



# Back to the origins of braneantibrane inflation

Gonzalo Villa de la Viña Based on 2409.xxxxx [hep-th] With M. Cicoli, C. Hughes, A. R. Kamal, F. Marino, F. Quevedo and M. Ramos-Hamud 04/09/2024, COST Action CA21106 2nd general meeting, Istanbul, Turkey

## Inflation and the UV

• Naturalness (small inflaton mass):

## Inflation and the UV

- Naturalness (small inflaton mass):
  - Shift symmetries
  - Scaling symmetries
  - SUSY

## Inflation and the UV

- Naturalness (small inflaton mass):
  - Shift symmetries
  - Scaling symmetries
  - SUSY

Higher-dimension operators:

 $\mathcal{O}_6 = V(\varphi) \left(\frac{\varphi}{\Lambda}\right)^2 \to \eta \sim \left(\frac{M_p}{\Lambda}\right)$ 

## Inflation and the UV

- Naturalness (small inflaton mass):
  - Shift symmetries
  - Scaling symmetries
  - SUSY
- Higher-dimension operators:

$$\mathcal{O}_6 = V(\varphi) \left(\frac{\varphi}{\Lambda}\right)^2 \to \eta \sim \left(\frac{M_p}{\Lambda}\right)^2$$

Inflation requires understanding of the UV!

#### Goals

#### A consistent description of brane-antibrane inflation.

Dvali, Tye'99 Burgess, Majumdar, Nolte, Quevedo, Rajesh, Zhang'01 Dvali, Shafi, Solganik'01 Kachru, Kallosh, Linde, Maldacena, McAllister, Trivedi'03

#### Goals

A consistent description of brane-antibrane inflation.
<u>Issues tackled in the paper:</u>

- Eta problem
- EFT consistency
- Experimental data
- Late-time vacuum

Today: eta problem

Dvali, Tye'99 Burgess, Majumdar, Nolte, Quevedo, Rajesh, Zhang'01 Dvali, Shafi, Solganik'01 KKLMMT: Kachru, Kallosh, Linde, Maldacena, McAllister, Trivedi'03

## **Brane-antibrane potential**

Brane positions are described by scalar fields

### Brane-antibrane potential

- Brane positions are described by scalar fields
- The potential perceived by a brane position  $\varphi$  caused by an antibrane is of the form:

$$V(\mathcal{V},\varphi) = C_0 \left(1 - \frac{D_0}{\varphi^4}\right)$$

#### **Brane-antibrane potential**

- Brane positions are described by scalar fields
- The potential perceived by a brane position  $\varphi$  caused by an antibrane is of the form:

$$V(\mathcal{V},\varphi) = C_0 \left(1 - \frac{D_0}{\varphi^4}\right)$$

When the antibrane sits in a warped throat:

 $C_0 \sim D_0 \sim rac{e^{-8\pi K/g_s M}}{\mathcal{V}^{4/3}} \qquad K, M \in \mathbb{Z}$  Naturally flat with O(1) inputs!

#### Brane-antibrane inflation: pros and cons

#### Brane-antibrane inflation: pros and cons

• Pros

- Naturally stringy ingredients.
- Order one tuning renders very flat potentials
- Favoured by observations.
- Plethora of predictions and stringy features after inflation.

Kofman, Yi'05 Frey, Mazumdar, Myers'05

#### Brane-antibrane inflation: pros and cons

#### Pros

- Naturally stringy ingredients.
- Order one tuning renders very flat potentials
- Favoured by observations.
- Plethora of predictions and stringy features after inflation.

Kofman, Yi'05 Frey, Mazumdar, Myers'05

#### Cons

- Antibrane breaks SUSY.
- Need to stabilize the WISP by excellence in string compactifications: the volume modulus.
  Eta problem.

#### WISPS in inflation: the volume

• Life is difficult because generically there is at least one further scalar: the volume modulus.

### WISPS in inflation: the volume

- Life is difficult because generically there is at least one further scalar: the volume modulus.
- Its potential is generically steep.
- So inflation can never occur far from a minimum of the potential in the volume direction.
- Need to find this minimum (moduli stabilization).

Kachru, Kallosh, Linde, Trivedi'03 Balasubramanian, Berglund, Conlon, Quevedo'05

- The volume receives a correction due to brane backreaction:  $\mathcal{V}^{2/3}=\rho-\bar{\varphi}\varphi$ 

• In the original proposal only ho is stabilized.

KKLMMT'03

- The volume receives a correction due to brane backreaction:  $\mathcal{V}^{2/3}=\rho-\bar{\varphi}\varphi$ 

KKLMMT'03

• In the original proposal only ho is stabilized. Recall:

$$V(\mathcal{V}, \varphi) \sim \frac{1}{\mathcal{V}^{4/3}} \left( 1 - \frac{D_0}{\varphi^4} \right)$$

- The volume receives a correction due to brane backreaction:  $\mathcal{V}^{2/3}=\rho-\bar{\varphi}\varphi$
- In the original proposal only  $ho\,$  is stabilized. This implies: KKLMMT'03

$$V(\mathcal{V}, \varphi) \sim \frac{1}{(\langle \rho \rangle - \bar{\varphi} \varphi)^2} \left( 1 - \frac{D_0}{\varphi^4} \right)$$

- The volume receives a correction due to brane backreaction:  $\mathcal{V}^{2/3}=\rho-\bar{\varphi}\varphi$
- In the original proposal only  $\rho$  is stabilized. This implies: KKLMMT'03  $V(\mathcal{V},\varphi) \sim \frac{1}{(\langle \rho \rangle - \bar{\varphi}\varphi)^2} \left(1 - \frac{D_0}{\varphi^4}\right) \sim \frac{1}{\langle \rho \rangle^2} \left(1 - \frac{D_0}{\varphi^4}\right) (1 + 2\bar{\varphi}\varphi)$

Which is a large mass term for the inflaton (eta problem)

- The volume receives a correction due to brane backreaction:  $\mathcal{V}^{2/3}=\rho-\bar{\varphi}\varphi$
- In the original proposal only  $\rho$  is stabilized. This implies: KKLMMT'03  $V(\mathcal{V},\varphi) \sim \frac{1}{(\langle \rho \rangle - \bar{\varphi}\varphi)^2} \left(1 - \frac{D_0}{\varphi^4}\right) \sim \frac{1}{\langle \rho \rangle^2} \left(1 - \frac{D_0}{\varphi^4}\right) (1 + 2\bar{\varphi}\varphi)$
- Which is a large mass term for the inflaton (eta problem)
- Solution: use corrections\* that stabilize the whole quantity  $\mathcal{V}$ .

\*For aficionados: perturbative corrections to the Kahler potential

#### **Nonlinearly realized SUSY**

• The antibrane breaks SUSY but its effects can be accommodated into an effective SUSY description.

Komargodski, Seiberg'09 Kallosh, Quevedo, Uranga'15

#### **Nonlinearly realized SUSY**

- The antibrane breaks SUSY but its effects can be accommodated into an effective SUSY description.
- A description in terms of SUSY variables allows for a systematic classification of possible corrections.

Komargodski, Seiberg'09 Kallosh, Quevedo, Uranga'15

#### **Nonlinearly realized SUSY**

- The antibrane breaks SUSY but its effects can be accommodated into an effective SUSY description.
- A description in terms of SUSY variables allows for a systematic classification of possible corrections.
- We discover a new perfect square structure in the scalar potential which stabilises the volume.

Komargodski, Seiberg'09 Kallosh, Quevedo, Uranga'15

 The (perturbative) scalar potential in the framework of nonlinearly realized SUSY is given by 5 ingredients:

 The (perturbative) scalar potential in the framework of nonlinearly realized SUSY is given by 5 ingredients:

Three functions of the volume

 $f(\mathcal{V}^{2/3}), g(\mathcal{V}^{2/3}), h(\mathcal{V}^{2/3})$ 

 The (perturbative) scalar potential in the framework of nonlinearly realized SUSY is given by 5 ingredients:

Three functions of the volume

 $f(\mathcal{V}^{2/3}), g(\mathcal{V}^{2/3}), h(\mathcal{V}^{2/3}) \xrightarrow{\text{Tree level}} f(x) = x, \ g(x) = 0, \ h(x) = 1$ 

 The (perturbative) scalar potential in the framework of nonlinearly realized SUSY is given by 5 ingredients:

Three functions of the volume

 $f(\mathcal{V}^{2/3}), g(\mathcal{V}^{2/3}), h(\mathcal{V}^{2/3}) \xrightarrow{\text{Tree level}} f(x) = x, \ g(x) = 0, \ h(x) = 1$ 

• A (for current purposes) tunable constant  $\,W_0$ 

• A function of  $\, arphi$  , but not  ${\mathcal V}$  which contains the Coulomb potential  $\, W_X$ 

 The most general scalar perturbative potential in the framework of nonlinearly realized SUSY reads\*:

$$V = \frac{1}{U} \left[ \left( f'W_X - 3g'W_0 \right)^2 - f'' \left( fW_X^2 - 6gW_X W_0 - 9hW_0^2 \right) \right]$$

• With U a function of f,g,h.

\*Neglecting (small) inflaton F-terms

 The most general scalar perturbative potential in the framework of nonlinearly realized SUSY reads\*:

$$V = \frac{1}{U} \left[ \left( f'W_X - 3g'W_0 \right)^2 - f'' \left( fW_X^2 - 6gW_X W_0 - 9hW_0^2 \right) \right]$$

- With U a function of f,g,h.
- Tree level reproduces standard uplift.

 Wx=g=0 reproduce standard breaking of no-scale structure (f" is BBHL)

\*Neglecting (small) inflaton F-terms

 $V = \frac{1}{U} \left[ \left( f'W_X - 3g'W_0 \right)^2 - f'' \left( fW_X^2 - 6gW_X W_0 - 9hW_0^2 \right) \right]$ 

• If the second term is small, perfect square that stabilises  $\langle \mathcal{V} \rangle$  at a Minkowski minimum, uplifted to dS by the second term.

 $V = \frac{1}{U} \left[ \left( f'W_X - 3g'W_0 \right)^2 - f'' \left( fW_X^2 - 6gW_X W_0 - 9hW_0^2 \right) \right]$ 

If the second term is small, perfect square that stabilises ⟨𝒱⟩ at a Minkowski minimum, uplifted to dS by the second term.
For g' ~ 1/𝒱<sup>a</sup> :

 $\langle \mathcal{V} \rangle = \left(\frac{W_0}{W_X}\right)^{1/a}$ 

 $V = \frac{1}{U} \left[ \left( f'W_X - 3g'W_0 \right)^2 - f'' \left( fW_X^2 - 6gW_X W_0 - 9hW_0^2 \right) \right]$ 

If the second term is small, perfect square that stabilises ⟨𝒱⟩ at a Minkowski minimum, uplifted to dS by the second term.
For g' ~ 1/𝒱<sup>a</sup> :

 $\langle \mathcal{V} \rangle = \left(\frac{W_0}{W_X}\right)^{1/a}$ 

Controlled SUGRA for small Wx!!

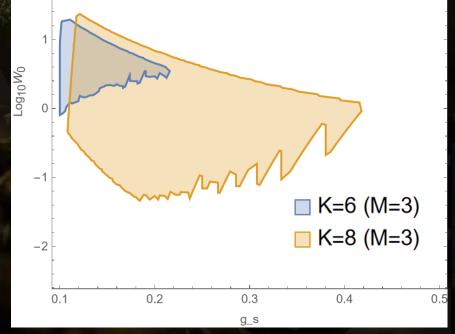
#### **Brane-antibrane inflation reloaded!**

- If g'>>f" (typically BBHL), then the volume is integrated leaving a Coulomb potential\*.
- Brane-antibrane inflation can now occur.

\*In the paper we also propose a stabilization mechanism in the opposite case

#### **Brane-antibrane inflation reloaded!**

- If g'>>f" (typically BBHL), then the volume is integrated leaving a Coulomb potential\*.
- Brane-antibrane inflation can now occur.
- After EFT and experimental constraints, for g~log[V]:



\*In the paper we also propose a stabilization mechanism in the opposite case

### **Conclusions and future directions**

- Perturbative corrections to the Kahler potential allow for braneantibrane inflation.
- Nonlinearly realized SUSY uncovers a perfect-square structure in the scalar potential.
- In this talk, g plays a prominent role. It is a kinetic mixing goldstinovolume. EFT arguments suggest it's there, the question is at what power in 1/V.
- We find a parameter space of microscopic data where braneantibrane inflation can take place.

## **Backup: Kahler and superpotential**

- Most general\* Kahler and superpotentials with constrained superfields X^2=0:
- $K = -3\log[f(\mathcal{V}^{2/3}) + g(\mathcal{V}^{2/3})(X + \bar{X}) + h(\mathcal{V}^{2/3})X\bar{X}]$  $W = W_0 + XW_X$

\*Contributions outside of the log can always be absorbed up to a higher order correction