

# Neural Quantum Monte Carlo for quantum simulators

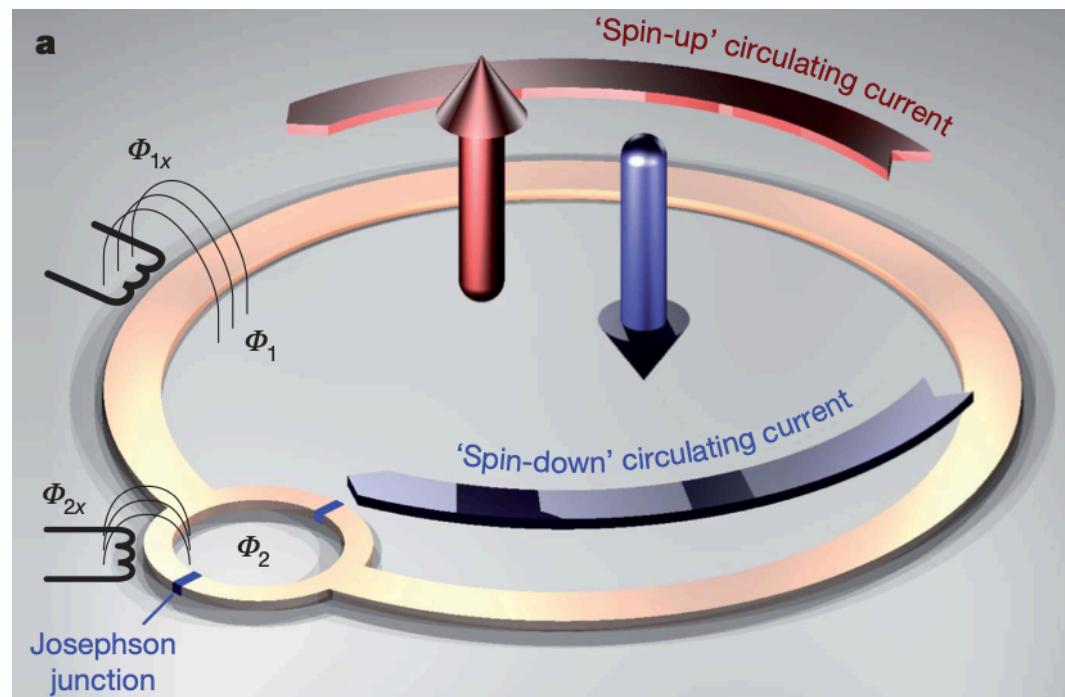
**Luca Brodolini - University of Camerino**

**Quantum Computing @ INFN - Padova, 28-31 October 2024**

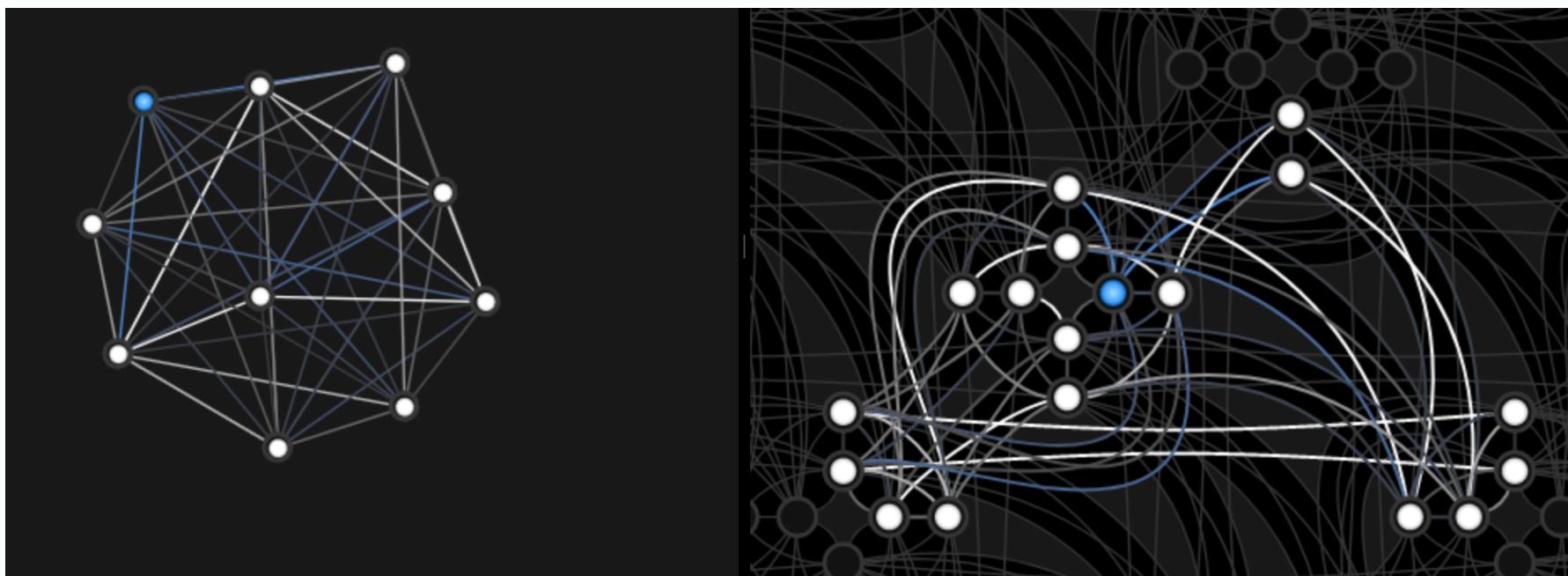
# Programmable quantum platforms

$$H = - \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z - \sum_i \Gamma_i \sigma_i^x - h \sum_i \sigma_i^z$$

## Quantum annealers

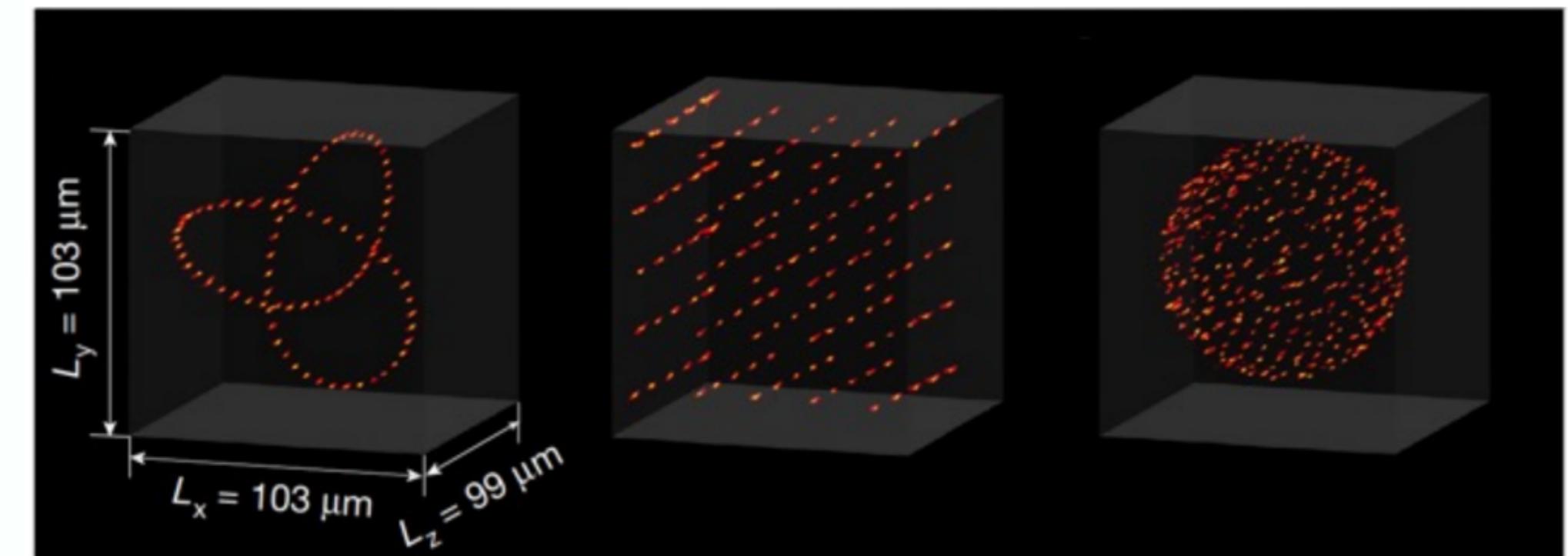
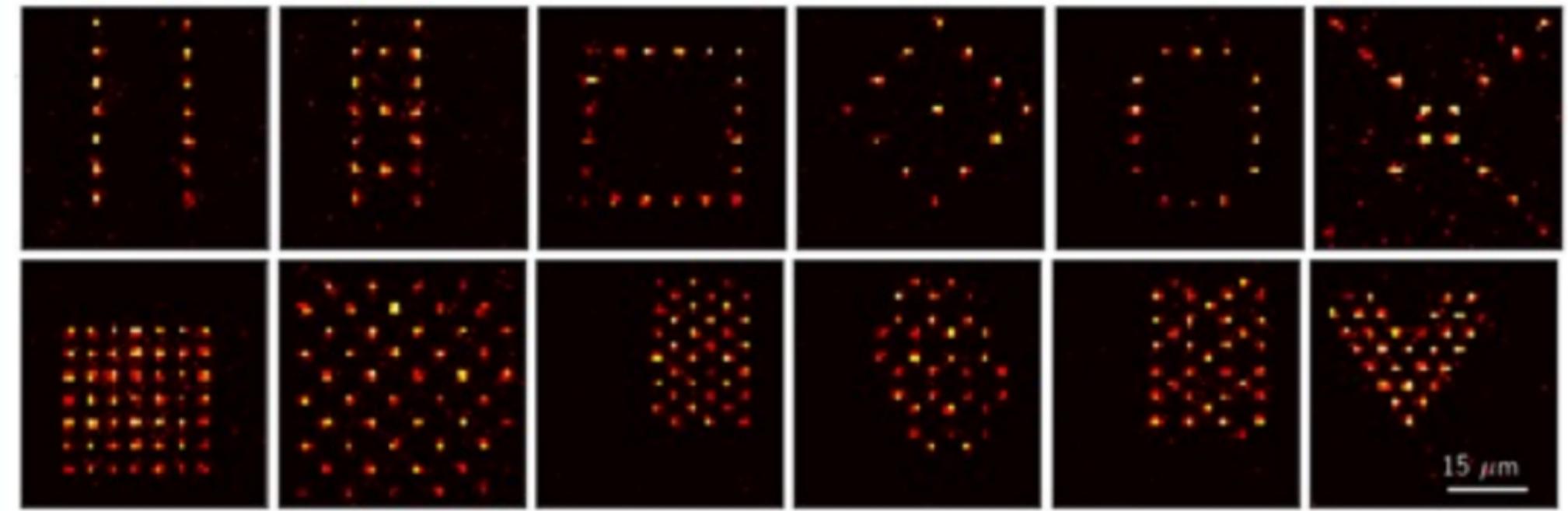


Nature 473, 194–198 (2011)



D-Wave QPU - 10 qubits all-to-all connections

## Rydberg atoms in optical tweezers

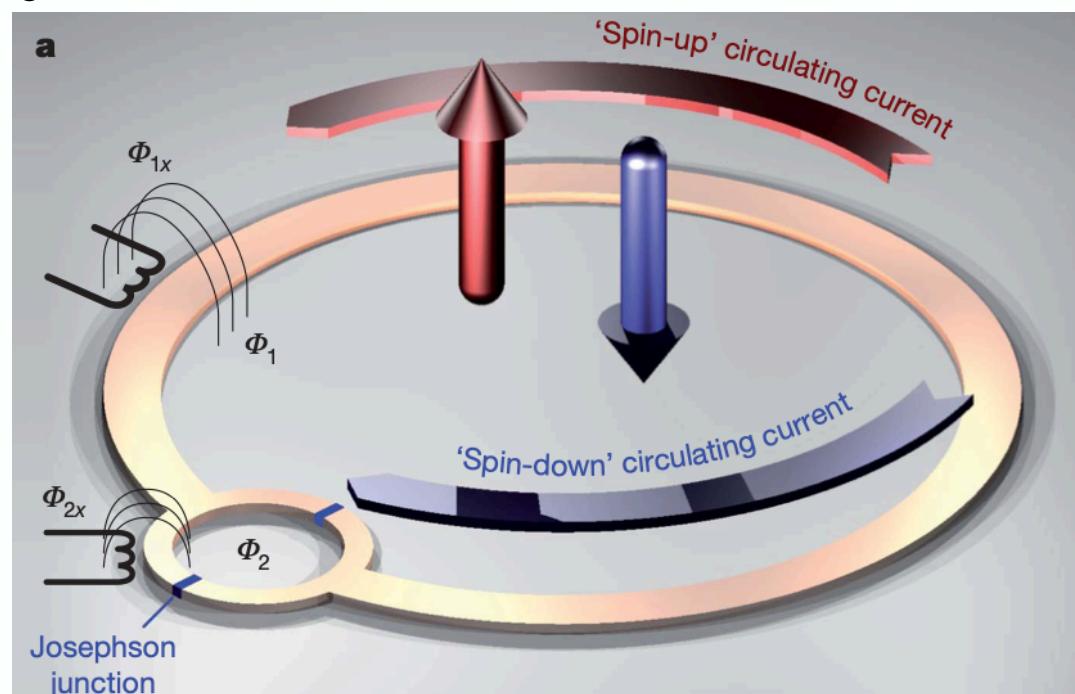


Nature 561, 79–82 (2018)

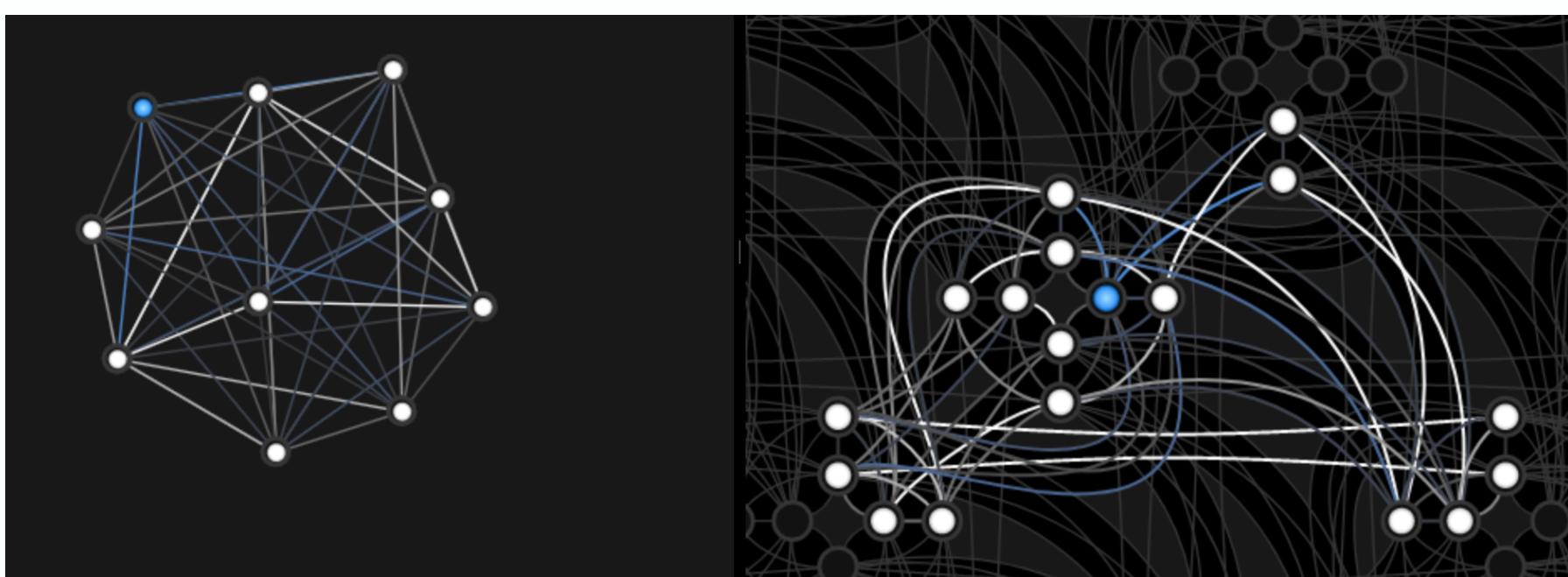
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## Quantum annealers



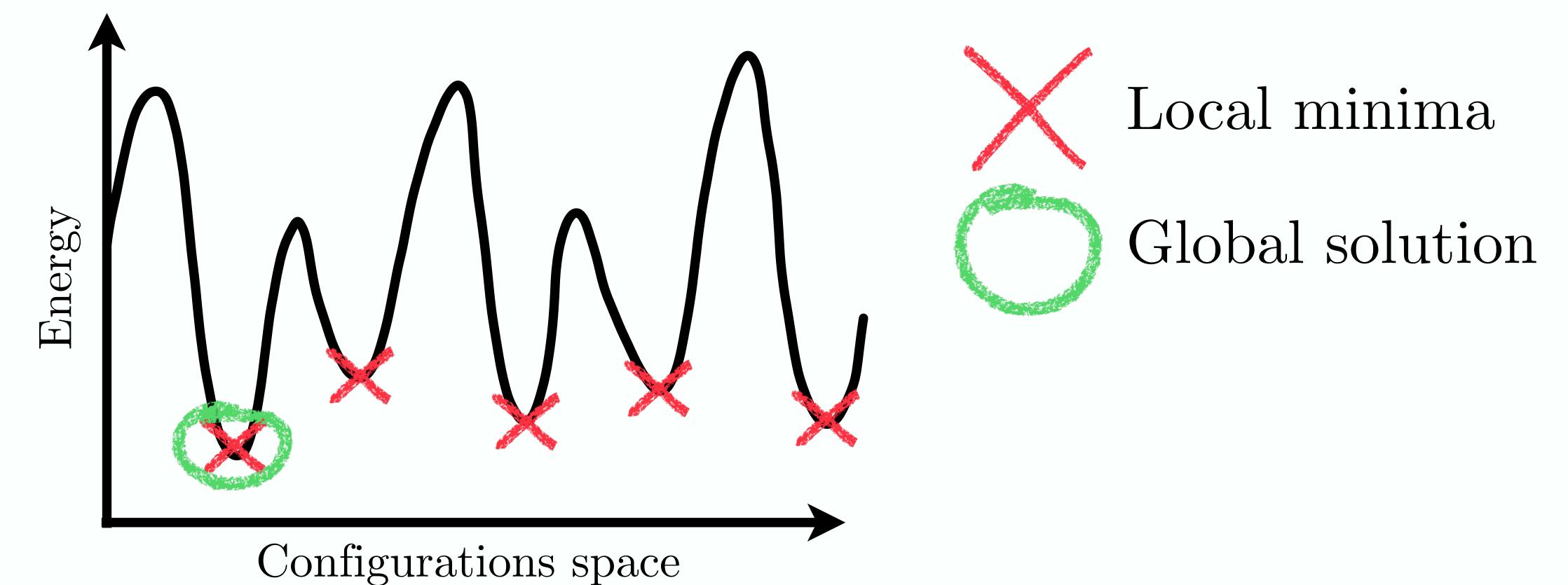
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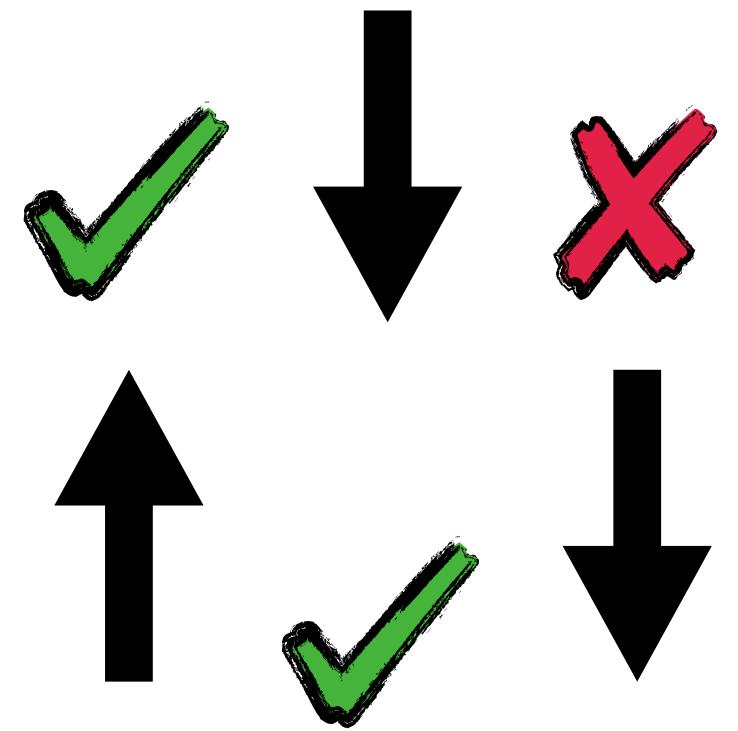
D-Wave QPU - 10 qubits all-to-all connections

## Problems:

- Optimization (energy minimization)
- Sampling (from many low energy states)



# Quantum spin glasses



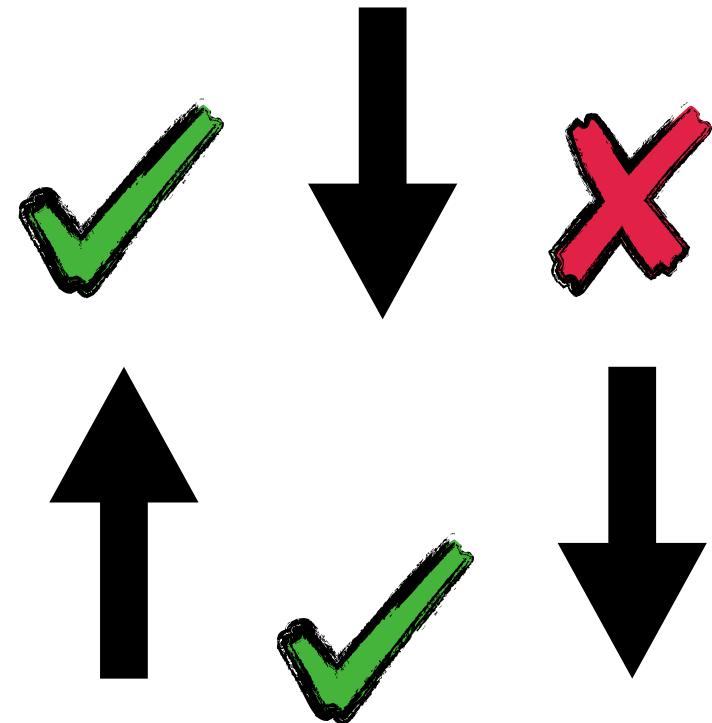
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Frustration in triangular lattice with anti-ferromagnetic interactions.

Frustration arises from:

- Geometry of the lattice
- Nature of the couplings  $J_{ij}$

# Quantum spin glasses



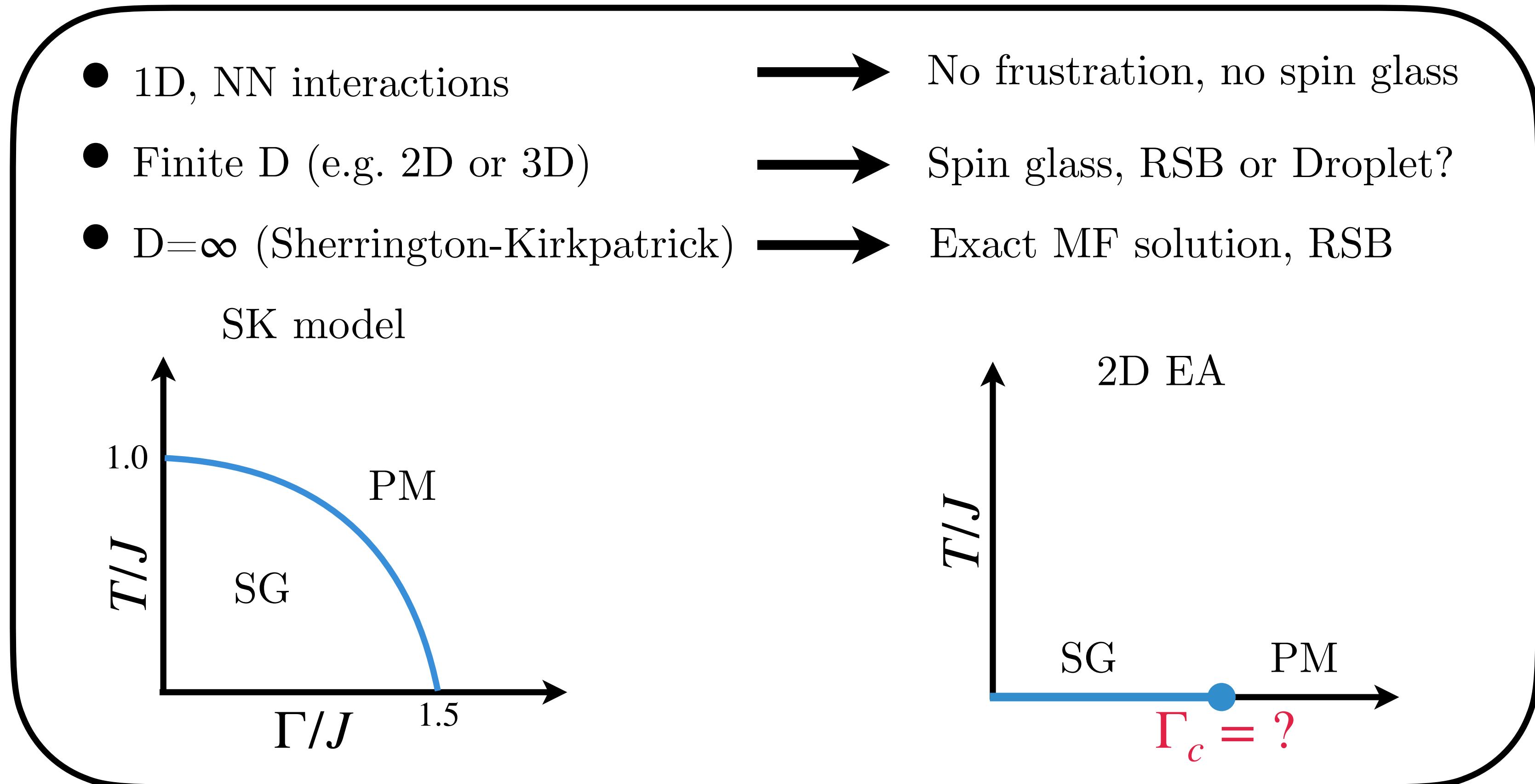
Frustration in triangular lattice with anti-ferromagnetic interactions.

Frustration arises from:

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$$H = - \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z - \sum_i \Gamma_i \sigma_i^x$$

- 1D, NN interactions → No frustration, no spin glass
- Finite D (e.g. 2D or 3D) → Spin glass, RSB or Droplet?
- $D=\infty$  (Sherrington-Kirkpatrick) → Exact MF solution, RSB



# Projective Quantum Monte Carlo for random Ising Model

Imaginary-time evolution:  $\psi(\mathbf{x}, \tau) = \exp(-\tau H)\psi(\mathbf{x}, 0) \underset{\tau \rightarrow \infty}{=} \psi_{\text{GS}}(\mathbf{x})$

Stochastic process:  $\psi(\mathbf{x}, \tau + \Delta\tau) = \sum_{\mathbf{x}'} G(\mathbf{x}, \mathbf{x}', \Delta\tau) \psi(\mathbf{x}', \tau)$

Green's function:  $G(\mathbf{x}, \mathbf{x}', \Delta\tau) = \left\langle \mathbf{x} \left| \exp \left( -\tau (H - E_{\text{ref}}) \right) \right| \mathbf{x}' \right\rangle$

Not a stochastic matrix!

- In general  $G(\mathbf{x}, \mathbf{x}', \Delta\tau)$  can be  $< 0 \implies$  sign problem! **NOT OUR CASE**
- $\sum_{\mathbf{x}} G(\mathbf{x}, \mathbf{x}', \Delta\tau) \neq 1 \implies$  Transition probability not normalized

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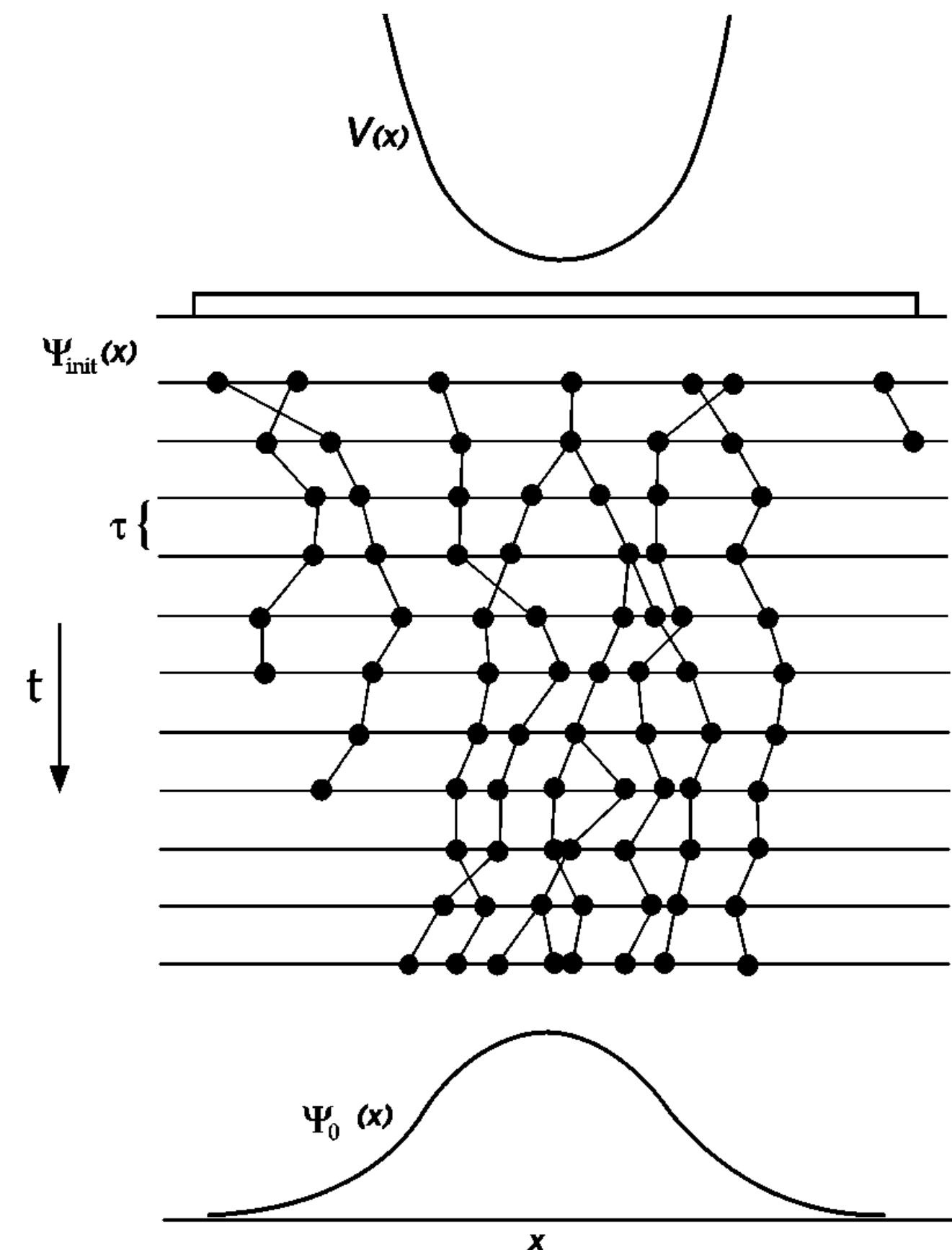
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Stochastic evolution of walker population  
branching process: kill or replicate



# Projective Quantum Monte Carlo for random Ising Model

## Importance sampling with a neural guiding wave-function

Imaginary-time evolution:  $\psi(\mathbf{x}, \tau)\psi_g(\mathbf{x}) = \exp(-\tau H)\psi(\mathbf{x}, 0)\psi_g(\mathbf{x}) \underset{\tau \rightarrow \infty}{=} \psi_{\text{GS}}(\mathbf{x})\psi_g(\mathbf{x})$

Common choices for  $\psi_g(\mathbf{x})$ :

Stochastic process:  $\psi(\mathbf{x}, \tau + \Delta\tau)\psi_g(\mathbf{x}) = \sum_{\mathbf{x}'} \tilde{G}(\mathbf{x}, \mathbf{x}', \Delta\tau)\psi(\mathbf{x}', \tau)\psi_g(\mathbf{x}')$

Modified Green's function:  $\tilde{G}(\mathbf{x}, \mathbf{x}', \Delta\tau) = \left\langle \mathbf{x} \middle| \exp \left( -\tau (H - E_{\text{ref}}) \right) \middle| \mathbf{x}' \right\rangle \frac{\psi_g(\mathbf{x})}{\psi_g(\mathbf{x}')}$

- Jastrow WF;
- Hartree-Fock WF;
- Variational ansatz;

# Projective Quantum Monte Carlo for random Ising Model

## Importance sampling with a neural guiding wave-function

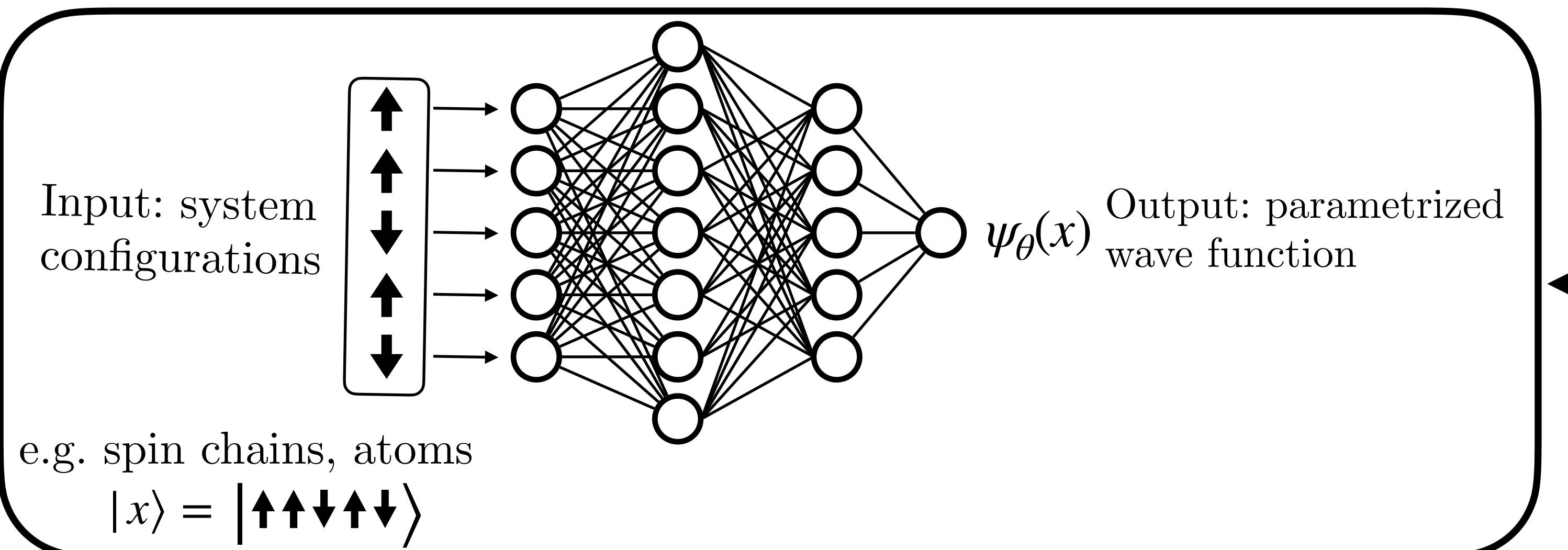
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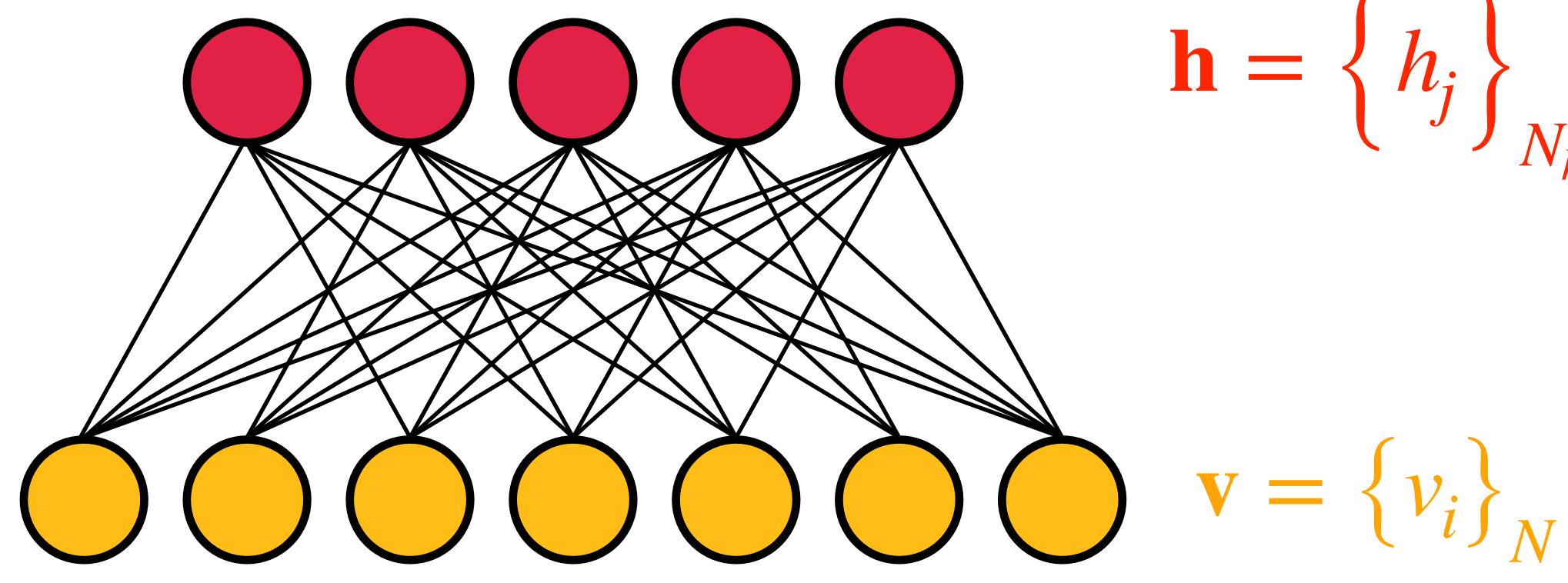
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Common choices for  $\psi_g(\mathbf{x})$ :

- Jastrow WF;
- Hartree-Fock WF;
- Variational ansatz;
- Neural Quantum States

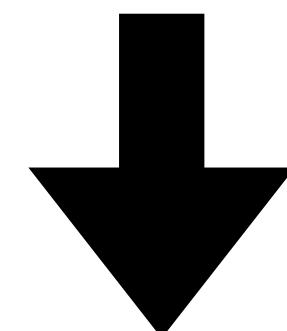


# Neural guiding function - Restricted Boltzmann Machine



RBM learns marginal distribution of visible variables

$$P_{RBM}(\mathbf{v}) = \sum_{\mathbf{h}} P(\mathbf{v}, \mathbf{h}) = \frac{1}{Z} \sum_{\mathbf{h}} \exp(-E_{RBM}(\mathbf{v}, \mathbf{h}))$$



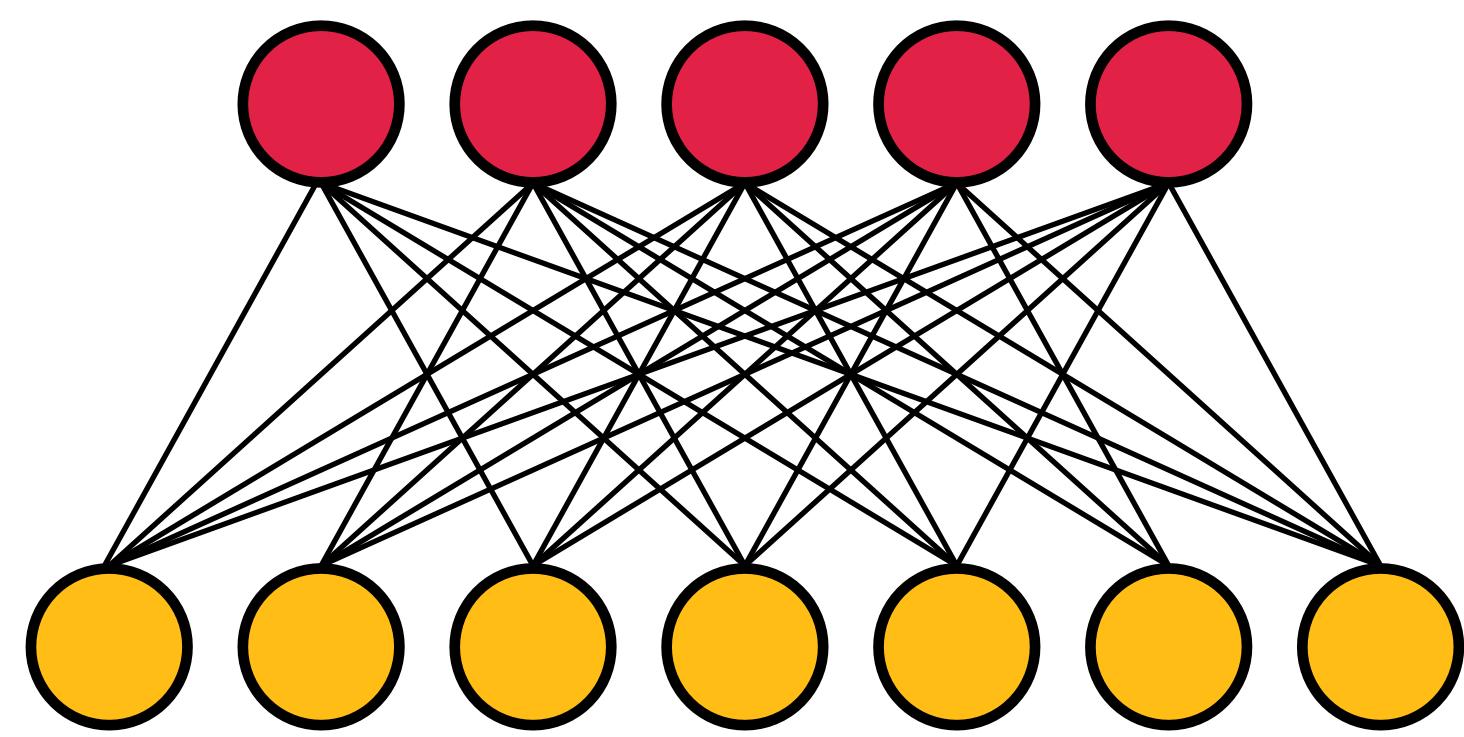
$$\psi_g = \psi_{RBM} \propto \sqrt{P_{RBM}}$$

Energy-based model

$$E_{RBM}(\mathbf{v}, \mathbf{h}) = - \sum_{j=1}^{N_h} \sum_{i=1}^N w_{ij} v_i h_j - \sum_{i=1}^N v_i b_i - \sum_{j=1}^{N_h} h_j c_j$$

$w_{ij}$  : weights       $b_i, c_j$  : biases

# Neural guiding function - Restricted Boltzmann Machine

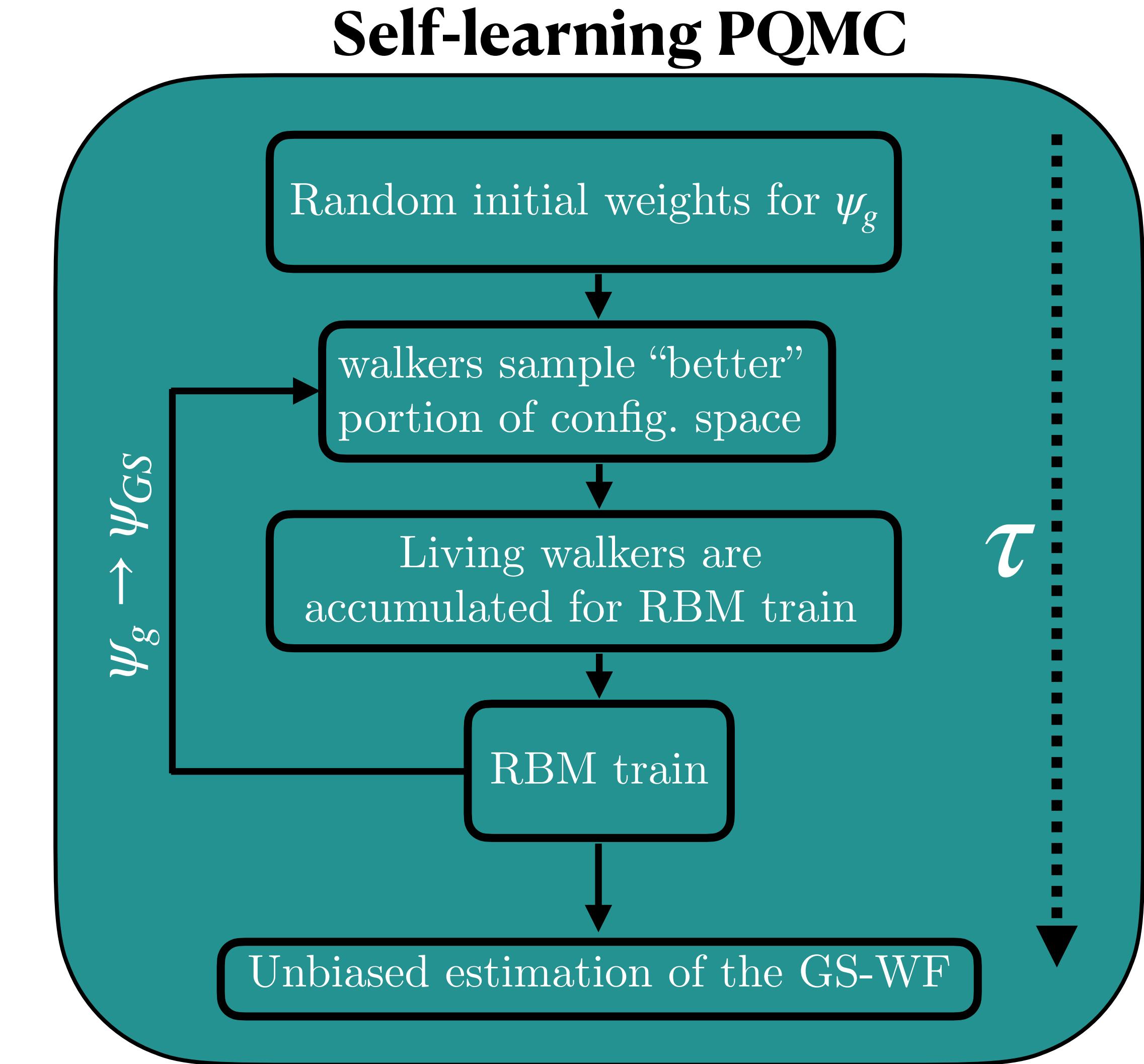


$$\mathbf{h} = \left\{ h_j \right\}_{N_h}$$

$$\mathbf{V} = \left\{ v_i \right\}_N$$

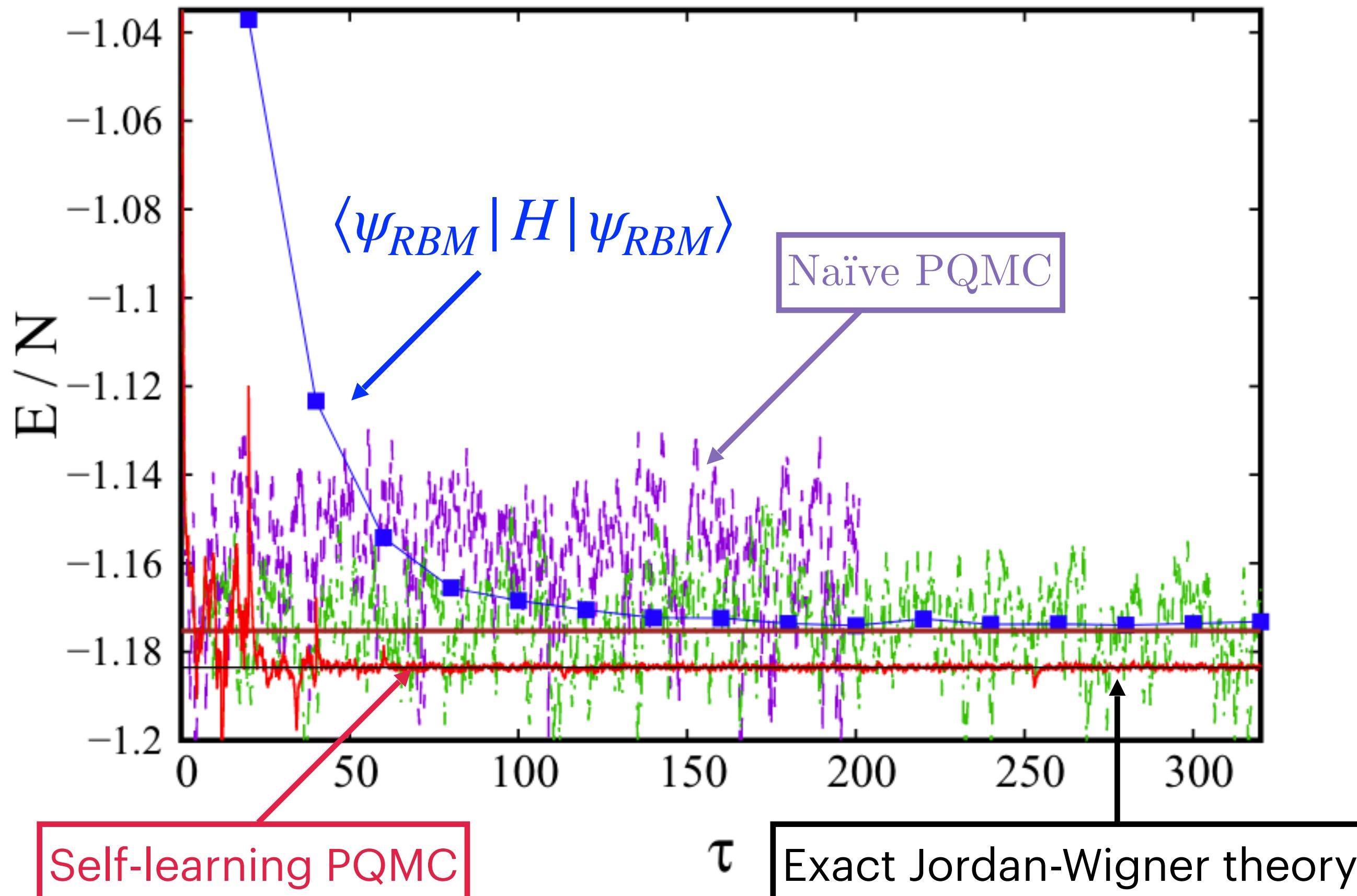
$$P_{RBM}(\mathbf{v}) = \sum_{\mathbf{h}} P(\mathbf{v}, \mathbf{h}) = \frac{1}{Z} \sum_{\mathbf{h}} \exp(-E_{RBM}(\mathbf{v}, \mathbf{h}))$$

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# Neural guiding function - Restricted Boltzmann Machine

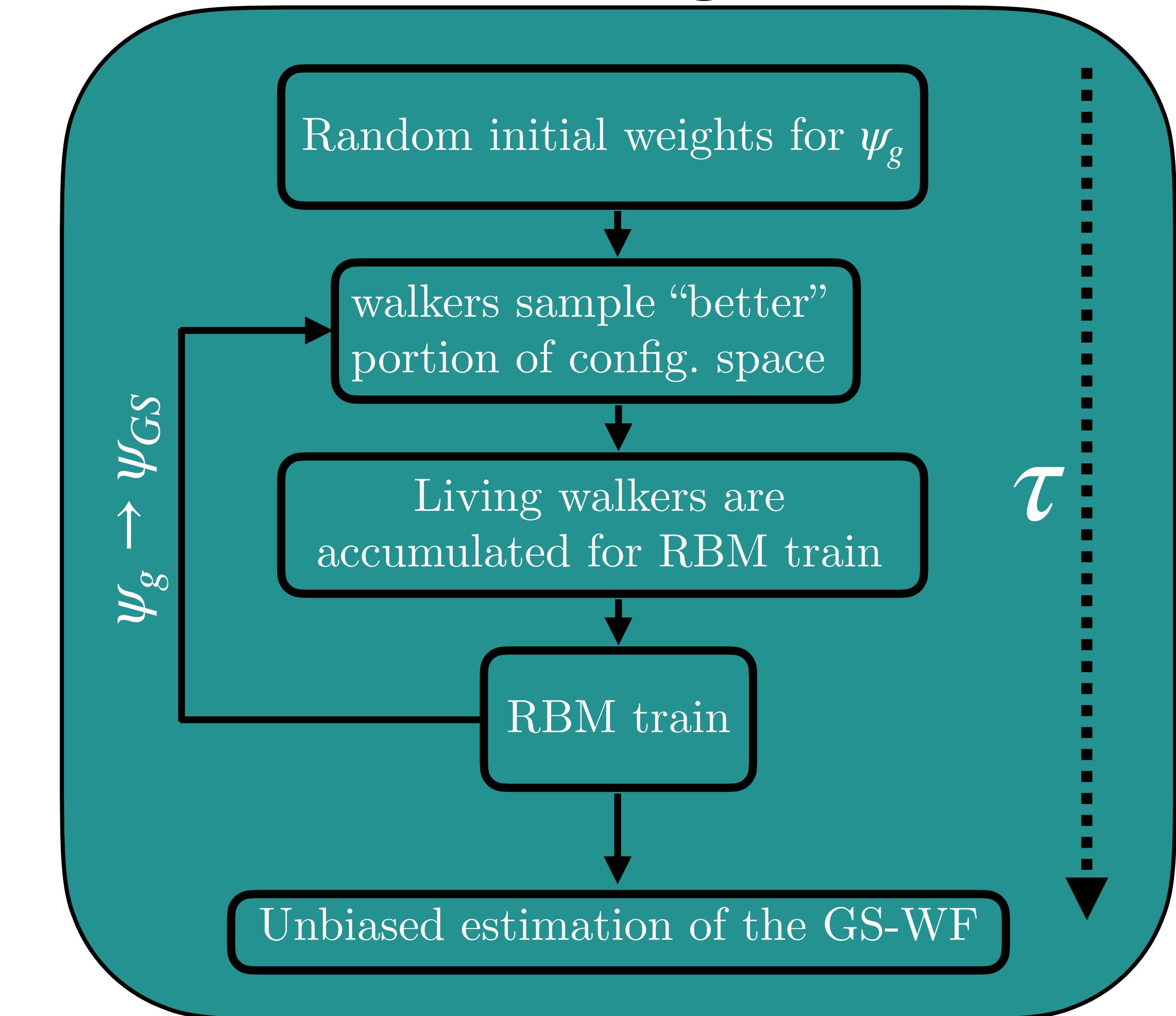
Ferrom. 1D, N=80



- No population control bias!

- Exact variational ansatz not needed!

## Self-learning PQMC



# 2D Edwards-Anderson model

$$H = - \sum_{\langle ij \rangle} J_{ij} \sigma_i^z \sigma_j^z - \sum_i \Gamma_i \sigma_i^x \quad J_{ij} \sim \text{Norm}[\mu = 0, \sigma = 1]$$

Two copies of the system are needed

EA order parameter:

$$q_{EA} = \frac{1}{N} \sum_{i=1}^N \langle \sigma_i^z \rangle^2$$

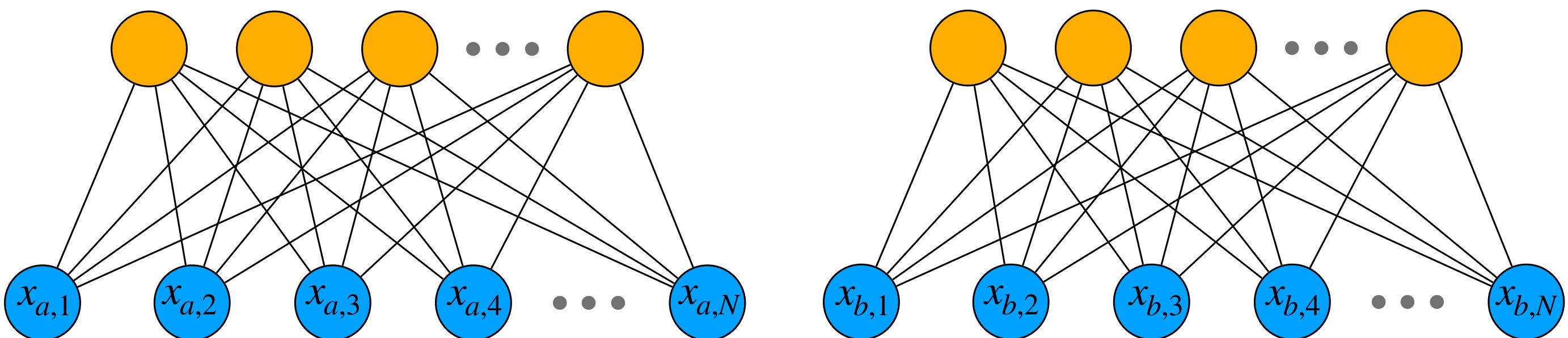
Spin overlap:

$$\hat{q} = \frac{1}{N} \sum_{i=1}^N \sigma_{iA}^z \sigma_{iB}^z$$

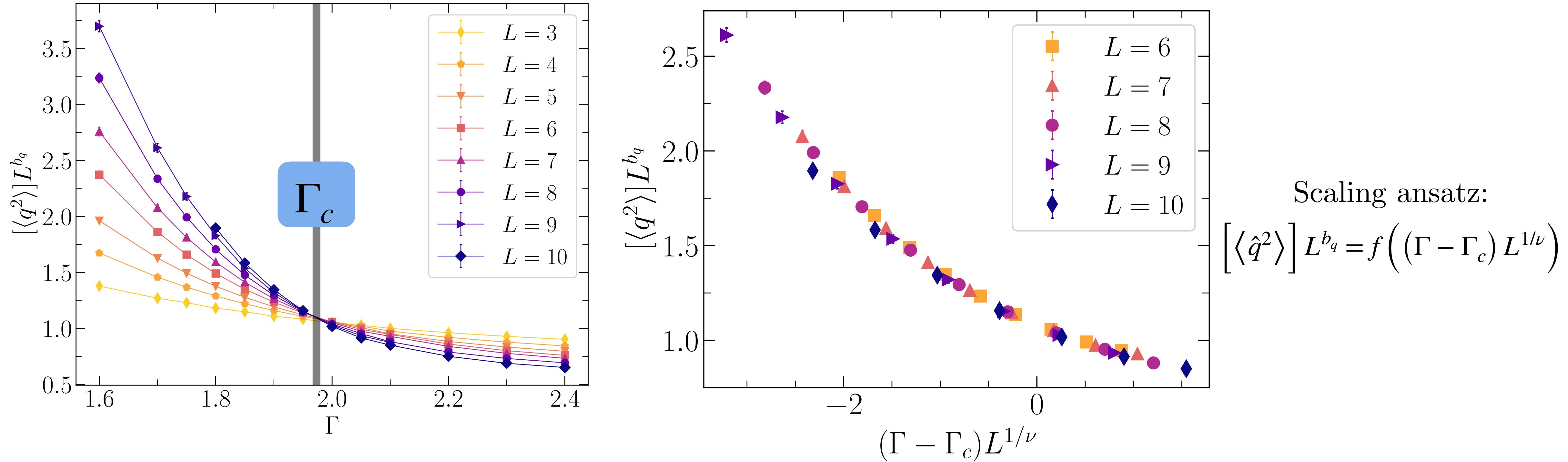
Spin glass susceptibility:  $\frac{\chi_{SG}}{N} = [\langle \hat{q}^2 \rangle]$

$$H_{TOT} = H_A + H_B$$

separable RBM



# 2D Edwards-Anderson model



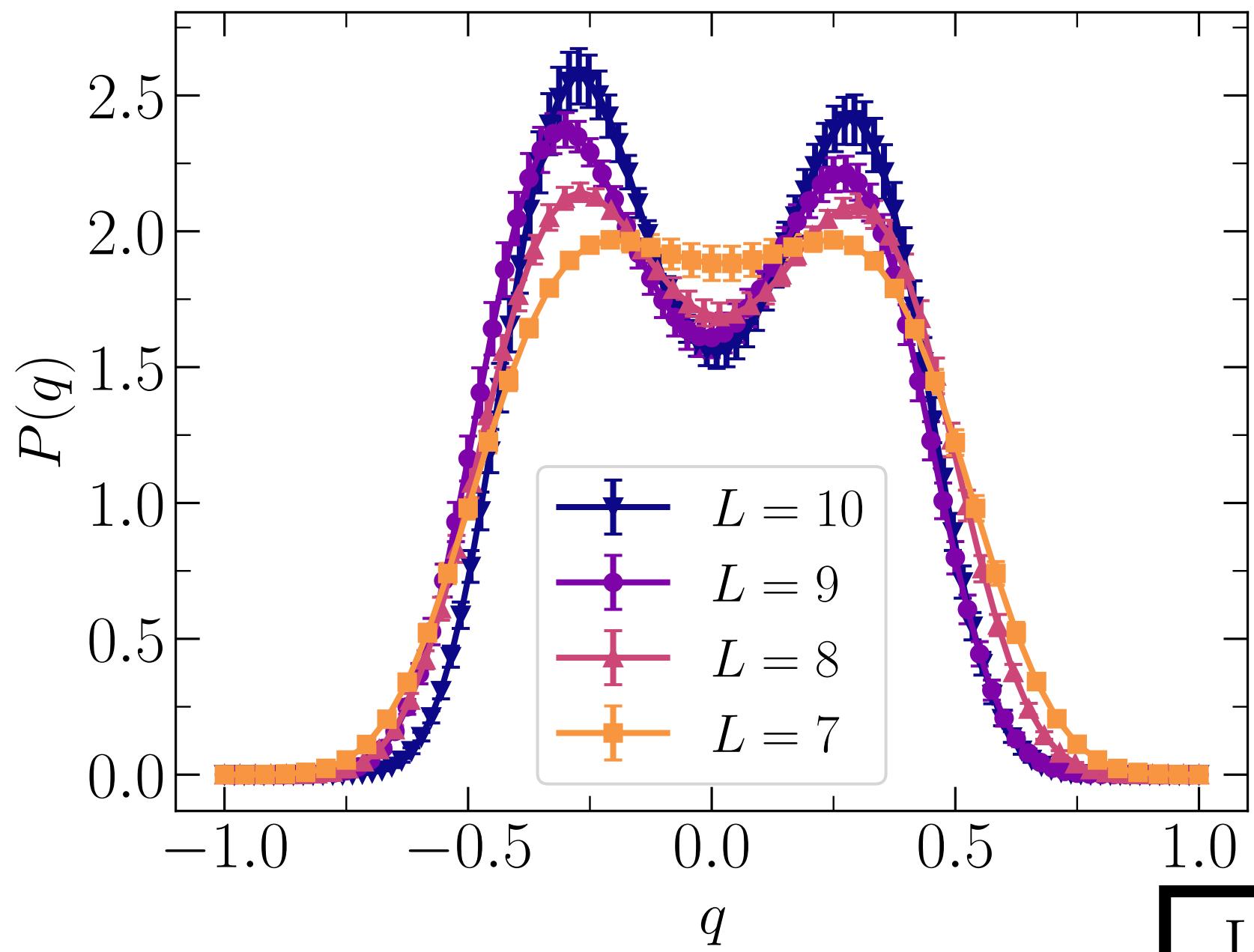
Couplings	$1/\nu$	$b_q$	$\Gamma_c$	Ref.
50%. $\pm 1$	1.02(16)	1.76(3)	2.11(1)	A.D. King <i>et al.</i> , Nature (2023)
50% $\pm 1$	0.71(24)(9)	1.73(8)(8)	2.18(1)	M. Bernaschi <i>et al.</i> , Nature (2024)
$\mathcal{N}(0, 1)$	1.11(22)	1.68(8)	1.98(7)	This work

# 2D Edwards-Anderson model

## Replica symmetry breaking

$$\hat{q} = \frac{1}{N} \sum_{i=1}^N \sigma_{iA}^z \sigma_{iB}^z$$

Overlap distribution:  $P(q) = \langle \delta(q - \hat{q}) \rangle$



Droplet picture

$$P(q) = \frac{1}{2}(\delta(q - q_{EA}) + \delta(q + q_{EA}))$$

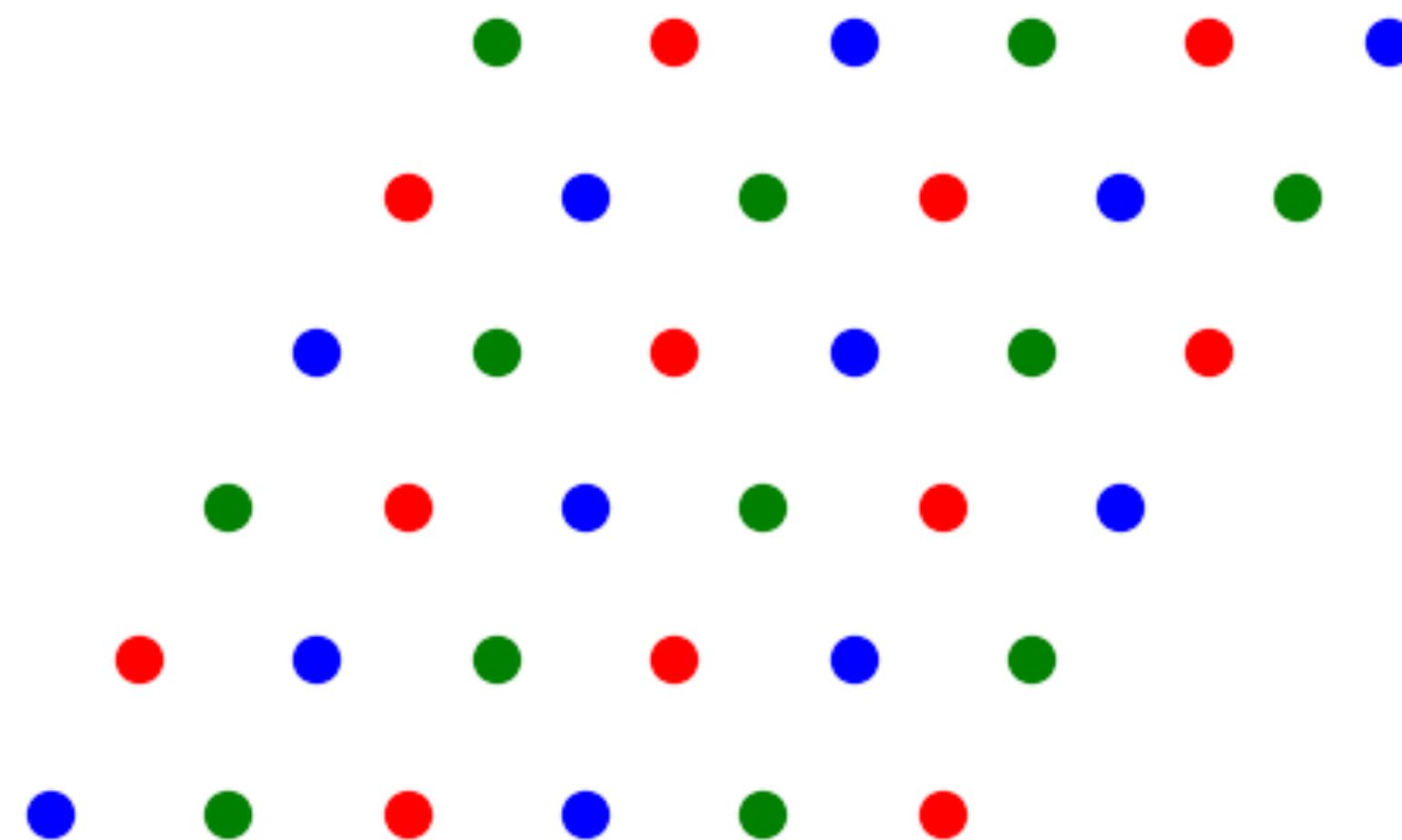
Replica Symmetry Breaking

$$P(q = 0) > 0$$

- Non-trivial distribution of feasible sizes, hints to RSB
- Larger sizes needed to take  $L \rightarrow \infty$  limit

# Preliminary results on triangular lattices

$J_{ij} < 0$  (antiferrom.)  $\rightarrow$  Intrinsic geometric frustration



### 3 sublattices (blue, red, green)

## “Order by disorder” mechanism:

# Small quantum fluctuations: Clock phase

# High quantum fluctuations: PM phase

Ordered phase: spins in the sublattices  
assumes 6 discrete preferred orientations

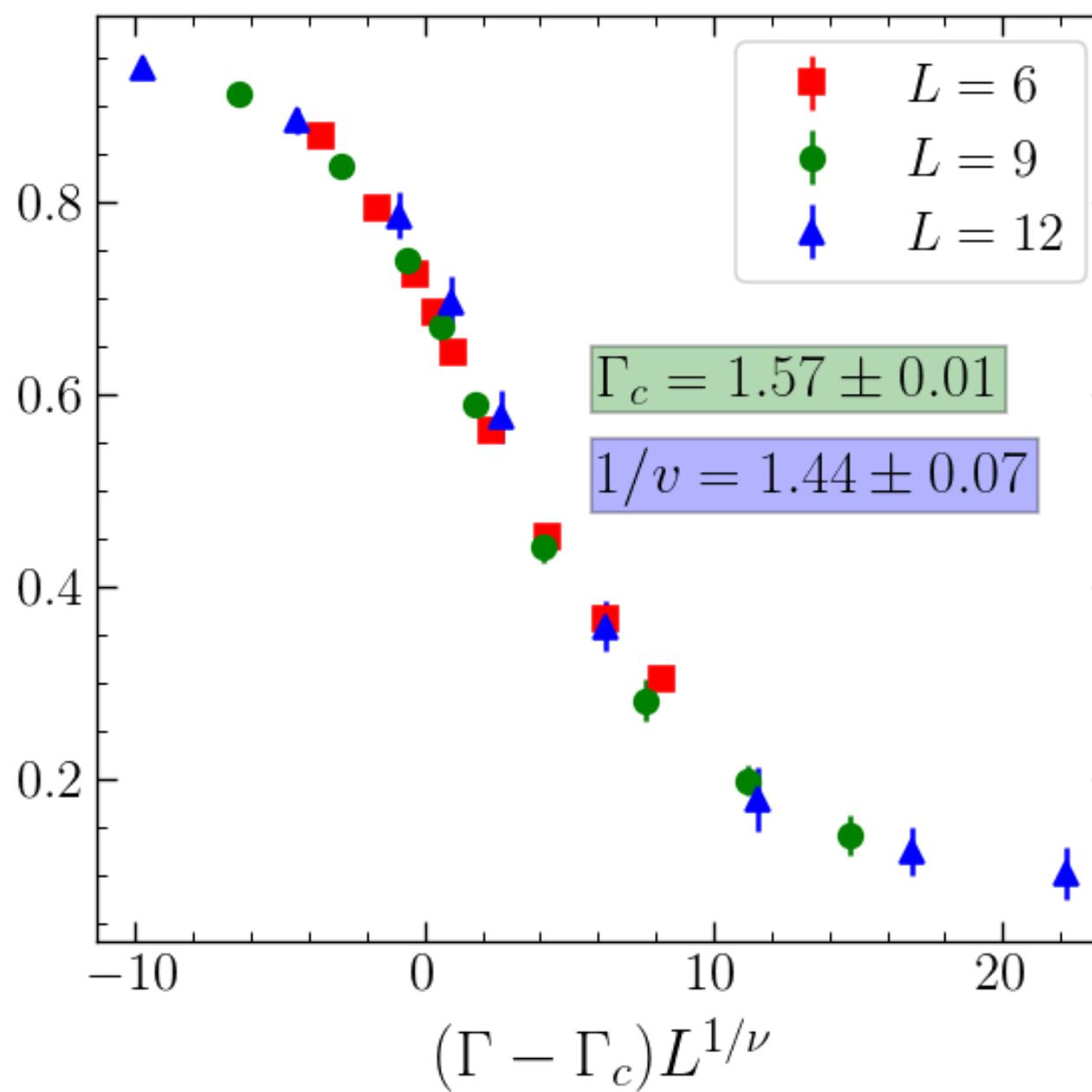
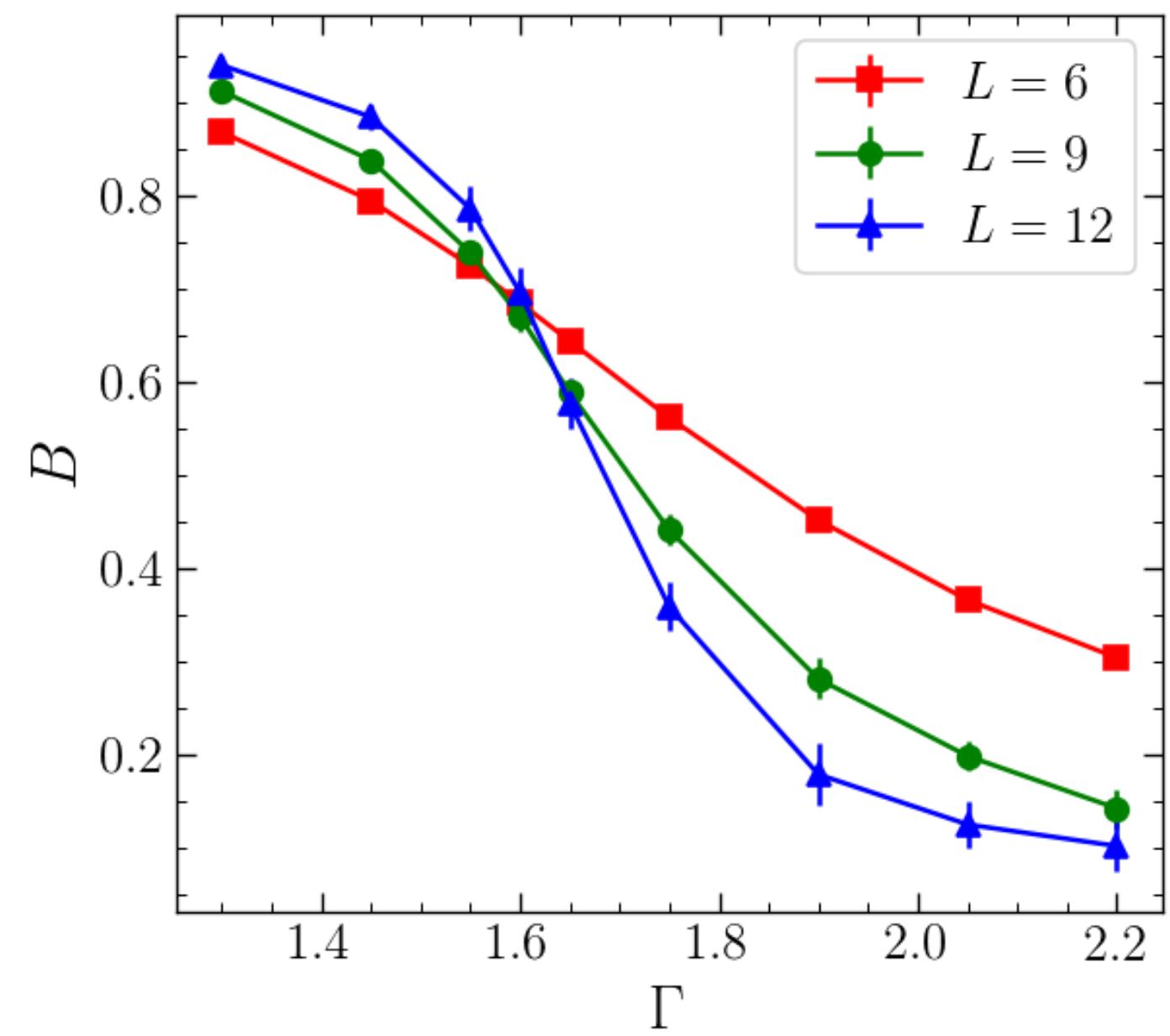
## XY complex magnetization:

$$me^{i\theta} = \textcolor{blue}{m_1} + \textcolor{red}{m_2}e^{i(4\pi/3)} + \textcolor{green}{m_3}e^{i(-4\pi/3)}/\sqrt{3}$$

# Preliminary results on triangular lattices

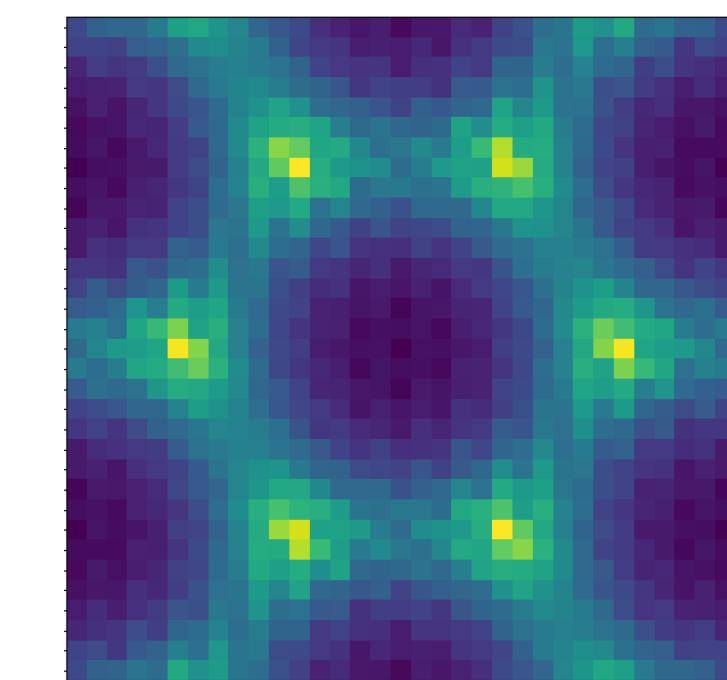
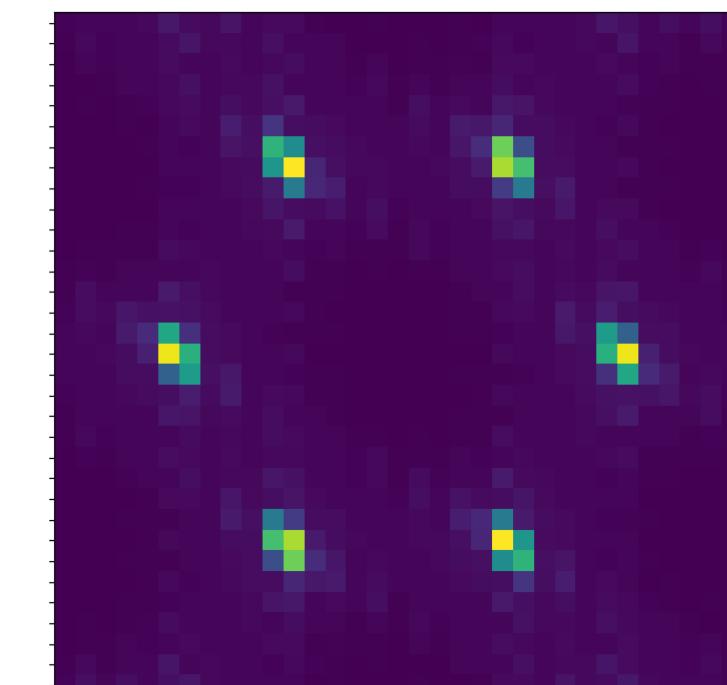
NN Model

Binder cumulant:  $B = 2 \left( 1 - \frac{\langle m^4 \rangle}{2\langle m^2 \rangle^2} \right)$



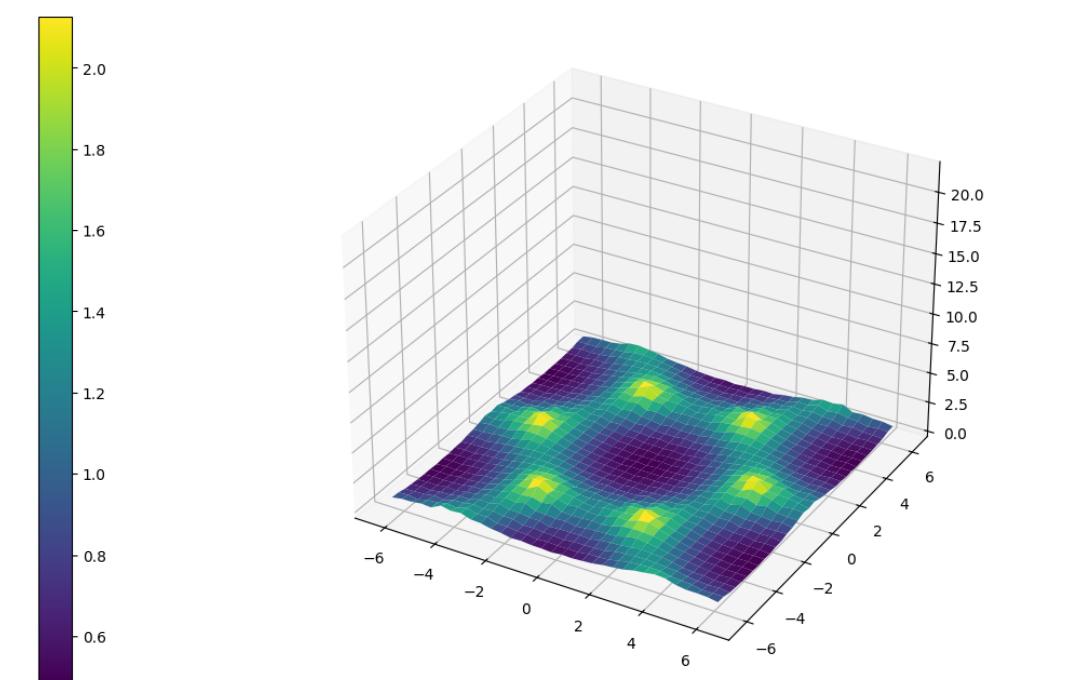
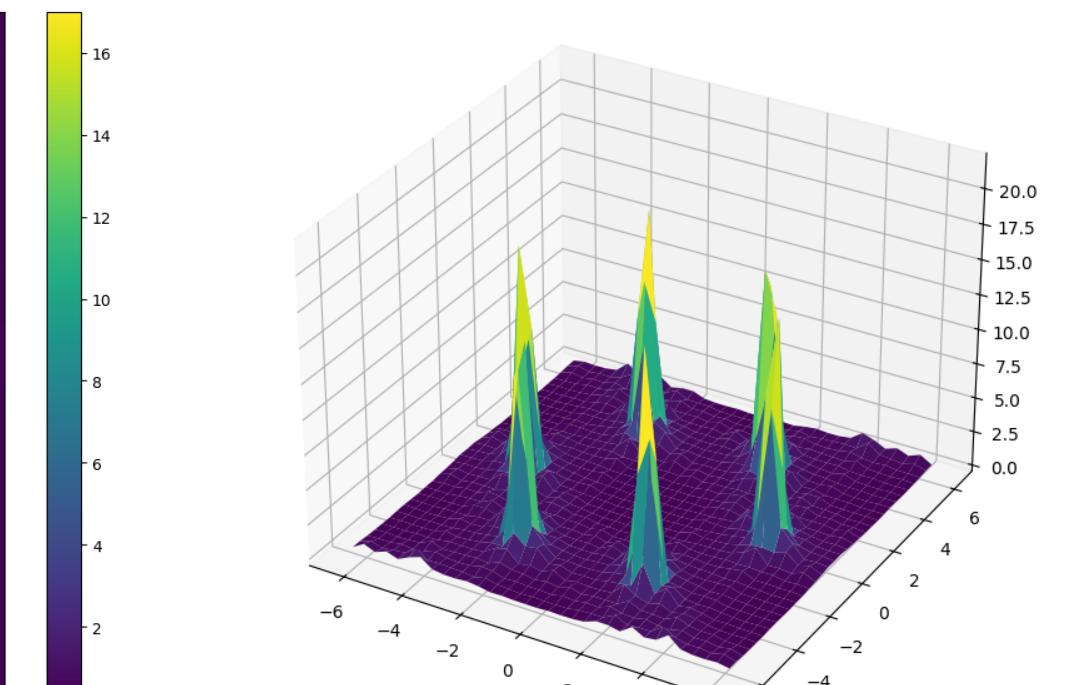
Spin structure factor:

$$S(\mathbf{q}) = \langle S_{\mathbf{q}}^z S_{-\mathbf{q}}^z \rangle = \frac{1}{N} \sum_{ij} \langle S_i^z S_j^z \rangle e^{i\mathbf{q} \cdot (\mathbf{x}_i - \mathbf{x}_j)}$$



$\Gamma = 1.0$

$\Gamma = 2.0$



# Conclusions

- Our Neural PQMC technique allow us to investigate the spin glass phase of the 2D EA model
- We determine the critical transverse field as well as the critical exponents in good agreement with state-of-the-art works
- Large scale PQMC simulation allow us to compute  $P(q)$ : hints to RSB

L.Brodoloni, S.Pilati arXiv:2407.05978(2024)

- Preliminari results on “order-by-disorder” phenomena enlight a promising route to explore different kinds of interactions (Rydberg)

## Complex Quantum Matter @ UNICAM



Left to right: Giovanni Midei, Victor Velasco, Luis Ardila, Sebastiano Pilati, Andrea Perali, Nicola Pinto, Simone Cantori, Luca Brodoloni, Andrea Della Valle and Verdiana Piselli. CQM-Unicam, May 30 2024.

Not Included: Carlo Lucheroni, Filippo Pascucci, Meenakshi Sharma and Gabriele Spada.

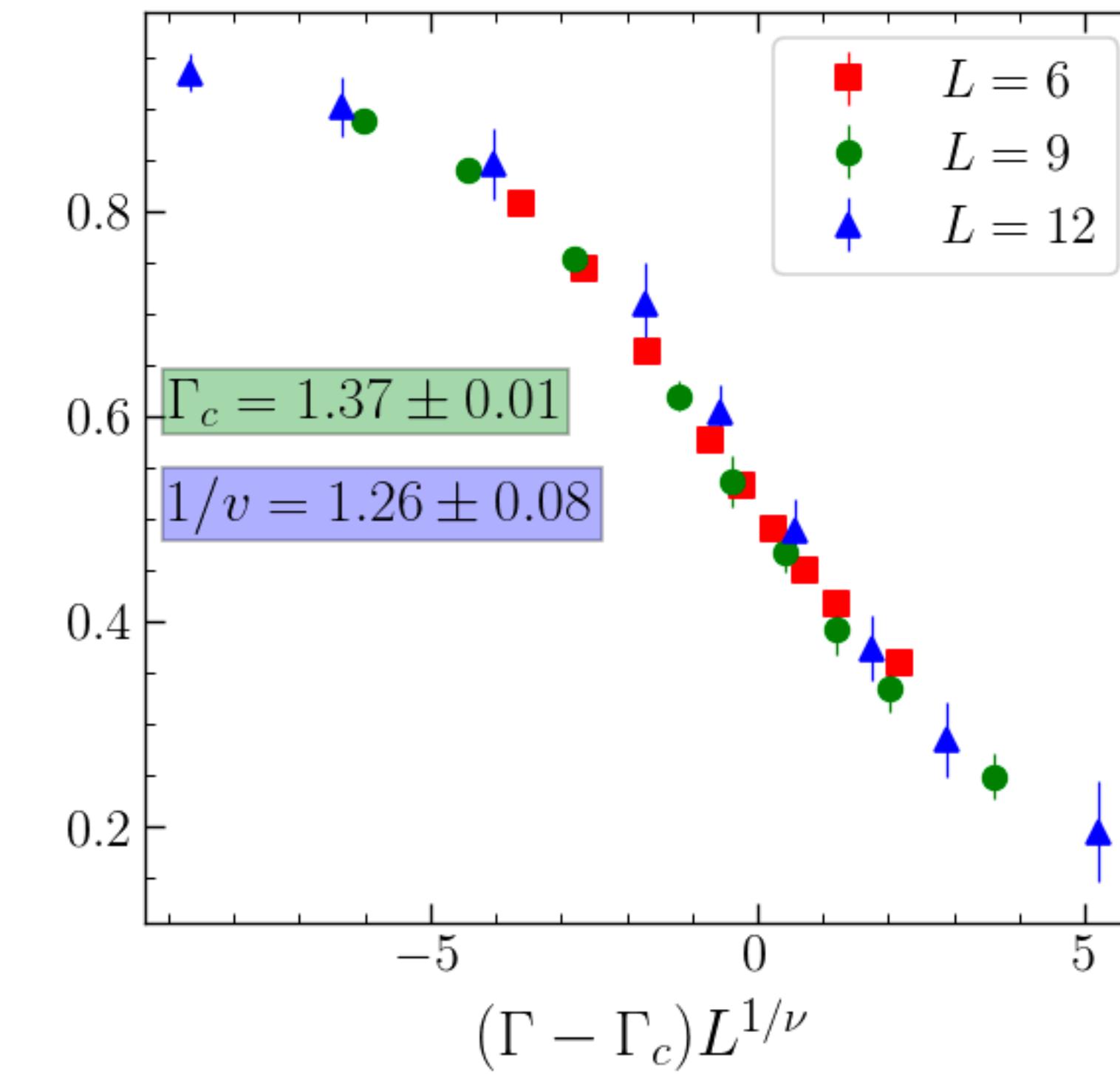
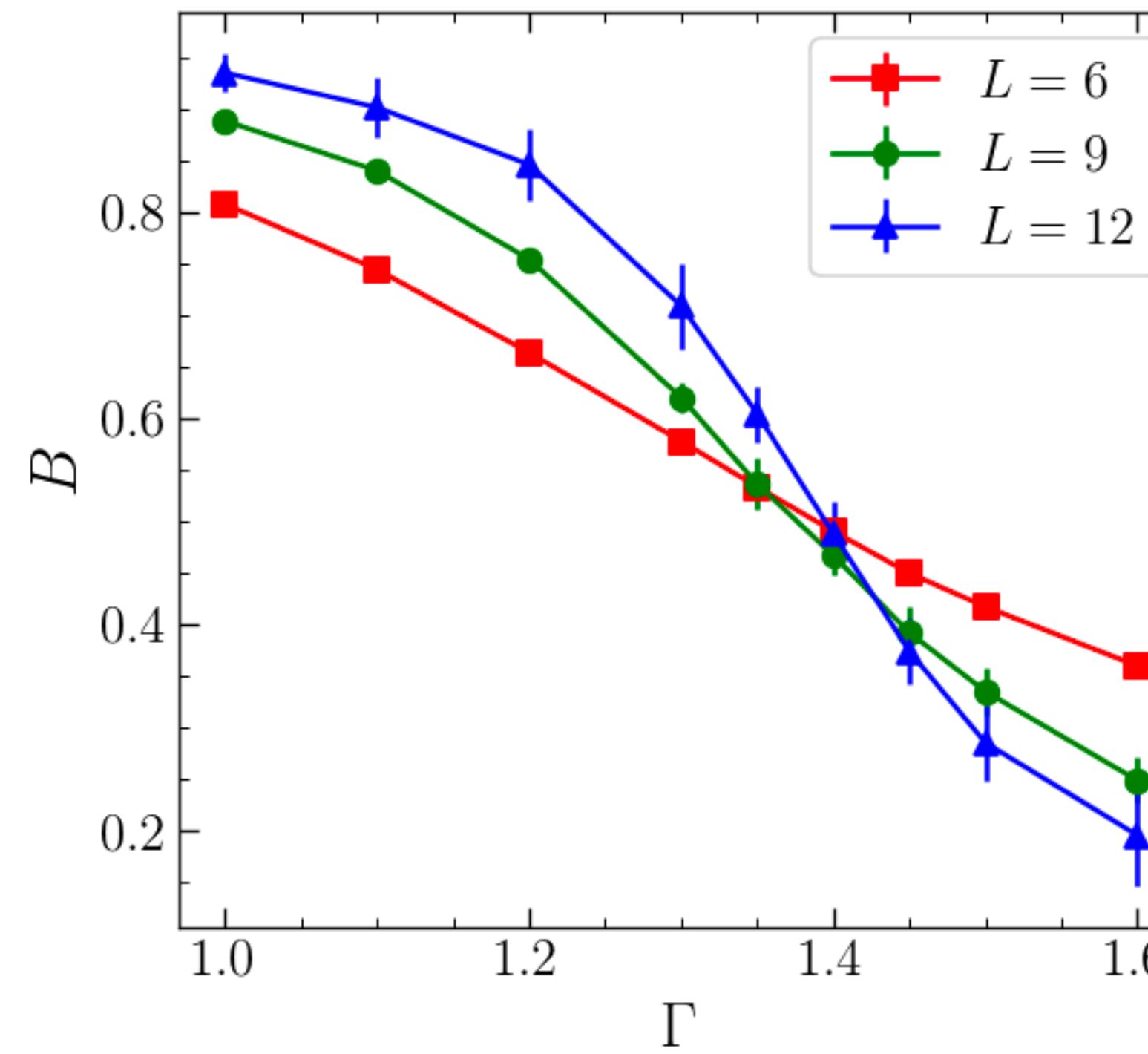
# Long range interactions model

$$H = - \sum_{i < j} J_{ij} \sigma_i^z \sigma_j^z - \sum_i \Gamma_i \sigma_i^x$$

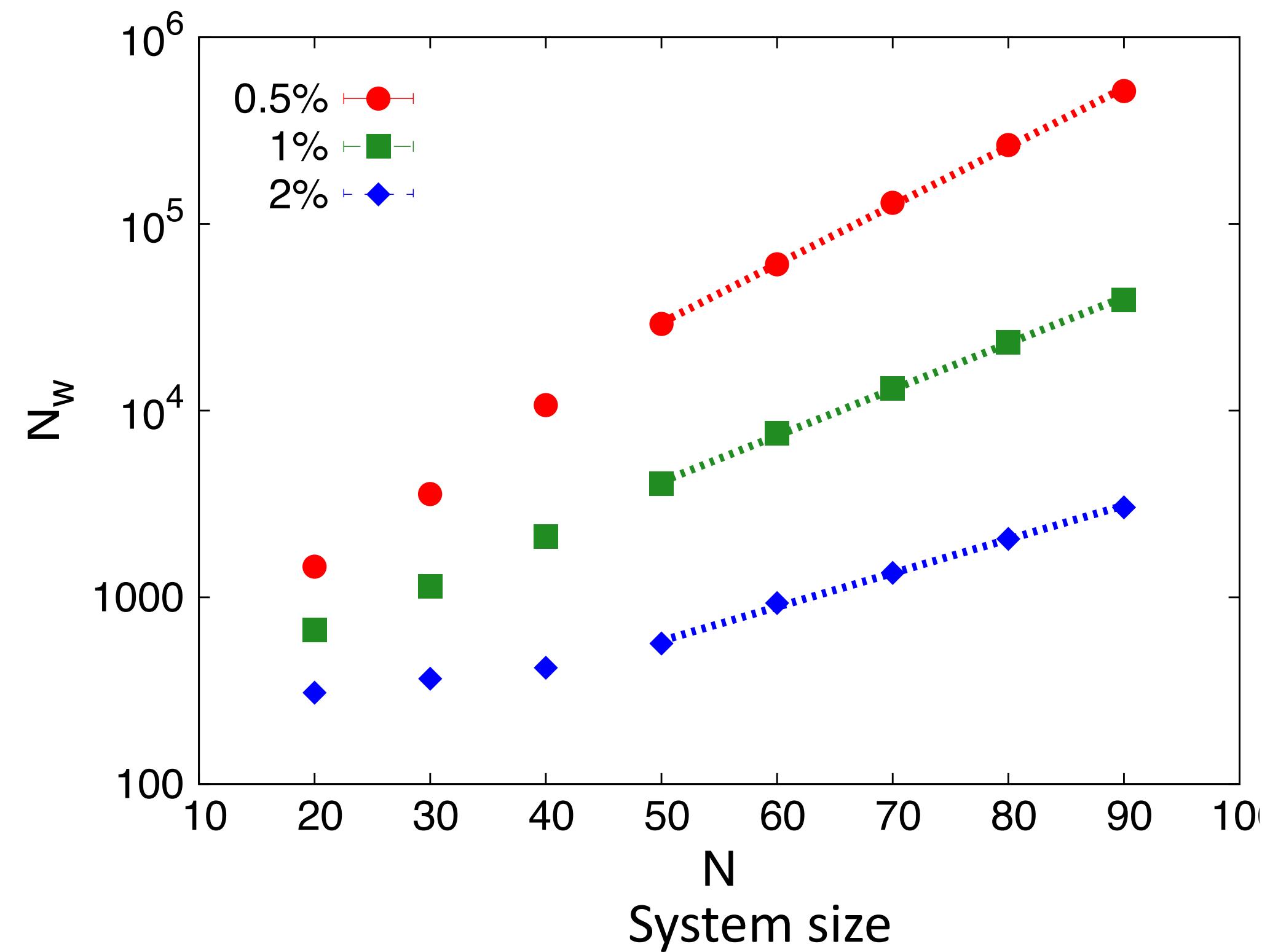
$$J_{ij} \propto \frac{1}{r_{ij}^6}$$

Rydberg interactions Ising model

$\alpha = 6$  Model



# of walkers required to keep relative err. fixed



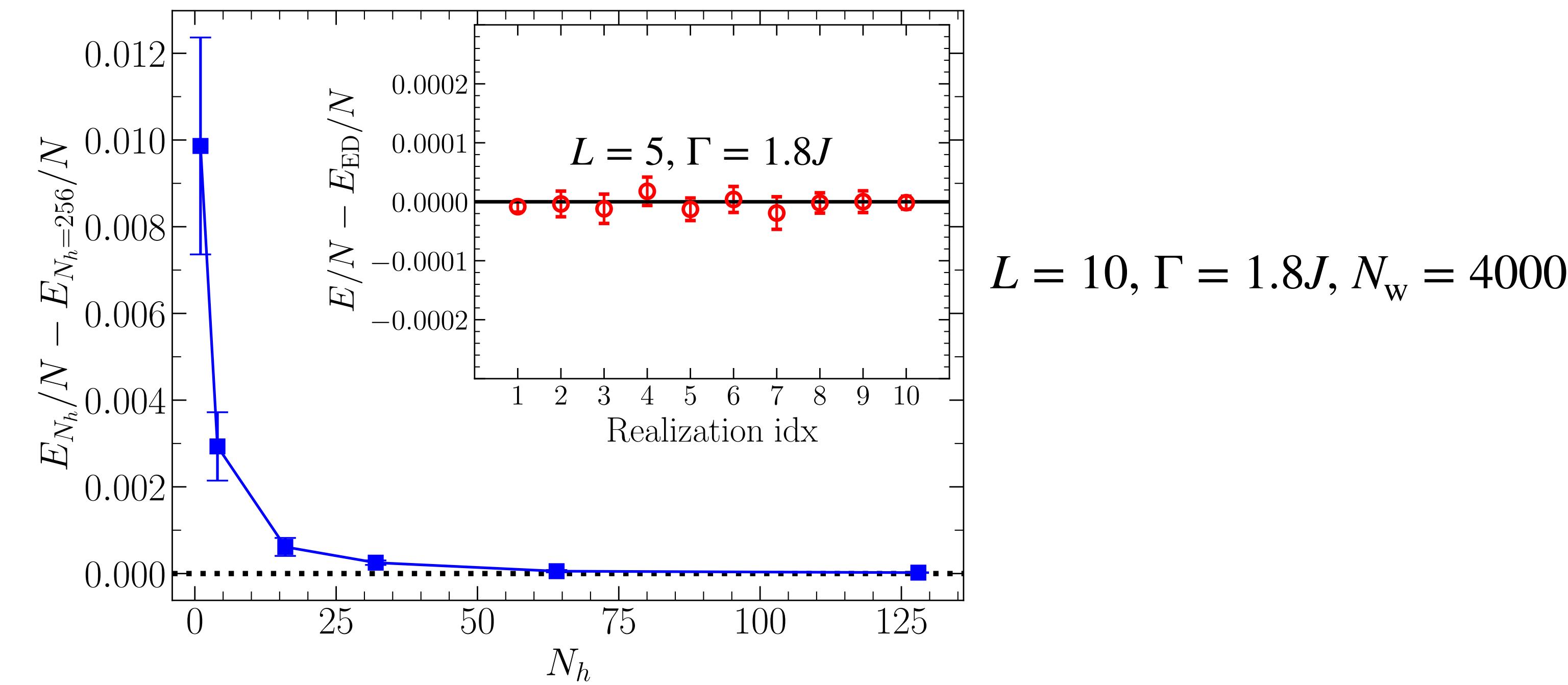
Exponential growth with the system size

No guiding wf

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Unbiased estimations of GS properties



$L = 10, \Gamma = 1.8J, N_w = 4000$

