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On (De Sitter) Vacua in Heterotic String Theory

Nicole Righi

based on work with

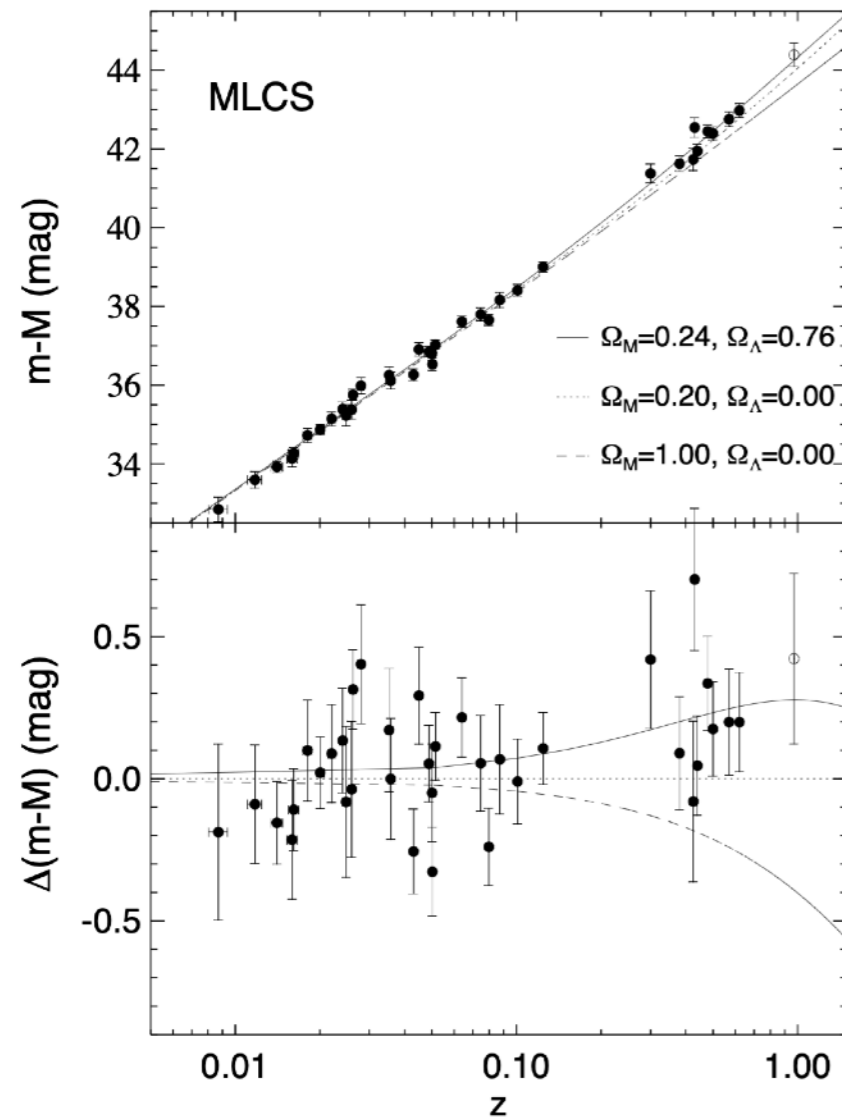
Rafael Álvarez-García, Christian Kneissl,

Jacob M. Leedom, Alexander Westphal

2nd COST Meeting, ISU Istanbul 2024

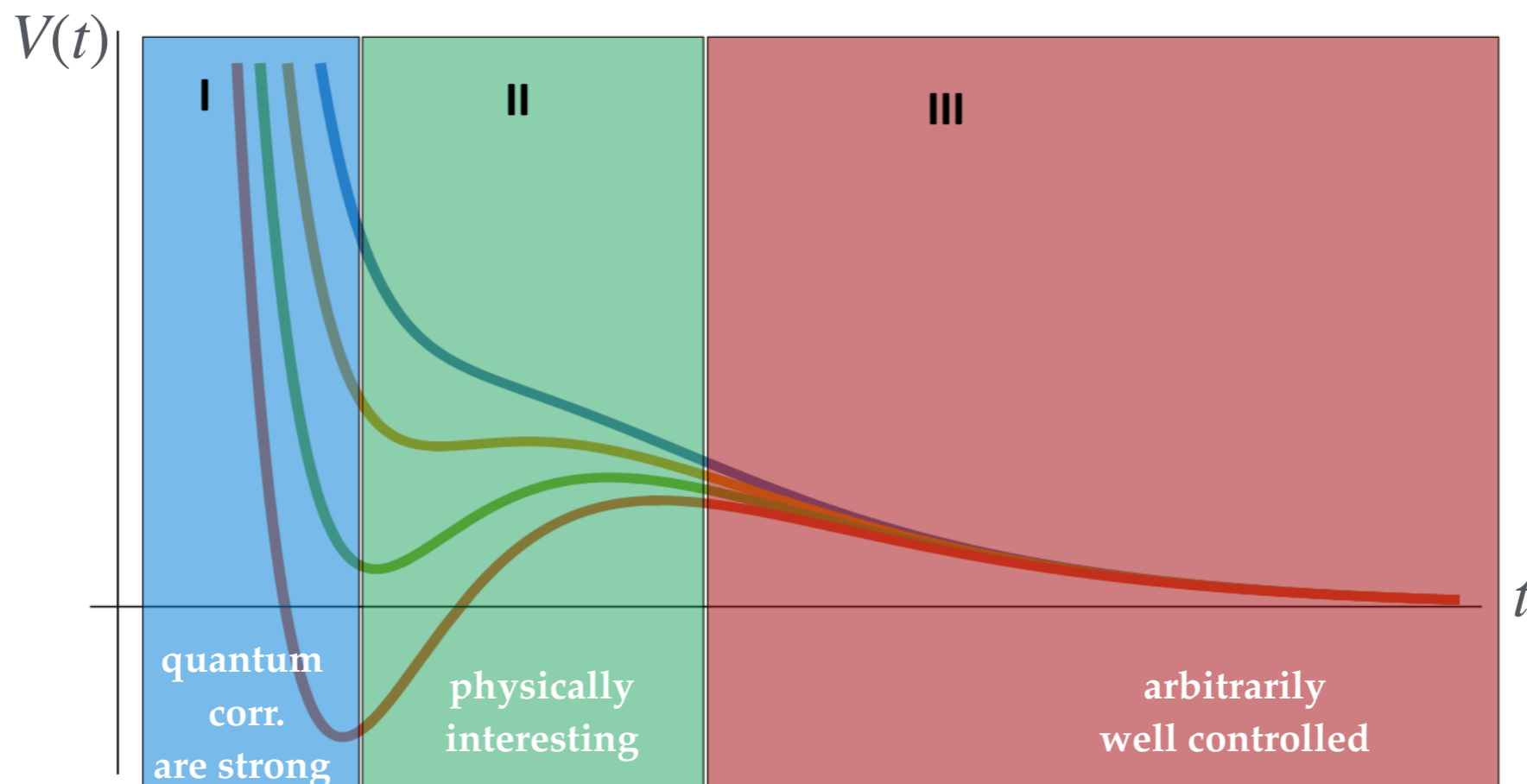
The positive Λ

[Riess et al, '98]



⇒ String theory should reproduce this

Dine-Sieberg Problem



[McAllister, Quevedo, *Moduli Stabilization in String Theory*]

e.g. in IIB on CY orientifolds:

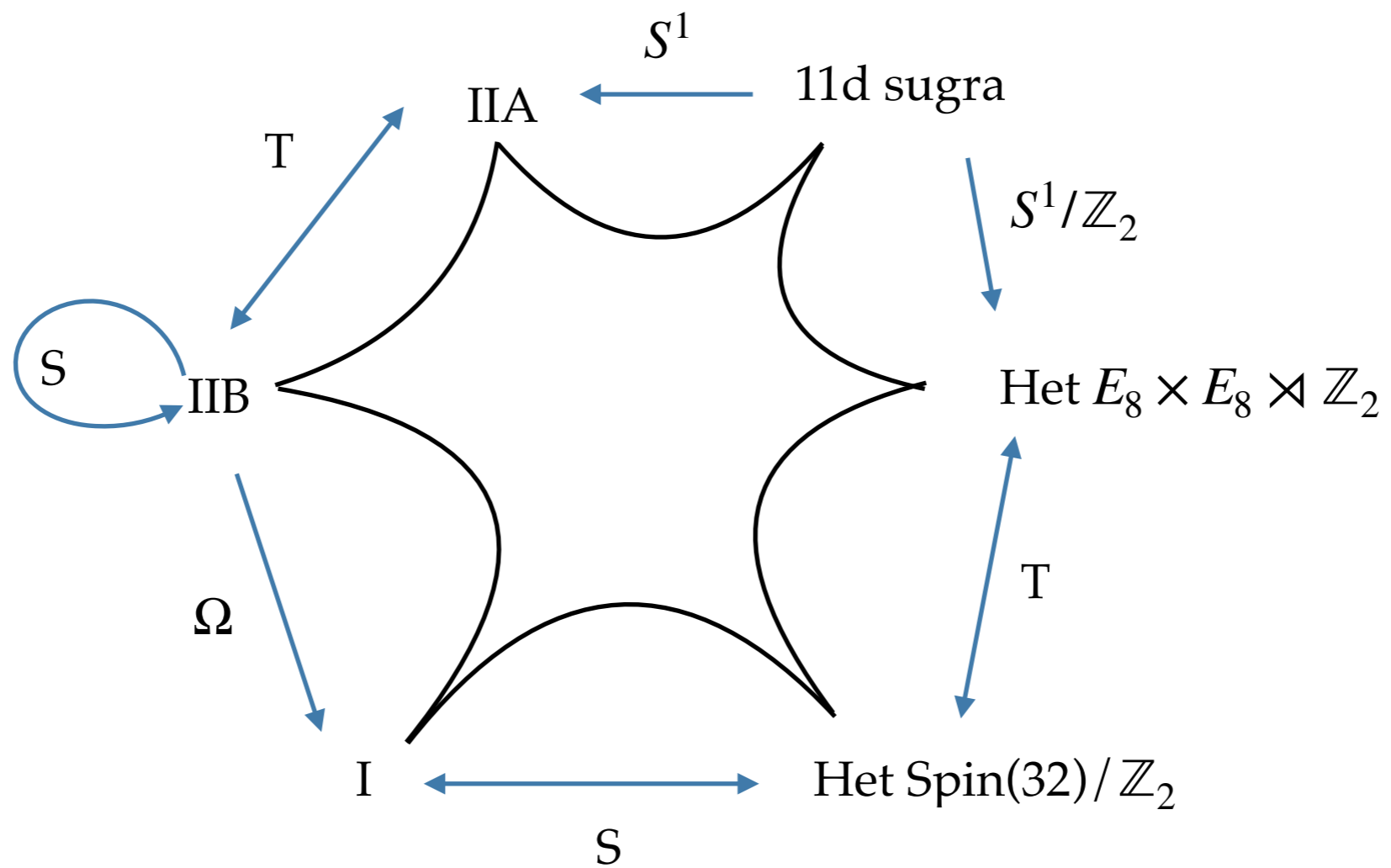
$$|W_0| \sim \delta W \gg \delta K \quad [\text{Kachru, Kallosh, Linde, Trivedi '03}] \quad \text{KKLT}$$

$$|W_0| \gg \delta K \sim \delta W \quad [\text{Balasubramanian, Berglund, Conlon, Quevedo '05}] \quad \text{LVS}$$

$$|W_0| \gg \delta K \gg \delta W \quad [\text{Berg, Haack, Kors '05}] \quad \text{most plagued by the problem}$$

[Antoniadis, Chen, Leontaris '18]

Motivation for Heterotic

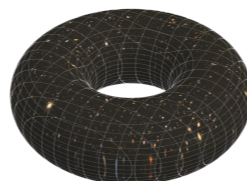


Motivation for Heterotic

Heterotic string is a remarkable playground:

- 4d $N = 2$ heterotic - type II duality

- presence of modular symmetries



- control of the $N = 1$ EFT from higher susy subsectors \longrightarrow Can use volumes $1 \lesssim t \lesssim 10$

- outstanding work on **vacua** and **SM** \longrightarrow No-go thms for de Sitter vacua

\Rightarrow Lessons for the Swampland program

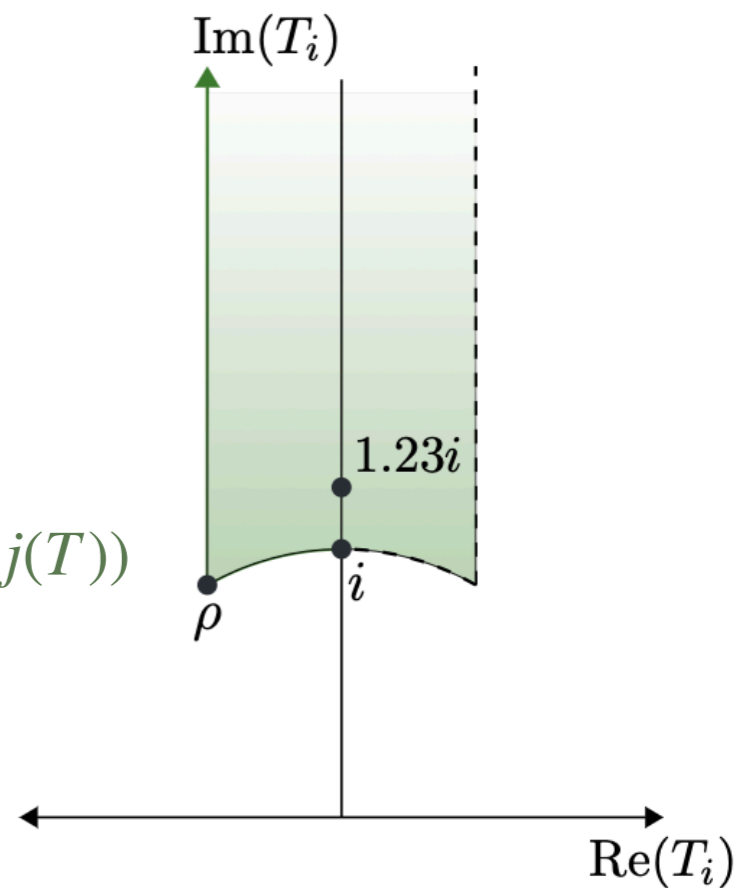
The Potential

Kahler potential $K = k(S, \bar{S}) - 3 \ln(-i(T - \bar{T}))$

Superpotential $W = \frac{\Omega(S)H(T)}{\eta^6(T)} \quad \Omega(S) \sim e^{-S}$
 $H(T) = \left(\frac{E_4(T)}{\eta^8(T)}\right)^n \left(\frac{E_6(T)}{\eta^{12}(T)}\right)^m \mathcal{P}(j(T))$

Scalar potential

$$V = e^{k(S, \bar{S})} Z(T, \bar{T}) |\Omega(S)|^2 \left[(A(S, \bar{S}) - 3) |H(T)|^2 + \widehat{V}(T, \bar{T}) \right]$$



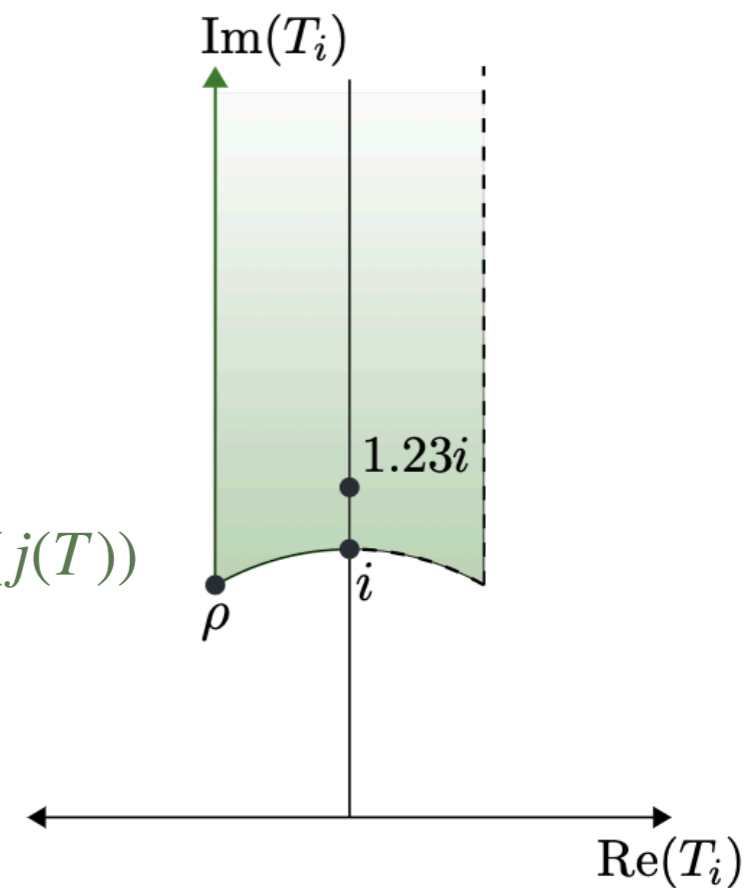
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Scalar potential

$$V = e^{k(S, \bar{S})} \underbrace{Z(T, \bar{T})}_{> 0} |\Omega(S)|^2 \left[\underbrace{(A(S, \bar{S}) - 3)}_{\sim |F_S|^2 \geq 0} |H(T)|^2 + \underbrace{\widehat{V}(T, \bar{T})}_{> 0} \right]$$



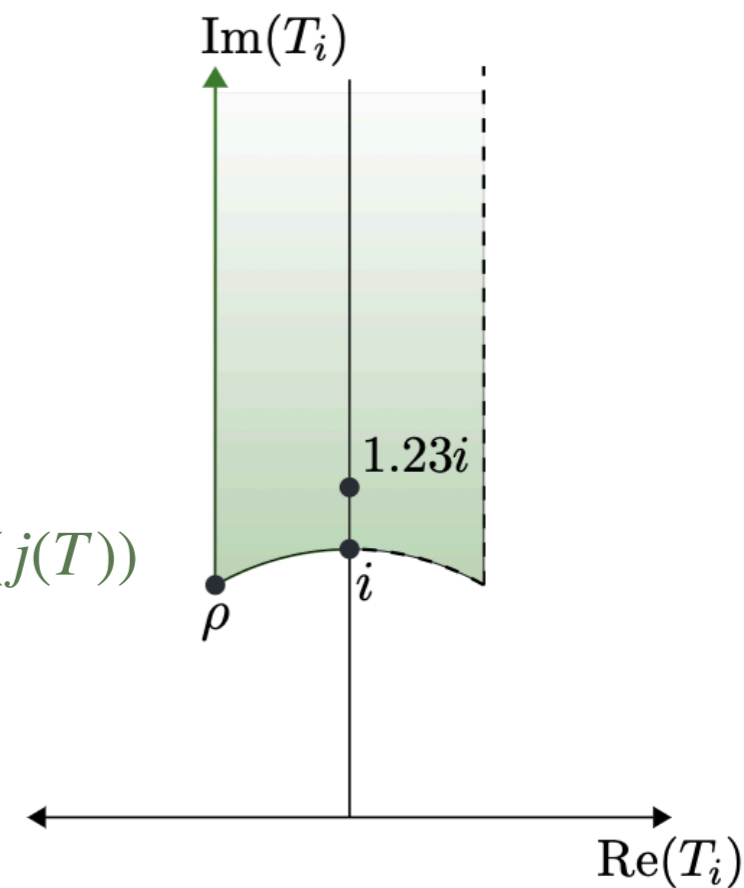
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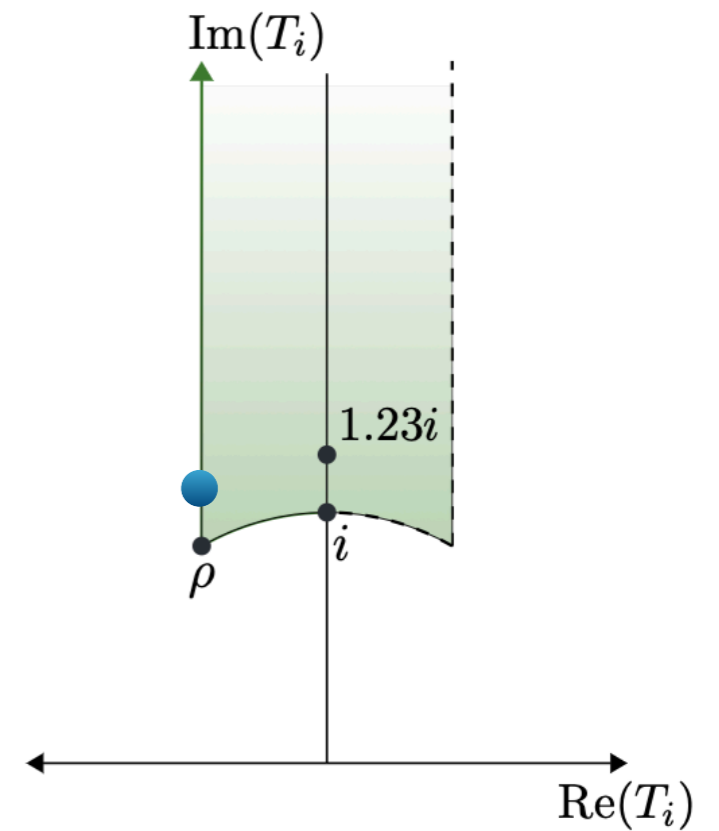
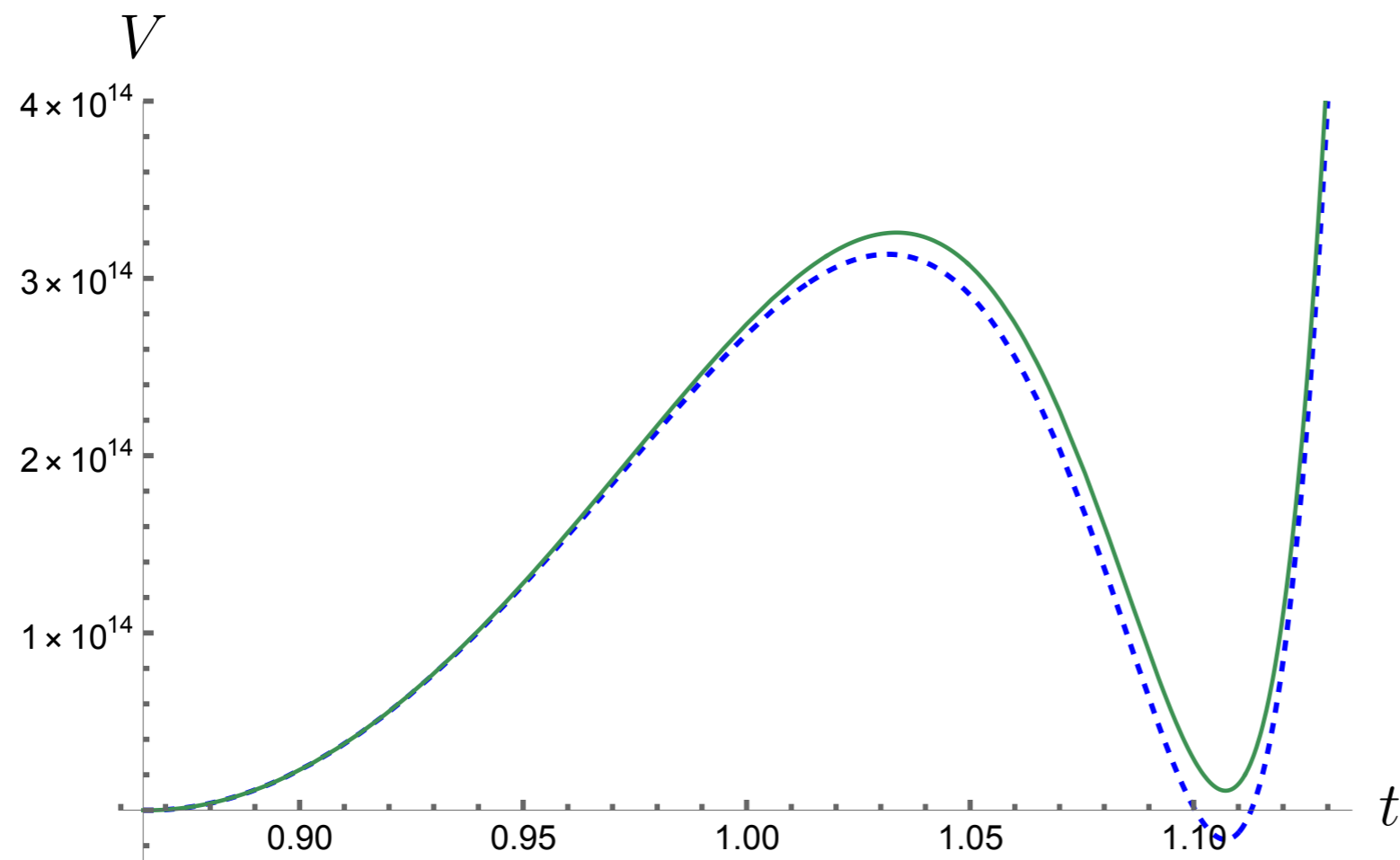
$$V = e^{k(S, \bar{S})} \underbrace{Z(T, \bar{T})}_{> 0} |\Omega(S)|^2 \left[\underbrace{(A(S, \bar{S}) - 3)}_{\sim |F_S|^2 \geq 0} |H(T)|^2 + \underbrace{\widehat{V}(T, \bar{T})}_{> 0} \right]$$



1) if $K = -\log(S + \bar{S}) \Rightarrow V \leq 0$

2) if $K = -\log(S + \bar{S}) + \delta K_{np} \Rightarrow V$ can be positive!

De Sitter Vacua



$\text{---} A(S, \bar{S}) = 0$
 $\text{---} A(S, \bar{S}) > 3 - \frac{\hat{V}(T, \bar{T})}{H(T)}$

De Sitter Vacua: Sanity Checks

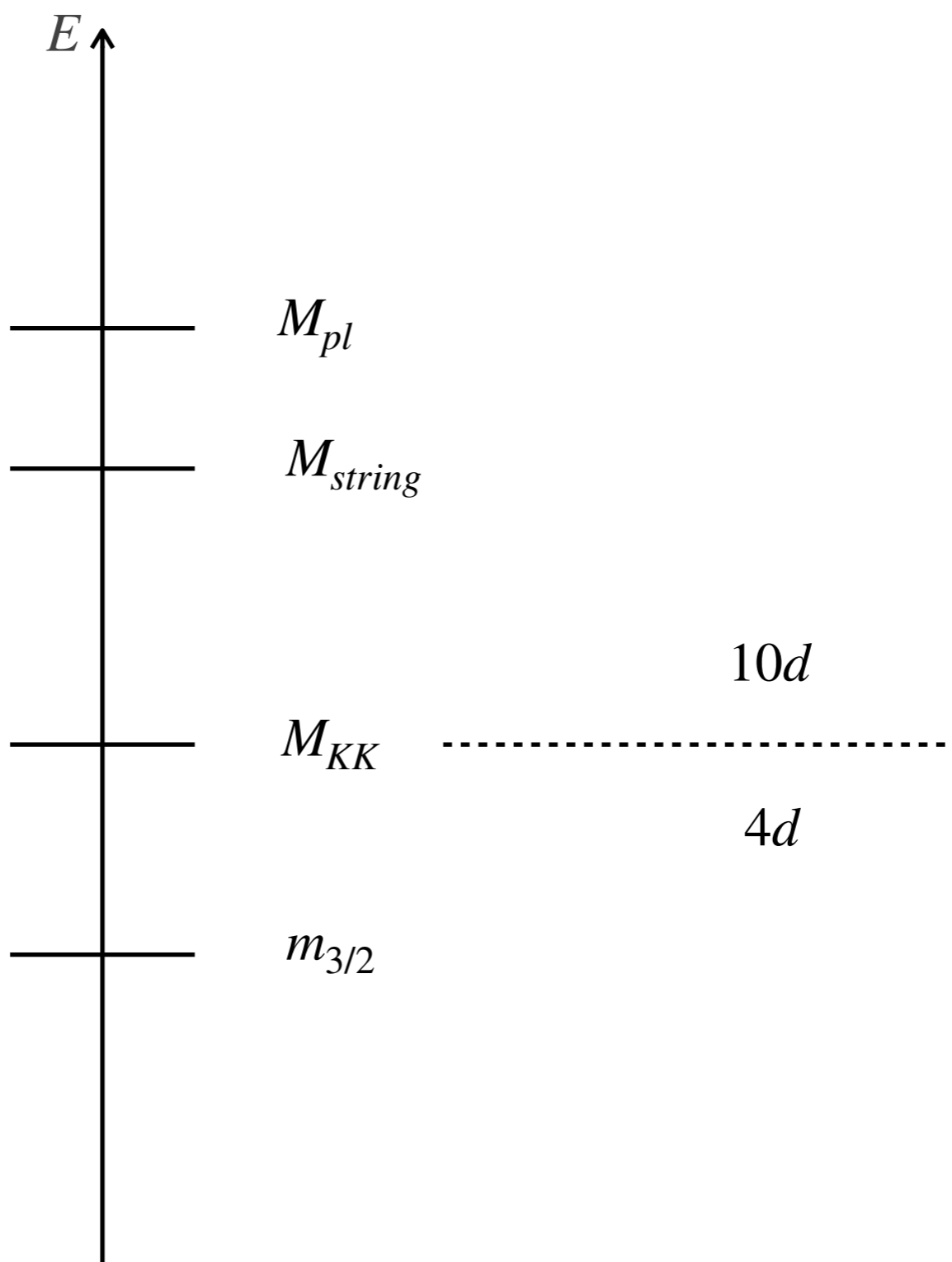
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- 2) can a δK_{np} of the right magnitude exist?

De Sitter Vacua: Sanity Checks

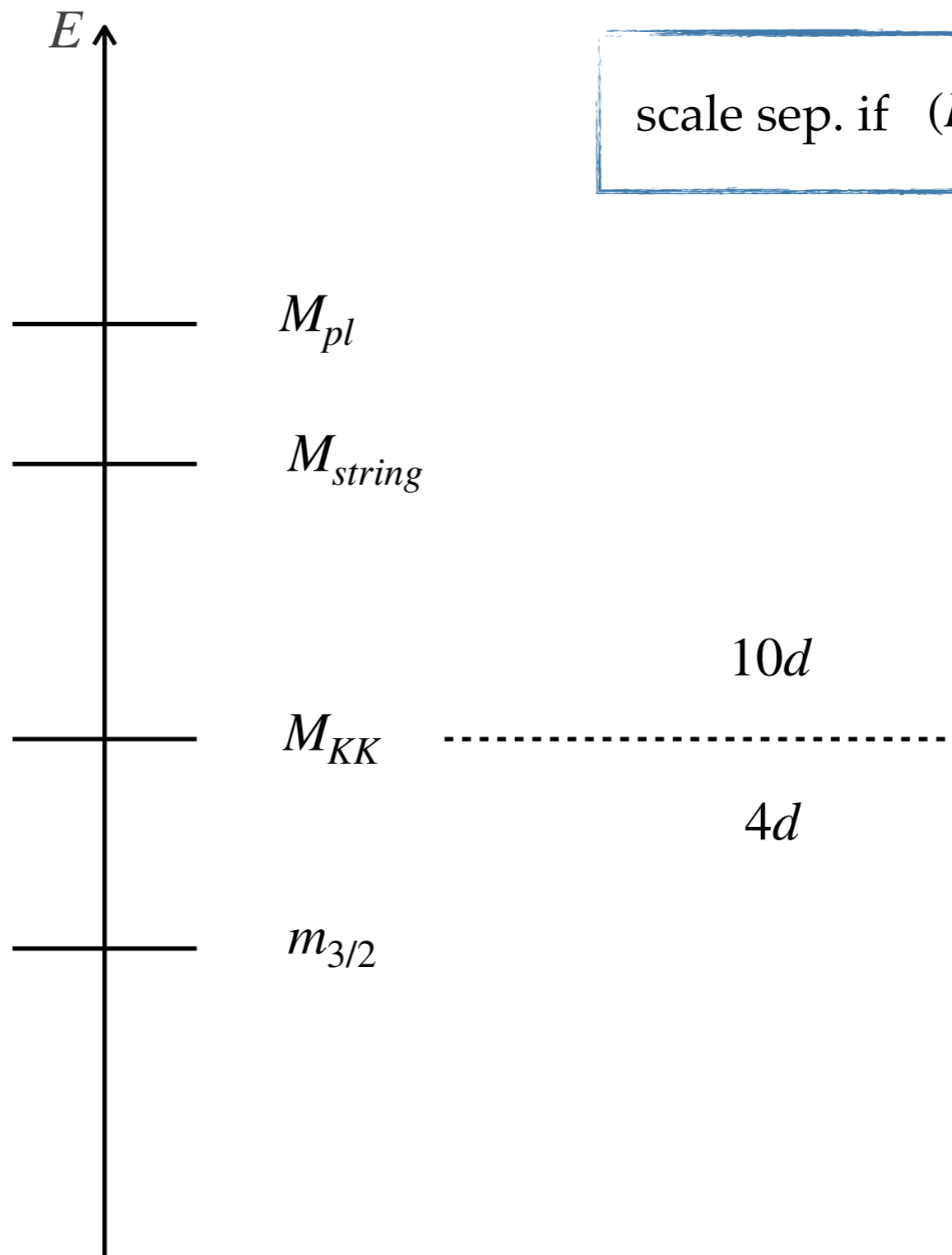
1) can we really ignore KK modes? i.e. do we have scale separation? [NR *WiP*]

2) can a δK_{np} of the right magnitude exist? [Álvarez-García, Kneissl, Leedom, NR, '24]

Scale Separation



Scale Separation



scale sep. if $(L_{AdS}M_{KK})^{-1} \ll 1$

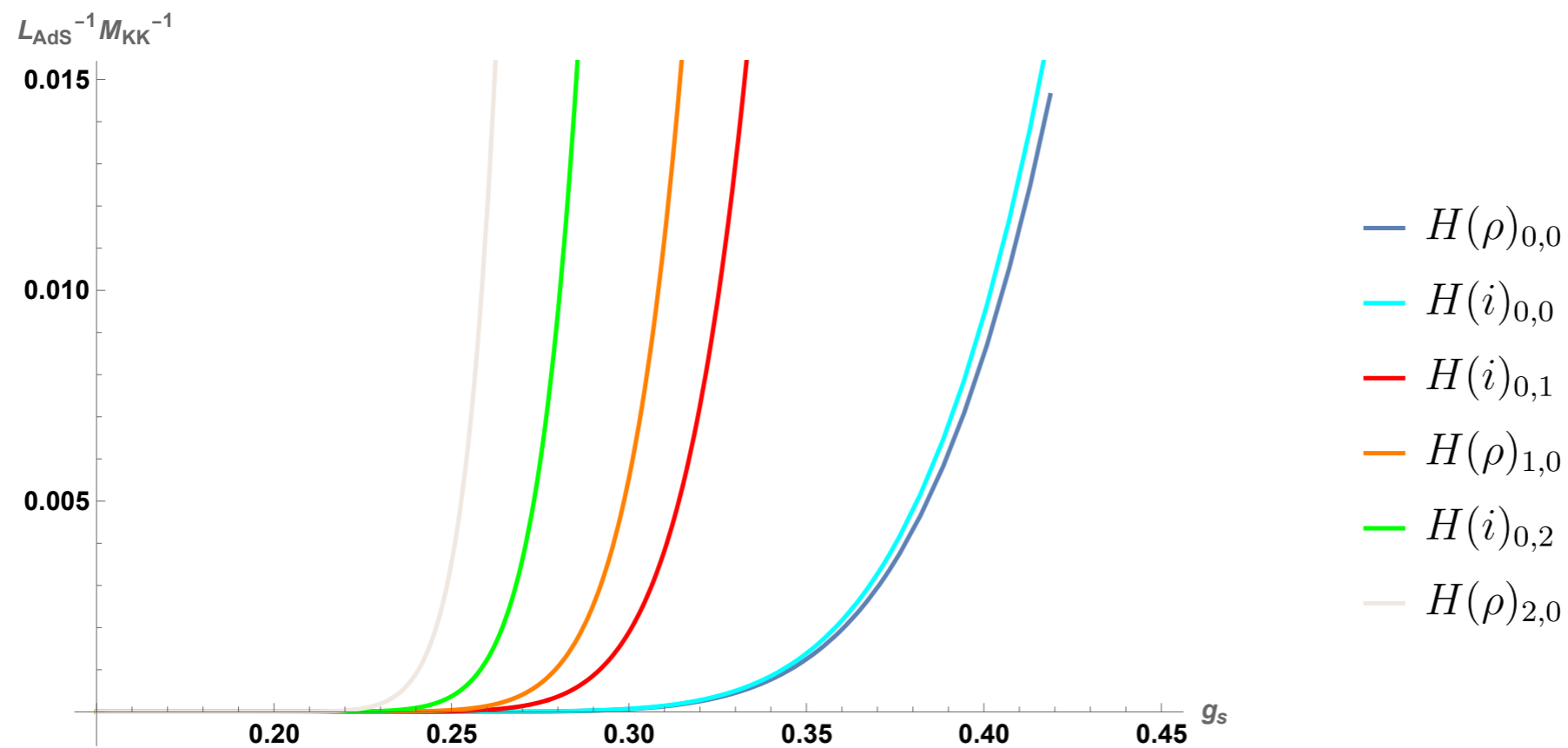
Scale Separation

[NR WiP]

$$\frac{L_{AdS}^{-1}}{M_{KK}} \simeq \frac{|H(T)| |\Omega(S)|}{t |\eta(T)^6|}$$

scale sep. if $(L_{AdS} M_{KK})^{-1} \ll 1$

susy-preserving vev

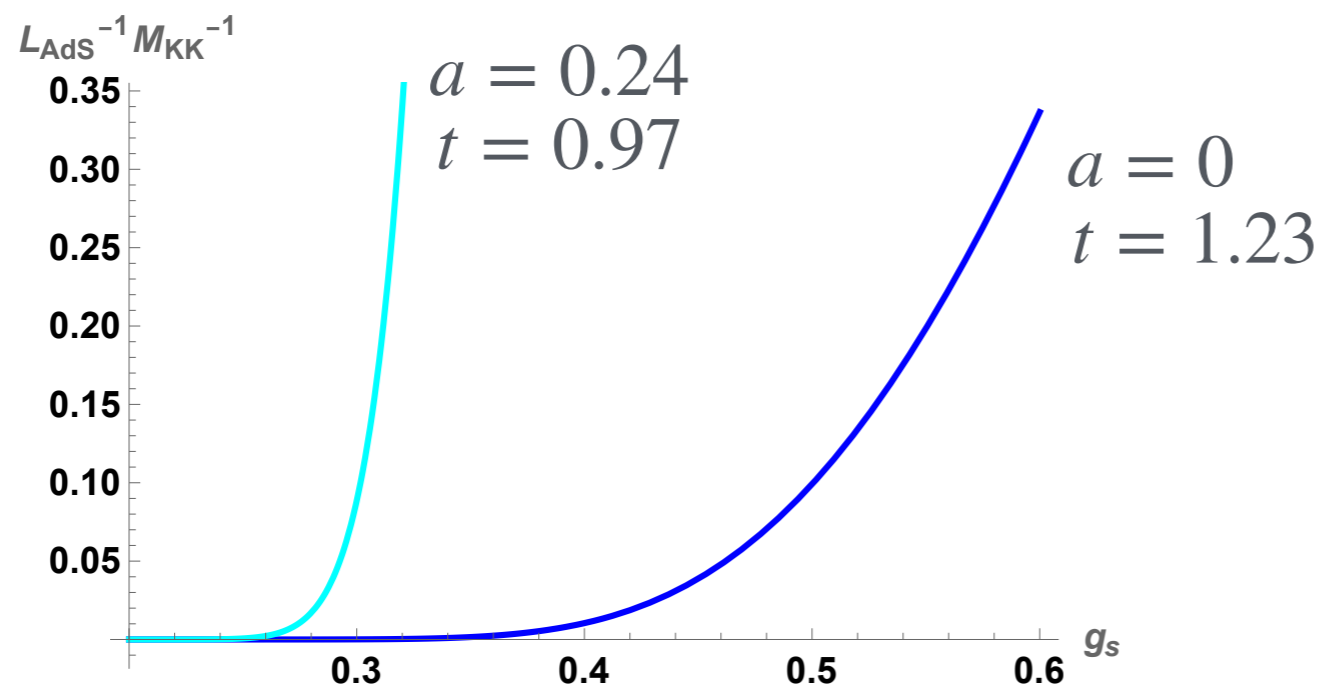


Scale Separation

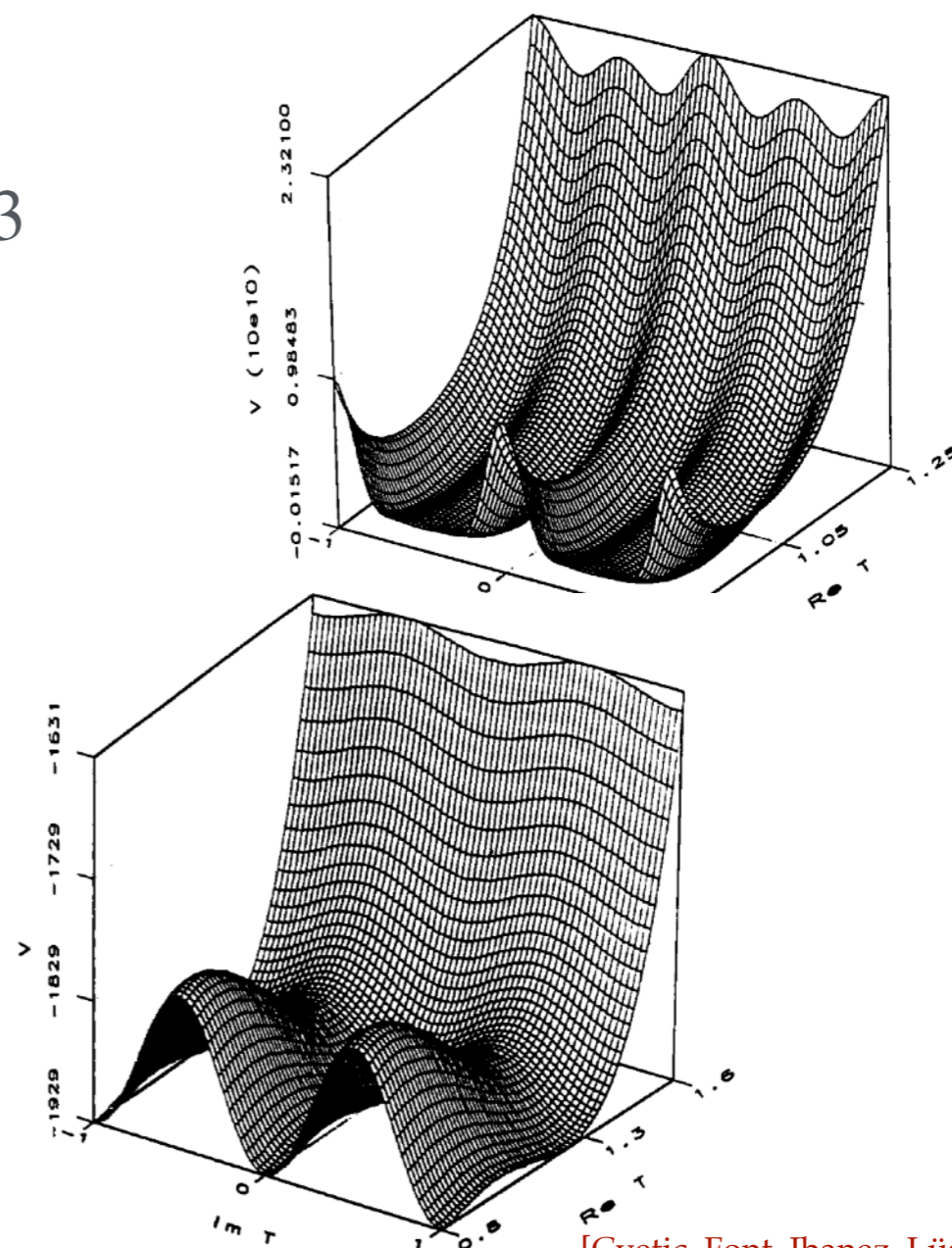
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susy-breaking vev



Heterotic Open Strings

$$K = -\ln(S + \bar{S}) + \delta K_{np}$$

All closed string theories have nonperturbative contributions of strength $\sim e^{-1/g_s}$

[Shenker '90]

Quite odd in heterotic — no D-branes!

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but there are “D-branes”!
the e^{-1/g_s} should come from a 10d instanton
(at least for $\text{Spin}(32)/\mathbb{Z}_2$)

[Hull, '97 & '98]
[Polchinski, '05]

[Álvarez-García, Kneissl, Leedom, NR, '24]

What have we found + What has to be done

- Constructed new de Sitter vacua in heterotic
- Found the critical ingredient: $\delta K_{np} \sim e^{-1/g_s}$
- Started checking its consistency via scale separation & 10d origin
 - scale separated dS vacua?
 - 10d instanton: need of an action
 - What about 10d $E_8 \times E_8$?
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Thank you

Heterotic Open Strings: Arguments

- Hull's orientifold introduces a D1-string
- Inflow of fermion d.o.f from spacetime to worldsheet
- Disk diagrams from the D1 with endpoints on the instanton giving e^{-1/g_s} contributions



Shenker's argument, explained in HO by a D(-1)-brane

- Confirmed by 4-graviton 10d scattering amplitudes
- $\pi_9(\text{SemiSpin}(32)) = \mathbb{Z}_2 + \text{Derrick's theorem} \quad + \quad \Omega_{10}^{\text{Spin}}(B\text{SemiSpin}(32)) = 10\mathbb{Z}_2$