

Toward ultra-low timing-jitters for beam-driven plasma accelerators

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On behalf of the SPARC_LAB collaboration



A particle accelerator is a machine that uses electromagnetic fields to propel charged particles to very high speeds and energies, and to contain them in well-defined beams

Wikipedia

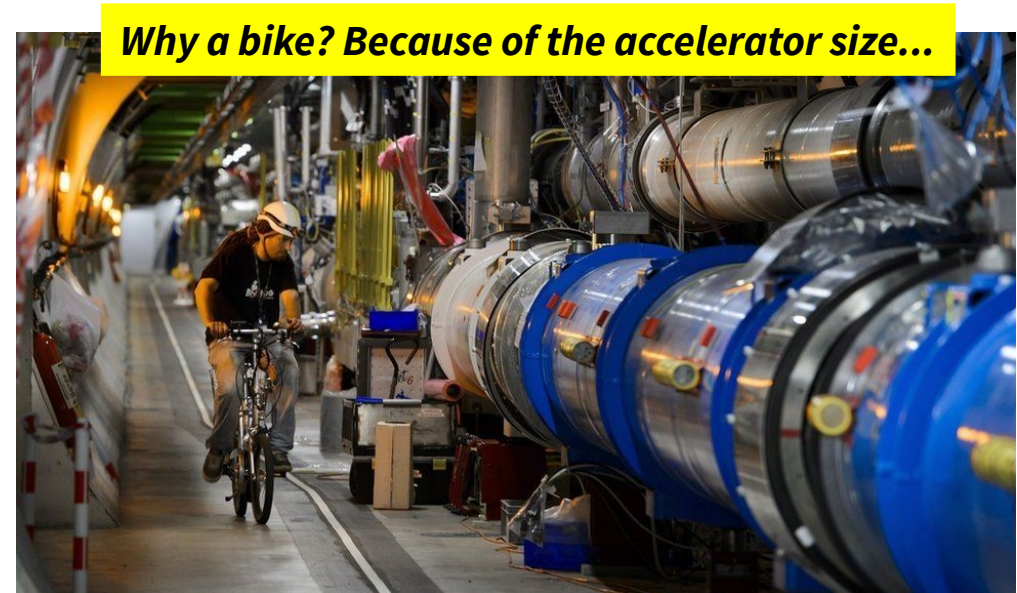
According to the De Broglie hypothesis, larger is the particle **energy** better is the **spatial resolution** at which matter can be investigated

Today use of particle accelerators

High-energy and nuclear physics

Sources of synchrotron and FEL radiation

Medical/industrial applications

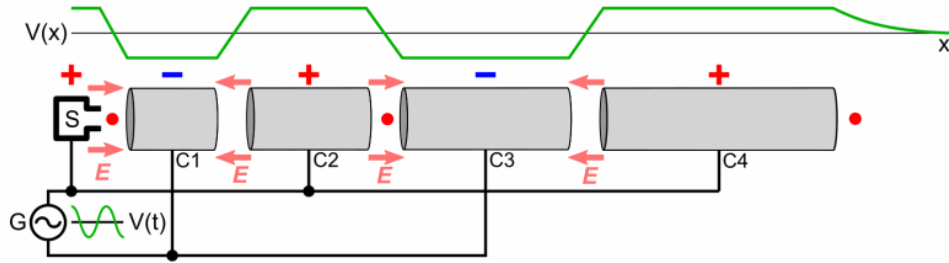


$$\lambda = \frac{h}{p}$$

Planck constant

Microscope resolution

Particle momentum (~energy)



RF technology uses high power microwaves in resonant cavities with metallic walls.

Typical RF frequencies are in GHz range.

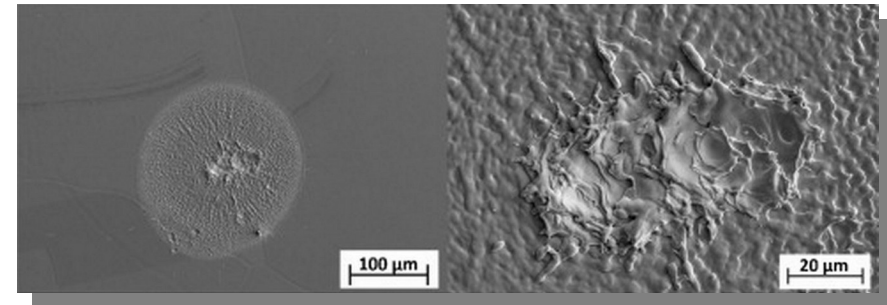
The cavities dimensions are of the order of the microwaves wavelength (1-60 cm)



The power stored in the cavity cannot grow to infinite

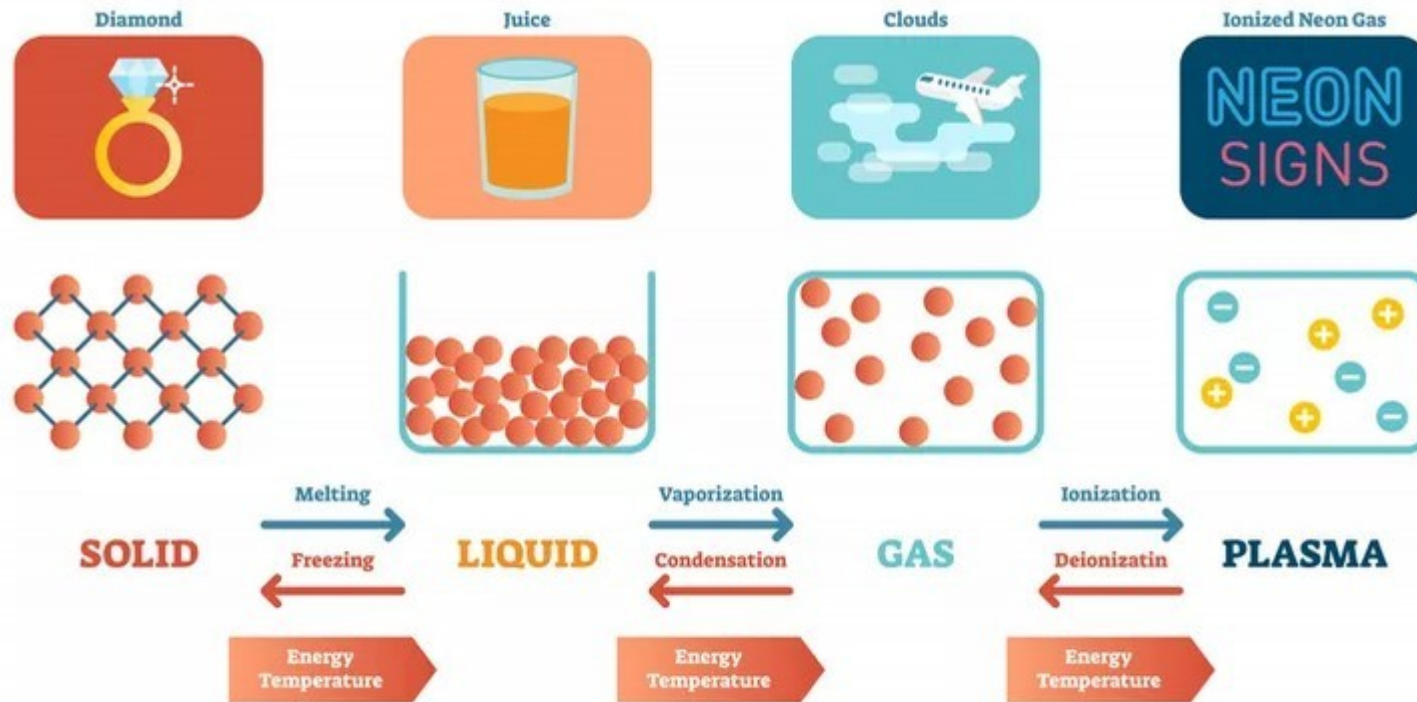
RF breakdown: imperfections on the cavities can trigger sparks and damage the structure

There is a maximum sustainable electric field



$$E_{max} \approx 100 \frac{MV}{m}$$

States of Matter



Plasma is the 4th state of matter and is made of free electrons and positively charged ions

It is typically made by heating a gas until its electrons have sufficient energy to escape from the (positive) nuclei

Being already ionized, the plasma cannot be “damaged” by any spark and can thus sustain huge fields

From Maxwell's equations, the electric field in a (positively) charged sphere with uniform density n_i at location \mathbf{r} is

$$\vec{E}(\mathbf{r}) = \frac{q_i n_i}{3 \epsilon_0} \mathbf{r}$$

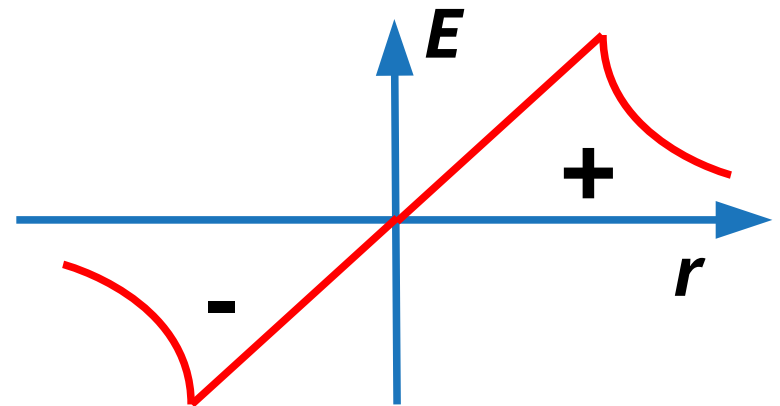
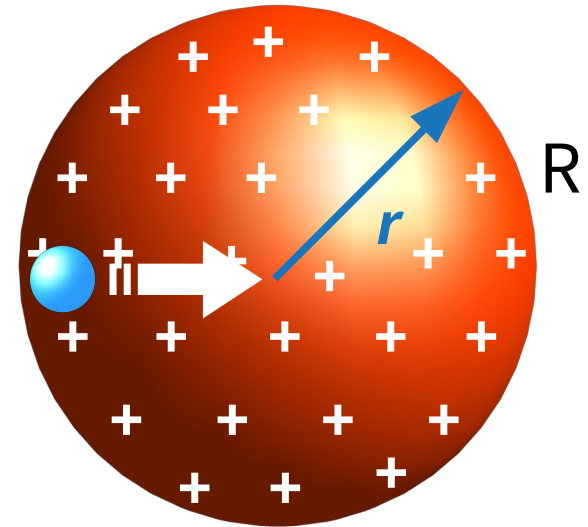
The field is **increasing** inside the sphere

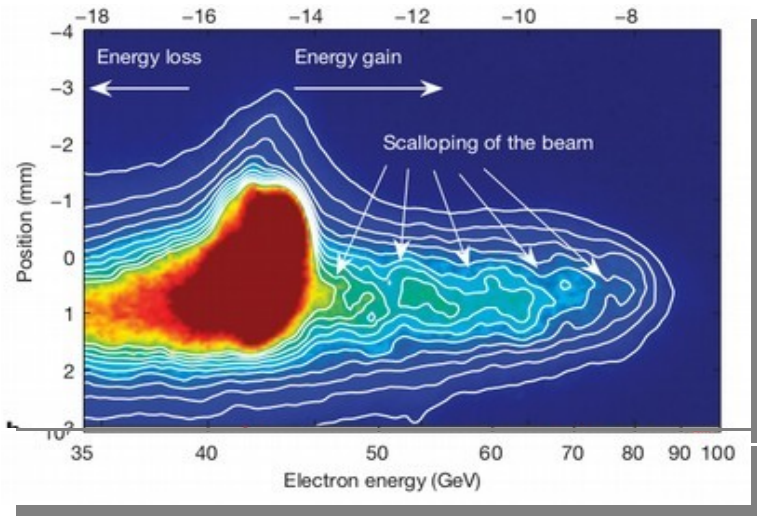
Let's put some numbers

$$n_i = 10^{16} \text{ cm}^{-3}$$

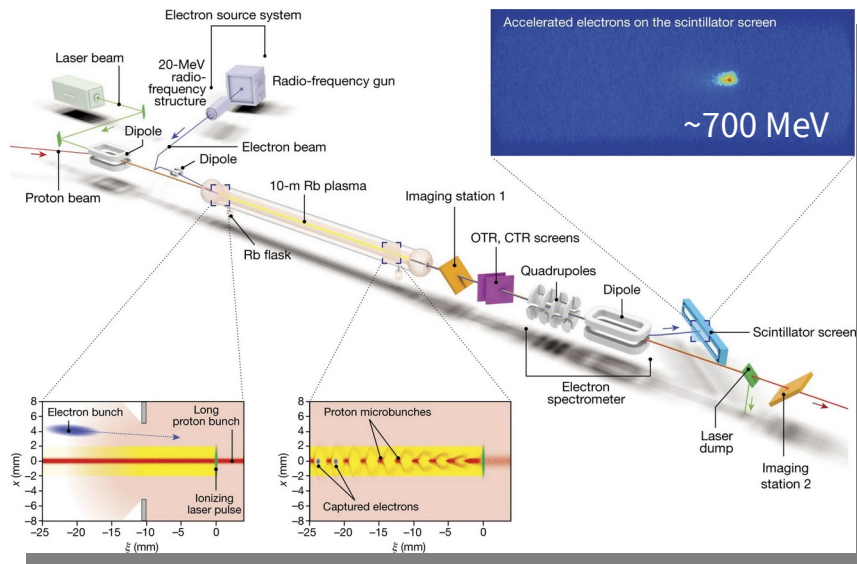
$$R = 0.5 \lambda_p = 150 \mu\text{m}$$

$$\Rightarrow E \approx 10 \frac{\text{GV}}{\text{m}}$$

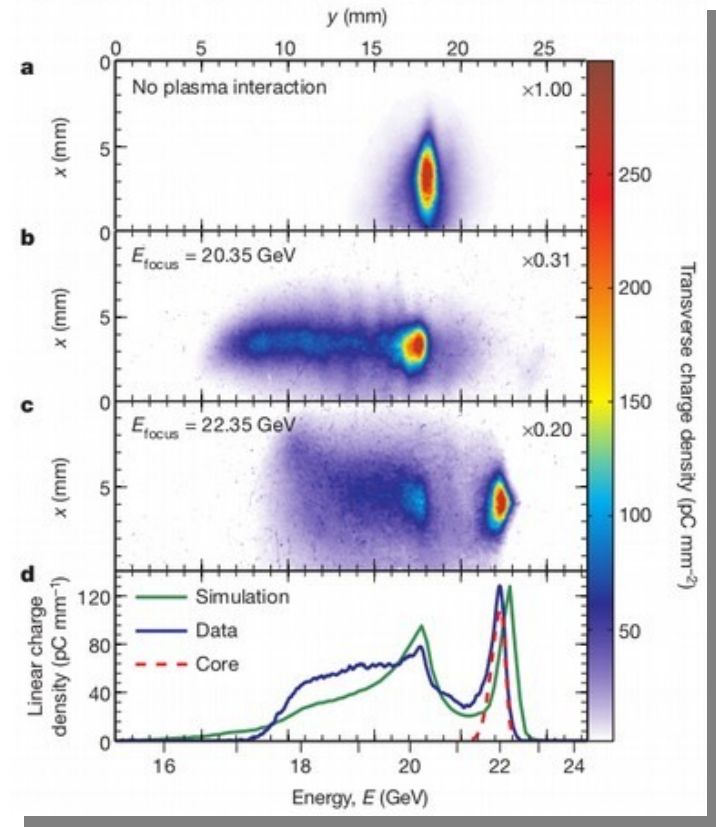




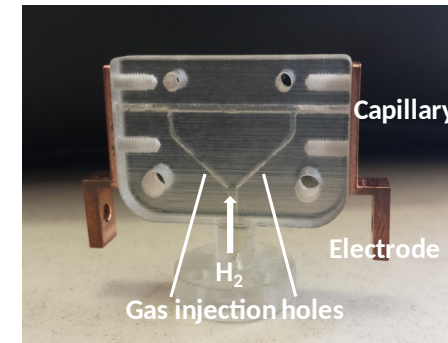
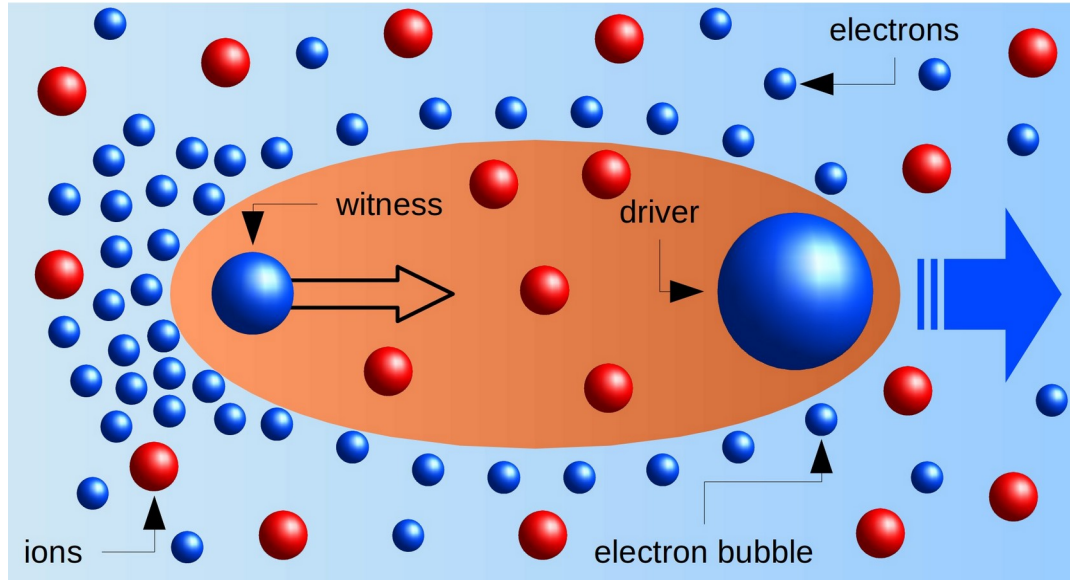
Blumenfeld, Ian, et al. "Energy doubling of 42 GeV electrons in a metre-scale plasma wakefield accelerator." Nature 445.7129 (2007): 741-744.



Adli, Erik, et al. "Acceleration of electrons in the plasma wakefield of a proton bunch." Nature 561.7723 (2018): 363-367.

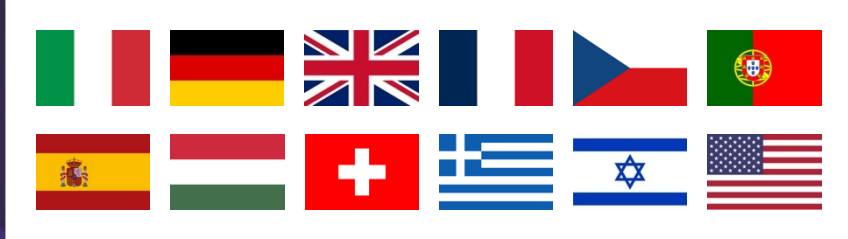


Litos, M., et al. "High-efficiency acceleration of an electron beam in a plasma wakefield accelerator." Nature 515.7525 (2014): 92-95.

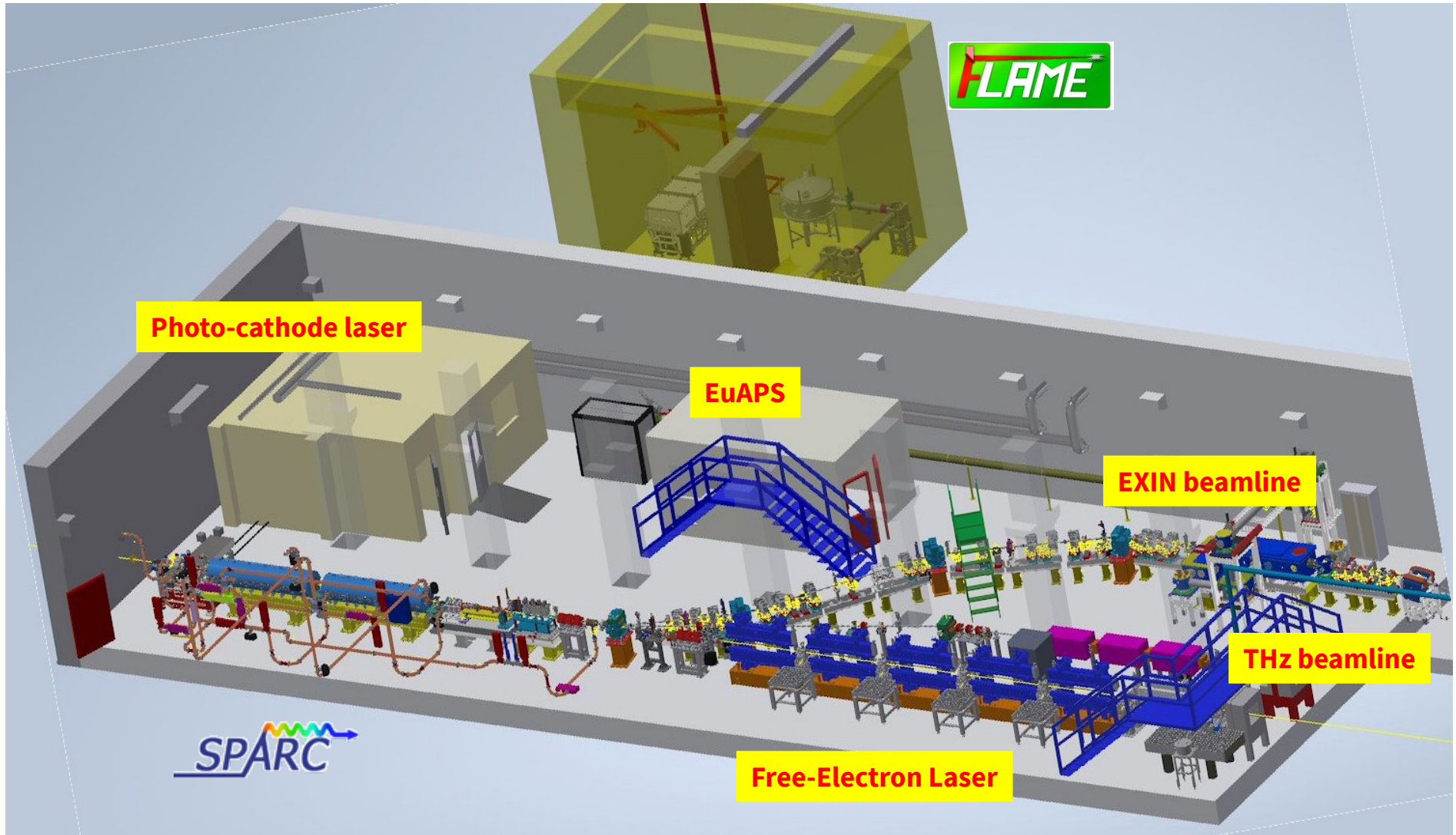


$$E_0 = \frac{m_e c \omega_p}{e} \approx 96 \sqrt{n_0 (\text{cm}^{-3})} \rightarrow E_0 \approx 10 \frac{\text{GV}}{\text{m}} @ n_0 = 10^{16} \text{cm}^{-3}$$

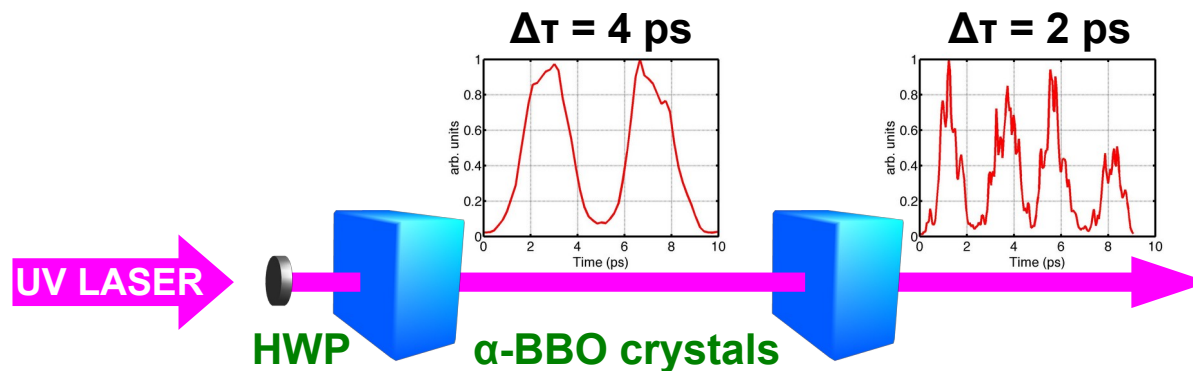
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PLASMA RESEARCH
ACCELERATOR WITH
EXCELLENCE IN
APPLICATIONS



This project has received funding from the European Union's Horizon Europe research and innovation programme under grant agreement No. 101079773



Ferrario, M., et al. "SPARC_LAB present and future." NIMB 309 (2013): 183-188.



Ferrario, M., et al "Laser comb with velocity bunching: Preliminary results at SPARC." NIM 637.1 2011 S43-S46.

To generate the **driver** and **witness** bunches necessary for plasma acceleration we always used the laser-comb technique

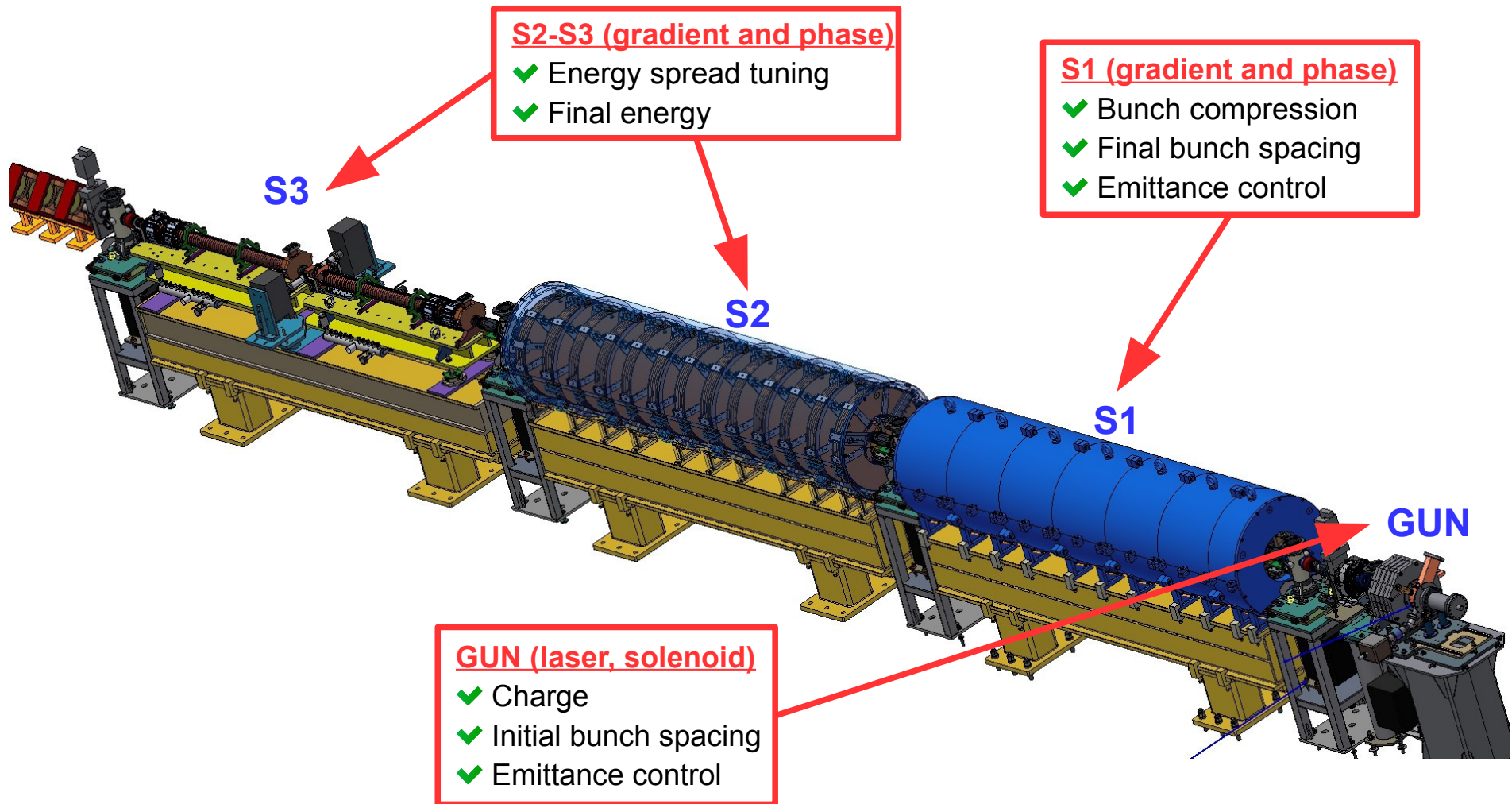
Multi-bunch structure directly generated on the cathode by using a multi-pulse laser system

The laser pulses can be tuned in energy, distance and durations

→ the corresponding electron bunches are tuned in charge, distance, durations!

It's a robust technique used so far for multiple purposes (plasma acceleration, THz radiation, test of advanced diagnostics, etc.)

Villa, F., et al. "Laser pulse shaping for multi-bunches photo-injectors." NIM A 740 (2014): 188-192.



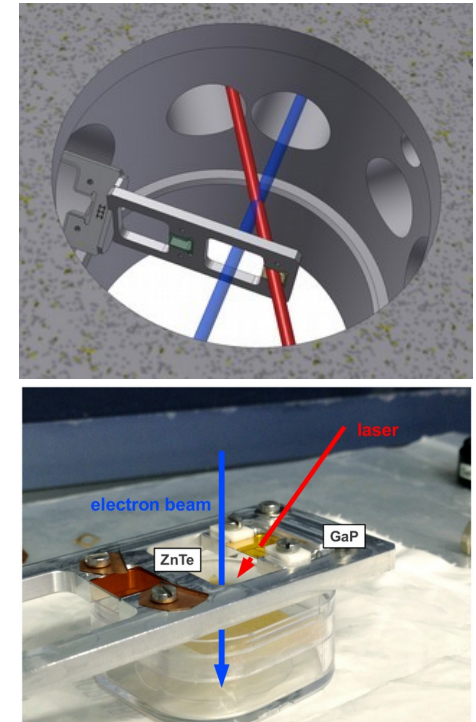
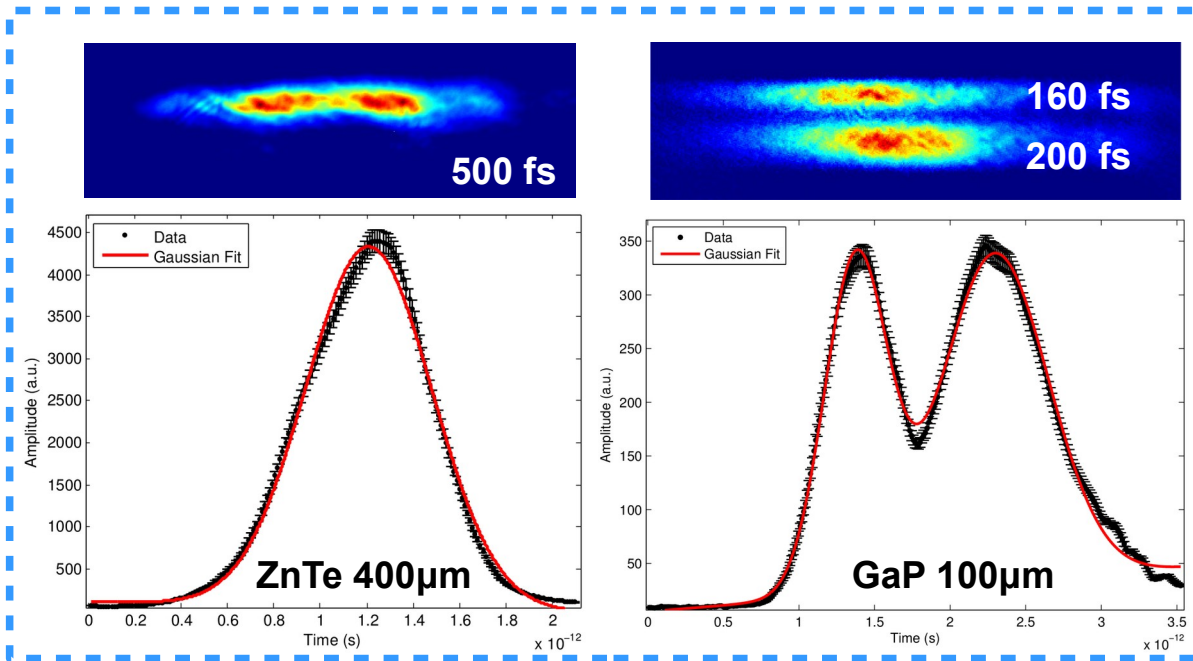
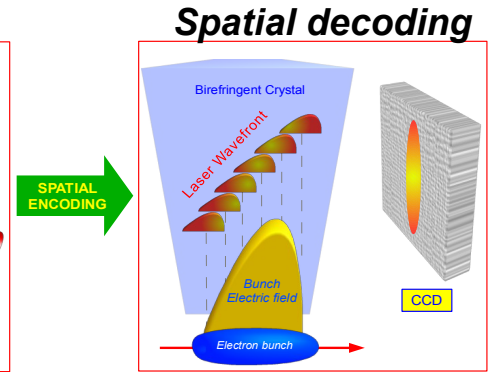
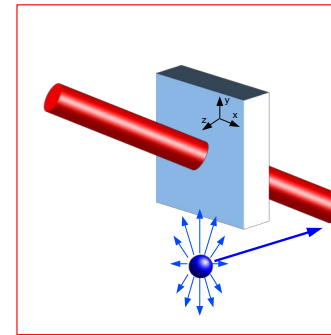
Chiadroni, E., et al. "Characterization of the THz radiation source at the Frascati linear accelerator." RSI 84.2 2013
 Mostacci, A., et al. "Advanced beam manipulation techniques at SPARC." Proceedings of IPAC2011

Multi-bunches trains have been measured with Electro-Optical Sampling

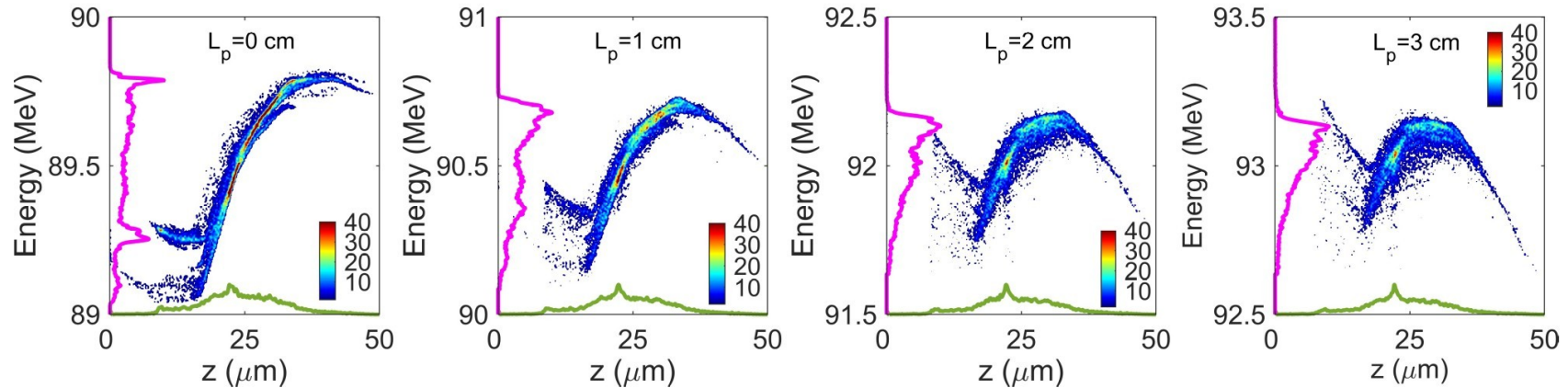
Single-shot, non-intercepting

80 fs (rms) temporal resolution

Goal: monitor beam injection in plasma

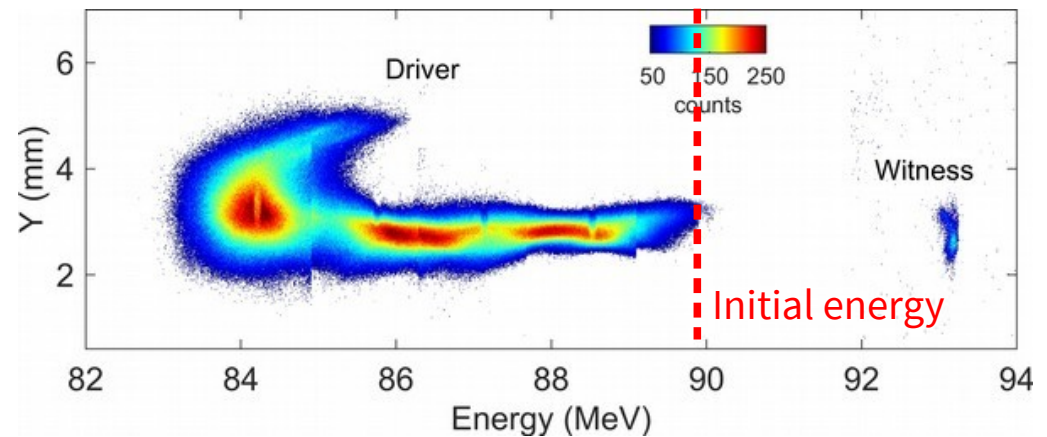


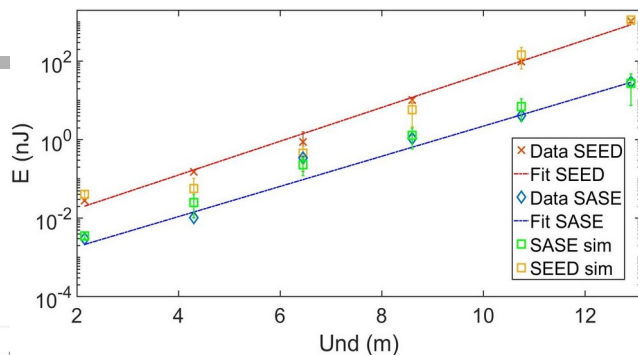
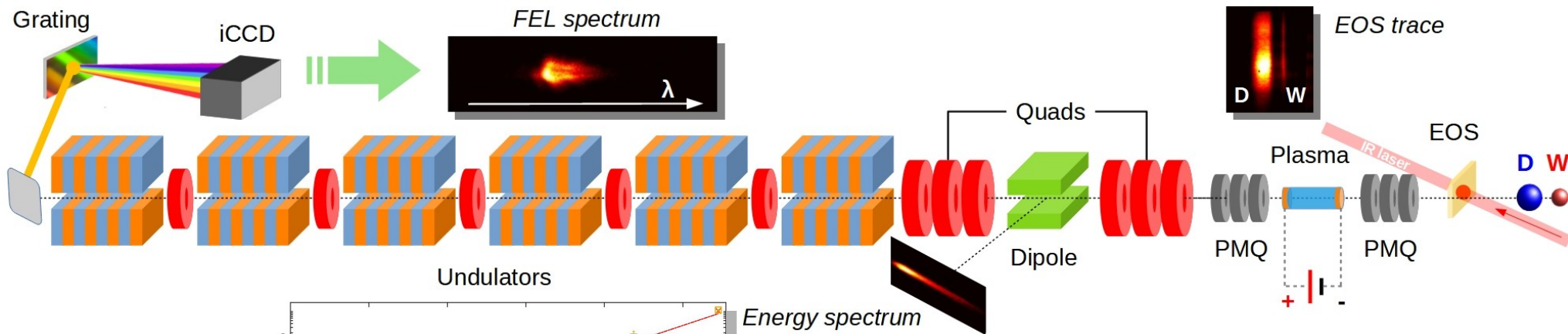
R. Pompili, et al., NIM A: Accelerators. 740, 216 (2014).



Energy spread minimization in a beam-driven plasma wakefield accelerator

R. Pompili¹, D. Alesini¹, M. P. Anania¹, M. Behtouei¹, M. Bellaveglia¹, A. Biagioni¹, F. G. Bisesto¹, M. Cesarini^{1,2}, E. Chiadroni¹, A. Cianchi³, G. Costa¹, M. Croia¹, A. Del Dotto¹, D. Di Giovenale¹, M. Diomede¹, F. Dipace¹, M. Ferrario¹, A. Giribono¹, V. Lollo¹, L. Magnisi¹, M. Marongiu¹, A. Mostacci², L. Piersanti¹, G. Di Pirro¹, S. Romeo¹, A. R. Rossi⁴, J. Scifo¹, V. Shpakov¹, C. Vaccarezza¹, F. Villa¹ and A. Zigler^{1,5}





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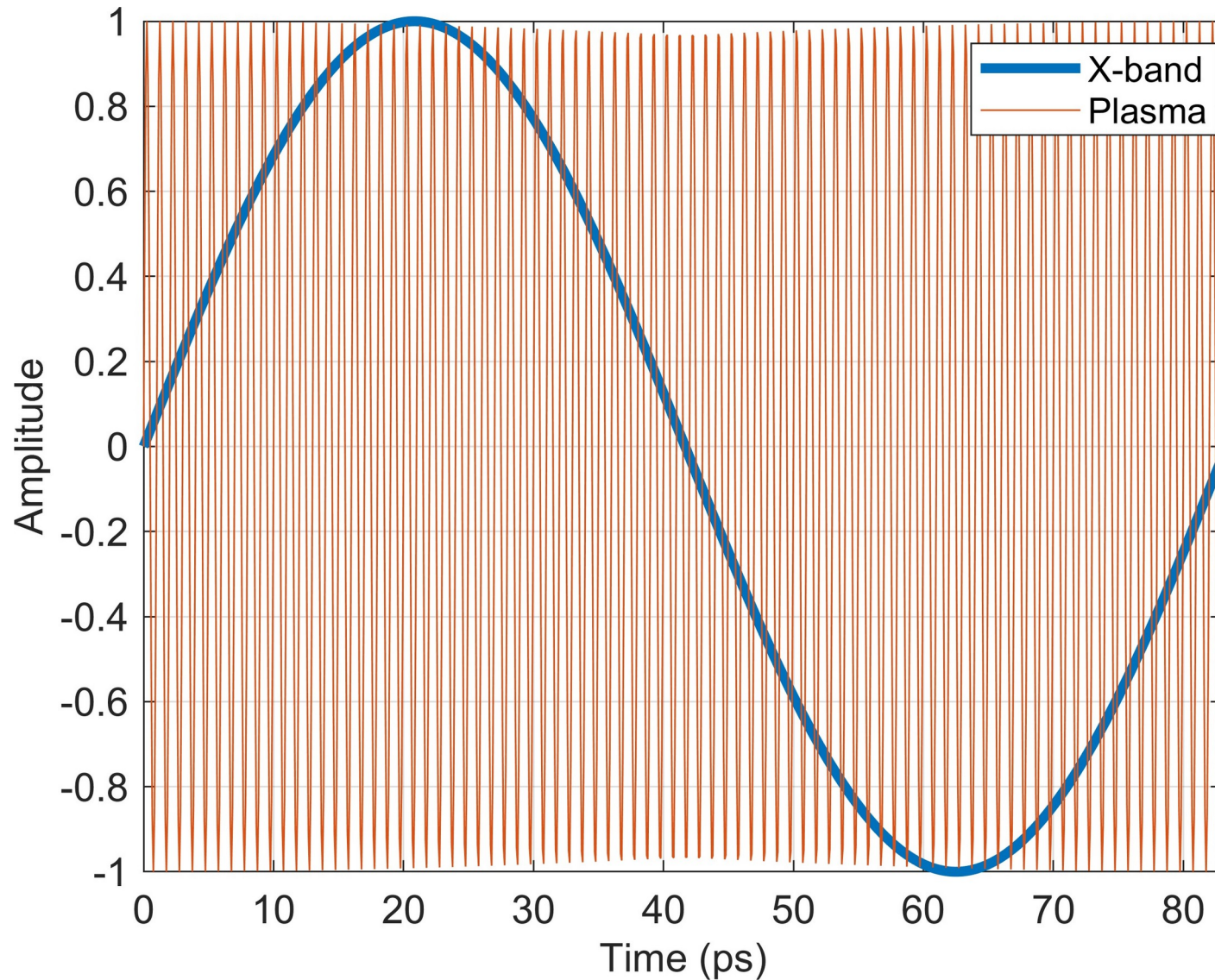
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Stable Operation of a Free-Electron Laser Driven by a Plasma Accelerator

M. Galletti *et al.*
Phys. Rev. Lett. **129**, 234801 – Published 29 November 2022



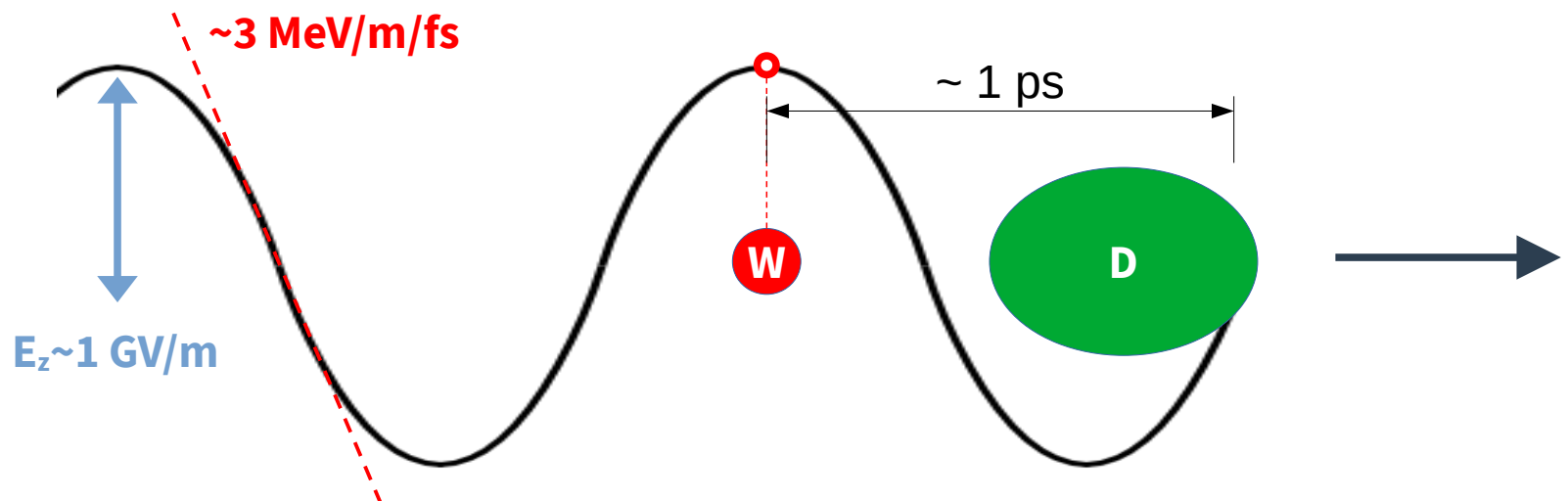
A typical configuration that is used at SPARC_LAB (and possibly will be implemented in EuPRAXIA) foresees the generation of a driver-witness beam configuration with the following parameters

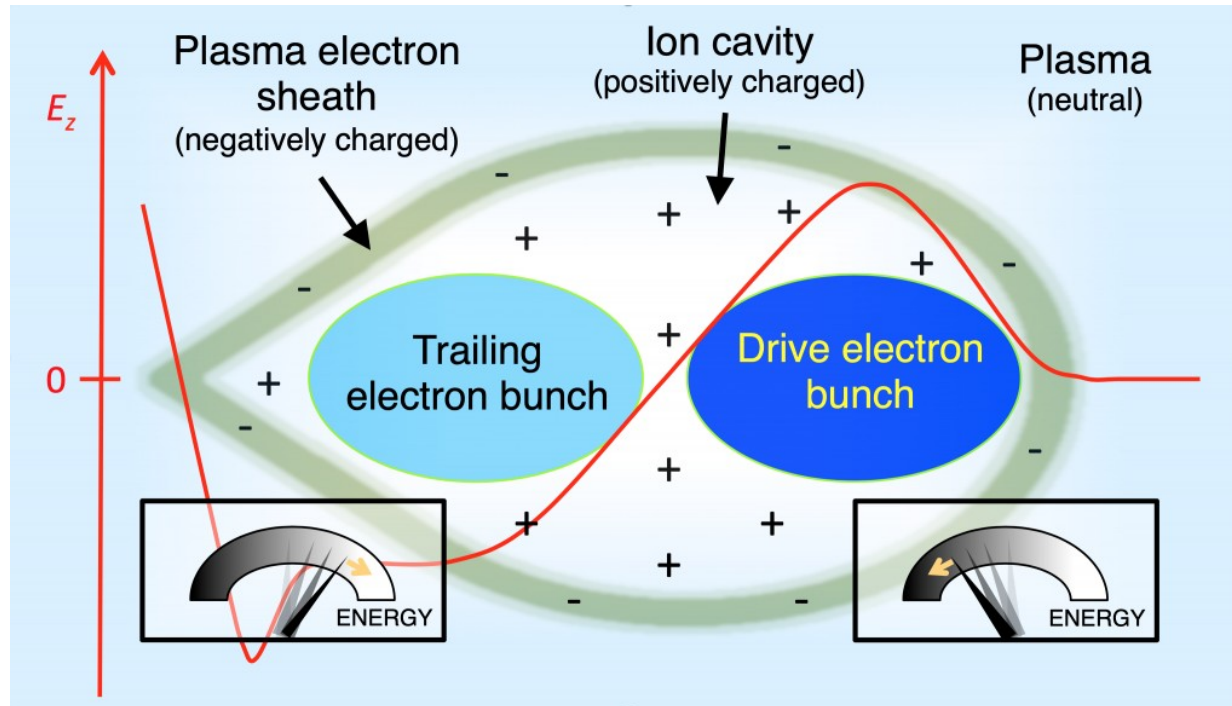
Charge: 400 pC (driver) + 50 pC (witness)

Distance: ~ 1 ps (corresponding to $\sim 3/4$ of the plasma oscillation period).

Durations: as short as possible (typically ~ 100 - 200 fs for the driver and < 30 fs for the witness)

Witness behind driver (in order to be accelerated by the plasma wakefield excited by the latter)



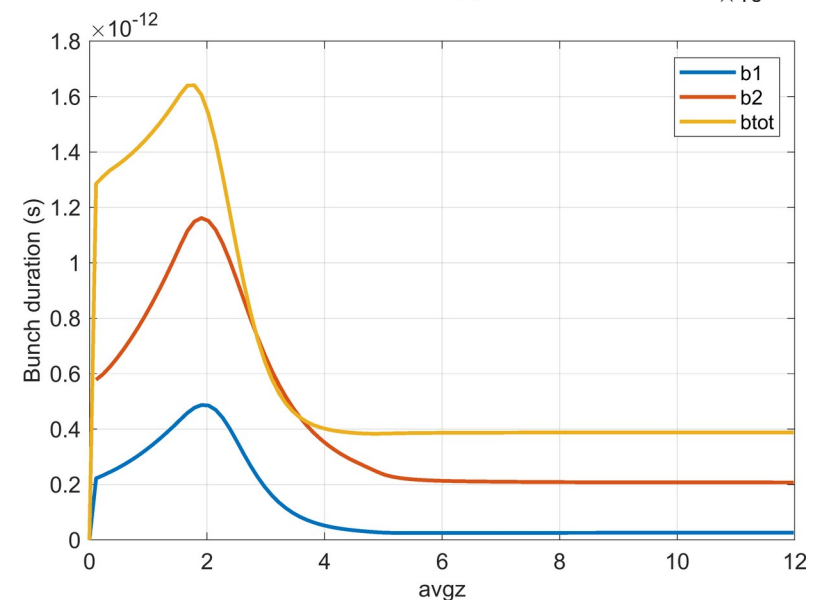
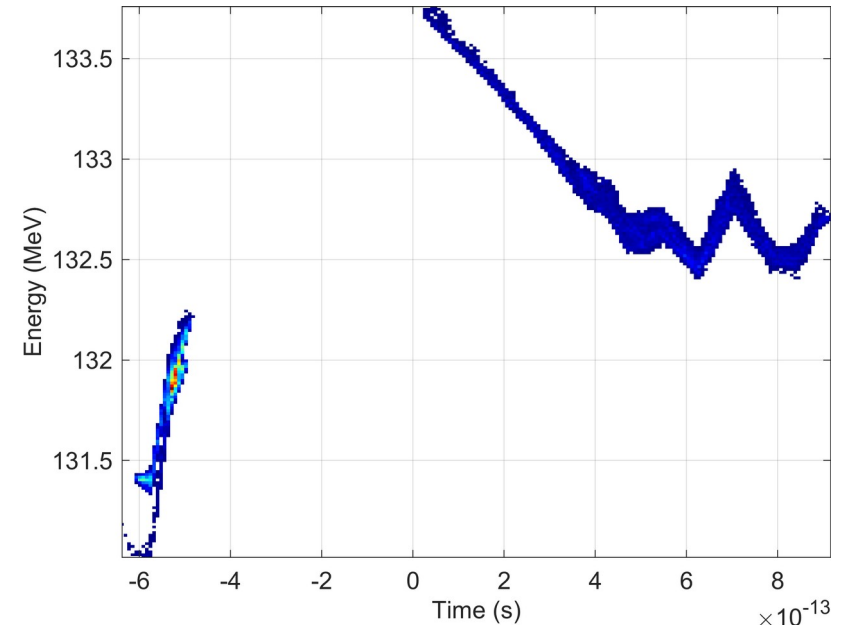
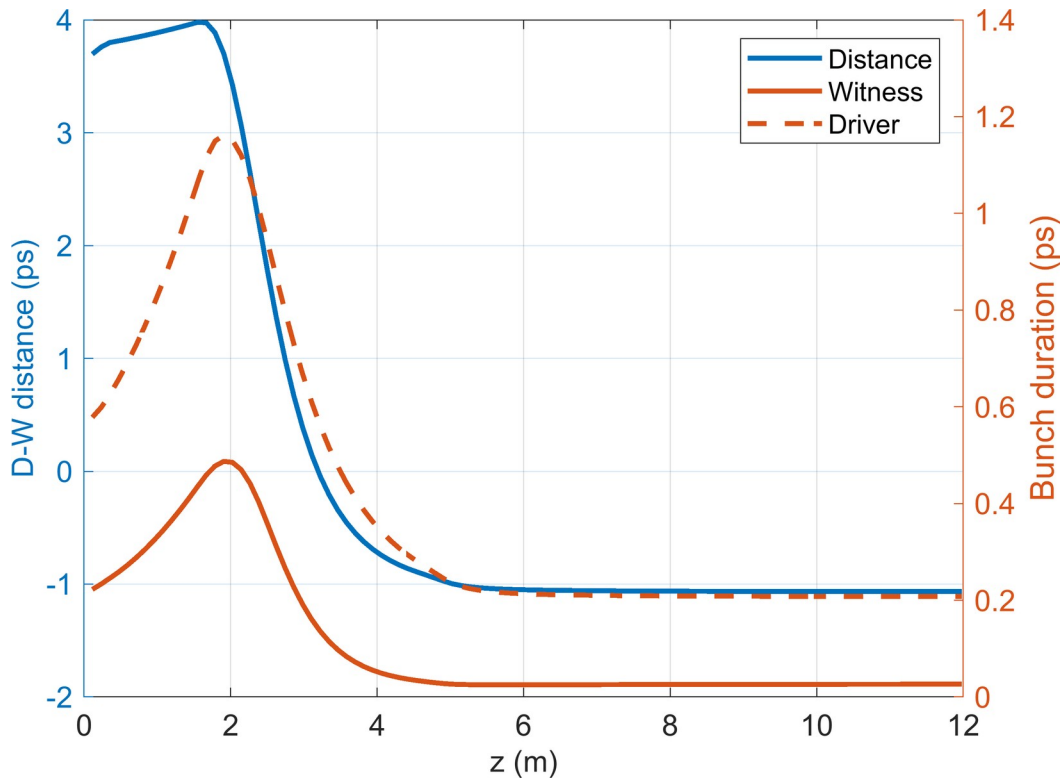


<https://loa.ensta-paris.fr/research/upx-research-group/particle-beam-driven-plasma-wakefield-acceleration/>

In nonlinear plasma wakefield regimes the behavior is not sinusoidal!

The slope becomes larger close to the plasma “bubble” tail → 3 MeV/m/fs can be considered a lower bound

Witness beam-loading effects makes the picture even more complicated...

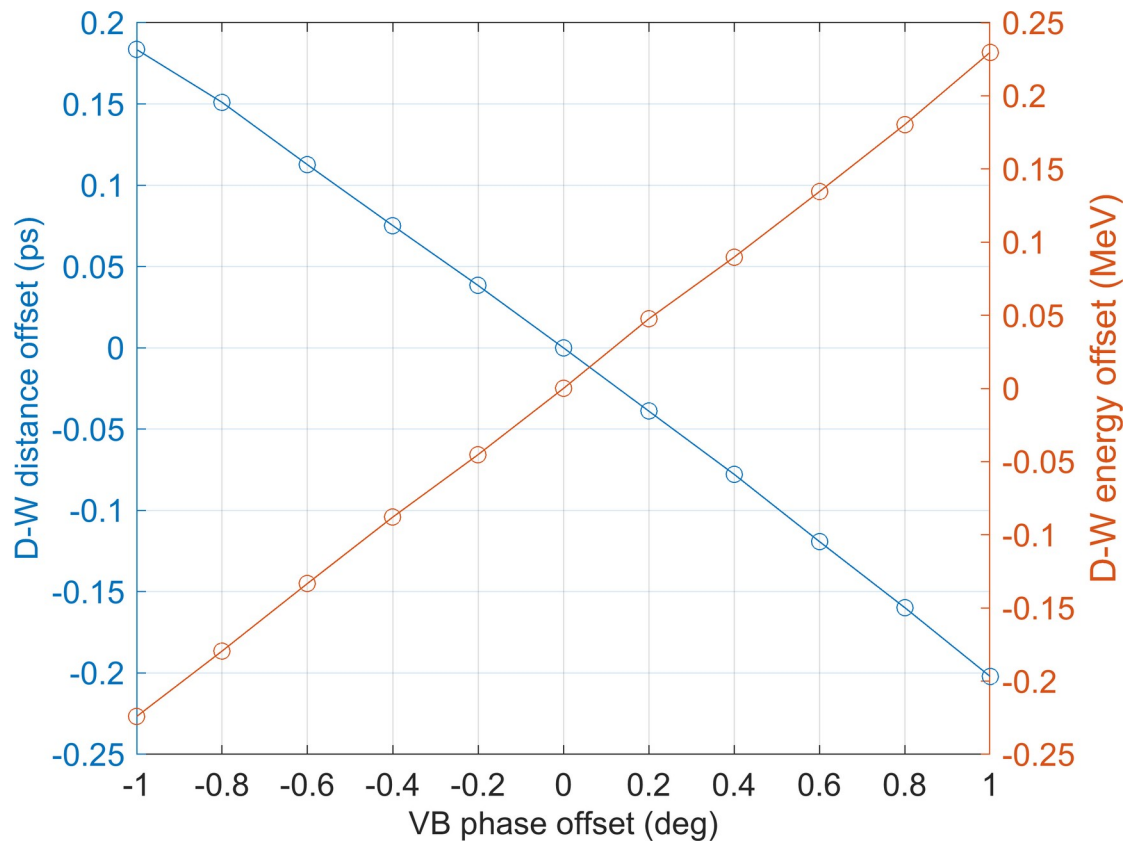


Two delayed (4 ps) laser pulses are sent to the cathode

Witness is photo-emitted in advance with respect to the driver

S1 RF phase is set at -90° (wrt on-crest phase) \rightarrow velocity-bunching longitudinal compression takes place

During the VB process, the driver-witness bunches are shortened and their positions are swapped \rightarrow witness behind the driver



Effects of RF phase-jitter on the driver-witness distance

1° offset → ~220 fs change in the D-W distance

Considering the plasma wakefield slope (~3 MeV/m/fs) and a state-of-the-art RF jitter of ~0.02° →

~5 fs D-W distance jitter → ~15 MeV energy jitter for a witness accelerated in 1 m long plasma

The RF timing-jitter also affects the bunch durations during the velocity-bunching process

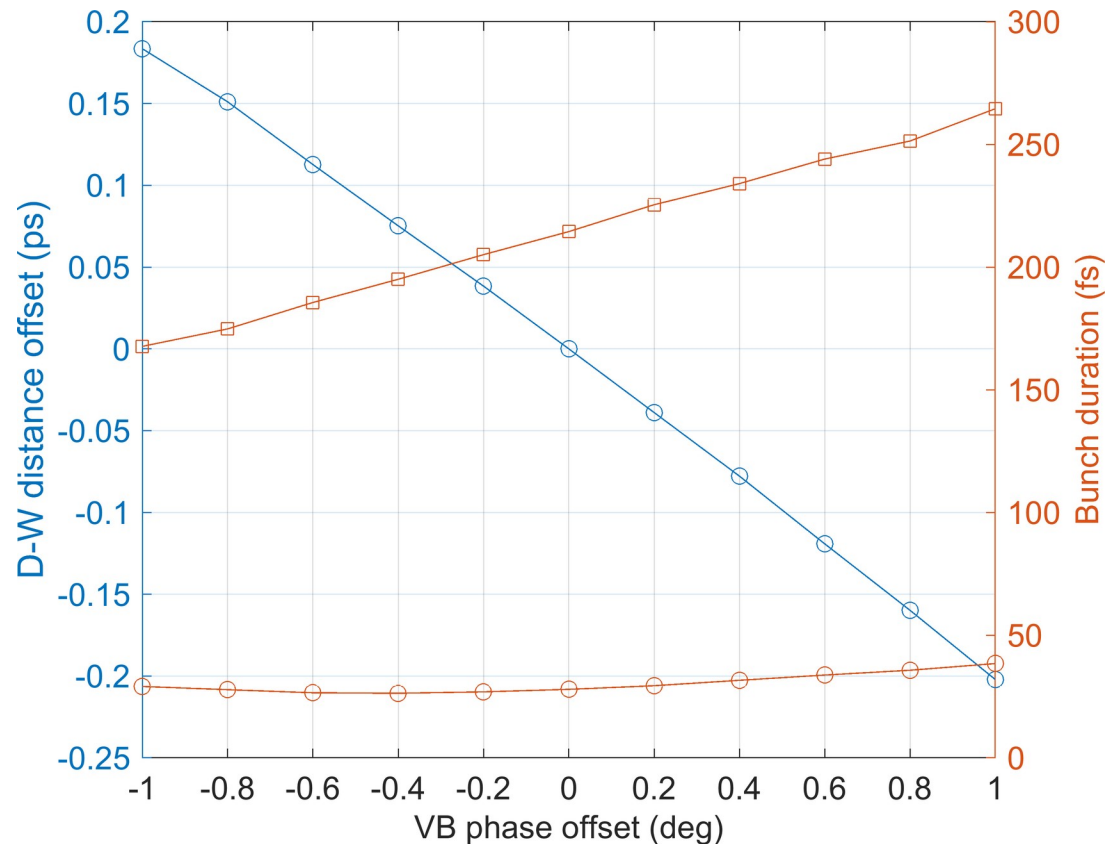
This translates in an additional plasma wakefield jitter since

$$E_z \propto I_{driver} \propto \frac{1}{\sigma_{t,D}^{1/3}}$$

In this case 1° RF phase offset increase/decrease the driver duration by ~ 50 fs

Considering again a state-of-the-art RF jitter of $\sim 0.02^\circ \rightarrow \sim 0.3\%$ plasma wake amplitude jitter

$\rightarrow 3$ MeV energy jitter with 1 m plasma



Plasma wakefield acceleration (PWFA) enables GV/m accelerating gradients overcoming RF accelerator technology by orders of magnitude

Proof-of-principle experiments demonstrated the huge gradients, recent developments demonstrated higher beam-quality

However, beam-driven PWFA still requires a conventional RF linac upstream of the plasma

→ very demanding requirements in terms of phase-stability and beam brightness

Few fs timing-jitters can produce tens of MeV energy jitter downstream of the plasma → too large if compared with currently operating facilities (e.g. for FEL)

Need to push phase-stability beyond state-of-the-art

Thanks!

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