## Dynamical Stabilization of the Fermi Scale

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LNF 2008





### Part A

Introducing Technicolor

Precision Data

Flavor versus Heavy Quark Masses

Walking Dynamics

Natural Dark Matter candidates

## Part B

Phase Diagrams of 4-Dimensional Gauge Theories

All orders (non)supersymmetric beta function

Ladder approximation

Lattice results

From Walking to Unparticle Physics

## Part C

### Minimal Walking Technicolor (MWT)

### The MWT Lagrangian

### "A Primer: LHC Phenomenology"

### Playing with Unification

# Part A

## Dynamical EW Breaking

$$L(H) \to -\frac{1}{4} F^{a\mu\nu} F^a_{\mu\nu} + i \bar{Q} \gamma^\mu D_\mu Q + \cdots$$

Dots are partially fixed by Anomalies as well as other principles

$$\dots \rightarrow L(\text{New SM Fermions})$$

## Technicolor

New Strong Interactions at ~ 250 GeV [Weinberg, Susskind]

Natural to use QCD-like dynamics.

 $SU(N)_{TC} \times SU(3)_C \times SU_L(2) \times U_Y(1)$ 

$$\langle Q^f \tilde{Q}_{f'} \rangle = \Lambda_{TC}^3 \qquad \qquad \Lambda_{TC} \simeq 250 \ GeV$$

## Electroweak Precision Measurements

Kennedy-Lynn, Peskin-Takeuchi, Altarelli-Barbieri, Bertolini- Sirlin, Marciano-Rosner



$$\Pi_{XY}^{\mu\nu}(q^2) = \Pi_{XY}(q^2)g^{\mu\nu} + \cdots$$



LEP EWG Summer 2006

## Large & Positive S from QCD-like TC

Peskin and Takeuchi, 90



# Masses to SM Fermions

#### The need to Extend Technicolor



## Different Approaches

#### Scalar-less New Gauge Interactions (Extended TC)

#### Marry SUSY and Technicolor

. . . . .

#### Add New Scalars in the Flavor Sector

#### Extended Technicolor





#### Beta - Terms



 $m_{\rm f} \approx \frac{g_{ETC}^2}{\Lambda_{ETC}^2} < \bar{Q}Q >_{ETC}$ 

Gamma - Terms

$$\frac{1}{\Lambda_{ETC}^2} (\bar{s}\gamma^5 d) (\bar{s}\gamma^5 d) + \frac{1}{\Lambda_{ETC}^2} (\bar{\mu}\gamma^5 e) (\bar{e}\gamma^5 e) + \dots$$

 $\Lambda_{ETC} \ge 10^3 \Lambda_{TC}$ 

Too small Top mass if

 $<\bar{Q}Q>_{ETC}\approx<\bar{Q}Q>_{TC}\sim\Lambda_{TC}^3$ 

# Walking versus Running

### Near Conformal Properties



## Why the walking can help?

$$\left\langle \bar{Q}Q_{ETC} \right\rangle = \exp\left(\int_{\Lambda_{TC}}^{\Lambda_{ETC}} d\ln(\mu) \ \gamma_m(\alpha(\mu))\right) \left\langle \bar{Q}Q_{TC} \right\rangle$$

#### <u>QCD-Like</u>

$$\exp\left(\int_{\Lambda_{TC}}^{\Lambda_{ETC}} d\ln(\mu) \, \gamma_m(\alpha(\mu))\right) \sim (\ln(\Lambda_{ETC}/\Lambda_{TC}))^{\gamma_m}$$

#### Near the conformal window

$$\exp\left(\int_{\Lambda_{TC}}^{\Lambda_{ETC}} d\ln(\mu) \ \gamma_m(\alpha(\mu))\right) \sim \left(\left(\Lambda_{ETC}/\Lambda_{TC}\right)^{\gamma_m(\alpha^{\star})}\right)$$

Fermion Mass Enhancement

$$m_{\rm f} \approx \frac{g_{ETC}^2}{\Lambda_{ETC}^2} < \bar{Q}Q >_{ETC} = \frac{g_{ETC}^2}{\Lambda_{ETC}^2} \left(\frac{\Lambda_{ETC}}{\Lambda_{TC}}\right)^{\gamma_m(\alpha^*)} < \bar{Q}Q >_{TC}$$

## S in Walking Technicolor

$$S_{\text{pert}} = \frac{1}{6\pi} \frac{N_f}{2} d(R)$$

$$S_{WTC} < S_{TC}$$

Appelquist, F.S.

However we will take:

 $S_{WTC} \approx S_{\text{pert}}$ 

## Besides



## S from extra Leptons

Fermions : 
$$\psi_L = \begin{pmatrix} \psi_{1L} \\ \psi_{2L} \end{pmatrix}$$
,  $\psi_{1R}$ ,  $\psi_{2R}$   
Hypercharge :  $Y$ ,  $Y + \frac{1}{2}$ ,  $Y - \frac{1}{2}$ 

$$S_{\text{Leptons}} = \frac{1}{6\pi} \left[ 1 - 2Y \ln\left(\frac{M_1}{M_2}\right)^2 + \frac{1 + 8Y}{20} \left(\frac{m_Z}{M_1}\right)^2 + \frac{1 - 8Y}{20} \left(\frac{m_Z}{M_2}\right)^2 + O\left(\frac{m_Z^4}{M_i^4}\right) \right]$$

$$M_{1,2}^2 \gg m_Z^2$$

# Rule:

Find Walking Theories with EW embedding minimizing S

### Identified Many EW Viable Walking and Custodial TC Models

Minimal Walking Technicolor (MWT)

Higher Dimensional Representations

**Beyond Minimal Walking Technicolor** 

Partially EW Gauged Technicolor

Split Technicolor

Additional Fermions in SM

Custodial TC

Similar to BESS models of Casalbuoni et al.

F.S. - Tuominen 04 Dietrich - F.S. - Tuominen 05 Dietrich - F.S. 06

Ryttov, F.S. 06

Appelquist, F.S. 98 Foadi, Frandsen, F.S. 07

# Dark Matter

 $\frac{\Omega_{DM}}{\Omega_B}\sim 5$ 

# What if DM is due to an asymmetry?

A particle similar to the nucleon

But electrically neutral

At most EW-type cross sections

Great if connected to the Electroweak Symmetry Breaking sector of the SM

# Technicolor is a natural example

Technibarion is similar to the nucleon TB number like the B number At most EW-type cross sections EW scale and interactions built in

> Nussinov, 86 Barr - Chivukula - Farhi 90 Gudnason - Kouvaris - F.S. 06

# Naive estimates

$$\frac{m_{TB}}{m_p} \approx \frac{1 \text{ TeV}}{1 \text{ GeV}} = 10^3$$

$$\frac{TB}{B} \propto \exp\left[-\frac{m_{TB}(T^*)}{T^*}\right] \sim 10^{-3} \qquad T^* \sim 200 \text{ GeV}$$

$$\frac{\Omega_{TB}}{\Omega_B} = \frac{TB}{B} \frac{m_{TB}}{m_p} \sim \mathcal{O}(1)$$

# Part B



## Progress in Strong Interactions

Phase Diagram of Gauge Theories

New Limits for Strongly Interacting Theories

All order beta function for nonsupersymmetric theories

## The 2 index Symmetric-Theory



Here Q and  $\widetilde{Q}$  are Weyl fermions.

The A-type is obtained by substituting  $\Box \Box$  with  $\Box$ .

#### Phase Diagram for the Symmetric Theory



F.S. - Tuominen 04

#### Using Ladder Approximation

Is this the minimal walking theory?

## Non-SUSY Phase Diagram for HDRs


## SUSY Phase Diagram for HDRs



Intriligator-Seiberg

### All orders beta function

$$\beta(g) = -\frac{g^3}{(4\pi)^2} \frac{\beta_0 - \frac{2}{3}T(r)N_f\gamma(g^2)}{1 - \frac{g^2}{8\pi^2}C_2(G)\left(1 + \frac{2\beta_0'}{\beta_0}\right)}$$

$$\gamma = -d\ln m/d\ln \mu$$

$$\beta_0 = \frac{11}{3}C_2(G) - \frac{4}{3}T(r)N_f$$

$$\beta_0' = C_2(G) - T(r)N_f$$

Ryttov and F.S. 07



### Bounds on the Conformal Window

# $\beta = 0 \qquad \longrightarrow \qquad \gamma = \frac{11C_2(G) - 4T(r)N_f}{2T(r)N_f}$

Unitarity of the Conformal Operators demands:

$$\gamma \le 2$$

Ryttov and F.S. 07

## Non-SUSY Phase Diagram Bound



### Landscape of the Unparticle World



Georgi 07

### Information on nonperturbative aspects of HDRs

# Part C

### The Minimal Walking Theory

2 Adj. Dirac Flavors of SU(2)

$$Q_L^a = \begin{pmatrix} U^a \\ D^a \end{pmatrix}_L, \quad U_R^a, \quad D_R^a \qquad a = 1, 2, 3$$

$$Y(Q_L) = \frac{y}{2}$$
  $Y(U_R, D_R) = \left(\frac{y+1}{2}, \frac{y-1}{2}\right)$ 

$$\mathcal{L}_{L} = \begin{pmatrix} N \\ E \end{pmatrix}_{L} \qquad N_{R} \qquad E_{R}$$
$$Y(\mathcal{L}_{L}) = -3\frac{y}{2} \qquad Y(N_{R}, E_{R}) = \left(\frac{-3y+1}{2}, \frac{-3y-1}{2}\right)$$

 $\mathcal{N} = 4$  super Yang-Mills

### The MWT Lagrangian

$$\mathcal{L}_{H} \rightarrow \left[ \frac{1}{4} \mathcal{F}^{a}_{\mu\nu} \mathcal{F}^{a\mu\nu} + i\bar{Q}_{L}\gamma^{\mu}D_{\mu}Q_{L} + i\bar{U}_{R}\gamma^{\mu}D_{\mu}U_{R} + i\bar{D}_{R}\gamma^{\mu}D_{\mu}D_{R} + i\bar{L}_{L}\gamma^{\mu}D_{\mu}L_{L} + i\bar{N}_{R}\gamma^{\mu}D_{\mu}N_{R} + i\bar{E}_{R}\gamma^{\mu}D_{\mu}E_{R} \right]$$

$$\mathcal{F}^{a}_{\mu\nu} = \partial_{\mu}\mathcal{A}^{a}_{\nu} - \partial_{\nu}\mathcal{A}^{a}_{\mu} + g_{TC}\epsilon^{abc}\mathcal{A}^{b}_{\mu}\mathcal{A}^{c}_{\nu} \qquad a, b, c = 1, \dots, 3$$

$$D_{\mu}Q_{L}^{a} = \left(\delta^{ac}\partial_{\mu} + g_{TC}\mathcal{A}^{b}_{\mu}\epsilon^{abc} - i\frac{g}{2}\vec{W}_{\mu}\cdot\vec{\tau}\delta^{ac} - ig'\frac{y}{2}B_{\mu}\delta^{ac}\right)Q_{L}^{c}$$

What you see is "not" what LHC will see

### LHC effective theory

 $\mathcal{L}(\text{Composites}) + \mathcal{L}(\text{Mixing with SM}) + \mathcal{L}(\text{New Leptons}) + \mathcal{L}(\text{SM} - \text{Higgs})$ 

#### Low Energy Scalar Sector

 $M_{ij} \sim Q_i^{\alpha} Q_j^{\beta} \varepsilon_{\alpha\beta}$ 



#### Low Energy Vector Sector

 $A_i^{\mu,j} \sim Q_i^{\alpha} \sigma^{\mu}_{\alpha\dot{\beta}} \bar{Q}^{\dot{\beta},j} - \frac{1}{4} \delta_i^j Q_k^{\alpha} \sigma^{\mu}_{\alpha\dot{\beta}} \bar{Q}^{\dot{\beta},k}$ 

$$A^{\mu} = \begin{pmatrix} \frac{a^{0\mu} + v^{0\mu} + v^{4\mu}}{2\sqrt{2}} & \frac{a^{+\mu} + v^{+\mu}}{2} \\ \frac{a^{-\mu} + v^{-\mu}}{2} & \frac{-a^{0\mu} - v^{0\mu} + v^{4\mu}}{2\sqrt{2}} \\ \frac{x^{\mu}_{UD} - s^{\mu}_{UD}}{2} & \frac{x^{\mu}_{DD}}{2} \\ \frac{x^{\mu}_{UD}}{\sqrt{2}} & \frac{x^{\mu}_{UD} - s^{\mu}_{UD}}{2} \\ \frac{x^{\mu}_{UD} + s^{\mu}_{UD}}{\sqrt{2}} & \frac{x^{\mu}_{DD} - s^{\mu}_{UD}}{2} \\ \frac{x^{\mu}_{UD} + s^{\mu}_{UD}}{2} & \frac{x^{\mu}_{DD}}{\sqrt{2}} \\ \frac{a^{+\mu} - v^{+\mu}}{2} & \frac{-a^{0\mu} + v^{0\mu} - v^{4\mu}}{2\sqrt{2}} \\ \end{pmatrix}$$

**Comprehensive Effective Technicolor Lagrangian** 

Vector Mesons

Yukawas

\*\* Link to MWT via Modified Weinberg Sum Rules \*\*

Written in a renormalizable form

With imposed constraints from Precision Data

A working technicolor benchmark

Foadi, Frandsen, Ryttov & F.S. 07

#### MWT versus EWPD



LEP EWG Summer 2006

 $100 \text{ Gev} < M_1 < 800 \text{ Gev}$ 

Y=-3/2 for Leptons

 $100 \text{ Gev} < M_2 < 1000 \text{ Gev}$ 

$$S_{\text{Leptons}} = \frac{1}{6\pi} \left[ 1 - 2Y \ln\left(\frac{M_1}{M_2}\right)^2 + \frac{1 + 8Y}{20} \left(\frac{m_Z}{M_1}\right)^2 + \frac{1 - 8Y}{20} \left(\frac{m_Z}{M_2}\right)^2 + O\left(\frac{m_Z^4}{M_i^4}\right) \right]$$

# LHC Phenomenology

In collaboration with

A. BelyaevR. FoadiM. T. FrandsenM. O. Jarvinen

A primer/work in progress

# Vector-Axial Mass Splitting



Within the LHC reach for small values of S if  $M_V > M_A$ 

## Vector-Axial Decay Widths



Vector wider than the axial 'cause of the decay V->2A

# Vector Resonances Branching Ratios



BRs for different Composite Higgs Masses BR in WW drops and rises (small S)

# Axial Resonances Branching Ratios



BRs for different Composite Higgs Masses

# Vector & Axial Direct production



# Associate Higgs Production



Vanishing for finite values of the axial mass: small S

# Playing with Unification

Farhi-Susskind, 79 Gudnason-Ryttov-F.S. 06 1 - loop running for SU(n)

$$\alpha_n^{-1}(\mu) = \alpha_n^{-1}(M_Z) - \frac{b_n}{2\pi} \ln\left(\frac{\mu}{M_Z}\right)$$

#### Degree of SU(3)xSU(2)xU(1) Unification

$$B_{Th} \leftarrow \frac{b_3 - b_2}{b_2 - b_1} = \frac{\alpha_3^{-1} - \alpha^{-1} \sin^2 \theta_w}{(1 + c^2)\alpha^{-1} \sin^2 \theta_w - c^2 \alpha^{-1}} \to B_{Exp}$$

Unification Scale

$$M_{GUT} = M_Z \exp\left[2\pi \frac{\alpha_2^{-1}(M_Z) - \alpha_1^{-1}(M_Z)}{b_2 - b_1}\right]$$



# Adding Adjoint SM Matter

### Minimal Walking Technicolor + SM Adjoint Matter

Colored Octet of Majorana Particles

Weak Triplet of Majorana Fermions

Extra SM Weyl Singlet

$$b_{3} = \frac{4}{3}N_{g} - 11 + 2$$

$$b_{2} = \frac{4}{3}(N_{g} + 1) - \frac{22}{3} + \frac{4}{3}$$

$$b_{1} = \frac{4}{3}(N_{g} + 1)$$

Gudnason-Ryttov-F.S. ph/0612230

Bajc-Senjanovic ph/0612029

 $B_{\rm Th} = 13/18 = 0.72(2)$  versus  $B_{\rm Exp} \simeq 0.72$ 



### What LHC can see and how may it deceive you

- Composite States: Technirho, composite (light) Higgs
- Detect Light Higgs: "Elementary or Composite ?"
- Study: pp -> HW, pp -> HZ. pp->HW (enhanced)
- 4th Family of Leptons no Quarks

### What LHC can see and how may it deceive you

- wino/bino/gluino-like produced. Have you seen SUSY, WTC or ... ?
- Study their couplings to SM fermions

## Unpleasant scenario:

• WW scattering is unitarized at the tree level up to 4 TeV with new vectors states of about 1.5 TeV. LHC and ILC will not dicover! (Foadi and FS 08)

#### Summary

- Introduced different types of viable technicolor theories
- Phase diagram of Higher Dimensional Representations
- Presented Minimal Walking Technicolor
- Benchmark for "any" serious model breaking dynamically the electroweak symmetry and passing LEP I & II precision tests.
- Dark Matter as a technibaryon
- Unification

### Unifying also TC Cartoon!



#### MWT versus EWPD



46 Gev <  $M_1 < 800$  Gev

Y=-1/2 for Leptons

 $100 \text{ Gev} < M_2 < 1000 \text{ Gev}$ 



#### DD - Technibaryon as DM



Amount of LTB dark matter as function of LTB mass with L' = 0, L = B

Gudnason, Kouvaris and F.S. 06

Constrained by CDMS to be at most only a 40-60% fraction of DM
# Cosmological Constant Problem

• Why is empty space so nearly empty?

$$\rho_{vac}{<}10^{\text{-}46}GeV\ ^{4}\approx 10^{\text{-}29}\ g\ cm^{\text{-}3}$$

• Standard Model sets the scale to:

 $\rho_{sm}$  >108GeV 4

• Mismatch by 54 order of magnitude!!



#### Dark Matter in MWT

a) Heavy Neutrino for y=1/3.

Kainulainen, Tuominen, Virkajärvi. 06

b) Dark Majorana

Kouvaris 07

#### TB Dark Matter Beyond MWT

SU(4), 8 flavors, only 2-flavors EW gauged

$$\epsilon_{t_1 t_2 t_3 t_4} Q_L^{t_1 t_2, f} Q_L^{t_3 t_4, f'} \epsilon_{ff'}$$

Dietrich-F.S. 06

Cannot be easily constrained by CDMS

## Higgsless versus Higgsfull



 $\frac{M_H}{M_V} > 1$ 



 $\frac{M_H}{M_V} \le 1$ 

#### Spectrum of Hadronic/Technihadronic States

Using 't Hooft Large N and Unitarity in Pion-Pion Scattering in QCD

Vector Meson is a quark-antiquark state:

 $\rho(770)$ 

Broad Sigma of multiquark nature

 $f_0(600)$ 

F.S. & Schechter, 95 Harada, F.S. and Schechter, 03 Caprini, Colangelo,Leutwyler 05 Maiani,Piccinini, Polosa, Riquer 04 F.S. and Schechter, 07

### Higgsless: 't Hooft Extension

$$M_{T\rho} = \frac{\sqrt{2}v_0}{F_{\pi}} \frac{\sqrt{3}}{\sqrt{N_D N_{TC}}} m_{\rho} \qquad v_0 \sim 250 \text{GeV}$$

$$M_{Tf_0} = \frac{\sqrt{2}v_0}{F_{\pi}\sqrt{N_D}} \left(\frac{N_{TC}}{\sqrt{3}}\right)^{\frac{p-1}{2}} m_{f_0} \qquad p \ge 1$$
ES.07
$$N_D = \frac{N_{TF}}{2} \qquad 1.4 \qquad M_{T\rho}(\text{TeV})$$

$$N_D = 2 \qquad 0.9 \qquad 0.8 \qquad 0.7 \qquad M_{Tf_0}(\text{TeV}) \qquad M_{Tf_0}(\text{TeV})$$

9

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$$\epsilon_{t_1 t_2 t_3 t_4} Q_L^{t_1 t_2, f} Q_L^{t_3 t_4, f'} \epsilon_{ff'}$$

Dietrich-F.S. 06

Cannot be easily constrained by CDMS

## S - T

S-measures the left - right type current correlator

$$S = -16\pi \frac{\Pi_{3Y}(m_Z^2) - \Pi_{3Y}(0)}{m_Z^2}$$

T-measures deviations from

$$m_{\rm W}^2 = \cos^2 \theta_{\rm W} \, m_{\rm Z}^2$$

$$T = 4\pi \frac{\Pi_{11}(0) - \Pi_{33}(0)}{s_W^2 c_W^2 m_Z^2}$$

# Conditions

Universe Electric Neutrality

Chemical Equilibrium

EW Sphaleron Processes, Kuzmin-Rubakov-Shaposhnikov

TB - B violated at High Energies & approx. conserved at EW

#### $M_H/M_V < 1$ in MWT theories

$$M_{T\rho} = \frac{\sqrt{2}v_0}{F_{\pi}} \frac{\sqrt{3}\sqrt{2}}{\sqrt{N_D N_{TC}(N_{TC} \mp 1)}} m_{\rho}$$

$$M_{Tf_0} = \frac{\sqrt{2}v_0}{F_{\pi}} \frac{\sqrt{3}\sqrt{2}}{\sqrt{N_D N_{TC}(N_{TC} \mp 1)}} m_{f_0}$$
E.S. 07

