Relic neutrino Background from Cosmic-Ray Reservoirs

Andrea Giovanni De Marchi University of Bologna and INFN Dark Matter and Cosmic Rays, Napoli November 11th 2024





Based on: AGDM, Granelli, Nava, Sala [2405.04568]

Not observed yet! $v_e + {}^{3}H \rightarrow {}^{3}He^+ + e^-$ Direct detection via capture on tritium (PTOLEMY)^{1,2}: ${}^{3}\underline{H} \rightarrow {}^{3}He^+ + e^- + \overline{v}_e$



• Well separated peak, but...

Not observed yet!

Direct detection via capture on tritium (PTOLEMY)^{1,2}: ${}^{3}H \rightarrow {}^{3}He^{+} + e^{-} + \overline{\nu}_{e}$



- Well separated peak, but...
- Heisenberg uncertainty broadens the distribution, now hidden under background

 $\nu_{e} + {}^{3}\text{H} \rightarrow {}^{3}\text{He}^{+} + e^{-}$

Not observed yet! $v_e + {}^{3}H \rightarrow {}^{3}He^+ + e^-$ Direct detection via capture on tritium (PTOLEMY)^{1,2}: ${}^{3}\underline{H} \rightarrow {}^{3}He^+ + e^- + \overline{v}_e$



- Well separated peak, but...
- Heisenberg uncertainty broadens the distribution, now hidden under background
- Some final states work, but suppressed

Not observed yet! $v_e + {}^{3}H \rightarrow {}^{3}He^+ + e^-$ Direct detection via capture on tritium (PTOLEMY)^{1,2}: ${}^{3}H \rightarrow {}^{3}He^+ + e^- + \overline{v}_e$



- Well separated peak, but...
- Heisenberg uncertainty broadens the distribution, now hidden under background
- Some final states work, but suppressed
- Needs new tech, problem solved if ~ unbound tritium

Cosmic rays upscatter the relic neutrinos

 $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR!³

Cosmic rays upscatter the relic neutrinos

 $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR!³

Ciscar-Monsalvatje, Herrera, Shoemaker 2402.0098:

Cosmic rays upscatter the relic neutrinos $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR! ³

Ciscar-Monsalvatje, Herrera, Shoemaker 2402.0098:

• CRs in our galaxy:



Cosmic rays upscatter the relic neutrinos $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR! ³

Ciscar-Monsalvatje, Herrera, Shoemaker 2402.0098:

• CRs in our galaxy: $\eta_{\nu} \sim 10^{13}$



Cosmic rays upscatter the relic neutrinos $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR! ³

Ciscar-Monsalvatje, Herrera, Shoemaker 2402.0098:

• CRs in our galaxy:

 $\eta_{\nu} \sim 10^{13}$

• Blazars (TXS 0506+056):



Cosmic rays upscatter the relic neutrinos $\sigma(CR - \nu) \propto E_{CR}$ most upscattered by UHECR! ³

Ciscar-Monsalvatje, Herrera, Shoemaker 2402.0098:

• CRs in our galaxy: $\eta_{\nu} \sim 10^{13}$

• Blazars (TXS 0506+056): $\eta_{\nu} \sim 10^{10}$







Where are UHECR produced? Fang, Murase (2017, Nature Phys.) 1704.00015:

• Look at galaxy clusters



Where are UHECR produced? Fang, Murase (2017, Nature Phys.) 1704.00015:

- Look at galaxy clusters
- AGN jets inject UHE particles in the cluster $(\frac{d\Phi}{dE} \sim E^{-\alpha}, \ \alpha \in [2, 2.5])$



Where are UHECR produced? Fang, Murase (2017, Nature Phys.) 1704.00015:

- Look at galaxy clusters
- AGN jets inject UHE particles in the cluster $\left(\frac{d\Phi}{dE} \sim E^{-\alpha}, \ \alpha \in [2, 2.5]\right)$
- Clusters have $B \sim \mu G$ magnetic fields



Where are UHECR produced? Fang, Murase (2017, Nature Phys.) 1704.00015:

- Look at galaxy clusters
- AGN jets inject UHE particles in the cluster $(\frac{d\Phi}{dE} \sim E^{-\alpha}, \alpha \in [2, 2.5])$
- Clusters have $B \sim \mu G$ magnetic fields
- UHECRs spend $\tau_{esc} \sim {\rm Gyr}$ in the cluster before escaping



Where are UHECR produced? Fang, Murase (2017, Nature Phys.) 1704.00015:

- Look at galaxy clusters
- AGN jets inject UHE particles in the cluster $\left(\frac{d\Phi}{dE} \sim E^{-\alpha}, \alpha \in [2, 2.5]\right)$
- Clusters have $B \sim \mu G$ magnetic fields
- UHECRs spend $\tau_{esc} \sim {\rm Gyr}$ in the cluster before escaping

A lot of time to upscatter neutrinos!

Flux from Cosmic Reservoirs

- Improves previous bounds by orders of magnitude^{4,5}
- Overdensities only on cluster scale, not diffuse
- Can tell apart from Cosmogenic neutrinos:
 - Spectral shape (DIS is crucial)
 - Flavour composition



4. Ciscar-Monsalvatje, Herrera, Shoemaker 2402.00985

5. Franklin, Martinez-Soler, Perez-Gonzalez, Turner 2404.02202

Flavour composition

vs are non-relativistic $\Rightarrow \sigma$ depends on m_{ν}

We computed the flux of mass eigenstates v_i , preserved during propagation

At detection, the flux of flavour eigenstate v_{α} is

$$\frac{d\Phi_{\alpha}}{dE_{\nu}} = \sum_{i} |U_{\alpha i}|^2 \frac{d\Phi_{i}}{dE_{\nu}}$$



Flavour composition

- Higher neutrino mass: degenerate neutrinos, 1:1:1 flavour ratio
- Lower neutrino mass: the heaviest neutrino(s) dictate the flavour composition
- NO/IO: less/more electron neutrinos



Do these overdensities make sense?

- Limit to overdensity in SM: Pauli blocking, needs BSM
- Smirnov, Xu 2201.00939 get close with new Yukawa interaction
- Limit on mass of the cluster: alleviated by non-homogeneous distribution



Conclusions

- We implemented a corrected cross section and CR composition and showed this has orders of magnitude impact wrt to the literature
- We set the most stringent bound on η_{ν} (can be even stronger with correct normalization of CR flux and non homogeneous distribution)
- We provided two ways to disentangle this signal from others (cosmogenic):
 - Energy dependence (correct cross section including DIS is crucial)
 - Flavour composition, which depends on absolute neutrino masses

What else?

- BSM mechanisms that induce neutrino overdensities? v self-interactions? v-DM interactions?
- Improved modeling of CR diffusion in the reservoir? Very naive treatment
- Hadronic resonances: the most important region to set the limit is also the most complex and not modeled in our analysis
- Other environments?
- What if the target is not the RvB? Upscattering DM

Thank you for your attention!

Backup slides

Effective distance D_{eff}

A measure of the effective distance traveled by Cosmic Rays through an equivalent homogeneous environment

$$D_{eff} = \mathcal{B}c\tau_{esc}$$

where

$$\mathcal{B} = \int d^3 \vec{r} \ f_{CR}(\vec{r}) \frac{\eta(\vec{r})}{\overline{\eta}}$$

Energy dependence of τ_{esc}

Trapping time τ_{esc} depends on the energy *E* of the particle and on the magnetic field *B* as⁶

$$\tau_{esc} \sim 1 \text{ Gyr} \times \begin{cases} (ZeB/E)^{1/3} & \text{if } ZeB/E < lc \\ \\ (ZeB/E)^2 & \text{if } ZeB/E > lc \end{cases}$$



Pauli blocking

vs with momentum higher than $p_{\rm esc} = m_{\rm v} v_{\rm esc}$ are not gravitationally bound to the cluster. There are therefore

$$N = \frac{1}{8} \left(\frac{4\pi}{3} p_{\rm esc}^3 \right) \frac{V}{\pi^3}$$

available states. NFW profile, compute $v_{\rm esc}(r) = \sqrt{\frac{2GM(< r)}{r}}$, thus $\eta(r)$ and average it.

Allowed overdensity scales as $m_{
m v}^3$ and as $M_{
m vir}{}^7$

^{7.} Ringwald, Wong (2004) hep-ph/0408241

Cluster mass limit

NFW:

$M_{\rm vir} \propto r_{\rm vir}^3$

so average density is (almost) the same for all halos, $\overline{
ho}_{
m halo}$

We impose

 $\overline{\eta}_{\nu} n_{\nu}^0 \sum_i m_i < \overline{\rho}_{\text{halo}}$

Cosmic-Rays composition

