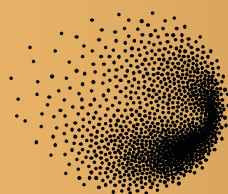


KU LEUVEN

Update on modern muonic x-ray experiments

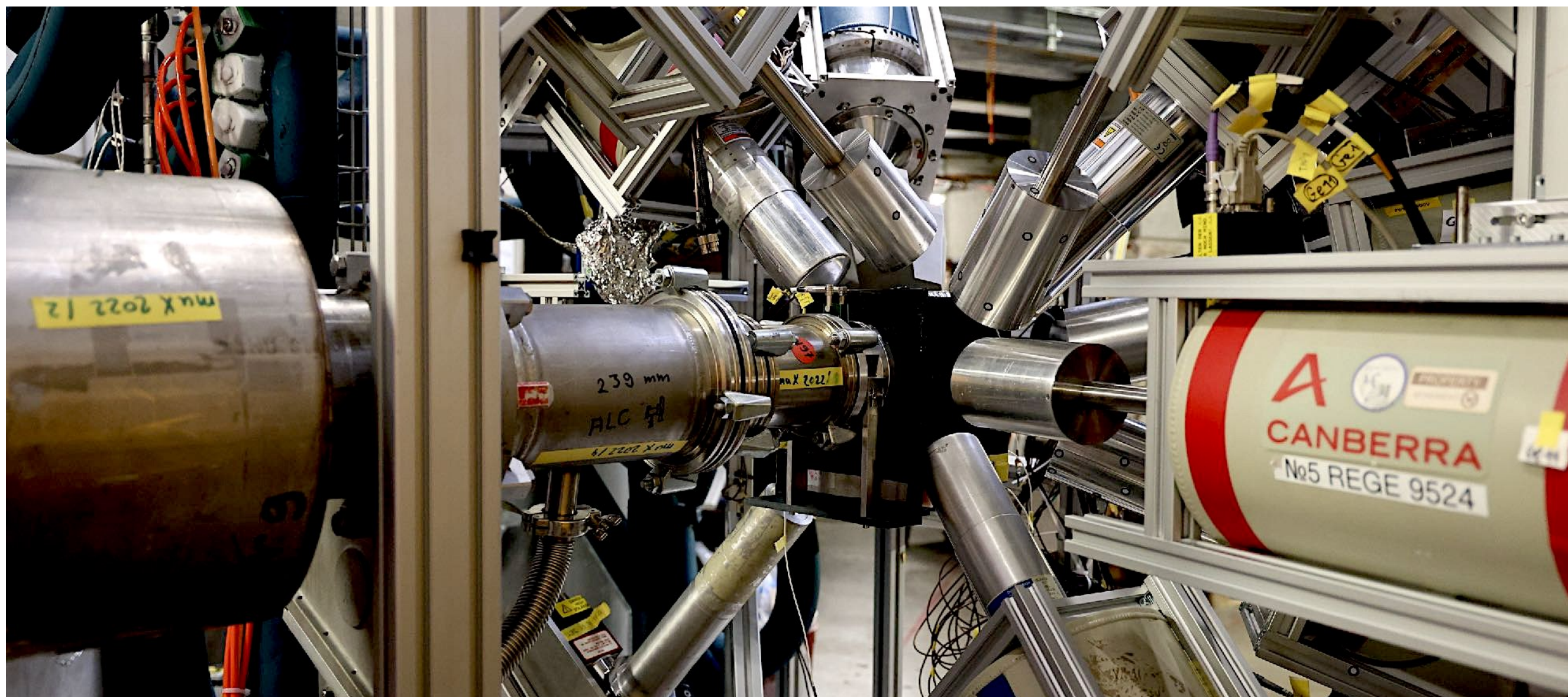


PSI

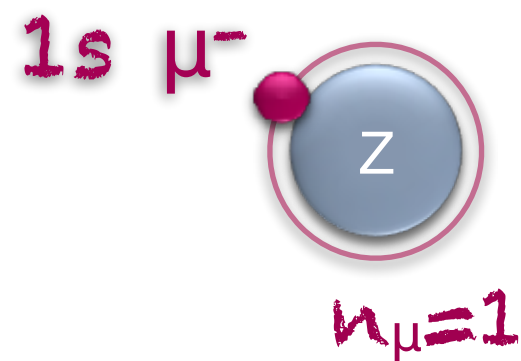
Stella Vogiatzi

On behalf of the RefRadii & muX collaborations

26-30 May, 2025



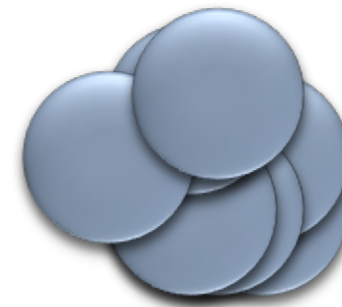
Our *exotic* muonic atom



$n_e = 1$

Our *exotic* muonic atom

Stopping negative muons ****directly**** to
targets of >100 mg



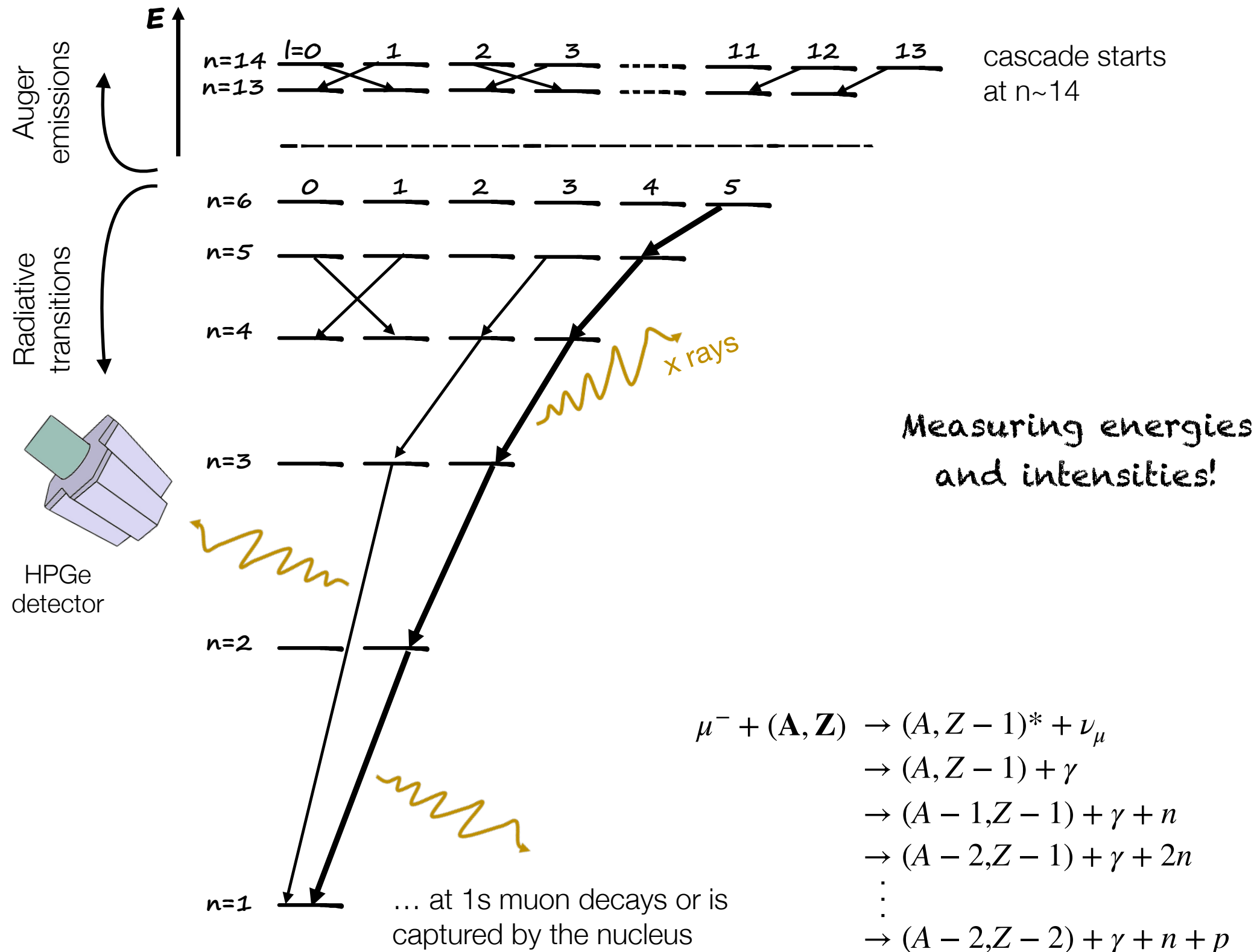
$30 \text{ MeV}/c \mu^-$



$n_e=1$

Achtung! ****...only when you have an abundant and stable target**

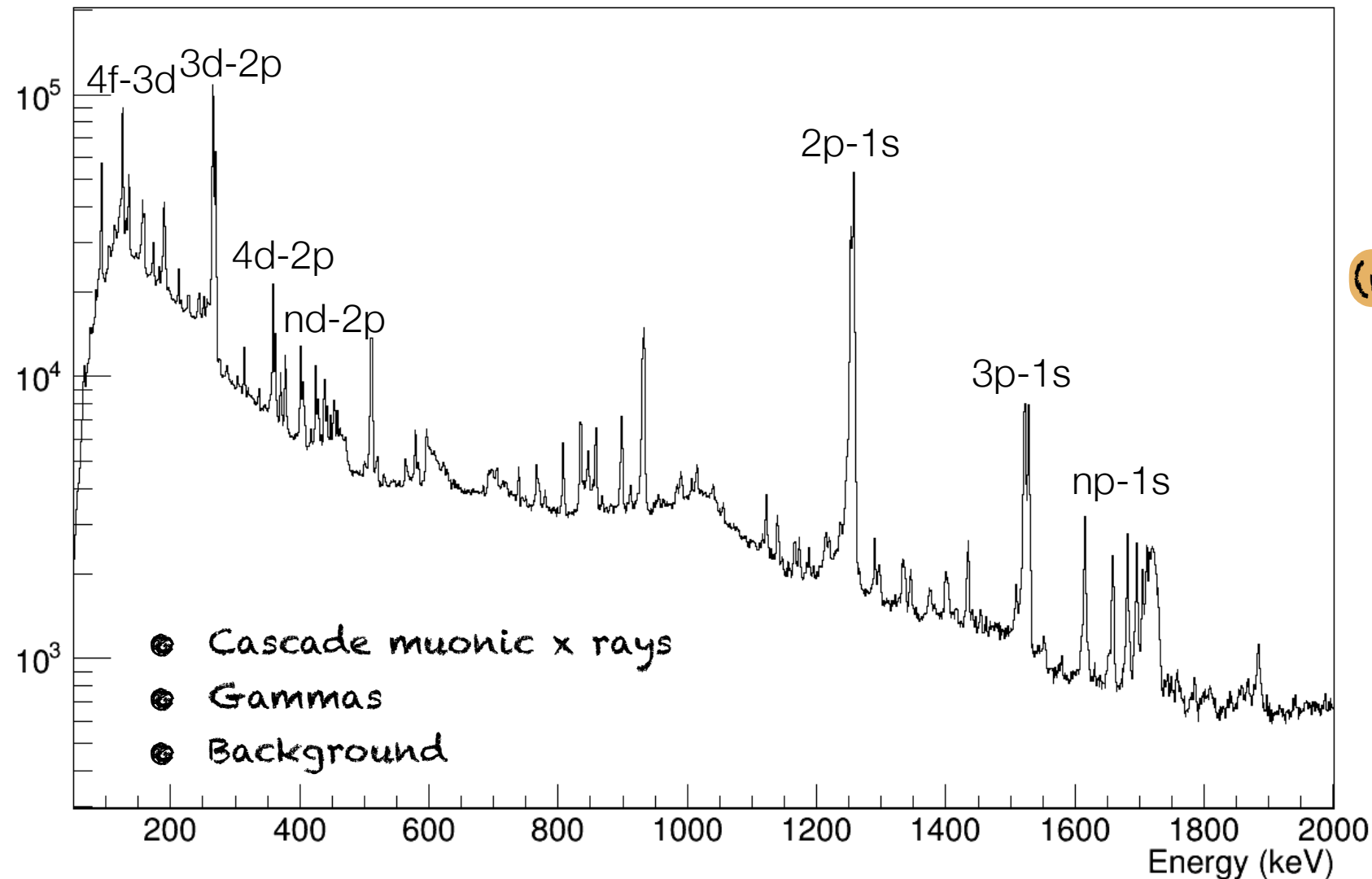
From muon capture to the ground state



Energy spectrum of a muonic atom *As seen by a HPGe detector*

Eur. Phys. J. Plus 135, 777 (2020)

muonic iron



Muonic cascade intensities for
elemental compositions

(MIXE)

Muonic cascade energies for
charge radii extraction

(muX & Reference Radii & QUARTET)

mid- and heavy-
isotopes (up to ^{96}Cm)
with HPGe

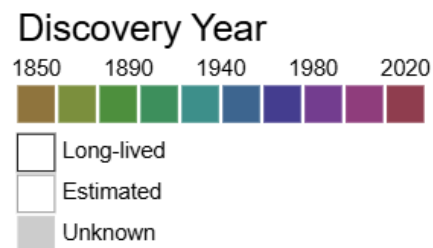
Low-masses (^3Li to
 ^{12}Mg) with MMC

Nuclear capture gammas
for capture rates
(MONUMENT)

Science case

By measuring absolute nuclear charge radii...

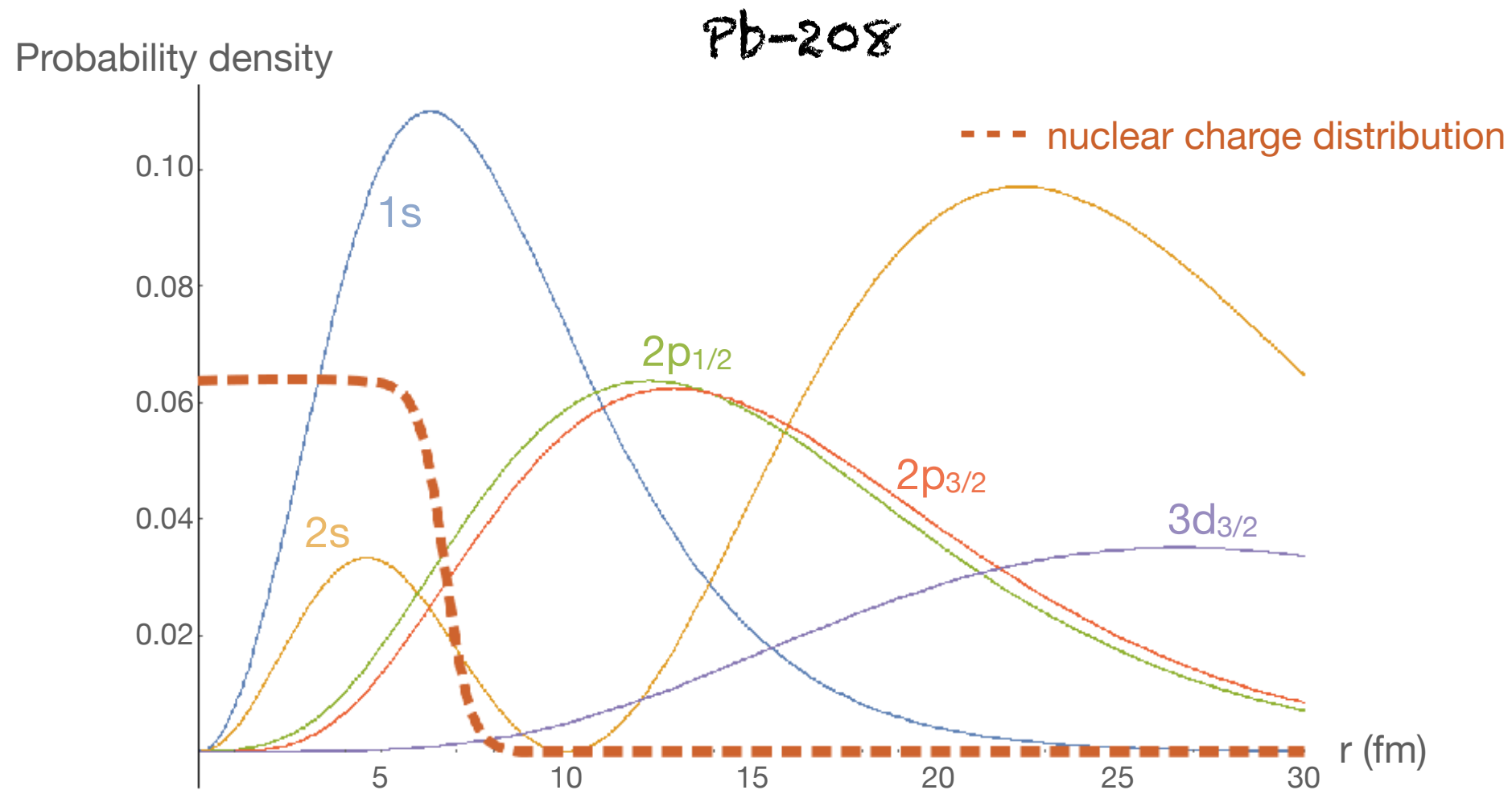
- Benchmark laser spectroscopy results & challenge state-of-art nuclear models
- Provide better understanding of the underlying nuclear structure



- Provide insight on the evolution of nuclear shapes across the nuclear landscape

- Probe Atomic Parity Violating effects & New Physics

Why not using electrons to probe charge radii?



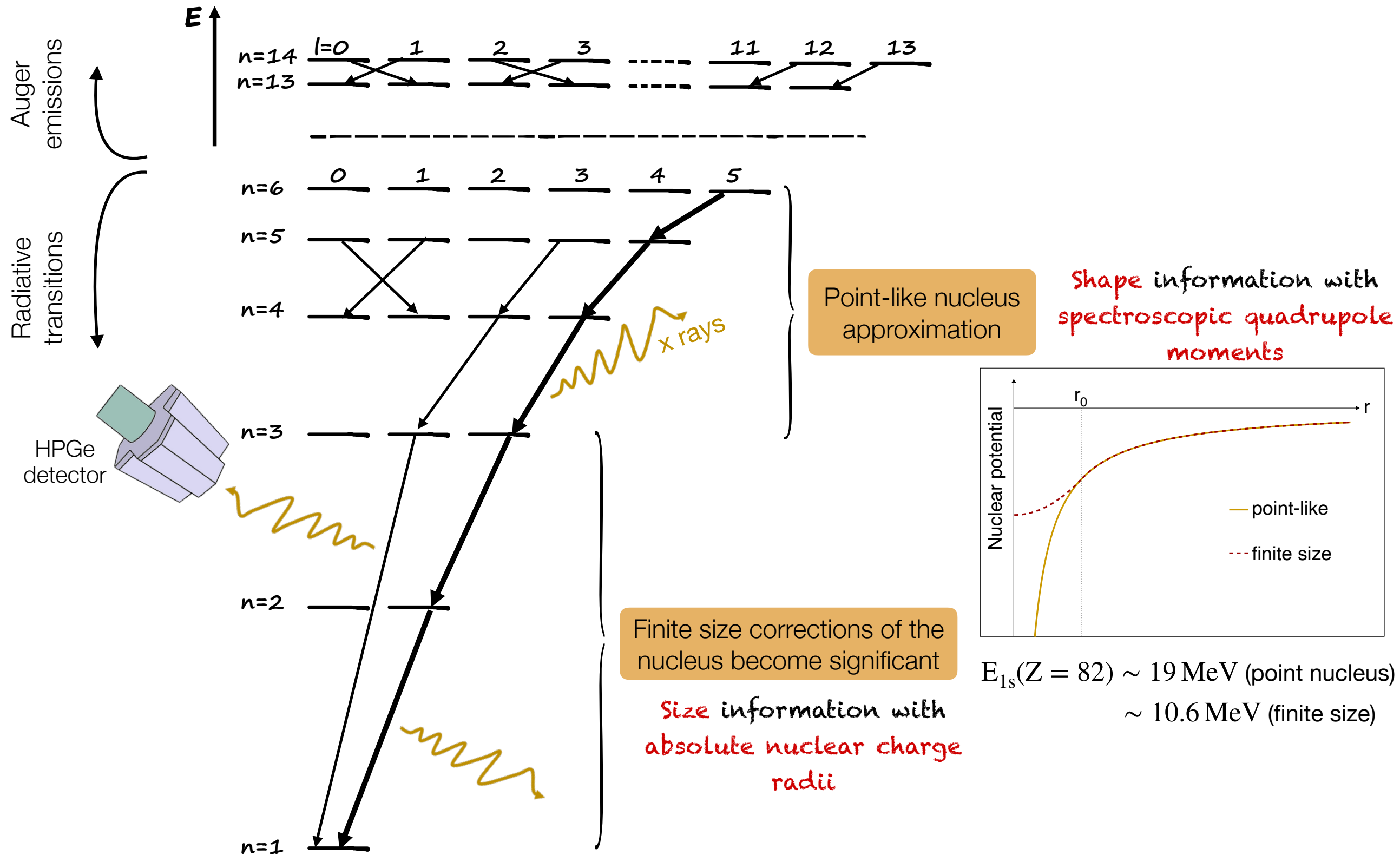
- Large overlap of the low-lying muonic states with the nuclear charge distribution.
- The measurement of the muonic energy levels allows to extract properties of the nucleus such as the charge radius

2p-1s transition most sensitive to R_N

$$R_N^2 = \langle r^2 \rangle = \frac{1}{Ze} \int d^3\vec{r} \rho(\vec{r}) r^2$$

nuclear charge distribution

Back to the muonic cascade

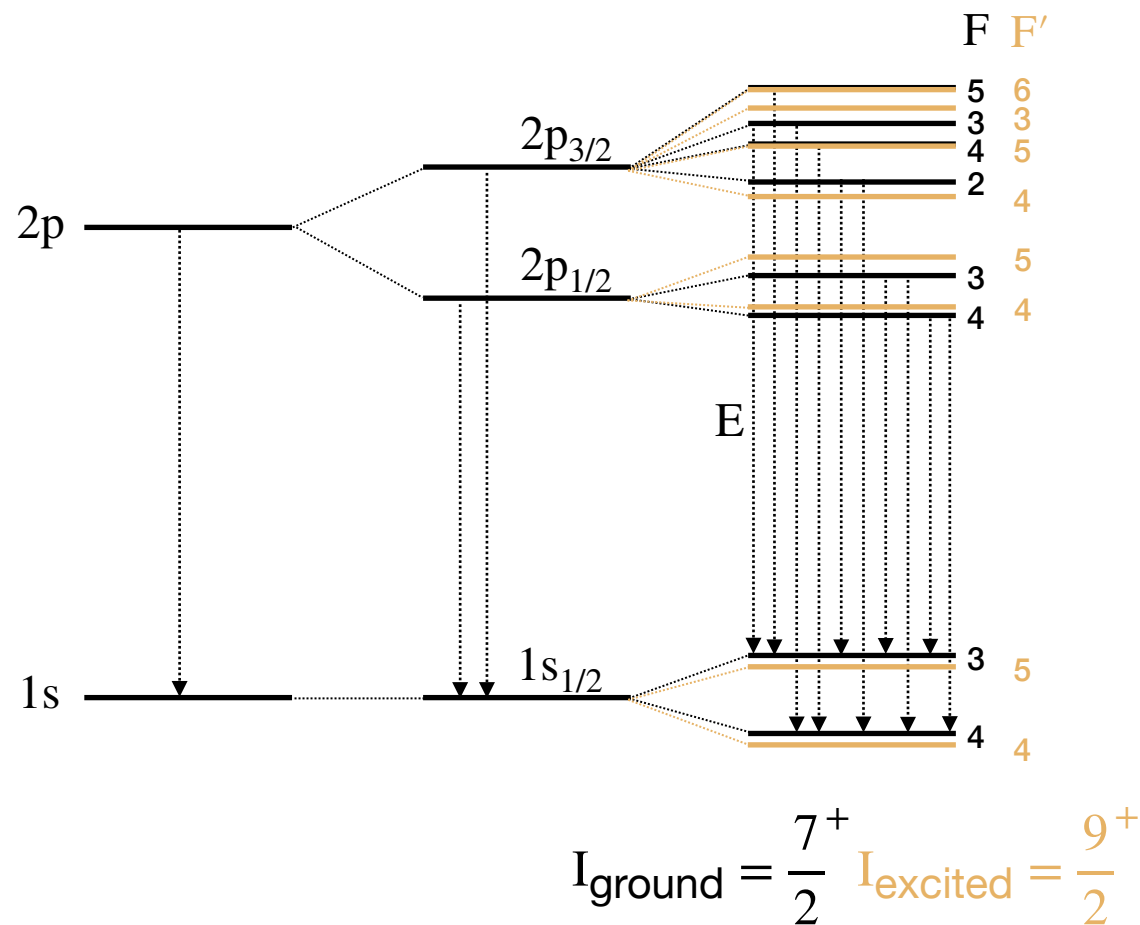


Muonic atoms & Hyperfine Splitting

The blessing and the curse

Fine splitting (FS): $\vec{J} = \vec{I} + \vec{s}$

Static hyperfine splitting (HFS): $\vec{F} = \vec{I} + \vec{J}$



- Energy displacement due to the E2 and M1 interactions:

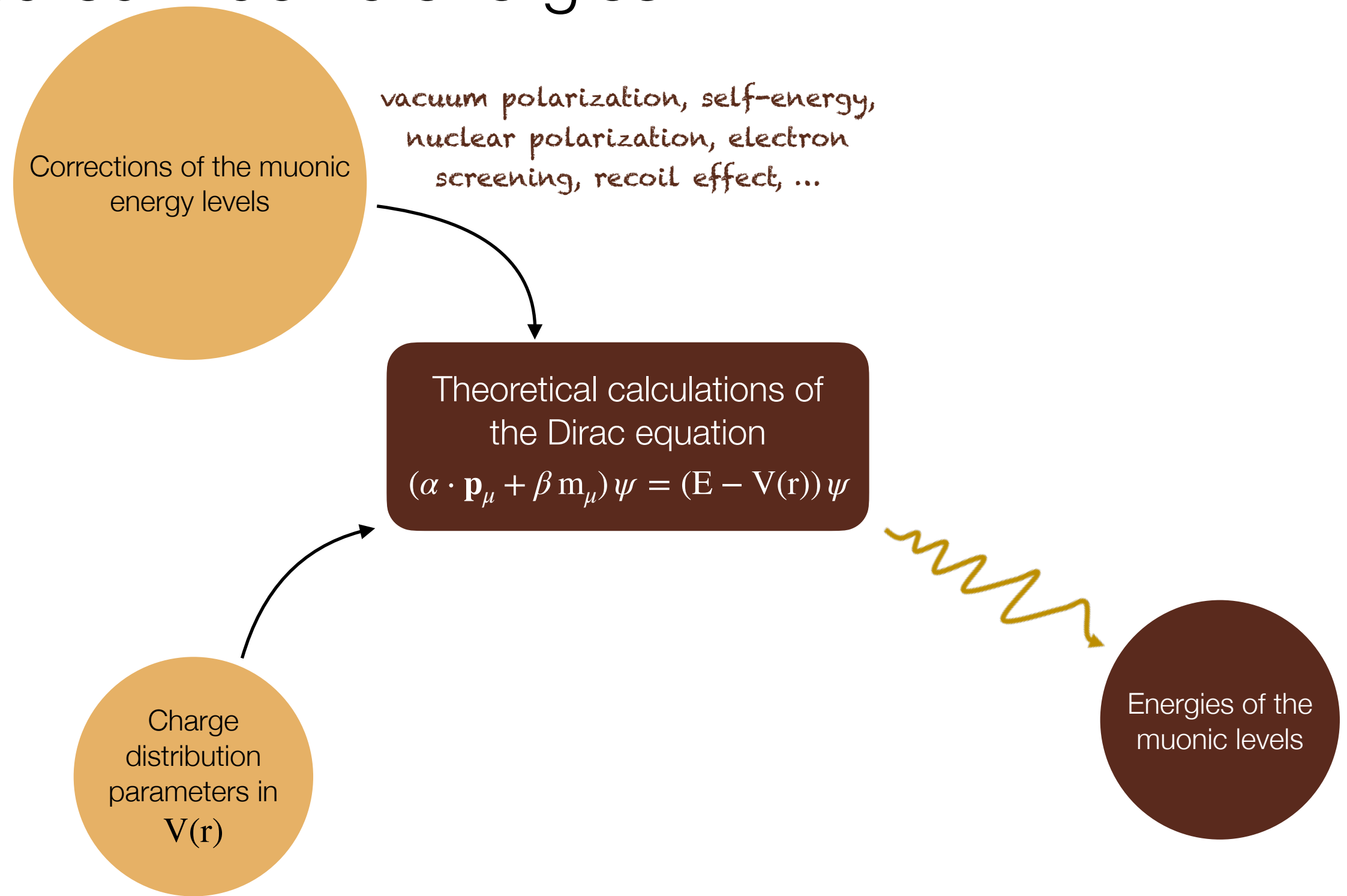
$$\Delta E_F(\text{E2}) = A_2 6 \frac{K(K+1) - 4/3 I(I+1)J(J+1)}{4I(2I-1)J(2J-1)}$$

$$\Delta E_F(\text{M1}) = \frac{A_1}{2} \{F(F+1) - I(I+1) - J(J+1)\}$$

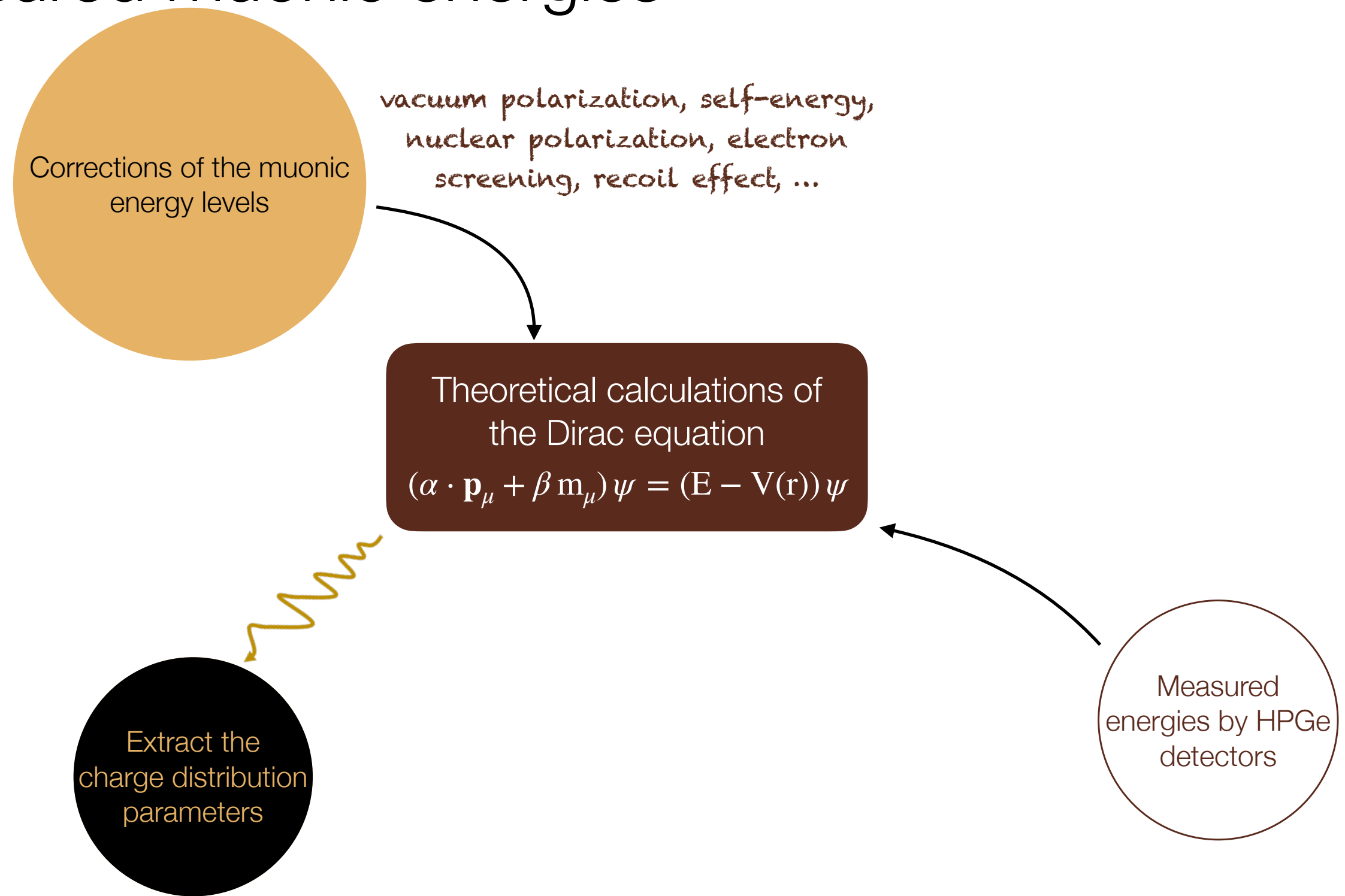
Dynamic HFS effect

- The hyperfine levels from ground and excited nuclear states are mixed due to the high energy of muonic transitions
- HFS also observed in even-even nuclei with zero spin in the ground state

How to extract nuclear charge parameters from measured muonic energies

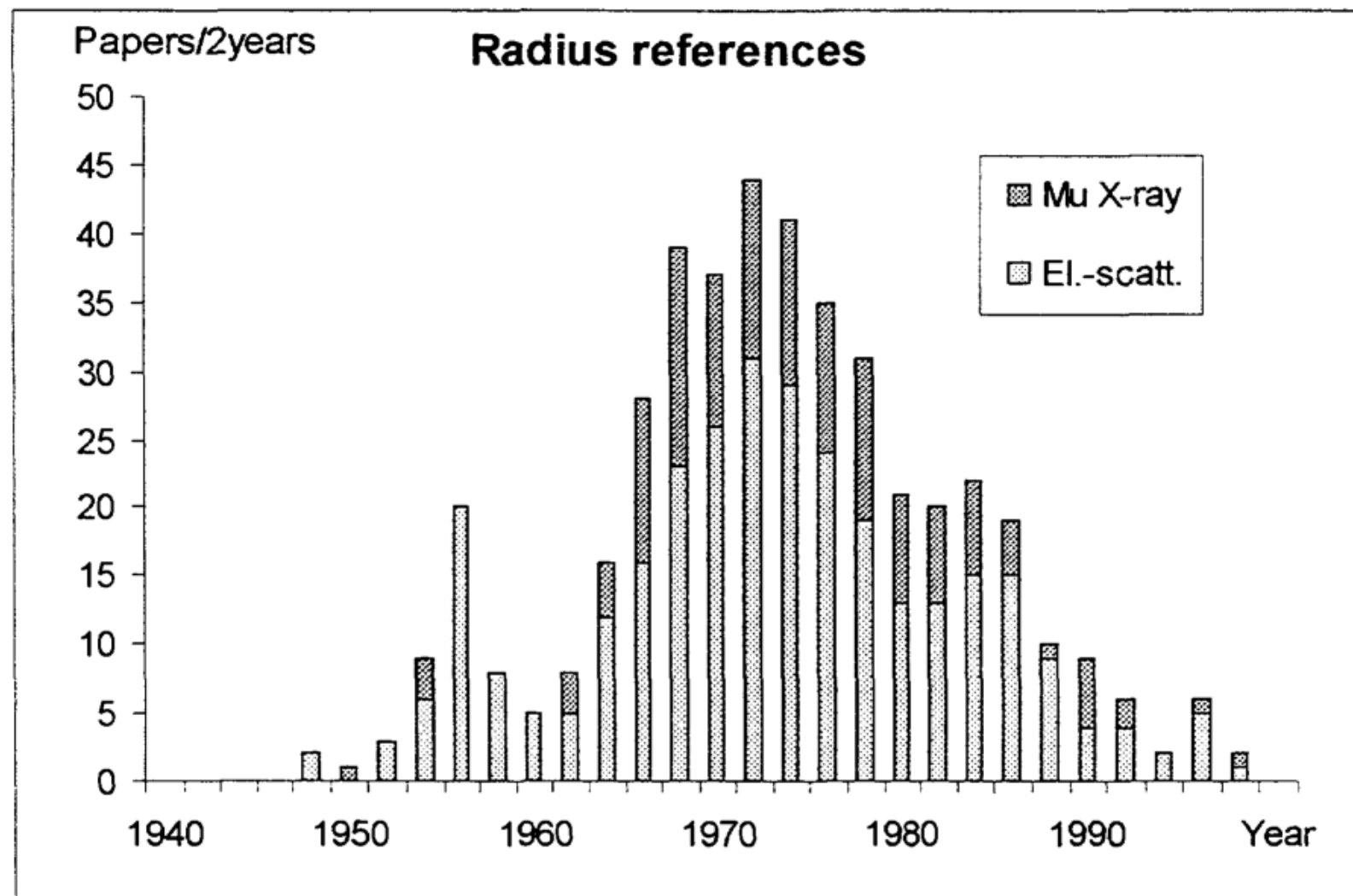


How to extract nuclear charge parameters from measured muonic energies



A technique used heavily in the past..

Angeli, I., & International Atomic Energy Agency, International Nuclear Data Committee, Vienna (Austria). (1999). Table of nuclear root mean square charge radii.



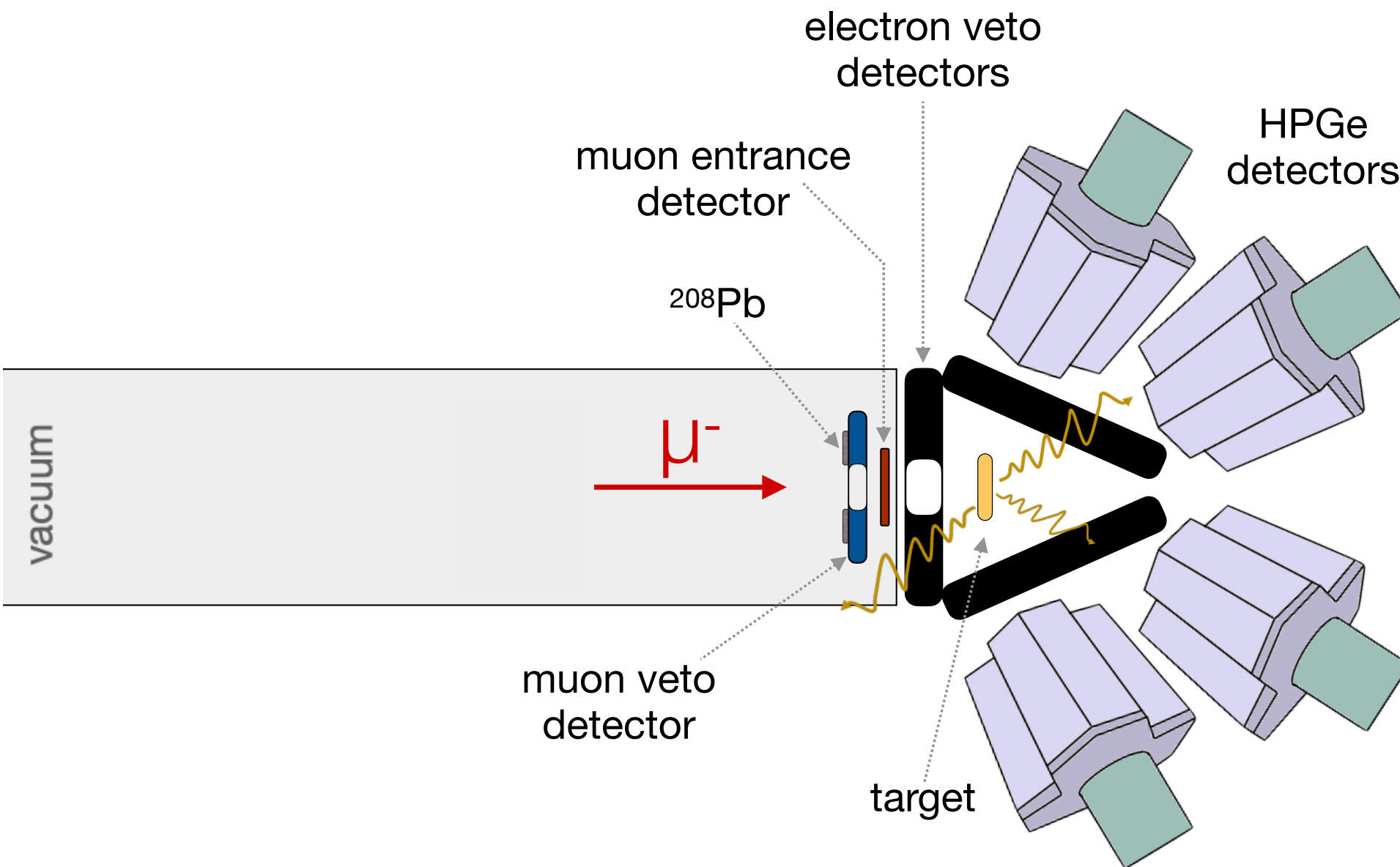
Down to a few
microgram
targets!!!

Extending measurements to rare and radioactive targets.

High Intensity Proton
Accelerator (HIPA) facility at Paul
Scherrer Institute



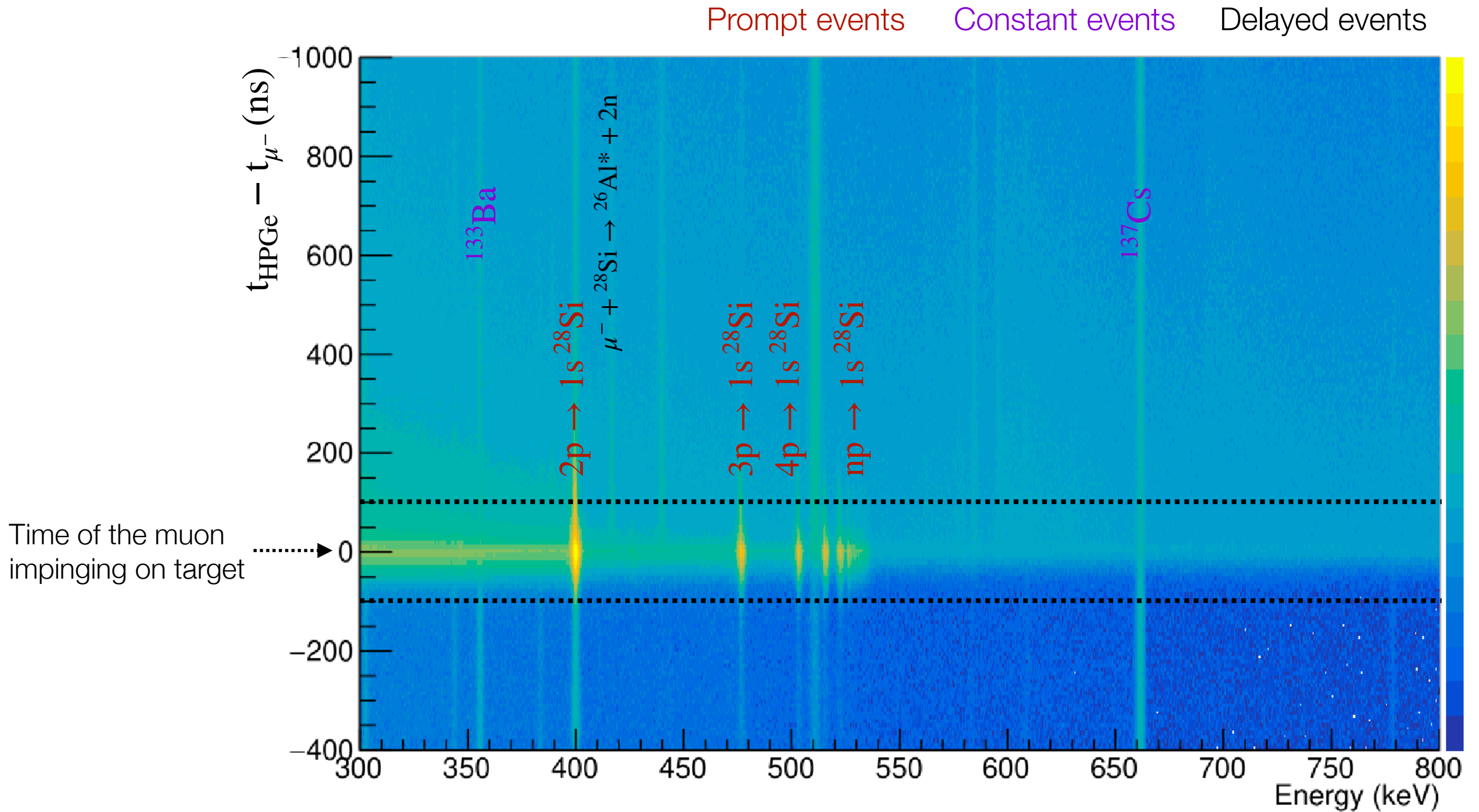
Detection system *Schematics*



Macroscopic
 ^{28}Si target

Detection system

Muon-time correlated analysis



Detection system

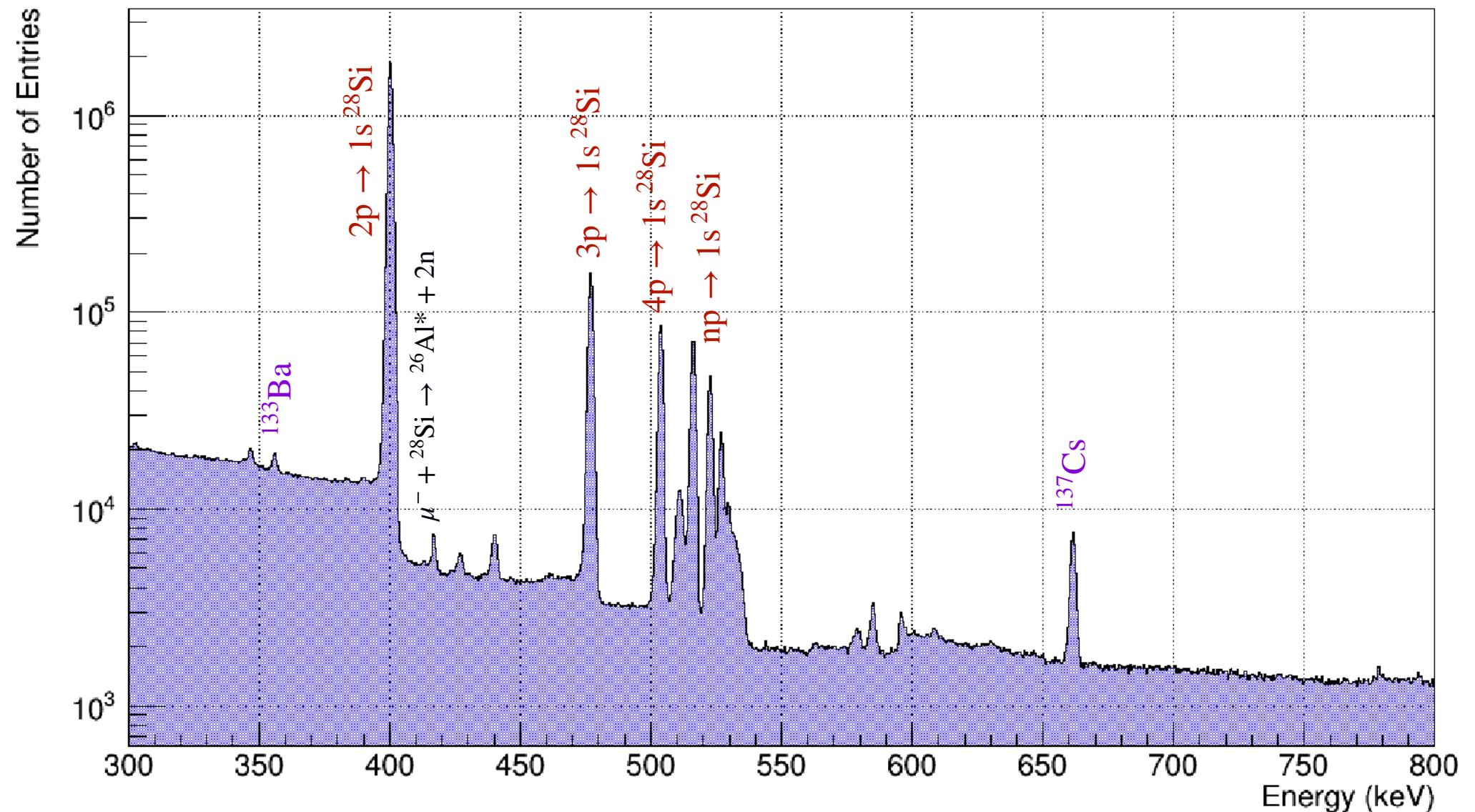
Muon-time correlated analysis

After ELET : $\text{FWHM}_{\text{dt}} \sim 14 \text{ nsec}$

Prompt events

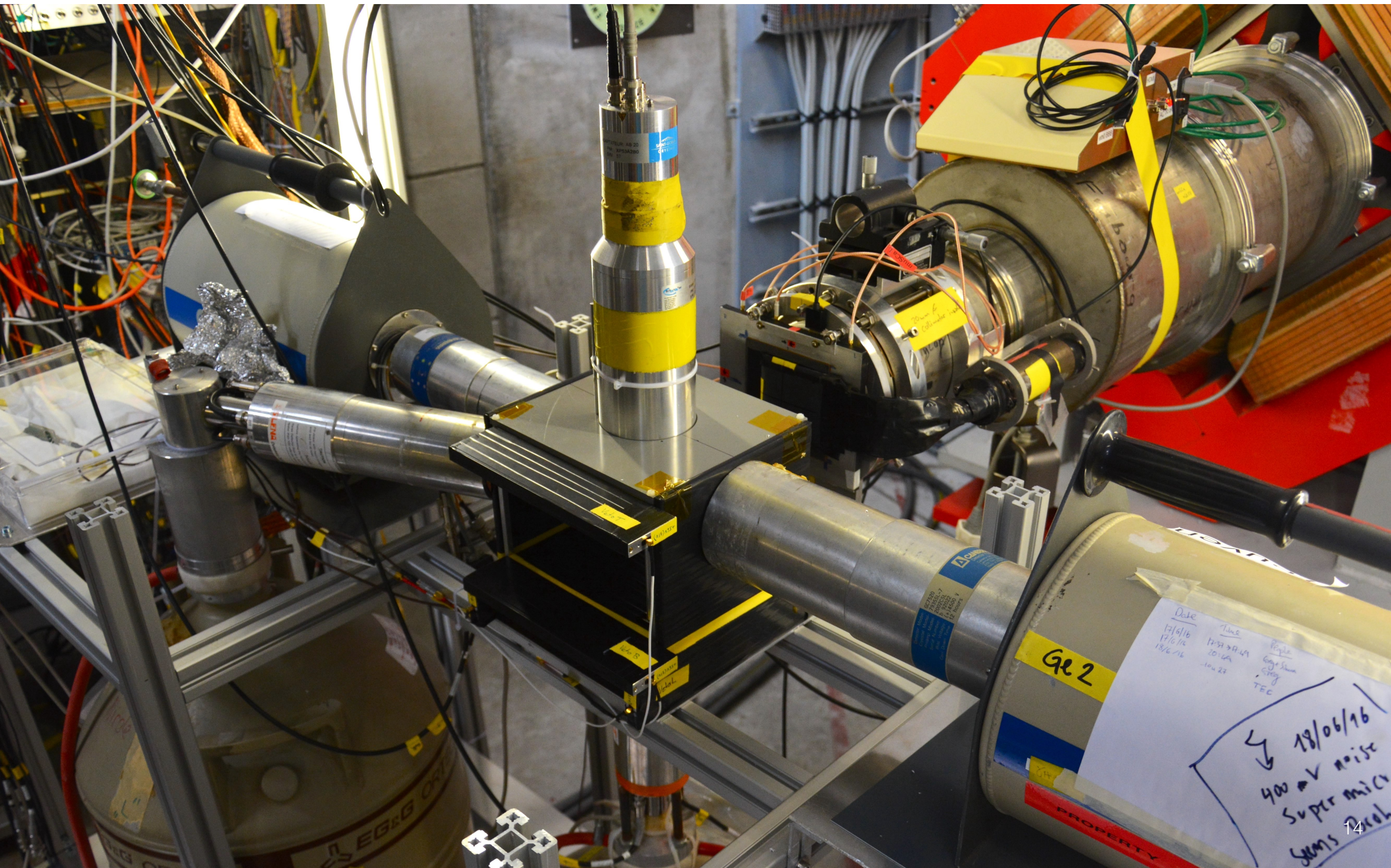
Constant events

Delayed events



Experimental timeline

2016 — ^{185}Re and ^{187}Re (500 mg)



^{185}Re and ^{187}Re 5g4f analysis

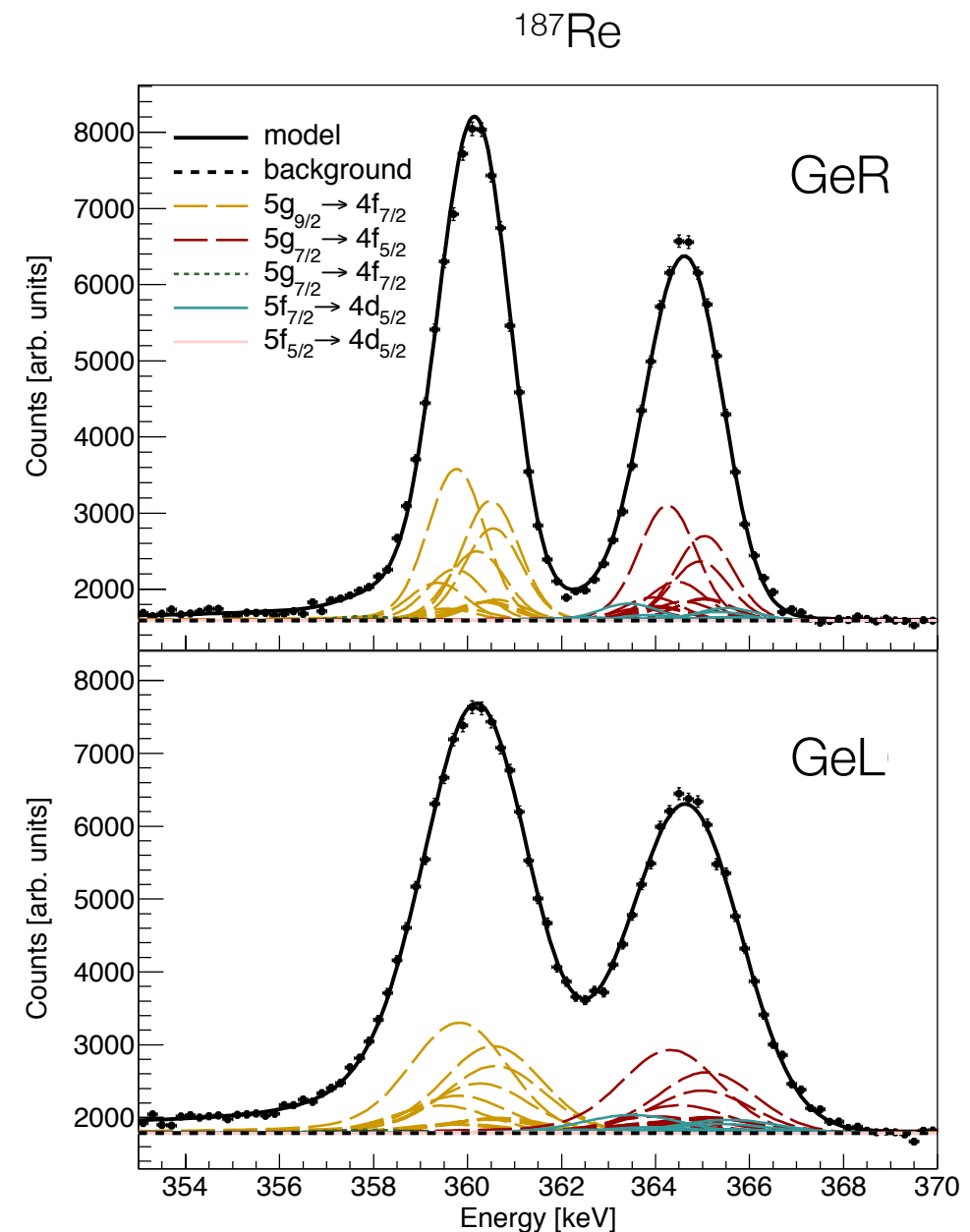
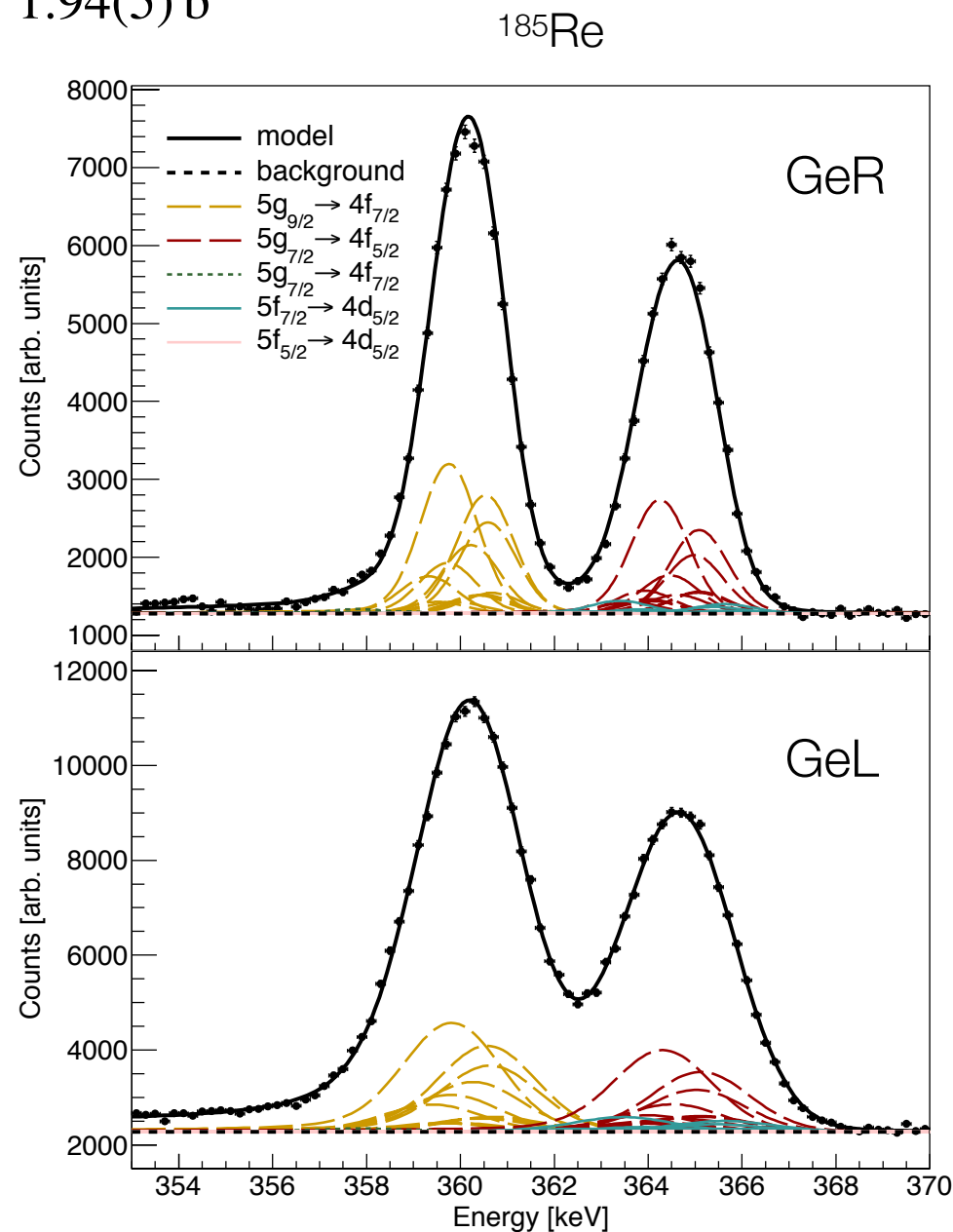
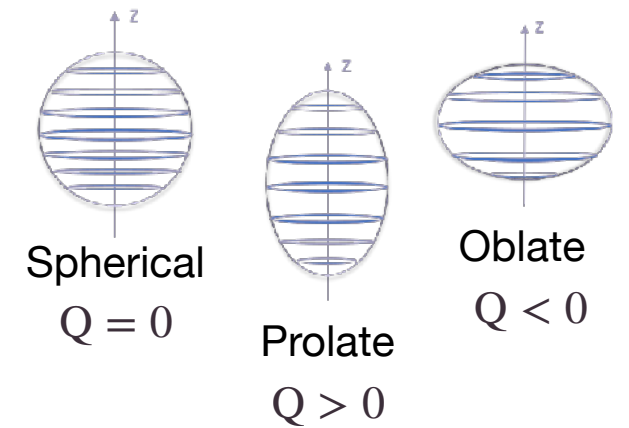
Spectroscopic quadrupole moment

Phys. Rev. C 101, 054313 (2020)

$$Q(^{185}\text{Re}) = 2.07(5) \text{ b}$$

$$Q(^{187}\text{Re}) = 1.94(5) \text{ b}$$

Our first physics
result

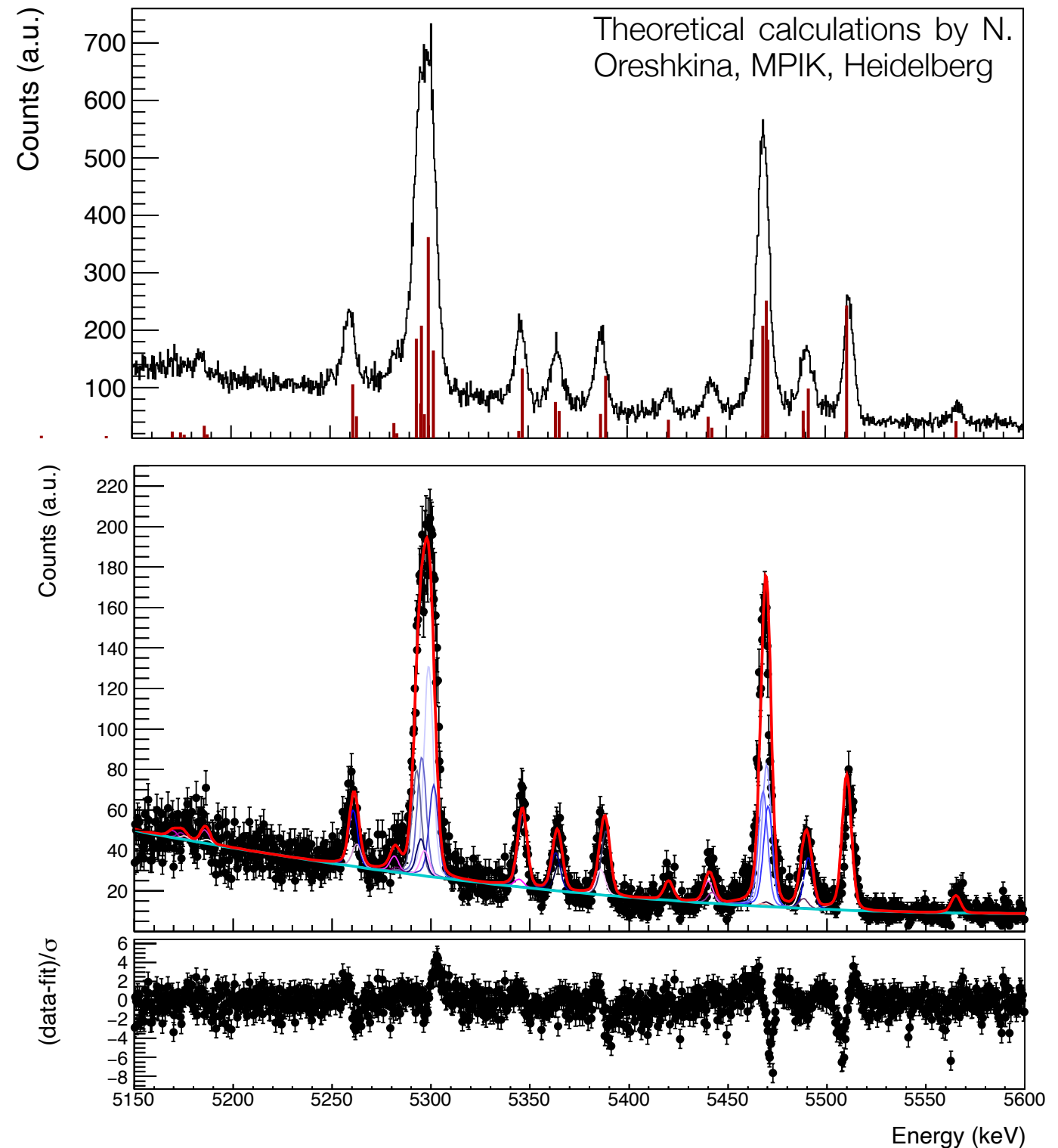


^{185}Re and ^{187}Re 2p1s analysis

Outlook: Nuclear charge radius

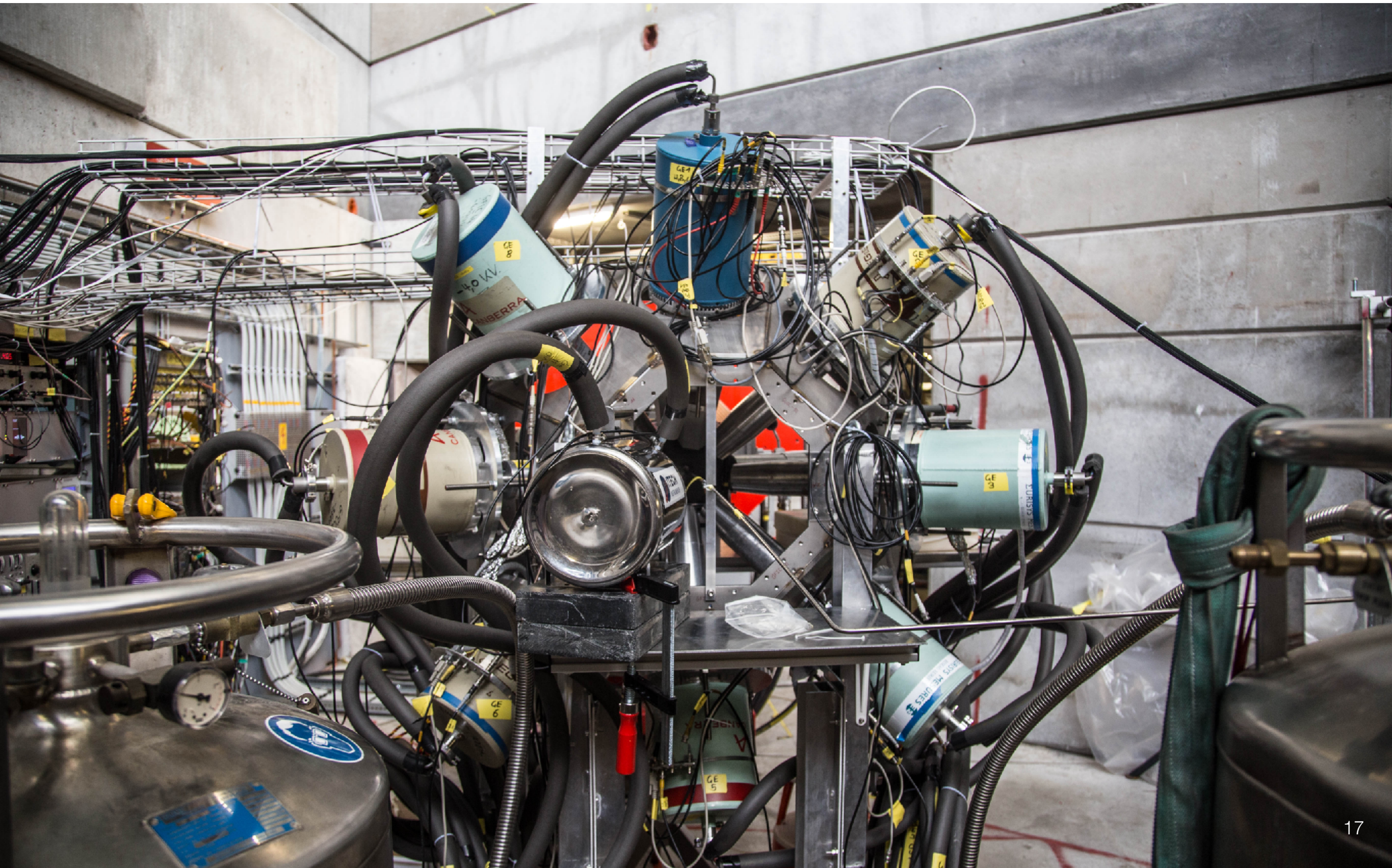
PhD thesis of S.M. Vogiatzi

- First time attempting to measure the last stable element with a non-measured absolute nuclear charge radius.
- Admixture of ground and excited hyperfine states due to the dynamic effect — very complicated muonic energy spectrum.



Experimental timeline

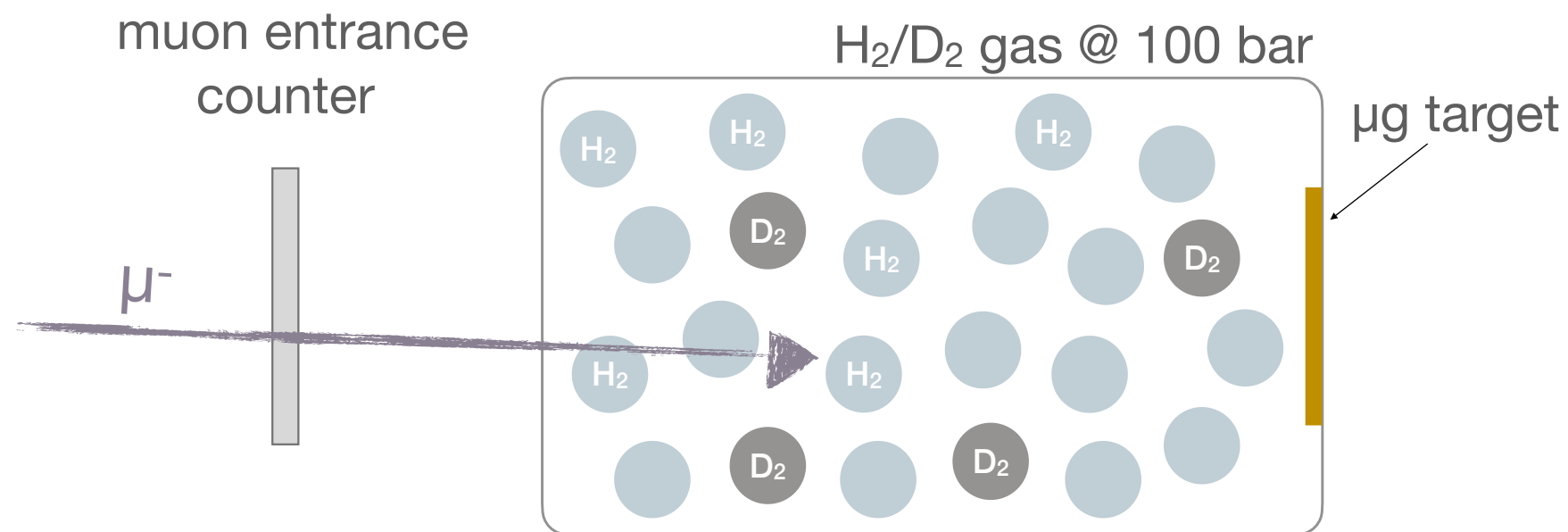
2017 — Towards microgram targets



Muon transfer to microgram targets

Developed by the muX collaboration [Eur. Phys. J. A 59, 15 \(2023\)](#)

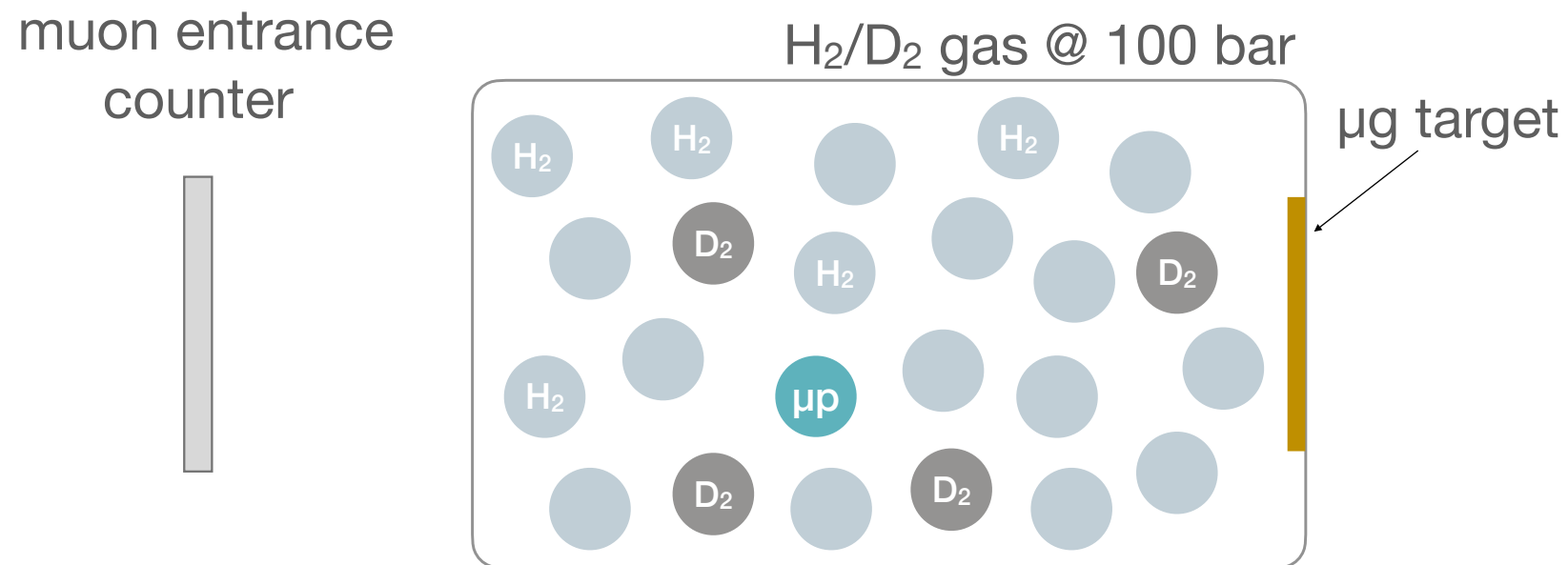
Inspired by the work of Strasser *et al.* and Kraiman *et al.*



Muon transfer to microgram targets

Developed by the muX collaboration [Eur. Phys. J. A 59, 15 \(2023\)](#)

Inspired by the work of Strasser *et al.* and Kraiman *et al.*

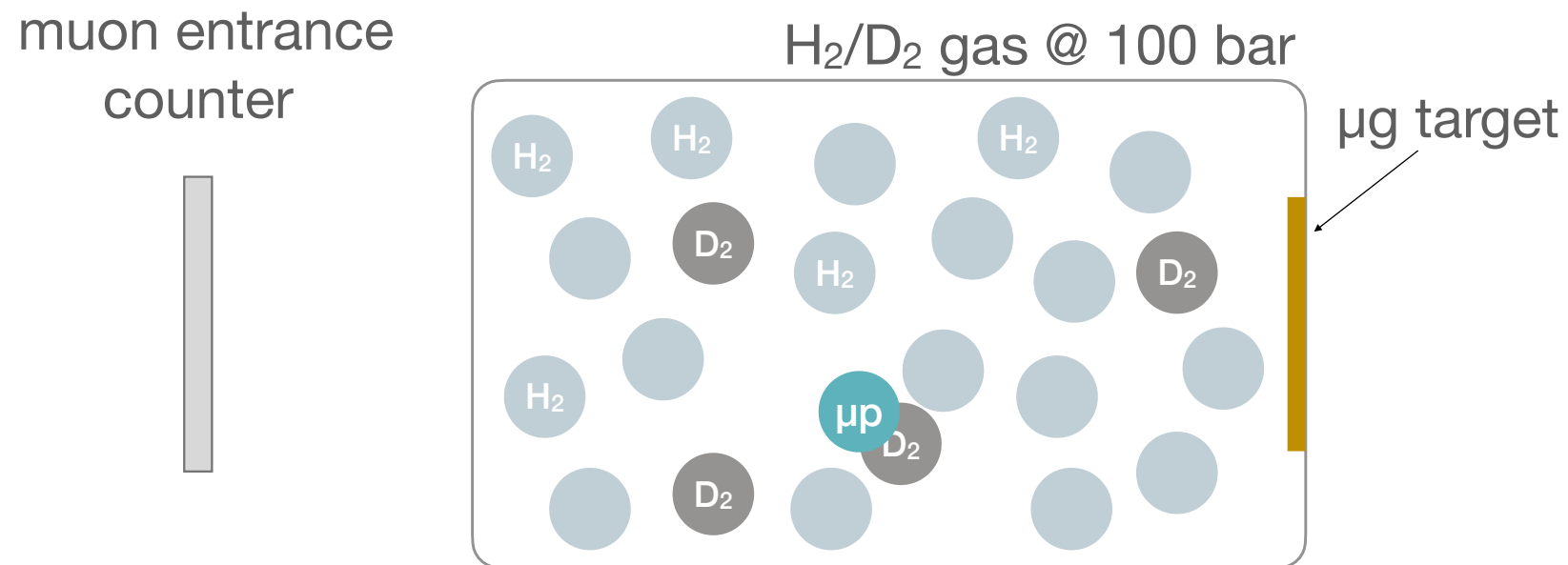


1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp

Muon transfer to microgram targets

Developed by the muX collaboration [Eur. Phys. J. A 59, 15 \(2023\)](#)

Inspired by the work of Strasser *et al.* and Kraiman *et al.*

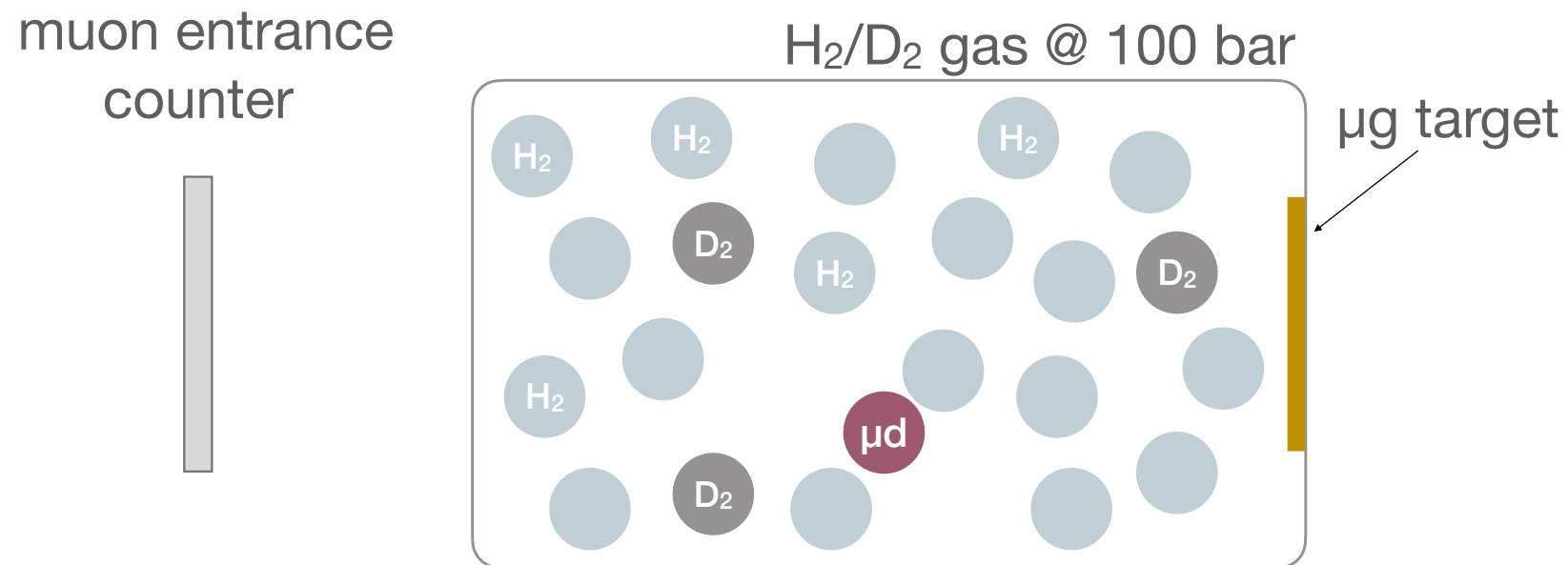


1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp
2. transfer to deuterium $\mu p \rightarrow \mu d$

Muon transfer to microgram targets

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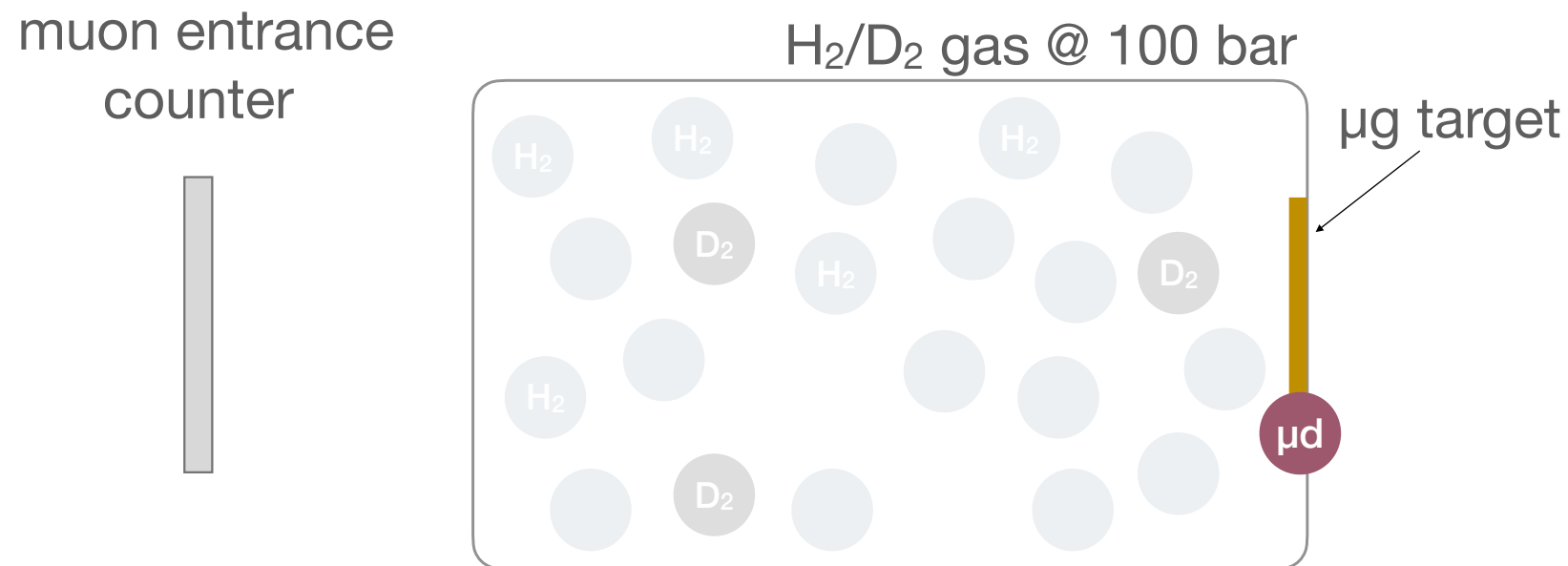


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Muon transfer to microgram targets

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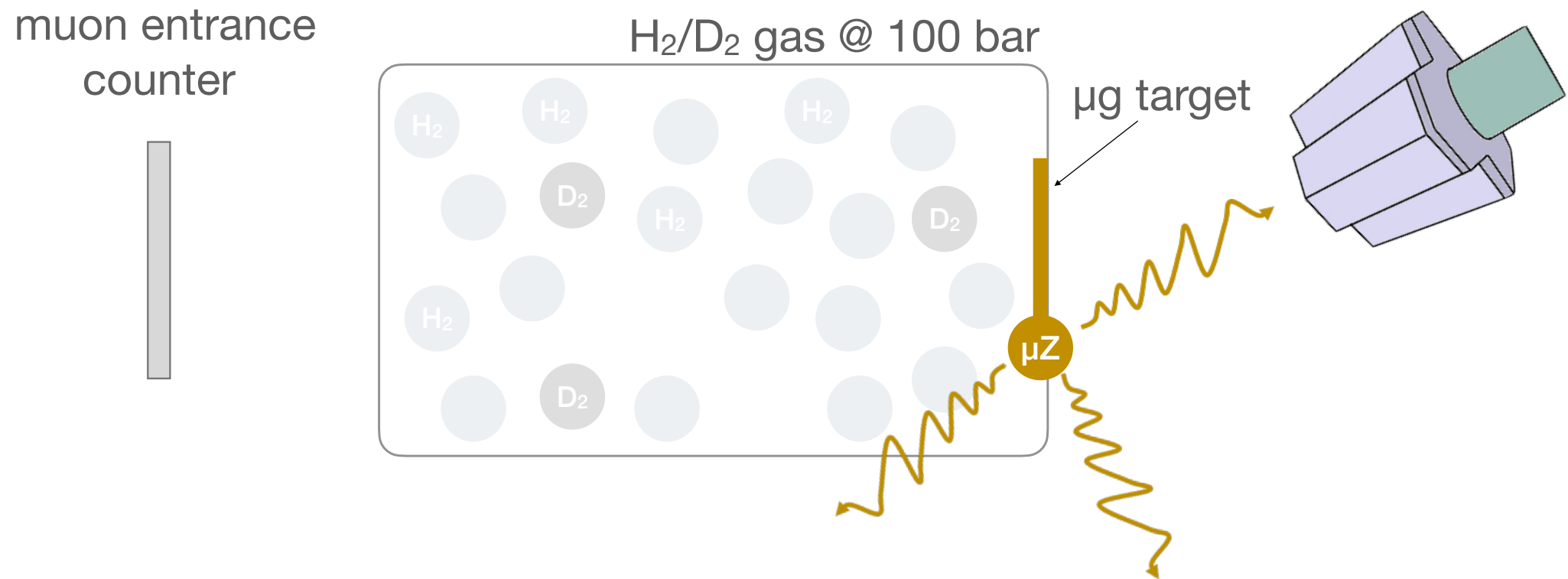
1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp
2. transfer to deuterium $\mu p \rightarrow \mu d$
3. μd moves almost freely in the H₂ gas

[Physical Review A 73, 034501 \(2006\)](#)

Muon transfer to microgram targets

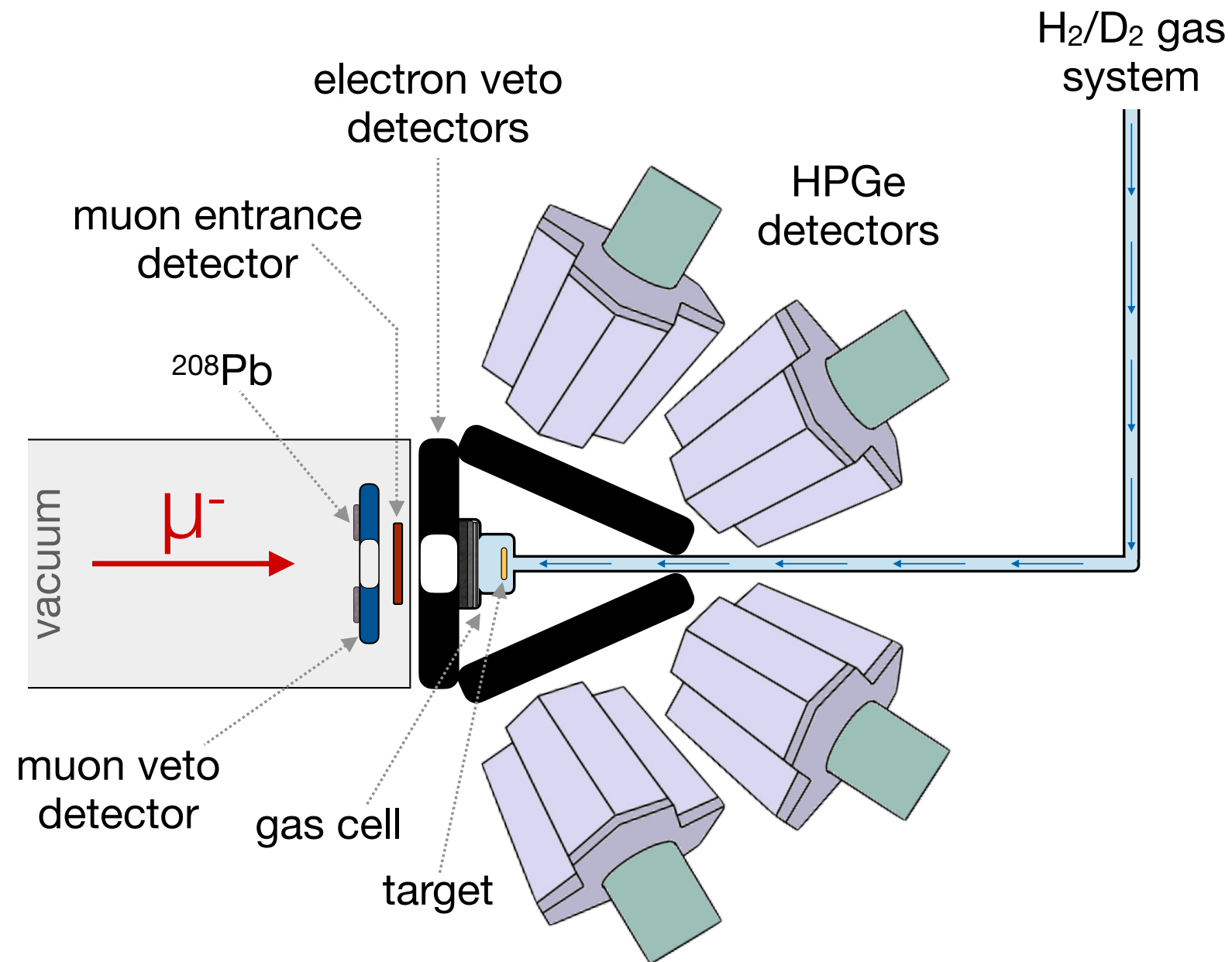
Developed by the muX collaboration [Eur. Phys. J. A 59, 15 \(2023\)](#)

Inspired by the work of Strasser *et al.* and Kraiman *et al.*

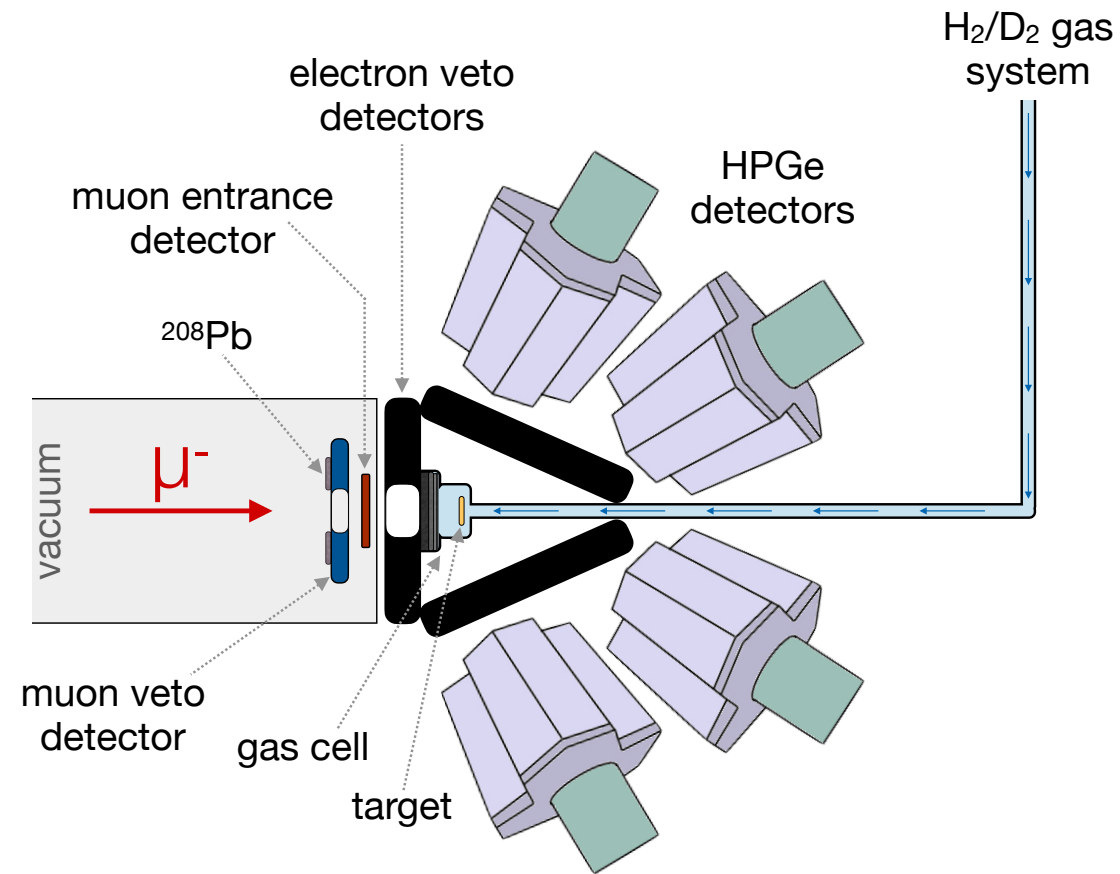


1. μ^- stops in 100 bar of H₂ + 0.25% D₂ & forms muonic hydrogen μp
2. transfer to deuterium $\mu p \rightarrow \mu d$
3. μd moves almost freely in the H₂ gas [Physical Review A 73, 034501 \(2006\)](#)
4. transfer to high-Z element $\mu d \rightarrow \mu Z$ when hitting target & emission of x rays during the atomic cascade

Gas cell measurements

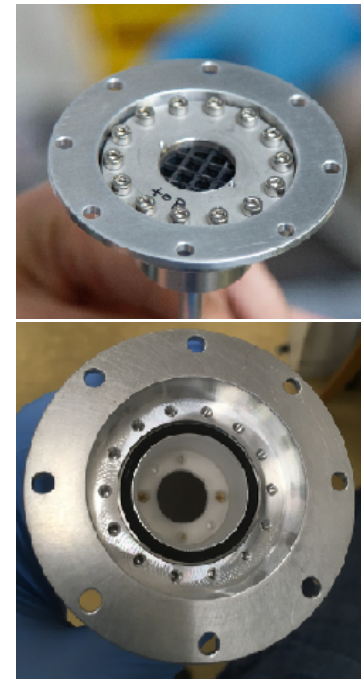
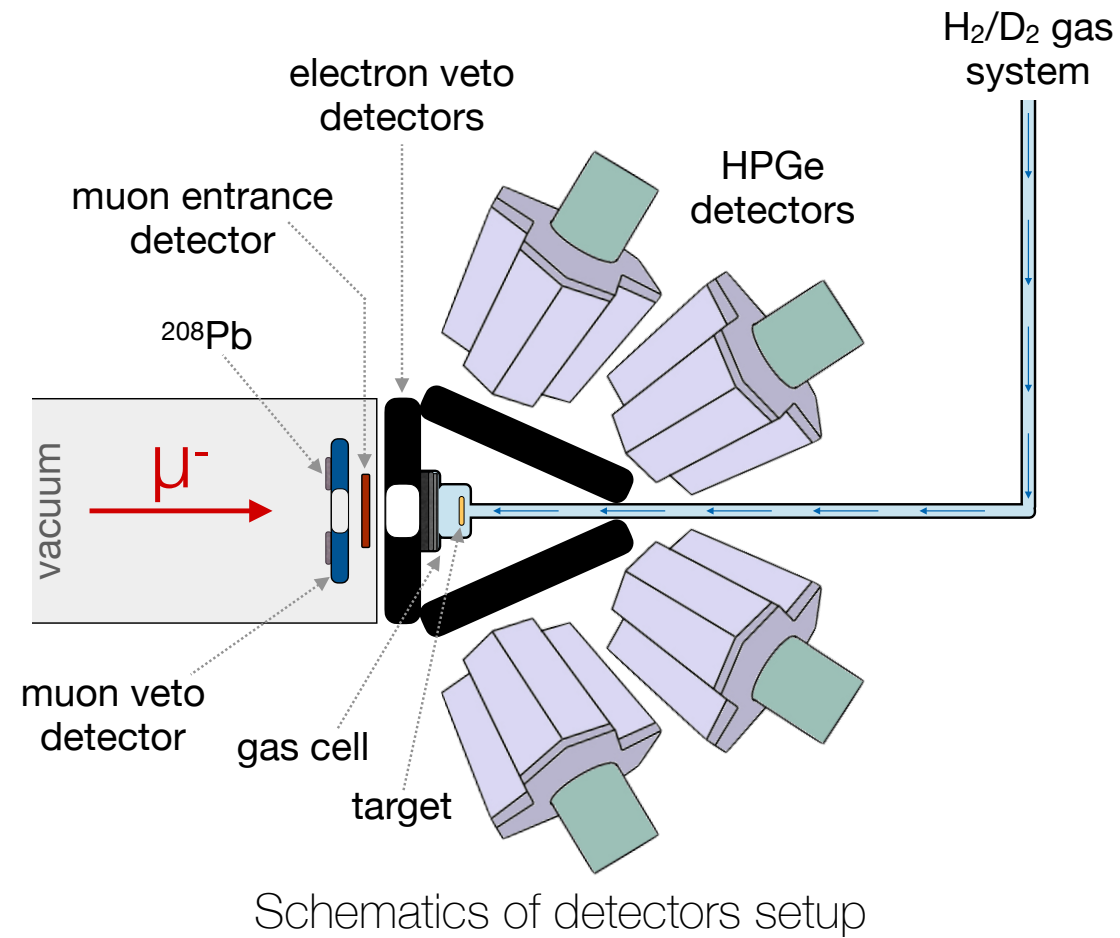


Gas cell measurements

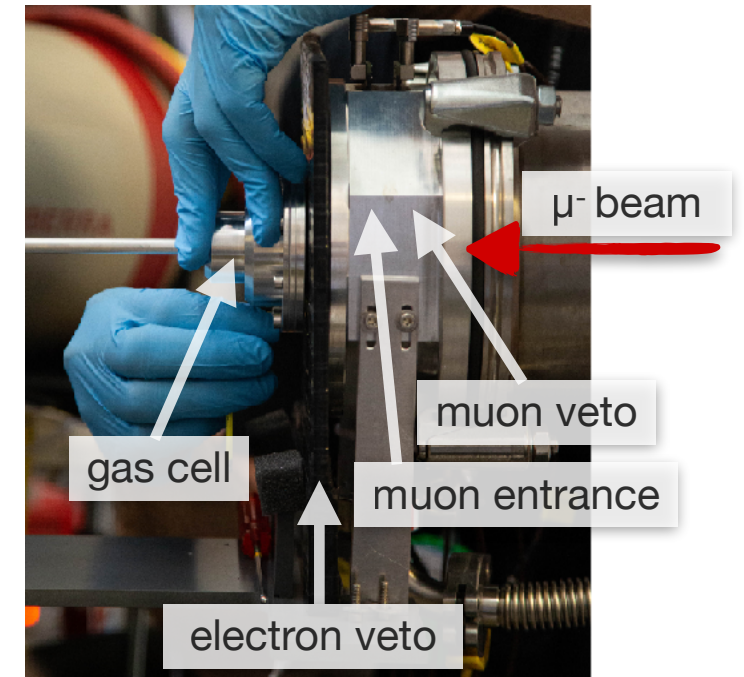


Schematics of detectors setup

Gas cell measurements

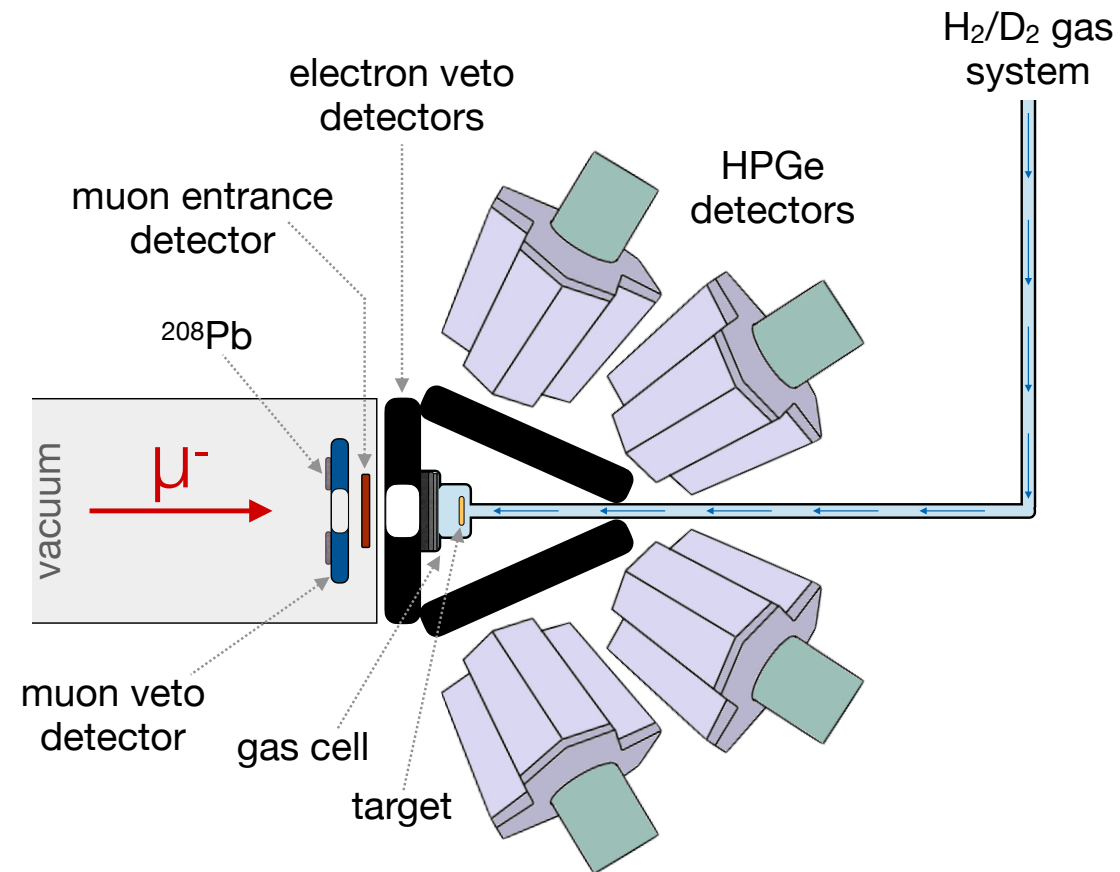


The gas cell

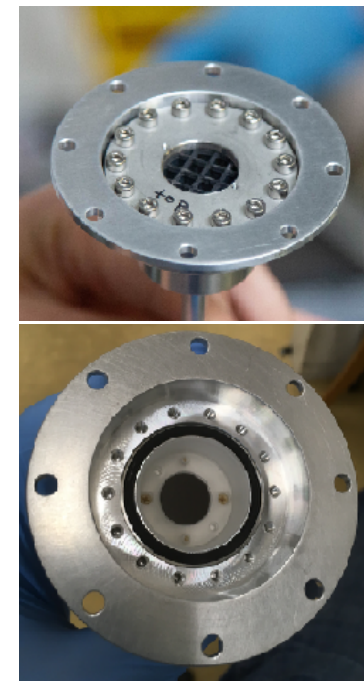


The muon and electron counters

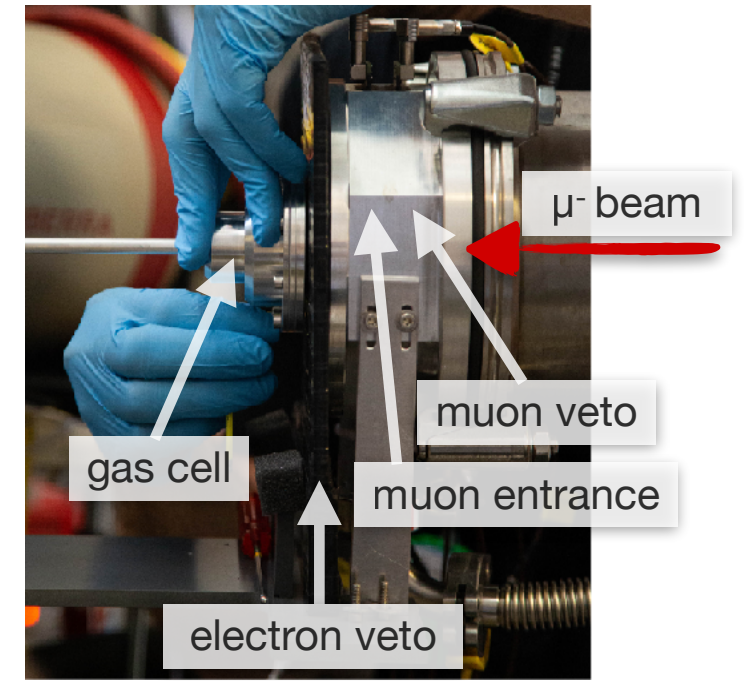
Gas cell measurements



Schematics of detectors setup



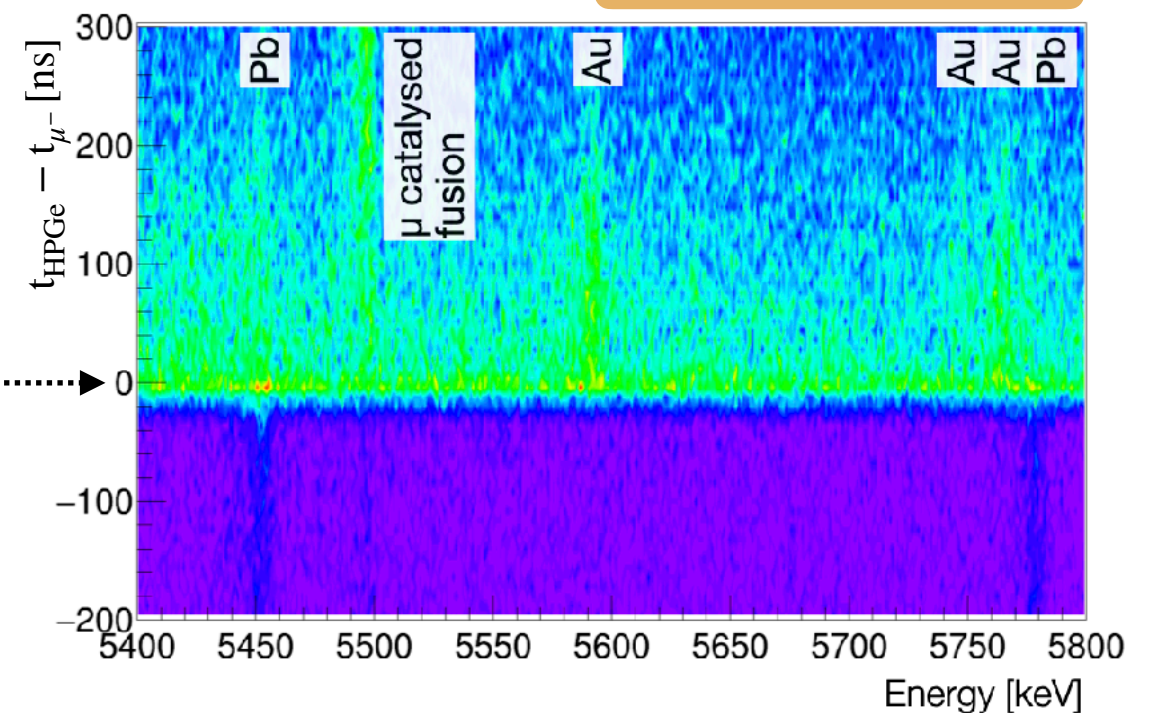
The gas cell



The muon and electron counters

Demonstration of principal in a 5 μg gold target in 2017, 18.5 h of measurement

the time the muon enters the gas cell



Experimental timeline

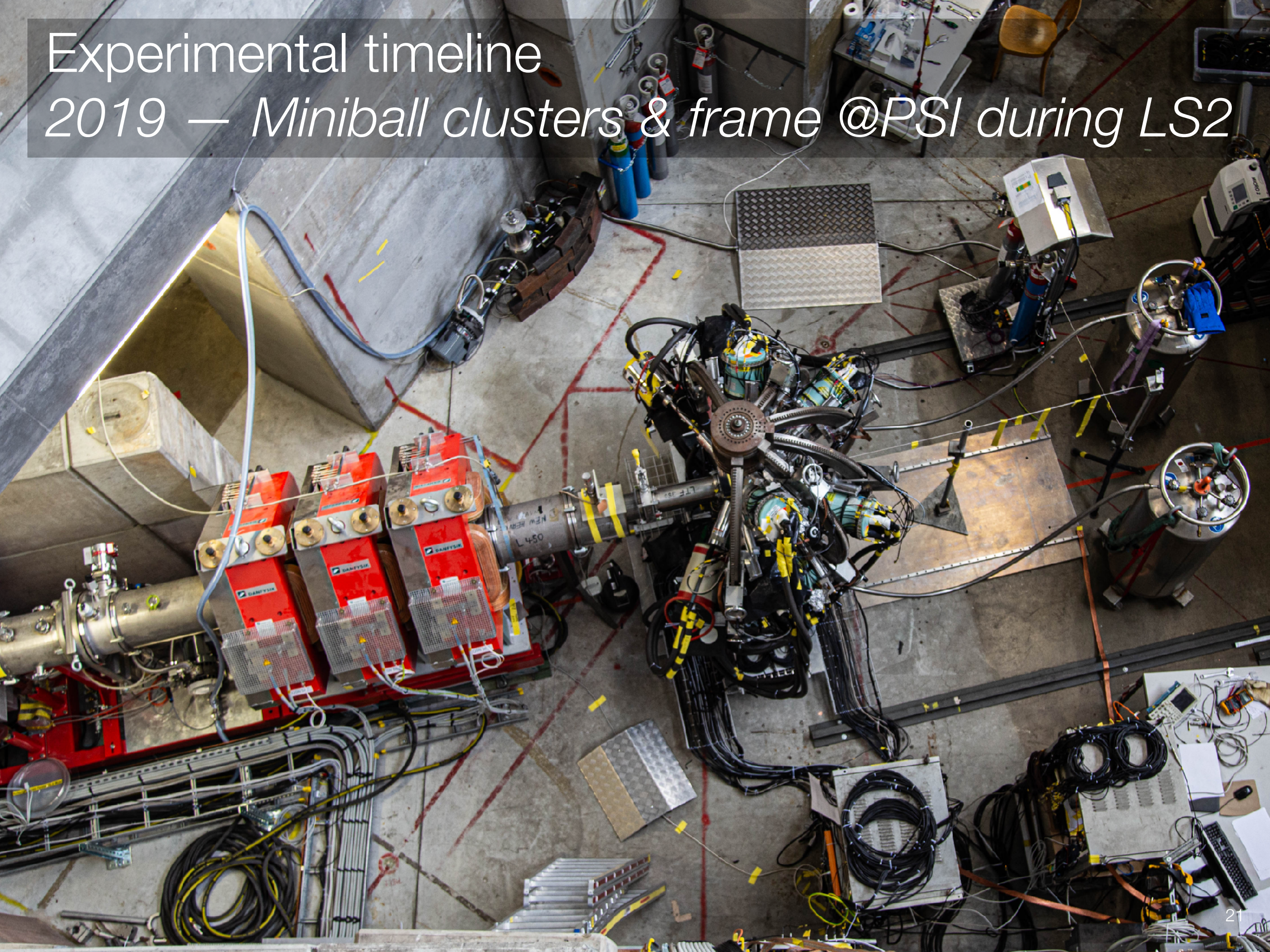
2019 — first radioactive targets ^{248}Cm & ^{226}Ra ($5\ \mu\text{g}$)



MINIBALL array

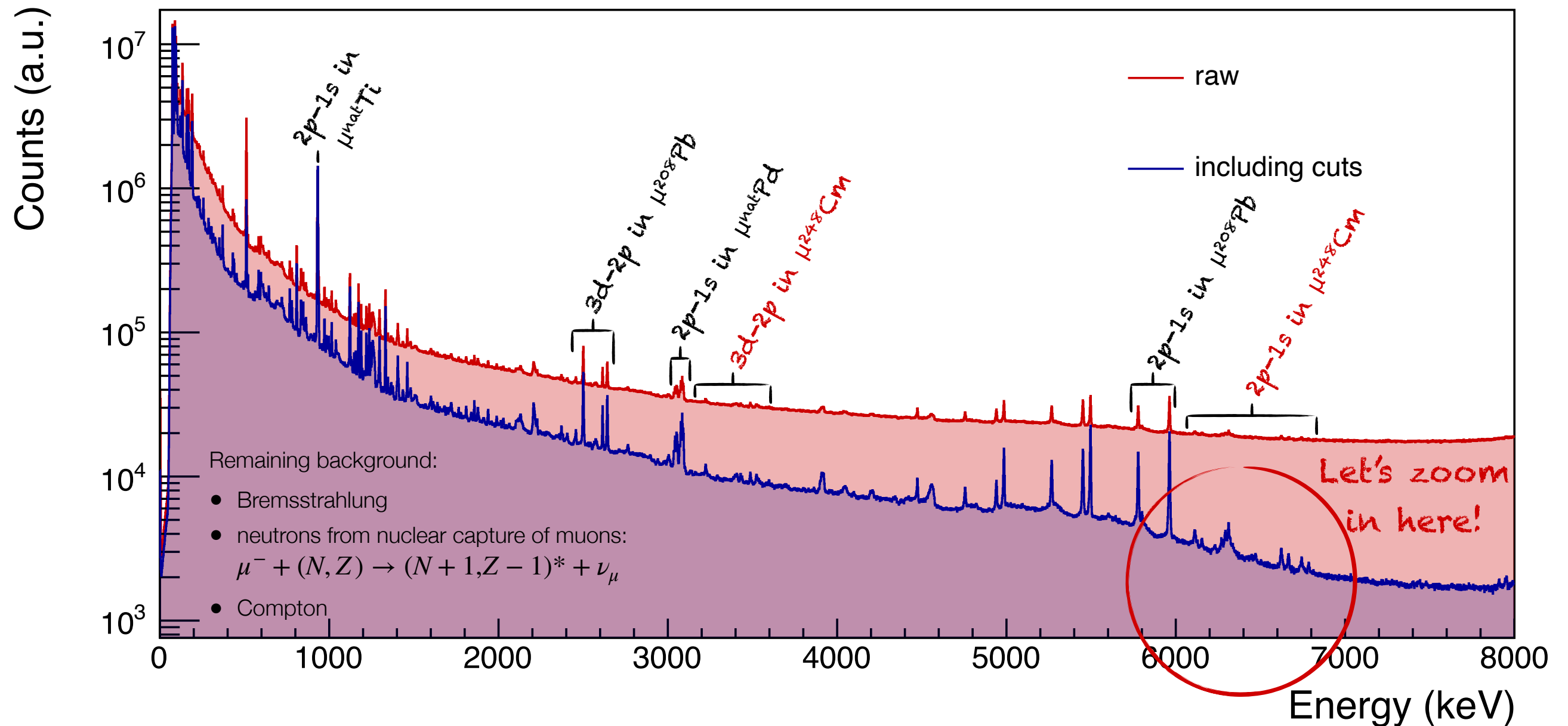
Experimental timeline

2019 — Miniball clusters & frame @PSI during LS2

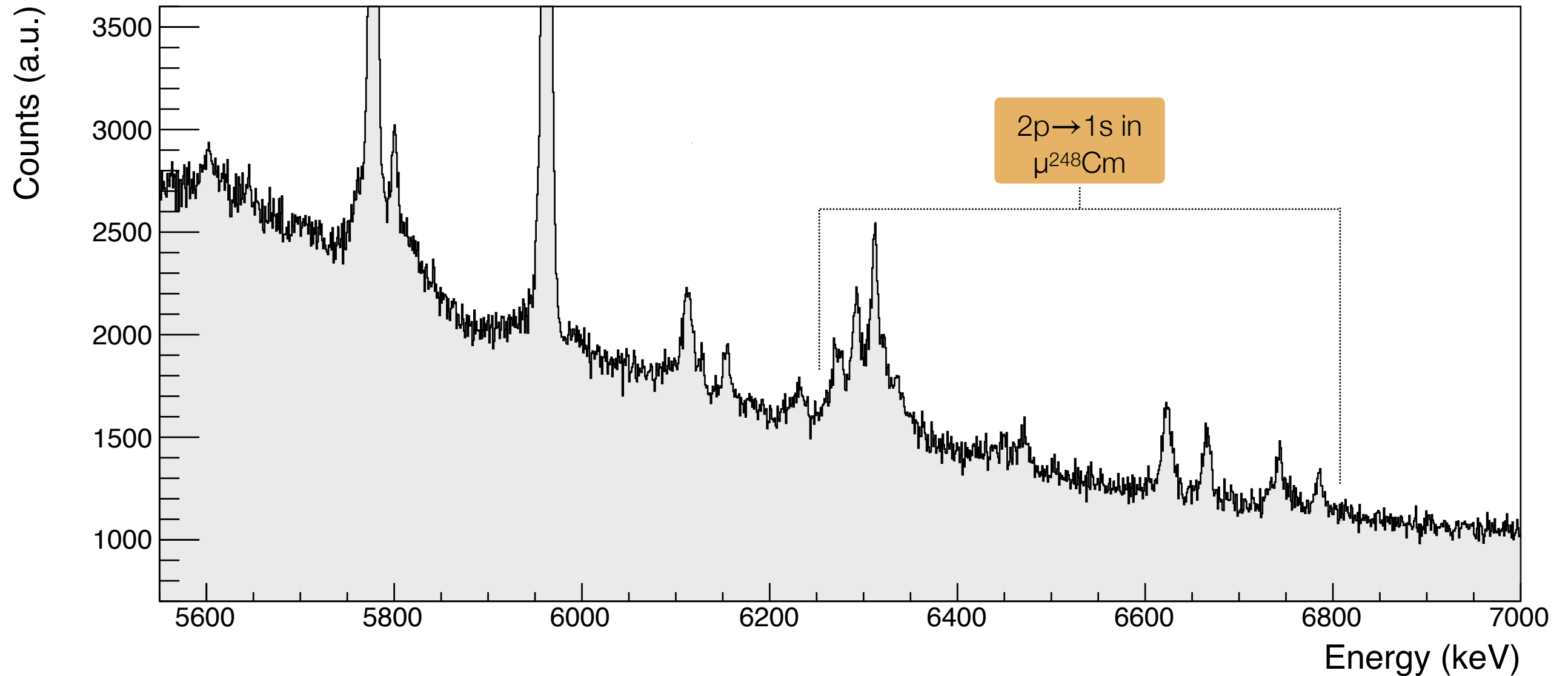


How to get to good S/B ratio at high energies?

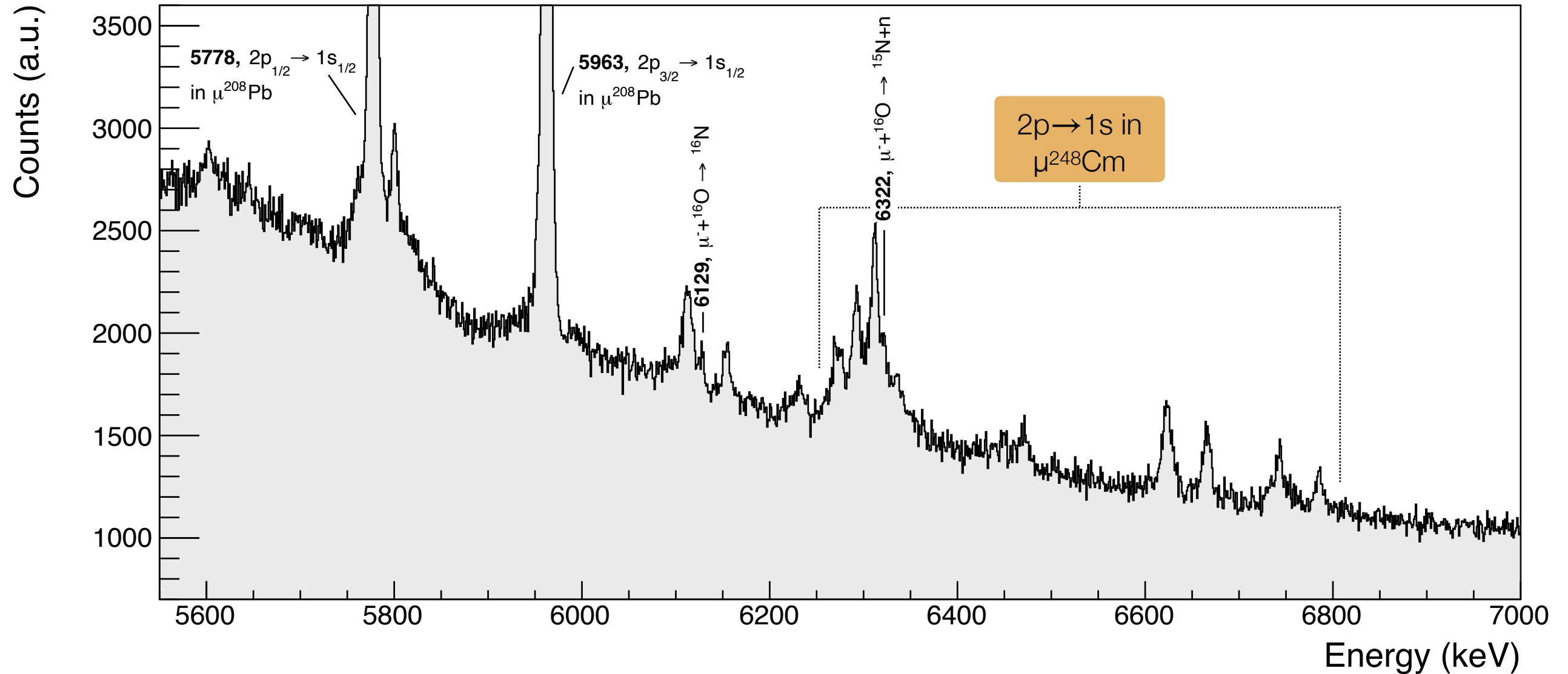
- Electron veto cuts
- Addback
- Gain drift corrections
- Baseline correction



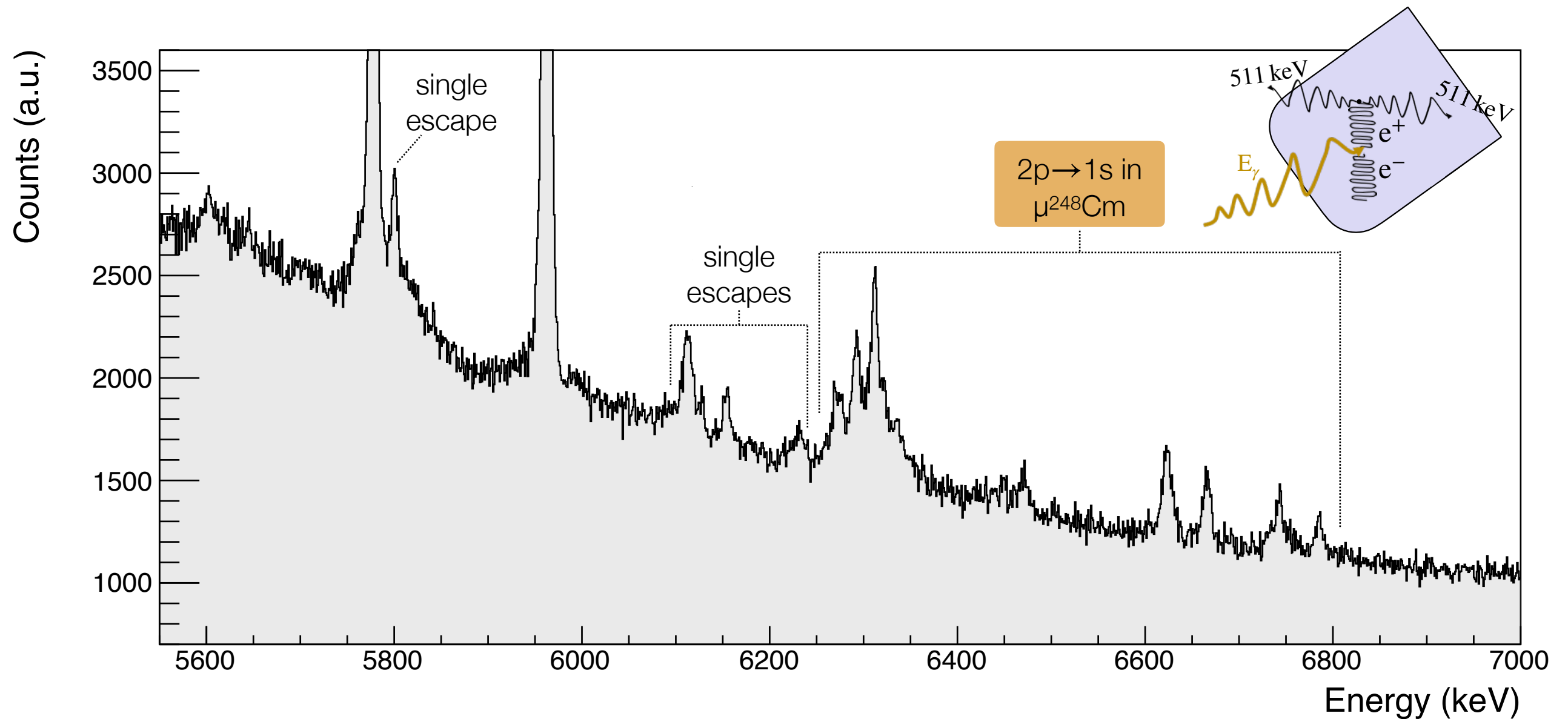
^{248}Cm 2p1s analysis



^{248}Cm 2p1s analysis

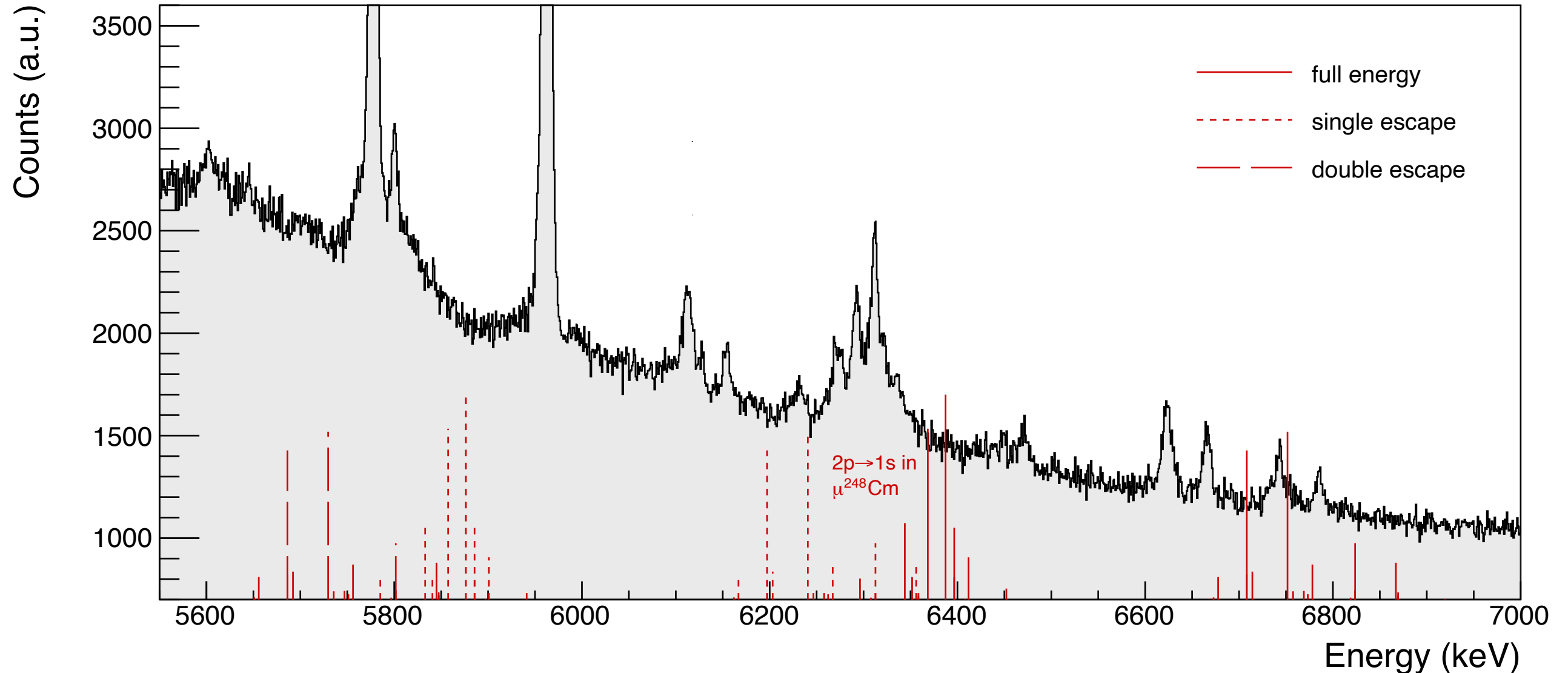


^{248}Cm 2p1s analysis



^{248}Cm 2p1s analysis

Theoretical calculations are performed by
N. Oreshkina & I. Valuev, MPIK, Heidelberg

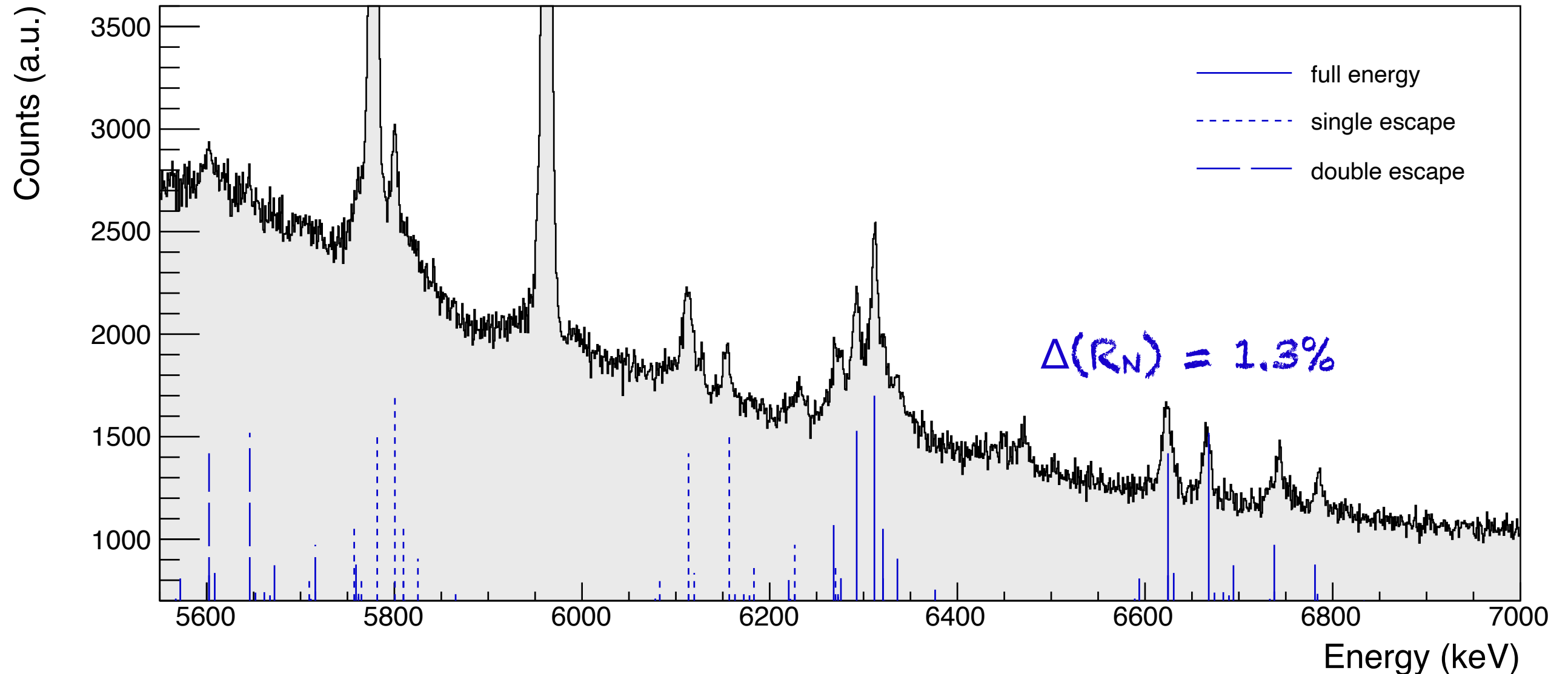


Using the nuclear charge radius
prediction by Angeli & Marinova

$$R = 5.8687 \text{ fm}$$

^{248}Cm 2p1s analysis

Theoretical calculations are performed by
N. Oreshkina & I. Valuev, MPIK, Heidelberg

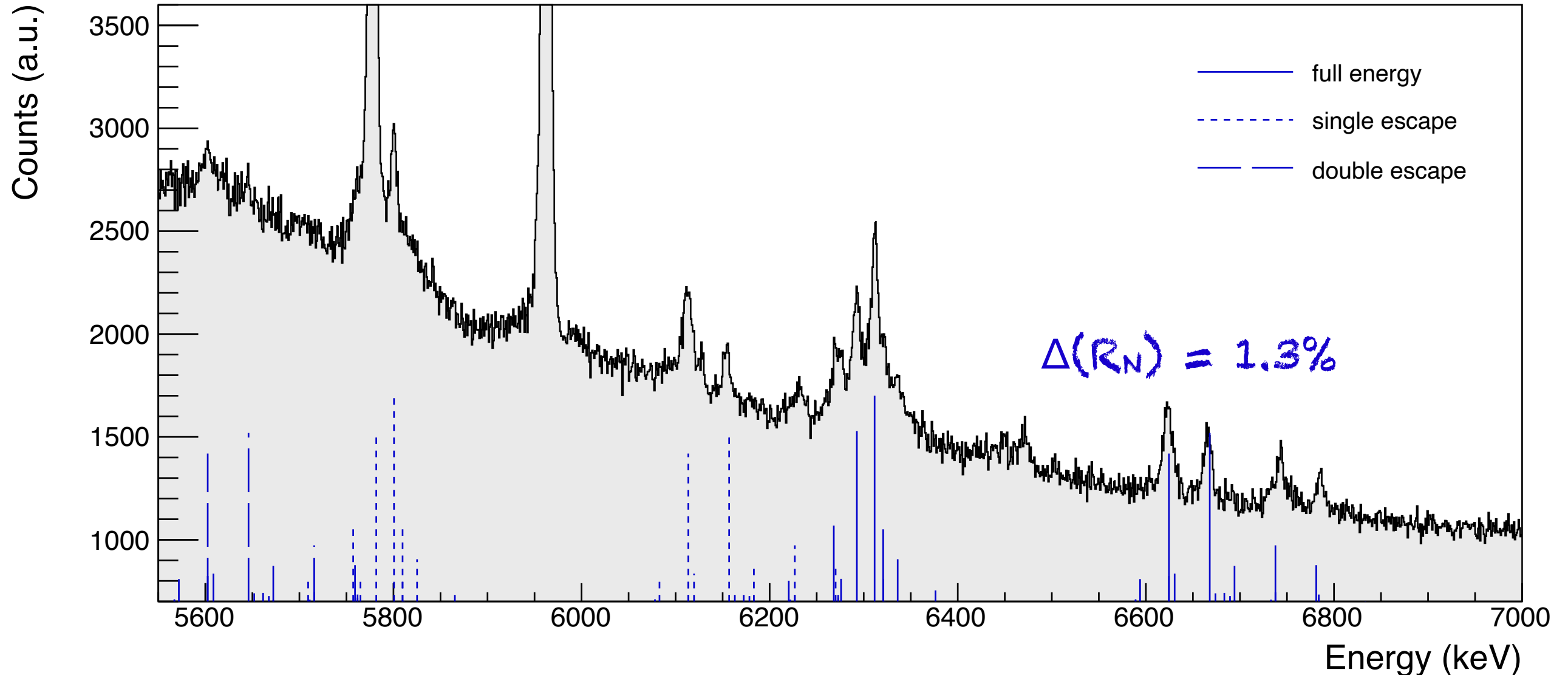


Using the nuclear charge radius
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$$R = 5.8687 \text{ fm}$$

^{248}Cm 2p1s analysis

Theoretical calculations are performed by
N. Oreshkina & I. Valuev, MPIK, Heidelberg



Using the nuclear charge radius
prediction by Angeli & Marinova $R=5.8687$ fm

Sensitivity on the quadrupole moment Q is also observed

^{248}Cm 2p1s analysis

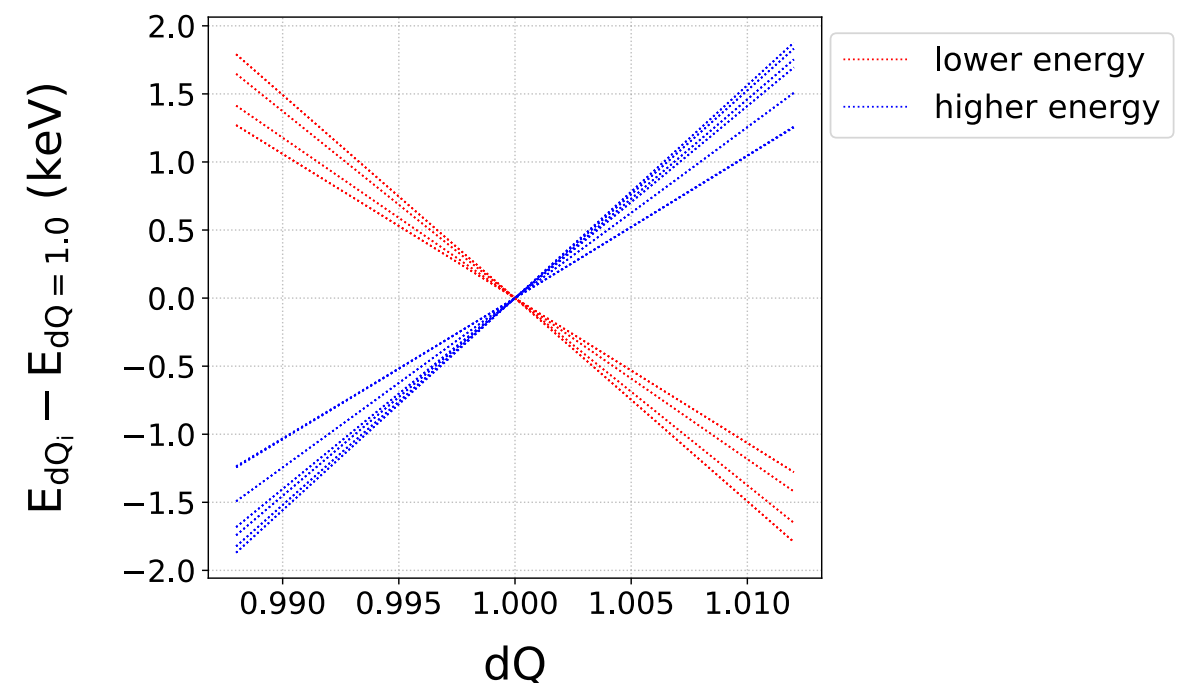
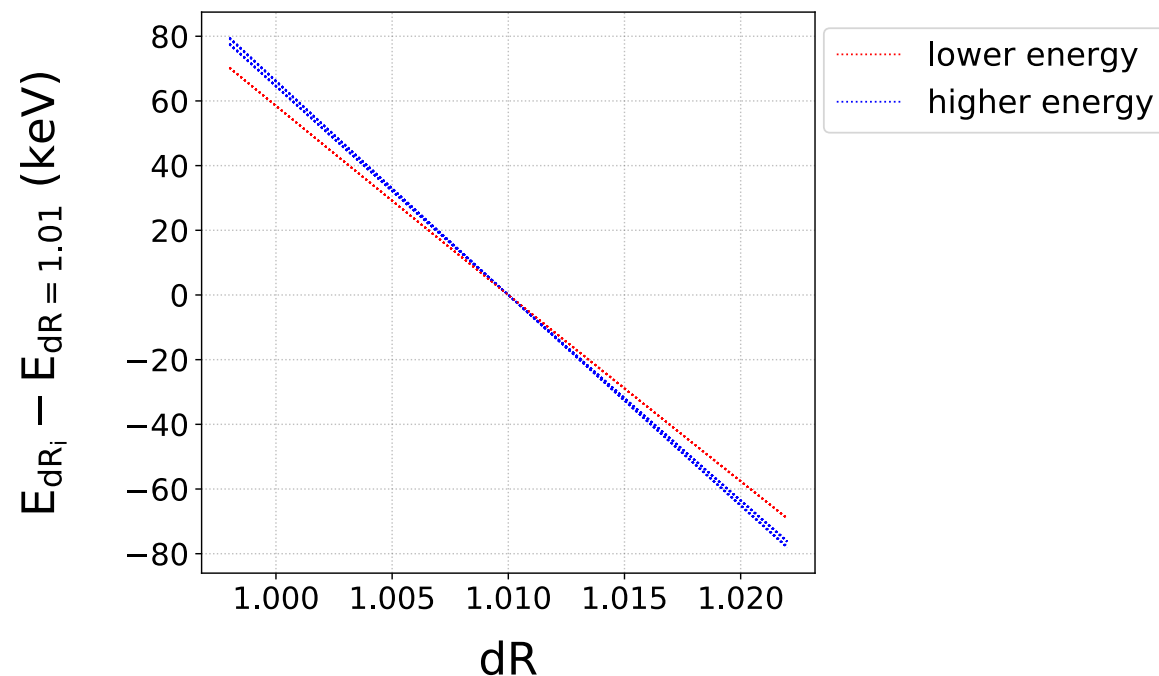
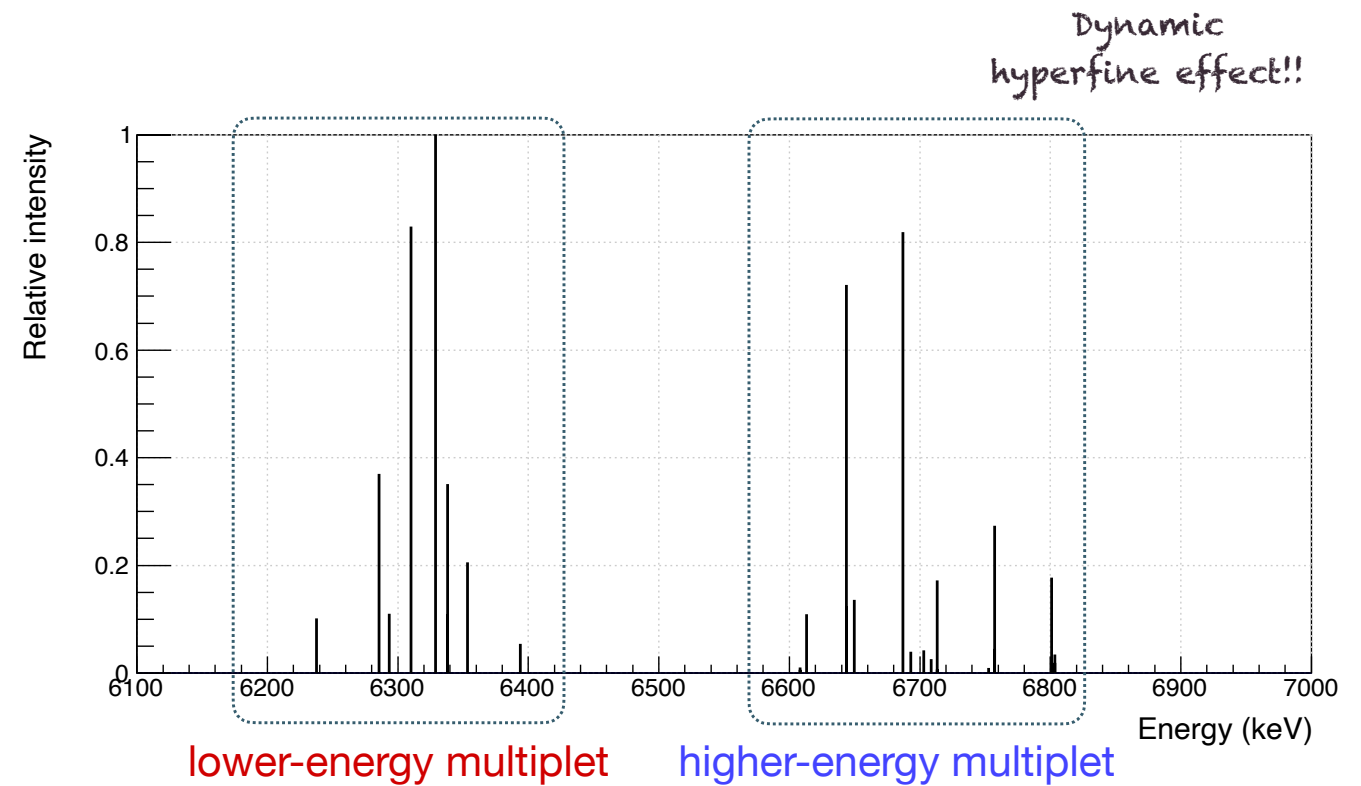
Transition energy dependence on (R_N, Q)

Theoretical calculations are performed by
N. Oreshkina & I. Valuev, MPIK, Heidelberg

$$dR = \frac{R}{R_N}, \text{ where } R_N = 5.8687 \text{ fm}$$

$$dQ = \frac{Q}{Q_N}, \text{ where } Q_N = 12.04 \text{ b}$$

^{248}Cm is even-even $\rightarrow$ This is the **intrinsic Q** calculated as $16\pi B(E2)_{0^+ \rightarrow 2^+} / 5e^2]^{1/2}$, integrated in all possible orientations $\rightarrow$ spherical density (is that right?)



^{248}Cm 2p1s analysis

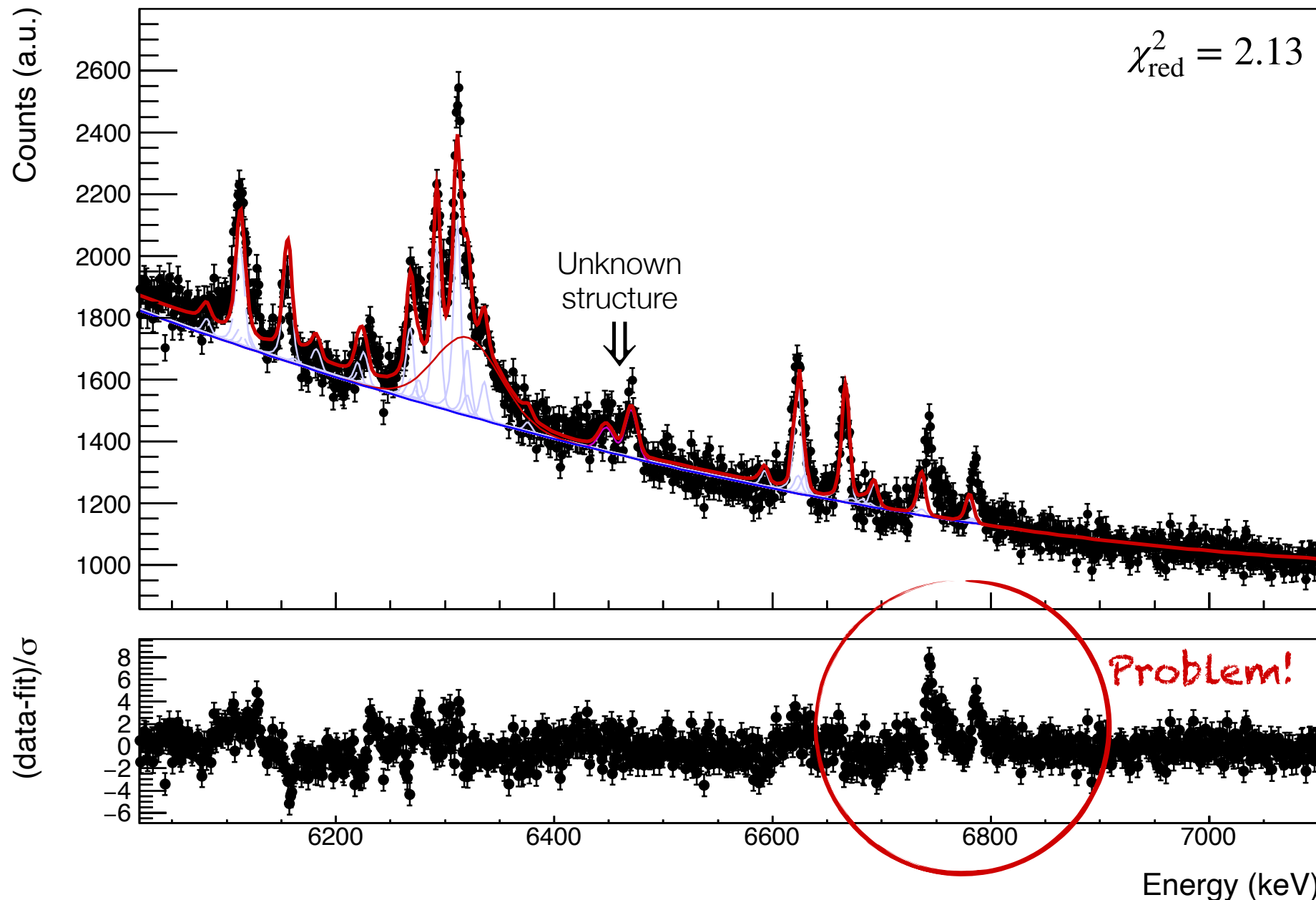
Nuclear charge radius

$$R = 5.9455(1)_{\text{stat}}(117)_{\text{sys}} \text{ fm}$$

Statistical result

$$Q = 12.003(8)_{\text{stat}}(361)_{\text{sys}} \text{ b}$$

My PhD thesis



Included corrections:

- Finite size (2-Fermi)
- Vacuum polarisation
- Nuclear polarisation

Not included corrections:

- Magnetic moment ($\sim$ sub-keV)
- Electron screening?
- Hadronic vacuum polarization
- Recoil effect

Sources of systematic uncertainties

Combined results

ed results

		Charge radius		Quadrupole moment	
Systematic effect	Description	$\Delta(\text{dR})$	σ_{dR}	$\Delta(\text{dQ})$	σ_{dQ}
Fitting features	Instrumental line-shape	0	0.000 032	0	0.001 4
	SE/FE ratio	0	0.000 015	0	0.001 3
	Binning	0	0.000 008	0	0.000 8
	Fitting energy range	0	0.000 033	0	0.001 3
	Background model	0	0.000 042	0	0.004 0
	Free intensities fit	0	0.000 028	0	0.000 4
	Combined	0	0.000 070	0	0.004 7
Energy calibration	Wrong energy of ^{16}N line	0.000 12	0.000 04	-0.000 82	0.000 55
	Energy calibration scheme	0	0.000 007	0	0.000 67
	Uncertainty of literature energy	0	0.000 018	0	0.002 25
	Line-shape for energy calibration	-0.000 068	0.000 038	0	0
	Combined	0.000 052	0.000 058	-0.000 82	0.002 4
Theory	Uncertainty of nuclear polarization correction	0	0.000 20	0	0.000 11
Charge distribution model	Change of the skin thickness parameter	0	0.002 2	0	0.001 91
Discrepancies of spectrum and fit	Free Gaussian fits	0	0.000 95	0	0.028 2
Total		0.000 052	0.002 4	-0.000 82	0.028 8

Combined result

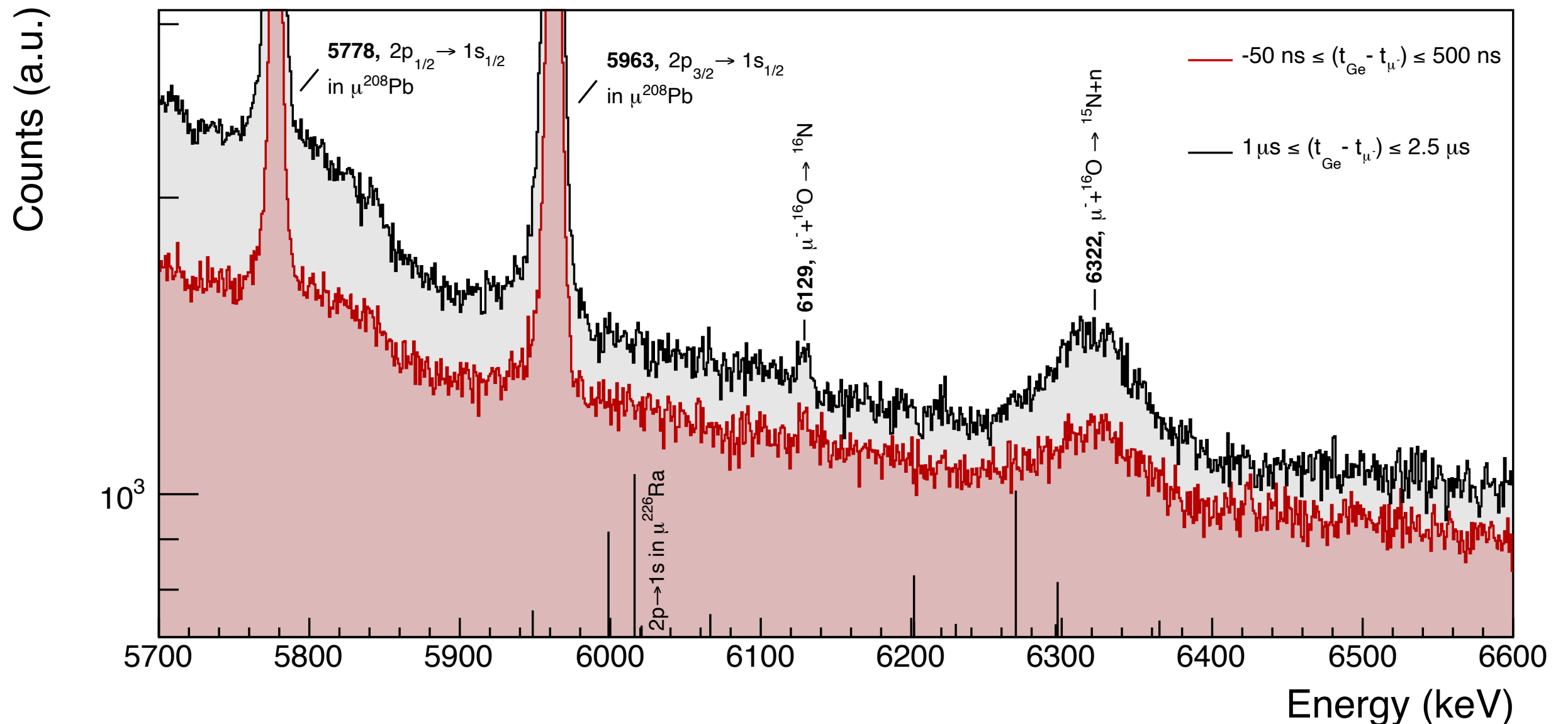
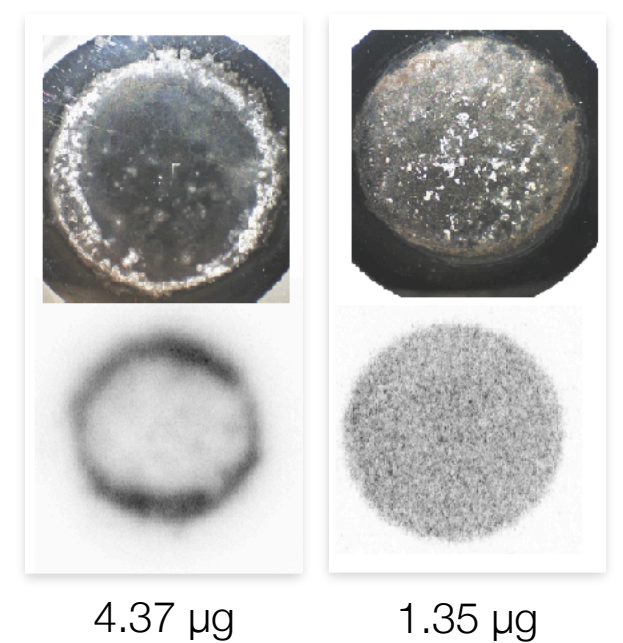
$$R = 5.9455(1)_{\text{stat}}(117)_{\text{sys}} \text{ fm}$$

$$Q = 12.003(8)_{\text{stat}}(361)_{\text{sys}} \text{ b}$$

Needing further theory input

^{226}Ra spectrum in the 2p1s region

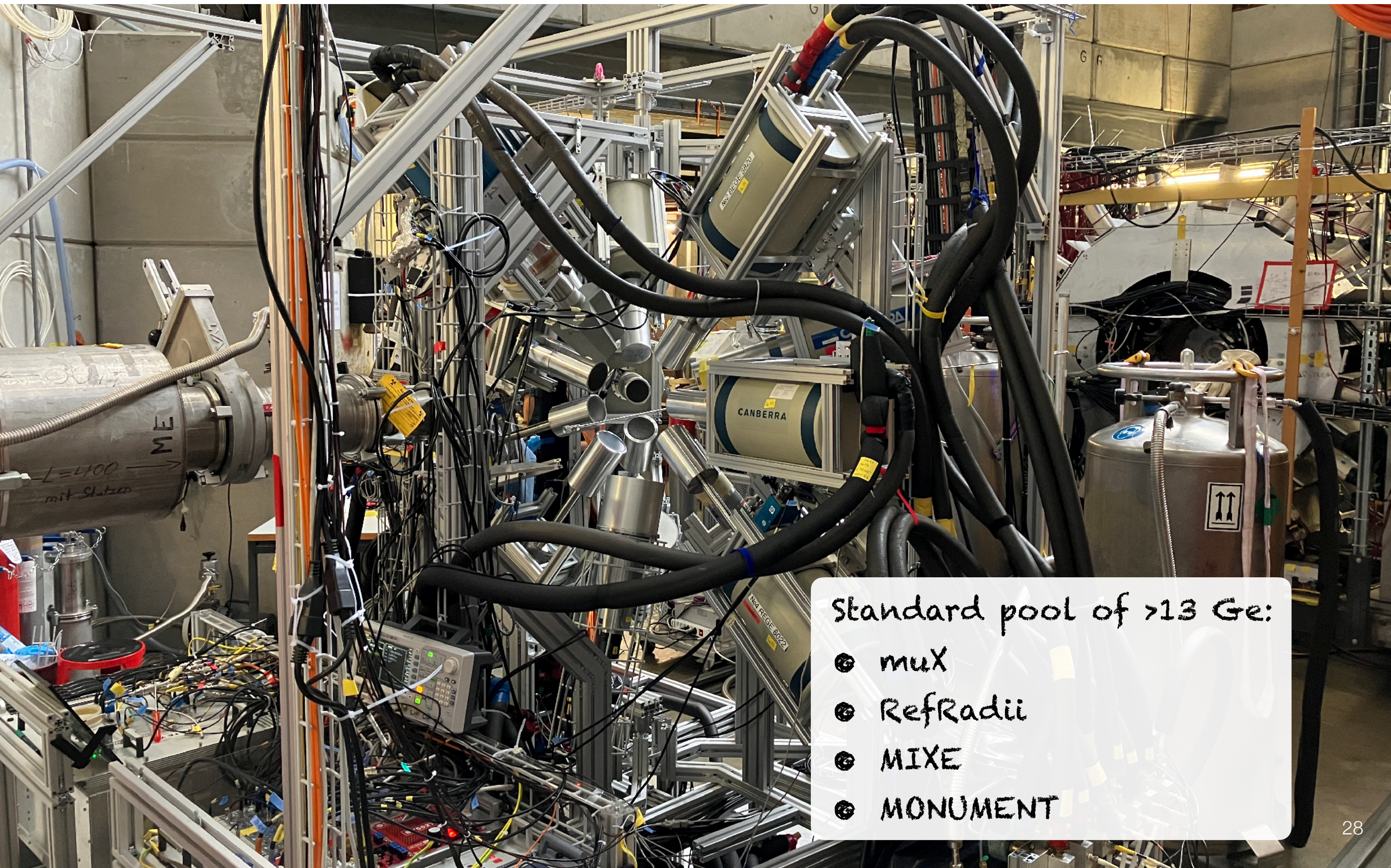
No 2p1s structure is observed



Experimental timeline

Sci Rep 15, 6939 (2025)

2022 — *Assessing the target fabrication techniques*

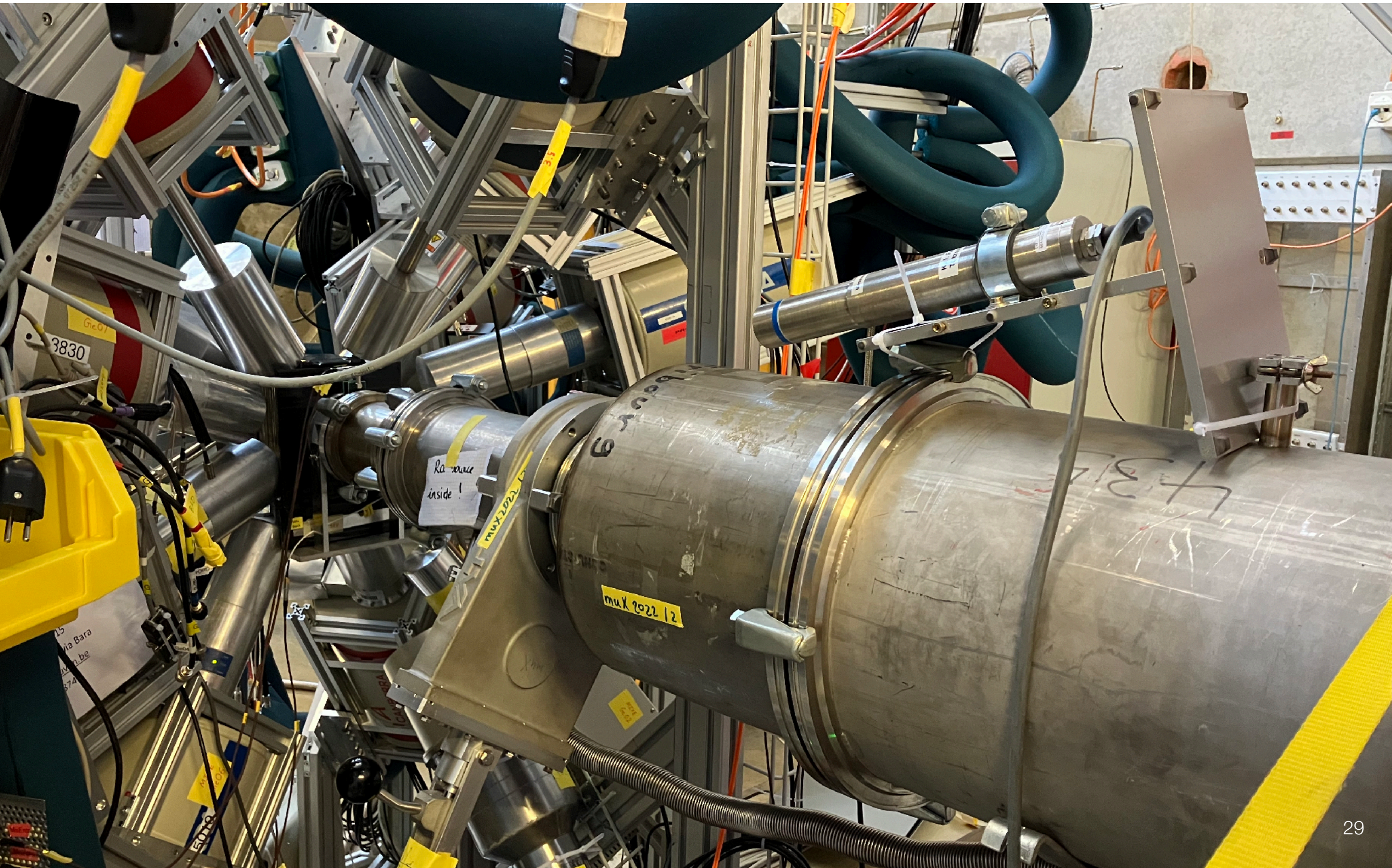


Standard pool of >13 Ge:

- muX
- RefRadii
- MIXE
- MONUMENT

Experimental timeline

2023 — $^{35,37}\text{Cl}$ & $^{39,40,41}\text{K}$ & $^{107,109}\text{Ag}$ ($\mathcal{O}(\text{mg to g})$ & $7.5\ \mu\text{g}$)

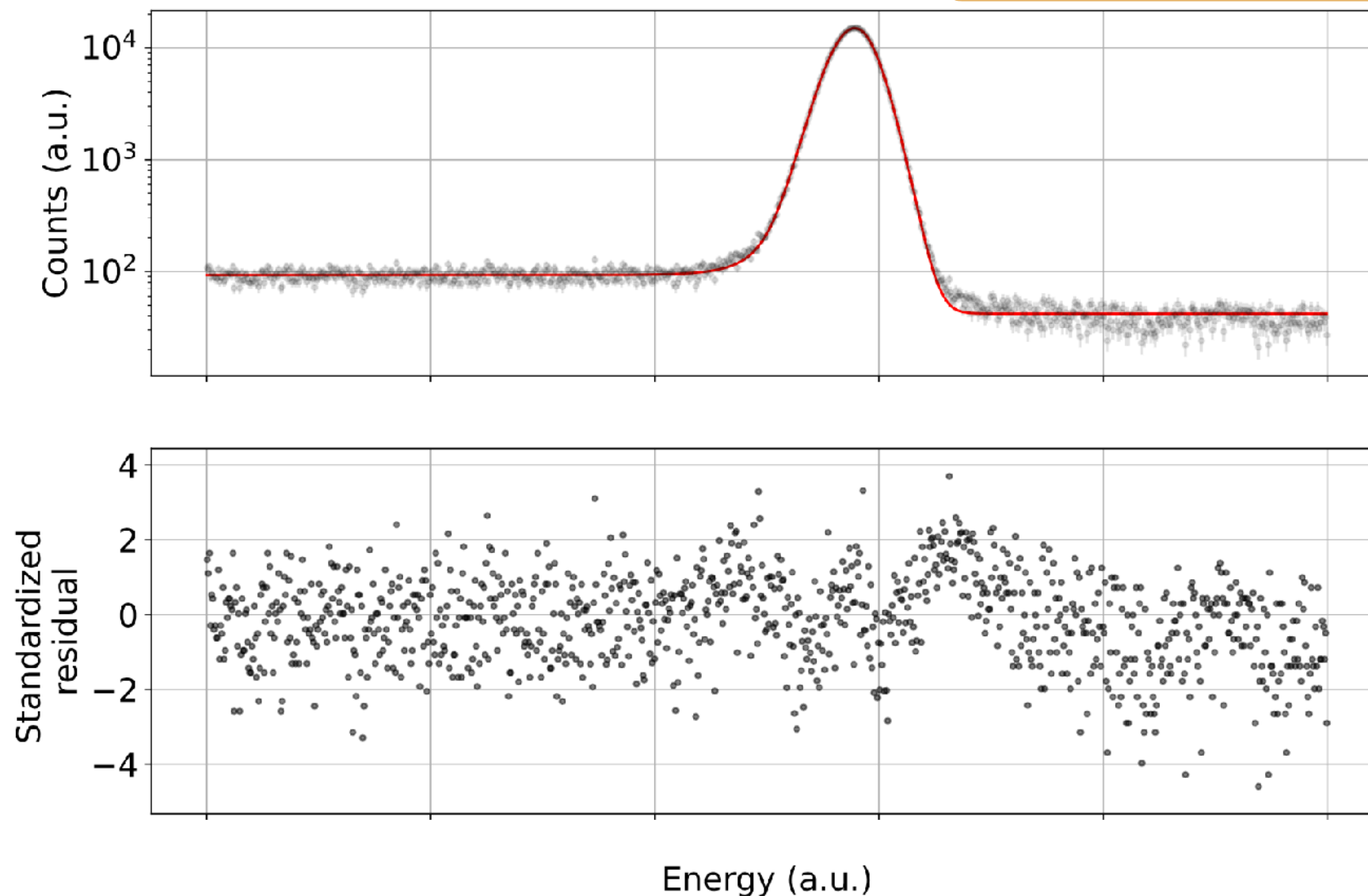


^{35}Cl 2p1s analysis (200 mg)

Nuclear charge radius

Publication in progress;
 $^{35,37}\text{Cl}$, $^{39,41}\text{K}$

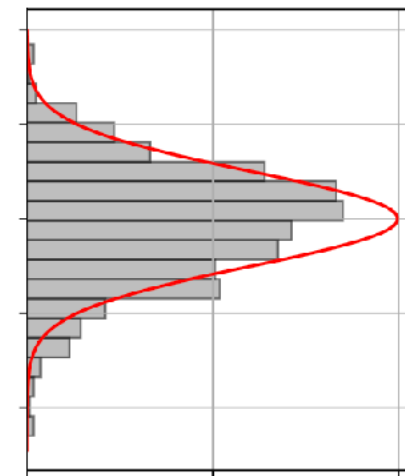
Analysis of data and image
courtesy of **Michael Heines**



Bunch of corrections for quality
data analysis pipeline:

- Thorough line-shape investigation
- Gain stabilisation
- Time algorithm correction
- Rise-time correction
- Combined detector analysis
-

Extensive analysis note written!

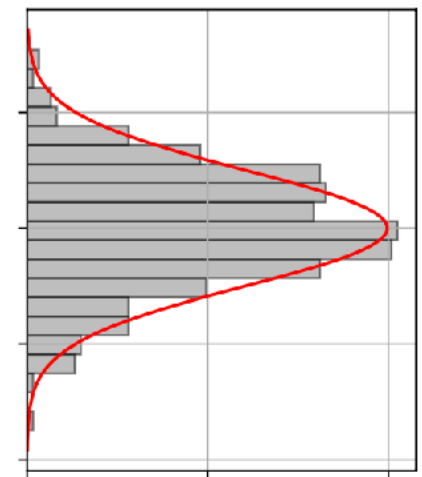
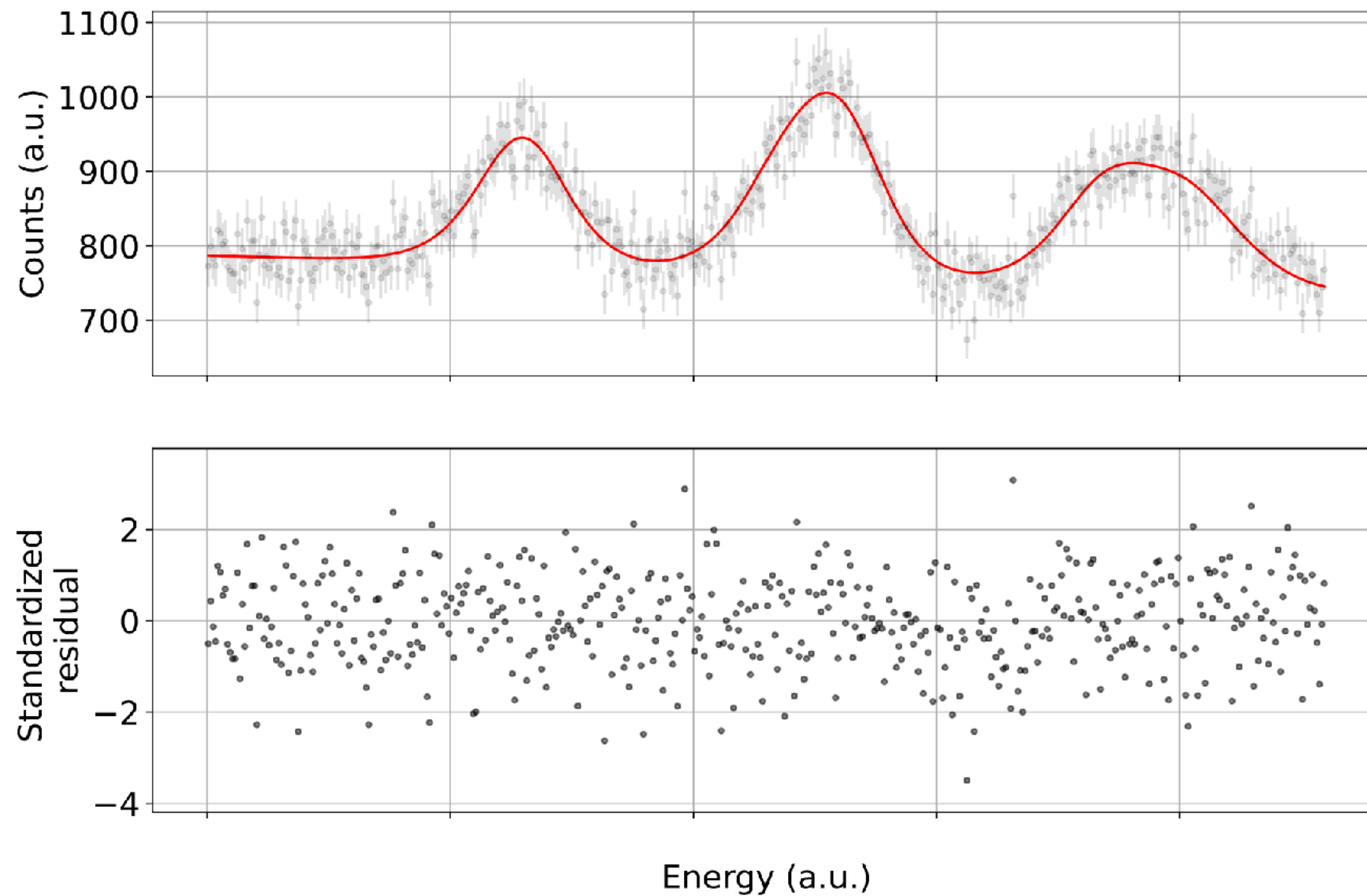


^{40}K 2p1s analysis (7.5 μg)

Nuclear charge radius

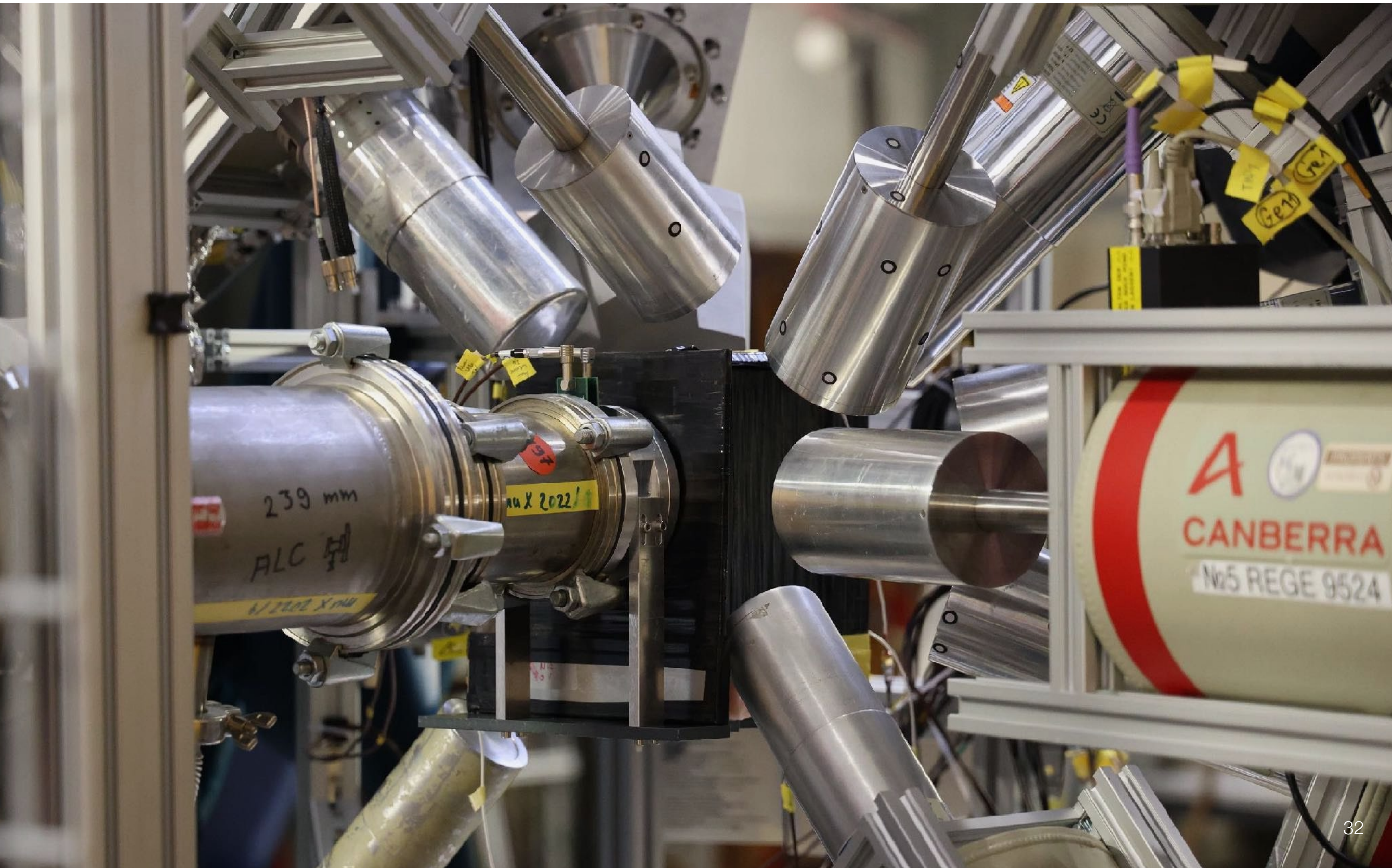
$$E_{2p-1s} = 712.\text{xxx}(\text{xx}) \text{ keV}$$

Analysis of data and image
courtesy of **Michael Heines**



Experimental timeline

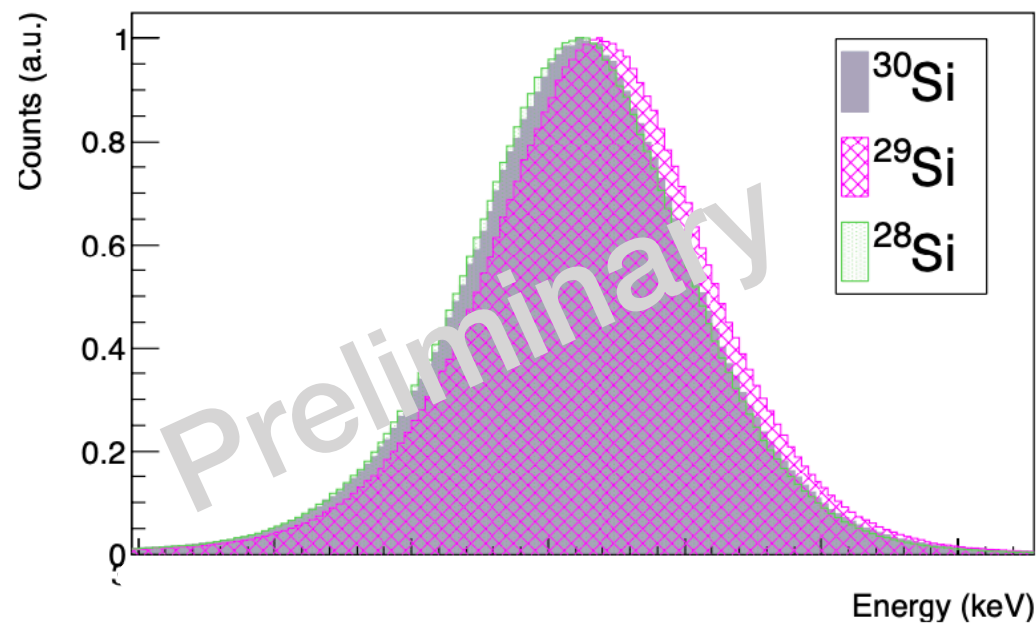
2024 — ^{27}Al , $^{28,29,30}\text{Si}$, ^{139}La , $^{175,176}\text{Lu}$ — $\mathcal{O}(\text{mg to g})$



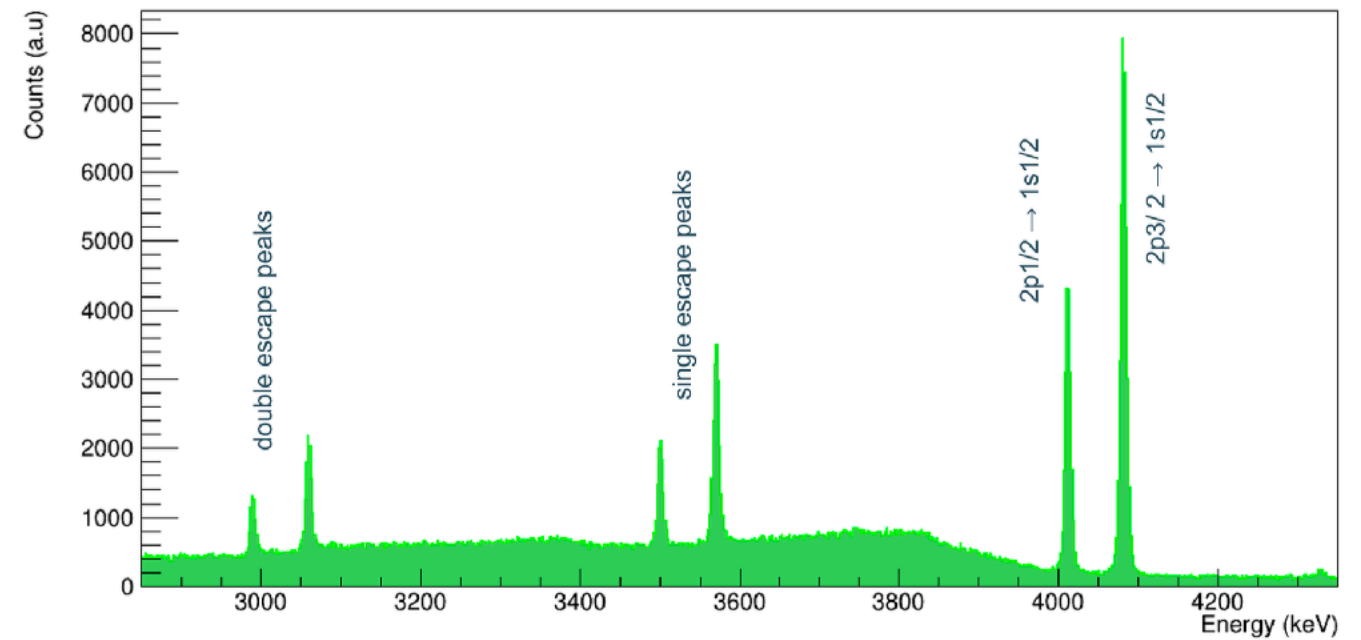
2024 data acquired

Analysis of data by Marie Deseyn & War War Myint Myat Phyo

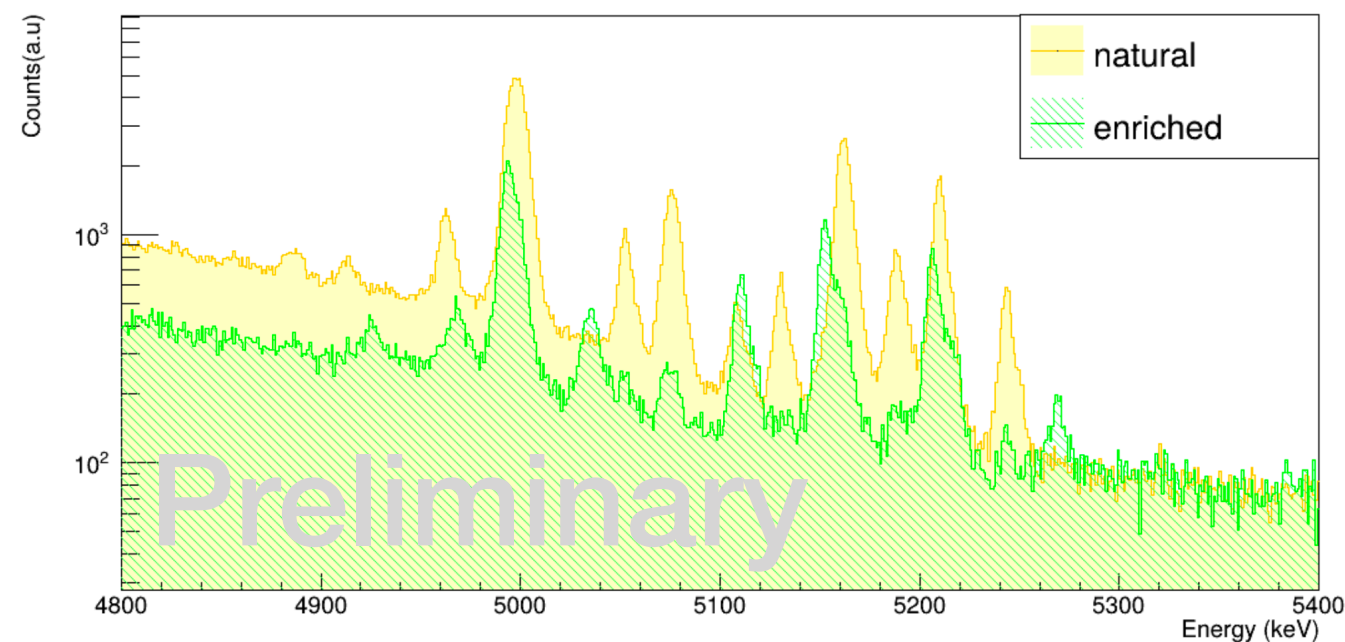
Towards nuclear charge radii extraction



Muonic 2p1s in $^{28,29,30}\text{Si}$



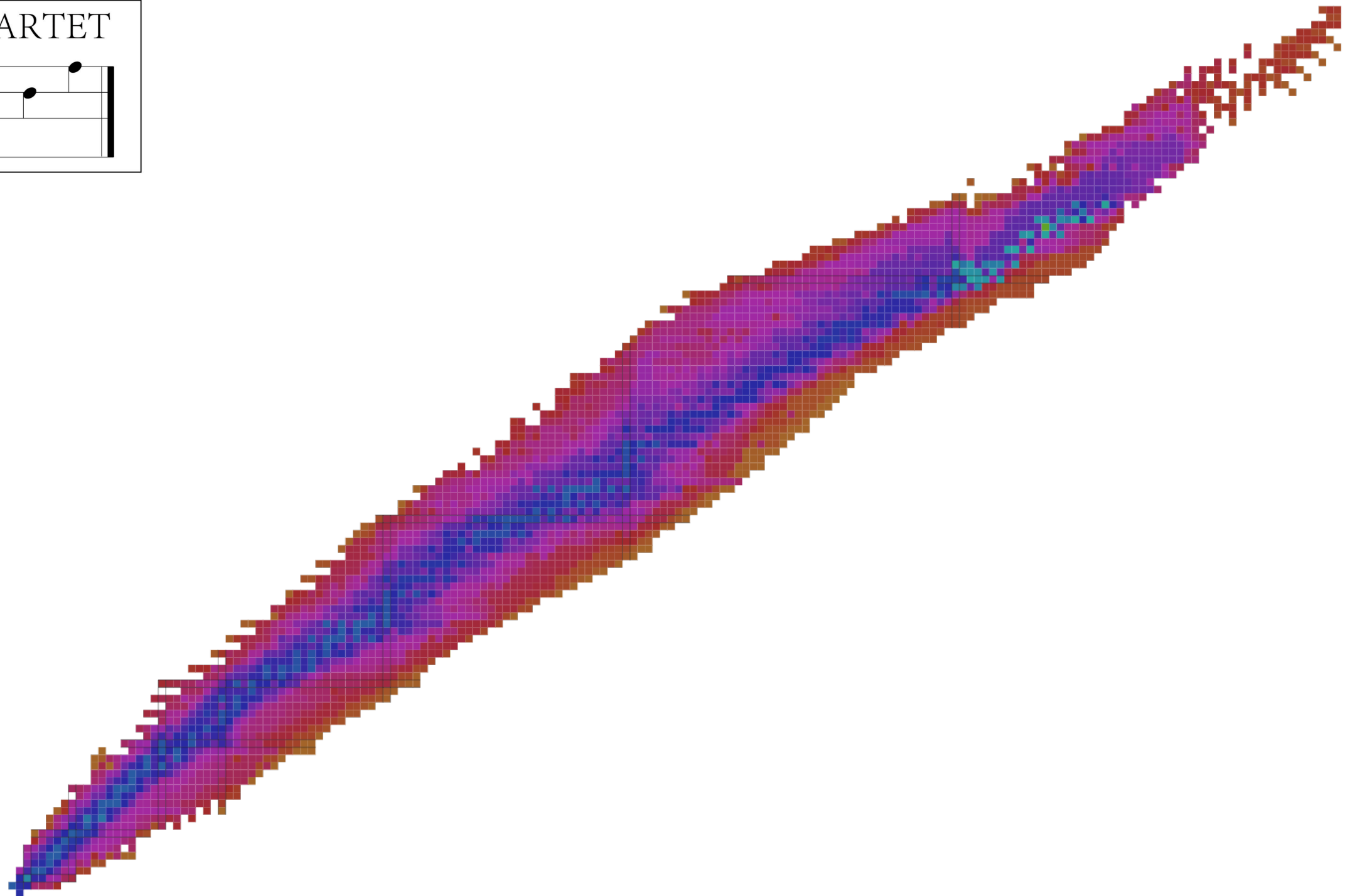
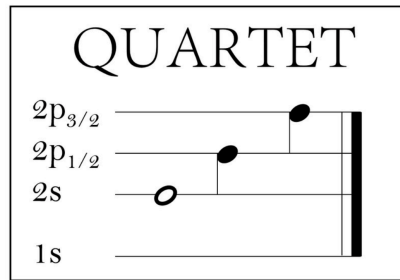
Muonic 2p1s in ^{139}La



Muonic 2p1s in $^{175,176}\text{Lu}$

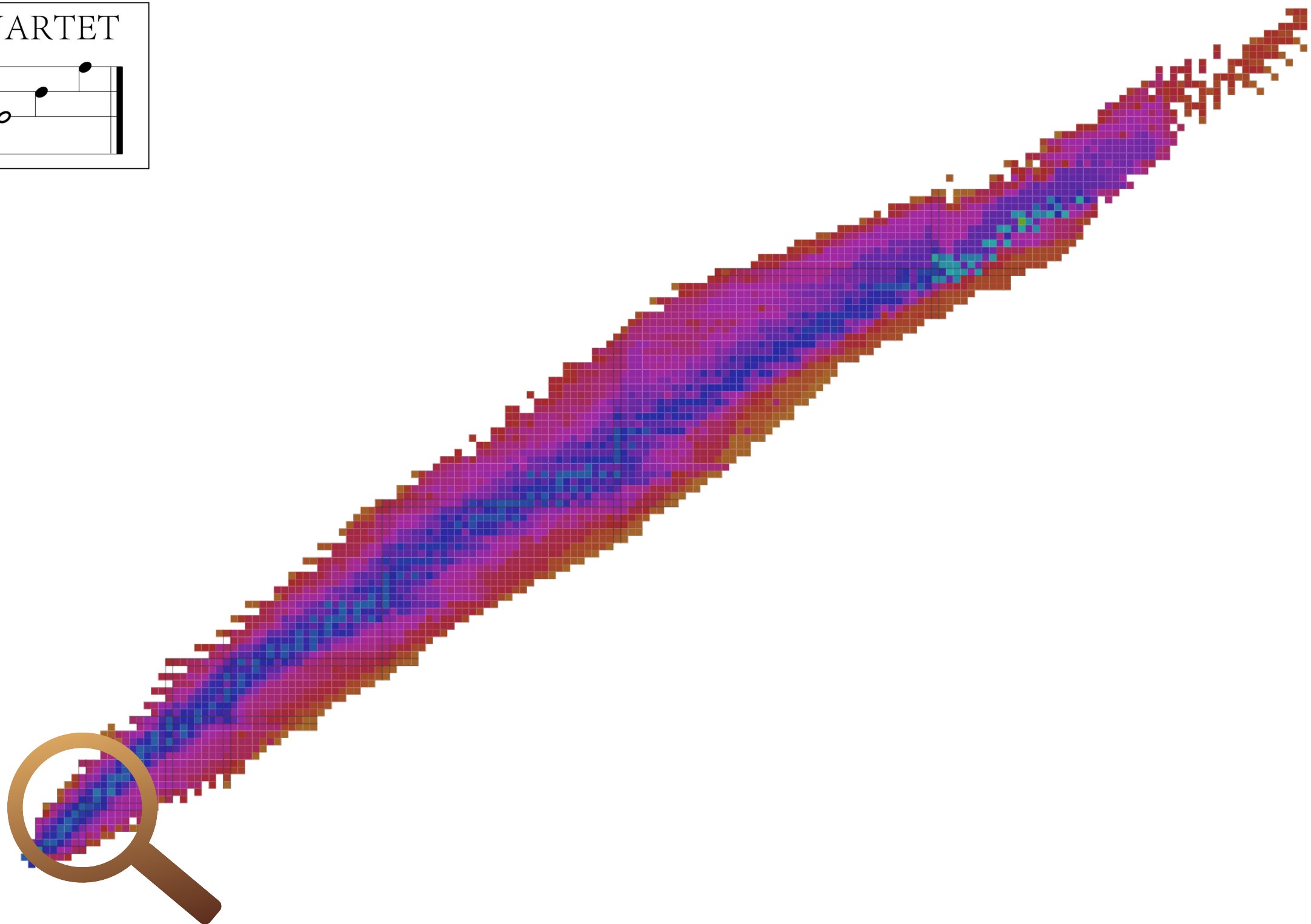
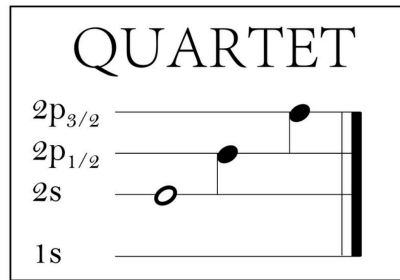
How to tackle the very low-mass regime

QUAntum inte**Rac**Tions with **E**xotic a**T**oms (QUARTET)



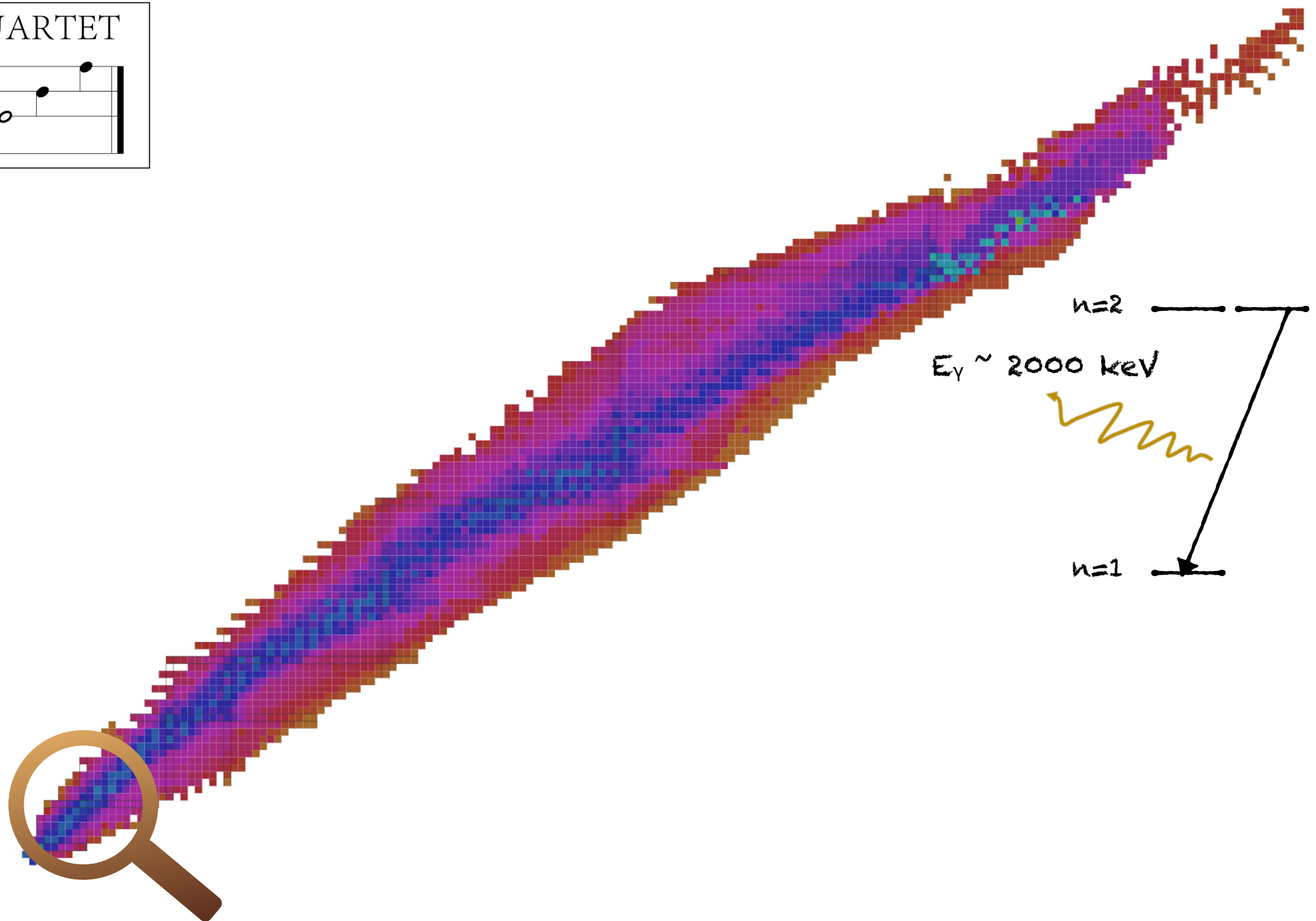
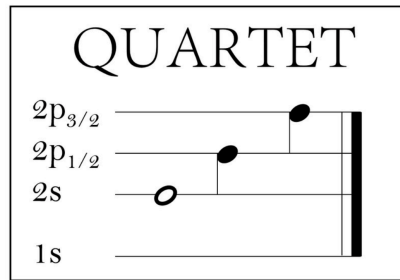
How to tackle the very low-mass regime

QUAntum inte**Rac**Tions with **E**xotic a**T**oms (QUARTET)



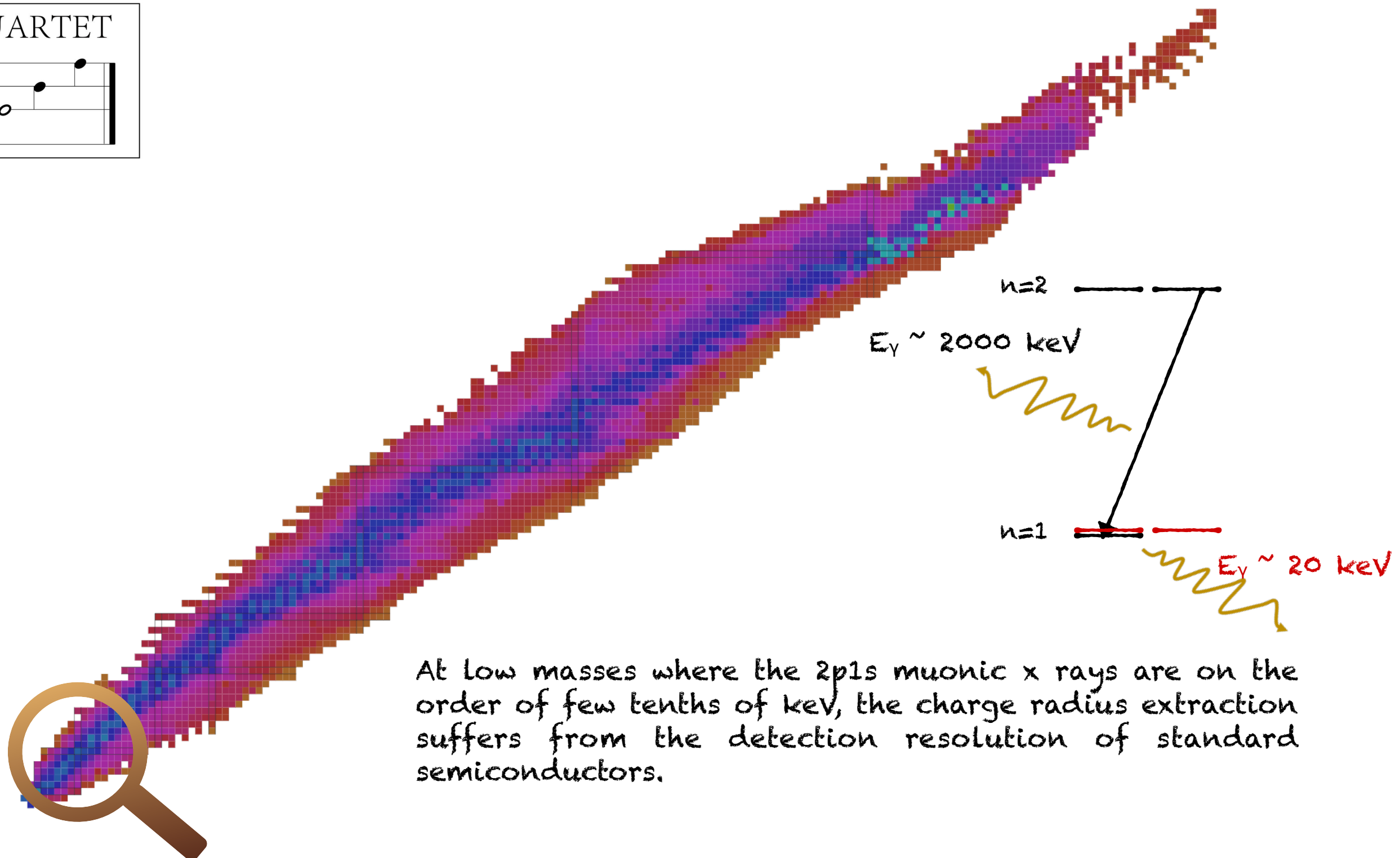
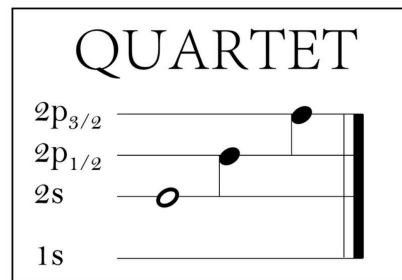
How to tackle the very low-mass regime

QUAntum inte**Rac**Tions with **E**xotic a**T**oms (QUARTET)



How to tackle the very low-mass regime

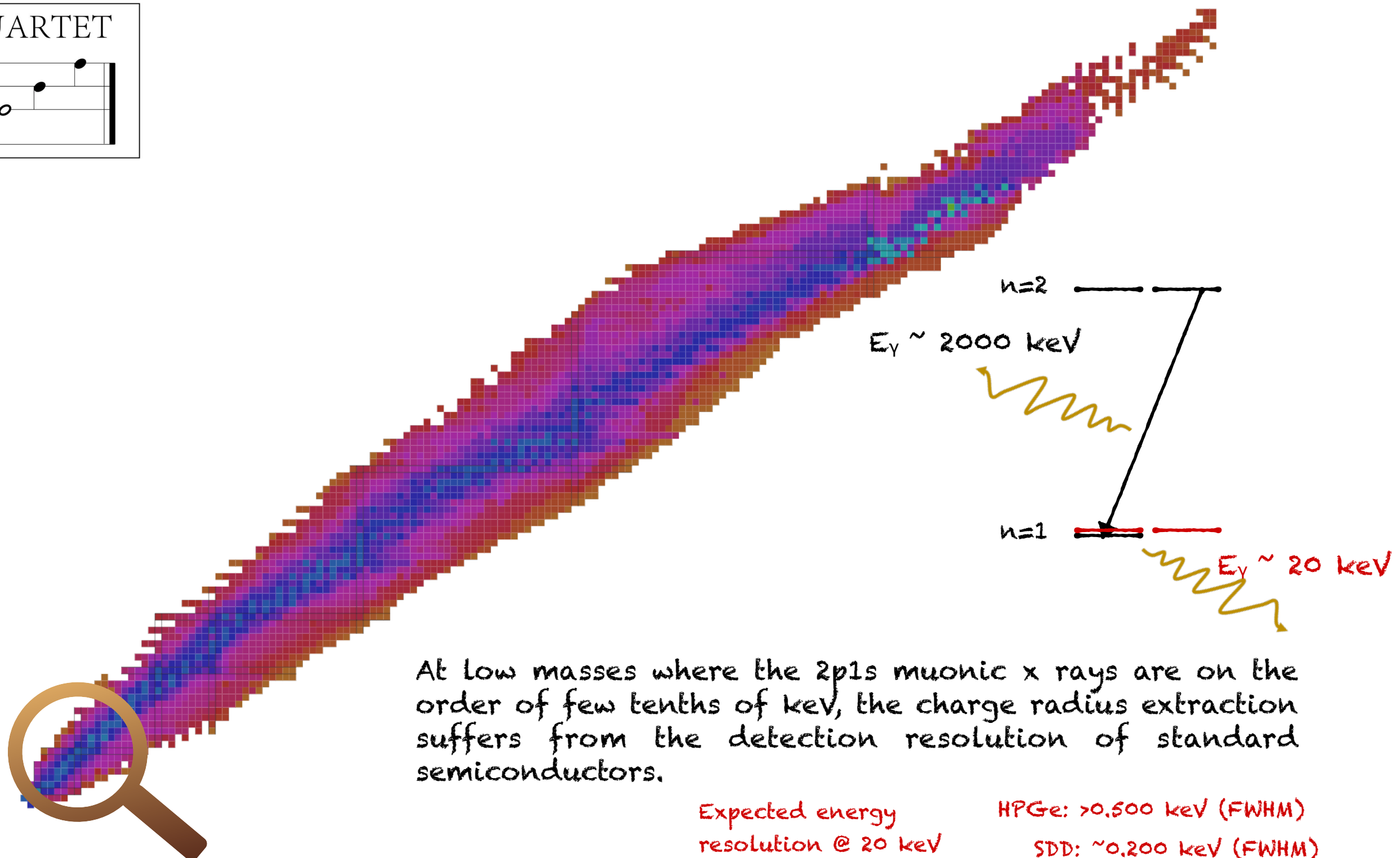
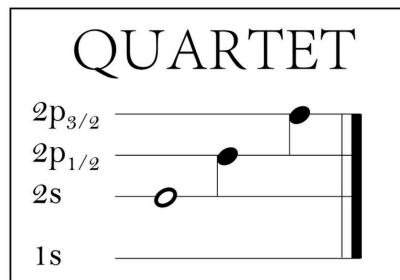
QUAntum inte**R**ac**T**ions with **E**xotic a**T**oms (QUARTET)



At low masses where the $2p1s$ muonic x rays are on the order of few tenths of keV, the charge radius extraction suffers from the detection resolution of standard semiconductors.

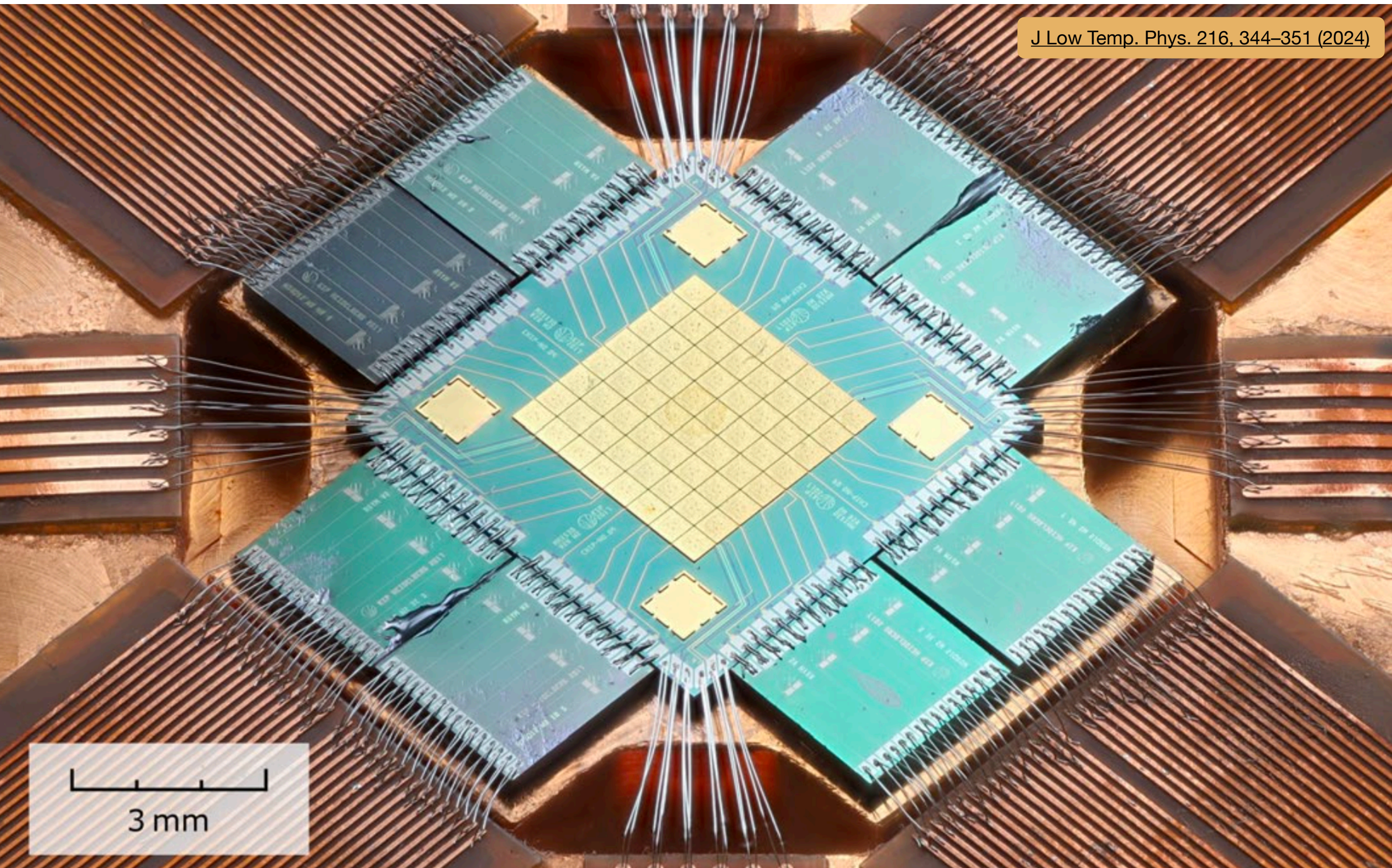
How to tackle the very low-mass regime

QUAntum inte**Rac**Tions with **E**xotic a**T**oms (QUARTET)



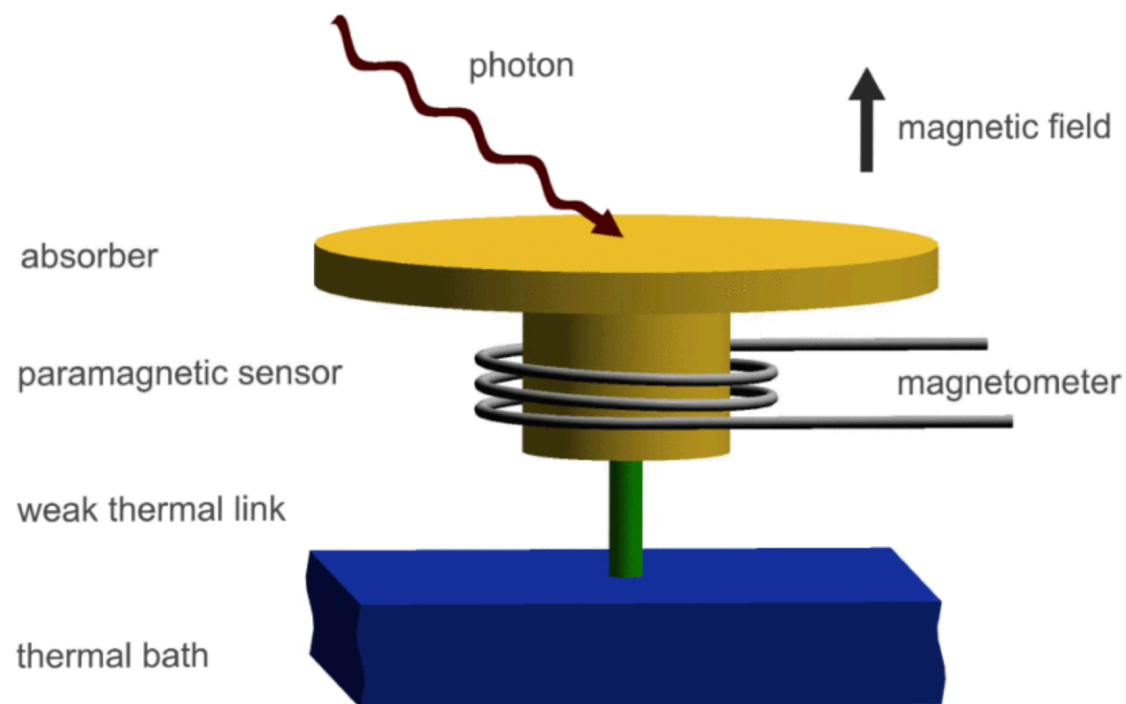
Probing muonic x-ray energies with high precision with metallic magnetic calorimeters (MMC)

J Low Temp. Phys. 216, 344–351 (2024)



Working principle of MMC

Achieving a high energy resolution



IEEE Transactions on Applied
Superconductivity 19, 63 (2009)

Expected energy
resolution @ 20 keV

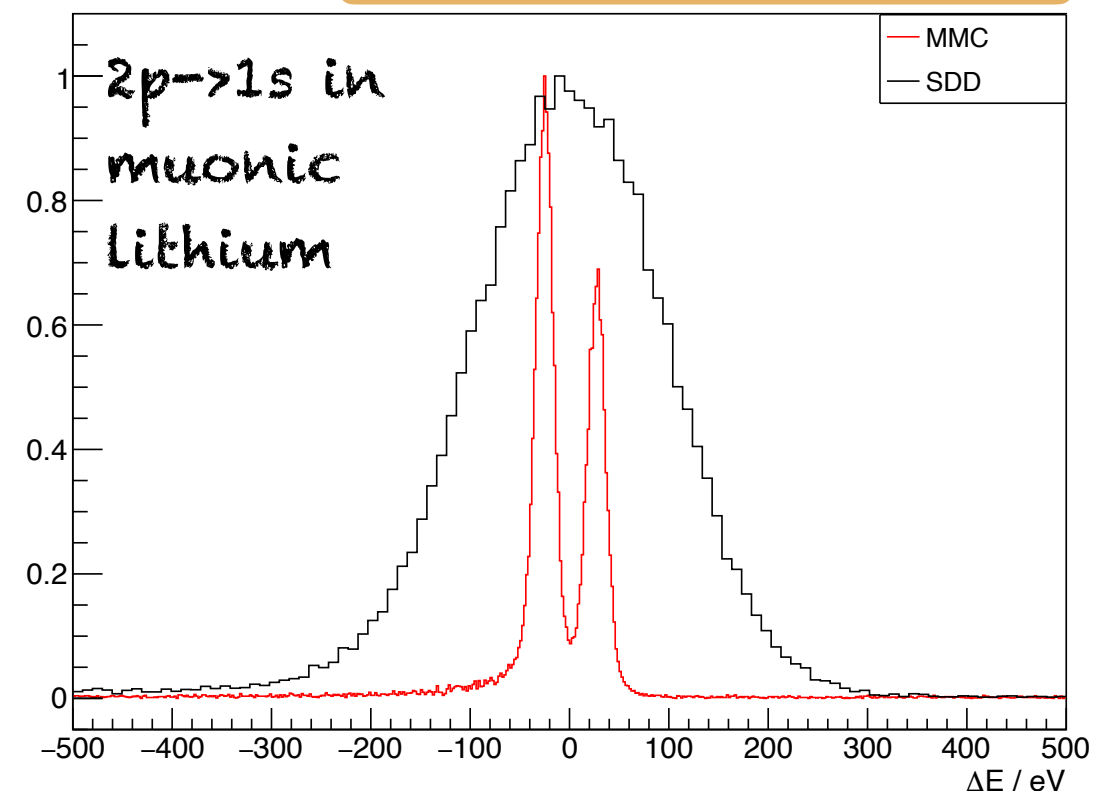
HPGe: >0.500 keV (FWHM)

SDD: ~0.200 keV (FWHM)

MMC: ~0.010 keV (FWHM)

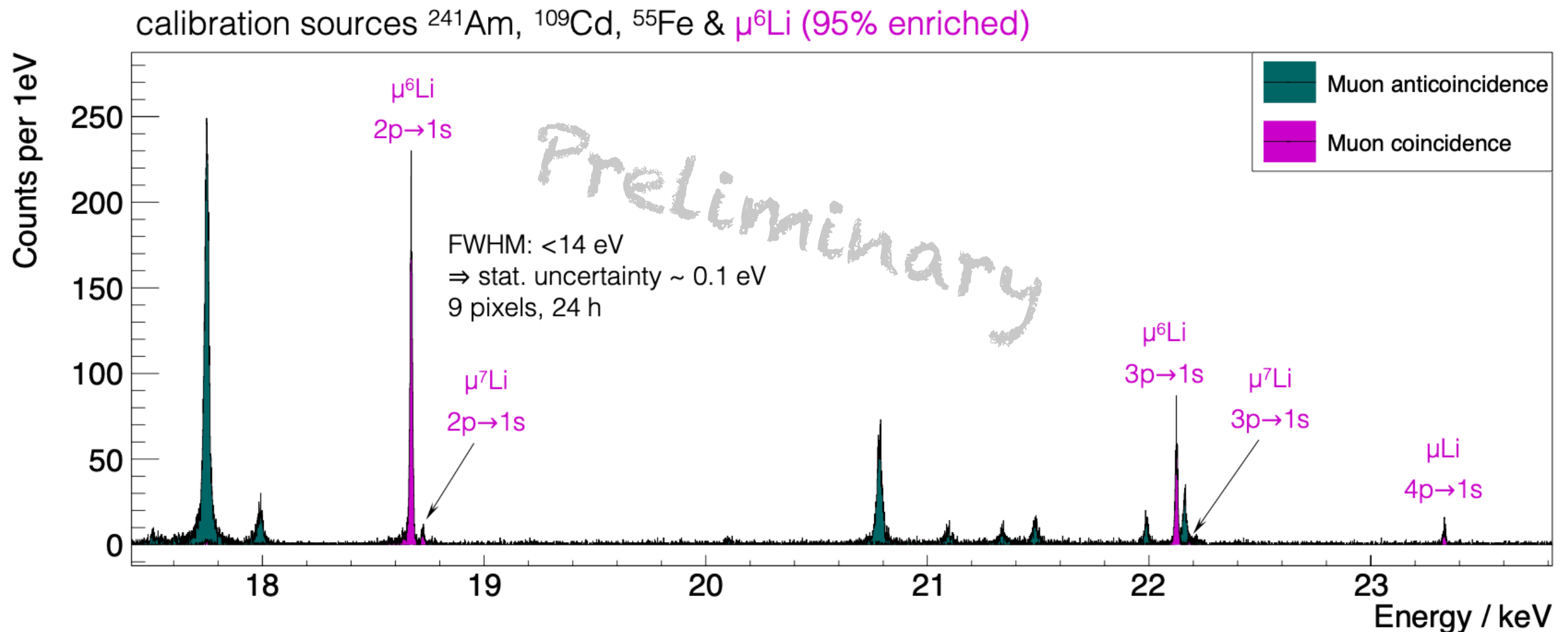
- Tiny pixel $500 \times 500 \mu\text{m}^2$
- Temperature increase of the detector due to photon's absorption
- Decrease of sensor's magnetization, read by magnetometer (SQUID)
- Connected to thermal bath, the detector operates in a cryostat.

Image courtesy of Katharina von Schoeler



Measurement demonstration of lithium *Energy spectrum*

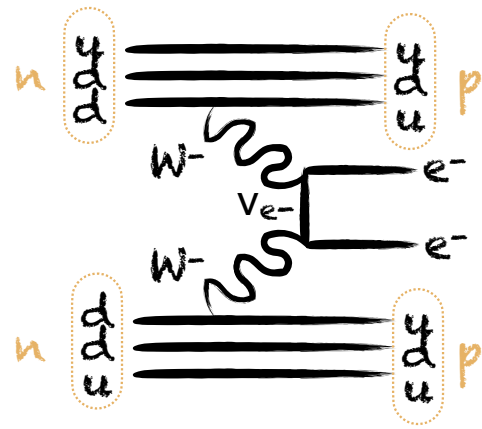
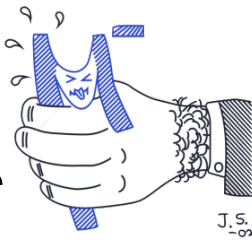
Analysis of data and image courtesy
of **Katharina von Schoeler**



More elements have been measured since then!!

Studying capture rates after muon's nuclear capture

Muon **O**rdinary capture for the **NU**clear **M**atrix elem**ENT**s in $\beta\beta$ decays (MONUMENT)

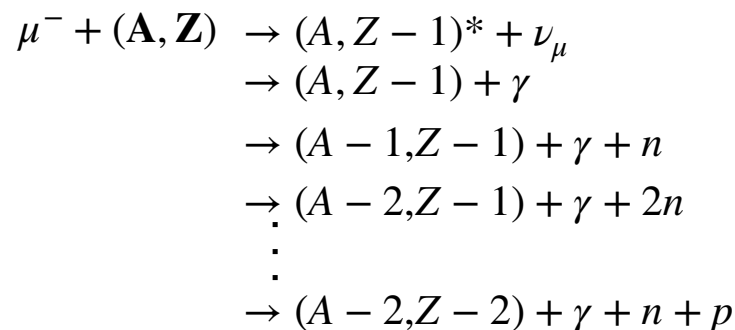
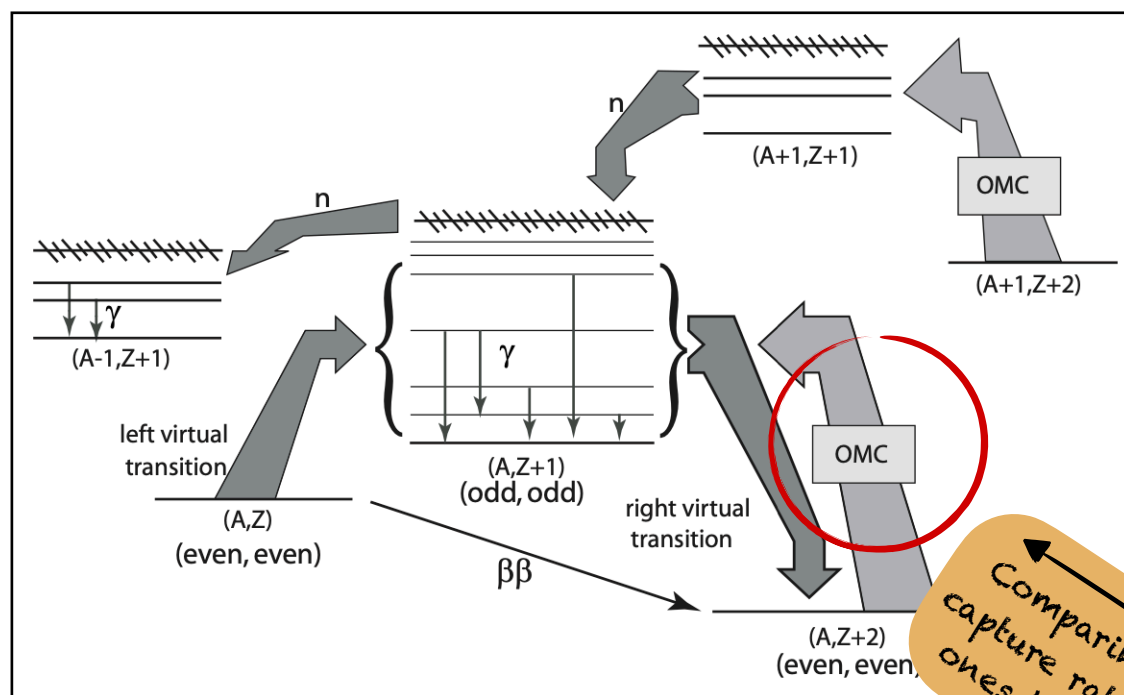


Decay rate of $0\nu\beta\beta$

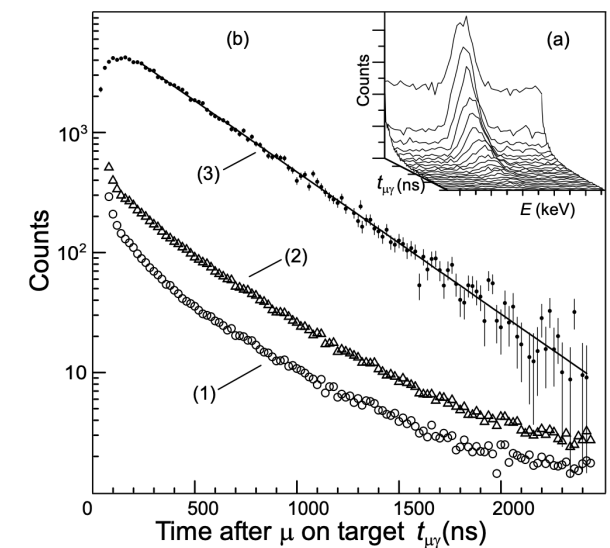
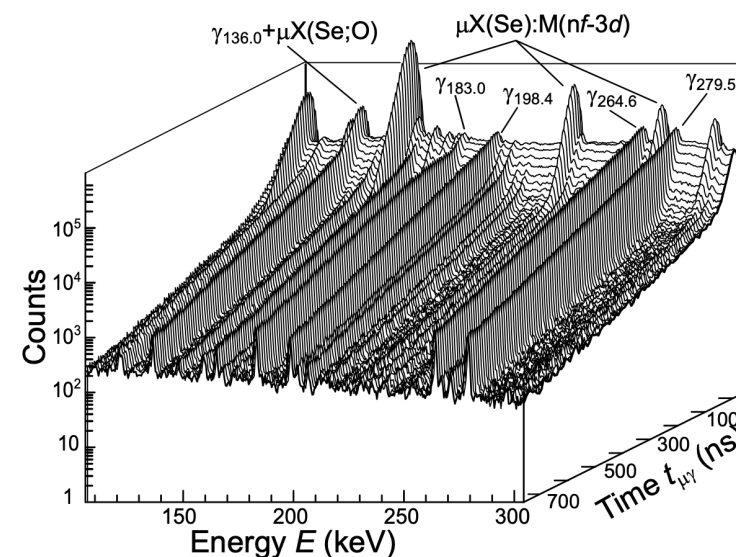
$$\Gamma_{\beta\beta}^{0\nu} = \frac{1}{T_{\beta\beta}^{0\nu}} = G^{0\nu} \left| M^{0\nu} \right|^2 \langle m_{\beta\beta} \rangle^2$$

NME squared
→ calculated with large theoretical uncertainty

Square of effective Majorana mass



Comparing experimental capture rates to calculated ones by the shell model



Phys. Rev. C 99, 024327 (2019)
Eur. Phys. J. C 84, 1188 (2024)
EPL 58, 666 (2002)

Total capture rate

$$\lambda_{\text{total}} = \lambda_{\text{decay}} + \lambda_{\text{absorbed}} = 1/2.2 \mu\text{s} + 1/T_{\text{capture}}$$

Partial capture rate

$$\lambda_{\text{capture-i}} = \lambda_{\text{total}} \gamma_i$$

Absolute γ -ray yield

Studying samples' elemental composition

Muon Induced X-ray Emissions (MIXE)

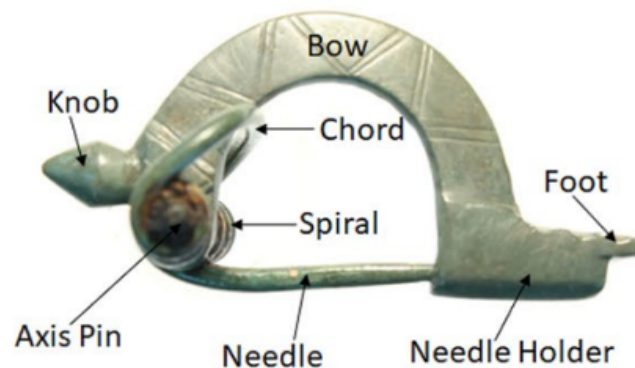
- Determining the elemental composition of different samples (such as archaeological artefacts, meteorites, batteries) from the cascade transition yields of the muonic x rays.
- The muonic energies are characteristic of the element.

Arrowhead made of meteoritic iron



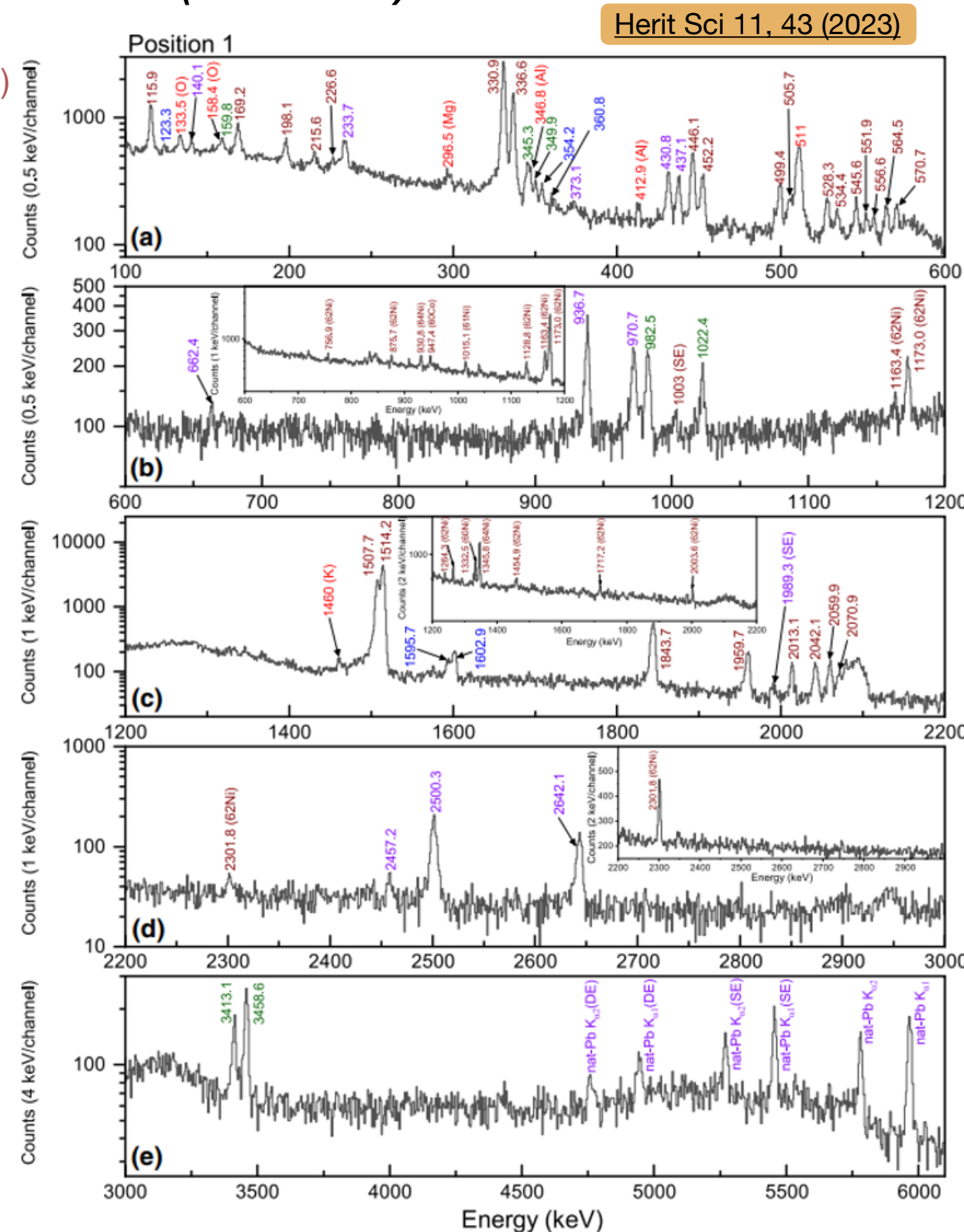
Journal of Archaeological Science 157, 105827 (2023)

Knob bow fibula



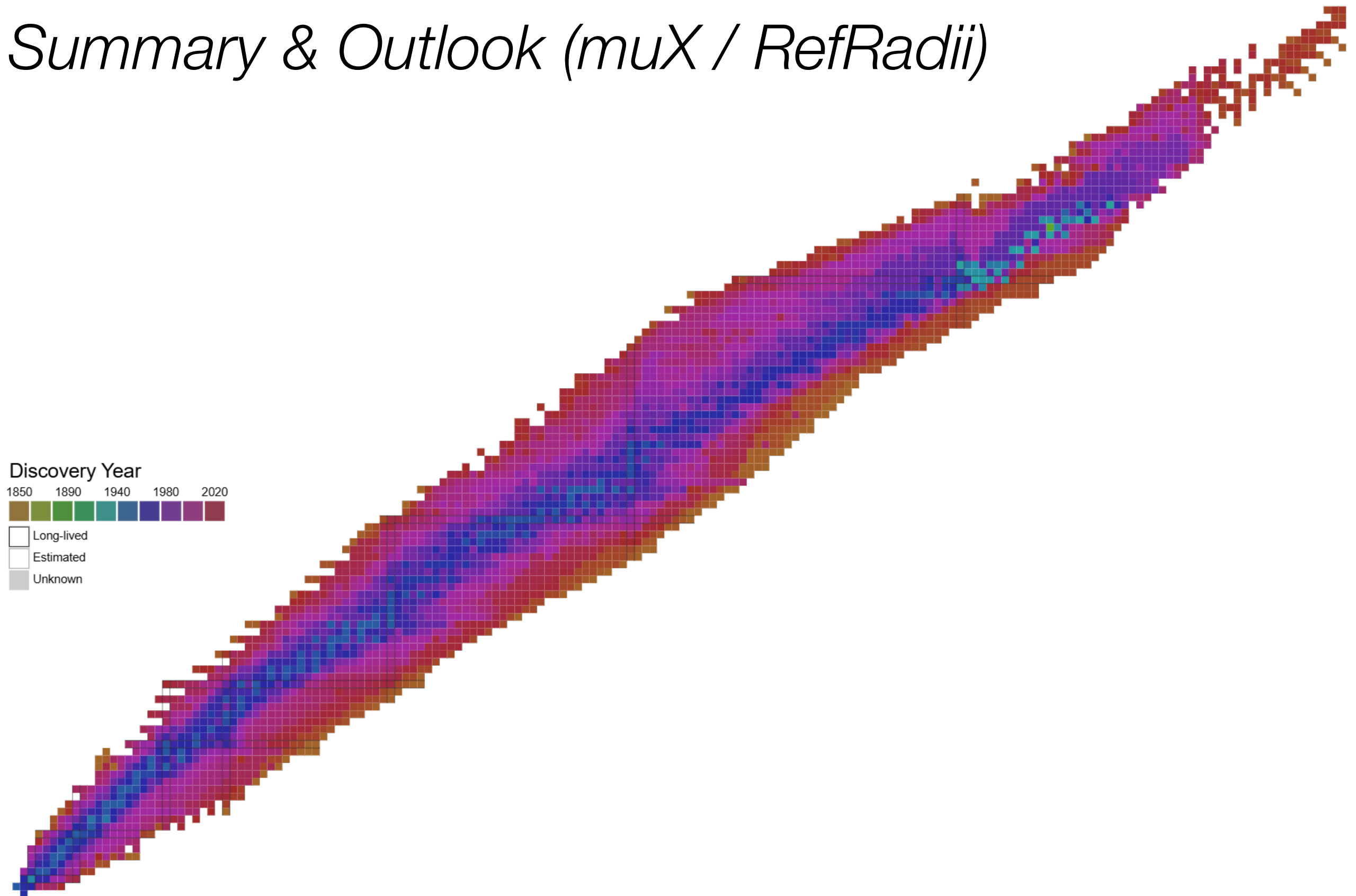
Herit Sci 11, 43 (2023)

Copper (Cu)
Zinc (Zn)
Tin (Sn)
Lead (Pb)



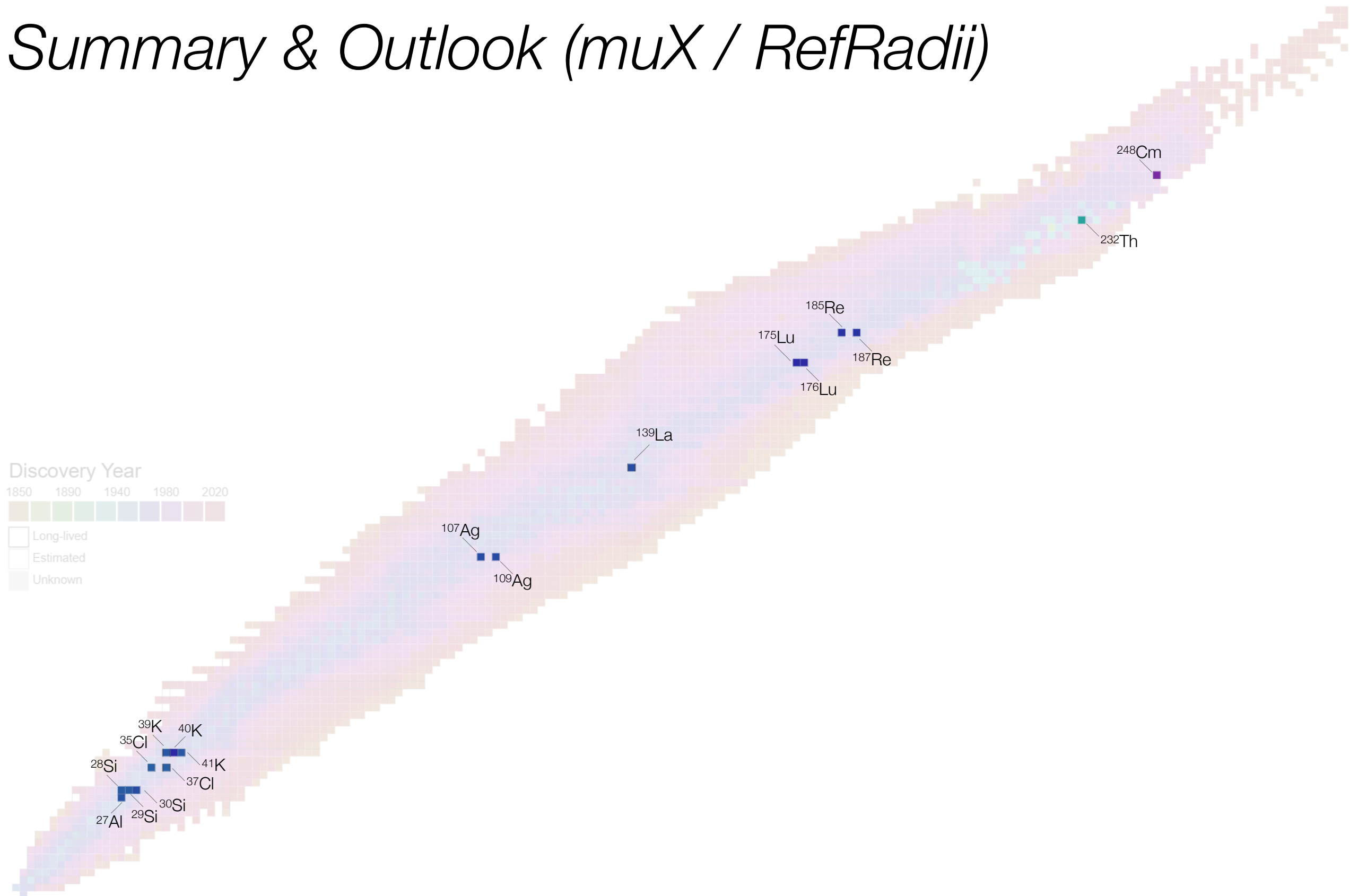
Science case

Summary & Outlook (muX / RefRadii)



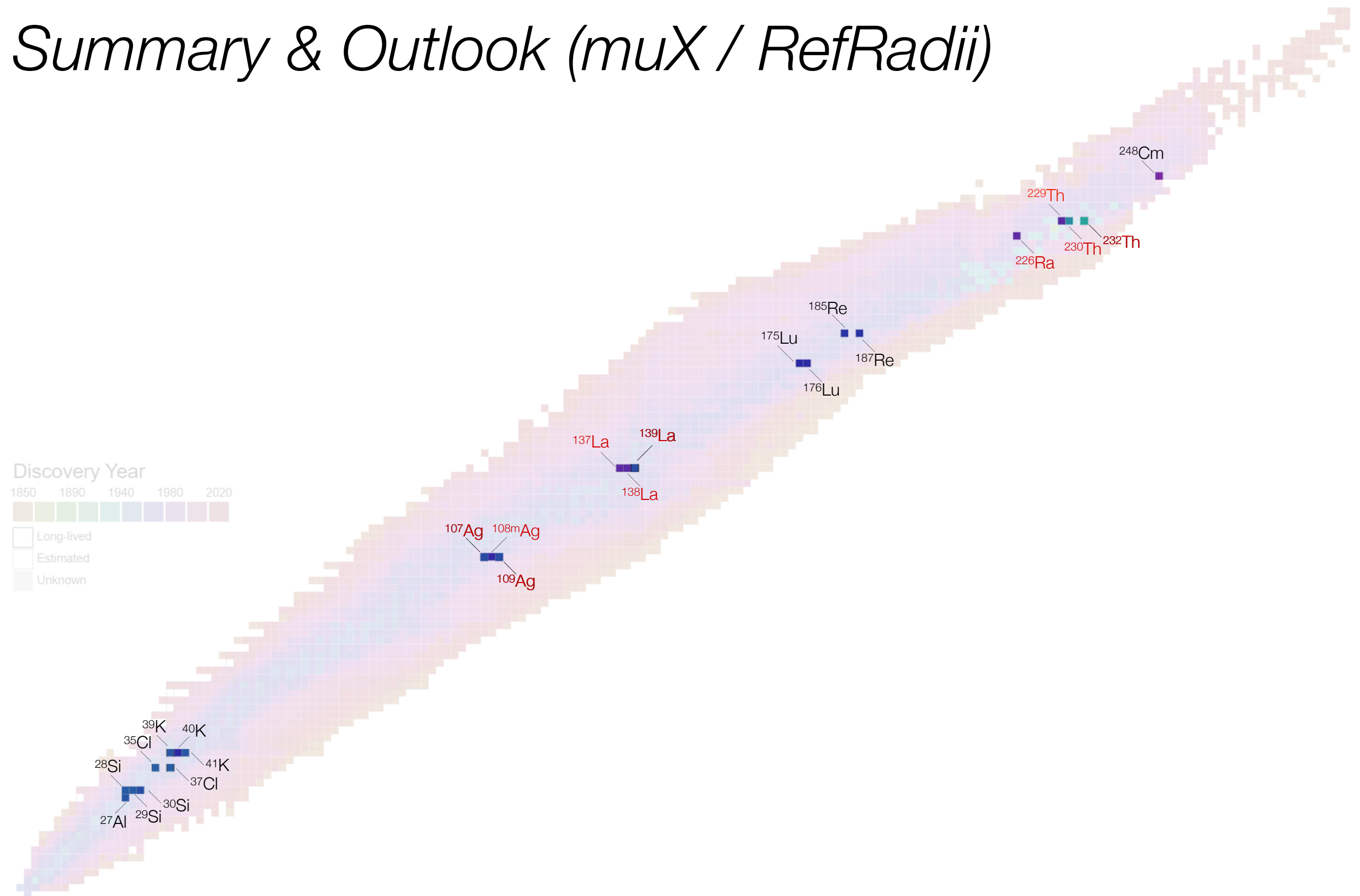
Science case

Summary & Outlook (muX / RefRadii)



Science case

Summary & Outlook (muX / RefRadii)



Behind the story told...

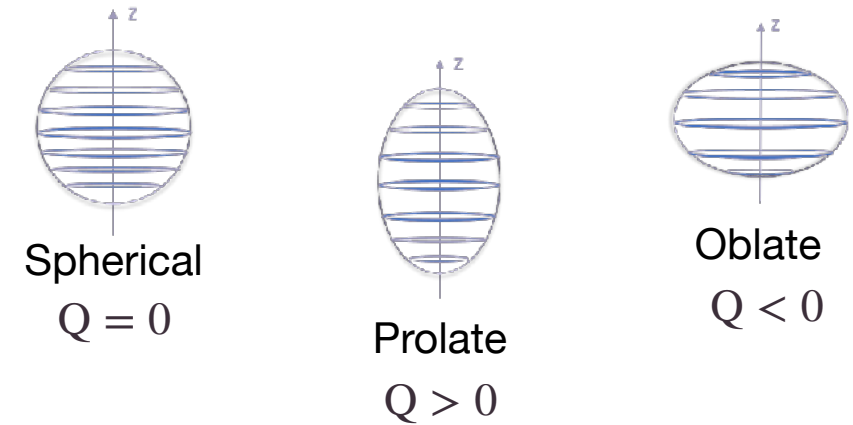


Backup Slides

^{185}Re and ^{187}Re 5g4f analysis

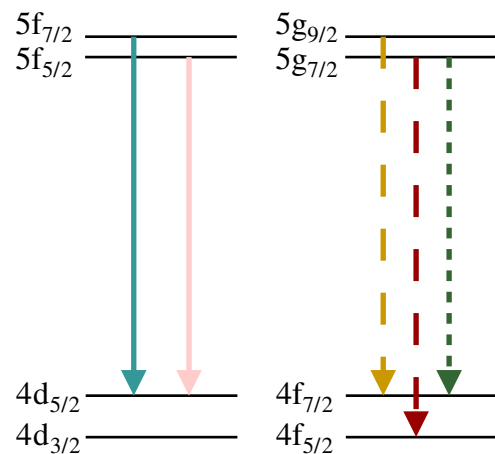
Spectroscopic quadrupole moment

The spectroscopic quadrupole moment Q is extracted from the 5 $\rightarrow$ 4 transitions being not sensitive to the details of the nuclear charge distribution

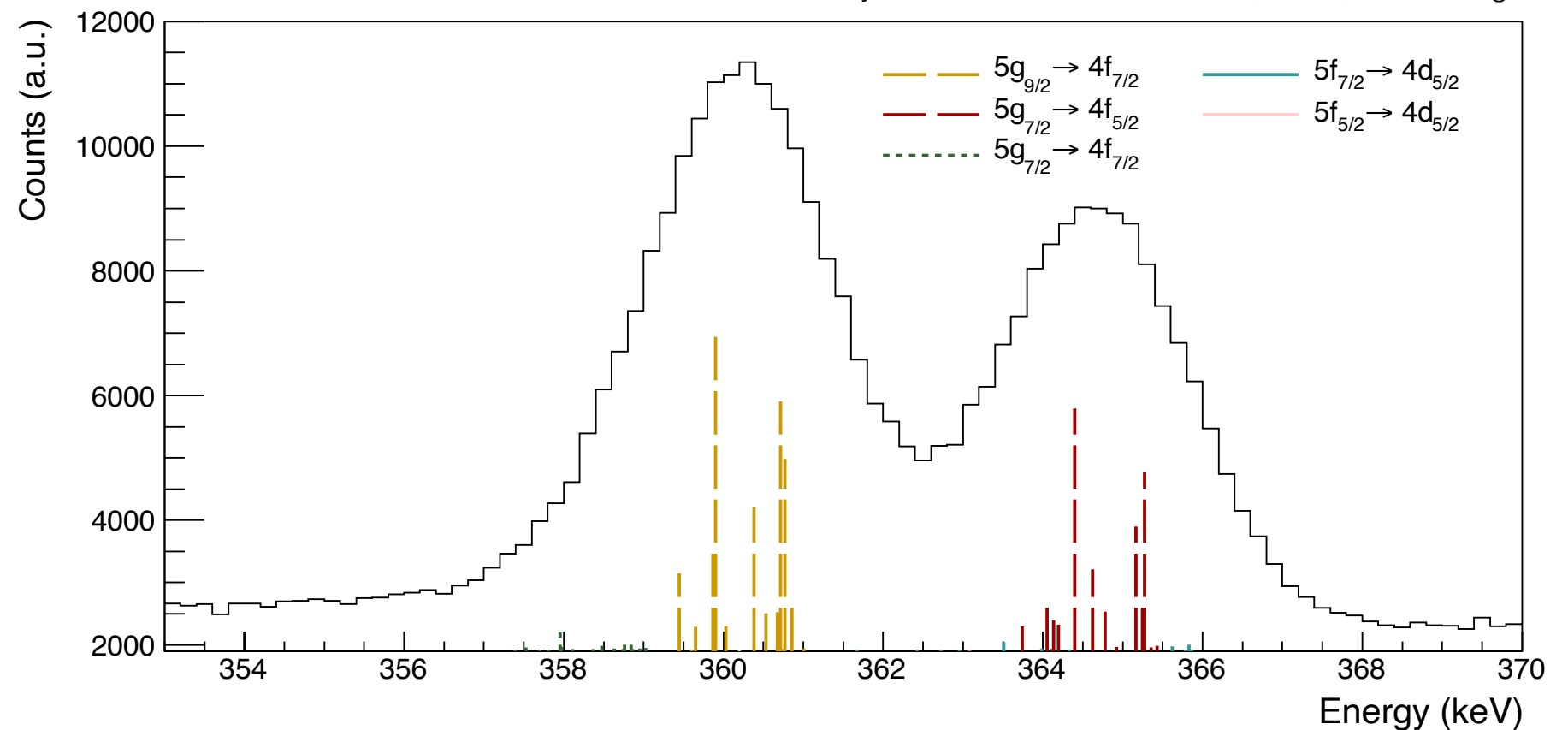


To extract the quadrupole moment:

- Theoretical predictions of hyperfine transitions
- Good knowledge of the instrumental response is required



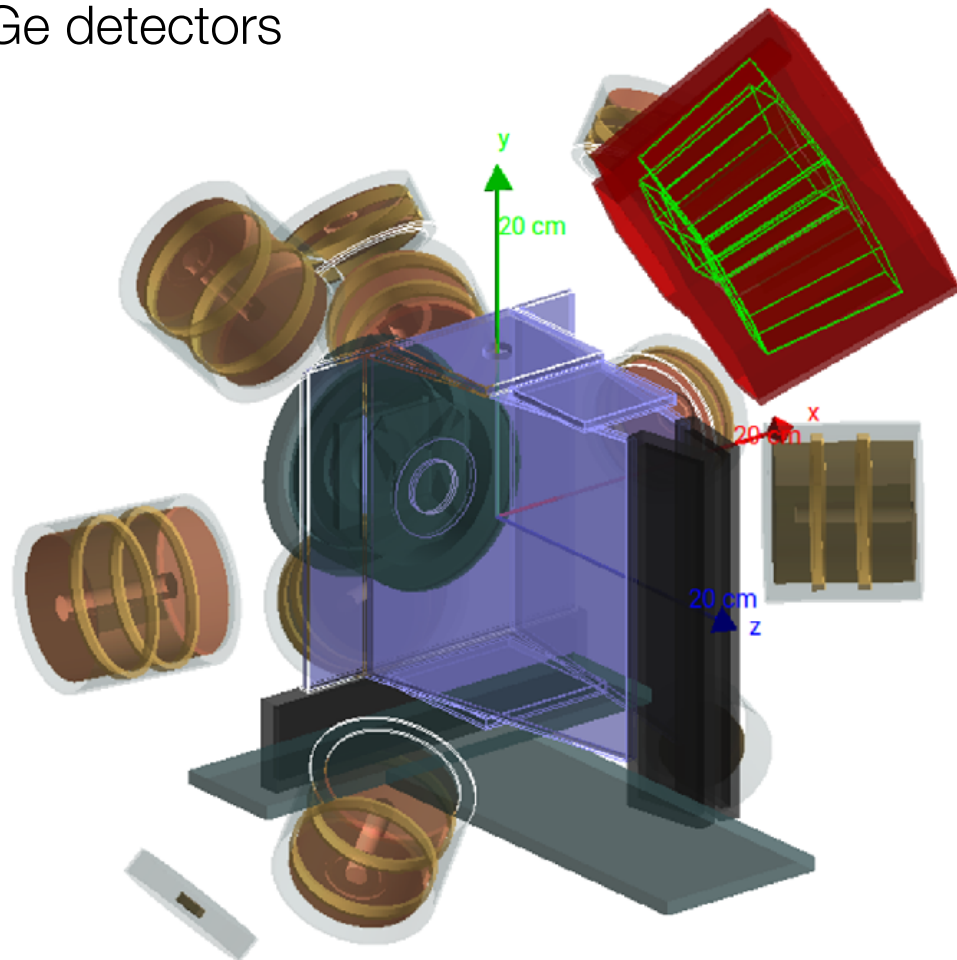
Theoretical calculations by **N. Michel & N. Oreshkina**, MPIK, Heidelberg



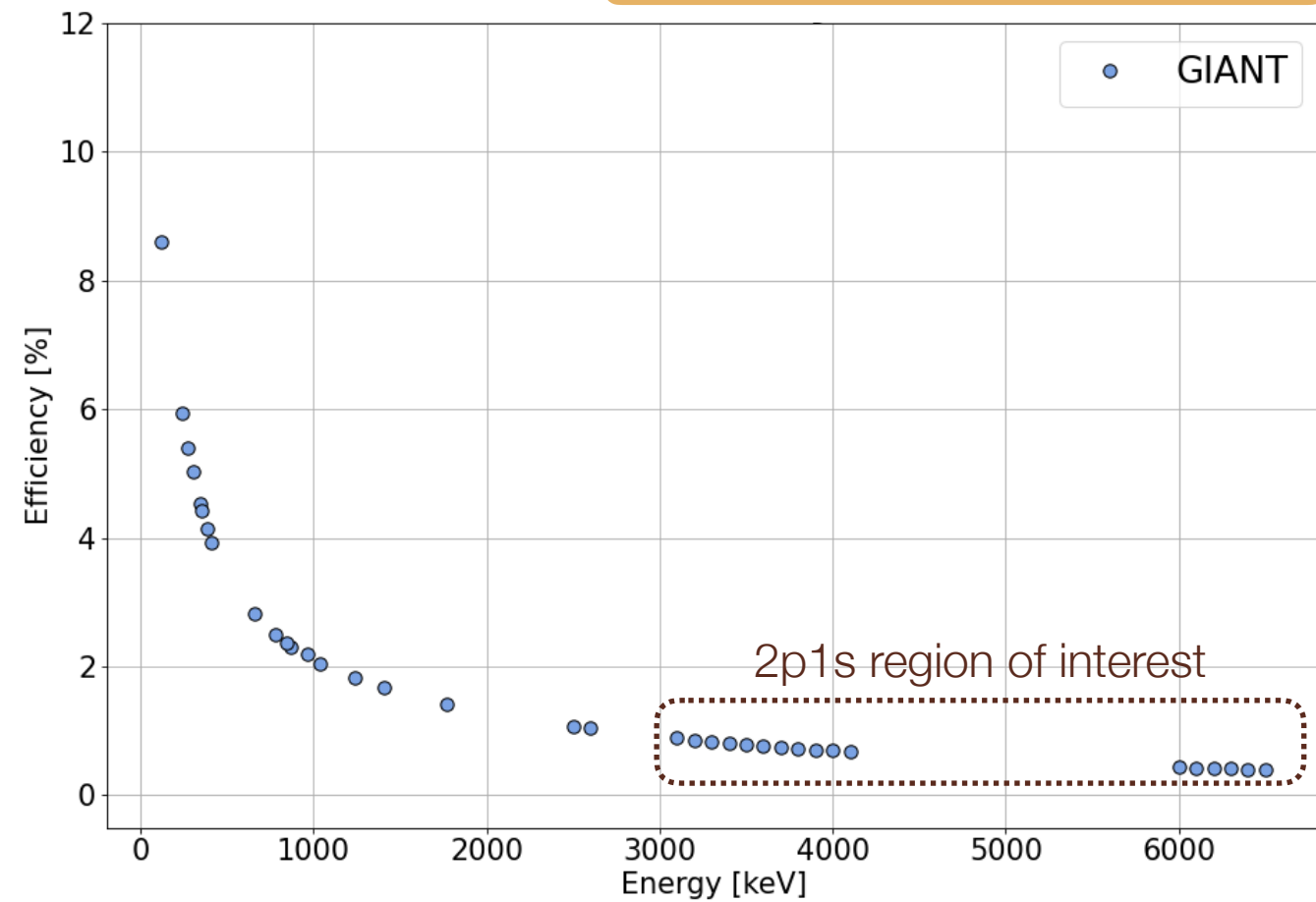
Optimizing the geometry for higher efficiency

Full GEANT4 simulations

- Full GEANT4 simulations of the detectors' position in the frame considering all physical limitations of the setup.
- Angle dependencies on single crystals and combine all HPGe detectors to reach the maximum possible efficiency
- Use wisely the electron veto-ed space for good HPGe detectors

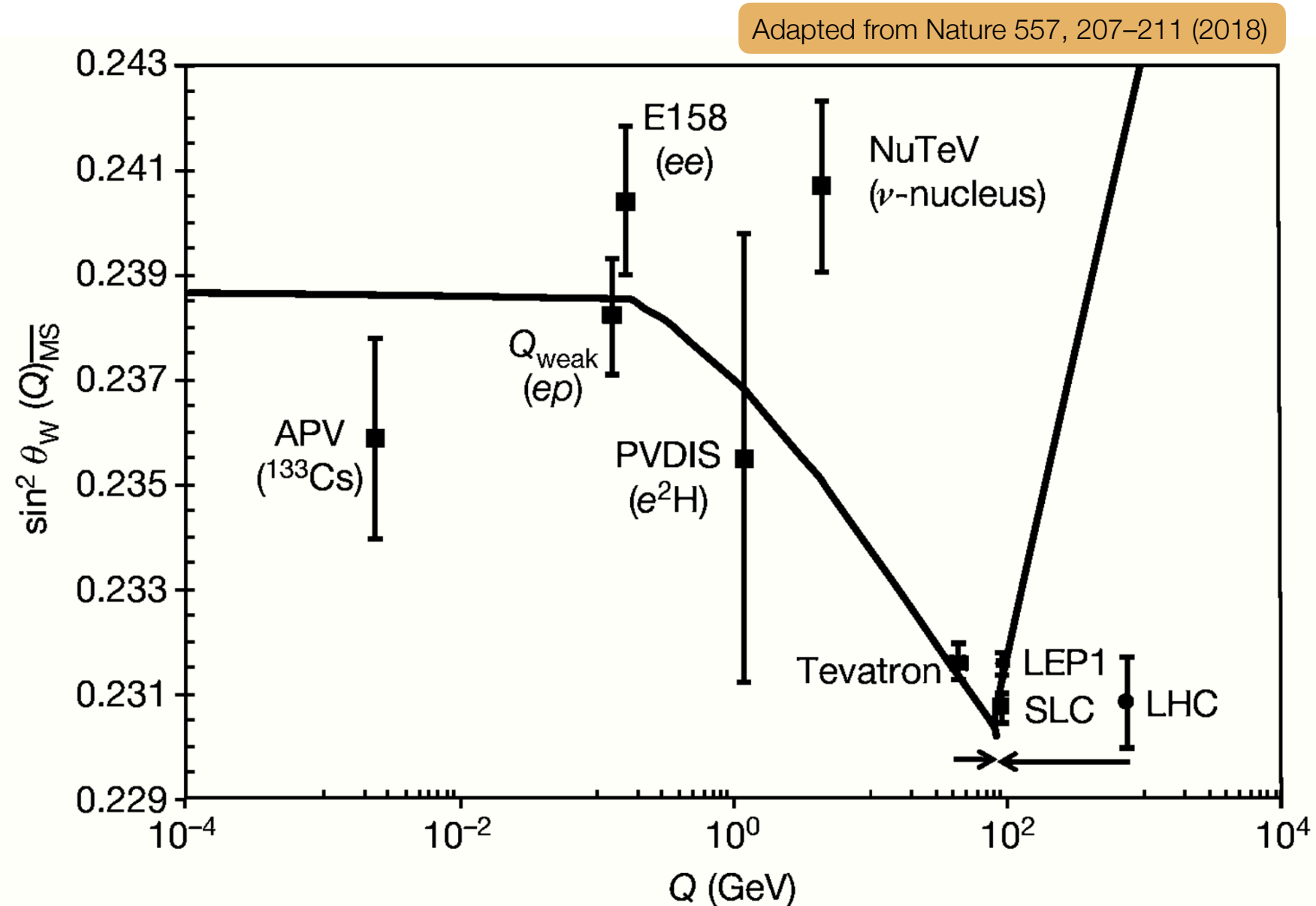


Geant4 simulations and visualisations created by **Marie Deseyn**



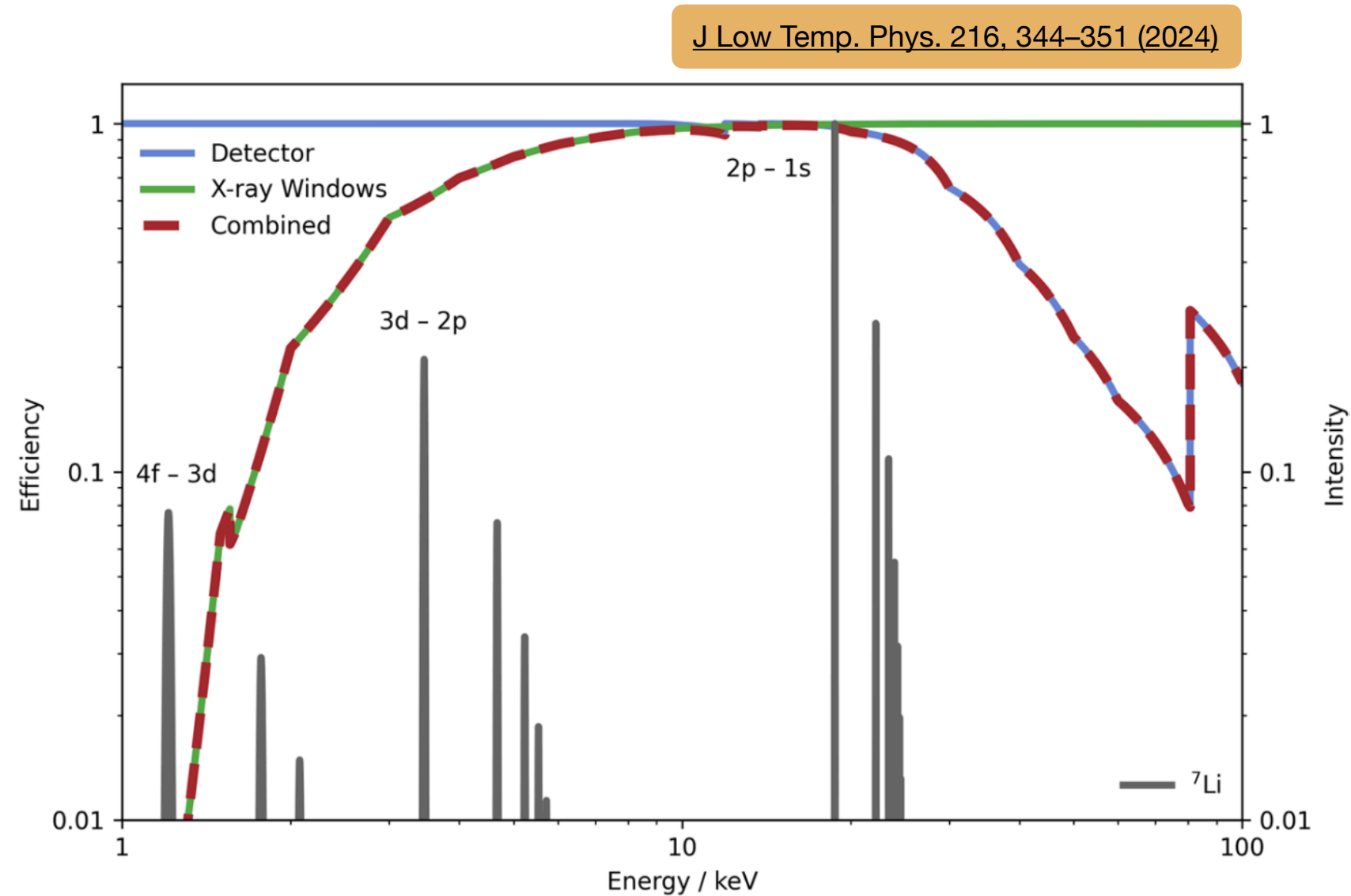
Motivation for ^{226}Ra

Atomic parity violation in radium



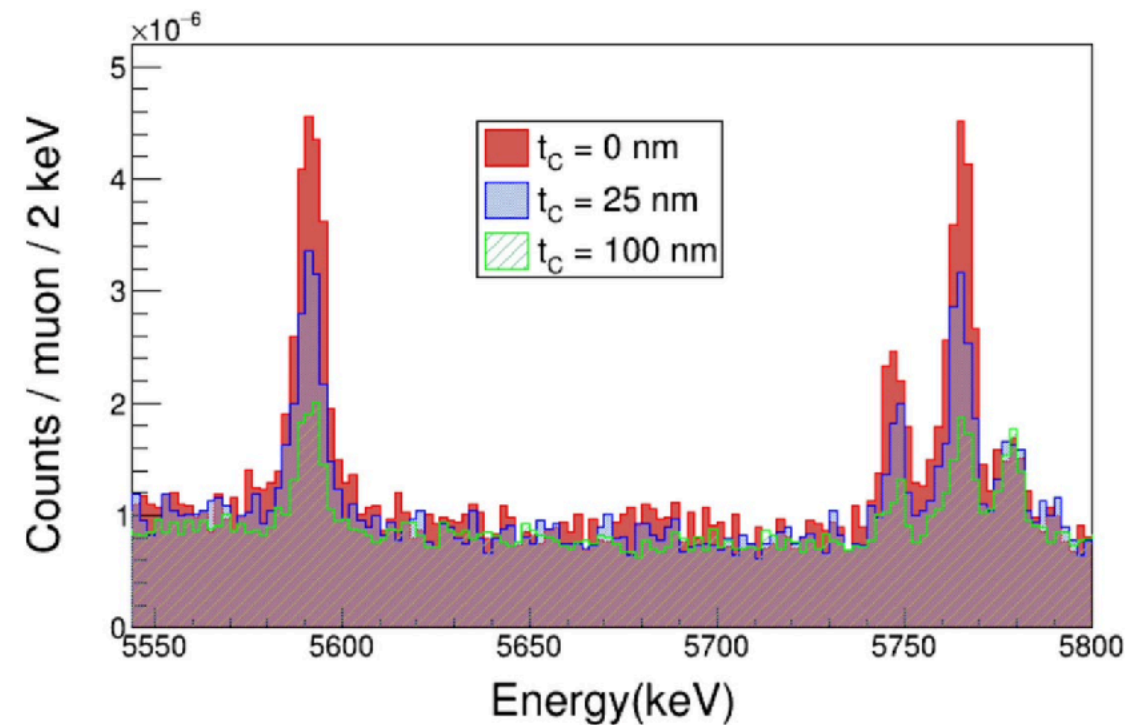
MMC detector

Detection efficiency



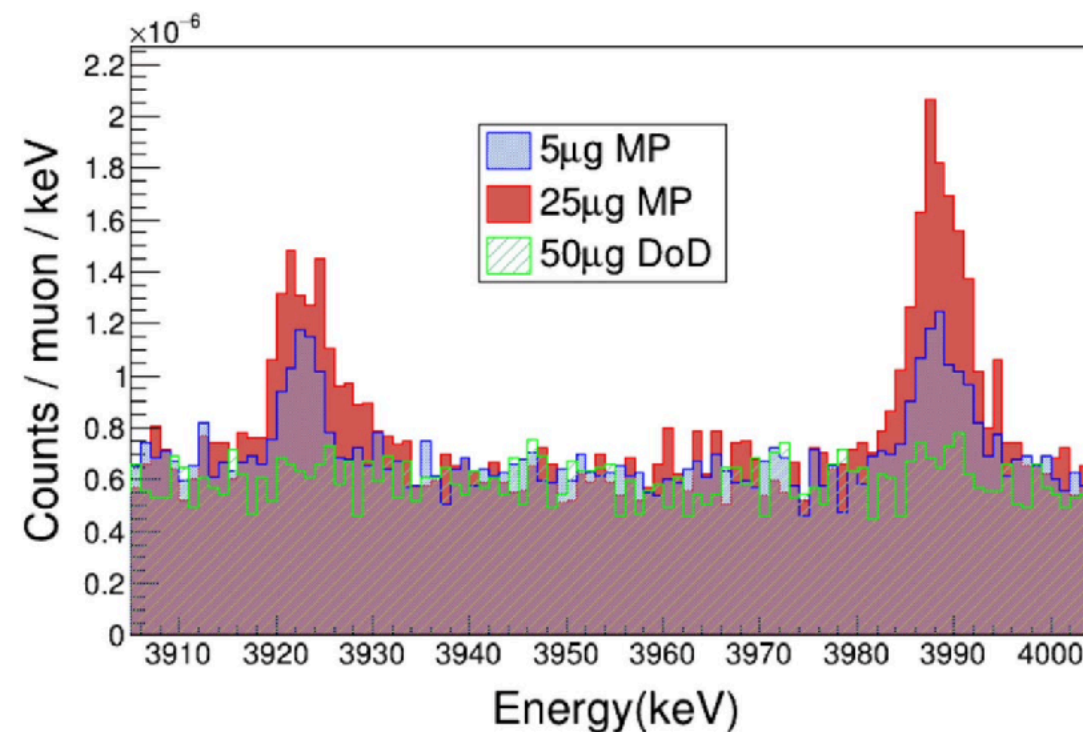
Comparing target fabrication techniques

Graphite targets

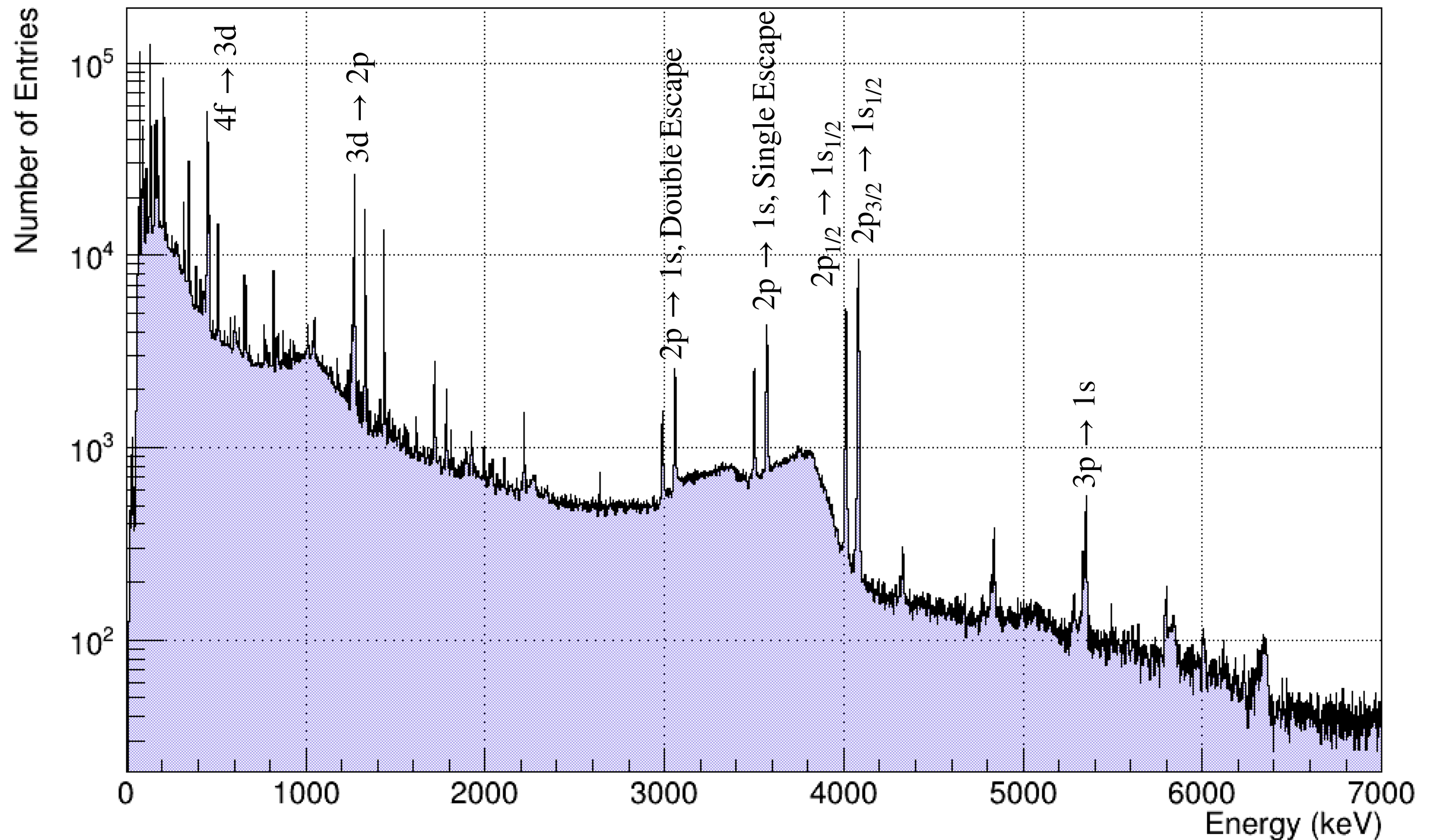


Sci Rep 15, 6939 (2025)

Molecular plating &
Drop-on-Demand



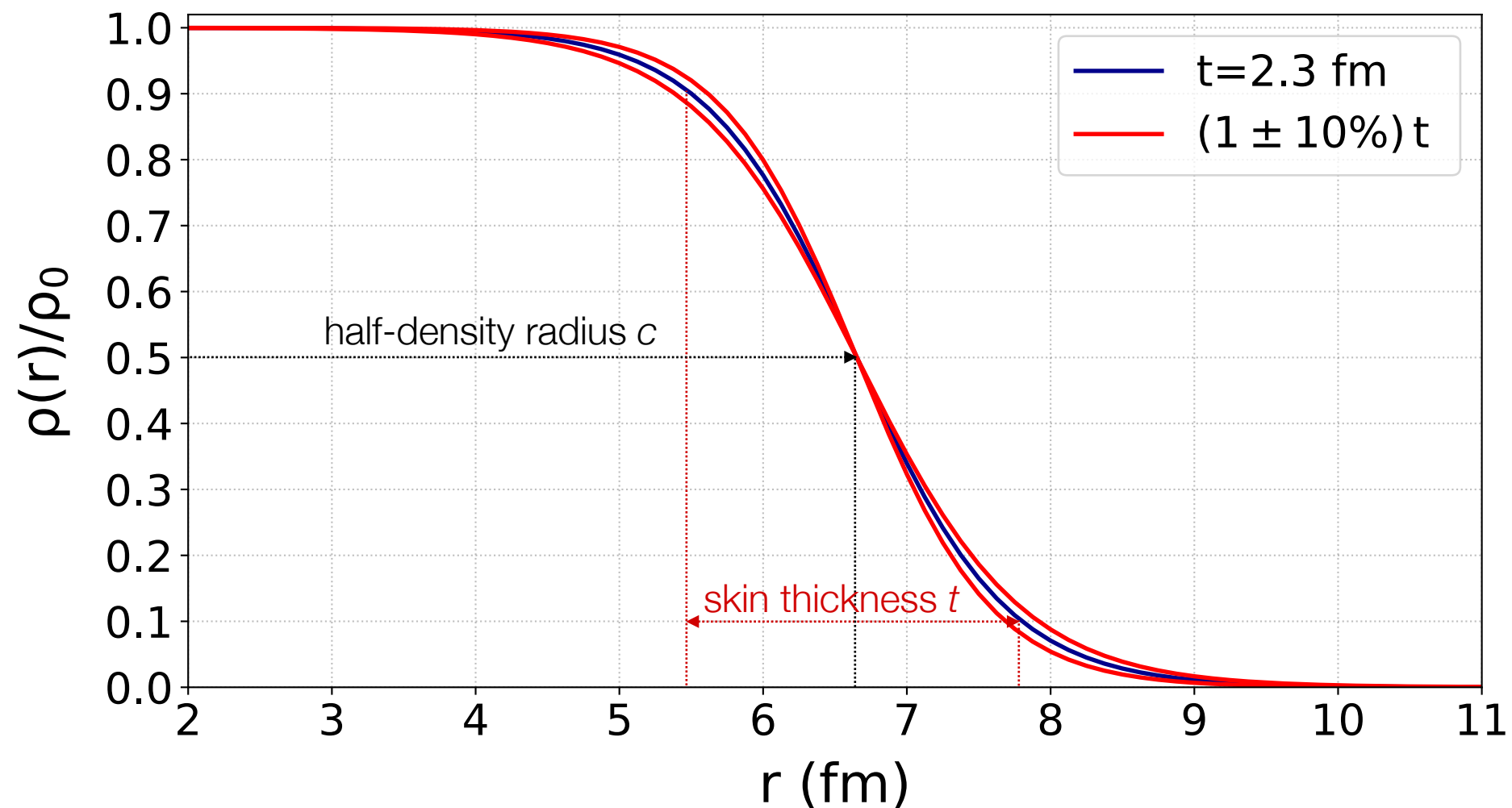
Muonic spectrum in $^{\text{nat}}\text{La}$



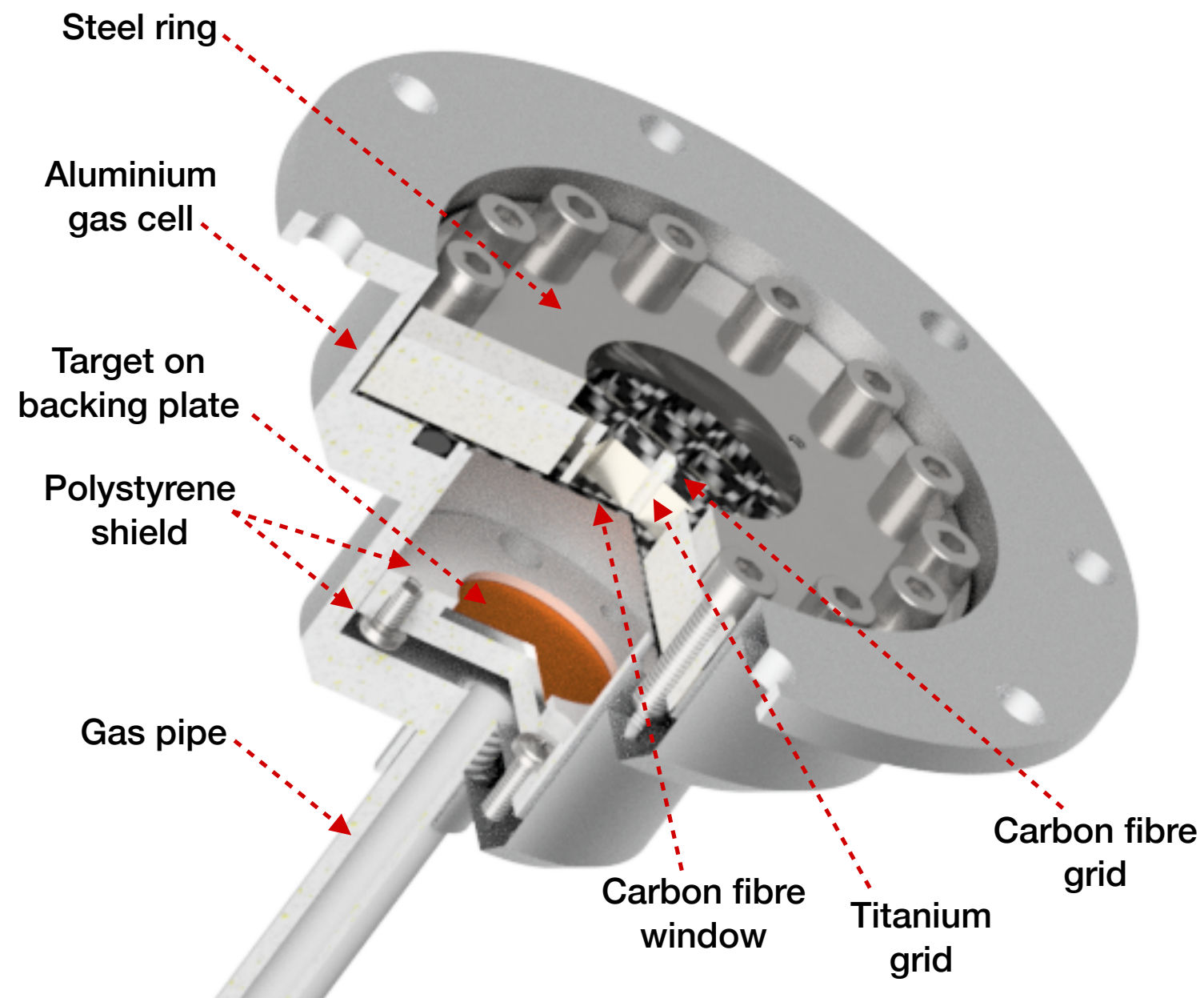
Charge distribution model

Nuclear charge distribution
for spherical nuclei

$$\rho(r) = \rho_0 \left(1 + e^{4 \ln 3 \frac{r-c}{t}} \right)^{-1}$$

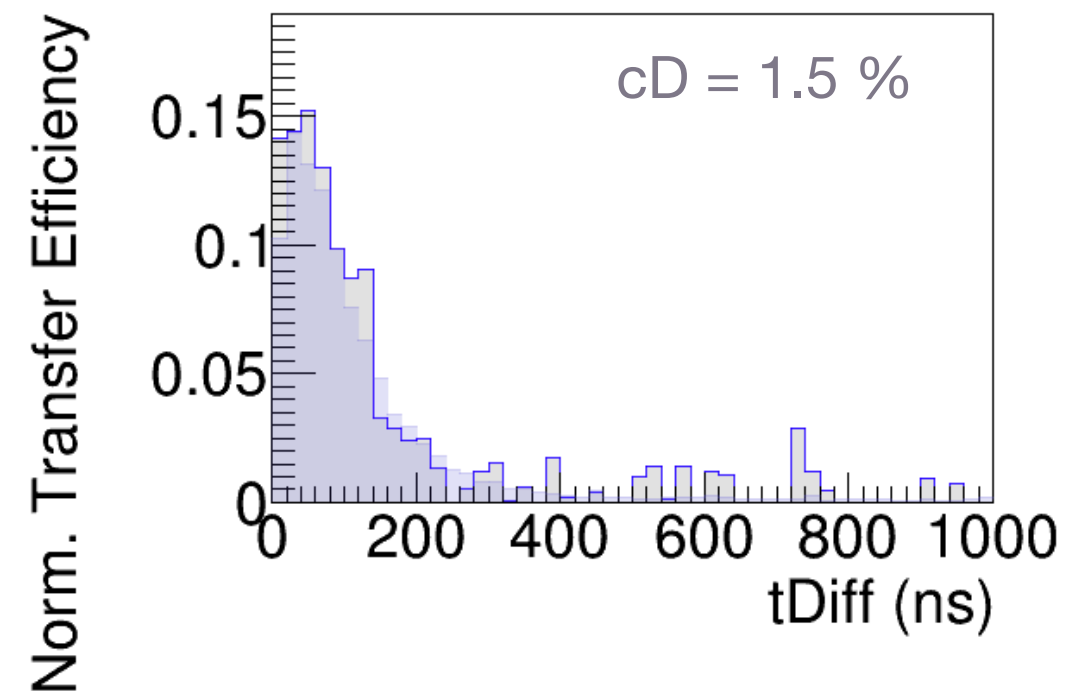
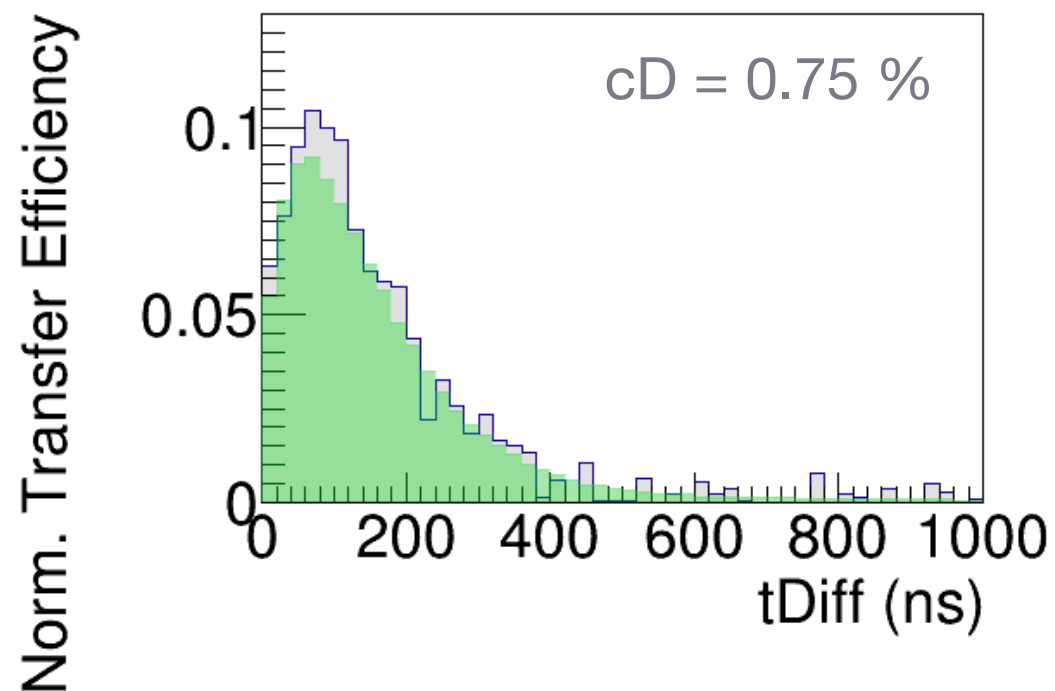
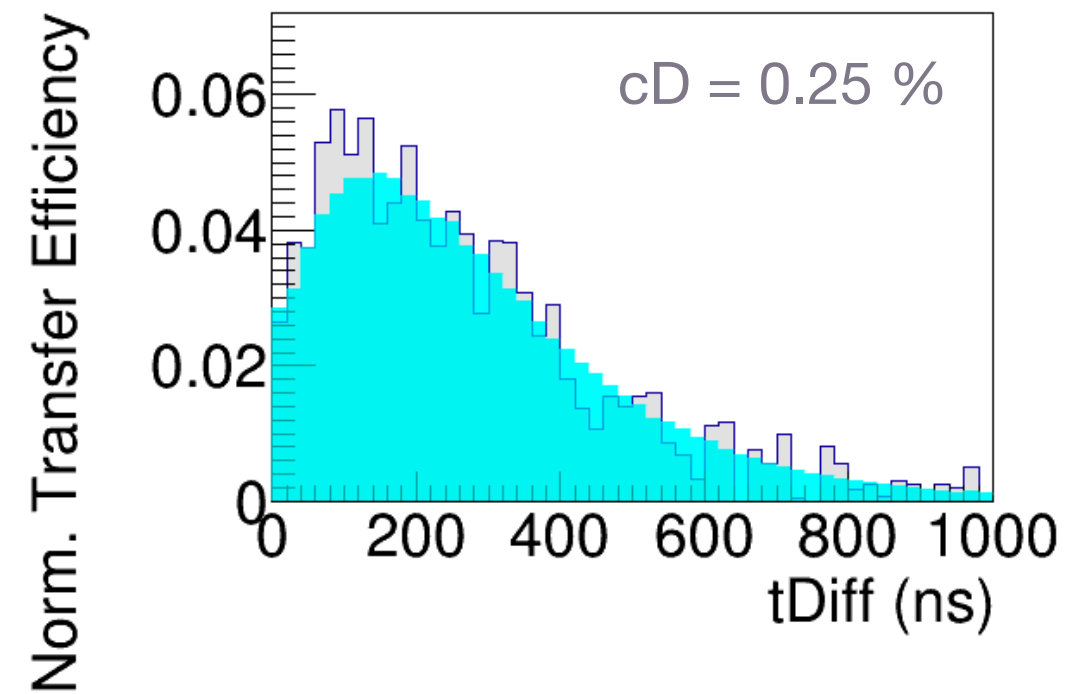
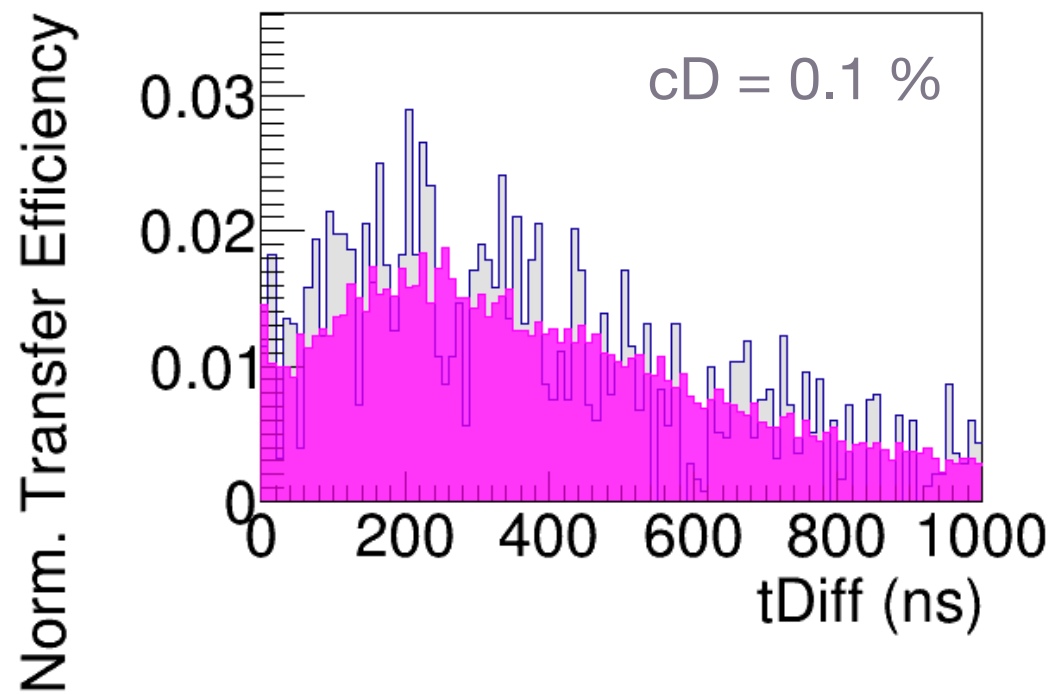


Gas cell structure

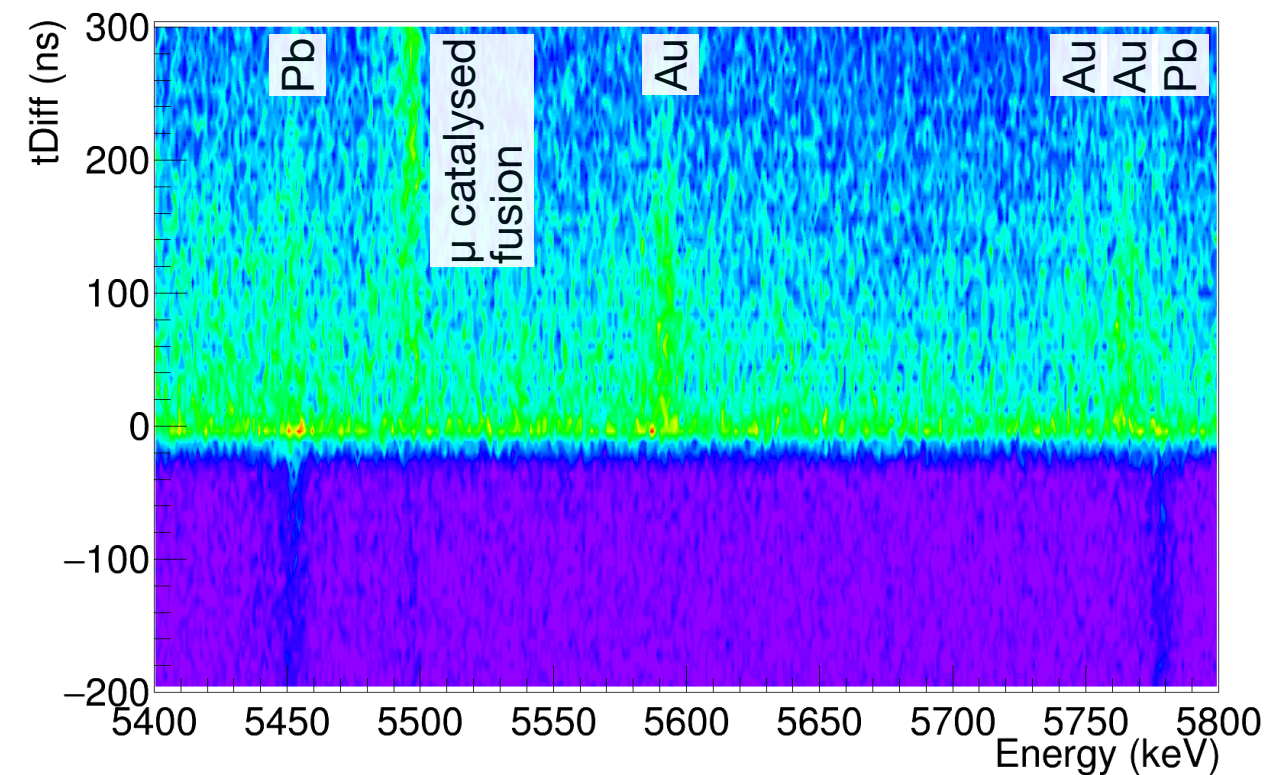
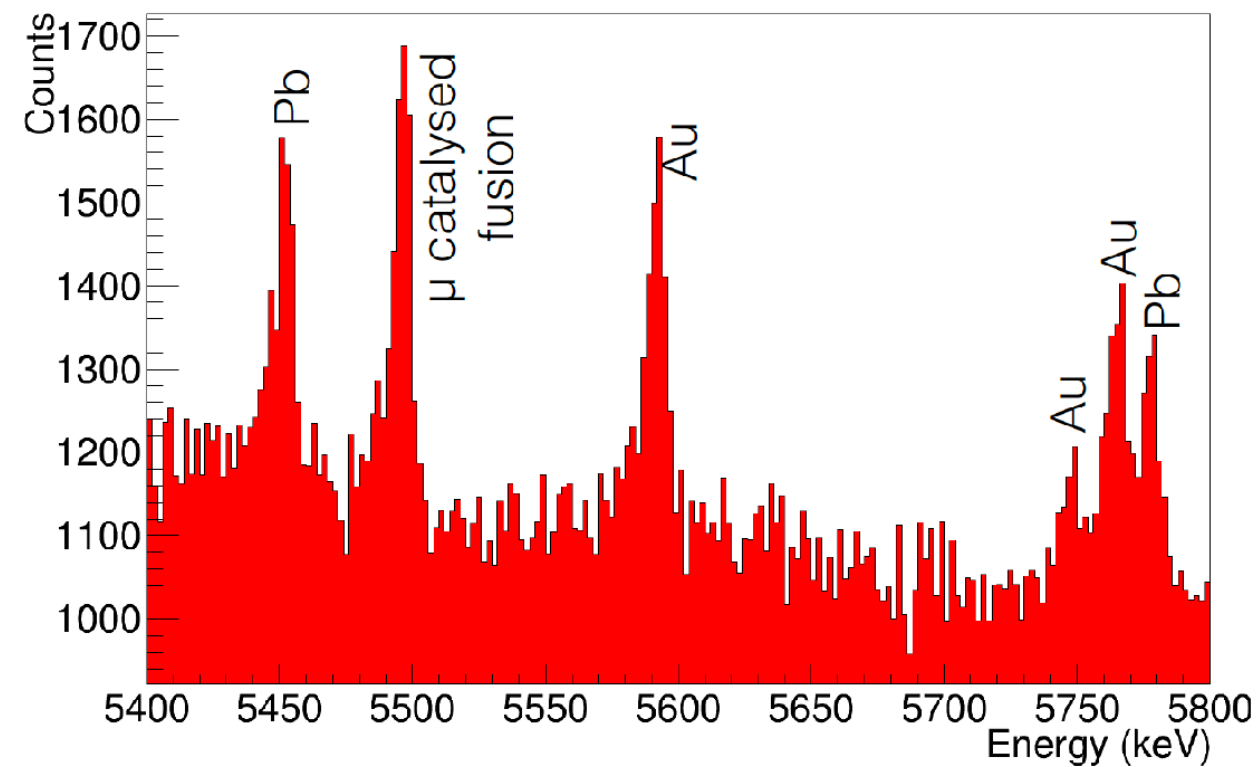


Time distributions after transfer

PhD thesis of A. Skawran
Simulations by J. Nuber



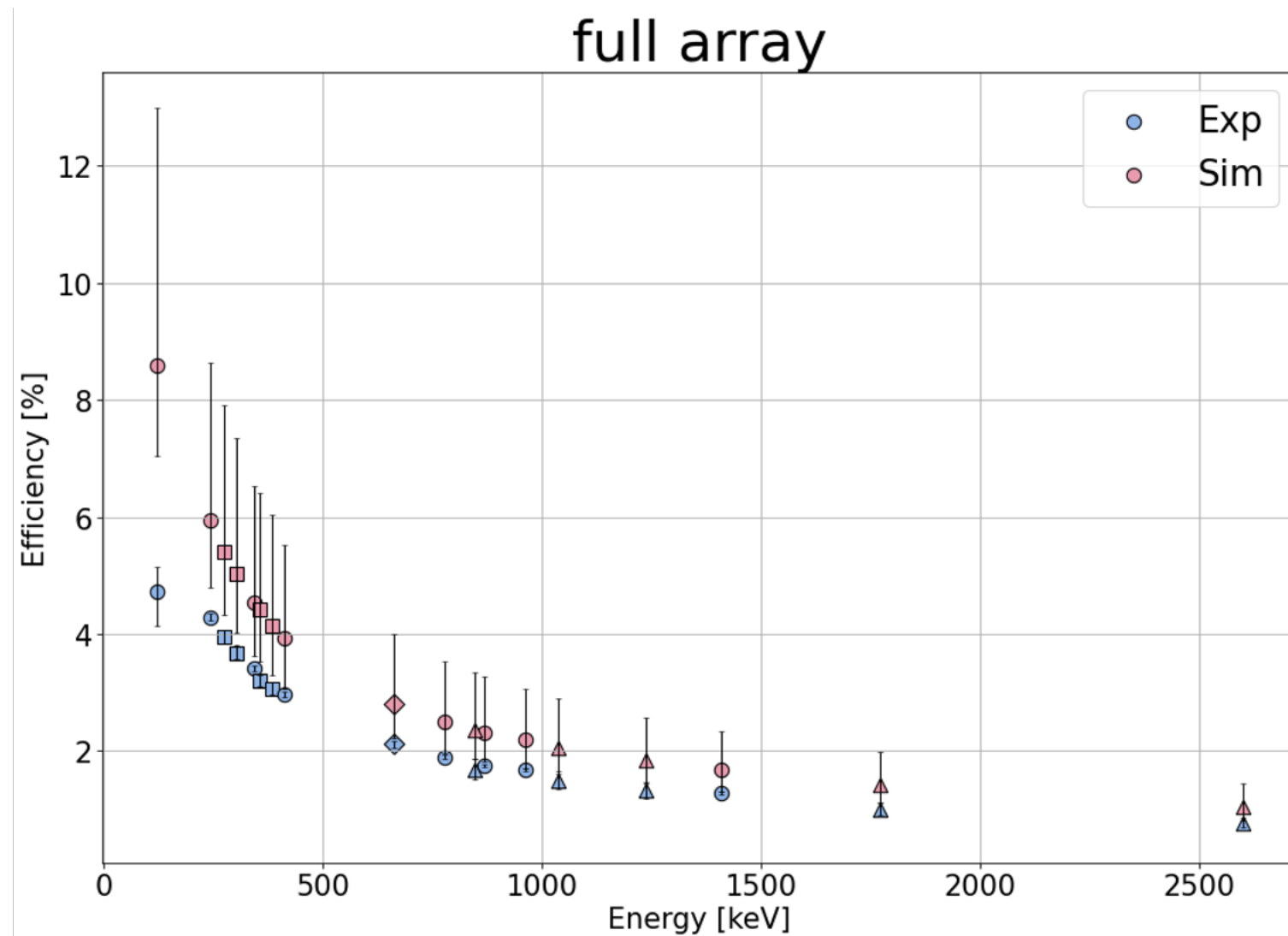
Demonstration of muon transfer reaction



GIANT efficiency comparison

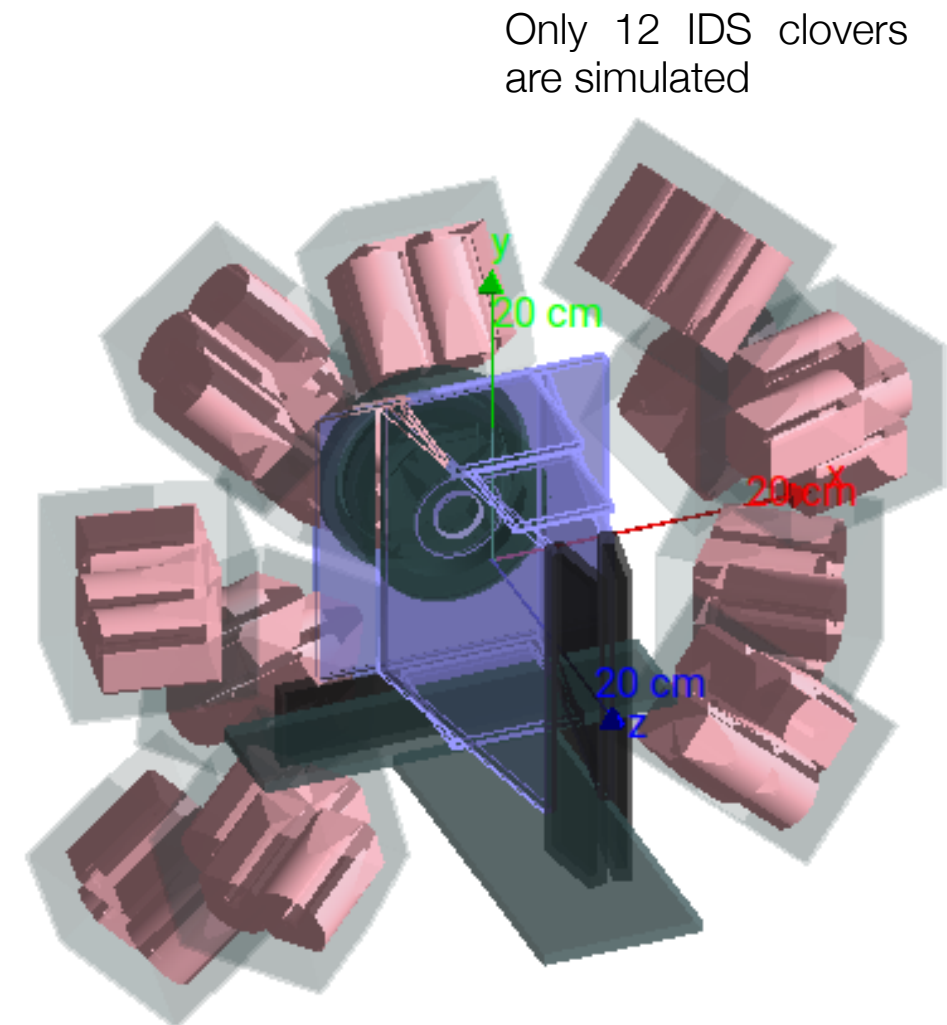
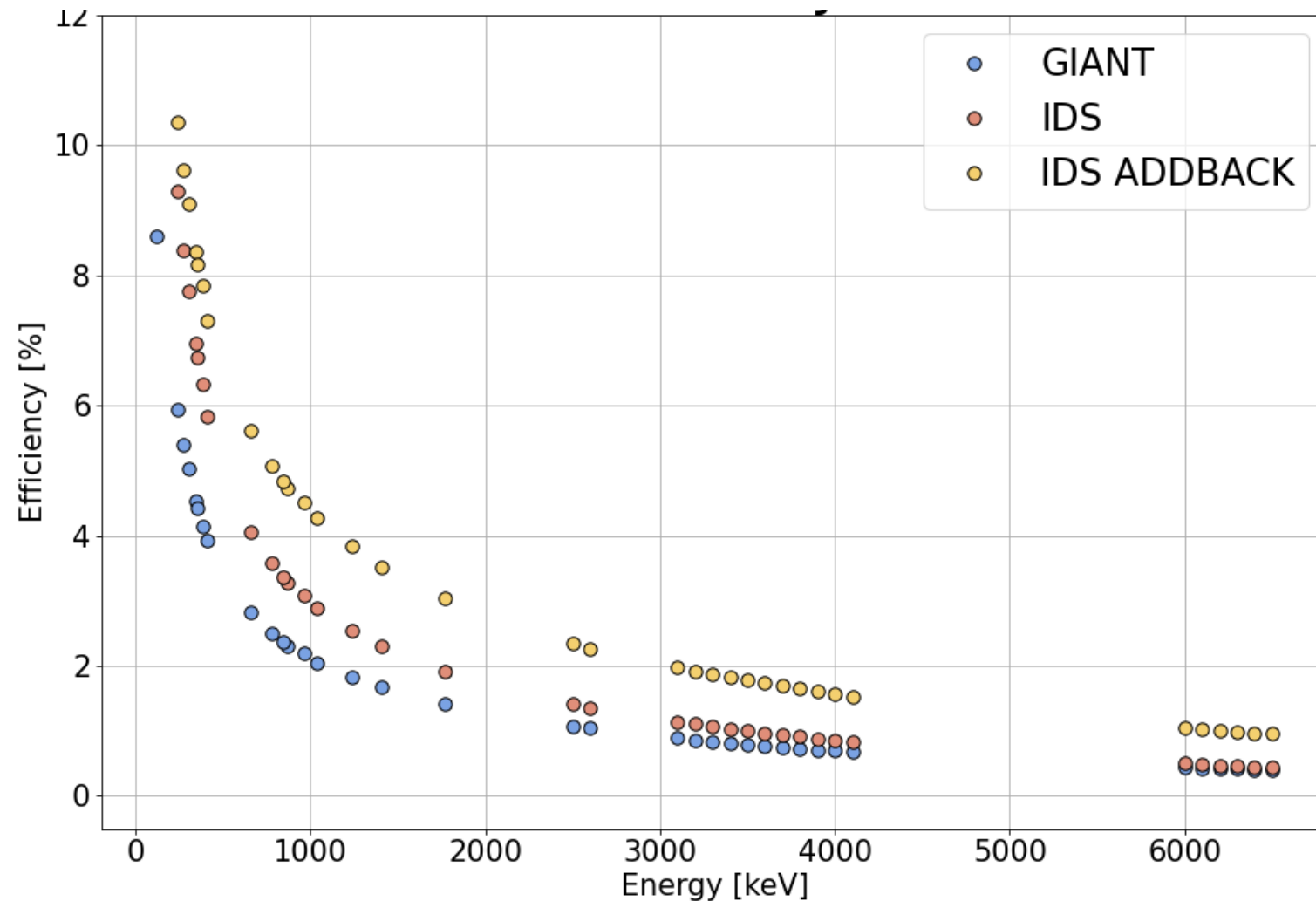
Experimental VS simulated

Image courtesy of **Marie Deseyn**



Challenging the microgram targets

Higher efficiency with the IDS array

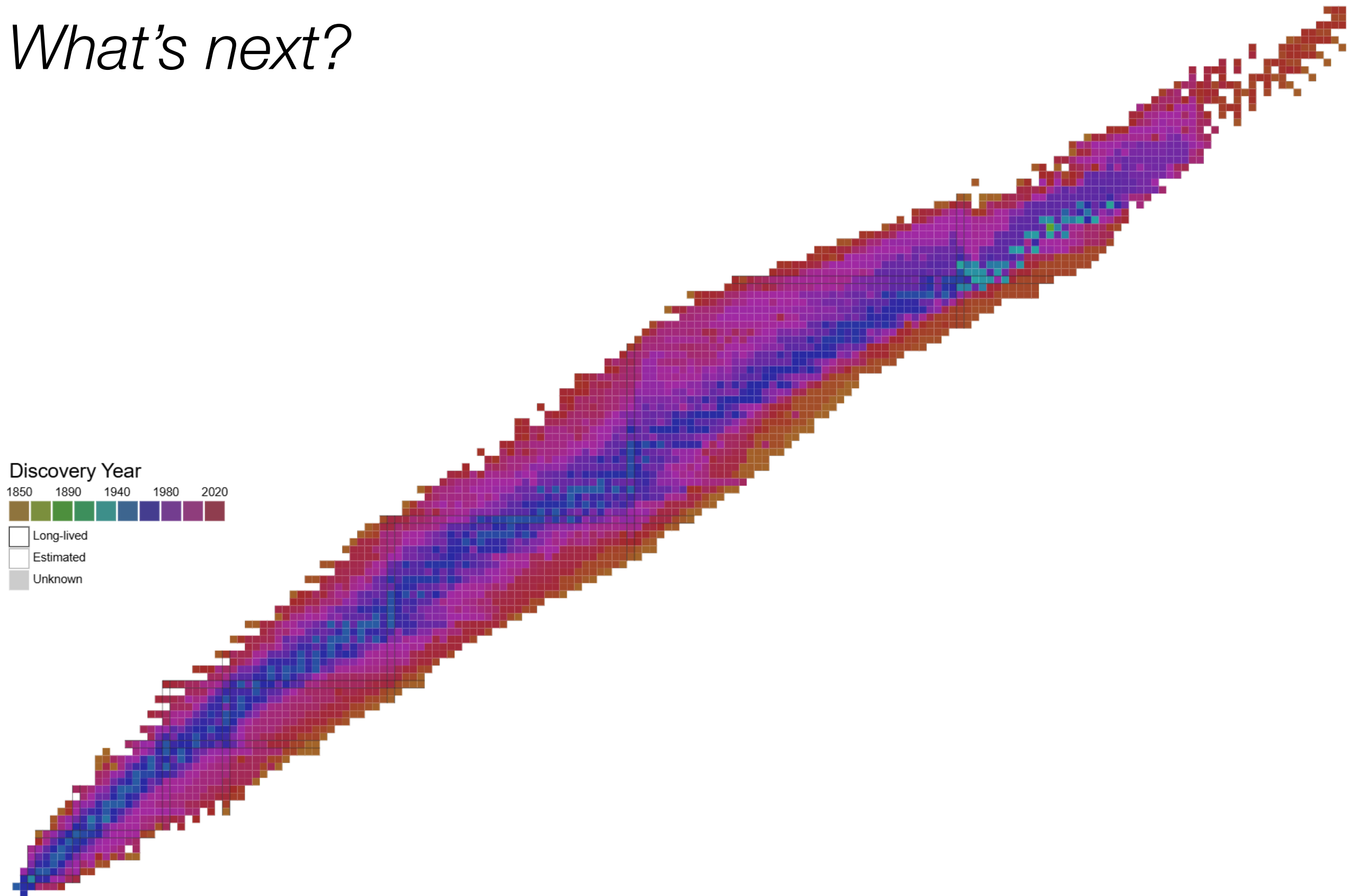


**IDS simulation assumes 20 mm between end cup and all crystals. Efficiency to be increased by ~25% for the corrected gap.

Geant4 simulations and visualisations created by **Marie Deseyn**.
IDS detectors' geometry by **War War Myint Myat Phyo**.

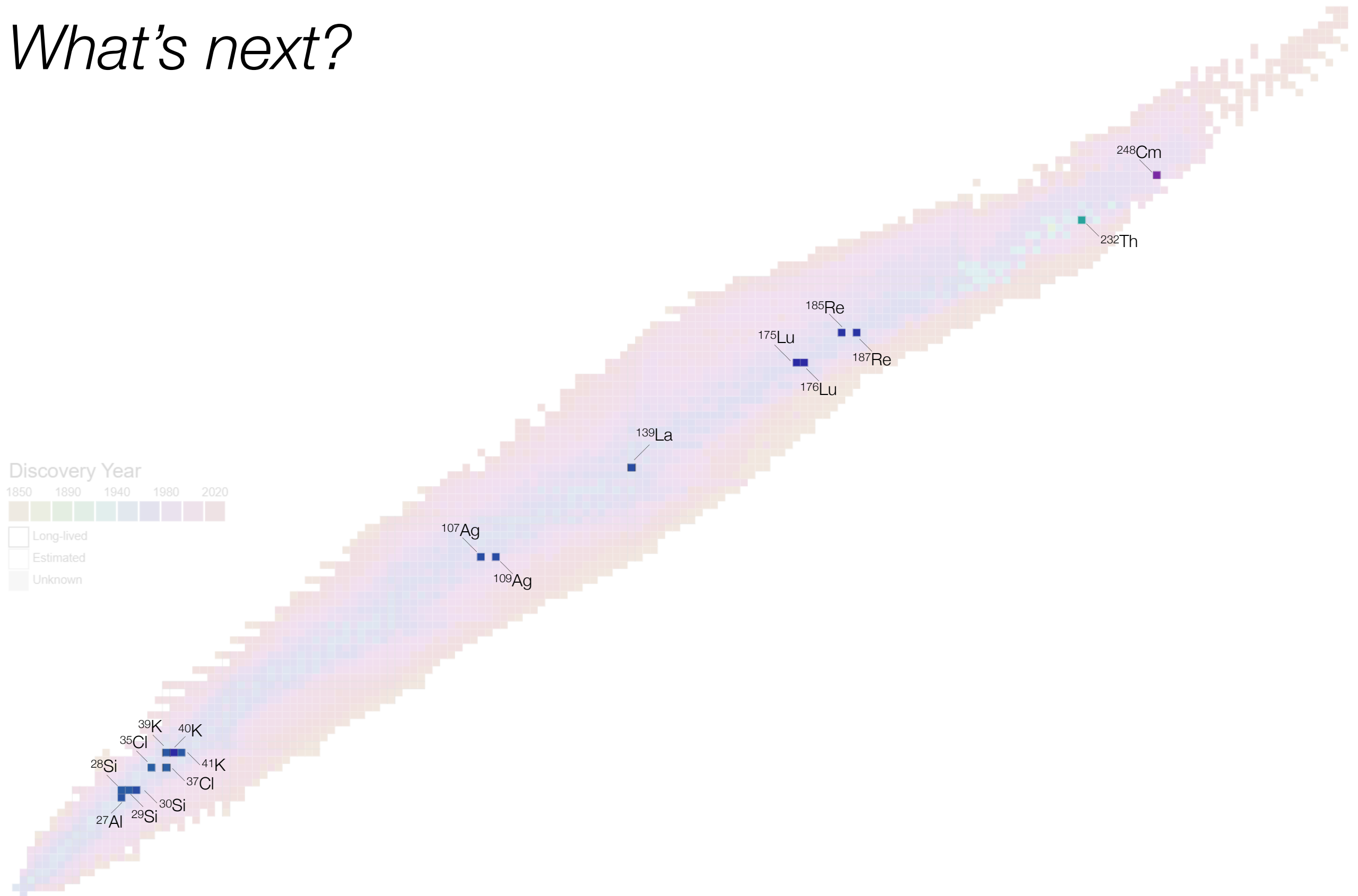
Science case

What's next?



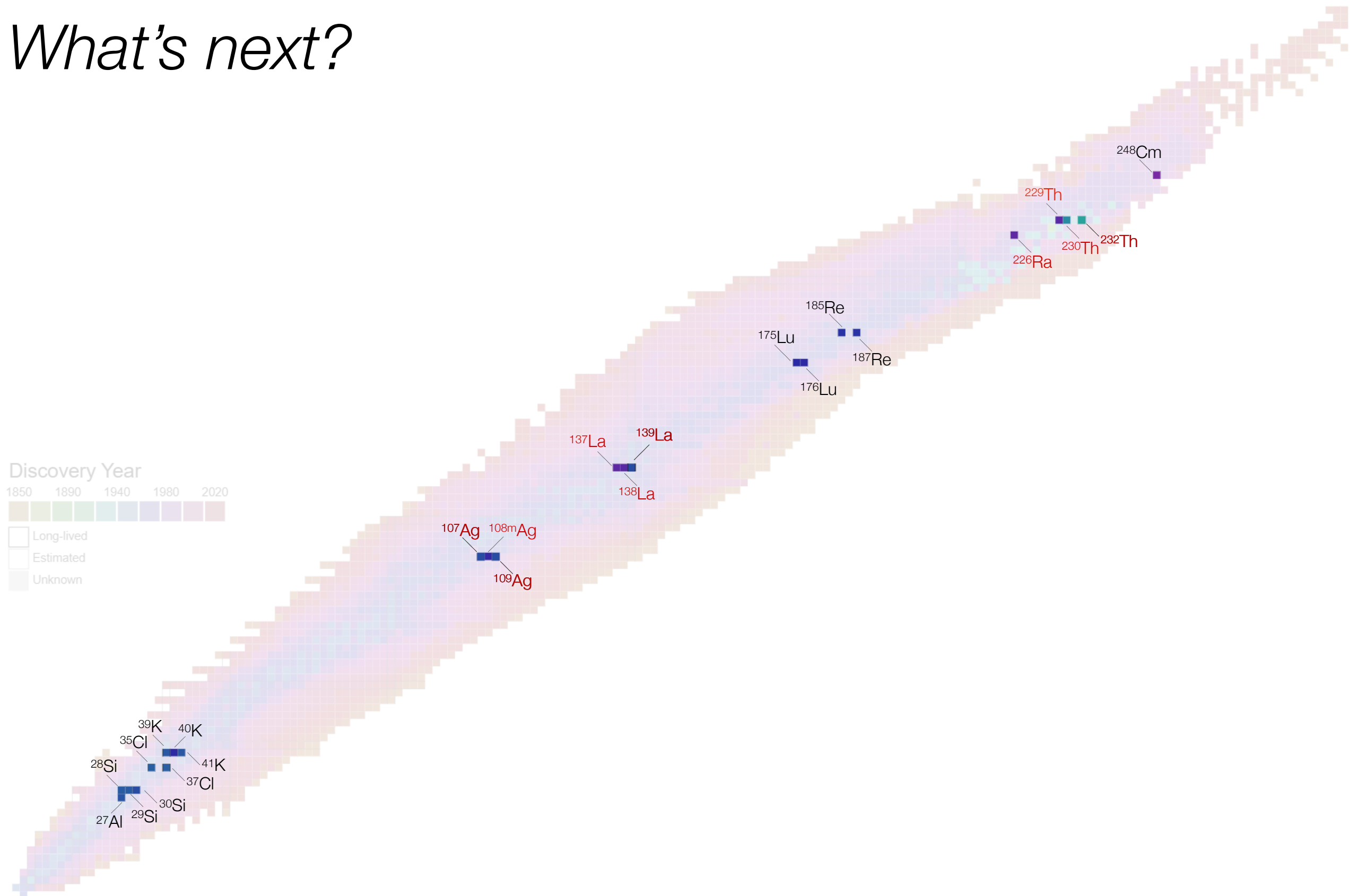
Science case

What's next?



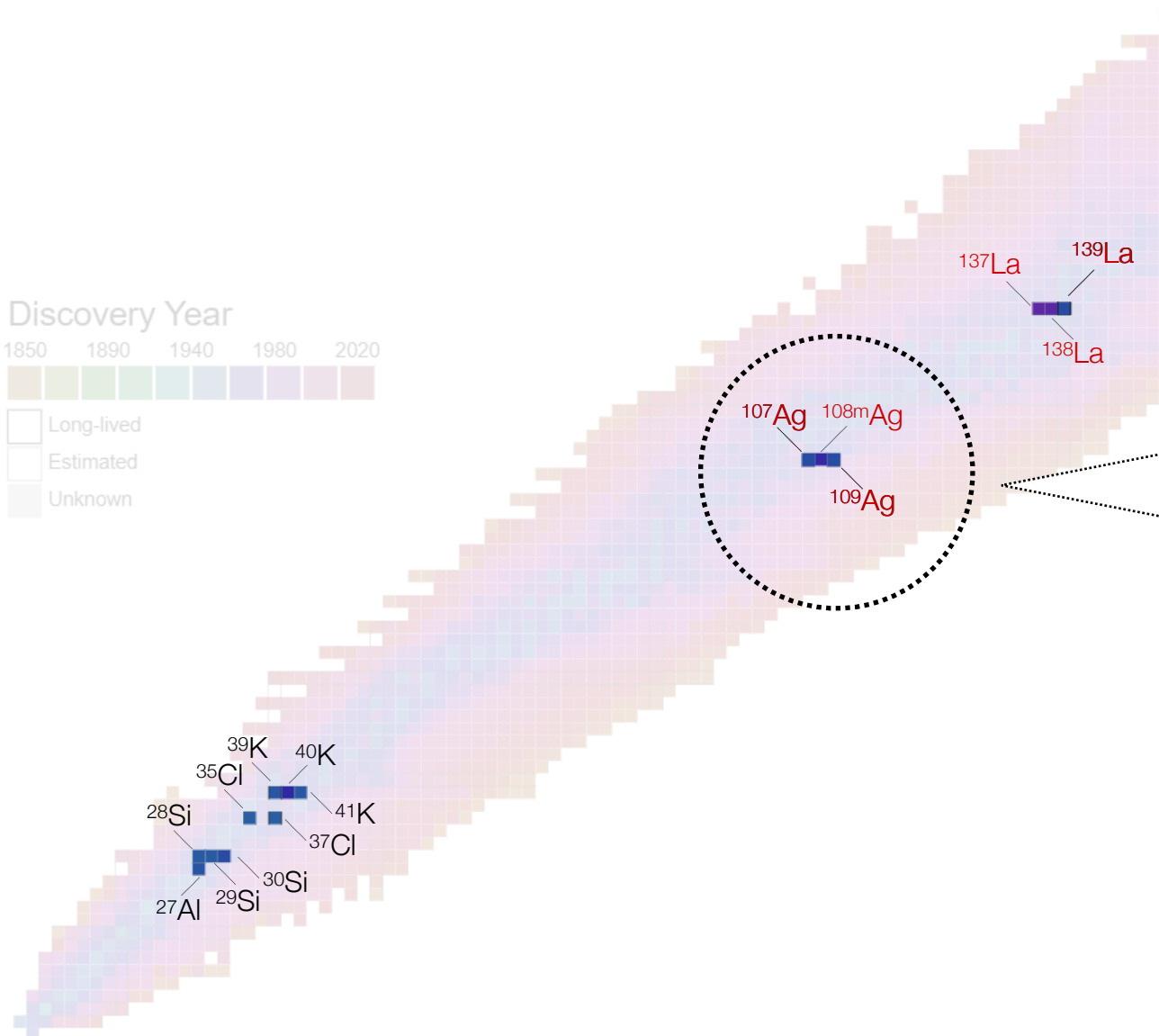
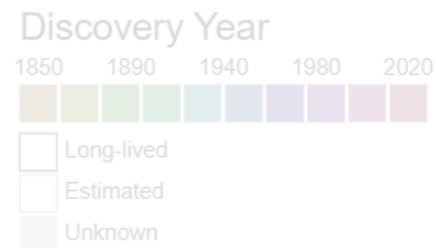
Science case

What's next?



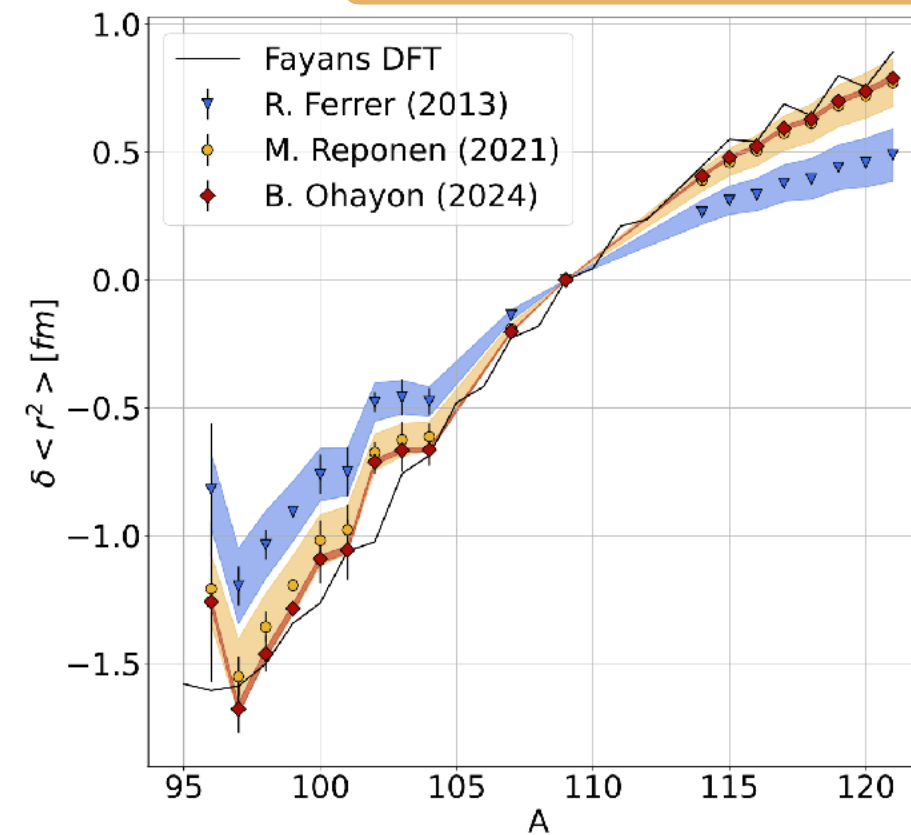
Science case

What's next?



Benchmarking the laser spectroscopy results & test state-of-art nuclear models

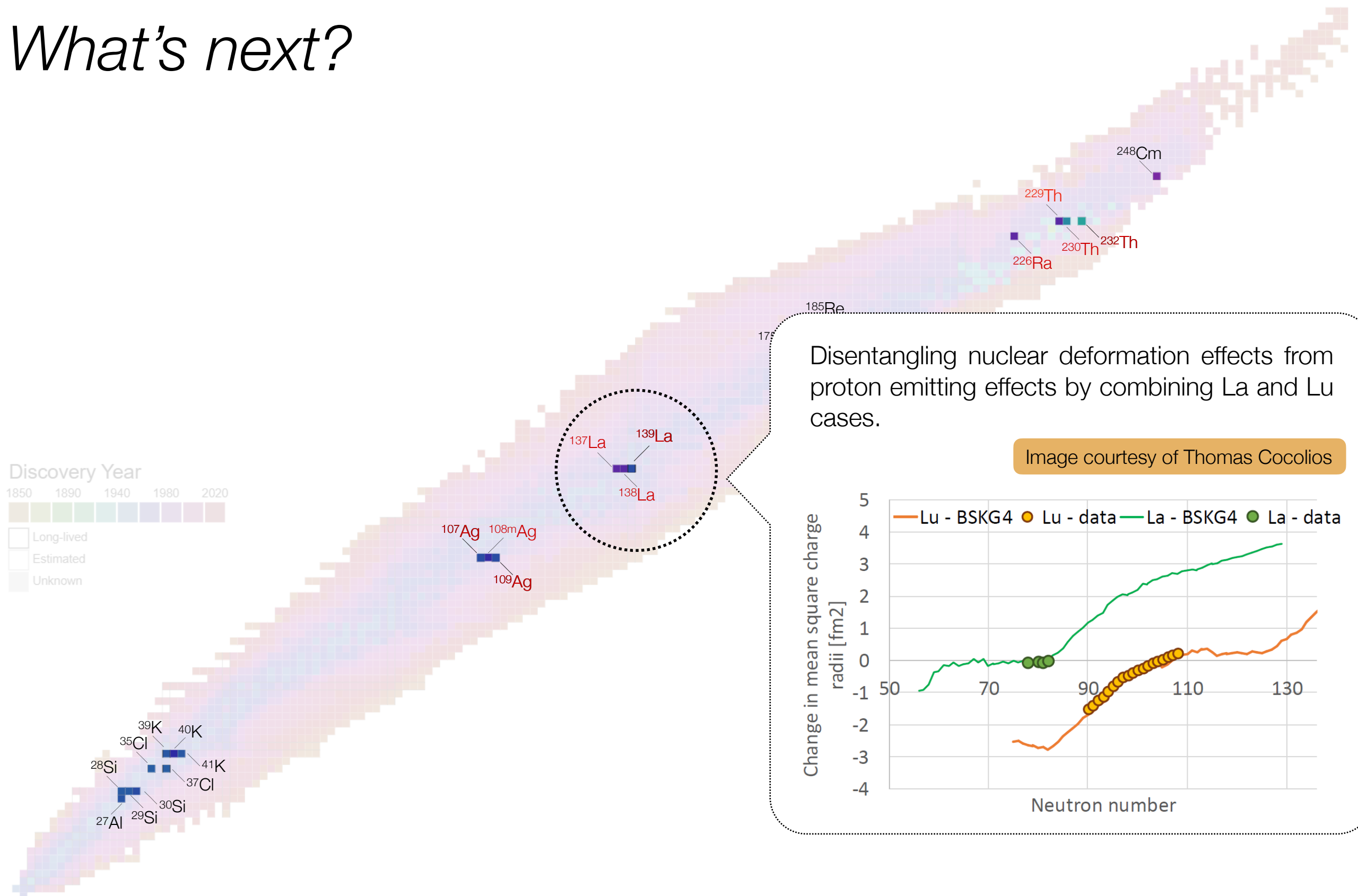
Image courtesy of Marie Deseyn



Phys. Rev. Research 6, 033040 (2024)
Nat. Commun. 12, 4596 (2021)
Phys. Lett. B 278, 191 (2014)

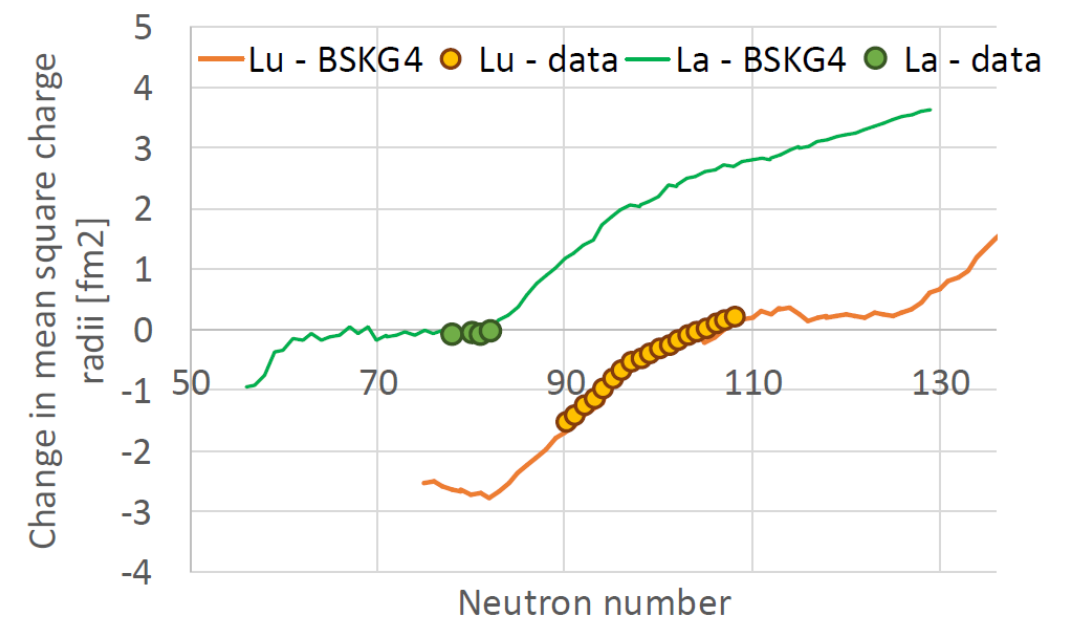
Science case

What's next?



Disentangling nuclear deformation effects from proton emitting effects by combining La and Lu cases.

Image courtesy of Thomas Cocolios

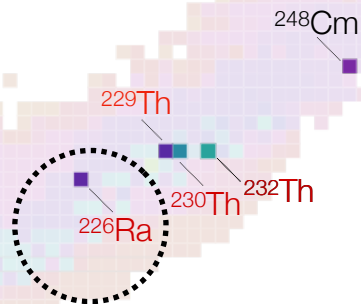
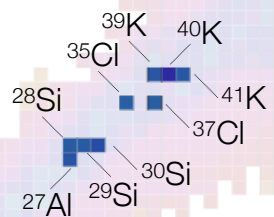
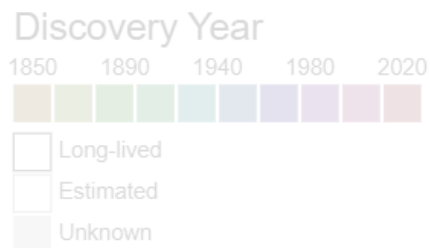
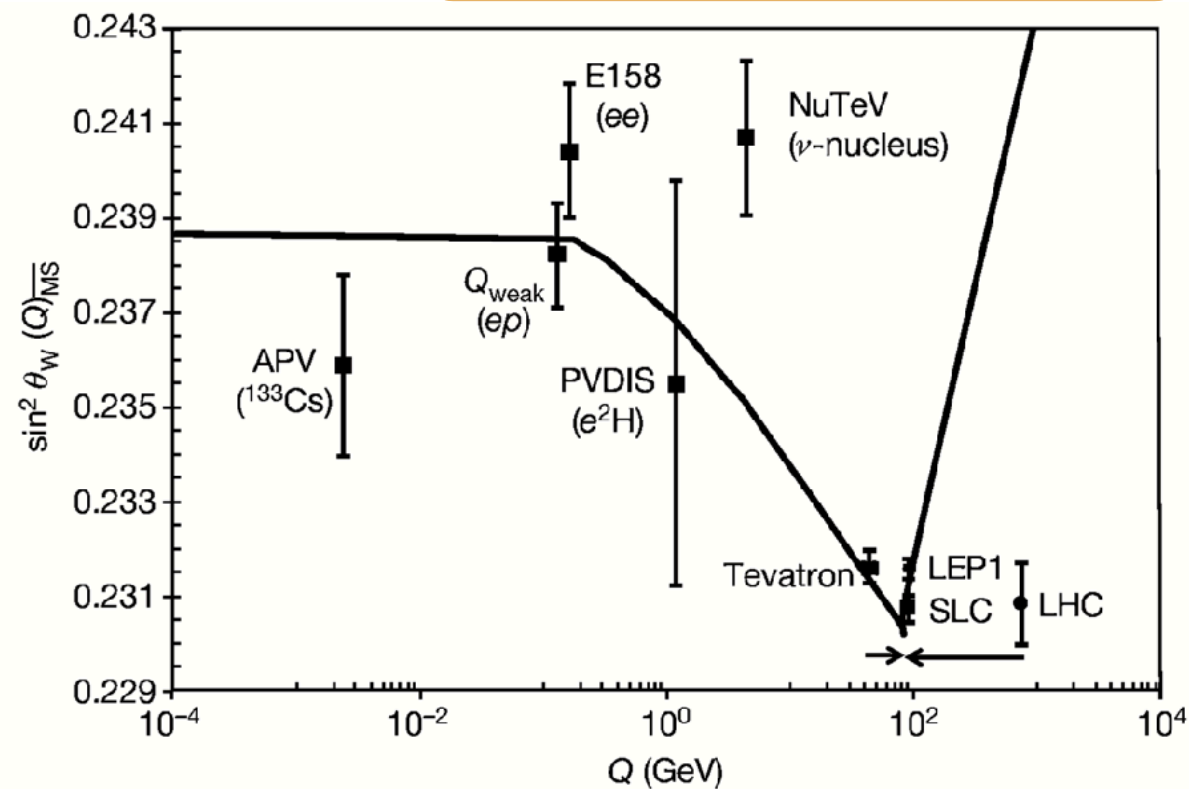


Science case

What's next?

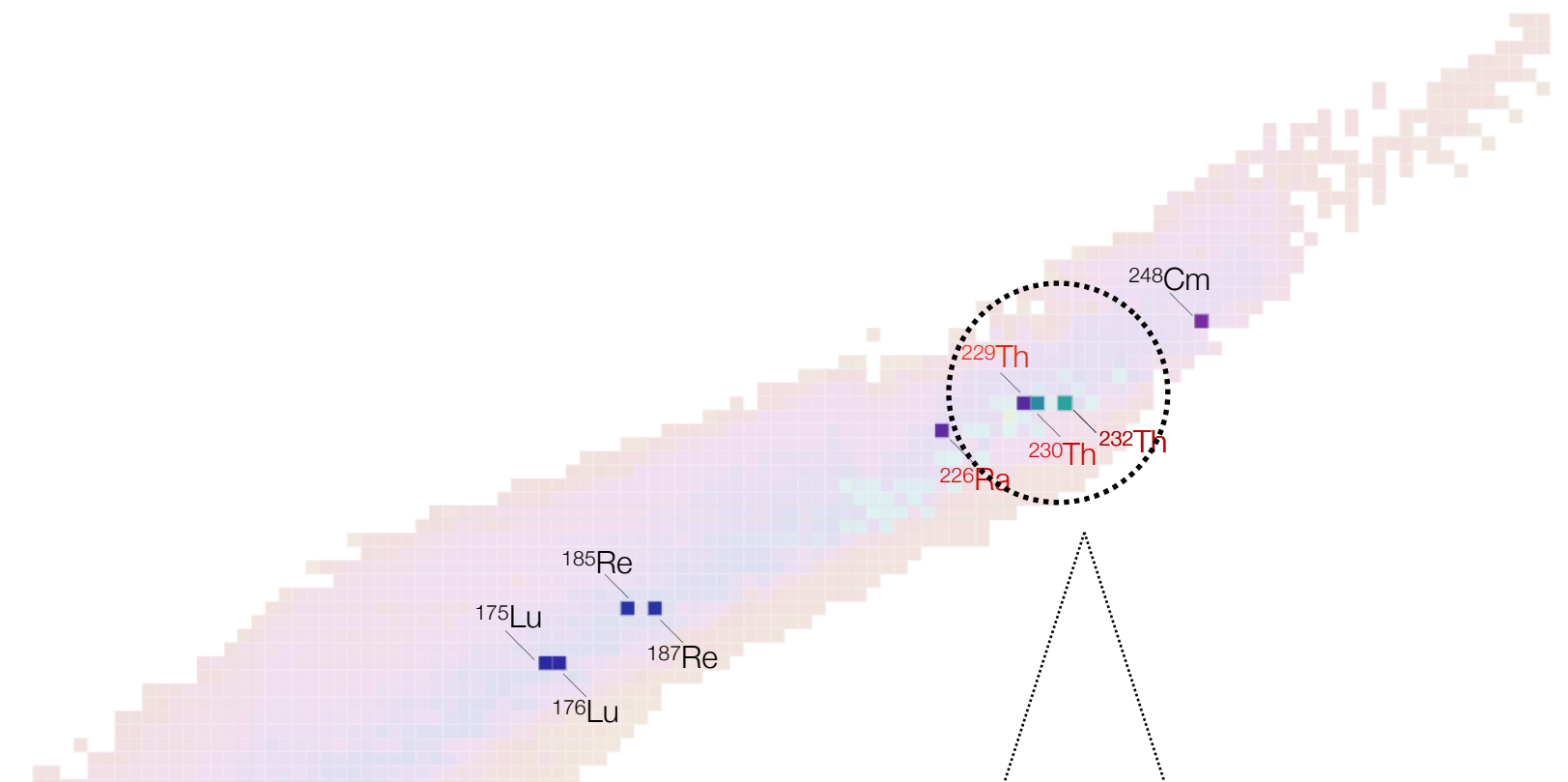
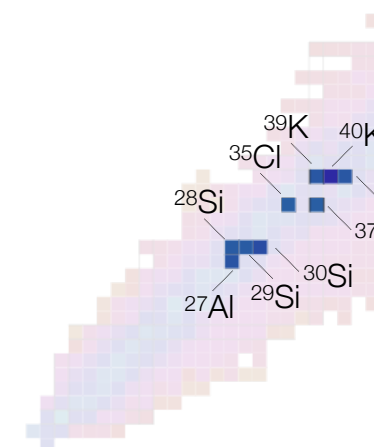
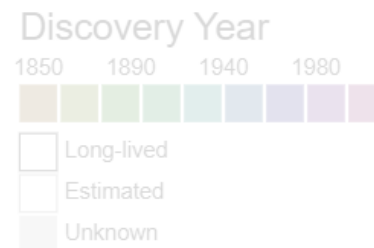
Extracting the Weinberg angle at low Q and probing new physics.

Adapted from Nature 557, 207–211 (2018)

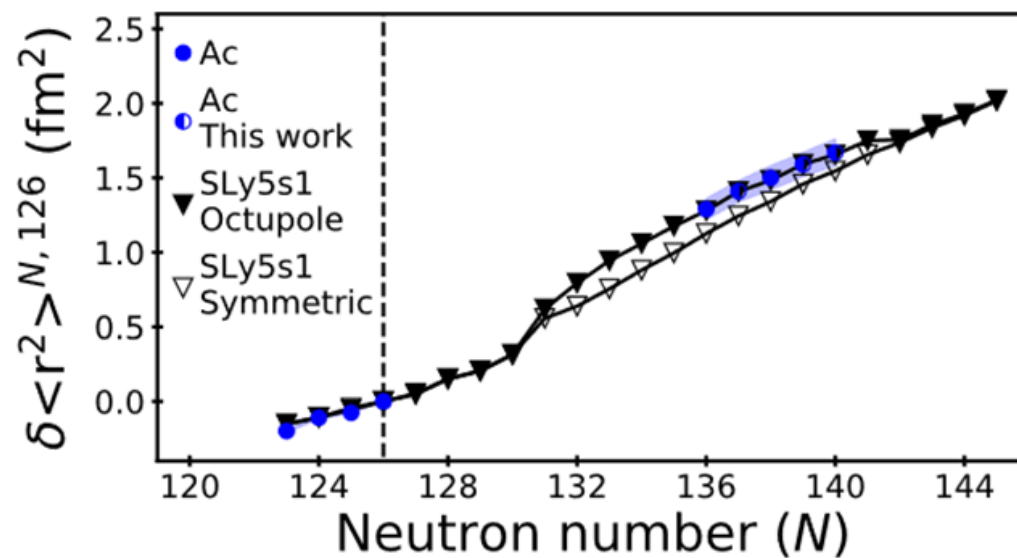


Science case

What's next?



PRC 100, 044321 (2019)



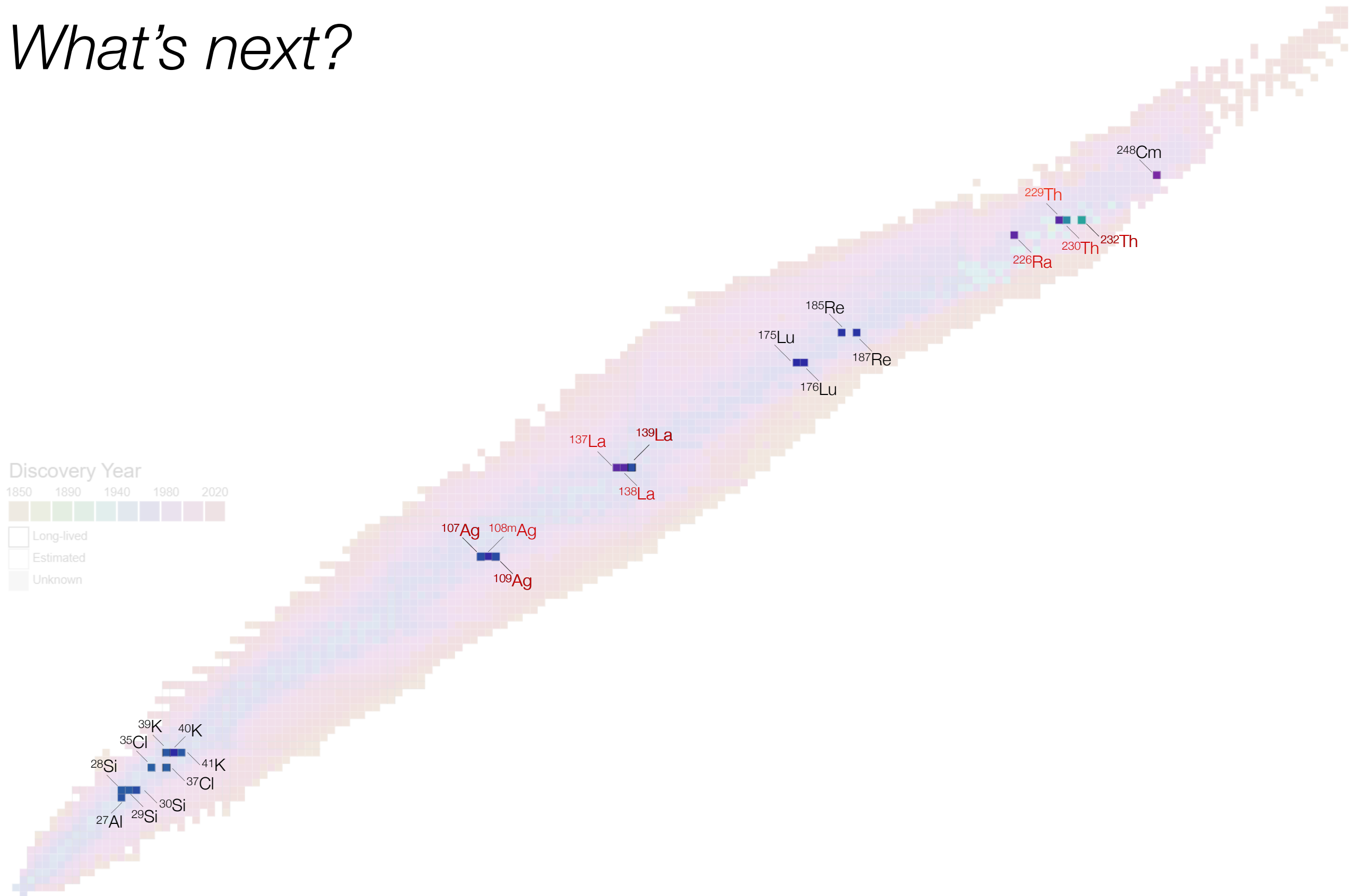
Strong octupole deformation in heavy mass isotopes can probe Beyond Standard Model (BSM) physics:

- Atomic Parity Violating (APV) effects in strongly octupole deformed isotopes are enhanced
- ^{229}Th nuclear clock is the best candidate to study variations of the fine structure constant

Phys. Rev. A 102, 052833 (2020)
arXiv:2407.17526

Science case

What's next?



Science case

What's next?

