

From boost to brake: Dark Matter upscattering by the Diffuse Supernova Neutrino Background

Tim Herbermann

NOW 2024

2403.15367 (with A. Das, M. Sen & V. Takhistov)

2408.02721 (with M. Lindner & M. Sen)

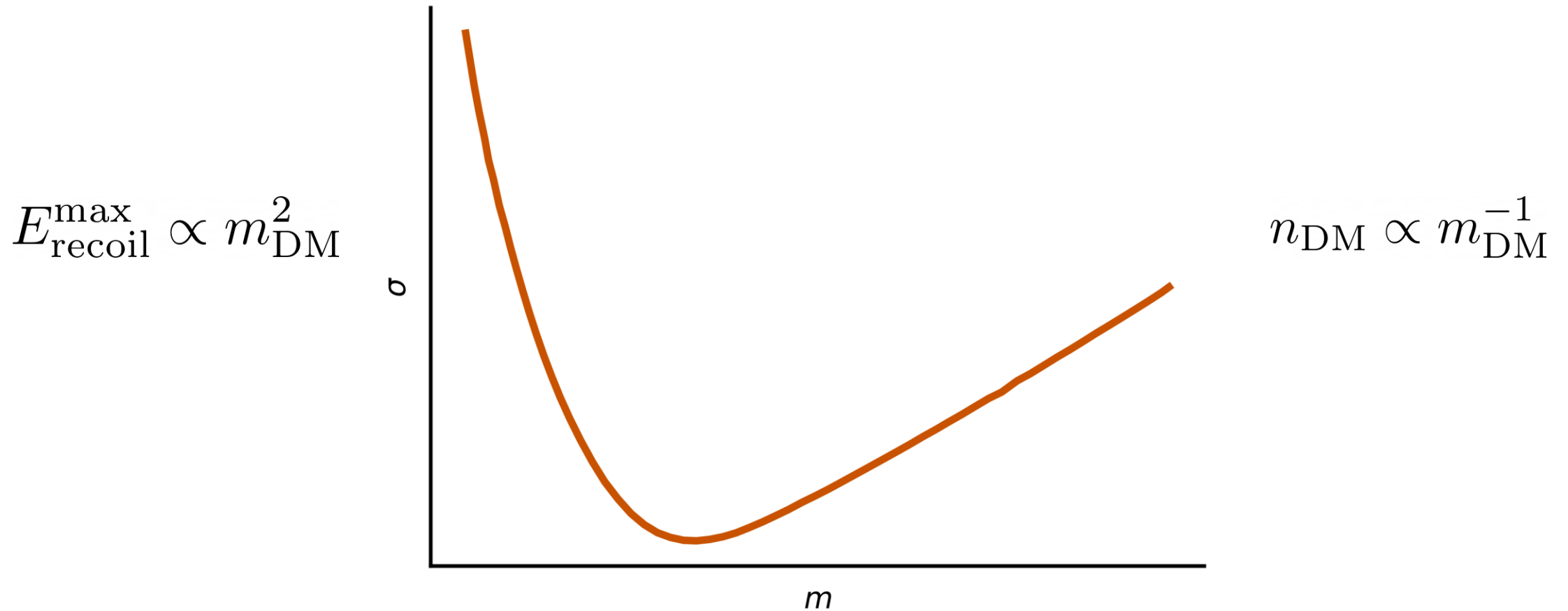


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The problem in direct detection



Boosted DM to the rescue!

- Fraction of DM with boosted kinetic energies
- Recoil energy no longer mass suppressed
- Often irreducible, since we already assume interactions for detection → Comes for free!

$$T_\chi \gg m_\chi v_{\text{gal}}^2$$



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We use neutrinos from Diffuse Supernova Neutrino Background (DSNB) for upscattering of leptophilic DM



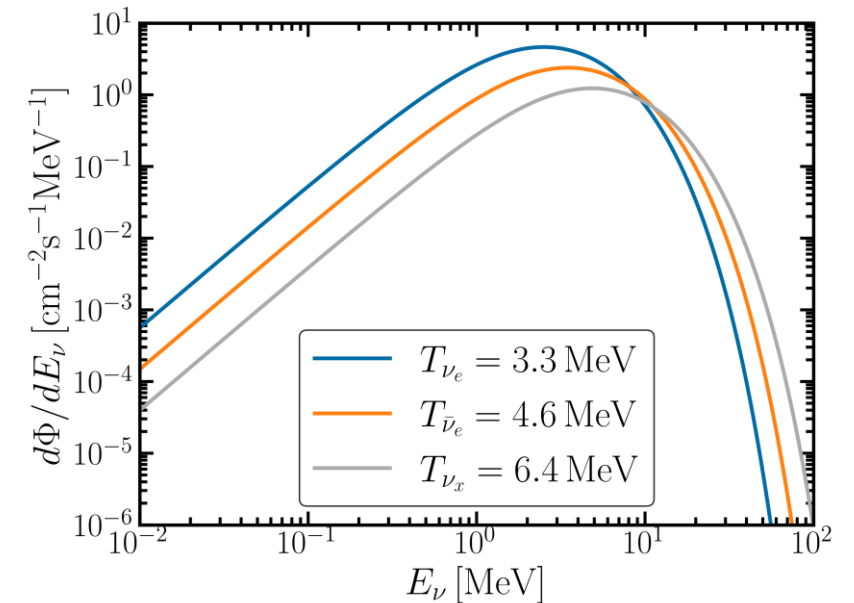
Why boosting with the DSNB?

- Abundant flux at MeV energies

$$T_{\chi}^{\max} = \frac{E_{\nu}^2}{E_{\nu} + m_{\chi}/2}$$

- ν -DM interactions can arise naturally in leptophilic models
- Probe of joint e -DM and ν -DM interaction

See talk by C. Volpe



$$\Phi_{\nu}(E) = \int_0^{z_{\max}} \frac{dz}{H(z)} R_{\text{CCSN}}(z) F_{\nu}(E') \Big|_{E'=E(1+z)}$$

$$F_{\nu}(E) = \frac{E_{\nu}^{\text{tot}}}{6} \frac{120}{7\pi^4} \frac{E^2}{T_{\nu}^4} \frac{1}{e^{E/T_{\nu}} + 1}$$

$$E_{\nu}^{\text{tot}} = 3 \times 10^{53} \text{erg}$$

A. Das & M. Sen (2021),
De Romeri et al (2024)

The boost...



Local upscattering rate
of DM by DSNB



Integrate all lines of sight



$$\frac{d\Phi_\chi}{dT_\chi} = \int \frac{d\Omega}{4\pi} \int_{\text{l.o.s.}} dl \int_{E_\nu^{\min}}^{E_\nu^{\max}} dE_\nu \frac{\rho_\chi(l)}{m_\chi} \frac{d\Phi_\nu}{dE_\nu} \frac{d\sigma_{\nu\chi}}{dT_\chi}$$

... and the brake

Underground direct detection sites

Energy loss from elastic scattering
before detection leads to flux attenuation

$$\frac{dT_\chi}{dx}(x) = - \sum_i n_i(x) \int_0^{T_i^{\max}} dT_i T_i \frac{d\sigma_{i\chi}}{dT_i}$$

Conservative: Purely elastic, no multiple scattering, no deflection



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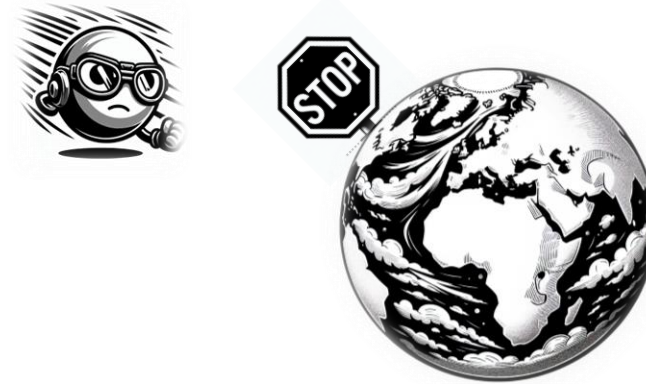


Previously: Analytic approximation

- Constant cross section only
- Assumptions on kinematics

Does not apply here!

... and the brake



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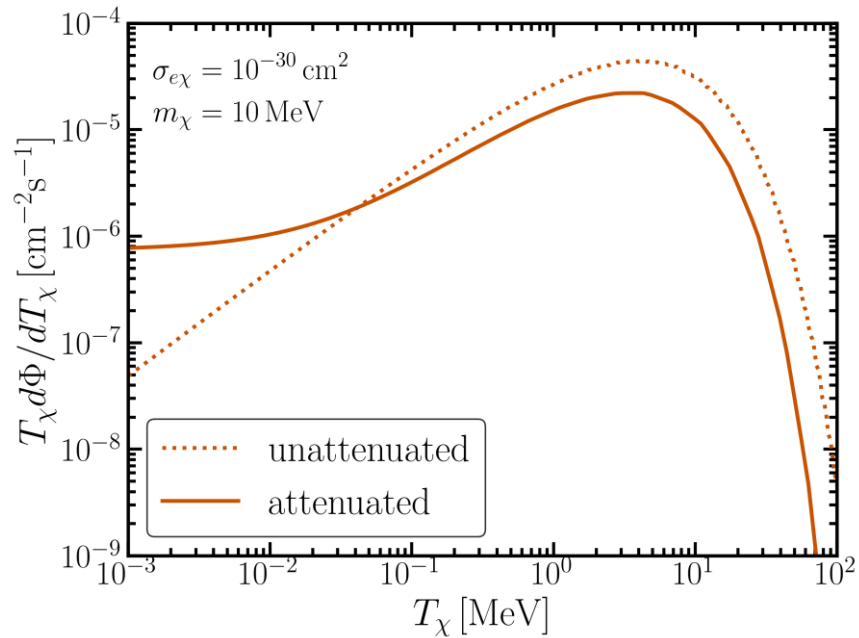
Conservative: Purely elastic, no multiple scattering, no deflection

NEW: Full numerical solution

- No additional approximations
- Enables model dependent studies
- Numerically challenging

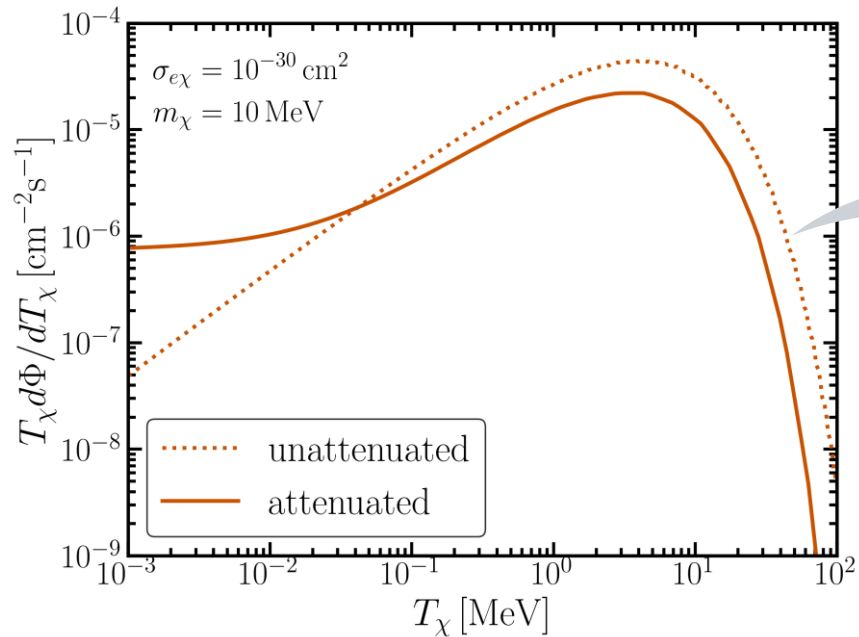
Effect of attenuation

Attenuation for constant cross section

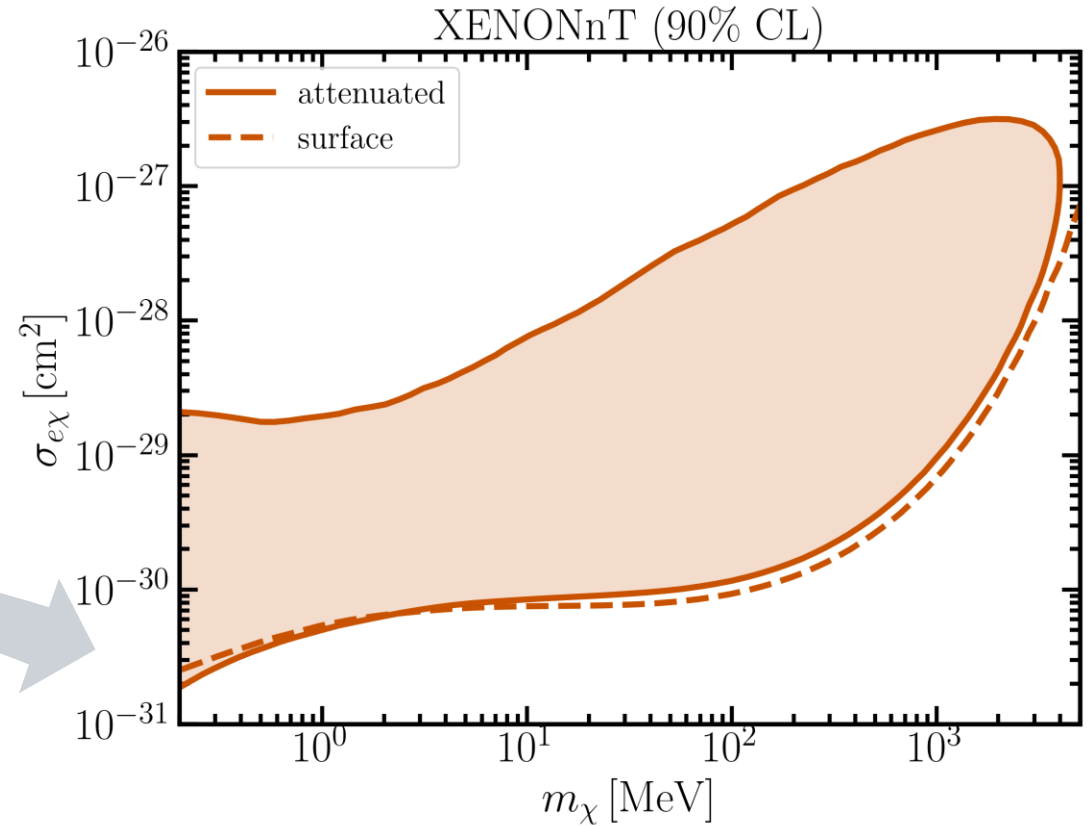


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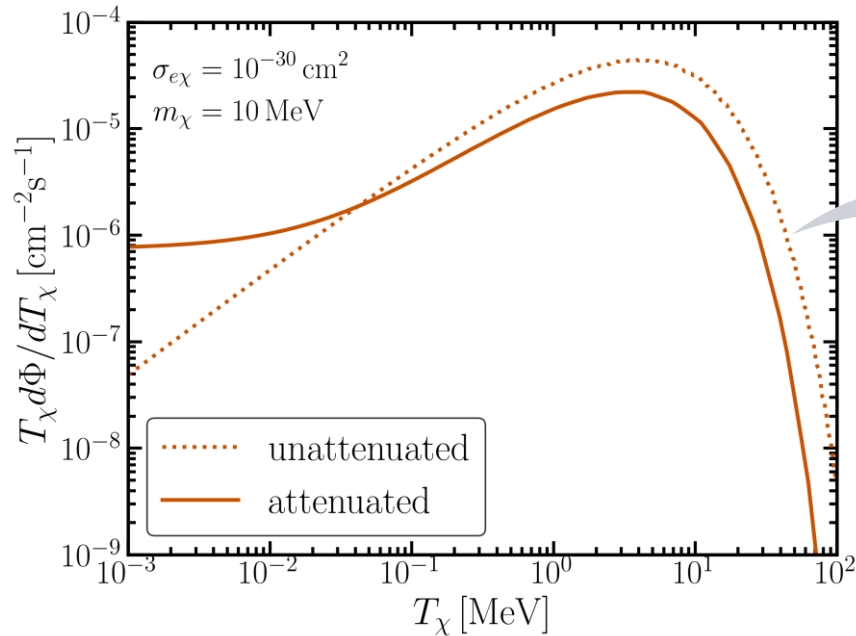


$$\text{+} \quad \frac{dR}{dT_e} = N_e \int dT_\chi \frac{d\Phi_\chi}{dT_\chi} \frac{d\sigma_{e\chi}}{dT_e}$$

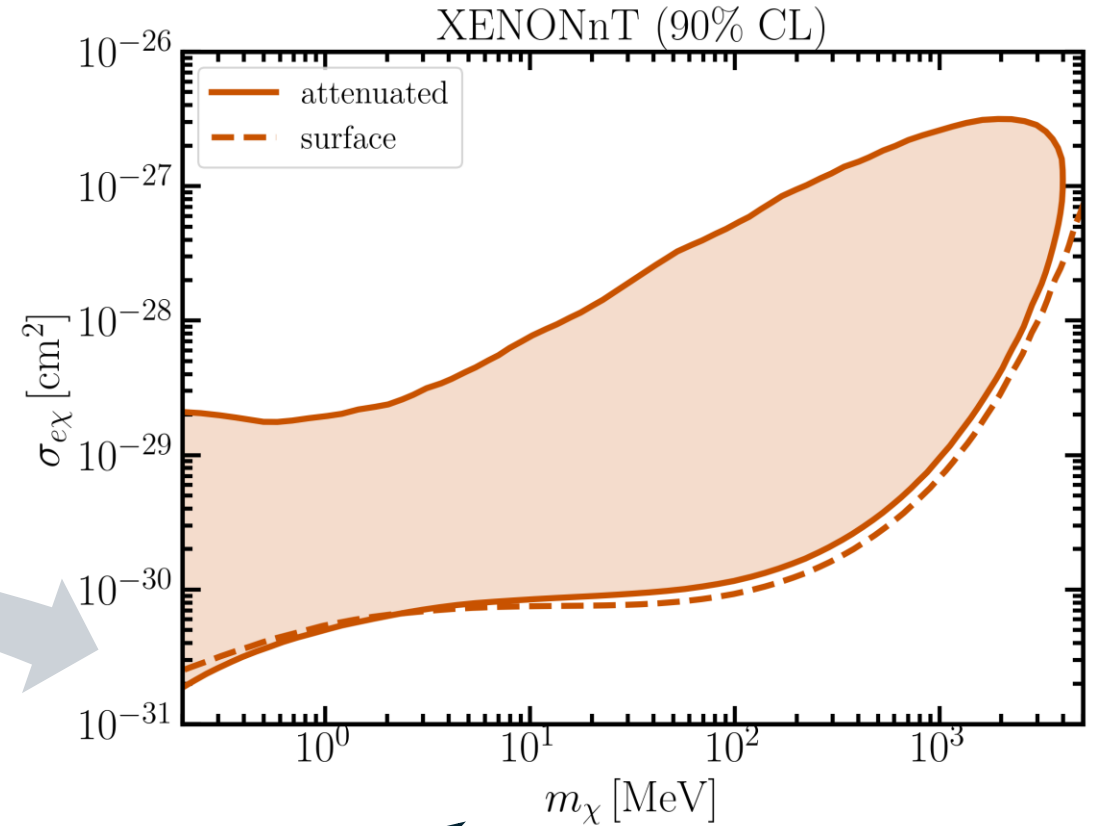


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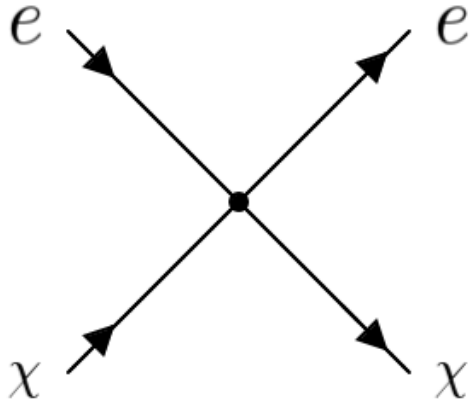


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Attenuation ceiling in tension with approximate solutions highlighting the necessity of a full solution.

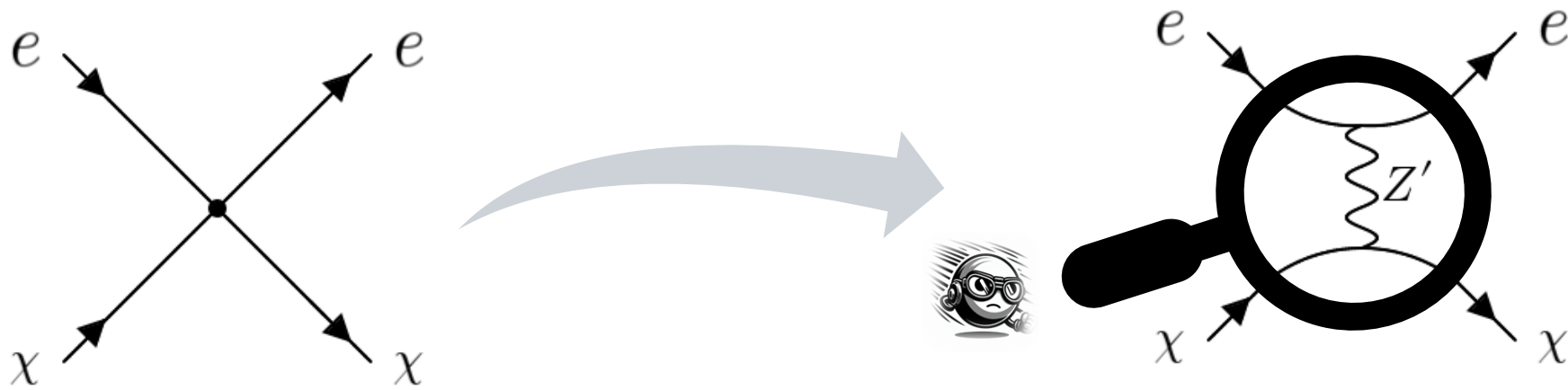
The role of energy dependence



Non relativistic DM mapped on
effective cross section

$$\bar{\sigma}_{e\chi} = \frac{g_e^2 g_\chi^2}{\pi} \frac{\mu_{e\chi}^2}{(q_{\text{ref}}^2 + M_{\text{med}}^2)^2}, \quad q_{\text{ref}} = \alpha m_e$$

The role of energy dependence



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Boosted DM (semi-)relativistic

$$q \gg \alpha m_e$$

Are we sensitive to underlying interaction?

Energy dependence matters!

Limits for leptophilic Z' toy-model

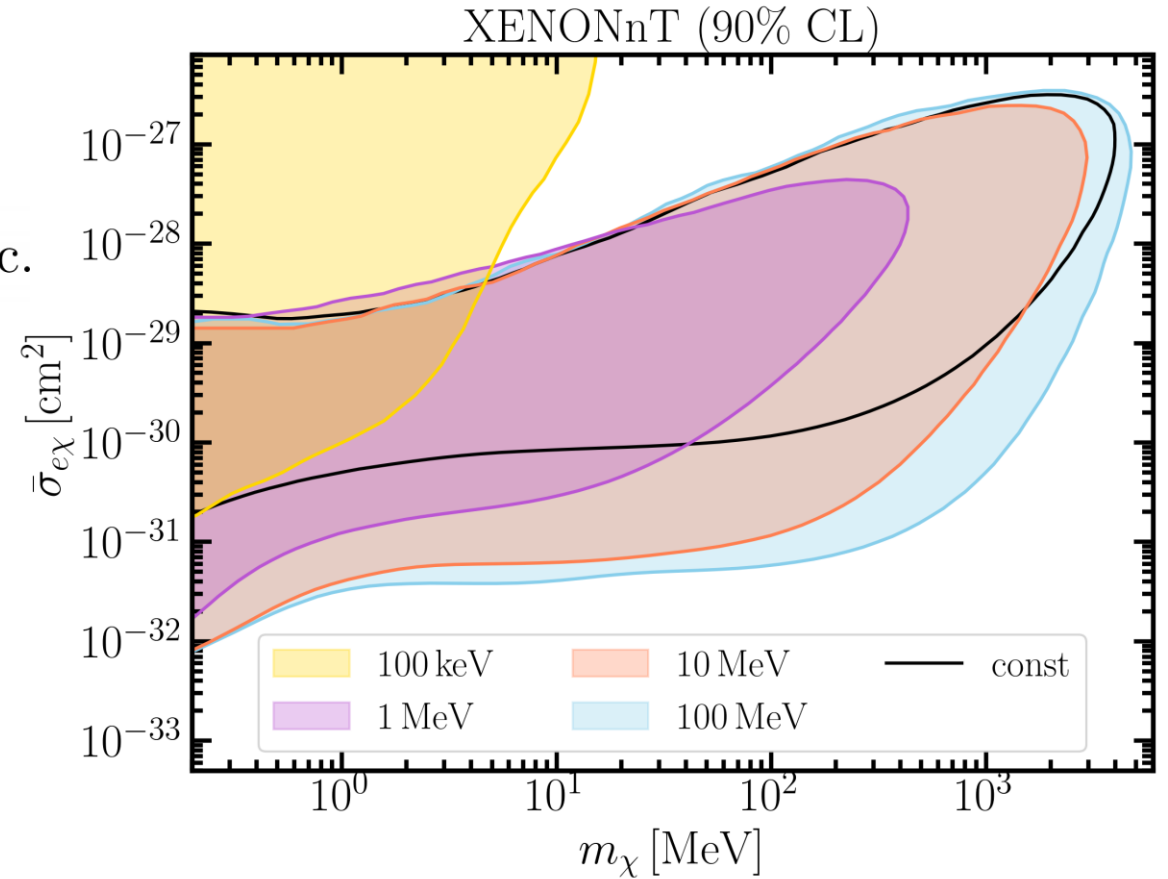
$$\mathcal{L} \supset \frac{1}{2} m_{Z'}^2 Z'_\mu Z'^\mu + g \bar{L} \gamma_\mu L Z'^\mu + g_\chi \bar{\chi} \gamma_\mu \chi Z'^\mu + \text{h.c.}$$

and map on effective cross section to compare

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07.09.2024

Tim Herbermann



JCAP 07 (2024), 045
(A. Das, TH, M. Sen & V. Takhistov)

15

Energy dependence matters!

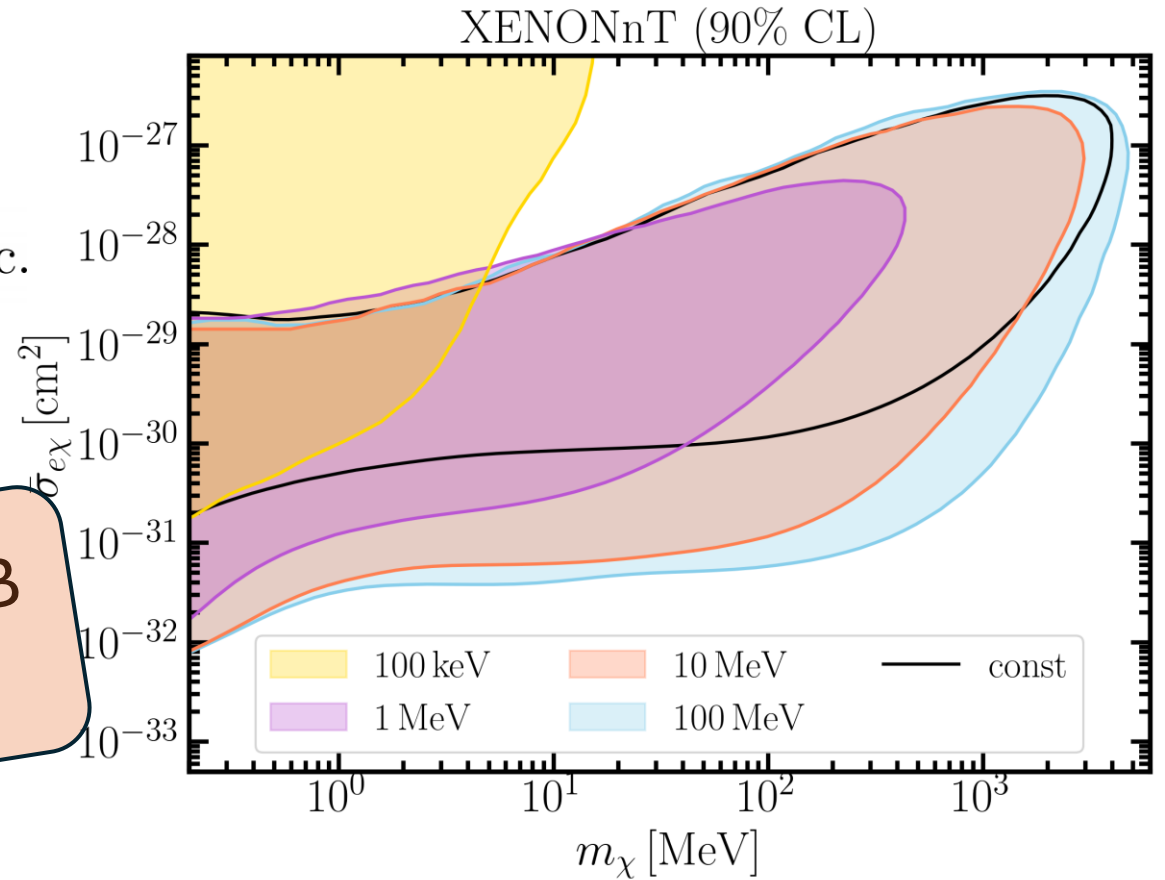
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This also applies beyond DSNB boosted leptophilic DM

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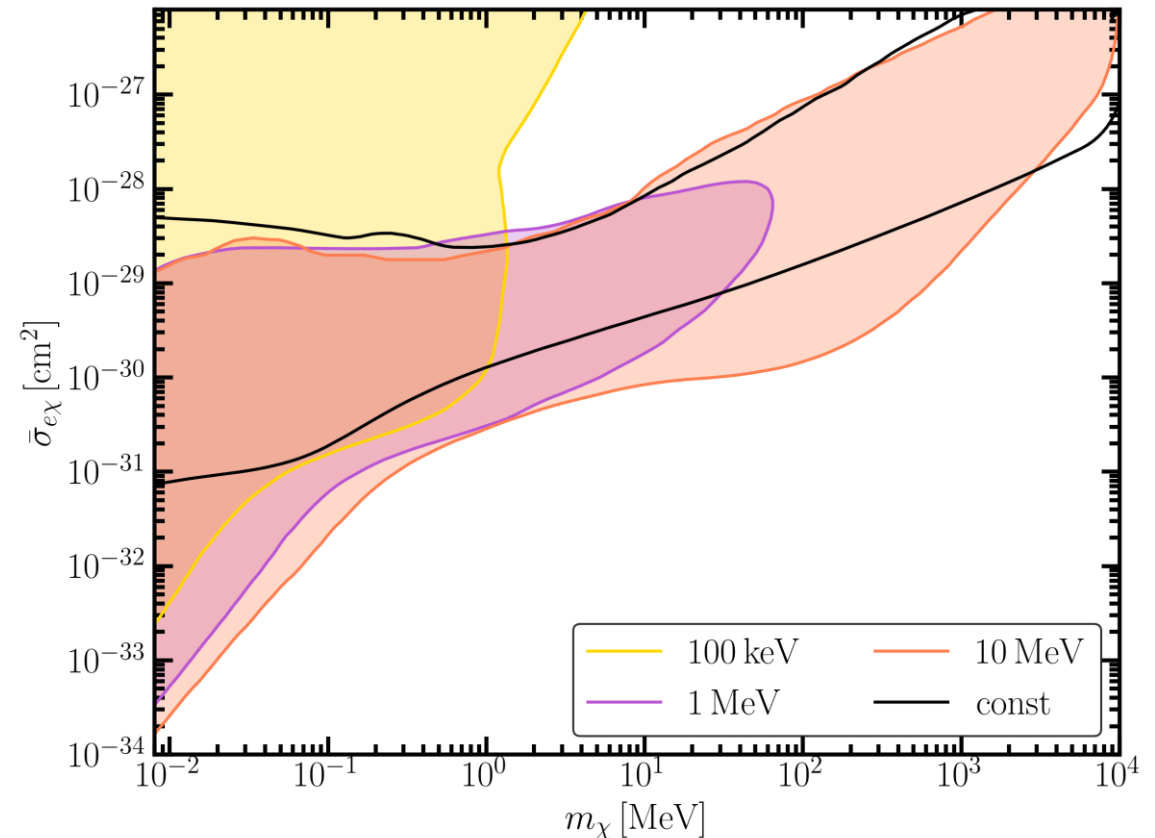


JCAP 07 (2024), 045
(A. Das, TH, M. Sen & V. Takhistov)

Beyond DSNB?

Energy dependence and
attenuation general features
of (boosted) DM

Indeed, we found similar
results for cosmic ray
electron boosted DM!



arXiv: 2408.02721
(TH, M. Lindner, M. Sen)

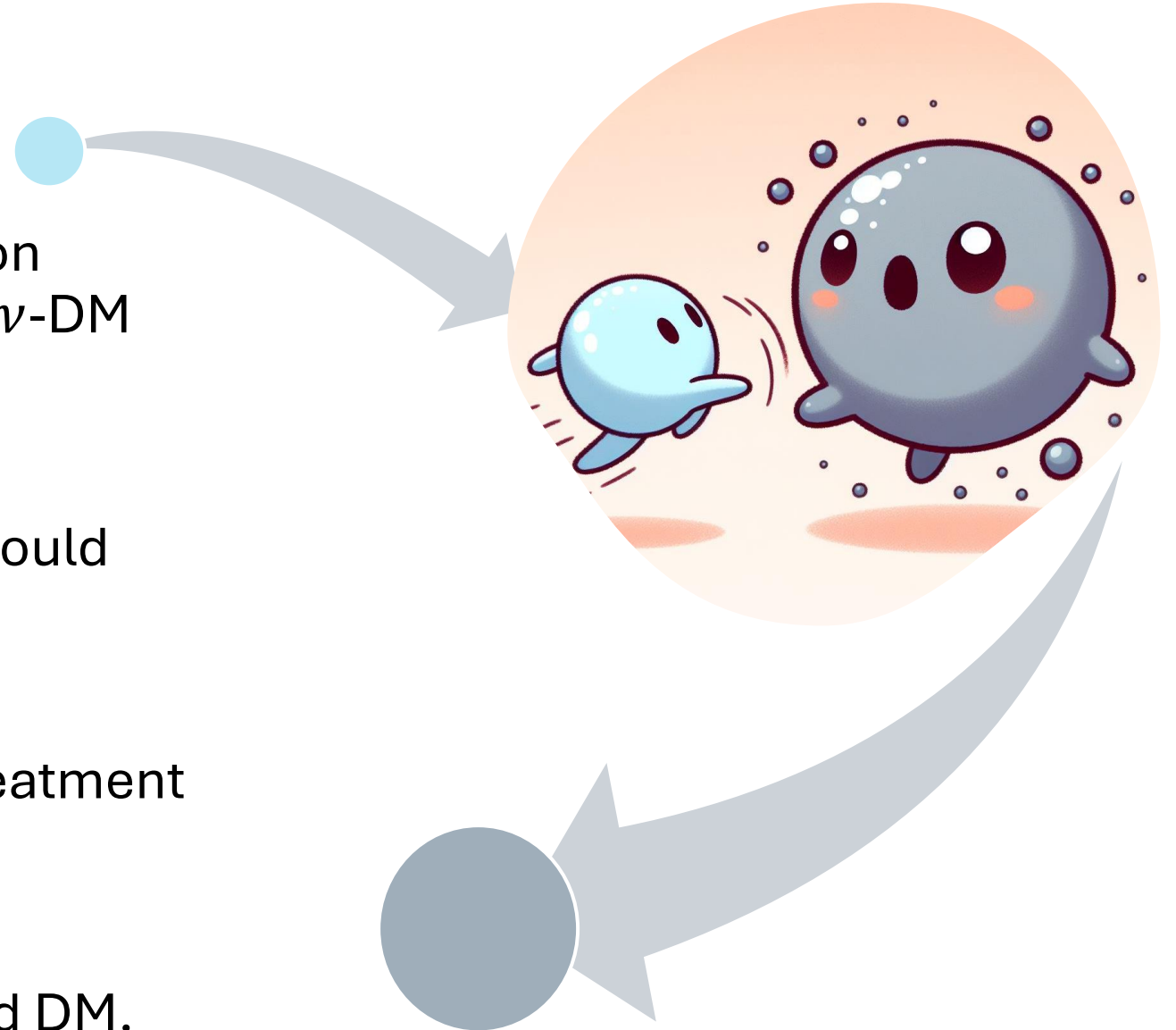
Summary

Unique testing ground for DM-lepton interactions complementing other ν -DM probes.

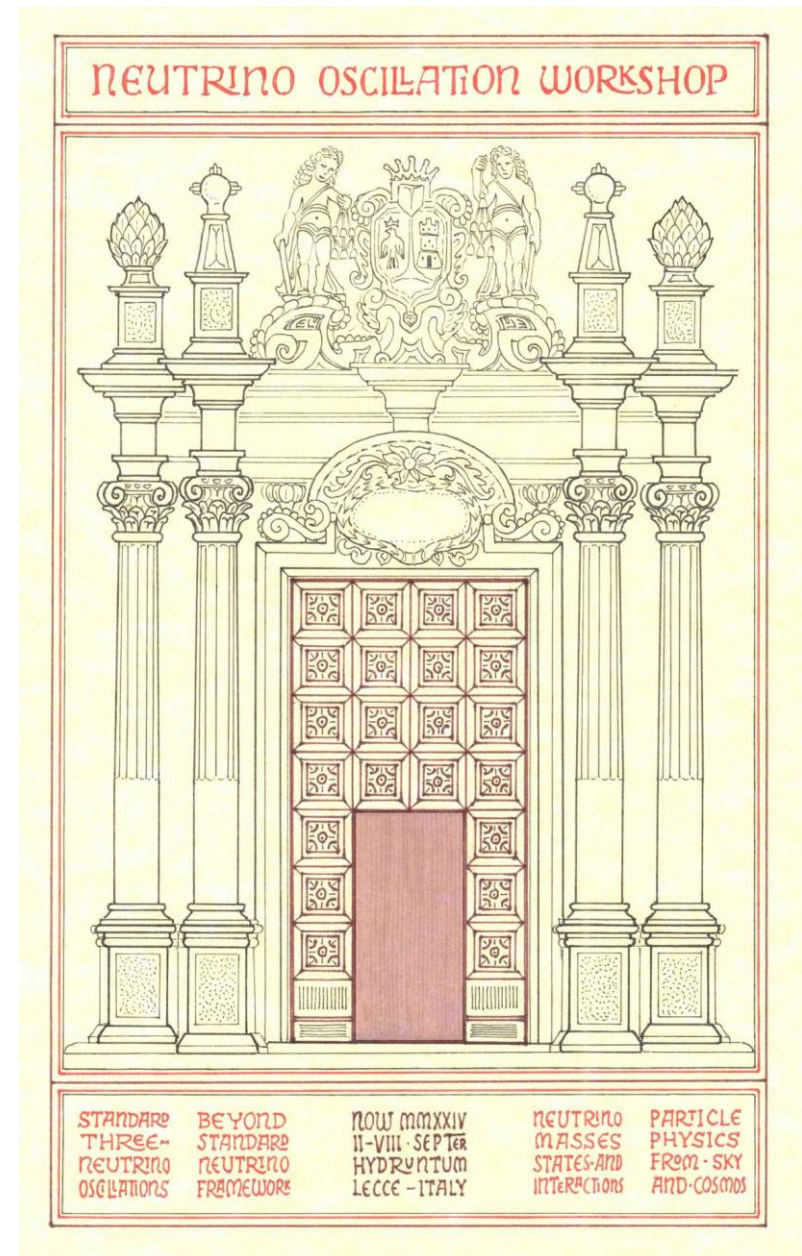
Boosted DM comes for free - we should exploit it.

Energy dependence and correct treatment of attenuation matter.

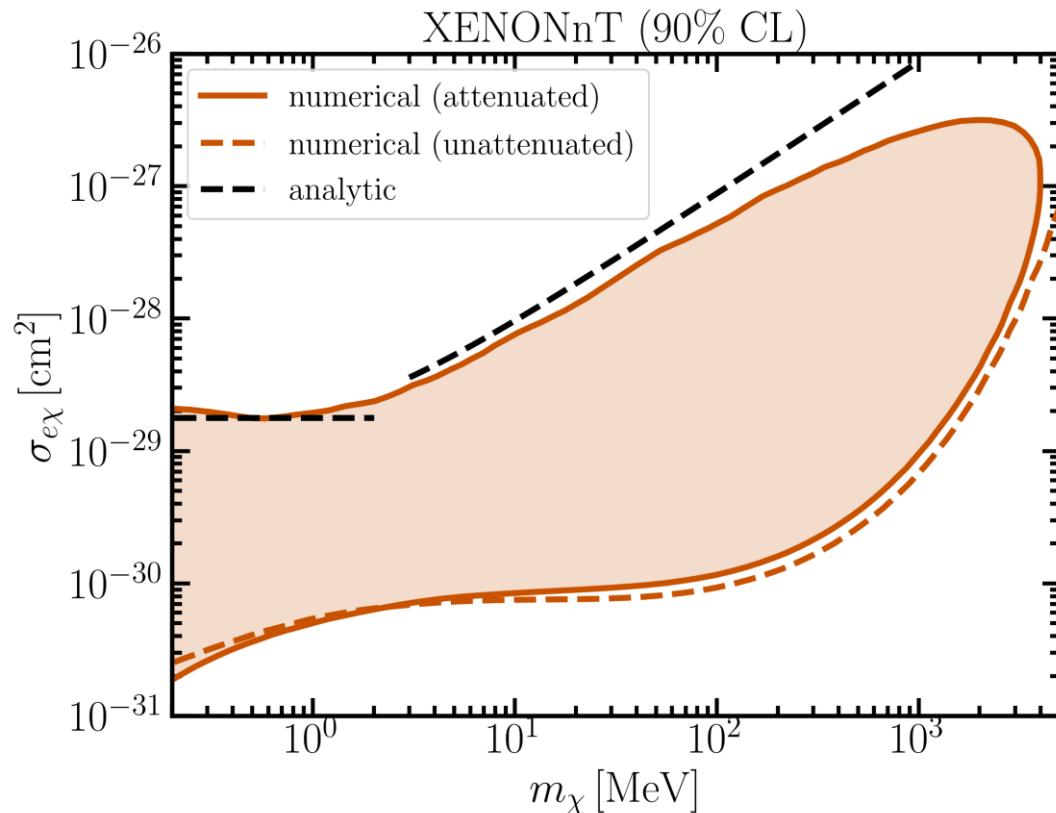
Implications beyond DSNB boosted DM.



Thank you!



Understanding attenuation



Approximate for constant cross section and detector at depth d

i. $m_\chi \gg m_e$

$$\frac{dT_\chi}{dx} = -\frac{1}{2m_\chi\lambda_{\text{eff}}}(T_\chi^2 + 2m_\chi T_\chi), \quad \lambda_{\text{eff}}^{-1} = 4n_e\bar{\sigma}_{e\chi}m_em_\chi/(m_e + m_\chi)^2 \propto m_\chi^{-1}$$

$$d/\lambda_{\text{eff}} \approx 1 \quad \Rightarrow \quad \sigma_{e\chi} \propto m_\chi$$

ii. $m_\chi \sim m_e$

$$\frac{dT_\chi}{dx} \approx -\frac{1}{2}n_e\bar{\sigma}_{e\chi}T_\chi \quad \Rightarrow \quad \sigma_{e\chi} \approx \text{const}$$