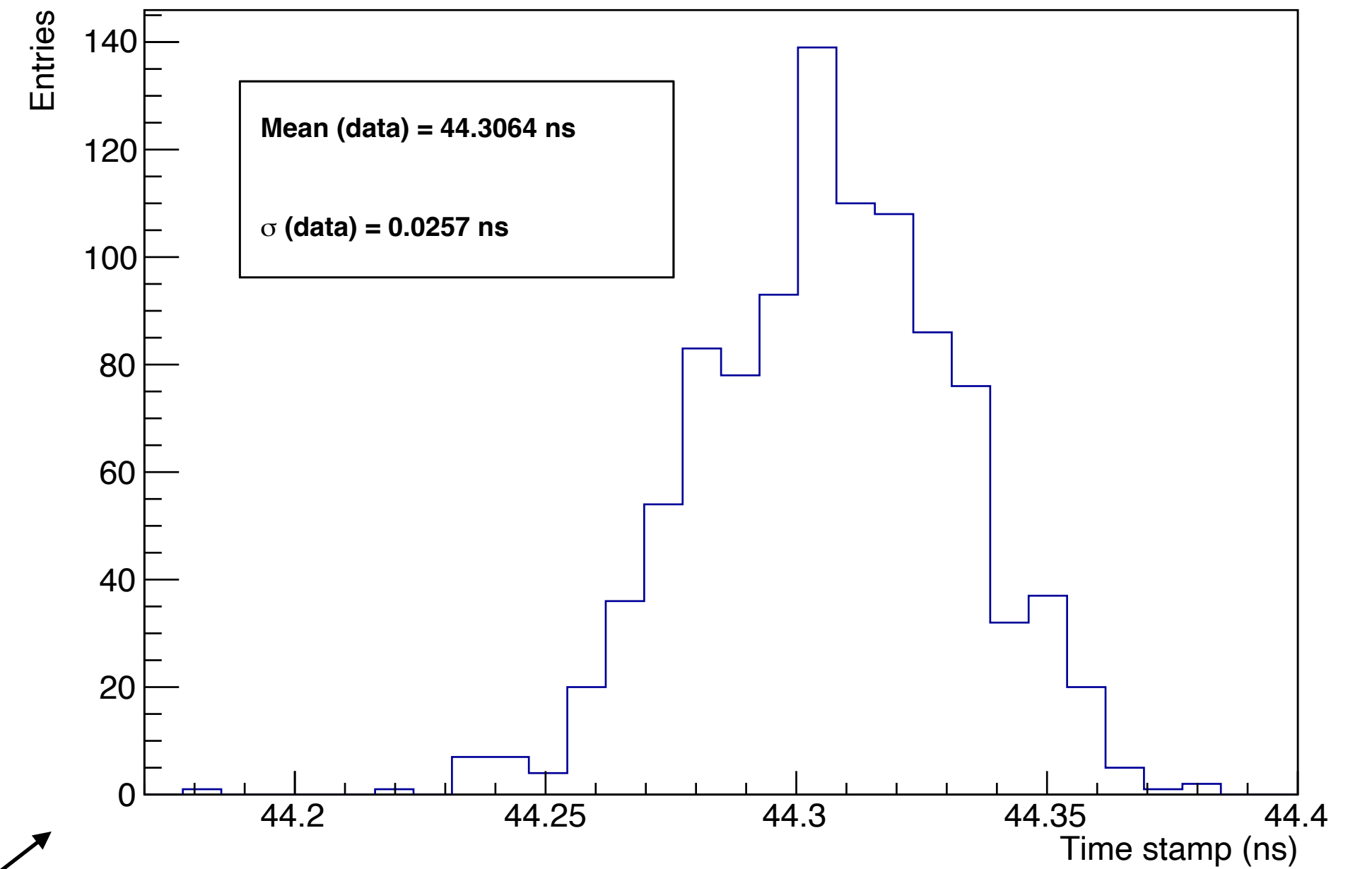
- Only the central PMT is considered
- **Time resolution = std. dev.** of the time stamps sample
- **Rise time:** time between the 10% and 90% of the pulse's amplitude
- In general the shapes are not gaussian
- Presence of outliers

Example here:

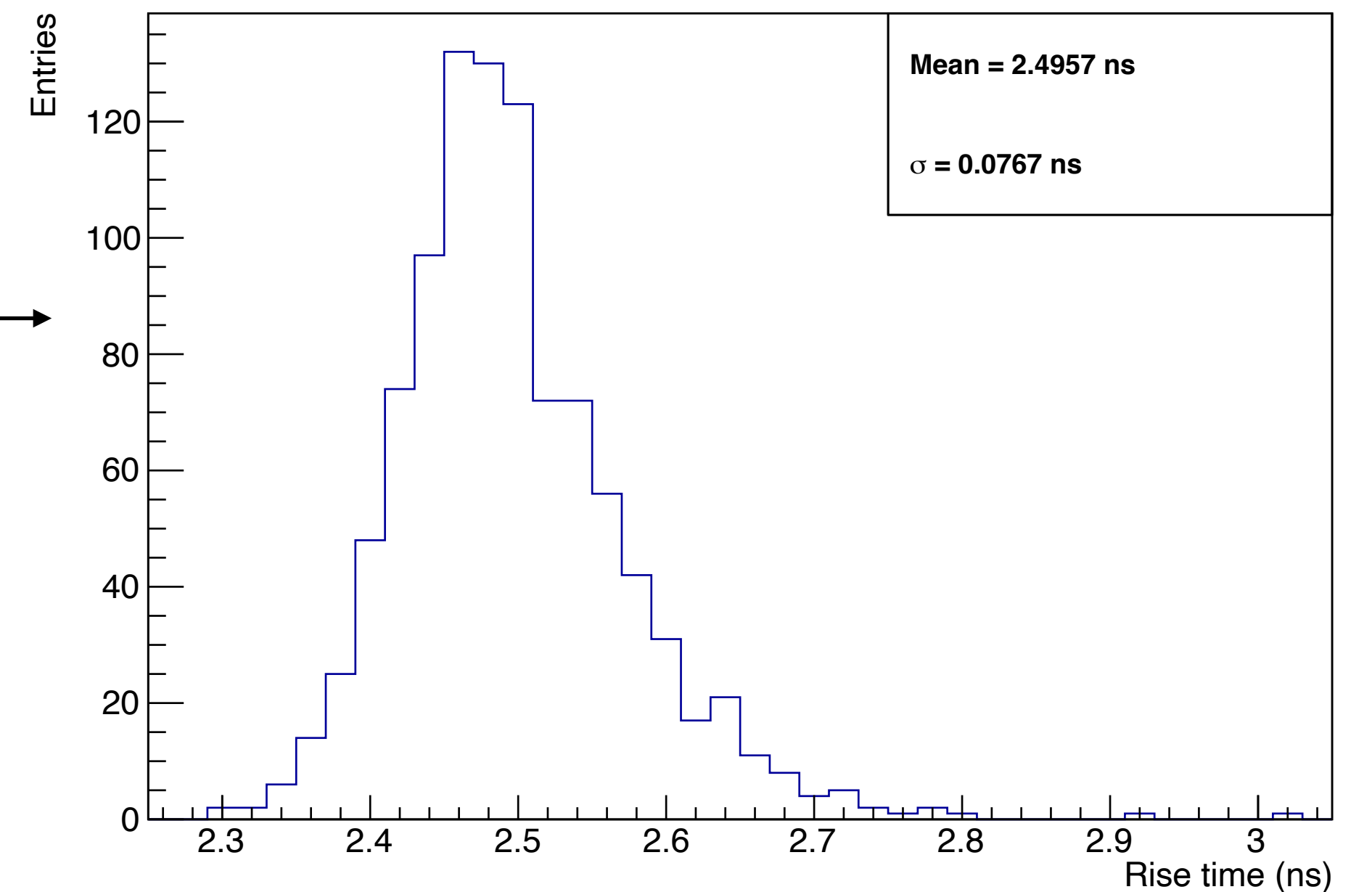
- $\tau = 1.0 \text{ ns}$ (FWHM = 3.4 ns)
- CFD fraction = 40%



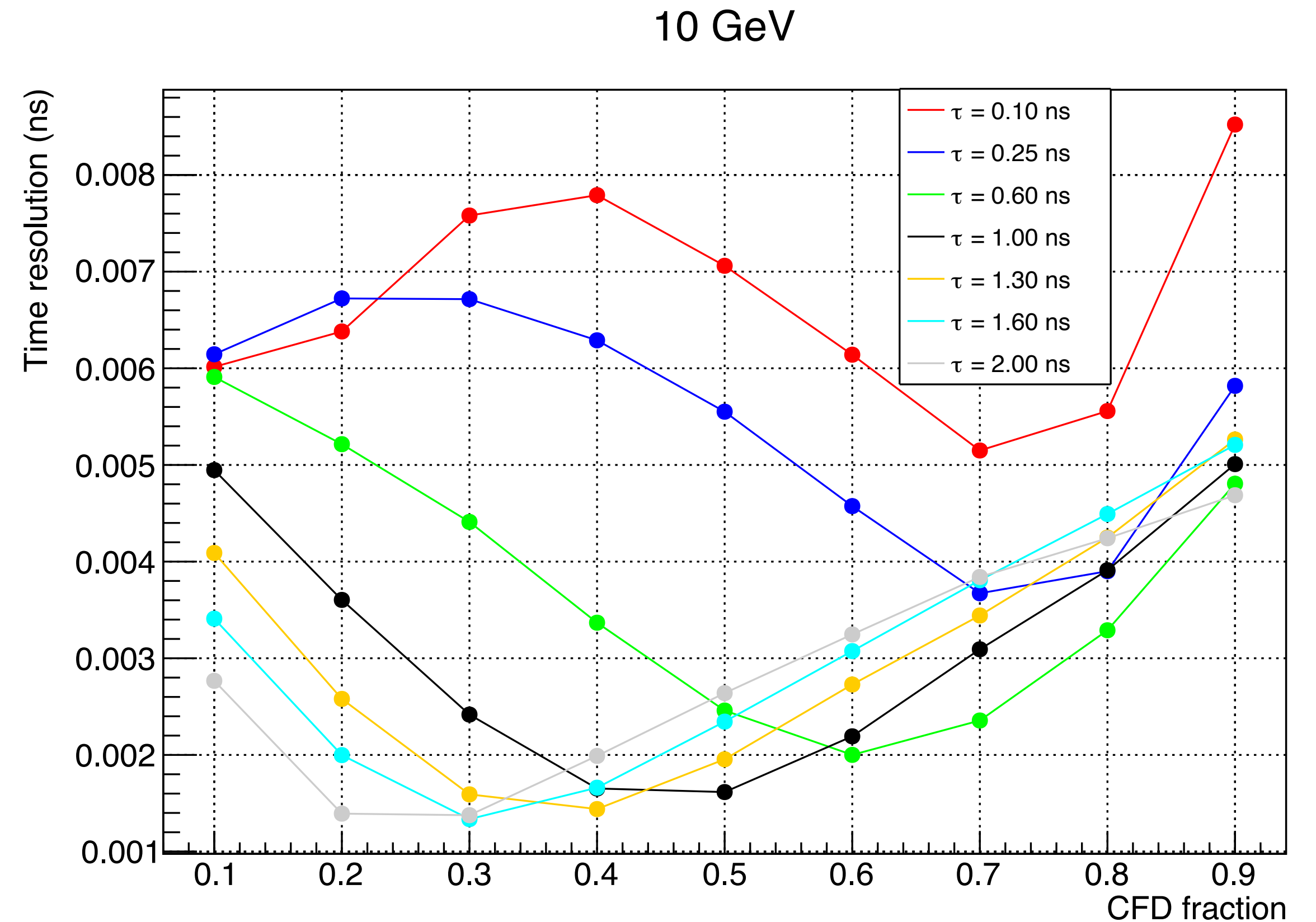
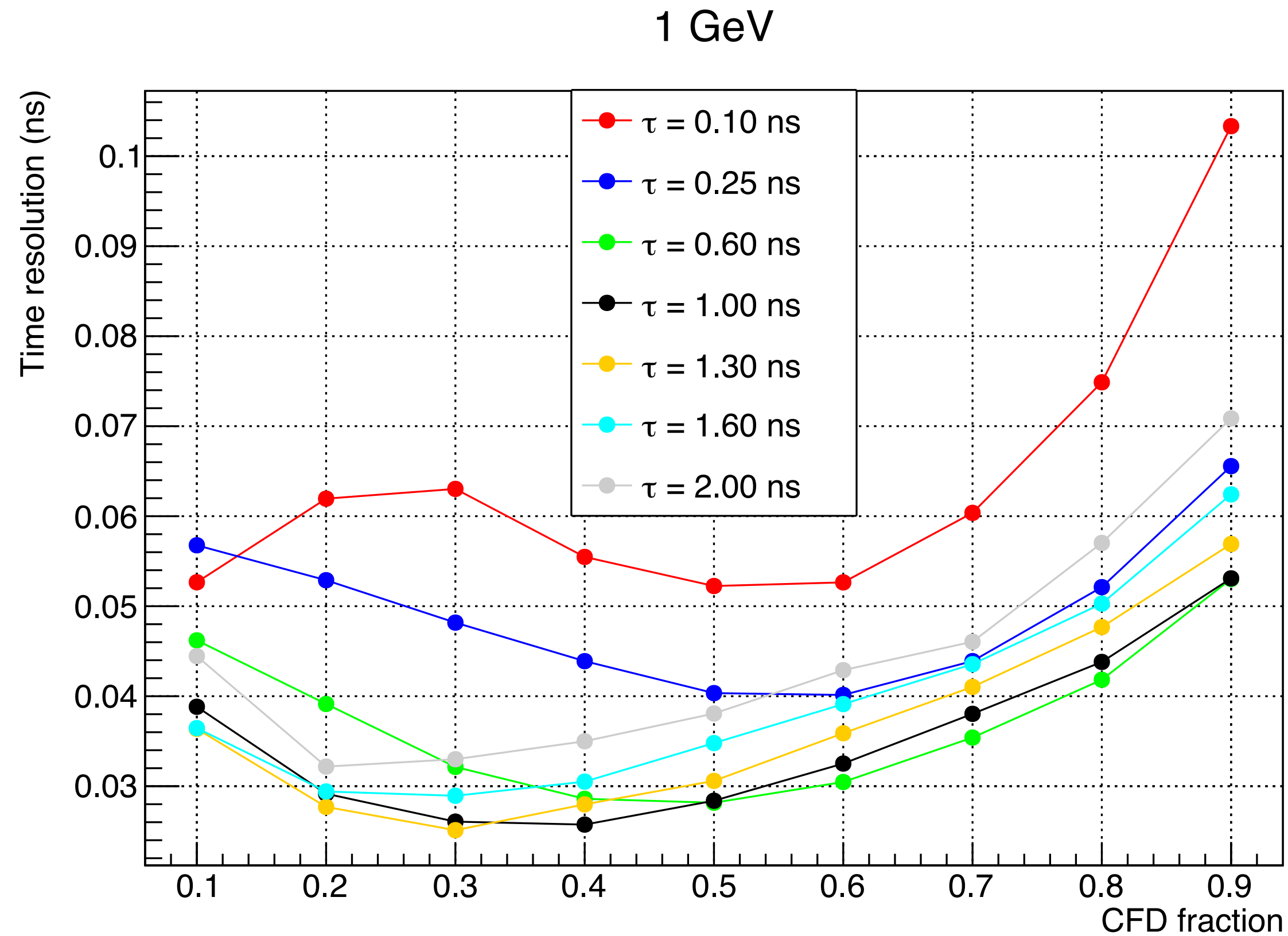
Time stamp distribution



Rise time distribution

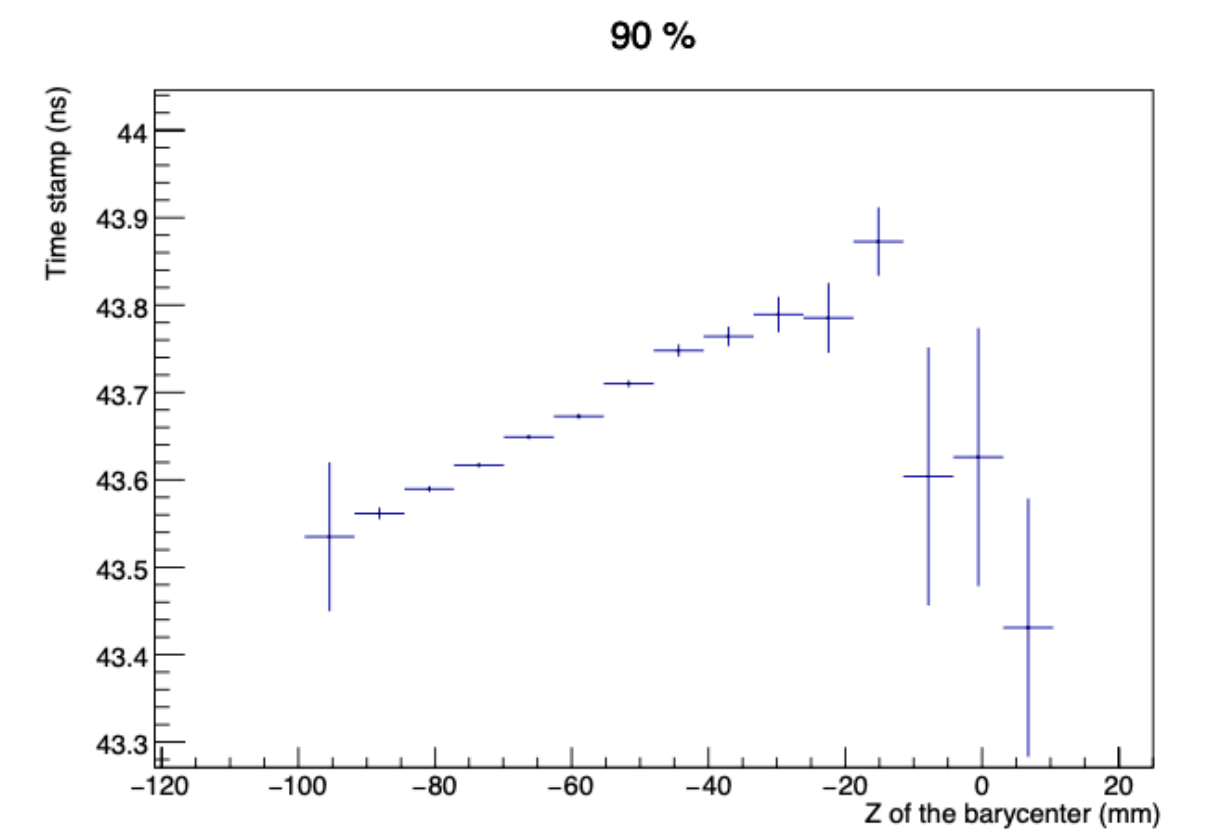
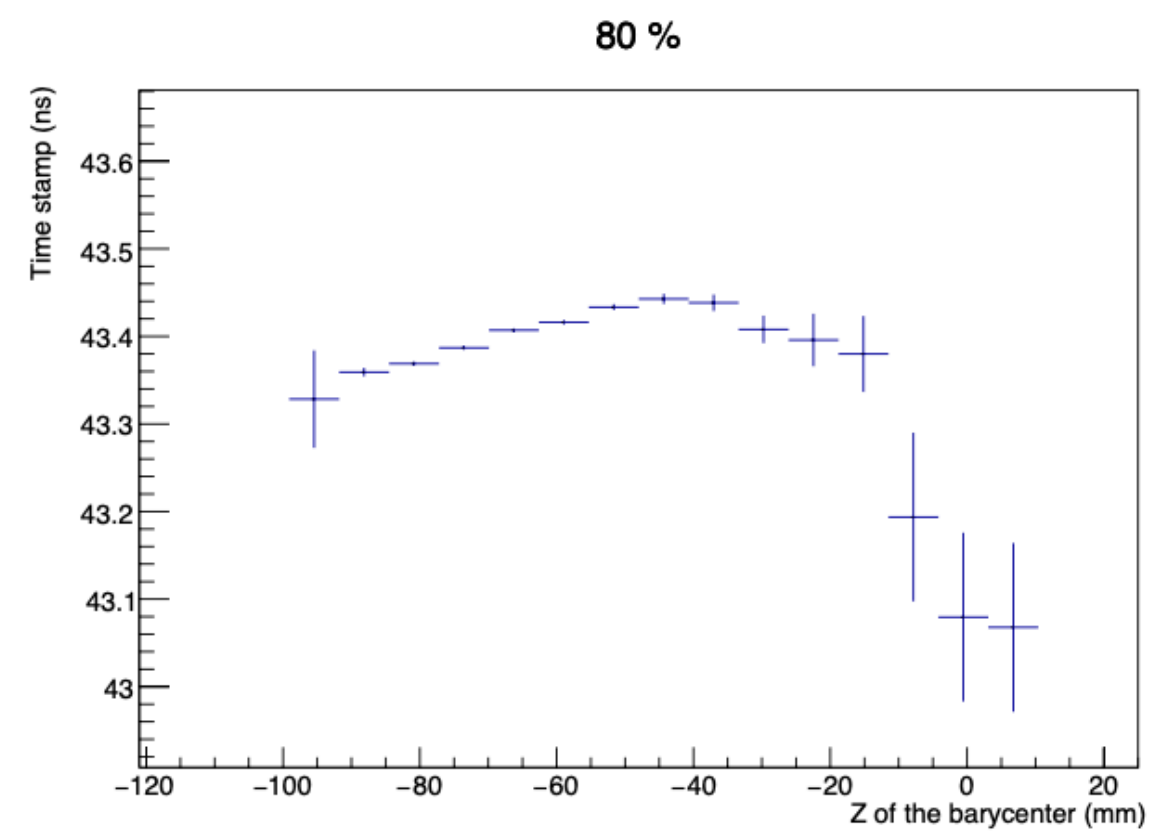
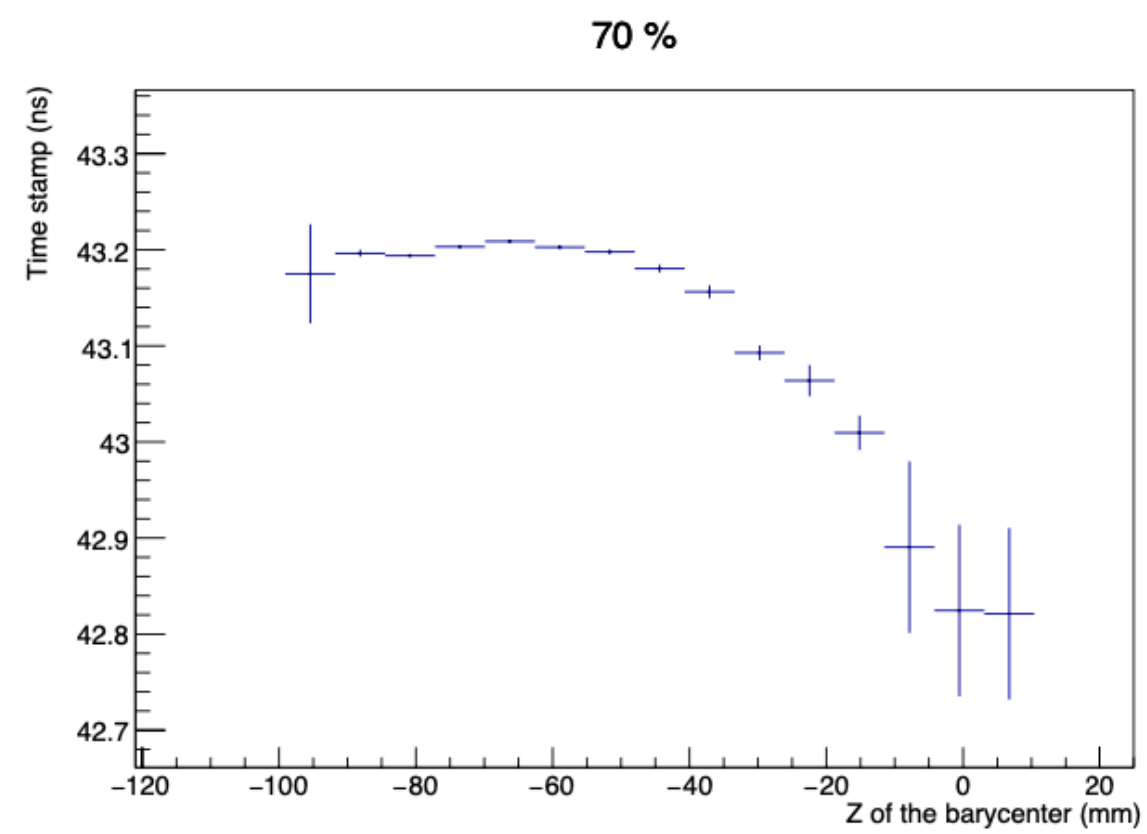
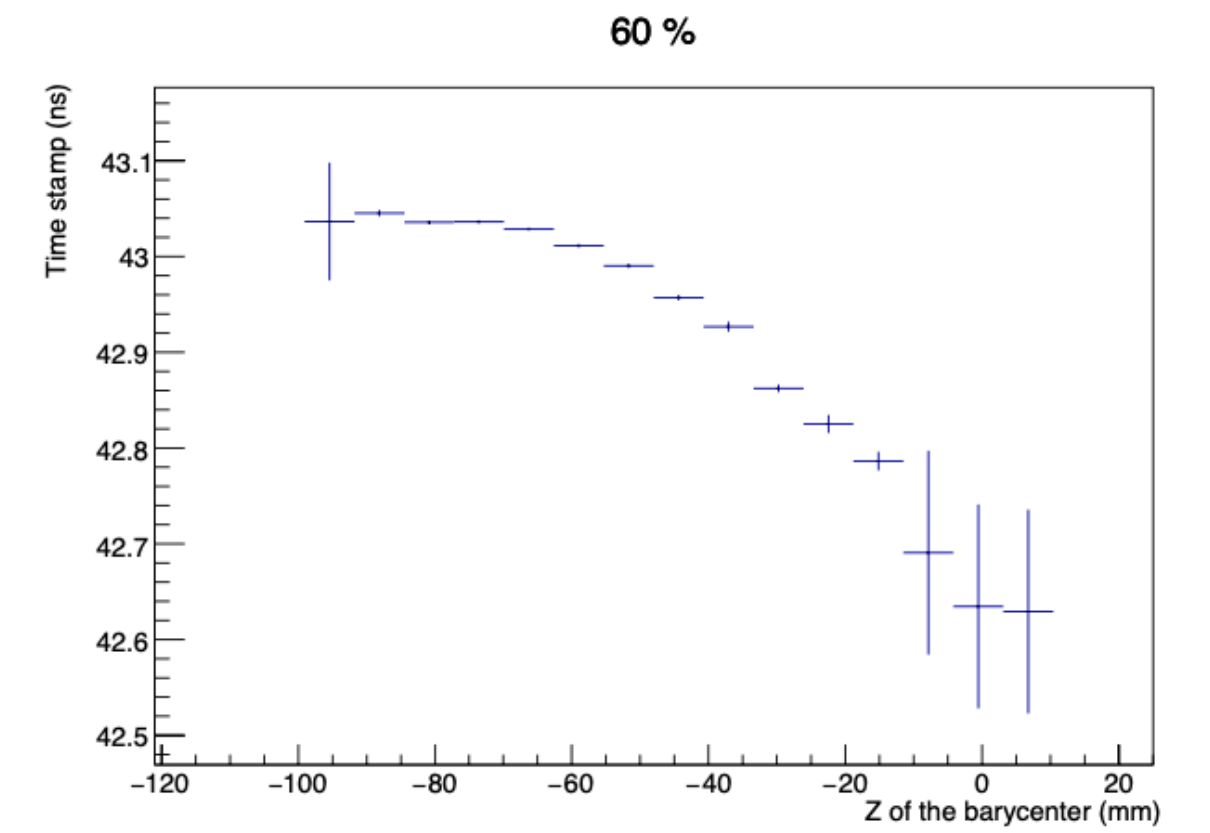
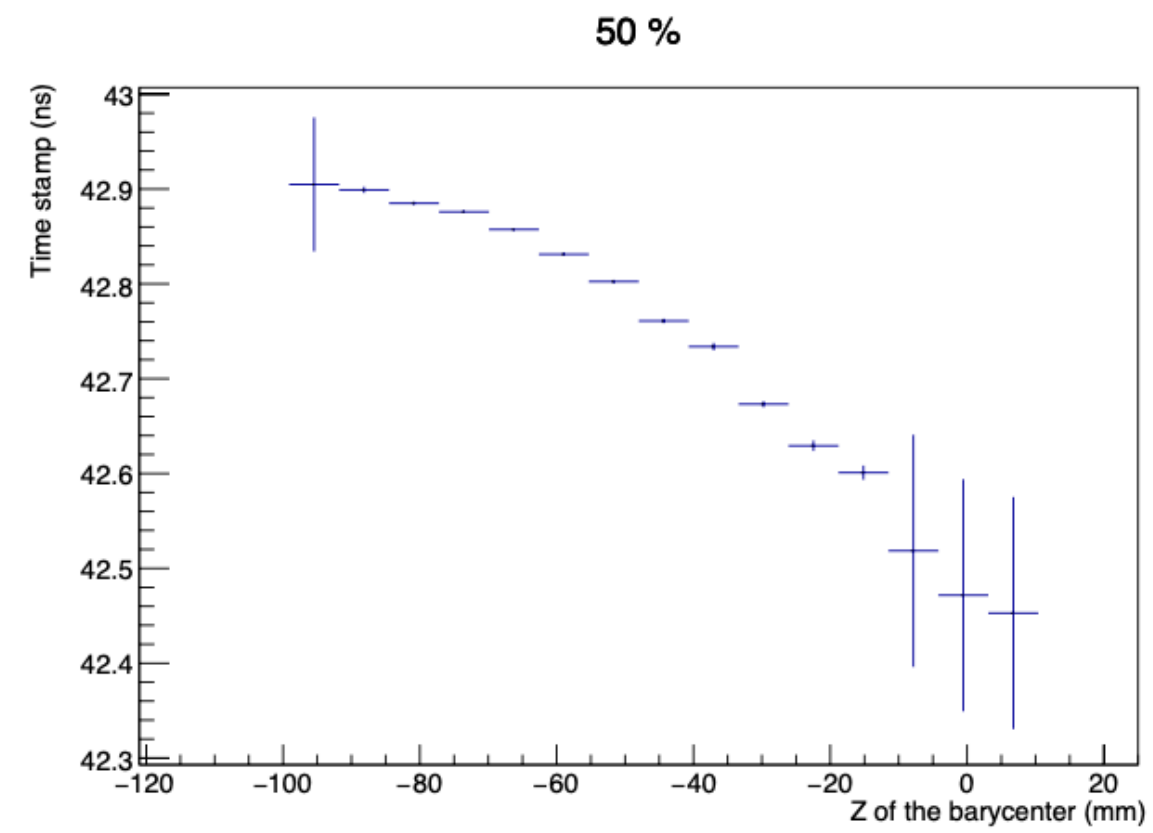
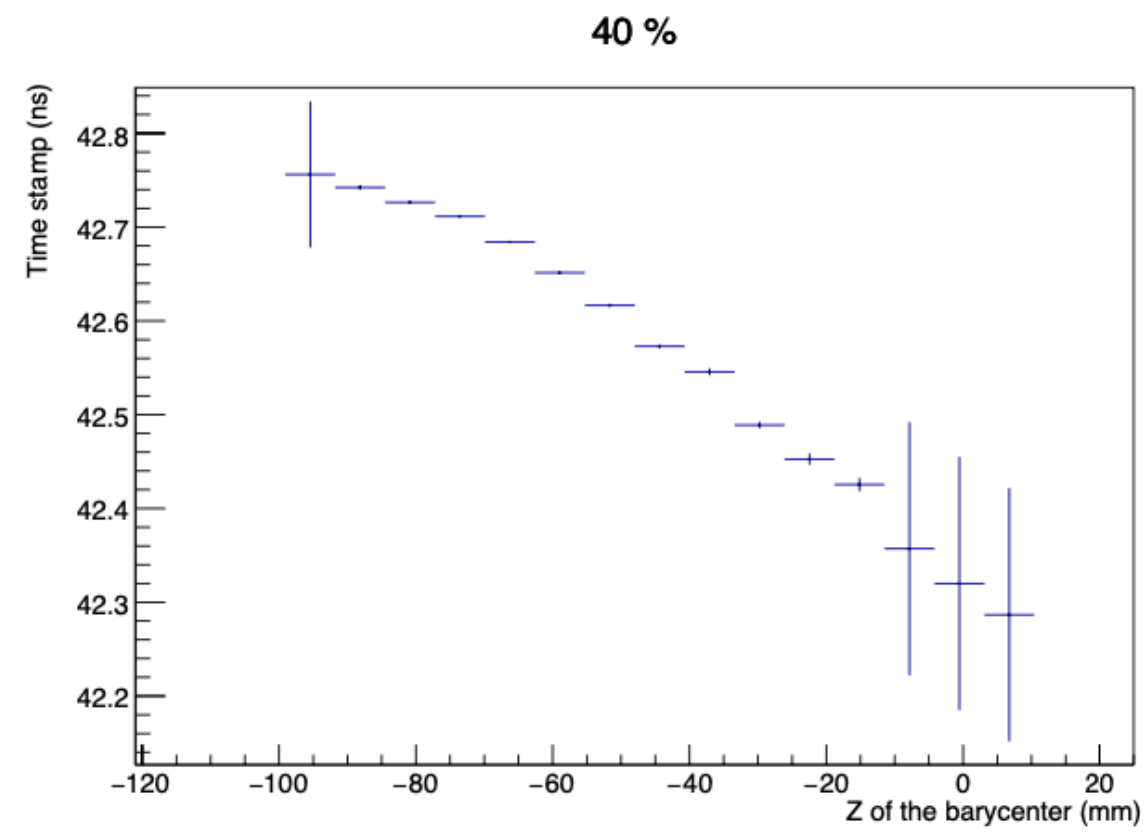
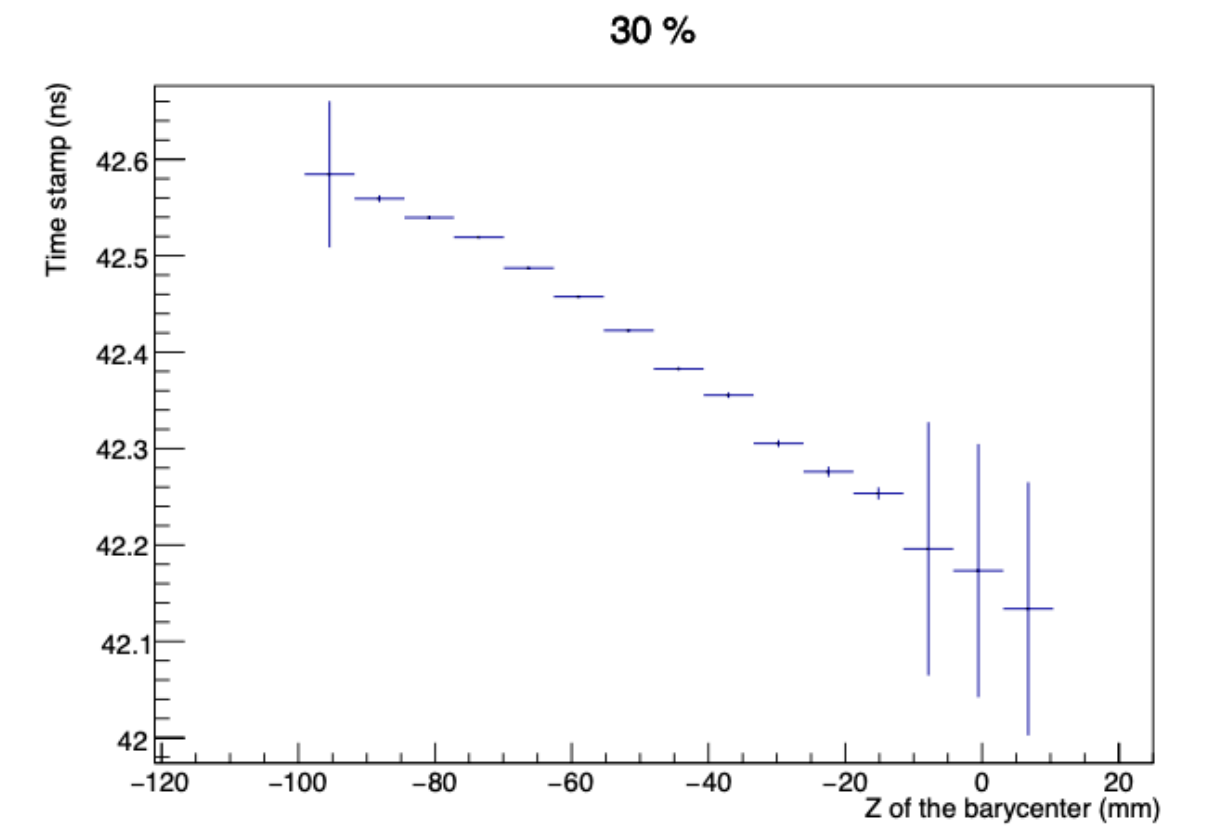
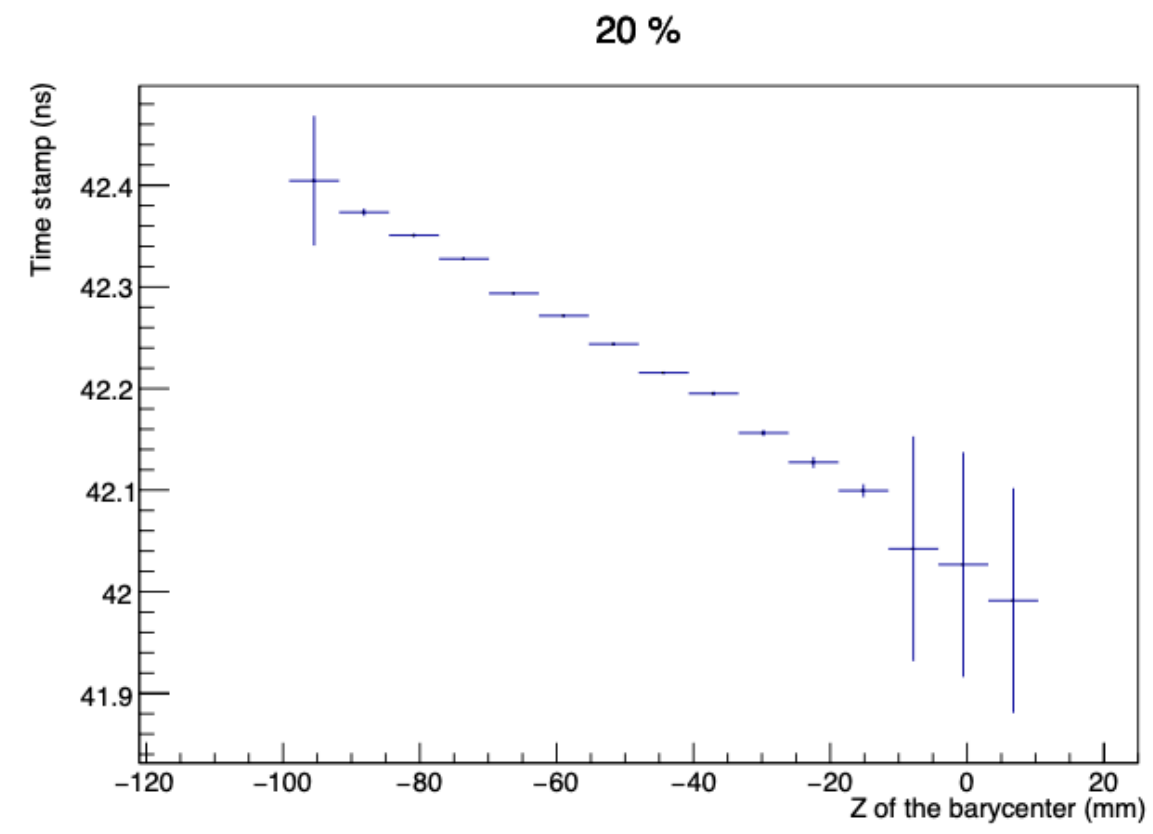
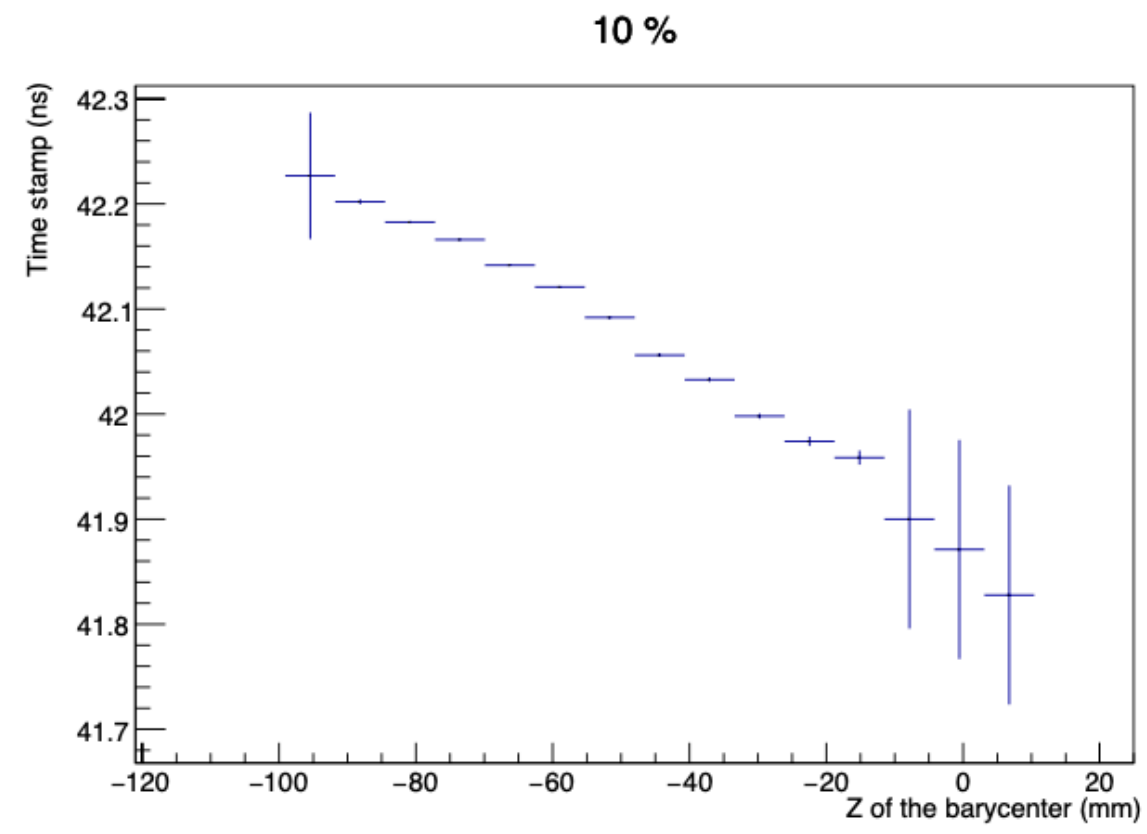


CFD fraction scan

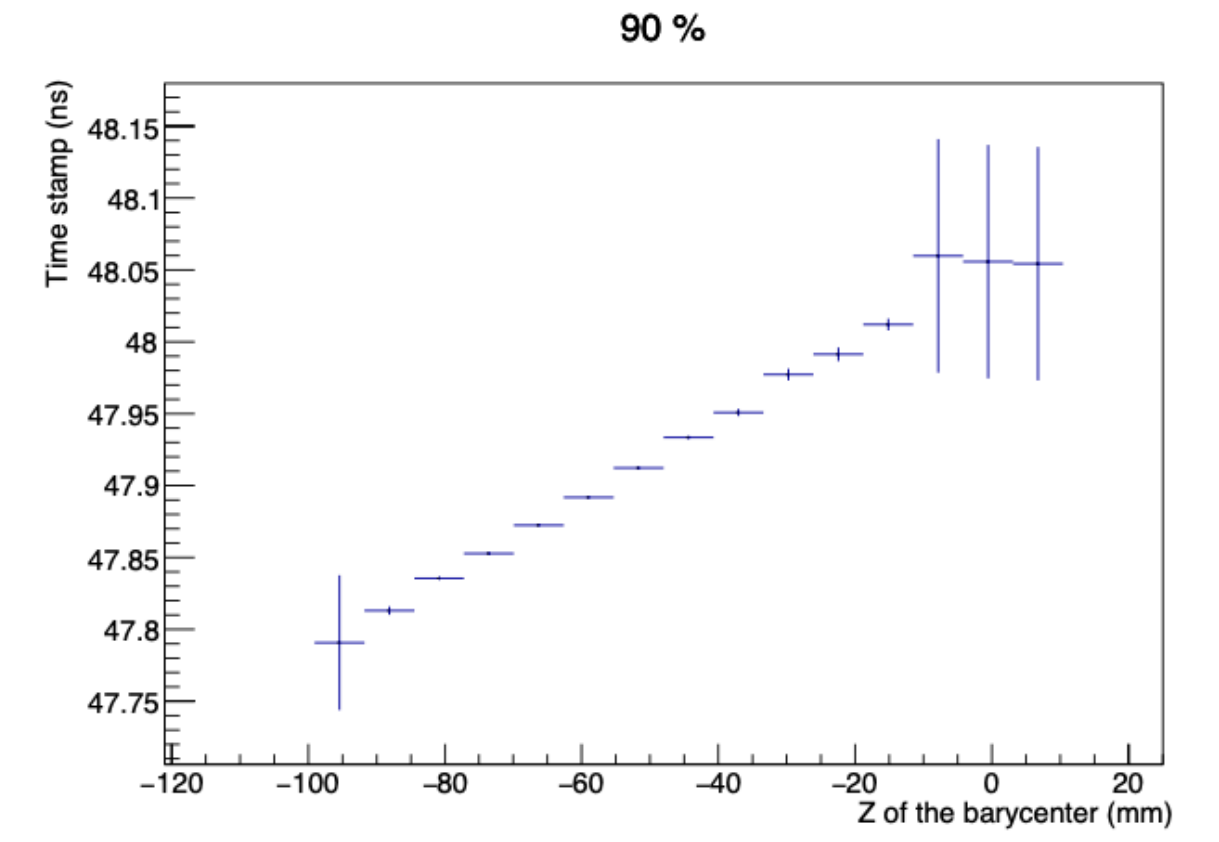
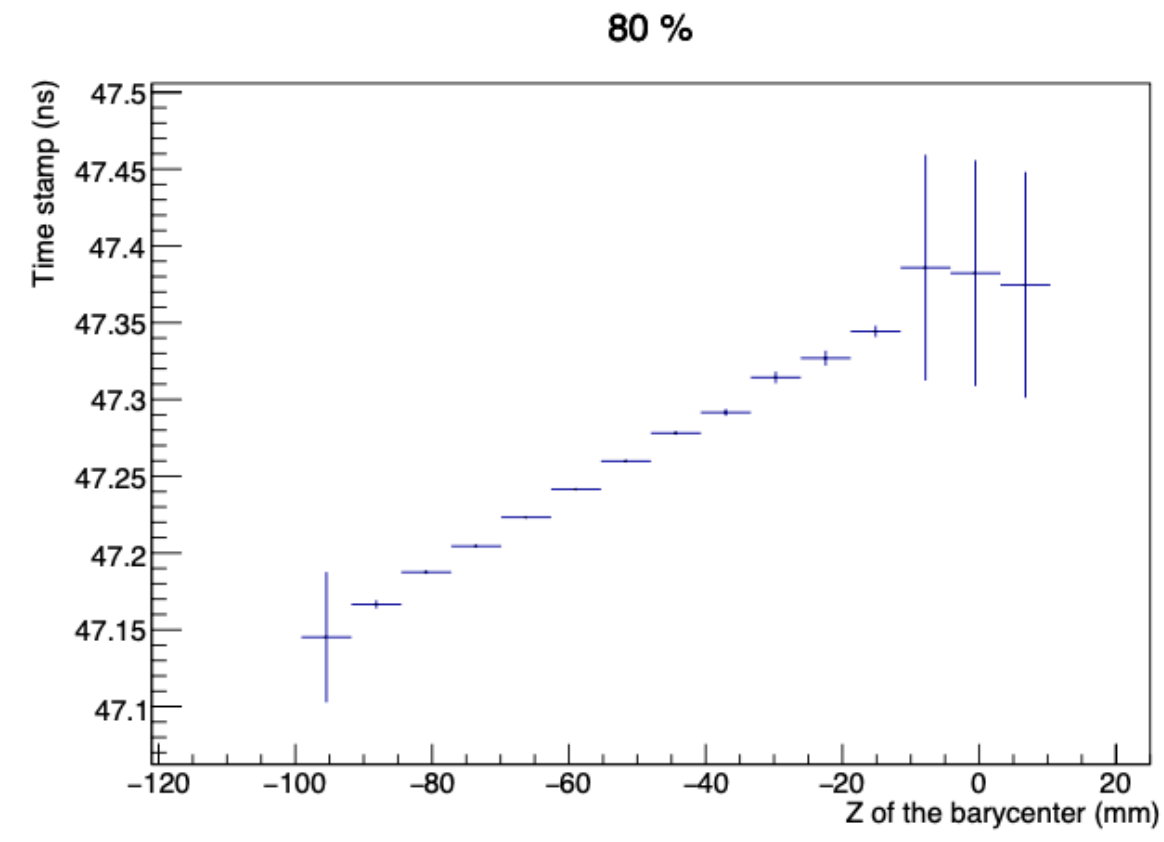
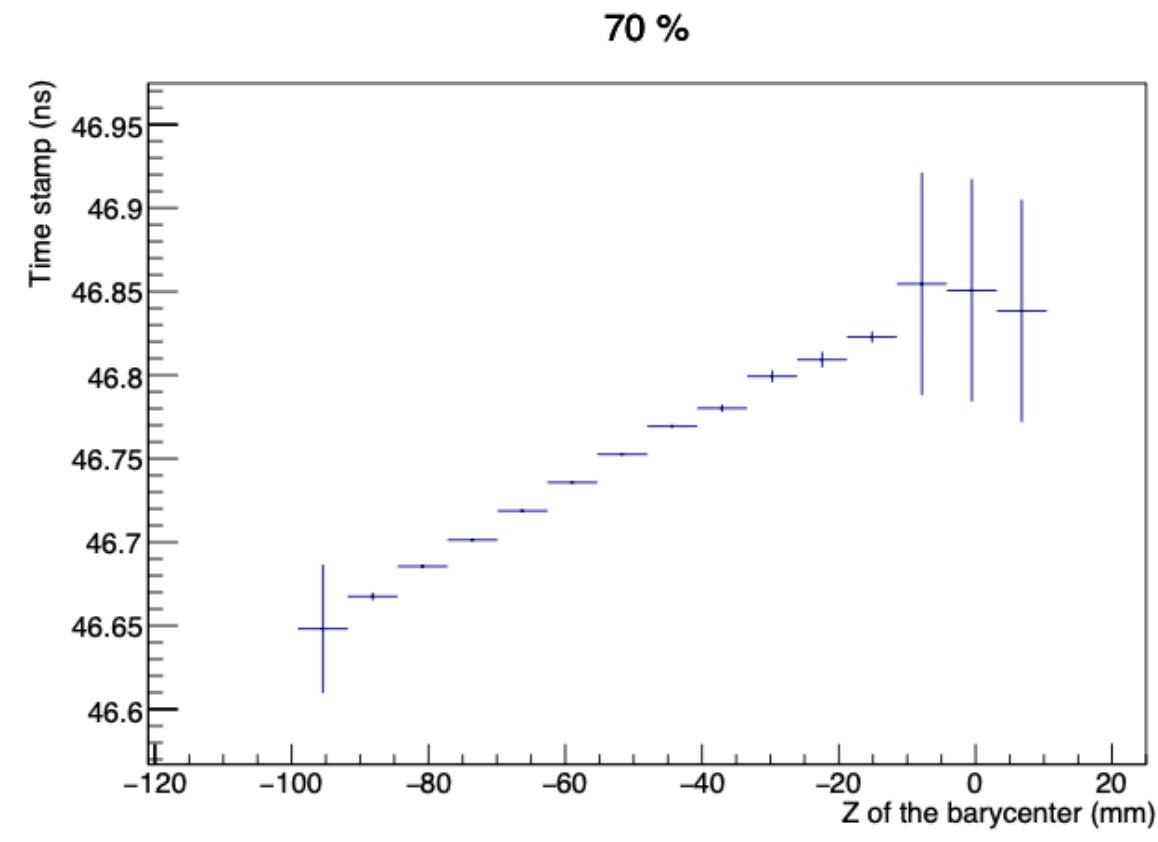
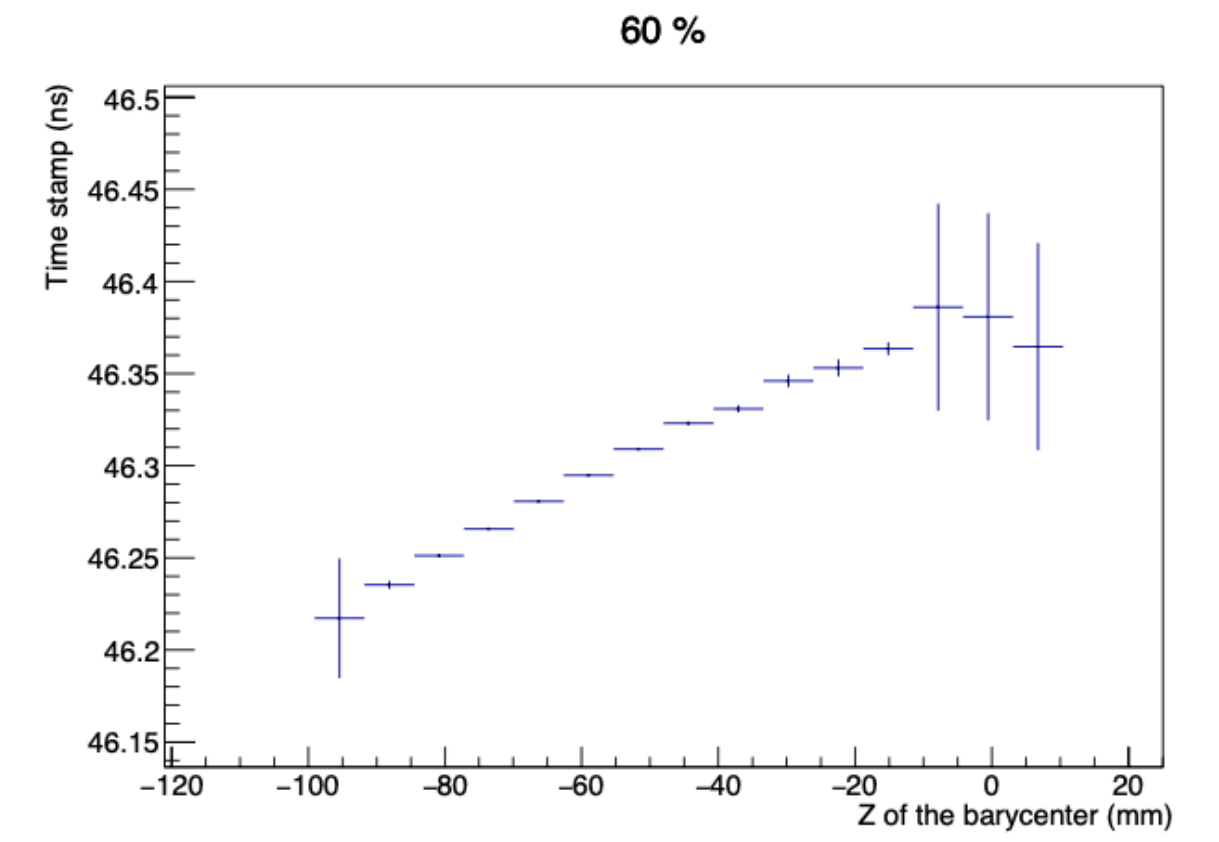
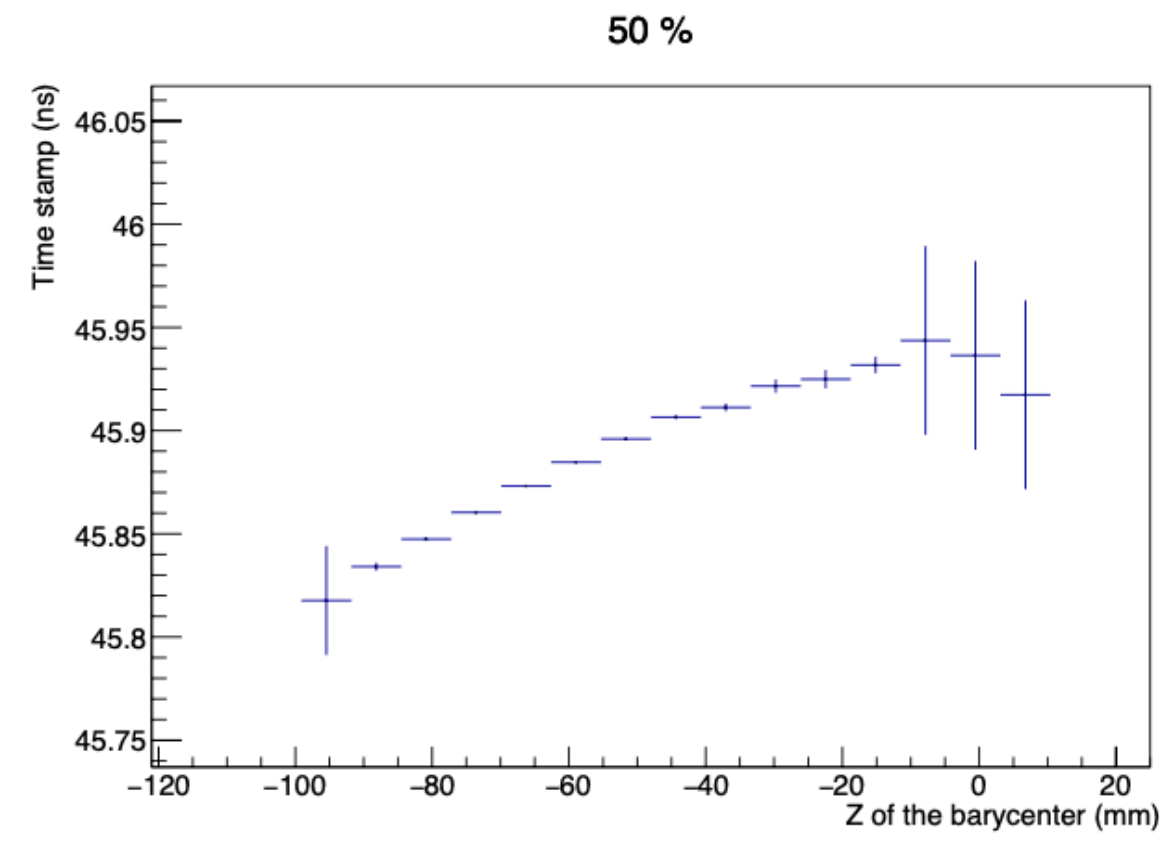
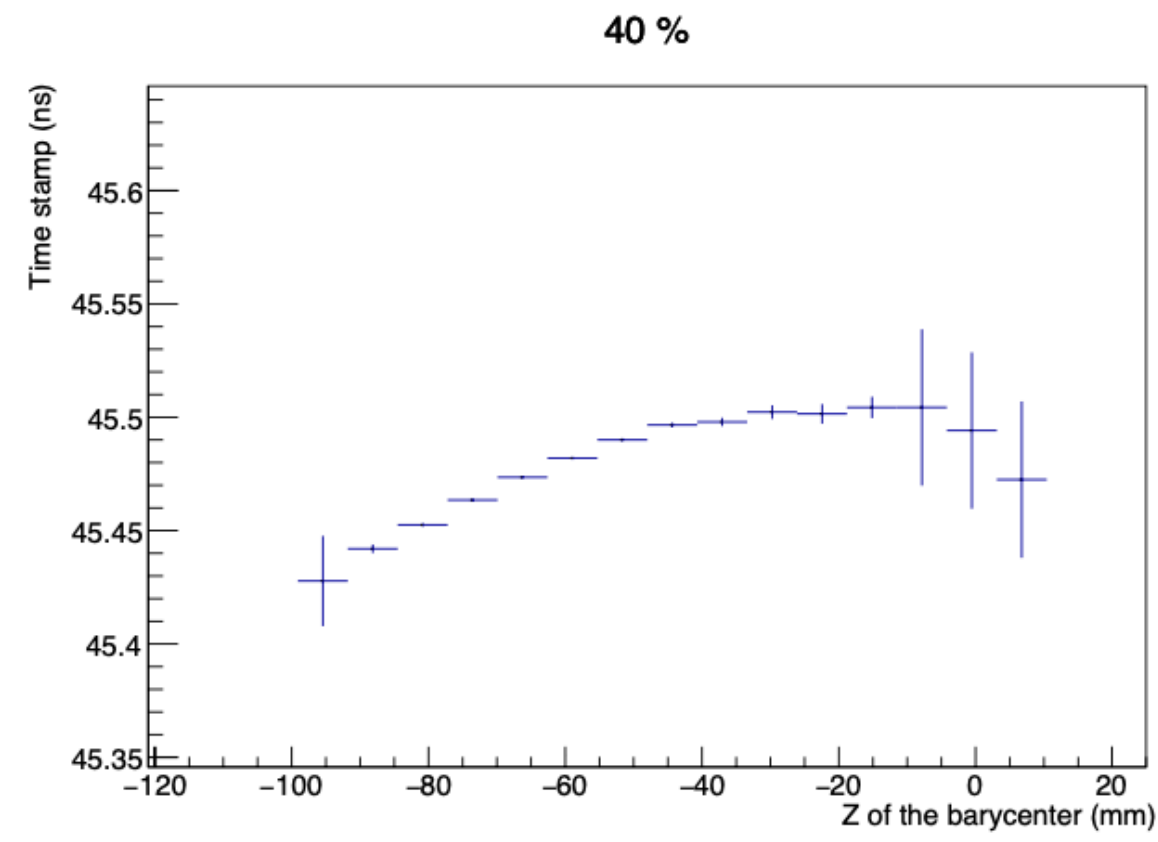
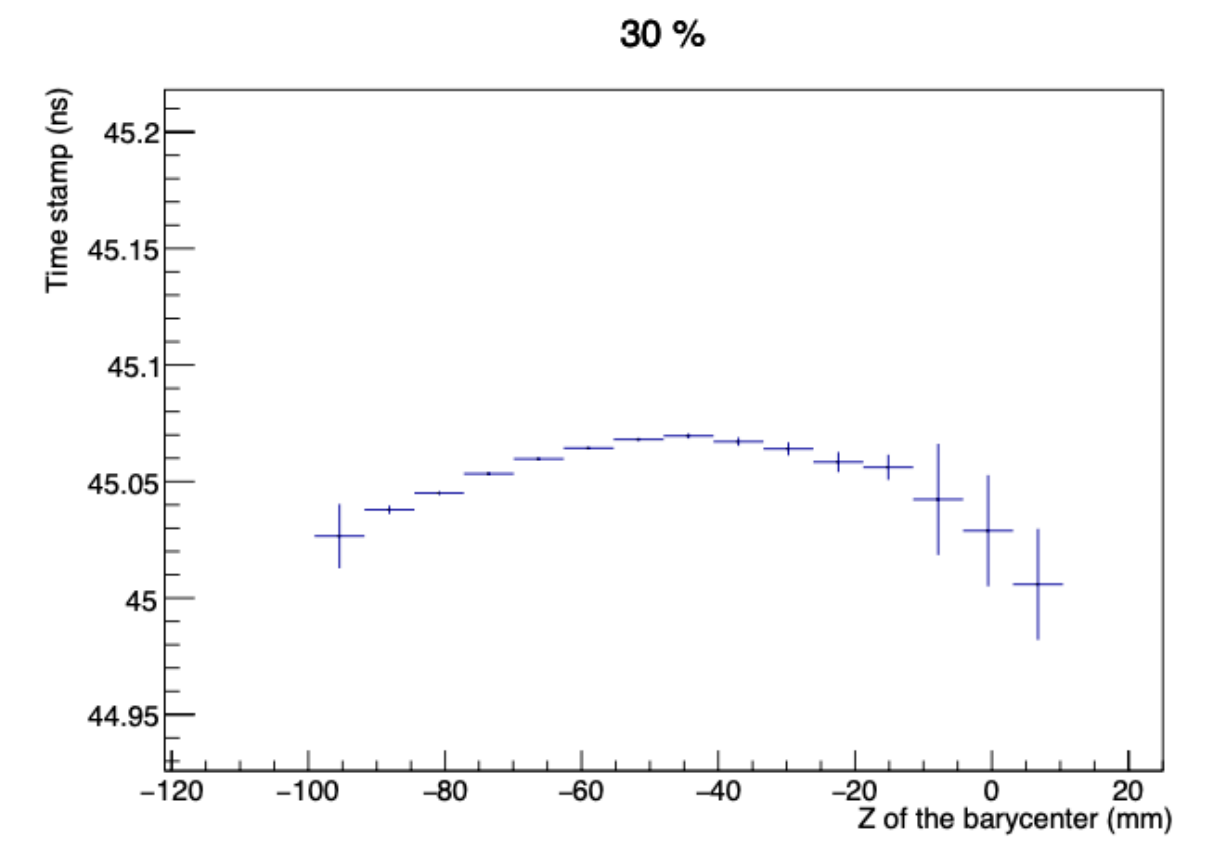
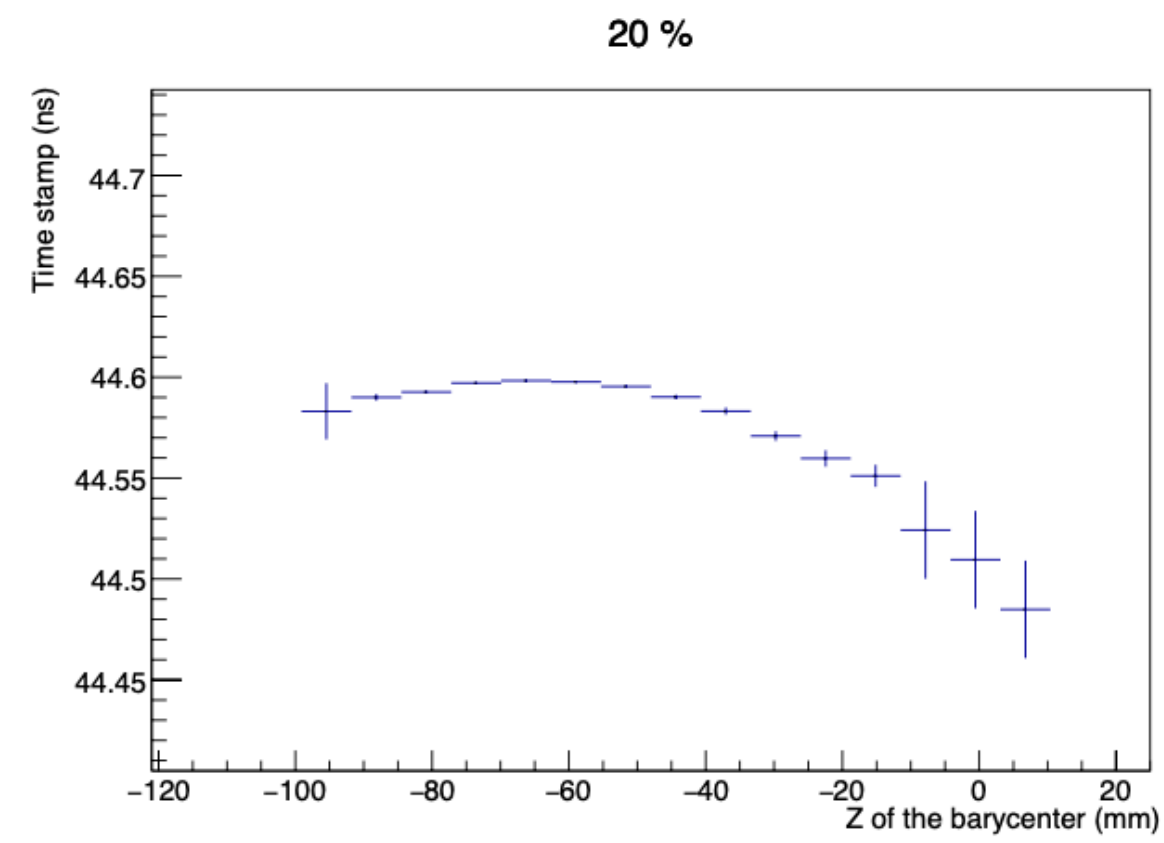
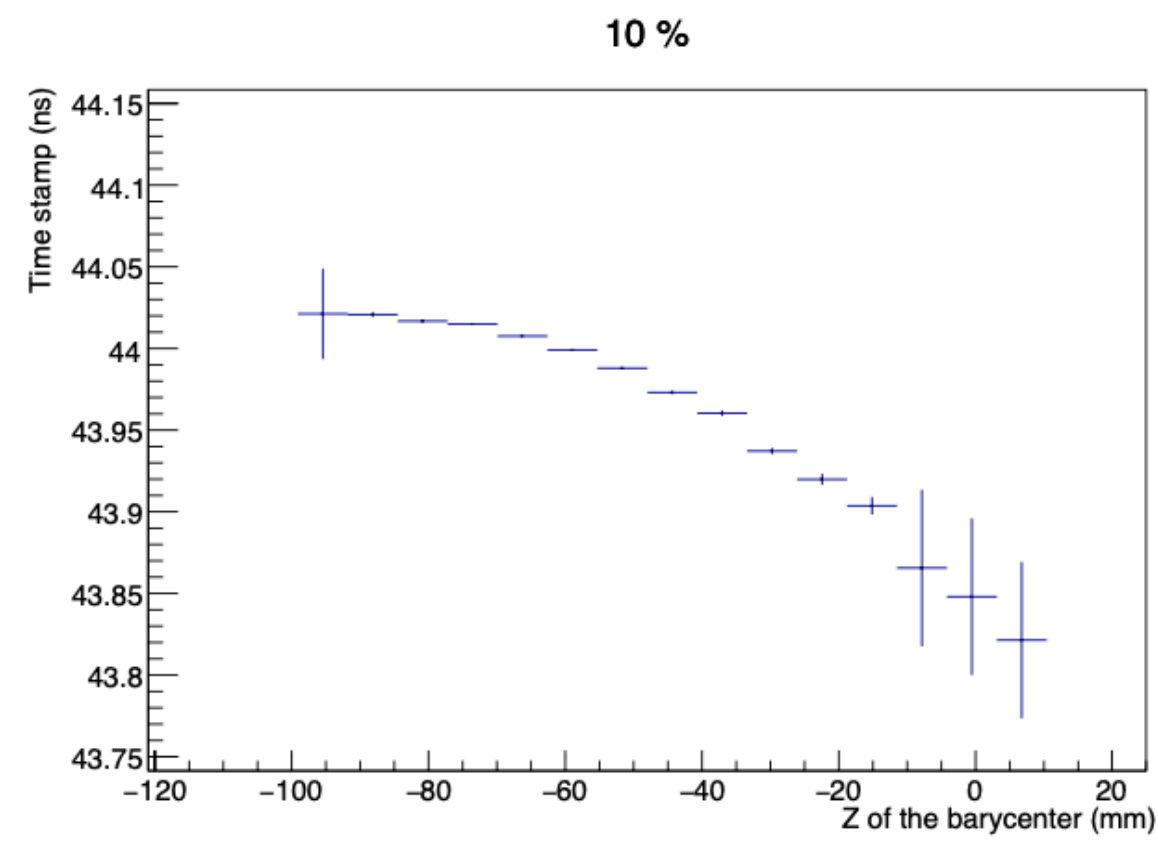


- The best CFD threshold tends to decrease for slower PMTs
- Best values always between 20% and 70%

**Tau = 0.1 ns
(FWHM = 0.35 ns)**

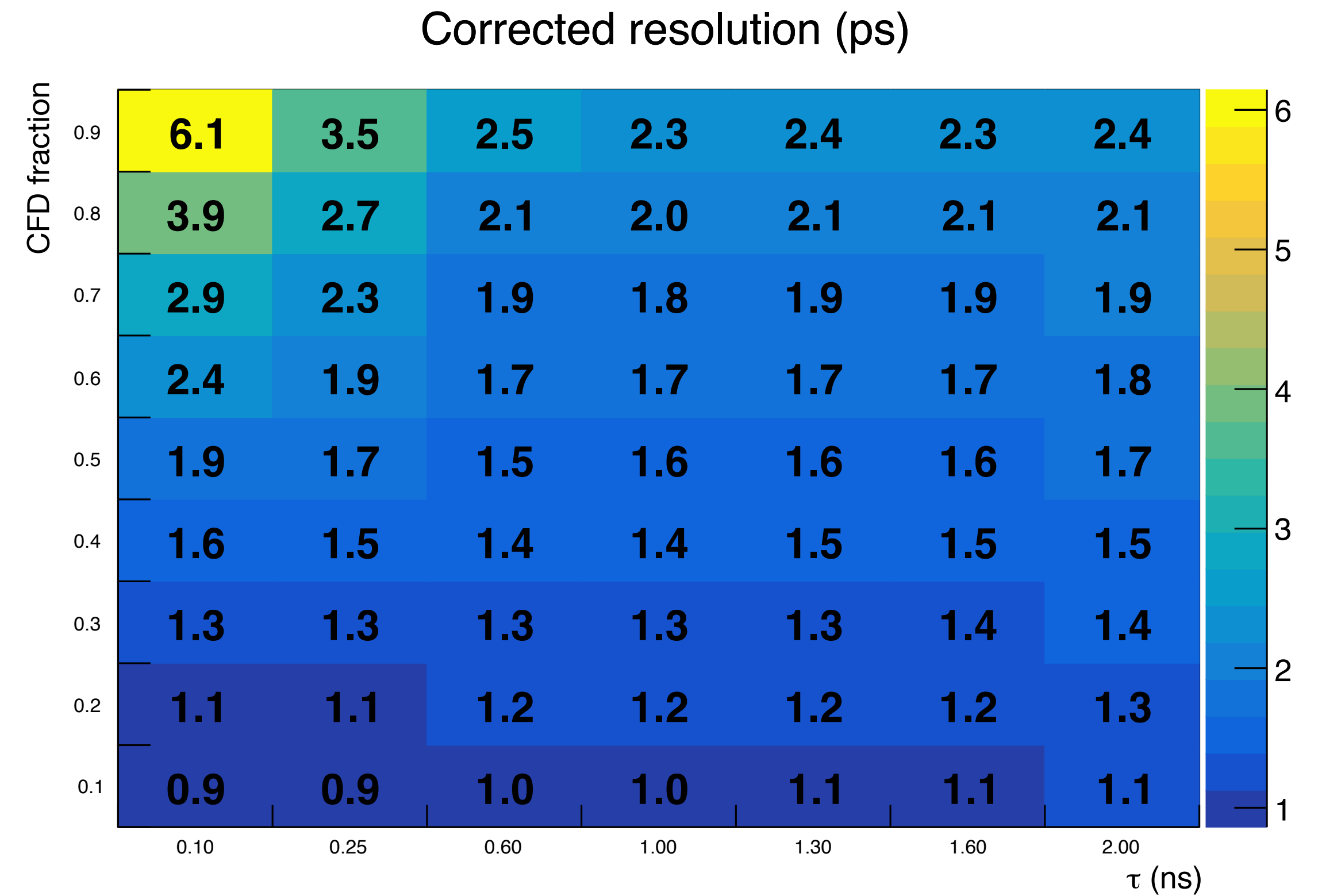
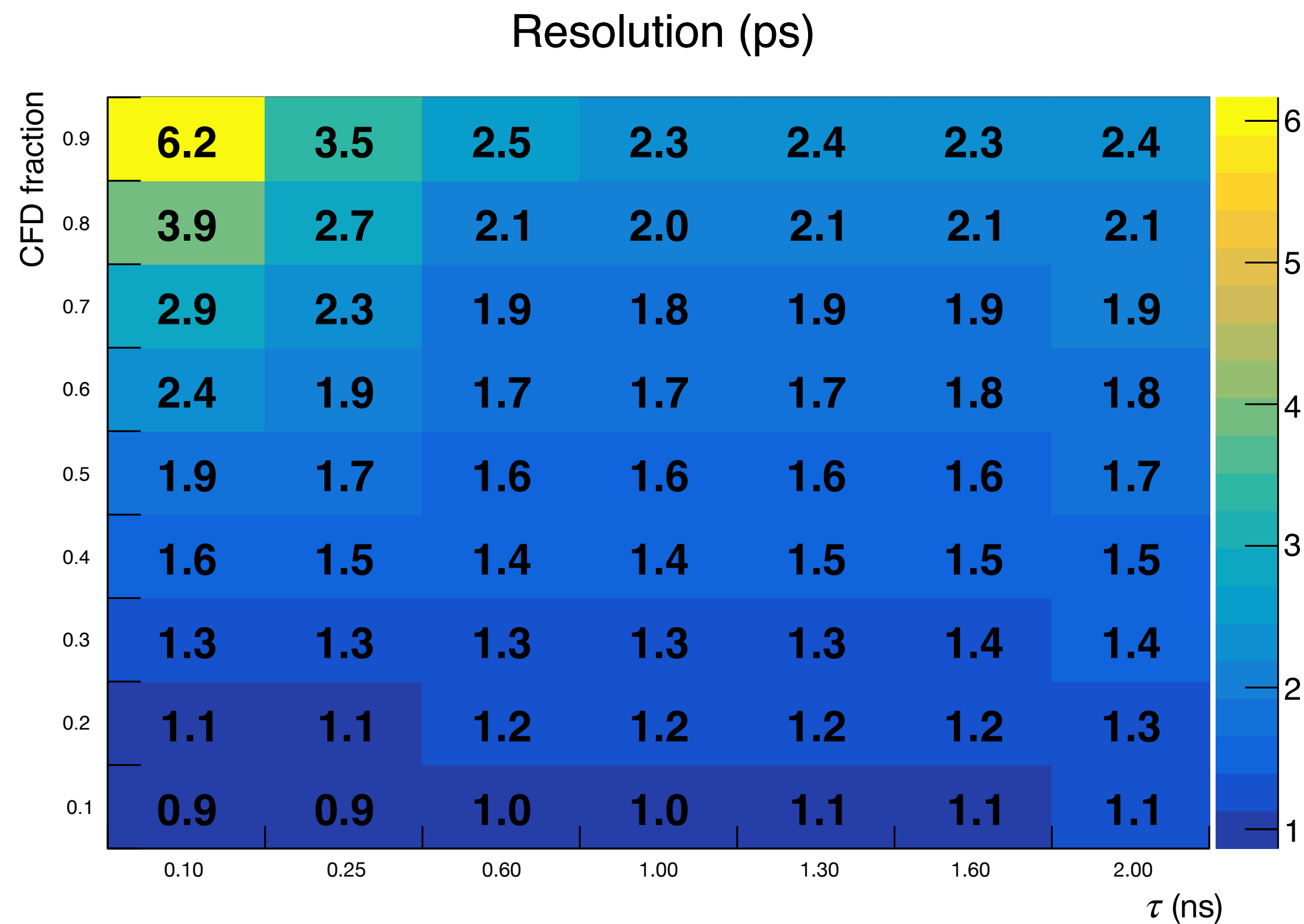


Tau = 2.0 ns
(FWHM = 7 ns)



Why are slower PMTs better ?

- Arrival time to the PMT window: $t_{total} = t_{generation}$
- Corrections to the time resolution are negligible



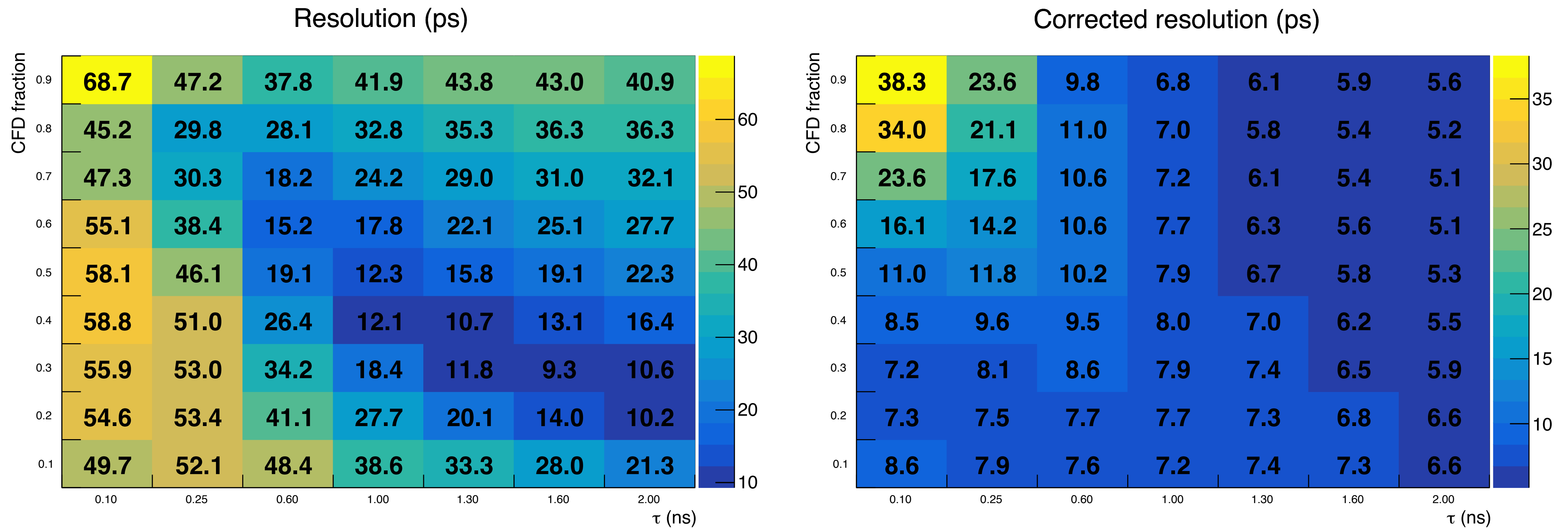
Why are slower PMTs better ?

- Arrival time to the PMT window: $t_{total} = t_{generation} + t_{deposition}$
- ~45 ps worsening of the biased resolution
- ~40 ps corrections, homogeneous behaviour in the parameter space



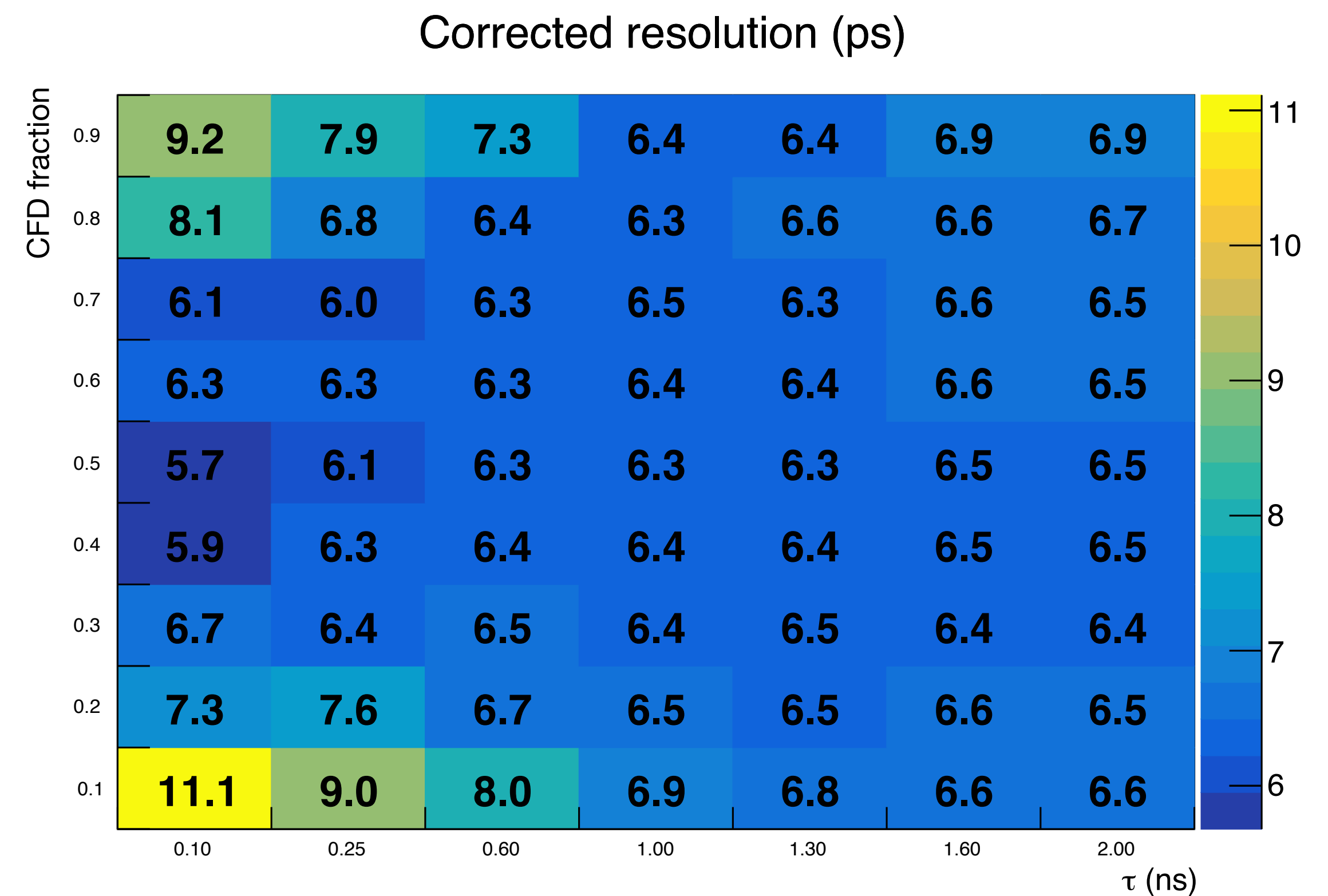
Why are slower PMTs better ?

- Arrival time to the PMT window: $t_{total} = t_{generation} + t_{deposition} + t_{propagation}$
- Best resolutions: when the two correlations (reflected/direct photons) cancel out each other
- After correction: more homogeneous behaviour in the parameter space

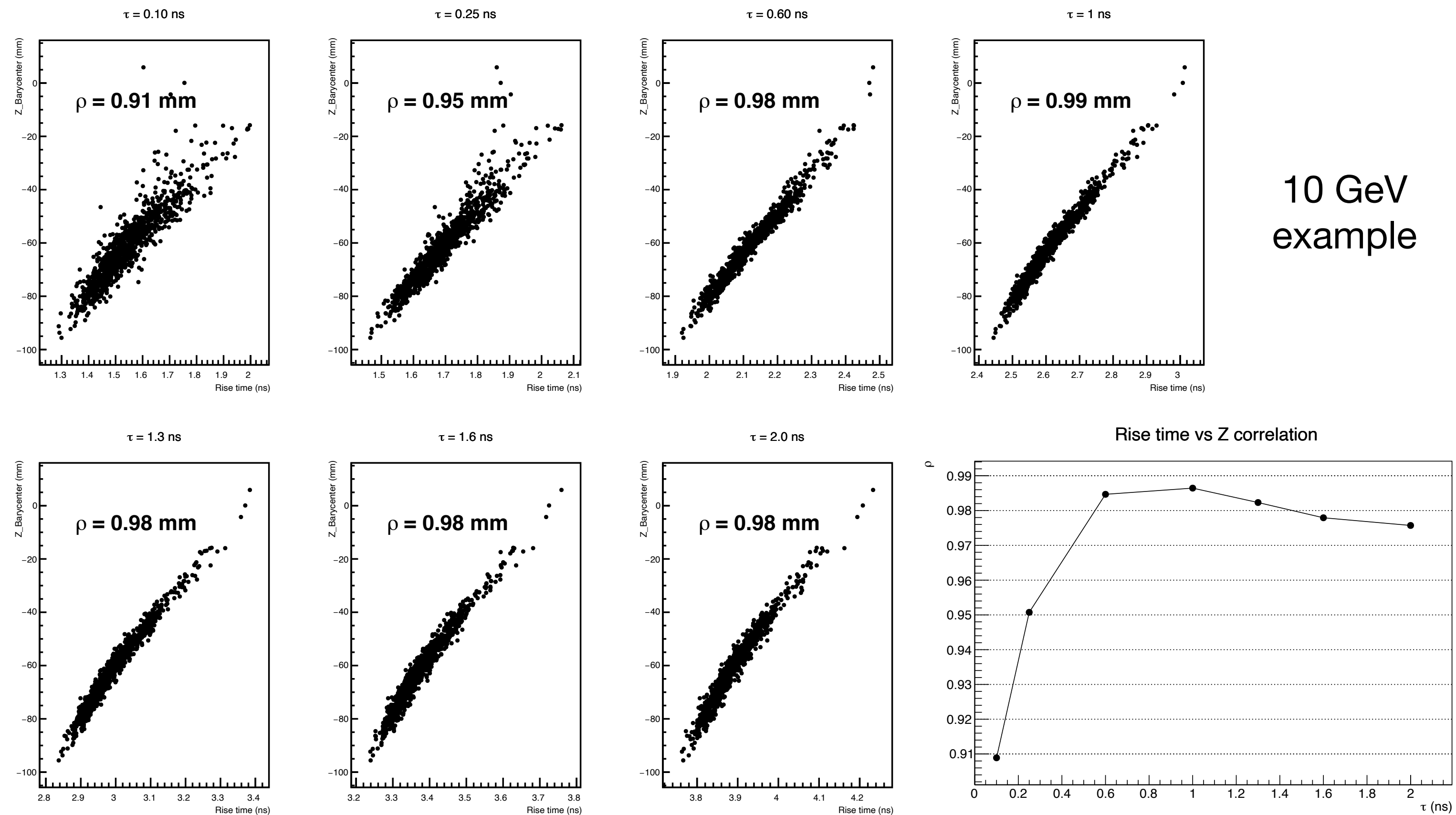


Why are slower PMTs better ?

- Arrival time to the PMT window: $t_{total} = t_{generation} + t_{deposition} + t_{propagation}$
- **Only direct photons considered**
- Homogeneous behaviour in both cases



Shower depth vs rise time



- High correlation between shower depth and rise time:
 - Deeper showers present a **higher spread** in the Z direction
 - Deeper showers feature **more separated peaks** in the propagation time distribution of the photons

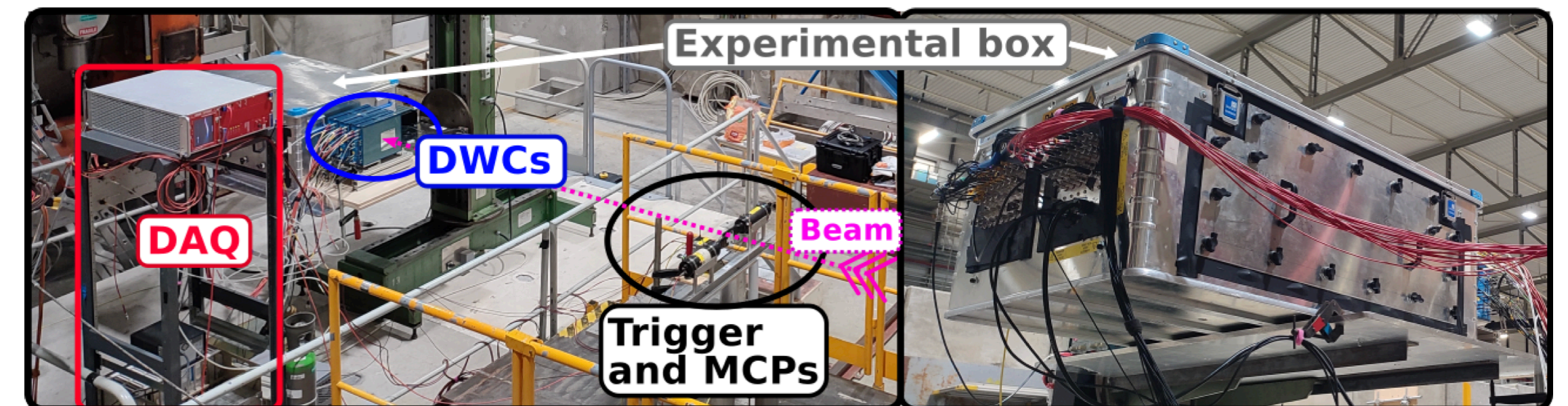
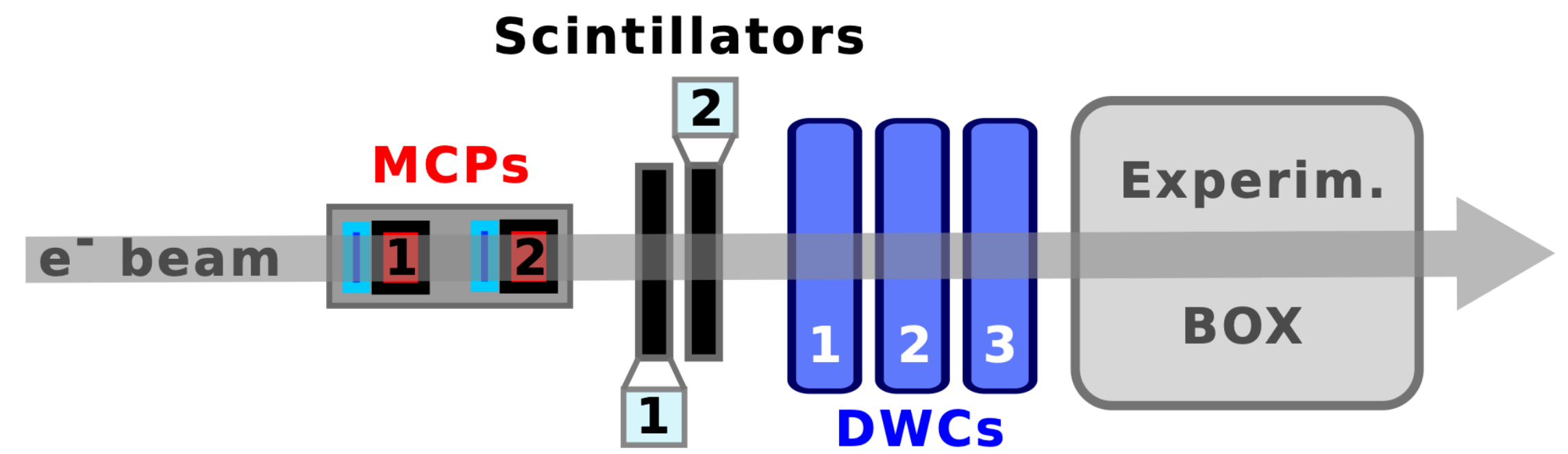
SPS testbeam setup

- June 2024 at the CERN Super Proton Synchrotron (SPS)
- ~ 2 weeks of data-taking

- Characterization of SpaCal and Shashlik prototypes
- e^- and hadrons beams (20 GeV - 100 GeV)

Experimental setup

- **Trigger:** scintillators + PMTs
- **Timing reference:** MCPs ($\sigma \simeq 15$ ps)
- **Tracking:** DWCs ($\sigma_{xy} \simeq 200$ μm)
- Prototype inside a thin and dark experimental box
- **Digitizer:** 5×10^9 Hz (200 ps sampling period)



Few technical details

- Spatial selection: 10x10 mm² around the centre of one cell
- **R7600U** and **R9880U** feature **Extended Red Multi Alkali (ERMA)** photocathodes
- **R9880U** and **R14755U** are the **fastest ones**

Quantum efficiency

	SCSF-78	3HF
R7600U	12%	12%
R9880U	15%	18%
R14755U	26%	10%
Tilecal	/	/

(Taken from data sheet)

Single photoelectron pulse

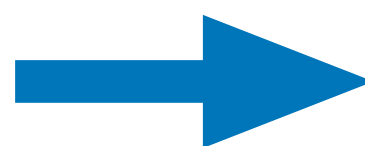
	FWHM (ns)	Rise time (ns)
R7600U	3.2	1.6
R9880U	1.25	0.57
R14755U	0.68	0.4
Tilecal	/	/

(Taken from data sheet)

Small module - fit parameters

SCSF-78 fibres

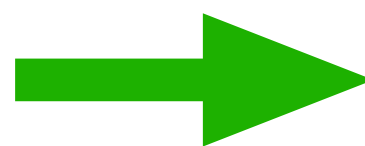
Before correction	Sampling term (ps x GeV ^{1/2})	Constant term (ps)
R7600U	113 ± 3	14.5 ± 0.5
R9880U	169 ± 4	18.8 ± 0.6
R14755U	148 ± 5	24.8 ± 0.6
Tilecal	95 ± 4	14.4 ± 0.5



After correction	Sampling term (ps x GeV ^{1/2})	Constant term (ps)
R7600U	103 ± 3	13.5 ± 0.4
R9880U	127 ± 3	16.9 ± 0.5
R14755U	87 ± 5	23.4 ± 0.4
Tilecal	92 ± 4	14.5 ± 0.5

3HF fibres

Before correction	Sampling term (ps x GeV ^{1/2})	Constant term (ps)
R7600U	174 ± 3	14.1 ± 0.8
R9880U	223 ± 7	19.6 ± 1.4
R14755U	116 ± 12	48.6 ± 0.6
Tilecal	303 ± 6	27.4 ± 1.2

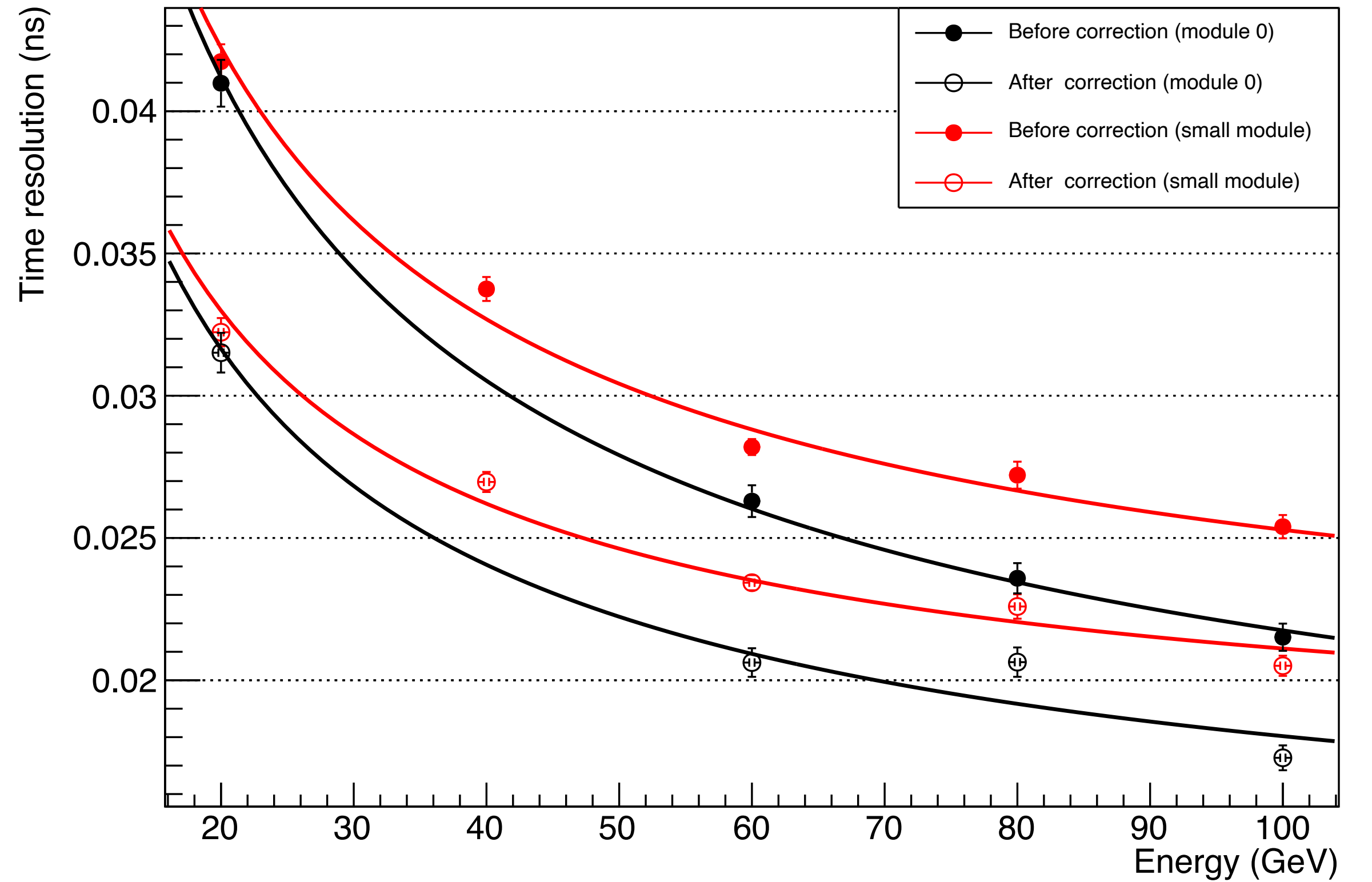


After correction	Sampling term (ps x GeV ^{1/2})	Constant term (ps)
R7600U	157 ± 4	14.5 ± 0.7
R9880U	179 ± 7	21.0 ± 1.0
R14755U	202 ± 5	20.5 ± 0.9
Tilecal	293 ± 6	29.5 ± 1.1

Small module vs module 0 - fit parameters

	Sampling term (ps x GeV ^{1/2})	Constant term (ps)
Small module	169 ± 4	18.8 ± 0.6
Small module (corrected)	127 ± 3	16.9 ± 0.5
Module 0	175 ± 5	12.9 ± 1.1
Module 0 (corrected)	130 ± 5	12.5 ± 0.8

W-Poly small module vs module 0



Uncertainties

- Several sources of uncertainties
 - Quantum efficiency not precisely known (taken from producer's data sheet)
 - Laser instability affecting the measurements
 - General electronic noise
- Fluctuations presenting a long period (~ 1 minute) when measuring I_{anode}

Physical quantity	Uncertainty	Type
Power	0.05 nW	absolute
Dark current	10%	relative
Anodic current	1%	relative
Quantum efficiency	0.05	absolute

N.B. These are not precision measurements and these uncertainty values are preliminary