Axions, Neutrinos, and Rare Decay Anomaly of Belle-II

Based on: Hati, Leite, NN, Valle, 24-08.00060 [hep-ph] 19th Patras Workshop on Axions, WIMPs and WISPs

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17-09-2024



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Question: Can we explain neutrino mass, strong CP problem and Belle I anomaly from a single framework? GENERALITAT VALENCIANA Based on: Hati, Leite, NN, Valle, 24-08.00060 [hep-ph] 19th Patras Workshop on Axions, WIMPs and WISPs



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Belle-II Anomaly • The latest Belle II measurement: B (${ m {\it B}^+} ightarrow { m K^+} + { m {\it E}_{miss}}$) = $(2.3 \pm 0.5) imes 10^{-5}$

- lacksquare
- The effective Lagrangian within the LEFT approach: •

$$\mathcal{L}_{\text{LEFT}}^{bs\nu\nu} \supset \frac{8G_F V_{tb} V_{ts}^*}{\sqrt{2}} \frac{\alpha}{4\pi} \sum_{\alpha\beta} \left[(C_L^{sb\alpha\beta} + C_L^{sb,\text{SM}} \delta^{\alpha\beta}) (\bar{s}_L \gamma_\mu b_L) (\bar{\nu}_L^\alpha \gamma^\mu \nu_L^\beta) + C_R^{sb\alpha\beta} (\bar{s}_R \gamma_\mu b_R) (\bar{\nu}_L^\alpha \gamma^\mu \nu_L^\beta) \right] \\ \text{NP} \qquad \text{NP} \\ \text{Define:} \ R_{K^{(*)}}^\nu = \frac{\text{Br}(B \to K^{(*)} \nu \bar{\nu})}{\text{Br}(B \to K^{(*)} \nu \bar{\nu})_{\text{SM}}} \\ R_K^\nu = 5.3 \pm 1.7 \qquad \text{; Only Belle || Adachi et al. 2311.14647 [hep-ex]} \\ R_K^\nu = 2.93 \pm 0.90 \quad \text{; Belle || + previous results} \\ \text{Marzocca et. al. 2404.06533 [hep-ph]} \\ R_{K^*}^\nu = 1.0 \pm 1.1 \quad \text{; Belle || expt. with combination of all channels} \\ \text{Grapher of the set of$$

Adachi et al. 2311.14647 [hep-ex]

Interpreting "missing energy" as a pair of neutrinos leads to 2.7 σ deviation from the SM

Grygler et al. 1702.03224 [nep-ex]

Belle-II Anomaly

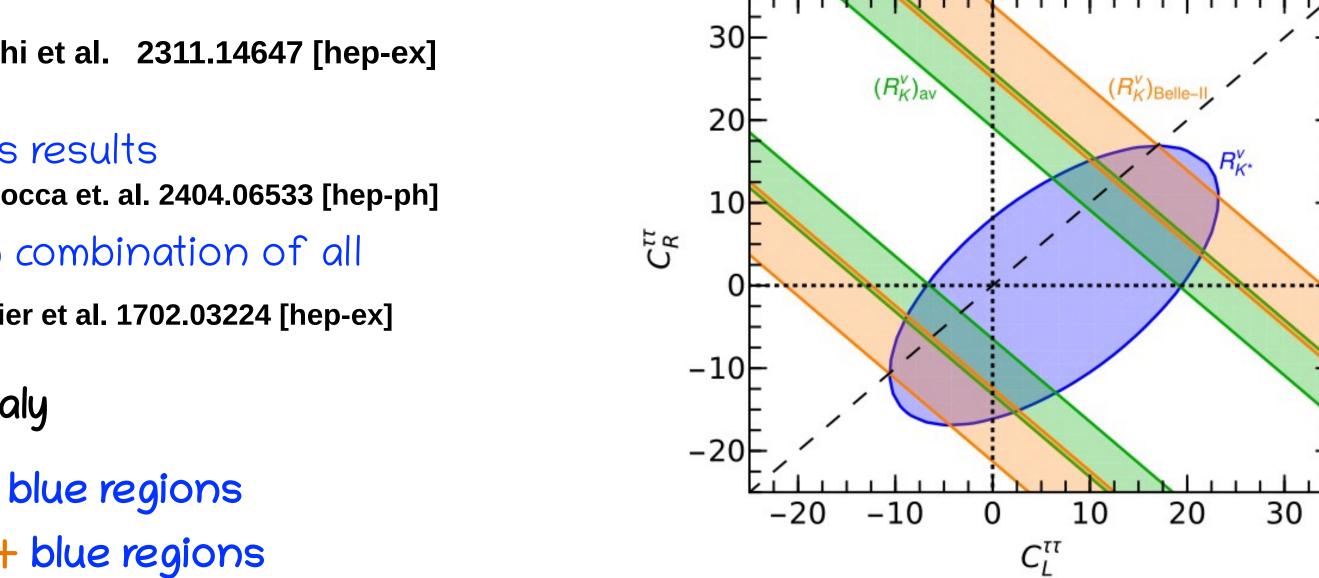
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Orange + blue regions [-12.41, -16.20]

Adachi et al. 2311.14647 [hep-ex]

Interpreting "missing energy" as a pair of neutrinos leads to 2.7 σ deviation from the SM





Framework: KSVZ-type QCD axion model

Fields content:

SM Fields				
Fields	\mathcal{SM}	$U(1)_{PQ}$	Families	
L_L	(1, 2, -1/2)	0	3	
e_R	(1,1,-1)	0	3	
Q_L	(3,2,1/6)	0	3	
d_R	(3,1,-1/3)	0	3	
u_R	(3,1,2/3)	0	3	
Φ	(1, 2, 1/2)	0	1	

Interaction Lagrangian:

$$-\mathcal{L}_{\text{Yuk.}} \supset Y_{ab}^{\Psi} \overline{\Psi_{aL}} \Psi_{bR} \sigma + \frac{Y_{ab}^{\Sigma}}{2} \Psi_{aR}^{T} C \Sigma \Psi_{bL}$$

The relevant Lagrangian for η -mediated $b \to s \nu_{\tau} \overline{\nu_{\tau}}$ transition

$$\mathcal{L}_{\eta}^{\text{mass}} = \overline{\nu_{\tau L}} \,\lambda_{33} \, b_R \, \eta^{1/3} + \overline{\nu_{\tau L}} \, \lambda_{14} \, b_R \, \eta^{1/3} + \overline{\nu_{\tau L}} \, \lambda_{14} \, b_R \, \eta^{1/3} + \overline{\nu_{\tau L}} \, \lambda_{14} \, b_R \, \eta^{1/3} + \overline{\nu_{\tau L}} \, b_R \, \eta^{1/3} \, b$$

Also see: Batra, Câmara, Joaquim, Srivastava, Valle, PRL 132 (2024) 5, 051801

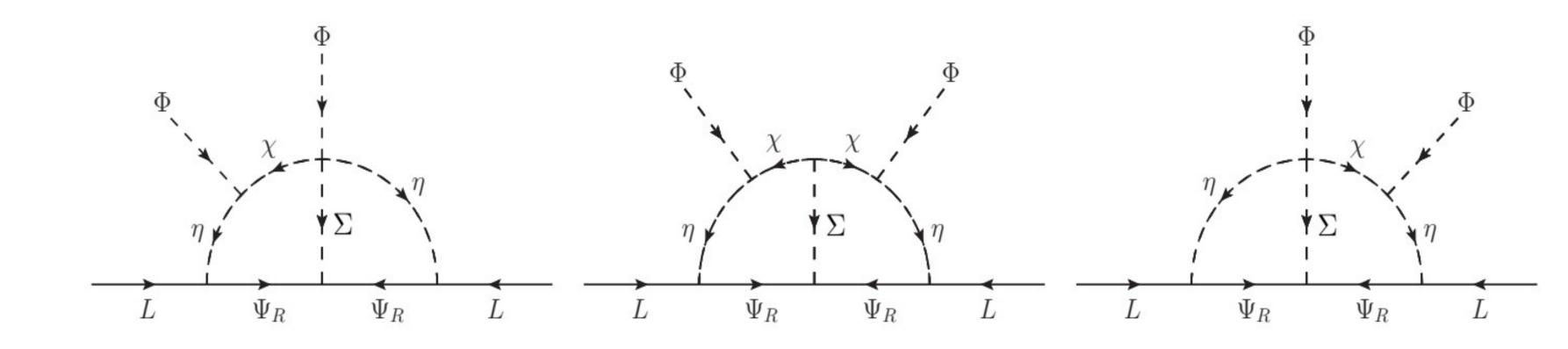
Fields	\mathcal{SM}	$U(1)_{PQ}$	Multiplicities
Ψ_L	(3, 1, -1/3)	0	2
Ψ_R	(3, 1, -1/3)	-1/2	2 Vector-like quarks
σ	(1, 1, 0)	1/2	1
η	(3,2,1/6)	-1/2	1 } Leptoquark-like
χ	(3, 1, -1/3)	-1/2	1
Σ	$(ar{6},1,2/3)$	1	1 Colored-scalars

PCM Eiolda

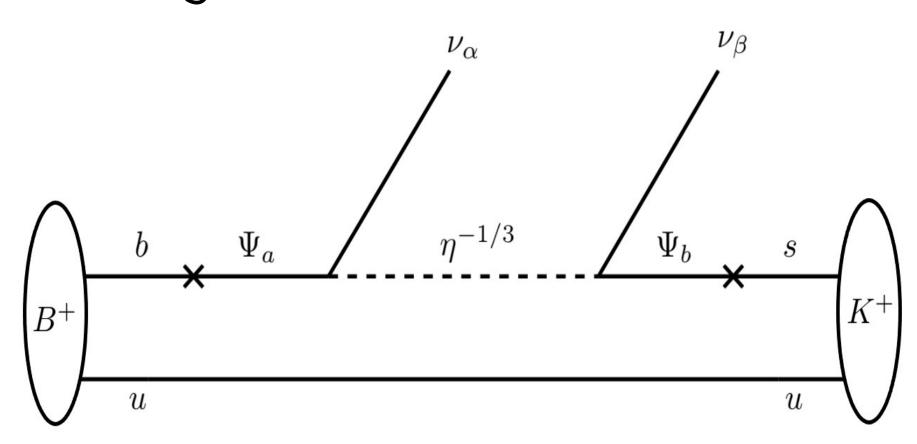
 $h_{R} + Y_{ia}\overline{L_{iL}}\,\tilde{\eta}\,\Psi_{aR} + M_{ai}^{\Psi d}\,\overline{\Psi_{aL}}d_{iR} + y_{ij}^d\,\overline{Q_{iL}}\,\Phi\,d_{jR} + \text{H.c.}$

 $\lambda_{32} s_R \eta^{1/3} + \overline{b_R} \lambda_{33}^* \nu_{\tau L} \eta^{-1/3} + \overline{s_R} \lambda_{32}^* \nu_{\tau L} \eta^{-1/3}$

Neutrino mass



The diagram for $b \to s \nu_{\tau} \overline{\nu_{\tau}}$ transition



Cont...

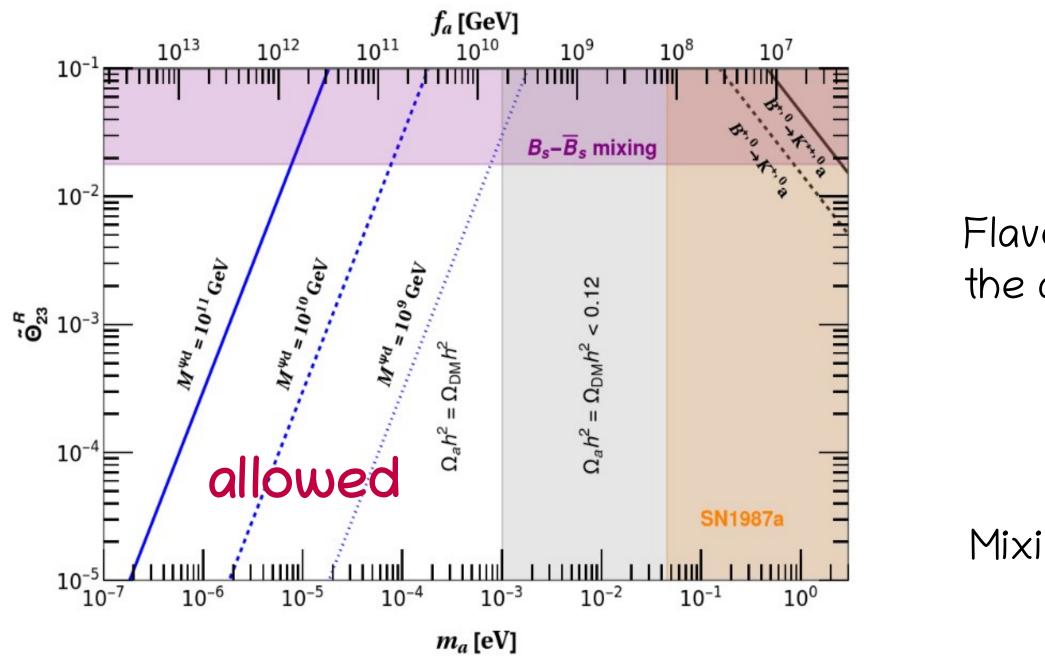
The effective Lagrangian: $\mathcal{L}_{\text{eff}}^{bs\nu\nu} \supset -\frac{\lambda_{32}^* \lambda_{33}}{2m_{\eta}^2} \left(\bar{s}_R \gamma_{\mu} b_R \right) \left(\bar{\nu}_{\tau L} \gamma^{\mu} \nu_{\tau L} \right)$

The relevant Wilson coefficient:

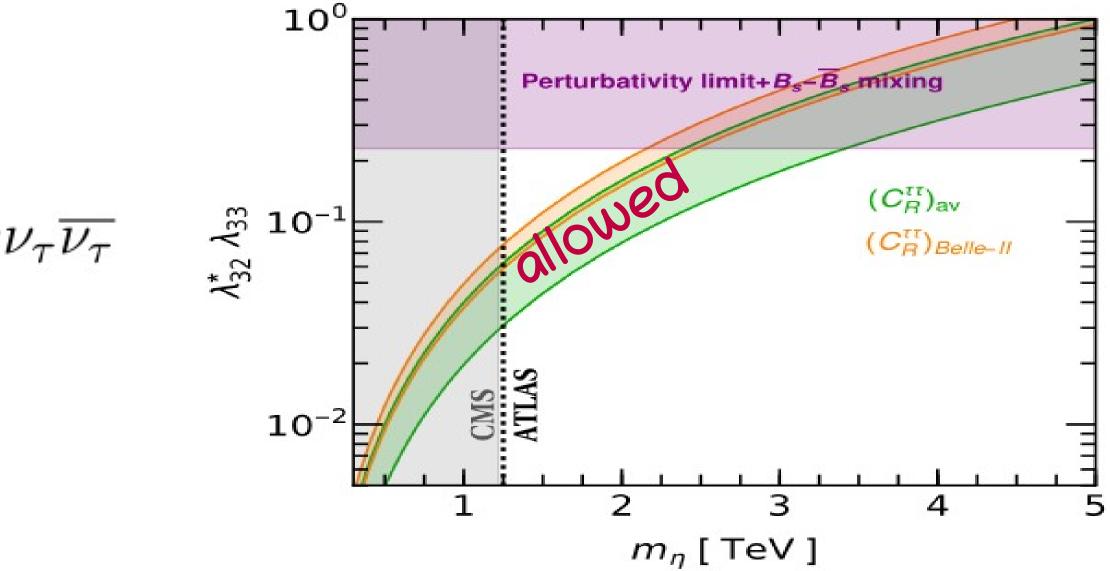
$$C_R^{sb\tau\tau} = -\left(\frac{8G_F V_{tb} V_{ts}^*}{\sqrt{2}} \frac{\alpha}{4\pi}\right)^{-1} \frac{\lambda_{32}^* \lambda_{33}}{2m_\eta^2}$$

We have
$$C_R^{sb au au} = -\left(rac{8G_FV_{tb}V_{ts}^*}{\sqrt{2}}rac{lpha}{4\pi}
ight)^{-1}rac{\lambda_{32}^*\lambda_{33}}{2m_\eta^2}$$

Allowed parameter space that could explain $b
ightarrow s
u_ au \overline{
u_ au}$



Results



Flavor-violating interaction terms between down-type quarks and the axion a

$$\mathcal{L}_{a}^{\mathrm{FV}} = \frac{\partial_{\mu}a}{f_{a}} \overline{d'_{i}} \gamma^{\mu} \left(c_{ij}^{V} + c_{ij}^{A} \gamma_{5} \right) d'_{jR}$$

Where, $c_{bs}^{V} = c_{bs}^{A} = \frac{1}{2} \mathcal{Q}_{R} \tilde{\Theta}_{23}^{R}$

Mixing parameter

$$\tilde{\Theta}^R_{23} = \frac{M_{22}^{\Psi d} M_{23}^{\Psi d}}{M_2^{\Psi^2}} = \frac{M_{22}^{\Psi d} M_{23}^{\Psi d}}{(Y_{22}^{\Psi} f_a)^2} < 0.018 \ \text{from} \ B_s - \bar{B}_s \ \text{mixin}$$



Final remarks

- The framework explains neutrino mass, strong CP problem and Belle II anomaly
- KSVZ-type QCD axion model has been adopted: \bullet

The SM is extended with a new pair of vector-like isosinglet quarks, one scalar leptoquark, and two coloured scalar mediators with non-zero PQ charges.

- Crucial: Mass-mixing of vector-like quarks with the SM quarks explain the anomaly
- $B_s \bar{B}_s$: Mixing is important to prove relevant parameter space
- The QCD axion in the model could explain the relic Dark matter abundance of the Universe

