

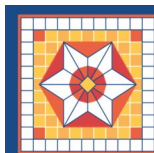
Long live the heavy ALP!

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FLASY

ROME 2025

11TH WORKSHOP

Flavor Symmetries
and Consequences
in Accelerators
and Cosmology



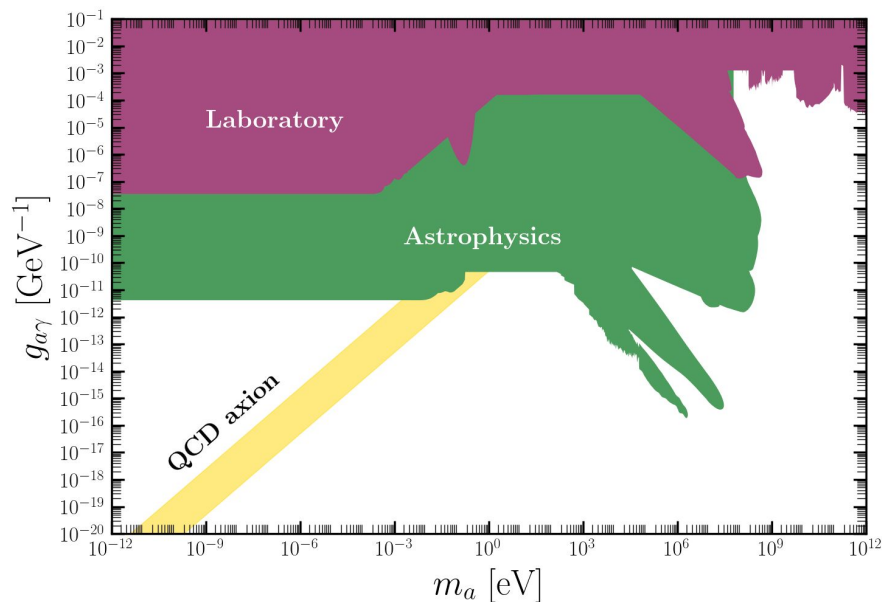
Motivation

Strong CP problem

Most elegant solution: **QCD axion**

[R.D. Peccei, H. R. Quinn, Phys. Rev. D, vol. 16. pp. 1791-1797, 1977]

From SSB of $U(1)_{PQ}$ a pNGB arises with mass generated dynamically from QCD sector, bounded to $m_a f_a \cong m_\pi f_\pi$



In general: $g_{a\gamma} \propto \frac{1}{f_a}$

For $f_a > f_\pi$ is guaranteed $m_a < f_a$ in the band

Outside? Still QCD axion?

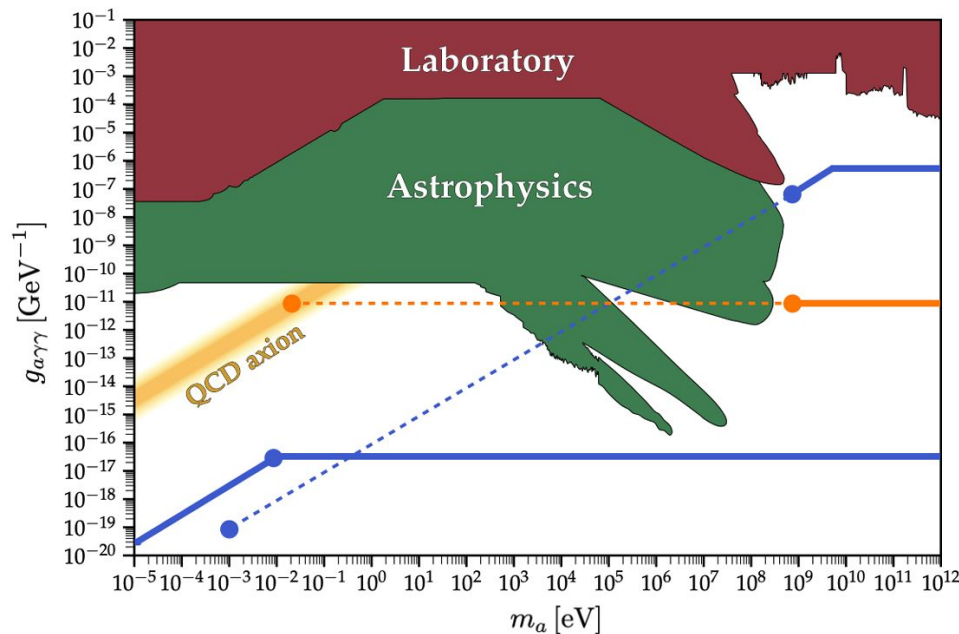
[C. O'Hare, cajohare.github.io/AxionLimits, July 2020]

Motivation

QCD axion besides $m_a f_a \cong m_\pi f_\pi$

Yes, there are non-canonical QCD axions allowed to not satisfy that relation:

Extra-dimensional Maxions: Axion field maximally mixing with N singlet scalars arising from extra dimensions.



Orange line: Canonical QCD axion where all but one modes are decoupled.

Blue line: Extra-dimensional QCD maxion, with no single axion close to canonical band.

[A. de Giorgi, M. Ramos, Phys. Rev. D 111 (2025) 7, 075006]

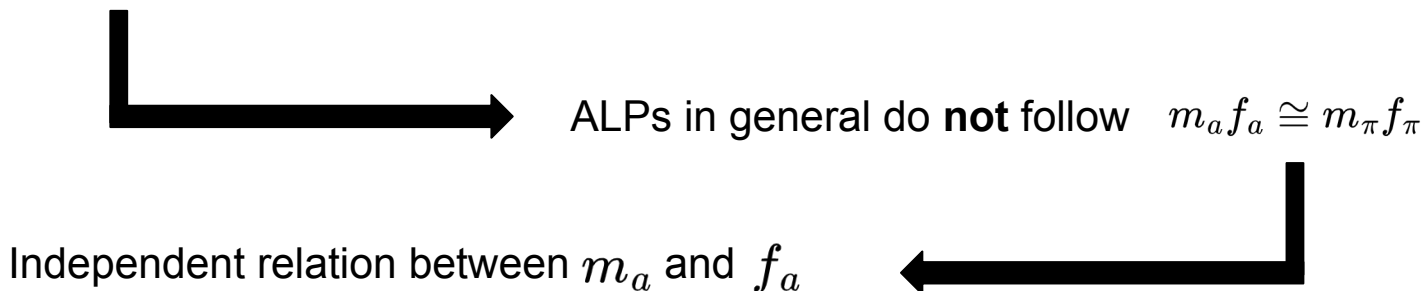
Axion like particles (ALPs)

Yes, ALPs also can solve strong CP problem.

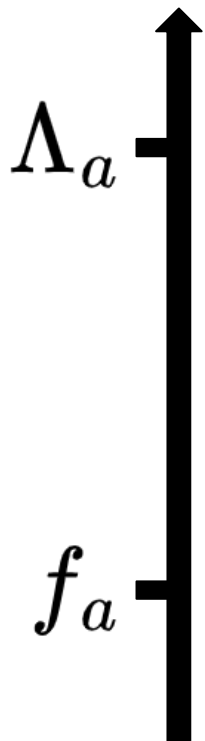
But what is an ALP?

- pNGB arising from the SSB of a **general** global symmetry with an **explicit** mass term.

The dynamical contribution like in the QCD axion is **different** to the explicit term.



Axion like particles (ALPs)



Validity of the EFT beyond f_a

Effective Lagrangian with explicit mass:

[H. Georgi, D. Kaplan, L. Randall, PLB 169B(1986)73]

$$\mathcal{L}_a = \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} m_a^2 a^2 + \frac{\partial_\mu a}{f_a} \sum_f \bar{f} \gamma^\mu \mathbf{c}_f f - \frac{a}{f_a} \sum_b c_b X_{\mu\nu}^b X^{\mu\nu, b}$$

Naive dimensional analysis: $\Lambda \leq 4\pi f_a$


[A. Manohar, H. Georgi, Nucl. Phys. B 234 (1984) 189-212]

General BIAS towards $m_a < f_a$

ALP EFT valid up to $\Lambda_a \leq 4\pi f_a$  $m_a f_a \not\approx m_\pi f_\pi$

Possible scenario where $f_a \leq m_a \leq \Lambda_a$?

YES

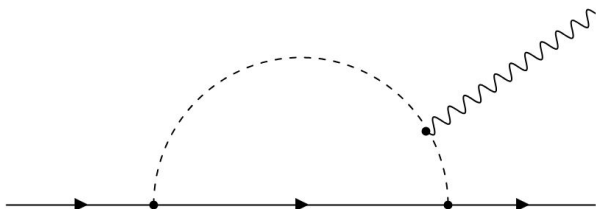
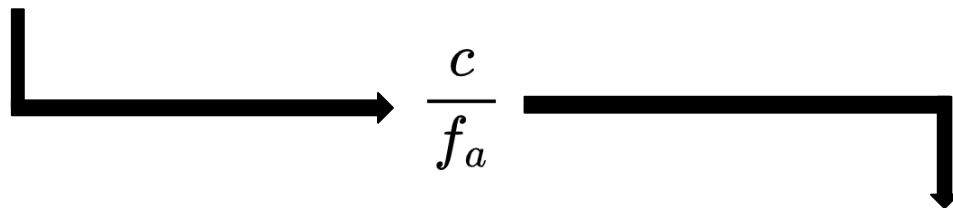
Well-known example: Chiral perturbation theory  $f_\pi < m_\pi < \Lambda_{\chi\text{pt}}$

[S. Weinberg, Physica A96, 327 (1979)]

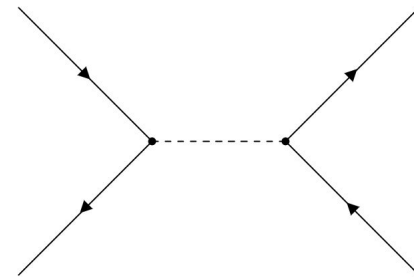
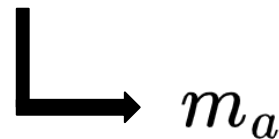
Does it have impact?

Two ways to test:

1) Clean observable in a 2-2 scattering for high energies



Clean 1-loop observable



Does it have impact?

Two ways to test:

1) Clean observation a 2-2 scale lies

Not enough
precision!

2) observation

m_a

Procedure

Does it have impact?

Two ways to test:

2) Available parameter space for 5 clean and very precise lepton observables

Obs.

Exp.

SM contribution

| | | |
|---------------------------|---|---|
| $(g - 2)_e$ | $a_e^{\text{exp}} = 0.00115965218062(12)$ [1] | $a_e^{\text{SM-Rb}} = 0.001159652180252(95)$ [2] |
| $(g - 2)_\mu$ | $a_\mu^{\text{exp}} = 0.00116592059(22)$ [3] | $a_\mu^{\text{SM-Lattice}} = 0.00116592033(62)$ [4] |
| $\mu \rightarrow e\gamma$ | $Br(\mu \rightarrow e\gamma) < 1.5 \cdot 10^{-13}$ (90% C.L.) [5] | $< 10^{-54}$ |
| $\mu \rightarrow 3e$ | $Br(\mu \rightarrow 3e) < 1.0 \cdot 10^{-12}$ (90% C.L.) [6] | |
| $\mu N \rightarrow e N$ | $Br(\mu^- \text{Au} \rightarrow e^- \text{Au}) < 7 \cdot 10^{-13}$ (90% C.L.) [7] | |

Procedure

ALP-Lepton Lagrangian

$$\mathcal{L}_a \supset \mathcal{L}_{\partial a}^l = \frac{\partial_\mu a}{f_a} \left[\bar{L}'_L \gamma^\mu \mathbf{c}'_L L'_L + \bar{e}'_R \gamma^\mu \mathbf{c}'_e e'_R \right] \supset \mathcal{L}_{c.l} = \frac{\partial_\mu a}{f_a} \sum_{ij} \left[\bar{l}_i \gamma^\mu c_{ij} l_j \right]$$

UV model independent

Only consider the first two generations

\mathbf{c}'_f are 2×2 matrix in flavour space and $c_{ij} = c_{ij}^L P_L + c_{ij}^R P_R$ for i, j electron and muon

Allowed both LFC and LFV couplings.

We also work in the range $m_a \gg m_\mu$

$$(g - 2)_i$$

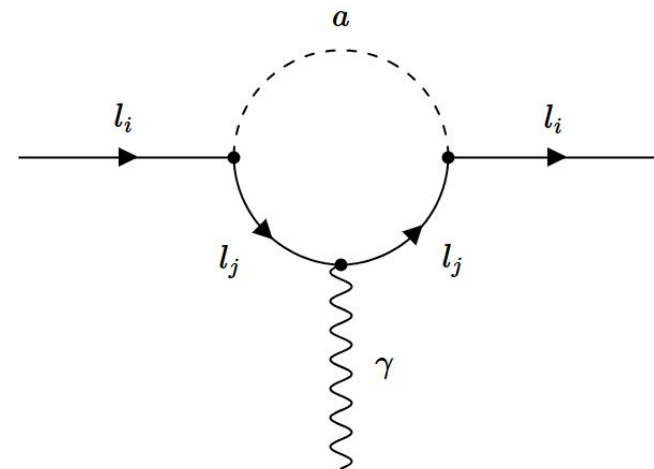
Up to one loop

Red: FV Blue: FC

Dependence:

$$a_e^{\text{ALP}} \propto \frac{m_e m_\mu^3}{f_a^2 m_a^2} \text{Re}[c_{e\mu}^L c_{\mu e}^R] \log\left(\frac{m_a^2}{m_\mu^2}\right)$$

$$a_\mu^{\text{ALP}} \propto \frac{m_\mu^4}{f_a^2 m_a^2} \left[- (c_{\mu\mu}^L - c_{\mu\mu}^R)^2 \log\left(\frac{m_a^2}{m_\mu^2}\right) + (|c_{\mu e}^L|^2 + |c_{\mu e}^R|^2) \right]$$



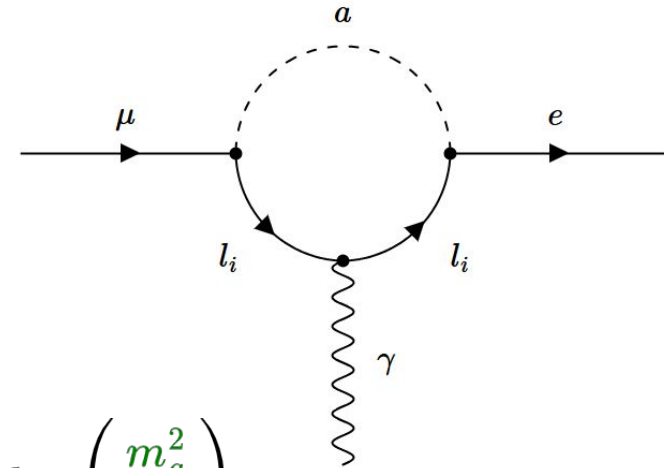
$$\mu \rightarrow e \gamma$$

Up to one loop

Red: FV Blue: FC

Dependence:

$$Br(\mu \rightarrow e \gamma) \propto \frac{\alpha_{\text{em}} m_\mu^2}{G_F^2 f_a^4 m_a^4} (c_{\mu\mu}^L - c_{\mu\mu}^R)^2 (|c_{\mu e}^L|^2 + |c_{\mu e}^R|^2) \log \left(\frac{m_a^2}{m_\mu^2} \right)$$



$$\mu N \rightarrow e N$$

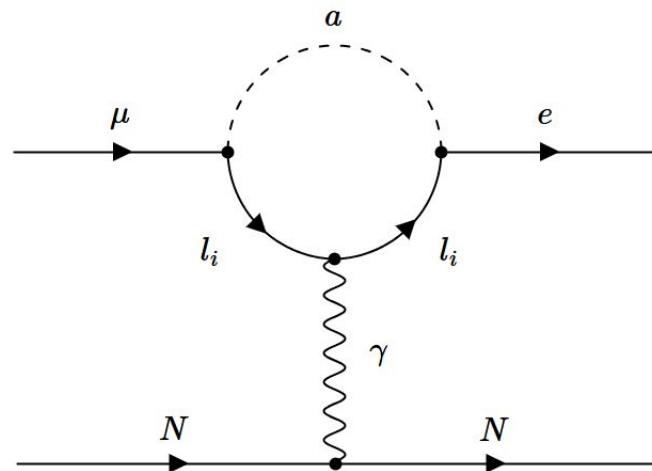
Up to one loop (2 more diagrams)

Red: FV Blue: FC

Dependence:

$$Br(\mu N \rightarrow e N) \propto \frac{m_\mu^8}{f_a^4 m_a^4} (c_{\mu\mu}^L - c_{\mu\mu}^R)^2 \log^2 \left(\frac{m_a^2}{m_\mu^2} \right) \\ \times \left[a(|c_{\mu e}^L|^2 + |c_{\mu e}^R|^2) - b \operatorname{Re}[c_{e\mu}^L c_{\mu e}^R] \right]$$

Besides a prefactor accounting for N . Also: $a/b \sim 1$



$$\mu \rightarrow 3e$$

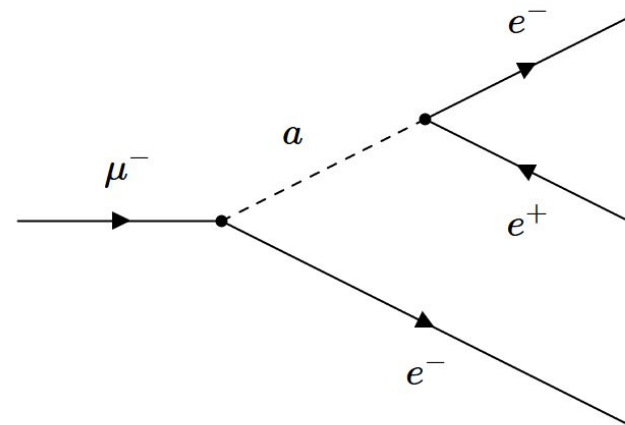
Only tree level (1 more diagram)

Red: FV Blue: FC

Penguin diagram suppressed by α_{em}

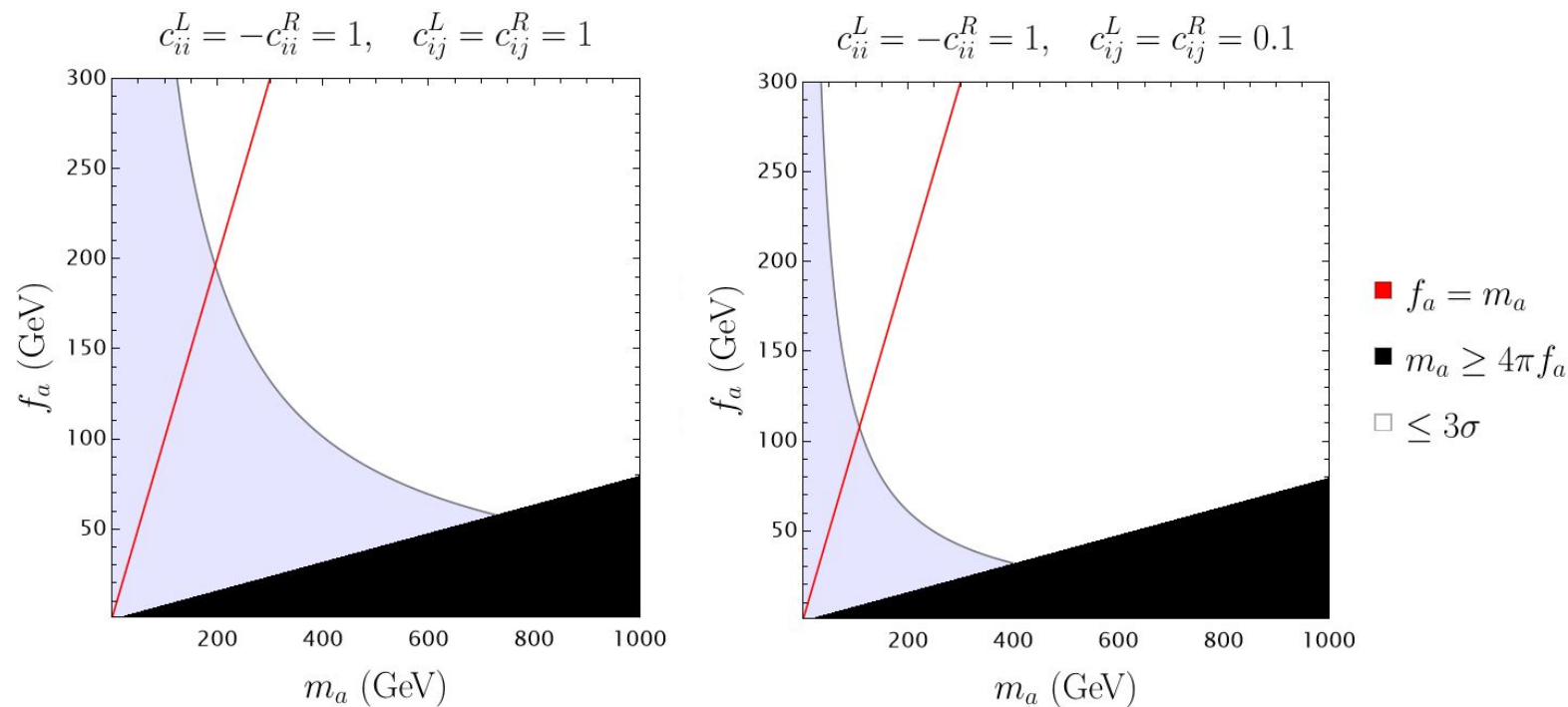
Dependence:

$$Br(\mu \rightarrow 3e) \propto \frac{m_e^2 m_\mu^2}{G_F^2 f_a^4 m_a^4} (c_{ee}^L - c_{ee}^R)^2 (|c_{\mu e}^L|^2 + |c_{\mu e}^R|^2)$$



f_a vs m_a

Relevant process: $\mu \rightarrow e\gamma$



Results

f_a vs m_a

Future prospects for $\mu \rightarrow e\gamma$ and

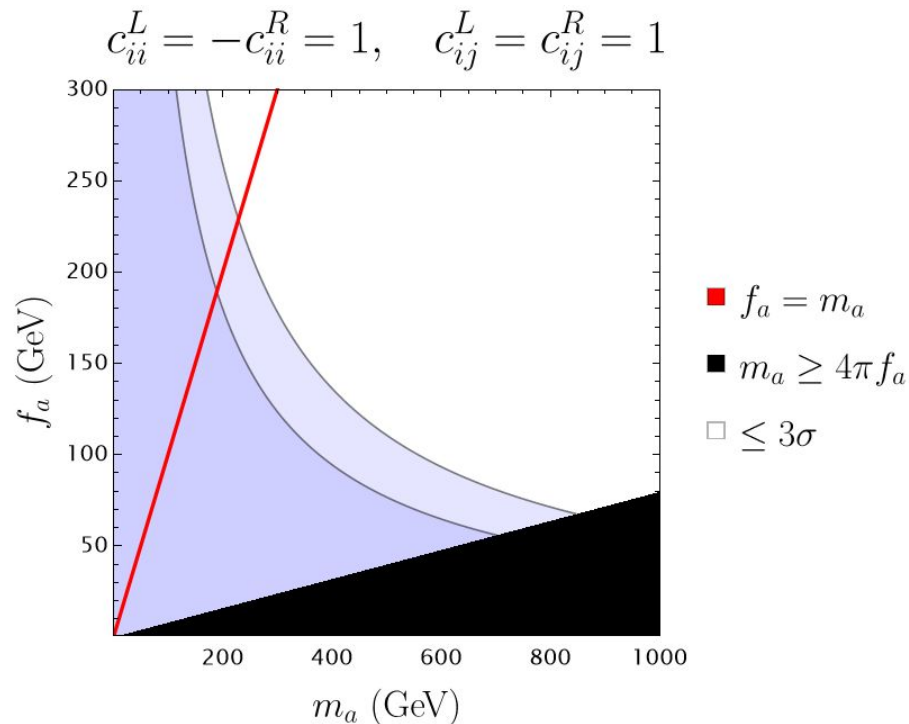
$\mu^- \text{Al} \rightarrow e^- \text{Al}$

$$Br(\mu \rightarrow e\gamma) < 6 \cdot 10^{-14} \text{ (90\% C.L)}$$

[K. Afanaciev et al., arXiv: 2504.15711]

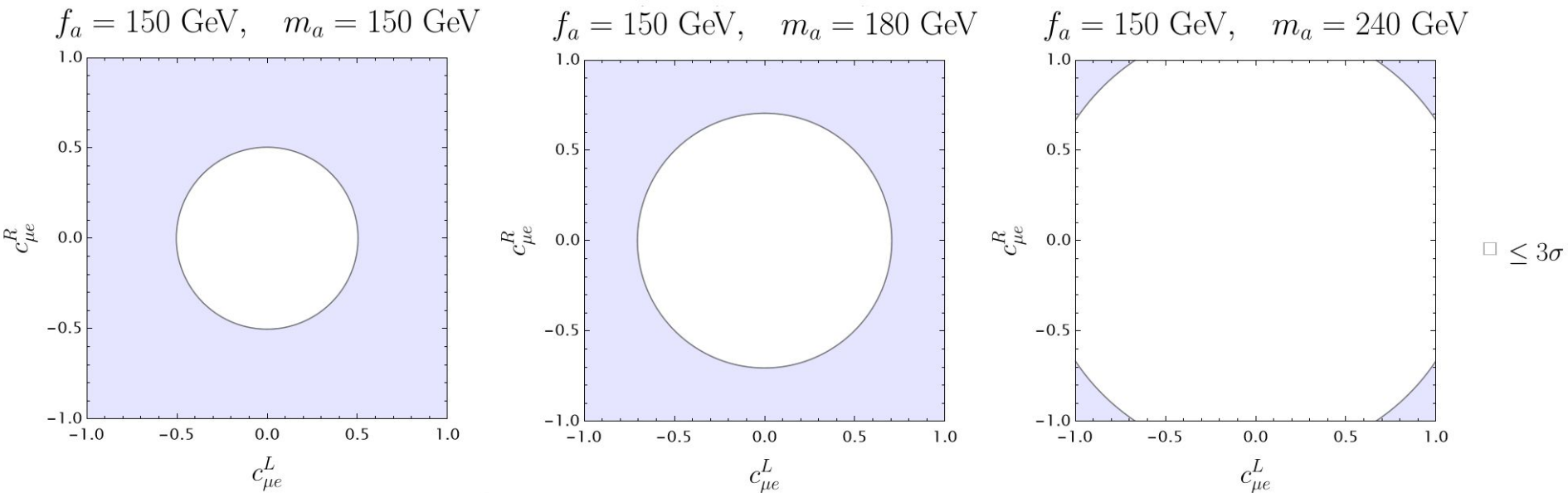
$$Br(\mu^- \text{Al} \rightarrow e^- \text{Al}) < 8 \cdot 10^{-17} \text{ (90\% C.L)}$$

[R. H. Bernstein, Front. in Phys. 7 (2019) 1]



Results

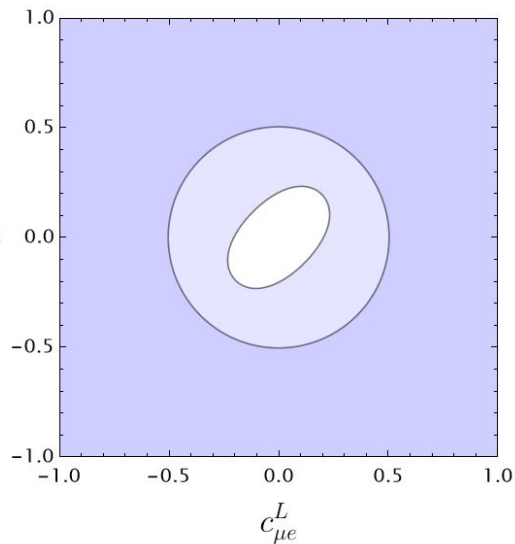
$$c_{\mu e}^R \text{ vs } c_{\mu e}^L, \quad c_{ii}^L = -c_{ii}^R = 1$$



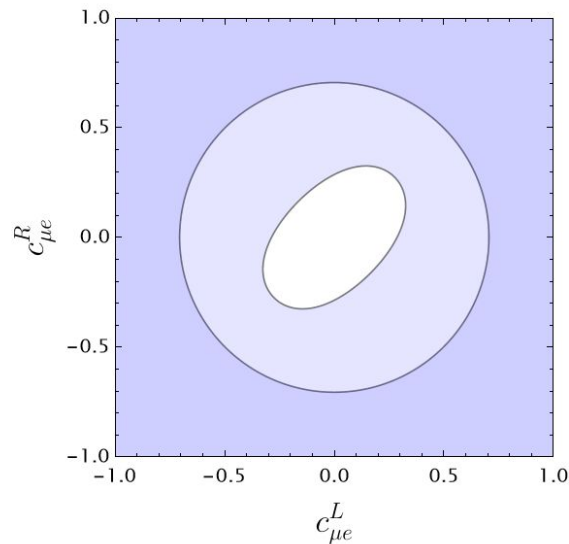
$$c_{\mu e}^R \text{ vs } c_{\mu e}^L, \quad c_{ii}^L = -c_{ii}^R = 1$$

For the previous future prospects:

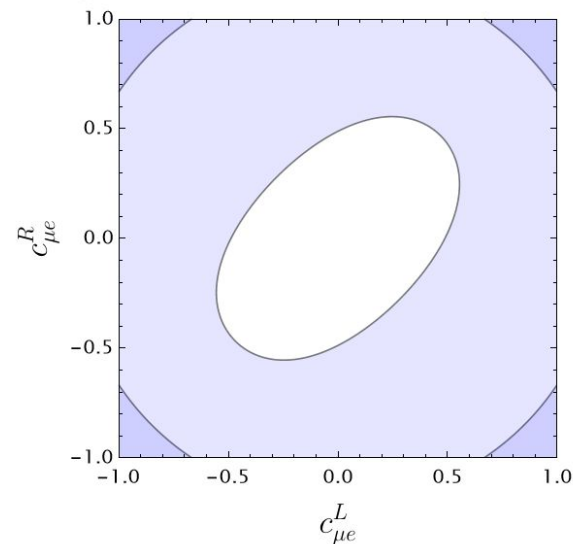
$f_a = 150 \text{ GeV}, \quad m_a = 150 \text{ GeV}$



$f_a = 150 \text{ GeV}, \quad m_a = 180 \text{ GeV}$



$f_a = 150 \text{ GeV}, \quad m_a = 240 \text{ GeV}$



$\square \leq 3\sigma$

- No reason to not consider $f_a \leq m_a \leq \Lambda_a$
- Maybe answers are hidden within that regime.
- The proof of concept depicted here can be translated to general EFT scenarios.

Thank you for your attention

References not shown:

- [1] - [PDG]
- [2] - [L. Morel et al., Nature 588 (2020) 7836, 61-65]
- [3] - [D. P. Aguillard et al., Phys. Rev. D 110 (2024) 3, 032009]
- [4] - [R. Aliberti et al., arXiv: 2002.12347]
- [5] - [K. Afanaciev et al., arXiv: 2504.15711]
- [6] - [U. Bellgardt et al., Nucl. Phys. B 299 (1988) 1-6]
- [7] - [W. H. Bertl et al., Eur. Phys. J. C 47 (2006) 337-346]

Work supported by:

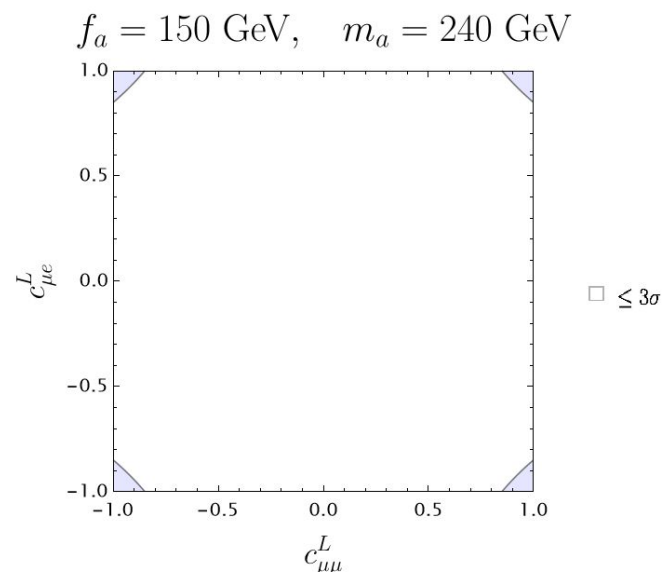
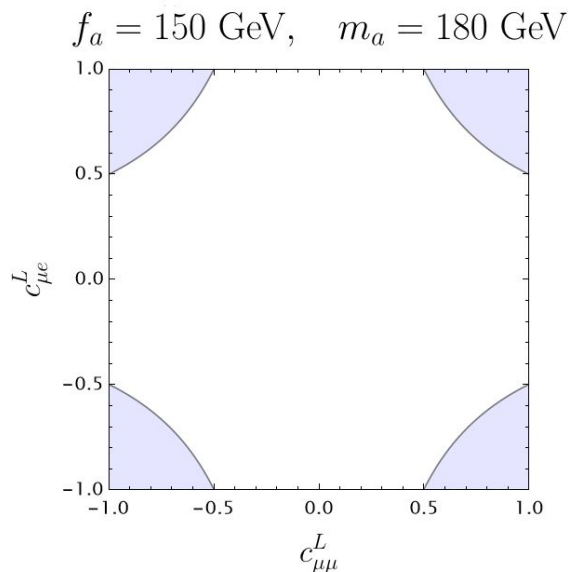
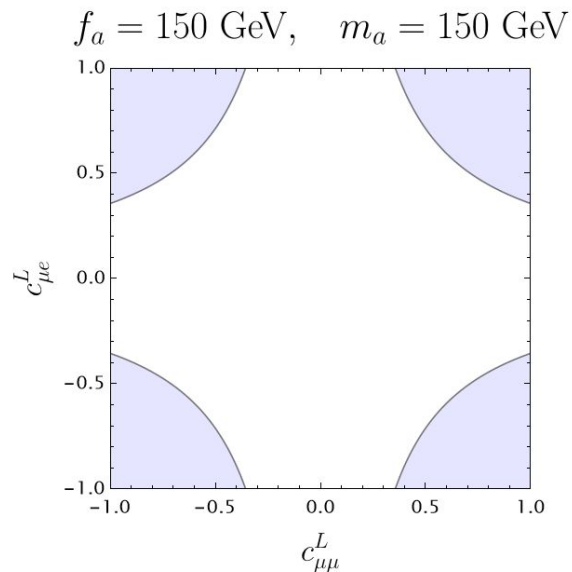
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EXCELENCIA
SEVERO
OCHOA

Parameter scan and fit

$$c_{\mu e}^L \text{ vs } c_{\mu\mu}^L$$



Parameter scan and fit

$$c_{\mu\mu}^R \text{ vs } c_{\mu\mu}^L, \quad c_{\mu e}^L = c_{\mu e}^R = 1$$

