#### **Gravitational Wave Probes of Dark Matter**



Gianfranco Bertone GRAPPA center of excellence, U. of Amsterdam

GEMMA 2 @ La Sapienza — September 16, 2024

### Plan of the talk:

•Why study DM in strong gravity?

•The DM-BH connection

Gravitational Wave probes of DM

# What is the Universe made of?



### How is DM distributed?





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#### Rotation curve of the Milky Way



locco, GB et al. 2015: <u>http://www.nature.com/nphys/journal/v11/n3/full/nphys3237.html</u>

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http://www.illustris-project.org/media/

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### Candidates

- No shortage of ideas..
- Tens of dark matter models, each with its own phenomenology
- Models span 90 orders of magnitude in DM candidate mass!



Dark Matter Candidate Mass [eV]

### Why study DM in Strong Gravity



GB, Tait, Nature (2018) 1810.01668

 Identifying DM = discriminating among hundreds of DM candidates

• DM candidates differ in terms of:

- small-scale distribution
- Scattering rate:  $\Gamma_{\chi n} \sim \sigma_{\chi n} n_{\chi} n_n$
- Self-annihilation rate:  $\Gamma_{\chi\chi} \sim < \sigma v > n_{\chi}^2$

 Idea: study DM phenomenology in strong gravity = very small scales, very high-densities

### The team



Pippa Cole



Adam Coogan



Bradley Kavanagh



Thomas Spieksma



Daniele Gaggero



Gimmy Tomaselli

+ Ismini Andrianou, Leon Kamermans, Theophanes Karydas, David Nichols, Renske Wierda, ...

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# **Black Holes**

- In GR, completely described by (M, L, q)
- BUT observed (M, L, q; z) drawn from probability distribution that carries information about history (PBHs..)



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- In GR, completely described by (M, L, q)
- BUT observed (M, L, q; z) drawn from probability distribution that carries information about history (PBHs..)
- Don't exist in vacuum. Environment:
  - Enables EM detection (direct imaging of accretion discs, dynamical M from stars, ..)
  - Affects P(M, L, q; z)(q=0, formation scenario, merger rate history, ...)
  - Alters GW signals (dephasing, caracteristic features,...)



Event Horizon Telescope, 2019

### **BH** environments



#### Adiabatic compression of DM around BHs



Conservation of adiabatic invariants:

$$I_i(E_i,L) = I_f(E_f,L) \quad \rightarrow \quad f_f(E_f,L) = f_i(E_i,L) \quad \rightarrow \quad \rho_f(r) = \frac{4\pi}{r^2} \int_{E_f^{\min}}^0 dE_f \int_{L_f^{\min}}^{L_f^{\max}} dL_f \frac{L_f}{v_r} f_f(E_f,L_f) \, .$$

(Peebles 1972, Young 1980, Quinlan, Hernquist and Sigurdsson 1995, Gondolo and Silk 2000, ...)

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#### DM 'spikes'



 $\rho_{\rm cusp}(r) \sim r^{-\gamma}$ 

 $(NFW: \gamma = 1)$ 

$$\rho_{\text{spike}}(r) \sim r^{-\gamma_{\text{sp}}}, \ \gamma_{\text{sp}} = \frac{9 - 2\gamma}{4 - \gamma}$$
$$(\gamma = 1 \rightarrow \gamma_{\text{sp}} = 7/3)$$

Gondolo and Silk 2000

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#### 3 steps: SMS growth, Collapse, BH growth



- I Adiabatic growth on extended stellar object Blumenthal 1986; Young 1980; Spolyar, Freese, Gondolo 2007; Freese et al. 2008
- II Collapse to BH on free-fall timescale E.g. Ullio, Zhao, Kamionkowski 2001 (circular orbits)
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#### Realistic dark matter overdensities around BHs



GB, Wierda, Gaggero, Kavanagh, Volonteri, Yoshida - 2404.0873 I

#### Y-rays from DM spikes in EAGLE simulations



Aschersleben, GB et al JCAP09(2024)005

#### Fermi-LAT, H.E.S.S. and CTAO sensitive to dark matter self-annihilation around IMBHs well below thermal relic cross section

# DM overdensities around PBHs



PBH

'Turnaround' point, when particles decouple from expansion

 $\rho_{\rm DM}(r) \sim r^{-9/4}$ 

Adamek+ 1901.08528, Boudaud+ 2106.07480

### If DM=WIMPs, large annihilation flux!



If (subdominant) PBHs discovered: Extraordinarily stringent constraints on new physics at the weak scale!



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• Detecting a subdominant PBHs with the Einstein Telescope would essentially rule out not only WIMPs, but entire classes of BSM models (even those leading to subdominant DM!)

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#### Phenomenology of DM in Strong Gravity





(Classical paper: Chandrasekhar 1931)



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Additional energy loss term:  $\dot{E}_{orb} = -\dot{E}_{GW} - \dot{E}_{DF}$ 

Evolution of binary separation:  $\dot{r}_2$ 

$$f_2 = - \frac{64 G^3 M m_1 m_2}{5 c^5 (r_2)^3} - \frac{8\pi G^{1/2} m_2 \log \Lambda r_2^{5/2} \rho_{\text{DM}}(r_2)}{\sqrt{M}}$$

$$\sqrt{Mm_1}$$

Easy, right?

(Eda+ 2013, 2014)



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#### Not so fast..

DM distribution is heated: 
$$\Delta E_{\rm DF}(r_{\rm i}, r_{\rm f}) = -\int_{r_{\rm i}}^{r_{\rm f}} \frac{\mathrm{d}E_{\rm DF}}{\mathrm{d}t} \left(\frac{\mathrm{d}r_2}{\mathrm{d}t}\right)^{-1} \mathrm{d}r_2$$



#### Not so fast II...

DM medium NOT homogenous

Scattered particles are in a  $\sim$  torus around the secondary object orbit

Ellipticity, high-v particles, relativistic corrections, accretion etc..



(Kavanagh, GB et al. 2002.12811, Becker+ 2112.09586, Dosopoulou 2305.17281, ...)

### Equal-mass 'Dressed' BH-BH merger



Kavanagh, Gaggero & GB, arXiv:1805.09034

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Kavanagh, Gaggero & GB, arXiv:1805.09034

#### EMRIs = Extreme Mass Ratio Inspirals



 $m_1 \gg m_2$ 

#### Co-evolution of binary and DM distribution

Energy losses due to dynamical friction:

$$\dot{E}_{\rm orb} = -\dot{E}_{\rm GW} - \dot{E}_{\rm DF}$$

Evolution of binary separation:

$$\dot{r}_{2} = -\frac{64 G^{3} M m_{1} m_{2}}{5 c^{5} (r_{2})^{3}} - \frac{8\pi G^{1/2} m_{2} \log \Lambda r_{2}^{5/2} \rho_{\rm DM}(r_{2}, t) \xi(r_{2}, t)}{\sqrt{M} m_{1}}$$



Time-dependent dark matter phase space density:

$$T_{\rm orb}\frac{\partial f(\mathcal{E},t)}{\partial t} = -p_{\mathcal{E}}f(\mathcal{E},t) + \int \left(\frac{\mathcal{E}}{\mathcal{E}-\Delta\mathcal{E}}\right)^{5/2} f(\mathcal{E}-\Delta\mathcal{E},t)P_{\mathcal{E}-\Delta\mathcal{E}}(\Delta\mathcal{E})\,\mathrm{d}\Delta\mathcal{E}$$

Kavanagh, GB et al. 2002. [28] [ [see also Trestini's talk]

#### Time-dependent dark matter density profile



Kavanagh, GB et al 2002.12811, https://doi.org/10.6084/m9.figshare.11663676

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#### Effect of the environment on the waveform



- Waveforms are dephased, with a characteristic  $\Delta\phi(f)$
- Additional energy loss  $\rightarrow$  shorter time to merger

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### Gravitational Waveform dephasing

• Calculate the number of cycles including the effect DM

$$N_{\text{cycles}}(t_{\text{f}}, t_{\text{i}}) = \int_{t_{\text{i}}}^{t_{\text{f}}} f_{\text{GW}}(t) \mathrm{d}t$$

Calculate difference wrt vacuum

 $\Delta N_{\rm cycles} = N_{\rm cycles}^{\rm vac}(f_{\rm GW,f}, f_{\rm GW,i}) - N_{\rm cycles}^{\rm DM}(f_{\rm GW,f}, f_{\rm GW,i})$ 

- Static: Assuming DM fixed (Eda+ 2013, 2014)
- Dynamic: including evolution of DM phase space (2002.12811)



Kavanagh, GB et al. 2002. I 28 I I

#### Detecting / discovering / Measuring DM with GWs

#### Coogan, GB, Gaggero, Kavanagh Nichols 2021



- Dark dresses within  $\sim 100$  Mpc are <u>detectable</u> with Lisa
- Can discover that fiducial systems are <u>not</u> <u>GR-in-vacuum</u> (in terms of Bayes factor)
- Can <u>measure</u>:
  DM density profile

  normalization
  slope

  mass ratio

# OK, but can we identify DM with GW observations?

Can we tell e.g. WIMPs from ultra-light DM, WDM, selfinteracting DM, etc..?

### Gravitational atoms



Y. Zel'Dovich (1971,1972); C. Misner (1972); A. Starobinsky (1973); Detweiler (1980); W. East and F. Pretorius (2017); and many many others, see e.g. the review by R. Brito, V. Cardoso, and P. Pani (2015)

- If ultra-light bosons exist, they can be produced around rotating black holes through Superradiance
- Extraction of mass and angular momentum
   → cloud of the bosonic field
- BH + boson cloud = gravitational atom. Bound states |nlm> in analogy with proton + electron structure in H atom

See talk by Cristina Mondino

### EMRIs in presence of Gravitational Atoms



Energy lost by the binary due to 'ionisation'

- 'Resonances' due to transitions between bound states  $< a \mid V_*(t) \mid b >$ Baumann, Chia, Porto, arXiv:1804.03208
- 'lonization', i.e. transitions to continuum  $< a | V_*(t) | klm >$ Baumann, GB, Stout, Tomaselli Phys.Rev.Lett. 128 (2022) 22, 221102
- Role of accretion on companion, eccentricity, inclination Baumann, GB, Stout, Tomaselli 2112.14777, Tomaselli, Spieksma, GB 2305.15460, 2403.03147

### Published yesterday:



- When inclination angle falls inside angular interval  $\chi_i$  around a counterrotating configuration, the cloud survives all the resonances, becoming observable late in the inspiral
- Otherwise, cloud is destroyed (red line), leaving a distinctive mark on the orbital parameters.
- Binaries that form at small radii are an exception: They may skip the destructive (hyper)fine resonances.

#### Density profiles depend on the DM properties

#### Self-annihilating DM



[GB & Merritt astro-ph/0504422, Shapiro & Shelton1606.01248]

#### $10^{32}$ CDM, $1/2 \le \gamma_{\rm sp} \le 7/3$ SIDM, contact interaction $10^{29}$ SIDM, massless mediator DM density $ho \, [{\rm M}_\odot/{\rm Mpc}^3]$ SIDM, massive mediator $10^{26}$ $10^{23}$ $10^{20}$ $10^{17}$ < Fd $10^{14}$ $10^{-4}$ $10^{-3}$ $10^{-2}$ $10^{-1}$ $10^{0}$ $10^{1}$ $10^{2}$ $10^{3}$ $10^{4}$ $10^{5}$ $10^{6}$ Radius r [pc]

[Alonso-Alvarez+ 2401.14450]

#### Self-interacting DM

# In case of detection, how well can we reconstruct parameters?



Cole, GB et al. Nature Astron. 7 (2023) 8, 943-950

### New results/Work in progress

- Realistic spike formation scenarios, via formation and collapse of Supermassive Stars (2404.08731)
- **Refined modeling** of eccentricity, accretion, torques, etc (2402.13053, 2402.13762, 2403.03147)
- Relativistic effects
- Fast statistical inference of environments w/ machine learning
- •Imprint of DM particle properties on the waveform
- •Population studies, Merger rates, etc

### Gravitational wave probes of DM



"Gravitational wave probes of dark matter: challenges and opportunities" GB, Croon, et al. 1907.10610

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# Conclusions

- Studying DM in strong gravity opens new opportunities to identify it
- DM can reach very high density around BHs
- $\bullet$  We can probe these very high densities with (Y-rays and) GWs



# Supplementary material

## **Further GW-DM connections:**



"Gravitational wave probes of dark matter: challenges and opportunities" GB, Croon, et al. 1907.10610

### **BH** environments



### Other environments



Pippa Cole, GB + <u>2302.03351</u>