

Cosmology with binary black holes

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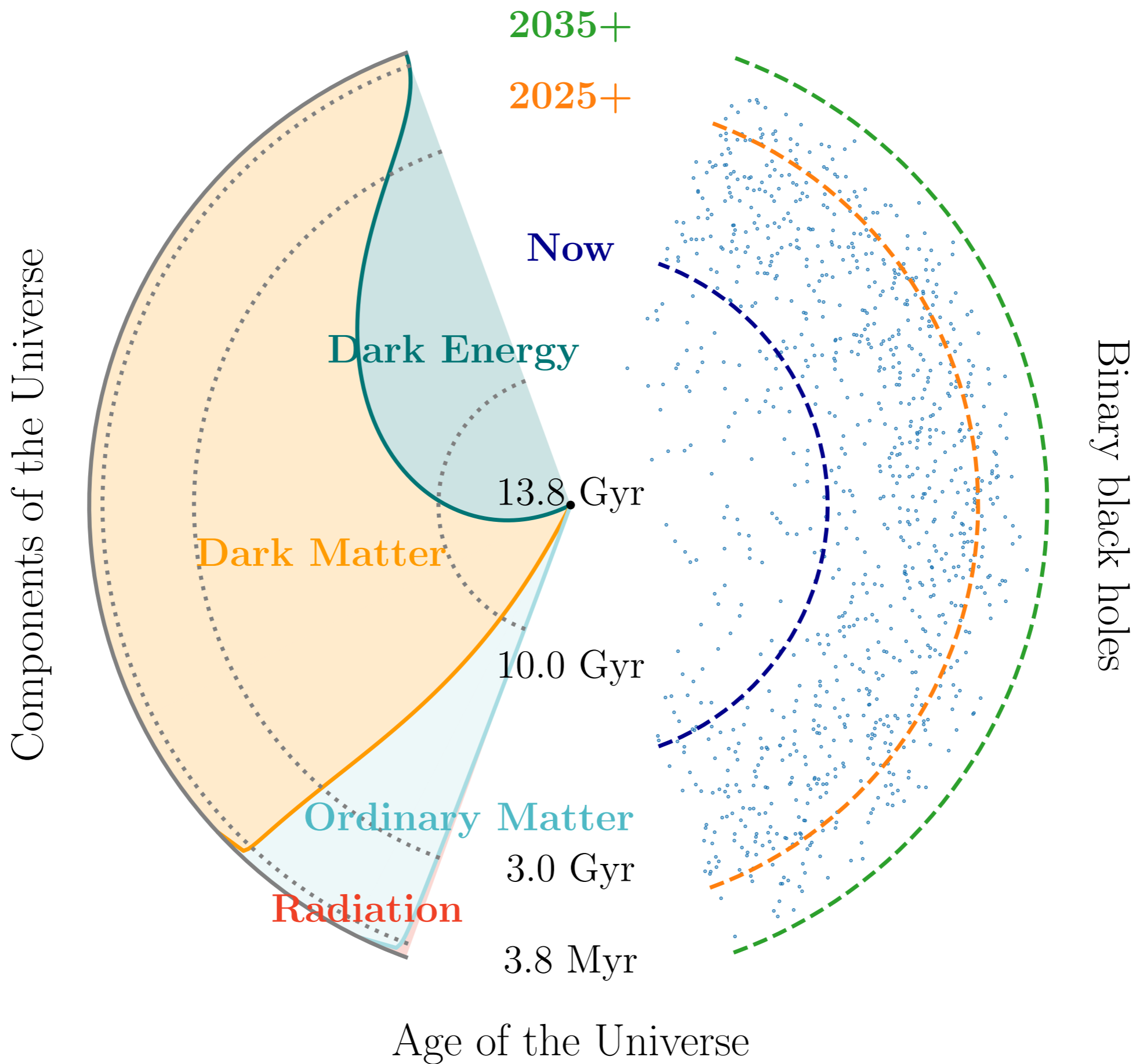
VILLUM FONDEN



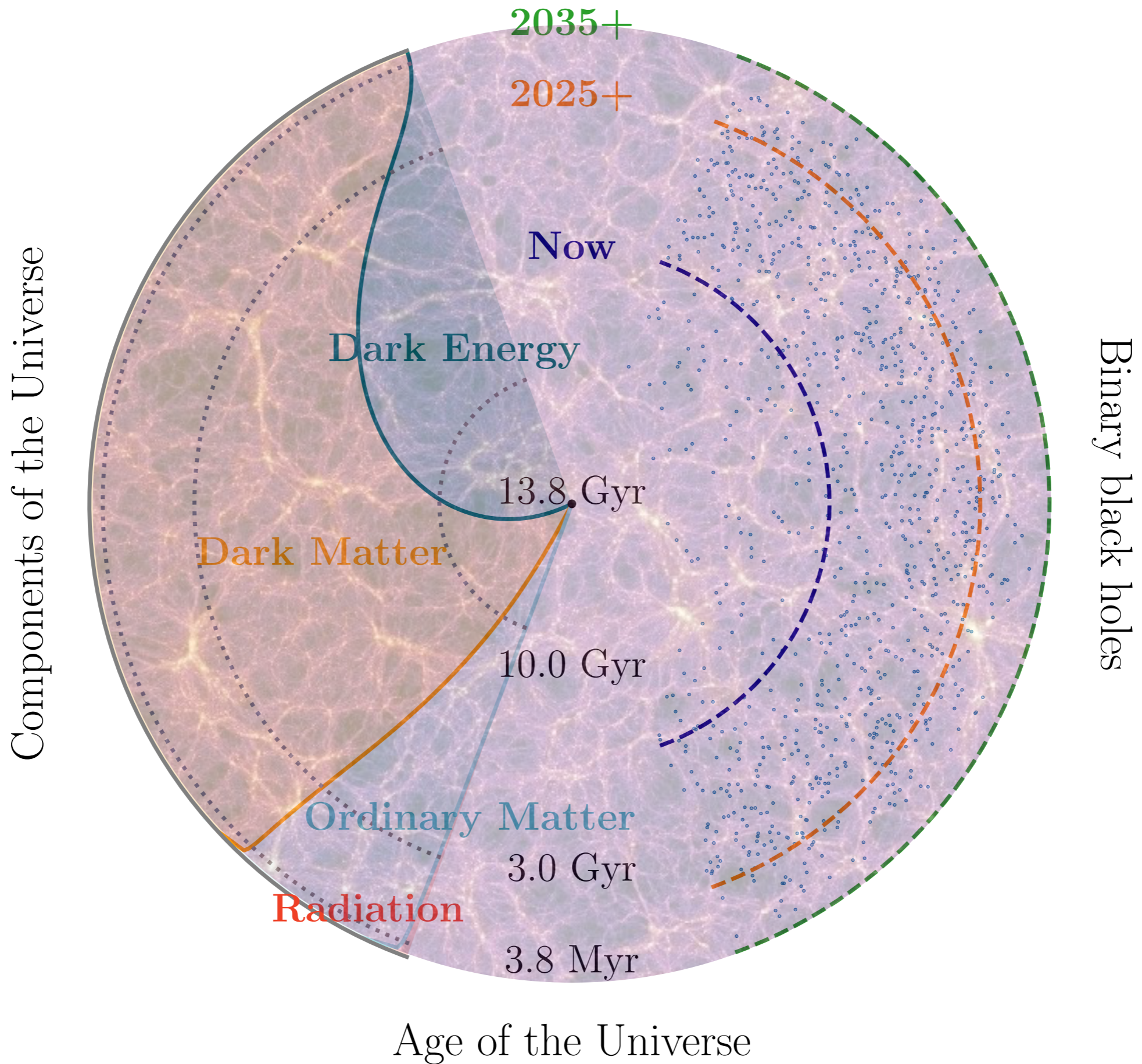
KØBENHAVNS
UNIVERSITET

[Pompei]

Gravitational Wave horizons

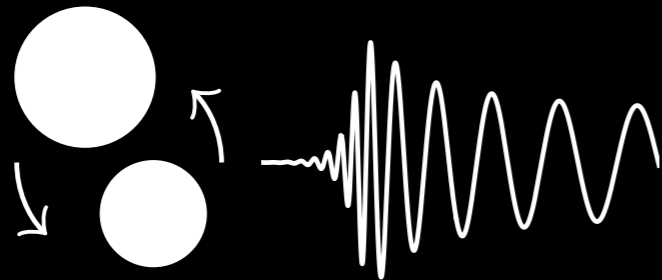


Gravitational Wave horizons



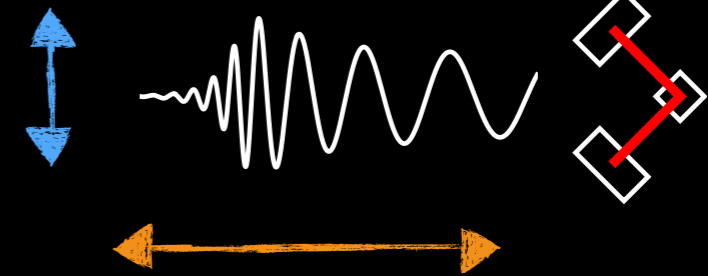
GWs are standard sirens

[well understood selection effects]



[GR predicts waveform]

[cosmo propagation]



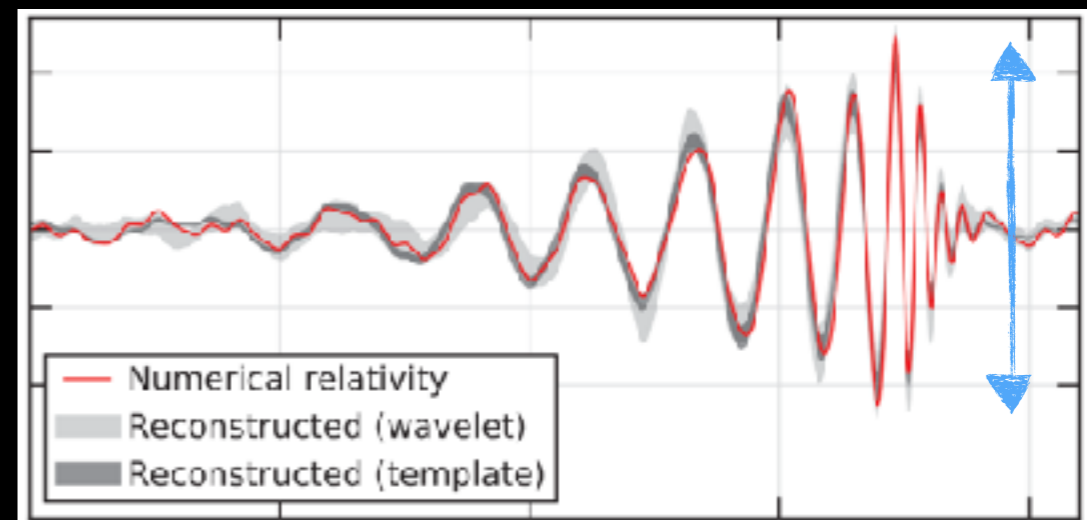
$$d_L(z)$$

[GW Hubble diagram]

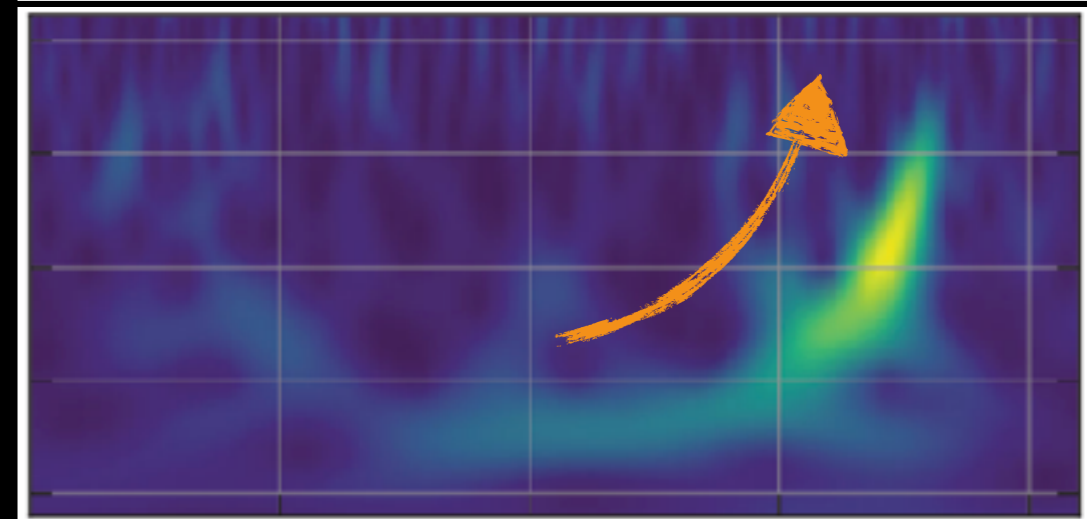
$$m_{\text{det}} = (1 + z)m$$

[Interplay with astrophysics]

strain



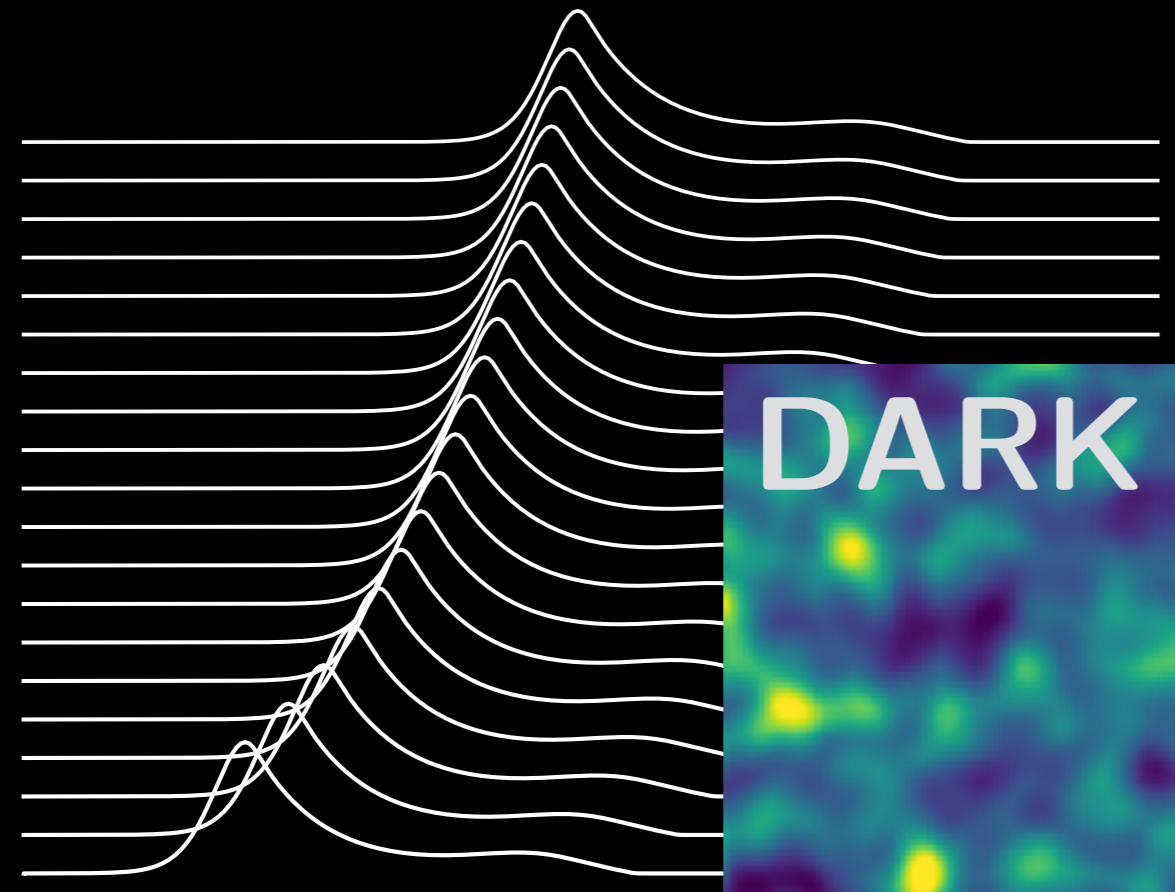
frequency



time

SPECTRAL SIRENS

Luminosity distance



log[Detector frame]



Love sirens



[Credit: D. Berry]

Cosmography with next-generation gravitational wave detectors

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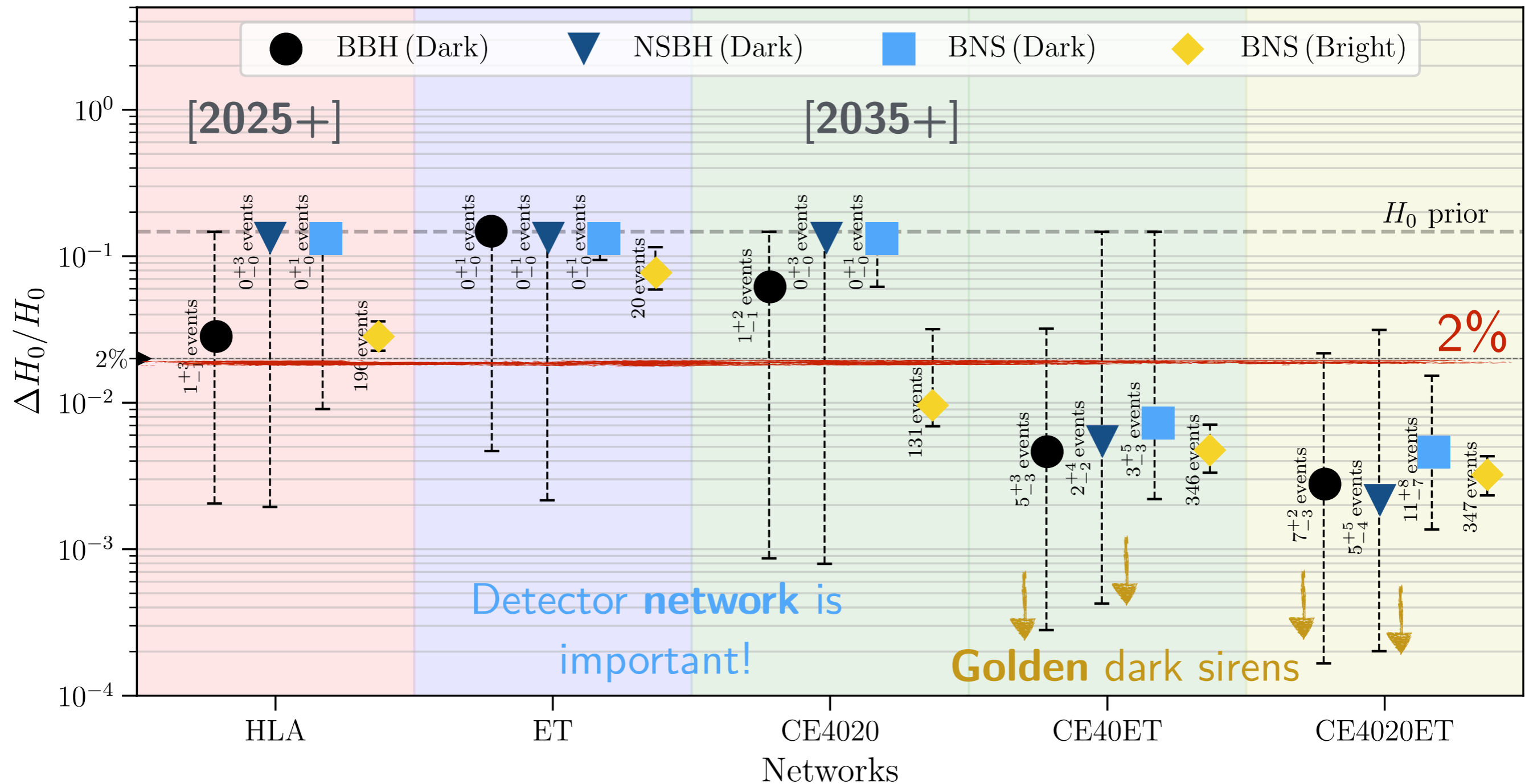
Abstract

Advancements in cosmology through next-generation (XG) ground-based gravitational wave (GW) observatories will bring in a paradigm shift. We explore the pivotal role that GW standard sirens will play in inferring cosmological parameters with XG observatories, not only achieving exquisite precision but also opening up unprecedented redshifts. We examine the merits and the systematic biases involved in GW standard sirens utilizing binary black holes, binary neutron stars, and neutron star-black hole mergers. Further, we estimate the precision of bright sirens, golden dark sirens, and spectral sirens for these binary coalescences and compare the abilities of various XG observatories (A[#], cosmic explorer, Einstein telescope, and their possible networks). When combining different sirens, we find sub-percent precision over more than 10 billion years of cosmic evolution for the Hubble expansion rate $H(z)$. This work presents a broad view of opportunities to precisely measure the cosmic expansion rate, decipher the elusive dark energy and dark matter, and potentially discover new physics in the uncharted Universe with XG GW detectors.

Keywords: gravitational wave, standard siren, cosmology [Chen, Ezquiaga & Gupta (CQG'24)] 6

H_0 (also) with dark sirens

H: Hanford (US)
 L: Livingston (US)
 A: Aundha (India)
 ET: Einstein Telescope (EU)
 CE: Cosmic Explorer (US)

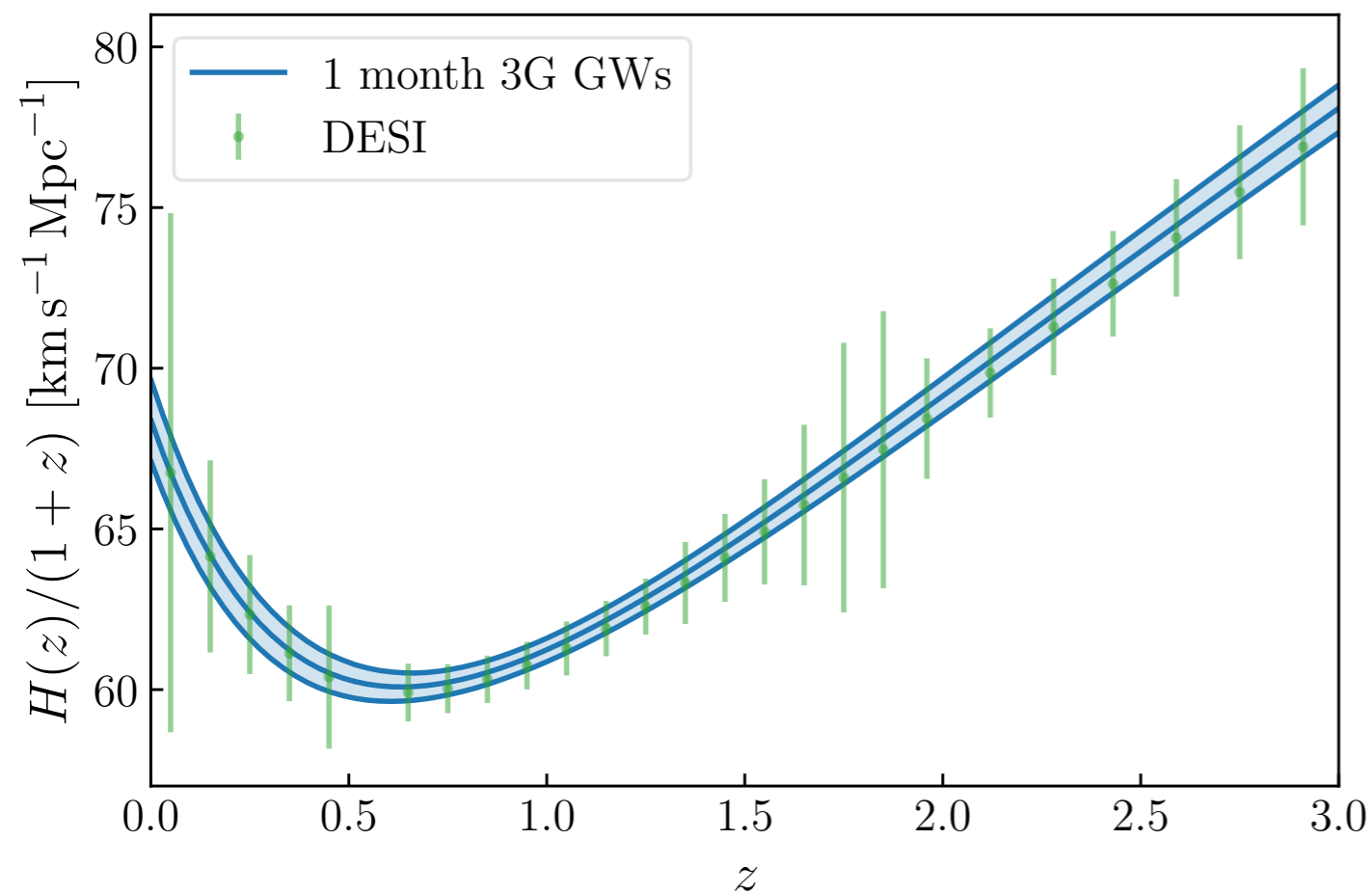


[Chen, Ezquiaga & Gupta (CQG'24)]

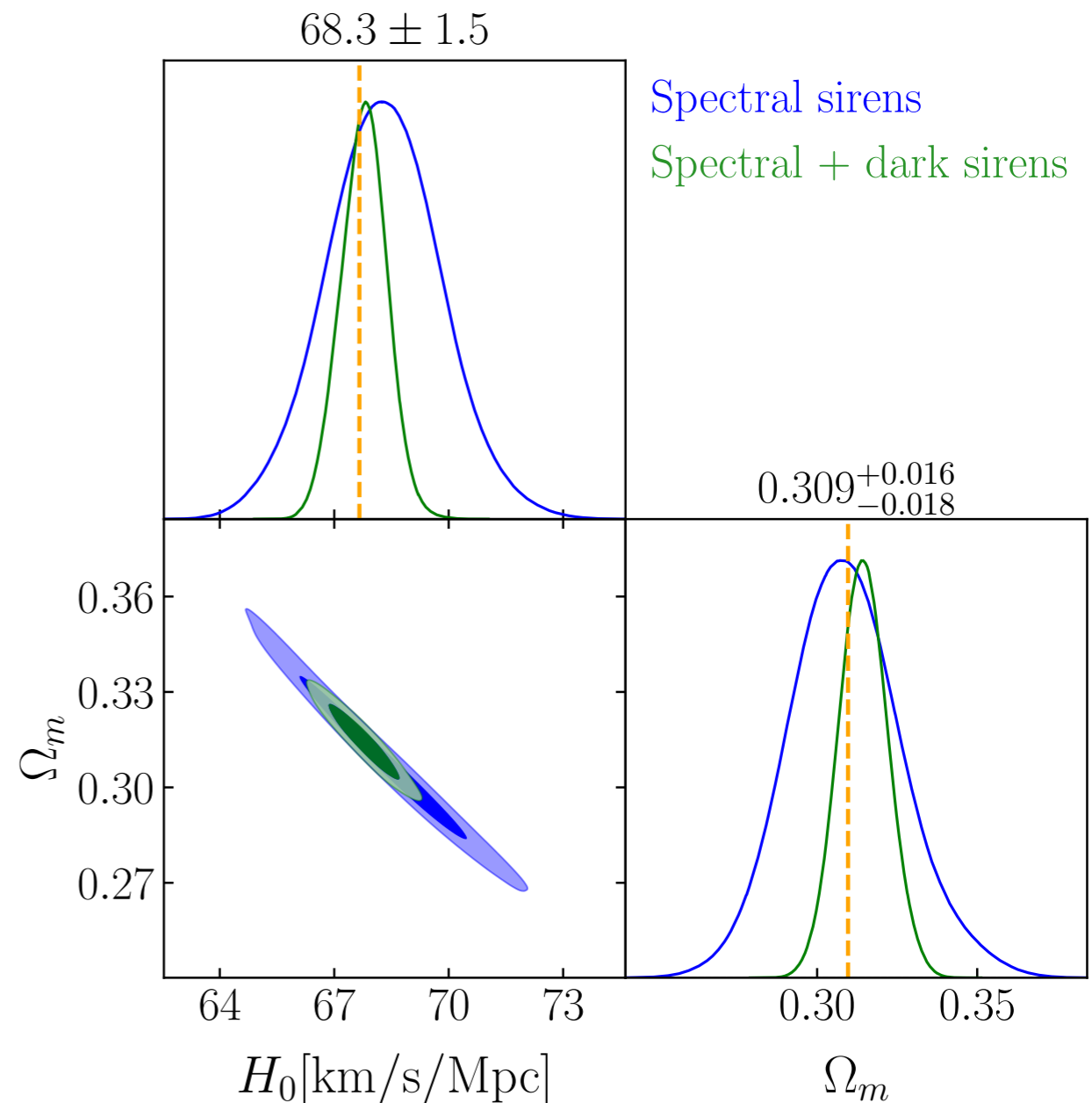
Expansion rate at high redshift $H(z)$

Combining sirens **sub-percent** precision across cosmic history!

Spectral sirens are competitive
with cosmic surveys



[Ezquiaga & Holz (PRL'22)]



[Chen, Ezquiaga & Gupta (CQG'24)]

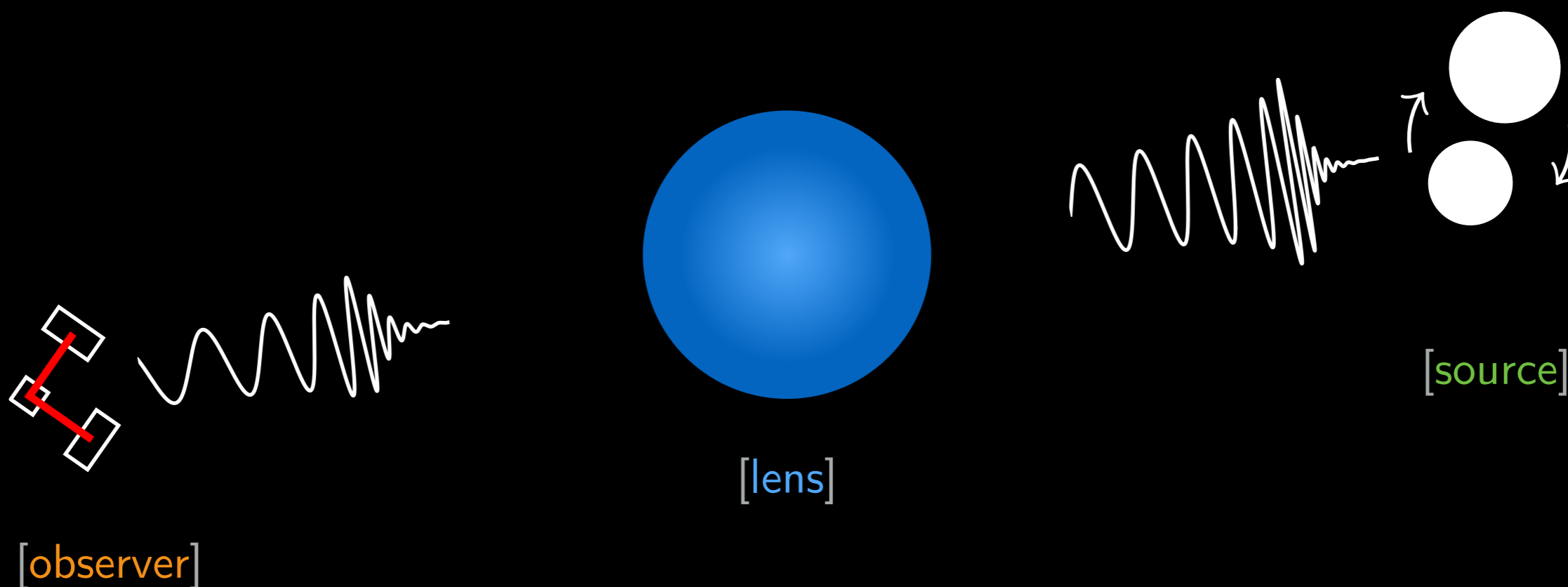


Gravitational lensing

- Solve GW propagation on a **curved** background

$$\square \bar{h}_{\mu\nu} + 2\bar{R}_{\alpha\mu\beta\nu} \bar{h}^{\alpha\beta} = 0$$

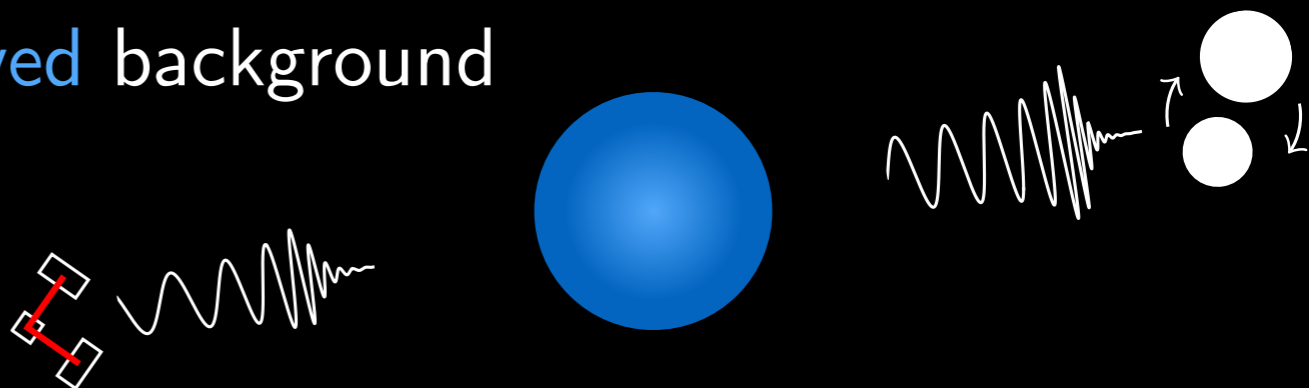
- We want to make a mapping between the **source** and the **observer** through the **lens**



Gravitational lensing

- Solve GW propagation on a **curved** background

$$\square \bar{h}_{\mu\nu} + 2\bar{R}_{\alpha\mu\beta\nu} \bar{h}^{\alpha\beta} = 0$$



- Within **weak-gravity** & **thin lens** approximations, in **Fourier** space:

$$h_L(\omega) = F(\omega, \theta_S) \cdot h(\omega)$$

$$F(\omega, \vec{y}) = \frac{\omega}{2\pi i} \int d^2x \exp[i\omega T_d(\vec{x}, \vec{y})]$$

[Dimensionless variables] $\vec{x} \equiv \vec{\theta}/\theta_*$, $\vec{y} \equiv \vec{\theta}_S/\theta_*$, $\omega \equiv \tau_D \theta_*^2 \omega$

$$T_d \equiv t_d/\tau_D \theta_*^2 \quad \tau_D \equiv (1 + z_L) D_L D_S / c D_{LS}$$

Stationary Phase Approximation

- Solve integral in the limit of highly oscillatory integrand

$$F(w, \vec{y}) = \frac{w}{2\pi i} \int d^2x \exp[iwT_d(\vec{x}, \vec{y})]$$

- Stationary points define the **images**:

$$\left. \frac{\partial t_d}{\partial \theta_a} \right|_{\vec{\theta}=\vec{\theta}_j} = 0$$

$$T_d(\vec{\theta}) \approx T_d(\vec{\theta}_j) + \frac{1}{2} \sum_{(a,b)=1}^2 \delta\theta_a \delta\theta_b \frac{\partial^2 T_d(\vec{\theta}_j)}{\partial \theta_a \partial \theta_b} + \dots$$

- Hessian matrix determines magnifications

$$\mu(\theta_j) = 1/\det(T_{ab}(\theta_j))$$

$$T_{ab} \equiv \tau_D^{-1} \partial^2 t_d / \partial \theta_a \partial \theta_b$$

Strong lensing

$$\Delta t_d \cdot \omega \gg 1$$

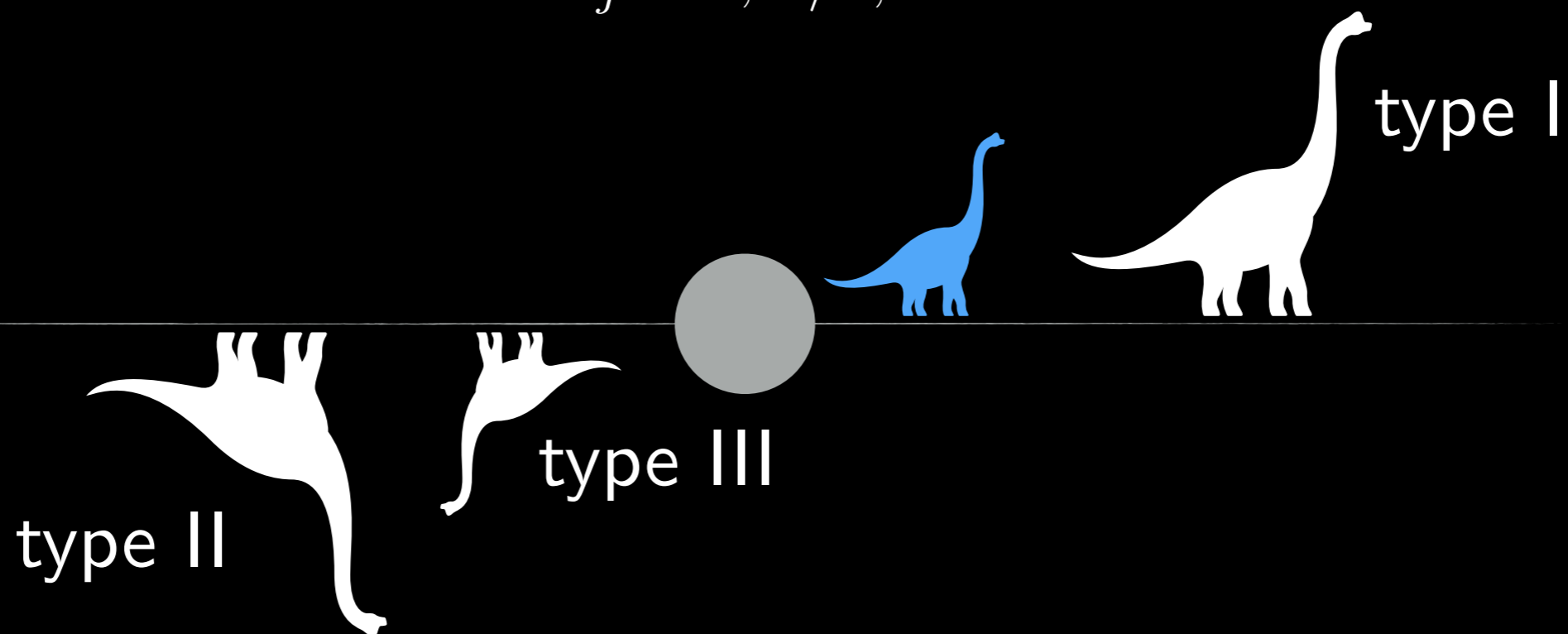
$$h_L(\omega) = F(\omega, \theta_S) \cdot h(\omega)$$

$$F \approx \sum_j |\mu_j|^{1/2} \exp(i\omega t_j - i\pi n_j)$$

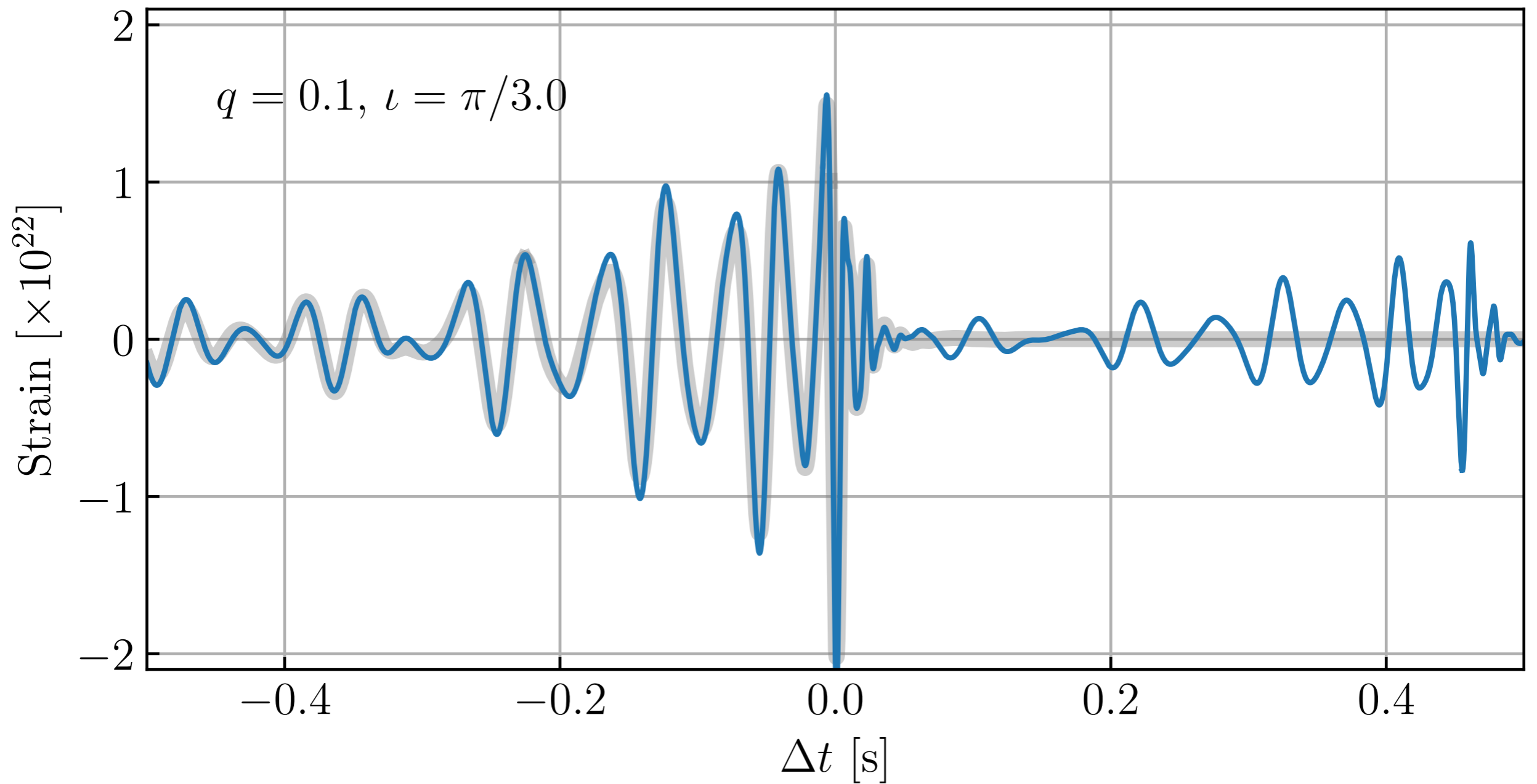
Magnification
Time delay
Phase shift

- Each image type (I, II and III) acquire a different phase shift

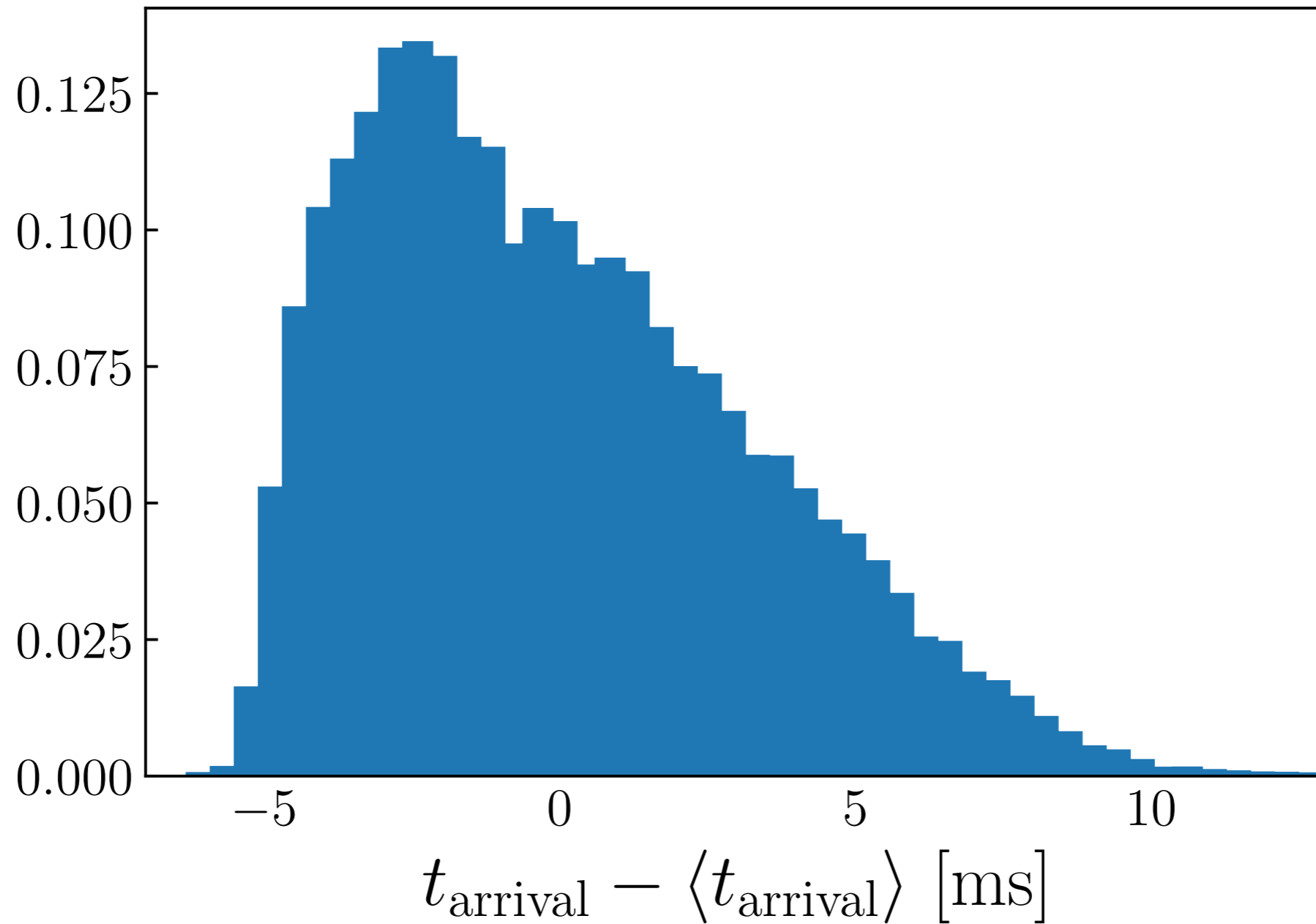
$$n_j = 0, 1/2, 1$$



Repeated, coherent signals

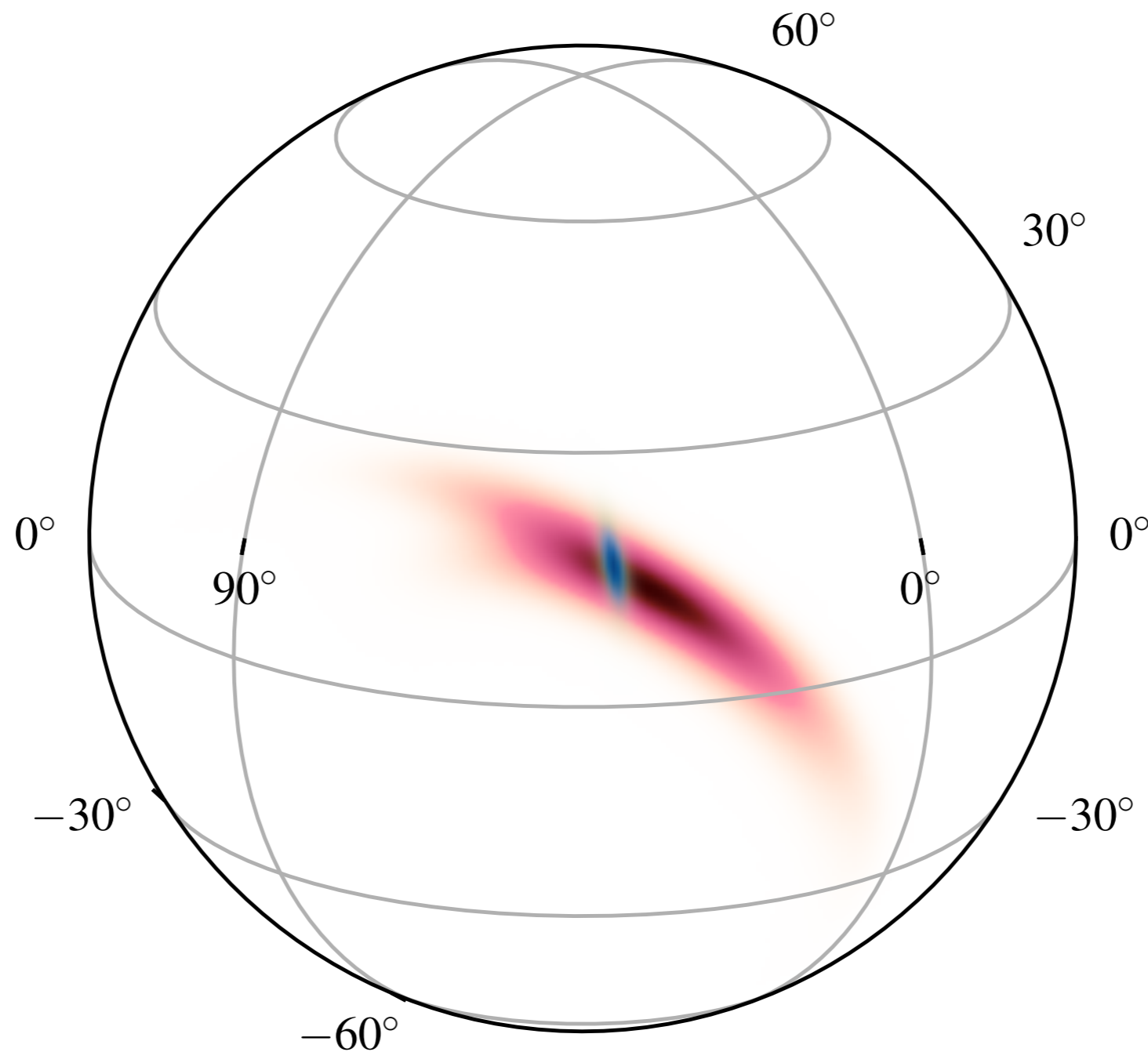


Precise timing



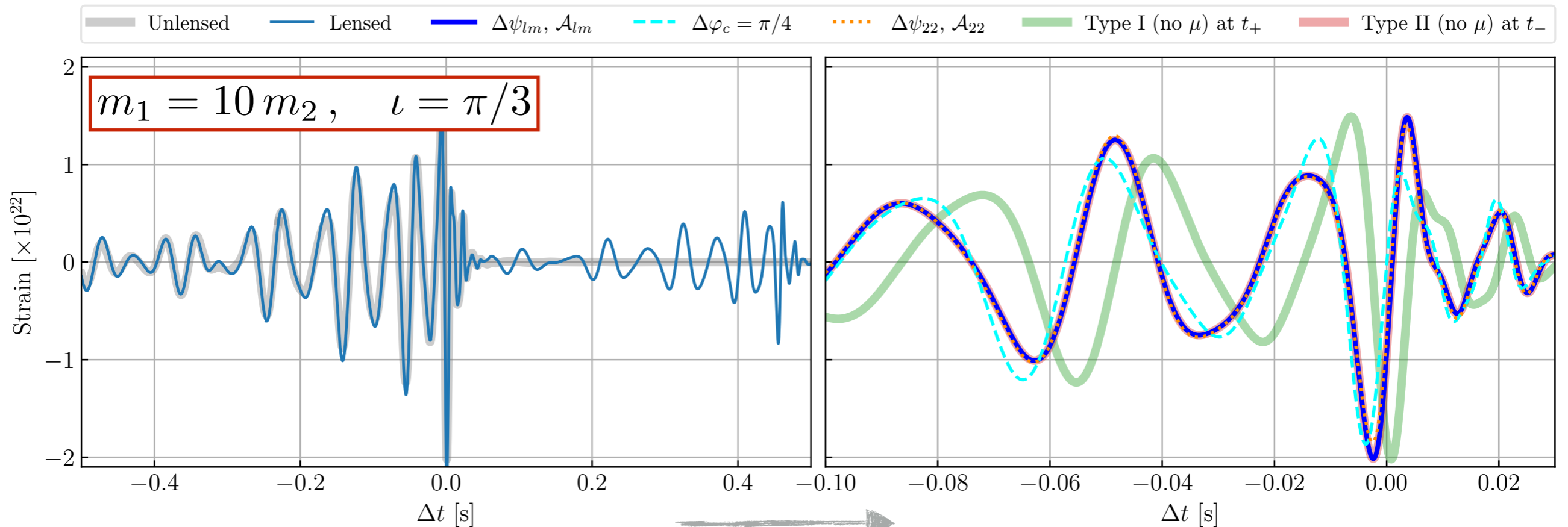
Poor sky localization

$$\theta_E \sim 1'' \sqrt{\frac{M}{10^{12} M_\odot}} \sqrt{\frac{1 \text{ Gpc}}{D}}$$



Waveform distortions in **type II** images

- Lensing imprints *small* but *characteristic* modifications in the signals that cannot be mapped to other astrophysical parameters



Zoom in the shape
of each image

Wave optics

$$\Delta t_d \cdot \omega$$

- Time delay scales with the lens mass

$$\Delta t_d(y = 1) \simeq 4 \left(\frac{(1 + z_L) M_L}{100 M_\odot} \right) \text{ ms} \quad [\text{point mass lens}]$$

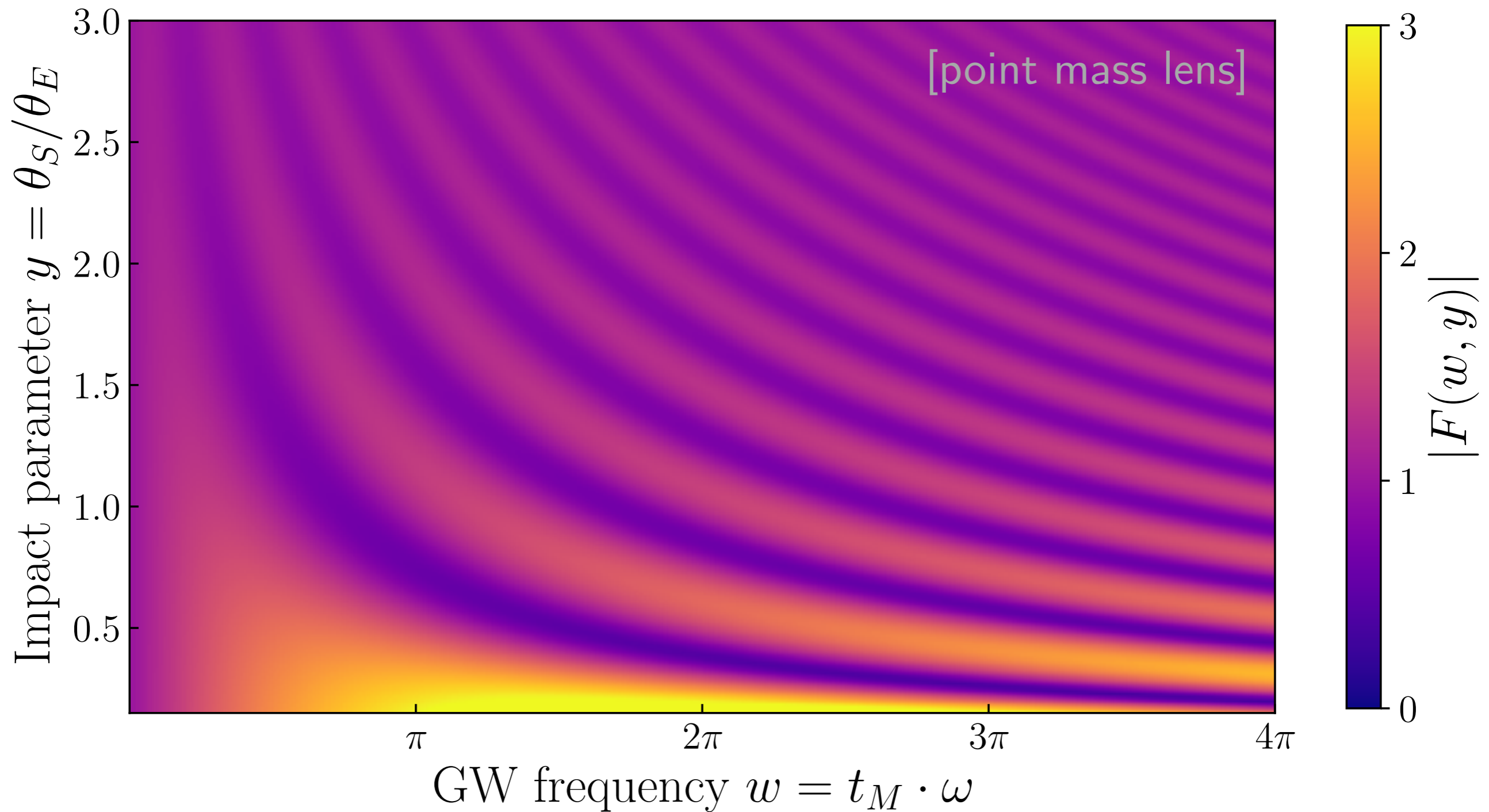
- GW frequency scales with binary mass (*has astrophysical size!*)

$$f \sim \frac{1}{2\pi} \frac{1}{2t_{\text{Sch}}} \sim 800 \text{ Hz} \left(\frac{10 M_\odot}{M} \right)$$

- Wave optics regime: $\Delta t_d \cdot \omega \sim 1$

- Low-frequency limit has small lensing $\omega \rightarrow 0 \Rightarrow F \rightarrow 1$

Wave optics: point lens

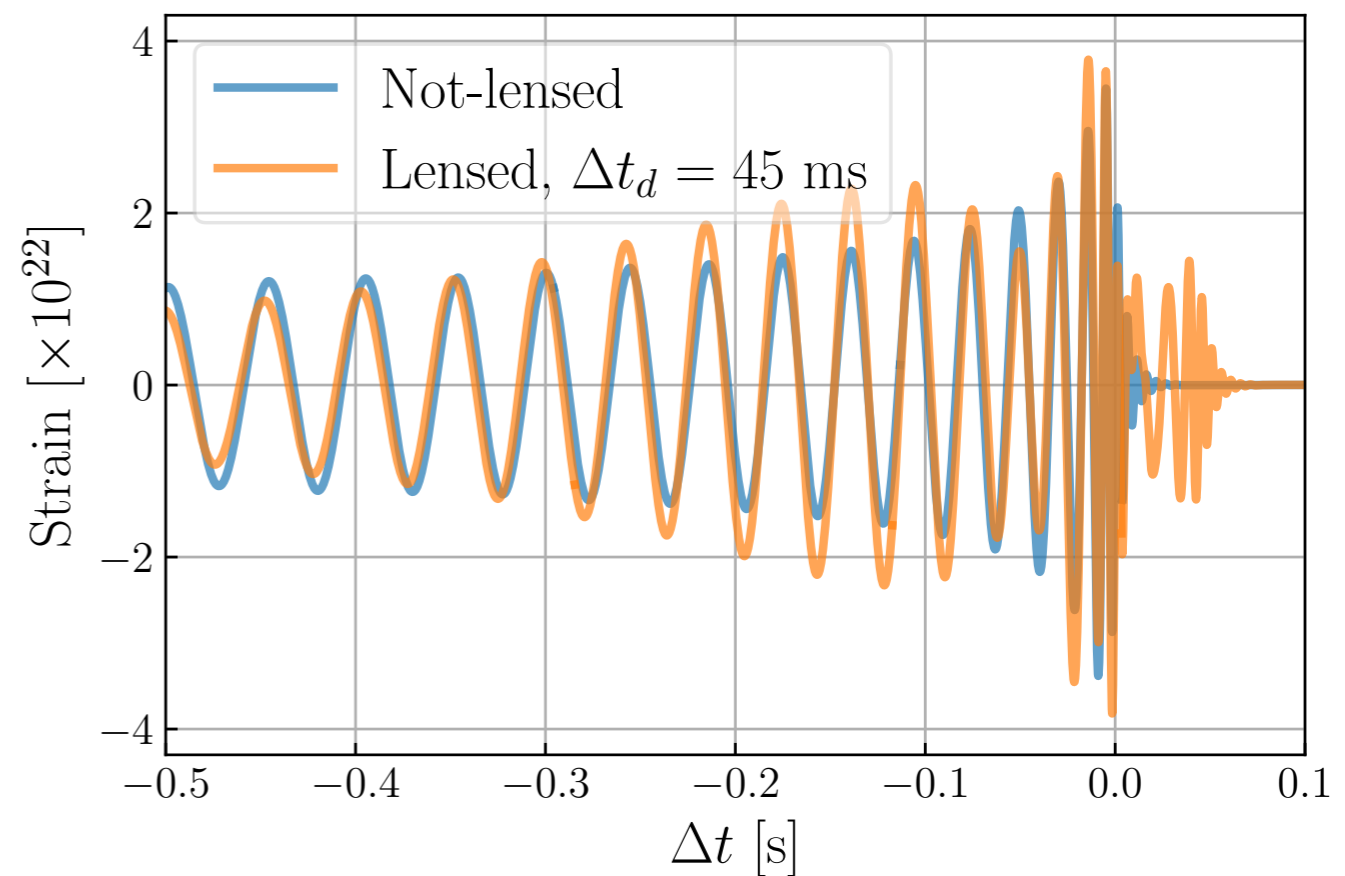
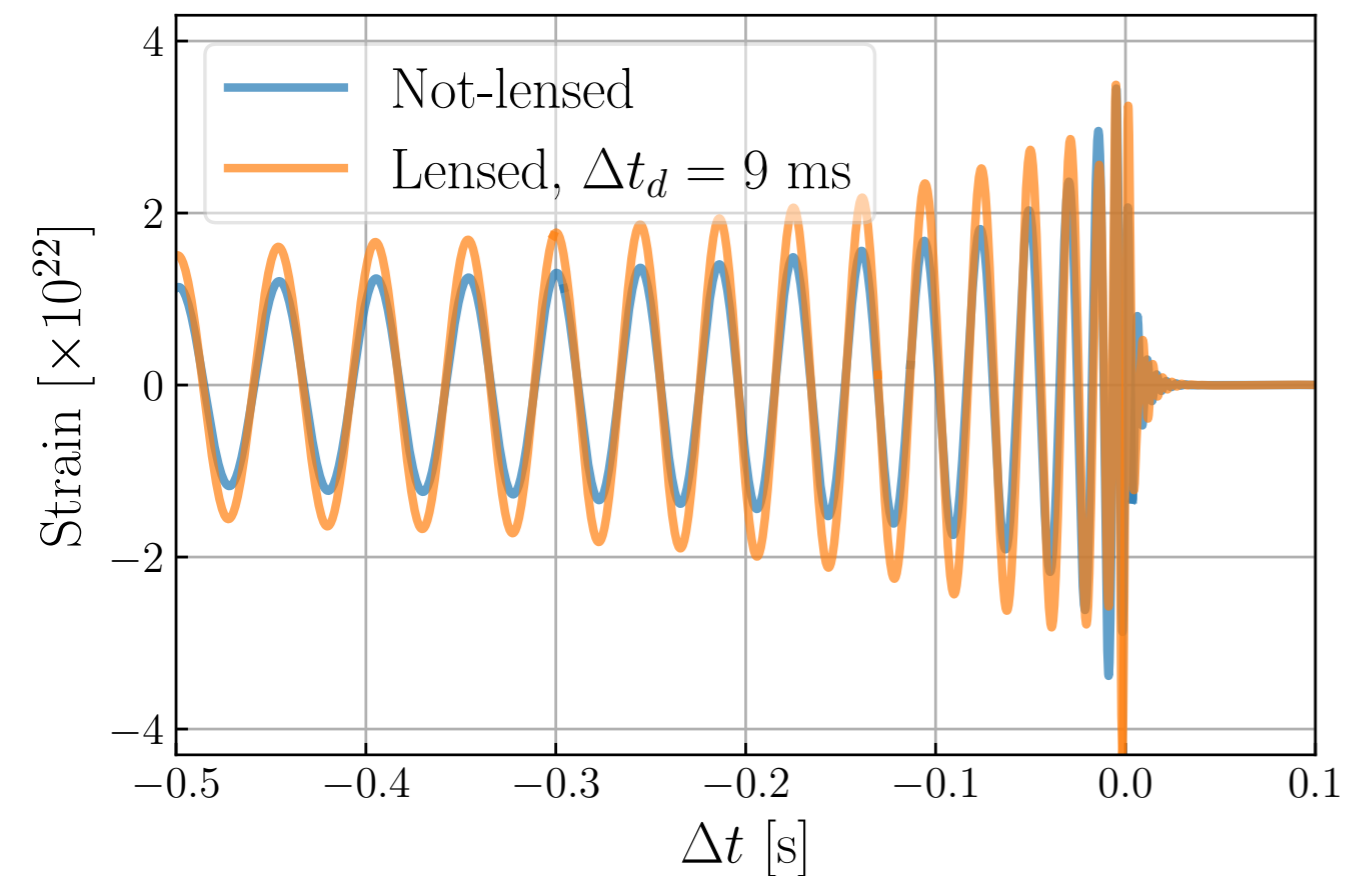


Wave optics: point lens

$$\Delta t_d(y = 1) \simeq 4 \left(\frac{(1 + z_L) M_L}{100 M_\odot} \right) \text{ ms}$$

Diffraction

Interference



Many recent works: probing compact objects ([Dai et al.'18](#), [Diego'19](#), [Tambalo et al.'22](#), ...), strong lensing + microlensing ([Seo et al.'21](#), [Mena et al.'22](#), ...), breaking mass-sheet degeneracy ([Cremonese, Ezquiaga, Salzano'21](#)), solving diffraction integral ([Feldbrugge&Turok'20](#), [Tambalo et al.'22](#)) ... +++

Caustics

- For point sources, there are **singular** points in the lens mapping

$$\det \left(\frac{\partial^2 T_d(\theta_j)}{\partial \theta_a \partial \theta_b} \right) \rightarrow 0 \quad \Rightarrow \quad \mu(\theta_j) \rightarrow \infty$$

- Caustics exhibit **universal** behaviors (described by catastrophe theory)

$$\mu_{\pm} \sim 1/\sqrt{\Delta\theta_S} \sim \Delta t^{-1/3}$$

- SPA is **broken** when approaching to a caustic
- **Maximum magnification** set by diffraction



Observational Signatures of Highly Magnified Gravitational Waves from Compact Binary Coalescence

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(Dated: July 26, 2024)

Gravitational lensing has empowered telescopes to discover astronomical objects that are otherwise out of reach without being highly magnified by foreground structures. While we expect gravitational waves (GWs) from compact binary coalescences to also experience lensing, the phenomenology of highly magnified GWs has not been fully exploited. In this letter, we fill this gap and explore the observational signatures of these highly magnified GWs. We find that these signatures are robust against modeling details and can be used as smoking-gun evidence to confirm the detection of lensing of GWs without any electromagnetic observation. Additionally, diffraction becomes important in some cases, which limits the maximum possible magnification and gives waveform signatures of lensing that can only be observed by GW detectors. Even with current-generation observatories, we are already sensitive to these highly magnified GWs and can use them to probe the high-redshift Universe beyond the usual horizon.

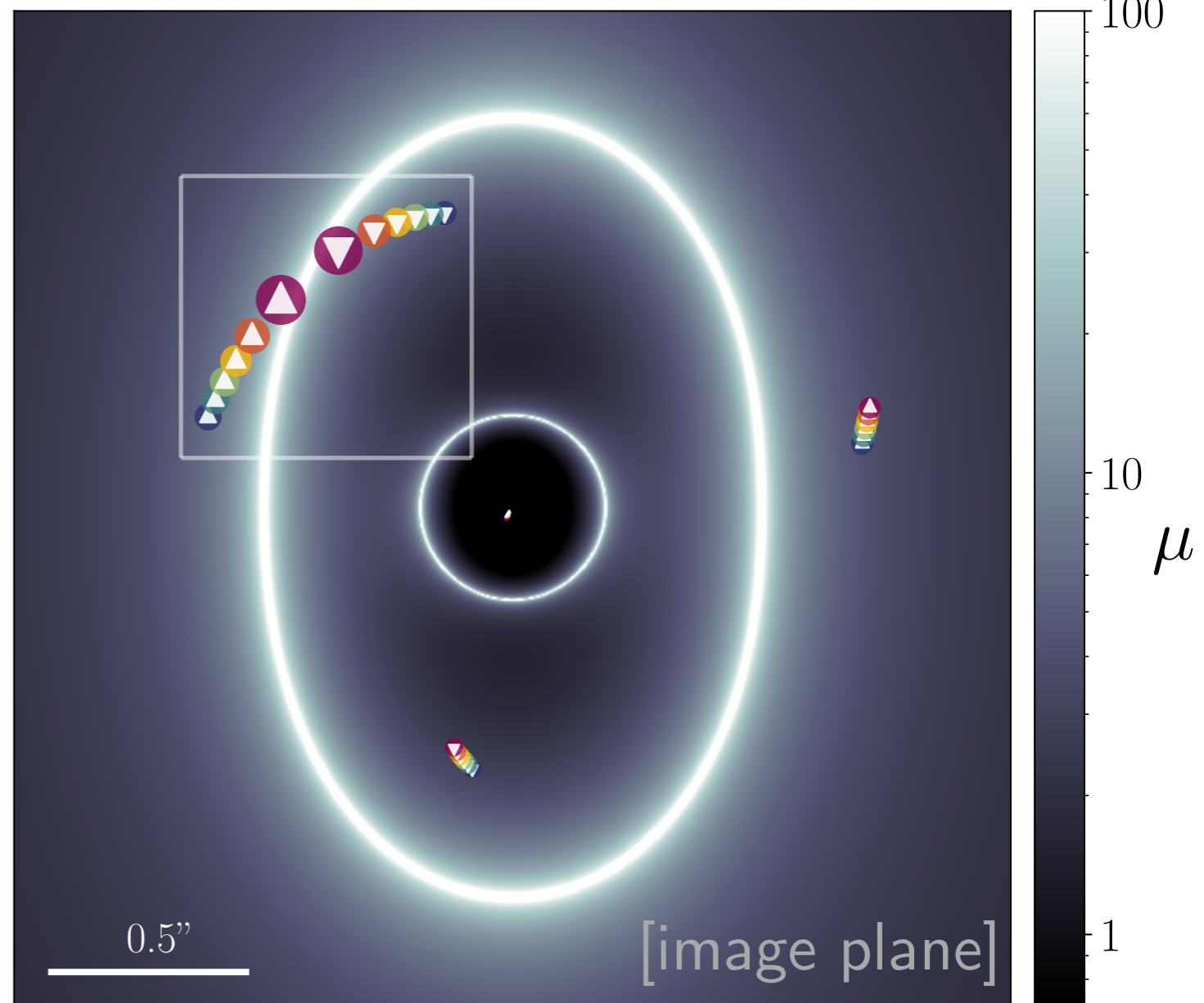
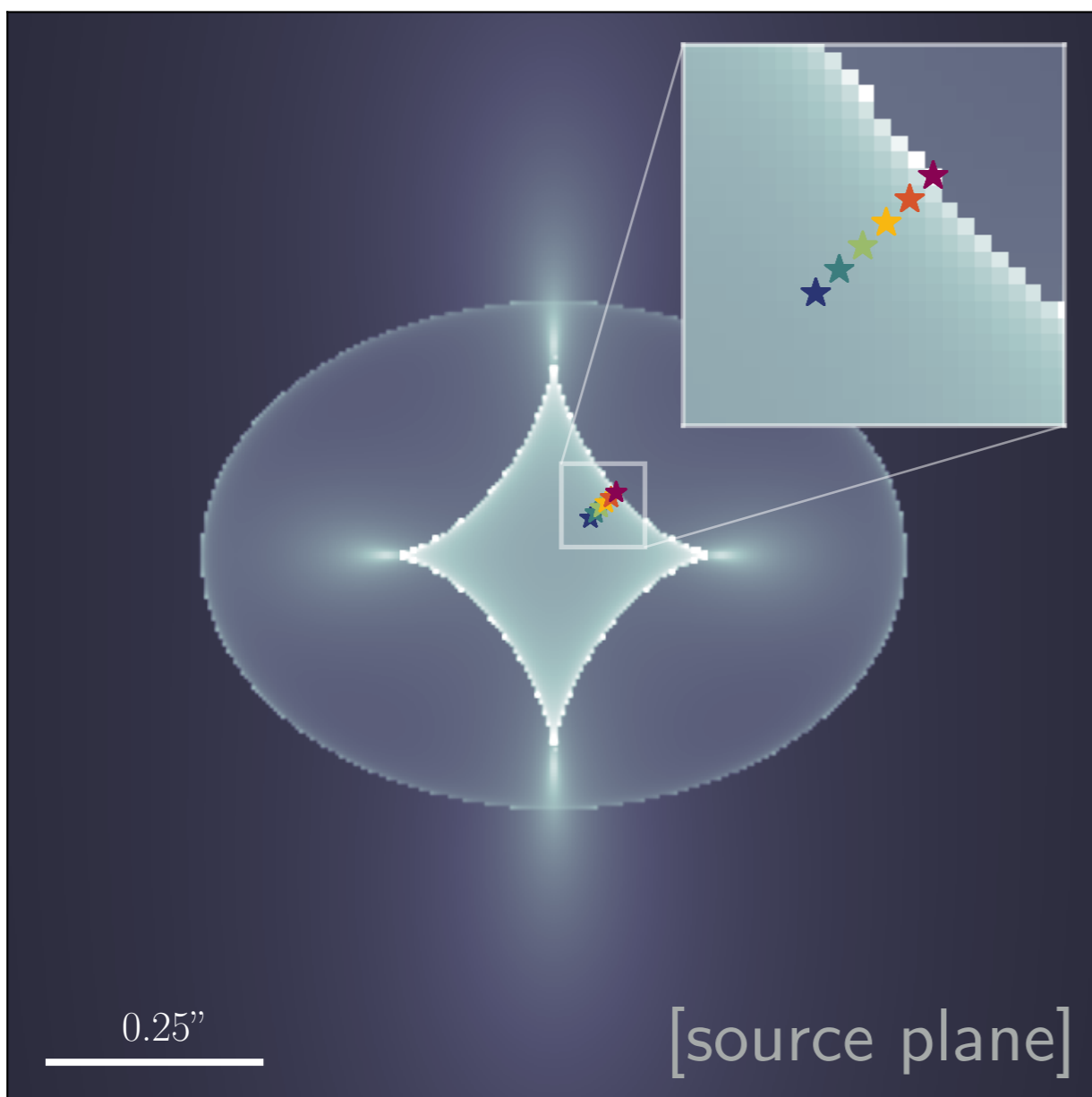


Rico Lo



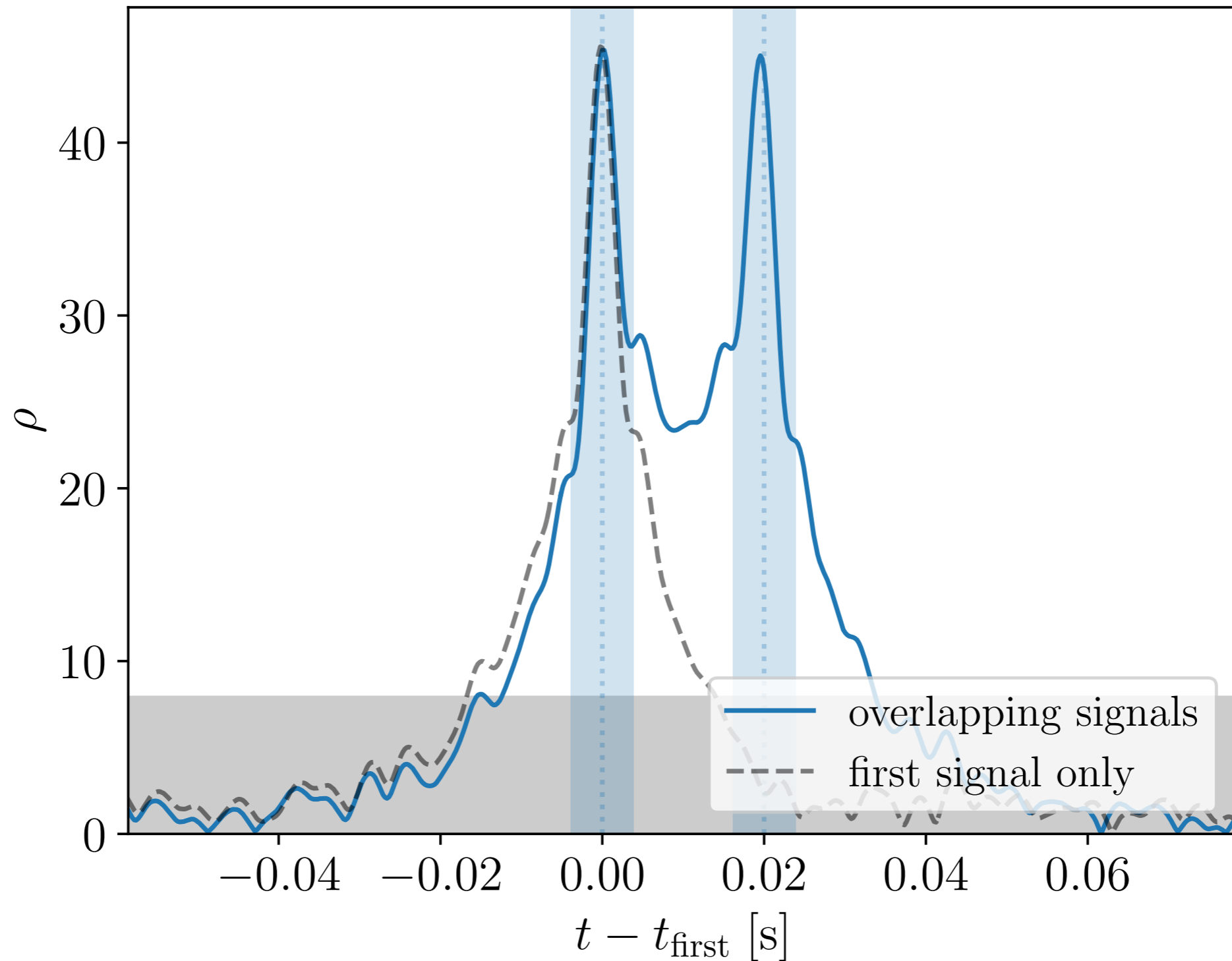
Luka Vujeva

Fold caustics around galaxies

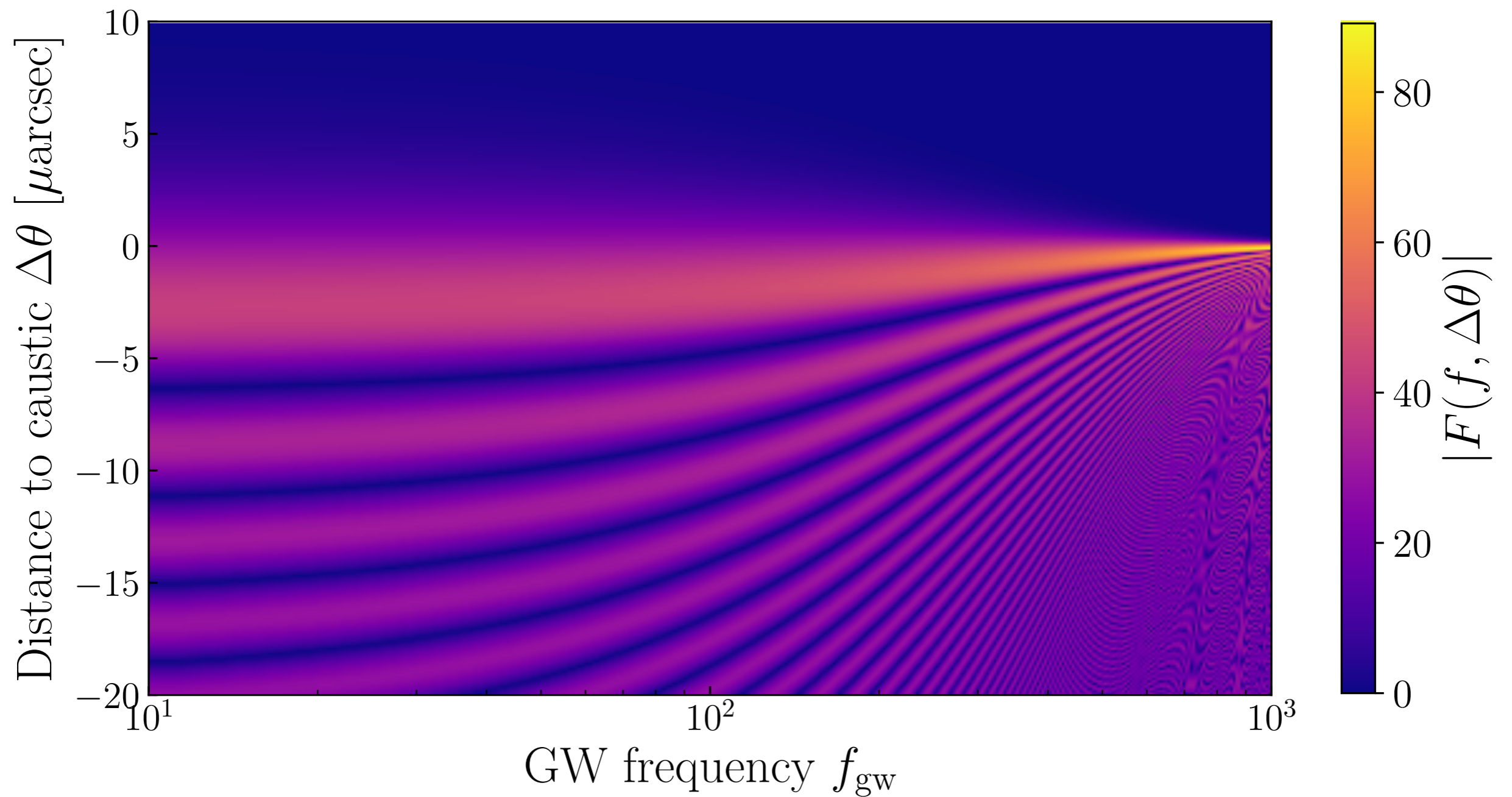


[galaxy lens with a cored singular isothermal ellipsoid density profile]

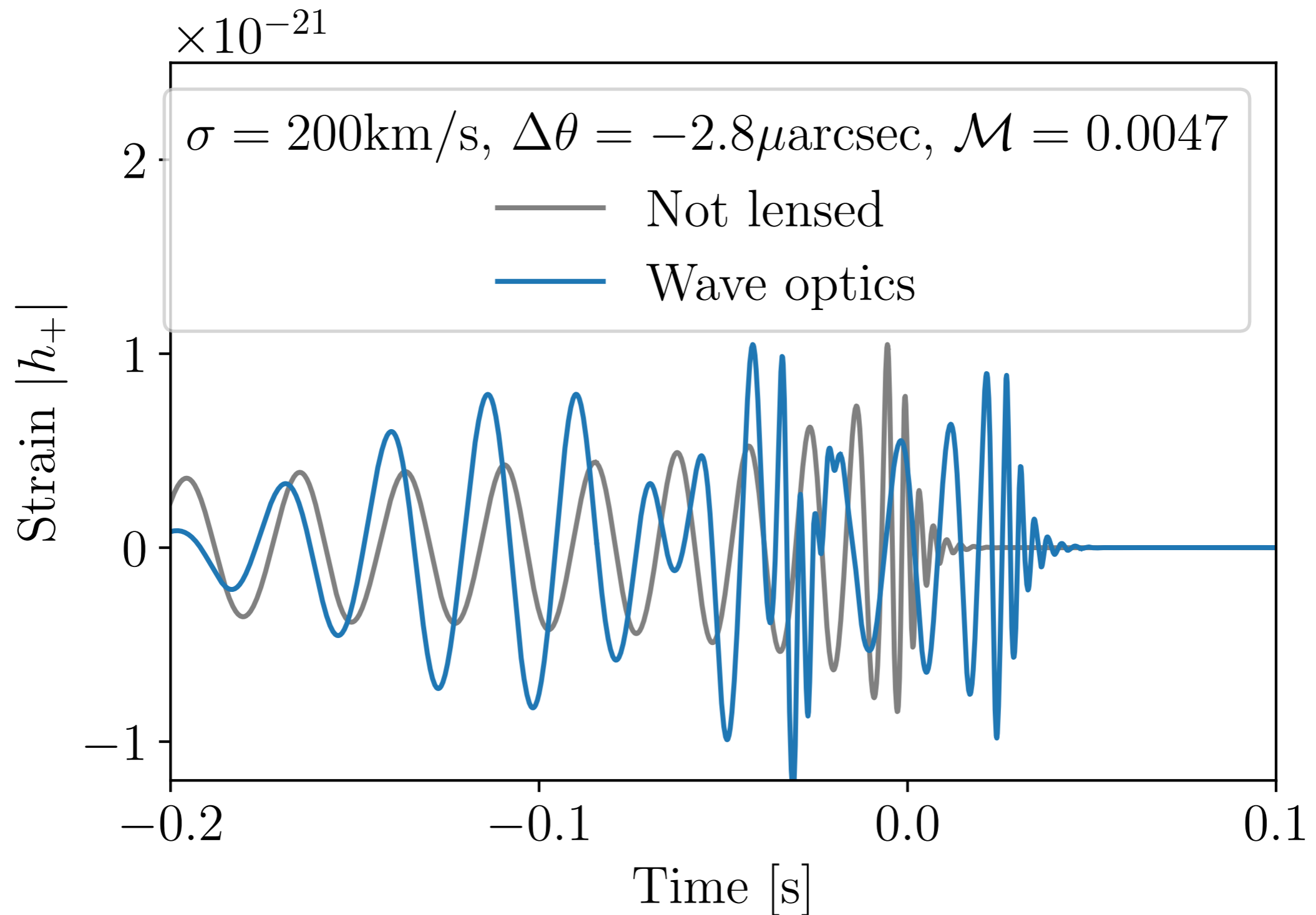
Overlapping signals



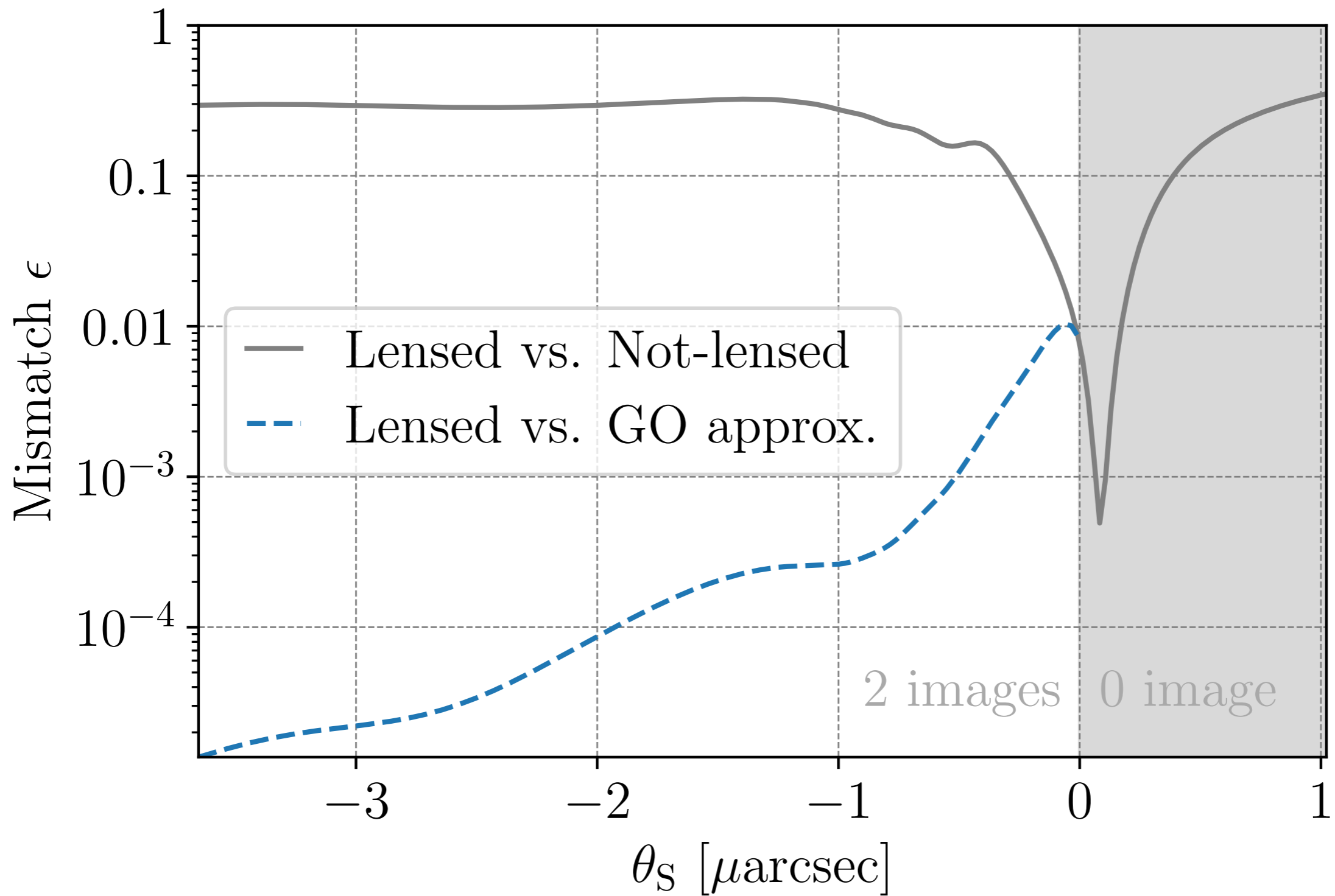
Diffraction around fold caustics



Waveform distortions

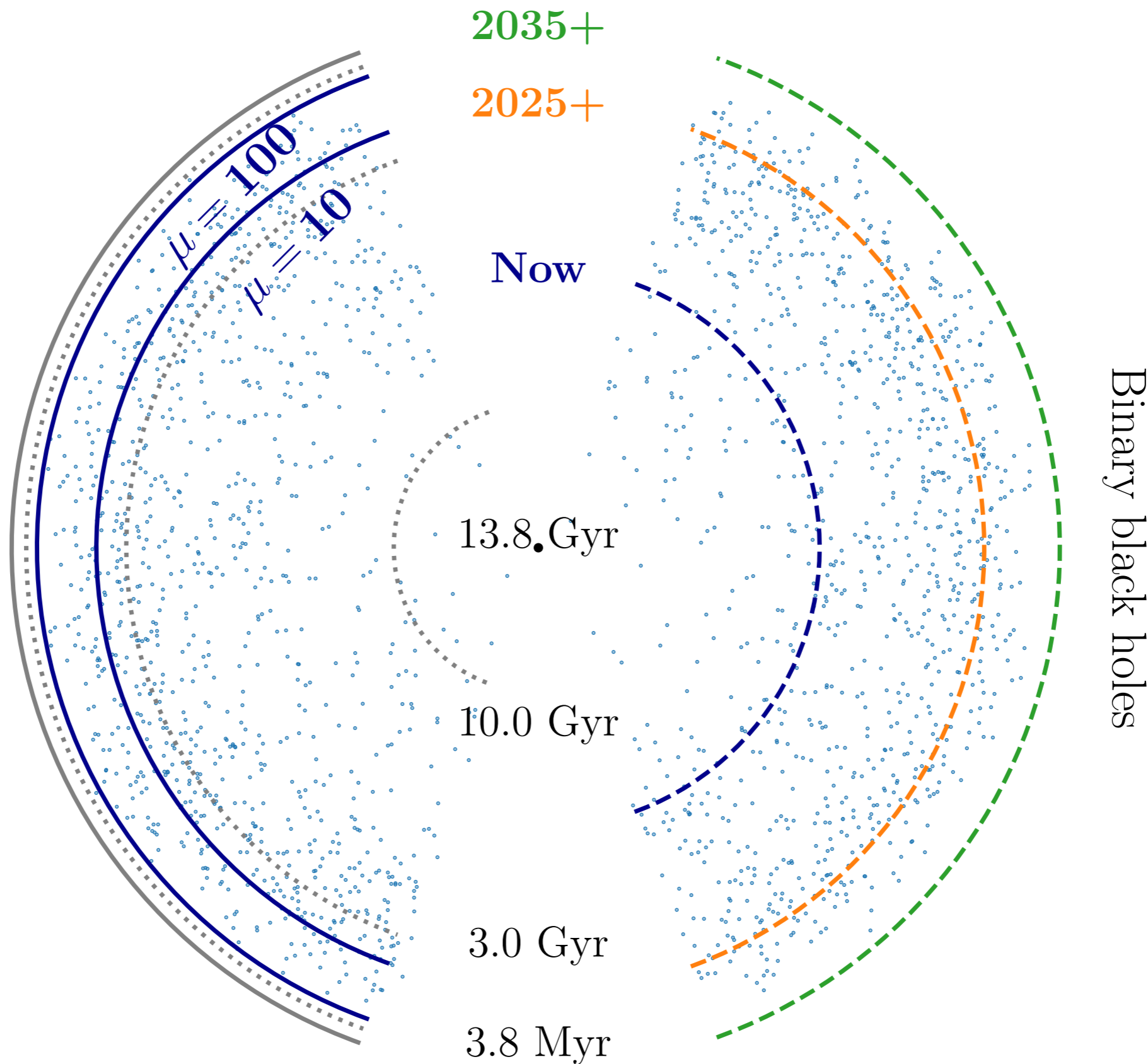


Observing GW diffraction



Gravitational Wave horizons

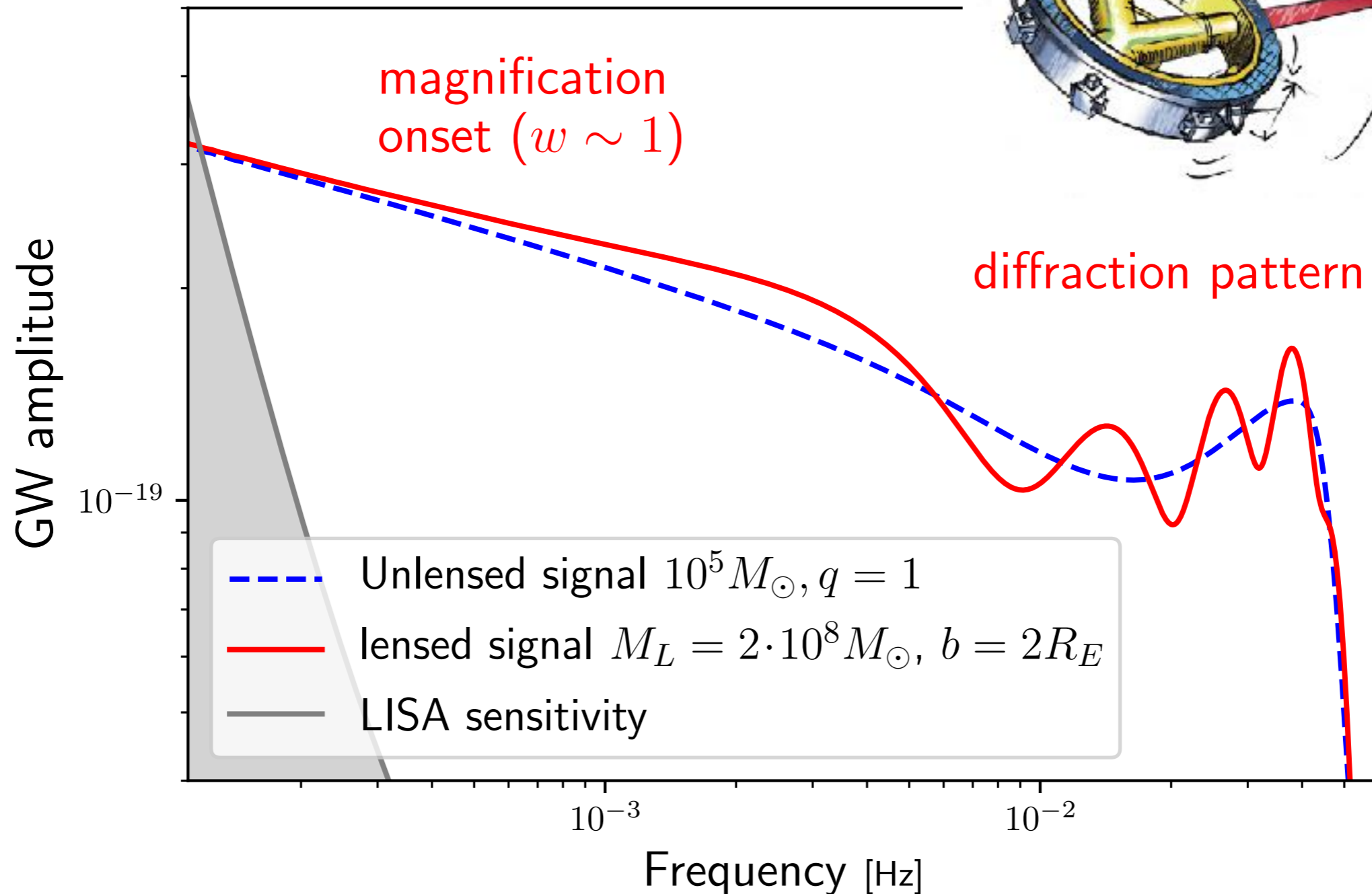
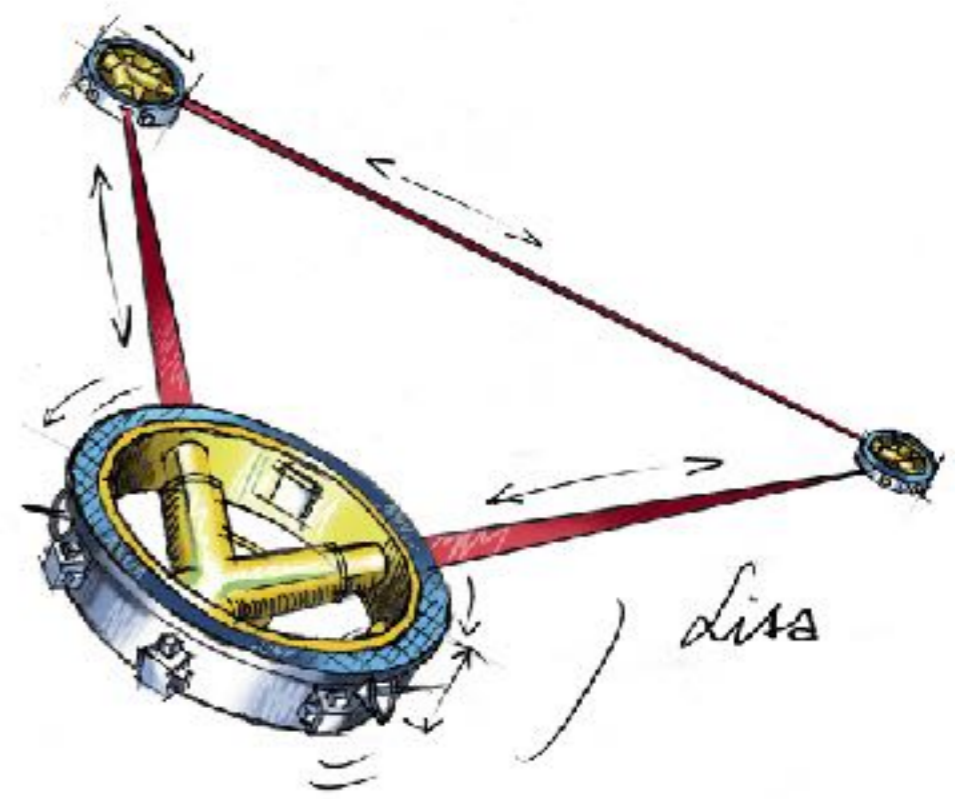
[a shortcut to
next-generation
detectors]



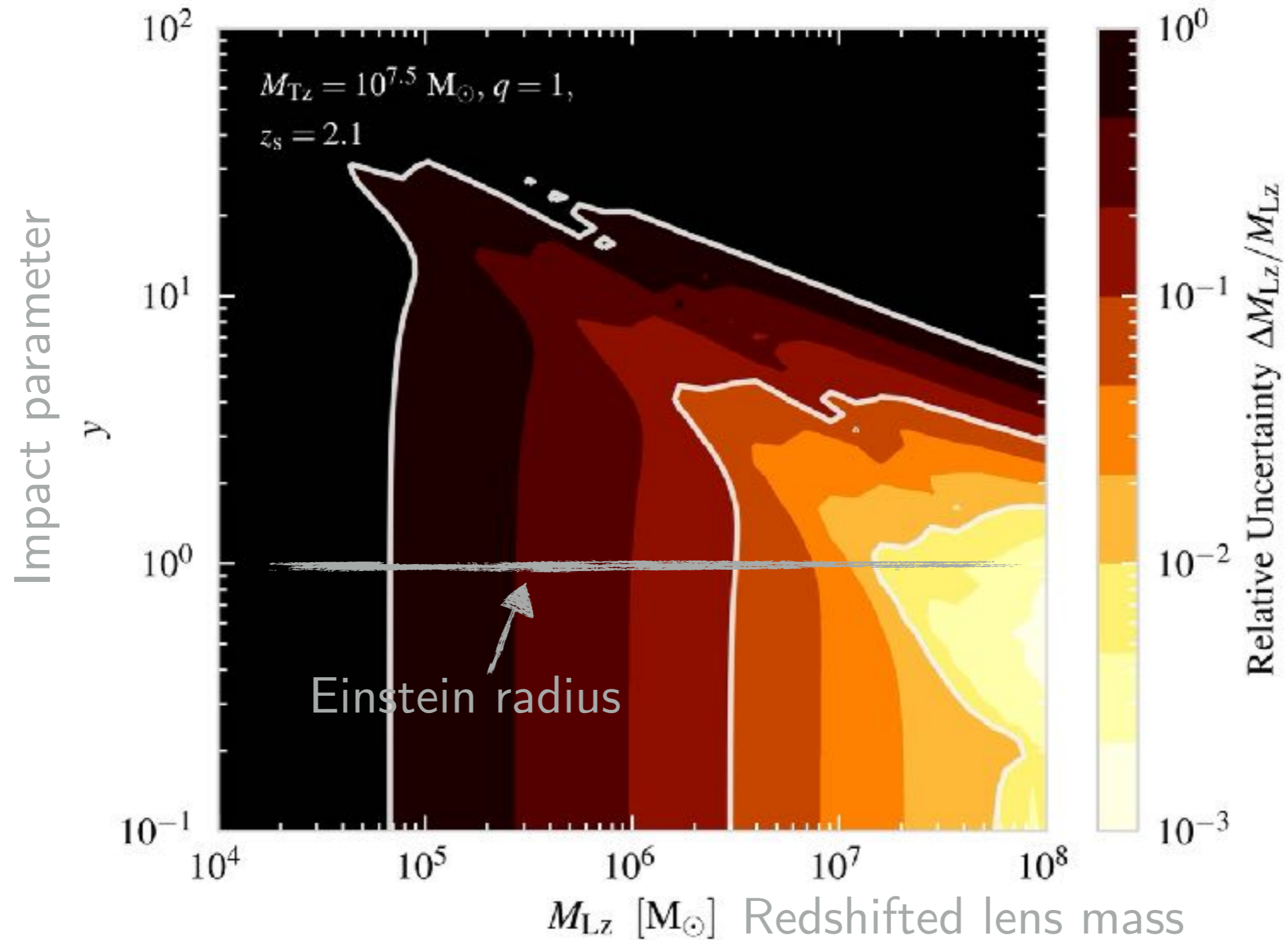
Age of the Universe

*stellar mass
binary black holes

Wave effects: LISA



Increased optical depth in wave optics

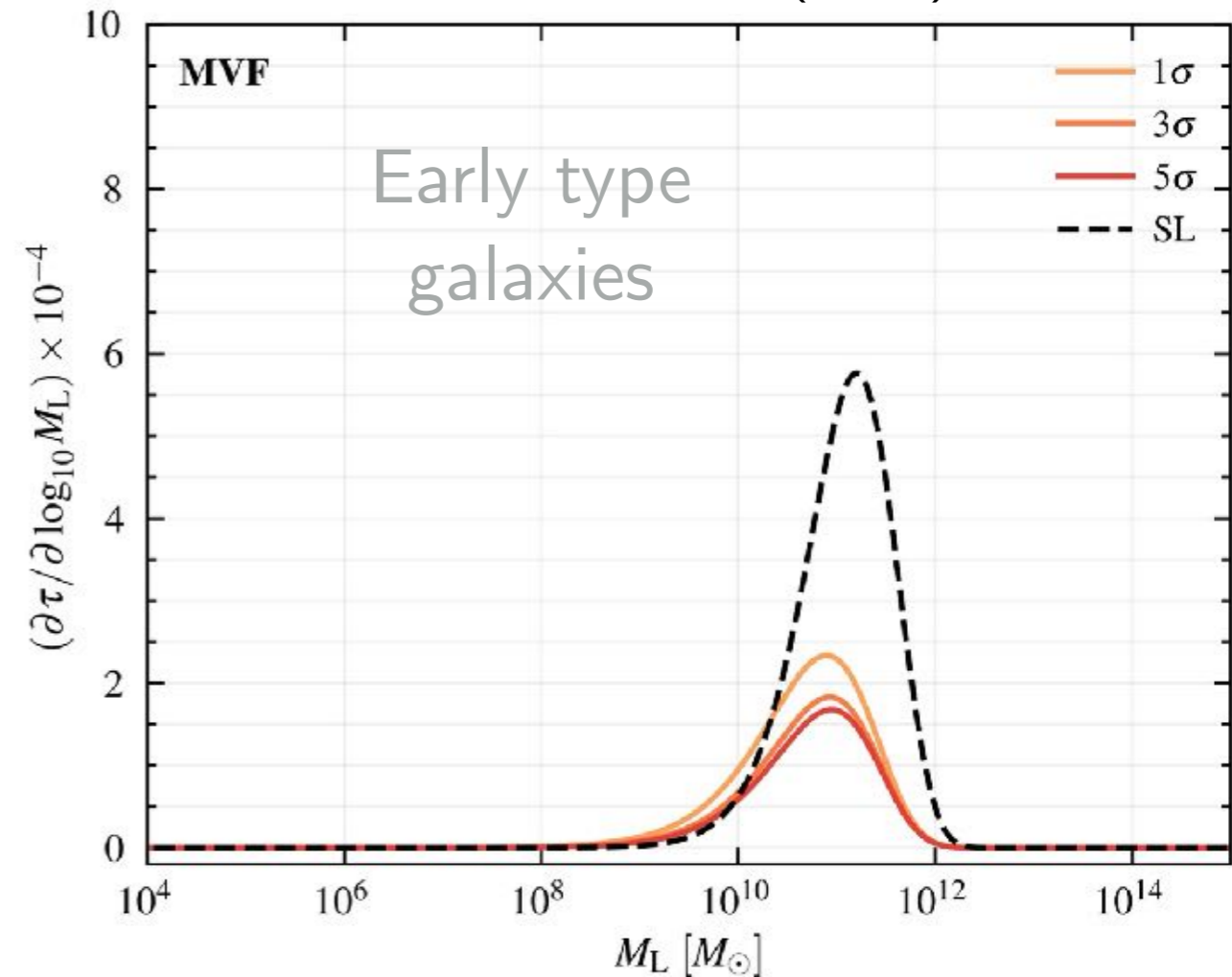
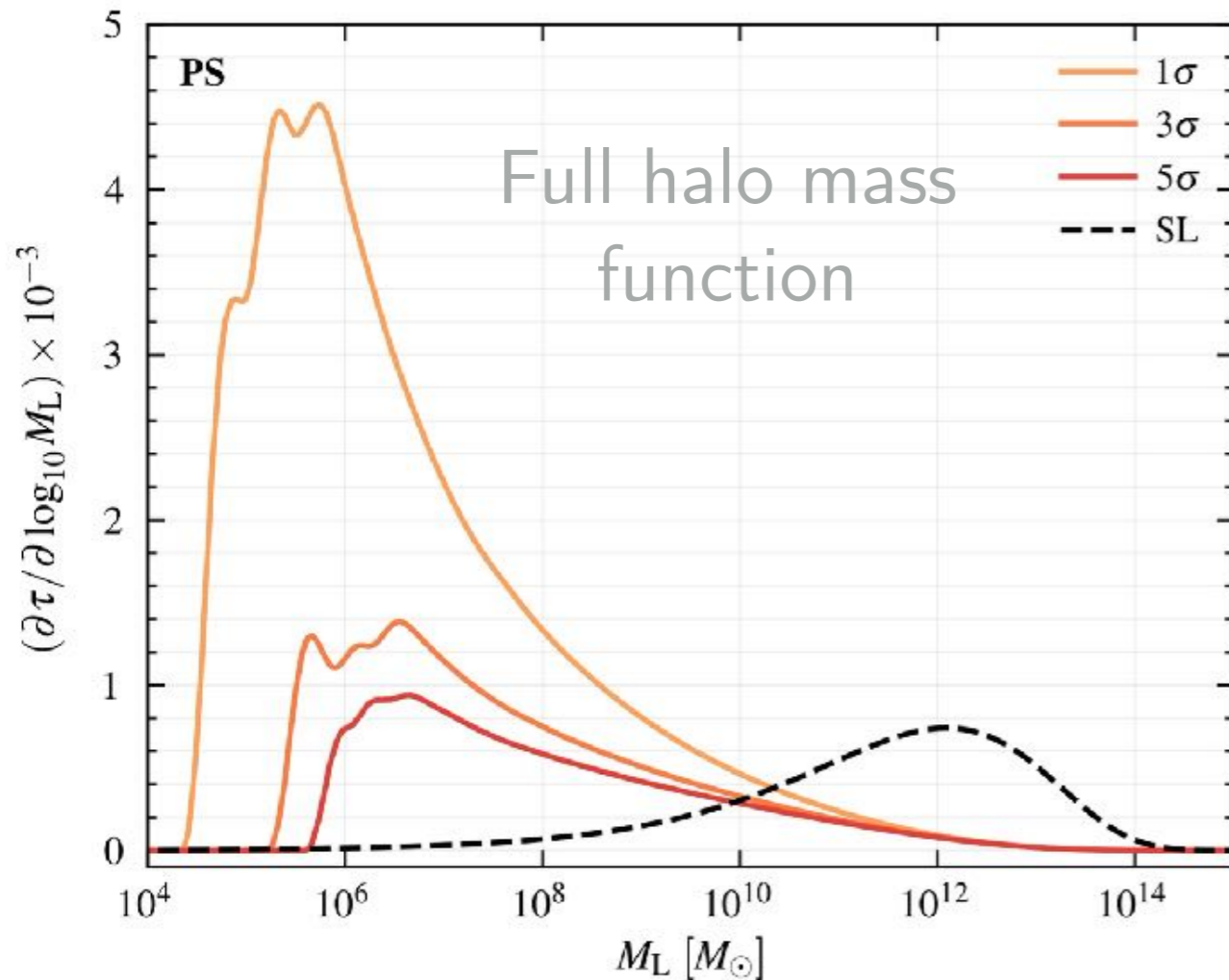


Probing low-mass dark matter halos

- Probability of wave optics lensing:
few percent and larger than strong lensing



Mesut Çalışkan
(JHU)



Conclusions

Binary black holes are precious cosmological probes:

1. H_0 (also) with **dark sirens**
2. Expansion rate at high redshift $H(z)$ with **spectral sirens**
3. Probing **inhomogeneous** Universe via **lensing**
4. **Highly magnified** GWs have **unique** observational signature
5. Probing **low-mass halos** with **LISA** binaries



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