

Heavy flavour production in $p\bar{p}$ and AA collisions at the LHC

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work done in collaboration with:

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Outline

- Heavy quark production in pQCD: problems and theoretical tools
- Our approach: a Montecarlo setup with a **pQCD event generator** (POWHEG, within the POWHEG-BOX set-up) followed by an implementation of **Langevin diffusion** of heavy quarks in QGP.
- Further refinements (especially in the pp sector) with respect to version of QM2011.
- Predictions for pp collisions at LHC: $\sqrt{s}=7 \text{ TeV}$ and $\sqrt{s}=2.76 \text{ TeV}$
- Predictions for Pb-Pb collisions at LHC: $\sqrt{s}=2.76 \text{ TeV/n}$

Heavy flavour production in pQCD

The large mass M of c and b quarks makes possible a pQCD calculation of $Q\bar{Q}$ production:

- it sets a minimal “off-shellness” of the intermediate propagators (diagrams don’t explode);
- it sets a hard scale for the evaluation of $\alpha_s(\mu)$ (fastening of the convergence of the perturbative series);
- it prevents collinear singularities (suppression of emission of small-angle gluon, the so-called dead cone effect)

Both the total cross section $\sigma_{Q\bar{Q}}^{tot}$ and the invariant single-particle spectrum $E(d\sigma_Q/d^3p)$ are well-defined quantities which can be calculated in pQCD.

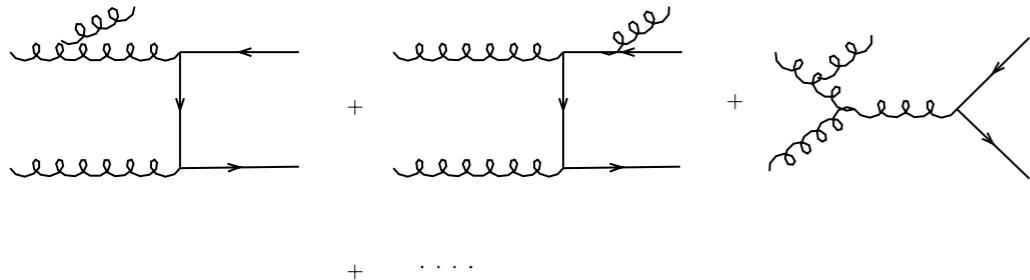
Factorization scheme (large Q^2):

$$\int G_{a/A}(x_a) G_{b/B}(x_b) \hat{\sigma}_{ab \rightarrow c\bar{c}}(\hat{s} = x_a x_b s) | D_{D/c}(z) = \sigma_{AB \rightarrow DX}$$

pQCD partonic cross section

Next to Leading Order (NLO) processes

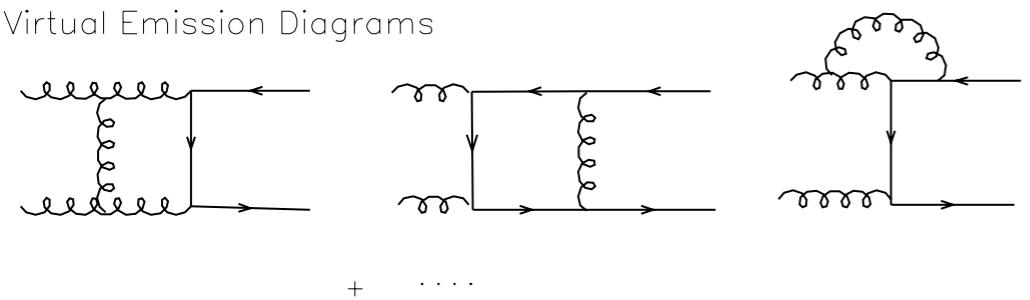
Real Emission Diagrams



- Real emission

$$|\mathcal{M}_{\text{real}}|^2 \sim \mathcal{O}(\alpha_s^3)$$

Virtual Emission Diagrams



- Virtual corrections

$$2\text{Re}\mathcal{M}_0\mathcal{M}_{\text{virt}}^* \sim \mathcal{O}(\alpha_s^3)$$

- Calculation at NLO accuracy gives the $\mathcal{O}(\alpha_s^3)$ result for $\sigma_{Q\bar{Q}}^{\text{tot}}$ and $E(d\sigma_Q/d^3 p)$
- It has been implemented for many years in several numerical codes (as MNR and others) and Montecarlo event generators (as MC@NLO, POWHEG, etc)
- However... **large terms of collinear origin** $\sim \alpha_s \ln(p_T/M)$ **can become large for ($p_T \gg m$)**
- Different schemes for **resummation** of these logs are possible

Some examples of pQCD tools for heavy flavour studies

- Fixed-order codes: **MNR** (calculation) or **MC@NLO, POWHEG** (MC generators).
- **FONLL**: fixed-order (NLO) calculation of hard processes + **next-to-leading log resummation** (with implementation also of non-perturbative input to model fragmentation and decays).
- alternative schemes, e.g. **GM-VFNS**

Montecarlo codes provide the advantage to save maximum information about the event (i.e. correlations in azimuth or rapidity between $Q\bar{Q}$)

POWHEG is a code interfaced to Shower Monte Carlo programs (like **HERWIG**, **PYTHIA**) that describe at the **Leading Log accuracy** how initial and final states partons evolve according to DGLAP.

Our approach: a multi-step simulation of heavy quark production in pp and AA

- Initial generation of QQ pairs with **POWHEG** (pQCD@NLO), with the addition *a posteriori* of an **intrinsic k_T kick**; added Parton Shower (**PYTHIA**)
- for AA collisions:** EPS09 nuclear corrections to parton distributions (both at NLO accuracy) have been implemented
- Heavy quark position are distributed in the transverse plane according to nuclear geometry (**Glauber**); Cronin effect (**k_T broadening**) included.
- Langevin evolution in the QGP:** at each step $u^\mu(x)$ and $T(x)$ are given by **hydro codes**, and used to evaluate transport coefficients of the expanding fluid and to update position and 4-momentum of the heavy quark.
- At T_c HQ are made **hadronize**. **Fragmentation** is performed by sampling hadron species from experimental branching-fractions, and by sampling momentum from appropriate **parametrizations of fragmentation functions**.
- Finally, heavy quark hadrons are made **decay into electrons**, by using the **PYTHIA decayer** with branching-ratios from Particle Data Group review.

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Relativistic Langevin equation

$$\frac{\Delta p^i}{\Delta t} = - \underbrace{\eta_D(p)p^i}_{\text{determ.}} + \underbrace{\xi^i(t)}_{\text{stochastic}}$$

with the properties of the noise encoded in

$$\langle \xi^i(\mathbf{p}_t) \xi^j(\mathbf{p}_{t'}) \rangle = b^{ij}(\mathbf{p}_t) \frac{\delta_{tt'}}{\Delta t} \quad b^{ij}(\mathbf{p}) \equiv \kappa_L(p) \hat{p}^i \hat{p}^j + \kappa_T(p) (\delta^{ij} - \hat{p}^i \hat{p}^j)$$

Transport coefficients to be calculated (HTL pQCD):

- **momentum diffusion** $\kappa_T \equiv \frac{1}{2} \frac{\langle \Delta p_T^2 \rangle}{\Delta t}$ and $\kappa_L \equiv \frac{\langle \Delta p_L^2 \rangle}{\Delta t}$

- **friction term** (dependent on the **discretization scheme!**):

$$\eta_D^{\text{Ito}}(p) = \frac{\kappa_L(p)}{2TE_p} - \frac{1}{E_p^2} \left[(1-v^2) \frac{\partial \kappa_L(p)}{\partial v^2} + \frac{d-1}{2} \frac{\kappa_L(p) - \kappa_T(p)}{v^2} \right]$$

fixed in order to insure approach to equilibrium (**Einstein relation**):

Langevin \Leftrightarrow Fokker Planck with steady solution $\exp(-E_p/T)$

Hydrodynamics

The fields $u^\mu(x)$ and $T(x)$ are taken from the output of two longitudinally boost-invariant (“Hubble-law” longitudinal expansion $vz = z/t$) hydro codes ^(1,2):

$$x^\mu = (\tau \cosh \eta, \mathbf{r}_\perp, \tau \sinh \eta) \quad \text{with} \quad \tau \equiv \sqrt{t^2 - z^2}$$
$$u^\mu = \bar{\gamma}_\perp (\cosh \eta, \bar{\mathbf{v}}_\perp, \sinh \eta) \quad \text{with} \quad \bar{\gamma} \equiv \frac{1}{\sqrt{1 - \bar{\mathbf{v}}_\perp^2}}$$

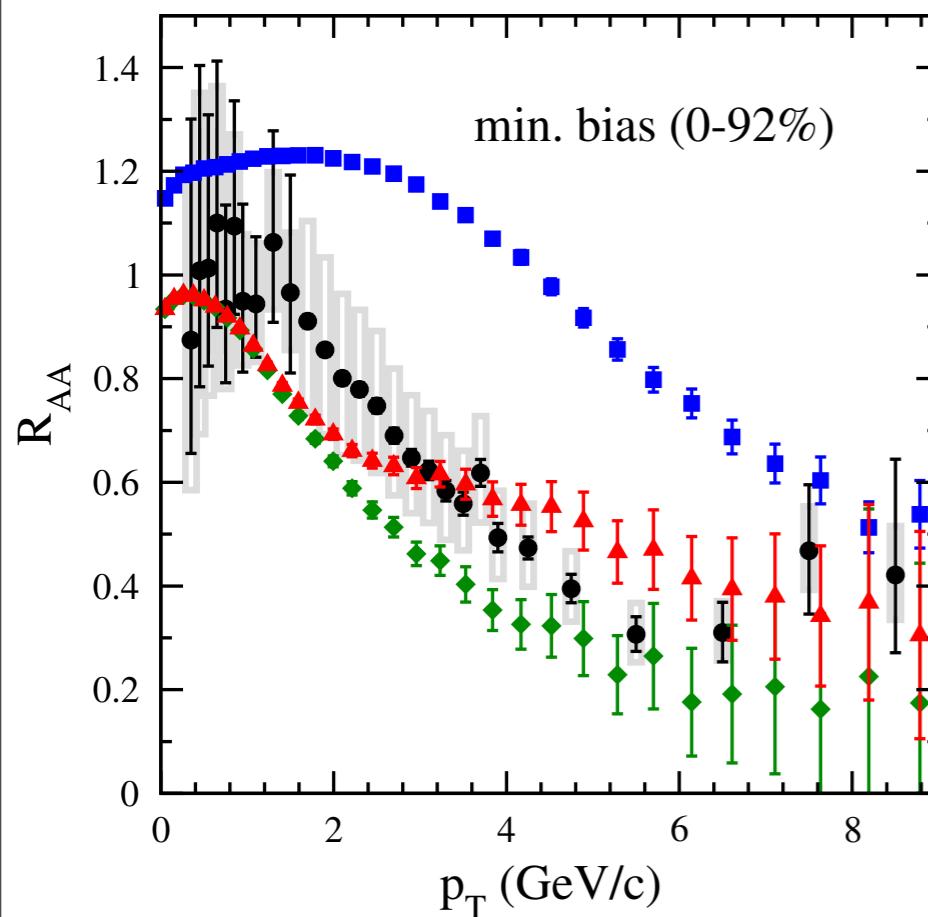
- $u^\mu(x)$ used to perform the update each time in the fluid rest-frame
- $T(x)$ allows to fix at each step the value of the transport coefficients.

[1] P.F. Kolb, J. Sollfrank and U. Heinz, Phys. Rev. C 62 (2000) 054909

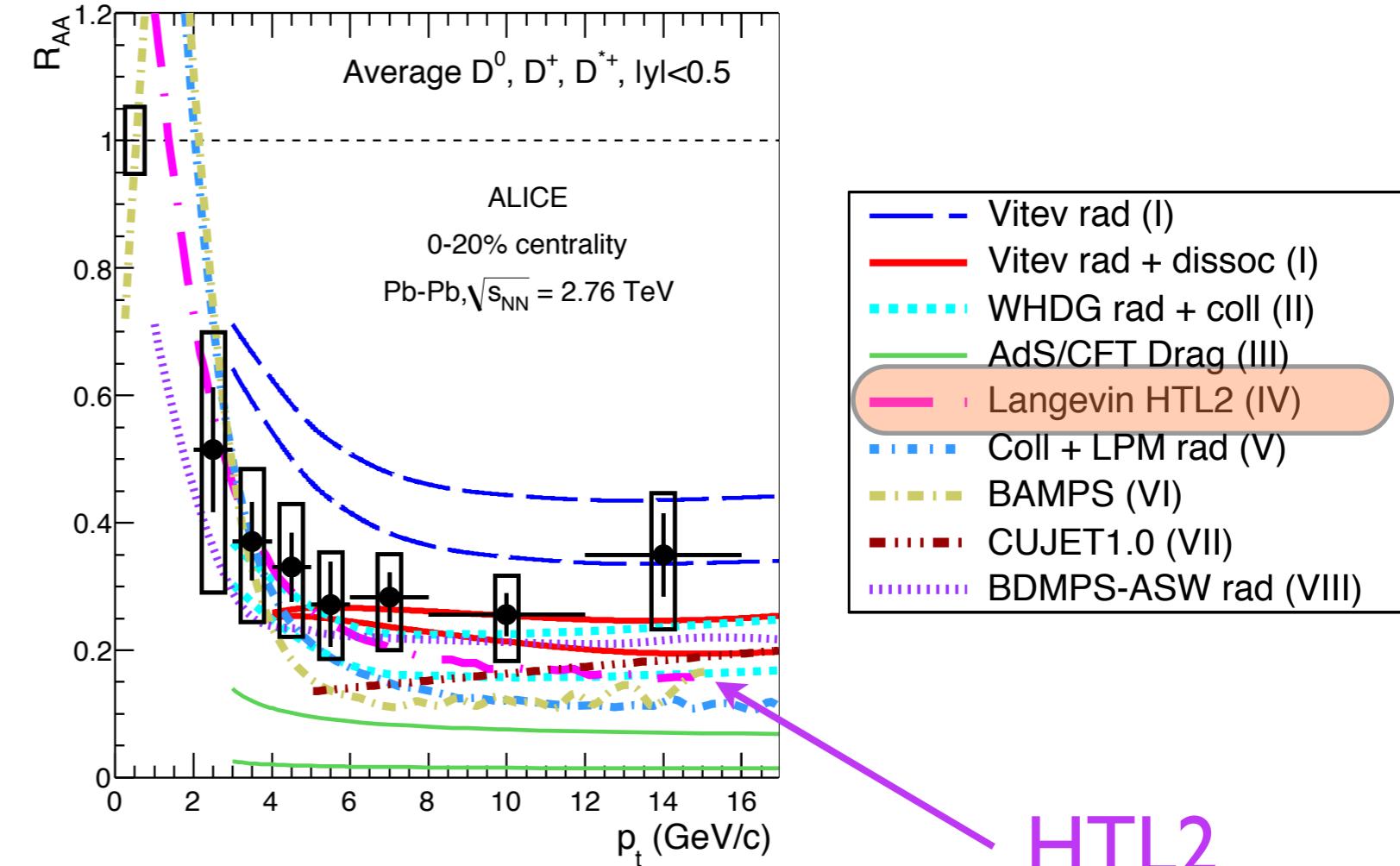
[2] P. Romatschke and U. Romatschke, Phys. Rev. Lett. 99 (2007) 172301

First results (released at QM2011)

PHENIX HFE



ALICE D mesons



- missing Parton Shower evolution (ISR and FSR)
- obsolete FF (Peterson)
- output files have poor information for more refined analysis (i.e. rapidity or angular correlations)
- needed study of systematic uncertainties

but...

Fragmentation Functions

Functional forms from Braaten et al. arXiv:hep-ph/9409316
with different functional forms for pseudoscalar and vector mesons:

$$D_{Q \rightarrow P}(z) = N \frac{rz(1-z)^2}{(1-(1-r)z)^6} [6 - 18(1-2r)z + (21-74r+68r^2)z^2 - 2(1-r)(6-19r+18r^2)z^3 + 3(1-r)^2(1-2r+2r^2)z^4], \quad (31)$$

$$D_{Q \rightarrow V}(z) = 3N \frac{rz(1-z)^2}{(1-(1-r)z)^6} [2 - 2(3-2r)z + 3(3-2r+4r^2)z^2 - 2(1-r)(4-r+2r^2)z^3 + (1-r)^2(3-2r+2r^2)z^4]. \quad (32)$$

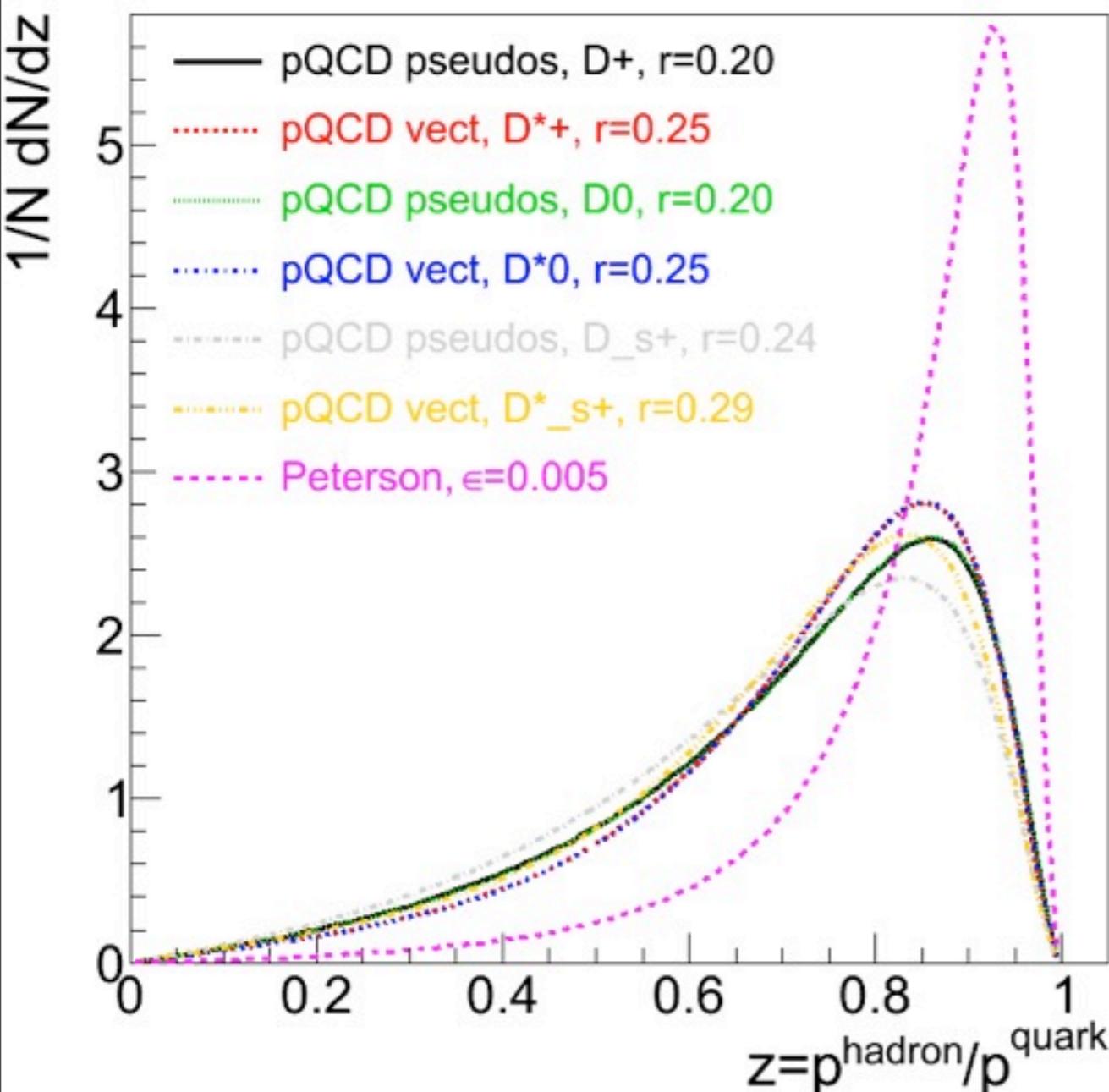
z=p_H / p_Q

One parameter: r (but different authors adopt different values)

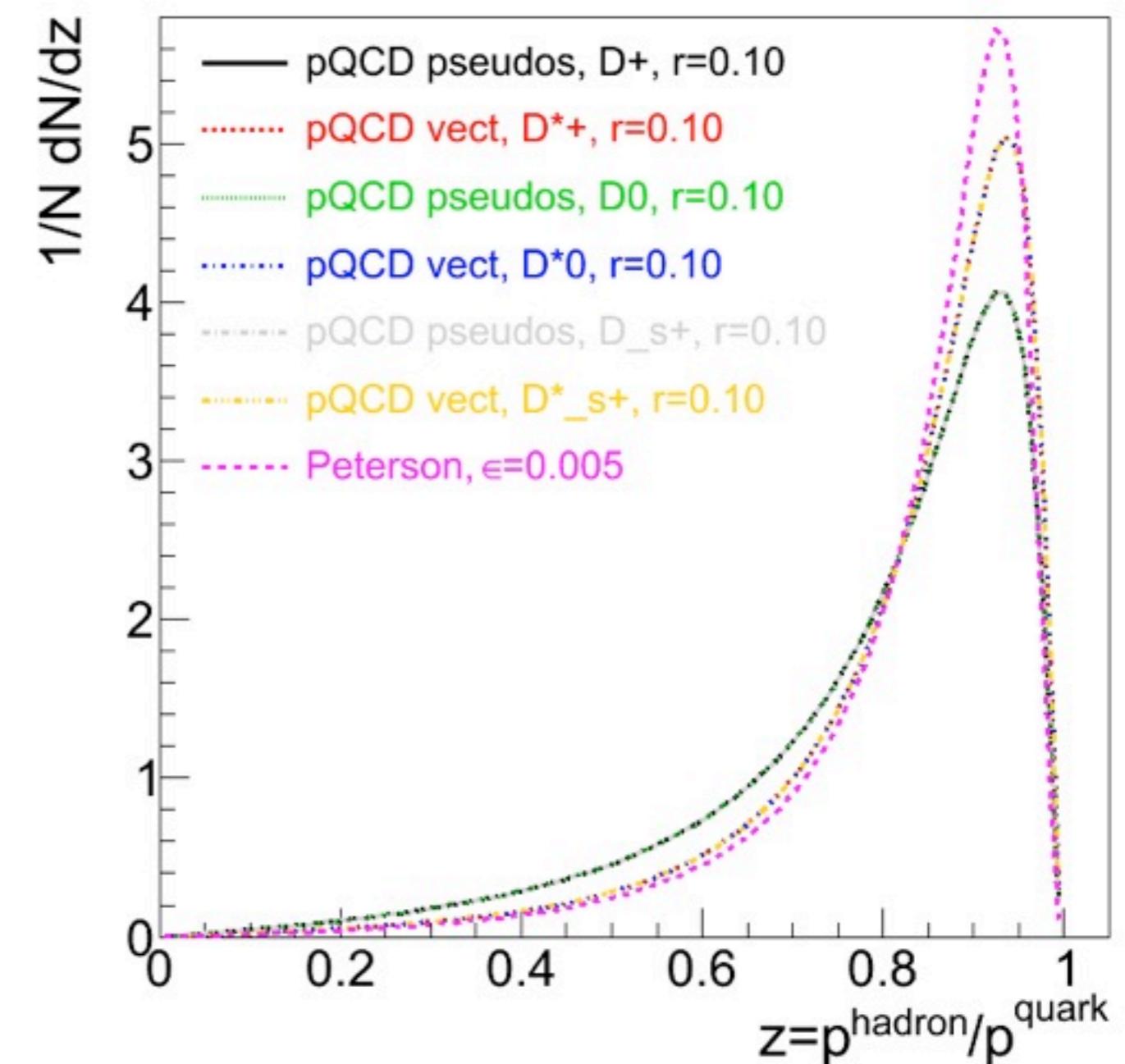
- in **Braaten et al.**: $r=(m_H-m_Q)/m_H$ corresponding to the contribution of light quark to hadron mass (e.g. with **m_Q=1.5 GeV, r=0.2 for D⁰ and D⁺**)
- In **FONLL** (arXiv:hep-ph/0502203 and arXiv:hep-ph/0306212):
r=0.1 for m_c=1.5 GeV (fitted on D* ALEPH data)

Which is the effect on syst.uncertainty of calculations?

Comparison between different FF (charm)

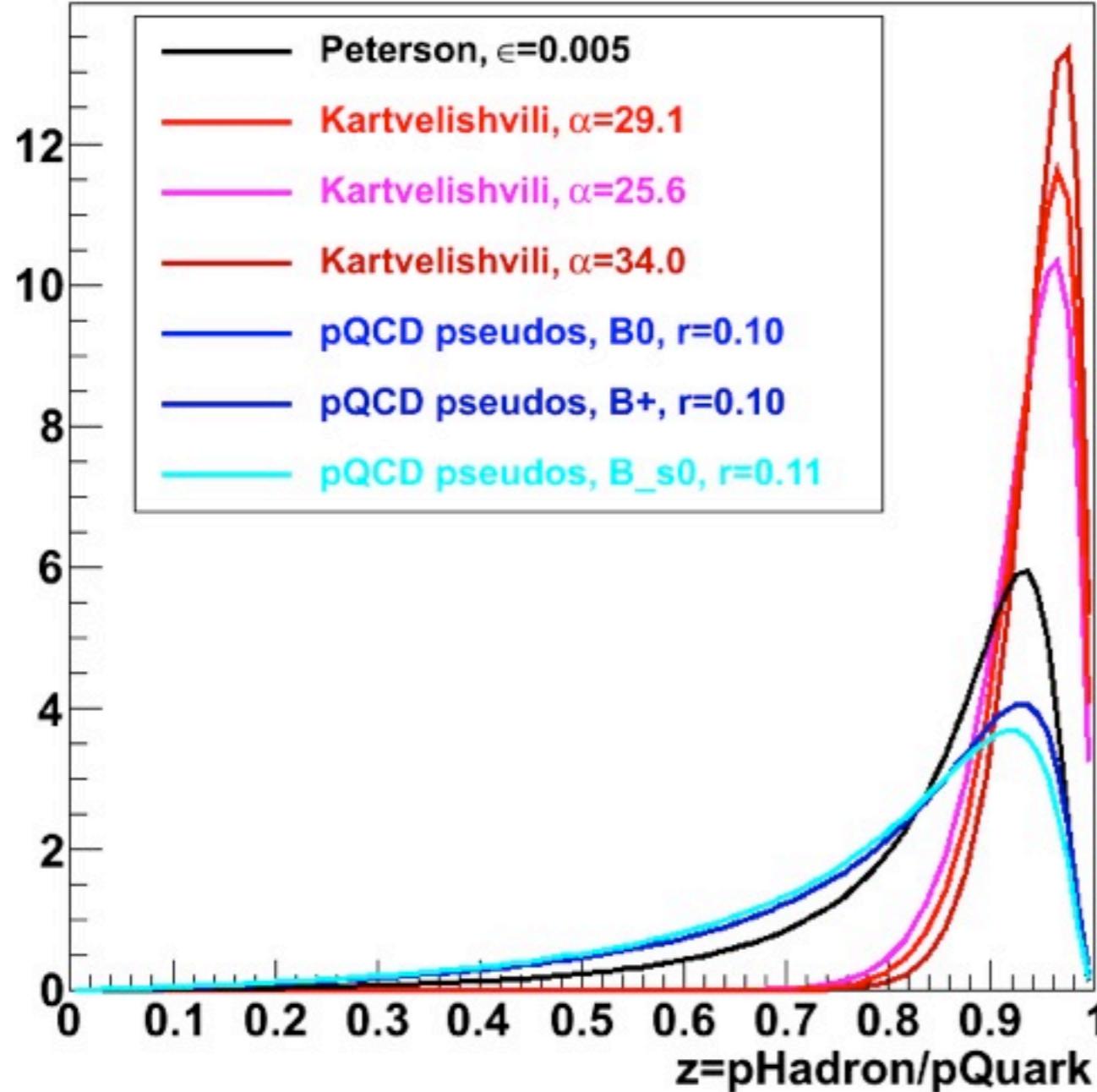


Braaten et al.



FONLL

Comparison between different FF (bottom)



Kartvelishvili et al: FONLL
pQCD: Braaten et al

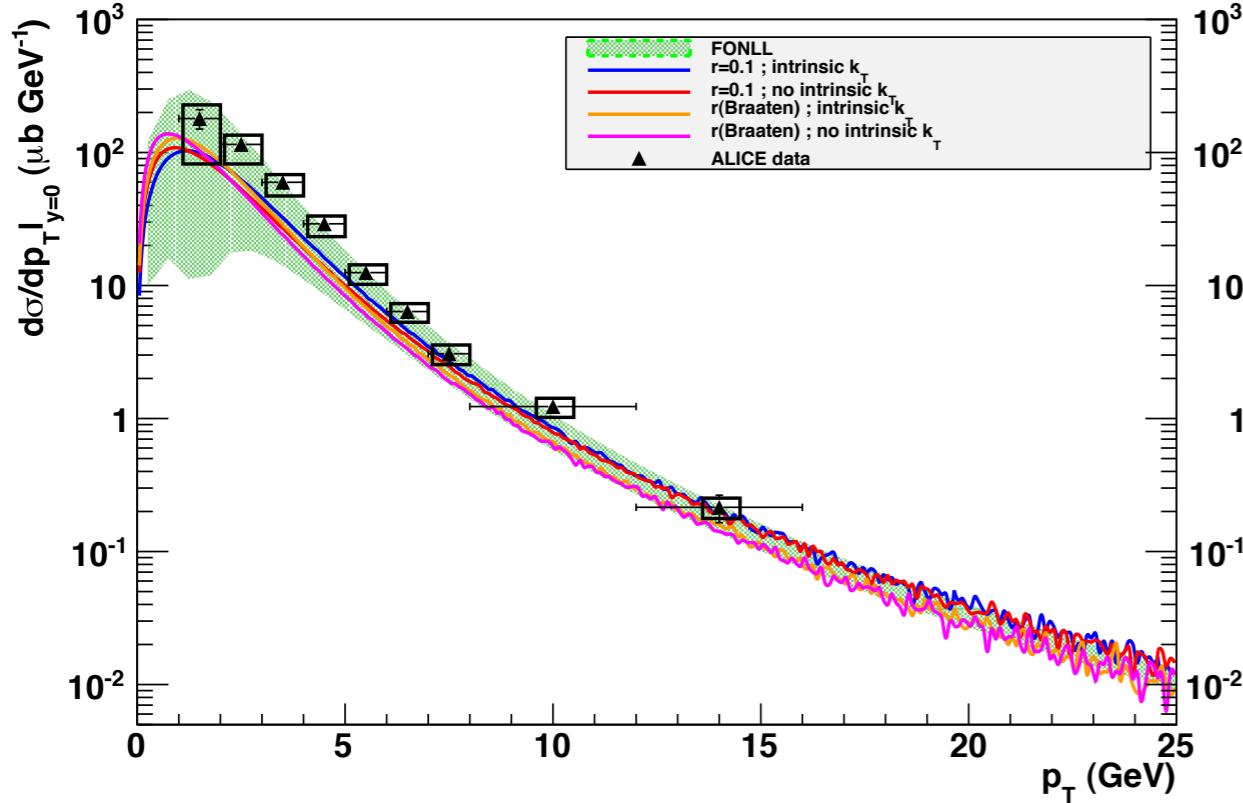
Some predictions for pp collisions at the LHC

p_T differential cross-sections

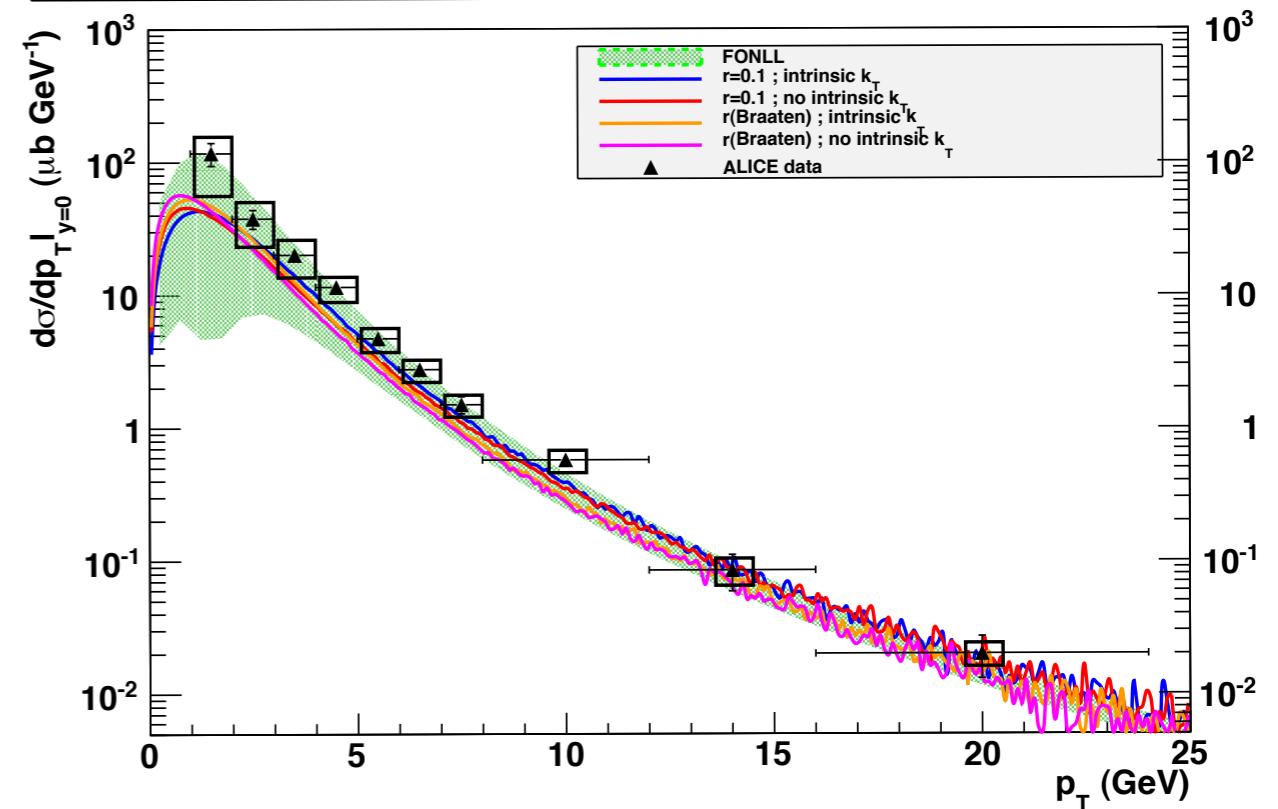
- prompt D, D^* mesons - ALICE $|y| < 0.5$
- B mesons - CMS $|y| < 2.2 - 2.4$

p_T spectra of D mesons in ALICE: $\sqrt{s} = 7$ TeV

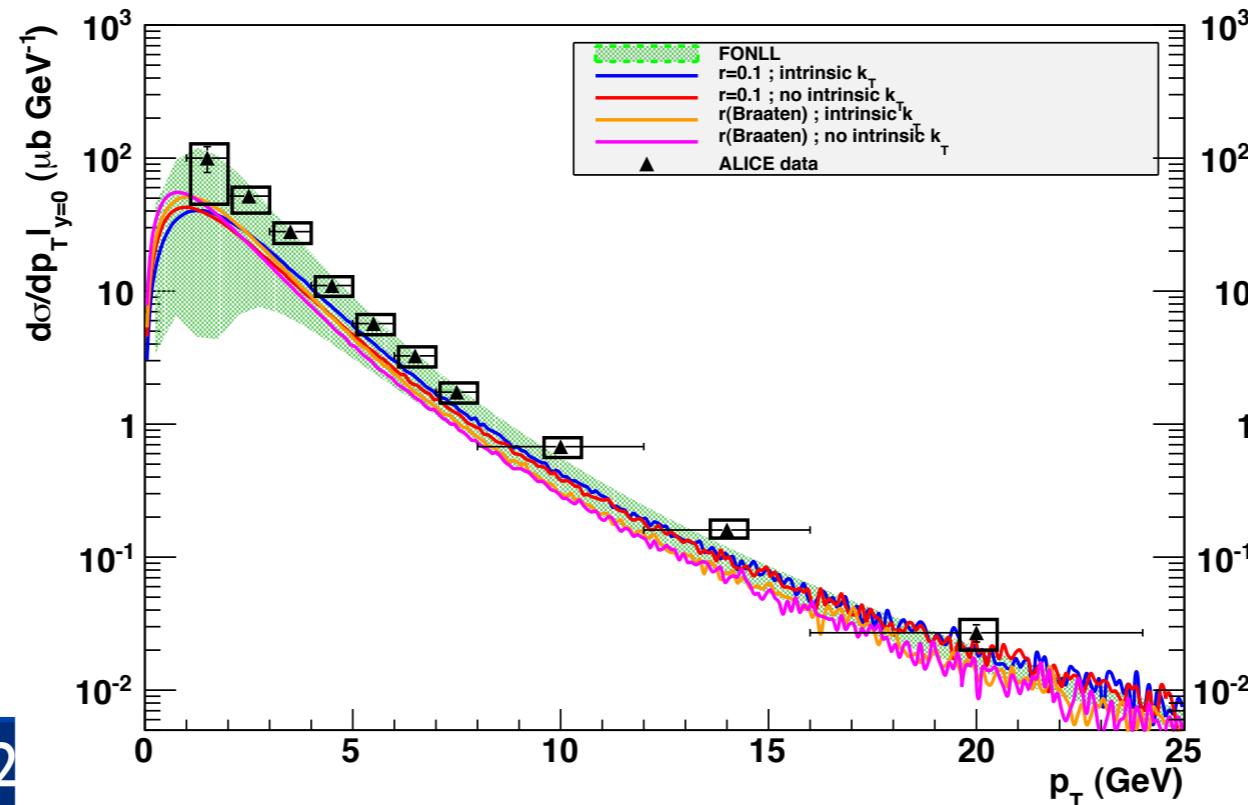
D^0 in pp at LHC ($\sqrt{s}=7$ TeV, $|y|<0.5$): POWHEG-BOX+PYTHIA Parton Shower



D^+ in pp at LHC ($\sqrt{s}=7$ TeV, $|y|<0.5$): POWHEG-BOX+PYTHIA Parton Shower

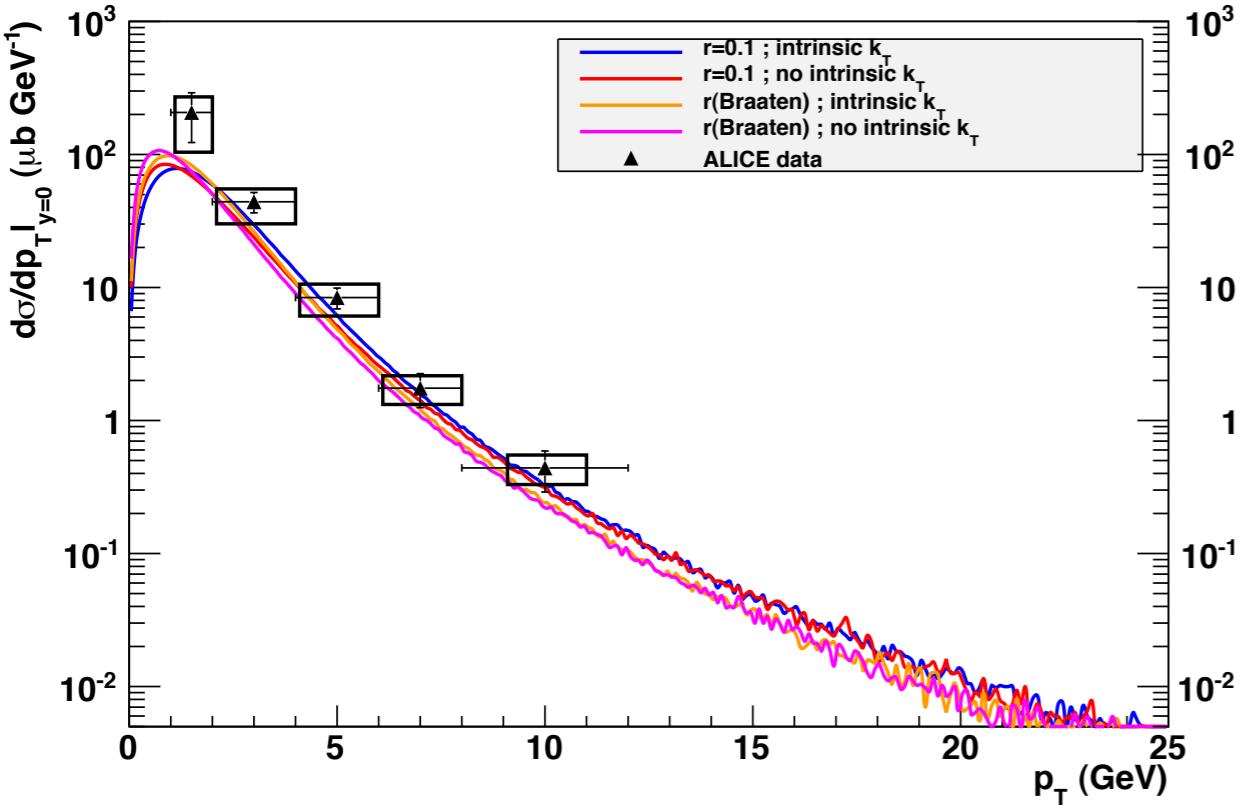


D^{*+} in pp at LHC ($\sqrt{s}=7$ TeV, $|y|<0.5$): POWHEG-BOX+PYTHIA Parton Shower

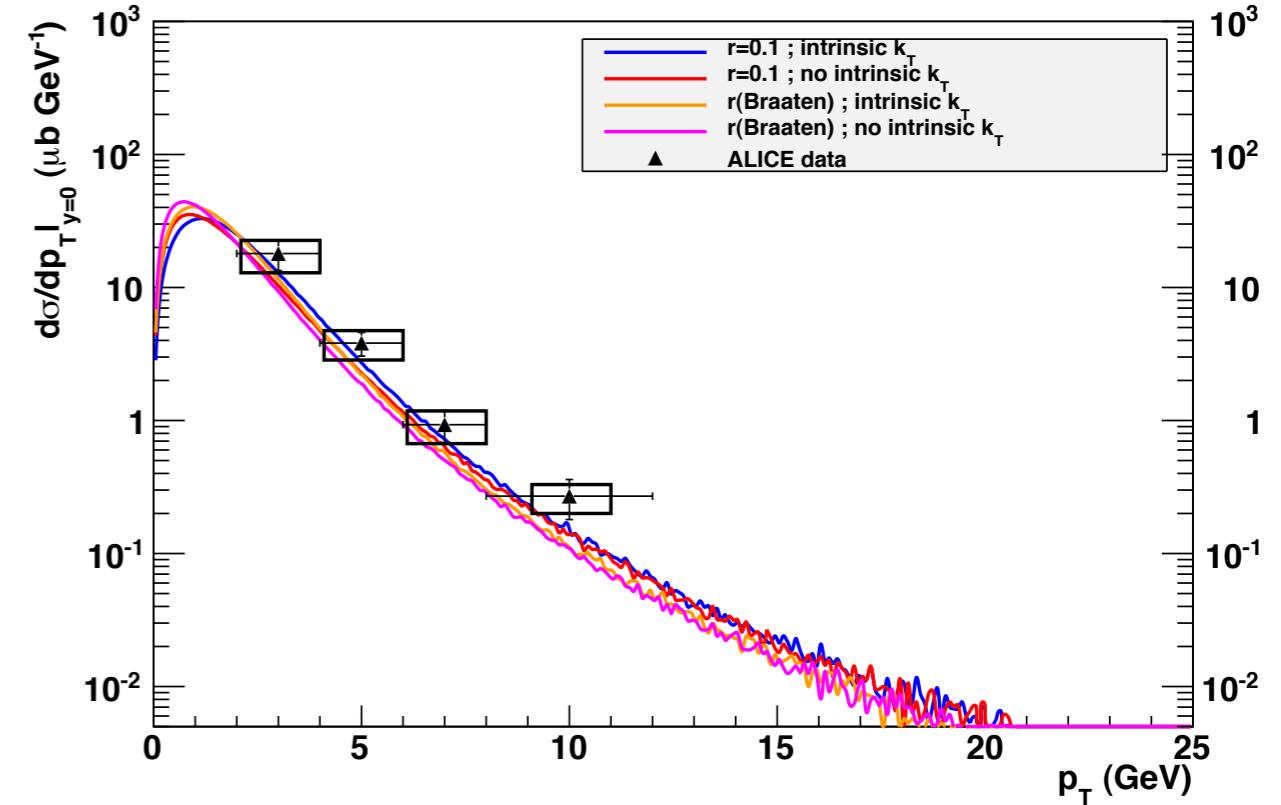


p_T spectra of D mesons in ALICE: $\sqrt{s} = 2.76$ TeV

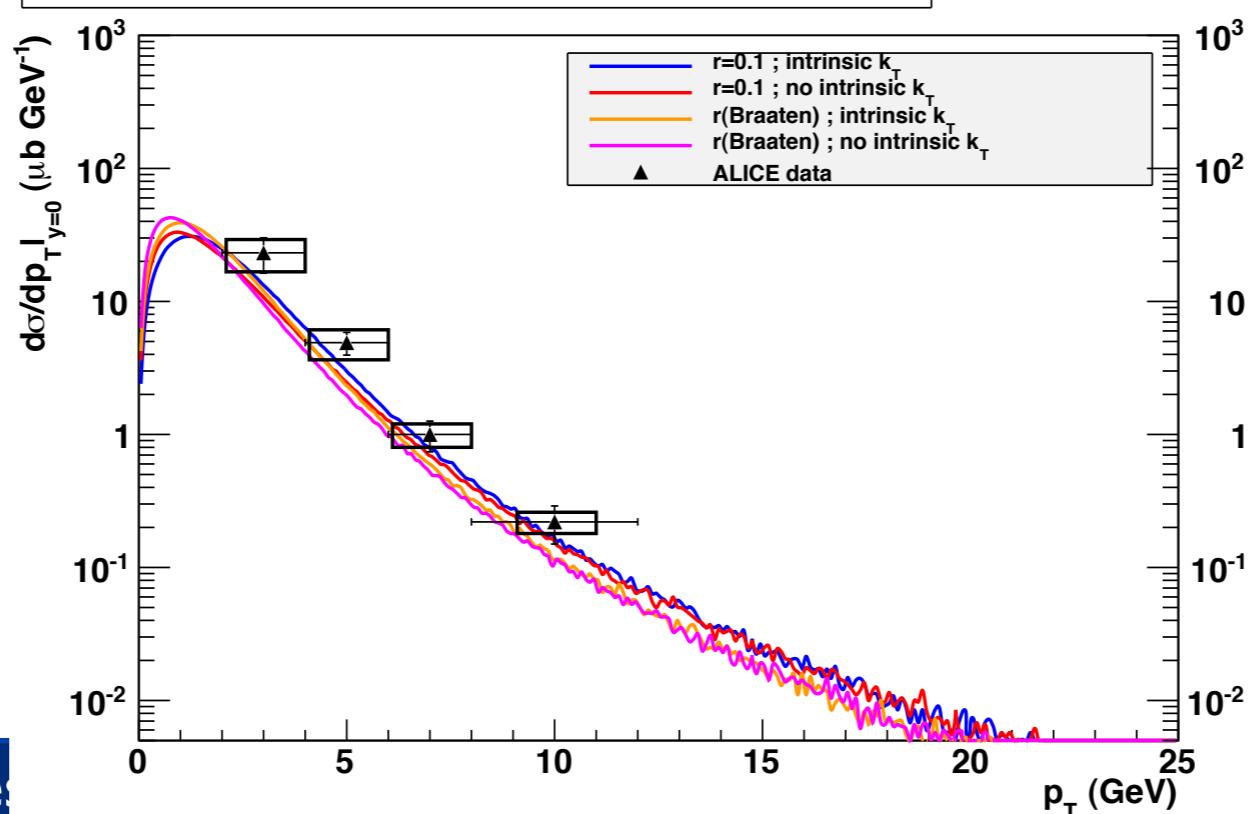
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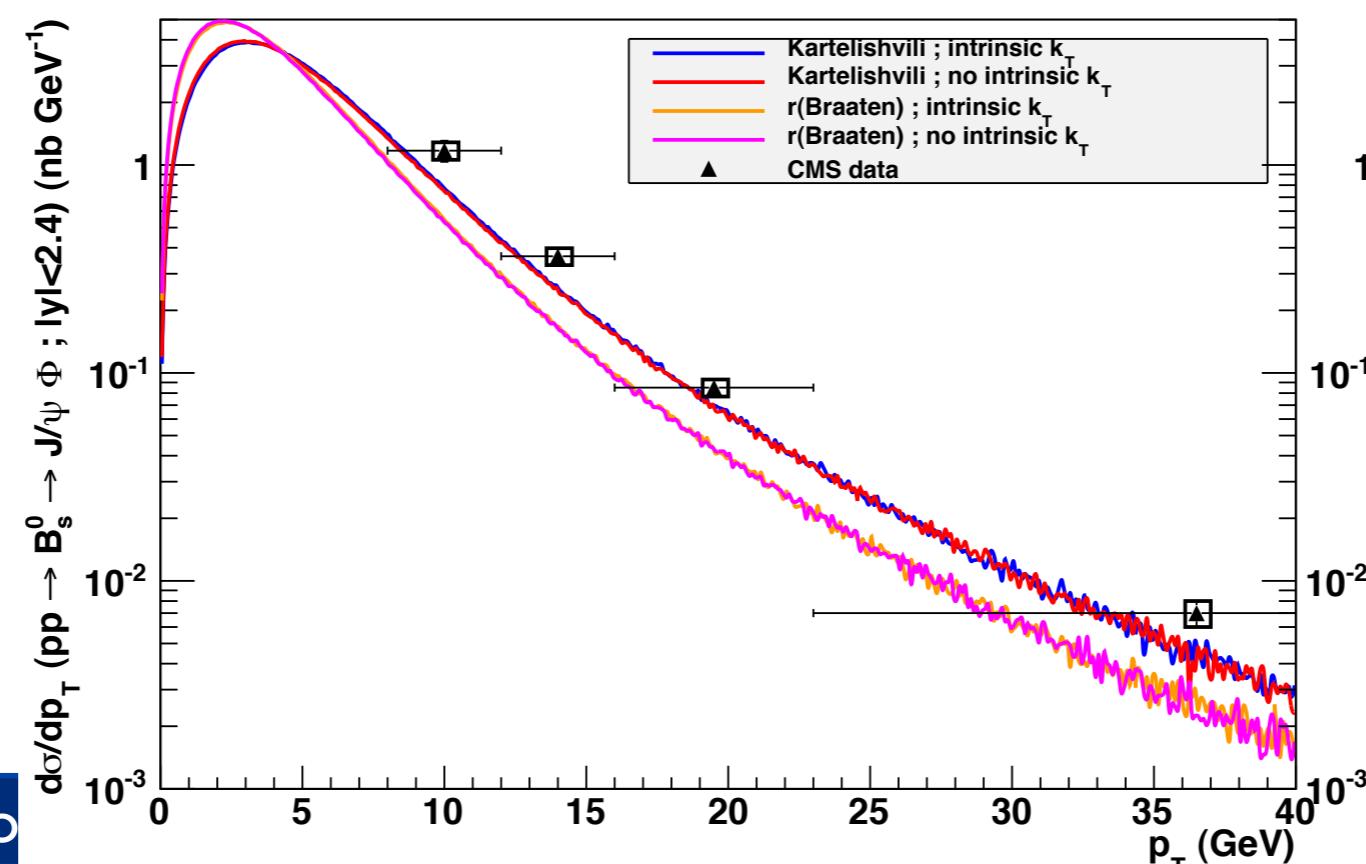
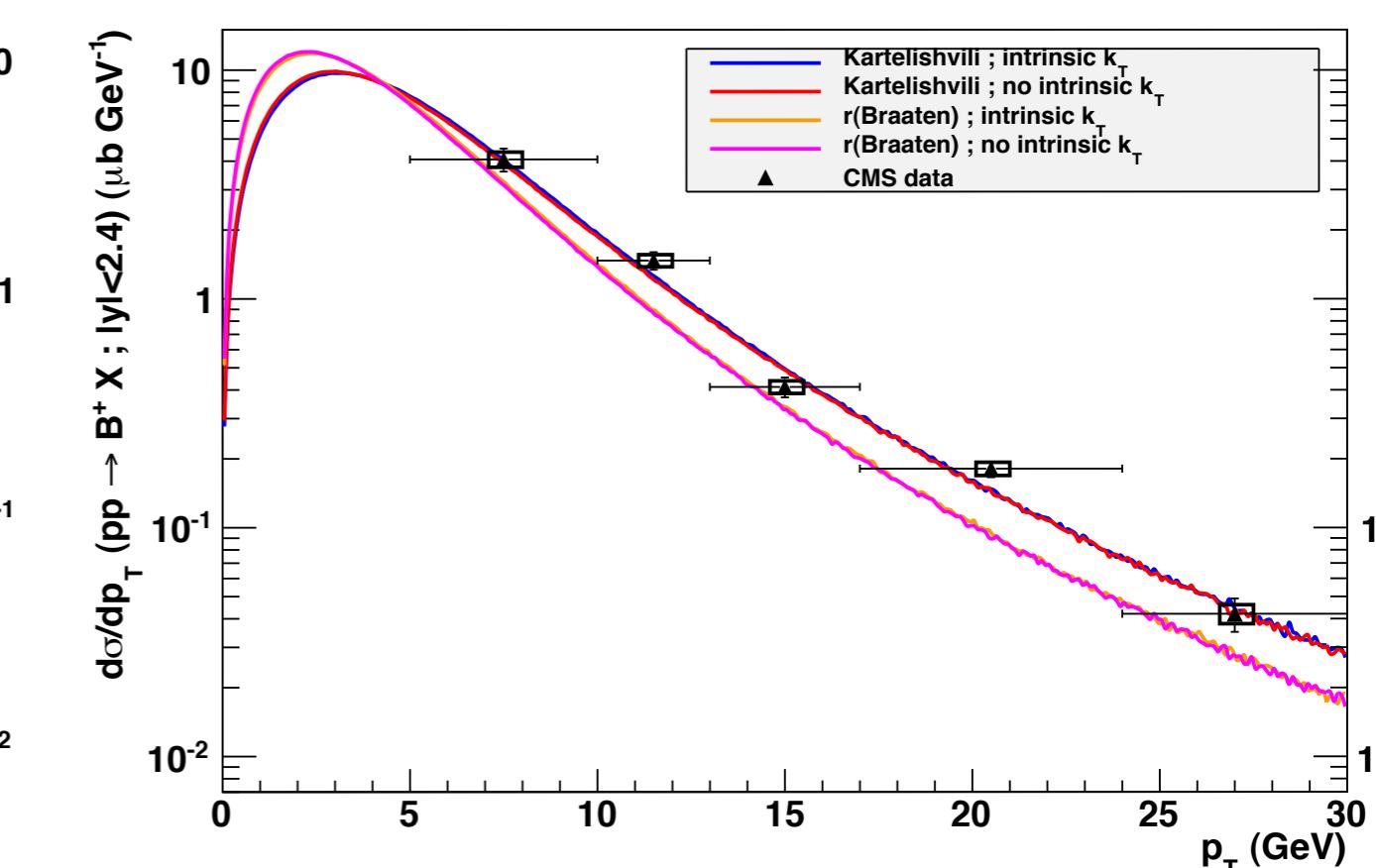
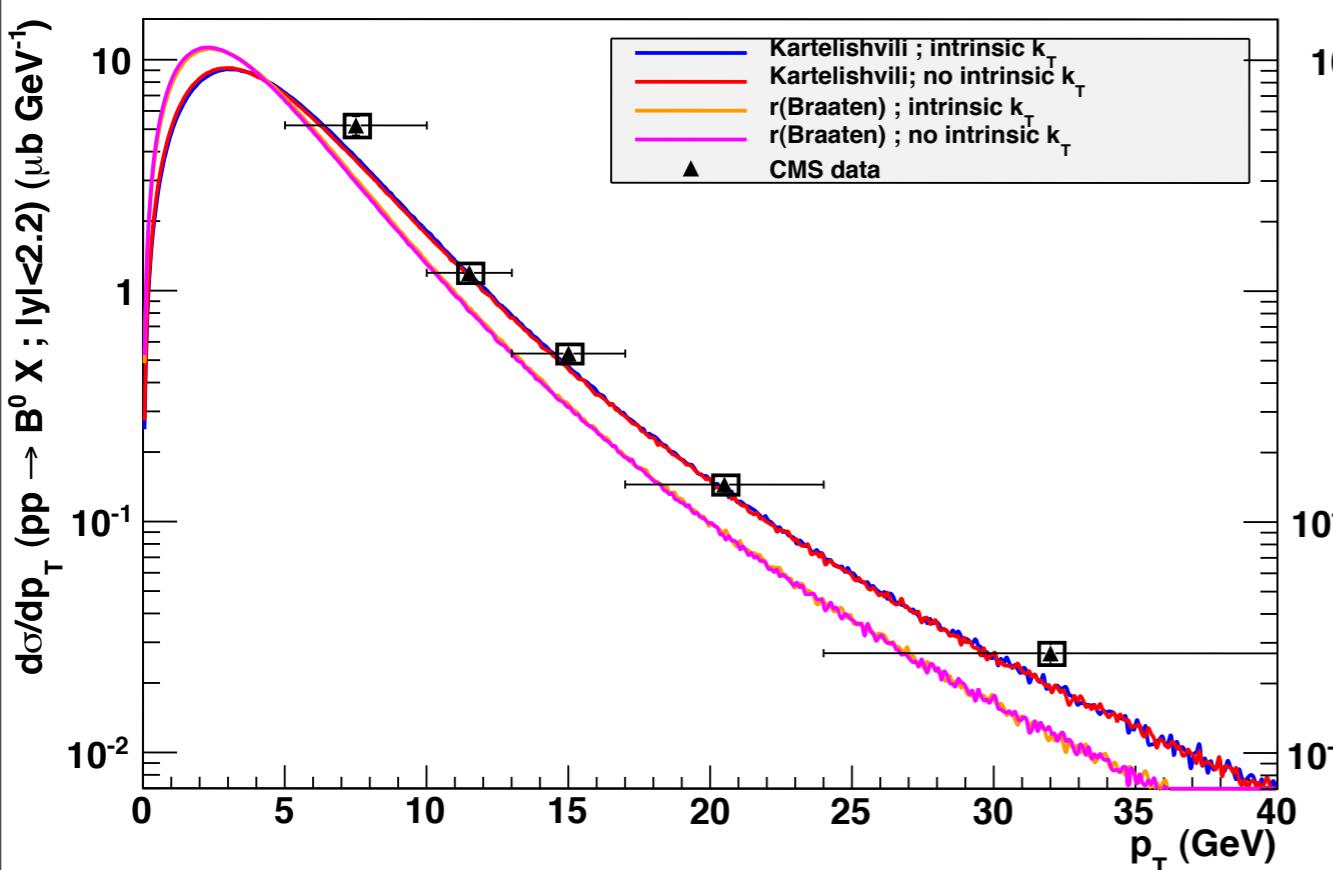
D^+ in pp at LHC ($\sqrt{s}=2.76$ TeV, $|y|<0.5$): POWHEG-BOX+PYTHIA Parton Shower



D^{*+} in pp at LHC ($\sqrt{s}=2.76$ TeV, $|y|<0.5$): POWHEG-BOX+PYTHIA Parton Shower



p_T spectra of B mesons in CMS $\sqrt{s} = 7$ TeV



Predictions for PbPb collisions at the LHC: $\sqrt{s} = 2.76$ TeV/nucleon

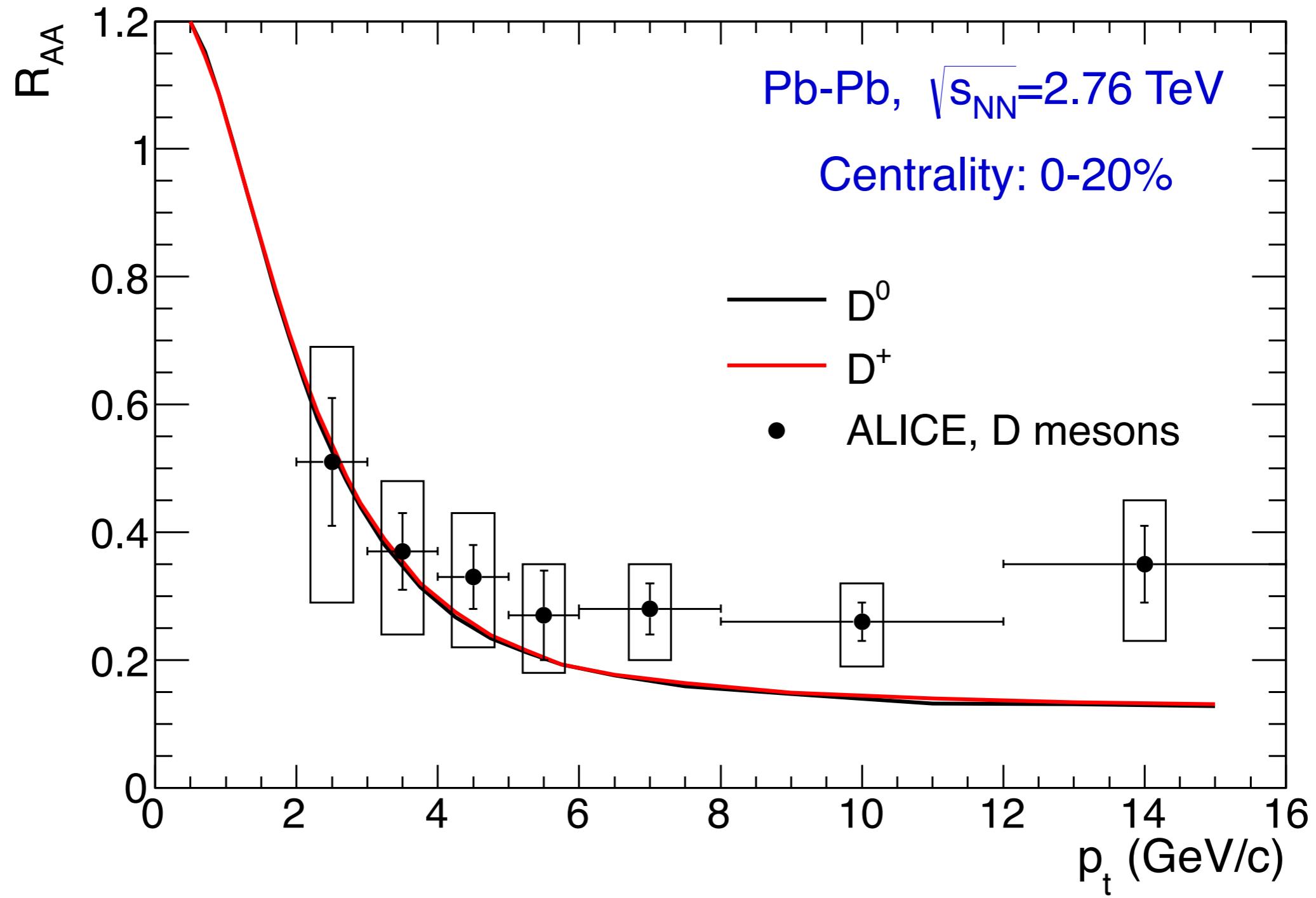
Nuclear modification factors R_{AA}

- R_{AA} D, D^* mesons in ALICE $|y| < 0.5$
centr. (0-20)%

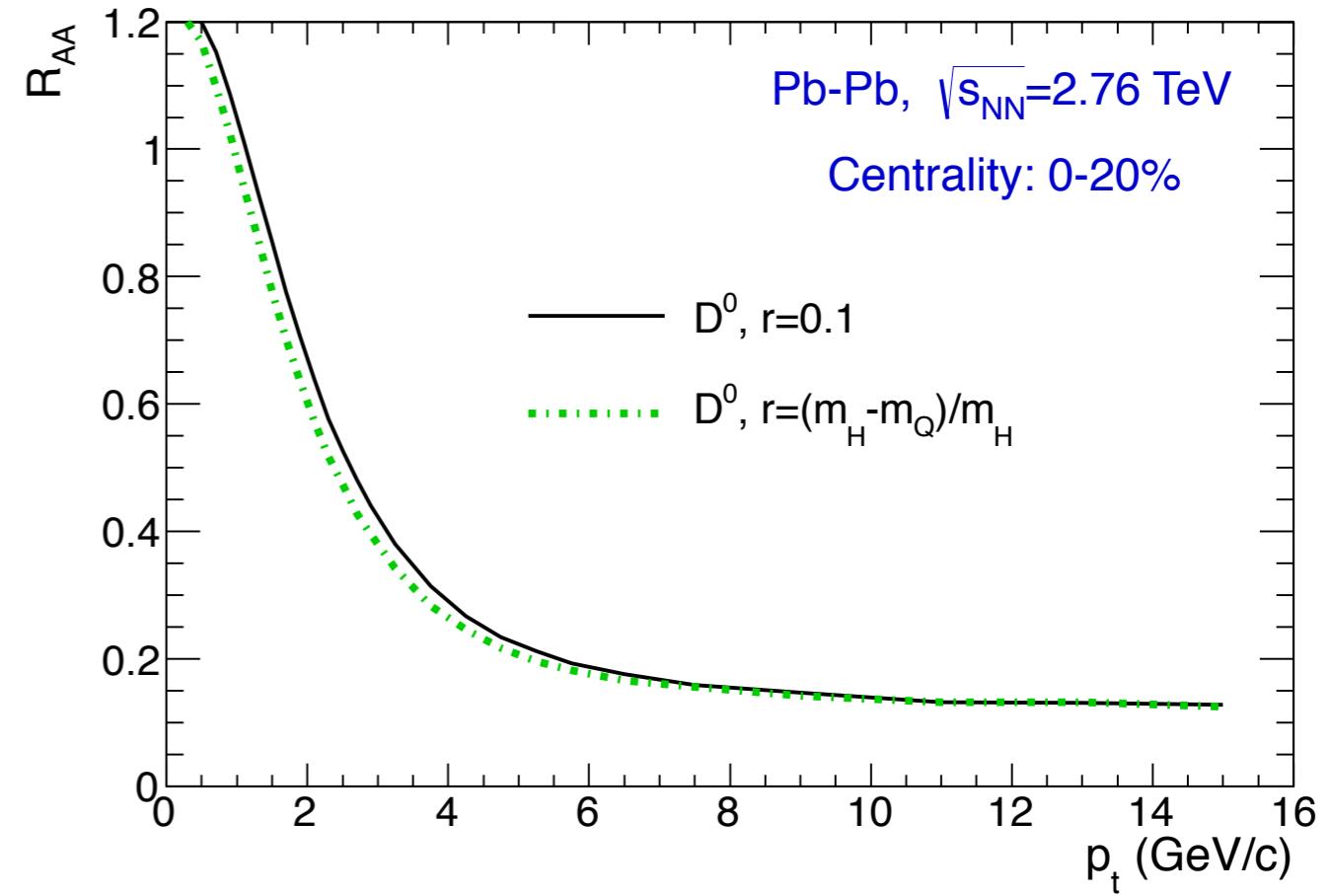
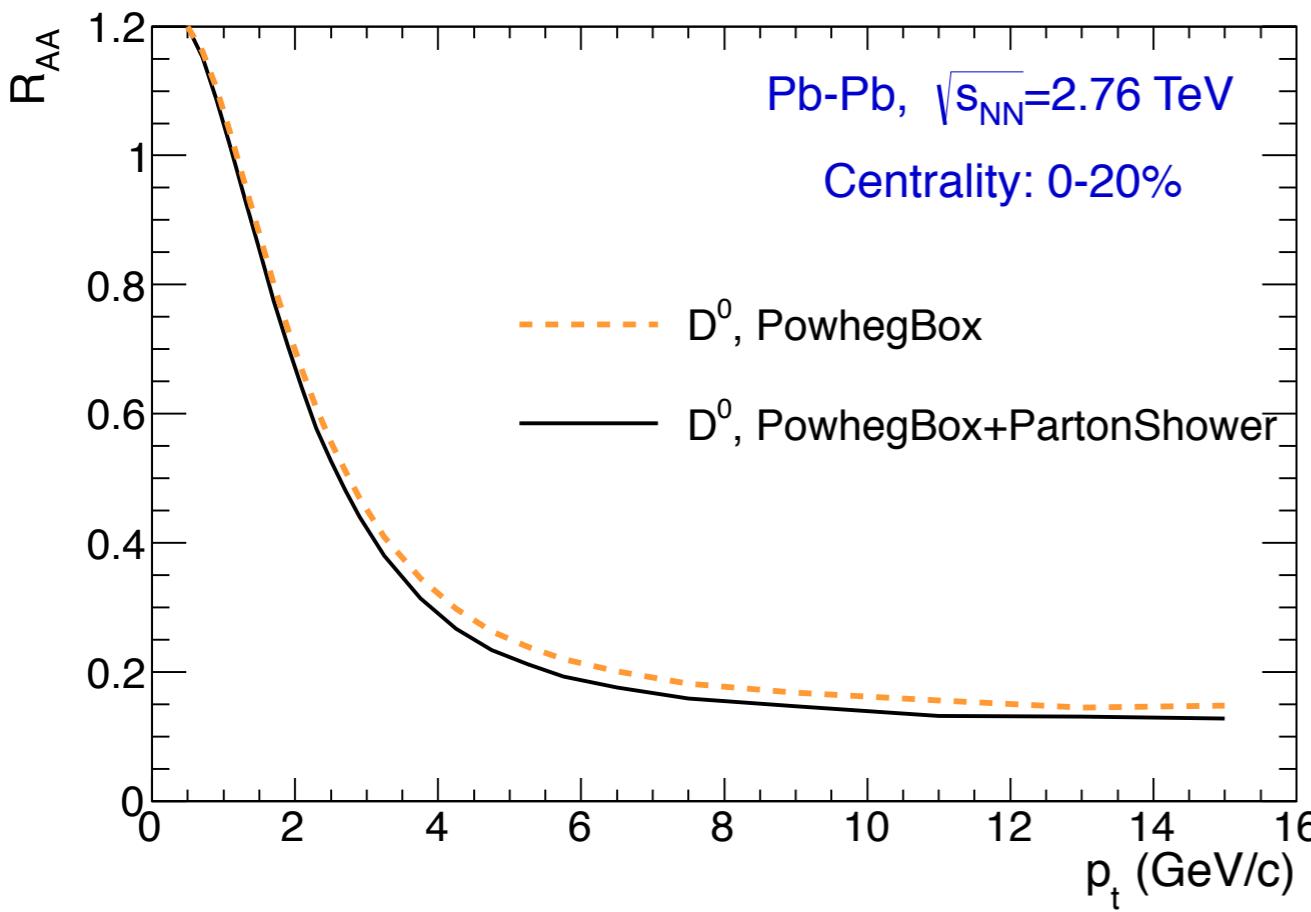
Elliptic flow v_2

- v_2 of D mesons in ALICE $|y| < 0.5$, centr. (30-50)%

R_{AA} of D mesons in ALICE

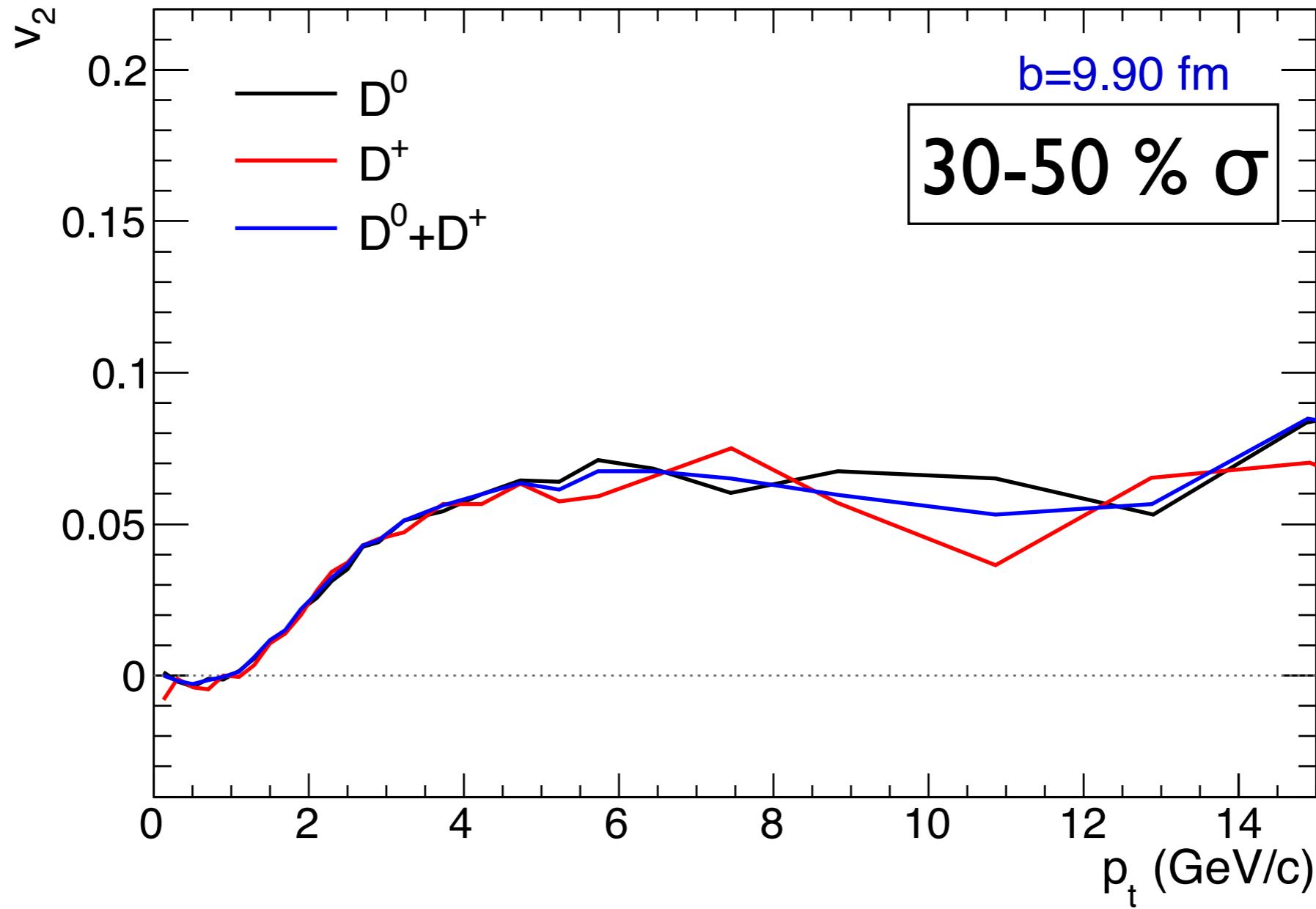


R_{AA} of D mesons in ALICE: systematics

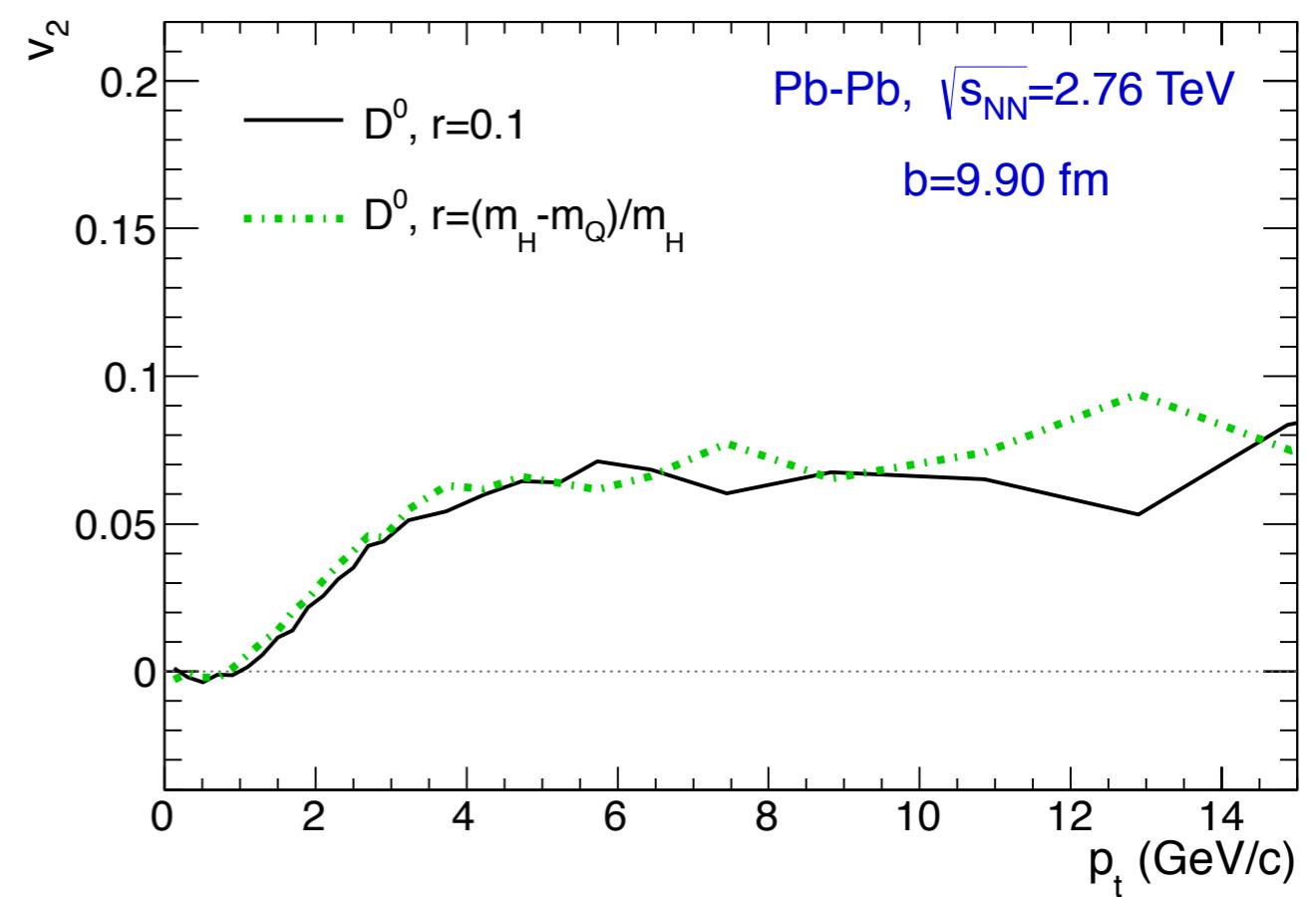
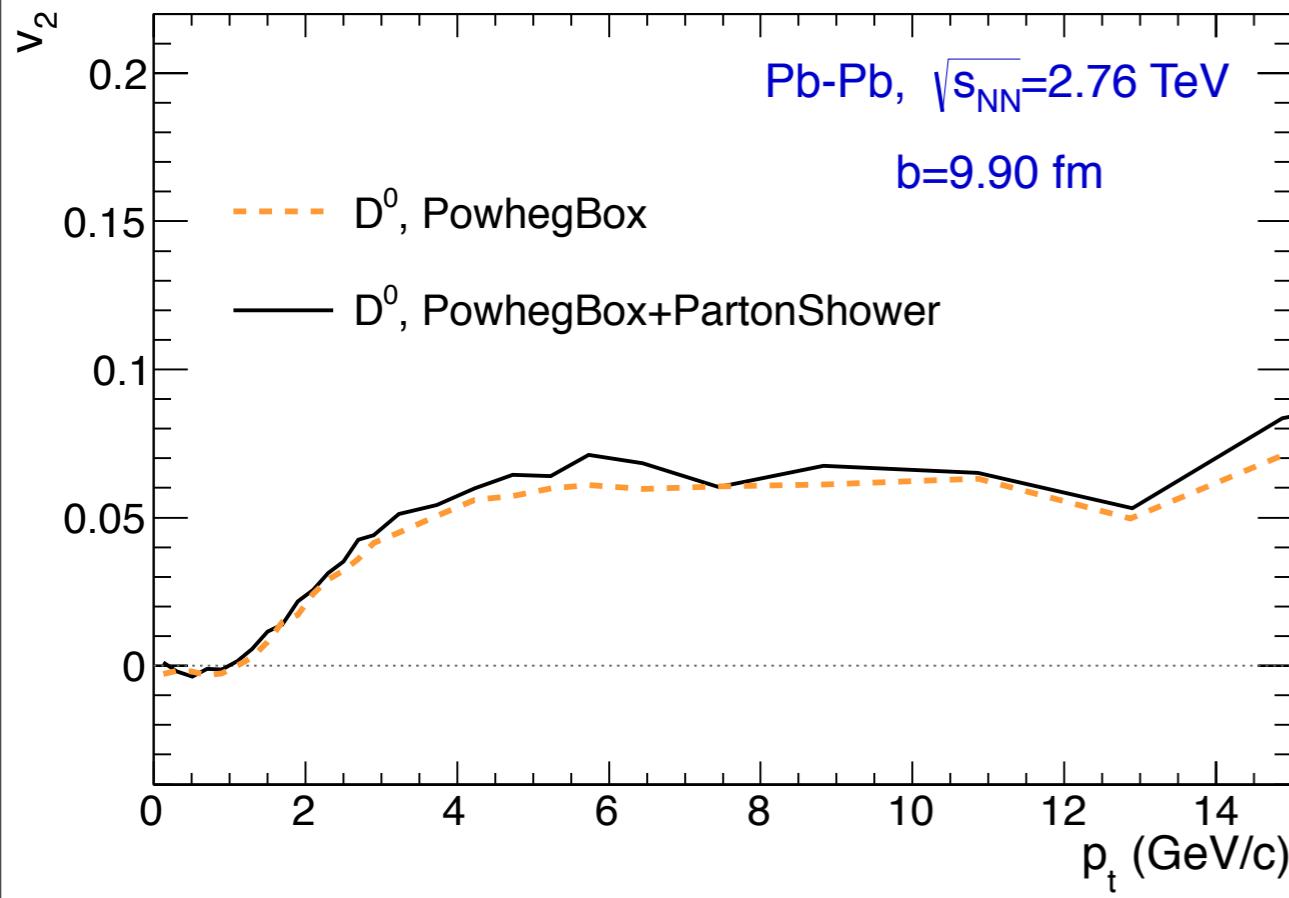


Comparison of changes in R_{AA} with/without PYTHIA Parton Shower
and for changes of the r parameter in the charm Braaten fragmentation function

Elliptic flow v_2 in ALICE



v_2 of D^0 mesons in ALICE: systematics



Comparison of changes in v_2 (D^0) with/without PYTHIA Parton Shower
and for changes of the r parameter in the charm Braaten fragmentation
function

Summary

- Improvements to our simulation of HQ production in pp collisions have been implemented, by interfacing the most recent version of POWHEG with PYTHIA (to perform Parton Showering).
- Systematic studies on the effect of an intrinsic k_T , or a change in FF have been performed.
- First check against data from LHC experiments on pp collisions at $\sqrt{s} = 7$ TeV show a good level of agreement, taking into account the different sources of theoretical uncertainties. Our setup provides a solid benchmark to be used as a reference for results obtained in nucleus-nucleus collisions.
- Preliminary comparisons with ALICE data show a good agreement with the pattern of R_{AA} in central collisions, whereas the elliptic flow v_2 in semi-peripheral collisions appears to be underestimated.

BACKUP SLIDES

Evaluation of transport coefficients $\kappa_T/\kappa_L(p)$

The interaction rate (from the squared matrix element of the process) must be weighted by the squared transverse/longitudinal exchanged momentum.

intermediate cutoff $|t|^* \sim m_D^2$ introduced to separate the contributions of

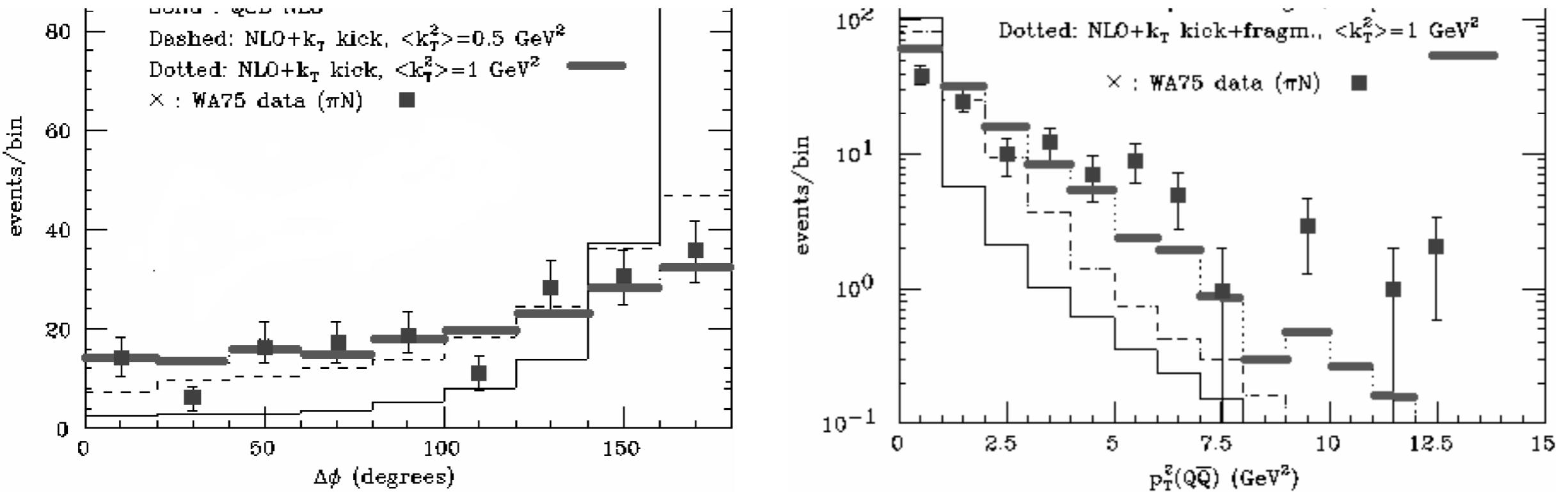
- **soft collisions** ($|t| < |t|^*$): Hard Thermal Loop (HTL) approximation in a weak-coupling scenario, with the running coupling constant $g(\mu)$ taken at a scale $\mu \sim T$, the **Debye screening mass m_D** preventing infrared divergencies.
- **hard collisions** ($|t| > |t|^*$): kinetic pQCD calculation

Two calculations,
with $g(\mu)$ evaluated at:

$\mu \sim T$ as for the soft component (HTL1)

$$\mu = |t| = -Q^2 \quad (\text{HTL2})$$

Effect of intrinsic k_T



WA75 (and WA92)

- Azimuthal correlations: $dN/d\Delta\phi$ with $\Delta\phi \equiv \phi_Q - \phi_{\bar{Q}}$
- Pair transverse momentum: $dN/d(p_T^{Q\bar{Q}})^2$ with $P_T^{Q\bar{Q}} \equiv P_T^Q + P_T^{\bar{Q}}$

LO: $\Delta\phi = \pi$

NLO: $\Delta\phi = \pi + \mathcal{O}(\alpha_s)$

$P_T^{Q\bar{Q}} = 0$ Q and \bar{Q} back-to-back
 $P_T^{Q\bar{Q}} = \mathcal{O}(\alpha_s)$ $Q\bar{Q}g$ final-state

Intrinsic k_T looks necessary to reproduce the data!

We can check what happens in our simulations with/without it.