



Axion search goes deep underground: θ-dependence of α-decay half-lives

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based on the work w/ Carlo Broggini, Giuseppe Di Carlo and Luca Di Luzio arxiv: 2404.18993, published on PLB

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Axion potentially addresses the Strong CP problem

$$\mathcal{L}_{\theta} = \frac{g^2 \theta}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$|\theta| \lesssim 10^{-10}$$

>
$$\theta$$
 promoted to a dynamical field: $\theta \rightarrow \frac{a}{f_a}$

> QCD potential relaxed *dynamically* to zero:



Oscillating axion DM

- An axion contribution to DM density is unavoidable so it is a natural DM candidate
- > In a misalignment mechanism the DM axion field induces a time varying θ -term

Preskill, Wise, Wilczek '83, Abbott, Sikivie '83, Dine, Fischler '83



$$\theta \simeq \sqrt{\frac{2\rho_{DM}}{m_a^2 f_a^2}} \cos(\omega t + \vec{p} \cdot \vec{x} + \phi)$$

α -radioactivity and axion DM

> Search for time modulation of α -radioactivity from axion DM

Broggini, Di Carlo, Di Luzio, CT arxiv:2006.12321

Lee, Meißner, Olive, Shifman, Vonk arxiv:2006.12321

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Non vanishing θ -term impacts nuclear physics

$$M_{\pi}^{2}(\theta) = M_{\pi}^{2} \cos \frac{\theta}{2} \sqrt{1 + \varepsilon^{2} \tan^{2} \frac{\theta}{2}}$$

For its impact anthropic context, see Ubaldi arxiv:0811.1599

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Time modulation of radiative decays

For a recent study of time modulation signals on β -decays, see Zhang, Houston, Li arxiv:2303.09865

Theory of α -decay

Gamow theory of α -decay

$$T_{1/2} = \frac{\ln 2}{\nu_0} \exp(K),$$

$$K = \frac{2}{\hbar} \int_{r_1}^{r_2} dr \, \sqrt{2\mu [V_{\rm tot}(r) - Q_{\alpha}]}$$



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$$K = Z_{\alpha} Z_{d} \alpha_{\text{QED}} \left(\frac{8\mu}{Q_{\alpha}} \right)^{1/2} F \left(\frac{Q_{\alpha} R_{\text{well}}}{Z_{\alpha} Z_{d} \alpha_{\text{QED}}} \right),$$

with

$$F(x) = \arccos \sqrt{x} - \sqrt{x} \sqrt{(1-x)} \approx \frac{\pi}{2} - 2\sqrt{x} + \dots,$$



Signal estimate

Halftimes are highly sensible to the Q-value, i.e. the energy released in the decay

Question: how the θ -term impacts the Q?

 $Q_{\alpha} = BE(A - 4, Z - 2) + BE(4, 2) - BE(A, Z),$

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> The θ -term changes the size of the scalar (attractive) and vector (repulsive) nuclear interaction that contributes to the BEs

$$H = G_S(\bar{N}N)(\bar{N}N) + G_V(\bar{N}\gamma_{\mu}N)(\bar{N}\gamma^{\mu}N),$$

$$\eta_S = \frac{G_S(\theta)}{G_S(\theta=0)}, \quad \eta_V = \frac{G_V(\theta)}{G_V(\theta=0)}.$$

Damour, Donoghue arxiv:0712.2968

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$$\eta_{S} = \frac{G_{S}(\theta)}{G_{S}(\theta=0)}, \quad \eta_{V} = \frac{G_{V}(\theta)}{G_{V}(\theta=0)}.$$

$$Q_{\alpha}(\theta) = Q_{\alpha}(\theta=0) - 97 \text{ MeV}(\eta_{S}(\theta)-1) \times ((A-4)^{2/3} + 4^{2/3} - A^{2/3}).$$

$$Damour, Donoghue arxiv:0712.2968$$

$$\eta_{S}(\theta) = 1.4 - 0.4 \frac{M_{\pi}^{2}(\theta)}{M_{\pi}^{2}}$$





$$\frac{\mathring{T}_{1/2}(0)}{T_{1/2}(0)} \approx \mathring{K} \approx \frac{\partial K}{\partial Q_{\alpha}} \mathring{Q}_{\alpha} \approx 8.45 \ (Z-2) \left[4^{2/3} - \frac{8}{3A^{1/3}} \right] \left(\frac{\text{MeV}}{Q_{\alpha}} \right)^{3/2}$$



Figure A.3: Contours of $\mathring{T}_{1/2}(0)/T_{1/2}(0)$ in the (Q_{α}, Z) plane for A = 241. The case of ²⁴¹Am is indicated by a white dot.

- 630

- 567

- 504

- 441

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Exp. set-up



- > A 3"x3" NaI crystal detects the gamma rays which are subproducts of the α -decays, primarily (85% of the time) at 59.5 keV
- > The crystal is closed inside a cooper-lead shield

Similar studies have already been performed, see for example: Bellotti, Broggini, Di Carlo, Laubenstein, Menegazzo arxiv:1802.09373

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A picture of our experiment, based on Americium-241, placed under Gran Sasso

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Figure 1: γ -spectrum (counts per second per keV) of the ²⁴¹Am source (upper curve) compared to the background (lower curve). The dominant contribution arises from the γ at 59.5 keV.

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Exp. prospects

Prospect limits from measurements of Americium-241 α -decay



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> Phase I, one month of data:

$$1 \,\mu s < \Delta t < 10 \text{ days}$$

 $I_{exp} = 2 \times 10^{-5} \text{ at } 2\sigma$



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> Phase I, one month of data:

 $1 \,\mu s < \Delta t < 10 \text{ days}$ $I_{exp} = 2 \times 10^{-5} \text{ at } 2\sigma$

> Phase II, three years of data:

 $1 \,\mu s < \Delta t < 1 \,yr$ $I_{exp} = 4 \times 10^{-6} \text{ at } 2\sigma$





