

Stringy Axion Dark Energy

Francisco Gil Pedro

Based on:

2002.02695, 2007.11011 and 2206.10649
M. Brinkman, M. Cicoli and G. Dibitetto



2112.10783, 2112.10779, and 2407.03405
M. Cicoli, C. Cunillera and A. Padilla



Outline

- Introduction
- Dynamical alternatives to the cosmological constant
 - Multifield quintessence
 - Axionic quintessence
 - Embedding axionic quintessence

The CC problem

1998 SN Ia -> Universe in accelerated expansion

$$H_0 = 100 h \text{ km/s/Mpc} \approx 10^{-60} M_{pl}$$

$h = 0.675$ PLANCK 2018

$h = 0.73$ SH0ES

Explained by a fluid with $p = \omega \rho$ $\omega \approx -0.99^{+0.15}_{-0.13}$ $\Omega \sim 0.7$

Compatible with a cosmological constant $\Lambda \approx 10^{-120} M_P^4$

See however DESI: 2404.03002 $\omega(a) = \omega_0 + (1 - a)\omega_a$

$-0.55^{+0.39}_{-0.21} < -1.32$

The CC problem, theoretically

What if $\omega \neq -1$?

(Too) Many alternatives:

Quintessence: Single field $\lambda = -\frac{V_\phi}{V}$ $\lambda < 0.6$ [Agrawal et al. '18]
[Akrami et al. '18]

No dS conjecture? $\lambda \geq \mathcal{O}(1)$

How to explain the DE scale?

Radiative stability?

Fifth forces?

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[Akrami et al. '18]

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Alternatives:

-multifield dynamics

-axionic quintessence

[(Brinkman), Cicoli, Dibitetto, FGP '20-22]

[Cicoli, Cunillera, Padilla, FGP '21]

Stringy axions and moduli

Universal features of string theories:

- gravity
- extra dimensions
- higher form antisymmetric tensors

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Freund-Rubin compactification

$$S_{(D)} = \frac{M^{D-2}}{2} \int d^D x \sqrt{g_{(D)}} (R_{(D)} - |F_p|^2)$$

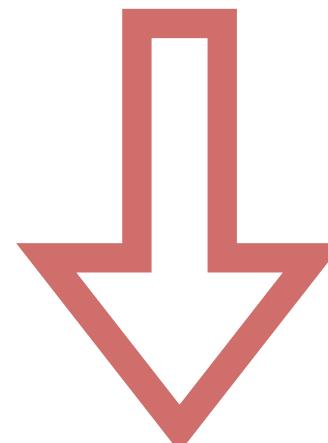
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$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu + \sigma^2 g_{mn} dx^m dx^n$$

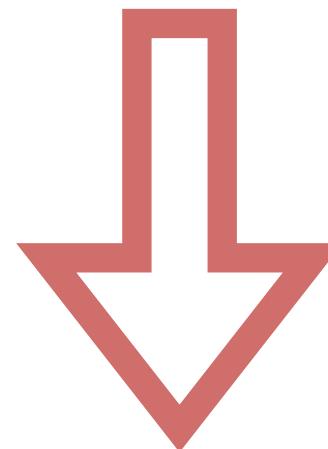
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$$S_{(D)} = \frac{M_P^2}{2} \int d^4 x \sqrt{g_{(4)}} \left(R_{(4)} - \frac{1}{2}(D-4)(D-2)(\partial \log \sigma)^2 + \frac{\chi}{2} M_P^2 \sigma^{2-D} - M_P^2 \sigma^{4-2p-D} f^2 \right)$$

where $M_P^2 = M^2 \sigma_0^{D-4}$ $\int d^{D-4} y \sqrt{g_{(D-4)}} R_{(D-4)} = \chi M^{6-D}$ $\int d^{D-4} y \sqrt{g_{(D-4)}} g^{m_1 n_1} \dots g^{m_p n_p} F_{m_1 \dots m_p} F_{n_1 \dots n_p} \equiv M^{2-D} f^2$

Stringy axions and moduli

[Cicoli et al. '23]

Theory	Dimension	Supercharges	Massless Bosons
Heterotic $E_8 \times E_8$	10	16	g_{MN}, B_{MN}, ϕ A_M^{ij}
Heterotic $SO(32)$	10	16	g_{MN}, B_{MN}, ϕ A_M^{ij}
Type I $SO(32)$	10	16	g_{MN}, ϕ, A_M^{ij} C_{MN}
Type IIA	10	32	g_{MN}, B_{MN}, ϕ C_M, C_{MNP}
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M-Theory	11	32	g_{MN}, C_{MNP}

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Kahler moduli
 g_{MN}
complex structure moduli

τ_i $h_{1,1}$
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(real) Kahler potential $K(\Phi, \bar{\Phi})$

(holomorphic) superpotential $W(\Phi)$

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(holomorphic) superpotential $W(\Phi)$

$$\mathcal{L} = -K_{i\bar{j}}\partial_\mu T_i \partial^\mu \bar{T}_{\bar{j}} - V.$$

$$V = e^K \left[K^{i\bar{j}} D_i W D_{\bar{j}} \bar{W} - 3|W|^2 \right]$$

Multifield Quintessence

Multifield modulus+axion

[Spintessence, Boyle et al.'01]
[Sonner and Townsend '06]
[Achucarro and Palma '18]

$$\frac{\mathcal{L}}{\sqrt{-g}} = \frac{1}{2} \gamma_{ij} \partial_\mu \phi^i \partial^\mu \phi^j - V(\phi) ,$$

[(Brinkman), Cicoli, Dibitetto, FGP '20-22]

$$\gamma_{ij} = \begin{pmatrix} 1 & 0 \\ 0 & f^2(\phi_1) \end{pmatrix}$$

$$f(\phi_1) = e^{-k_1 \phi_1}$$

$$V = V_0 e^{-k_2 \phi_1}$$

$$R_{fs} = -2k_1^2$$

Exploit difference between

$$\epsilon_V = \frac{1}{2} \left(\frac{V'}{V} \right)^2$$

and

$$\epsilon_H = -\frac{\dot{H}}{H^2}$$

accelerated expansion in steep potentials: $\frac{k_2^2}{2} \gtrsim \mathcal{O}(1)$

Multifield Quintessence

$$\begin{cases} \ddot{\phi}_1 + 3H\dot{\phi}_1 - f f_1 \dot{\phi}_2^2 + V_1 = 0 , \\ \ddot{\phi}_2 + 3H\dot{\phi}_2 + 2\frac{f_1}{f}\dot{\phi}_2\dot{\phi}_1 = 0 , \end{cases}$$

● Geodesic:

$$\dot{\phi}_2 = 0 \quad \dot{\phi}_1 = -\frac{V_1}{3H}$$

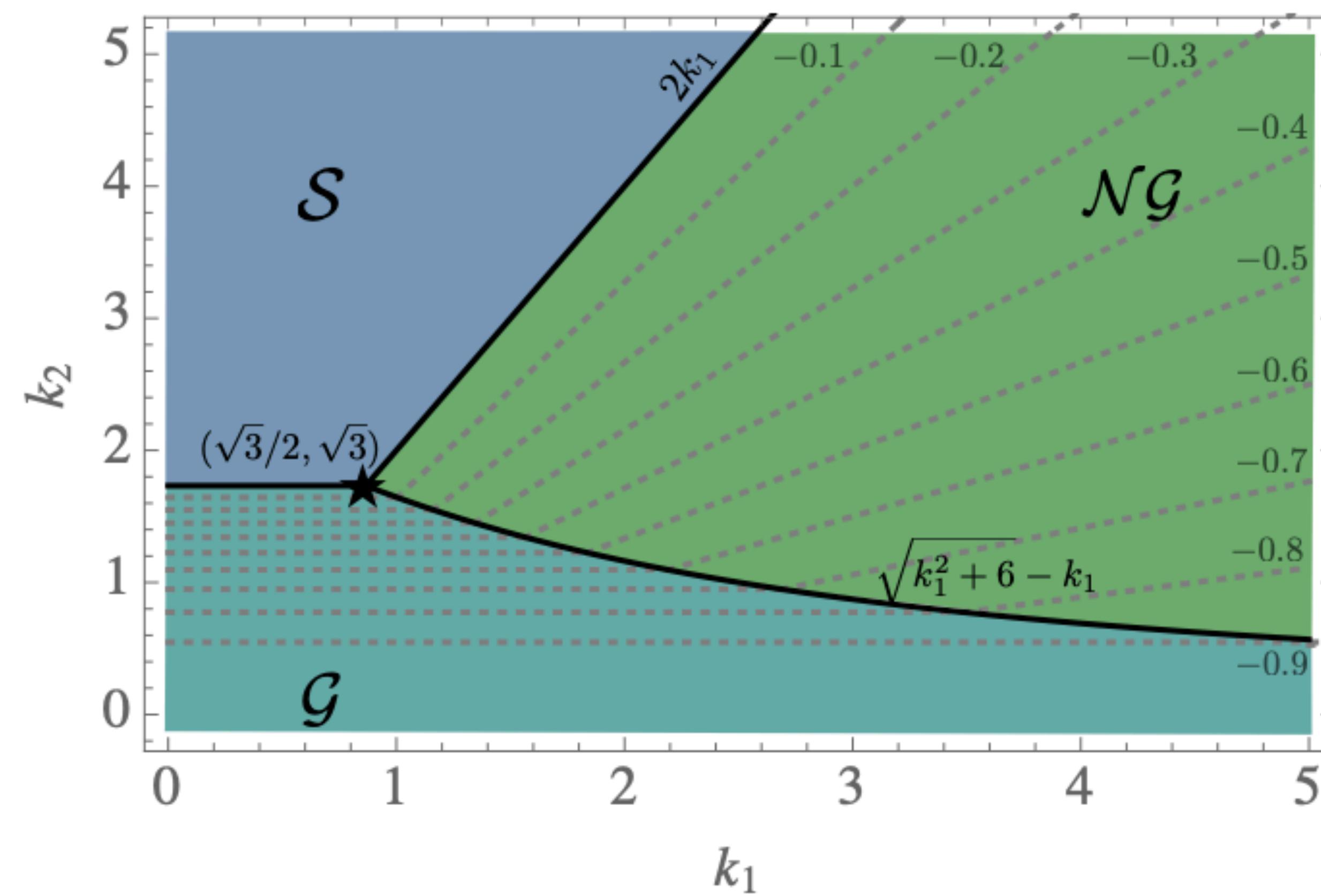
● Non Geodesic:

$$\dot{\phi}_2 \neq 0 \quad \dot{\phi}_1 = -\frac{3Hf}{2f_1}$$

	x_1	x_2	y_1	Ω_ϕ	ω_ϕ	existence
\mathcal{K}_+	1	0	0	1	1	all k_1, k_2, γ
\mathcal{K}_-	-1	0	0	1	1	all k_1, k_2, γ
\mathcal{F}	0	0	0	0	undefined	all k_1, k_2, γ
\mathcal{S}	$\frac{\sqrt{3/2}\gamma}{k_2}$	0	$\frac{\sqrt{3/2}\gamma(2-\gamma)}{k_2}$	$\frac{3\gamma}{k_2^2}$	$\gamma - 1$	$0 < \gamma < 2 \wedge k_2^2 \geq 3\gamma$
\mathcal{G}	$\frac{k_2}{\sqrt{6}}$	0	$\sqrt{1 - \frac{k_2^2}{6}}$	1	$-1 + \frac{k_2^2}{3}$	$k_2 < \sqrt{6}$
\mathcal{NG}	$\frac{\sqrt{6}}{(2k_1+k_2)}$	$\frac{\pm\sqrt{k_2^2+2k_2k_1-6}}{2k_1+k_2}$	$\sqrt{\frac{2k_1}{2k_1+k_2}}$	1	$\frac{k_2-2k_1}{k_2+2k_1}$	$k_2 \geq \sqrt{6+k_1^2} - k_1$

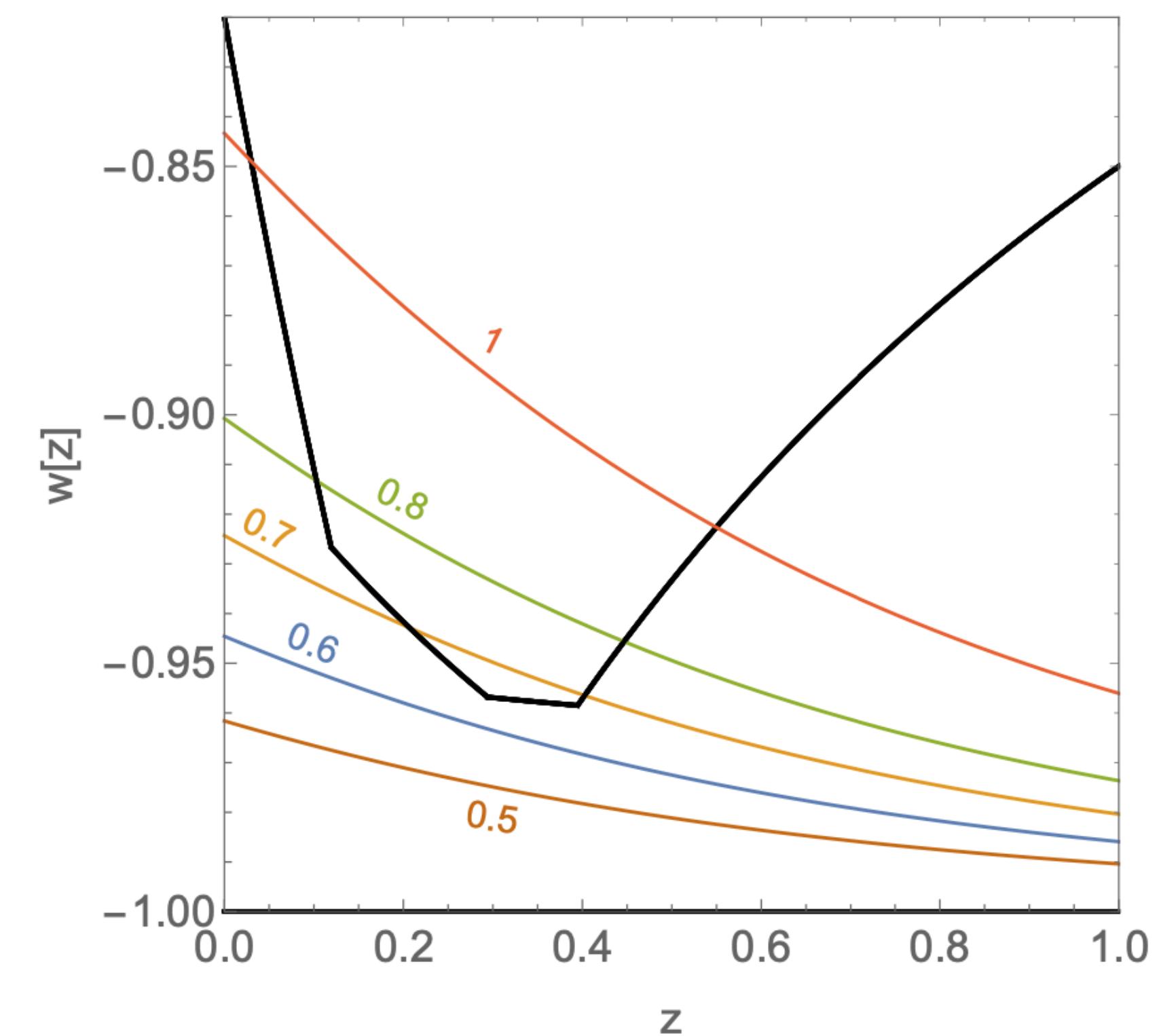
Multifield Quintessence

Single field regime $k_1 = 0$



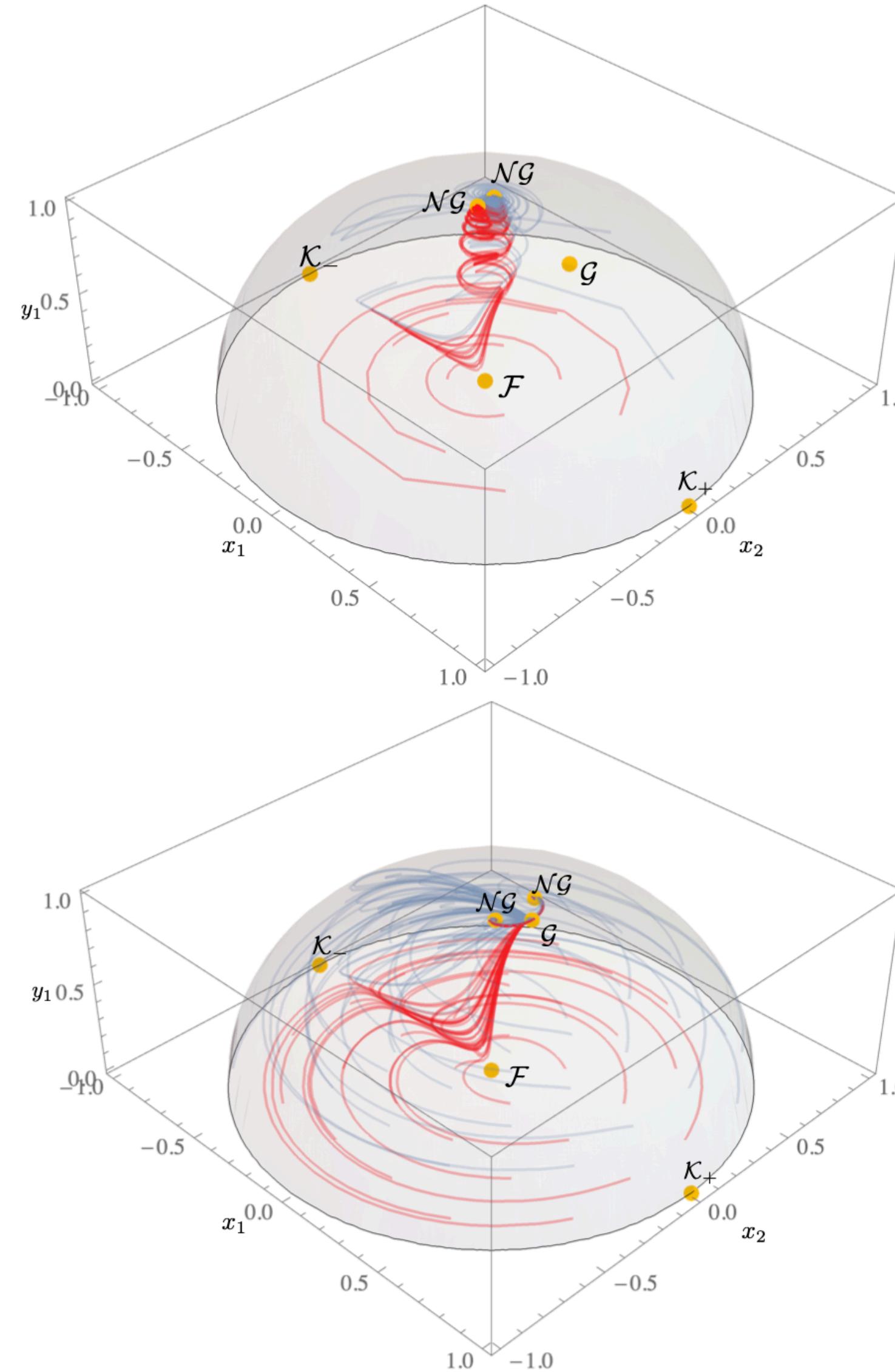
Look for observationally viable quintessence

$$\omega_{DE} \sim -1 \quad \Omega_{DE} = 0.7$$

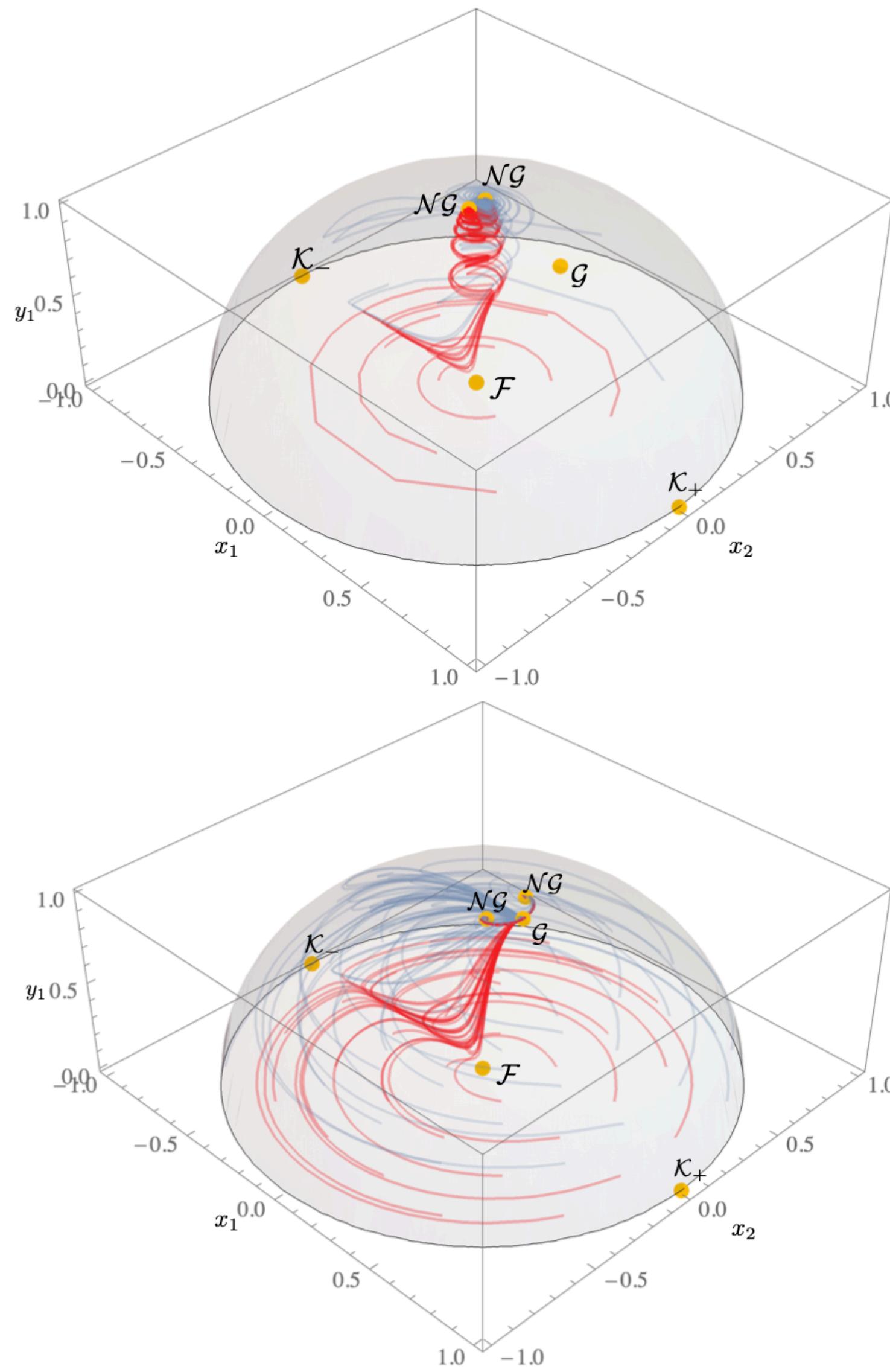


[Agrawal et al., '18]

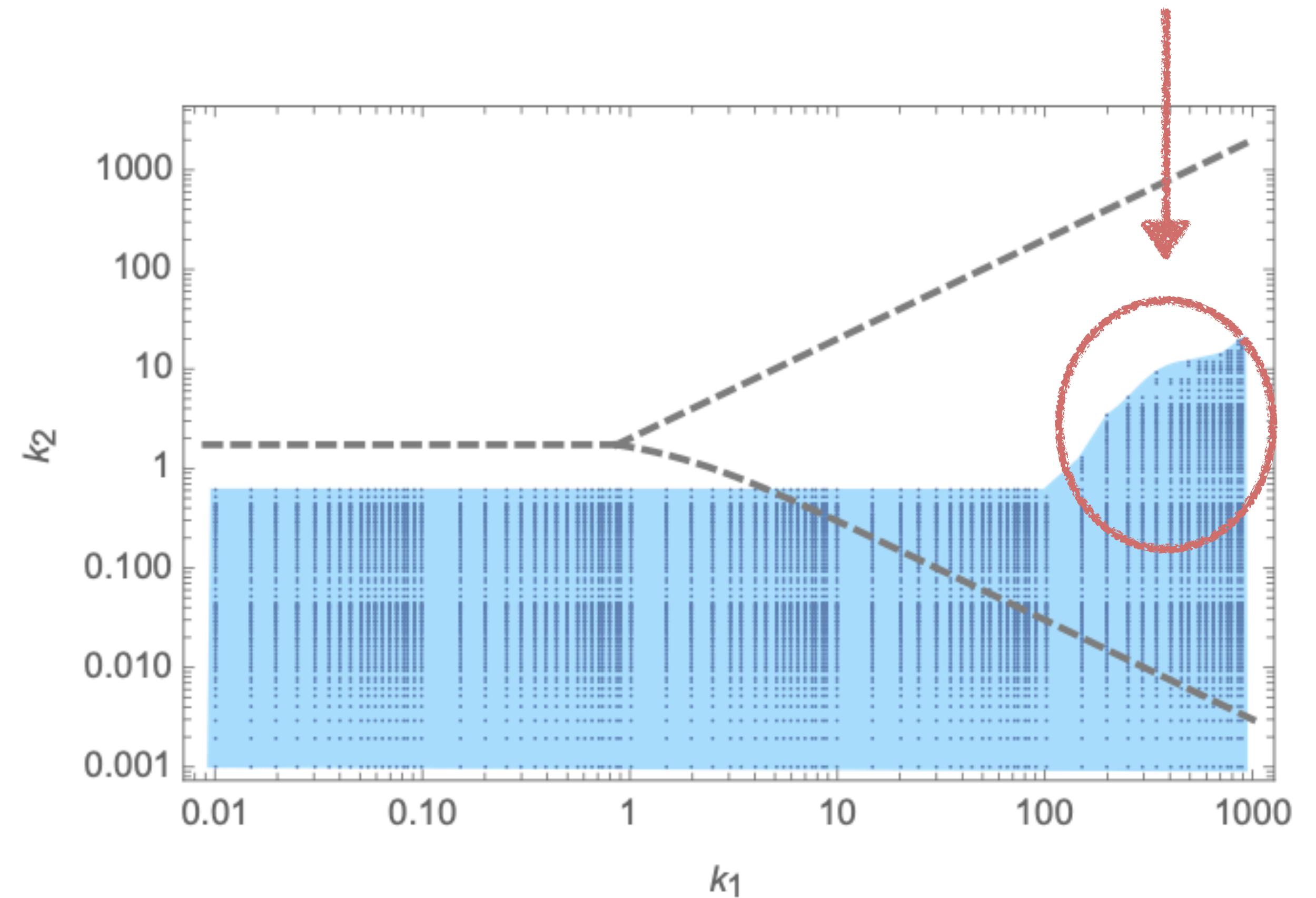
Multifield Quintessence



Multifield Quintessence



Steep potentials require LARGE field space curvatures



Multifield Quintessence

[Saltman,Silverstein, '04]

[(Brinkman),Cicoli,Dibitetto,FGP 2020-22]

?Embeddable into string theory?

$$K = -p \ln[X + \bar{X}] \quad V_K \propto e^K = (X + \bar{X})^{-p}$$

$$k_1 = \sqrt{2/p} \quad \text{and} \quad k_2 = \sqrt{2p}.$$

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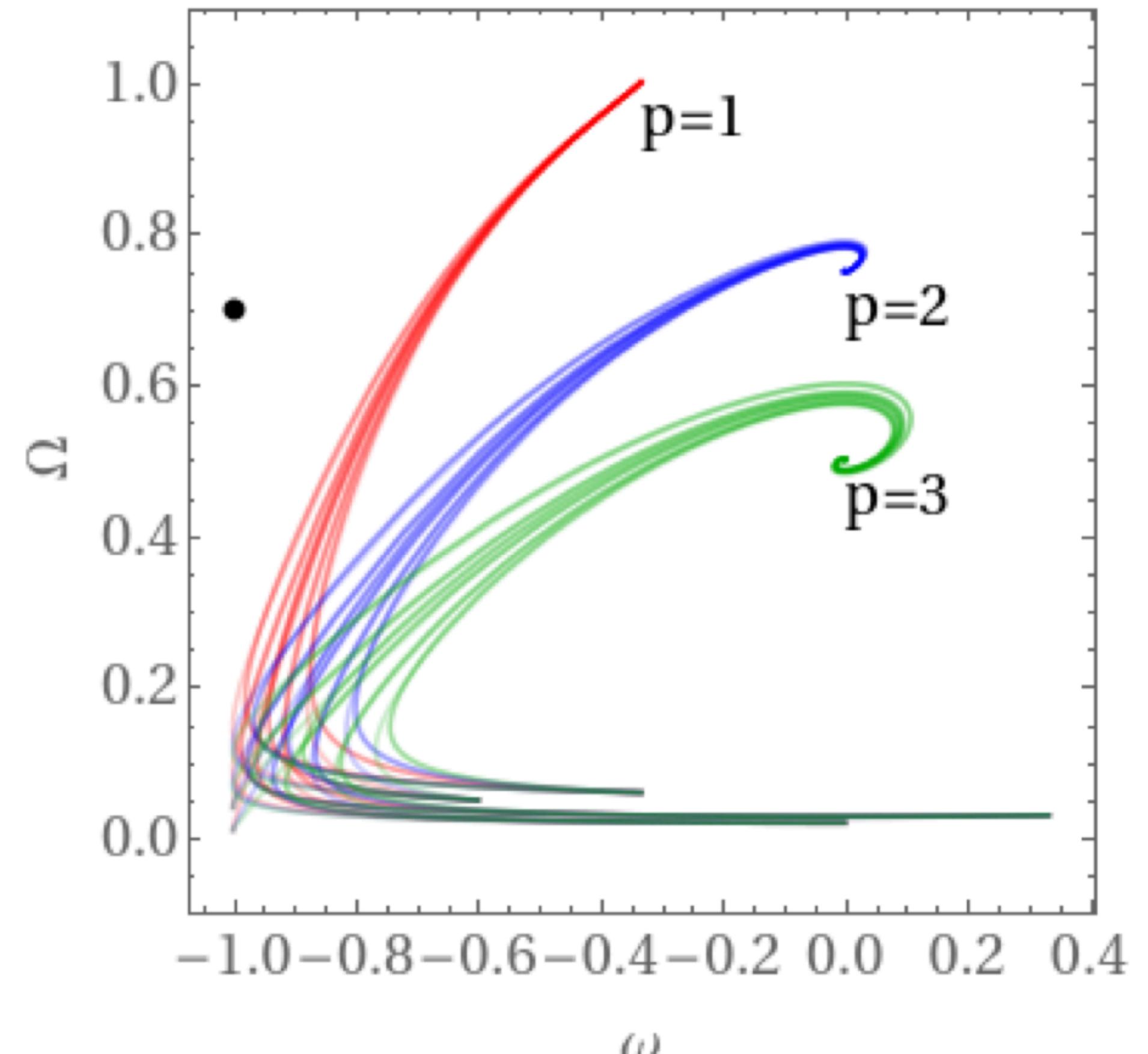
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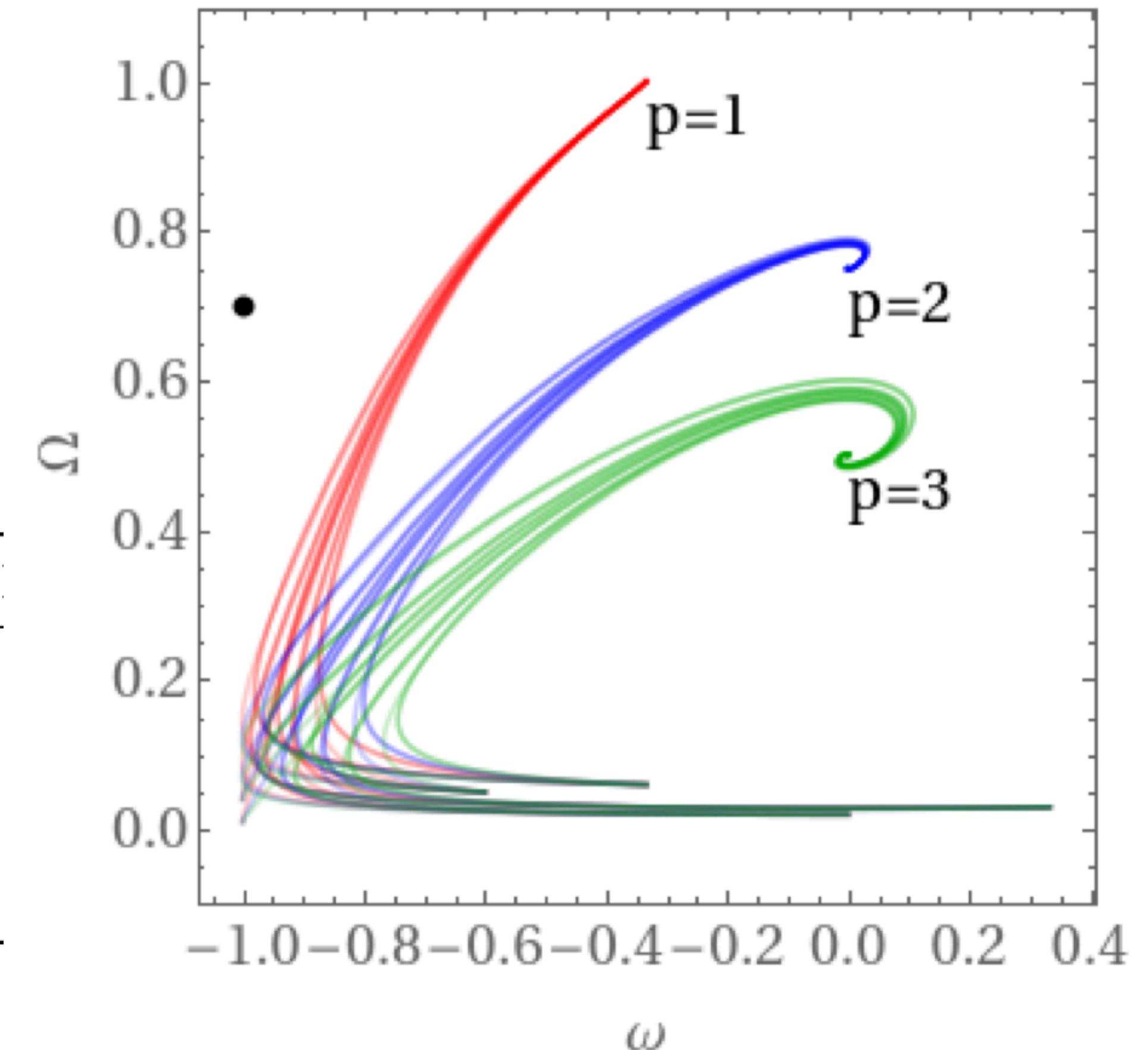
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p	X	Theory	Sources	$\mathcal{M}_{\text{internal}}$
1	$S = e^{-\varphi} + i a$	Heterotic	—	SU(3) str.
2	$T_2 = \text{Vol}(\Sigma_4^{(2)}) + i \int_{\Sigma_4^{(2)}} C_{(4)}$	Type IIB	D3/D7, O3/O7	K3-fibered CY ₃
3	$T = \text{Vol}(\Sigma_4) + i \int_{\Sigma_4} C_{(4)}$	Type IIB	D3/O3	CY ₃
7	$Z = \text{Vol}(\Sigma_3) + i \int_{\Sigma_3} A_{(3)}$	M-theory	KK6/KKO6	G ₂ str.



Axionic DE

[Kaloper,Sorbo '05]

Axionic DE: -Potential generated by non perturbative effects

- Derivative couplings to matter
- Radiative stability

scale suppression

fifth force suppression

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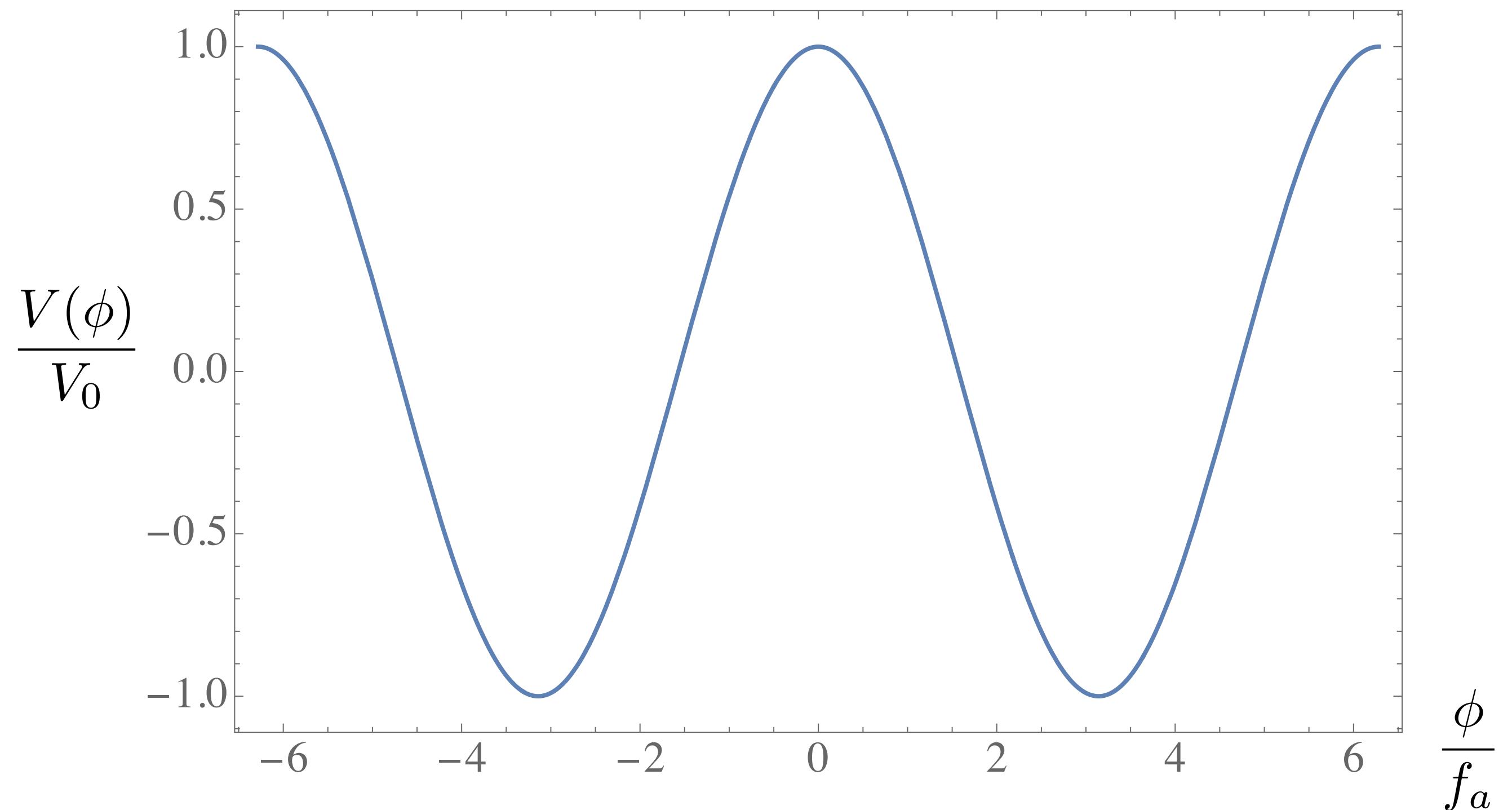
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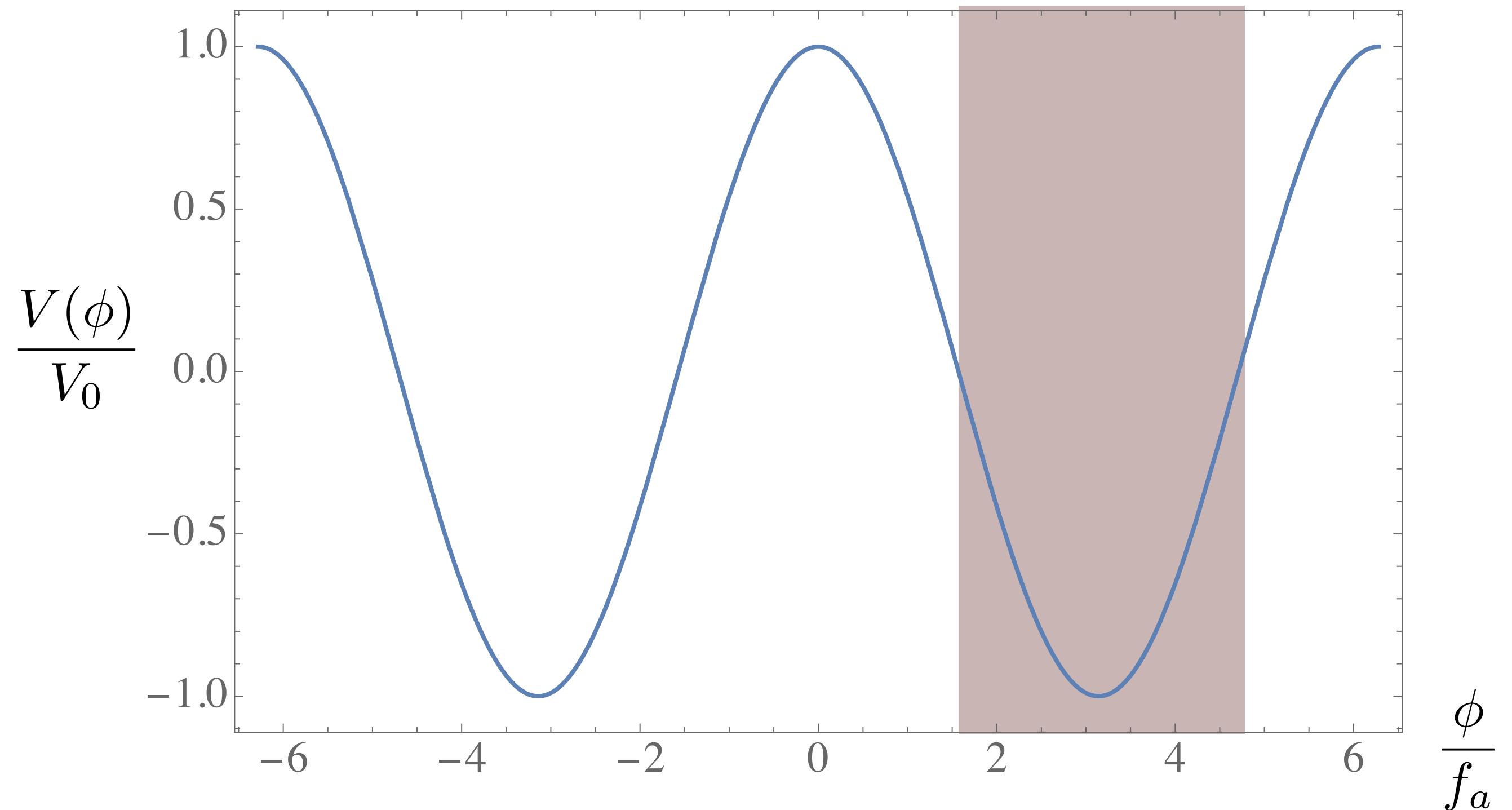
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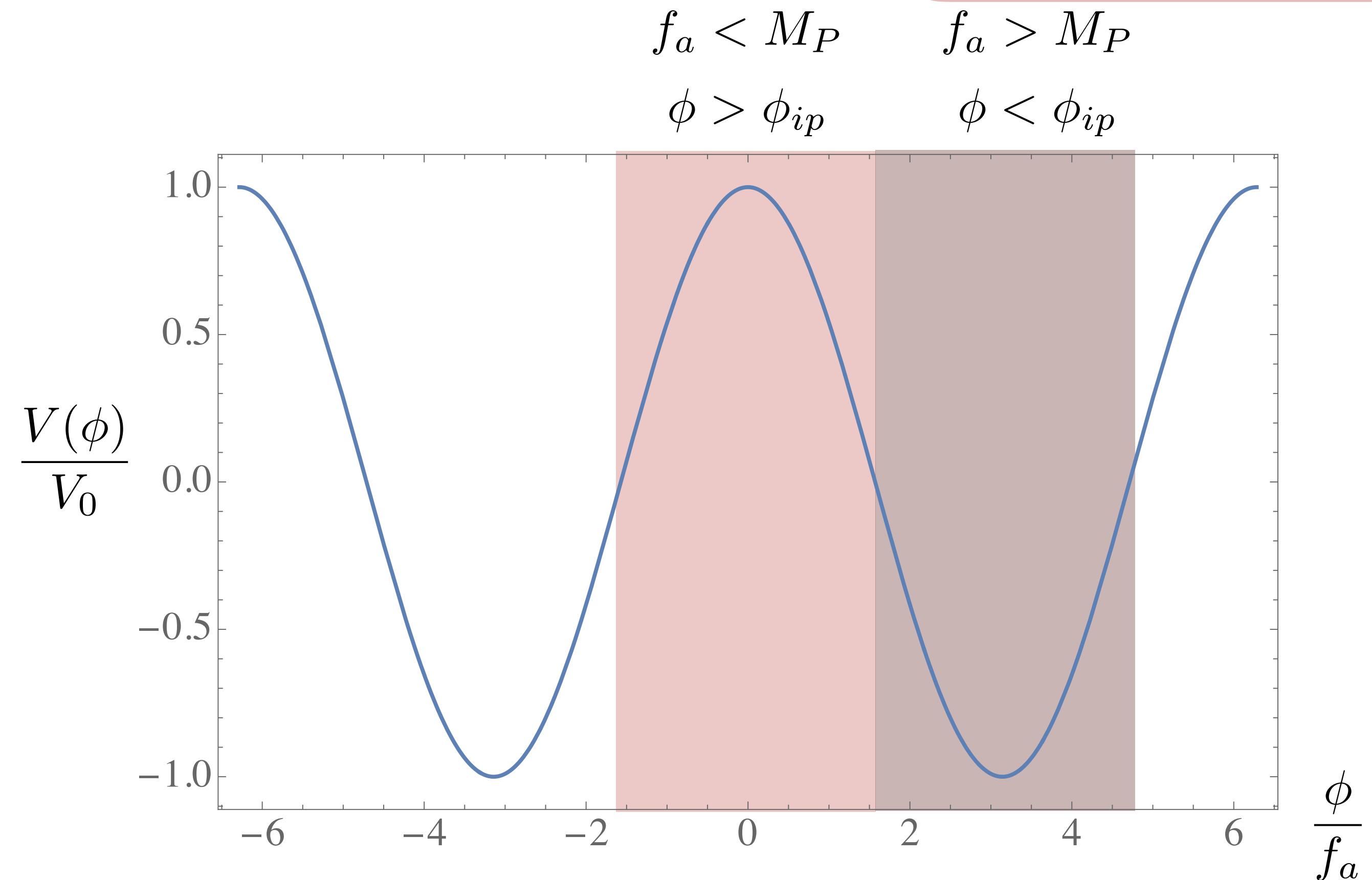
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$$f_a < M_P$$

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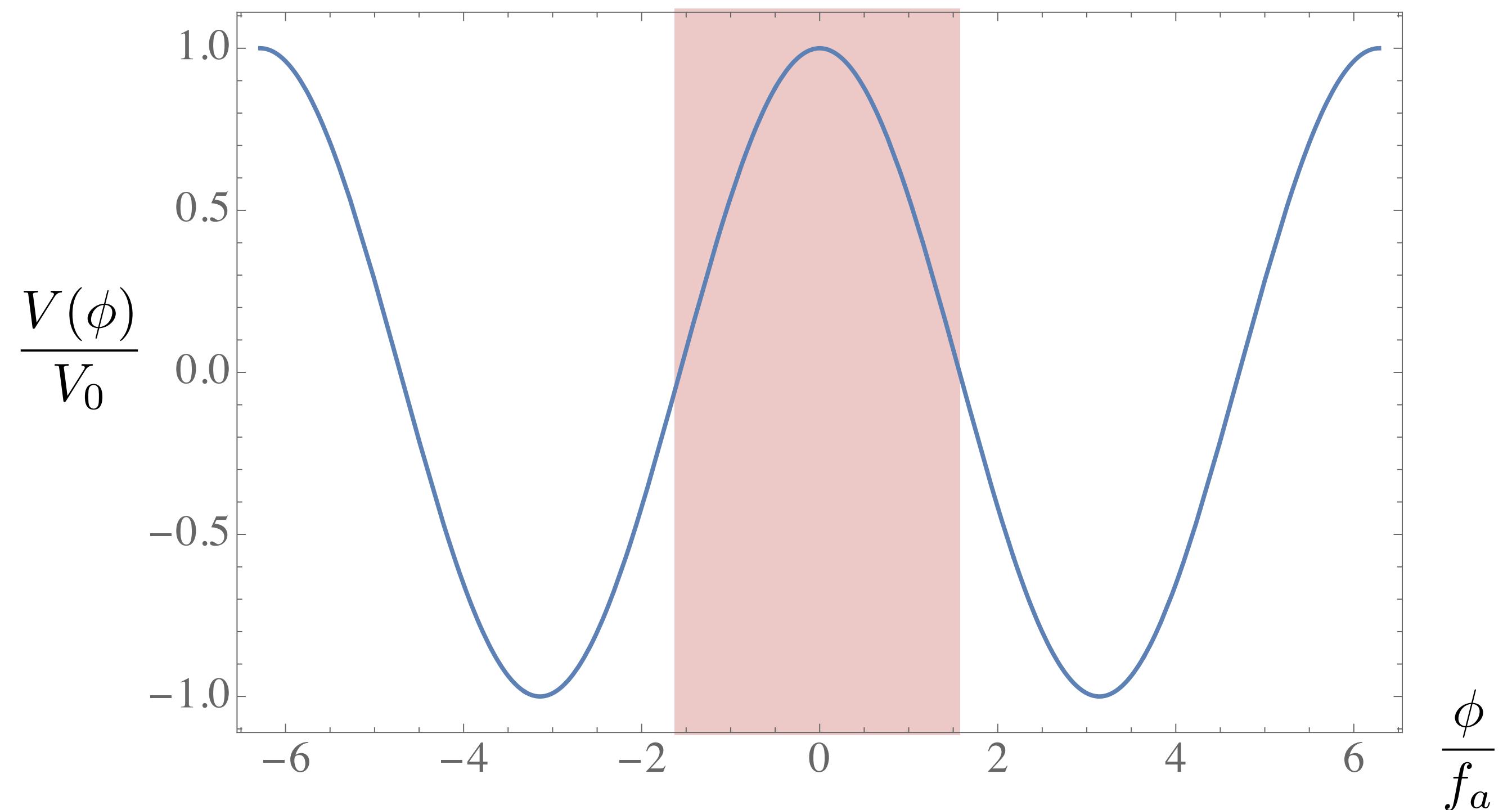
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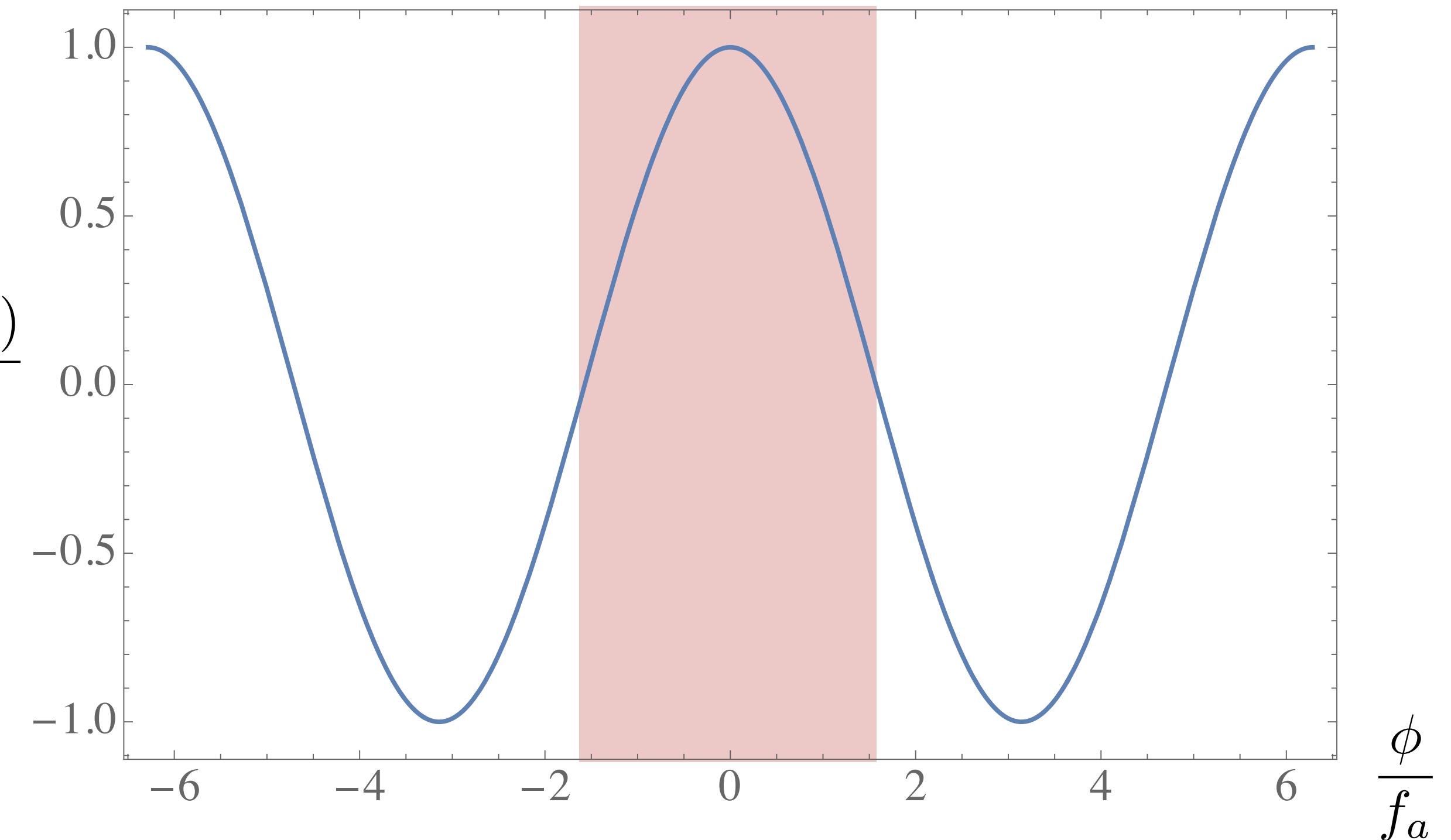
$$V = V_0 \left(1 - \cos \frac{\phi}{f_a} \right)$$

[Banks et al. '03]

$$\frac{f_a}{M_P} \sim \frac{(g_s)^p}{vol^q}$$

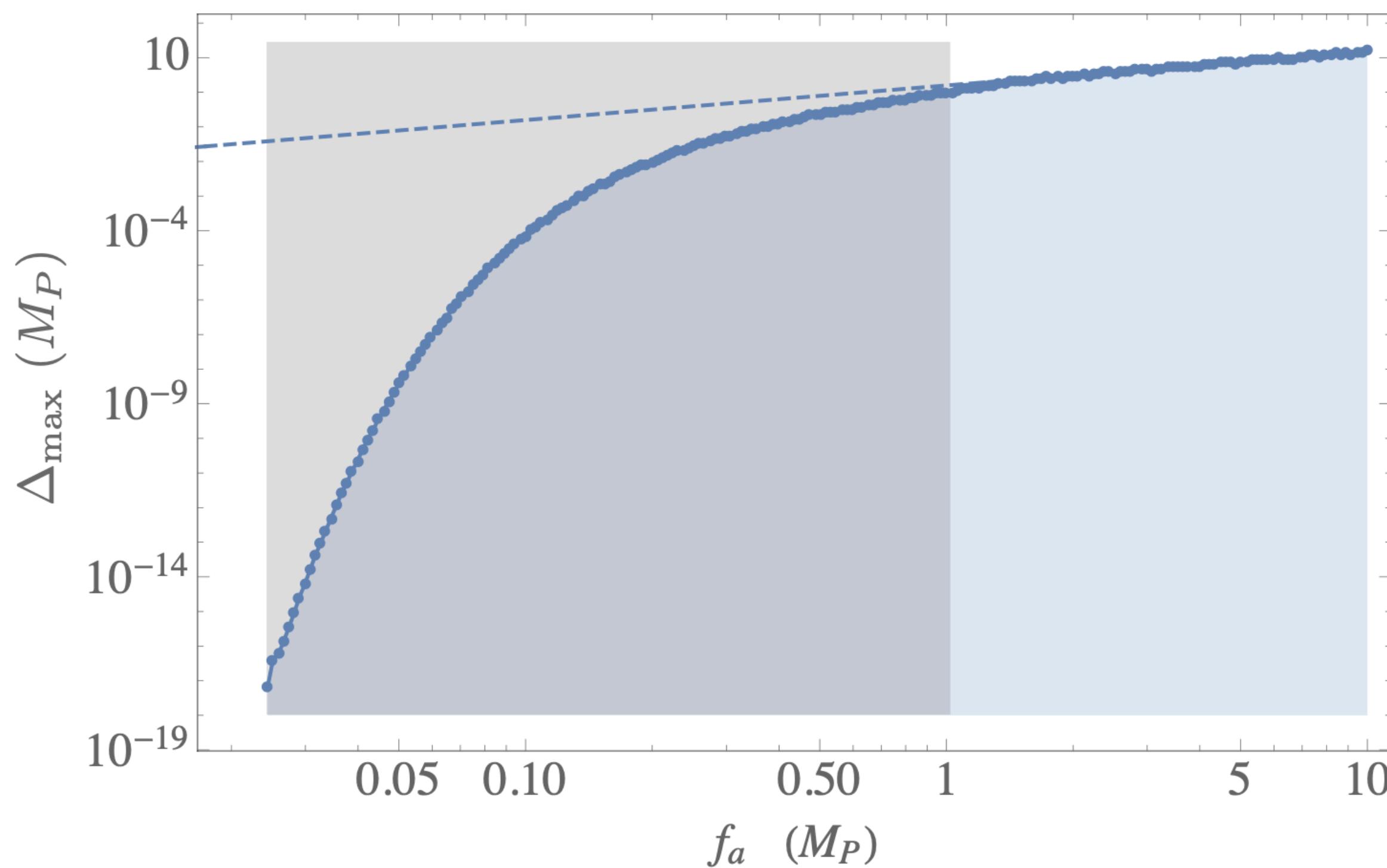
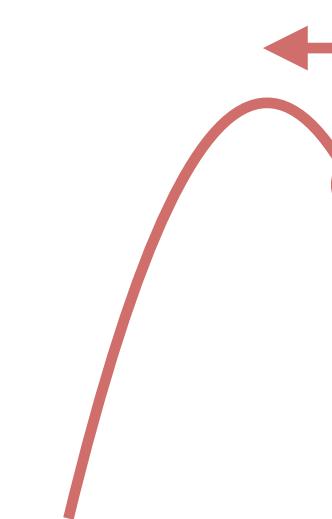
axionic WGC: $S_a f_a < 1$

$$\begin{aligned} f_a &< M_P \\ \phi &> \phi_{ip} \end{aligned}$$



Axionic DE

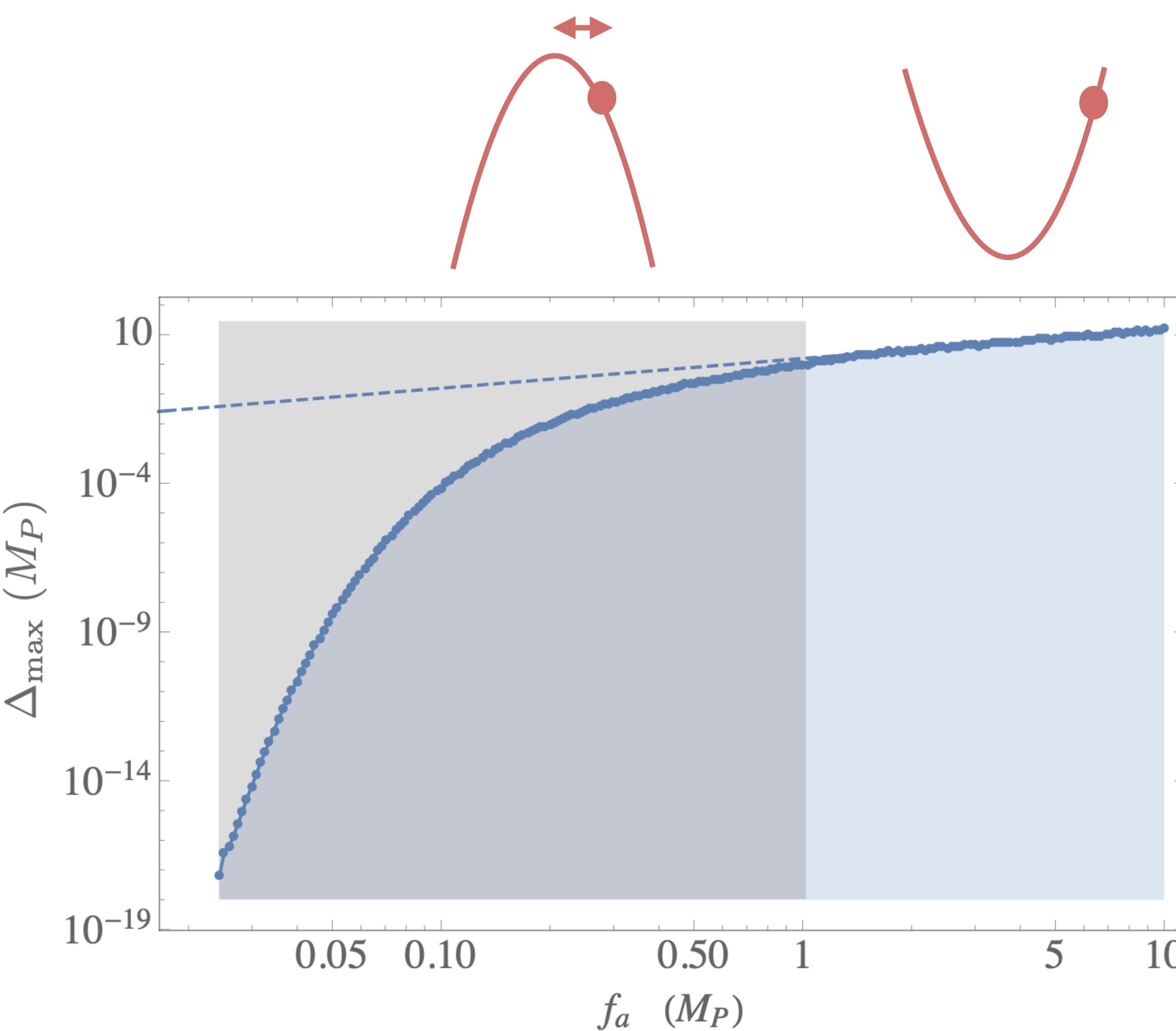
$$\Delta_{max} = |\phi_i - \phi_{max}|$$



[Dutta and Scherrer '08]
[Cicoli, Cunillera, Padilla, FGP '21]

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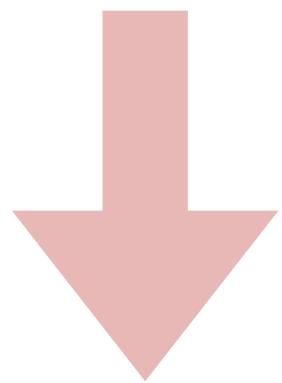
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Small decay constant



Start very close to top of the hill

tuning of initial conditions

Axionic DE

Axion is classically frozen until today

[Cicoli, Cunillera,Padilla,FGP '21]

$$\frac{\partial \phi}{\partial N} = -\frac{V_\phi}{3H_{\text{inf}}^2}$$

$$V_\phi \ll H_{\text{inf}}^2 \quad \frac{\partial \phi}{\partial N} = 0$$

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Quantum diffusion during inflation

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[Cicoli, Cunillera,Padilla,FGP '21]

$$\frac{\partial \phi}{\partial N} = -\frac{V_\phi}{3H_{\text{inf}}^2} + \frac{H_{\text{inf}}}{2\pi} \xi$$

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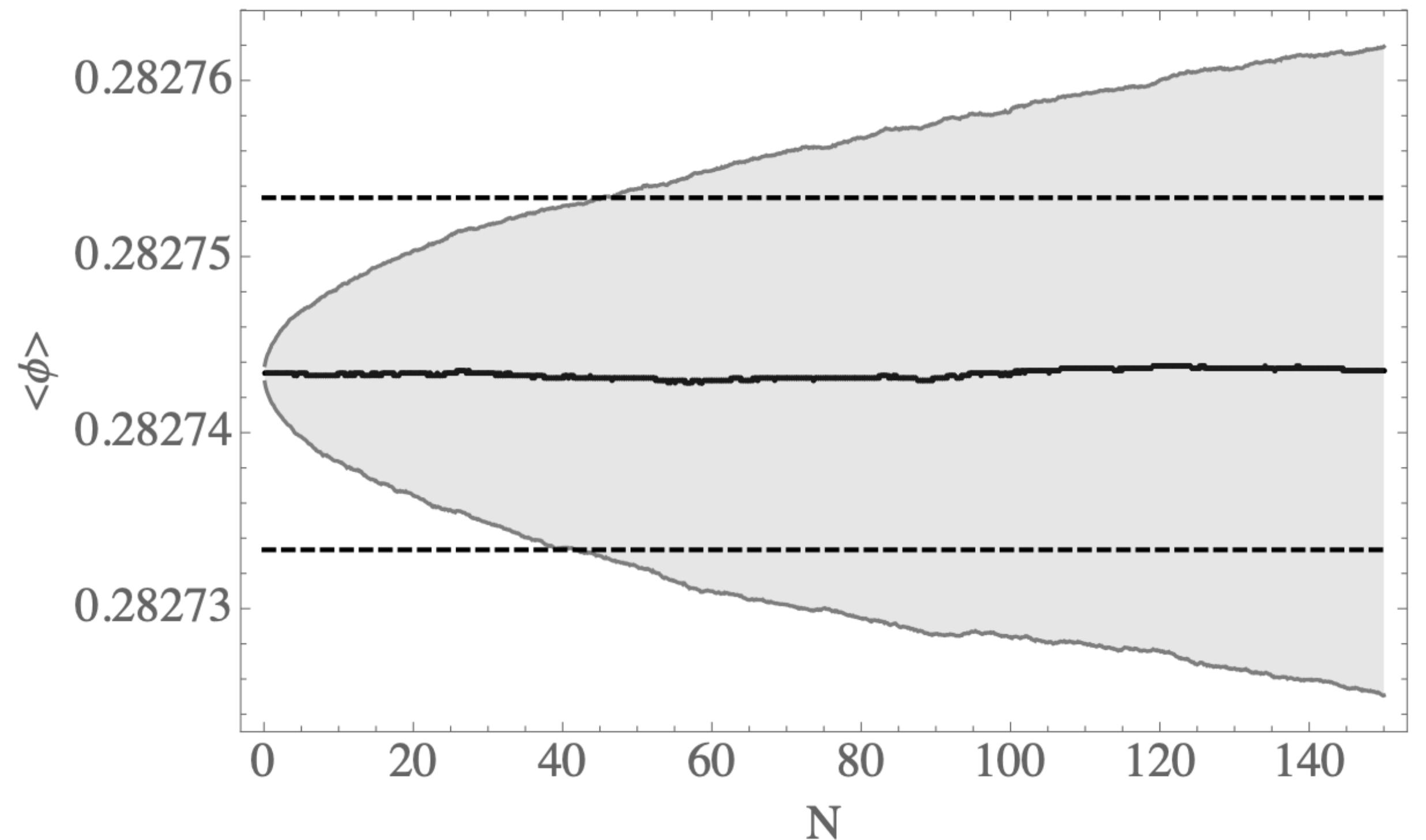
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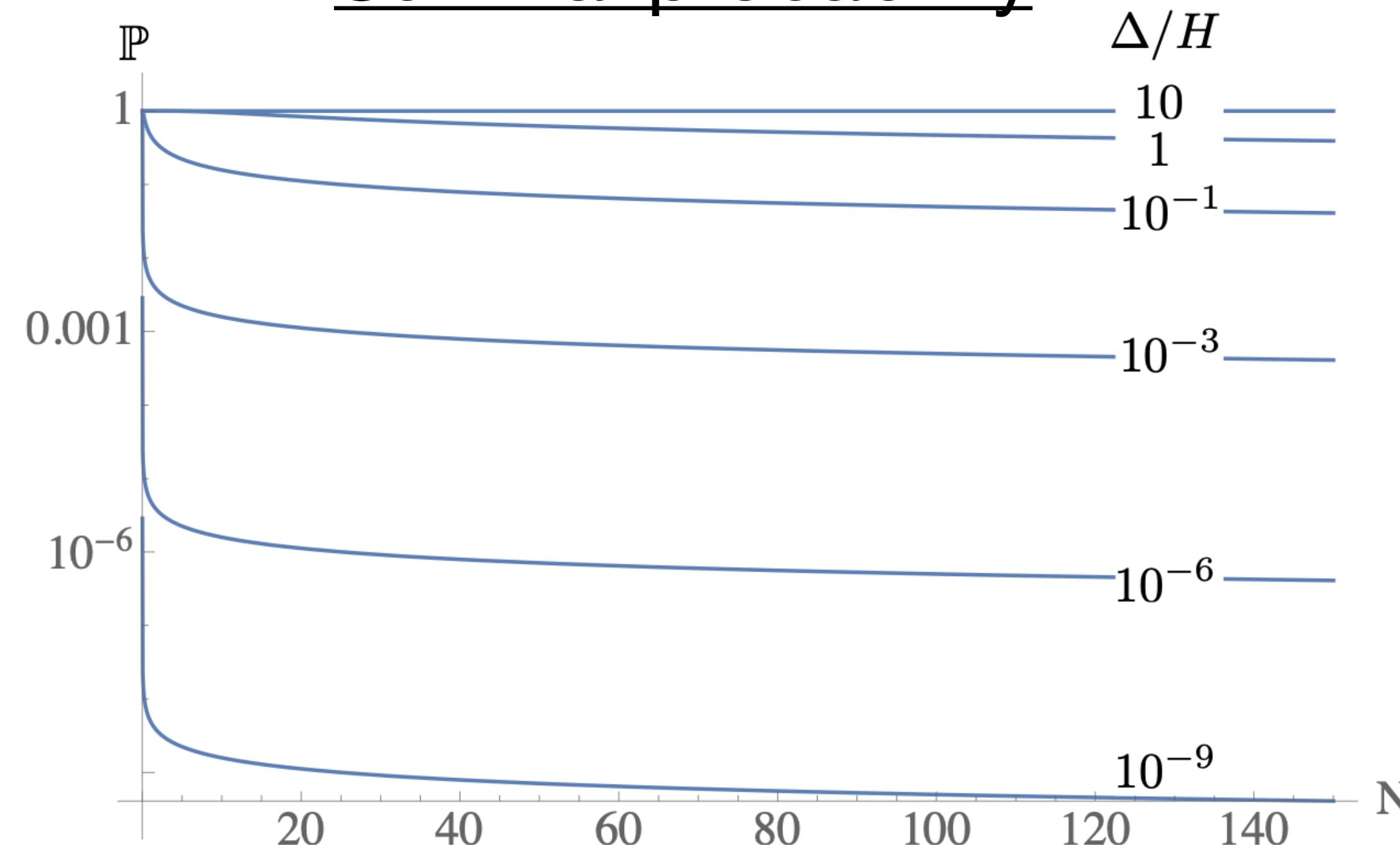
Choice of ics gets blurred during inflation



Axionic DE

[Cicoli, Cunillera,Padilla,FGP '21]

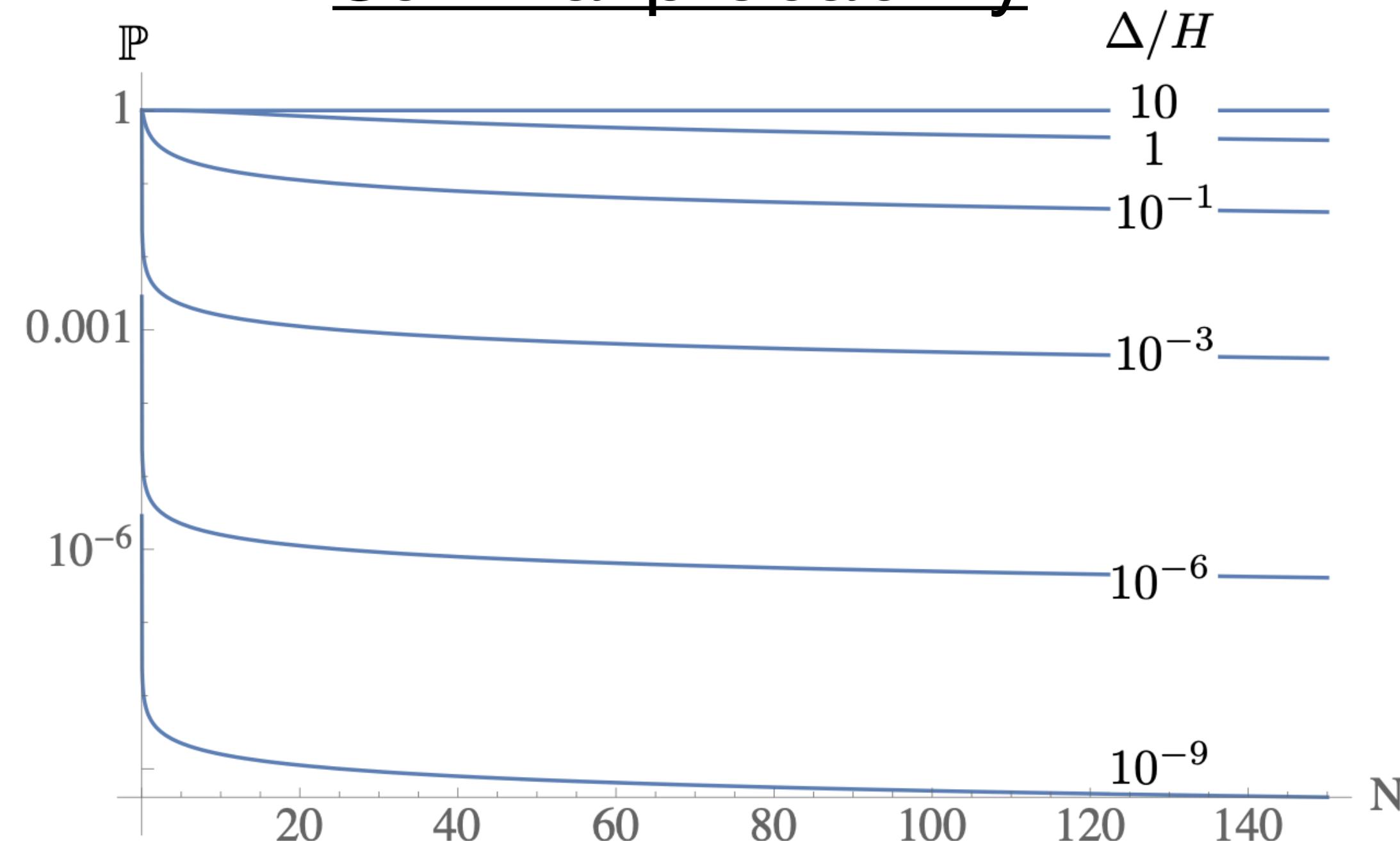
Survival probability



Axionic DE

[Cicoli, Cunillera,Padilla,FGP '21]

Survival probability



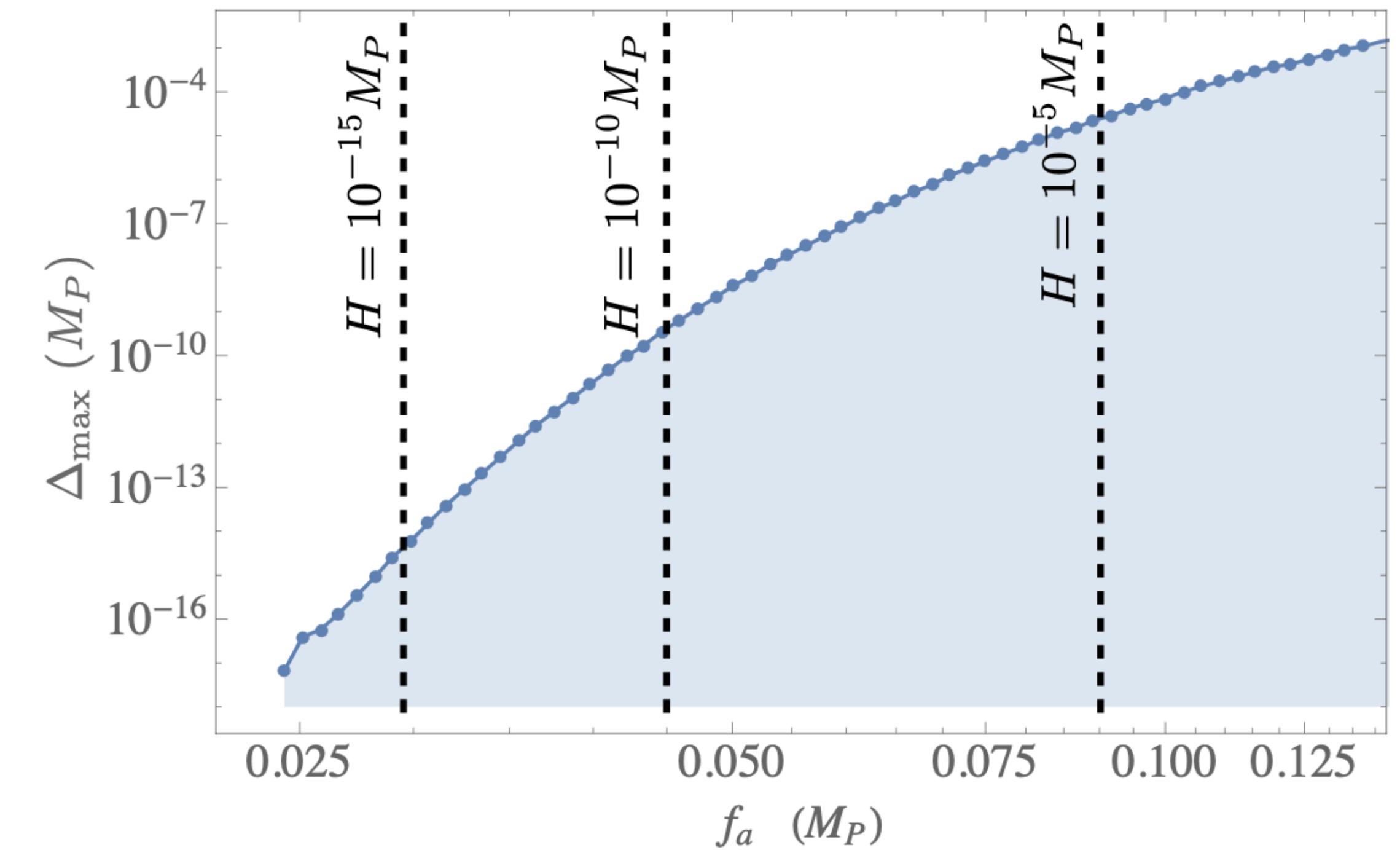
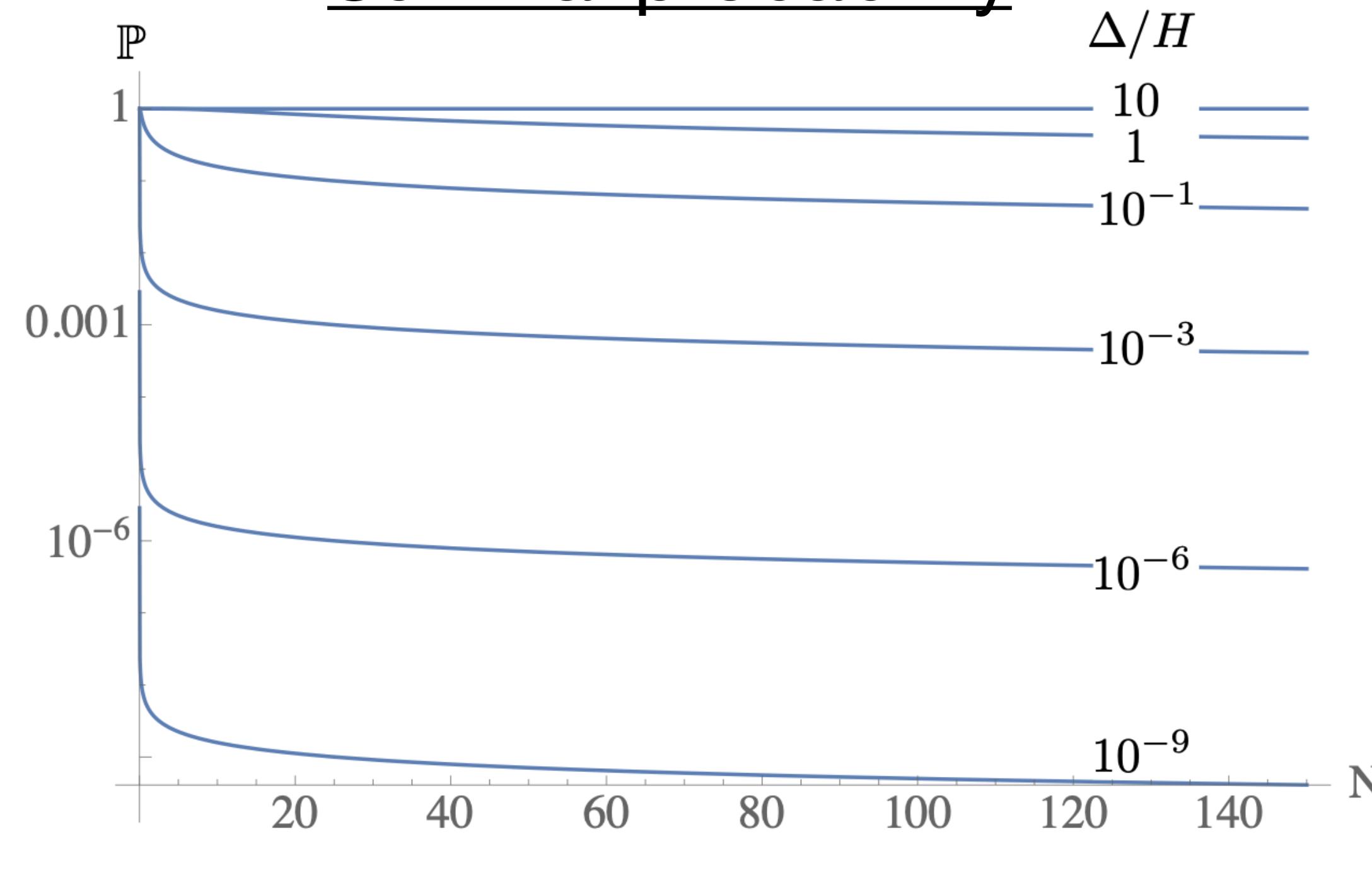
Safe from diffusion if

$$\Delta_{max} > H_{inf}$$

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[Cicoli, Cunillera,Padilla,FGP '21]

Survival probability



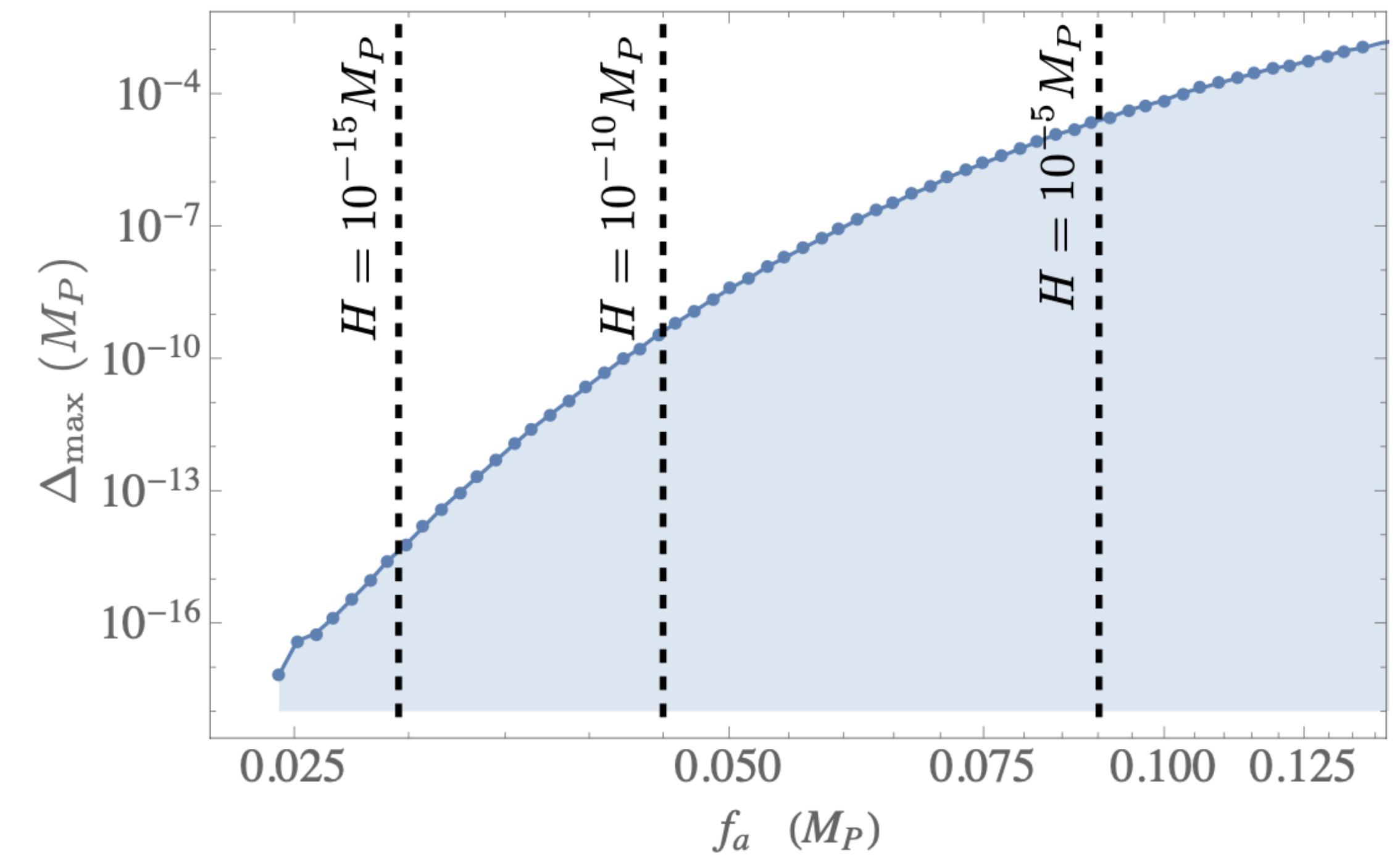
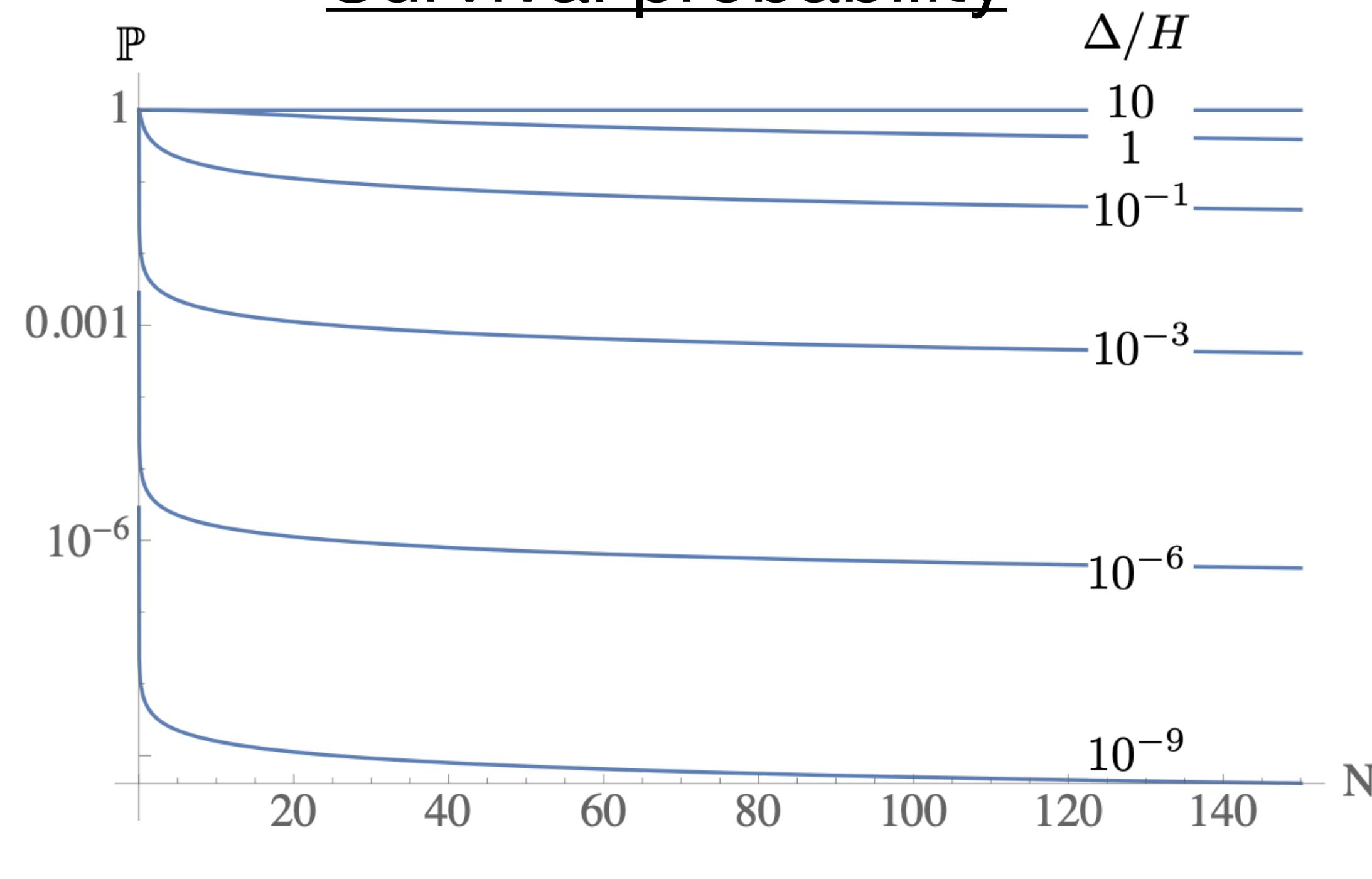
Safe from diffusion if

$$\Delta_{\max} > H_{inf}$$

Axionic DE

[Cicoli, Cunillera,Padilla,FGP '21]

Survival probability



Safe from diffusion if

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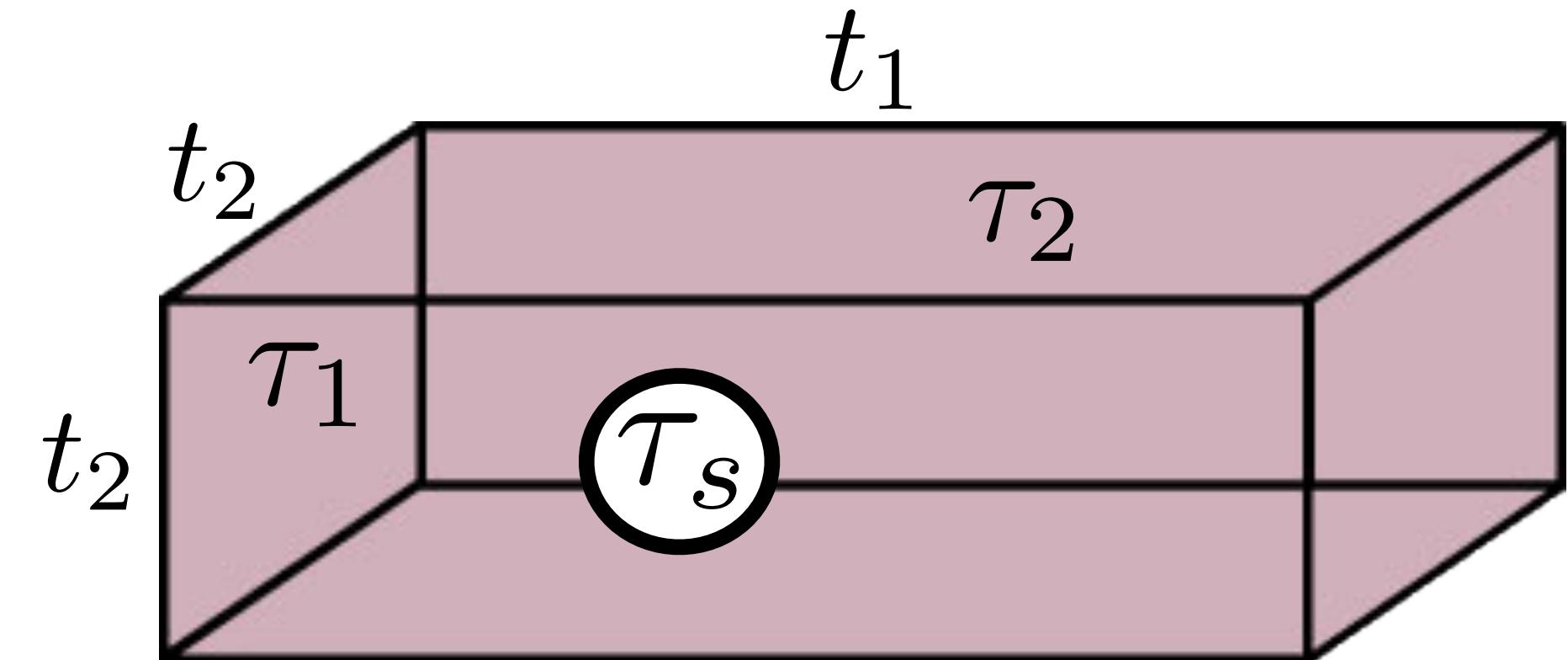
example: $H_{\inf} \sim 10^{-5} M_P \rightarrow f_a > 0.08 M_P$

Axionic DE embedding

[Cicoli, Cunillera,Padilla,FGP '24]

Embedding into type IIB compactifications

$$\mathcal{V} = t_1 t_2^2 - t_s^3 = \sqrt{\tau_1 \tau_2} - \tau_s^{3/2}$$



$$T_j = \tau_j + i\theta_j \quad \text{3 moduli + 3 axions}$$

Theory defined by

$$K = -2 \ln \left(\mathcal{V} + \frac{\xi}{2g_s^{3/2}} \right) + K_{g_s}$$

$$W_{\text{np}} = A_s e^{-a_s T_s} + A_2 e^{-a_2 T_2 + A_1 e^{-a_1 T_1}}$$

Axionic DE embedding

Hierarchical structure

$$V = V_{LVS}(\mathcal{V}, \tau_s, \theta_s) + V_{inf}(\tau_1/\tau_2) + V_{late}(\theta_1, \theta_2)$$

→ smaller →

LVS stabilisation α'^3 corrections + n.p. effects

[Conlon et al. '05]

$$V_{\text{vol}} = \frac{\kappa}{\mathcal{V}^n} + \frac{3\xi|W_0|^2}{4g_s^{3/2}\mathcal{V}^3} - 4|W_0|A_s a_s \tau_s \frac{e^{-a_s \tau_s}}{\mathcal{V}^2} + 4\sqrt{2\hat{k}} A_s^2 a_s^2 \sqrt{\tau_s} \frac{e^{-2a_s \tau_s}}{\mathcal{V}}.$$

Axionic DE embedding

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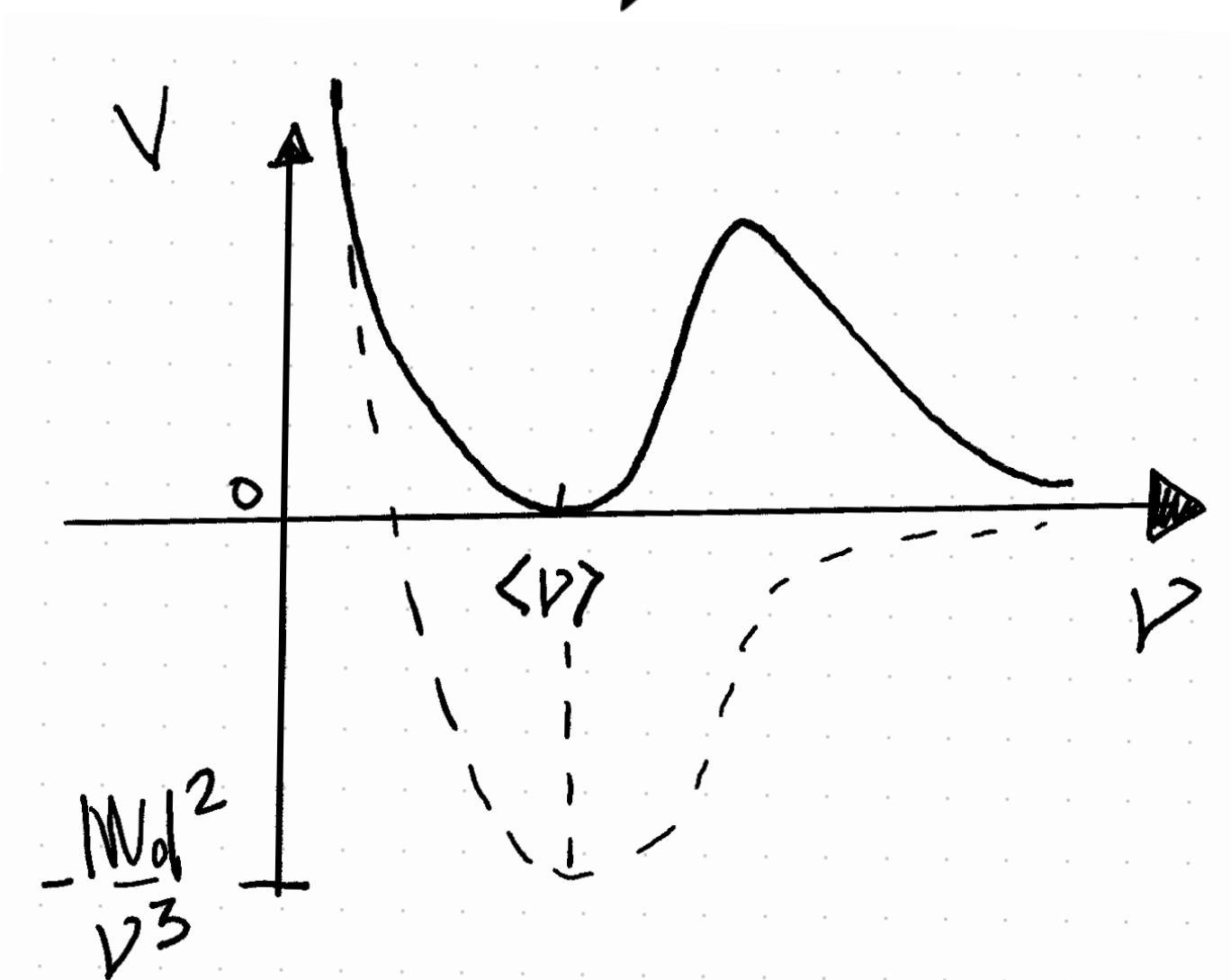
uplift

non SUSY AdS

$$\langle V \rangle \sim -\frac{|W_0|^2}{\mathcal{V}^3}$$

$$\langle \tau_s \rangle \propto 1/g_s$$

$$\langle \mathcal{V} \rangle \propto e^{1/g_s}$$



Axionic DE embedding

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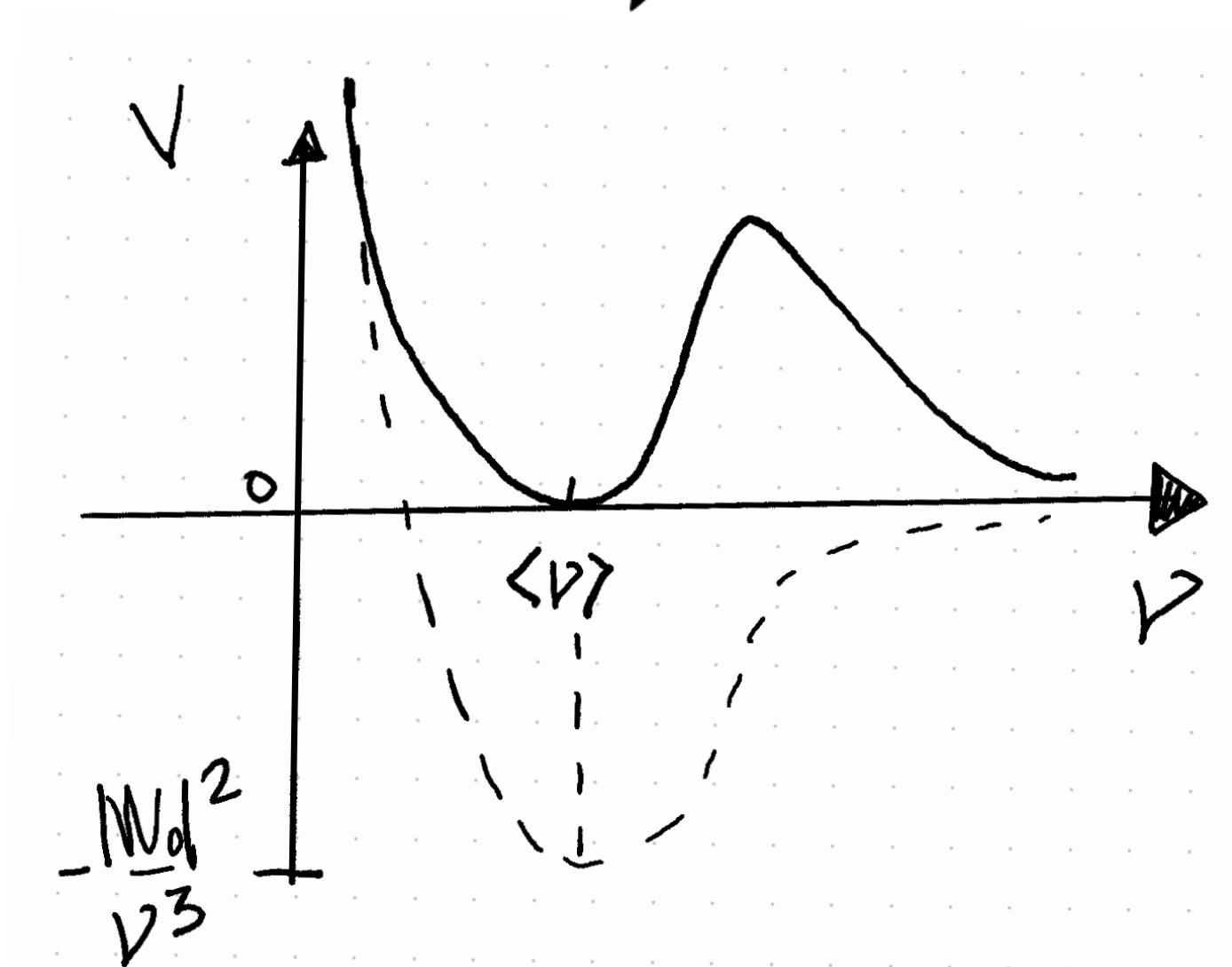
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$$\langle \tau_s \rangle \propto 1/g_s$$

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Mass hierarchy

$$m_{\tau_s} \sim m_{\theta_s} \sim \frac{M_P}{\mathcal{V}} \gg m_{\mathcal{V}} \sim \frac{M_P}{\mathcal{V}^{3/2}}$$



Axionic DE embedding

Fibre inflation: string loops (W) + HD

[Cicoli et al. '08]

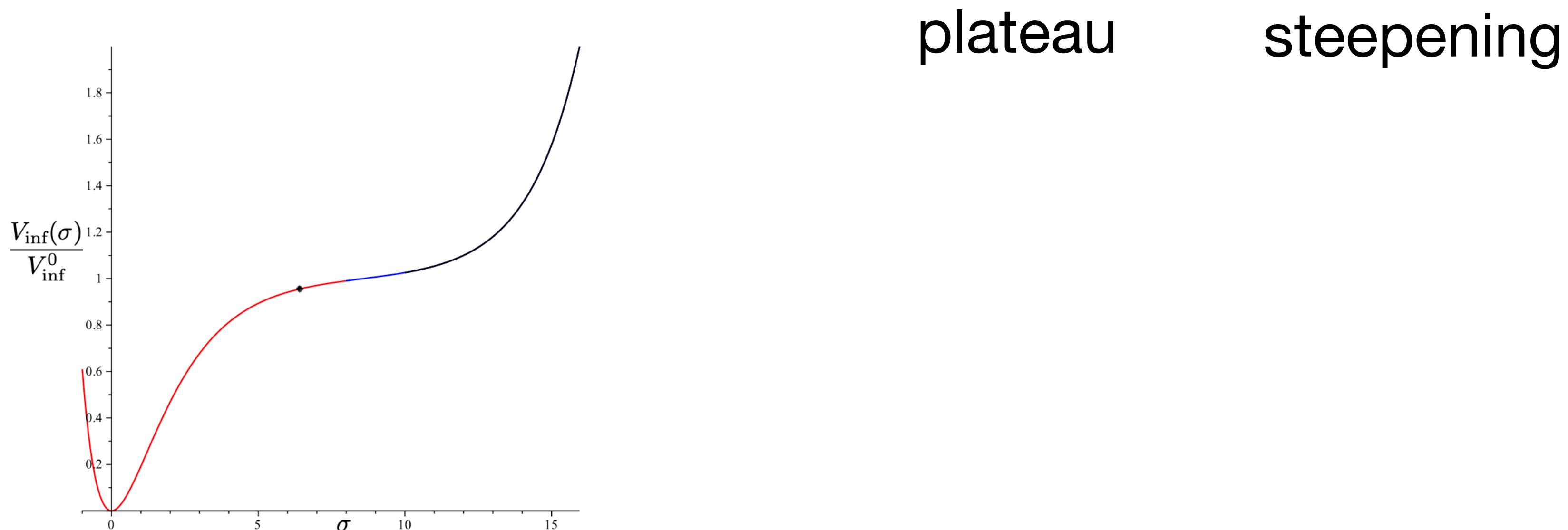
$$K_{g_s} = \sum_i g_s \frac{C_i(U, \bar{U}) t_i^\perp}{\mathcal{V}} + \sum_i \frac{\tilde{C}_i(U, \bar{U})}{t_i^\cap \mathcal{V}}.$$

$$V_{\text{hd}} = -\frac{\lambda}{g_s^{3/2}} \frac{|W_0|^4}{\mathcal{V}^4} \Pi_i t_i$$

$$\sigma \sim \log(\tau_1/\tau_2)/\sqrt{3}$$

$$V_{\text{inf}}(\sigma) = V_0 \left[\left(1 - e^{-\frac{\sigma}{\sqrt{3}}}\right)^2 - 2\mathcal{R} \left(1 - \cosh\left(\frac{\sigma}{\sqrt{3}}\right)\right) \right]$$

$$V_0 \propto \frac{|W_0|^4}{\mathcal{V}^{11/3}}$$



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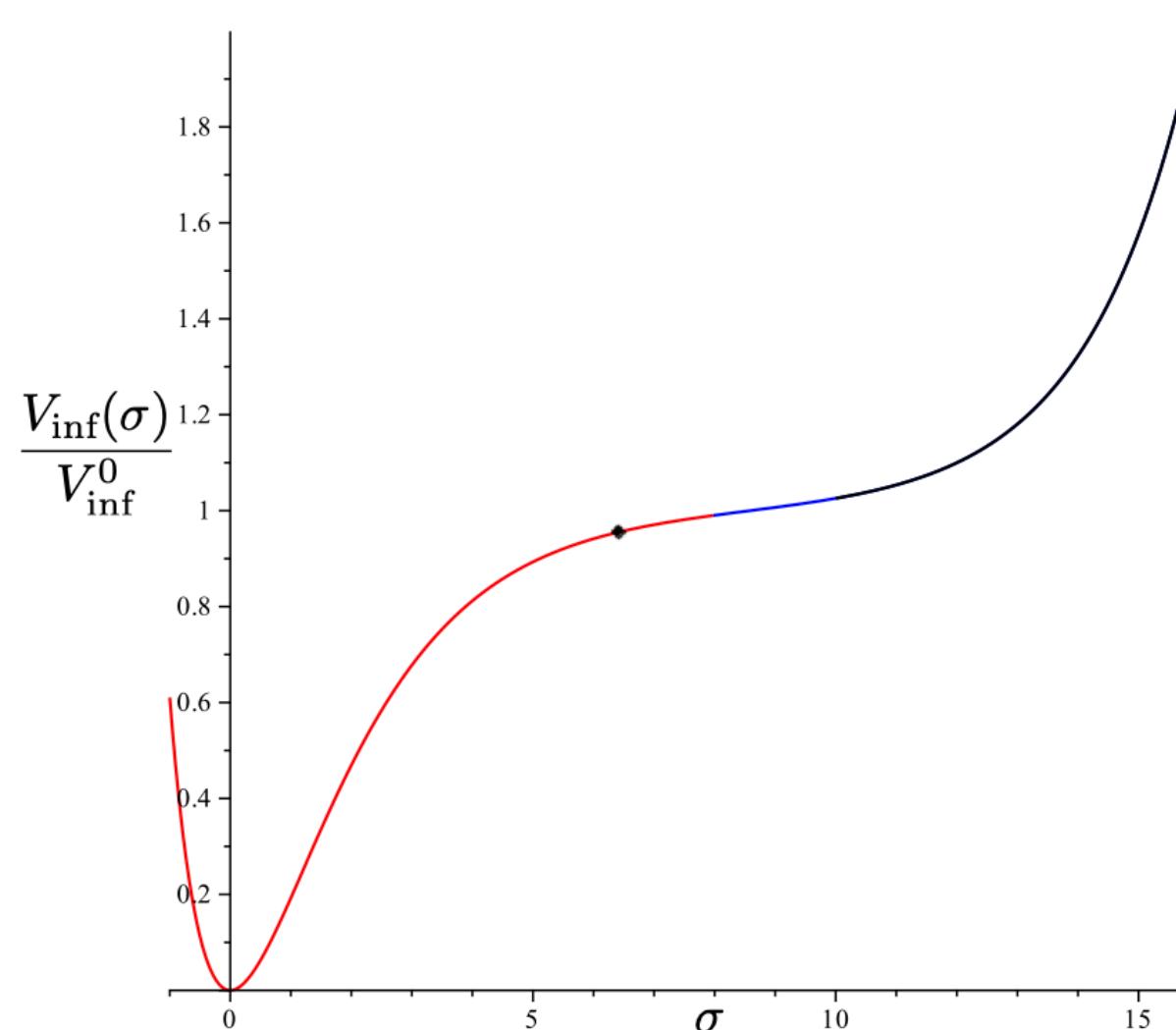
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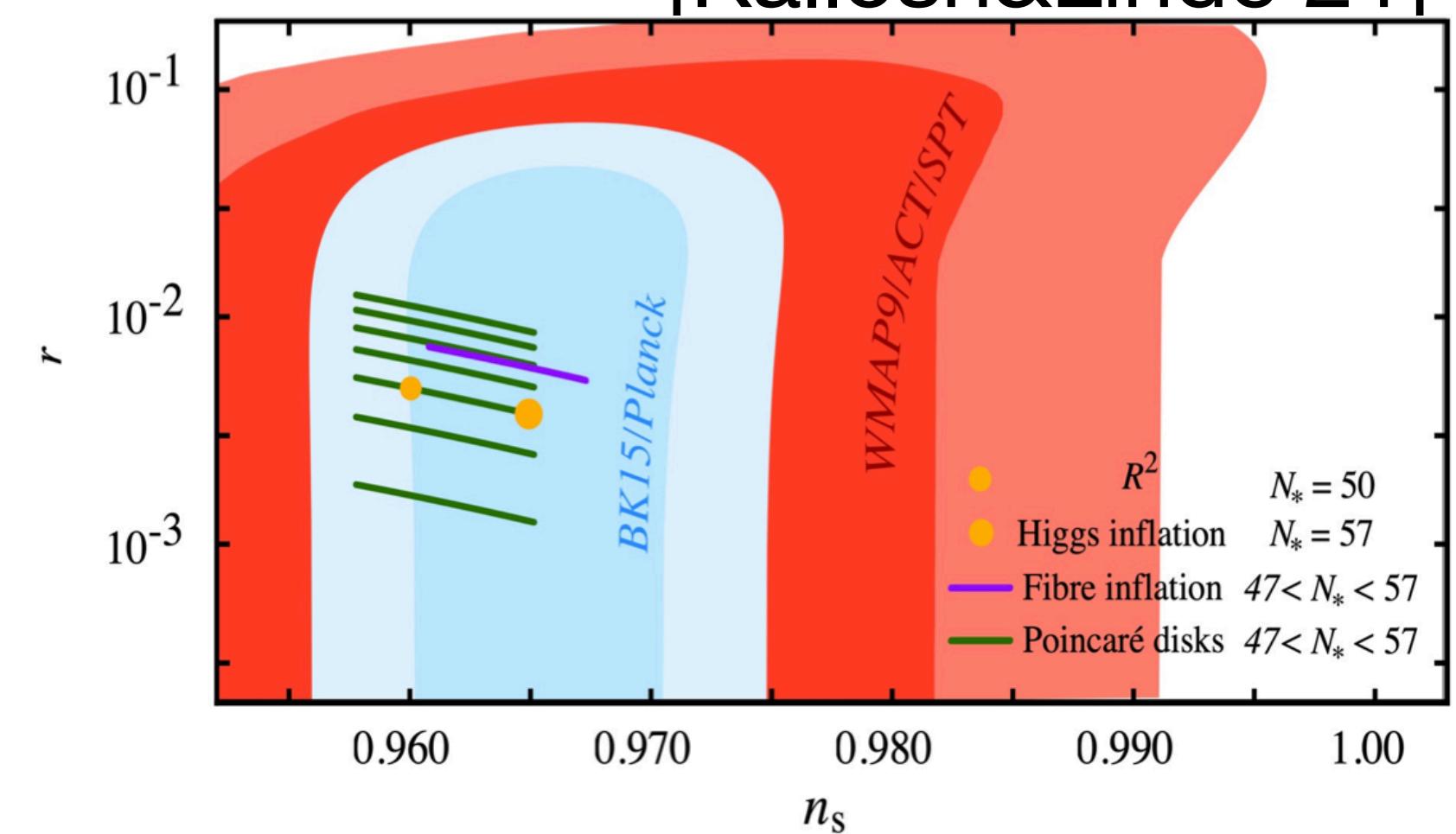


plateau steepening

$$H_{inf} \sim 10^{-5} M_P$$

$$\mathcal{V} \sim \mathcal{O}(10^3)$$

[Kallosh&Linde'21]



Axionic DE embedding

Fibre inflation: string loops (W) + HD

[Cicoli et al. '08]

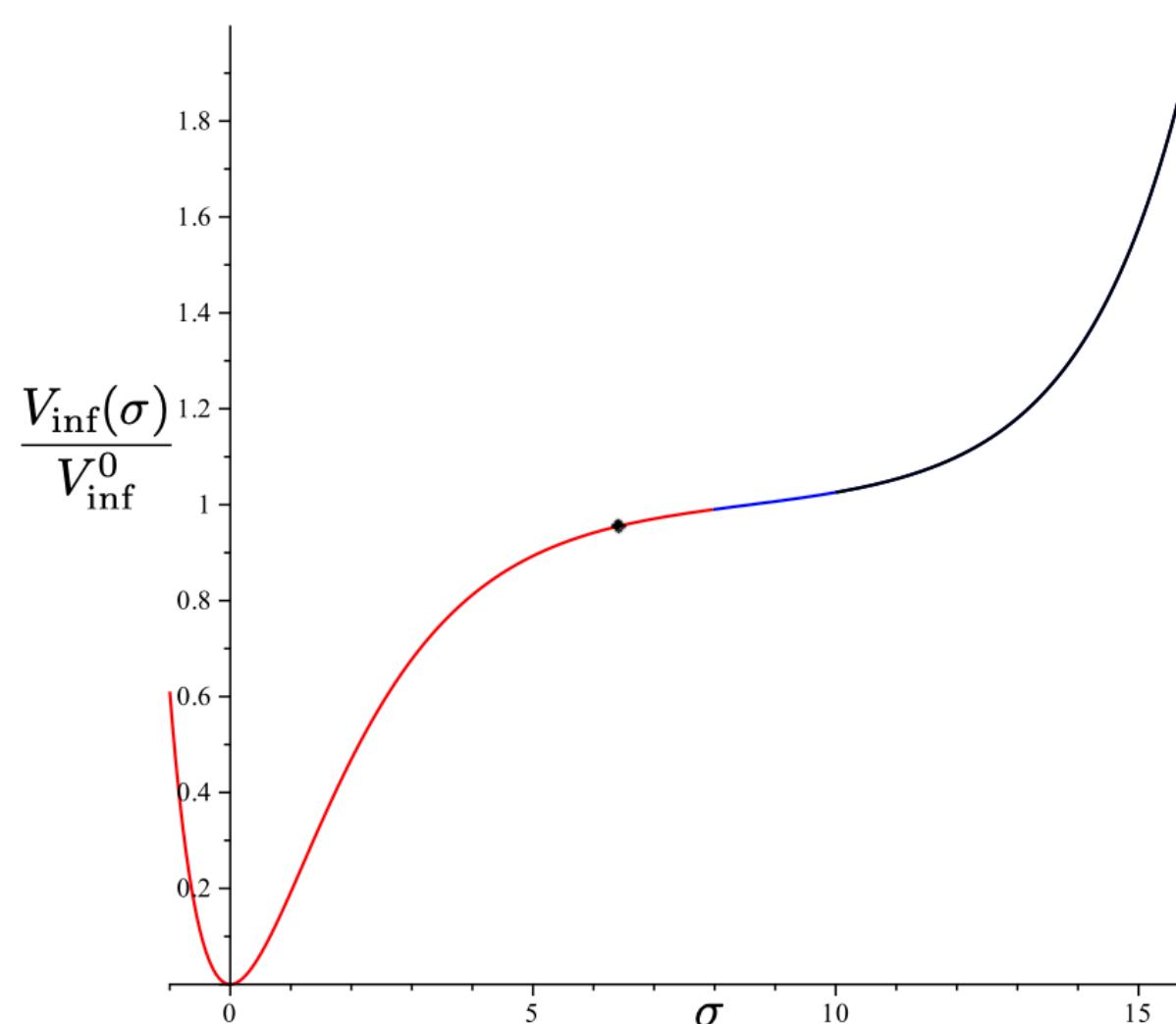
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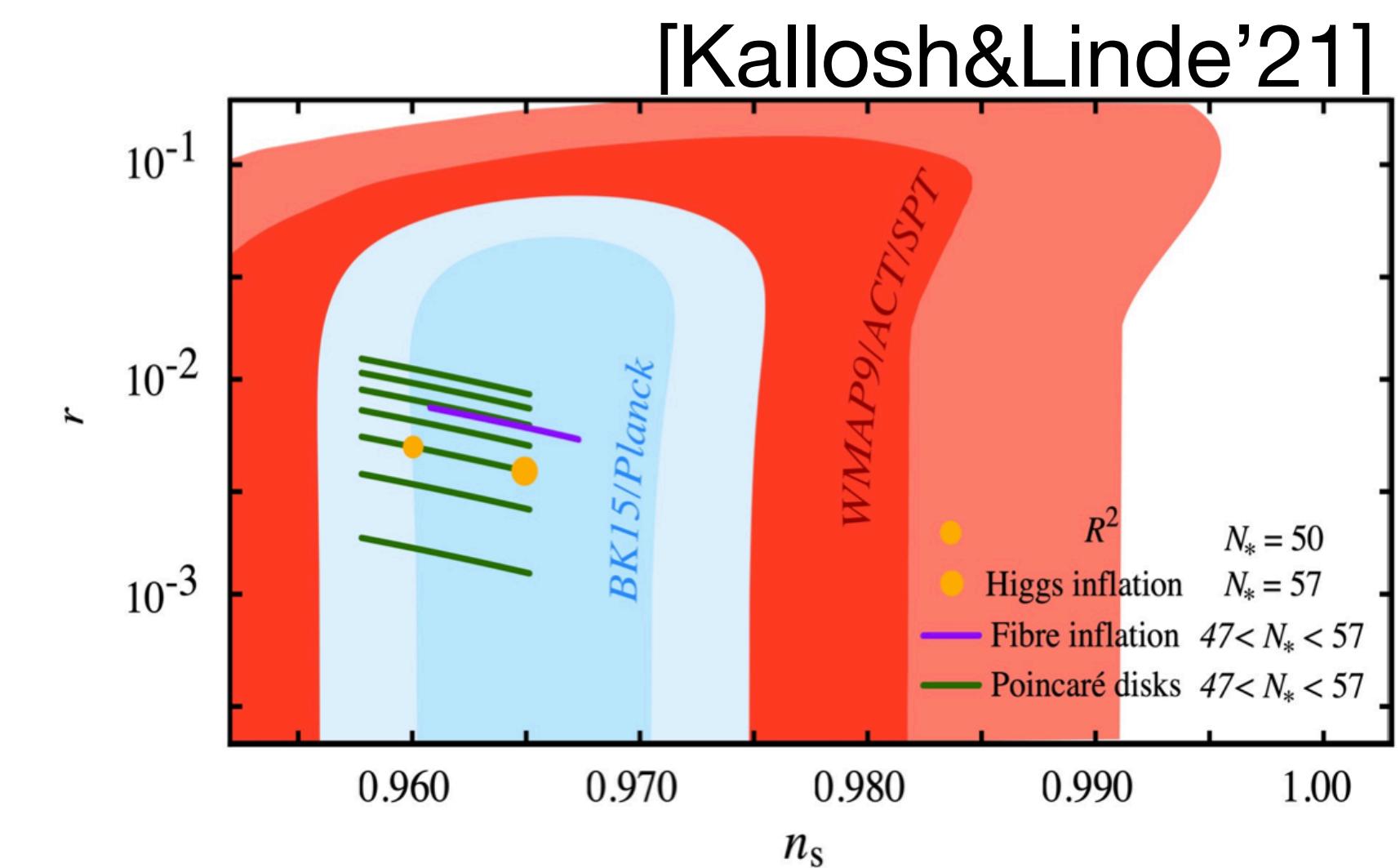
plateau steepening

$$H_{\text{inf}} \sim 10^{-5} M_P$$

$$\mathcal{V} \sim \mathcal{O}(10^3)$$

single field inflation

$$m_{\theta_s} \sim m_{\tau_s} \sim m_{3/2} \gg m_V > m_\sigma \sim H_{\text{inf}}$$



[Kallosh&Linde'21]

Axionic DE embedding I

Late time potential generated by poly-instanton

$$\begin{aligned} W_{\text{np}} &= A_s e^{-a_s T_s} + \boxed{A_2 e^{-a_2 T_2 + A_1 e^{-a_1 T_1}}} & [\text{Blumenhagen et al. '08-'12}] \\ &= A_s e^{-a_s T_s} + A_2 e^{-a_2 T_2} + A_2 A_1 e^{-(a_2 T_2 + a_1 T_1)} + \dots & [\text{Lüst, Zhang. '13}] \end{aligned}$$

$$V_{\text{late}} = \Lambda_2^4 [1 - \cos(a_2 \theta_2)] + \Lambda_1^4 [1 - \cos(a_2 \theta_2 + a_1 \theta_1)]$$

$$\Lambda_2^4 = \frac{4W_0 A_2}{\mathcal{V}^2} (a_2 \langle \tau_2 \rangle) e^{-a_2 \langle \tau_2 \rangle} \quad \Lambda_1^4 = \frac{4W_0 A_1 A_2}{\mathcal{V}^2} (a_1 \langle \tau_1 \rangle + a_2 \langle \tau_2 \rangle) e^{-a_1 \langle \tau_1 \rangle - a_2 \langle \tau_2 \rangle}$$

Hierarchy: $a_i \tau_i > 1$ $\Lambda_1^4 \ll \Lambda_2^4$

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Hierarchy:

$$a_i \tau_i > 1$$

$$\Lambda_1^4 \ll \Lambda_2^4$$

$$f_1 = \frac{1}{\sqrt{2} a_1 \langle \tau_1 \rangle} \quad f_2 = \frac{1}{a_2 \langle \tau_2 \rangle}$$

$$V_{DE} = \Lambda_1^4 \left(1 - \cos \frac{\phi_1}{f_1} \right)$$

$$f_1, f_2 \ll 1$$

Axionic DE embedding I

Model's structure:

$$\Lambda_2^4 = \Lambda_2^4(f_2) \quad \Lambda_1^4 = \Lambda_1^4(f_1, f_2)$$

Example: safe ics $\leftrightarrow f_1 = 0.085$ DE scale $\Lambda_1^4 \sim 10^{-120} \leftrightarrow f_2 = 0.0038$

$\Lambda_2^4 \sim 10^{-117}$

Axionic DE embedding I

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axionic mass hierarchy

$$m_2 \sim 10^{-56} = 10^{-29} eV$$

$$m_1 \sim 10^{-59} = 10^{-32} eV$$

$$< H_{eq} \sim 10^{-27} eV$$

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ϕ_2 can only be a subdominant component of DM

$$\frac{\Omega_2}{\Omega_m} \simeq \frac{3}{2} \left(\frac{\delta\phi_2}{M_p} \right)^2 \rightarrow \leq 0.0002$$

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ϕ_2 can only be a subdominant component of DM $\frac{\Omega_2}{\Omega_m} \simeq \frac{3}{2} \left(\frac{\delta\phi_2}{M_p} \right)^2 \rightarrow \leq 0.0002$

$$a_i = \frac{2\pi}{N_i}$$

$$N_1 = 1$$

$$\langle \tau_1 \rangle = 1.3$$

$$N_2 = 27$$

$$\langle \tau_2 \rangle \sim 10^3$$

$$N_1 = 5$$

$$\langle \tau_1 \rangle \sim 5$$

$$N_2 = 10$$

$$\langle \tau_2 \rangle \sim 500$$

Summary

- From a bottom up perspective dynamical DE is no easier than dS
- But if observationally $\omega \neq -1$ or theoretically no dS in QG
- Two alternatives: multifield quintessence
 - ▶ acceleration in steep potentials
 - ▶ string embedding?
- axionic quintessence
 - ▶ radiative stability (shift symmetry)
 - ▶ scale suppression
 - ▶ evades fifth force
 - ▶ $f_a < M_P$ hilltop

Thank you

Backup slides

Axionic DE

Simplest models

$$V = e^{-S} \cos\left(\frac{\phi}{f}\right)$$

where

$$S = \frac{M_P}{f}$$

DE constrains

$$V \sim 10^{-120} M_P^4$$

$$S \sim 276$$

$$f \sim 3 \times 10^{-3} M_P$$

Requires VERY low scale inflation

Need to be more sophisticated: break the link between V_0 and f

Alignment
Poly-instantons

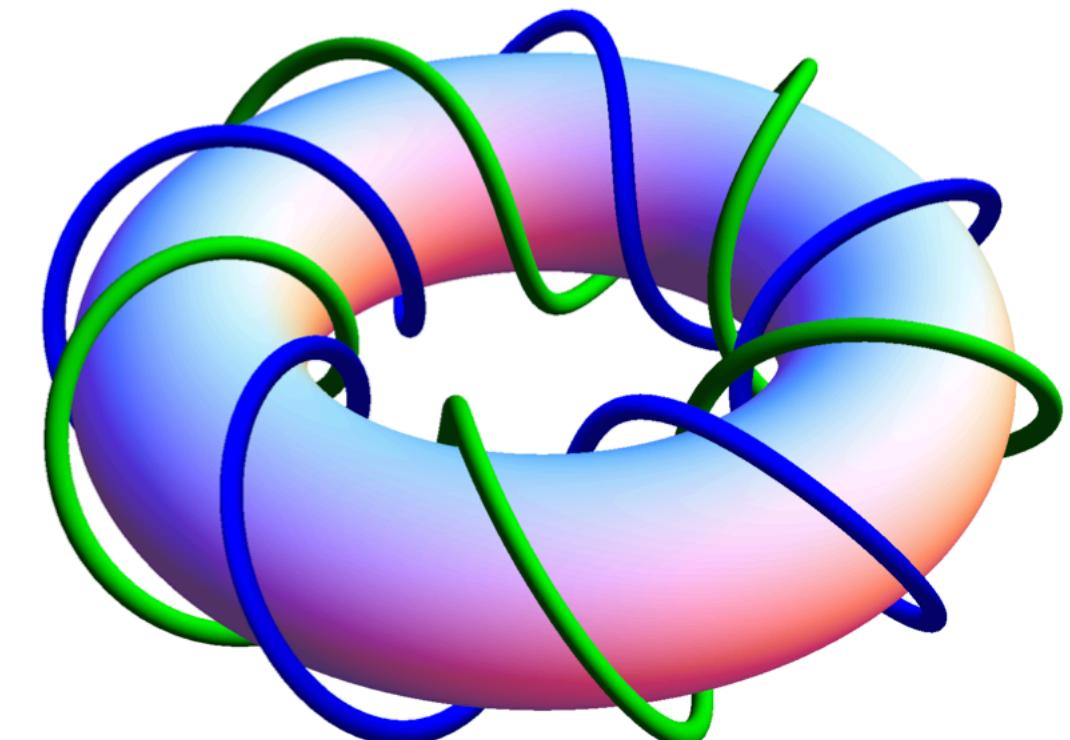
Axionic DE embedding II

Late time potential generated by KNP alignment

$$W_{\text{np}} = A_s e^{-a_s T_s} + \sum_{i=1}^3 A_i e^{-a_i(q_{i1} T_1 + q_{i2} T_2)}$$

$$V_{\text{late}} \simeq \hat{\Lambda}_1^4 \left[1 - \cos \left(\frac{q_{11}}{N_1} \frac{\phi_1}{f_1} + \frac{q_{12}}{N_1} \frac{\phi_2}{f_2} \right) \right] + \hat{\Lambda}_2^4 \left[1 - \cos \left(\frac{q_{21}}{N_2} \frac{\phi_1}{f_1} + \frac{q_{22}}{N_2} \frac{\phi_2}{f_2} \right) \right]$$

q_{ij} winding numbers



[Long et al, 2014]

$$\hat{\Lambda}_i^4 \equiv \frac{4|W_0|A_i}{\mathcal{V}^2} a_i (q_{i1} \tau_1 + q_{i2} \tau_2) e^{-a_i(q_{i1} \tau_1 + q_{i2} \tau_2)}$$

Axionic DE embedding II

Late time potential generated by KNP alignment

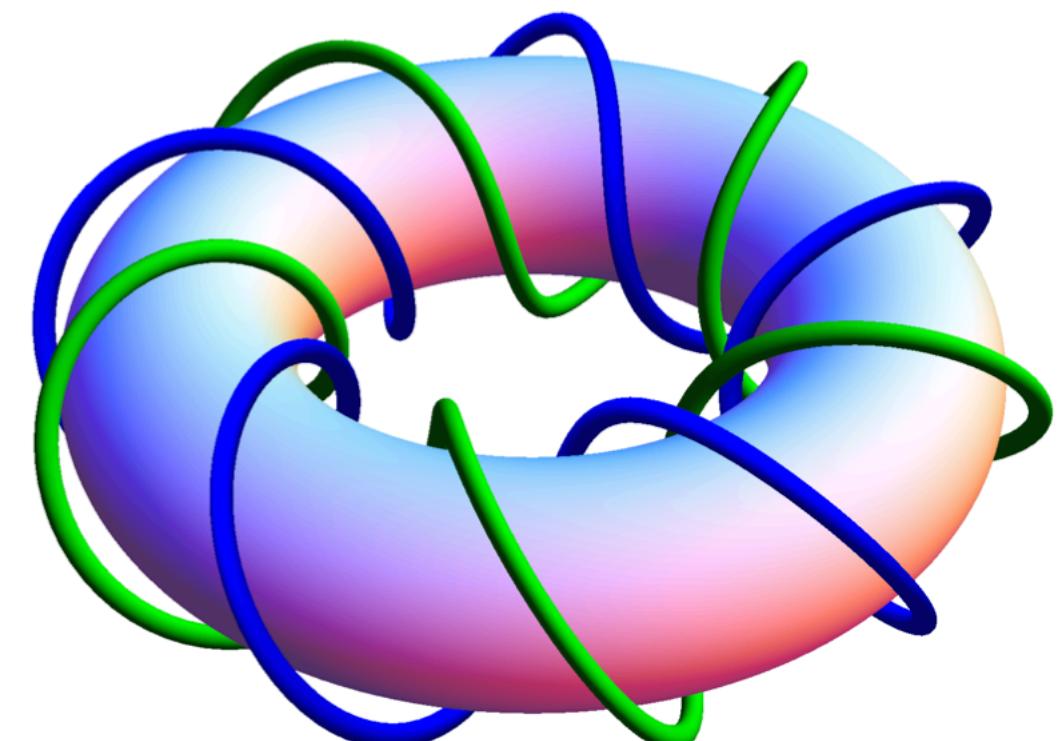
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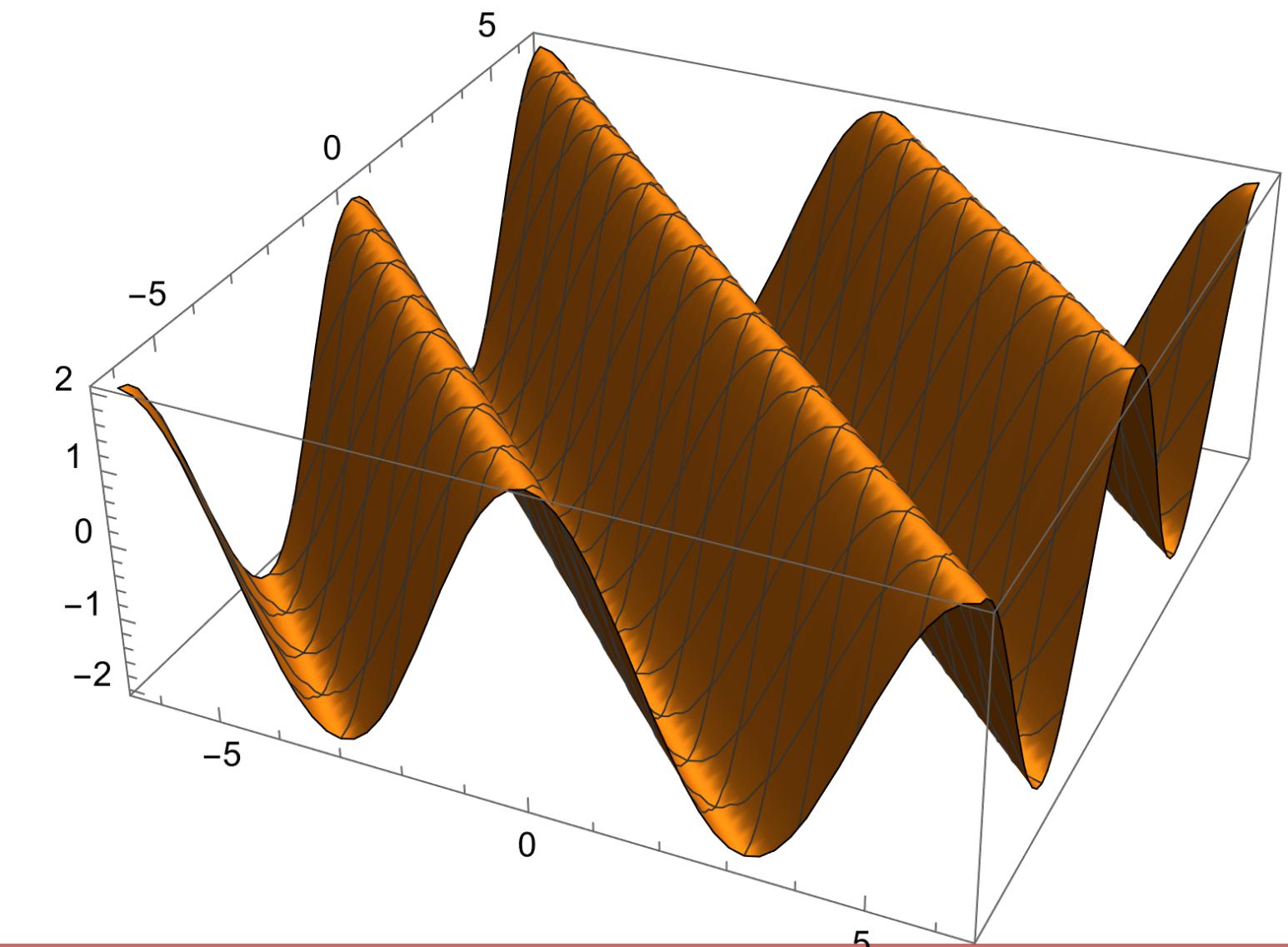
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one massless axion if $q_{11}q_{22} = q_{12}q_{21}$

q_{ij} winding numbers



[Long et al, 2014]



Axionic DE embedding II

Late time potential generated by KNP alignment

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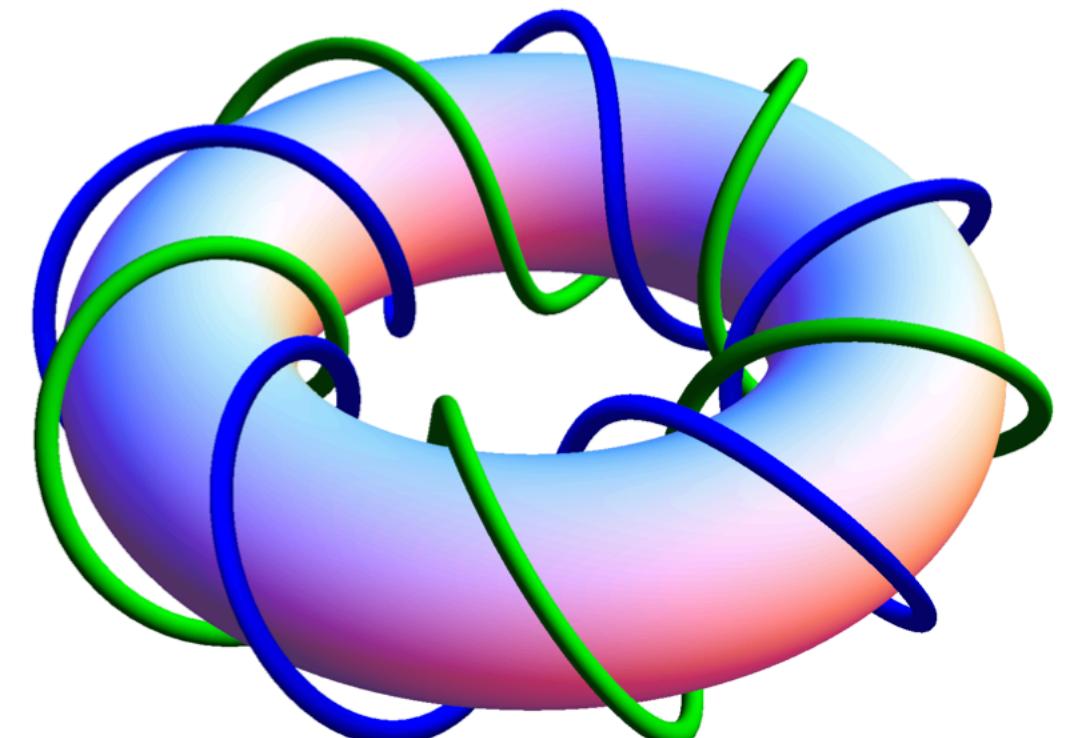
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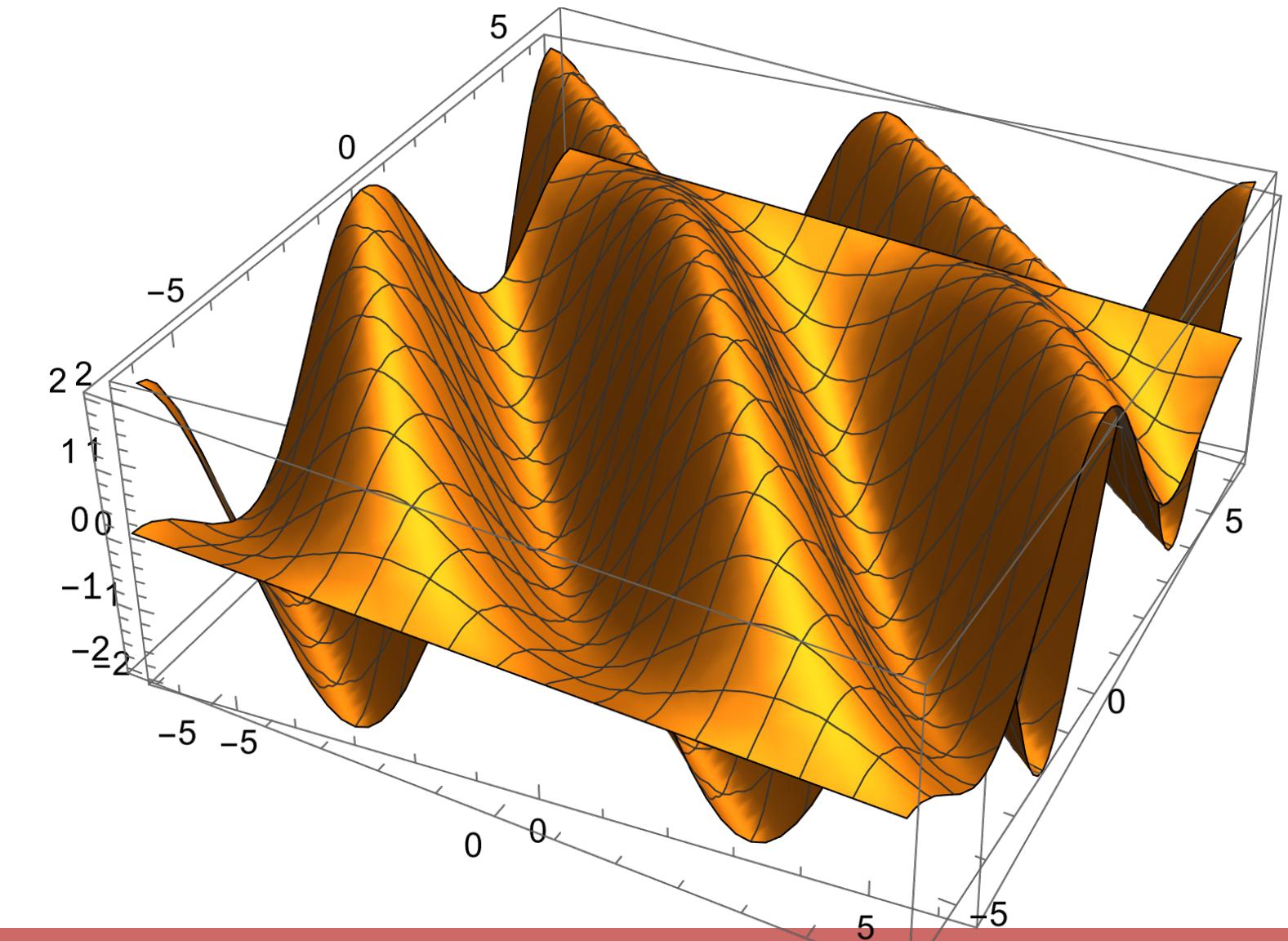
one massless axion if $q_{11}q_{22} = q_{12}q_{21}$

mass hierarchy if $q_{11}q_{22} = (1 + \delta)q_{12}q_{21}$ $\delta \ll 1$

q_{ij} winding numbers



[Long et al, 2014]



Axionic DE embedding II

In terms of mass eigenstates

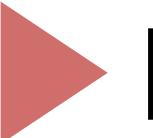
$$V_{\text{late}} \simeq \hat{\Lambda}_2^4 \left[1 - \cos \left(q_{2H} \frac{\phi_H}{\mathfrak{f}_H} \right) \right] + \hat{\Lambda}_1^4 \left[1 - \cos \left(q_{1H} \frac{\phi_H}{\mathfrak{f}_H} + q_L \frac{\phi_L}{\mathfrak{f}_L} \right) \right]$$

$$\mathfrak{f}_H \equiv \sqrt{\frac{q_{21}^2 + q_{22}^2}{q_{21}^2/\mathfrak{f}_1^2 + q_{22}^2/\mathfrak{f}_2^2}} \quad \mathfrak{f}_L \equiv \frac{\mathfrak{f}_1 \mathfrak{f}_2}{\mathfrak{f}_H}$$

$$m_H \simeq q_{2H} \frac{\hat{\Lambda}_2^2}{\mathfrak{f}_H} \gg |m_L| \simeq q_L \frac{\hat{\Lambda}_1^2}{\mathfrak{f}_L}$$

$$V(\phi_L) \simeq \hat{\Lambda}_1^4 \left[1 - \cos \left(\frac{\phi_L}{f_L} \right) \right]$$

$$f_L \equiv \frac{\mathfrak{f}_L}{q_L} = \frac{N_1 \mathfrak{f}_1 \mathfrak{f}_2}{|\det q|} \sqrt{q_{21}^2/\mathfrak{f}_1^2 + q_{22}^2/\mathfrak{f}_2^2}$$

More freedom: DE sector does not fix mass of heavy axion  DM

Axionic DE embedding II

assuming $m_H > H_{eq} = 10^{-27} eV$

[Marsh.'15]

$$\Omega_H = \frac{1}{6} (9\Omega_r)^{3/4} \sqrt{\frac{m_H}{10^{-60} M_p}} \left(\frac{\delta\phi_H}{M_p} \right)^2$$

isocurvature constraint:

$$\mathcal{A}_I = \left(\frac{\Omega_H}{\Omega_m} \right)^2 \left(\frac{H_{\inf}}{\pi \delta\phi_H} \right)^2 < 0.038 \mathcal{A}_s$$

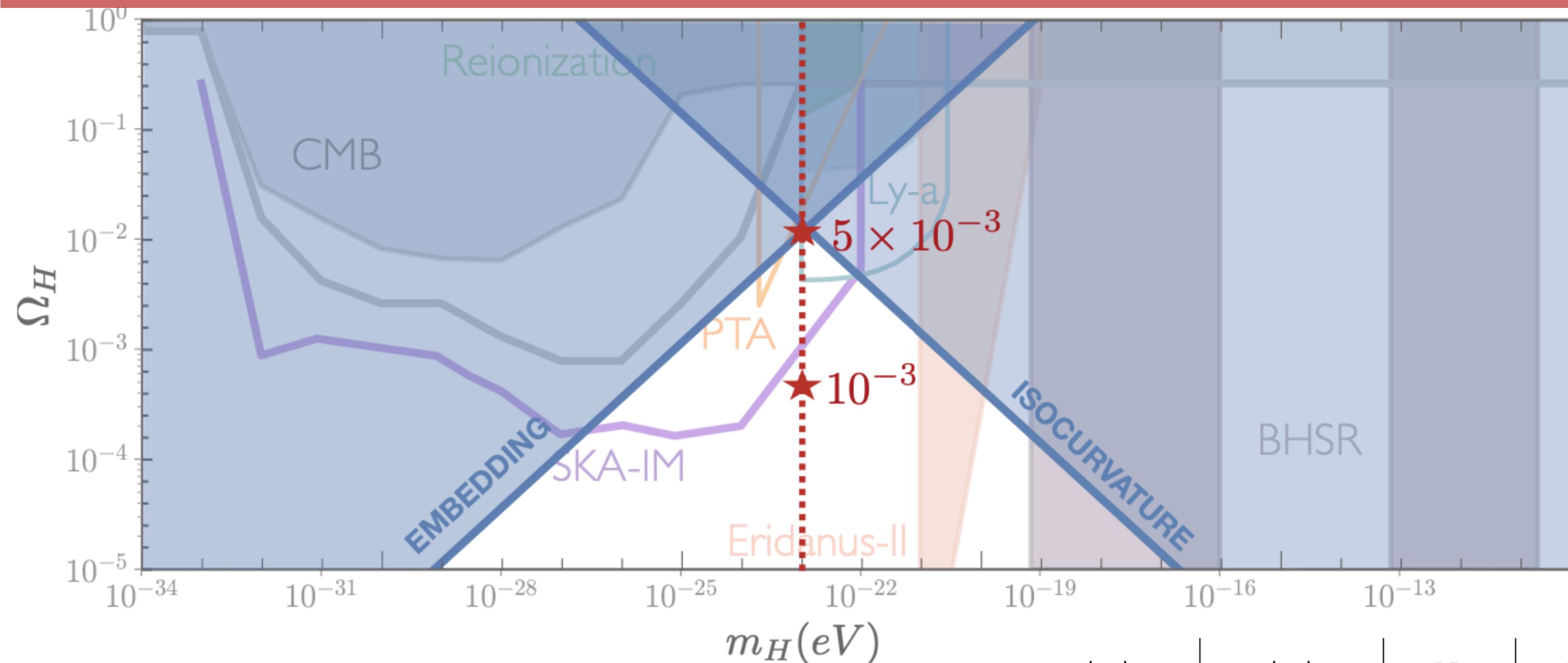
$$\frac{\Omega_H}{\Omega_m} \lesssim 3395 \sqrt{\frac{10^{-60} M_p}{m_H}} \left(\frac{0.0005}{\epsilon} \right)$$

embedding constraint:

$$f_H < 0.005 M_P$$

$$\frac{\Omega_H}{\Omega_m} \lesssim 6.95 \times 10^{-7} \sqrt{\frac{m_H}{10^{-60} M_p}}$$

Axionic DE embedding II



$$m_H \sim 10^{-23} \text{ eV}$$

$$\frac{\Omega_H}{\Omega_M} < 0.045$$

$\langle \tau_1 \rangle$	$\langle \tau_2 \rangle$	N_1	N_2	q_{ij}	f_L	f_H	Ω_H^{\max}/Ω_m
106.13	97.07	23	42	$\begin{pmatrix} 2 & 8 \\ 3 & 13 \end{pmatrix}$	0.17	0.005	0.045
95.08	102.56	20	16	$\begin{pmatrix} -19 & 26 \\ -13 & 18 \end{pmatrix}$	0.15	0.001	0.002
94.13	103.07	11	24	$\begin{pmatrix} -122 & 116 \\ -231 & 220 \end{pmatrix}$	0.11	0.0001	2×10^{-5}