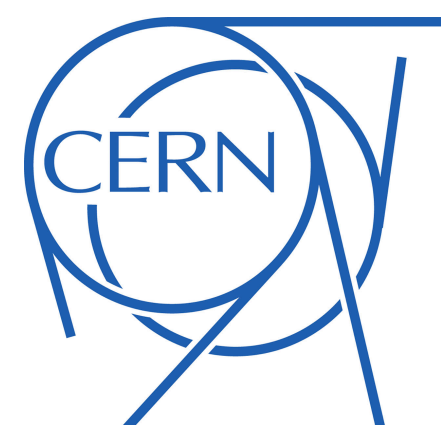


# Lattice QCD and hadronic uncertainties: status and prospects

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Workshop Italiano sulla Fisica ad Alta Intensità (WIFAI 2024)  
Bologna, 12-15 November 2024

- The **Standard Model (SM)** works very well, but it leaves many questions still unanswered

why three generations?

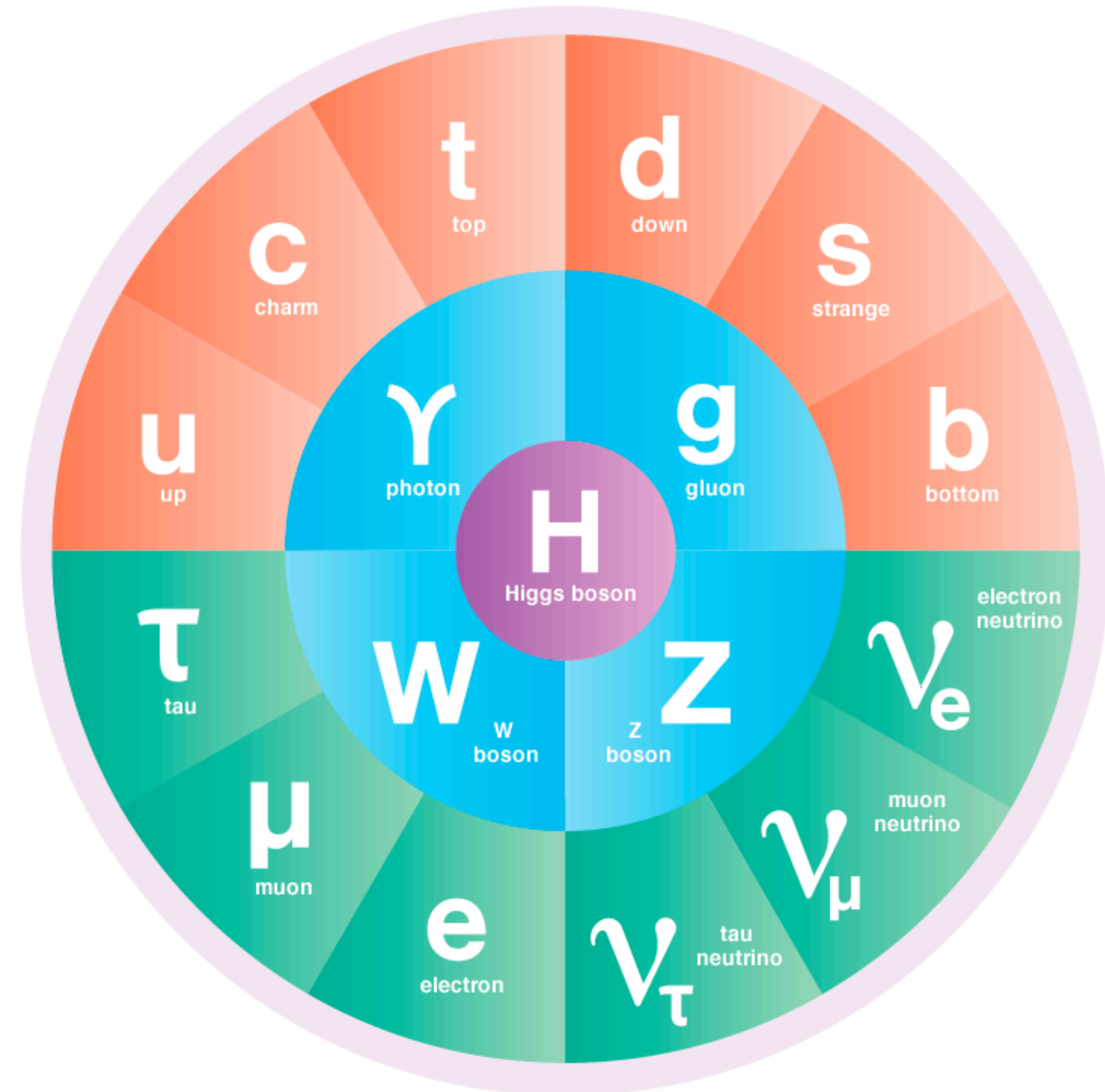
why this quark mass hierarchy?

why this pattern in quark mixing?

why such matter-antimatter asymmetry?

...

- We want to look for **new physics effects**: deviations from SM expectations could be a signal of new physics!
- Need for precise and controlled **experimental measurements** and **theoretical predictions**



# Flavour physics

Flavour physics offers opportunities to test the Standard Model and probe new physics effects

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

in the Standard Model:

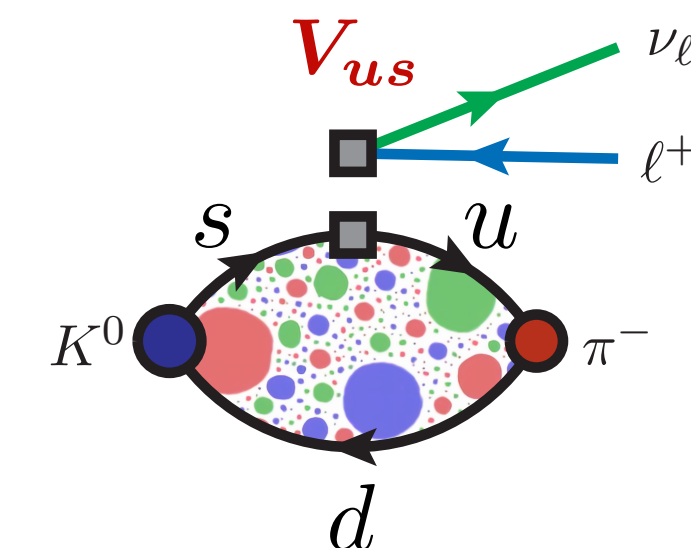
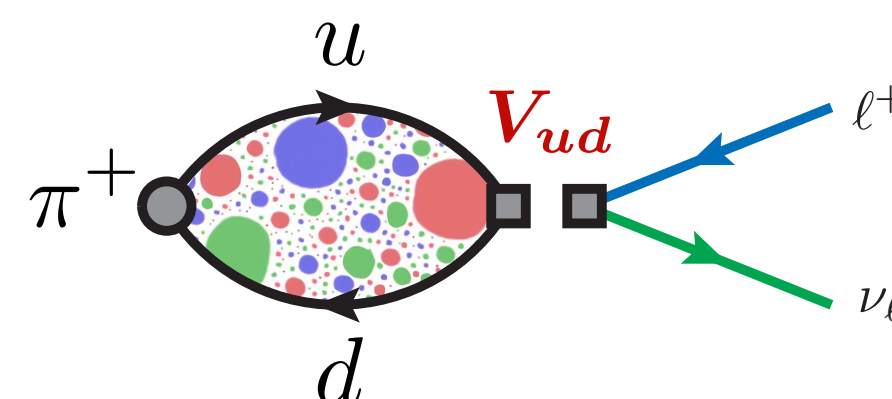
3 mixing angles + 1 CPV phase

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

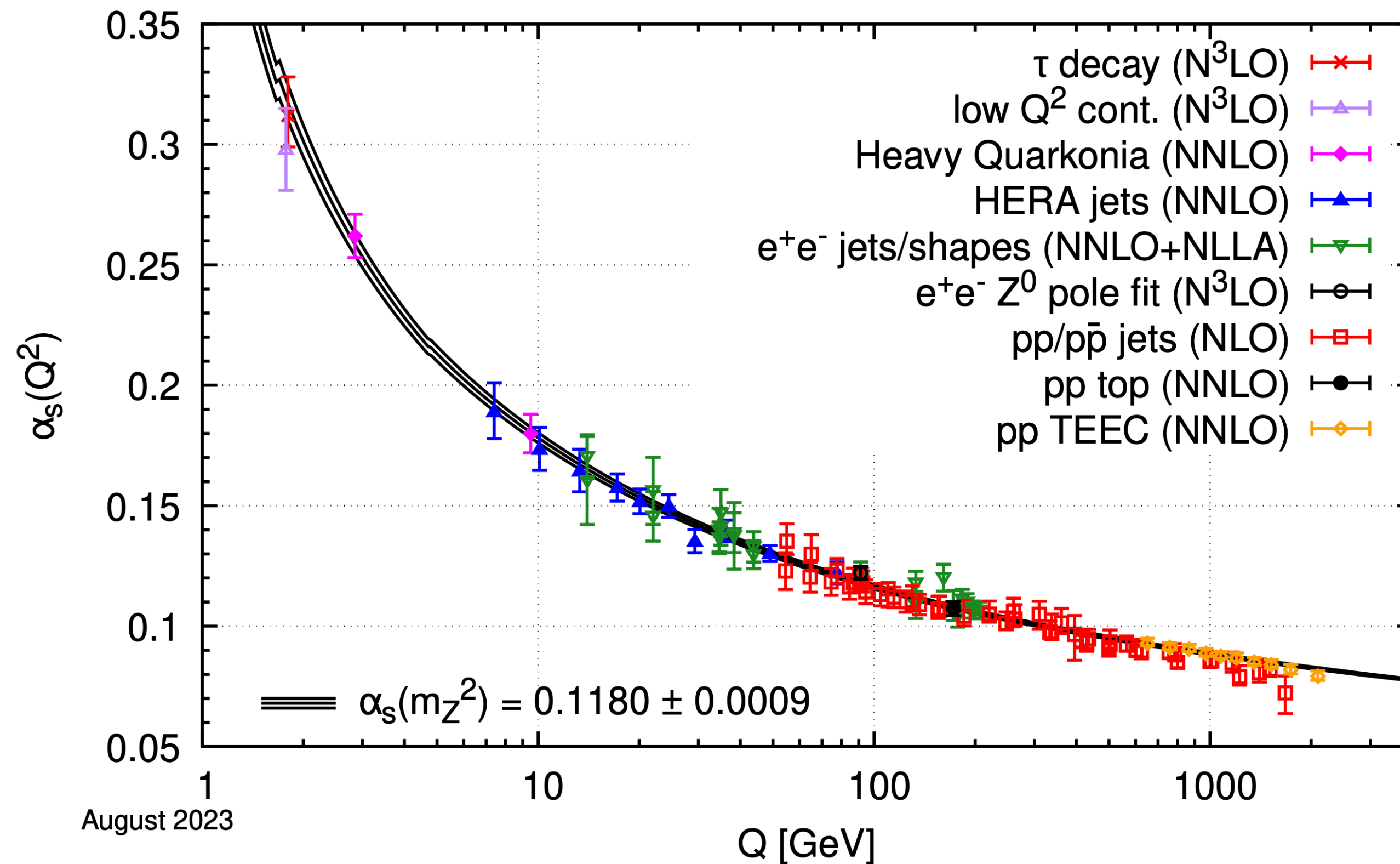
Matrix elements can be extracted e.g. from **leptonic** and **semileptonic** decays of hadrons

$$\underbrace{\frac{\Gamma[K \rightarrow l\nu_l(\gamma)]}{\Gamma[\pi \rightarrow l\nu_l(\gamma)]}}_{\text{experiments}} \propto \underbrace{\left| \frac{V_{us}}{V_{ud}} \right|^2}_{\text{QCD}} \underbrace{\left( \frac{f_K}{f_\pi} \right)^2}_{\text{QCD}}$$

$$\underbrace{\Gamma[K \rightarrow \pi l\nu_l(\gamma)]}_{\text{experiments}} \propto \underbrace{|V_{us}|^2}_{\text{QCD}} \underbrace{|f_+^{K\pi}(0)|^2}_{\text{QCD}}$$



# The strong coupling constant



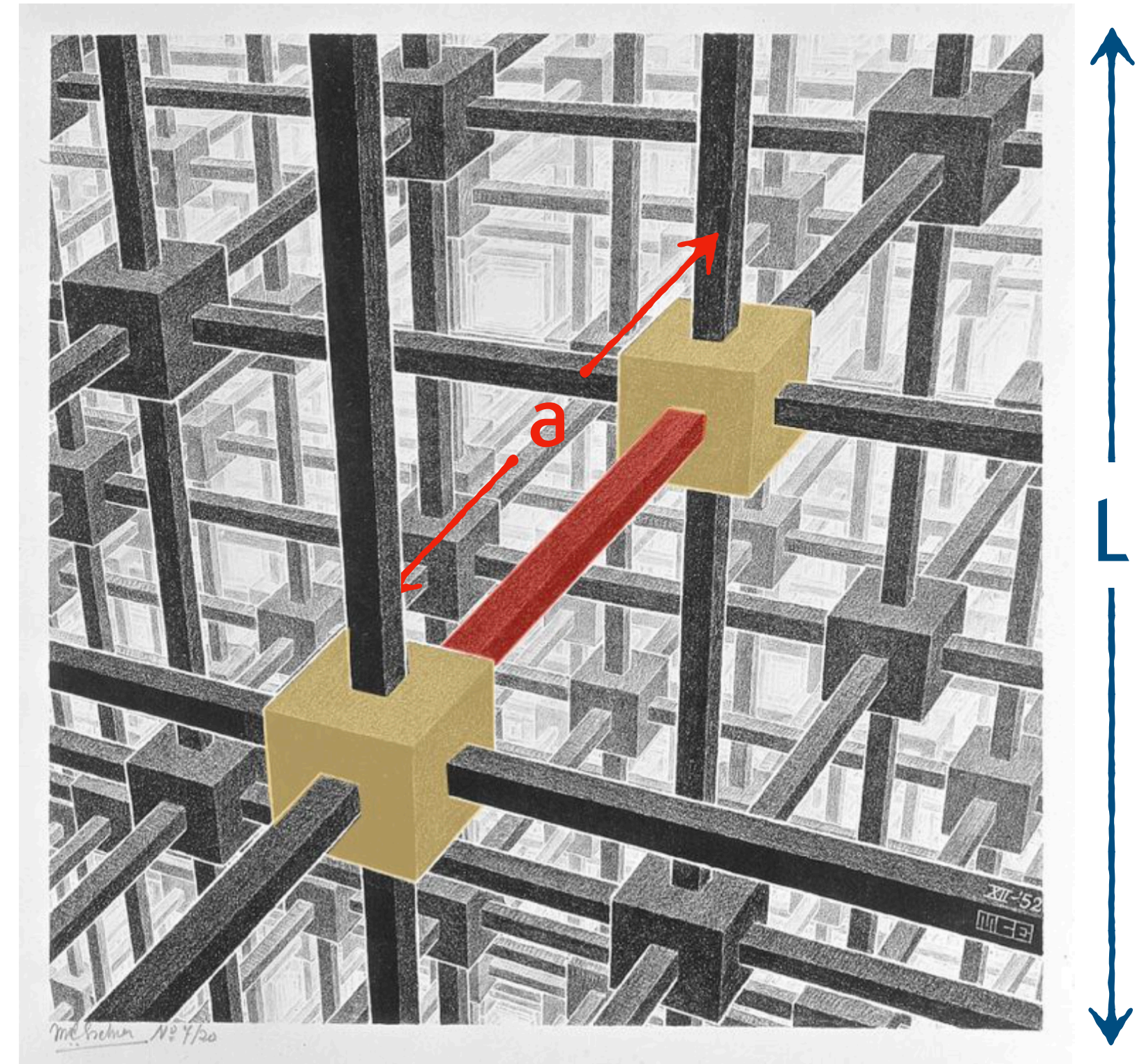
- The strong coupling constant  $\alpha_s(Q^2)$  runs with the energy  $Q$
- At high energies  $Q \sim m_Z$  the coupling is small:
  - ▶ perturbative expansion
  - ▶ quarks are asymptotically free
- At small energies  $Q \sim \Lambda_{\text{QCD}}$  the coupling is strong:
  - ▶ non-perturbative
  - ▶ quarks are confined

# Lattice QCD in a [small] nutshell

- QCD on a discrete and finite Euclidean space-time
- Based on Feynman path integrals

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}[U] \mathcal{D}[\psi, \bar{\psi}] \mathcal{O}[\psi, \bar{\psi}, U] e^{-S[\psi, \bar{\psi}, U]}$$

- Path integral solved using Monte Carlo methods
- Physical QCD results obtained, after renormalization, by taking the continuum & infinite-volume limit

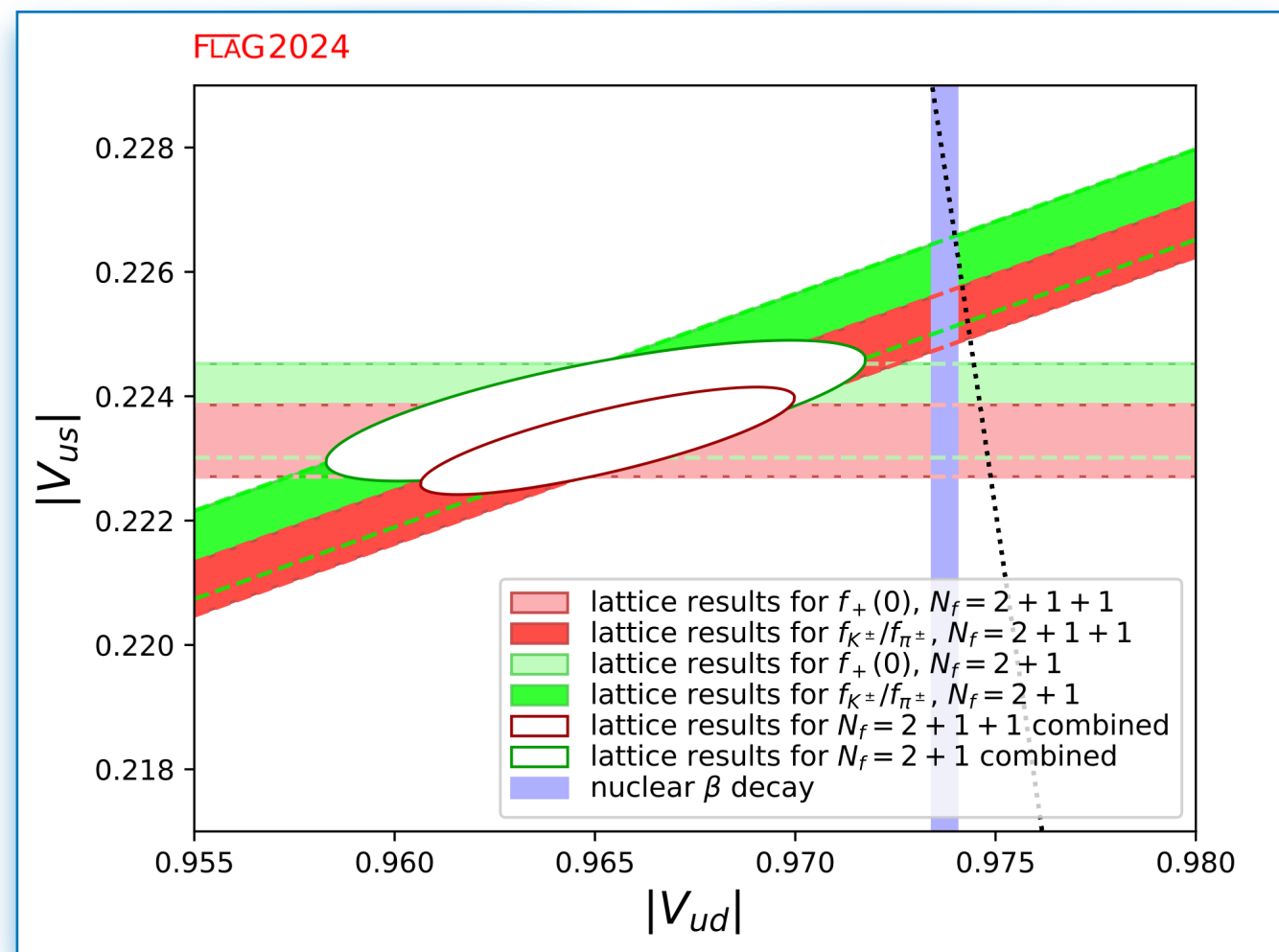


M.C. Escher, "Cubic space division" (1953)

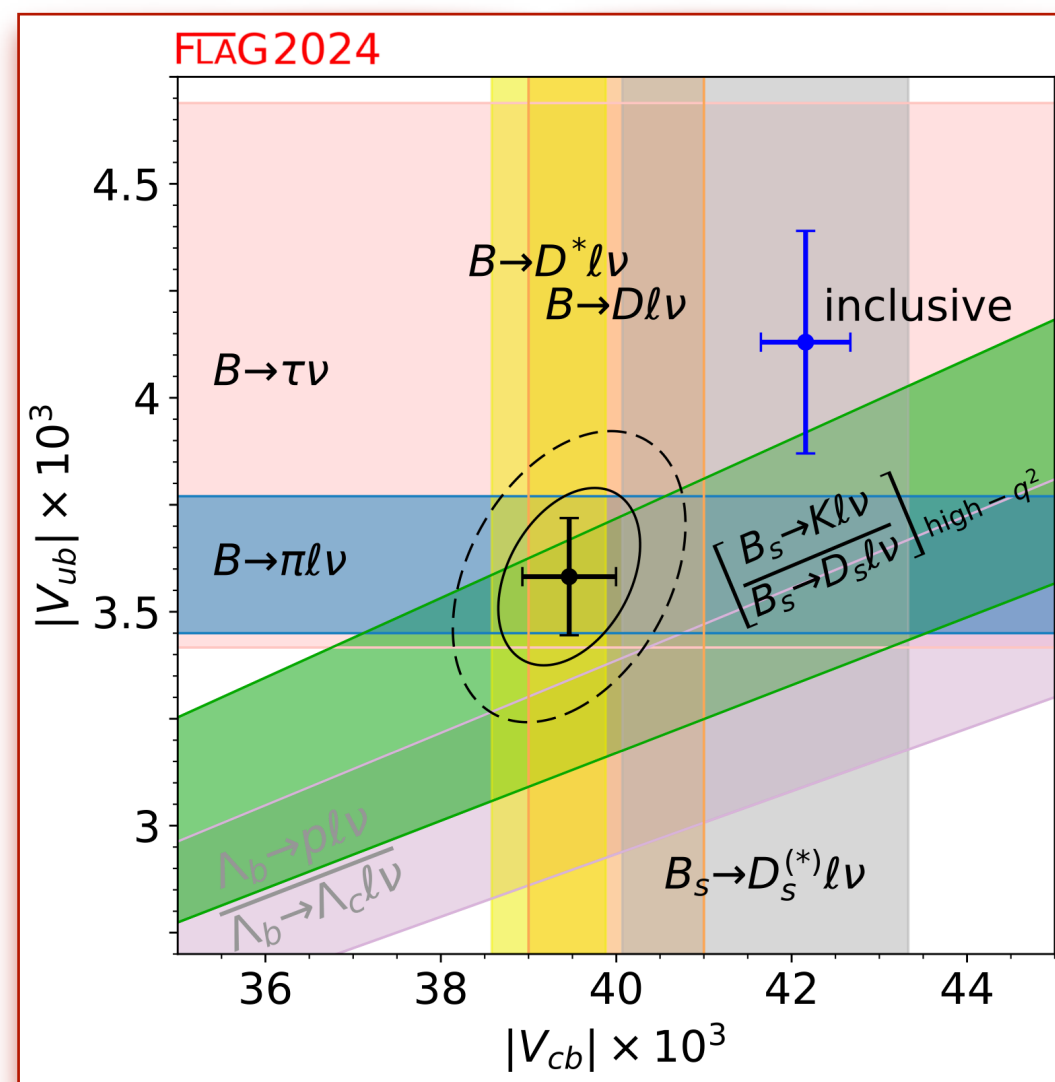
# Plan of the talk

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

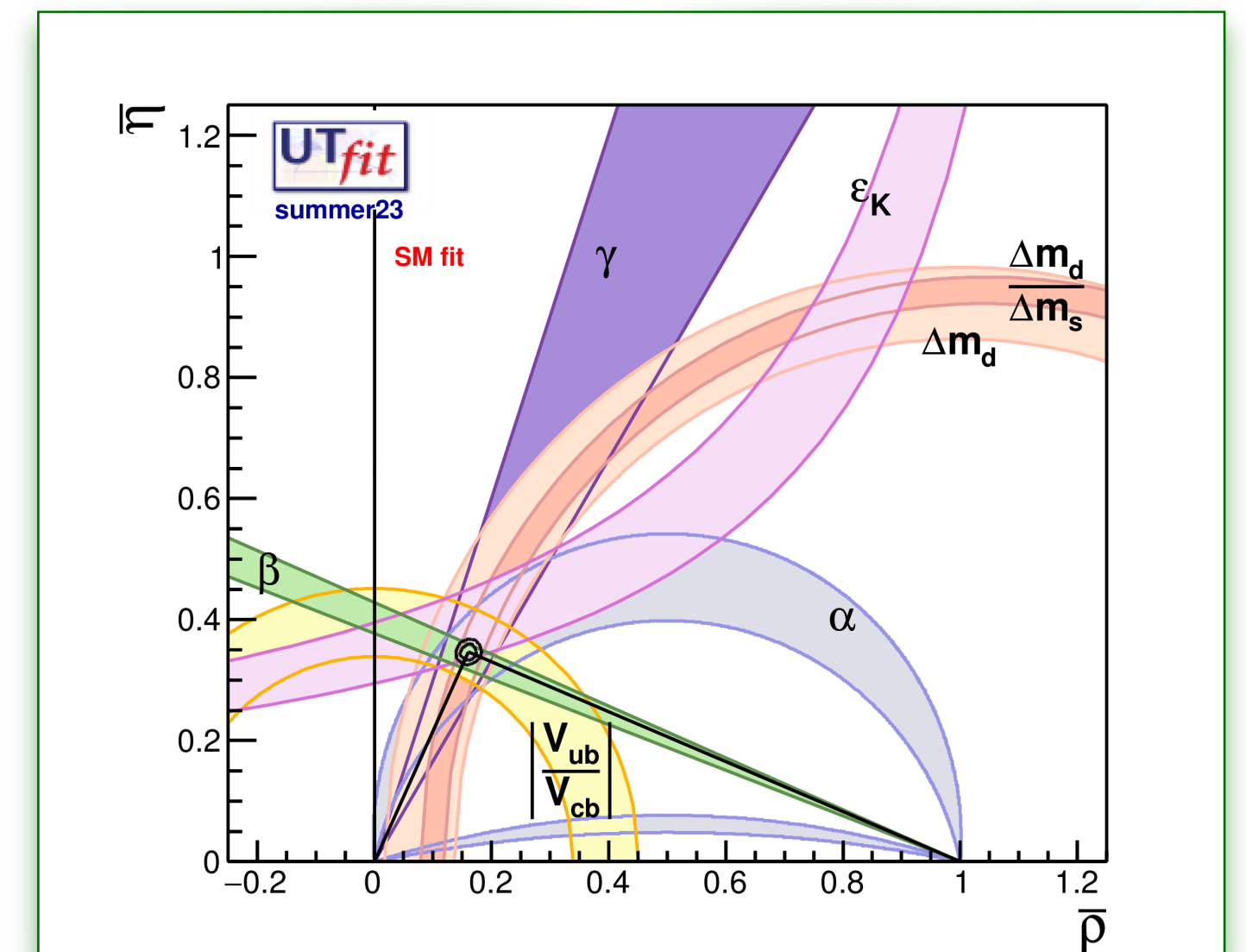
## Cabibbo anomaly



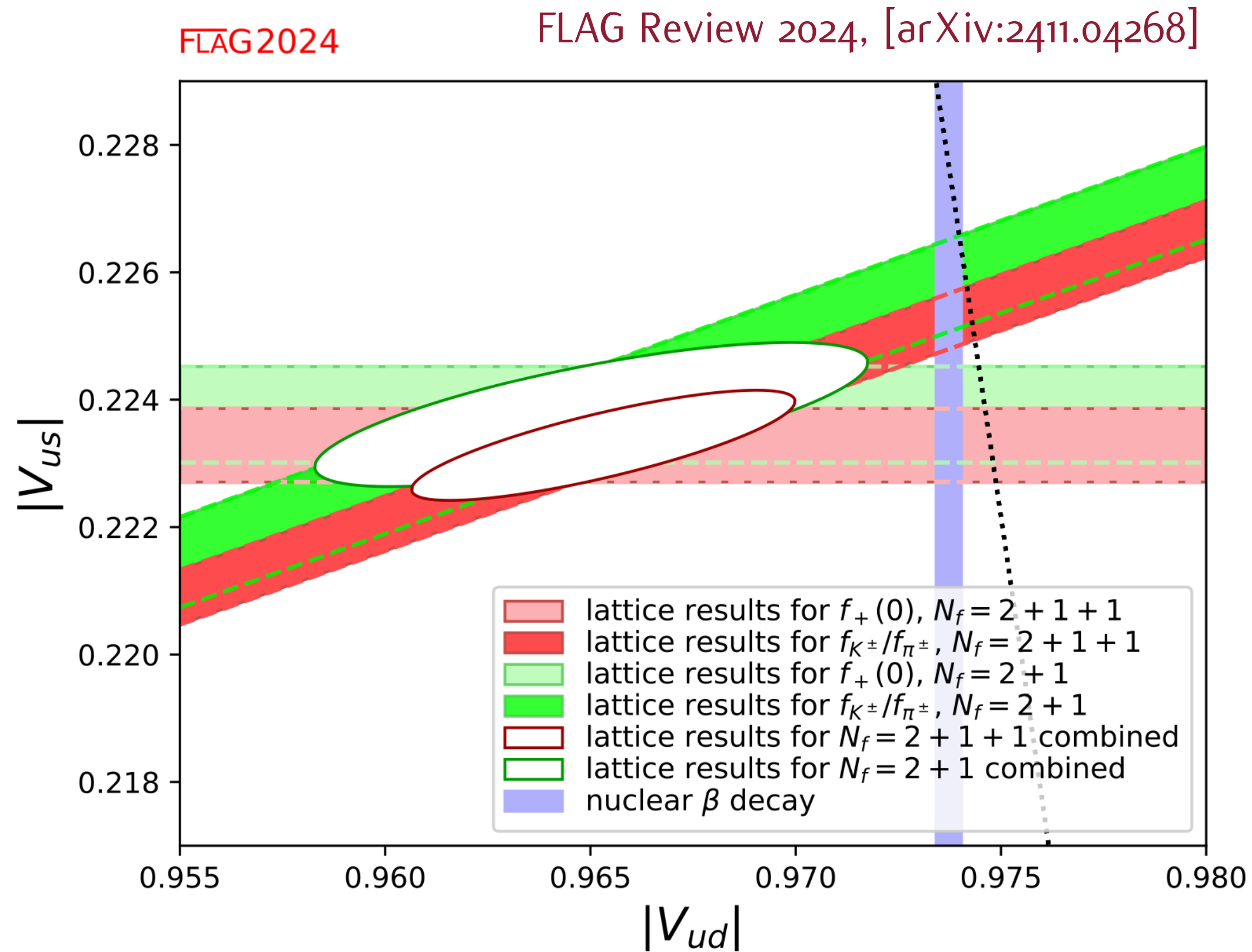
## $|V_{cb}|$ puzzle



## CPV in neutral mesons



# The Cabibbo anomaly



$$\frac{|V_{us}|}{|V_{ud}|} \frac{f_{K^\pm}}{f_{\pi^\pm}} = 0.27599(41)$$

M.Moulson, PoS CKM2016 (2017)

PDG, PTET 2022 (2022)

$$|V_{us}| |f_+^{K^0\pi^-}(0)| = 0.21654(41)$$

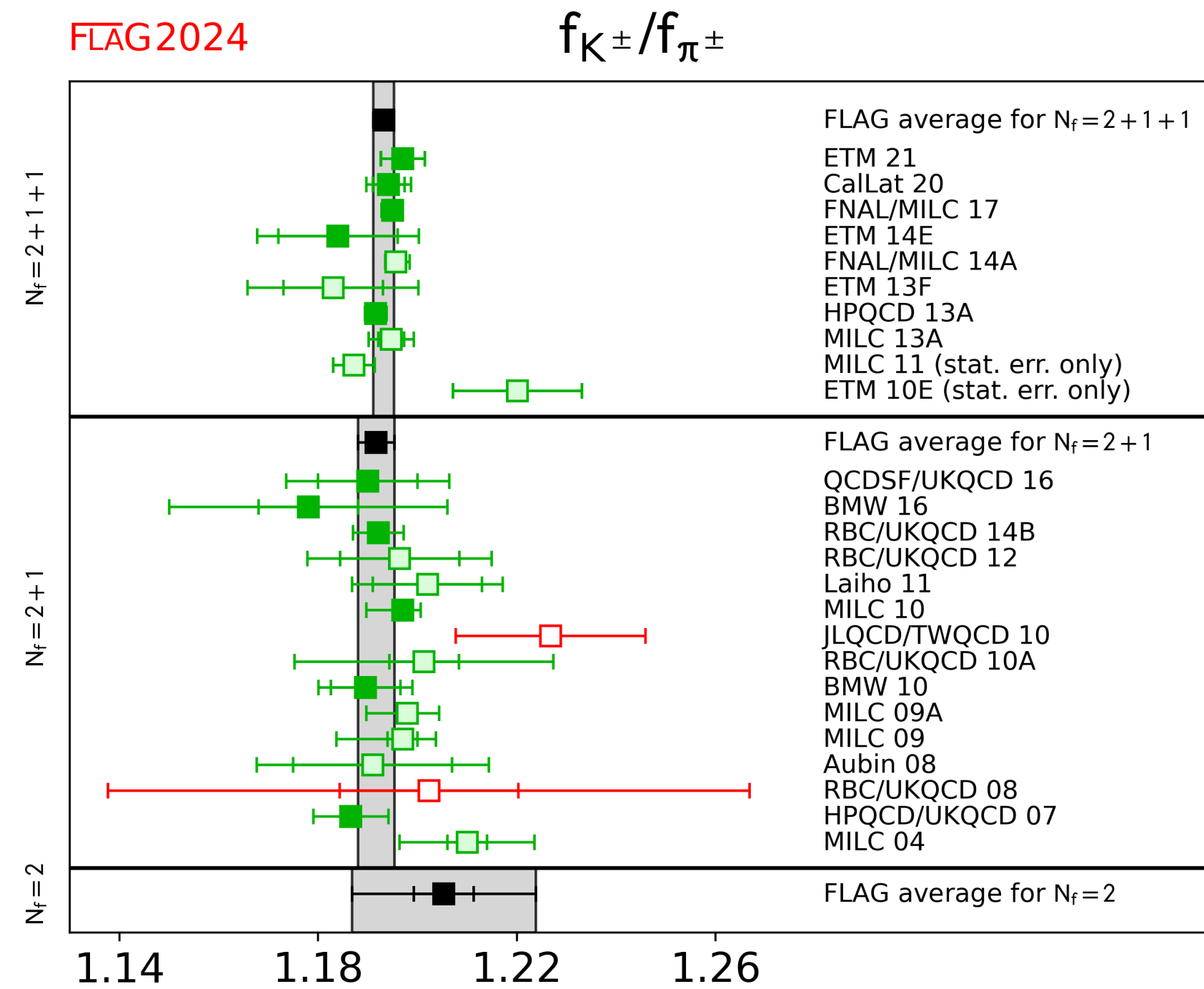
Different tensions in the  $V_{us}$ - $V_{ud}$  plane:

$$|V_u|^2_{\text{red oval}} - 1 = 2.8\sigma$$

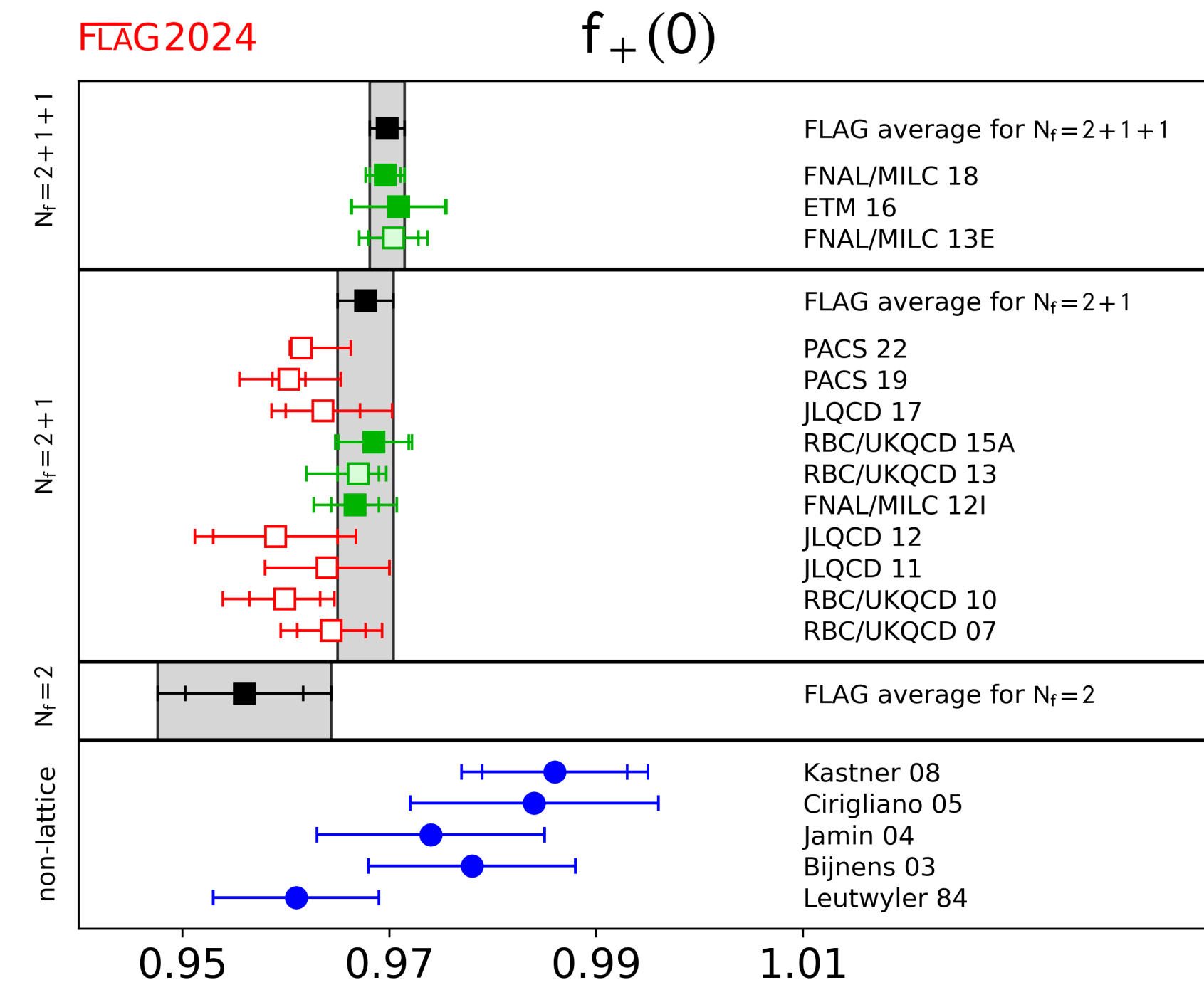
$$|V_u|^2_{\text{blue square}} - 1 = 3.1\sigma \quad |V_u|^2_{\text{blue square, red square}} - 1 = 1.7\sigma$$

Experimental and theoretical control of these quantities is of crucial importance to solve the issue

# Lattice QCD inputs



$$f_{K^\pm}/f_{\pi^\pm} = 1.1934(19)$$



$$f_+^{K^\pi}(0) = 0.9698(17)$$

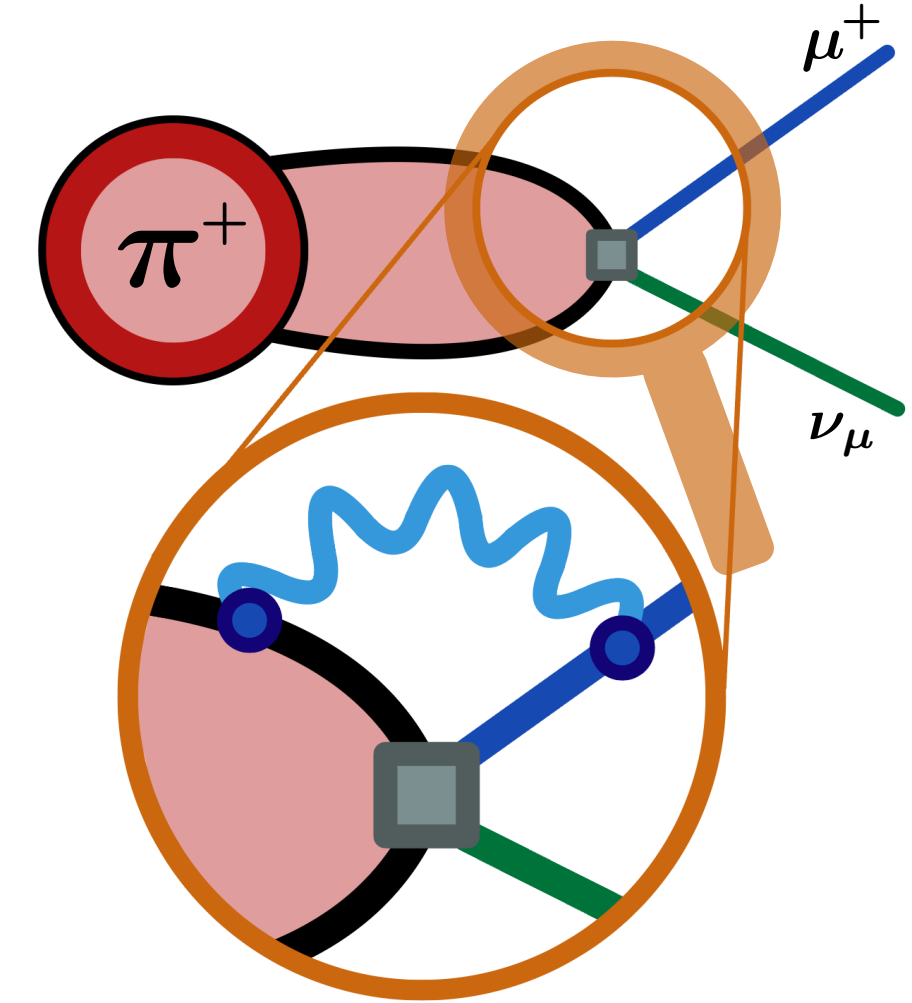
**FLAG2024** *Flavour Lattice Averaging Group*  $f_K/f_\pi$  and  $f_+^{K^\pi}(0)$  determined from lattice QCD with sub percent precision!



# QED and isospin-breaking effects

Current level of precision requires the inclusion of isospin-breaking corrections due to

- strong effects  $[m_u - m_d]_{\text{QCD}} \neq 0$
  - electromagnetic effects  $\alpha \neq 0$
- $\sim \mathcal{O}(1\%)$



$$\frac{\Gamma(K \rightarrow l\nu_l)}{\Gamma(\pi \rightarrow l\nu_l)} \propto \frac{|V_{us}|^2}{|V_{ud}|^2} \left(\frac{f_K}{f_\pi}\right)^2 (1 + \delta R_{K\pi}) \quad \Gamma(K \rightarrow \pi l\nu_l) \propto |V_{us}|^2 |f_+^{K\pi}(0)|^2 \mathcal{I}_{K\pi}^l (1 + \delta R_{K\pi}^l)$$

- ▶ results currently quoted in the PDG come from  $\chi$ PT
- ▶ fully non-perturbative (structure dependent) quantities
- ▶ first-principle lattice calculations are possible!

V.Cirigliano & H.Neufeld, PLB 700 (2011)

# Lattice QCD + QED

A conceptual challenge: how to define QED in a finite periodic box?

- ▶ need to circumvent Gauss' law: no charged states in a periodic box
- ▶ finite-volume effects can be sizeable and power-like
- ▶ logarithmic infrared divergences arise when studying decays

Problems well studied. Different lattice QED formulations proposed and used.

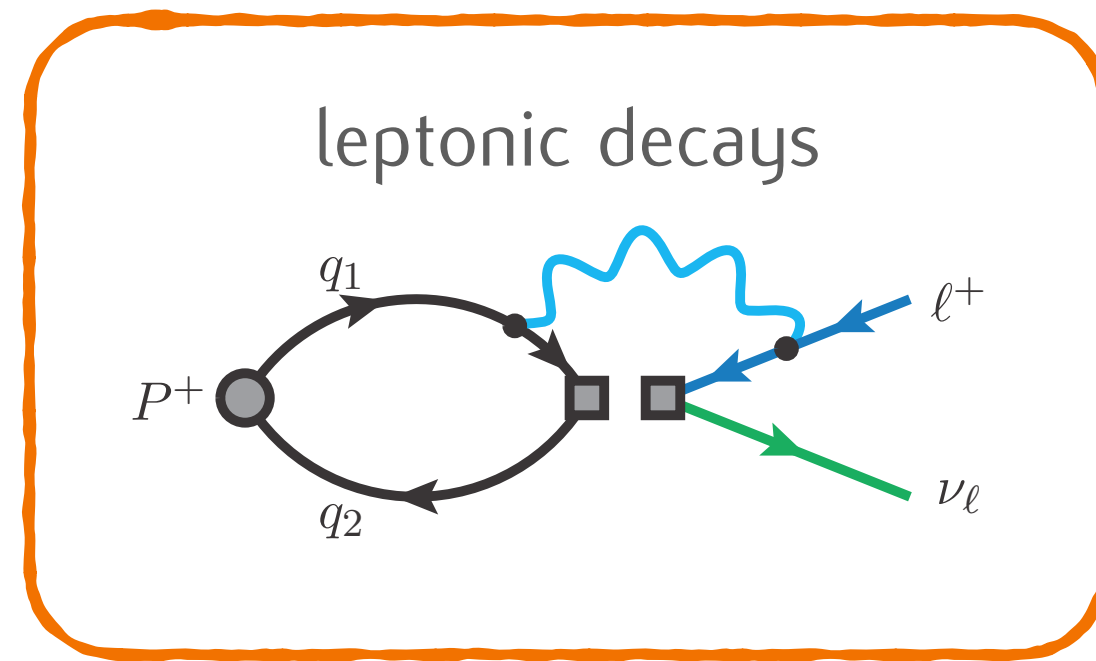
RM123 approach:

G.M.de Divitiis et al. [RM123], PRD 87 (2013)

$$\langle \mathcal{O} \rangle = \int \mathcal{D}\Phi \mathcal{O} e^{-S_{\text{iso}} - \Delta S} = \langle \mathcal{O} \rangle_{\text{iso}} + \langle \Delta S \mathcal{O} \rangle_{\text{iso}} + \dots$$

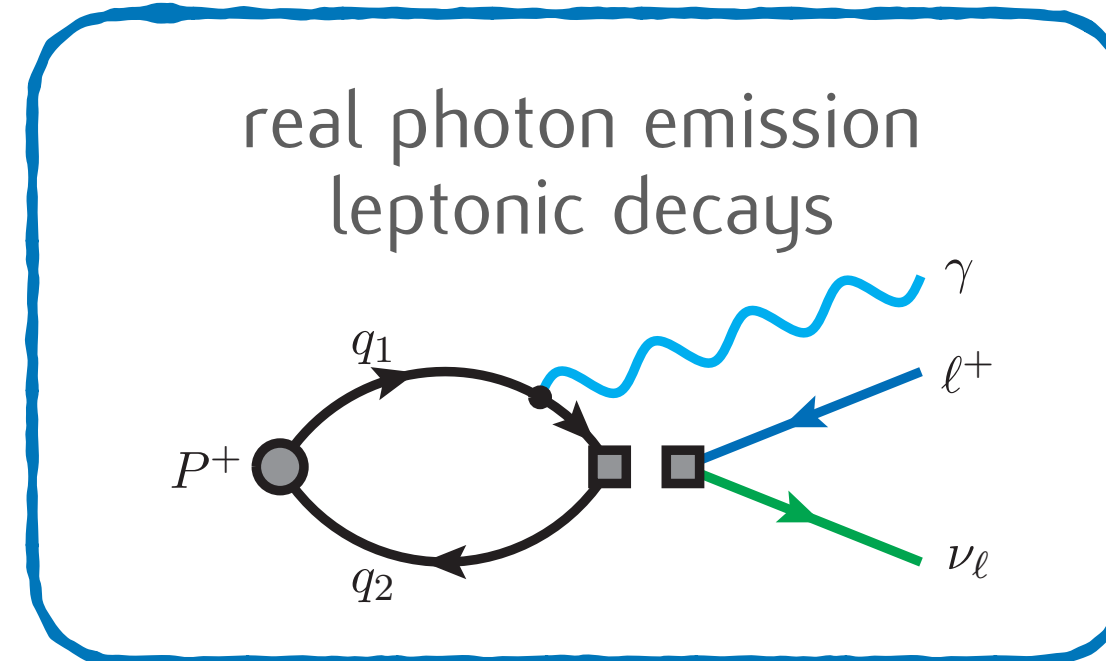
$$\text{“iso”} = \begin{cases} m_u = m_d \\ \alpha_{\text{em}} = 0 \end{cases}$$

# Weak decays — some recent works



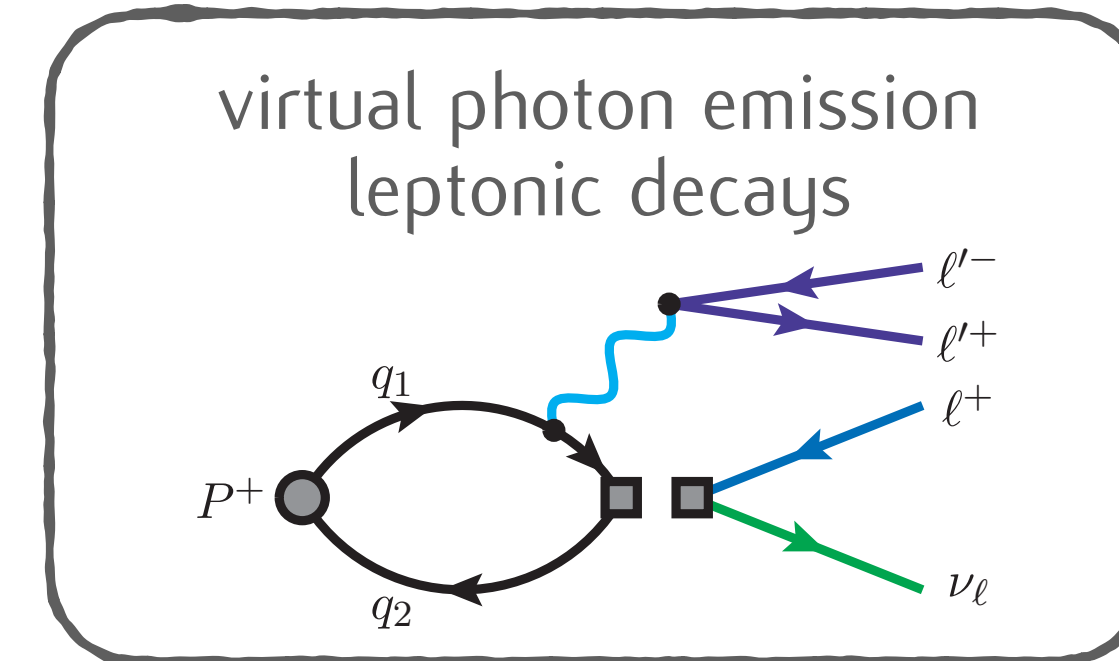
N.Carrasco et al., PRD 91 (2015)  
 V.Lubicz et al., PRD 95 (2017)  
 N.Tantalo et al., [1612.00199v2]  
 D.Giusti et al., PRL 120 (2018)  
 MDC et al., PRD 100 (2019)  
 MDC et al., PRD 105 (2022)  
 P.Boyle, MDC et al., JHEP 02 (2023)  
 N.Christ et al., [2304.08026]

R.Frezzotti et al., [2402.03262]



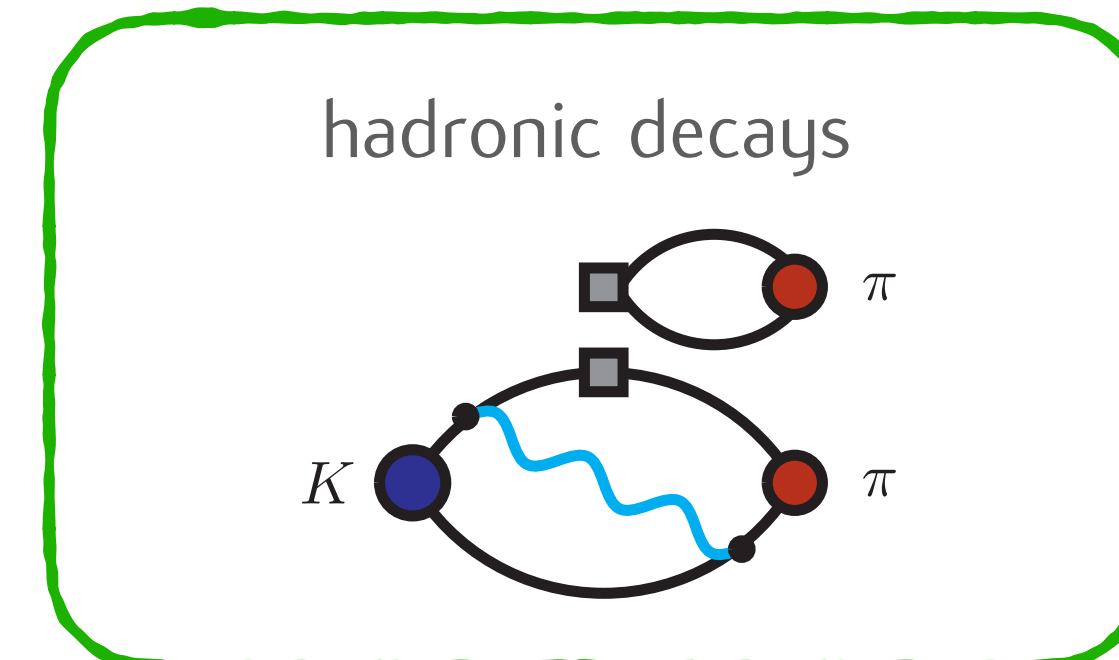
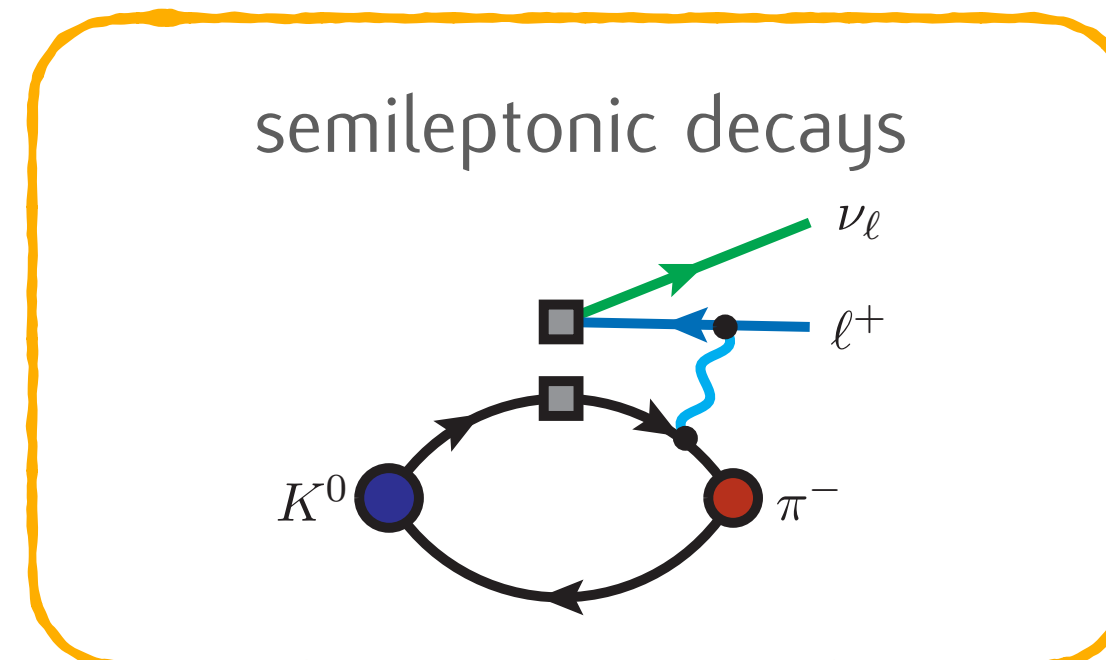
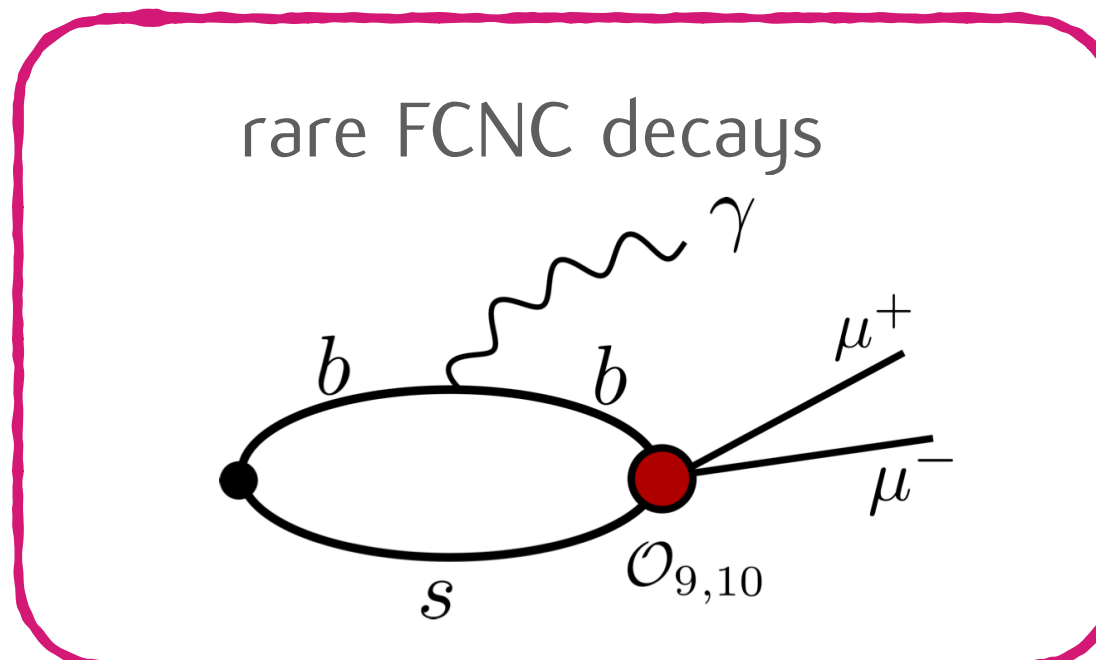
G.M.de Divitiis et al., [1908.10160]  
 C.Kane et al., [1907.00279 & 2110.13196]  
 R.Frezzotti et al., PRD 103 (2021)  
 A.Desiderio et al., PRD 102 (2021)  
 D.Giusti et al., [2302.01298]  
 R.Frezzotti et al., [2306.05904]

C.Sachrajda et al., [1910.07342]  
 N.Christ et al., PRD 108 (2023)  
 N.Christ et al., [2402.08915]



G.Gagliardi et al., Phys. Rev. D 105 (2022)  
 R.Frezzotti et al., [2306.07228]

R.Abbott et al., PRD 102 (2020)  
 Z.Bai et al., PRL 115 (2015)  
 N.Christ et al., PRD 106 (2022)  
 N.Christ & X.Feng, EPJ Web Conf. 175 (2018)  
 Y.Cai & Z.Davoudi, [1812.11015]



# Leptonic decay rate at $\mathcal{O}(\alpha)$

F.Bloch & A.Nordsieck, PR 52 (1937)  
 N.Carrasco et al., PRD 91 (2015)  
 V.Lubicz et al., PRD 95 (2017)  
 N.Tantalo et al. [1612.00199v2]  
 MDC et al., PRD 105 (2022)

The RM123+Soton recipe (2015)

$$\Gamma(P_{\ell 2}) = \lim_{\Lambda_{\text{IR}} \rightarrow 0} \left\{ \text{IR finite} \left[ \text{Diagram 1} - \text{Diagram 2} \right] \right\} + \lim_{\Lambda_{\text{IR}} \rightarrow 0} \left\{ \text{Diagram 3} + \text{Diagram 4} \right\} \\
 + \lim_{\Lambda_{\text{IR}} \rightarrow 0} \left\{ \text{Diagram 5} - \text{Diagram 6} \right\}$$

The diagrams are Feynman diagrams for leptonic decay at  $\mathcal{O}(\alpha)$ . Each diagram features a yellow circle labeled  $\mathcal{P}$  on the left. A blue wavy line (photon) is emitted from the vertex. A green arrow (lepton) is shown on the right. Diagram 1 is a loop diagram with a black line and a blue line. Diagram 2 is a tree-level diagram with a black line. Diagram 3 is a tree-level diagram with a blue line. Diagram 4 is a tree-level diagram with a blue line. Diagram 5 is a loop diagram with a black line and a blue line. Diagram 6 is a tree-level diagram with a blue line. The labels "IR finite" are placed below the first and fifth diagrams.

# Leptonic decay rate at $\mathcal{O}(\alpha)$

F.Bloch & A.Nordsieck, PR 52 (1937)  
 N.Carrasco et al., PRD 91 (2015)  
 V.Lubicz et al., PRD 95 (2017)  
 N.Tantalo et al. [1612.00199v2]  
 MDC et al., PRD 105 (2022)

The RM123+Soton recipe (2015)

$$\Gamma(P_{\ell 2}) = \lim_{L \rightarrow \infty} \left\{ \text{on the lattice} \right\} + \lim_{m_\gamma \rightarrow 0} \left\{ \text{in perturbation theory} \right\}$$

$$+ \lim_{L \rightarrow \infty} \left\{ \text{on the lattice} \right\}$$

enough for  $K_{\mu 2}$  and  $\pi_{\mu 2}$

D. Giusti et al., PRL 120 (2018)  
 MDC et al., PRD 100 (2019)  
 P.Boyle, MDC et al., JHEP 02 (2023)

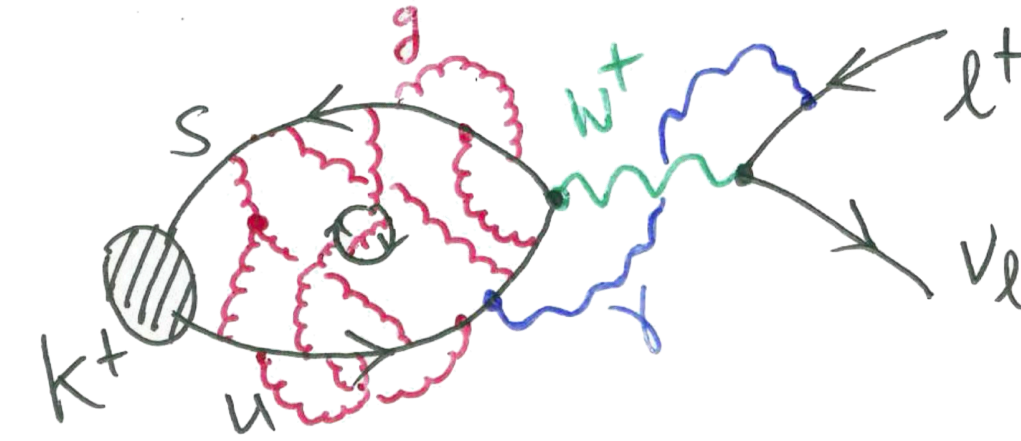
relevant for  $K_{e 2}$  and  $\pi_{e 2}$   
 & decays of heavier mesons

G.M.de Divitiis et al., [1908.10160]  
 R.Frezzotti et al., PRD 103 (2021)  
 A.Desiderio et al., PRD 102 (2021)

C.Kane et al., [1907.00279 & 2110.13196]  
 D.Giusti et al., [2302.01298]  
 R.Frezzotti et al., [2306.05904]

# Results for $\delta R_{K\pi}$

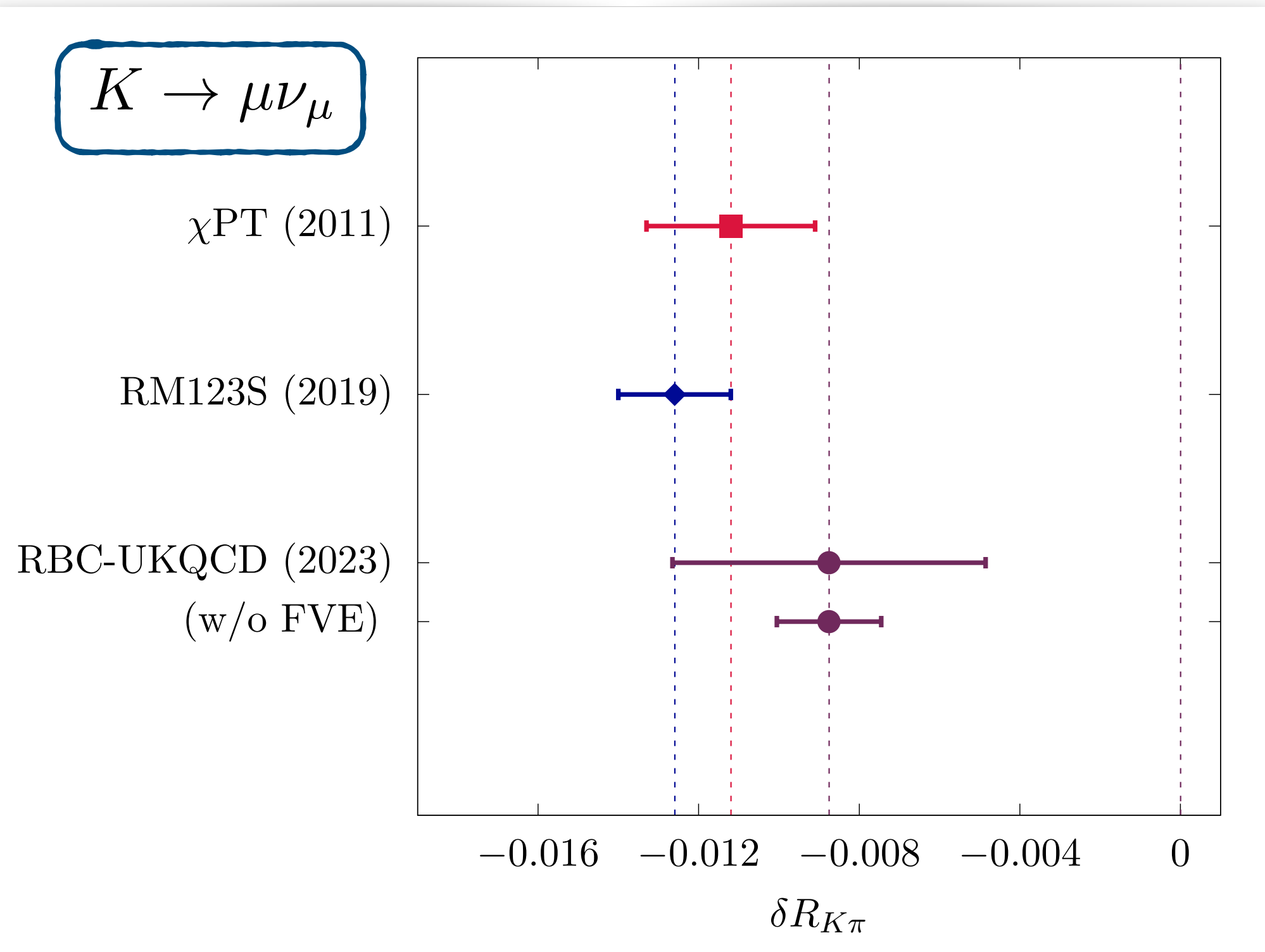
- $\delta R_{K\pi} = -0.0112 (21)$
- ◆  $\delta R_{K\pi} = -0.0126 (14)$
- $\delta R_{K\pi} = -0.0086 (13)(39)_{\text{vol.}}$



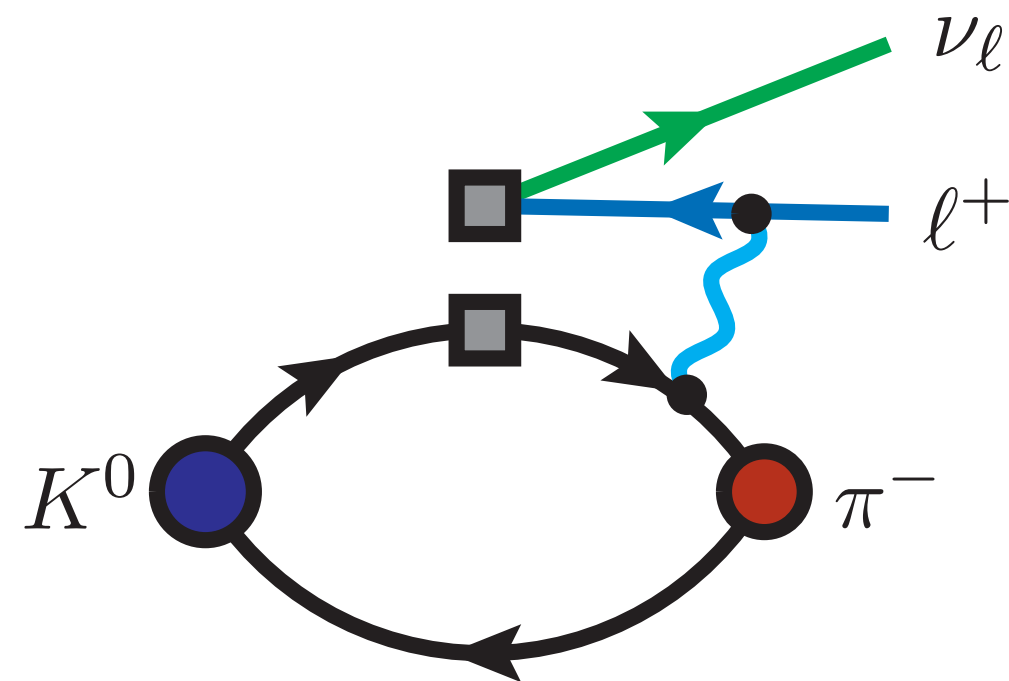
V.Cirigliano et al., PLB 700 (2011)  
 MDC et al., PRD 100 (2019)  
 P.Boyle, MDC et al., JHEP 02 (2023)

$$\frac{\Gamma(K \rightarrow l\nu_l)}{\Gamma(\pi \rightarrow l\nu_l)} \propto \frac{|V_{us}|^2}{|V_{ud}|^2} \left(\frac{f_K}{f_\pi}\right)^2 (1 + \delta R_{K\pi})$$

- Good evidence that  $\delta R_{K\pi}$  can be computed from first principles non-perturbatively on the lattice!
- RBC-UKQCD error dominated by a large systematic uncertainty related to finite-volume effects (!)  
 Work in progress to improve the result.
- Errors on  $|V_{us}|/|V_{ud}|$  from theoretical inputs could become comparable with those from experiments



# Semileptonic kaon decays



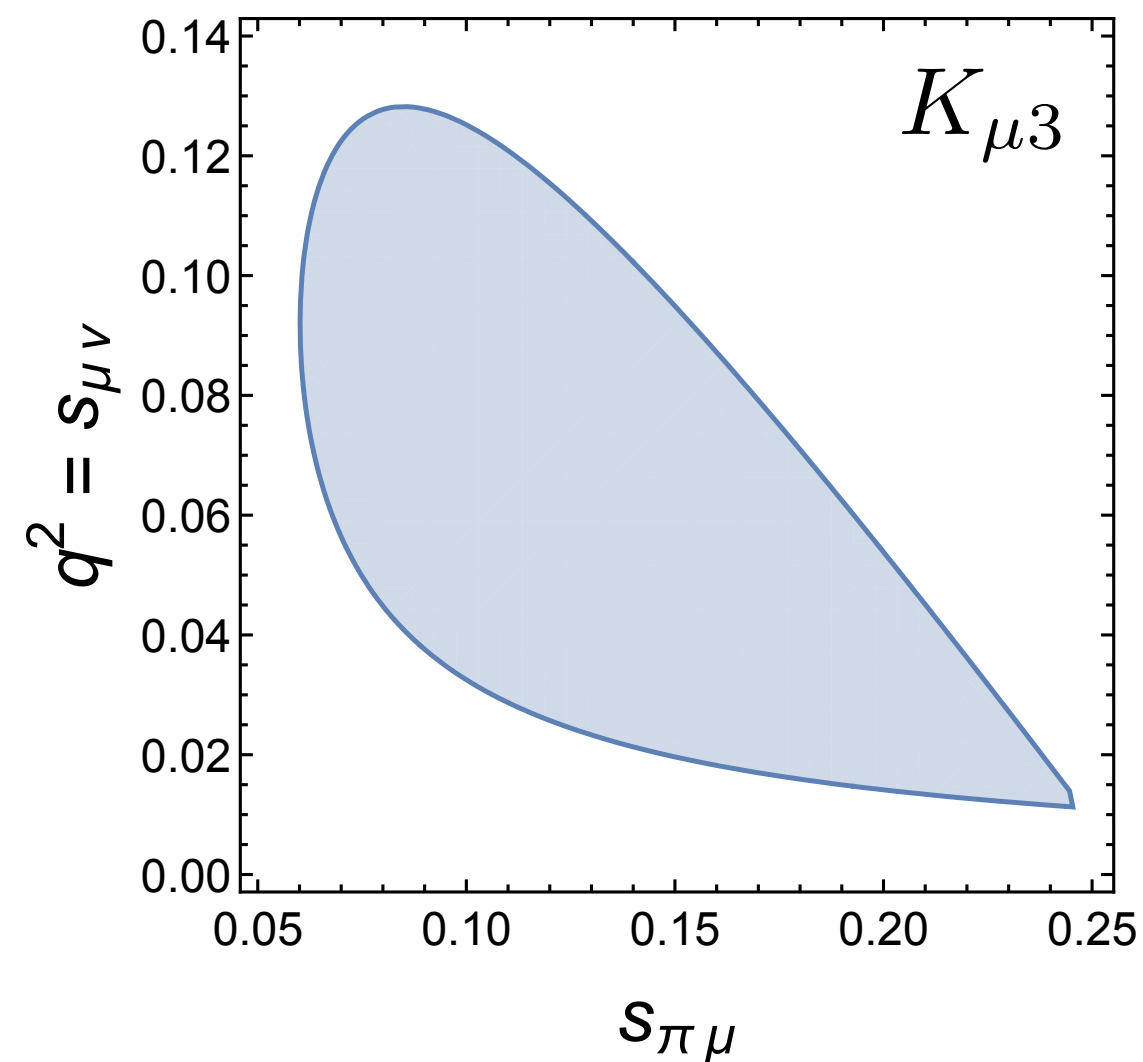
To go beyond current precision, we need to include **isospin-breaking effects** computed from first principles lattice QCD+QED.

Significant additional **difficulties** compared to leptonic decays:

- integration over **three-body phase-space**
- problems of **analytical continuation** from Euclidean to Minkowski when intermediate on shell  $\pi\text{-}\ell$  states are lighter than external ones

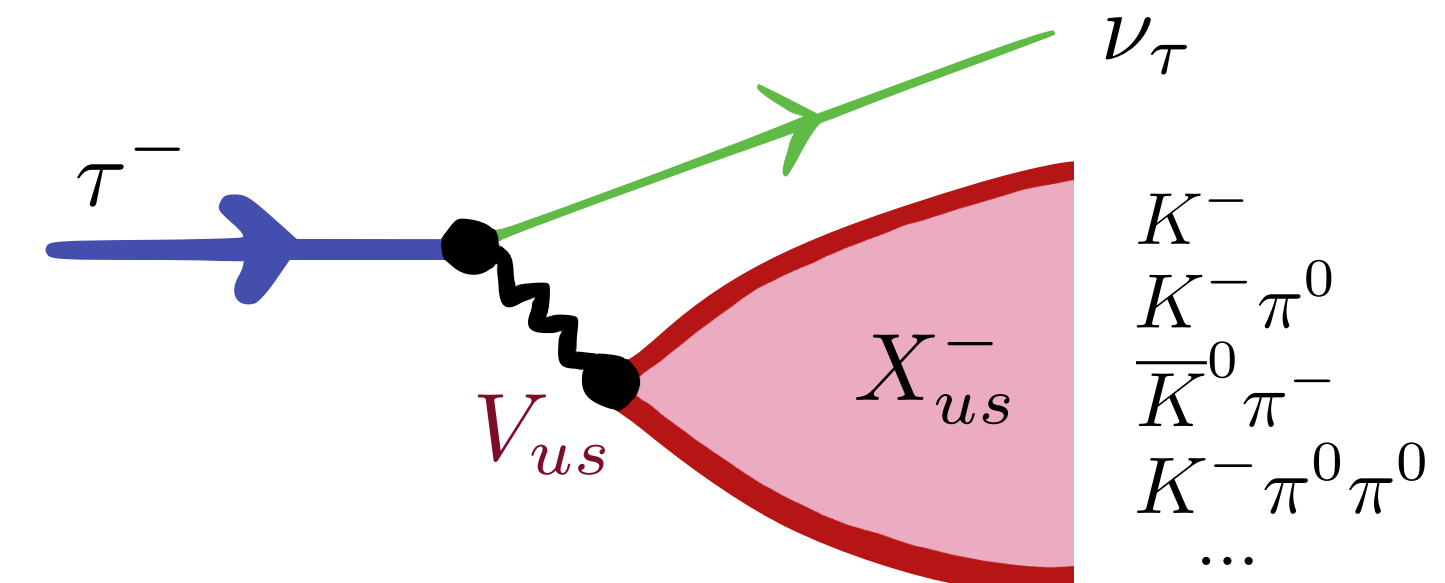
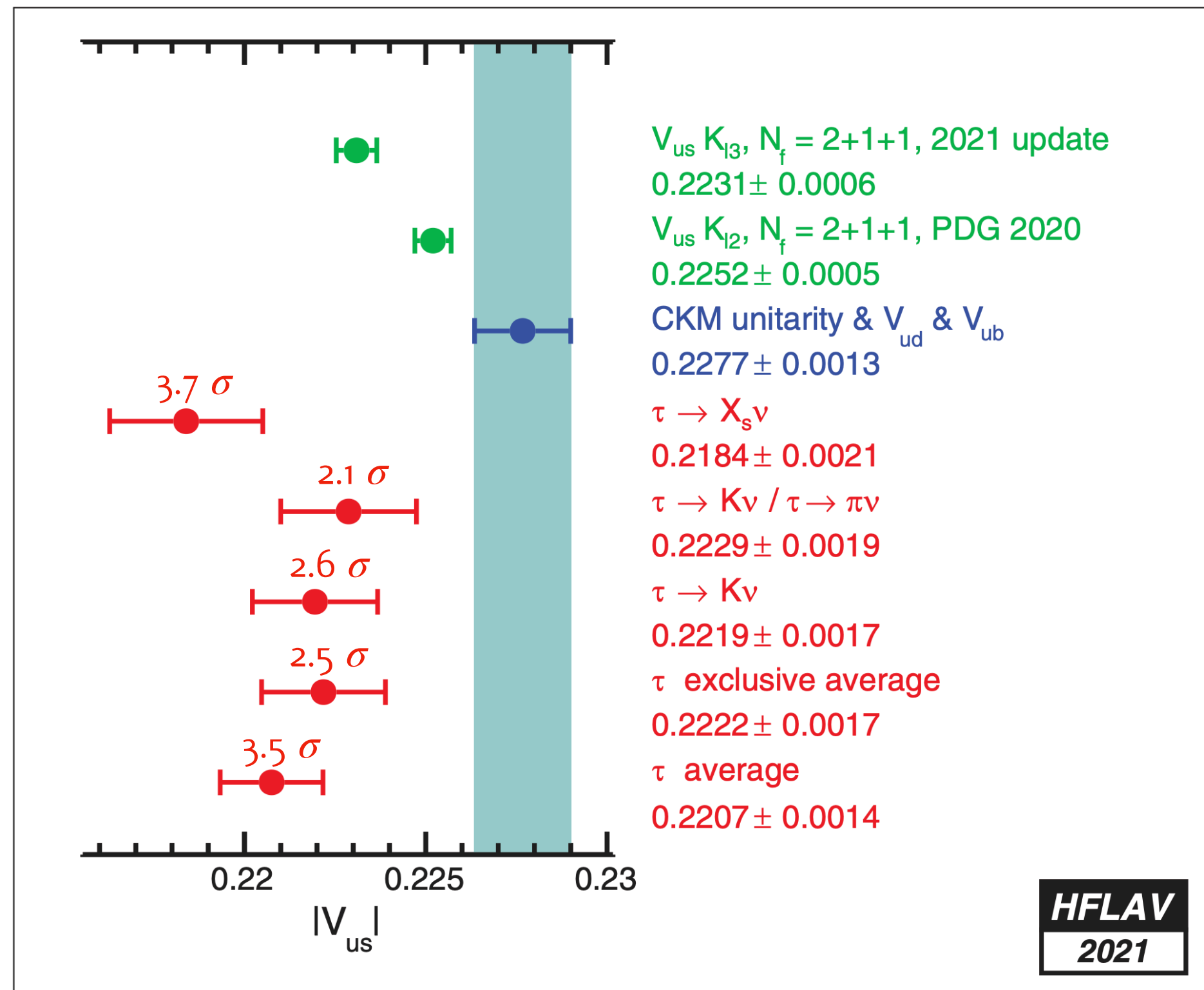
Proper **finite-volume QED formalism** is still missing, but solutions are under study by different groups.

Recent  $\text{QED}_\infty$  proposal: [N.Christ et al., PRD 108 \(2023\)](#) & [PoS LATTICE2023 \(2024\) 266](#)



# Inclusive hadronic $\tau$ decays

Alternative determinations of  $|V_{us}|$  can be obtained from inclusive hadronic  $\tau$  decays



- And yet another puzzle: lower value of  $|V_{us}|_{\tau\text{-incl.}}$
- Inclusive  $\tau \rightarrow X_{us} \nu_\tau$  result in HFLAV plot obtained using truncated operator product expansion (OPE)
- Exclusive channels give results larger than  $|V_{us}|_{\tau\text{-incl.}}$  but smaller than that obtained imposing CKM unitarity



# Inclusive hadronic $\tau$ decays

A.Evangelista et al. (ETMC), PRD 108 (2023)

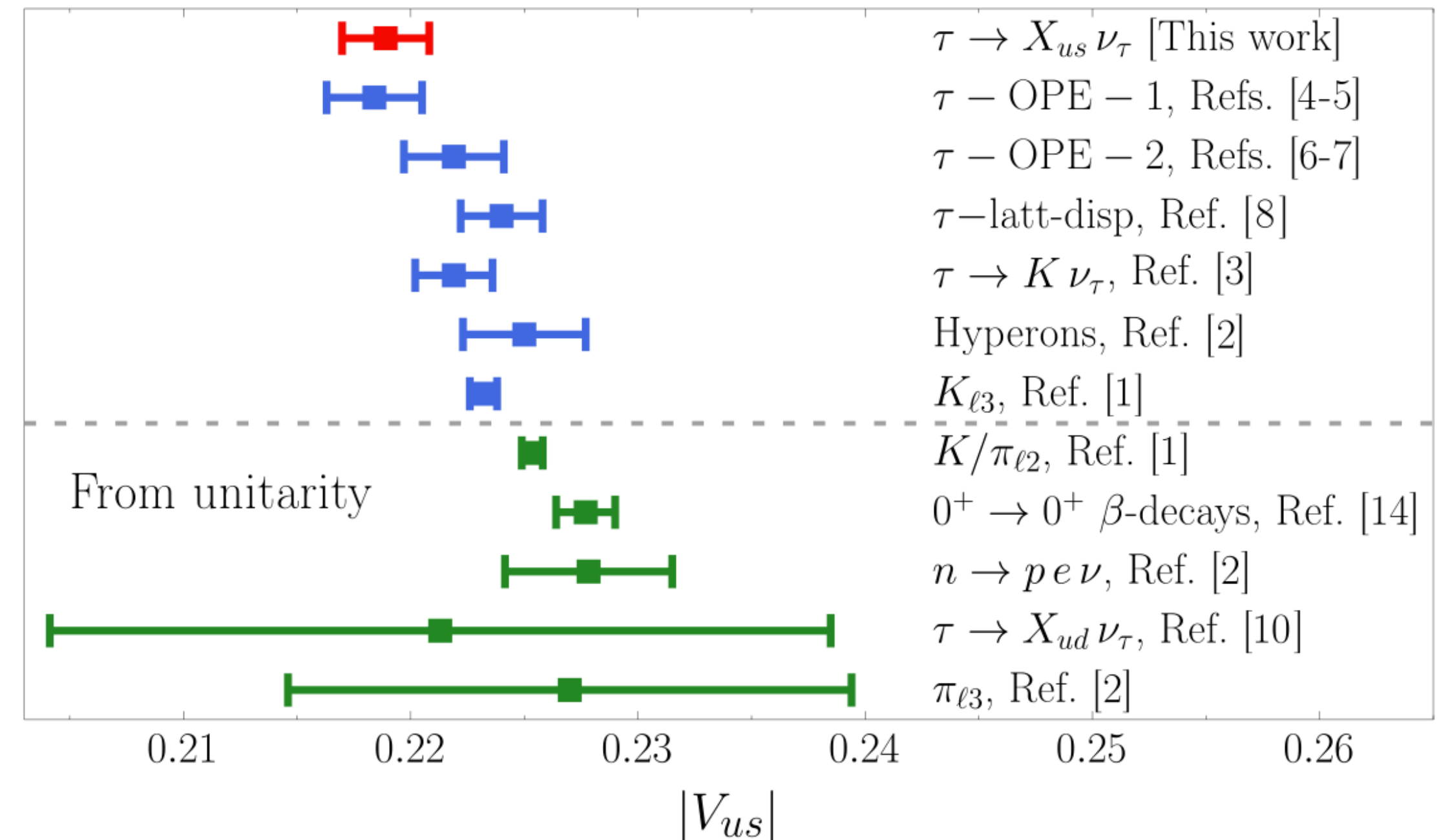
C.Alexandrou et al. (ETMC), PRL 132 (2024)

Recent calculation obtains inclusive decay rate using smeared spectral densities reconstructed from finite-volume Euclidean lattice correlators M.Hansen, A.Lupo & N.Tantalo, PRD 99 (2019)

$$\rho(\omega) = \langle \tau^- | H_W^{us} (2\pi) \delta(\mathbb{H} - \omega) H_W^{us} | \tau^- \rangle$$

$$\begin{aligned} \hat{\rho}_L(E, \epsilon) &= \int_0^\infty \frac{d\omega}{2\pi} \Delta_\epsilon(E, \omega) \rho_L(\omega) \\ &= \sum_{t=0}^T g_t(E, \epsilon) C_L(t) \end{aligned}$$

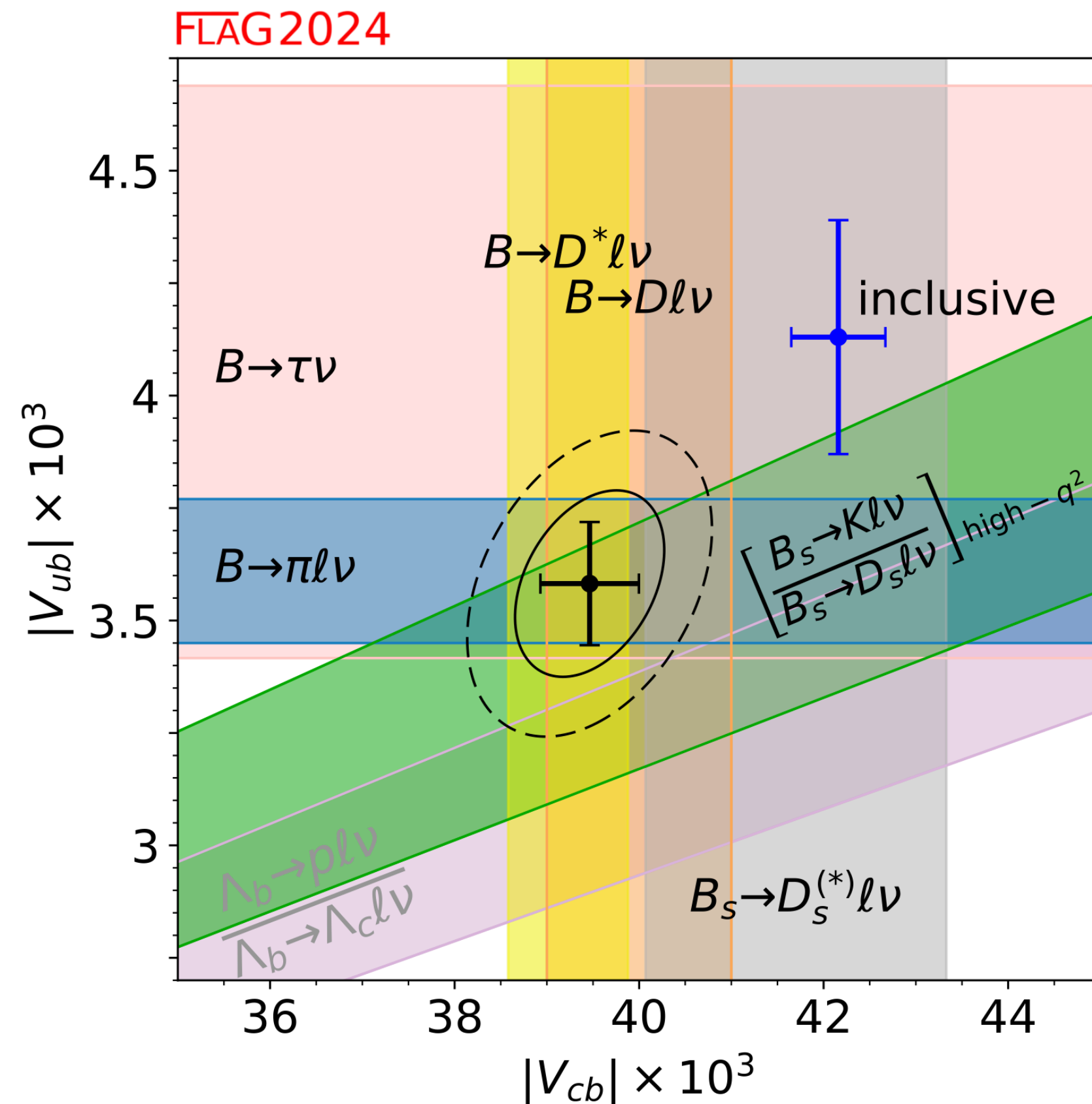
$$\Rightarrow \Gamma(\tau \rightarrow X_{us} \nu_\tau) = \lim_{\epsilon \rightarrow 0} \lim_{L \rightarrow \infty} \frac{\hat{\rho}_L(m_\tau, \epsilon)}{2m_\tau}$$



> **Next steps:** inclusion of QED and strong isospin-breaking effects

Intense activity on inverse problems in lattice QCD recently discussed at workshop Lattice@CERN 2024

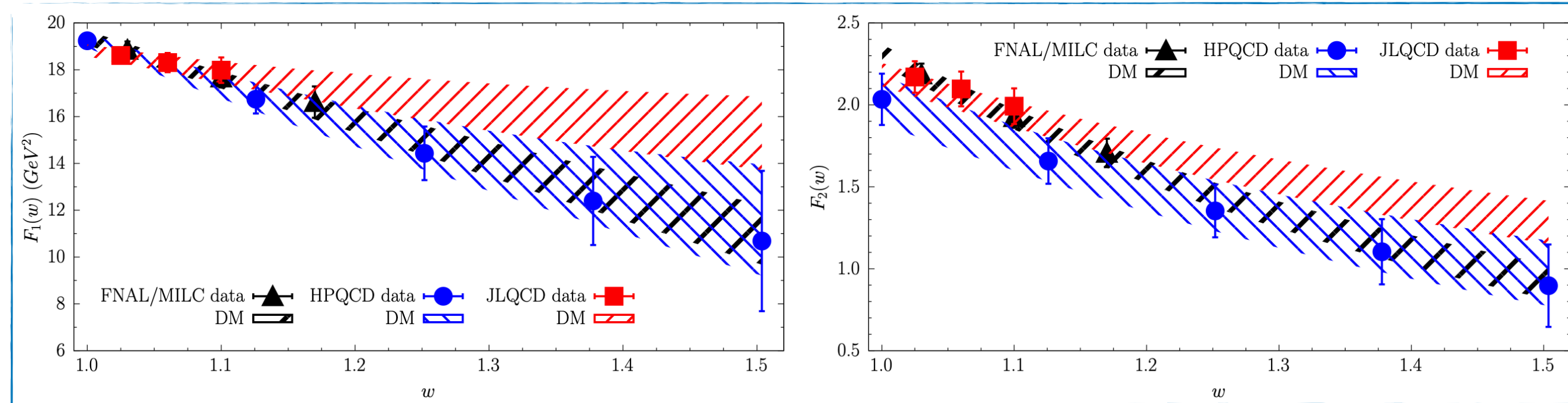
# The $|V_{cb}|$ puzzle



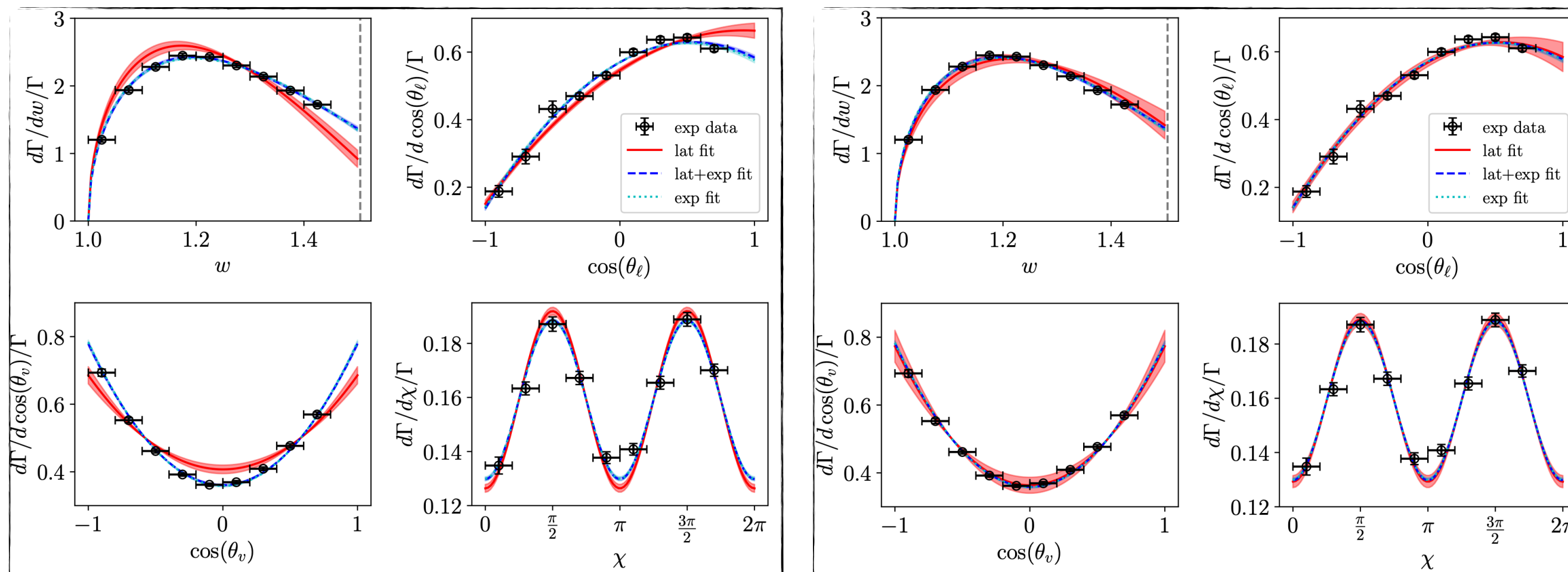
- Long-standing tension between inclusive and exclusive determinations of  $|V_{cb}|$
- Most precise experimental results come from  $B \rightarrow D^* \ell \nu$  decay (Belle and Belle II)
- Form factors are computed in lattice QCD by 3 collaborations: FNAL/MILC, HPQCD and JLQCD
- There are some tensions in the form-factor shapes, which are currently under investigation

# Comparison of lattice results

G.Martinelli et al., EPJC 84 (2024) 4

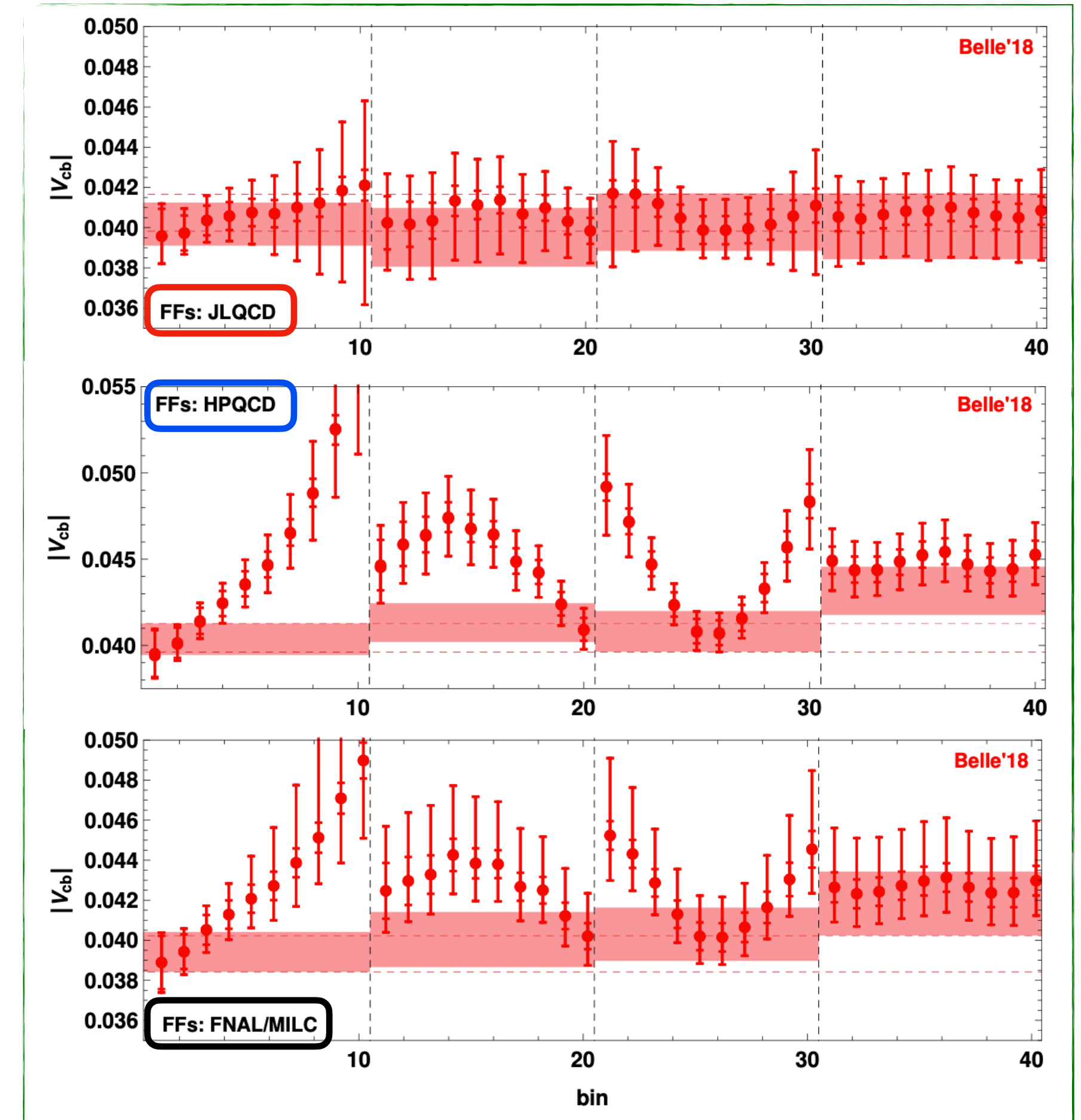


M.Bordone & A.Jüttner, arXiv:2406.10074



FNAL/MILC & HPQCD

JLQCD



M.Jung @ Siegen 2024

# Inclusive decays on the lattice

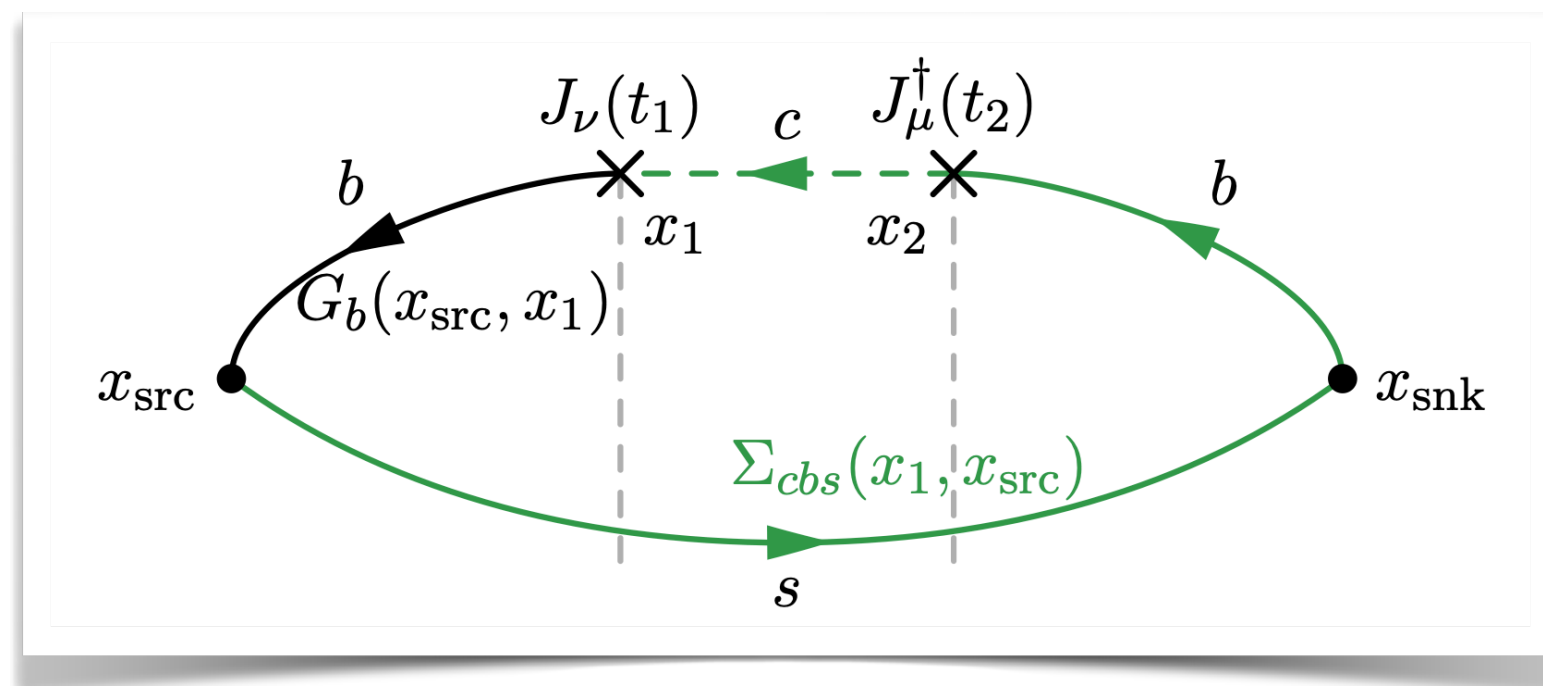
- Interesting proposal by [P.Gambino & S.Hashimoto PRL 125 \(2020\)](#) for lattice calculations of **inclusive semileptonic decays**: alternative way to get independent determination of  $V_{cb}$ .

- Use spectral reconstruction techniques to evaluate

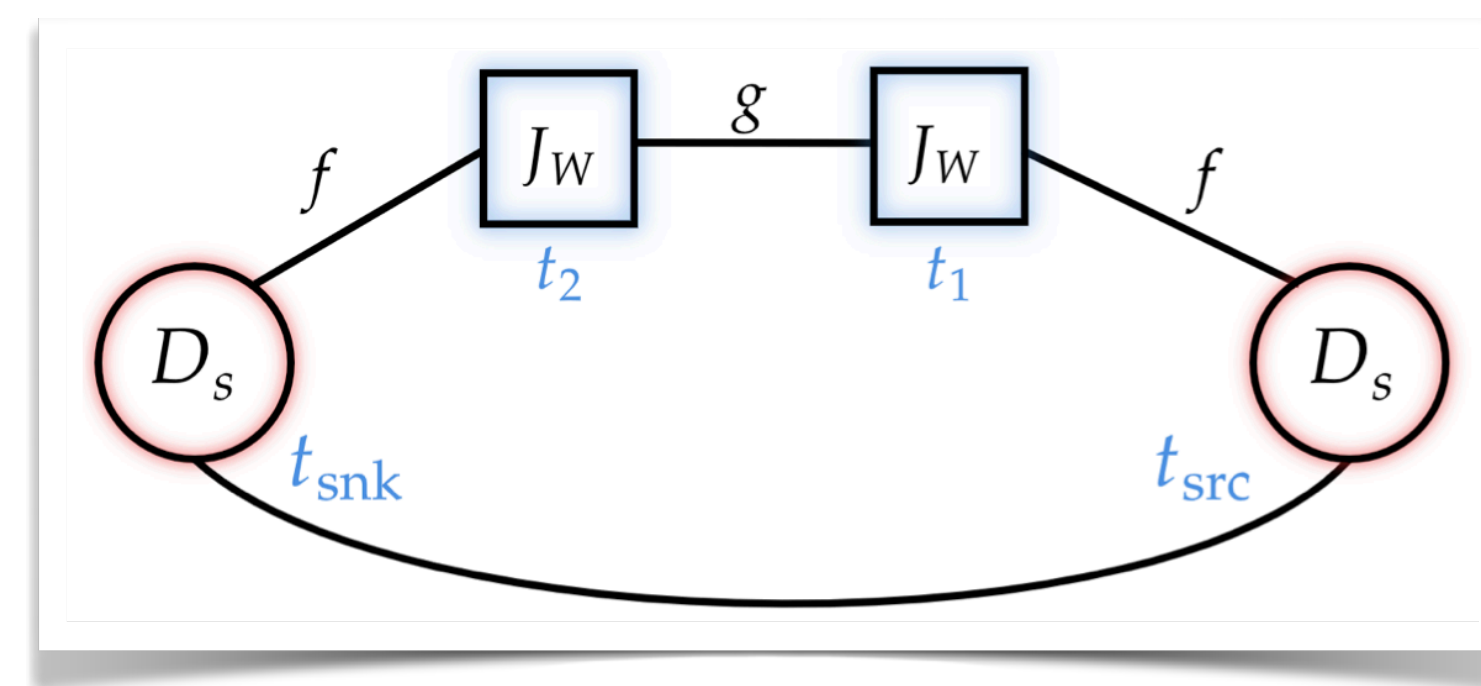
$$\frac{d\Gamma}{dq^2 dq_0 dE_l} = \frac{G_F^2 |V_{cb}|^2}{8\pi^3} L_{\mu\nu} W^{\mu\nu} \longrightarrow W^{\mu\nu} \sim \sum_X \langle B | J^\mu | X \rangle \langle X | J^{\nu\dagger} | B \rangle$$

- Ongoing work using different approaches on B and D meson decays

A.Barone et al., JHEP 07 (2023) 145



A.De Santis & C.Groß (ETMC) @ Lattice 2024



# CP violation in neutral kaons

flavour eigenstates

$$i\frac{\partial}{\partial t} \begin{pmatrix} K^0(t) \\ \bar{K}^0(t) \end{pmatrix} = \left( \mathbf{M} - \frac{i}{2}\mathbf{\Gamma} \right) \begin{pmatrix} K^0(t) \\ \bar{K}^0(t) \end{pmatrix}$$

weak eigenstates

$$|K_{L,S}\rangle = \frac{1}{\sqrt{1+|\bar{\epsilon}|^2}} (\bar{\epsilon}|K_{\pm}\rangle + |K_{\mp}\rangle)$$

CP eigenstates

$$|K_{\pm}\rangle = \frac{1}{\sqrt{2}} (|K^0\rangle \pm |\bar{K}^0\rangle)$$

Neutral mesons can mix because the flavour eigenstates are different from the weak eigenstates

- "Indirect" CP violation in the mixing

$$|\epsilon| = (2.228 \pm 0.011) \times 10^{-3}$$

PDG, PTET 2022 (2022)

- "Direct" CP violation in the decay

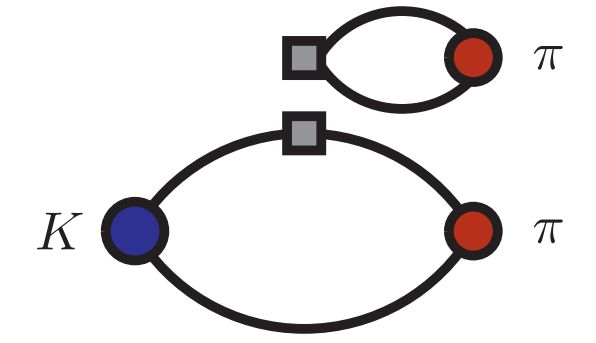
$$\eta_{+-} = \frac{\langle \pi^+ \pi^- | \mathcal{H}_w | K_L \rangle}{\langle \pi^+ \pi^- | \mathcal{H}_w | K_S \rangle} = \epsilon + \epsilon'$$

$$\eta_{00} = \frac{\langle \pi^0 \pi^0 | \mathcal{H}_w | K_L \rangle}{\langle \pi^0 \pi^0 | \mathcal{H}_w | K_S \rangle} = \epsilon - 2\epsilon'$$

$$\text{Re}(\epsilon'/\epsilon) = (1.66 \pm 0.23) \times 10^{-3}$$

PDG, PTET 2022 (2022)

# Direct CP violation in $K \rightarrow \pi\pi$



If isospin-symmetry is conserved, then the CP violation parameters can be expressed as

$$\frac{\epsilon'}{\epsilon} = \frac{ie^{i(\delta_2 - \delta_0)}}{\sqrt{2}\epsilon} \frac{\text{Re}(A_2)}{\text{Re}(A_0)} \left[ \frac{\text{Im}(A_2)}{\text{Re}(A_2)} - \frac{\text{Im}(A_0)}{\text{Re}(A_0)} \right]$$

$$A_I = \langle (\pi\pi)_I | H_W^{\Delta S=1} | K \rangle$$

$$\delta_I = \pi\pi \text{ scattering phase shifts}$$

( $I$  = isospin)

1. RBC-UKQCD performed first calculation of  $\epsilon'$  in 2015 Z.Bai et al., PRL 115 (2015)
2. Improved result in 2020: 3.5x more statistics + improved systematics R.Abbott et al., PRD 102 (2020)

$$\text{lattice: } \text{Re}(\epsilon'/\epsilon) = 21.7 (2.6)_{\text{stat.}} (8.0)_{\text{sys.}} \times 10^{-4}$$

$$\text{experiments: } \text{Re}(\epsilon'/\epsilon) = 16.6 (2.3) \times 10^{-4}$$

# Direct CP violation in $K \rightarrow \pi\pi$

## Systematic error budget

(from C.Kelly @Lattice2023)

- (~12%) Perturbation theory in Wilson coeffs to match 3f – 4f weak EFT at  $m_c$ 
  - Improve with 4f calculation (active charm) : computationally infeasible?
  - Non-perturbative calculation of matching matrix : investigation underway

[M.Tomii, PoS LATTICE2018 (2019) 216]
- (~23%) Lack of EM+isospin-breaking contributions in lattice calculation
  - Lattice measurement of these effects extremely challenging but approach is being formulated.

[Phys.Rev.D 106 (2022) 1, 014508] V.Cirigliano et al.,  
[Christ, PoS LATTICE2021 (2022) 312] JHEP 02 (2020)
- (~12%) Use of single lattice spacing to compute  $I=0$  amplitude
  - Repeat calculation with multiple, finer lattice spacings: **my current focus**

- Intense work by RBC-UKQCD to reduce ~12% error due single lattice spacing (C.Kelly @Lattice2024)  
+ parallel ongoing project using different computational approach (M.Tomii @Lattice2024)
- **IB correction will soon become relevant!** (but very tricky to compute on the lattice)  
Usually  $O(1\%)$ , but the " $\Delta I = 1/2$  rule" can give a ~20x enhancement in  $\epsilon'/\epsilon$ .

# Indirect CP violation in kaon mixing

$$\epsilon = e^{i\phi_\epsilon} \sin(\phi_\epsilon) \left( \frac{-\text{Im}M_{\bar{0}0}}{\Delta M_K} + \frac{\text{Im}(A_0)}{\text{Re}(A_0)} \right)$$

$$\tan(\phi_\epsilon) = \frac{\Delta M_K}{\Delta\Gamma_K/2}$$

$$\Delta\Gamma_K = \Gamma_{K_S} - \Gamma_{K_L}$$

$$\Delta M_K = M_{K_L} - M_{K_S}$$

The quantity  $M_{\bar{0}0}$  splits into a **short** and **long** distance parts

$$\begin{aligned} M_{\bar{0}0} &= \langle \bar{K}^0 | \mathcal{H}_w | K^0 \rangle = \langle \bar{K}^0 | \mathcal{H}_w | K^0 \rangle_{\text{SD}} + \langle \bar{K}^0 | \mathcal{H}_w | K^0 \rangle_{\text{LD}} \\ &= \langle \bar{K}^0 | \mathcal{H}_w^{\Delta S=2} | K^0 \rangle + \mathcal{P} \sum_n \frac{\langle \bar{K}^0 | \mathcal{H}_w^{\Delta S=1} | n \rangle \langle n | \mathcal{H}_w^{\Delta S=1} | K^0 \rangle}{M_K - E_n} \end{aligned}$$

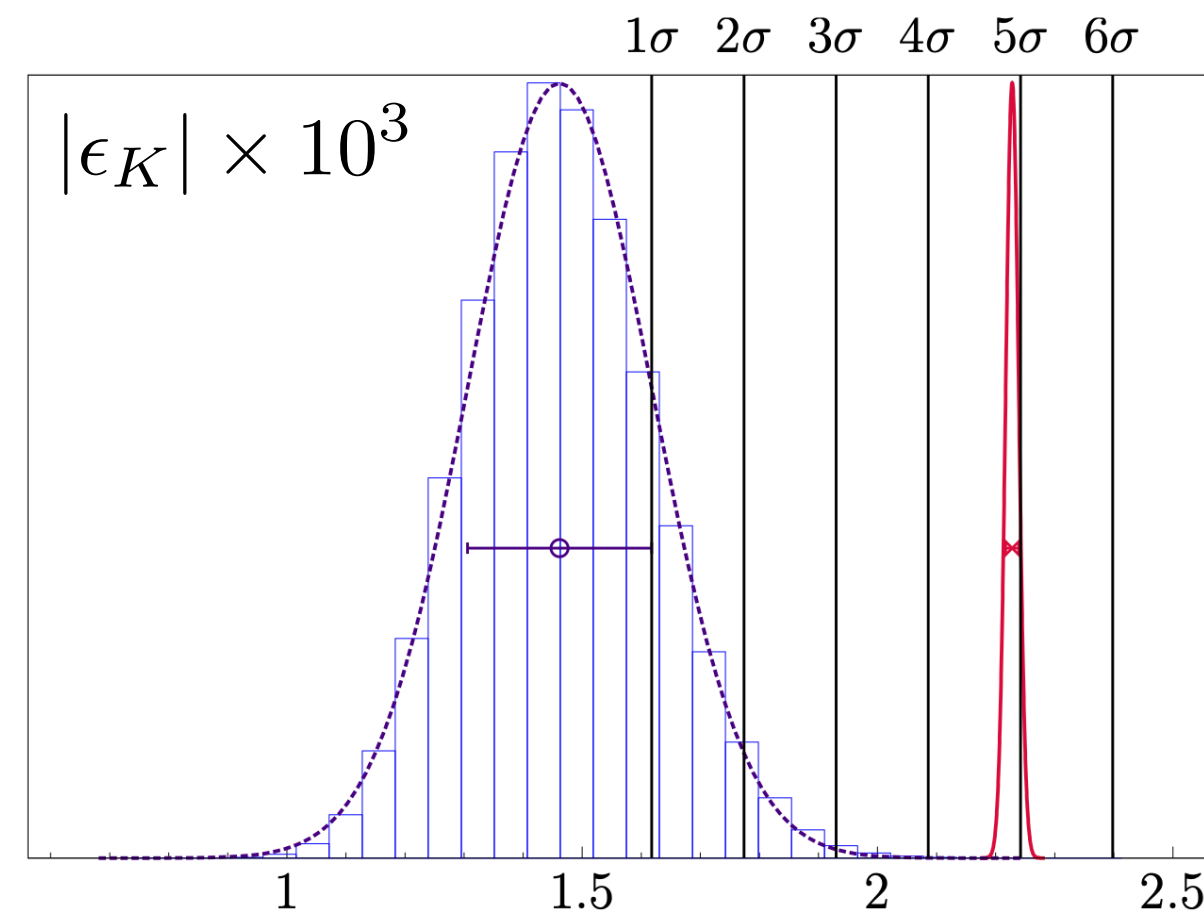
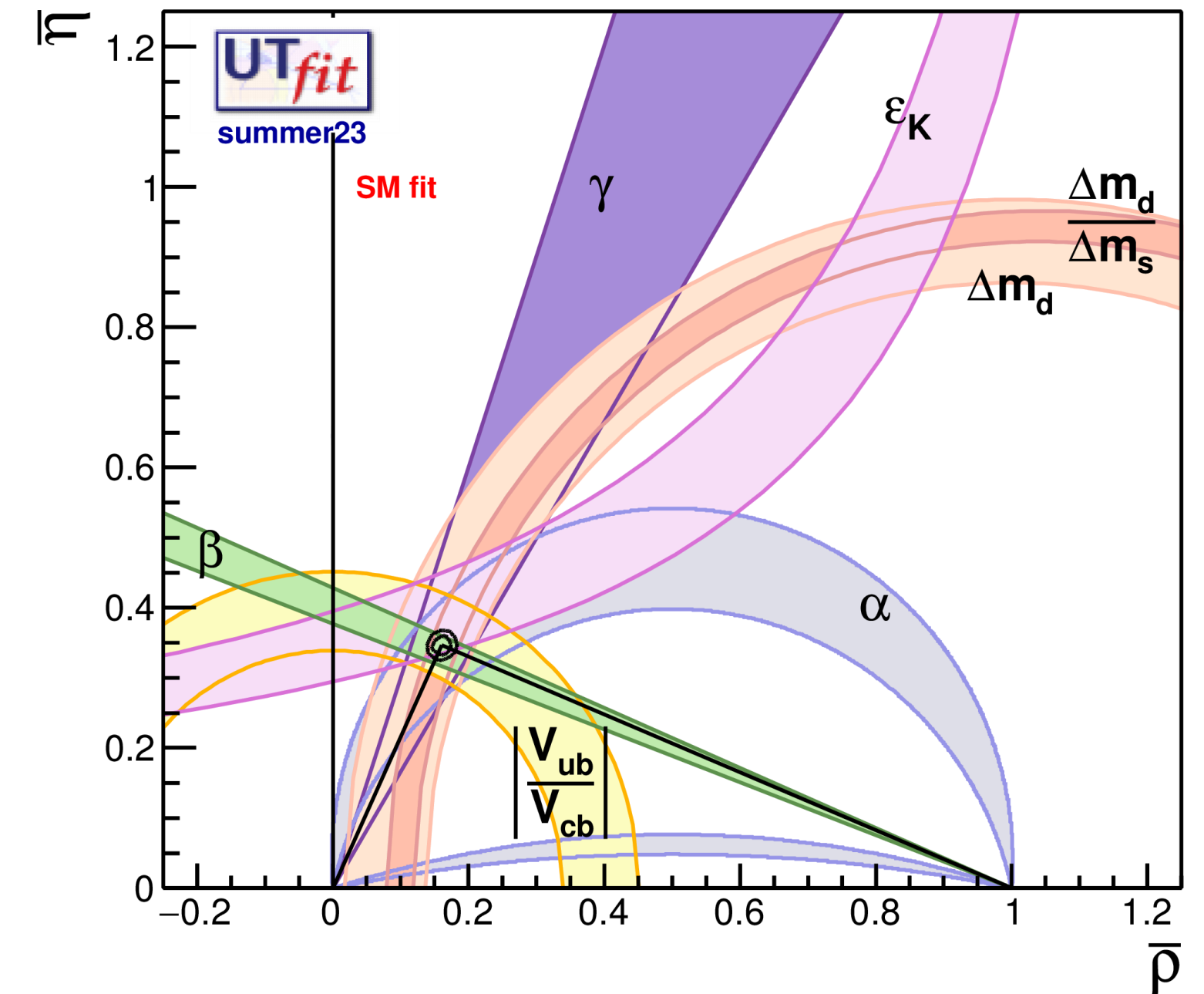
On the **lattice**, we can compute both:

- $\langle \bar{K}^0 | \mathcal{H}_w | K^0 \rangle_{\text{SD}}$  e.g. P.A.Boyle et al. PRD 110 (2024) ← dominating contribution
- $\langle \bar{K}^0 | \mathcal{H}_w | K^0 \rangle_{\text{LD}}$  Z.Bai et al. PRD 109 (2024)

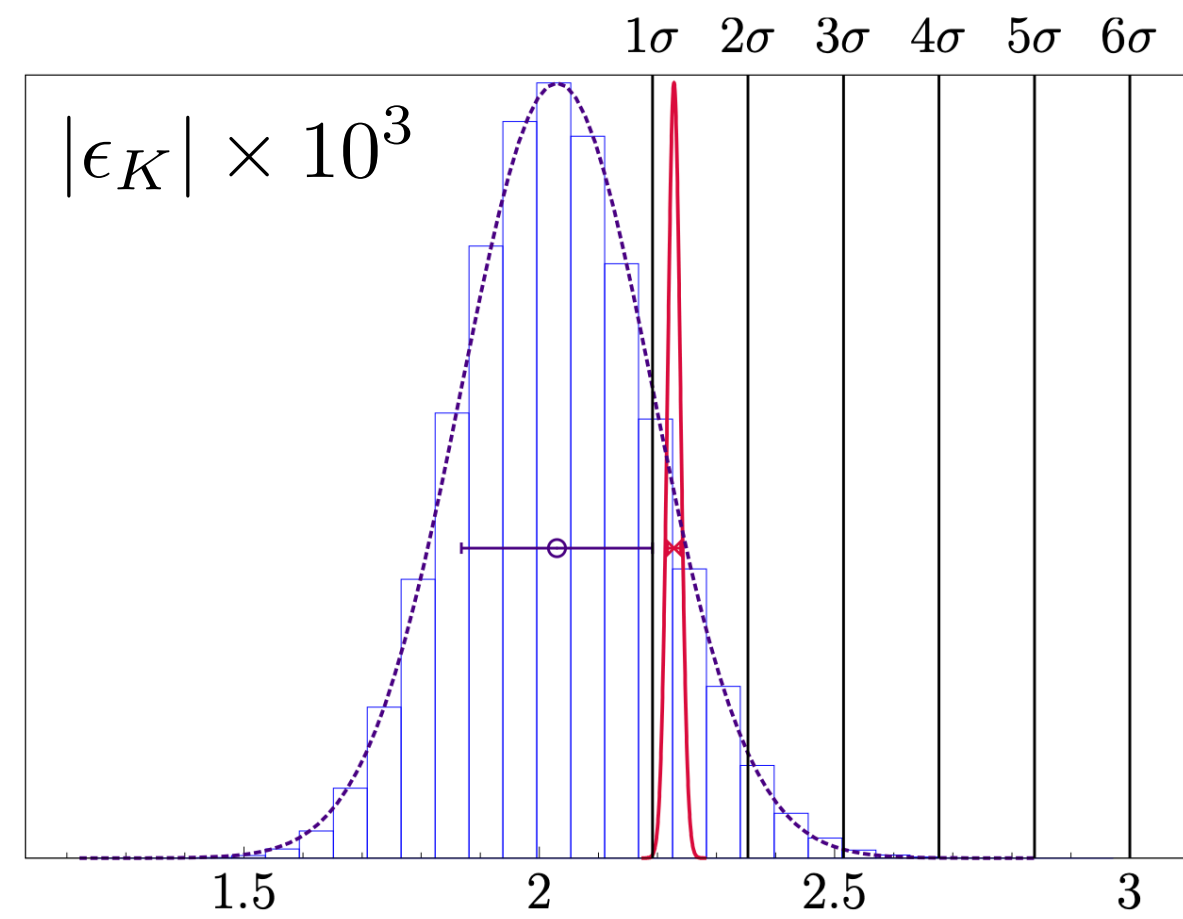


# CP violation in neutral kaons

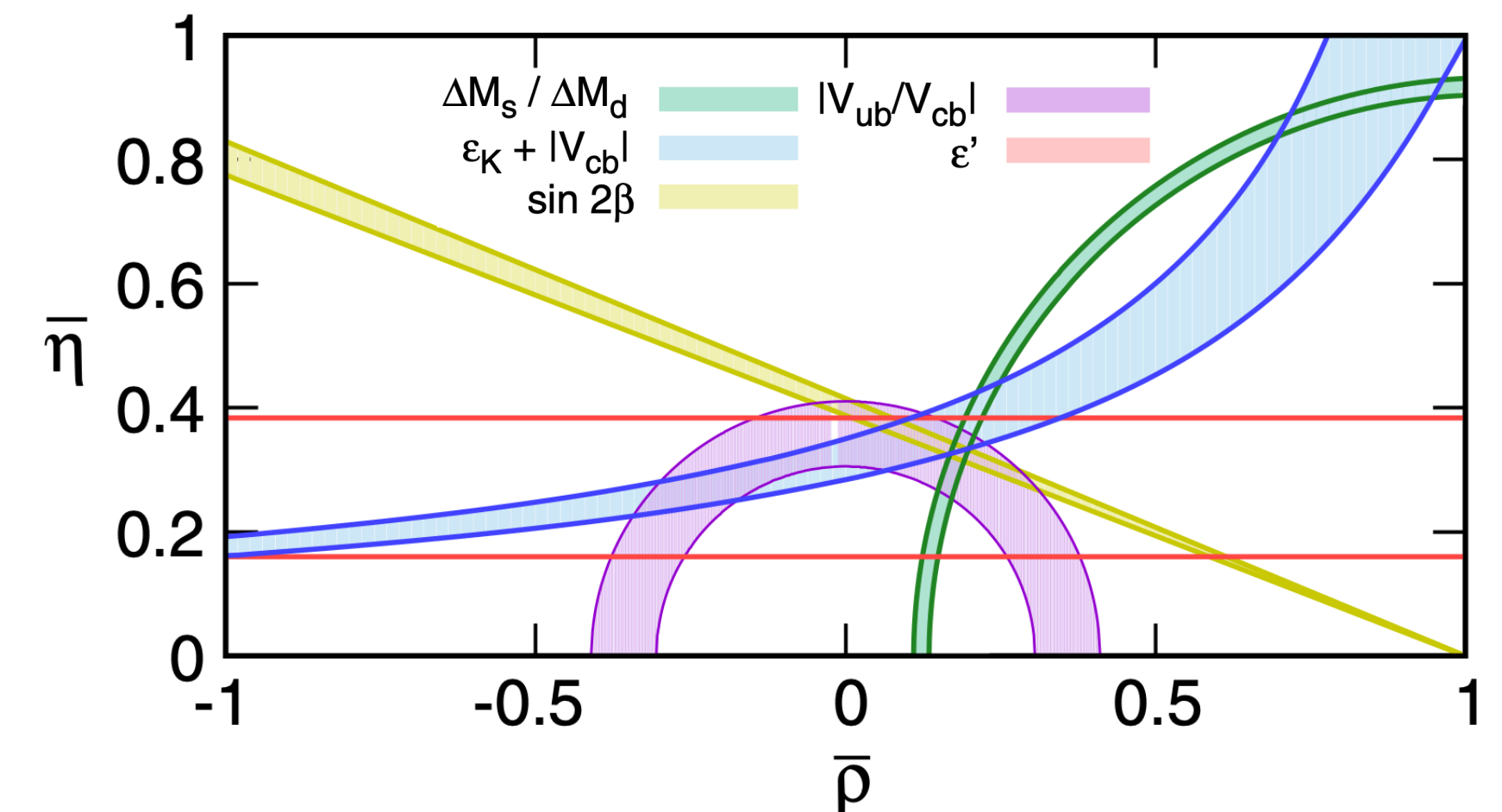
- $\text{Re}(\epsilon'/\epsilon)$  currently at **40% precision**.  
Significant improvements expected in next couple of years
- $(\epsilon_K)_{\text{LD}}$  also at **40% precision**.  
Errors of  $\sim 10\%$  can be achieved on the long term
- Lattice inputs to  $(\epsilon_K)_{\text{SD}}$  can be computed with high precision, but **overall uncertainty is dominated by  $|V_{cb}|$**
- $|V_{cb}|$ -puzzle affects the SM prediction for  $|\epsilon_K|$



(a) Exclusive  $|V_{cb}|$  (FNAL/MILC-22, BGL)



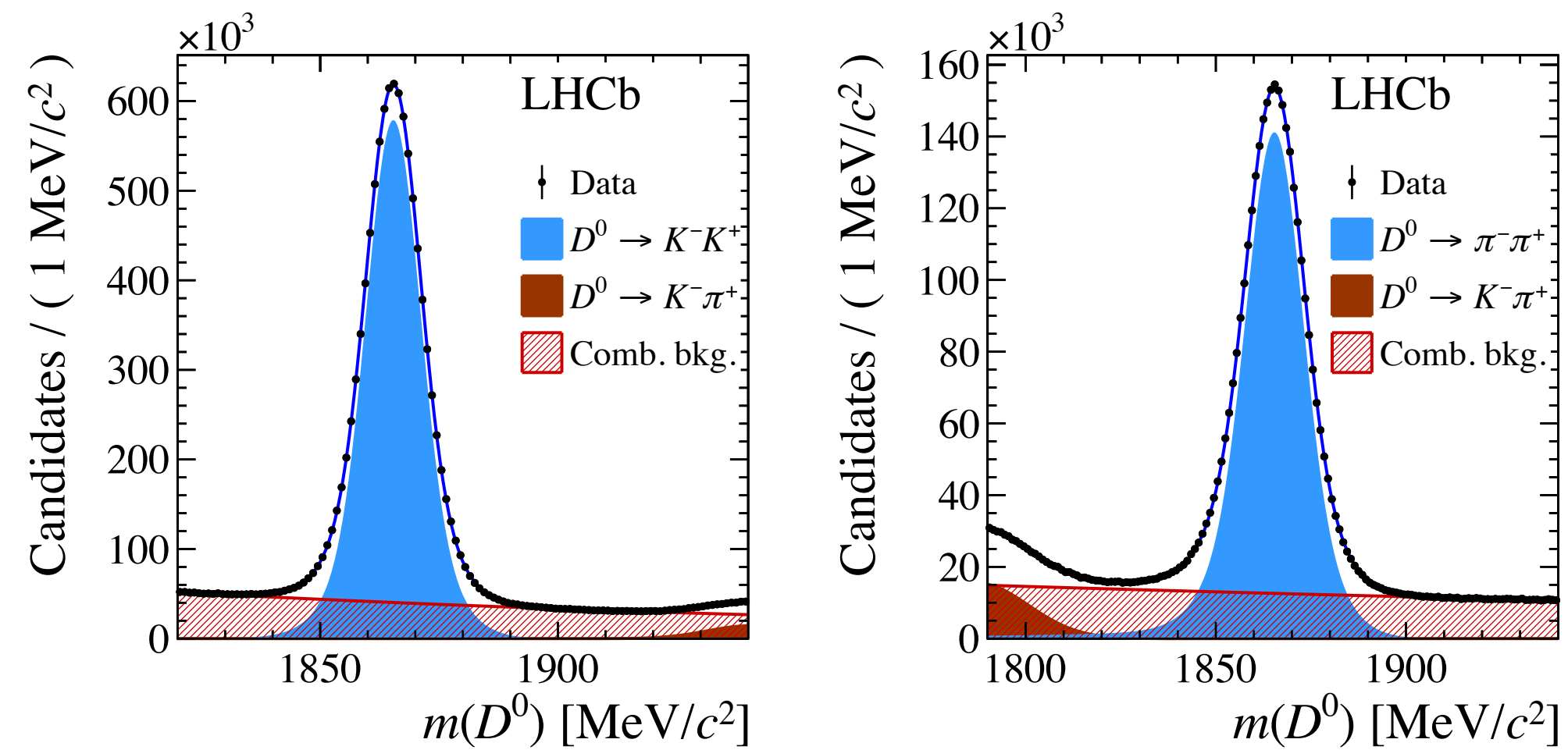
(b) Inclusive  $|V_{cb}|$  (HFLAV-23, 1S scheme)



R.Abbott et al., PRD 102 (2020)

# The charm sector

## Hadronic $D$ decays

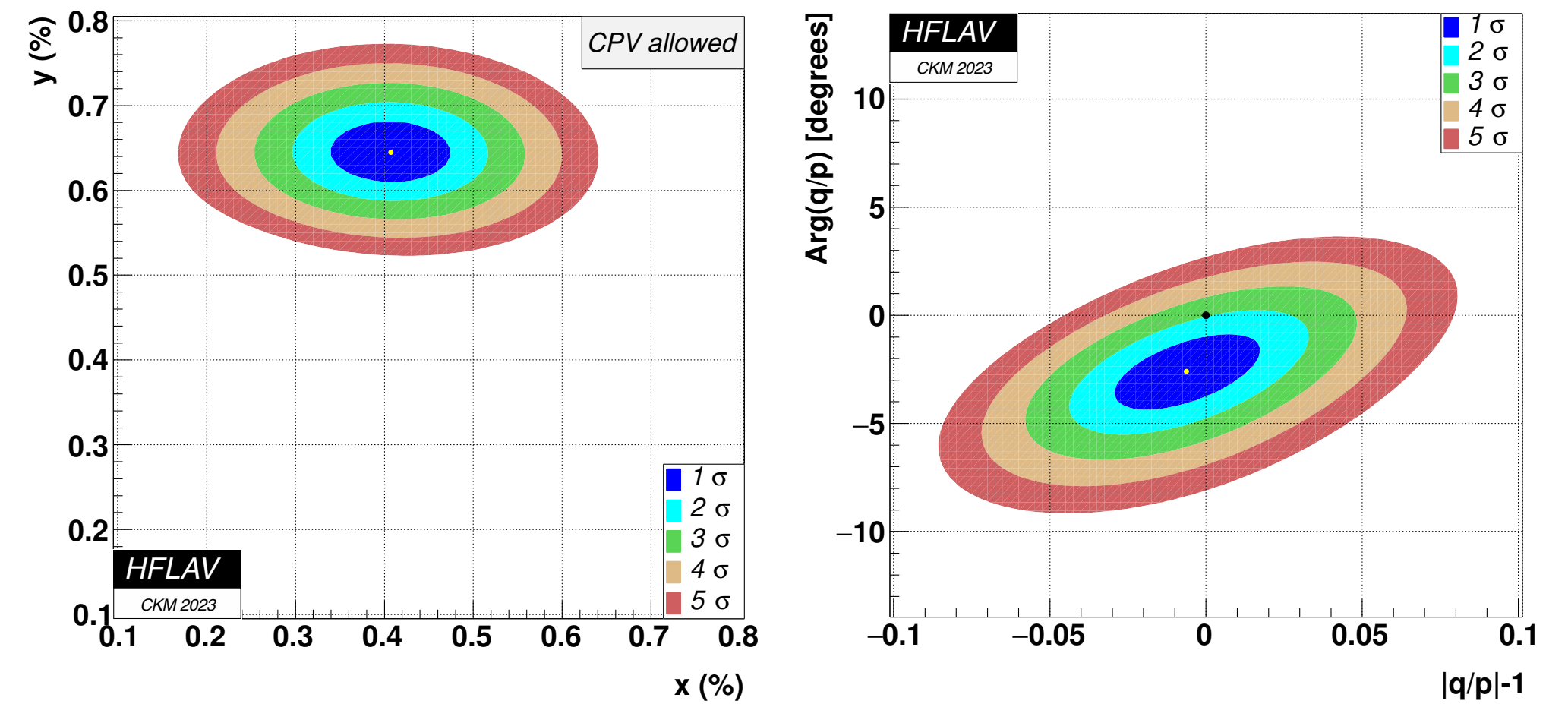


LHCb, PRL 122 (2019)

- $$\Delta A_{CP} = A_{CP}(K^- K^+) - A_{CP}(\pi^- \pi^+)$$

$$= (-15.4 \pm 2.9) \times 10^{-4}$$

## $D^0$ - $\bar{D}^0$ mixing

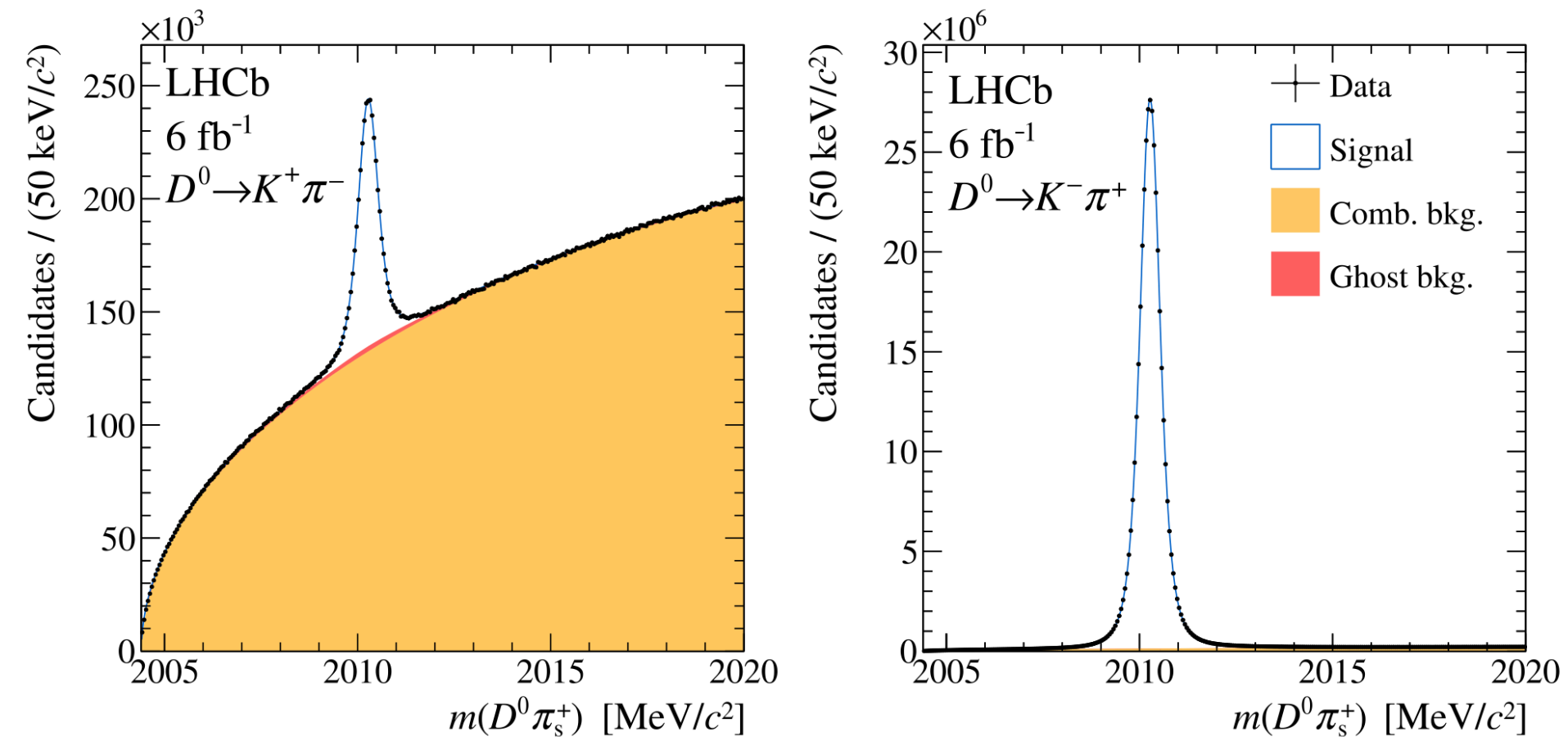


HFLAV, PRD 107 (2023)

- no-mixing point excluded at  $> 11.5 \sigma$
- no evidence for indirect CP violation

# The charm sector

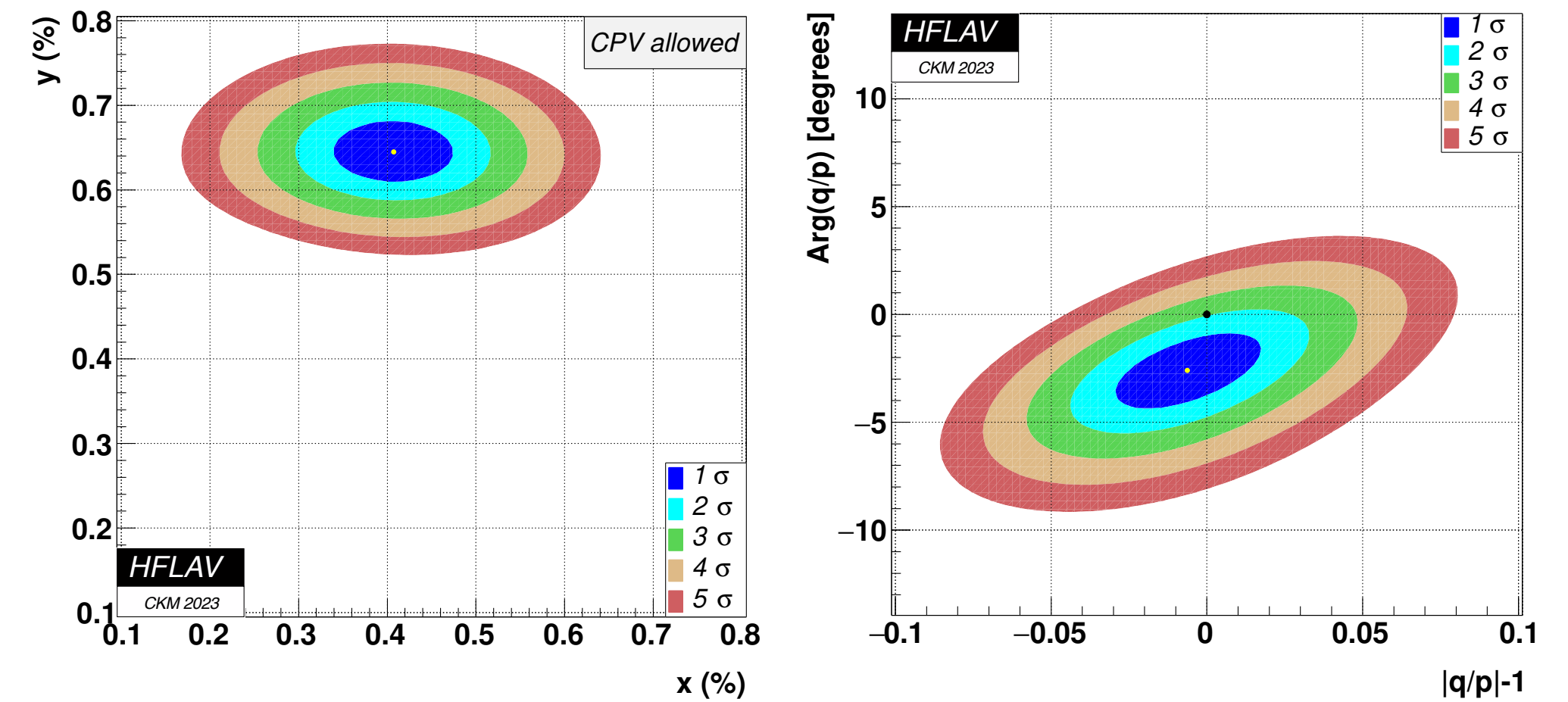
## Hadronic $D$ decays



LHCb, [arXiv:2407.18001]

- Measure of CP violating parameters in  $D^0 \rightarrow K^\pm \pi^\mp$

## $D^0$ - $\bar{D}^0$ mixing



HFLAV, PRD 107 (2023)

- no-mixing point excluded at  $> 11.5 \sigma$
- no evidence for indirect CP violation

# Towards hadronic D decays on the lattice

- First exploratory calculation of  $D \rightarrow K\pi$  at  $SU(3)_F$  symmetric point F.Joswig et al., PoS LATTICE2022 (2023)
- Difficult calculation with many challenges

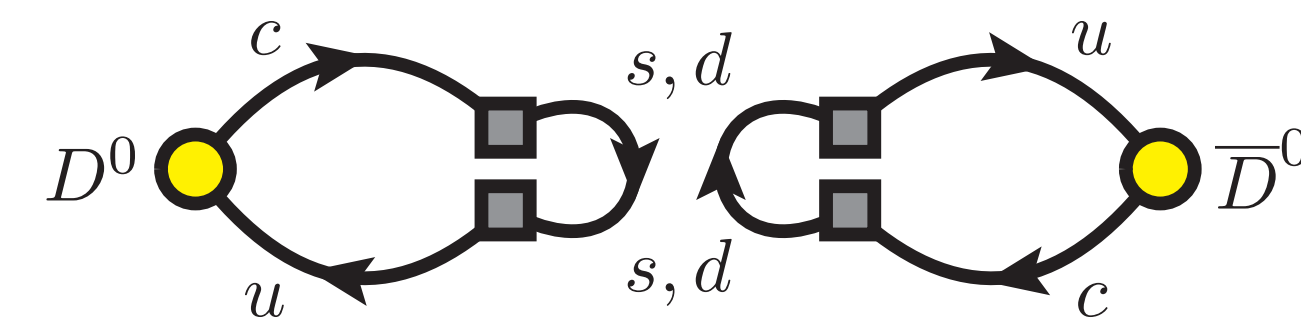
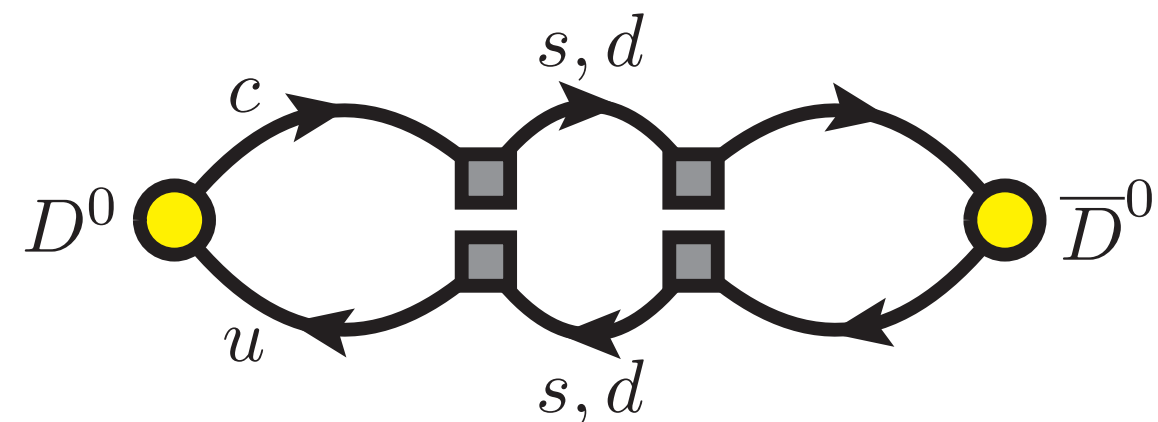
$$A(D \rightarrow h_1 h_2) = C_{n,L,h_1 h_2}^{\text{LL}} \left[ \lim_{a \rightarrow 0} Z^{\overline{\text{MS}}} \langle n, L | \mathcal{H}_W | D, L \rangle \right]$$

- Non-perturbative renormalization of four-quark operators ○ ○ ○ ○ ○
- Reliable creation of excited multi-hadron final states ○ ○ ○ ○ ○
- Removal of discretization effects (enhanced by the charm mass) ○ ○ ○ ○ ○
- Formalism to relate finite-volume matrix elements to the amplitudes ○ ○ ○ ○ ○
- Extraction of the matrix element from three-point functions ○ ○ ○ ○ ○

from M.T.Hansen @Lattice2023

# $D^0$ - $\bar{D}^0$ mixing on the lattice?

- long distance contributions dominate
- huge number of intermediate states makes direct calculations impossible
- we can take advantage of recent developments in spectral density reconstruction techniques
- use  $i\epsilon$ -prescription as a "smearing" of spectral densities related to the mixing parameters



$$\rho(\omega) = \langle \bar{D}^0 | H_w (2\pi) \delta(\mathbb{H} - \omega) H_w | D^0 \rangle$$

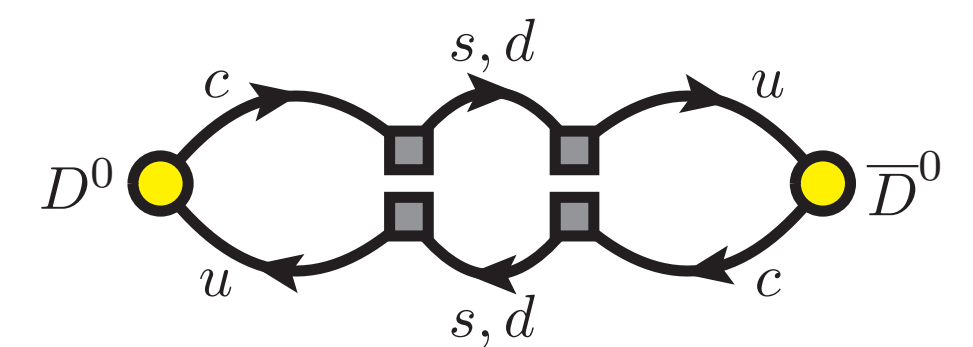
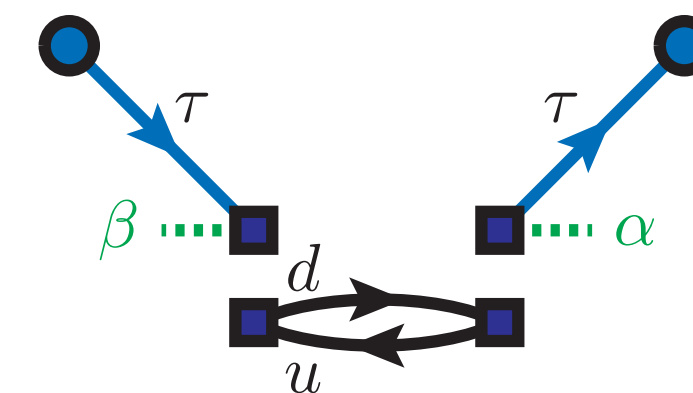
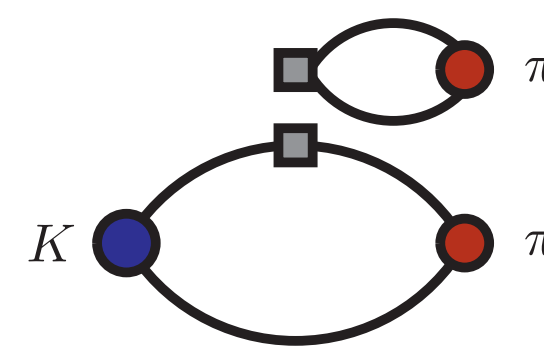
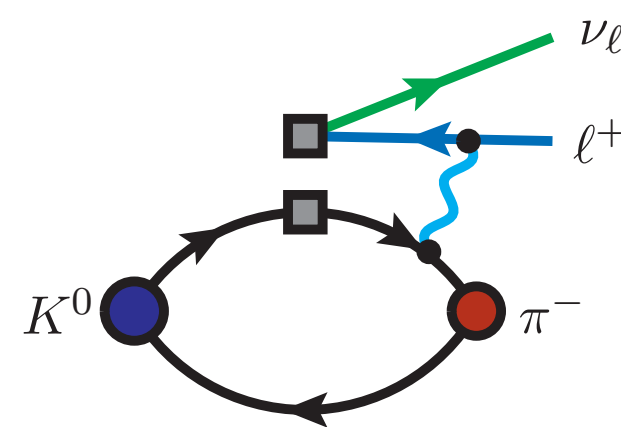
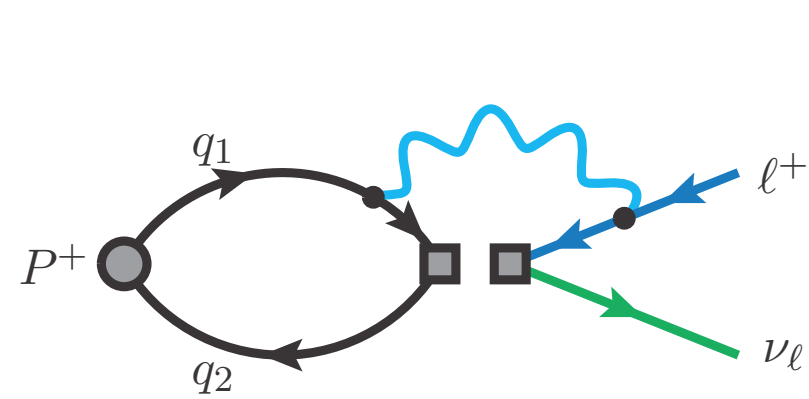
$$\begin{aligned} \hat{\rho}_L(E, \epsilon) &= \int_0^\infty \frac{d\omega}{2\pi} \frac{\rho_L(\omega)}{\omega - E - i\epsilon} \\ &= \sum_{t=0}^T \left\{ g_t^R(E, \epsilon) + i g_t^I(E, \epsilon) \right\} C_L(t) \end{aligned}$$

$$M_{12} \propto \lim_{\epsilon \rightarrow 0} \lim_{L \rightarrow \infty} \text{Re}[\hat{\rho}_L(m_D, \epsilon)]$$

$$\Gamma_{12} \propto \lim_{\epsilon \rightarrow 0} \lim_{L \rightarrow \infty} \text{Re}[\hat{\rho}_L(m_D, \epsilon)]$$

# Conclusions

- Flavour physics offers unique opportunities for indirect searches of New Physics
- Lattice QCD is at a mature stage on many flavour observables and is now in the precision era: towards physical pion masses, QED & isospin-breaking effects, physical b quarks, ...
- Progress expected on Cabibbo anomaly,  $V_{cb}$  puzzle and CPV in kaons in next years. Maybe some of the tensions will be clarified?
- Developments of new techniques & algorithmic advances make it possible to study long-distance observables considered before inaccessible!



# Thank you



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**Backup slides**



# Prospects for $|V_{us}/V_{ud}|$

An exercise on the error budget

$$\left| \frac{V_{us}}{V_{ud}} \right|^2 = \left[ \frac{\Gamma(K_{\ell 2}) M_{K^+}^3 (M_{K^+}^2 - M_{\mu^+}^2)^2}{\Gamma(\pi_{\ell 2}) M_{\pi^+}^3 (M_{\pi^+}^2 - M_{\mu^+}^2)^2} \right]_{\text{exp}} \cdot \left[ \frac{f_{K,0}}{f_{\pi,0}} \right]^2 (1 + \delta R_{K\pi})$$

- Using our new result

$$\delta R_{K\pi} = -0.0086 (13)(39)_{\text{vol.}}$$

	$[f_{K,0}/f_{\pi,0}]$	$ V_{us}/V_{ud} $
FLAG21 2+1 average	1.1930 (33)	0.23154 (28) <sub>exp</sub> (15) <sub><math>\delta R</math></sub> (45) <sub><math>\delta R, \text{vol.}</math></sub> (65) <sub><math>f_P</math></sub>

- Using RM123S result

$$\delta R_{K\pi} = -0.0126 (14)$$

	$[f_{K,0}/f_{\pi,0}]$	$ V_{us}/V_{ud} $
FLAG19 2+1+1 average	1.1966 (18)	0.23131 (28) <sub>exp</sub> (17) <sub><math>\delta R</math></sub> (35) <sub><math>f_P</math></sub>