Heavy hadron production with a coalescence plus fragmentation approach from AA to pp collisions

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***In collaboration with*:**
S. Plumari, M.L. Sambataro,Y.Sun, V.Greco

<u>Meeting SIM e PRIN2022</u>

Outline

Hadronization:

- **Fragmentation**
- **Coalescence model**

Heavy hadrons in AA collisions:

- **Λ_c , D spectra and ratio: RHIC and LHC**
  S. Plumari, et al. Eur. Phys. J. C 78, no.4, 348 (2018)

Heavy hadrons in small systems (pp @ 5.02 TeV):

- **Λ_c/D^0 , Ξ_c/D^0, Ω_c/D^0** *V. Minissale, et al., Phys. Lett. B 821, 136622 (2021)*
- **Λ_b/B^0, Ξ_b/B^0 , Ω_b/B^0** *(new) V. Minissale et al, arXiv:2405.19244 [hep-ph]*

Multicharm production *(new) V. Minissale, et al., Eur. Phys. J. C 84, no.3, 228 (2024)*

Specific of Heavy Quarks

- $m_{c,b} >> \Lambda_{QCD}$

produced by pQCD process (out of equilibrium)

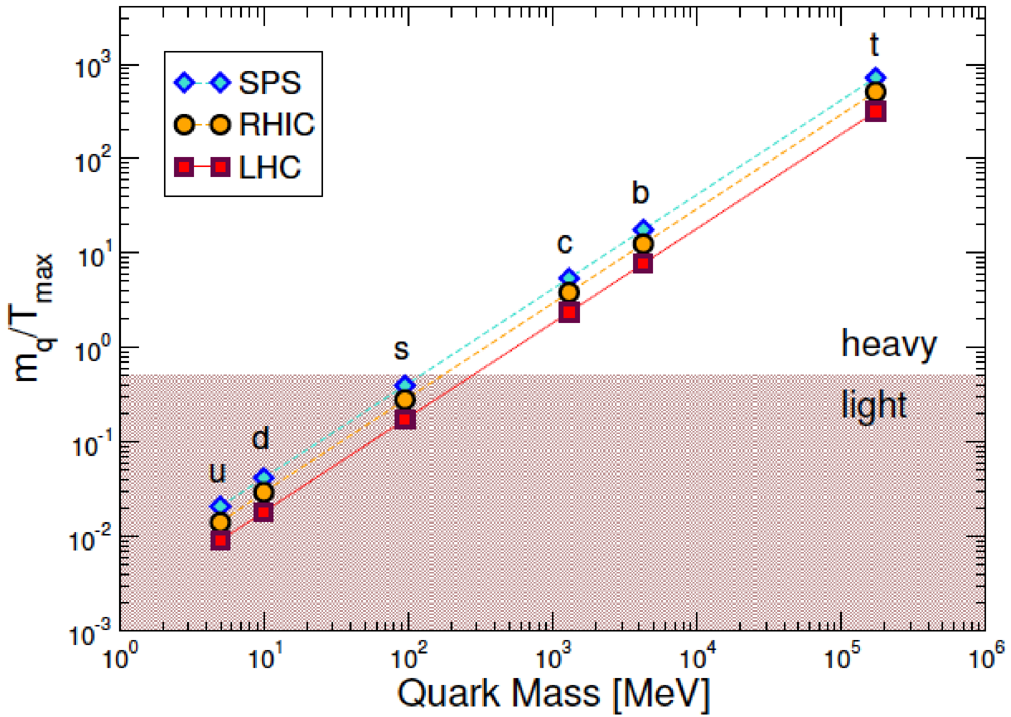
- $m_{c,b} >> T_0$

negligible thermal production

- $\tau_0 << \tau_{QGP}$
- $\tau_{therm.} \approx \tau_{QGP} >> \tau_{g,q}$

HQs experience the full QGP evolution

Carry informations about initial stages, more than light quarks



Recent reviews:

1) X.Dong, V. Greco Prog. Part. Nucl. Phys. 104 (2019)
2) A.Andronic Eur.Phys.J.C 76 (2016) 3, 107
3) F.Prino, R.Rapp, J.Phys.G 43 (2016) 9, 093002

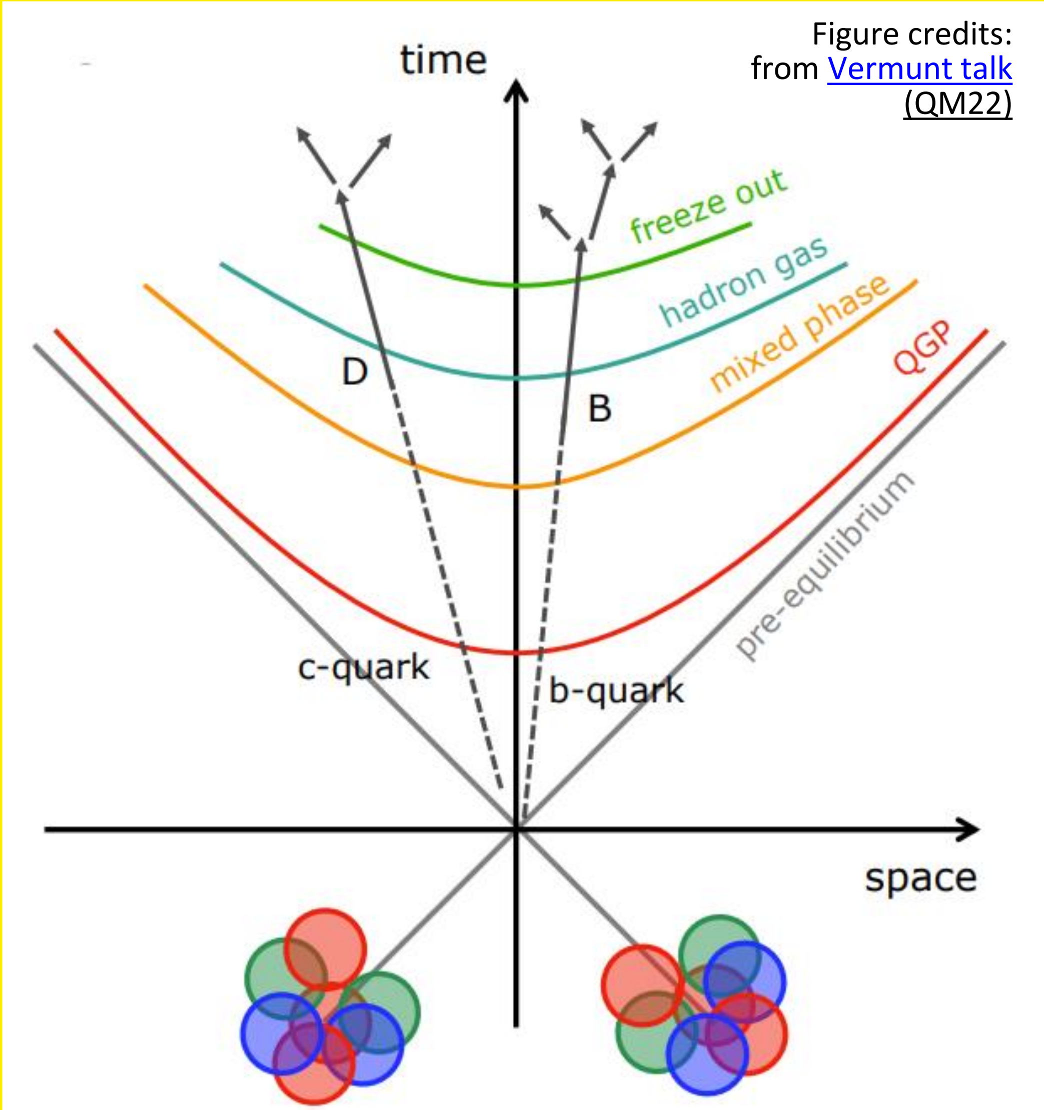


Figure credits: from Vermunt talk (QM22)

Relativistic Boltzmann transport at finite η/s

Bulk evolution

$$\underbrace{p^{\mu}\,\partial_{\mu} f_{q,g}(x,p)}_{\text{free-streaming}} + \underbrace{M(x)\,\partial^{x}_{\mu} M(x)\,\partial^{\mu}_{p} f_{q,g}(x,p)}_{\text{field interaction } \varepsilon-3p\neq 0} = \underbrace{C_{22}[f_{q,g}]}_{\text{collisions } \eta\neq 0}$$

free-streaming

field interaction
ε-3p≠0

collisions
η≠0

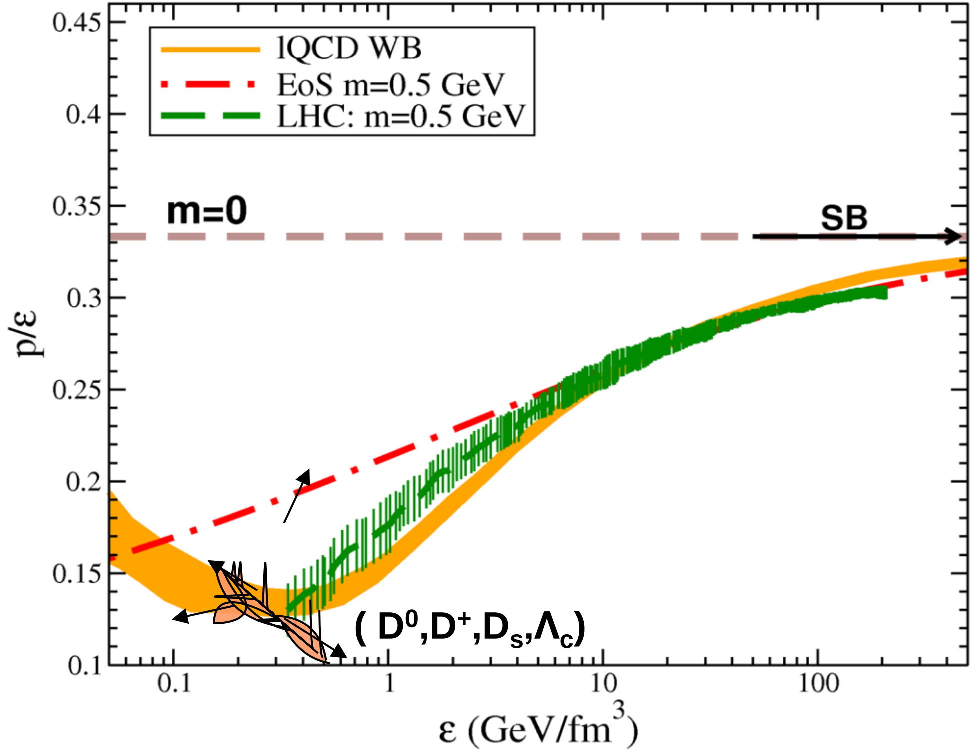
Heavy quark evolution

$$p^{\mu}\partial_{\mu} f_Q(x,p) = C[f_q, f_g, f_Q]$$

Describes the evolution of the one body distribution function f(x,p)

It is valid to study the evolution of both bulk and Heavy quarks

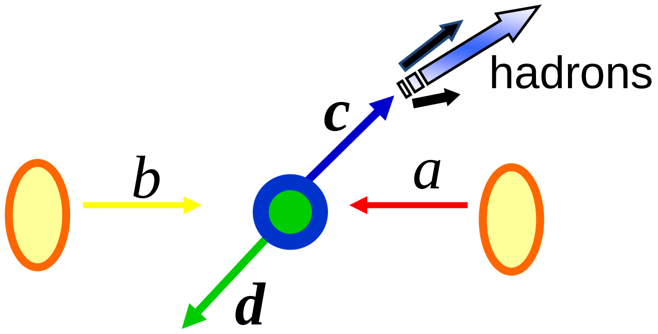
Possible to include f(x,p) out of equilibrium



S. Plumari et al.,J.Phys.Conf.Ser. 981 012017 (2018).

Heavy flavour Hadronization

Microscopic



Fragmentation:

production from hard-scattering processes (PDF+pQCD).
Fragmentation functions: data parametrization, assumed "universal"

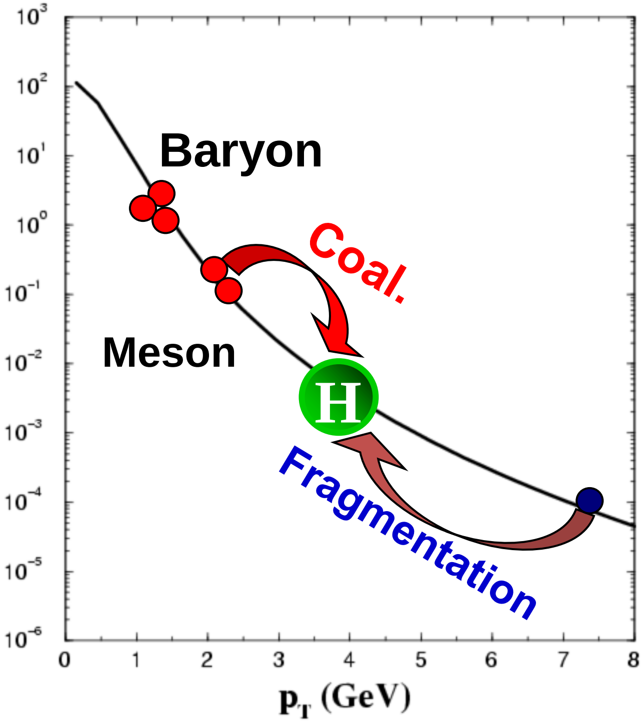
$$\sigma_{pp\to h} = PDF(x_a, Q^2)PDF(x_b, Q^2)\otimes\sigma_{ab\to q\bar{q}}\otimes D_{q\to h}(z, Q^2)$$

Parton shower: String fragmentation(Lund model – PYTHIA)
+colour reconnection(interaction from different scattering)
Cluster decay (HERWIG)

Coalescence: recombination of partons in QGP close in phase space

$$\frac{dN_{Hadron}}{d^2p_T} = g_H \int \prod_{i=1}^{n} p_i \cdot d\sigma_i \frac{d^3p_i}{(2\pi)^3} f_q(x_i, p_i) f_W(x_1, \dots, x_n; p_1, \dots, p_n)\,\delta\left(p_T - \sum_i p_{Ti}\right)$$

Have described first AA observations in light sector for the enhanced baryon/meson ratio and elliptic flow splitting



Macroscopic

Statistical hadronization:

Equilibrium + hadron-resonance gas + freeze-out temperature.
Production depends on hadron masses and degeneracy, and on system properties.
pQCD Charm production + total yield from charm cross section (not Temp.)
charm hadrons according to thermal weights

Catania Model: Coalescence + Fragmentation

Statistical factor colour-spin-isospin

Parton Distribution function

Hadron Wigner function

$C_H = \mathcal{N} f_H$

$$\frac{dN_{Hadron}}{d^2p_T} = g_H \int \prod_{i=1}^{n} p_i \cdot d\sigma_i \frac{d^3p_i}{(2\pi)^3} f_q(x_i, p_i) C_H(x_1, \ldots, x_n; p_1, \ldots, p_n) \delta\left(p_T - \sum_i p_{iT}\right)$$

Wigner function – Wave function

Wigner function **width** fixed by root-mean-square charge radius from **quark model**

$\sigma_{ri} = 1/\sqrt{\mu_i \omega}$

C.-W. Hwang, EPJ C23, 585 (2002)
C. Albertus et al., NPA 740, 333 (2004)

$$\Phi_M^W(\boldsymbol{r}, \boldsymbol{q}) = \int d^3r' \; e^{-i\boldsymbol{q}\cdot\boldsymbol{r}'} \phi_M\left(\boldsymbol{r} + \frac{\boldsymbol{r}'}{2}\right) \phi_M^*\left(\boldsymbol{r} - \frac{\boldsymbol{r}'}{2}\right)$$

$\phi_M(\boldsymbol{r})$ *meson wave function*

Assuming gaussian wave function

$$f_H(\ldots) = \prod_{i=1}^{N_q - 1} 8 \exp\left(\frac{-x_{ri}^2}{\sigma_{ri}^2} - p_{ri}^2 \sigma_{ri}^2\right)$$

only one width coming from $\phi_M(\mathbf{r})$, constraint $\sigma_r \sigma_p = 1$

Meson	$\langle r^2 \rangle_{ch}$	σ_{p1}	σ_{p2}
$D^+ = [c\bar{d}]$	0.184	0.282	—
$D_s^+ = [\bar{s}c]$	0.083	0.404	—
Baryon	$\langle r^2 \rangle_{ch}$	σ_{p1}	σ_{p2}
$\Lambda_c^+ = [udc]$	0.15	0.251	0.424
$\Xi_c^+ = [usc]$	0.2	0.242	0.406
$\Omega_c^0 = [ssc]$	-0.12	0.337	0.53

$$\langle r^2 \rangle_{ch} = \frac{3}{2}\left(\frac{m_2^2 Q_1 + m_1^2 Q_2}{(m_1+m_2)^2}\sigma_{r1}^2 + \frac{m_3^2(Q_1+Q_2) + (m_1+m_2)^2 Q_3}{(m_1+m_2+m_3)^2}\sigma_{r2}^2\right)$$

Meson	$\langle r^2 \rangle_{ch}$	σ_{p1}	Baryon	$\langle r^2 \rangle_{ch}$	σ_{p1}	σ_{p2}
$B^-[b\bar{u}]$	-0.378	0.302	$\Lambda_b^0[udb]$	0.13	0.264	0.5
$\bar{B}^0[b\bar{d}]$	0.187	0.303	$\Xi_b^0[usb]$	0.16	0.279	0.527
$\bar{B}_s^0[b\bar{s}]$	0.119	0.374	$\Xi_b^-[dsb]$	-0.21	0.295	0.557
$B_c^-[b\bar{c}]$	-0.043	0.74	$\Omega_b^-[ssb]$	-0.18	0.318	0.592

Catania Model: Coalescence + Fragmentation

Statistical factor colour-spin-isospin

Parton Distribution function

Hadron Wigner function $C_H = \mathcal{N} f_H$

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Wigner function width fixed by root-mean-square charge radius from **quark model**

C.-W. Hwang, EPJ C23, 585 (2002)
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Wigner function – Wave function

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Assuming gaussian wave function

$$f_H(\dots) = \prod_{i=1}^{N_q-1} 8 \exp\left(\frac{-x_{ri}^2}{\sigma_{ri}^2} - p_{ri}^2\sigma_{ri}^2\right)$$

only one width coming from $\phi_M(\mathbf{r})$, constraint $\sigma_r\sigma_p = 1$

Normalization $\mathcal{N}$ of $C_H(\dots)$ requiring that $P_{coal}=1$ at $p_T=0$

The charm that does not coalesce undergo fragmentation

Catania Model: Coalescence + Fragmentation

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LIGHT

Thermal+flow for **u,d,s** ($p_T < 3$ GeV)

$$\frac{dN_{q,\bar{q}}}{d^2p_T} \sim \exp\left(-\frac{\gamma_T - p_T \cdot \beta_T \mp \mu_q}{T}\right)$$

$$\beta(r) = \frac{r}{R}\beta_{max}$$

$$V = \pi R^2 \tau \cosh(y_z), R(\tau_f) = R_0(1 + \beta_{max}\tau_f)$$

PbPb@5TeV(0-10%): $\tau_f = 8{,}5 \frac{fm}{c} \rightarrow V_{|y|<0{,}5} = 4500\ fm^3$

+quenched minijets for u,d,s ($p_T > 3$ GeV)

CHARM

In AA collisions charm distribution from the studies of R_{AA} and v_2 of **D-meson** to determine the Space Diffusion coefficient:

parton simulations solving relativistic Boltzmann transport equation

In pp collisions the charm distribution are the FONLL distribution

Coalescence simulation in a fireball with radial flow for light quarks → dimension set by experimental constraints

Catania Model: Coalescence + Fragmentation

Statistical factor colour-spin-isospin — *Parton Distribution function* — *Hadron Wigner function* — $C_H = \mathcal{N} f_H$

$$\frac{dN_{Hadron}}{d^2p_T} = g_H \int \prod_{i=1}^{n} p_i \cdot d\sigma_i \frac{d^3p_i}{(2\pi)^3} f_q(x_i, p_i) C_H(x_1, \ldots, x_n; p_1, \ldots, p_n) \, \delta\left(p_T - \sum_i p_{iT}\right)$$

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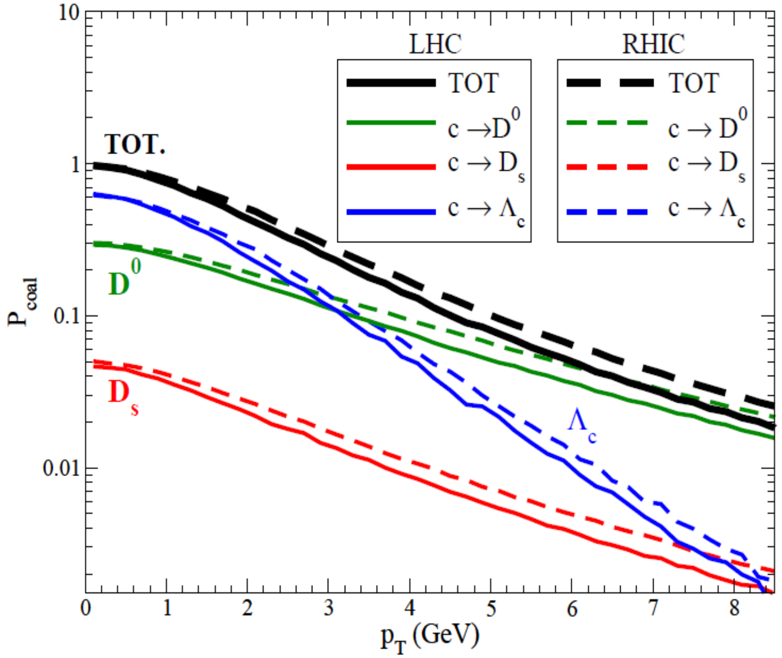
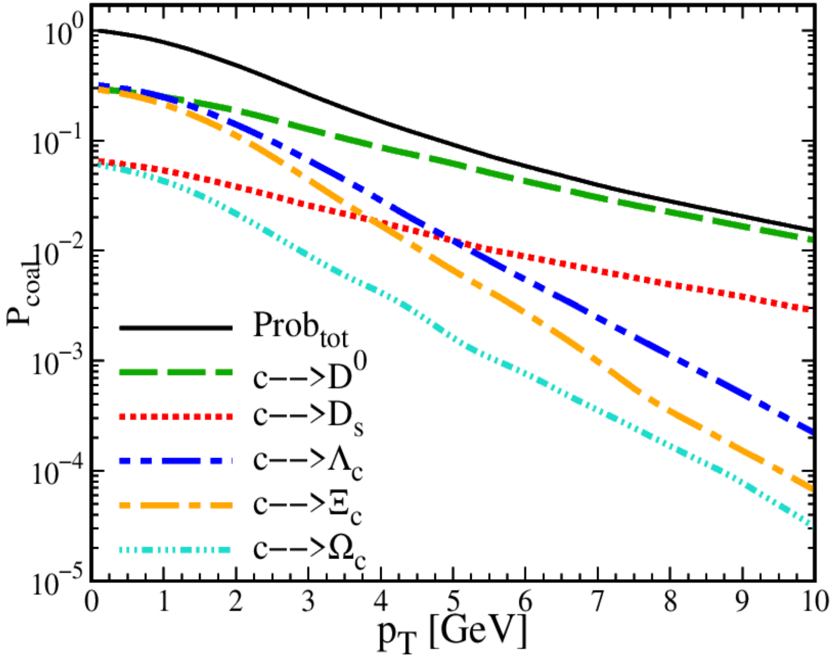
The charm that does not coalesce undergo fragmentation

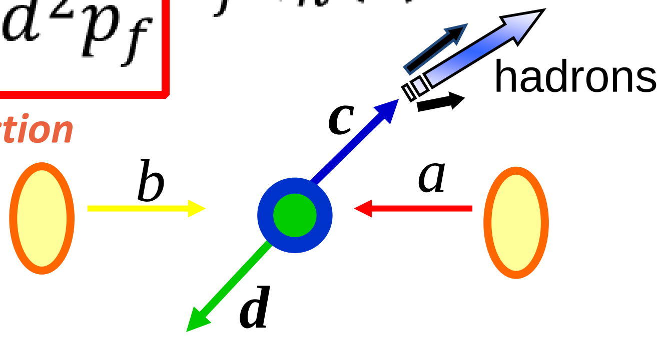




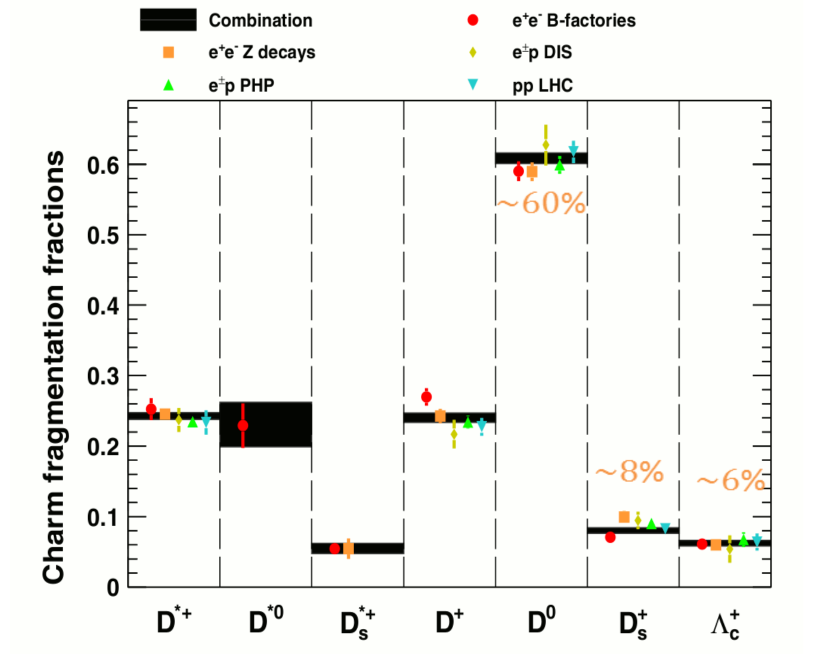
Catania Model: Coalescence + Fragmentation

$$\frac{dN_h}{d^2p_h} = \sum_f \int dz \frac{dN_f}{d^2p_f} D_{f\to h}(z)$$

Parton Distribution Function



The distribution function is evaluated at the Fixed-Order plus Next-to-Leading-Log (FONLL)
M. Cacciari, P. Nason, R. Vogt, PRL 95 (2005) 122001
In AA: bulk+charm evolution with Relativistic Transport Boltzmann Equation



M. Lisovyi, et al. EPJ C76 (2016) no.7, 397

We use the Peterson fragmentation function

C. Peterson, D. Schalatter, I. Schmitt, P.M. Zerwas PRD 27 (1983) 105

$$D_{f\to h}(z) \propto \frac{1}{z\left[1 - \frac{1}{z} - \frac{\epsilon}{1-z}\right]^2}$$

Sligthly modified to reproduce tail of the Λ_c/D^0

Charm Fragmentation Fraction (c→h)
Measurement in $e^{\pm} p$, $e^+ e^-$ *and old pp data*

$$\left(\frac{\Lambda_c^+}{D^0}\right)_{e^+e^-} \simeq 0.1 \qquad \left(\frac{D_s^+}{D^0}\right)_{e^+e^-} \simeq 0.13$$

Heavy flavour: Resonance decay

Meson	Mass(MeV)	I (J)	Decay modes	B.R.
$D^+ = \bar{d}c$	1869	$\frac{1}{2}\,(0)$		
$D^0 = \bar{u}c$	1865	$\frac{1}{2}\,(0)$		
$D_s^+ = \bar{s}c$	2011	$0\,(0)$		
Resonances				
D^{*+}	2010	$\frac{1}{2}\,(1)$	$D^0\pi^+$; D^+X	68%,32%
D^{*0}	2007	$\frac{1}{2}\,(1)$	$D^0\pi^0$; $D^0\gamma$	62%,38%
D_s^{*+}	2112	$0\,(1)$	D_s^+X	100%
Baryon				
$\Lambda_c^+ = udc$	2286	$0\,(\frac{1}{2})$		
$\Xi_c^+ = usc$	2467	$\frac{1}{2}\,(\frac{1}{2})$		
$\Xi_c^0 = dsc$	2470	$\frac{1}{2}\,(\frac{1}{2})$		
$\Omega_c^0 = ssc$	2695	$0\,(\frac{1}{2})$		
Resonances				
Λ_c^+	2595	$0\,(\frac{1}{2})$	$\Lambda_c^+\pi^+\pi^-$	100%
Λ_c^+	2625	$0\,(\frac{3}{2})$	$\Lambda_c^+\pi^+\pi^-$	100%
Σ_c^+	2455	$1\,(\frac{1}{2})$	$\Lambda_c^+\pi$	100%
Σ_c^+	2520	$1\,(\frac{3}{2})$	$\Lambda_c^+\pi$	100%
$\Xi_c^{\prime +,0}$	2578	$\frac{1}{2}\,(\frac{1}{2})$	$\Xi_c^{+,0}\gamma$	100%
Ξ_c^+	2645	$\frac{1}{2}\,(\frac{3}{2})$	$\Xi_c^+\pi^-$,	100%
Ξ_c^+	2790	$\frac{1}{2}\,(\frac{1}{2})$	$\Xi_c^{\prime}\pi$,	100%
Ξ_c^+	2815	$\frac{1}{2}\,(\frac{3}{2})$	$\Xi_c^{\prime}\pi$,	100%
Ω_c^0	2770	$0\,(\frac{3}{2})$	$\Omega_c^0\gamma$,	100%

In our calculations we take into account hadronic channels including the ground states + first excited states

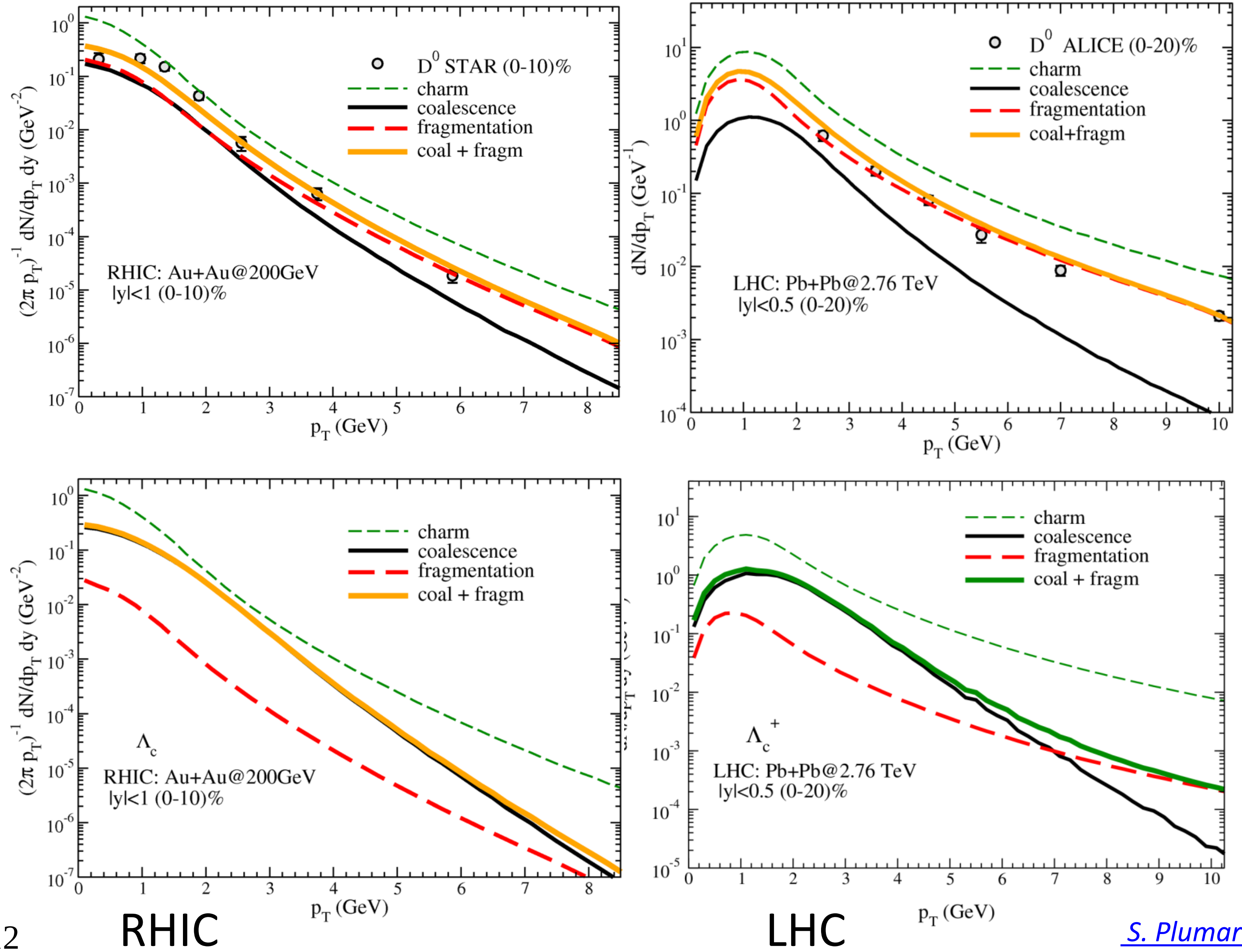
Statistical factor suppression for resonances

$$\frac{|(2J+1)(2I+1)|_{H^*}}{|(2J+1)(2I+1)|_{H}}\left(\frac{m_{H^*}}{m_H}\right)^{3/2} e^{-(m_{H^*}-m_H)/T}$$

AA @ RHIC & LHC

wave function widths σ_p of baryon and mesons are the same at RHIC and LHC!

Data from: STAR Coll. PRL 113, 142301 (2014), ALICE Coll. JHEP 09 (2012) 112



D^0

Coalescence lower at LHC than at RHIC

main contribution: Fragmentation

Λ_c

Coalescence lower at LHC than at RHIC

main contribution: Coalescence

S. Plumari, V. Minissale et al., Eur. Phys. J. C78 no. 4, (2018) 348

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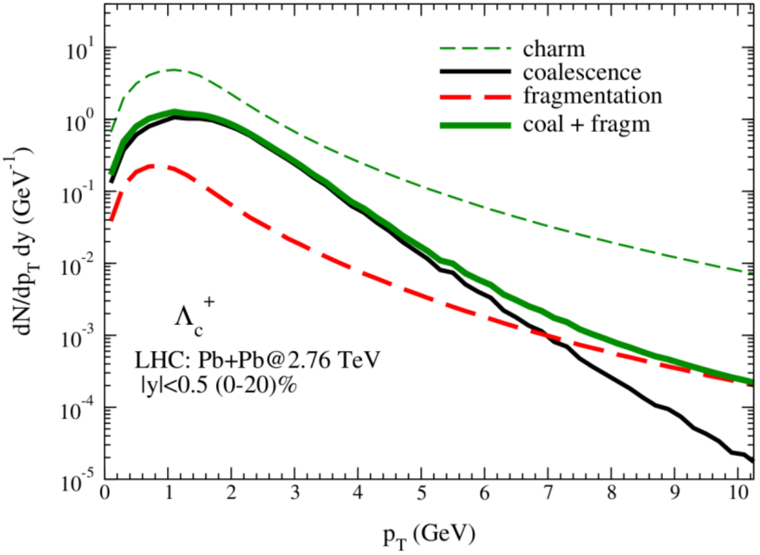
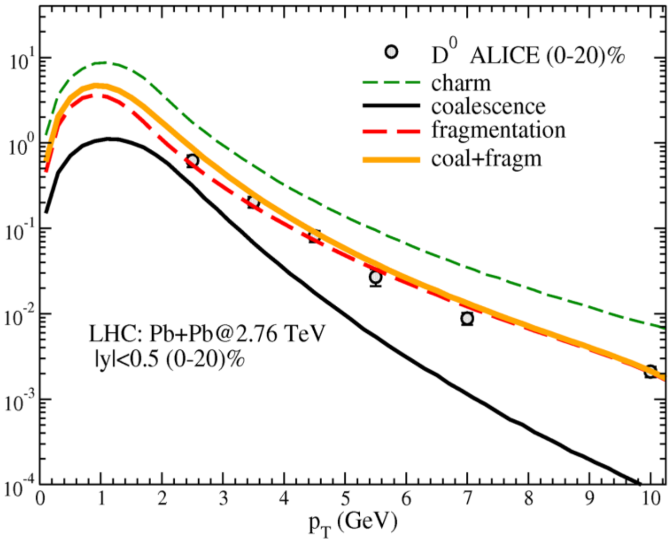
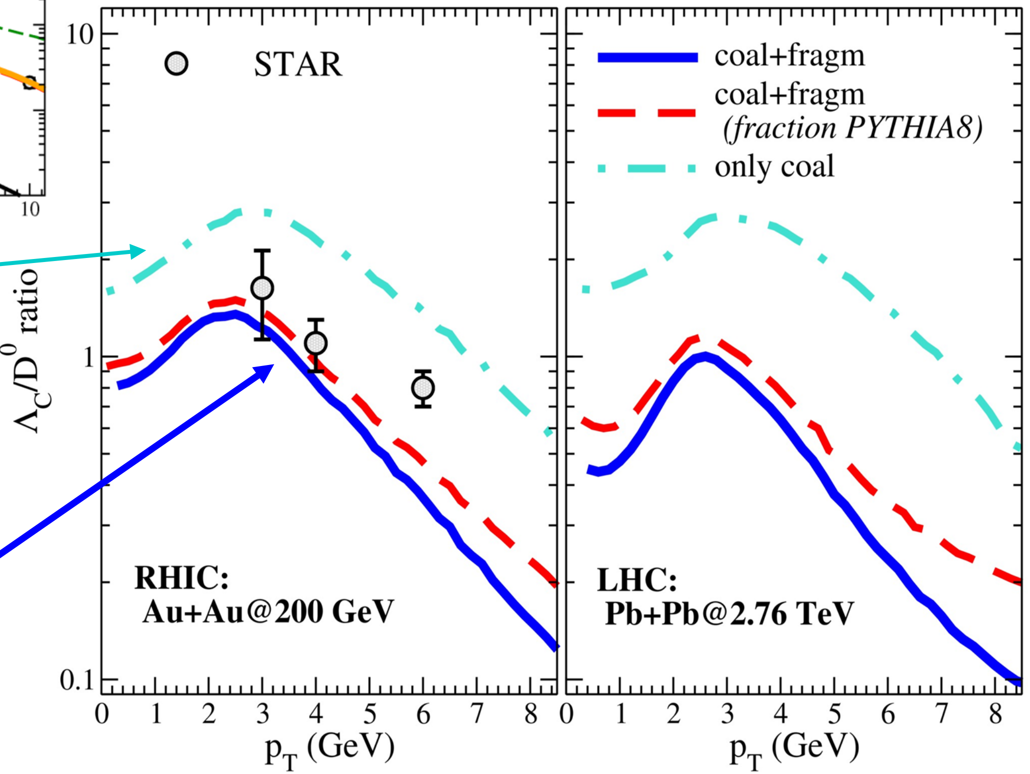




Coalescence lower at LHC than at RHIC



Only Coalescence ratio is similar at both energies.

Fragmentation ~ 0.1 at both energies.

the **combined ratio is different** because the <u>coalescence over fragmentation ratio at LHC is smaller than at RHIC</u>

Therefore at LHC the larger contribution in particle production from fragmentation leads to a final ratio that is smaller than at RHIC.

STAR Coll., Phys.Rev.Lett. 124 (2020) 17, 172301

S. Plumari, V. Minissale et al., Eur. Phys. J. C78 no. 4, (2018) 348

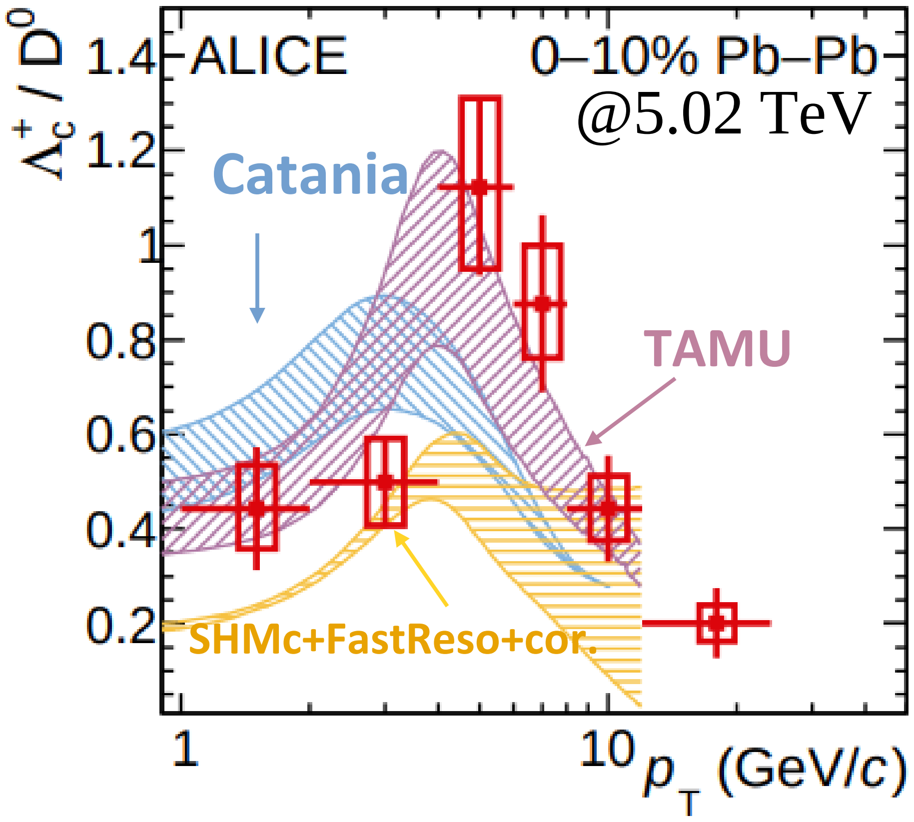
AA @ RHIC & LHC

wave function widths σ_p of baryon and mesons are the same at RHIC and LHC!

Results for 0-10% in PbPb @5.02TeV:

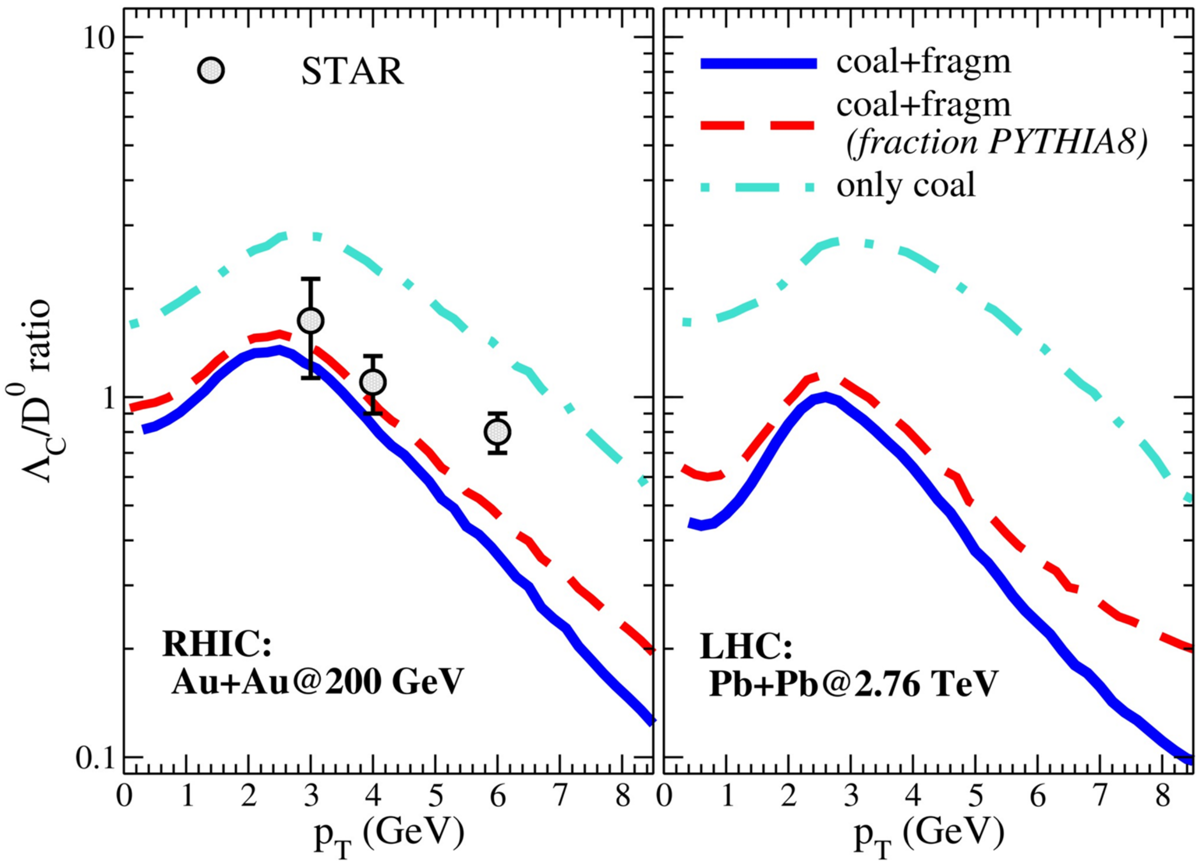
Consistent with the trend shown at RHIC and LHC @2.76TeV

Available data at low p_T → differences recombination vs SHM



ALICE Coll. Phys. Lett. B 839 (2023) 137796

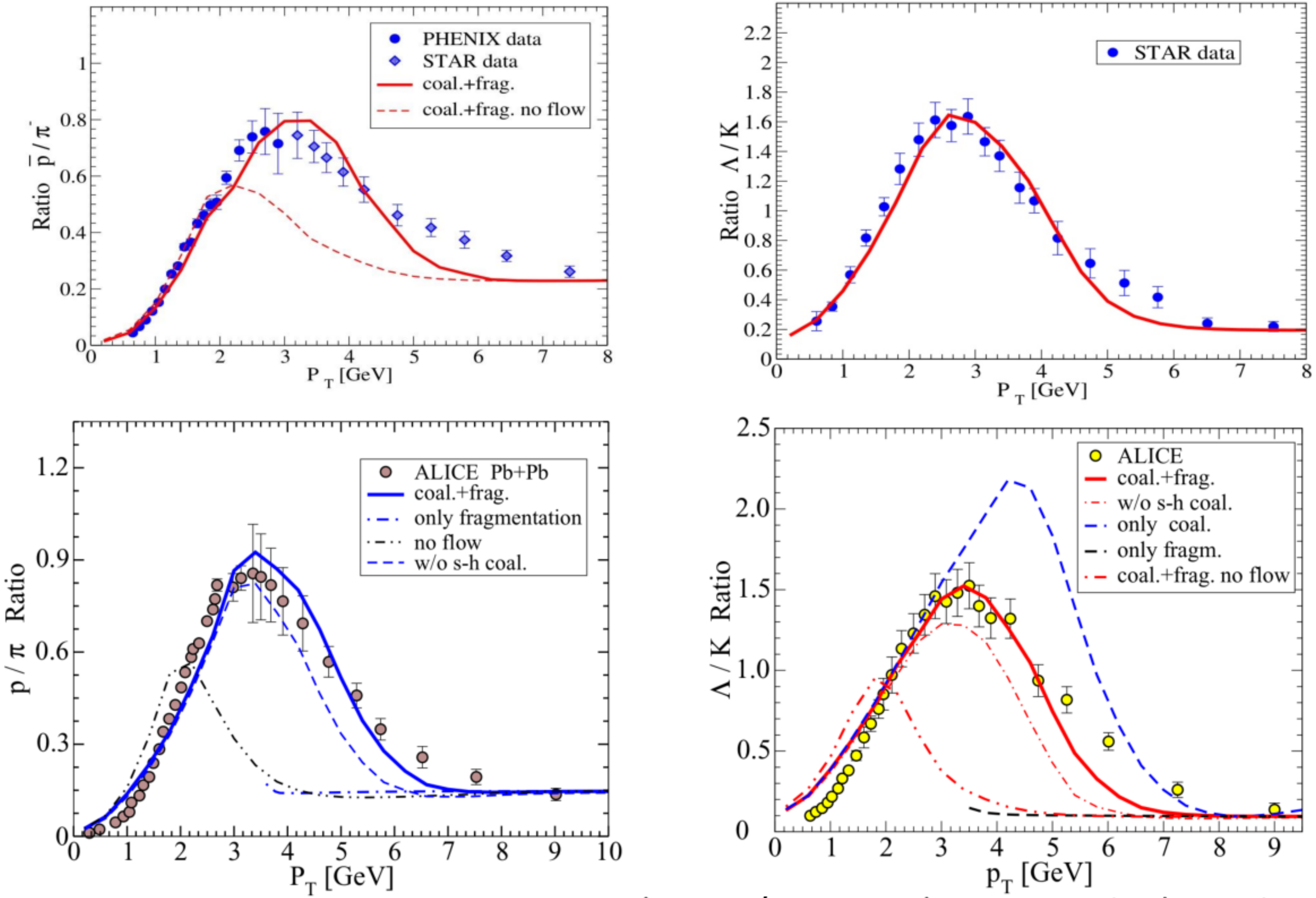
STAR Coll., Phys.Rev.Lett. 124 (2020) 17, 172301



S. Plumari, V. Minissale et al., Eur. Phys. J. C78 no. 4, (2018) 348

Baryon to meson ratio at RHIC & LHC

Minissale, Scardina, Greco, Phys.Rev. C 92 (2015) 5,054904



RHIC

LHC

- coalescence naturally predict a baryon/meson enhancement in the region $p_T \simeq$ 2–4GeV with respect to pp collisions
- Lack of baryon yield in the region $p_T \simeq$ 5–7GeV

Heavy flavour Hadronization

Fragmentation:
production from hard-scattering processes (PDF+pQCD).
Fragmentation functions: data parametrization, assumed "universal"

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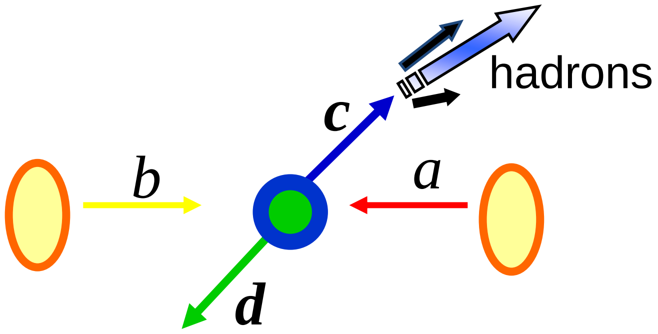
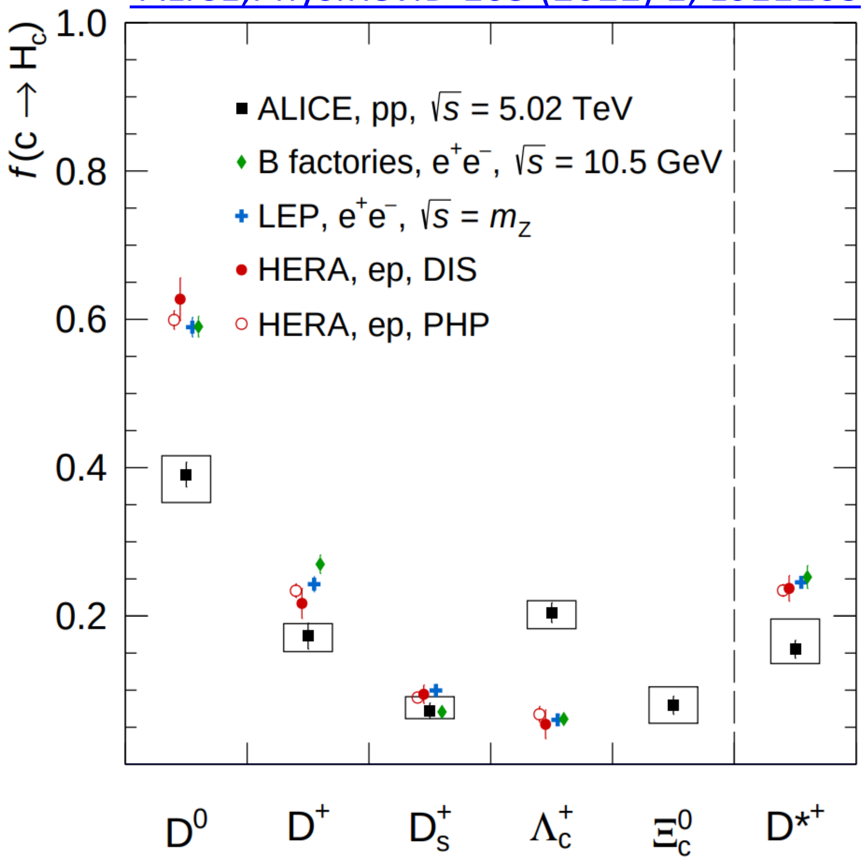


ALICE,Phys.Rev.D 105 (2022) 1, L011103



Things get more complicated after experimental evidence in pp@5TeV:

Fragmentation fractions (c→h) depends on collision system...*and QGP presence*?

No more Universality?

Heavy flavour Hadronization

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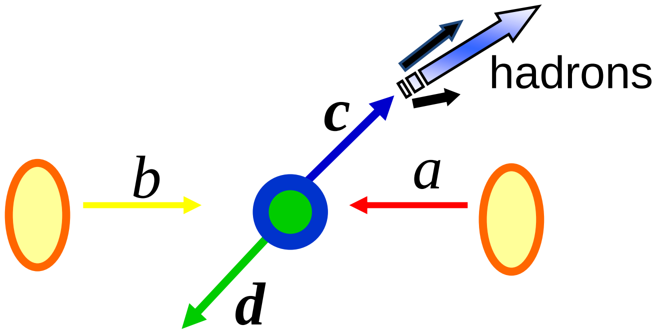


ALICE, PRL 127 202301 (2021)
ALICE, PRC 104 054905 (2021)

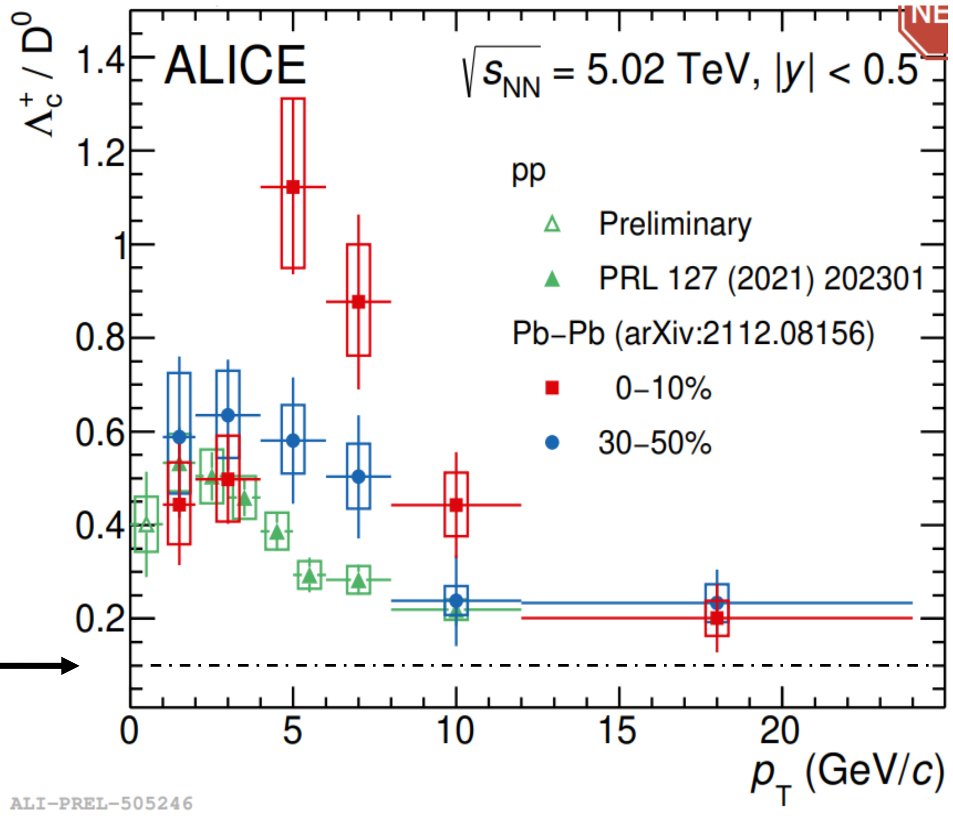
Things get more complicated after experimental evidence in pp@5TeV:

Fragmentation fractions (c→h) depends on collision system...*and QGP presence*?

No more Universality?

Baryon/meson ratio is underestimated, and no p_T dependence

$$\left(\frac{\Lambda_c^+}{D^0}\right)_{e^+e^-} \simeq 0.1$$



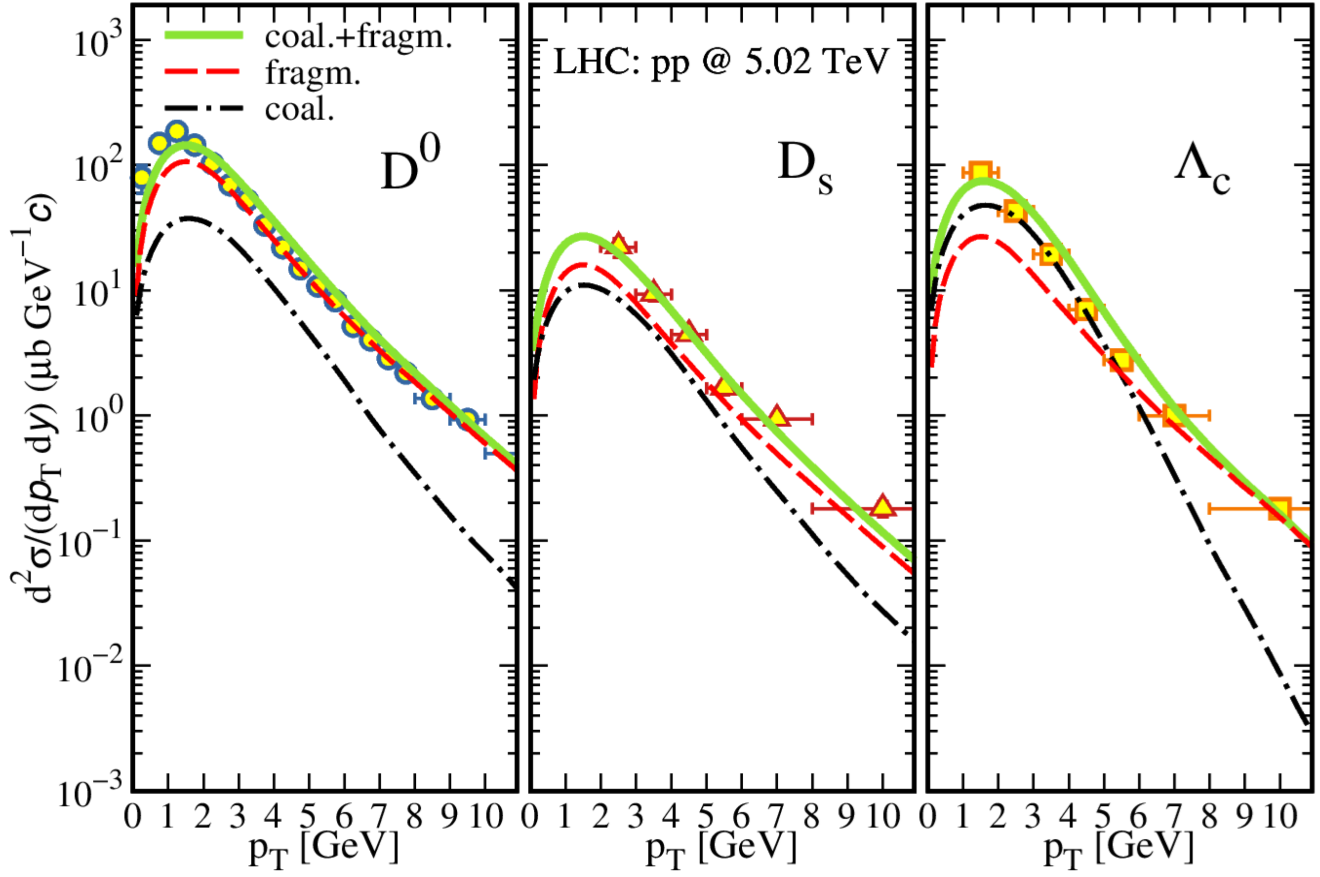
Small systems: Coalescence in pp?

Common consensus of possible presence of QGP in smaller system.

What if:

- **Assuming QGP formation also in pp?**
- **What coalescence+fragmentation predicts in this case?**

Data from:
S. Acharya et al. (ALICE), Eur. Phys. J. C 79, 388 (2019)
ALICE Coll., Phys.Rev.Lett. 127 (2021) 20, 202301 - Phys.Rev.C 104 (2021) 5, 054905



V. Minissale, S. Plumari, V. Greco, Physics Letters B 821 (2021) 136622

If we assume in p+p @ 5 TeV a medium similar to the one simulated in hydro:

p+p @ 5 TeV

- τ_{pp}=2 fm/c
- β_0=0.4
- R=2.5 fm
- V~30 fm^3

LIGHT

- Thermal Distribution (p_T< 2 GeV)

$$\frac{dN_q}{d^2r_T d^2p_T} = \frac{g_g \tau m_T}{(2\pi)^3} \exp\left(-\frac{\gamma_T(m_T - p_T \cdot \beta_T)}{T}\right)$$

- Minijet Distribution (p_T> 2 GeV)
  NO QUENCHING

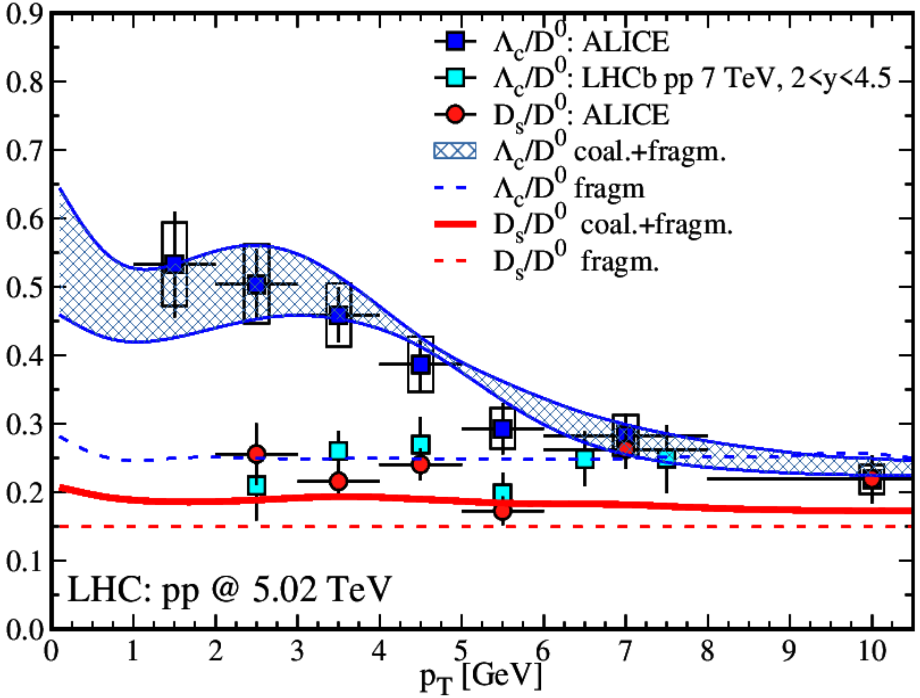
CHARM

FONLL Distribution

wave fuction widths σ_p of baryon and mesons kept the same from AA to pp

Small systems: Coalescence in pp?

V. Minissale, S. Plumari, V. Greco, Physics Letters B 821 (2021) 136622



Error band correspond to <r^2> uncertainty in quark model

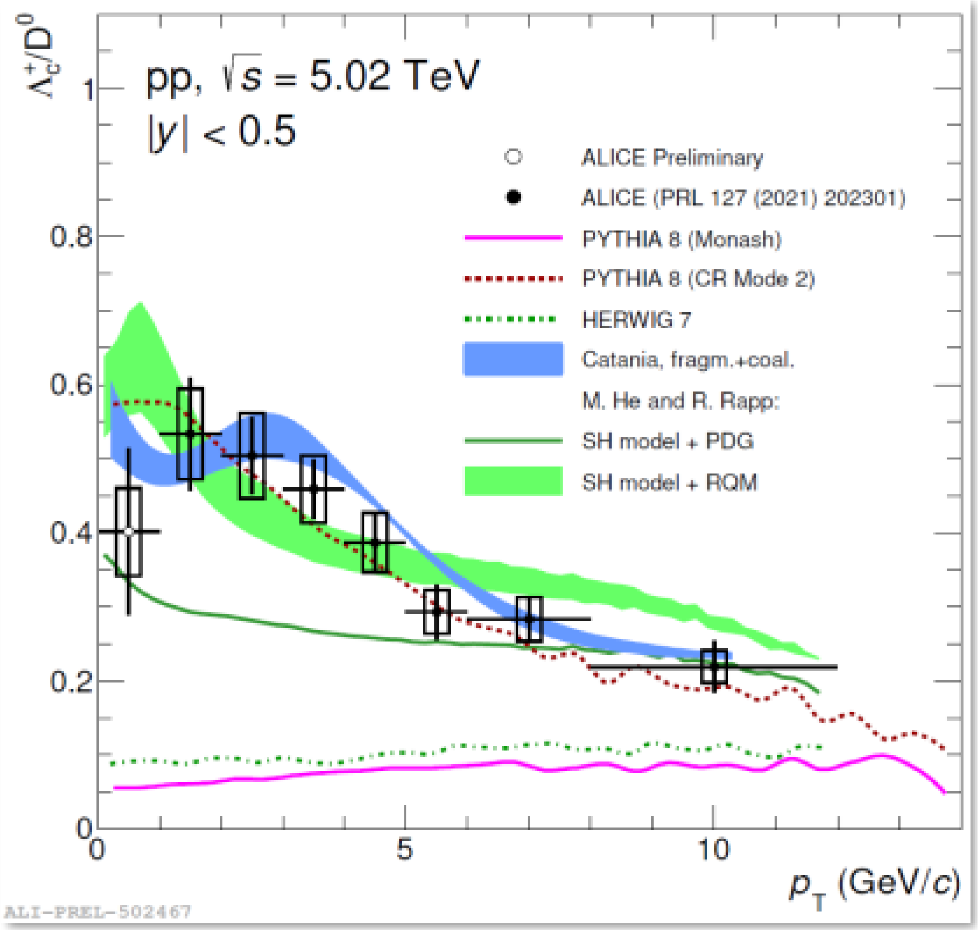
Reduction of rise-and-fall behaviour in Λ_c / D^0 ratio:

-Confronting with AA: Coal. contribution smaller w.r.t. Fragm.

-FONLL distribution flatter w/o evolution trough QGP

-Volume size effect

The increase of Λ_c production in pp have effect on R_{AA} of Λ_c

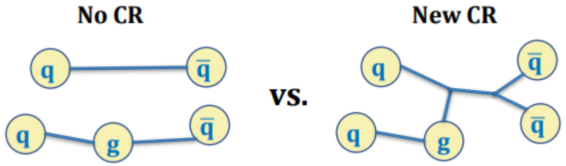


Other models:

He-Rapp, Phys.Lett.B 795 (2019) 117-121: Increase≈2 to Λ_c production: SHM with resonance not present in PDG

PYTHIA8 + color reconnection

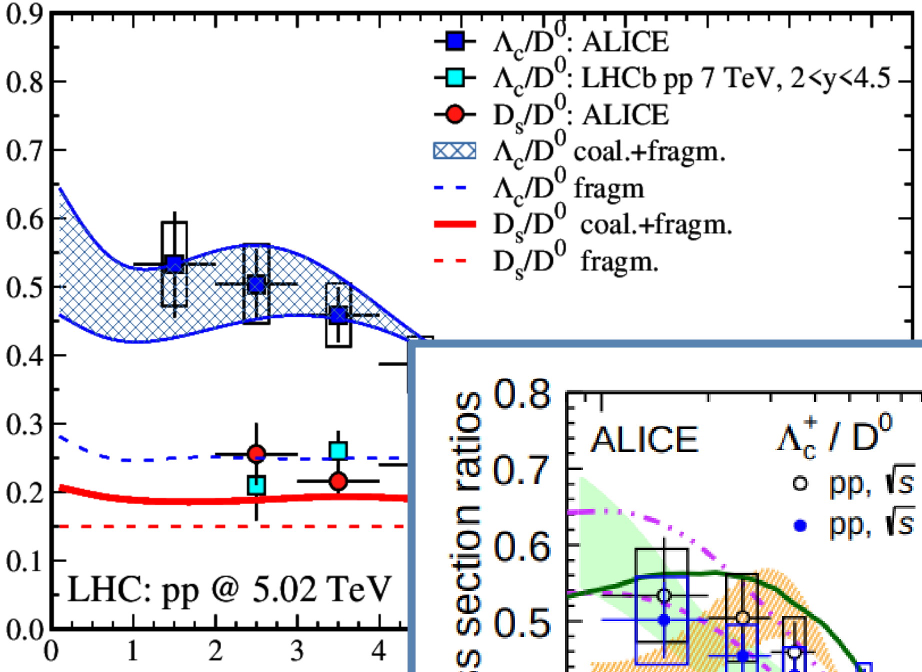
CR with SU(3) weights and string length minimization



ALICE, Phys.Rev.Lett. 127 (2021) 20, 202301

Small systems: Coalescence in pp?

V. Minissale, S. Plumari, V. Greco, Physics Letters B 821 (2021) 136622



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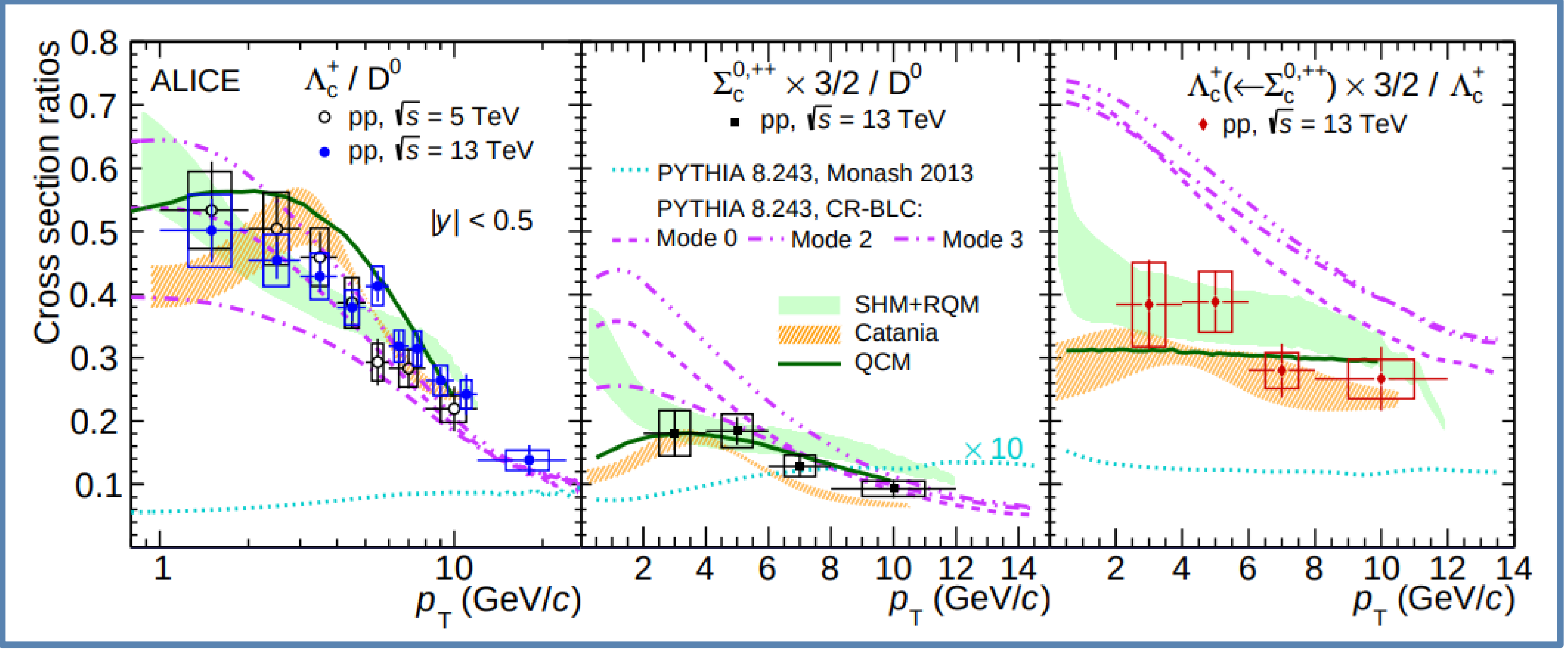
-FONLL distribution flatter w/o evolution trough QGP

-Volume size effect

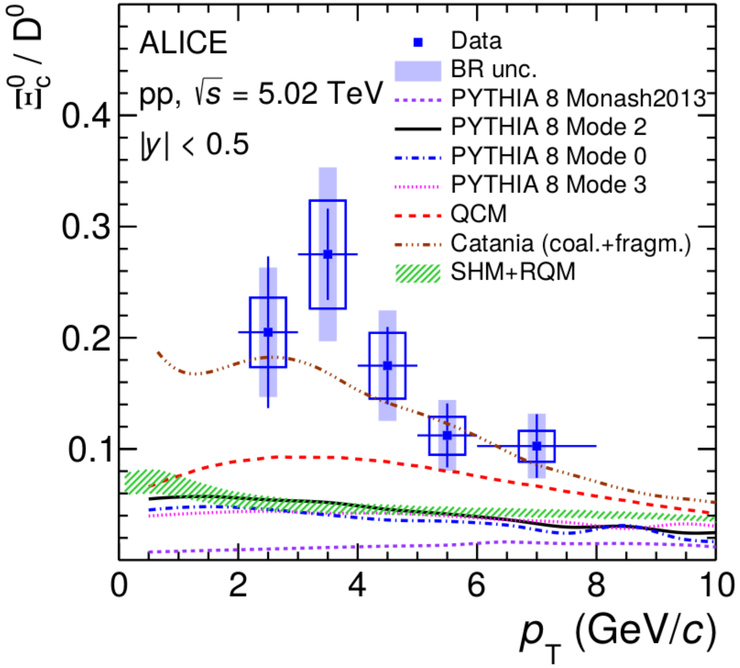
The increase of Λ_c production in pp have effect on R_{AA} of Λ_c

Error band correspond to



ALICE Coll., Physical Review Letters 128, 012001 (2022)

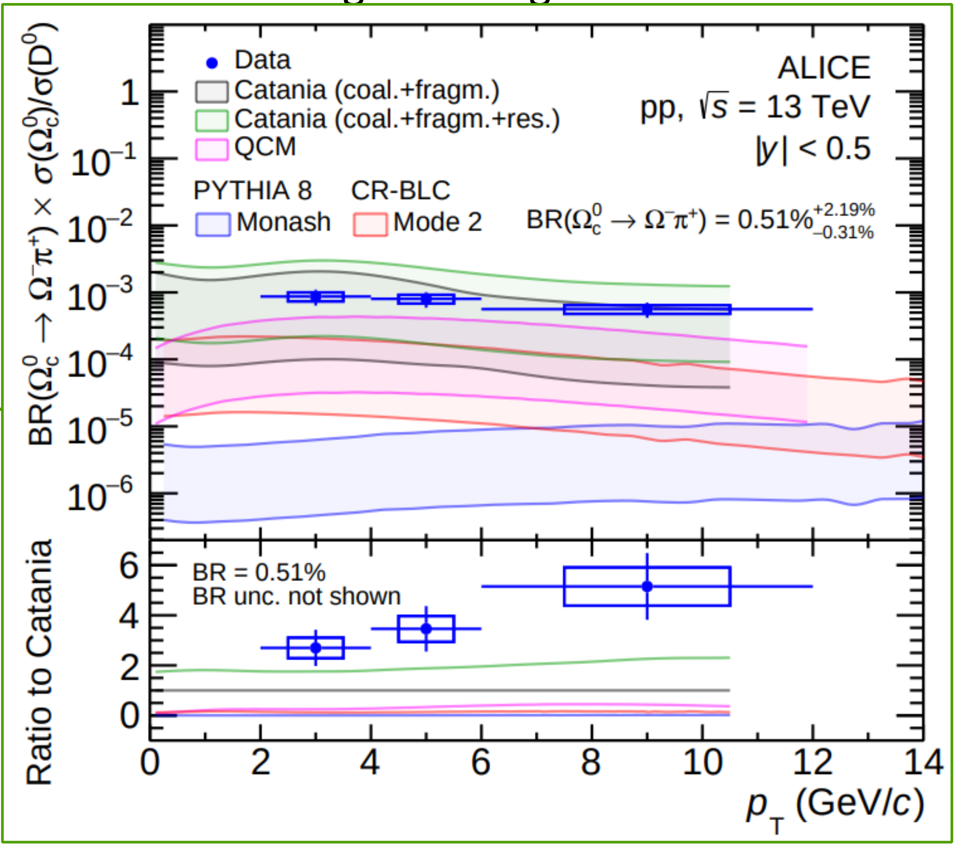
Small systems: Coalescence in pp?



New measurements of heavy hadrons at ALICE:

- Ξ_c/D^0 ratio, same order of Λ_c/D^0: coalescence gives enhancement
- very large Ω_c/D^0 ratio, our model does not get the big enhancement

Uncertainties bands coming from the Branching Ratio error

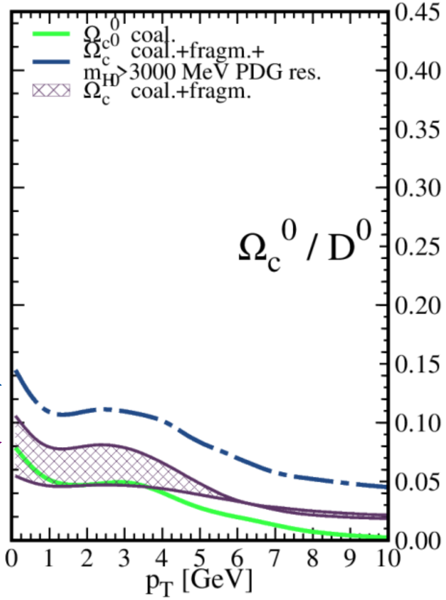


Assuming additional PDG resonances with J=3/2 and decay to Ω_c additional to $\Omega_c^0(2770)$

$\Omega_c^0(3000), \Omega_c^0(3005), \Omega_c^0(3065), \Omega_c^0(3090), \Omega_c^0(3120)$

supply an idea of how these states may affect the ratio

Error band correspond to $\langle r^2 \rangle$ uncertainty in quark model

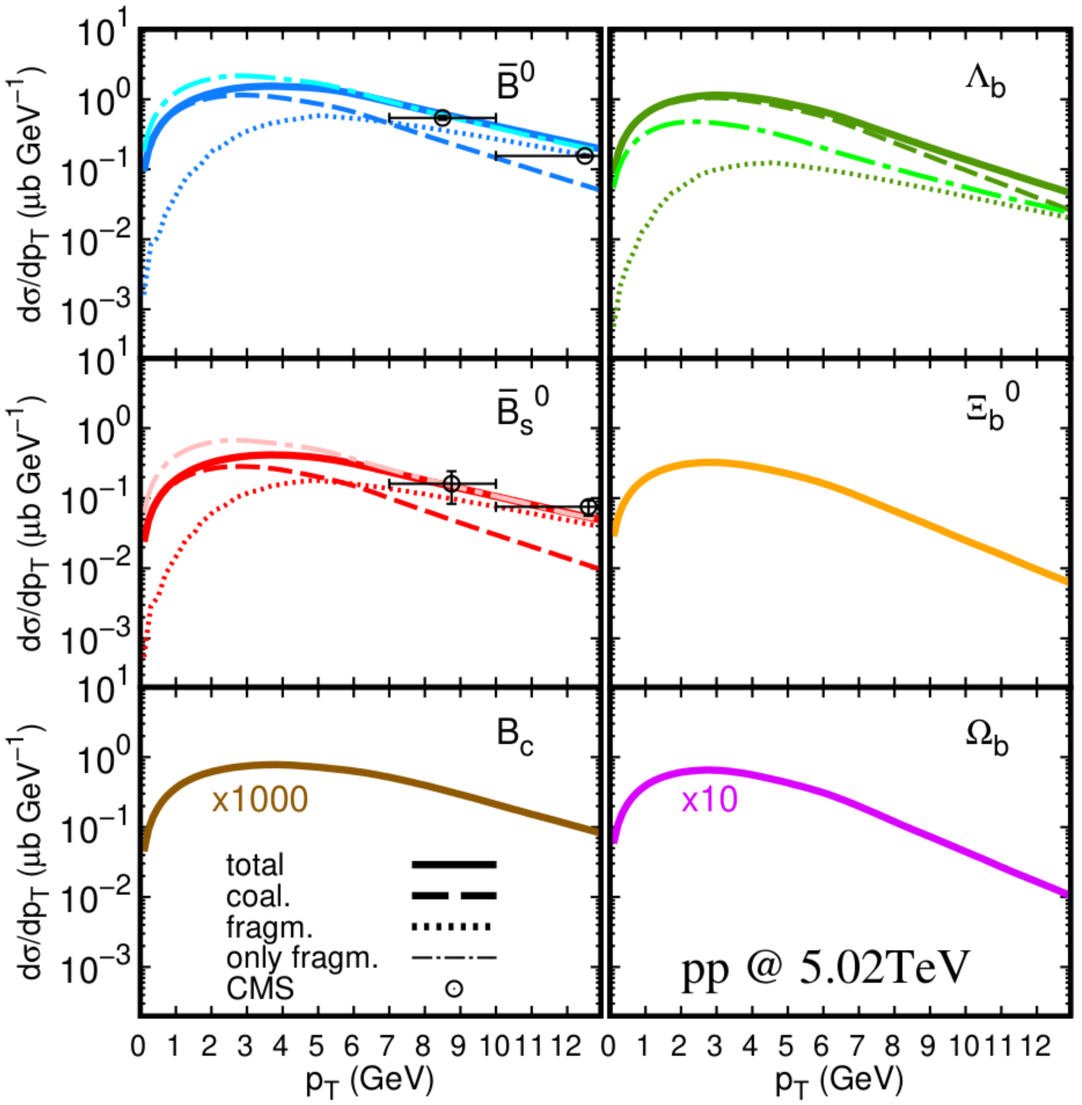


ALICE Coll. JHEP 10 (2021) 159
ALICE Coll. arXiv:2205.13993
V. Minissale, S. Plumari, V. Greco, Physics Letters B 821 (2021) 136622

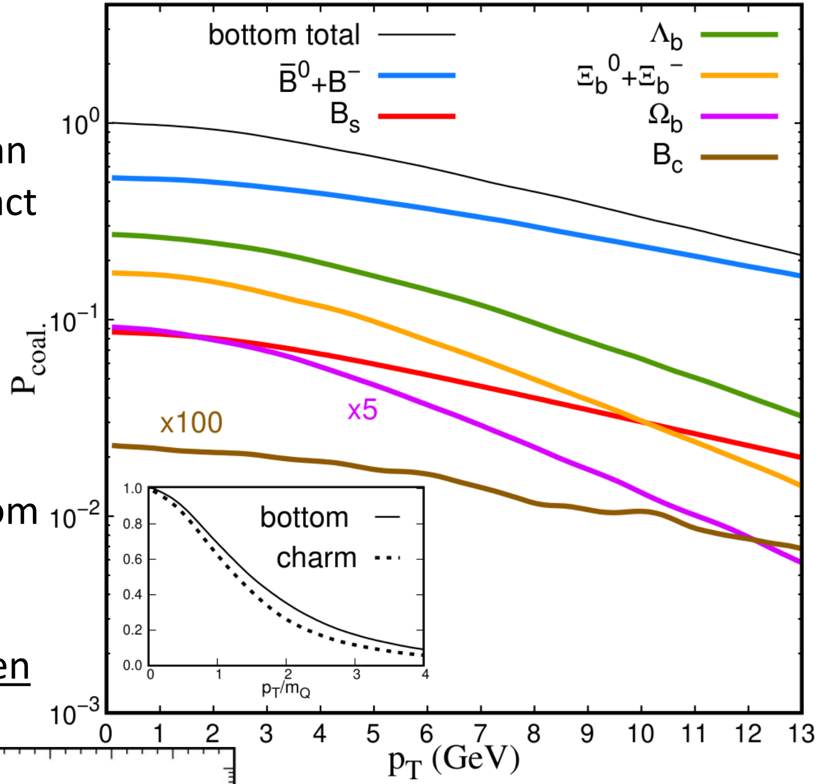
Small systems: Coalescence in pp for bottom hadron?

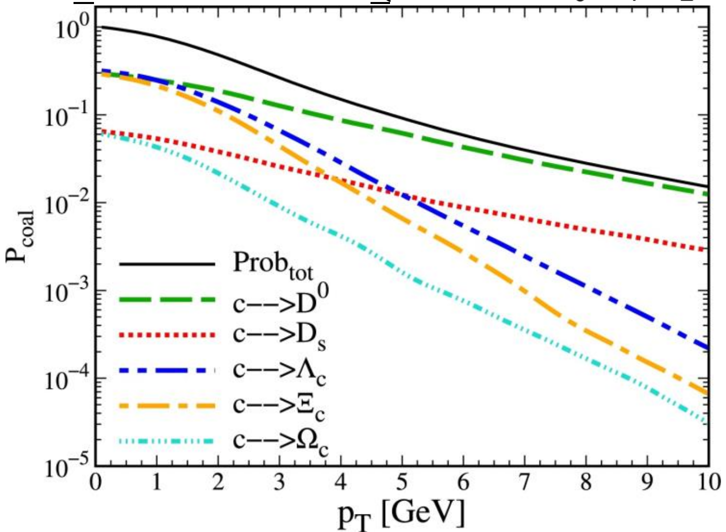
V. Minissale et al, arXiv:2405.19244

Data from: A. M. Sirunyan et al. (CMS), PRL 119, 152301 (2017).



- Bottom P_{coal} is flatter than charm → Coal. big impact on bottom hadron production in a wider p_T region
- B meson at $p_T < 5$ GeV mainly from Coal
- Λ_b production mainly from Coal. for $p_T < 10$ GeV
- Similar probabilities when p_T is scaled with m_Q

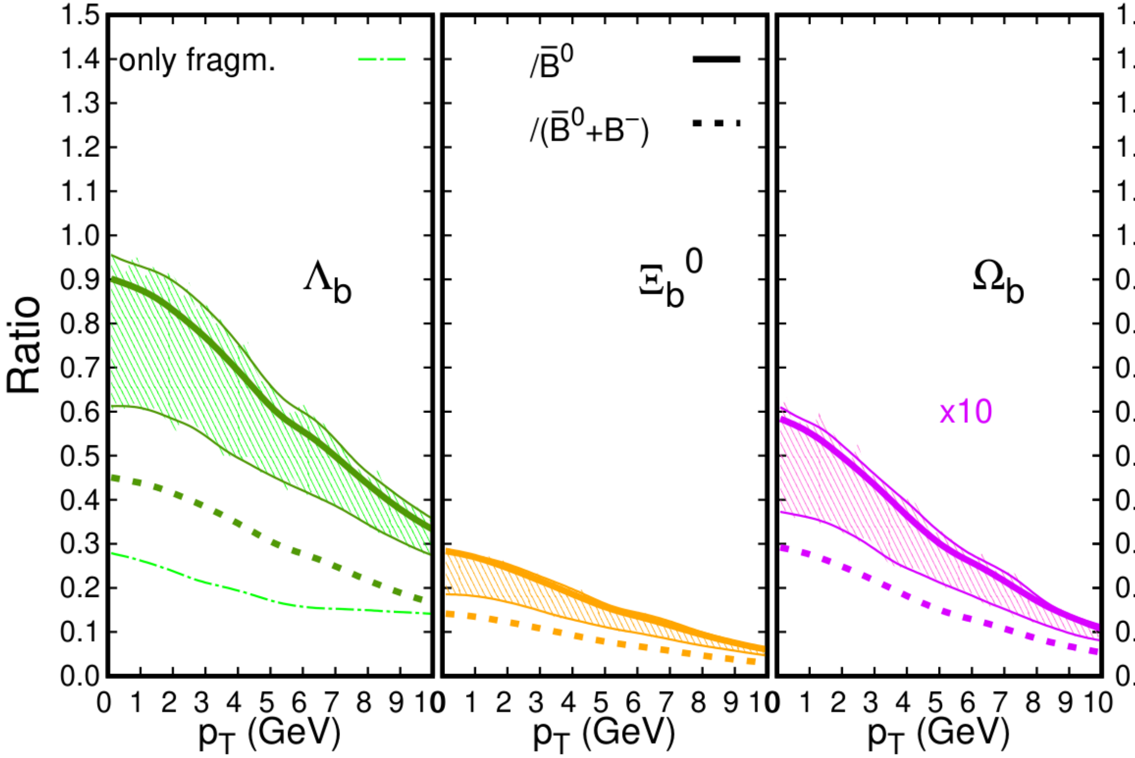




 Kartvelishvili fragm. function Phys.Lett.B 78 (1978) 615-617

Small systems: Coalescence in pp for bottom hadron?

V. Minissale et al, arXiv:2405.19244



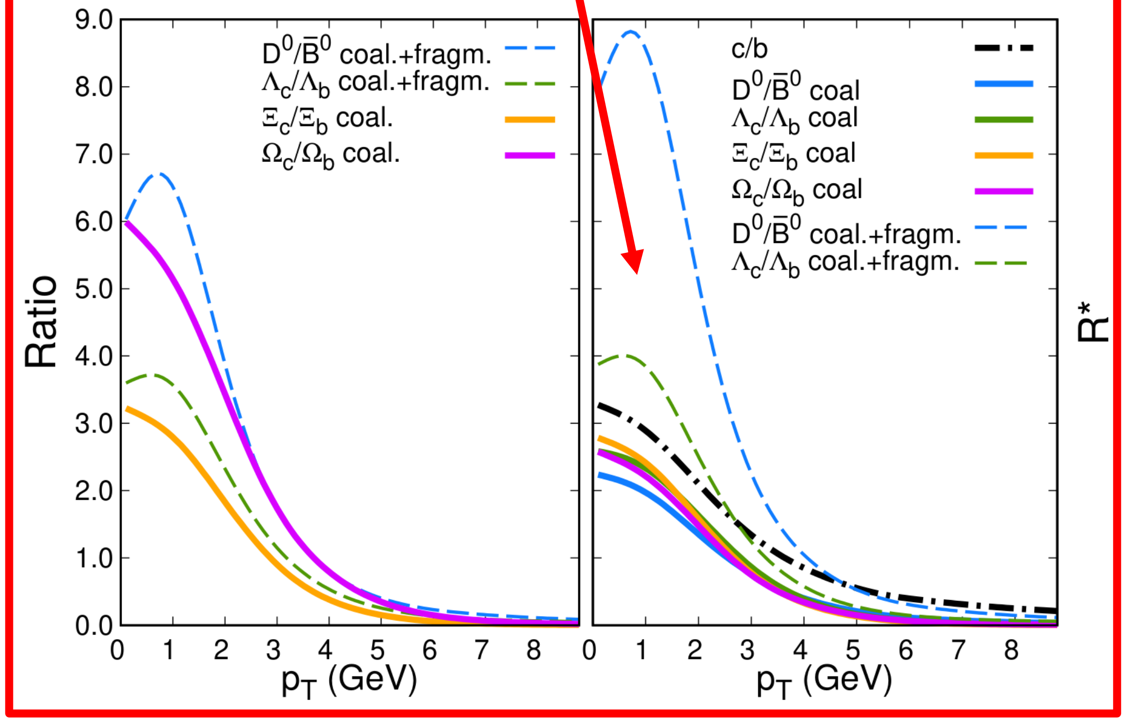
Error band correspond to $<r^2>$ uncertainty in quark model

Coal gives enhancement of Baryon/meson ratio

$$g_{H^*} = \frac{|(2J+1)(2I+1)|_{H^*}}{|(2J+1)(2I+1)|_H}\left(\frac{m_{H^*}}{m_H}\right)^{3/2} e^{-(m_{H^*}-m_H)/T}$$

$$\mathcal{F}_{\rm H}^{\rm c,b} = 1 + \sum_{Res} g_{H^*}^{c,b}$$

$$R^*_{H_{c,b}} = \left(\frac{\mathcal{F}_{\rm H}^{\rm c}}{\mathcal{F}_{\rm H}^{\rm b}}\right)^{-1} \frac{d\sigma^{H_c}/dp_T}{d\sigma^{H_b}/dp_T}$$

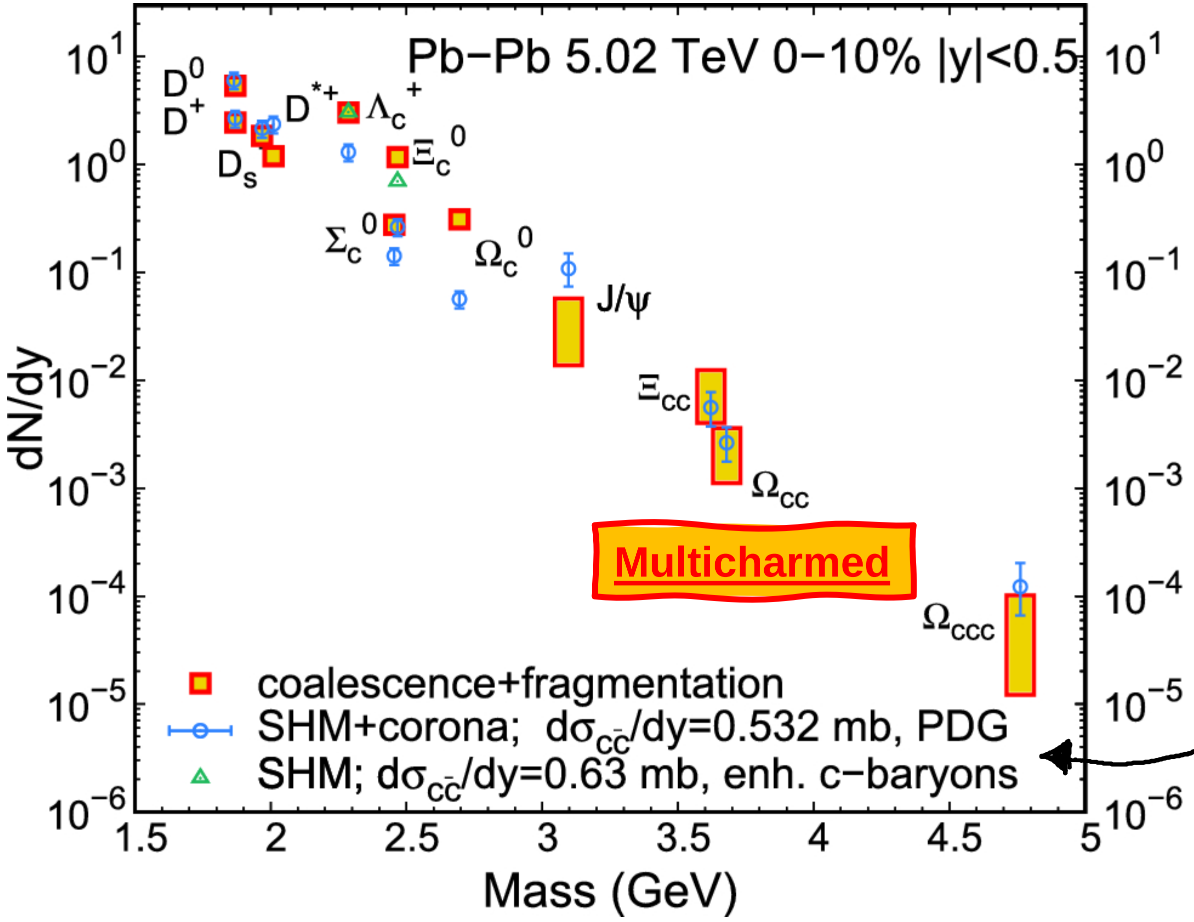


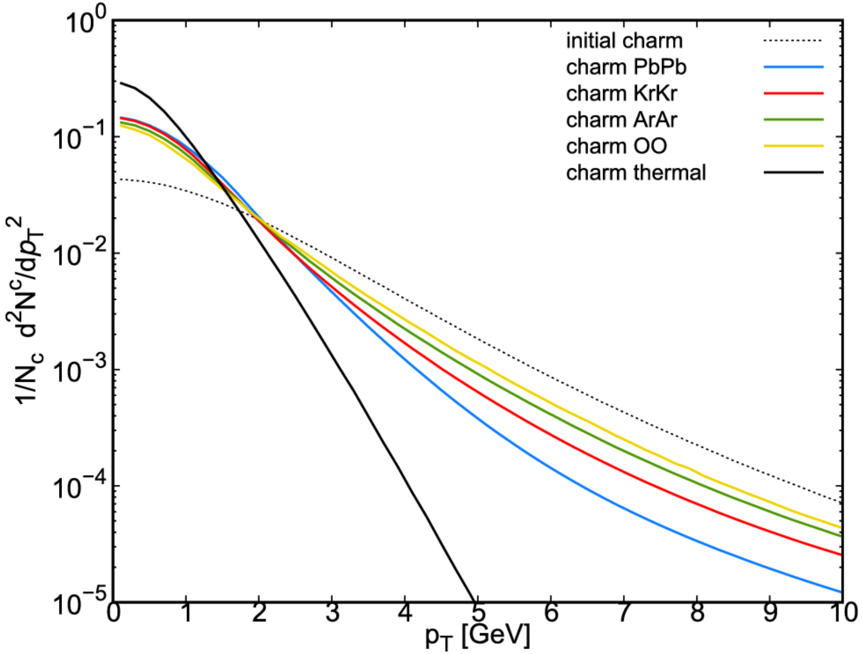
D/B, Λ_c/Λ_b, Ξ_c/Ξ_b, Ω_c/Ω_b provide information about hadronization and f(c)/f(b)

Scaling when only coal. is assumed, considering only the g.s.

Yields in PbPb from coalescence vs SHM

V. Minissale, S. Plumari, Y. Sun and V. Greco, *Eur. Phys. J. C 84, no.3, 228 (2024)*





We employ same volume in SHM
A. Andronic et al., JHEP 07 (2021) 035

	OO	ArAr	KrKr	PbPb
$R_0(fm)$	2.76	3.75	4.9	6.5
$R_{max}(fm)$	5.2	7.65	10.1	14.1
$\tau(fm)$	4	5	6.2	8
β_{max}	0.55	0.6	0.64	0.7
$V_{\lvert y\rvert<0.5}(fm^3)$	345	920	2000	5000

→ upper limit: charm thermal distribution

→ lower limit: PbPb distribution with widths rescaled as standard Harm. Oscill. (ω from Ω_c^0)

$\Sigma_c^0, \Xi_c^0, \Omega_c^0$, widths from quark model

Ξ_{cc}, Ω_{cc} widths obtained rescaling with harm. oscillator

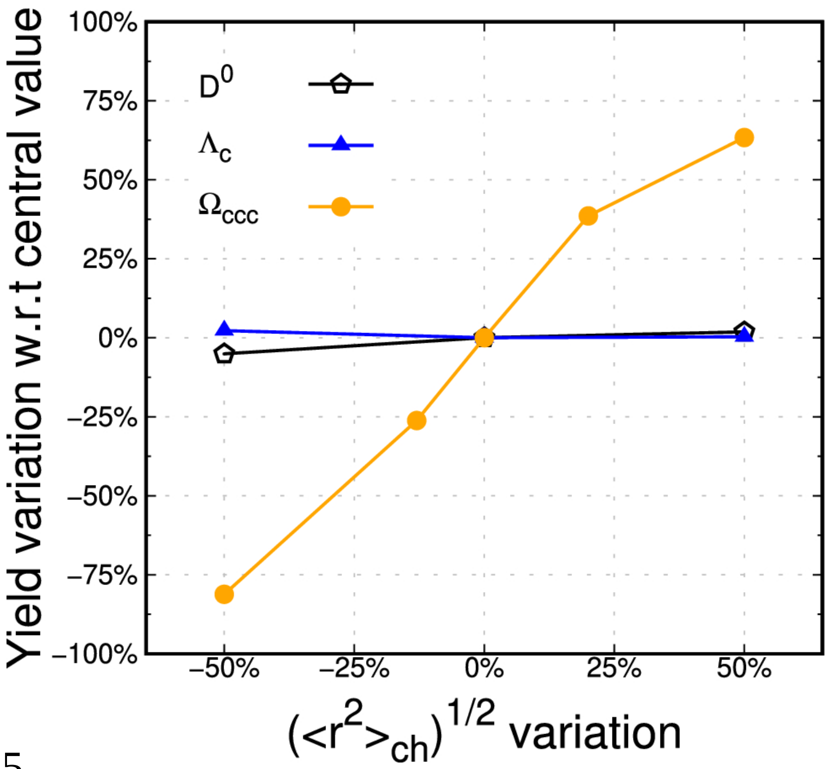
Widths and system size effects on multicharmed hadrons

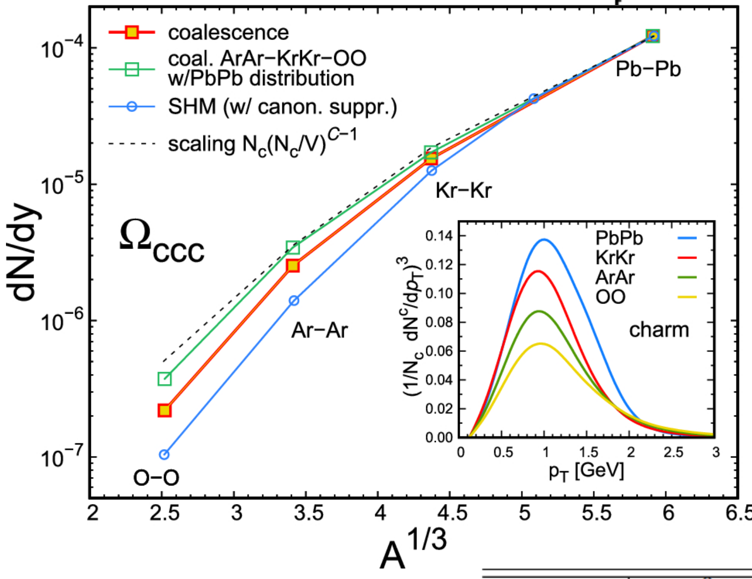
Widths effects on production:

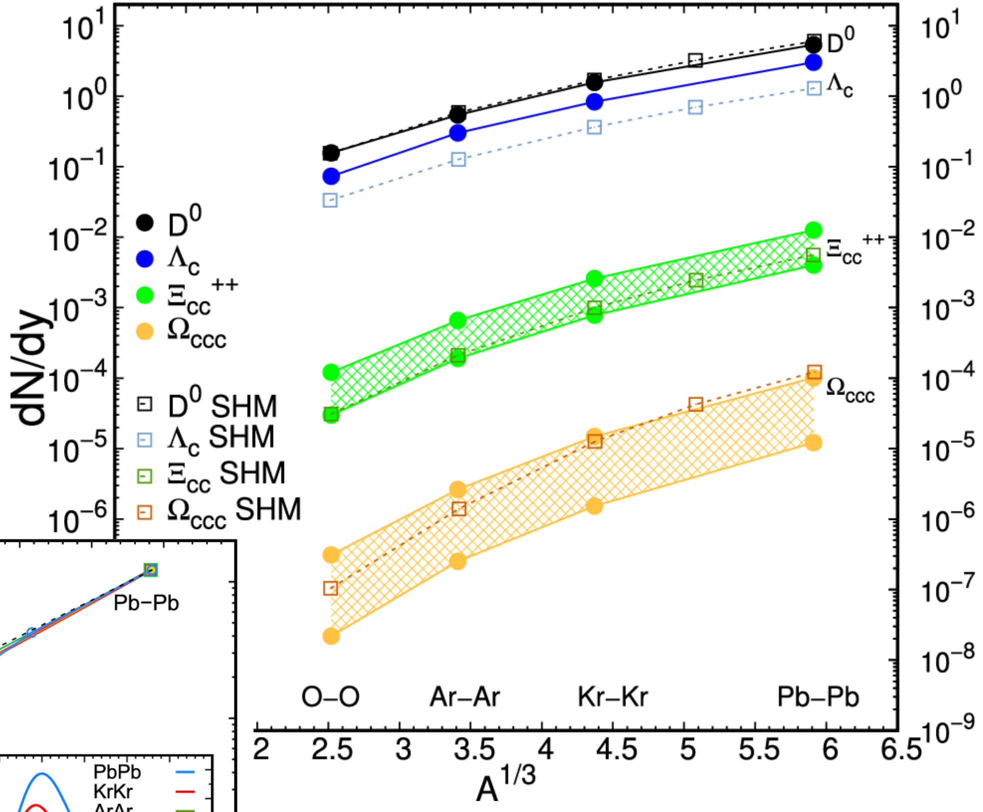
D^0 and Λ_c determine the yield, the radius variation is compensated by the constraint on the charm hadronization

A ± 50% in the radius of Ω_{ccc} induces **a change in the yield of about ± 70%**

V. Minissale, S. Plumari, Y. Sun and V. Greco, *Eur. Phys. J. C 84, no.3, 228 (2024)*







Follow the $N_c \left(\frac{N_c}{V}\right)^{C-1}$ scaling at fixed distribution

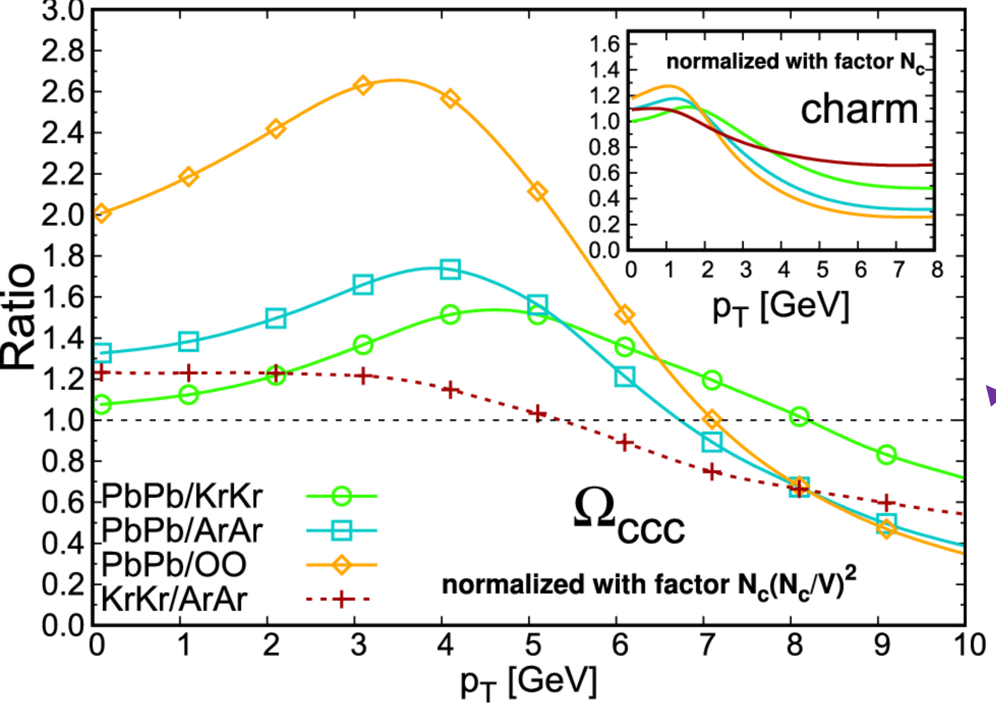
System size scan production

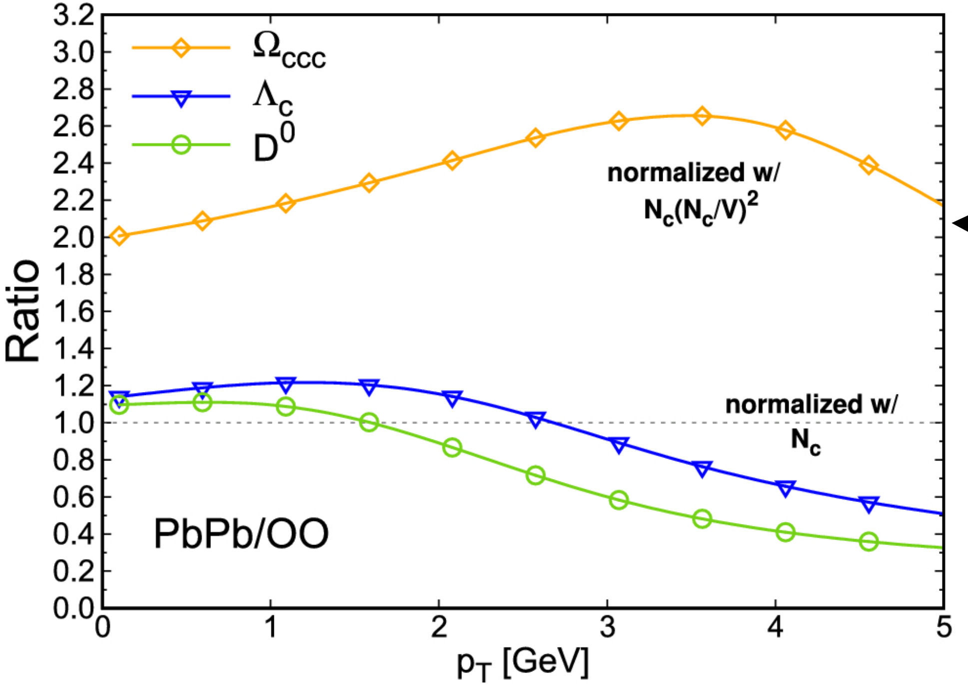
- upper limit: charm thermal distr.
- lower limit: PbPb distr
- Ω_{ccc} production follow c^3 distr.

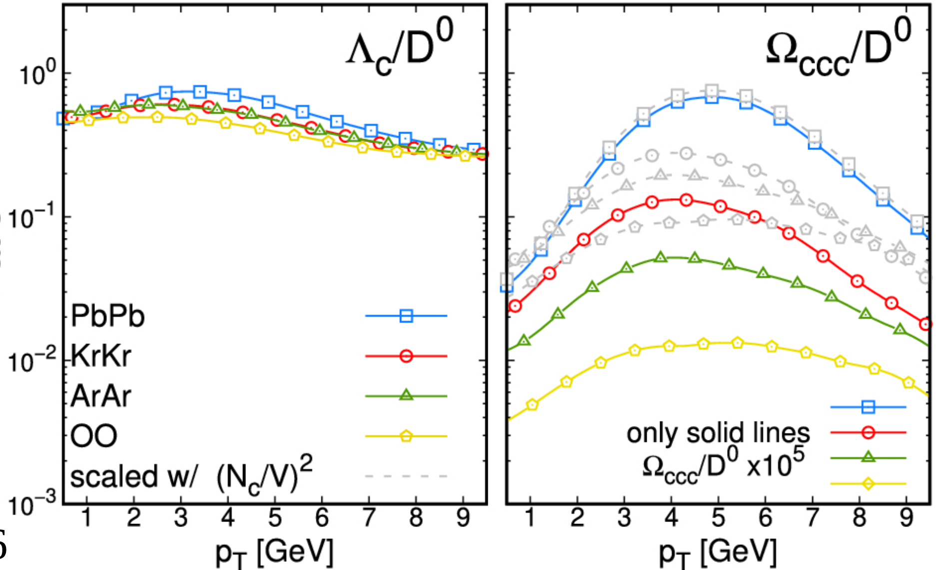
	D^0	Λ_c	$\Xi_{cc}^{+,++}$	Ω_{ccc}
OO	0.156	0.0732	$3-12.1\cdot10^{-5}$	$2.2-29.2\cdot10^{-8}$
$ArAr$	0.543	0.301	$1.9-6.6\cdot10^{-4}$	$2.5-26.3\cdot10^{-7}$
$KrKr$	1.564	0.835	$0.78-2.6\cdot10^{-3}$	$1.5-14.9\cdot10^{-6}$
$PbPb$	5.343	3.0123	$4-12.5\cdot10^{-3}$	$0.12-1.01\cdot10^{-4}$

Ratios of p_T distribution Ω_{ccc} in PbPb/KrKr/ArAr/OO

[V. Minissale, S. Plumari, Y. Sun and V. Greco, *Eur. Phys. J. C 84, no.3, 228 (2024)*]







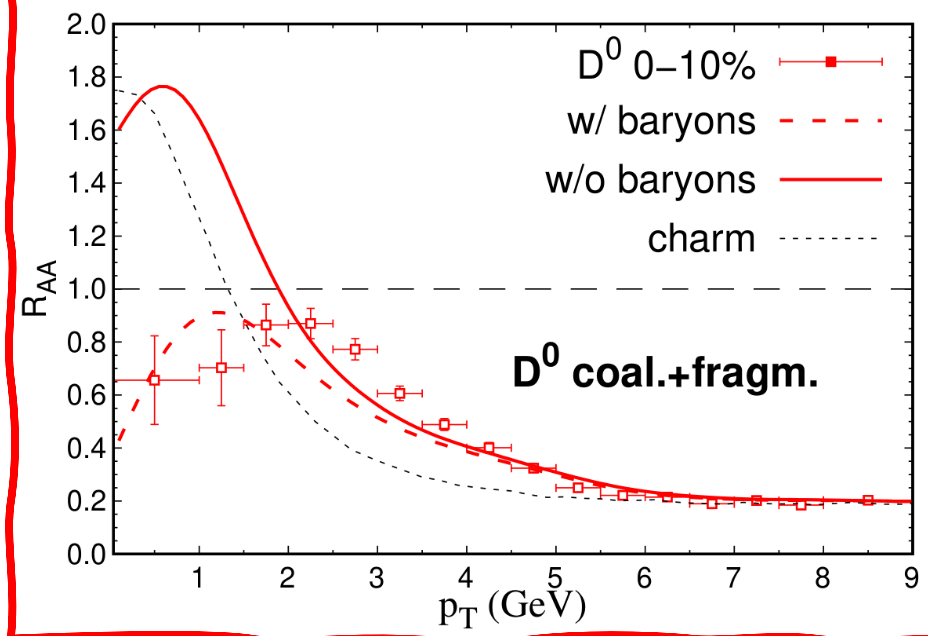
It can be a meter of non-equilibrium.

Translation of features of charm spectra at low p_T to hadron high momentum region.

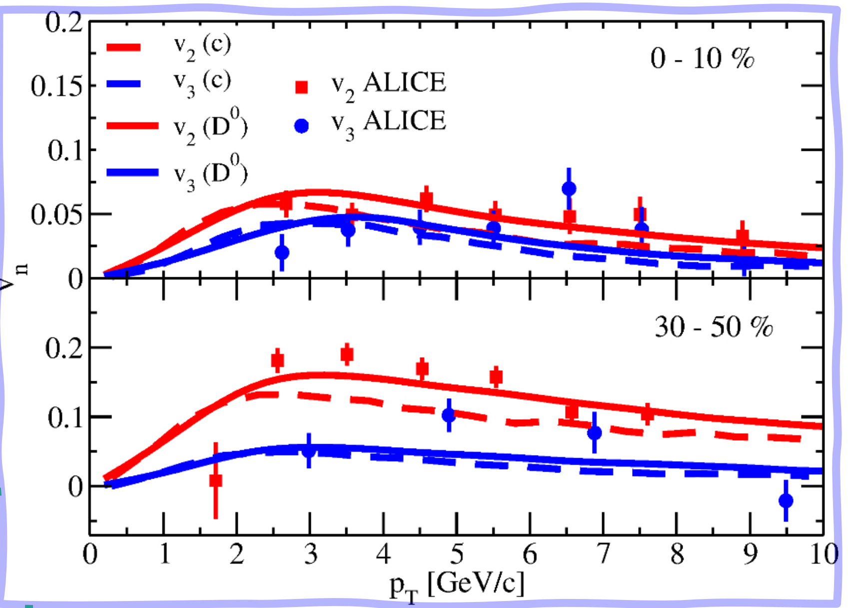
More sensitive multicharm Ω_{ccc}/D^0 respect to Λ_c/D^0

More sensitive to system change w.r.t. D^0 Λ_c (fragm. Effect)

Implications and developments:

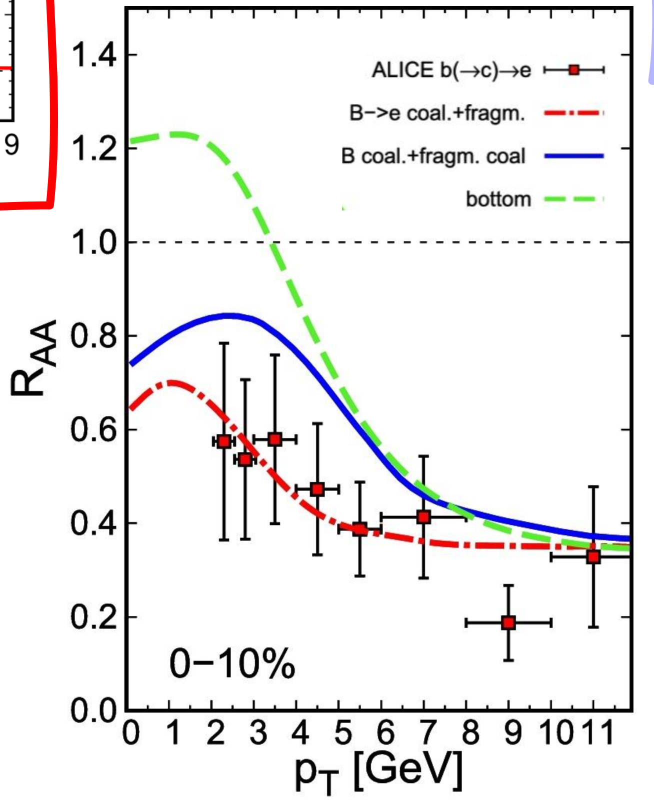


The large Λ_c, Ξ_c production has effects on the R_{AA} of D^0, because of the charm conservation



Coalescence give an enhancement to the $v_n(p_T)$ of final hadrons compared to the charm $v_n(p_T)$.

Sambataro,Sun,Minissale,Plumari,Greco, Eur.Phys.J.C 82 (2022) 9, 833



Electrons from semileptonic B meson decay with a coal + fragm model for B meson production

Sambataro,Minissale,Plumari,Greco *Phys.Lett.B* 849 (2024) 138480

Conclusions

- **Charm hadronization in AA:**

Coalescence+fragmentation gives enhancement of Λ_c production at intermediate momentum region: $\Lambda_c/D^0 \sim 1$ for $p_T \sim 3\ GeV$

- ***In p+p assuming a medium:***

Charm:

Coal.+fragm. good description of heavy baryon/meson ratio (closer to Λ_c/D^0 , Ξ_c/D^0, Ω_c/D^0 data)

Bottom:

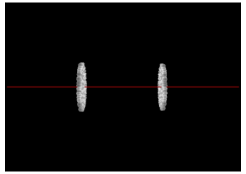
- ✓ B, B_s good agreement with exp. data.
- ✓ Coal.+fragm. Enhancement of Λ_b
- ✓ Predictions for Λ_b/B^0 , Ξ_b/B^0, Ω_b/B^0
- ✓ D/B, Λ_c/Λ_b , Ξ_c/Ξ_b provide information about hadronization and f(c)/f(b)

- ***Multicharm hadrons:*** - role of non-equilibrium distribution function

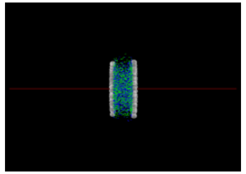
  -in accord with SHM predictions

Backup Slides

Quark Gluon Plasma in Ultra-Relativistic Heavy-Ion Collisions

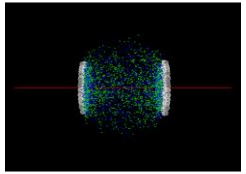


Initial Stage



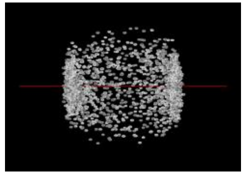
Pre-equilibrium stage

Expansion

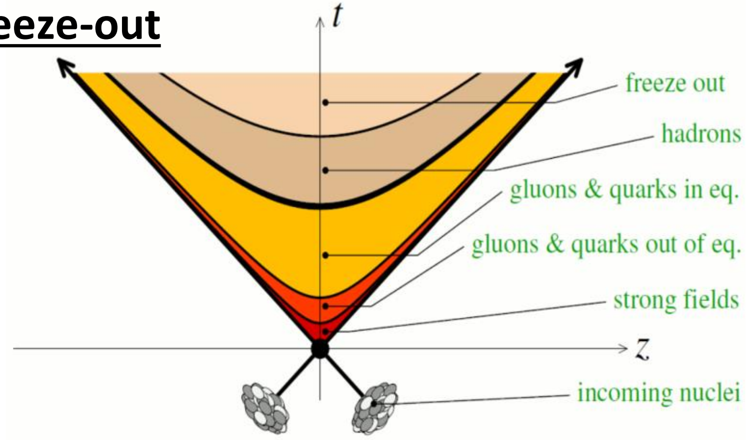


QGP

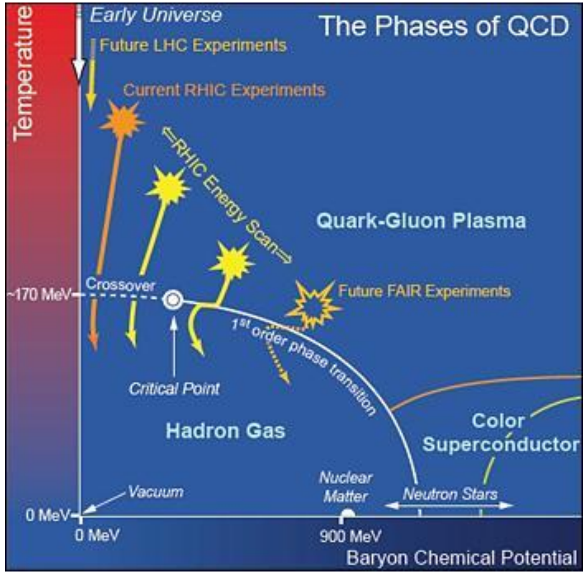
Hadronization

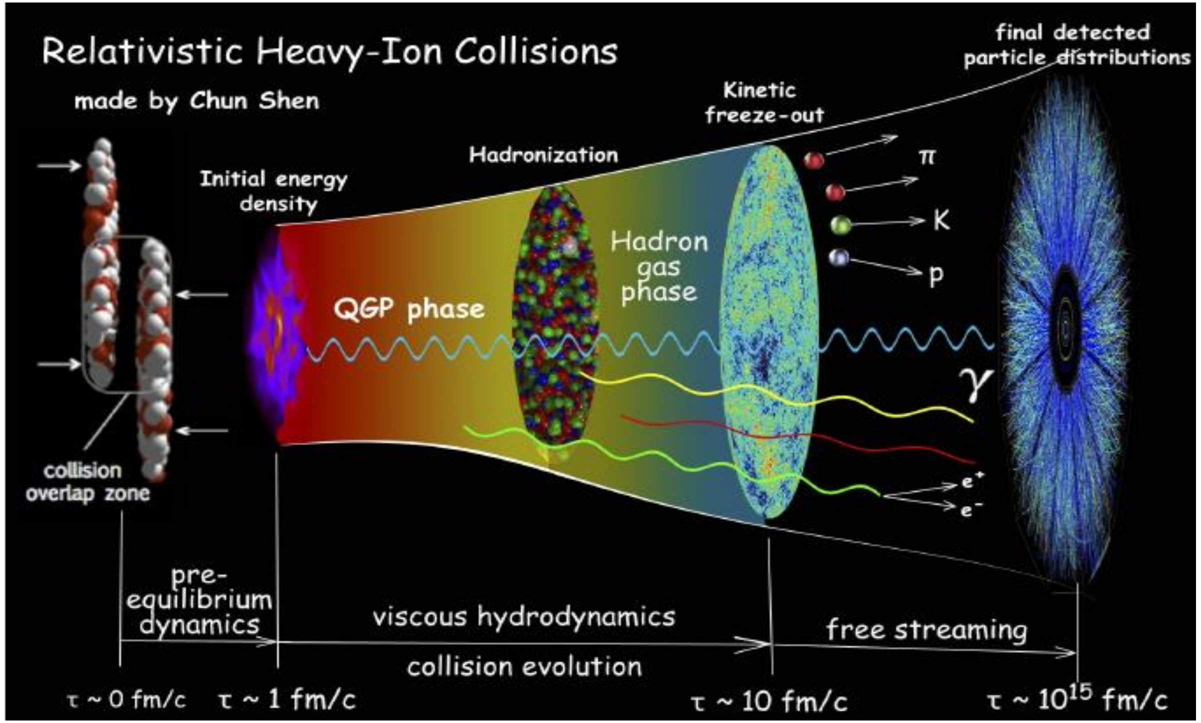


Chemical and kinetic freeze-out



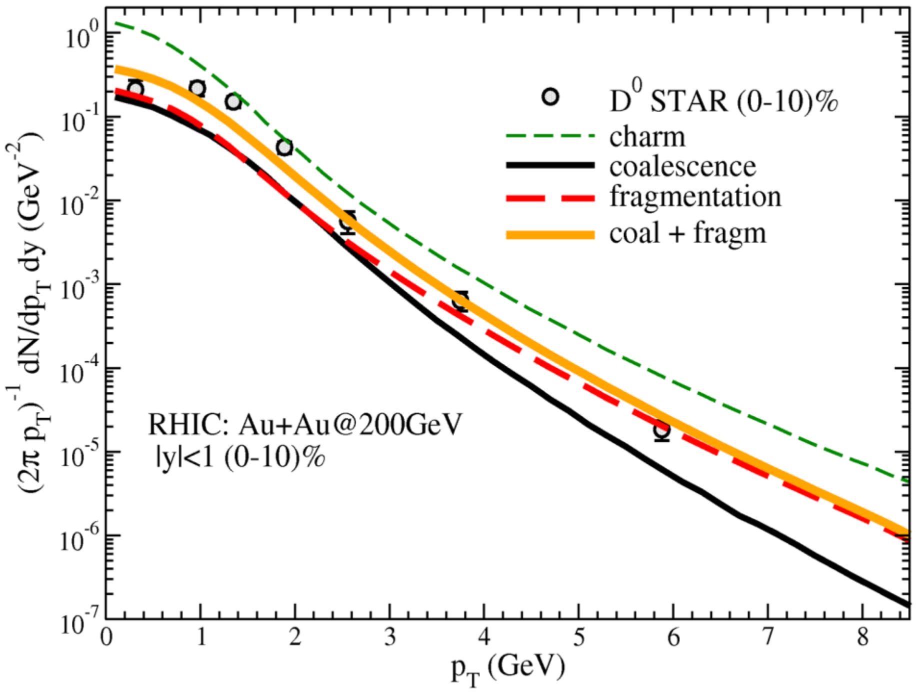
- Nuclear matter: Critical Energy and Temperature in the transition between confined and deconfined phase
- If $T>T_C$ colour charges are deconfined in a Quark Gluon Plasma (QGP)
- Different value of T and ρ for deconfinement → Phase Diagram





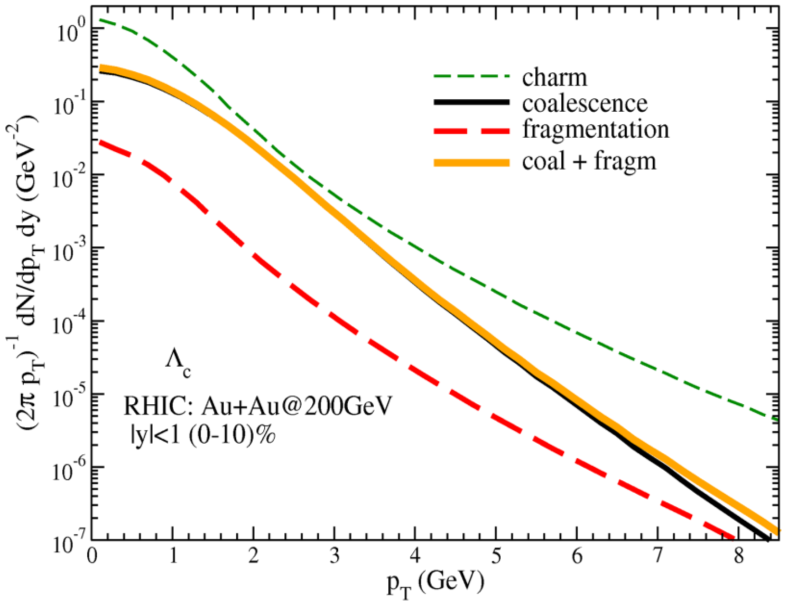
RHIC: results

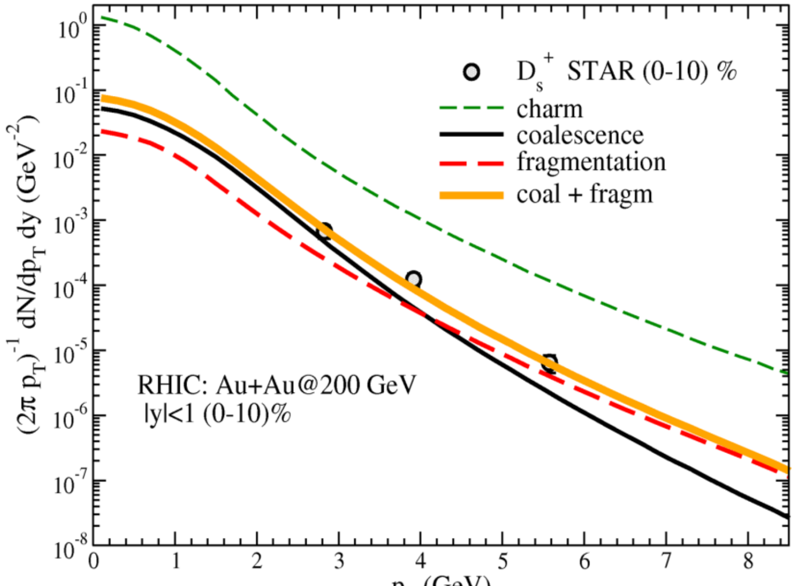
S. Plumari, V. Minissale et al., Eur. Phys. J. ***C78*** *no. 4, (2018) 348*



Data from STAR Coll. PRL **113** (2014) no.14, 142301

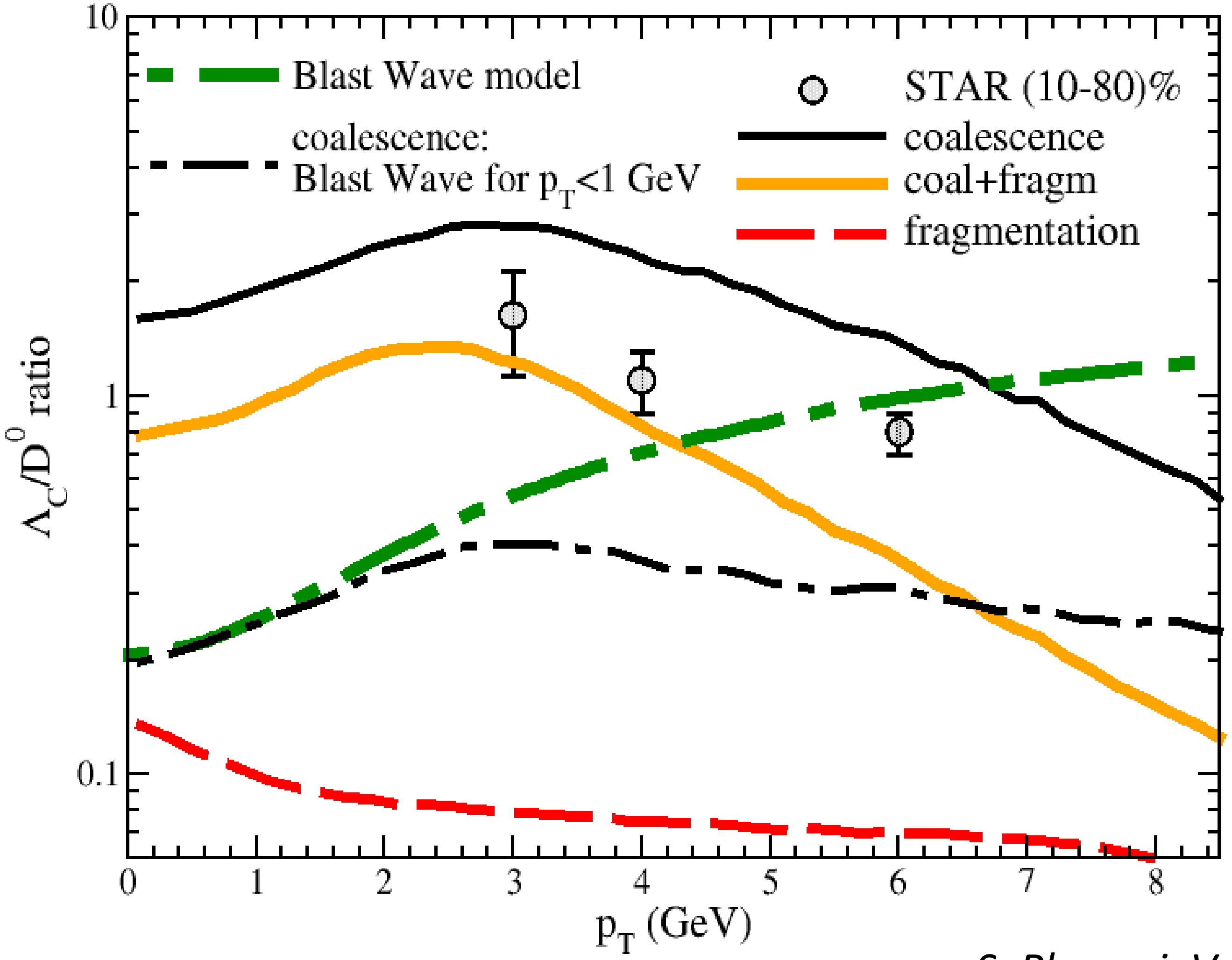
- For D^0 coalescence and fragmentation comparable at 2 GeV
- fragmentation fraction for D^+_s are small and less than about 8% of produced total heavy hadrons
- Λ_c^+ fragmentation is even more smaller, coalescence gives the dominant contribution





RHIC: Baryon/meson

STAR, Phys.Rev.Lett. 124 (2020) 17, 172301



Compared to light baryon/meson ratio the Λ_c/D^0 ratio has a larger width (flatter)

More flatter → should coalescence extend to higher pt? Indication also in light sector

V. Minissale, F. Scardina, V. Greco **PRC 92**,054904 (2015)

Cho, Sun, Ko et al.,**PRC 101 (2020)** 2, 024909

Needed data at low p_T

S. Plumari, V. Minissale et al., Eur. Phys. J. ***C78*** *no. 4, (2018) 348*

Multicharm production Pb-Pb, Kr-Kr, Ar-Ar, O-O

V. Minissale, S. Plumari, Y. Sun and V. Greco, *Eur. Phys. J. C 84, no.3, 228 (2024)*

Baryon			
$\Xi_{cc}^{+,++} = dcc, ucc$	3621	$\frac{1}{2}\,(\frac{1}{2})$	
$\Omega_{scc}^{+} = scc$	3679	$0\,(\frac{1}{2})$	
$\Omega_{ccc}^{++} = ccc$	4761	$0\,(\frac{3}{2})$	
Resonances			
Ξ_{cc}^{*}	3648	$\frac{1}{2}\,(\frac{3}{2})$	$1.71 \times g.s$
Ω_{scc}^{*}	3765	$0\,(\frac{3}{2})$	$1.23 \times g.s$

like S.Cho and S.H. Lee, PRC101 (2020)
from R.A. Briceno et al., PRD 86(2012)

Strengths of the approach:

- Does not rely on distribution in equilibrium for charm
  → useful for small AA down to pp collisions and at p_T> 3-4 GeV

- Provide a p_T dependence of spectra and their ratios vs p_T

Widths from harmonic oscillator rescaling and from<r> of Tsingua approach

	$\sigma_{p_1}(GeV)$	$\sigma_{p_2}(GeV)$	$\sigma_{r_1}(fm)$	$\sigma_{r_2}(fm)$
Ξ_c	0.262	0.438	0.751	0.450
Ω_c	0.345	0.557	0.572	0.354
Ξ_{cc}^{ω}	0.317	0.573	0.622	0.344
$\Omega_{ccc}^{\sigma_r\sigma_p=3/2}$	0.522	0.522	0.566	0.566

Elliptic Flow – Quark Number Scaling

Fourier expansion of the azimuthal distribution

$$f(\varphi, p_T) = 1 + 2\sum_{n=1}^{\infty} v_n(p_T)\cos n\varphi$$

momentum anisotropy in the transverse plane

n=2 Elliptic flow

coalescence brings to

$$v_{2,M}(p_T) \approx 2v_{2,q}(p_T/2)$$
$$v_{2,B}(p_T) \approx 3v_{2,q}(p_T/3)$$

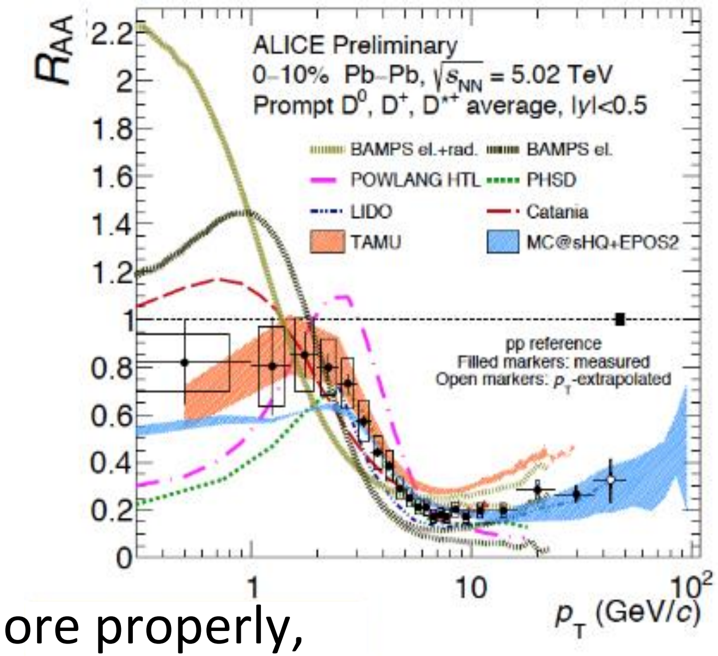
Partonic elliptic flow

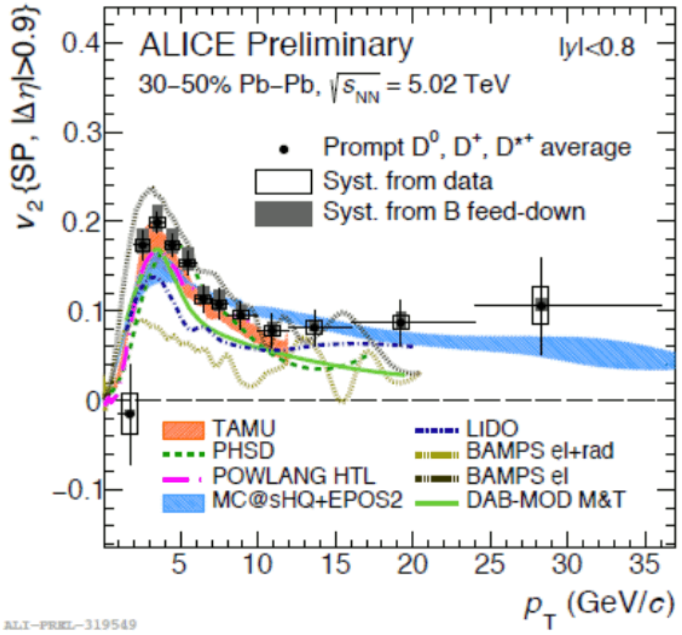
Hadronic elliptic flow

Assumption

- one dimensional
- Dirac delta for Wigner function
- isotropic radial flow
- not including resonance effect

Transport approaches





Models not really tested at p→0
The new data → determine D_s(T) more properly,
i.e. p→0 where it is defined and computed in lQCD

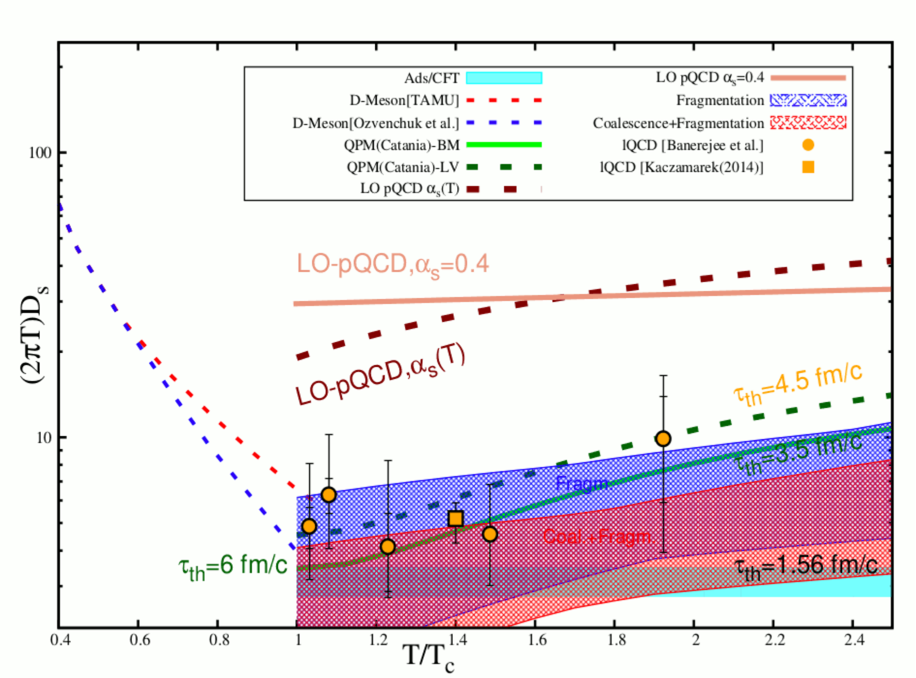
	Catania	Duke	Frankfurt(PHSD)	LBL	Nantes	TAMU
Initial HQ (p)	FONLL	FONLL	pQCD	pQCD	FONLL	
Initial HQ (x)	binary coll.	binaryy coll.	binary coll.	binary coll.		binary coll.
Initial QGP	Glauber	Trento	Lund		EPOS	
QGP	Boltzm.	Vishnu	Boltzm.	Vishnu	EPOS	2d ideal hydro
partons	mass	m=0	m(T)	m=0	m=0	m=0
formation time QGP	0.3 fm/c	0.6 fm/c	0.6 fm/c (early coll.)	0.6 fm/c	0.3 fm/c	0.4 fm/c
interactions in between	HQ-glasma	no	HQ-preformed plasma	no		no

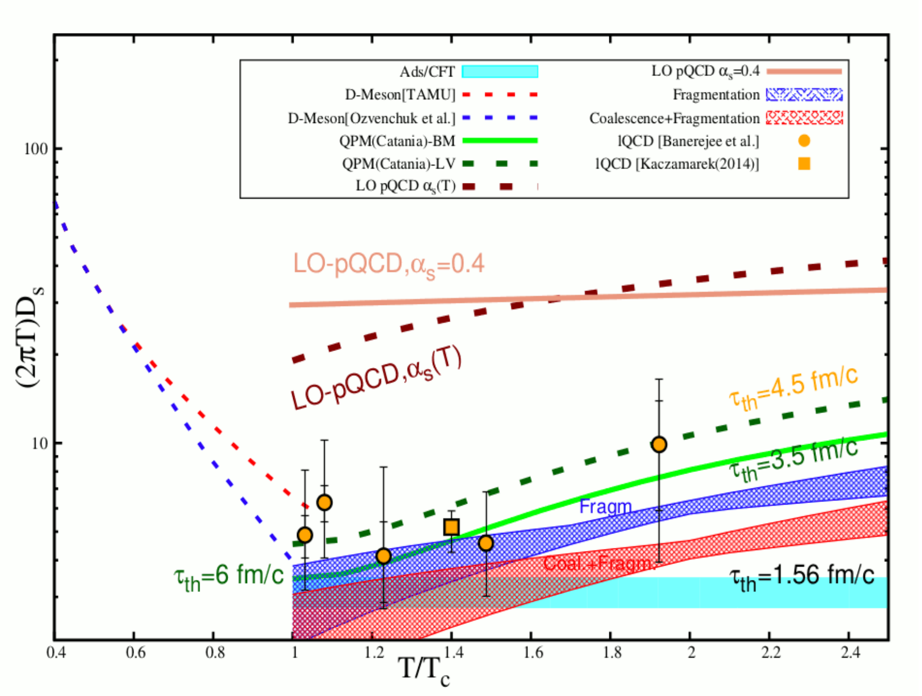
2018-2019

Several Collab. in joint activities:

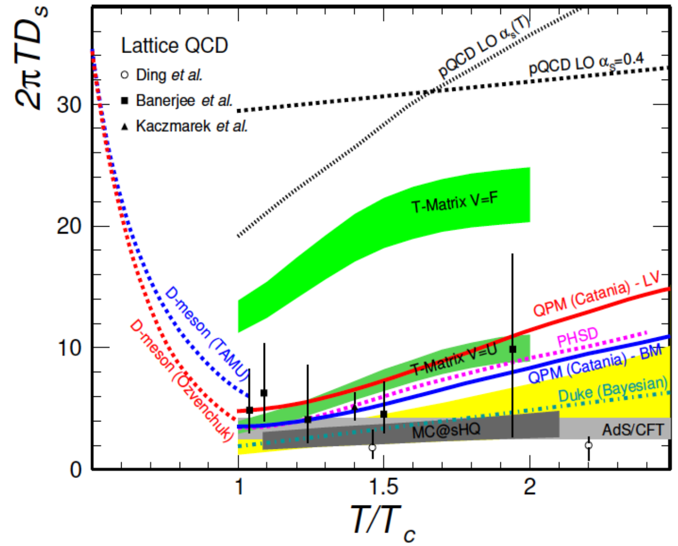
- EMMI-RRTF:
  R. Rapp et al., Nucl. Phys. A 979 (2018)
- HQ-JETS:
  S. Cao et al.,Phys. Rev. C 99 (2019)
- Y. Xu et al., Phys. Rev. C 99 (2019)

Transport coefficient





X. Dong and V. Greco, PPNP(2019)



Different hadronization models can affect the extraction of the charm quark diffusion coefficient

2018-2019

Several Collab. in joint activities:

- EMMI-RRTF: R. Rapp et al., Nucl. Phys. A 979 (2018)
- HQ-JETS: S. Cao et al.,Phys. Rev. C 99 (2019)
- Y. Xu et al., Phys. Rev. C 99 (2019)