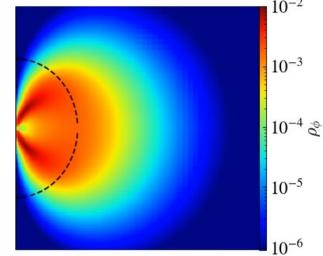
Non-linear dynamics in modified gravity



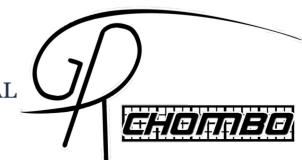
Llibert Aresté Saló



Pollica, 19th September 2024







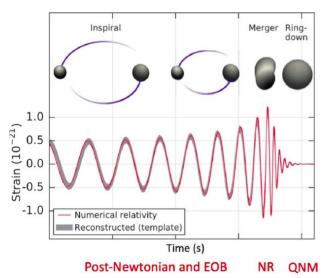
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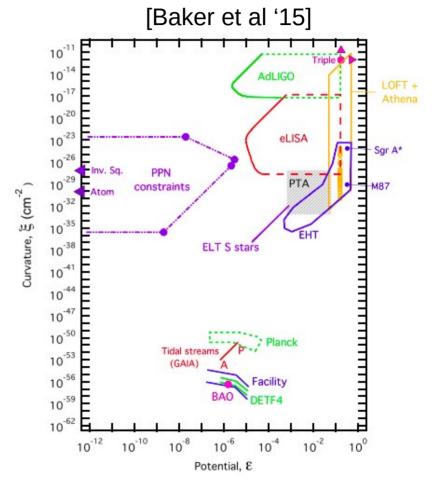
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Motivation

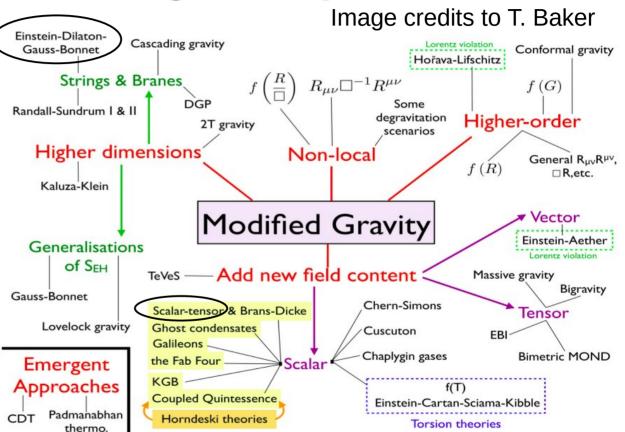
- Detection of gravitational waves → Testing of the strong field regime.
- Numerical Relativity enables us to compute those waveforms.





Modified gravity

- Multiple possibilities for modifying gravity.
- Effective field theory.



The theory

We consider the Four-Derivative Scalar Tensor Theory of Gravity:

$$S = \int d^4x \sqrt{-g} (R + X - V(\phi) + g_2(\phi) X^2 + \lambda(\phi) \mathcal{L}^{GB}),$$

$$2^{\text{nd deriv.}} \qquad 4^{\text{th deriv.}}$$

$$X = -\frac{1}{2} (\partial_\mu \phi)^2, \mathcal{L}^{GB} = R^2 - 4 R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma},$$

which yields Einstein-scalar-Gauss-Bonnet (EsGB) when $V(\phi)=g_2=0$.

- It violates the No-Hair Theorem.
- It is an EFT of gravity in the weak coupling regime:

$$\lambda(\phi), g_2(\phi) \ll L^2$$

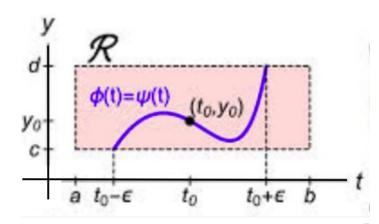
where L is the characteristic length of the system.

Well-posedness

• An initial value (or Cauchy) problem

$$\partial_t u = F(x^i, u, \partial_i u, \dots, \partial_{i_1} \dots \partial_{i_m} u, \dots)$$

$$u|_{t=0} = f(x^i)$$



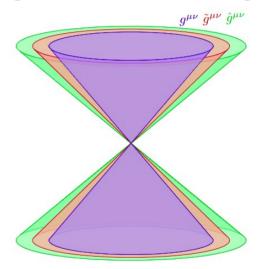
is (locally) well-posed if there exists a unique (local) solution for $t \in [0, T]$ for a given T>0 which depends smoothly on the initial data.

 Only well-posed initial value problems can lead to stable Numerical Relativity simulations.

Modified Harmonic Gauge

- Proposed in [Kovács and Reall '20].
- Well-posed formulation of weakly coupled Lovelock and Horndeski theories of gravity.
- Different propagation speeds for the unphysical modes.
- Implemented in EsGB in GHC coordinates [East and Ripley '21, Corman, Ripley and East '23, Corman and East '24].

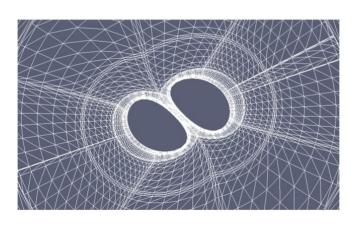
[Kovács and Reall '20]

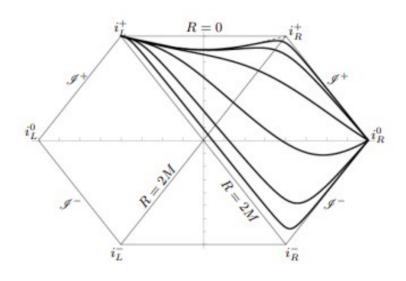


$$\begin{split} \tilde{g}^{\mu\nu} &= g^{\mu\nu} - a(x) n^{\mu} n^{\nu} \\ \hat{g}^{\mu\nu} &= g^{\mu\nu} - b(x) n^{\mu} n^{\nu} \\ 0 &< a(x) < b(x) \end{split}$$

Our formulation

- Moving puncture gauge (singularityavoiding coordinates).
- It does not need excision of the inner horizon





[Hannam et al '08]

Results

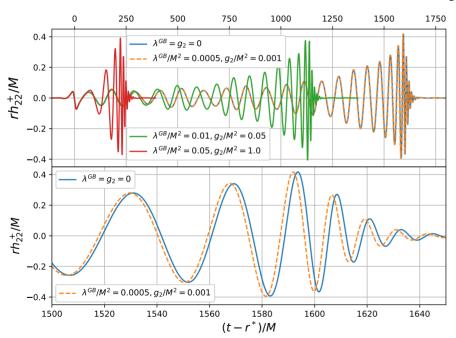
$$S = \int d^{4}x \sqrt{-g} (R + X + g_{2}(\phi) X^{2} + \lambda(\phi) \mathcal{L}^{GB})$$

$$X = -\frac{1}{2} (\partial_{\mu} \phi)^{2}, \mathcal{L}^{GB} = R^{2} - 4 R_{\mu\nu} R^{\mu\nu} + R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma}$$

- Well-posed in our modified CCZ4 formulation in the weakly coupled regime.
- Implemented in its full non-linear form in GRFolres [LAS et al '23], an extension of GRChombo, https://github.com/GRTLCollaboration/GRFolres.

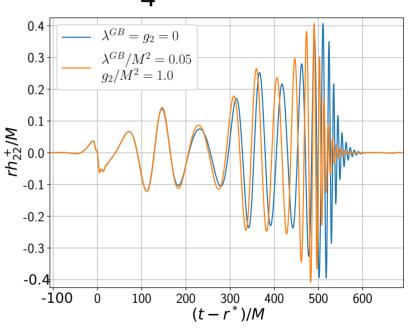
Shift-symmetric 4∂ST theory

• All Black Hole solutions are hairy.



Non-spinning Black Hole binaries [LAS, Clough and Figueras '22]

$$\lambda(\phi) = \frac{1}{4} \lambda^{GB} \phi$$
 $g_2(\phi) = g_2$

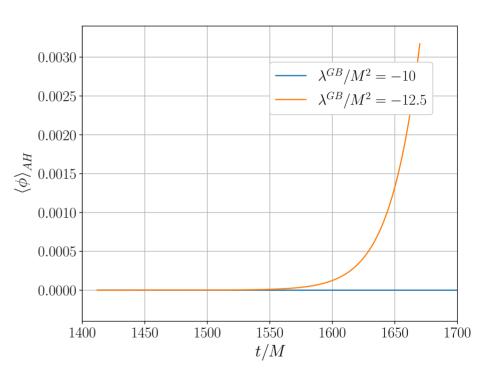


Initially spinning Black Hole binaries [LAS, Clough and Figueras '23]

Quadratic EsGB

$$\lambda(\phi) = \frac{1}{4} \lambda^{GB} \phi^2$$

- Studied in [Silva et al '21, Elley et al '22] without backreaction.
- Hairy and non-hairy Black Holes.
- Spin-induced tachyonic instability that triggers spontaneous scalarisation [Silva, Sakstein, Gualtieri, Sotiriou and Berti '18].



[LAS, Clough and Figueras '23]

Exponential quadratic EsGB

$$\lambda(\phi) = \frac{1}{4\sigma} \lambda^{GB} (1 - e^{-\sigma\phi^2}) \text{ Grav. strain } \frac{\frac{8}{2} \frac{0.2}{0.0}}{\frac{1}{2} \frac{0.2}{0.2}}$$

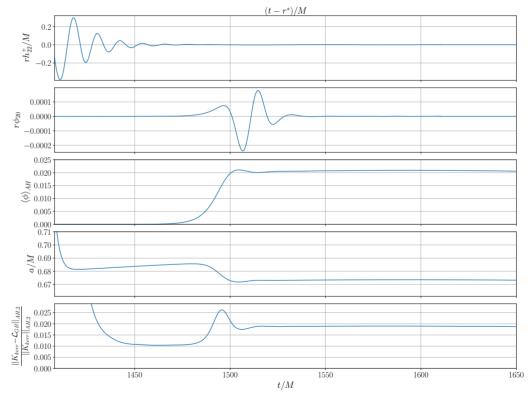
- Proposed in [Doneva et al '22].
- Spin-induced scalarisation.
- Stable hairy Black Hole Merger.

Scalar radiation

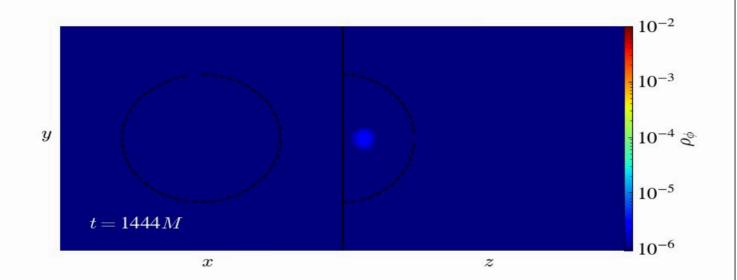
Exp. value of ϕ at the Merger's AH

Spin

Deviation from Kerr's Kretschmann

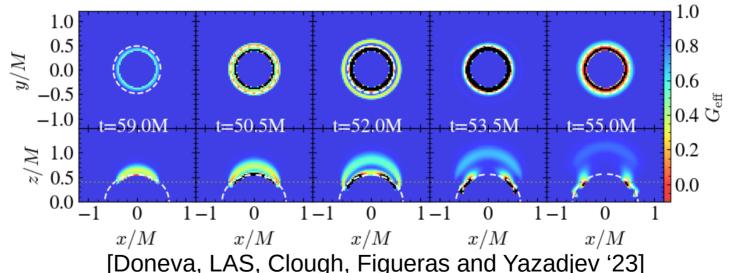


[LAS, Clough and Figueras '23]



Hyperbolicity loss

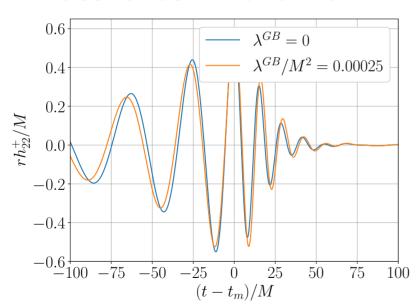
- Some of the physical modes lie on the null cone of an effective metric.
- The change of sign of its determinant determines the transition from a hyperbolic system to an elliptic system.

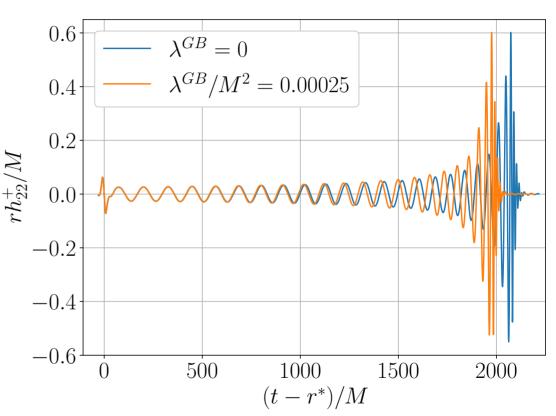


See also [Ripley and Pretorius '19]

Non-equal mass binaries

 We have been able to evolve non-equal mass binaries with mass ratios 1:2 and 1:3



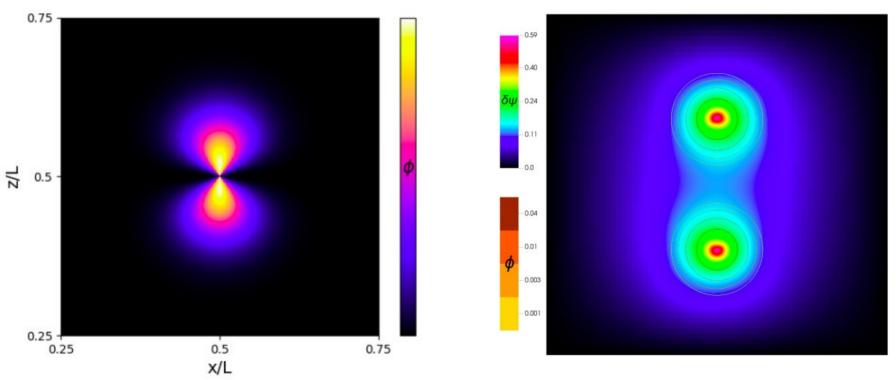


See also [Corman, Ripley and East '23]

[Doneva, LAS, Clough, Figueras and Yazadjev '24] (in preparation)

Initial conditions

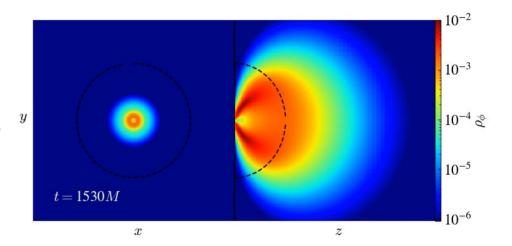
CTTK method [Aurrekoetxea, Clough and Lim '22]



[Brady, LAS, Clough, Figueras and P.S. '23]

Summary

- Well-posed formulation in singularity-avoiding coordinates of the 4∂ST theory.
- Equal and non-equal mass ^y binaries simulations in the full non-linear theory.
- Non-trivial dynamics for the scalar field.



Further work

- Make our modified gravity waveforms available to the community.
- Understand whether modified gravity effects are degenerate with environmental effects, a shift in GR parameters...
- Explore other high derivative theories (see [Held, Figueras, Kovács '24])

Further work

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THANK YOU FOR YOUR ATTENTION!