



Ultralight Dark Matter Detection with Gravitational-Wave Interferometers

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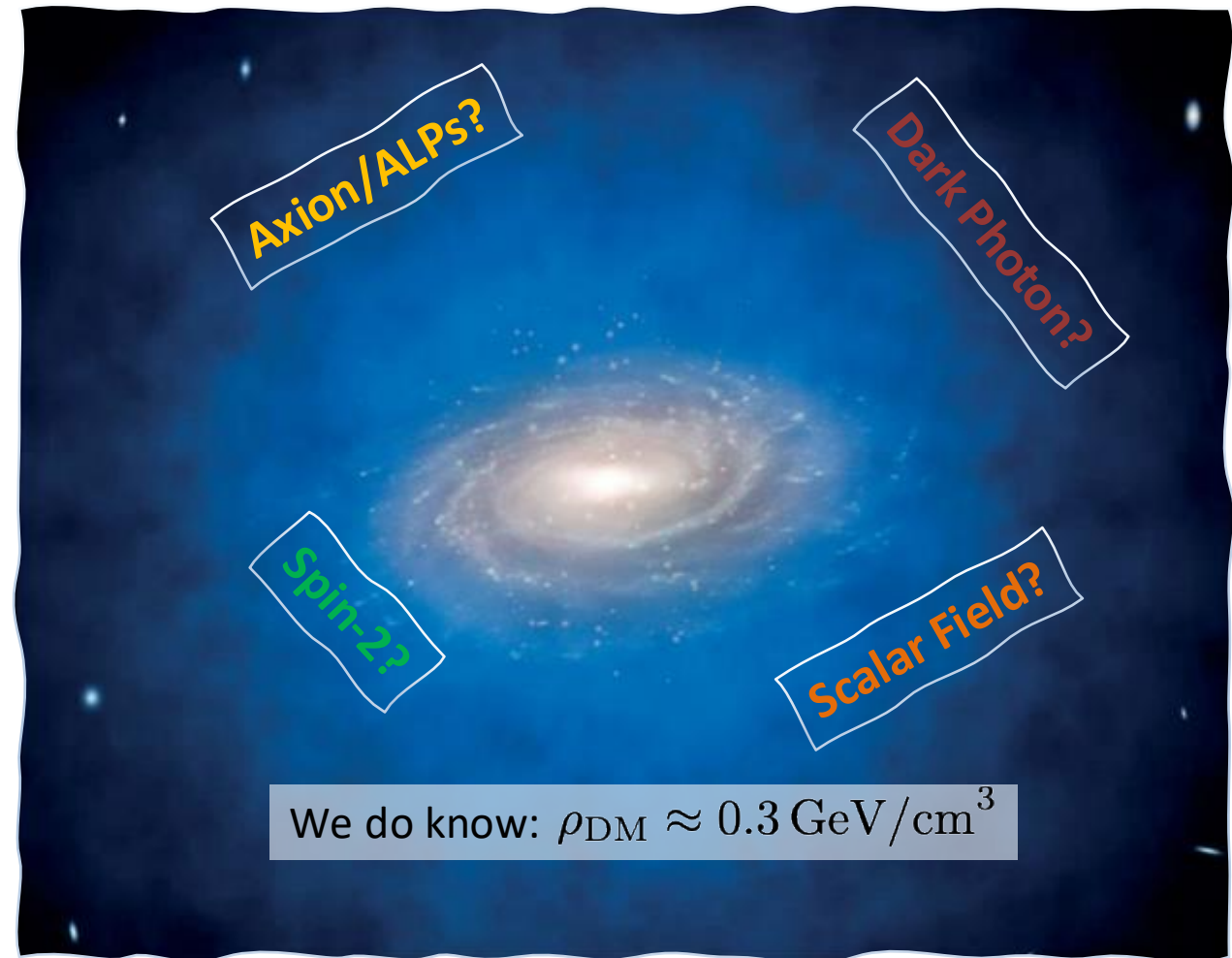
Introduction to Scalar Field Dark Matter



- Dark matter is one of the mysteries of modern physics.
- Scalar field dark matter is a theoretical candidate that could account for dark matter through low-mass fields.
- Why scalar fields?

Scalar Field Coupling

$$\mathcal{L}_{\text{int}} = \frac{\phi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} - \frac{\phi}{\Lambda_e} m_e \bar{\psi}_e \psi_e$$





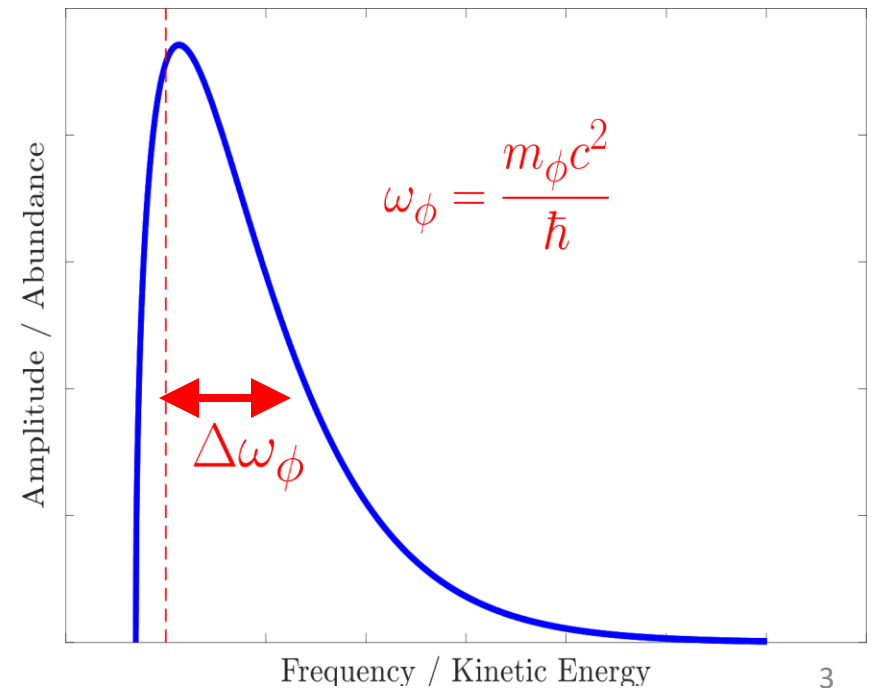
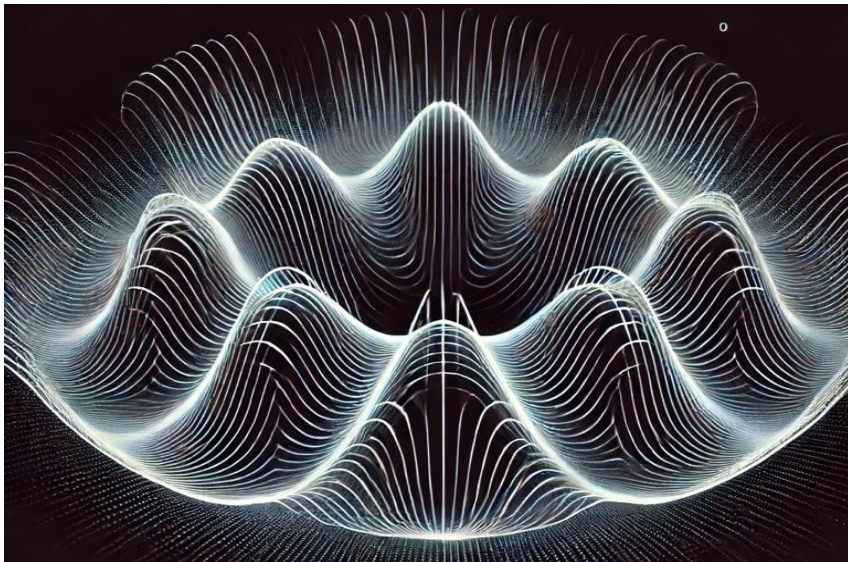
What is Scalar Field Dark Matter?



- Scalar field dark matter may have been created in the early universe.
- Coherently oscillating fields that can cause observable effects in physical systems.

$$\phi(t, \vec{r}) = \left[\frac{\hbar \sqrt{2} \rho_{\text{local}}}{m_\phi c} \right] \cos \left(\omega_\phi t - \vec{k}_\phi \cdot \vec{r} \right)$$

$$\frac{\delta \omega_\phi}{\omega_\phi} \approx 10^{-6}$$





Expected Effects of Scalar Field Dark Matter



- Scalar field dark matter affects fine structure constant and electron mass.
- Modulations in the size and refractive index of materials.

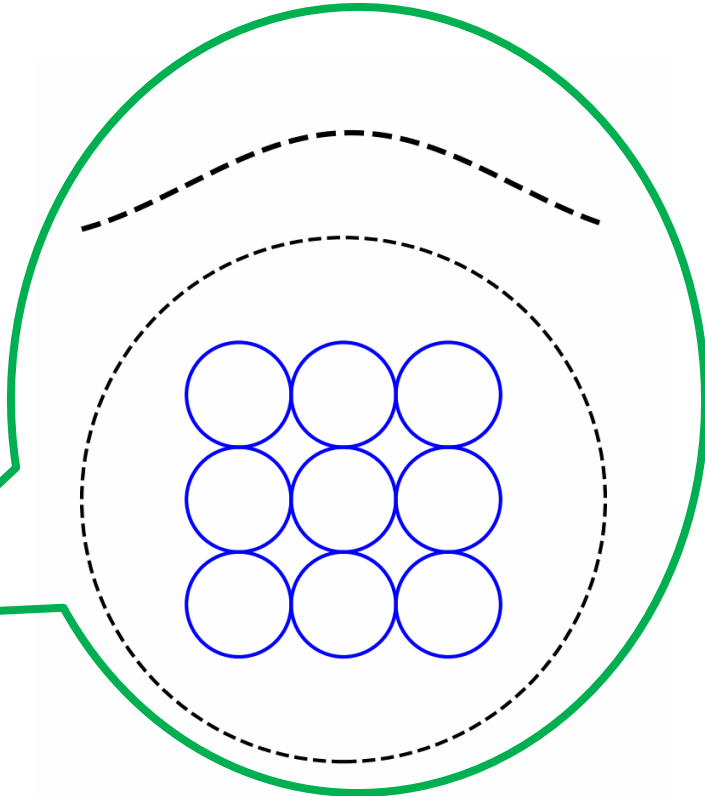
$$\mathcal{L}_{\text{int}} = \frac{\phi}{\Lambda_\gamma} \frac{F_{\mu\nu} F^{\mu\nu}}{4} - \frac{\phi}{\Lambda_e} m_e \bar{\psi}_e \psi_e$$

$$\frac{\delta\alpha}{\alpha} = \frac{\phi}{\Lambda_\gamma}$$

$$\frac{\delta m_e}{m_e} = \frac{\phi}{\Lambda_e}$$

$$\phi(t, \vec{r}) = \left[\frac{\hbar \sqrt{2} \rho_{\text{local}}}{m_\phi c} \right] \cos(\omega_\phi t - \vec{k}_\phi \cdot \vec{r})$$

$$r_0 = \frac{\hbar}{c} \frac{1}{m_e \alpha}$$

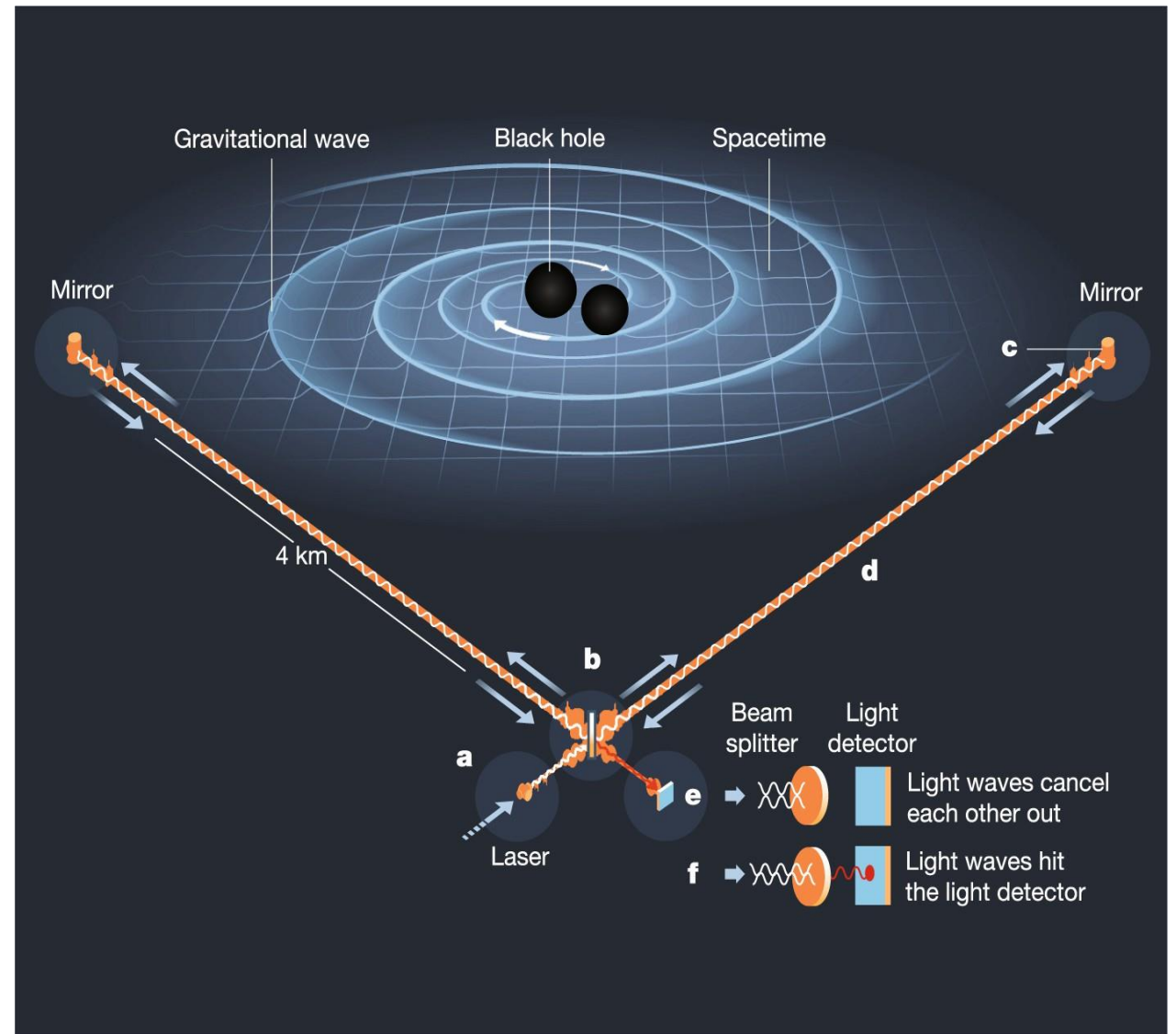




LIGO as a Detection Tool



- LIGO is a dual-recycled Fabry-Pérot Michelson interferometer designed to detect gravitational waves.
- Its extreme sensitivity to length changes makes it an ideal instrument for detecting scalar field dark matter.





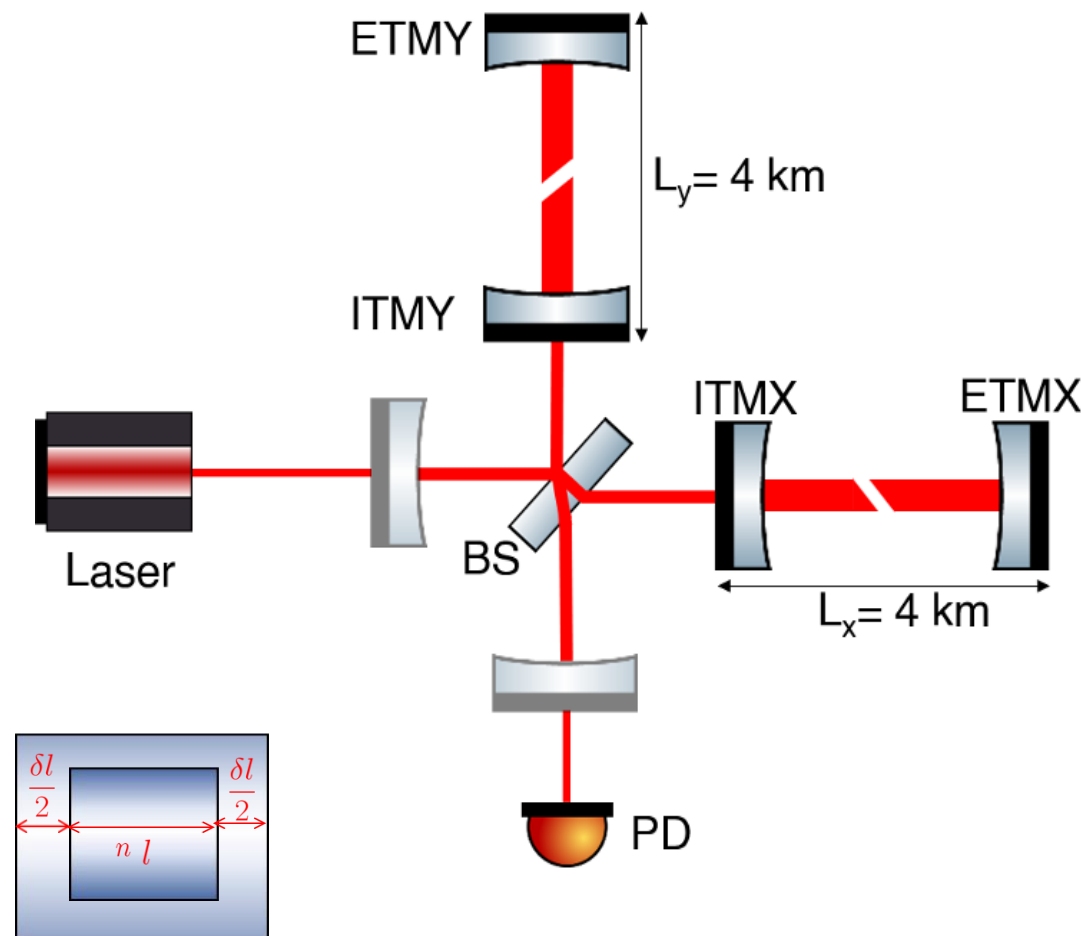
How Scalar Field Dark Matter Affects LIGO Measurements



- Oscillations in LIGO's beamsplitter and test masses, leading to detectable path length differences.
- These effects are measured through changes in the interferometer's output signal.

DM “size” effect only:

- **Beamsplitter**
- Splitting occurs far from the center of mass
- **Test masses**
- Asymmetry from thickness differences





LIGO's Dual Detectors





Calibrated Strain

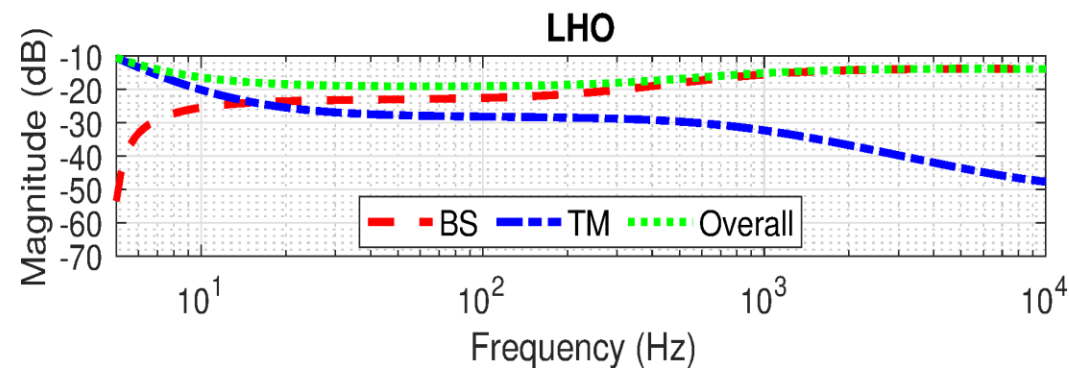
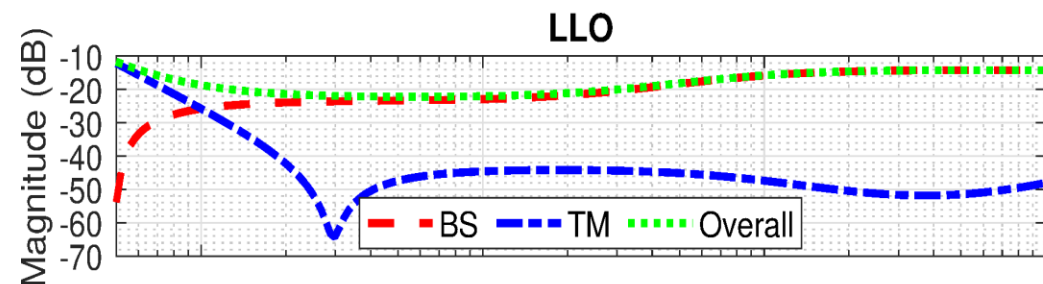
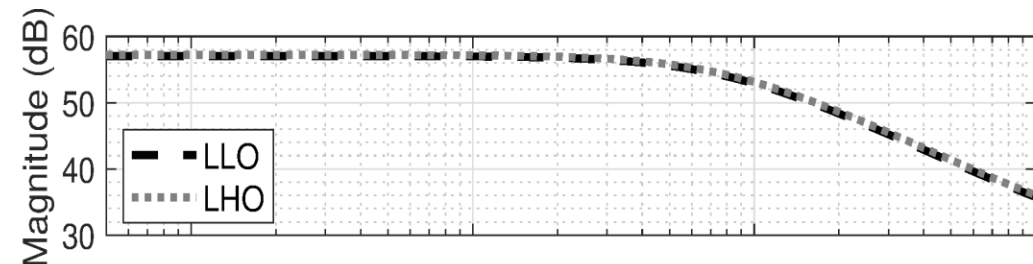


- Available strain data is calibrated to gravitational waves:

$$h(\omega) = \frac{I_{\text{PD}}(\omega)}{L T_{\text{GW}}(\omega) e^{i\phi_{\text{GW}}}}$$

- We are instead interested in DM-induced strain:

$$s_{\text{DM}}(\omega) = \frac{I_{\text{PD}}(\omega)}{|n t_{\text{B}} T_{\text{B}} e^{i\phi_{\text{B}}} + t_{\text{M}} T_{\text{M}} e^{i\phi_{\text{M}}}|}$$

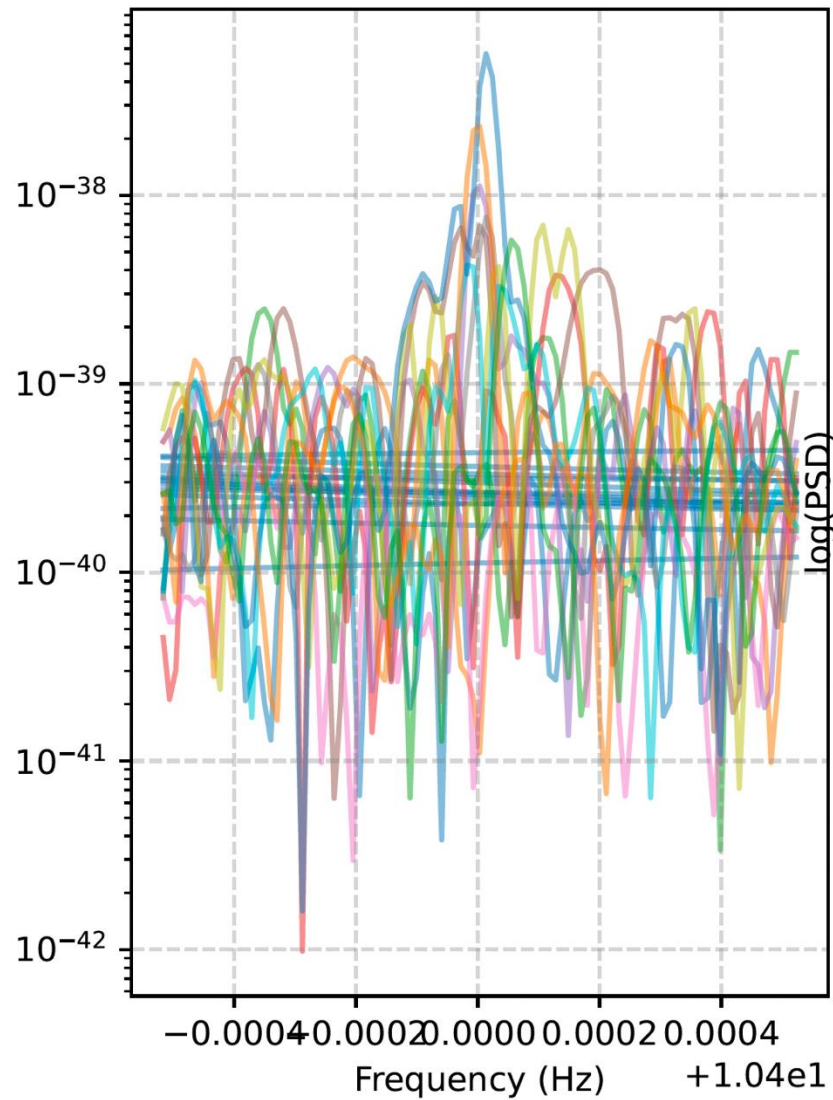




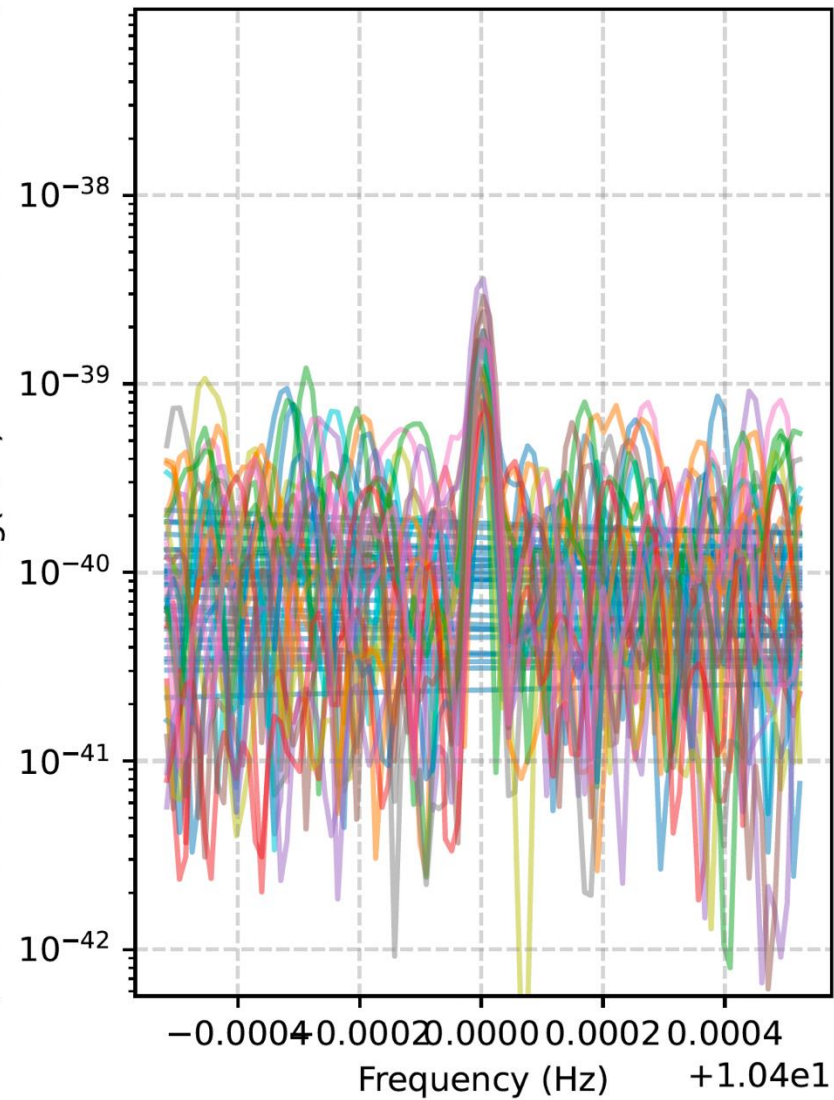
10.4 Hz: our best candidate



Hanford Observatory



Livingston Observatory





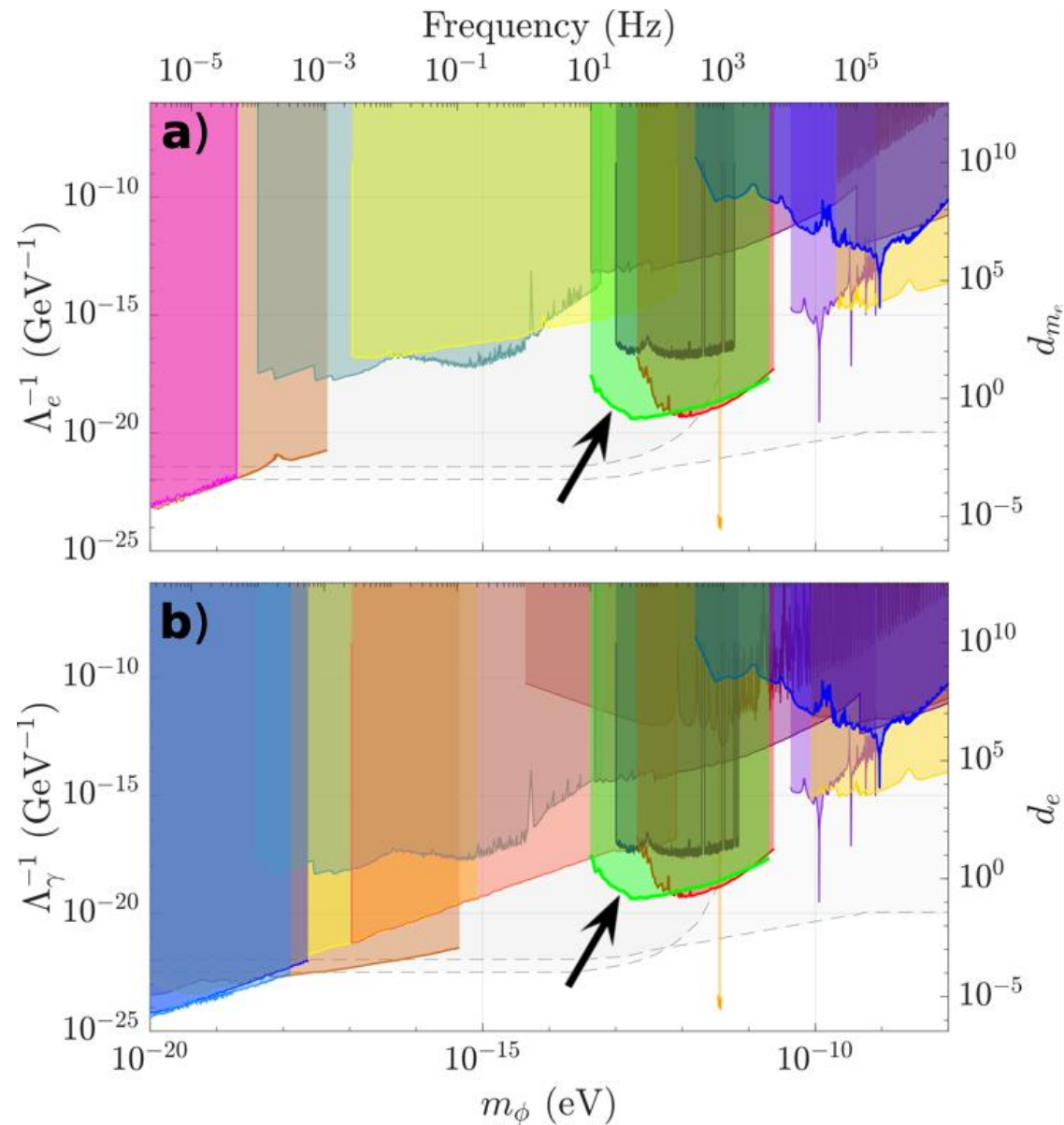
Upper Limits Results



- Our results in green
- Up to x1000 improvement below 180 Hz
- Competitive up to 5 kHz

A.S. Göttel, A. Ejlli, K. Karan, S.M. Vermeulen, L. Aiello V. Raymond, H. Grote

<https://arxiv.org/abs/2401.18076>



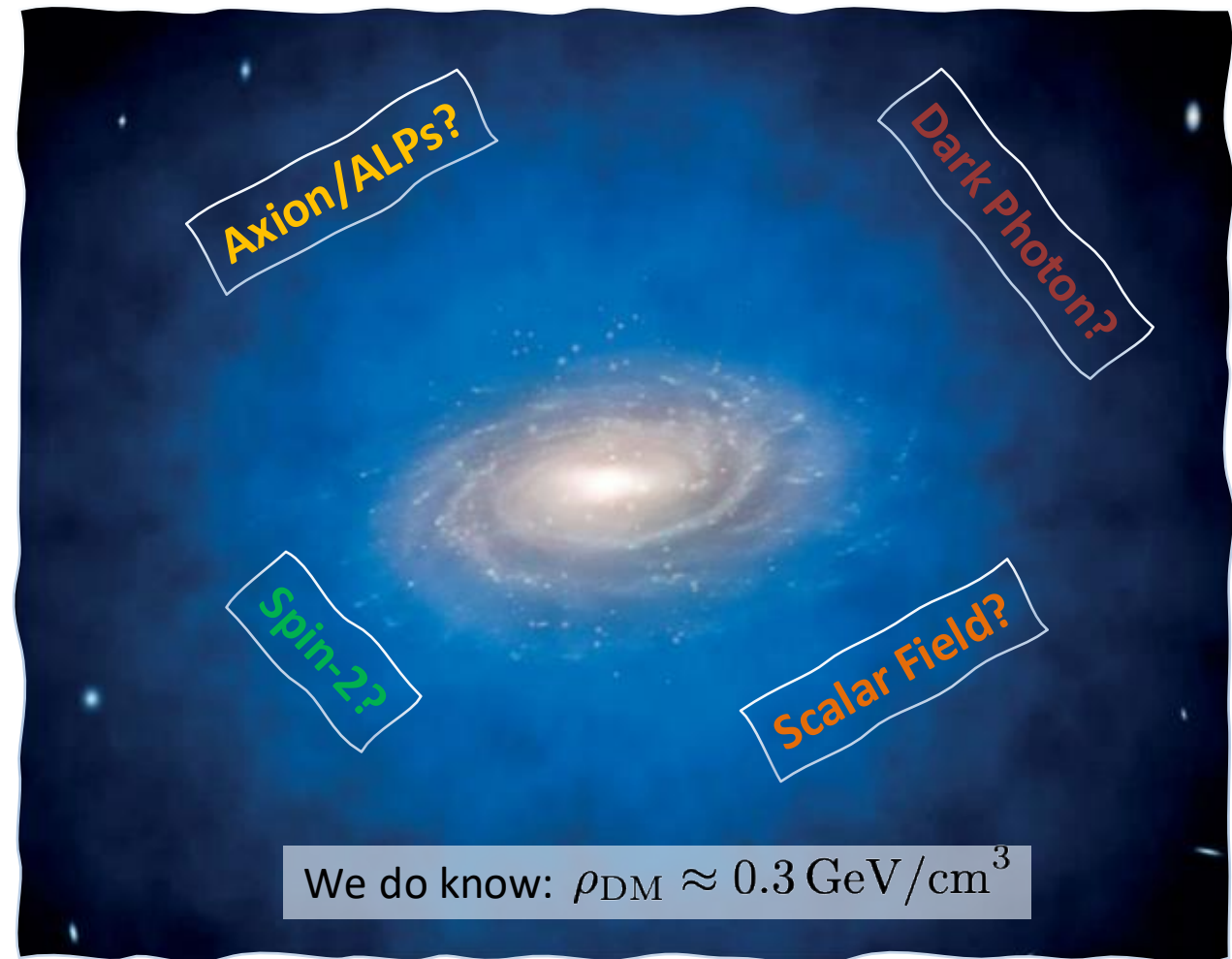


Axion Field Dark Matter



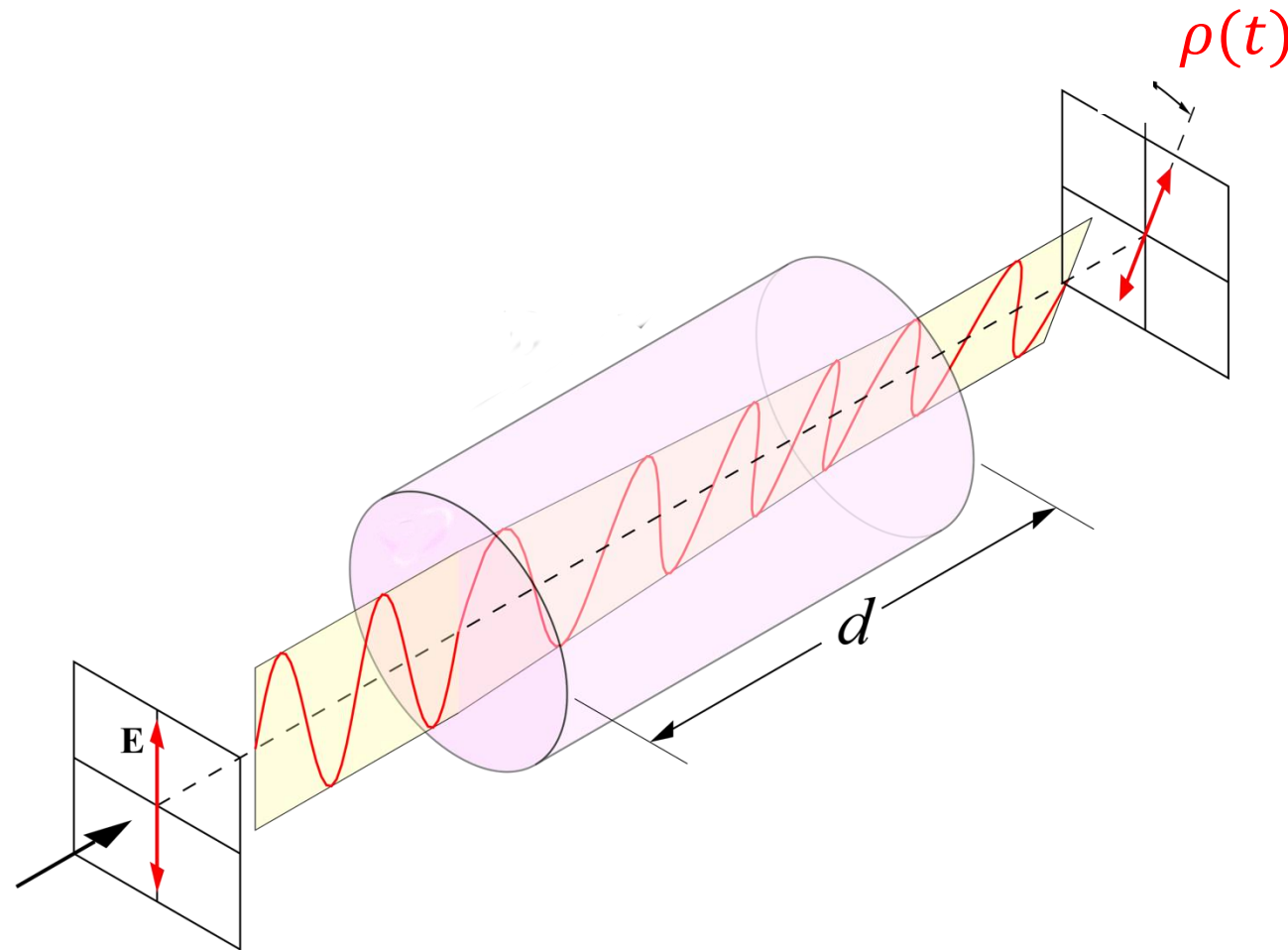
ALPs Coupling

$$\mathcal{L}_{\text{int}} = \frac{ag_{a\gamma}}{4} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

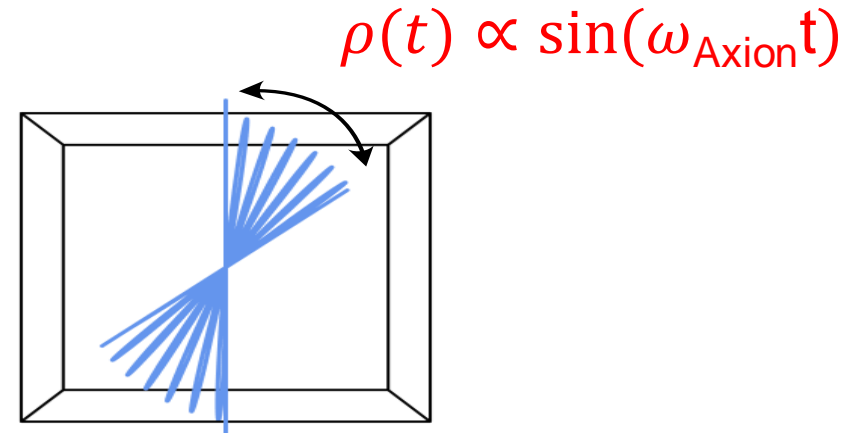




Effect of Axion Field on Polarization of Light

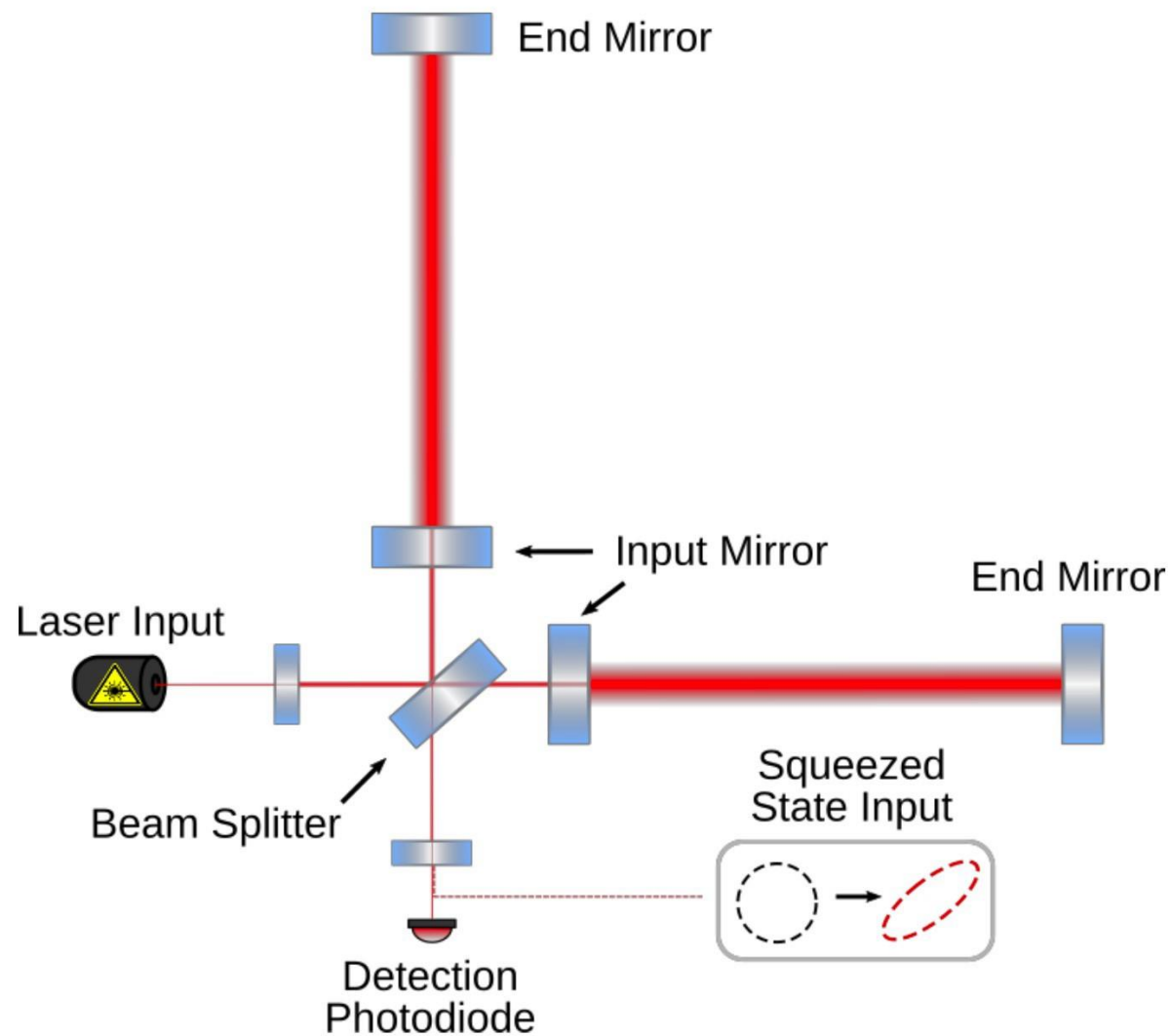


- Axion field rotates the polarization of linearly polarized light
- Angle of rotation oscillates with the frequency of Axion field

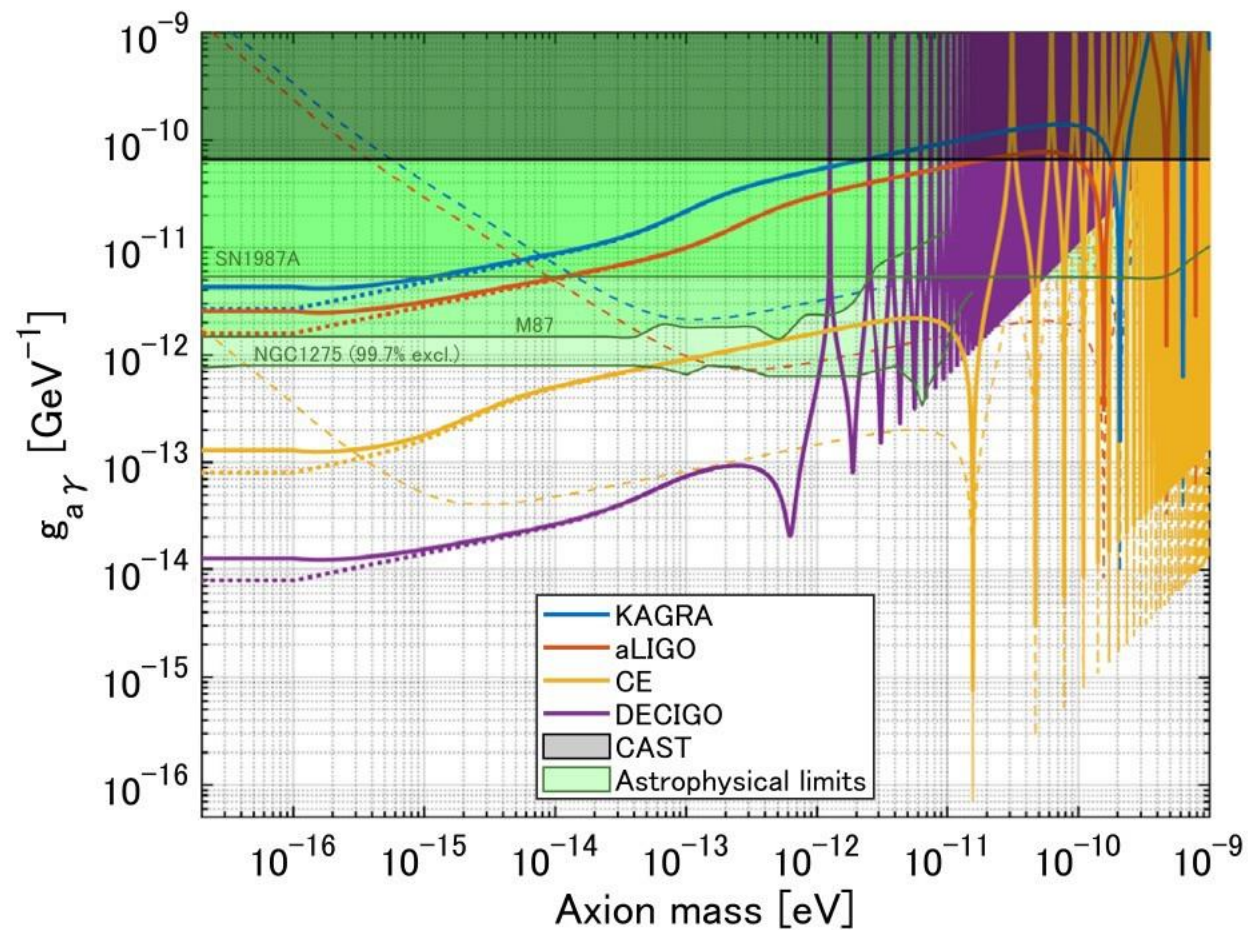




GW Detectors and Axion Search

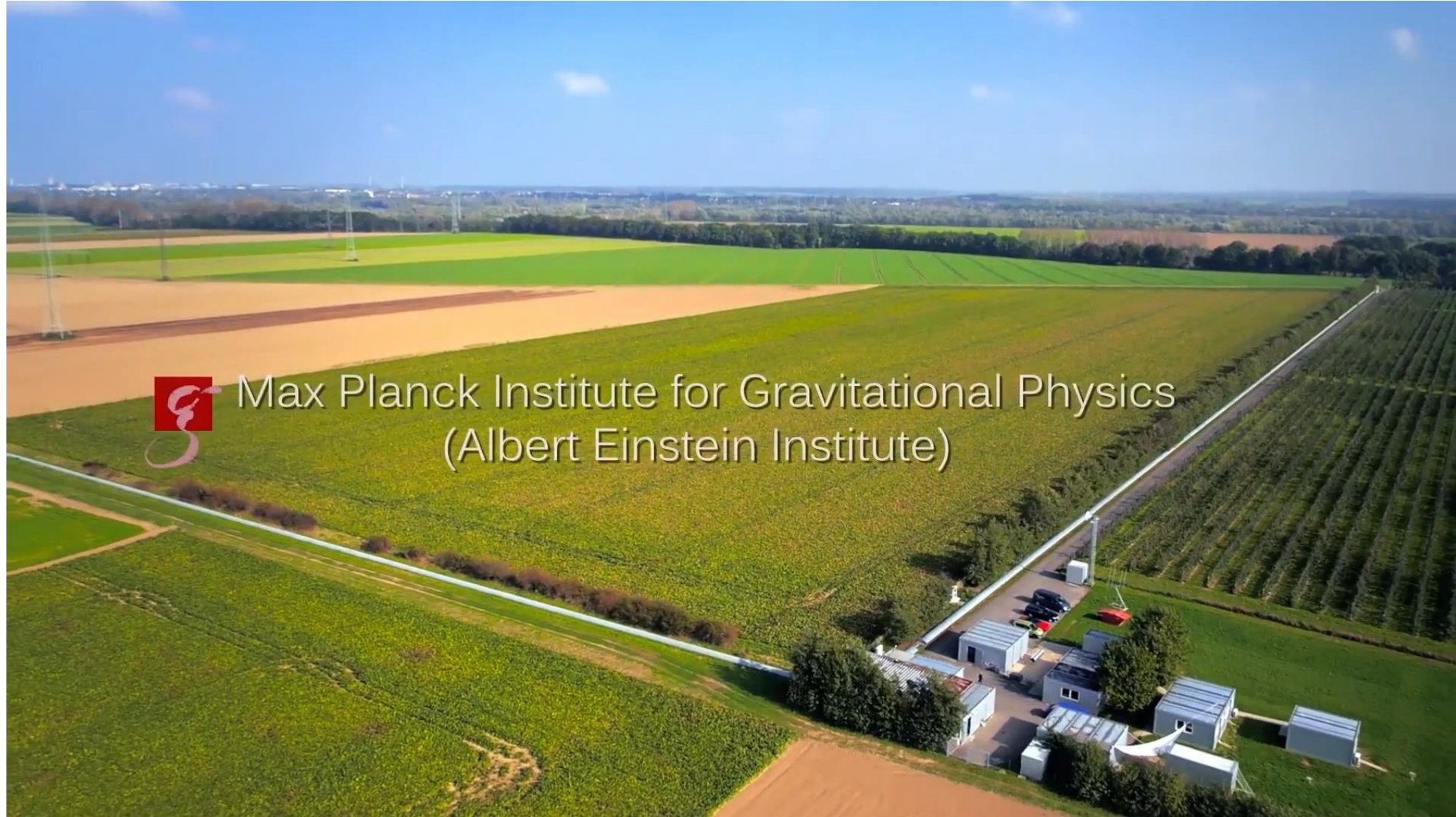


K. Nagano, *et al*
PRD 104, 062008 (2021)





GEO 600 – Gravitational Wave Detector





PAPER

DarkGEO: a large-scale laser-interferometric axion detector

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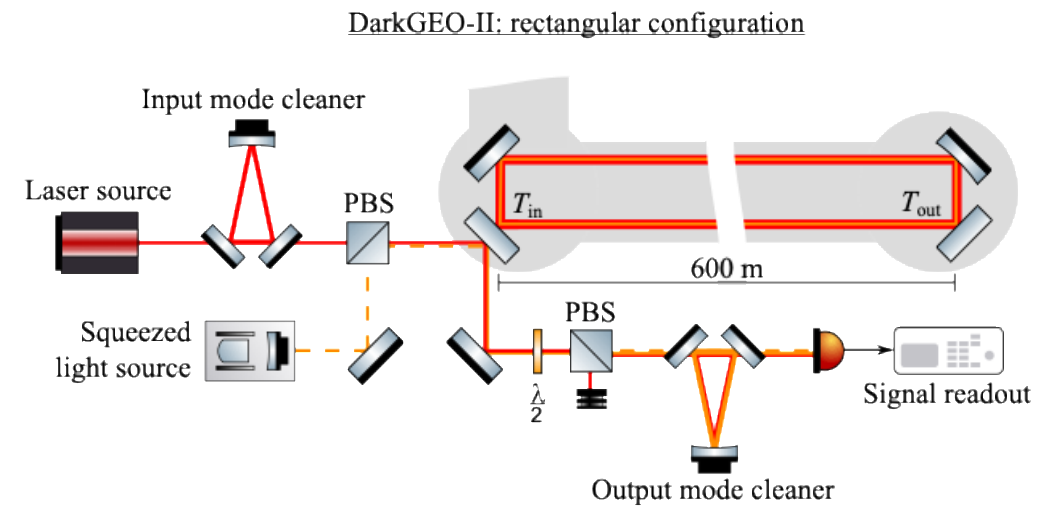
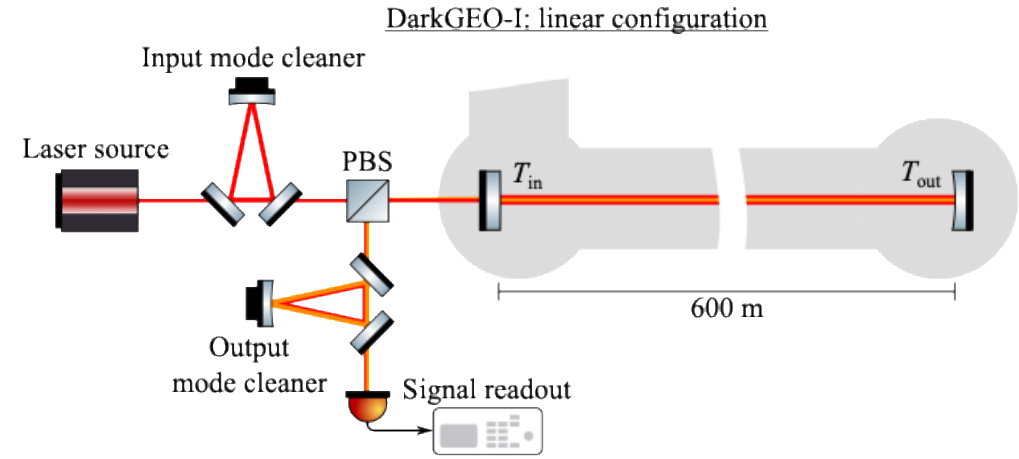
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- **DarkGEO-I: Linear Configuration**
- **DarkGEO-II: Rectangular Configuration**
- **DarkGEO-III: DarkGEO-III: Full Coincidence Search**

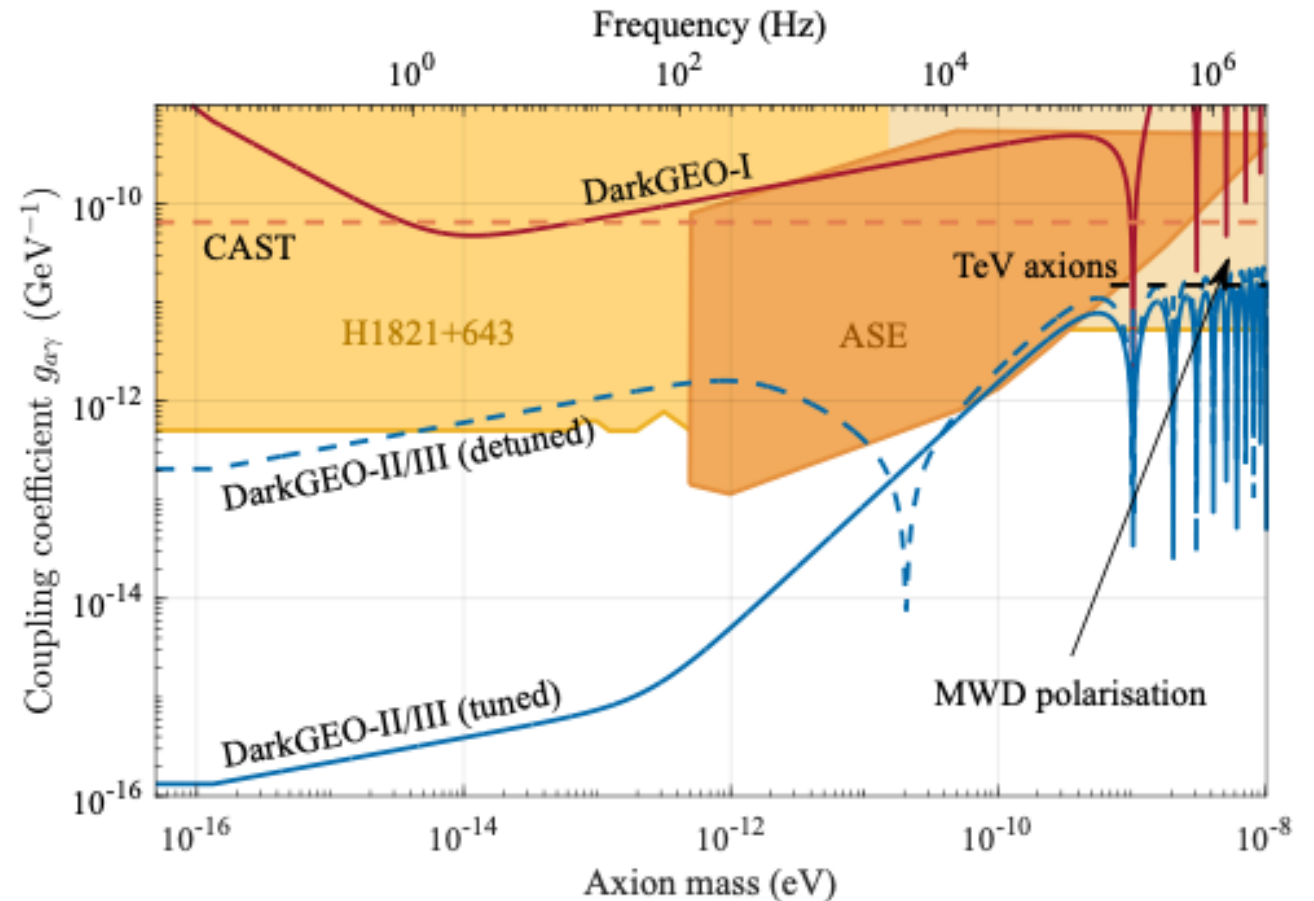




Potential of Large Scale Axion Polarimetry



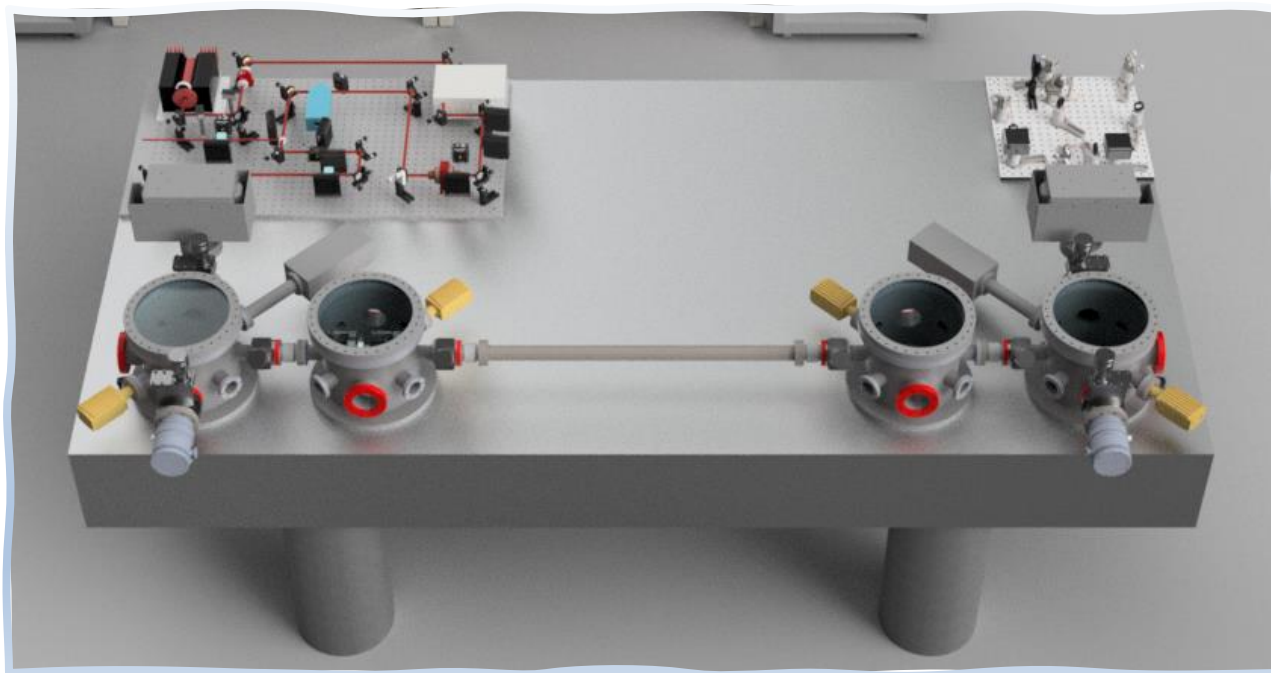
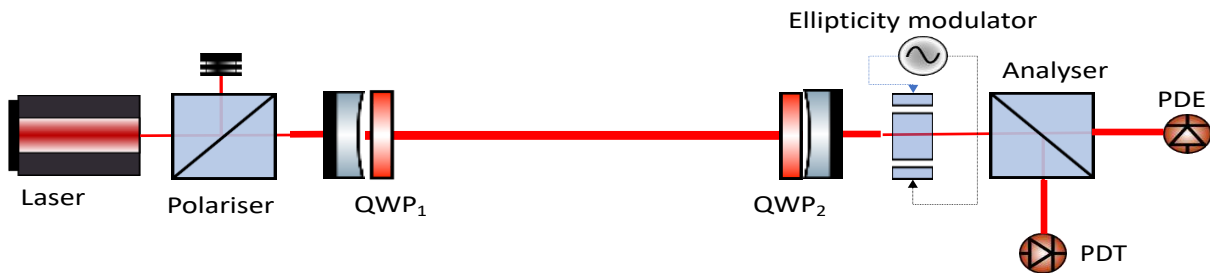
- High Intra-Cavity Power 10 MW
- Squeezed Light at 10 dB to reduce quantum noise and enhance detection.
- Assumes shot-noise-limited performance across the measurement band for optimal signal-to-noise ratio.



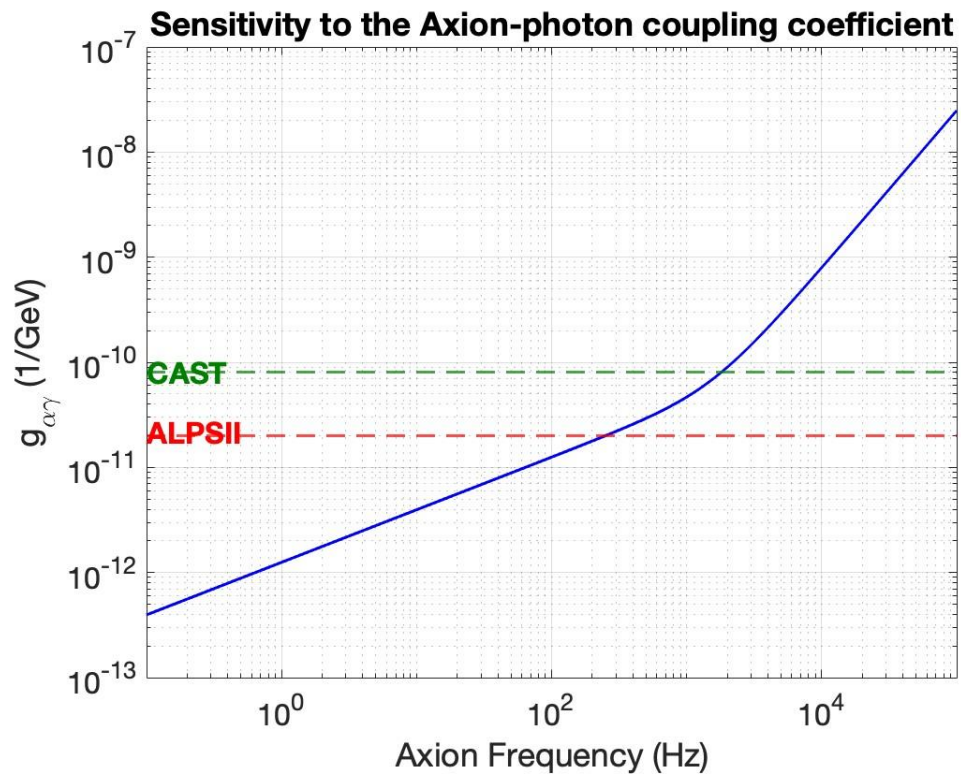
J. Heinze *et al*, [New J. Phys. 26 \(2024\) 055002](#)



Polarimetry-Prototype under construction at AEI MPG



$L = 1.5 \text{ m}$



➤ Sensitivity is proportional to the length!

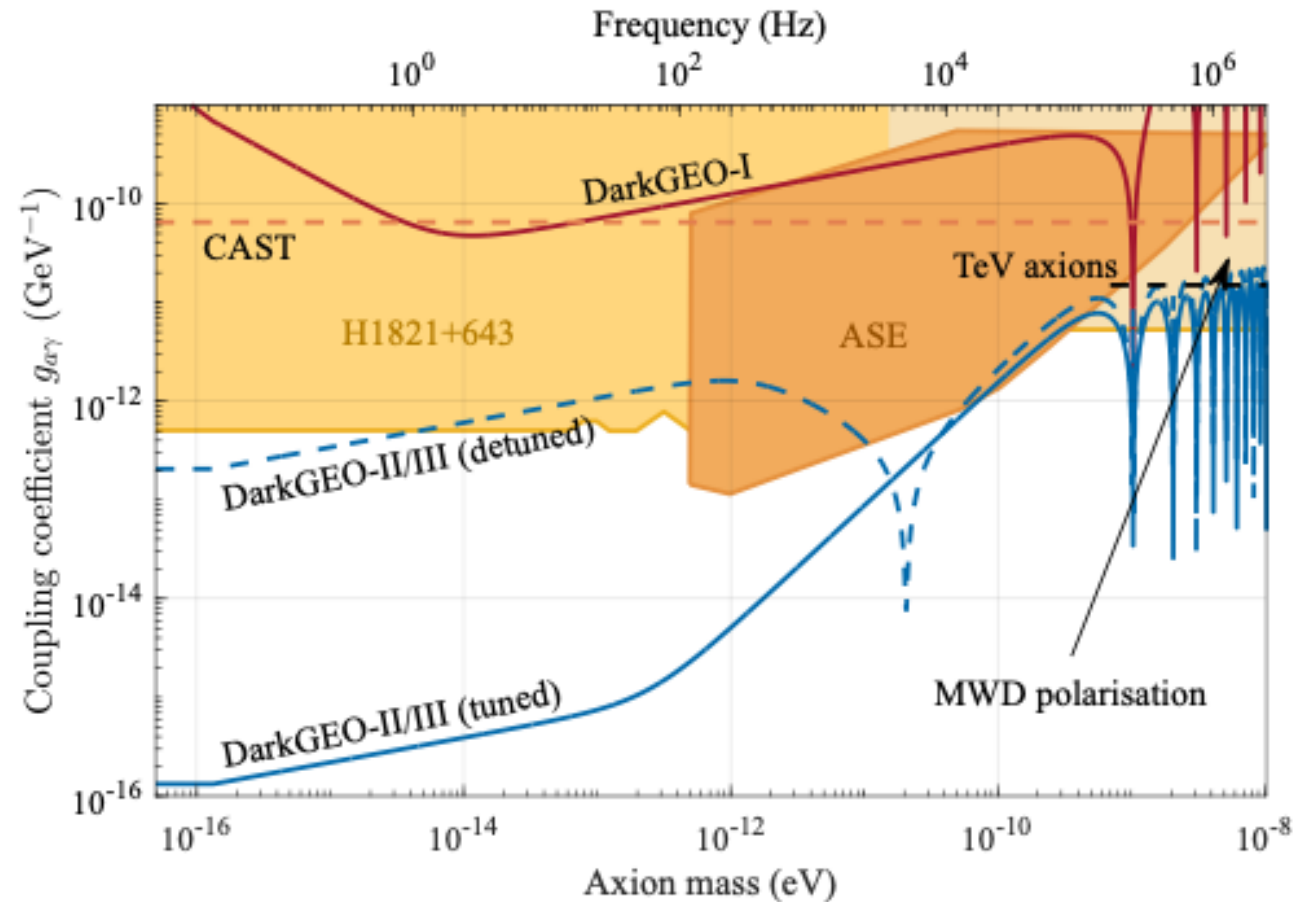
$$g_{\alpha\gamma} = \frac{S_P^{(\text{tot})}}{2\tau} \sqrt{\frac{P_{\text{AR}}^2 + 4 \sin^2(\pi\nu_a\tau)}{2\rho_{\text{local}}}}$$



Discussion: Why targeting 10^{-16} - 10^{-10} eV?



- Could this mass range validate multiple theoretical models of axions?
- Could the anomalies in cosmology and astrophysics be due to axions in this specific mass window?
- Could this experiment fill the critical gap left by other axion searches, making it a unique contribution to the field?





Thank you for your attention