

UNIVERSITÄ SEIT 1386

# Breaking black-hole uniqueness at supermassive scales

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#### **Black-hole uniqueness**

- A central prediction of GR is that <u>all</u> vacuum black holes are described exactly by the Kerr metric
- Black holes all have the same spacetime structure, rescaled only by their mass
- The same theoretical model describes astrophysical objects ranging, at least, <u>10 orders of magnitude in mass</u>, from  $M \sim M_{\odot}$  to  $M \sim 10^{10} M_{\odot}$
- Any deviation from Kerr would signal <u>new physics</u>



Refer to Thomas' talk on motivations to look for new physics in gravity



#### Where can we potentially observe new physics?

- Cosmological scales
- <u>Strong-gravity regime</u>

New physics should manifest only in these regimes, and agree with GR in all other regimes where it is well-tested (up to experimental accuracy).

Finding an alternative model fulfilling these requirements is highly non-trivial!



# Scalarization

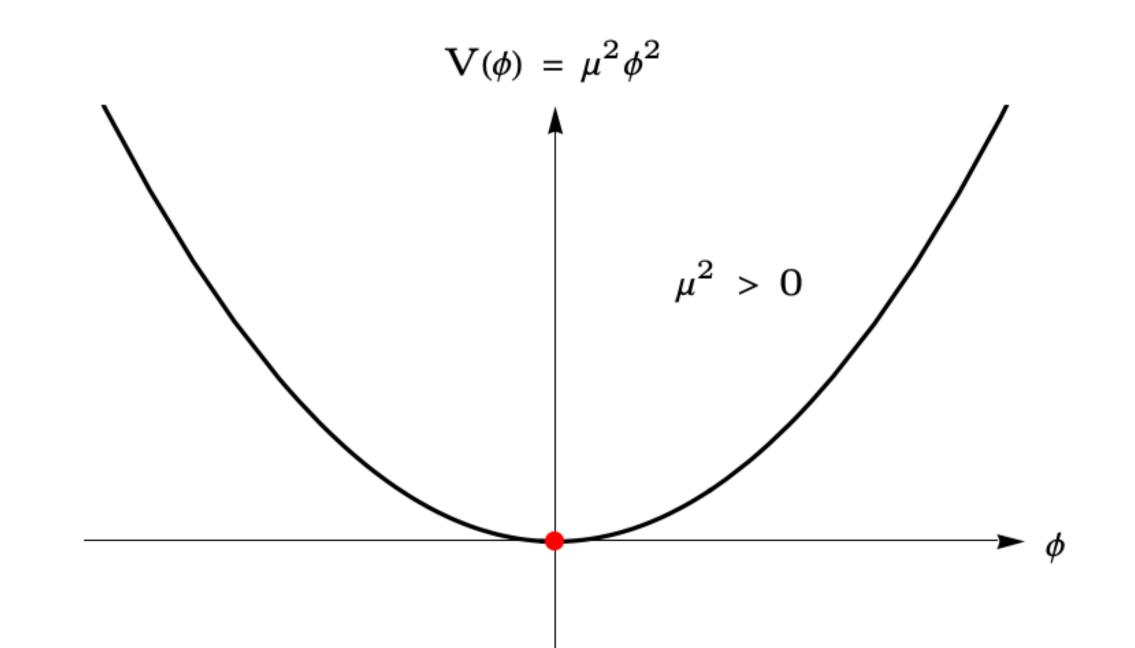
#### Scalarization

- One way to excite new physics <u>exclusively</u> in certain regimes is through scalarization
- Scalarization occurs in gravitational theories obeying the following conditions: 1. There is (at least) one new scalar field
- - 2. The theory allows for (at least) two different branches of solutions:
    - (i) a scalar-free solution which coincides exactly with GR
    - (ii) a scalarized solution, different from those of GR, which is excited only in certain regimes

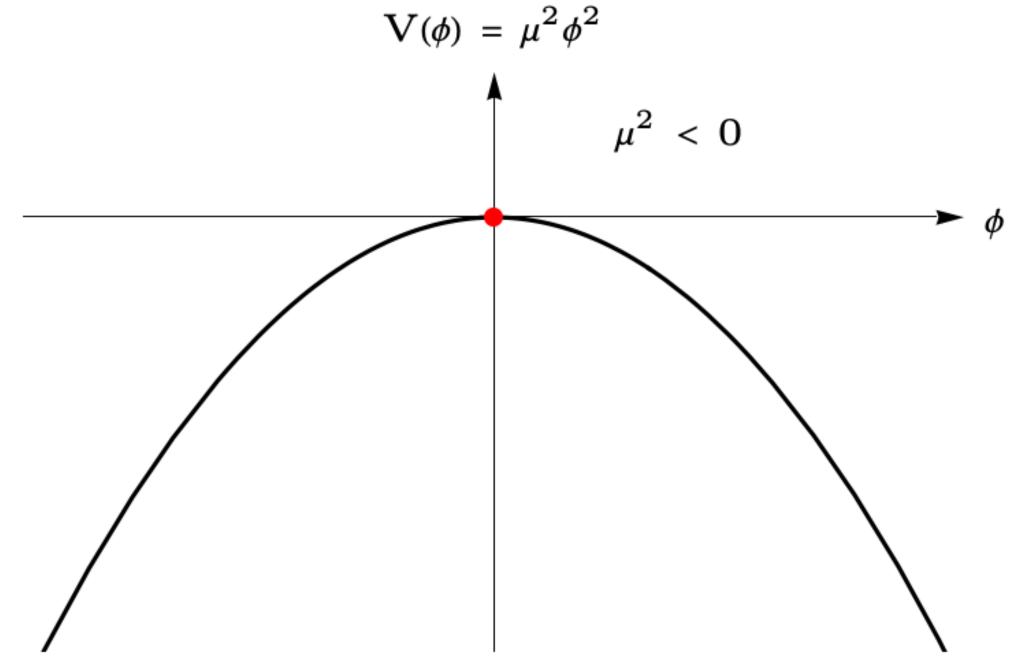
Damour et. al, 1993 Doneva et. al, 2017 Antoniou et. al. 2017 Silva et. al, 2017

. . .

#### **Tachyonic Instabilities** Prelude to scalarization

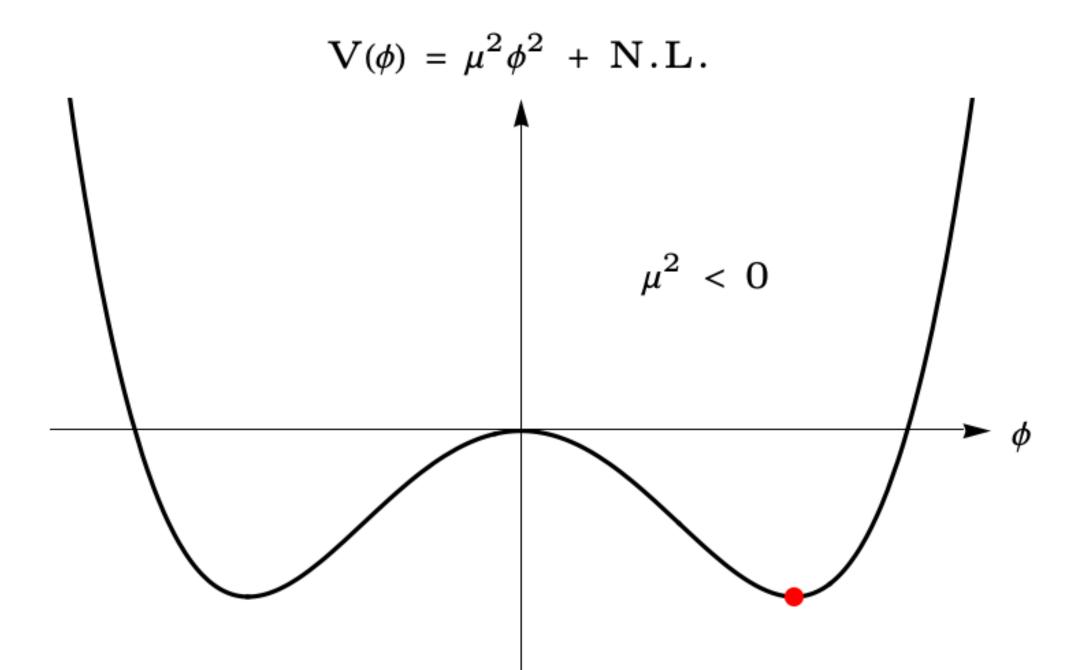


#### **Tachyonic Instabilities Prelude to scalarization**





#### **Tachyonic Instabilities Prelude to scalarization**





#### Scalarization

• New scalar,  $\phi$ , that obeys the following Klein-Gordon "wave" equation  $\Box \phi = \mu_{\text{eff}}^2 \phi$ 

- Suppose  $\mu_{\rm eff}^2 \ll 0$  near a compact object. The scalar becomes a tachyon near the compact object, inducing an instability
- Then new physics can be excited near compact objects in certain regimes, even if in all others the scalar is dormant (GR +  $\phi$  = 0)

where  $\mu_{\text{eff}} \equiv \mu_{\text{eff}}(x)$  is a (potentially) spacetime dependent effective mass.

#### Scalarization An example

concept:

$$S = \frac{1}{16\pi} \int d^4x \sqrt{-g}$$

• The scalar field equation is

 $\Box \phi = (\mu^2 - \alpha \mathscr{G}) \phi \equiv \mu_{\text{eff}}^2 \phi$ 

The field equations allow GR vacuum solutions together with  $\phi = 0$ 

#### Consider the following theory belonging to the Horndeski class as a proof-of- $\left(R - (\partial \phi)^2 - \mu^2 \phi^2 + \alpha \phi^2 \mathscr{G}\right)$ $\mathscr{G} = R^2 - 4R_{\alpha\beta}R^{\alpha\beta} + R_{\alpha\beta\mu\nu}R^{\alpha\beta\mu\nu}$ $[\alpha] = \text{length}^2$

#### Scalarization of black holes How does the scalar behave near a Schwarzschild black hole?

- The closer we are to the horizon ( $r_H = 2M$ ), the larger the GB curvature  $\mathcal{G} = \frac{48M^2}{r^6} > 0$
- The smaller the black hole, the larger the curvature

 $\mu_{eff}^2 = \mu^2 - \alpha \mathscr{G} < 0$  in its vicinity

 $\mathscr{G}(r_H) \propto 1/M^4$ 

• If the black hole is small enough (i.e., the curvature sufficiently large), we can have

• In these conditions there is a tachyonic instability around sufficiently small black holes

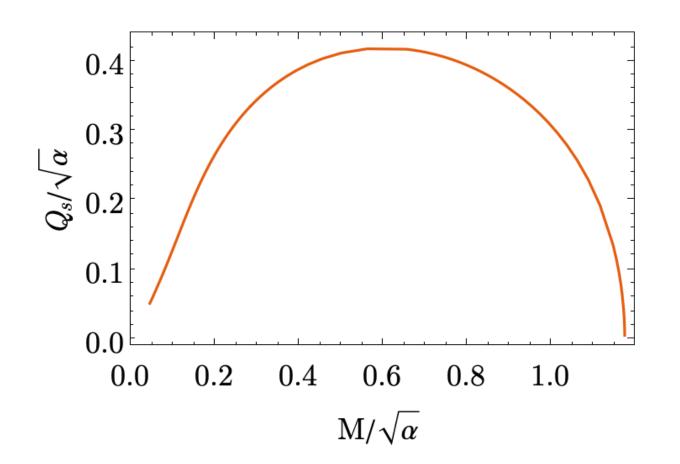


#### Scalarization of black holes

Scalarization occurs when

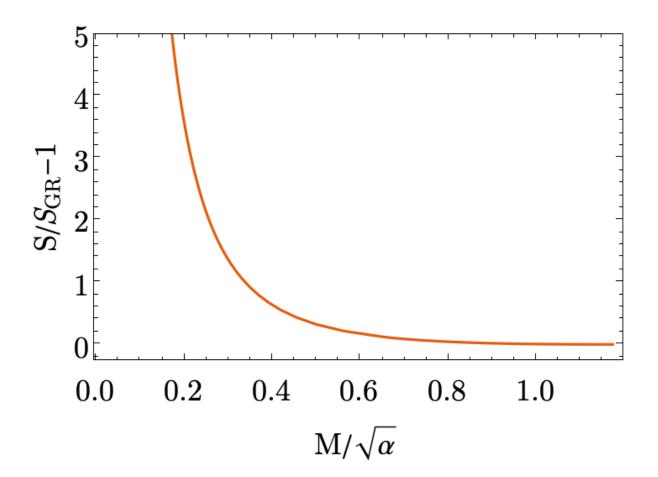
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hole uniqueness is lost in this regime!



$$M \sim \sqrt{\alpha}$$

# • A new branch of scalarized black holes branches of Schwarzschild. Black-



#### Scalarization of black holes Which black holes scalarize?

with masses  $M \sim \sqrt{\alpha}$  scalarize)

• Can  $\alpha$  be arbitrarily large? The answer is "no"



#### • The coupling constant $\alpha$ defines the scale of new physics (only black holes)

#### Scalarization of black holes Which black holes scalarize?

- Neutron stars are also highly compact objects, and therefore susceptible to scalarization
- However, binary pulsar data we have is in great agreement with GR • The scale of new physics can be at most  $\sqrt{\alpha} \sim \mathcal{O}(km)$  [see e.g. Doneva et.
- al, 2112.03869]
- We can only potentially see new physics in solar-mass black holes  $M_{\odot} \sim 1 \mathrm{km}$

# Breaking black-hole uniqueness at supermassive scales

#### **Supermassive scales**

- Supermassive scales: black holes with  $M \gtrsim 10^5 M_{\odot}$
- new physics
- SMBHs have very low horizon curvature
- theory must be of order of the gravitational radius of the SMBH
- be insensitive to new physics



• Common expectation: the higher the curvature, the more likely we are to see

Intuitively, to significantly impact SMBHs, new length scales/couplings of our

Experiments probing larger black holes such as LISA or the EHT would (likely)



### Scalarizing supermassive black holes

- Is it really impossible to have new physics affect SMBHs, while the weak-field limit, neutron stars etc... all remain well-described by GR?
- For SMBHs to scalarize in some model, it must be that

- $\Box \phi = \mu_{\rm eff}^2 \phi$
- $\mu_{\text{eff}}^2 \sim \begin{cases} \text{significantly negative around supermassive black holes} \\ \text{positive or possibly negative but very small in other regimes} \end{cases}$

# **Scalarizing supermassive black holes** We considered the following expression for the effective mass $\mu_{\text{eff}}^2 = -\alpha_1 \,\mathcal{G} + \,\alpha_2^3 \,\mathcal{G}^2$

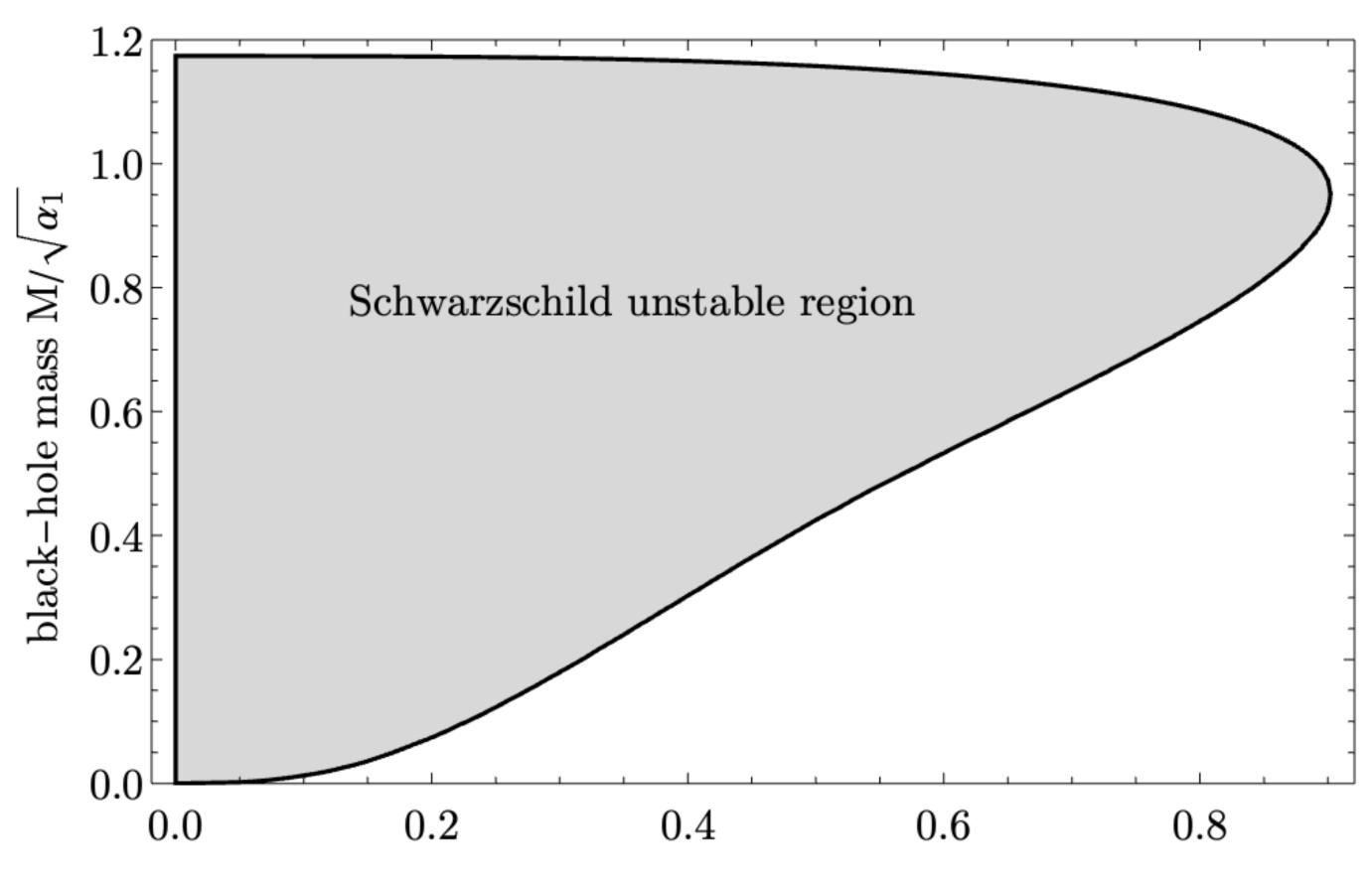
- the effective mass
- possibly leading to scalarization

• When  $\mathscr{G} \gg 1$ , the second term dominates giving a positive contribution to

• When curvatures are intermediate, there is an interplay of the two terms,

• When  $\mathscr{G} \ll 1$ , so is the effective mass, and we do not expect instabilities

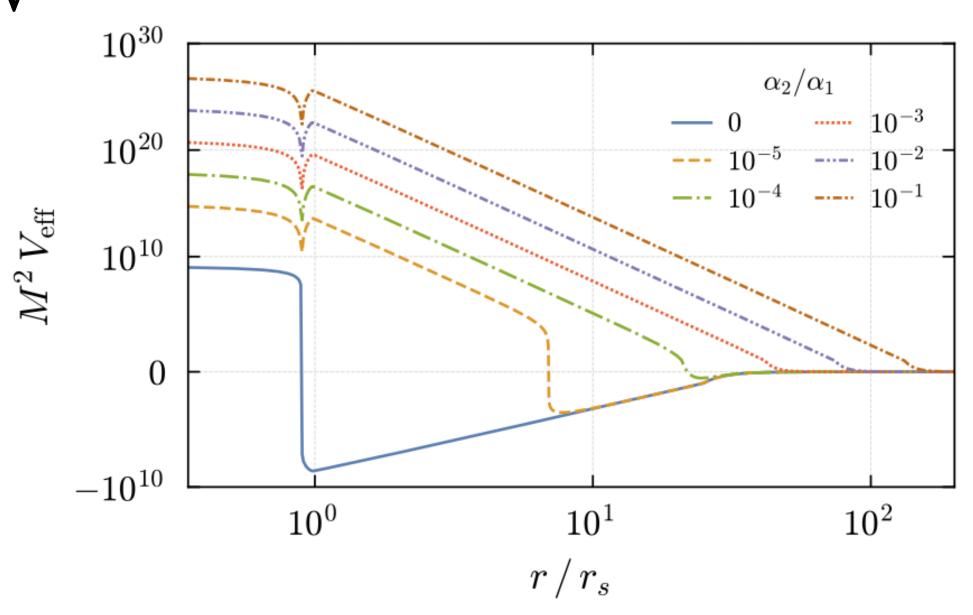
#### **Scalarizing supermassive black holes** Instability of the Schwarzschild black hole



coupling ratio  $\alpha_2/\alpha_1$ 

#### **Evading scalarization of other astrophysical objects** Neutron stars

- We do not expect scalarization to occur in neutron stars if both couplings are taken to be in the supermassive scale (curvature is very high)
- Example: Effective potential of scalar perturbations for a neutron star with  $M \approx 1.4 M_{\odot}$ ,  $r_s \approx 11 {\rm km}$ , and taking  $\sqrt{\alpha_1} \sim 10^6 M_{\odot}$



$$\frac{d^2u}{dr_*^2} + \left(\omega^2 - V_{\text{eff}}\right)u = 0$$



)

#### Scalarizing supermassive black holes The model

So far we have only considered the scalar-field equation

 $\Box \phi = (-\alpha_1 \mathscr{G})$ 

What reasonable model could produce such an equation?

$$(1+\alpha_2^3 \mathscr{G}^2)\phi \equiv \mu_{\text{eff}}^2\phi$$

#### Scalarizing supermassive black holes The model

Consider the following theory

$$S = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left[ \frac{1}{16\pi} \int d^4x \sqrt{-g} \right]$$

By introducing an auxilliary scalar  $\psi$ 

$$S = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left[ R - (\partial x) \right]$$

- The model belongs to a bi-scalar extension of Horndeski

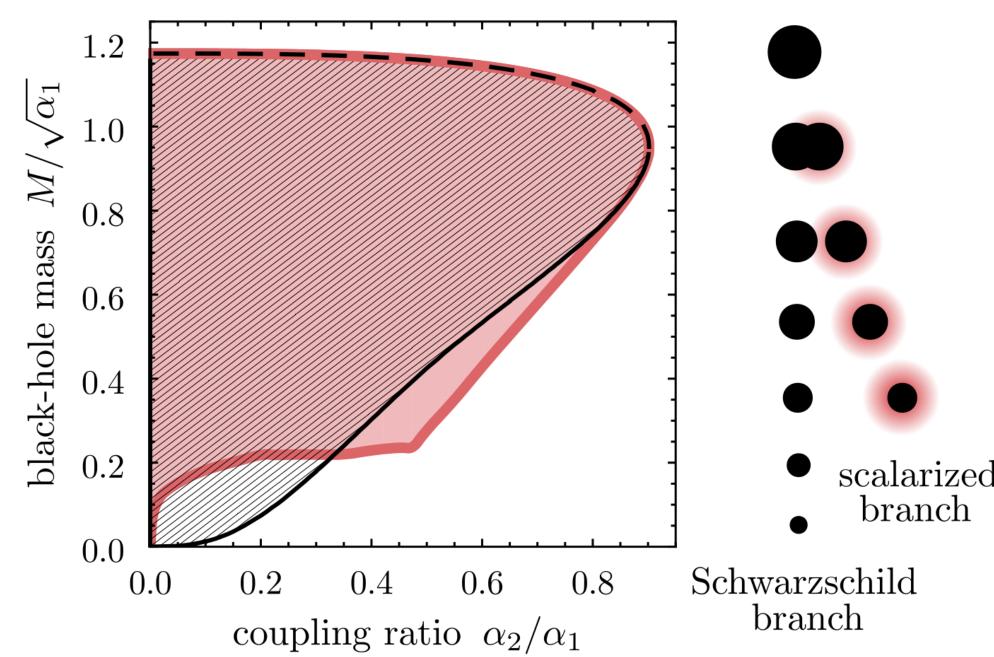
 $\left[R - (\partial \phi)^2 + (\alpha_1 \mathcal{G} - \alpha_2^3 \mathcal{G}^2)\phi^2\right]$ 

 $(\partial \phi)^2 + \alpha_1 \phi^2 \mathcal{G} - 2\alpha_2^3 \left( \psi \mathcal{G} - \frac{\psi^2}{2} \right) \phi^2$ 

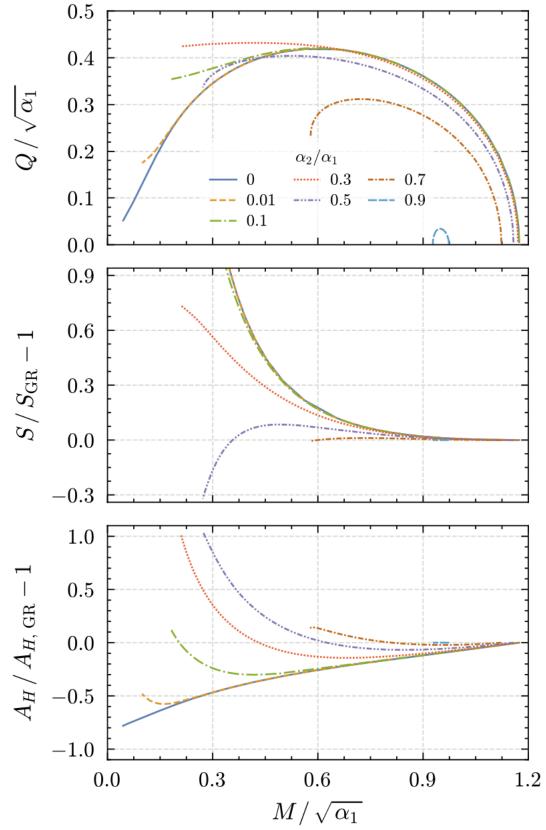
All equations of motion are second-order in all involved fields,  $g_{\mu\nu}$ ,  $\phi$  and  $\psi$ 

#### **Scalarizing supermassive black holes Scalarized black holes**

 We have numerically solved the field equations in spherical symmetry to obtain scalarized black hole solutions



# 0.5



## **Cosmological stability of the theory**

- during inflation
- that prevents inflation [Charmousis et. al, 1903.02399]
- In our case  $\mu_{eff}^2 = -\alpha_1 \mathcal{G} + \alpha_2^3 \mathcal{G}^2$ . The new term dominates during inflation and prevents any instability
- Compatible with all cosmological evolution described by GR

• In the standard scalarization scenario  $\mu_{eff}^2 = -\alpha_1 \mathscr{G}$ , which is highly negative

• Assuming the Universe starts in a GR vacuum, there is catastrophic instability

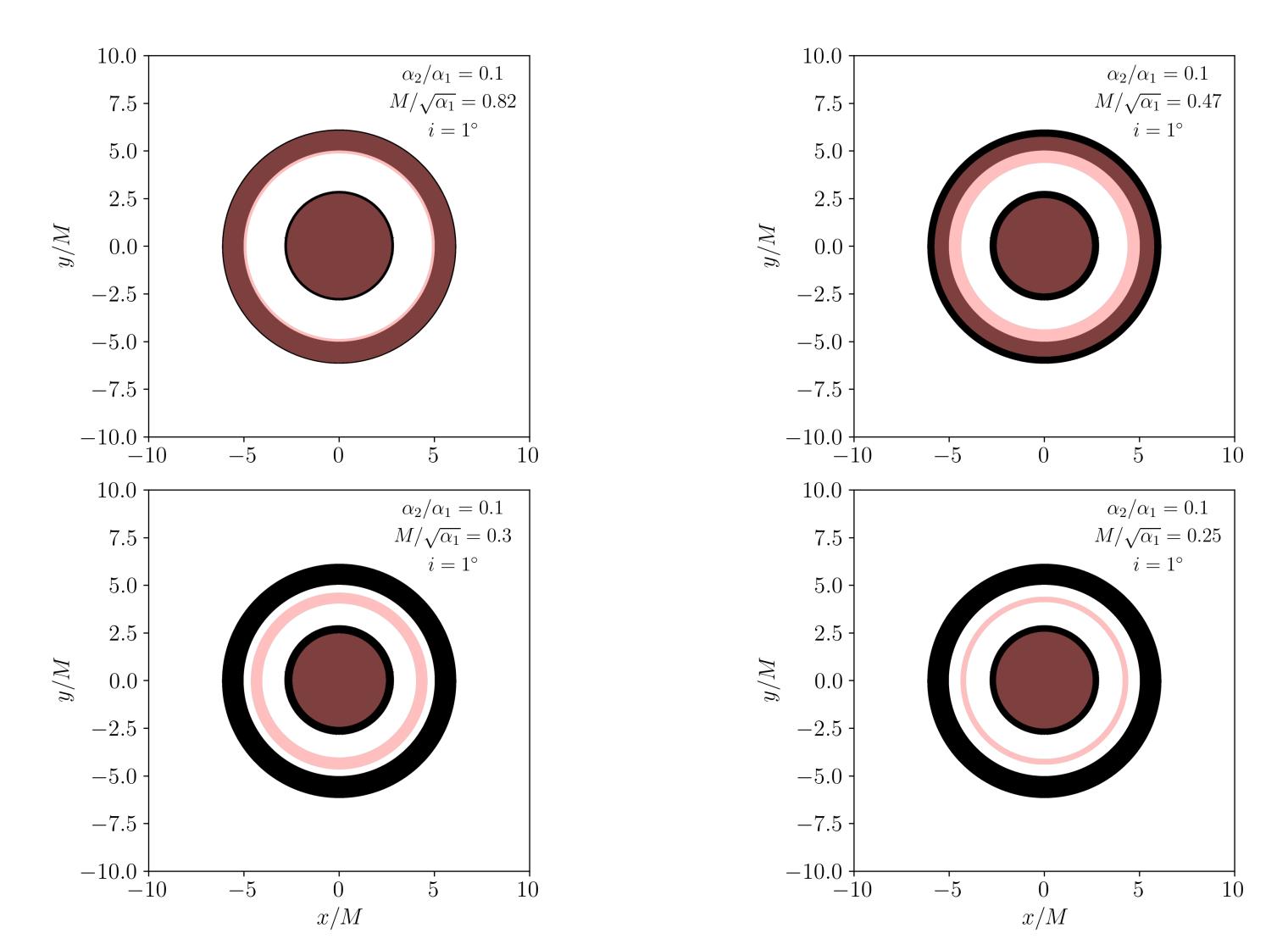
# Conclusions

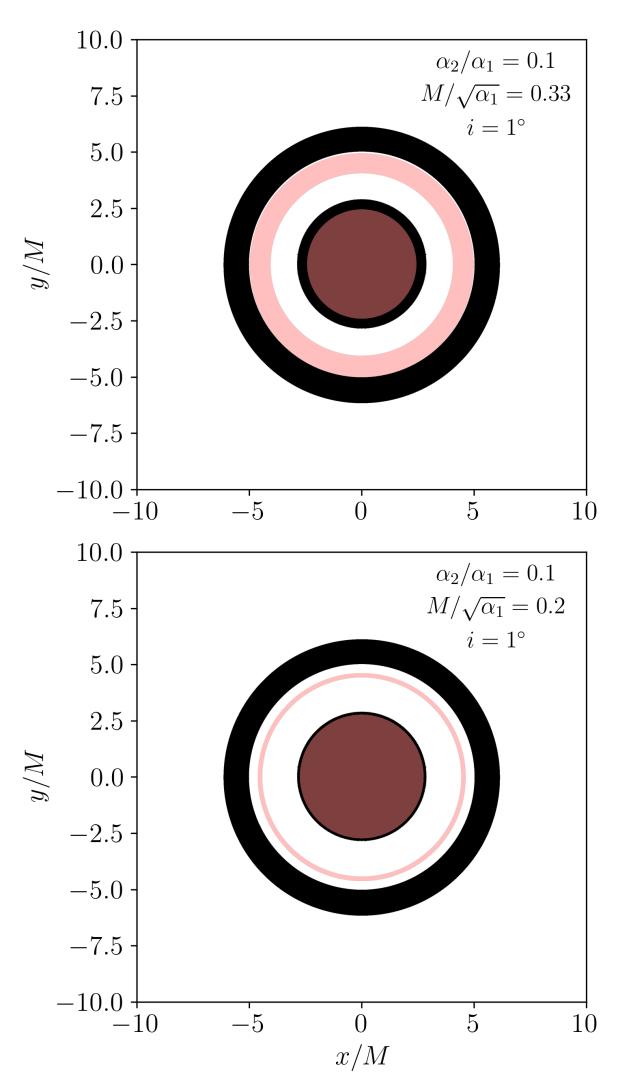
#### Conclusions

- We introduced a novel scalarization model where new physics manifests only for supermassive black holes
- Consistent cosmological evolution
- To constraint the theory and probe supermassive scales we need experiments targeting those scales
- Potential observational consequences for: EHT, LISA, pulsar-timing arrays...
- This theory serves as a proof-of-concept that supermassive black holes can significantly deviate from Kerr



#### **Ongoing and Future Work** Black-Hole Imaging





## **Ongoing and Future Work**

- 0.20
- 0.15• Rotating black holes
  - ·~ 0.10
    - 0.05
    - 0.00

SMBH scalarization from an EFT point of view



