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# Solving Beautiful Puzzles

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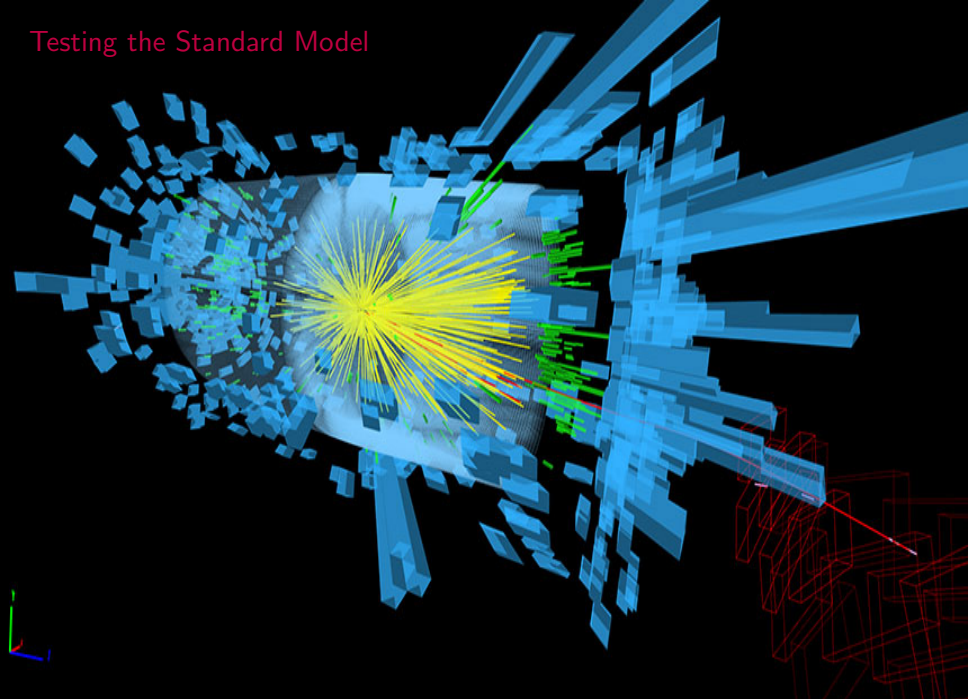
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K. Keri Vos

Maastricht University & Nikhef

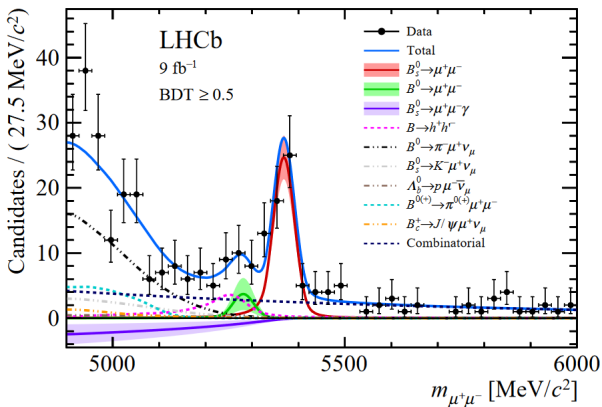
= Laboratori Nazionali di Frascati 2024 =

# Testing the Standard Model



# Testing the Standard Model: Indirect

LHCb Collaboration [Phys. Rev. Lett. 128, (2022) 041801]



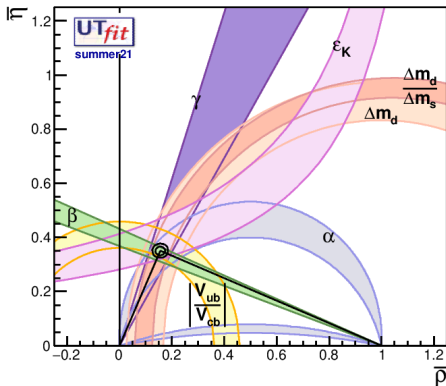
## Precision frontier

Tiny deviations from SM predictions constrain effects of New Physics

# The Flavour Puzzle

Thanks to Marcella Bona for providing the 2021 plots

$$\bar{\rho} + i\bar{\eta} = -\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}$$



Huge amounts of data + theory advances = Precision frontier

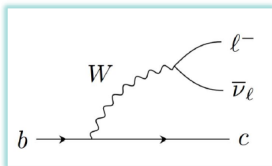
Tiny deviations from SM predictions constrain effects of New Physics

# SM or beyond?

## Challenge:

Disentangle SM long-distances effects from the effects of new interactions

Quark level process

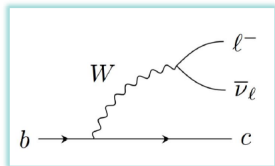


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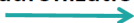
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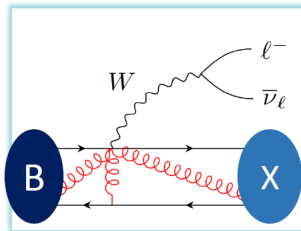
Quark level process



hadronization



Reality: Bound state

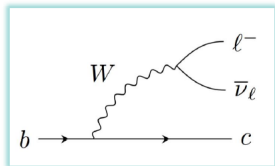


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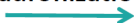
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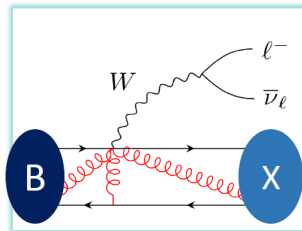
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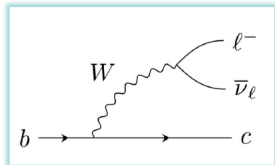
- Reliable theory uncertainties are essential!

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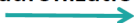
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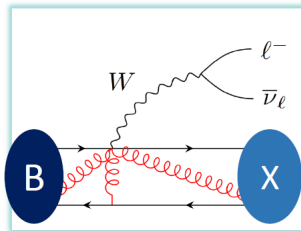
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Reality: Bound state



- Reliable theory uncertainties are essential!
- Look for the cleanest observables/methods

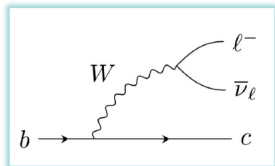


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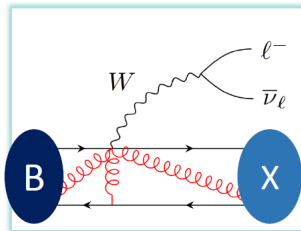
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hadronization



Reality: Bound state



- Reliable theory uncertainties are essential!
- Look for the cleanest observables/methods
- Some anomalies already spotted

# Puzzles in Flavour Physics

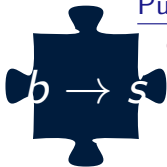
## Challenge:

Disentangle SM long-distance effects from NP effects from CPV effects



## Puzzles in semileptonic decays

- Inclusive versus Exclusive
- $V_{cb}$  and  $V_{ub}$
- LFUV in  $R_D$  and  $R_{D^*}$



## Puzzles in rare decays

- Anomalies in  $b \rightarrow sll$



# Inclusive versus Exclusive Decays

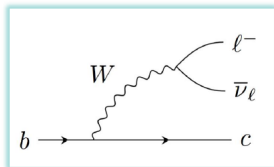
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# The Beauty of Semileptonic Decays

## Motivation:

- Theoretically relatively easy to describe: factorization of strong interaction effects

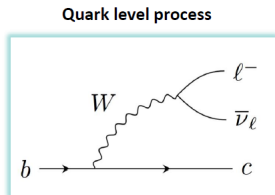
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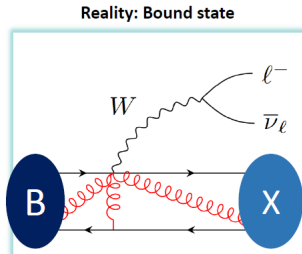
# The Beauty of Semileptonic Decays

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hadronization  
→



# The Beauty of Semileptonic Decays

## Motivation:

- Theoretically relatively easy to describe: factorization of strong interaction effects



## Two options:

- Exclusive decays: pick one final state with the desired quarks ( $V_{cb} \rightarrow D^{(*)}$  and  $V_{ub} \rightarrow \pi$ )
- Inclusive decays: everything you can think of! (denoted with  $X_c$  or  $X_u$ )

# The Beauty of Semileptonic Decays

## Motivation:

- Theoretically relatively easy to describe: factorization of strong interaction effects



## Challenge:

- Dealing with QCD at large distances/small scales
- Parametrize fundamental mismatch in non-perturbative objects
  - Calculable: Lattice or Light-cone sumrules = **Exclusive Decays**
  - Measurable: from data = **Inclusive Decays**

# Why inclusive decays?

- Set up OPE and heavy quark expansion
- Well established framework
- Extract important CKM parameters  $|V_{cb}|, |V_{ub}|$  (and  $|V_{cs}|$ ?)
- Extract power corrections from data
- Cross check of exclusive decays



# The Heavy Quark Expansion

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# Inclusive Decays = Heavy Quark Expansion

- $b$  quark mass is large compared to  $\Lambda_{\text{QCD}}$
- Setting up the HQE: momentum of  $b$  quark:  $p_b = m_b v + k$ , expand in  $k \sim iD$
- Optical Theorem  $\rightarrow$  (local) Operator Product Expansion (OPE)

$$d\Gamma = d\Gamma_0 + \frac{d\Gamma_1}{m_b} + \frac{d\Gamma_2}{m_b^2} + \dots \quad d\Gamma_i = \sum_k C_i^{(k)} \langle B | O_i^{(k)} | B \rangle$$

- $C_i^{(k)}$  perturbative Wilson coefficients
- $\langle B | \dots | B \rangle$  non-perturbative matrix elements  $\rightarrow$  string of  $iD$
- operators contain chains of covariant derivatives

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- Standard tool for inclusive  $B \rightarrow X_c \ell \nu$  decays

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HQE elements:  $\langle B | \mathcal{O}_i^{(n)} | B \rangle = \langle B | \bar{b}_v (iD_\mu) \dots (iD_{\mu_n}) b_v | B \rangle$

- Extracted from kinematic moments of the data
- Ideas for the lattice Juetner et al. [202305.14092]

$\Gamma_i$  are power series in  $\mathcal{O}(\alpha_s)$

$$\Gamma = \Gamma_0 + \frac{1}{m_b} \Gamma_1 + \frac{1}{m_b^2} \Gamma_2 + \frac{1}{m_b^3} \Gamma_3 \dots$$

- $\Gamma_0$ : decay of the free quark (partonic contributions),  $\Gamma_1 = 0$
- $\Gamma_2$ :  $\mu_\pi^2$  kinetic term and the  $\mu_G^2$  chromomagnetic moment

$$2M_B \mu_\pi^2 = - \langle B | \bar{b}_v i D_\mu i D^\mu b_v | B \rangle$$

$$2M_B \mu_G^2 = \langle B | \bar{b}_v (-i \sigma^{\mu\nu}) i D_\mu i D_\nu b_v | B \rangle$$

- $\Gamma_3$ :  $\rho_D^3$  Darwin term and  $\rho_{LS}^3$  spin-orbit term

$$2M_B \rho_D^3 = \frac{1}{2} \langle B | \bar{b}_v [i D_\mu, [i v D, i D^\mu]] b_v | B \rangle$$

$$2M_B \rho_{LS}^3 = \frac{1}{2} \langle B | \bar{b}_v \{ i D_\mu, [i v D, i D_\nu] \} (-i \sigma^{\mu\nu}) b_v | B \rangle$$

- $\Gamma_4$ : 9 parameters Mannel, Turczyk, Uraltsev, JHEP 1010 (2011) 109
- $\Gamma_5$ : 18 parameters Mannel, Turczyk, Uraltsev, JHEP 1010 (2011) 109

# Powercounting in the HQE

I:  $m_Q \sim m_q \gg \Lambda_{\text{QCD}}$  OPE for  $b \rightarrow c l \bar{\nu}$

- $q$  is treated as a heavy degree of freedom
- two-quarks operators:  $\bar{Q}_v(iD^\alpha \dots iD^\beta)Q_v$
- IR sensitivity to mass  $m_q$

$$\Gamma \Big|_{1/m_Q^3} = \left[ \frac{34}{3} + 8 \log \rho + \dots \right] \frac{\rho_D^3}{m_Q^3}, \quad \text{with } \rho = (m_q/m_Q)^2$$

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II:  $m_Q \gg m_q \gg \Lambda_{\text{QCD}}$  start with  $q$  dynamical

- four-quark operators  $(\bar{Q}_v \Gamma q)(q \bar{\Gamma} Q_v)$
- removed when matching onto two-quark operators
- RGE running gives  $\log(m_q/m_Q)$

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III:  $m_Q \gg m_q \sim \Lambda_{\text{QCD}}$  OPE for  $c \rightarrow s \ell \bar{\nu}$  Fael, Mannel, KKV [1910.05234]

- $q$  dynamical degree of freedom
- four-quark operators remain in OPE
- no explicit  $\log(m_q/m_Q)$ : hidden inside new non-perturbative HQE parameters



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IV:  $m_Q \gg \Lambda_{\text{QCD}} \gg m_q$  for  $b \rightarrow u$  and  $c \rightarrow d$  transitions

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III and IV have four-quark (weak annihilation) effects

# Weak Annihilation

Uraltsev, Bigi, Voloshin, Mannel, Turczyk; Ligeti, Luke, Manohar, Phys. Rev.D82 (2010) 033003  
Gambino, Kamenik, Nucl.Phys.B840 (2010) 424

- IR sensitivity to light quark gives additional four-quark non-pert. parameters:

$$\langle B | (\bar{b}_v \gamma^\nu P_L q) (\bar{q} \gamma^\mu P_L b_v) | B \rangle = 2M_B [T_1(\mu) g^{\mu\nu} + T_2(\mu) v^\mu v^\nu]$$

- Starting at  $\mathcal{O}(1/m_b^3)$  and mix with  $\rho_D^3$  under renormalization
- Challenging to study non-perturbatively
- Very important to achieve precise  $B \rightarrow X_{d,s} \ell \ell$  predictions  
Hurth, Huber, Lunghi, Jenkins, Qin, KKV [2007.04191,2404.03517]

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- Can be obtained from  $D$  and  $D_s$  semileptonics using HQE?

$$B_{WA}^{bq} = \frac{m_B f_B^2}{m_D f_D^2} B_{WA}^{cq}$$

- Effect is  $(m_b/m_c)^3$  enhanced compared to  $B$  decays

# Heavy quark expansion for charm?

- Expansion parameters  $\alpha_s(m_c)$  and  $\Lambda_{\text{QCD}}/m_c$  less than unity, but not so small ...
- Turn vice into virtue: more sensitive to higher  $1/m_Q$  corrections
- Exploit the full physics potential of BES III, LHCb ...
- Constrain Weak Annihilation (WA) contributions

$$\rightarrow B_d \rightarrow s\ell\ell$$

$$\rightarrow V_{ub}$$

[Huber, Hurth, Lunghi, Jenkins, KKV, Qin]

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- Extraction of  $|V_{cs}|$  and  $|V_{cd}|$ ?

[Huber, Hurth, Lunghi, Jenkins, KKV, Qin]

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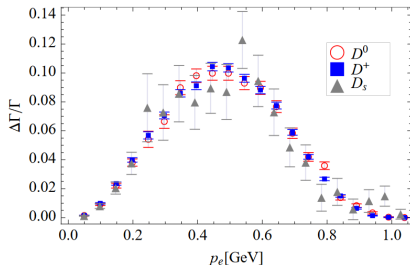
[Huber, Hurth, Lunghi, Jenkins, KKV, Qin]

## Challenges:

- Valence and non-valence WA operators at higher orders
- Scale for radiative corrections
- Charm mass definition

# Extracting weak annihilation from data

CLEO data, Gambino, Kamenik [1004.0114]

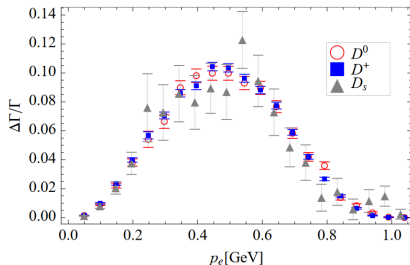


- Lepton energy moments extracted from spectrum
- Kinetic mass for charm at  $\mu = 0.5$  GeV threshold, HQE parameters as input
- Max 2% weak annihilation (WA) contribution to  $B \rightarrow X_u \ell \nu$



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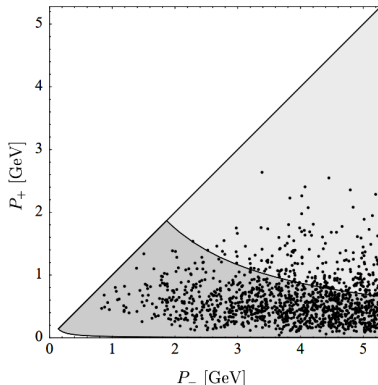


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- **In progress:** Feasibility study to measure  $q^2$  moments at BESIII Bernlochner, Gilman, Malde, Prim, KKV, Wilkinson

# Inclusive $B \rightarrow X_u$ semileptonic decays

# Modified Heavy Quark Expansion

- Cuts needed to suppress large charm background
- Pushes towards specific corner of the phase space
  - Local OPE as in  $b \rightarrow c$  cannot work
  - Sensitivity to  $b$ -quark wave function properties (Fermi motion)
  - Deal with energetic light degrees of freedom
  - **More than two scales involved!**
- Expansion parameter  $\Lambda_{\text{QCD}}/(m_b - 2E_\ell)$
- Use light-cone OPE with light-cone directions  $n$  and  $\bar{n}$



# Factorization of scales

- Separates the different scales in the problem

$$d\Gamma = H \otimes J \otimes S$$

- H: Hard scattering kernel at  $\mathcal{O}(m_b)$
- J: universal Jet function at  $\mathcal{O}(\sqrt{m_b\Lambda_{\text{QCD}}})$
- S: Shape function at  $\mathcal{O}(\Lambda_{\text{QCD}})$
- Framework to include radiative corrections (+ NNLL resummation)
- Introduces 3 subleading shape functions

- Leading order shape functions

$$2m_B f(\omega) = \langle B(v) | \bar{b}_v \delta(\omega + i(n \cdot D)) b_v | B(v) \rangle$$

- Leading order shape functions

$$2m_B f(\omega) = \langle B(v) | \bar{b}_v \delta(\omega + i(n \cdot D)) b_v | B(v) \rangle$$

- Universal
- Similar to parton distribution in DIS

- Leading order shape functions

$$2m_B f(\omega) = \langle B(v) | \bar{b}_v \delta(\omega + i(n \cdot D)) b_v | B(v) \rangle$$

- Charged Lepton Energy Spectrum (at leading order)

$$\frac{d\Gamma}{dy} \sim \int d\omega \theta(m_b(1-y) - \omega) f(\omega)$$

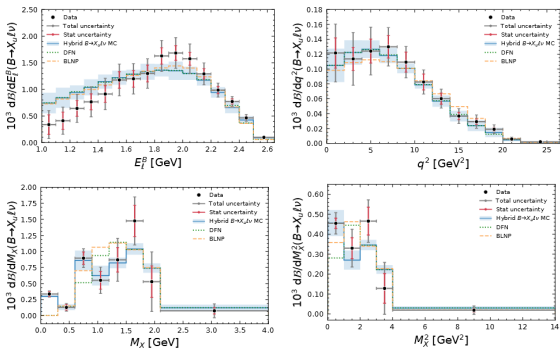
- Moments of the shapefunction are related to HQE ( $b \rightarrow c$ ) parameters:

$$f(\omega) = \delta(\omega) + \frac{\mu_\pi^2}{6m_b^2} \delta''(\omega) - \frac{\rho_D^3}{m_b^3} \delta'''(\omega) + \dots$$

- Shape function is non-perturbative and cannot be computed

# Shape function parametrization

Differential spectra from BelleII [2107.13855]



- Often linked to  $B \rightarrow X_s \gamma$
- Updated experimental measurements could constrain SFs further



Different frameworks for inclusive  $B \rightarrow X_u$ :

- **BLNP**: Bosch, Lange, Neubert, Paz uses Soft Collinear Effective Theory (SCET)
- **GGOU** Gambino, Giordano, Ossola, Uraltsev
  - OPE with hard-cutoff
  - No subleading SFs

Approaches to calculate the SF perturbatively:

- **DGE**: Dressed Gluon Exponentiation Andersen, Gardi
- **ADFR** Aglietti, Di Lodovico, Ferrerar, Ricciardi

Average of all four approaches:

$$|V_{ub}|_{\text{incl}} = \sqrt{\frac{\Delta\mathcal{B}(B \rightarrow X_u \ell \nu)}{\tau_B \delta\Gamma(B \rightarrow X_u \ell \nu)}} = (4.10 \pm 0.28) \cdot 10^{-3}$$

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Exclusive world average:  $|V_{ub}|_{\text{excl}} = (3.44 \pm 0.12) \cdot 10^{-3}$

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- **ADFR** Aglietti, Di Lodovico, Ferrerar, Ricciardi

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$$|V_{ub}|_{\text{incl}} = \sqrt{\frac{\Delta\mathcal{B}(B \rightarrow X_u \ell \nu)}{\tau_B \delta\Gamma(B \rightarrow X_u \ell \nu)}} = (4.10 \pm 0.28) \cdot 10^{-3}$$

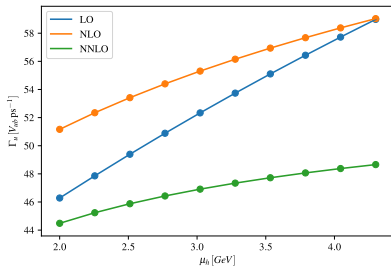
Inclusive determinations need to be scrutinized

## Update of BLNP approach

- Systematic framework: Soft Collinear Effective Theory (SCET)
- **In progress:** include known  $\alpha_s^2$  corrections

# Shape function parametrization

Preliminary! Olschewsky, Lange, Mannel, KKV [240x..xxxx]



- $\alpha_s^2$  corrections give large corrections [see also Pezszak 2019]
- Required to make precision predictions

## Update of BLNP approach

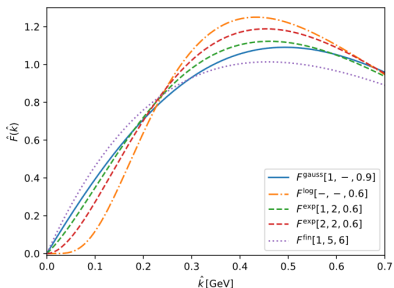
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## Update of BLNP approach

- Systematic framework: Soft Collinear Effective Theory (SCET)
- **In progress:** include known  $\alpha_s^2$  corrections
- Moments of shape functions can be linked to HQE parameters in  $b \rightarrow c$ 
  - **In progress:** include higher-moments
  - kinetic mass scheme as in  $b \rightarrow c$
- Shape function is non-perturbative and cannot be computed
  - **In progress:** new flexible parametrization

# Shape function parametrization

Olschewsky, Lange, Mannel, KKV [240x.xxxx]

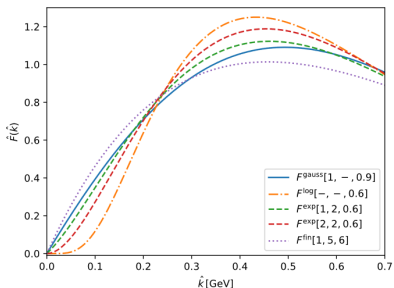


- All moments of shape functions are linked to HQE parameters
- Allows for a range of different shapes  $\rightarrow$  systematic uncertainty



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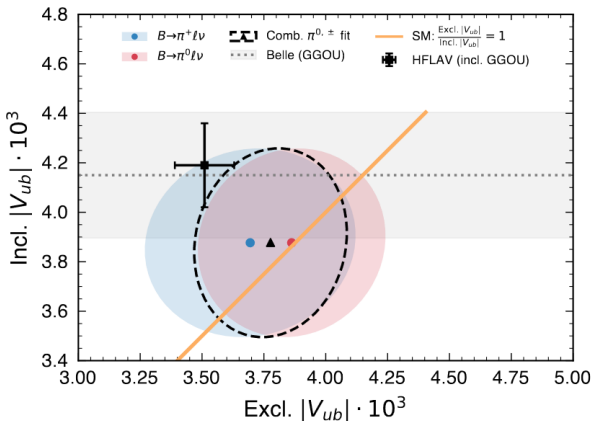
In progress:

Lange, Mannel, Olschewsky, KKV [in progress]

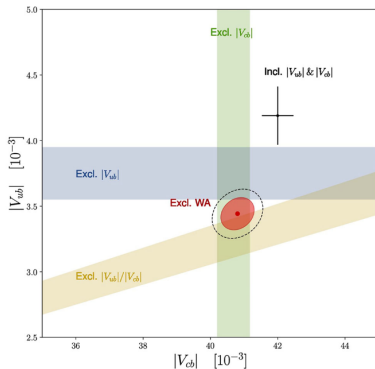
$$|V_{ub}|_{\text{incl}} = \text{Stay Tuned!}$$

# Inclusive versus Exclusive

BelleII [2303.17309]



- First simultaneous measurement
- Experimental advantages due to common backgrounds and modeling



- Includes also ratio measurements of  $|V_{ub}/V_{cb}|$  from LHCb ( $B_s$  and  $\Lambda$ )

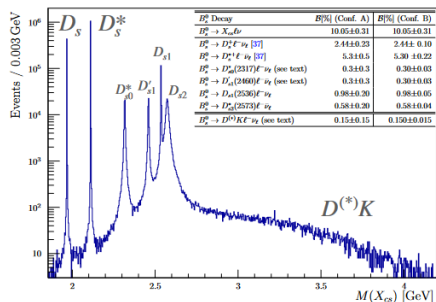
# Inclusive Measurements at LHCb?

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# Inclusive $B_s$ decays?

De Cian, Feliks, Rotondo, KKV [2312.05147]. Pic from M. Fael

First study of the possibilities using sum-over-exclusive technique

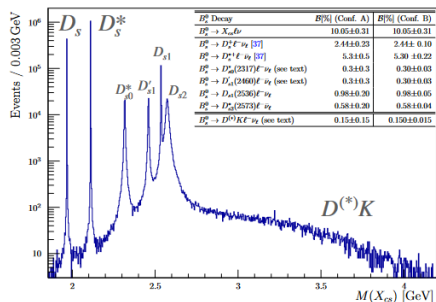


- $B_s$  spectrum well-separated
- Only  $M_X^2$  moments available
- Study  $SU(3)$  breaking of HQE

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- $B_s$  spectrum well-separated
- Only  $M_X^2$  moments available
- Study  $SU(3)$  breaking of HQE

- Improve knowledge  $D_S^{**}$  states
- Understand non-resonant contribution
- $|V_{cb}|$  extraction requires Branching ratio from Belle II!

Set up an OPE as for  $B \rightarrow X_u$

- Power-corrections give large uncertainties
- Normalizing to  $B \rightarrow X_u$  may reduce uncertainty:

$$\mathcal{R}(q_0^2) = \int_{q_0^2}^{M_B^2} dq^2 \frac{d\Gamma(\bar{B} \rightarrow X_s \ell^+ \ell^-)}{dq^2} \bigg/ \int_{q_0^2}^{M_B^2} dq^2 \frac{d\Gamma(\bar{B} \rightarrow X_u \ell \bar{\nu})}{dq^2}$$

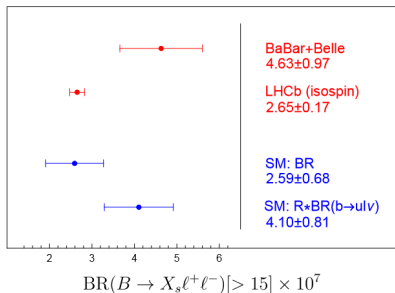
# Inclusive $B \rightarrow X_s \ell \ell$ decays

Hurth, Huber, Lunghi, Jenkins, Qin, KKV [2007.04191, 2404.03517]

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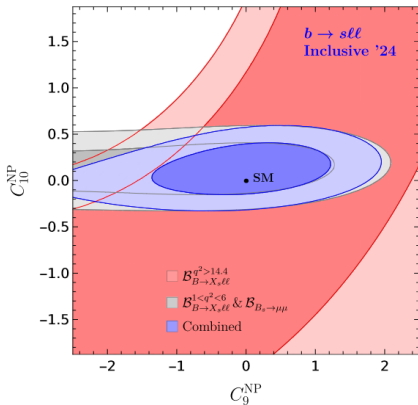
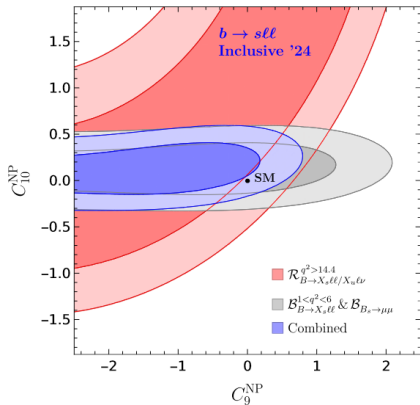
- At high- $q^2$   $X_s = K, K^*, K\pi, \dots$
- Use sum-over-exclusives from LHCb measurements!



# Inclusive $B \rightarrow X_s \ell \ell$ decays

Hurth, Huber, Lunghi, Jenkins, Qin, KKV [2007.04191,2404.03517]

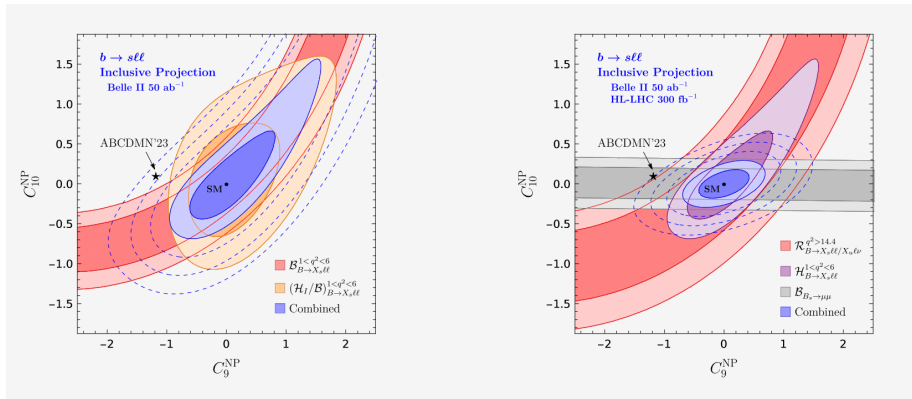
Important to cross-check the exclusive channels!



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Important to cross-check the exclusive channels!



- New measurements by Belle [2311.00458]
- **New!**  $\alpha_s^3$  corrections for  $b \rightarrow u$  Fael, Usovitsch [2310.03685]
  
  
  
  
  
  
  
  
  
  
- We can predict the  $B \rightarrow X\ell\nu$  rate in local OPE

- New measurements by Belle [2311.00458]
- **New!**  $\alpha_s^3$  corrections for  $b \rightarrow u$  Fael, Usovitsch [2310.03685]
- **In progress:** Direct calculation of the ratio

$$C \equiv \left| \frac{V_{cb}}{V_{ub}} \right|^2 \frac{\mathcal{B}(B \rightarrow X_u \ell \nu)}{\mathcal{B}(B \rightarrow X_c \ell \nu)}$$

- Either in shapefunction region or in local OPE (see also Mannel, Rahimi, KKV [2105.02163])
- We can predict the  $B \rightarrow X \ell \nu$  rate in local OPE

# Solving beautiful puzzles

Many exciting Puzzles remain

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- Comparing Inclusive and Exclusive important to test QCD description
- Need to revise previous assumptions and ensure reliable systematic uncertainties

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Close collaboration between theory and experiment necessary!



# Backup

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# What mass to use?

Bigi, Shifman, Uraltsev, Vainshtein, hep-ph/9704245, hep-ph/9405410; Czarnecki, Melnikov, Uraltsev, hep-ph/9708372.

- Renormalon issues require short-distance mass
- Kinetic mass: relating hadron versus quark mass  
QCD corrections using hard cut off  $\mu$

$$m_Q(\mu)^{\text{kin}} = m_Q^{\text{Pole}} - [\bar{\Lambda}]_{\text{pert}} + \left[ \frac{\mu_\pi^2}{2m_Q} \right]_{\text{pert}} + \dots$$

$$[\bar{\Lambda}]_{\text{pert}} = \frac{4}{3} C_F \frac{\alpha_s(m_c)}{\pi} \mu \quad [\mu_\pi^2]_{\text{pert}} = C_F \frac{\alpha_s(m_c)}{\pi} \mu^2$$

- Higher-order terms in the HQE generate corrections  $(\alpha_s/\pi)\mu^n/m_Q^n$ .

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- Higher-order terms in the HQE generate corrections  $(\alpha_s/\pi)\mu^n/m_Q^n$ .
- $\Lambda_{\text{QCD}} < \mu < m_Q$ : expansion parameters  $\mu/m_Q$ 
  - Well established for  $m_B$ :  $\mu/m_B \simeq 0.2$
  - Charm??
    - $\rightarrow \mu = 1 \text{ GeV} \rightarrow \mu/m_c \simeq 1$
    - $\rightarrow \mu = 0.5 \text{ GeV} \rightarrow \mu/m_c \simeq 0.4$

Challenge:  $\mu = 0.5 \text{ GeV}$  touches upon the non-perturbative regime?

# Ratios of $V_{cb}$ and $V_{ub}$ : a $B_s$ puzzle

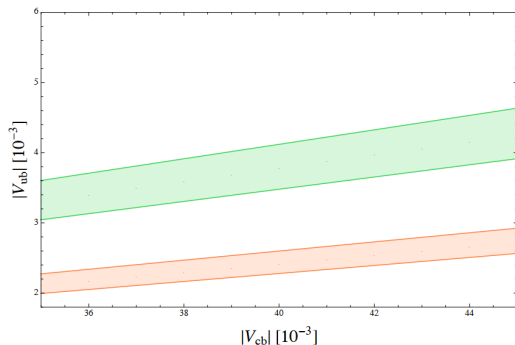
Bolognani, van Dyk, KKV [2308.0437]  
LHCb [2012.05143], Khodjamirian, Rusov [2017]

- Also  $B_s \rightarrow K\mu\nu$  is sensitive to  $|V_{ub}|$
- Only accessible at LHCb, but normalization needed
- Using  $B \rightarrow D\mu\nu$  gives access to the ratio

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LHCb high  $q^2$  ratio:  $FF_K$  determined with LQCD

LHCb low  $q^2$  ratio:  $FF_K$  determined with LCSR

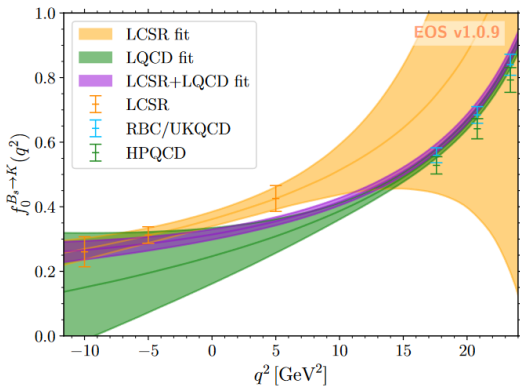
$$\left| \frac{V_{ub}}{V_{cb}} \right|_{\text{low } q^2} = 0.061 \pm 0.004$$
$$\left| \frac{V_{ub}}{V_{cb}} \right|_{\text{high } q^2} = 0.095 \pm 0.008$$

**3.8 $\sigma$**

# A puzzle in $B_s$ decays?

Bolognani, van Dyk, KKV [2308.0437]  
LHCb [2012.05143], Khodjamirian, Rusov [2017]

- **Recent update:** New form factor predictions combining lattice and light-cone sumrule information



# A puzzle in $B_s$ decays?

Bolognani, van Dyk, KKV [2308.0437]  
LHCb [2012.05143], Khodjamirian, Rusov [2017]

- **Recent update:** New form factor predictions combining lattice and light-cone sumrule information
- Puzzle becomes less:  $1.9\sigma$  difference

$$q^2 < 7\text{GeV}^2 \rightarrow \frac{|V_{ub}|}{|V_{cb}|} = 0.0681 \pm 0.004 \quad q^2 > 7\text{GeV}^2 \rightarrow \frac{|V_{ub}|}{|V_{cb}|} = 0.0801 \pm 0.005$$

# Spectral-Density Mass

Chetyrkin, Kuehn, Steinhauser hep-ph/9705254, Penin, Pivovarov hep-ph/9805344 Boushmelev, Mannel, KKV [2301.05607]

- $m_c$  not observable  $\rightarrow$  no physical meaning
- Extracted from data: moments of the spectral density in  $e^+e^- \rightarrow$  hadrons

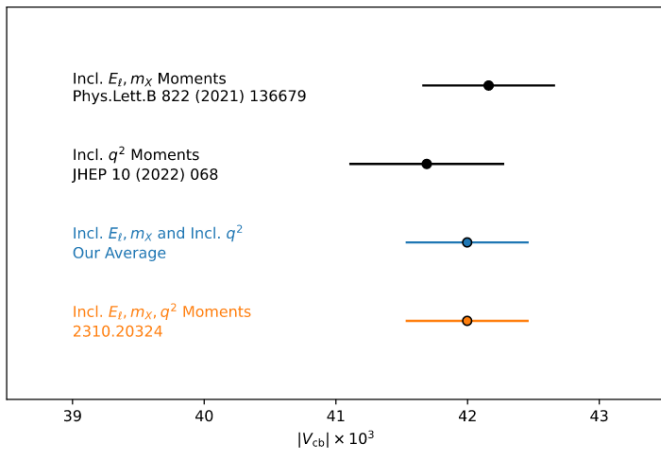
$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

- Replace  $m_c$  by moments of the spectral density!
- First study shows small improvement in pert. series
- **In progress:** Similar approach for the charm + power corrections



# Summary of $|V_{cb}|$ inclusive

Fael, Prim, KKV, Eur. Phys. J. Spec. Top. (2024). <https://doi.org/10.1140/epjs/s11734-024-01090-w>



- Need new branching ratio measurements!

# Contamination of the $B \rightarrow X_c \ell \nu$ signal

Rahimi, Mannel, KKV [arXiv: 2105.02163]

Avoid background subtraction by calculating the full inclusive width:

$$d\Gamma(B \rightarrow X\ell) = d\Gamma(B \rightarrow X_c \ell \bar{\nu}) + d\Gamma(B \rightarrow X_u \ell \bar{\nu}) + d\Gamma(B \rightarrow X_c (\tau \rightarrow \ell \bar{\nu} \nu) \bar{\nu})$$

- $b \rightarrow u \ell \nu$  contribution: suppressed by  $V_{ub}/V_{cb}$
- $b \rightarrow c (\tau \rightarrow \mu \nu \bar{\nu}) \bar{\nu}$  contribution: phase space suppressed
- QED effects
- Quark-hadron duality violation?

## Goal:

provide theoretical description and compare with Monte-Carlo data used by Belle (II)

## Challenge:

estimate how much this description would improve  $V_{cb}$  determination

# $b \rightarrow ul\nu$ contribution: Local OPE

Neubert (1994); Bosch, Paz, Lange, Neubert (2004,2005)

- Can be analyzed in local OPE as  $B \rightarrow X_c l \nu$  by taking  $m_c \rightarrow 0$  limit
- For  $V_{ub}$  determination
  - large charm background requires experimental cuts
  - reduces the inclusivity and local OPE no longer converges
  - spectrum described by non-local OPE
  - convolution of pert. coefficients with shape function

## Goal:

provide theoretical description and compare with Monte-Carlo data used by Belle (II)

- NLO +  $1/m_b^2 + 1/m_b^3$
- In agreement with partonic calc of DFN De Fazio, Neubert (1999); Gambino, Ossola, Uraltsev (2005)
- First study: no  $\alpha_s$  for  $1/m_b^2$ , no additional uncert. due to missing higher orders
- Inputs HQE parameters from  $B \rightarrow X_c l \nu$  study Gambino, Schwanda [2014]; Gambino, Healey, Turczy [2016]

# Monte Carlo versus HQE

Rahimi, Mannel, KKV [arXiv: 2105.02163]; De Fazio, Neubert 1999; Bosch, Lange, Neubert, Paz 2005

Compare local OPE with generator level Monte-Carlo data provided by Cao, Bernlochner

## Monte Carlo:

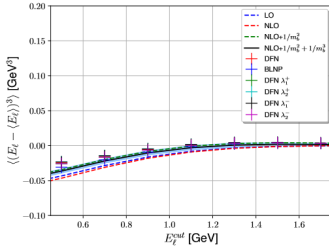
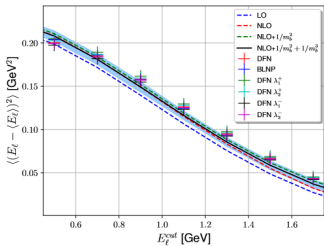
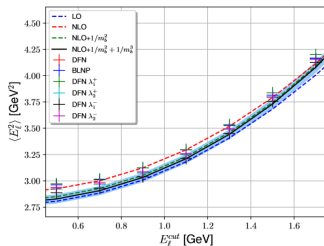
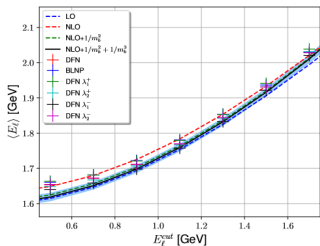
- BLNP: specific shape function input parameters shape function parameters  $b = 3.95$  and  $\Lambda = 0.72$
- DFN:  $\alpha_s$  corrections convoluted with the exponential shape function model
  - Inputs from  $B \rightarrow X_c \ell \nu$  and  $B \rightarrow X_s \gamma$  data using KN-scheme Kagan, Neubert 1998
  - $(\lambda_1^+, \lambda_2^+, \lambda_1^-, \lambda_2^-)$  are obtained by varying  $\bar{\Lambda}$  and  $\mu_\pi^2$  within  $1\sigma$  Buchmuller, Flacher, 2006

Hadronic contributions: “hybrid Monte Carlo” Belle Collaboration [arXiv:2102.00020.]

- convolution with hadronization simulation based on PYTHIA
- plus explicit resonances:  $\bar{B} \rightarrow \pi \ell \bar{\nu}$  and  $\bar{B} \rightarrow \rho \ell \bar{\nu}$

# Monte Carlo versus HQE

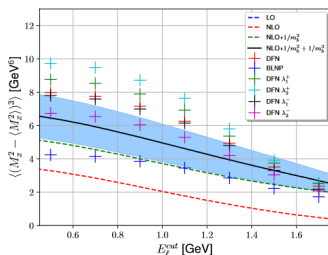
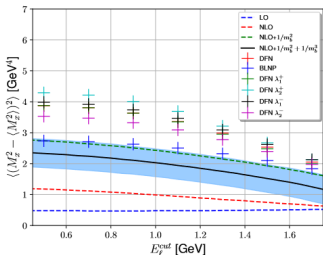
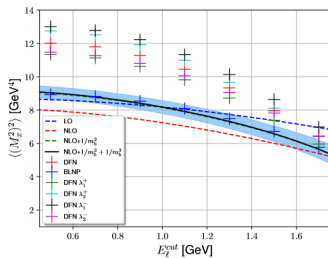
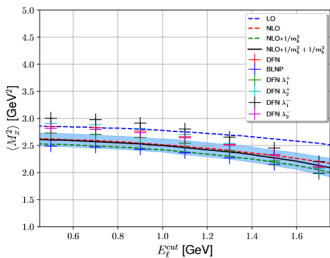
Rahimi, Mannel, KKV [arXiv: 2105.02163]; MC data by Lu Cao and Florian Bernlochner



MC-results are in good agreement with the HQE results

# Monte Carlo versus HQE

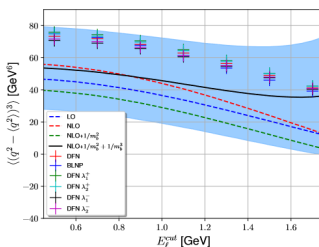
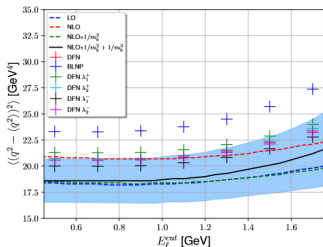
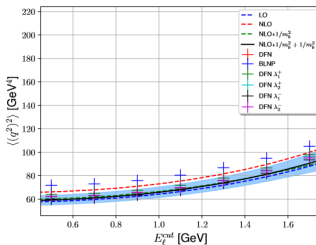
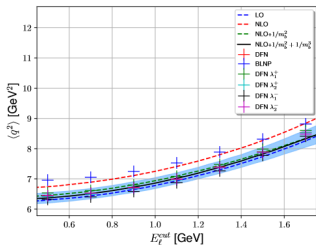
Rahimi, Mannel, KKV[arXiv: 2105.02163]; MC data by Lu Cao and Florian Bernlochner



Wide spread between MC for higher moments

# Monte Carlo versus HQE

Rahimi, Mannel, KKV[arXiv: 2105.02163]; MC data by Lu Cao and Florian Bernlochner



## Remarks:

- DFN: Smearing corresponding to a shape function, mimicking some non-perturbative effects; may not capture all
- BLNP: should reproduce the HQE, with parameters adjusted to local HQE prediction
  - should include higher moments of the shape-function model?
  - include subleading shape functions?
- our HQE: interesting to include  $\alpha_s$  to HQE parameters,  $\alpha_s^2$ ?



# Spectral-Density Mass

Chetyrkin, Kuehn, Steinhauser hep-ph/9705254, Penin, Pivovarov hep-ph/9805344

- $m_c$  not observable  $\rightarrow$  no physical meaning
- Extracted from data: moments of the spectral density in  $e^+e^- \rightarrow$  hadrons

$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

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- Start from vacuum correlator

$$\int d^4x e^{-iqx} \langle 0 | T[j_\mu(x)j_\nu(0)] | 0 \rangle = (g_{\mu\nu}q^2 - q_\mu q_\nu) \Pi(q^2)$$

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- Expand around  $q^2 = 0$ : ( $\bar{C}_n = \bar{C}_n^{(0)} + \frac{\alpha_s(\mu)}{\pi} \bar{C}_n^{(1)} + \dots$ )

$$\Pi(q^2) = \Pi(0) + \frac{4}{9} \frac{3}{16\pi^2} \sum_{n=1}^{\infty} \bar{C}_n \left( \frac{q^2}{4m_c^2} \right)$$

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- $\bar{C}_n$  known up to  $\alpha_s^2$  and related to moments

$$\bar{C}_n = (4m_c^2)^n M_n \quad \text{with} \quad M_n = \int \frac{ds}{s^{n+1}} R(s) \quad (1)$$

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$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

- Expand around  $q^2 = 0$ : ( $\bar{C}_n = \bar{C}_n^{(0)} + \frac{\alpha_s(\mu)}{\pi} \bar{C}_n^{(1)} + \dots$ )

$$\Pi(q^2) = \Pi(0) + \frac{4}{9} \frac{3}{16\pi^2} \sum_{n=1}^{\infty} \bar{C}_n \left( \frac{q^2}{4m_c^2} \right) = \Pi(0) + \frac{q^2}{12\pi^2} \int \frac{ds}{s} \frac{R(s)}{s - q^2}$$

- $\bar{C}_n$  known up to  $\alpha_s^2$  and related to moments

$$\bar{C}_n = (4m_c^2)^n M_n \quad \text{with} \quad M_n = \int \frac{ds}{s^{n+1}} R(s) \quad (1)$$

- Replace  $m_c$ :

$$m_c = \frac{1}{2} \left( \frac{\bar{C}_n}{M_n} \right)^{1/(2n)}$$

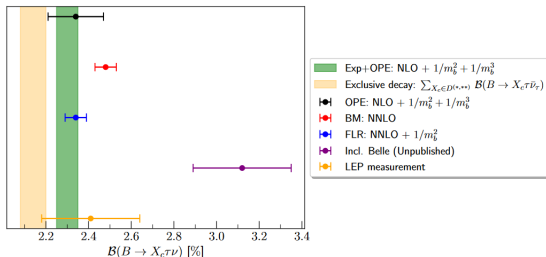
$$R_{e/\mu}(X) \equiv \frac{\Gamma(B \rightarrow X_c e \bar{\nu}_e)}{\Gamma(B \rightarrow X_c \mu \bar{\nu}_\mu)}$$

- Belle II result:  $R(X_{e/\mu}) = 1.033 \pm 0.022$  with cut, see H. Junkerkalefeld [ICHEP]
- In agreement with new SM predictions:  $1.006 \pm 0.001$  at  $1.2\sigma$

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