



Energy reconstruction and final-state interactions in proton knockout

Alexis Nikolakopoulos

NOW 2024,

7th September 2024

Neutrino-nucleus cross section in oscillation analysis

A measurement in a neutrino experiment:

$$\frac{dN}{dX} = \int dE_\nu \Phi(E_\nu) K(E_\nu, X)$$

Is an integral transform with complicated kernel

$$\frac{dN}{d|\vec{p}_1|} = \int \underbrace{d\Omega_1 d\vec{p}_2 \cdots d\vec{p}_n}_{\text{Phase space(E)}} \underbrace{\epsilon(\vec{p}_1, \cdots, \vec{p}_n)}_{\text{efficiency}} \int \underbrace{dE_\nu \Phi(E_\nu)}_{\text{Flux}} \underbrace{\frac{d\sigma(E_\nu)}{d\vec{p}_1 \cdots d\vec{p}_n}}_{\text{Cross section}}$$

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Oscillation analysis

Constrain and determine uncertainty on $K(E, X)$

Compare predictions with different $\phi(E)$ to measurement

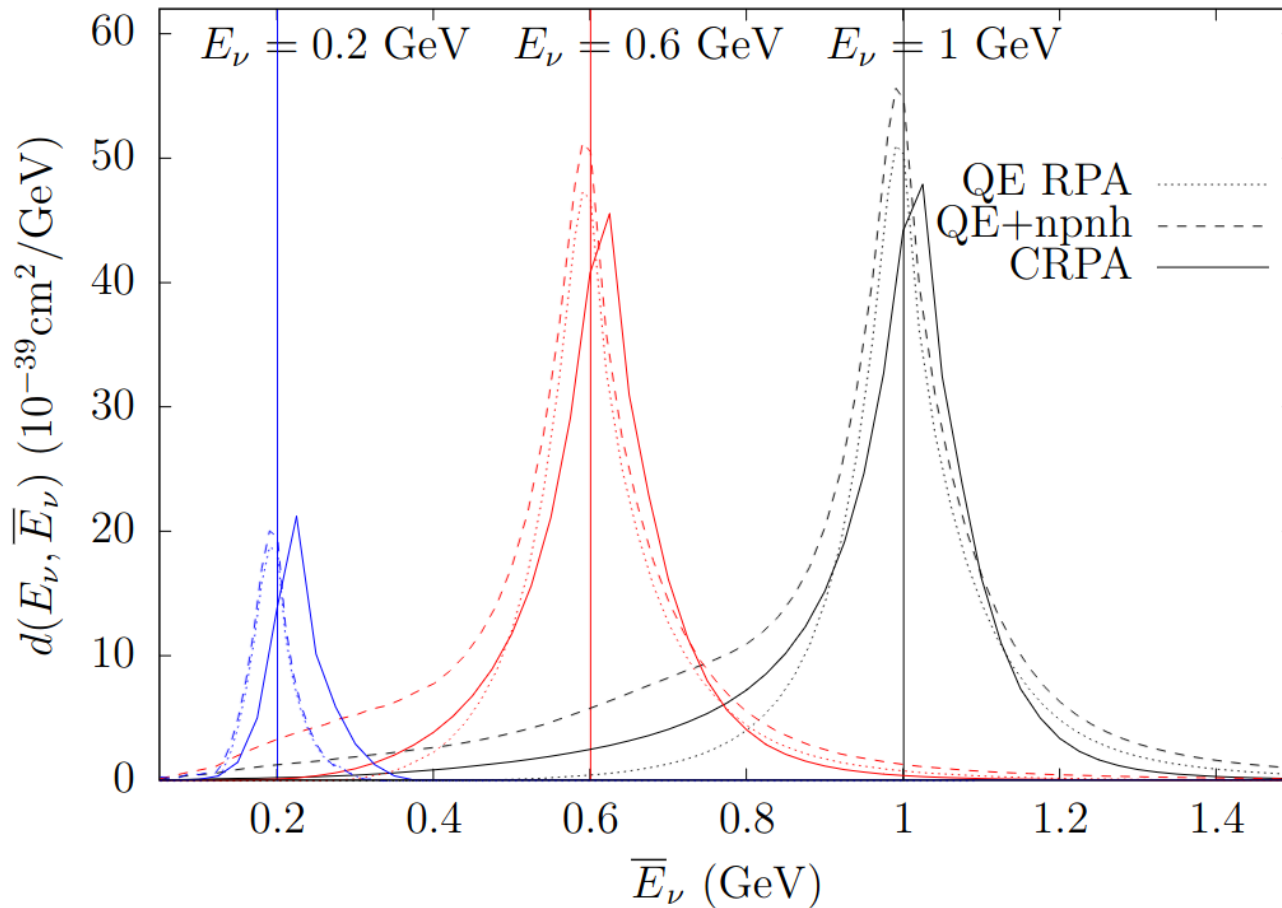
Uncertainty, constraints, and flexibility in $K(E, X)$
Determine possible precision of analysis

Neutrino-nucleus cross section in oscillation analysis

MiniBooNE use a 1-d observable

$$\frac{dN}{dX} = \int dE_\nu \Phi(E_\nu) K(E_\nu, X) \quad X = \bar{E}_\nu = \frac{2M'_n E_l - (M_n'^2 + m_l^2 - M_p^2)}{2(M'_n - E_l + P_l \cos \theta)}$$

[arxiv:1808.07520]



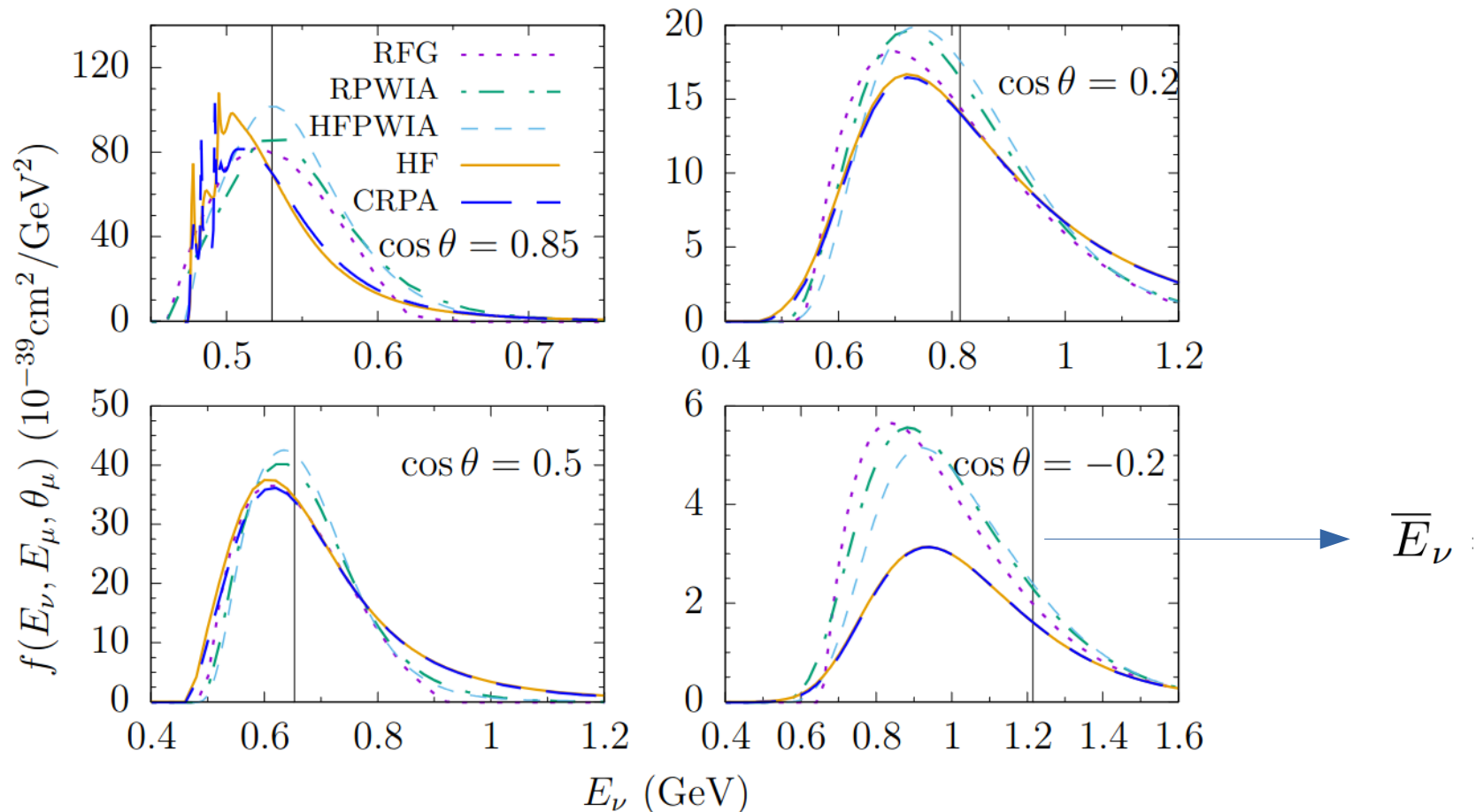
Width and model-dependence of distribution affect oscillation analyses

More observables → more information

[arxiv:1808.07520]

$$\frac{dN}{dX} = \int dE_\nu \Phi(E_\nu) K(E_\nu, X) \quad X = (T_\mu, \theta_\mu)$$

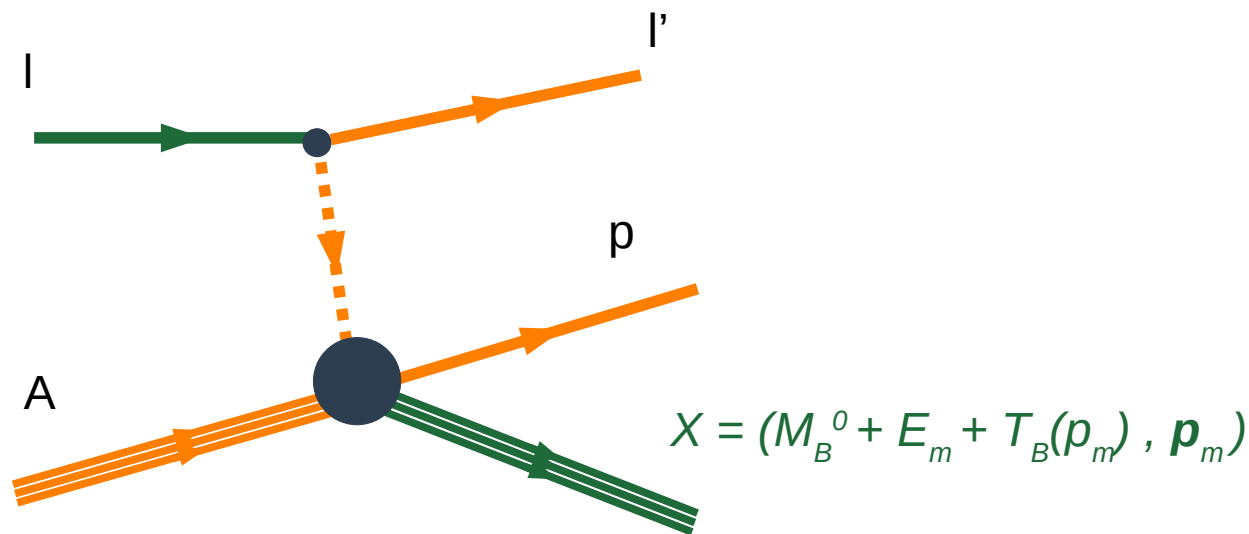
$T_\mu = 0.35$ GeV with MiniBooNE flux



Smearing determined by inclusive double differential CS

From inclusive to semi-inclusive with LArTPCs → 1μ1p events

$$\left\langle \frac{d^6\sigma}{dk_l d\Omega_l dp_N d\Omega_N} \right\rangle = \int dE \phi(E) \frac{d^6\sigma(E)}{dk_l d\Omega_l dp_N d\Omega_N} \quad X = (K_\mu, K_N)$$



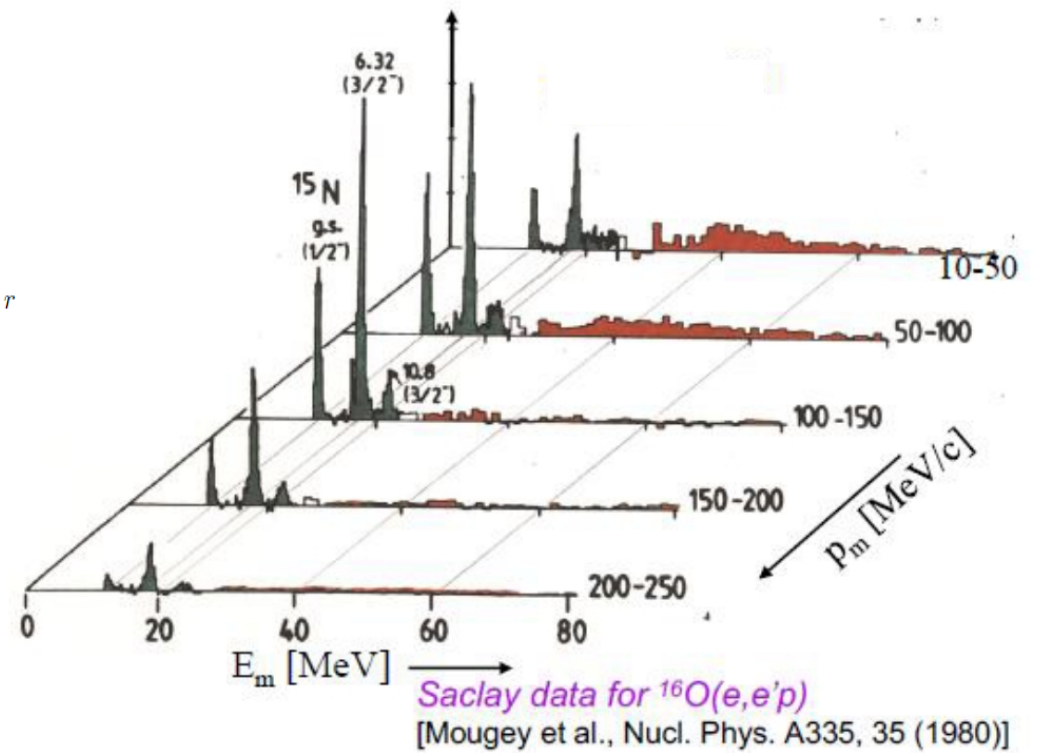
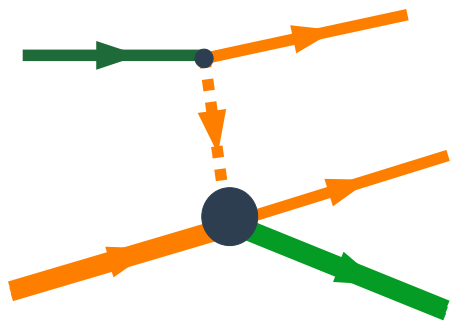
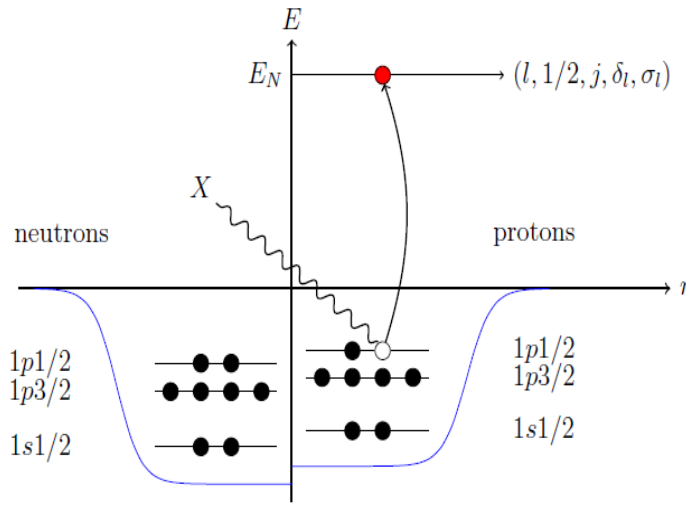
Missing energy

$$E_m = E_l - E_{l'} - T_N - T_B$$

Missing momentum

$$|\mathbf{p}_m| = |\mathbf{k}_l - \mathbf{k}_{l'} - \mathbf{k}_N|$$

Exclusive electron scattering: Missing energy distributions

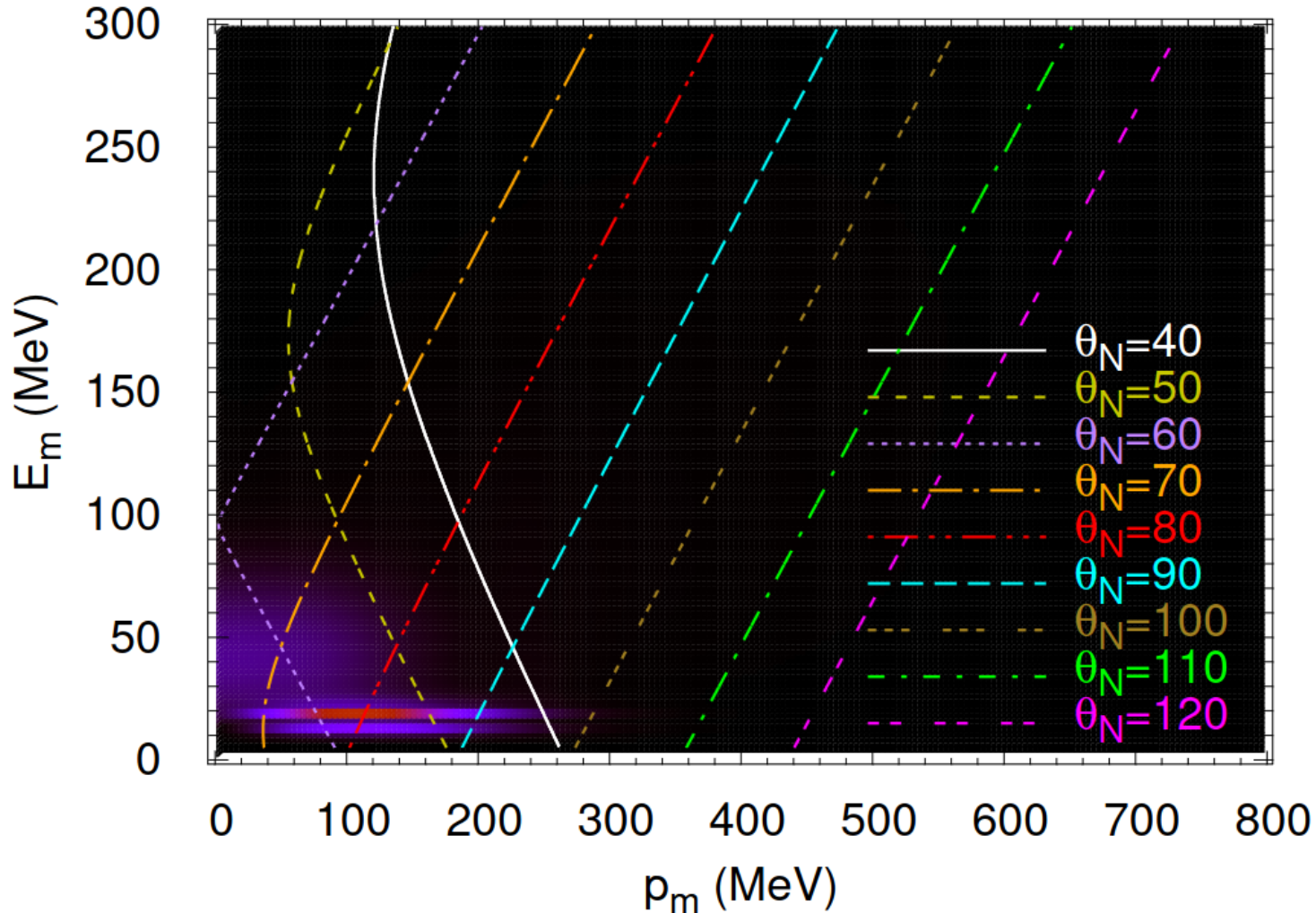


Direct nucleon knockout dominated by distinct shell-model peaks
 → Residual system with low excitation energy E_m

Neutrino scattering: energy reconstruction for $1\mu 1p$ events

[R. Gonzalez-Jimenez et al. arxiv:2104.01701]

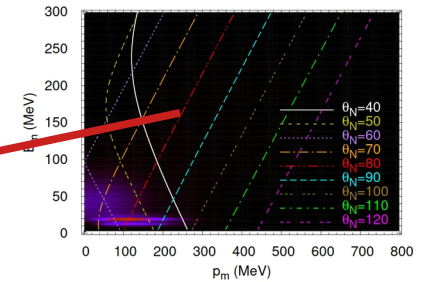
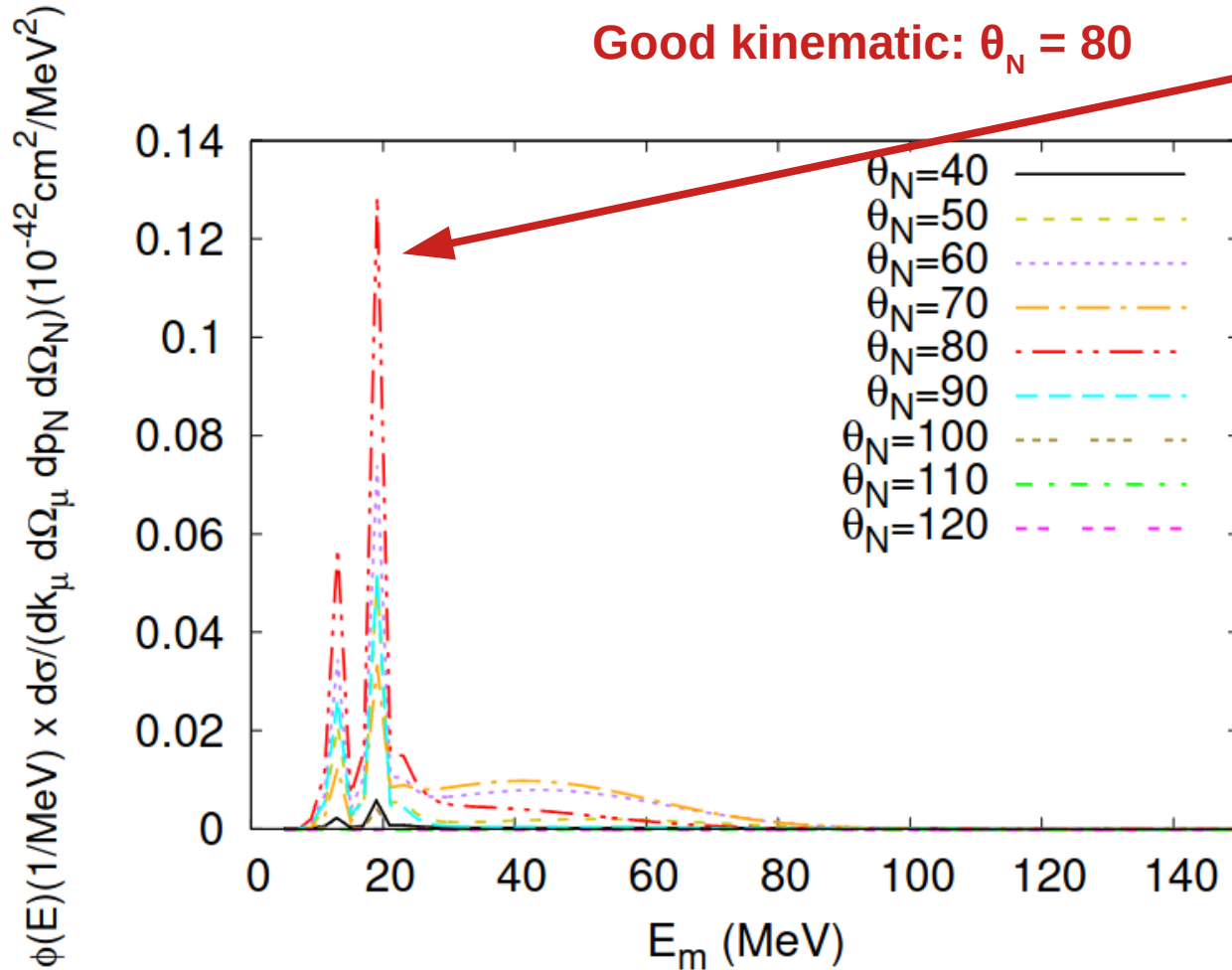
Fixed $k_\mu, k_N \rightarrow E_\nu$ defines a trajectory in E_m, p_m space



'Good kinematics' ($\theta_N = 80$) \rightarrow cross low- E_m peaks at small p_m

Neutrino scattering: energy reconstruction for $1\mu 1p$ events

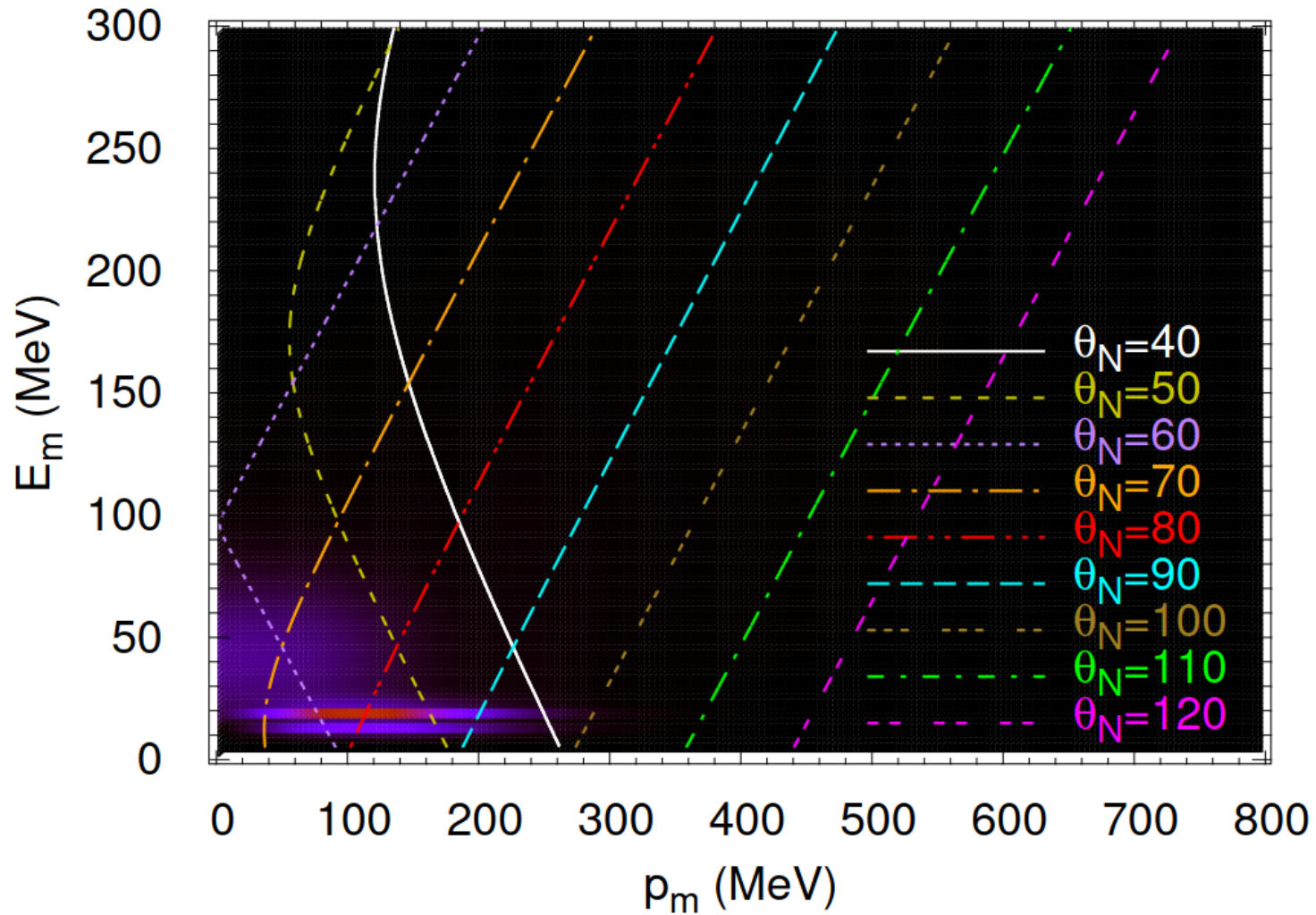
[R. Gonzalez-Jimenez et al. arxiv:2104.01701]



Good kinematics → Energy estimator with narrow uncertainty!
Sub-percent energy reconstruction → [arxiv:2104.017101]

Neutrino scattering: energy reconstruction for $1\mu 1p$ events

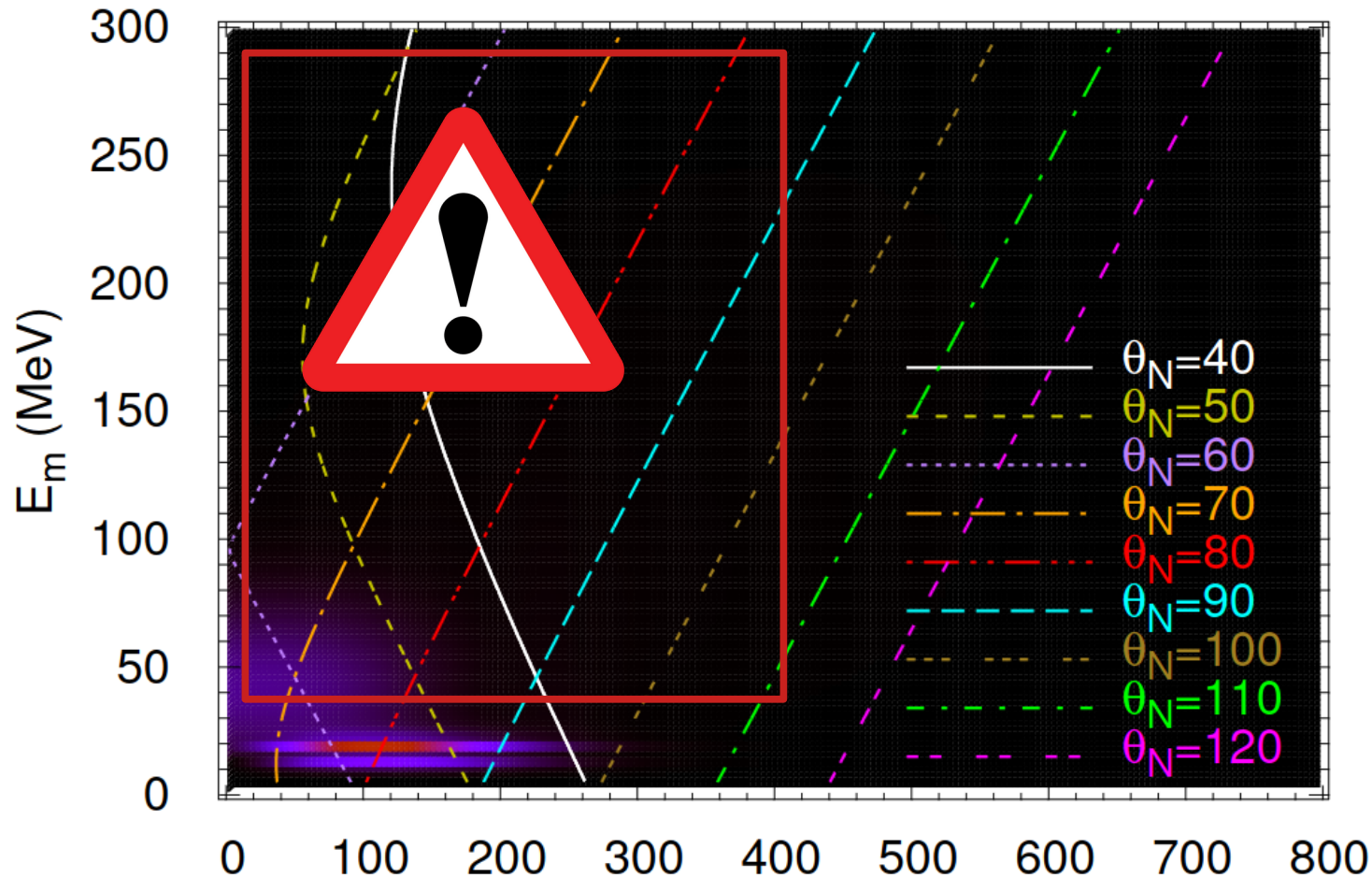
[R. Gonzalez-Jimenez et al. arxiv:2104.01701]



‘Good kinematics’ → percent-level energy reconstruction [arxiv:2104.017101]

Neutrino scattering: energy reconstruction for $1\mu 1p$ events

[R. Gonzalez-Jimenez et al. arxiv:2104.01701]



!Problem!

2 – nucleon knockout, pion-production, final-state interactions, ...
Populate the high- E_m region

Neutrino scattering: energy reconstruction for $1\nu 1n$ events

This talk: Final-state interactions in nucleon knockout

Based on

[A. Nikolakopoulos , A. Ershova, R. Gonzalez-Jimenez, J. Isaacson,
A.M. Kelly, K. Niewczas, N. Rocco, F. Sanchez, **arxiv:2406.09244**]

[A. Nikolakopoulos , R. Gonzalez-Jimenez, N. Jachowicz,
K. Niewczas, F. Sanchez, J.M. Udias, **PRC 105, 054603**]

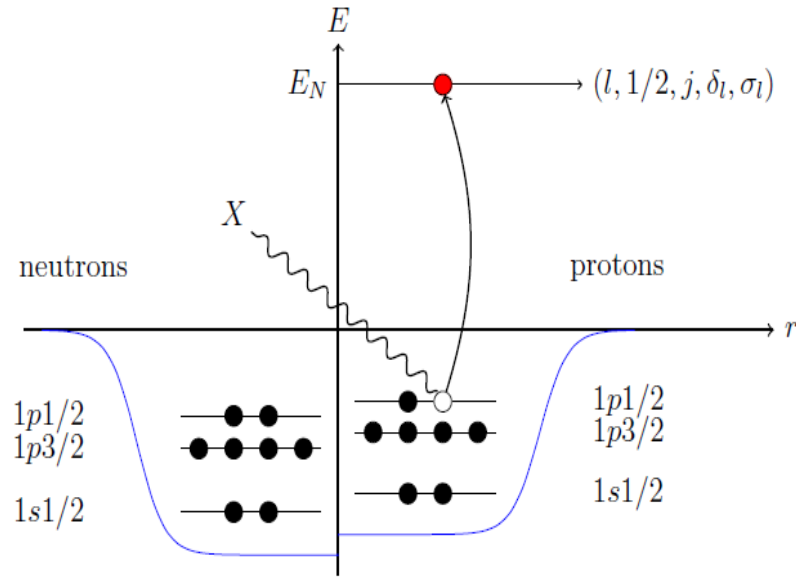
(Other efforts, e.g. electron-muon neutrino interactions in extra slides)

0 100 200 300 400 500 600 700 800

!Problem!

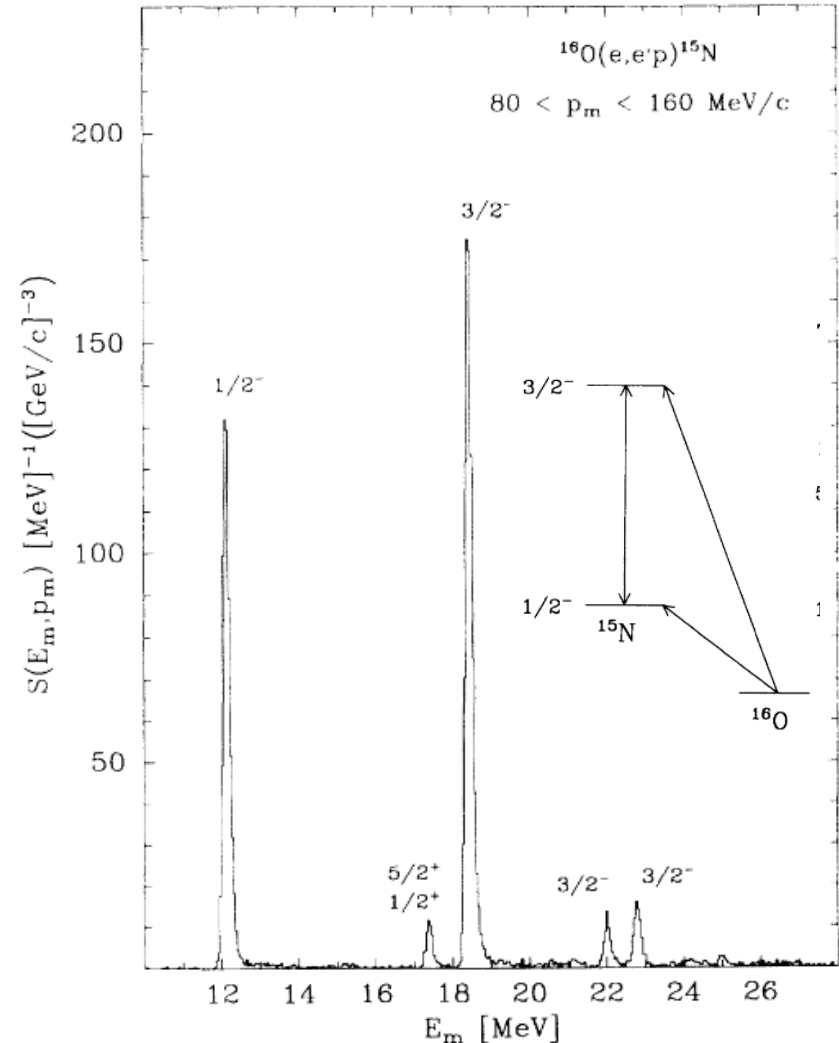
2 – nucleon knockout, pion-production, final-state interactions, ...
Populate the high- E_m region

Final-state interactions in exclusive (e,e'p) : optical potential



$$\mathcal{M} = j_{lep,\mu} \langle \Psi_f | \mathcal{O}^\mu | \Psi_i \rangle$$

$$\langle \Psi_f | = \langle \phi_p | \langle {}^{15}N^* |$$

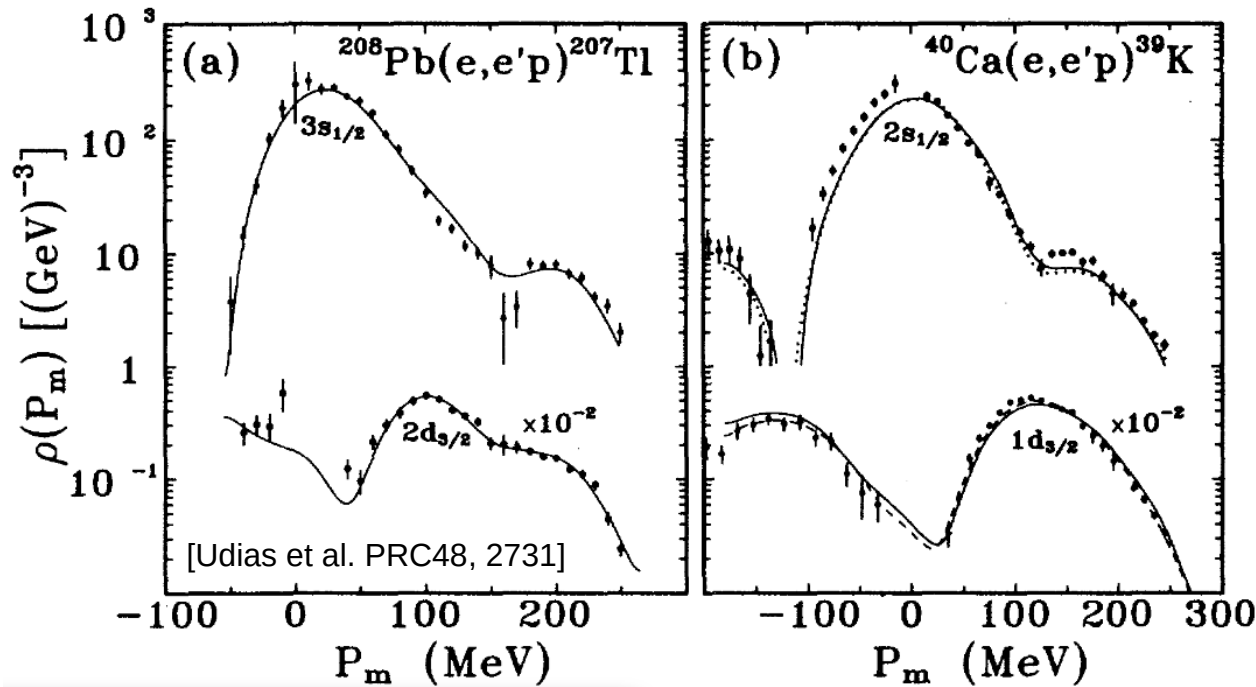


→ **Outgoing proton does not exchange energy with residual system**

Final-state interactions in exclusive (e,e'p) : optical potential

$$\mathcal{J}_{\kappa}^{m_j}(Q, P_N) = \int d\mathbf{p} \bar{\psi}(\mathbf{p} + \mathbf{q}, \mathbf{k}_N, s_N) \mathcal{O}^{\mu} \psi_{\kappa}^{m_j}(\mathbf{p})$$

Final-state : solution of Dirac equation in optical potential



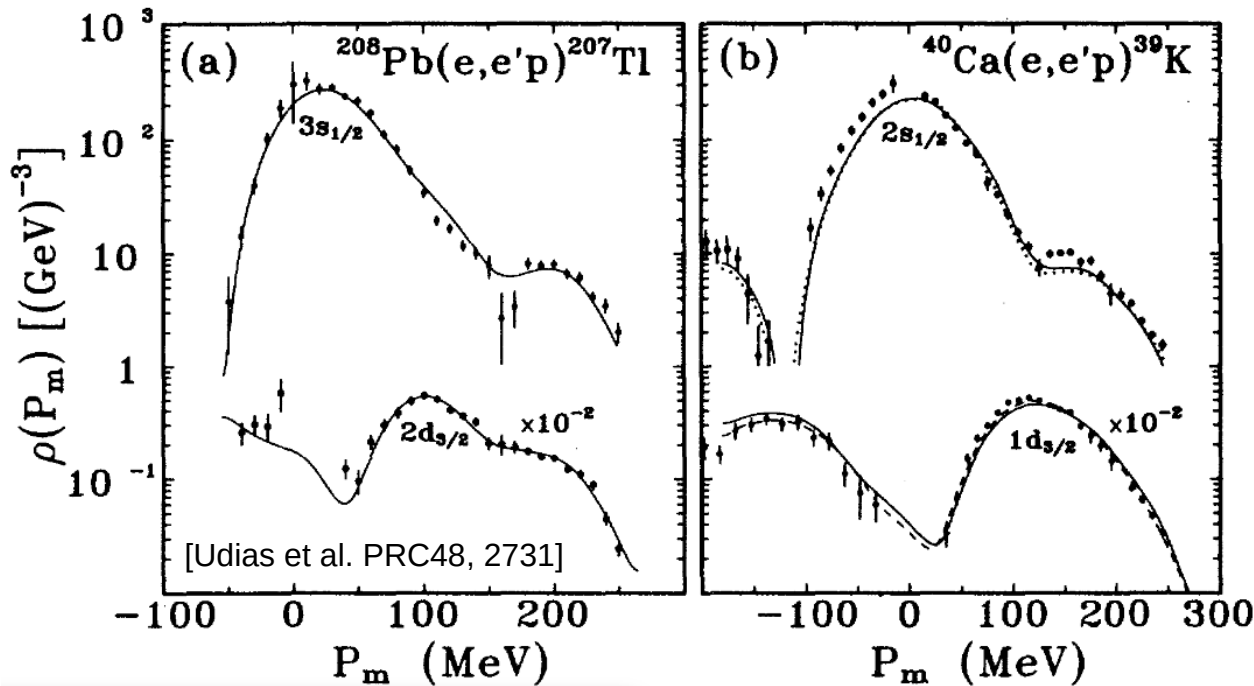
- 'Standard' approach for FSI in exclusive (e,e'p) analysis

e.g.
Jlab analyses of ^{40}Ar , ^{48}Ti
[PRD 107, 012005]
[PRD 105, 112002]

Final-state interactions in exclusive (e,e'p) : optical potential

$$\mathcal{J}_{\kappa}^{m_j}(Q, P_N) = \int d\mathbf{p} \bar{\psi}(\mathbf{p} + \mathbf{q}, \mathbf{k}_N, s_N) \mathcal{O}^{\mu} \psi_{\kappa}^{m_j}(\mathbf{p})$$

Final-state : solution of Dirac equation in optical potential



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e.g.
 Jlab analyses of ^{40}Ar , ^{48}Ti
 [PRD 107, 012005]
 [PRD 105, 112002]

The optical potential removes nucleons that undergo inelastic FSI



In neutrino experiments want to know where the nucleon goes

Where do the protons go?: Intranuclear Cascade model (INC)

- ED-RMF

FSI in inclusive interactions

-INC

FSI for (semi-)exclusive channels

- ROP

FSI in single exclusive channel

Production of final-state $|X\rangle = |p\rangle|^{39}\text{Ar}^*\rangle$

$$|\mathcal{M}|^2 \approx \left| \sum_{\alpha} \langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle \langle \psi_{\alpha} | X \rangle \right|^2, \quad \longrightarrow \quad \text{Restrict to 1-body operator}$$

$$\approx \sum_{\alpha} |\langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle|^2 |\langle \psi_{\alpha} | X \rangle|^2 \quad \longrightarrow \quad \text{Classical approximation}$$

$$\approx \sum_{\alpha} |\langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle|^2 P(X|\alpha). \quad \longrightarrow \quad \text{Intranuclear Cascade}$$

Where do the protons go?: Intranuclear Cascade model (INC)

- ED-RMF

FSI in inclusive

-INC

FSI for relevant (semi-)exclusive channels

- ROP

FSI in single exclusive channel





Production of final-state $|X\rangle = |p\rangle|^{39}\text{Ar}^*\rangle$

$$\begin{aligned}
 |\mathcal{M}|^2 &\approx \left| \sum_{\alpha} \langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle \langle \psi_{\alpha} | X \rangle \right|^2, && \longrightarrow \text{Restrict to 1-body operator} \\
 &\approx \sum_{\alpha} |\langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle|^2 |\langle \psi_{\alpha} | X \rangle|^2 && \longrightarrow \text{Classical approximation} \\
 &\approx \sum_{\alpha} |\langle \Psi_0 | T_{1b} | \psi_{\alpha} \rangle|^2 P(X|\alpha) && \longrightarrow \text{Intranuclear Cascade}
 \end{aligned}$$

ROP
ED-RMF
INC

Can benchmark the INC with Optical potential with same nuclear model
For direct proton knockout

**Benchmarking intranuclear cascade models for neutrino scattering
with relativistic optical potentials**

A. Nikolakopoulos ^{1,2,*} R. González-Jiménez ³ N. Jachowicz,¹ K. Niewczas,^{1,4} F. Sánchez ⁵ and J. M. Udías ³

Input to the INC

Fully differential events from RDWIA or RPWIA
For $1\mu 1p$





Cuts on the INC results

Single proton events where proton does not lose
Energy \rightarrow no inelastic FSI

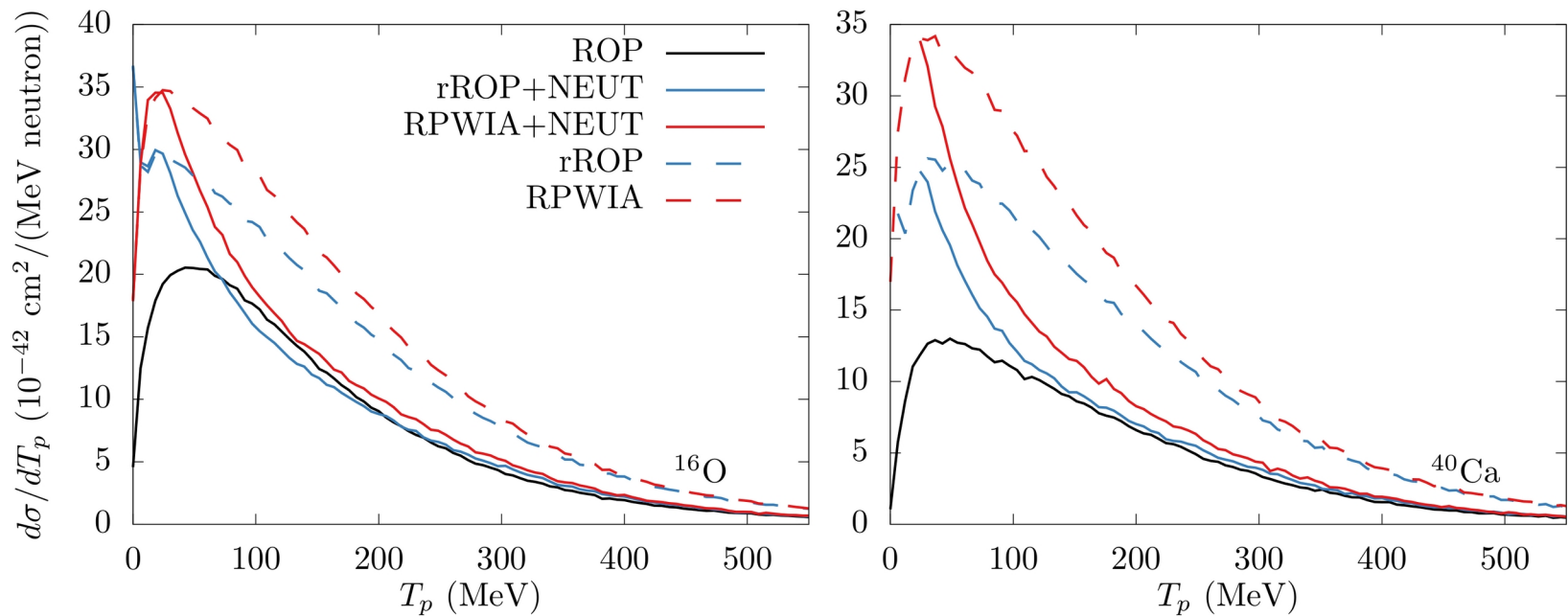


Can be compared to ROP results

Benchmarking intranuclear cascade models for neutrino scattering with relativistic optical potentials

A. Nikolakopoulos ^{1,2,*} R. González-Jiménez ³ N. Jachowicz,¹ K. Niewczas,^{1,4} F. Sánchez ⁵ and J. M. Udías ³

Flux-folded with T2K ND flux: NEUT INC

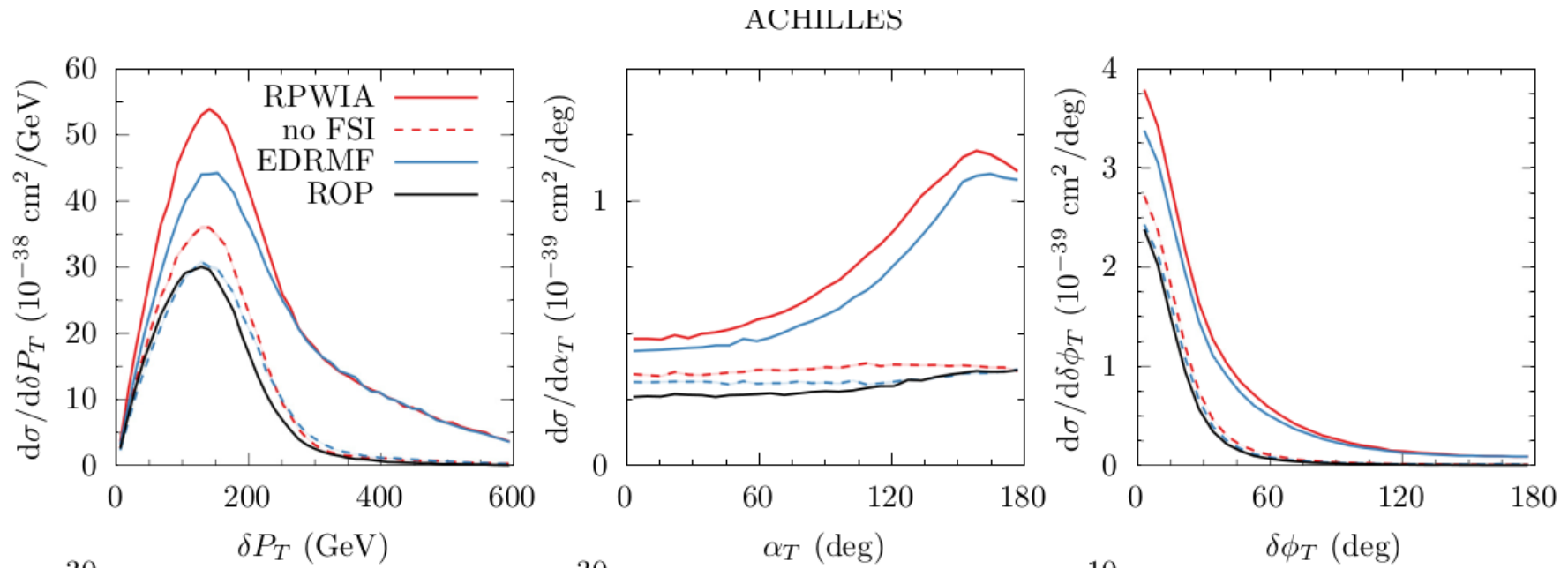


ROP and INC agree at large T_p but large disagreement for small T_p

Neutrino-induced proton knockout from argon in MicroBooNE

[Arxiv:2406.09244]

- Flux-folded results for MicroBooNE
- ACHILLES, INCL, NEUT, and NuWro INC models
- Large set of kinematic distributions for comparison
- Detailed comparisons in backup slides

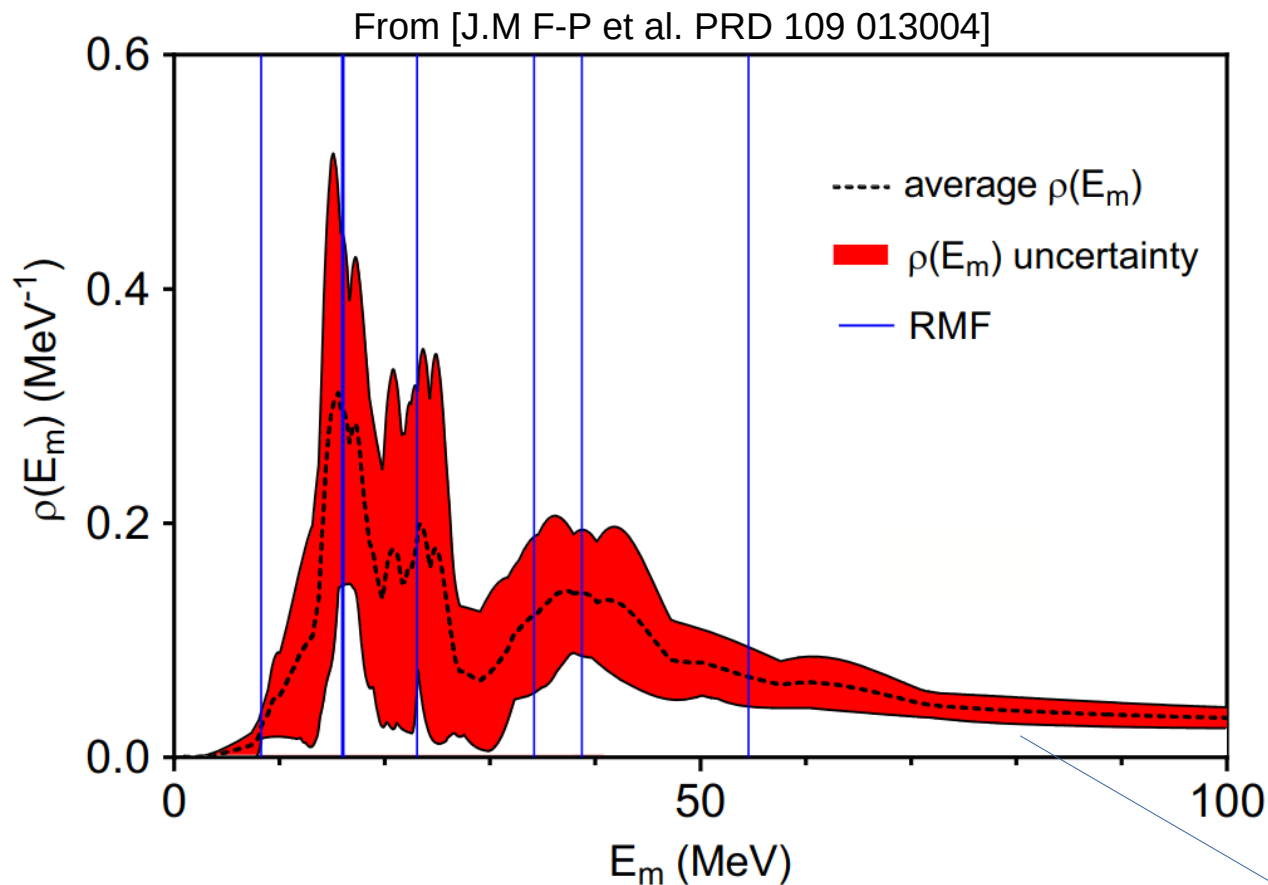


Some findings:

- Agreement depends on input calculation (ED-RMF \leftrightarrow RPWIA)
- Large differences between INCs (low T_p & treatment of correlations)
- No full agreement between any INC and ROP

RDWIA calculations with realistic spectral functions for MicroBooNE

See: [J. M. Franco-Patino et al. PRD 109, 013004] & [R. Gonzalez-Jimenez et al. PRC 105, 025502]



mean field

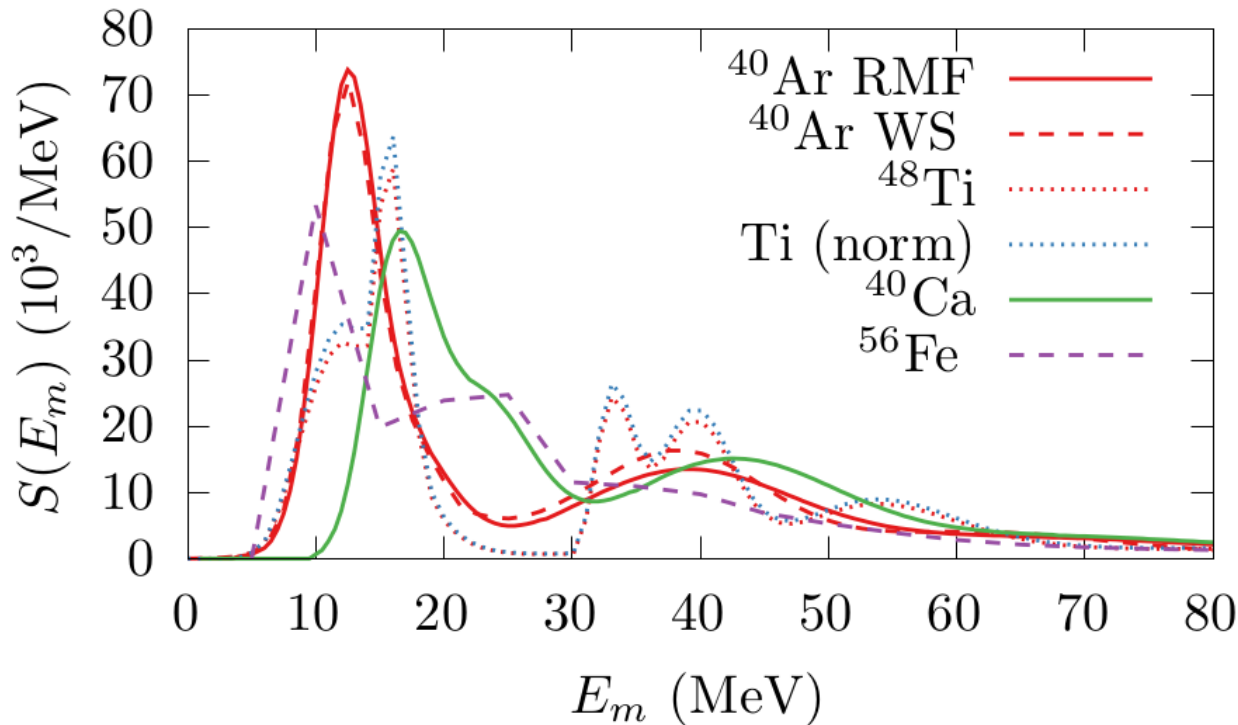
mean field + src

$$\delta(E_i - E_f) \sum_f |\mathcal{M}_{if}| = L_{\mu\nu} \sum_{\kappa} (2J_{\kappa} + 1) \delta(E_m - E_{\kappa}) H_{\kappa}^{\mu\nu}(Q, P_N) \rightarrow L_{\mu\nu} \left\{ \sum_{\kappa} N_{\kappa} \rho_{\kappa}(E_m) H_{\kappa}^{\mu\nu}(Q, P_N) + \rho_{corr}(E_m) H_{corr}^{\mu\nu}(Q, P) \right\}$$

RDWIA calculations with spectral functions for MicroBooNE

$$L_{\mu\nu} \left\{ \sum_{\kappa} N_{\kappa} \rho_{\kappa}(E_m) H_{\kappa}^{\mu\nu}(Q, P_N) + \rho_{corr}(E_m) H_{corr}^{\mu\nu}(Q, P) \right\}$$

Choices of N_{κ} and $\rho(E_m)$



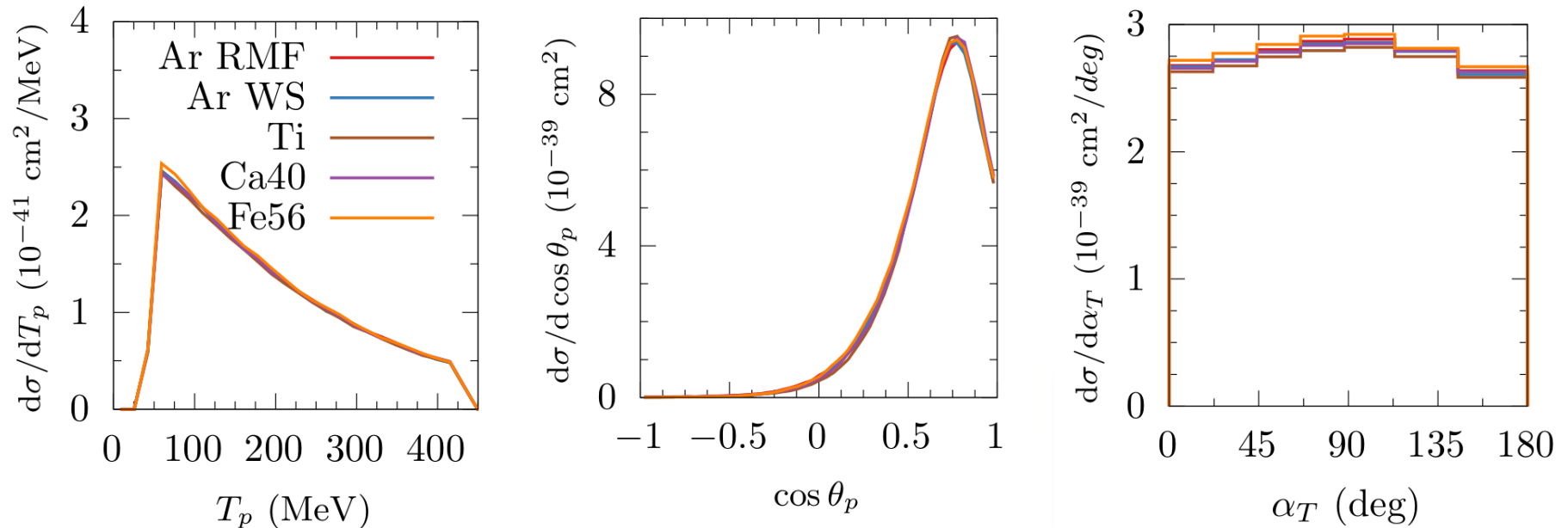
- ^{40}Ar spectral functions
[Butkevich PRC 85, 065501]
& [Jlab, PRD 107, 012005]
- ^{48}Ti from Jlab
[PRD 107, 012005]
- ^{56}Fe
[Benhar et al. NPA 579, 493]
- ^{40}Ca
[Butkevich PRC 85, 065501]

Large variation in E_m profiles to check sensitivity of observables

Sensitivity to variations in the spectral functions: PWIA calculations

[Arxiv:2406.09244]

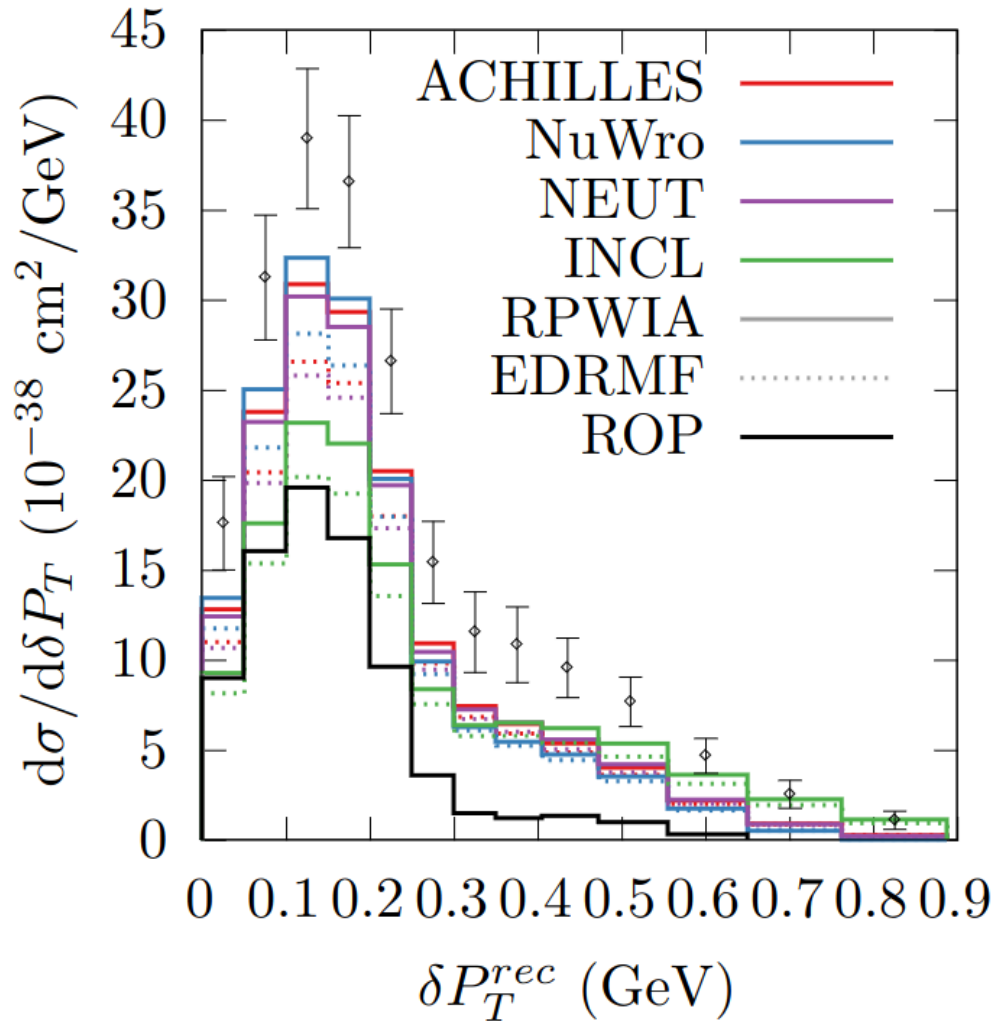
Observables for MicroBooNE flux-averaged signal



- Negligible differences between different spectral-functions for observables that do not correlate p_p and p_μ
- Mild sensitivity only to p_m distributions
- Current MicroBooNE data not sensitive to missing energy distributions

Comparisons to MicroBooNE data

[Arxiv:2406.09244]

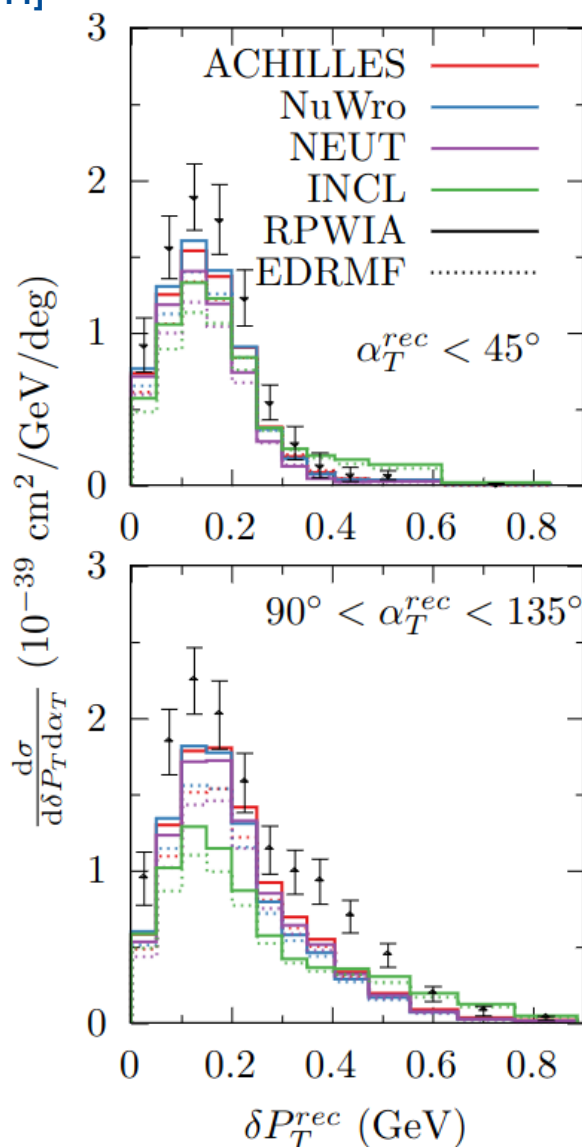


Isolate the effect of the INC
→ exact same inputs for each

Realistic final-state nucleon wave-functions
→ 10% effect
(Can generate events + Implementation in NEUT underway)

Comparisons to MicroBooNE data

[Arxiv:2406.09244]



Isolate the effect of the INC
→ exact same inputs for each

Realistic final-state nucleon wavefunctions
→ 10% effect

Cut on α_T
→ reduce rescattering for small α_T
→ Still sizeable contribution of resc
→ ROP alone cannot reproduce data

Low α_T dP_T underpredicted
→ Increase axial coupling ??
→ **Need to include the interference with 2-body currents!**

- **1-proton knockout could provide excellent energy-resolution**
→ **Backgrounds non-trivial !**
- **RDWIA + realistic spectral function + INC**
= most comprehensive description of single-nucleon knockout
→ **No full agreement with recent MicroBooNE data**

Where to go from here?

Theory

- * **Interference with 2-body currents needs to be included !**
See e.g. [T. Franco Munoz et al. PRC108, 064608] [Lovato et al. , 2312.12545]
- * **Two-nucleon and single pion production (SPP) contributions**
 - **ACHILLES: will soon include 2-body interference + SPP with full FSI**
 - **NEUT: will include RDWIA calculations with SF**
 - **NuWro: new SPP [2405.0512] and 2-N [K. Niewczas Phd] implementations**

Experiment

- * **High statistics in SBND can test 'good event' selection in $1\mu 1p$ in LarTPC**
- * **New electron scattering data in non-trivial kinematic regions**
→ **A1 at MAMI, e4v at Jlab take electron data for ν program (+ more facilities ?!)**

Other stuff



Terminology : RDWIA, RPWIA and PWIA & ED-RMF and ROP

-Relativistic Distorted Wave Impulse Approximation (RDWIA)

$$\mathcal{J}_{\kappa}^{m_j}(Q, P_N) = \int d\mathbf{p} \bar{\psi}(\mathbf{p} + \mathbf{q}, \mathbf{k}_N, s_N) \mathcal{O}^{\mu} \psi_{\kappa}^{m_j}(\mathbf{p})$$

- Relativistic Plane Wave Impulse Approximation (RPWIA)

$$\mathcal{J} = (2\pi)^{3/2} \bar{u}(\mathbf{k}_N, s_N) \mathcal{O}^{\mu} \psi_{\kappa}^{m_j}(\mathbf{k}_N - \mathbf{q})$$

- Plane-Wave Impulse Approximation (PWIA)

The initial state is assumed proportional to a positive-energy spinor:

$$\psi_{\kappa}^{m_j}(\mathbf{p}) \propto f(|\mathbf{p}|)u(\mathbf{p})$$

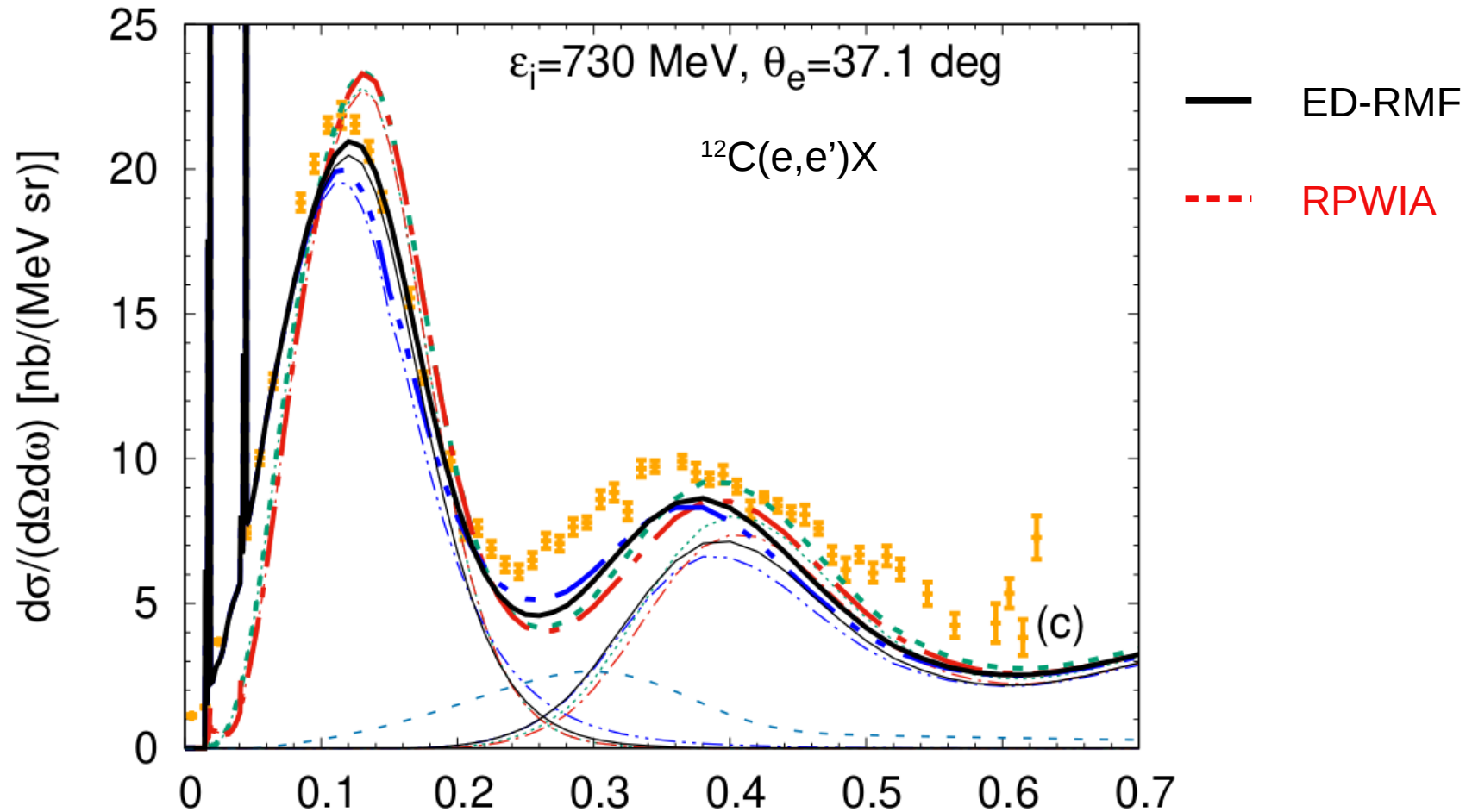
One obtains a factorized expression ('spectral function approach')

$$\frac{d\sigma(E_{\nu})}{dp_{\mu}d\Omega_{\mu}d\Omega_p dp_N} = \frac{G_F^2 \cos^2 \theta_c p_{\mu}^2 p_N^2}{(2\pi)^2} \frac{M_N^2}{E_{\nu} E_{\mu} E_N \bar{E}} L_{\mu\nu} h_{s.n.}^{\mu\nu} S(E_m, p_m)$$

Terminology : RDWIA, RPWIA and PWIA & ED-RMF and ROP

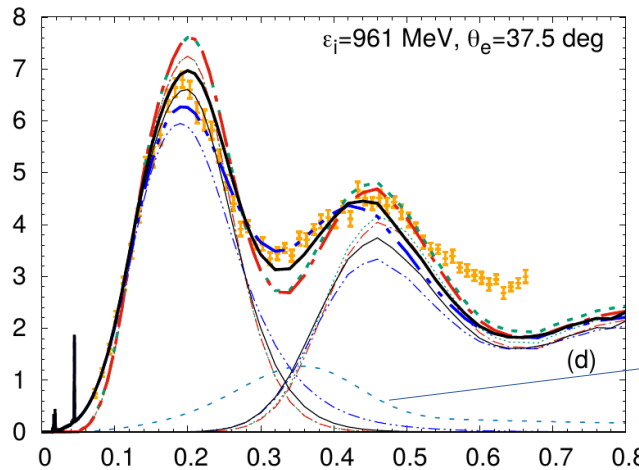
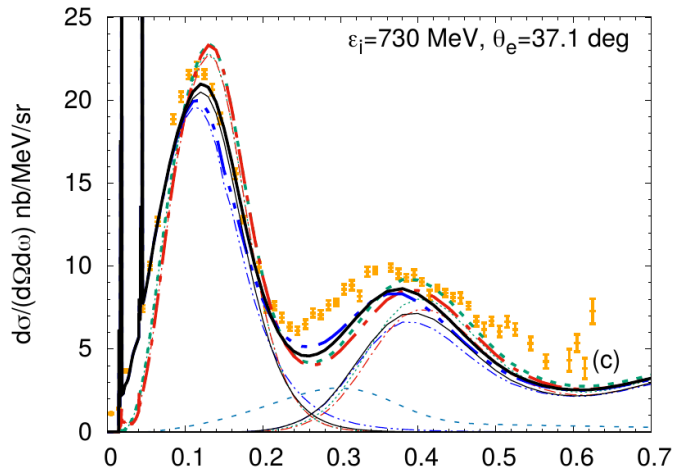
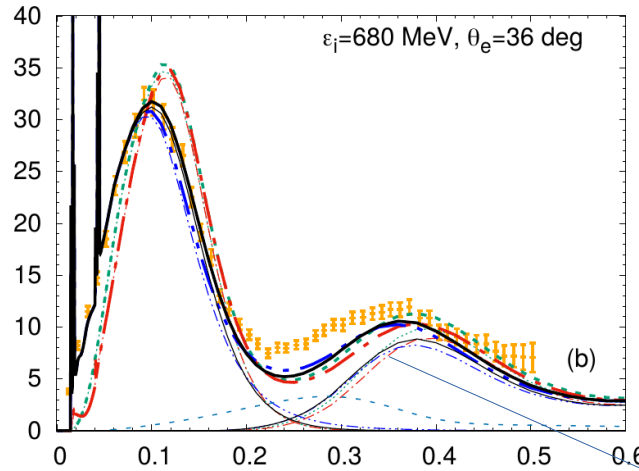
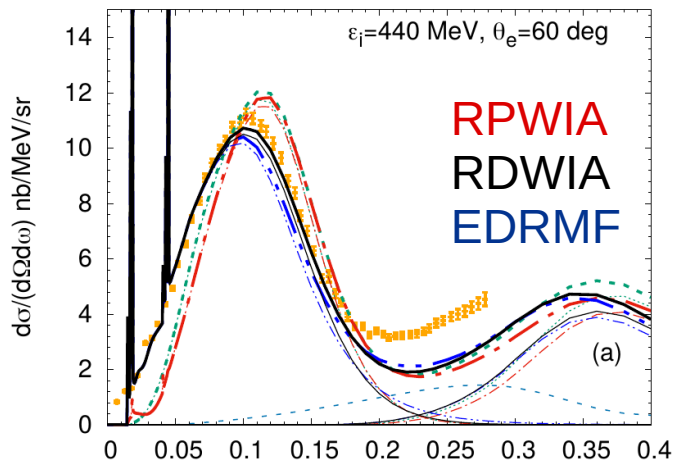
- Energy-Dependent Relativistic Mean-Field (ED-RMF)

$\bar{\psi}(\mathbf{p} + \mathbf{q}, \mathbf{k}_N, s_N)$ \longrightarrow Final-state in **real** Energy-Dependent potential
 \rightarrow suitable for **FSI in inclusive** cross section



Inclusive electron scattering off a RMF nucleus

[R. Gonzalez-Jimenez, A. Nikolakopoulos, N. Jachowicz, J.M. Udias PRC 100, 045501 (2019)]

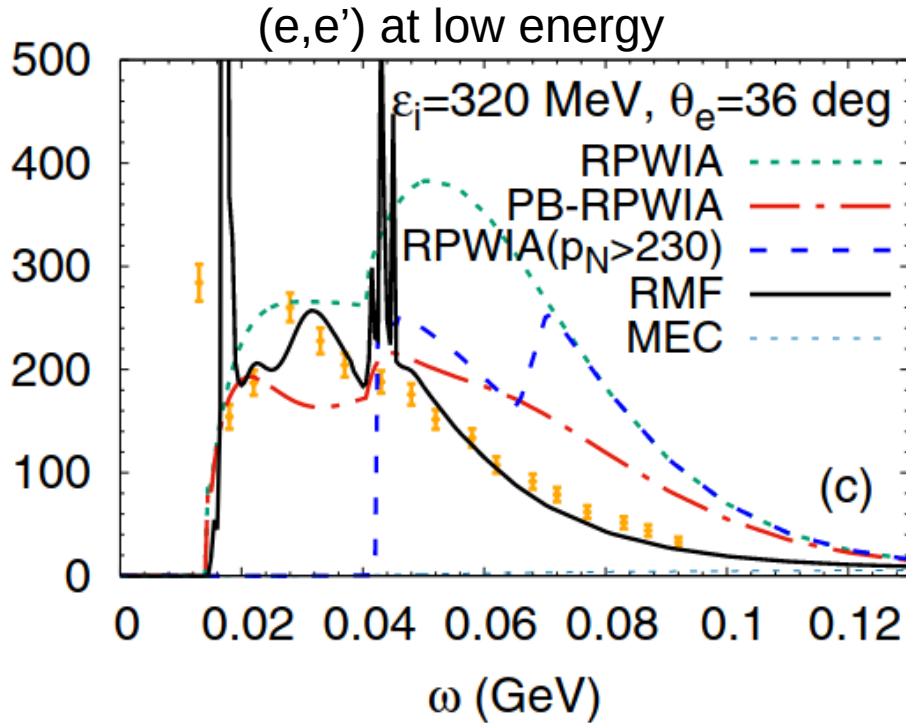


Single π
Production

2-body currents
[Megias et al.
PRD 91, 073004]

Nucleon knockout from a RMF nucleus: consistent initial & final-states

[R. Gonzalez-Jimenez, A. Nikolakopoulos, N. Jachowicz, J.M. Udias PRC 100, 045501 (2019)]



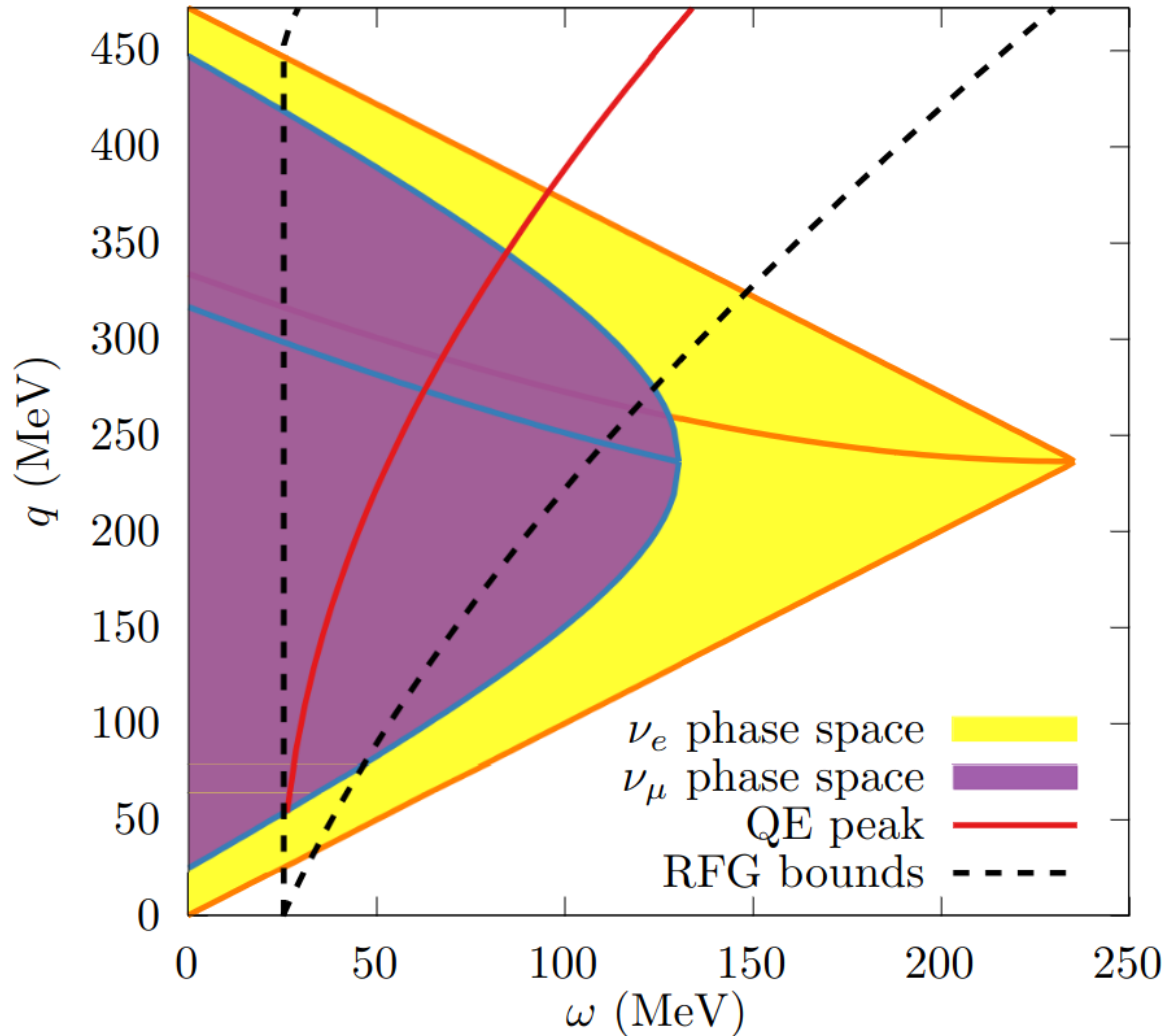
- Consistent states are crucial at low energies
 - provided by RMF
 - ED-RMF = RMF by construction at low-energy
- Consistent states are orthogonal ⇒ Pauli-Blocking

‘Pauli-Blocked’ RPWIA (PB-RPWIA): orthogonalize with respect to bound-states

$$|\Psi^{S_N}(\mathbf{p}_N)\rangle = |\psi_{pw}^{S_N}(\mathbf{p}_N)\rangle - \sum_{\kappa, m_j} [C_{\kappa}^{m_j, S_N}(\mathbf{p}_N)]^{\dagger} |\psi_{\kappa}^{m_j}\rangle \quad C_{\kappa}^{m_j, S_N}(\mathbf{p}_N) \equiv \langle \psi_{pw}^{S_N}(\mathbf{p}_N) | \psi_{\kappa}^{m_j} \rangle.$$

Electron- and muon neutrino interactions

$$E_\nu = 236 \text{ MeV}$$



Neutrino interactions are constrained in near-detector:
Mostly ν_μ measurements

In ν_e appearance or CP violation measurements:

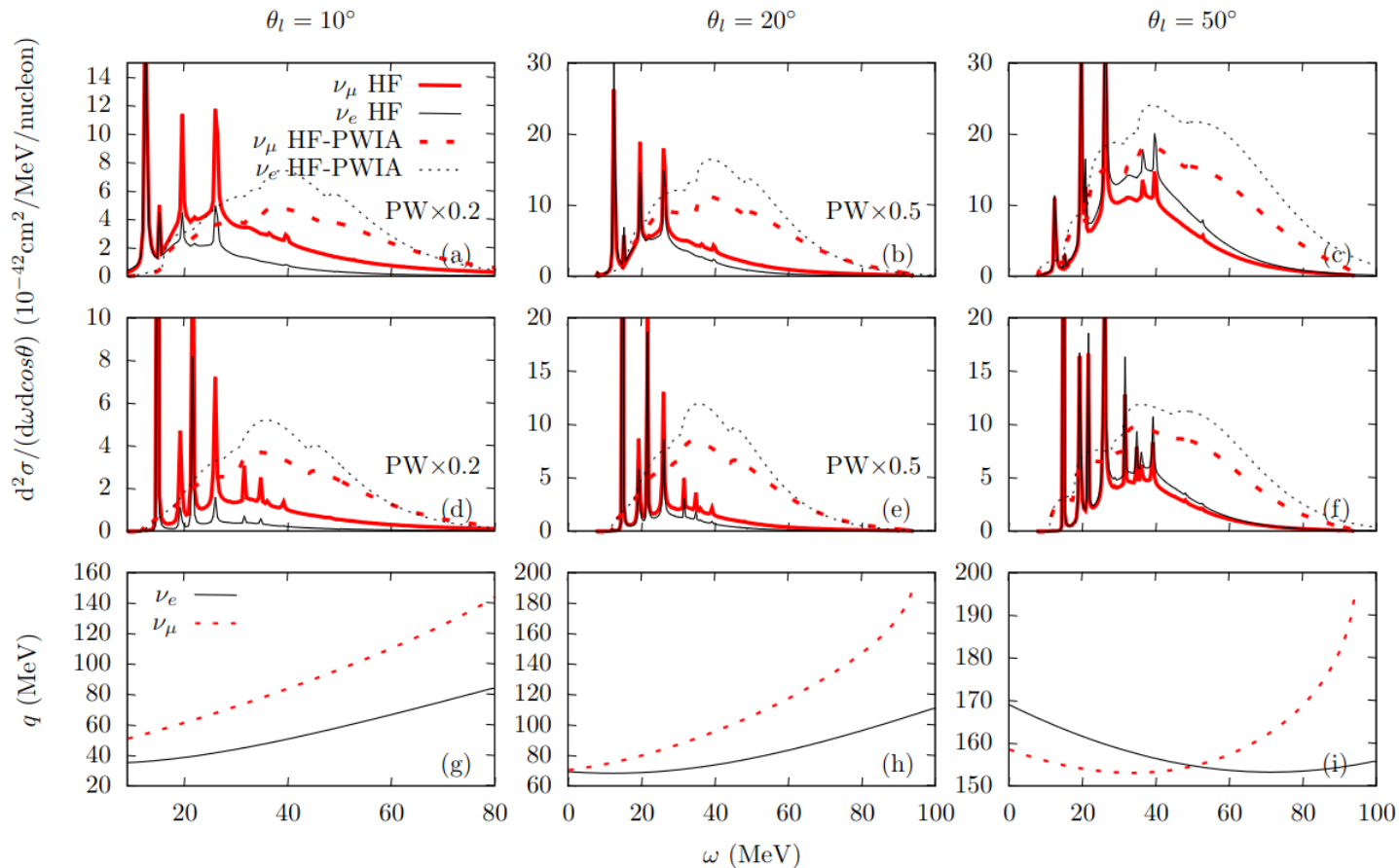
Physics that affects (anti)- ν interactions of different flavours differently can lead to systematic uncertainty

At high-E:
Cross sections are \sim the same

Low-E:
Naive expectation: electron CS larger because of phase space

Electron- and muon neutrino interactions at low energy

[A. Nikolakopoulos et al. PRL 123, 123, 052501 (2019)]



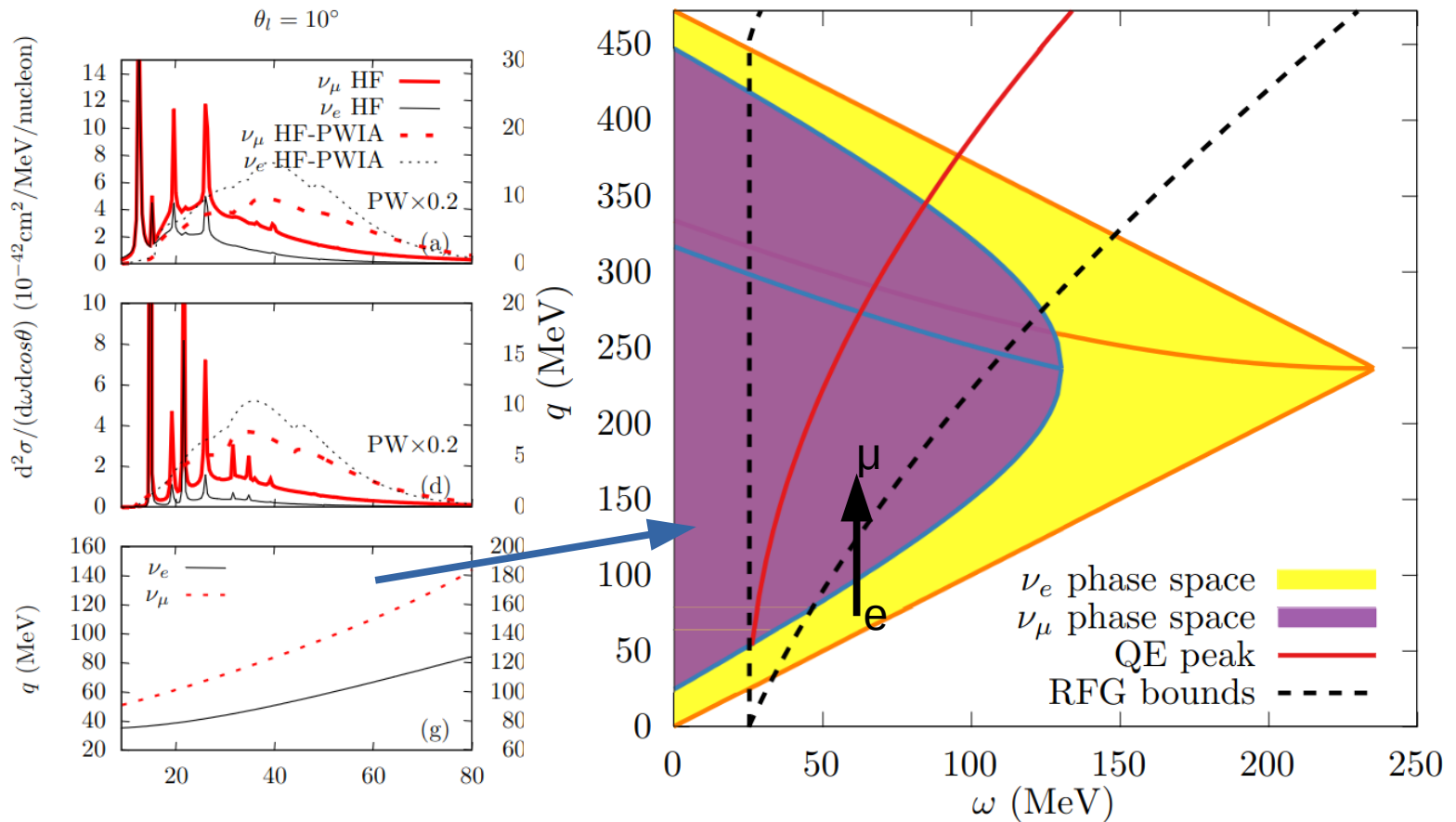
At low energy-momentum transfer ν_μ cross sections are larger than ν_e

$$q = \sqrt{E_\nu^2 + P_l^2 - 2 \cos \theta_l E_\nu P_l} \approx E_\nu - \sqrt{(E_\nu - \omega)^2 - m_l^2}$$

Electron- and muon neutrino interactions at low energy

[A. Nikolakopoulos e

$E_\nu = 236 \text{ MeV}$



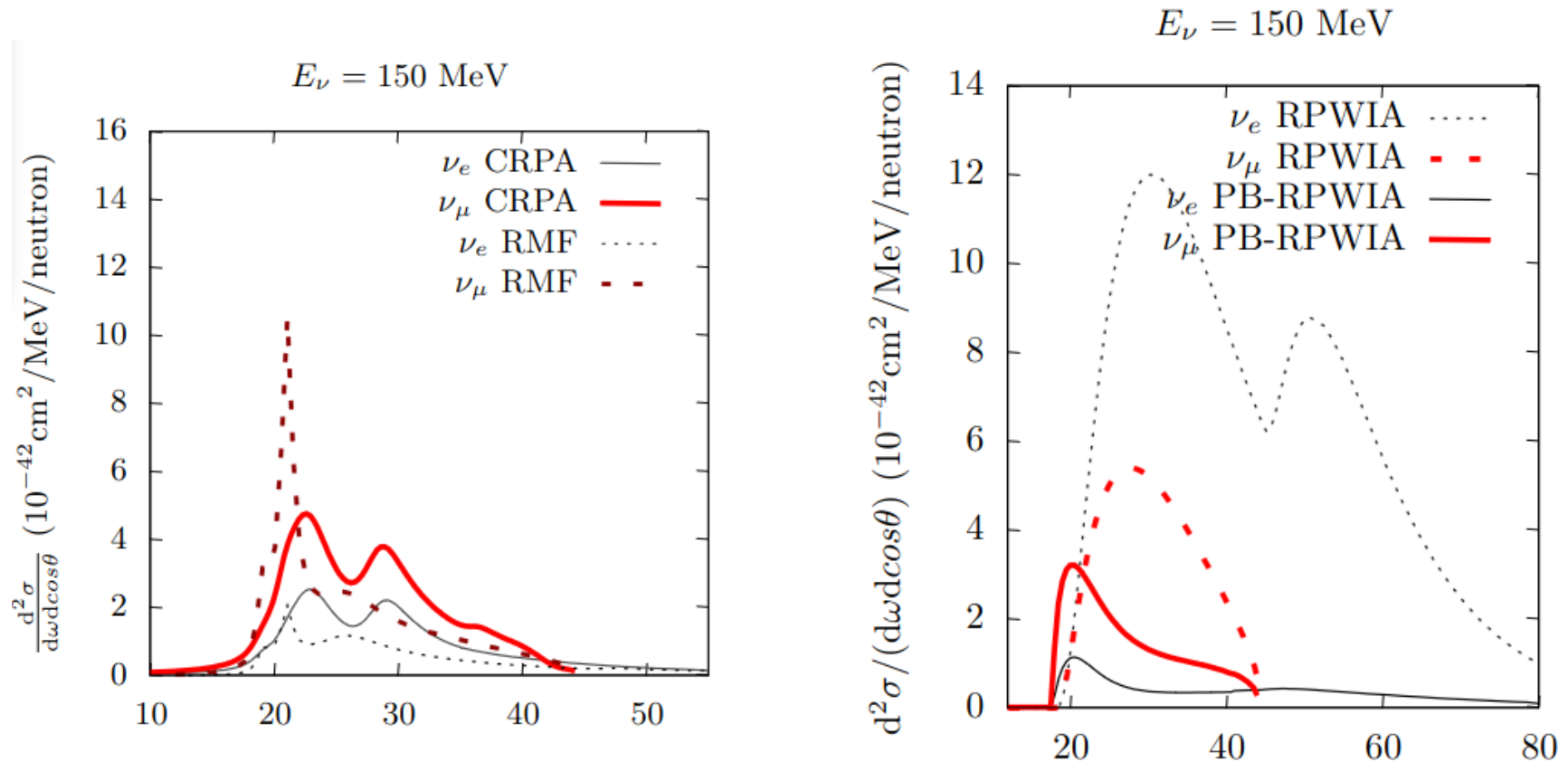
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Electron- and muon neutrino interactions at low energy

[A. Nikolakopoulos et al. PRL 123, 123, 052501 (2019)]

Can be understood from the orthogonality with bound states



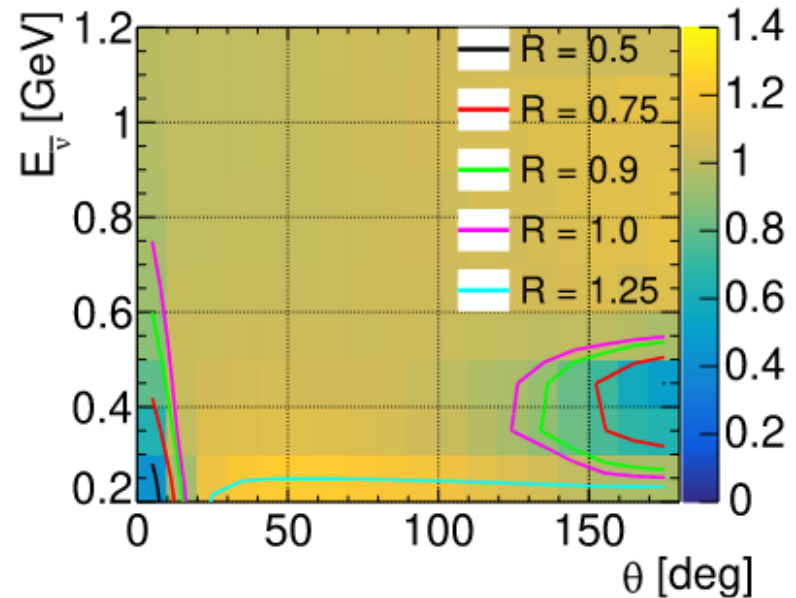
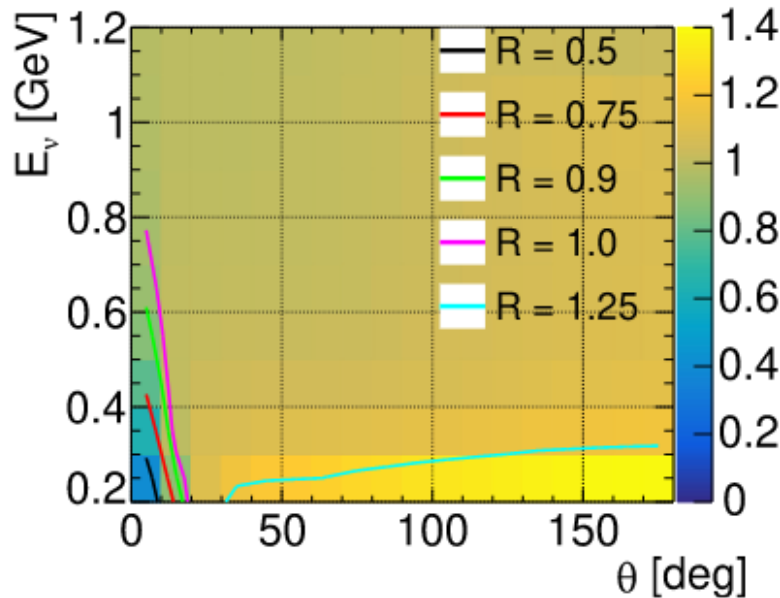
Uncertainties on the $\langle \bar{\nu} \rangle_e / \langle \bar{\nu} \rangle_\mu$ and $\nu_e / \bar{\nu}_e$ cross-section ratio from the modelling of nuclear effects at 0.2 to 1.2 GeV neutrino energies and their impact on neutrino oscillation experiments

T. Dieminger,^{1,*} S. Dolan,^{2,†} D. Sgalaberna,^{1,‡} A. Nikolakopoulos,³
 T. Dealtry,⁴ S. Bolognesi,⁵ L. Pickering,⁶ and A. Rubbia¹

arXiv:2301.08065v3 [hep-ph] 5 Apr 2023

$$\left[\frac{d\sigma_{\nu_e}(E_\nu)}{d\cos\theta} / \frac{d\sigma_{\nu_\mu}(E_\nu)}{d\cos\theta} \right]$$

$$\left[\frac{d\sigma_{\bar{\nu}_e}(E_\nu)}{d\cos\theta} / \frac{d\sigma_{\bar{\nu}_\mu}(E_\nu)}{d\cos\theta} \right]$$



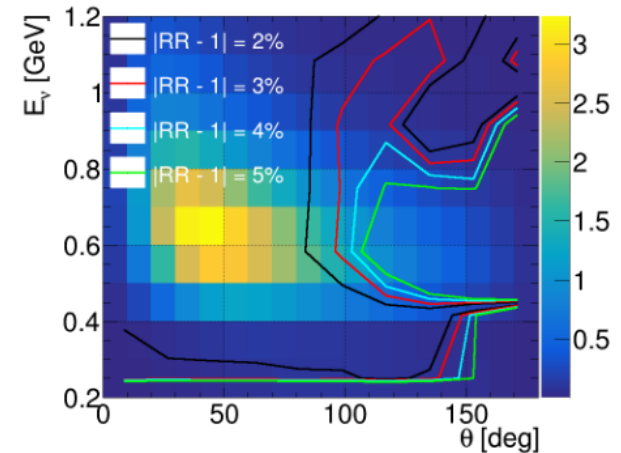
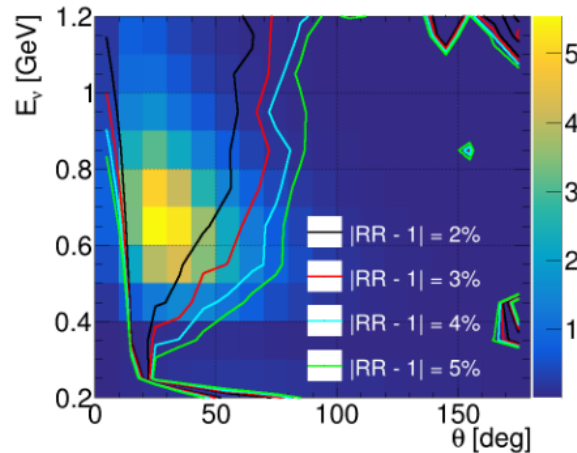
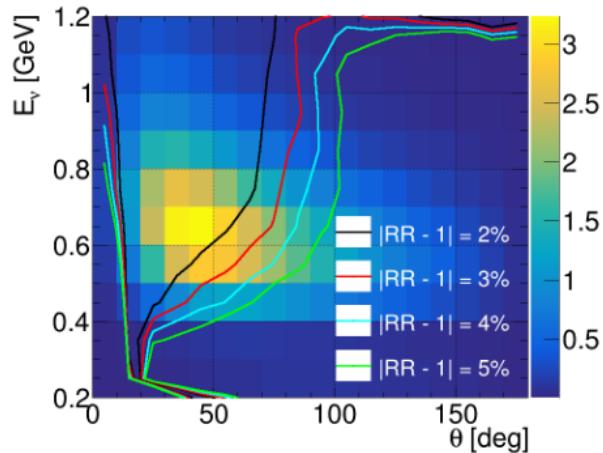
Assessing model-dependence: double ratios

$$RR_{\nu_\alpha/\nu_\beta}^{\text{Model 1/Model 2}}(E_\nu, \theta) = \frac{R_{\nu_\alpha/\nu_\beta}^{\text{Model 1}}(E_\nu, \theta)}{R_{\nu_\alpha/\nu_\beta}^{\text{Model 2}}(E_\nu, \theta)}$$

$$RR_{\nu_e/\nu_\mu}^{\text{CRPA/SF}}$$

$$RR_{\bar{\nu}_e/\bar{\nu}_\mu}^{\text{CRPA/SF}}$$

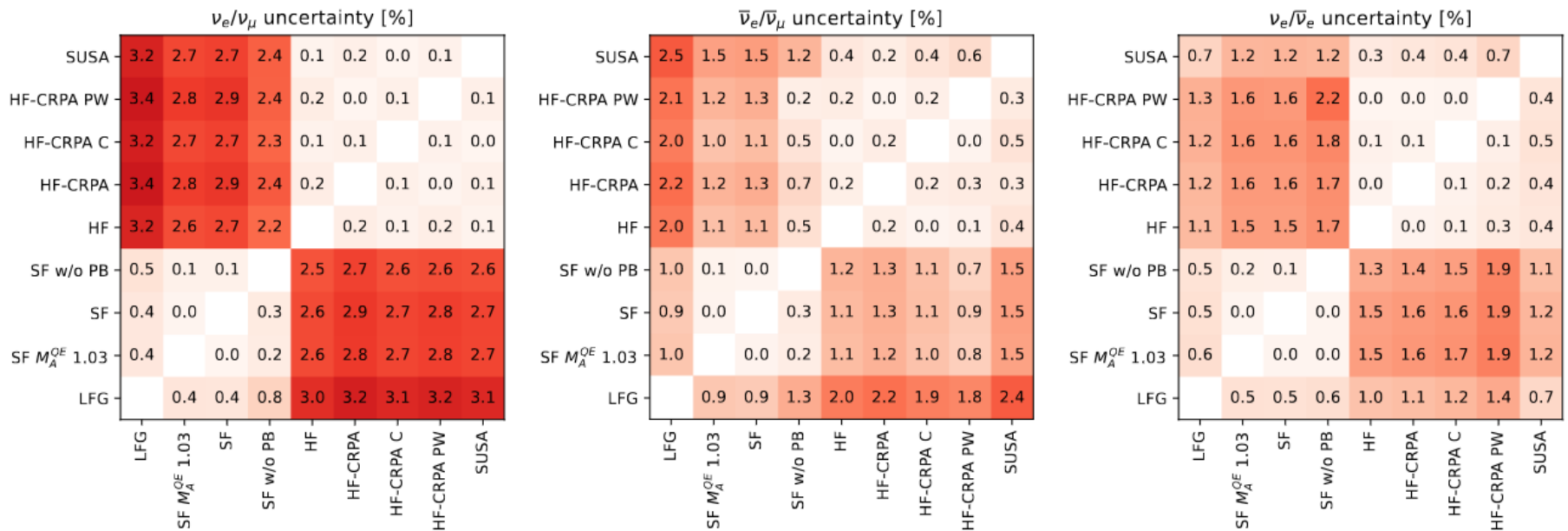
$$RR_{\nu_e/\bar{\nu}_e}^{\text{CRPA/SF}}$$



High-angle region in T2K is most relevant

Assessing model-dependence

Uncertainty from averaging model-differences *in Ratio* over predicted event rate

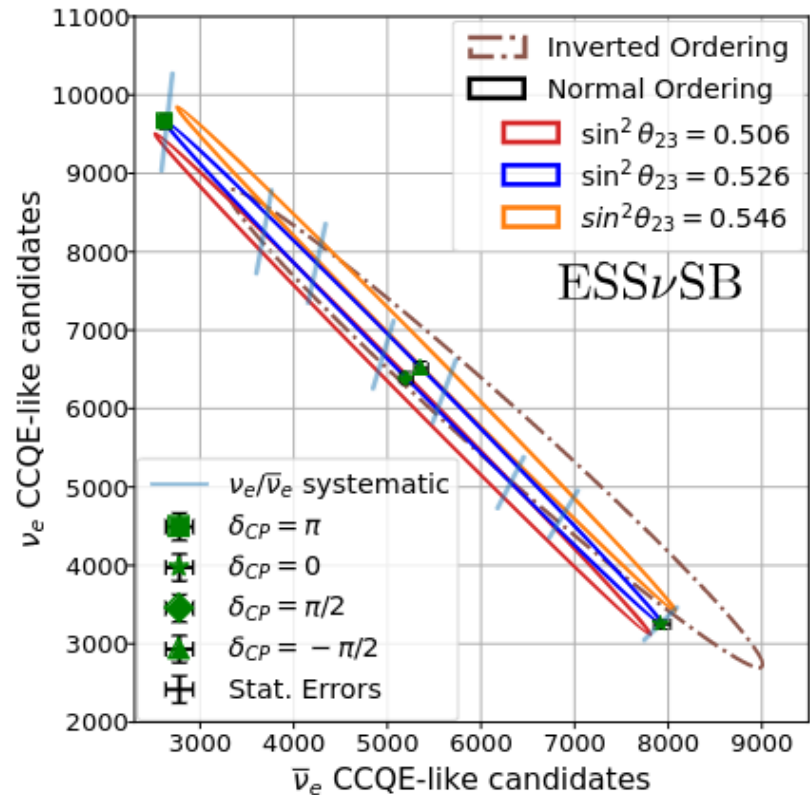
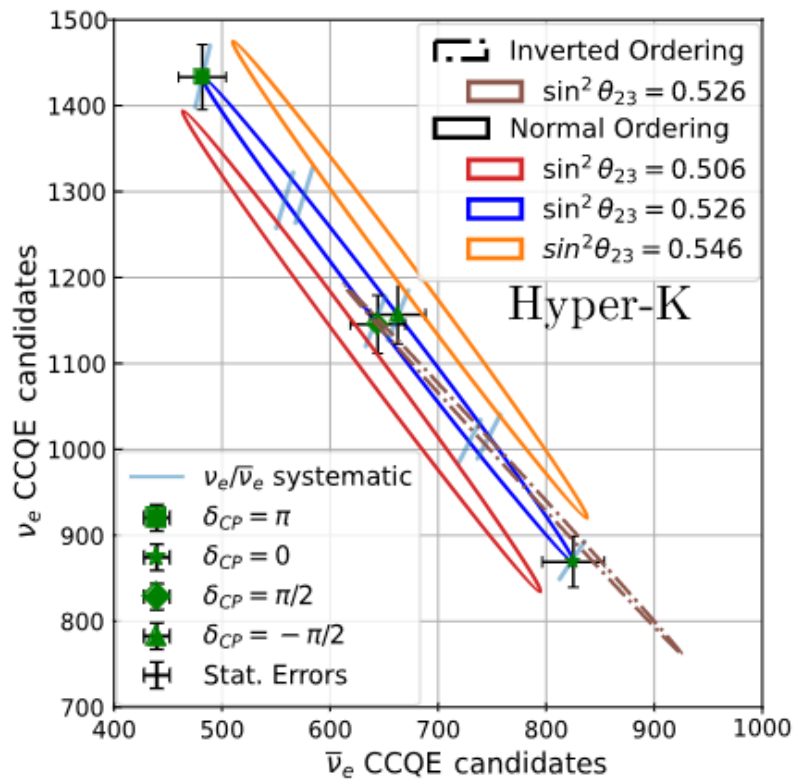


Quite robust when varying parameters within one model

Two 'sets': LFG and PWIA+SF and mean-field based models (or NEUT generator vs. other calculations ?)

Assessing impact on appearance experiment: bi-event plots

$$\nu_\mu \rightarrow \nu_e \quad \text{and} \quad \bar{\nu}_\mu \rightarrow \bar{\nu}_e$$



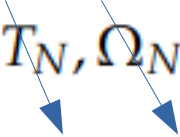
1) Not significant for $\sin \delta_{CP}$

2) Increases degeneracy in $\sin^2 \theta_{23}$

Magnitude 2-4% percent: uncertainty in line with values used in experiment

From inclusive to semi-inclusive one-nucleon knockout in neutrino event generators

Alexis Nikolakopoulos^{1,*}, Steven Gardiner¹, Afroditi Papadopoulou², Stephen Dolan³ and Raúl González-Jiménez⁴

$$P(E_l, \theta_l, T_N, \Omega_N)$$


Replace by
'factorized approach'

We get the GENIE version
based on **the same** inclusive
cross section!

We generate events for (e,e'p) in
RDWIA with **real potential**

- Full consistent description of exclusive kinematics 1e1p
- Integrate over the proton → get the correct inclusive cross section (=includes 'elastic' FSI!)
- For every event we **replace** the nucleon kinematics by the GENIE prediction (SuSAv2 implementation)

Nucleon kinematics from inclusive cross section in GENIE

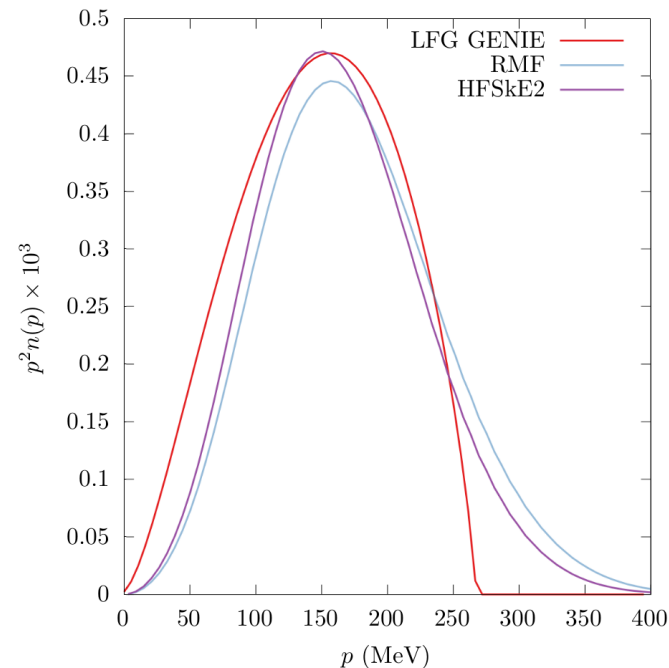
$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

1. Draw initial nucleon \mathbf{p}_m from $p^2 n(p)$ (e.g. LFG)

!! 2. Compute $E_m^2 = p_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$



Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

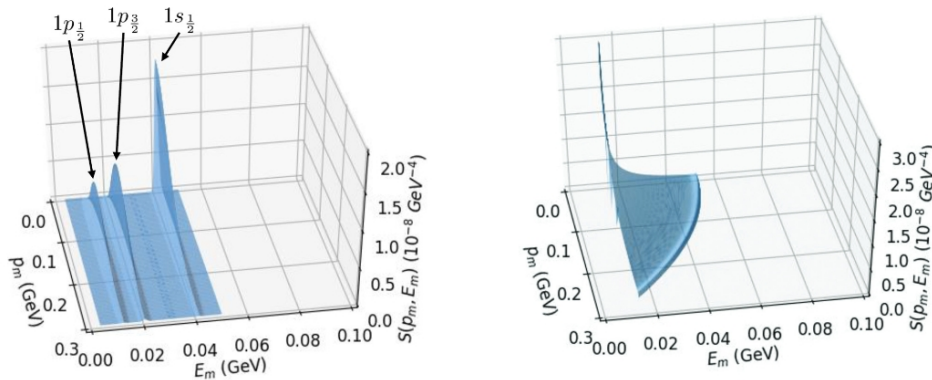
1. Draw initial nucleon \mathbf{p}_m from $p^2 n(p)$ (e.g. LFG)

!! 2. Compute $E_m^2 = \mathbf{p}_m^2 + M_N^2$

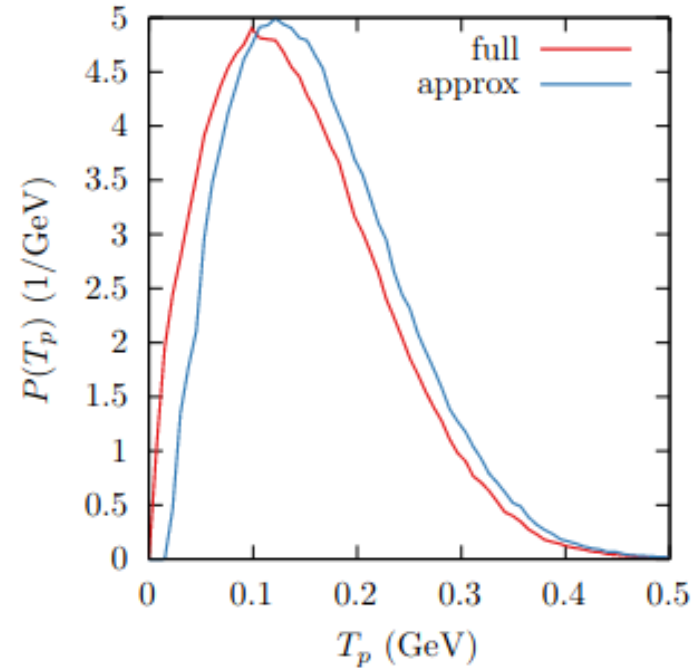
3. $E_N = E_m + \omega - E_b(q)$

Get a shift in total energy spectrum →

LFG $E_m^2 = \mathbf{p}_m^2 + M_N^2$ is not realistic!



[V. Orden & Donnelly PRC 100 044620]



E4nu kinematics $E = 1.15$ GeV

Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information \rightarrow Need to generate it in GENIE

1. Draw initial nucleon \mathbf{p}_m from $p^2 n(p)$ (e.g. LFG)

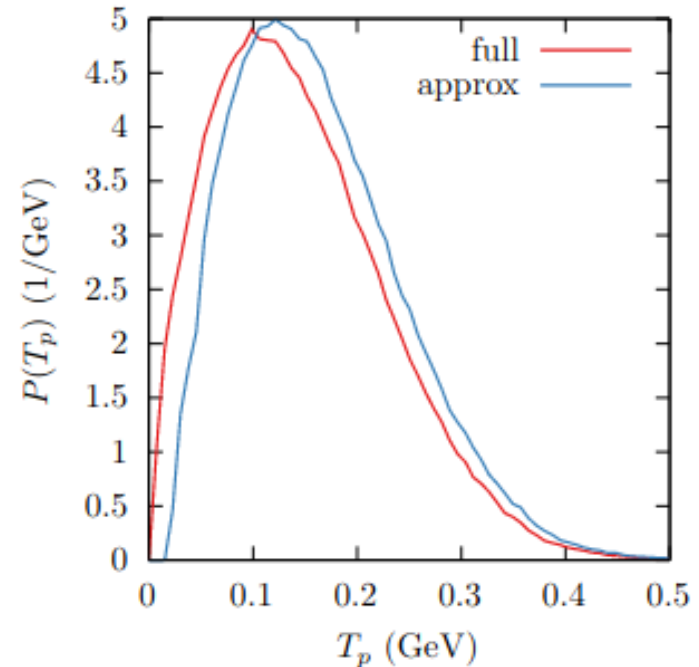
!! 2. Compute $E_m^2 = \mathbf{p}_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$

4. $k_N^2 = E_N^2 - M_N^2$

!! $|\mathbf{p}_m + \mathbf{q}| \neq k_N = \sqrt{E_N^2 - M_N^2}$

$\rightarrow \mathbf{k}_N = \frac{k_N}{|\mathbf{p}_m + \mathbf{q}|} (\mathbf{p}_m + \mathbf{q})$



Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

1. Draw initial nucleon \mathbf{p}_m from p^2 $n(p)$ (e.g. LFG)

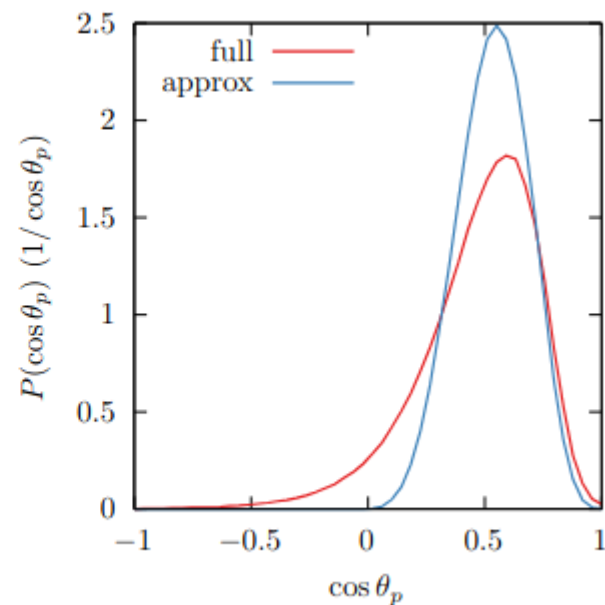
!! 2. Compute $E_m^2 = \mathbf{p}_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$

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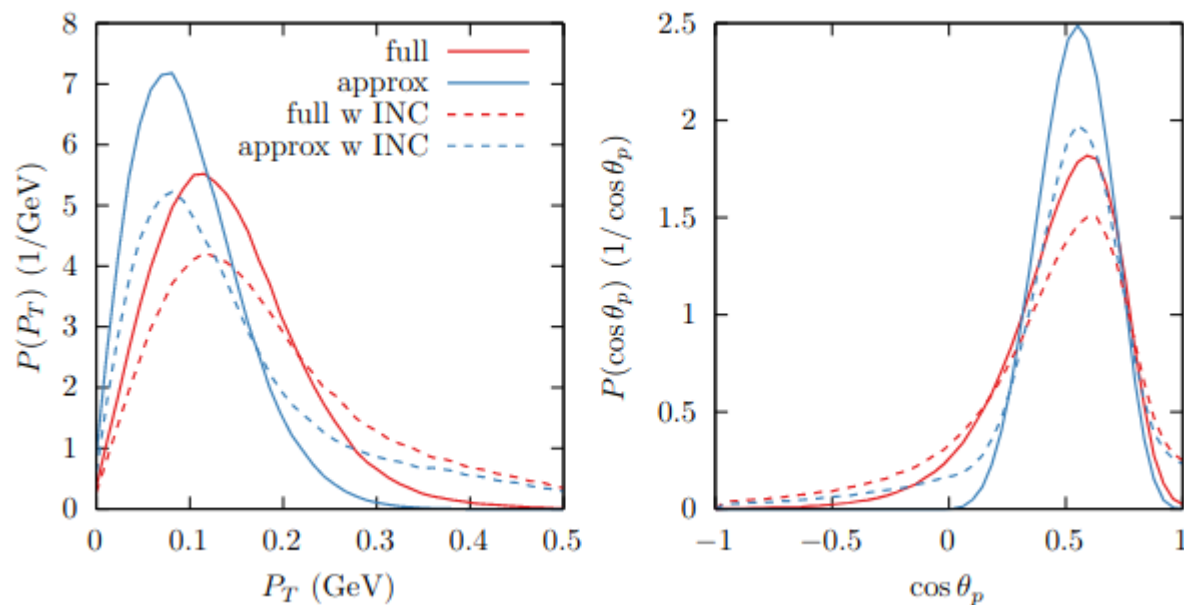


Serious differences in angular distributions!

Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

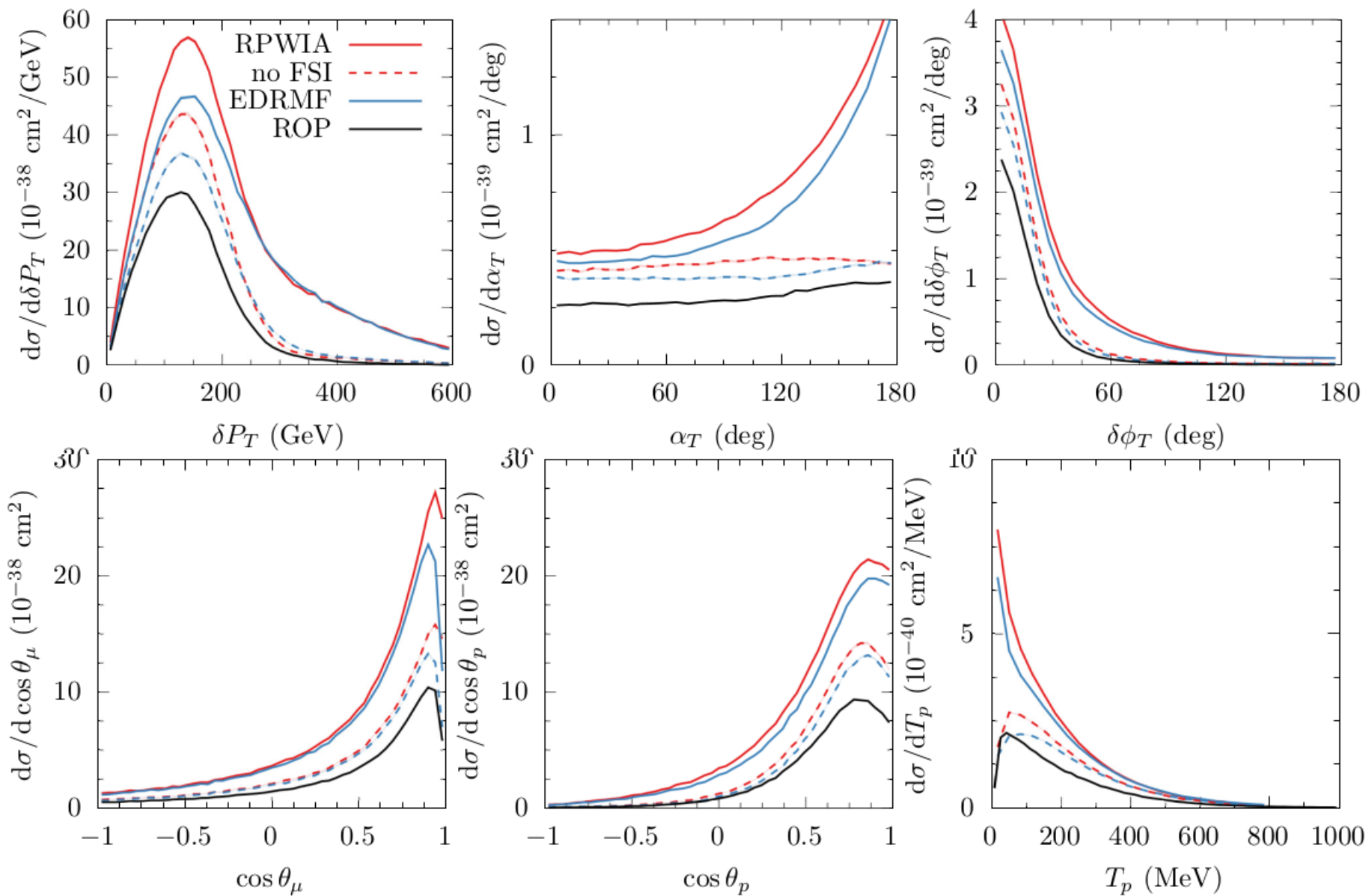


Results for e4nu kinematics $E=1.159$ including the GENIE cascade!

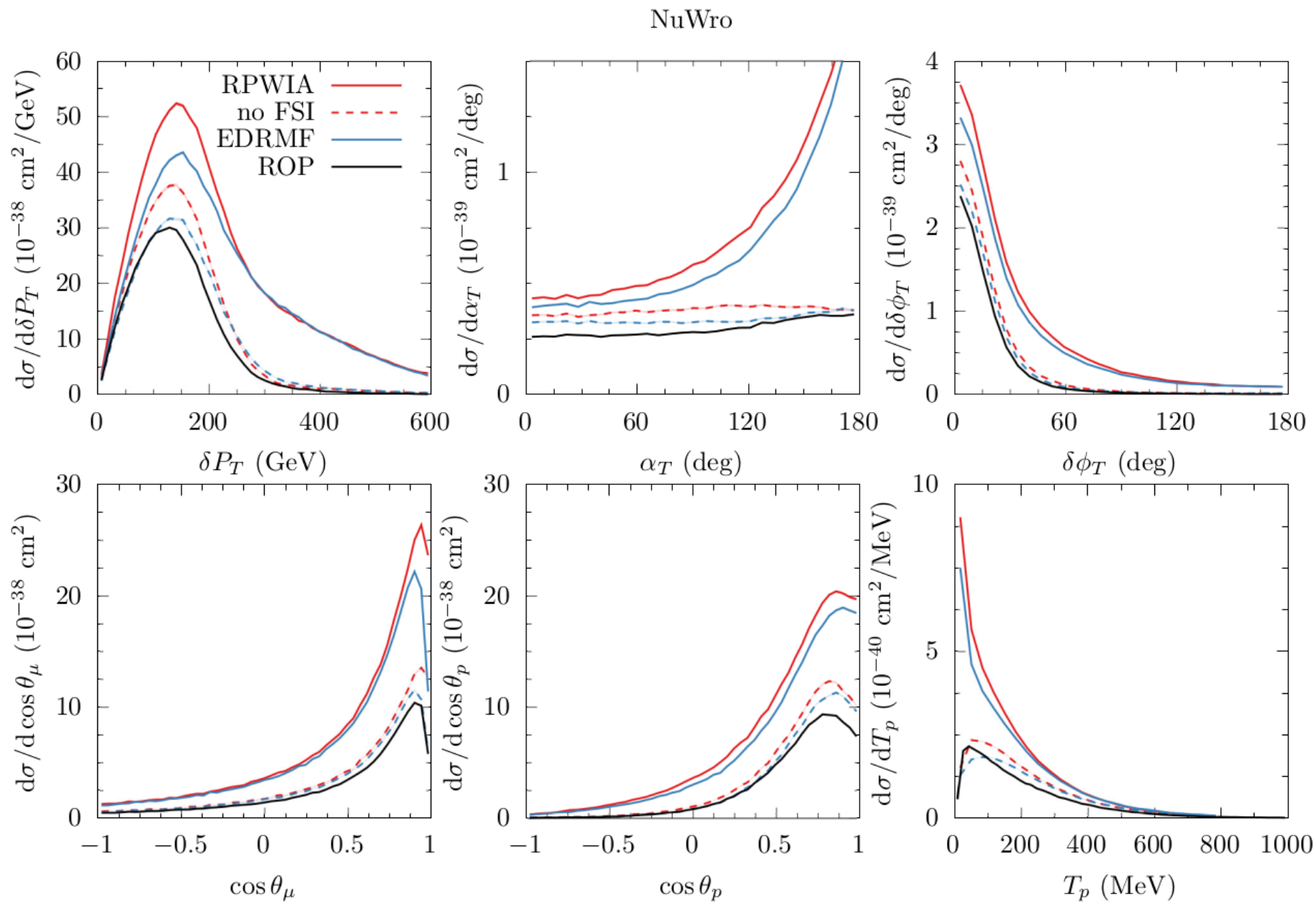
Shape differences biggest in P_T and angular distributions

NuWro with SRC effect in Mean-free path

NuWro

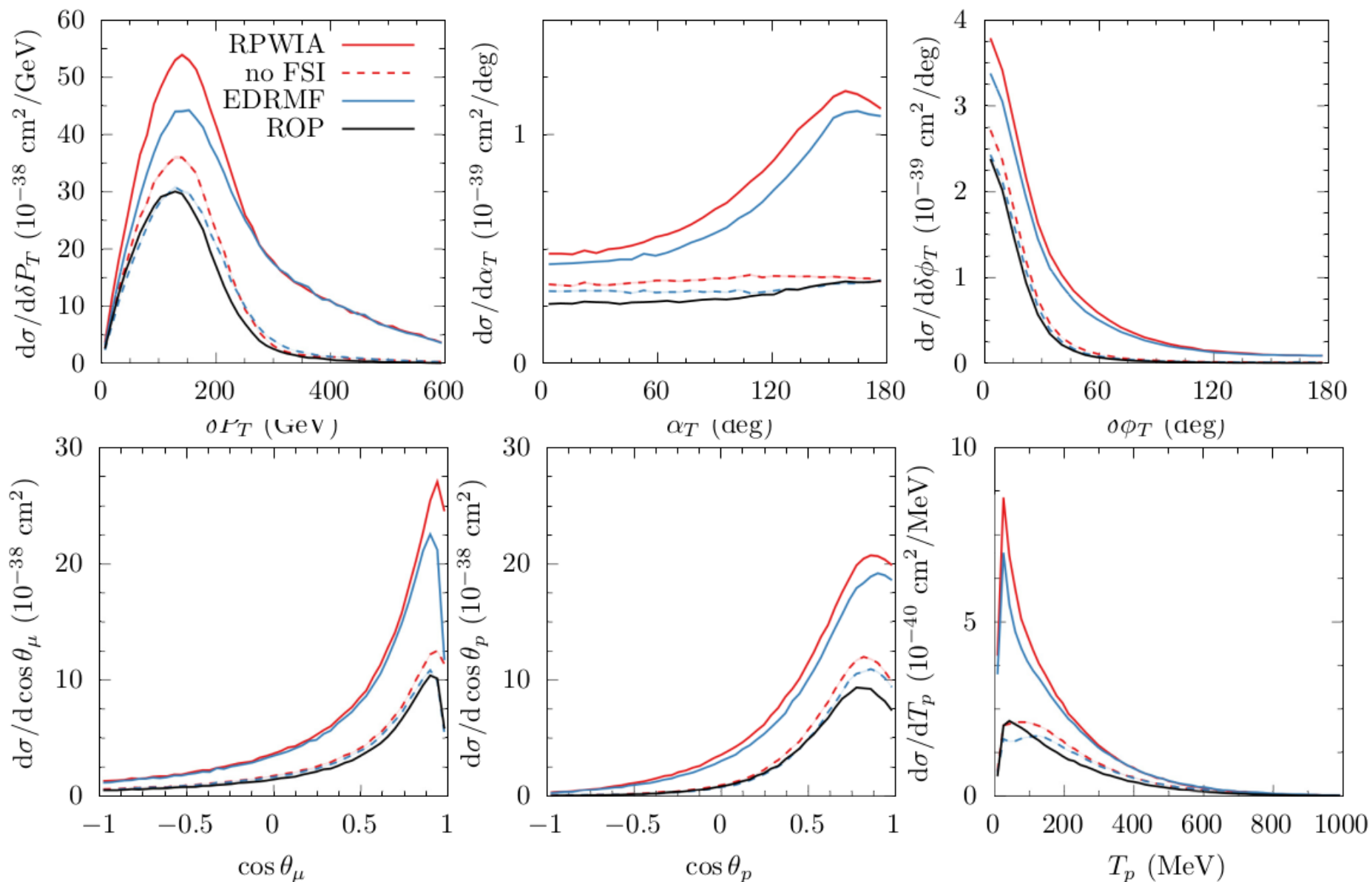


NuWro without SRC effect in Mean-free path



ACHILLES with Formation time

ACHILLES



ACHILLES without Formation time

ACHILLES

