



# Performance of the Muon Detector for the LHCb Upgrade II

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PhD presentation  
XXXVI cycle



# Outline



- Studies on the design of the LHCb muon detector at high luminosity
  - Rate estimate
  - Proposal of a new readout scheme
  - Dead time and muon ID
  
- Search for the rare decay mode  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-$ 
  - Optimisation of the  $D_s$  selection
  - $D_s^+ \mu^+ \mu^-$  selection
  - Observation of new decay modes



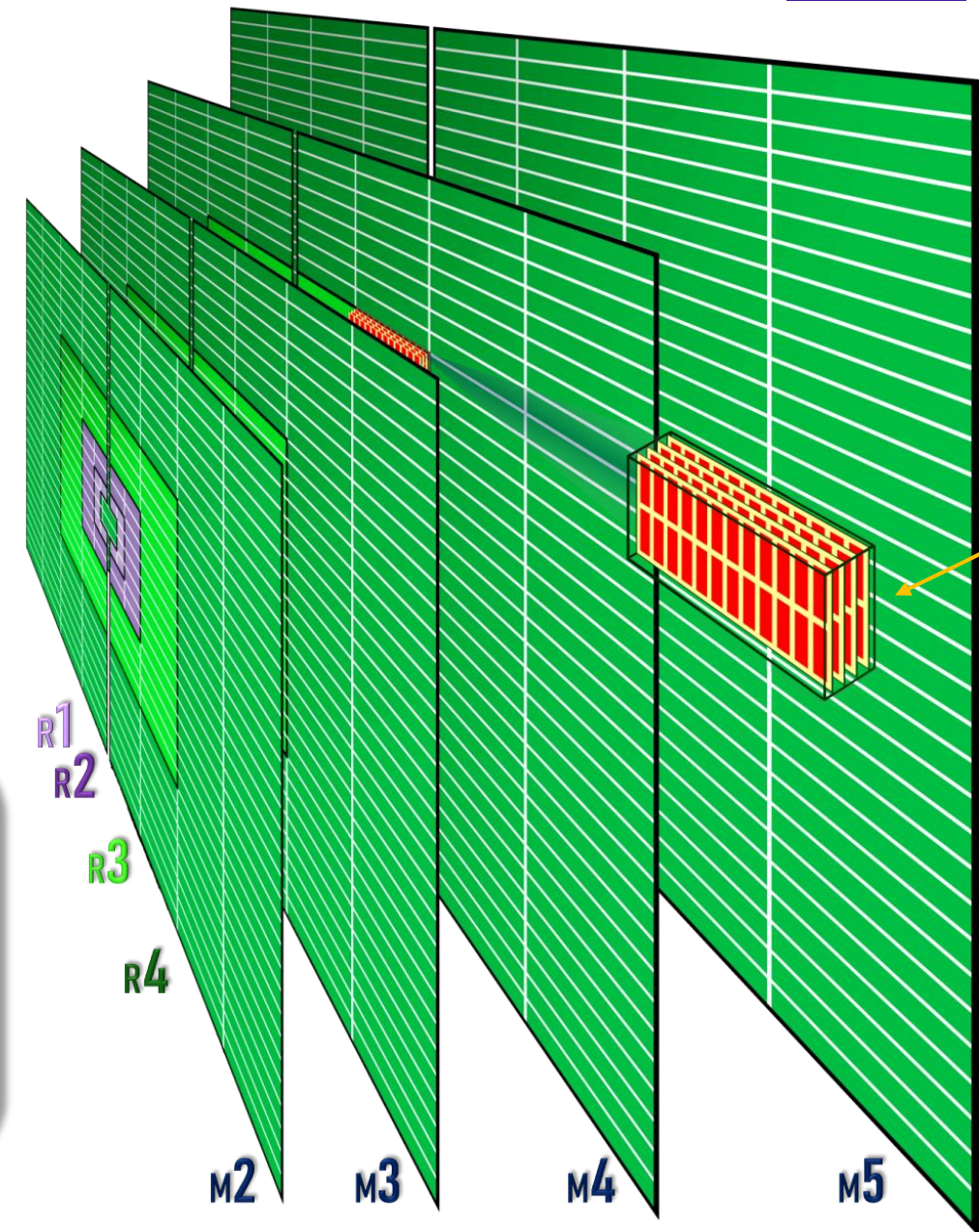
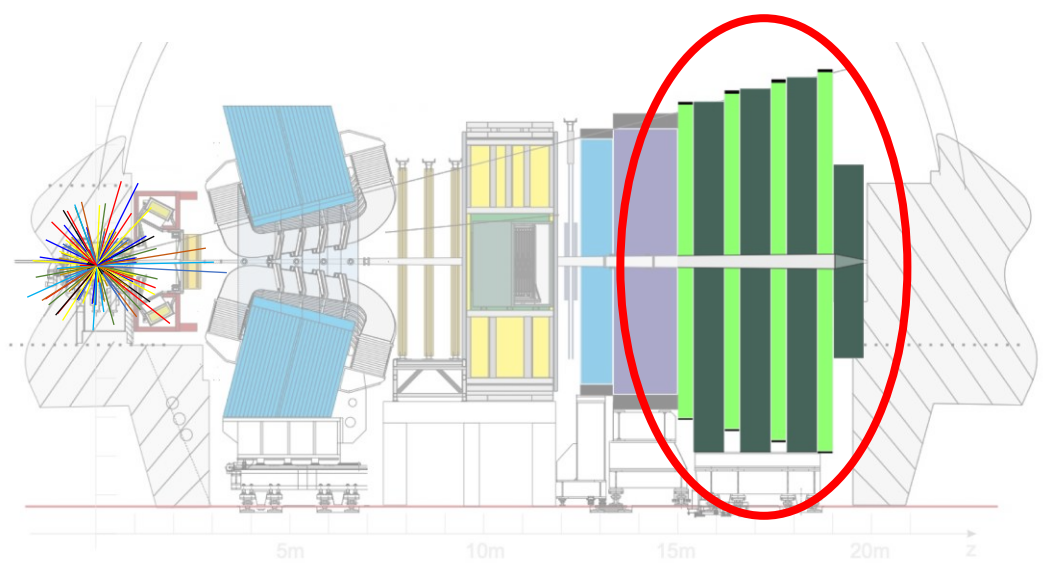
# Outline



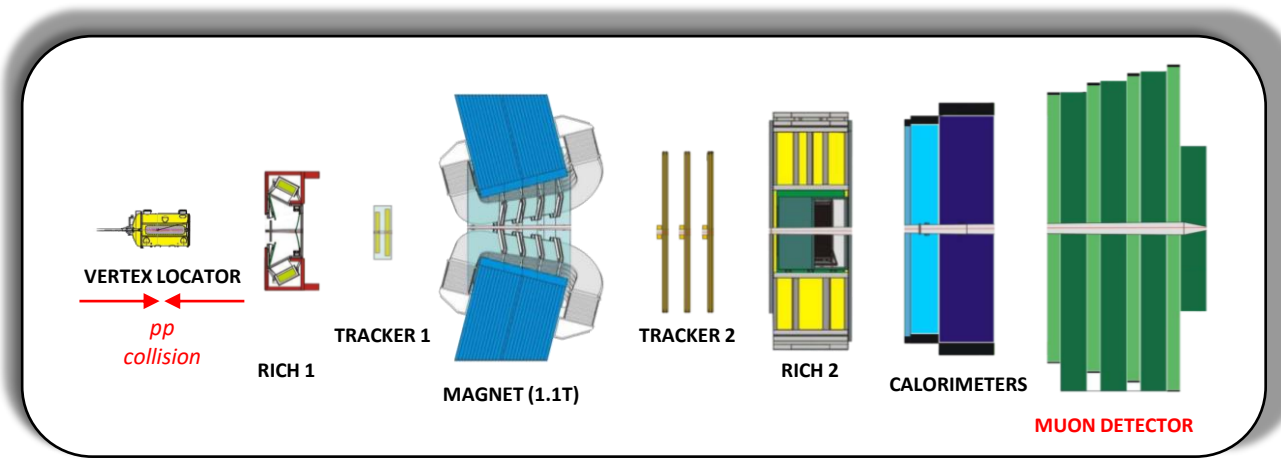
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# LHCb muon detector at Run 3

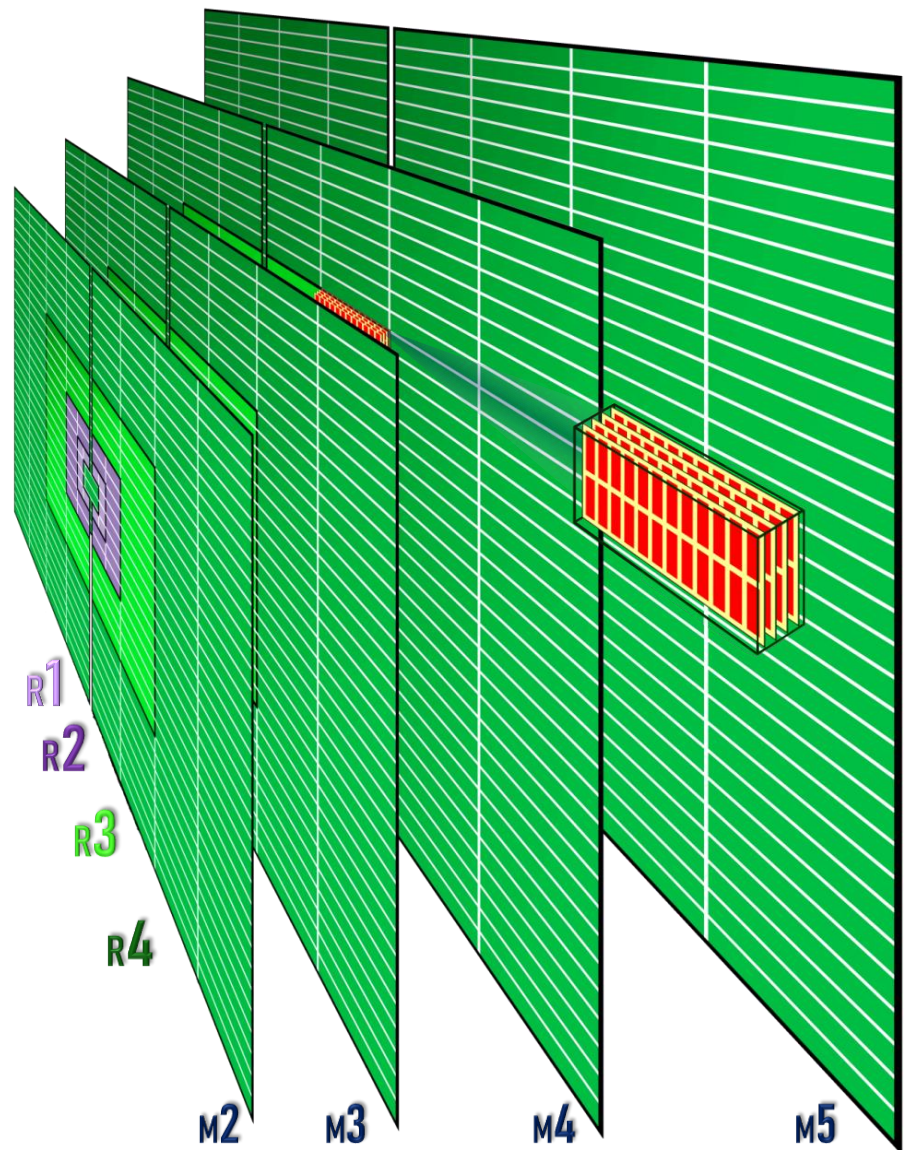


4-gaps  
MWPC  
of region  
R3M4

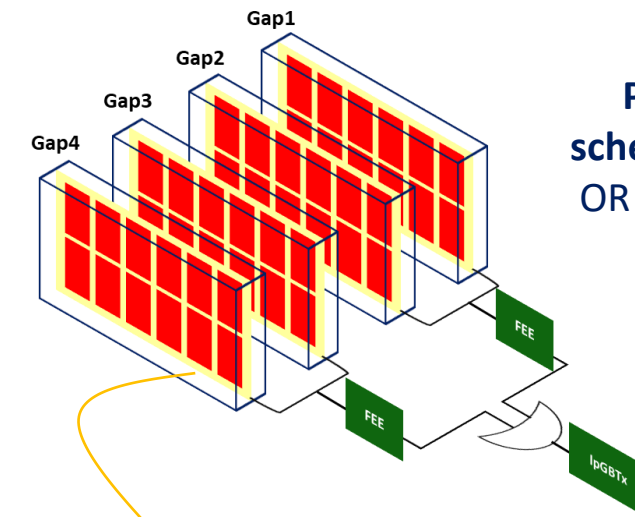
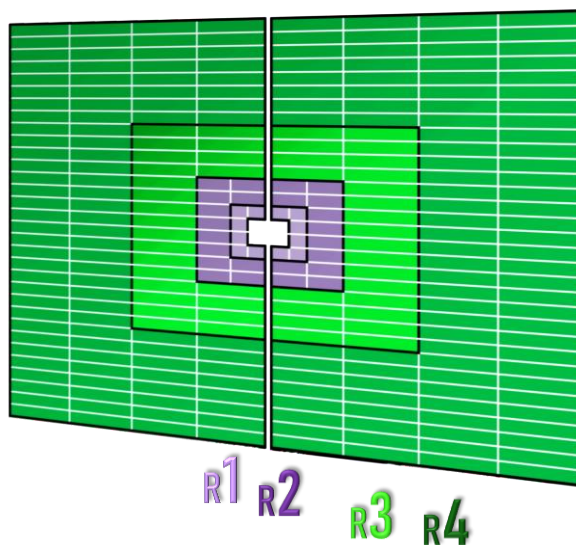




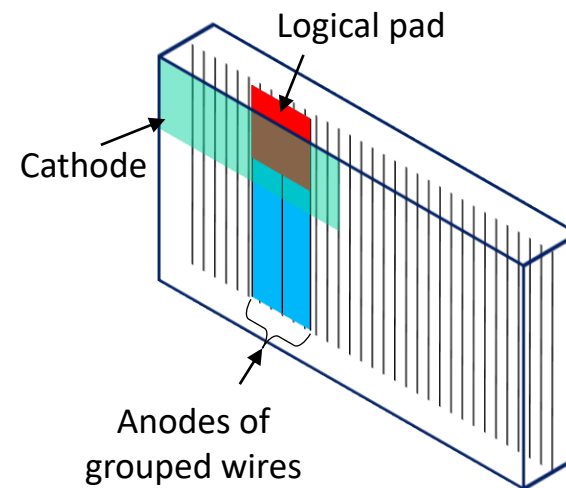
# LHCb muon detector at Run 3



4 stations of 4-gaps multi-wire proportional chambers (MWPCs) with gas mixture Ar:CO<sub>2</sub>:CF<sub>4</sub> (40:55:5). Iron absorbers 80 cm thick between stations



**Present readout scheme of a chamber:**  
OR of 2 gaps, again in OR at FEE level



In each gap logical pads are logic AND between grouped wires and strips (or cathode pads)



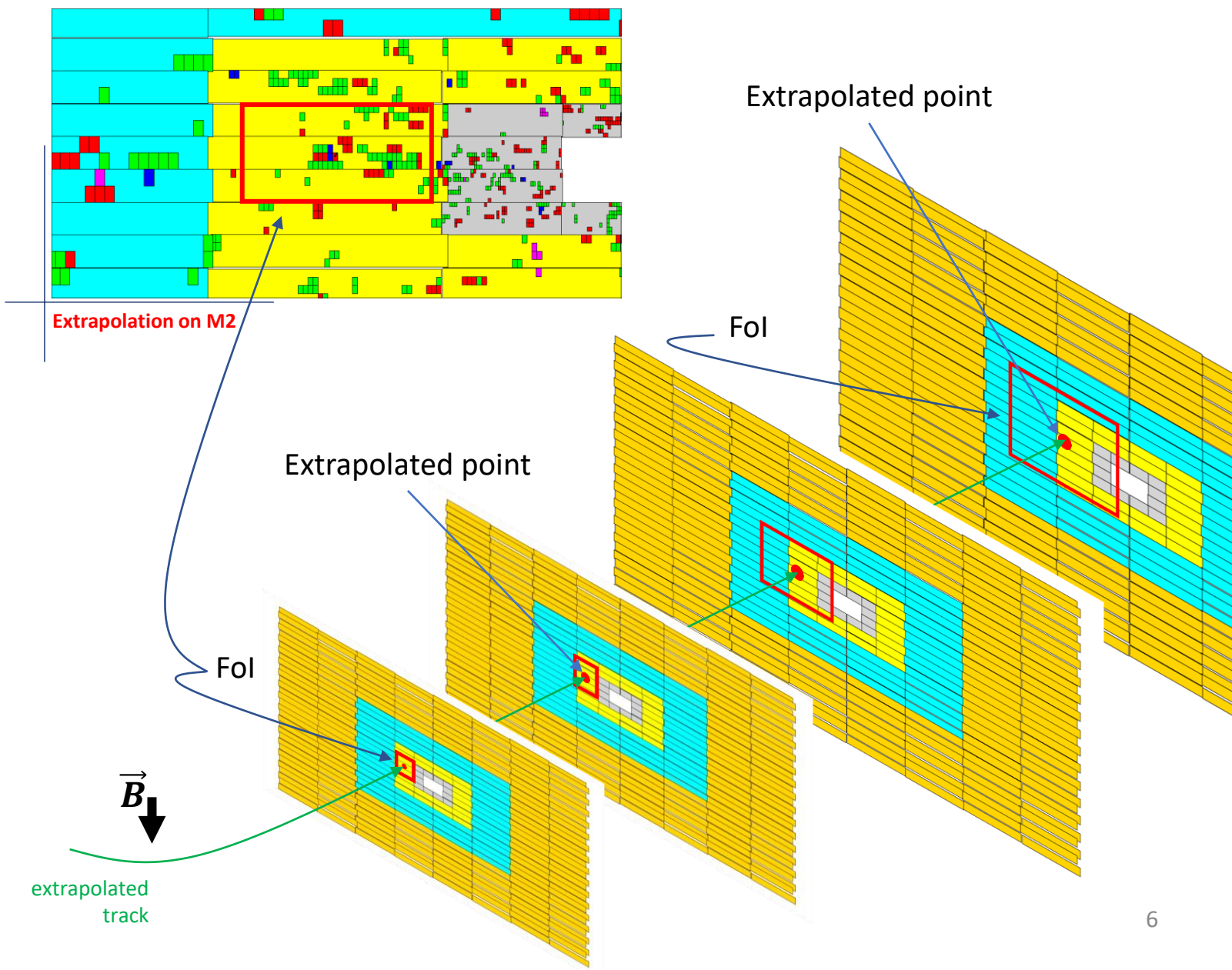
# LHCb muon detector at Run 3



The **muon identification** acts at different levels. The first one is the *IsMuon* request such that, given an extrapolated track:

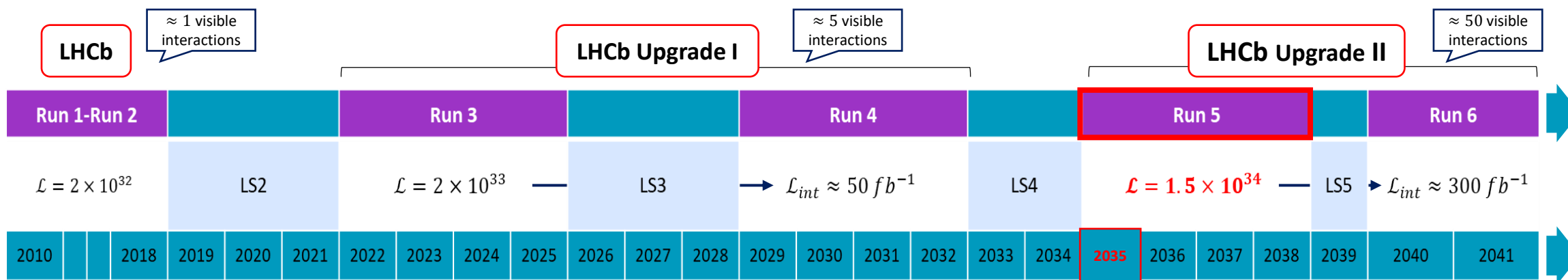
1. The nearest track hit on each station is searched in a *FoI*: area around the extrapolated point on the station
2. The presence of hits on stations is required, according to the momentum:

$p$ [GeV/c]	The track is a muon:
$p < 3$	never
$3 < p < 6$	M2 & M3
$6 < p < 10$	M2 & M3 & (M4    M5)
$p > 10$	M2 & M3 & M4 & M5





# Upgrade II of LHCb apparatus



Observable	Current LHCb (up to $9 \text{ fb}^{-1}$ )	Upgrade I ( $50 \text{ fb}^{-1}$ )	Upgrade II ( $300 \text{ fb}^{-1}$ )
<b>CKM tests</b>			
$\gamma$ ( $B \rightarrow DK$ , etc.)	$4^\circ$	$1^\circ$	$0.35^\circ$
$\phi_s$ ( $B_s^0 \rightarrow J/\psi\phi$ )	32 mrad	10 mrad	4 mrad
$ V_{ub} / V_{cb} $ ( $A_b^0 \rightarrow p\mu^-\bar{\nu}_\mu$ , etc.)	6%	2%	1%
$a_{\text{sl}}^d$ ( $B^0 \rightarrow D^-\mu^+\nu_\mu$ )	$36 \times 10^{-4}$	$5 \times 10^{-4}$	$2 \times 10^{-4}$
$a_{\text{sl}}^s$ ( $B_s^0 \rightarrow D_s^-\mu^+\nu_\mu$ )	$33 \times 10^{-4}$	$7 \times 10^{-4}$	$3 \times 10^{-4}$
<b>Charm</b>			
$\Delta A_{CP}$ ( $D^0 \rightarrow K^+K^-, \pi^+\pi^-$ )	$29 \times 10^{-5}$	$8 \times 10^{-5}$	$3.3 \times 10^{-5}$
$A_\Gamma$ ( $D^0 \rightarrow K^+K^-, \pi^+\pi^-$ )	$11 \times 10^{-5}$	$3.2 \times 10^{-5}$	$1.2 \times 10^{-5}$
$\Delta x$ ( $D^0 \rightarrow K_s^0\pi^+\pi^-$ )	$18 \times 10^{-5}$	$4.1 \times 10^{-5}$	$1.6 \times 10^{-5}$
<b>Rare Decays</b>			
$\mathcal{B}(B^0 \rightarrow \mu^+\mu^-)/\mathcal{B}(B_s^0 \rightarrow \mu^+\mu^-)$	69%	27%	11%
$S_{\mu\mu}$ ( $B_s^0 \rightarrow \mu^+\mu^-$ )	—	—	0.2
$A_\Gamma^{(2)}$ ( $B^0 \rightarrow K^{*0}e^+e^-$ )	0.10	0.043	0.016
$A_\Gamma^{\text{Im}}$ ( $B^0 \rightarrow K^{*0}e^+e^-$ )	0.10	0.043	0.016
$\mathcal{A}_{\phi\gamma}^{\Delta\Gamma}$ ( $B_s^0 \rightarrow \phi\gamma$ )	$+0.41$ $-0.44$	0.083	0.033
$S_{\phi\gamma}$ ( $B_s^0 \rightarrow \phi\gamma$ )	0.32	0.062	0.025
$\alpha_\gamma(A_b^0 \rightarrow A\gamma)$	$+0.17$ $-0.29$	0.097	0.038
<b>Lepton Universality Tests</b>			
$R_K$ ( $B^+ \rightarrow K^+\ell^+\ell^-$ )	0.044	0.017	0.007
$R_{K^*}$ ( $B^0 \rightarrow K^{*0}\ell^+\ell^-$ )	0.12	0.022	0.009
$R(D^*)$ ( $B^0 \rightarrow D^{*-}\ell^+\nu_\ell$ )	0.026	0.005	0.002

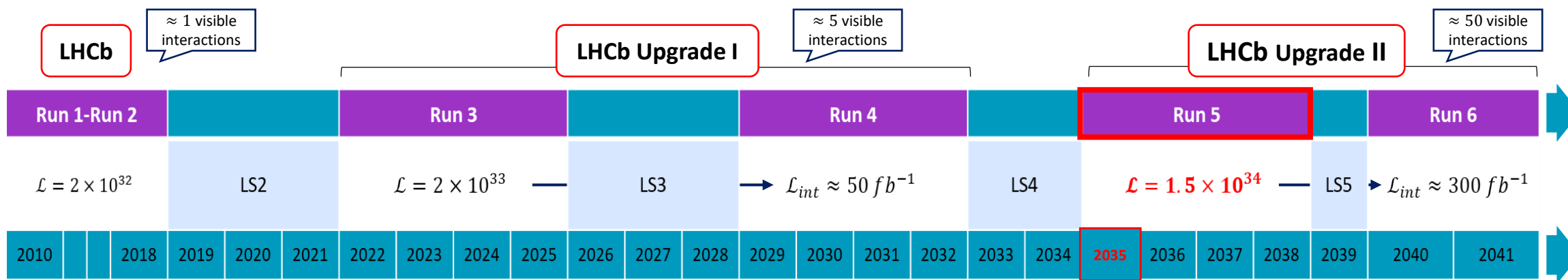
LHCb-TDR-023

A great opportunity for LHCb:

- Measurements **precision** will improve
- More sensibility to **new physics** phenomena



# Upgrade II of muon detector



- ❖ Insufficient rate capability of present detector
- ❖ Chambers granularity inadequate for the future high rates
- ❖ High pile-up conditions will cause:
  - inefficiency effects by electronics dead time
  - inefficiency effects of muon mis-identification



- **New readout scheme simulation**
- New detector ( $\mu$ RWELL) for inner regions
- **Electronics dead time effects**
- **mis-Identification effects**





# Rates estimate at Run 5

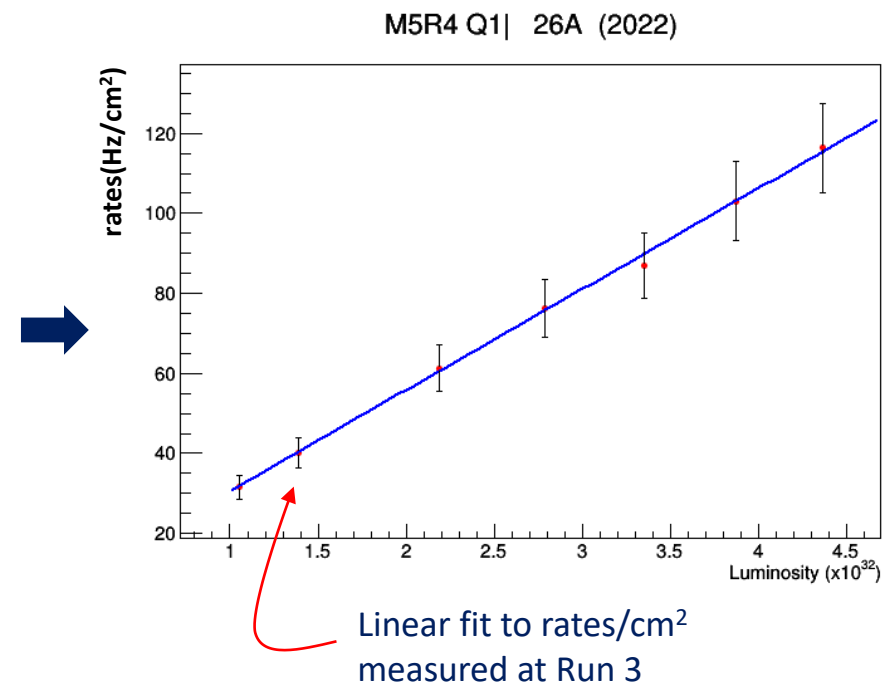


## At the LHC Run 3 start:

rates measurements of each chambers gap, at seven luminosity configurations

A linear fit for the rates extrapolation at Run 5 luminosity:  
 $L = 1.5 \times 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$

Number of visible $pp$ interactions per bunch crossing	
$\mu = 1.10$	$L = 1.05 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 1.45$	$L = 1.39 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 2.29$	$L = 2.19 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 2.92$	$L = 2.79 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 3.50$	$L = 3.35 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 4.05$	$L = 3.87 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 4.56$	$L = 4.36 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$



Maximum rates (M2) of muon detector design (**Run 1**):

	M2 (kHz/cm <sup>2</sup> )
R1	37.5
R2	26.5
R3	6.5
R4	1.2

Maximum rates extrapolated at **Run 5**, in red, reduction in the hypothesis of a **filter** upstream wrt M2

Maximum chamber rate (kHz/cm <sup>2</sup> )				
	M2	M3	M4	M5
R1	594.0 -> 344.5	274.5	203.5	232.7
R2	255.6 -> 79.2	64.2	34.1	39.0
R3	53.4 -> 19.2	8.9	6.2	8.9
R4	9.9	3.0	1.7	6.8

$\langle \frac{\sigma_R}{R} \rangle \%$

10.1%  
10.2%  
8.5%  
9.6%



# Rates estimate at Run 5



### Rates of M2 gaps (Hz/cm<sup>2</sup>)

1001	910	1247	1893	2512	1403	1025	1096
917	804	1165	1776	2110	1280	849	1036
1389	784	1307	1931	2106	1299	841	1533
1559	747	1322	2153	2456	1378	811	1743
2347	911	1878	3049	3357	1832	1028	2520
2149	964	2212	3923	4189	2213	1117	2377
3008	1346	3158	6023	6585	3145	1505	3162
2771	1266	3808	8021	9170	3865	1686	2988
3612	2067	5441	12527	13263	5342	2258	3696
6452	2411	6748	17733	20786	7056	2598	3711
4870	3228	10223	29439	32666	10271	3561	6056
4546	3825	12713	47328	53362	12902	4084	4784
6446	5028	17773	66493	148811	19378	5429	6676
5886	5787	23113	99470	255560	23768	6213	6244
8355	7272	31962	147585	321062	32190	7757	8502
7282	7473	34047	187623	594044	573691	7902	7647
9783	7957	39218	193571	496249	549110	8619	9938
6942	6921	30431	143561	341093	546084	7337	7299
7764	6280	25658	103585	209874	248696	8609	8103
5358	4802	17515	65005	122387	135696	5125	5503
5779	4111	14025	49121	48987	14323	4421	5783
3889	3116	9632	30130	31072	9905	3232	4095
4183	2658	8154	19704	21454	8021	2835	4469
5035	1872	4899	12240	13814	5210	2104	3253
3301	1674	2767	9149	9554	4351	1880	3542
2431	1283	2943	5789	6430	2771	1397	2646
2470	1138	2533	4445	4789	2448	1250	2568
1776	891	1725	2909	3107	1702	927	1901
1861	821	1525	2350	2474	1492	876	2060
1200	719	1195	1661	1873	1173	760	1451
1210	864	1387	1765	1854	1316	922	1351
1009	1161	1857	2117	2322	1755	1188	1094

### Rates of M2 R1-R2 gaps (Hz/cm<sup>2</sup>)

66493	120583	148811	77788
99470	217584	255560	107048
147585	321062	538980	508077
187623	594044	340550	170105
193571	496249	573691	205862
143561	341093	549110	217988
103585	209874	546084	344551
65005	122387	248696	114114
		135696	73421

Maximum rates extrapolated at **Run 5**,  
in red, reduction in the hypothesis of a **filter** upstream wrt M2

Maximum chamber rate (kHz/cm <sup>2</sup> )				
	M2	M3	M4	M5
R1	594.0 -> 344.5	274.5	203.5	232.7
R2	255.6 -> 79.2	64.2	34.1	39.0
R3	53.4 -> 19.2	8.9	6.2	8.9
R4	9.9	3.0	1.7	6.8

$\langle \frac{\sigma_R}{R} \rangle$  %

10.1%  
10.2%  
8.5%  
9.6%



# Background study at Run 5



**Minimum bias  
Low threshold:**

- Sample of 2600 events at  $L = 1.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

**97% are background:**  
defined as all except muons with  $E_\mu > 3 \text{ GeV}$

Fraction of signals of the present readout scheme by  
single-gap cases of background particles

	<b>M2</b> (%)	<b>M3</b> (%)	<b>M4</b> (%)	<b>M5</b> (%)
R1	$49.5 \pm 0.1$	$56.7 \pm 0.1$	$58.6 \pm 0.2$	$59.4 \pm 0.2$
R2	$49.7 \pm 0.1$	$70.4 \pm 0.2$	$75.0 \pm 0.2$	$76.0 \pm 0.2$
	<b><math>49.60 \pm 0.05</math></b>	<b><math>61.1 \pm 0.1</math></b>	<b><math>63.6 \pm 0.1</math></b>	<b><math>64.3 \pm 0.1</math></b>
R3	$57.5 \pm 0.1$	$73.0 \pm 0.2$	$76.8 \pm 0.3$	$77.1 \pm 0.3$
R4	$49.7 \pm 0.1$	$66.6 \pm 0.3$	$76.1 \pm 0.4$	$69.0 \pm 0.4$
	<b><math>53.6 \pm 0.1</math></b>	<b><math>70.8 \pm 0.2</math></b>	<b><math>76.6 \pm 0.2</math></b>	<b><math>74.0 \pm 0.3</math></b>



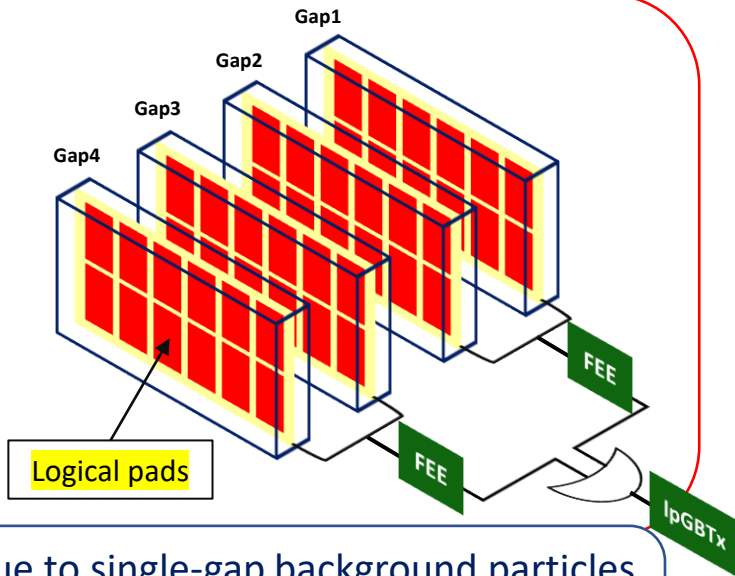
# Studies for a new readout scheme



## Current readout scheme

### 4-gaps OR-ed

If a charged particle crosses one out of four gaps, a signal is generated



Rates due to single-gap background particles

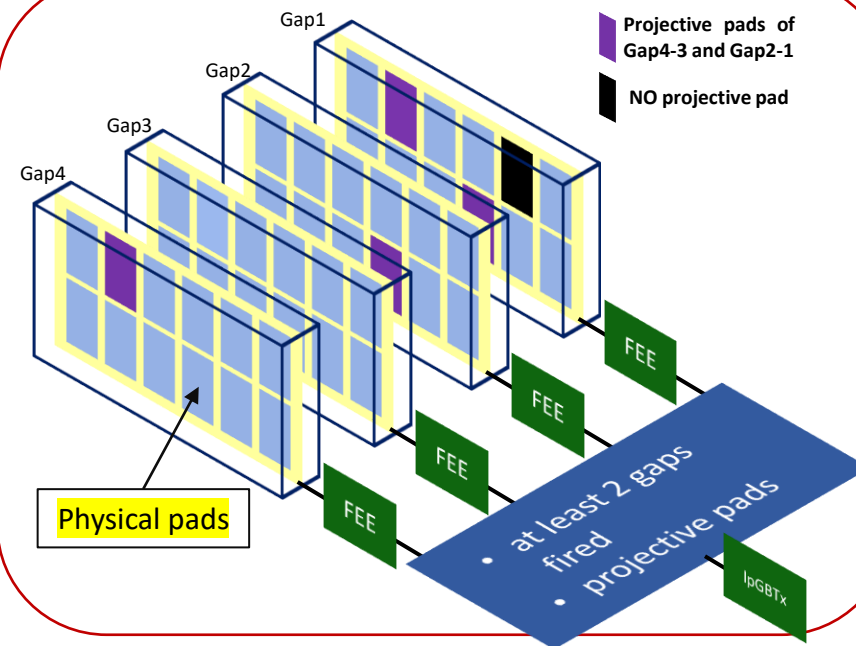
**up to 80%:**

**The new readout scheme is necessary**

Presented at:

[PoS LHCP2021 \(2021\) 192](#)

## New readout scheme



## New readout scheme for Upgrade II:

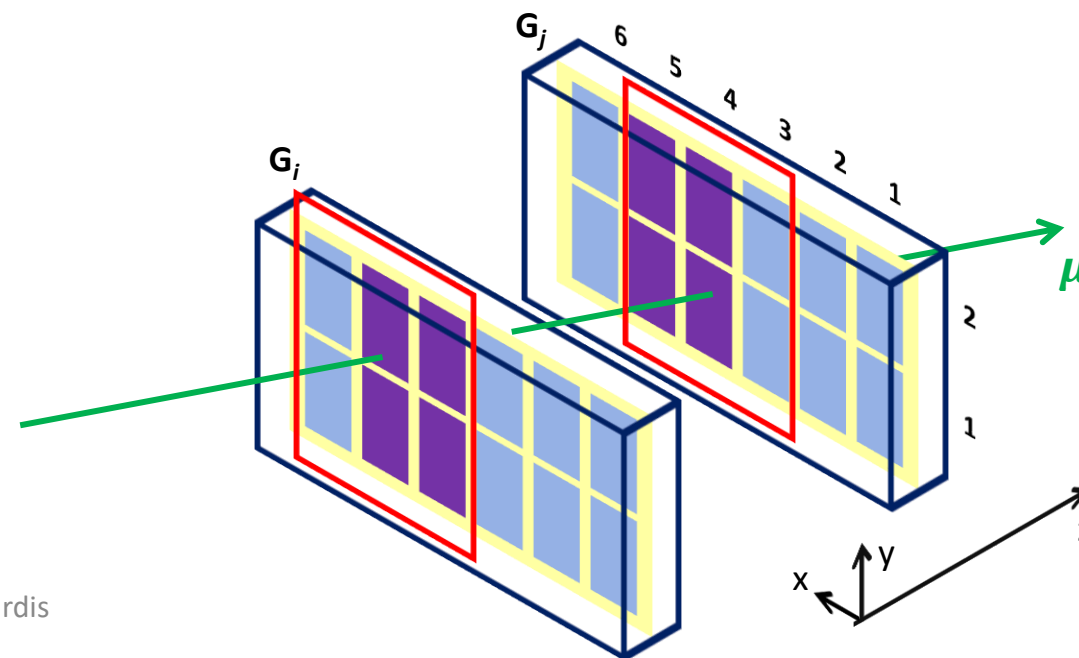
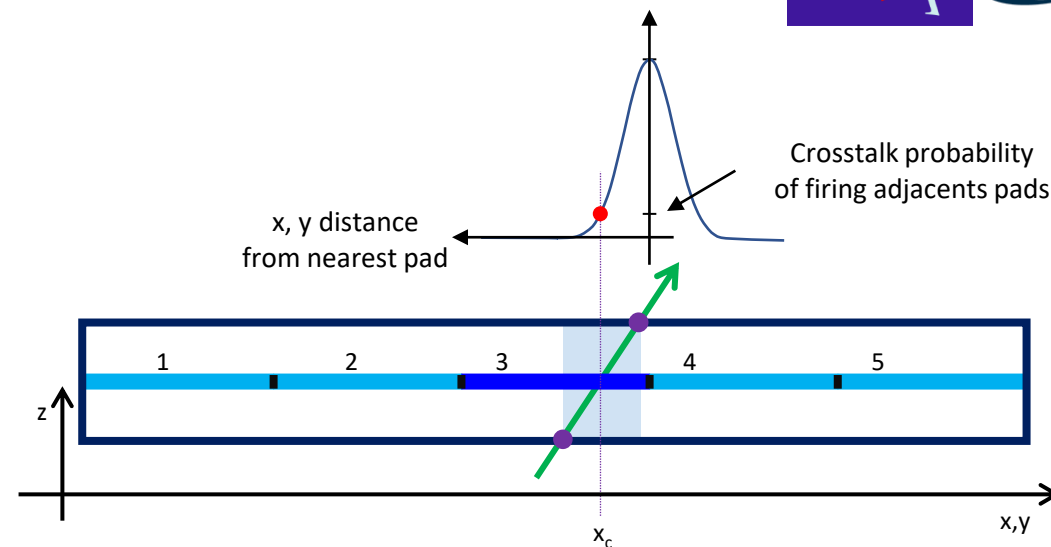
- readout for each gap
- signal in at least 2 gaps out of 4 at the same time
- at least two projective physical pads fired in the corresponding gaps.



# Muon Detector simulation



- Physical pads (anodes and cathodes) simulated in chambers at present geometrical configuration (Run 3)
- Electrical and geometrical crosstalk simulated for physical pad firing, at particles chambers crossing
- $\mu$ -RWELL in inner regions R1-R2 in the 4-gaps configuration
- Signal generation delay, into 25 ns time window, simulated for present MWPCs and  $\mu$ -RWELL
  - Efficiency of **85%** for MWPCs and **95%** for  $\mu$ -RWELLS
- Projectivity condition, between not only consecutive physical pads, simulated for signal generation (not in R4 of M2, M3 and M4 where chambers are fixed at the present OR-ed Bi-gap scheme)





# Background reduction with the new readout scheme



$\mu$ -RWELL time-eff= 95%  
MWPCs time-eff= 85%

The new readout scheme has a very significant effect on the background signals removal in the chambers.

A great merit is given by the AND request.

What is the effect on signal muons?

R1  
R2

R3  
R4

	<b>M2</b> (%)	<b>M3</b> (%)	<b>M4</b> (%)	<b>M5</b> (%)
R1	$66.40 \pm 0.07$	$77.0 \pm 0.1$	$78.20 \pm 0.13$	$78.6 \pm 0.1$
R2	$60.70 \pm 0.07$	$82.10 \pm 0.13$	$86.5 \pm 0.2$	$86.8 \pm 0.2$
	<b><math>63.50 \pm 0.05</math></b>	<b><math>78.60 \pm 0.08</math></b>	<b><math>80.7 \pm 0.1</math></b>	<b><math>81.0 \pm 0.1</math></b>
R3	$64.40 \pm 0.09$	$79.0 \pm 0.2$	$80.8 \pm 0.3$	$81.0 \pm 0.3$
R4	0.0	0.0	0.0	$72.30 \pm 0.43$
	<b><math>32.10 \pm 0.05</math></b>	<b><math>50.9 \pm 0.1</math></b>	<b><math>56.0 \pm 0.1</math></b>	<b><math>77.70 \pm 0.11</math></b>



# Efficiency loss with the new readout scheme



$J/\psi \rightarrow \mu^+ \mu^-$ :

- Sample of 200 000 of  $J/\psi \rightarrow \mu^+ \mu^-$  candidate.

**Inefficiency less than 2%, against the large background reduction up to 80%**

The effects of:  
-dead time and  
-muons mis-identification  
are not included

$\mu$ -RWELL time-eff= 95%  
MWPCs time-eff= 85%

CERN/LHCC 2021-012 LHCb TDR 23

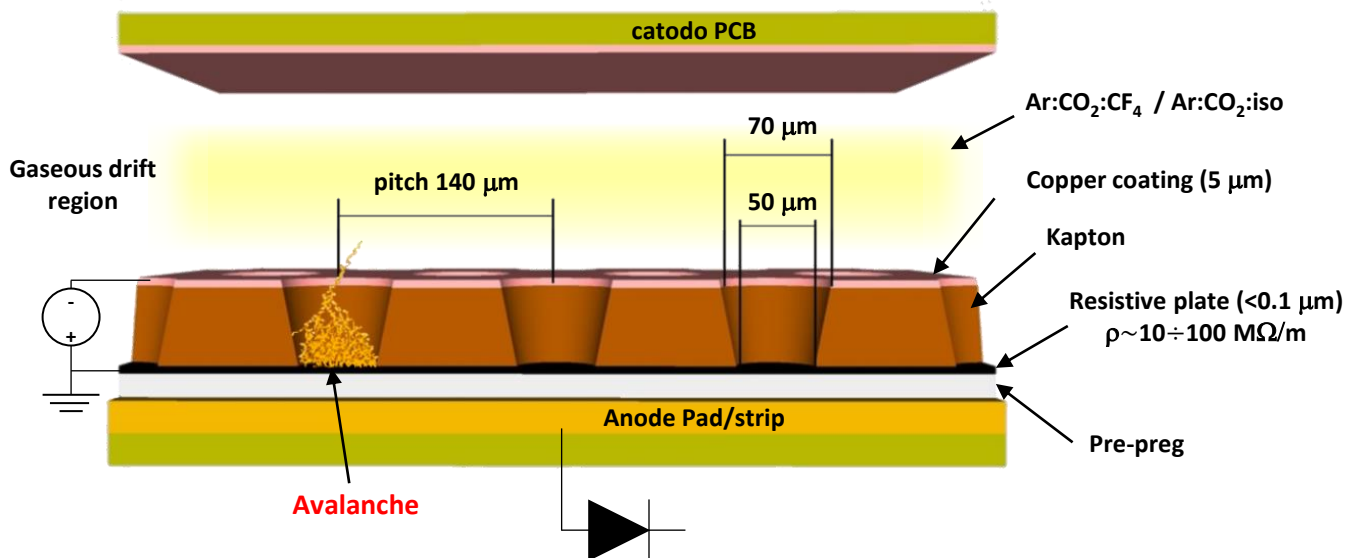
Stations to be crossed for current $\mu$ -identification			
$p$ [GeV/c]	overall detector (%)	inner regions R1-R2 (%)	outer regions R3-R4 (%)
$3 < p < 6$ (M2 & M3)	$1.13 \pm 0.08$	$1.73 \pm 0.35$	$1.09 \pm 0.08$
$6 < p < 10$ (M2 & M3 & (M4    M5))	$1.06 \pm 0.06$	$0.70 \pm 0.16$	$1.10 \pm 0.07$
$p > 10$ (M2 & M3 & M4 & M5)	$2.23 \pm 0.04$	$0.52 \pm 0.03$	$4.04 \pm 0.08$

# $\mu$ RWELL for Upgrade II

JINST (2019) 289 14 P05014

- Cathode plan
- Drift region
- $\mu$ RWELL system:
  - Wells to copper clad foil
  - Resistive plate
  - Readout PCB of pad/stip

The ionization avalanche in the drift region is amplified in the wells and limited by the resistive plate.



Present pads size [mm]

$\mu$ -RWELL pads size [mm]

## Performances:

- High rate capability (up to 1 MHz/cm<sup>2</sup>)
- High granularity, fundamental for the dead time inefficiency reduction

Reg / Sta	M2	M3	M4	M5
R1	38x31 -> 9x9	41x34 -> 10x10	29x36 -> 11x10	31x39 -> 12x11
R2	76x31 -> 9x18	82x34 -> 10x19	58x73 -> 11x21	62x77 -> 12x22
R3	25x125	27x135	58x145	62x155
R4	50x250	54x270	58x290	62x309





# Inefficiency induced by dead time



From Run 5 rates, the inefficiency by the electronics dead time is calculated

- The FEE dead time is assumed conservatively of 100 ns

JINST (2016) 11 P04010

A software has been developed to evaluate the dead time effects on decay modes of special interest

Maximum deadtime inefficiency % HCAL - MWPC				
	M2	M3	M4	M5
R1	17.14	6.65	7.50	8.66
R2	17.81	4.62	5.69	7.34
R3	7.21	1.72	3.49	5.68
R4	8.24	3.37	2.30	8.55

Maximum deadtime inefficiency % HCAL - $\mu$ RWELL				
	M2	M3	M4	M5
R1	1.18	0.48	0.79	0.95
R2	1.22	0.32	0.31	0.41
R3	7.21	1.72	3.49	5.68
R4	8.24	3.37	2.30	8.55

Moving to high-granularity  $\mu$ RWell chambers

Presented at:

[IFAE 2023 – Catania](#)

[6th Workshop on LHCb upgrade II - Barcellona](#)

SCENARIO		$B_s^0 \rightarrow \mu^+ \mu^-$ (%)	$D^0 \rightarrow \mu^+ \mu^-$ (%)	$K_s^0 \rightarrow \mu^+ \mu^-$ (%)	$B_s^0 \rightarrow J/\psi(\mu^+ \mu^-)\phi$ (%)
HCAL	MWPC	30.50 ± 0.06	32.10 ± 0.25	25.80 ± 0.35	31.0 ± 0.1
	$\mu$ -RWELL	10.30 ± 0.04	9.50 ± 0.16	8.50 ± 0.22	9.60 ± 0.05
SHIELD	MWPC	19.00 ± 0.05	19.40 ± 0.21	14.0 ± 0.3	18.80 ± 0.07
	$\mu$ -RWELL	8.60 ± 0.04	7.60 ± 0.14	6.4 ± 0.2	7.80 ± 0.05
w/o M2	MWPC	13.40 ± 0.04	17.2 ± 0.2	10.70 ± 0.24	16.4 ± 0.1
	$\mu$ -RWELL	6.00 ± 0.03	5.30 ± 0.12	3.20 ± 0.14	5.30 ± 0.04

The **muon identification** acts at different levels:

2. For each IsMuon track candidate, the  $\chi_{corr}^2$  is calculated, from distances between the extrapolated points and hits on stations

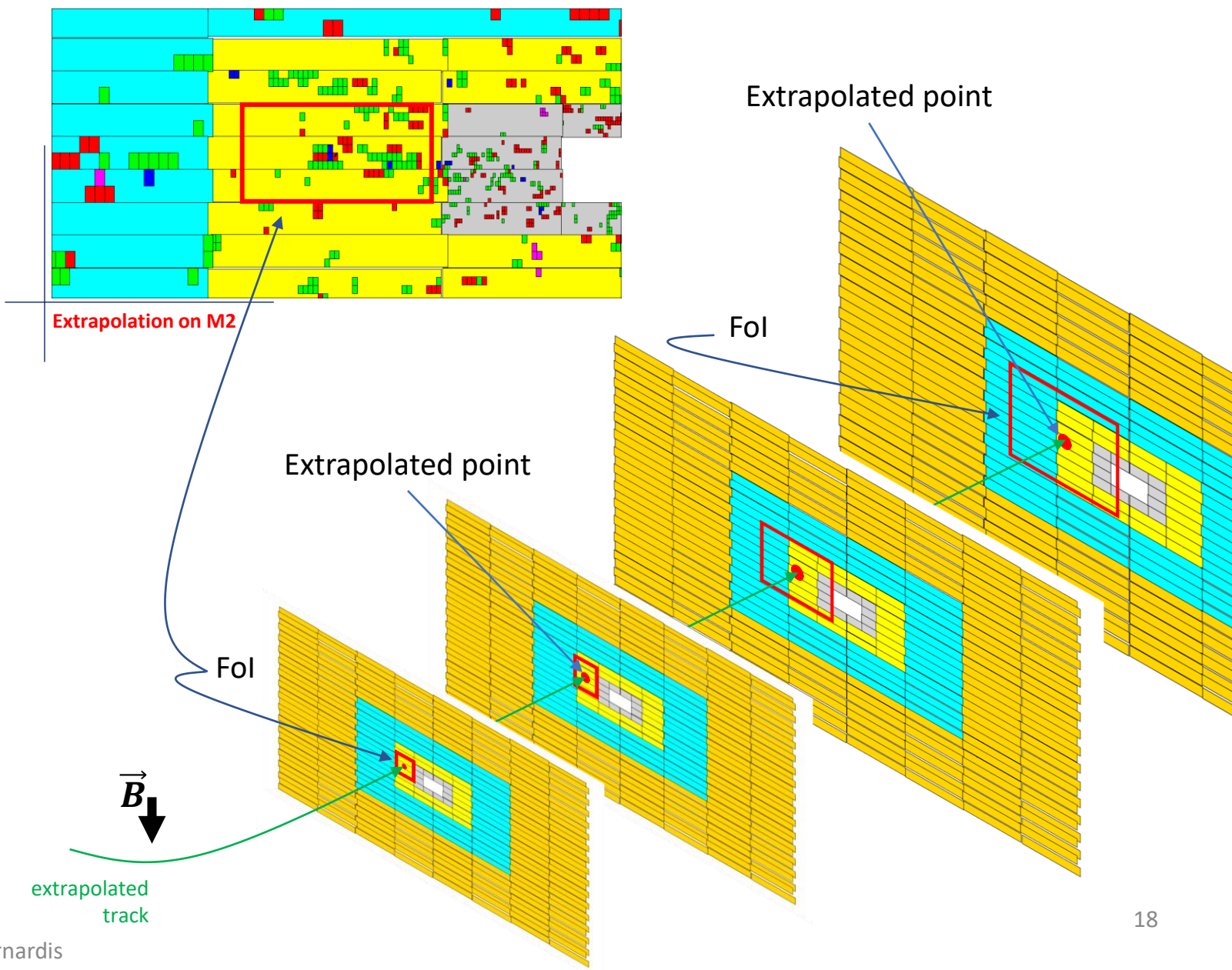
JINST 15 (2020) T12005

3. Cut on  $\chi_{corr}^2$  is applied

**False muon hits** can be identified on stations at very crowded environment.

Development of a software to simulate the full hits generation at Run 5, in 25 ns time window :

- New *FoI* definitions
- Different  $\chi_{corr}^2$  cut





# Outline



- Studies on the design of the LHCb muon detector at high luminosity
  - Rate estimate
  - Proposal of a new readout scheme
  - Dead time and muon ID

- Search for the rare decay mode  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-$ 
  - Optimisation of the  $D_s$  selection
  - $D_s^+ \mu^+ \mu^-$  selection
  - Observation of new decay modes



# The puzzle of the $D_{s1}(2460)$ meson



## Theoretical predictions

- Broad  $c\bar{s}$  meson state
- Massive enough that dominant strong decays would have been isospin-conserving in  $D^*K$  final states.

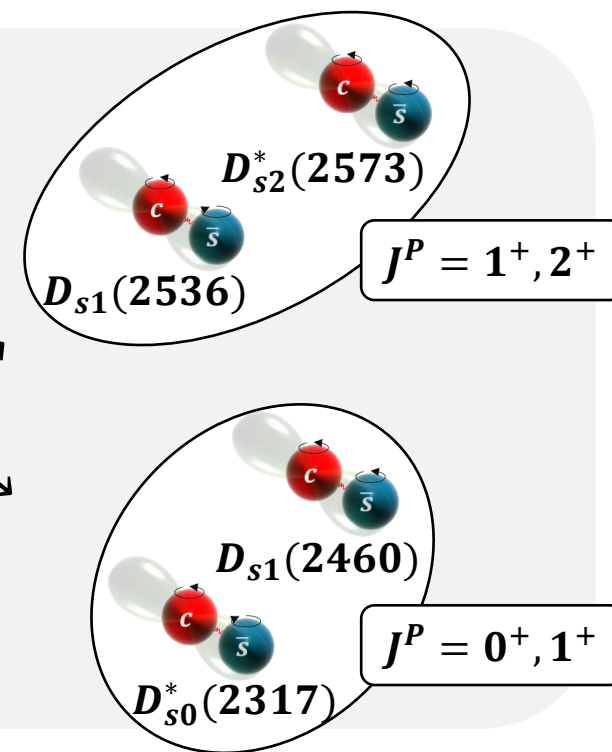
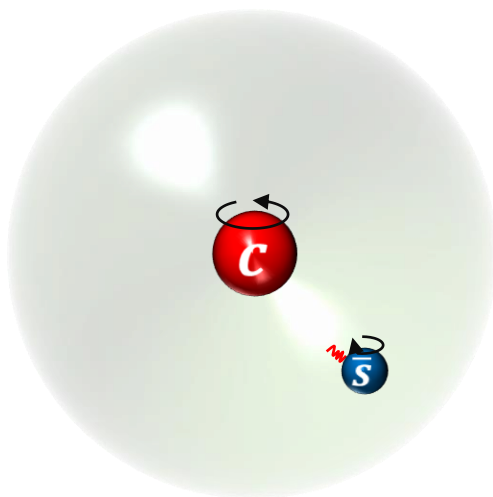
Phys. Rev. Lett. 66 (1991) 1130  
 Phys. Rev. D 12 (1975) 147

Excited  $P$ -states ( $\vec{L} = 1$ )

$$L + S_{\bar{s}} \begin{cases} \rightarrow j_{\bar{s}} = 3/2 \\ \rightarrow j_{\bar{s}} = 1/2 \end{cases} + S_c = 1/2$$

### In the heavy quark limit model

Bounding force depends on the spin-orbit coupling only of the light  $\bar{s}$  quark





# The puzzle of the $D_{s1}(2460)$ meson



## Surprising observation:

$$D_{s1}(2460)^+ \rightarrow D_s^{*+} \pi^0$$

**CLEO (2003)**

- mass smaller ( $\approx 150$  MeV) than all theoretical predictions
- observed in the isospin-violating  $D_s^* \pi^0$  channel
- narrow state

$D_{s1}(2460)^+$

decay modes	
$D_s^{*+} \pi^0$	$(48 \pm 11)\%$
$D_s^+ \gamma$	$(18 \pm 4)\%$
$D_s^+ \pi^+ \pi^-$	$(4.3 \pm 1.3)\%$
<b>TOT</b>	<b><math>(70 \pm 16)\%</math></b>

In a  $c\bar{s}$  scenario:

isospin-violating decay is suppressed.

For the  $D_{s1}(2460)$ :

isospin-violating decay  $D_s^* \pi^0$  is much frequent, suggesting a 4-quarks system ( $c\bar{s}u\bar{u}$  or  $c\bar{s}d\bar{d}$ )

## Dalitz decays

- $\pi^0 \rightarrow \gamma \gamma^* \rightarrow \gamma e^+ e^-$  decay process predicted by R.H. Dalitz. The virtual photon directly produces an electron-positron pair.

(nowadays)

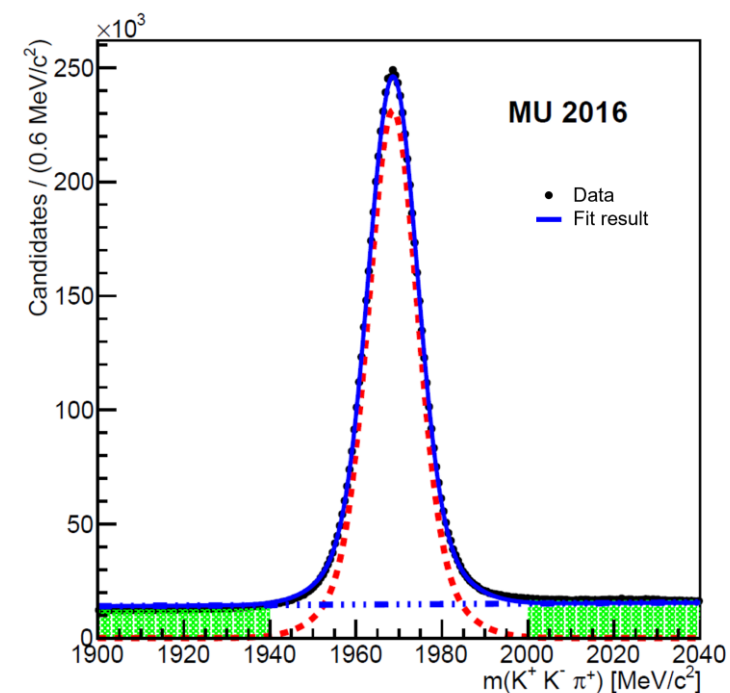
- Any decay process with two leptons in the final state, produced by a virtual photon.



# Optimisation of the $D_s^+$ selection



BEFORE



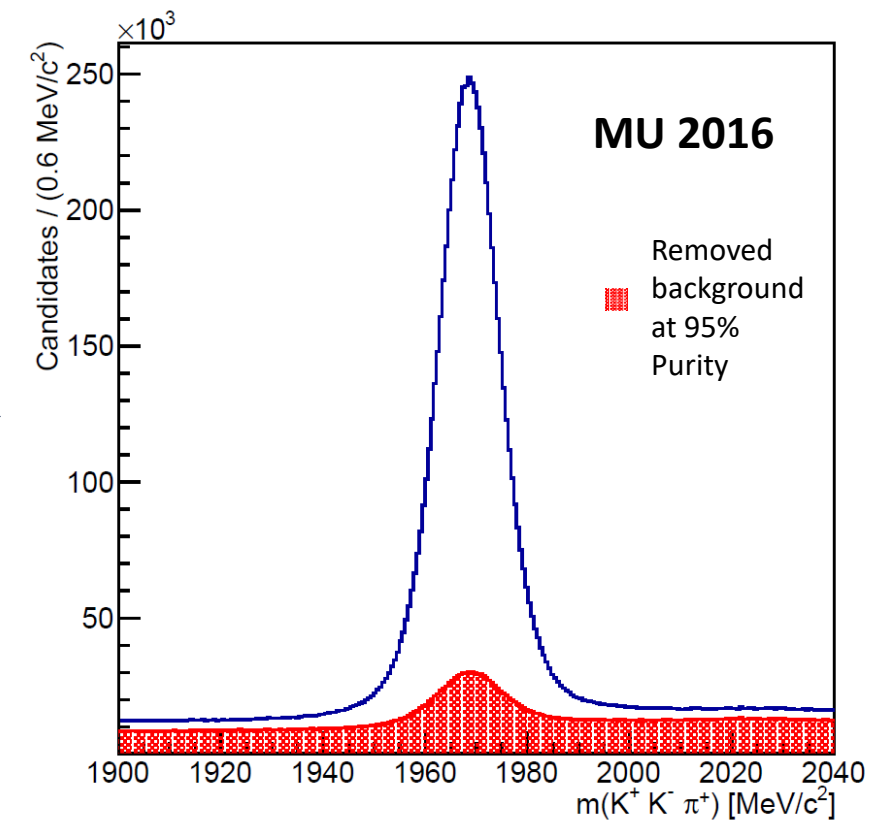
sFit to the  $K^+K^-\pi^+$  mass spectrum

BDT multivariate analysis to classify:

- signal** candidates (s-weighted  $D_s^+$  distribution)
- and
- background** candidates (sidebands).



AFTER



# $D_s^+ \mu^+ \mu^-$ selection

RUN 2  
(2015-2018)  
 $6 \text{ fb}^{-1}$



$$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^- + c. c.$$

$$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^+ + c. c.$$

$$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^- \mu^- + c. c.$$

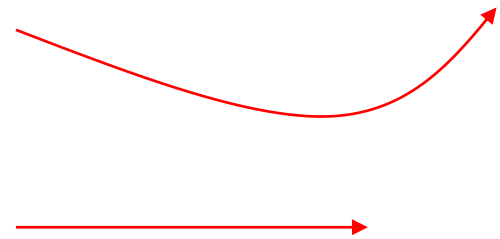
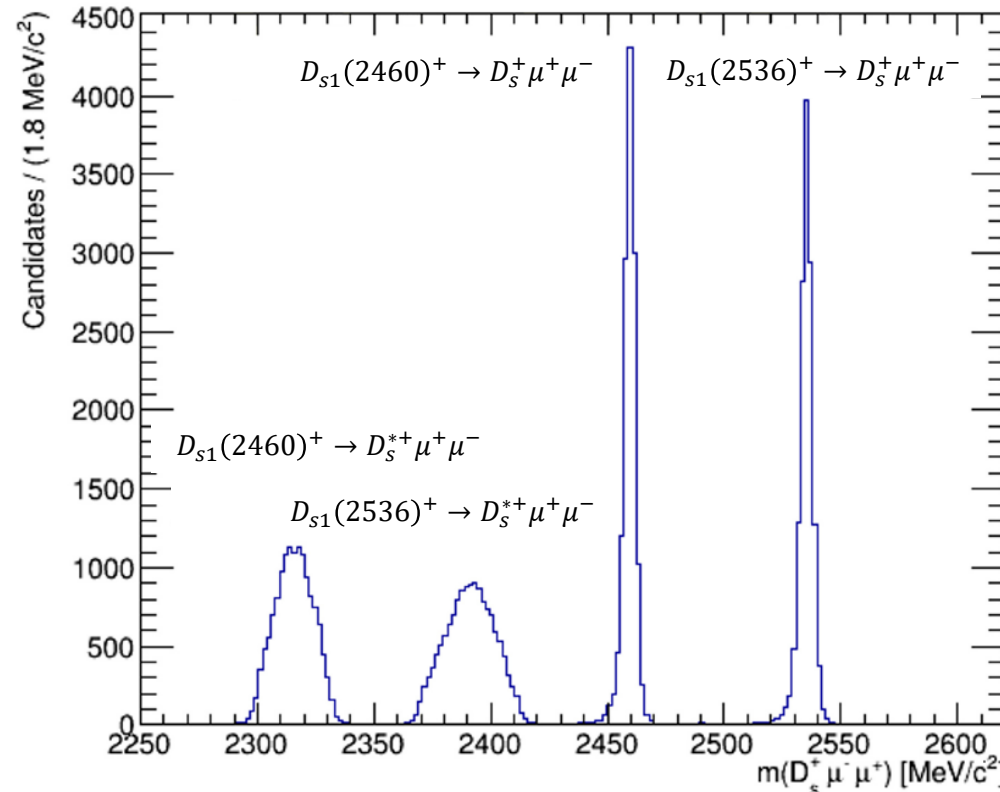
- $D_s$  candidates:  
multivariate selection application

- $D_{s1}(2460)$  candidates:

-BDT and BDTG algorithms for  
signal and background events  
classifications.

- MC samples as signal input  
( $\frac{1}{2}$  for training and  $\frac{1}{2}$  for testing)

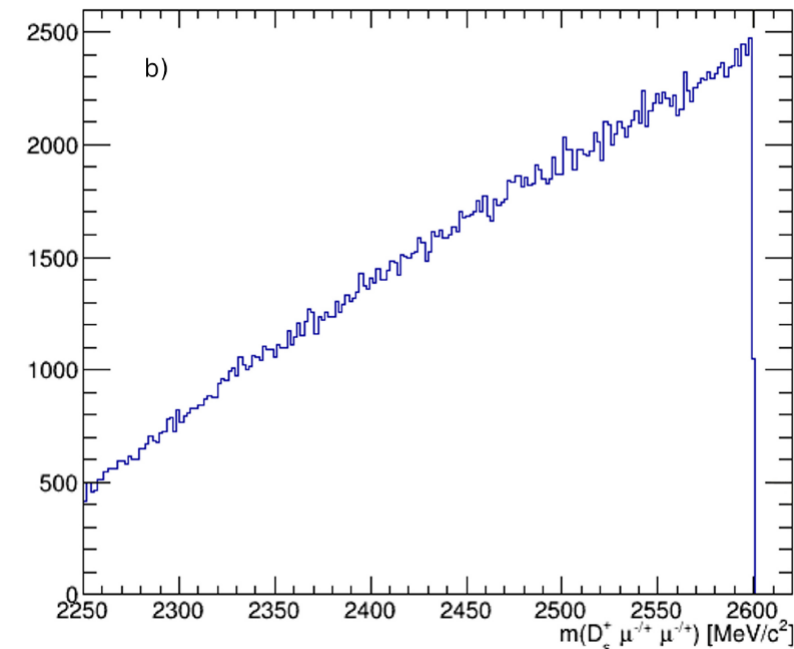
- Wrong Signs sample  
as a proxy of **combinatorial  
background**  
( $\frac{1}{2}$  for training and  $\frac{1}{2}$  for testing)



Presented at:

[EPJ Web Conf. 270 \(2022\) 00013](#)

[LHCb Italia 2022 Collaboration Meeting](#)



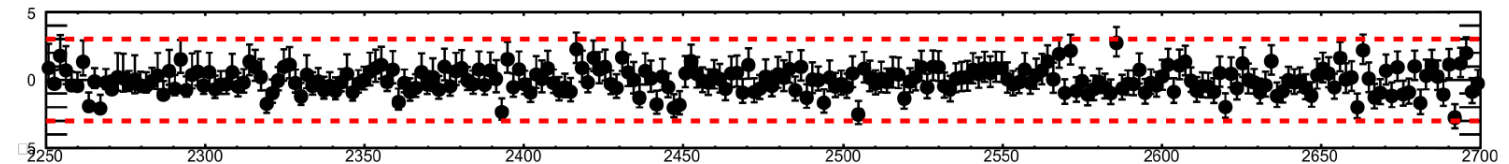
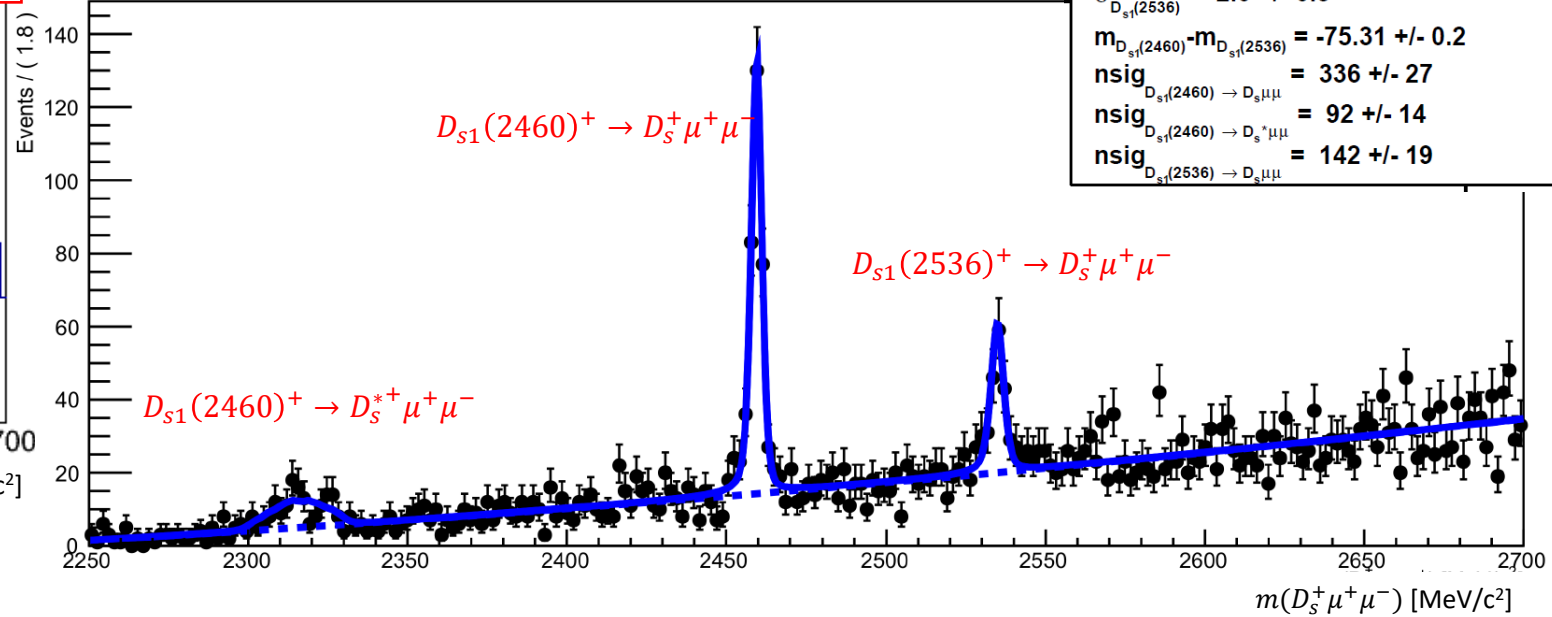
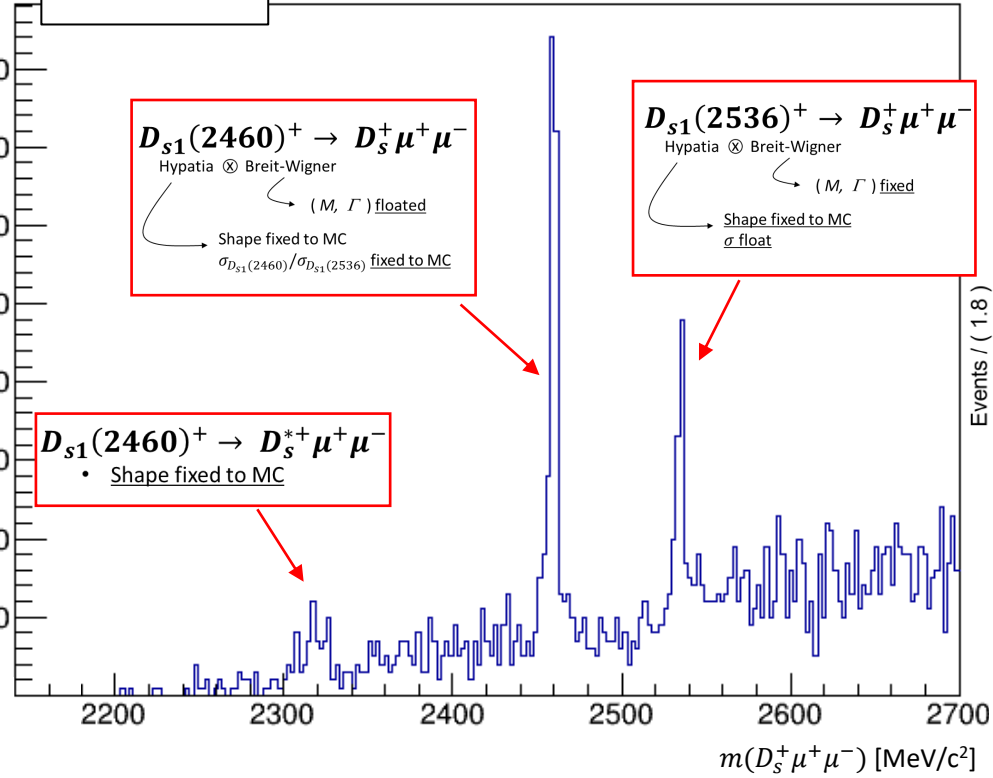
# Fit result:

RUN 2  
(2015-2018)  
 $6 \text{ fb}^{-1}$



## Fit model

$(PDF_{total})$






# New observations:

- $m_{D_{s1}(2460)} = 2459.9 \pm 0.2 \text{ MeV}/c^2$
- $\Gamma_{D_{s1}(2460)} = 0.5 \pm 0.7 \text{ MeV}/c^2$

upper limit evaluated at 95% CL:


$$\Gamma_{D_{s1}(2460)} < 1.77 \text{ MeV}/c^2$$



$D_{s1}(2460)^\pm$	<b>MASS</b>	$2459.5 \pm 0.6 \text{ MeV}/c^2$
$D_{s1}(2460)^\pm$	<b>WIDTH</b>	$< 3.5 \text{ MeV}/c^2$

**PDG**

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PDG

$$\frac{\mathfrak{B}(D_{s1}(2460) \rightarrow D_s^* \mu^+ \mu^-)}{\mathfrak{B}(D_{s1}(2460) \rightarrow D_s \mu^+ \mu^-)} = \text{BR}_{2460} \cdot \frac{\epsilon_1}{\epsilon_2} = 0.48 \pm 0.10$$

$D_{s1}(2460) \rightarrow D_s^* \mu^+ \mu^-$

peak significance :  $n_\sigma = 7.7\sigma$

**NEW DECAY OBSERVATION**

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NEW DECAY OBSERVATION

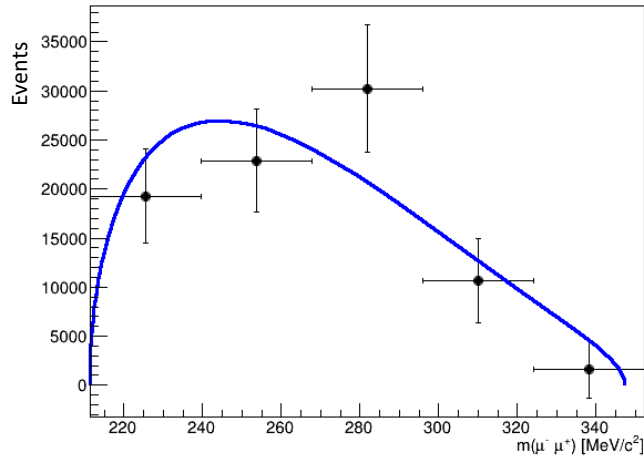
First observation of the decay:  $D_{s1}(2536) \rightarrow D_s^+ \mu^+ \mu^-$

Unexpected observation,  
given that the corresponding radiative decay is not seen

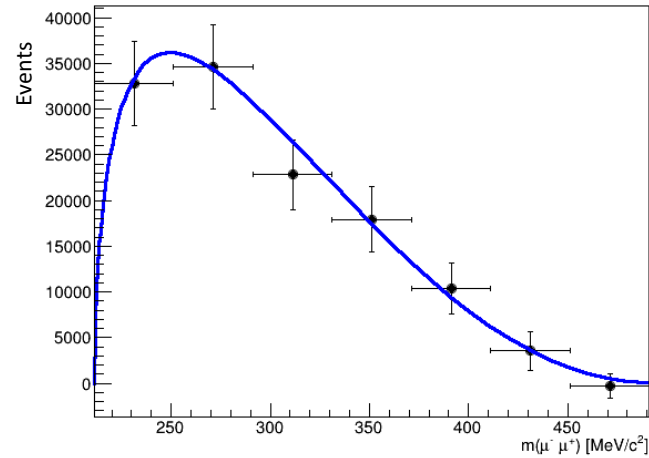
Is it a  
Dalitz  
decay?

# Study of the $\mu^+ \mu^-$ spectrum:

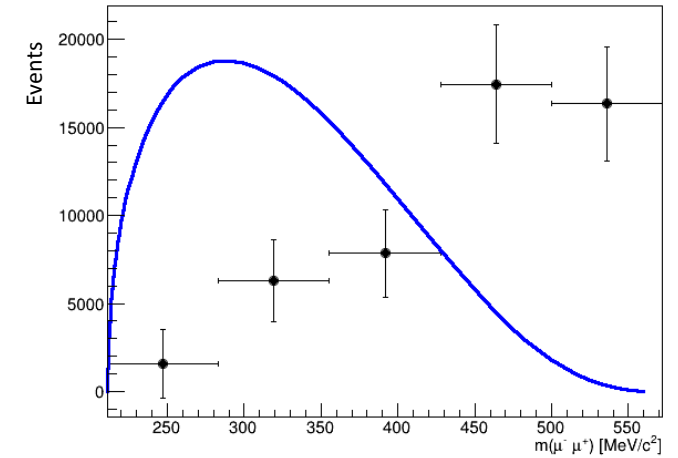
$D_{s1}(2460)^+ \rightarrow D_s^{*+} \mu^+ \mu^-$



$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-$



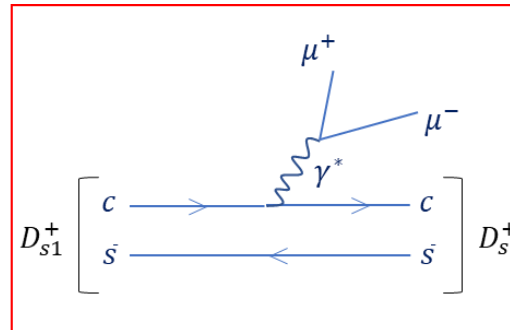
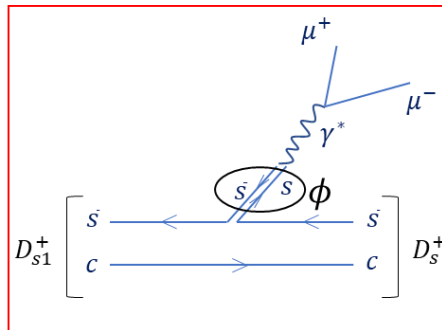
$D_{s1}(2536)^+ \rightarrow D_s^+ \mu^+ \mu^-$



## Fits using models built for the Dalitz processes in the QCD framework

[Phys. Rev. D 108, 074027 - 2023]

A virtual photon couples to a  $c$  and  $\bar{s}$  quark. The latter through an intermediate  $\phi$  meson, according to the Vector Meson Dominance model.



**Dalitz nature** can be supported for decays:

- $D_{s1}(2460)^+ \rightarrow D_s^{*+} \mu^+ \mu^-$
- $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-$

Dalitz nature is disfavoured for decay:

- $D_{s1}(2536)^+ \rightarrow D_s^+ \mu^+ \mu^-$



# Summary



- Studies for the muon detector at high luminosity
  - New readout scheme and  $\mu$ -RWELL in the inner regions
  - Large background hits reduction up to 80%
  - Inefficiency induced by dead time under study. New scenarios proposed
  - Preliminary results on  $\mu$  mis-ID studies pave the way to deeper investigations and new ideas
- Search for the rare decay mode  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-$ 
  - Precise measurements of  $m_{D_{s1}(2460)}$  and  $\Gamma_{D_{s1}(2460)}$
  - New observation of the  $D_{s1}(2460)^+ \rightarrow D_s^{*+} \mu^+ \mu^-$  decay
  - Precise measurement of  $\mathcal{B}(D_{s1}(2460) \rightarrow D_s^* \mu^+ \mu^-) / \mathcal{B}(D_{s1}(2460) \rightarrow D_s \mu^+ \mu^-)$
  - New observation of the  $D_{s1}(2536)^+ \rightarrow D_s^+ \mu^+ \mu^-$  decay



Thank you  
for your attention

**BACKUP**



# Background study at Run 5



## Minimum bias Low threshold:

- Sample of 2600 events at  $L = 1.5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

$\approx 800$  charged particles cross the detector at each bunch crossing:

$e^\pm$	69%
$\pi^\pm$	12%
$p^\pm$	12%
$\mu^\pm$	6% (all)
	3% (bckg)
$K^\pm$	0.6%
$nucl$	0.4%

## Background particles distribution in regions:

R1:	37%	} 66%
R2:	29%	
R3:	18%	} 34%
R4:	16%	

### Electrons:

R1:	36 %
R2:	28 %
R3:	18 %
R4:	18 %

### Pions:

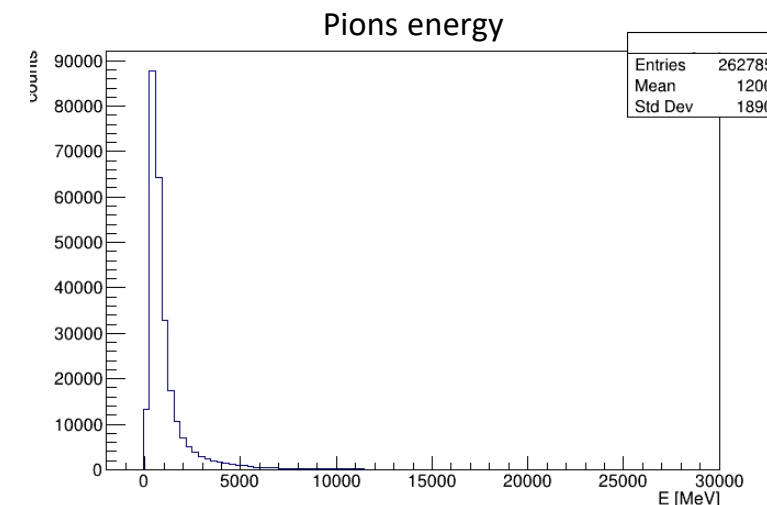
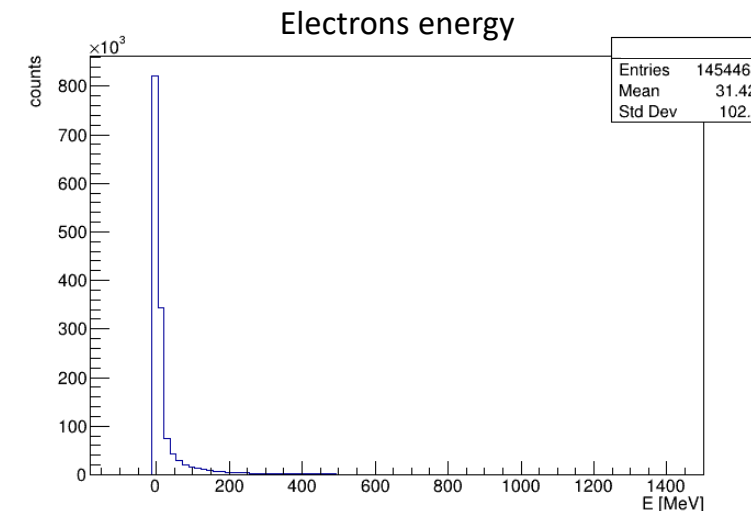
R1:	48 %
R2:	32 %
R3:	13 %
R4:	7 %

### Protons:

R1:	41 %
R2:	32 %
R3:	17 %
R4:	10 %

### Muons ( $E_\mu < 3 \text{ GeV}$ ):

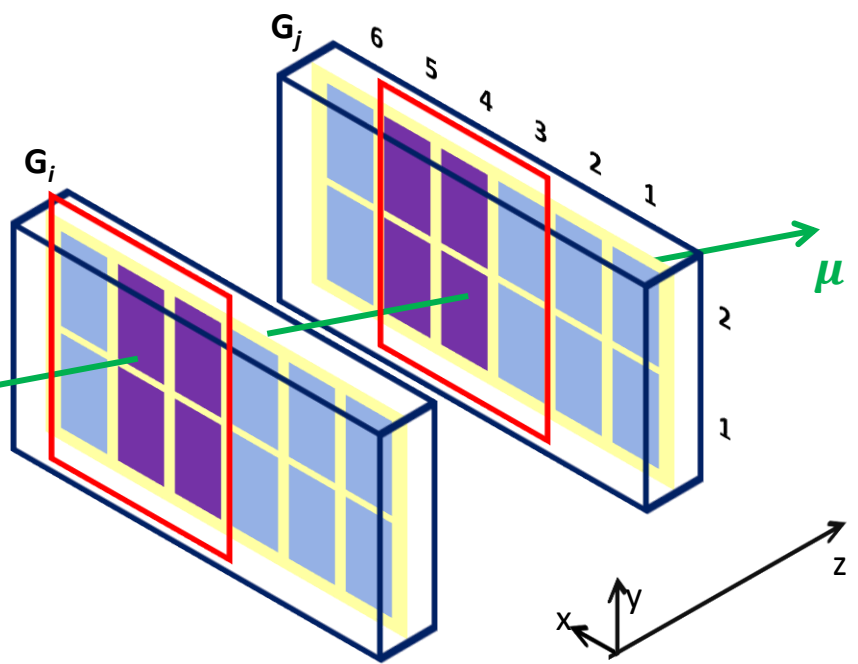
R1:	18 %
R2:	23 %
R3:	21 %
R4:	38 %







# Projectivity and time inefficiency improvements

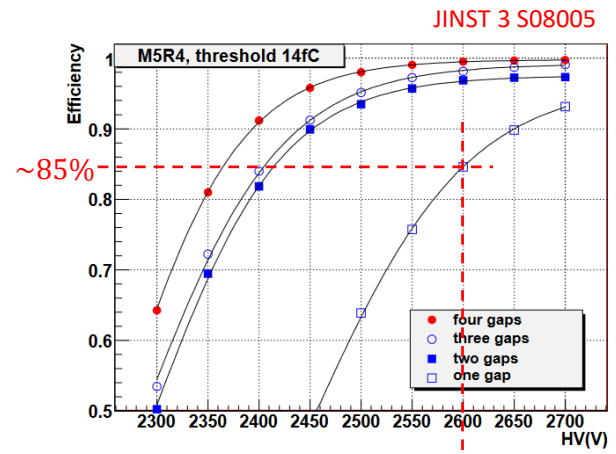


A chamber is fired if at least one couple of **not only consecutive** gaps, have at least two fired projective pads.

Pads that could be fired ( $P=1$ ) or not fired ( $P=0$ ), in gaps  $G_i$  and  $G_j$ . Of them, violet pads in a gap are projective with respect to those in the other gap.

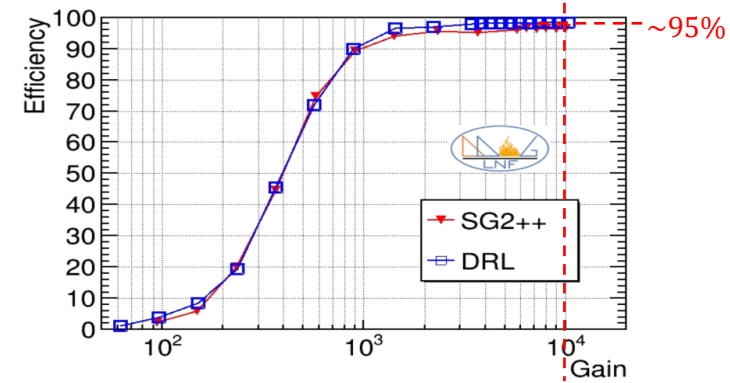
## MWPCs

As preliminary evaluation, **Single gap Time Efficiency (SgTE)** equal to: **85%**



## μRWELL

As preliminary evaluation, **Single gap Time Efficiency (SgTE)** equal to: **95%**

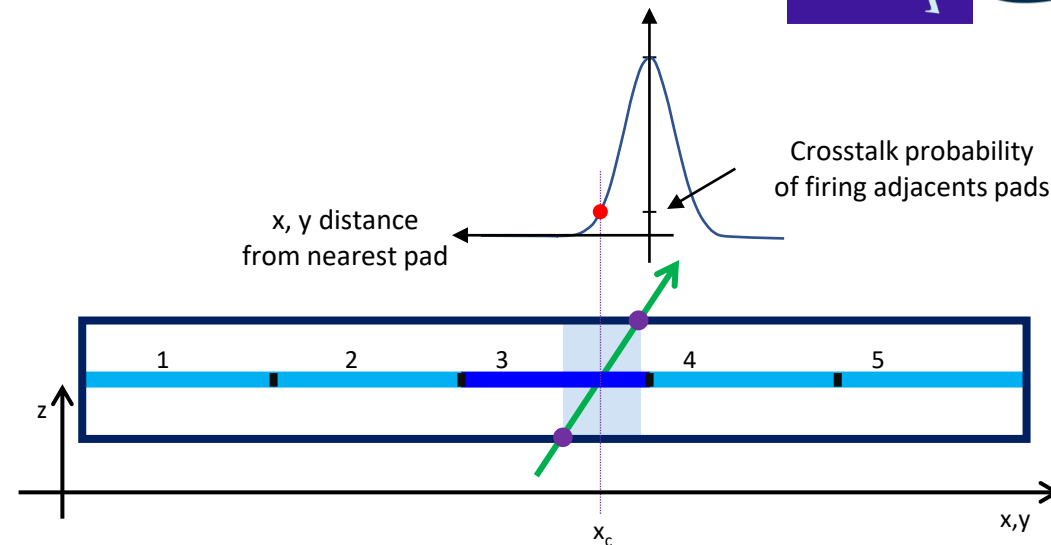




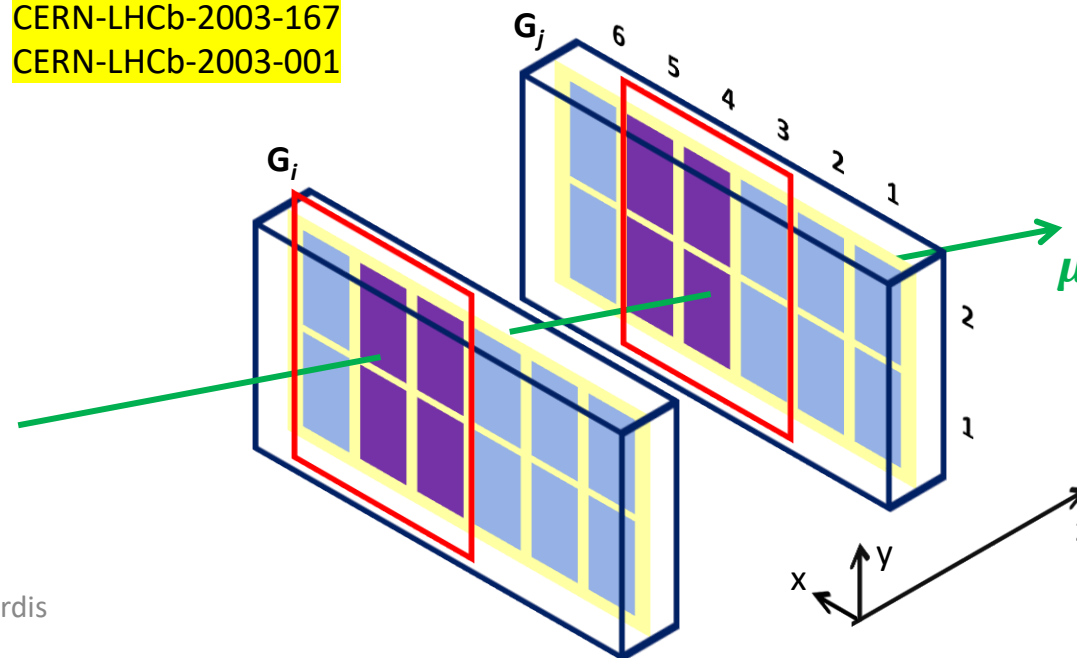
# Muon Detector simulation



- Physical pads (anodes and cathodes) simulated in chambers at present geometrical configuration (Run 3)
- Electrical and geometrical crosstalk simulated for physical pad firing, at particles chambers crossing
- $\mu$ -RWELL in inner regions R1-R2 in the 4-gaps configuration
- Signal generation delay, into 25 ns time window, simulated for present MWPCs and  $\mu$ -RWELL  
 Efficiency of **85%** for MWPCs and **95%** for  $\mu$ -RWELLS
- Projectivity condition, between not only consecutive physical pads, simulated for signal generation (not in R4 of M2, M3 and M4 where chambers are fixed at the present **OR-ed Bi-gap scheme**)



LHCB-MUON 2001-120 V2.0  
 CERN-LHCb-2001-150  
 CERN-LHCb-2003-167  
 CERN-LHCb-2003-001

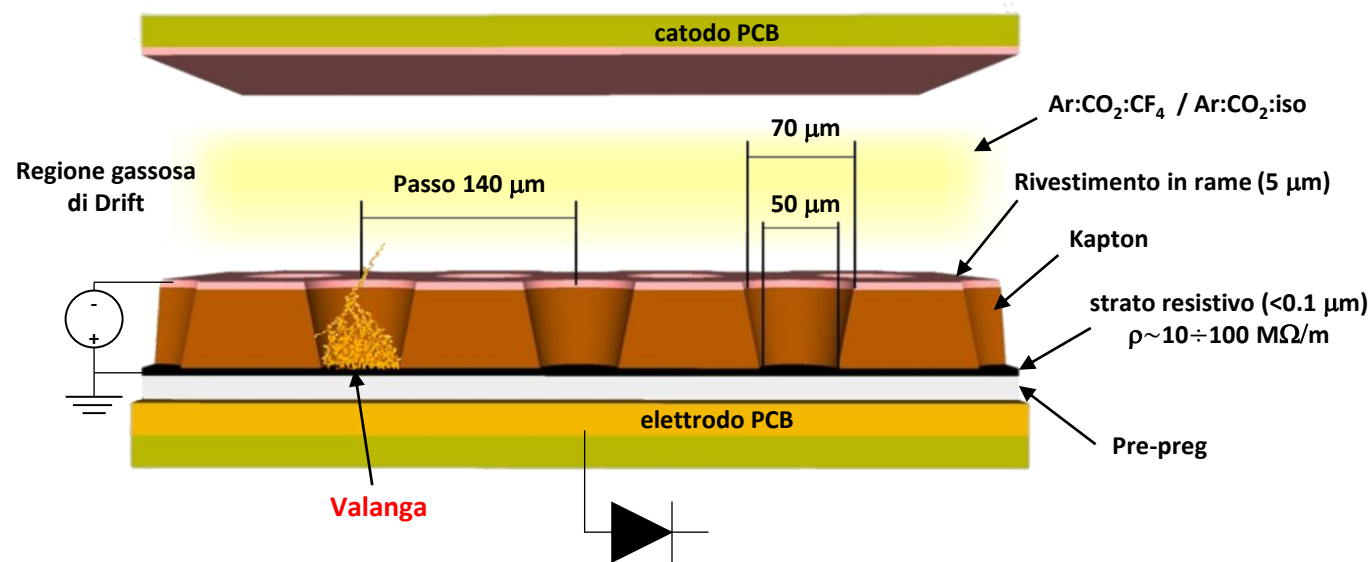


# Le $\mu$ RWELL

JINST (2019) 289 14 P05014

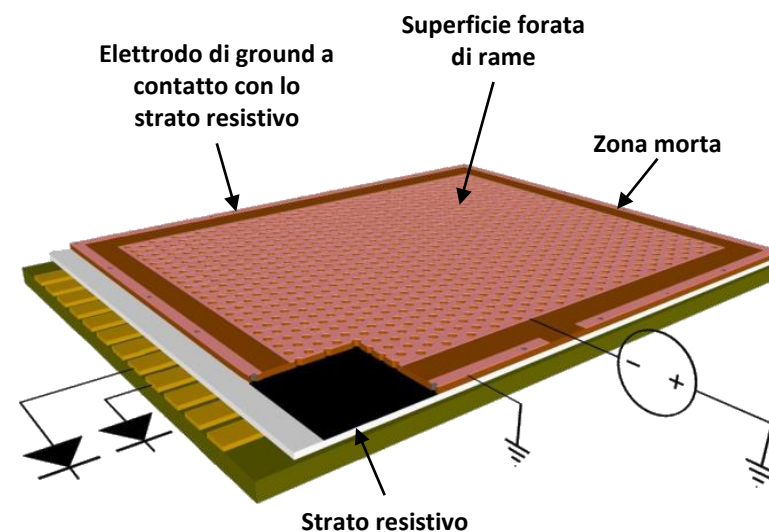
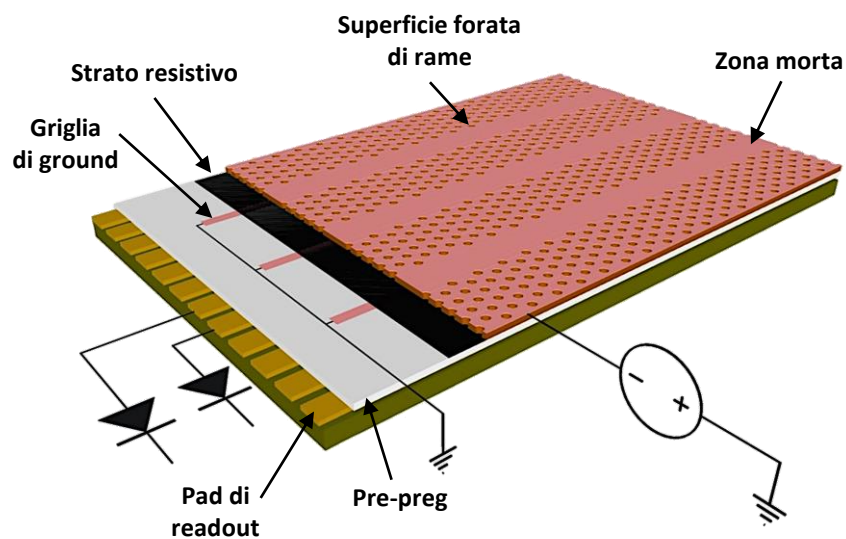
Un rivelatore  $\mu$ RWELL è un MicroPattern Gas Detector (MPGD) che consiste principalmente dei seguenti elementi:

- Piano catodico
- Regione di deriva
- Scheda  $\mu$ RWELL, compattata con:
  - foglio di kapton forato e rivestito in rame
  - Strato resistivo
  - PCB di readout con segmentazione a pad/stip



La ionizzazione a valanga, generata nella regione di drift, è amplificata nei fori ed è controllata grazie allo strato resistivo.

La perdita di efficienza ad alti rates (Run 5) è risolta con una fitta rete di elettrodi di messa a terra che evacuano rapidamente le valanghe.





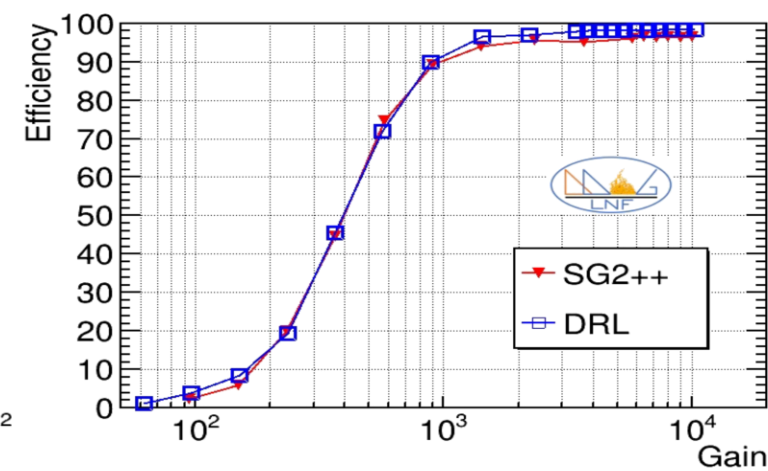
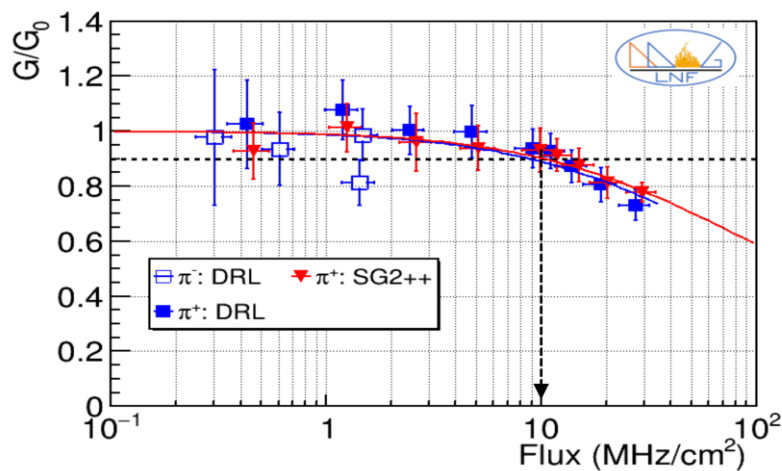
# Le $\mu$ RWELL



JINST (2019) 289 14 P05014

## Prestazioni:

- elevata rate capability (**fino a 1 MHz**)
- elevata efficienza in risposta temporale  
 $\epsilon = 96\%$  con miscela Ar:CO<sub>2</sub>:CF<sub>4</sub> = 45:15:40  
 in fase di studio con miscela Ar:CO<sub>2</sub>:iso = 68:30:02
- Dimensioni fino a 1.2x0.5 m
- elevata granularità**  
 che ha un effetto molto significativo  
 sull'inefficienza da tempo morto (dopo ->)



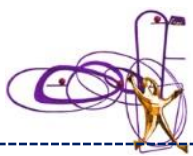
## Confronto tra le dimensioni (mm<sup>2</sup>):

- Pads fisiche attuali delle MWPCs
- Pad delle  $\mu$ RWELL (in rosso)



Reg / Sta	M2	M3	M4	M5
R1	38x31 -> 9x9	41x34 -> 10x10	29x36 -> 11x10	31x39 -> 12x11
R2	76x31 -> 9x18	82x34 -> 10x19	58x73 -> 11x21	62x77 -> 12x22
R3	25x125	27x135	58x145	62x155
R4	50x250	54x270	58x290	62x309

LHCb-TDR-023



# The muon mis-IDentification



The **muon identification** acts at different levels:

- For each IsMuon track candidate, the  $\chi^2_{corr}$  is calculated, from distances between the extrapolated points and hits on stations

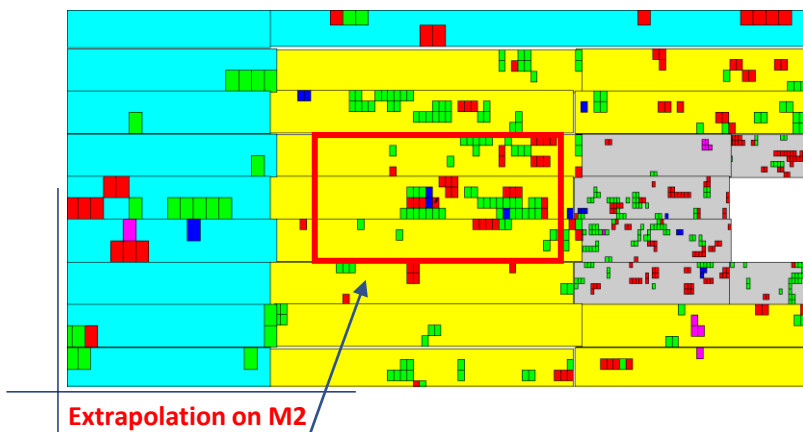
JINST 15 (2020) T12005

$$\chi^2_{corr} \approx \frac{1}{N} \sum_i \left\{ \left( \frac{x_i - x_{track}}{pad_x} \right)^2 + \left( \frac{y_i - y_{track}}{pad_y} \right)^2 \right\}$$

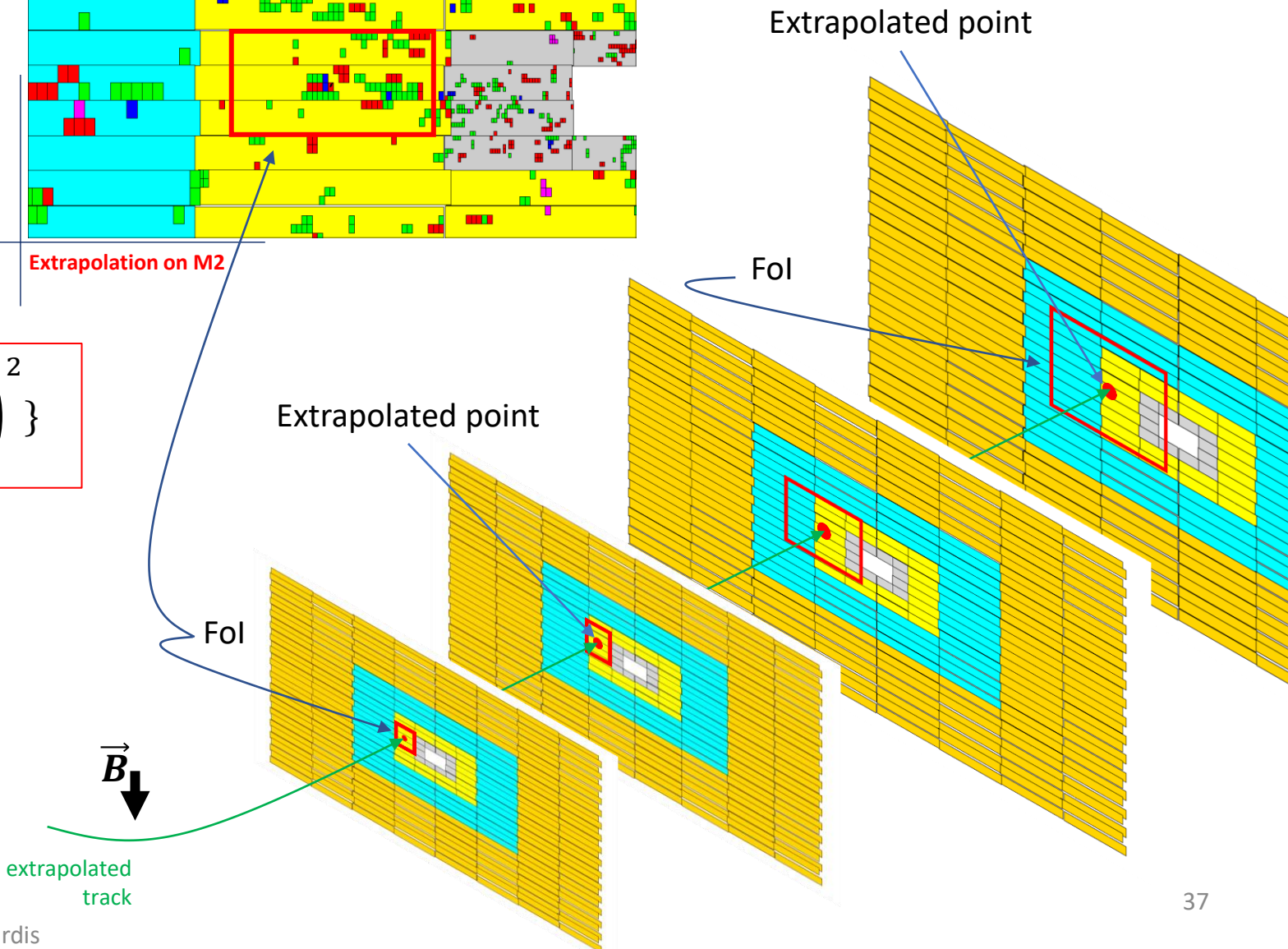
**False muon hits can be identified on stations at very crowded environment.**

Development of a software to simulate the full hits generation at Run 5, in 25 ns time window :

- New *FoI* definitions
- Different  $\chi^2_{corr}$  cut



Extrapolation on M2

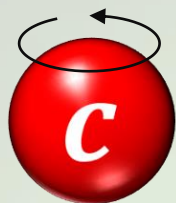


# $D_s$ SPECTROSCOPY

## In the heavy quark limit model



At high energy scale,  $\alpha_s$  is small: perturbative interactions and at short distance scales  $\ll R_{had}$



Strong coupling constant

$$\alpha_s(Q^2) \propto \frac{12\pi}{\ln(Q^2/\Lambda_{\text{QCD}}^2)}$$

$\Lambda_{\text{QCD}} \approx 0.2 \text{ GeV}$  is the energy scale that separates small from large distance scale.

At small energy scale,  $\alpha_s$  is large:

interactions with non perturbative confinement phenomena of quarks and gluons at large length scale

$$R_{had} \sim \frac{1}{\Lambda_{\text{QCD}}} \approx 1 \text{ fm}$$

- The **heavy quark** is much smaller than the hadron size and it is surrounded by a strongly interacting cloud of light quarks, antiquarks, and gluons with which the **light quark** interacts.
- The heavy-quark spin effect is not seen by the light quarks, such that the **heavy-quark spin decouples**.

$$m_c \approx 1.3 \text{ GeV}/c^2$$

$$m_s \approx 93 \text{ MeV}/c^2$$

# The $P$ -wave ( $L = 1$ ) states

Phys. Rev. Lett. 66 (1991) 1130

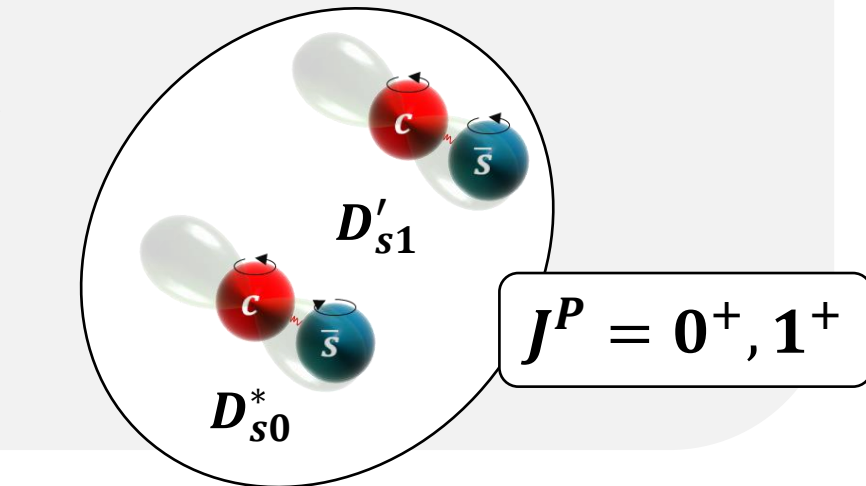
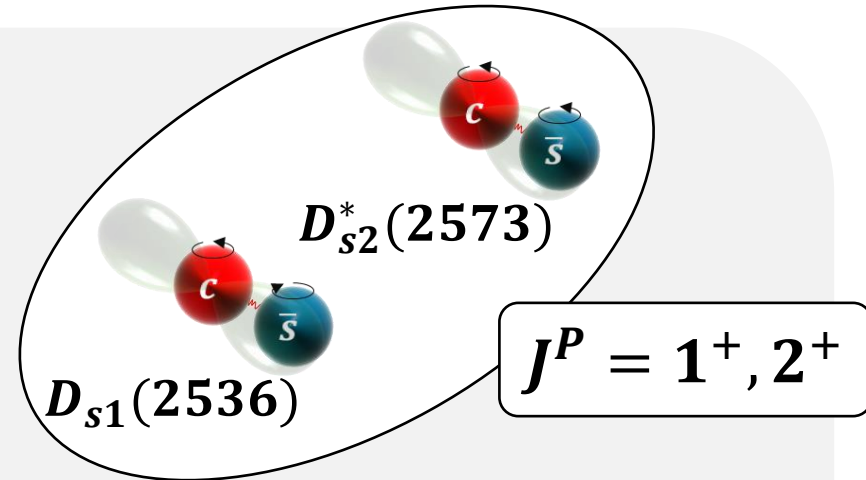
Phys. Rev. D 12 (1975) 147

## In the heavy quark limit model

Bounding force depends on the spin-orbit coupling only of the light  $\bar{s}$  quark

Excited  $P$ -states ( $\vec{L} = 1$ )

$$L + S_{\bar{s}} \begin{cases} \nearrow & j_{\bar{s}} = 3/2 \\ \searrow & j_{\bar{s}} = 1/2 \end{cases} + S_c = 1/2$$



# $D_s$ SPECTROSCOPY

MESON STATE  
 $2S+1 L_J$

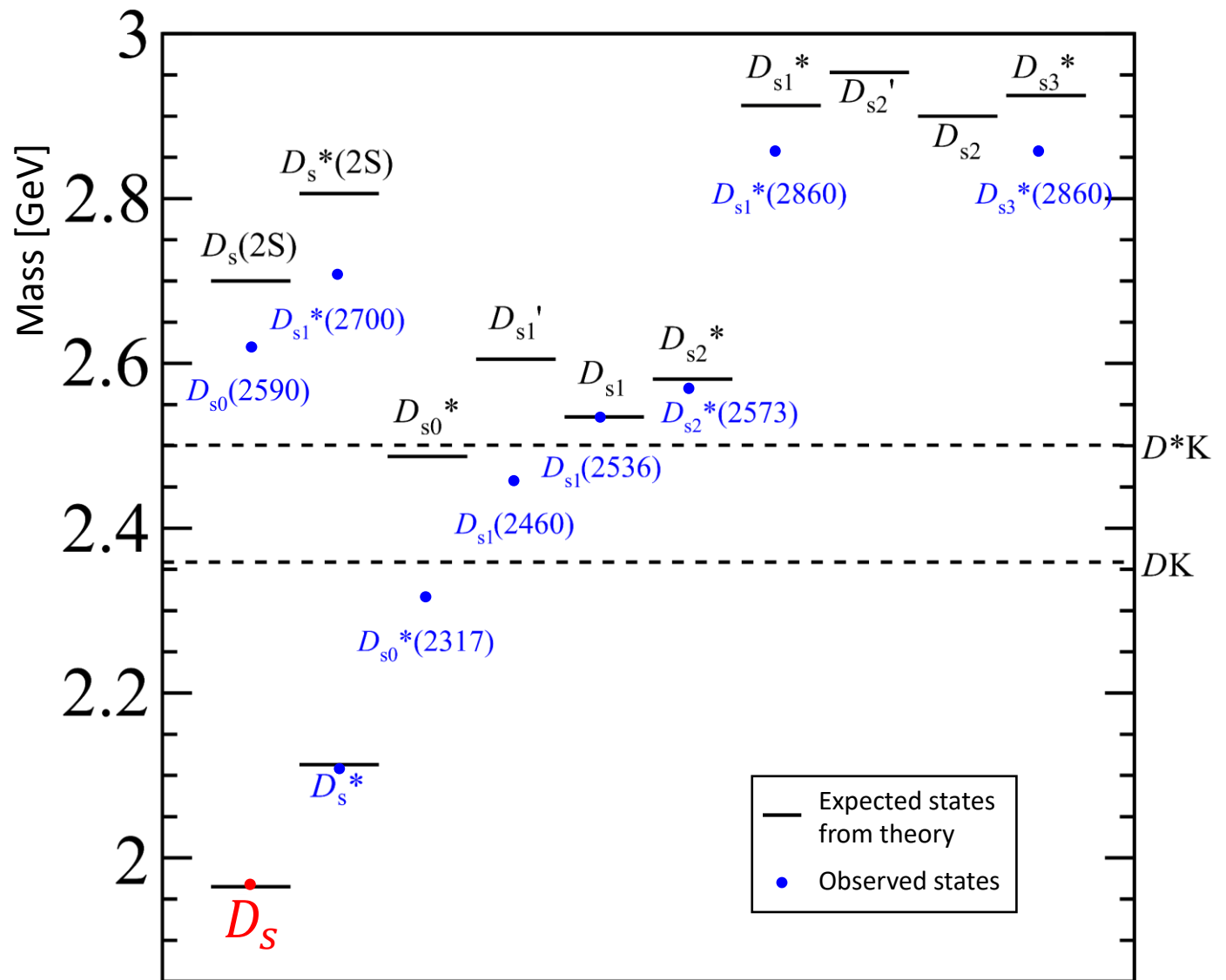
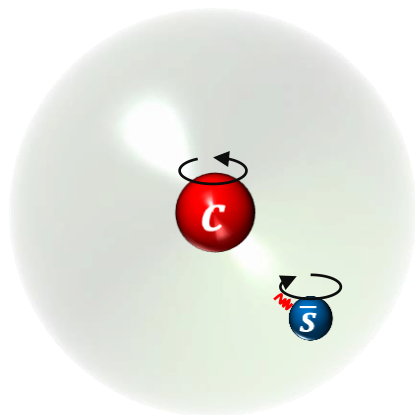
$L$ : orbital angular momentum between constituent quarks  
 ( $S, P, D$  correspond to  $L = 0, 1, 2$ )

$S = 0, 1$ : sum of quark spins

$J$ : total spins of meson state

$j_{\bar{s}}$ : sum of  $L$  with the spin of light  $\bar{s}$  quark

$P$ : parity



$2S+1 L_J$	$^1S_0$	$^3S_1$	$^1P_0$	$^3P_1$	$^1P_1$	$^3P_2$	$^1D_1$	$^3D_2$	$^1D_2$	$^3D_3$
$j_{\bar{s}}$	1/2	1/2	1/2	1/2	3/2	3/2	3/2	3/2	5/2	5/2
$J^P$	$0^-$	$1^-$	$0^+$	$1^+$	$1^+$	$2^+$	$1^-$	$2^-$	$2^-$	$3^-$



# $D_s^+ \mu^+ \mu^-$ selection

Signal and background samples

**RUN 2**  
(2015-2018)  
 $6 \text{ fb}^{-1}$

DATA sample at disposal:

$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^- + c. c.$

$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^+ + c. c.$

$D_{s1}(2460)^+ \rightarrow D_s^+ \mu^- \mu^- + c. c.$

- $D_s$  candidates:**

multivariate selection application

- $D_{s1}(2460)$  candidates:**

-requirements on  $D_s$  mass and PV constraints

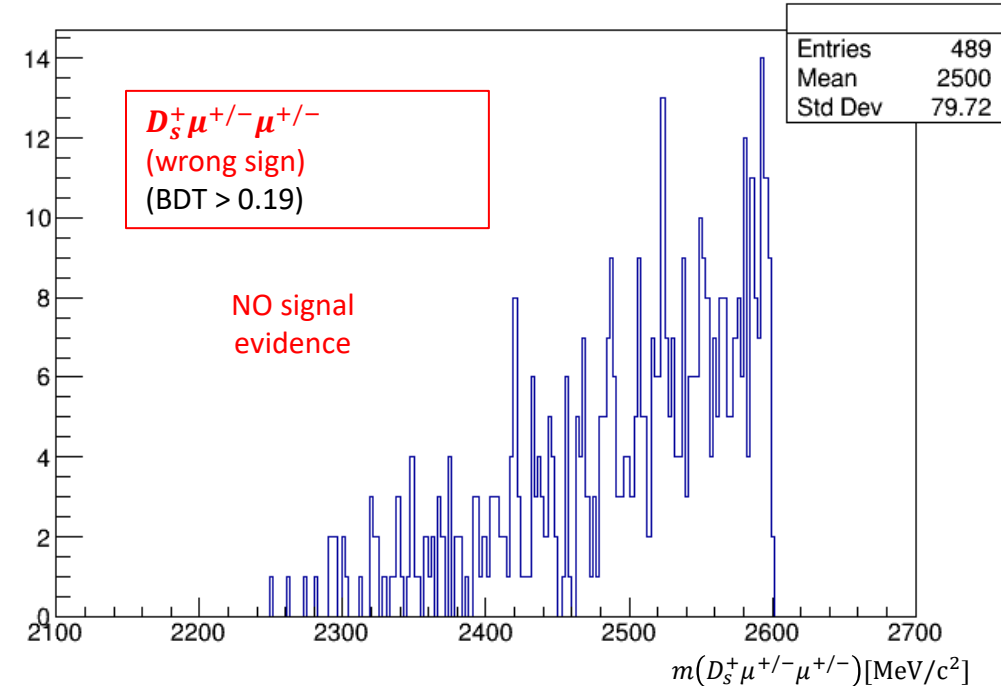
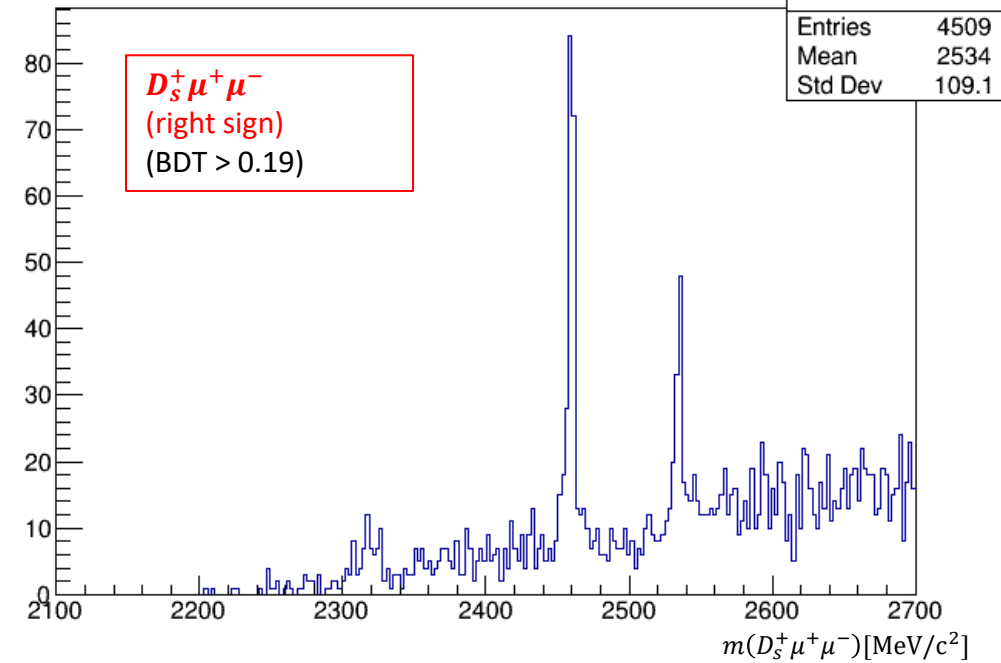
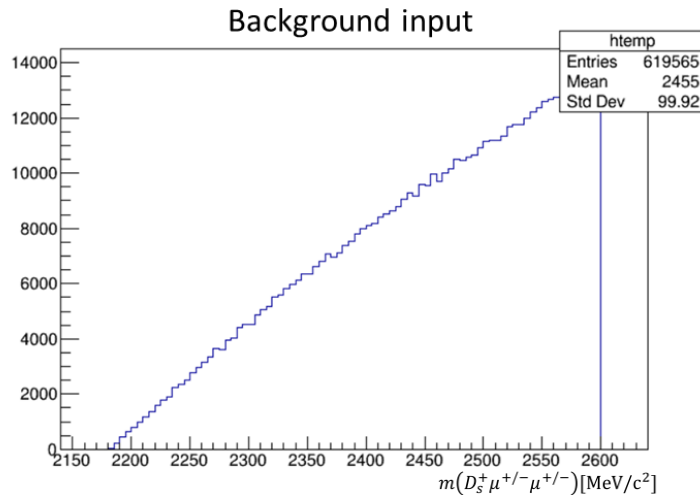
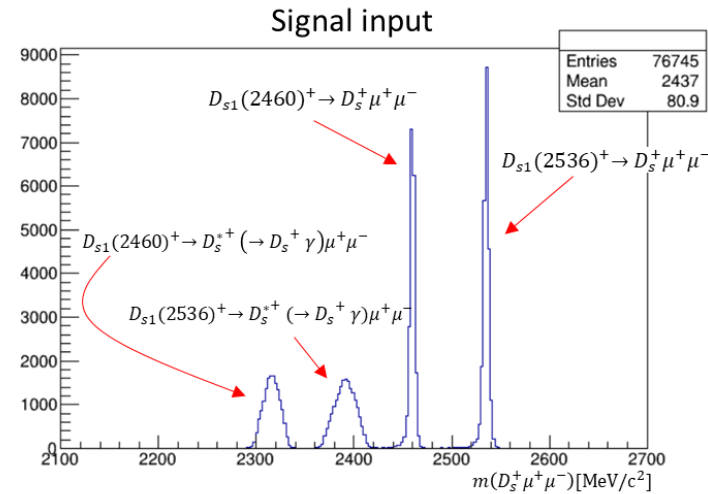
-BDT and BDTG algorithms for signal and background events classifications.

- **A mixture of four MC samples as signal input**  
( $\frac{1}{2}$  for training and  $\frac{1}{2}$  for testing)

- **Wrong Signs sample**

$m(D_s^+ \mu^+ \mu^-) < 2600 \text{ MeV}/c^2$  used as a proxy of **combinatorial background**  
( $\frac{1}{2}$  for training and  $\frac{1}{2}$  for testing)

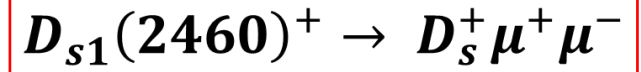
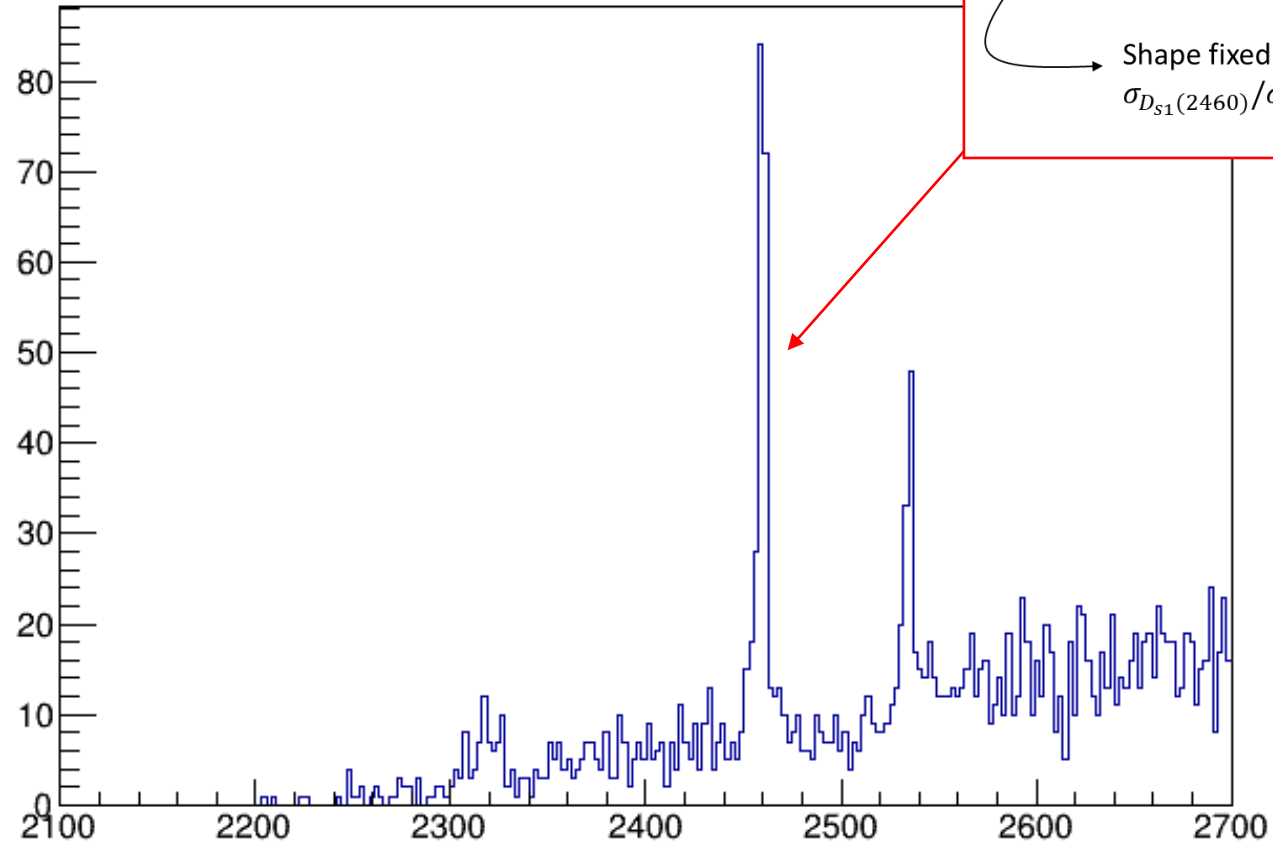
-PUNZI function used for BDT cut optimisation



# Fit model

( $PDF_{total}$ )

RUN 2  
(2015-2018)  
 $6 fb^{-1}$



Hypatia  $\otimes$  Breit-Wigner

$(M, \Gamma)$  floated

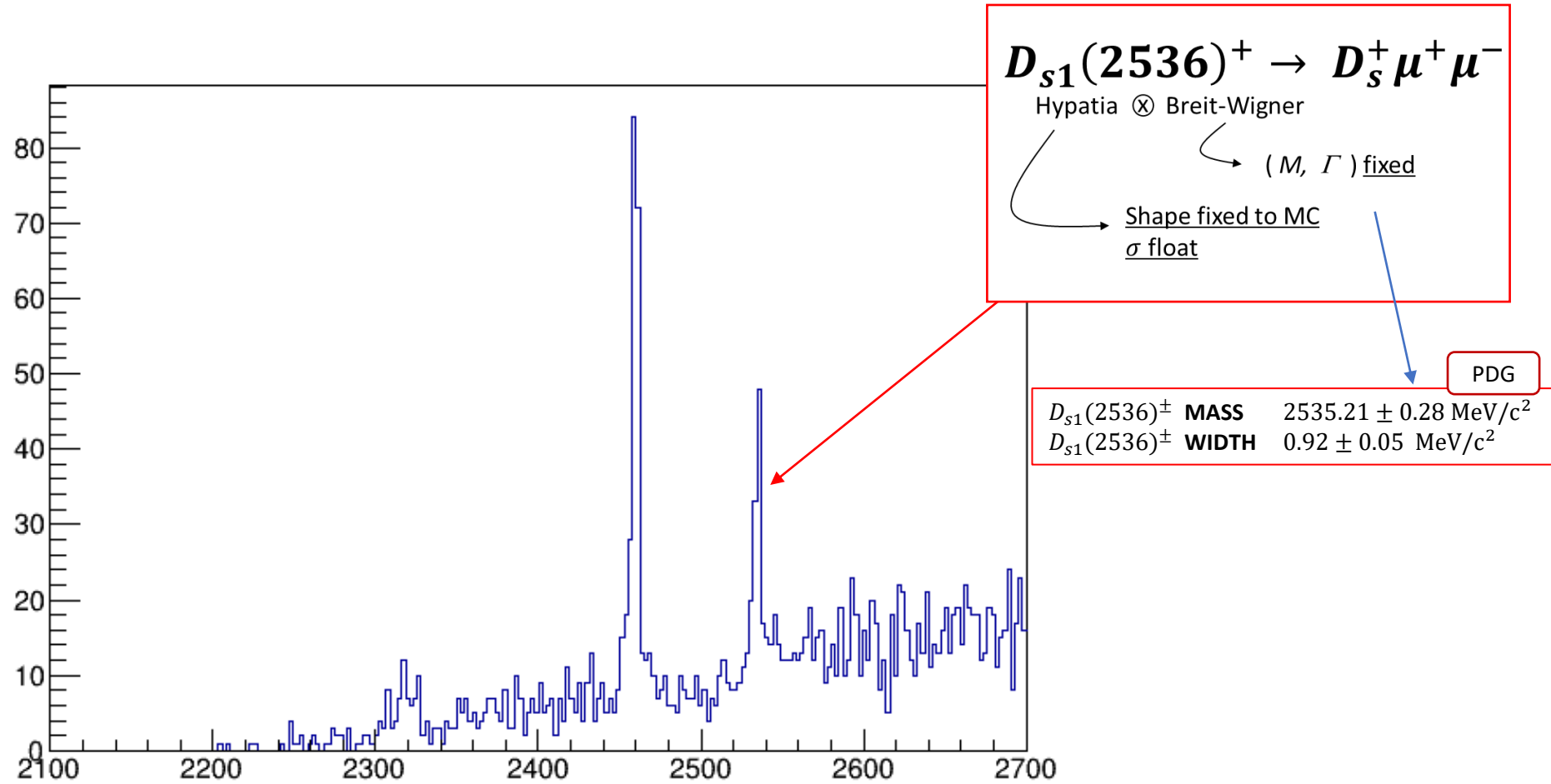
Shape fixed to MC

$\sigma_{D_{s1}(2460)}/\sigma_{D_{s1}(2536)}$  fixed to MC

# Fit model

( $PDF_{total}$ )

RUN 2  
(2015-2018)  
 $6 fb^{-1}$



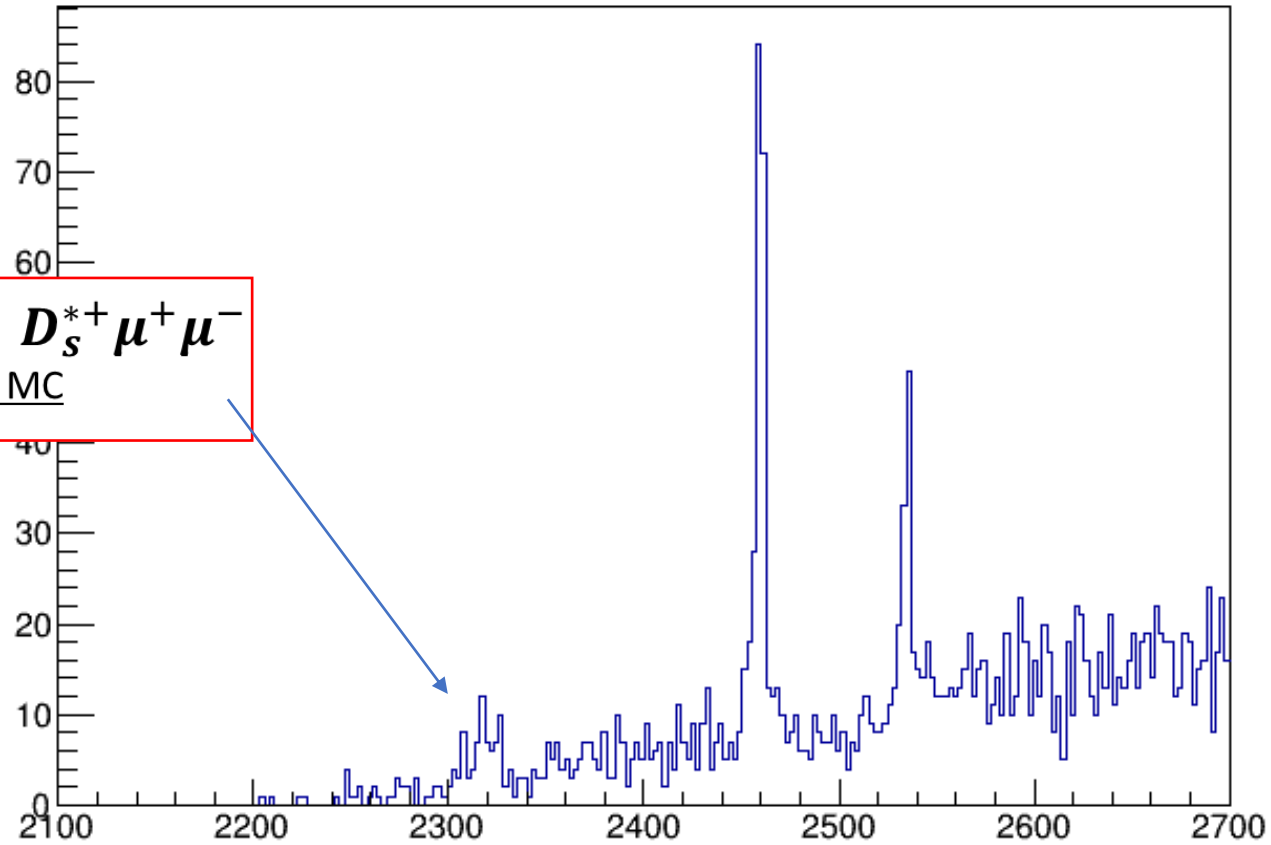
# Fit model

( $PDF_{total}$ )

RUN 2  
(2015-2018)  
 $6 fb^{-1}$



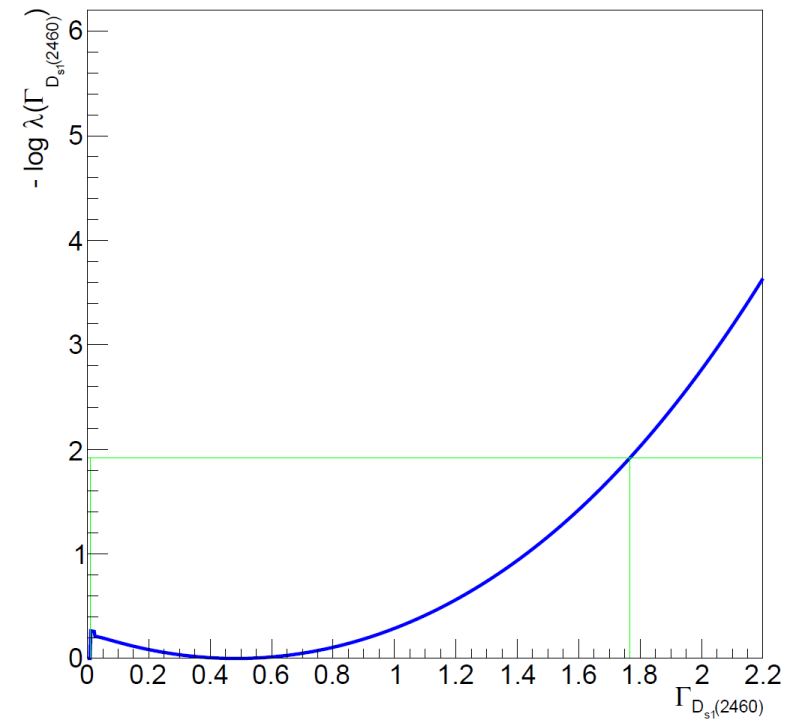
$D_{s1}(2460)^+ \rightarrow D_s^{*+} \mu^+ \mu^-$   
• Shape fixed to MC



# Measurements:

- $m_{D_{S1}(2460)} = 2459.9 \pm 0.2 \text{ MeV}/c^2$
- $\Gamma_{D_{S1}(2460)} = 0.5 \pm 0.7 \text{ MeV}/c^2 \sim 0$

log profile likelihood scan at the  $\Gamma_{D_{S1}(2460)}$  variation, such to evaluate the upper limit at 95% CL:  $\Gamma_{D_{S1}(2460)} < 1.77$

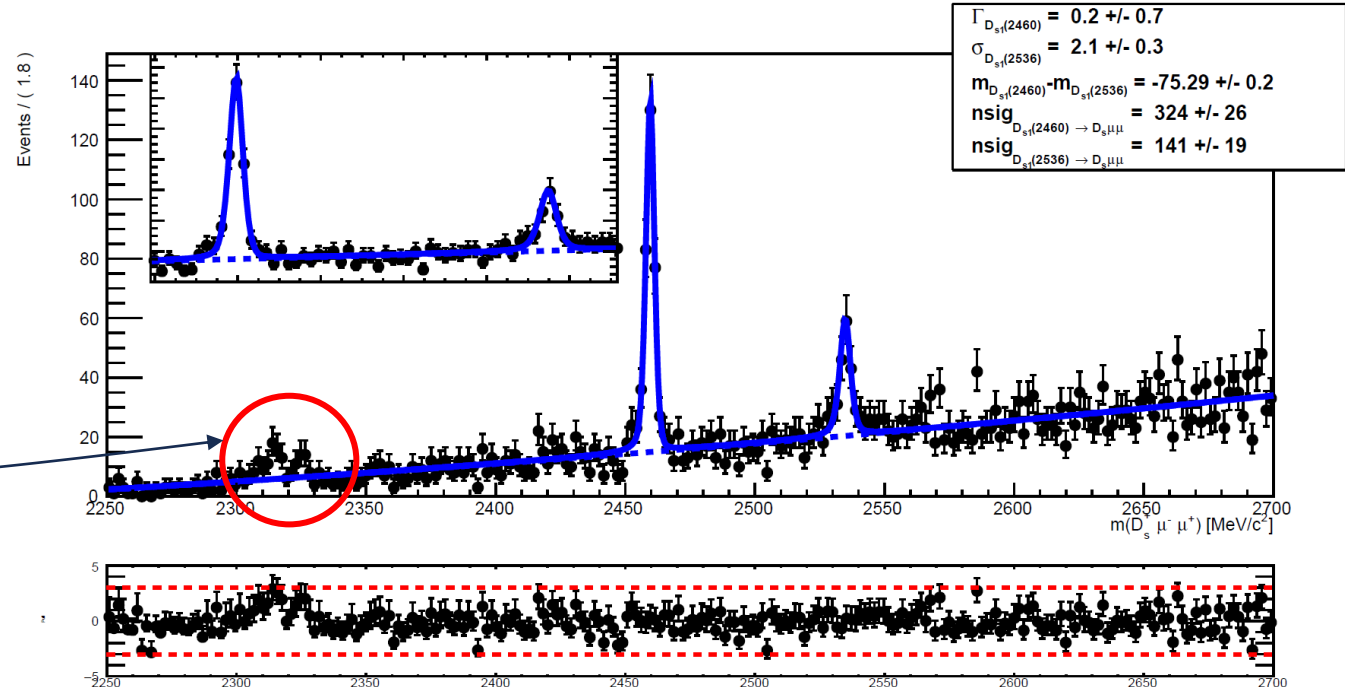


Peak presence associated to the decay events:



Peak significance estimation is necessary

It is done with the difference of the log-likelihood at minimum between the previous best-fit and the one without the PDF contribution.



Then:

$$n_\sigma = 7.7\sigma > 5\sigma$$

**NEW DECAY OBSERVATION**

# Measurements:

Relative branching ratio measurement of:  
 $D_{S1}(2460)^+ \rightarrow D_s^{*+} (\rightarrow D_s^{*+} \gamma / \pi^0) \mu^+ \mu^-$

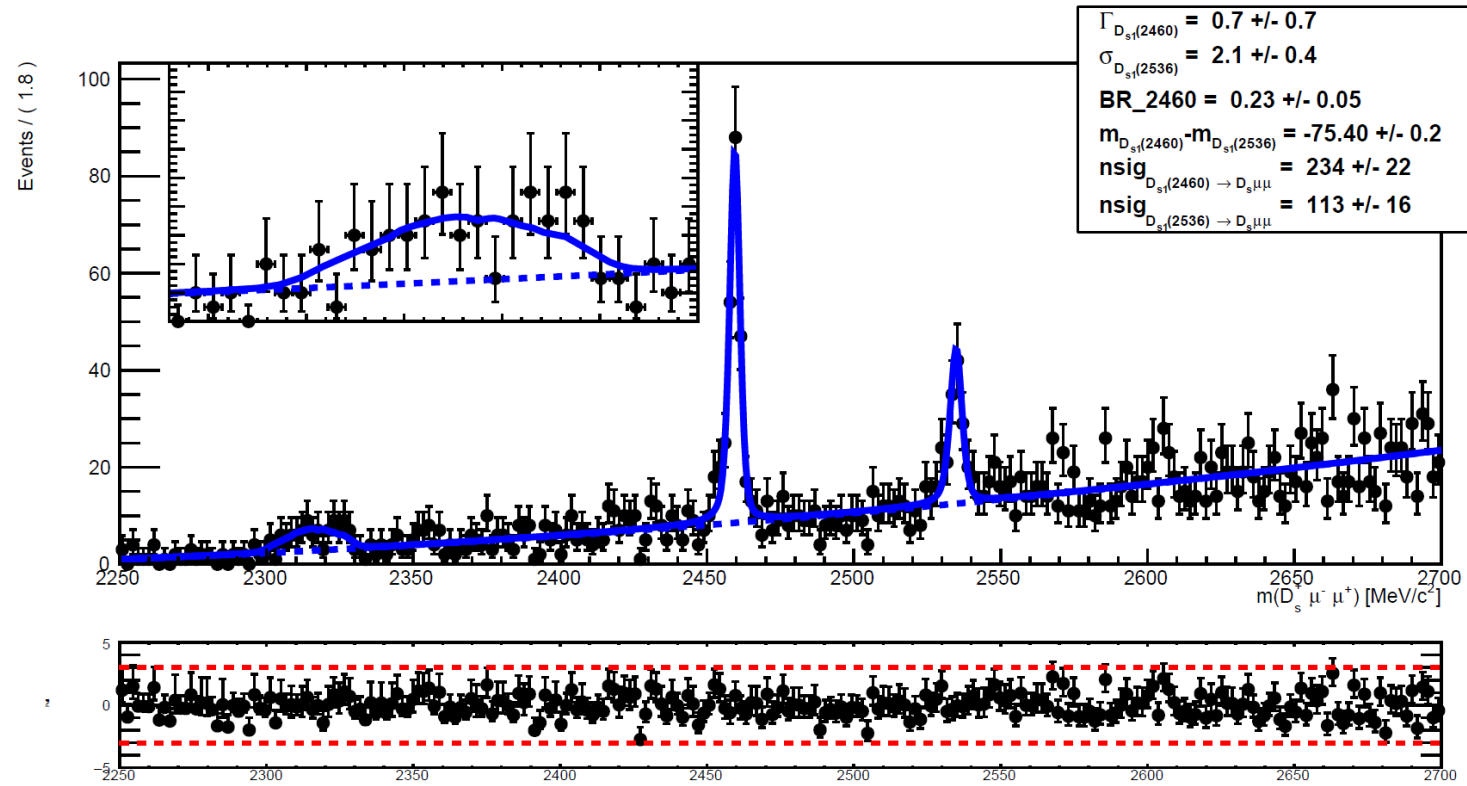
- Trigger constraint application:  
only the Run 2 events satisfying the decay criteria are selected

- Best-fit to triggered data,  
to measure the ratio BR\_2460:

$$B_{2460} = \frac{nsig(D_{S1}(2460)^+ \rightarrow D_s^{*+} \mu^+ \mu^-)}{nsig(D_{S1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^-)} = 0.23 \pm 0.05 \%$$

- Evaluations of the LHCb reconstruction efficiencies,  $\epsilon_1, \epsilon_2$  for both  $D_{S1}(2460)^+$  decays, by **MC samples**:

$$\epsilon_{1(2)}(D_{S1}(2460)^+ \rightarrow D_s^{(*)+} \mu^+ \mu^-) = \frac{\#event\ reconstructed}{\#event\ generated}$$



$$\frac{\mathcal{B}(D_{S1}(2460) \rightarrow D_s^* \mu^+ \mu^-)}{\mathcal{B}(D_{S1}(2460) \rightarrow D_s \mu^+ \mu^-)} = BR_{2460} \cdot \frac{\epsilon_1}{\epsilon_2} = 0.48 \pm 0.10$$

Encouraged to search for the decay signal:  
 $D_{S1}(2536) \rightarrow D_s^* \mu^+ \mu^-$

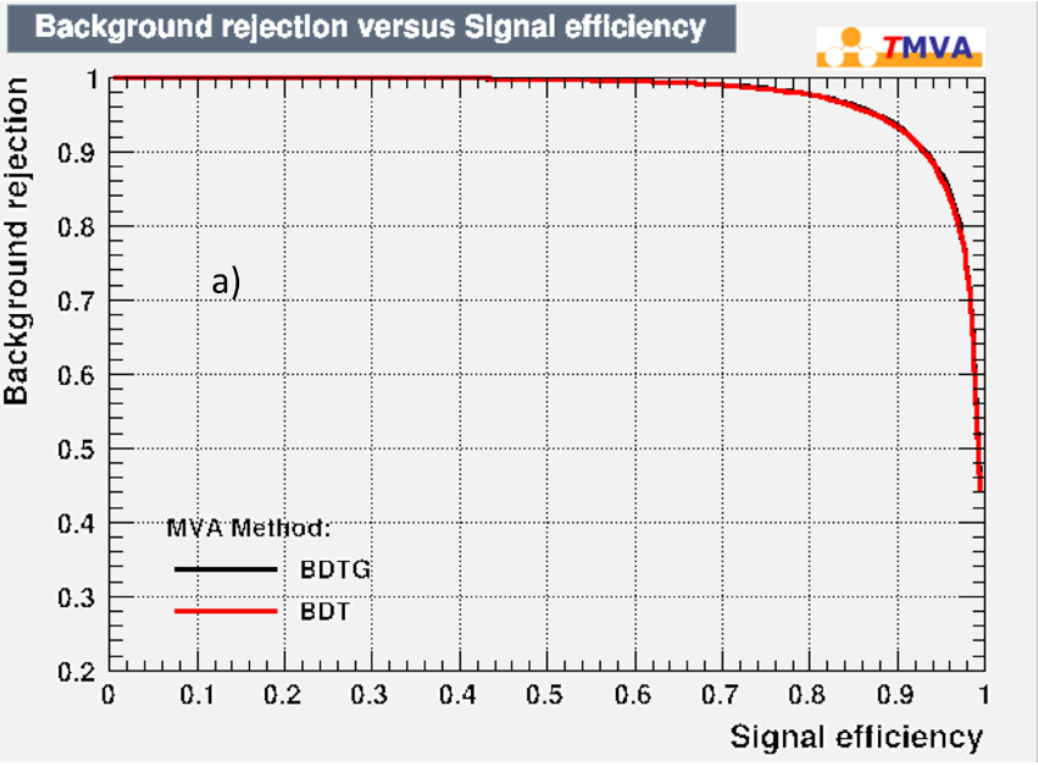
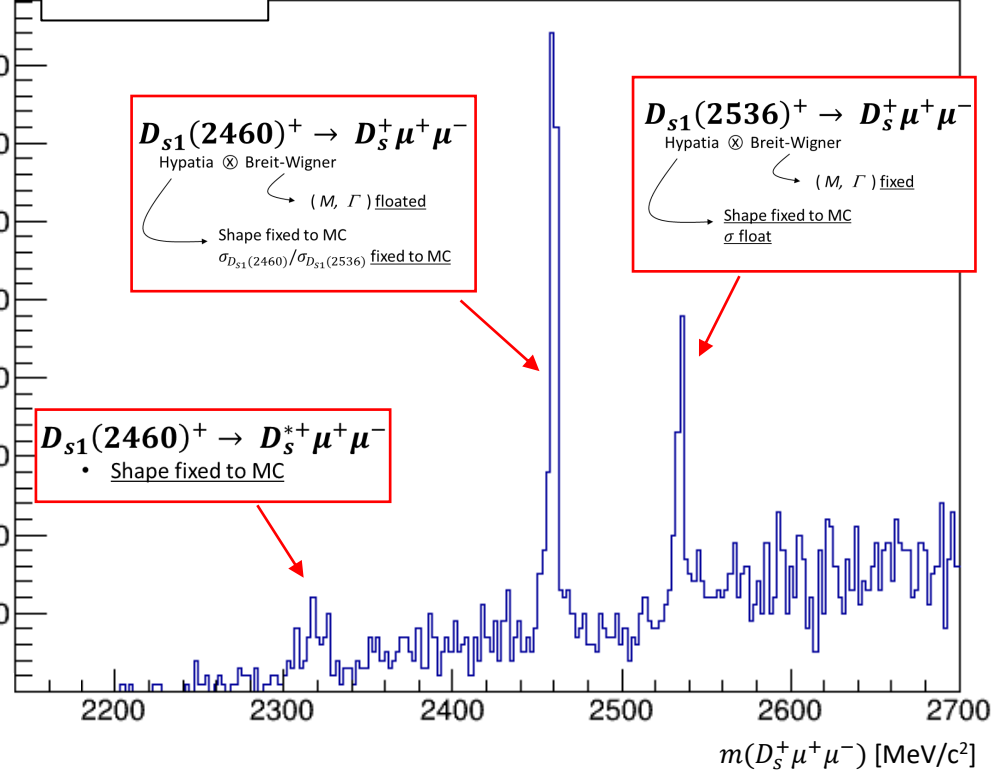
# Fit result:

RUN 2  
(2015-2018)  
 $6 \text{ fb}^{-1}$



## Fit model

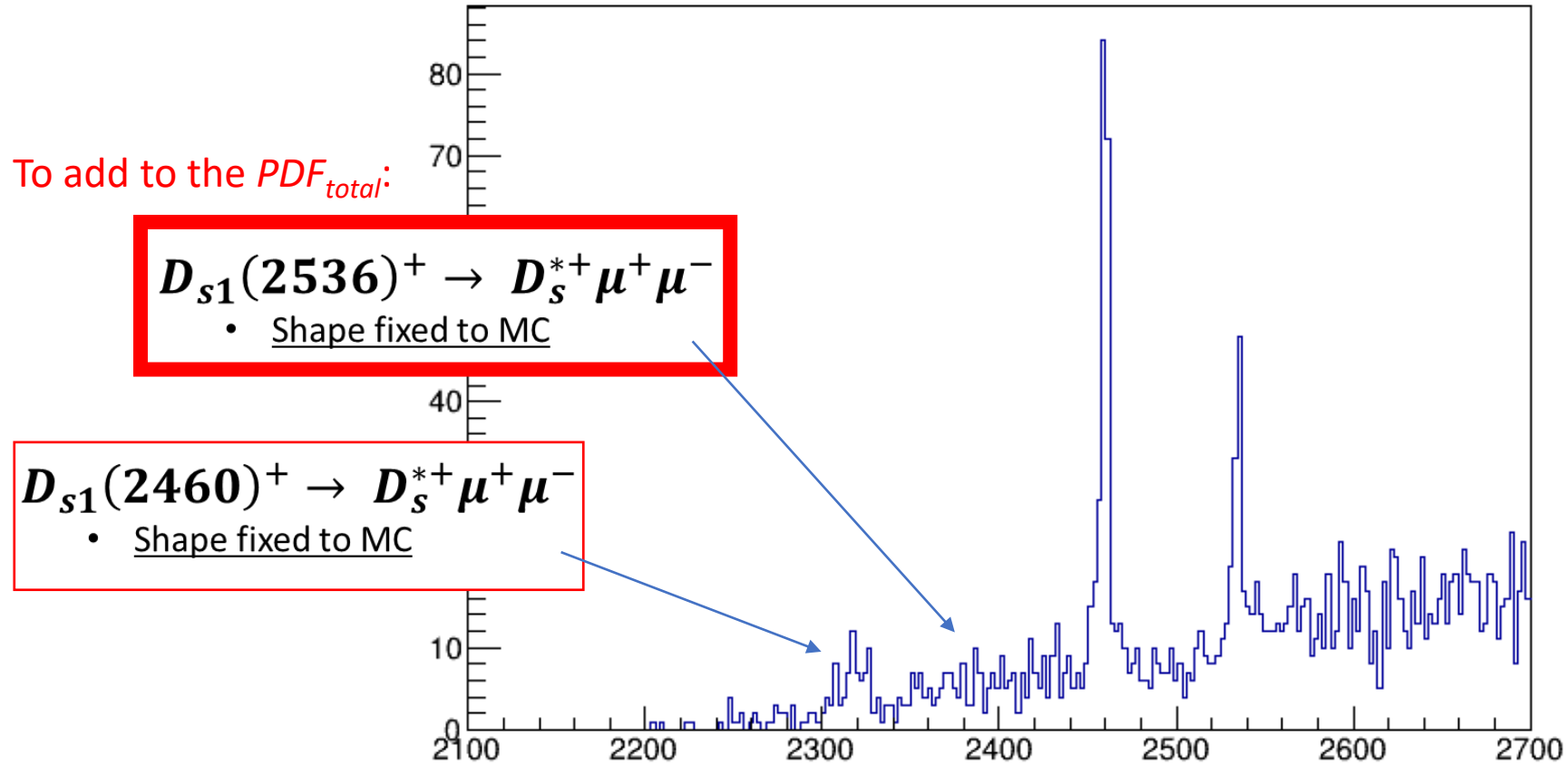
( $PDF_{total}$ )



ROC curves from the BDT and the BDTG algorithms used for the multivariate  $D_s^+ \mu^+ \mu^-$  selection. It is evident that the algorithms have about same performance (integral ROC  $\sim 0.97$ ). However, the BDT classification algorithm is chosen.

# Fit model

RUN 2  
(2015-2018)  
 $6 \text{ fb}^{-1}$





# Measurements:

Relative branching ratio measurement of:  
 $D_{s1}(2536)^+ \rightarrow D_s^{*+} (\rightarrow D_s^{*+} \gamma / \pi^0) \mu^+ \mu^-$

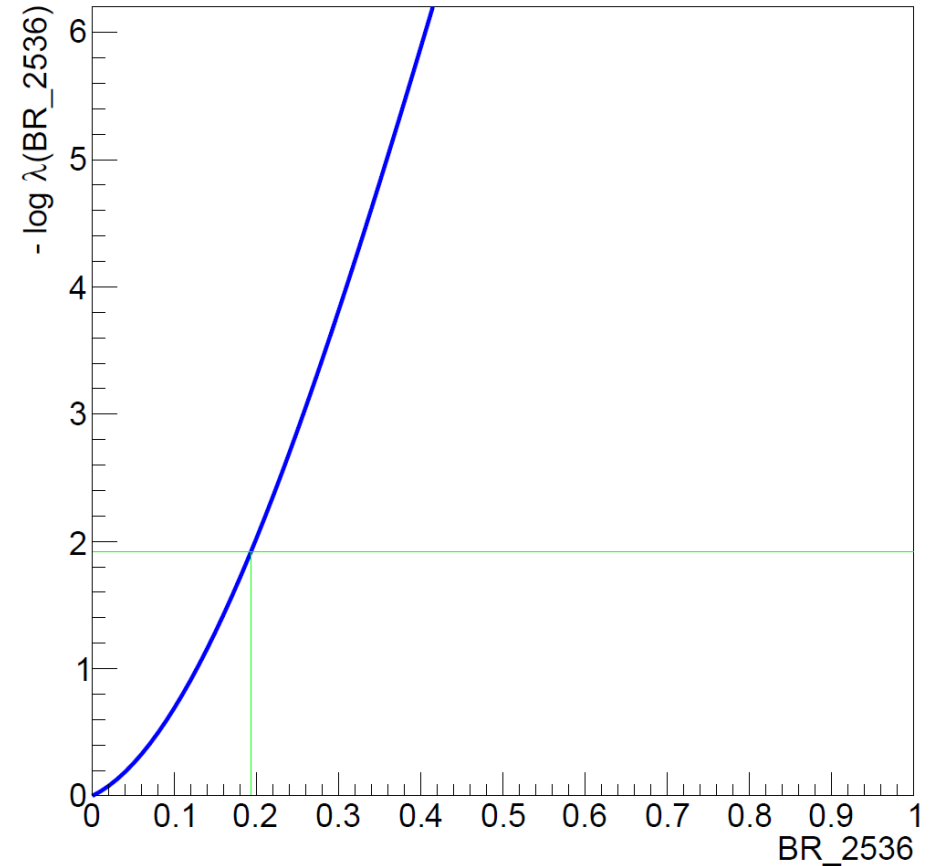
- Best-fit to triggered data, to measure the ratio BR\_2536:

$$B_{2536} = \frac{nsig(D_{s1}(2536)^+ \rightarrow D_s^{*+} \mu^+ \mu^-)}{nsig(D_{s1}(2536)^+ \rightarrow D_s^+ \mu^+ \mu^-)} = 0.0 \pm 0.1 \% \sim 0$$

- log profile likelihood scan for the upper limit at 95% CL:  
 $BR_{2536} < 0.193$

- Evaluations of the LHCb reconstruction efficiencies,  $\epsilon'_1, \epsilon'_2$  for both  $D_{s1}(2536)^+$  decays, by **MC samples**:

$$\epsilon'_{1(2)} (D_{s1}(2536)^+ \rightarrow D_s^{(*)+} \mu^+ \mu^-) = \frac{\#event\ reconstructed}{\#event\ generated}$$



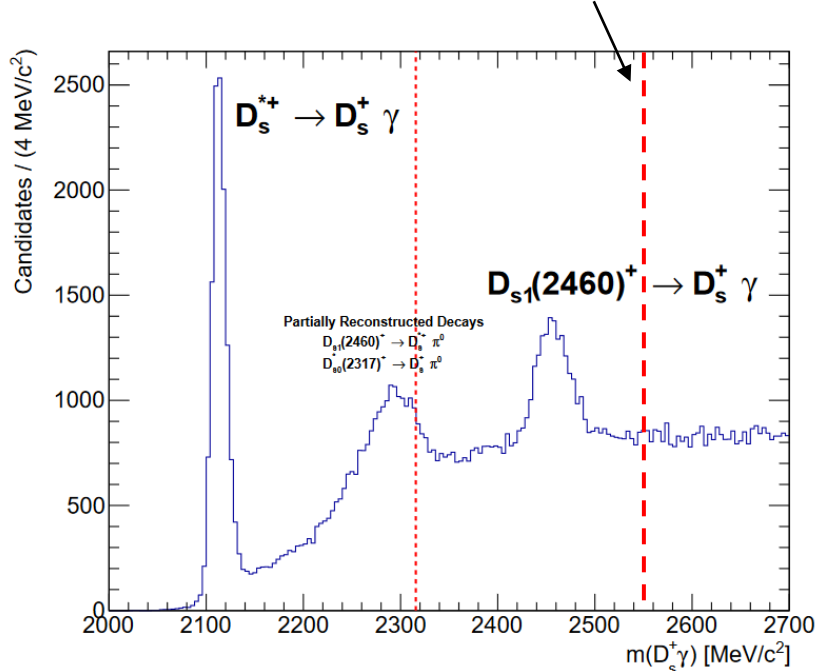
$$\frac{\mathfrak{B}(D_{s1}(2536) \rightarrow D_s^* \mu^+ \mu^-)}{\mathfrak{B}(D_{s1}(2536) \rightarrow D_s \mu^+ \mu^-)} = BR_{2536} \cdot \frac{\epsilon'_1}{\epsilon'_2} < 0.133$$

# Summary:

**RUN 2**  
(2015-2018)  
 $6 \text{ fb}^{-1}$

## Observation of 3 signals/structures:

- First observation of  $D_{s1}(2460) \rightarrow D_s \mu^+ \mu^-$  decays.
- Evidence/observation of  $D_{s1}(2460) \rightarrow D_s^* \mu^+ \mu^-$  decays.
- First observation of the puzzling decay of  $D_{s1}(2536) \rightarrow D_s \mu^+ \mu^-$  despite the **lack of  $D_s \gamma$  decay observation.**



- Radiative decays largely suppressed due to dominance of strong decays to  $D^{*+} K^0$  decay mode [PR D 2016, 93.3: 034035]

- $D_{s1}(2536)^+ \rightarrow D_s^+ \eta$  ( $BR(\eta \rightarrow \mu^+ \mu^-) \cong 6 \cdot 10^{-6}$ ) forbidden for the angular momentum conservation:

$$\vec{J}_{D_{s1}(2536)} = \vec{1} = \vec{J}_{final} = \vec{L} + \vec{S}_{D_s} + \vec{S}_\eta = \vec{L}$$

$$P_{D_{s1}(2536)} = +1 \neq P_{final} = (-1)(-1)(-1)^L = -1$$

- $D_{s1}(2536)^+ \rightarrow D_s^+ \omega$  ( $BR(\omega \rightarrow \mu^+ \mu^-) \cong 7 \cdot 10^{-5}$ ) possible, according to the angular momentum, parity and isospin conservation:

$$\left. \begin{aligned} \vec{J}_{D_{s1}(2536)} = 1 = \vec{J}_{final} = \vec{L} + \vec{S}_{D_s} + \vec{S}_\omega = \vec{L} + 1 \\ P_{D_{s1}(2536)} = +1 \neq P_{final} = (-1)(-1)(-1)^L = +1 \\ I_{D_{s1}(2536)} = 0 = I_{final} = 0 + 0 \end{aligned} \right\} \vec{L} = 0, 2$$

- $D_{s1}(2536)^+ \rightarrow D_s^+ \rho$  ( $BR(\rho \rightarrow \mu^+ \mu^-) \cong 5 \cdot 10^{-5}$ ) possible, according to the angular momentum and parity but the isospin conservation:

$$I_{D_{s1}(2536)} = 0 \neq I_{final} = 0 + 1 = 1$$

N.B.  
 $D_{s1}(2460)^+ \rightarrow D_s^{*+} \pi^0$   
is also an isospin violating decay but  $BR \sim 50\%$



# The puzzle of the $D_{s1}(2460)$ meson



**Surprising observation:**  
 $D_{s1}(2460)^+ \rightarrow D_s^{*+} \pi^0$   
**CLEO (2003)**

- mass smaller than all theoretical predictions
- observed in the isospin-violating  $D_s^* \pi^0$  channel
- narrow state

## Theoretical predictions

- Broad  $c\bar{s}$  meson state
- Massive enough that dominant strong decays would have been isospin-conserving in  $D^*K$  final states.

Presented at:

[EPJ Web Conf. 270 \(2022\) 00013](https://arxiv.org/abs/2201.00013)

**In a  $c\bar{s}$  scenario:**  
isospin-violating decay is suppressed.

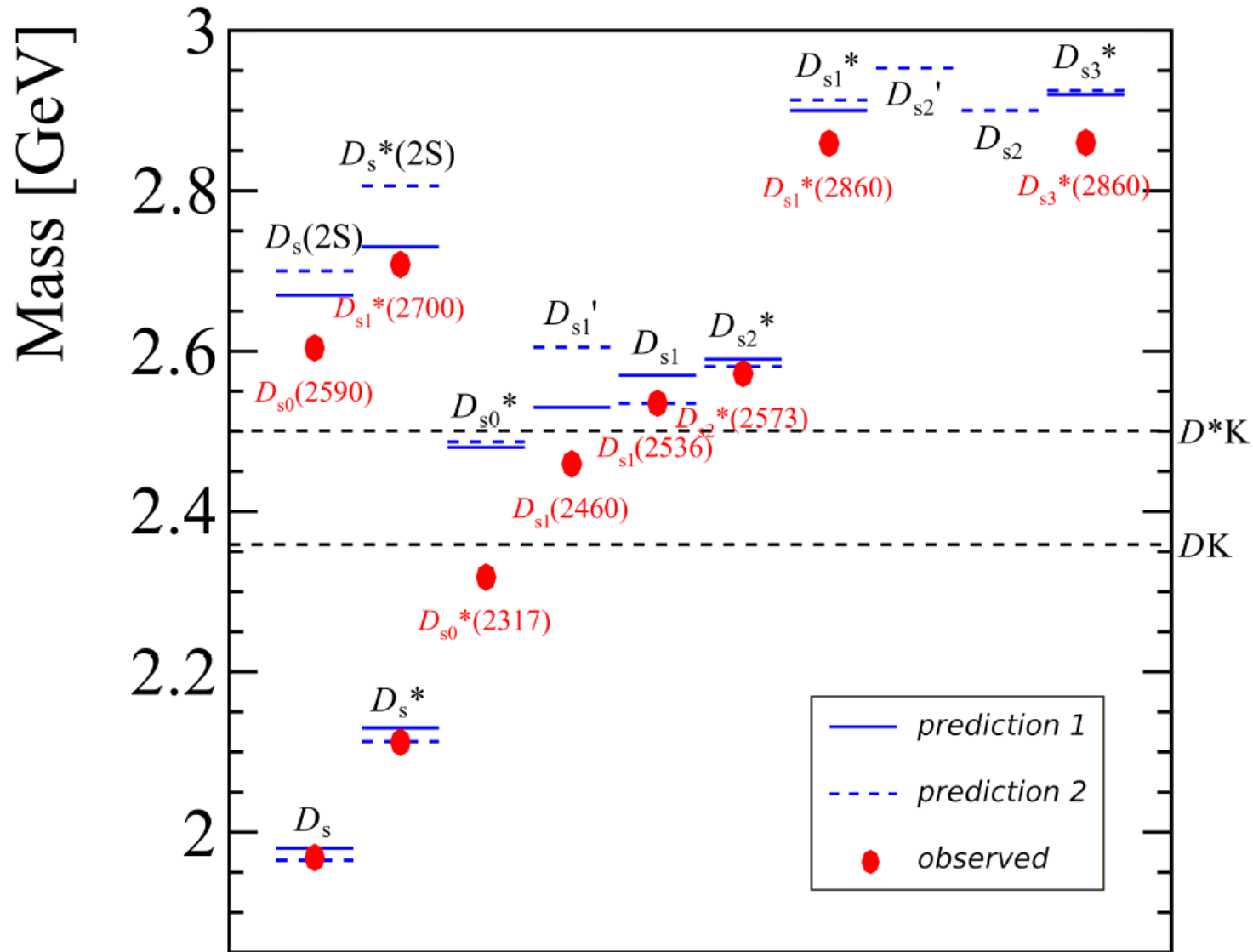
**For the  $D_{s1}(2460)$ :**  
isospin-violating decay  $D_s^* \pi^0$  is much frequent, suggesting a 4-quarks system ( $c\bar{s}u\bar{u}$  or  $c\bar{s}d\bar{d}$ )

$D_{s1}(2460)^+$

decay modes	
$D_s^{*+} \pi^0$	$(48 \pm 11)\%$
$D_s^+ \gamma$	$(18 \pm 4)\%$
$D_s^+ \pi^+ \pi^-$	$(4.3 \pm 1.3)\%$
<b>TOT</b>	<b><math>(70 \pm 16)\%</math></b>

$D_{s1}(2460)^\pm$ <b>MASS</b>	$2459.5 \pm 0.6 \text{ MeV}/c^2$
$D_{s1}(2460)^\pm$ <b>WIDTH</b>	$< 3.5 \text{ MeV}/c^2$

PDG

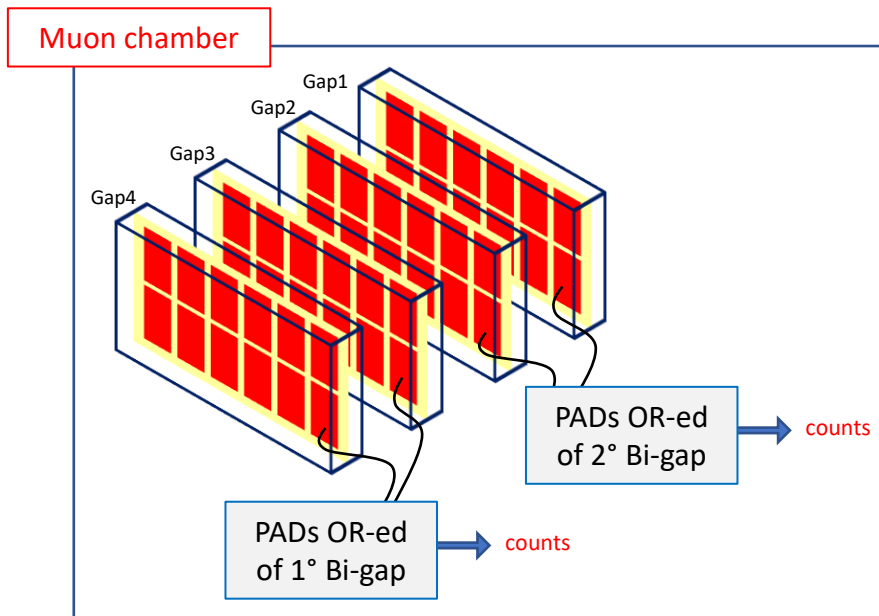


$2S+1$	$L_J$	$1S_0$	$3S_1$	$1P_0$	$3P_1$	$1P_1$	$3P_2$	$1D_1$	$3D_2$	$1D_2$	$3D_3$
$j$	1/2	1/2	1/2	1/2	3/2	3/2	3/2	3/2	3/2	5/2	5/2
$J^P$	$0^-$	$1^-$	$0^+$	$1^+$	$1^+$	$2^+$	$1^-$	$2^-$	$2^-$	$3^-$	



# Estimate of rates at Upgrade II

- Run 3 scalers data collected during November 2022, used to extrapolate rate values at the U2 luminosity
- Scaler data: counts of OR-ed bi-gaps
  - collected in 20 sec
  - at 7 different values of luminosity



Protons colliding per bunch ( $\mu$  real):

$\mu = 1.10$	• $L = 1.05 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 1.45$	• $L = 1.39 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 2.29$	• $L = 2.19 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 2.92$	• $L = 2.79 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 3.50$	• $L = 3.35 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 4.05$	• $L = 3.87 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
$\mu = 4.56$	• $L = 4.36 \times 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$

# Absorber instead of HCAL

## Rates reduction

(~80% of which is background):

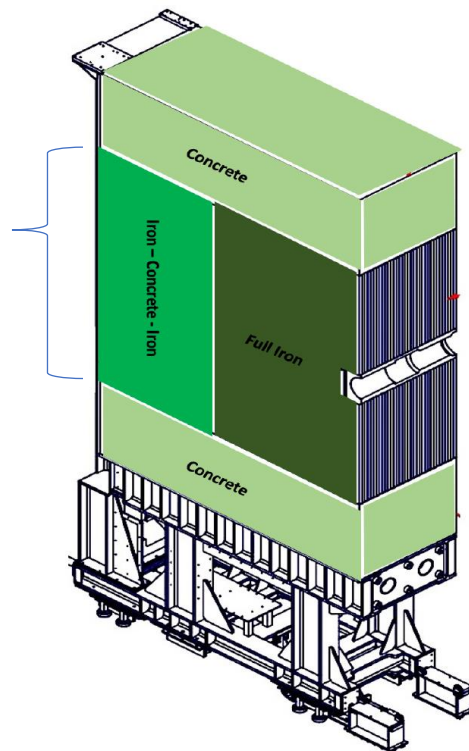
- M2R1: -42%
- M2R2: -69%
- M2R3: -64%

$\mu$ -ID efficiency loss negligible.

Evaluated with  $K_S^0 \rightarrow \mu^+ \mu^-$  events

	HCAL	Shielding
$1\mu$ , all regions	$97.5 \pm 0.2\%$	$97.7 \pm 0.2\%$
R1	$93.1 \pm 0.9\%$	$93.4 \pm 0.8\%$
R2	$98.2 \pm 0.3\%$	$98.7 \pm 0.2\%$
R3	$99.1 \pm 0.2\%$	$97.4 \pm 0.3\%$
R4	$96.9 \pm 0.4\%$	$98.8 \pm 0.2\%$
$2\mu$ , all regions	$94.8 \pm 0.4\%$	$95.4 \pm 0.3\%$

ferro-calcestruzzo-ferro  
( $6.2 \lambda_I$ ), agente su R4 di M2



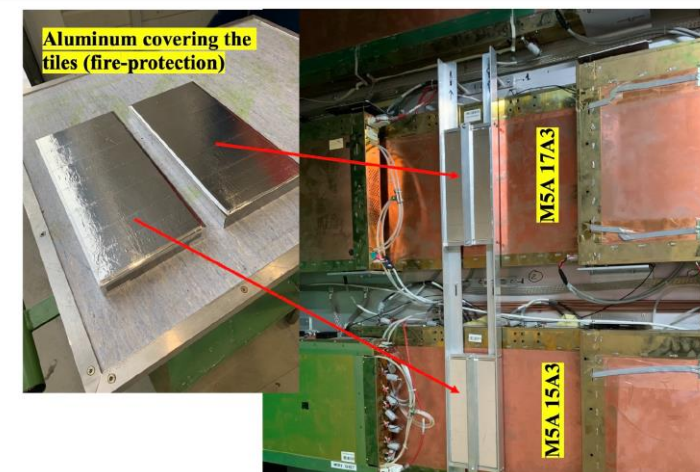
LHCB-INT-2019-008

solo ferro ( $10.1 \lambda_I$ ),  
agente su R1-R3 di M2

Calcestruzzo ( $4 \lambda_I$ ), agente su  
R4 di M2

Mirrorbor  
(under study)

For thermal neutrons  
absorption

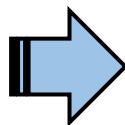




# Data usage in the analysis



**RUN 2**  
(2015-2018)  
 $6 \text{ fb}^{-1}$



1.  $D_s^+ \rightarrow K^+ K^- \pi^+ + c. c.$

Used for  $D_s^+$  decay identification with a multivariate technique

2.  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^- + c. c.$

Used for  $D_{s1}(2460)^+$  search

3.  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^+ \mu^+ + c. c.$

4.  $D_{s1}(2460)^+ \rightarrow D_s^+ \mu^- \mu^- + c. c.$

Decay channels with wrong signs used for background studies



# CHAMBERS



Table 3.1: Parameters of the present (Run 3) LHCb muon detector. Focusing on the chambers readout system, the logical (*L*) and the physical pads (anodic *A.*, and/or cathodic *C.*) granularity are also indicated. The latter consisting in anodes (*A.*) and cathodes (*C.*). The chambers areas are expressed in  $\text{cm}^2$ .

	M2	M3	M4	M5
R1 surface	0.9	1.0	1.2	1.4
N. chambers	12	12	12	12
Ch. area	30×25	32.4×27	34.8×29	37.1×30.9
L. granularity	48×8	48×8	12×8	12×8
Readout	combined	combined	Cathodes	Cathodes
A. granularity	48×1	48×1	-	-
C. granularity	8×8	8×8	12×8	12×8
R2 surface	3.6	4.2	4.8	5.5
N. chambers	24	24	24	24
Ch. area	60×25	64.8×27	69.5×29	74.3×30.9
L. granularity	48×4	48×4	12×4	12×4
Readout	combined	combined	Cathodes	Cathodes
A. granularity	48×1	48×1	-	-
C. granularity	8×8	8×8	12×4	12×4
R3 surface	14.4	16.8	19.3	22.1
N. chambers	48	48	48	48
Ch. area	120×25	129.6×27	139×29	148.5×30.9
L. granularity	48×2	48×2	12×2	12×2
Readout	Cathodes	Cathodes	Cathodes	Cathodes
A. granularity	-	-	-	-
C. granularity	48×2	48×2	24×2	24×2
R4 surface	57.7	67.2	77.4	88.3
N. chambers	192	192	192	192
Ch. area	120×25	129.6×27	139×29	148.5×30.9
L. granularity	24×1	24×1	6×1	6×1
Readout	Anodes	Anodes	Anodes	Anodes
A. granularity	24×1	24×1	24×1	24×1
C. granularity	-	-	-	-

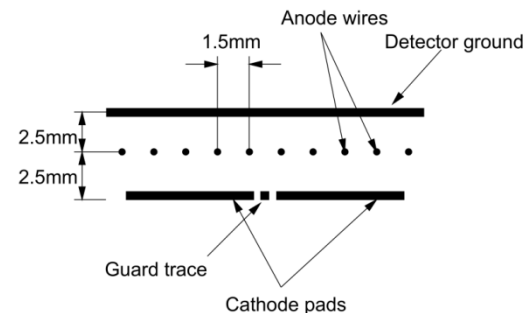
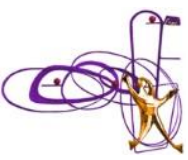


Table 3.2: Cathodes and anodes size, in  $\text{mm}^2$ , in the present (Run 3) muon detector where present.

		M2	M3	M4	M5
R1	Anodes	6.3 × 250	6.7 × 270	-	-
	Cathodes	38 × 31	41 × 34	29 × 36	31 × 39
R2	Anodes	12.5 × 250	13.5 × 270	-	-
	Cathodes	76 × 31	82 × 34	58 × 73	62 × 77
R3	Anodes	-	-	-	-
	Cathodes	25 × 125	27 × 135	58 × 145	62 × 155
R4	Anodes	50 × 250	54 × 270	58 × 290	62 × 309
	Cathodes	-	-	-	-





## Detector performances till now:

- *Geometrical acceptance:*  $10 < \theta < 300/250$  mrad (bending/non bending plane)
- *p* measurement resolution:  $\Delta p/p = 0,5\% / 1,0\%$  at low momentum / at 200 GeV/c
- *Impact parameter* resolution:  $(15 + 29/p_T[\text{GeV}]) \mu\text{m}$
- *Track reconstruction* resolution: 96% for long tracks
- *Muon identification* efficiency:  $\approx 97\%$