

FUN WITH PTA INTERPRETATIONS

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Fundamental physics and GW detectors

Pollica

September 13, 2024

New era in fundamental (particle) physics

Standard model is valid to high energies

No clear hint where new physics is



New era in fundamental (particle) physics

At the same time:

SM is **incomplete**, does not explain dark matter, baryon asymmetry, inflation, ...

All linked to very early Universe dynamics



New era in fundamental (particle) physics

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SM is **incomplete**, does not explain dark matter, baryon asymmetry, inflation, ...

All linked to very early Universe dynamics

Gravitational waves are messengers from this era

- ▶ great opportunity



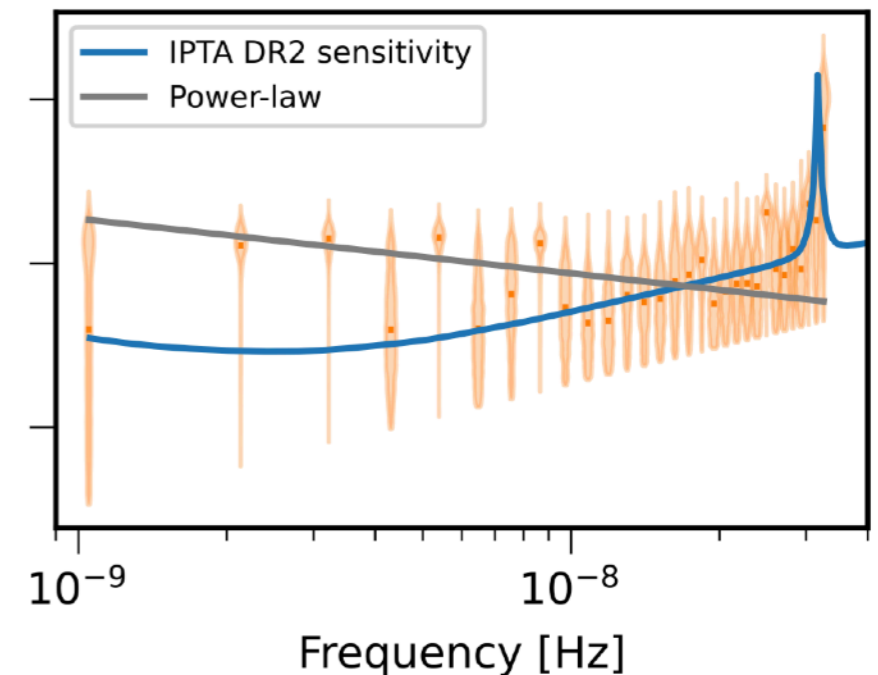
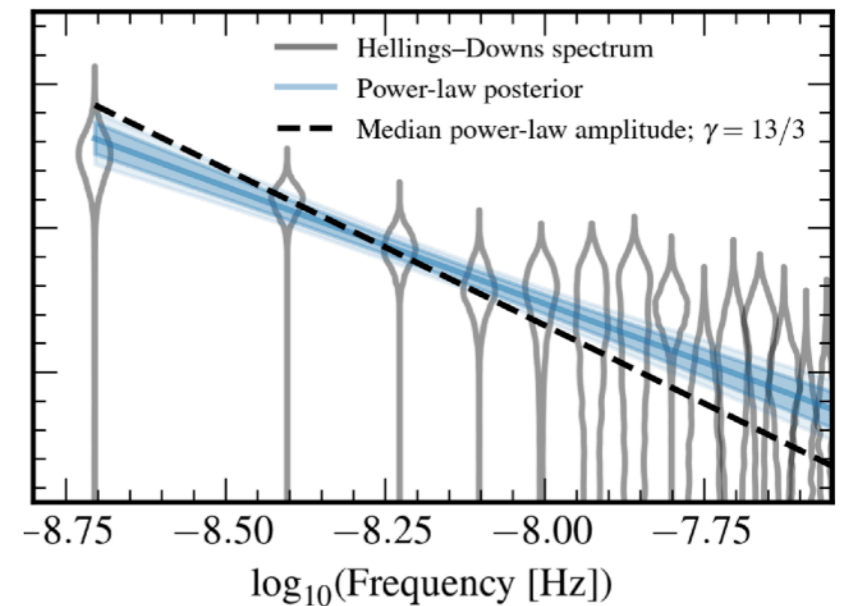
PTA: First observation of stochastic GW background

Could be of primordial origin, though a large astrophysical contribution (SMBHB) is likely

In any case:

Testing ground for model building, parameter reconstruction, ...

- ▶ what can we learn from GW data
- ▶ how else can we probe and distinguish models



Outline

Audible axions: GWs from rolling axions & PTA

GWs from domain walls

- ▶ spectral distortions as complementary probe of GW sources

PTA GWs from supermassive pBH

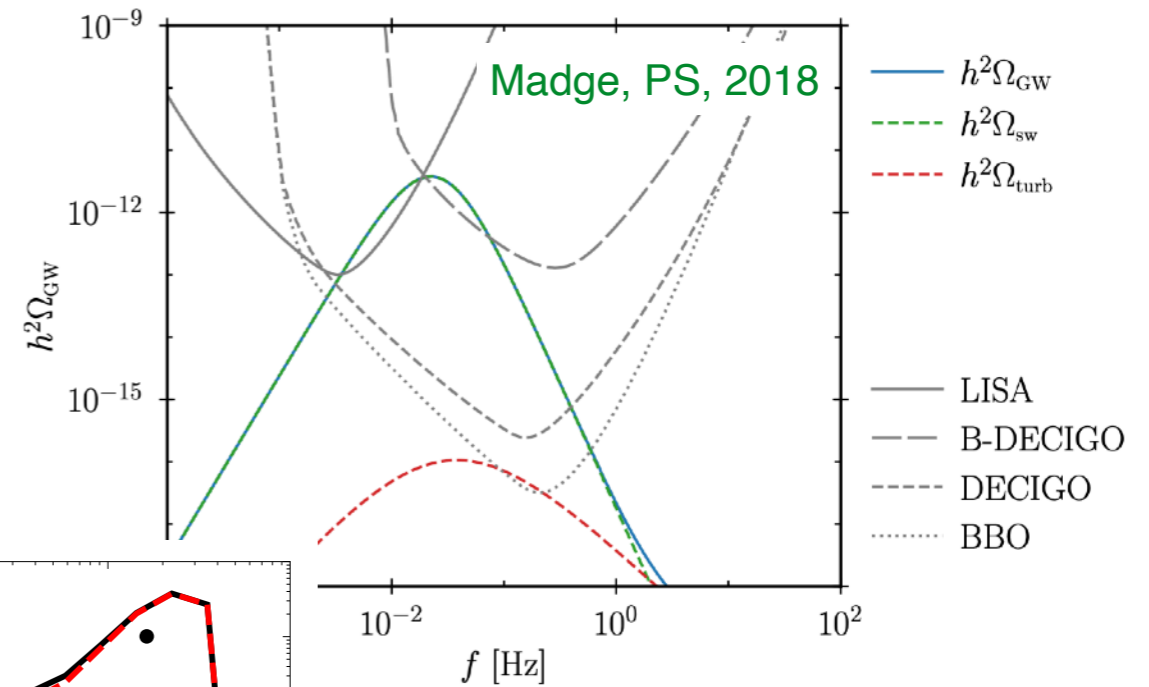
Not today:

- ▶ Ultra-high frequency GWs
- ▶ Supercooled PTs
- ▶ Strongly coupled PTs from holography

GW sources and their signal shapes

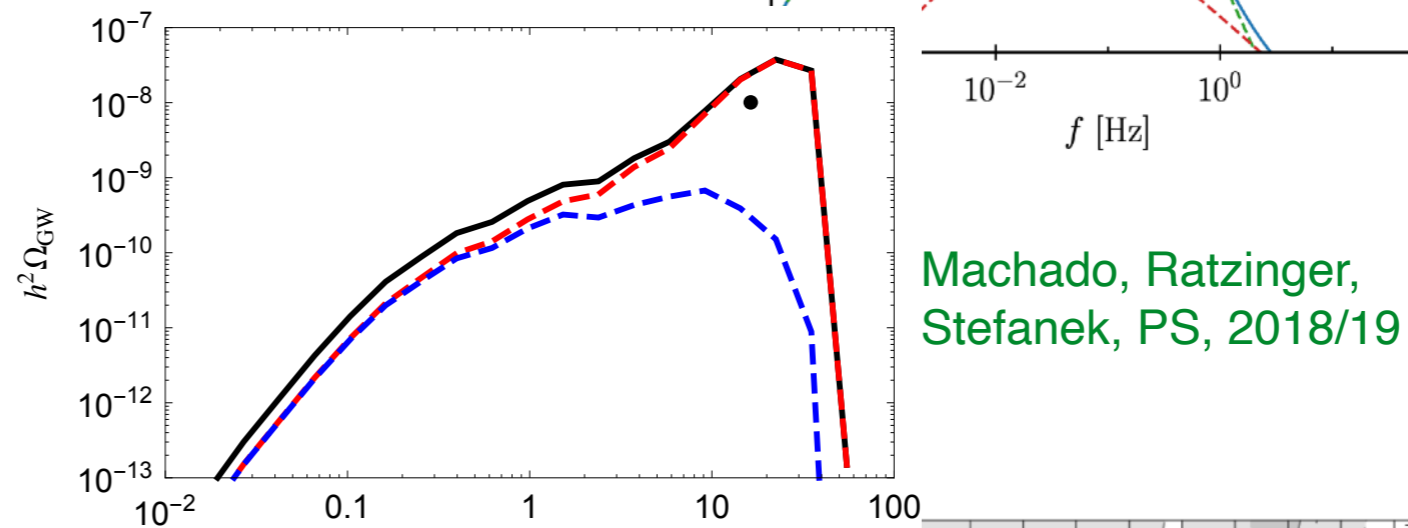
Phase transition

- ▶ Peak position depends on critical temperature



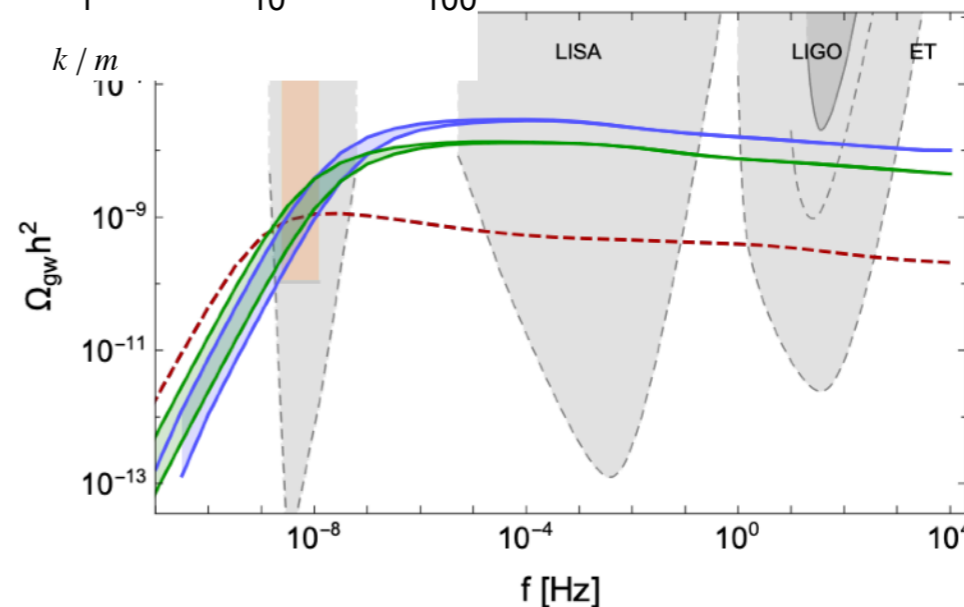
Audible axions:

- ▶ Peaked but chiral



Cosmic strings

- ▶ Flatter spectrum



Audible Axioms

Axion/ALP with dark photon

Take an axion ϕ

Dark photon X

$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{\alpha}{4f} \phi X_{\mu\nu} \tilde{X}^{\mu\nu} \right]$$

In radiation domination,
i.e. after inflation!

Axion/ALP with dark photon

Take an axion ϕ

Dark photon X

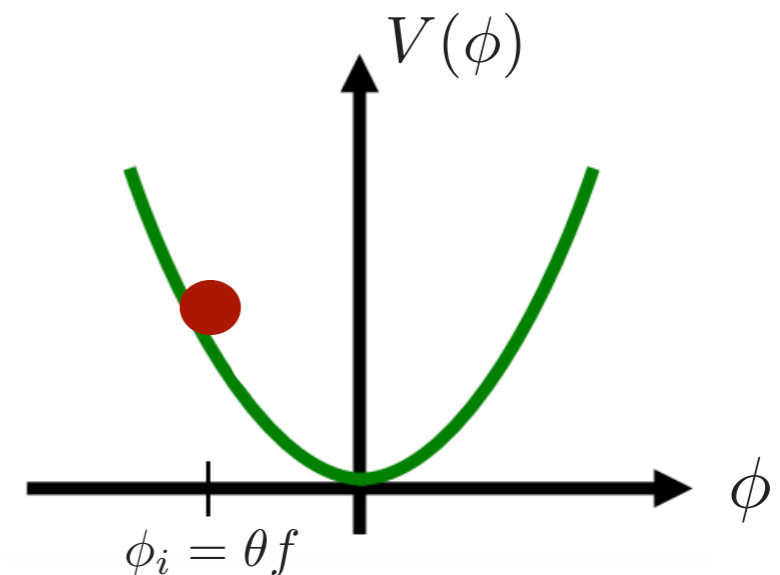
$$\mathcal{S} = \int d^4x \sqrt{-g} \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{\alpha}{4f} \phi X_{\mu\nu} \tilde{X}^{\mu\nu} \right]$$

In radiation domination,
i.e. after inflation!

$$V(\phi) = m^2 f^2 \left[1 - \cos\left(\frac{\phi}{f}\right) \right]$$

Initial ϕ_i , starts rolling when

$$H \sim m_\phi$$

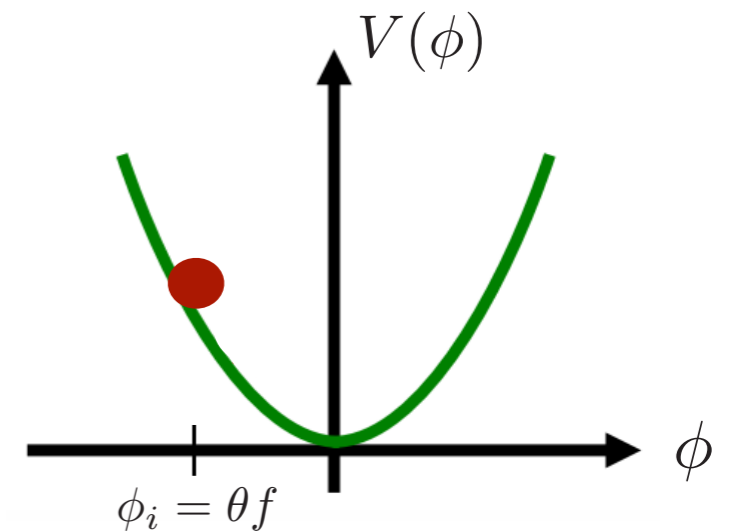


ALP dynamics - with dark photon

Equation of motion

$$\phi'' + 2aH\phi' + a^2V'(\phi)$$

$$\cancel{-\nabla^2\phi} - \frac{\alpha}{fa^2}\mathbf{X}' \cdot (\nabla \times \mathbf{X}) = 0$$



ALP starts rolling when $H \sim m_\phi$

ALP is damped due to exponential production of dark photons

- ▶ Reduced relic abundance - enlarge natural DM parameter space
- ▶ Or production of vector DM

Agrawal, Marques-Tavares, Xue, 2018
And others...

How does this work?

Equation of motion (in momentum space)

$$X''_{\pm}(\tau, \mathbf{k}) + \left(k^2 \pm k \frac{\alpha}{f} \phi'(\tau) \right) X_{\pm}(\tau, \mathbf{k}) = 0$$

The rolling ALP induces a tachyonic instability

$$X''_{\pm} + \omega_{\pm}(\tau) X_{\pm} = 0 \quad \text{with} \quad \omega_{\pm} = k^2 \mp k \frac{\alpha}{f} \phi'$$

Exponential growth of a range of dark photon modes

$$X(\tau) \propto e^{|\omega|\tau} \quad \text{for} \quad k \sim \frac{\alpha \phi'}{2f}$$

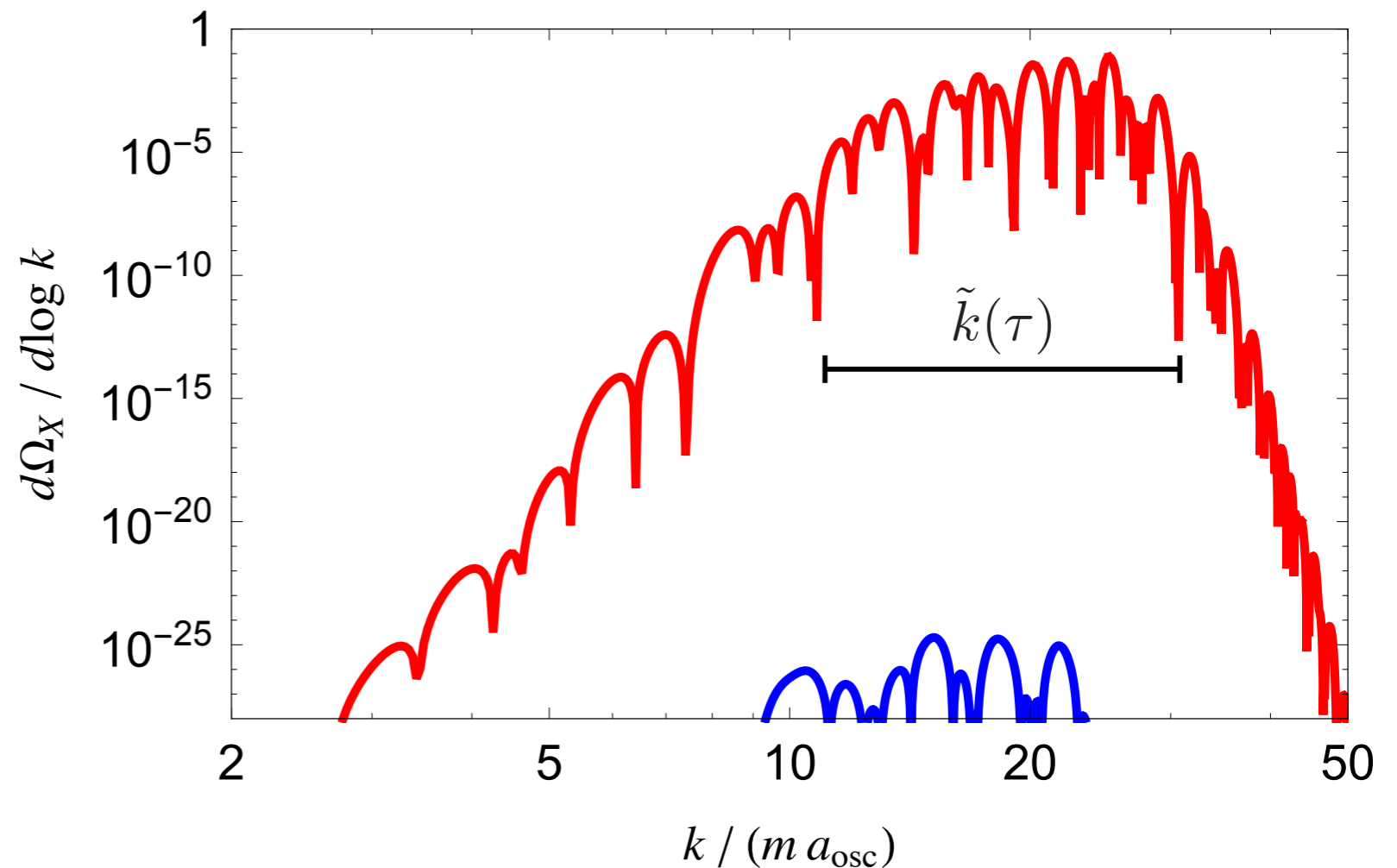
Dark photon spectrum

Initial condition violates parity (field rolls to the left or to the right)

One dark photon helicity dominates

A certain range of modes undergoes growth

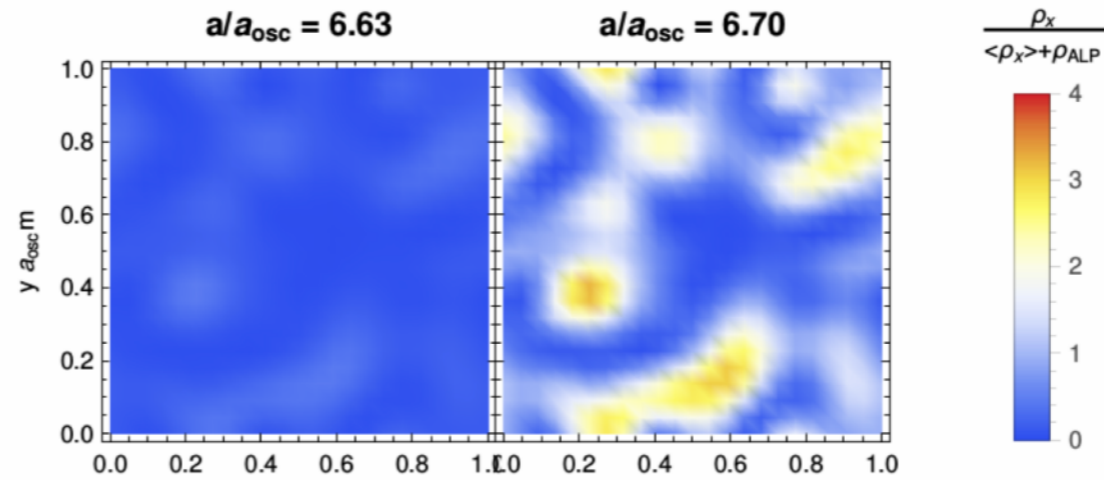
$$0 < k < \frac{\alpha\phi'}{f}, \quad \frac{k}{m} \lesssim \alpha\theta$$



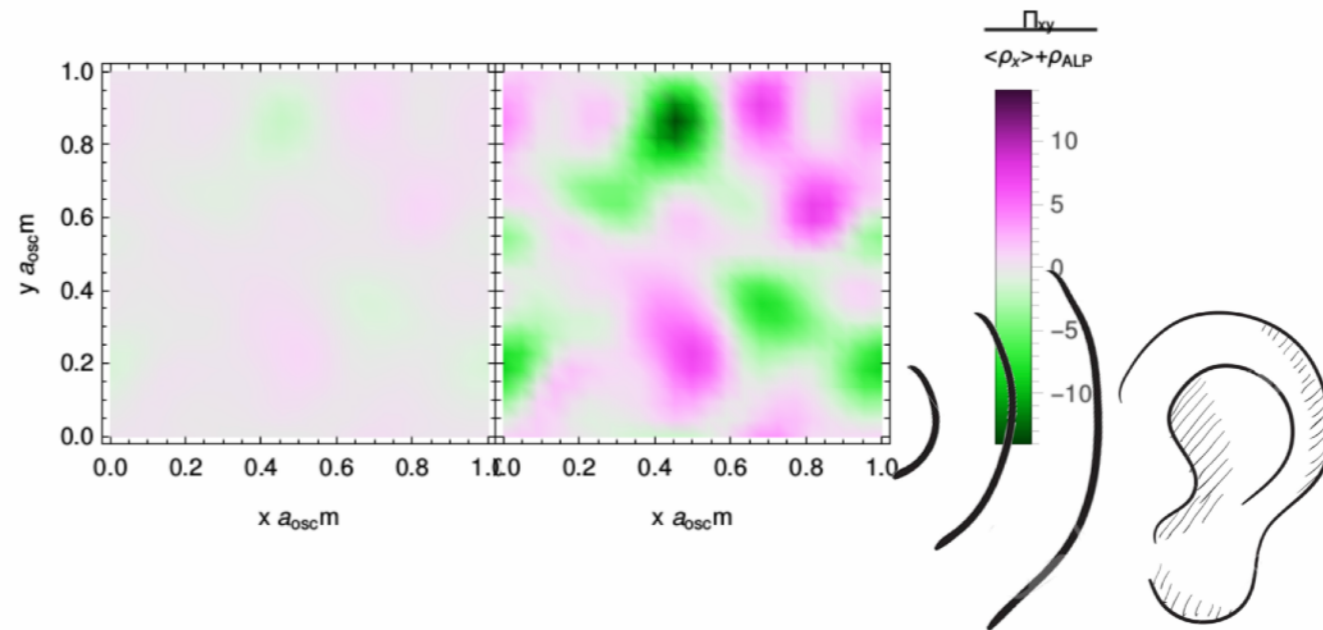
Machado, Ratzinger, Stefanek, PS, 1811.01950

GW production

Energy Density
of Dark Photon



Anisotropic
Stress



Gravity Waves

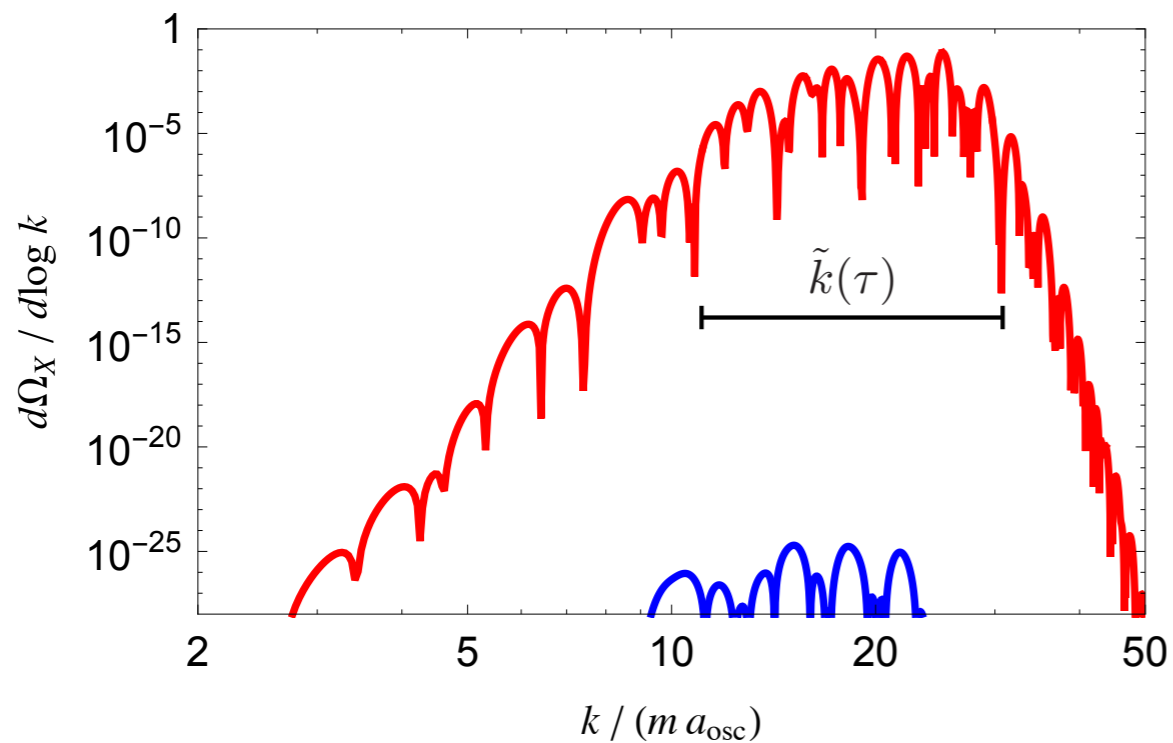
$$h''_{ij}(\mathbf{k}, \tau) + k^2 h_{ij}(\mathbf{k}, \tau) = \frac{2}{M_P^2} \Pi_{ij}(\mathbf{k}, \tau),$$

Anisotropic
stress

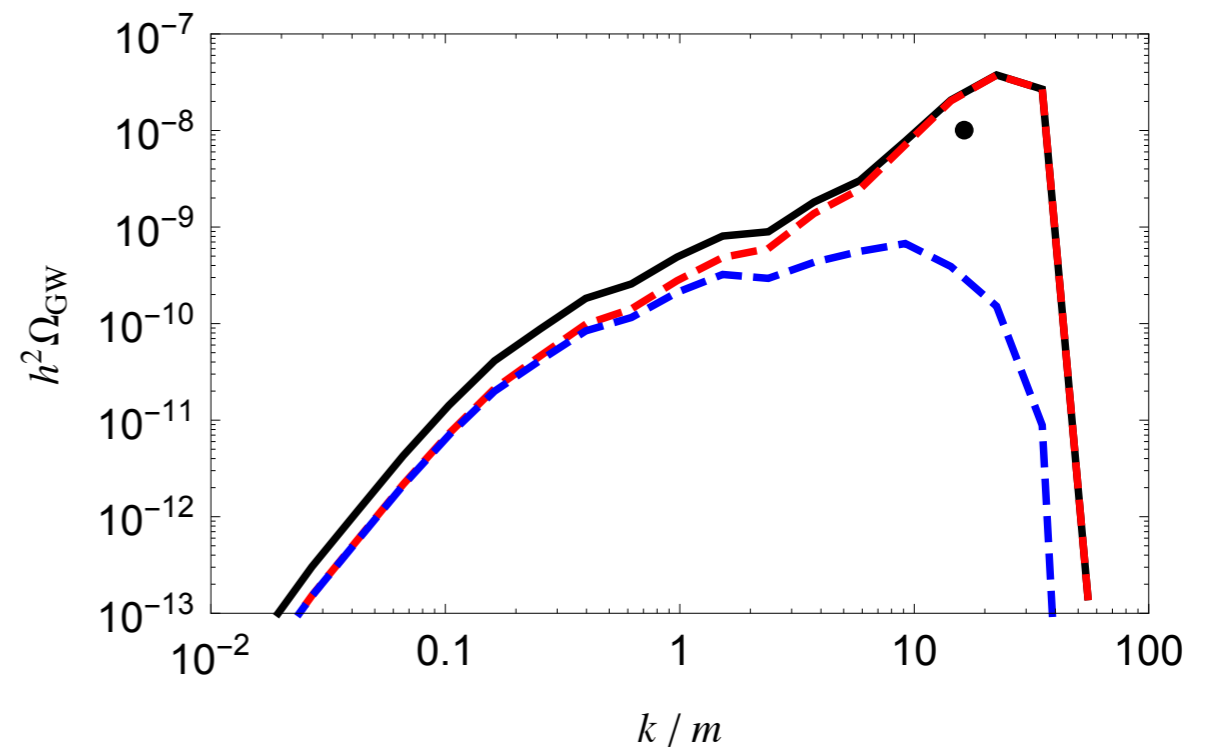
$$\hat{\Pi}_{ij}(\mathbf{k}, \tau) = \frac{\Lambda_{ij}^{kl}}{a^2} \int \frac{d^3 q}{(2\pi)^3} \left[\hat{E}_k(\mathbf{q}, \tau) \hat{E}_l(\mathbf{k} - \mathbf{q}, \tau) + \hat{B}_k(\mathbf{q}, \tau) \hat{B}_l(\mathbf{k} - \mathbf{q}, \tau) \right].$$

From dark photons to GWs

The exponential growth amplifies quantum fluctuations in the dark photon fields which source a **chiral** gravitational wave background



Dark photon spectrum

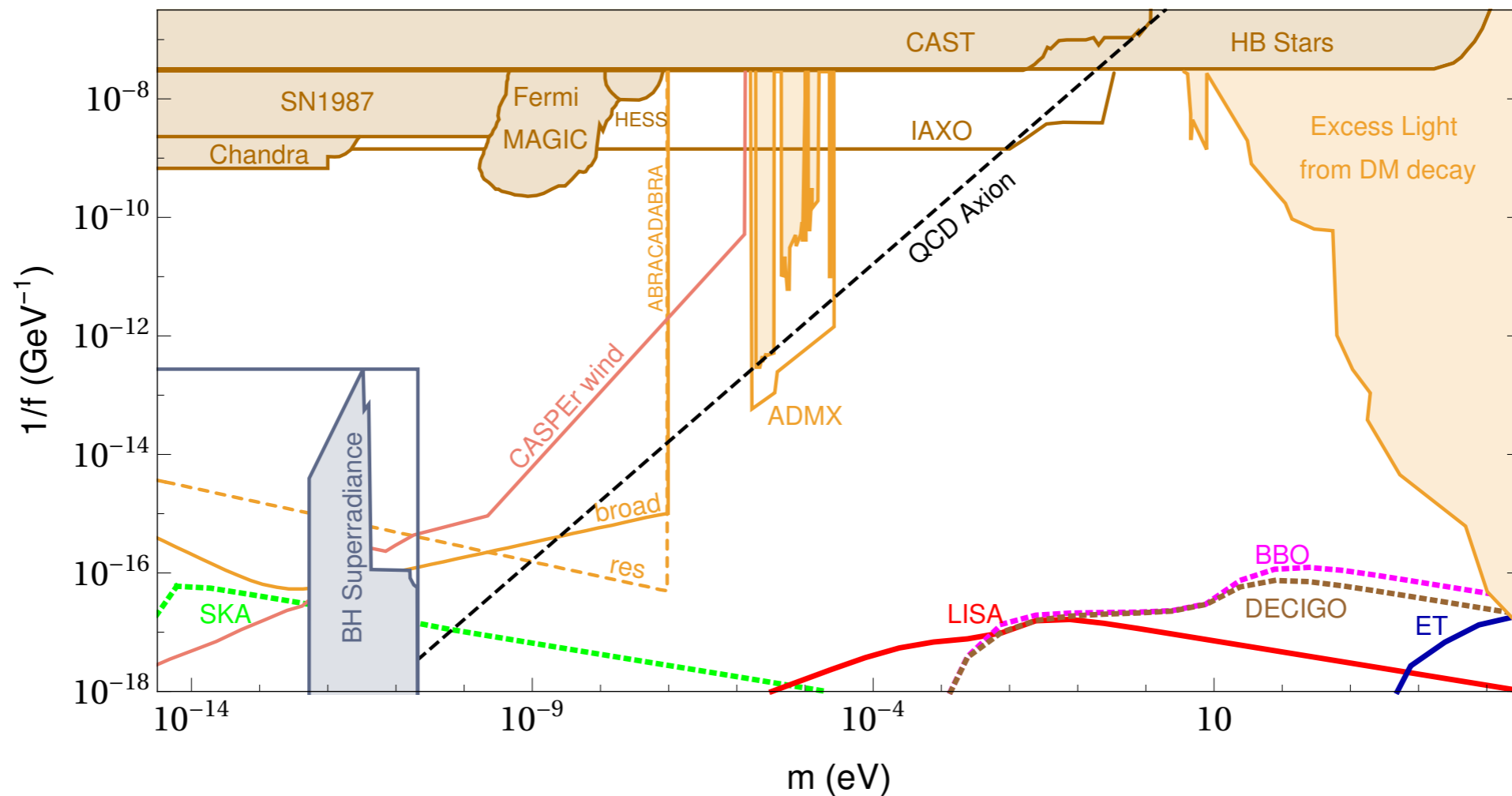


GW spectrum

Machado, Ratzinger, Stefanek, PS, 1811.01950

GW probes of audible ALPs

Machado, Ratzinger, Stefanek, PS, 1912.01107

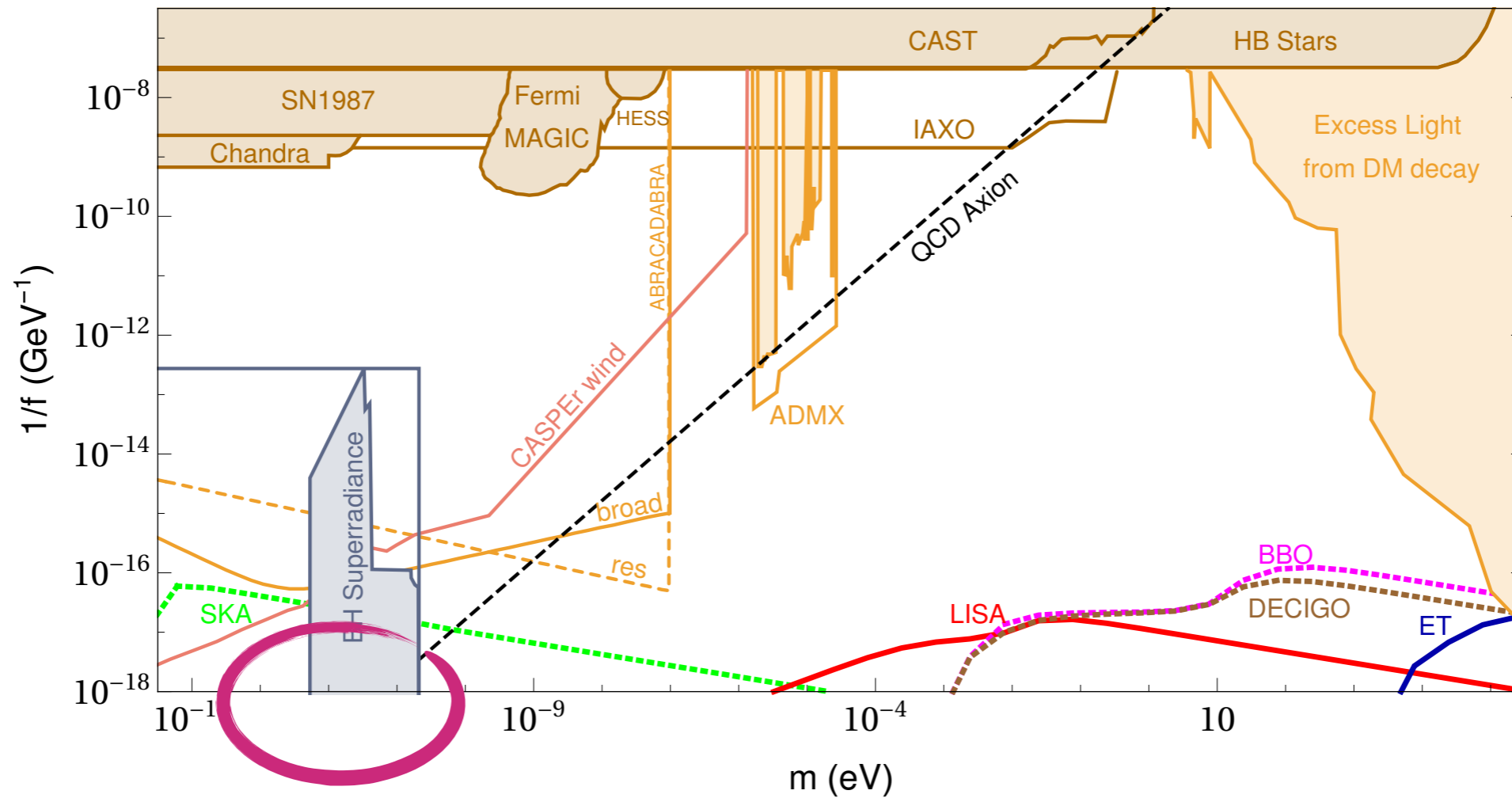


Mainly sensitive to high scale ALPs, since

$$f_0 \approx m \left(\frac{T_0}{T_*} \right) (\alpha\theta)^{2/3} = \sqrt{\frac{m}{M_P}} T_0 (\alpha\theta)^{2/3}, \quad \Omega_{\text{GW}}^0 \approx \Omega_\gamma^0 \left(\frac{f}{M_P} \right)^4 \left(\frac{\theta^2}{\alpha} \right)^{4/3}$$

GW probes of audible ALPs

Machado, Ratzinger, Stefanek, PS, 1912.01107

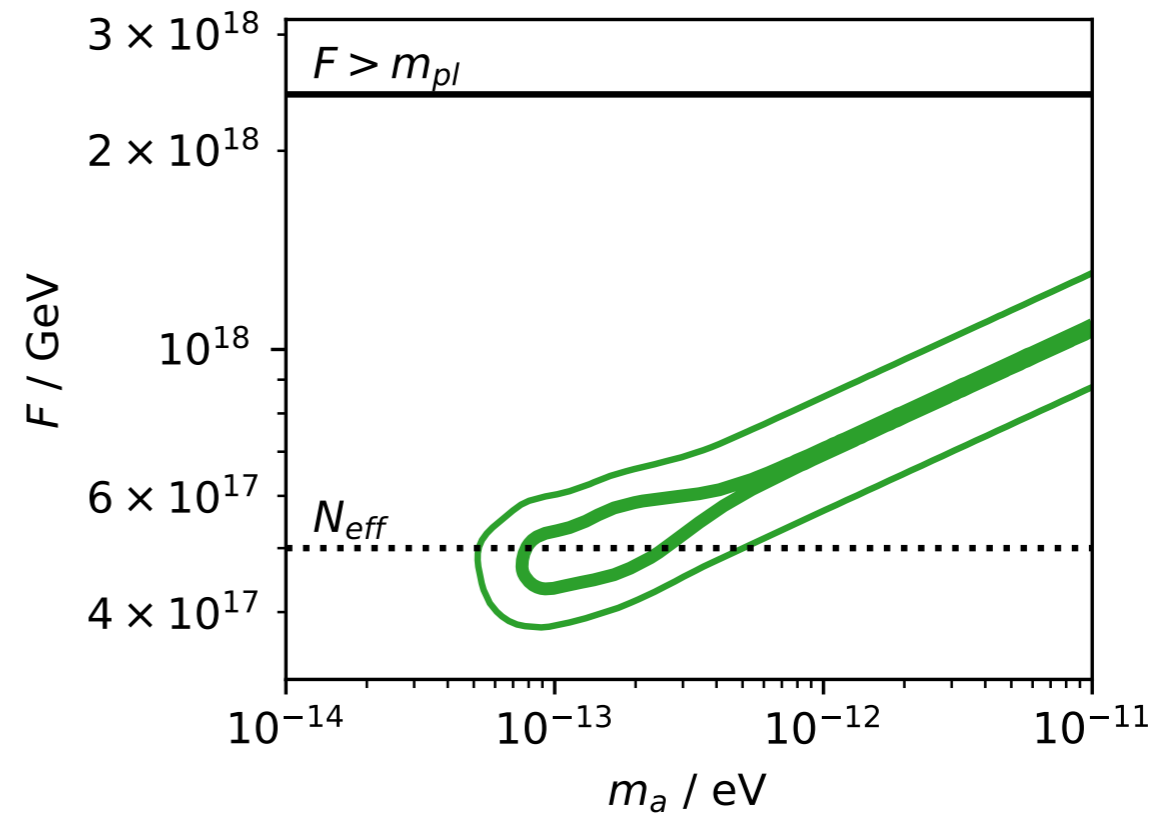


PTA region

Did PTAs hear the audible axion?

2020: Maybe

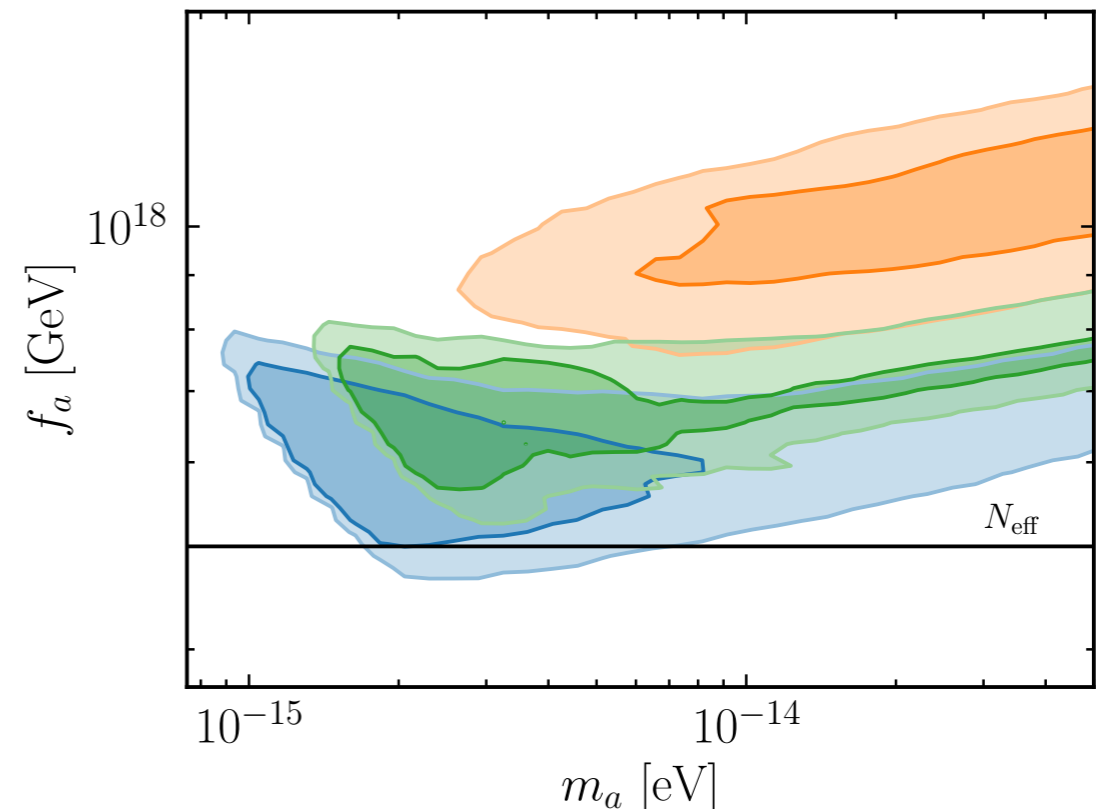
Wolfram Ratzinger & PS, 2009.11875



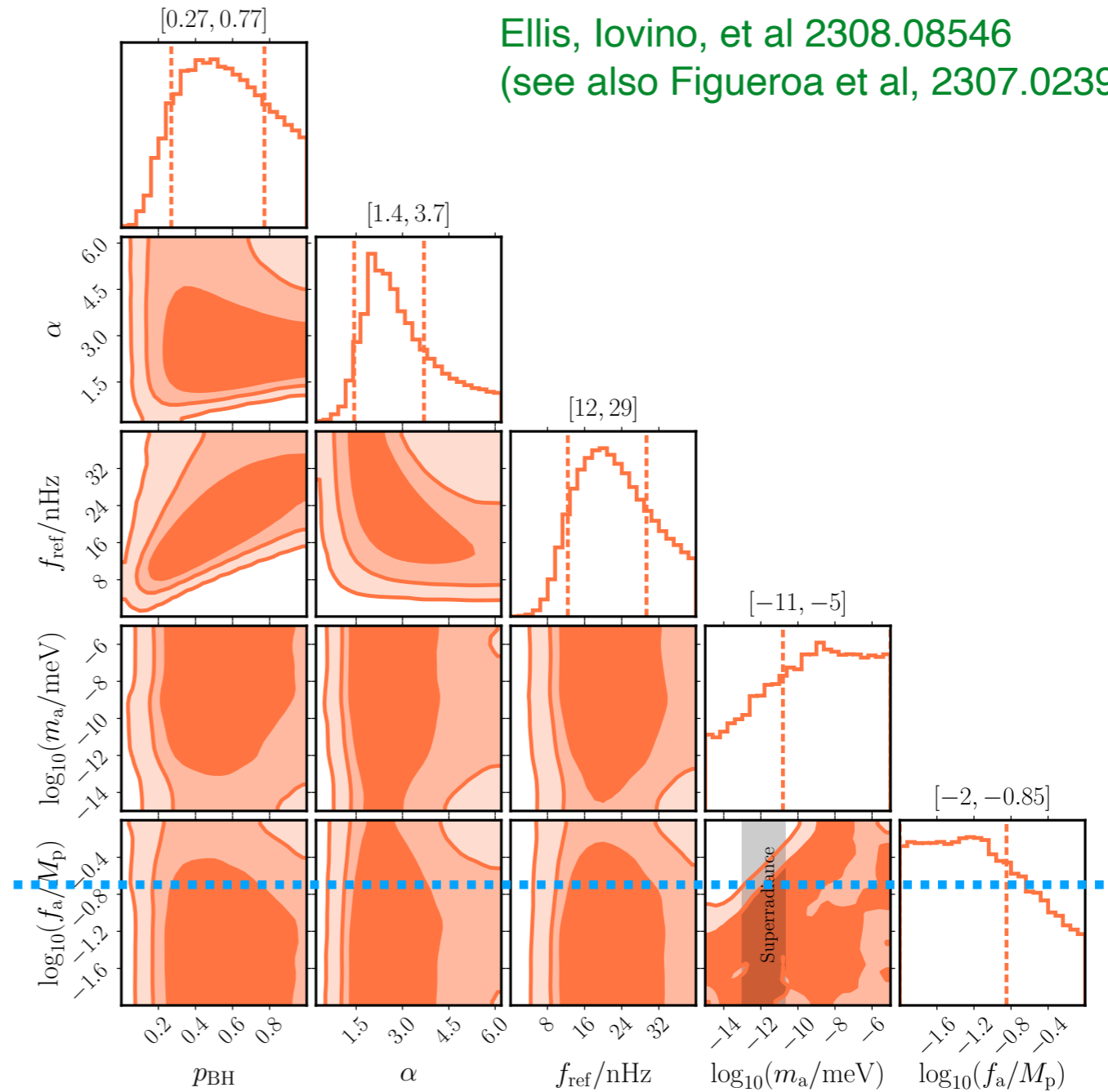
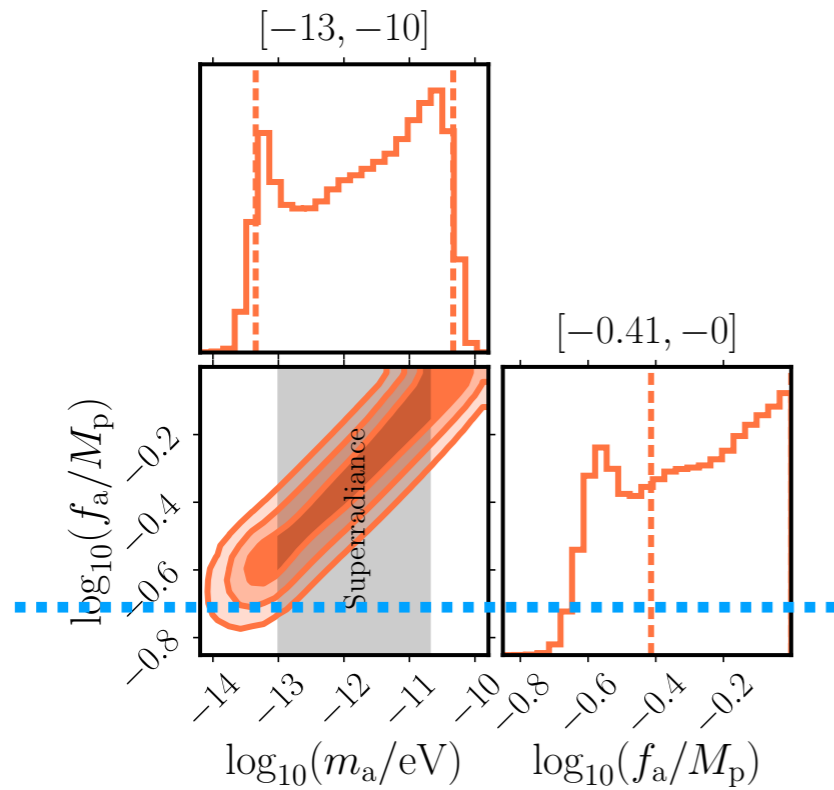
2023:

- ▶ Barely consistent with N_{eff}
- ▶ Not all of signal from AA
- ▶ OK since we also expect an astrophysical contribution

Madge et al,
[2306.14856](#)



Fits including SMBHB (from 2308.08546)



Fit moves to lower f_a

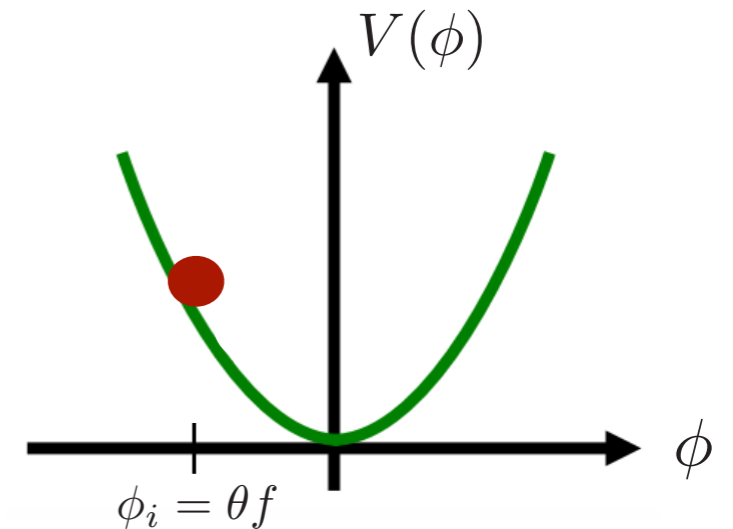
Consistent with N_{eff}

Using template w/o
backreaction

ALP dynamics - once more

Equation of motion

$$\begin{aligned} \phi'' + 2aH\phi' + a^2V'(\phi) \\ - \nabla^2\phi - \frac{\alpha}{fa^2}\mathbf{X}' \cdot (\nabla \times \mathbf{X}) = 0 \end{aligned}$$

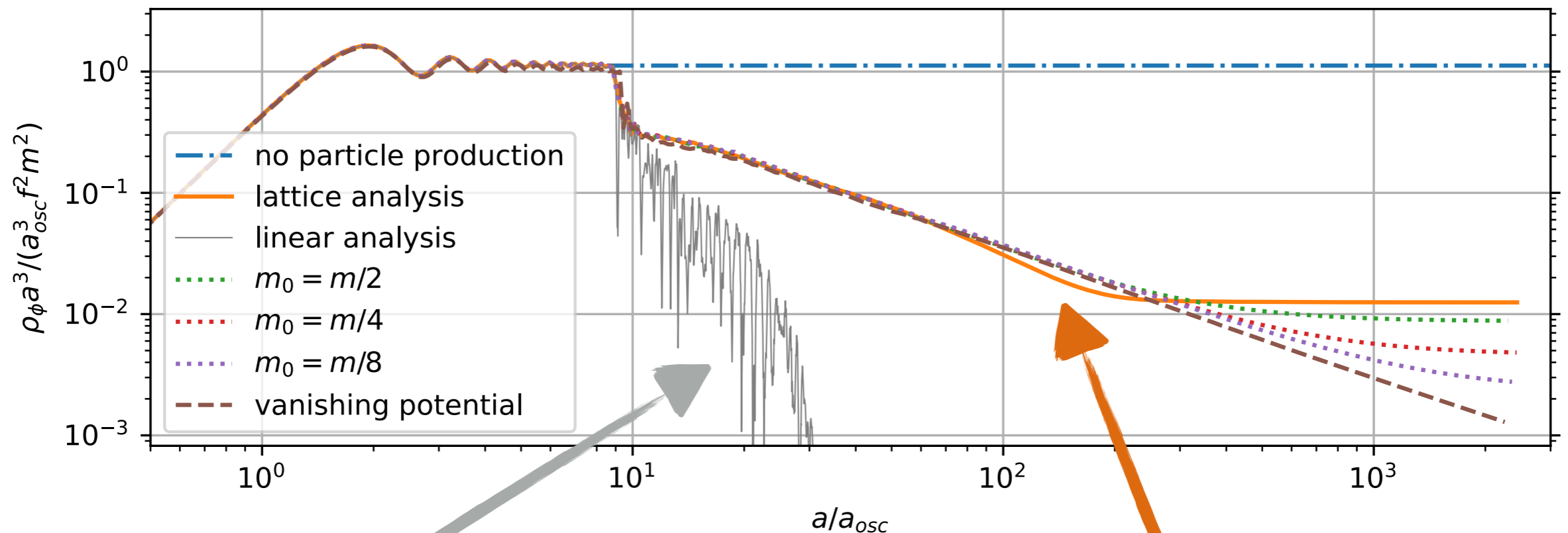


Once a significant population of dark photons is produced, the back-scattering into ALP fluctuations becomes non-negligible

Requires fully numerical treatment on the lattice

Important to get correct relic abundance prediction

From 2012.11584 with W. Ratzinger, B. Stefanek

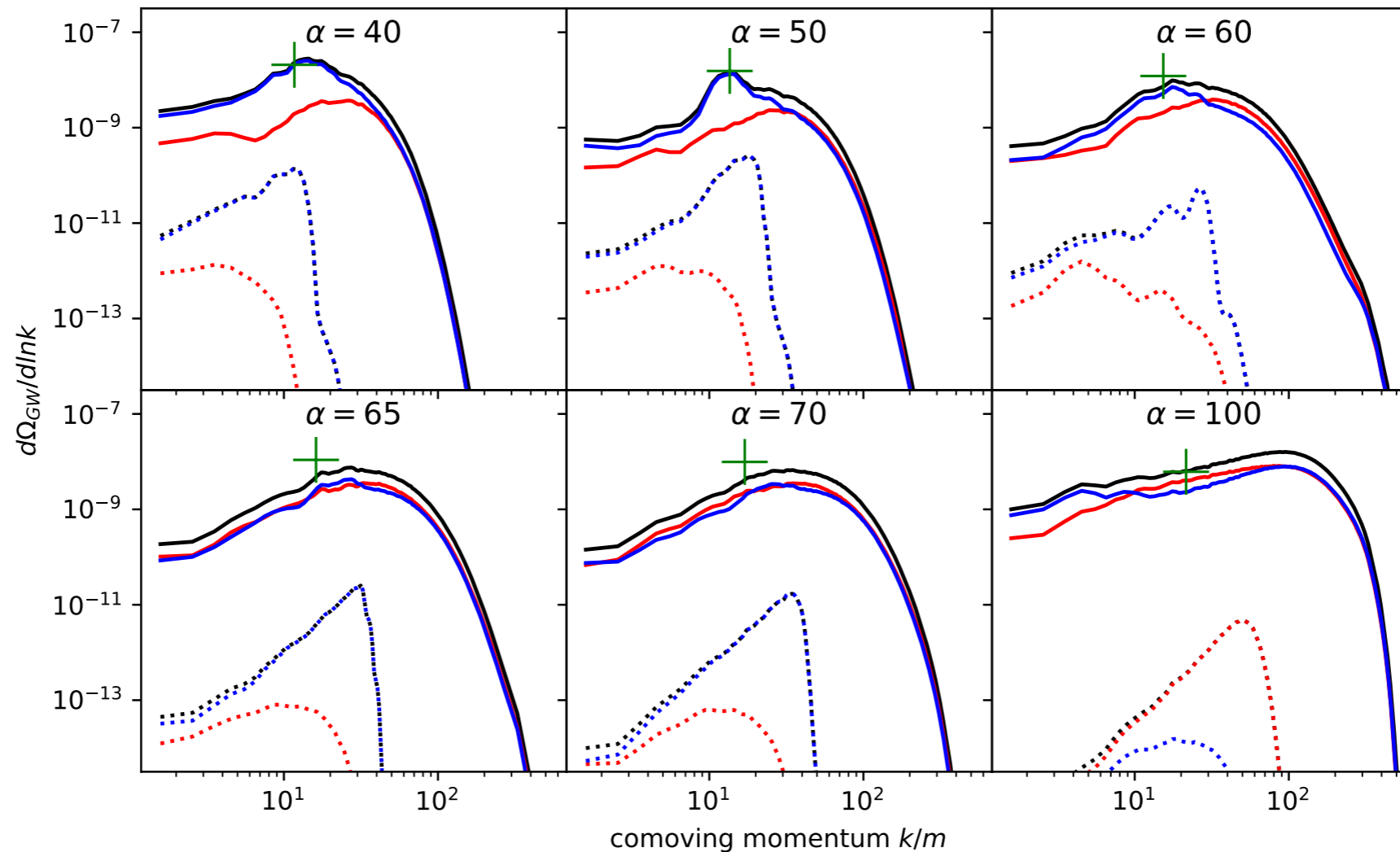


without back-scattering

Lattice result

See also Kitajima, Sekiguchi, Takahashi, 2018
Agrawal, Kitajima, Reece et al, 2020

Corrections to GW signal



Qualitative features unchanged, but polarisation is washed out at large couplings

From 2012.11584 with W. Ratzinger, B. Stefanek
see also 2010.10990 by (Kitajima, Soda, Urakawa)

Notes

Model variations: Audible Relaxion, Axion kinetic misalignment (see extra slides)

Also works for:

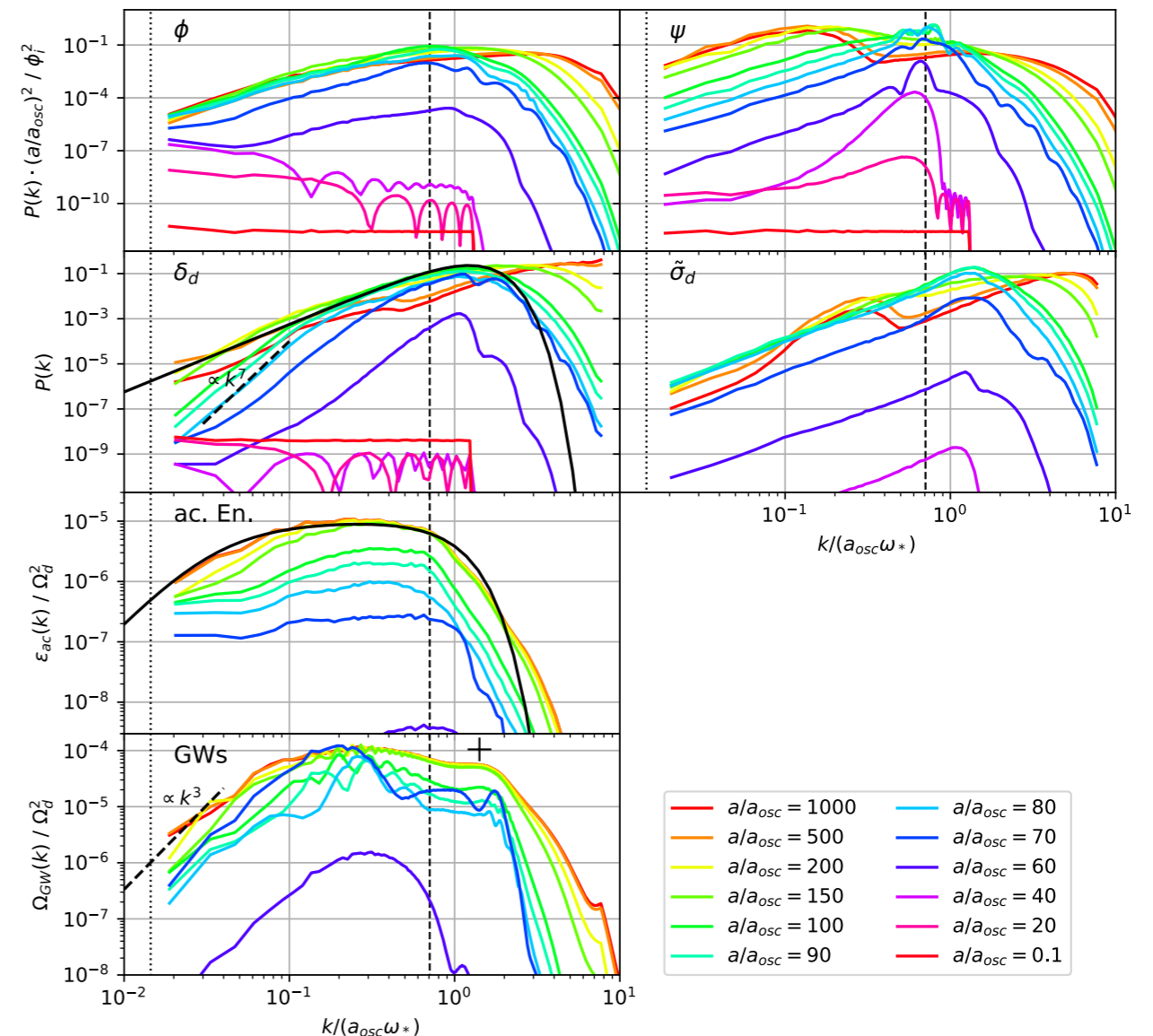
- Scalar dark sectors, e.g.

$$V(\phi, \psi) = \frac{1}{4} \lambda \phi^4 + \frac{1}{2} g^2 \phi^2 \psi^2$$

Ramberg, Ratzinger & PS, 2209.14313
see also Cui et al, 2310.13060

- single field models (e.g. axion fragmentation)

see e.g. Chatrchyan, Jaeckel, 2004.07844, Fonseca, Morgante, Sato, Servant, 1911.08472, ...

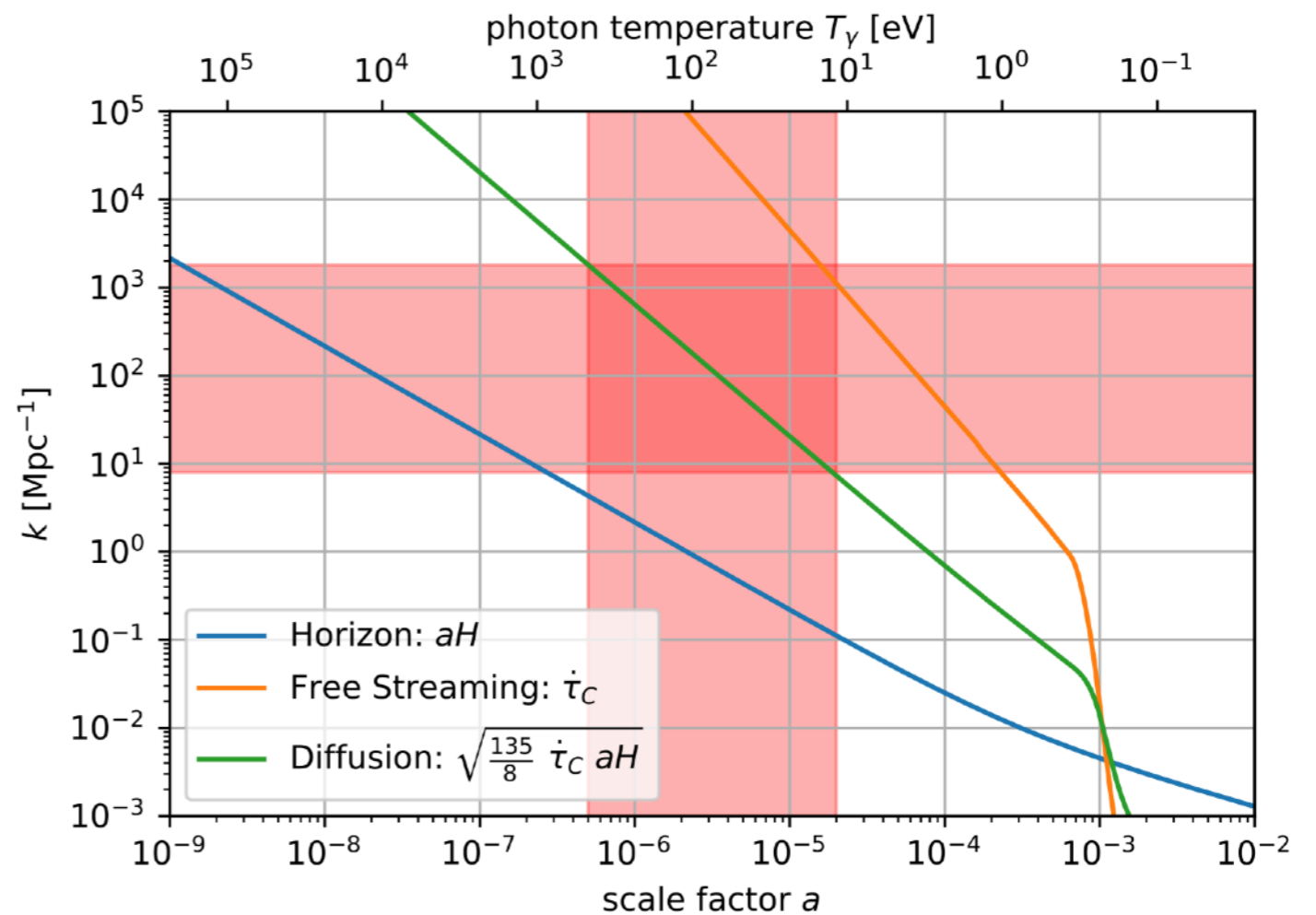


Spectral distortions?

Around $10^4 \lesssim z \lesssim 10^6$,
photon number is frozen

Any energy added to the
photons leads to a so
called μ distortion

Energy source we
consider here:
Gravitational damping of
dark sector fluctuations



Spectral distortions from dark sector anisotropies

Assume decoupled dark sector, $\Omega_d \ll 1$

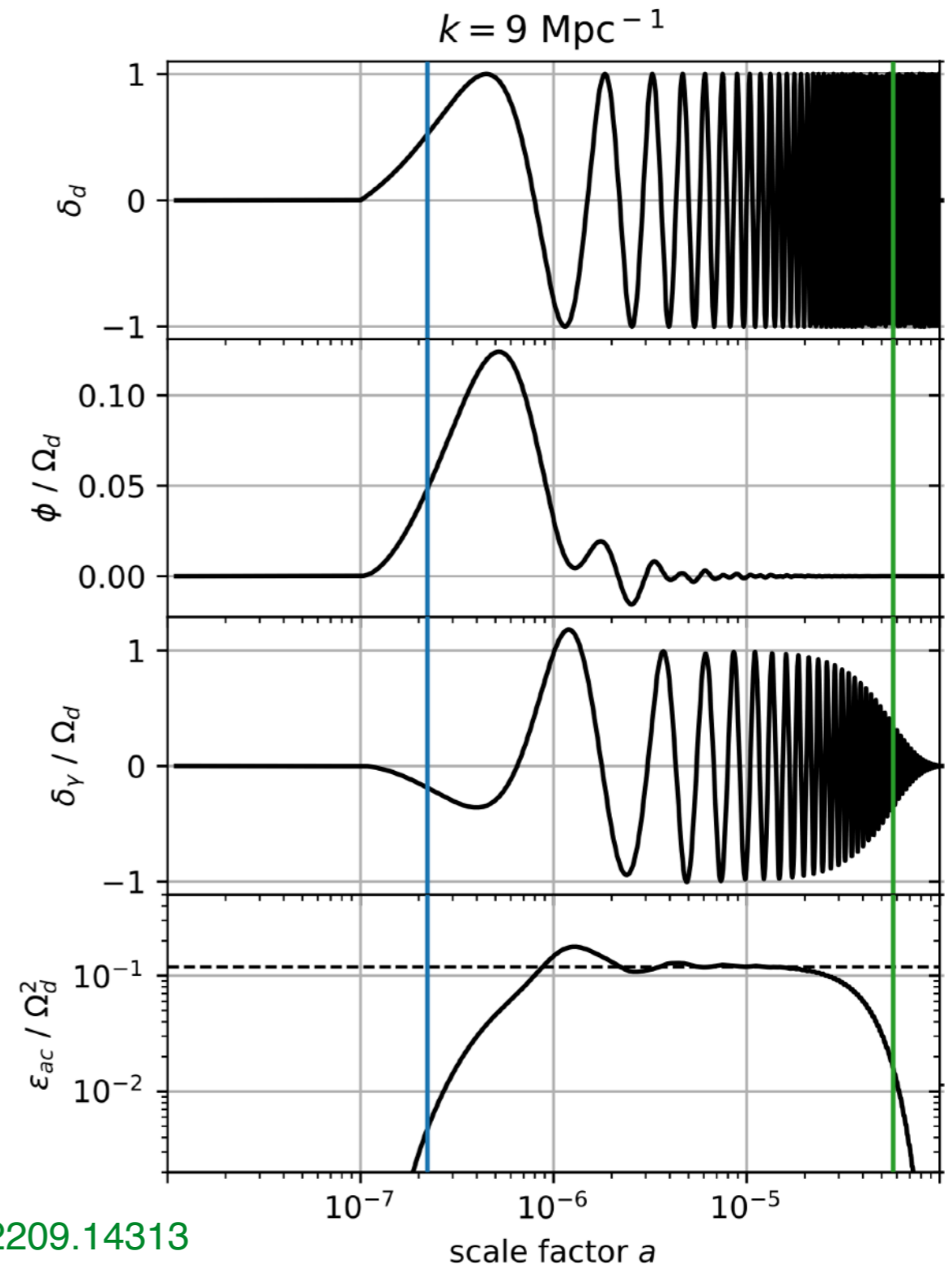
Large fluctuations

$$\delta_d = \delta\rho_d / \rho_d \sim 1$$

- ▶ Gravitationally induced sound waves in photons ϵ_{ac}

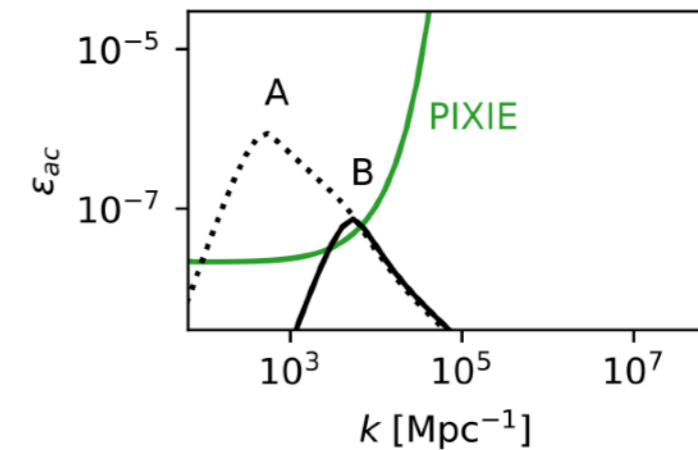
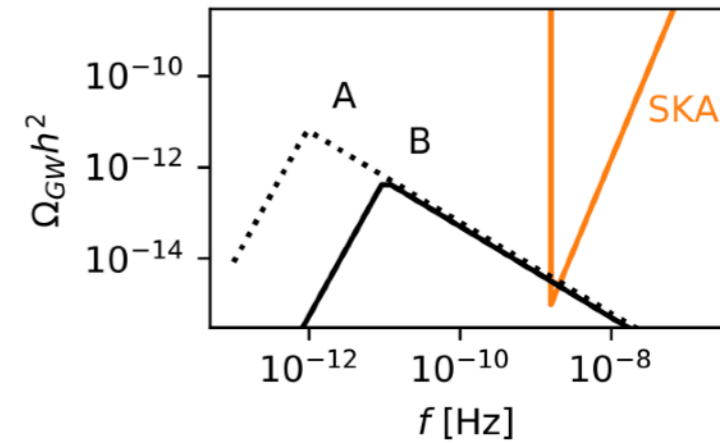
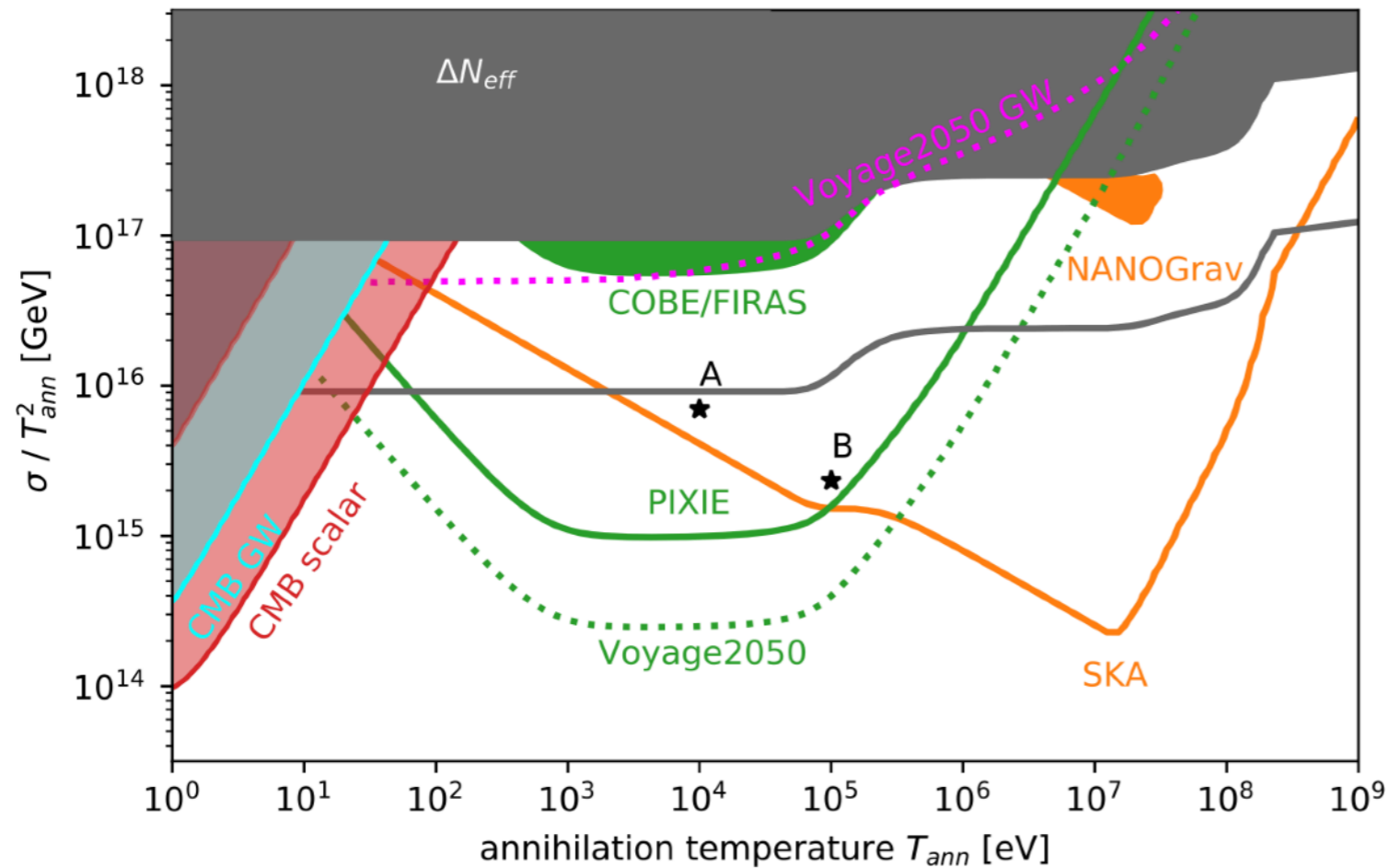
Resulting μ distortions

$$\mu = \int d \log k \epsilon_{ac}^{\text{lim}}(k) \mathcal{W}(k),$$



Ramberg, Ratzinger & PS, 2209.14313

Example source: Annihilating domain walls



Already probes allowed parameter space

Complementary to GW probes, can break degeneracy

- For all low scale sources (PTs, strings, AA,...)

Ramberg, Ratzinger & PS, 2209.14313

Axion/ALP domain walls

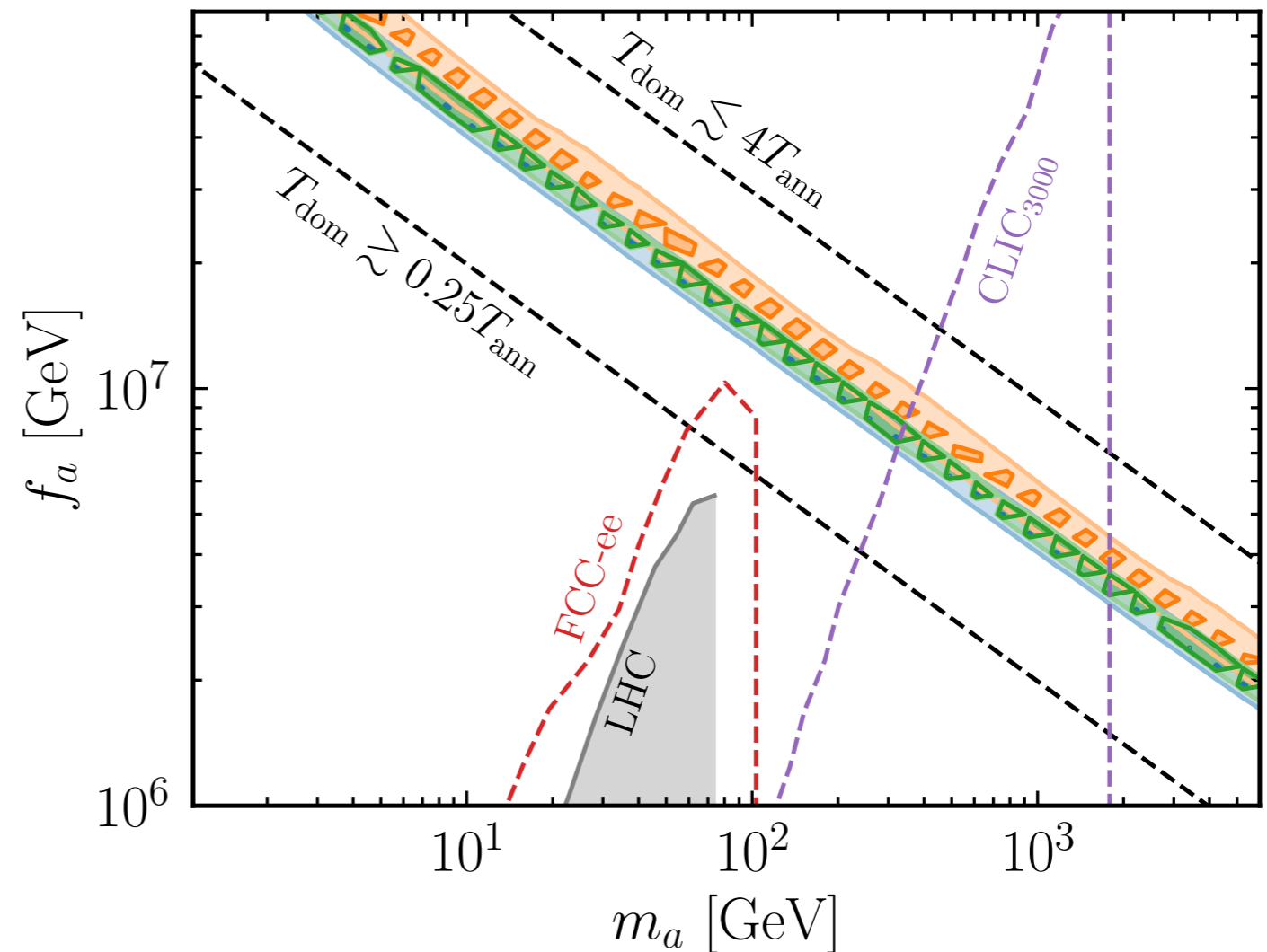
Madge et al,
[2306.14856](#)

Concrete model:
Aligned/clockwork
Axions [Higaki et al, 1606.05552](#)

Heavy axion
“partners” at weak
scale

In reach of future
colliders [Bauer et al, 1808.10323](#)

Visible, no spectral distortions



Invisibly decaying DWs

Madge et al,
[2306.14856](https://arxiv.org/abs/2306.14856)

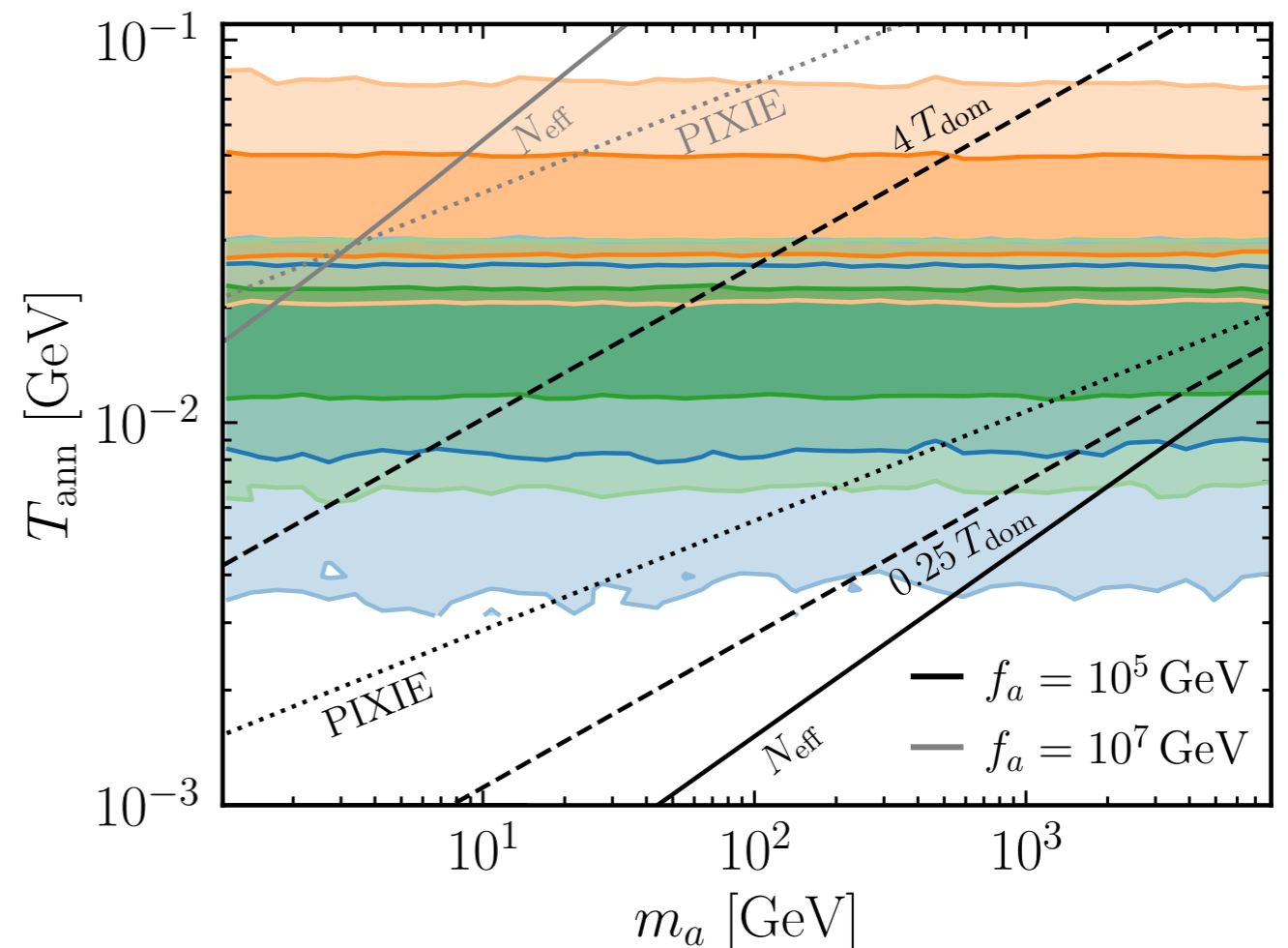
DWs annihilate to
dark radiation

- ▶ N_{eff} ok mostly

Dark sector anisotropies
induce CMB spectral
distortions

- ▶ In reach of future
experiments (PIXIE)

Also: PBH formation (Y. Gouttenoire 2023)



One more: Primordial black holes

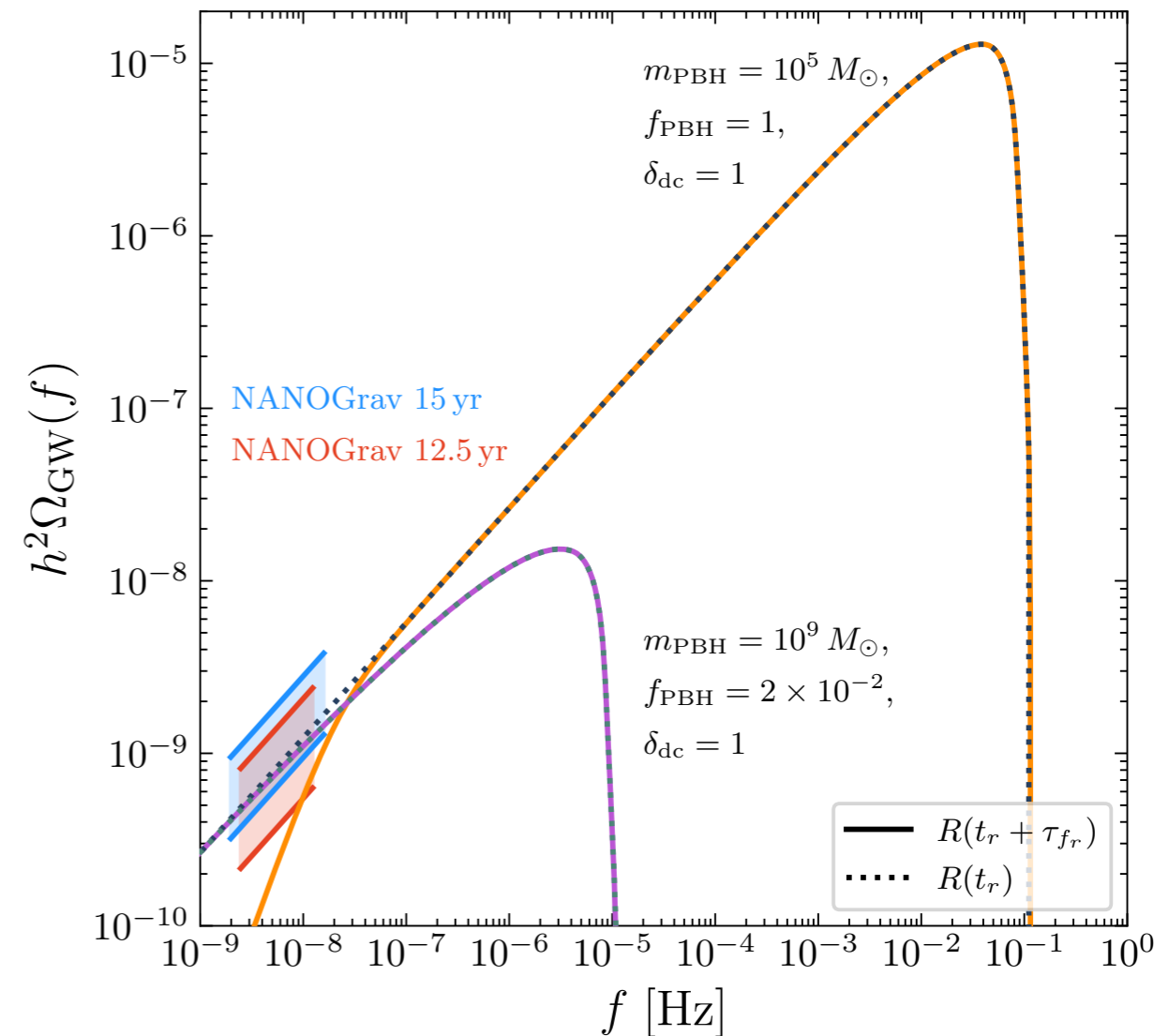
Binaries of supermassive PBH
produce a stochastic GWB

Parameters:

- ▶ Mass m_{PBH}
- ▶ DM fraction f_{PBH}
- ▶ Clustering δ_c

Approximations:

- ▶ Circular orbits, monochromatic mass distribution, no environmental effects



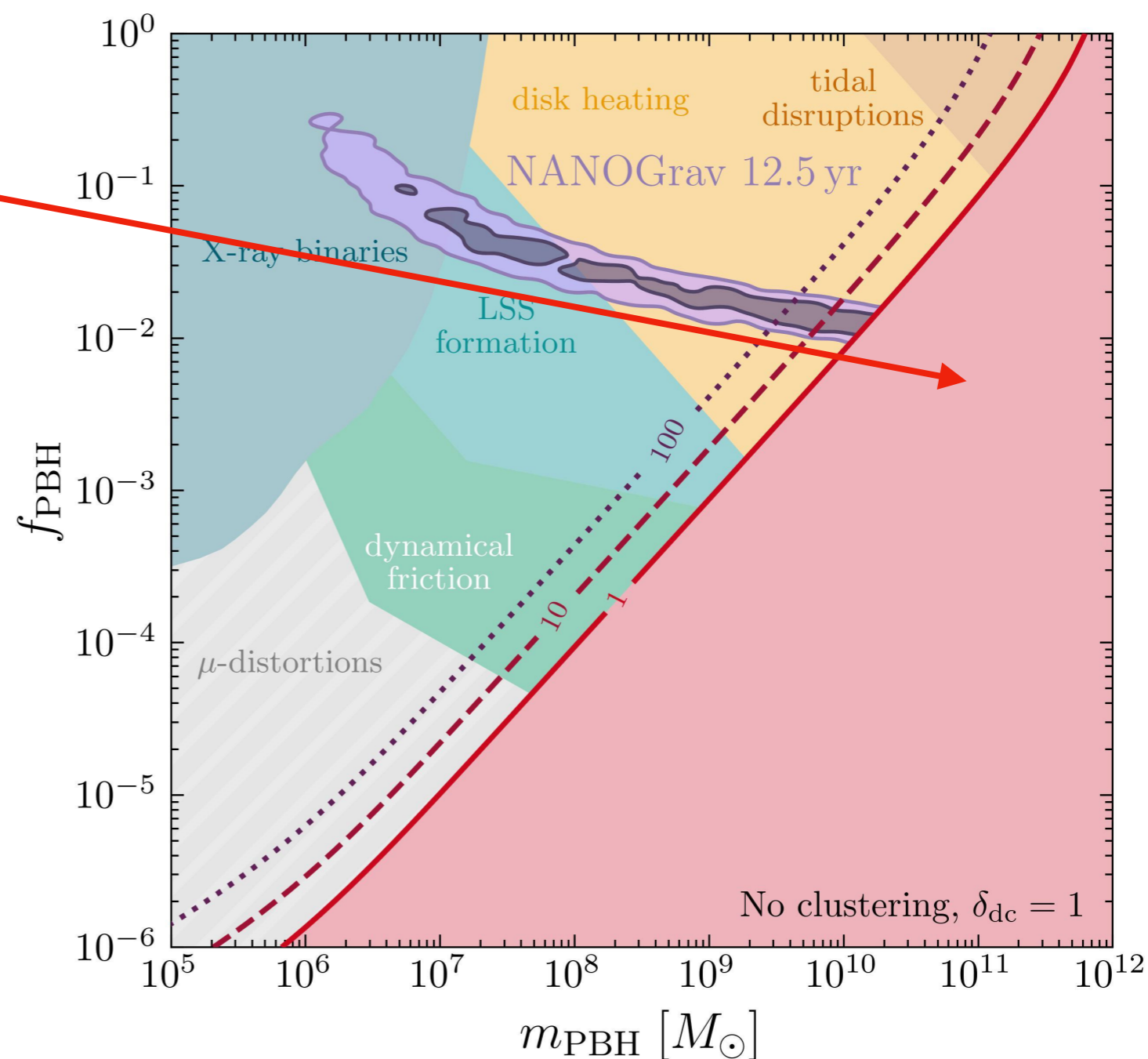
PBH: No clustering

Viable region at
very large masses

already pointed out by Atal, Sanglas,
Triantafyllou, 2012.14721

However: Fewer than
one pBH contributes
to signal on average
there

- ▶ Not a stochastic BG
- ▶ Not even a signal
most of the time



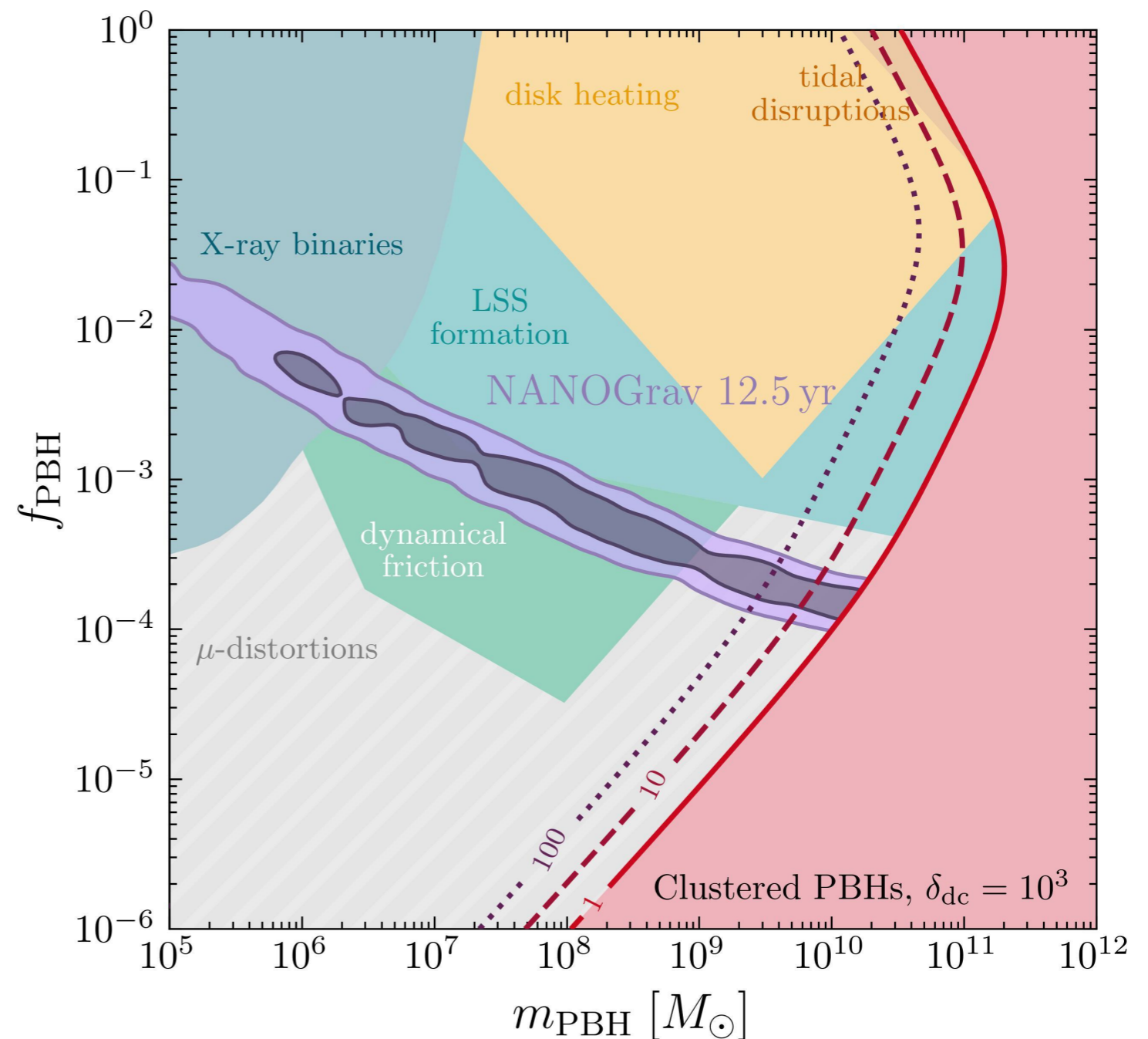
PBH: With clustering

Depta et al,
2306.17836

Now an actually viable region emerges

Assuming a suitable production mechanism

- ▶ Needs to evade mu distortion bounds
- ▶ Non-gaussian!

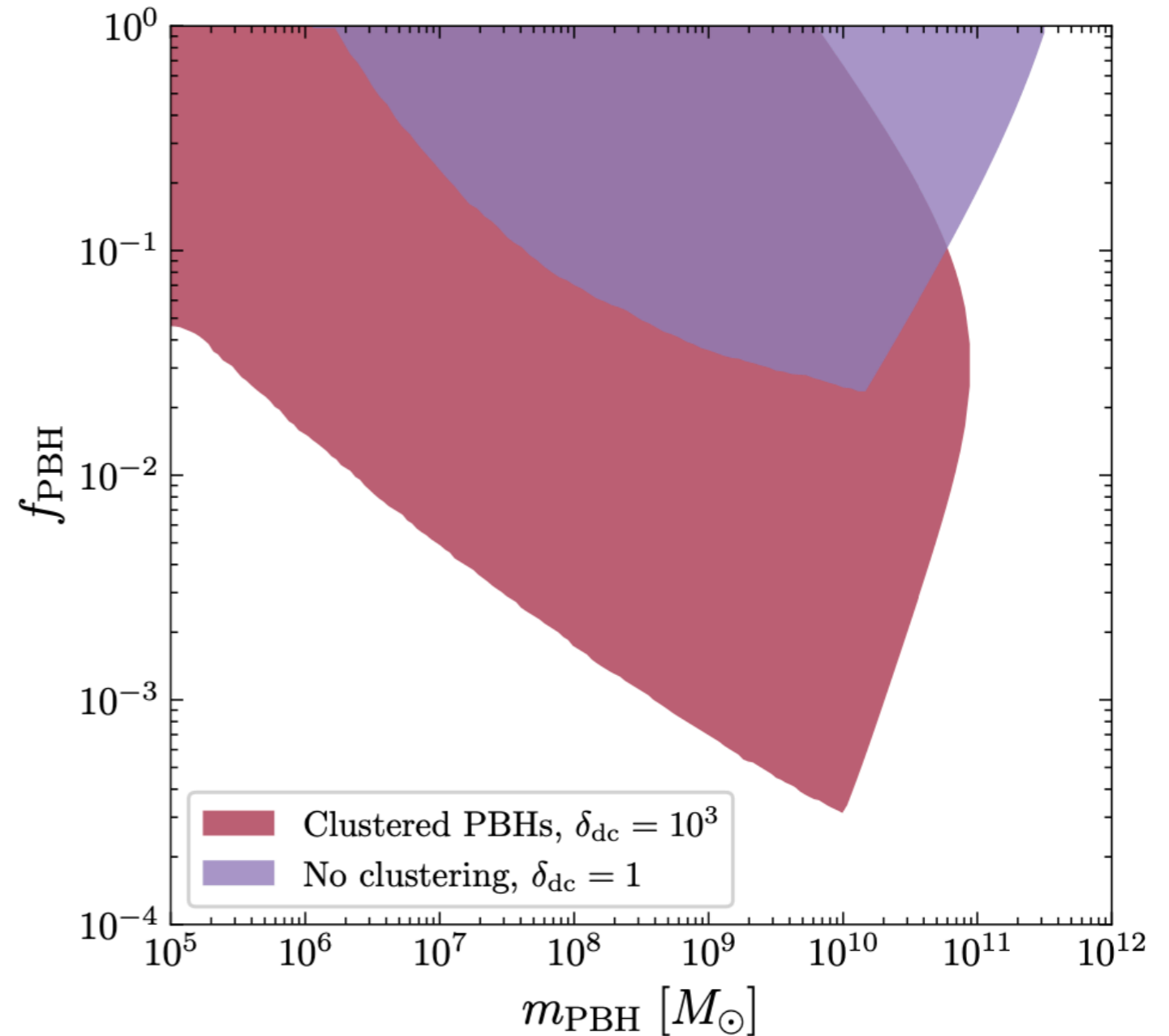


Superlarge PBH

Depta et al,
2306.17836

Less crazy: Astro-independent bounds on clustered pBH from PTAs

Expect anisotropies in GW background :)



Summary

GWs are new window to early, dark Universe

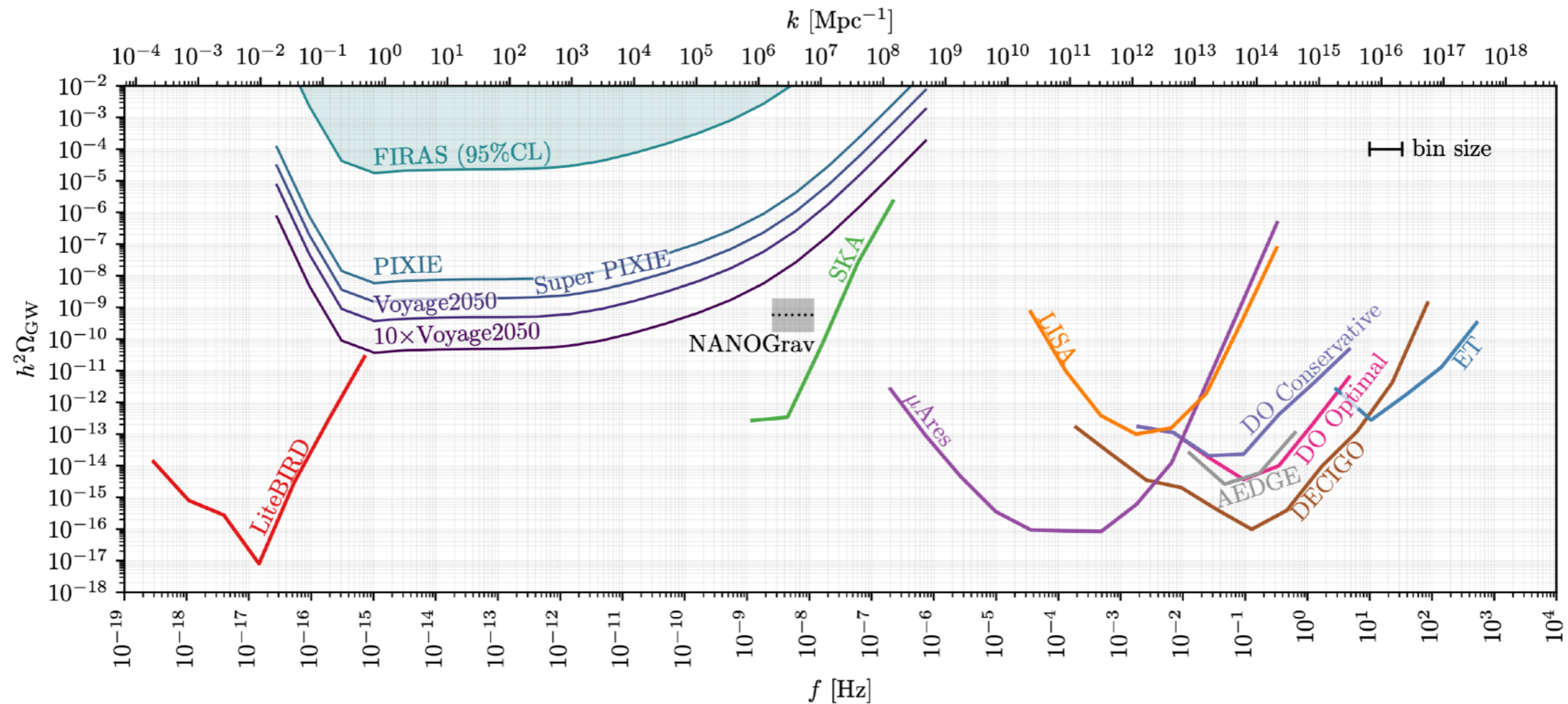
Today:

- ▶ Audible axions are cool
- ▶ Spectral distortions are cool
- ▶ Supermassive pBH are also cool, but maybe a bit crazy ;)

Many things to be done (simulations!), much data will come in the future -> Exciting times!

Extra slides :)

Spectral distortions as probes of low scale GWs



From Kite, Ravenni, Patil, Chluba, MNRAS 2021

Tensor fluctuations (GWs) also source μ distortions

- But difficult to test. Better to directly go for the scalar fluctuations (that also source the GWs)

High frequency GW searches

Higher Frequency \rightarrow shorter wavelength

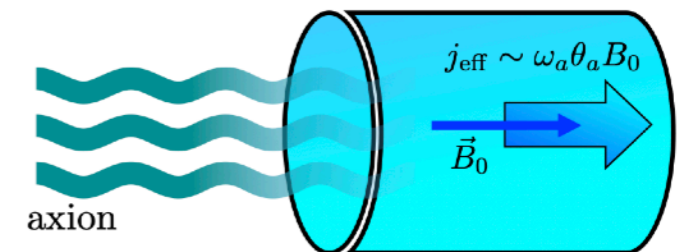
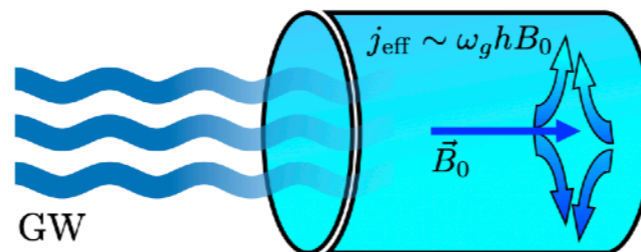
- ▶ Experiment may fit in your laboratory

Gravity couples to everything

- ▶ Any very sensitive device could potentially be a detector

Current interest:

- ▶ Cavities for axion searches
- ▶ Gertsenshtein effect:
GWs convert to photons in strong magnetic field



Berlin et al, 2112.11465

Sources? Primordial BH, superradiance, or...?

E&M on curved backgrounds is confusing however

E and B fields not uniquely defined everywhere in detector, depend on chosen coordinate frame

$$E_a = F_{a0}?$$

Observables should be independent!

$$\underline{E}_a = \hat{e}_a^\mu F_{\mu\nu} u^\nu !$$

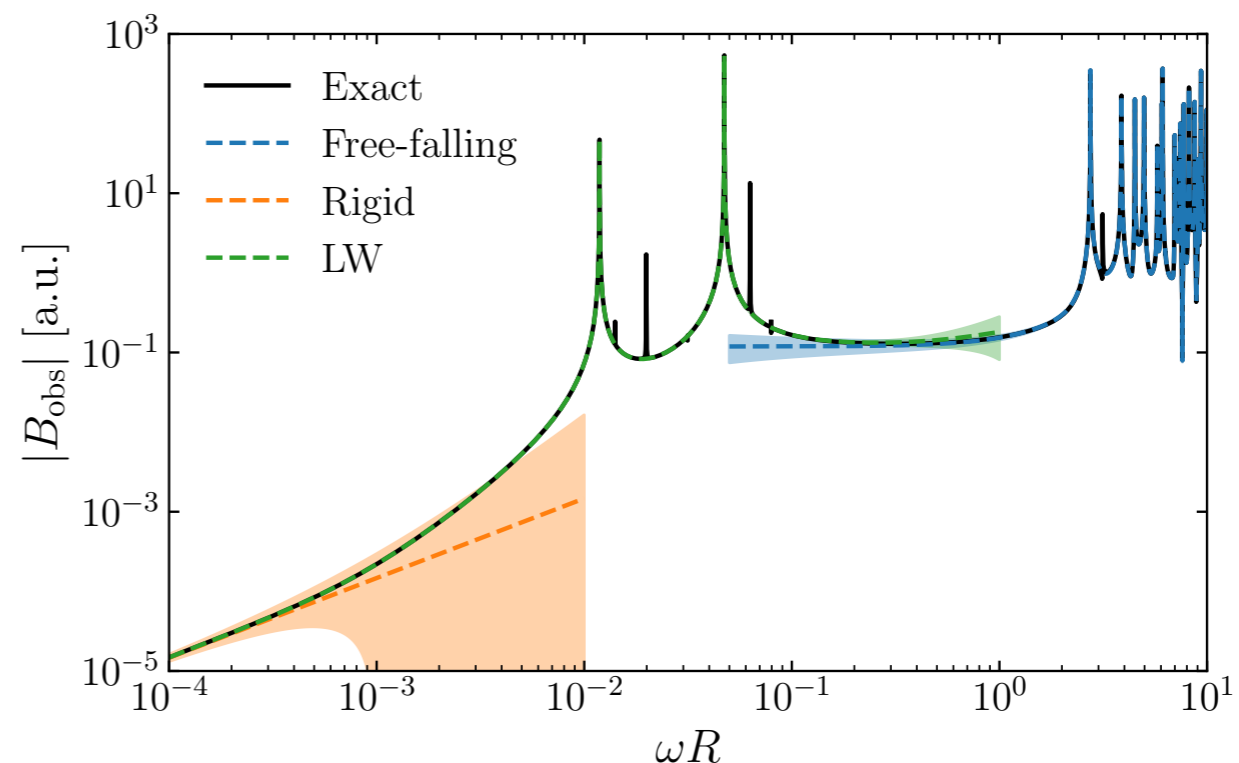
Proposed coordinate independent perturbation scheme

Applied to:

- ▶ Thin rod
- ▶ Sphere

Including mechanical deformations

Compared with commonly used approximations → can identify range of validity and provide error estimate



Wolfram Ratzinger, Sebastian Schenk, PS, 2404.08572

Audible relaxion

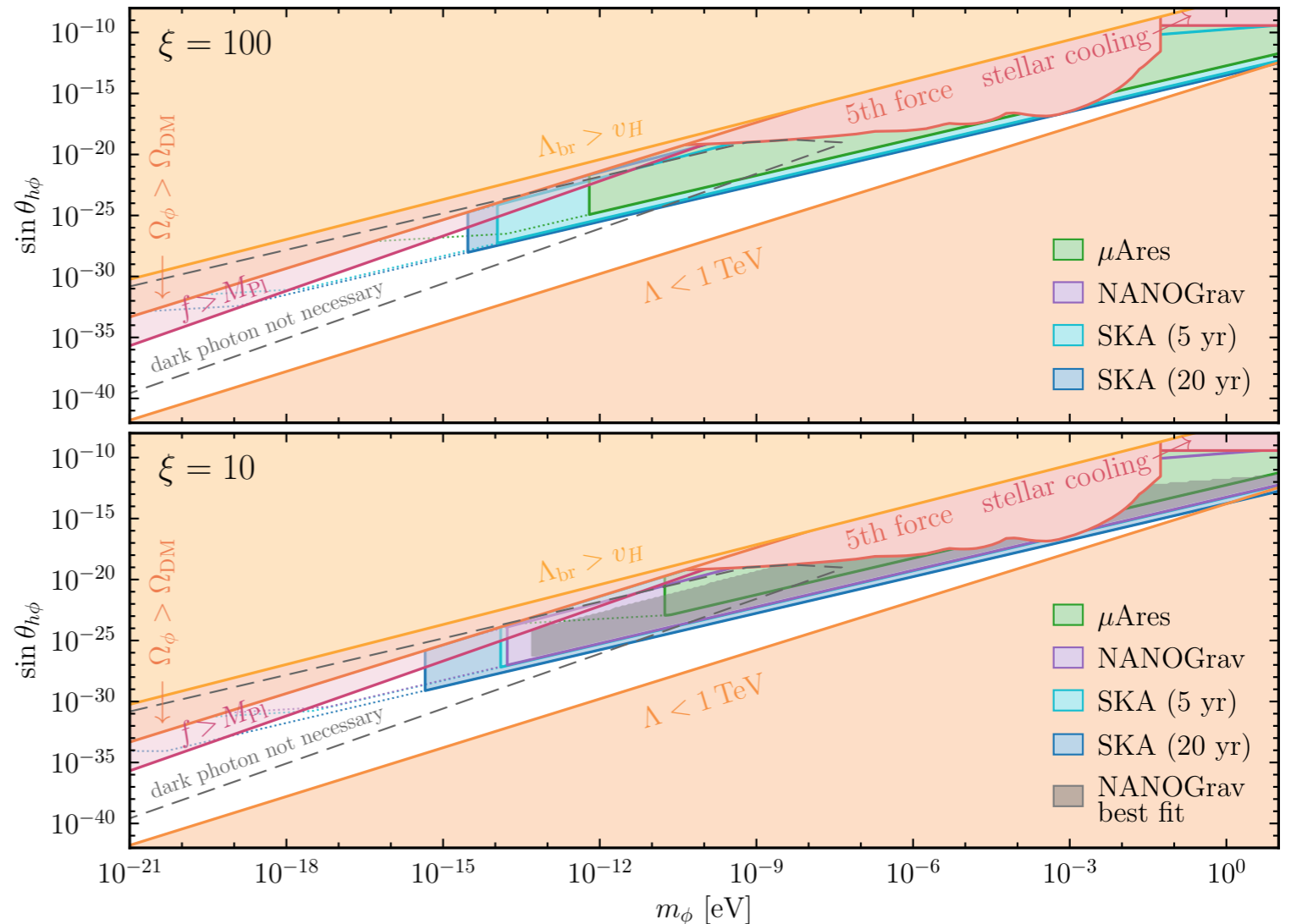
Audible relaxion

$$-\mathcal{L} \supset V(H, \phi) + \frac{r_X}{4} \frac{\phi}{f_\phi} X_{\mu\nu} \tilde{X}^{\mu\nu}$$

$$V(H, \phi) = V_{\text{roll}}(\phi) + \mu_H^2(\phi) |H|^2 + \lambda |H|^4 + V_{\text{br}}(H, \phi)$$

Dark photon
friction essential
for trapping
relaxion after reheating

→ Potentially observable GW signal



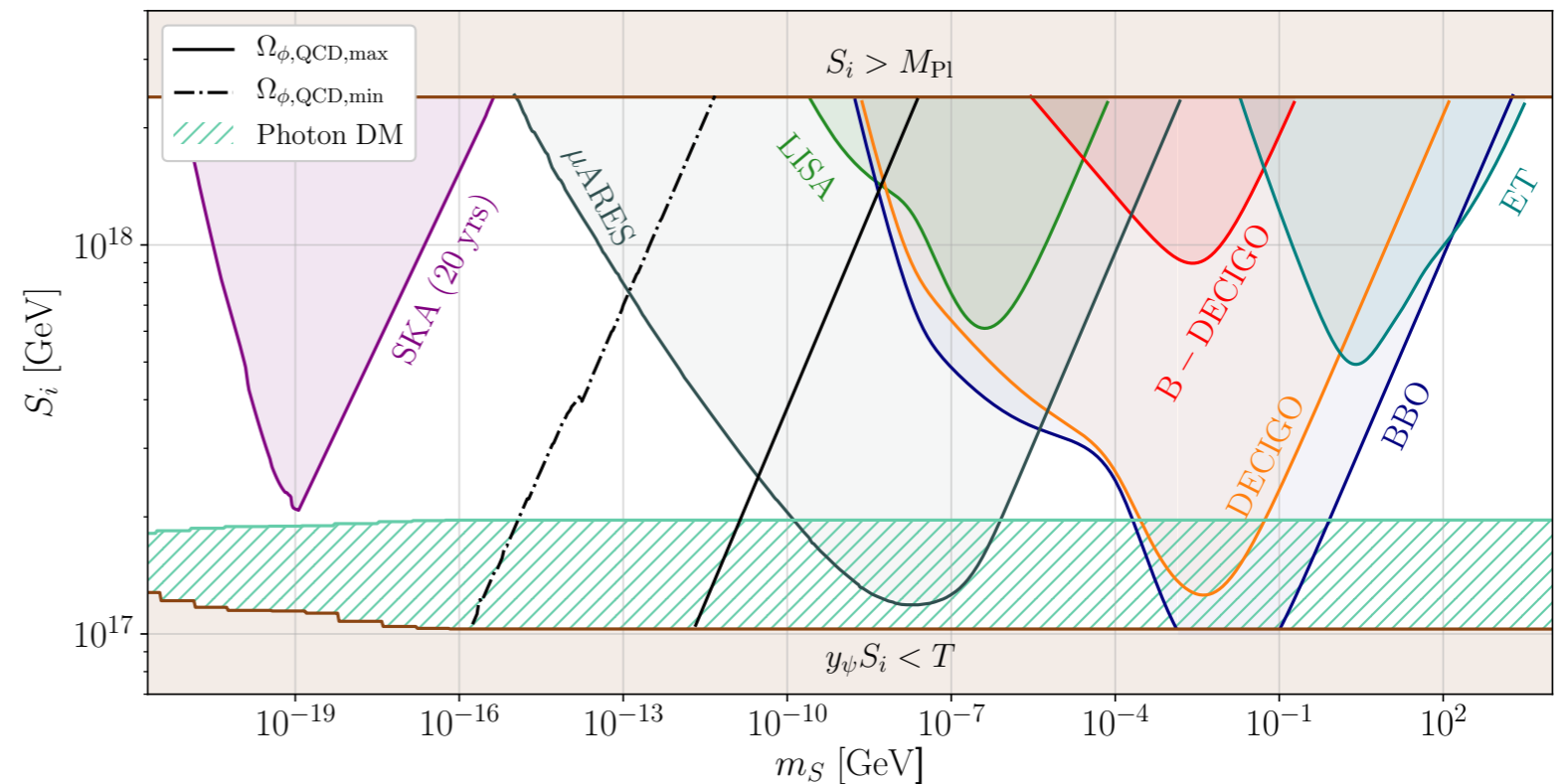
GWs from kinetic misalignment

Consider the case of large initial $\dot{\phi}$

Detectable signal also for smaller decay constants

Fix ALP mass to fit DM relic abundance

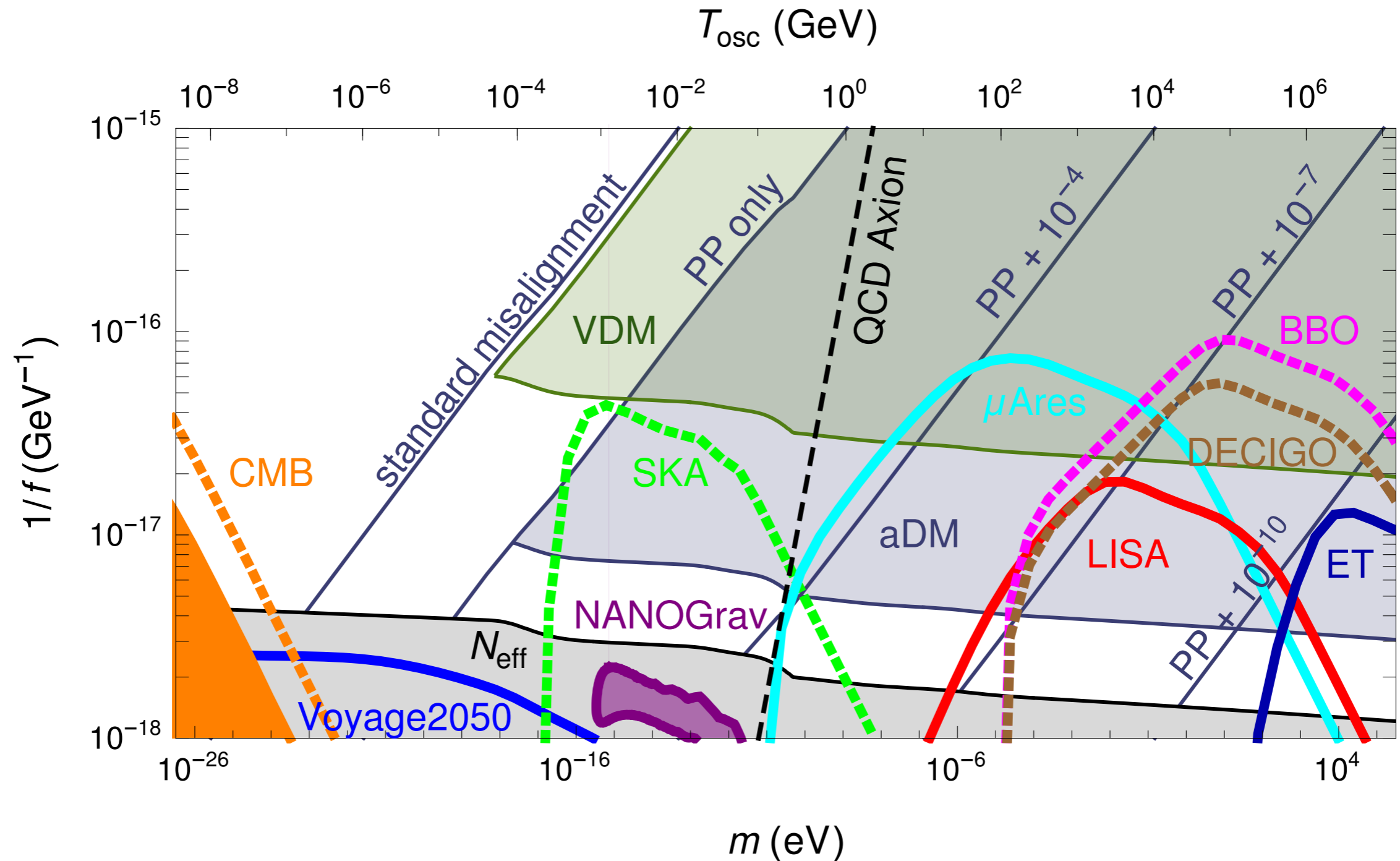
Also consistent with Axiogenesis!



From Madge, Ratzinger, Schmitt, PS, 2111.12730

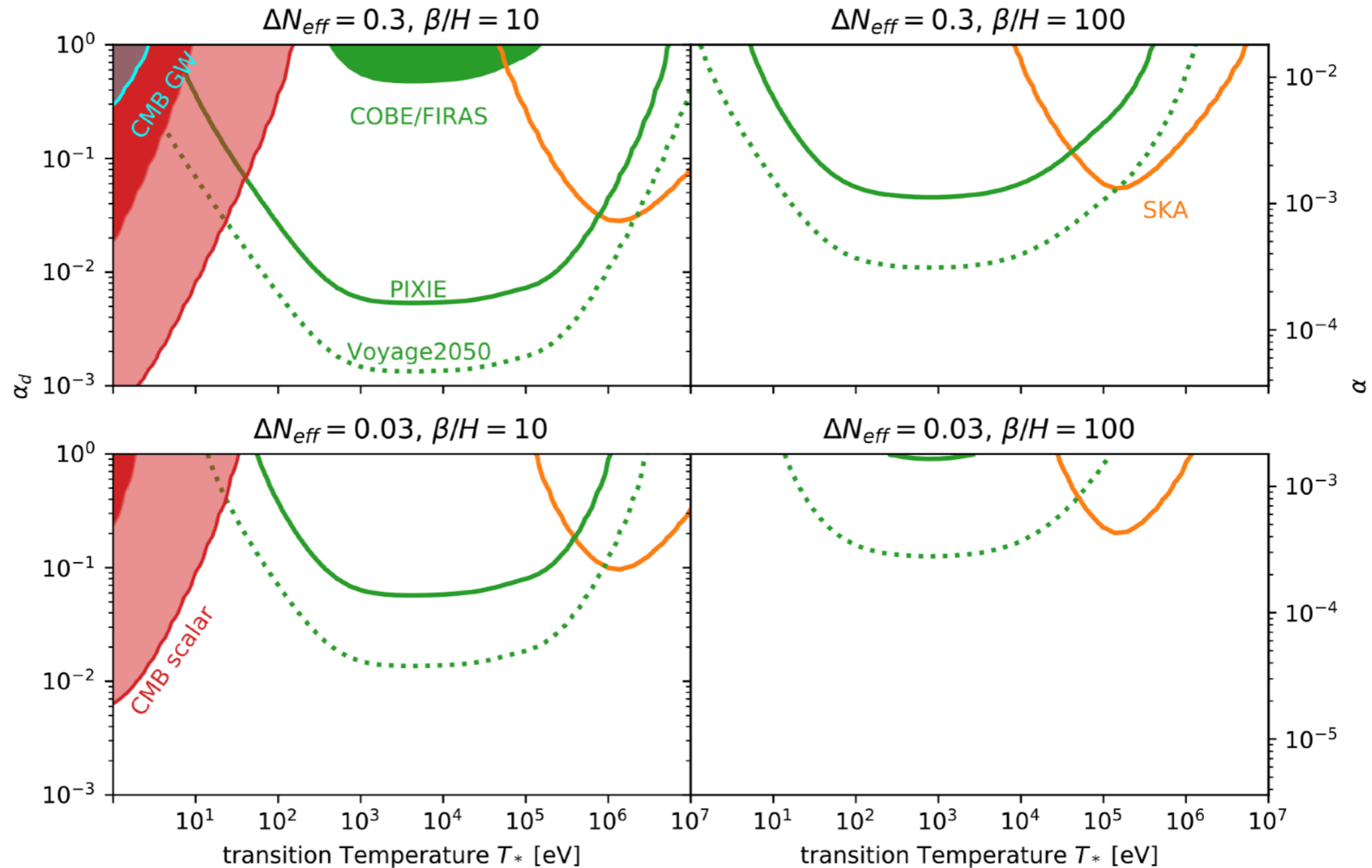
See also Co, Harigaya, Pierce, 2104.02077

Detectable region - update



From 2012.11584 with W. Ratzinger, B. Stefanek

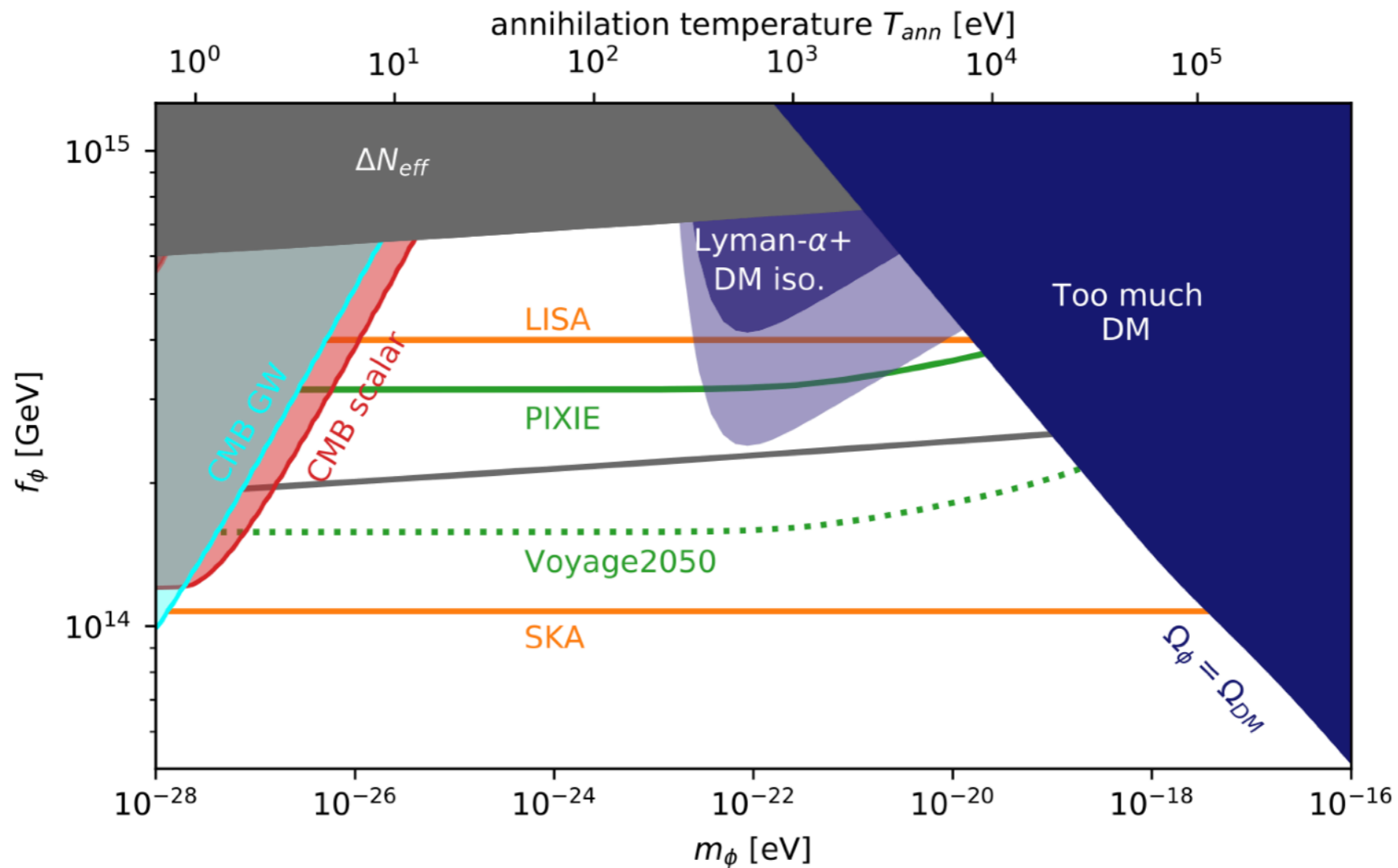
Example source I: Dark sector phase transition



Note: Ω_d fixed to satisfy N_{eff} constraints

Ramberg, Ratzinger & PS, 2209.14313

Source III: (global) cosmic strings



Note: Local strings mainly radiate from small loops and are thus NOT an efficient source of spectral distortions

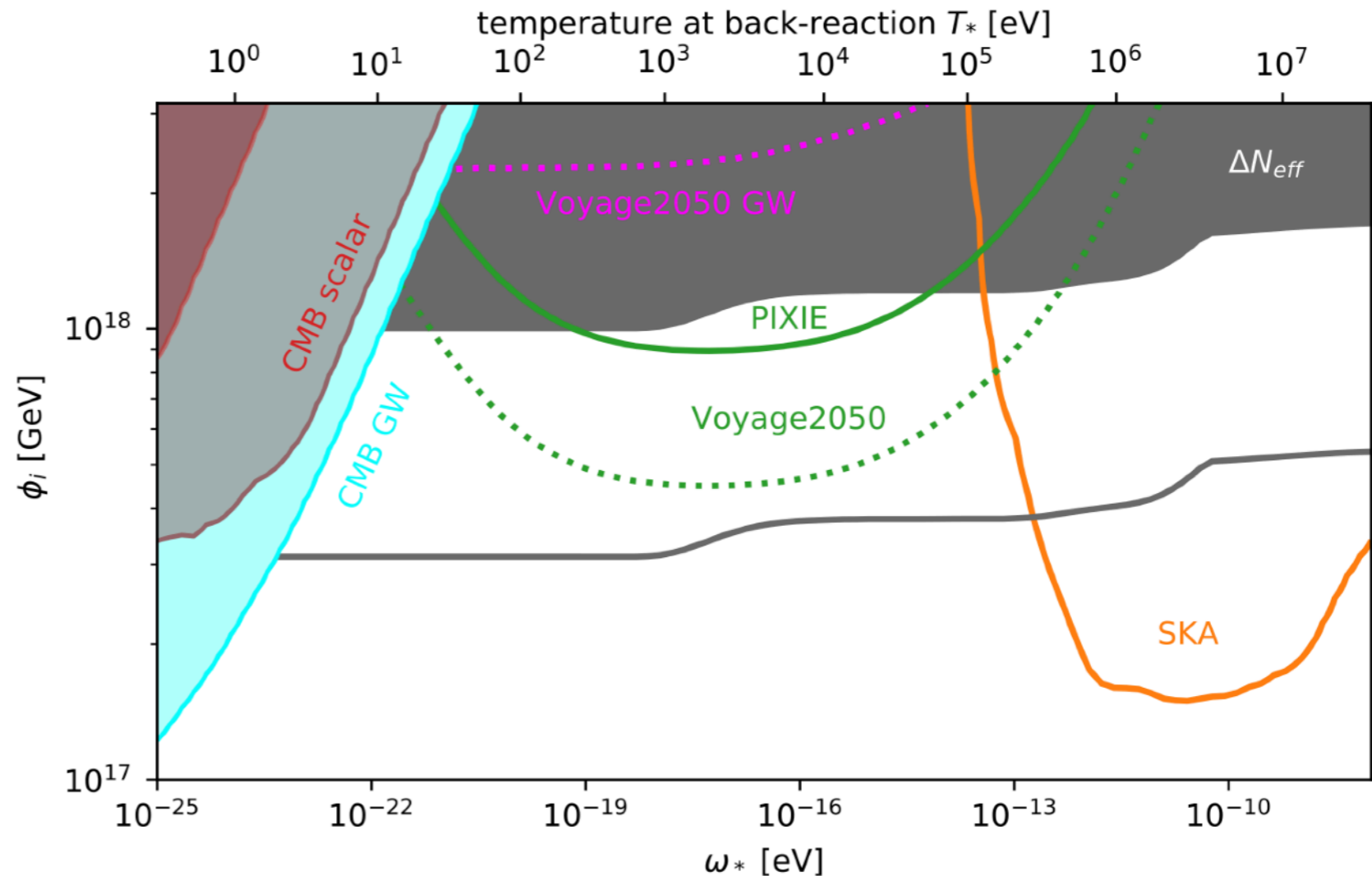
Example source IV: Audible axions...

Not yet...

Results for scalar toy model

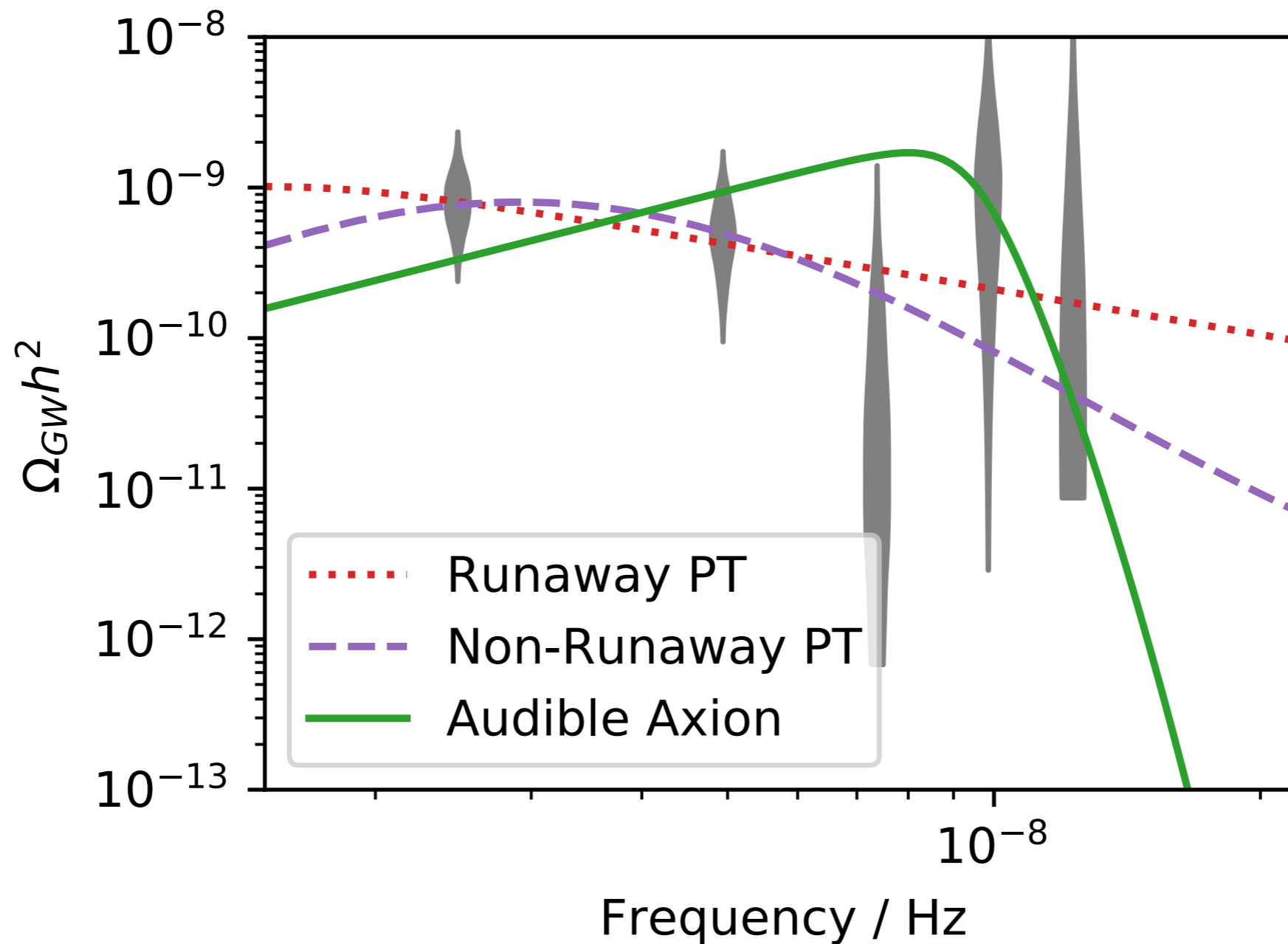
Constraints not as strong since fluctuations are not horizon size

Expect better sensitivity for axion fragmentation



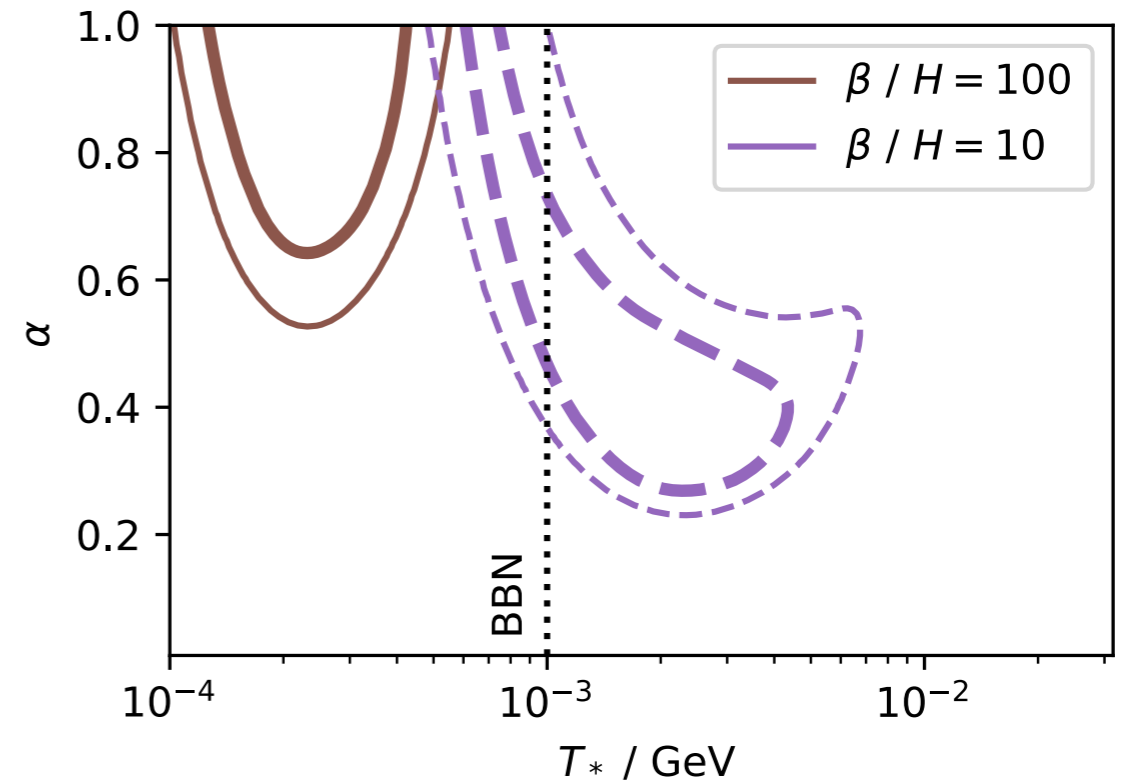
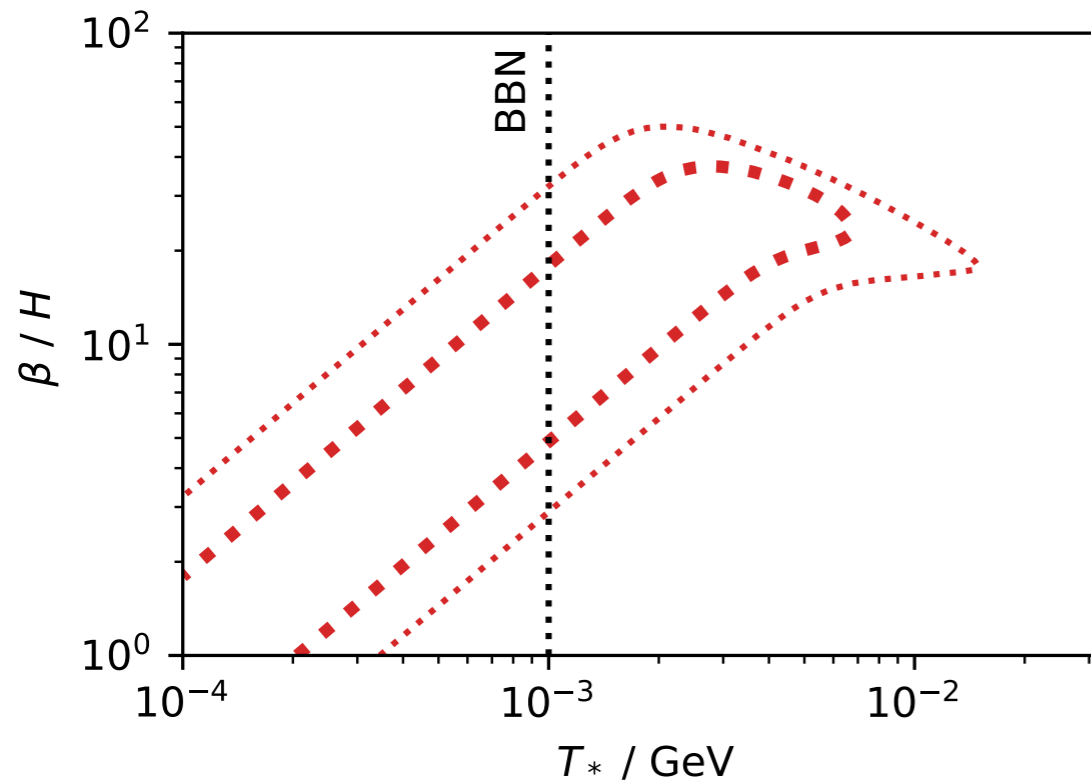
Ramberg, Ratzinger & PS, 2209.14313

Fit with broken power law signals



Wolfram Ratzinger & PS, 2009.11875

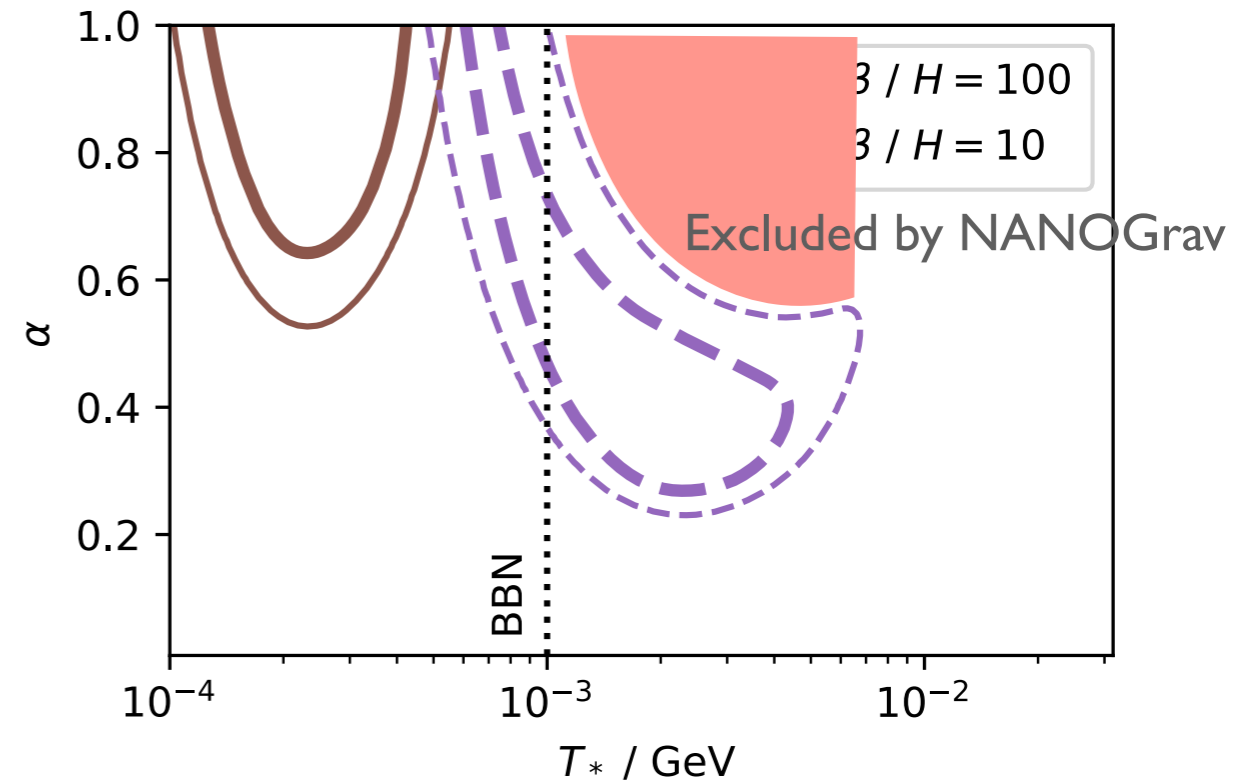
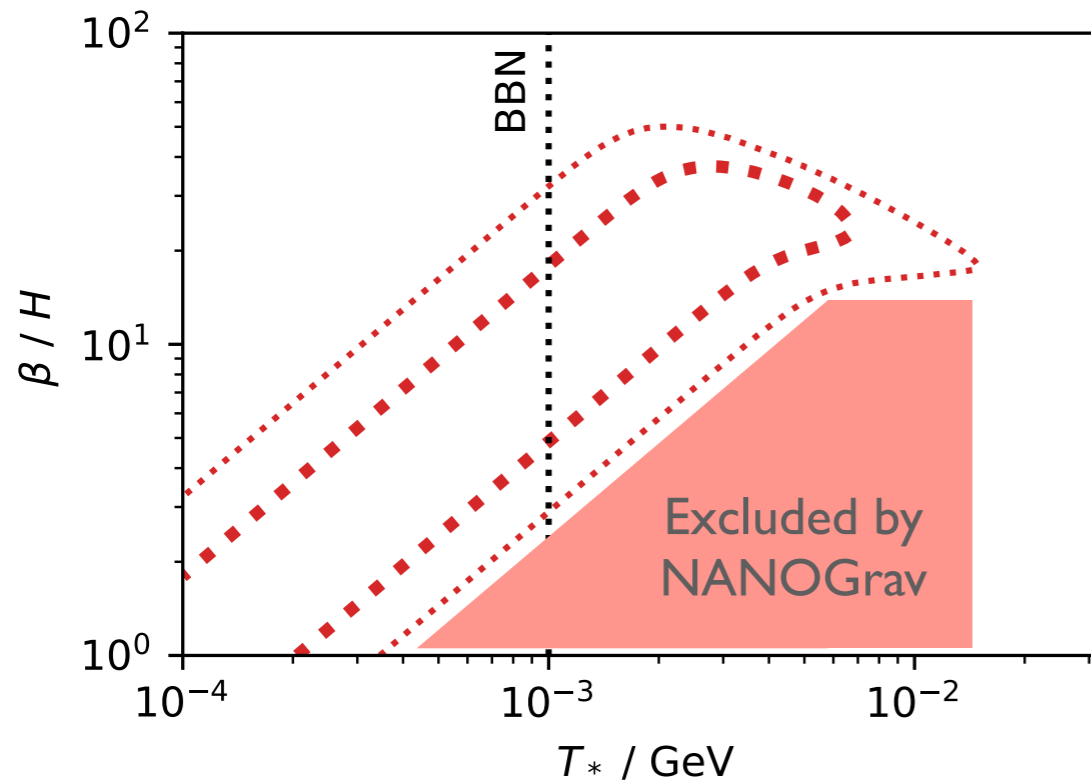
Fit with Phase Transition



Generic PT parameterisation, best fit with PT at temperatures in few MeV range

Challenge for model building \rightarrow Hint for dark sector

Fit with Phase Transition

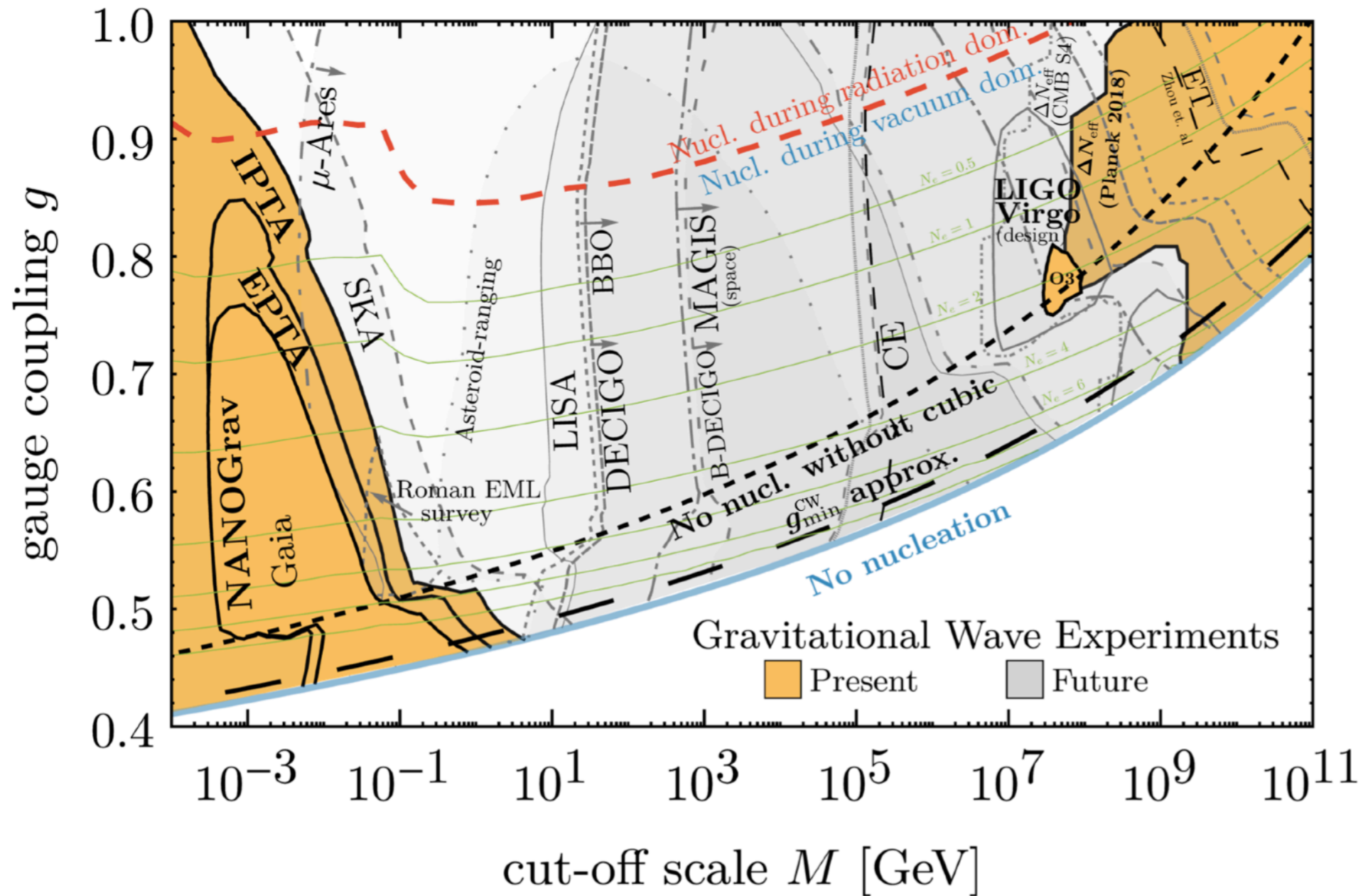


Generic PT parameterisation, best fit with PT at temperatures in few MeV range

Some model parameters excluded by PTA data now!

At higher frequencies

Levi, Opferkuch, Redigolo, 2212.08085



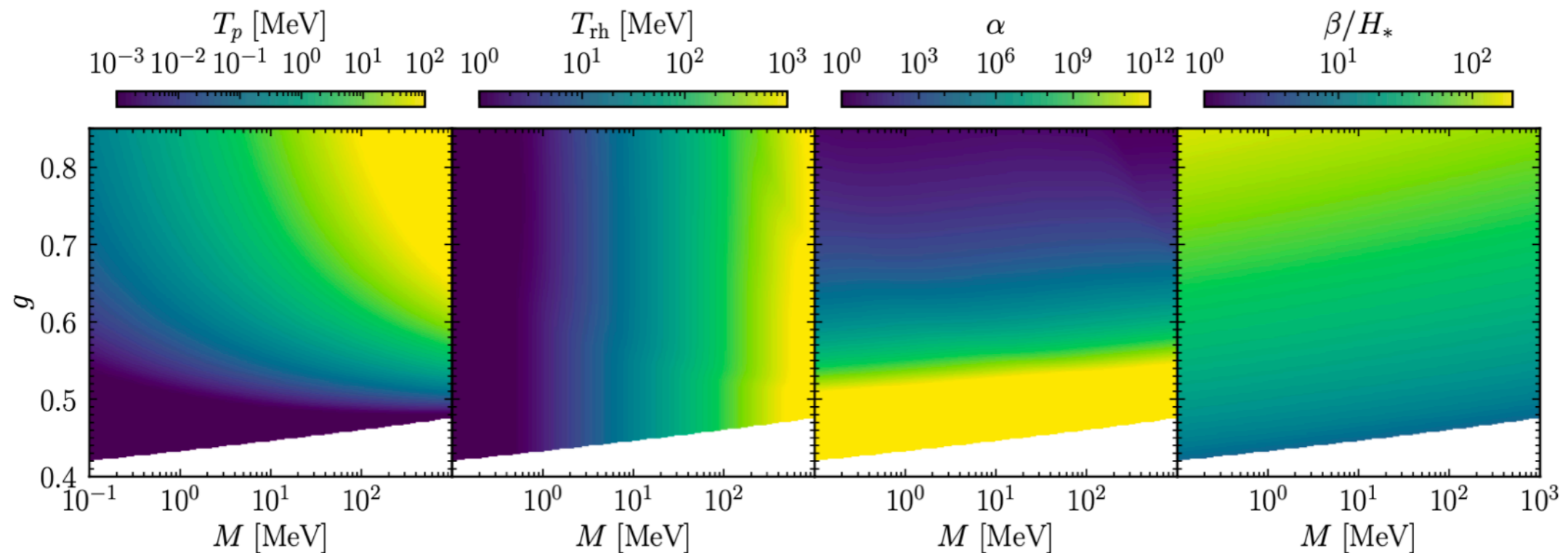
LISA will probe above 10 GeV, colliders could fill gap

Supercooled phase transitions

Benchmark model: Coleman-Weinberg model with vanishing tree level potential

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^2 + D_\mu\Phi^\dagger D^\mu\Phi - V(\Phi, T)$$

Two parameter model: Mass scale M and coupling g



Madge et al,
[2306.14856](https://arxiv.org/abs/2306.14856)

Signal dominated by colliding bubbles and sound shells

Simulated by Lewicki and Vaskonen, 2208.11697

Supercooled phase transitions

Madge et al,
2306.14856

Comparison with
12 year data

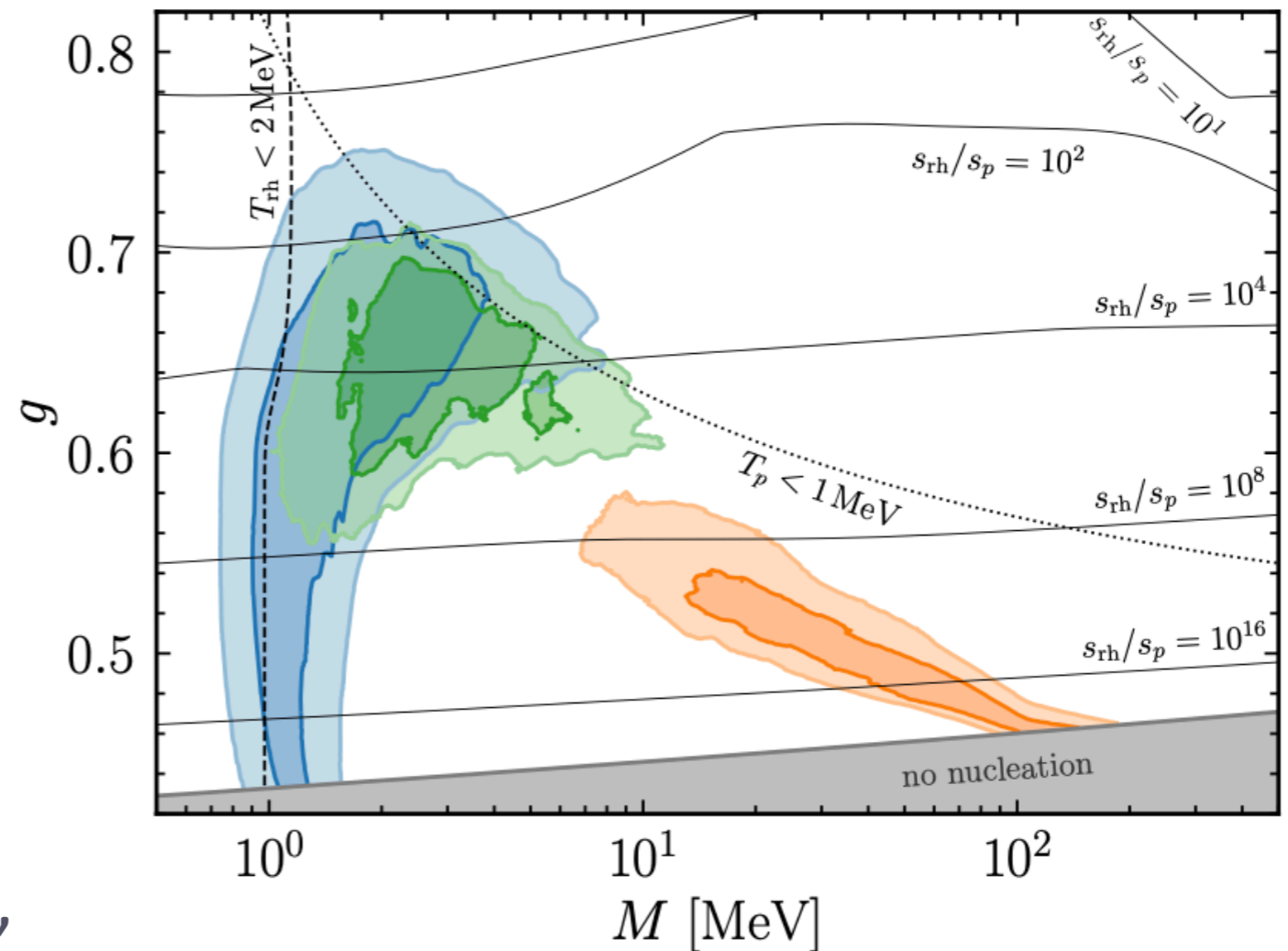
Large supercooling
and reheating

- ▶ Dilution of baryons,
dark matter
- ▶ Two BBNs

Pheno: Light scalar $m_\phi \approx M$,
decay to electrons and photons

Higgs portal not viable, instead

FCC? Or low energy e+e- machine (e.g. MESA in Mainz)



$$\mathcal{L} \supset c_{ee} \frac{|\Phi|^2}{\Lambda^2} LH\bar{e} + c_{\gamma\gamma} \frac{|\Phi|^2}{\Lambda^2} F_{\mu\nu} F^{\mu\nu}$$

Axion/ALP domain walls

Domain walls appear when discrete symmetries are spontaneously broken to degenerate ground states

Long lasting GW source, until DWs annihilate, before dominating the Universe ideally

Review:
Saikawa,
[1703.02576](#)

Axion DW: $U(1)_{\text{PQ}} \rightarrow Z_N$

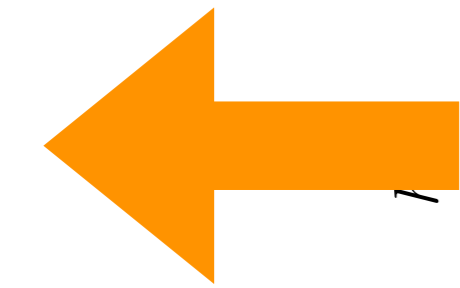
Surface tension $\sigma = 8m_a f_a^2$

Annihilation triggered by QCD instantons

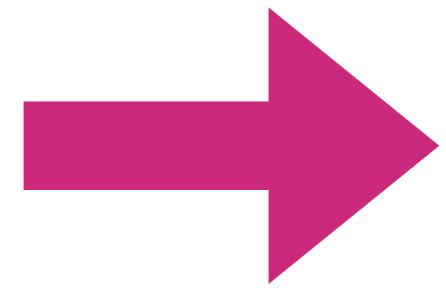
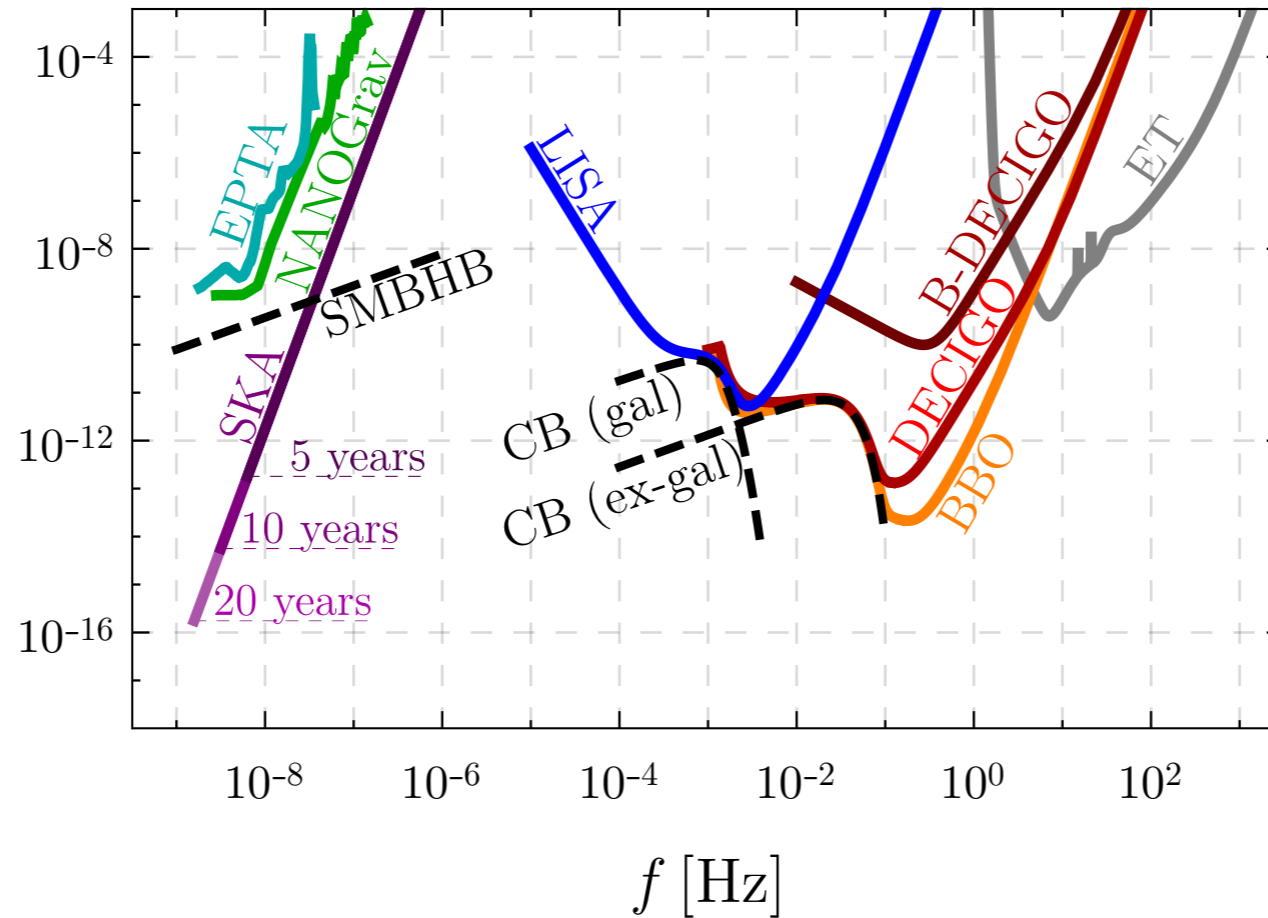
$$T_{\text{ann}} \sim 1 \text{ GeV} \left(\frac{g_*(T_{\text{ann}})}{80} \right)^{-\frac{1}{4}} \left(\frac{\Lambda_{\text{QCD}}}{400 \text{ MeV}} \right)^2 \left(\frac{10^7 \text{ GeV}}{f_a} \right) \sqrt{\frac{10 \text{ GeV}}{m_a}}$$

Madge et al,
[2306.14856](#)

Pushing the limits



CMB
spectral
distortions



High frequency
GW detectors