X-ray constraints on Dark Matter

Jordan Koechler

New postdoc!

INFN Welcome Day

November 3rd, 2024

INFN



- Step 1: Home region near the German/Switzerland border of France
- Did my first two years of B.Sc. there







• Step 2: Finished the B.Sc. and M.Sc. in Montpellier







• Step 3: M.Sc. thesis in Lyon with Giacomo Cacciapaglia @ IP2I









• Step 4: Ph.D. in Paris with Marco Cirelli @ LPTHE (Sorbonne)









• Step 5: Here I am in Turin! Postdoc with Mattia di Mauro









- You could probably see me running near the Po river most days of the week.
- A bit of a foodie.
- Big fan of video games. Enjoying board games too.
- Would partake in a book club if some people are into reading.
- Really into visiting museums too...

About this talk

- Dark matter 101
- Constraining dark matter candidates from diffuse X-rays
 - Sub-GeV DM
 - Primordial black holes

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- Evidence 1: Flat spiral galaxy rotation curves
- Newton's second law: $v_{circ}(r) \propto 1/\sqrt{r}$ for $r \to \infty$
- Observed: $v_{circ}(r) = const$
- New matter type to explain the discrepancy?



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 - In pink: X-ray mapping of the intracluster gas



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 - In pink: X-ray mapping of the intracluster gas
 - In blue: weak lensing mapping of the clusters' mass



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- Evidence 2: Extra invisible mass in galaxy clusters
- Example: The Bullet Cluster
 - In pink: X-ray mapping of the intracluster gas
 - In blue: weak lensing mapping of the clusters' mass
- Most of the the clusters' mass is not composed of gas



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$$T(\theta,\phi) = T_0 \sum_{\ell=0}^{+\infty} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\theta,\phi)$$
$$T_0 = 2.725 \text{ K}$$
$$C_\ell = \langle a_{\ell m} a_{\ell m}^* \rangle = \frac{1}{2\ell+1} \sum_{m=-\ell}^{\ell} |a_{\ell m}|^2$$





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• Evidence 3: Temperature anisotropies of the Cosmic Microwave Background





[Aghanim et al., A&A 641 (2020) A1]

A Universe filled with baryonic matter, dark matter, radiation and dark energy provides an excellent fit to the measurements!

Measurements of the CMB temperature anisotropies (+ other things) provide



• Known DM properties:

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Spans over 90 orders of magnitude!

A way to potentially probe DM: indirect detection





Candidate 1: DM as sub-GeV elementary particles



- Candidate 1: DM as sub-GeV elementary particles
- Candidate 2: DM as primordial black holes



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 <u>Solar screening</u> suppresses the flux
- What to do?
 - Look at Voyager 1 & 2 data!



[Boudaud, Lavalle & Salati, *Phys.Rev.Lett.* 119 (2017) 2, 021103]

- <u>Issue 1</u>: when DM produces e[±]
 <u>Solar screening</u> suppresses the flux
- What to do?
 - Look at Voyager 1 & 2 data!
 - Look for DM-produced γ -rays



[Boudaud, Lavalle & Salati, *Phys.Rev.Lett.* 119 (2017) 2, 021103]

 <u>Issue 2</u>: when DM produces γ
 <u>No sensitive enough observatories</u> at MeV energies



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No sensitive enough observatories at MeV energies

 Secondary emissions allow to circumvent the issue → study X-rays signals from light DM

[Cirelli et al., *Phys.Rev.D* 103 (2021) 6, 063022]

Adapted from [De Angelis et al. *Exper.Astron.* 44 (2017) 1, 25]



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Inverse-Compton scattering!



Adapted from [De Angelis et al. *Exper.Astron.* 44 (2017) 1, 25]





Local number density of DM-produced e^\pm

Local number density of ambient photons

DM

DM

Local number density of DM-produced e^\pm



Local number density of DM-produced e^{\pm}



Local number density of DM-produced e^\pm

Local number density of DM-produced e^{\pm}

$$\vec{\nabla} \left(D \, \vec{\nabla} f_{e^{\pm}} - \vec{v}_{c} f_{e^{\pm}} \right) + \frac{\partial}{\partial K_{e}} \left(b_{loss} f_{e^{\pm}} + \beta^{2} D_{pp} \frac{\partial f_{e^{\pm}}}{\partial K_{e}} \right) + Q_{e^{\pm}}^{DM} = 0$$
spatial
convection
energy loss
momentum space
source

diffusion convection energy loss momentum space sou

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spatial diffusion

energy loss momentum space source diffusion

Solve this equation by using DRAGON2

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source

Solve this equation by using DRAGON2

$$D = D_0 \beta^{\eta} \frac{(R/R_0)^{\delta}}{[1 + (R/R_0)^{\Delta\delta/s}]^s} \qquad D_{pp} = \frac{4}{3} \frac{1}{\delta(4 - \delta^2)(4 - \delta)} \frac{v_A^2 p^2}{D}$$

Transport parameters (D_0 , η , R_0 , δ , $\Delta\delta$, s, v_c , v_A , L) are set using CR fits



https://github.com/bsafdi/XMM_BSO_DATA

Datasets + Instrument response functions

Conservative approach



• Impose a 2σ bound whenever $\chi^2_{>}(p, m_{DM}) \ge 4$









Halo height	H	$8.00^{+2.35}_{-1.96}~{ m kpc}$
Norm. of Diffusion coeff.	D_0	$1.02^{+0.12}_{-0.10}\times10^{29}~{\rm cm^{2}s^{-1}}$
Norm. rigidity	R_0	$4 \mathrm{GV}$
Diffusion spectral index	δ	0.49 ± 0.01
$\beta \ {\rm exponent}$	η	$-0.75\substack{+0.06 \\ -0.07}$
Alfvén velocity	v_A	$13.40^{+0.96}_{-1.02}~{ m km/s}$
Break rigidity	R_b	$312\pm31~{\rm GV}$
Index break	$\Delta\delta$	0.20 ± 0.03
Smooth. param.	s	0.04 ± 0.0015

DM profiles:

- NFW
- Burkert
- cNFW with $\gamma = 1.26$

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[Hawking, Commun. Math. Phys. 43 (1975) 199]

$$T = \frac{1}{4\pi GM} \frac{\sqrt{1 - a^{\star^2}}}{1 + \sqrt{1 - a^{\star^2}}} \qquad a^{\star} = J/(GM)^2$$
$$\frac{d^2 N_i}{dt dE_i} = \frac{1}{2\pi} \sum_{\text{d.o.f.}} \frac{\Gamma_i(E_i, M, a^{\star})}{e^{E'_i/T} \pm 1}$$
$$E'_i = E_i - m\Omega \qquad \Omega = \frac{a^{\star}}{2GM \left(1 + \sqrt{1 - a^{\star^2}}\right)}$$



[De la Torre Luque, **JK** & Balaji, arXiv:2406.11949 (accepted in *Phys.Rev.D)*] BlackHawk [Arbey, Auffinger, *Eur.Phys.J.C* 79 (2019) 8, 693], [Arbey, Auffinger, *Eur.Phys.J.C* 81 (2021) 10]





Propagation of DM-produced e^{\pm}



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Conclusion

- Considering secondary emissions can help us circumventing the MeV gap in γ -ray observatories
- As a bonus: it has a great constraining power
- Robustness is debatable, due to numerous astro uncertainties
- Possible improvement: Astrophysical background modeling

Thank you for your attention!