

04.09.2024

NOW 2024

T2K-NOvA Tension and Physics BSM



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Outstanding progress in ν physics in ~ 25 years

Discoveries	Interpretation	known knowns
<p>Zenith angle dependence (Multi-GeV)</p> <p> $\chi^2(\text{shape}) = 2.8/4 \text{ dof}$ $U_P = 0.93 +0.13 -0.12$ </p> <p> $\chi^2(\text{shape}) = 30/4 \text{ dof}$ $U_P = 0.54 +0.06 -0.05$ (6.2σ !!) </p> <p> $\phi_{\mu} (\times 10^6 \text{ cm}^{-2} \text{ s}^{-1})$ $\phi_e (\times 10^8 \text{ cm}^{-2} \text{ s}^{-1})$ </p> <p> + many other ones: solar, KamLAND, θ_{13} at reactors & T2K ... </p>	<p> $\Delta m^2 [\text{eV}^2]$ $\tan^2 \theta$ </p> <p> $\nu_e \leftrightarrow \nu_\mu$ $\nu_\mu \leftrightarrow \nu_\tau$ $\nu_e \leftrightarrow \nu_\tau$ $\nu_e \leftrightarrow \nu_\mu$ </p> <p> <small>All limits are at 90%CL unless otherwise noted</small> </p> <p> http://hitoshi.berkeley.edu/nuetrino </p>	<p>known knowns</p> <p> $\delta m^2/\text{eV}^2 \sim 7.34 \times 10^{-5} \pm 2.2\%$ $\Delta m^2/\text{eV}^2 \sim 2.48 \times 10^{-3} \pm 1.3\%$ $\sin^2 \theta_{12} \sim 0.303 \pm 4.4\%$ $\sin^2 \theta_{13} \sim 0.0225 \pm 3.8\%$ $\sin^2 \theta_{23} \sim 0.545 \pm 5.0\%$ </p> <p>known unknowns</p> <p> $\delta(\text{CP})$ $\text{sign}(\Delta m^2)$ $\text{octant}(\theta_{23})$ absolute ν mass Dirac/Majorana NSI, sterile states, non-unitarity, ... </p> <p>unknown unknowns</p>

3-flavor scheme now established as the standard framework...

The 3ν mixing matrix

$$|\nu_\alpha\rangle = \sum_{i=1}^3 U_{\alpha i}^* |\nu_i\rangle \quad U = O_{23} \Gamma_\delta O_{13} \Gamma_\delta^\dagger O_{12}$$

$$\Gamma_\delta = \text{diag}(1, 1, e^{+i\delta})$$

$$\delta \in [0, 2\pi]$$

Dirac CP-violating phase δ

U is non-real if $\delta \neq (0, \pi)$

Explicit
form

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

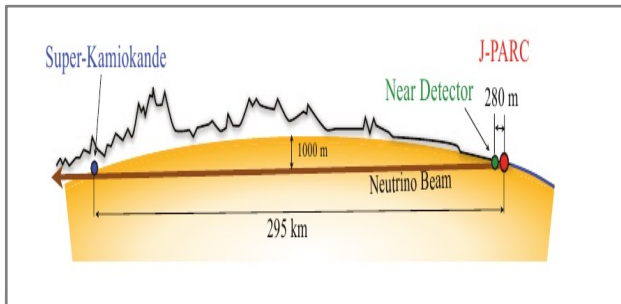
$$\theta_{23} \sim 45^\circ$$

$$\theta_{13} \sim 9^\circ$$

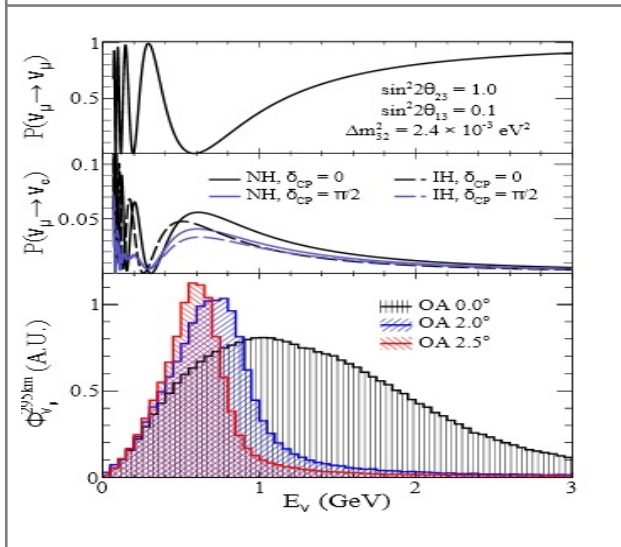
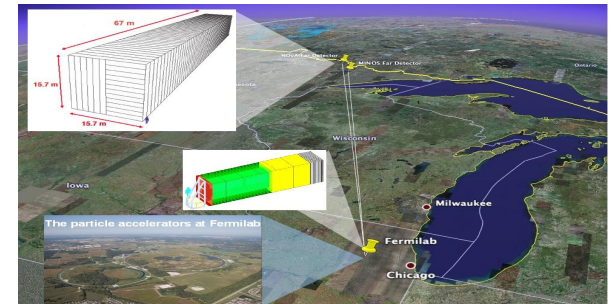
$$\theta_{12} \sim 34^\circ$$

Three non-zero θ_{ij} : Way open to CPV searches...

Two main actors: T2K & NO_vA



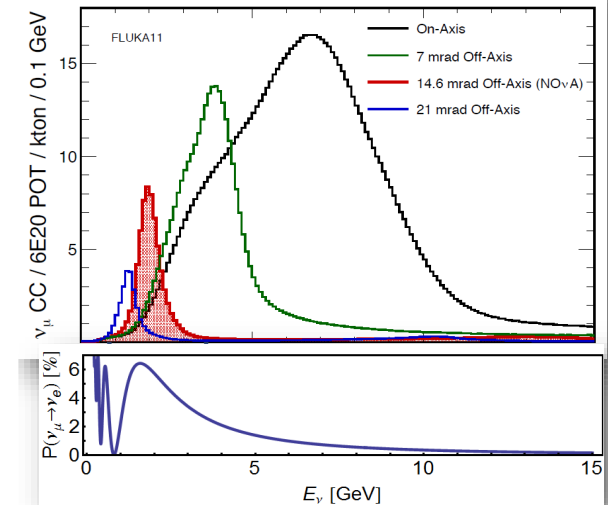
off-axis
beam



$$\Delta = \frac{\Delta m_{13}^2 L}{4E} \approx \frac{\pi}{2}$$

First
oscillation
maximum

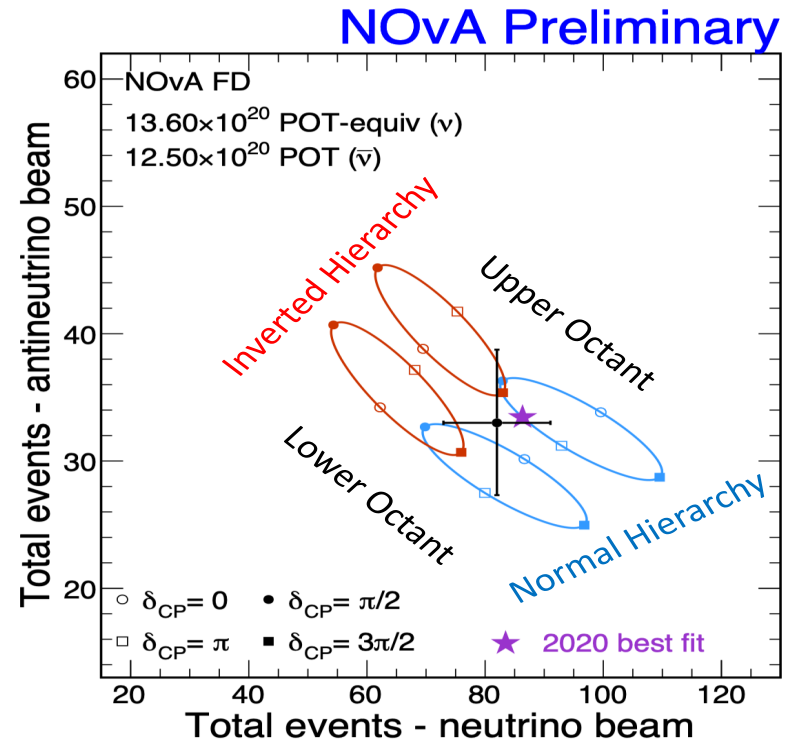
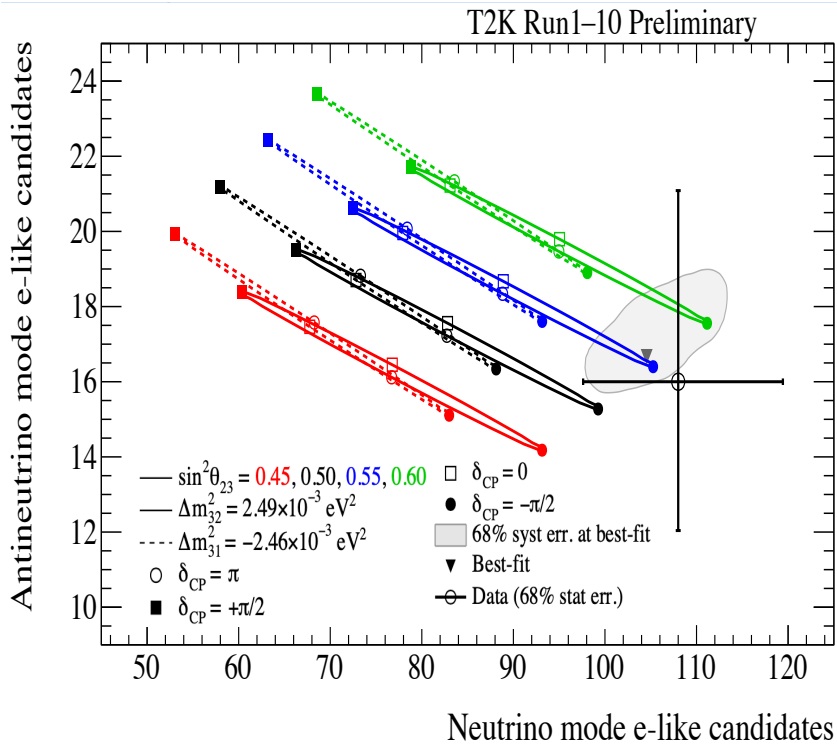
Far Detector flux NO_vA Simulation



E = 0.6 GeV
L = 295 km

E = 2 GeV
L = 810 km

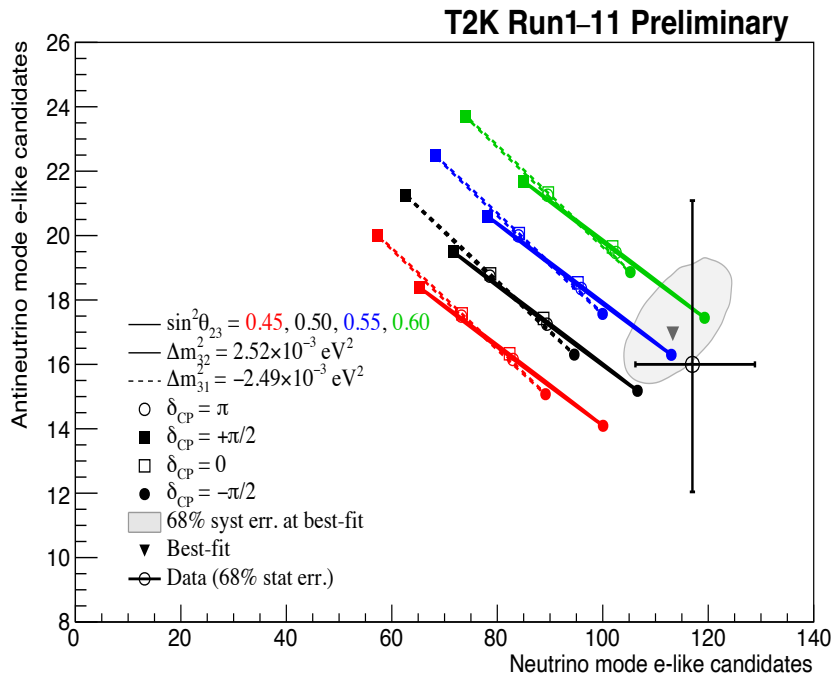
Bird's-eye view: 2020



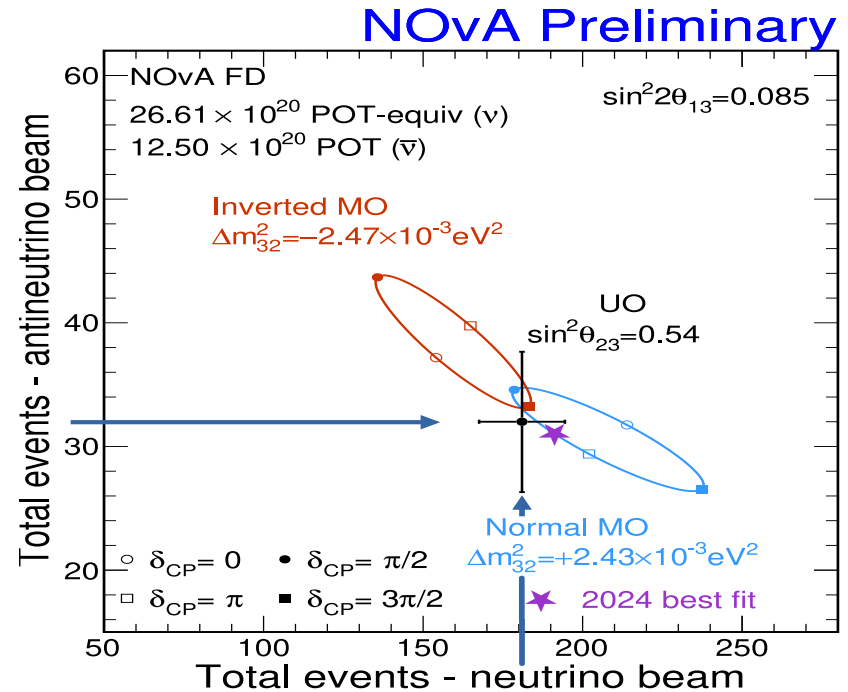
TENSION
for Normal Ordering:

$\left\{ \begin{array}{l} \text{T2K prefers } \delta_{CP} \sim 1.5\pi \\ \text{NOvA prefers } \delta_{CP} \sim 0.8\pi \end{array} \right.$

Bird's-eye view: 2024



T2K almost unaltered statistics



NOvA doubled $\bar{\nu}$ statistics

TENSION INCREASED
for Normal Ordering:

T2K prefers $\delta_{CP} \sim 1.5\pi$
NOvA prefers $\delta_{CP} \sim 0.9\pi$

Almost unchanged

Soon after Neutrino 2020

PHYSICAL REVIEW LETTERS **126**, 051802 (2021)

Nonstandard Neutrino Interactions as a Solution to the NO ν A and T2K Discrepancy

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(Received 17 August 2020; accepted 3 December 2020; published 4 February 2021)

The latest data of the two long-baseline accelerator experiments NO ν A and T2K, interpreted in the standard three-flavor scenario, display a discrepancy. A mismatch in the determination of the standard CP phase δ_{CP} extracted by the two experiments is evident in the normal neutrino mass ordering. While NO ν A prefers values close to $\delta_{CP} \sim 0.8\pi$, T2K identifies values of $\delta_{CP} \sim 1.4\pi$. Such two estimates are in disagreement at more than 90% C.L. for 2 degrees of freedom. We show that such a tension can be resolved if one hypothesizes the existence of complex neutral-current nonstandard interactions (NSIs) of the flavor changing type involving the $e - \mu$ or the $e - \tau$ sectors with couplings $|e_{e\mu}| \sim |e_{e\tau}| \sim 0.2$. Remarkably, in the presence of such NSIs, both experiments point towards the same common value of the standard CP phase $\delta_{CP} \sim 3\pi/2$. Our analysis also highlights an intriguing preference for maximal CP violation in the nonstandard sector with the NSI CP phases having best fit close to $\phi_{e\mu} \sim \phi_{e\tau} \sim 3\pi/2$, hence pointing towards imaginary NSI couplings.

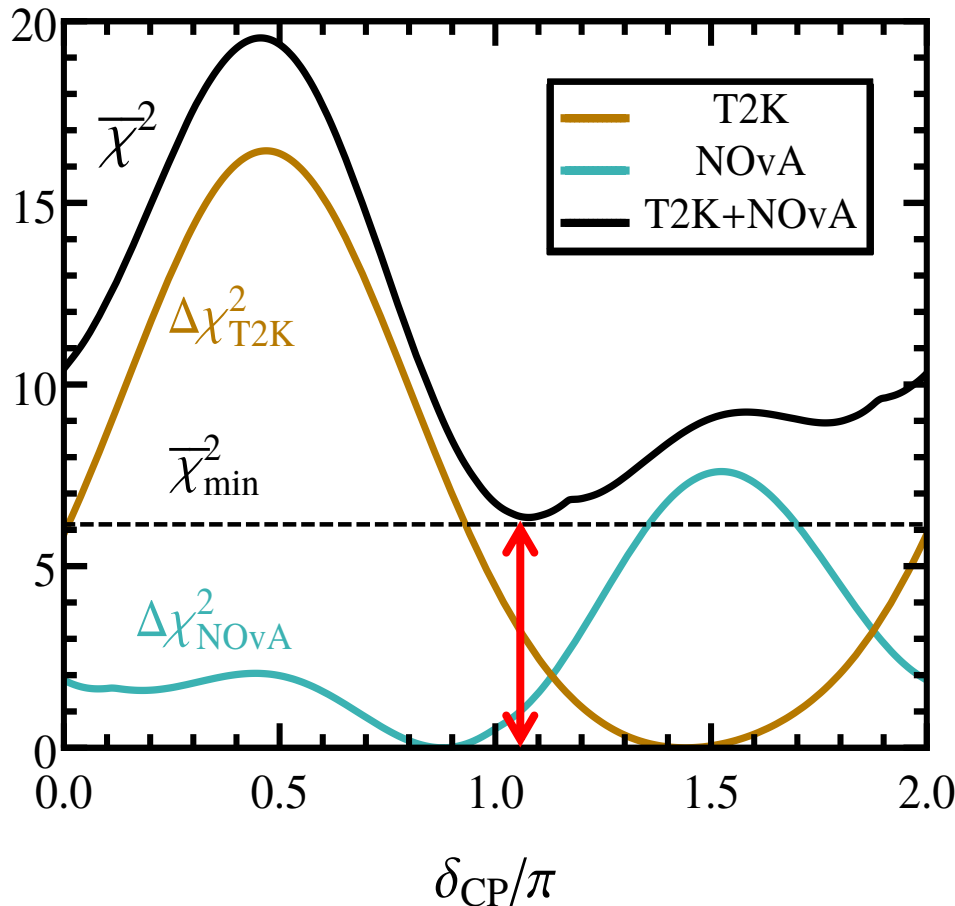
DOI: [10.1103/PhysRevLett.126.051802](https://doi.org/10.1103/PhysRevLett.126.051802)

See also Denton, Gehrlein & Pestes, PRL 126 051801 (2021)

After Neutrino 2024

Chatterjee & Palazzo arXiv:2409.XXXXX

$$\bar{\chi}^2(\delta_{\text{CP}}) = \chi_{\text{T2K+NOvA}}^2(\delta_{\text{CP}}) - (\chi_{\text{T2K,min}}^2 + \chi_{\text{NOvA,min}}^2)$$



Maltoni & Schwetz Criterion
hep-ph/0304176 PRD
(2003)

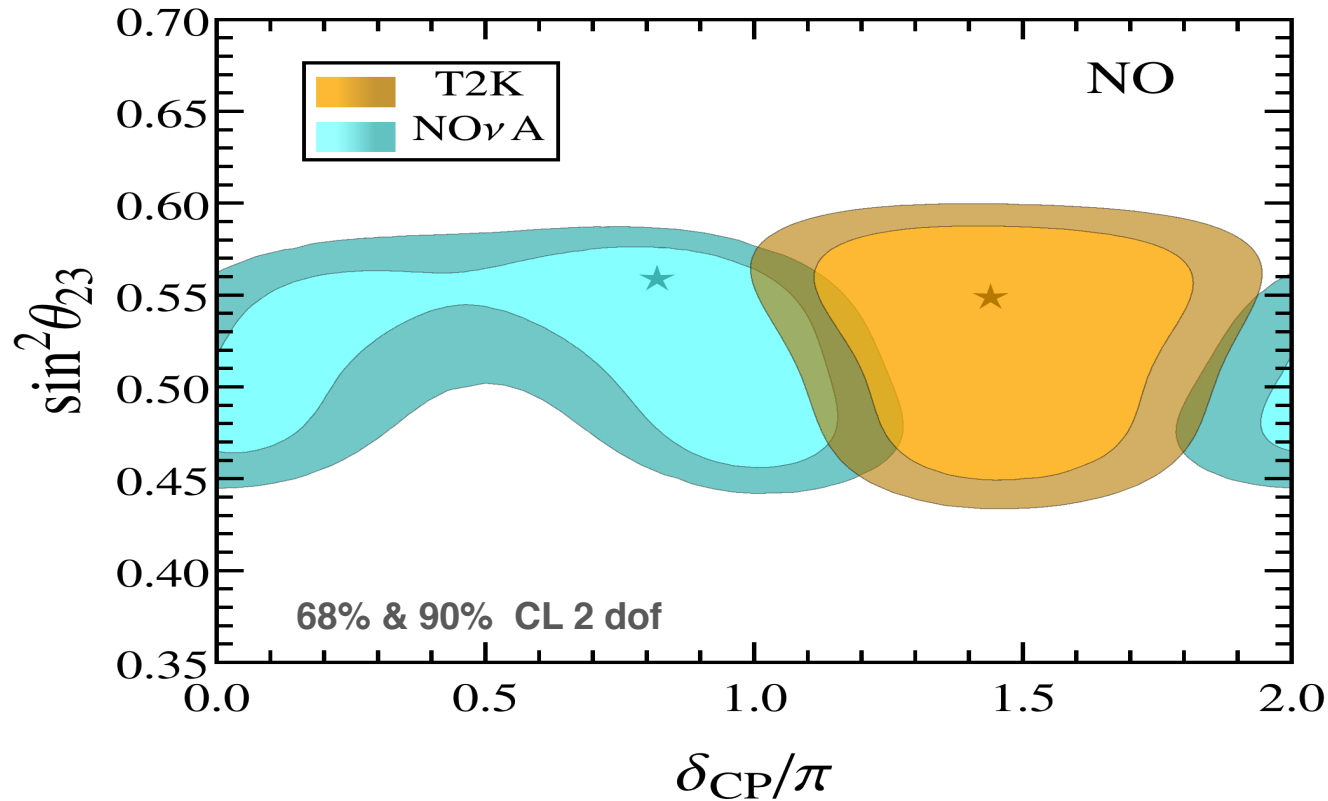
$$\bar{\chi}_{\text{min}}^2 = 6.3$$

$$\text{GoF} = 1.4 \times 10^{-2}$$

Alert for:

- Experimentalists
- Phenomenologists
- Theorists

View in two-parameter space (2024)



Maybe a statistical fluctuation or a systematic error

But interesting to consider alternative explanations...

Why to consider non-standard interactions

T2K and NOvA have different baselines and peak energies ($L/E = \text{constant}$)

Matter effects depend on the ratio $v = \frac{2V_{CC}E}{\Delta m_{31}^2} = 0.18 \left[\frac{E}{2.0 \text{ GeV}} \right]$

T2K	$v \sim 0.05$
NOvA	$v \sim 0.17$

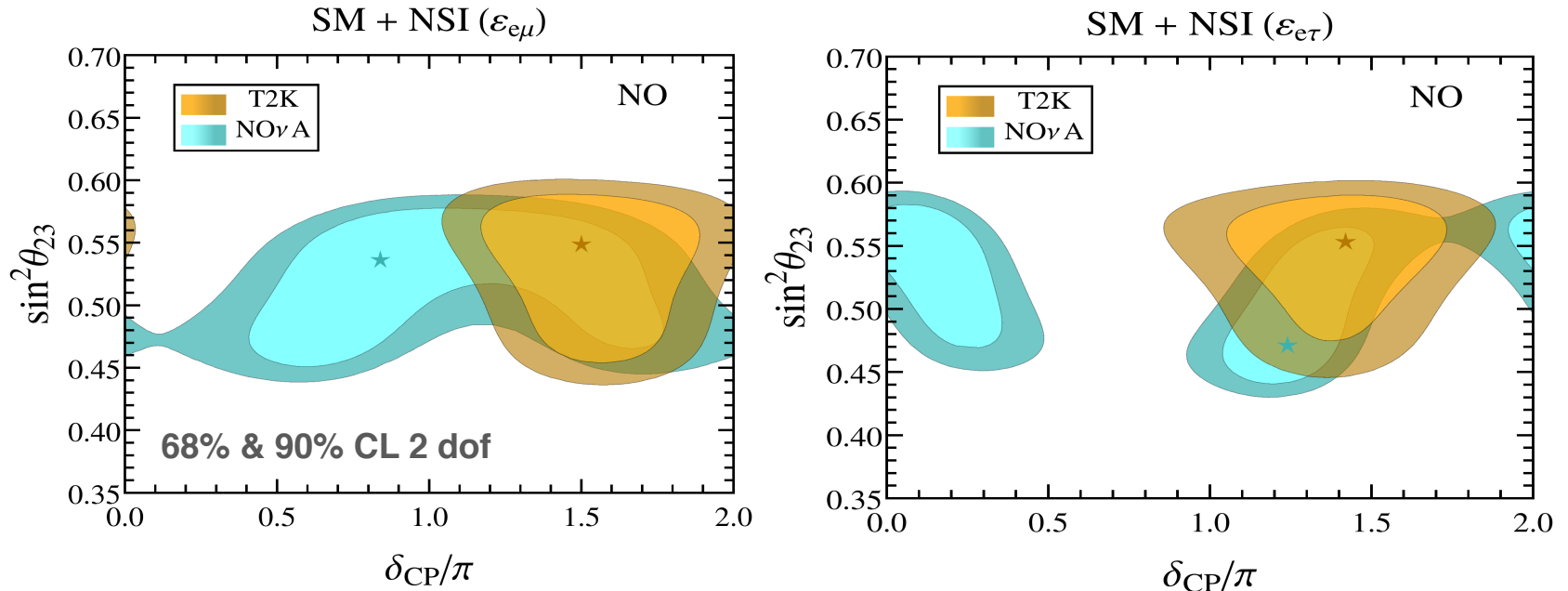
New matter effects encoded by NSI are also proportional to v

Basic Idea: suppose NSI exist, then:

T2K is a “quasivacuum” experiment. Its estimate of δ_{CP} is independent of NSI.

NOvA is a “matter dominated” experiment. The extracted value of δ_{CP} is affected by NSI. If NSI are taken into account, the estimate of δ_{CP} should return in agreement with that of T2K.

NSI bring the estimates of δ_{CP} in agreement



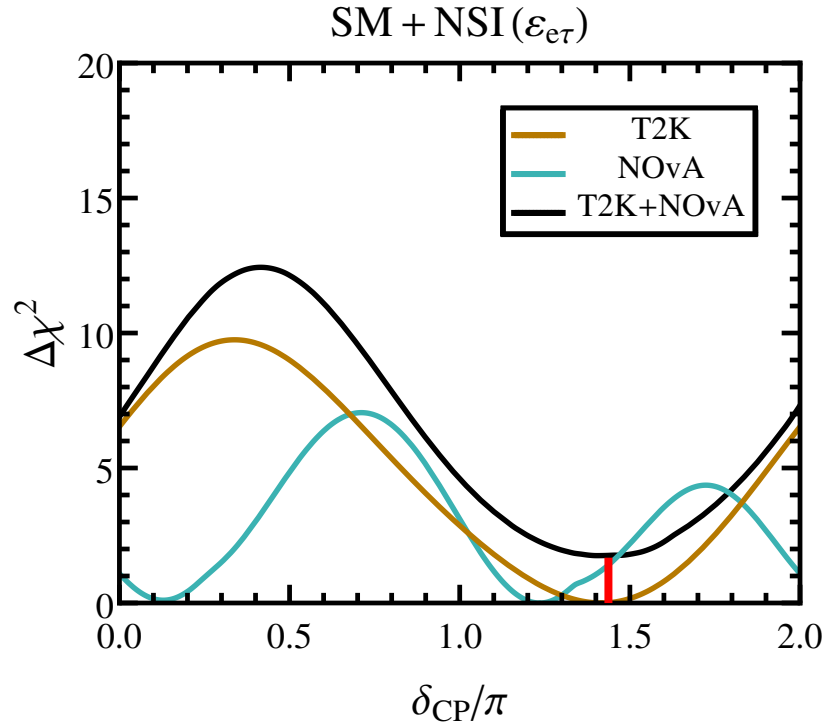
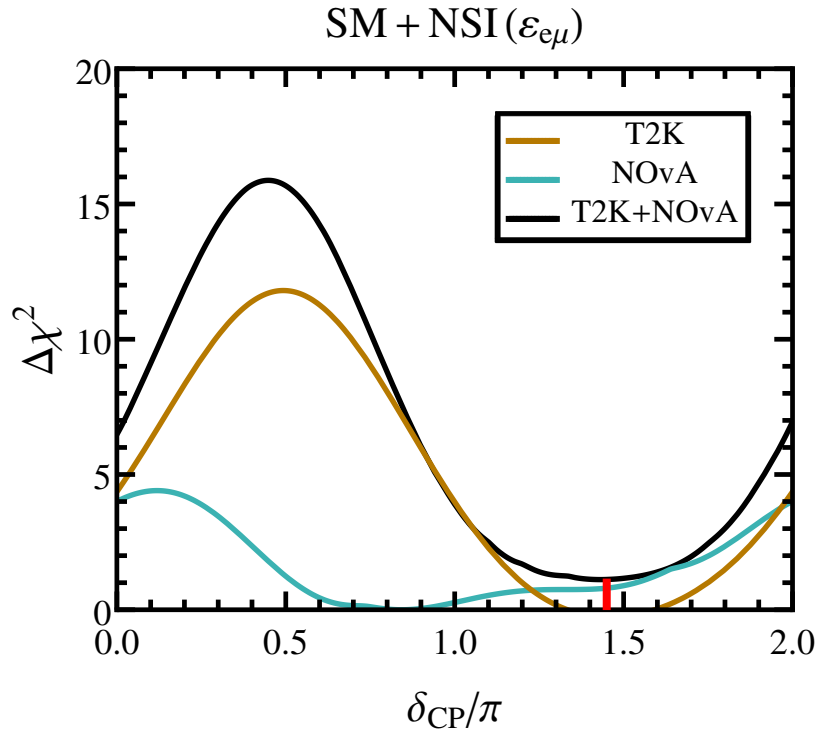
Contours obtained for the best fit of T2K + NOvA:

$$\begin{cases} [\epsilon_{e\mu} = 0.13, \phi_{e\mu} = 1.35\pi] \\ [\epsilon_{e\tau} = 0] \end{cases}$$

$$\begin{cases} [\epsilon_{e\mu} = 0] \\ [\epsilon_{e\tau} = 0.22, \phi_{e\tau} = 1.70\pi] \end{cases}$$

T2K regions almost unaltered **NOvA regions strongly modified**

NSI substantially reduce the tension

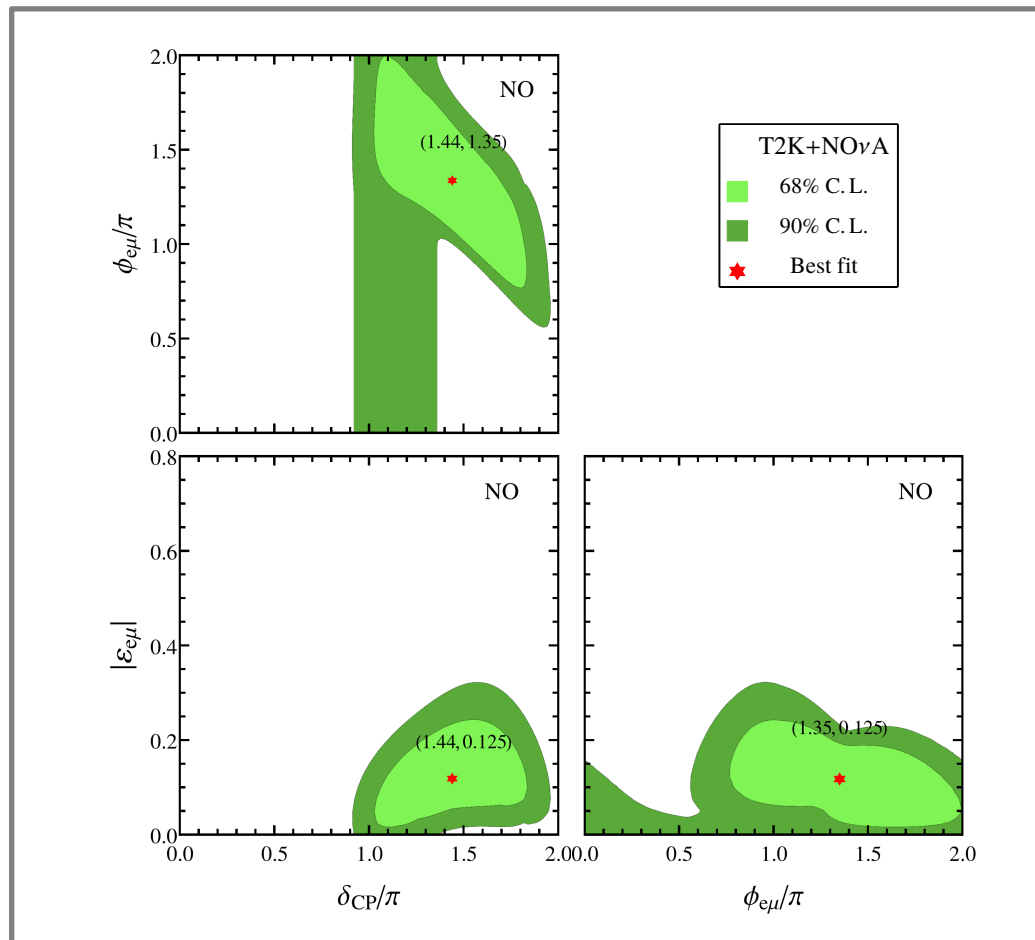


Single $\Delta\chi^2$ obtained for the best fit of T2K + NOvA:

$$\begin{cases} [\varepsilon_{e\mu} = 0.13, \phi_{e\mu} = 1.35\pi] \\ [\varepsilon_{e\tau} = 0] \end{cases}$$

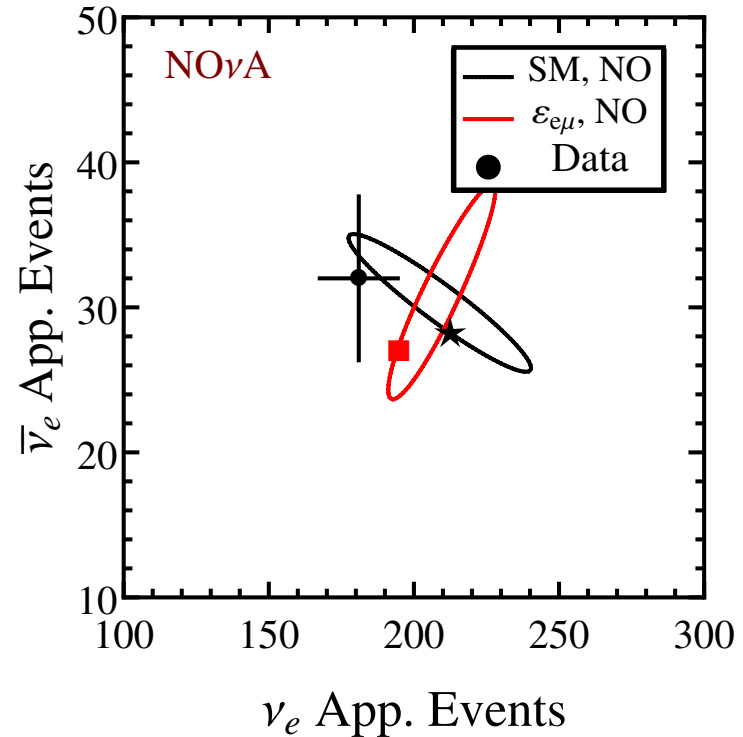
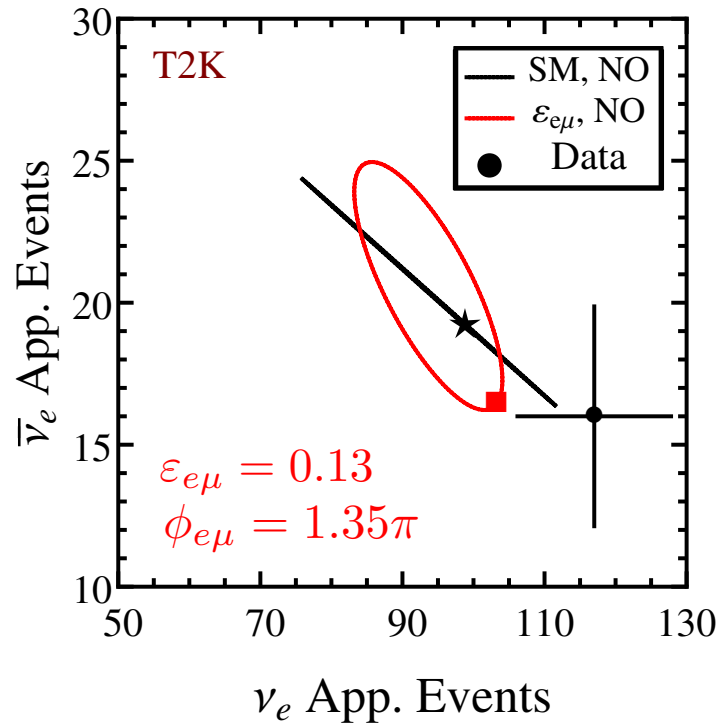
$$\begin{cases} [\varepsilon_{e\mu} = 0] \\ [\varepsilon_{e\tau} = 0.22, \phi_{e\tau} = 1.70\pi] \end{cases}$$

Hint of non-zero $\varepsilon_{e\mu}$ from T2K + NO ν A



~ 1.8 sigma preference for NSI

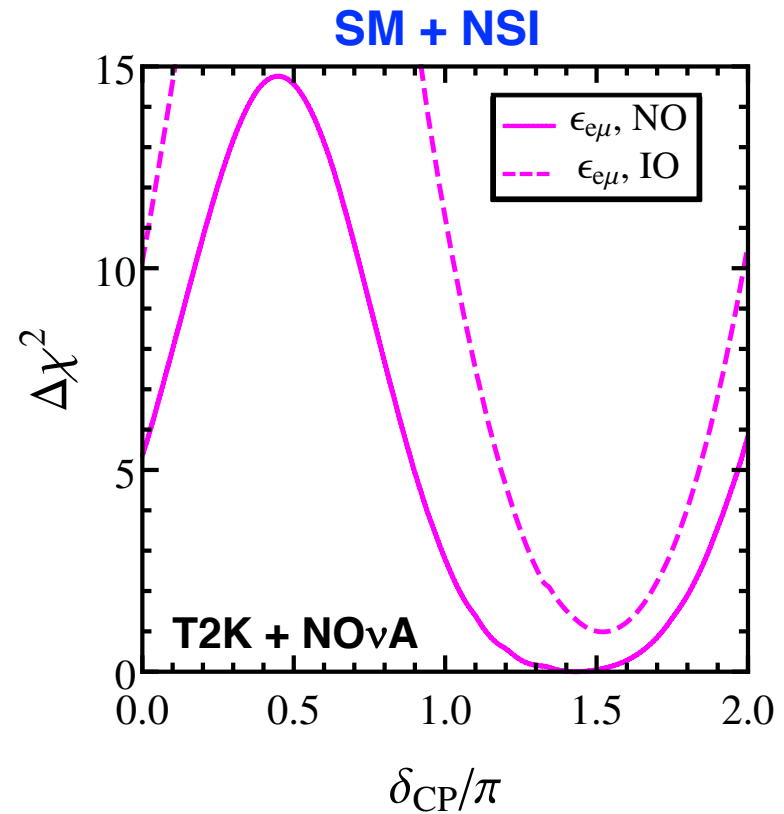
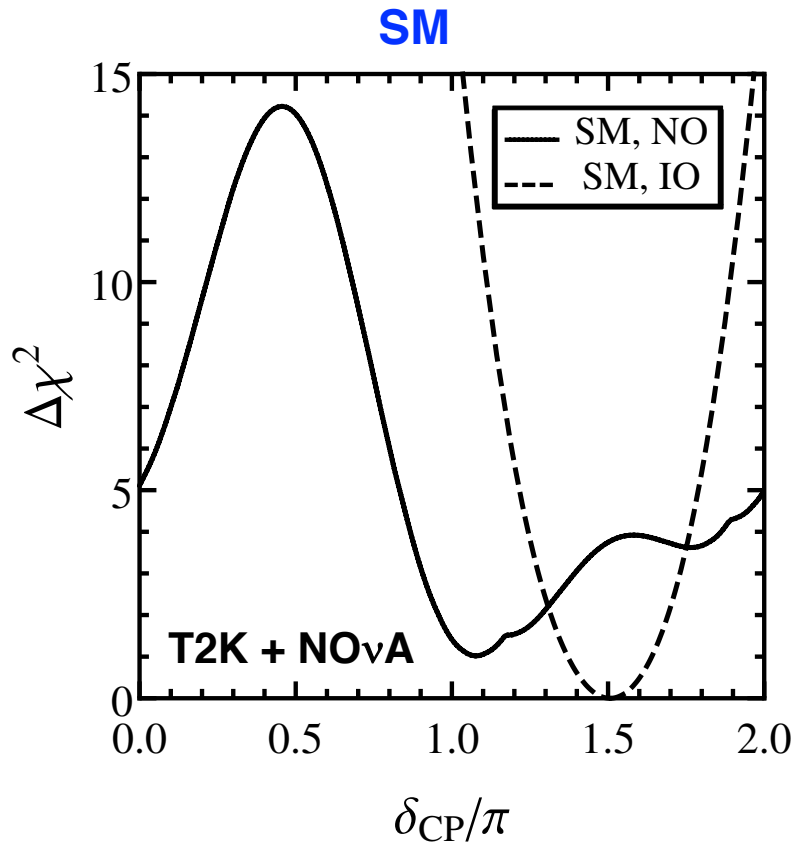
Bievents plots in the presence of NSI



Best fit of $[\delta_{CP}, \phi_{e\mu}, \varepsilon_{e\mu}]$ of combination of T2K + NOVA

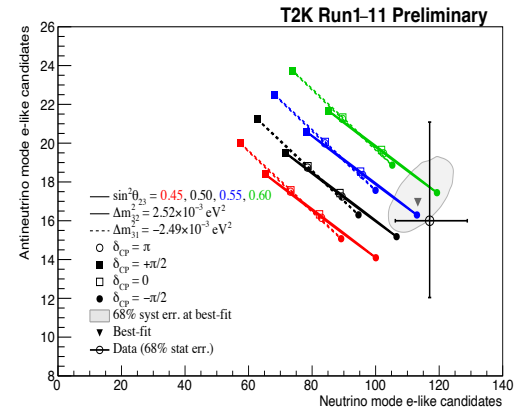
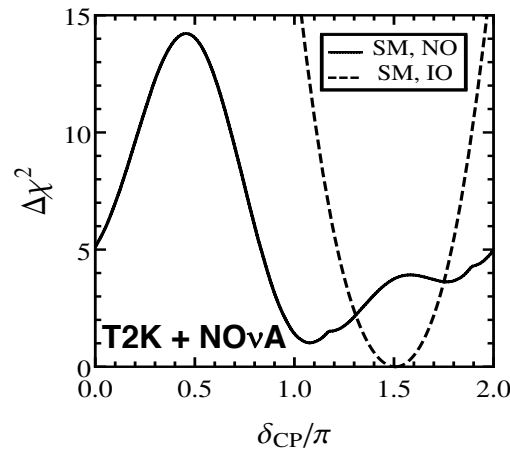
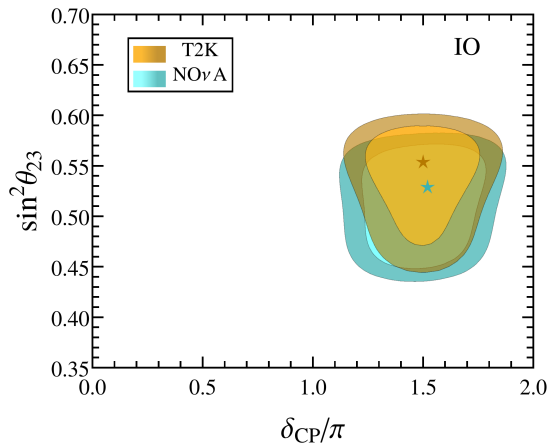
- ★ $\delta_{CP} = 1.1\pi$ compromise between 1.5π and 0.9π
- $\delta_{CP} = 1.44\pi$ almost no need of compromise

NSI restore the preference for NO



Better agreement with all the other data

Can the tension be resolved assuming IO?



For IO the best fit of δ_{CP} is the same in T2K and NOvA (left panel).

However, IO gains only $\chi^2_{IO} - \chi^2_{NO} \sim -1$ in T2K + NOvA combination (middle panel).

The reason is that T2K disfavors IO (dotted ellipses) (right panel).

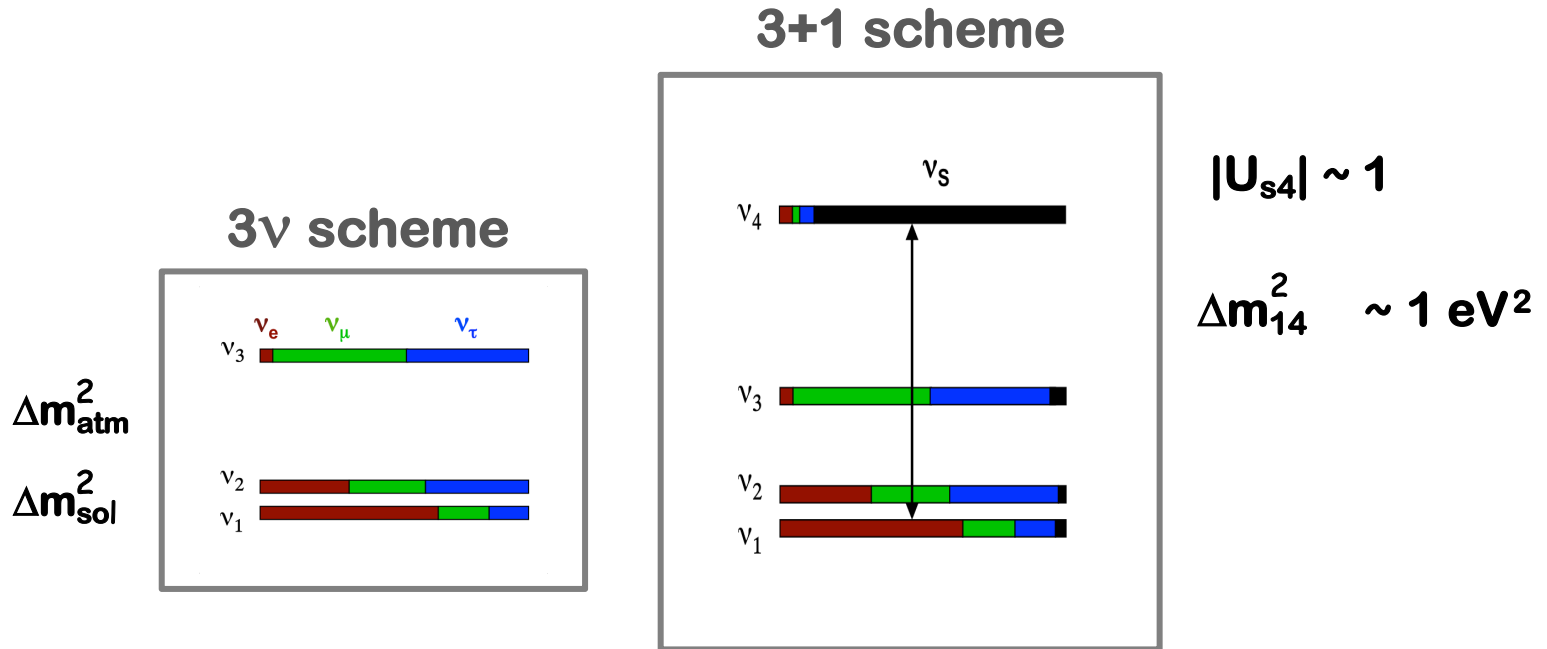
T2K and NOvA disappearance channel + Reactors prefer NO ($\chi^2_{IO} - \chi^2_{NO} \sim 4$).

SK atmospheric data (ν 2020) prefer NO ($\chi^2_{IO} - \chi^2_{NO} \sim 3$).

Therefore, IO seems not to be the favored solution
but I think it is still premature to discard it

Can sterile neutrinos resolve the tension?

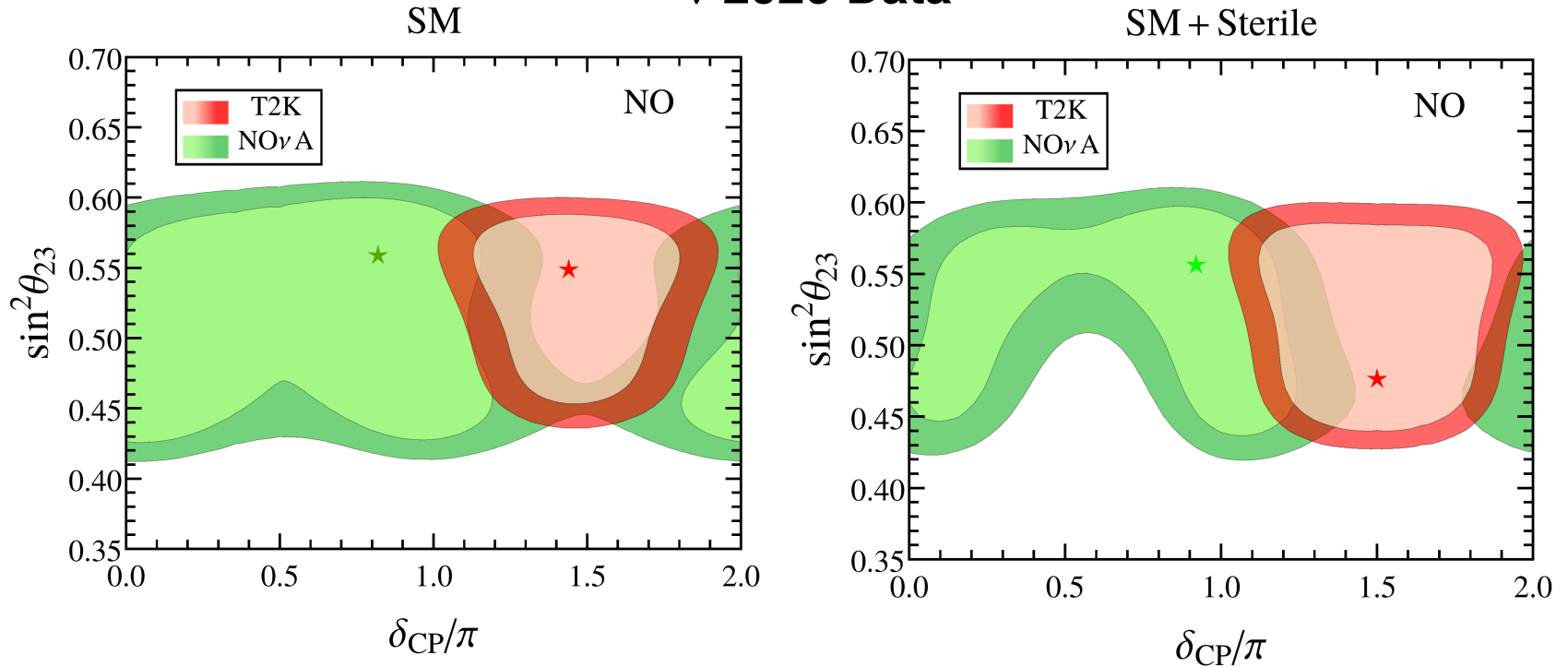
Chatterjee & Palazzo arXiv:2005.10338



At LBL the effective 2-flavor SBL description is no more valid and calculations should be done in the 3+1 (or 3+N_s) scheme

Sterile don't work: tension is still there!

v 2020 Data



Why?

Sterile vs NSI:

Two different kinds of interference

Sterile
Kinematical
Same
amplitude in
NOvA / T2K

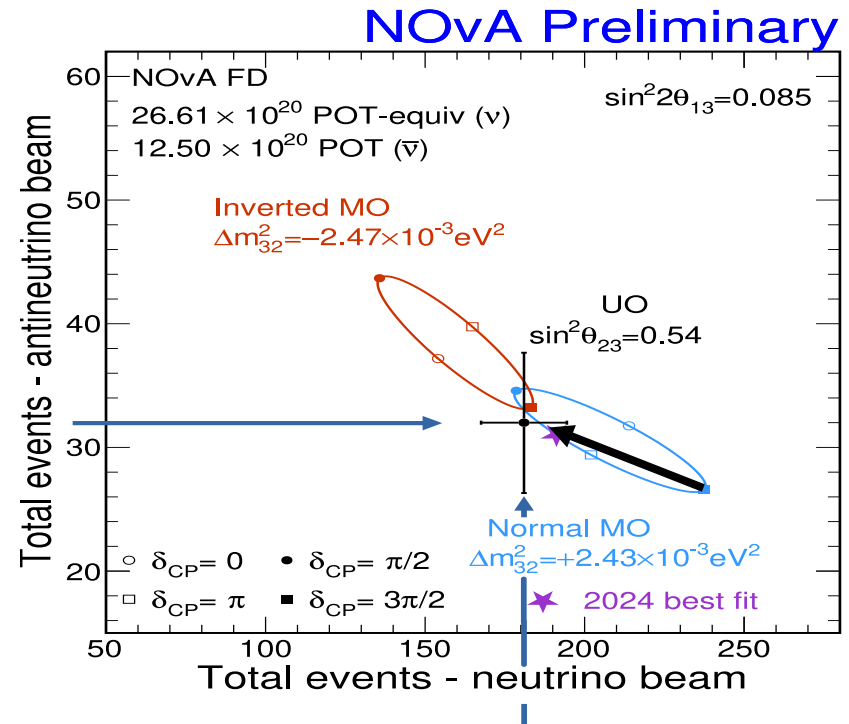
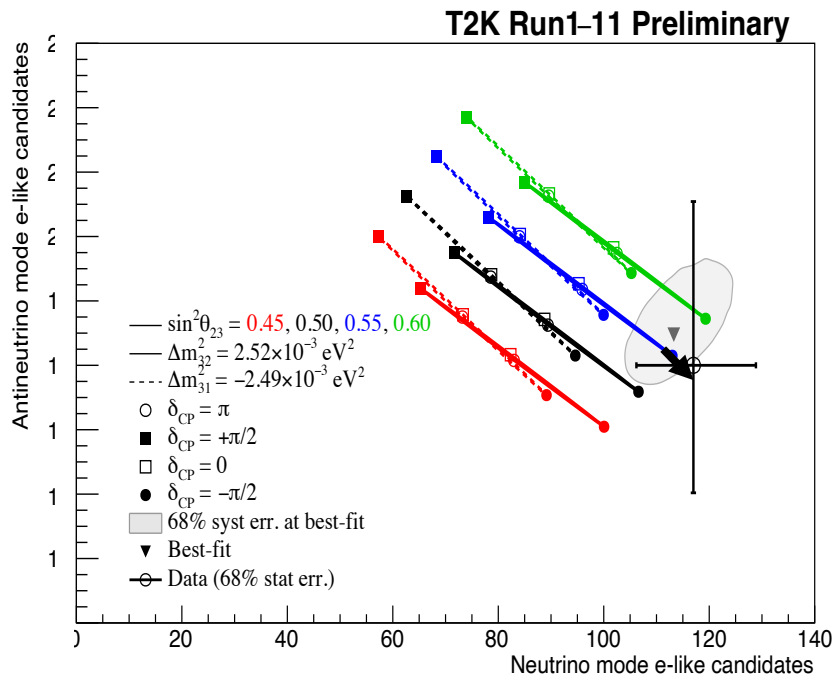
$$\left\{ \begin{array}{l} P^{\text{ATM}} \simeq 4s_{23}^2 s_{13}^2 \sin^2 \Delta \\ P_{\text{I}}^{\text{INT}} \simeq 8s_{13}s_{23}c_{23}s_{12}c_{12}(\alpha\Delta) \sin \Delta \cos(\Delta + \delta_{13}) \\ P_{\text{II}}^{\text{INT}} \simeq 4\underline{s_{14}}\underline{s_{24}}\underline{s_{13}}s_{23} \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14}) \end{array} \right.$$

NSI
Dynamical
Different
amplitude in
NOvA / T2K

$$\left\{ \begin{array}{l} P_0 \simeq 4s_{13}^2 s_{23}^2 f^2 \\ P_1 \simeq 8s_{13}s_{12}c_{12}s_{23}c_{23}\underline{\alpha}fg \cos(\Delta + \delta_{\text{CP}}) \\ P_2 \simeq 8\underline{s_{13}}\underline{s_{23}}\underline{v}|\underline{\varepsilon}_{\alpha\beta}|[af^2 \cos(\delta_{\text{CP}} + \phi_{\alpha\beta}) + bfg \cos(\Delta + \delta_{\text{CP}} + \phi_{\alpha\beta})] \end{array} \right.$$

$$v = \frac{2V_{\text{CC}}E}{\Delta m_{31}^2} = 0.18 \left[\frac{E}{2.0 \text{ GeV}} \right]$$

Gaining insight with bievents plots



We need a noticeable displacement in NOvA

while no displacement is needed in T2K

This is possible only with NSI (different values of ν)

Conclusions

T2K and NO_vA display a tension at ~2.5 sigma level

Complex flavor-changing NSI can solve the tension for $\varepsilon \sim 0.1$

Sterile neutrinos are not able to do the same job

Dynamical (NSI) vs kinematical (sterile) mechanism

Various other mechanisms have been investigated (non-unitarity, LV, ultra-light DM, ...) none satisfactory

New information may come in a few years from T2K, NO_vA, ANTARES and Icecube

If the NSI indication persists, DUNE and HK will definitely confirm/disconfirm it.

Back up slides

Thank you for your attention

Analytical expectations with NSI

$P_{\mu e}$ involves 4 small quantities

$$s_{13} = 0.15 \quad \epsilon \quad v = \frac{2V_{CC}E}{\Delta m_{31}^2} = 0.18 \left[\frac{E}{2.0 \text{ GeV}} \right] \quad \epsilon$$

$$\alpha = 0.03 \quad \epsilon^2 \quad |\epsilon_{\alpha\beta}| \sim 0.2 \quad \epsilon$$

$P_{\mu e}$ is the sum of three terms

$$P_{\mu e} \simeq \underbrace{P_0 + P_1}_{\text{SM}} + \underbrace{P_2}_{\text{NSI}}$$

T2K $v \sim 0.05$
NOvA $v \sim 0.18$

$$P_0 \simeq 4s_{13}^2 s_{23}^2 f^2 \quad \epsilon^2$$

$$P_1 \simeq 8s_{13}s_{12}c_{12}s_{23}c_{23}\alpha fg \cos(\Delta + \delta_{CP}) \quad \epsilon^3$$

$$P_2 \simeq 8s_{13}s_{23}v|\epsilon_{\alpha\beta}|[af^2 \cos(\delta_{CP} + \phi_{\alpha\beta}) + bfg \cos(\Delta + \delta_{CP} + \phi_{\alpha\beta})] \quad \epsilon^3$$

$$f \equiv \frac{\sin[(1-v)\Delta]}{1-v}, \quad g \equiv \frac{\sin v\Delta}{v}$$

$$a = s_{23}^2, \quad b = c_{23}^2 \quad \text{if } \alpha\beta = e\mu$$

$$a = s_{23}c_{23}, \quad b = -s_{23}c_{23} \quad \text{if } \alpha\beta = e\tau$$

$$\nu \rightarrow \bar{\nu} \quad [v, \delta_{CP}, \phi_{\alpha\beta}] \rightarrow [-v, -\delta_{CP}, -\phi_{\alpha\beta}]$$

P_2 brings one additional CP-phase $\phi_{\alpha\beta}$

Parametric curve in biprobability plot:

$$[x, y] = [P_{\mu e}, \bar{P}_{\mu e}]$$

- For fixed $\phi_{\alpha\beta} \rightarrow$ ellipse for varying δ_{CP}
- For fixed $\delta_{CP} \rightarrow$ ellipse for varying $\phi_{\alpha\beta}$

What theory says about NSI?

T2K and NO ν A point to effective couplings of about 0.2. These can be obtained with fundamental couplings on electrons, u and d quarks of a few %. This is still a large number from a theoretical perspective.

Neutrinos are components of an SU(2)_L doublet. Gauge invariance at high energies implies that NSI operators come together with operators involving charged leptons, on which there are strong constraints from CLFV.

So, it is very difficult to build models with large NSI [Gavela et al. [0809.3451](#)]

Some possibilities:

Heavy mediators { **Tree-level see-saw** [Forero & Huang [1608.04719](#)]
Radiative see-saw [Babu et al. [1907.09498](#)]

Light mediators are an appealing alternative

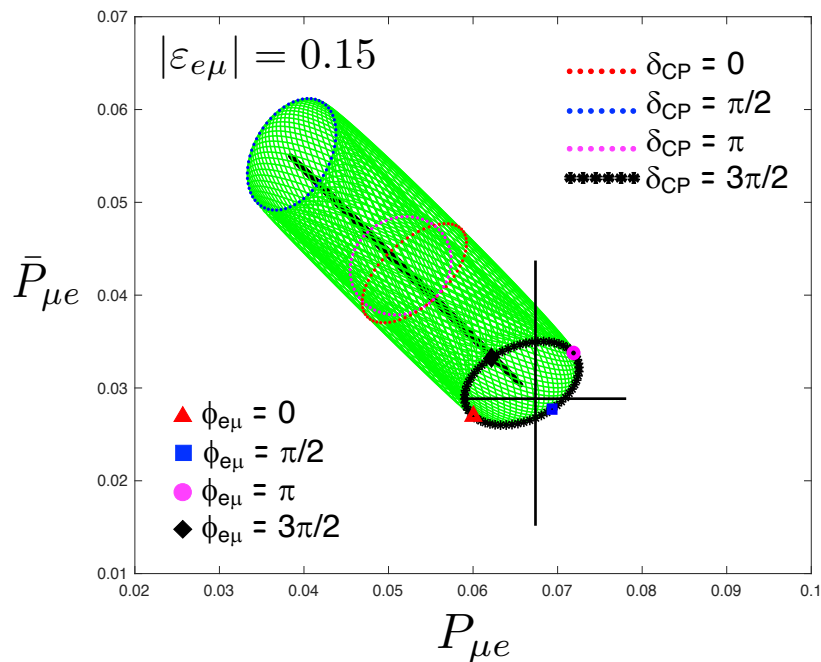
Farzan, Heeck [1607.07616](#) Farzan [1912.09408](#)

Note that forward scattering probes $q^2 = 0$ and a light mediator is felt as an heavy one. Hence, also in this case it is legitimate to describe NSI by an effective dim-6 operator.

Biprobability plots in the presence of NSI

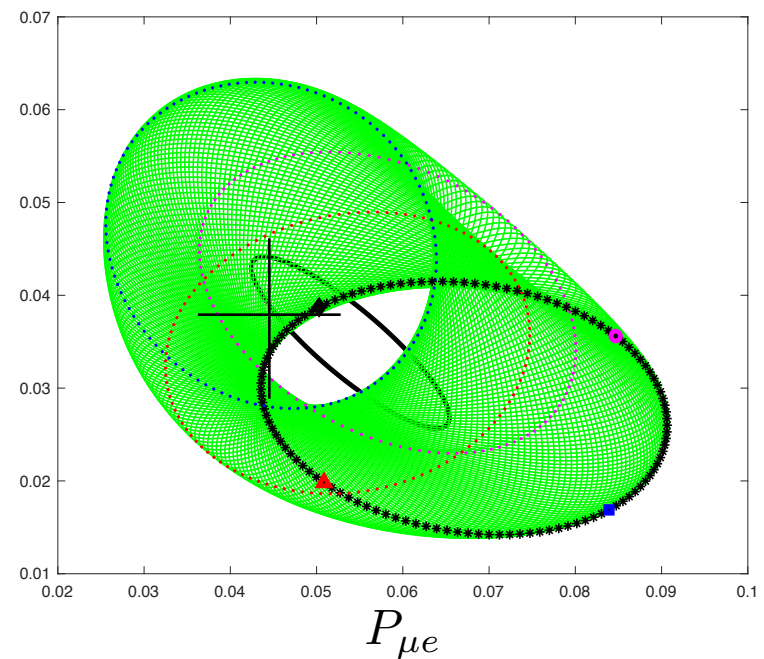
T2K

Strongly favors $\delta_{CP} \sim 3\pi/2$ ellipse
(almost no sensitivity to $\phi_{e\mu}$)



NOvA

In agreement with $\delta_{CP} \sim 3\pi/2$ ellipse.
On this ellipse it pins down $\phi_{e\mu} \sim 3\pi/2$



Mixing Matrix in the 3+1 scheme

$$U = \widetilde{R}_{34} R_{24} \widetilde{R}_{14} R_{23} \underbrace{\widetilde{R}_{13} R_{12}}_{3\nu}$$

$$R_{ij} = \begin{bmatrix} c_{ij} & s_{ij} \\ -s_{ij} & c_{ij} \end{bmatrix}$$

$$\widetilde{R}_{ij} = \begin{bmatrix} c_{ij} & \widetilde{s}_{ij} \\ -\widetilde{s}_{ij}^* & c_{ij} \end{bmatrix}$$

$$\begin{aligned} s_{ij} &= \sin \theta_{ij} \\ c_{ij} &= \cos \theta_{ij} \\ \widetilde{s}_{ij} &= s_{ij} e^{-i\delta_{ij}} \end{aligned}$$

$$3\nu \begin{cases} 3 \text{ mixing angles} \\ 1 \text{ Dirac phase} \\ 2 \text{ Majorana phases} \end{cases}$$

$$3+1 \begin{cases} 6 \\ 3 \\ 3 \end{cases}$$

$$3+N \begin{cases} 3+3N \\ 1+2N \\ 2+N \end{cases}$$

In general, we have additional sources of CPV

LBL transition probability in 3-flavor

$$P_{\nu_{\mu} \rightarrow \nu_e}^{3\nu} = P^{\text{ATM}} + P^{\text{SOL}} + P^{\text{INT}}$$

In vacuum:

$$P^{\text{ATM}} = 4s_{23}^2 s_{13}^2 \sin^2 \Delta$$

$$P^{\text{SOL}} = 4c_{12}^2 c_{23}^2 s_{12}^2 (\alpha \Delta)^2$$

$$P^{\text{INT}} = 8s_{23}s_{13}c_{12}c_{23}s_{12}(\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{CP})$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}, \quad \alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

$$\Delta \sim \pi/2$$

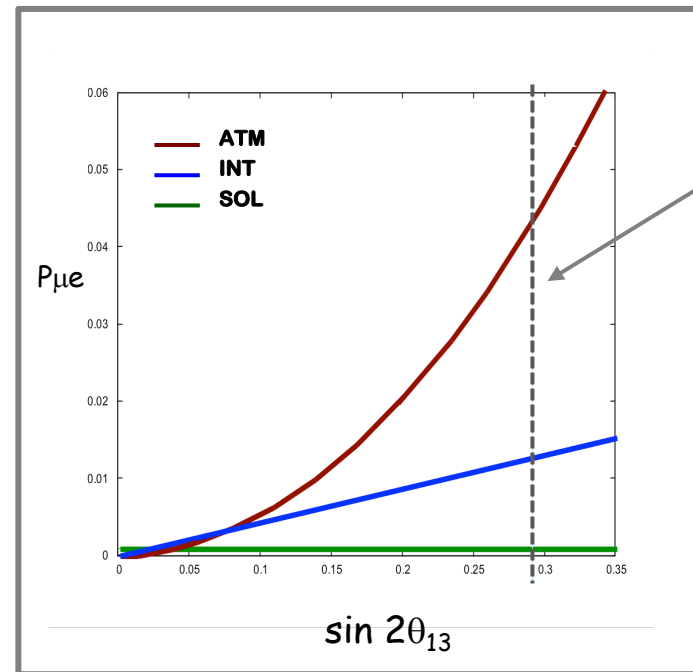
$$\alpha \sim 0.03$$

P^{ATM} leading $\rightarrow \theta_{13} > 0$

P^{INT} subleading \rightarrow dependency on δ

P^{SOL} negligible

T2K osc. maximum E = 0.6 GeV



A new interference term in the 3+1 scheme

N. Klop & A.P., PRD (2015)

- $\Delta_{14} \gg 1$: fast oscillations are averaged out

- But interference of Δ_{14} & Δ_{13} survives and is observable

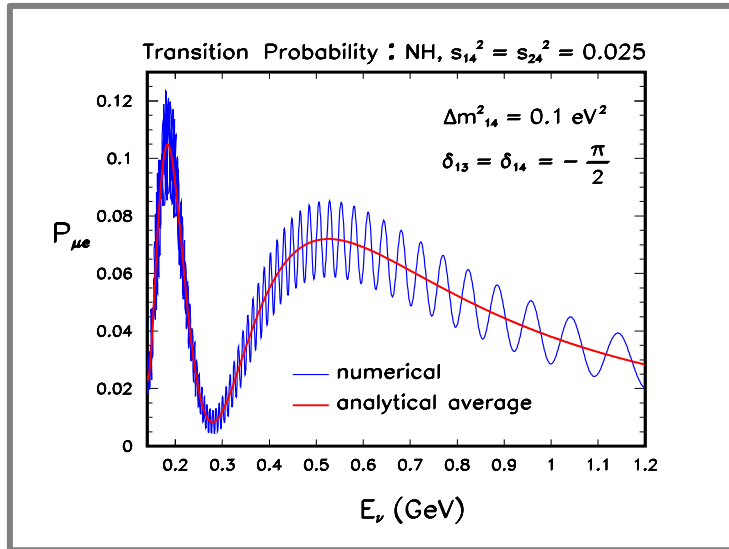
$$P_{\mu e}^{4\nu} \simeq P^{\text{ATM}} + P_{\text{I}}^{\text{INT}} + P_{\text{II}}^{\text{INT}}$$

$$\begin{aligned} S_{13} \sim S_{14} \sim S_{24} &\sim 0.15 \sim \epsilon \\ \alpha = \delta m^2 / \Delta m^2 &\sim 0.03 \sim \epsilon^2 \end{aligned}$$

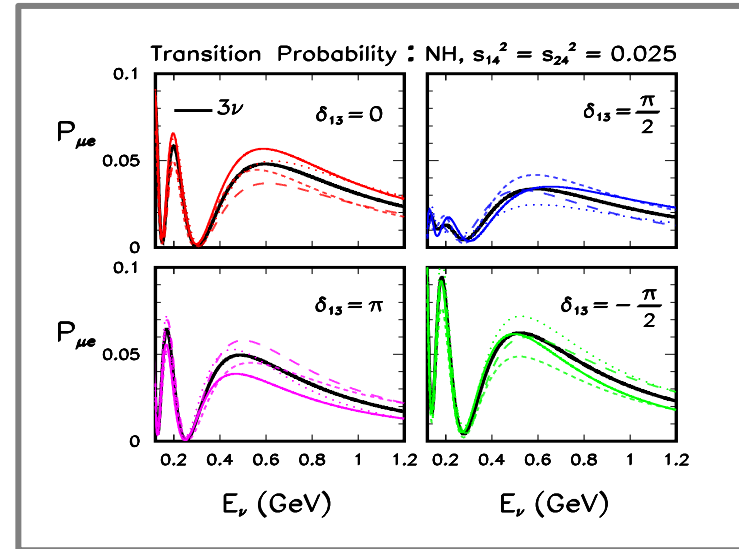
$$\left\{ \begin{aligned} P^{\text{ATM}} &\simeq 4s_{23}^2 \underline{s_{13}^2} \sin^2 \Delta && \sim \epsilon^2 \\ P_{\text{I}}^{\text{INT}} &\simeq 8 \underline{s_{13}} s_{23} c_{23} s_{12} c_{12} (\underline{\alpha \Delta}) \sin \Delta \cos(\Delta + \delta_{13}) && \sim \epsilon^3 \\ P_{\text{II}}^{\text{INT}} &\simeq 4 \underline{s_{14}} \underline{s_{24}} \underline{s_{13}} s_{23} \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14}) && \sim \epsilon^3 \end{aligned} \right.$$

Sensitivity to the new CP-phase δ_{14}

Numerical examples of 4ν probability



The fast oscillations get averaged out due to the finite energy resolution



Different line styles



Different values of δ_{14}

The modifications induced by δ_{14} are almost as large as those induced by the standard CP-phase δ_{13}

In principle we can try to explain the tension with 4ν