



Search for new particles decaying to a pair of Higgs bosons with the ATLAS experiment

Nicolas Cartalade 30/09/2024 Dipartimento di Fisica UNIBO, Bologna



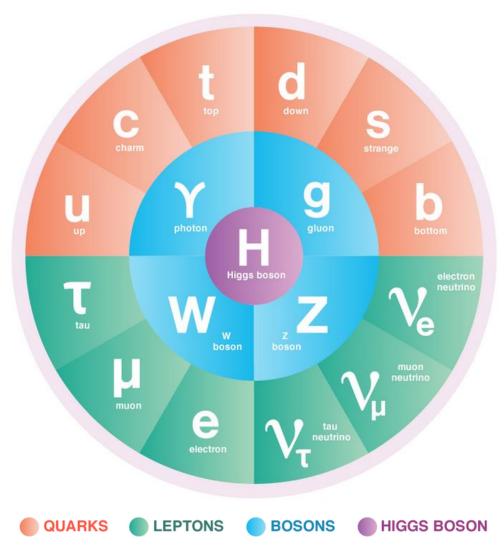


- Introduction
- Why are we studying the Higgs pair production ?
- Production process and decay mode HH
- Selection
- Modeling
- Final Fit
- Conclusions and next steps

The Standard Model (SM)



- Currently the most experimentally validated theory, it fails to explain some observations, suggesting that a more fundamental theory exists
- 12 **fermions** (6 quarks 6 leptons) : elementary constituents of matter
- 4 gauge **bosons**: mediators of fundamental interactions $(g, \gamma, Z^0, W^{\pm})$
- **Higgs boson** (2012): Brout-Englert-Higgs mechanism (1964) that **gives mass** to elementary particles



The Experimental Framework aboratoire de Physique de Clermon 44m The LHC (Large Hadron Collider) at CERN World's largest accelerator Proton-proton collider 25m - Genève Cern ixel detector LAr electromagnetic calorimeters Toroid ma Transition radiation tracker Solenoid magnet Muon chambers Semiconductor tracker ALICE 100 m **ATLAS** experiment LHC CMS

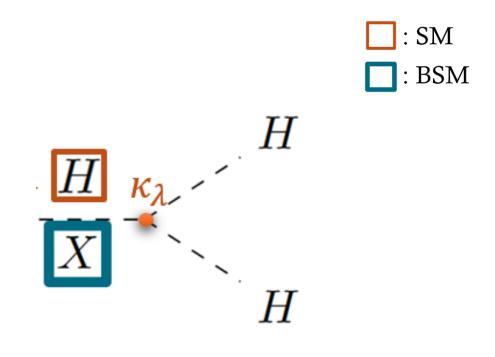
27 km in circumference

The **detector** (almost 4π coverage) of the Trackers and Calorimeters: designed to measure **trajectory** and **energy** of emitted particles

Tile calorimeters LAr hadronic end-cap and forward calorimeters



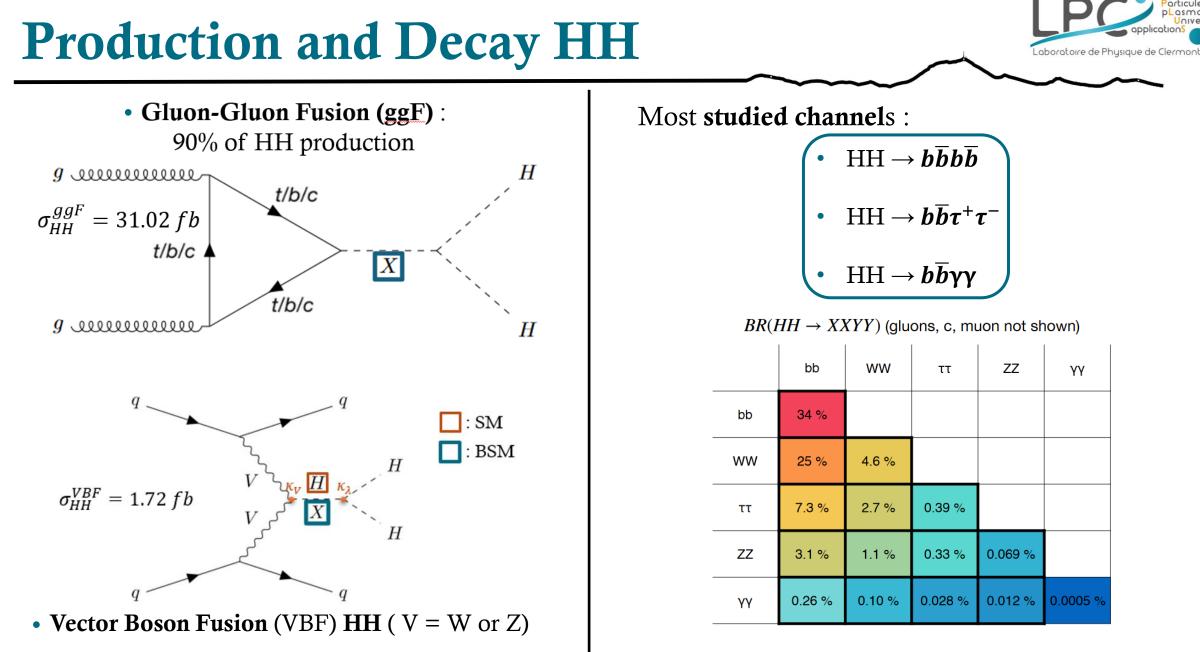
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Resonant Production mode: X is a scalar particle

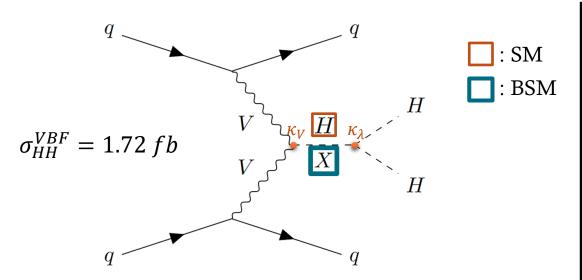
- Higgs mechanism gives mass to particles so we can probe the potential existence of a new massive X scalar boson through its interaction with the Higgs boson
- Possibility to discover a new particle or set limits on the production of a new scalar boson through this analysis and to characterize the Higgs boson selfcoupling (κ_λ)

Internship supervised by Louis D'Eramo



Production and Decay HH





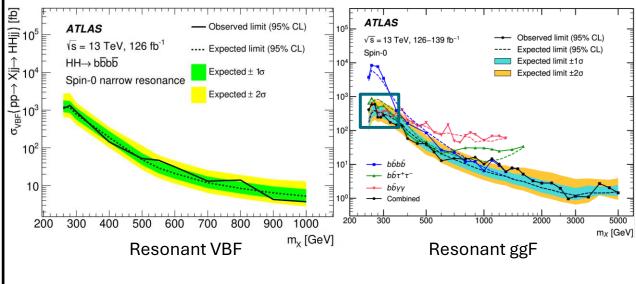
Vector Boson Fusion (VBF) **HH** (V = W or Z) **Resonant** : X is a spin-0 particle

- Cross-section for the single Higgs Boson VBF production which is a rare process is around $\sigma_H^{VBF} =$ 3.78 pb (2000 times the di-Higgs one)
- Detection only of signatures (final states) because the lifetime of the Higgs is too low (10⁻²² s)

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$\mathrm{HH} ightarrow b \overline{b} \gamma \gamma$

- $\mathbf{H} \rightarrow b\overline{b}$: High Branching Ratio (57%)
- H → γγ: photons allows to trigger events at a lower energy and to obtain a good resolution for m_{γγ} and m_X thanks to the electromagnetic calorimeter (Lar)

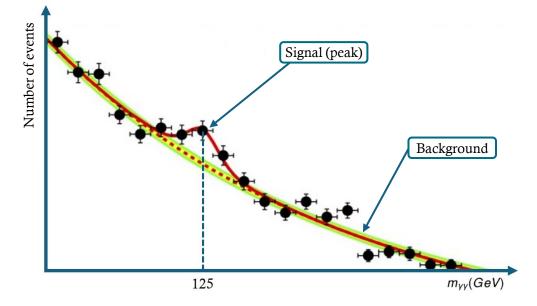


Objective : Final Fit

where $f_{signal}(x)$ denotes the **fitted** function for **the signal** (smallest χ^2), while $f_{bckg}(x)$ represents all the background, modeled by a function for the **continuous background**, and includes **single H** and **double Higgs** production, both described by the **same function** as the one of the **signal**.

 $f_{final}(x) = f_{signal}(x) + f_{bckg}(x)$

- Perform the **FIT** adding the **fitted pdf** of the **signal** and the one of the **background** (smallest spurious signal) to know if the **BDT** used in **SM HH** production is **adapted** to the **VBF Resonant** HH data.
- Establish an upper limit (exclusion zone for new models) on the cross-section $\sigma_{VBF}(X \rightarrow HH)$.

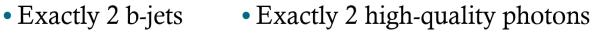




Selection





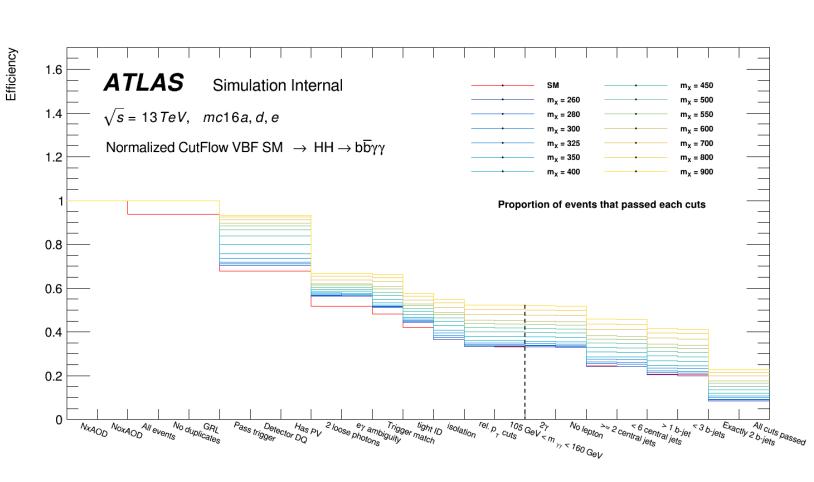


- < 6 central jets No lepton $(t \to bW \to b\ell^- \bar{\nu}_\ell)$
- Identification of 2 high-quality photons
 - with transverse momentum $p_T^{\gamma} > 35$ GeV and $|\eta^{\gamma}| < 2.5$ (η , the pseudo-rapidity, is an angle that describe the direction of a particle in the detector with respect to the beam pipe)
 - isolated to reduce background contamination (the energy within a certain area around the photon's energy deposit must remain below a defined threshold)
- The invariant mass of the two selected photons must be in the range $105 < m_{\gamma\gamma} < 160 \text{ GeV}$
- B-jet selection: Each b-tagged jet must have a transverse momentum $p_T > 30$ GeV and $|\eta| < 2.5$

Selection CutFlow : Signal all m_X

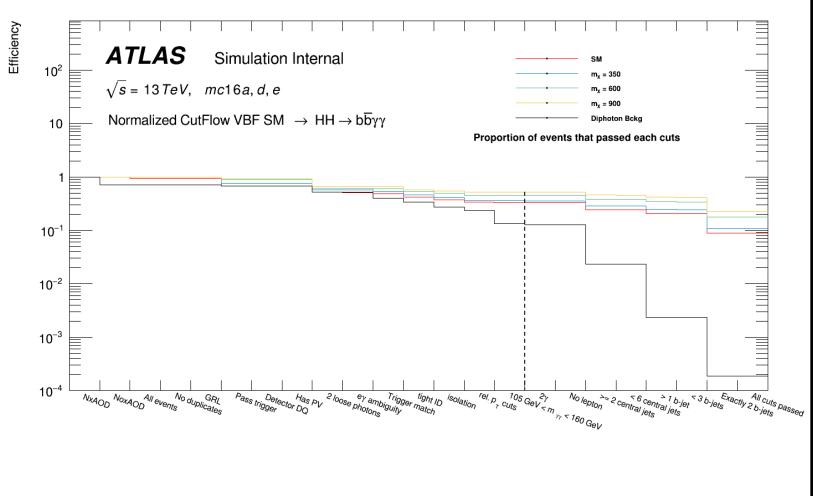


- Reconstruct the **CutFlow** applying **cuts after cuts** in order to **select signal-like events** which correspond to the final state $b\overline{b}\gamma\gamma$.
- CutFlow **renormalised** divided into two parts, the $\mathbf{H} \rightarrow \gamma \gamma$ cuts (selecting good data taking conditions and 2 good photons) and the $\mathbf{HH} \rightarrow b\bar{b}\gamma\gamma$ (b-jet selection) ones.
- Development of an automatic algorithm that scan all the **cuts efficiency** (number of events after a cut / number of events before this cut) and select the ones that **remove the most data** to **plots them** in function of m_X (backup).



Selection : Signal and Background

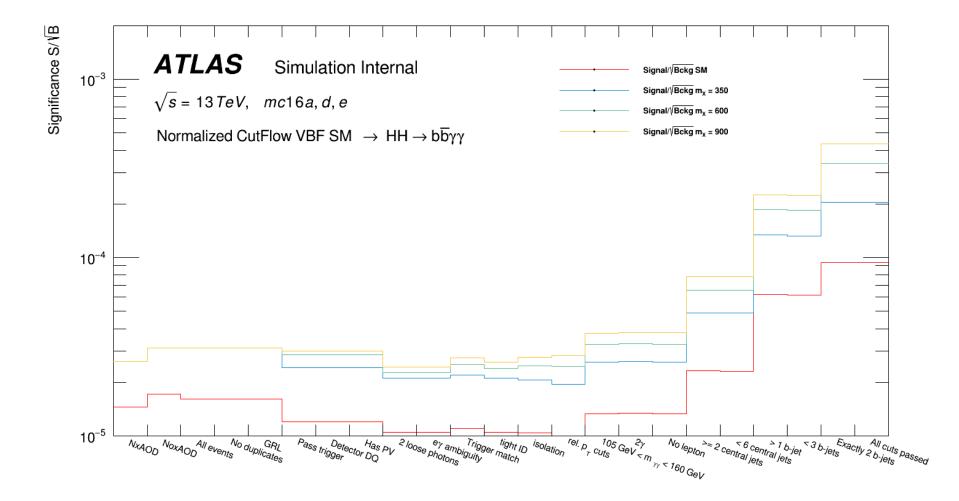


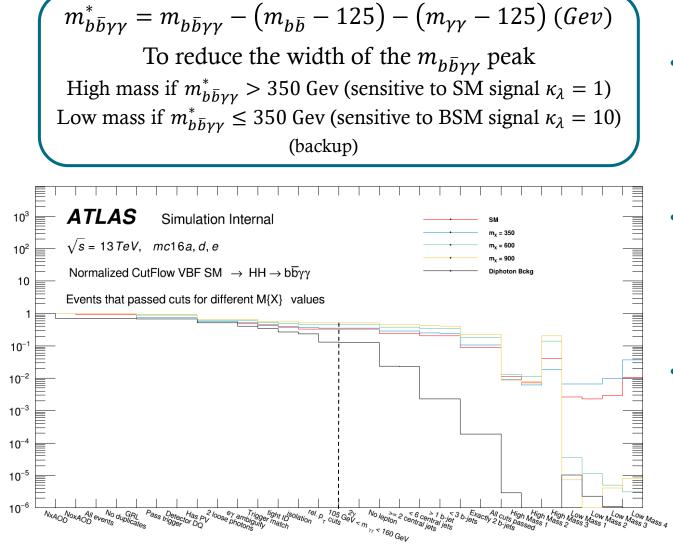


- Until now, we have only talked about the study of the signal so let's see how the background reacts to this selection.
- For that I decide to work with only the SM and 3 samples m_X and the diphoton background.
- As we were expecting the cuts are removing the background much more than the signal and $m_X \nearrow \Rightarrow$ Efficiency \nearrow (more boosted objects so better identification).
- Plot the **Significance** which is the Signal over the square root of the background $(\frac{S}{\sqrt{B}})$. Significance increasing cuts after cuts and $m_X \nearrow \Rightarrow$ Significance \nearrow .

Selection : Significance





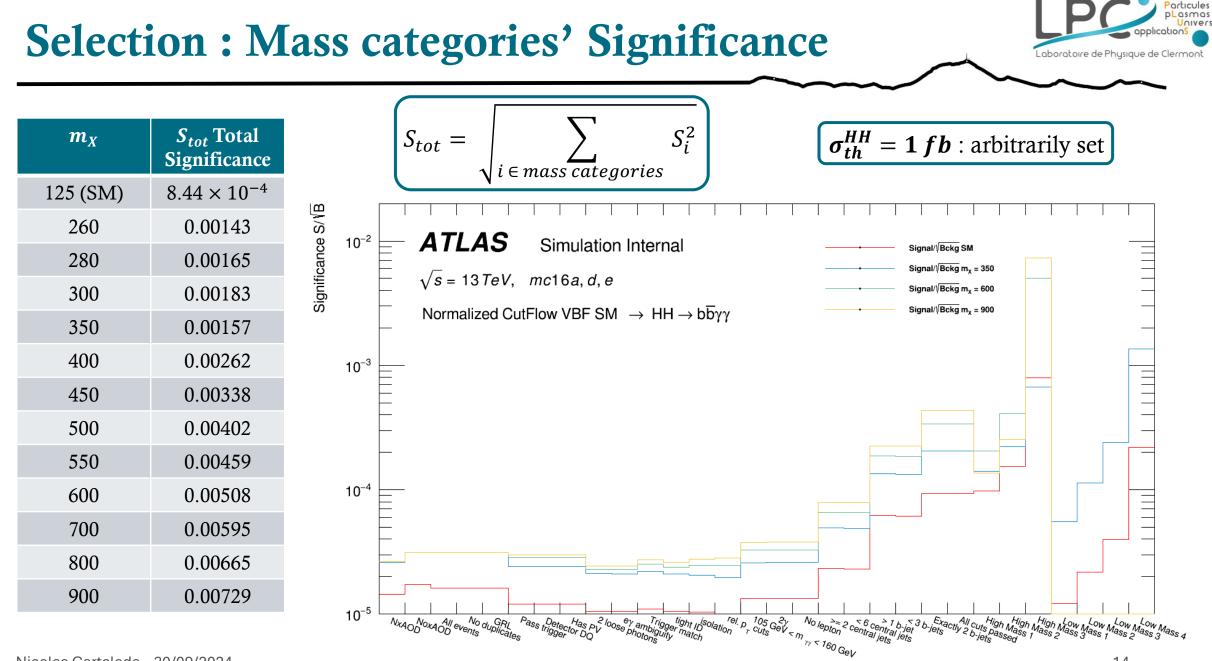


bbyy Selection : Mass categories



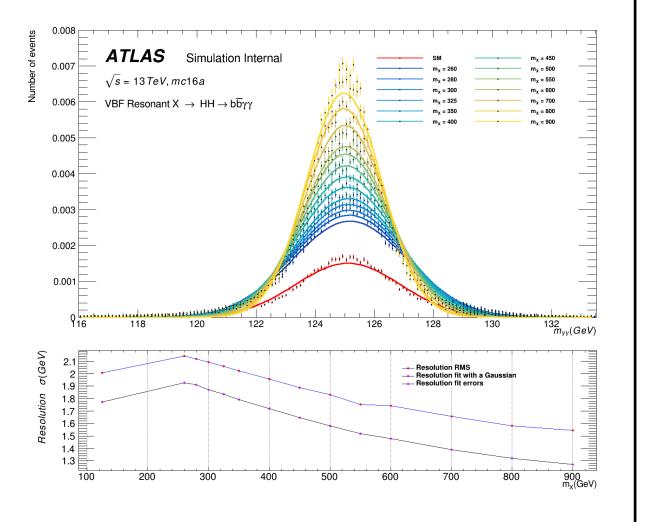
- The analysis follows the **selection methodology** from a previously published study on **nonresonant HH** $\rightarrow b\overline{b}\gamma\gamma$ (JHEP 01 2024 066)
- Worked with a BDT (Boosted Decision Tree) algorithm and the mass categories created thanks to BDT score ∈ [0, 1[(signal purity indicator).
- 3 High Mass categories and 4 Low Mass ones defined to optimize the significance considering the various possible κ_λ values and designed to improve the selection of events

Efficiency



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Modeling Resonant signal

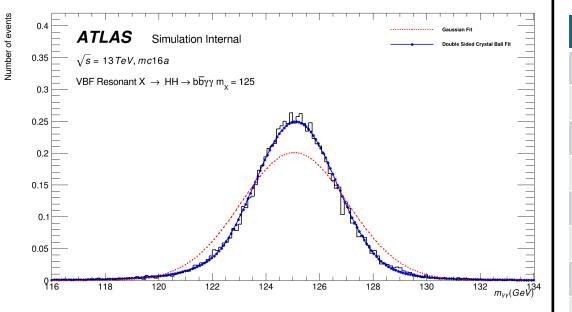


- On the upper plot we can see that when $m_X \nearrow \Rightarrow$ *Number of events* \nearrow (more **boosted objects so** better identification)
- Code to plot the $m_{\gamma\gamma}$ branch applying **pre-selection** and **weights** for different Monte-Carlo generated root files with $m_X \in [260, 900]$ *GeV*
- Graphs of resolution σ of gaussian which fit the data but this function is not optimal as we can see for the peak and the tails \rightarrow Double Sided Crystal Ball function and applied to both a χ^2 test to compare them

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Goodness of fit





- Fit the data with a DoubleSidedCrystalBall (DSCB) function instead of fit it with a gaussian to compare their goodness of fit (ChiSquare/ndof)
- As we can see in this table the **goodness of fit of the DSCB** function is quite **better than the gaussian** one, thanks to the **adaptable tails** of the DSCB function

(backup)

m_X	Gaussian $\chi^2/ndof$	DSCB $\chi^2/ndof$
125 (SM)	55,7	7,22
260	76,3	4,41
280	94,7	4,94
300	84,9	5,14
325	102,5	4,95
350	105,7	5,18
400	95,2	5,02
450	101,6	6,05
500	113,2	5,88
550	104,3	7,23
600	110	6,82
700	133,7	7,61
800	113,9	8,04
900	127,9	7,99



Different types of background

- **Continuum diphoton** production with jets ($\gamma\gamma$ + jets)
 - > $t\bar{t}\gamma\gamma$ production, where top quark pairs are produced alongside two photons
- **Single and pair Higgs boson** production through : ggF, VBF, ttH, bbH

Fit functions of the continuous background

• **Decreasing Exponential** Function

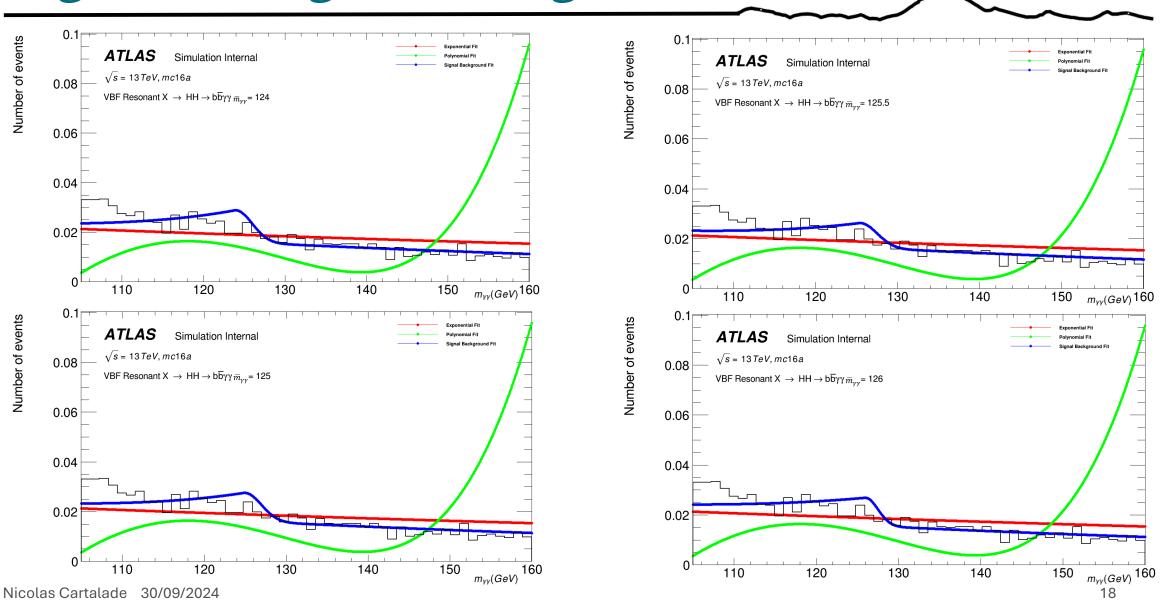
 $f_{bckg}(x) = Ae^{-Bx}$

where *A* and *B* determined by fitting the background data in regions outside of the signal window $\in [120, 130]$ Gev (**blinding** the signal region)

• **Polynomial** Function

 $f_{bckg}(x) = a_0 + a_1x + a_2x^2 + \dots + a_nx^n$ where the coefficients $a_0, a_1, a_2, \dots, a_n$ determined through the fitting process. The choice of polynomial degree *n* is balanced to ensure sufficient **flexibility** without overfitting

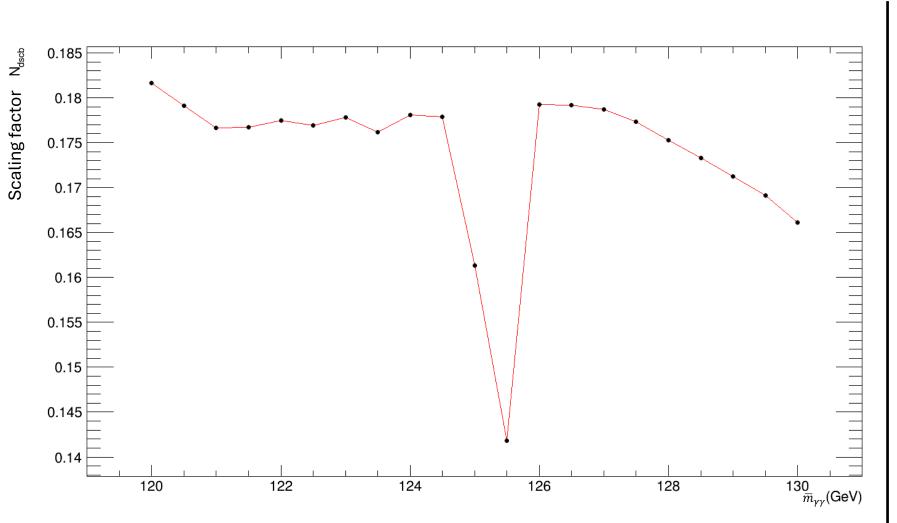
Signal + Bckg over Bckg



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Spurious Signal





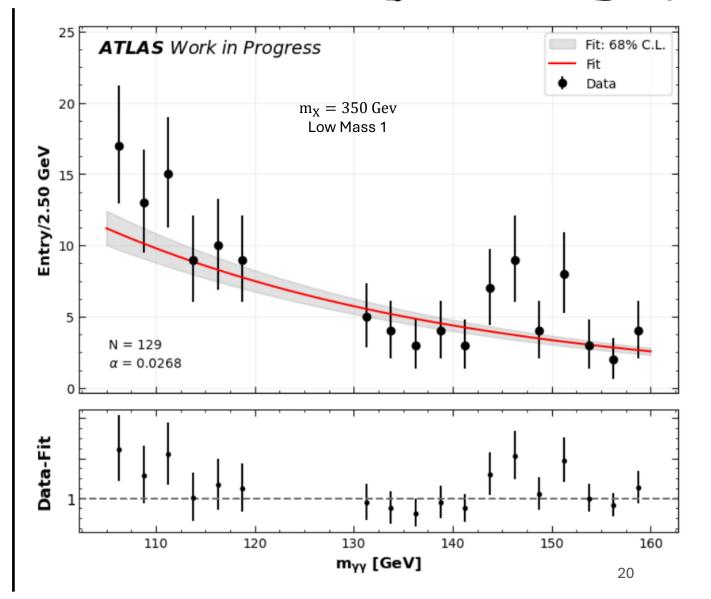
- $\overline{m}_{\gamma\gamma} \in [120, 130]$ *Gev* as a free parameter
- Spurious signal (N_{dscb}) : Maximum of 'fakesignal' find out when fitting signal +background function onto the background only sample.
- Nuisance parameter so we want to **find** the **function** fitting the background with the **lower spurious signal**

Data continuum background fit



• Finally using the **decreasing exponential** function to fit the continuous background component (often preferred for **smoothly falling distributions**)

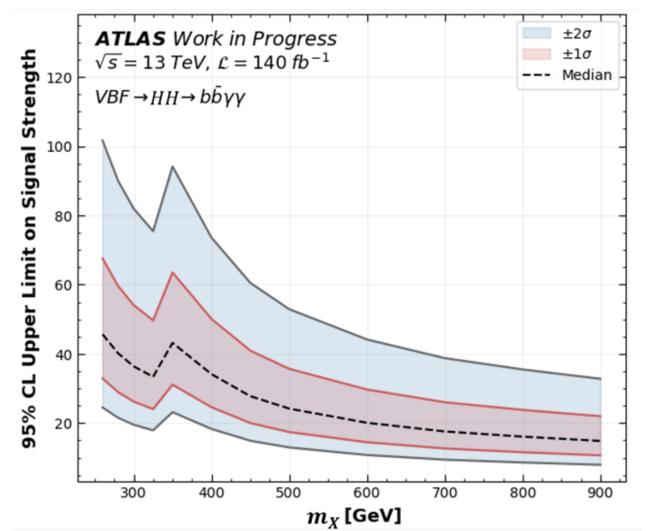
• As an **example**, we have on the right the **background fit of the data** for $m_X = 350$ Gev and for the Low Mass 1.



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Final Fit : signal strength

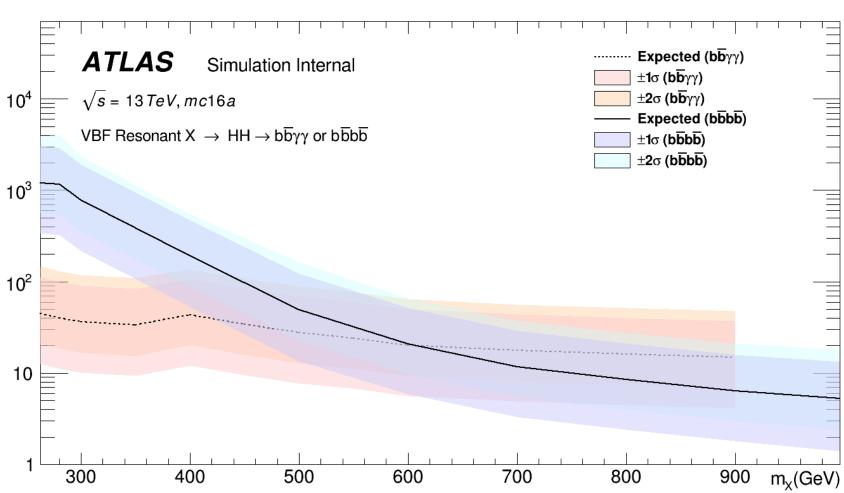
- Performed a **blinded** Asimov (data equal to expected value) **fit** using only **sideband data** to **avoid bias** from potential real signals.
- Compute at the end the **signal strength** : $\mu_{HH} = \frac{\sigma_{obs}^{HH} (upper limit)}{\sigma_{th}^{HH}}$
- Peak (350 GeV) could be explained from the choice of discrimination variable $(m_{b\overline{b}\gamma\gamma}^*)$ and the transition between Low and High Mass categories, redistributing signal events across more categories and reducing total significance for $m_X \sim 350$ Gev





Upper limit : Cross-section





- Upper limit cross-section for resonant HH model improved for m_X below 600 GeV thanks to this study focused on $b\overline{b}\gamma\gamma$ final state
- **Combining** $b\overline{b}\gamma\gamma$ and $b\overline{b}b\overline{b}$ channels enhances the upper limit and is crucial for constraining the resonant HH VBF process





- This internship represent the **first Resonant HH VBF** into $b\overline{b}\gamma\gamma$ analysis
- Used Double-Sided Crystal Ball (DSCB) for signal, exponential for background with the study of the spurious signal and the χ^2 function to optimize the fits
- Applied and study **Selection** and **Categorization** in order to optimize the **Significance**
- Improved upper limit on cross-section (for masses $m_X < 600$ GeV) of VBF Resonant HH process thanks to the $b\bar{b}\gamma\gamma$ final state





Current and following steps for the PhD :

Selection

- Adapt the BDT's structure and the definition of the mass categories to the resonant signal to improve cuts' efficiency and significance
- Explore neural networks for enhanced mass categorization and better signal-background separation

Modeling

• Future analyses should **incorporate systematic uncertainties** to provide more comprehensive and realistic results. While these had **minimal impact** in the **non-resonant analysis** due to **limited data**, they will become **crucial for improving fit accuracy as datasets grow**



Thank you for your attention

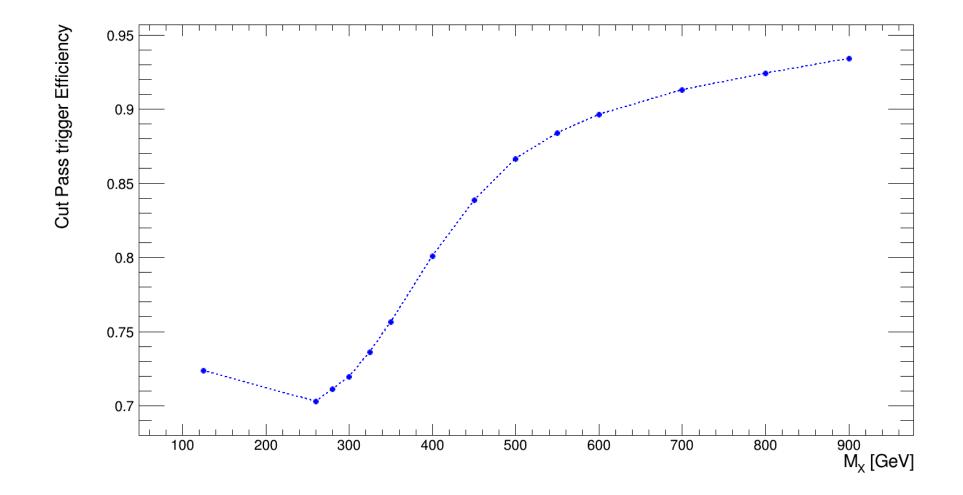
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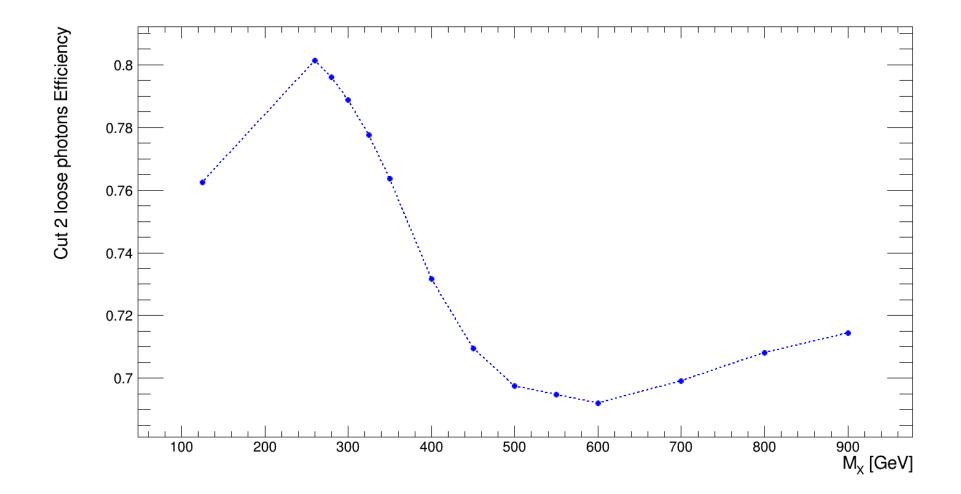
Backup slides

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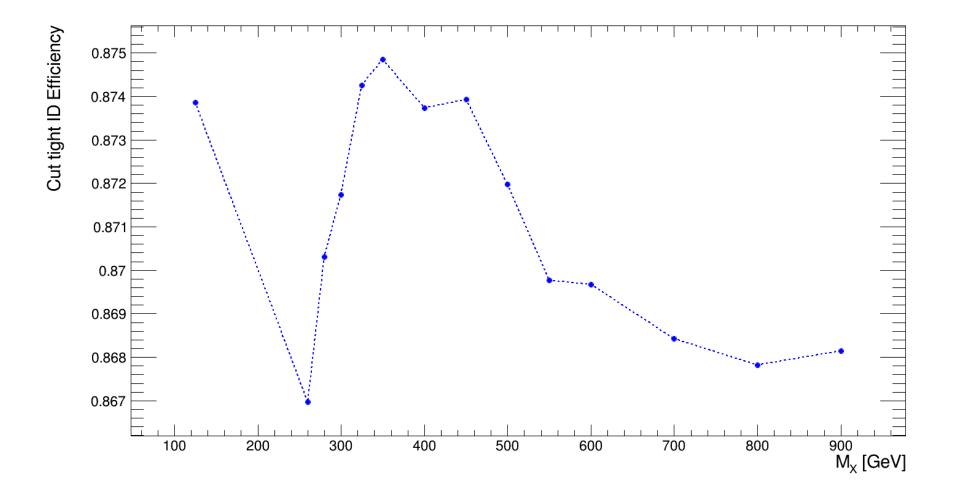




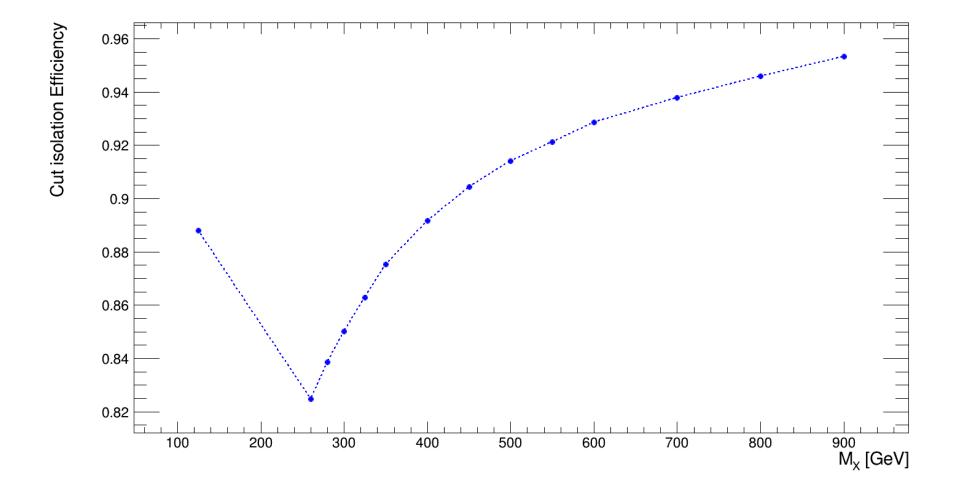




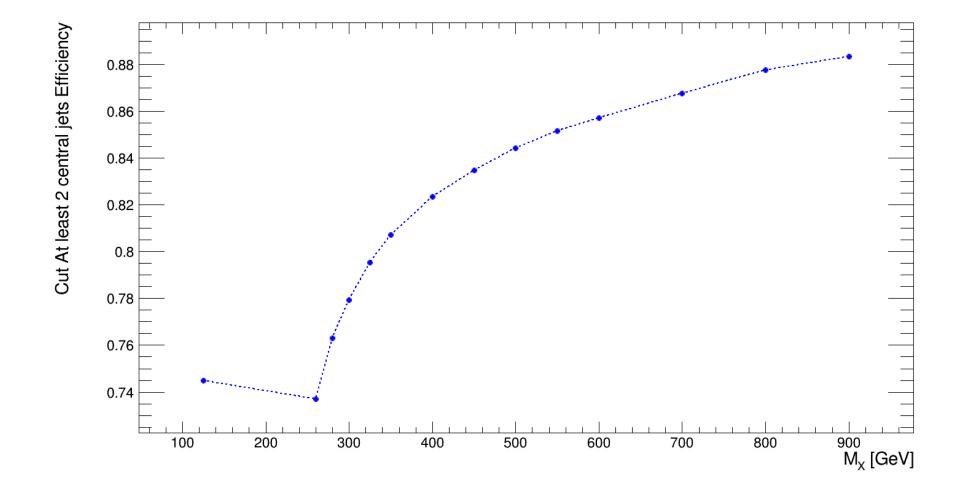




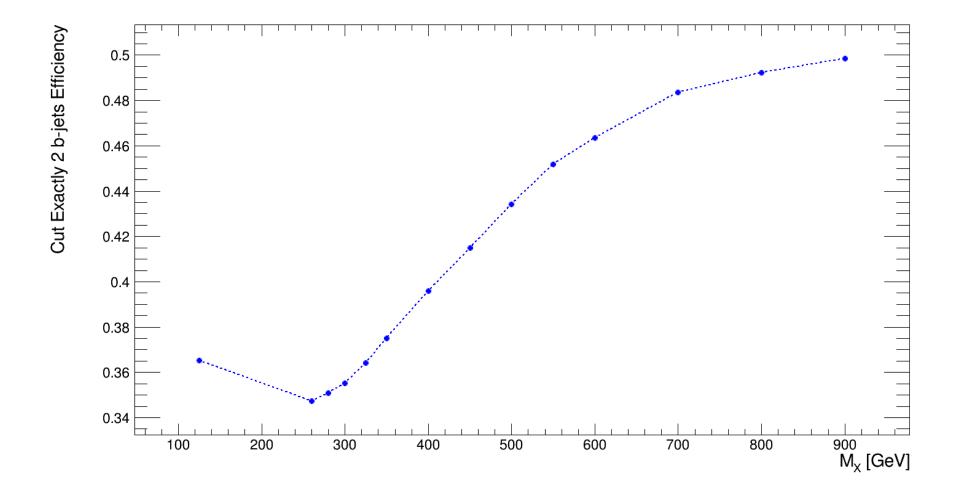






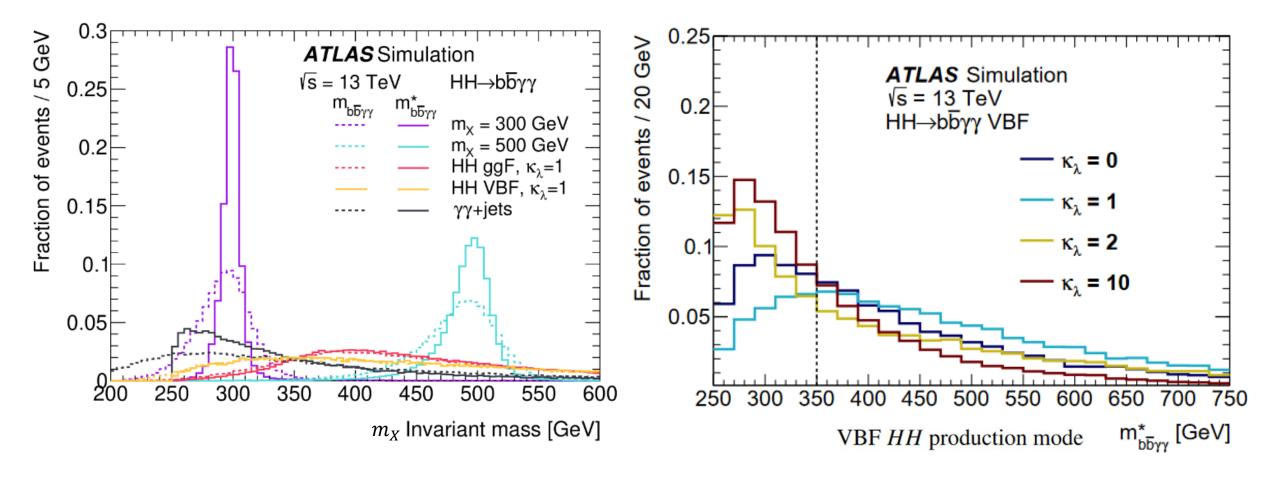


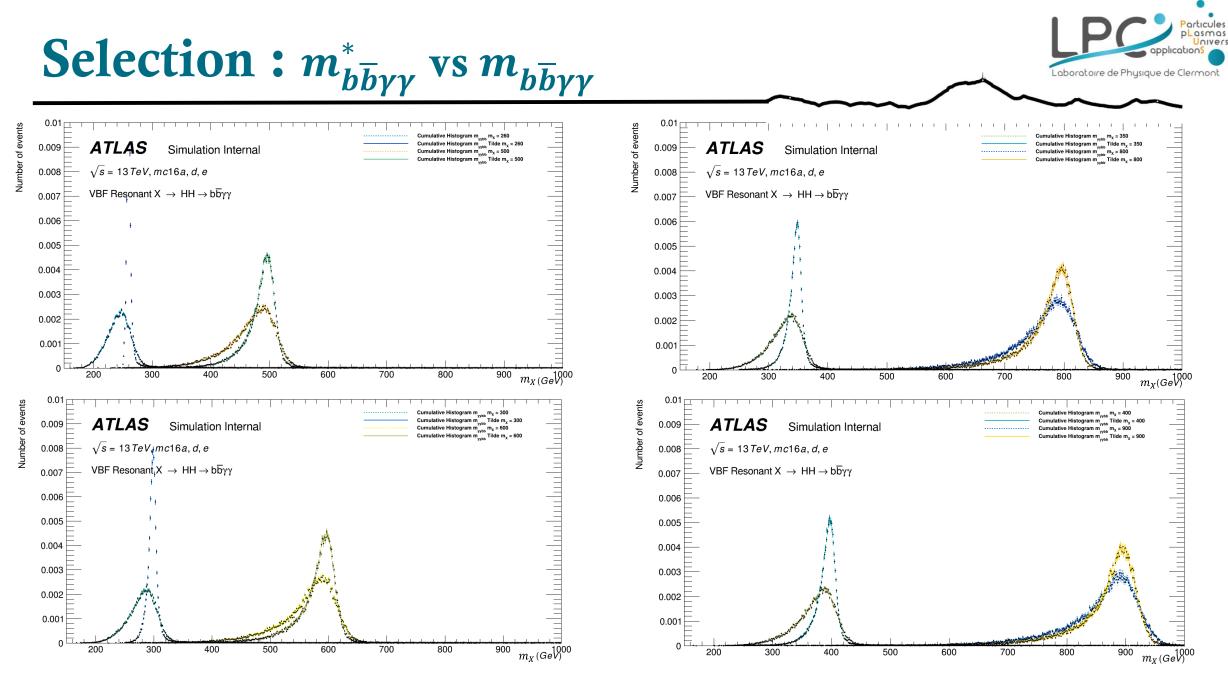




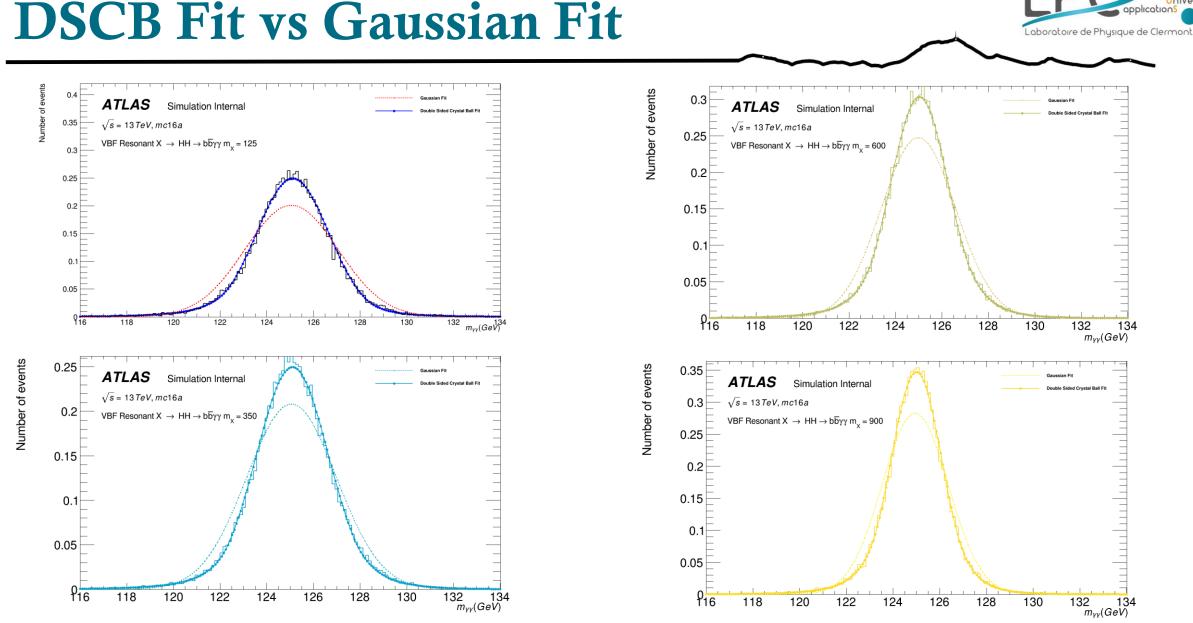
CutFlow : BDT mass categories







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Formulas



$${f Gaussian} ext{ function :} \ f(x) = rac{1}{\sigma\sqrt{2\,\pi}}\,{
m e}^{-rac{(x-\mu)^2}{2\sigma^2}}$$

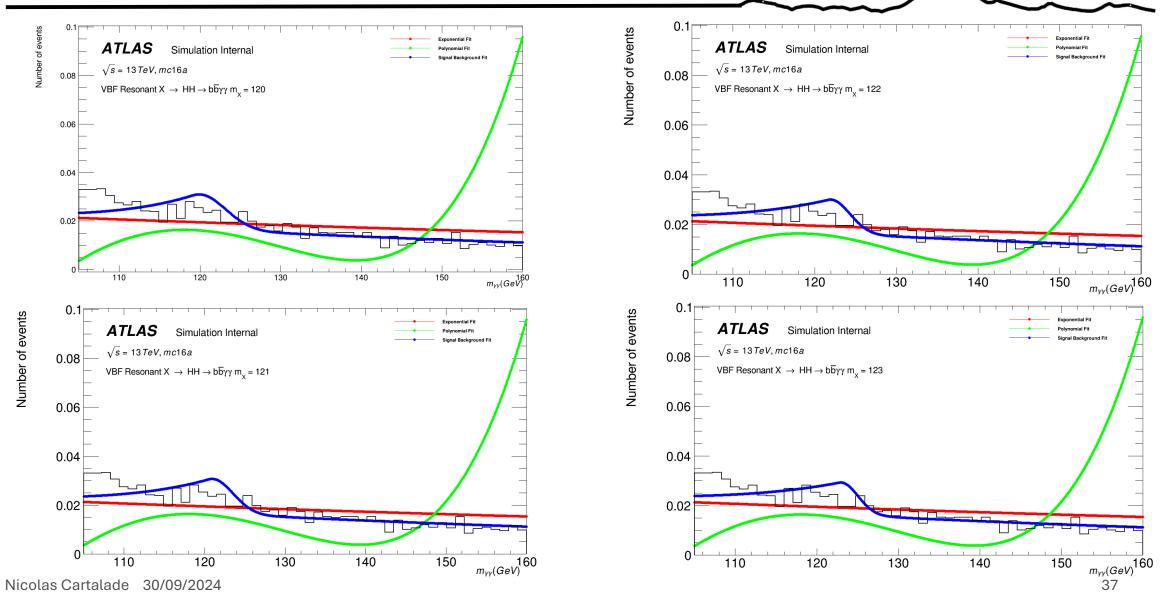
Exponential function :
$$f(x; \lambda) = \begin{cases} \lambda e^{-\lambda x} & x \geq 0, \\ 0 & x < 0. \end{cases}$$

 $\begin{array}{l} \textbf{Double Sided Crystal Ball function :} \\ f(m;m_0,\sigma,\alpha_L,n_L,\alpha_R,n_R) = \begin{cases} A_L \cdot (B_L - \frac{m-m_0}{\sigma_L})^{-n_L}, & \text{for } \frac{m-m_0}{\sigma_L} < -\alpha_L \\ \exp\left(-\frac{1}{2} \cdot \left[\frac{m-m_0}{\sigma_L}\right]^2\right), & \text{for } \frac{m-m_0}{\sigma_L} \leq 0 \\ \exp\left(-\frac{1}{2} \cdot \left[\frac{m-m_0}{\sigma_R}\right]^2\right), & \text{for } \frac{m-m_0}{\sigma_R} \leq \alpha_R \\ A_R \cdot (B_R + \frac{m-m_0}{\sigma_R})^{-n_R}, & \text{otherwise,} \end{cases} \\ \end{array}$ $\begin{array}{l} \textbf{where} \\ A_i = \left(\frac{n_i}{|\alpha_i|}\right)^{n_i} \cdot \exp\left(-\frac{|\alpha_i|^2}{2}\right) \end{cases}$

 $B_i = rac{n_i}{|lpha_i|} - |lpha_i|$

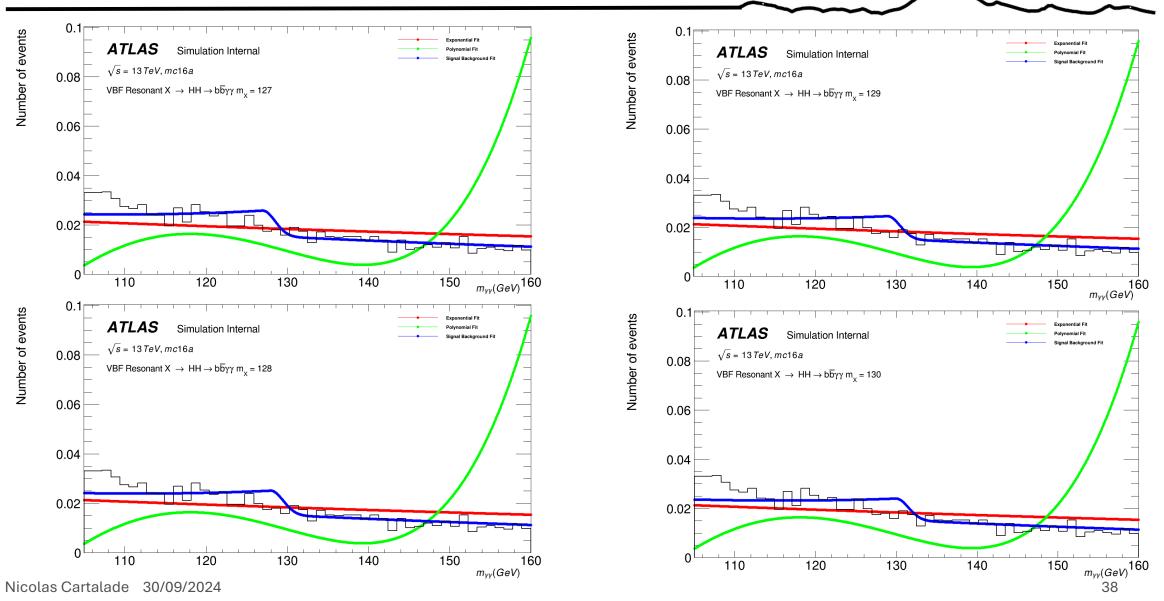
Signal and Bckg over Bckg





Signal and Bckg over Bckg









m_X	Cuts passed	High Mass 1	High Mass 2	High Mass 3	Low Mass 1	Low Mass 2	Low Mass 3	Low Mass 4
125 (SM)	0.08867	0.01151	0.00775	0.04049	0.00269	0.00230	0.00291	0.01076
260	0.08384	7.47×10^{-5}	1.60×10^{-5}	1.86×10^{-5}	0.01373	0.01097	0.01256	0.03752
280	0.08872	$1,31 \times 10^{-4}$	5.19×10^{-5}	3.28×10^{-5}	0.01202	0.01087	0.01355	0.04345
300	0.09271	$2,36 \times 10^{-4}$	6.60×10^{-5}	5.23×10^{-5}	0.01100	0.01028	0.01402	0.04873
350	0.10674	0.00916	0.00617	0.01884	0.00677	0.00659	0.00971	0.03715
400	0.09920	0.01846	0.01449	0.07197	6.59×10^{-4}	3.26×10^{-4}	3.41×10^{-4}	$4.50 imes 10^{-4}$
450	0.13597	0.01593	0.01309	0.09423	1.88×10^{-4}	6.24×10^{-5}	5.07×10^{-5}	4.80×10^{-5}
500	0.15060	0.01486	0.01218	0.11257	1.27×10^{-4}	2.50×10^{-5}	1.97×10^{-5}	1.52×10^{-5}
550	0.16489	0.01420	0.01155	0.12877	7.73×10^{-5}	1.58×10^{-5}	1.36×10^{-5}	1.72×10^{-5}
600	0.17644	0.01327	0.01141	0.14253	3.53×10^{-5}	1.16×10^{-5}	4.96×10^{-6}	3.08×10^{-6}
700	0.19796	0.01209	0.01030	0.16749	2.39×10^{-5}	0	2.39×10^{-6}	5.67×10^{-6}
800	0.21402	0.01071	0.00904	0.18737	2.04×10^{-5}	0	0	0
900	0.22720	0.00886	0.00705	0.20541	7.28×10^{-6}	0	4.05×10^{-6}	8.07×10^{-6}
Background	1.87×10^{-4}	2.90×10^{-6}	5.25×10^{-7}	5.48×10^{-7}	1.03×10^{-5}	2.33×10^{-6}	1.12×10^{-6}	5.10×10^{-6}





m_X	Cuts passed	High Mass 1	High Mass 2	High Mass 3	Low Mass 1	Low Mass 2	Low Mass 3	Low Mass 4
125 (SM)	9.39×10^{-5}	9.79×10^{-5}	1.55×10^{-4}	7.93×10^{-4}	1.21×10^{-5}	2.18×10^{-5}	3.98×10^{-5}	2.19×10^{-4}
260	1.61×10^{-4}	1.15×10^{-6}	5.81×10^{-7}	6.62×10^{-7}	1.12×10^{-4}	1.89×10^{-4}	3.11×10^{-4}	0.00138
280	1.70×10^{-4}	2.01×10^{-6}	1.88×10^{-6}	1.16×10^{-6}	9.81×10^{-5}	1.87×10^{-4}	3.35×10^{-4}	0.00160
300	1.78×10^{-4}	3.63×10^{-6}	2.39×10^{-6}	1.86×10^{-6}	8.98×10^{-5}	1.77×10^{-4}	3.47×10^{-4}	0.00179
350	2.05×10^{-4}	1.41×10^{-4}	2.23×10^{-4}	6.68×10^{-4}	5.52×10^{-5}	1.13×10^{-4}	2.40×10^{-4}	0.00137
400	2.32×10^{-4}	2.84×10^{-4}	5.25×10^{-4}	0.00255	5.38×10^{-6}	5.60×10^{-6}	8.45×10^{-6}	1.65×10^{-5}
450	2.61×10^{-4}	2.46×10^{-4}	4.74×10^{-4}	0.00334	1.54×10^{-6}	1.07×10^{-6}	1.26×10^{-6}	1.76×10^{-6}
500	2.88×10^{-4}	2.29×10^{-4}	4.41×10^{-4}	0.00399	1.04×10^{-6}	4.31×10^{-7}	4.88×10^{-7}	5.58×10^{-7}
550	3.16×10^{-4}	2.19×10^{-4}	4.18×10^{-4}	0.00456	6.31×10^{-7}	2.71×10^{-7}	3.38×10^{-7}	6.35×10^{-7}
600	3.38×10^{-4}	2.04×10^{-4}	4.13×10^{-4}	0.00505	2.88×10^{-7}	2.00×10^{-7}	1.23×10^{-7}	1.13×10^{-7}
700	3.80×10^{-4}	1.86×10^{-4}	3.73×10^{-4}	0.00594	1.95×10^{-7}	0	5.92×10^{-8}	2.08×10^{-7}
800	4.10×10^{-4}	1.65×10^{-4}	3.27×10^{-4}	0.00664	1.67×10^{-7}	0	0	0
900	4.36×10^{-4}	1.37×10^{-4}	2.55×10^{-4}	0.00728	5.94×10^{-8}	0	1.00×10^{-7}	2.97×10^{-7}

Opportunities aboratoire de Physique de Clermont Internship Work: PhD work: PhD Work: Maintenance RUN 3 RUN 2 HL-LHC and Upgrades 2015 2018 2022 2025 2040 Instrumentation on the High Granularity Time Fusion Gluon-Gluon (ggF) : **Detector (HGTD)** : 90% de la production HH : SM t/b/c : BSM $\sigma_{HH}^{ggF} = 31.02 \, fb$ $\sigma_{HH}^{VBF} = 1.72 \ fb$ t/b/c Xt/b/c g g**Fusion de Bosons Vecteurs HH** (V = W or Z) First study of VBF Resonant signals HH $\rightarrow b\overline{b}\gamma\gamma$: Complementarity and combinations of results $HH \rightarrow b\bar{b}b\bar{b}$ **VBF quarks** forward, to **differentiate** them from **b** Particle **identification** and **differentiation** using quarks from $H \rightarrow b\bar{b}$: HGTD detection time with a temporal resolution \leq 50 ps