

$\mathcal{O}(a^2)$ Corrected Renormalization Constants

$$N_f = 4$$

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C. Alexandrou, M. Constantinou, C. Kallidonis

Physics Department, University of Cyprus

OUTLINE

- A** Operators under study
- B** Non-Perturbative evaluation
- C** Statistics
- D** Perturbative computation
- E** Renormalization Conditions (RI'-MOM)
- F** Results
- G** Summary

NON-PERTURBATIVE RENORMALIZATION

Operators

$$\mathcal{O}^S = \hat{1}$$

$$\mathcal{O}^P = \gamma^5$$

$$\mathcal{O}^V = \gamma^\mu$$

$$\mathcal{O}^A = \gamma^5 \gamma^\mu$$

$$\mathcal{O}^T = \gamma^5 \sigma^{\mu\nu}$$

$$\mathcal{O}^{DV} = \gamma^{\{\mu} \overleftrightarrow{D}^{\nu\}}$$

$$\mathcal{O}^{DA} = \gamma^5 \gamma^{\{\mu} \overleftrightarrow{D}^{\nu\}}$$

- No mixing with lower dimension operators
- No disconnected diagrams
- symmetrization and subtraction of the traces

Momentum Source Method

- ★ Dirac equation solved with momentum source
- ★ # of inversion depends on the # of momenta considered
- ★ Application of any operator
- ★ High statistical accuracy is achieved

Statistics

| | β | k_{cr} | μ_0 | $L^3 \times T$ |
|-----|---------|----------|---------|------------------|
| 3p | 1.95 | 0.160826 | 0.0180 | $24^3 \times 48$ |
| 3m | 1.95 | 0.161229 | 0.0180 | $24^3 \times 48$ |
| 2p | 1.95 | 0.160826 | 0.0085 | $24^3 \times 48$ |
| 2m | 1.95 | 0.161229 | 0.0085 | $24^3 \times 48$ |
| 8p | 1.95 | 0.160524 | 0.0020 | $24^3 \times 48$ |
| 8m | 1.95 | 0.161585 | 0.0020 | $24^3 \times 48$ |
| 5p | 2.10 | 0.155949 | 0.0078 | $32^3 \times 64$ |
| 5m | 2.10 | 0.156291 | 0.0078 | $32^3 \times 64$ |
| 4m | 2.10 | 0.156250 | 0.0064 | $32^3 \times 64$ |
| 3p | 2.10 | 0.156017 | 0.0046 | $32^3 \times 64$ |
| 2ap | 2.10 | 0.166042 | 0.0030 | $32^3 \times 64$ |
| 2am | 2.10 | 0.166157 | 0.0030 | $32^3 \times 64$ |

| $\beta = 1.95$ ($\mu = 0.0180, 0.0085$) | $\beta = 1.95$ ($\mu = 0.002$) | $\beta = 2.10$ |
|---|----------------------------------|--------------------------------|
| $(n_t, 2, 2, 2), n_t : 1 - 7$ | $(n_t, 2, 2, 2), n_t : 10 - 12$ | $(n_t, 2, 2, 2), n_t : 3 - 12$ |
| $(n_t, 3, 3, 3), n_t : 1 - 10$ | $(n_t, 3, 3, 3), n_t : 5 - 10$ | $(n_t, 3, 3, 3), n_t : 1 - 13$ |
| $(3, 3, 3, 2)$ | $(3, 3, 3, 2), n_t : 8 - 10$ | $(n_t, 4, 4, 4), n_t : 4 - 15$ |
| $(3, 3, 3, 4)$ | $(3, 3, 3, 4)$ | $(n_t, 5, 5, 5), n_t : 2 - 11$ |
| $(3, 3, 4, 4)$ | $(3, 3, 4, 4)$ | $(3, 3, 3, 2), (3, 3, 2, 2)$ |
| | | $(3, 3, 3, 4), (3, 3, 4, 4)$ |

~ 10-50 configurations per momentum

Renormalization Conditions (RI'-MOM scheme):

$$Z_q = \frac{1}{12} \text{Tr}[(S^L(p))^{-1} S^{(0)}(p)] \Big|_{p^2=\mu^2}$$

$$Z_q^{-1} Z_{\mathcal{O}}^{\mu\nu} \frac{1}{12} \text{Tr}[\Gamma_{\mu\nu}^L(p) \Gamma^{(0)-1}_{\mu\nu}(p)] \Big|_{p^2=\mu^2} = 1$$

$$S^{(0)}(p) = \frac{-i \sum_{\rho} \gamma_{\rho} \sin(p_{\rho})}{\sum_{\rho} \sin(p_{\rho})^2}$$

$$\Gamma_{\mu\nu}^{(0)}(p) = -i \tilde{\Gamma}_{\{\mu\nu\}} \sin(p_{\nu})$$

$$\frac{1}{L^2} \ll \Lambda_{\text{QCD}}^2 \ll \bar{\mu}^2 \ll \frac{1}{a^2}$$

Reliable perturbation theory (Wilson coefficients)

Small $\mathcal{O}(a)$ lattice effects

In practice: $0.5 \leq (ap)^2 \leq 4.5$

Example of Perturbative Results:

- Iwasaki gluons , $c_{\text{SW}} = 0$
- Landau gauge
- $m = 0, \mu_0 = 0$

$$\begin{aligned} \frac{1}{4} \text{Tr} \left[(S^L(p))^{-1} S_{\text{tree}}^{(0)}(p) \right] \Big|_{p_\rho = \mu_\rho} &= 1 + \tilde{g}^2 \left\{ -8.11657 \right. \\ &+ a^2 \left[\mu^2 (0.620227 - 0.0748167 \ln(a^2 \mu^2)) \right. \\ &+ \left. \left. \frac{\sum_\rho \mu_\rho^4}{\mu^2} (1.85344 - 0.963033 \ln(a^2 \mu^2)) \right] \right\} \\ &+ \mathcal{O}(a^4 g^2, g^4) \end{aligned}$$

- Ambiguity on the choice of the momentum direction
- **(3,3,3,3):** $(ap)^2 = 2.06045, (ap)^4 = 1.18557$
- **(8,2,2,2):** $(ap)^2 = 2.06045, (ap)^4 = 1.75809$

Criterion for choosing the momenta

arXiv:1004.1115

$$\text{local : } \frac{\sum_{\rho} p_{\rho}^4}{\left(\sum_{\rho} p_{\rho}^2\right)^2} \leq 0.28 \quad \text{twist - 2 : } \frac{\sum_{\rho} p_{\rho}^4}{\left(\sum_{\rho} p_{\rho}^2\right)^2} \leq 0.4$$

Minimizes the $\sum_{\rho} p_{\rho}^4$ contributions of the perturbative expressions

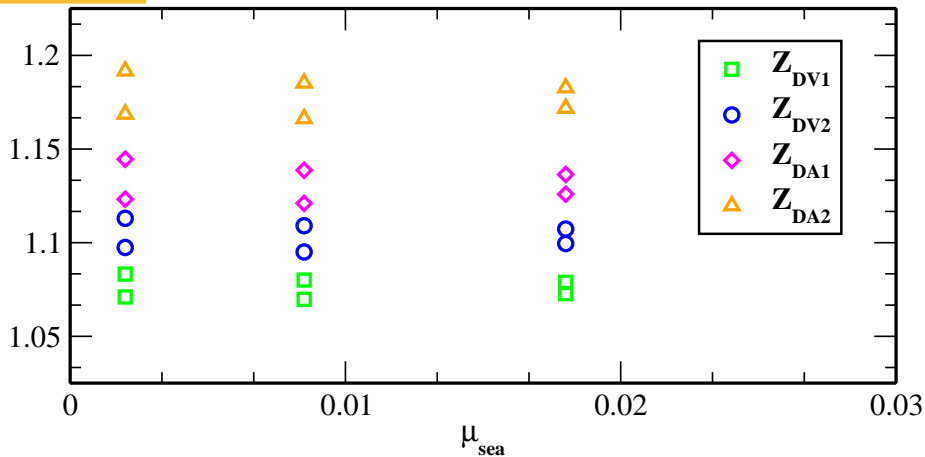
$$\begin{aligned} Z_{\text{DV1}}^{\text{impr}} = Z_{\text{DV1}} \\ -a^2 \tilde{g}^2 \left[c_1 p^2 + c_2 p_{\nu_1}^2 + \frac{c_3 p_{\nu_1}^4}{p^2} + \frac{c_4 p_{\nu_1}^4}{8p_{\nu_1}^2 + p^2} + \frac{c_5 p_{\nu_1}^6}{(8p_{\nu_1}^2 + p^2)^2} + \left(c_6 p^2 + c_7 p_{\nu_1}^2 + \frac{c_8 p_{\nu_1}^4}{(8p_{\nu_1}^2 + p^2)} \right) \ln[a^2 p^2] \right. \\ \left. + p^4 \left(c_9 \frac{1}{p^2} + c_{10} \frac{p_{\nu_1}^2}{p^2} + c_{11} \frac{1}{8p_{\nu_1}^2 + p^2} + c_{12} \frac{p_{\nu_1}^2}{(8p_{\nu_1}^2 + p^2)^2} \right. \right. \\ \left. \left. + \left(c_{13} \frac{1}{p^2} + c_{14} \frac{1}{(8p_{\nu_1}^2 + p^2)} \right) \ln[a^2 p^2] \right) \right] \end{aligned}$$

No momentum criterion was applied here.

Mass dependence

$$\beta = 1.95$$

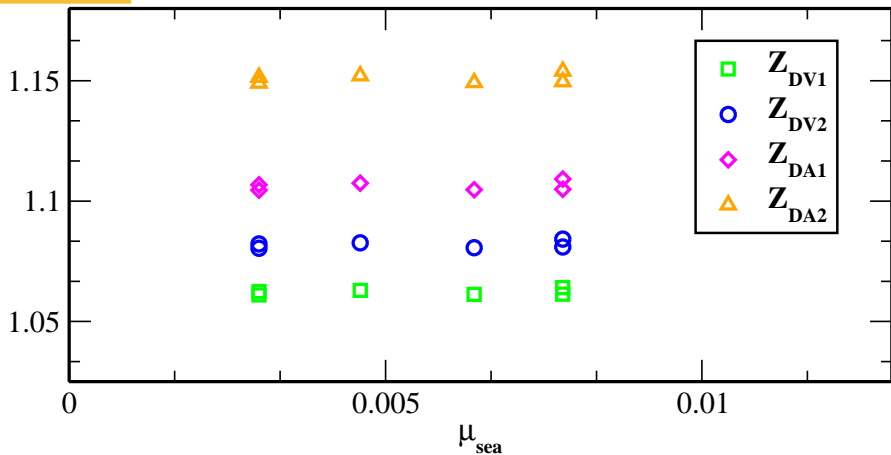
$$24^3 \times 48$$



Mass dependence

$$\beta = 2.10$$

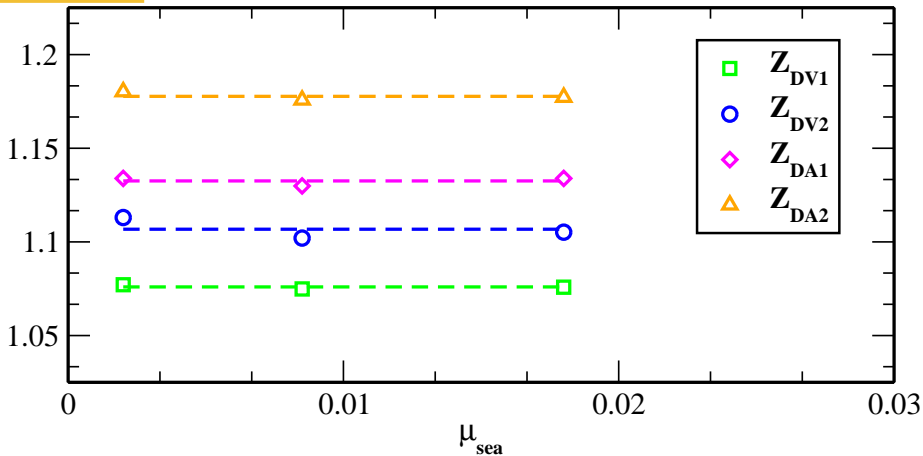
$$32^3 \times 64$$



Theta average: Mass dependence

$$\beta = 1.95$$

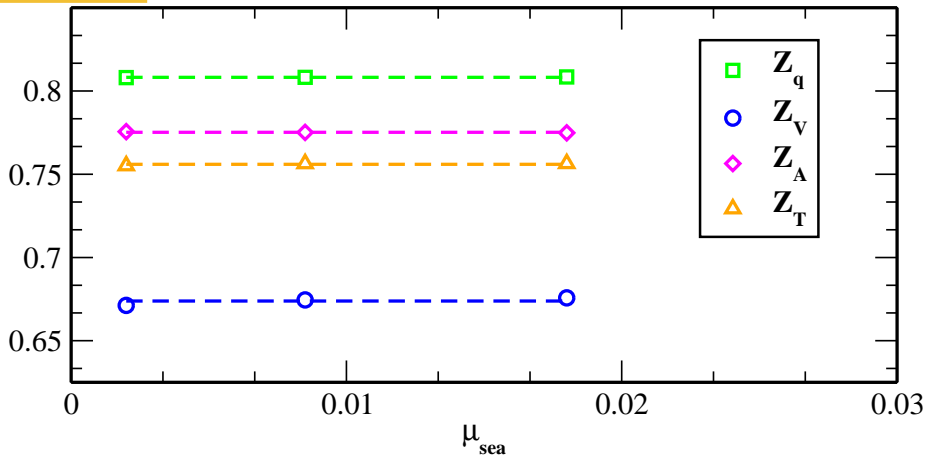
$$24^3 \times 48$$



Theta average: Mass dependence

$$\beta = 1.95$$

$$24^3 \times 48$$



Volume effects

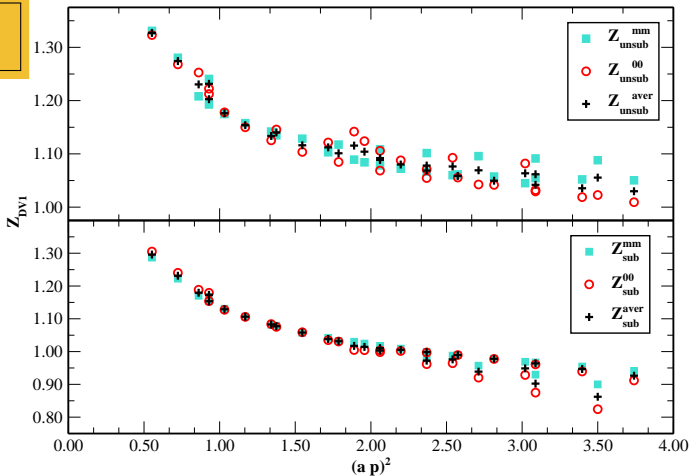
Not studied yet

We don't expect significant dependence on volume

Averaging over different directions

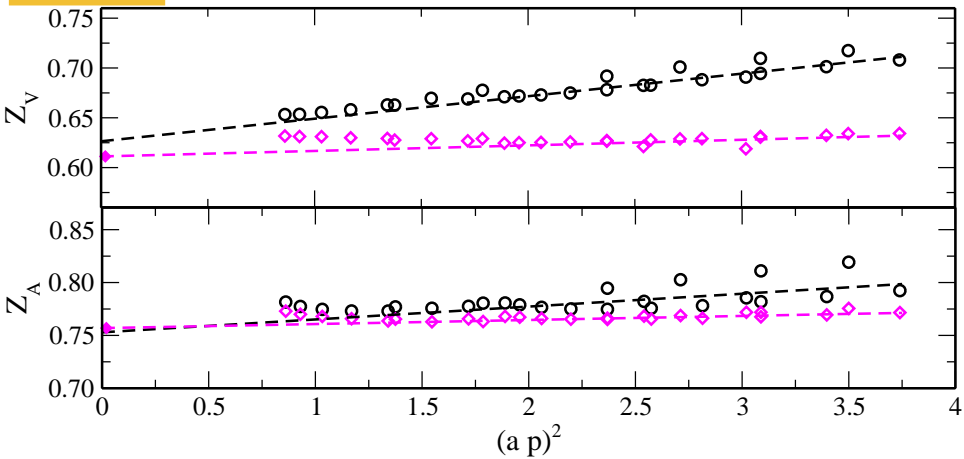
$$\beta = 1.95$$

$$24^3 \times 48$$



$$Z_{DV1}^{aver} = \left(Z_{DV1}^{00} + \sum_{i=1}^3 Z_{DV1}^{ii} / 3 \right) / 2$$

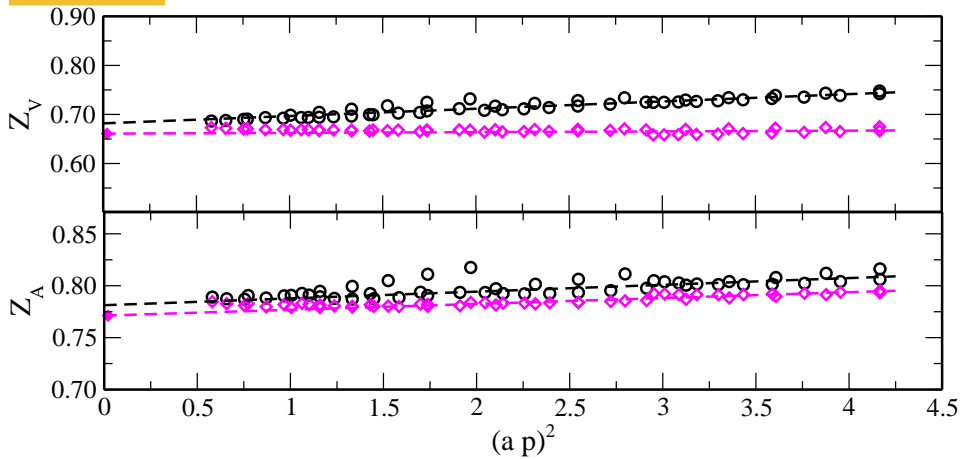
Conversion to \overline{MS} : $Z_{DV1}^{\overline{MS}} = C^{aver} Z_{DV1}^{aver}$

Z_V, Z_A $\beta = 1.95$ $24^3 \times 48$ 

Z_V, Z_A

$\beta = 2.10$

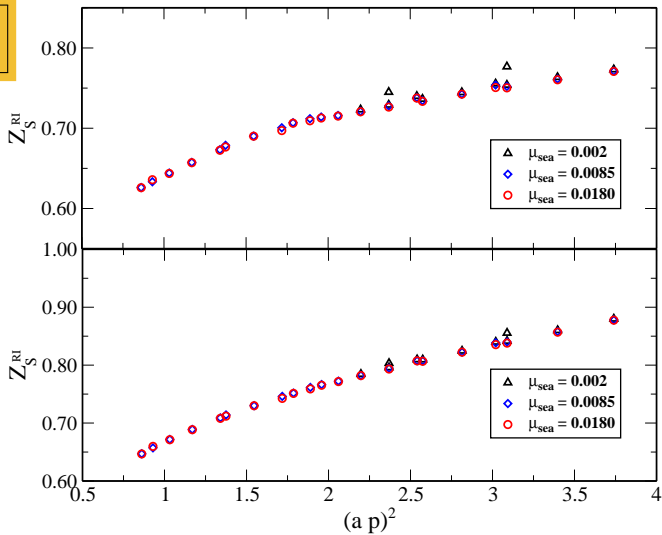
$32^3 \times 64$



Pion pole for Z_S

$\beta = 1.95$

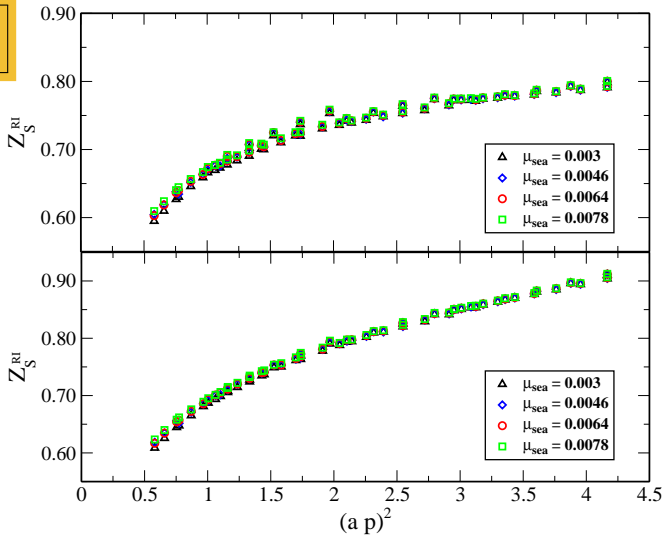
$24^3 \times 48$



Pion pole for Z_S

$\beta = 2.10$

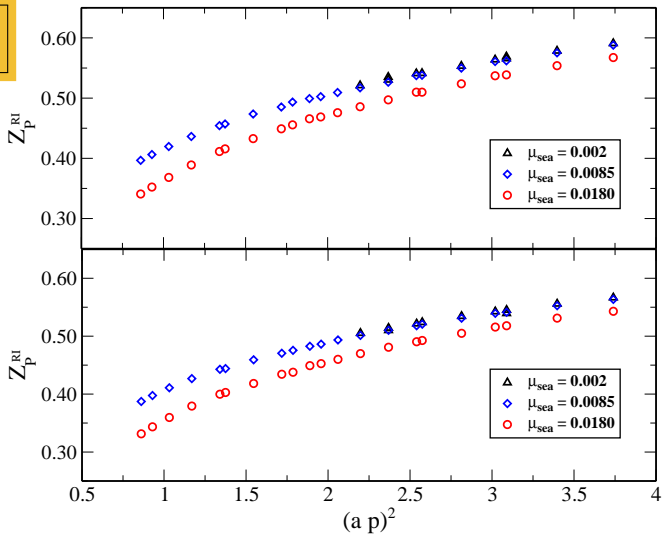
$32^3 \times 64$



Pion pole for Z_P

$\beta = 1.95$

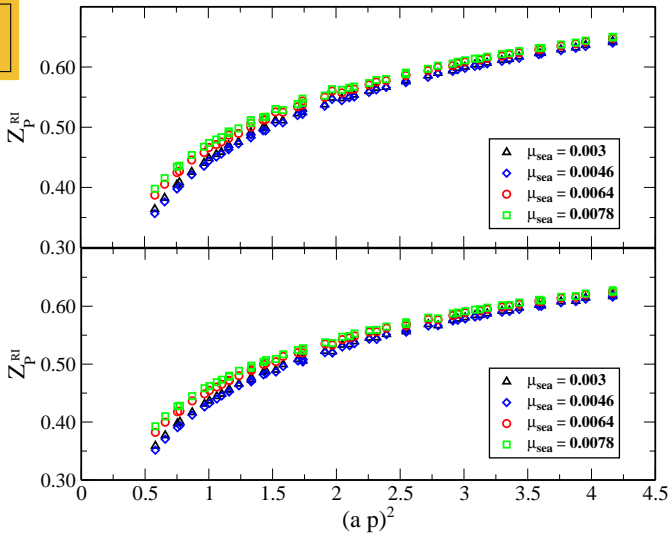
$24^3 \times 48$



Pion pole for Z_P

$\beta = 2.10$

$32^3 \times 64$



Conversion to $\overline{\text{MS}}$

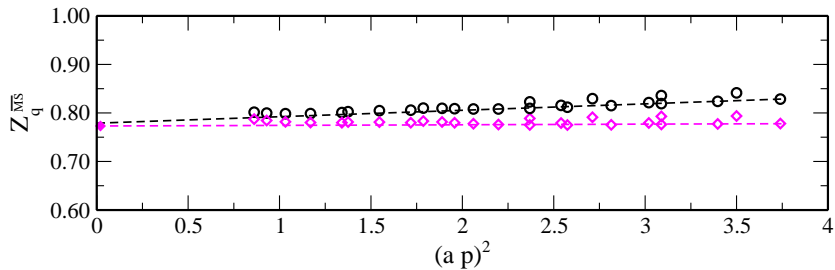
ultra-local operators: 2-loop formulae

one-derivative operators: 3-loop formulae

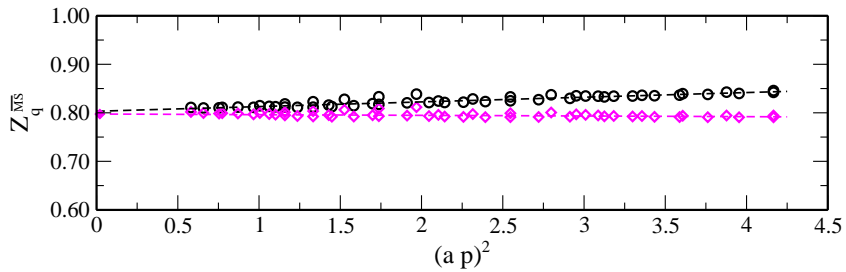
★ bare coupling

$$\beta = 1.95$$

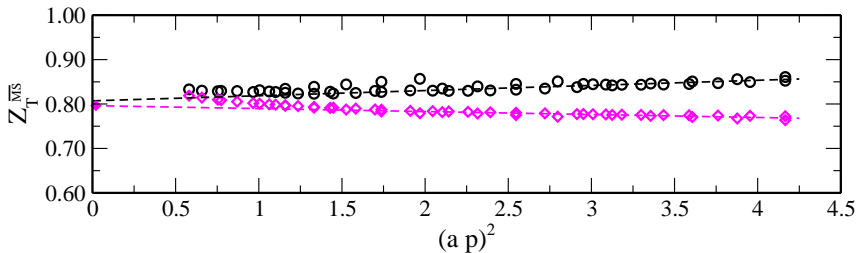
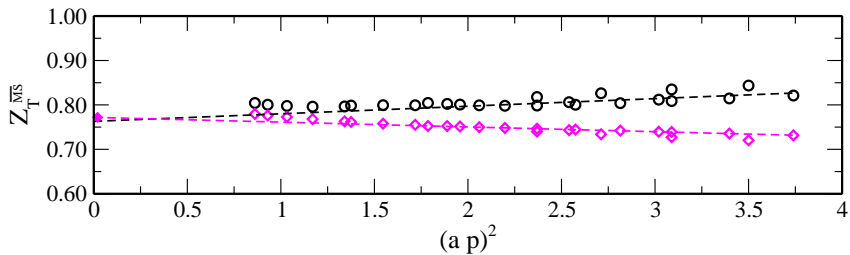
$$Z_a^{\overline{\text{MS}}}(2\text{GeV})$$



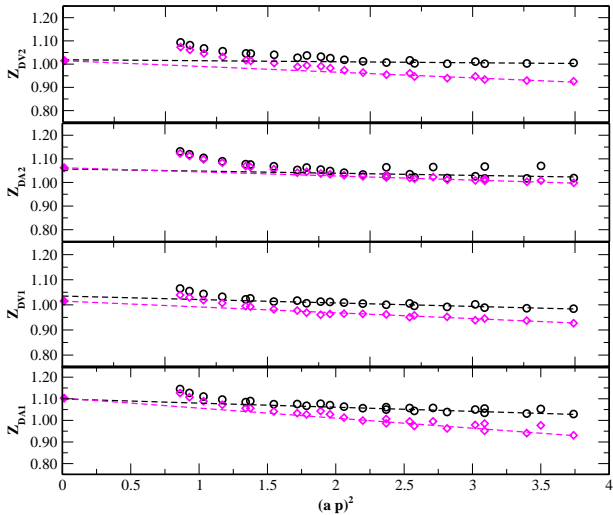
$$\beta = 2.10$$



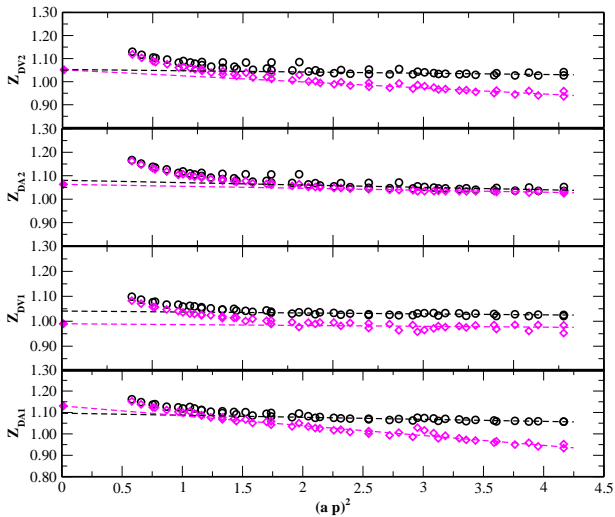
$$Z_T^{\overline{\text{MS}}}(2\text{GeV})$$



$$Z_{\text{DV}}^{\overline{\text{MS}}}(2\text{GeV}), Z_{\text{DA}}^{\overline{\text{MS}}}(2\text{GeV})$$



$$Z_{DV}^{\overline{\text{MS}}}(2\text{GeV}), Z_{DA}^{\overline{\text{MS}}}(2\text{GeV})$$



Fitting ($\overline{\text{MS}}$ at 2GeV)

Very Preliminary Results

| β | Z_V | Z_A | Z_T | Z_q |
|---------|----------|----------|----------|-----------|
| 1.95 | 0.611(9) | 0.757(3) | 0.773(2) | 0.7719(9) |
| 2.10 | 0.661(5) | 0.771(2) | 0.796(3) | 0.798(3) |

| β | Z_{DV1} | Z_{DV2} | Z_{DA1} | Z_{DA2} |
|---------|-----------|-----------|-----------|-----------|
| 1.95 | 1.015(6) | 1.015(9) | 1.103(9) | 1.062(3) |
| 2.10 | 0.990(9) | 1.052(9) | 1.13(1) | 1.063(3) |

Systematic errors: Not computed yet

Summary

- $\mathcal{O}(a^2)$ subtraction are crucial
- Mass dependence is insignificant after θ -averaging

Future Work

- Extend momentum range
- Explore volume dependence
- Pion pole subtraction for Z_S and Z_P
- Apply criterion for momentum selection (each operator separately)
- Systematic and statistical errors

THANK YOU