

Light Colored Scalars as Messengers of Up-Quark, Down-Quark and Charged Lepton Flavor Dynamics in Grand Unified Theories*

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LES RENCONTRES DE PHYSIQUE DE LA VALLEE D'AOSTE, La Thuile

March 2nd, 2012

*I.D., Svjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, *Phys. Lett. B* 682 (2009) 67-73; *Phys. Rev. D* 81 (2010) 055009, *Phys. Rev. D* 82 (2010) 094015.

*I.D., Jure Drobnak, Svjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, *JHEP* (2011) 1111:002.

OUTLINE

•MOTIVATION

•EXPERIMENTAL STATUS

• $d = 6$ PROTON DECAY OPERATORS
SCALAR CONTRIBUTION

•LIGHT SCALARS IN $SU(5)$

•CONCLUSIONS

MOTIVATION

Leptoquarks[#] are inherent to any theory that treats quarks and leptons on the same footing.



- UNIFICATION THEORIES (PATI-SALAM[#], $SU(5)$, $SO(10)$, $E_6\dots$)



Leptoquarks can generate proton decay.



LEPTOQUARKS \equiv QUALITATIVELY NEW PHYSICS!

[#]J. C. Pati and A. Salam, *Phys. Rev. D* 10 275-289, 1974.

MOTIVATION

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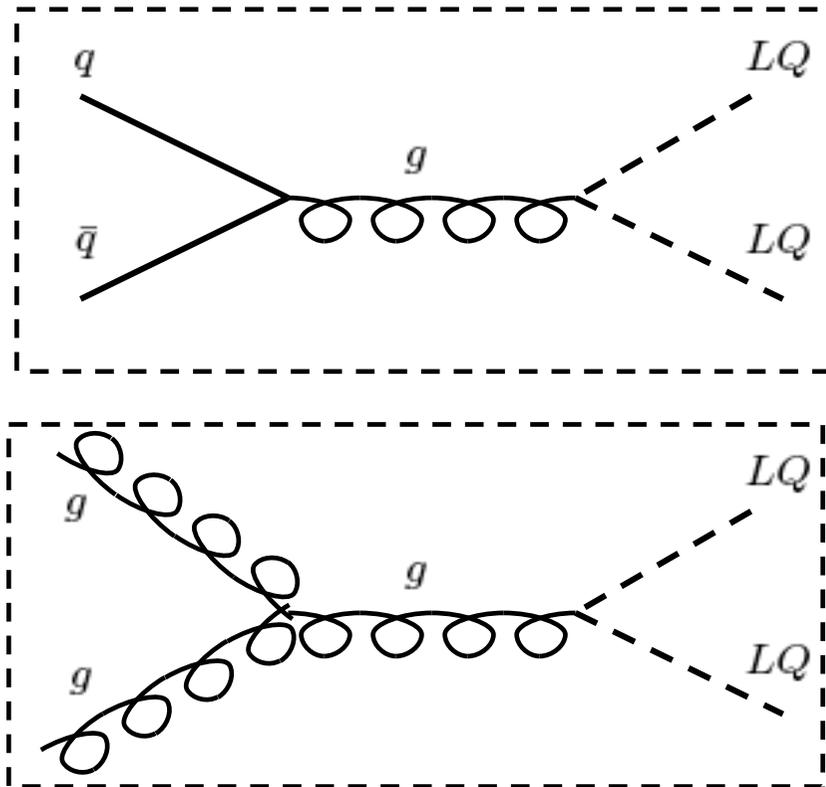


LEPTOQUARKS \equiv QUALITATIVELY NEW PHYSICS!

[#]J. C. Pati and A. Salam, *Phys. Rev. D* 10 275-289, 1974.

EXPERIMENTAL STATUS

Leptoquarks (LQ) can be produced directly in colliders.



EXPERIMENTAL LIMITS:

| | CMS | ATLAS |
|-----------------|-------------------------------|--------------------------------------|
| 1 st | $m_{LQ} > 384 \text{ GeV}^\#$ | $m_{LQ} > 660 \text{ GeV}^*$ |
| 2 nd | | $m_{LQ} > 422 \text{ GeV}^\emptyset$ |
| 3 rd | $m_{LQ} > 350 \text{ GeV}^\&$ | |

[#]V. Khachatryan et al. (CMS), Phys. Rev. Lett. 106, 201802 (2011), 1012.4031.

[∅]G. Aad et al. (ATLAS) (2011), 1104.4481.

*ATLAS (2011), arXiv:1112.4828.

&CMS PAS EXO-11-030.

EXPERIMENTAL STATUS

(PROTON DECAY)

| PROCESS | τ_p (10^{33} years) | |
|---------------------------------|-----------------------------|---|
| $p \rightarrow \pi^0 e^+$ | 8.2 | * |
| $p \rightarrow \pi^0 \mu^+$ | 6.6 | |
| $p \rightarrow K^+ \bar{\nu}$ | 2.3 | @ |
| $p \rightarrow K^0 e^+$ | 1.0 | |
| $p \rightarrow K^0 \mu^+$ | 1.3 | |
| $p \rightarrow \eta e^+$ | 0.313 | |
| $p \rightarrow \eta \mu^+$ | 0.126 | |
| $p \rightarrow \pi^+ \bar{\nu}$ | 0.025 | |
| ⋮ | ⋮ | |
| $p \rightarrow \pi^0 e^+$ | 13.0 | ¶ |
| $p \rightarrow \pi^0 \mu^+$ | 11.0 | |
| $p \rightarrow K^+ \bar{\nu}$ | 4.0 | |



*[Super-Kamiokande Collaboration], arXiv:0903.0676.

@[Super-Kamiokande Collaboration], arXiv:hep-ex/0502026.

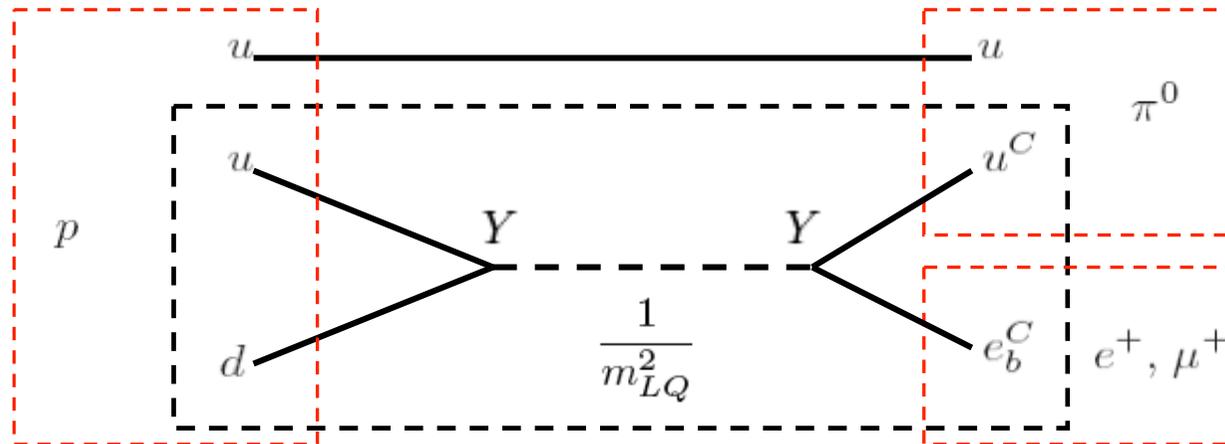
¶ <http://www.phys.utk.edu/blv2011/sessions01-06.html> (Makoto Miura)

$d=6$ PROTON DECAY OPERATORS

(SCALAR CONTRIBUTIONS*)

PROTON DECAY MEDIATING LEPTOQUARKS SHOULD BE VERY HEAVY!

$$\Gamma_6 \sim \frac{Y^4}{m_{LQ}^4} m_p^5 \quad \rightarrow \quad m_{LQ} > 10^{12} \text{ GeV}$$

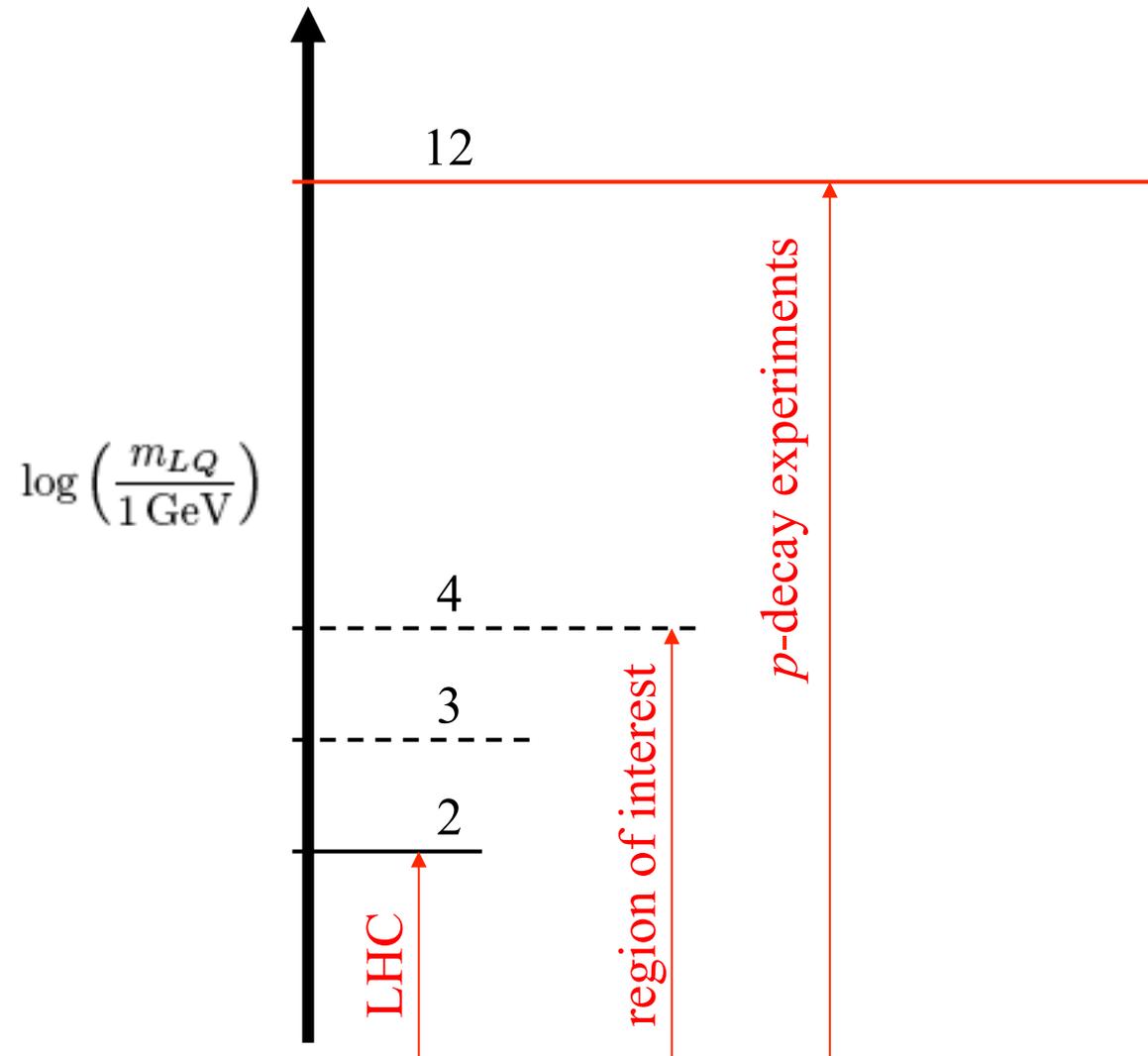


$Y \equiv$ Yukawa coupling(s)

$m_{LQ} \equiv$ Leptoquark mass

*S. Weinberg, *Phys. Rev. D* 22:1694, 1980.

RELEVANT SCALES



CASE STUDY: AN $SU(5)$ SCENARIO*

FERMIONS OF THE STANDARD MODEL (SM $\equiv SU(3) \times SU(2) \times U(1)$):

$$L_a \equiv (\mathbf{1}, \mathbf{2}, -1/2)_a = (\nu_a \quad e_a)^T$$

$$e_a^C \equiv (\mathbf{1}, \mathbf{1}, 1)_a$$

LEPTONS

$$Q_a \equiv (\mathbf{3}, \mathbf{2}, 1/6)_a = (u_a \quad d_a)^T$$

$$u_a^C \equiv (\bar{\mathbf{3}}, \mathbf{1}, -2/3)_a$$

$$d_a^C \equiv (\bar{\mathbf{3}}, \mathbf{1}, 1/3)_a$$

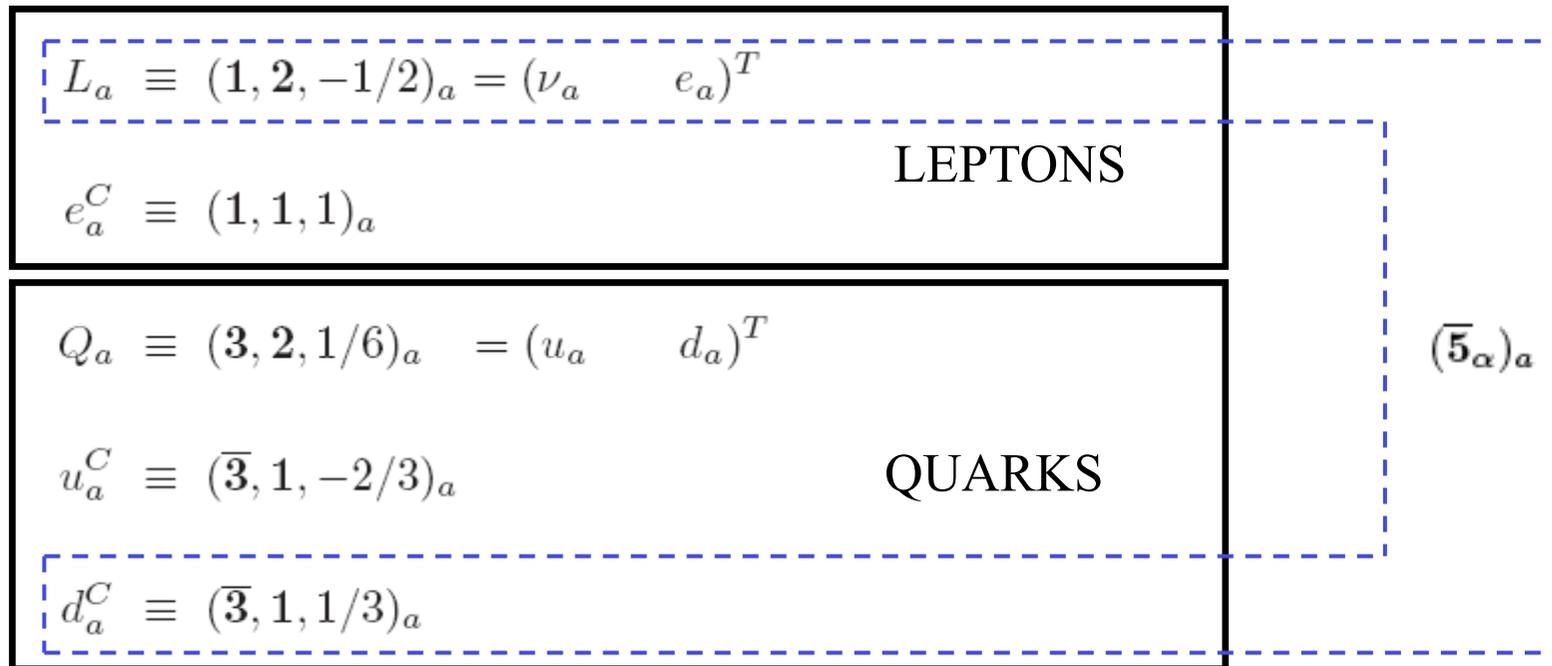
QUARKS

$a = 1, 2, 3$
FAMILY INDEX

*H. Georgi and S.L. Glashow (1974).

CASE STUDY: AN $SU(5)$ SCENARIO*

FERMIONS OF THE STANDARD MODEL:



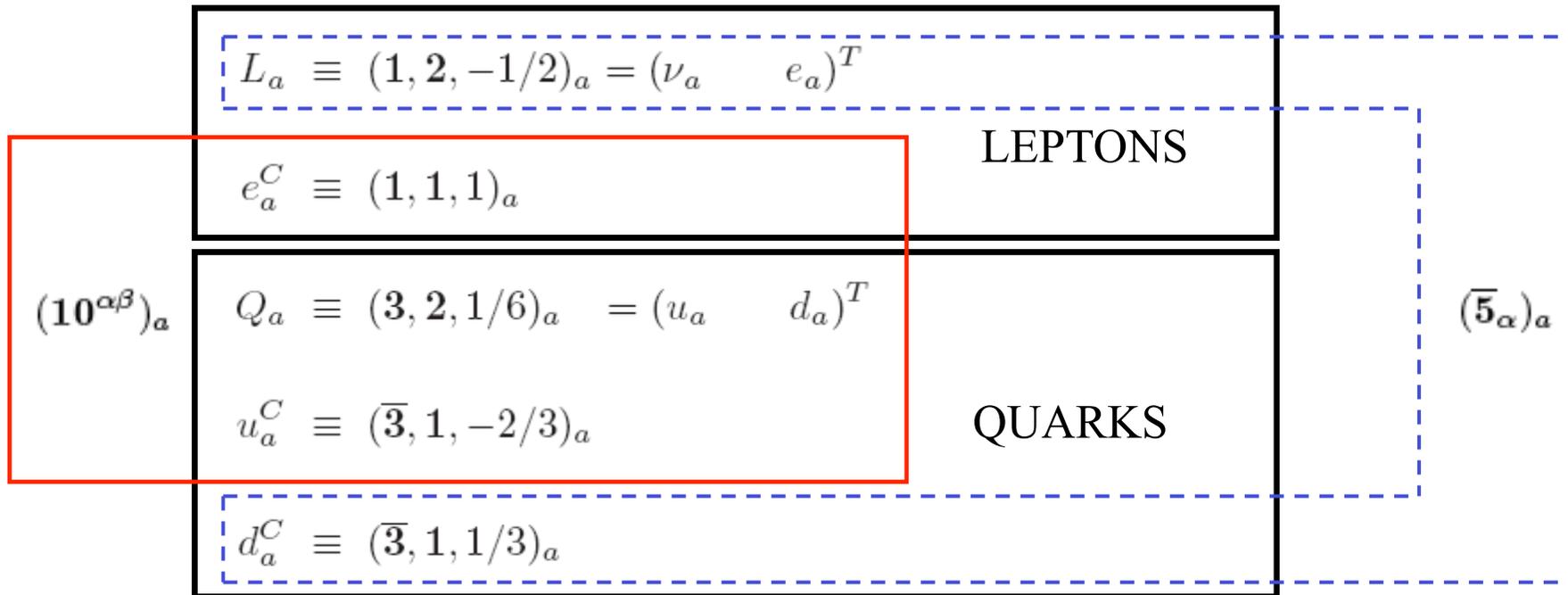
$a = 1, 2, 3$
 FAMILY INDEX

$\alpha, \beta = 1, 2, 3, 4, 5$
 $SU(5)$ GROUP INDICES

*H. Georgi and S.L. Glashow (1974).

CASE STUDY: AN $SU(5)$ SCENARIO*

FERMIONS OF THE STANDARD MODEL:



$a = 1, 2, 3$
FAMILY INDEX

$\alpha, \beta = 1, 2, 3, 4, 5$
 $SU(5)$ GROUP INDICES

*H. Georgi and S.L. Glashow (1974).

FERMION MASSES

(SCALAR REPRESENTATIONS IN $SU(5)$)

$$10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$$

$$\bar{5} \times \bar{5} = \bar{10} \oplus \bar{15} : M_N$$

$$10 \times 10 = \bar{5} \oplus \bar{45} \oplus \bar{50} : M_U$$

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5

10

15

45

50

FERMION MASSES

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$$10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$$

$$\bar{5} \times \bar{5} = \cancel{10} \oplus \bar{15} : M_N$$

$$10 \times 10 = \bar{5} \oplus 45 \oplus \cancel{50} : M_U$$

5

~~10~~

15

45

~~50~~

FERMION MASSES

(SCALAR REPRESENTATIONS IN $SU(5)$)

$$10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$$

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$$10 \times 10 = \bar{5} \oplus \bar{45} \oplus \cancel{50} : M_U$$

5

~~10~~

15

45

~~50~~



$$(10)_i(\bar{5})_j 5^*$$

$$(\bar{5})_i(\bar{5})_j 15$$

$$(10)_i(\bar{5})_j 45^*$$

$$(10)_i(10)_j 5$$

$$(10)_i(10)_j 45$$

$i, j = 1, 2, 3$
FAMILY INDICES

FERMION MASSES

(SCALAR REPRESENTATIONS IN $SU(5)$)

$$10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$$

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5

~~10~~

15

45

~~50~~



| | | | |
|-------------------------|-----------------------------|--------------------------|------------|
| $(10)_i(\bar{5})_j 5^*$ | | $(10)_i(\bar{5})_j 45^*$ | M_E, M_D |
| | $(\bar{5})_i(\bar{5})_j 15$ | | M_N |
| $(10)_i(10)_j 5$ | | $(10)_i(10)_j 45$ | M_U |

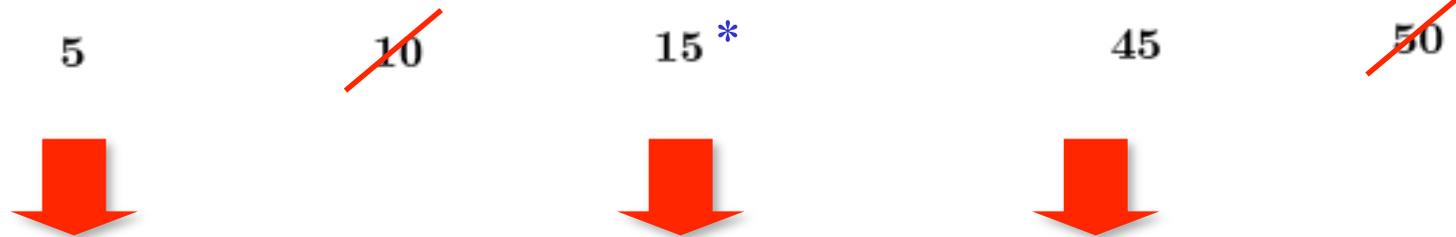
$$i, j = 1, 2, 3$$

FAMILY INDICES

FERMION MASSES

(SCALAR REPRESENTATIONS IN $SU(5)$)

| | | |
|----------------------------------------------|--------------------------------------------------------------|-------------------------------------------------------------------|
| $10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$ | $\bar{5} \times \bar{5} = \cancel{10} \oplus \bar{15} : M_N$ | $10 \times 10 = \bar{5} \oplus \bar{45} \oplus \cancel{50} : M_U$ |
|----------------------------------------------|--------------------------------------------------------------|-------------------------------------------------------------------|



| | | | |
|-------------------------|-----------------------------|--------------------------|------------|
| $(10)_i(\bar{5})_j 5^*$ | $(\bar{5})_i(\bar{5})_j 15$ | $(10)_i(\bar{5})_j 45^*$ | M_E, M_D |
| $(10)_i(10)_j 5$ | $(\bar{5})_i(\bar{5})_j 15$ | $(10)_i(10)_j 45$ | M_U |

| |
|------------------------------------|
| $i, j = 1, 2, 3$ FAMILY INDICES |
|------------------------------------|

*I.D., Pavel Fileviez Pérez, *Nucl. Phys. B* 723 (2005) 53-76.

FERMION MASSES

(SCALAR REPRESENTATIONS IN $SU(5)$)

$$10 \times \bar{5} = 5 \oplus 45 : M_E, M_D$$

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$$10 \times 10 = \bar{5} \oplus \bar{45} \oplus \cancel{50} : M_U$$

5



45



| | | | |
|-------------------------|--|--------------------------|------------|
| $(10)_i(\bar{5})_j 5^*$ | | $(10)_i(\bar{5})_j 45^*$ | M_E, M_D |
| | | | |
| $(10)_i(10)_j 5$ | | $(10)_i(10)_j 45$ | M_U |

LEPTOQUARKS IN $SU(5)$

$$\mathbf{5} = (D, T)$$

$$D = (1, 2, 1/2)$$

$$T = (\mathbf{3}, 1, -1/3)$$

$$\mathbf{45} = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

$$\Delta_1 = (8, 2, 1/2)$$

$$\Delta_2 = (\bar{\mathbf{6}}, 1, -1/3)$$

$$\Delta_3 = (\mathbf{3}, \mathbf{3}, -1/3)$$

$$\Delta_4 = (\bar{\mathbf{3}}, 2, -7/6)$$

$$\Delta_5 = (\mathbf{3}, 1, -1/3)$$

$$\Delta_6 = (\bar{\mathbf{3}}, 1, 4/3)$$

$$\Delta_7 = (1, 2, 1/2)$$

LEPTOQUARKS IN $SU(5)$

$$\mathbf{5} = (D, T)$$

$$D = (1, 2, 1/2)$$

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--- \equiv Higgs doublet

$$\mathbf{45} = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

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LEPTOQUARKS IN $SU(5)$

$$\mathbf{5} = (D, T)$$

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$$T = (3, 1, -1/3)$$

--- \equiv Higgs doublet

— \equiv “genuine” leptoquark

$$\mathbf{45} = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

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$$\Delta_3 = (3, 3, -1/3)$$

$$\Delta_4 = (\bar{3}, 2, -7/6)^*$$

$$\Delta_5 = (3, 1, -1/3)$$

$$\Delta_6 = (\bar{3}, 1, 4/3)$$

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*I.D., Sjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, *Phys. Lett. B* 682 (2009) 67-73.

LEPTOQUARKS IN $SU(5)$

(p -DECAY MEDIATING LEPTOQUARK)

$$5 = (D, T)$$

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$$T = (3, 1, -1/3)$$

$$45 = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

$$\Delta_1 = (8, 2, 1/2)$$

$$\Delta_2 = (\bar{6}, 1, -1/3)$$

$$\Delta_3 = (3, 3, -1/3)$$

$$\Delta_4 = (\bar{3}, 2, -7/6)$$

$$\Delta_5 = (3, 1, -1/3)$$

$$\Delta_6 = (\bar{3}, 1, 4/3)$$

$$\Delta_7 = (1, 2, 1/2)$$

--- \equiv Higgs doublet

— \equiv “genuine” leptoquark

□ \equiv p -decay mediating leptoquark

LEPTOQUARKS IN $SU(5)$

(p -DECAY MEDIATING LEPTOQUARK)

$$5 = (D, T)$$

$$D = (1, 2, 1/2)$$

$$T = (3, 1, -1/3)$$



$$45 = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

$$\Delta_1 = (8, 2, 1/2)$$

$$\Delta_2 = (\bar{6}, 1, -1/3)$$

$$\Delta_3 = (3, 3, -1/3)$$

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$$\Delta_5 = (3, 1, -1/3)$$

$$\Delta_6 = (\bar{3}, 1, 4/3)$$

$$\Delta_7 = (1, 2, 1/2)$$



--- \equiv Higgs doublet

— \equiv “genuine” leptoquark

□ \equiv p -decay mediating leptoquark

LEPTOQUARKS IN $SU(5)$

(p -DECAY MEDIATING SCALAR LEPTOQUARKS)

$$\mathbf{5} = (D, T)$$

$$D = (1, 2, 1/2)$$

$$T = (\mathbf{3}, 1, -1/3)$$

$$\mathbf{10} = (\Psi_a, \Psi_b, \Psi_c)$$

$$\Psi_a = (1, 1, 1)$$

$$\Psi_b = (\bar{\mathbf{3}}, 1, -2/3)$$

$$\Psi_c = (\mathbf{3}, 2, 1/6)$$

$$\mathbf{45} = (\Delta_1, \Delta_2, \Delta_3, \Delta_4, \Delta_5, \Delta_6, \Delta_7)$$

$$\Delta_1 = (8, 2, 1/2)$$

$$\Delta_2 = (\bar{\mathbf{6}}, 1, -1/3)$$

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$$\Delta_4 = (\bar{\mathbf{3}}, 2, -7/6)$$

$$\Delta_5 = (\mathbf{3}, 1, -1/3)$$

$$\Delta_6 = (\bar{\mathbf{3}}, 1, 4/3)$$

$$\Delta_7 = (1, 2, 1/2)$$

\square $\equiv p$ -decay mediating leptoquark

LEPTOQUARKS IN $SU(5)$

(p -DECAY MEDIATING SCALAR LEPTOQUARKS)

ALL IN ALL, THERE ARE EIGHTEEN (FIFTEEN) PROTON DECAY
MEDIATING SCALARS IF NEUTRINOS ARE DIRAC (MAJORANA)!*

*I.D., Svjetlana Fajfer and Nejc Košnik, work in progress.

LEPTOQUARKS IN THE 5 OF $SU(5)$ *

(p -DECAY MEDIATING SCALAR LEPTOQUARKS)

*I.D., Svjetlana Fajfer and Nejc Košnik, work in progress.

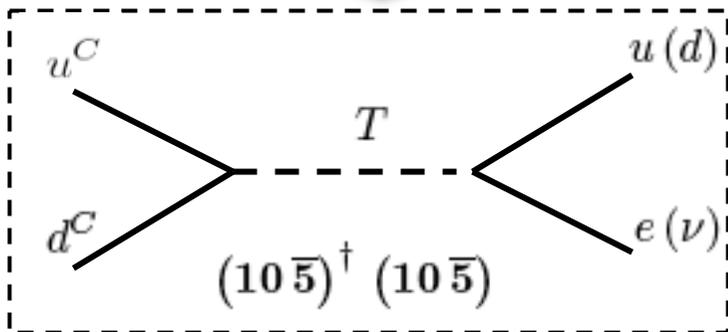
$$(10^{\alpha\beta})_i (\bar{5}_\alpha)_j 5_\beta^*$$



$$M_E = M_D^T$$



$$\begin{array}{l} u_{ai}^{CT} C d_{bj}^C T_c^* \\ u_{ai}^T C e_j T_a^* \\ d_{ai}^T C \nu_j T_a^* \end{array}$$



$$\underline{T}$$

$$(3, 1, -1/3)$$

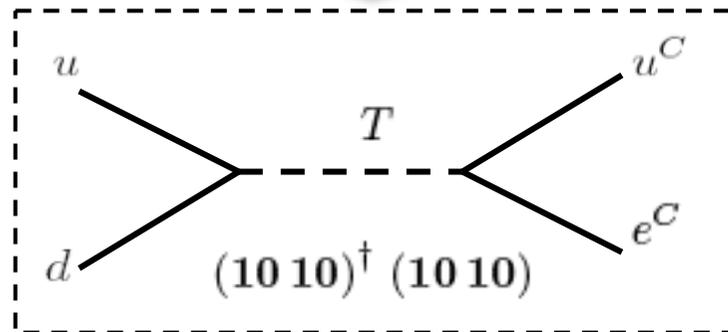
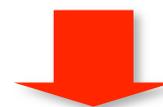
$$\epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\gamma\delta})_j (5)^\epsilon$$

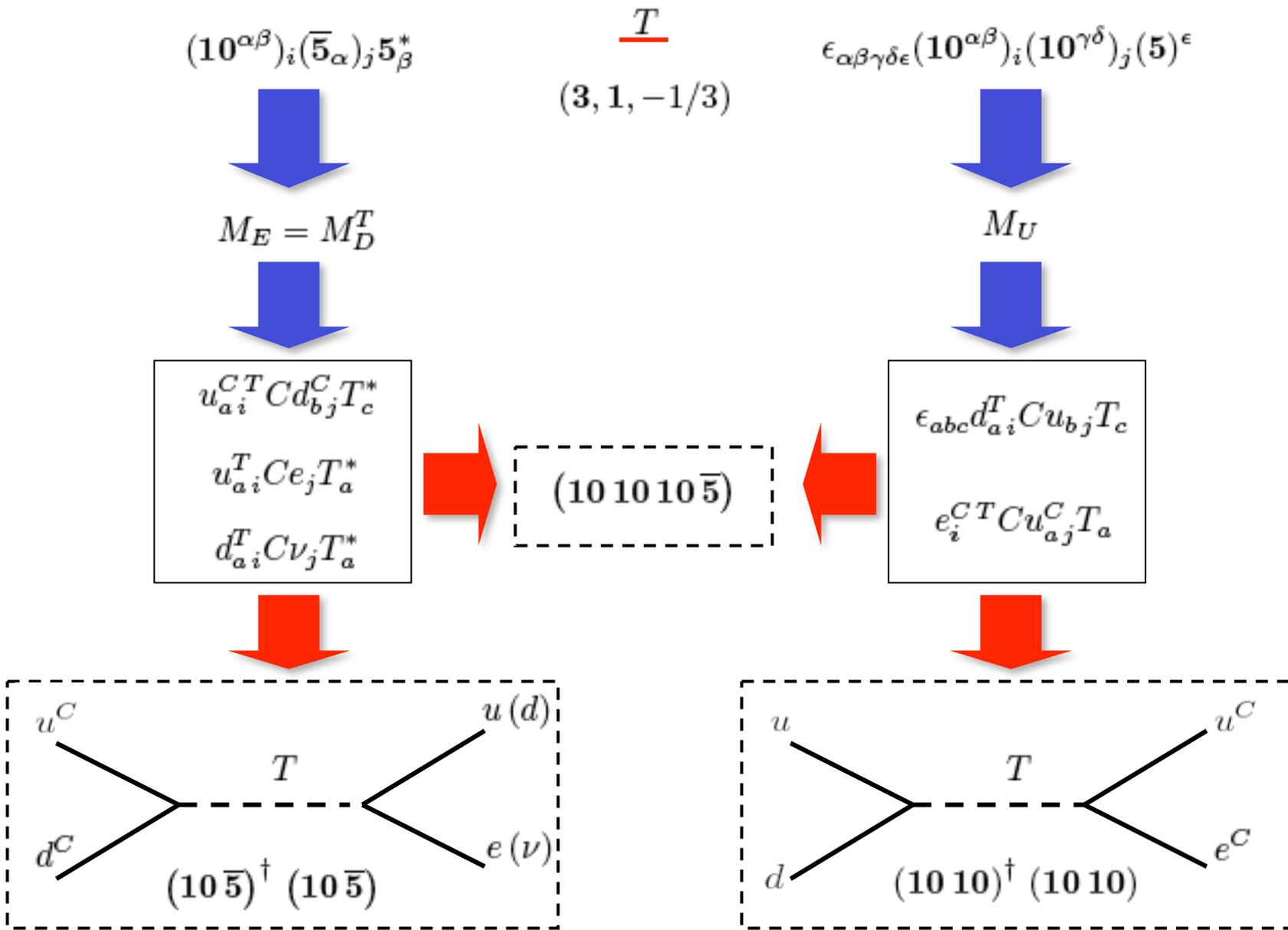


$$M_U$$



$$\begin{array}{l} \epsilon_{abc} d_{ai}^T C u_{bj} T_c \\ e_i^{CT} C u_{aj} T_a \end{array}$$





LEPTOQUARKS IN THE 5 OF $SU(5)$ (p -DECAY MEDIATING SCALAR LEPTOQUARKS)

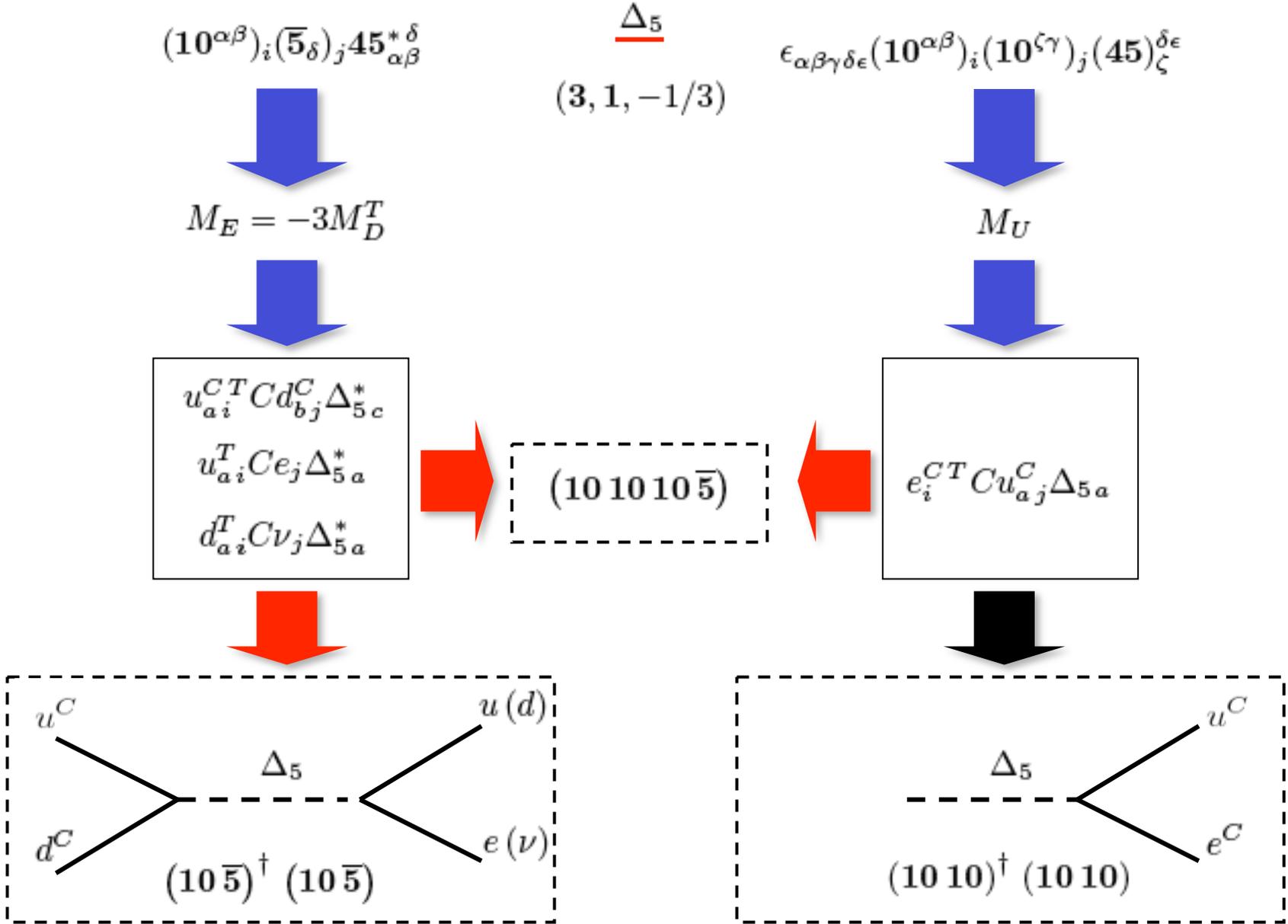
$$\Gamma(p \rightarrow e^+ \pi^0) \sim \frac{\alpha^2}{v_5^4 m_\Delta^4} \left| \frac{3}{8} (V_{UD})_{11} (V_{UD})_{13} m_\tau m_b \right|^2$$

$$\Gamma(p \rightarrow \mu^+ \pi^0) \sim \frac{\alpha^2}{v_5^4 m_\Delta^4} \left| \frac{3}{8} (V_{UD})_{11} (V_{UD})_{12} m_\tau m_s \right|^2$$

*I.D., Svjetlana Fajfer and Nejc Košnik, work in progress.

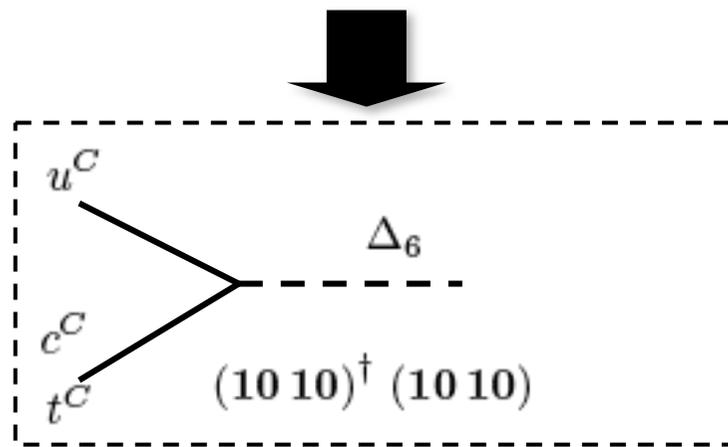
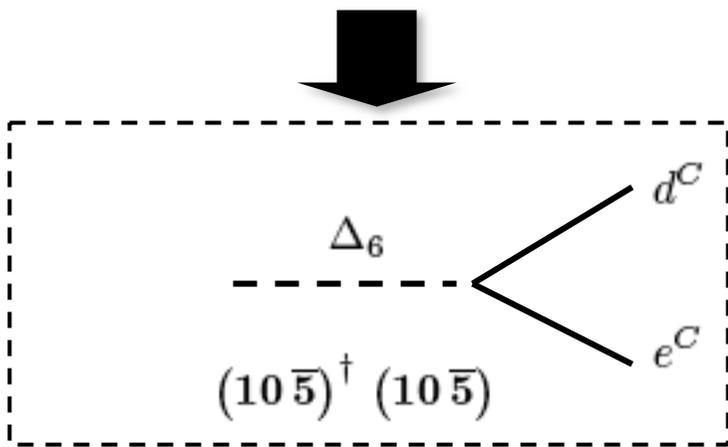
LEPTOQUARKS IN THE 45 OF $SU(5)$ *
(p -DECAY MEDIATING SCALAR LEPTOQUARKS)

*I.D., Svjetlana Fajfer and Nejc Košnik, work in progress.



$$\begin{array}{ccc}
 (10^{\alpha\beta})_i (\bar{5}_\delta)_j 45_{\alpha\beta}^{*\delta} & \begin{array}{c} \underline{\Delta_6} \\ (\bar{3}, 1, 4/3) \end{array} & \epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\zeta\gamma})_j (45)_\zeta^{\delta\epsilon} \\
 \downarrow & & \downarrow \\
 M_E = -3M_D^T & & M_U \\
 \downarrow & & \downarrow
 \end{array}$$

$$\boxed{e_i^{CT} C d_{aj}^C \Delta_{6a}^*} \quad \rightleftarrows \quad (10 \ 10 \ 10 \bar{5}) \quad \rightleftarrows \quad \boxed{\epsilon_{abc} u_{ia}^{CT} C u_{bj}^C \Delta_{6c}}$$



$$(10^{\alpha\beta})_i (\bar{5}_\delta)_j 45_{\alpha\beta}^{*\delta}$$



$$M_E = -3M_D^T$$



$$e_i^{CT} C d_{aj}^C \Delta_{6a}^*$$



$(g-2)_\mu$ ANOMALY

$$\underline{\Delta_6}$$

$(\bar{3}, 1, 4/3)$

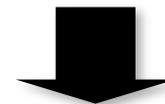
$$\epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\zeta\gamma})_j (45)_\zeta^{\delta\epsilon}$$



$$M_U$$



$$\epsilon_{abc} u_{ia}^{CT} C u_{bj}^C \Delta_{6c}$$



$t\bar{t}$ ASYMMETRY

$$\underline{\Delta_6} \quad \epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\zeta\gamma})_j (45)_{\zeta}^{\delta\epsilon}$$

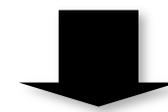
$$(\bar{3}, 1, 4/3)$$



M_U



$$Y = -Y^T \quad \epsilon_{abc} u_{ia}^{CT} C u_{bj}^C \Delta_{6c} *$$



$t\bar{t}$ ASYMMETRY

*I.D., Svjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, *Phys. Rev. D* 81 (2010) 055009, *Phys. Rev. D* 82 (2010) 094015.

Δ_6 LEPTOQUARK: UP-QUARK SECTOR

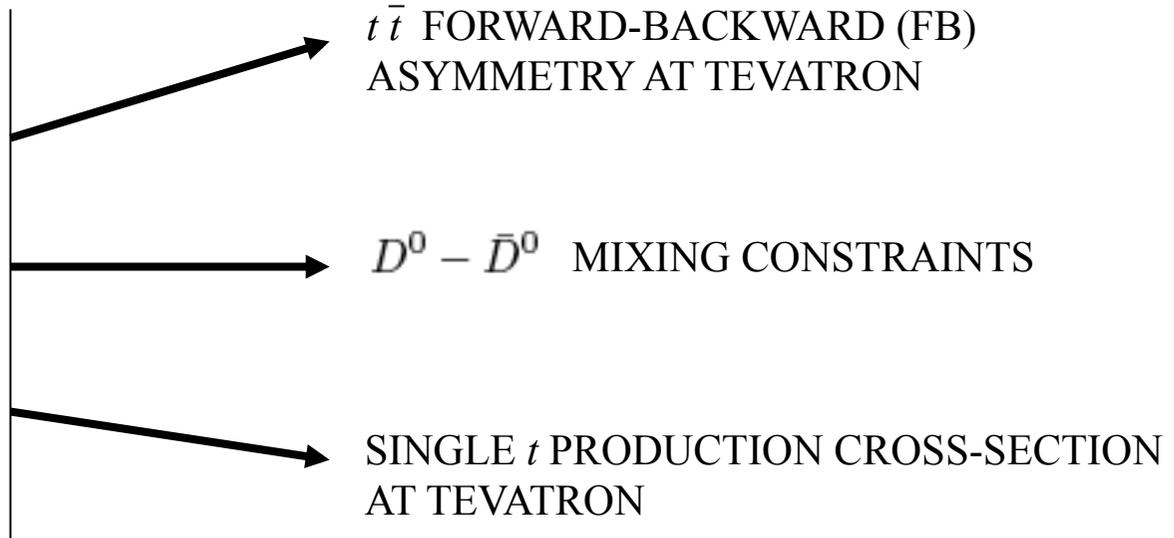
CONSTRAINTS ON Y ORIGINATE FROM
THE UP-QUARK PHENOMENOLOGY!

$$m_{\Delta_6} = 400 \text{ GeV}$$

$$|Y_{13}| = 1.9$$

$$|Y_{23}| \leq 0.0033$$

$$|Y_{12}| \leq 0.042$$



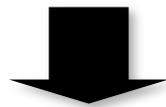
$$(10^{\alpha\beta})_i (\bar{5}_\delta)_j 45_{\alpha\beta}^{*\delta}$$



$$M_E = -3M_D^T$$



$$e_i^{CT} C d_{aj}^C \Delta_{6a}^*$$



$(g-2)_\mu$ ANOMALY

$$\underline{\Delta_6}$$

$(\bar{3}, 1, 4/3)$

$$\epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\zeta\gamma})_j (45)_\zeta^{\delta\epsilon}$$



$$M_U$$

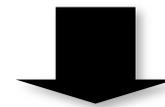


$SU(5)$

$SO(10)$

$45 \in 120$

$$\epsilon_{abc} u_{ia}^{CT} C u_{bj}^C \Delta_{6c}$$



$t\bar{t}$ ASYMMETRY

$$(10^{\alpha\beta})_i (\bar{5}_\delta)_j 45_{\alpha\beta}^{*\delta}$$



$$M_E = -3M_D^T$$



$$e_i^{CT} C d_{aj}^C \Delta_{6a}^*$$



$(g-2)_\mu$ ANOMALY

$$\underline{\Delta_6}$$

$(\bar{3}, 1, 4/3)$

$$\epsilon_{\alpha\beta\gamma\delta\epsilon} (10^{\alpha\beta})_i (10^{\zeta\gamma})_j (45)_\zeta^{\delta\epsilon}$$



$$M_U$$



$SU(5)$

$SO(10)$

$45 \in 126$

$$\epsilon_{abc} u_{ia}^{CT} C v_{bj}^C \Delta_{6c}$$



$t\bar{t}$ ASYMMETRY

$$(10^{\alpha\beta})_i (\bar{5}_\delta)_j 45_{\alpha\beta}^{*\delta}$$

$$\underline{\Delta_6} \in 126$$

$$(\bar{3}, 1, 4/3)$$



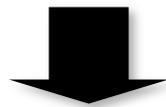
$$M_E = -3M_D^T$$



Y

$$e_i^{CT} C d_{aj}^C \Delta_{6a}^*$$

*



$(g - 2)_\mu$ ANOMALY

*I.D., Jure Drobnak, Svjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, , JHEP (2011) 1111:002.

Δ_6 LEPTOQUARK: THE DOWN-QUARK AND CHARGED LEPTON SECTORS

$$|Y^{(1\sigma)}| \in \left(\begin{array}{ccc} < 1.4 \times 10^{-6} & < 8.7 \times 10^{-5} & < 4.1 \times 10^{-4} \\ < 3.6 \times 10^{-3} \cup [2.1, 2.9] & < 3.6 \times 10^{-3} \cup [2.1, 2.9] & < 6.2 \times 10^{-4} \cup [2.2, 2.8] \\ < 5.6 \times 10^{-3} & < 8.1 \times 10^{-3} & < 9.6 \times 10^{-3} \end{array} \right)$$

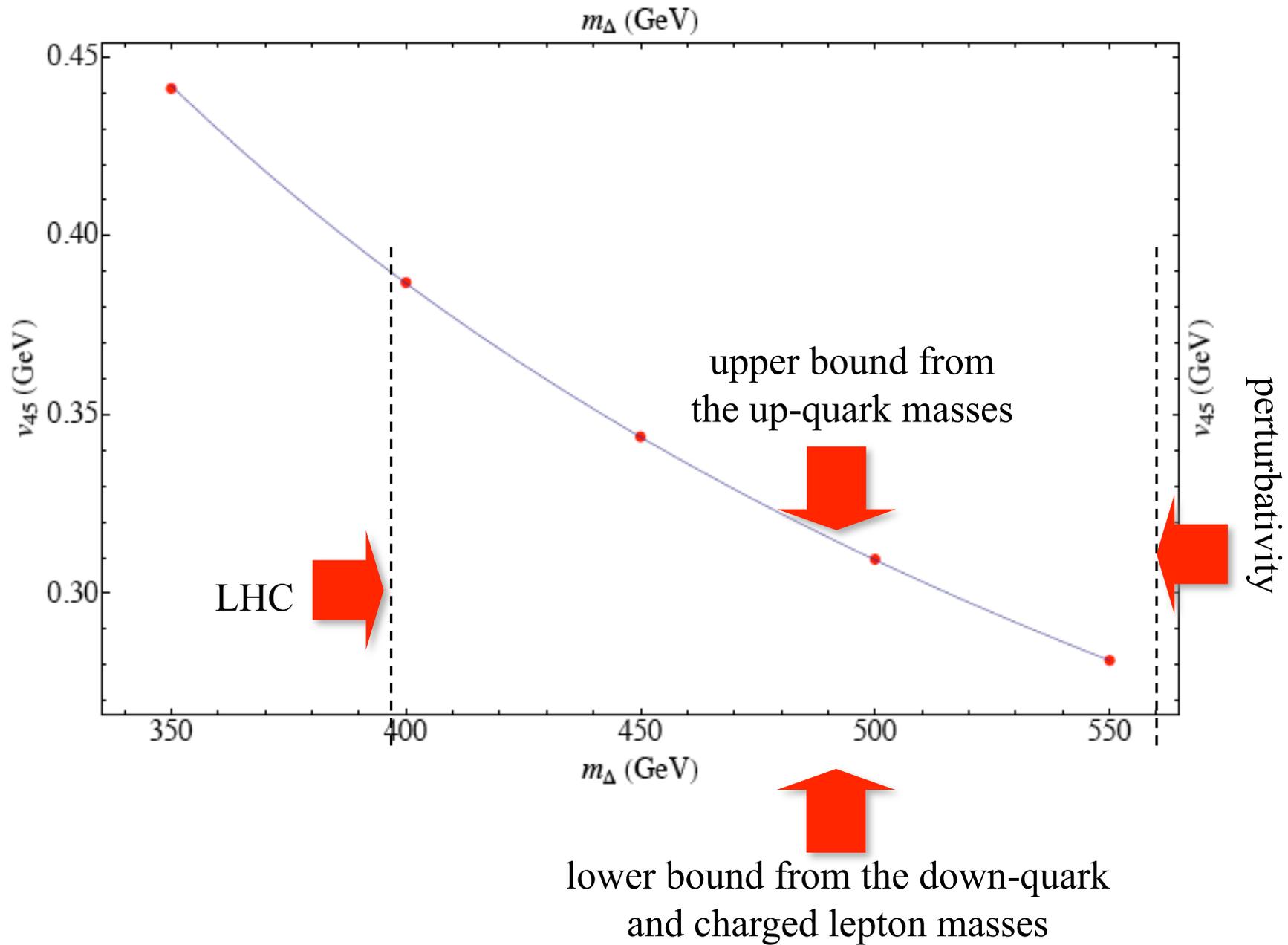
*

III

$$\begin{pmatrix} 0 & 0 & 0 \\ \blacksquare & 0 & 0 \\ \bullet & \bullet & \bullet \end{pmatrix}, \quad \begin{pmatrix} 0 & 0 & 0 \\ 0 & \blacksquare & 0 \\ \bullet & \bullet & \bullet \end{pmatrix}, \quad \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \blacksquare \\ \bullet & \bullet & \bullet \end{pmatrix}$$

*I.D., Jure Drobnak, Svjetlana Fajfer, Jernej F. Kamenik and Nejc Košnik, , JHEP (2011) 1111:002.

IMPLICATIONS FOR THE UP-QUARK SECTOR



CONCLUSIONS

Leptoquark states represent qualitatively new physics.

Proton decay operators induced via scalar leptoquark exchanges exhibit strong model dependence.

That feature opens up possibilities for existence of light leptoquark states with interesting phenomenological consequences without conflict with proton stability.

CONCLUSIONS

Light scalar leptoquarks can, for example address the issue of $(g - 2)_\mu$ anomaly or $t\bar{t}$ asymmetry.

Scenarios that incorporate light leptoquarks could thus be directly probed at colliders.

THANK YOU!

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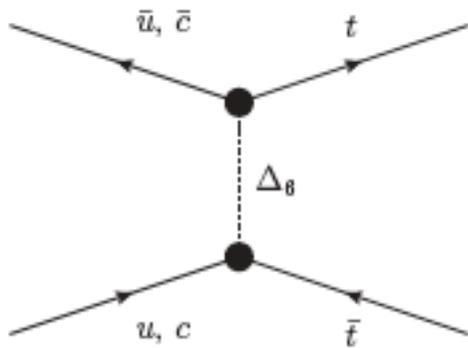
FORWARD-BACKWARD ASYMMETRY

SIMULTANEOUS FIT TO THE INTEGRATED CROSS SECTION σ^{exp} AND A_{FB}

$$A_{FB}^{\text{exp}} - A_{FB}^{\text{SM}} = (14.2 \pm 6.9)\%$$

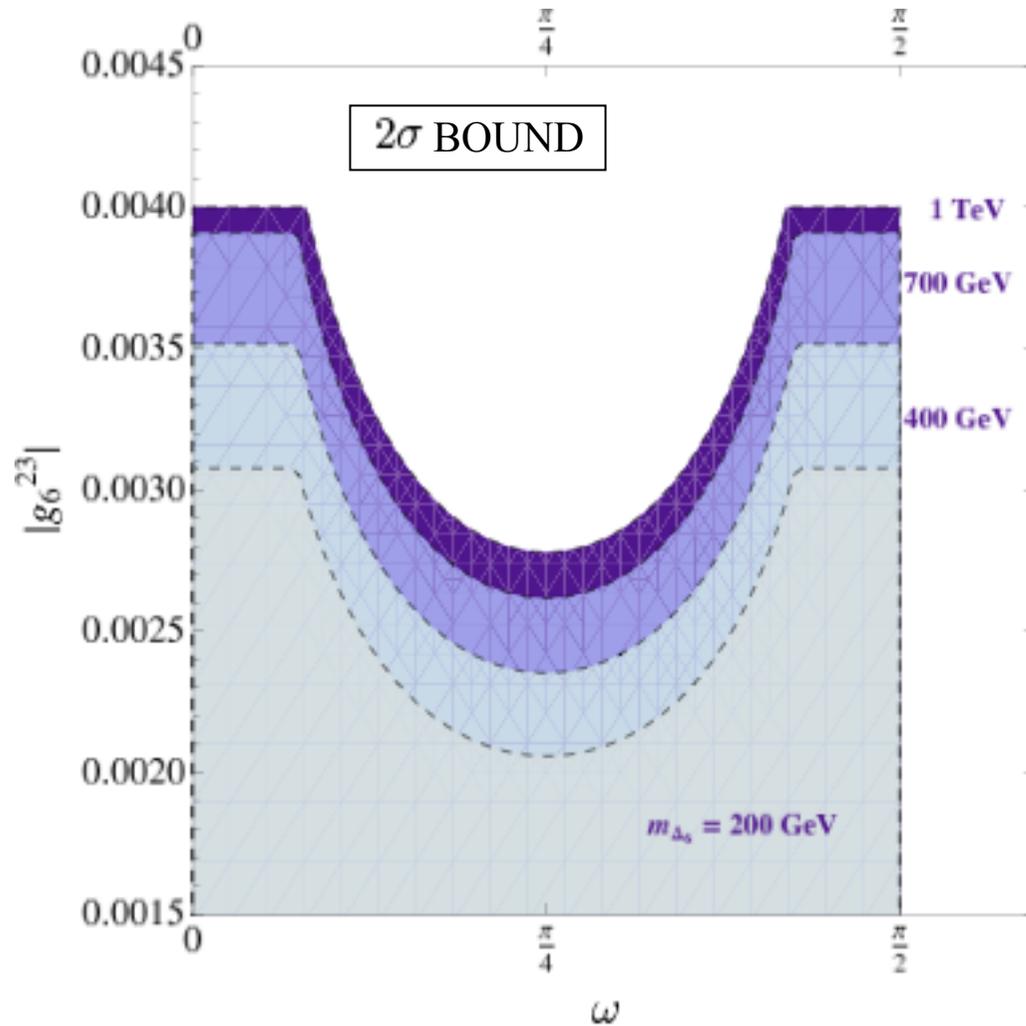
$$\sigma^{\text{exp}} = 7.0 \pm 0.6 \text{ pb}$$

$$|g_6^{13}| = 0.9(2) + 2.5(4) \frac{m_{\Delta_6}}{1 \text{ TeV}}$$



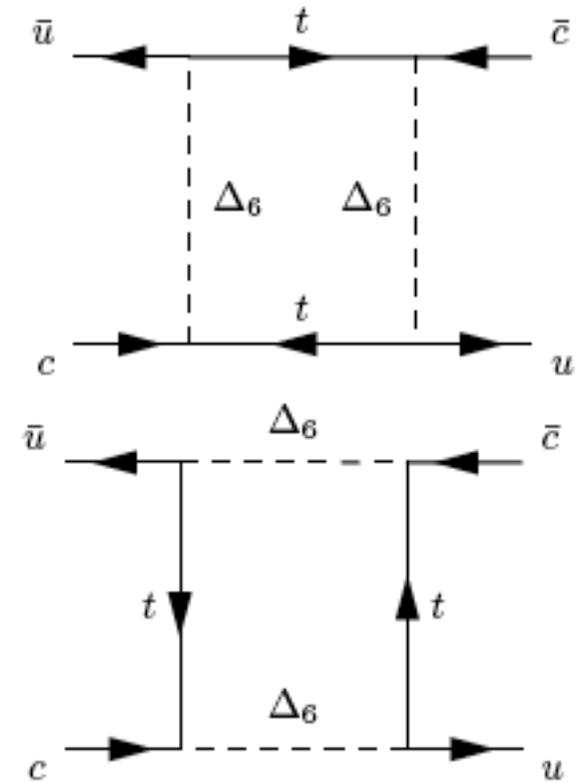
LEADING CONTRIBUTIONS TO $t \bar{t}$
PRODUCTION CROSS SECTION AND A_{FB}
AT TEVATRON.

$D^0 - \bar{D}^0$ MIXING CONSTRAINTS



$\omega \equiv$ RELATIVE PHASE BETWEEN g_6^{13} AND g_6^{23}

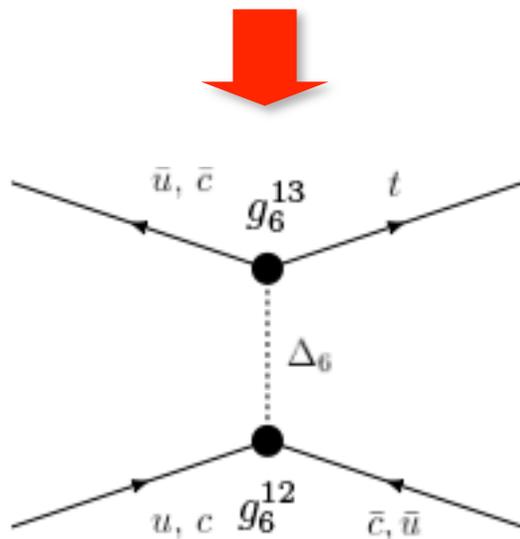
$(g_6^{13} g_6^{23*})^2$ CONTRIBUTIONS TO $|\Delta C| = 2$:



$m_{\Delta_6} = 400 \text{ GeV}$
 $|g_6^{23}| \leq 0.0033$

SINGLE t PRODUCTION CROSS-SECTION

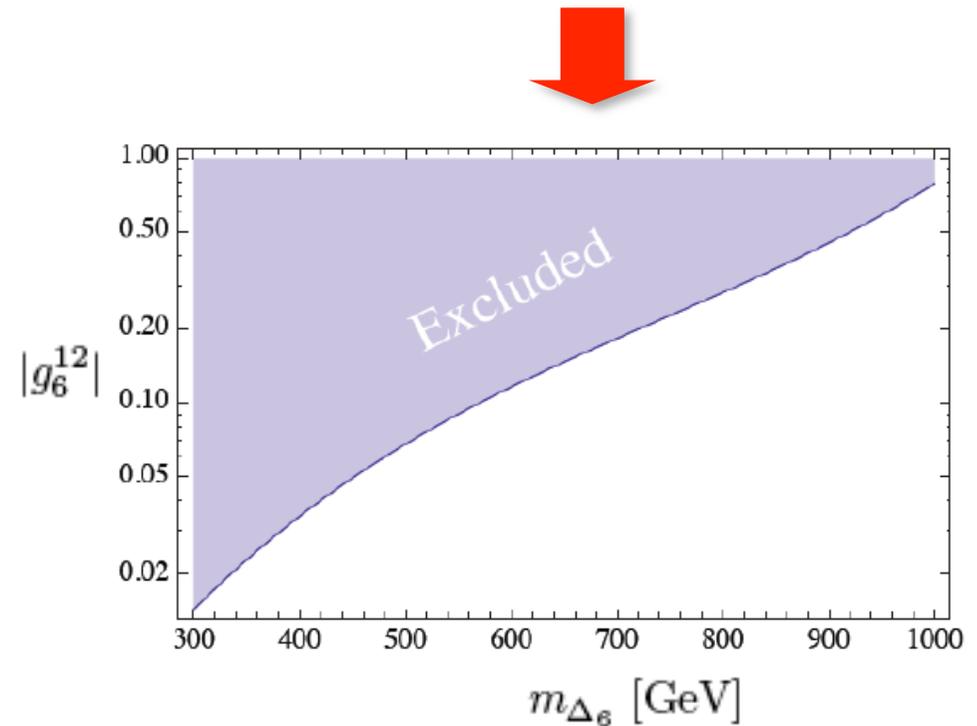
THE SINGLE t PRODUCTION IS SENSITIVE TO THE PRODUCT $|g_6^{12} g_6^{13*}|$.



$$m_{\Delta_6} = 400 \text{ GeV}$$
$$|g_6^{12}| \leq 0.042$$

TEVATRON RESULT: $\sigma_{1t} = 2.76_{-0.47}^{+0.58} \text{ pb}$

WE REQUIRE: $\sigma_{1t}^{\Delta_6} \leq 1 \text{ pb}$



| decay mode | 90 % C.L. exp. bound on \mathcal{B} | 1σ upper bound in units $(m_\Delta/400 \text{ GeV})^4$ |
|------------------------------------|---------------------------------------|------------------------------------------------------------------------------|
| $B_d \rightarrow e^- e^+$ | 8.3×10^{-8} | $ Y_{eb}Y_{ed}^* ^2 < 4.4$ |
| $B_d \rightarrow \mu^- \mu^+$ | 4.2×10^{-9} | $ Y_{\mu b}Y_{\mu d}^* ^2 < 5.0 \times 10^{-6}$ |
| $B_d \rightarrow \tau^- \tau^+$ | 4.1×10^{-3} | $ Y_{\tau b}Y_{\tau d}^* ^2 < 1.3 \times 10^{-2}$ |
| $B_s \rightarrow e^- e^+$ | 2.8×10^{-7} | $ Y_{eb}Y_{es}^* ^2 < 10.1$ |
| $B_s \rightarrow \mu^- \mu^+$ | 1.2×10^{-8} | $ Y_{\mu b}Y_{\mu s}^* ^2 < 1.1 \times 10^{-5}$ |
| $B_d \rightarrow e^\mp \mu^\pm$ | 6.4×10^{-8} | $ Y_{eb}Y_{\mu d}^* ^2 + Y_{\mu b}Y_{ed}^* ^2 < 1.6 \times 10^{-4}$ |
| $B_d \rightarrow \mu^\mp \tau^\pm$ | 2.2×10^{-5} | $ Y_{\mu b}Y_{\tau d}^* ^2 + Y_{\tau b}Y_{\mu d}^* ^2 < 2.2 \times 10^{-4}$ |
| $B_d \rightarrow \tau^\mp e^\pm$ | 2.8×10^{-5} | $ Y_{\tau b}Y_{ed}^* ^2 + Y_{eb}Y_{\tau d}^* ^2 < 2.7 \times 10^{-4}$ |
| $B_s \rightarrow e^\mp \mu^\pm$ | 2.0×10^{-7} | $ Y_{eb}Y_{\mu s}^* ^2 + Y_{\mu b}Y_{es}^* ^2 < 3.4 \times 10^{-4}$ |

$$\delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}} = (2.87 \pm 0.93) \times 10^{-9}$$

| | | |
|-------------------|----------------------------------------|----------|
| $ \epsilon_K $ | $2.228(11) \times 10^{-3}$ | [23] |
| Δm_K | $3.483(6) \times 10^{-15} \text{ GeV}$ | [23] |
| ϕ_ϵ | $43.5(7)^\circ$ | [23] |
| f_K | $0.1560(11) \text{ GeV}$ | [27] |
| \hat{B}_K | $0.725(26)$ | [27] |
| κ_ϵ | $0.94(2)$ | [48] |
| η_1 | $1.31^{(+25)}_{(-22)}$ | [49] |
| η_2 | $0.57(1)$ | [46, 50] |
| η_3 | $0.496(47)$ | [51] |

| decay mode | 90% C.L. exp. bound on \mathcal{B} | 1σ upper bound in units $(m_\Delta/400 \text{ GeV})^4$ |
|----------------------------------------|--------------------------------------|------------------------------------------------------------------------------|
| $B^+ \rightarrow \pi^+ \ell^- \ell^+$ | 4.9×10^{-8} | $ Y_{eb}Y_{ed}^* ^2 + Y_{\mu b}Y_{\mu d}^* ^2 < 3.0 \times 10^{-7}$ |
| $B^+ \rightarrow \pi^+ e^\pm \mu^\mp$ | 1.7×10^{-7} | $ Y_{eb}Y_{\mu d}^* ^2 + Y_{\mu b}Y_{ed}^* ^2 < 1.1 \times 10^{-6}$ |
| $B^+ \rightarrow K^+ e^\pm \mu^\mp$ | 9.1×10^{-8} | $ Y_{eb}Y_{\mu s}^* ^2 + Y_{\mu b}Y_{es}^* ^2 < 4.3 \times 10^{-7}$ |
| $B^+ \rightarrow K^+ \tau^\pm \mu^\mp$ | 7.7×10^{-5} | $ Y_{\tau b}Y_{\mu s}^* ^2 + Y_{\mu b}Y_{\tau s}^* ^2 < 5.7 \times 10^{-4}$ |

| decay mode | 90% C.L. exp. bound on \mathcal{B} | 1σ upper bound in units $(m_\Delta/400 \text{ GeV})^4$ |
|-----------------------------|--------------------------------------|-------------------------------------------------------------------------------|
| $\tau \rightarrow e\pi^0$ | 8.0×10^{-8} | $ Y_{ed}Y_{\tau d}^* ^2 < 1.9 \times 10^{-4}$ |
| $\tau \rightarrow \mu\pi^0$ | 1.1×10^{-7} | $ Y_{\mu d}Y_{\tau d}^* ^2 < 2.7 \times 10^{-4}$ |
| $\tau \rightarrow eK_S$ | 3.3×10^{-8} | $ Y_{ed}Y_{\tau s}^* - Y_{es}Y_{\tau d}^* ^2 < 3.2 \times 10^{-5}$ |
| $\tau \rightarrow \mu K_S$ | 4.0×10^{-8} | $ Y_{\mu d}Y_{\tau s}^* - Y_{\mu s}Y_{\tau d}^* ^2 < 4.0 \times 10^{-5}$ |
| $\tau \rightarrow \mu\eta$ | 6.5×10^{-8} | $ 0.69 Y_{\mu d}Y_{\tau d}^* - Y_{\mu s}Y_{\tau s}^* ^2 < 1.3 \times 10^{-4}$ |

| decay mode | 90% C.L. exp. bound on \mathcal{B} | 1σ upper bound in units $(m_\Delta/400 \text{ GeV})^4$ |
|------------------------------|--------------------------------------|-----------------------------------------------------------------|
| $\mu \rightarrow e\gamma$ | 2.4×10^{-12} | $ \sum_{i=d,s,b} Y_{ei}Y_{\mu i}^* ^2 < 4.6 \times 10^{-8}$ |
| $\tau \rightarrow \mu\gamma$ | 4.4×10^{-8} | $ \sum_{i=d,s,b} Y_{\mu i}Y_{\tau i}^* ^2 < 4.8 \times 10^{-3}$ |
| $\tau \rightarrow e\gamma$ | 3.3×10^{-8} | $ \sum_{i=d,s,b} Y_{ei}Y_{\tau i}^* ^2 < 3.6 \times 10^{-3}$ |

