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Weak cosmic censorship and the rotating quantum BTZ black hole

Andrea Pierfrancesco Sanna





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Wald's test

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Throwing **test** particles at an extremal Kerr black hole

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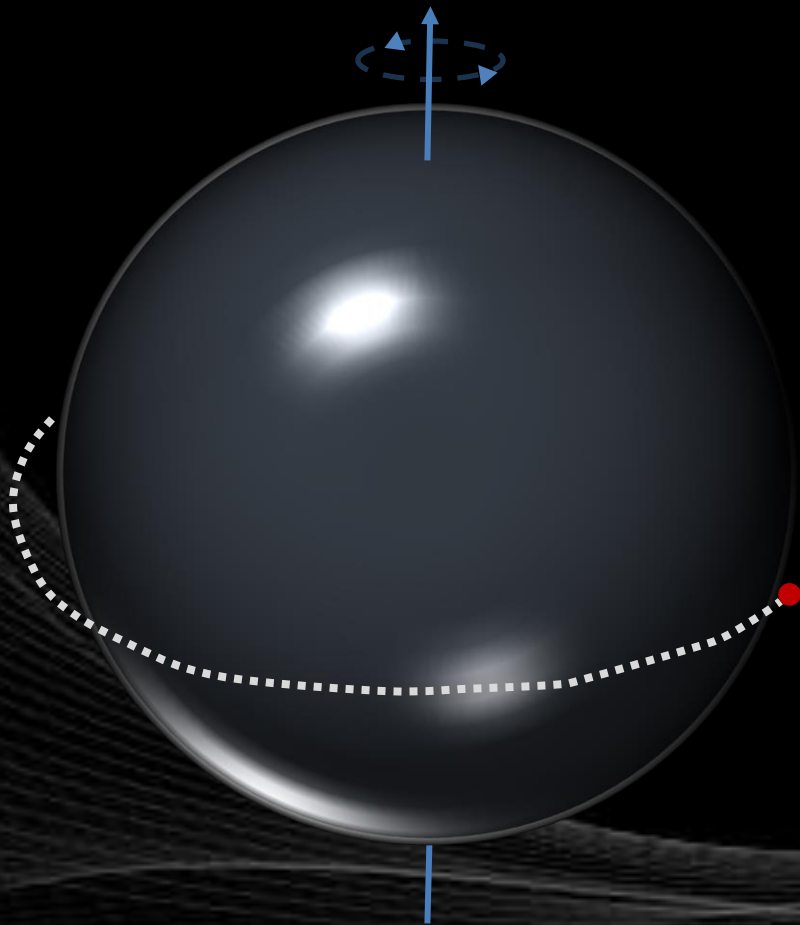
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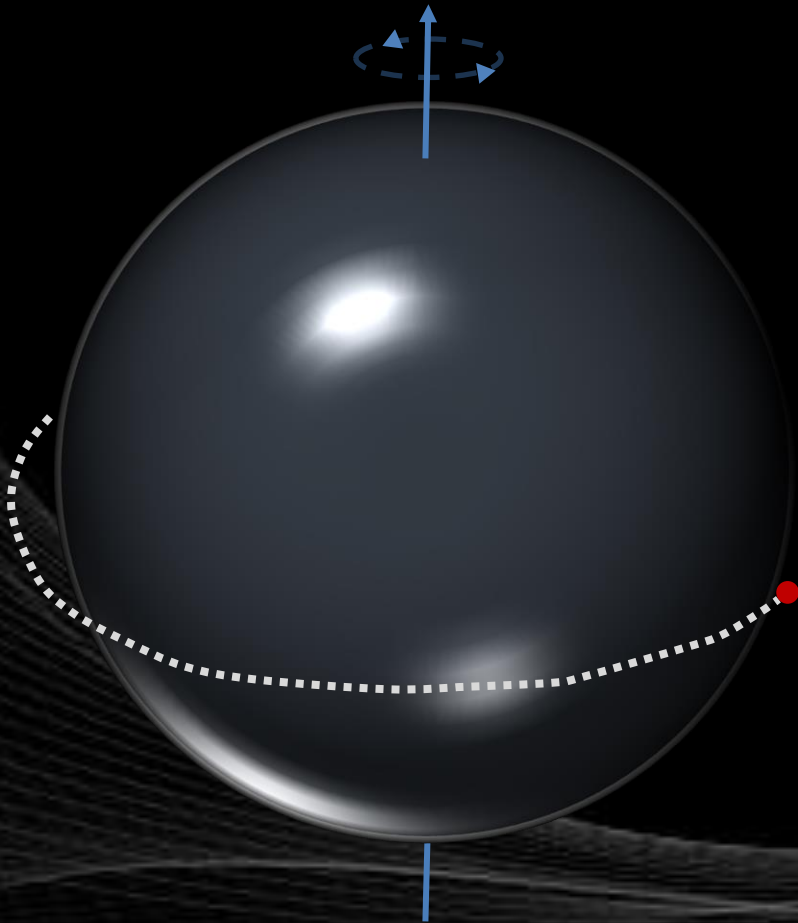


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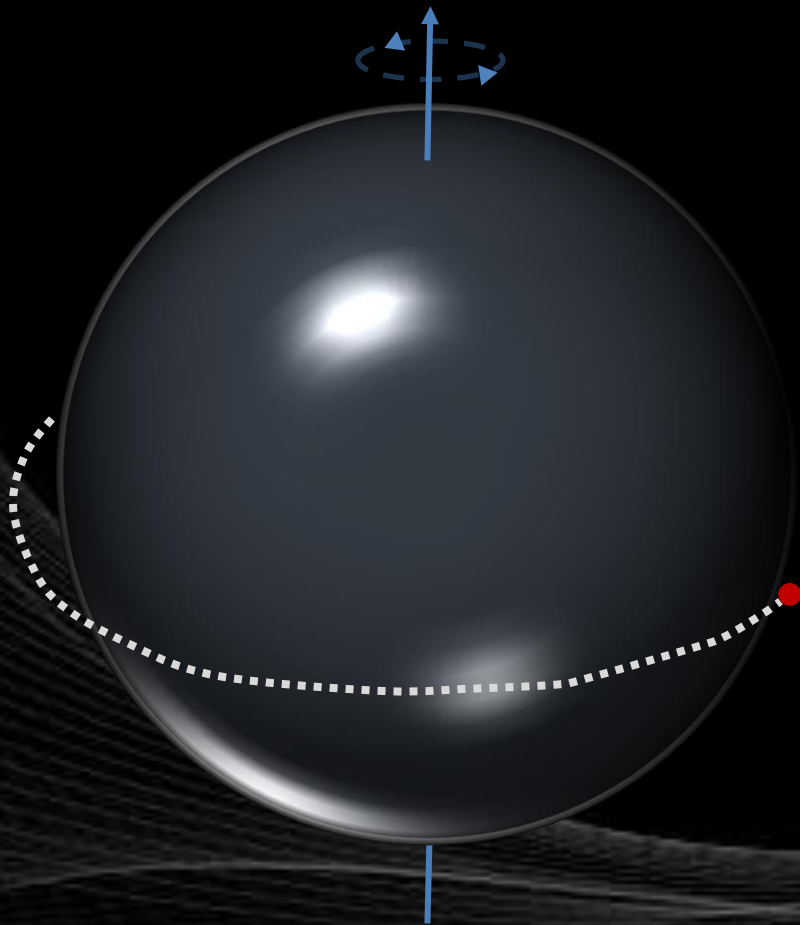
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Particles with dangerously high angular momenta are **not captured** by the black hole

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No naked singularity is formed

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Can one spin up the black hole past extremality?



The diagram shows a black hole with a vertical blue arrow indicating its spin axis. A dashed blue circle with arrows around the axis represents rotation. A dotted white line shows the path of a test particle orbiting the black hole, with a red dot at the point of impact. The background features a grid of lines representing spacetime curvature.

What is the impact of quantum corrections?

No naked singularity is formed

quantum black holes

Semiclassical gravity

$$G_{\mu\nu}(g_{\alpha\beta}) = 8\pi G \langle T_{\mu\nu}(g_{\alpha\beta}) \rangle$$

Classical geometry of a black hole modified by quantum fields

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Holographic formulation

Braneworld Holography



AdS_{D+1}

∂AdS

Braneworld Holography

Limitation of conventional AdS/CFT: boundary of AdS is **fixed**



AdS_{D+1}

∂AdS

Braneworld Holography

Limitation of conventional AdS/CFT: boundary of AdS is **fixed** \longrightarrow Introducing the brane

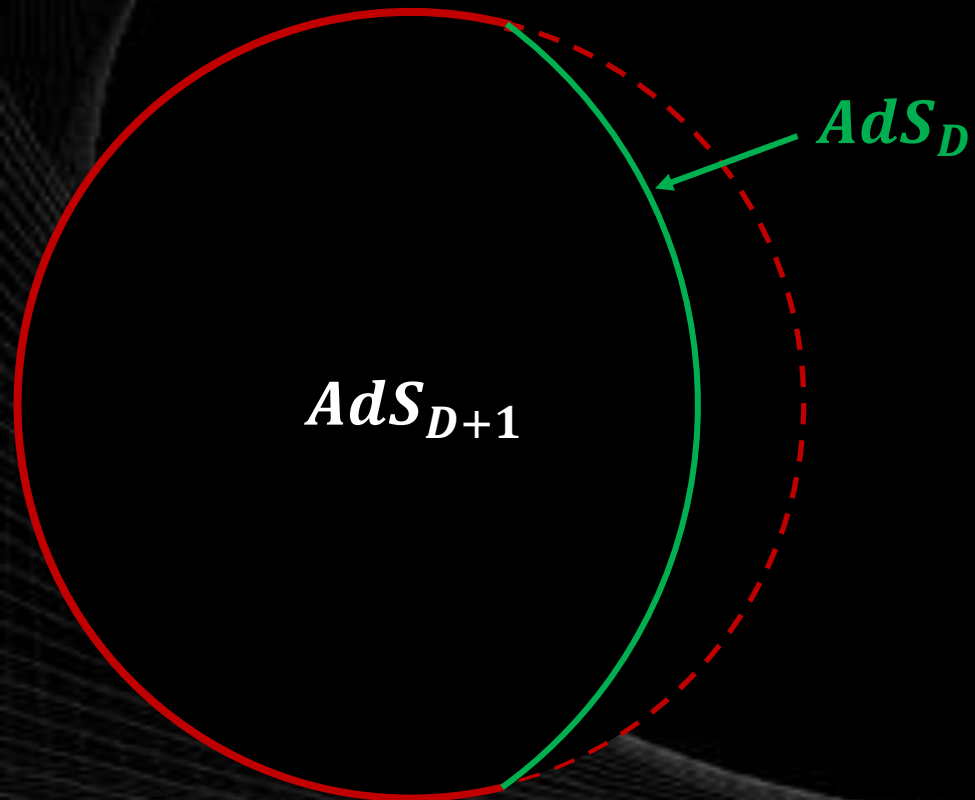


The diagram shows a large red circle representing the boundary of AdS_{D+1} . Inside this circle, a green circle represents a brane. The green circle is smaller than the red one and is positioned such that it is tangent to the red circle at two points. The text AdS_{D+1} is written in white inside the red circle.

AdS_{D+1}

Braneworld Holography

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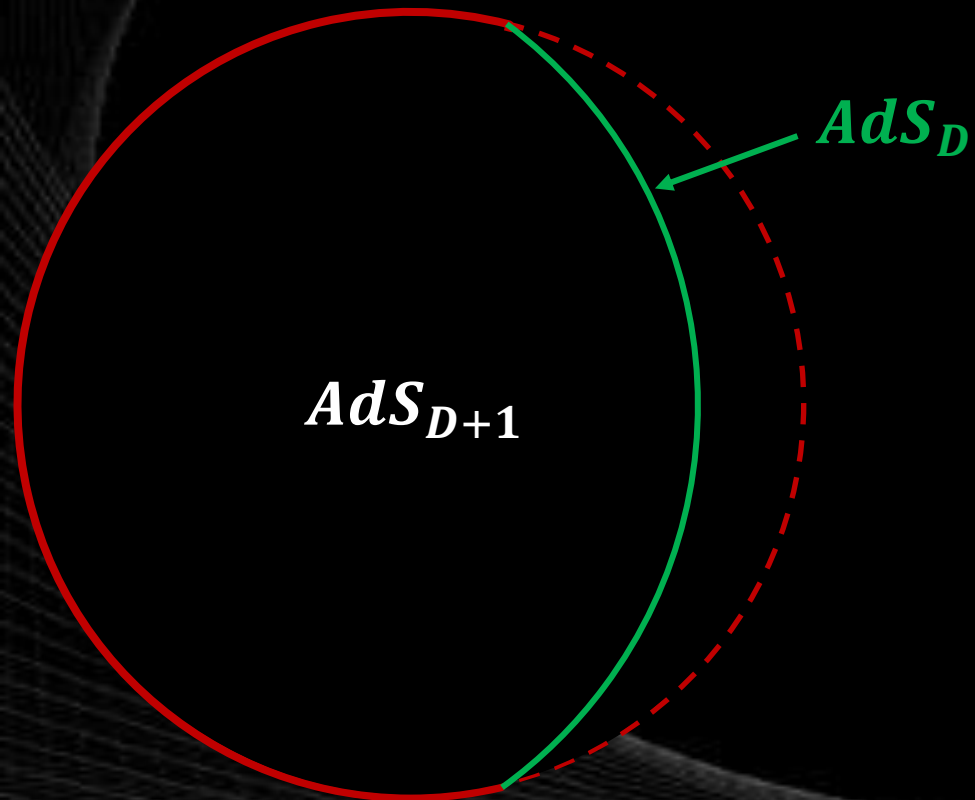


The geometry on the brane is dynamical!
Graviton modes are induced on the brane

Randall and Sundrum, PRL **83** (1999)
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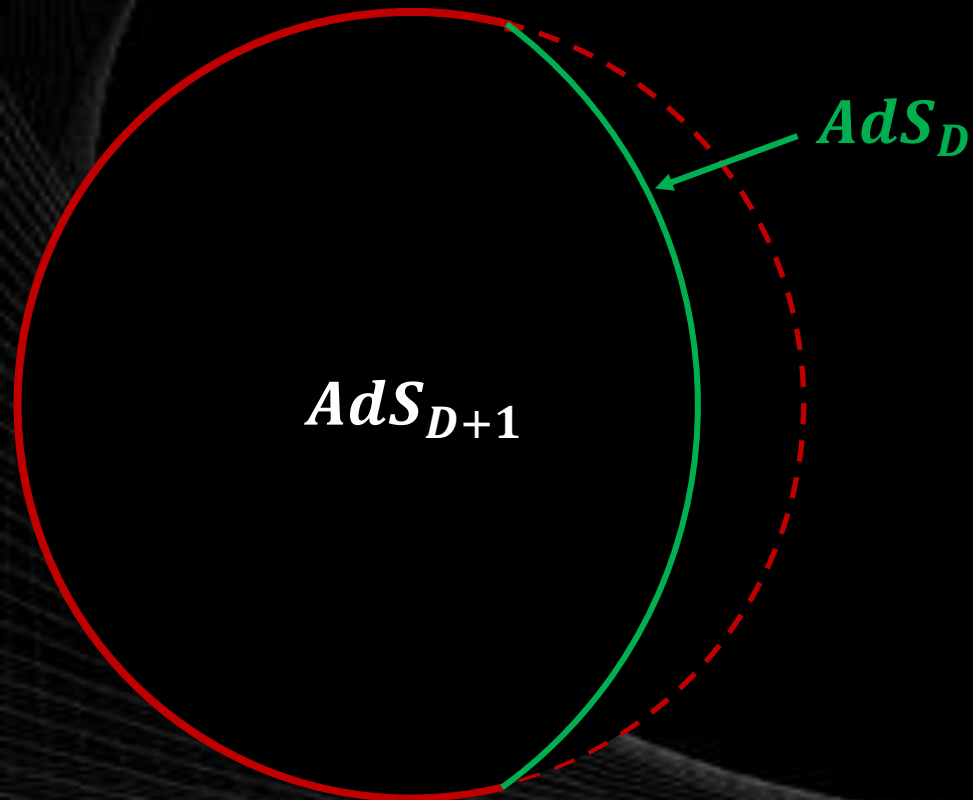
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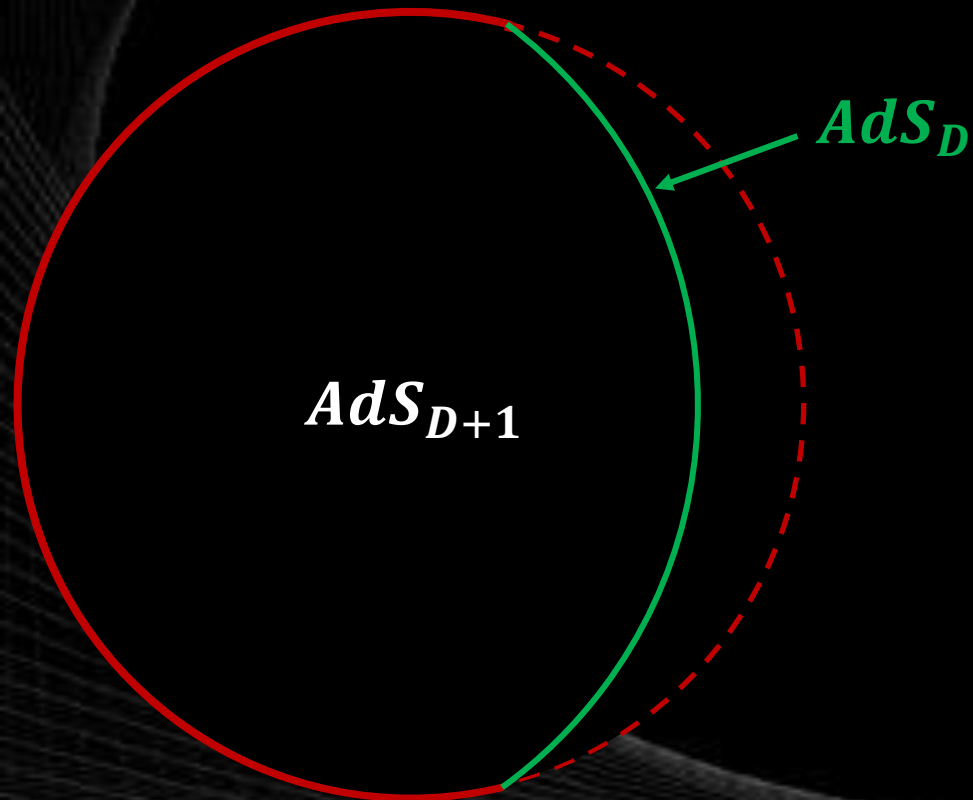
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$$\mathcal{S} = \frac{\ell_4}{8\pi G_4} \int d^3x \sqrt{-h} \left[\frac{4}{\ell_4^2} \left(1 - \frac{\ell_4}{\ell} \right) + R + \ell_4^2 \left(\frac{3}{8} R^2 - R_{\mu\nu} R^{\mu\nu} \right) + \dots \right] + \mathcal{S}_{\text{CFT}}$$

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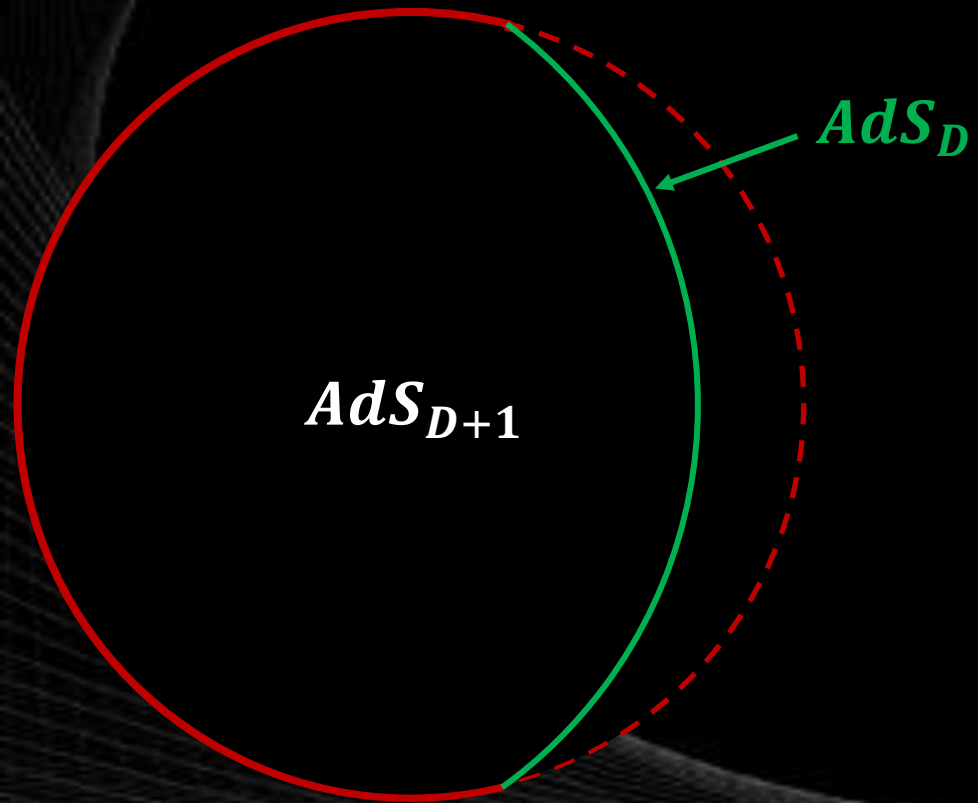
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$$\ell^{-1} \sim \text{Brane tension} \sim \text{Cutoff}$$

Braneworld Holography

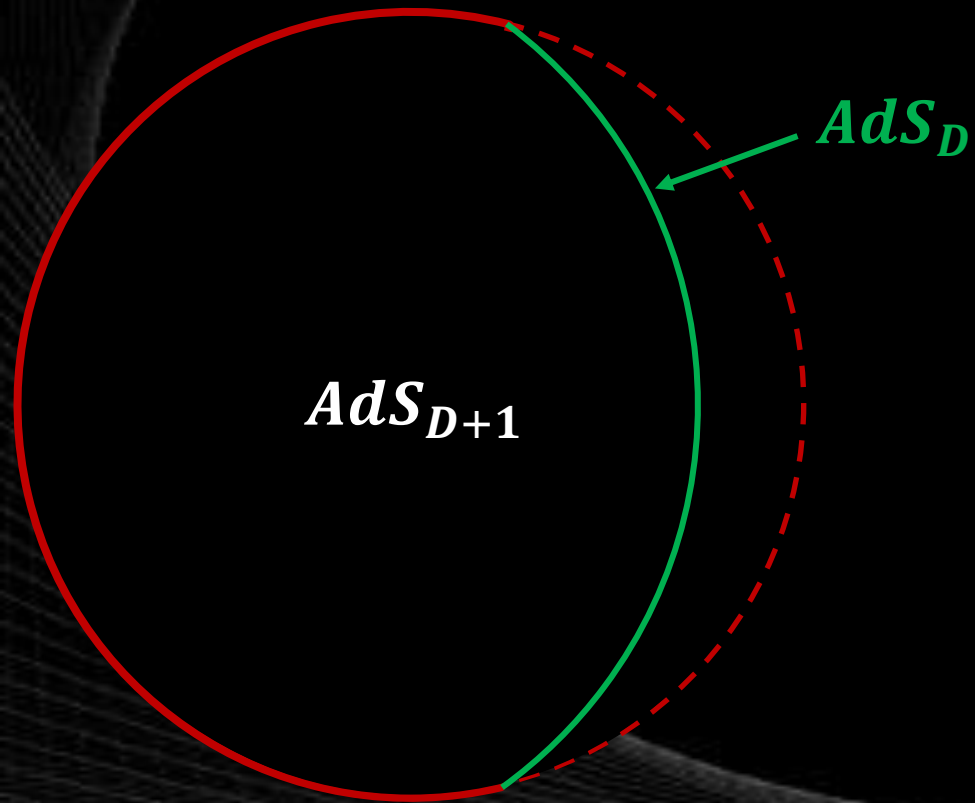
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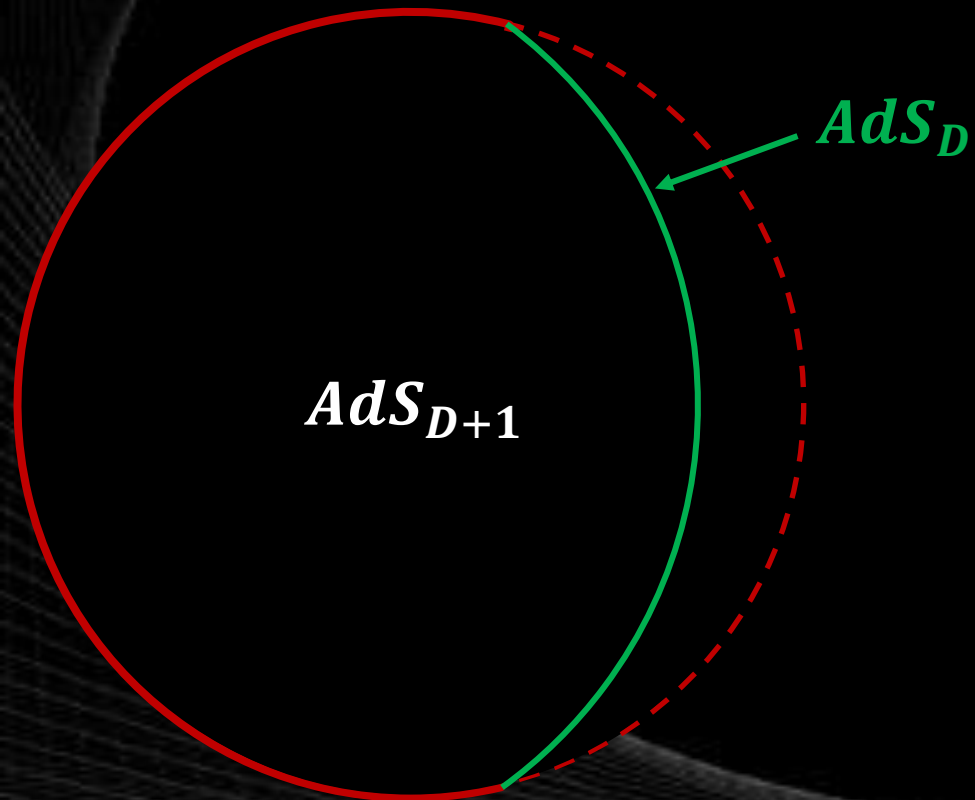


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Einstein + higher curvature terms

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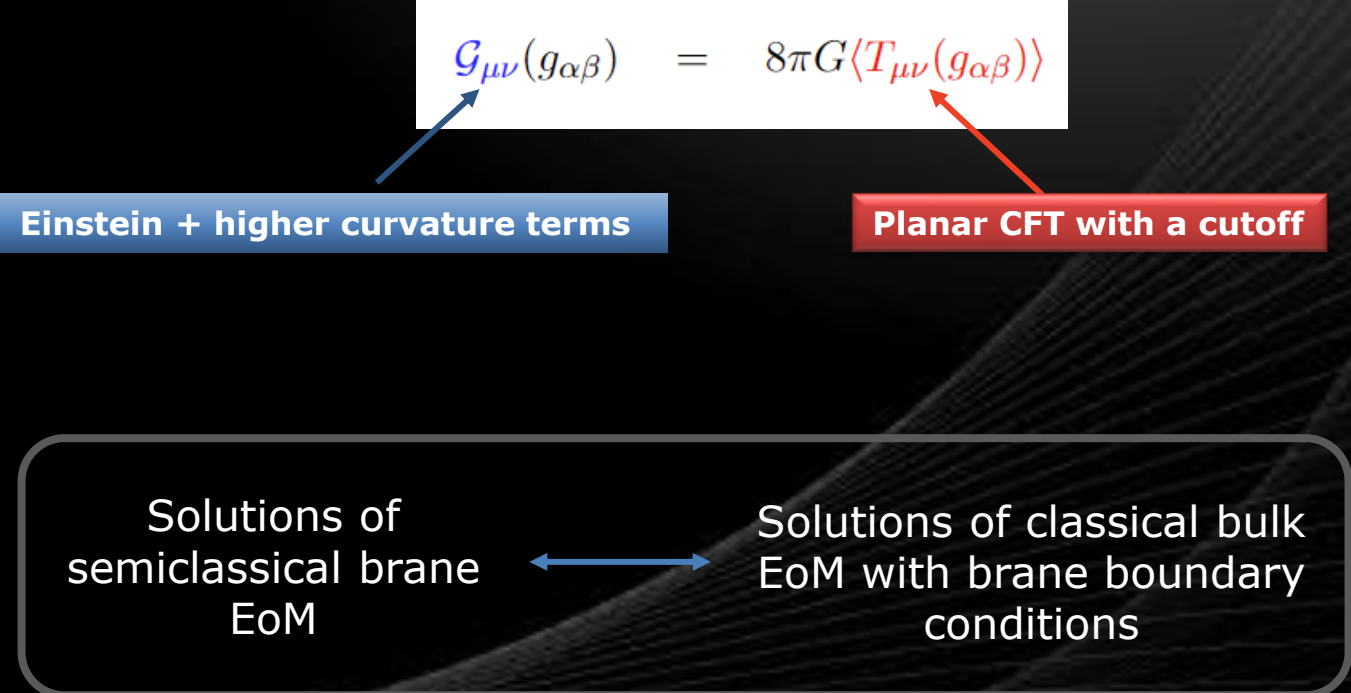
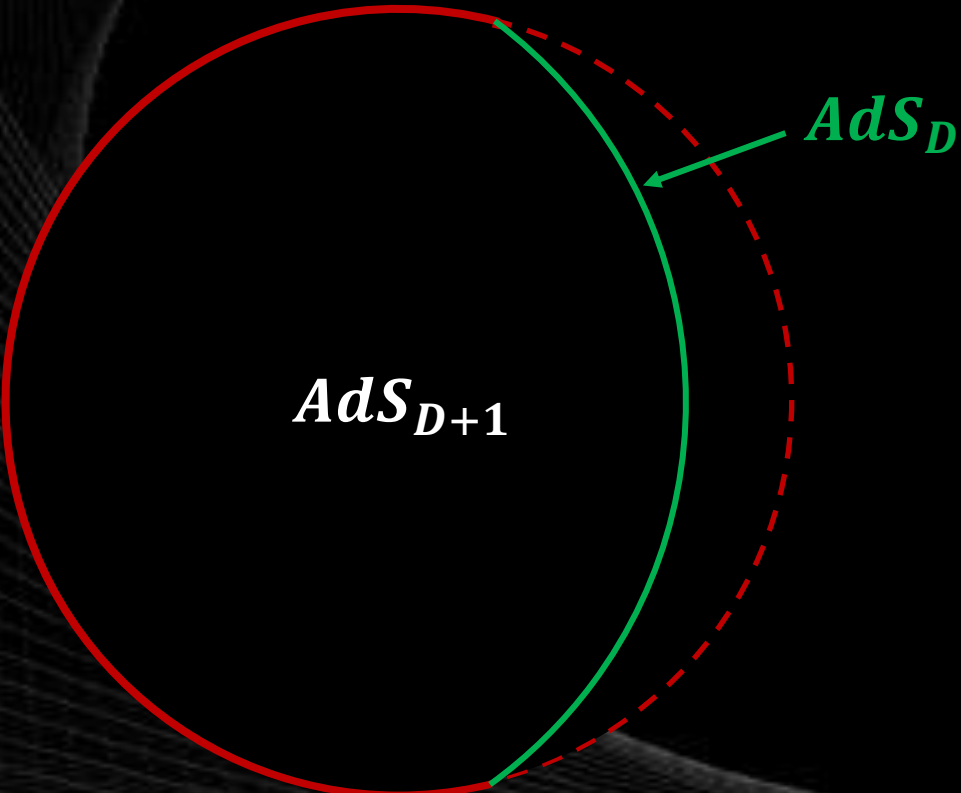


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Einstein + higher curvature terms \longrightarrow Planar CFT with a cutoff

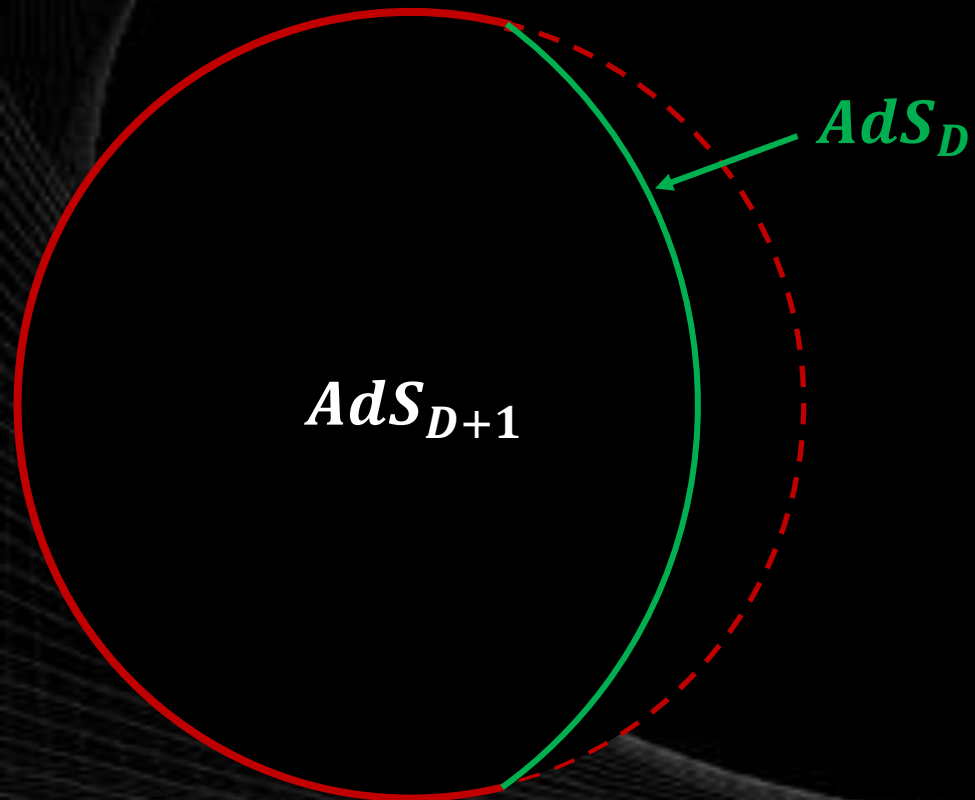
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Planar CFT with a cutoff

The black hole solutions localized on the brane in the AdS_{D+1} braneworld which are found by solving the classical bulk equations in AdS_{D+1} with the brane boundary conditions, correspond to quantum-corrected black holes in D dimensions, rather than classical ones.

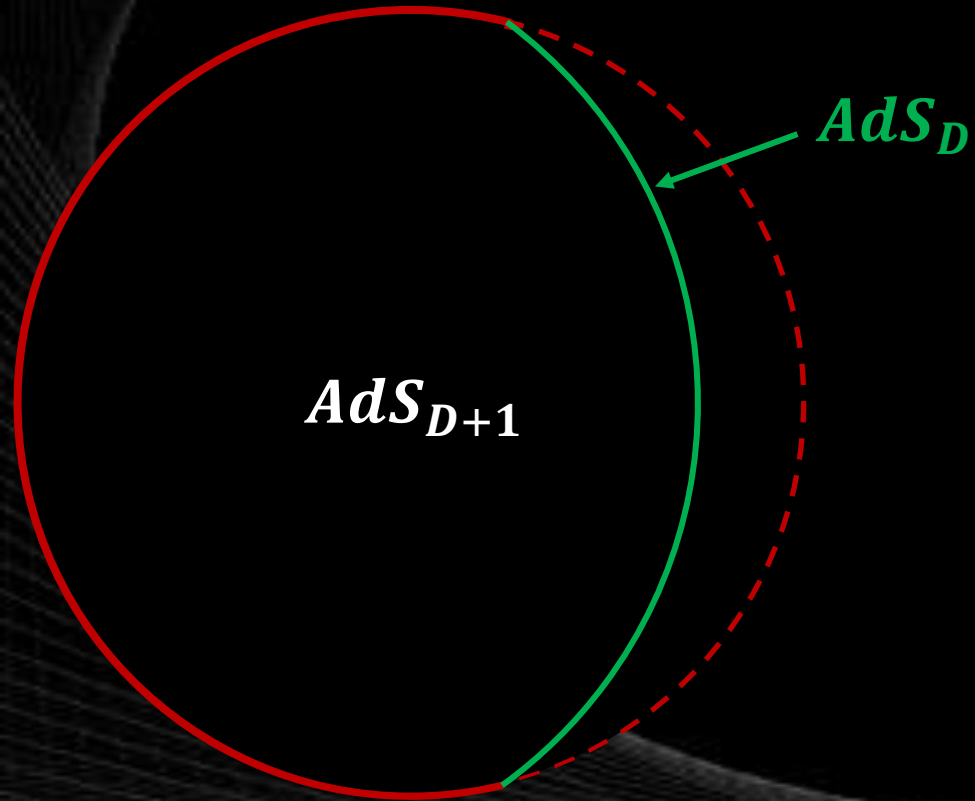
Emparan, Fabbri, Kaloper, JHEP **08** (2002)

Emparan, Horowitz, Myers, JHEP **01** (2000) 007

Emparan, Horowitz, Myers, JHEP **01** (2000) 021

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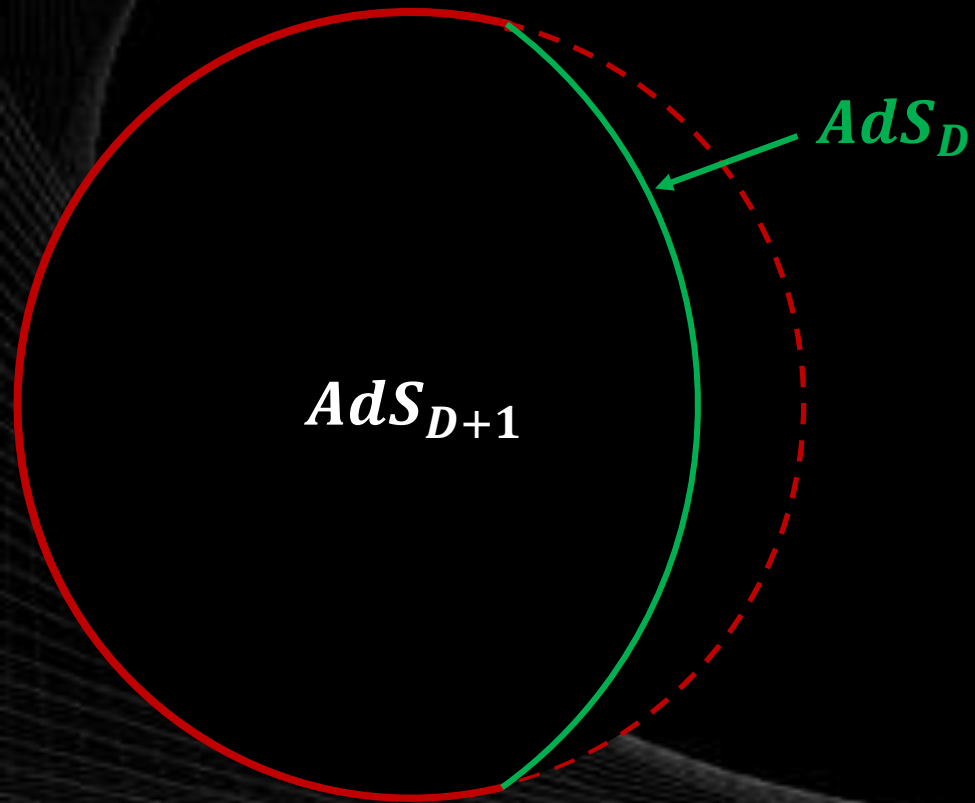
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Consequences

Braneworld Holography

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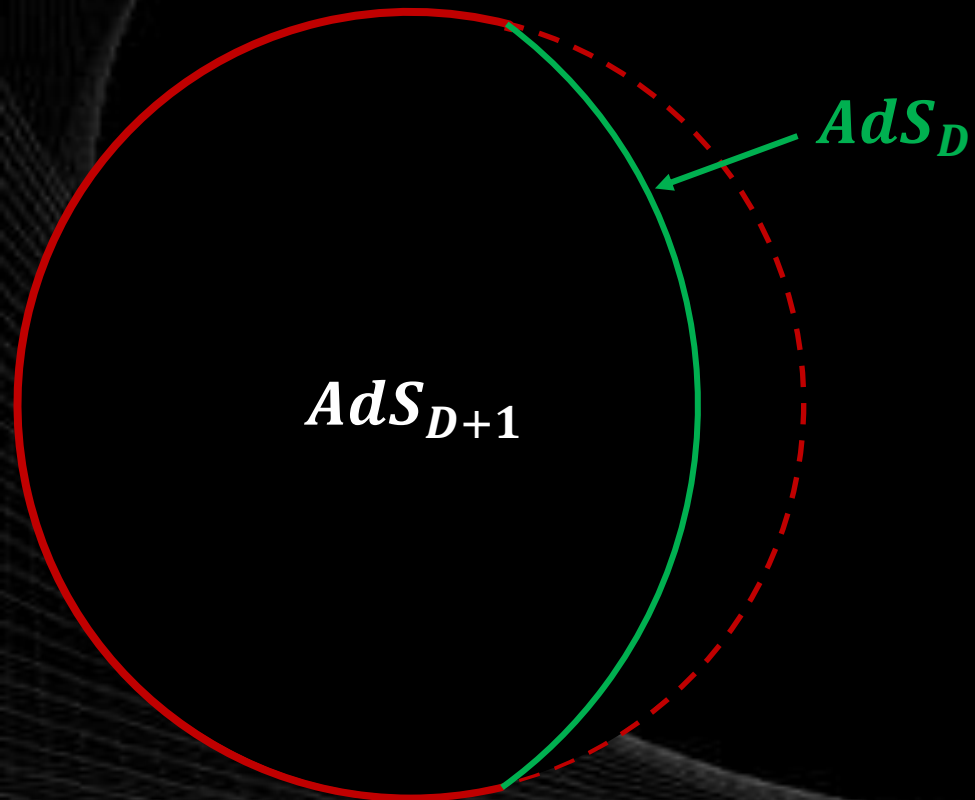
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- Existence of new BH solutions

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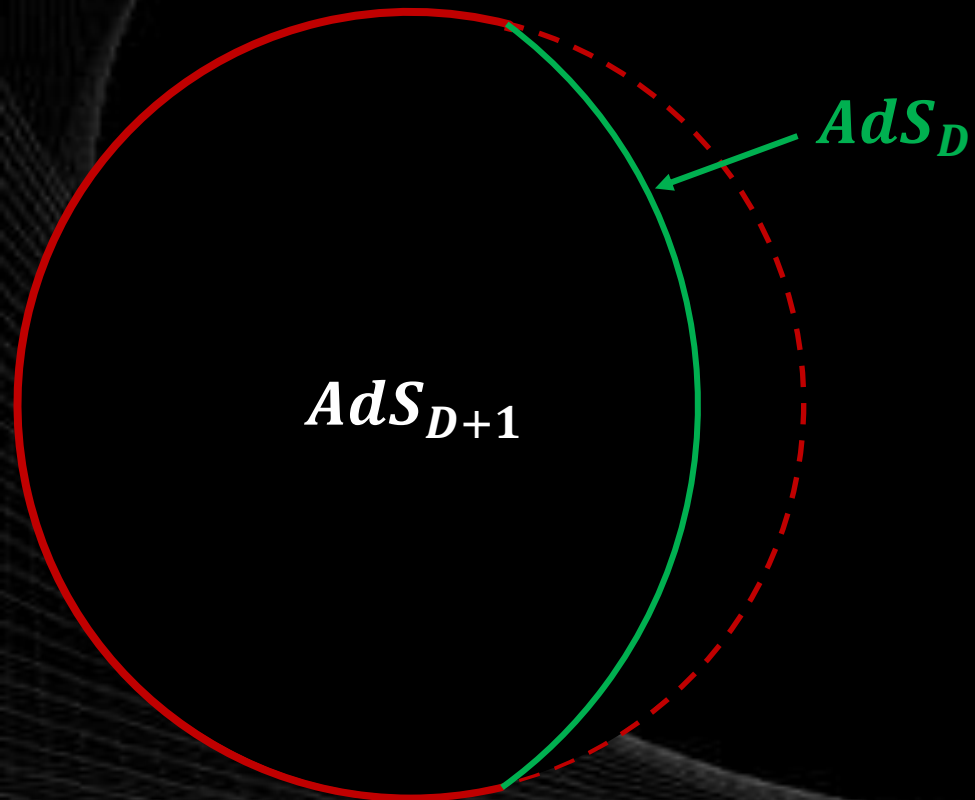
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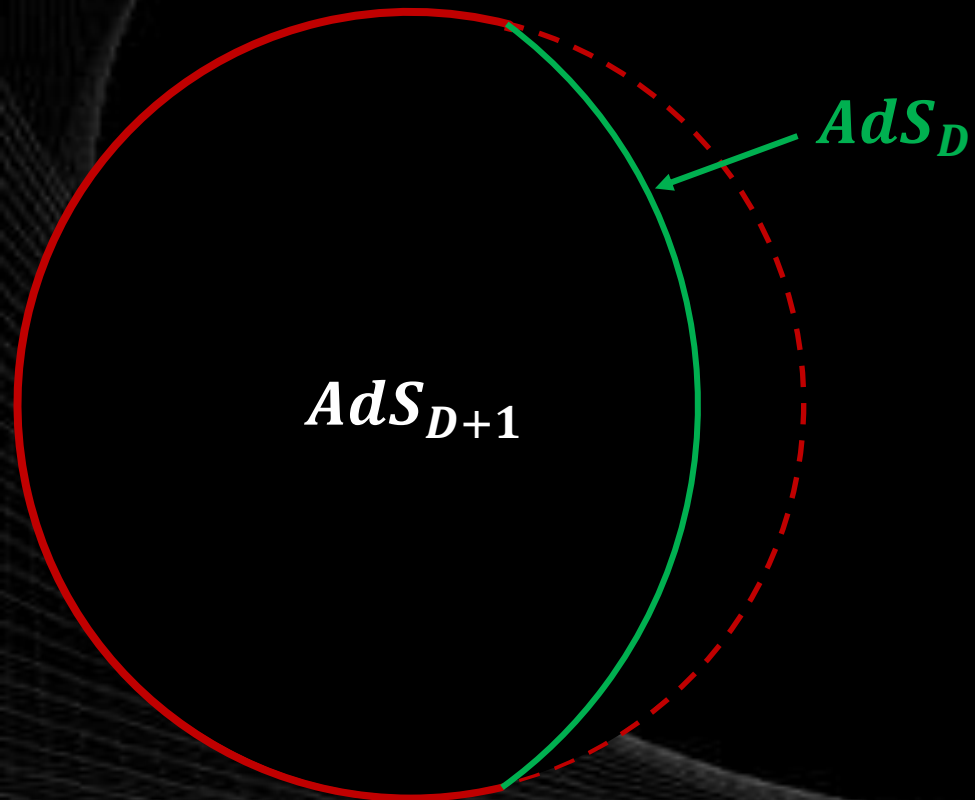
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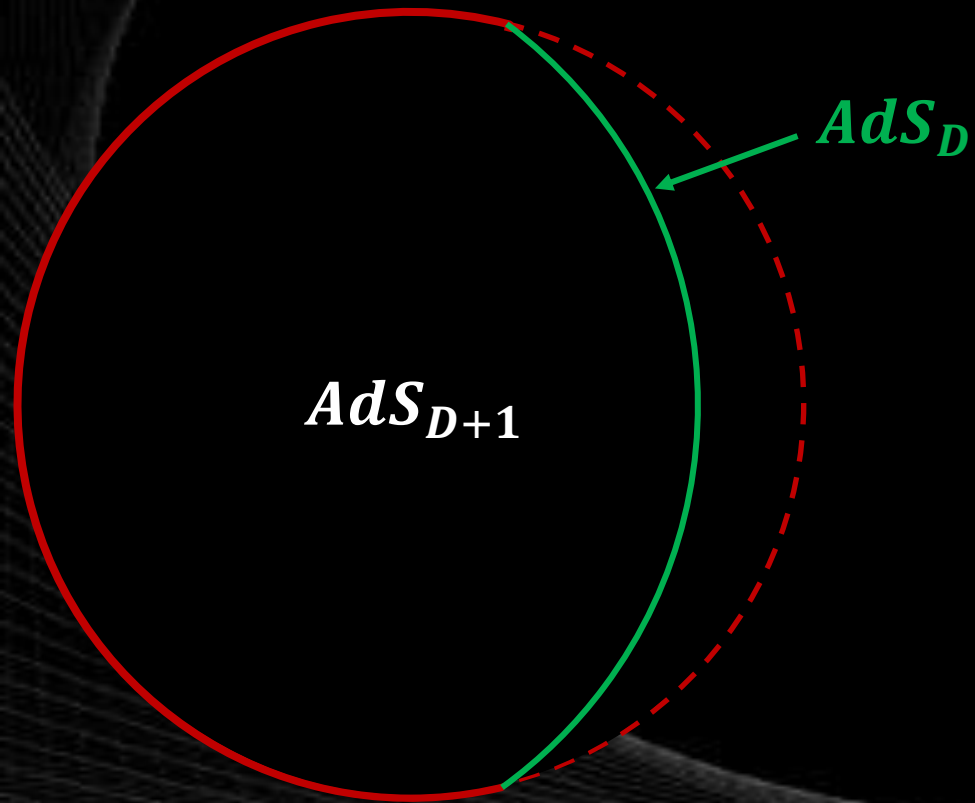
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- Study quantum corrections onto BHs

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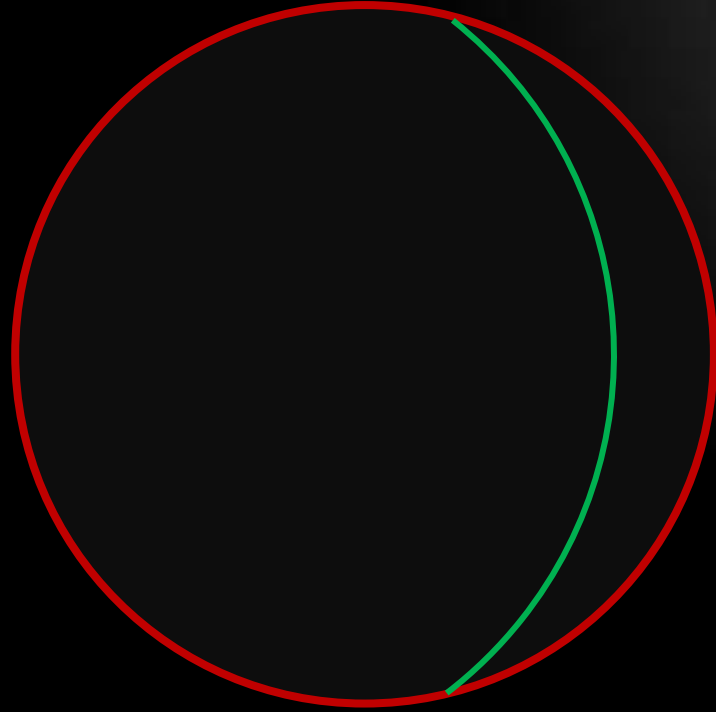
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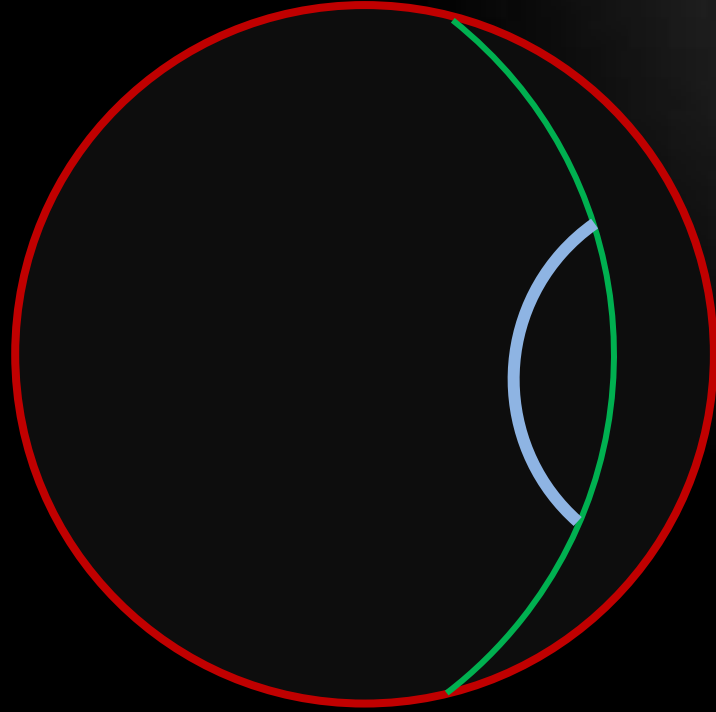
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quantum BTZ

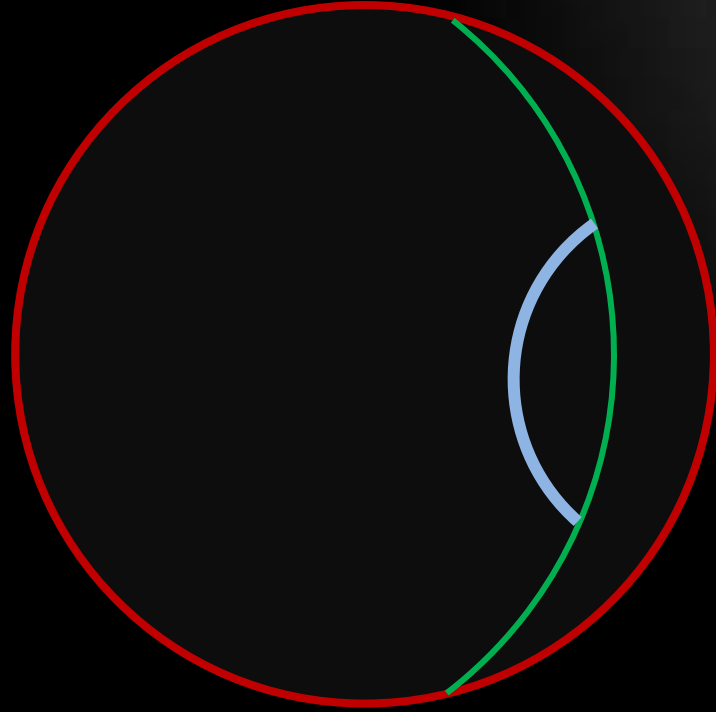
Black holes on branes



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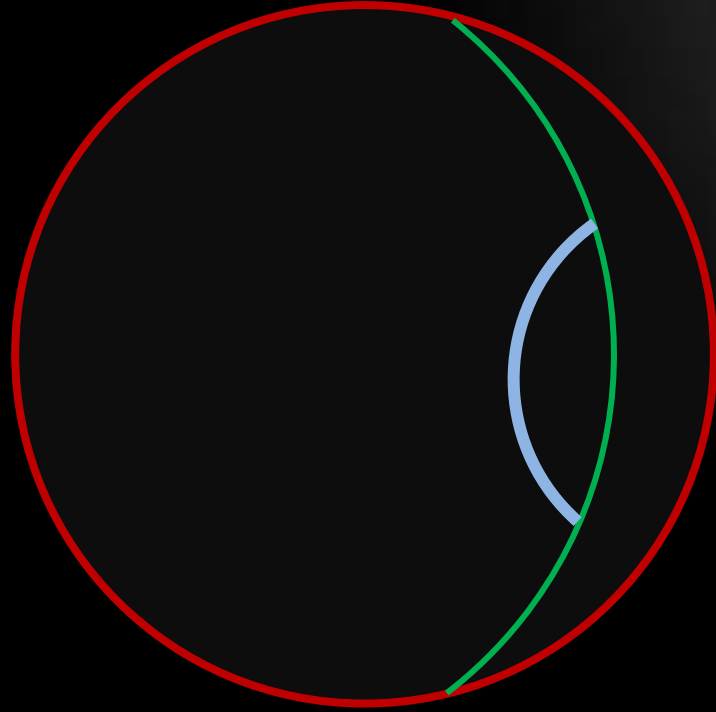


Black holes on branes



A bulk black hole whose horizon intersects the brane is accelerating towards the boundary!

Black holes on branes



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4D **classical** bulk solution: rotating AdS C-metric sliced with a brane Plebanski & Demianski, Annals Phys. **98** (1976)

Rotating qBTZ metric

Empanan, Frassino, Way, JHEP **11** (2020)

$$ds^2 = g_{tt} dt^2 + g_{\phi\phi} d\phi^2 + 2g_{t\phi} dt d\phi + g_{rr} dr^2$$

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$$g_{tt} = -\frac{8\sqrt{1-\tilde{a}^2}\nu\ell_3(\tilde{a}^2 - \kappa x_1^2 + 1)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^3 \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}} + \frac{16\tilde{a}^2 - 4(\tilde{a}^2 + 1)\kappa x_1^2}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} - \frac{r^2}{\ell_3^2}$$

$$g_{\phi\phi} = r^2 - \frac{8\tilde{a}^2\sqrt{1-\tilde{a}^2}\nu\ell_3^3(\tilde{a}^2 - \kappa x_1^2 + 1)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^3 \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}}$$

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Relevant parameters

$$\tilde{a}, x_1, \kappa, \ell_3, \nu$$

$$\kappa = 0, \pm 1$$

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Empanan, Frassino, Way, JHEP **11** (2020)

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Relevant parameters

$$\tilde{a}, x_1, \kappa, \ell_3, \nu$$

$$\kappa = 0, \pm 1$$

$\nu \equiv \ell/\ell_3$ strength of the backreaction

$\nu \rightarrow 0$: classical BTZ limit

Rotating qBTZ metric

Empanan, Frassino, Way, JHEP **11** (2020)

$$ds^2 = g_{tt} dt^2 + g_{\phi\phi} d\phi^2 + 2g_{t\phi} dt d\phi + g_{rr} dr^2$$

$$g_{tt} = -\frac{8\sqrt{1-\tilde{a}^2}\nu\ell_3(\tilde{a}^2 - \kappa x_1^2 + 1)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^3 \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}} + \frac{16\tilde{a}^2 - 4(\tilde{a}^2 + 1)\kappa x_1^2}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} - \frac{r^2}{\ell_3^2}$$

$$g_{\phi\phi} = r^2 - \frac{8\tilde{a}^2\sqrt{1-\tilde{a}^2}\nu\ell_3^3(\tilde{a}^2 - \kappa x_1^2 + 1)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^3 \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}}$$

$$g_{t\phi} = -\frac{4\tilde{a}\ell_3(\tilde{a}^2 - \kappa x_1^2 + 1)}{(3 - \tilde{a}^2 - \kappa x_1^2)^2} \left(1 + \frac{2\sqrt{1-\tilde{a}^2}\nu\ell_3}{(3 - \tilde{a}^2 - \kappa x_1^2) \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}} \right)$$

$$g^{rr} = \frac{r^2}{\ell_3^2} - \frac{8(1-\tilde{a}^2)^{3/2}\nu\ell_3(\tilde{a}^2 - \kappa x_1^2 + 1) \sqrt{\frac{4\tilde{a}^2\ell_3^2(\kappa x_1^2 - 2)}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2} + r^2}}{r^2(3 - \tilde{a}^2 - \kappa x_1^2)^3} + \frac{16\tilde{a}^2\ell_3^2(\tilde{a}^2 - \kappa x_1^2 + 1)^2}{r^2(\tilde{a}^2 + \kappa x_1^2 - 3)^4} + \frac{4[(\tilde{a}^2 + 1)\kappa x_1^2 - 4\tilde{a}^2]}{(\tilde{a}^2 + \kappa x_1^2 - 3)^2}$$

Relevant parameters

$$\tilde{a}, x_1, \kappa, \ell_3, \nu$$

$$\kappa = 0, \pm 1$$

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$\nu \rightarrow 0$: classical BTZ limit

Charges

$$M = \frac{1}{2\mathcal{G}_3} \frac{-\kappa x_1^2 + \tilde{a}^2(4 - \kappa x_1^2)}{(3 - \kappa x_1^2 - \tilde{a}^2)^2}$$

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Empanan, Frassino, Way, JHEP **11** (2020)

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Relevant parameters

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$$\mathcal{G}_3 \propto G_4/\ell$$

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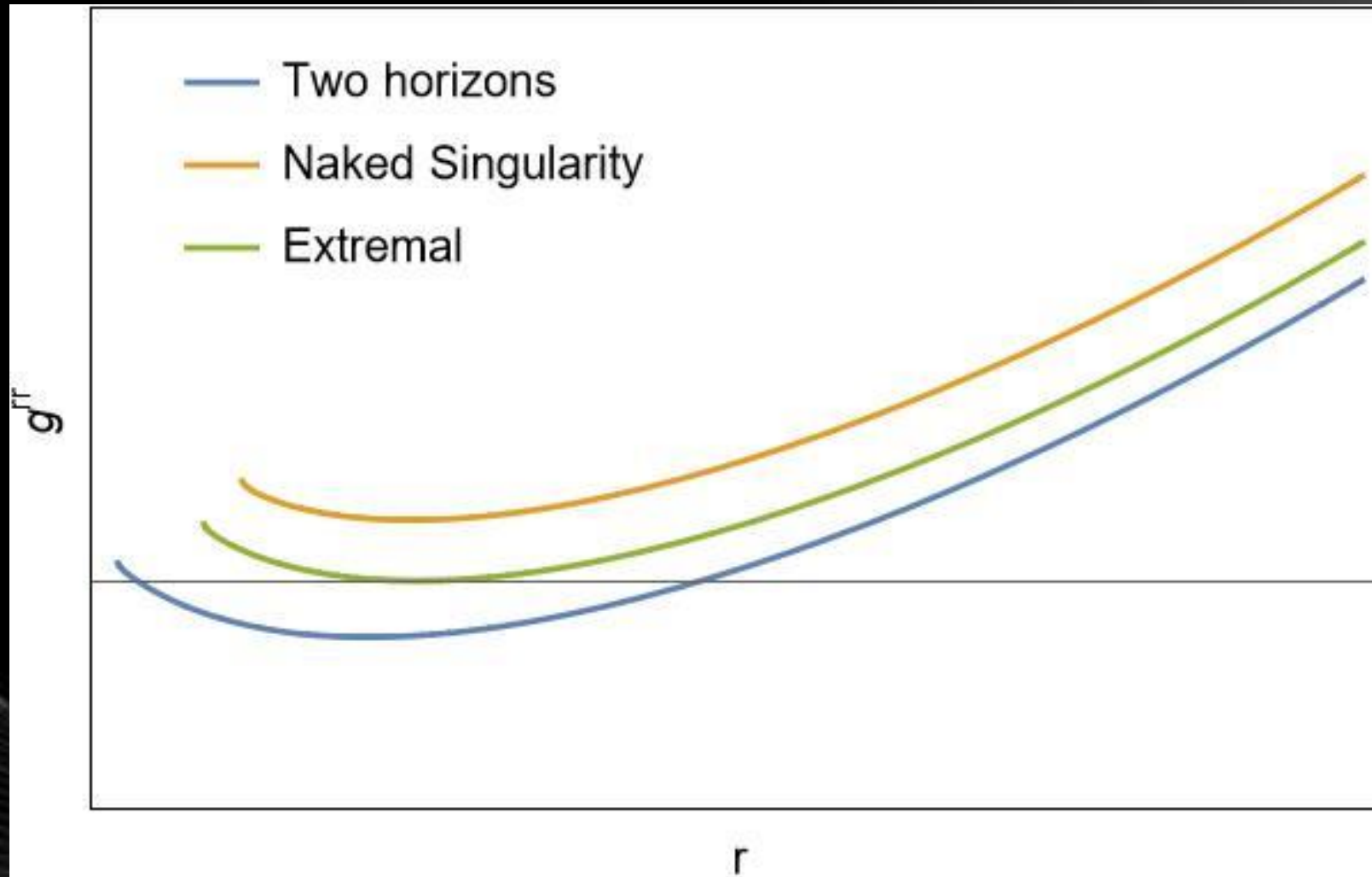
Charges

$$M = \frac{1}{\mathcal{G}_3} \frac{-\kappa x_1^2 + \tilde{a}^2(4 - \kappa x_1^2)}{(3 - \kappa x_1^2 - \tilde{a}^2)^2}$$

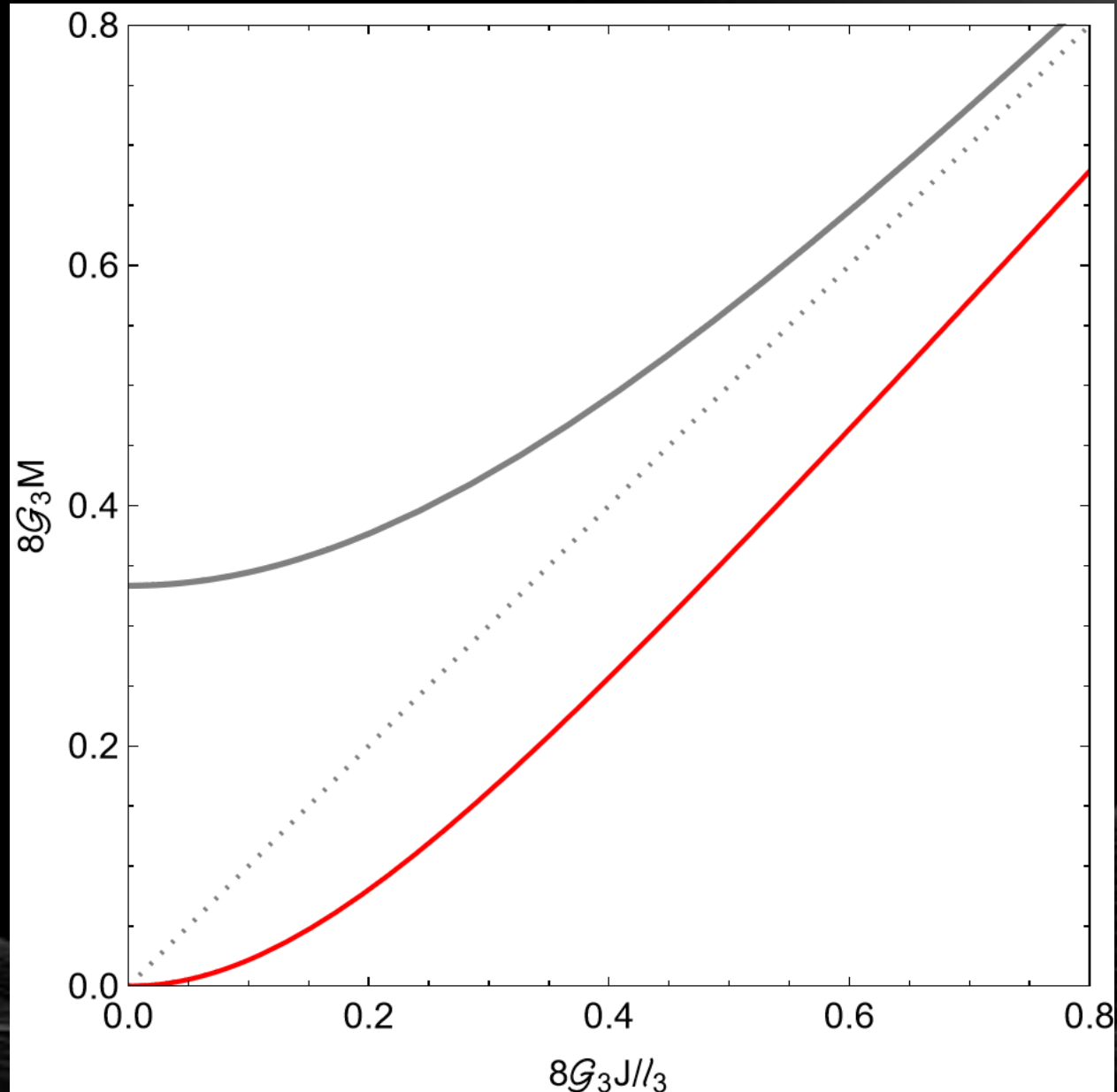
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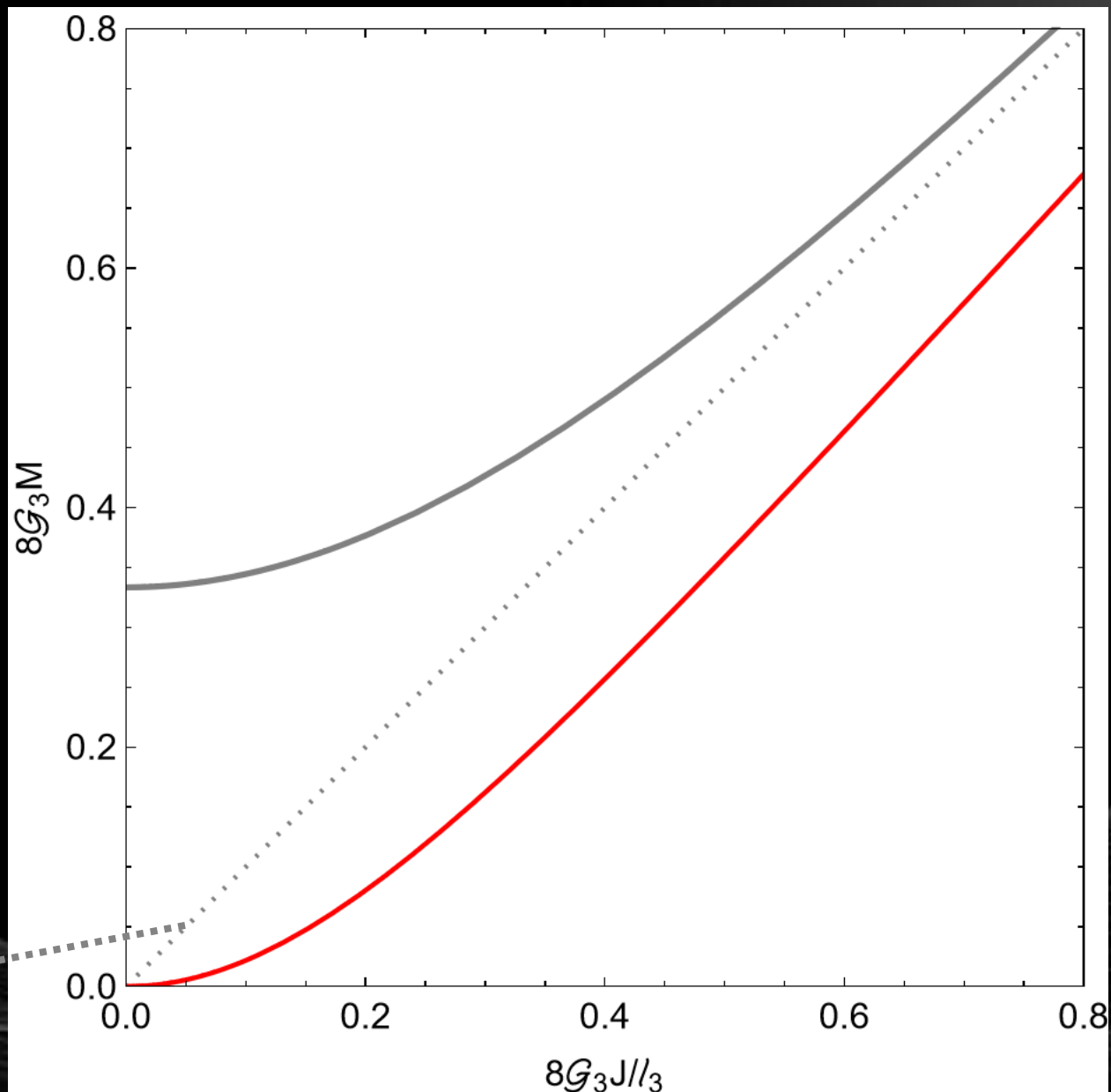
Metric structure



M vs J diagram

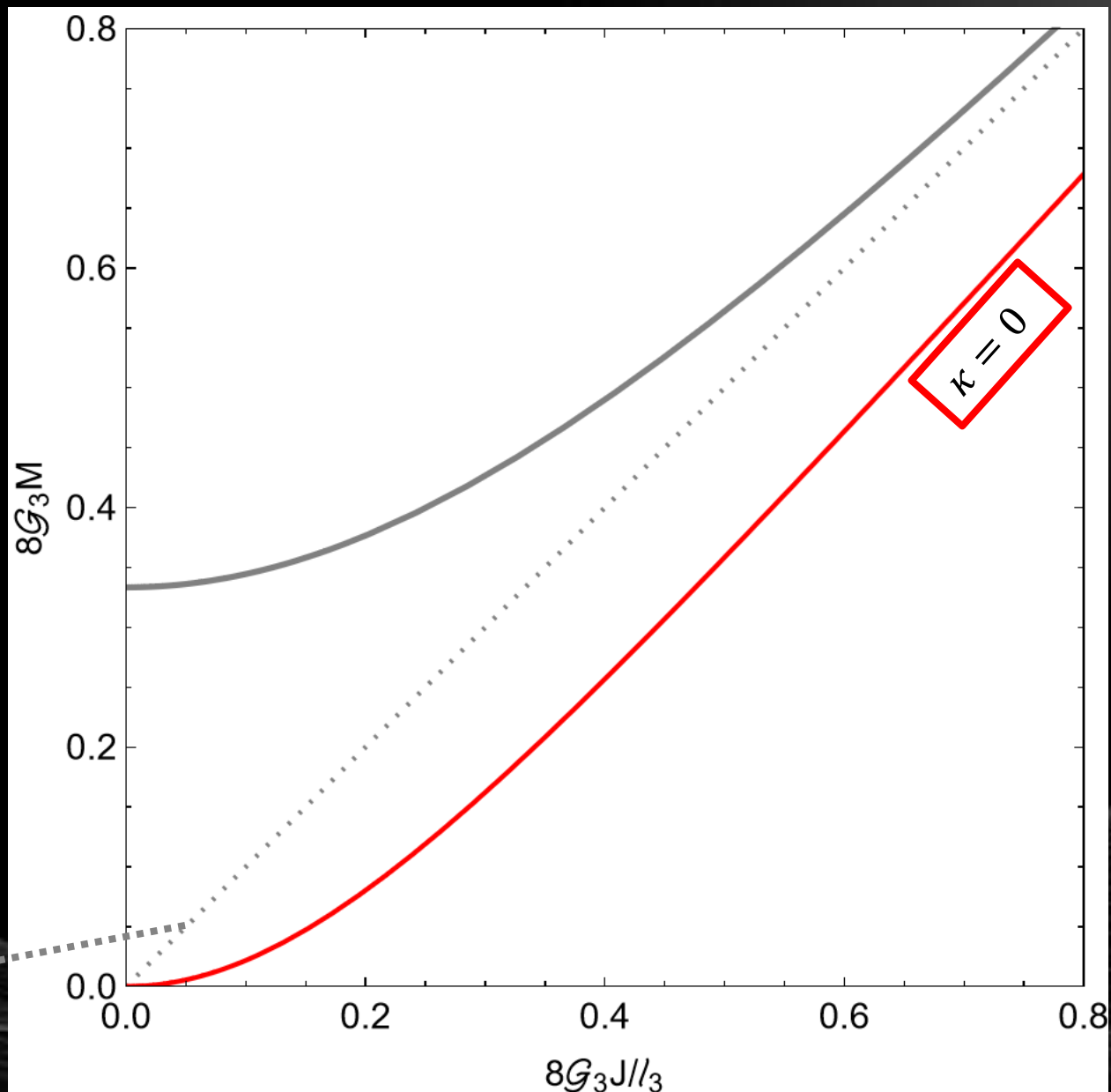


M vs J diagram



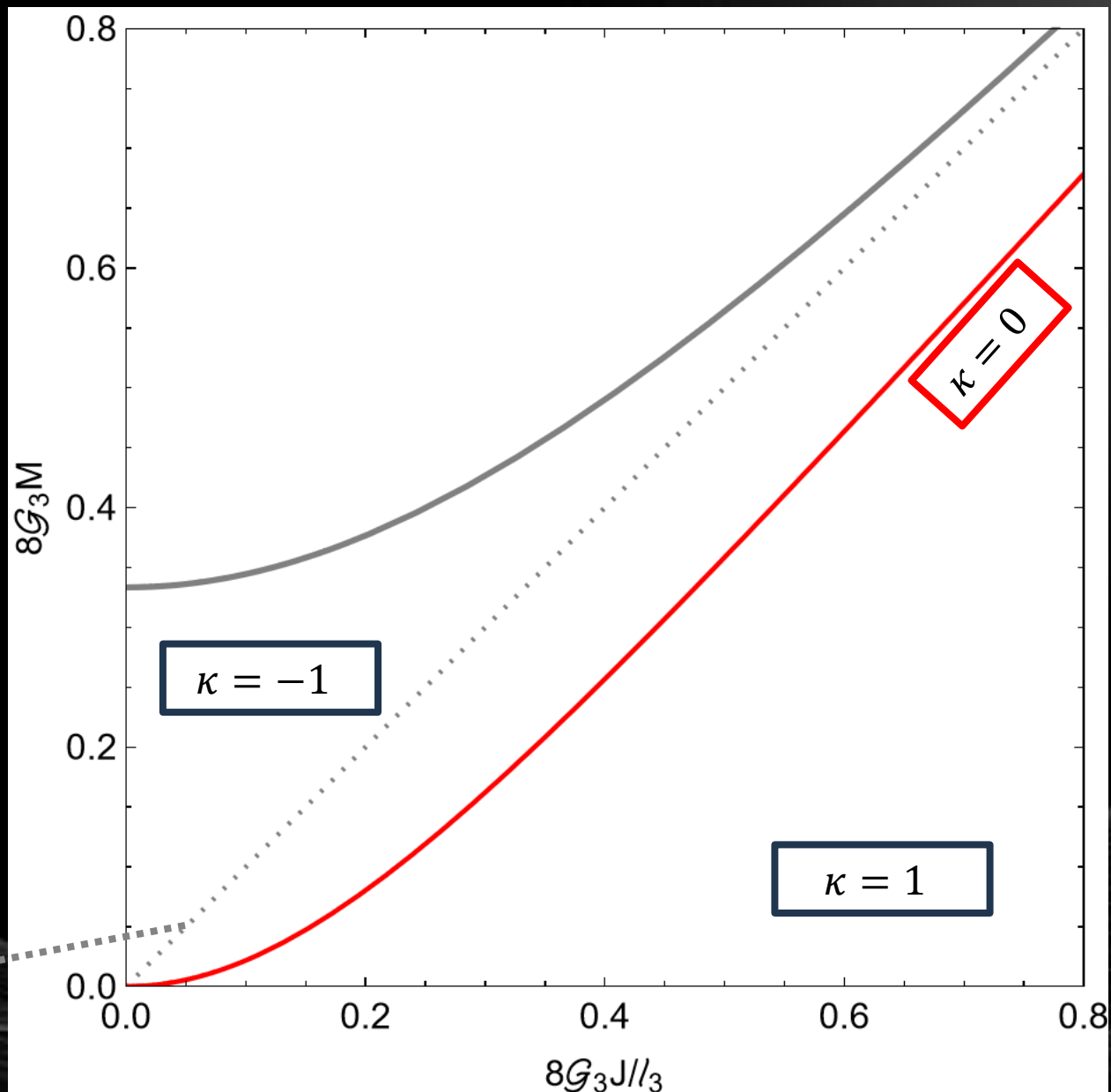
Classical BTZ extremality

M vs J diagram



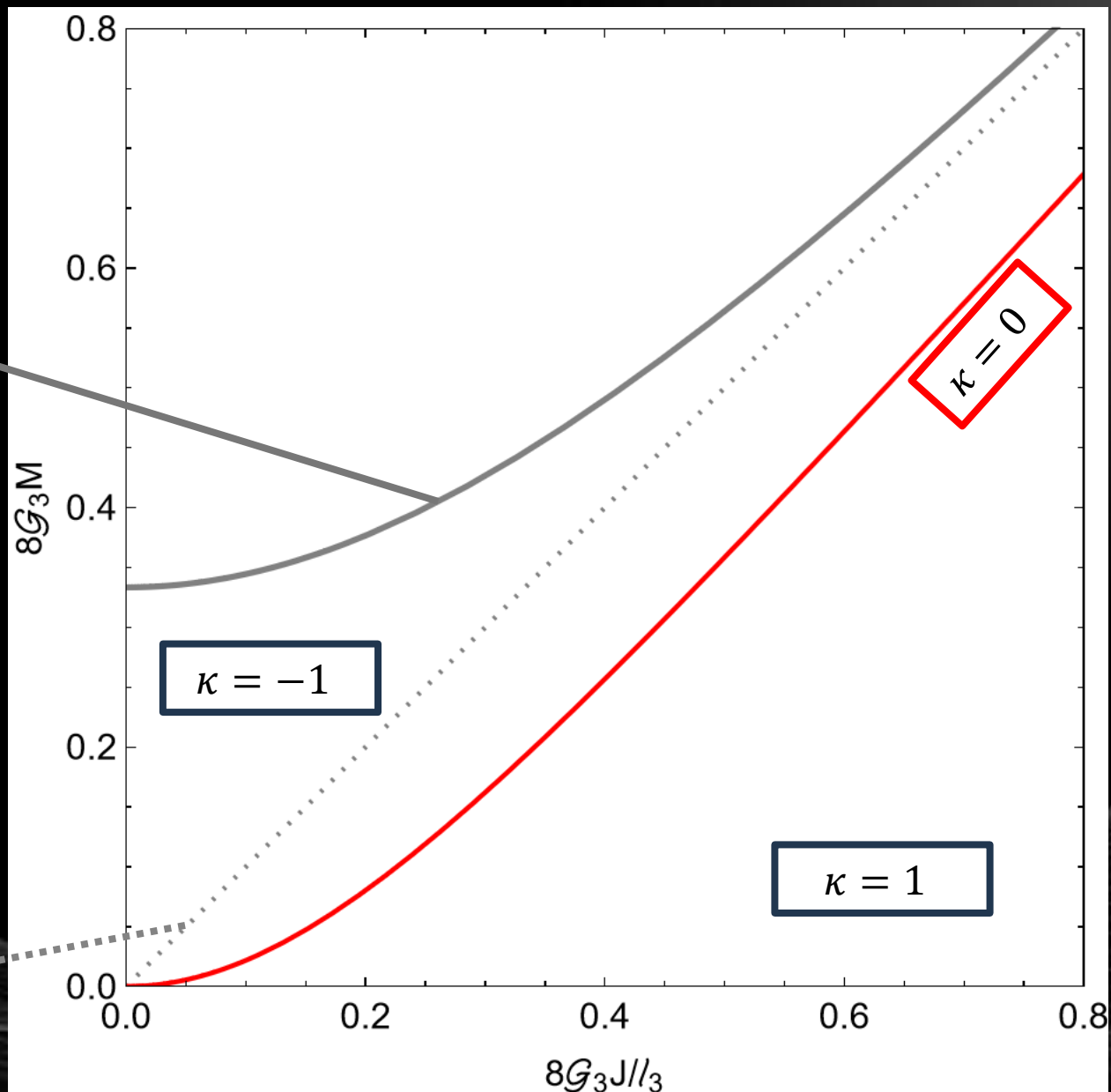
Classical BTZ extremality

M vs J diagram



Classical BTZ extremality

M vs J diagram



Maximum Mass

From holographic construction

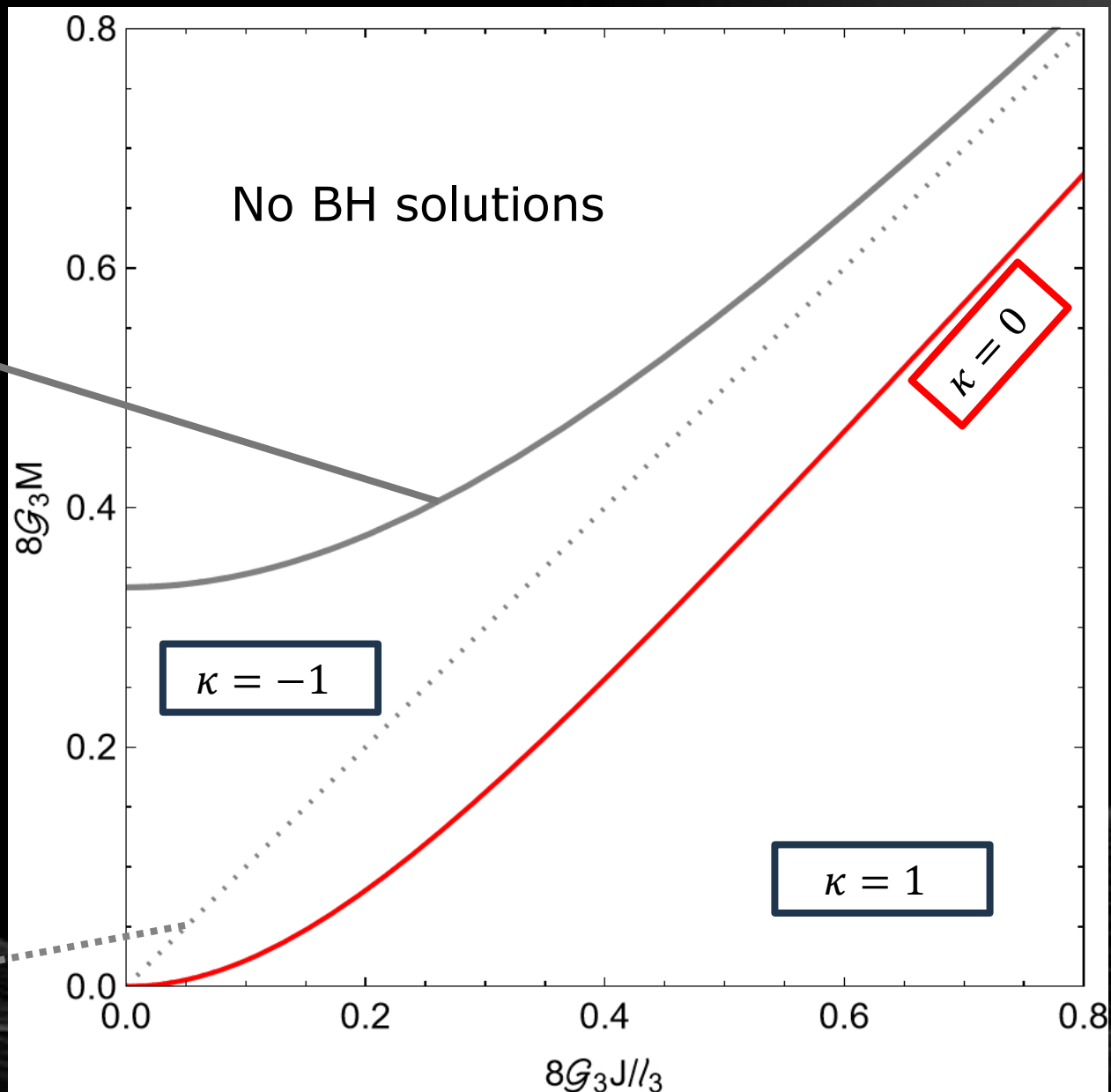
$\kappa = -1$

$\kappa = 0$

$\kappa = 1$

Classical BTZ extremality

M vs J diagram



Maximum Mass

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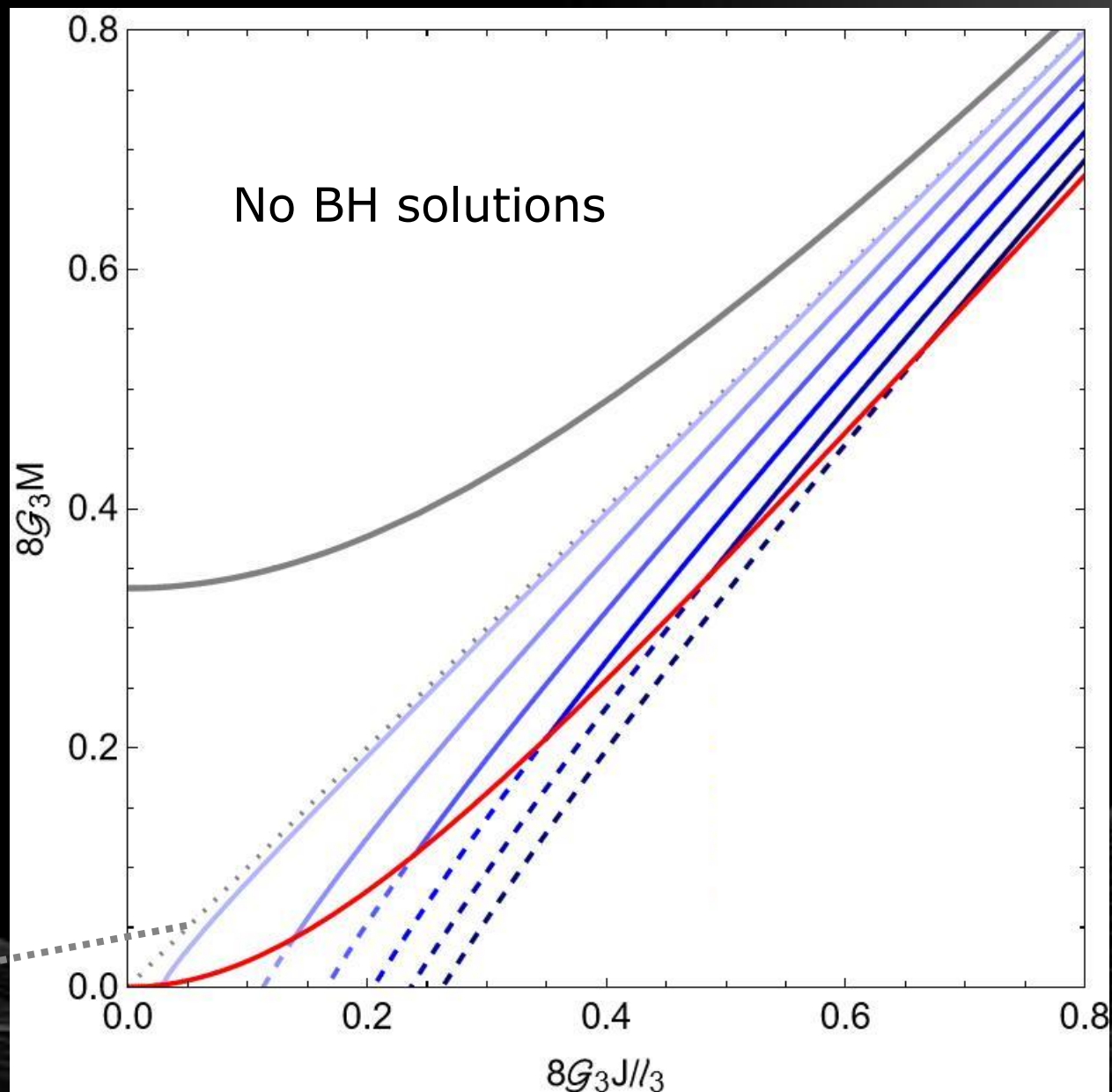
$\kappa = -1$

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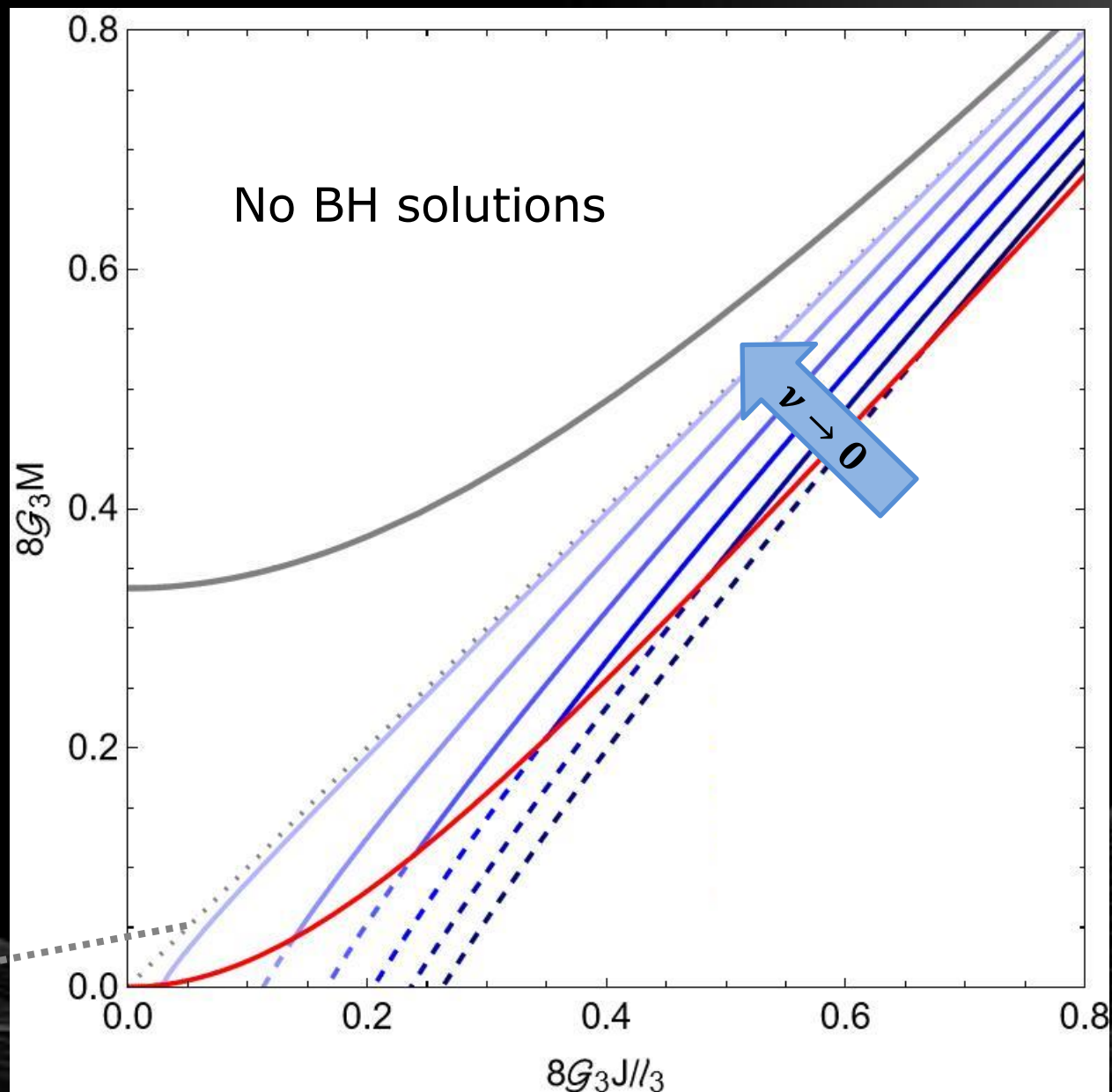
Classical BTZ extremality

M vs J diagram



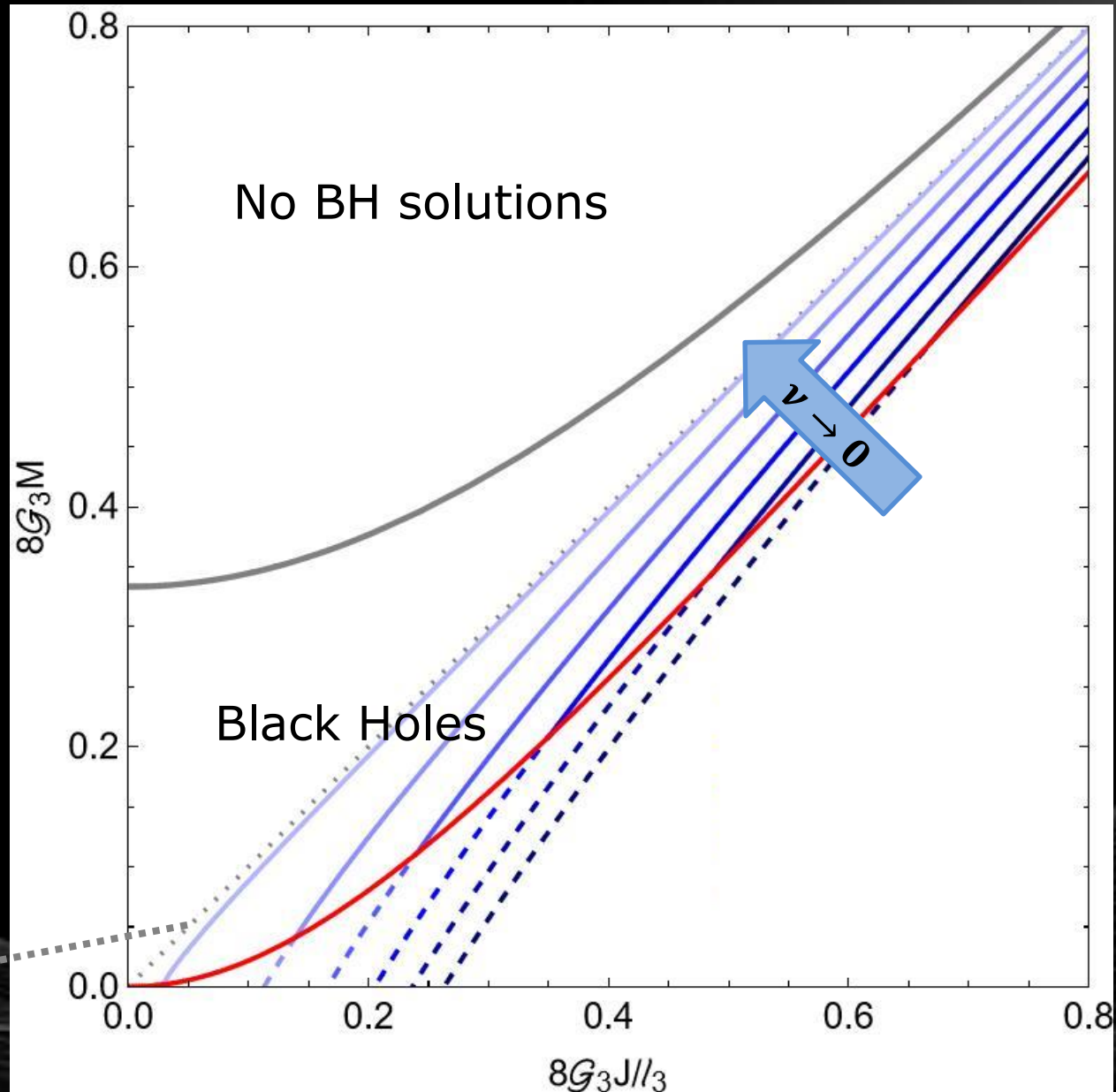
Classical BTZ extremality

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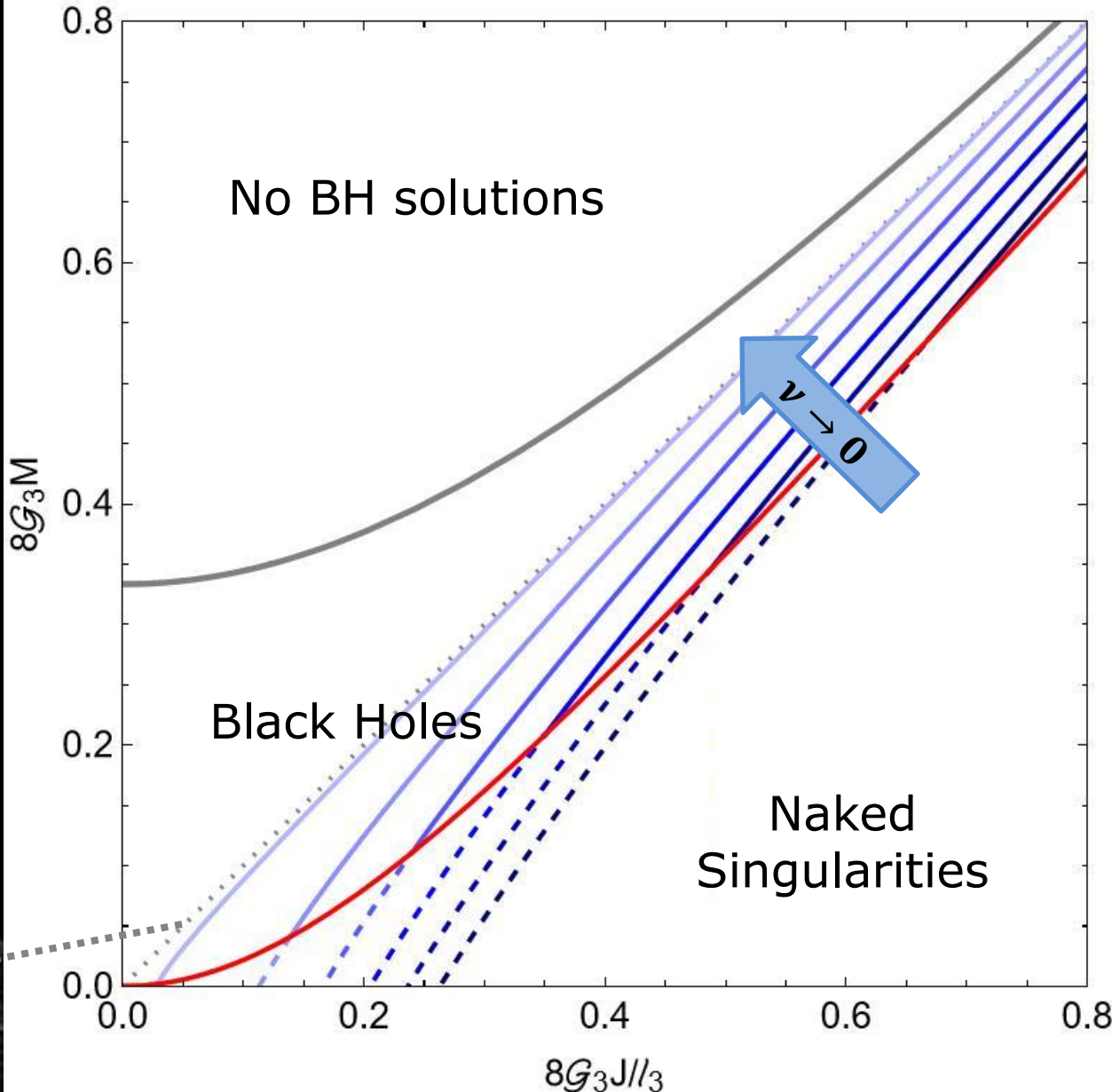
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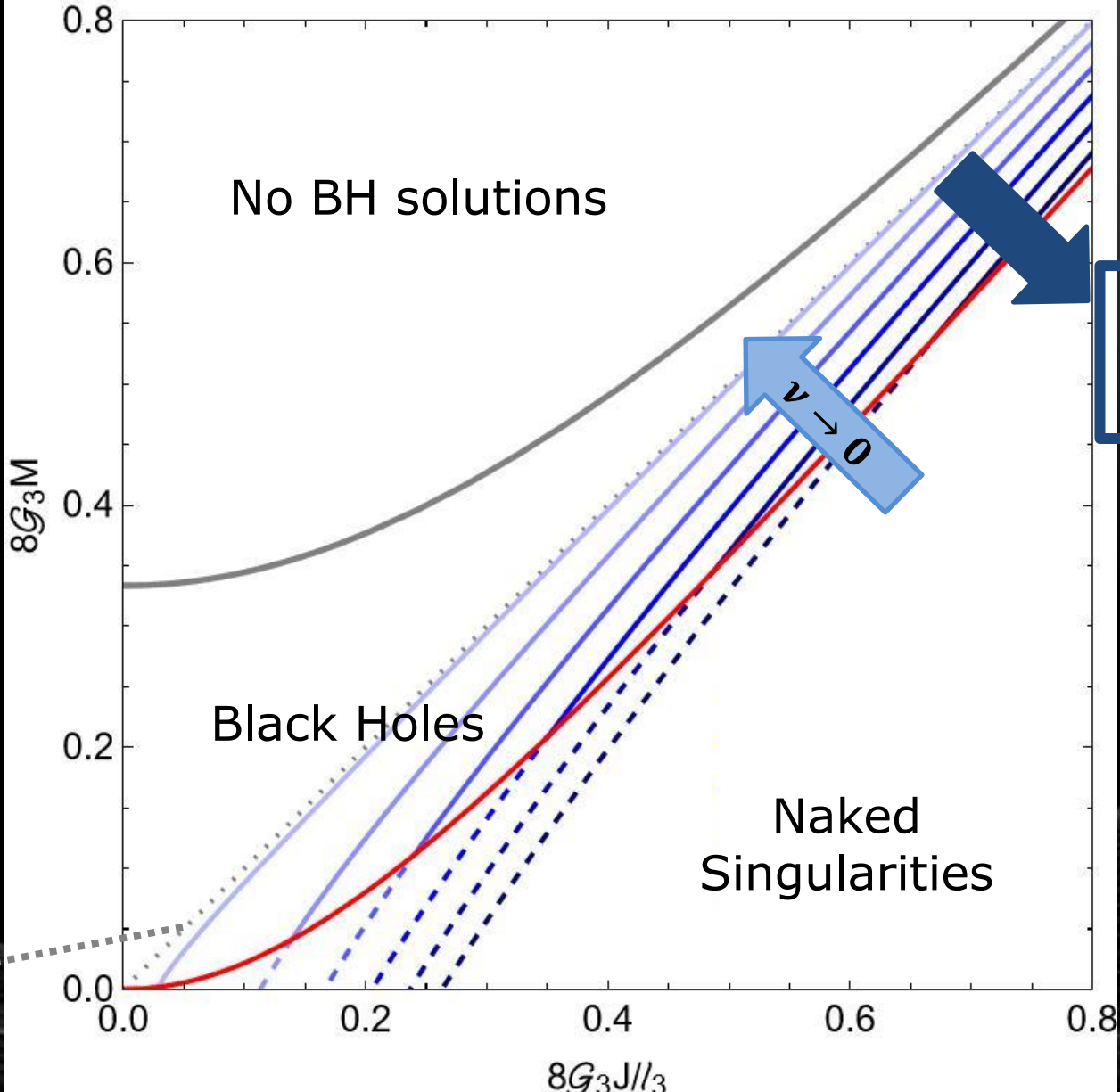
Classical BTZ extremality

M vs J diagram



Classical BTZ extremality

M vs J diagram



Classical BTZ extremality

BH solution parameter space grows with ν

wCCC tests w/ qBTZ

qBTZ metric: tests of wCCC

Main Idea

qBTZ metric: tests of wCCC

Main Idea

Start with the extremal configuration

qBTZ metric: tests of wCCC

Main Idea

Start with the extremal configuration



Particle Absorption

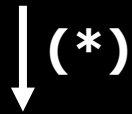
qBTZ metric: tests of wCCC

Main Idea

Start with the extremal configuration



Particle Absorption



Check the sign of g^{rr} at the minimum of the final state

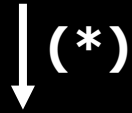
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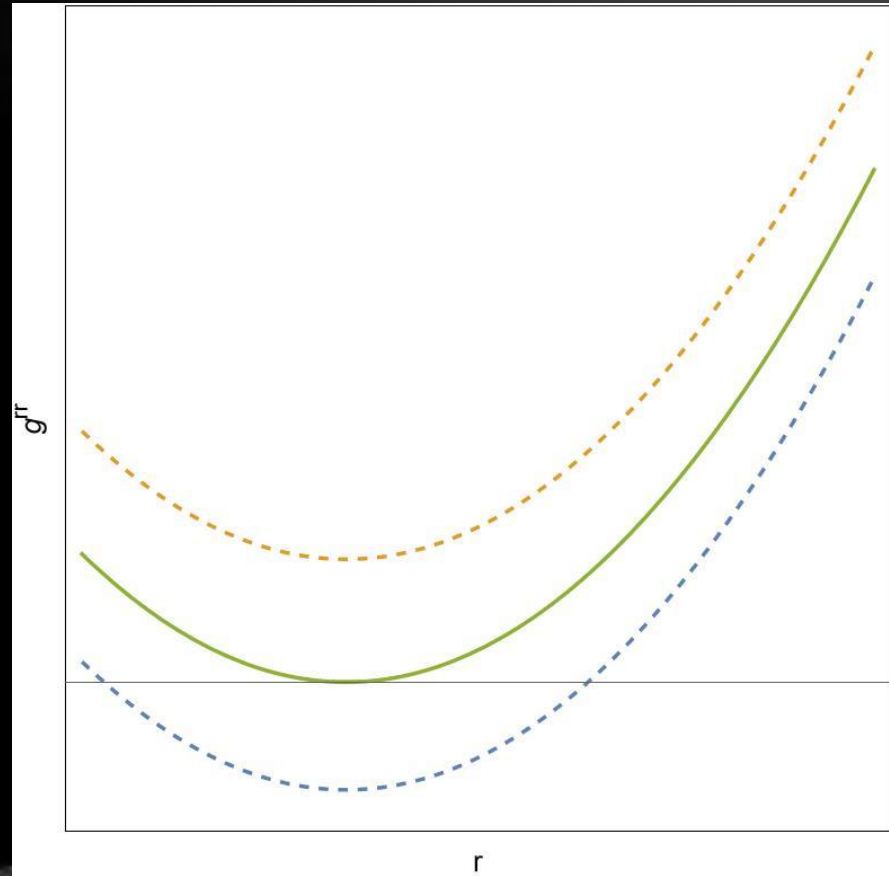
Start with the extremal configuration



Particle Absorption



Check the sign of g^{rr} at the minimum of the final state



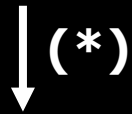
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Main Idea

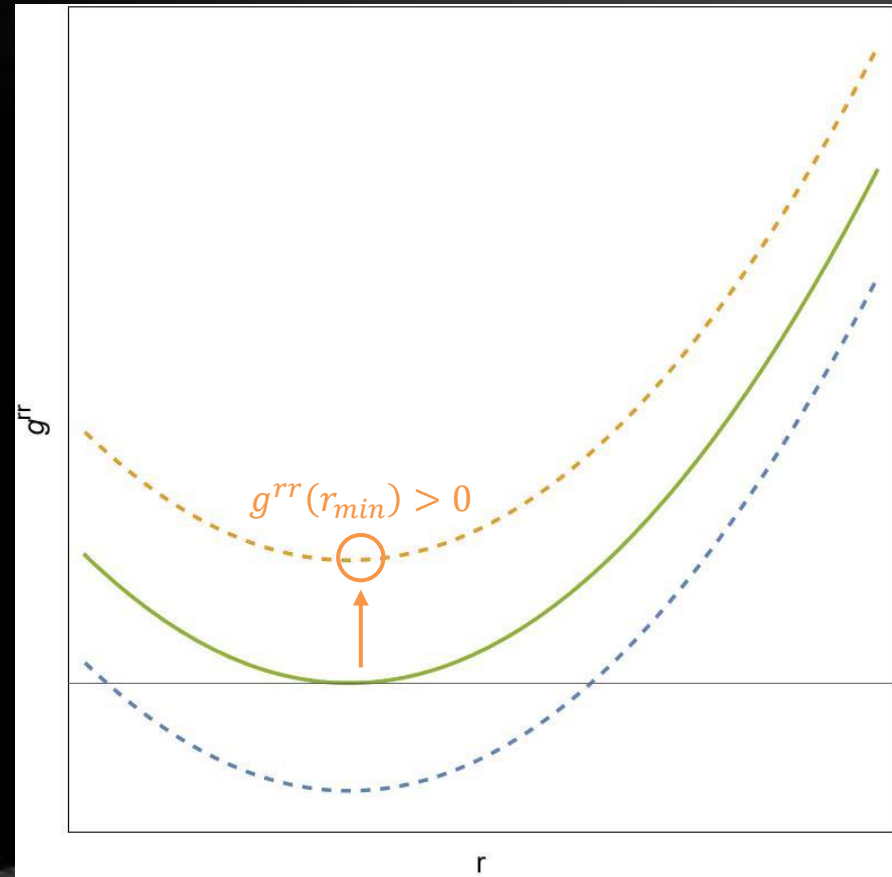
Start with the extremal configuration



Particle Absorption



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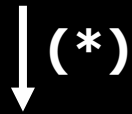
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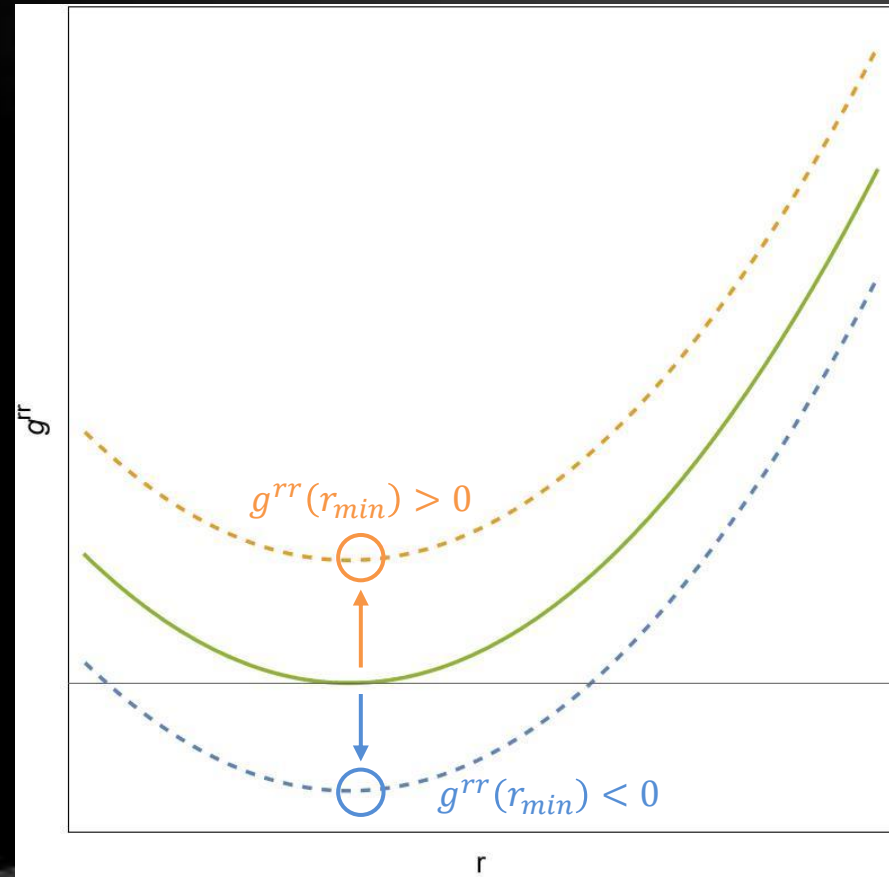
Start with the extremal configuration



Particle Absorption



Check the sign of g^{rr} at the minimum of the final state



Perturbative analysis

Perturbative analysis

Linear response after absorption of a test particle

Perturbative analysis

Linear response after absorption of a test particle

Strategy

Perturbative analysis

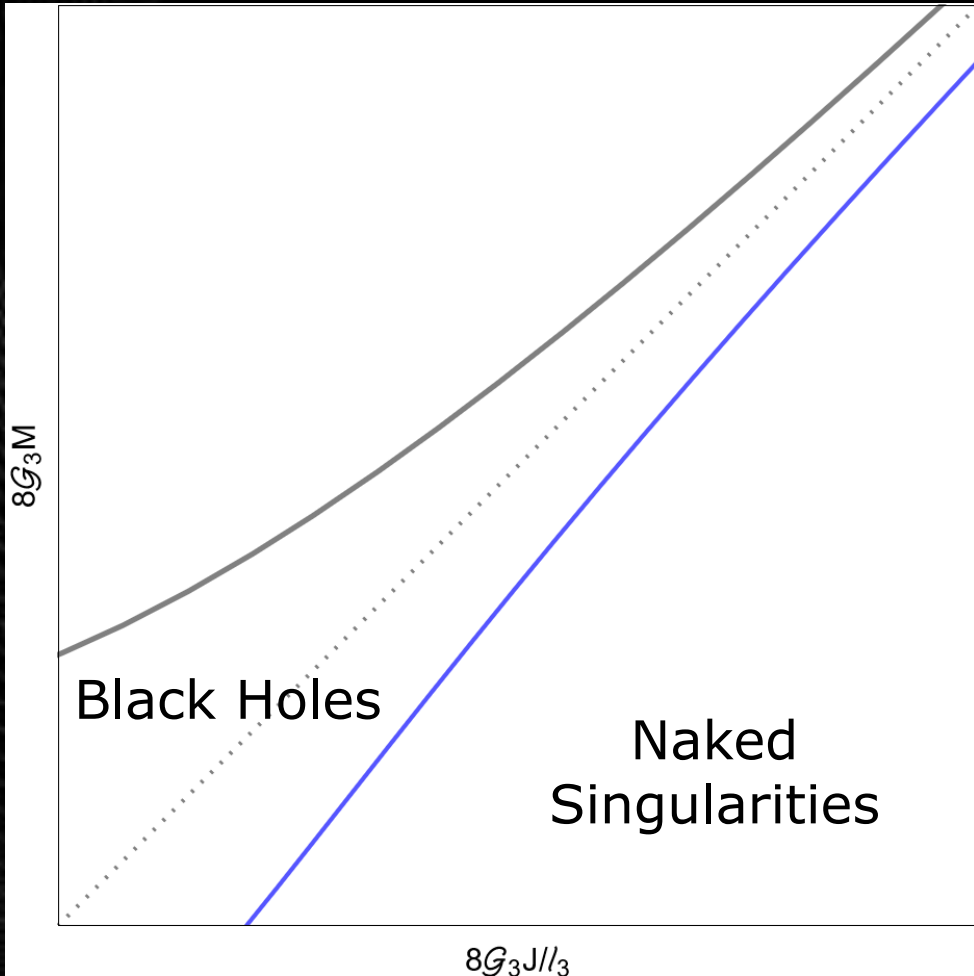
Linear response after absorption of a test particle

Strategy

- Choose an initial extremal configuration: fix ℓ_3 , κ , x_1 and \tilde{a}

Perturbative analysis

Linear response after absorption of a test particle



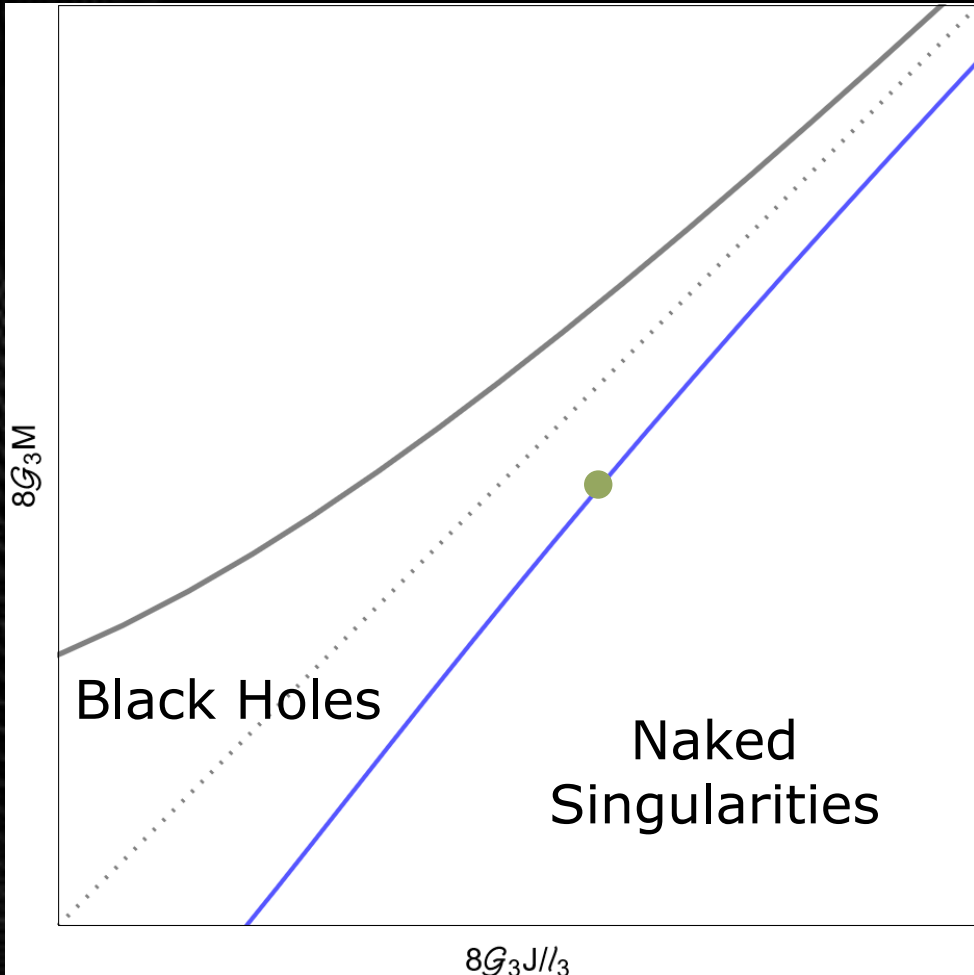
Strategy

- Choose an initial extremal configuration: fix ℓ_3 , κ , x_1 and \tilde{a}

$$v = v_{ext}(\ell_3, x_1, \tilde{a})$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

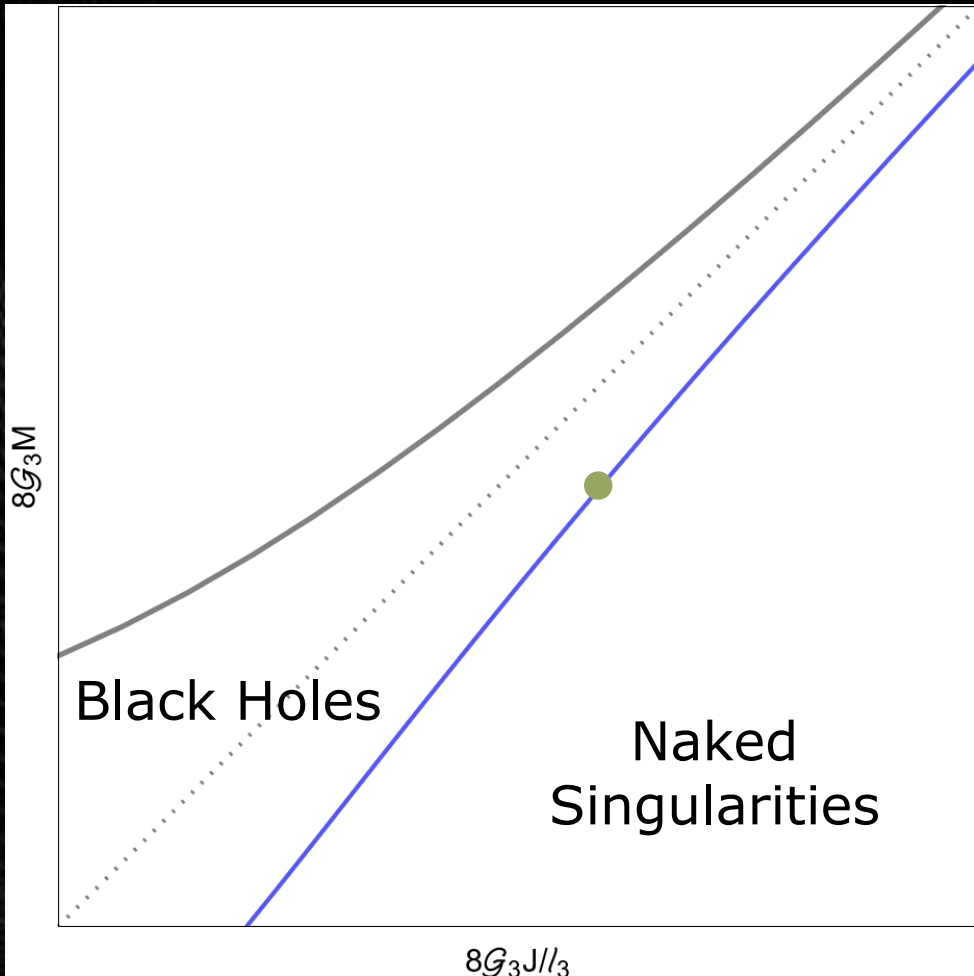
- Choose an initial extremal configuration: fix ℓ_3 , κ , x_1 and \tilde{a}

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- Initial values of M and J fixed

Perturbative analysis

Linear response after absorption of a test particle



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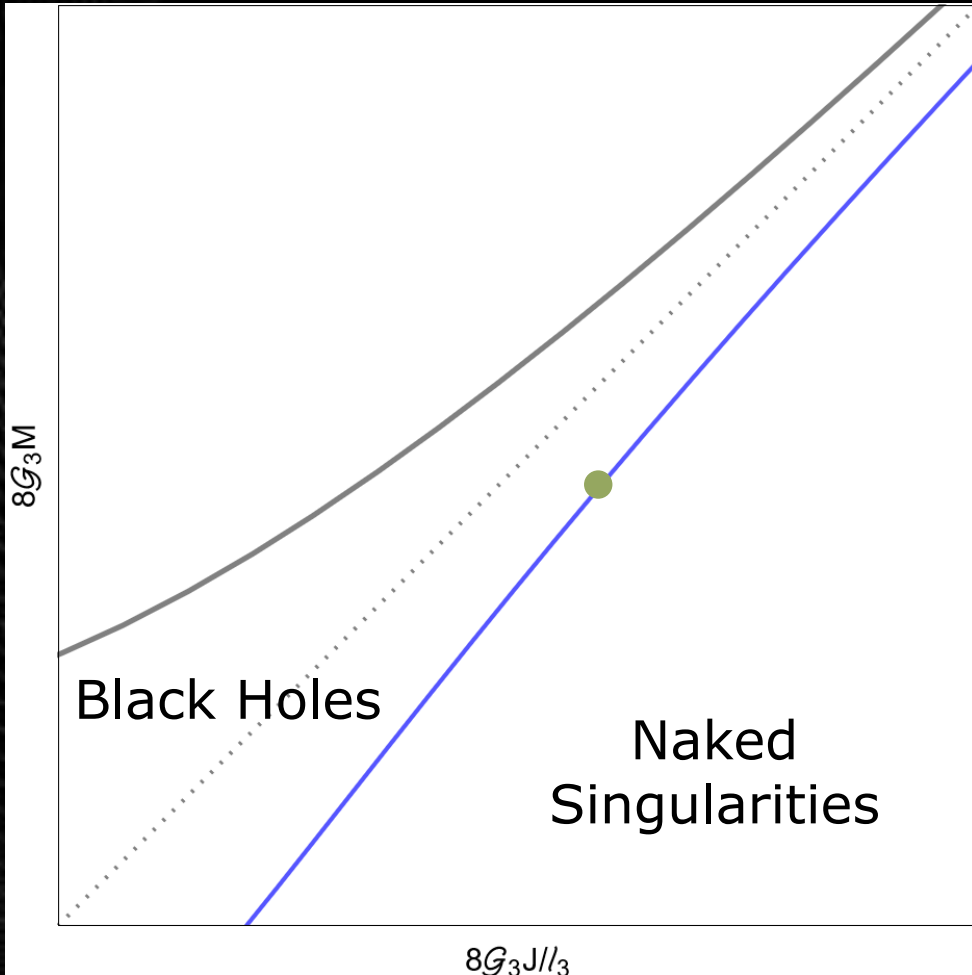
- Initial values of M and J fixed

- Particle absorption changes the charges

$$\begin{aligned} M &\rightarrow M + \delta M \\ J &\rightarrow J + \delta J \end{aligned}$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

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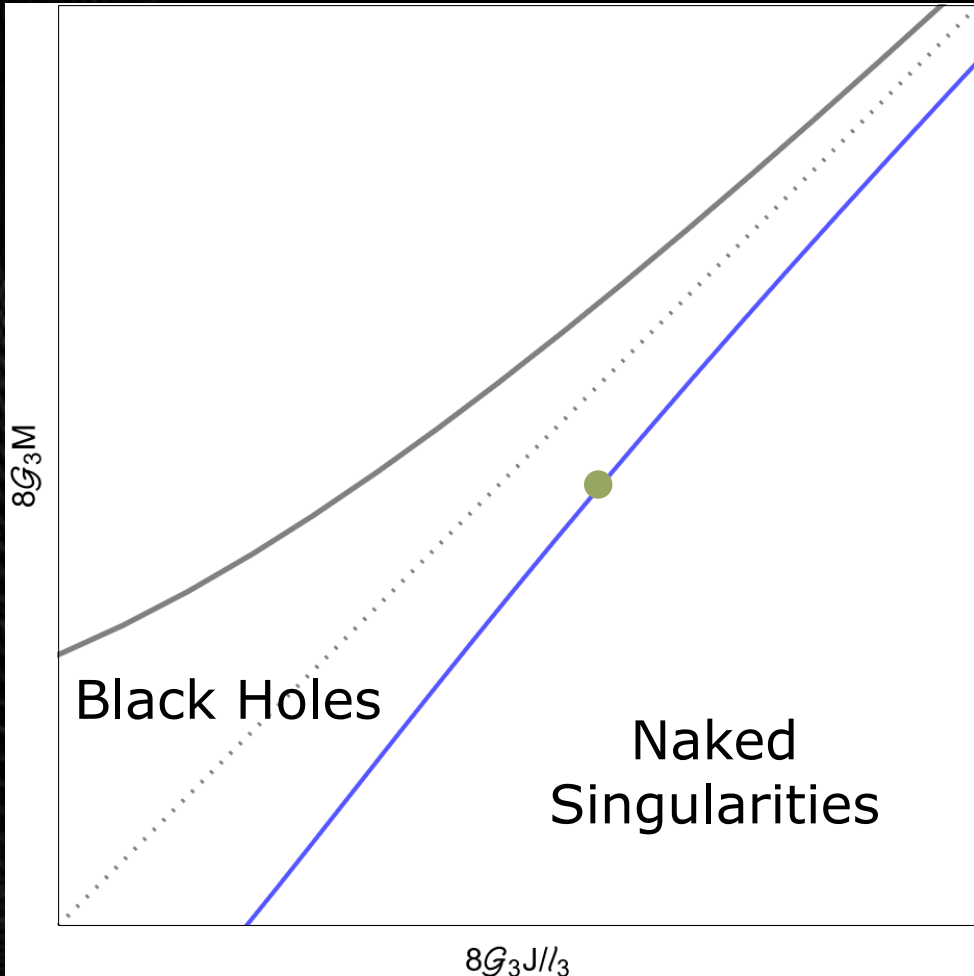
- Particle absorption changes the charges

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- $(\delta M, \delta J) \rightarrow (\delta \tilde{a}, \delta x_1)$

Perturbative analysis

Linear response after absorption of a test particle

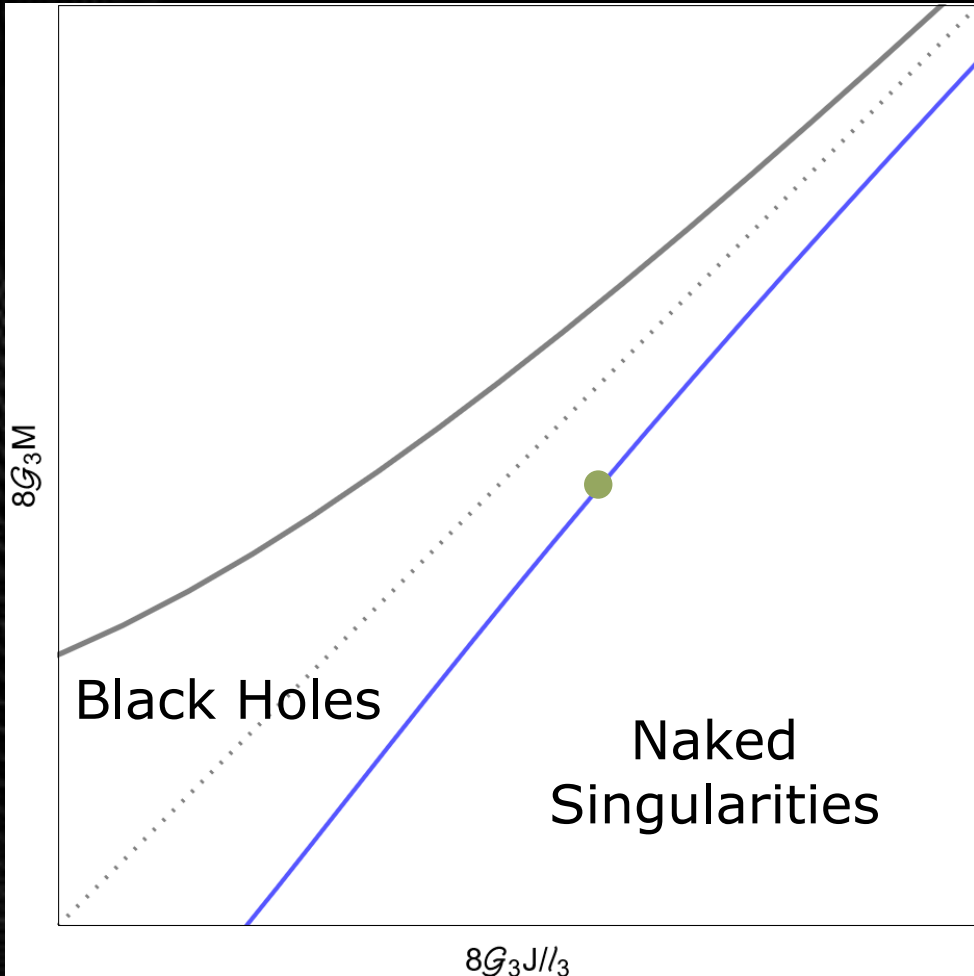


Strategy

$$g^{rr}(r_{\min} + \delta r, \tilde{a} + \delta \tilde{a}, x_1 + \delta x_1) = g^{rr}(r_{\min}, \tilde{a}, x_1) + \delta g^{rr}$$

Perturbative analysis

Linear response after absorption of a test particle

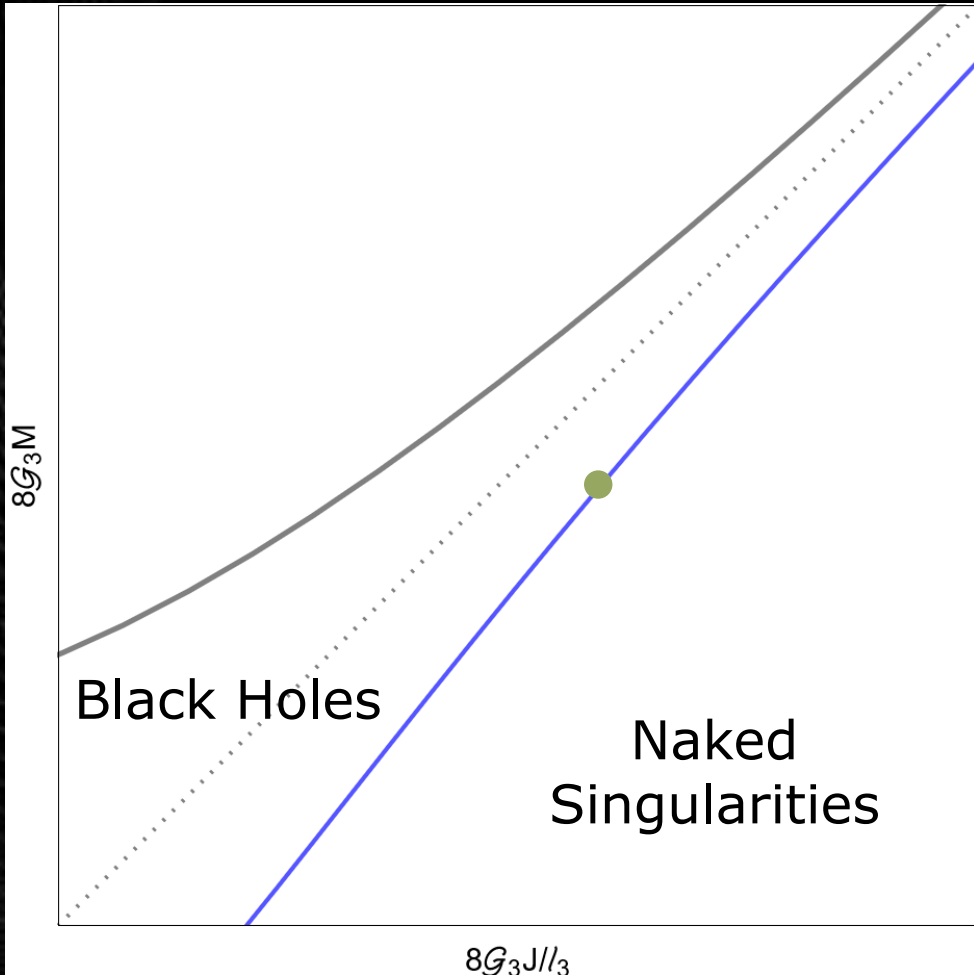


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Perturbative analysis

Linear response after absorption of a test particle



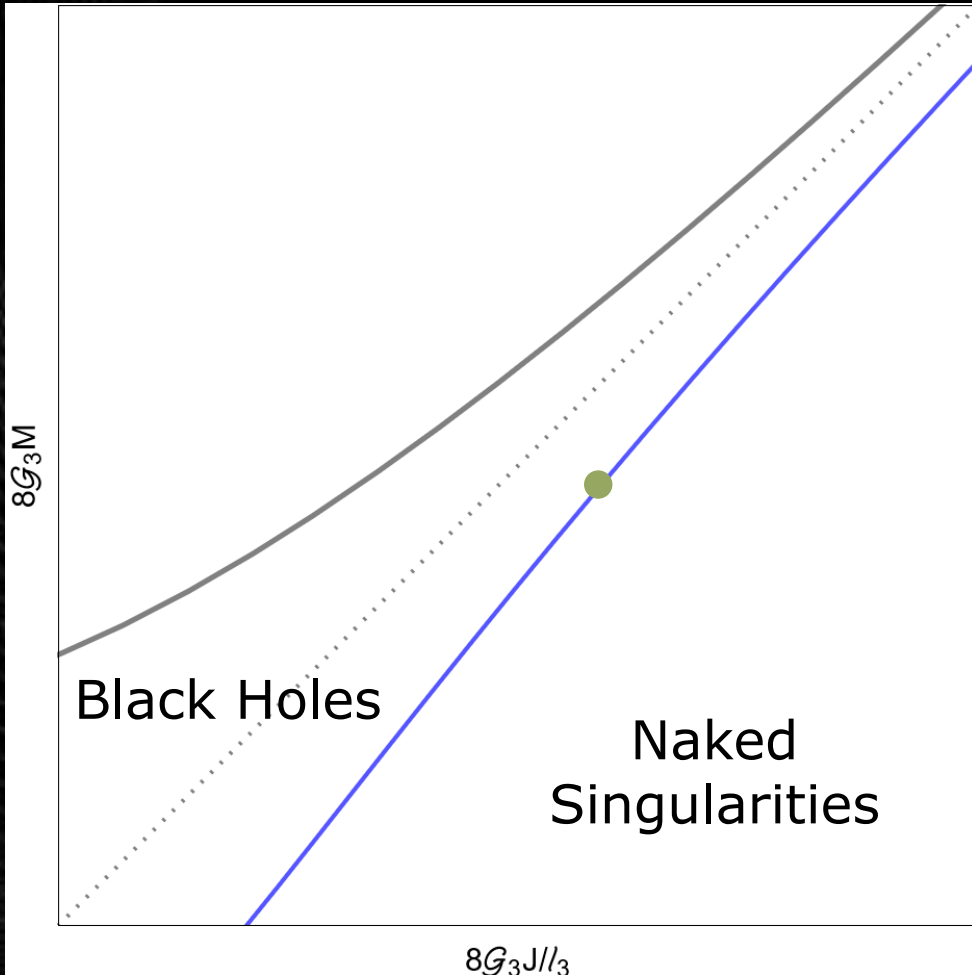
Strategy

$$g^{rr}(r_{\min} + \delta r, \tilde{a} + \delta \tilde{a}, x_1 + \delta x_1) = g^{rr}(r_{\min}, \tilde{a}, x_1) + \delta g^{rr}$$

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Perturbative analysis

Linear response after absorption of a test particle



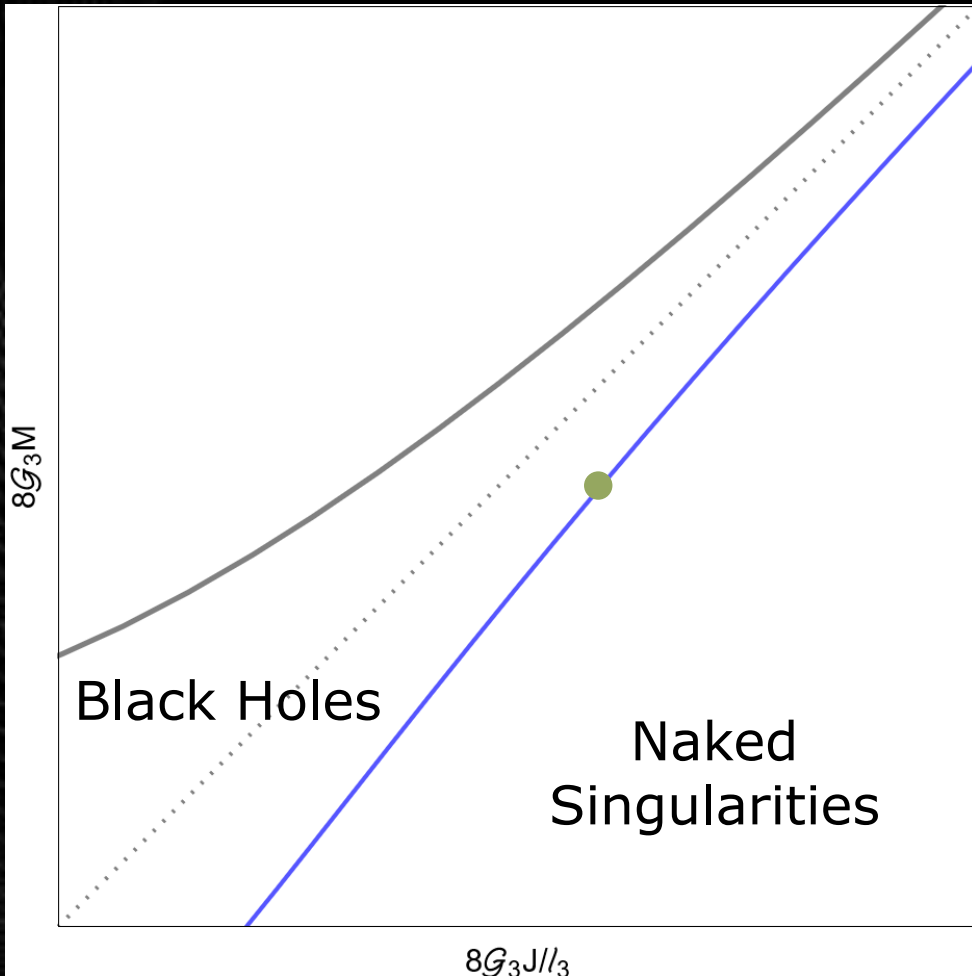
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Perturbative analysis

Linear response after absorption of a test particle



Strategy

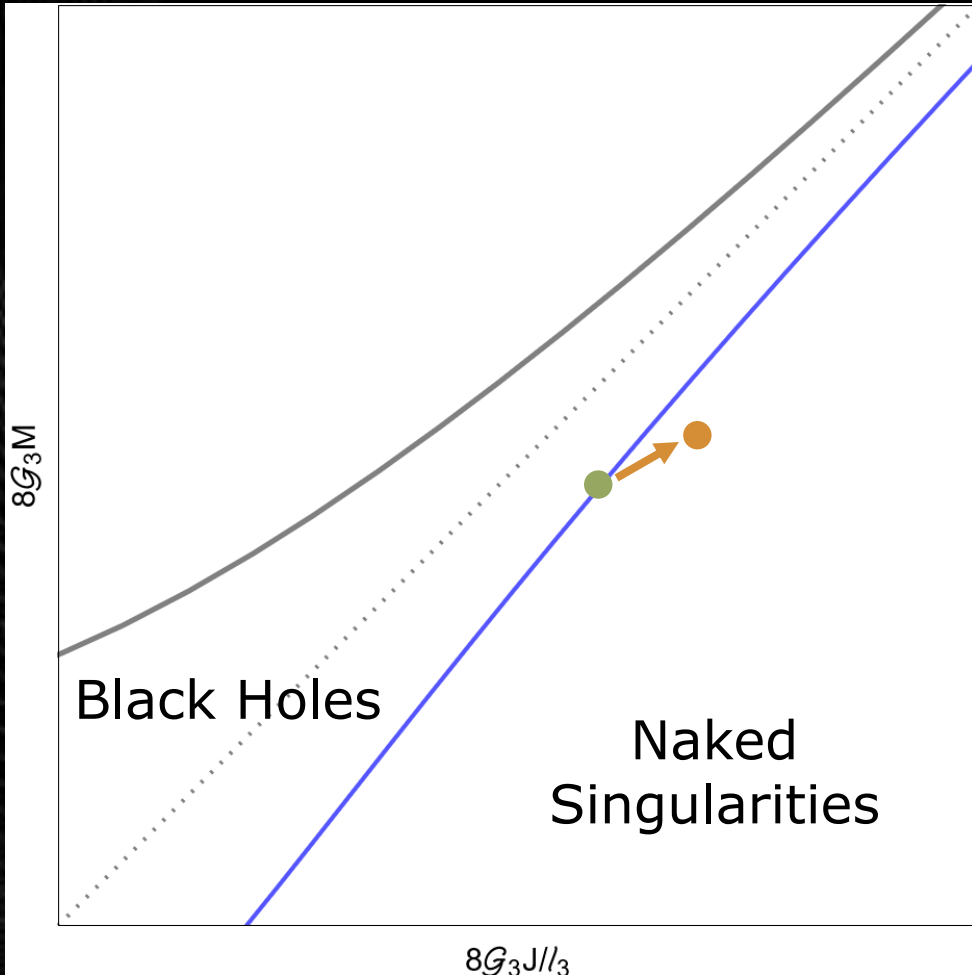
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$$\delta g^{rr} > 0$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

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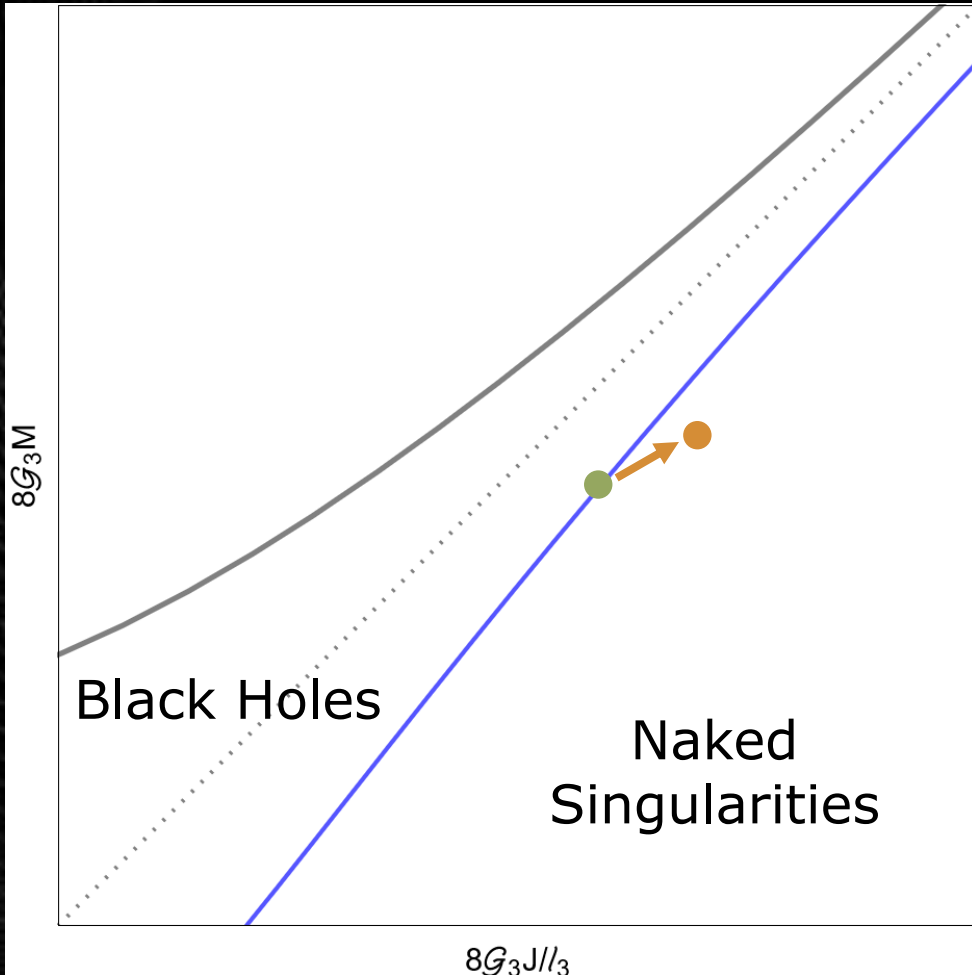
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$$\delta g^{rr} > 0$$

Naked singularity forms

Perturbative analysis

Linear response after absorption of a test particle



Strategy

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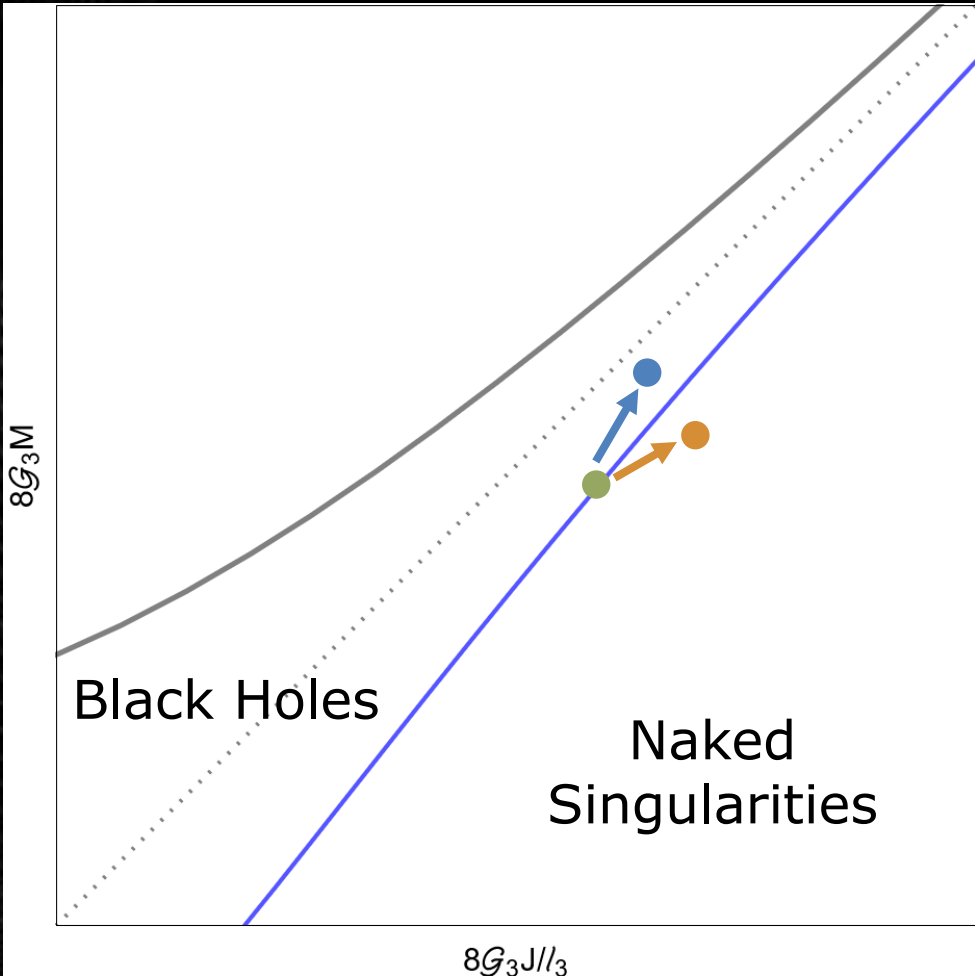
$$\delta g^{rr} > 0$$

Naked singularity forms

$$\delta g^{rr} < 0$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

$$g^{rr}(r_{\min} + \delta r, \tilde{a} + \delta \tilde{a}, x_1 + \delta x_1) = g^{rr}(r_{\min}, \tilde{a}, x_1) + \delta g^{rr}$$

$$\delta g^{rr} = \left. \frac{\partial g^{rr}}{\partial r} \right|_{r=r_{\min}} \delta r + \left. \frac{\partial g^{rr}}{\partial \tilde{a}} \right|_{r=r_{\min}} \delta \tilde{a} + \left. \frac{\partial g^{rr}}{\partial x_1} \right|_{r=r_{\min}} \delta x_1$$

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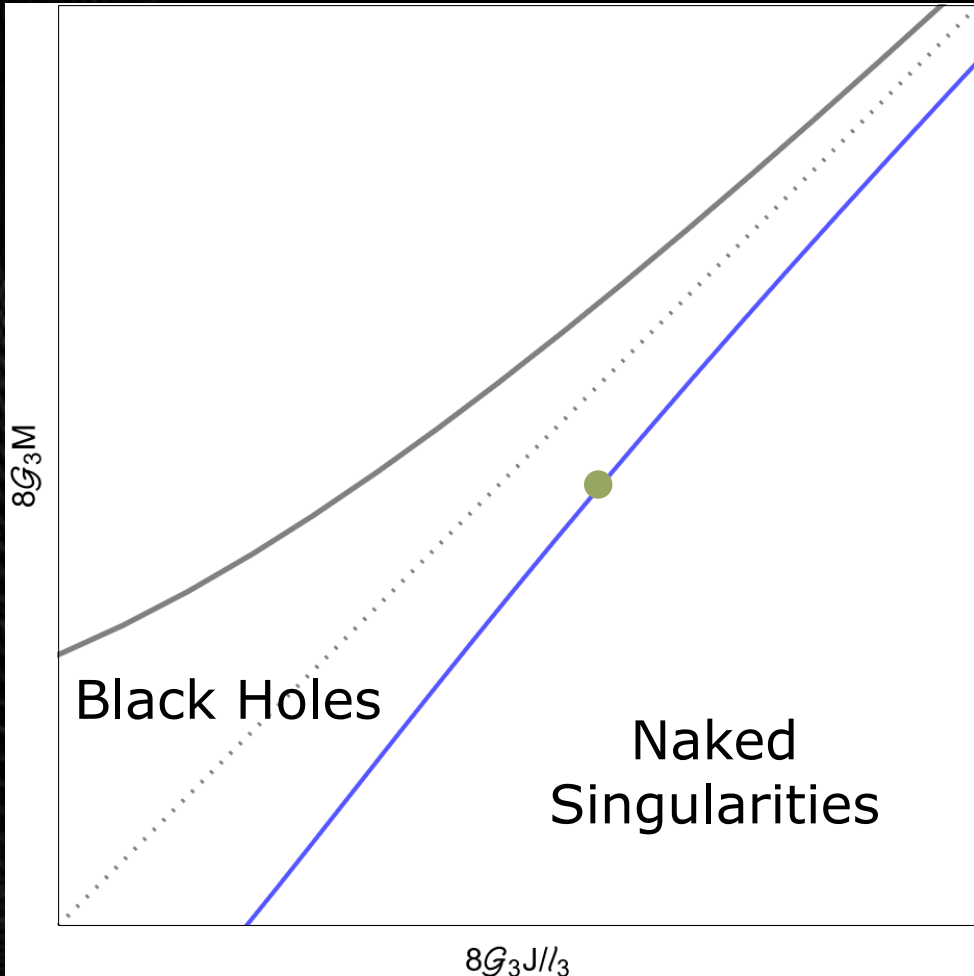
Naked singularity forms

$$\delta g^{rr} < 0$$

Horizon forms

Perturbative analysis

Linear response after absorption of a test particle



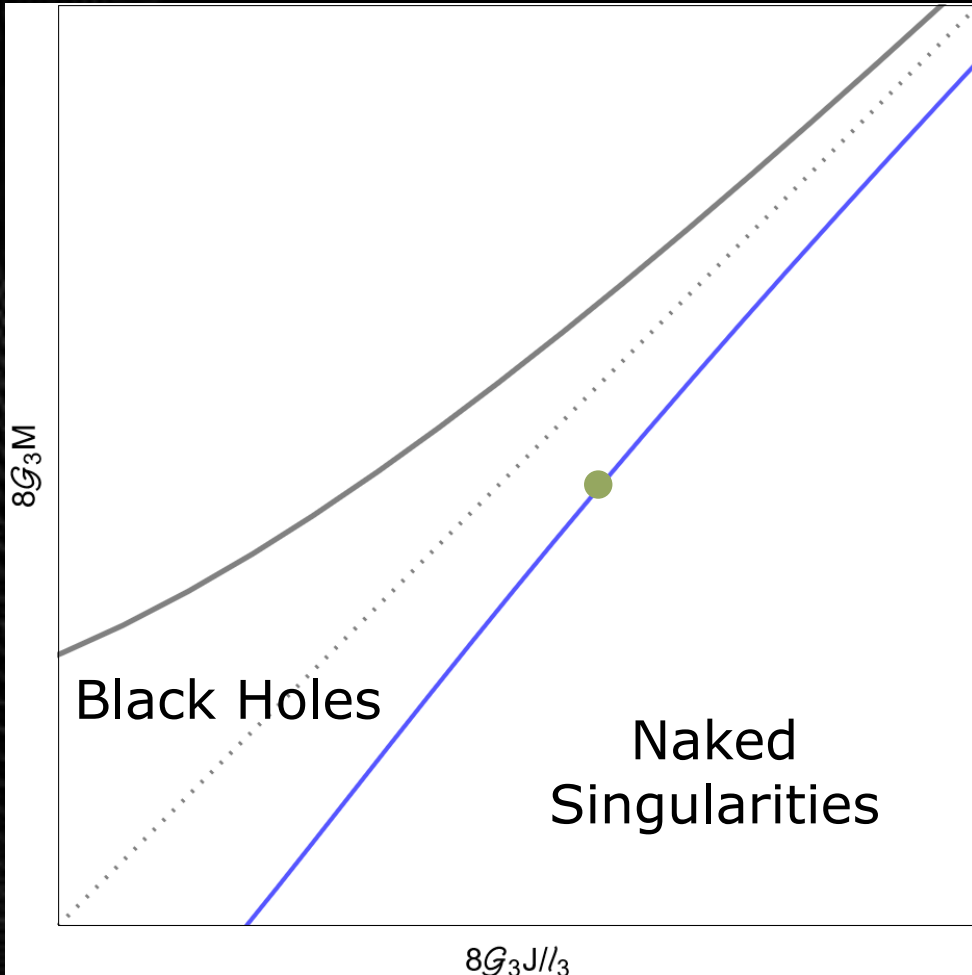
Strategy

Only particles with sufficiently low angular momentum are captured

$$\frac{\delta J}{\delta M} \leq L_{\max}$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

Only particles with sufficiently low angular momentum are captured

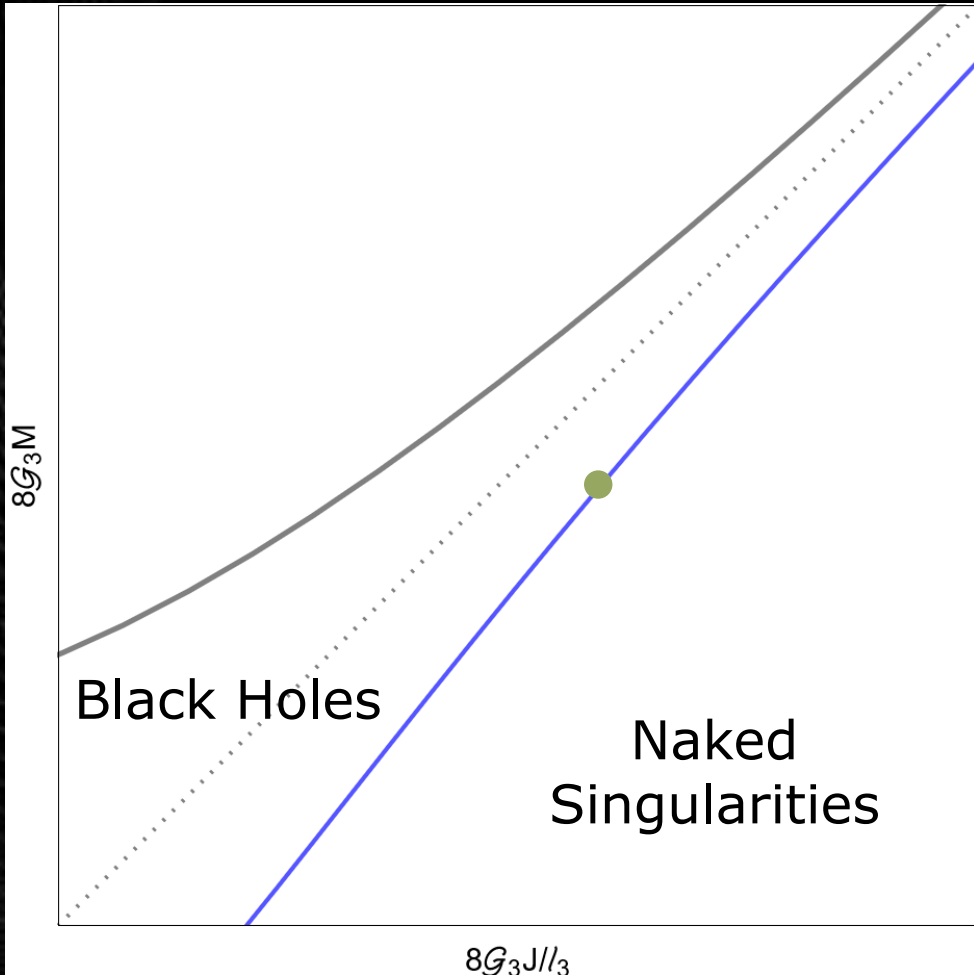
$$\frac{\delta J}{\delta M} \leq L_{\max}$$

The worst case

$$\frac{\delta J}{\delta M} = L_{\max}$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

Only particles with sufficiently low angular momentum are captured

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The worst case

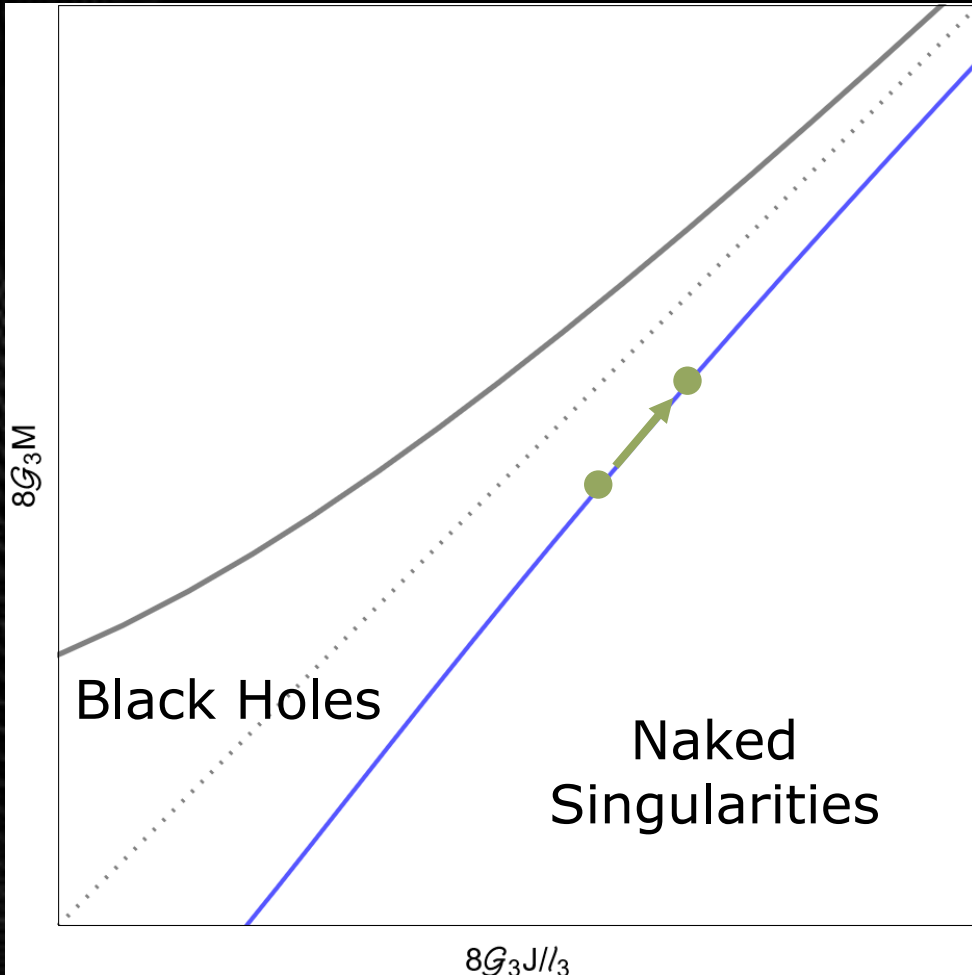
$$\frac{\delta J}{\delta M} = L_{\max}$$



$$\delta g^{rr} = 0$$

Perturbative analysis

Linear response after absorption of a test particle



Strategy

Only particles with sufficiently low angular momentum are captured

$$\frac{\delta J}{\delta M} \leq L_{\max}$$

The worst case

$$\frac{\delta J}{\delta M} = L_{\max}$$



$$\delta g^{rr} = 0$$

Perturbative analysis

Linear response after absorption of a test particle

Strategy

The event horizon is not destroyed!

Black Holes

Naked
Singularities

$8G_3J/l_3$

The worst case

$$\frac{\delta J}{\delta M} = L_{\max}$$



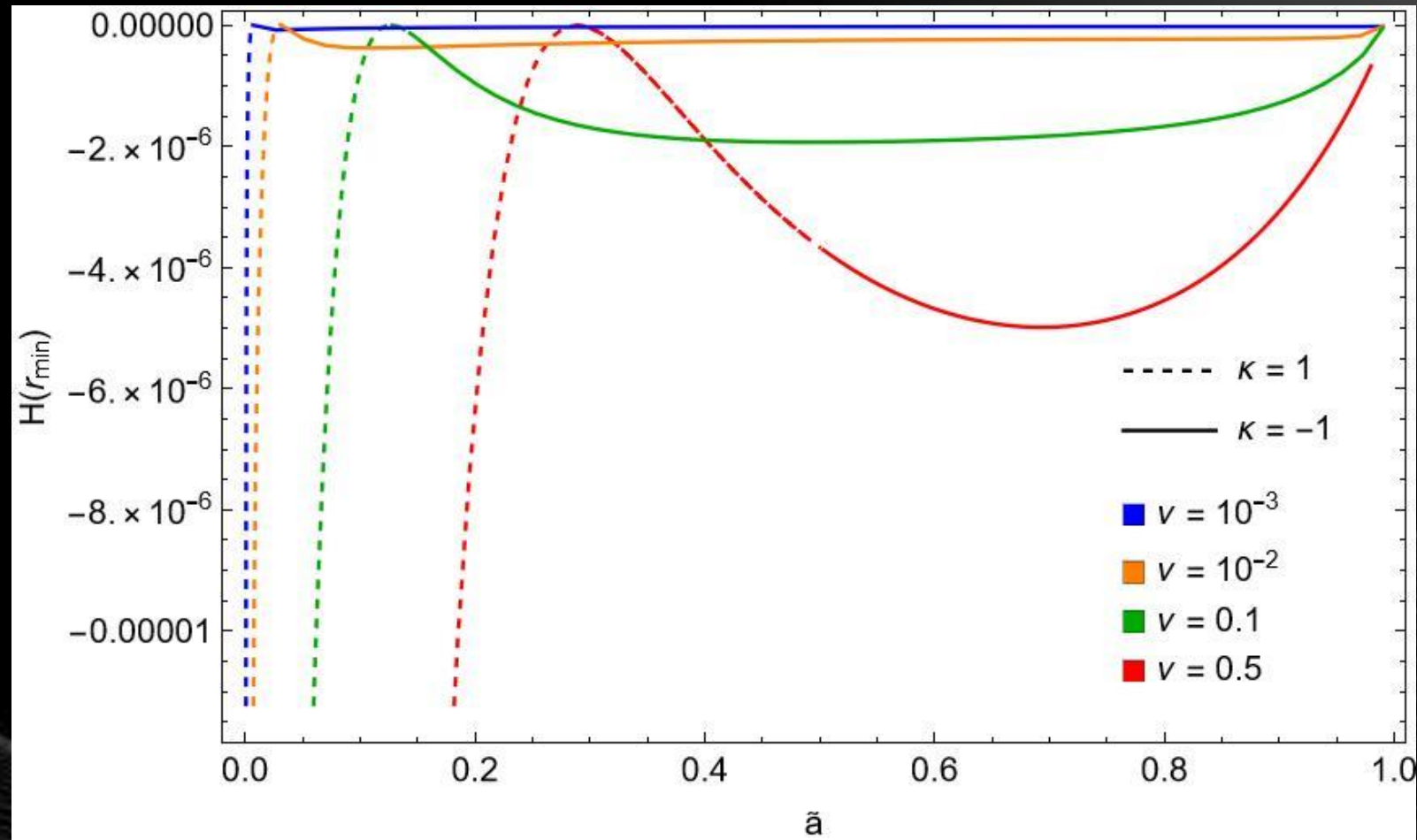
$$\delta g^{rr} = 0$$

Numerical assessment

Test particles with finite mass ($\delta M \ll M$)

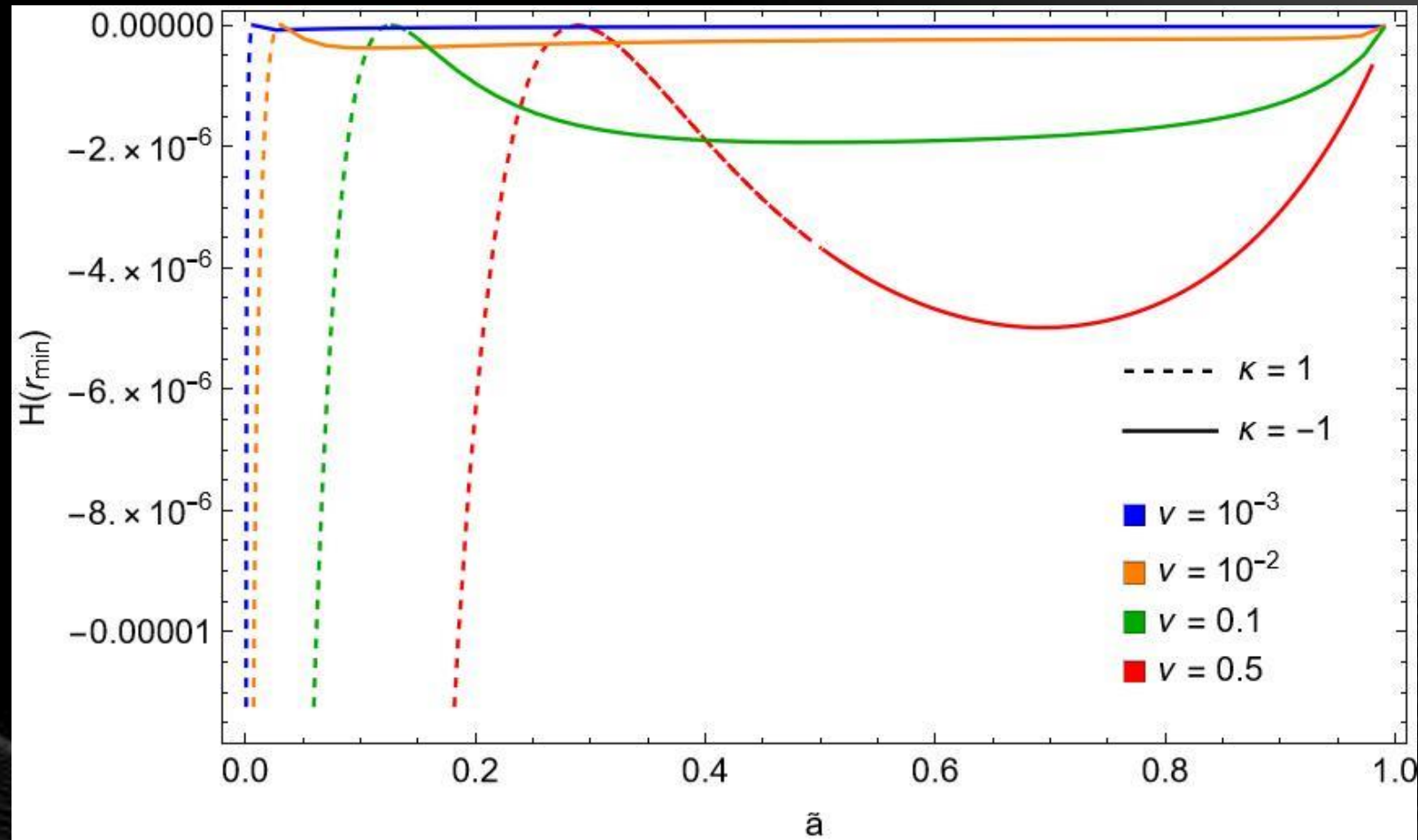
Numerical assessment

Test particles with finite mass ($\delta M \ll M$)



Numerical assessment

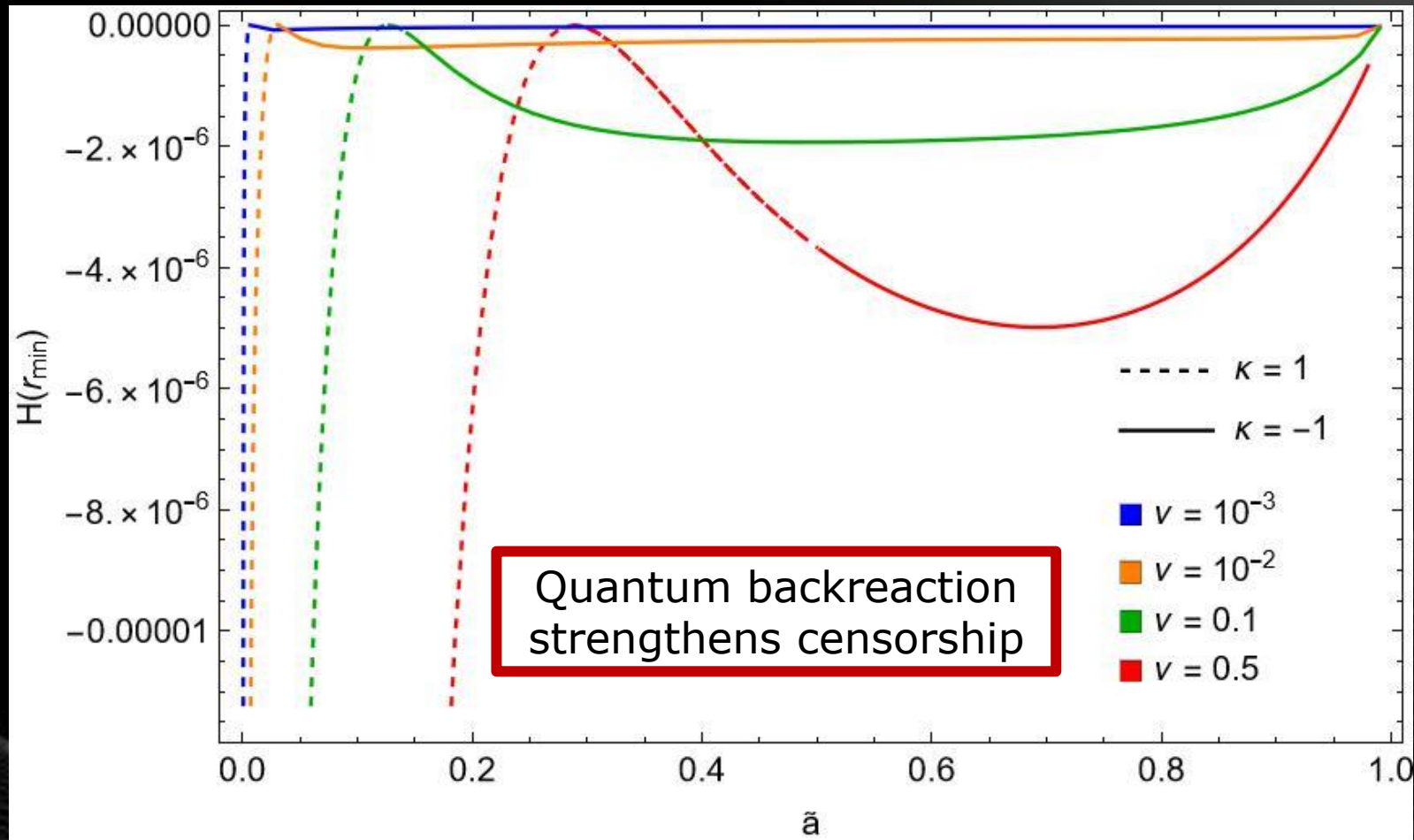
Test particles with finite mass ($\delta M \ll M$)



wCCC respected

Numerical assessment

Test particles with finite mass ($\delta M \ll M$)



wCCC respected



THANKS FOR THE ATTENTION!