

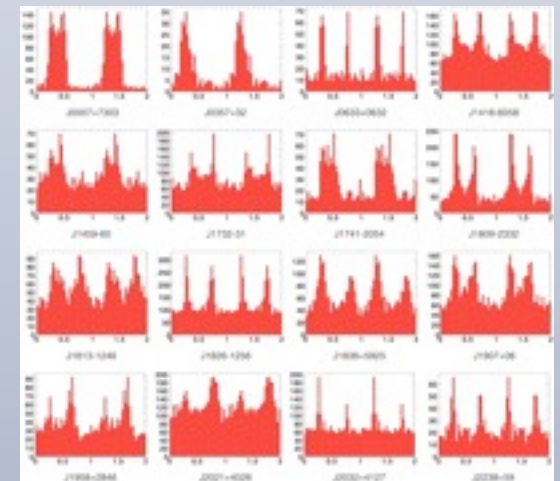
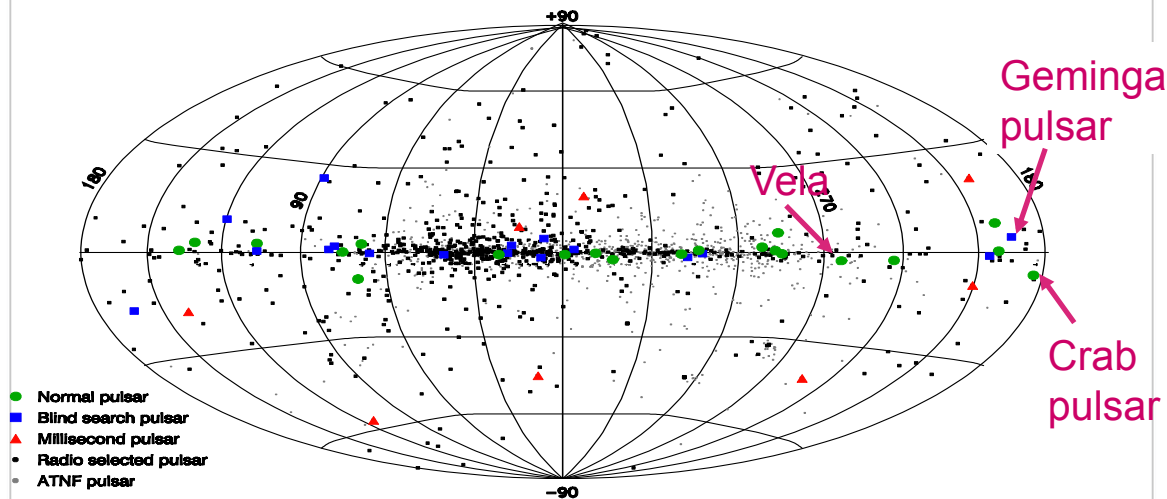
Gamma-ray emission from Crab pulsar and the nebula: paradigm shifts?

Maxim Lyutikov (Purdue U., Osservatorio Arcetri)

1. Gamma-ray emission of pulsars

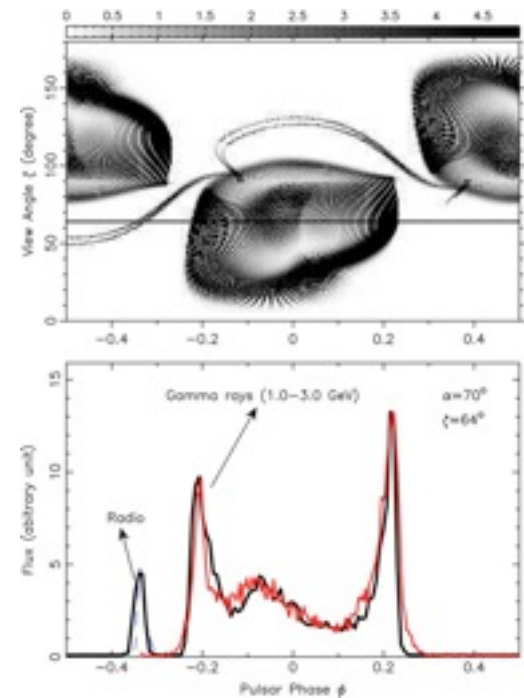
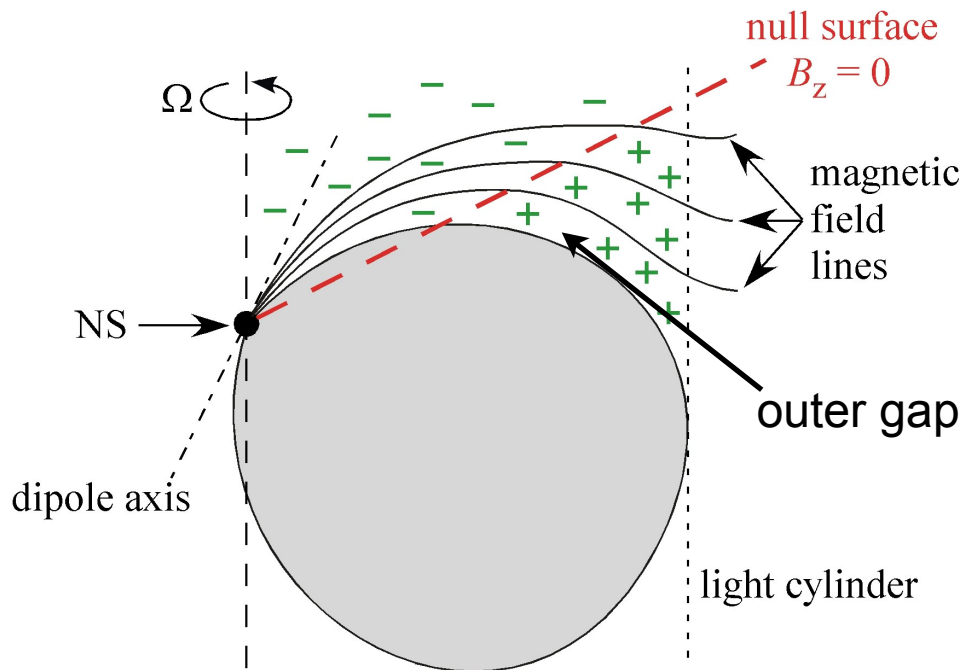
Lively time in high energy astrophysics and pulsar astrophysics

- *Fermi*: *LAT* has detected > 100 pulsars above 100 MeV



Outer gap models: geometrical - good fits

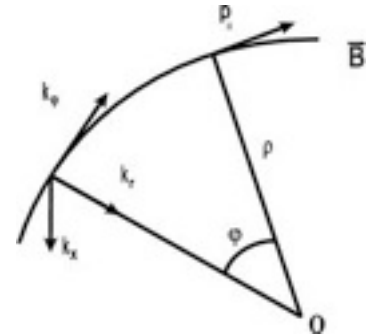
- Rotation induces charge density in the magnetosphere
- null line: outer gap
- E_{\parallel} accelerates particles which emit



Emit what?

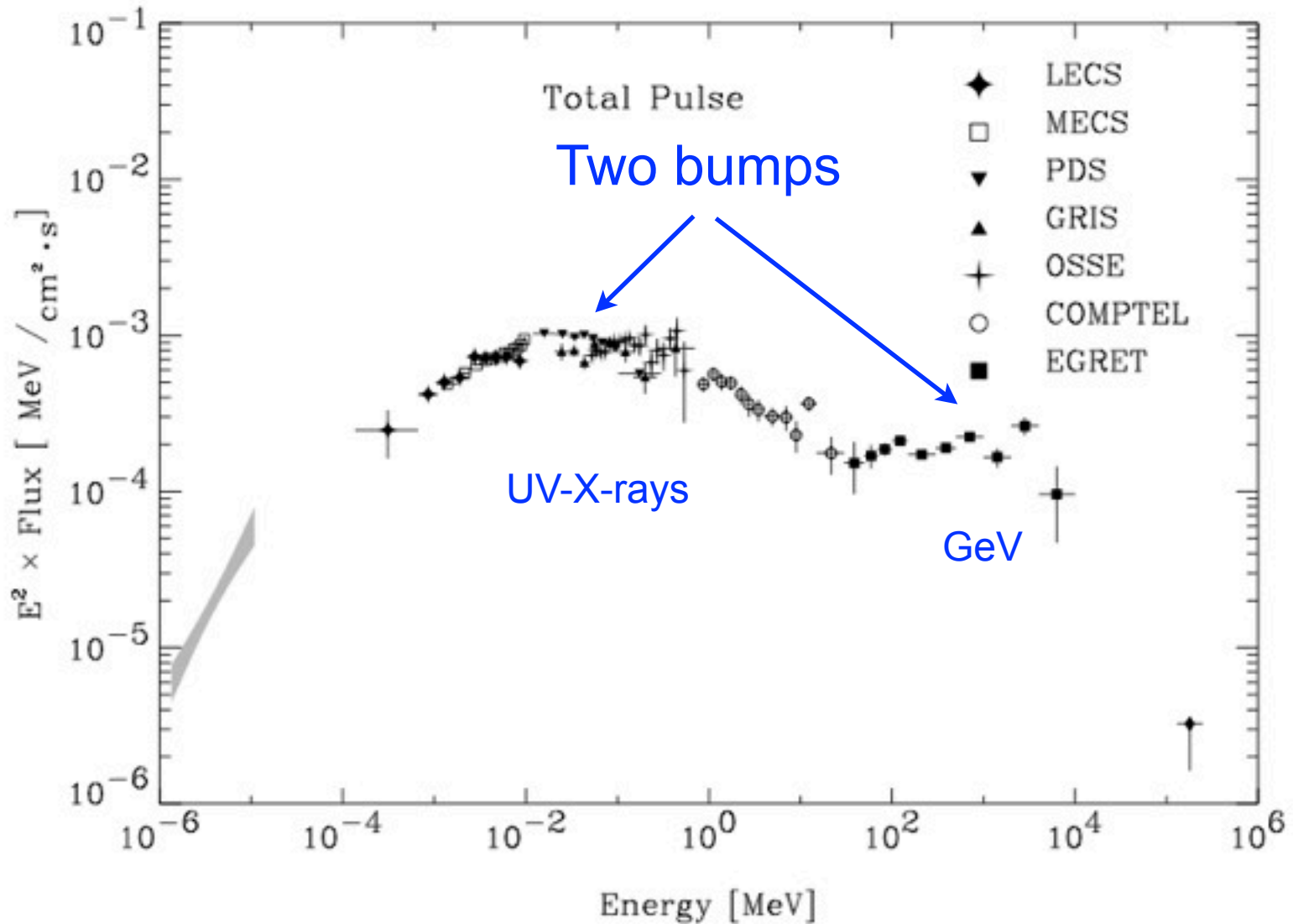
- Curvature emission (Chen & Ruderman 1986)

$$\epsilon_{ph} = \frac{\hbar \gamma^3 c}{R_C}$$

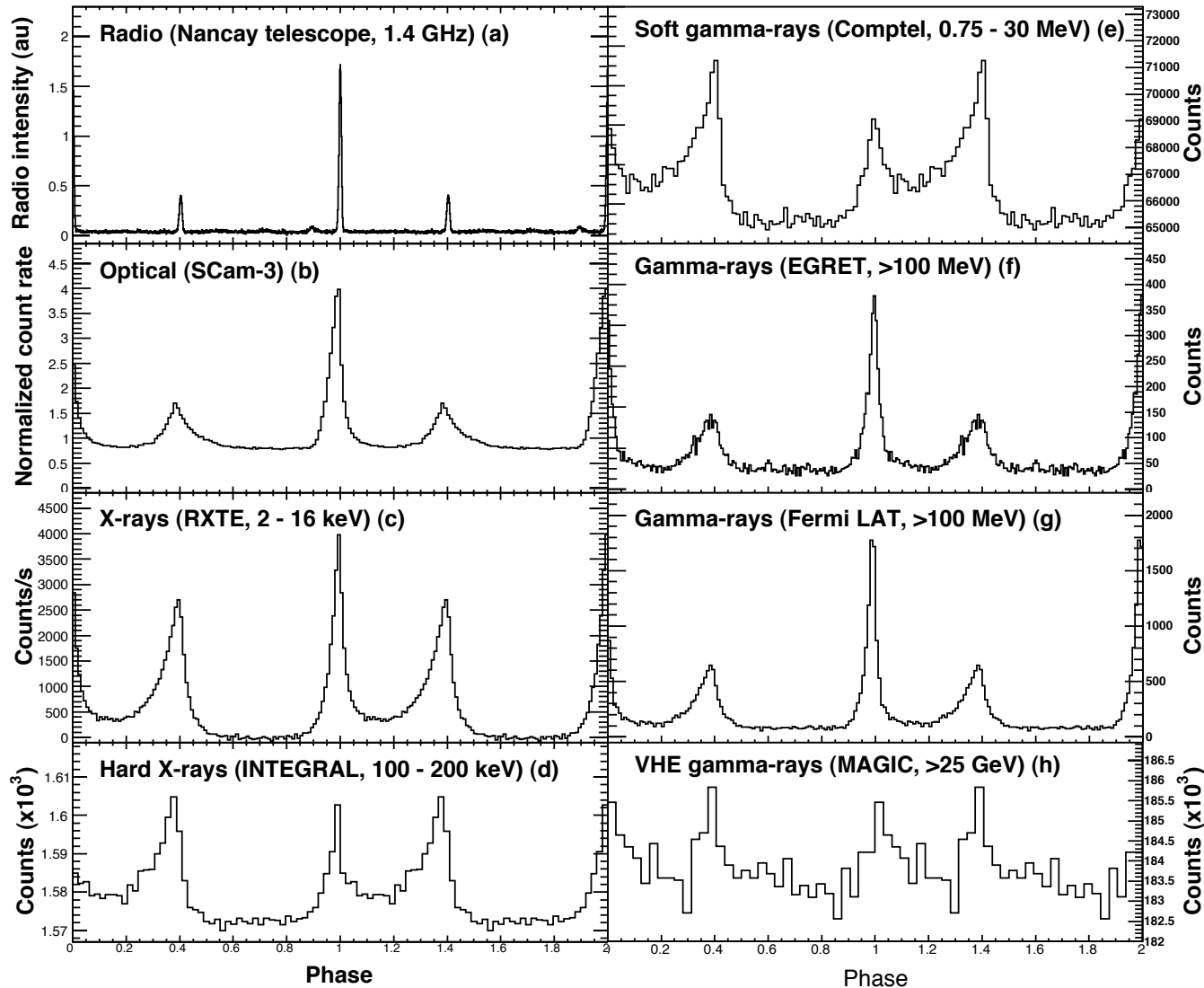


- Hard to solve the full electrodynamic picture:
- E_{\parallel} accelerates particles, produce pairs and currents, pairs screen E_{\parallel} , currents distort B-field, changing E_{\parallel} , non-local radiative transfer.
- Typically $E_{\parallel} \sim 10^{-2}-10^{-1}$ B (Hirota)
- **Above the break the spectrum must be exponentially suppressed**

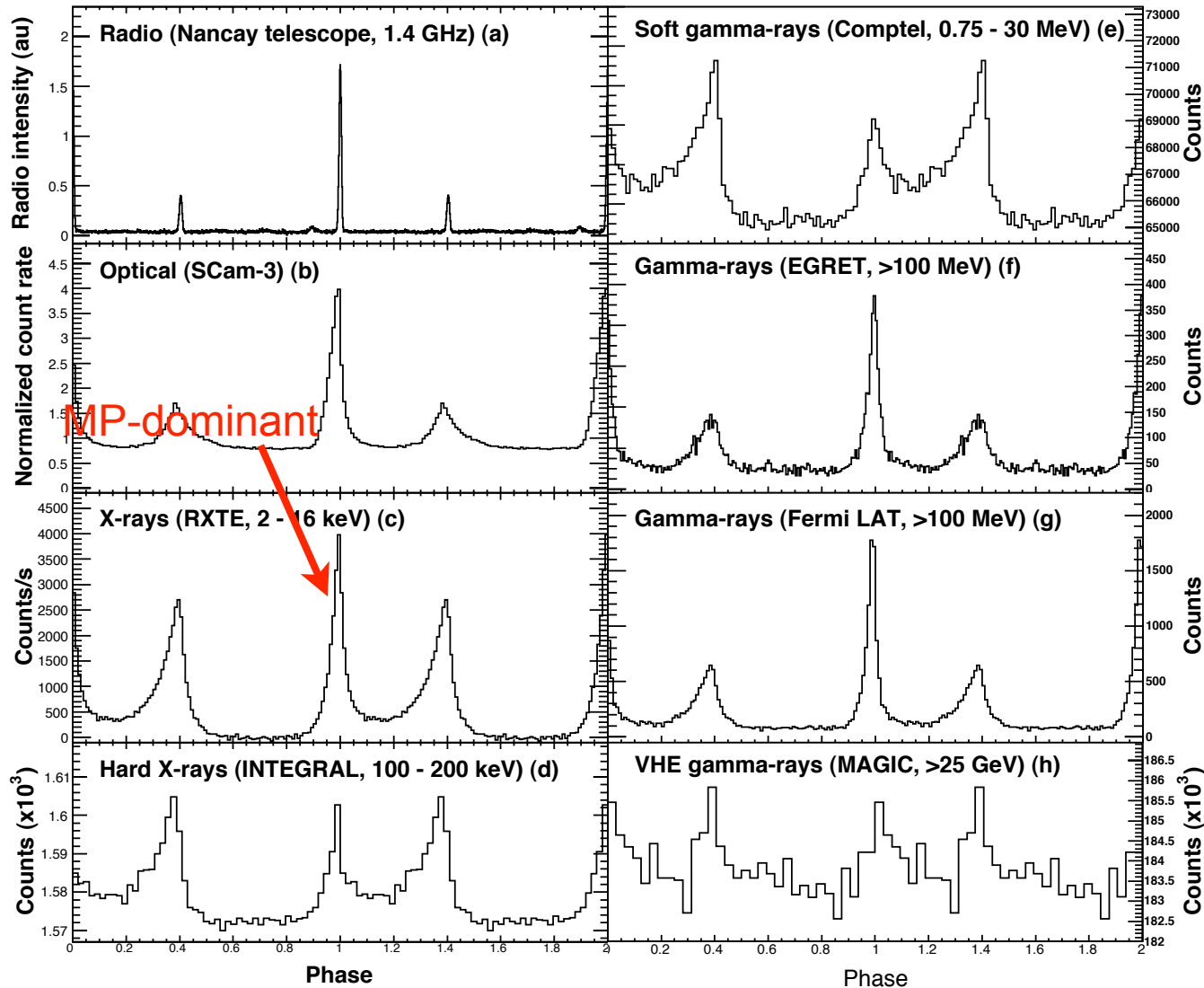
Crab pulsar spectrum: two bumps



Crab profiles

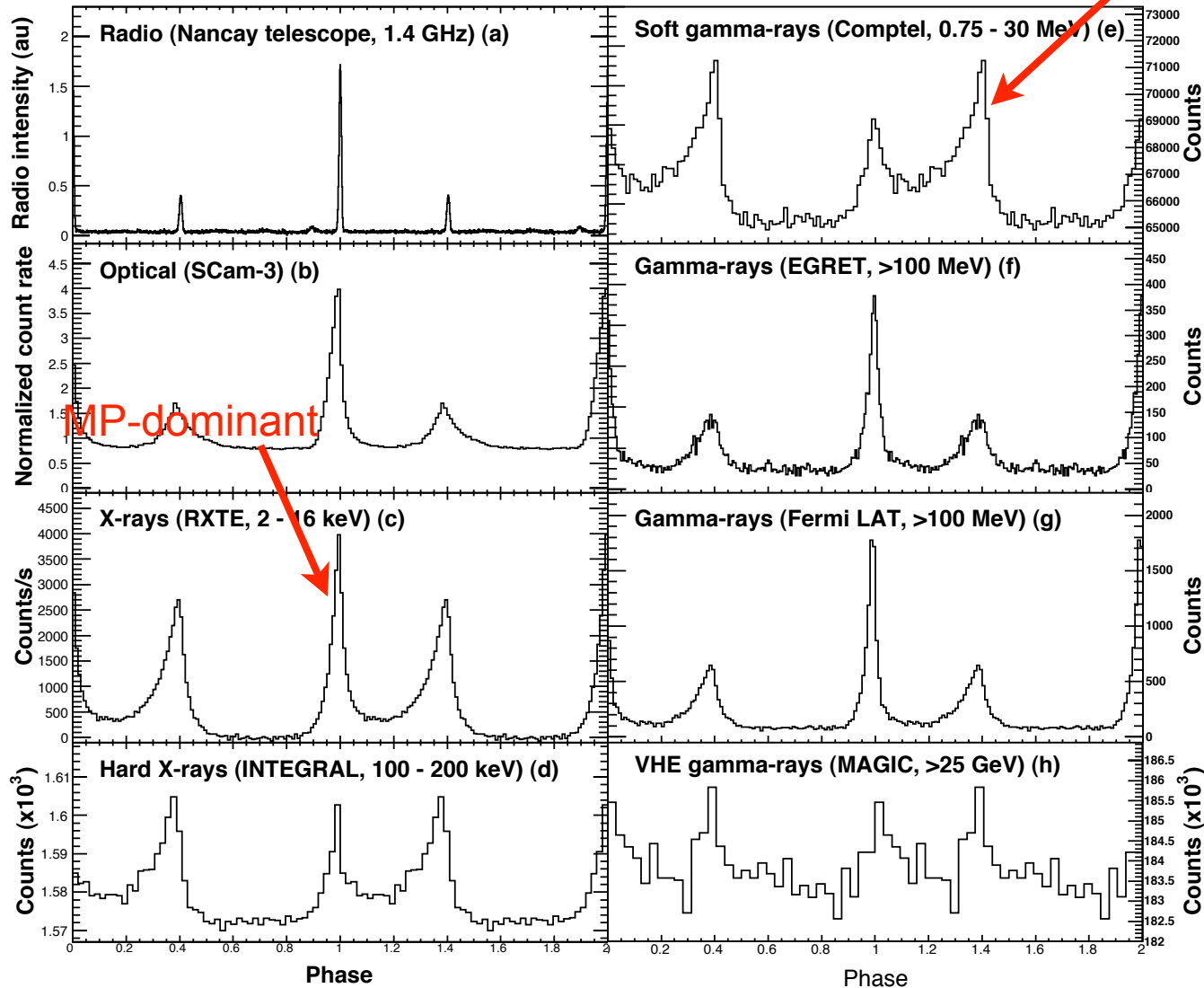


Crab profiles



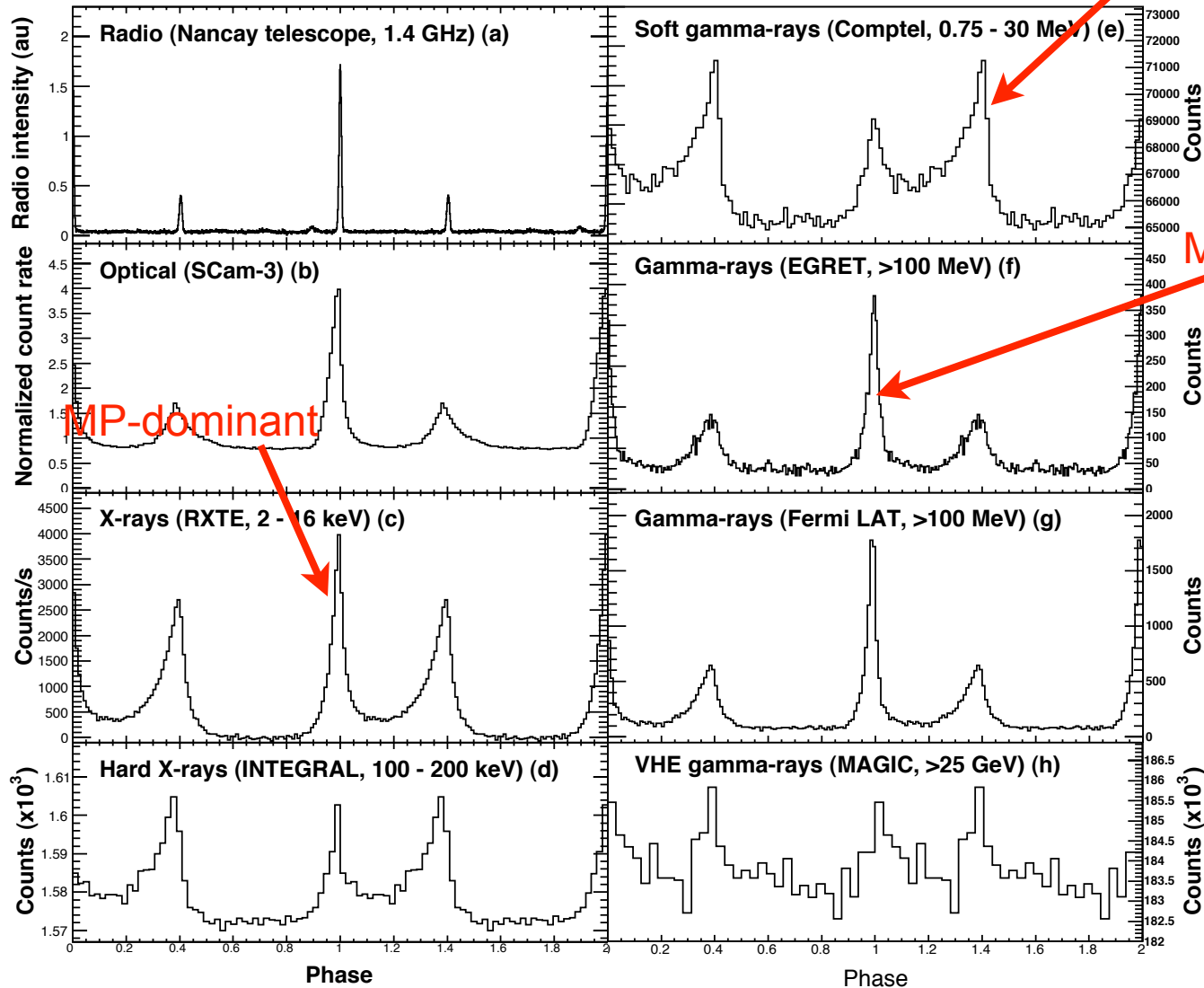
Crab profiles

IP-dominant



Crab profiles

IP-dominant

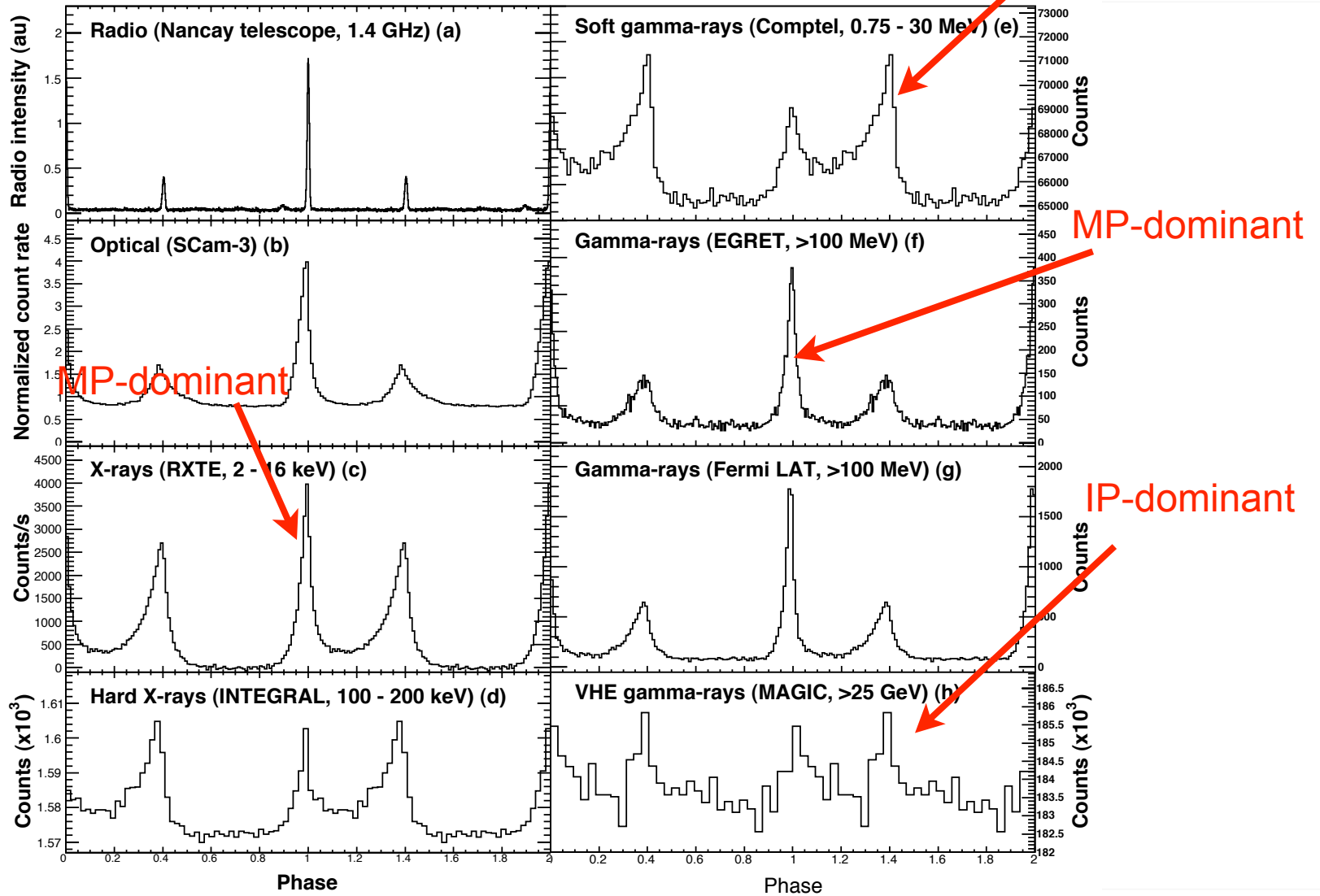


MP-dominant

MP-dominant

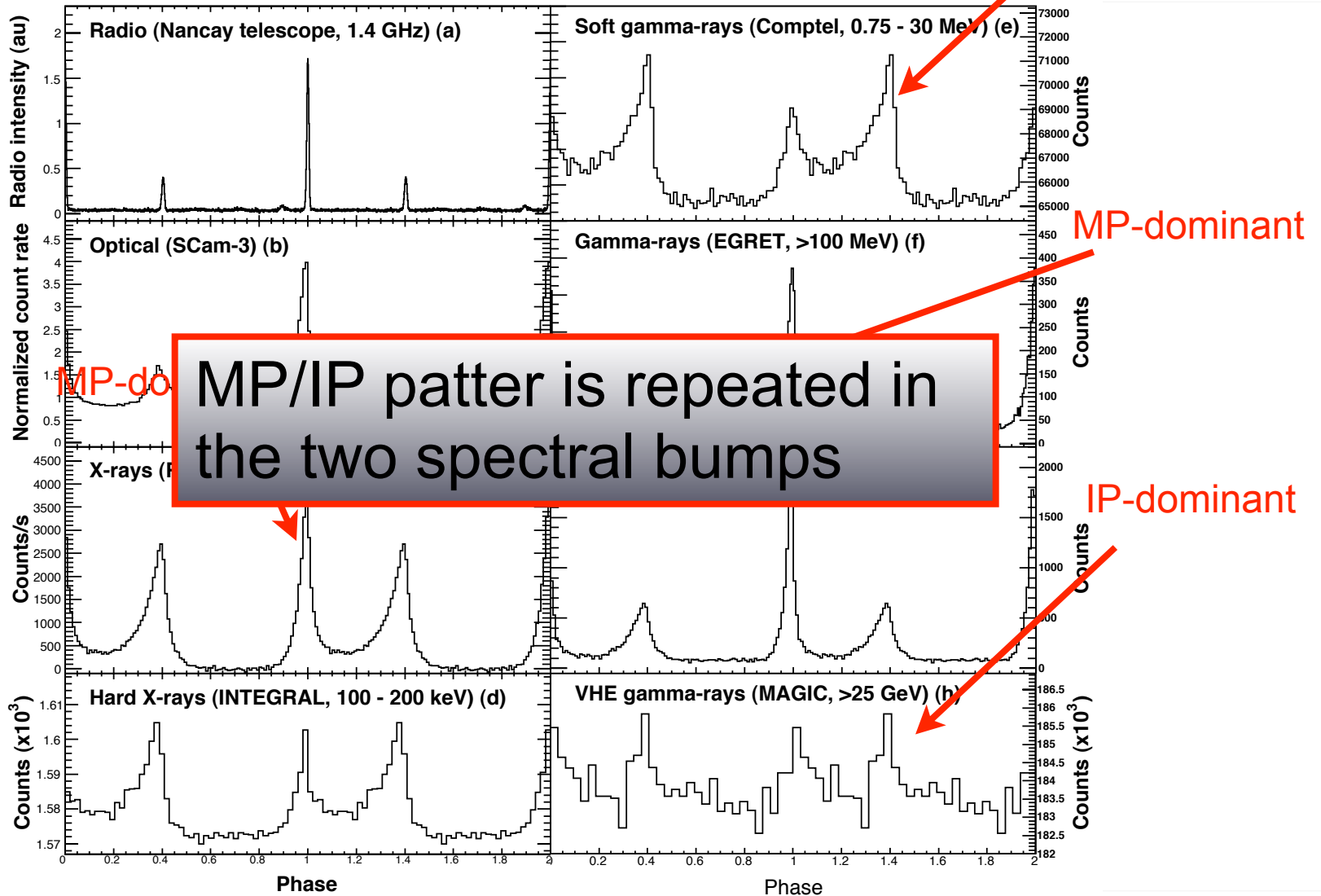
Crab profiles

IP-dominant



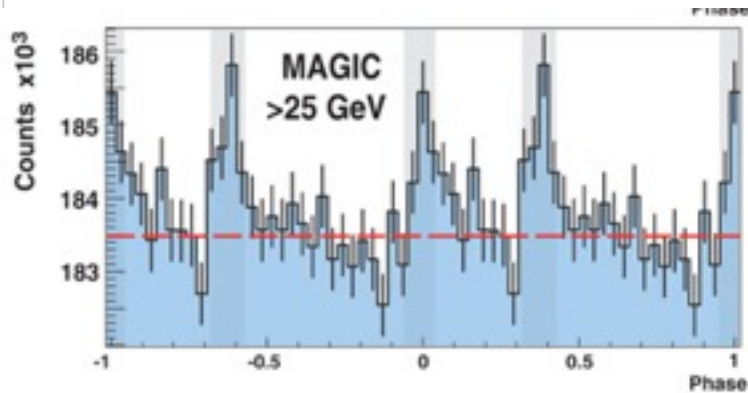
Crab profiles

IP-dominant



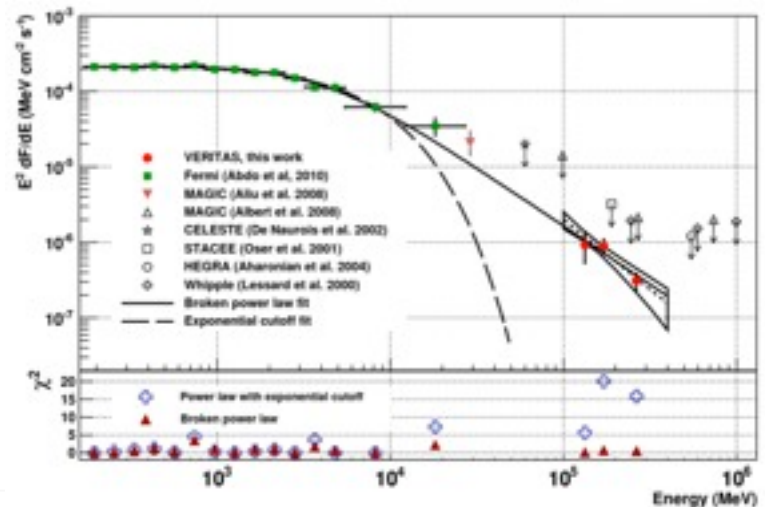
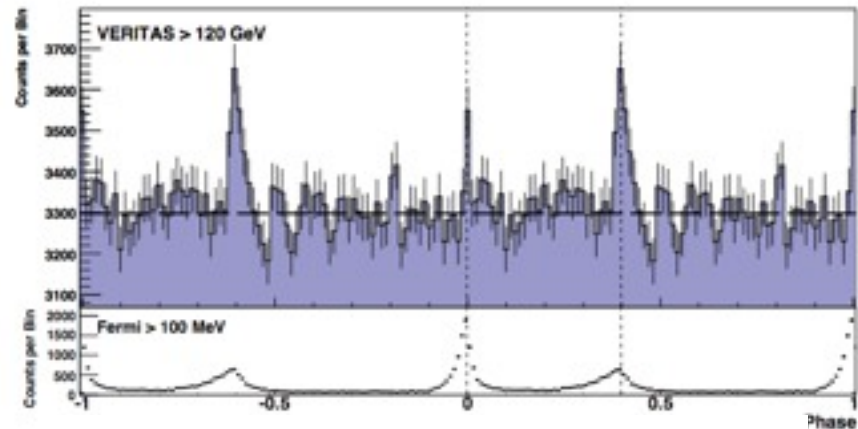
New results

- MAGIC sees Crab at ~ 25 GeV



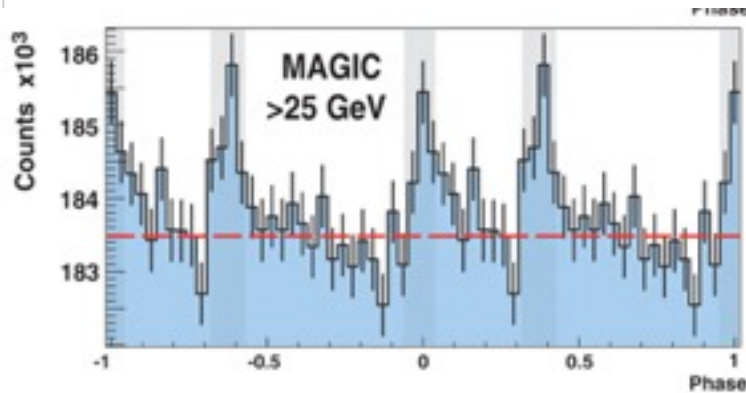
- With indication of non-exponential cut-off (2011)

- VERITAS sees Crab at > 150 GeV!



New results

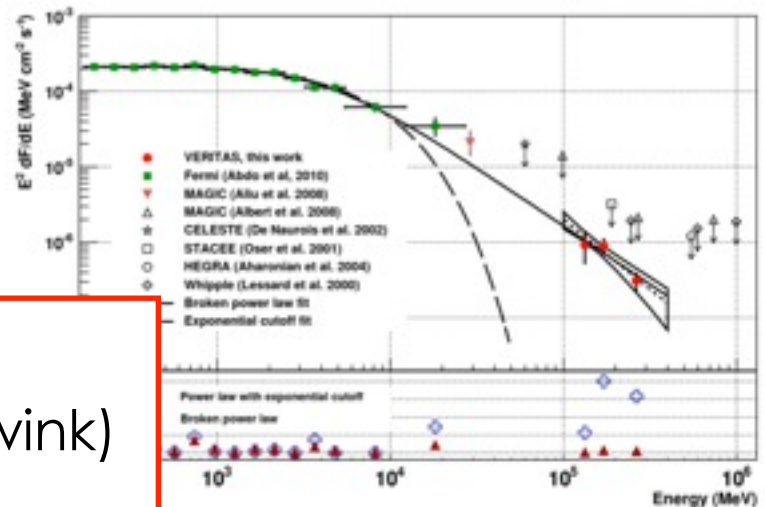
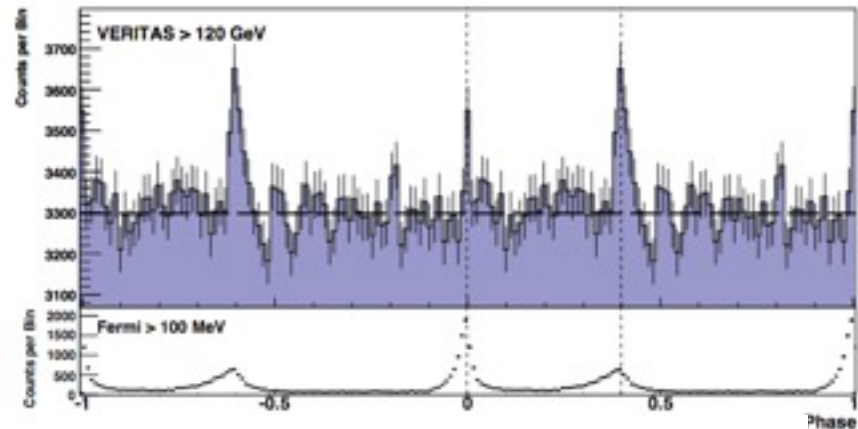
- MAGIC sees Crab at ~ 25 GeV



- With indication of non-exponential cut-off (2011)

- Cut-off is non-exponential!
- IP is brighter than MP (wink, wink)

- VERITAS sees Crab at > 150 GeV!



Curvature emission near light cylinder is excluded

- Astrophysical E-fields < B-field
- Equate acceleration by $E_{\parallel} = \eta B$ to curvature losses in $R_C = \xi R_{LC}$

Maximum possible energy break due to curvature emission

$$\epsilon_{br} = (3\pi)^{7/4} \frac{\hbar}{(ce)^{3/4}} \eta^{3/4} \sqrt{\xi} \frac{B_{NS}^{3/4} R_{NS}^{9/4}}{P^{7/4}} \left(\frac{r_{em}}{R_{LC}} \right)^{-9/4}$$

For Crab, assuming E=B

$$\approx 150 \text{ GeV}$$

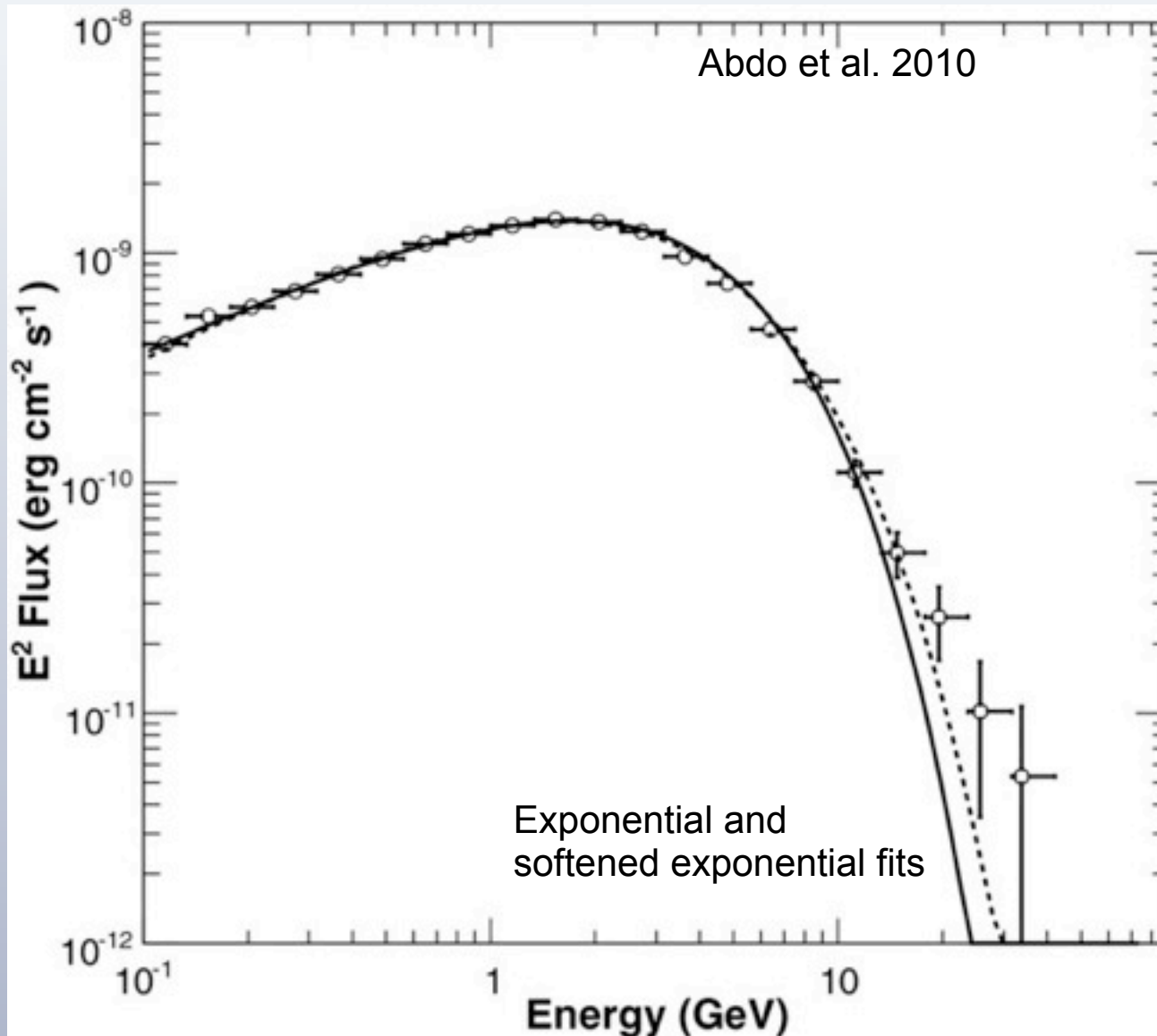
- Detection of Crab above 150 GeV (with non-exponential cut-off) exclude curvature emission from near the light cylinder (Lyutikov et al. 2011)
- E.g. Bai & Spitkovsky: emission within 5% of LC

Implications of Crab detection by VERITAS:

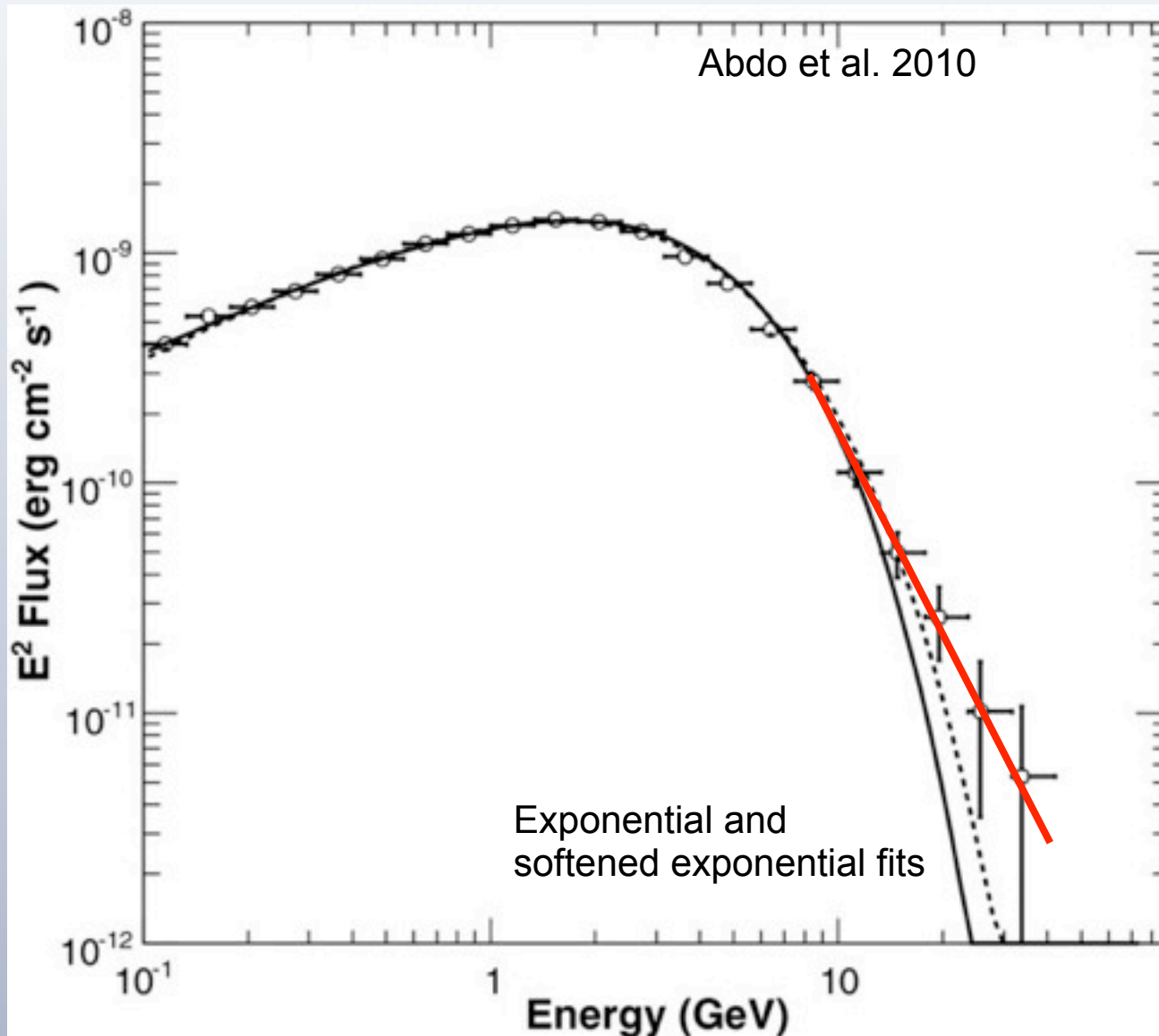
- Spectral break in Crab is not due to curvature emission of the maximal energy of particles
- Alternative possibility: IC scattering
- Break due to the details of particle distribution and scattering cross-section (in the KN regime)

- Is Crab special (e.g. high level of soft photons)?
- What about other pulsars?
 - Vela
 - Geminga

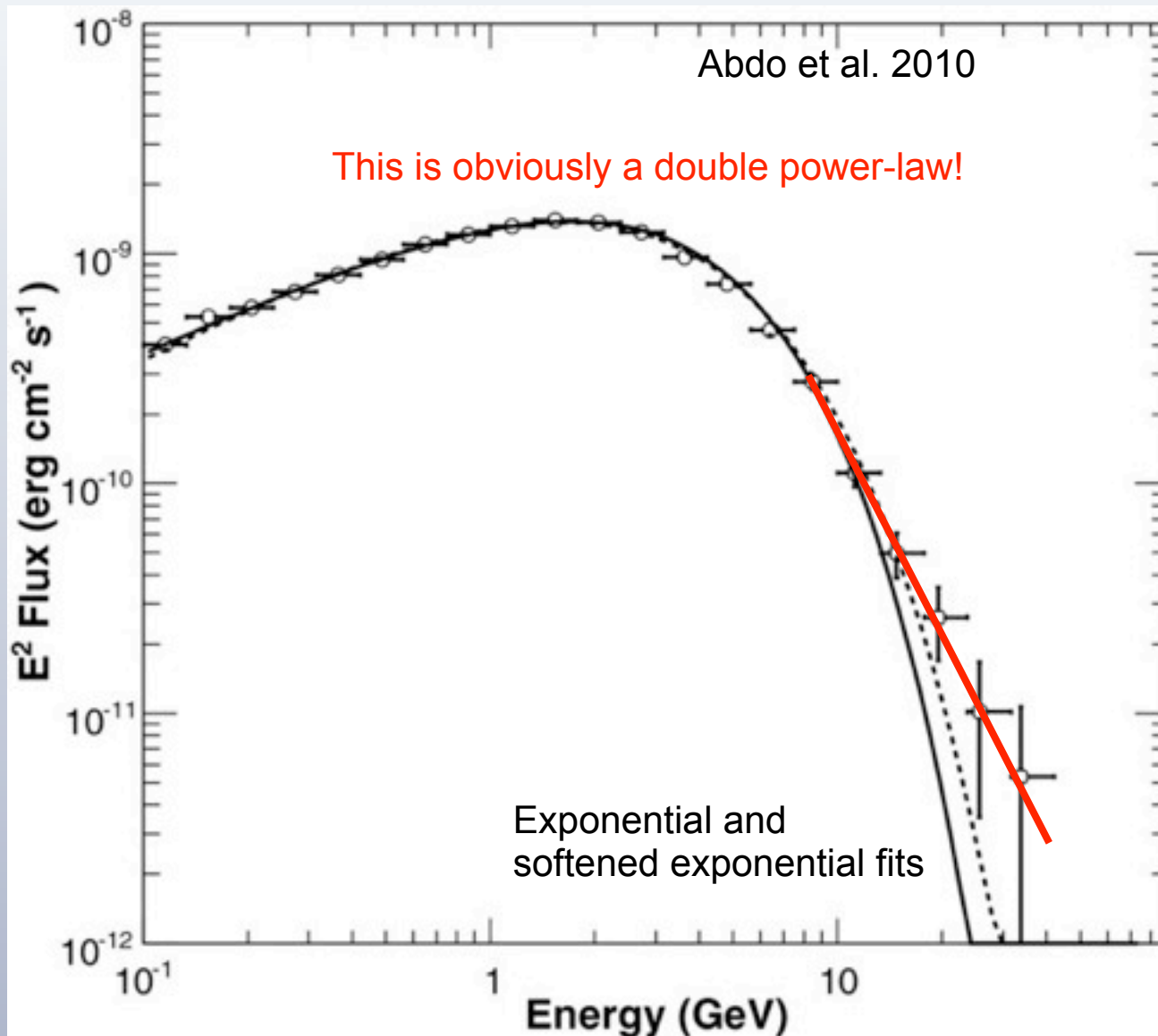
Fermi spectrum of Geminga



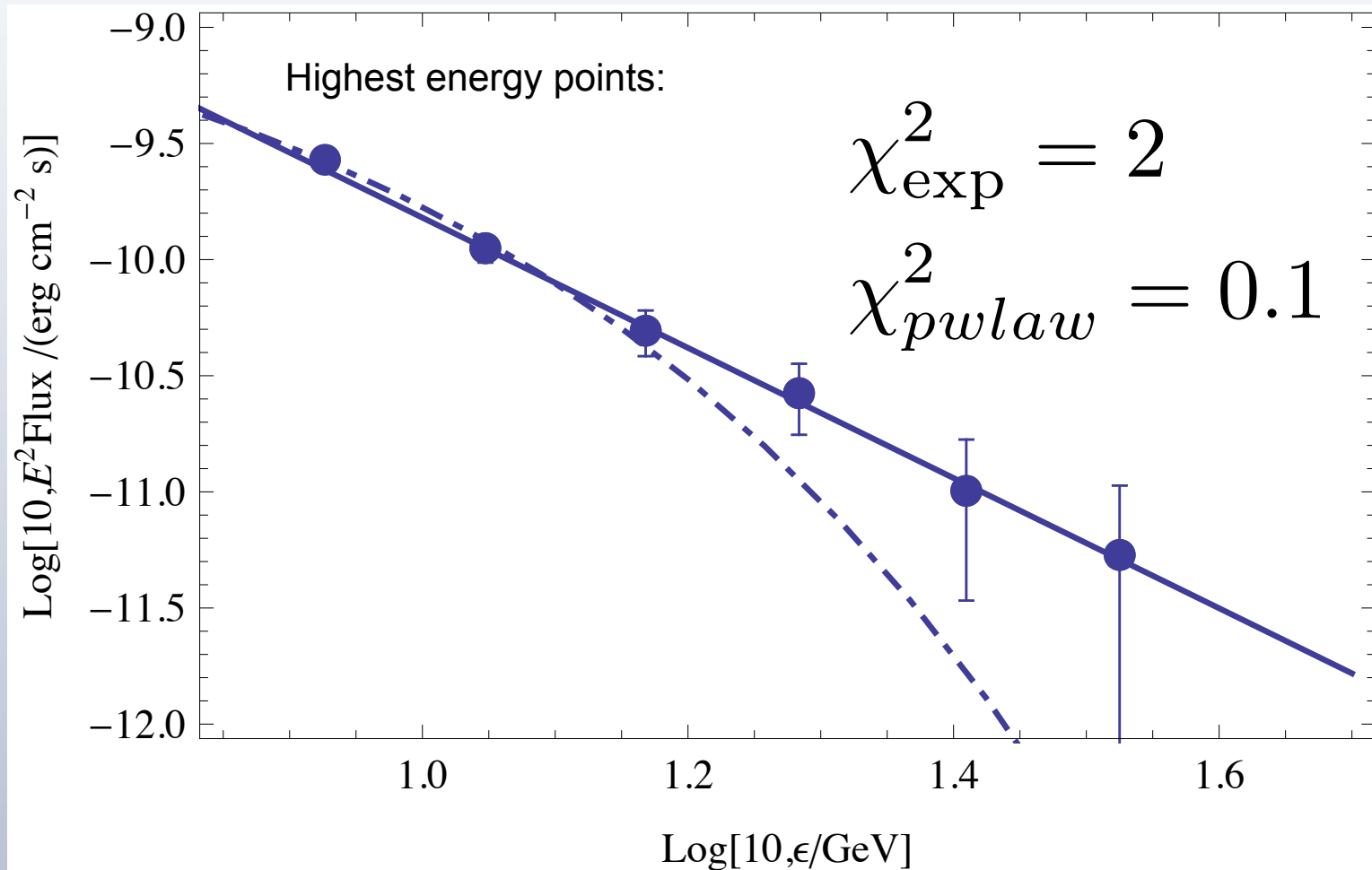
Fermi spectrum of Geminga



Fermi spectrum of Geminga



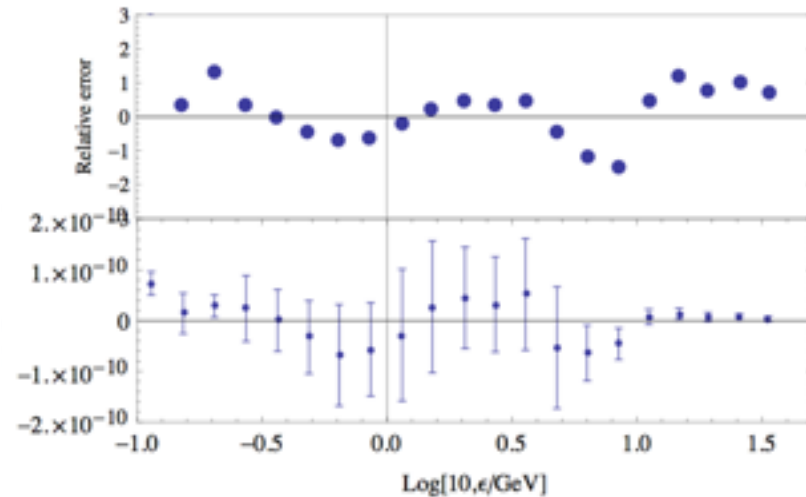
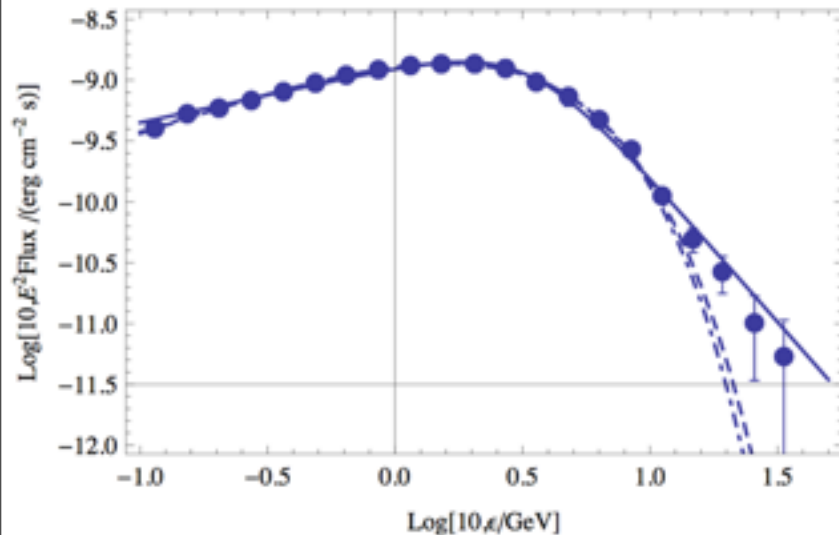
Geminga: fits



- The highest energy data points actually have the **smallest** error bars.
- Too broad energy bins?

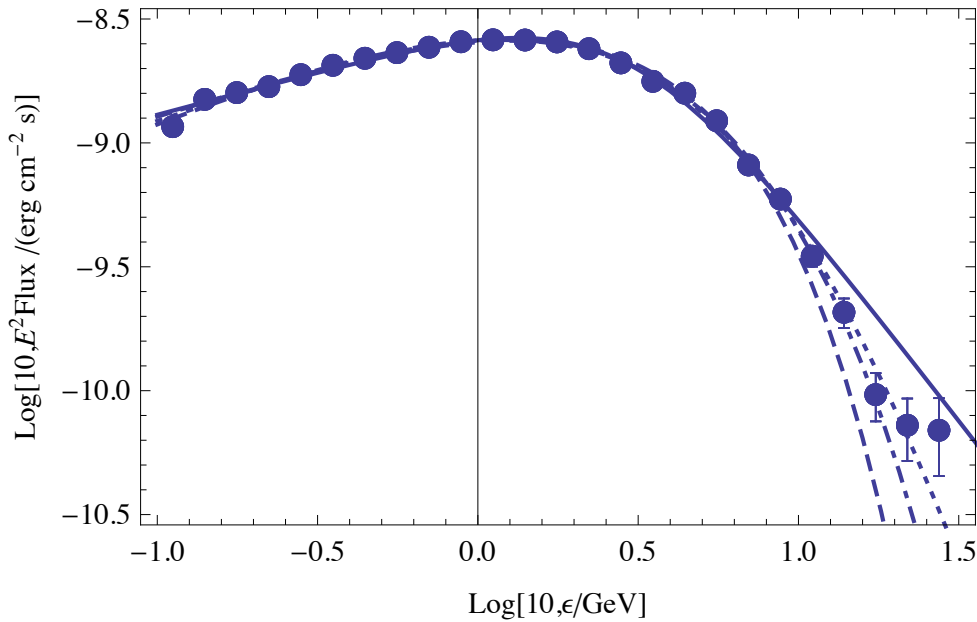
Geminga: broad band fits

errors χ^2 for double pwlaw



- The errors are not random
- **Most of the χ^2 is accumulated near the break energy due to the ARBITRARY parametrization of the spectral roll-off**
- Similar results for phase-resolved spectra

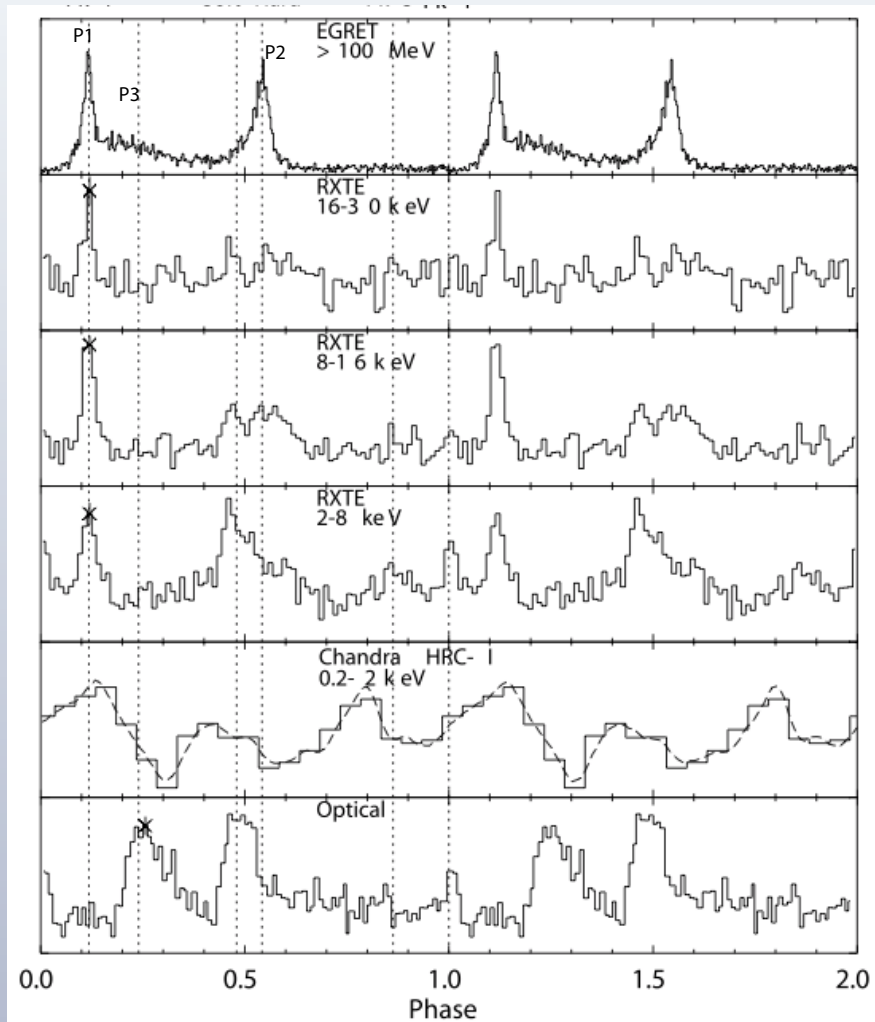
Vela



- Double power-law are as good as exp.
- Various parametrizations of roll-offs reduce χ^2

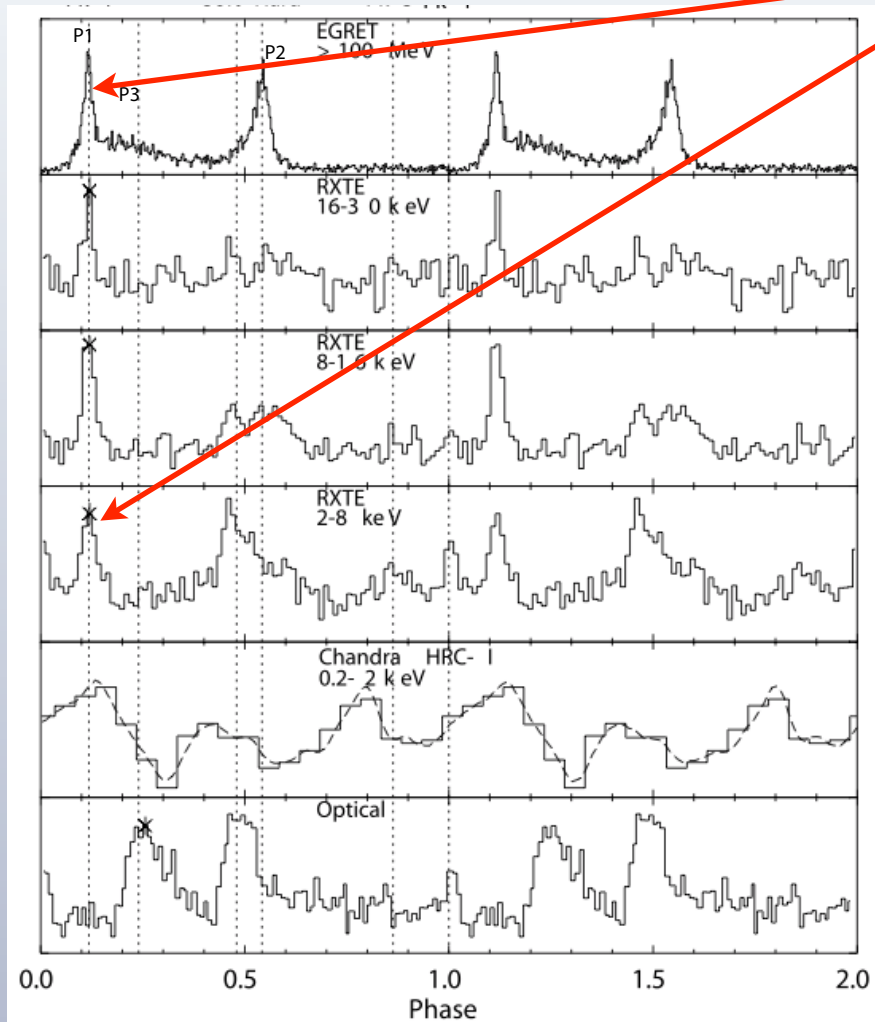
model	Fit function	α	β	ϵ_{br}	reduced χ^2	b	dof
<i>a</i>	$\left(\left(\frac{\epsilon}{\epsilon_{br}} \right)^\alpha + \left(\frac{\epsilon}{\epsilon_{br}} \right)^{-\beta} \right)^{-1}$	1.68	0.34	2.91	2.2	—	21
<i>b</i>	$\left(\left(\frac{\epsilon}{\epsilon_{br}} \right) + \left(\frac{\epsilon}{\epsilon_{br}} \right)^{-\beta} \right)^{-\alpha}$	4.11	0.11	8.8	1.3	—	21
<i>c</i>	$\epsilon^\beta e^{-\frac{\epsilon}{\epsilon_{br}}}$	—	.42	3.1	2.0	—	22
<i>d</i>	$\epsilon^\beta e^{-\left(\frac{\epsilon}{\epsilon_{br}} \right)^b}$	—	0.59	1.58	1.4	0.73	21

Vela profiles:



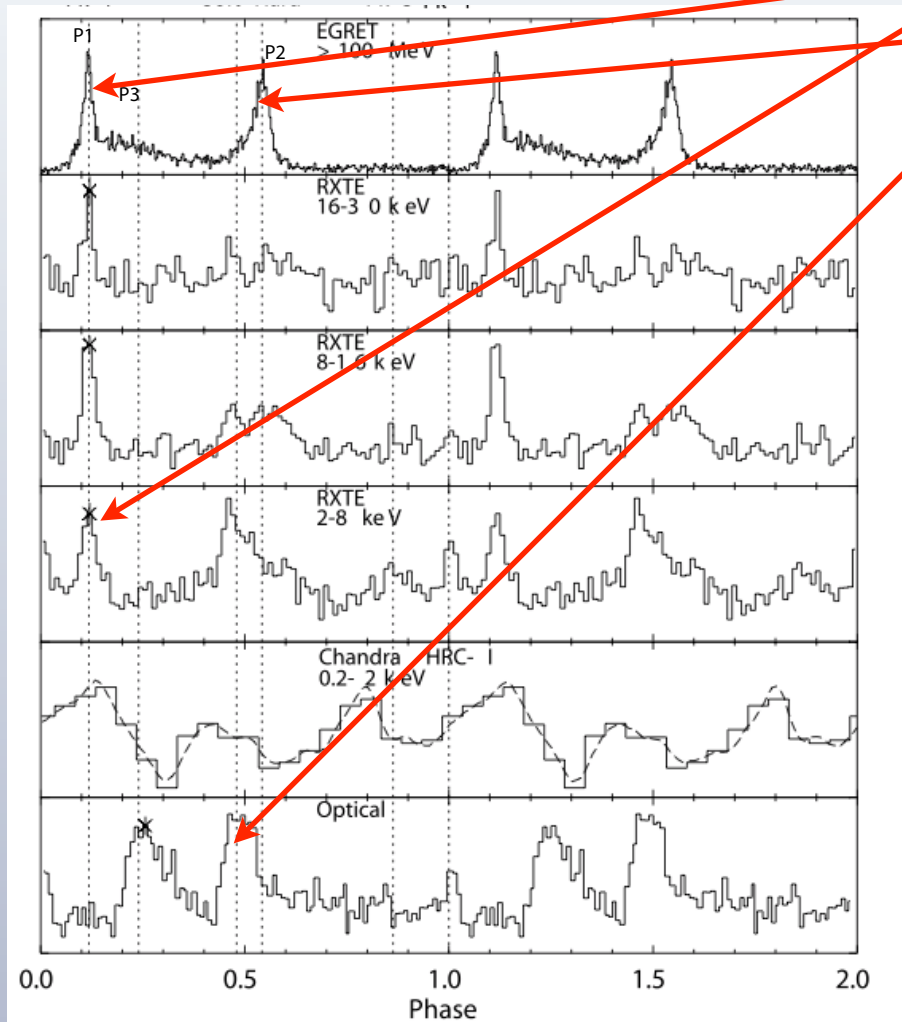
- All gamma-ray peaks have mirror images at lower energies
- No simple 1:1 correspondence: anisotropy and KN effects

Vela profiles:



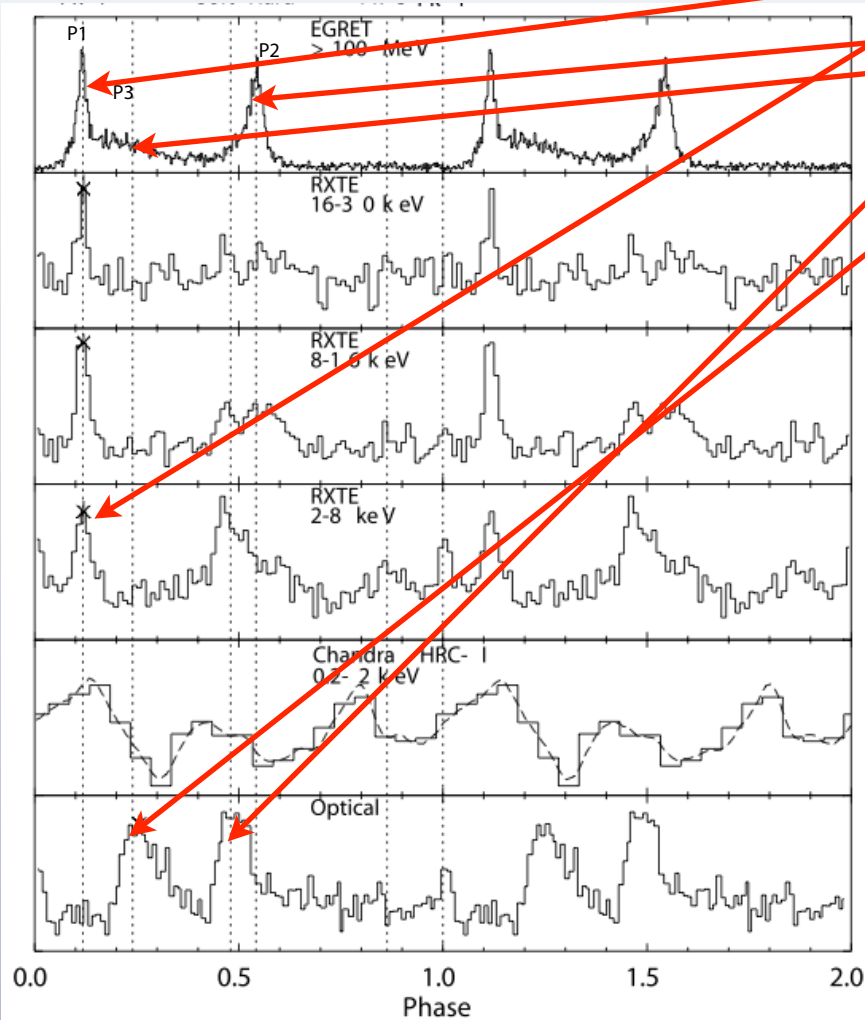
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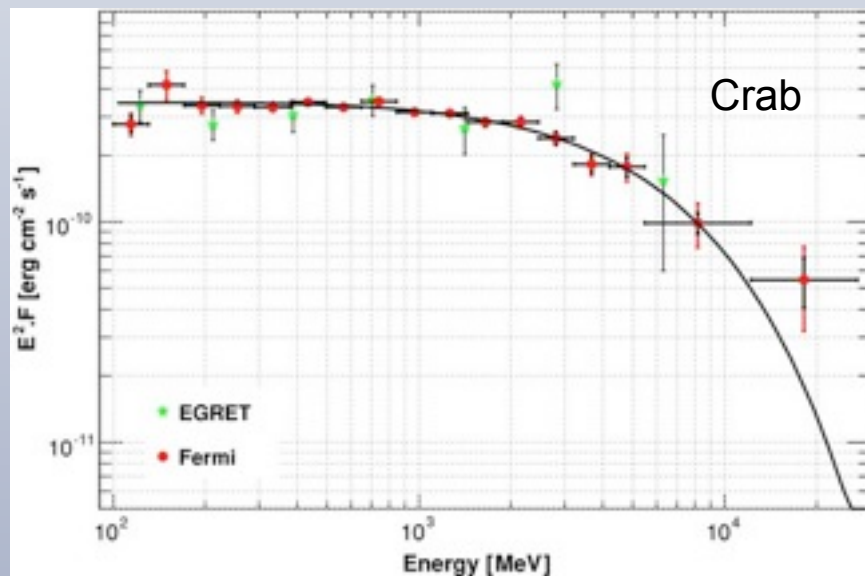
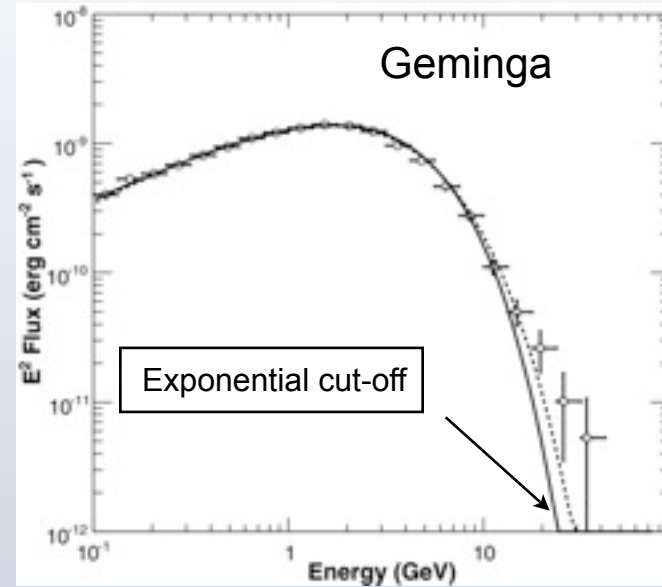
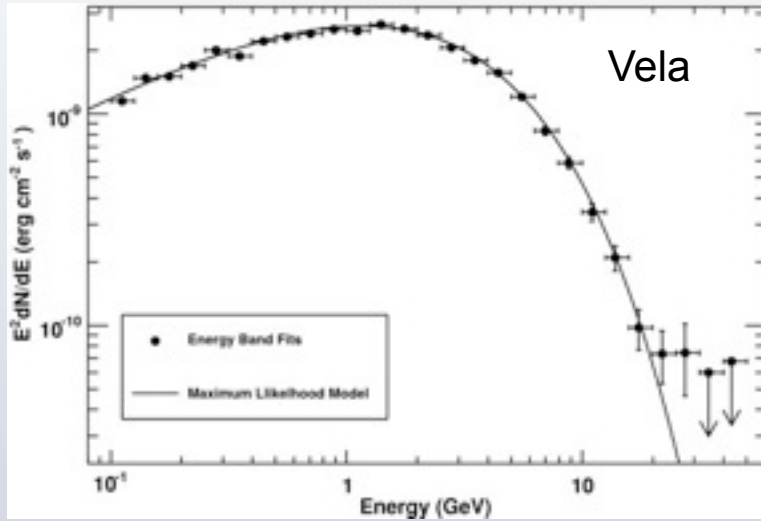
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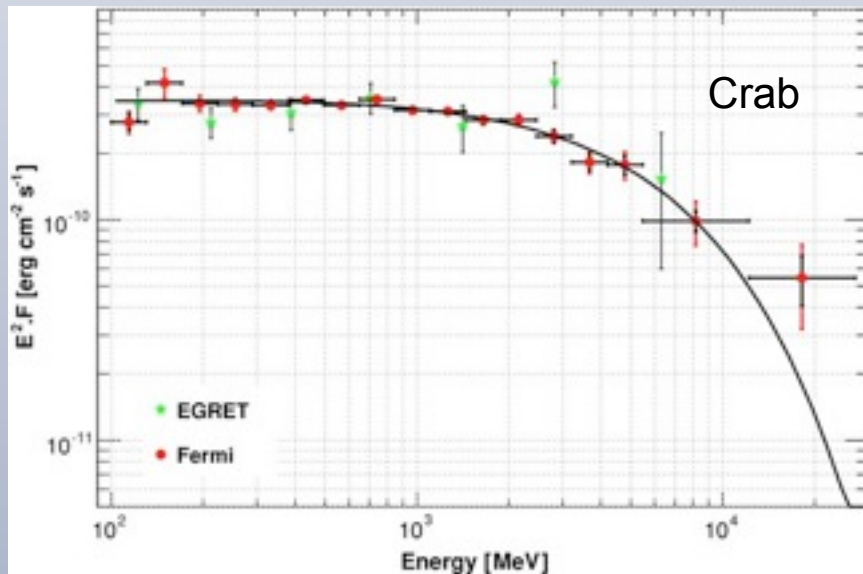
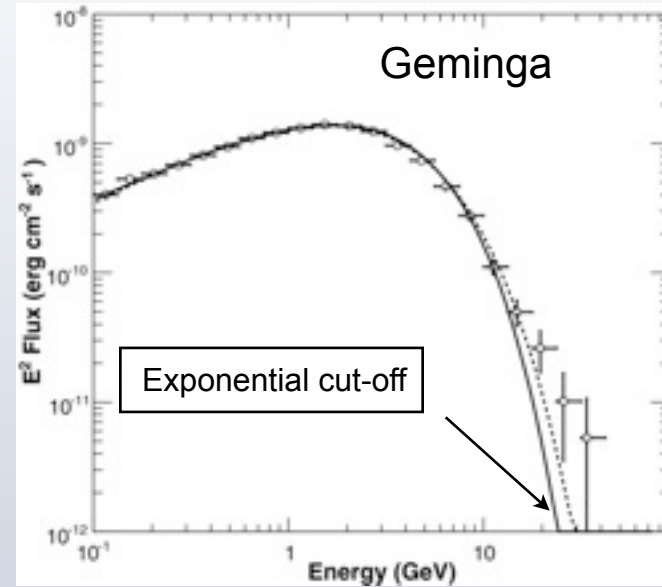
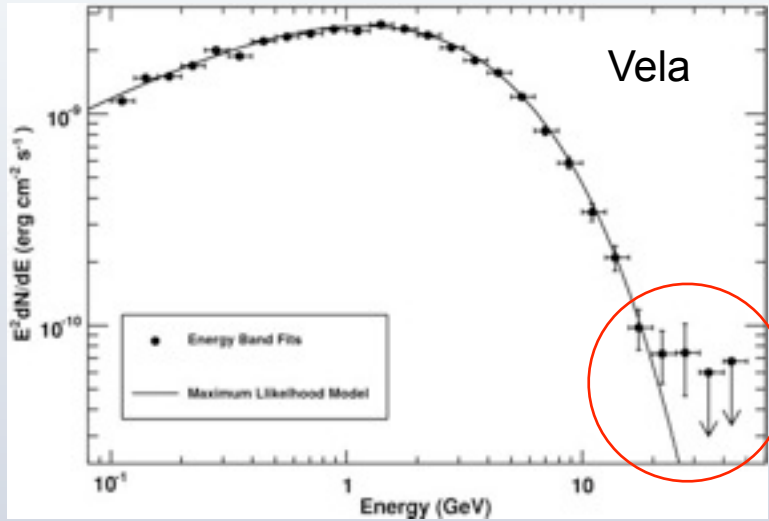


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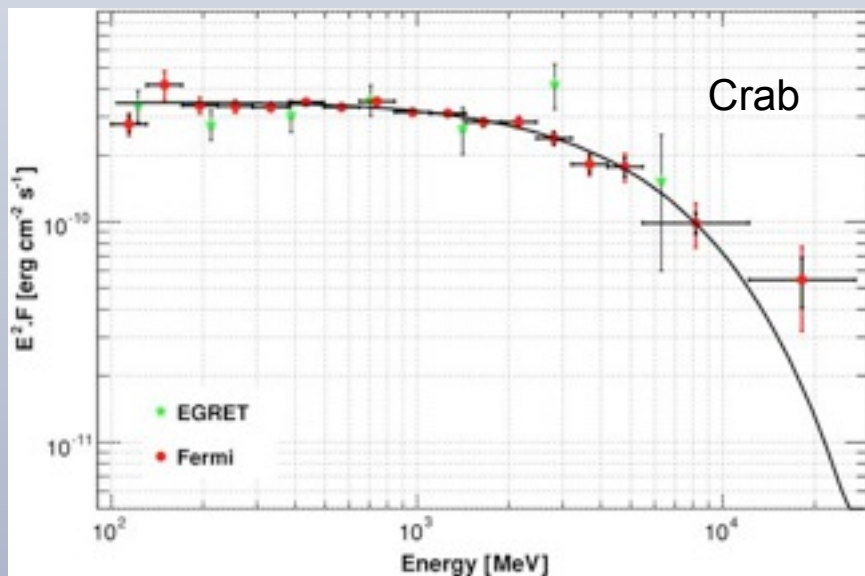
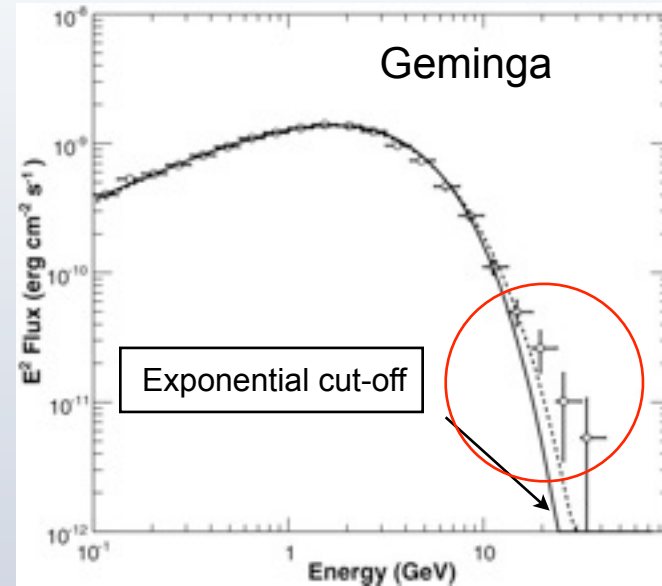
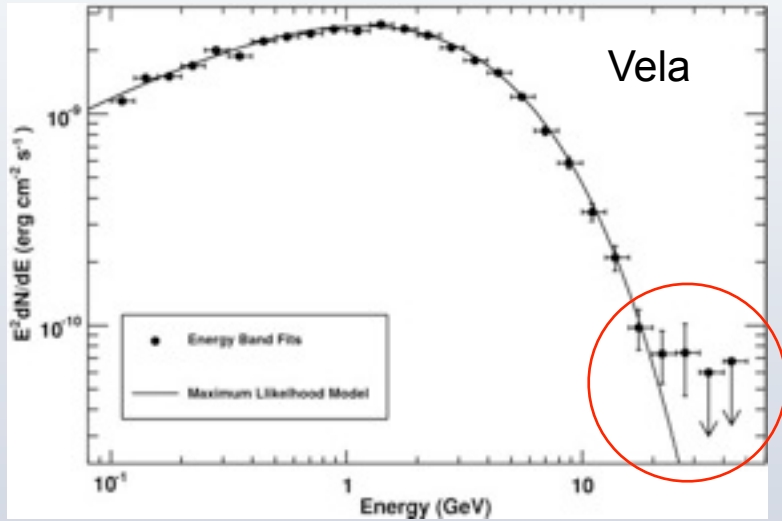
Looking back at Fermi spectra:



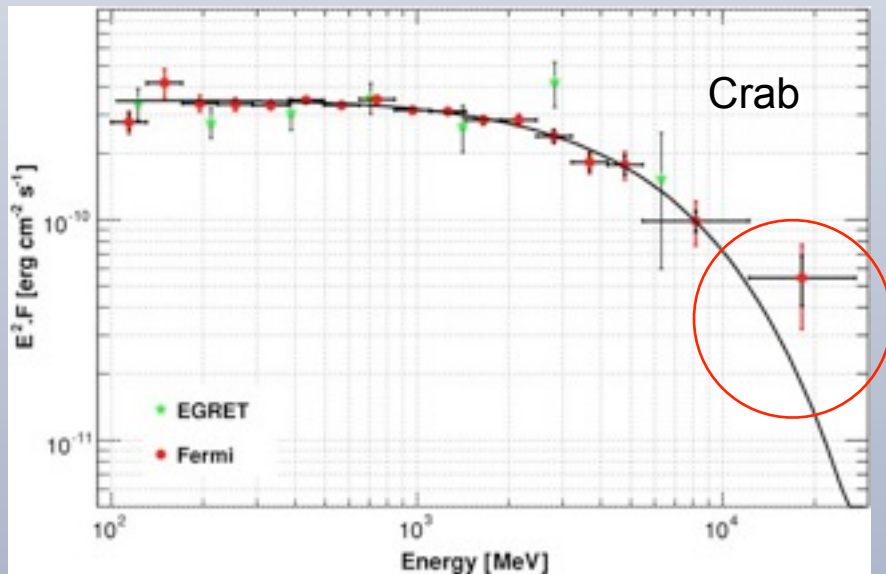
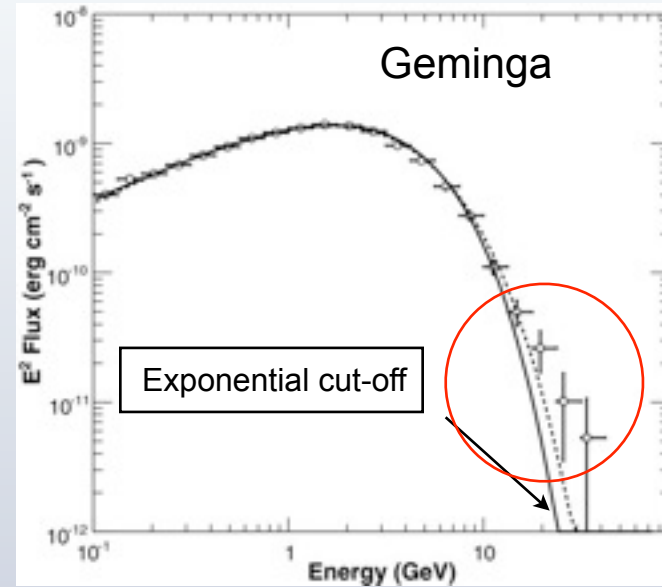
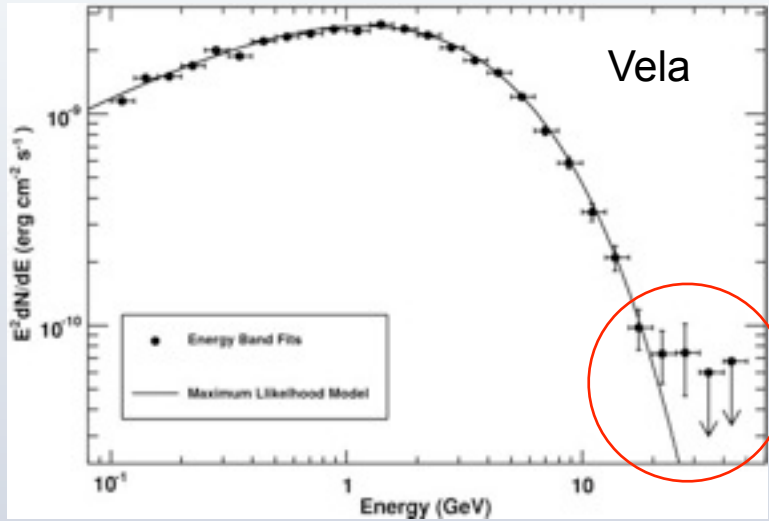
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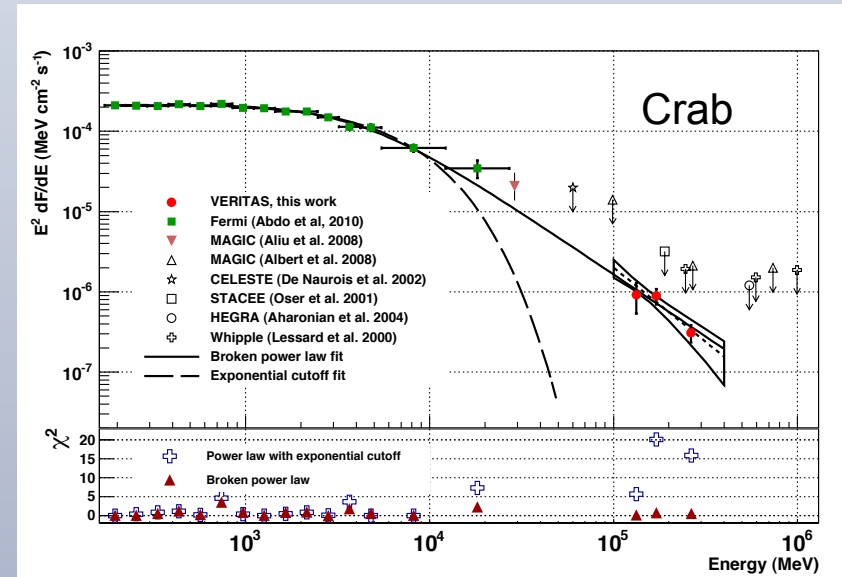
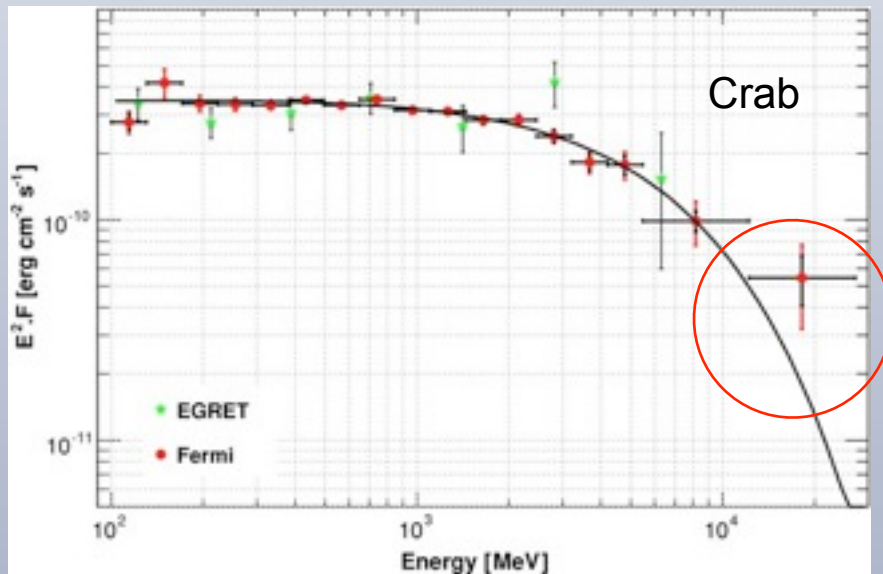
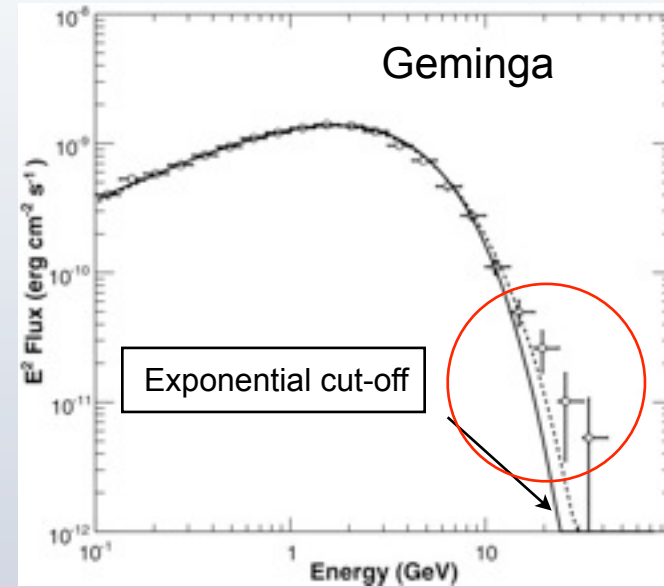
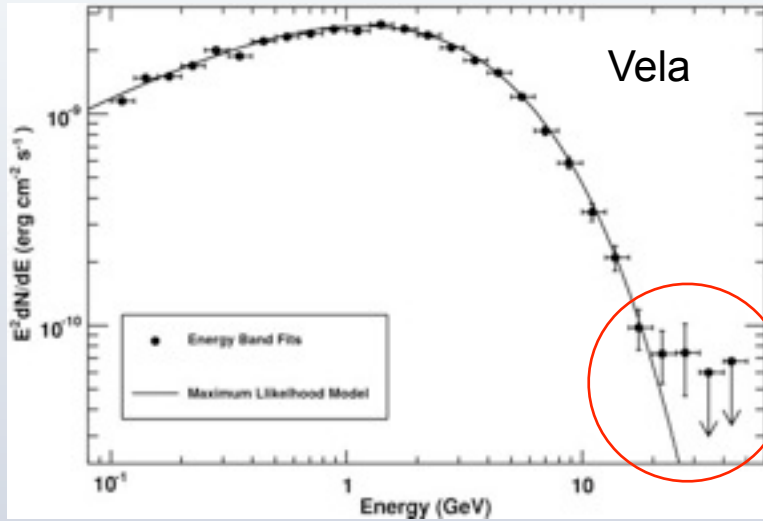
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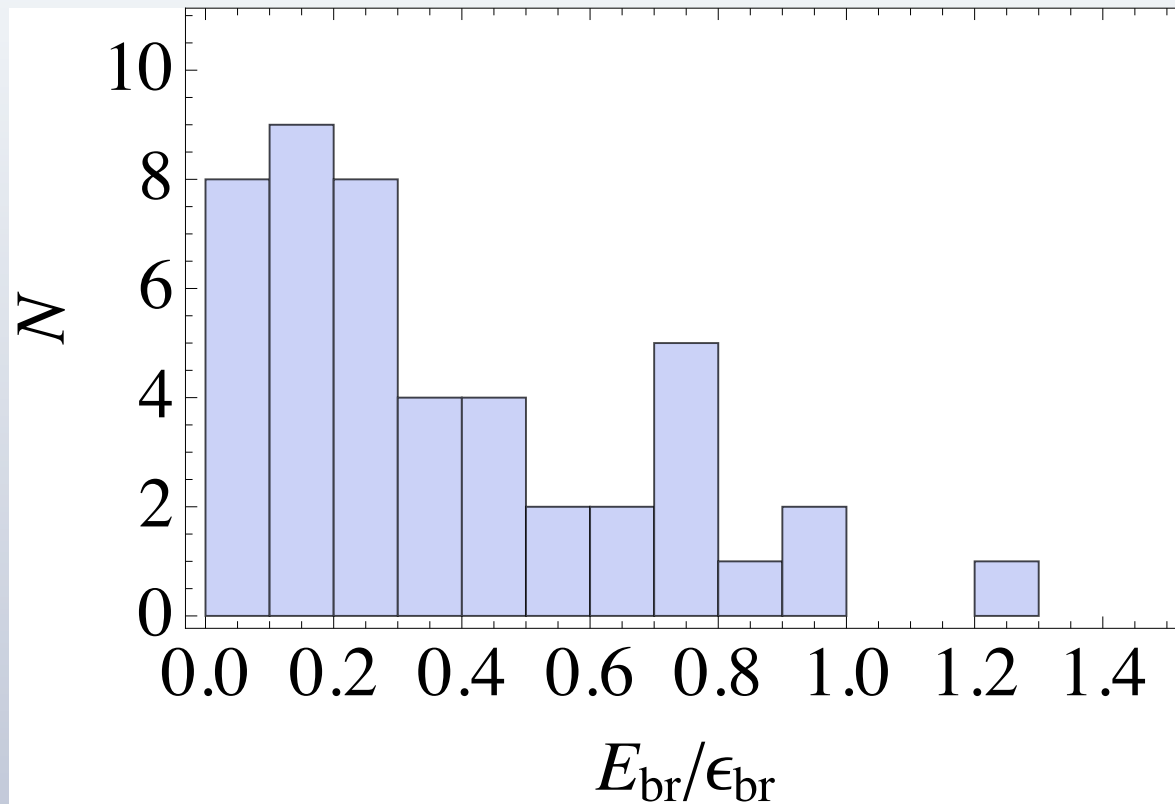
Looking back at Fermi spectra:



Looking back at Fermi spectra:



Other pulsar: maximal curvature energy at light cylinder



- Ratio of the observed break energies E_{br} for 46 pulsars to the maximum predicted for curvature radiation ϵ_{br}
- For Crab $E_{br}/\epsilon_{br} \sim 0.05$ seemed OK, but not OK \rightarrow Lower limits

Scaling of L_γ with E_{sd}

- Acceleration by $E_{||} =$ radiation losses

$$L_\gamma = ec \times \eta B \times n_{GJ} \times \eta_G R_{LC}^3$$

$$L_\gamma \approx \eta \eta_G \dot{E}_{SD}$$

- **Linear!**
- (Everyone uses $L_\gamma \propto \sqrt{\dot{E}_{SD}}$ - potential-limited, from polar caps)
- Known problem of too many observed pulsars
- uncertainty of the geometrical parameter η_G and accelerating E-field η

Upshot: pulsar spectral modeling

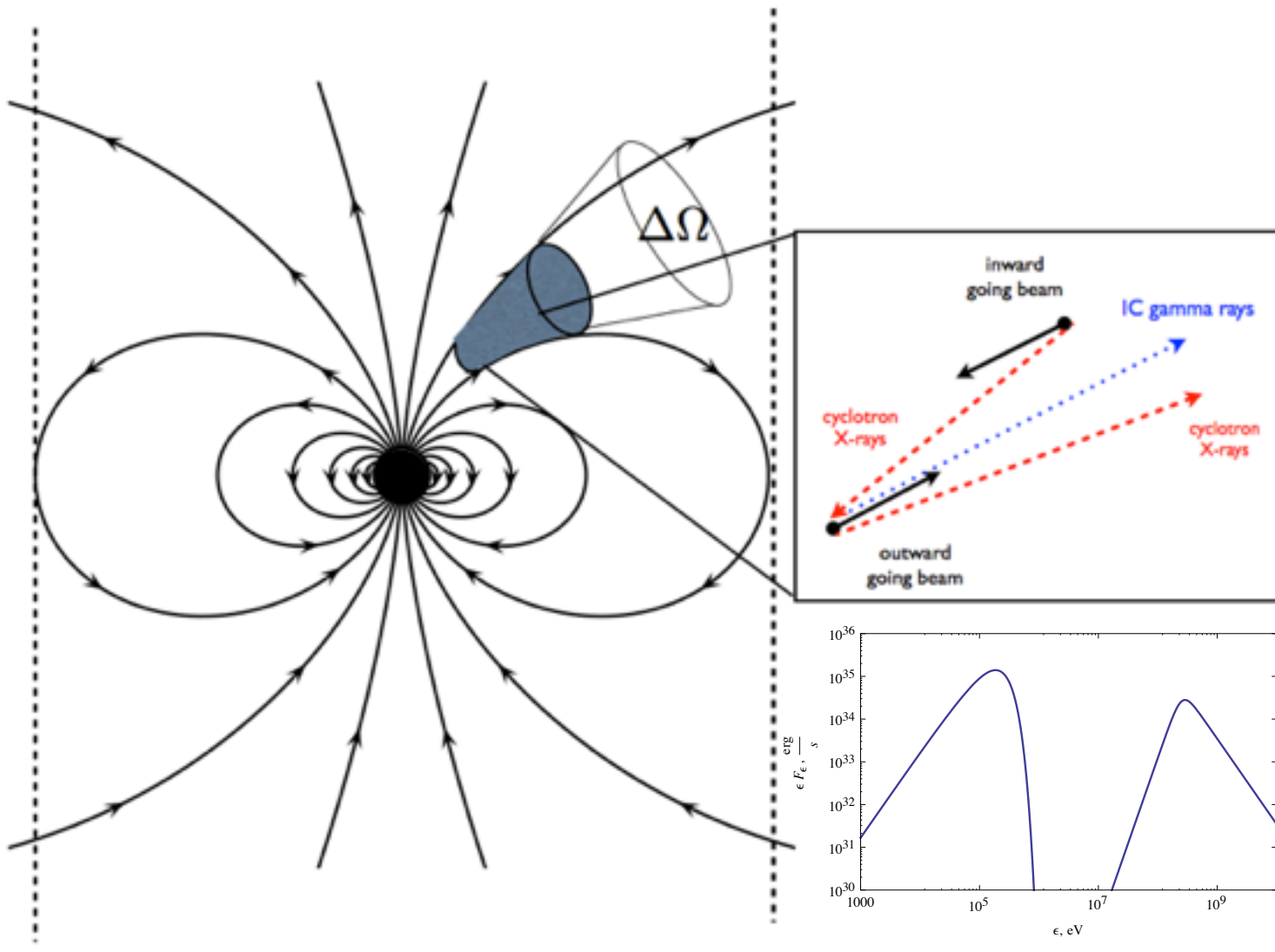
- Broadband pulsar high energy spectra of the brightest gamma-ray pulsars are generally consistent with the double power law shapes.
 - Crab and Geminga are inconsistent with exp. cut-off
 - Vela is consistent with double power law
- In Geminga most errors in the broadband double power law fit are accumulated due to the **arbitrary** parametrization of the spectral roll-off between the two asymptotic power laws
 - Due to low background above ~ few GeV Fermi is especially sensitive there
- **This favors the inverse Compton (IC) scattering for the production of gamma-ray photons**

IC model (in progress)

- ``Off-the shelf'' SSC models not applicable
 - Random B-field of given value
 - Isotropic particle distribution
 - single value for bulk motion
- Regular B-field, changing sharply in the emission region
- Strong radiative damping: non-isotropic distribution
- Continuous $v_{||}$

- Need two distinct components:
 - optical-X-ray: boosted cyclotron-synchrotron
 - gamma rays: IC scattering of the ``first'' cyclotron photons by the ``second'' particles

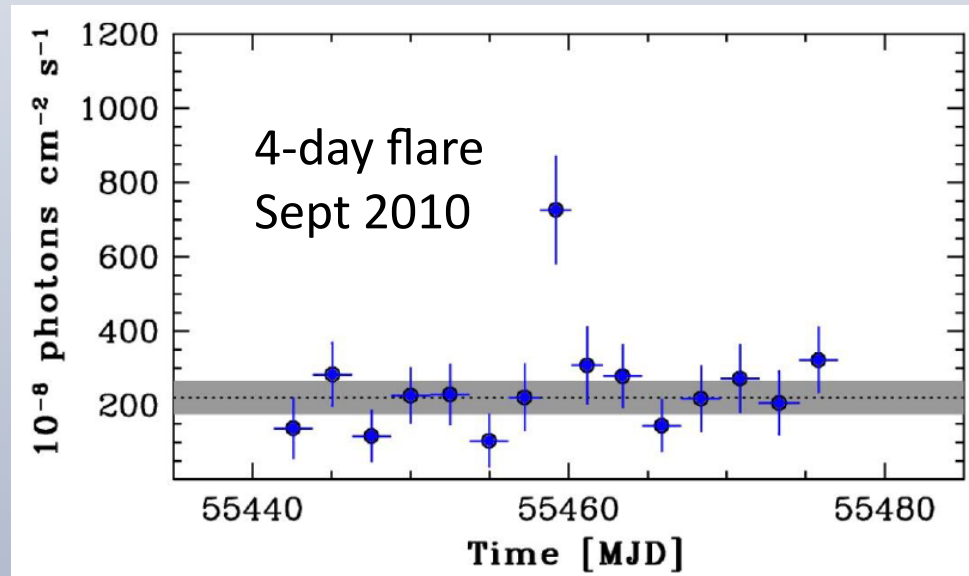
IC model (in progress)



$$\beta_0 = 3 \cdot 10^{-6}$$
$$\lambda = 5 \cdot 10^6$$

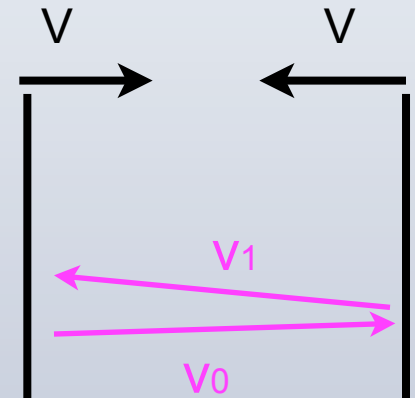
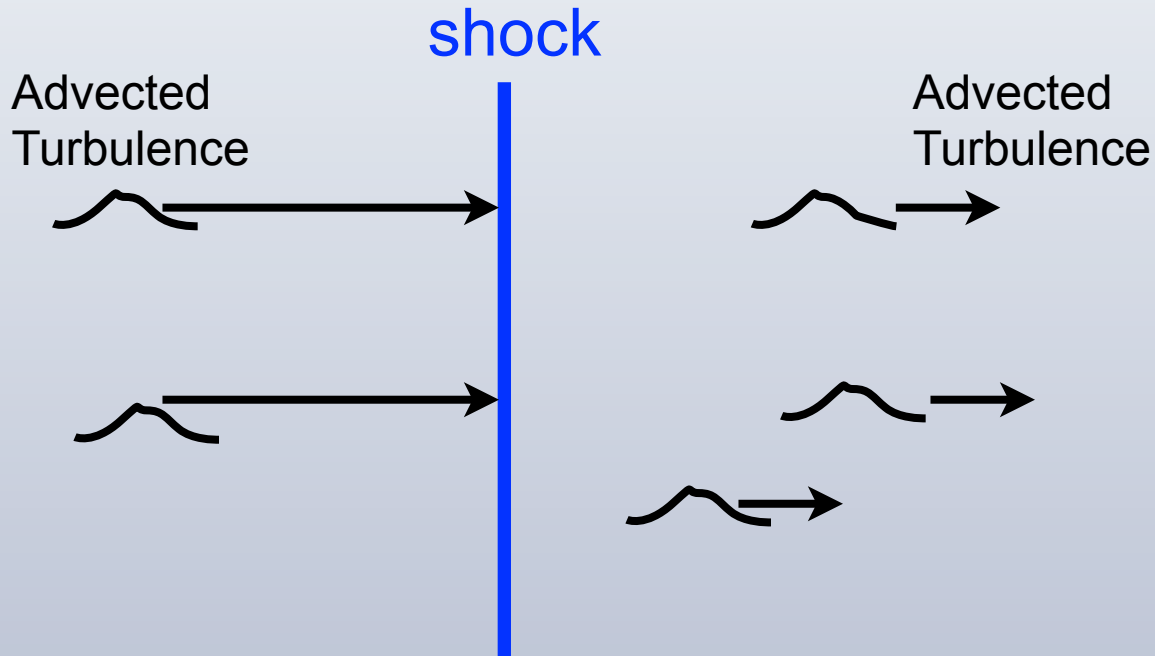
2. Crab nebula flares seem inconsistent with the stochastic shock acceleration

(with Eric Clausen-Brown)



Tavani et al. 2011

The accepted model of Cosmic Ray acceleration: Fermi-I at shocks

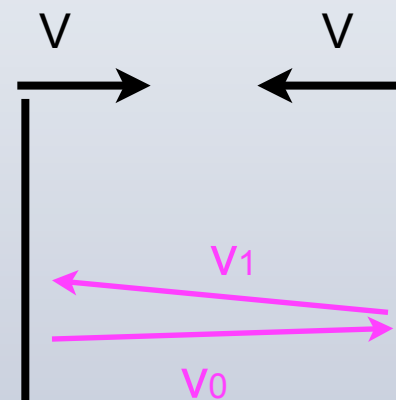
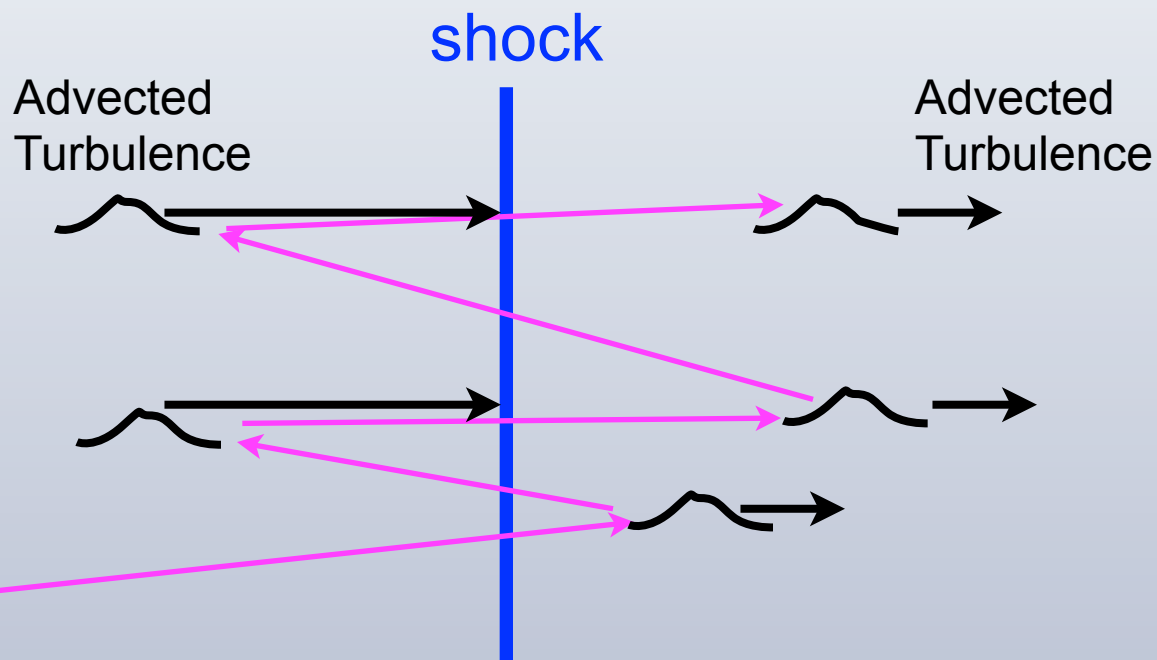


$$v_0 + V = v_1 - V$$

$$v_1 = v_0 + 2V$$

$$\frac{\delta\epsilon}{\epsilon} \sim \frac{v}{c}$$

The accepted model of Cosmic Ray acceleration: Fermi-I at shocks



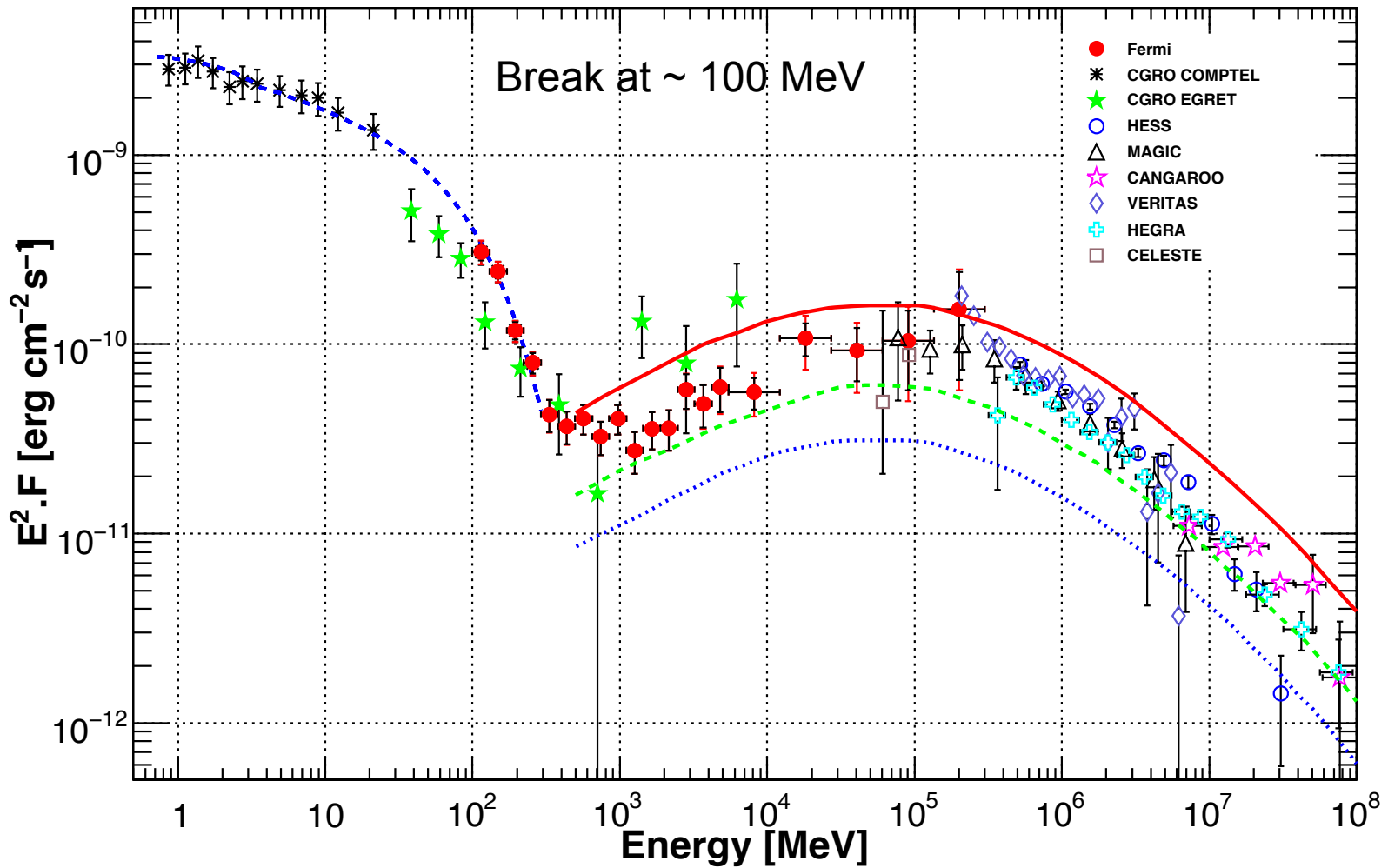
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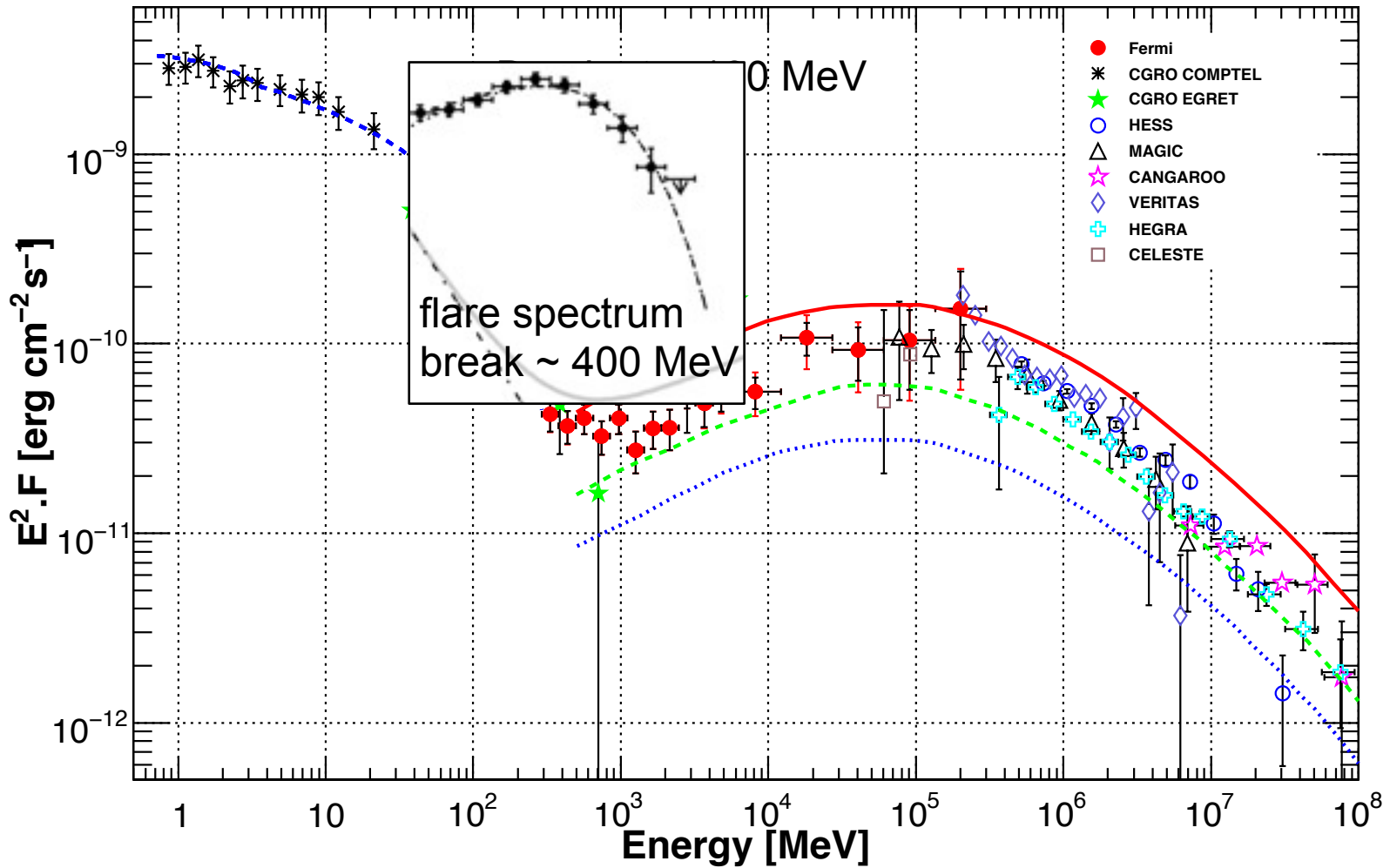
Crab nebula

Tavani et al. 2010
Beuhler et al., 2011



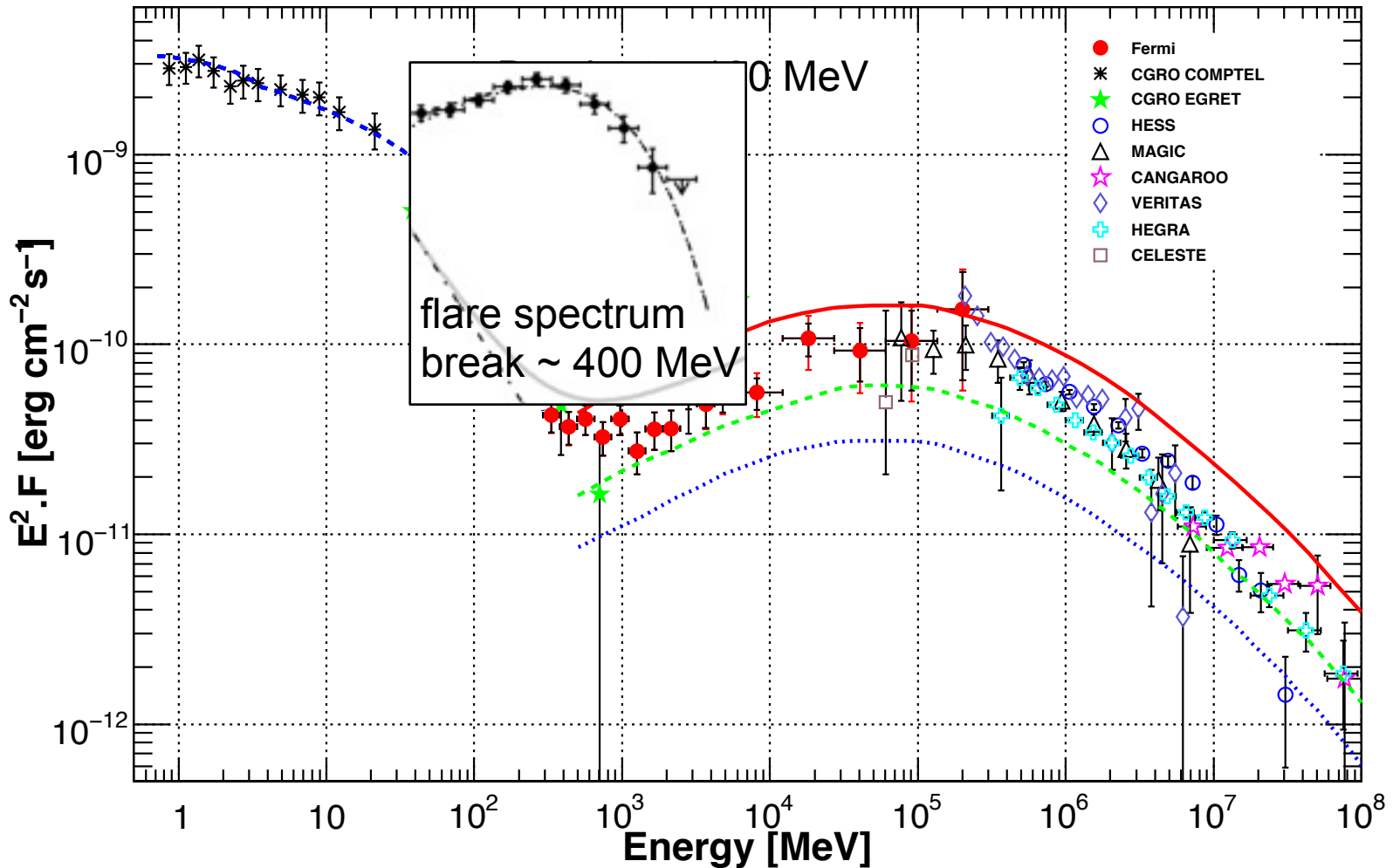
Crab nebula

Tavani et al. 2010
Beuhler et al., 2011



Crab nebula

Tavani et al. 2010
Beuhler et al., 2011



Flare time scales \ll nebular dynamical time, termination shock size

Upper limit to synchrotron frequency

Accelerating E-field < B-field

$$eEc = \eta eBc = \frac{4e^4}{9m^2c^3} B^2 \gamma^2$$

$$E_p = \frac{27}{16\pi} \eta \frac{mhc^3}{e^2} = 236 \eta \text{ MeV.}$$

Lyutikov '10,
Komissarov & Lyutikov '11
de Jager '98 (for shocks)

- Same as Fermi acceleration on inverse gyroscale
(requires very efficient scattering, stochastic
acceleration: $\eta \ll 1$)

I am not aware of any stochastic acceleration model
that gives $\eta \sim 1$. Typically $\eta \sim 10^{-2}$.

Possible resolutions.

Doppler boosting

$$\omega \rightarrow \delta\omega$$

$$I \propto \delta^{2+\alpha}, \delta^{3+\alpha} \rightarrow \delta^{3-4}$$

$$\delta = \frac{1}{\Gamma(1 - \beta \cos \theta)}$$

Even delta ~ 2 can increase the flux by 10-20 times

Doppler boosting

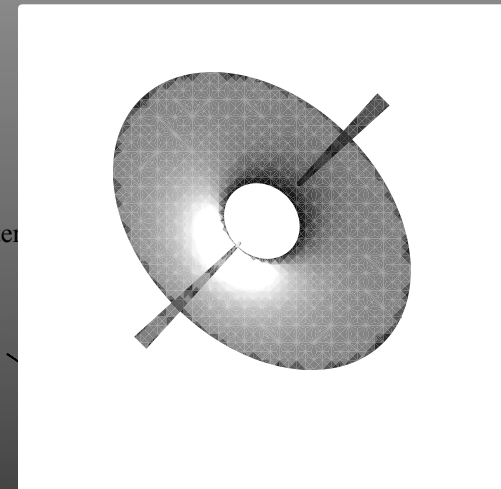
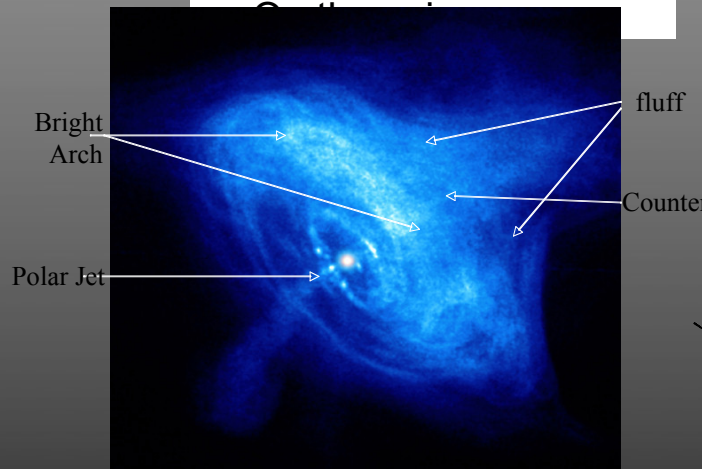
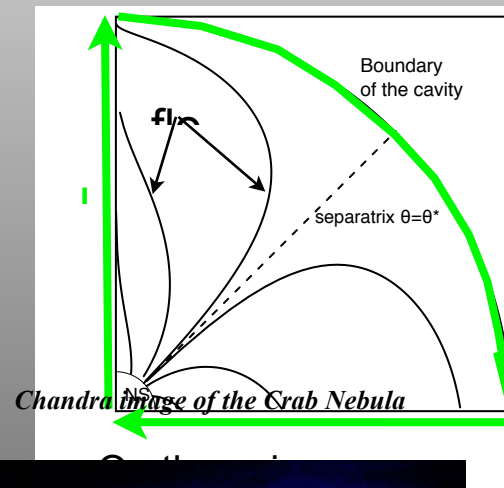
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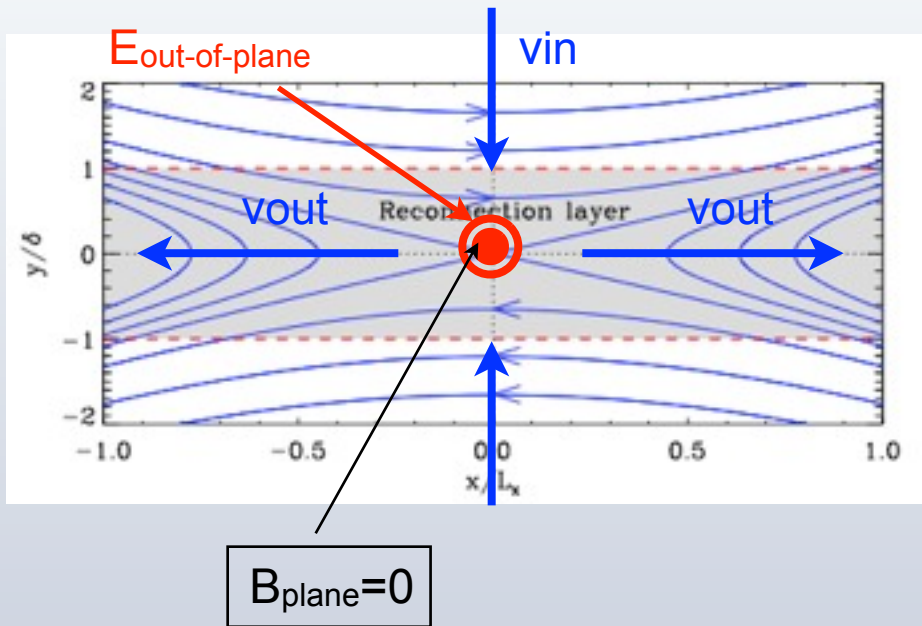
Even delta ~ 2 can increase the flux by 10-20 times

Lyutikov (2010): Ideal flow in the bulk, dissipation on boundary



***Paradigm shift: some (most?)
particles are accelerated by
magnetic reconnection (and
not shocks)***

Reconnection: $E \sim B$

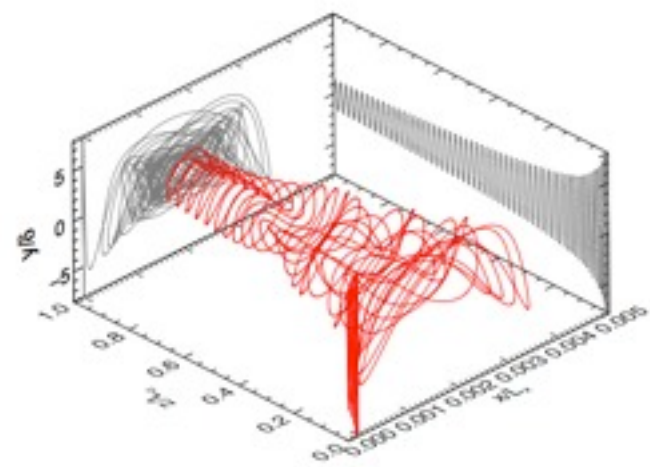


- Relativistic reconnection: Lyutikov & Uzdensky, Lyubarsky, Hoshino
- $E \sim (v_{in}/c) B$.
- Inflow velocity $v \sim c$, $E \sim B$

$$\beta_{in} \sim \sqrt{\frac{\sigma}{S}}$$

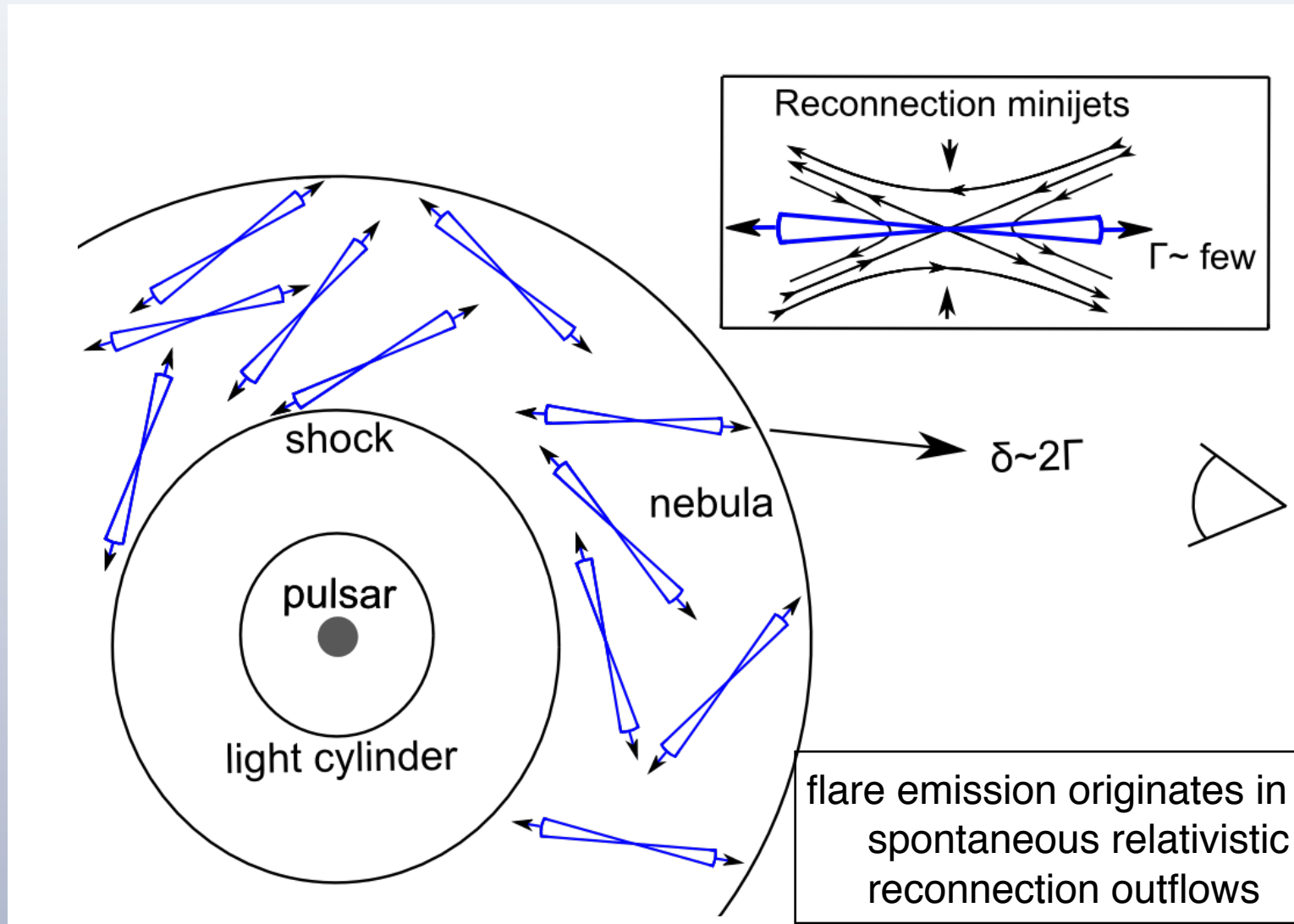
$$\gamma_{out} = \gamma_{in} (1 + \sigma)$$

Uzdensky et al.: Accelerate in a region where B is small, with $E > B$, emit where B is large.

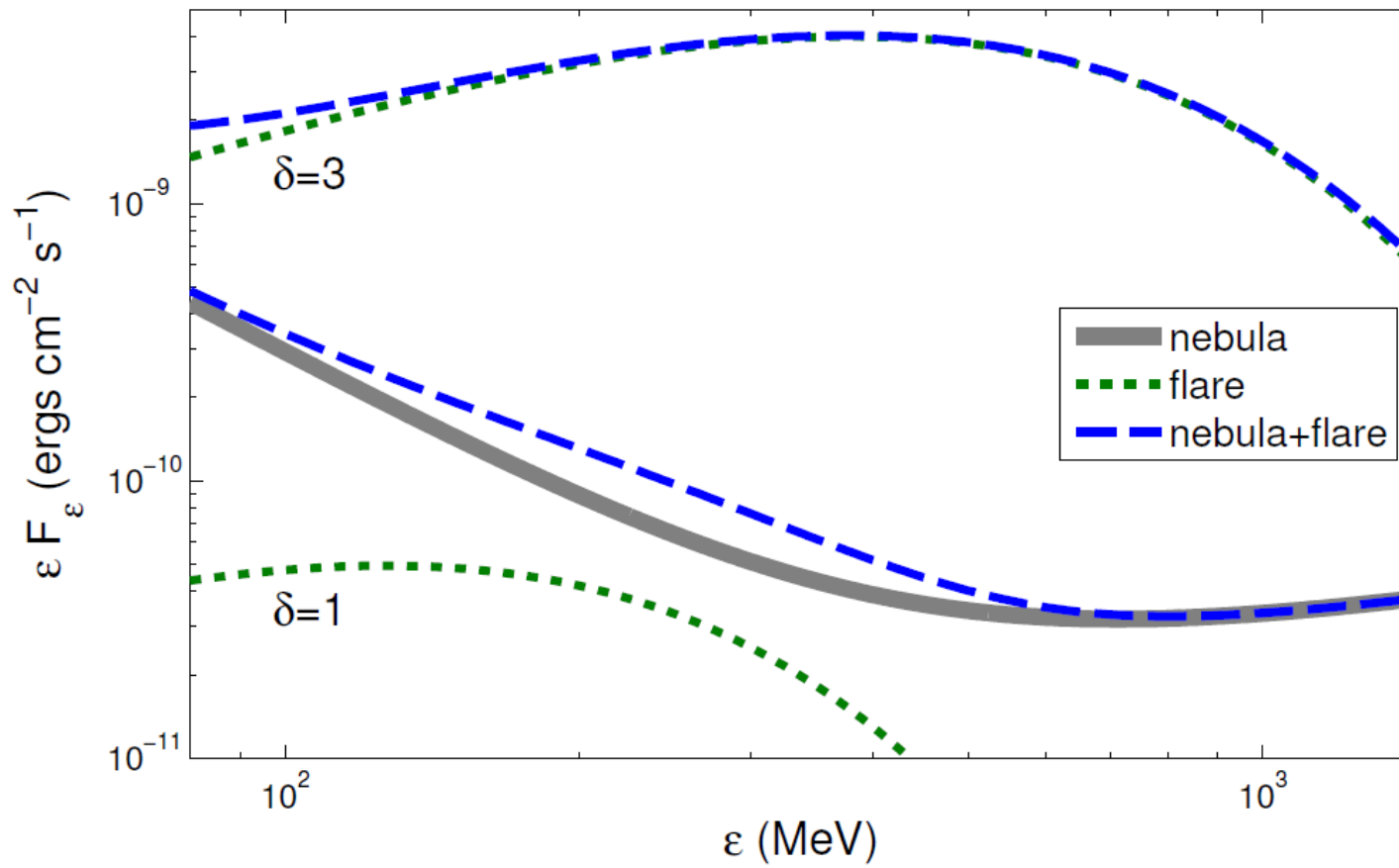


Reconnection + bulk motion

Clausen-Brown & Lyutikov 2012

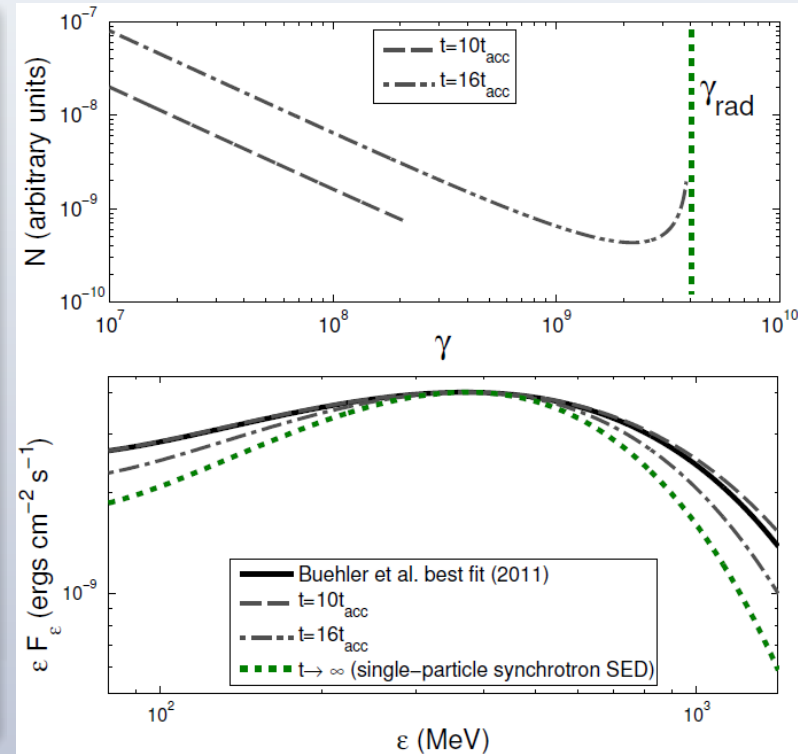


$\Gamma \sim$ few increases flux and peak energy



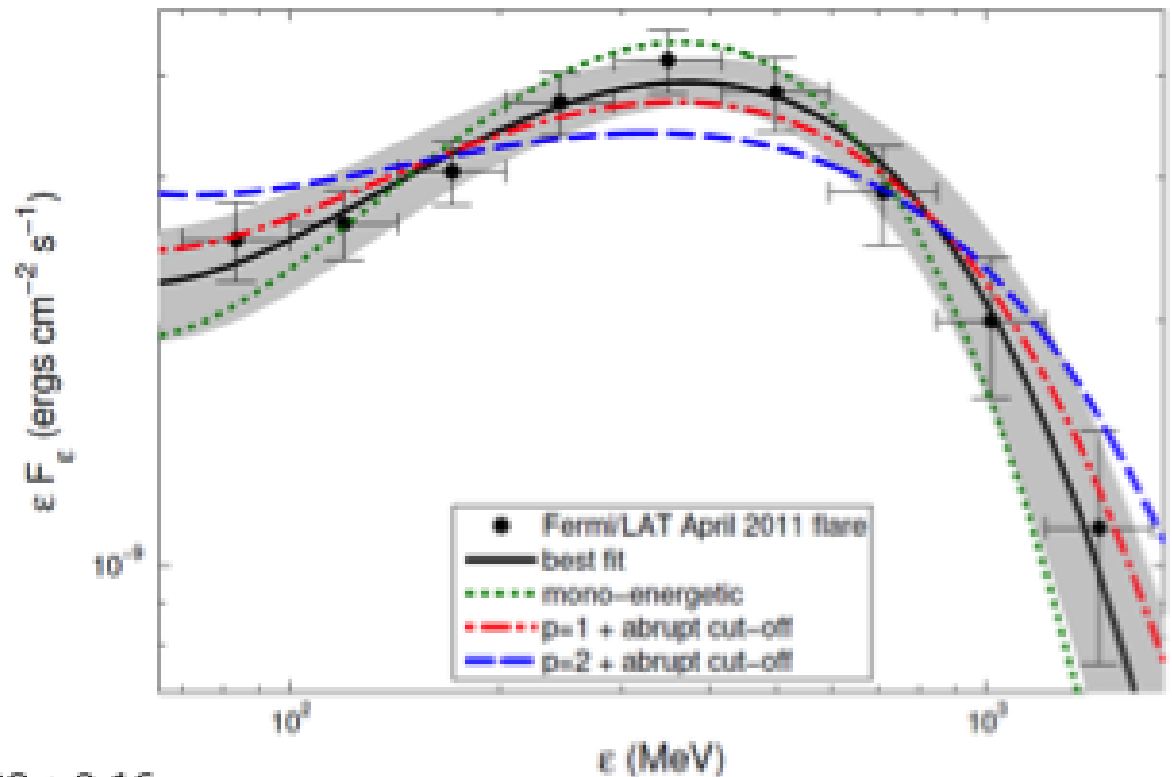
Radiation reaction-limited acceleration: pile-up

- Many reconnection models predict a particle spectrum of power-law $p=1$
- When this particle spectrum approaches radiation limit, a “pile-up” distribution is produced
- As pile-up grows, the SED approaches that of a single-particle synchrotron SED



Gray area is one-sigma error region from fitting data to:

$$\Phi = A\epsilon^{-\gamma_f} \exp\left(-\frac{\epsilon}{\epsilon_{cut}}\right)$$



- Best fit for this SED is $\gamma_f = 1.08 \pm 0.16$
 - $\gamma_f = 0.7$ approximates single-particle synchrotron SED
 - $\gamma_f = 1.3$ approximates $p=1$ particle distribution
 - $\gamma_f = 1.56$ approximates $p=2$ particle distribution (shock acceleration)
- This is consistent with power-law of index $p = 1$ with pile-up (previous slide), or with distribution of power law index $p = -0.2$
- Observed SED inconsistent with shocks, but is in line with reconnection particle spectrum + pile-up

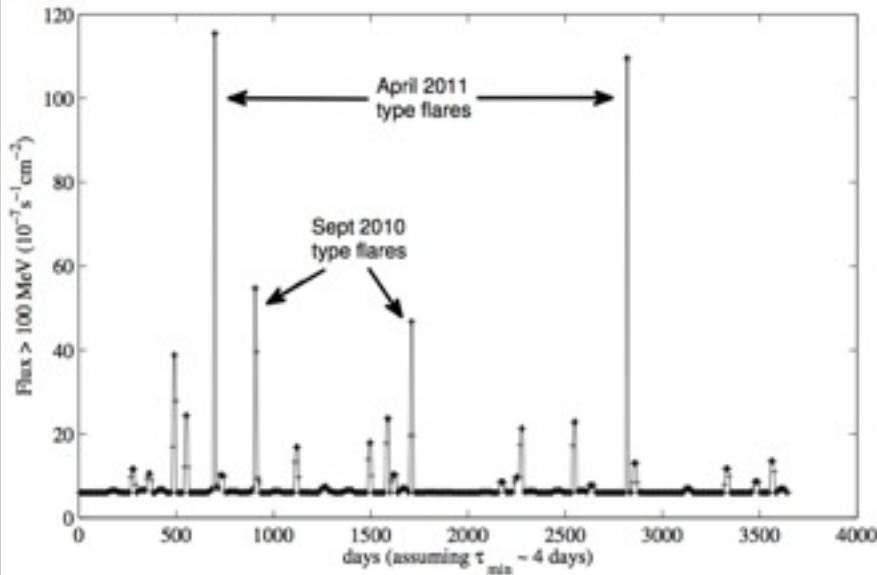
Flare statistics

- Model A:**

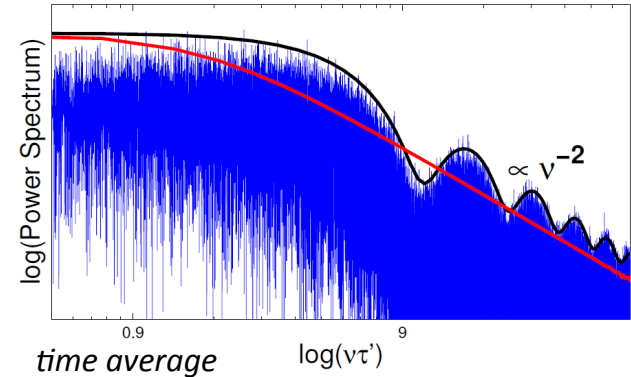
- shock gives average emission
- rare reconnection events
- probability of flare flux

$$\rho(F)dF \propto F^{-\frac{q+1}{q}} dF \approx \frac{1}{F} dF$$

- average flare flux is dominated by bright rare flares.

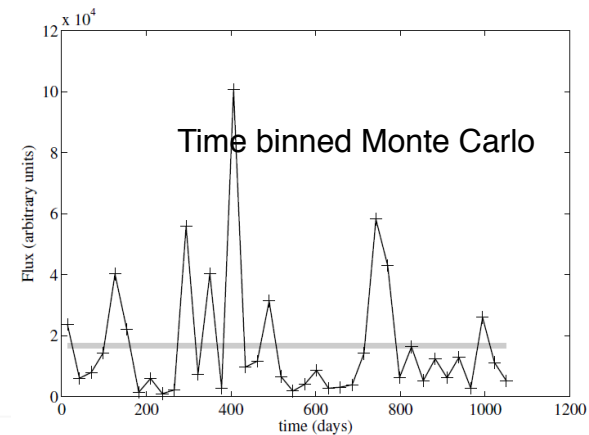


- Model B: overlapping flares**



$$\langle F(t) \rangle \approx n\tau' (2\Gamma)^{2+\alpha} F'$$

$$\text{std. dev.} \longrightarrow \frac{\sigma}{\langle F(t) \rangle} \approx \gamma_1 \approx \frac{\Gamma^{3/2}}{(n\tau')^{1/2}}$$



Conclusion

- **Reconnection is an important, perhaps dominant, mechanism of particle acceleration in PWNs.**
- Observed SED implies hard electron spectrum as predicted by reconnection models.
- Can reproduce intermittency