

QUANTIFYING THE COMPLEXITY OF LEARNING QUANTUM FEATURES

LEONARDO BANCHI UNIVERSITY OF FLORENCE (ITALY)

EASY VS COMPLEX QUANTUM LEARNING PROCESSES

- Classical input *x*, classical feature y
- By varying x we can explore different states
- We want to extract the feature y e.g. via measurements

EXAMPLES

- Classical machine learning or QML with classical data
- Sample complexity in supervised learning: number of training data *N*
- (Validation) Error

What we want to learn

EXAMPLES

- Quantum sensing / Quantum hypothesis testing
- shots) *S*
- Continuous y (metrology): mean square error $\epsilon = \mathscr{O}\left(S^{-1}\right)$ or
- Discrete y : probability of misclassification $\epsilon = 1 2^{-H_{\text{min}}(Q|Y)} \approx e^{-\alpha S}$

$$
quare error c = O(S^{-1}) or c = O(S^{-2})
$$

$$
p(x, y) \longrightarrow y
$$

What we want to learn

Sample complexity is the number of copies of ρ (e.g. number of measurement

OPEN QUEST

-
- for all possible x ? Generalisation
- Error: $\epsilon = f(N, S)$?

Open question

How many values N of x does our learning strategy need to be able to extract y

How many copies S (e.g. measurement shots) of each $\rho(x,y)$ do we need?

Statistical Complexity of Quantum Learning

Leonardo Banchi, Jason Luke Pereira, Sharu Theresa Jose, and Osvaldo Simeone*

1.1. Scope

systems through the lens of information

on the amount of available data—which may be classical or

WORST CASE

 $\epsilon = \mathcal{O} \left(\left(N S \right) \right)$ $-1/3$)

iographysical system in microscopic physical system in \mathbb{R}^n) with tomography

$$
\epsilon = \mathcal{O}\left(d\left(NS\right)^{-1/2}\right) \quad \text{with tomography}
$$

LEARNING WITH QUANTUM STATES

L BANCHI, J PEREIRA, S PIRANDOLA PRX QUANTUM 2, 040321 (2021)

CLUSTERING IN HILBERT SPACES

- ▸ **Generalisation depends on both data and algorithm**
- ▸ Good generalisation when the states are **clustered** in the Hilbert space
- ▸ Large *I*(*Y* : *Q*), small *I*(*X* : *Q*)

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LEARNING TO CLASSIFY QUANTUM PHASES OF MATTER

RECOGNISING QUANTUM PHASES OF MATTER IS "EASY"

- $x = (h_1, h_2)$ is the set of classical Hamiltonian parameters (magnetic fields, couplings)
- $\rho(x,y)$ is the ground state
- is the phase (paramagnetic, *y* topological…)

SEE MEHRAN'S TALKS

WHY LEARNING QUANTUM PHASES OF MATTER ?

The learnt observable requires few measurements, e.g. single-qubit measurements

Generalisation bound grows (linearly?) with the number of qubits

See arXiv:2409.05188

LEARNING TO CLASSIFY ENTANGLEMENT WITHOUT KNOWING ENTANGLEMENT THEORY

LEARNING TO CLASSIFY ENTANGLEMENT IS "HARD"

Simple analytical strategy: measure the purity (~ Renyi-2 entropy) of either A or B subsystems. Outcomes 1 vs *O*(d^{−1}) variance $\left(1 - d^{-2}\right)$ /*S*

Toy problem: binary classification of separable (class 0) vs maximally entangled (class 1) states

 $|\psi_0(x)\rangle = U_A(x) |0\rangle \otimes U_B(x) |0\rangle$ $|\psi_1(x)\rangle \propto U_A(x) \otimes U_B(x)$

The states from both classes uniformly populate the Hilbert space, the two classes are not clustered

Linear observables cannot discriminate the two states.

$$
|\psi_1(x)\rangle \propto U_A(x) \otimes U_B(x) \sum_{i=1}^d |i, i\rangle_{AB}
$$

$$
\int dx |\psi_0(x)\rangle\langle\psi_0(x)| = \int dx |\psi_1(x)\rangle\langle\psi_1(x)| = \frac{\log 1}{d^2}
$$

CLASSIFICATION WITH SUPPORT VECTOR MACHINES

Theorem: if we estimate kernels via the SWAP test, the optimal observable learnt from data converges to the purity measurement when $N\to\infty, S\to\infty$

$$
K(x, x') = |\langle \psi(x) | \psi(x') \rangle|^4
$$

= Tr $\left[\rho(x)^{\otimes 2} \rho(x')^{\otimes 2} \right]$
Feature space

Error *e* with
$$
N = \frac{1}{\epsilon} \left(1 + \frac{d^4}{S} \right)^2
$$
 data

Banchi et al. arXiv (soon)

PARAMETERS WITH ACCURACY HIGHER THAN 99% Banchi et al. arXiv (soon)

TEMPORAL STOCHASTIC PROCESSES

▸ EUR vs 1 SGD

▸ What will happen in the future?

?

TEMPORAL STOCHASTIC PROCESSES

Not all hidden variables are relevant to predict future observations Replace the huge hidden states with a finite memory, which contains all the relevant information to predict future outcomes

QUANTUM SIMULATORS

K are Kraus operators

Matrix Product States for Quantum Stochastic Modeling

 \mathcal{L} , received 2018; revised manuscript received manuscript received 10

Chengran Yang,^{1,2,*} Felix C. Binder,^{1,2,†} Varun Narasimhachar,^{1,2} and Mile Gu^{1,2,3,‡} ¹School of Physical and Mathematical Sciences, Nanyang Technological University, 637371 Singapore, Singapore ²Complexity institute, Nanyang Technological University, 639798 Singapore, Singapore ³ Centre for Quantum Technologies, National University of Singapore, 3 Science Drive 2, 117543 Singapore, Singapore

Ivan Glasser^{1,2}*, Ryan Sweke³, Nicola Pancotti^{1,2}, Jens Eisert^{3,4}, J. Ignacio Cirac^{1,2} ¹Max-Planck-Institut für Quantenoptik, D-85748 Garching ²Munich Center for Quantum Science and Technology (MCQST), D-80799 München 3Dahlem Center for Complex Quantum Systems, Freie Universität Berlin, D-14195 Berlin 4Department of Mathematics and Computer Science, Freie Universität Berlin, D-14195 Berlin

PHYSICAL REVIEW LETTERS 121, 260602 (2018)

Expressive power of tensor-network factorizations for probabilistic modeling

In general

Quantum: Classical:

In some cases (*ε*-machines) the quantum model uses less memory

ANTUM SIMULATORS

LEARNING FRAMEWORK

- : vector of past observations *x*
- : vector of future observations *y*

- memory $\mathbf{r} \cap \mathbf{r} \cap \mathbf{r}$
- Does the quantum model display an advantage in generalisation?

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MACHINE

Department of Physics and Astronomy, University of Florence, via G. Sansone 1, I-50019 Sesto Fiorentino, (FI), Italy INFN Sezione di Firenze, via G. Sansone 1, I-50019 Sesto Fiorentino, (FI), Italy), (FI), Italy

PAPER

Leonardo Banchi

E-mail: leonardo.banchi@unifi.it

1. Introduction

Accuracy vs memory advantage in the quantum simulation of stochastic processes 29 December 2023

The quantum model better compresses past observations into a quantum $_$ $I(X: \mathcal{Q}_{\text{quantum mean}} \leq I(X: \mathcal{Q}_{\text{classical memory}})$ m eddes t_{max} a direct mapping between its exact classical and \mathbf{r}_{max} past observations into a quantum we was a $p_{\text{max}} = \frac{1}{2}$ \sim \sim chassical memory, or alternatively, better accuracy with less memory, or alternatively, better accuracy with less memory, \sim \mathbf{m} citation and DOI.

OPEN ACCESS

ARNING Science and Technology

EXAMPLE: STOCHASTIC CLOCK EVAMDI E. STAPHASTIP PLAP

EXAMPLE: STOCHASTIC CLOCK

\hat{W} \bar{E} ⁴ \hat{R} \hat{H} AL+W \hat{W} \hat{R} \hat{S} \hat{S} DATA \hat{C} \hat{C} \hat{S} \hat{S} MPHOGRAPHY)

bl. of by extra lymph. pass S C hence using 100 million samples. Table (c) shows the Bhattachary samples of the resulting quantum and resulting $\frac{1}{\sqrt{2}}$ 1 $1 \quad 3$ $\overline{2}$ $\overline{2}$ 2 $\mathbf 1$ $\mathbf{1}$ $2 \quad 3$ $\overline{2}$ $\overline{2}$ $\overline{2}$ $\overline{2}$ $\overline{2}$ $\overline{2}$ 1 $\mathbf{1}$ quantum and classical models. Each table entry shows the fitting procedure for $\mathbf{1}$ random initial $\mathbf{1}$ 1 $\mathbf{1}$ \cdots ... \cdots $143 \t3$ $\overline{2}$ $\overline{2}$ $\overline{2}$ 1. $\mathbf{1}$ 144 2 1 1 1 1 1 $\overline{2}$ 145 2 $\overline{2}$ 1. 1 $\mathbf{1}$ 146 2 1 1 1 1 1 1472 $\overline{2}$ $\mathbf{2}$ 2 2

 $X =$ lymphography["data"]["features"]

148 rows \times 19 columns

 $N=2$ (comp) 0.842 0.652 Probability of observing the **last ten** outcomes in **Figure 5.** Distribution of the distance between consecutive ones in the exact and fitted classical (c-fit) and quantum (q-fit) the Lymphography dataset, given the previous simulators, and the complete computer computer computer computer computer computer computer the exact simulation of the exact simulation of the exact simulation uses \mathcal{L} is a simulated computation uses \mathcal{L} is a s the fitted ones experience observations and the sequence of \mathcal{L} \mathcal{L}

Each entry shows the best value over 100 and the next ten functions on the next ten functions. Parenthesis. The last ten outcomes in the parenthesis. The parenthesis in the previous observations, $\frac{3}{6}$, $\frac{3}{6}$, $\frac{3}{6}$ repetitions of the fitting procedure for random initial configurations, and the median in

LEARNING WITH QUANTUM DATA

nature reviews physics

Abstract Communication

Nature Reviews Physics 5, 141-156 (2023)

Learning quantum systems

Leonardo Banchi ^{®,9}, Florian Marquardt^{10,11}, Luca Pezzè^{1,2}⊠ & Cristian Bonato ^{® 5}⊠

The embedded input can be expressed as a Fourier-like series

Generalisation error from the Fourier matrix $F_{\omega\omega'}$ $=$ 1 *N*Ω *N* ∑ $n =$ $e^{ix_n(\omega-\omega')}$

$2^{I_{1/2}(X:Q)} < 2^{H_{1/2}(f)} \leq |\Omega|$

with spectrum *f* e mbedding circuit $x \mapsto \rho(x)$ decision via POVM Π_c

See also Caro et al. Quantum (2021)

class/label $c =$ "cat"

(b) Quantum classifier $|0\rangle$ $|0\rangle$. .
.
. .
.
. $|0\rangle$ $U_1(x)$ V_1 $U_2(x)$ *...* $\left| U_L(x) \right| \left| V_L \right| \qquad \left| \Pi_c \right|$

$$
|\psi(x)\rangle = \frac{1}{\sqrt{|\Omega|}} \sum_{\omega \in \Omega} e^{i\omega x} |\phi_{\omega}\rangle
$$

*[|]c*ih*c[|] ^c* optimization of POVM Π_c

EXAMPLE: PARAMETRIC QUANTUM CIRCUIT EMBED

(a) Data distribution

predicted class *c* of *x*

(c) Dilated measurements

= optimization of unitary *U^M*

BARCODES AND PATTERN REC

- ▸ Barcode classification must identify each pixel correctly
- ▸ Handwriting classification is easier as errors are tolerated!

 $error \simeq F(\rho_{black}, \rho_{white})$ Hamming_{4→9}

▸

- ‣ L. Banchi, Q. Zhuang, S. Pirandola, Phys. Rev. Applied 14, 064026 (2020)
- ‣ C Harney, L Banchi, S Pirandola, Phys. Rev. A 103, 052406 (2021)

‣ JL Pereira, L Banchi, Q Zhuang, S Pirandola, Phys. Rev. A 103, 042614 (2021)

CONCLUSION

▸ Generalisation favoured by **discarding information** about the input that is irrelevant for predicting the output **C**uantum compression?

-
- ▸ **Few measurement shots** may be a dominant source of generalisation error
- ▸ Quantitative data-dependent bounds. Greater physical insight.
- ▸ Different applications:
	- ▸ Quantum pattern recognition with entanglement-enhanced quantum sensor
	- ▸ Classification of quantum states and phases of matter
	- ▸ Quantum machine learning of classical data

STATISTICAL LEARNING THEORY

