



SPACECRAFT CHARGING: INCOMING AND OUTGOUNG ELECTRONS

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OUTLINE



- Introduction to Spacecraft Charging
- Incoming currents
- Outgoing currents
- Critical temperature: Theory and observations
- Outgoing electron yields depend on surface conditions
- Conclusions



Spacecraft Charging Adverse Effects



Spacecraft Charging is Harmful to the Health of Onboard Electronics

Charging affects

- Scientific measurements
- Telemetry signals
- Electronic communications





Discharges degrade

- Solar cells
- Controls
- Navigation





Spacecraft Charging Leading Cause of Spacecraft Failures



Missions Terminated *Due to the Space Environment*



Most Mission Terminated were due to Dielectric and Surface Electrostatic Discharges



Push and Pull of the Magnetosphere

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Local Time Dependence of Spacecraft Charging



Solar wind compresses the magnetosphere on the day side and elongates it on the night side. The magnetotail snaps back from to time. Energetic plasmas from the magnetotail may reach the geosynchronous altitudes at midnight. Electrons drift to the dawnside. As a result, **spacecraft charging occurs mostly during 0-6 AM**.



ELECTRON AND ION FLUXES AT GEOSYNCHRONOUS ALTITUDES



(Measurements on LANL-1984, March, Eclipse Periods, 2000)







• Consider a plasma in thermal equilibrium

$$\frac{1}{2}m_{e}v_{e}^{2} = \frac{1}{2}M_{i}V_{i}^{2}$$

• The electrons are much lighter and faster than the ions

 $v_e >> V_i$

• Therefore, the flux of electrons is much higher than that of ions

$$n_e q_e v_e >> n_i q_i V_i$$

- This is why spacecraft often charge to negative potentials in a plasma
- This is true not only in space but also in the laboratory.



$$f(E) = n(m/2\pi kT)^{3/2} \exp\left(-E/kT\right)$$

 $\log f(E) = C - E / kT$



- Shift of the energy distribution.
- The shift eφ gives the charging potential φ.

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ECLIPSE CHARGING OF LANL-97A TO -3 kV

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Charging of DSCS Geosynchronous Satellite





Ambient Ion Flux Spectrum measured on DSCS. Note the ion energy at 2-4 kV during the period of 11000-15000 UT.

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Lai [1999]



AMBIENT

IONS





CURRENT

 $I_e - I_i = I_s + I_b + I_{ph}$

 $\int_0^\infty dEE f_e(E) \exp\left(-q_e \phi/kT_e\right) - \int_0^\infty dEE f_i(E) \left(1 - q_i \phi/kT_i\right)^\alpha = \int_0^\infty dEE f_e(E) \left[\delta(E) + \eta(E)\right] + I_{ph}$





- Incoming currents can be measured onboard.
- Outgoing currents can hardly be measured.
- To calculate outgoing currents, one needs the yield functions δ(E), η(E), and Y(ω) for secondary electrons, backscattered electrons, and photoelectrons, respectively.
- Traditionally, one adopts the yield functions from some standard sources, such as materials encyclopedia.
- The functions depend very much on surface conditions!



ELECTRON INDUCED ELECTRONS FROM A SPACECRAFT SURFACE



• Typical secondary electron yield $\delta(E)$ and backscattered electron yield $\eta(E)$.

• $\delta(E)$ features a maximum at an energy range of 60 - 2000 eV for typical materials. Beyond the max, it falls monotonically.

• $\eta(E)$ is much smaller and flatter compared with $\delta(E)$.

- The sum, $\delta(E) + \eta(E)$, is the total number of outgoing electrons for an incoming electron of energy E.
- Recent ly (Cimino, 2004) : η(E)
 rises at below about 20-50 eV and
 reaches unity at E = 0.



ELECTRON DISTRIBUTION AND SECONDARY EMISSION YIELD





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Fig.2 (Upper) Temperature T of a Maxwellian electron distribution. (Lower) Secondary electron coefficient & as a function of primary electron energy. • Imagine 2 camps of electrons.

• The low energy camp favors charging to positive voltage, albeit a few V only.

- The high energy camp favors charging to negative voltage, which can be -kV.
- As temperature increases, there must exist a critical temperature T*, above which the high energy camp wins.
- Above T*, negative charging occurs.



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- For studying charging onset, one can ignore the ions because the ion flux is two orders of magnitude lower than that of electrons.
- The Maxwellian distribution function f(E) is a function of E, n, and T $f(E) = n(m/2\pi kT)^{3/2} \exp(-E/kT)$

Substitute f(E) into the following equation of definite integrals: $\int_{0}^{\infty} dEE f(E) = \int_{0}^{\infty} dEE \left[\delta(E) + \eta(E)\right] f(E)$

- After the integration over *E*, the equation is a function of *n* and *T* only.
- Since *n* is a multiplicative factor on both sides of the equation, *n* cancels out.
- Therefore, the equation is a function of *T* only. The value *T** of *T* satisfying this equation is the critical temperature.



• A $\delta(E)$ formula [*Sanders and Inouye*, 1979] is as follows:

$$\delta(E) = c \left[\exp\left(-E/a\right) - \exp\left(-E/b\right) \right]$$

where a = 4.3 E_{max} , b = 0.367 E_{max} , and c = 1.37 δ_{max} . The two parameters E_{max} and δ_{max} depend on the surface material.

• An $\eta(E)$ formula [*Prokopenko and Laframboise*, 1980] is as follows:

$$\eta(E) = A - B\exp(-CE)$$

where A, B, and C depend on the surface material.

Massachusetts Institute of Technology Inset of Spacecraft Charging TABLE 1



CRITICAL TEMPERATURE T* (keV)		
MATERIAL	ISOTROPIC	NORMAL
Mg	0.4	
AL	0.6	
Kapton	0.8	0.5
Al Oxide	2.0	1.2
Teflon	2.1	1.4
Cu-Be	2.1	1.4
Glass	2.2	1.4
SiO ₂	2.6	1.7
Silver	2.7	1.2
Mg Oxide	3.6	2.5
Indium Oxide	3.6	2.0
Gold	4.9	2.9
Cu-Be (Activated)	5.3	3.7
MgF ₂	10.9	7.8



EVIDENCE OF CRITICAL TEMPERATURE FOR THE ONSET OF SPACECRAFT CHARGING



Rubin et al. [1980] St

EVIDENCE OF CRITICAL TEMPERATURE LANL-97A

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Which one is the best?



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CHARGING









TO STUDY THE BIG BANG OF THE UNIVERSE, HIGG'S BOSON, AND EXTRA DIMENSIONS.

The world's most expensive physics experiments.

The accelerator tubes are miles long.

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Requirement - Extremely accurate trajectories of charged particle beams !!







ACCELERATOR WALL

A CERN CONCERN: Can low energy electron clouds affect fast ion trajectories?





• The Furman formula is of the form:

$$\delta(E,\theta,s) = \delta_{\max}(\theta) \frac{s(E/E_{\max})}{s-1+(E/E_{\max})^{s}}$$
$$\delta_{\max}(\theta) = \delta_{\max}(0) \exp\left[0.5(1-\cos\theta)\right]$$
$$E_{\max}(\theta) = E_{\max}(0)\left[1+0.7(1-\cos\theta)\right]$$

where s, the surface condition parameter, equals 1.35 for the LHC copper samples.



A PEDAGOGICAL EXAMPLE

A saw-tooth surface has low secondary electron yield.













CHARGING OF MIRRORS

The photoelectron flux emitted from a surface is given by

 $J_{ph}(\omega,\alpha) = J_0(\omega)Y(\omega,\alpha) \left[1 - R(\omega,\alpha)\right]$

where J_0 is the incident light intensity at frequency ω and incidence angle α , $Y(\omega, \alpha)$ is the photoelectron yield per absorbed photon, and $R(\omega, \alpha)$ is the reflectance.

- If $R \rightarrow 1$, the photoelectron flux J_{ph} must be small.
- CONJECTURE:

A mirror should charge in sunlight as if it were in eclipse!

(Lai, JGR, 110, pp.1104-1115, 2005)

• One should do a laboratory experiment to prove or refute this conjecture!







MIRRORS FLANKING SOLAR CELLS





A Differential Charging Hazard









CONCLUSIONS



- SPACECRAFT CHARGING AFFECTS SCIENTIFIC MEASUREMENTS ONBOARD.
- CURRENT BALANCE GOVERNS THE SPACECRAFT POTENTIAL.
- CRITICAL TEMPERATURE T* CONTROLS THE ONSET OF SPACECRAFT CHARGING.
- INCOMING CURRENTS CAN BE MEASURED; OUTGOING CURRENTS ARE CALCULATED USING THE YIELD FUNCTIONS.
- THE YIELD FUNCTIONS DEPEND ON THE TYPE OF MATERIAL, PRIMARY ELECTRON ENERGY, INCIDENCE ANGLE, AND SURFACE CONDITIONS SUCH AS ROUGHNESS, DOSE, MONOLAYERS.
- **REFLECTANCE AFFECTS PHOTOELECTRON YIELD.**
- **REFLECTIVE SURFACES CHARGE IN SUNLIGHT AS IF IN ECLIPSE.**
- SOLAR PANEL FLANKED BY MIRRORS IS A HAZARDOUS SITUATION.



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The Los Alamos Magnetospheric Plasma Analyzer (MPA) measurements were obtained from the CDAWeb data service at NASA Goddard Space Flight Center.