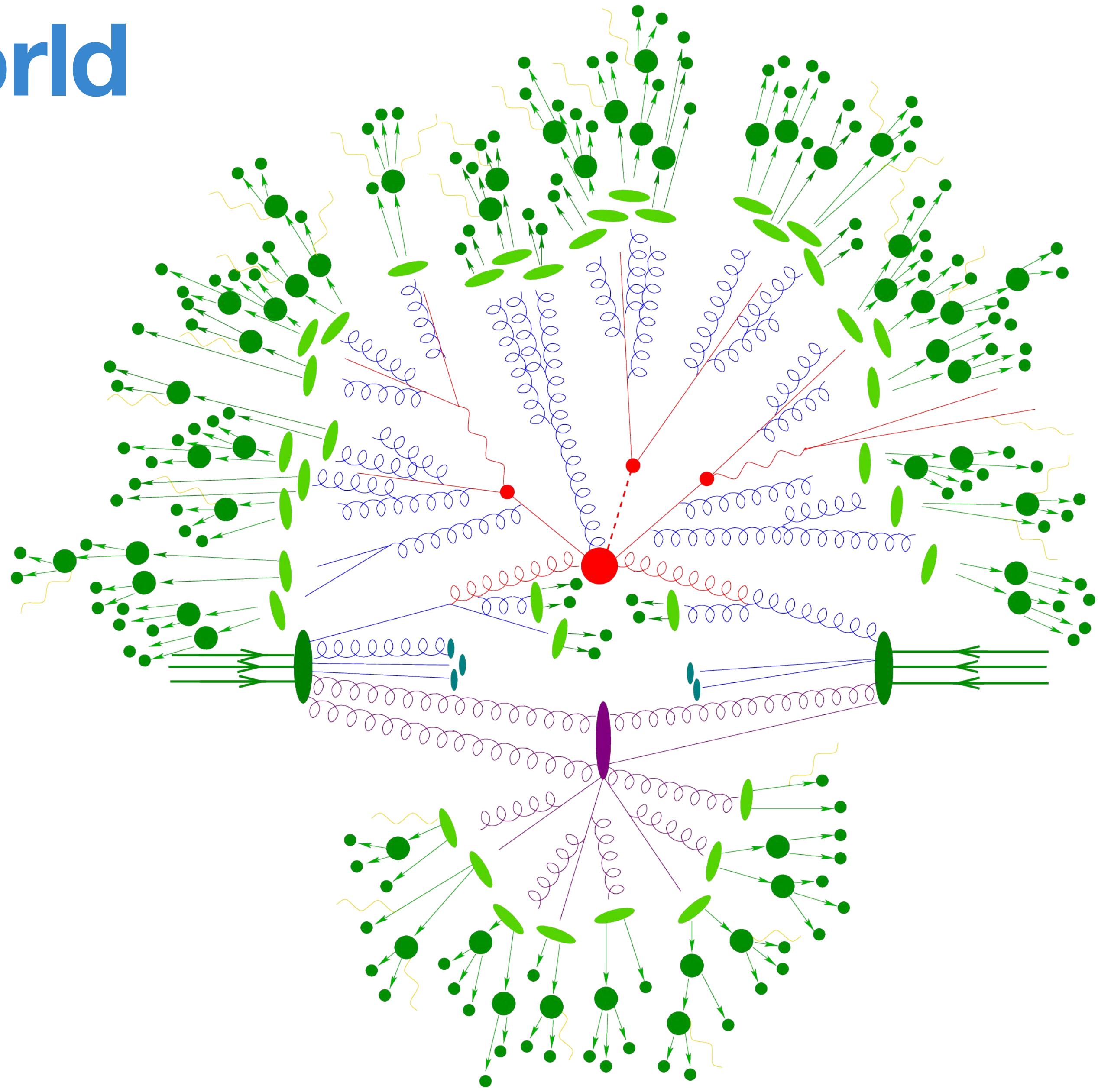
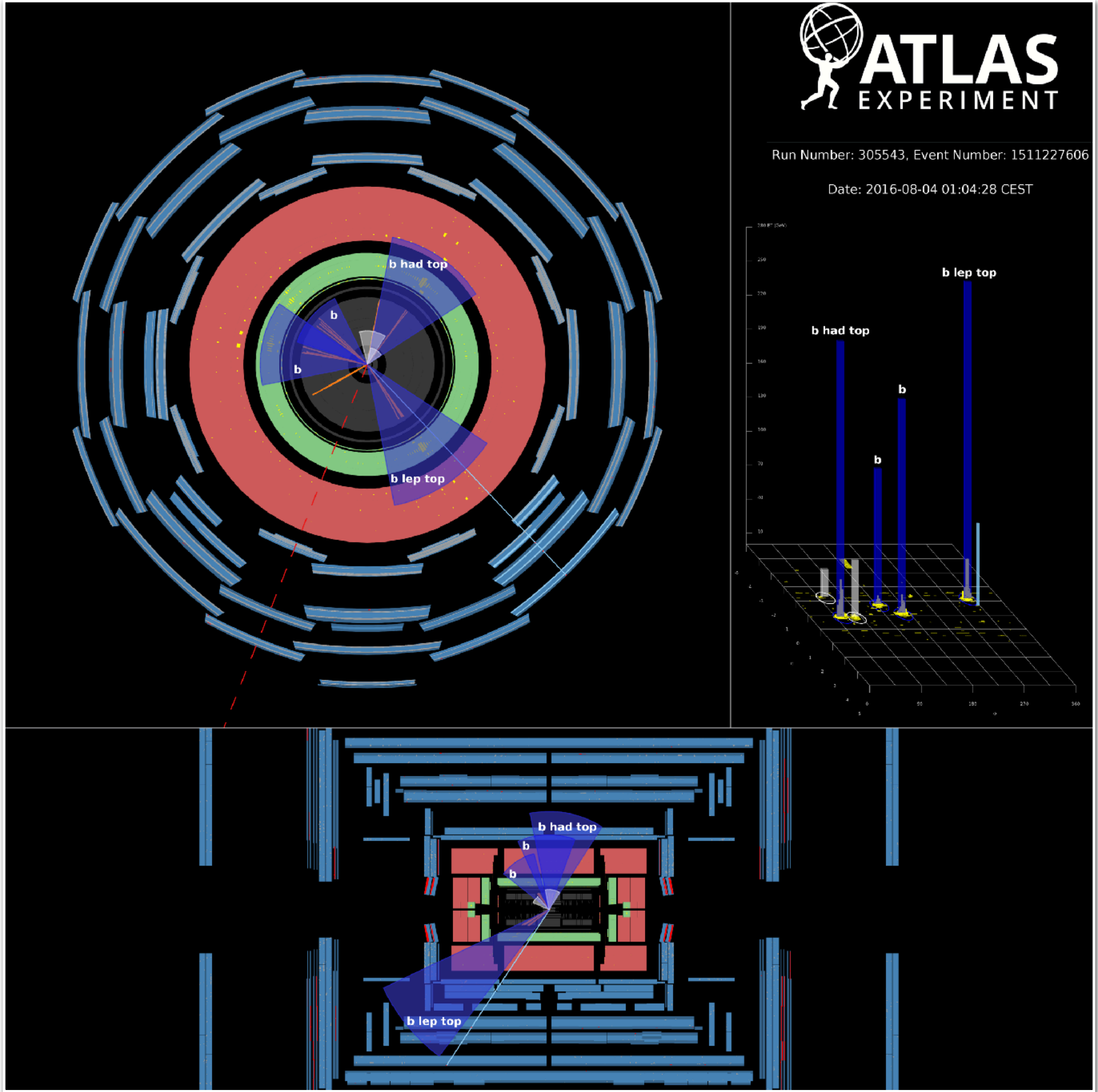


Parton shower accuracy (NLL and beyond)

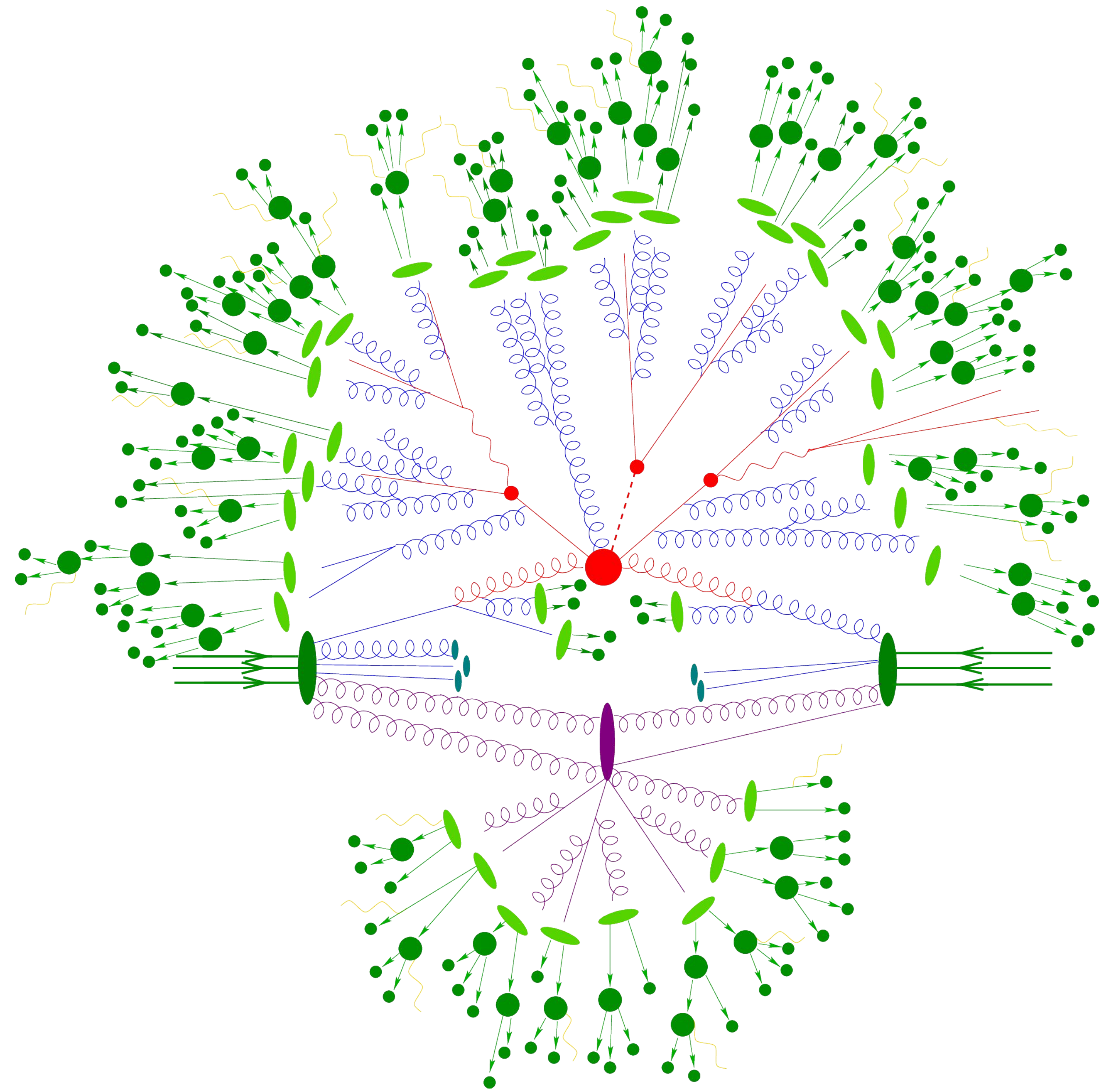
Daniel Reichelt, 2 October 2024

Colliders in the real world



Colliders for theorists

- Event simulation factorised into
 - **Hard Process**
 - **Parton Shower**
 - **PDF/Underlying event**
 - **Hadronisation**
 - **QED radiation**
 - **Hadron Decays**



Colliders for theorists

- Event simulation factorised into

- Hard Process

- Parton Shower

- Underlying event

- Hadronisation

- QED radiation

- Hadron Decays

This Talk:

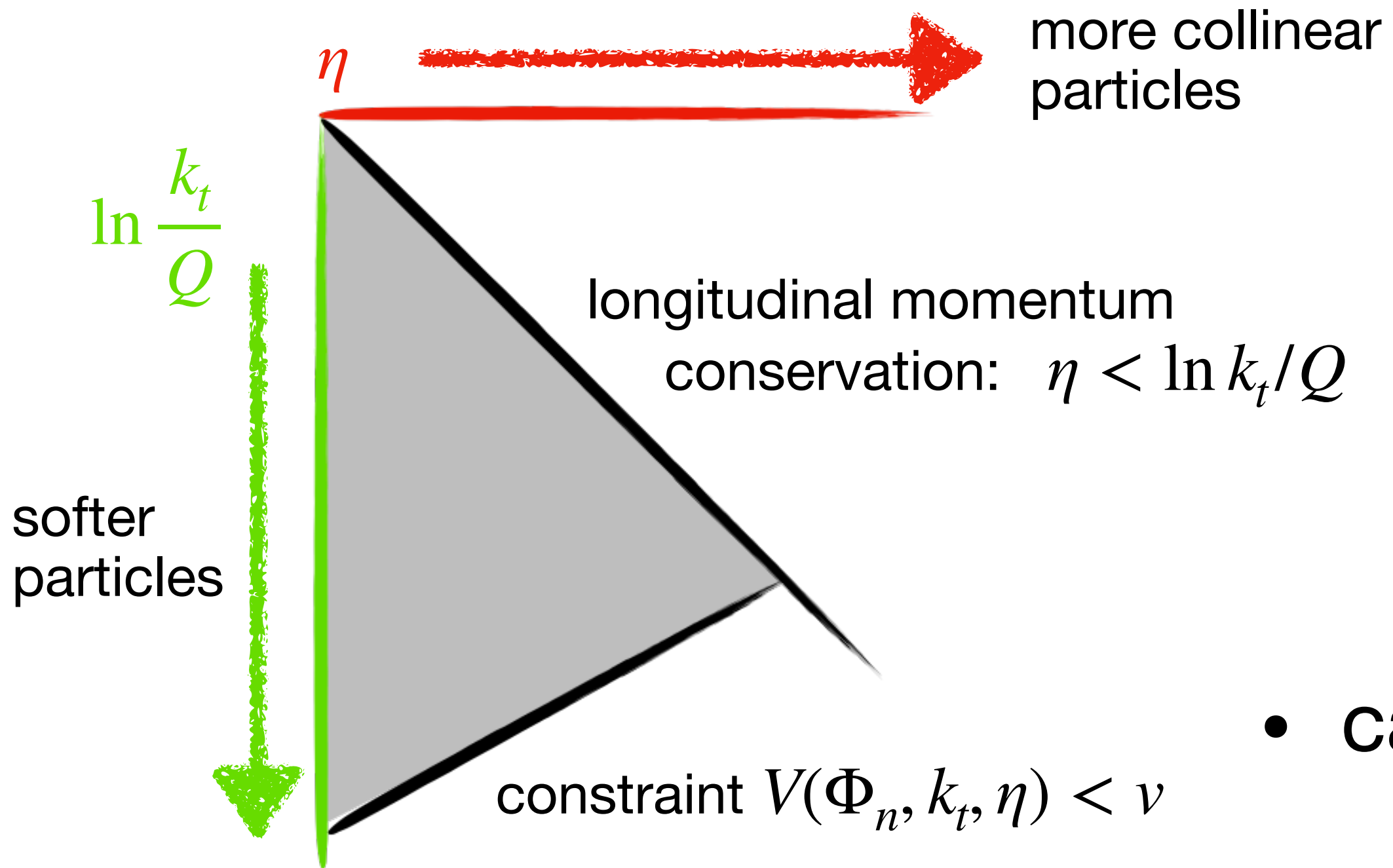
- Well established tools, used for decades to model collider physics
- Also connection to fixed order via matching/merging techniques well established, at least up to NLO
- But: basic shower picture based on leading-log approximation, some simple adjustments to get “at least most of” NLL

Factorisation and Logarithms

- QCD matrix elements factorise in the soft limit as ('Eikonal')

$$d\sigma_{n+1} = d\sigma_n \otimes d\Phi_{+1} \frac{\alpha_s}{2\pi} \sum_{k,i} \mathbb{T}_k \mathbb{T}_i \frac{p_k p_i}{(p_k q)(p_i q)}$$

$$\frac{p_k p_i}{(p_k q)(p_i q)} \sim 1/k_t^2$$



single emission phase space

transverse momentum

rapidity

azimuthal angle

$$d\Phi_{+1} \sim dk_t^2 d\eta d\phi$$

- calculate cross section, cut on $V(k_t, \eta) = k_t/Q > \nu$

$$\rightarrow \frac{\alpha_s}{2\pi} \int_{\nu Q}^Q \frac{dk_t}{k_t} \int_0^{\ln k_t/Q} d\eta \sim \frac{\alpha_s}{2\pi} \ln^2 1/\nu = \frac{\alpha_s}{2\pi} L^2$$

- “probability” for soft gluon emission above ν

Parton branchings

- In toy case of constant probability for one emission between two scales

$$P = \int_{t_c}^{t_0} dt' \lambda = \lambda \Delta t$$

- “No emission” probability given by unitarity
 - $\Delta(t_0, t_c) = \exp[-\lambda \Delta t]$
- Poisson-type distribution familiar from radioactive decay
- In reality not constant (see last slide), but Monte-Carlo methods available to generate emissions to corresponding “no-emission” factor (Veto-algorithm)

Missing ingredients for real (NLL) showers

- Precise choice of scale “ordering” variable $t \rightarrow$ I will mostly talk about $t \sim k_t^2$ ordered showers
- More accurate shower kernels
 - match to collinear part of Altarelli-Parisi splitting kernels
 - include CMW scheme (maybe not the Pythia default, but no conceptual question)
 - including additional effects on color, spin, generic higher order splitting kernels
- prescription to construct $n + 1$ parton final state (aka recoil scheme)

Parton showers - Cliff notes version

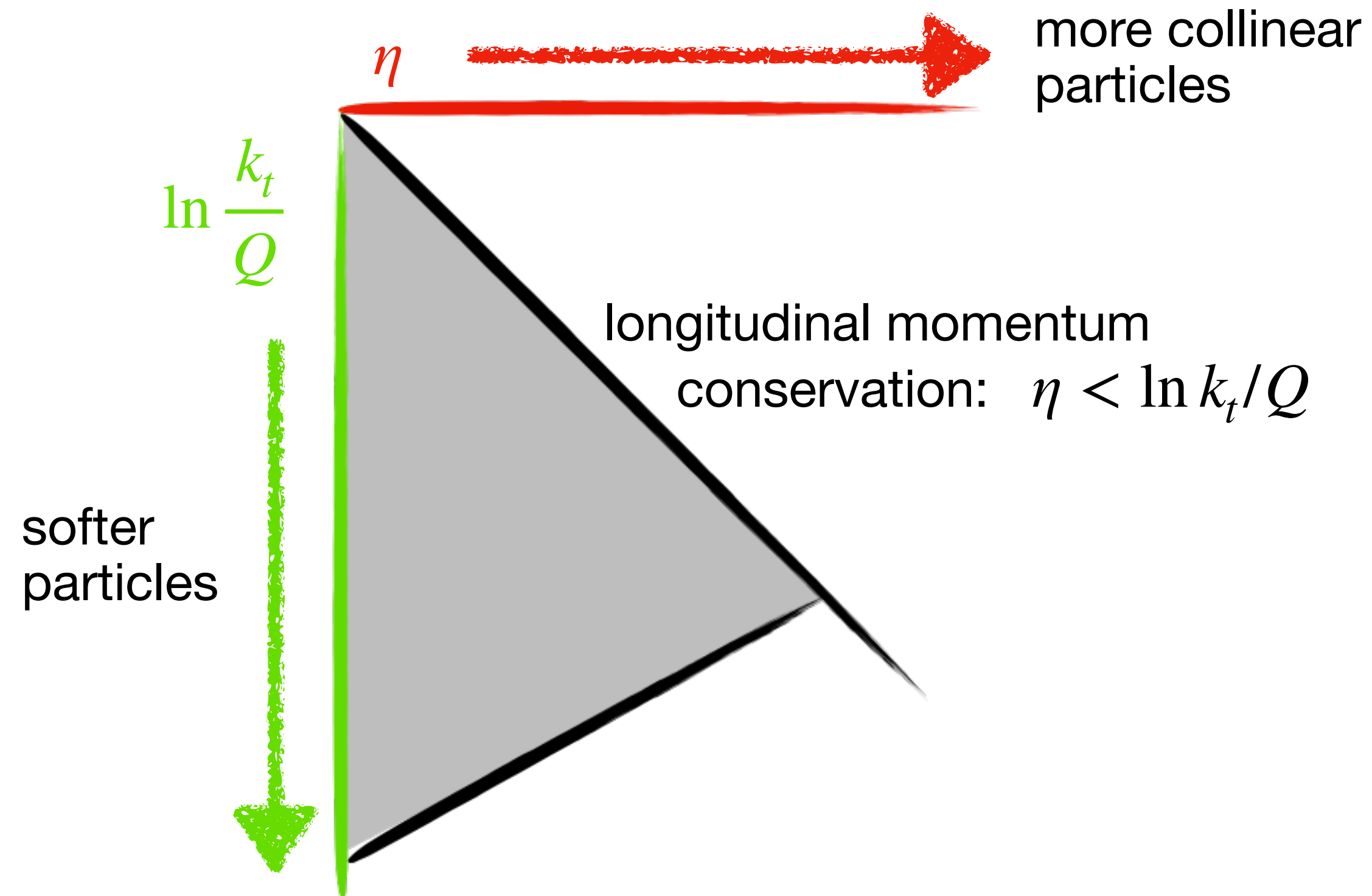
- no-emission probability (sudakov factor)
- Main ingredients to a shower:

$$\sim \exp \left[- \int_{t_0}^{t_1} \frac{dk_t}{k_t} dz \frac{\alpha_S}{2\pi} P(z) \right]$$

1. splitting kernels $P(z)$ captures soft and collinear limits of matrix elements

2. fill phase space ordered in evolution variable $(k_t, \theta, q^2, \dots) \Rightarrow$ here k_t ordered shower

3. generate new final state after emission according to recoil scheme



Colliders for theorists

- Event simulation factorised into

- Hard Process

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- See also large amount of effort dedicated to colour accuracy specifically, e.g. [Forshaw, Holguin, Plätzer '19,'20,'21], [De Angelis, Forshaw, Plätzer '20], [Nagy, Soper '19]

What I will not (so much) talk about:

- Issues with colour assignment:

- inherited from mismatch between PS evolution and resummed observable (different identification of “hardest” emission) [Dasgupta,Dreyer,Hamilton,Monni,Salam '18], [Hamilton, Medves, Salam, Scyboz, Soyez '20]
- for rest of the talk: assume suppression of effect with N_c is sufficient (whether you agree or not, we only have 30 min)

Colliders for theorists

- Event simulation factorised into

- Hard Process

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- Hadron Decays

What I will not (so much) talk about:

- Spin correlations:

- effective solution known in principle ([Collins '88], [Knowles '88,'90]), with application to angular-ordered and dipole showers [Richardson, Webster '18]

- see PanScales studies on implications for resummation properties for specific observables [Karlberg, Salam, Scyboz, Verheyen '21], [Hamilton, Karlberg, Salam, Scyboz, Verheyen '21]

Colliders for theorists

- Event simulation factorised into

- Hard Process

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What I will not (so much) talk about:

- Fixed-order inputs:

- See Emanuelle Re's talk yesterday about NNLOPS methods

- interplay with log accuracy issues in some points, in particular if NLO emission is performed separately a la Powheg-Box [Hamilton, Karlberg, Salam, Scyboz, Verheyen '21]

Colliders for theorists

- Event simulation factorised into

- Hard Process

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What I will (try to) talk about:

- NLL accurate parton showers

- PanScales studies on recoil schemes and solutions

- Pheno with NLL parton showers

- Towards NNLL

Treatment of multiple emissions e.g. in CAESAR

- factorisation of matrix elements in soft collinear limit well known (see last slide)
- how to extract NLL observable independent (i.e. without additional information)?
- method from [Banfi, Salam, Zanderighi '05]: need explicit implementation of soft-collinear limit*:

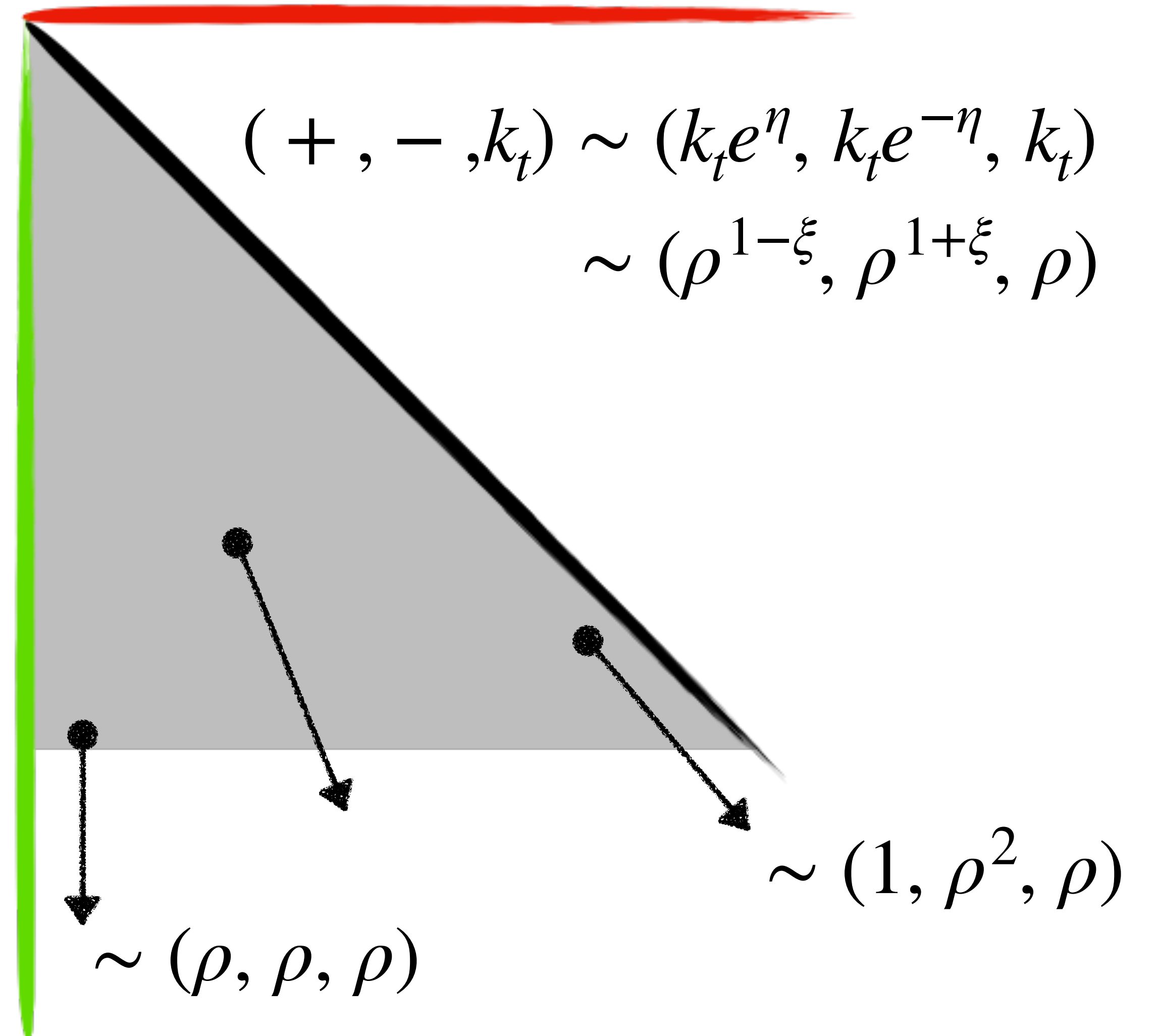
$$k_t^\rho = k_t \rho \quad \xi = \frac{\eta}{\eta_{\max}}$$

$$\eta^\rho = \eta - \xi \ln \rho$$

and assume

$$V(k_i^\rho) = \rho V(k_i)$$

→ numerically evaluate phase space integrals in this limit



* example assuming $V(k_t, \eta) \sim k_t/Q$ for brevity 13

Effect of recoil on accuracy

- question: do recoil effects indeed vanish in soft limit (i.e. $\rho \rightarrow 0$)?*

[Dasgupta,Dreyer,Hamilton,Monni,Salam '18]

- consider situation where we first emit \tilde{p}_{ij} from p_a, p_b , then emit p_j ,

$$\tilde{p}_{ij} \rightarrow p_i, p_j$$

- transverse momentum of p_i will be

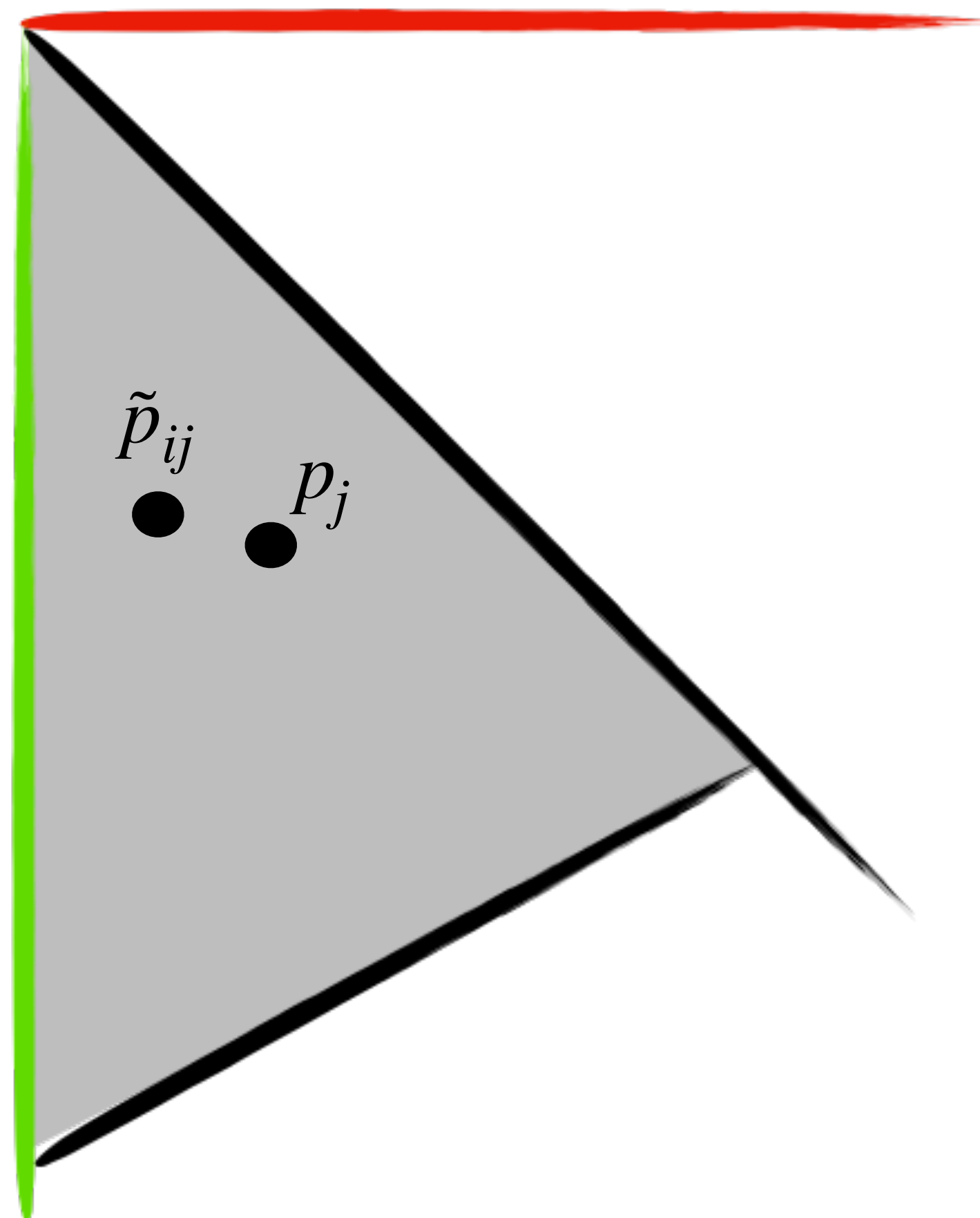
$$k_t^i \sim k_t^{ij} + k_t^j \rightarrow k_t^{ij} \text{ as } \frac{k_t^j}{k_t^i} \rightarrow 0$$

- but, relevant limit is $\frac{\Delta k_t^i}{k_t^i} \rightarrow \frac{\rho k_t^j}{\rho k_t^i} = \mathcal{O}(1)$

$$p_i = z\tilde{p}_{ij} + (1-z)y\tilde{p}_k + k_\perp$$

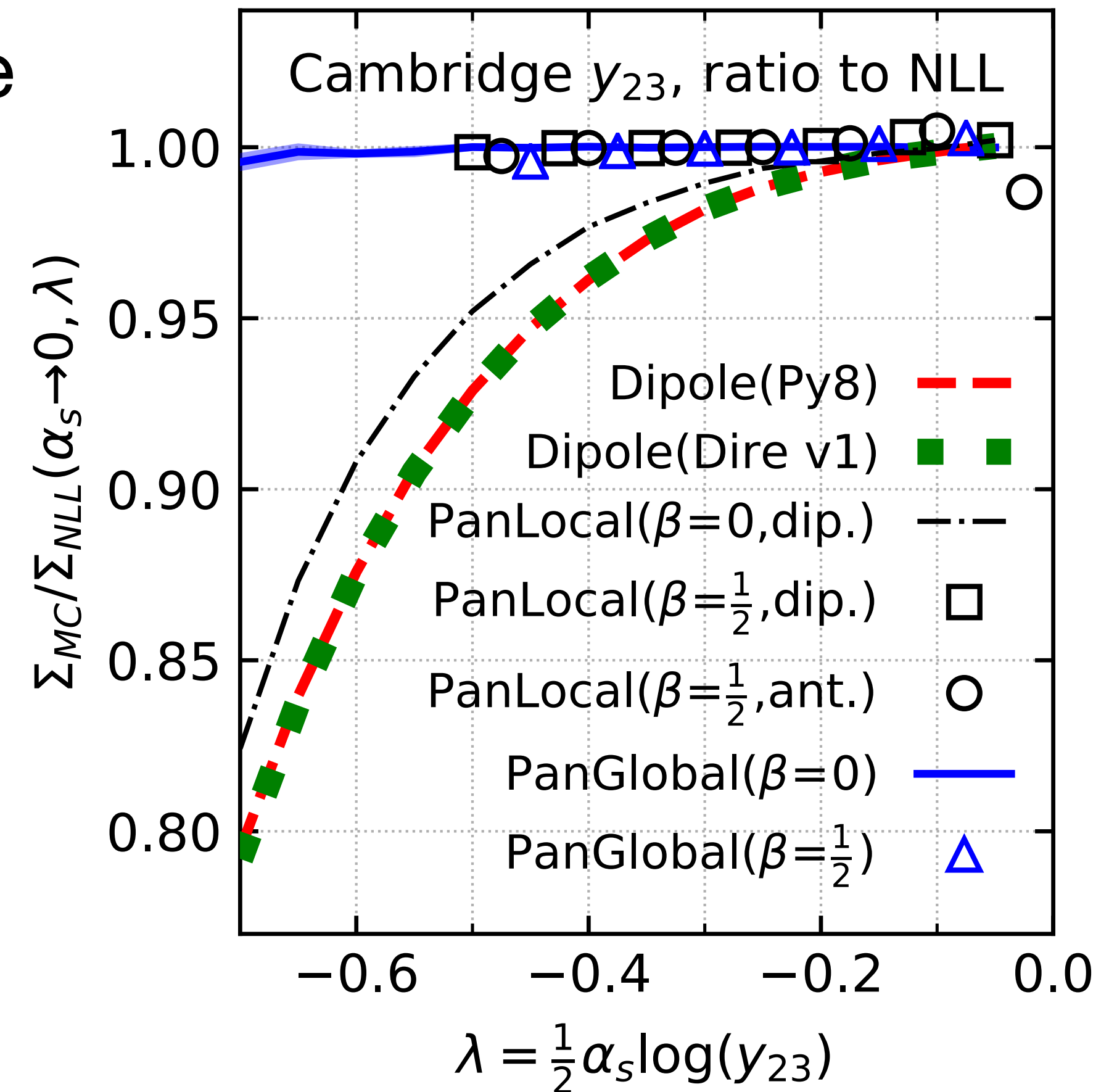
$$p_j = (1-z)\tilde{p}_{ij} + zy\tilde{p}_k - k_\perp$$

$$p_k = (1-y)\tilde{p}_k .$$



New Parton Showers - NLL accuracy

- typical claim based on accuracy of splitting functions etc.
- parton showers \sim NLL accurate if CMW scheme for strong coupling is used
- observation in [Dasgupta, Dreyer, Hamilton, Monni, Salam '18] [Dasgupta, Dreyer, Hamilton, Monni, Salam '20] (PanScales collaboration):
 - subtleties arise in distribution of recoil for subsequent emissions \Rightarrow phase space where accuracy is spoiled if soft gluon absorbs recoil
 - apparently restricts k_t ordered showers to global recoil schemes

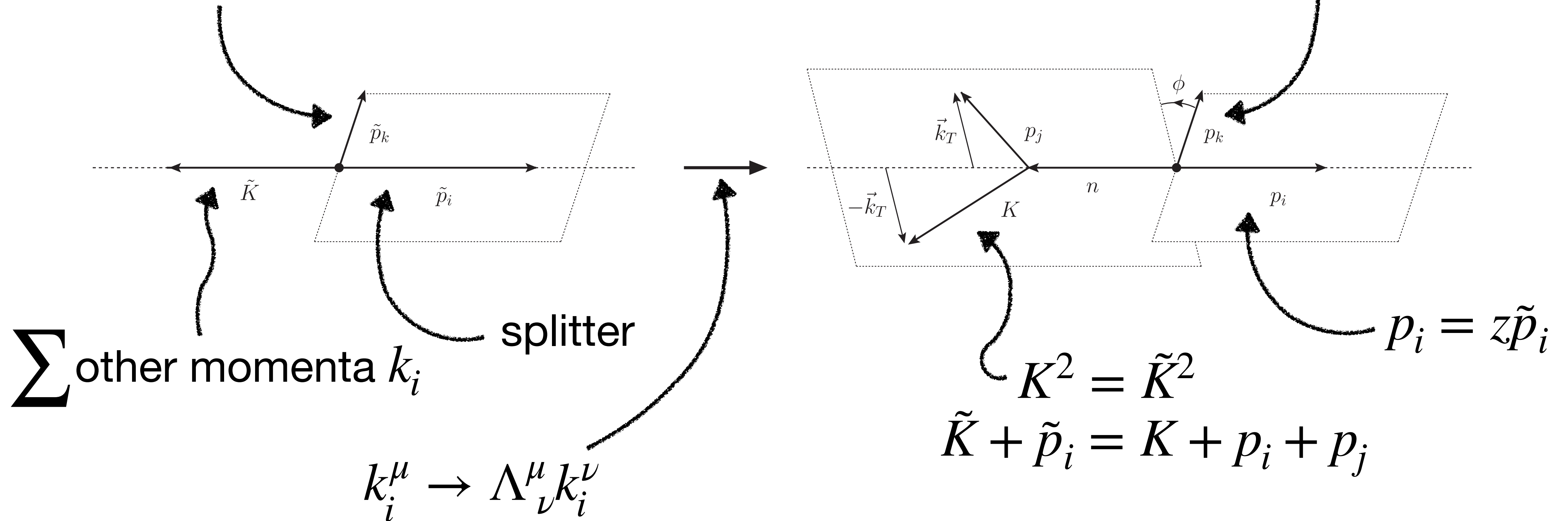


Kinematics - global recoil scheme (Alaric example)

[Herren, Höche, Krauss, DR, Schönherr, '22]

- Before splitting:
colour spectator

- After splitting:



[Catani, Seymour '97]

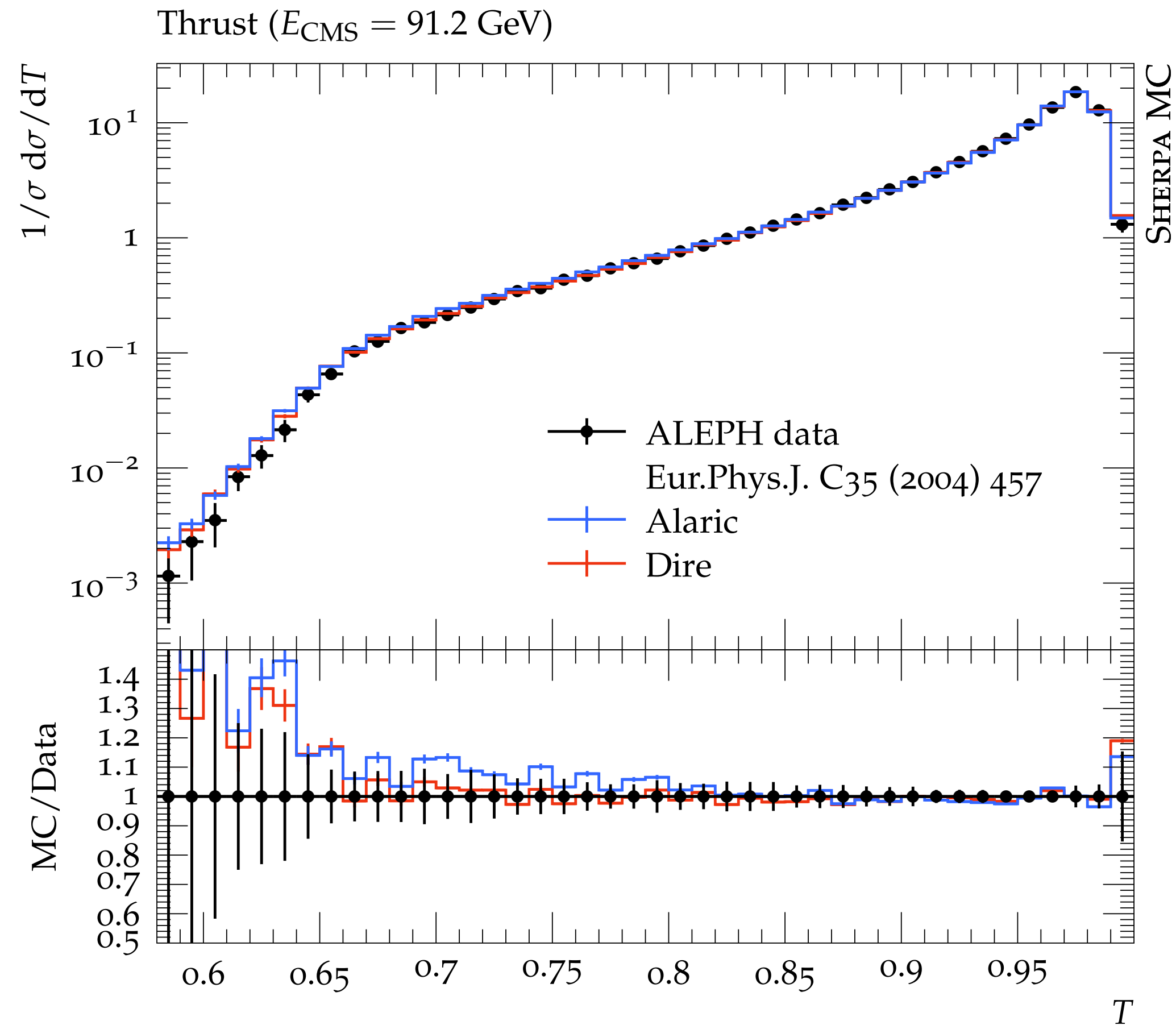
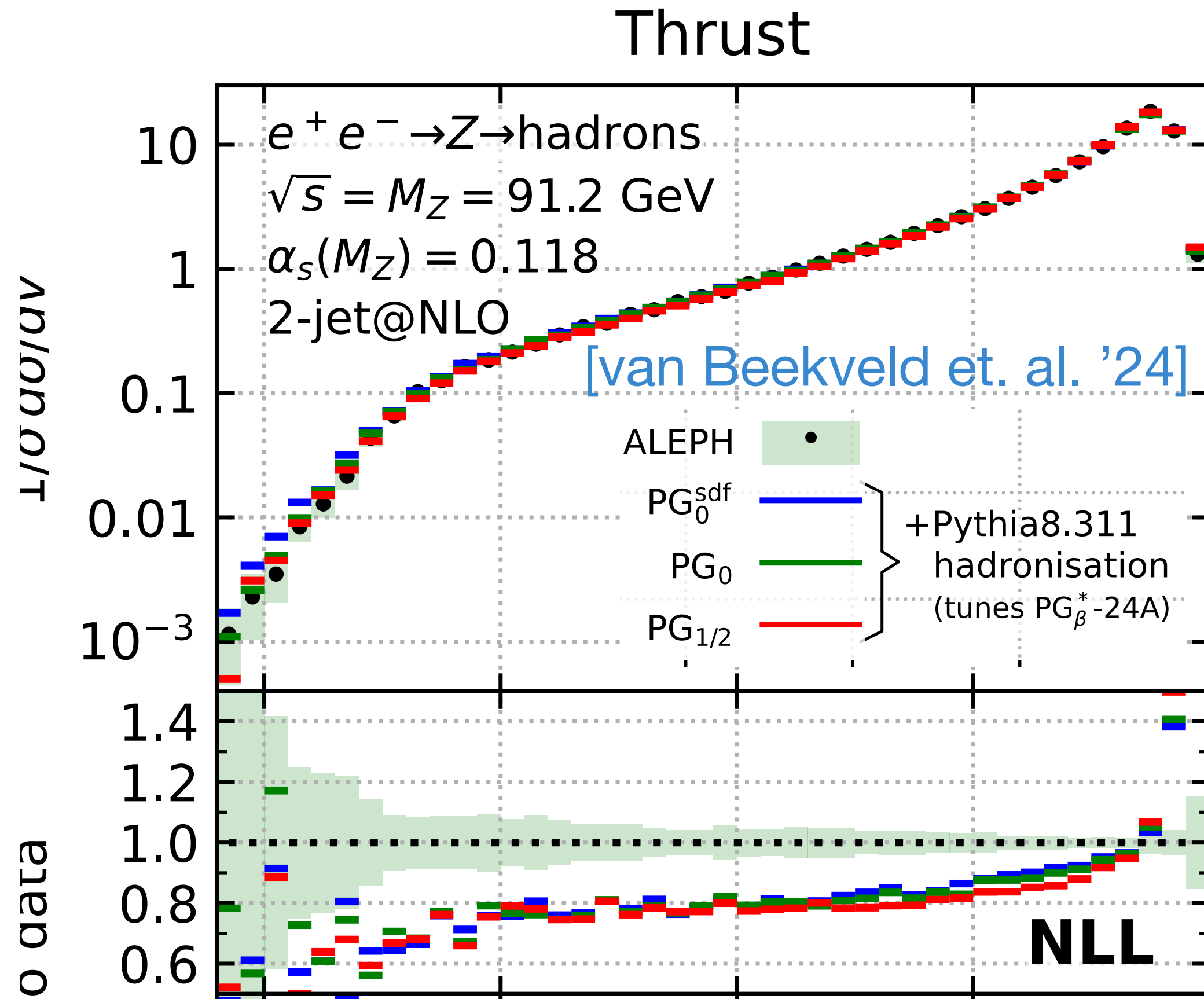
$$\Lambda_{\nu}^{\mu} = g_{\nu}^{\mu} - \frac{(K + \tilde{K})^{\mu}(K + \tilde{K})_{\nu}}{K \cdot \tilde{K} + \tilde{K}^2} + 2 \frac{K^{\mu} \tilde{K}_{\nu}}{\tilde{K}^2} \rightarrow \Lambda_{\nu}^{\mu} \tilde{K}^{\nu} = K^{\mu}$$

Lorentz transformation distributes recoil to hardest particles!

New Parton Showers - NLL accuracy

- Several solutions/re-evaluations of parton shower concepts:
- [Dasgupta, Dreyer, Hamilton, Monni, Salam, Soyez '20], [vanBeekveld, Ferrario Ravasio, Hamilton, Salam, Soto-Ontoso, Soyez '22] ...
 - partitioning of splitting functions and appropriate choice of evolution variable can lead to NLL accurate shower for local and global recoil strategies
- [Forshaw, Holguin, Plätzer '20]
 - Connections between angular ordered and dipole showers
- [Nagy, Soper '11]
 - local transverse, global longitudinal recoil
- [Herren, Höche, Krauss, DR, Schönherr, '22], [Höche, Asse '23], [Höche, Krauss, DR '24]
 - global recoil, enables analytic comparison to resummation and proof of NLL accuracy
- [Preuss '24]
 - global recoil in antenna shower Vinca

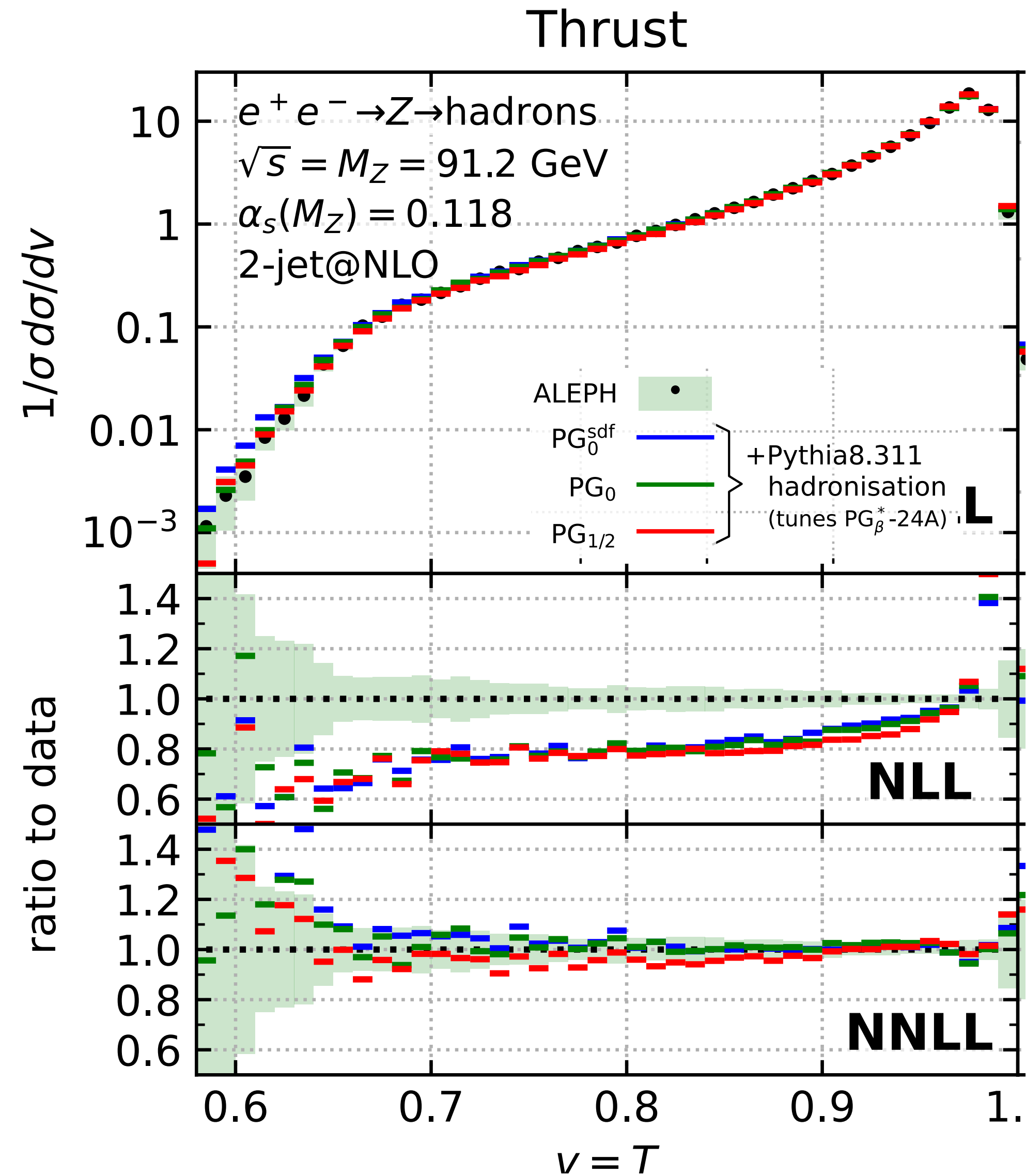
Pheno with NLL showers



- PanScales shower and Alaric @ NLL accuracy

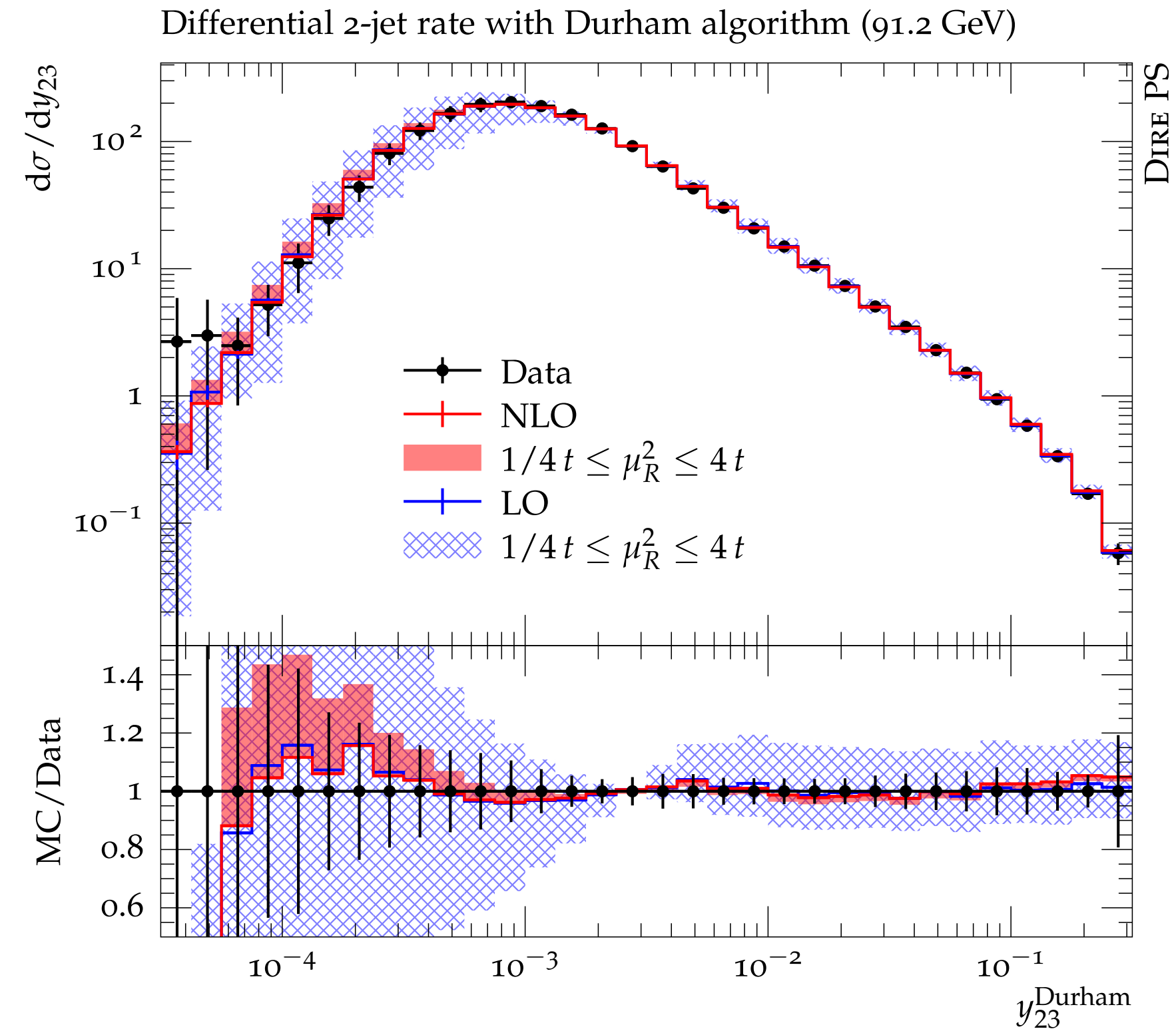
- significantly different conclusion about ability of NLL shower to describe data (similar level of tuning, $\alpha_s = 0.118$ is fixed, string fragmentation parameters in Pythia 8 tuned to LEP data)

Towards NNLL



[van Beekveld et. al. '24]

- Conclusion from PanScales studies: NNLL needed to describe even simple observables
- Achieved by multiplicative matching of NLO splitting kernels via + correction terms capturing effect of inclusive gluon emissions

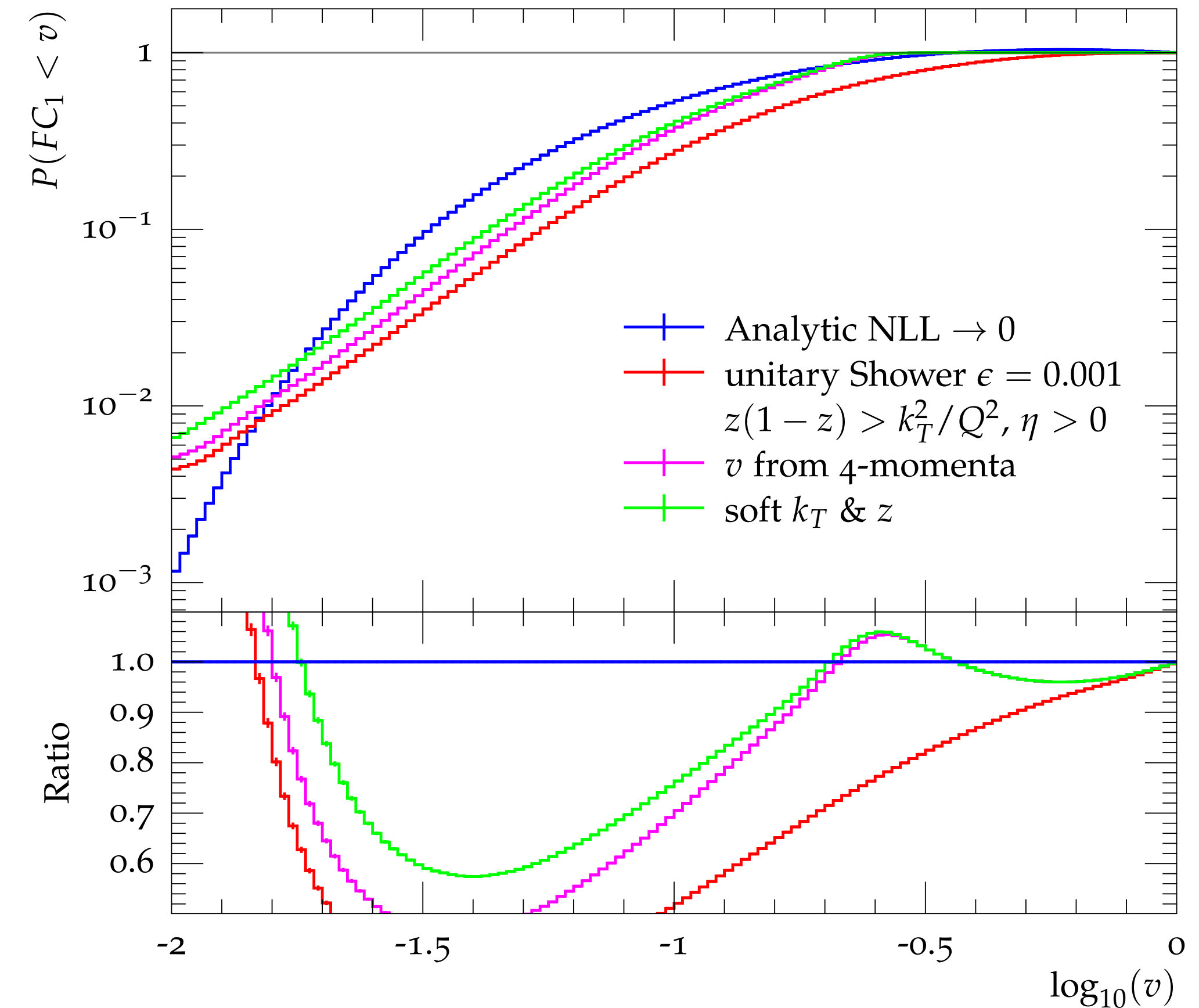


[Höche, Prestel '17]

- Appears to be in contrast with small effects found so far in implementing higher order splitting functions (though not in complete NNLL framework yet) [Höche, Prestel '17], [Dulat, Höche, Prestel '18], [Gellersen, Höche, Prestel]

Beyond logarithmic accuracy

- Observations
 - LL and NLL accurate showers can be very similar (e.g. failing of NLL accuracy numerically undetectable for Dire in prominent observables like Thrust)
 - NLL accurate showers can differ significantly from NLL result away from strict limit
 - \Rightarrow subleading effects play a significant role in phenomenological successful parton showers, more systematic understanding desirable, see also [Höche, Siegert, DR '17]



Alaric beyond NLL - subleading effects

assume Sudakov decompose like

$$p_i^\mu = z_i \hat{p}_{ij}^\mu + \frac{-k_t^2}{z_i 2p_{ij}\bar{n}} \bar{n}^\mu + k_t^\mu ,$$

$$p_j^\mu = z_j \hat{p}_{ij}^\mu + \frac{-k_t^2}{z_j 2p_{ij}\bar{n}} \bar{n}^\mu - k_t^\mu$$

derivation of splitting functions leads to:

$$P_{qq\parallel}^{(F)}(p_i, p_j, \bar{n}) = C_F (1 - \varepsilon)(1 - z_i)$$

$$P_{gg\parallel}^{(F)}(p_i, p_j, \bar{n}) = 2C_A z_i z_j ,$$

$$P_{gq\parallel}^{(F)}(p_i, p_j, \bar{n}) = T_R \left[1 - \frac{2 z_i z_j}{1 - \varepsilon} \right] .$$

actual shower kinematics:

$$p_i = z \tilde{p}_i ,$$

$$p_j = (1 - z) \tilde{p}_i + v(\tilde{K} - (1 - z + 2\kappa) \tilde{p}_i) - k_\perp ,$$

$$K = \tilde{K} - v(\tilde{K} - (1 - z + 2\kappa) \tilde{p}_i) + k_\perp ,$$

$$p_i = \frac{z}{1 - v(1 - z + \kappa)} \hat{p}_{ij} + \frac{z}{1 - v(1 - z + \kappa)} k_\perp + \mathcal{O}\left(\frac{k_\perp^2}{2\tilde{p}_i \tilde{K}}\right) ,$$

$$p_j = \frac{(1 - z)(1 - v) - v\kappa}{1 - v(1 - z + \kappa)} \hat{p}_{ij} - \frac{z}{1 - v(1 - z + \kappa)} k_\perp + \mathcal{O}\left(\frac{k_\perp^2}{2\tilde{p}_i \tilde{K}}\right)$$

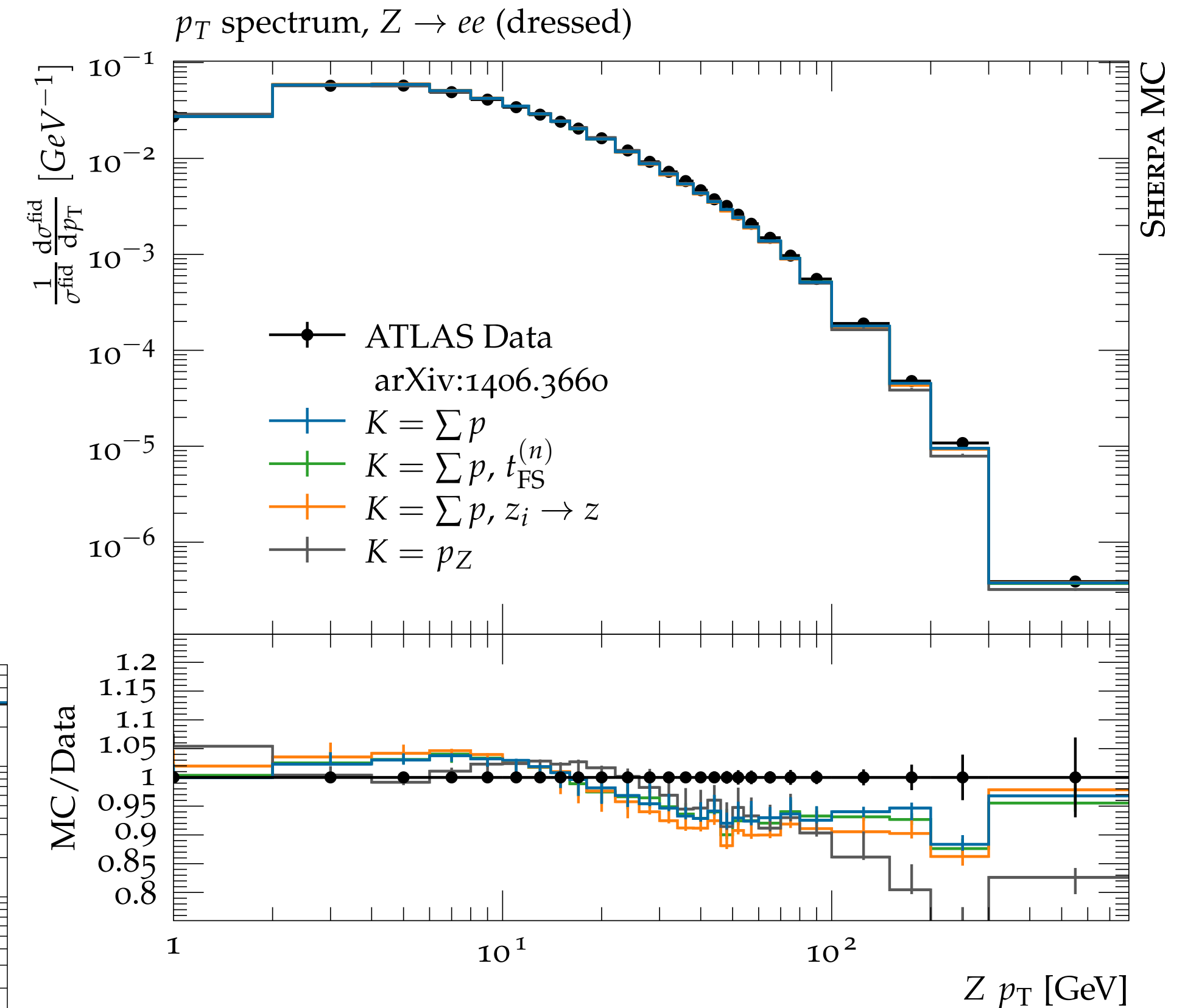
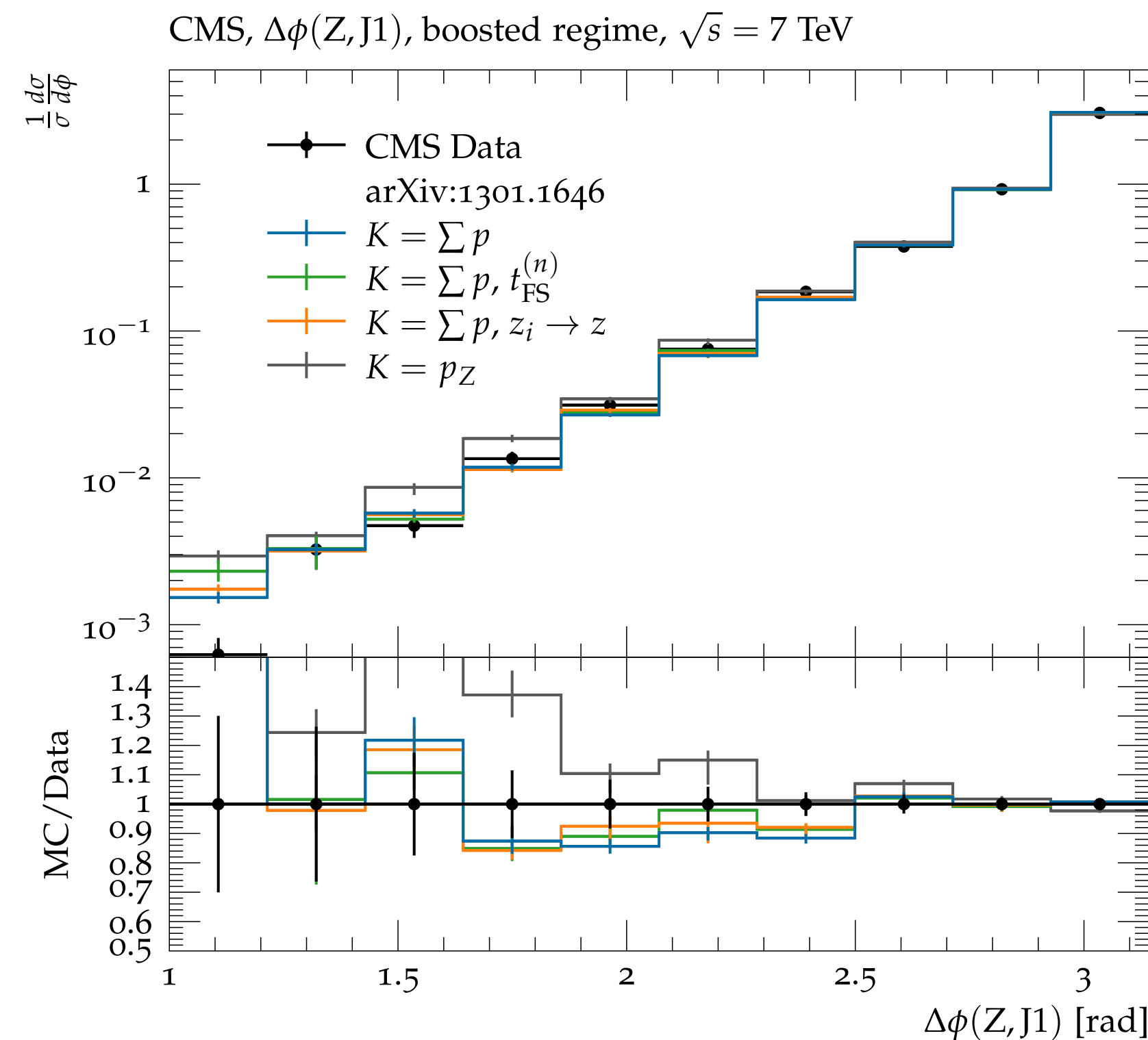
ultimately, “proper”
splitting variables:

$$z_i = \frac{z}{1 - v(1 - z + \kappa)} ,$$

$$z_j = 1 - \frac{z}{1 - v(1 - z + \kappa)}$$

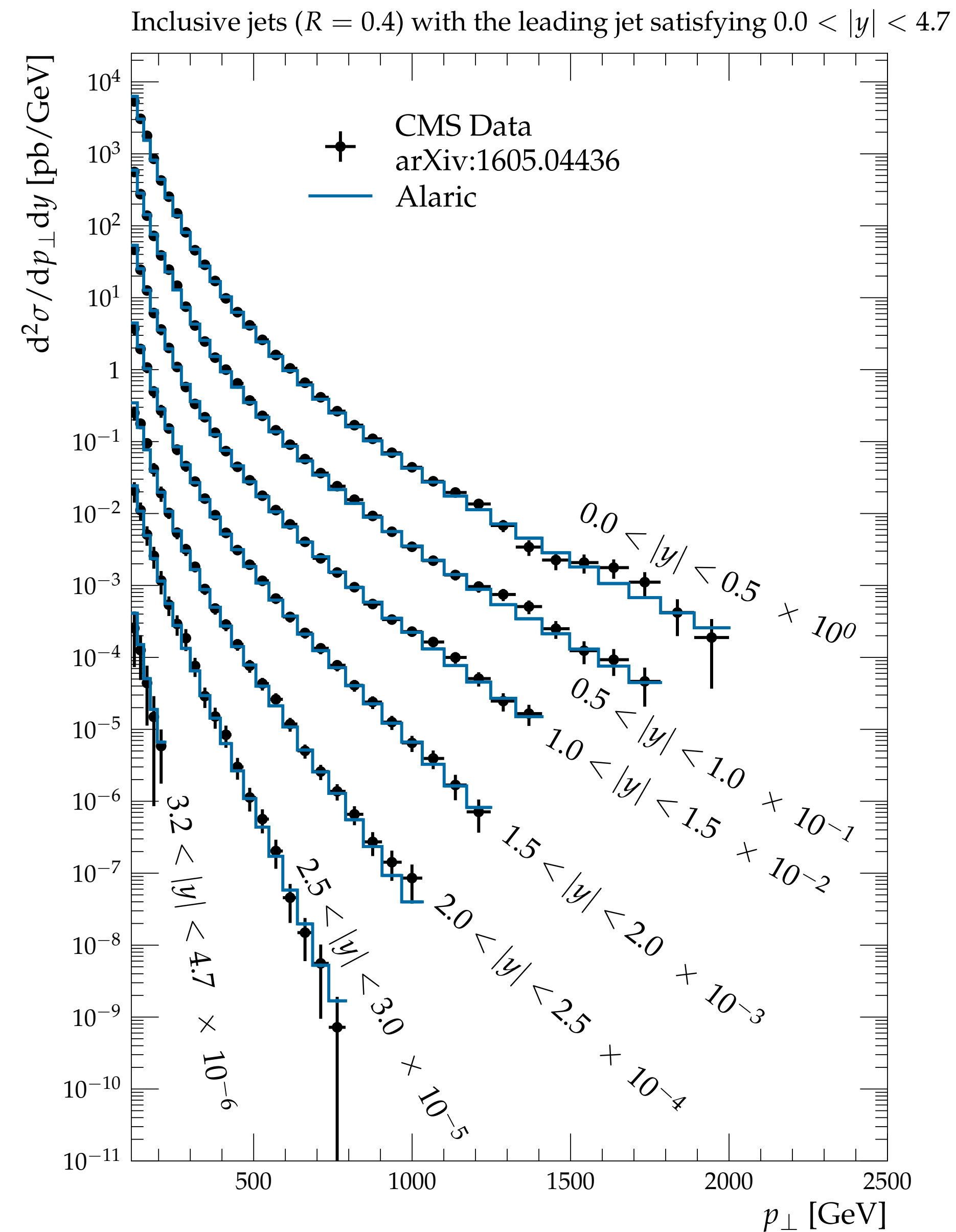
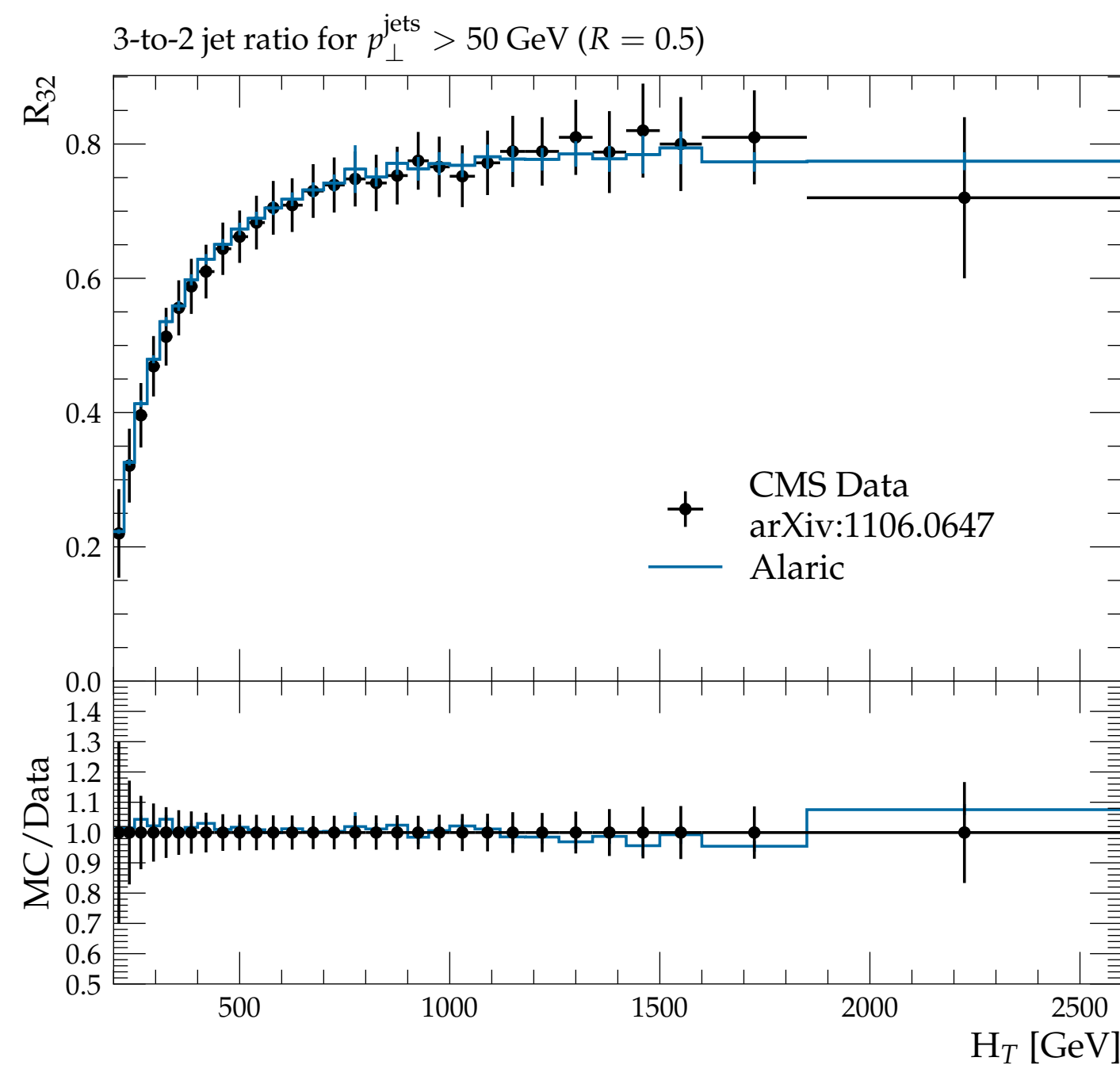
Alaric — subleading effects in Z+jets

- effects/choices beyond NLL accuracy:
 - choice of evolution variable (up to factors of $z \sim 1$)
- identify PS parameter z with z_i, z_j
- choice of recoil momentum K (NLL accuracy needs “hard” K)



Alaric at the LHC — jets

- [Höche, Krauss, DR '24] extend Alaric method to IS evolution
- satisfactory description of inclusive and dijet events
- transverse momentum spectrum of leading jet and ratio 3-to-2 jet rate



Conclusion

- Progress on logarithmic accuracy of parton showers (as compared to resummed calculations)
- Effect on “general-purpose” nature to be seen
 - reminder to Paolo Nason’s talk yesterday,
“ ‘best’ theory framework [has] not always [been] successful in SMC land ”
- Outlook:
 - Probably NNLL PS matched with NNLO fixed order in near future (at least on time scale of future collider)
 - Non-perturbative corrections/soft physics effect might become limiting factors