

Unitarity Triangle Angles Explained: a Predictive New Quark Mass Matrix Texture

Plus, un antipasto on the house:
A Brief Historical Perspective on TriBimaximal Mixing

Paul Harrison
University of Warwick

Talk at FLASY 2025
30th June 2025

arXiv: 2501.18508 with Bill Scott
(to be published in JHEP)

- Historical Perspective on TriBimaximal Mixing (TBM) for leptons (a story of serendipity and following the data)
 - ▶ ν oscillation landscape pre-1994
 - ▶ Trimaximal Mixing (TMX)
 - ▶ CHOOZ, SNO and TBM
 - ▶ Legacy

- For a longer version:

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- A Predictive New Quark Mass Matrix Texture

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- ▶ The new quark mass matrix texture
- ▶ Confronting the data
- ▶ Symmetries of the texture
- ▶ Discussion and conclusions

- Solar ν problem had been around since 1970s:
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- Atmospheric ν anomaly was known as a simple deficit of ν_μ c.f. ν_e : IMB (1985) and Kamiokande (1988) :
 - ▶ Was initially a simple deficit $\sim 50\%$, with no L/E dependence
 - ▶ Suggested **large mixing angles**

1994 Provided Step-change in Perspective

Solar Data: P. Darriulat's ICHEP Conference Summary talk:

Experiment	Data/SSM (BP) %	Data/SSM (TCL) %
GALLEX	$60 \pm 8 \pm 5$	$64 \pm 8 \pm 5$
SAGE	$52 \pm 8 \pm 5$	$56 \pm 9 \pm 5$
Kamiokande	$51 \pm 4 \pm 6$	$66 \pm 5 \pm 8$
Homestake	(Pro memoria) $29 \pm 3 \pm 9$	

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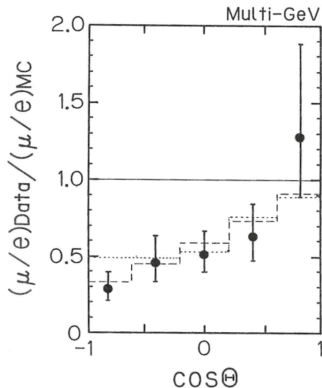
Atmospheric Neutrinos:

M. Nakahata, Glasgow ICHEP Conference 1994

Kamiokande Multi-GeV data

A clear zenith angle dependence!

Strongly favoured ν oscillations with large mixing angles



Meanwhile, also in 1994:

- Harrison and Scott propose TriMaximal mixing (TMX) for **quarks**(!):
 - ▶ “Generation permutation symmetry and the quark mixing matrix”, PLB 333 (1994) 471-475, hep-ph/9406351

$$U = \begin{bmatrix} \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{3}} \\ \frac{\omega}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{\bar{\omega}}{\sqrt{3}} \\ \frac{\bar{\omega}}{\sqrt{3}} & \frac{1}{\sqrt{3}} & \frac{\omega}{\sqrt{3}} \end{bmatrix} \Rightarrow (|U_{i\alpha}|^2) = \begin{bmatrix} \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \\ \frac{1}{3} & \frac{1}{3} & \frac{1}{3} \end{bmatrix}$$

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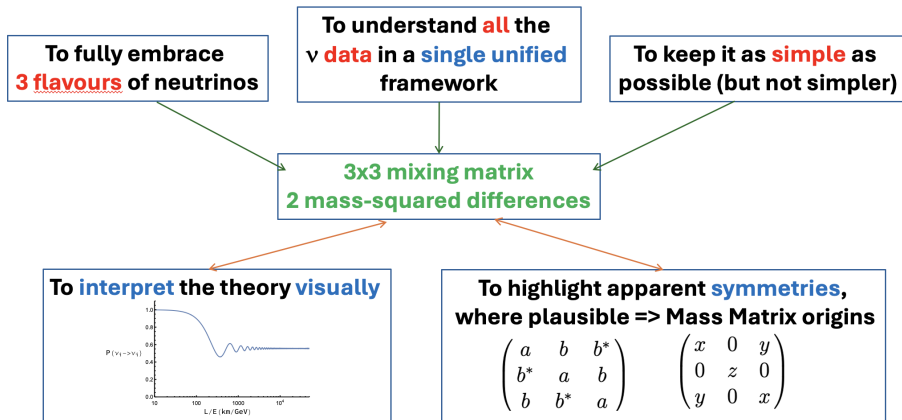
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- July '94, H&S realise that latest solar and atmospheric ν data are well-fitted by TMX!

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- Goals:

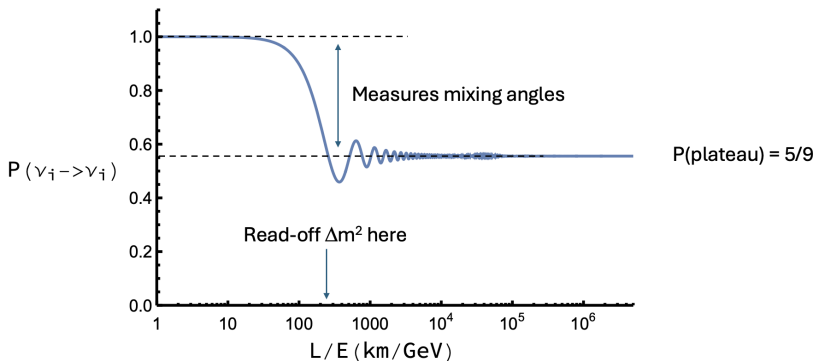


Threefold maximal lepton mixing (TMX) and the solar and atmospheric neutrino deficits (1995)

- PLB 349 (1995) 137 (“HPS1”)
- TMX $\Rightarrow$ universal survival (and appearance) probabilities (in vacuo) - eminently testable

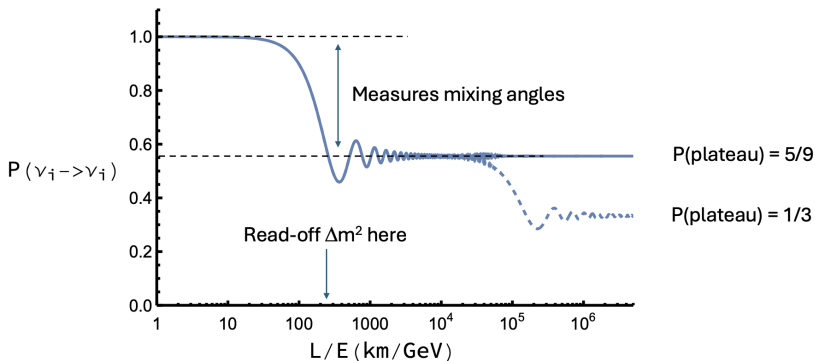
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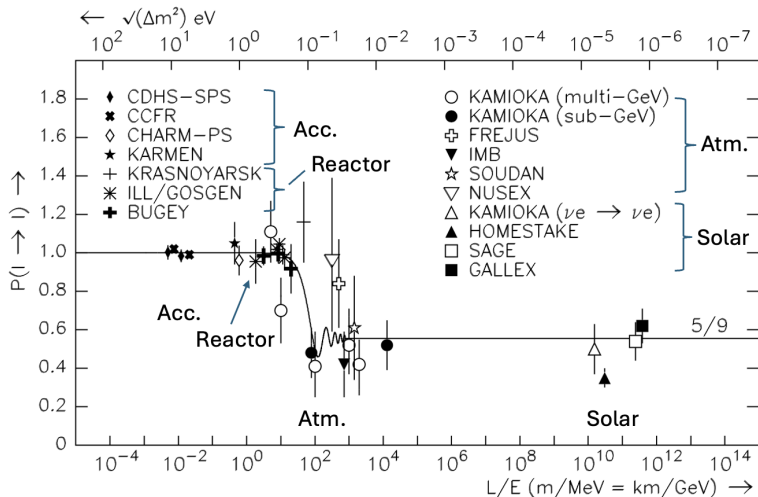


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Plotted with Data (taken from HPS1)



$$\chi^2/dof = 20.5/28;$$

$$\Delta m_{31}^2 = (7.8 \pm 1.2) \times 10^{-3} \text{ eV}^2$$

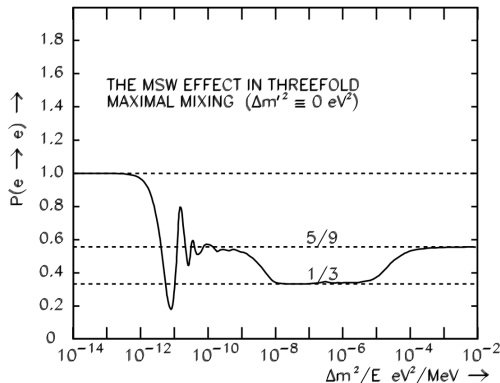
$$\Delta m_{21}^2 \text{ unresolved}$$

Included

A prediction:

The HPS “5/9-1/3-5/9” Bathtub

NB. required $\Delta m_{31}^2 \sim 10^{-8} - 10^{-5} \text{ eV}^2$



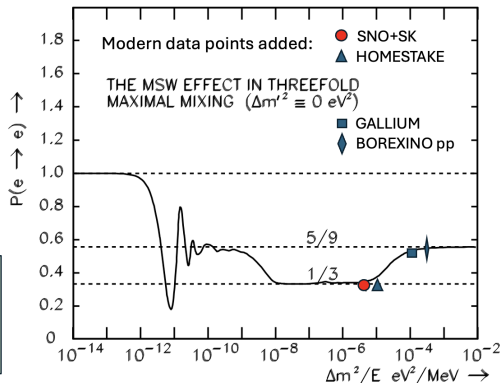
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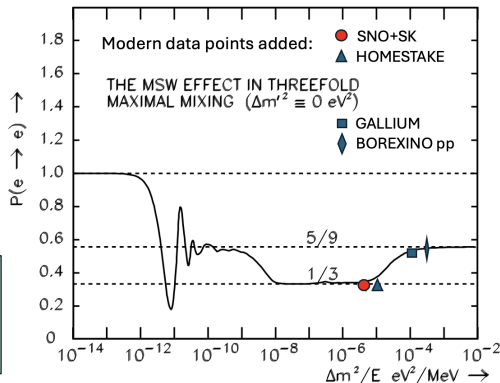
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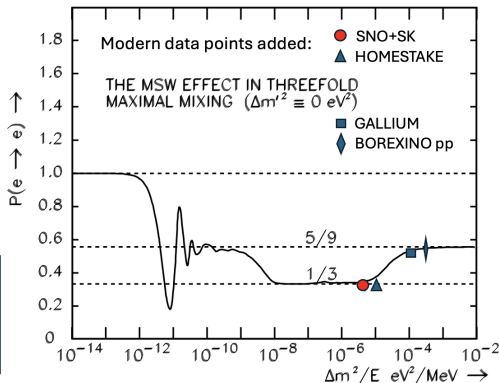
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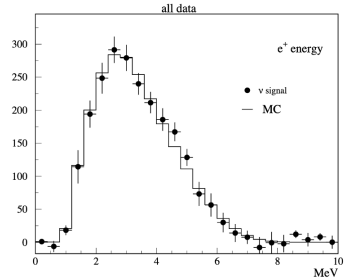


- HPS4 (PLB 458 (1999) 79) predicted for TMX “spectacular effects expected in long-baseline reactor and accelerator experiments”! viz. CHOOZ, PALO VERDE, MINOS etc.

BUT...**The CHOOZ
Reactor result 1999:**

$R = 1.01 \pm 0.03 \text{ (stat)} \pm 0.03 \text{ (sys)}$
($L \sim 1 \text{ km}$, $E \sim 3 \text{ MeV}$)
No Spectacular Effect!

Conclude: $|U_{e3}|^2$ small if $\Delta m^2 > 10^{-3} \text{ eV}^2$



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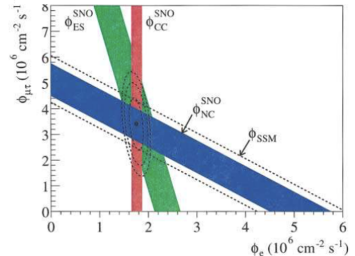
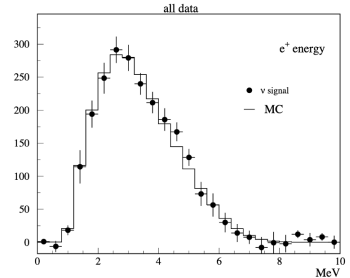
The SNO NC result 2002:

$\phi_{CC}/\phi_{NC} = 0.35 \pm 0.04 \sim 1/3!$

So...

Solar results energy-dependent

Trimaximal Mixing is Dead!



So, TriBiMaximal Mixing (TBM)

Long live
Tri(Bi)Maximal Mixing!

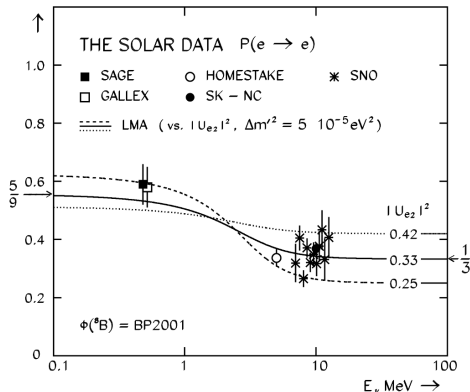
Tri-bimaximal mixing and the neutrino
oscillation data **PLB 530 (2002) 167**

“HPS5”

Return of the HPS “5/9-1/3-5/9” Bathtub!
(this time at the second threshold)

Conclude:

$$|U_{e2}|^2 \sim 1/3$$



So, TriBiMaximal Mixing (TBM)

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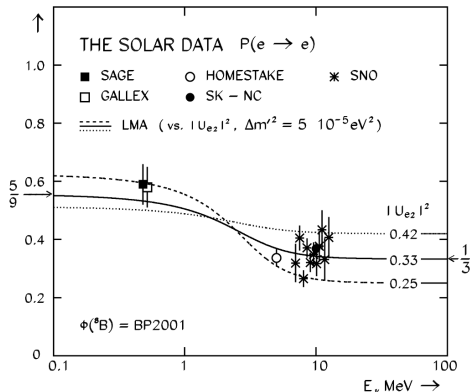
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Also...

**SUPERK Atmospheric neutrino
results pointing strongly to
twofold maximal $\nu_\mu - \nu_\tau$ mixing”**

Conclude:

$$|U_{\mu 3}|^2 \sim |U_{\tau 3}|^2 \sim 1/2$$



“The totality of the data clearly point to a particular form for the lepton mixing matrix”

$$U_{PMNS} = \begin{matrix} & \nu_1 & \nu_2 & \nu_3 \\ \begin{matrix} e \\ \mu \\ \tau \end{matrix} & \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} \end{matrix}$$

↔
Atmospheric
Scale
 Δm^2_{32}

Reactor
Mixing Angle
Atmospheric
Mixing Angle

“The totality of the data clearly point to a particular form for the lepton mixing matrix”

The diagram shows the PMNS lepton mixing matrix U_{PMNS} with its elements and associated physical parameters:

$$U_{PMNS} = \begin{pmatrix} e & \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ \mu & -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ \tau & -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

Annotations and physical parameters:

- Solar Scale Δm^2_{12}** : Indicated by a blue double-headed arrow between ν_1 and ν_2 .
- Solar Mixing Angle**: Indicated by a red arrow pointing to the $\frac{1}{\sqrt{3}}$ element in the first row, second column.
- Reactor Mixing Angle**: Indicated by a red arrow pointing to the 0 element in the first row, third column.
- Atmospheric Mixing Angle**: Indicated by a red arrow pointing to the $-\frac{1}{\sqrt{2}}$ element in the second row, fourth column.
- Atmospheric Scale Δm^2_{32}** : Indicated by a blue double-headed arrow between ν_2 and ν_3 .

“The totality of the data clearly point to a particular form for the lepton mixing matrix”

The diagram shows the PMNS matrix U_{PMNS} with rows labeled e , μ , and τ , and columns labeled ν_1 , ν_2 , and ν_3 . The matrix elements are:

$$U_{PMNS} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & -\frac{1}{\sqrt{2}} \\ -\sqrt{\frac{1}{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

Annotations with red arrows point to specific elements:

- Solar Scale Δm^2_{12}** : Points to the ν_1 and ν_2 columns.
- Solar Mixing Angle**: Points to the $\frac{1}{\sqrt{3}}$ element in the e row.
- Reactor Mixing Angle**: Points to the 0 element in the e row.
- Atmospheric Mixing Angle**: Points to the $-\frac{1}{\sqrt{2}}$ element in the μ row.
- Atmospheric Scale Δm^2_{32}** : Points to the ν_2 and ν_3 columns.

“Small non-zero values of U_{e3} and/or somewhat different values of $|U_{\mu 3}|^2$... are more-or-less equally acceptable experimentally.”

TBM was finally **excluded 10 years later** since (as anticipated) $U_{e3} = 0$ was incompatible with new results from the **Daya Bay** and **RENO** reactor experiments (2012) and later the **T2K** accelerator experiment (2013)

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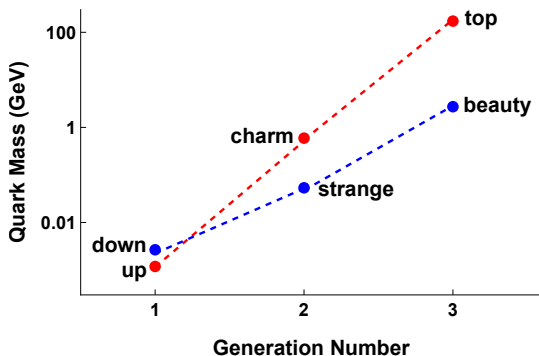
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Currently, $|U_{e3}|^2 \sim 0.02$. Thus **TBM** remains a **useful zeroth-order approximation** to U_{PMNS} , while allowing the exciting prospect that **leptonic CP violation may be accessible** in the future

Mystery of Quark Mass Spectra

- Quark masses show marked hierarchical structure:



- Is quasi-“geometric”:

$$\frac{m_c}{m_t} \simeq 0.0035$$

$$\frac{m_u}{m_c} \simeq 0.0020$$

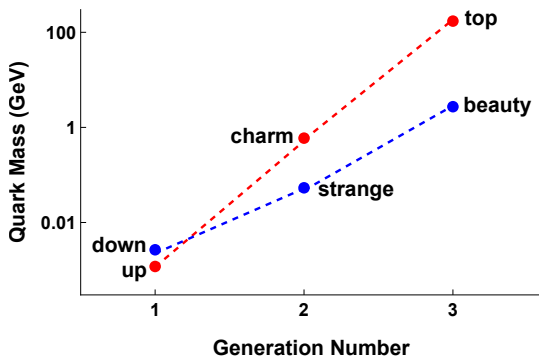
$$\frac{m_s}{m_b} \simeq 0.020$$

$$\frac{m_d}{m_s} \simeq 0.050$$

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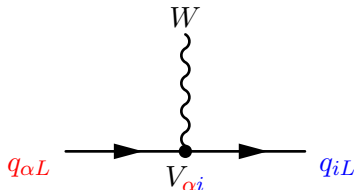
$$\frac{m_d}{m_s} \simeq 0.050$$

- Noted by very many authors
- Masses not predicted in the SM
- Hierarchy certainly not explained within SM
- BSM, Froggatt-Neilsen mechanism has had some success

Mystery of Quark Mixing Spectrum

- CKM quark mixing matrix:

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



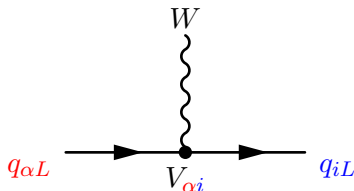
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$$\sim \begin{pmatrix} 1 & \lambda & A\lambda^3(\bar{\rho} - i\bar{\eta}) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4),$$

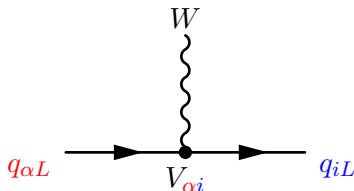
where $\lambda \equiv |V_{us}| \simeq 0.22$

A , $\bar{\rho}$ and $\bar{\eta} \lesssim \mathcal{O}(1)$



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$$\sim \begin{pmatrix} 1 & \lambda & A\lambda^3(\bar{\rho} - i\bar{\eta}) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4),$$

where $\lambda \equiv |V_{us}| \simeq 0.22$

A , $\bar{\rho}$ and $\bar{\eta} \lesssim \mathcal{O}(1)$

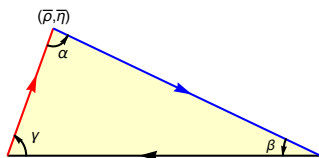
- Elements not predicted by the SM
- Strong hierarchy certainly not explained within SM
- But masses and mixings *both* arise in the Yukawa/Mass matrices

- Sides/Angles of UT are arbitrary in SM.
- But measured angles:

$$\alpha = (91.6 \pm 1.4)^\circ$$

$$\beta = (22.6 \pm 0.4)^\circ$$

$$\gamma = (65.7 \pm 1.3)^\circ$$



consistent with “special” values:

$$(\alpha, \beta, \gamma) \simeq \left(\frac{\pi}{2}, \frac{\pi}{8}, \frac{3\pi}{8}\right) \equiv (\alpha_0, \beta_0, \gamma_0).$$

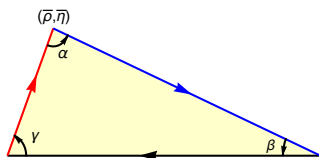
- Seems striking!

- Sides/Angles of UT are arbitrary in SM.
- But measured angles:

$$\alpha = (91.6 \pm 1.4)^\circ$$

$$\beta = (22.6 \pm 0.4)^\circ$$

$$\gamma = (65.7 \pm 1.3)^\circ$$



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- Seems striking!
- Coincidence or smoking gun?
- $\rightarrow$ Test as clue to what lies behind.

⇒ Build Special Angles into a Texture:

$$M_q^{HS} \equiv n_q \begin{pmatrix} c_q \lambda_q^4 & b \lambda_q^3 & 0 \\ b \lambda_q^{*3} & b \lambda_q^2 & A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix}, \quad \begin{array}{l} q = u, d, \\ \lambda_q \text{ complex} \\ \arg(\lambda_q) \text{ unobservable} \end{array}$$

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- Complex ratio is *fixed constant*:

$$\frac{\lambda_u}{\lambda_d} \equiv -i \tan \frac{\pi}{8}.$$

- ▶ $\arg \lambda_u / \lambda_d = -i$, is sole source of CP violation
- ▶ $|\lambda_u / \lambda_d| \simeq 0.41$ controls relative strength of “ u ” and “ d ” mass hierarchies
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- Describes 10 observables with 7 real parameters

- Diagonalise $\rightarrow$ masses:

$$D_q = U_q M_q^{HS} U_q^\dagger = m_3^q \begin{pmatrix} (c_q - b)\lambda_q^4 & 0 & 0 \\ 0 & b\lambda_q^2 & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad q = u, d,$$

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- Fits any m_u, m_d ✓ (no prediction here).

Leading-order Solution (Quark Mixing)

- Diagonalised by 2×2 (complex) rotations in 23 and 12 spaces.
- Small entries induced in the 13 elements of U_q :

$$U_q \simeq \begin{pmatrix} 1 & \pm \lambda_q & A_0 \lambda_q^3 \\ \mp \lambda_q^* & 1 & -A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix}, \quad q = u, d.$$

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$$\Rightarrow V_{CKM} = U_u U_d^\dagger \simeq \begin{pmatrix} 1 & \lambda_0 & A_0 \lambda_0^2 \lambda_u \\ -\lambda_0 & 1 & A_0 \lambda_0^2 \\ A_0 \lambda_0^2 \lambda_d^* & -A_0 \lambda_0^2 & 1 \end{pmatrix}$$

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- C.f. Wolfenstein form:

$$V_{CKM} = \begin{pmatrix} 1 & \lambda & A\lambda^3(\bar{\rho} - i\bar{\eta}) \\ -\lambda & 1 & A\lambda^2 \\ A\lambda^3(1 - \bar{\rho} - i\bar{\eta}) & -A\lambda^2 & 1 \end{pmatrix} \Rightarrow \begin{cases} \lambda \simeq \lambda_0 \checkmark \\ A \simeq A_0 \checkmark \\ (\bar{\rho} + i\bar{\eta}) \simeq \frac{\lambda_u^*}{\lambda_0} \end{cases}$$

The UT Angles

- Have deduced that:

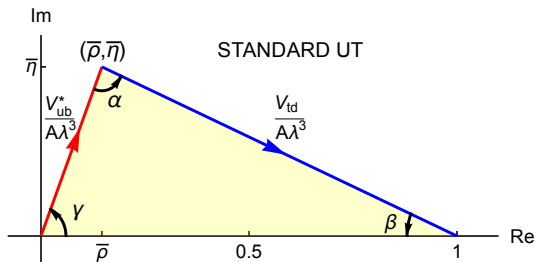
$$\lambda_u^* \simeq \lambda(\bar{\rho} + i\bar{\eta}) = \frac{V_{ub}^*}{A\lambda^2}$$

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$$\beta \simeq \arg \lambda_d$$

$$\text{and } \alpha \simeq \arg(-\frac{\lambda_u}{\lambda_d})$$



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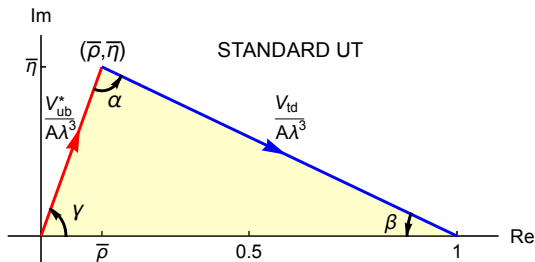
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- Recall, HS texture asserts $\frac{\lambda_u}{\lambda_d} = -i \tan \frac{\pi}{8}$.

► $\Rightarrow \alpha \simeq \frac{\pi}{2}$ ✓

► $\Rightarrow \tan \beta \simeq \left| \frac{\lambda_u}{\lambda_d} \right|$ (see Figure).

► $\Rightarrow \beta \simeq \frac{\pi}{8}$ ✓

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- Not necessarily!
- Because (exactly) these quantities run with renormalisation scale
- $\sim 14\%$ from weak to GUT scales: $A(\uparrow)$, $m_c/m_t(\uparrow)$ and $m_s/m_b(\downarrow)$.
- While λ , α , β , m_u/m_c and m_d/m_s are $\sim$ invariant.
- $\Rightarrow$ vary μ

Precision Fit to the Data

- Fit $\chi^2/\text{d.o.f} \simeq 1.01/2$

Observable	Input Renormali- sed to $\mu = 10^4$ TeV	Fitted Value at $\mu = 10^4$ TeV
$ m_u/m_c (\times 10^3)$	2.00 ± 0.05	2.00
$ m_d/m_s (\times 10^2)$	4.97 ± 0.06	4.97
$m_c/m_t (\times 10^3)$	3.46 ± 0.03	3.46
$m_s/m_b (\times 10^2)$	1.968 ± 0.008	1.968
λ	0.2250 ± 0.0007	0.2250
A	0.88 ± 0.02	0.88
$\bar{\rho}$	0.159 ± 0.009	0.152
$\bar{\eta}$	0.352 ± 0.007	0.348
UT Angles	—	Prediction from Fit
$\alpha(^{\circ})$	91.6 ± 1.4	91.30 ± 0.02
$\beta(^{\circ})$	22.6 ± 0.4	22.3 ± 0.1
$\gamma(^{\circ})$	65.7 ± 1.3	66.4 ± 0.1

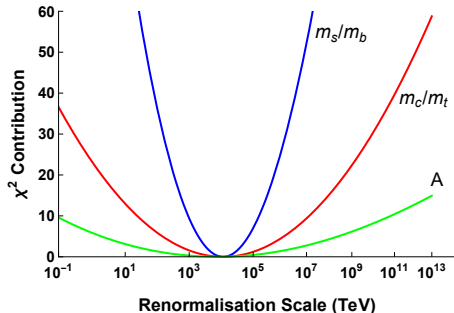
- Best fit renormalisation scale:
 $\mu \sim (0.3 \rightarrow 3) \times 10^4 \text{ TeV}.$
- Fitted values of the free parameters:
 - ▶ $\lambda_0 = 0.22646$
 - ▶ $A_0 = 0.854$
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- Three curves minimise at common scale $\sim 10^4 \text{ TeV}$
- Fitted values of observables in table represent predictions



The Leading Order UT (LO-UT)

- Define useful complex constants:

$$z_0 \equiv \lambda_u^* / \lambda_0 = i s_0 e^{-i\beta_0} = \rho_0 + i\eta_0,$$

$$\bar{z}_0 \equiv \lambda_d^* / \lambda_0 = c_0 e^{-i\beta_0} = 1 - z_0,$$

where

$$s_0 \equiv \sin \frac{\pi}{8}; \quad c_0 \equiv \cos \frac{\pi}{8}; \quad \eta_0 = s_0 c_0 = \frac{1}{2\sqrt{2}} \quad \text{and} \quad \rho_0 = \sin^2 \frac{\pi}{8}.$$

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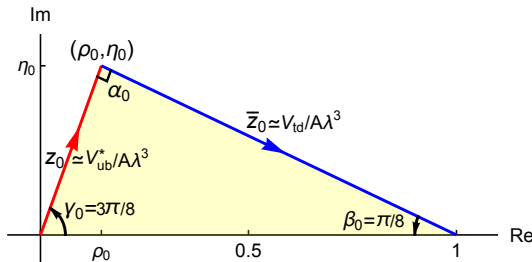
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- Use to construct LO-UT

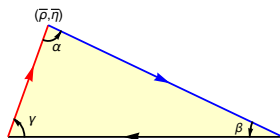


Symmetries of the MMs

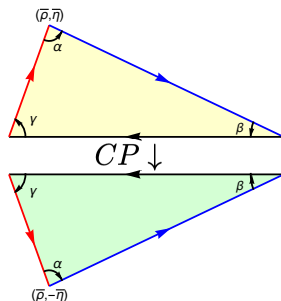
- Properties of the *paired system* (M_u , M_d), rather than of either in isolation

Recall: CP Transformation and Rephasing

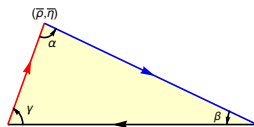
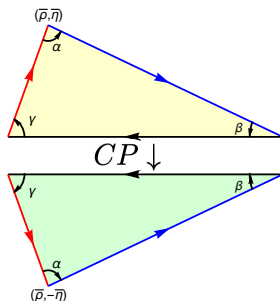
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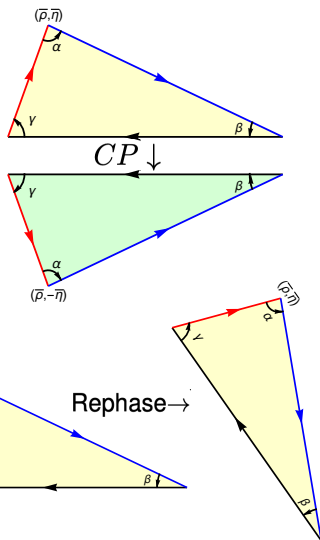


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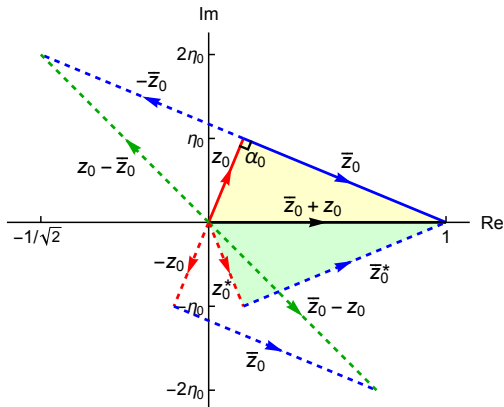
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- UT simply rotates in complex plane
- (Physical) shape and size invariant



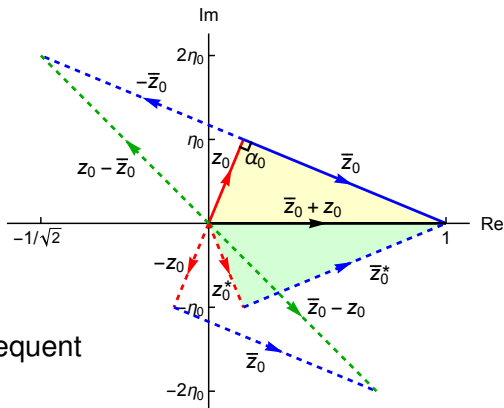
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- But iff $\alpha = \pm \frac{\pi}{2}$
- It is equivalent to a CP transformation
- Can be reversed by a subsequent actual CP transformation
- Symmetry is good to all orders

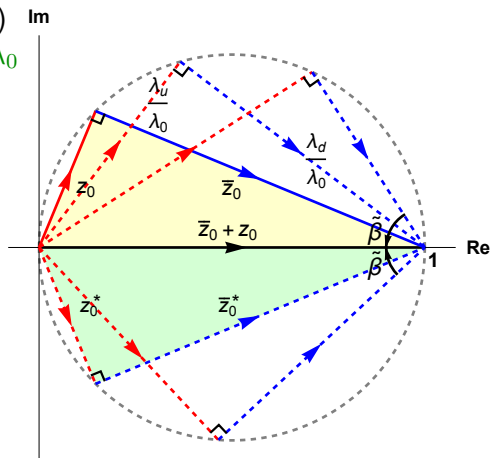


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- First consider $\beta_0 = \tilde{\beta} \neq \frac{\pi}{8}$ (fig→) keeping $\alpha = \frac{\pi}{2}$ and $\lambda_u + \lambda_d = \lambda_0$
- Clearly now

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$$\text{and } -\frac{\pi}{2} < \tilde{\beta} < \frac{\pi}{2}$$



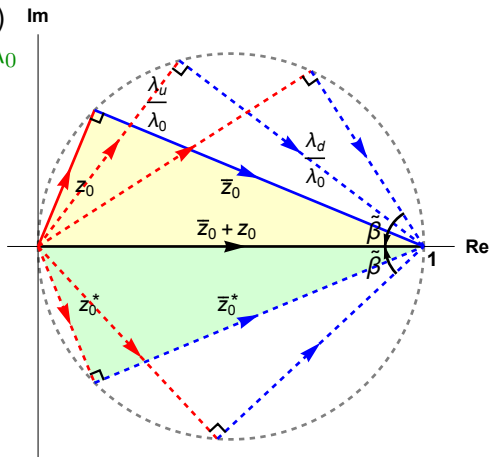
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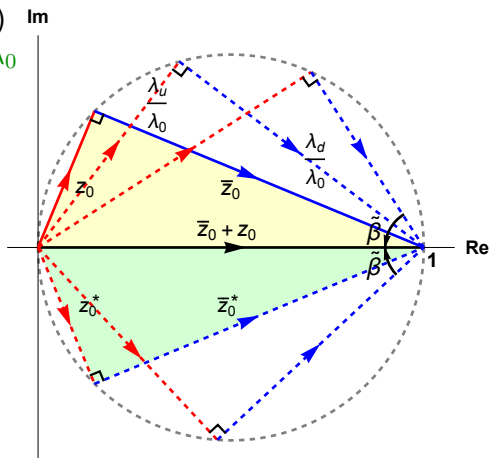
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- $\Rightarrow$ to fix $\beta_0 = \frac{\pi}{8}$ require symmetry under transformation $(*)$ followed by CP flip



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- Much to be done to find concrete BSM model(s) to implement/embed these symmetries.

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- Precise prediction of quark mass double ratio:

$$\frac{m_c}{m_t} \frac{m_b}{m_s} = \left| \frac{\lambda_u}{\lambda_d} \right|^2 = \tan^2 \frac{\pi}{8} (1 + \mathcal{O}(\lambda_0^2)) = 0.176 \pm 0.001$$

c.f. 0.177 ± 0.002 (exp)

- Proposed geometric-hierarchical MM texture
- Mass hierarchy “slopes” are related to UT sides
- Symmetries constrain UT angles $\rightarrow \alpha \simeq \frac{\pi}{2}$ and $\beta \simeq \frac{\pi}{8}$
- Origin of hierarchy not explained explicitly, but standard model-building methods can achieve that (e.g. F-N Mechanism)
- M_u and M_d exploit 7 pars to fit 10 observables with $\chi^2/\text{d.o.f} \simeq 1/2$
- Precise prediction of quark mass double ratio:

$$\frac{m_c}{m_t} \frac{m_b}{m_s} = \left| \frac{\lambda_u}{\lambda_d} \right|^2 = \tan^2 \frac{\pi}{8} (1 + \mathcal{O}(\lambda_0^2)) = 0.176 \pm 0.001$$

c.f. 0.177 ± 0.002 (exp)

- Precise predictions of UT angles:

$$\triangleright \alpha - \frac{\pi}{2} = (1.30 \pm 0.02)^\circ \text{ c.f. } (1.6 \pm 1.4)^\circ \text{ (exp)}$$

$$\triangleright \beta - \frac{\pi}{8} = (-0.2 \pm 0.1)^\circ \text{ c.f. } (0.1 \pm 0.4)^\circ \text{ (exp)}$$

$$\triangleright \gamma - \frac{3\pi}{8} = (-1.1 \pm 0.1)^\circ \text{ c.f. } (-1.8 \pm 1.3)^\circ \text{ (exp)}$$

Backup Slides

- Can re-write texture:

$$M_q^{HS} \equiv n_q \begin{pmatrix} c' \lambda_q^4 & b' \lambda_q^3 & 0 \\ b' \lambda_q^{*3} & b' \lambda_q^2 & A_0 \lambda_q^2 \\ 0 & A_0 \lambda_q^{*2} & 1 \end{pmatrix} \pm d \lambda_q^4 I$$

$b' \simeq b$. Still get good fit to data.

- First (leading) matrix solely responsible for quark mass differences and mixing parameters.
- Second (small) matrix is I_z -dependent “pedestal” on quark masses. Symmetric under a family- $SU(3)$ symmetry.
- All coefficients (λ_0 , A_0 , b' , c' , d) symmetric under isospin reflection operator $u \leftrightarrow d$.
- Symmetry broken (only) by λ_q , n_q and the sign of d .

- We give here the algebraic NLO solutions of the texture:

$$\lambda = \lambda_0 (1 + f_\lambda \lambda_0^2) + \mathcal{O}(\lambda_0^5)$$

$$A = A_0 \left\{ 1 + \left[\frac{1}{4}(3b - 2\rho_0) - 2f_\lambda \right] \lambda_0^2 \right\} + \mathcal{O}(\lambda_0^4)$$

$$\bar{\rho} = \rho_0 (1 + c_0 f_\rho \lambda_0^2) + \mathcal{O}(\lambda_0^4)$$

$$\bar{\eta} = \eta_0 \left\{ 1 + \left[s_0 f_\rho + \frac{1}{2}(1 - 5b) \right] \lambda_0^2 \right\} + \mathcal{O}(\lambda_0^4),$$

where $f_\lambda = \frac{1}{4}(3f_A + 4\eta_0 f_c - 5),$

$$f_A = \frac{1}{b} \left[A_0^2 + \frac{1}{2}(c_d + c_u) \right], \quad f_c = \frac{1}{b}(c_d - c_u)$$

and $f_\rho = s_0(1 + f_c) - \frac{1}{4s_0}[2f_A + 2f_c - 7b].$

- NLO corrections above, as fractions of LO terms are respectively: $-5.8 \times 10^{-3}, +2.6\%, +3.6\%$ and -1.8% (using fitted param values from table).

- For the quark mass ratios, we find:

$$\frac{m_1^q}{m_2^q} = -\lambda_q^2(1 - r_q) \left\{ 1 + \left[r_A \frac{(2-r_q)}{(1-r_q)} - 2 \right] \lambda_q^2 \right\} + \mathcal{O}(\lambda_q^6)$$

$$\frac{m_2^q}{m_3^q} = b\lambda_q^2 [1 + (1 - r_A)\lambda_q^2] + \mathcal{O}(\lambda_q^6),$$

where $r_q = \frac{c_q}{b}$ and $r_A = \frac{A_0^2}{b}$.

- NLO corrections to mass ratios m_c/m_t , m_s/m_b , m_u/m_c , m_d/m_s as fractions of LO terms are (resp.) -4.3×10^{-3} , -2.5% , $+4.3\%$, and $+4.5\%$ (using fitted param values from table).
- All results compatible with full numerical results reported in table.