

# COSMOLOGICAL BOUNDS ON THREE SCENARIOS OF AXION-LIKE PARTICLES AND CONDENSATES FROM NON-EQUILIBRIUM QFT

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Based on work in progress with R. Passante & L. Rizzuto

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# GOALS AND AN OUTLINE



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## GOALS AND A SUMMARY

We consider two methods for non-equilibrium QFT



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$$S = \int d^4x \sqrt{g} \left[ \frac{1}{2} \partial_\mu \Phi \partial^\mu \Phi - \frac{1}{2} m_\phi^2 \Phi^2 - \frac{\lambda_\phi}{4!} \Phi^4 + K_\chi + g_\chi \Phi O_\chi \right]$$



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## GOALS AND AN OUTLINE

We take  $\Phi = \varphi + \phi$  with  $\varphi(x) = \langle \Phi \rangle$



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We take  $\Phi = \varphi + \phi$  with  $\varphi(x) = \langle \Phi \rangle$  and extend the work done by [Ai, WY., Beniwal, A., Maggi, A. \(2024\) \[1\]](#) and [S. Cao and D. Boyanovsky PRD 107, 063518 \(2023\)\[2\]](#) in both the formalism and considering three ALPs models in a freeze-in scenario:



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1. Photophilic  $\rightarrow K_\chi = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} + K_f$  and  $O_\chi = -\frac{1}{4}\tilde{F}F$



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3. QCD  $\rightarrow K_\chi = -\frac{1}{4}G_{\mu\nu}^a G^{a\mu\nu}$  and  $O_\chi = -\frac{1}{4}\tilde{G}G$



## GOALS AND AN OUTLINE

From the quantum EoMs of  $\varphi$  and  $\Delta\phi$ , along with the help of the Liouville-Von Neumann equation, we get the following fully relativistic coupled Boltzmann equations:



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Such complete equations can be useful to handle a more complex case



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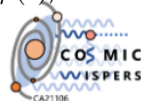
Objective: Get the relic abundances and  $\Delta N_{\text{eff}} (\varphi = \varphi(t))$ .



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# BOLTZMANN TRANSPORT EQUATIONS OF ALPS AND CONDENSATES



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$$\frac{df}{dt} - H k \frac{df}{dk} + M(t) \dot{M} \frac{df}{dk} = C[f]$$

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- We are solving them numerically by MicrOMEGAs.



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## PRELIMINARY RESULTS

- As one could have expected from qualitative arguments, the presence of the condensate is not significant for QCD ALPs  $\rightarrow m_a^2 \sim m_{a0}^2 + c_g^2 m_{QCD}^2$



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(this last is the same result obtained in the recent preprint [Jain, M., Maggi, A., Ai, W. Y., & Marsh, D. J. \(2024\) \[4\]](#))



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# SUMMARY AND CONCLUSIONS



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1. We have developed a full formalism for studying the evolution equation of the ALP distribution function, for both the "particle" and the "condensate" components.



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4. Adopt our formalism to study the distribution function of the topological defects of QCD axion and use it as an alternative method to (numerically) evaluate the axion spectrum and its spectral index  $q$ .



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***Thanks  
for your attention!***



# ADDITIONAL SLIDES



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## SECOND CASE: PHOTOPHOBIC ALPS

$$\Gamma_{2PI}[\psi, \Delta_\phi, \Delta_\gamma] = S[\psi] + i/2 \text{Tr} \ln \Delta_\phi^{-1} + i/2 \text{Tr} \ln \Delta_f^{-1} + i/2 \text{Tr}[G_\phi^{-1} \Delta_\phi] + i/2 \text{Tr}[G_f^{-1} \Delta_f] \quad (1)$$

where we have explicitly

$$\Delta_{f,mn} = \langle \bar{\Psi}_m(x) \Psi_n(y) \rangle \quad (2)$$

$$G_{f,mn}^{ab,-1} = ic^{ab}(i\gamma^\mu \partial_\mu - m_f - g_{aff} \partial_\mu \psi \gamma^\mu \gamma_5) \quad (3)$$



## QUANTUM EOMS

Precisely, we get from both the two methods the following quantum EoMs for  $\varphi(x) = \langle \phi \rangle$  and  $\Delta_\phi(x, y) = \langle T\phi(x)\phi(y) \rangle$ :

$$-(\square + \tilde{m}_\phi^2)\Delta^{ab}(x_1, x_2) - \sum_c c \int d^4x_3 \Pi_\phi^{ac}(x_1, x_3)\Delta^{cb}(x_3, x_2) = ic^{ab}\delta(x_1 - x_2) \quad (4)$$

$$(\square + \tilde{m}_\phi^2)\varphi + \frac{\lambda_\phi}{2}\Delta_\phi^{++}(x, x)\psi - \frac{\delta\Gamma_2}{\delta\phi^+}|_{\phi^+=\phi^-=\psi} + g_\chi\Delta_{\chi, \tilde{F}F} = 0 \quad (5)$$

With both methods, we assume the SM particles in our models are at thermal equilibrium.





EXAMPLE OF  $\Gamma_2$ 

$$\Gamma_2 \sim -i \left( \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \right)$$

FIG. 6.  $\Gamma_2$  for the photophilic case with the relevant 2PI diagrams