

The Modular Blueprint of Flavor: Quarks and Leptons Entangled

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Consider the “Principal congruence group of level N”

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The Quotient group gives the “Finite Modular Group”

$$\Gamma_N \equiv SL(2, \mathbb{Z}) / \pm \Gamma(N) \text{ or } \Gamma'_N \equiv SL(2, \mathbb{Z}) / \Gamma(N)$$

N	2	3	4	5
Γ_N	S_3	A_4	S_4	A_5
Γ'_N	S_3	$A'_4 \equiv T'$	S'_4	A'_5

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Fields transform as $\phi(\tau) \rightarrow (c\tau + d)^{-k_\phi} \rho_\phi(\gamma) \phi(\tau)$
 Yukawa couplings are in multiplets transforming as $Y(\tau) \rightarrow (c\tau + d)^{k_Y} \rho_Y(\gamma) Y(\tau)$

Weights $\swarrow \searrow$ Irreps of Γ_N (Γ'_N)

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Derive the model's predictions for the observables and compare them with experimental data to constrain those parameters

Experimental inputs

For instance in the following examples

Quark sector

Observable	Central value	$\pm 1\sigma$
$m_u/m_c (10^{-3})$	1.93	0.60
$m_c/m_t (10^{-3})$	2.82	0.12
$m_d/m_s (10^{-2})$	5.05	0.62
$m_s/m_b (10^{-2})$	1.82	0.10
$m_t [\text{GeV}]$	87.5	2.1
$m_b [\text{GeV}]$	0.97	0.01

Observable	Central value	$\pm 1\sigma$
$\sin^2 \theta_{12}^q (10^{-2})$	5.08	0.03
$\sin^2 \theta_{13}^q (10^{-5})$	1.22	0.09
$\sin^2 \theta_{23}^q (10^{-3})$	1.61	0.05
δ_{CKM}/π	0.385	0.017

Values of quark masses and mixings renormalized around 2×10^{16} GeV, assuming the supersymmetry breaking scale $M_{\text{SUSY}} = 10$ TeV and $\tan \beta = 10$

Lepton sector

Observable	Central value	$\pm 1\sigma$
$m_e/m_\mu (10^{-3})$	4.74	0.04
$m_\mu/m_\tau (10^{-2})$	5.88	0.05
$\delta m^2/ \Delta m^2 (10^{-2})$ (NO)	2.95	0.06
$\delta m^2/ \Delta m^2 (10^{-2})$ (IO)	2.99	0.07
$m_\tau [\text{GeV}]$	1.293	0.007
$\delta m^2 (10^{-5} \text{ eV}^2)$	7.37	0.17
$ \Delta m^2 (10^{-3} \text{ eV}^2)$ (NO)	2.495	0.020
$ \Delta m^2 (10^{-3} \text{ eV}^2)$ (IO)	2.465	0.020

We choose NO, currently preferred by global analysis

Observable	Central value	$\pm 1\sigma$
$\sin^2 \theta_{12}^\ell (10^{-1})$	3.03	0.14
$\sin^2 \theta_{13}^\ell (10^{-2})$	2.23	0.05
$\sin^2 \theta_{23}^\ell (10^{-1})$ (NO)	4.73	0.24
$\sin^2 \theta_{23}^\ell (10^{-1})$ (IO)	5.45	0.23
δ_{PMNS}/π (NO)	1.20	0.22
δ_{PMNS}/π (IO)	1.48	0.12

We prefer to use $\sin^2 \theta_{23}^\ell = 0.5$ with 1σ error given by 1/6 of the 3σ allowed region because of the octant ambiguity

effects on the neutrino mass ratios and mixing angles known to be small to a good approximation

We prefer not to use δ_{CP} because of the large uncertainty

S. Antusch et al., J. High Energy Phys. 03 (2005) 024

S. Antusch and V. Maurer, J. High Energy Phys. 11 (2013) 115.

Parameters and Observables

Parameters: N_{par}

The modulus τ

$$\mathbf{P}_q^{up} \quad \mathbf{P}_q^{down}$$

To construct the Yukawa matrices (Y^u, Y^d)

$$\mathbf{P}_{leptons}^{charged} \quad \mathbf{P}_\nu$$

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Observables: N_{obs}

$$(\sin^2 \theta_{12}^q, \sin^2 \theta_{13}^q, \sin^2 \theta_{23}^q, \delta_{CP}^q)$$

$$(m_u, m_d, m_c, m_s, m_t, m_b)$$

$$(\sin^2 \theta_{12}^\ell, \sin^2 \theta_{13}^\ell, \sin^2 \theta_{23}^\ell, \Delta m^2, \delta m^2)$$

$$(m_e, m_\mu, m_\tau)$$

Unknowns $\longrightarrow (\delta_{CP}^\ell, m_1, m_{\beta\beta}, m_\beta)$

To test the model

$$\chi_{\text{leptons}}^2(\mathbf{P}_{\text{leptons}}) = \sum_{i=1}^8 \left(\frac{O_i^{\text{theo}}(\mathbf{P}_{\text{leptons}}) - O_i^{\text{exp}}}{\sigma_i^{\text{exp}}} \right)^2$$

$$\chi_{\text{quarks}}^2(\mathbf{P}_{\text{quarks}}) = \sum_{i=9}^{18} \left(\frac{O_i^{\text{theo}}(\mathbf{P}_{\text{quarks}}) - O_i^{\text{exp}}}{\sigma_i^{\text{exp}}} \right)^2$$

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The χ^2 distribution undefined,
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It is not hard to find in the literature sentences such as “the model is fine because all the observables are within their 3σ experimentally allowed regions” or where only the $\Delta\chi^2$ is mentioned without also looking at the minimum χ^2_{min} or where the number of d.o.f. is completely wrong

For instance

$$\left[\begin{array}{l} N_{\text{obs}} = 18 \\ N_{\text{par}} = 17 \end{array} \right] \longrightarrow N_{\text{dof}} = 1$$

p-value can be calculated depending on the χ^2_{min}

- $\Delta\chi^2 = 1$ corresponds to the **68.27% confidence region (1σ)**
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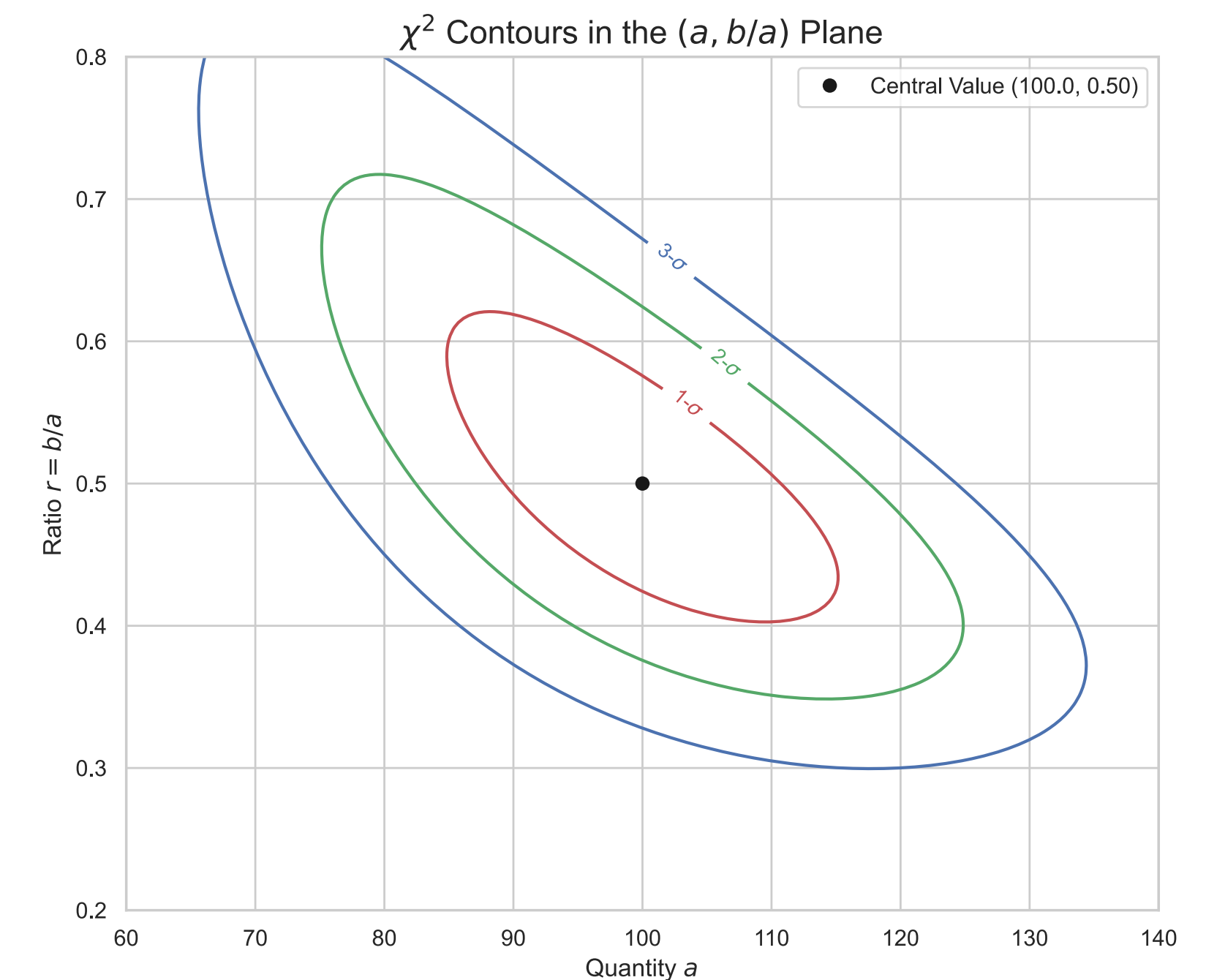
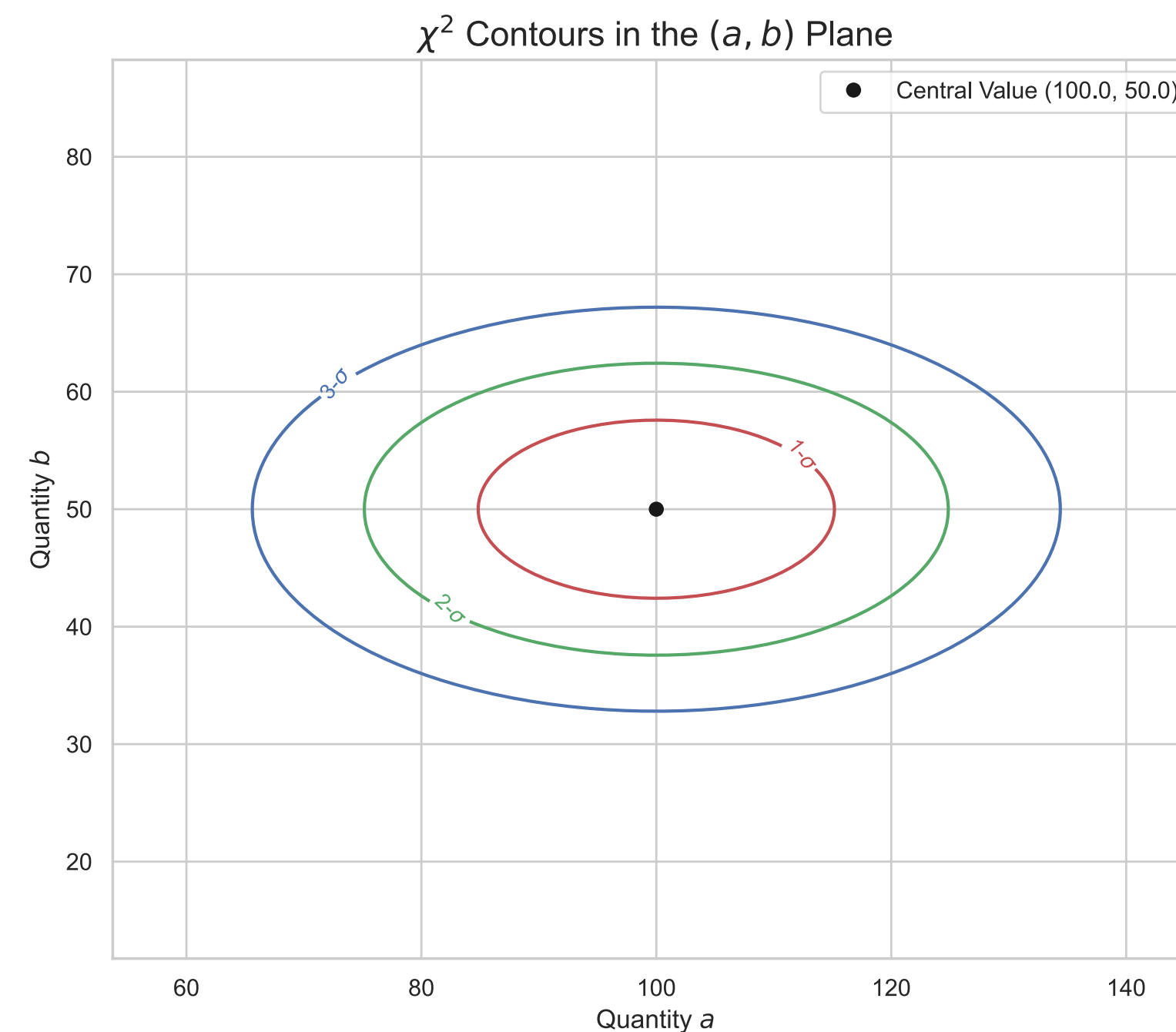
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Quark-lepton correlations (in a non-GUT framework) are a distinctive feature of models with modular symmetries and can help help clarify how the model works

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F. Feruglio et al., e-Print: [2505.20395](#)

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2O - the binary
octahedral group

14

$$N_{\text{dof}} > 0$$

S_3 modular invariant model in a $N = 1$ rigid SUSY

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Field	S_3 irrep	Modular weight k
Leptonic sector		
E_1^c	1	6
E_2^c	1'	4
E_3^c	1'	2
D_ℓ	2	4
ℓ_3	1'	2
H_d	1	0
H_u	1	0
Quark sector		
$Q_D = (Q_1, Q_2)^T$	2	4
Q_3	1'	2
u_1^c	1	6
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Tensor product decomposition rules for the group S_3 .

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$$W_\nu = \frac{1}{\Lambda} \Big[g_{8s} (D_\ell D_\ell)_1 Y_1^{(8)} + g'_8 (D_\ell D_\ell)_2 \cdot D_{2a}^{(8)} + \tilde{g}'_8 (D_\ell D_\ell)_2 \cdot D_{2b}^{(8)} \\ + g_{p6} (D_\ell \ell_3)_2 \cdot D_2^{(6)} + g'_w (\ell_3 \ell_3) Y_1^{(4)} \Big] H_u H_u .$$

S_3 modular invariant model in a $N = 1$ rigid SUSY

Field	S_3 irrep	Modular weight k
Leptonic sector		
E_1^c	1	6
E_2^c	1'	4
E_3^c	1'	2
D_ℓ	2	4
ℓ_3	1'	2
H_d	1	0
H_u	1	0
Quark sector		
$Q_D = (Q_1, Q_2)^T$	2	4
Q_3	1'	2
u_1^c	1	6
u_2^c	1'	4
u_3^c	1'	2
d_1^c	1	6
d_2^c	1'	4
d_3^c	1'	2
Higgs sector		
H_u	1	0
H_d	1	0

Follow the irrep structure of *JHEP* 09 (2023) 043

Tensor product decomposition rules for the group S_3 .

$\otimes$	1	1'	2
1	1	1'	2
1'	1'	1	2
2	2	2	$\mathbf{1} \oplus \mathbf{1}' \oplus \mathbf{2}$

Parameters

$$\tau = \text{Re}(\tau) + \text{Im}(\tau) \qquad 2$$

$$(\alpha_u, \tilde{\alpha}_u, \beta_u, \tilde{\beta}_u, \alpha_{uD}, \beta_{uD}, \gamma_{uD}) \qquad 7$$

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 S_3 -invariant scalar product between two doublets

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chosen as 10^{-16} eV

$$W_\nu = \frac{1}{\Lambda} \left[g_{8s} (D_\ell D_\ell)_1 Y_1^{(8)} + g'_8 (D_\ell D_\ell)_2 \cdot D_{2a}^{(8)} + \tilde{g}'_8 (D_\ell D_\ell)_2 \cdot D_{2b}^{(8)} \right. \\ \left. + g_{p6} (D_\ell \ell_3)_2 \cdot D_2^{(6)} + g'_w (\ell_3 \ell_3) Y_1^{(4)} \right] H_u H_u .$$

$\nearrow$
 S_3 -invariant scalar product between two doublets

Parameters

$\tau = \text{Re}(\tau) + \text{Im}(\tau)$

2

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7

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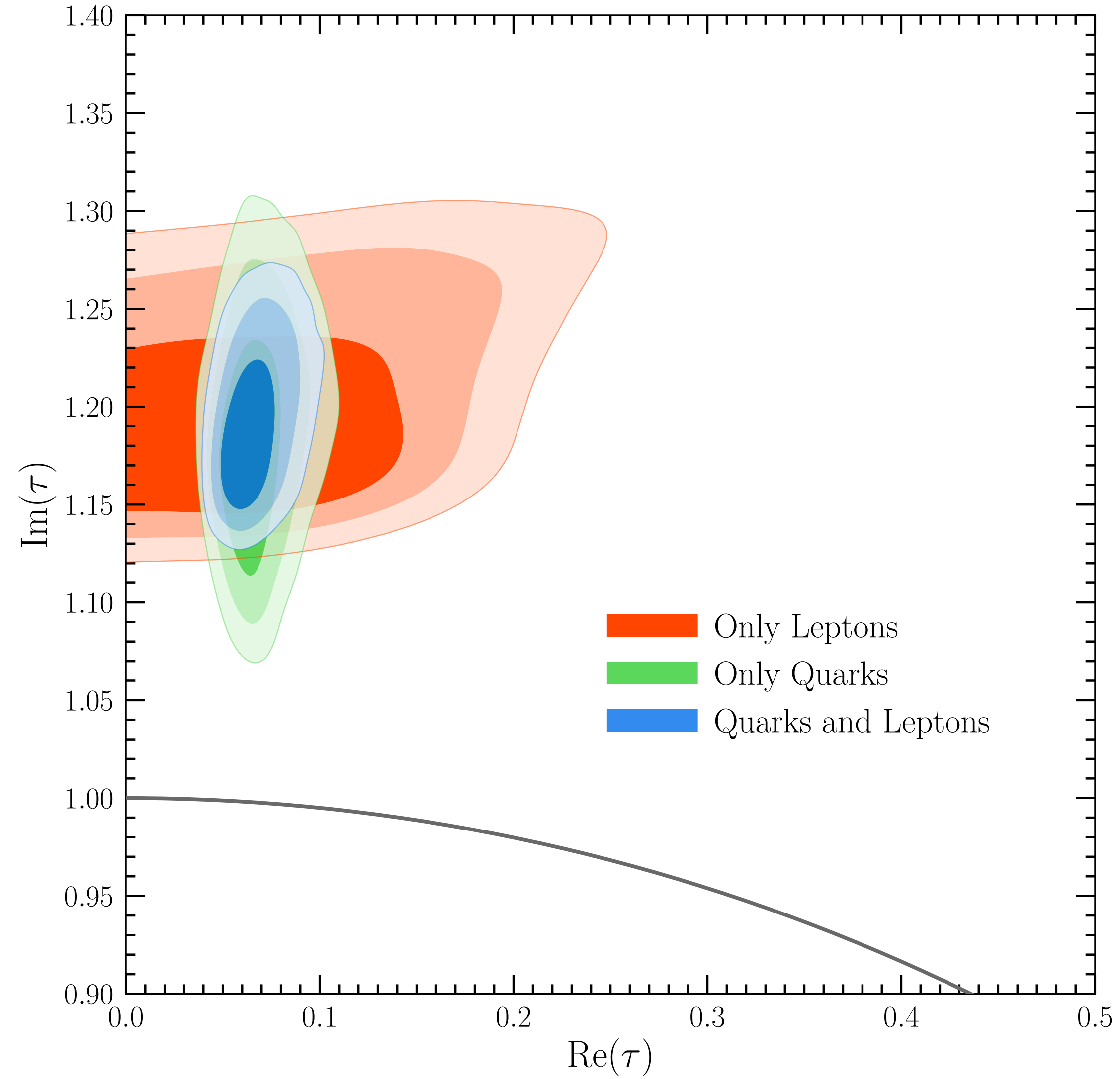
5

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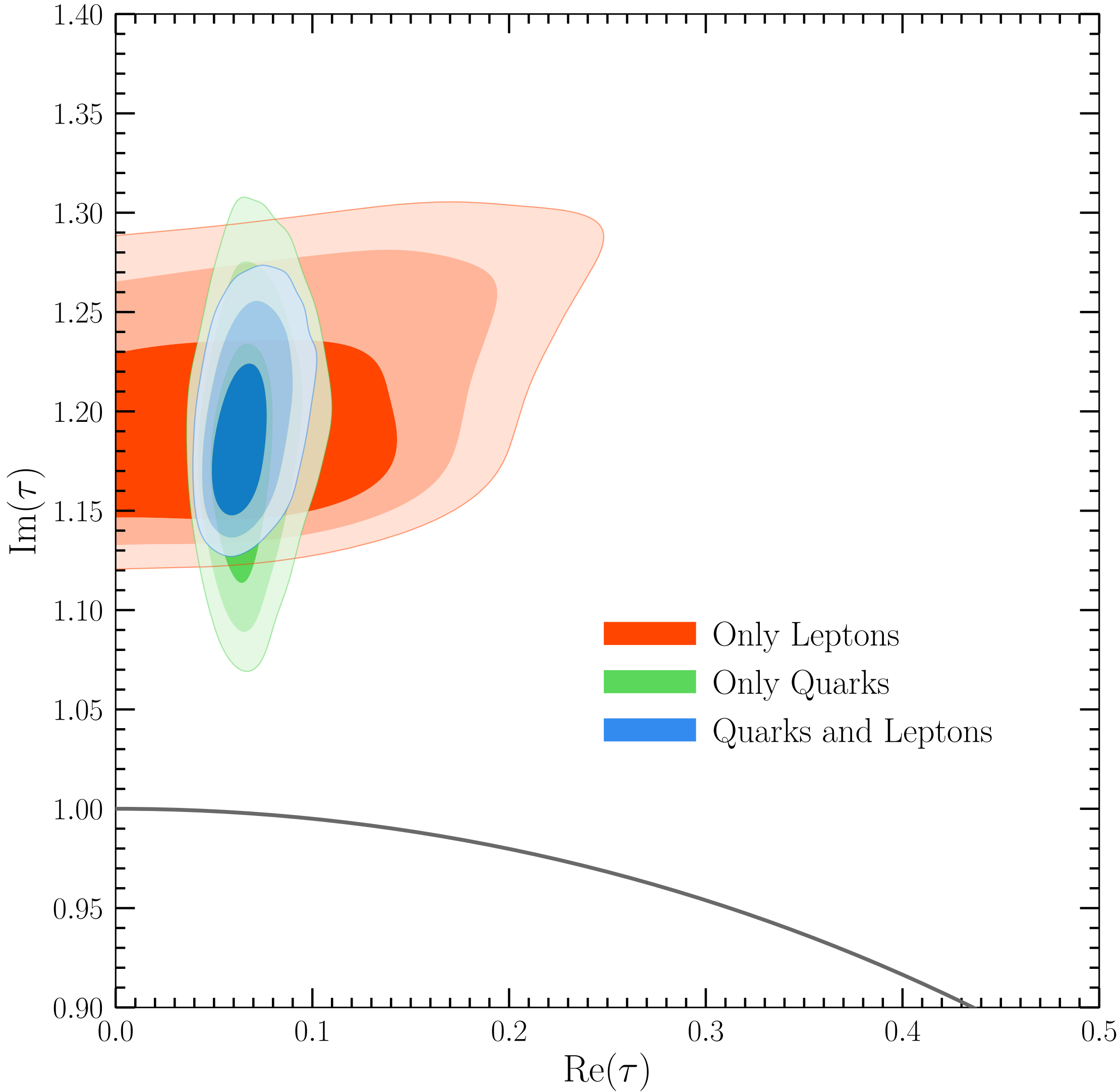
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S_3 unified quark-lepton model

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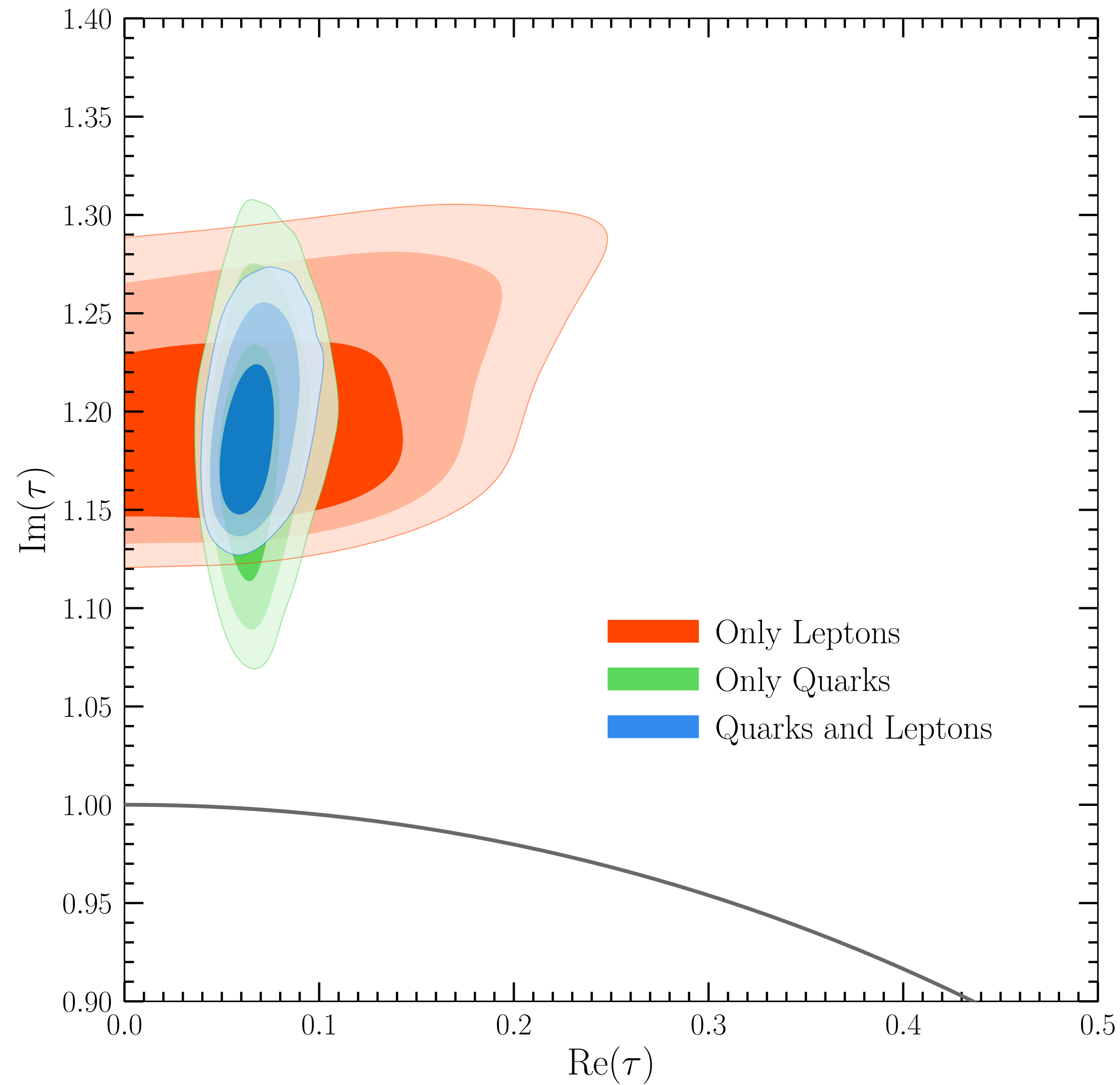


S_3 unified quark-lepton model



Modulus τ			
$\text{Re } \tau$	6.730×10^{-2}	$\text{Im } \tau$	1.210
Charged-lepton sector			
α	3.548×10^{-4}	α_T	9.634×10^{-3}
β	9.034×10^{-3}	β_T	7.819×10^{-2}
α_D	-2.192	β_D	2.847×10^{-1}
γ_D	7.466×10^{-1}		
Neutrino sector			
g_{8s}	8.446×10^{-8}	g_{p8}	1.202×10^{-7}
g_{pt8}	-2.837×10^{-7}	g_{p6}	2.100×10^{-8}
g_{pw}	-1.653×10^{-7}		
Up-quark sector			
α_u	-4.121×10^{-3}	α_{uT}	5.808×10^{-3}
β_u	-9.367×10^{-2}	β_{uT}	5.769×10^{-2}
α_{uD}	1.231×10^{-1}	β_{uD}	-7.267
γ_{uD}	6.706		
Down-quark sector			
α_d	6.584×10^{-3}	α_{dT}	-3.118×10^{-2}
β_d	5.120×10^{-2}	β_{dT}	-5.264×10^{-2}
α_{dD}	-1.143×10^{-1}	β_{dD}	-1.361
γ_{dD}	-6.709×10^{-1}		

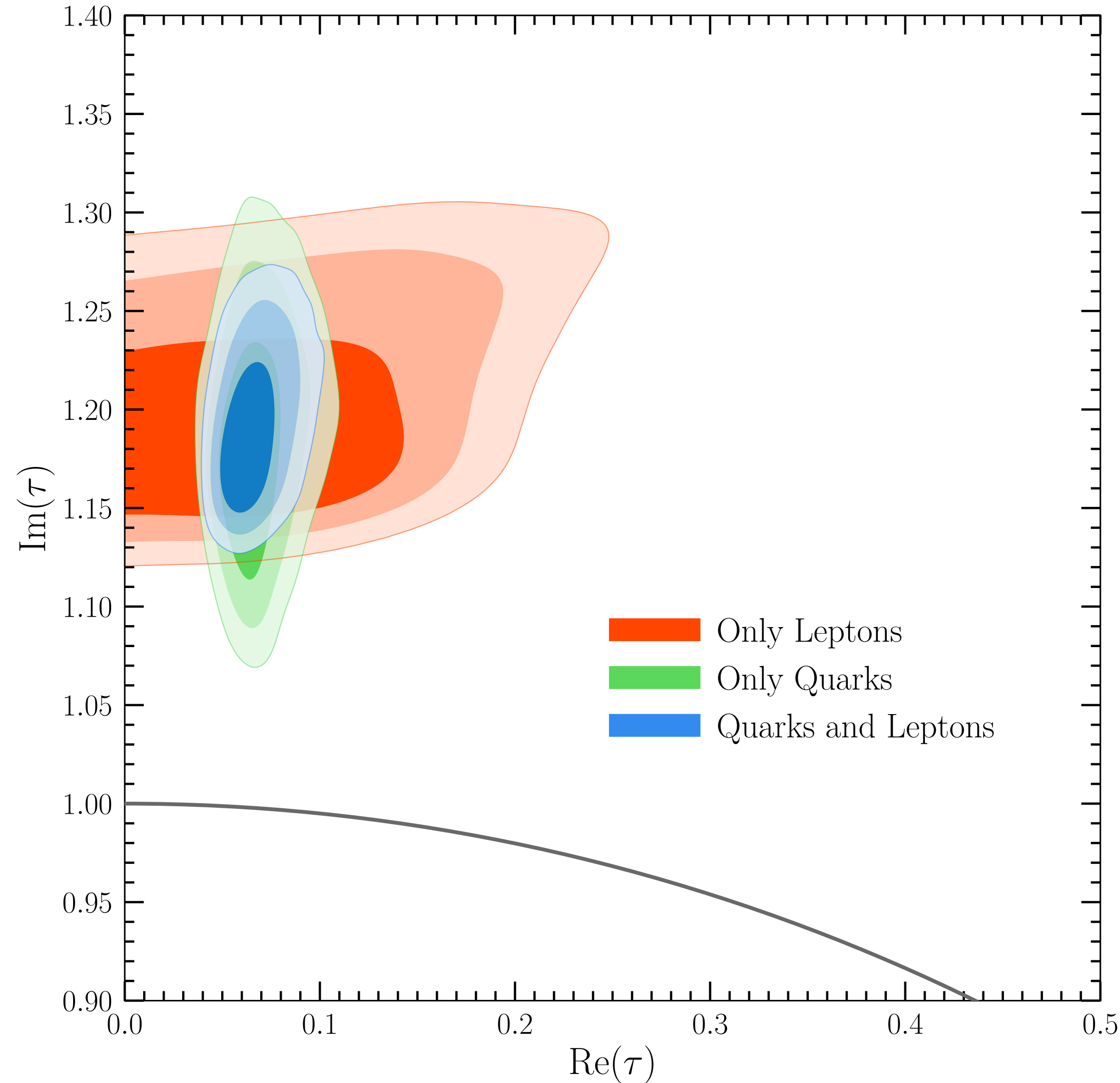
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The model fits the data but allowed regions are not well defined: they merely denotes $\Delta\chi^2 = 1, 4, 9$ and give an idea of where the modulus can be located and of the gradient of the χ^2 around the “only acceptable $\chi^2 = 0$ ” point

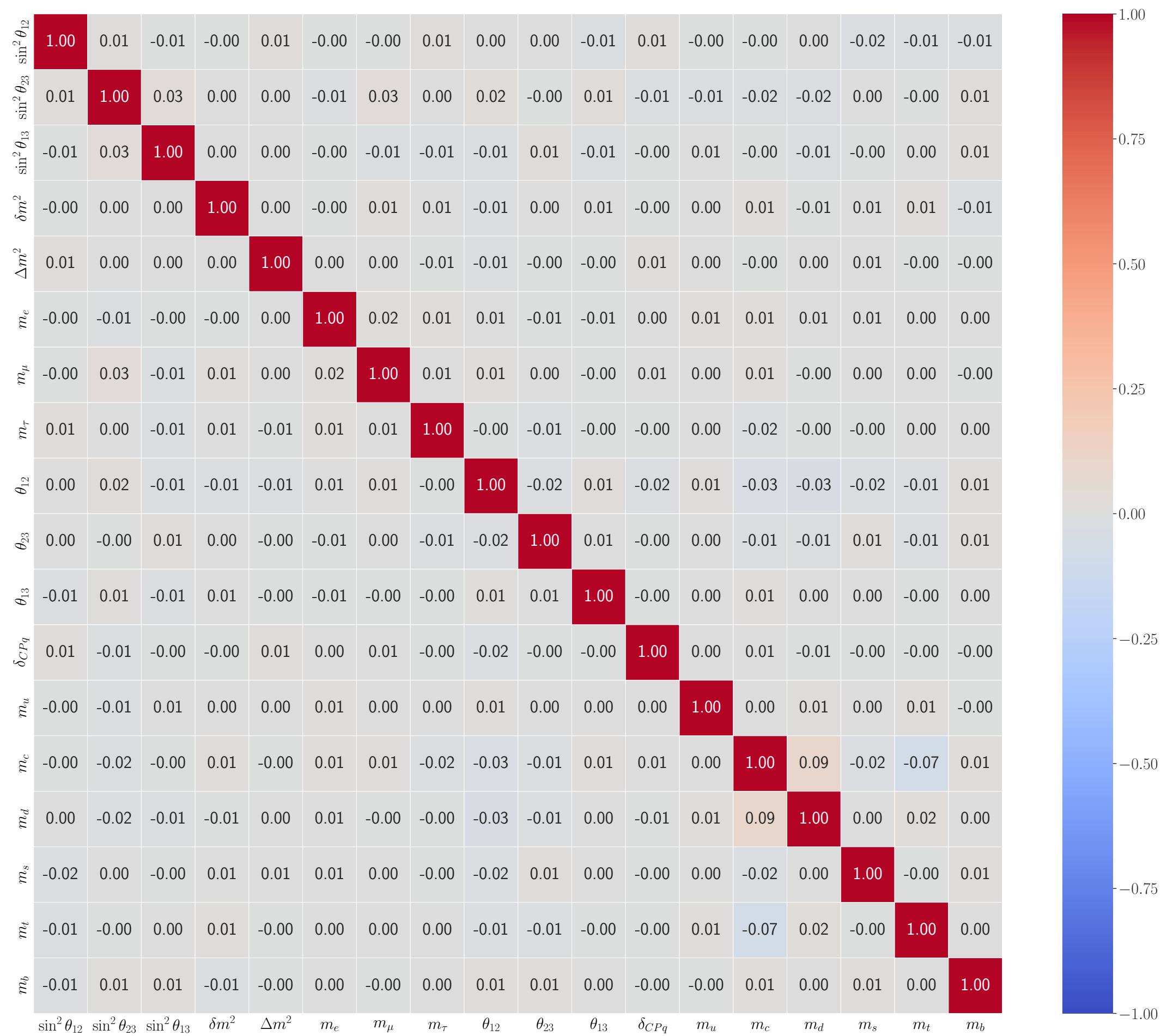
Interesting (the model being overfitted)

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- 1) As one might expect, the observables are almost completely uncorrelated

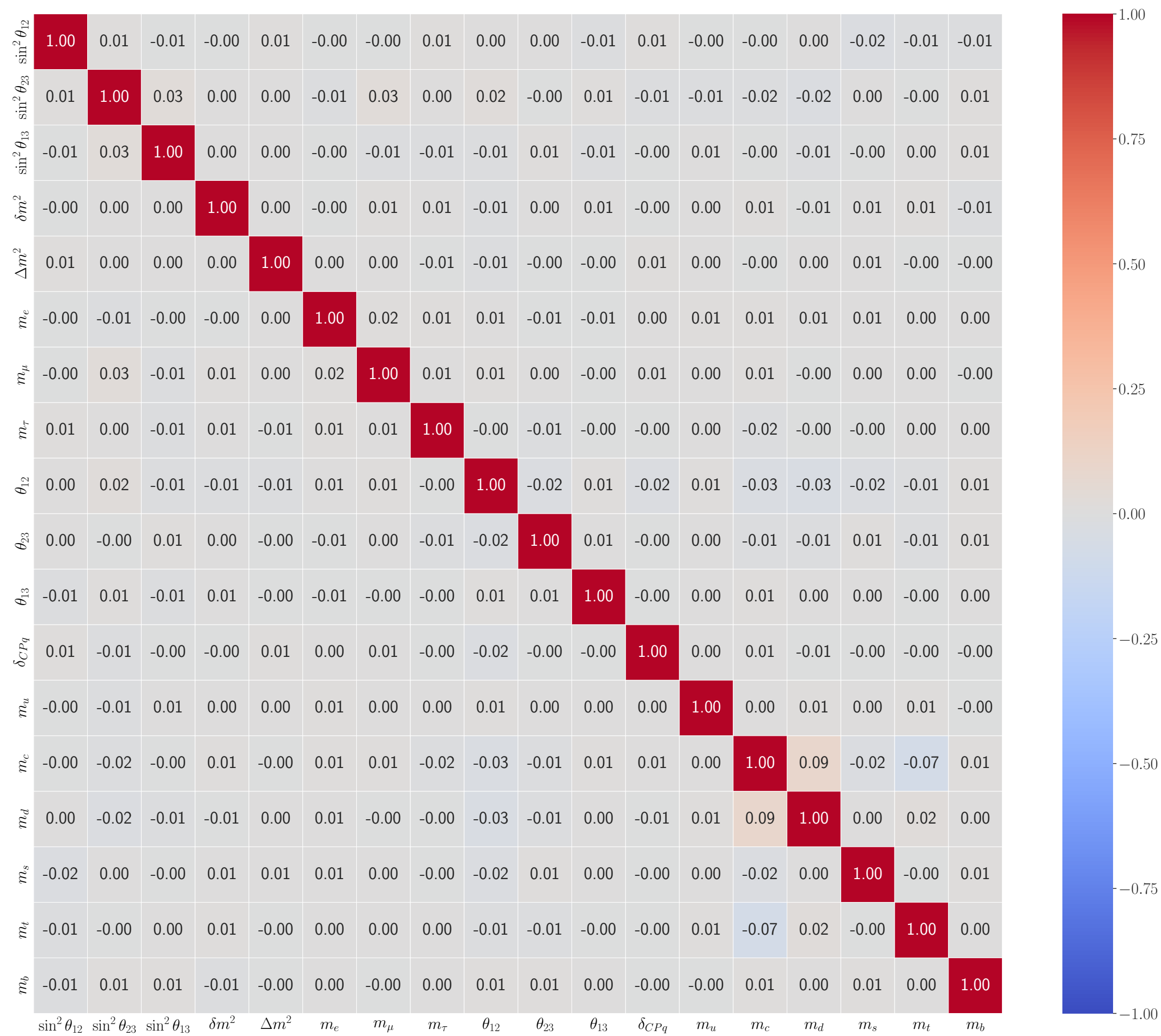
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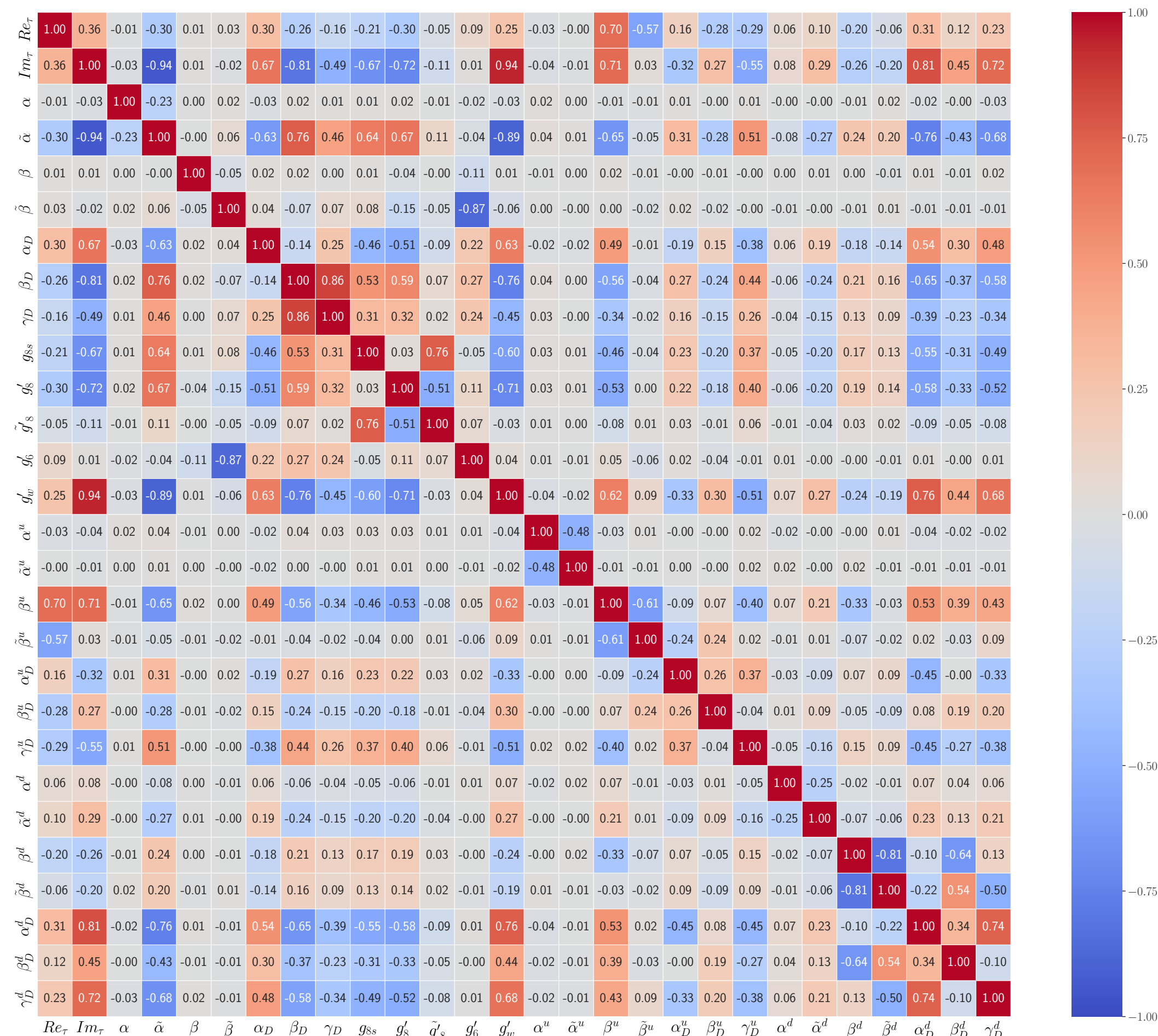
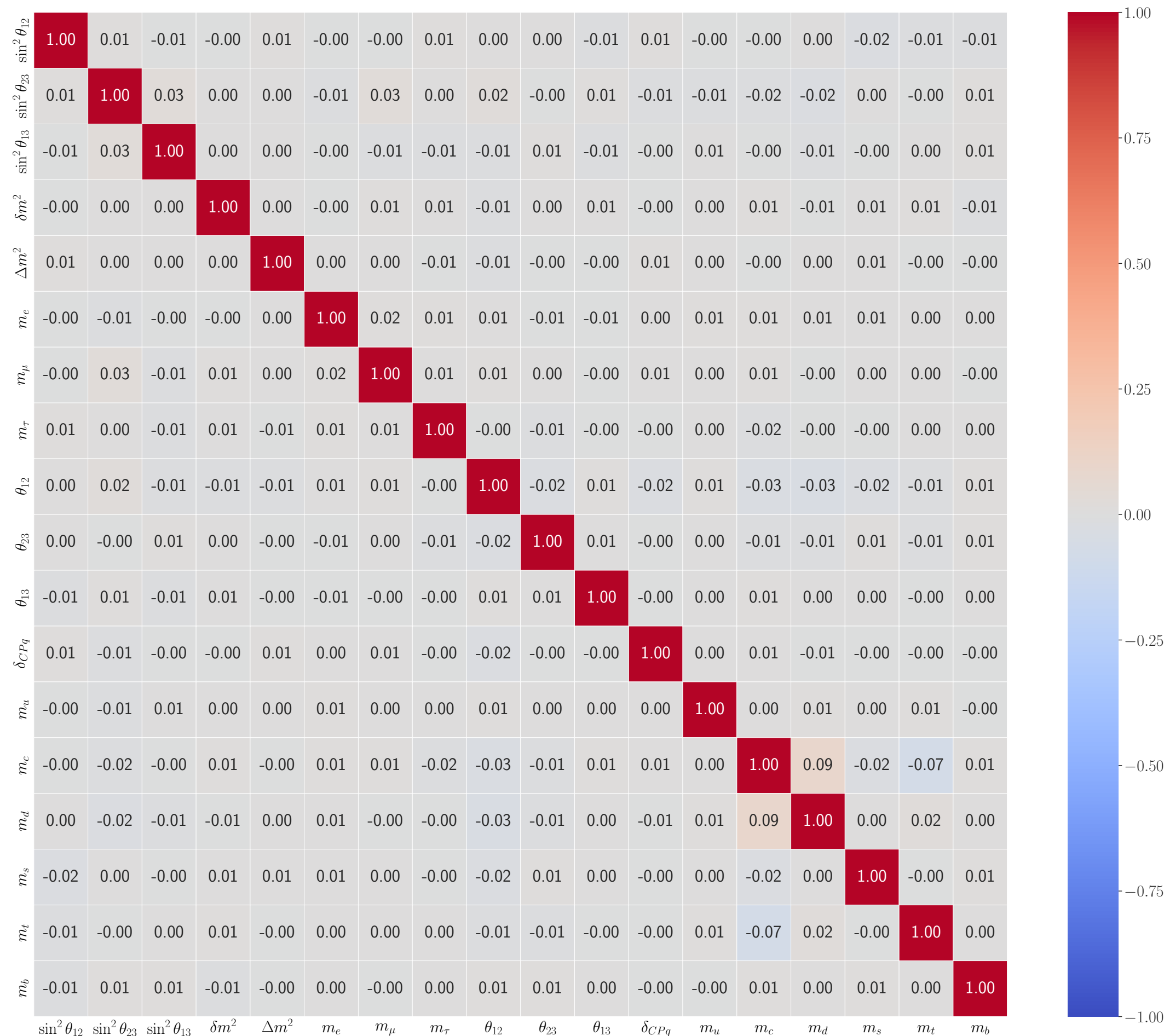
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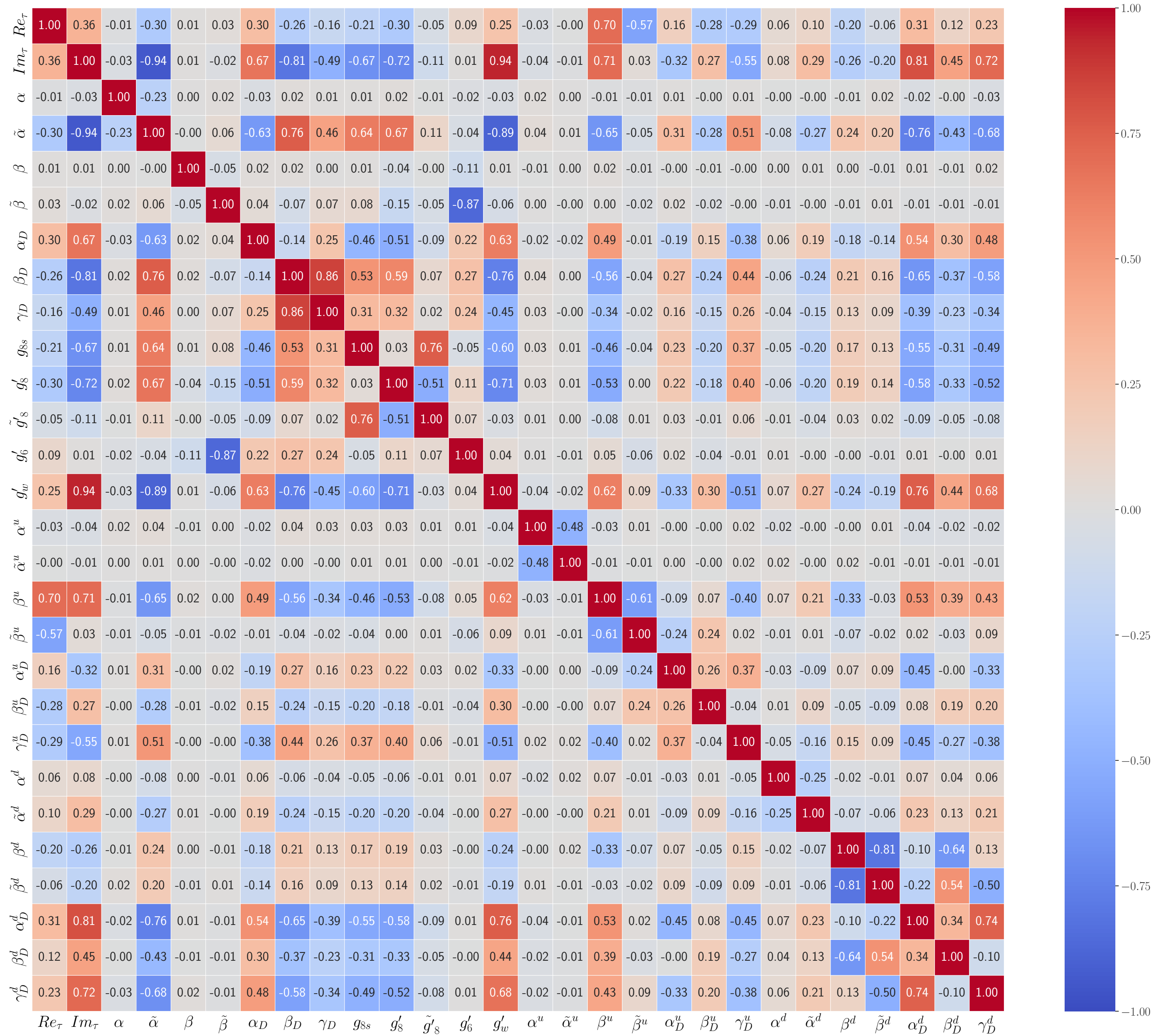
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- 2) Model parameters, however, show significant correlations



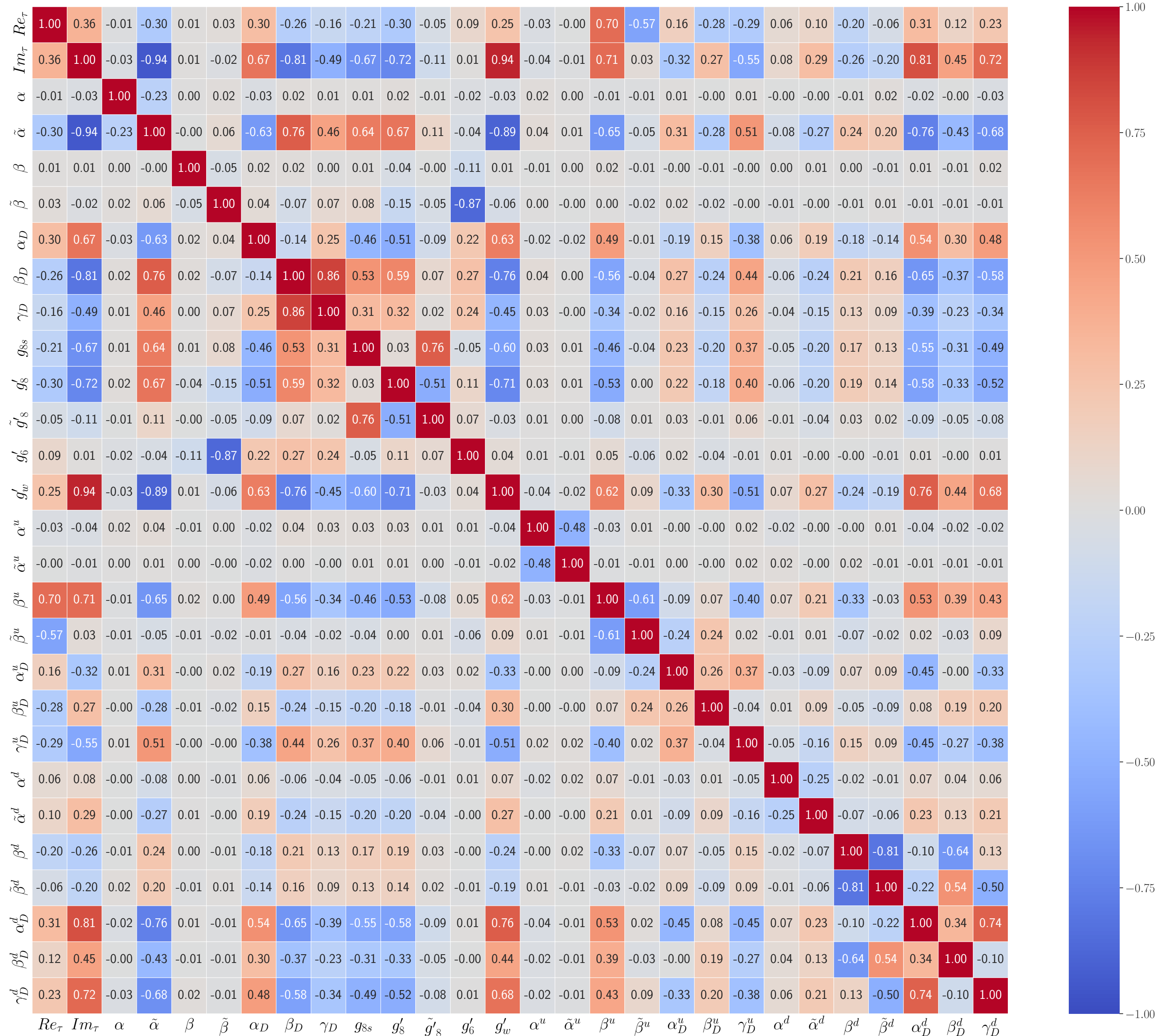
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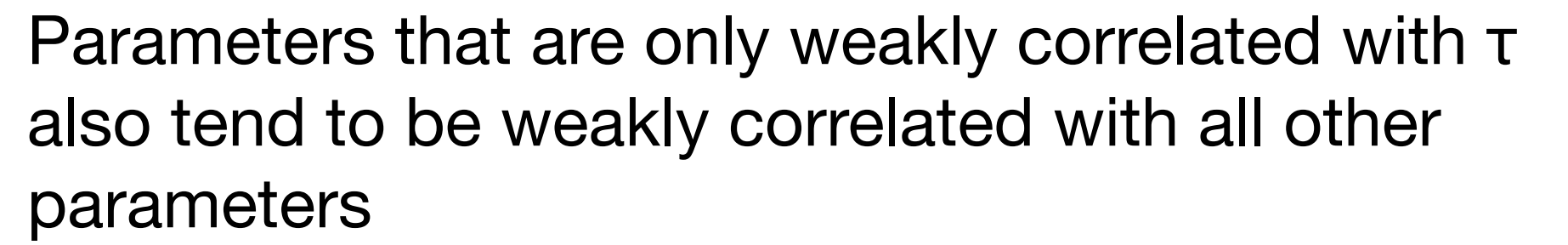
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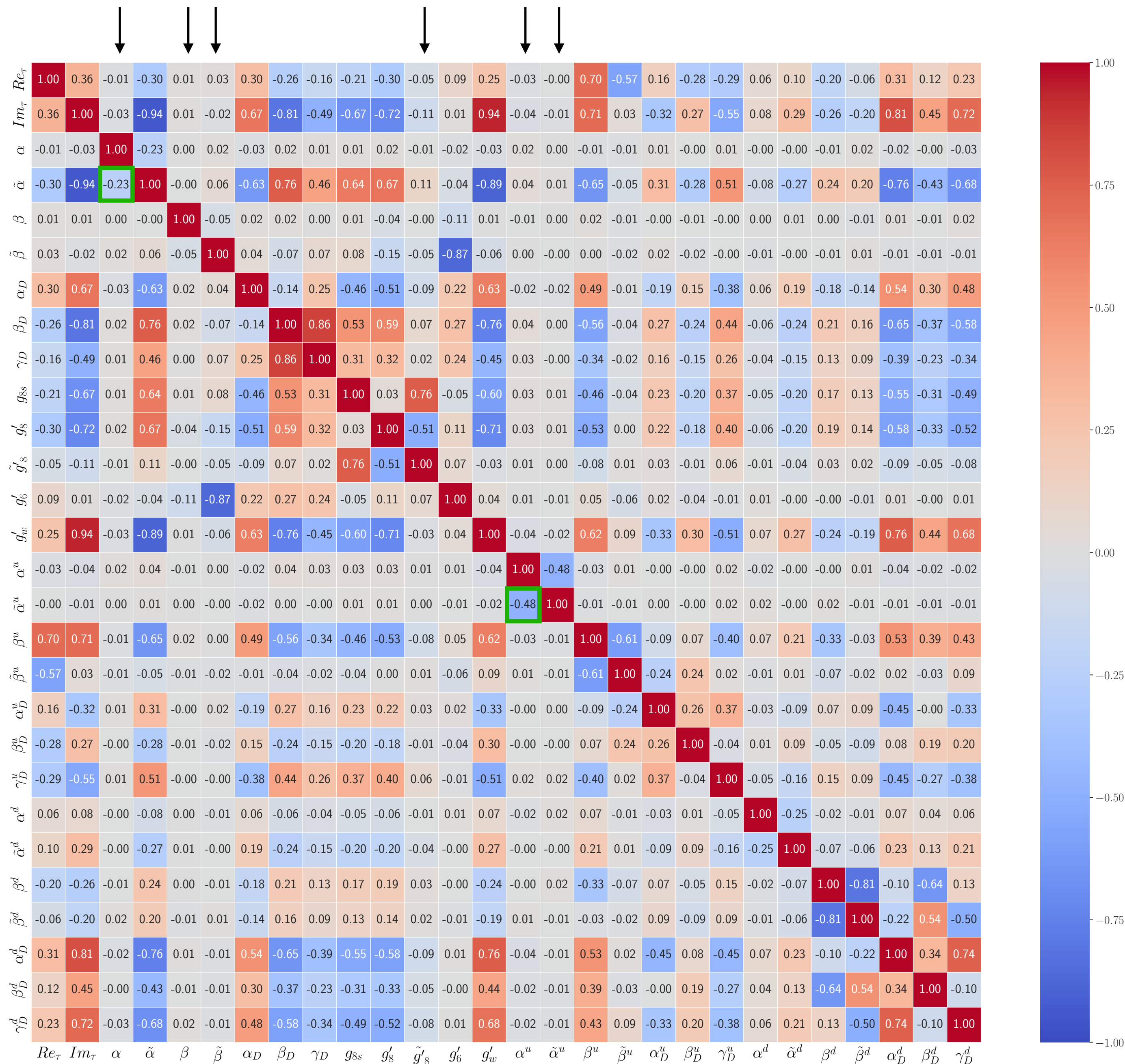


Correlations tend to be governed by the modulus



[illegible]

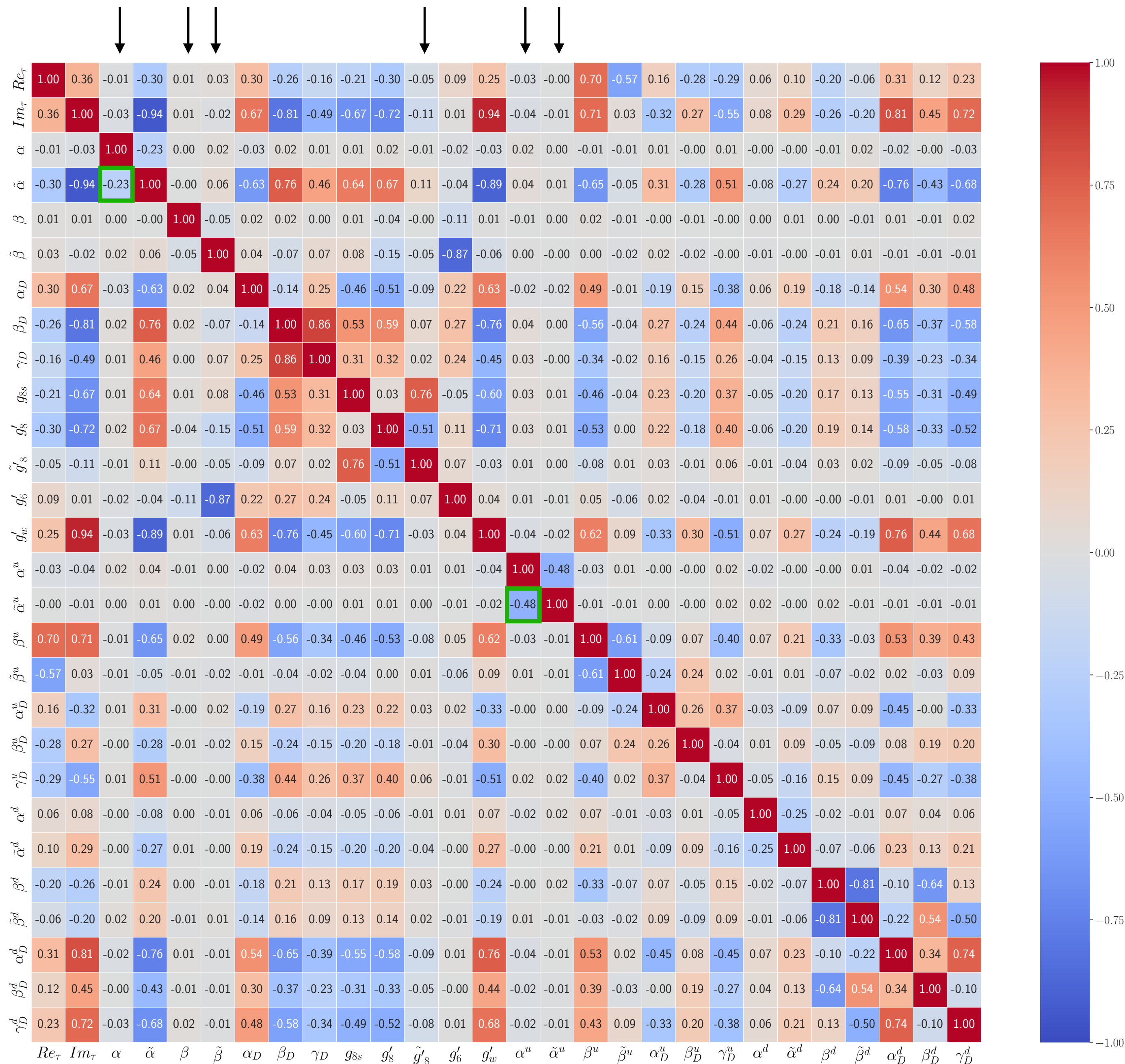
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Parameters that are only weakly correlated with τ also tend to be weakly correlated with all other parameters

$(\alpha, \tilde{\alpha})$ and $(\alpha_u, \tilde{\alpha}_u)$ are anti correlated because they appear in the same matrix element

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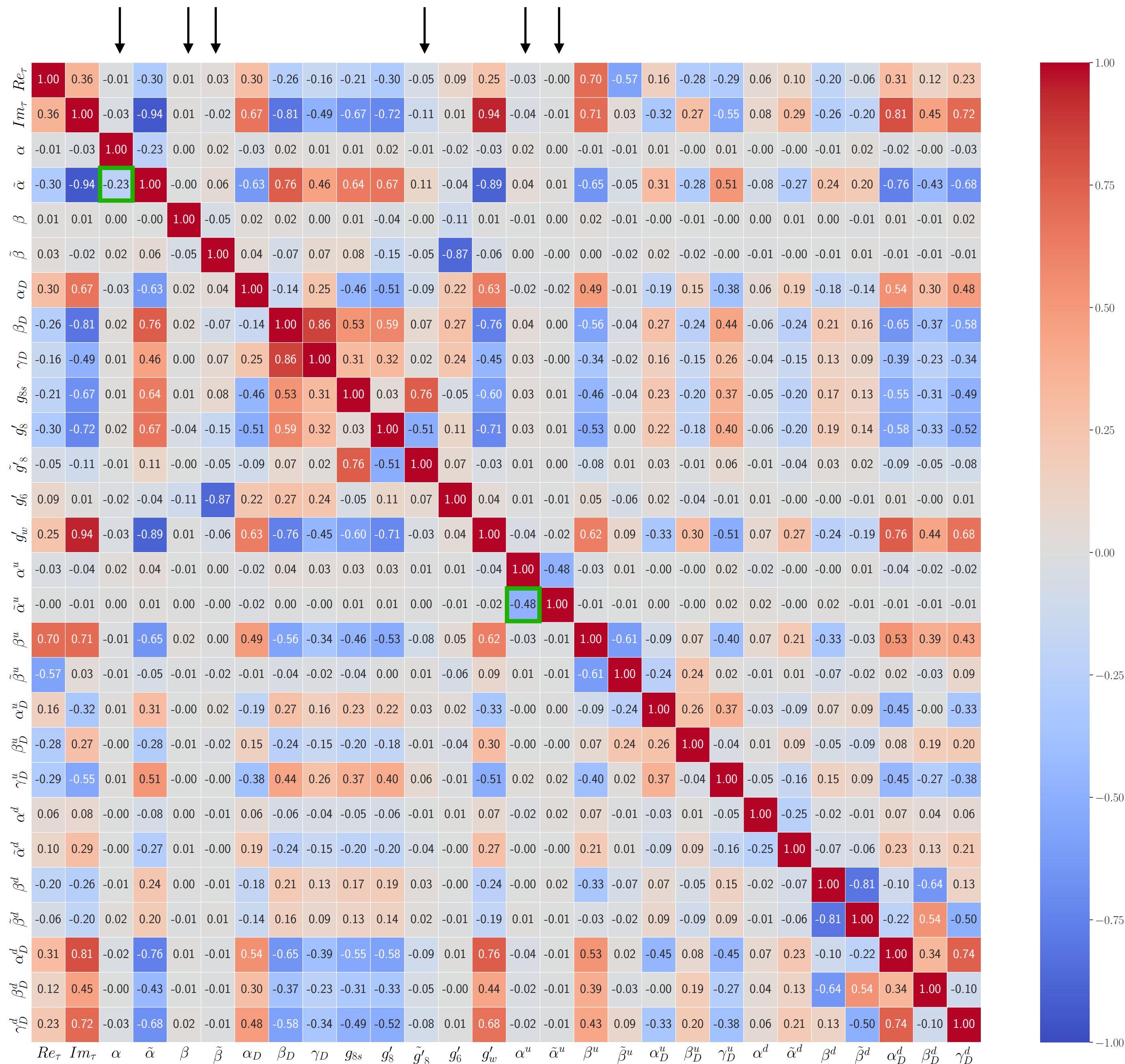


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These correlations reflect the non-linear modular constraints and the internal geometry of the parameter space, and are not a sign of fine-tuning

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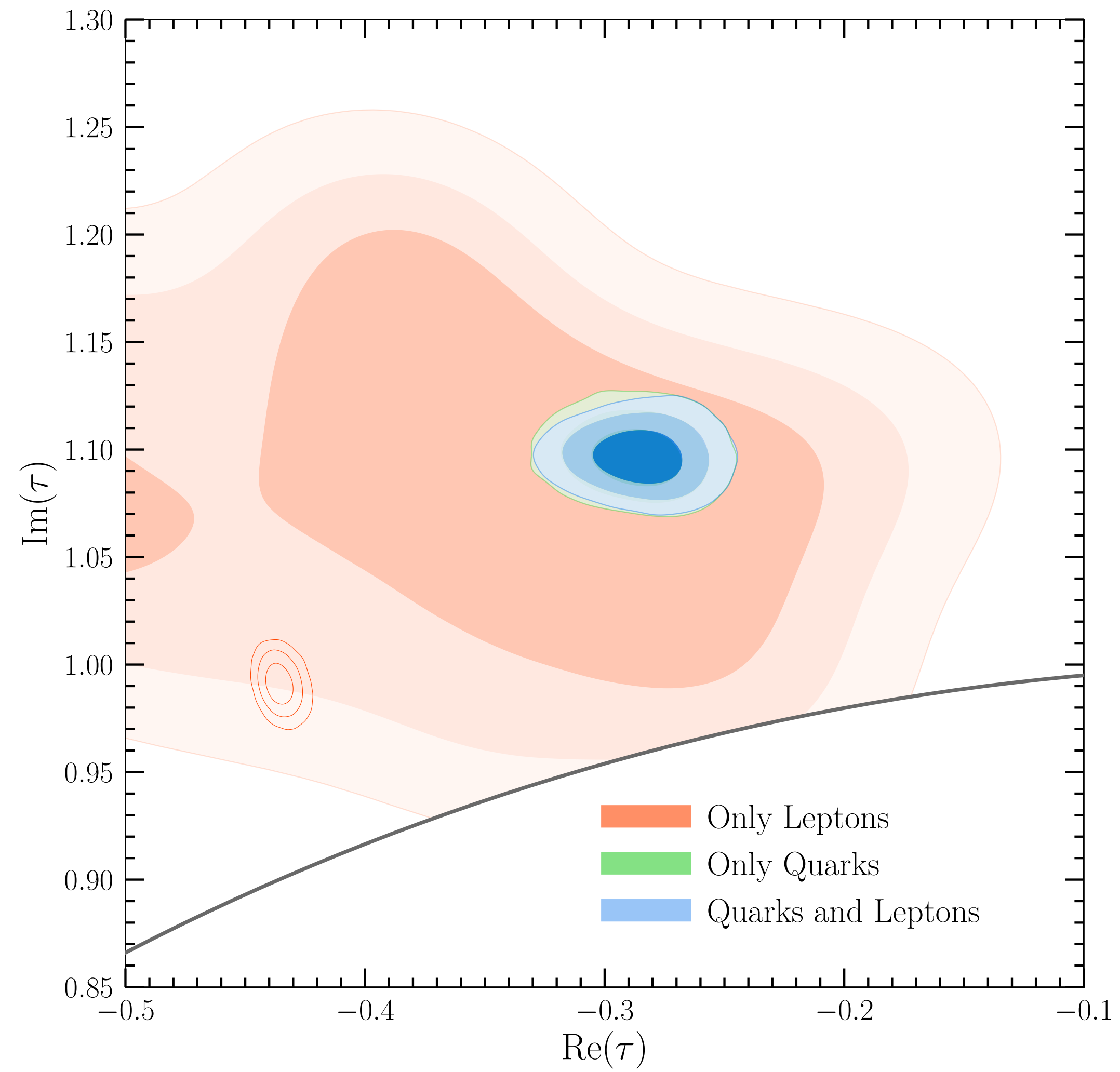
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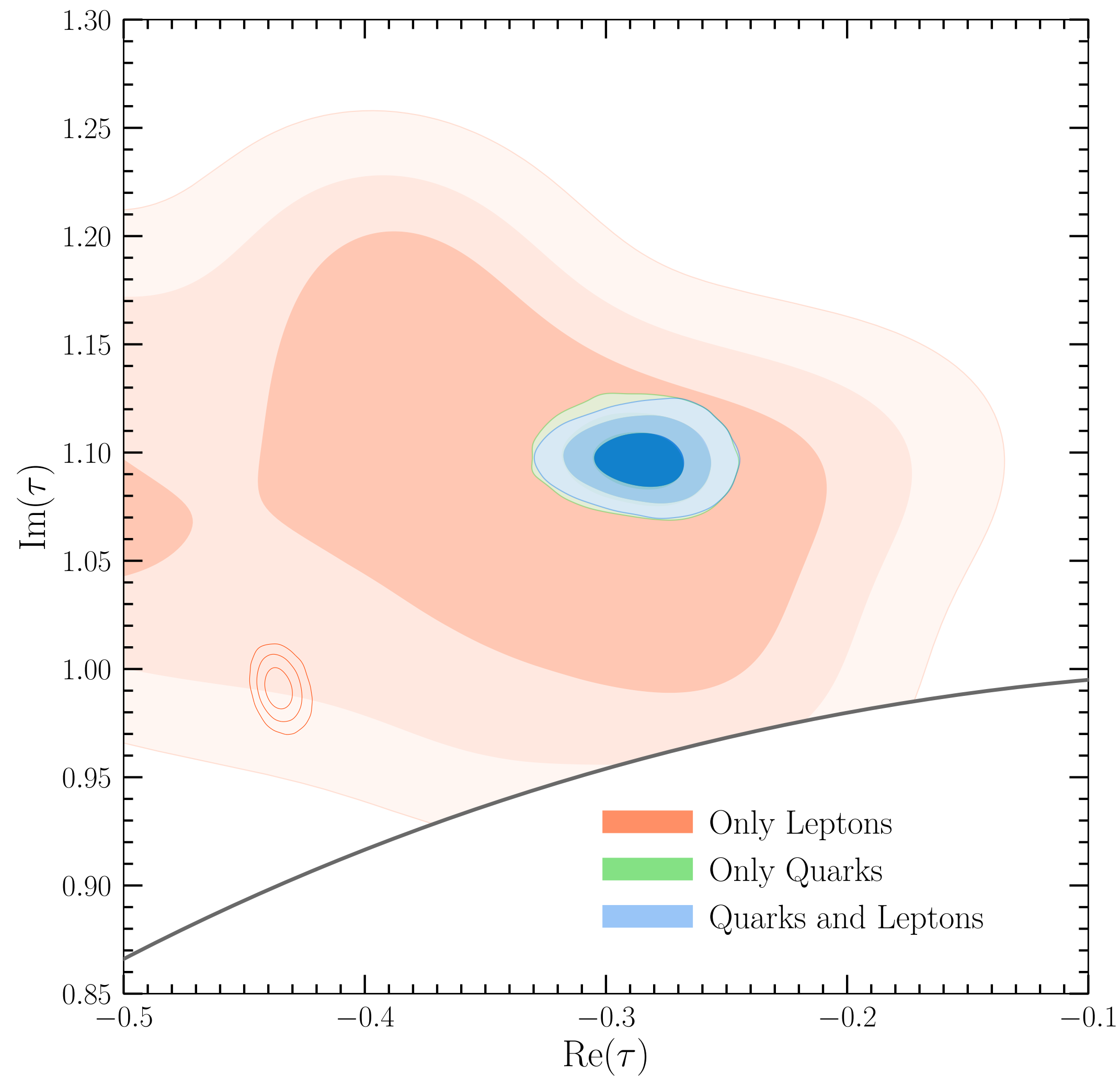


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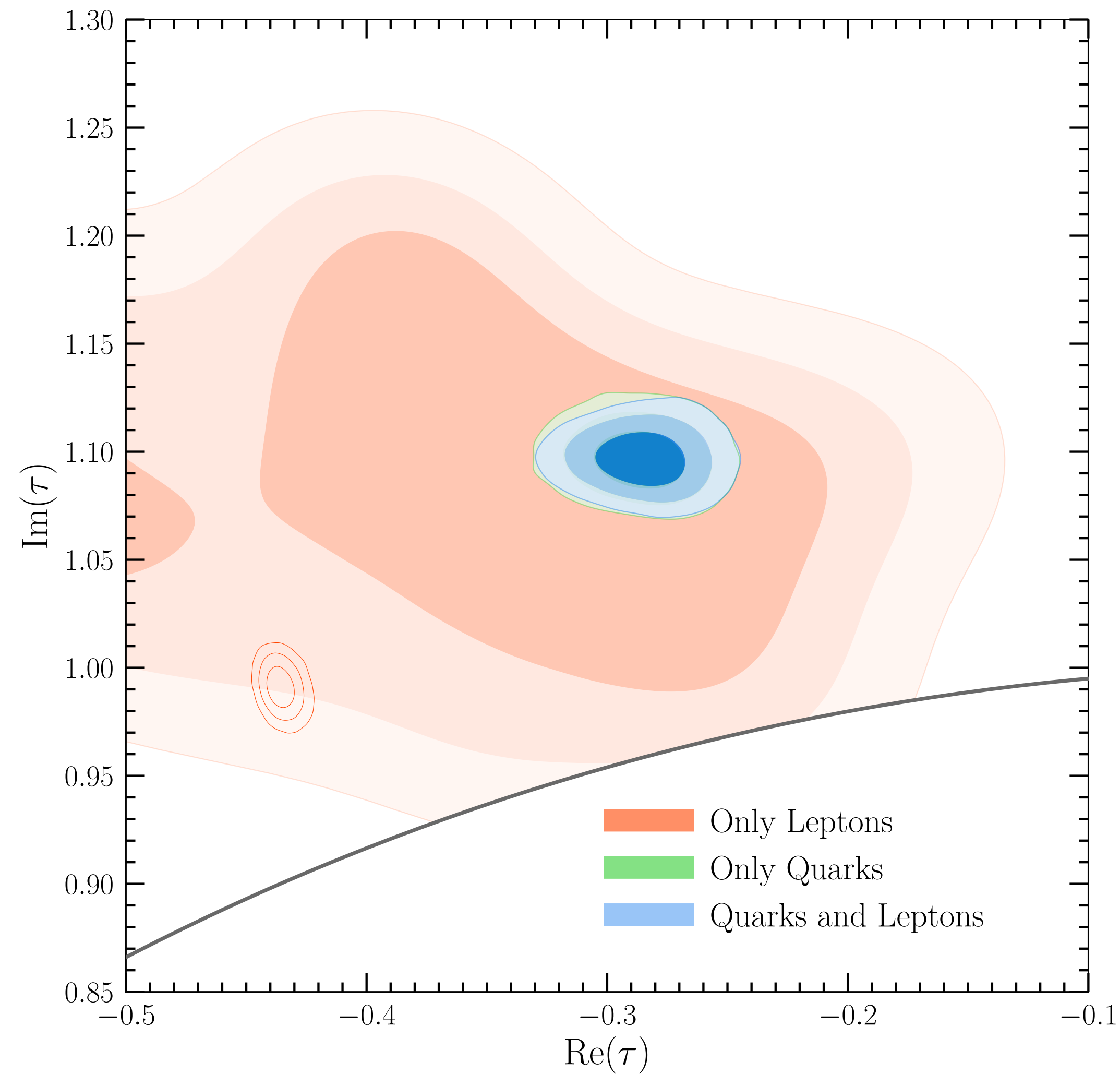
Also in the limiting case $N_{\text{dof}} = 0$ the χ^2 analysis is ill-defined and we can only draw qualitative conclusions from the analysis. The model can be "accepted" only if $\chi^2_{\text{min}} = 0$



Lepton allowed region for τ is much larger than the quark corresponding one, as expected and the combination in this plane essentially coincides with the quark allowed region

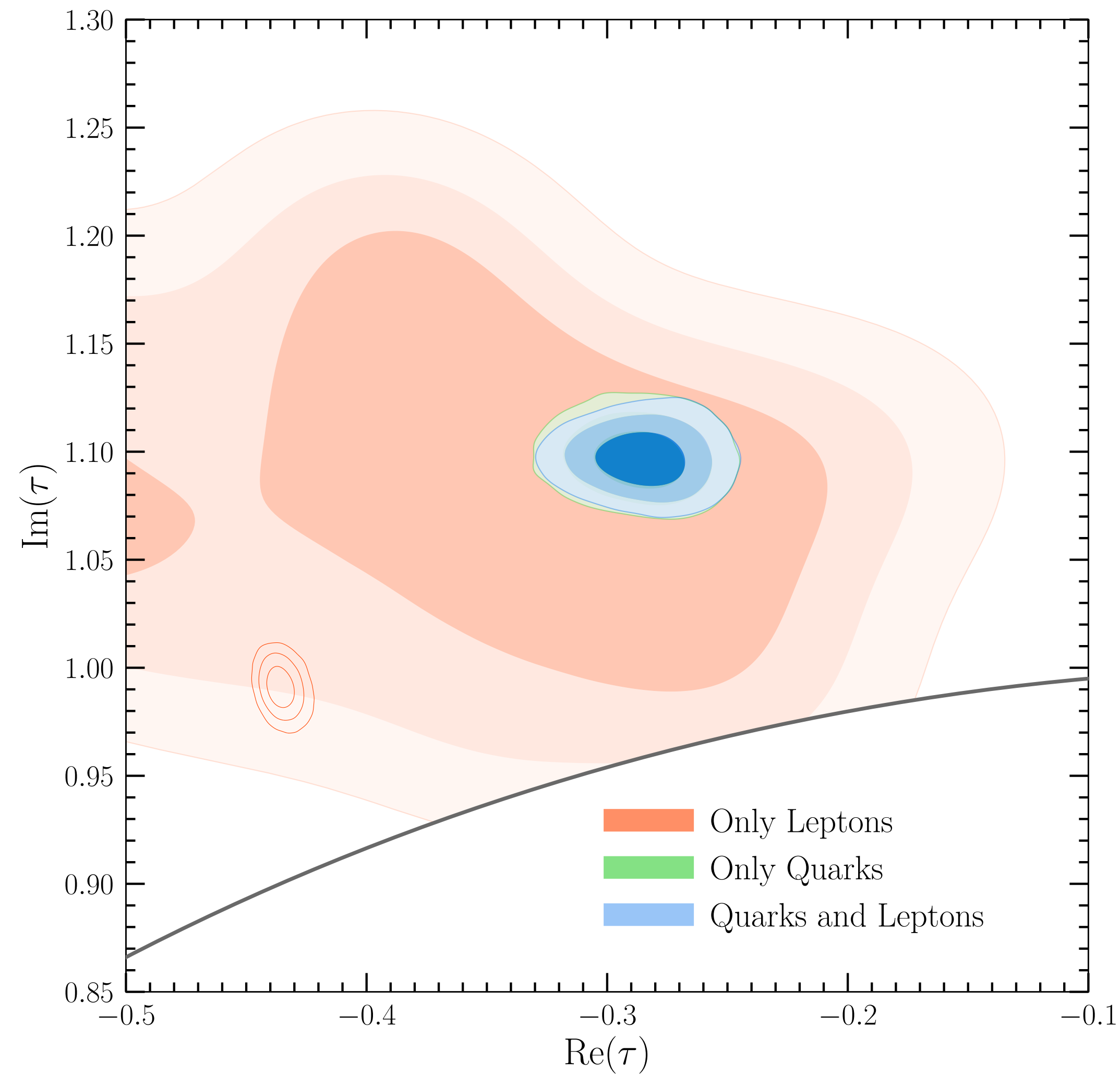


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Parameter	Value
Lepton sector	
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ℓ_{23}	1.94
e_{13}	2.21
e_{23}	7.9×10^{-3}
c_{33}^e	5.61
c_{33}^ν	7.6×10^{-2}
$c_{E3} c_{L3}$	7.66×10^{-5}
$\frac{c_{L3}^2}{2\Lambda_L}$	$6.60 \times 10^{-2}/10^{16} \text{ GeV}$
Quark sector	
q_{13}	3.7×10^{-2}
q_{23}	7.5×10^{-2}
u_{13}	3.5×10^{-2}
u_{23}	19.98
d_{13}	3.44
d_{23}	2.03×10^{-1}
$c_{U3} c_{Q3}$	4.15×10^{-4}
$c_{D3} c_{Q3}$	4.76×10^{-4}
Modulus	
τ	$-0.286 + 1.096 i$

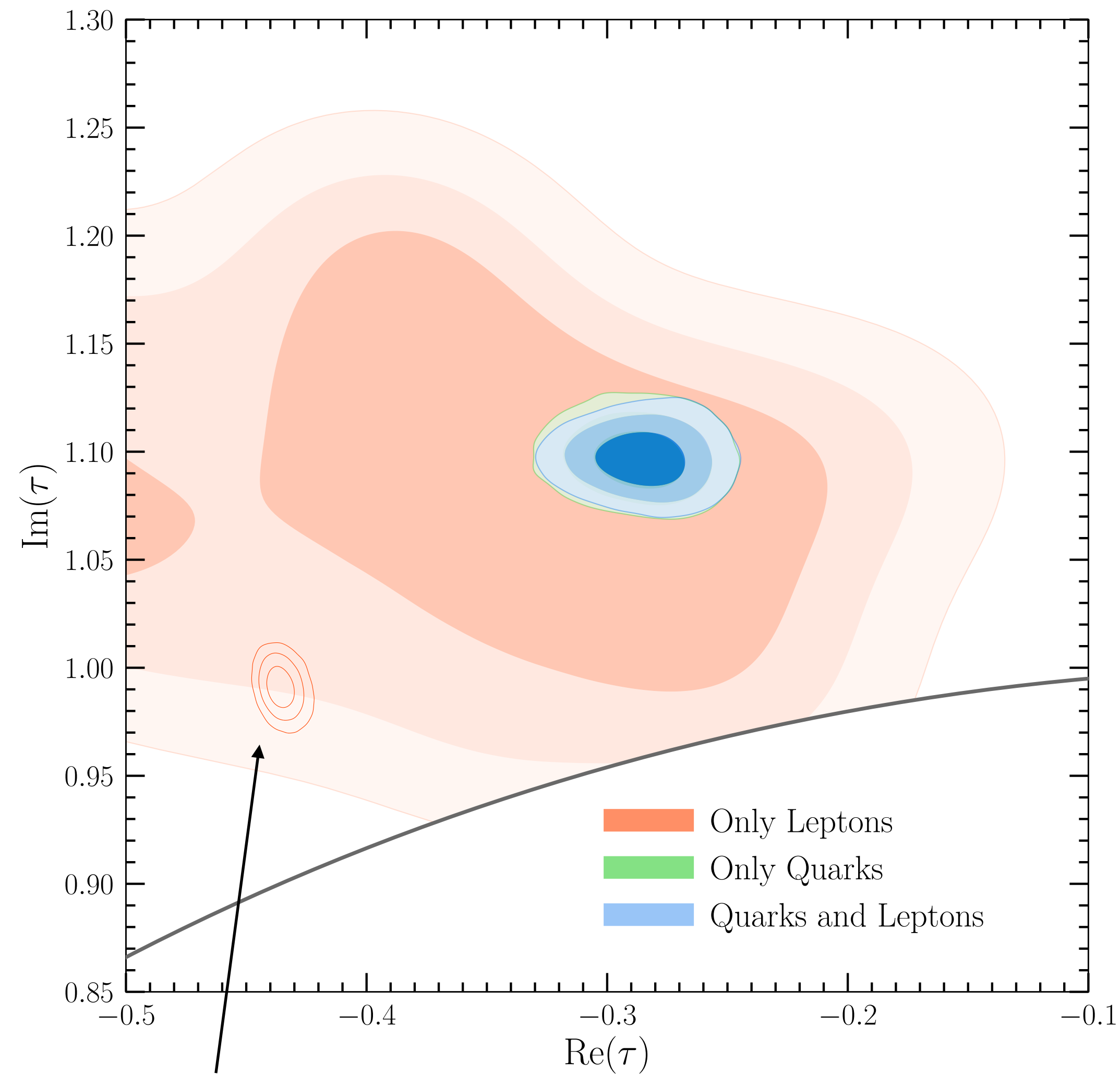
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See Feruglio's talk for the explanation of the solid orange lines

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Among the most correlated observables couples

$(\sin^2 \theta_{23}^\ell, \sin^2 \theta_{13}^\ell)$

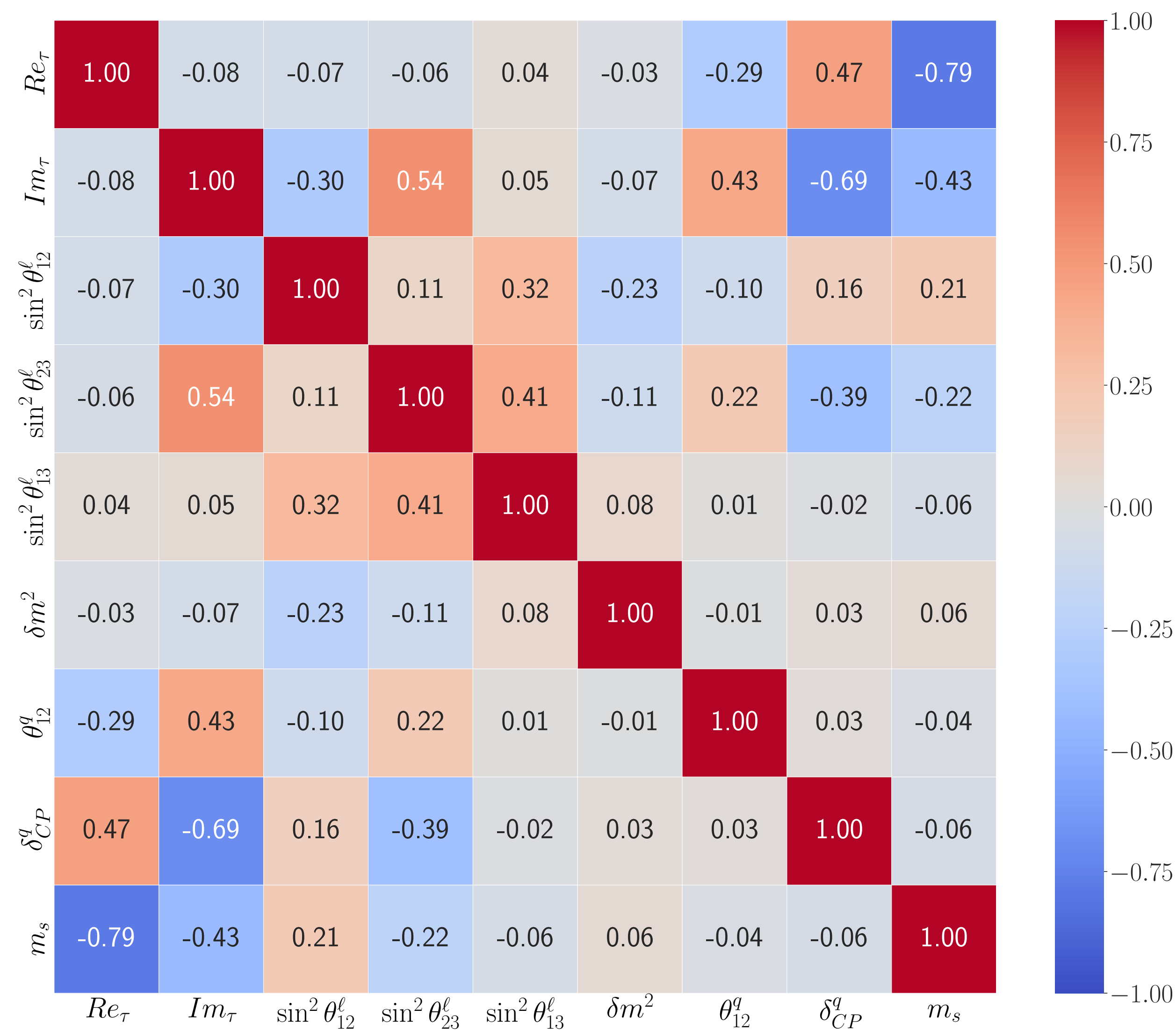
$(\sin^2 \theta_{23}^\ell, \delta_{CP}^q)$

$(\sin^2 \theta_{23}^\ell, m_s)$

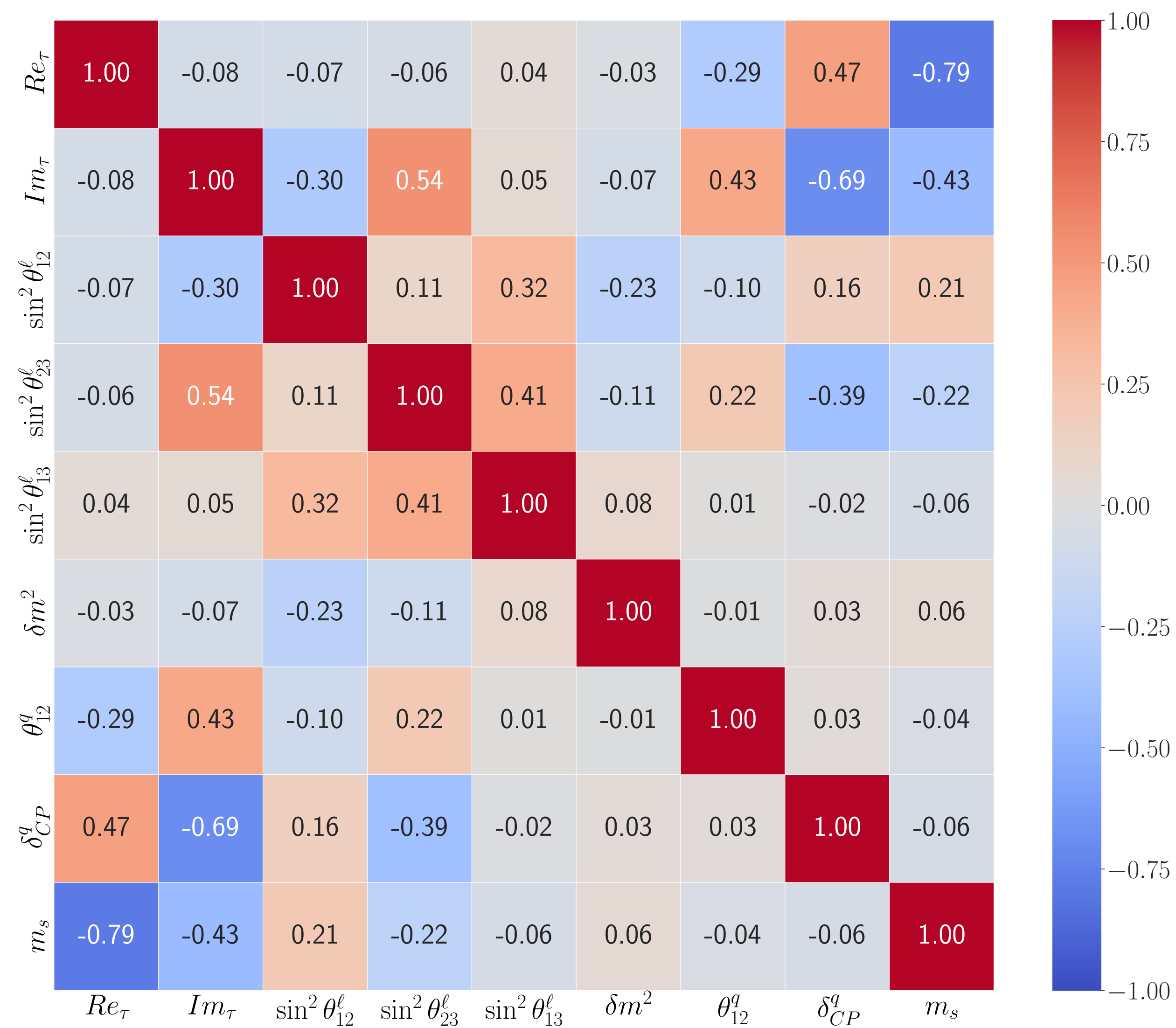
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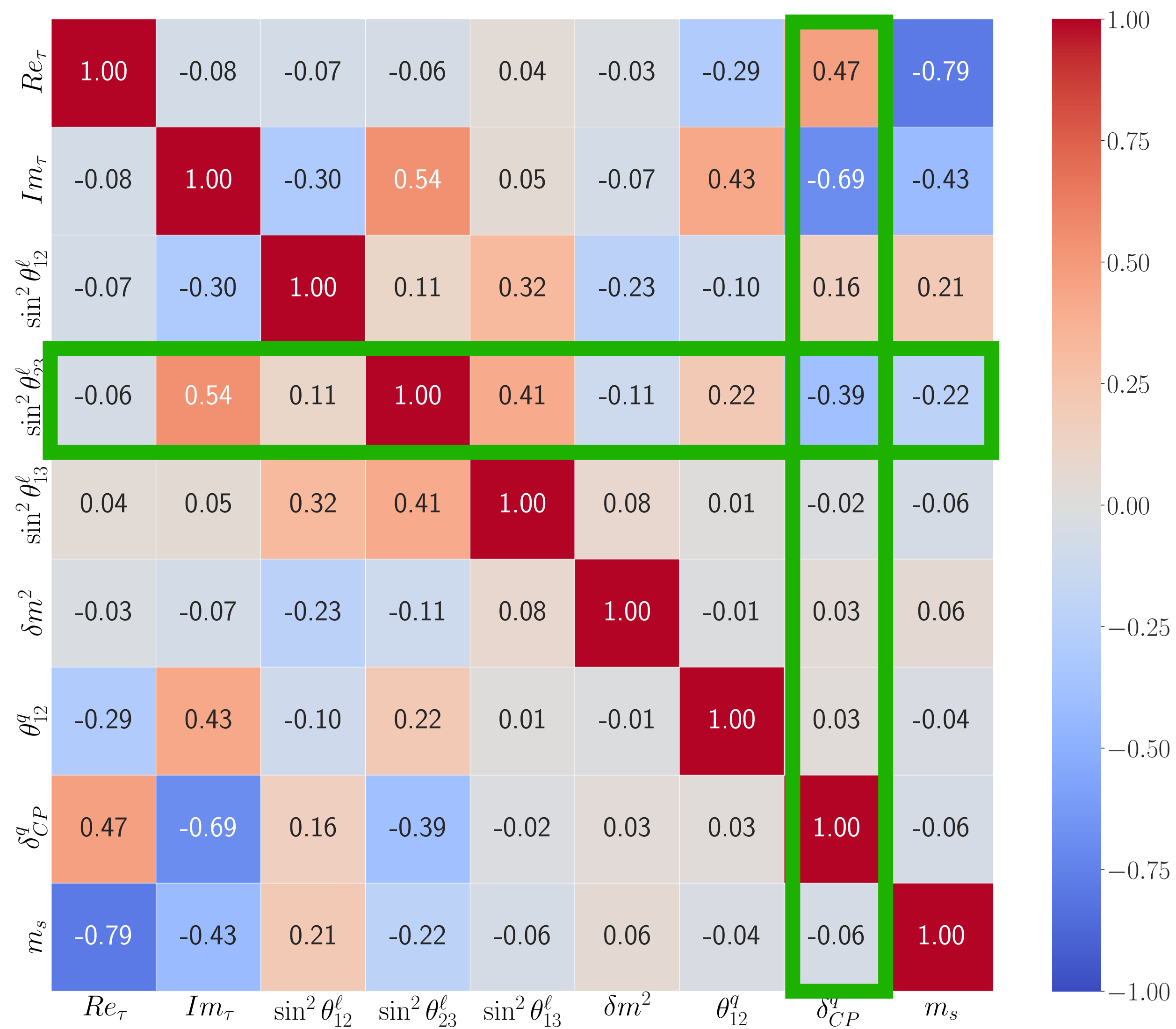


Most correlated Observables



Correlations must be governed by the modulus for lepton-quark correlated pairs of observables

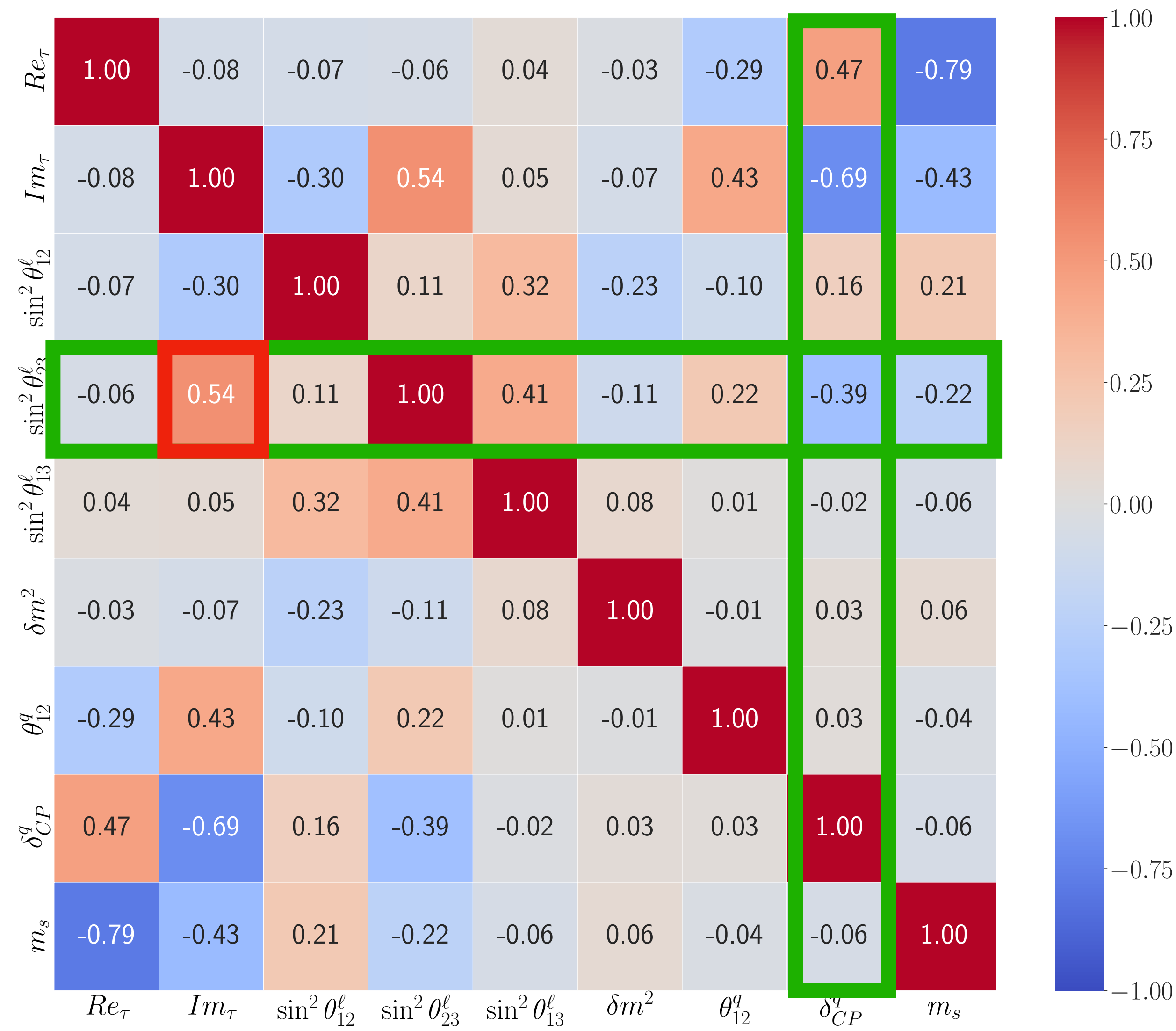
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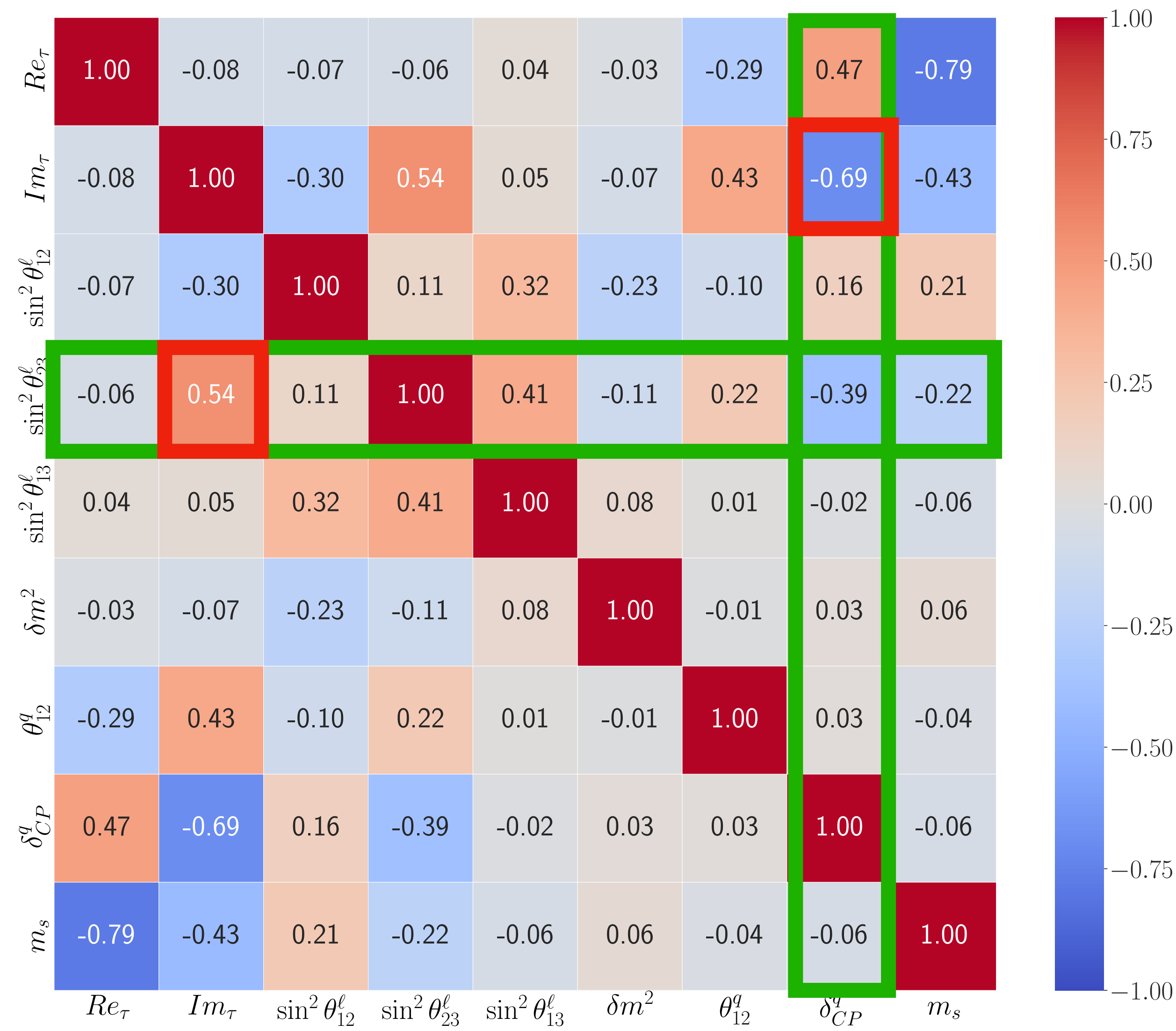


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$\sin^2 \theta_{23}^\ell$ positively correlated to $Im(\tau)$

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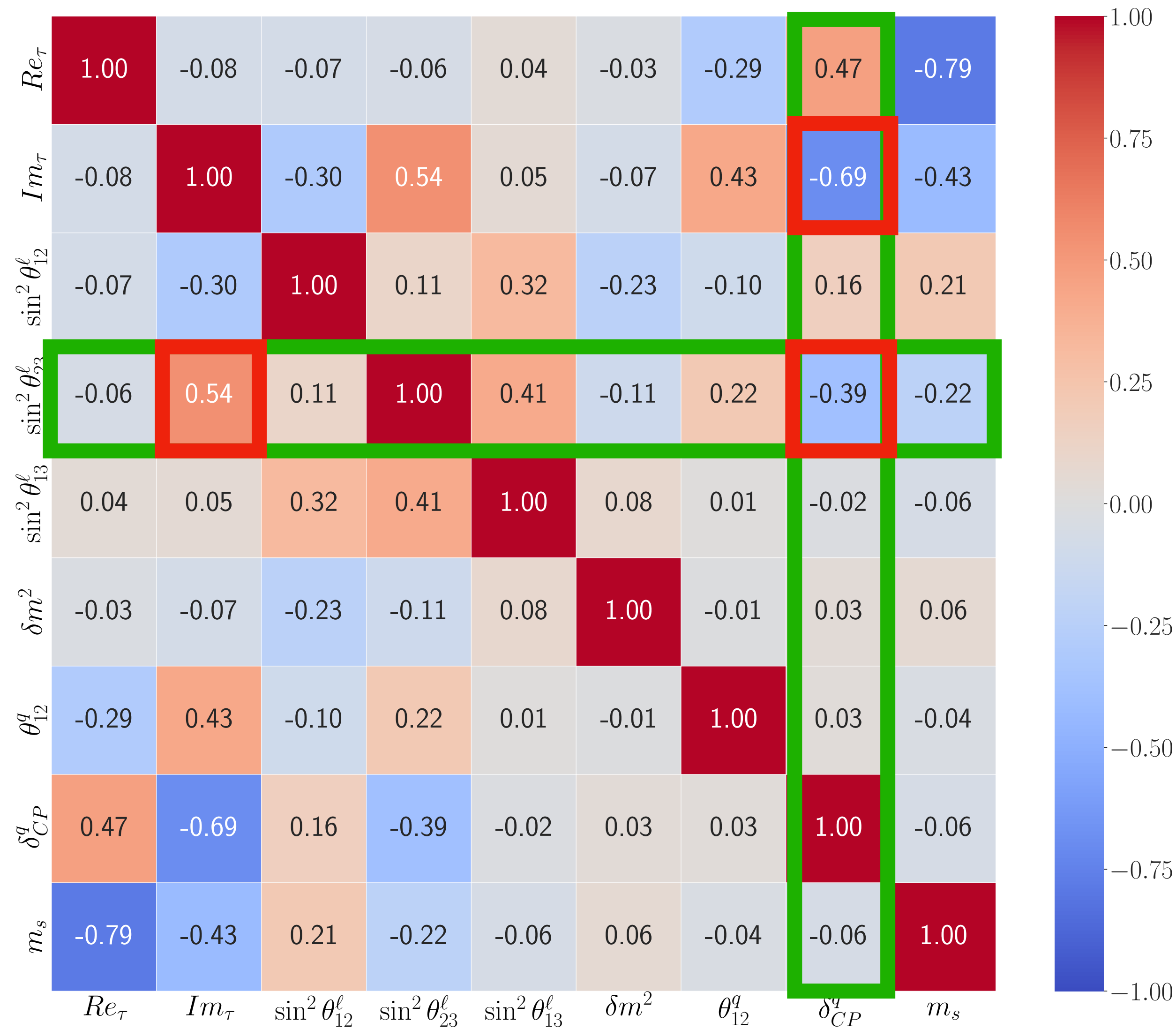
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δ_{CP}^q negatively correlated to $Im(\tau)$

Most correlated Observables



Correlations must be governed by the modulus for lepton-quark correlated pairs of observables

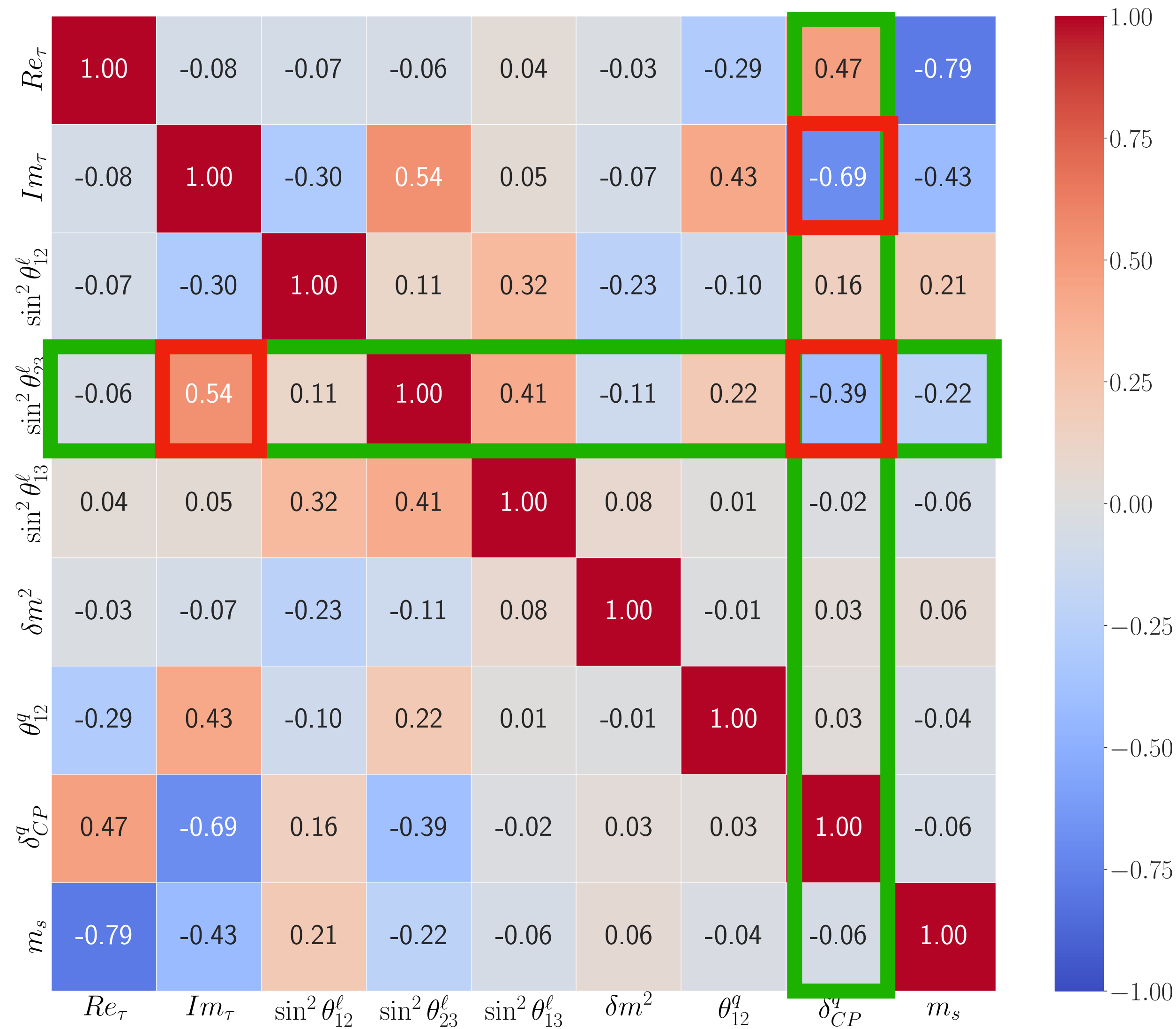
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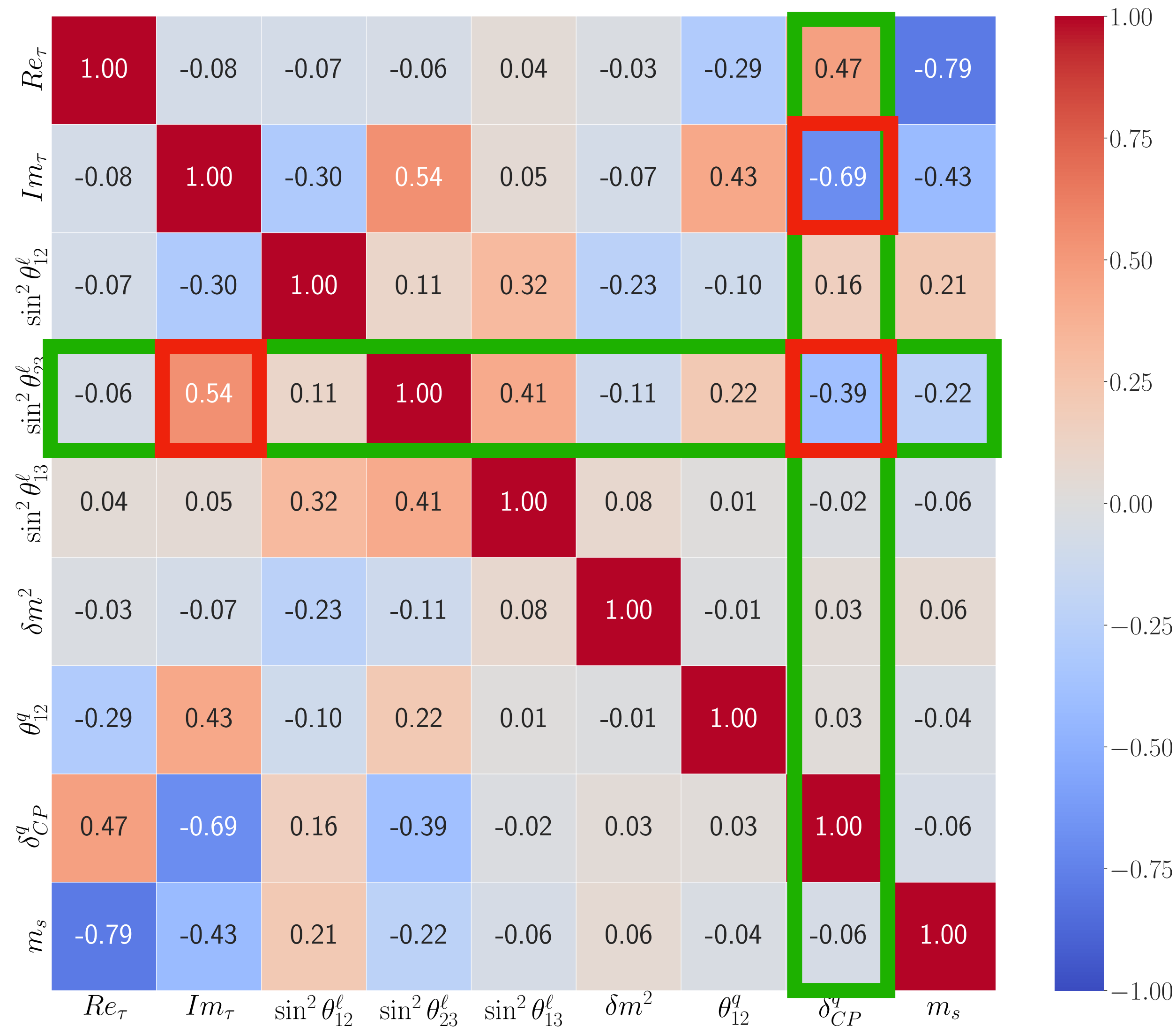
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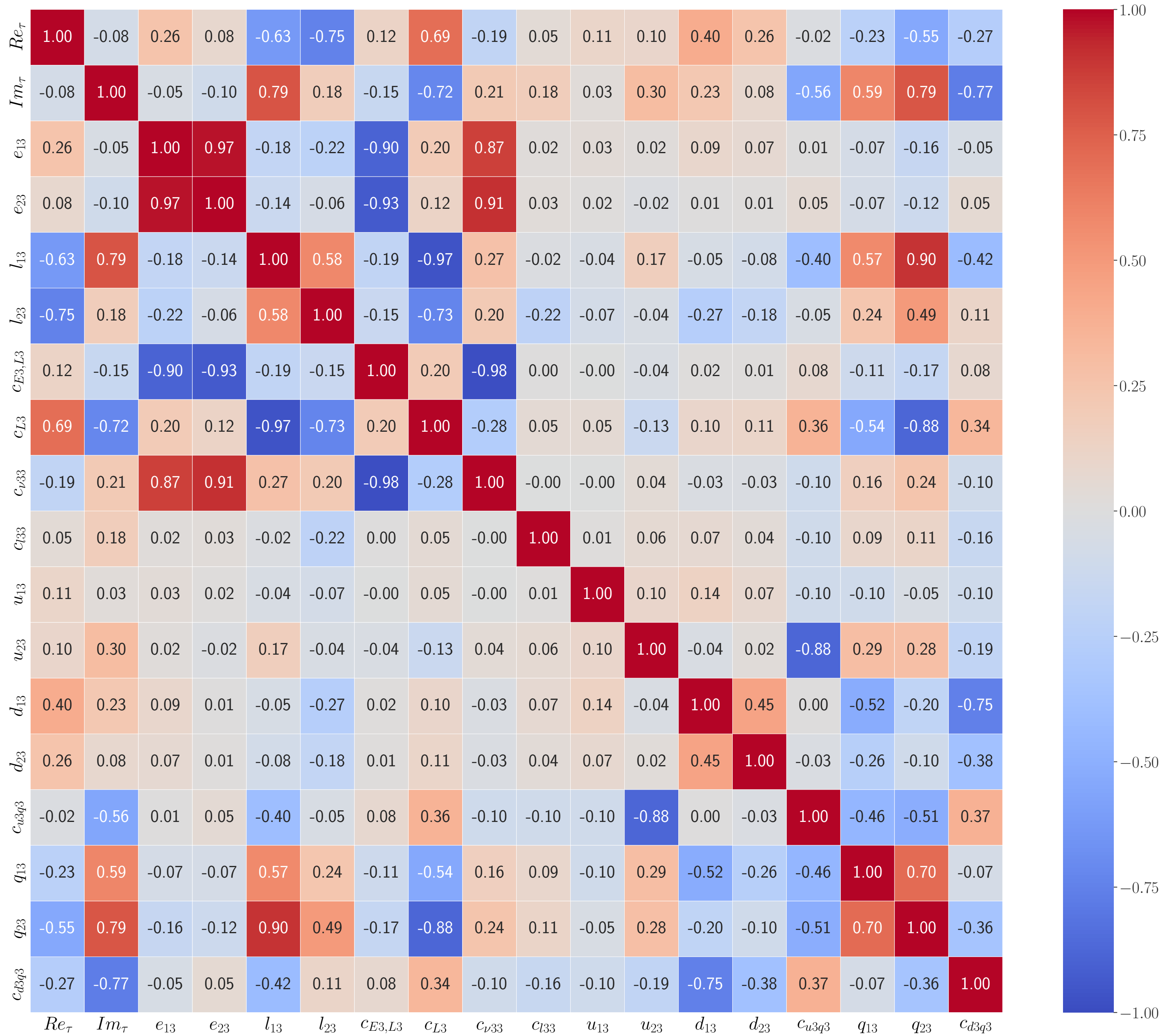
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Things are different when observables are in the same sector, for instance $(\sin^2 \theta_{23}^\ell, \sin^2 \theta_{12}^\ell)$ that are positive correlated but with opposite correlations to $Im(\tau)$. In this case one should look at the correlations with the other model parameters

Correlations between Parameters



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Also, parameters that are weekly correlated with the modulus can be strongly correlated between them, as for instance (e_{13}, e_{23}) that appear in the elements (11) and (21) of the charged lepton mass matrix



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Once again, since we are in the case of an overfitted model, the presence of strong correlations between parameters is related to the non-linear modular constraints and the internal geometry of the parameter space, and is not a sign of fine-tuning

Model based on the $2O$ group

The group $2O$, the binary octahedral group, is a group of order 48 closely related to the octahedral group O , the group of rotational symmetries of a cube (or an octahedron). Specifically, $2O$ is the double cover of the octahedral group and is isomorphic to S_4

Ding, Liu, Yao, *JHEP* 01 (2023); Ding, Liu, Lu, Weng, *JHEP* 11 (2023) 083, Ding, Lisi, AM, Petcov, *Phys.Rev.D* 111 (2025) 7, 075024

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2O has 8 irreducible representations

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Particle content of the model

Three $SU(2)$ lepton doublets $L_i = \begin{pmatrix} \nu_i \\ \ell_i \end{pmatrix}$ $i = e, \mu, \tau$

Three right-handed fermions $E^c = (e^c, \mu^c, \tau^c)$ that are $SU(2)$ singlets

Three right-handed neutrinos N_i^c $i = 1, 2, 3$ that are $SU(2)$ singlets

Six right-handed quarks $U^c = (u^c, c^c, t^c)$ and $D^c = (d^c, s^c, b^c)$ that are $SU(2)$ singlets

Three $SU(2)$ quark doublets $Q_i = \begin{pmatrix} u_i \\ d_i \end{pmatrix}$ $u_i = u, c, t$ $d_i = d, s, b$

Two Higgs fields, H_u and H_d that are invariant singlets of 2O with zero modular weights.

Representation under the 2O modular group and weights

$$L \sim \mathbf{3}, \quad E_D^c \equiv (e^c, \mu^c) \sim \hat{\mathbf{2}}', \quad \tau^c \sim \mathbf{1}', \quad N^c \sim \mathbf{3}$$

$$k_L = -1, \quad k_{E_D^c} = 6, \quad k_{\tau^c} = 5, \quad k_{N^c} = 1$$

$$Q_D \equiv (Q_1, Q_2)^T \sim \mathbf{2}, \quad Q_3 \sim \mathbf{1}', \quad U_D^c \equiv (u^c, c^c) \sim \hat{\mathbf{2}}'$$

$$t^c \sim \mathbf{1}', \quad D_D^c \equiv (d^c, s^c) \sim \mathbf{2}, \quad b^c \sim \mathbf{1}'$$

$$k_{Q_D} = 3 - k_{U_D^c} = k_{Q_3} = 6 - k_{t^c} = 6 - k_{D_D^c} = -k_{b^c}$$

The Lagrangian of the model contains the most general superpotential with all possible singlets under $2O$ built with the modular forms organized in various multiplets of different weights

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$$\mathcal{W}_\nu = g H_u (N^c L)_1 + \Lambda (N^c N^c)_2 Y_2^{(2)} \quad \text{neutrinos}$$

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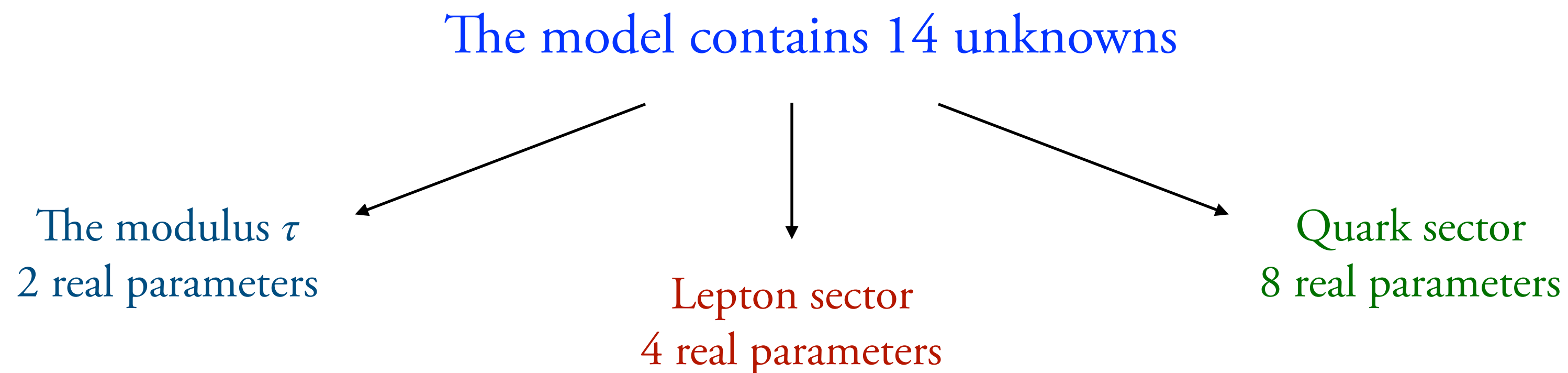
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Minimal model in the modular form literature (before arXiv:2506.19822)

Experimental inputs and model parameters

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Model should satisfy 18 experimental constraints (only Normal Ordering of neutrino masses allowed)

$$O_i = \left\{ \begin{array}{lll} \sin^2 \theta_{12}, \sin^2 \theta_{23}, \sin^2 \theta_{13} & i = 1, 2, 3 & \nu \text{ mixing,} \\ \delta m^2, \Delta m^2 & i = 4, 5 & \nu \text{ masses,} \\ r_{e\mu}, r_{\mu\tau}, m_\tau & i = 6, 7, 8 & \text{charged lepton masses,} \\ \theta_{12}^q, \theta_{23}^q, \theta_{13}^q, \delta_{CP}^q & i = 9, 10, 11, 12 & \text{quark mixing,} \\ r_{uc}, r_{ct}, r_{ds}, r_{sb}, m_t, m_b & i = 13, \dots, 18 & \text{quark masses.} \end{array} \right. \left. \begin{array}{l} \text{8 for leptons} \\ \text{10 for quarks} \end{array} \right\}$$

charged fermion masses and quark mixing
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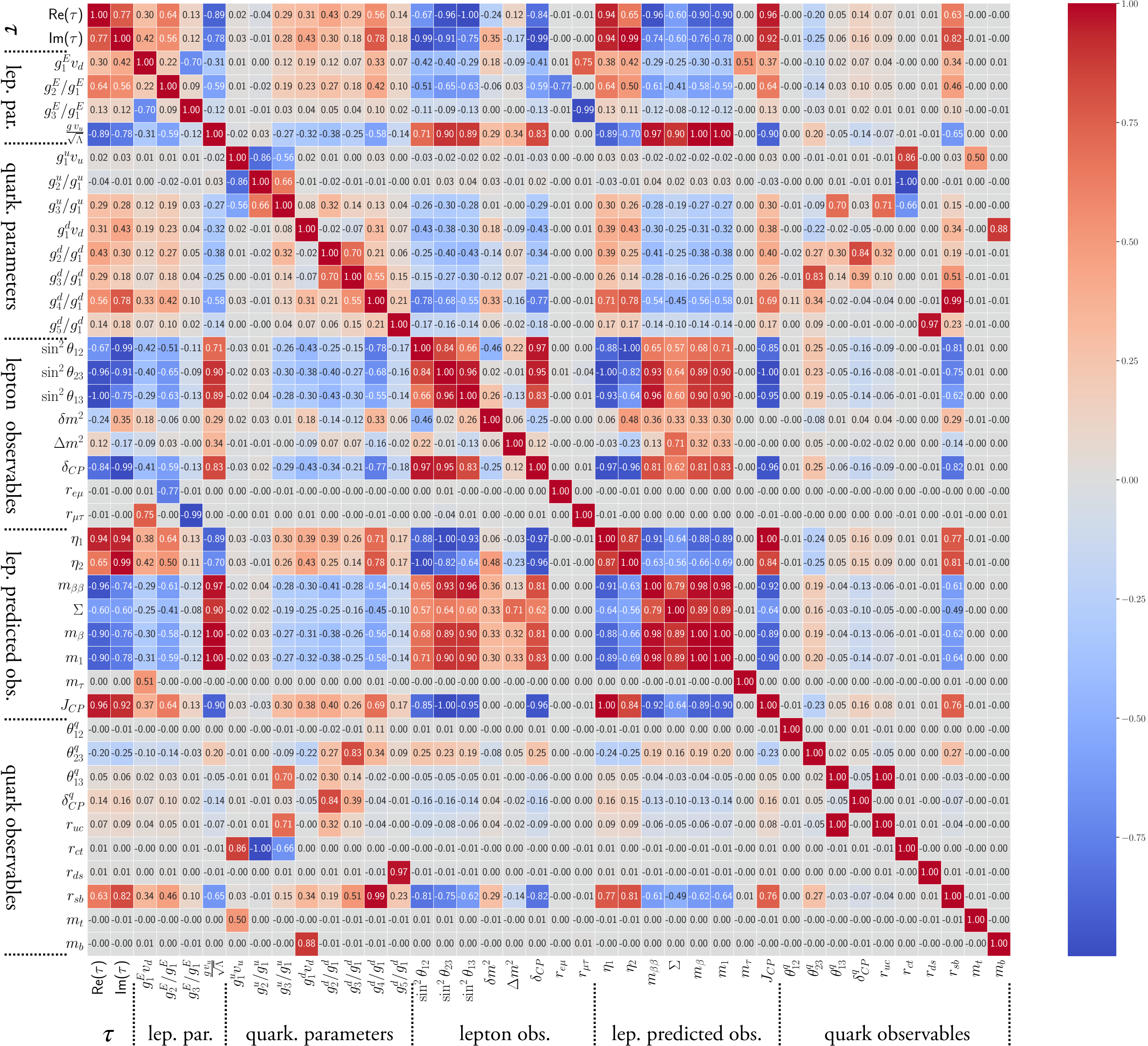
The Lagrangian of the model contains 14 real parameters, 6 for leptons, 10 for quarks, with τ in common

$$\mathbf{P}_{\text{leptons}} = (\tau, g_1^E v_d, \frac{g_2^E}{g_1^E}, \frac{g_3^E}{g_1^E}, \frac{g v_u}{\sqrt{\Lambda}}),$$

$$\mathbf{P}_{\text{quarks}} = (\tau, g_1^u v_u, \frac{g_2^u}{g_1^u}, \frac{g_3^u}{g_1^u}, g_1^d v_d, \frac{g_2^d}{g_1^d}, \frac{g_3^d}{g_1^d}, \frac{g_4^d}{g_1^d}, \frac{g_5^d}{g_1^d})$$

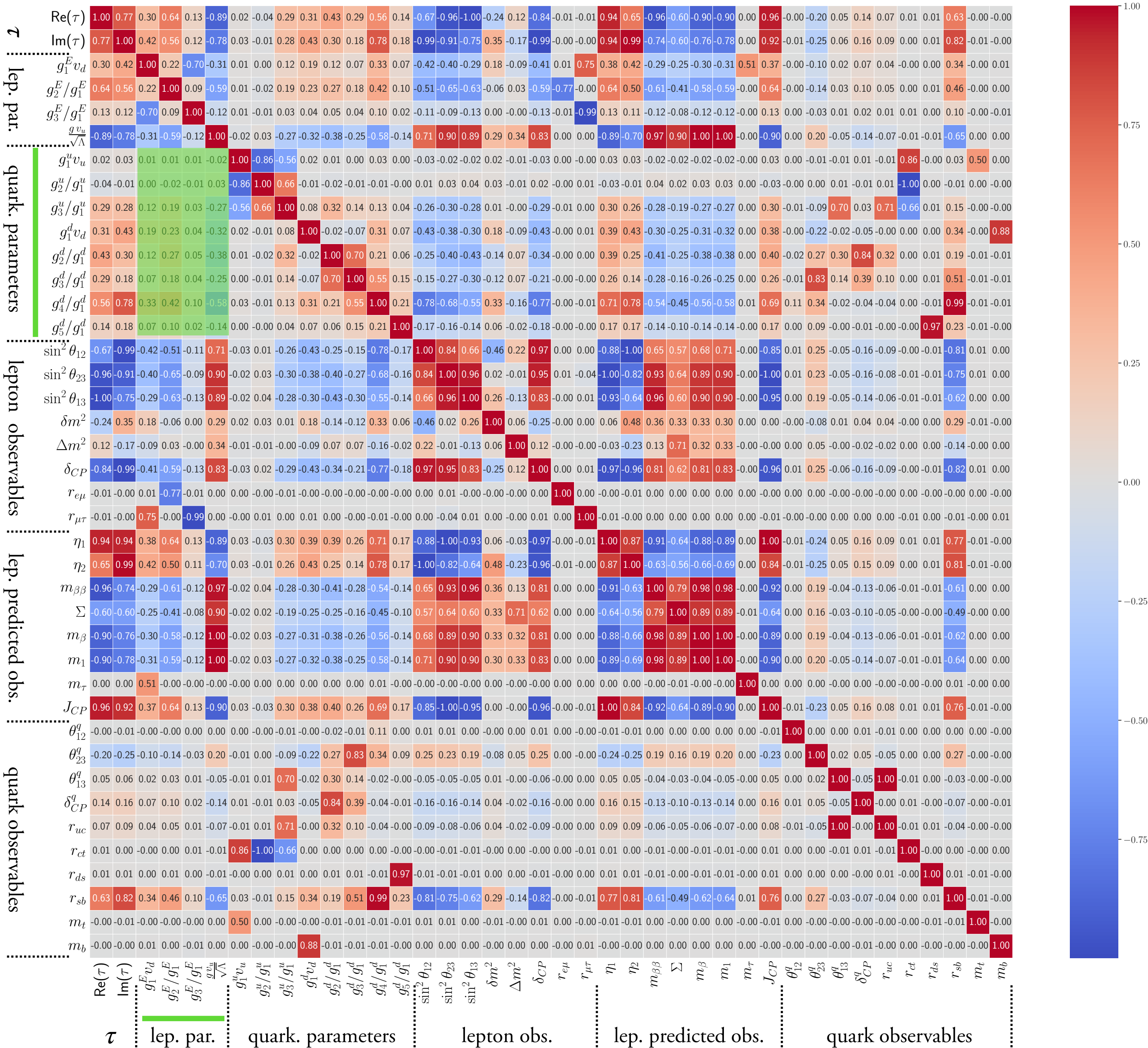
Complete correlations matrix

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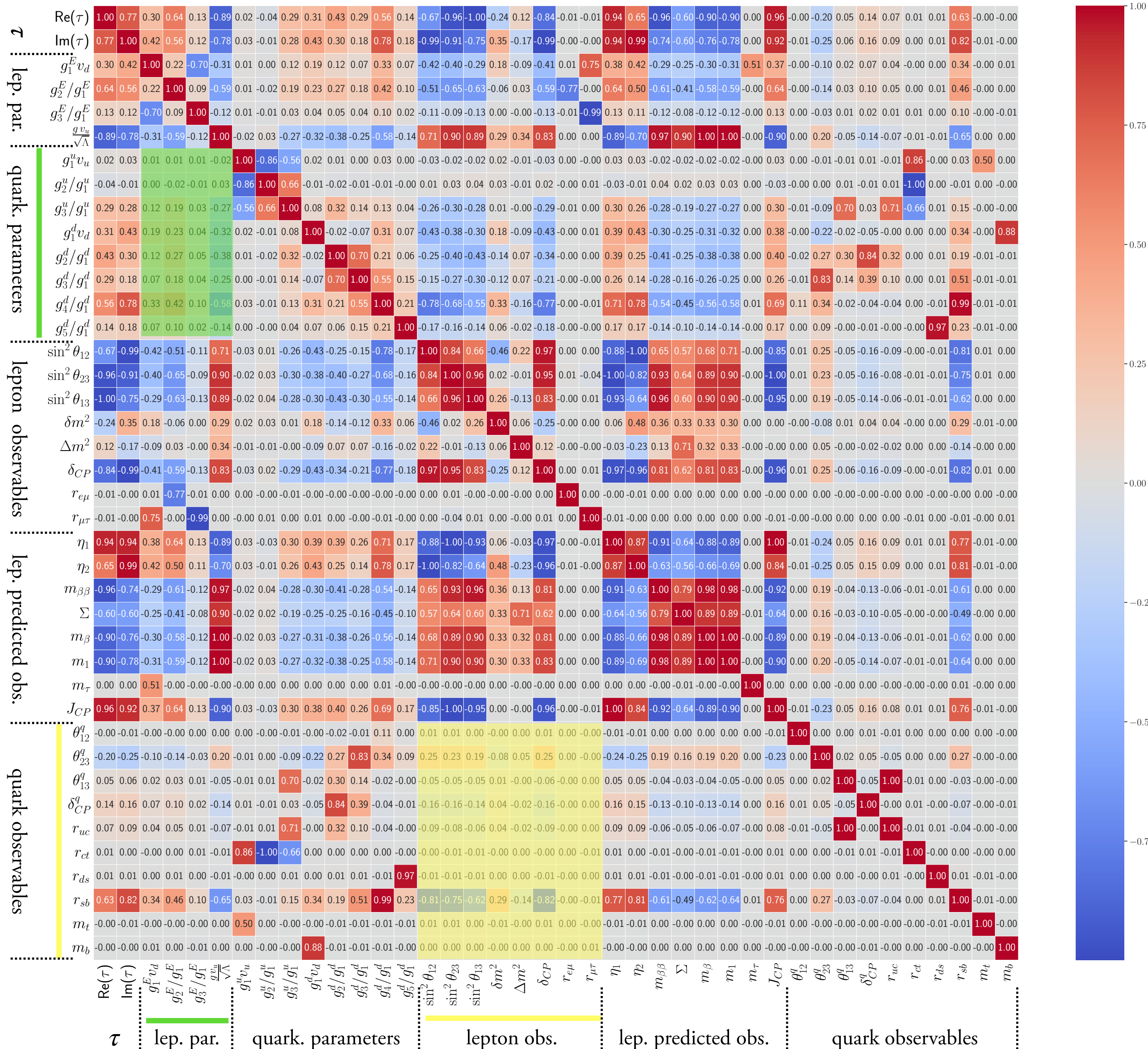
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Complete correlations matrix

Correlations between lepton and quark parameters

Correlations between lepton and quark observables

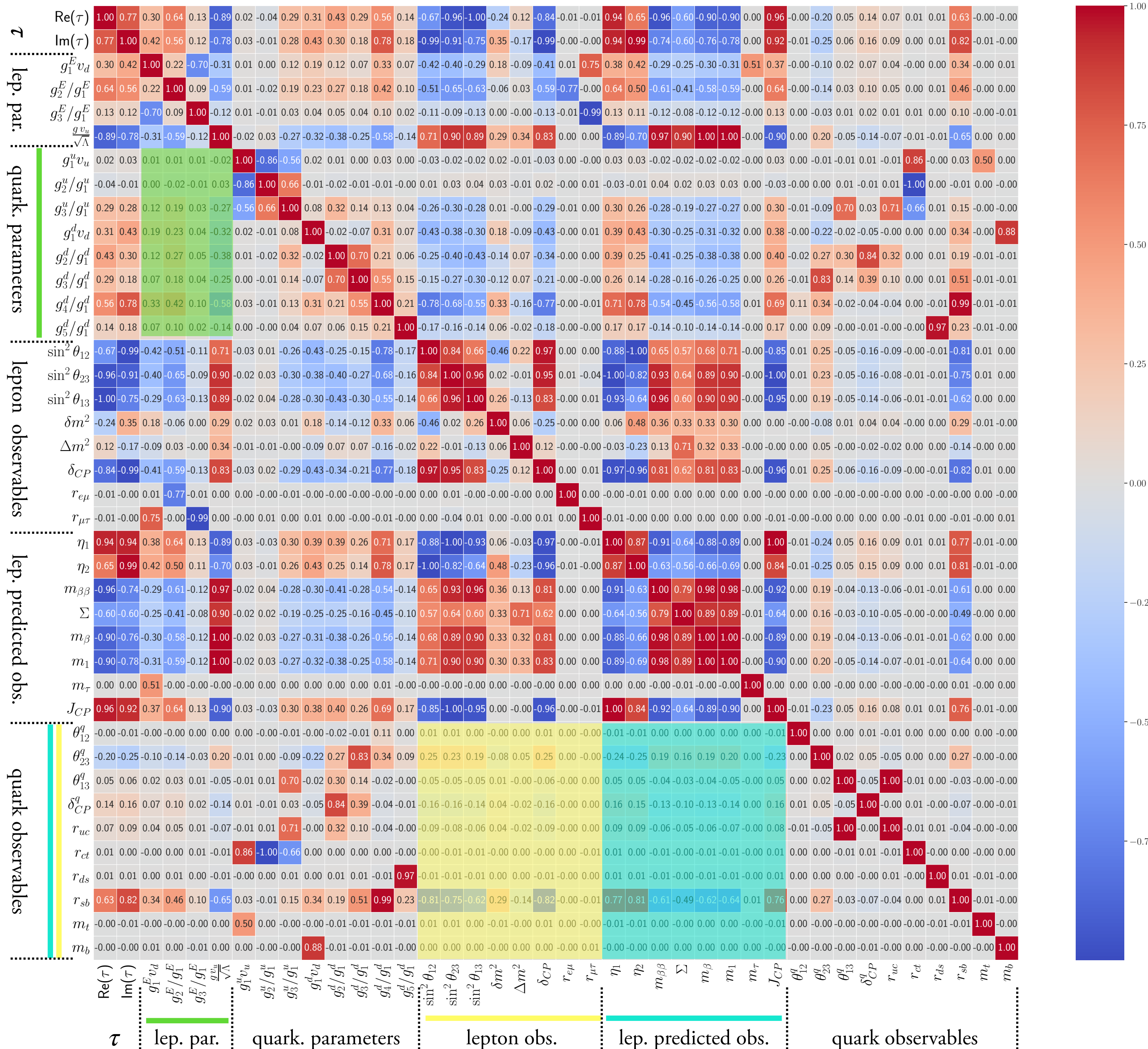


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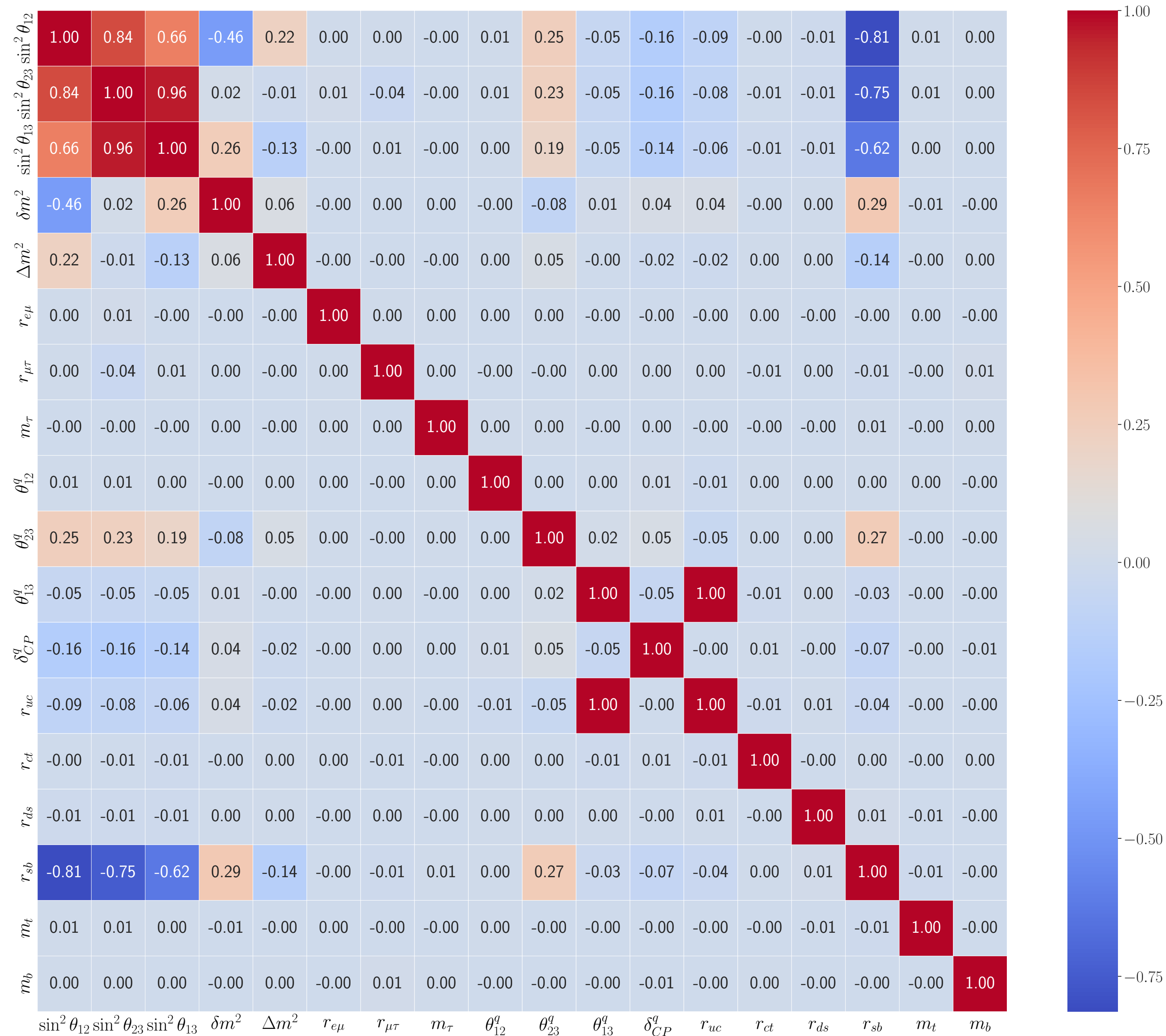
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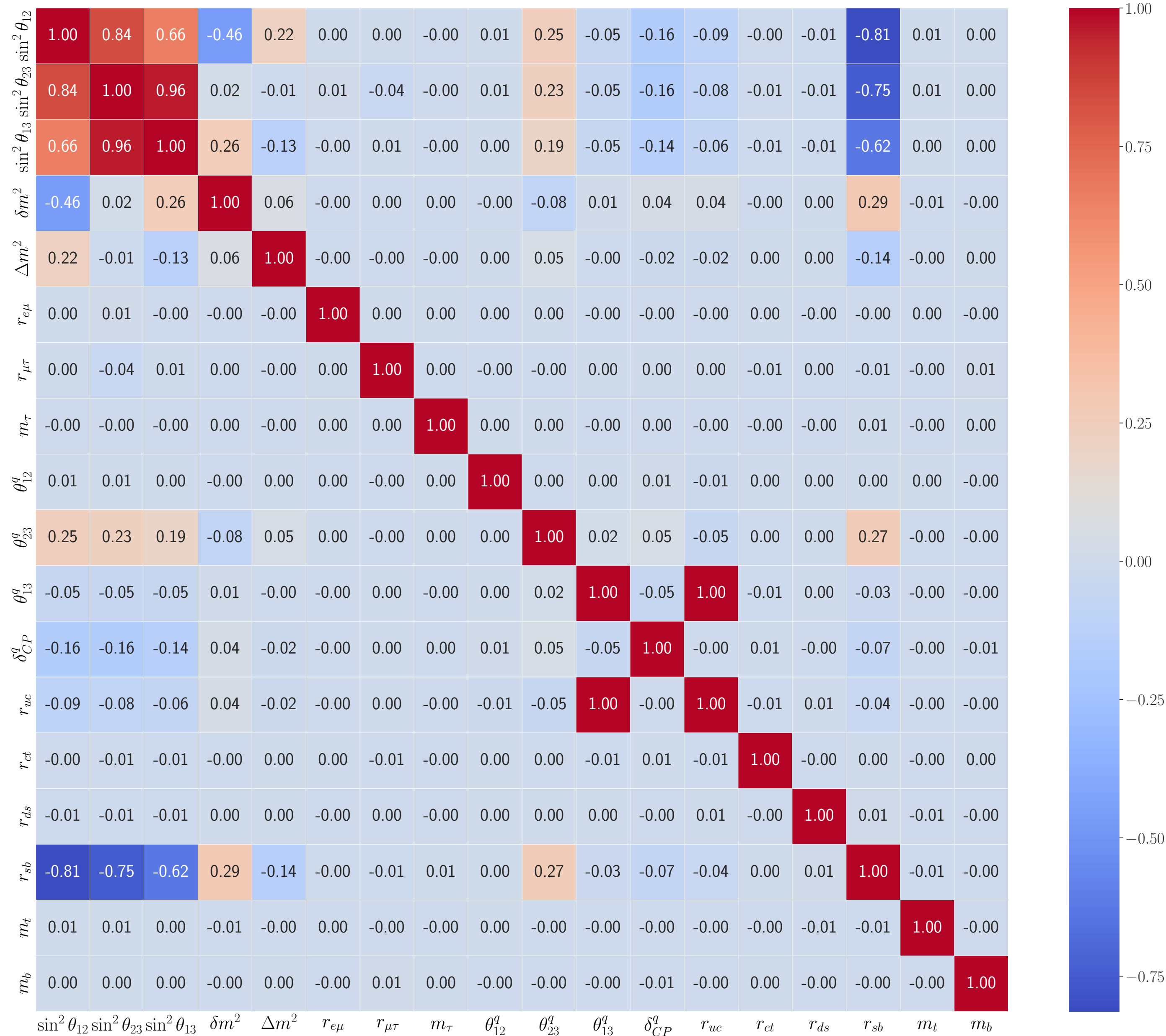
Correlations between predicted neutrino unknowns and quark observables



Correlations among Observables

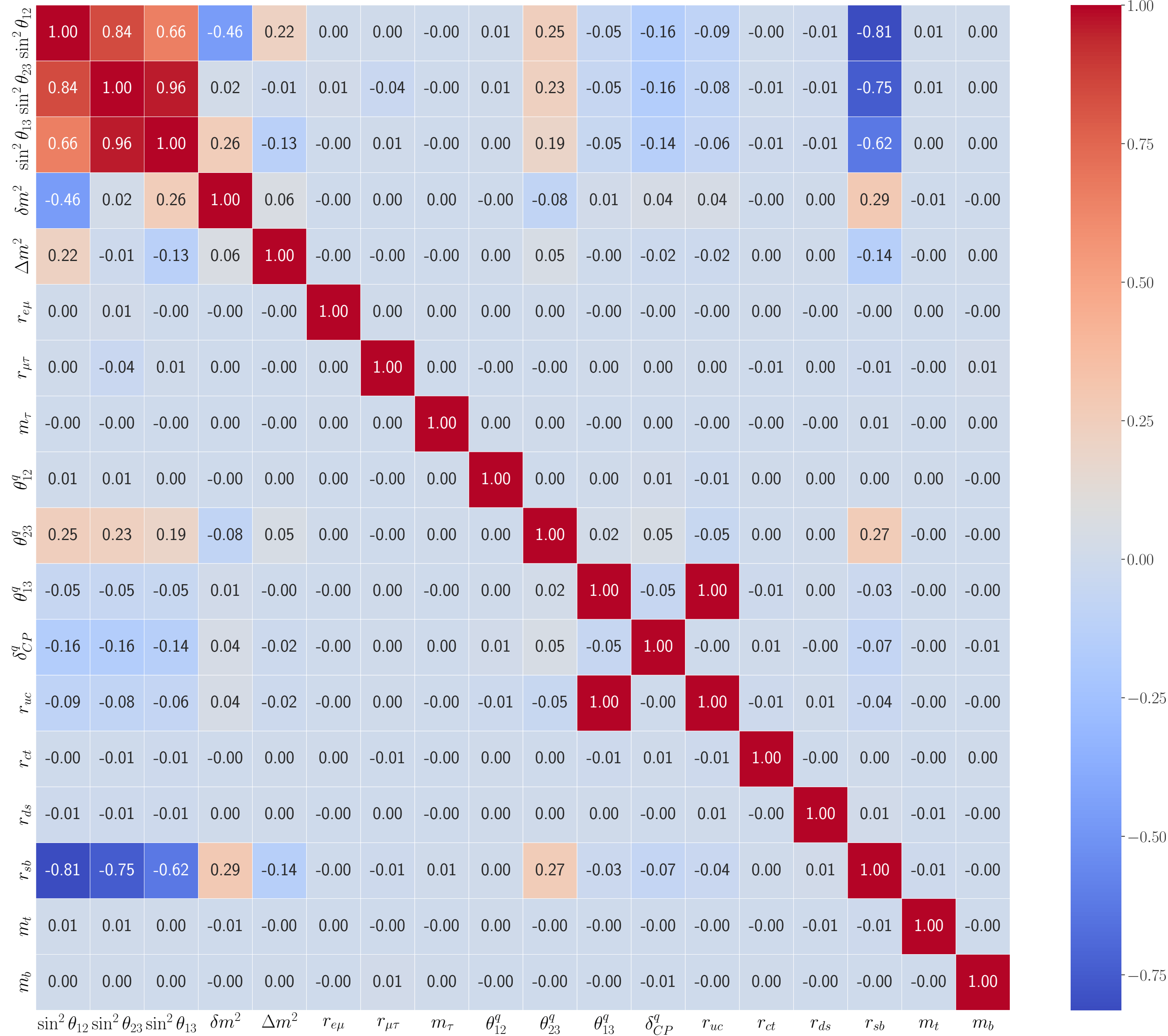


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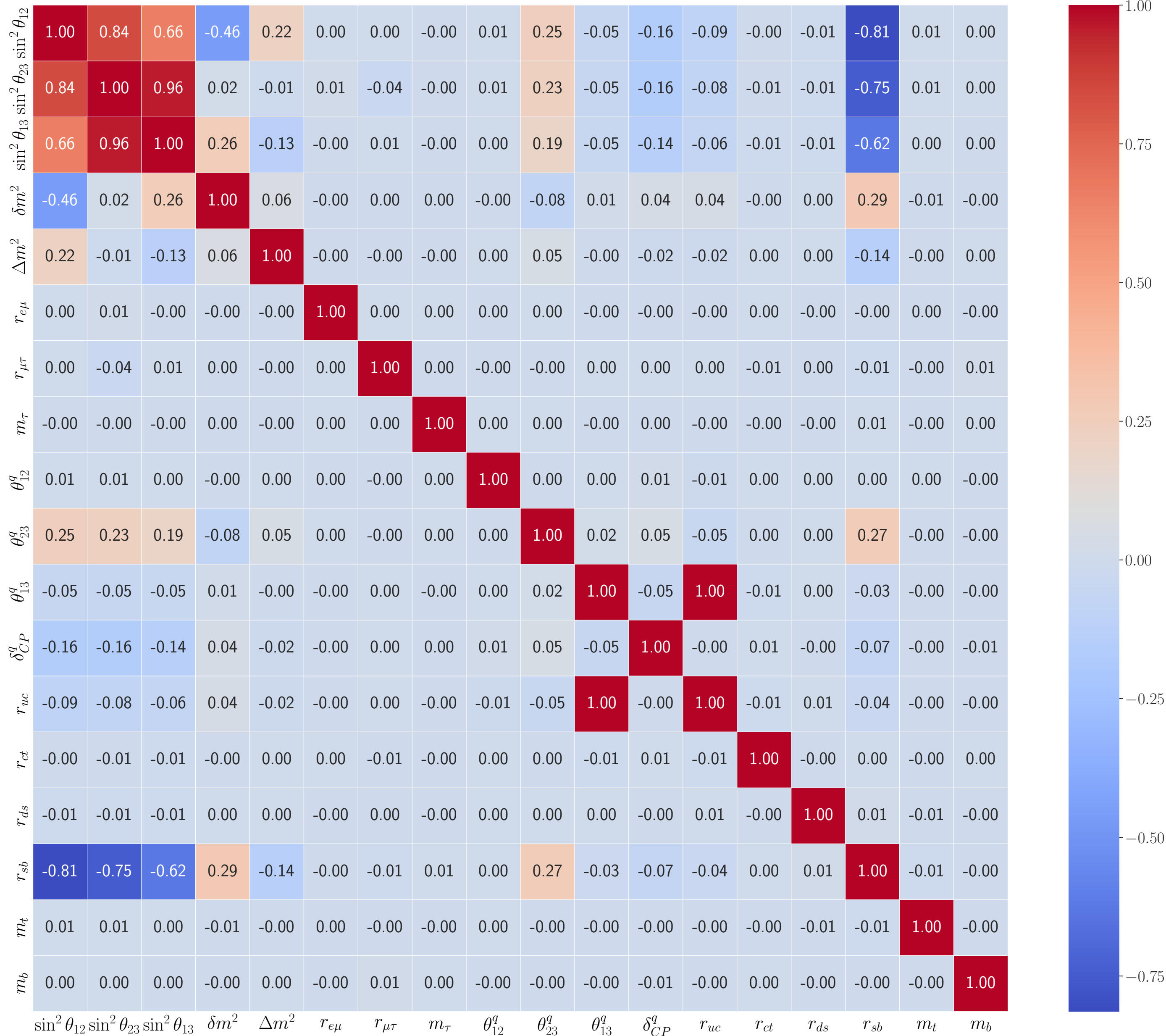
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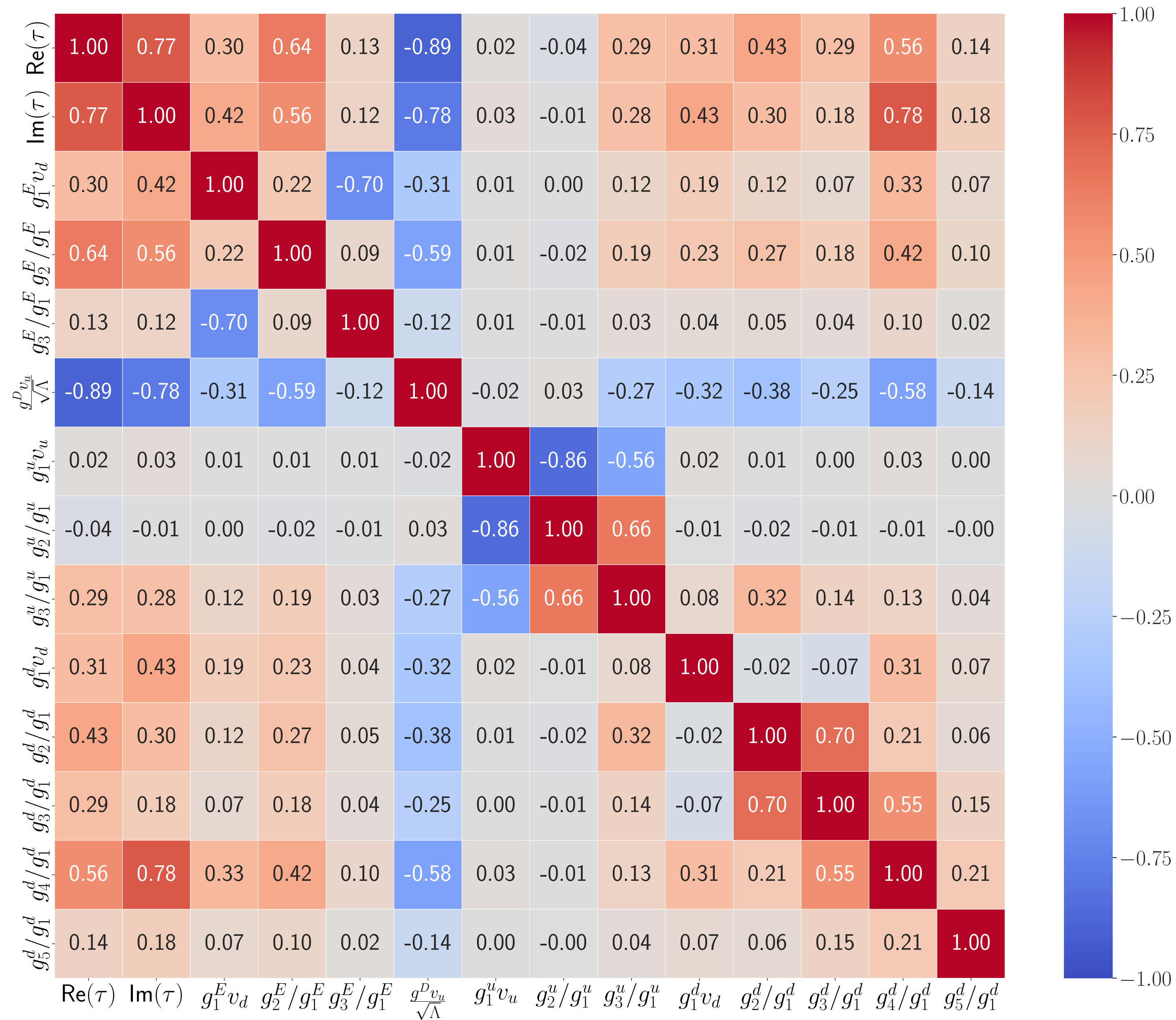
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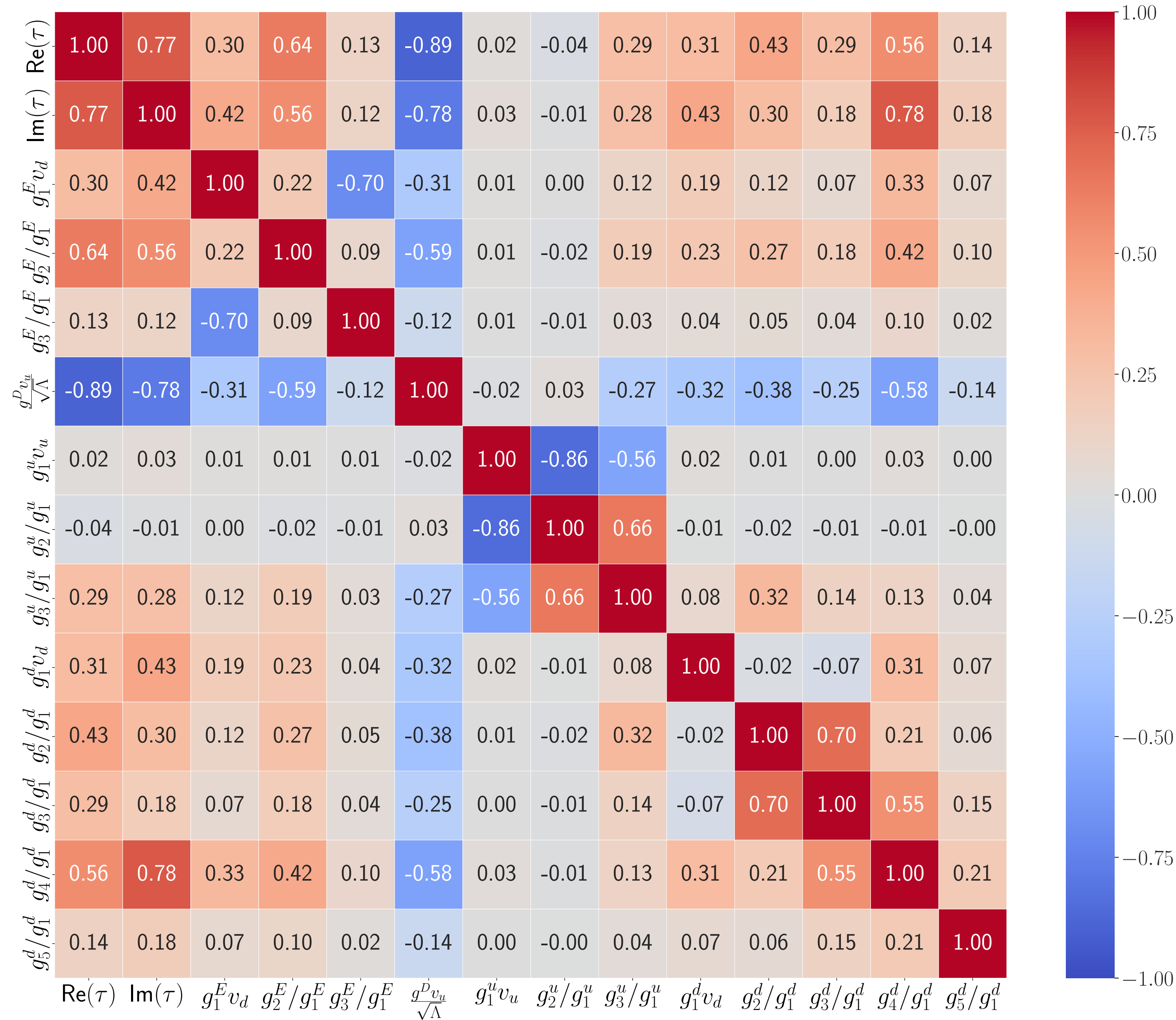
In most of the cases, correlations seem to be governed by the modulus τ

Interesting case of the (θ_{13}^q, r_{uc}) pair, for which the correlation is ~ 1 and is not related to τ

Correlations between Parameters

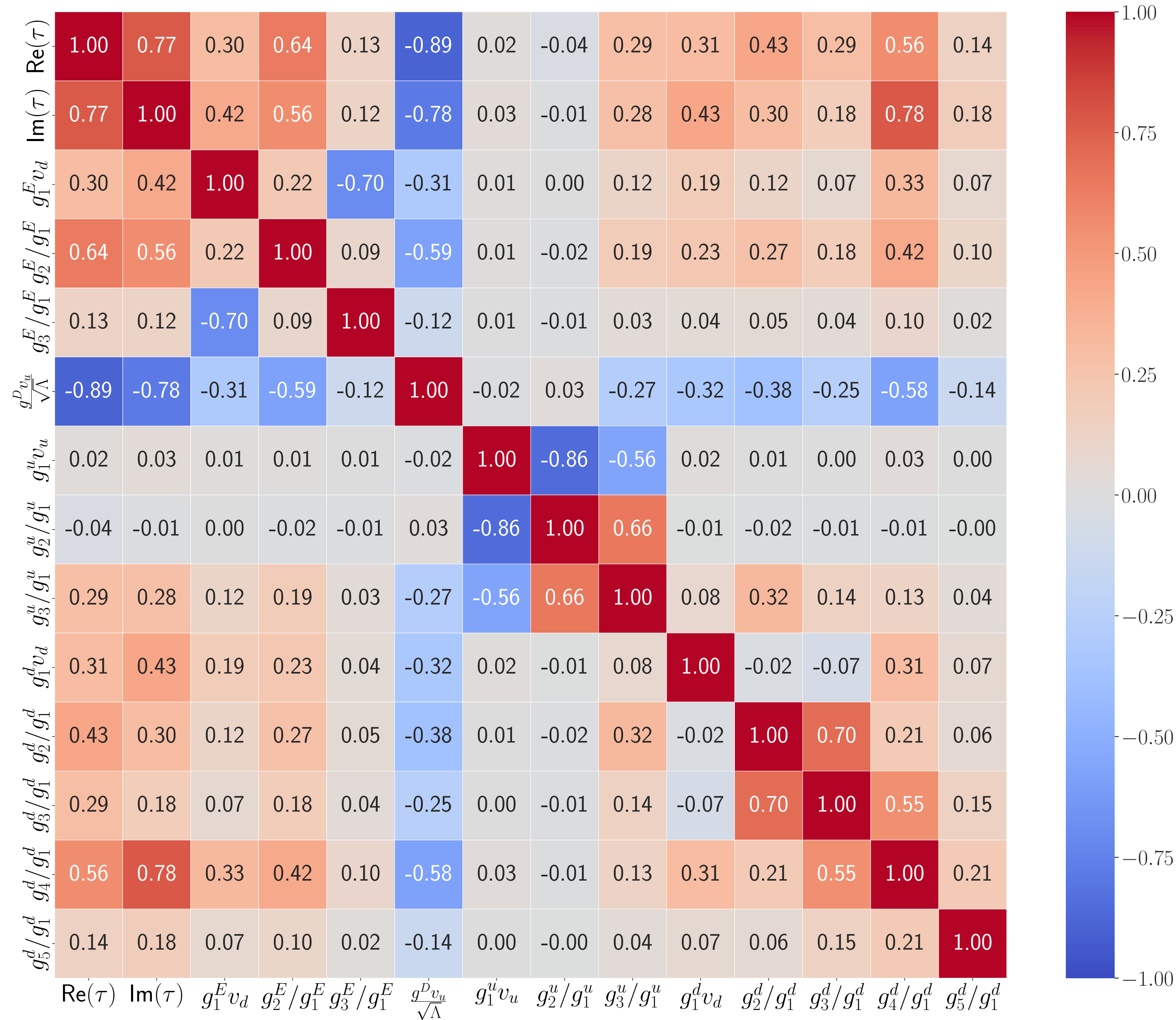


Correlations between Parameters



Correlations between Parameters

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Correlations between Parameters

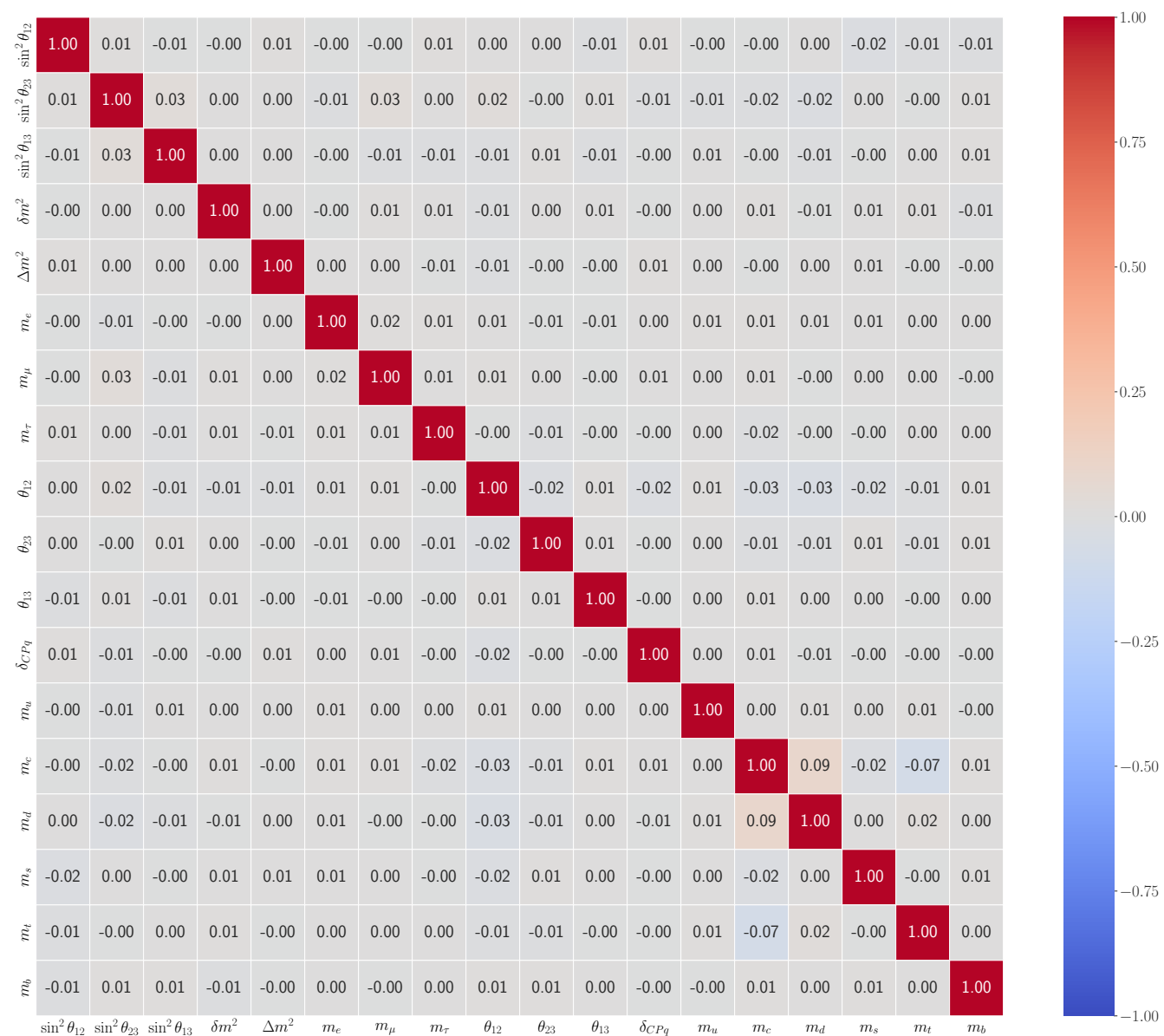
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Apparently, there is not a significant difference with respect to the case of overfitting models

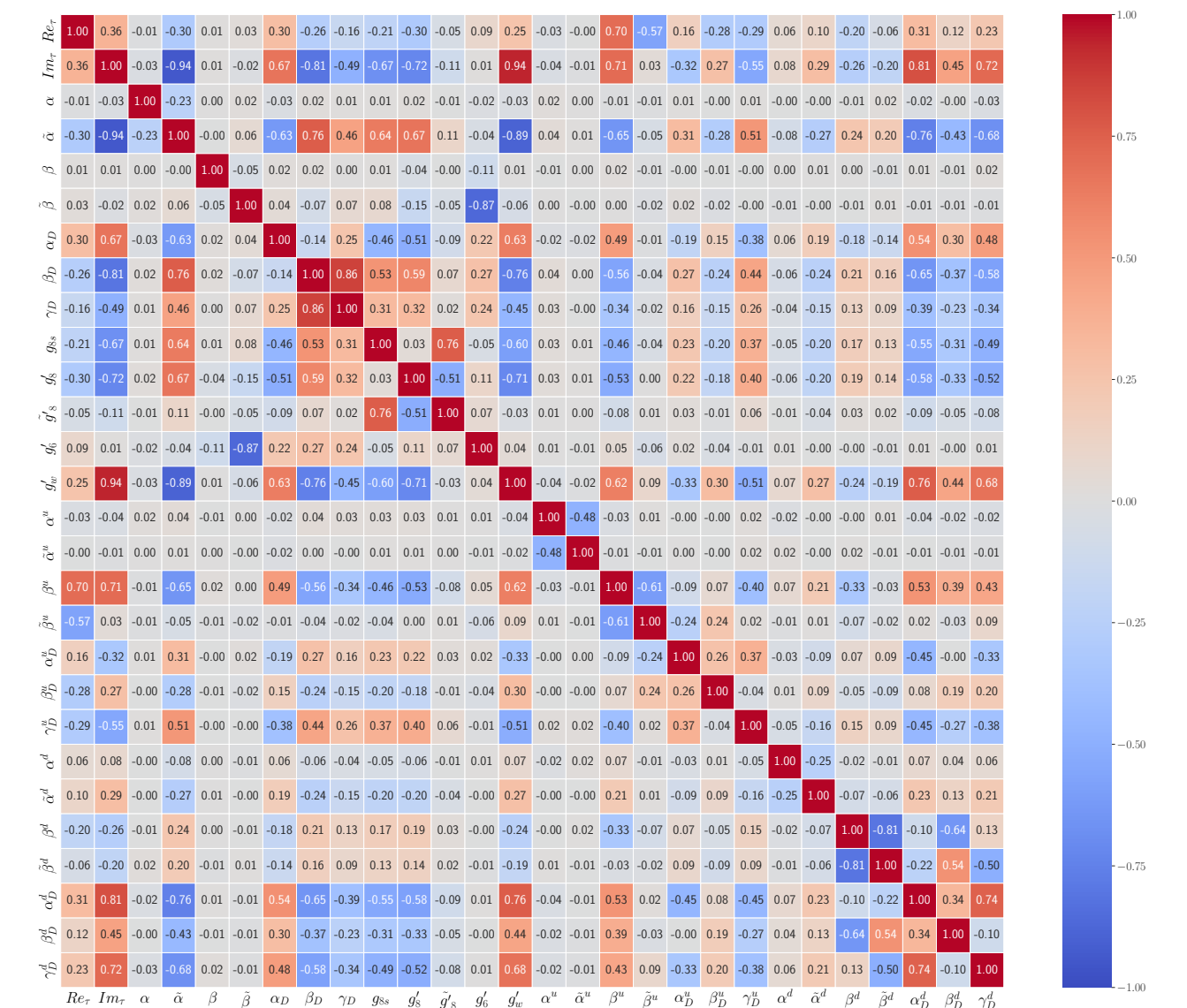
Observables

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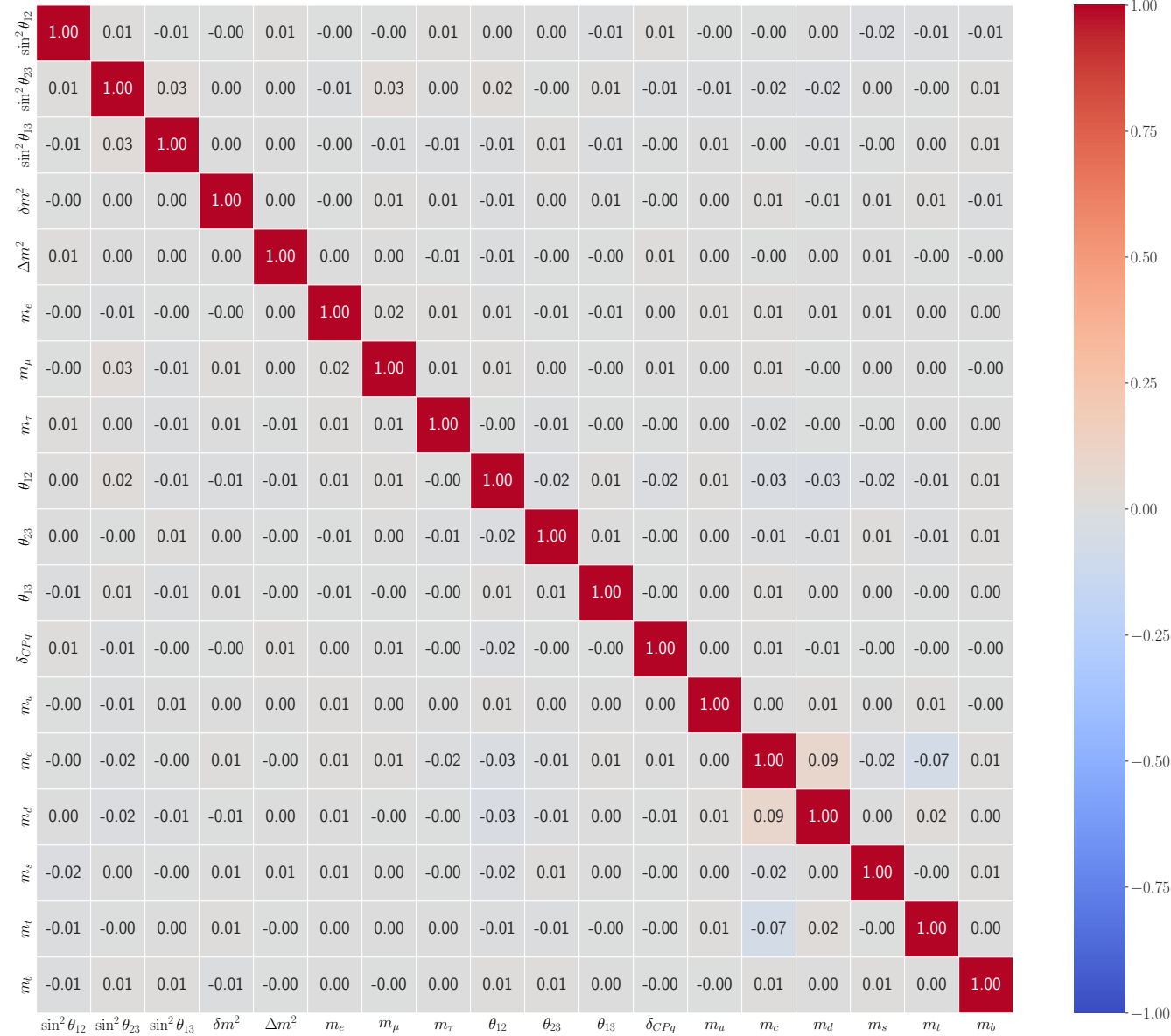
Observables



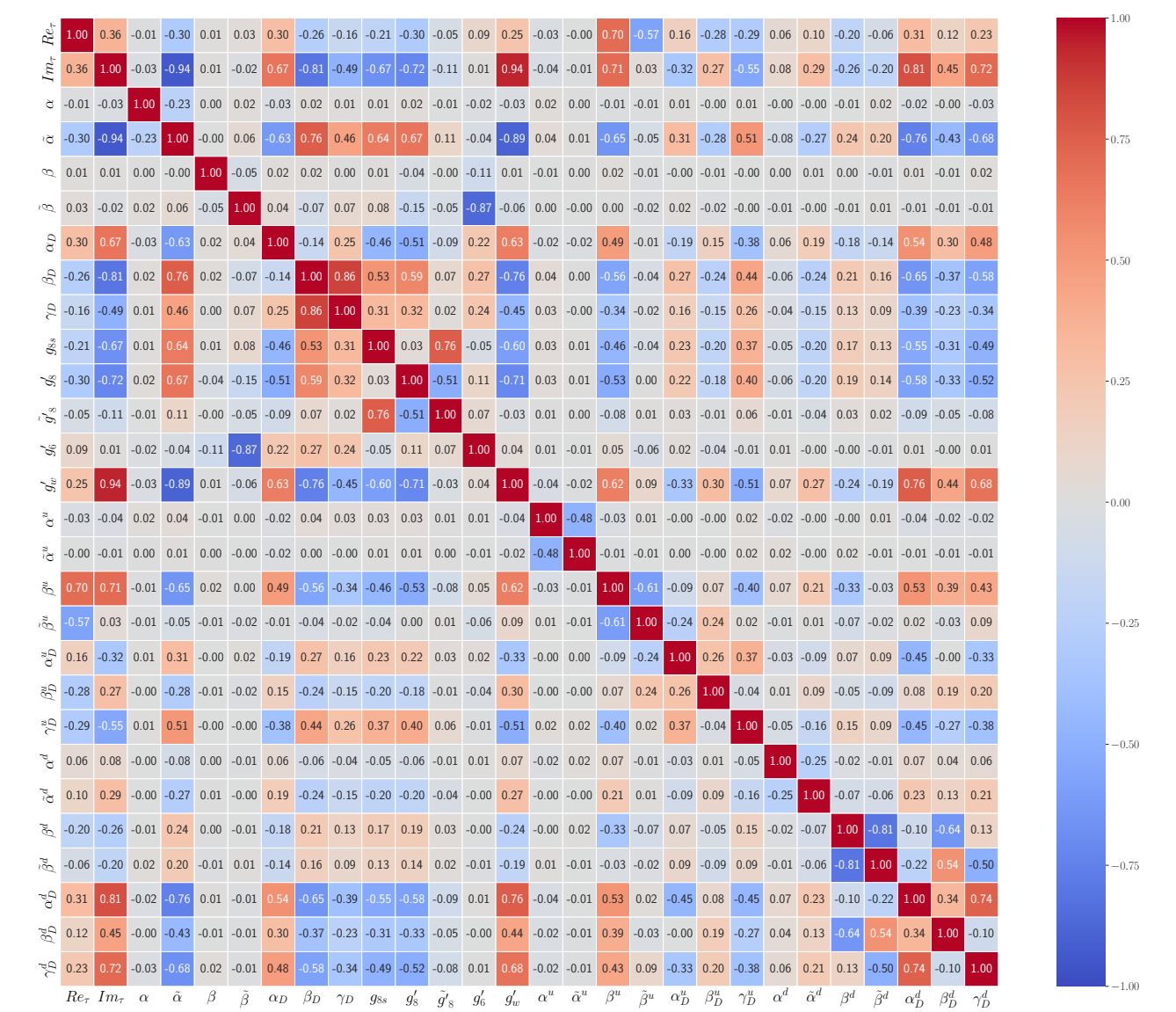
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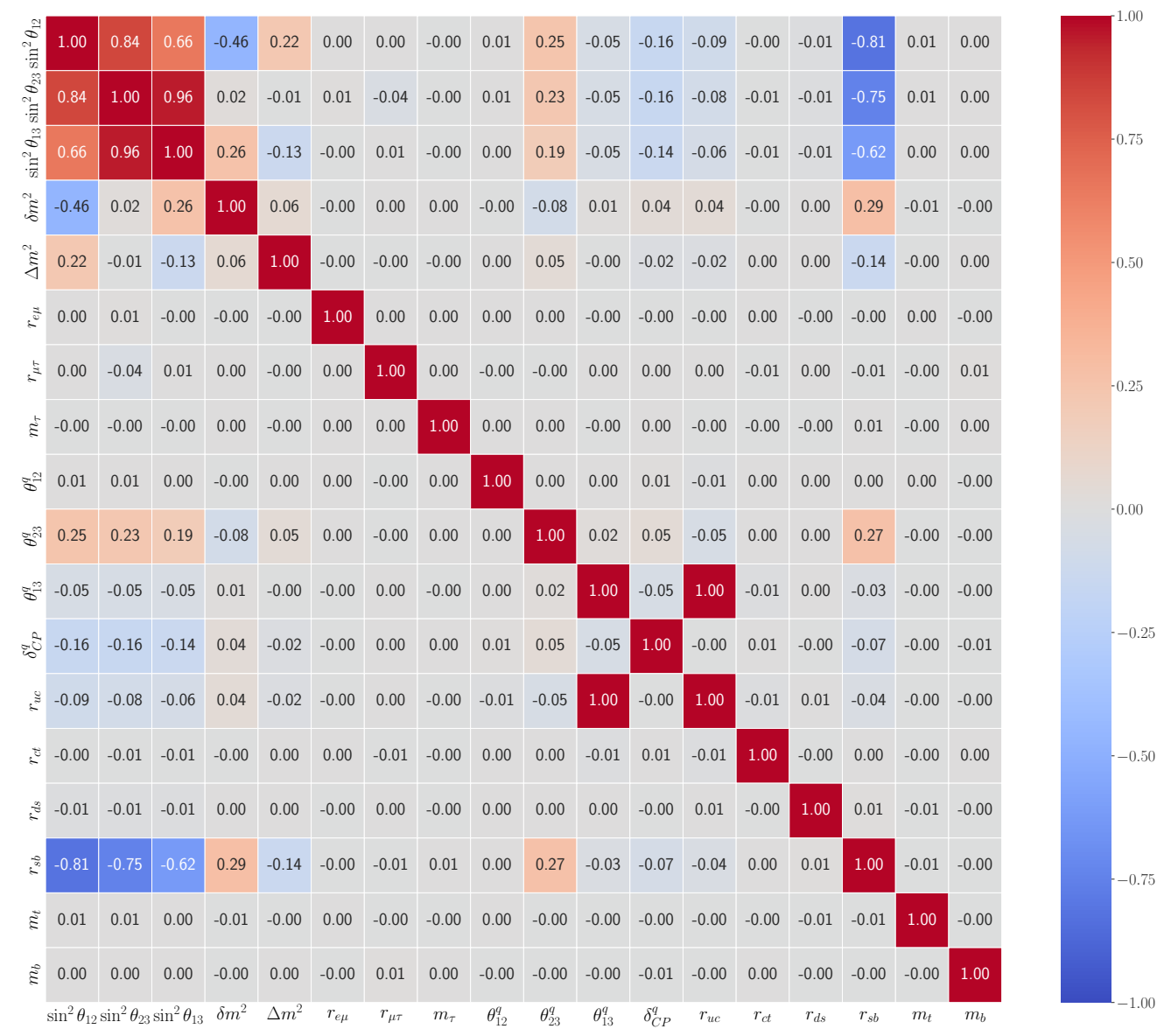
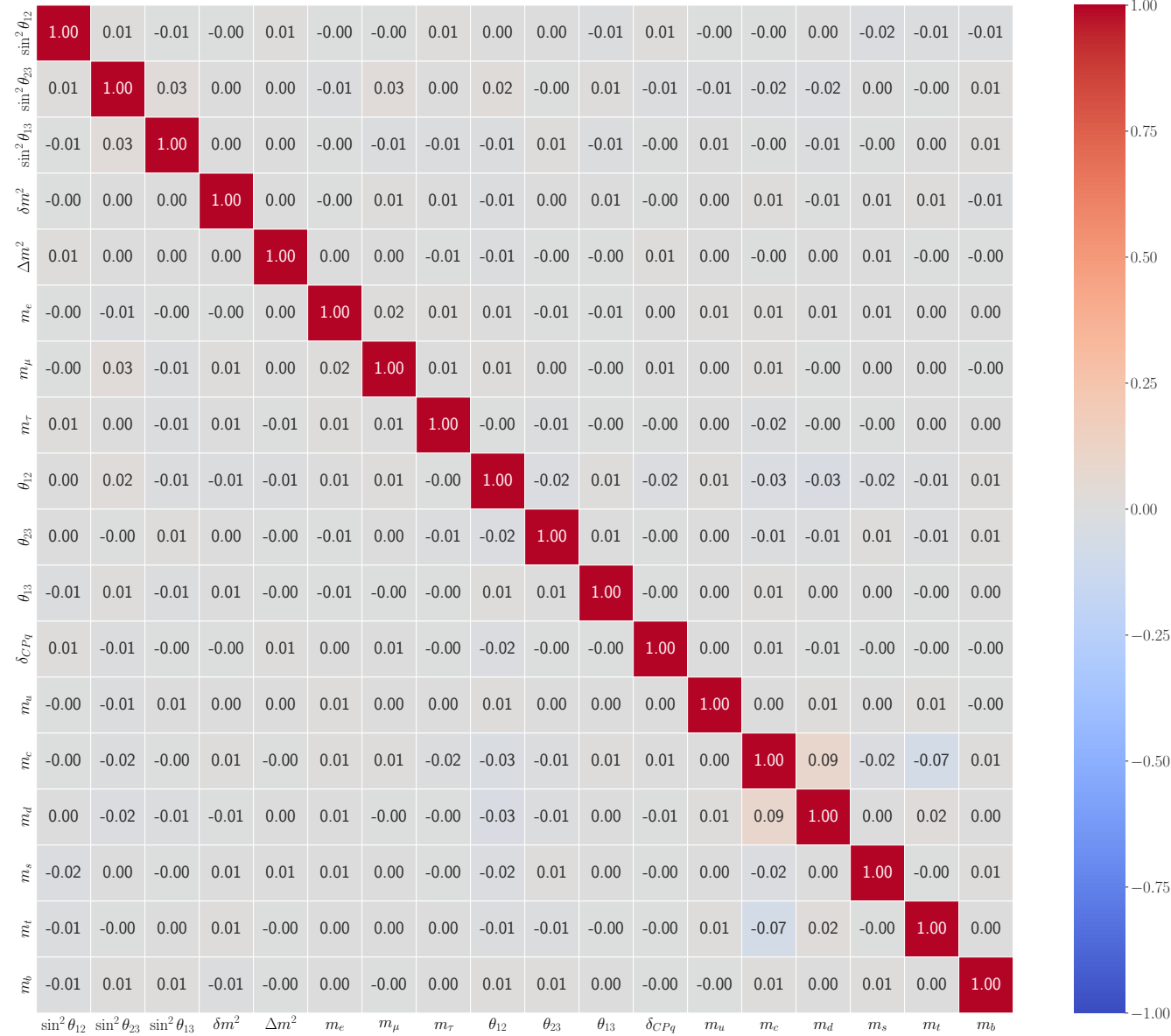
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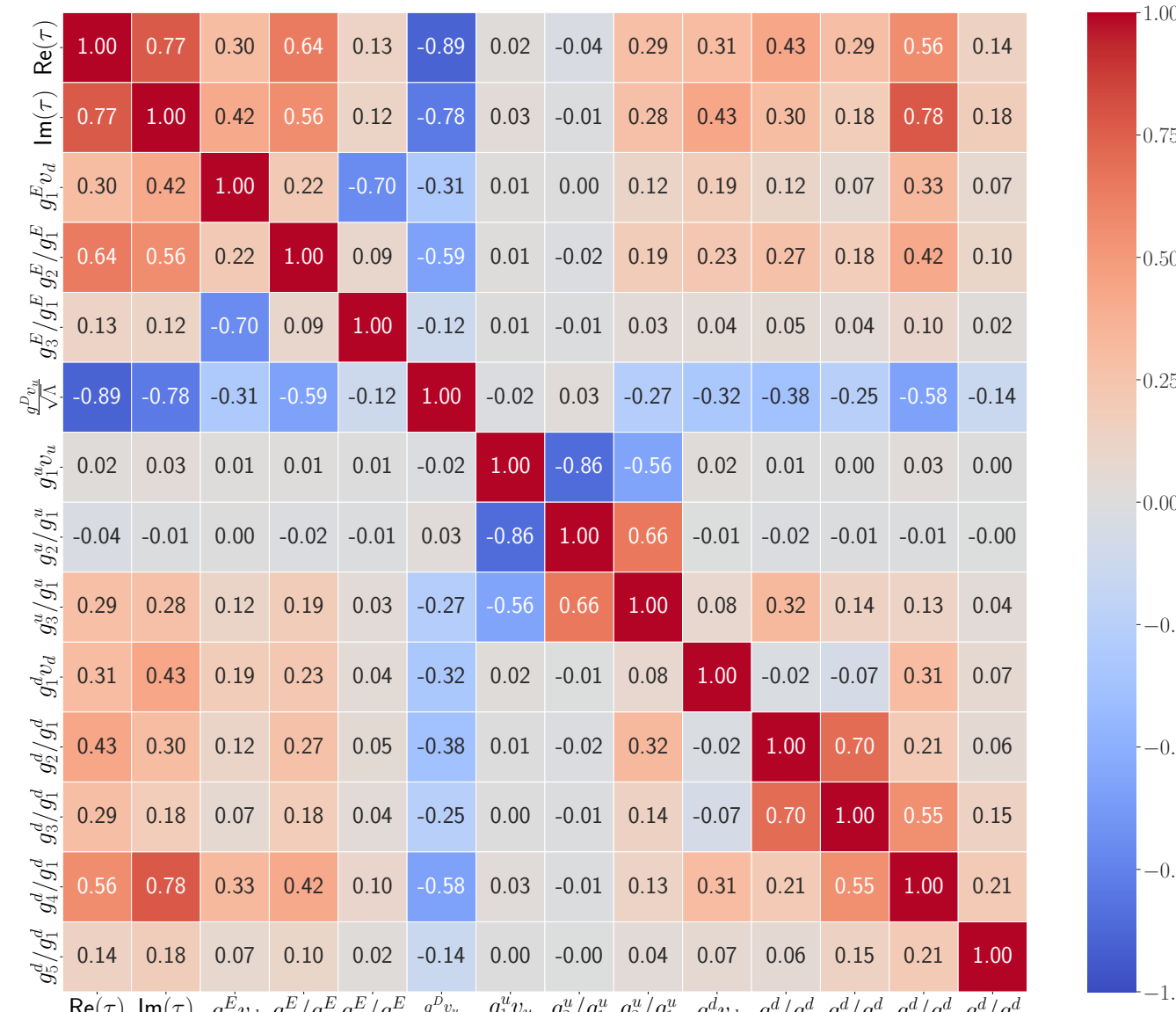
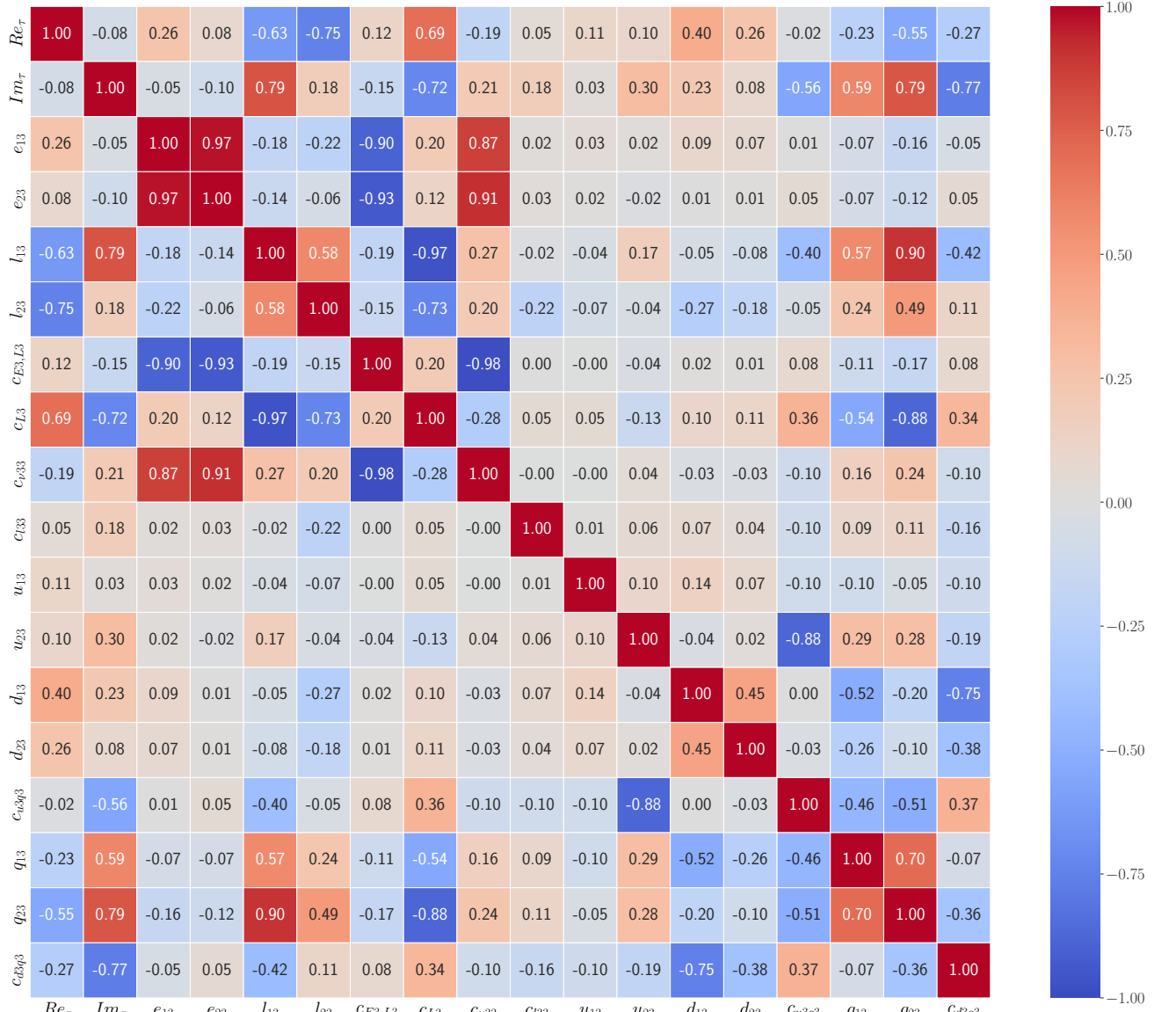
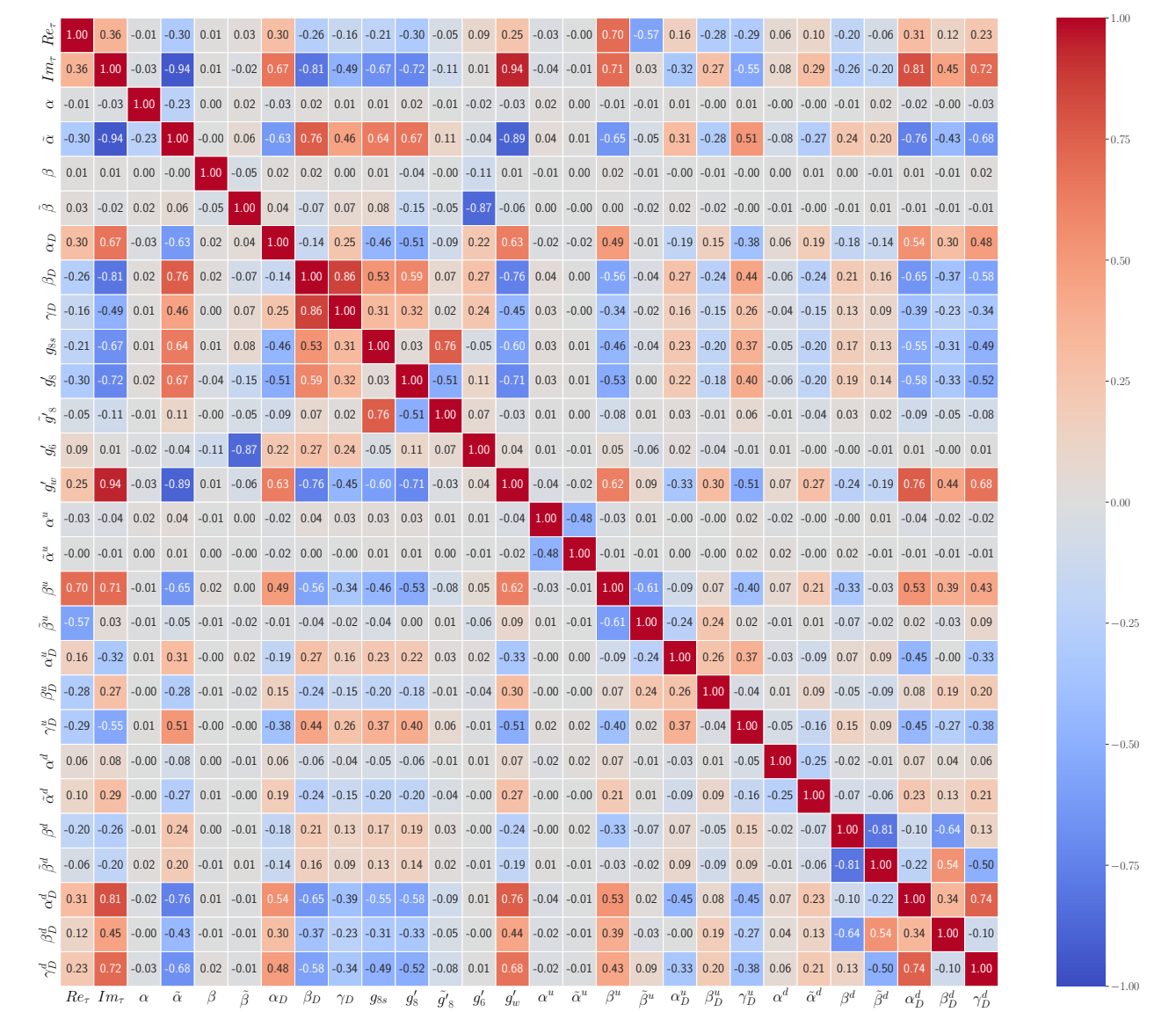
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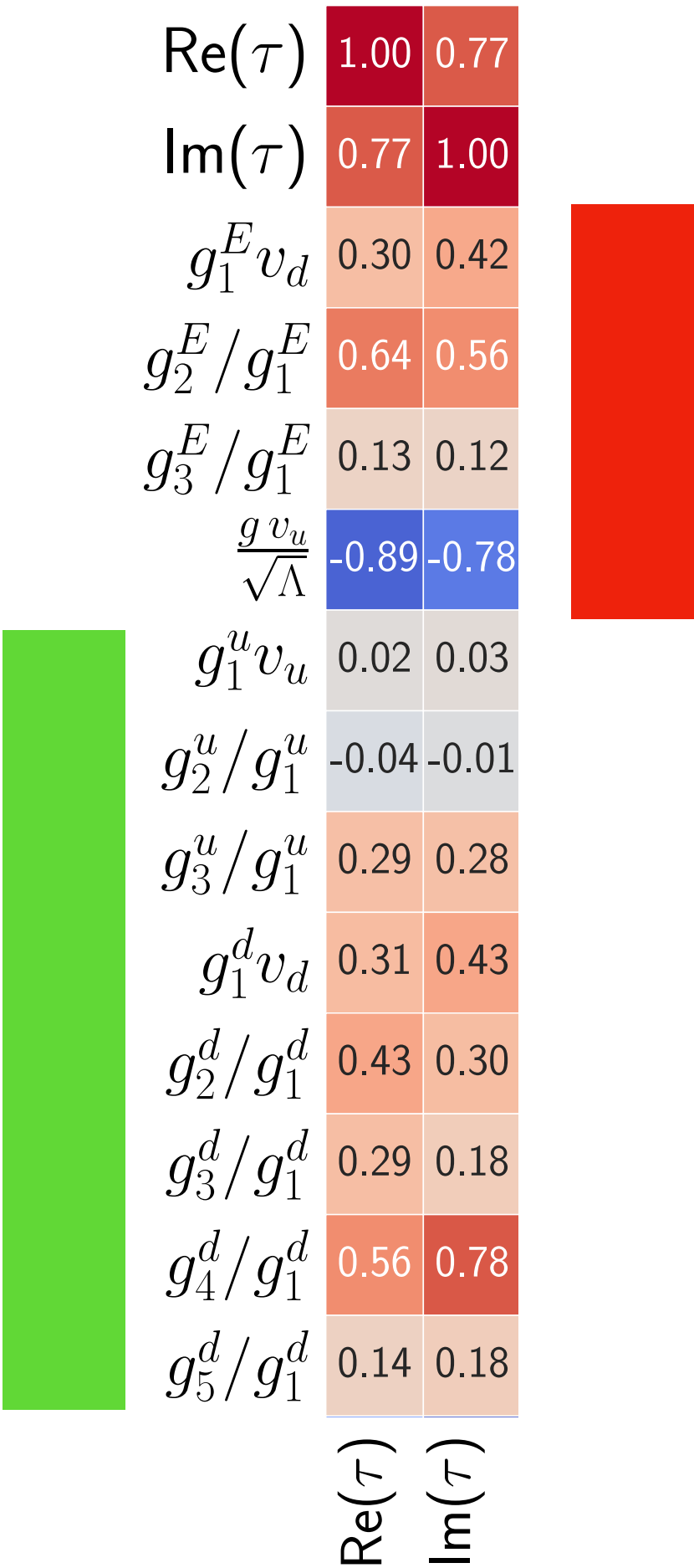
Example of correlations

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For instance $\rightarrow$ model parameters most correlated to τ in the two sectors

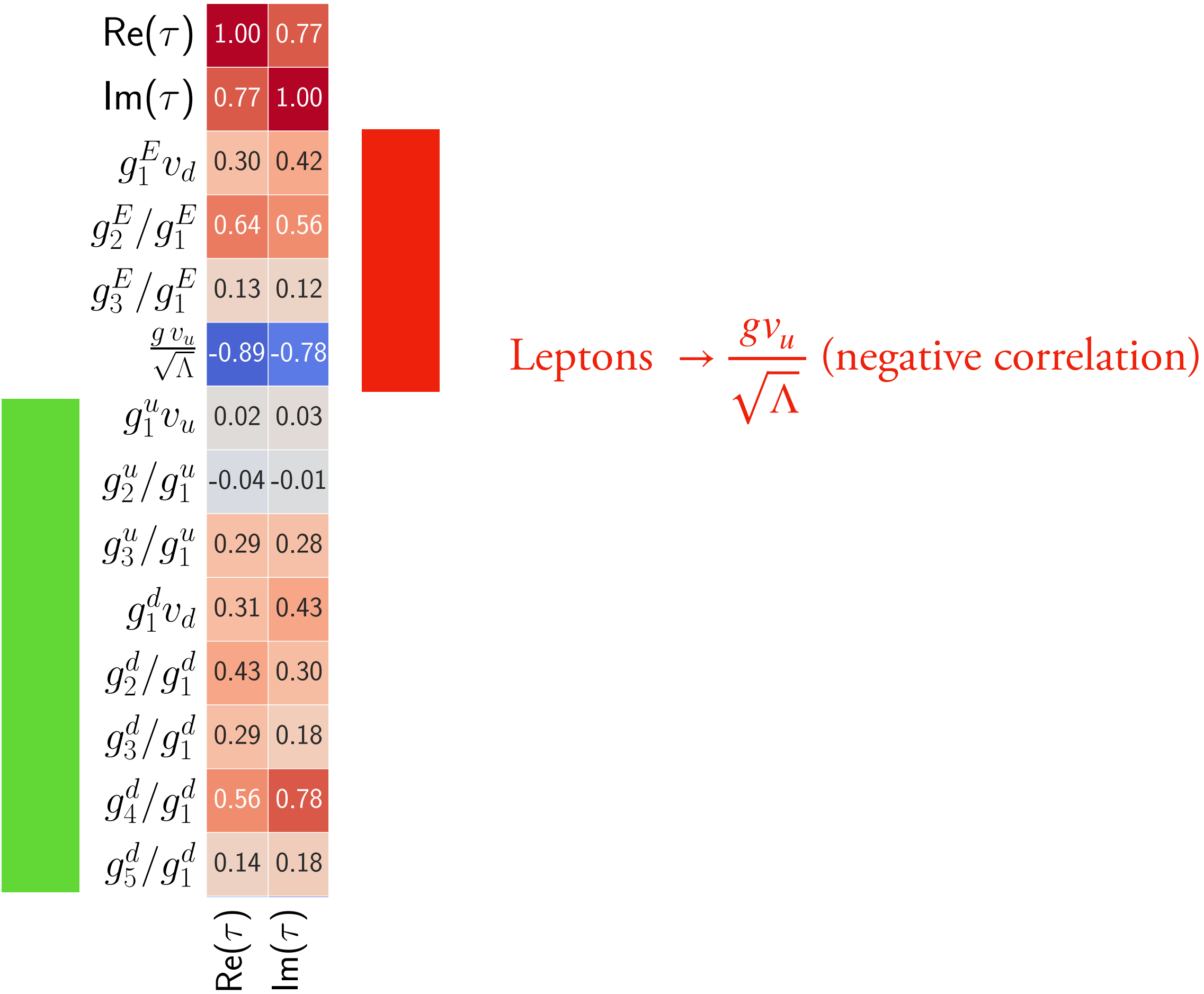
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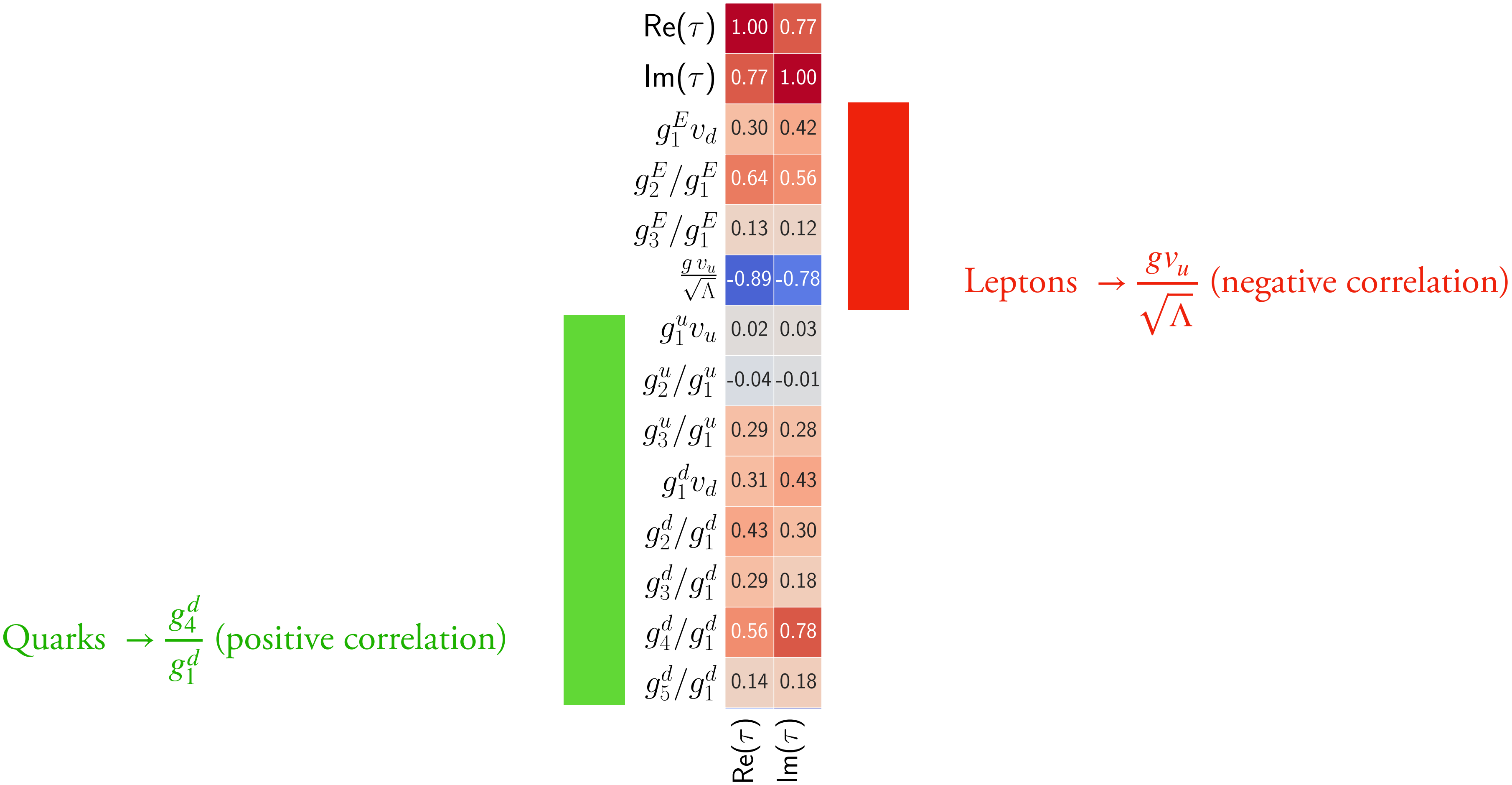
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We can directly verify the correlation among quark and lepton parameters

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$\frac{g_4^d}{g_1^d}$ and $\frac{g v_u}{\sqrt{\Lambda}}$ are, indeed the most (anti)correlated parameters among the two different sectors

$g_1^E v_d$	1.00	0.22	-0.70	-0.31	0.01	0.00	0.12	0.19	0.12	0.07	0.33	0.07
g_2^E / g_1^E	0.22	1.00	0.09	-0.59	0.01	-0.02	0.19	0.23	0.27	0.18	0.42	0.10
g_3^E / g_1^E	-0.70	0.09	1.00	-0.12	0.01	-0.01	0.03	0.04	0.05	0.04	0.10	0.02
$\frac{g v_u}{\sqrt{\Lambda}}$	-0.31	-0.59	-0.12	1.00	-0.02	0.03	-0.27	-0.32	-0.38	-0.25	-0.58	-0.14
$g_1^u v_u$	0.01	0.01	0.01	-0.02	1.00	-0.86	-0.56	0.02	0.01	0.00	0.03	0.00
g_2^u / g_1^u	0.00	-0.02	-0.01	0.03	-0.86	1.00	0.66	-0.01	-0.02	-0.01	-0.01	-0.00
g_3^u / g_1^u	0.12	0.19	0.03	-0.27	-0.56	0.66	1.00	0.08	0.32	0.14	0.13	0.04
$g_1^d v_d$	0.19	0.23	0.04	-0.32	0.02	-0.01	0.08	1.00	-0.02	-0.07	0.31	0.07
g_2^d / g_1^d	0.12	0.27	0.05	-0.38	0.01	-0.02	0.32	-0.02	1.00	0.70	0.21	0.06
g_3^d / g_1^d	0.07	0.18	0.04	-0.25	0.00	-0.01	0.14	-0.07	0.70	1.00	0.55	0.15
g_4^d / g_1^d	0.33	0.42	0.10	-0.58	0.03	-0.01	0.13	0.31	0.21	0.55	1.00	0.21
g_5^d / g_1^d	0.07	0.10	0.02	-0.14	0.00	-0.00	0.04	0.07	0.06	0.15	0.21	1.00
$g_1^E v_d$												
g_2^E / g_1^E												
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g_4^d / g_1^d												
g_5^d / g_1^d												

Example of correlations

We can directly verify the correlation among quark and lepton parameters

$\frac{g_4^d}{g_1^d}$ and $\frac{g v_u}{\sqrt{\Lambda}}$ are, indeed the most (anti)correlated parameters among the two different sectors

We trace back their anti correlation going back to their correlation to the real and imaginary parts of τ

$g_1^E v_d$	1.00	0.22	-0.70	-0.31	0.01	0.00	0.12	0.19	0.12	0.07	0.33	0.07
g_2^E / g_1^E	0.22	1.00	0.09	-0.59	0.01	-0.02	0.19	0.23	0.27	0.18	0.42	0.10
g_3^E / g_1^E	-0.70	0.09	1.00	-0.12	0.01	-0.01	0.03	0.04	0.05	0.04	0.10	0.02
$\frac{g v_u}{\sqrt{\Lambda}}$	-0.31	-0.59	-0.12	1.00	-0.02	0.03	-0.27	-0.32	-0.38	-0.25	-0.58	-0.14
$g_1^u v_u$	0.01	0.01	0.01	-0.02	1.00	-0.86	-0.56	0.02	0.01	0.00	0.03	0.00
g_2^u / g_1^u	0.00	-0.02	-0.01	0.03	-0.86	1.00	0.66	-0.01	-0.02	-0.01	-0.01	-0.00
g_3^u / g_1^u	0.12	0.19	0.03	-0.27	-0.56	0.66	1.00	0.08	0.32	0.14	0.13	0.04
$g_1^d v_d$	0.19	0.23	0.04	-0.32	0.02	-0.01	0.08	1.00	-0.02	-0.07	0.31	0.07
g_2^d / g_1^d	0.12	0.27	0.05	-0.38	0.01	-0.02	0.32	-0.02	1.00	0.70	0.21	0.06
g_3^d / g_1^d	0.07	0.18	0.04	-0.25	0.00	-0.01	0.14	-0.07	0.70	1.00	0.55	0.15
g_4^d / g_1^d	0.33	0.42	0.10	-0.58	0.03	-0.01	0.13	0.31	0.21	0.55	1.00	0.21
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Correlation between Parameters and Observables

Correlation between Parameters and Observables

All neutrino observables are positively correlated to $\frac{g_{\nu_u}}{\sqrt{\Lambda}}$, (in particular $\sin^2 \theta_{12}$ and the other two mixing angles)

$\frac{g_{\nu_u}}{\sqrt{\Lambda}}$	0.71	0.90	0.89	0.29	0.34	0.83	0.00	0.00
$\sin^2 \theta_{12}$								
$\sin^2 \theta_{23}$								
$\sin^2 \theta_{13}$								
δm^2								
Δm^2								
δ_{CP}								
$r_{e\mu}$								
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The quark parameter most correlated to quark observables is $\frac{g_4^d}{g_1^d}$

$\frac{g_4^d}{g_1^d}$	0.11	0.34	-0.02	-0.04	-0.04	0.00	-0.01	0.99	-0.01	-0.01
θ_{12}^q										
θ_{23}^q										
θ_{13}^q										
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r_{uc}										
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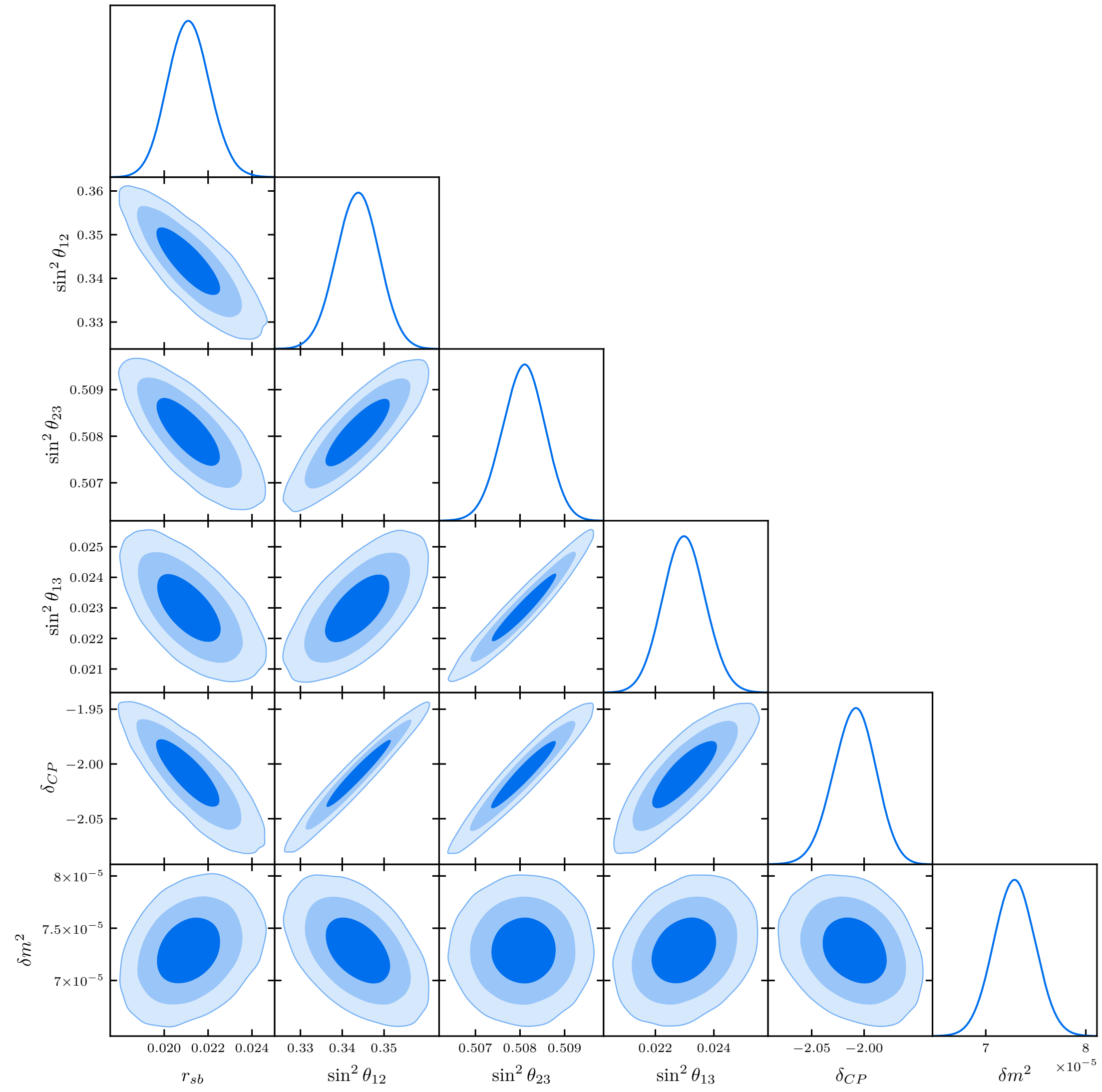
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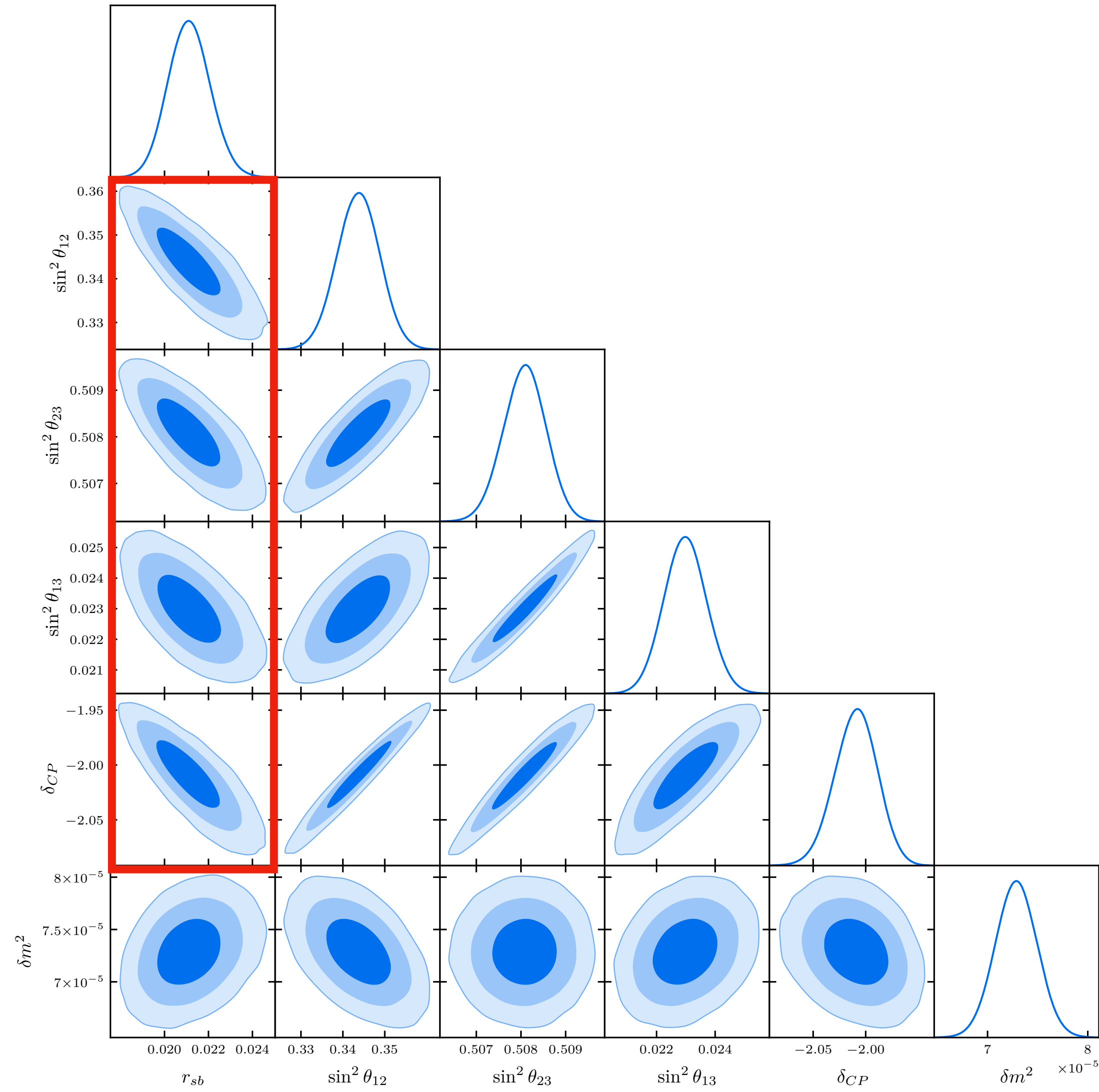


We can presume the presence of a significant anticorrelation between r_{sb} and all neutrino mixing angles

Correlations between quark and lepton observables

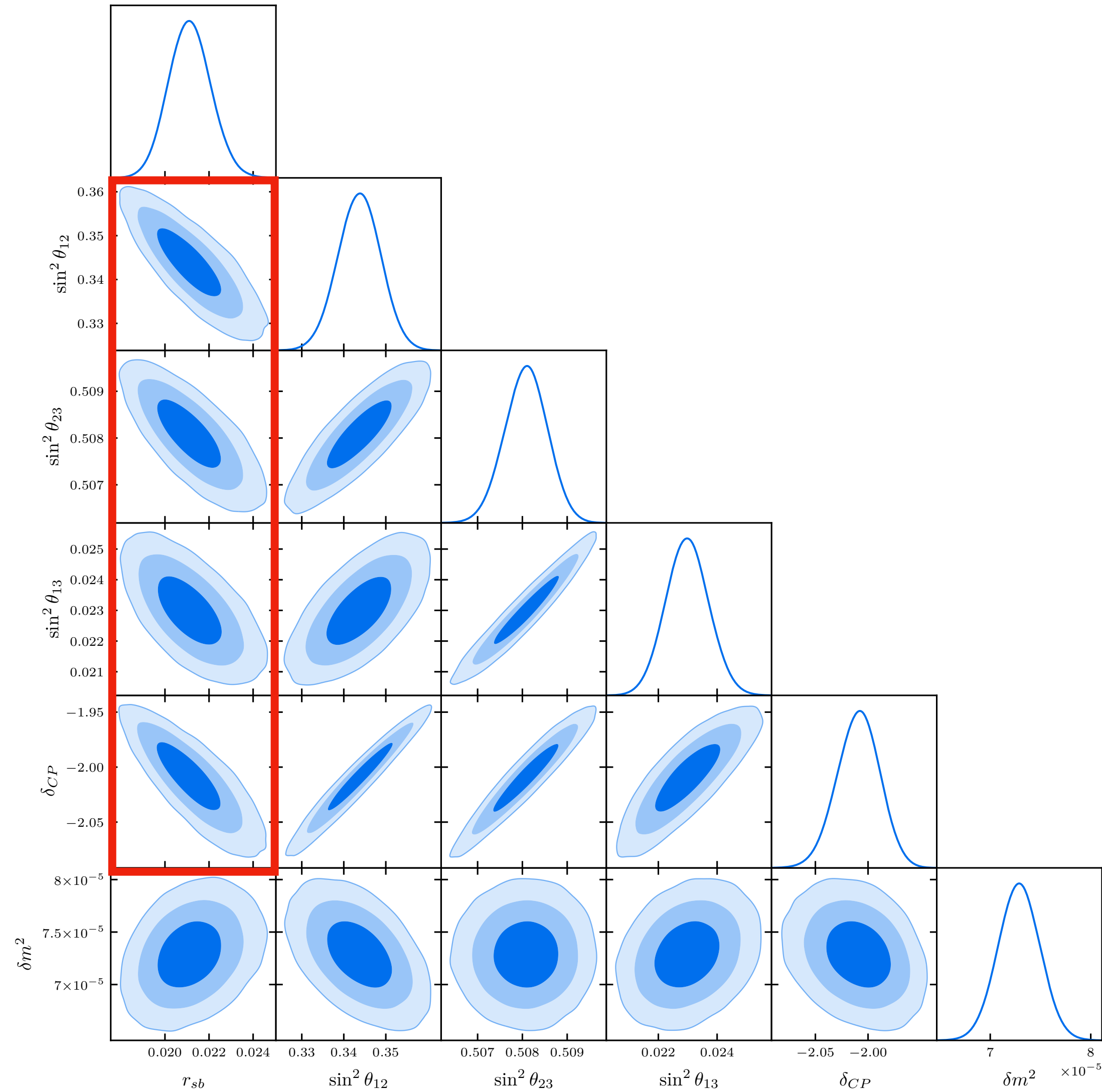


Correlations between quark and lepton observables



Significative anticorrelation between r_{sb} and $\sin^2 \theta_{12}$, $\sin^2 \theta_{23}$, $\sin^2 \theta_{13}$, δ_{CP}

Correlations between quark and lepton observables



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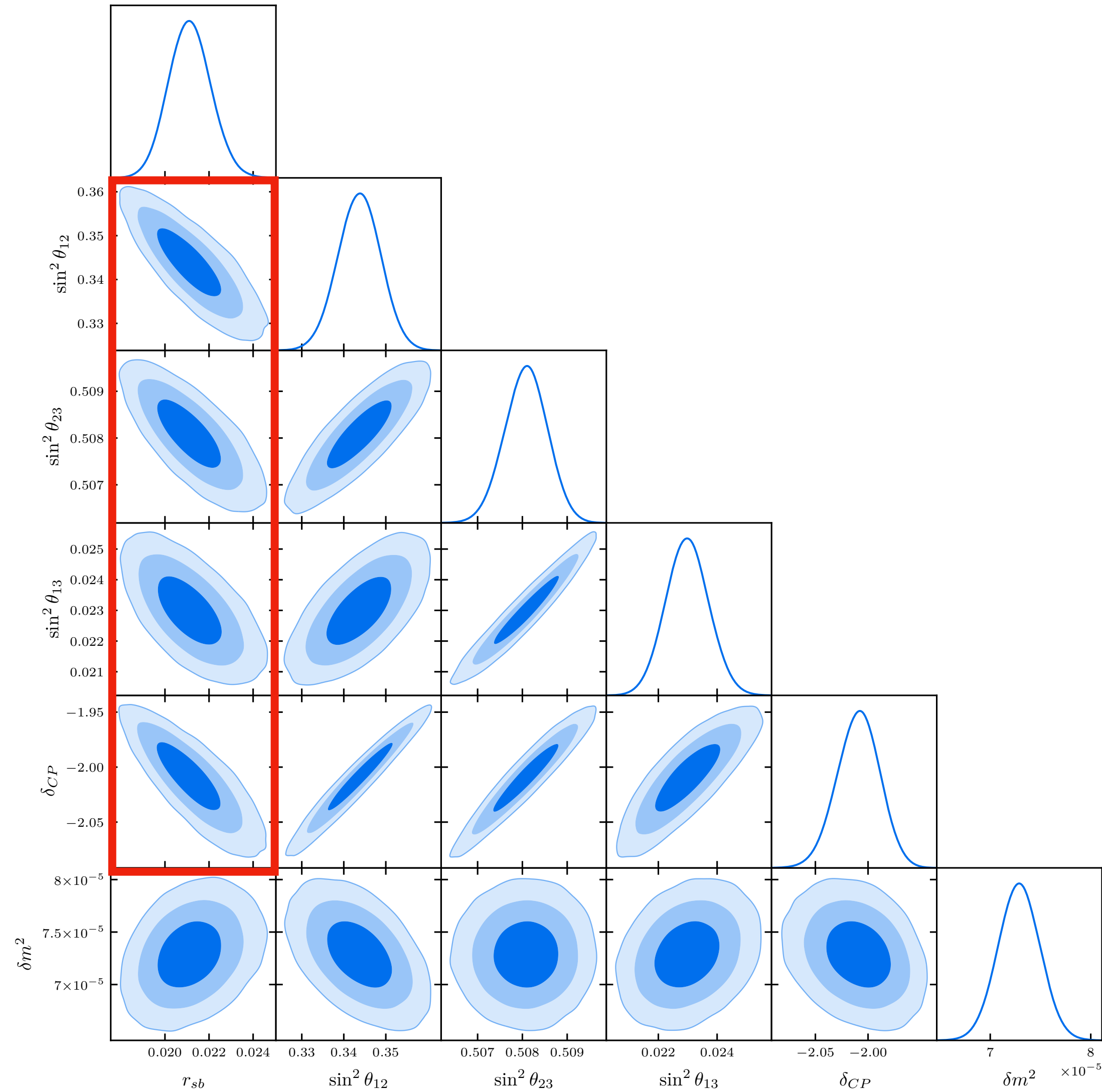
θ_{12}^q	0.01	0.01	0.00	-0.00	0.00	0.01	0.00	-0.00
θ_{23}^q	0.25	0.23	0.19	-0.08	0.05	0.25	0.00	-0.00
θ_{13}^q	-0.05	-0.05	-0.05	0.01	-0.00	-0.06	-0.00	0.00
δ_{CP}^q	-0.16	-0.16	-0.14	0.04	-0.02	-0.16	-0.00	0.00
r_{uc}	-0.09	-0.08	-0.06	0.04	-0.02	-0.09	-0.00	0.00
r_{ct}	-0.00	-0.01	-0.01	-0.00	0.00	-0.00	0.00	-0.01
r_{ds}	-0.01	-0.01	-0.01	0.00	0.00	-0.01	-0.00	0.00
r_{sb}	-0.81	-0.75	-0.62	0.29	-0.14	-0.82	0.00	-0.01
m_t	0.01	0.01	0.00	-0.01	-0.00	0.01	0.00	-0.00
m_b	0.00	0.00	0.00	-0.00	0.00	0.00	-0.00	0.01
	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$	$\sin^2 \theta_{13}$	δm^2	Δm^2	δ_{CP}	$r_{e\mu}$	$r_{\mu\tau}$

Going back to the results of the fit we understand how the 3σ tension on the values of $\sin^2 \theta_{12}$ predicted by the model (0.344) and the measured one (0.303) induces a new tension on r_{sb} in the combined fit



Overestimating $\sin^2 \theta_{12}$ at 3σ causes an underestimation of r_{sb} at $\sim 2.7\sigma$

Correlations between quark and lepton observables



Significative anticorrelation between r_{sb} and $\sin^2 \theta_{12}$, $\sin^2 \theta_{23}$, $\sin^2 \theta_{13}$, δ_{CP}

θ_{12}^q	0.01	0.01	0.00	-0.00	0.00	0.01	0.00	-0.00
θ_{23}^q	0.25	0.23	0.19	-0.08	0.05	0.25	0.00	-0.00
θ_{13}^q	-0.05	-0.05	-0.05	0.01	-0.00	-0.06	-0.00	0.00
δ_{CP}^q	-0.16	-0.16	-0.14	0.04	-0.02	-0.16	-0.00	0.00
r_{uc}	-0.09	-0.08	-0.06	0.04	-0.02	-0.09	-0.00	0.00
r_{ct}	-0.00	-0.01	-0.01	-0.00	0.00	-0.00	0.00	-0.01
r_{ds}	-0.01	-0.01	-0.01	0.00	0.00	-0.01	-0.00	0.00
r_{sb}	-0.81	-0.75	-0.62	0.29	-0.14	-0.82	0.00	-0.01
m_t	0.01	0.01	0.00	-0.01	-0.00	0.01	0.00	-0.00
m_b	0.00	0.00	0.00	-0.00	0.00	0.00	-0.00	0.01
	$\sin^2 \theta_{12}$	$\sin^2 \theta_{23}$	$\sin^2 \theta_{13}$	δm^2	Δm^2	δ_{CP}	$r_{e\mu}$	$r_{\mu\tau}$

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The existence of such correlations between the lepton and quark observables could be a strong indicator of an underlying modular flavour symmetry in particle physics

Conclusions

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Modular symmetries provide a compelling framework for flavor physics, offering an origin for mass hierarchies and CP violation from a discrete symmetry principle

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The choice of the modular group is critical for predictive power. Models with fewer parameters and positive degrees of freedom (like the 2O model, where $N_{\text{dof}} > 0$) allow for robust statistical tests and generate more complex correlation patterns among observables. In contrast, over-parameterized models (like the S_3 example, where $N_{\text{dof}} < 0$) can always achieve a perfect fit but lack true predictive power

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Future Outlook: The tension observed between the predicted value of $\sin^2 \theta_{12}$ and its experimental value in the 2O model highlights how these frameworks can be concretely tested and potentially falsified by future precision measurements. This demonstrates that modular symmetry models are not just theoretical curiosities but are becoming increasingly constrained by experimental data