

Muon-decay parameters at COHERENT

New horizons for CE ν NS experiments

Sergio Cruz-Alzaga

based on 2502.18175 in collaboration with
Martín González-Alonso (IFIC), Suraj Prakash (IFIC) and Victor Bresó-Pla (ITP Heidelberg)



EXCELENCIA
SEVERO
OCHOA



GENERALITAT
VALENCIANA

Gen=T

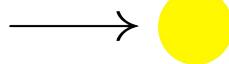


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NextGenerationEU



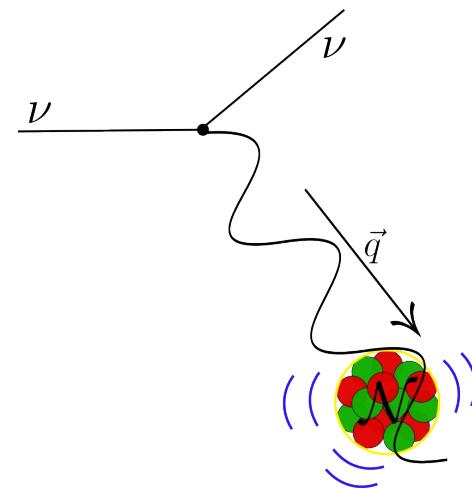
CE ν NS

(Coherent Elastic Neutrino-Nucleus Scattering)



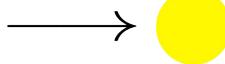
$$\propto Q_W = Z(1 - 4 \sin^2 \theta_W) - N \simeq -N$$

$$s_W^2 = 0.23$$



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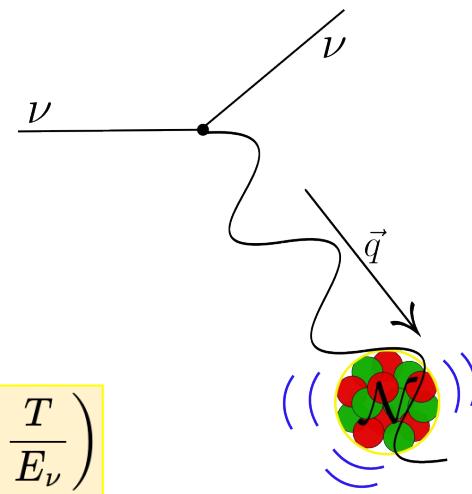
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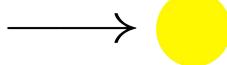
$$s_W^2 = 0.23$$

$$\frac{d\sigma}{dT} = \frac{m_N \mathcal{F}^2(q^2)}{8\pi v^4} Q_W^2 \left(1 - \frac{m_N T}{2E_\nu^2} - \frac{T}{E_\nu} \right)$$



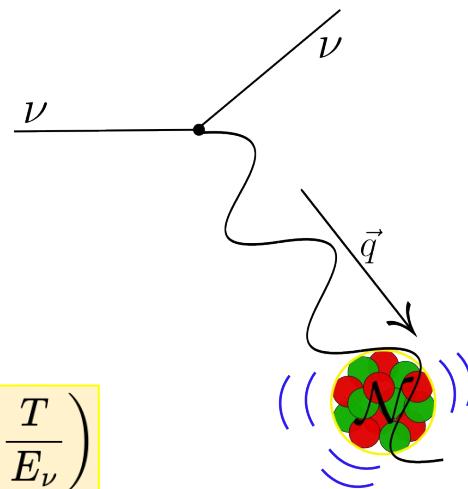
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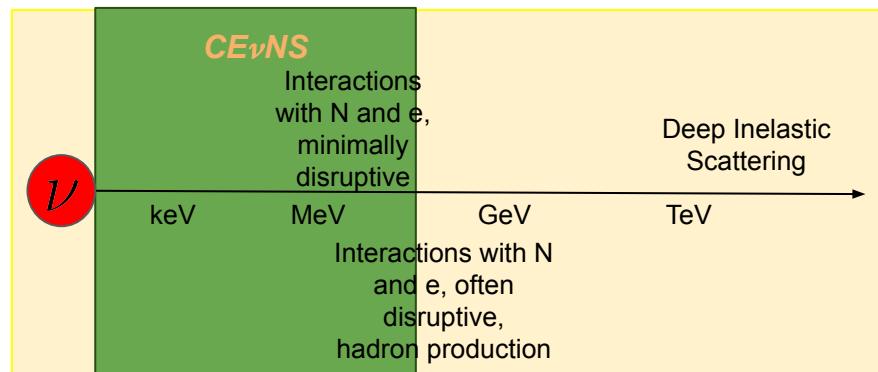
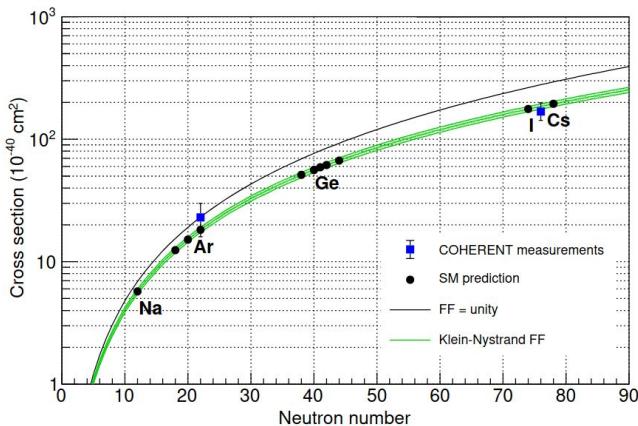


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Adapted from De Romeri

The COHERENT Experiment

*Spallation Neutron Source
(Oak Ridge National Laboratory)*



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- First observation of $CE\nu NS$

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- Neutrino detection

$CE\nu NS$

Targets → CsI, LAr, Ge

D. Akimov et al. (COHERENT) Science 357, 1123–1126 (2017)
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S. Adamski et al. (COHERENT), arXiv: 2406.13806

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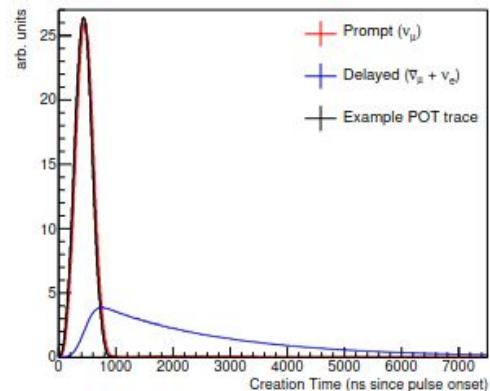
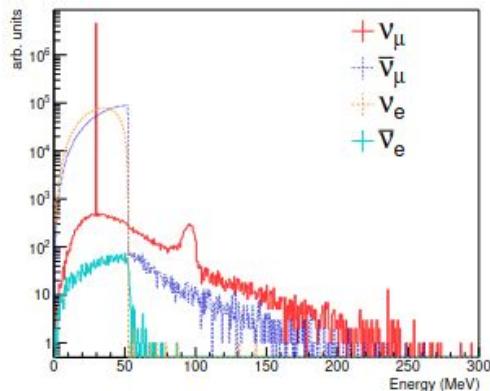
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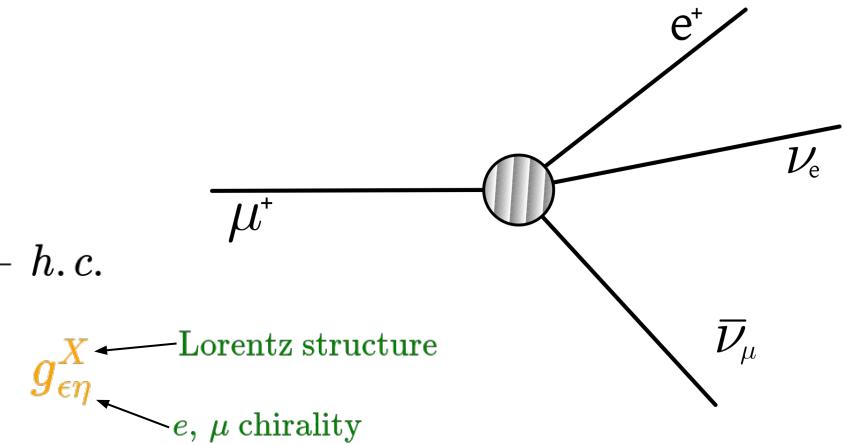


from P. S. Barbeau, et al. Ann. Rev. Nucl. Part. Sci. 73, 41–68 (2023)

Muon decay parameters

$$\mathcal{L} \supset -\frac{4G_F}{\sqrt{2}} \sum_{X,\eta,\epsilon} g_{\epsilon\eta}^X (\bar{e}_\epsilon \Gamma^X (\nu_e)_\rho) ((\bar{\nu}_\mu)_\gamma \Gamma_X \mu_\eta) + h.c.$$

$$\begin{aligned}\Gamma_X &= \{1, \gamma_\lambda, \sigma_{\lambda\omega}/\sqrt{2}\} \\ X &= \{S, V, T\}\end{aligned}$$

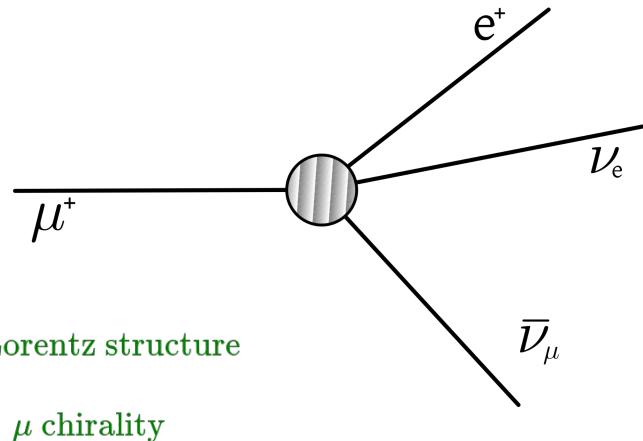


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- $\Gamma_\mu \xrightarrow[\text{measurement}]{} G_F \implies \sum_{\epsilon,\eta} \left[\frac{1}{4} |g_{\epsilon\eta}^S|^2 + |g_{\epsilon\eta}^V|^2 + 3|g_{\epsilon\eta}^T|^2 \right] = 1$
- SM limit $\xrightarrow{\text{red arrow}} g_{LL}^V = 1, \text{ rest} = 0$



(PDG collaboration), Review of particle physics, vol. 110, pp. 826–828, Phys. Rev. D (2024)

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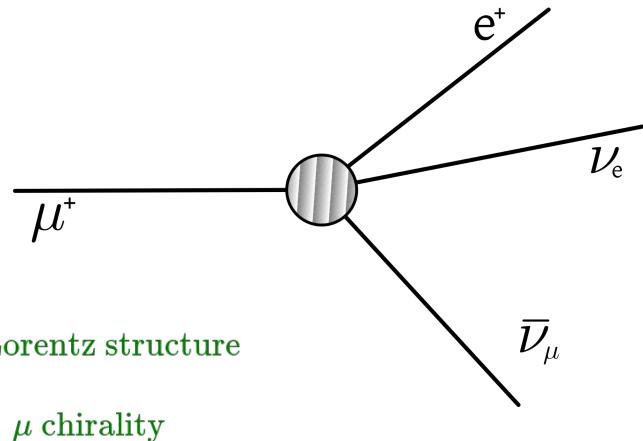
Usually measured from
the e^+/e^- detection

Flat direction

$$(g_{LL}^V)^2 + \frac{1}{4}(g_{LL}^S)^2$$

(PDG collaboration), Review of particle physics, vol. 110, pp. 826–828, Phys. Rev. D (2024)

B. Balke et al., Phys. Rev. D 37, 587 (1988)
R. P. MacDonald et al. (TWIST), Phys. Rev. D 78, 032010 (2008)
A. Hillairet et al. (TWIST), Phys. Rev. D 85, 092013 (2012)



Muon decay parameters

- For left-handed neutrinos:

$$\frac{d\Gamma_{\nu_L}}{dE_\nu} = \frac{24\Gamma_\mu}{m_\mu} P_{\nu_L} \left[y^2 (1-y) + \frac{8}{9} w_{\nu_L} y^2 \left(y - \frac{3}{4} \right) \right],$$

- For left-handed anti-neutrinos

$$\frac{d\Gamma_{\bar{\nu}_L}}{dE_{\bar{\nu}}} = \frac{24\Gamma_\mu}{m_\mu} P_{\bar{\nu}_L} \left[y^2 \left(\frac{1}{2} - \frac{y}{3} \right) + \frac{8}{9} w_{\bar{\nu}_L} y^2 \left(\frac{3}{4} - y \right) \right],$$

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$$P_{\nu_L} = |g_{LL}^V|^2 + |g_{LR}^V|^2 + \frac{1}{4} |g_{RR}^S|^2 + \frac{1}{4} |g_{RL}^S|^2 + 3 |g_{RL}^T|^2$$

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$$(g_{LL}^V)^2 + \frac{1}{4}(g_{LL}^S)^2 ???$$

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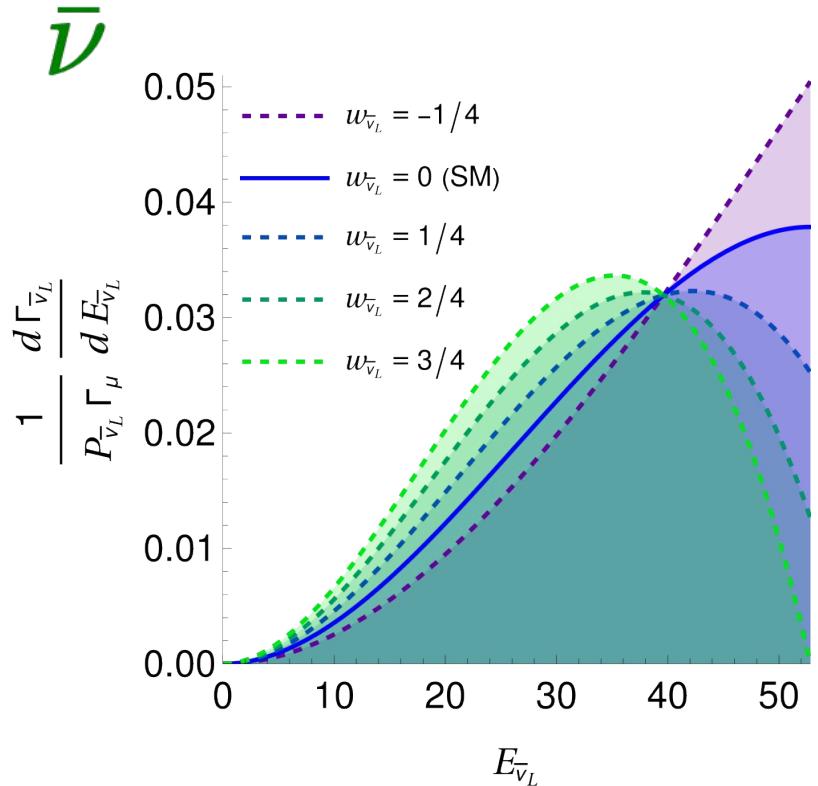
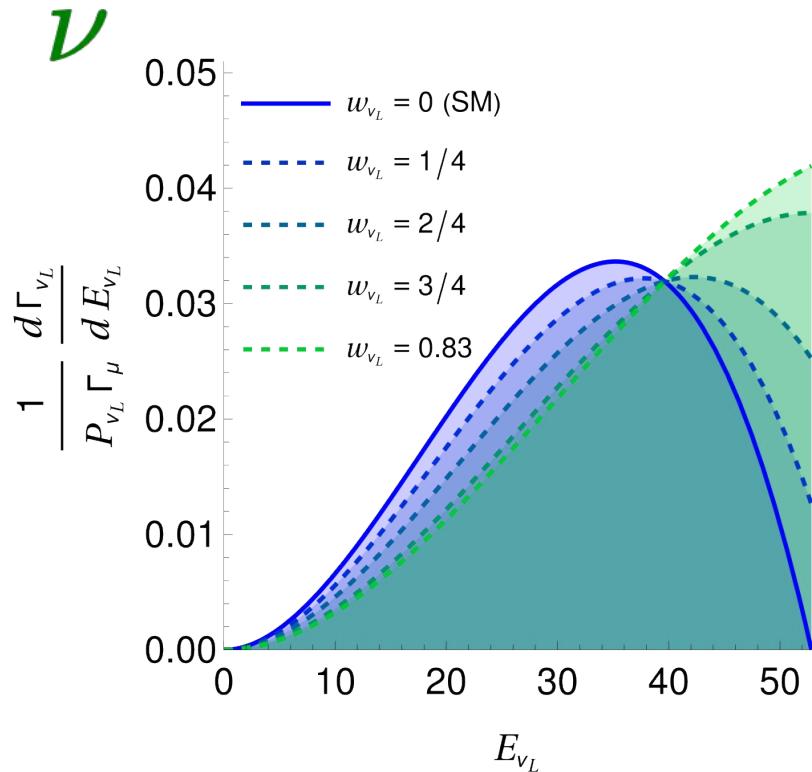
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In our work, we show the sensitivity of COHERENT-like experiments to these parameters and propose their first phenomenological extraction.

The (anti)neutrino flux



Rates

The rate per recoil energy T , per time t is:

$$\frac{dN}{dt dT} = g_\pi(t) \frac{dN^{\text{prompt}}}{dT} + g_\mu(t) \frac{dN^{\text{delayed}}}{dT},$$

where

$$\frac{dN^{\text{prompt}}}{dT} = N_T \int dE_\nu \frac{d\phi_{\nu_\mu}}{dE_\nu} \frac{d\sigma_{\nu_\mu}}{dT}, \quad \frac{dN^{\text{delayed}}}{dT} = N_T \int dE_\nu \left(\frac{d\phi_{\nu_e}}{dE_\nu} \frac{d\sigma_{\nu_e}}{dT} + \frac{d\phi_{\bar{\nu}_\mu}}{dE_\nu} \frac{d\sigma_{\bar{\nu}_\mu}}{dT} \right).$$

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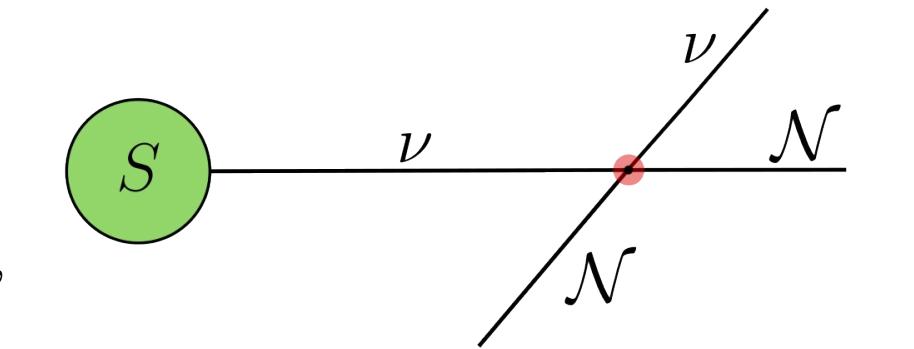
$$\frac{dN}{dt dT} = g_\pi(t) \frac{dN^{\text{prompt}}}{dT} + g_\mu(t) \frac{dN^{\text{delayed}}}{dT},$$

time dependence of the source

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ϕ_ν — neutrino flux



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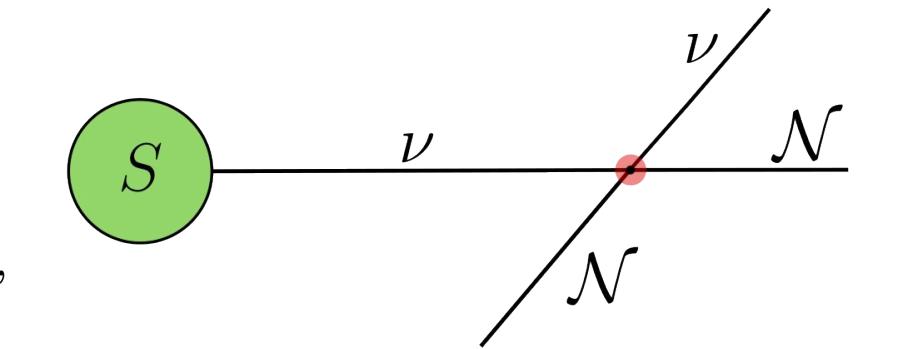
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σ — CE ν NS cross-section

This setup has been widely used in many different CE ν NS analysis (electroweak precision proves, nuclear studies and NP searchers).

Liao and Marfatia, Phys. Lett. B 775 (2017) 54
 C. Giunti, Phys. Rev. D 101 (2020) 035039
 Papoulias and Kosmas, Phys. Rev. D 97 (2018) 033003
 Aristizabal-Sierra, et. al., Phys. Rev. D 98 (2018) 075018

Skiba and Xia, JHEP 10 (2022) 102
 De Romeri, et. al., JHEP 04 (2023) 035
 Bresó-Pla, et.al., JHEP 05 (2023) 074

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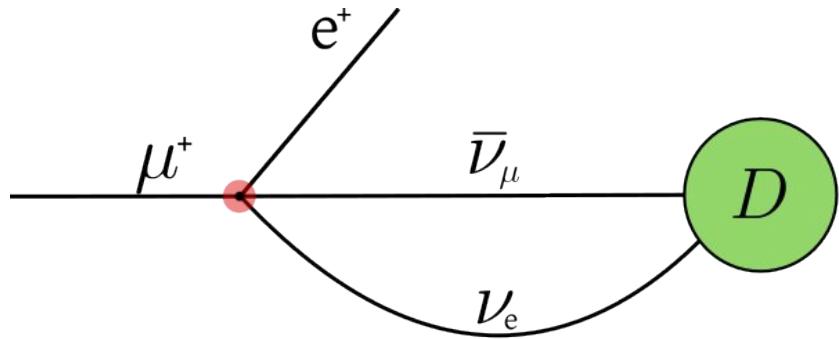
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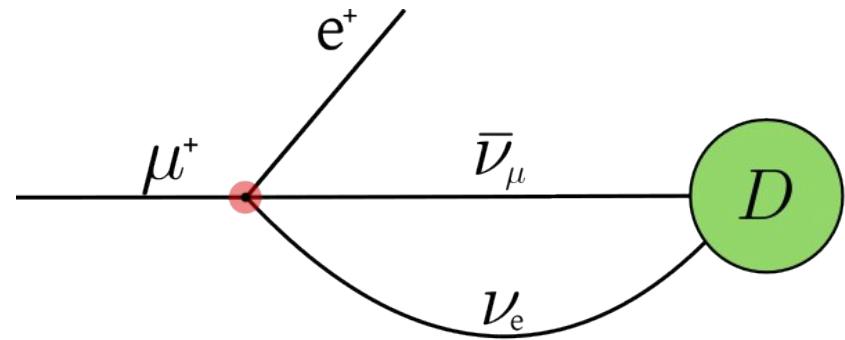


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$$x_{\nu_\mu} = 1,$$

$$x_{\bar{\nu}_\mu} = P_{\bar{\nu}_L} - \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L} + \frac{4}{3} P_{\nu_L} w_{\nu_L},$$

$$x_{\nu_e} = P_{\nu_L} - \frac{4}{3} P_{\nu_L} w_{\nu_L} + \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L}.$$

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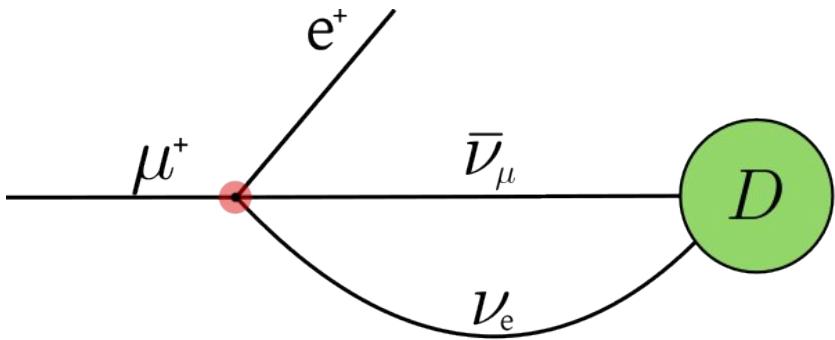
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prompt $\left\{ \begin{array}{l} x_{\nu_\mu} = 1, \\ x_{\bar{\nu}_\mu} = P_{\bar{\nu}_L} - \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L} + \frac{4}{3} P_{\nu_L} w_{\nu_L}, \\ x_{\nu_e} = P_{\nu_L} - \frac{4}{3} P_{\nu_L} w_{\nu_L} + \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L}. \end{array} \right.$

delayed



$$\frac{d\sigma^{\text{SM}}}{dT} = \frac{m_N G_F^2 \mathcal{F}^2(T)}{4\pi} Q_W^2 \left(1 - \frac{m_N T}{2E_\nu^2} - \frac{T}{E_\nu} \right)$$

$$\frac{d\phi_{\nu_\mu}^{\text{SM}}}{dE_\nu} = \frac{N_{\nu_\mu}}{4\pi L^2} \delta(E_\nu - E_{\nu,\pi}),$$

$$\frac{d\phi_{\nu_e}^{\text{SM}}}{dE_\nu} = \frac{N_{\nu_e}}{4\pi L^2} \frac{192 E_\nu^2}{m_\mu^3} \left(\frac{1}{2} - \frac{E_\nu}{m_\mu} \right),$$

$$\frac{d\phi_{\bar{\nu}_\mu}^{\text{SM}}}{dE_\nu} = \frac{N_{\bar{\nu}_\mu}}{4\pi L^2} \frac{64 E_\nu^2}{m_\mu^3} \left(\frac{3}{4} - \frac{E_\nu}{m_\mu} \right).$$

Muon decay

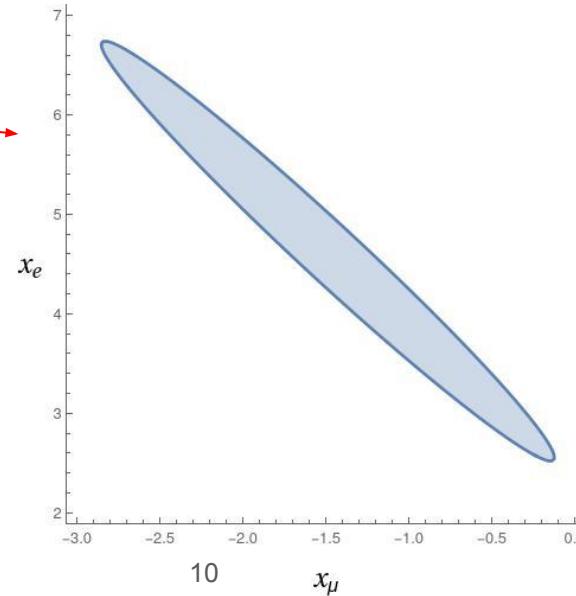
$$\frac{dN^{\text{pr./del.}}}{dT} = \sum_f N_T \int dE_\nu x_f \frac{d\phi_f^{\text{SM}}}{dE_\nu} \frac{d\sigma^{\text{SM}}}{dT} \quad \text{with} \quad \begin{aligned} x_{\nu_\mu} &= 1 , \\ x_{\bar{\nu}_\mu} &= P_{\bar{\nu}_L} - \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L} + \frac{4}{3} P_{\nu_L} w_{\nu_L} , \\ x_{\nu_e} &= P_{\nu_L} - \frac{4}{3} P_{\nu_L} w_{\nu_L} + \frac{4}{3} P_{\bar{\nu}_L} w_{\bar{\nu}_L} . \end{aligned}$$

$$\begin{aligned} x_{\bar{\nu}_\mu} &= -1.5(1.3) , \\ x_{\nu_e} &= 4.6(2.1) , \\ \rho &= -0.98 . \end{aligned}$$

Muon decay

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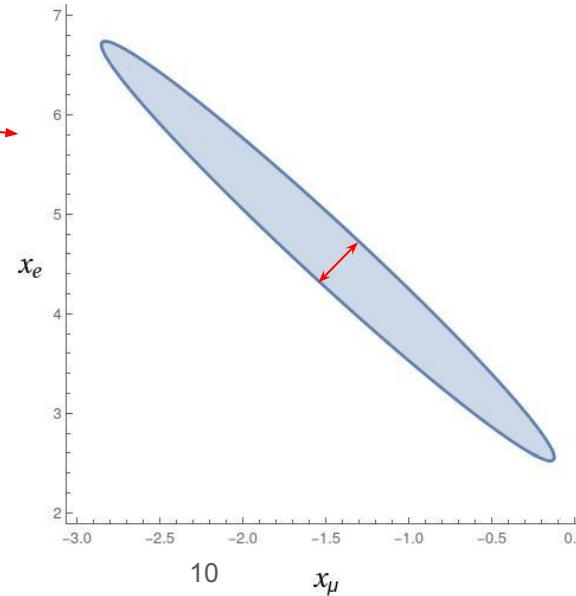
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Muon decay

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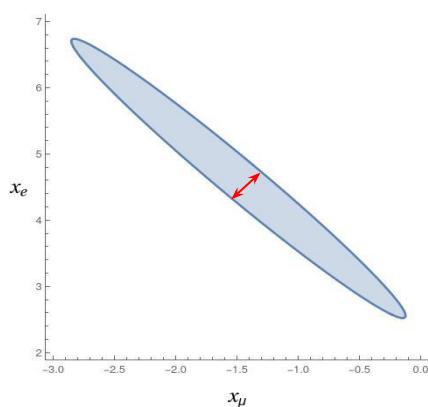
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Muon decay

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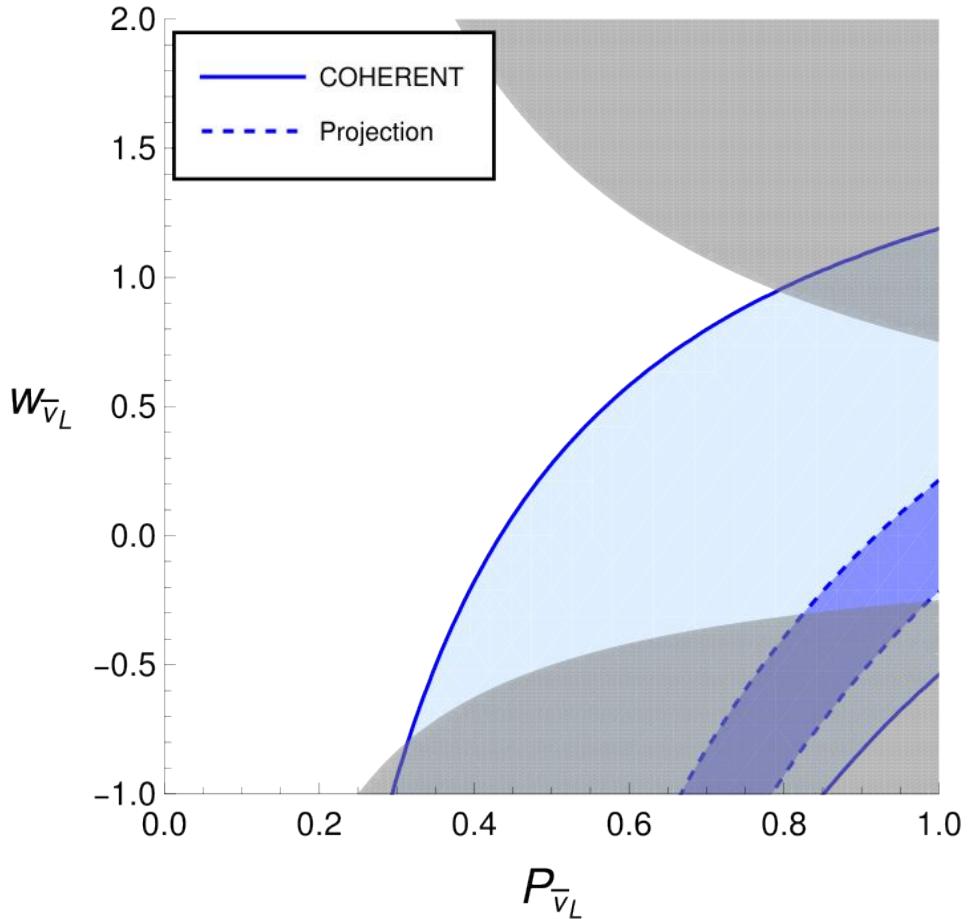
$$0.84 x_{\bar{\nu}_\mu} + 0.54 x_{\nu_e} = 1.25(21) ,$$

$$0.54 P_{\nu_L} + 0.84 P_{\bar{\nu}_L} + 0.40 (P_{\nu_L} w_{\nu_L} - P_{\bar{\nu}_L} w_{\bar{\nu}_L}) = 1.25 \pm 0.21 .$$

Muon decay

$$\begin{aligned} P_{\nu_L} &= 0.98 \left(\begin{array}{l} +02 \\ -37 \end{array} \right), & w_{\nu_L} &= 0.00 \left(\begin{array}{l} +23 \\ -00 \end{array} \right), \\ P_{\bar{\nu}_L} &= 0.90 \left(\begin{array}{l} +10 \\ -24 \end{array} \right), & w_{\bar{\nu}_L} &= 0.51 \left(\begin{array}{l} +24 \\ -51 \end{array} \right). \end{aligned}$$

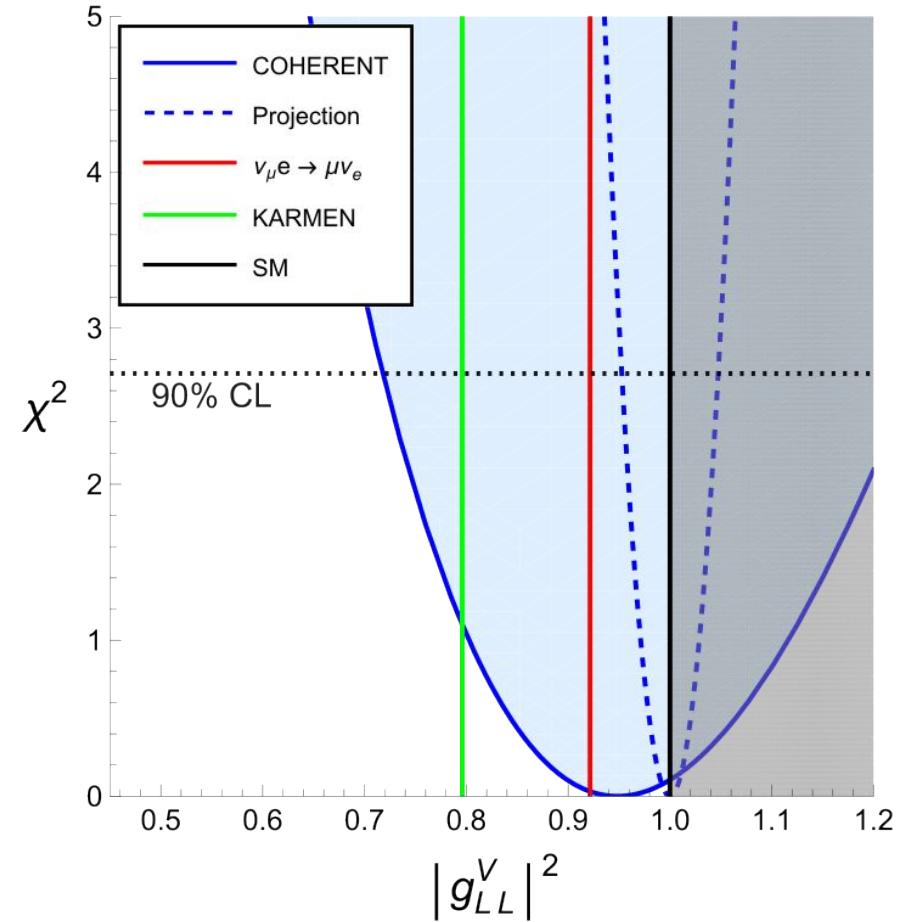
first constraint



Muon decay full χ^2

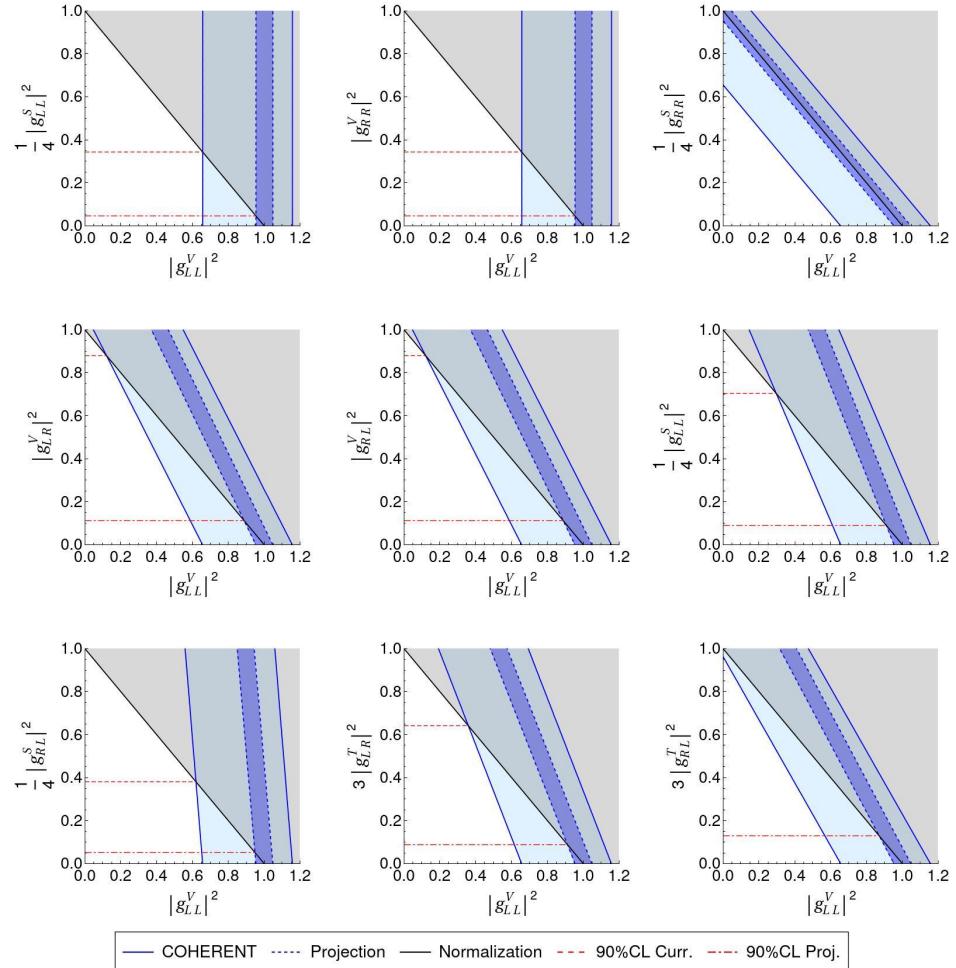
$$|g_{LL}^V|^2 = 0.95^{+0.05}_{-0.15}$$

independent of g_{LL}^S



Muon decay

Coefficient	Bound (90% C.L.)		
	Current	Projected	PDG ($\alpha = \mu, \beta = e$)
$ g_{LL}^V $	>0.848	>0.976	>0.960
$ g_{LR}^V $	0.79	0.34	0.023
$ g_{RL}^V $	0.79	0.34	0.105
$ g_{LR}^S $	1.4	0.57	0.050
$ g_{RL}^S $	1.6	0.67	0.420
$ g_{LR}^T $	0.47	0.21	0.015
$ g_{RL}^T $	0.42	0.17	0.105
$ g_{RR}^V (*)$	0.53	0.22	0.017
$ g_{LL}^S (*)$	1.07	0.44	0.550



A more general case (2505.01275)

- Flavour general interactions $\longrightarrow g_{\mu\epsilon}^X \longrightarrow [h_{\mu\epsilon}^X]_{\alpha\beta}$ ν flavour
 - NP in detections and in π decay \longrightarrow

$(\bar{u}\gamma^\mu d)(\bar{l}_\alpha\gamma_\mu P_H\nu_\beta)$	$(\bar{q}\gamma^\mu q)(\bar{\nu}_\alpha\gamma_\mu P_L\nu_\beta)$
$(\bar{u}\gamma^\mu\gamma^5 d)(\bar{l}_\alpha\gamma_\mu P_H\nu_\beta)$	$(\bar{q}\gamma^\mu\gamma^5 q)(\bar{\nu}_\alpha\gamma_\mu P_L\nu_\beta)$
$(\bar{u}d)(\bar{l}_\alpha P_H\nu_\beta)$	$(\bar{q}\gamma^\mu q)(\bar{\nu}_\alpha\gamma_\mu P_R\nu_\beta)$
$(\bar{u}i\gamma^5 d)(\bar{l}_\alpha P_H\nu_\beta)$	$(\bar{q}\gamma^\mu\gamma^5 q)(\bar{\nu}_\alpha\gamma_\mu P_R\nu_\beta)$
$(\bar{u}\sigma^{\mu\nu} P_H d)(\bar{l}_\alpha\sigma_{\mu\nu} P_H\nu_\beta)$	$(\bar{q}q)(\bar{\nu}_\alpha P_R\nu_\beta) + h.c.$
	$(\bar{q}i\gamma^5 q)(\bar{\nu}_\alpha P_R\nu_\beta) + h.c.$
	$(\bar{q}\sigma^{\mu\nu} q)(\bar{\nu}_\alpha\sigma_{\mu\nu} P_R\nu_\beta) + h.c.$
	$(\bar{\nu}_\alpha\sigma^{\mu\nu} P_R\nu_\beta)F_{\mu\nu} + h.c.$
 - A consistent EFT framework
- $\frac{1}{N_T} \frac{dN_\alpha^S}{dt dE_\nu dT} = \frac{N_S(t)}{32\pi L^2 m_S m_N E_\nu} \sum_{j,k,l} e^{-i \frac{L\Delta m_{kl}^2}{2E_\nu}} \int d\Pi_{P'} \mathcal{M}_{\alpha k}^P \bar{\mathcal{M}}_{\alpha l}^P \int d\Pi_{D'} \mathcal{M}_{jk}^D \bar{\mathcal{M}}_{jl}^D$

A. Falkowski et al., JHEP 11 (2020) 048 and V. Bresó-Pla et al., JHEP 05 (2023) 074

A more general case (2505.01275)

- Flavour general interactions
- NP in detections and in π decay
- A consistent EFT framework



$$\frac{dN^{\text{prompt/delayed}}}{dT} = N_T \int dE_\nu \frac{d\phi_{\nu_f}^{\text{SM}}}{dE_\nu} \frac{d\tilde{\sigma}_{\nu_f}}{dT}$$

$$\begin{aligned} \frac{d\tilde{\sigma}_{\nu_f}}{dT} = & \frac{m_N \mathcal{F}^2(q^2)}{8\pi v^4} \left\{ (\tilde{Q}_V^f)^2 \left(1 - \frac{m_N T}{2E_\nu^2} - \frac{T}{E_\nu} \right) \right. \\ & + (\tilde{Q}_S^f)^2 \frac{m_N T}{2E_\nu^2} + (\tilde{Q}_F^f)^2 \frac{v^2}{2m_N T} \left(1 - \frac{T}{E_\nu} \right) \\ & \left. + (\tilde{Q}_{SF}^f)^2 \frac{v}{2E_\nu} \left(1 - \frac{T}{2E_\nu} \right) \right\} \end{aligned}$$

$$\begin{aligned} \frac{d\phi_{\nu_\mu}^{\text{SM}}}{dE_\nu} &= \frac{N_{\nu_\mu}}{4\pi L^2} \delta(E_\nu - E_{\nu,\pi}) , \\ \frac{d\phi_{\nu_e}^{\text{SM}}}{dE_\nu} &= \frac{N_{\nu_e}}{4\pi L^2} \frac{192E_\nu^2}{m_\mu^3} \left(\frac{1}{2} - \frac{E_\nu}{m_\mu} \right) , \\ \frac{d\phi_{\bar{\nu}_\mu}^{\text{SM}}}{dE_\nu} &= \frac{N_{\bar{\nu}_\mu}}{4\pi L^2} \frac{64E_\nu^2}{m_\mu^3} \left(\frac{3}{4} - \frac{E_\nu}{m_\mu} \right) . \end{aligned}$$

A more general case

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Coefficient	Bound (90% C.L.)		Coefficient	Bound (90% C.L.)		
	Current	Projected		Current	Projected	
(I)				(II)		
$ \epsilon_{ee}^{uu} $	0.078	0.035	$ \epsilon_{e\mu}^{uu} $	0.12	0.016	
$ \epsilon_{\mu\mu}^{uu} $	0.049	0.016	$ \epsilon_{e\tau}^{uu} $	0.17	0.030	
$ \epsilon_{ee}^{dd} $	0.071	0.033	$ \epsilon_{\mu\tau}^{uu} $	0.15	0.019	
$ \epsilon_{\mu\mu}^{dd} $	0.043	0.015	$ \epsilon_{e\mu}^{dd} $	0.11	0.015	
(III)				$ \epsilon_{e\tau}^{dd} $	0.15	0.028
$ [\tilde{\epsilon}_S^{uu}]_{e\alpha} $	1.6×10^{-2}	9.1×10^{-3}	$ \epsilon_{\mu\tau}^{dd} $	0.13	0.018	
$ [\tilde{\epsilon}_S^{uu}]_{\mu\alpha} $	1.4×10^{-2}	5.8×10^{-3}	(IV)			
$ [\tilde{\epsilon}_S^{dd}]_{e\alpha} $	1.5×10^{-2}	8.9×10^{-3}	$ [\tilde{\epsilon}_L^{ud} - \tilde{\epsilon}_R^{ud}]_{\mu\alpha} $	0.51	0.30	
$ [\tilde{\epsilon}_S^{dd}]_{\mu\alpha} $	1.3×10^{-2}	5.7×10^{-3}	$ [\tilde{\epsilon}_P^{ud}]_{\mu\alpha} $	0.019	0.011	

Summary

- We have performed the first extraction of the anti-neutrino emission probability and the corresponding spectrum shape parameter.
- EFTs (Low Energy in our case) provide a suitable frame to study NP phenomenology in a model independent way.
- Neutrino experiments are required inputs in the EW precision fits.
- COHERENT-like experiments are entering the precision era and will become relevant precision EW inputs in the near future.

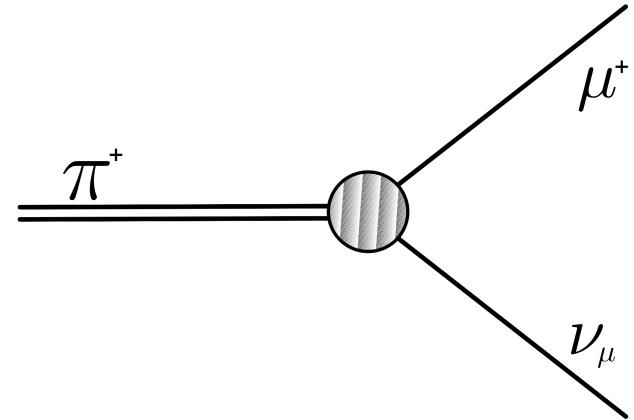
Summary

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Thank you!

Pion decay (pseudoscalar coupling)

$$\Delta\mathcal{L} = \sqrt{2} G_F V_{ud} \left\{ \epsilon_P (\bar{u} \gamma^5 d) (\bar{\mu} P_L \nu_\mu) + \tilde{\epsilon}_P (\bar{u} \gamma^5 d) (\bar{\mu} P_R \nu_\mu) \right\} + h.c.,$$



$$\frac{dN^{\text{pr./del.}}}{dT} = \sum_f N_T \int dE_\nu x_f \frac{d\phi_f^{\text{SM}}}{dE_\nu} \frac{d\sigma^{\text{SM}}}{dT}$$

$$\begin{aligned} \text{prompt} & \left\{ x_{\nu_\mu} \approx 1 - |\tilde{\epsilon}_P|^2 \frac{m_{\pi^\pm}^4}{m_\mu^2(m_u + m_d)^2}, \right. \\ \text{delayed} & \left\{ x_{\bar{\nu}_\mu} = 1, \right. \\ & \quad \left. x_{\nu_e} = 1. \right. \end{aligned}$$

- COHERENT data (projection) $\rightarrow \tilde{\epsilon}_P = 0.000 \pm 0.012 \text{ (0.009)} \quad 2.3 \text{ (2.6) TeV}$