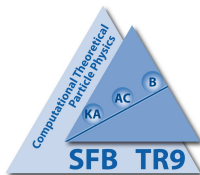


# Nucleon matrix elements and disconnected diagrams

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# Outline

**Nucleon  $\sigma$ -terms**

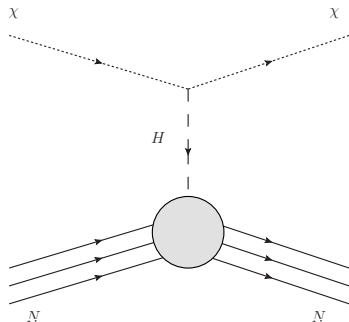
$$\langle X \rangle_{u+d-2s, \mu^2}$$

**Conclusion**

# Strange content of the Nucleon : Motivations

## Motivations

- Experimental direct detection of dark matter put bounds on the WIMP-Nucleon cross section
- Results are interpreted using various models (including SUSY) : systematic uncertainty due to  $\langle N(p) | \bar{q}q | N(p) \rangle$   
 → non-perturbative computation is required.



- sigma term :

$$\sigma_{\pi N} \equiv m_l \langle N(p) | \bar{u}u + \bar{d}d | N(p) \rangle$$

- dimensionless ratio :

$$Y_N \equiv \frac{2 \langle N(p) | \bar{s}s | N(p) \rangle}{\langle N(p) | \bar{u}u + \bar{d}d | N(p) \rangle}$$

- dimensionless coupling :

$$f_{Tq} = \frac{m_q}{m_N} \langle N | \bar{q}q | N \rangle$$

# Lattice setup (I)

## $N_f = 2 + 1 + 1$ dynamical simulations

- $N_f = 2 + 1 + 1$  configurations generated by ETMC
- Fixed lattice spacing  $a = 0.078$  fm
- $32^3 \times 64$  lattices :  $L = 2.5$  fm
- $m_{\text{PS}}L > 3.5$
- 3 pion masses [300, 450] MeV

## Mixed action setup

- Mixed action setup : introduce a doublet of twisted mass fermions  $\chi_q$  with a mass  $\mu_q$
- $\mu_q$  can be tuned to reproduce the  $K, D$  meson masses in the unitary setup
- Multimass solver inversions from the light to the charm quark mass

## Lattice setup (II)

- Rotation to the physical basis :

$$\bar{\chi}_q \gamma_5 \tau^3 \chi_q \longrightarrow \bar{q}q + \bar{q}'q' = \begin{cases} \bar{u}u + \bar{d}d & \text{if } q = u, d \\ \bar{s}s + \bar{s}'s' & \text{if } q = s \\ \bar{c}c + \bar{c}'c' & \text{if } q = c \end{cases}$$

- Three-points function definition (disconnected part only):

$$R(t_s, t_{\text{op}}) = \frac{\langle C_{2\text{pt}}(t_s)D(t_{\text{op}}) \rangle - \langle C_{2\text{pts}}(t_s) \rangle \langle D(t_{\text{op}}) \rangle}{C_{2\text{pts}}(t_s)}$$

- Connected pieces : fixed source-sink separation  $t_s/a = 12$

# Renormalization

- Discretization that break chiral symmetry suffer from mixing under renormalization of the light and strange content
- Twisted mass property in the mixed action setup : cancellation of chiral violating effects between “loops” of valence quarks regularized with opposite values of  $r$ .

$$(\bar{s}s)_R = \frac{Z_P}{2} [(\bar{s}s)_{r=+1} + (\bar{s}'s')_{r=-1}].$$

(R. Frezzotti, G. Herdoiza, G. Rossi)

- The quantity :

$$f_{T_s} = \frac{1}{2} \frac{\mu_s}{m_N} \langle N(p) | \bar{s}s + \bar{s}'s' | N(p) \rangle$$

is thus renormalization free.

# Twisted mass fermions and disconnected diagrams

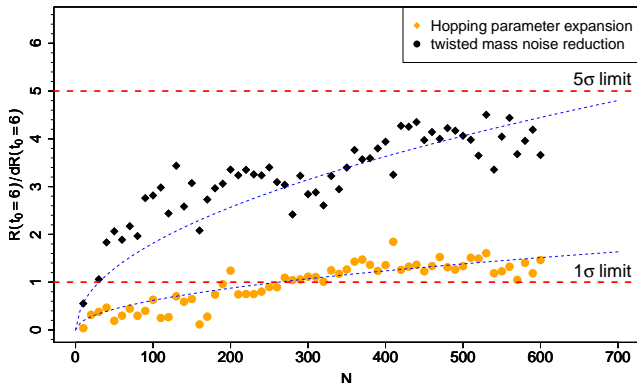
- Twisted mass Wilson Dirac operator :  $D_{\pm}[U] = D_W[U] + am_0 \pm ia\mu_q\gamma_5$
- Noise reduction technique specific for twisted mass fermions :
  - ◆ Exact relation : (C. Michael, C. Urbach, 2007)

$$\frac{1}{D_-} - \frac{1}{D_+} = 2ia\mu_q \frac{1}{D_-} \gamma_5 \frac{1}{D_+}$$

- r.h.s can be evaluated stochastically using “one-end-trick”
- Useful for correlators containing the insertion of the operator  $\bar{\chi}(x)\Gamma\tau^3\chi(x)$  (twisted basis)

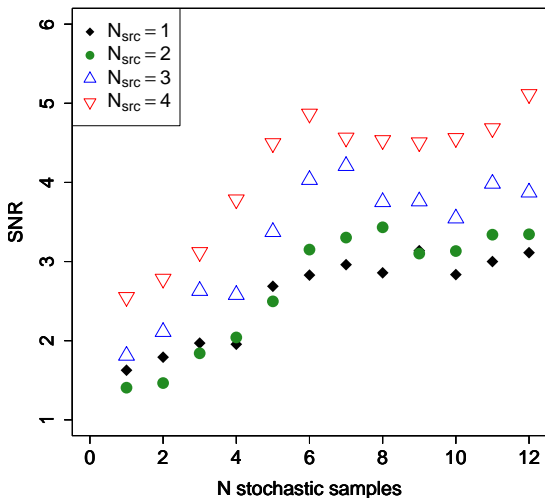
# Performances

- Signal-over-noise ratio of a typical three point function vs  $N_{\text{conf}}$
- In comparison with hopping parameter expansion noise reduction save a factor  $\sim 10$  in statistics.

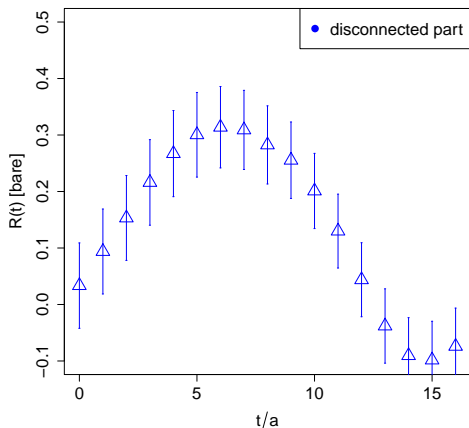
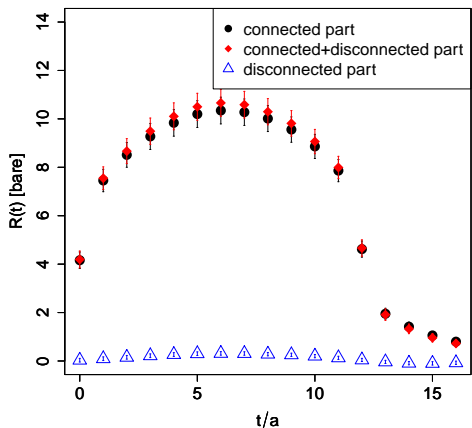




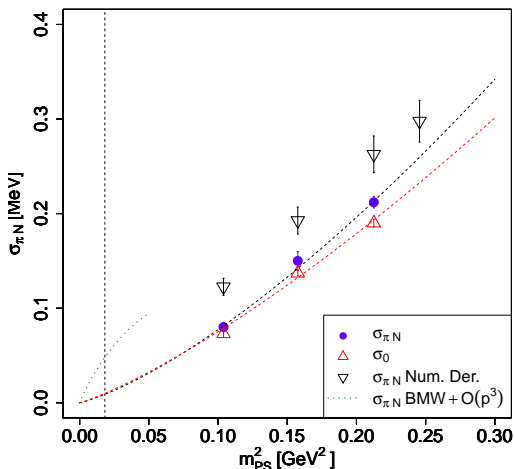
# Performances



# Result : light sector

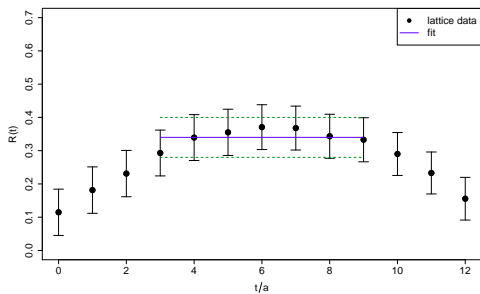


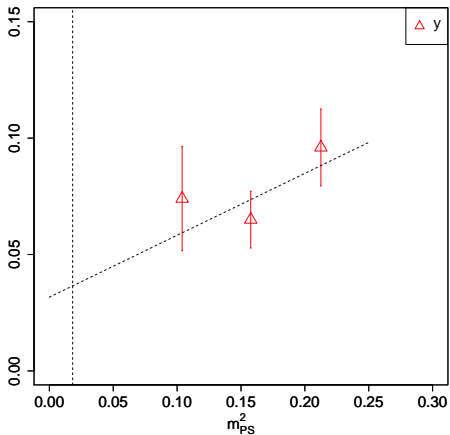
- $m_{PS} = 380$  MeV
- statistics :  $\sim 845$  gauge configurations
- 12 volume sources (no dilution), 4 sets of two point
- Disconnected piece is different from 0 at  $4\sigma$

chiral behaviour :  $\sigma_{\pi N}$ 

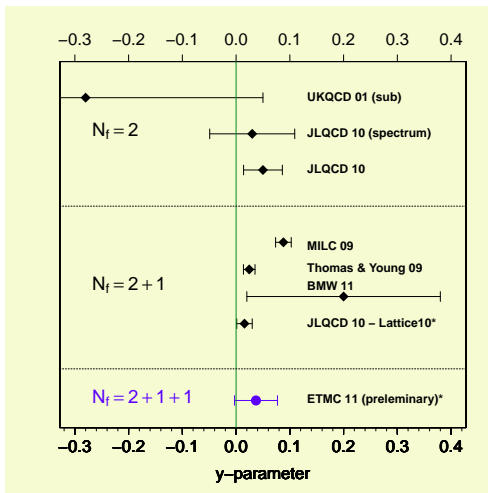
- $\sigma_{\pi N} = m_q \frac{\partial m_N}{\partial m_q} = m_q \langle N | \bar{u}u + \bar{d}d | N \rangle$
- $\sigma_0 = m_q \langle N | \bar{u}u + \bar{d}d - 2\bar{s}s | N \rangle$

# Result : strange sector



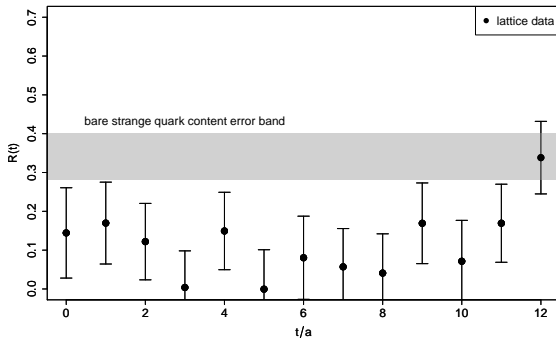
chiral behaviour :  $y_N$ 

# Preliminary results



- Systematic errors not estimated

# Result : charm sector



- $\langle N|\bar{c}c|N\rangle < \langle N|\bar{s}s|N\rangle$  but not really interesting
- we want to compare  $f_{T_S}$  and  $f_{T_C}$

# Motivations

Consider

$$\langle p, s | \bar{\psi} \gamma^{\{\mu} i D^{\nu\}} \lambda_8 \psi | p, s \rangle \Big|_{\mu^2} = 2 \langle X \rangle_{u+d-2s, \mu^2} p^{\{\mu} p^{\nu\}}, \quad \text{with} \quad \psi = \begin{pmatrix} u \\ d \\ s \end{pmatrix}$$

where  $\lambda_8 = \text{diag}(1, 1, -2)$ .

- $\langle X \rangle_{u+d-2s, \mu^2}$  has never been estimated on the lattice
- known experimentally
- Cancellation of renormalizations factor in  $\frac{\langle X \rangle_{u-d, \mu^2}}{\langle X \rangle_{u+d-2s, \mu^2}}$
- Problem : (quark)-disconnected diagrams



# Techniques

$$J^8 = \bar{\psi} \Gamma \lambda_8 \psi = \frac{1}{\sqrt{3}} (\bar{u} \Gamma u + \bar{d} \Gamma d - 2\bar{s} \Gamma s) \quad (1)$$

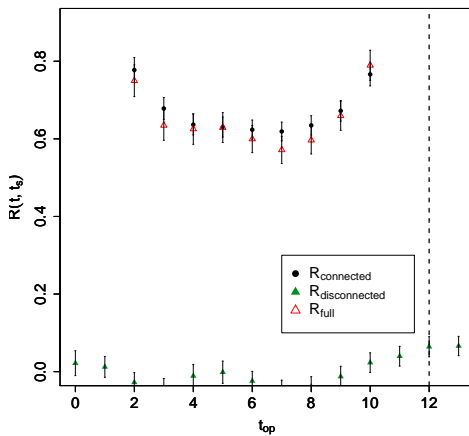
in the twisted basis :

$$\tilde{J}^8 = \bar{\chi} \Gamma \chi - \bar{\chi}_s \Gamma \chi_s \quad (2)$$

$$D_{q,\pm} - D_{q',\pm} = \pm 2i\kappa\gamma_5 a (\mu_q - \mu_{q'}) \quad (3)$$

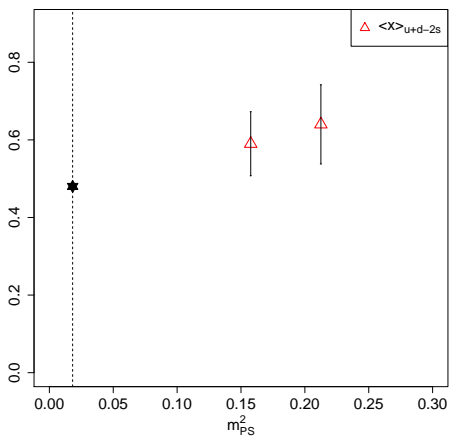
$$\frac{1}{D_{q',\pm}} - \frac{1}{D_{q,\pm}} = \frac{1}{D_{q,\pm}} (D_{q,\pm} - D_{q',\pm}) \frac{1}{D_{q',\pm}} = \pm 2i\kappa a (\mu_q - \mu_{q'}) \frac{1}{D_{q,\pm}} \gamma_5 \frac{1}{D_{q',\pm}}, \quad (4)$$

# Results



$t_s = 12a$  fixed

# Chiral behaviour : $\langle X \rangle_{u+d-2s, \mu^2}$



(Phenomenological value from S. Alekhin)

# Conclusion

- Direct computation of the nucleon sigma terms
- Crucial use of the twisted mass noise reduction
- Mixing problems are avoided in our setup
- Disconnected contribution  $4 - 5\sigma$  away from zero in the light and strange sector
- First computation of  $\langle X \rangle_{u+d-2s, \mu^2}$
- Use a generalization of the twisted mass noise reduction
- disconnected contribution compatible with 0, but with relatively small error

# Chiral condensate

