

# Summary: transverse momenta of hadrons

---

Conveners: Harut Avakian, Nobuo Sato

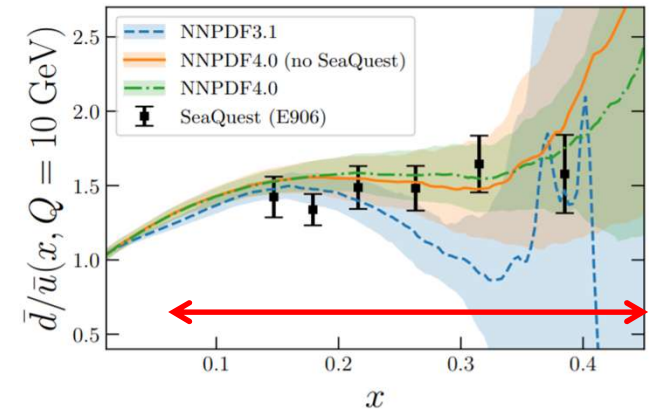
*Frascati, Dec 13, 2024*

- Main goals: relevant kinematics and observables
- What is SIDIS2.0?
  - Understanding hadronic correlations from  $ep \rightarrow e' \pi X$  to  $ep \rightarrow e' \pi \pi X$  and  $ep \rightarrow e' p \pi X$
  - Understanding “diffractive” contributions and impact on phenomenology
- Summary

# Understanding the QCD: from observables to QCD dynamics

## Moving from testing QCD to understanding it

→ Detailed studies of non-perturbative QCD dynamics in 3D space through measurements of multiplicities, and spin-azimuthal modulations



## JLAB uniqueness:

- Superior luminosity of CEBAF
- High resolutions of detectors
- Ability for multidimensional and multiparticle detection in the kinematics where non-perturbative effects are significant

Note: eptoproduction, in simplest case of a single hadron detected in the final state, is a measurement of observables in 5D space  $(x, Q^2, z, P_T, \phi)$ , 6D for transverse target,  $+ \phi_S$

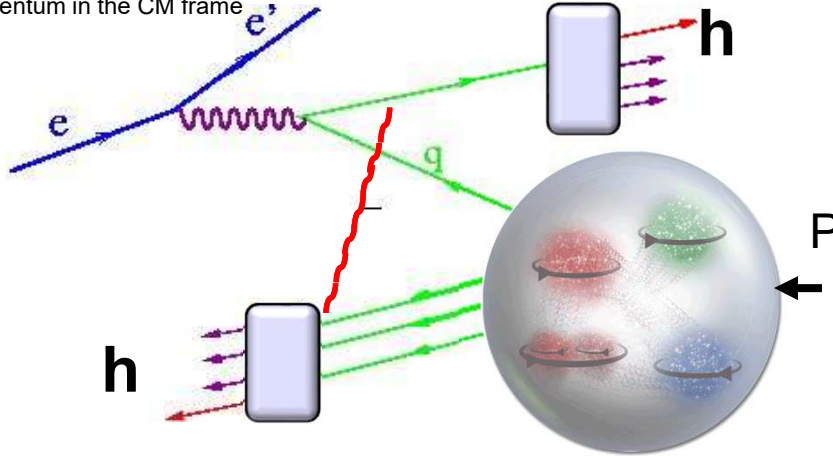
- Collinear SIDIS (last 50 years), is just the proper integration of observables, over  $P_T, \phi, \phi_S$
- Dominant fraction of hard scattering data available at  $Q^2 < 16 \text{ GeV}^2$

All the above makes JLab unique in disentangling the genuine intrinsic transverse structure of hadrons encoded in 3D partonic distributions with **controlled systematics in the kinematics dominated by valence quarks**

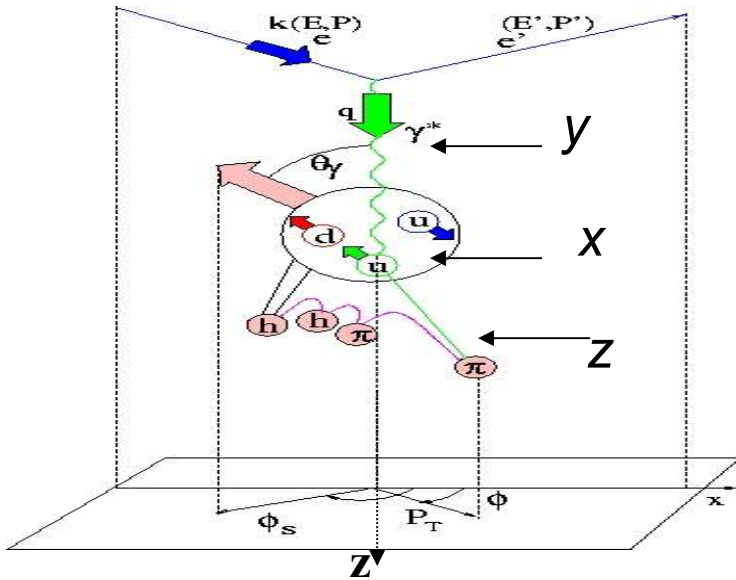
# Polarized Leptoproduction

$x_F$  – fractional longitudinal momentum in the CM frame

$x_F > 0$  (current fragmentation, CFR)



$x_F < 0$  (target fragmentation, TFR)



$$y = (qP)/(kP)$$

$$x = Q^2/2(qP)$$

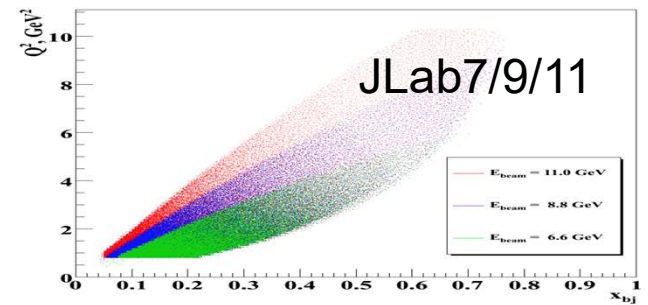
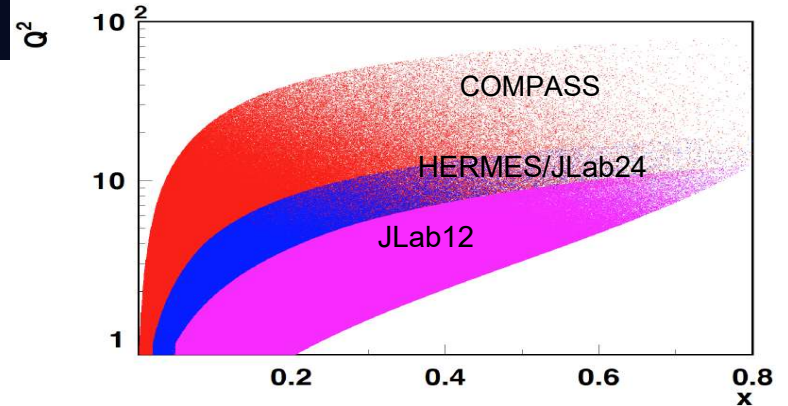
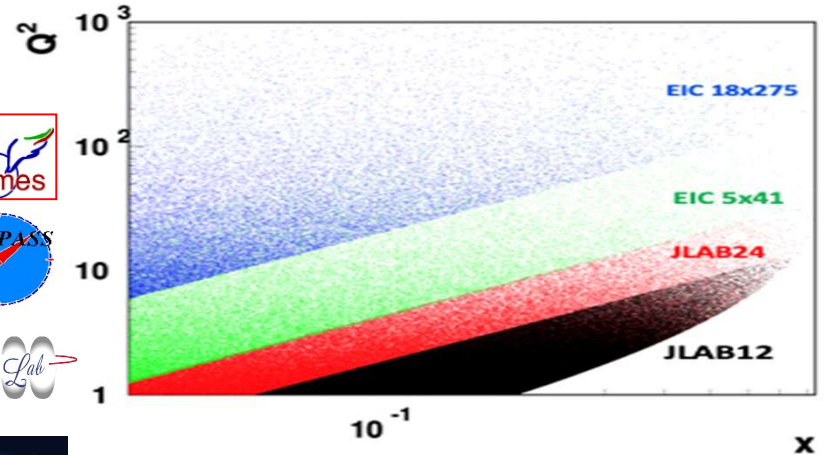
$$z = (qP_h)/(qP)$$

in  $eN \rightarrow e'hX$   
6D  $(x, Q^2, z, P_T, \phi, \phi_s)$

In exclusive limit

$$t = (P-P')^2$$

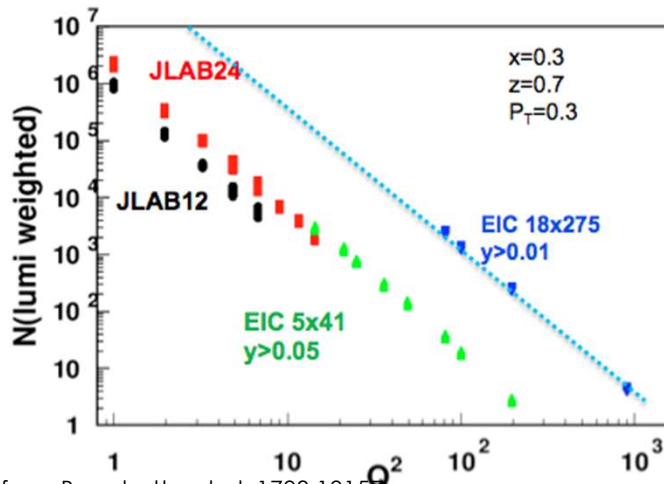
$$-t = (1-z)Q^2/x$$



# Structure functions and depolarization factors

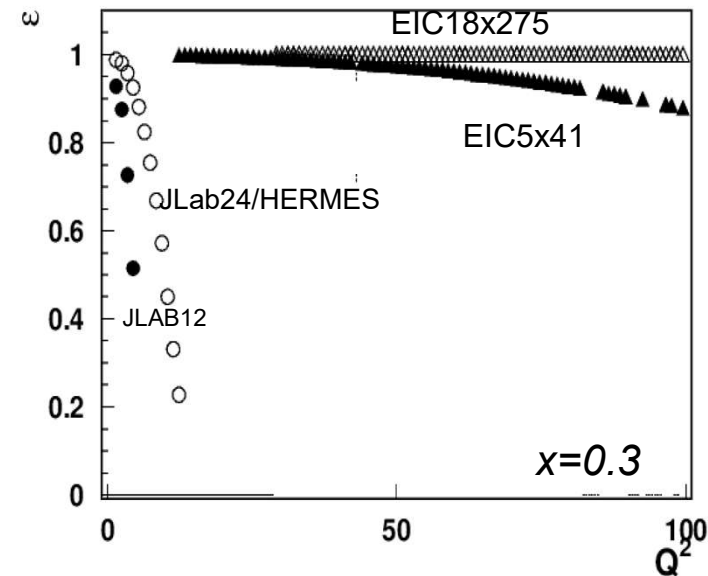
## x-section for eN→e'hX

$$\begin{aligned}
 \frac{d\sigma}{dx dy d\phi_S dz d\phi_h dP_{h\perp}^2} &= \frac{\alpha^2}{x y Q^2} \frac{y^2}{2(1-\varepsilon)} \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} \right. \\
 &+ \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin\phi_h F_{LU}^{\sin\phi_h} + S_L \left[ \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_h F_{UL}^{\sin\phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\
 &+ S_L \lambda_e \left[ \sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_h F_{LL}^{\cos\phi_h} \right] \\
 &+ S_T \left[ \sin(\phi_h - \phi_S) \left( F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} \right. \\
 &+ \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} + \sqrt{2\varepsilon(1+\varepsilon)} \sin\phi_S F_{UT}^{\sin\phi_S} \\
 &+ \left. \left. \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right] + S_T \lambda_e \left[ \sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} \right. \right. \\
 &+ \left. \left. \sqrt{2\varepsilon(1-\varepsilon)} \cos\phi_S F_{LT}^{\cos\phi_S} + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}
 \end{aligned}$$

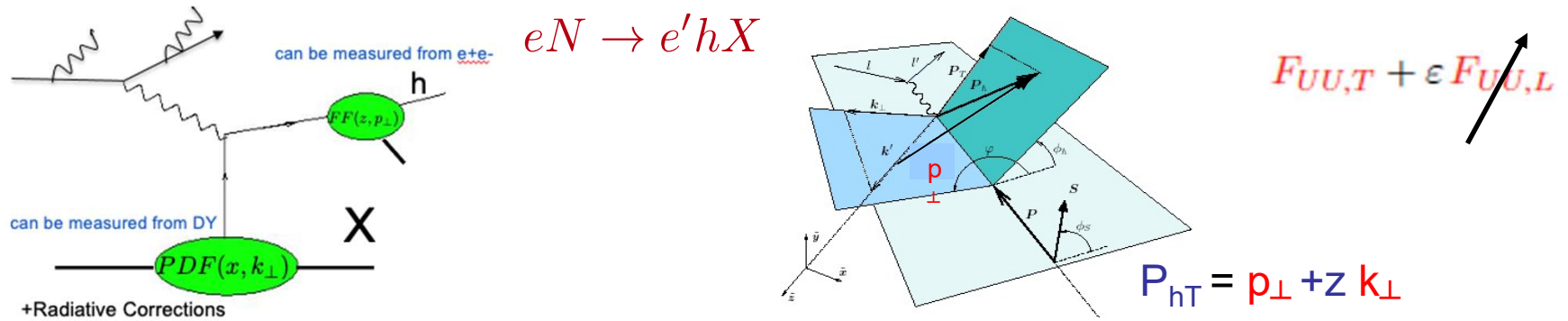


x-section from Bacchetta et al, 1703.10157

- Measurements of correlations of spin and transverse momenta, encoded in the Structure Functions (SFs), provides direct access to details of QCD dynamics
- Where are contributions from diffractive processes?
- Combination of statistics and depolarization factors defines measurable SFs
- At higher energies (EIC), observables surviving the  $\varepsilon \rightarrow 1$  limit (FUU, FUL, Transversely pol. FUT)



# SIDIS as THE theory describes it



Probability to produce 1 or 2 hadrons in single photon exchange:

$$\frac{d\sigma}{dx dy d\phi_S dz d\phi_h dP_{h\perp}^2} = x \sum_q \mathcal{H}_{UU,T}^q(Q^2, \mu^2) \int d^2\mathbf{k}_\perp d^2\mathbf{P}_\perp f_1^q(x, \mathbf{k}_\perp^2; \mu^2) D_1^{q \rightarrow h}(z, \mathbf{P}_\perp^2; \mu^2) \delta(z\mathbf{k}_\perp - \mathbf{P}_{hT} + \mathbf{P}_\perp) + Y_{UU,T}(Q^2, \mathbf{P}_{hT}^2) + \mathcal{O}(M^2/Q^2)$$

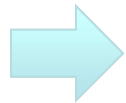
$F_{UU,T}(x, z, \mathbf{P}_{hT}^2, Q^2)$       TMD Parton Distribution Functions      TMD Parton Fragmentation Functions

- Factorization allows description using distribution functions (TMD-PDF) and fragmentation functions (TMD FF)
- $X \rightarrow$  multiplicity of unobserved hadrons **LARGE**, and x-section doesn't depend on  $X$

Conclusions in case of apparent disagreement:

“much bigger/smaller” defined in comparison with experiment

- Factorization is broken?
- Unaccounted terms may contribute



**Data has it all!!! Dealing with unaccounted terms:**

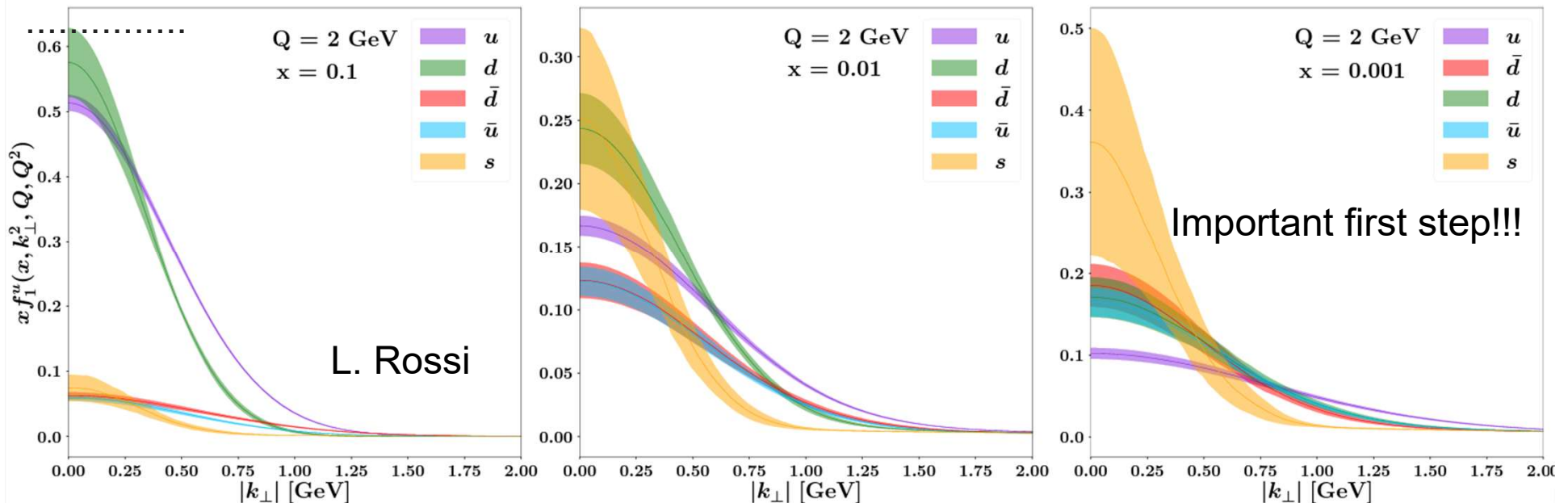
- Theory accounts for them (ex. VMs)
- Experiment measures and excludes them!!! (ex.VMs)**

# $P_T$ distributions provide access to $k_T$ -distributions

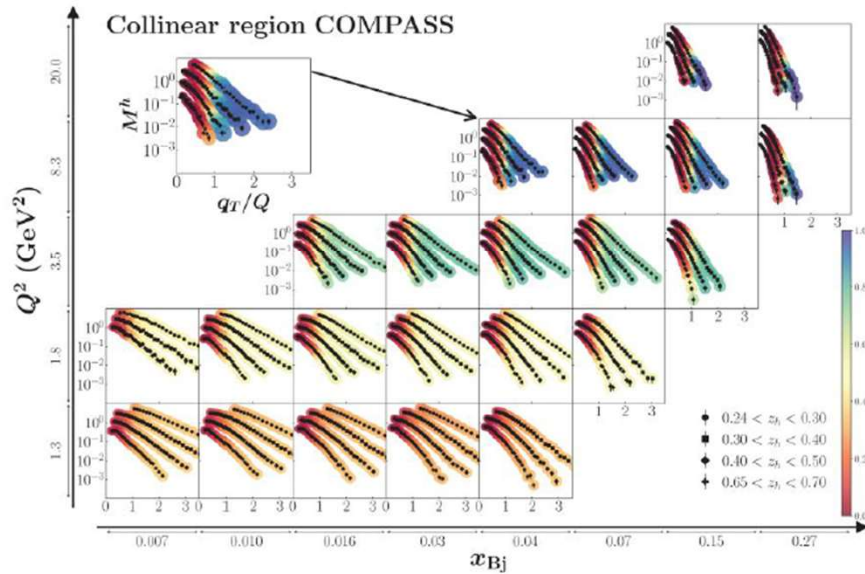
Understanding of  $P_T$ -distributions of hadrons in SIDIS, most critical for TMD measurements in the multidimensional space, providing access to QCD dynamics!!!

Expected:

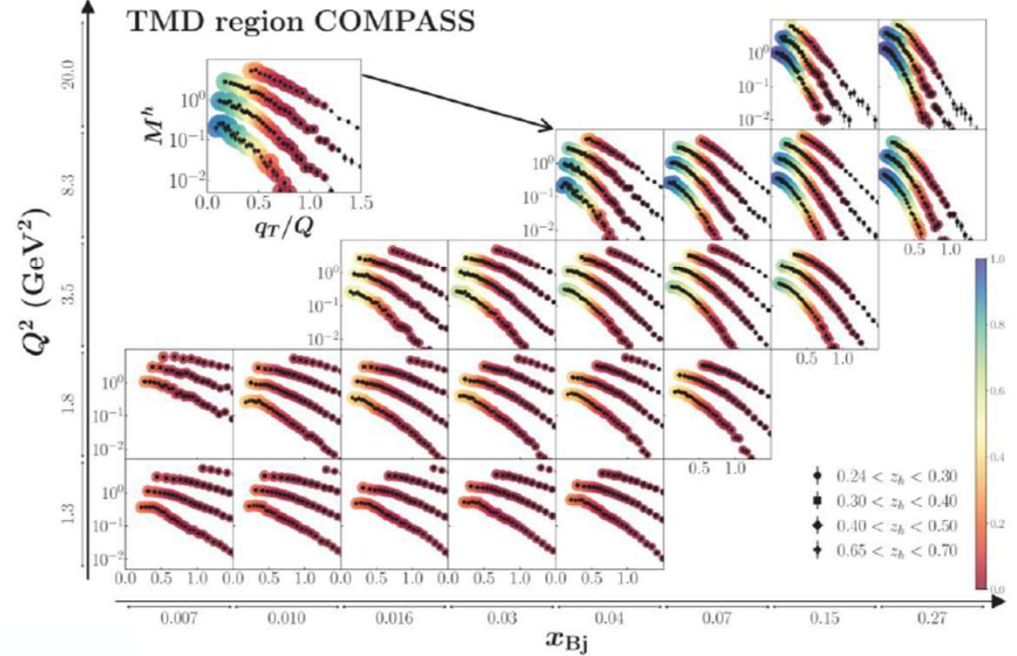
- 1) Perturbative contributions in the  $p_T$ -range covered by polarized SIDIS likely minor
- 2) Significantly wider in  $k_T$  distributions of u-quarks with spin opposite to proton spin (possible sign flips in asymmetries related to polarization of partons)
- 3) Significantly wider in  $k_T$  distributions of d-quarks (possible sign flips in asymmetries related to polarization of partons)
- 4) Significantly wider in  $k_T$  sea quark distributions (study contributions dominated by sea, K-,...)



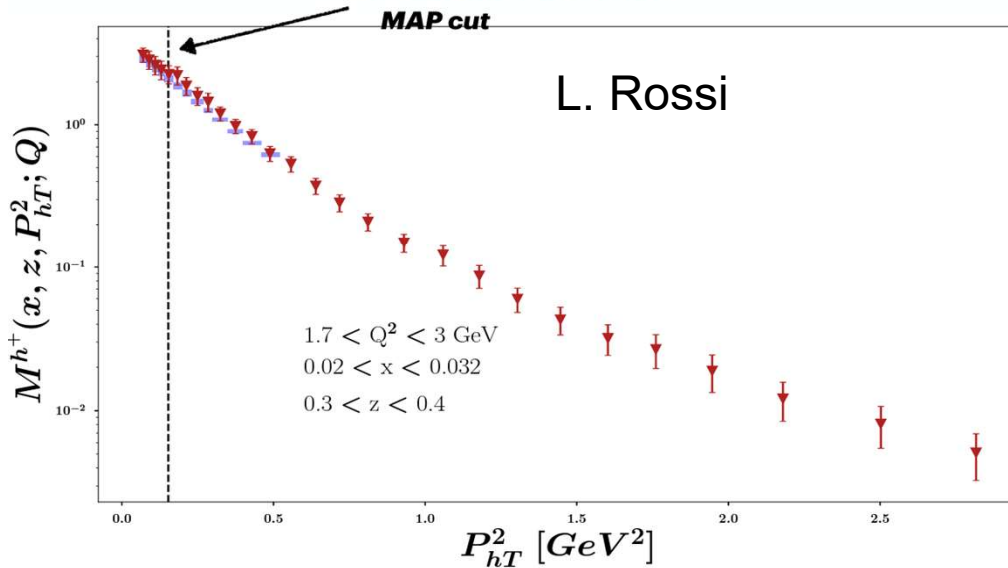
## B. Parsamyam



## JAM, JHEP 04 (2022) 084



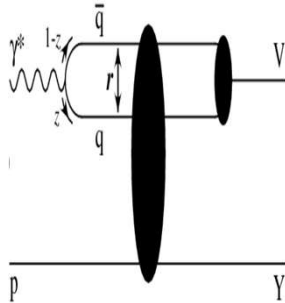
## What's next?



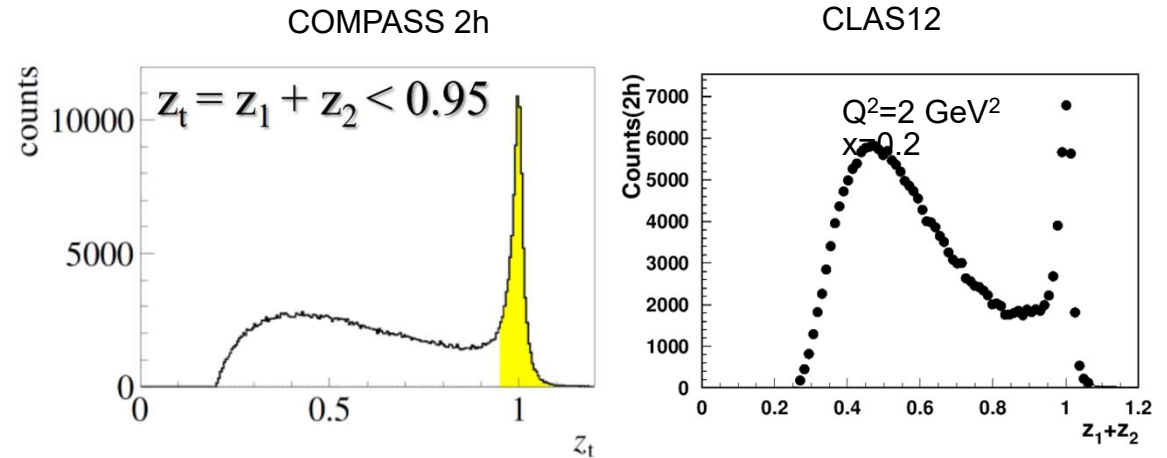
- The Theory works only for low  $P_T$
- Need major changes in the phenomenology to capitalize on the JLab22 upgrade, increasing significantly the  $P_T$  coverage

# Longitudinal photons and diffractive rhos

## Diffractive VMs ( $\rho^0$ )



$$\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = K(x, Q^2, y) [F_{UU,L} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} F_{UU}^{\cos} \dots]$$



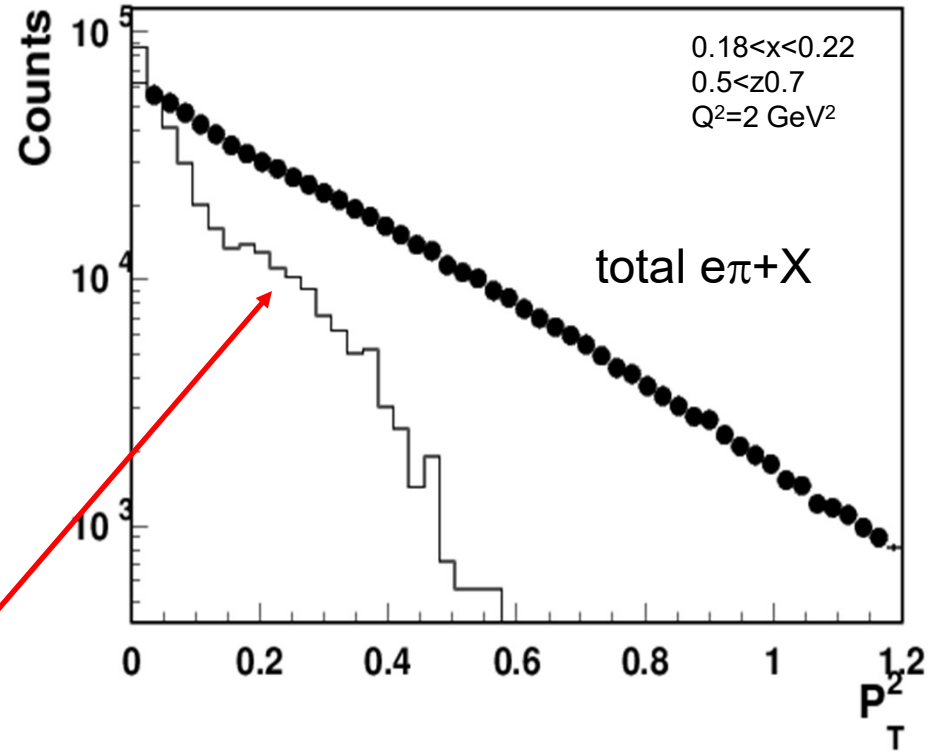
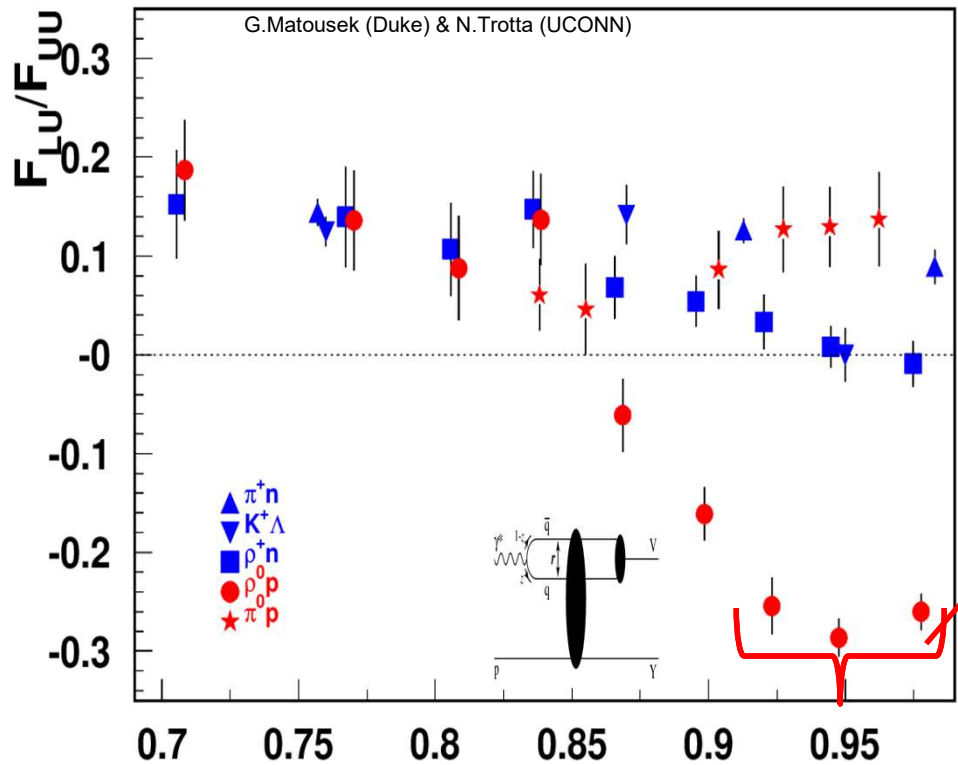
- CLAS12 measurements indicate the 2hadron exclusive sample is dominated by “diffractive  $\rho^0$ ” produced at very small  $t$
- JLab provides possibility of detailed studies of those rhos, crucial for interpretation in terms of TMDs of SIDIS data in general, and for EIC in particular.

- Estimated ~20% contributions from  $\rho$  to charged pion SIDIS, consistent with ~10% of diffractive DIS in inclusive DIS
- Indication: most longitudinally polarized  $\rho^0$  (note: higher the  $Q^2$  lower is  $\epsilon$ )

**Studies of exclusive processes require high resolution and multidimensional measurements !!!**



# Contributions of “diffractive $\rho^0$ s” in SIDIS



- Indications for diffractive:
- Major differences in observables dominated by quarks
  - Major differences with charged rho+

Estimated ~20% contributions from  $\rho$ , mainly show up at low  $P_T$  in SIDIS

The “diffractive”  $\rho$  will bias extractions of TMDs, unless properly subtracted in multidimensional space of SIDIS measurements.

# Excluding the “diffractive” rho from SIDIS

Depending on how we exclude the exclusive rho we can have several versions of experimental samples of inclusive hadrons, each with their own bias:

1) Standard SIDIS ( $eN \rightarrow ehX$ ,  $h=\pi, K, \dots$ ) within the full accessible kinematics, corrected for acceptance and RC, measured in the multidimensional space

→  $e\pi X$  biased with respect to theory by presence of contributions from diffractive rho, contributing to ~20% of counts, in low  $P_T$ , with contributions to SSA ~10 times higher

2) Standard SIDIS ( $eN \rightarrow e\pi X$ ) within the full accessible kinematics, corrected for acceptance and RC, measured in the multidimensional space, with subtracted in multi-D bins for rho0 contributions (“rho-subtracted SIDIS”)

→ requires measurements of pions from diffractive rho in multidimensional space, means detailed studies of SDMEs of rhos, requiring good precisions and huge statistics, also for all polarization observables, extensive validation needed, little known RC

3) SIDIS subsamples ( $eN \rightarrow ep\pi X$ ,  $eN \rightarrow e\pi\pi X$ ) within the full accessible kinematics, allowing clear elimination of rho0 contributions using cuts on missing masses of  $epX$  or  $e\pi\pi X$

(“rho-free SIDIS”)

→ biased by the presence of additional hadron in TFR ( $epX$ ) or CFR ( $eppX$ ), may need a new phenomenology

requires measurements of dependence on  $M_X$  to understand the bias,

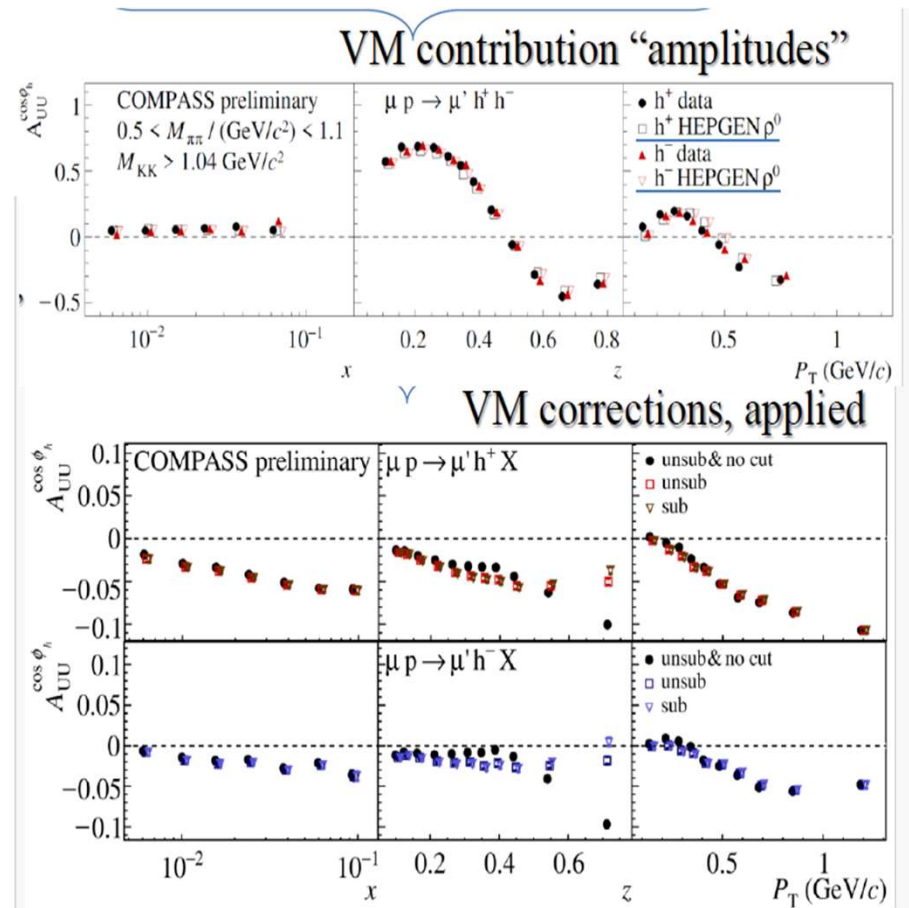
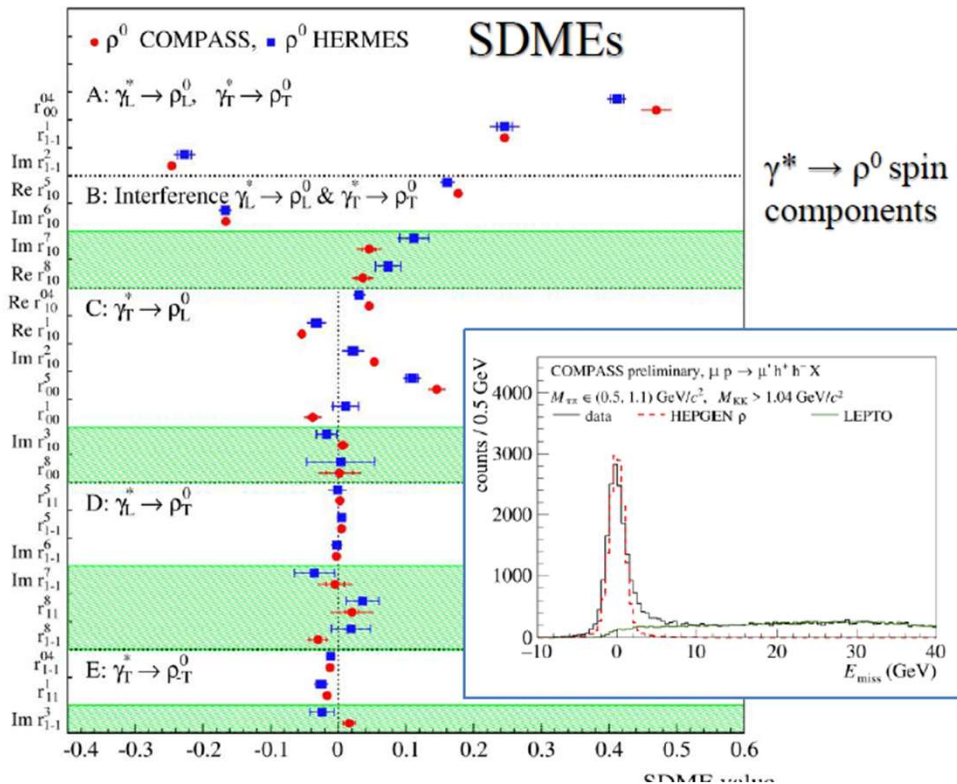
Theory should be able to evaluate the bias from the presence of an additional hadron

# Attempts to separate the “diffractive” contributions

Procedure: “rho-subtracted SIDIS”(B. Parsamyan)

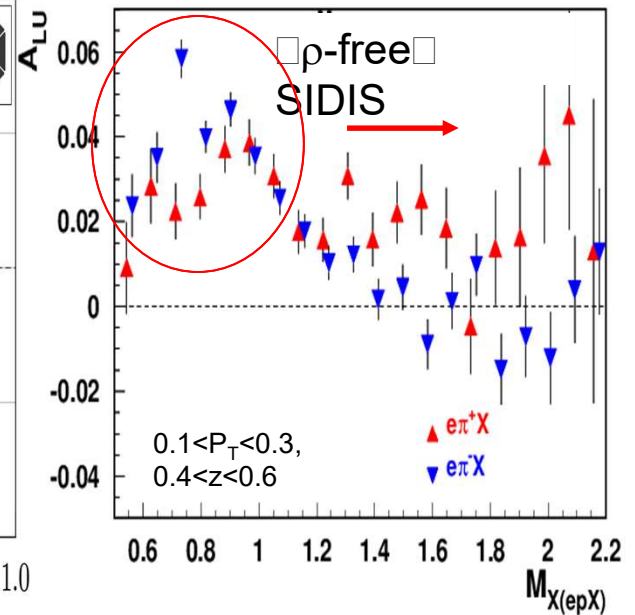
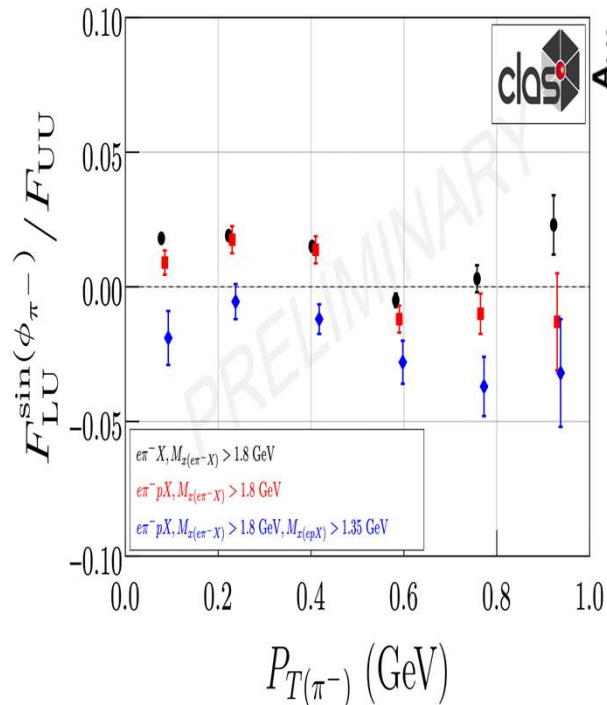
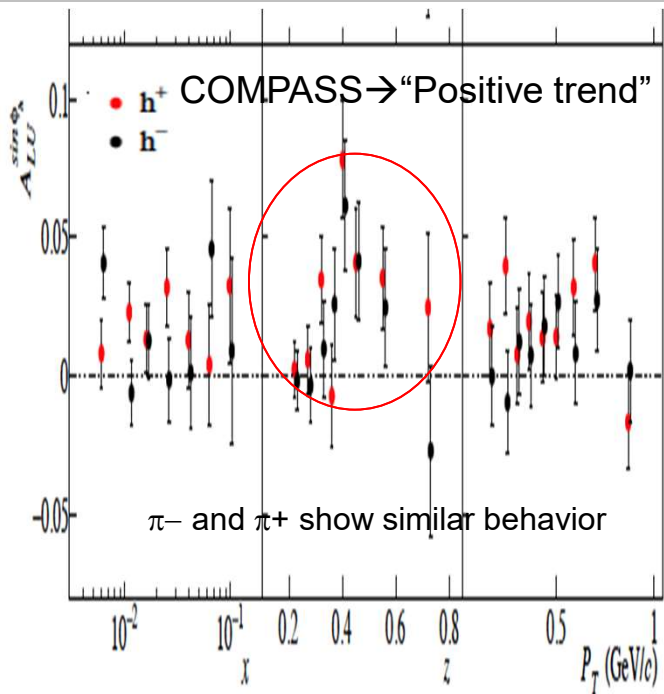
- 1) Measure exclusive rhos
- 2) Extract SDMEs (>20 azimuthal modulations)
- 3) Develop MC
- 4) Evaluate the correction

COMPASS, EPJC (2023) 83 924



Extracted modulations have weird  $P_T$  and most importantly  $Q^2$  dependence

# Exclusive $\rho$ contributions to $\pi$ : $P_T$ -dependence



COMPASS  $\rightarrow$  "Positive trend" also reproduced when additional proton in TFR detected (red)

- The same sign and size of  $\pi^+$  and  $\pi^-$  SSA indicates the  $\rho$  may not be properly subtracted (require detailed MC studies, which require proper SDMEs)
- While VM contributions are  $\sim 20\%$  in multiplicities **in SSA they can be  $> 100\%$**

Need proper MC for  $\rho$  for all polarization states  $\rightarrow$



"The greatest enemy of knowledge is not ignorance, it is the illusion of knowledge."

— Daniel J. Boorstin

# First attempts to account for rho (M. Cerutti)

## $\rho$ -subtraction exercise

see Harut's talk

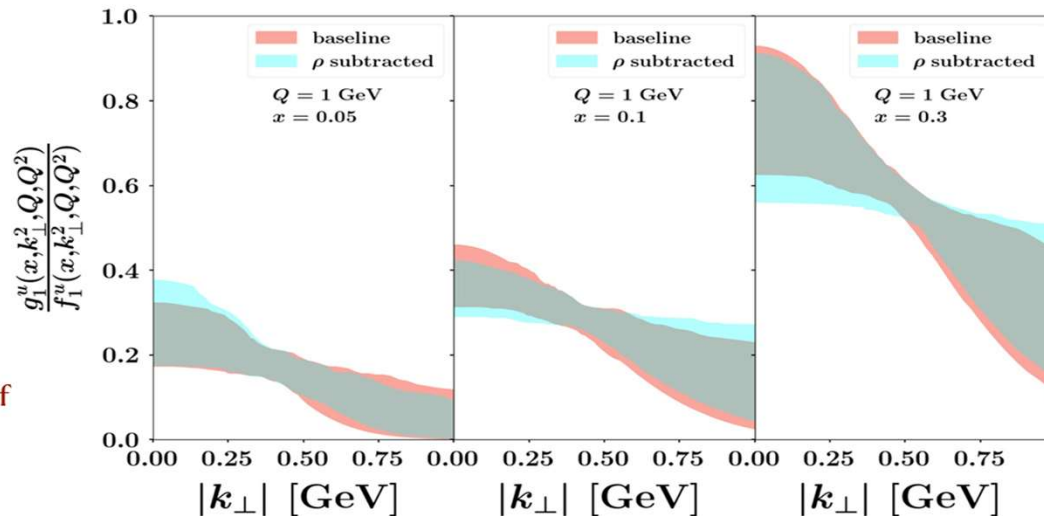
“Effective” subtraction of  $\rho$ -meson (diffractive) contribution

$$A_1^{\rho-sub} = \frac{A_1}{\Omega_\rho}$$

“deRHOification” factor  $\Omega_\rho = 1 + e^{-6P_{hT}^2}$

**Only for final-state  $\pi$**

- baseline fit (only  $\pi$ )  
 $\chi^2/N = 1.10$
- $\rho$ -subtracted fit  
 $\chi^2/N = 1.86$
- Impact on the shape of  $g_1/f_1$  ratio



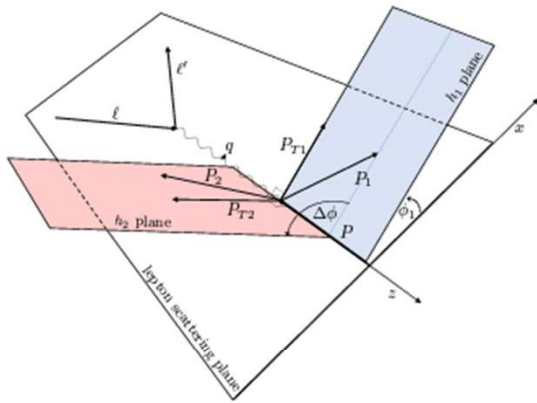
Marco Zaccheddu

Need a self consistent procedure with proper rho subtraction

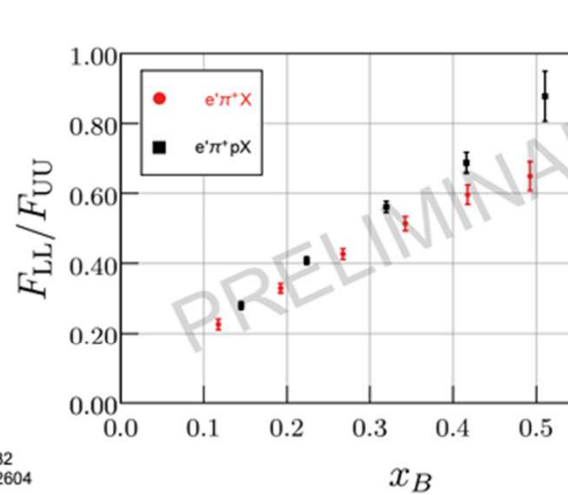
Pixelization of TMDs

# Back-to-back (dSIDIS) Formalism

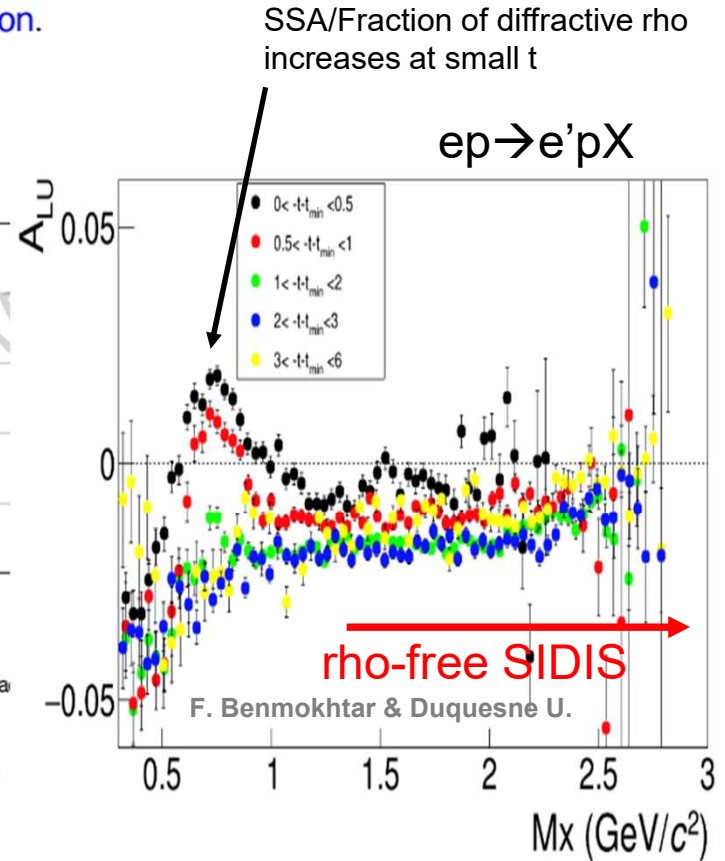
- When two hadrons are produced “back-to-back”<sup>1,2</sup> with one in the CFR and one in the CFR contains a convolution of a **fracture function** and a **fragmentation function**.
- Leading twist access to all quark-nucleon polarization combinations.



1. M. Anselmino et al., Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132  
 2. M. Anselmino et al., Phys. Lett. B. 713 (2012), 317-320, [hep-ph] 1112.2604



higher  $A_{LL}$  in b2b SIDIS, possibly due to exclusion of non-fa



Timothy Hayward

12/10/24

8

- Detection of the target fragment proton allows “rho free” SIDIS measurements
- Signal from rho clearly increases at small  $t$ , consistent with large negative SSA ( $r_{00}^8$  SDME, arising from longitudinal rhos produced from transverse photons)

# SUMMARY

- The superior luminosity of CEBAF, high resolutions of detectors, and ability for multidimensional and multiparticle detection, makes the JLab unique facility to study the non-perturbative QCD dynamics in 3D space in the kinematics dominated by valence quarks
- Interpretation of observables in charged pions SIDIS, and spin and azimuthal asymmetries, require detailed understanding of longitudinal photon contributions, vector mesons, and exclusive  $\rho^0$  in particular.
- Exclusive “diffractive rhos” due to sensitivity to gluon contributions, have very different from any other (quark dominated) processes spin and azimuthal asymmetries, and may provide understanding of “diffractive” phenomena
- Measurements of exclusive rhos at 22 GeV will be critical for interpretation of high energy data, also providing a bridge between JLab and EIC
- Development of “rho subtracted” SIDIS, will require detailed measurements of SDMEs in multidimensional space, validated by “rho free” measurements in multi-D space.
- Need a major upgrade for SIDIS phenomenology to extract TMDs with controlled systematics (SIDIS2.0 or reforming SIDIS) to use in most efficient way major improvements of  $P_T$  and  $Q^2$  coverage at 22 GeV

- 
- Support slides



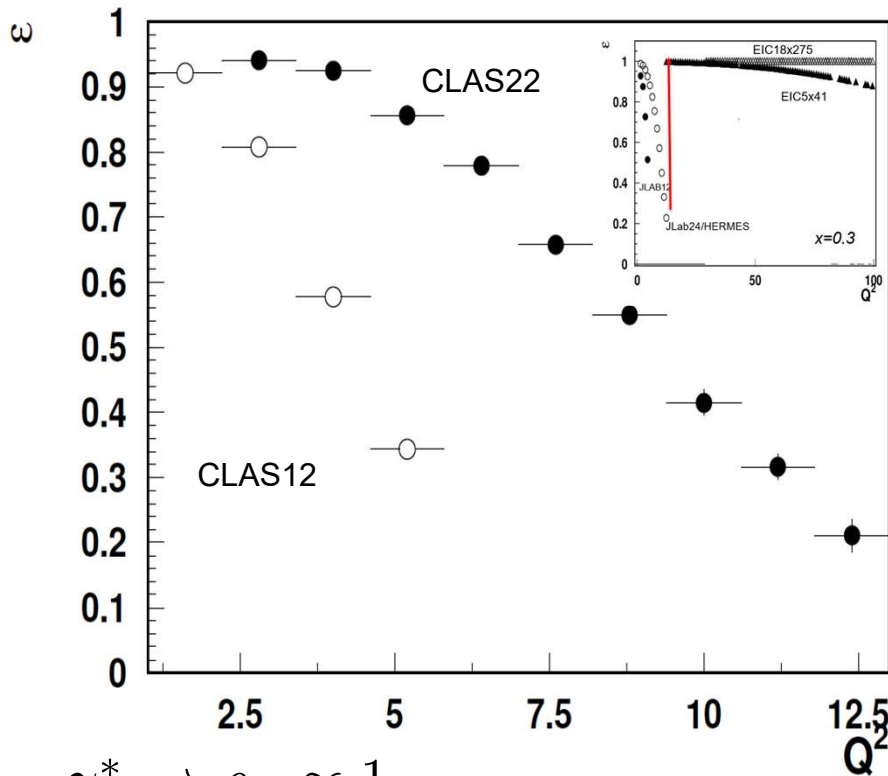
# JLab22 white paper: Hadronization and Transverse Momentum

---

Studies of Transverse Momentum Distributions (TMDs) require studies of transverse momentum dependences of SIDIS observables (multiplicities/asymmetries) in multidimensional space

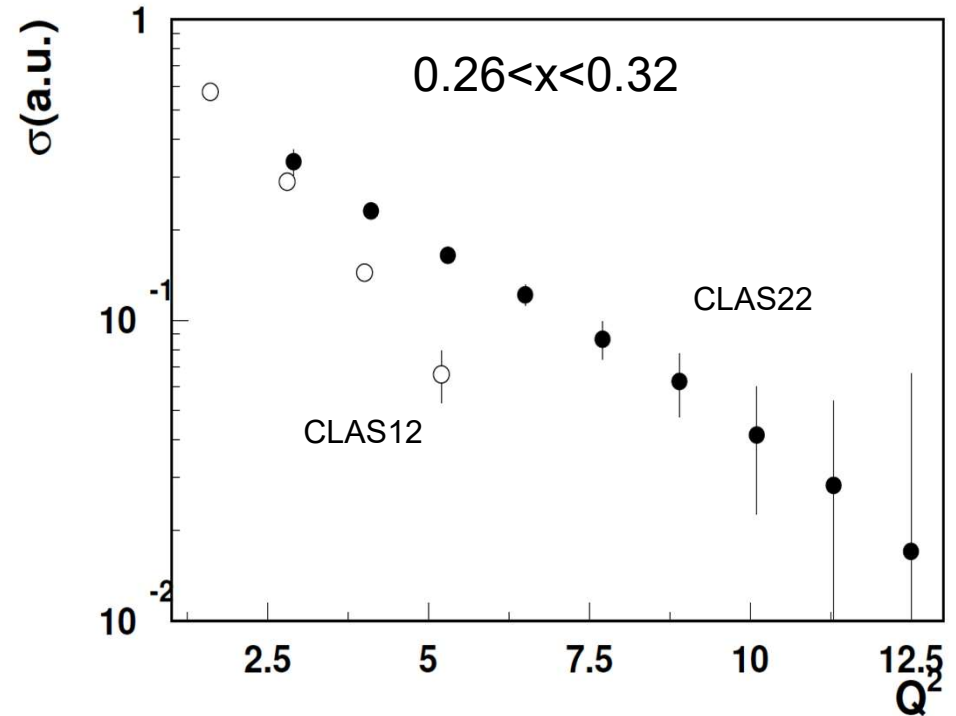
- Understanding the systematics in SIDIS studies
- physics backgrounds:
  - What is “diffractive DIS” and “diffractive SIDIS”?
- what we’ need to apply THE theory with controlled systematics?
  - separate different contributions to x-section (locate SF of interest)
  - separate different contributions to a given SF from different mechanisms (ex. longitudinal photon contributions)
  - separating the kinematics of current and target fragmentation
  - understanding the role of hadron correlations in SIDIS (impact of VMs)
  - use  $Q^2$ -dependent measurements as a unique tool to validate the interpretation of results

# Exclusive $\rho^0$ : extending the $Q^2$ with JLab22



$$\gamma_L^* \rightarrow \rho_L \propto 1$$

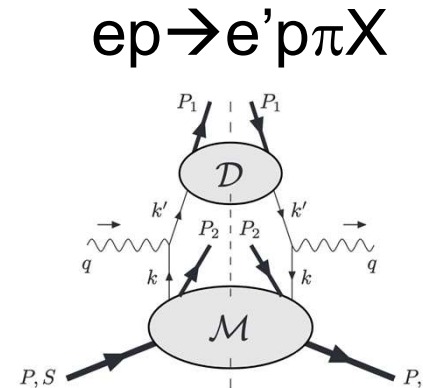
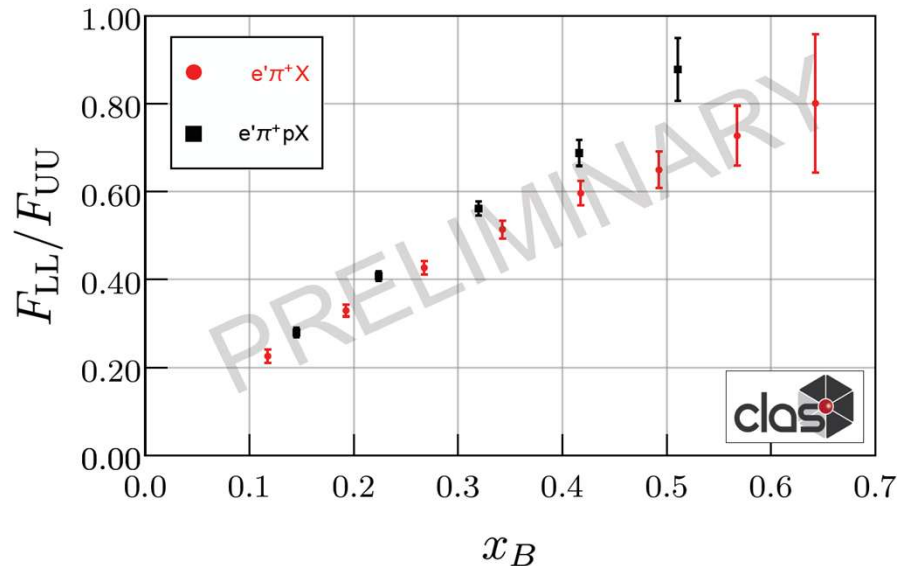
$$\gamma_T^* \rightarrow \rho_L \propto \sqrt{-t}/Q$$



Longitudinal photon contributions kinematically enhanced at higher energies

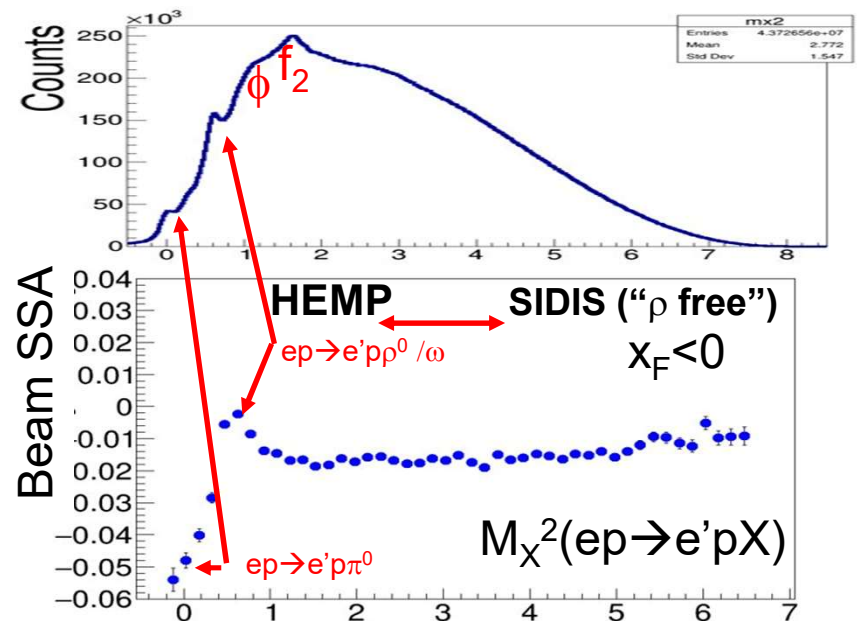
- Beam SSA provides important info on production of longitudinal rho from transverse photons
- Range in  $Q^2$  increases significantly with 22 GeV upgrade, allowing detailed studies at beyond 10 GeV<sup>2</sup> and providing a bridge to EIC

# Longitudinally polarized quarks in B2B SIDIS



N/q	U	L	T
U	$\hat{u}_1$	$\hat{l}_1^{\perp h}$	$\hat{l}_1^h, \hat{l}_1^{\perp}$
L	$\hat{u}_{1L}^{\perp h}$	$\hat{l}_{1L}$	$\hat{l}_{1L}^h, \hat{l}_{1L}^{\perp}$
T	$\hat{u}_{1T}^h, \hat{u}_{1T}^{\perp}$	$\hat{l}_{1T}^h, \hat{l}_{1T}^{\perp}$	$\hat{l}_{1T}, \hat{l}_{1T}^h, \hat{l}_{1T}^{\perp}, \hat{l}_{1T}^{\perp h}$

Detection of proton allows elimination of exclusive rho!



Possible theory formalisms:

- Formalism based on fracture functions (Anselmino, Barone, Kotzinian (back-to-back, b2b, hadron production, DSIDIS))
- Semi-exclusive processes, involving GPDs/GTMDs on proton side (TFR) and FFs on pion side (CFR) Yuan and Guo
- Differences in  $A_{LL}$ , due to different weights on PDFs can provide additional info on impact of possible ingredients
- Measurements of  $A_{LL}$  for  $\rho^0$  indicate very small values, and can be one of the reasons for higher  $A_{LL}$  with protons with a  $M_X$  cuts above 1.5 GeV (excluding exclusive  $\rho^0$ )
- Higher  $A_{LL}$  will change the phenomenology used last 40 years in DIS and SIDIS studies!!!

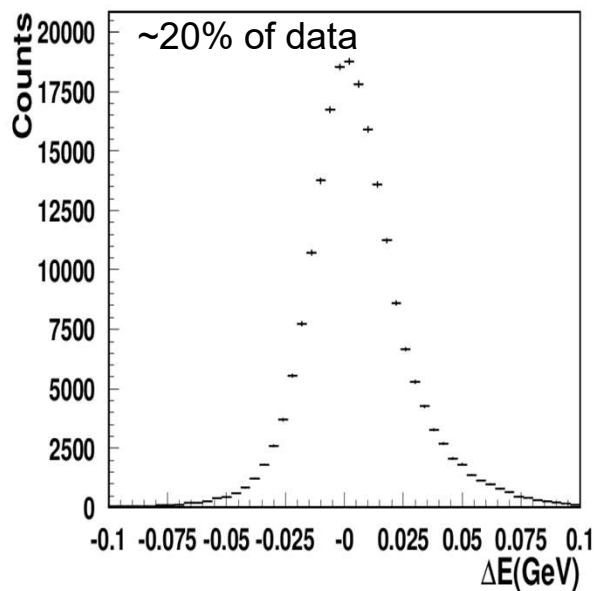
# Understanding exclusive rhos and SDME validations

Exclusivity condition defined by the missing Energy:

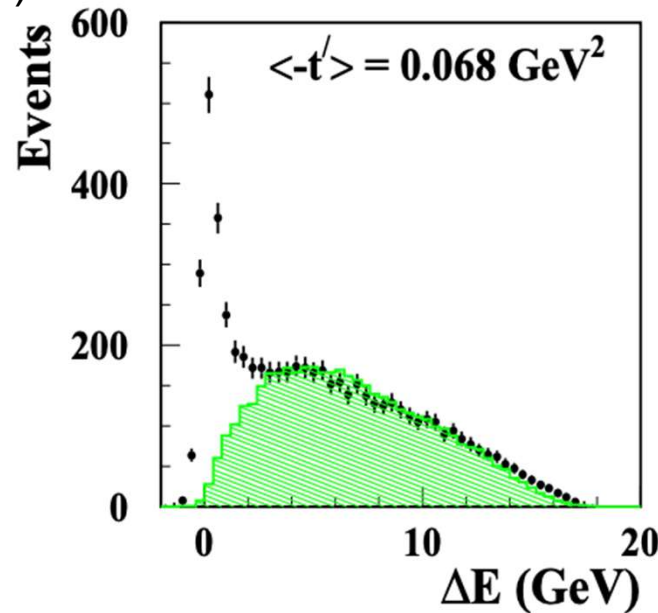
$$M_X^2 = (p + q - p_{\pi^+} - p_{\pi^-})^2$$

$$E_{\text{miss}} = \frac{M_X^2 - M^2}{2M}$$

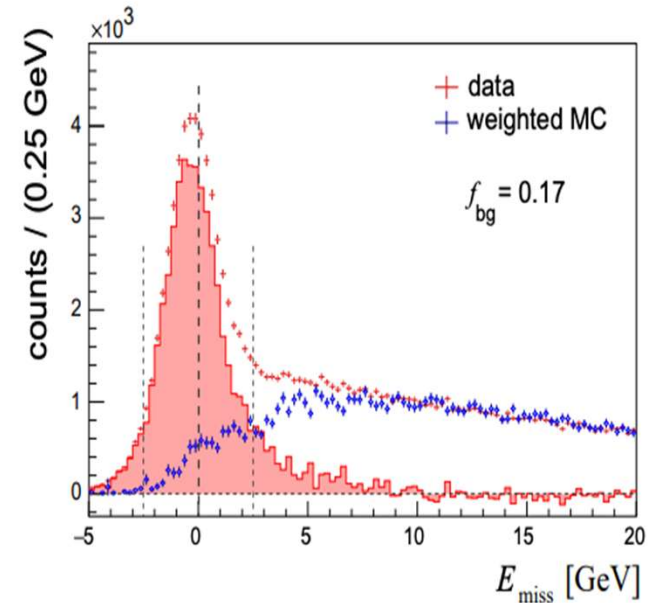
CLAS12 (width <0.1GeV)



HERMES (width ~0.6GeV)

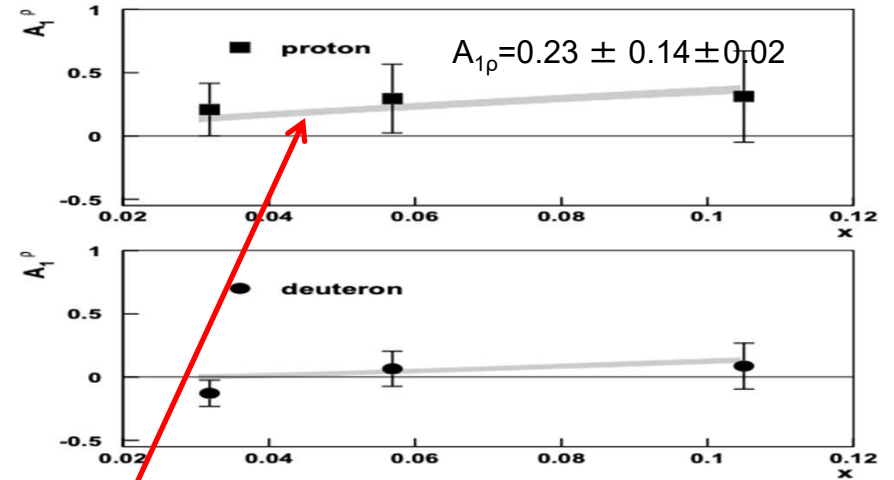
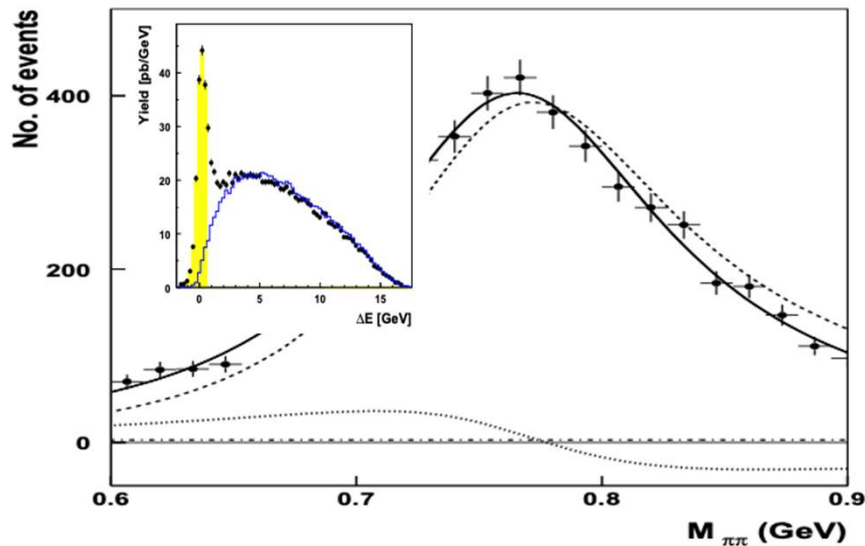


COMPASS (width ~2GeV)



- Guarantying the “exclusivity” requires good resolutions (get worse at higher energies)
- Subtraction procedure relays on normalization, based on exclusive limit of LUND-MC
- All distributions have have tails, indicating the RC may not be negligible
- Extraction of SDMEs, will require validation in the multi-D space (significant samples)

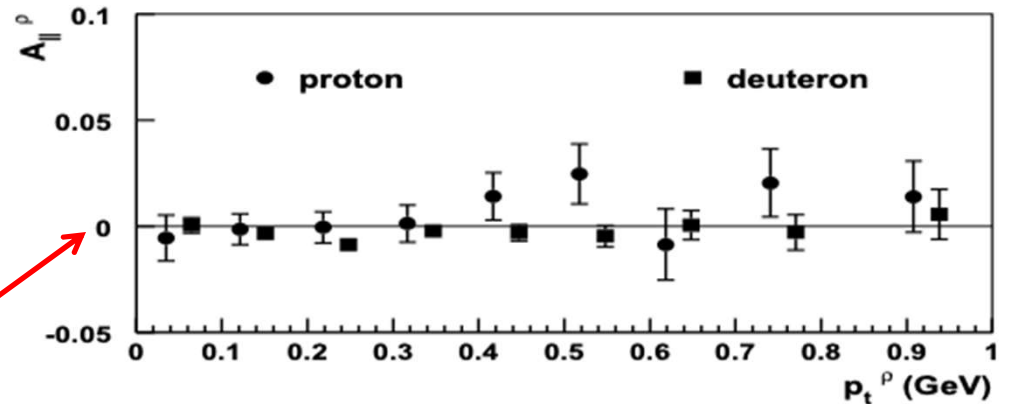
# $A_{LL}$ studies of exclusive $\rho^0$ : HERMES



1D plots can be really misleading need multi-D

For a proper extraction of multiplicities and spin-azimuthal modulations of exclusive  $\rho$ s, clean separation is needed for  $\rho^0$ , and longitudinally polarized  $\rho^0$  signal, in particular

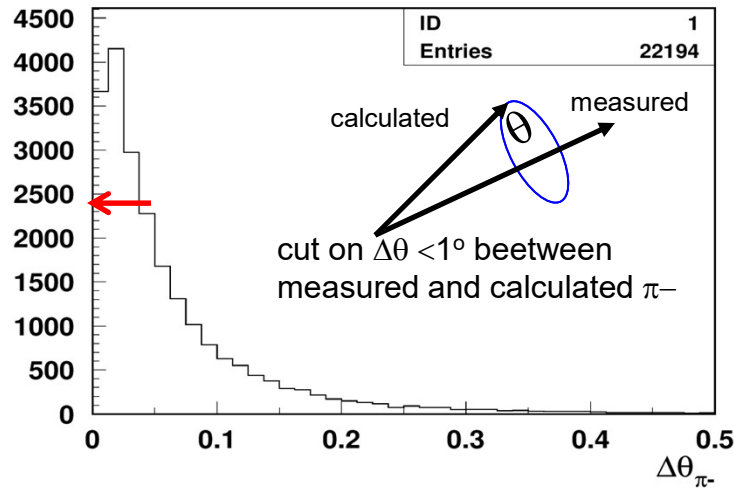
At low  $P_T$ , where the background is smaller, the asymmetry indeed tend to be negative



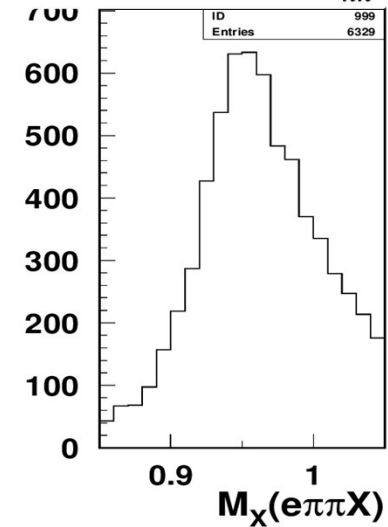
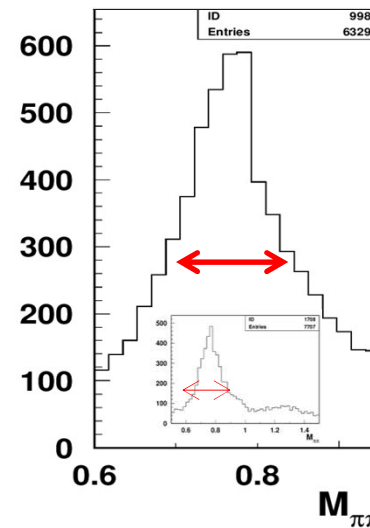
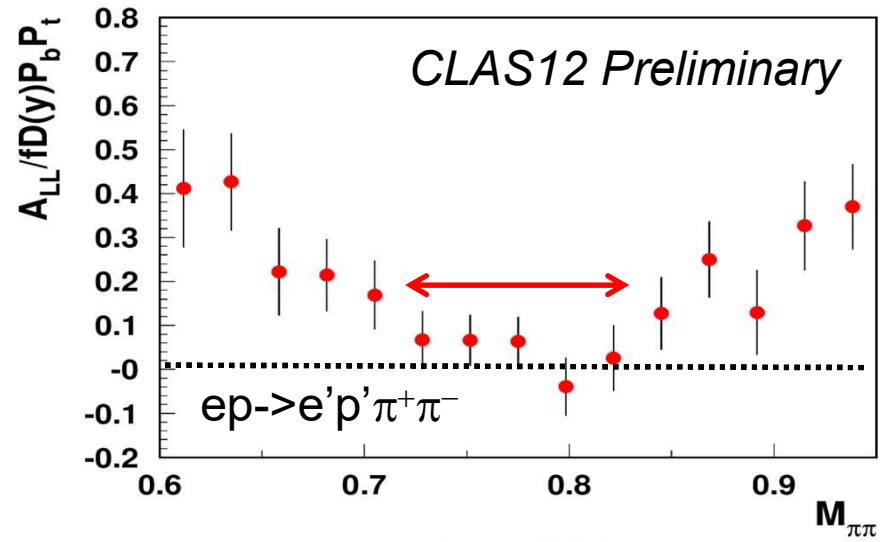
Accounting of  $\rho^0$  will change the phenomenology of helicity distributions

# Studies of $\rho^0$ impact with longitudinally polarized $\text{NH}_3$ target

## Separating exclusive dihadrons



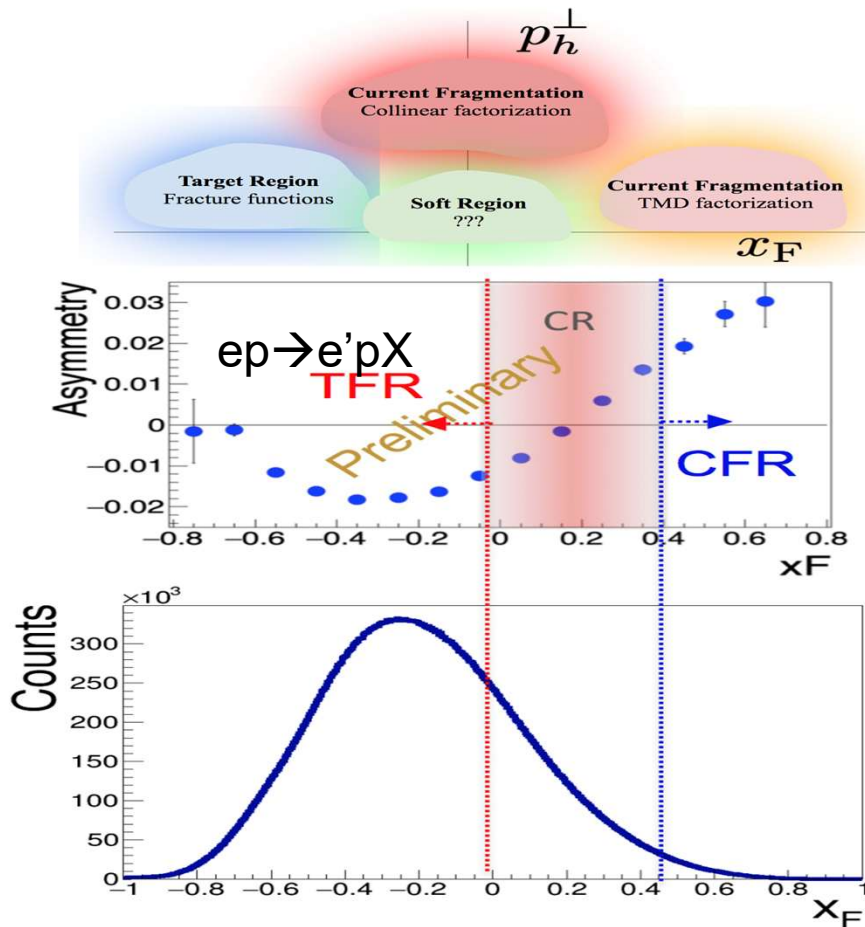
- Require the angle of negative pions is within a degree from calculated from  $e', p, \pi^+$  assuming exclusive  $e', p, \pi^+\pi^-$  event.
- Measurements of  $A_{LL}$  for  $\rho^0$  indicate very small values (with  $\sim 10\text{-}20\%$  bckg, likely negative  $\sim -2\text{-}10\%$ ), and can be one of the reasons for higher  $A_{LL}$  with protons with a  $M_X$  cuts above 1.35 GeV (excluding exclusive  $\rho^0$ )



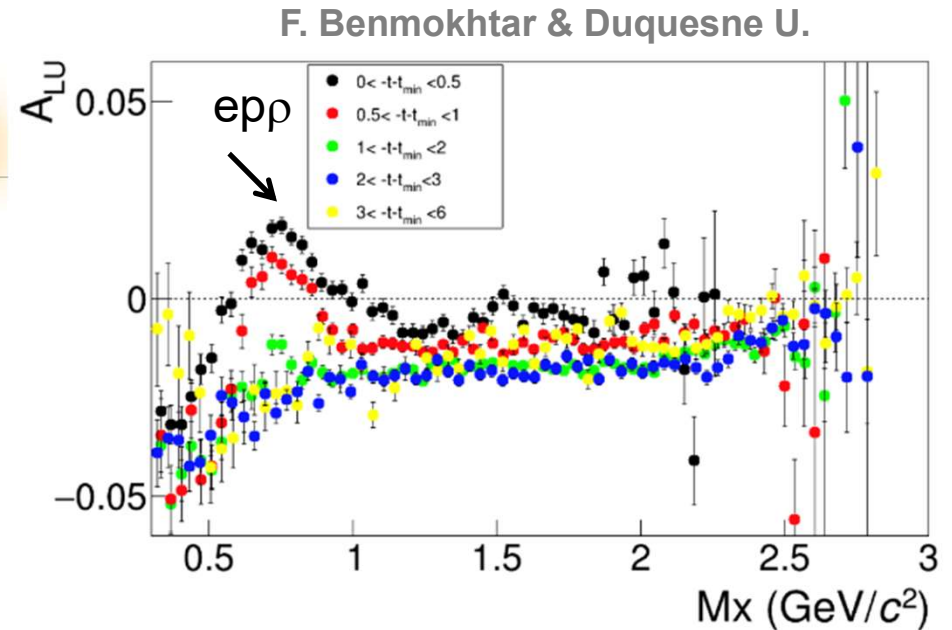
Need clear separation of hydrogen from  $\text{NH}_3$  and diffractive exclusive  $\rho^0$ s from exclusive  $\pi^+\pi^-$

# Beam SSAs as a tool to separate regions and contributions

Separating Target Fragmentation Region TFR from Current fragmentation region (CFR)



Negative sign of the SSA (plateau) defines the TFR dominance

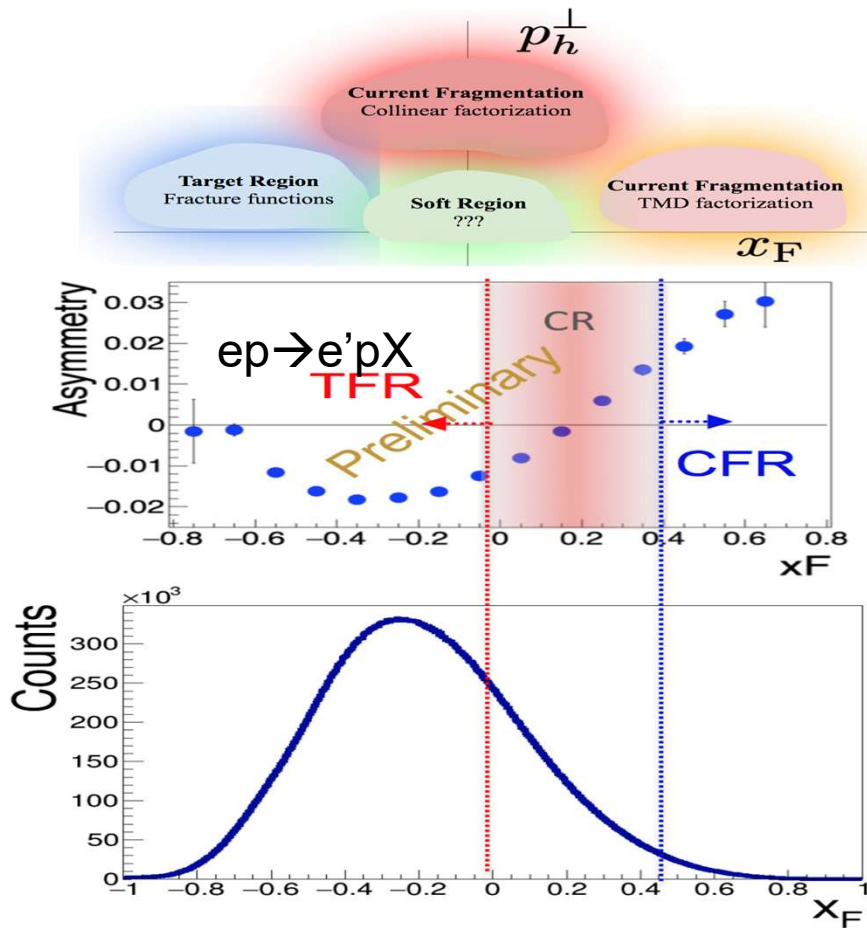


Major difference only for protons at small  $t$ !

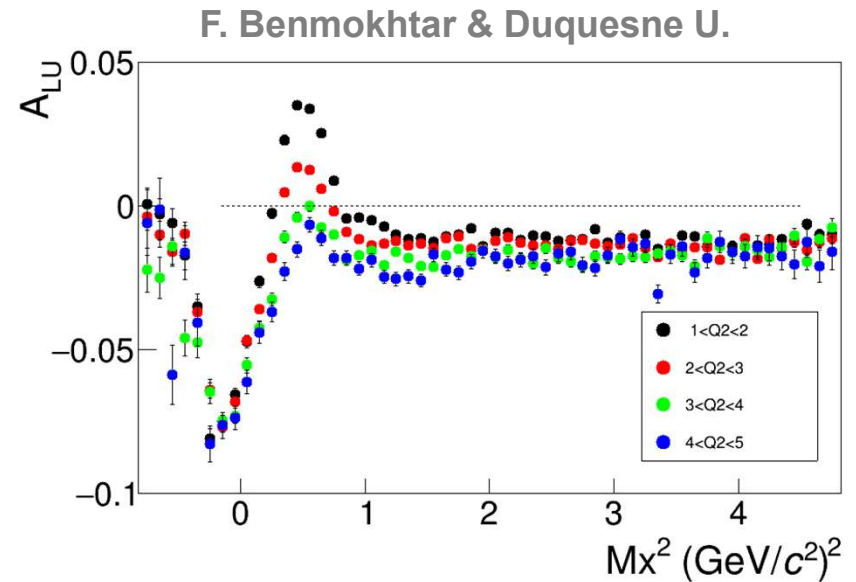
SIDIS beam SSA can serve as a tool to separate :  
1) kinematical regions (CFR/TFR)

# Beam SSAs as a tool to separate regions and contributions

Separating Target Fragmentation Region TFR from Current fragmentation region (CFR)



Negative sign of the SSA (plateau) defines the TFR dominance



SIDIS beam SSA can serve as a tool to separate:

- 1) kinematical regions (CFR/TFR)
- 2) dynamical contributions
- 3) cut on  $M_x$  eliminate exclusive VMs



# Understanding exclusive rhos and SDME validations

$$\mathcal{W}^U(\Phi, \phi, \cos \Theta)$$

$$+ \sqrt{2\epsilon(1+\epsilon)} \cos \Phi (r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta)$$

$$\mathcal{W}^L(\Phi, \phi, \cos \Theta)$$

$$+ \sqrt{2\epsilon(1-\epsilon)} \sin \Phi (r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta)$$

$$\gamma_T^* \rightarrow \rho_L^0 \quad \tau_{01} \approx \sqrt{\epsilon} \frac{\sqrt{(r_{00}^5)^2 + (r_{00}^8)^2}}{\sqrt{2r_{00}^{04}}}$$

Since the decay angle is correlated with the polarization of the rho, then  $r_{00}^8$  and  $r_{00}^5$  will be responsible for longitudinal rho, so tiny beam SSA expected for longitudinal rho

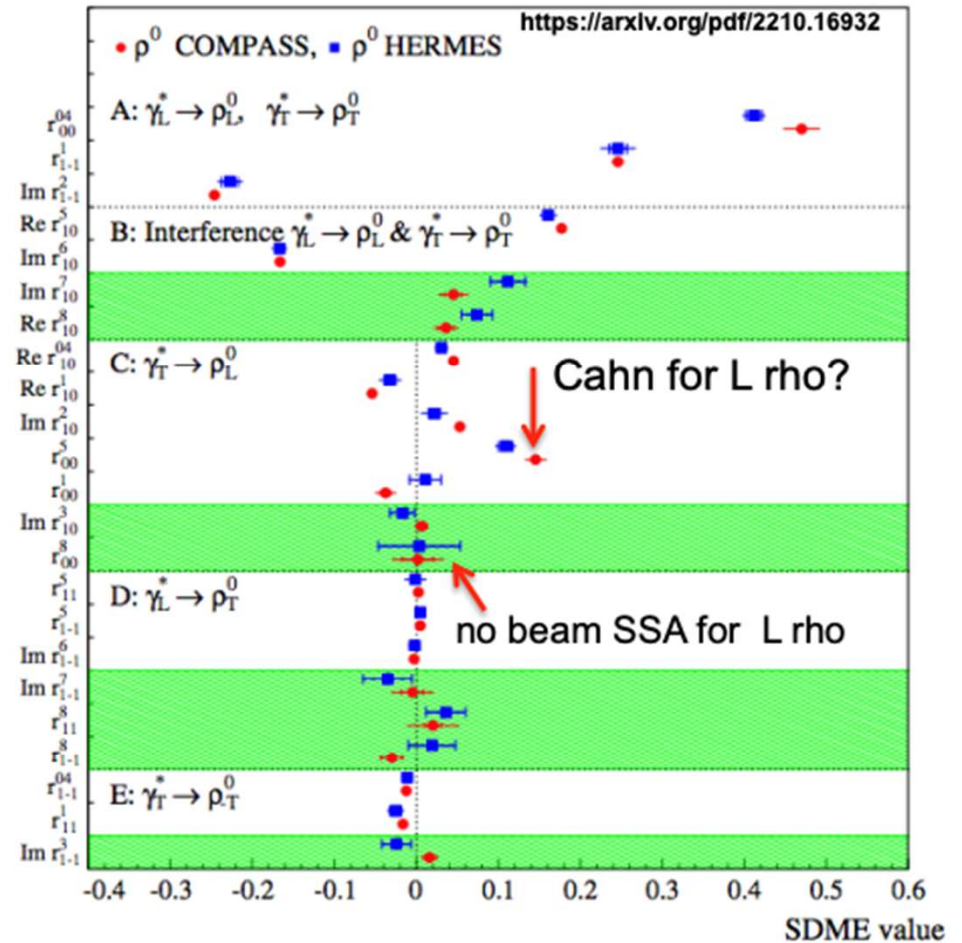


Fig. 12: Comparison of the 23 SDMEs for exclusive  $\rho^0$  lepton production on the proton extracted in the entire kinematic regions of the HERMES and COMPASS experiments. For HERMES the average kinematic values are  $\langle Q^2 \rangle = 1.96$  (GeV/c) $^2$ ,  $\langle W \rangle = 4.8$  GeV/c $^2$ ,  $\langle |t'| \rangle = 0.13$ , while those for COMPASS are  $\langle Q^2 \rangle = 2.40$  (GeV/c) $^2$ ,  $\langle W \rangle = 9.9$  GeV/c $^2$ ,  $\langle p_T^2 \rangle = 0.18$  (GeV/c) $^2$ . Inner error bars represent statistical uncertainties and outer ones statistical and systematic uncertainties added in quadrature. Unpolarised (polarised) SDMEs are displayed in unshaded (shaded) areas.

# SIDIS at JLab: the theorists' comments

---

## Statement:

“... SIDIS data has shown that there are basic open questions concerning the semi-inclusive pion/kaon production mechanisms at few-GeV energies, regarding e.g vector mesons and longitudinal photons....”

## Meaning:

JLab has problems specific for low energies, which should be solved, before THE theory of TMDs could be applied

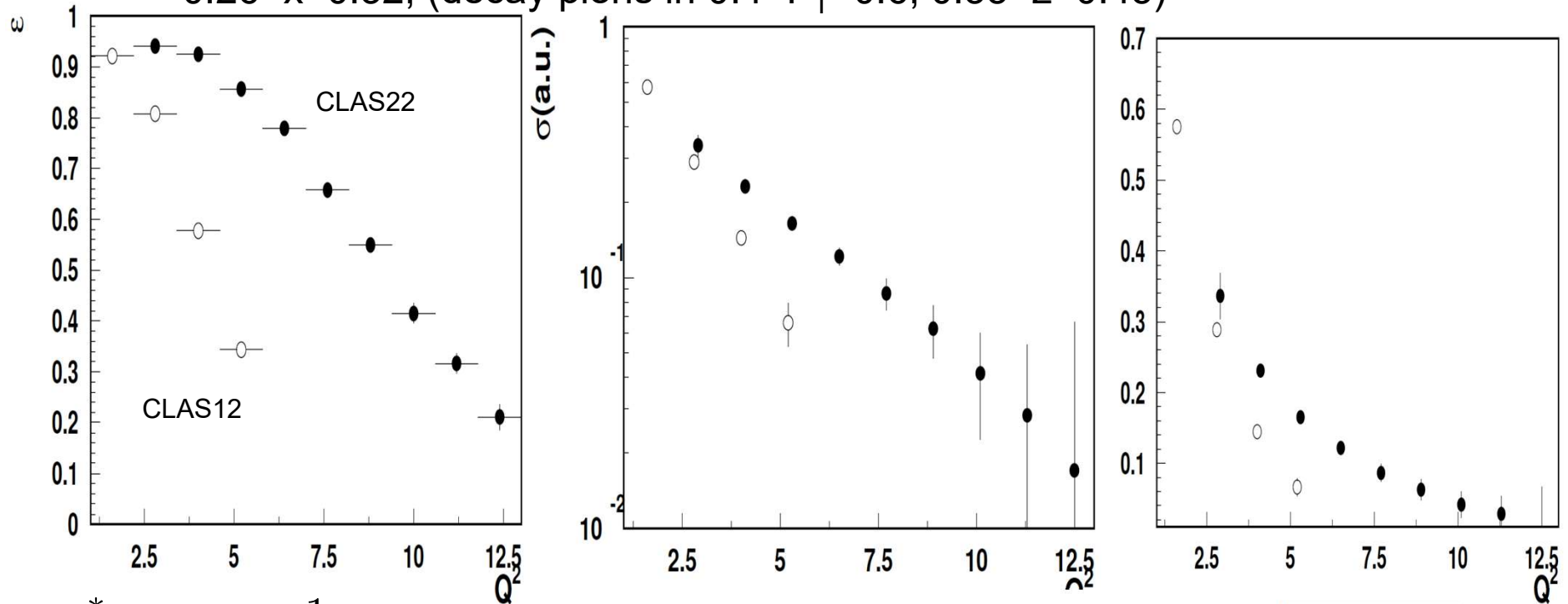
## Possible conclusion:

All problems are due to “few-GeV”, will magically vanish at higher energies, and TMDs can be studied in the valence region [in multidimensional space] at higher  $Q^2$  using THE theory [no need to deal with higher twists/correlations of quarks/hadrons/.....]



# Exclusive $\rho^0$ : extending the $Q^2$ with JLab22

0.26 < x < 0.32, (decay pions in 0.4 < P<sub>T</sub> < 0.6, 0.35 < z < 0.45)



$$\gamma_L^* \rightarrow \rho_L \propto 1$$

$$\gamma_T^* \rightarrow \rho_L \propto \sqrt{-t}/Q$$

Longitudinal photon contributions  
enhanced at higher energies

$$\gamma_T^* \rightarrow \rho_L^0 \quad \tau_{01} \approx \sqrt{\epsilon} \frac{\sqrt{(r_{00}^5)^2 + (r_{00}^8)^2}}{\sqrt{2r_{00}^4}}$$

$$W^L(\Phi, \phi, \cos\Theta)$$

$$+ \sqrt{2\epsilon(1-\epsilon)} \sin\Phi \left( r_{11}^8 \sin\Theta + r_{00}^8 \cos^2\Theta \right)$$

Range in  $Q^2$  increases significantly allowing  
detailed studies at beyond 10 GeV<sup>2</sup>

# Possible sources of large $P_T$ behaviour

- 1) Perturbative contributions and  $p_T$ -dependence of unpolarized FFs (so far unlikely...)
  - 2) Significantly wider in  $k_T$  distributions of u-quarks with spin opposite to proton spin (possible sign flips in asymmetries related to polarization of partons)
  - 3) Significantly wider in  $k_T$  distributions of d-quarks (possible sign flips in asymmetries related to polarization of partons)
  - 4) Significantly wider in  $k_T$  sea quark distributions (study contributions dominated by sea, K-,...)
- 
- 5) Increasing fraction of hadrons due to  $F_{UU,L}$  (needed for proper interpretation → separation of  $F_{UU,L}$  from total)
  - 6) Significant contributions from VMs to low  $P_T$  pion multiplicities, with direct pions showing up at large  $P_T$  (needed for proper interpretation → much wider in  $k_T$  original parton distributions)
  - 7) Radiative corrections (need the full x-section, typically applied to pions, while may be needed for underlying VMs,...)
  - 8) Two photon exchange (will need positron beam)

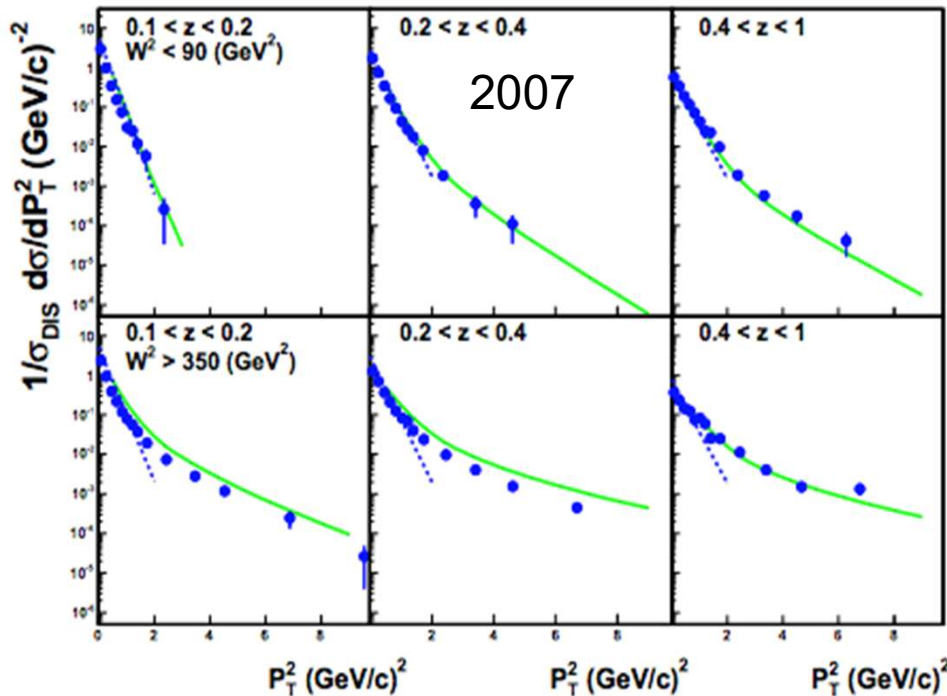
.....

# Multiplicities of hadrons in SIDIS

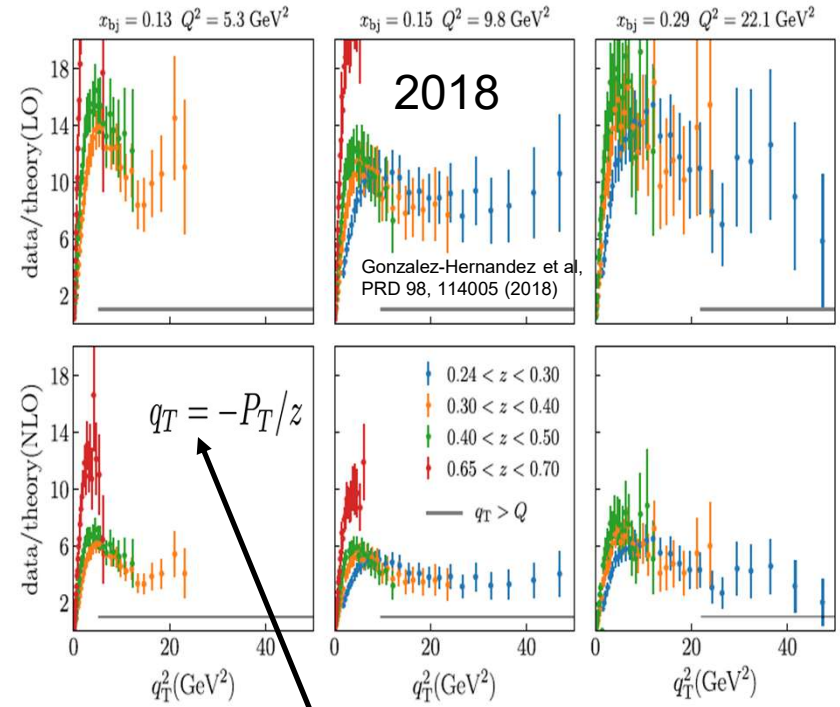
Anselmino et. al: hep-ph/0606286

$$F_{XY}^h(x, z, P_T, Q^2) \propto \underbrace{\sum H^q \times f^q(x, k_T, \dots)}_{\text{non-perturbative}} \otimes \underbrace{D^{q \rightarrow h}(z, p_T, \dots)}_{\text{perturbative}} + Y(Q^2, P_T) + \mathcal{O}(M/Q)$$

$$\int d^2 \vec{k}_T d^2 \vec{p}_T \delta^{(2)}(z \vec{k}_T + \vec{p}_T - \vec{P}_T)$$



Early attempts to use collinear and pQCD contributions (solid line) to fit EMC data

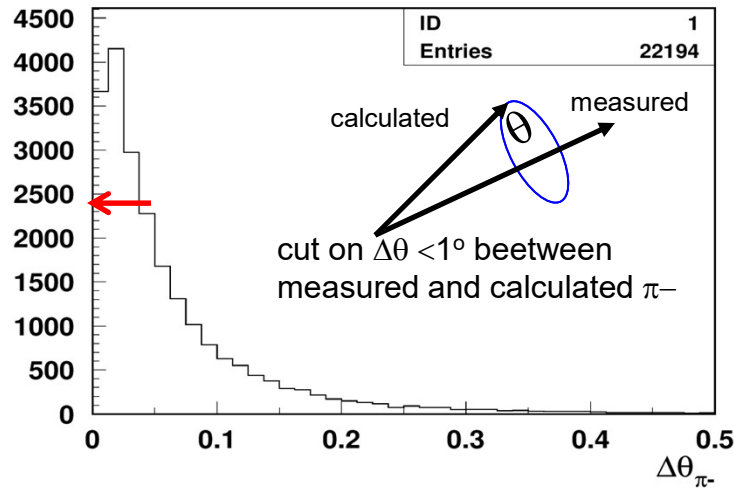


quark transverse momentum, assuming a direct link with pion transverse momentum

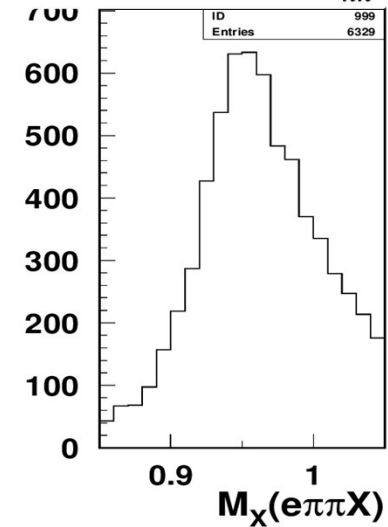
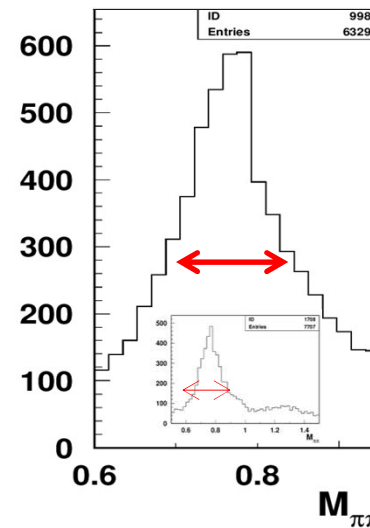
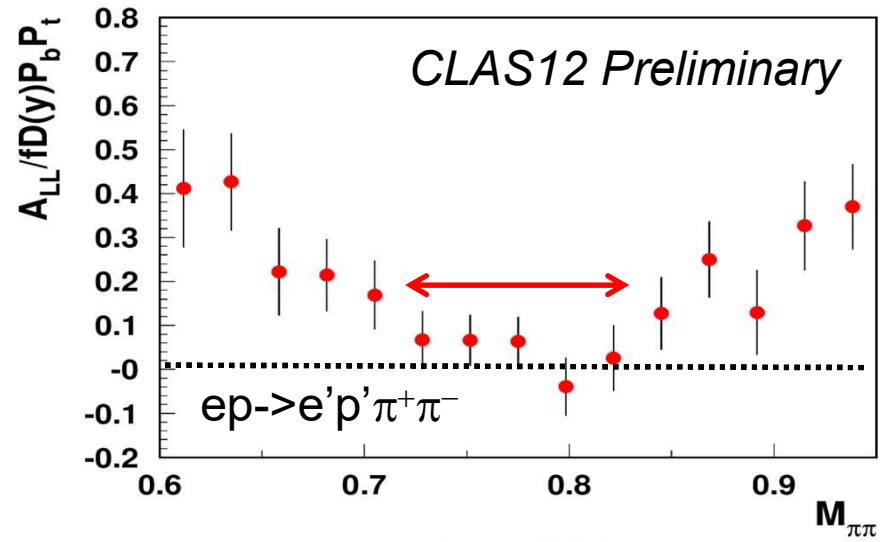
**Perturbative contributions underestimate the multiplicities by an order of magnitude for all accessible kinematics at COMPASS**

# Studies of $\rho^0$ impact with longitudinally polarized $\text{NH}_3$ target

## Separating exclusive dihadrons



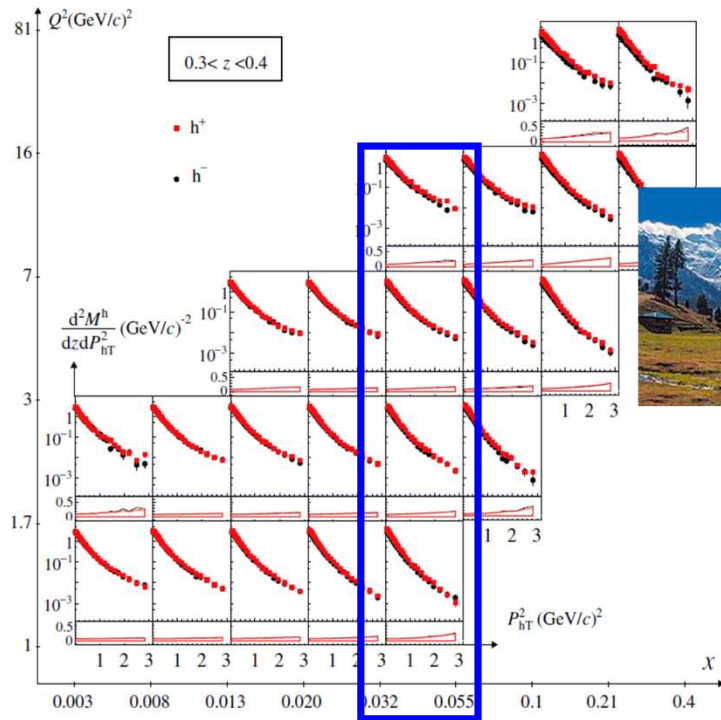
- Require the angle of negative pions is within a degree from calculated from  $e', p, \pi^+$  assuming exclusive  $e', p, \pi^+ \pi^-$  event.
- Measurements of  $A_{LL}$  for  $\rho^0$  indicate very small values (with  $\sim 10\text{-}20\%$  bckg, likely negative  $\sim -2\text{-}10\%$ ), and can be one of the reasons for higher  $A_{LL}$  with protons with a  $M_X$  cuts above 1.35 GeV (excluding exclusive  $\rho^0$ )



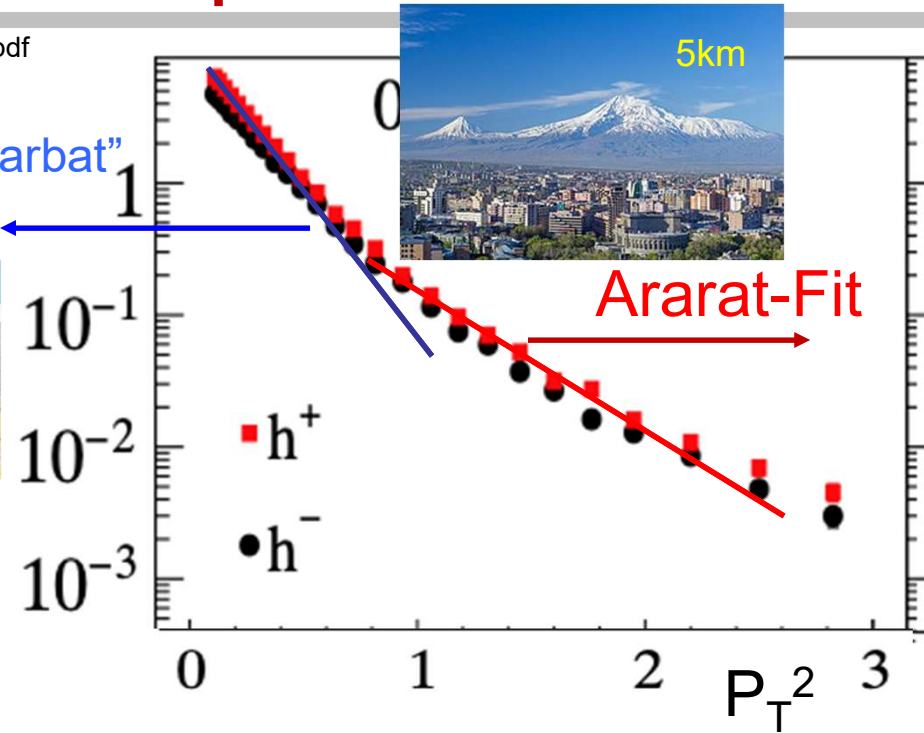
Need clear separation of hydrogen from  $\text{NH}_3$  and diffractive exclusive  $\rho^0$ s from exclusive  $\pi^+ \pi^-$

# q<sub>T</sub>-crisis or misinterpretation

<https://arxiv.org/pdf/1709.07374.pdf>



“Nanga Parbat”  
Fit



at higher  $Q^2$  the slope in  $P_T$  changes, why?

Higher the  $Q^2$  lower the  $\varepsilon$

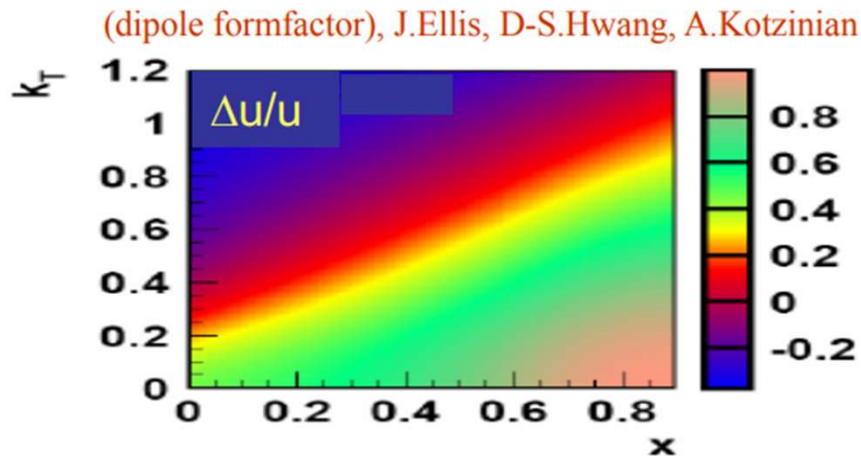
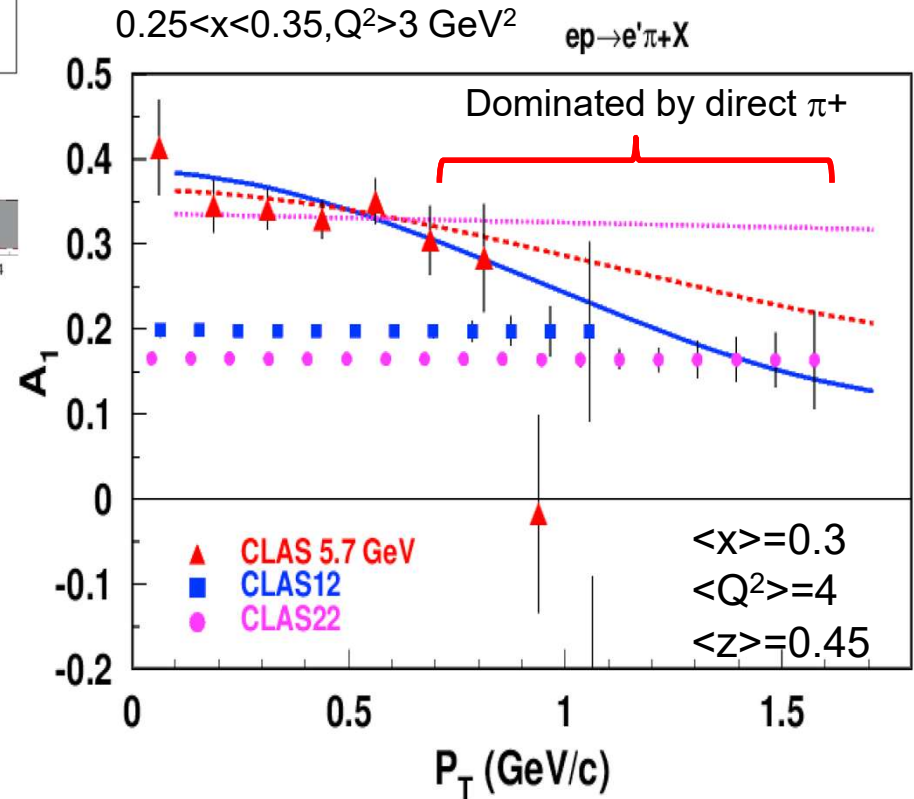
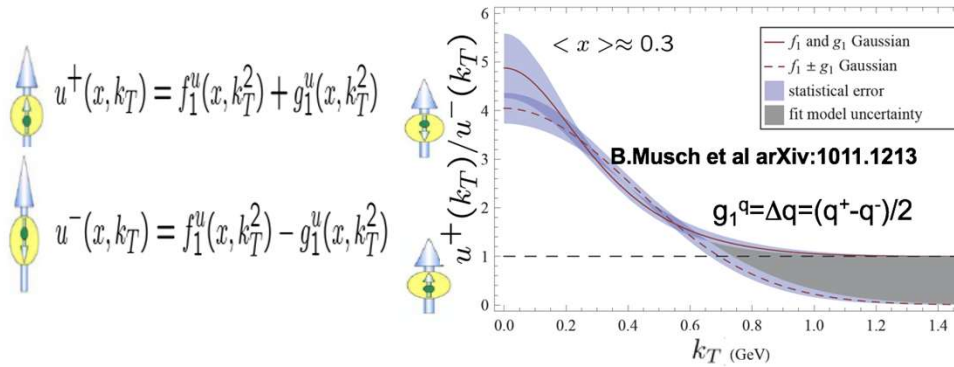
→ less diffractive rho at higher  $Q^2$  filling the low  $P_T$  in pion SIDIS.

**New procedure:** Fit from  $P_{Tmin}$  up  
 $P_{Tmin}$  can be lower at higher  $Q^2$ ,  
 as the contributions from diffractive rho decreases with  $Q^2$

Challenging for theory to explain the correlation of  $P_T$  and  $Q$   
 need experimental subtraction of rhos (proton detection will help)

# Unknown “known” $f_1, g_1$ TMDs

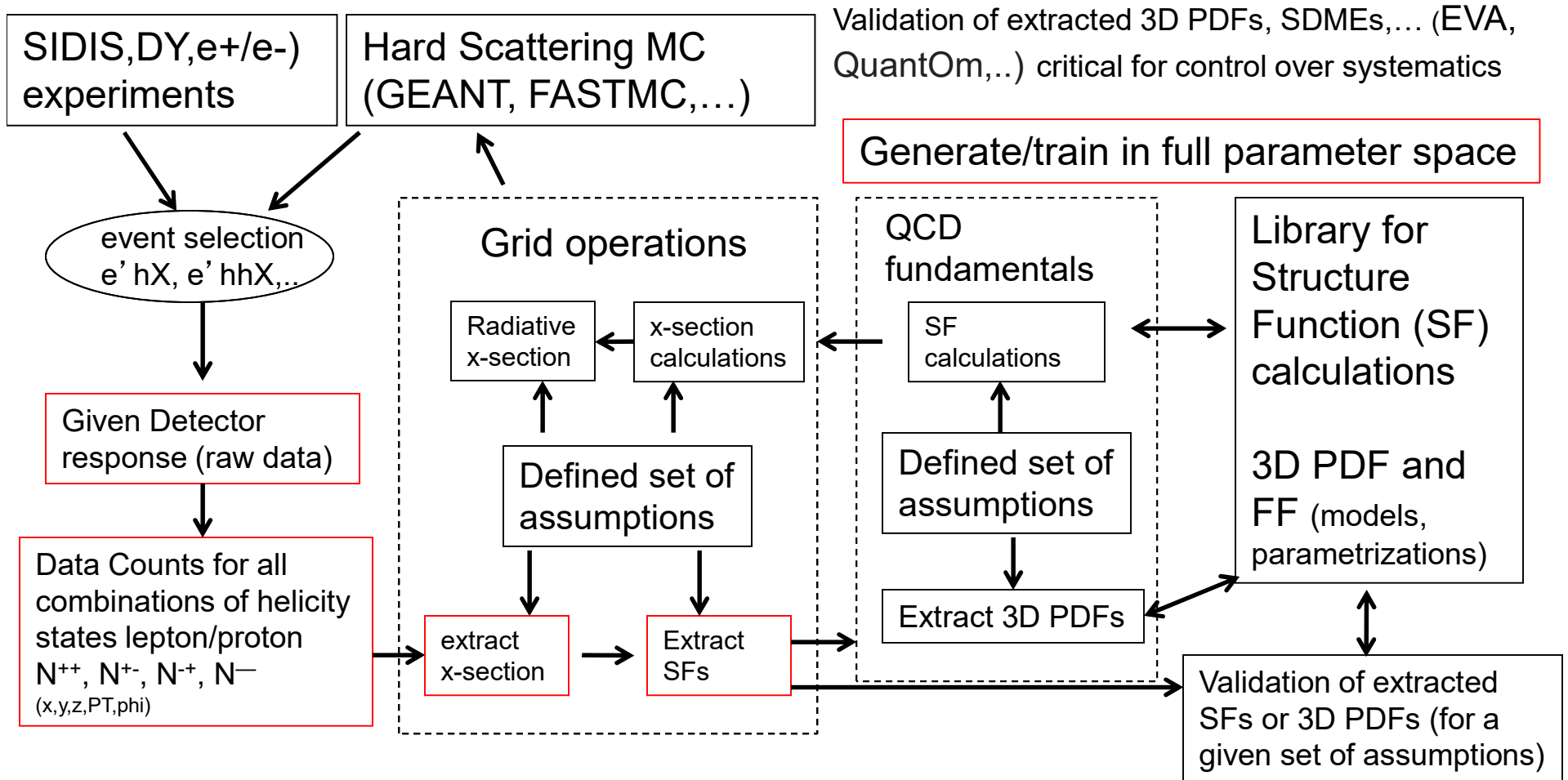
	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_{1L}$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}$	$h_1, h_{1T}^\perp$



- Models and lattice predict very significant spin and flavor dependence for TMDs
- Large transverse momenta are crucial to access the large  $k_T$  of quarks
- Several CLAS12 proposals dedicated to  $g_1(x, k_T)$ -studies CLAS12
- Understanding of  $k_T$ -dependence of  $g_1$  will help in modeling of  $f_1$



# 3D PDF Extraction and Validation (EVA) framework



Validation of extracted 3D PDFs, SDMEs, ... (EVA, QuantOm, ...) critical for control over systematics

Generate/train in full parameter space

Direct extraction of a given parameter sets from all steps (marked red) using AI tools techniques for the extraction of 3D PDFs and fragmentation functions from the **multidimensional** experimental observables with controlled systematics requires close collaboration of experiment, theory and computing

# x-section

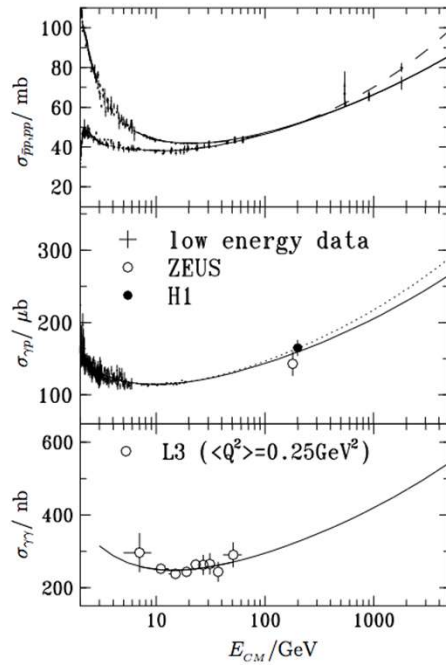


Figure 1.9: Total cross sections for pp ( $p\bar{p}$ ),  $\gamma p$  and  $\gamma\gamma$  scattering as a function of the center of mass energy  $E_{CM}$ . The curves represent the DL parameterization with  $\alpha_{IP}(0) = 1.0808$  (solid),  $= 1.112$  (dashed) and  $= 1.088$  (dotted).

Total hadron-hadron scattering can conveniently be described by the sum of a Reggeon and a Pomeron contribution. Donnachie and Landshoff [36] fitted all available hadronic data to the parameterization

$$\sigma_{tot} = A s^{\alpha_{RR}(0)-1} + B s^{\alpha_{IP}(0)-1}. \quad (1.38)$$

The parameters  $A$  and  $B$  depend on the particular process while global values for  $\alpha_{RR}(0) \approx 0.55$  and  $\alpha_{IP}(0) \approx 1.08$  are able to fit all considered data. A recent fit including newer data yielded  $\alpha_{IP}(0) \approx 1.096$  [37].

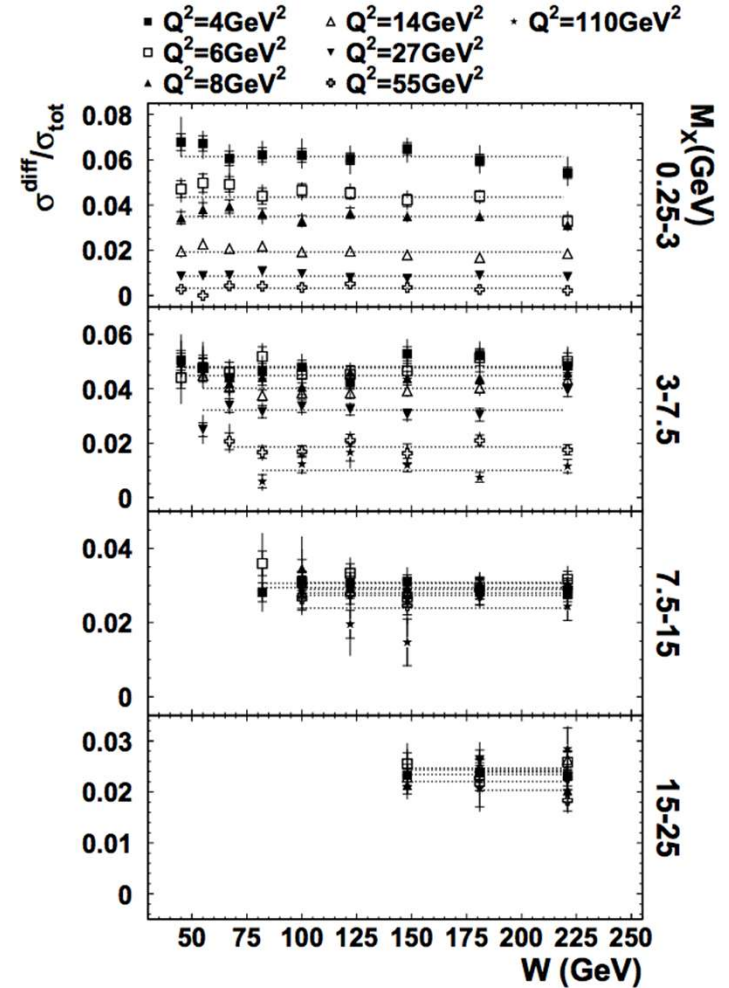
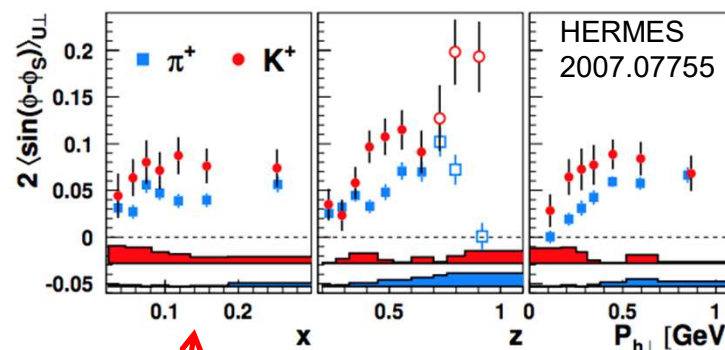
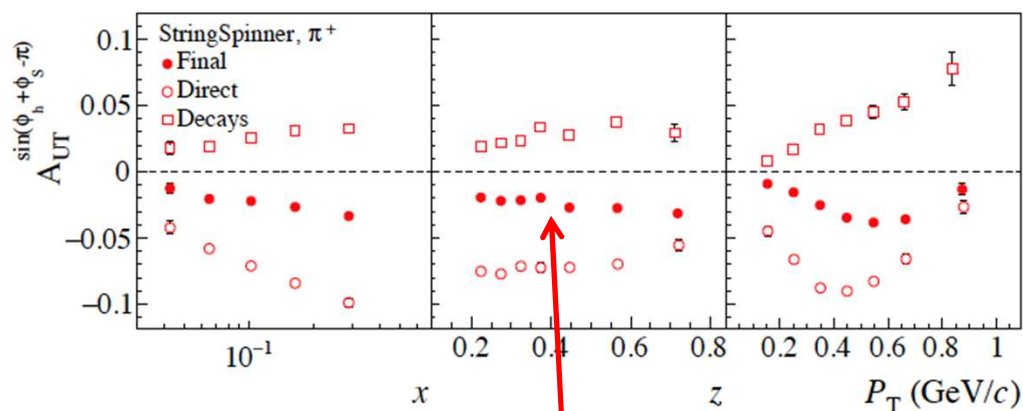
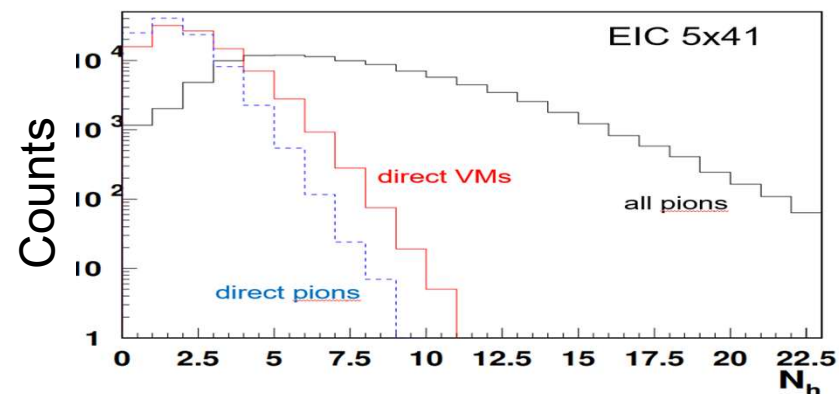
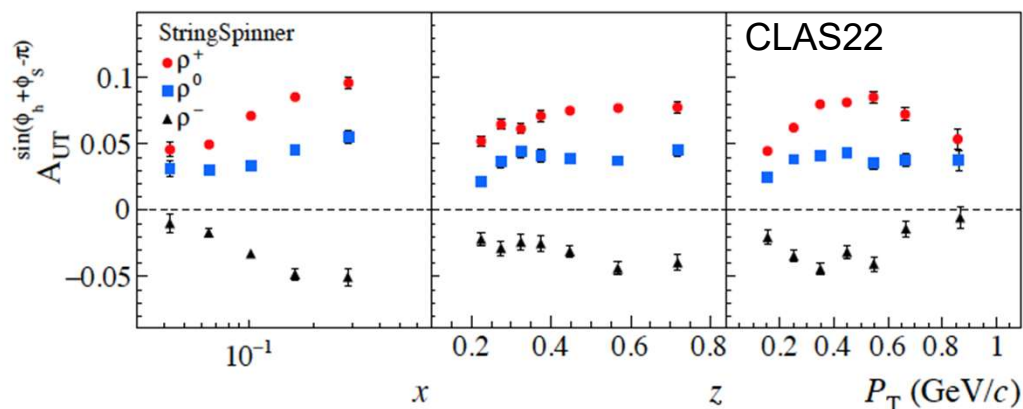


Figure 11.6: The ratio of the diffractive cross section  $\sigma^{diff}$ , integrated over the bin width  $M_a < M_X < M_b$ , and the total  $\gamma^*p$  cross section  $\sigma^{tot}$  is shown as a function of  $W$  for different bins of  $M_X$  and  $Q^2$ . The dotted lines indicate the average values of  $\sigma^{diff}/\sigma^{tot}$  in the measured  $W$  region for each bin in  $Q^2$  and  $M_X$ .

# VM contributions

A. Kerbizi (Trieste U.)



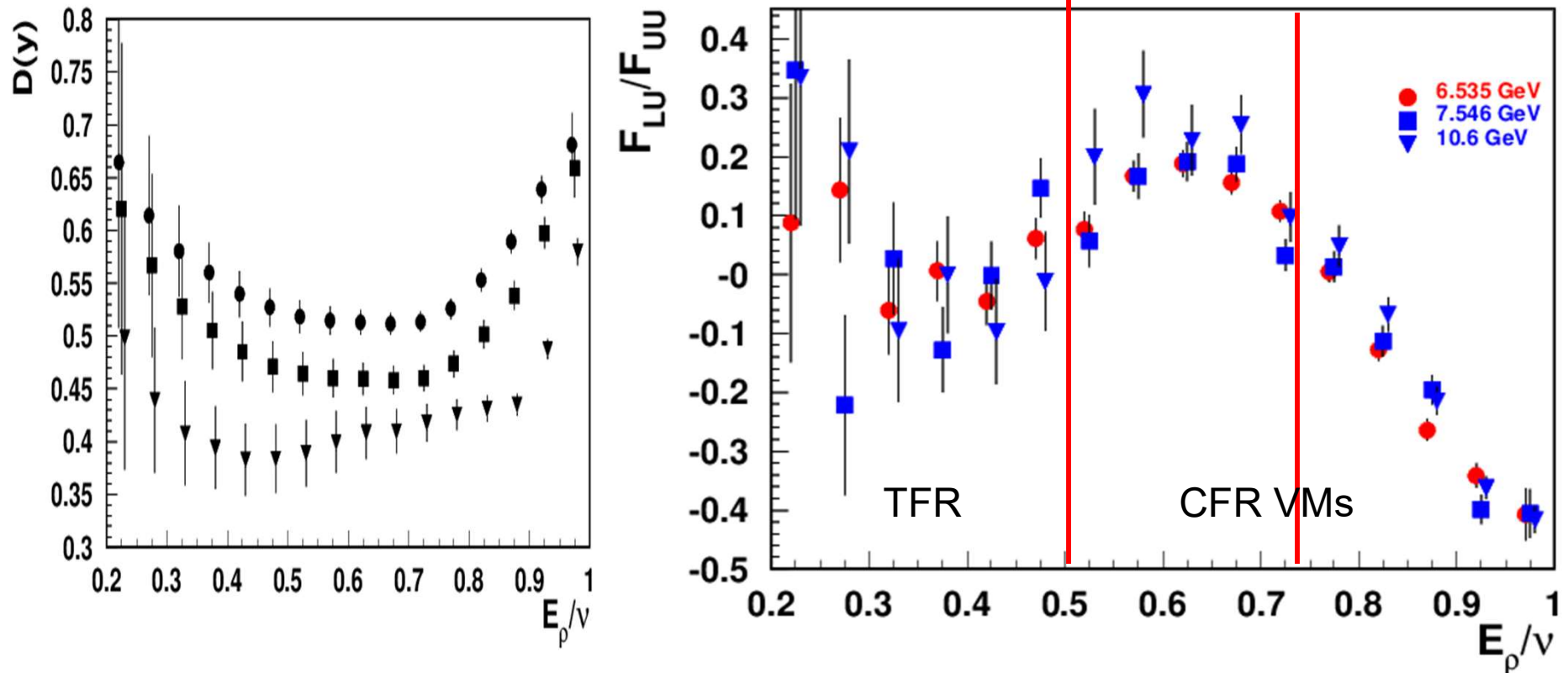
Strong dilution of SSAs(3-5 for pions 2-3 for Kaons) due to VM decays

The differences in pions vs Kaons SSAs may be coming from VMs???

Understanding VMs is critical for interpretation  
What is  $q_T$ ?

JLab can measure the SSA of VMs(also  $K^*$ ), and separate contributions

# VM contributions: z-dependence



SSA significant, and changing sign in 2 points separating 3 distinct regions, which have completely different impact on  $\pi^+$  SSAs

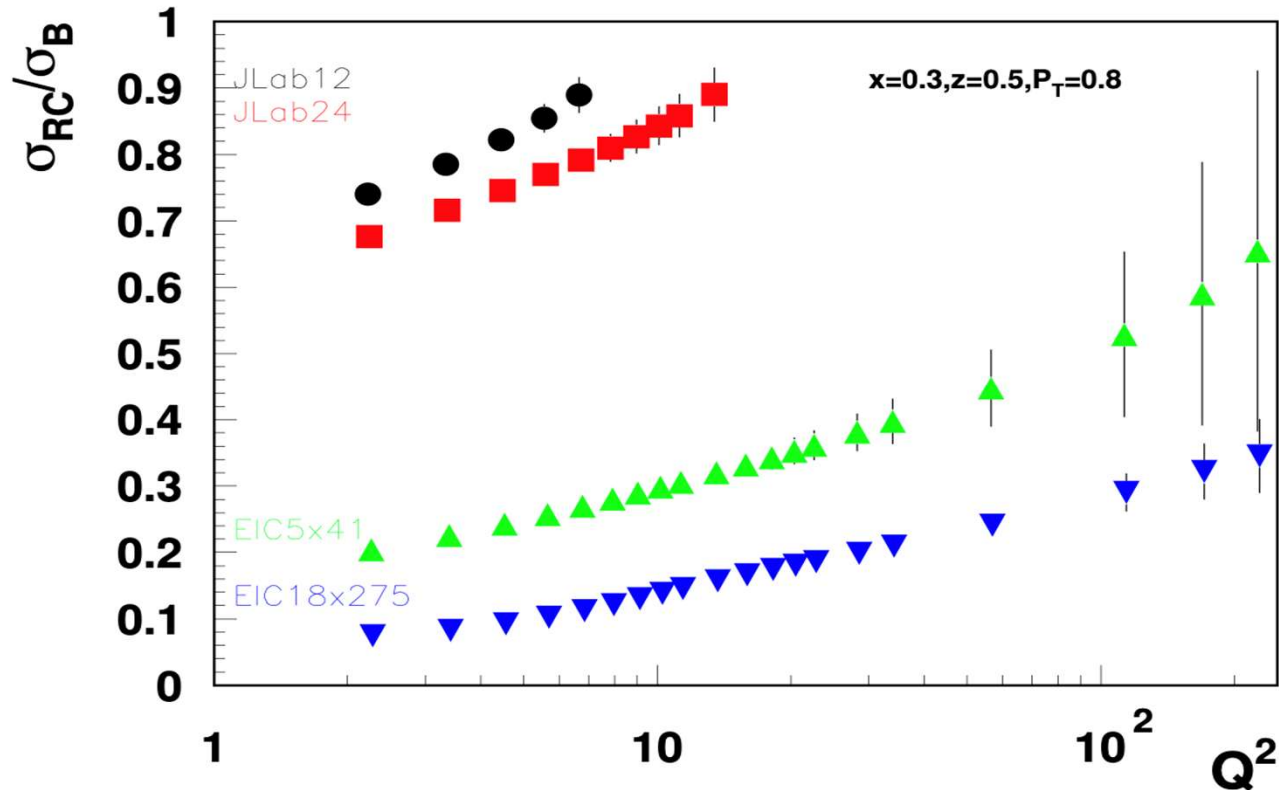
# From JLab to EIC: complementarity

The ratio of radiative cross ( $\sigma_{RC}$ ) section to Born ( $\sigma_B$ ) in SIDIS

*T. Liu et al*

*JHEP 11 (2021) 157*

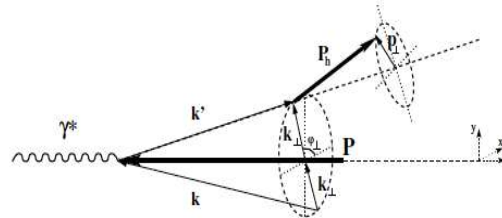
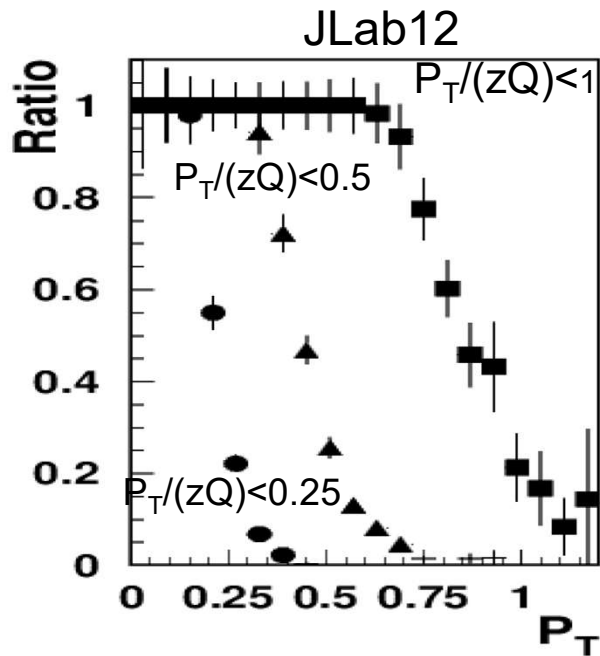
Gaussian  $F_{UU}(\phi_h=0)$



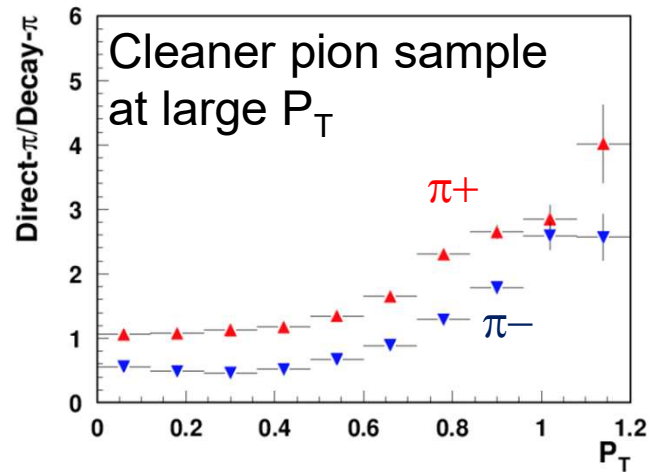
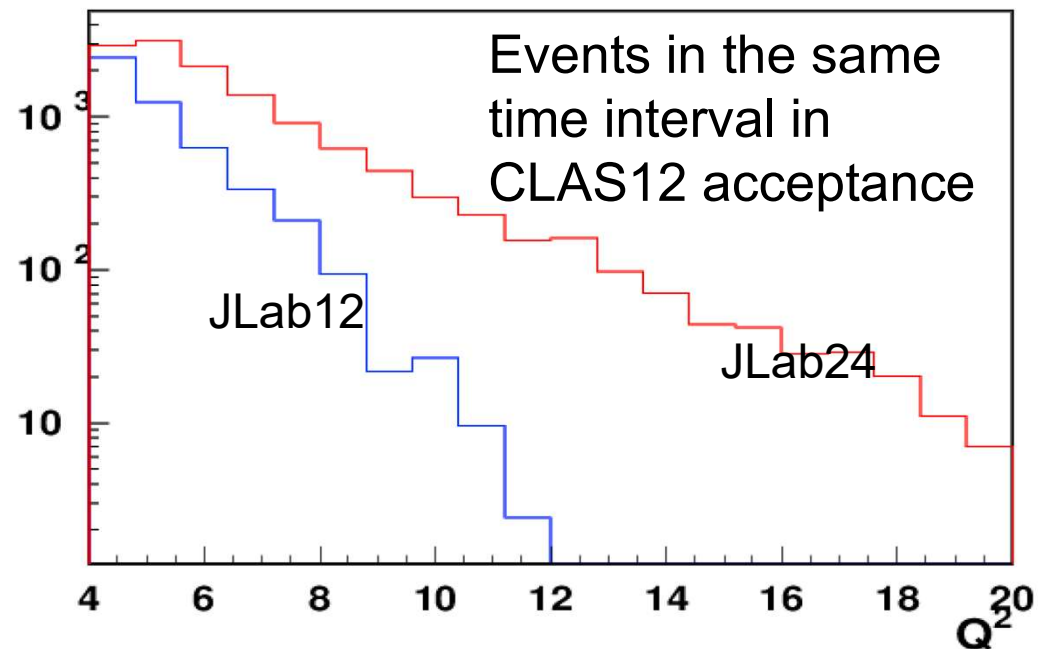
Cross section at low  $Q^2$   
suppressed at higher  
CM energies

- The radiative effects in SIDIS may be very significant and measurements in multidimensional space at different facilities will be crucial for understanding the systematics in evolution studies.
- Most sensitive to RC will be all kind of azimuthal modulations sensitive to cosines

# From JLa12 to JLab24 Larger $Q^2$ at large $P_T$



$$f_1^q(x, k_T) \otimes D_1^{q \rightarrow h}(z, p_T)$$



JLab24 will significantly increase the  $Q^2$  range, allowing detailed separation of higher twist SFs, needed for understanding the QCD