

Gravitational-wave interferometric detectors and the Virgo experiment

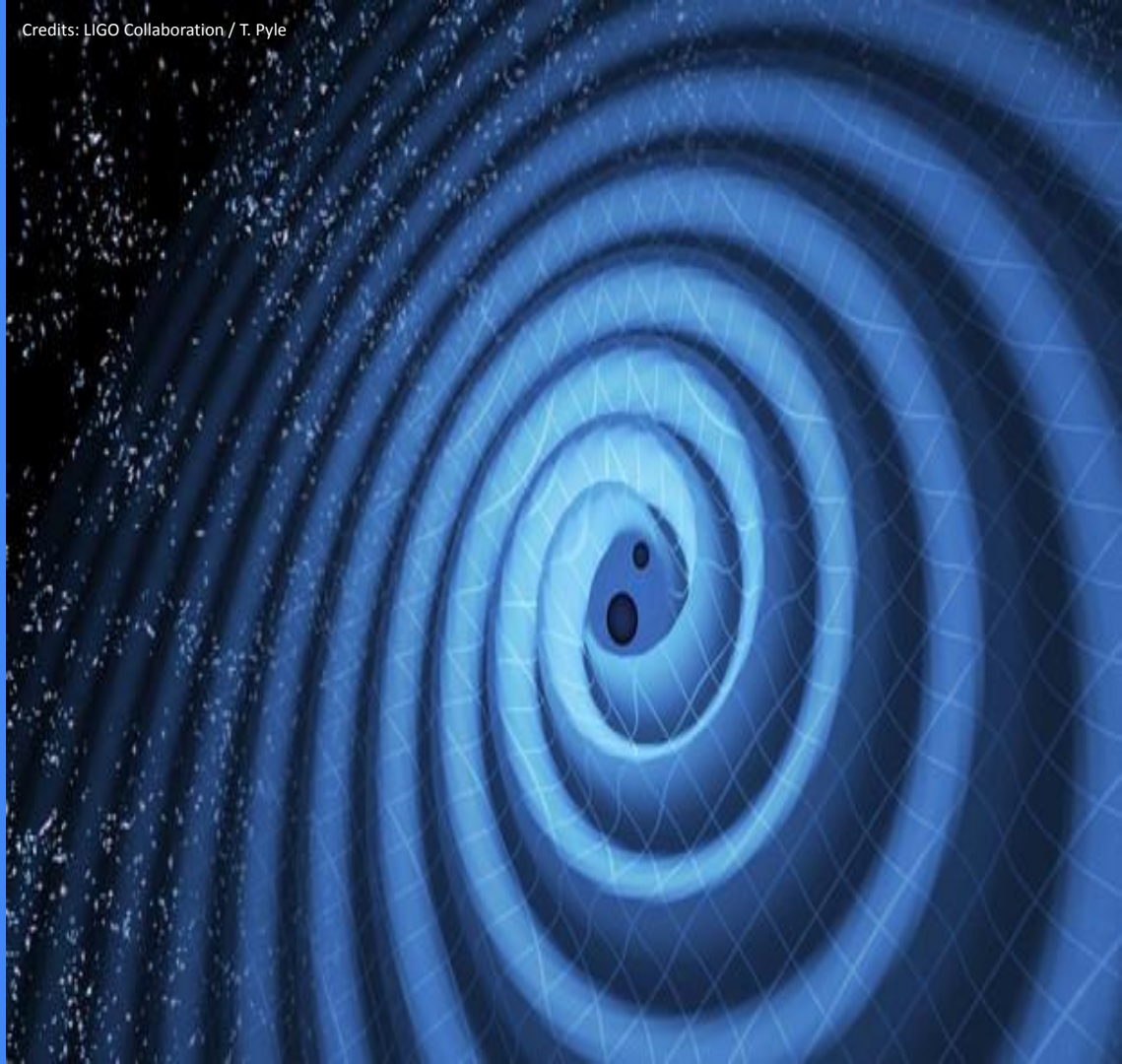
[F. De Marco](#)

Sapienza University and INFN Roma1

Young@INFN - Sapienza University, 11th April 2024



1. Gravitational Waves



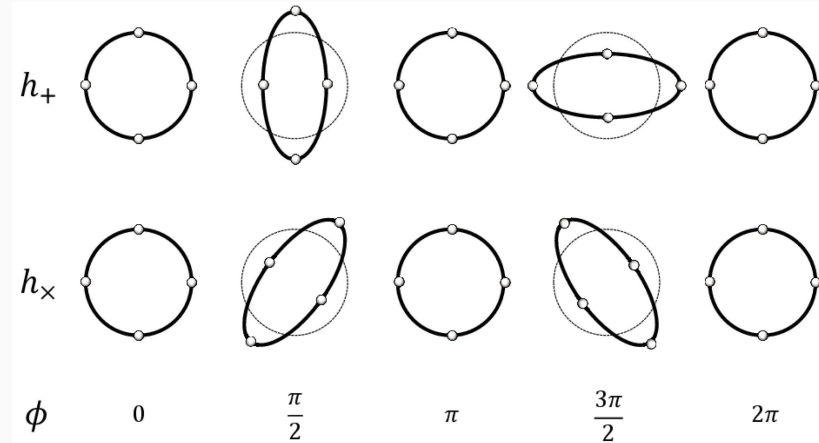
Gravitational-waves (GWs)

- Predicted by Einstein (1916): GR equations do have a wave solution!
- A GW can be seen as a *space-time strain* $h = \Delta L/L$
- In the Transverse-Traceless (TT) gauge:

$$h_{ij}^{TT}(t, \mathbf{x}) = \frac{4G}{c^4} \frac{1}{r} \int d^3\mathbf{x}' T_{ij} \left(t - \frac{r}{c} + \frac{\mathbf{x}' \cdot \hat{\mathbf{n}}}{c}, \mathbf{x}' \right)$$

- with no sources (masses):

$$h^{TT}(t, z) = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & h_+ & h_\times & 0 \\ 0 & h_\times & -h_+ & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \cos[\omega(t - z/c)]$$



Credits: [Królak, Patil - Universe 2017, 3\(3\), 59](#)

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HUGE suppression factor: $h \lesssim 10^{21}$ on Earth

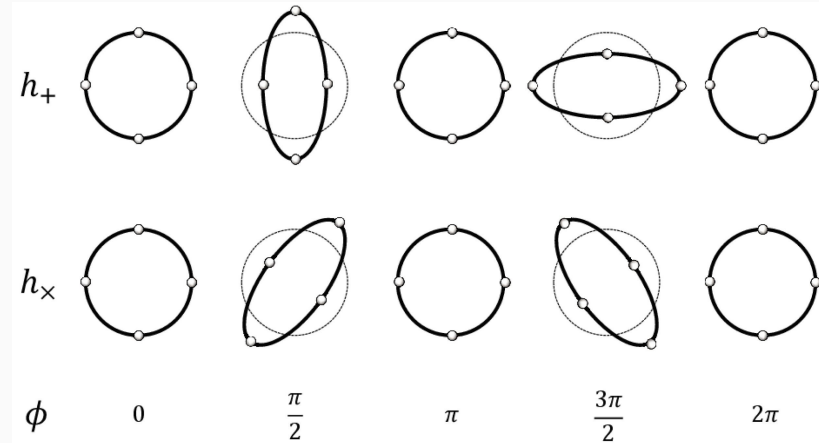
amplitude $\propto r^{-1}$

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travels at c

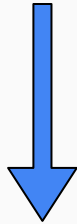
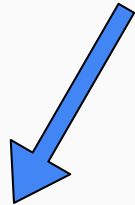
2 transverse polarizations: h_+ and h_\times



Credits: [Królak, Patil - Universe 2017, 3\(3\), 59](#)

Gravity at its strongest regime

- GWs are (almost) not perturbed by matter
- Generated in very strong regime (remind the suppression factor!)



Astrophysics

- source study
- multi-messenger observations

Cosmology

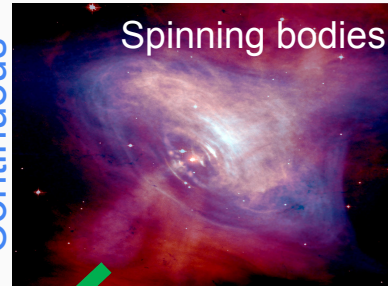
- BH population
- H_0 tension
- Primordial GW

Fundamental physics

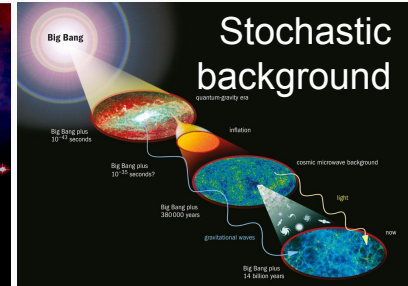
- Neutron Star EOS

GW sources

Modeled

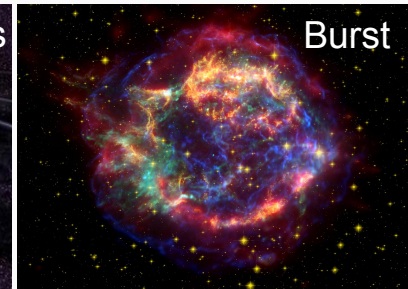


Unmodeled



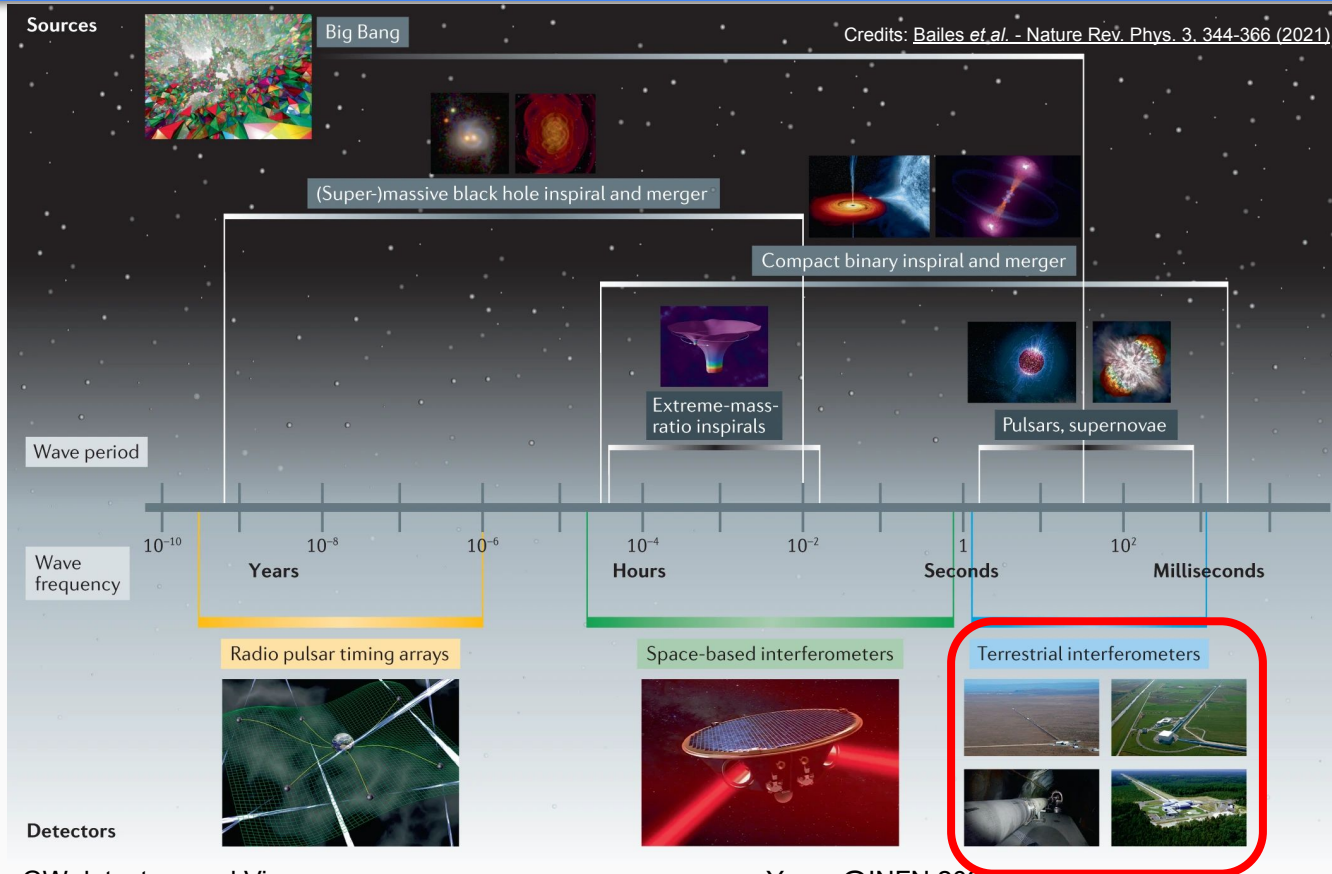
Continuous

Transient



See next talk by F. Attadio...

GW spectrum

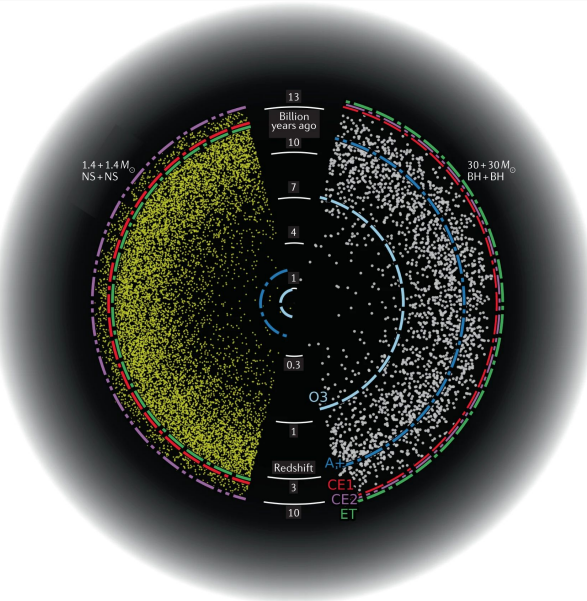


As with e.m. observations, different wavelength regimes require very different approaches!

State-of-the-art of GW detections

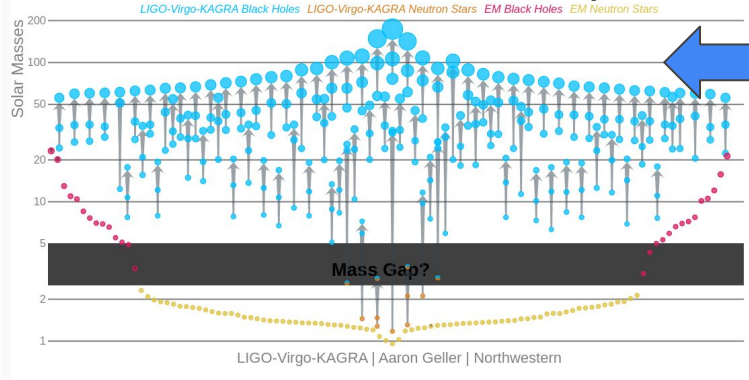
- A better sensitivity improves our astrophysical reach
- Figure of merit: Binary Neutron Star Horizon (today $\sim 55\text{-}60$ Mpc for Virgo)

Worldwide network of GW detectors



Credits: Bailes *et al.* - Nature Rev. Phys. 3, 344-366 (2021)

Masses in the Stellar Graveyard



And the observations?

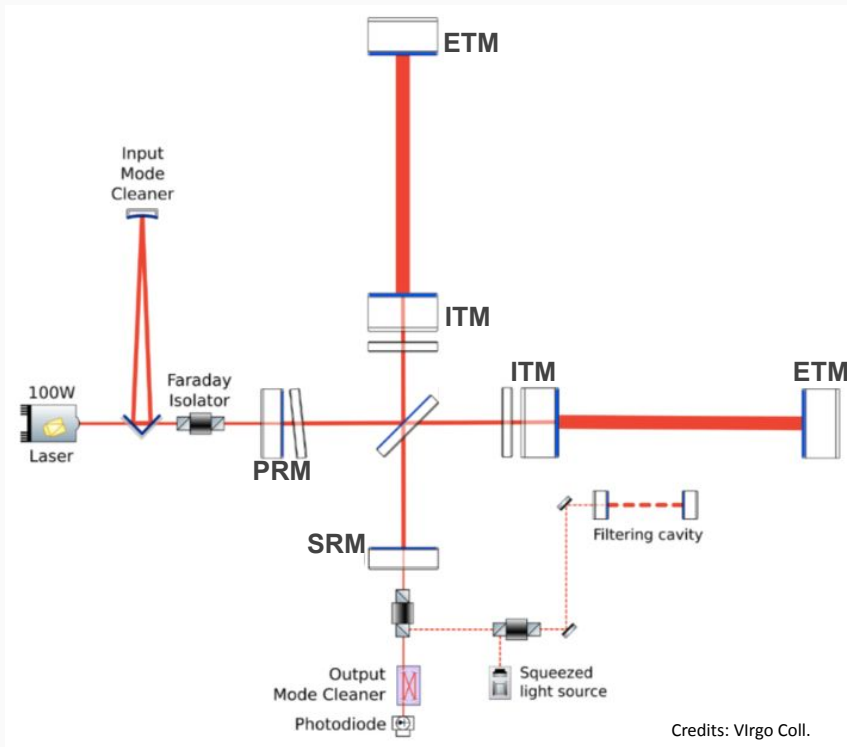
- 90 events in the past observing runs (O1, O2, O3)
- O4a started on **May 24th, 2023** with LIGO detectors
- O4b started **yesterday!**
- **82 event candidates** at April 11th, 2024

Check out the [public alerts web page](#) and keep counting!

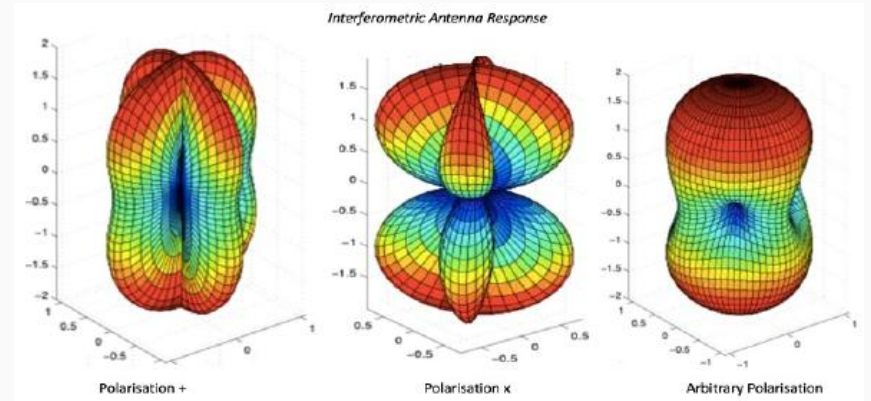
2. Ground-based detectors



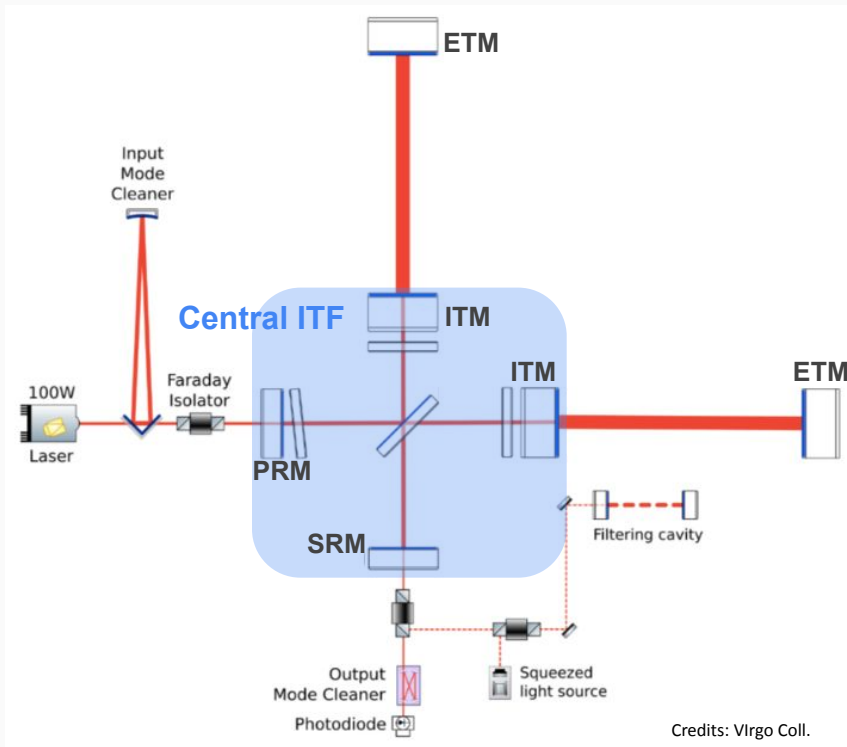
Earth-based interferometers



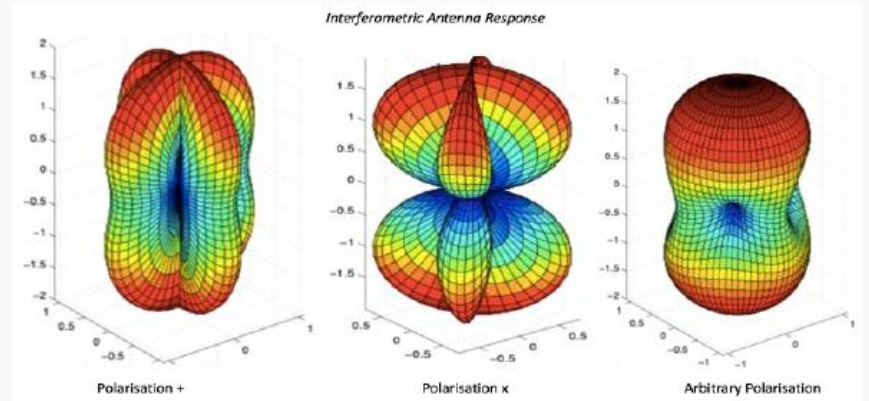
- Dual-recycled Michelson interferometer (ITF) with km-long Fabry-Perot (FP) arm cavities
- Test masses (mirrors) are displaced by the passage of a GW
- Differential measurement of **GW amplitude**
$$\Delta\phi = \omega_L \Delta t \simeq \frac{4\pi L}{\lambda_L} h_+$$
- Broad frequency and angular response, with some blind spots: antenna pattern



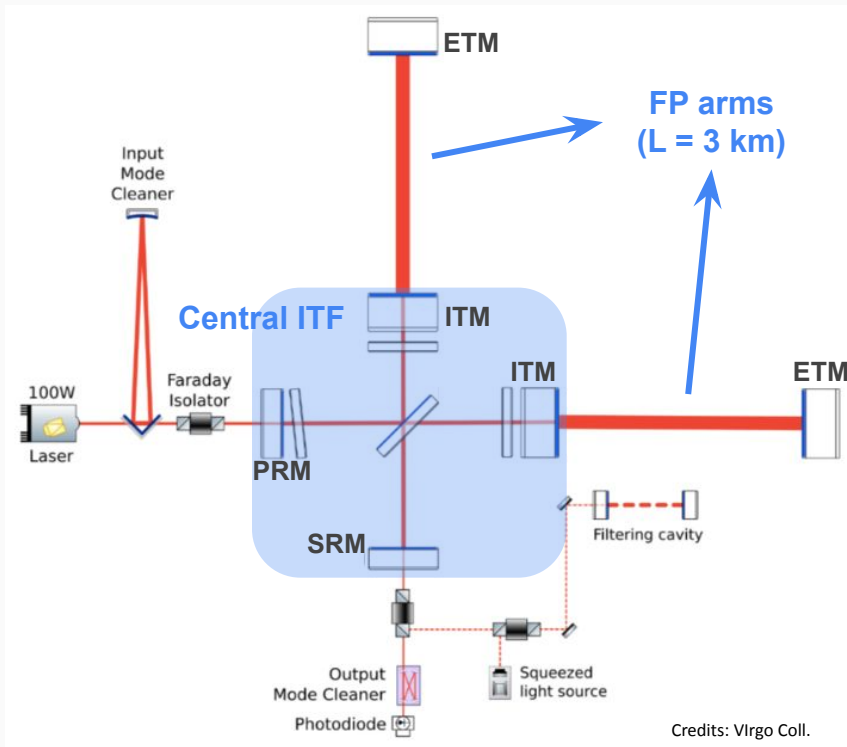
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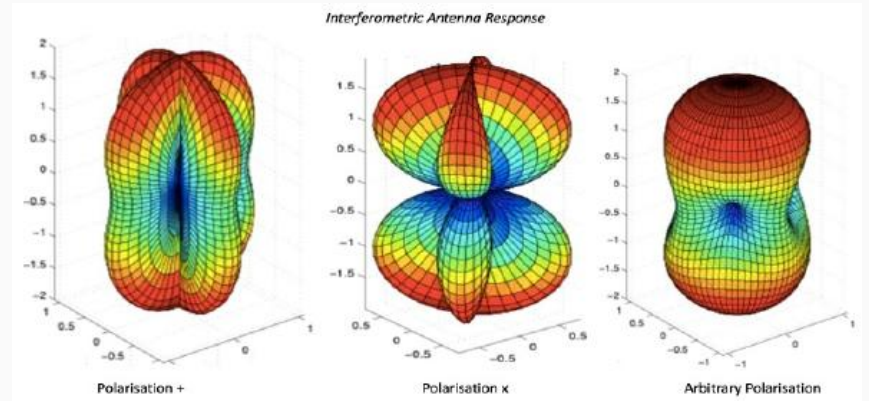
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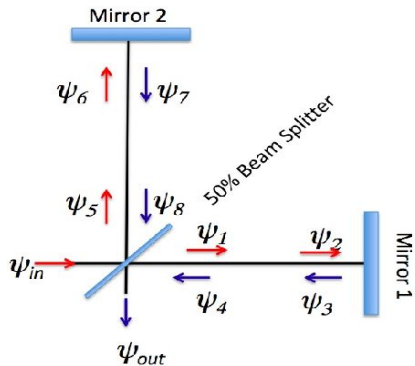
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Why Fabry-Perot (FP) arms?



Michelson ITF

$$\Delta\phi \simeq \frac{2\pi}{\lambda_L} (L_x + L_y) h$$

Differential measurement

Approximations!



- GW much longer than ITF
- best directional sensitivity
- no modulation sidebands in the laser
- equal ITF arms length

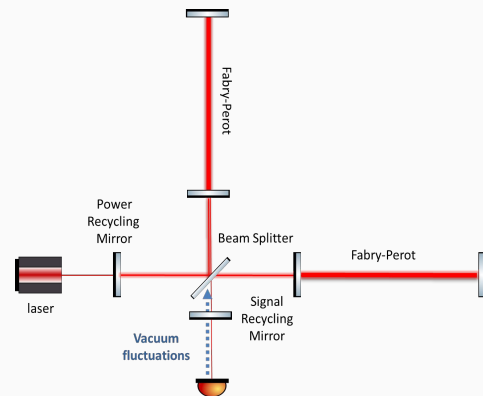
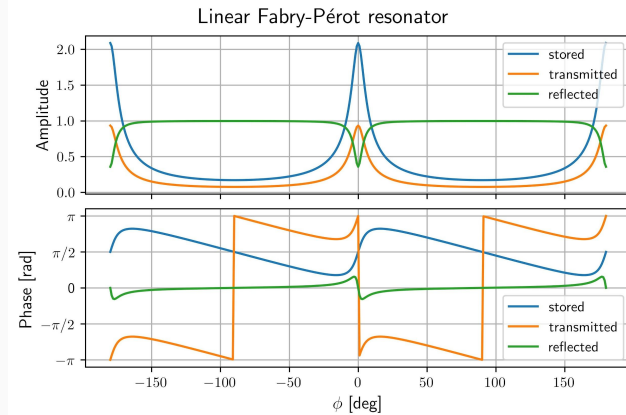
Michelson ITF with FP arms

$$\Delta\phi(\omega) \simeq \frac{2\pi}{\lambda_L} (L_x + L_y) \frac{2\mathcal{F}}{\pi} \frac{1}{\sqrt{1 + (\omega/\omega_p)^2}} h$$

FP gain FP cut-off

$$\mathcal{F} = 460$$

$$\omega_p \approx 2\pi \cdot 50\text{Hz}$$



Sensitivity curve

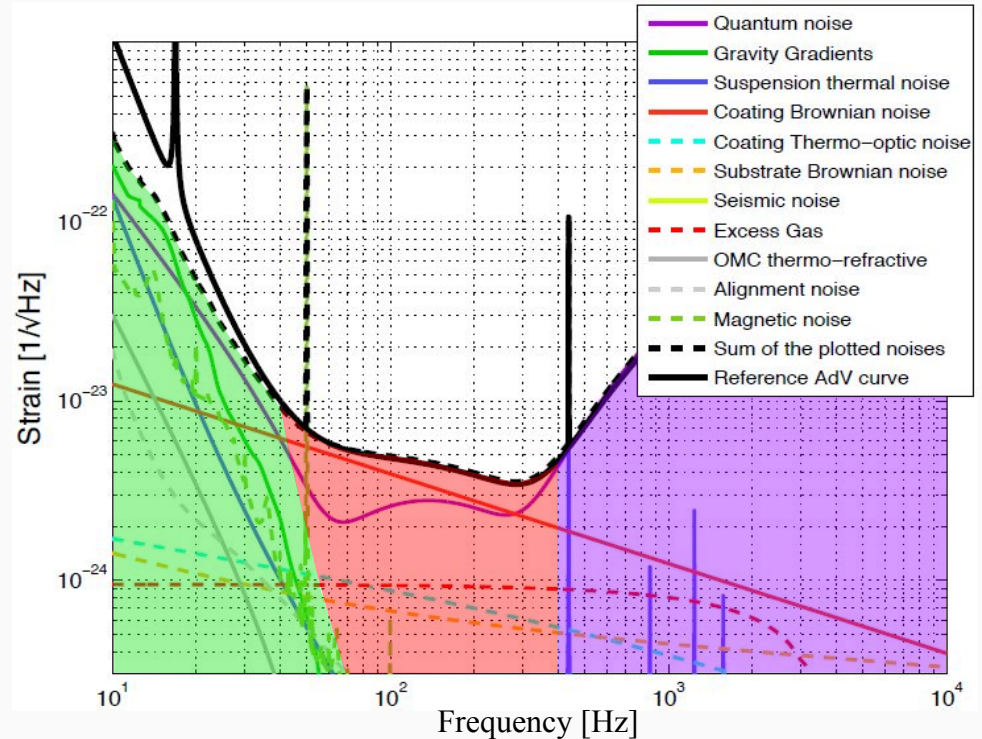
ASD of (independent and stationary!) noise sources:

$$\mathcal{N}(f) = \sqrt{2 \lim_{T \rightarrow \infty} \frac{1}{T} \left| \int_{-T/2}^{T/2} n(t) e^{-i2\pi f t} dt \right|^2}$$

“Fundamental” noise sources:

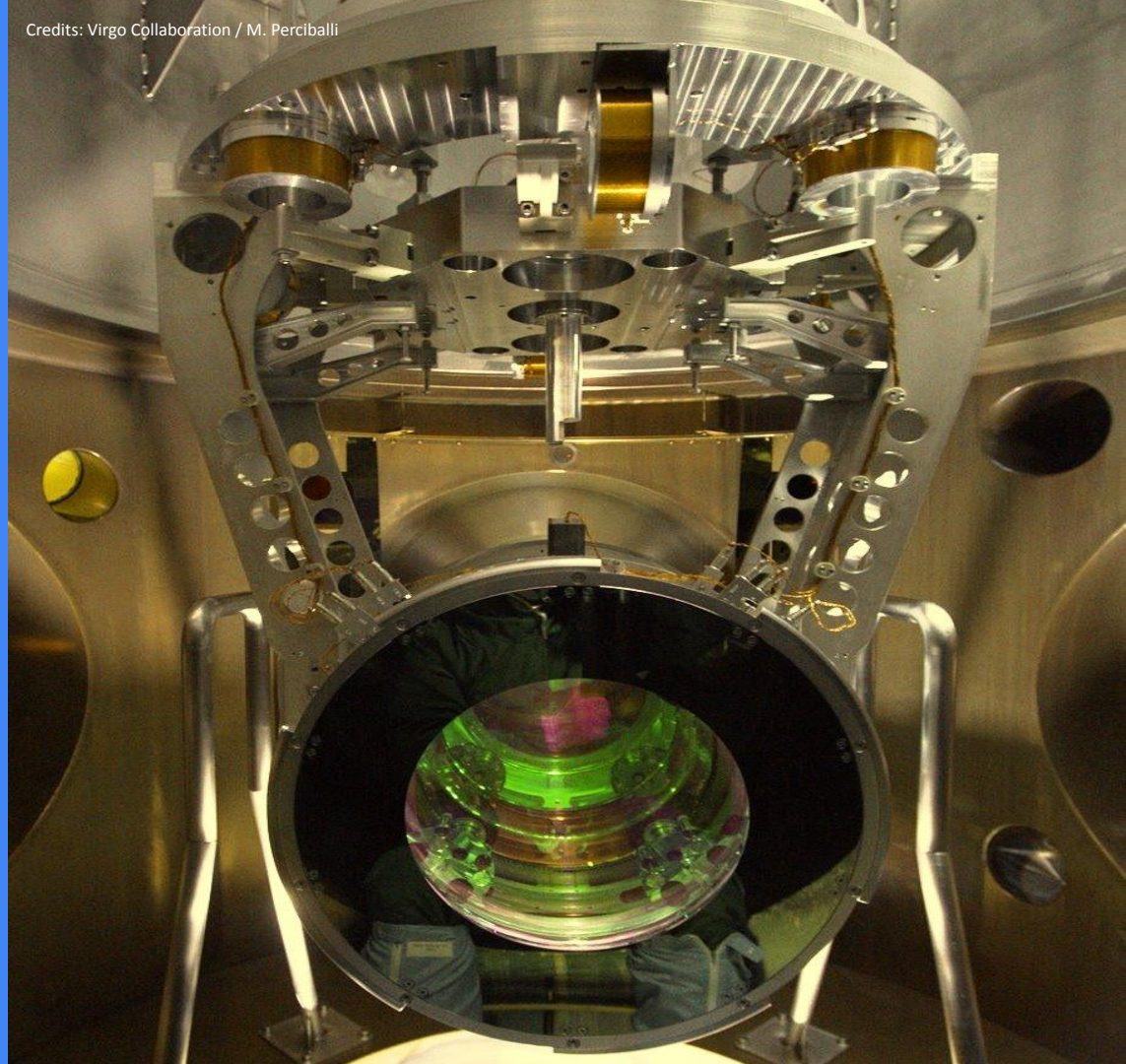
- **Seismic noise**
- **Thermal noise**
- **Quantum noise**

...but there are many other noises :(
Reducing their contribute is the main effort
of experimental physicists in GW!



Credits: [Virgo Coll. - Class. Quant. Grav. 32 024001 \(2015\)](#)

3. Noise sources



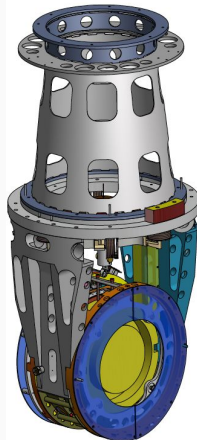
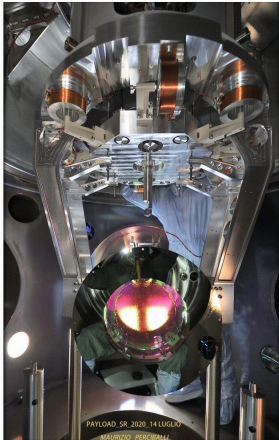
Seismic and environmental noise (<10 Hz)

- Relevant at low frequencies (micro-earthquakes, wind, sea activity...)
- Inverted pendulum + chain of pendulums with low proper frequency (0.6 Hz) to make a “cascade filter”: **Superattenuator**
- Heavier end test masses

42 kg
now

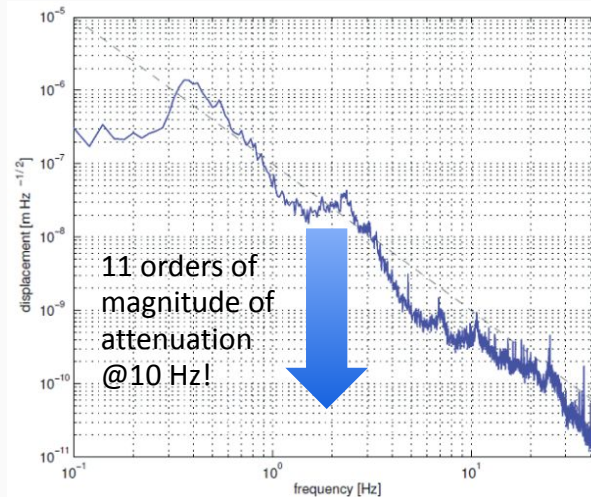


105 kg
O5

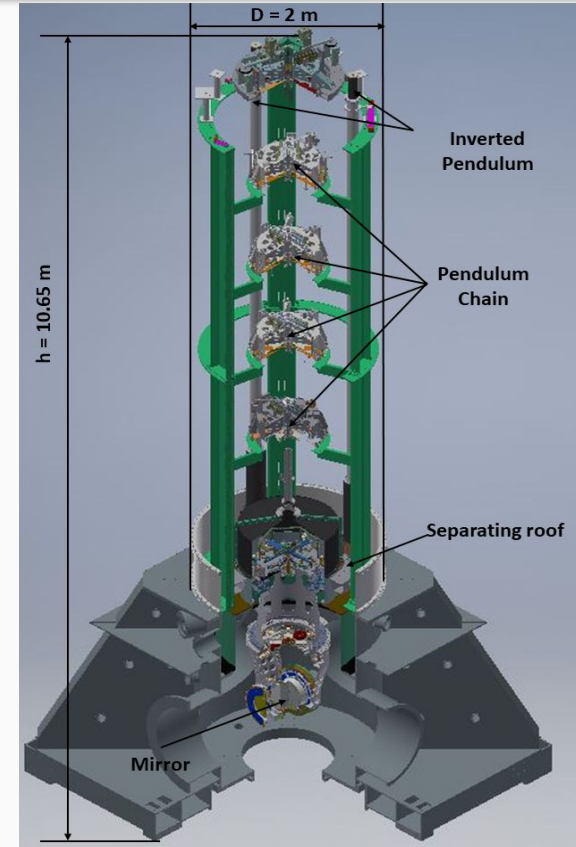


Credits: M. Perciballi

Seismic noise @Virgo site



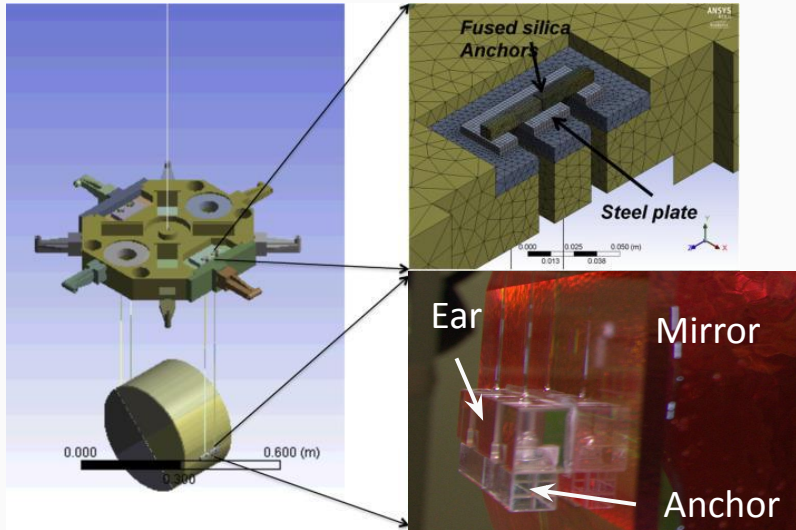
Credits: [Virgo Coll. - Jour. of Low Frequency Noise, Vibration and Active Control 30\(1\):63-79 \(2011\)](#)



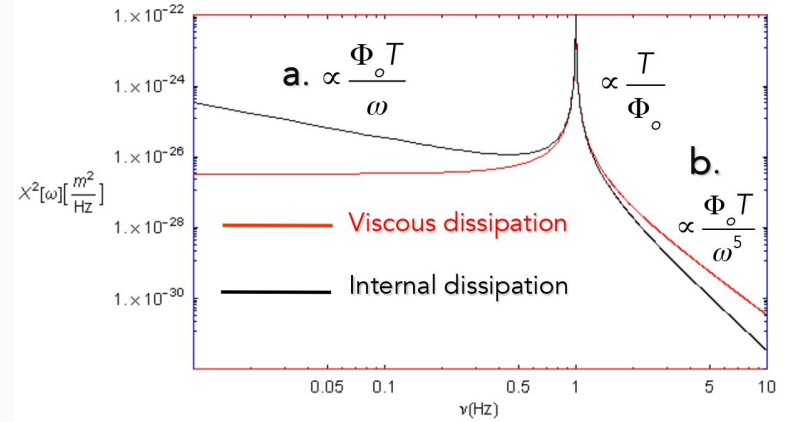
Thermal noise (<300 Hz)

- Any dissipation at non-zero temperature brings vibrational noise (Fluctuation-Dissipation Theorem)
- Relevant to ALL precision measurements!

$$S_x(\omega) \sim \begin{cases} \omega^{-1} & \text{if } \omega \ll \omega_0 \\ \omega^{-5} & \text{if } \omega \gg \omega_0 \end{cases}$$



Thermal noise of mirrors' pendulum motion



Many sources...

- Suspensions
- Bulk of the mirrors
- **Coatings**: very good from optical point of view, not from thermal one (active field of R&D)

...and solutions!

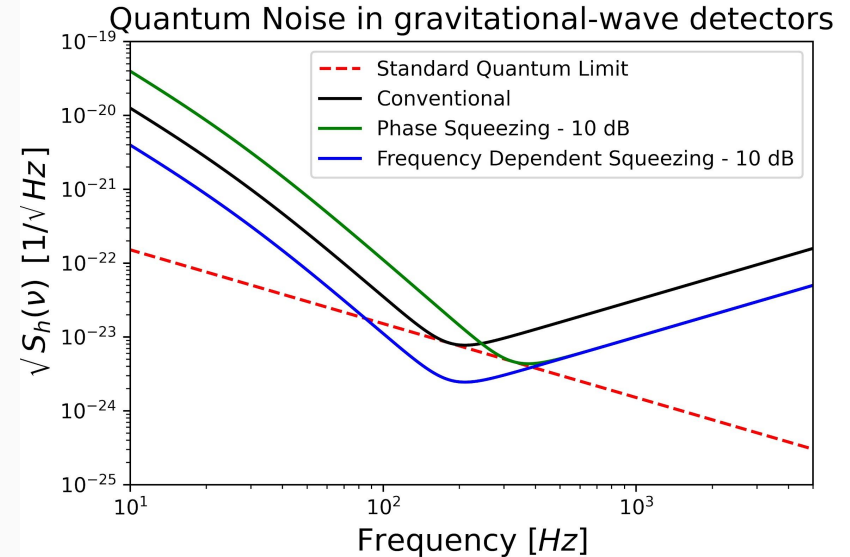
- monolithic suspensions
- better quality of materials
- lower optical power density (larger beams)
- cryogenics!! (KAGRA, ET...)

Quantum noise (>300 Hz)

- Macroscopic manifestation of the discreteness of laser light
- Originated by vacuum field entering the dark port of the ITF
- 2 contributions
 - back-action: Radiation-Pressure Noise (RPN)
 - detection: Shot Noise (SN)
- They meet at the **Standard Quantum Limit (SQL)**, lowest frequency of quantum noise

$$S_h(\omega) = \frac{h_{SQL}^2}{2} \left(\frac{1}{\mathcal{K}(\omega)} + \mathcal{K}(\omega) \right)$$

Standard Quantum Limit (SQL) ITF optomechanical coupling

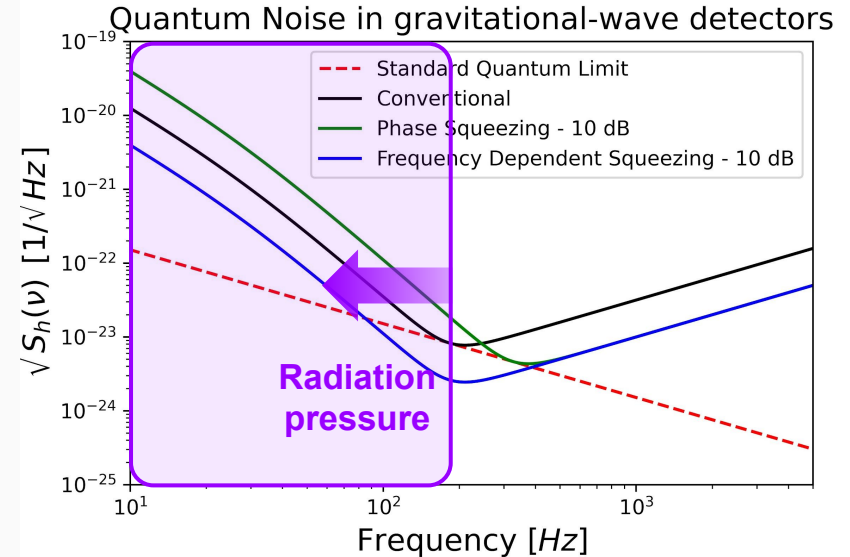


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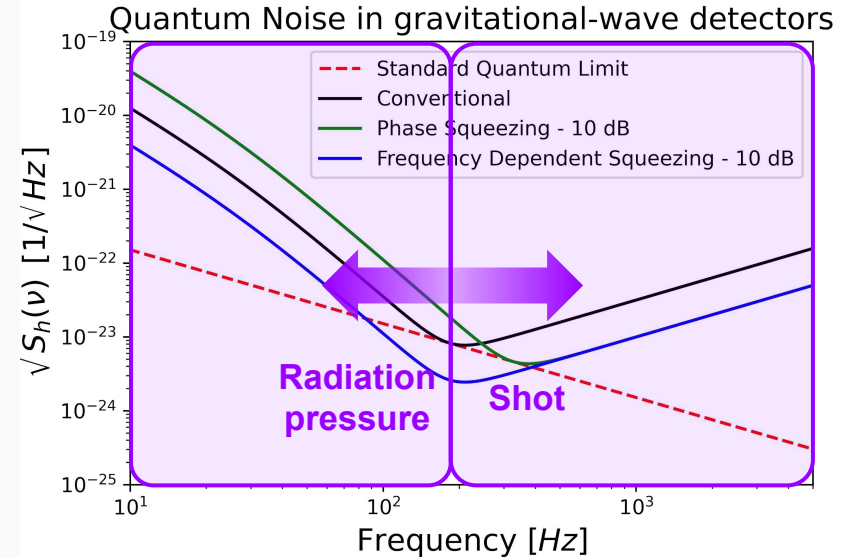


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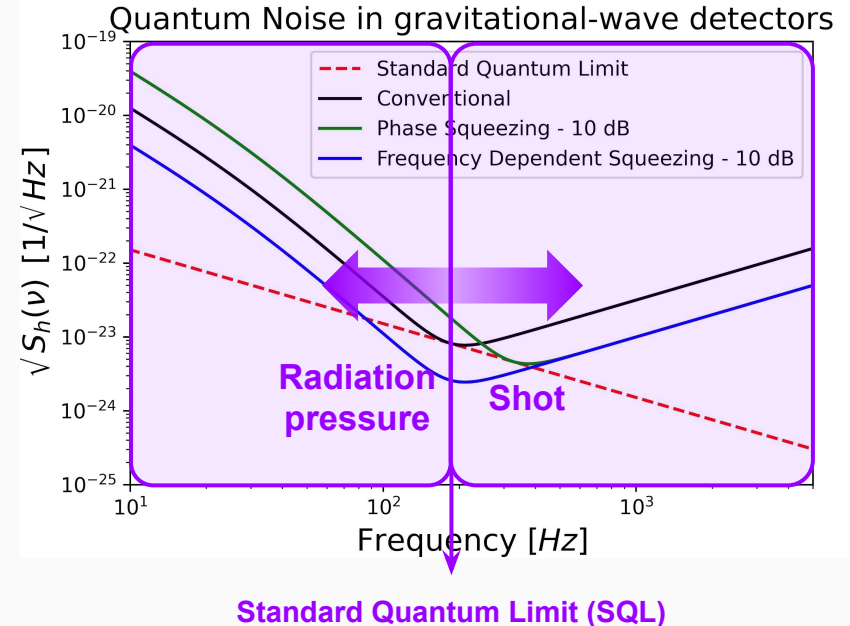


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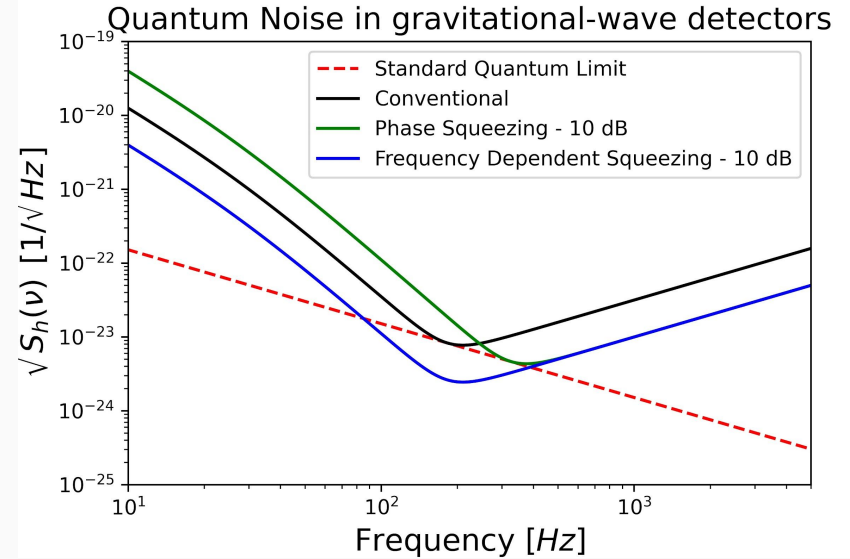
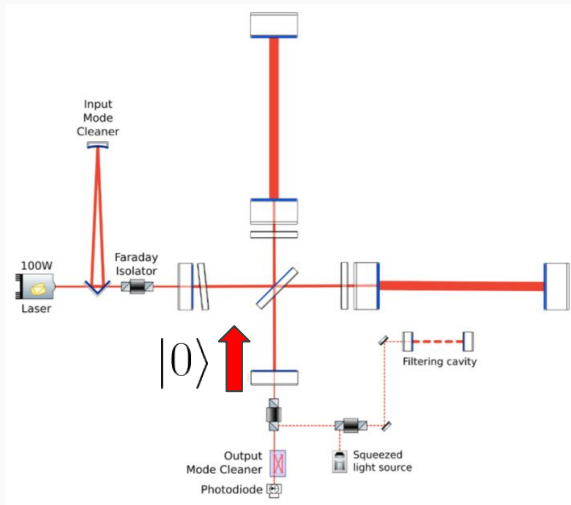


Quantum noise reduction with squeezing

Power spectrum of Quantum Noise:

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Standard Quantum Limit (SQL)
ITF optomechanical coupling

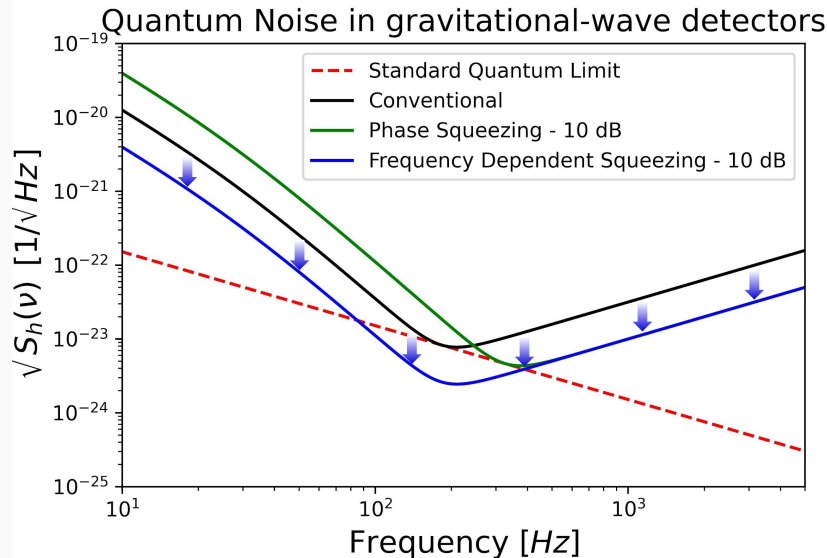
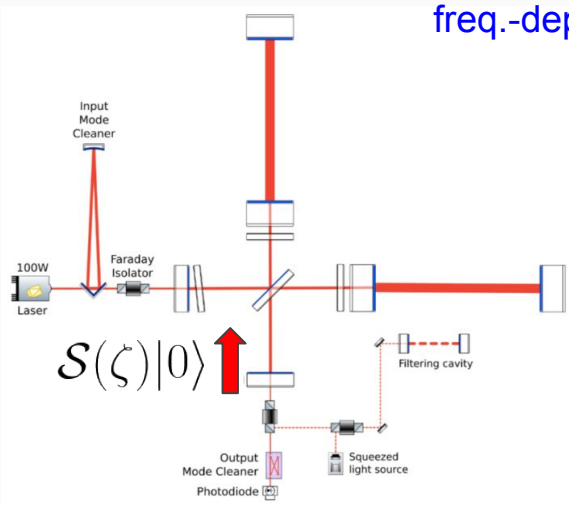


Quantum noise reduction with squeezing

Power spectrum of Quantum Noise:

$$S_h(\omega) = \frac{h_{SQL}^2}{2} \left(\frac{1}{\mathcal{K}(\omega)} + \mathcal{K}(\omega) \right) e^{-2r}$$

Suppression with
freq.-dependent squeezing

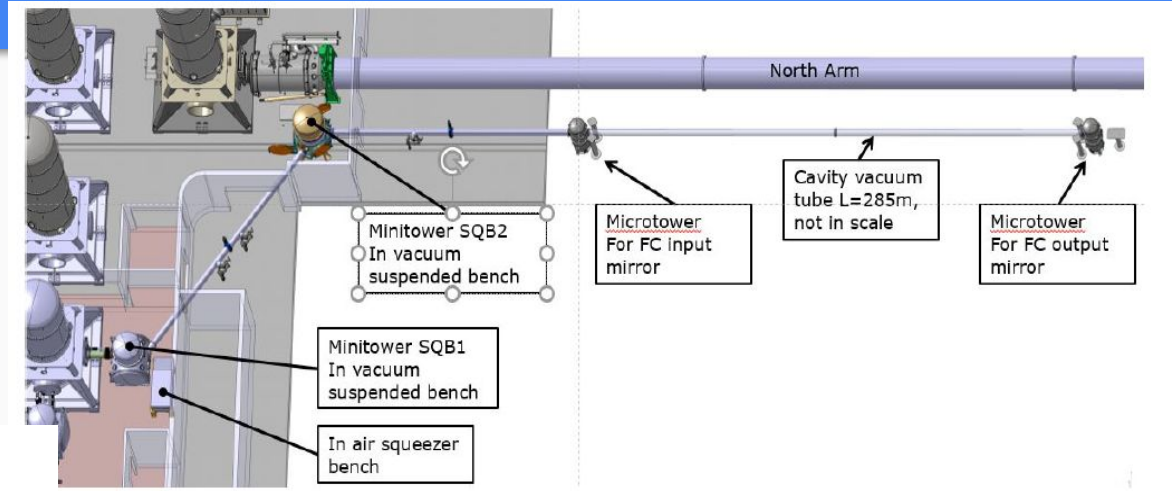
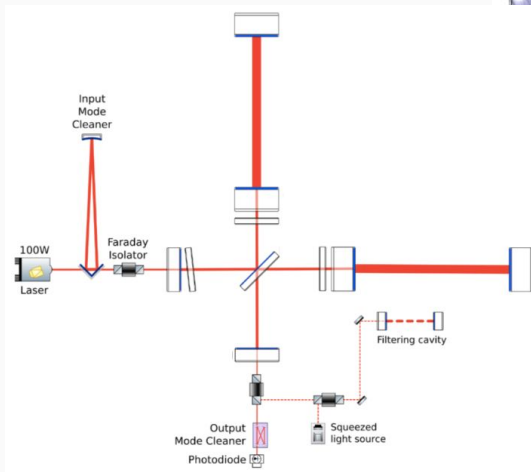


Provided a **specific tuning of the squeezing angle**, *Frequency-Dependent Squeezed* light reduces QN along the whole detection band!

Filter Cavity

- Squeezing ellipse can be rotated in a frequency-dependent manner with a detuned linear cavity (**Filter Cavity**)
- Central rotation angle @ ~ 50 Hz implies
 - Long cavity $L=285$ m
 - High finesse $F=10000$
- Round-trip losses in AdV+ FC: **50-90 ppm**

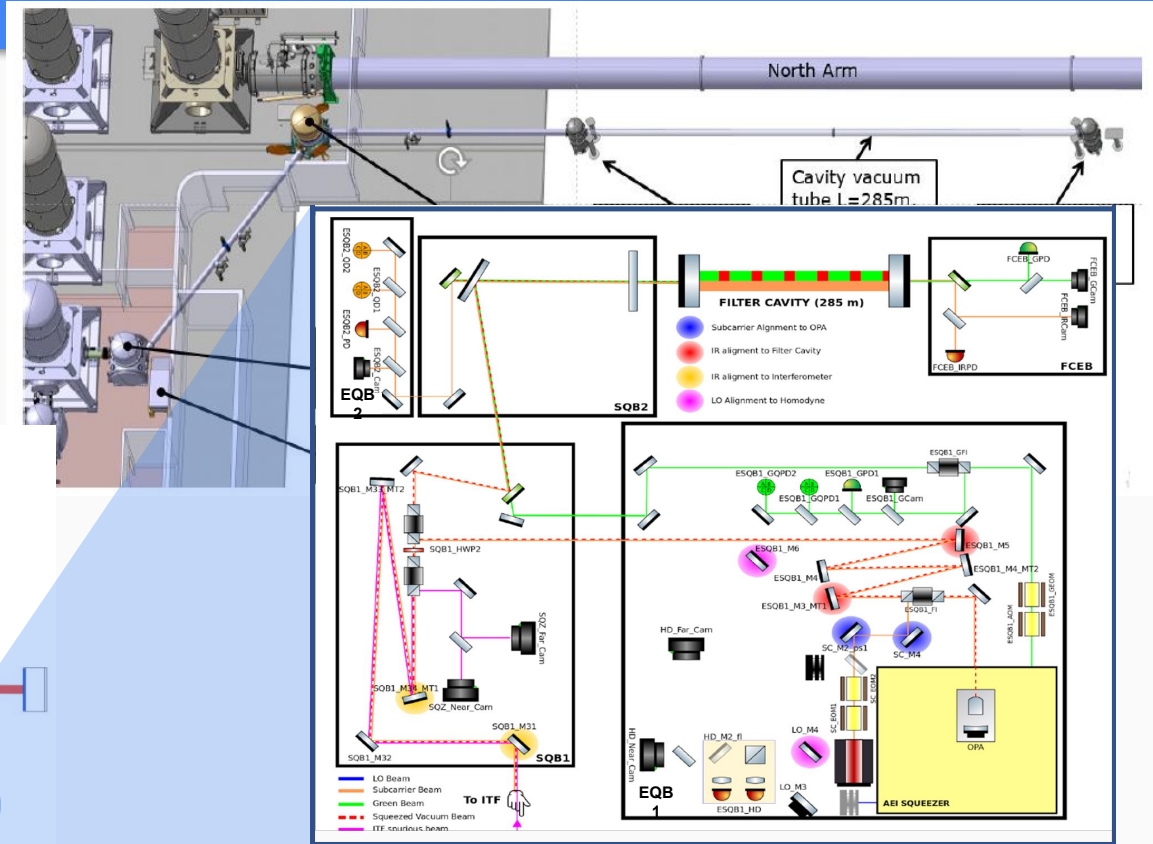
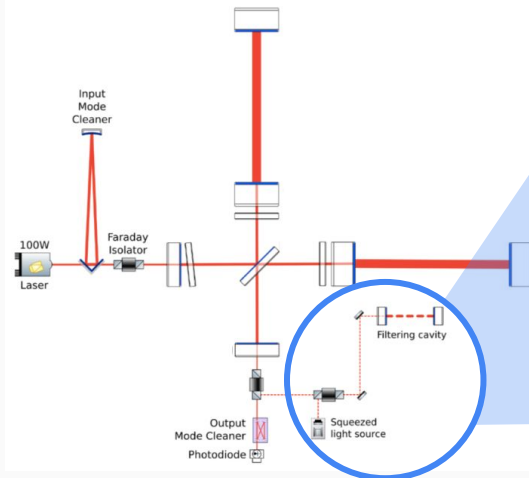
[Virgo Coll. - Phys. Rev. Lett. 131.041403 \(2023\)](#)



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[Virgo Coll. - Phys. Rev. Lett. 131 041403 \(2023\)](#)



Credits: [Di Pace - Phys. Scr. 96 124054](#)

To wrap up...

1. GWs are “ripples of space-time” caused by strong acceleration of huge masses
2. They can be revealed on the Earth with long-baseline Michelson interferometers, such as LIGO, VIRGO and KAGRA
3. The most important noise sources are seismic, thermal and quantum noise. Higher laser power, heavier mirrors, low-dispersion materials and squeezing of light are the main measures to counter them
4. Frequency-dependent squeezing of light can be implemented through an external detuned cavity coupled to the main interferometer
5. Improving the sensitivity of GW detectors allow to build up a consistent catalog of observations and make new physics!

Thank you for your attention

Questions??

Email:

francesco.demarco@roma1.infn.it



Backup slides



GW interaction with the detector

- Michelson ITF only, and GW much longer than the ITF
- GW distorts space-time metric, affecting the propagation time in the arms

$$t_x \simeq \frac{2}{c} \int_0^L \left(1 + \frac{1}{2} h_+ \right) dx$$

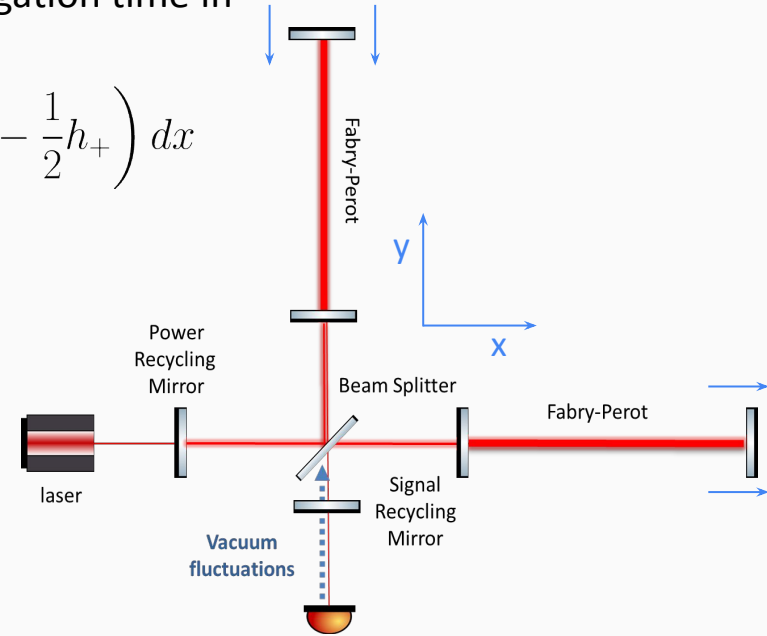
$$t_y \simeq \frac{2}{c} \int_0^L \left(1 - \frac{1}{2} h_+ \right) dx$$

- Phase shift at the output dark port

$$\Delta\phi = \omega_L \Delta t \simeq \frac{4\pi L}{\lambda_L} h_+$$

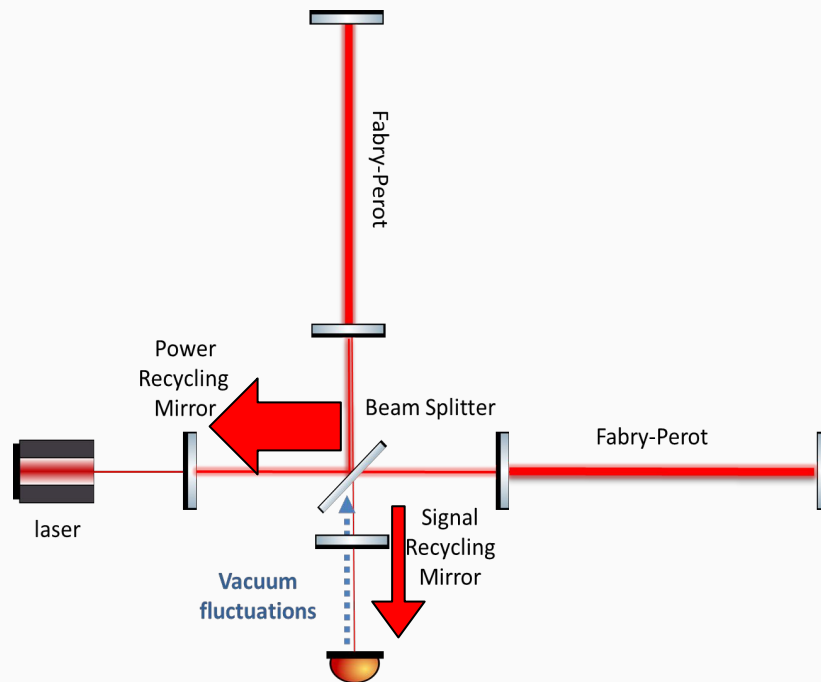
Typical values:

- FP arm length $L = 3$ km
(4 km for LIGO, 3 km for KAGRA)
- GW strain $h \approx 10^{-21}$



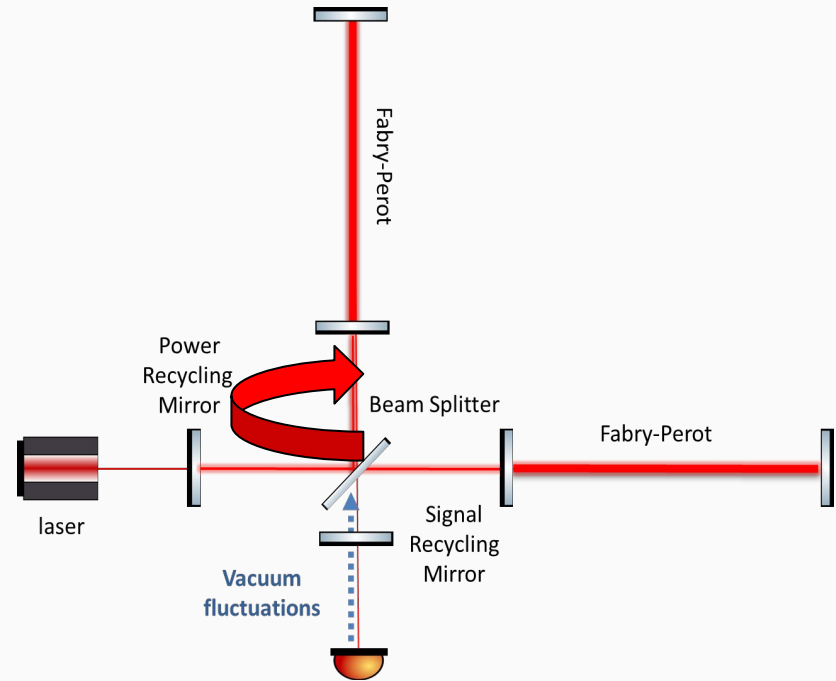
Power Recycling

- The ITF is always kept close to *dark fringe condition*, i.e. destructive interference at the output (dark) port
- Power is almost entirely reflected back towards the laser source



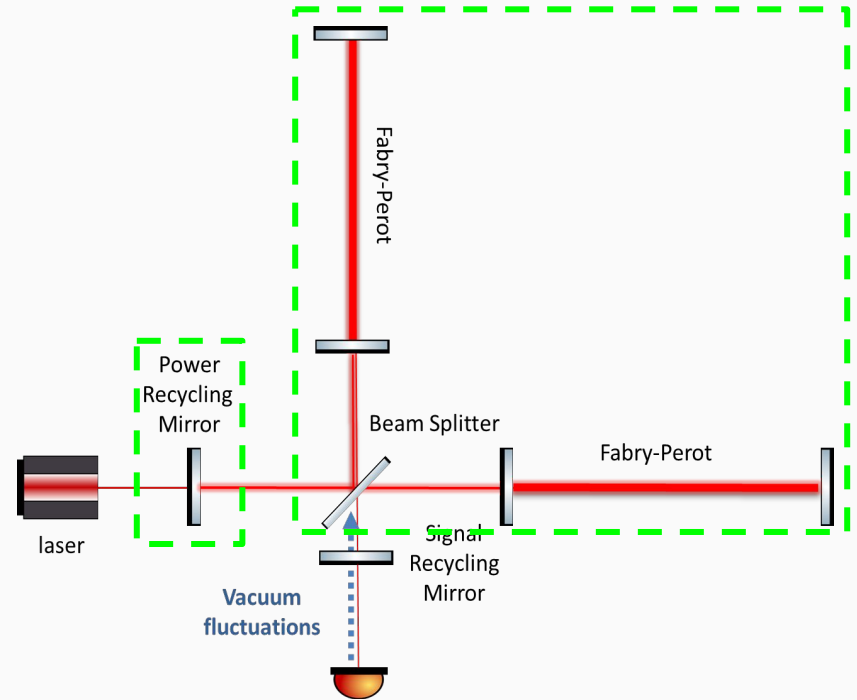
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- It creates another cavity, with the ITF as equivalent end mirror

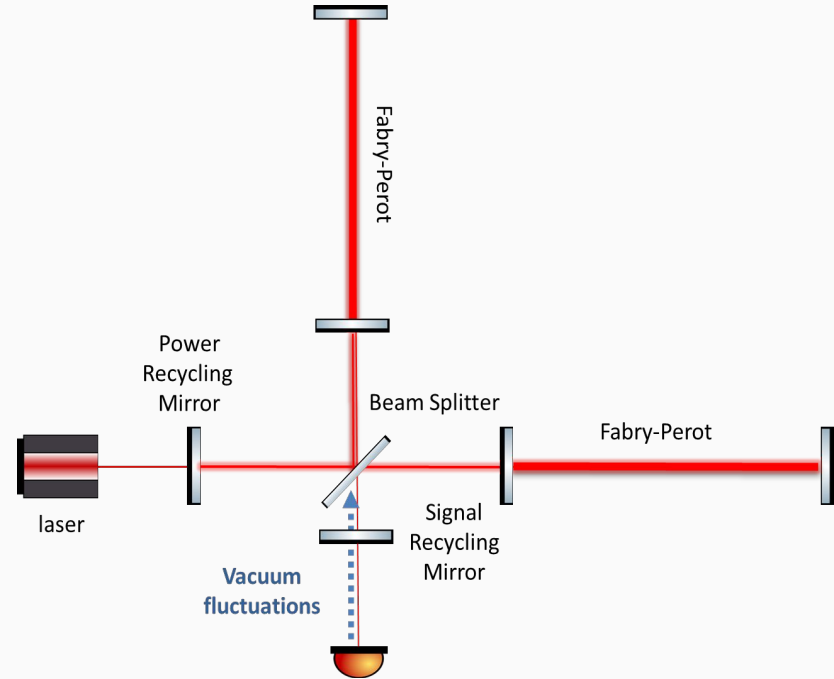


Signal Recycling

- Name is not so intuitive...
- FP cavities reduce the bandwidth

$$\gamma \simeq \frac{ct_{in}^2}{4L}$$

2π 50 Hz



Signal Recycling

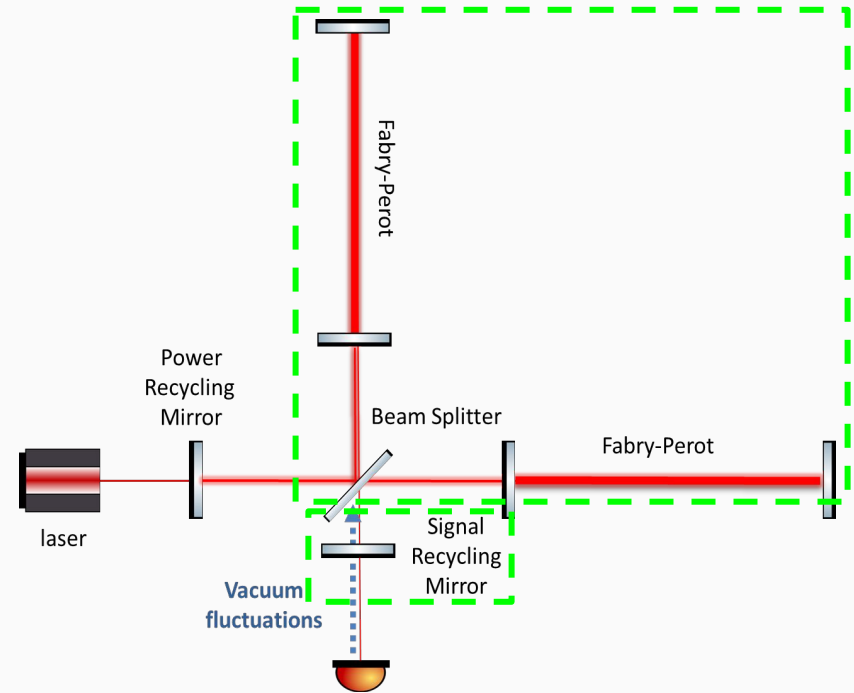
- Name is not so intuitive...
- FP cavities reduce the bandwidth
- You want something to keep the optical power high, while recovering a better bandwidth: the Signal Recycling Mirror (SRM)!
- As the PRM, it adds an optical cavity which closes with the ITF itself

$$\gamma \simeq \frac{ct_{in}^2}{4L}$$

$$2\pi 50 \text{ Hz}$$

$$\gamma_{SR} \simeq \frac{ct_{in}^2}{4L} \frac{1 + r_{SR}}{1 - r_{SR}}$$

$$2\pi 430 \text{ Hz}$$



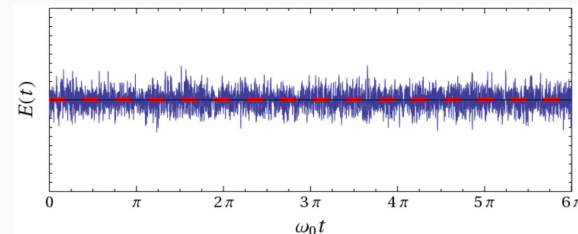
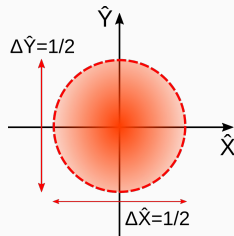
Squeezed states of light

Quantized EM field: $\hat{E}(t) \propto \int_0^\infty d\omega (\hat{a}_\omega e^{-i\omega t} + \hat{a}_\omega^\dagger e^{i\omega t}) \longrightarrow \hat{E}(t) \propto \int_0^\infty d\omega (\hat{X}_\omega \cos(\omega t) + \hat{Y}_\omega \sin(\omega t))$

Cosine (or amplitude) quadrature
Radiation Pressure Noise
Sine (or phase) quadrature
Shot Noise

Vacuum state:

$$|0\rangle$$



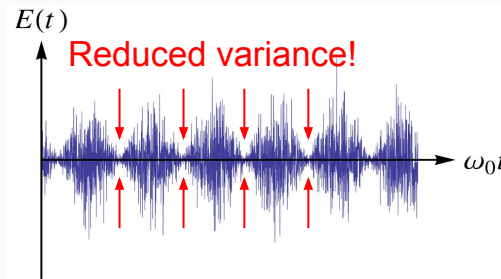
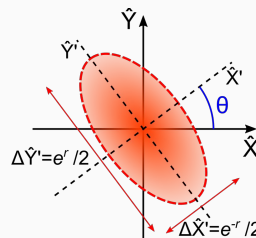
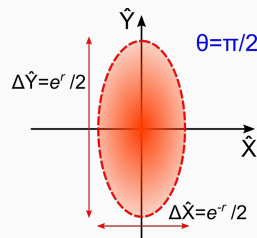
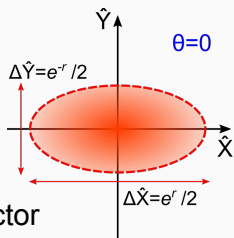
Credits: [Danilishin, Khalili - arXiv:1203.1706 \(2012\)](https://arxiv.org/abs/1203.1706)

Squeezed vacuum state:

$$\mathcal{S}(\zeta)|0\rangle$$

$\zeta = r e^{i\theta}$ Squeezing parameter

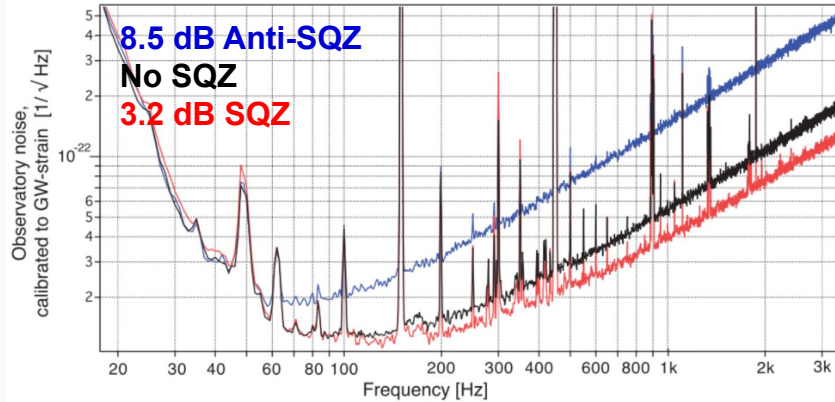
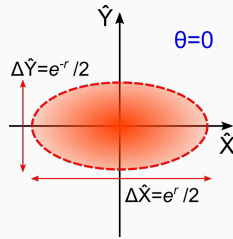
$r[\text{dB}] = 20 \log_{10}(e^r)$ Squeezing factor



Squeezing in O3 and O4

O3: Freq.-Independent Squeezing

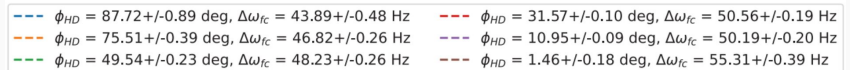
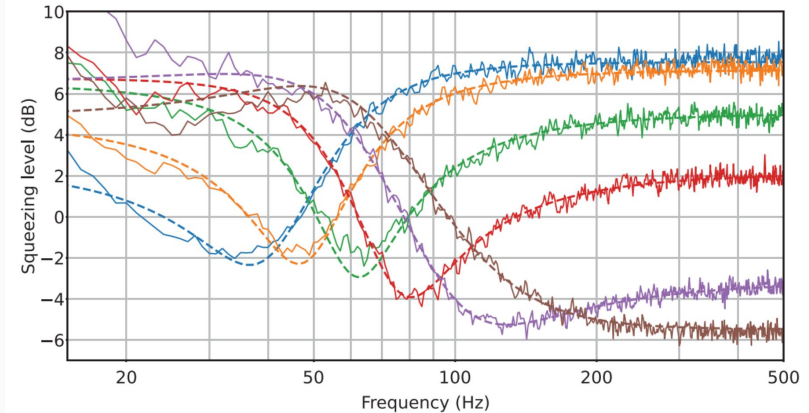
- Phase-squeezed light **3.2 dB**
- BNS range improved by 5 - 8 %
- Detection rate increased by 16 - 26 %



Credits: [Virgo Coll. - Phys. Rev. Lett. 123 231108 \(2019\)](#)

O4: Commissioning of a Freq.-Dependent Squeezing apparatus with Filter Cavity

- 5.6 dB at high frequencies and 2 dB around FC resonance
- Performances in the ITF similar to O3 due to ITF configuration

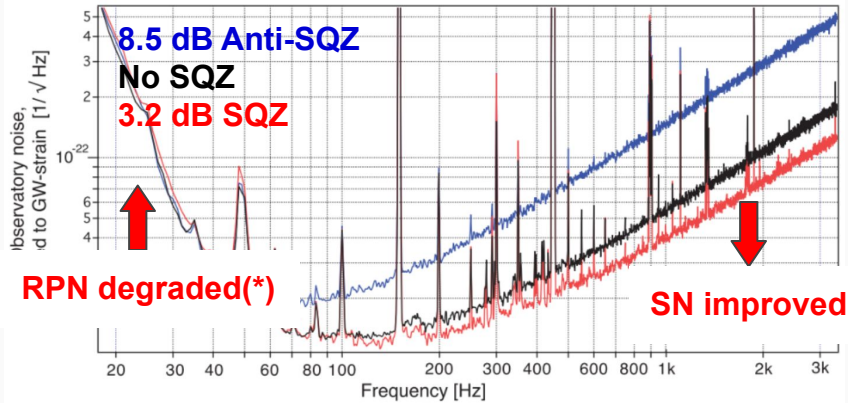
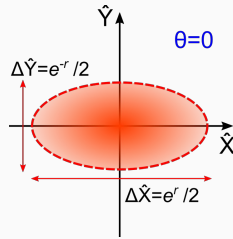


Credits: [Virgo Coll. - Phys. Rev. Lett. 131 041403 \(2023\)](#)

Squeezing in O3 and O4

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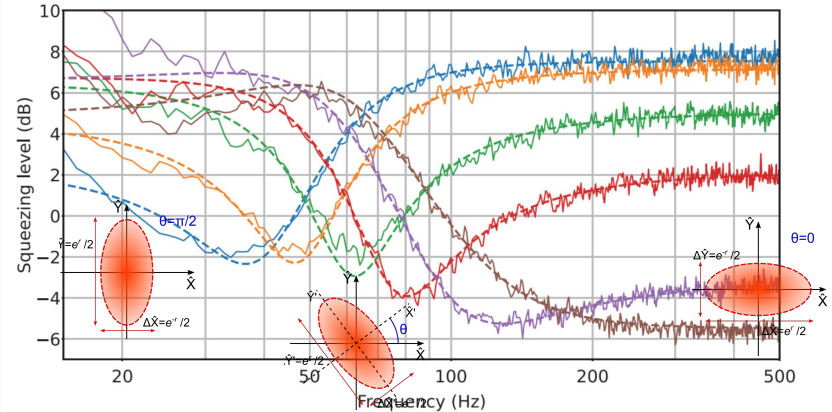


Credits: [Virgo Coll. - Phys. Rev. Lett. 123 231108 \(2019\)](#)

(*) Covered by other low-freq. noise sources

O4: Commissioning of a Freq.-Dependent Squeezing apparatus with Filter Cavity

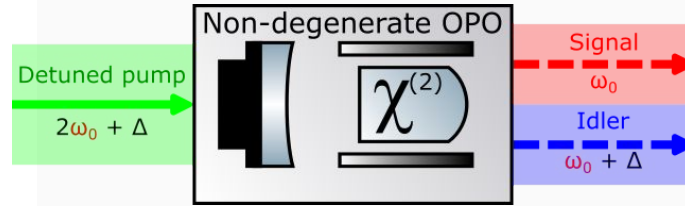
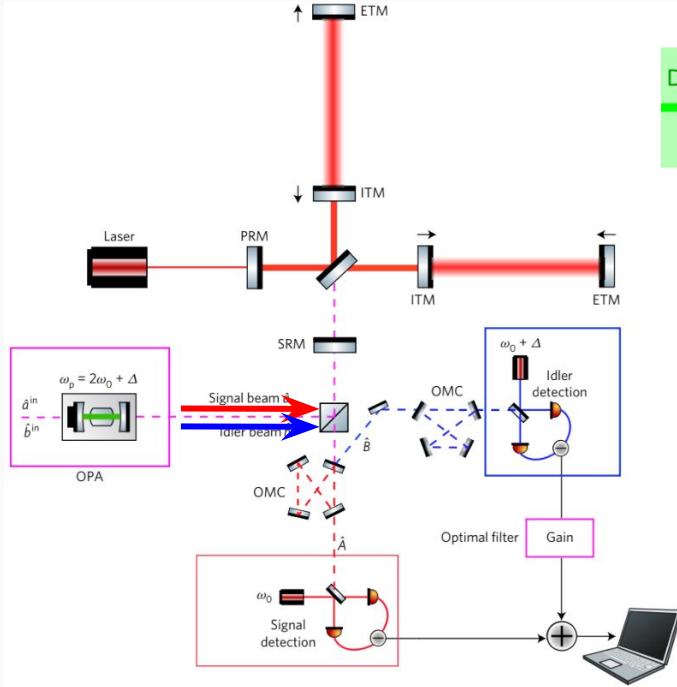
- 5.6 dB at high frequencies and 2 dB around FC resonance
- Performances in the ITF similar to O3 due to ITF configuration



--- $\phi_{HD} = 87.72 \pm 0.89$ deg, $\Delta\omega_{fc} = 43.89 \pm 0.48$ Hz	--- $\phi_{HD} = 31.57 \pm 0.10$ deg, $\Delta\omega_{fc} = 50.56 \pm 0.19$ Hz
--- $\phi_{HD} = 75.51 \pm 0.39$ deg, $\Delta\omega_{fc} = 46.82 \pm 0.26$ Hz	--- $\phi_{HD} = 10.95 \pm 0.09$ deg, $\Delta\omega_{fc} = 50.19 \pm 0.20$ Hz
--- $\phi_{HD} = 49.54 \pm 0.23$ deg, $\Delta\omega_{fc} = 48.23 \pm 0.26$ Hz	--- $\phi_{HD} = 1.46 \pm 0.18$ deg, $\Delta\omega_{fc} = 55.31 \pm 0.39$ Hz

Credits: [Virgo Coll. - Phys. Rev. Lett. 131 041403 \(2023\)](#)

A novel approach: EPR squeezing



1. **Signal** and **idler** are vacuum squeezed beams, EPR-entangled and detuned by Δ

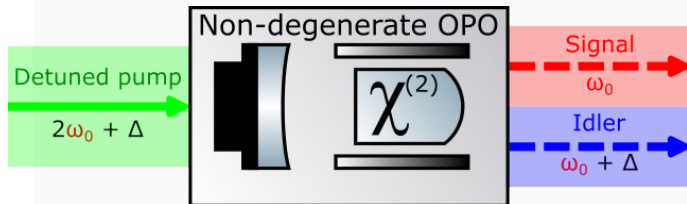
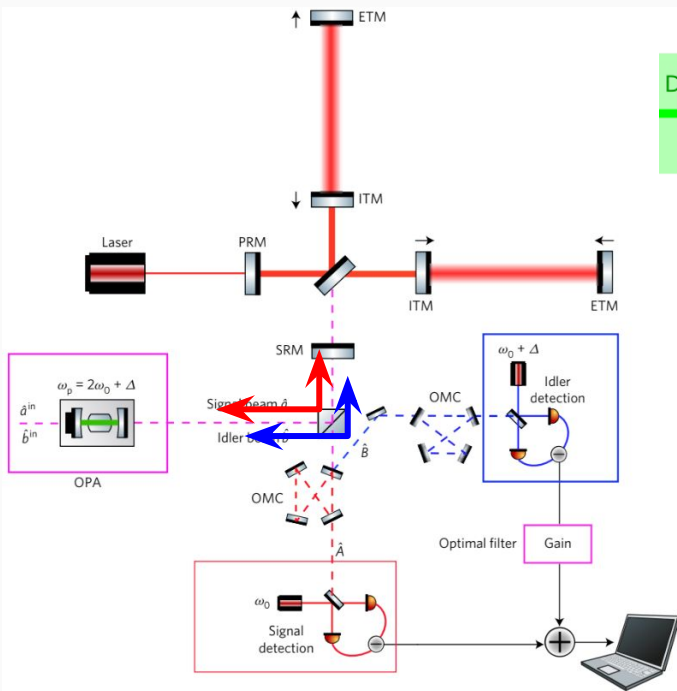
Pros:

- More compact and cheaper
- Avoids some optical losses

Cons:

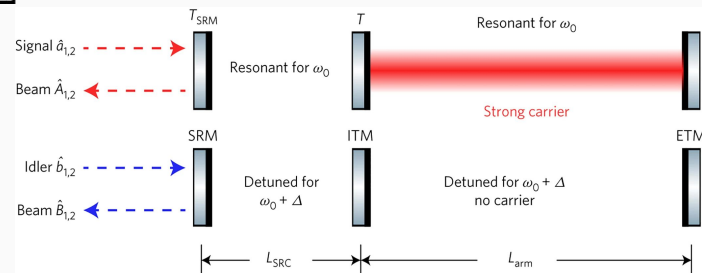
- 2 squeezed beams to be handled

A novel approach: EPR squeezing



1. **Signal** and **idler** are vacuum squeezed beams, EPR-entangled and detuned by Δ

2. The **idler** acquires frequency-dependence in the ITF due to its detuning



Pros:

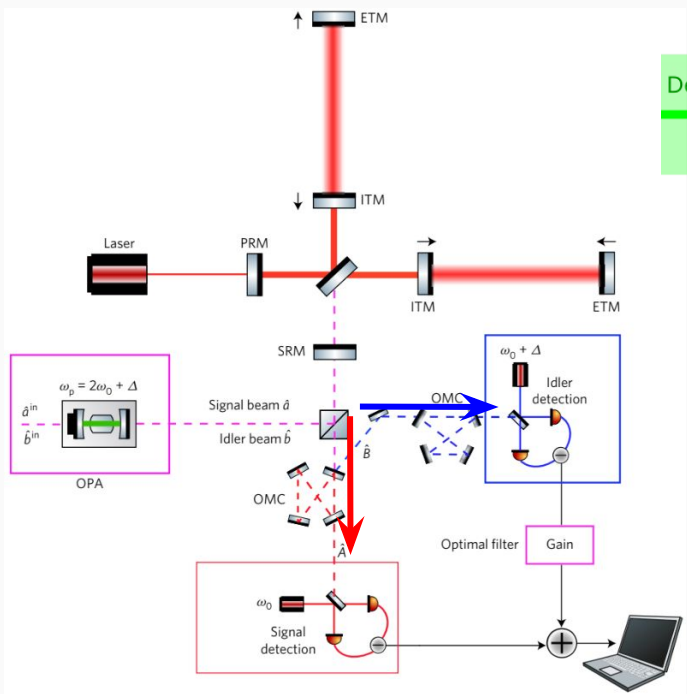
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Cons:

- 2 squeezed beams to be handled

Credits: [Ma et al. - Nature Phys 13, 776–780 \(2017\)](#)

A novel approach: EPR squeezing

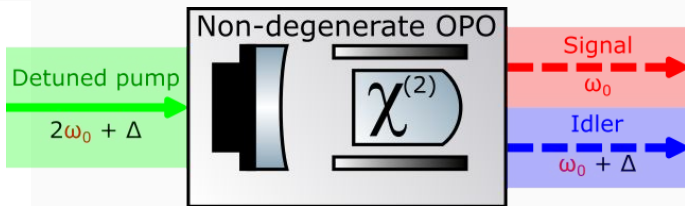


Pros:

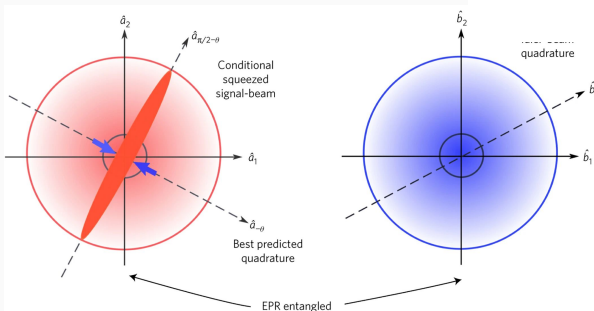
- More compact and cheaper
- Avoids some optical losses

Cons:

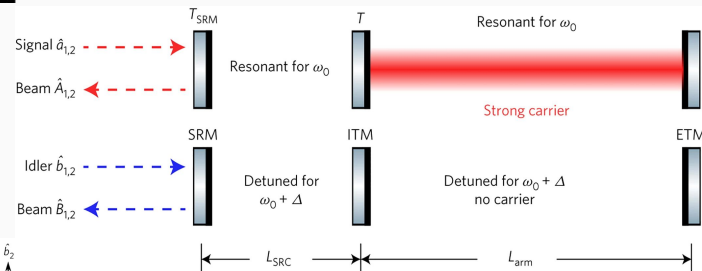
- 2 squeezed beams to be handled



2. The **idler** acquires frequency-dependence in the ITF due to its detuning



1. **Signal** and **idler** are vacuum squeezed beams, EPR-entangled and detuned by Δ



3. Combined measurement transfers the frequency dependence to the **signal** via **EPR entanglement**

Credits: [Ma et al. - Nature Phys 13, 776–780 \(2017\)](#)