# Probe QCD Confinement via $\pi^0\text{, }\eta$ and $\eta^\prime$

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#### Outline:

- QCD at low energies
- Theoretical methods
- Neutral sector
- $\bullet \ \pi^0 \to \gamma \gamma$
- Summary

### QCD at low energies



From the 22 GeV white paper: 2306.09360

on problem of masses see previous talk by Khépani Raya

#### Amplitudes

- important in particle physics: Lagrangian  $\rightarrow$  Feynman rules  $\rightarrow$  amplitudes  $\rightarrow$  cross-section
- new initiative to study these objects more deeply
- annual conferences: ..., Prague 22, CERN 23, IAS 24, Seoul 25
- amplitudes are the key object of theoretical studies
- example  $\rightarrow$  next page

### QCD: gluon amplitudes

- important in high-energy collider experiments (LHC)
- using conventional methods: complicated already at the tree-level



- intermediate steps are complicated, but the final result "nice"
- standard methods hard/impossible for higher multiplicities
- surprisingly some results super simple and closed for all multiplicities

$$A_n(--+\ldots+) = \frac{\langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \ldots \langle n1 \rangle}$$

(so called MHV, [Parke, Taylor '86])

#### pion amplitudes

[KK, Novotny, Trnka '13]

- We want to study low-energy QCD
- focus on dynamics of pions, kaons, etas
- very complicated already at the tree-level for large n
- simplify the problem: massless, large  $N_c$  (one trace  $\rightarrow$  cyclic ordering)
- 4pt:  $A = s_{13}$
- 6pt:



[Arkani-Hamed et al '23-'24]

The simplest model:  $Tr(\phi^3)$  only one vertex:

e.g. the 4pt amplitude:



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 $X_{ij} = (p_i + \ldots + p_{j-1})^2$ 6/22

[Arkani-Hamed et al '23-'24]

The magic:

$$A = \frac{1}{X_{13}} + \frac{1}{X_{24}}$$

odd/even shifts:

$$X_{ee} \to X_{ee} + \delta, \qquad X_{oo} \to X_{oo} - \delta$$
  
 $X_{eo} \to X_{eo}$ 

Do it in  $Tr(\phi^3)$  amplitude and expand in small momenta for large  $\delta$ :

$$A \to \frac{1}{X_{13} - \delta} + \frac{1}{X_{24} + \delta} \sim -X_{13} - X_{24} = s_{13}$$

which is the 4pt NLSM!

[Arkani-Hamed et al '23-'24]

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which is the 4pt NLSM!

#### True up to all multiplicities!!

and can be extended to the loop level

#### Theoretical methods

- amplitudes methods are important to uncover hidden structures
- especially true for the low-energy QCD
- we hope that the above miracles have some footprint in the low-energy data
- BSM? double copy of YM  $\rightarrow$  gravity double copy of NLSM  $\rightarrow$  special galileon

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#### problem: masses

- ongoing investigation (in collaboration with N. Arkani-Hamed)
- possible to extend it to SU(3) limit  $(m_u = m_d = m_s \neq 0)$
- very hard beyond this limit, for the physical masses

the most complicated is the neutral sector, i.e.  $\pi^0,\eta,\eta'$ 

it is necessary to study and understand all details of  $\pi^0,\eta$  and  $\eta'$  important properties/channels:

- $\eta, \eta' \to 3\pi$
- their lifetimes

### Decays of $\pi^0$

neutral pion: the lightest hadron rich experimental activity: Hades, KLOE-2, A2), JLab, NA62, ... accessible also at lattice

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$$\pi^0 \to \gamma\gamma$$
 (98.8%)  $\leftarrow$  this talk  
•  $\pi^0 \to e^+ e^- \gamma$  (1.17%)  
•  $\pi^0 \to$  positronium  $\gamma$  (2×10<sup>-9</sup>  
•  $\pi^0 \to e^+ e^- e^+ e^-$  (3×10<sup>-5</sup>)  
•  $\pi^0 \to e^+ e^-$  (6×10<sup>-8</sup>)

[Bernstein, Holstein: Neutral Pion Lifetime Measurements and the QCD Chiral Anomaly]

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[Bernstein, Holstein: Neutral Pion Lifetime Measurements and the QCD Chiral Anomaly] other possible decay modes

• 
$$\pi^0 \rightarrow 4\gamma$$
 (theory vs. exp:  $2 \times 10^{-11}$  vs  $< 2 \times 10^{-8}$ )  
•  $\pi^0 \rightarrow \nu \bar{\nu}$   
helicity supression – limit:  $< 4 \times 10^{-9}$  (NA62 '21, 60x improvement), theory:  
 $< 5 \times 10^{-10}$  (direct) or  $< 10^{-24}$  (cosmology)  
•  $\pi^0 \rightarrow \nu \bar{\nu} \gamma$  (Wolfram mode)

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SM:  $\sim 10^{-18}$ , exp.limit  $< 2 \times 10^{-7}$  (NA62 '19)

## $\mathsf{EFT} \to \mathsf{ChPT}$

EFT

- separated degrees of freedom (simplification)
- building the most general Lagrangian
- ordering principle (powercounting)

example: ChPT

- goldstone bosons (spontaneous symmetry breakdown of chiral symmetry)
- momentum power counting
- even sector: Lagrangian up to NNNLO
- odd sector?  $\Rightarrow$

symmetry pattern of QCD must be studied more carefully

#### Odd sector

• first we need to add EM interaction:  $\partial_{\mu}U \rightarrow D_{\mu}U = \partial_{\mu}U + i[U, v_{\mu}], \qquad v_{\mu} \sim QA_{\mu}$ 

 $\bullet$  and add by hands monomial to  $\mathcal{L}:$ 

$$UF_{\mu\nu}\tilde{F}^{\mu\nu}$$

U can be transformed out: we have to add (at least two) derivatives on U – vanishes in chiral limit (Sutherland theorem) way out: anomaly, in fact two anomalies

$$\partial^{\mu}A^{i}_{\mu} \neq 0$$

(non-trivial for i = 0, 3, 8, or for  $\pi^0$ ,  $\eta$ ,  $\eta'$  states)

incorporated to the action by Wess, Zumino and Witten (WZW) two-flavour case: [Kaiser'01]

 $\pi^0 \rightarrow \gamma \gamma$ : LO

- $\pi^0$  lightest hadron
  - $\Rightarrow$  primary decay mode  $\pi^0 o \gamma\gamma$
- in chiral limit exact due to QCD axial anomaly:

$$\Gamma(\pi^0 \to \gamma\gamma) = \frac{m_{\pi^0}^3}{64\pi} \left(\frac{\alpha N_C}{3\pi F_\pi}\right)^2 = 7.73 \, eV$$

 $\pi^0 \rightarrow \gamma \gamma$ : Correction to the current algebra prediction

- using [Pagels and Zepeda '72] sum rules in [Kitazawa '85]
- NLO corrections are hidden in  $F \to F_{\pi}$  and  $O(p^6)$  LECs [Donoghue, Holstein, Lin '85] [Bijnens, Bramon, Cornet '88]



• in 3-flavour case we can study  $\pi^0, \eta, \eta'$  mixing, resulting to [Goity, Bernstein, Holstein '02]:

 $\Gamma^{\mathsf{NLO}} = 8.1 \pm 0.08 \, \mathrm{eV}$ 

in 2-flavour case EM corrections [Ananth., Moussallam '02]:  $\Gamma^{\rm NLO}=8.06\pm0.02\pm0.06\,{\rm eV}$ 

• study based on dispersion relations [loffe, Oganesian '07]:  $\Gamma^{\rm NLO} = 7.93 \pm 0.11 \, {\rm eV}$ 

## $\pi^0 ightarrow \gamma \gamma$ at NNLO in 2 flavour ChPT [KK, Moussallam'09]

- NLO: a) One-loop diagrams with one vertex from L<sup>WZ</sup>, b) tree diagrams with one vertex from L<sup>WZ</sup> and one vertex from O(p<sup>4</sup>) Lagrangian, c) tree diagrams with one vertex from O(p<sup>6</sup>) anomalous-parity sector
- $O(p^6)$  anomalous-parity sector from [Bijnens, Girlanda, Talavera '02]
- representation of chiral field:  $U = \sigma + i \frac{\tau \cdot \pi}{F}$ ,  $\sigma = \sqrt{1 \vec{\pi}^2/F^2}$  (no  $\gamma 4\pi$  vertex at LO)



- verification of Z-factor,  $F_{\pi}/F$  [Bürgi '96], [Bijnens, Colangelo, Ecker, Gasser, Sainio '02]
- double log checked by Weinberg consistency relations

 $\pi^0 \to \gamma \gamma$  at NNLO, result

$$\begin{split} A_{NNLO} &= \frac{e^2}{F_{\pi}} \Big\{ \frac{1}{4\pi^2} \\ &+ \frac{16}{3} m_{\pi}^2 (-4c_3^{Wr} - 4c_7^{Wr} + c_{11}^{Wr}) + \frac{64}{9} B(m_d - m_u) (5c_3^{Wr} + c_7^{Wr} + 2c_8^{Wr}) \\ &- \frac{M^4}{24\pi^2 F^4} \left( \frac{1}{16\pi^2} L_{\pi} \right)^2 + \frac{M^4}{16\pi^2 F^4} L_{\pi} \Big[ \frac{3}{256\pi^4} + \frac{32F^2}{3} \left( 2c_2^{Wr} + 4c_3^{Wr} + 2c_6^{Wr} + 4c_7^{Wr} - c_{11}^{Wr} \right) \Big] \\ &+ \frac{32M^2 B(m_d - m_u)}{48\pi^2 F^4} L_{\pi} \Big[ -6c_2^{Wr} - 11c_3^{Wr} + 6c_4^{Wr} - 12c_5^{Vr} - c_7^{Wr} - 2c_8^{Wr} \Big] \\ &+ \frac{M^4}{F^4} \lambda_+ + \frac{M^2 B(m_d - m_u)}{F^4} \lambda_- + \frac{B^2(m_d - m_u)^2}{F^4} \lambda_- \Big\} \\ \lambda_+ &= \frac{1}{\pi^2} \Big[ -\frac{2}{3} d_{+}^{Wr}(\mu) - 8c_6^r - \frac{1}{4} (l_4^r)^2 + \frac{1}{512\pi^4} \left( -\frac{983}{288} - \frac{4}{3} \zeta(3) + 3\sqrt{3} \operatorname{Cl}_2(\pi/3) \right) \Big] \\ &+ \frac{16}{9} F^2 \Big[ 8l_3^r (c_3^{Wr} + c_7^{Wr} + 12c_4^{Wr} - 4c_7^{Wr} + c_{11}^{Wr}) \Big] \\ \lambda_- &= \frac{64}{9} \Big[ d_{-}^{Wr}(\mu) - 128F^2 l_7 (c_3^{Wr} + c_7^{Wr} + 2c_8^{Wr}) \Big] \\ \lambda_- &= d_{--}^{W-}(\mu) - 128F^2 l_7 (c_3^{Wr} + c_7^{Wr}) \,. \end{split}$$

- 4 LECs in 2 combinations of NLO
- additional 4 LECs in 3 combinations of NNLO

Is it at all possible to make some reliable prediction?

## $\pi^0 \rightarrow \gamma \gamma$ : modified counting

Use of SU(3) phenomenology via c<sup>Wr</sup><sub>i</sub> ↔ C<sup>Wr</sup><sub>i</sub> connection (based on [Gasser, Haefeli, Ivanov, Schmid '07,'08])

$$c_i^{Wr} = \frac{\alpha_i}{m_s} + \left(\beta_i + \gamma_{ij}C_j^{Wr} + \delta_i \ln \frac{B_0 m_s}{\mu^2}\right) + O(m_s)$$

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• implementation of modified counting

$$m_u, m_d \sim O(p^2)$$
 and  $m_s \sim O(p)$ 

#### Result:

$$\begin{split} A_{NNLO}^{mod} &= \frac{e^2}{F_{\pi}} \Biggl\{ \frac{1}{4\pi^2} - \frac{64}{3} m_{\pi}^2 C_7^{Wr} + \frac{1}{16\pi^2} \frac{m_d - m_u}{m_s} \Biggl[ 1 - \frac{3}{2} \frac{m_{\pi}^2}{16\pi^2 F_{\pi}^2} L_{\pi} \Biggr] \\ &+ 32B(m_d - m_u) \Biggl[ \frac{4}{3} C_7^{Wr} + 4C_8^{Wr} \Bigl( 1 - 3 \frac{m_{\pi}^2}{16\pi^2 F_{\pi}^2} L_{\pi} \Bigr) \\ &- \frac{1}{16\pi^2 F_{\pi}^2} \Bigl( 3L_7^r + L_8^r - \frac{1}{512\pi^2} (L_K + \frac{2}{3} L_{\eta}) \Bigr) \Biggr] - \frac{1}{24\pi^2} \left( \frac{m_{\pi}^2}{16\pi^2 F_{\pi}^2} L_{\pi} \Biggr)^2 \Biggr\} \end{split}$$

#### $\pi \rightarrow \gamma \gamma$ : Phenomenology inputs

- $F_{\pi} = 92.22 \pm 0.07$  MeV (using updated value of  $V_{ud}$  [Towner, Hardy'08]). using quark mass ratio (from lattice), pseudo-scalar meson masses, R from  $\eta \rightarrow 3\pi$
- $\frac{m_d m_u}{m_s} = (2.29 \pm 0.23) \, 10^{-2}$
- $B(m_d m_u) = (0.32 \pm 0.03) M_{\pi^0}^2$
- $3L_7 + L_8^r(\mu) = (0.10 \pm 0.06) \, 10^{-3}$   $(\mu = M_\eta)$  (from pseudo-scalar meson masses formula [Gasser, Leutwyler '85])
- $C_7^W = 0$  (more precisely  $C_7^W \ll C_8^W$ , motivated by simple resonance saturation)
- $C_8^W = (0.58 \pm 0.2) 10^{-3} {\rm GeV}^{-2}$  (from  $\eta \to 2\gamma$ )

result [KK, Moussallam'09]:

$$\Gamma_{\pi^0 \to 2\gamma} = (8.09 \pm 0.11) \text{ eV}$$

 $\pi^0 \to \gamma \gamma$ : comments

- 1)  $F_{\pi}$  is a crucial ingredient
  - $F_{\pi}$  vs  $\hat{F}_{\pi}$  [Bernard, Oertel, Passemar, Stern '08]
  - using  $\pi^0 \to \gamma \gamma$ :

$$F_{\pi}=93.85\pm1.4~{
m MeV}$$
 cf with  $\hat{F}_{\pi}=92.22(7)$   $(1.2\sigma~{
m difference})$ 

our F<sub>π</sub> from PDG is based on π<sub>l2</sub> and SM using [Marciano, Sirlin'93]
important input V<sub>ud</sub>: new update by [Hardy,Towner '20]

 $0.97418(26) \rightarrow 0.97373(31)$ 

we should change  $F_{\pi} \rightarrow 92.3$  ...

 $\pi^0 
ightarrow \gamma\gamma$ : comments

2)  $\eta/\eta'$  decays

 $\eta' \to \gamma \gamma$ :

• interconnected with  $\pi^0 \rightarrow \gamma \gamma$  (via constant  $C_7$ )

 $\eta \to \gamma \gamma$  :

- interconnected with  $\pi^0 \to \gamma \gamma$  ( $\to$  constant  $C_8$ )
- ongoing project to calculate it also at two loop (with J.Bijnens)
- $\bullet\,$  much complicated than  $\pi^0\to\gamma\gamma$



## $\pi^0 \to \gamma \gamma$ : comments

3) still interesting from purely theoretical point of view

a) Leading logarithm contribution of individual orders in percent of the leading order: [Bijnens,KK,Lanz'12]



b) surprising connections [Bijnens,KK,Lanz'12]

Adler-Lee-Treiman-Zee-Terentev theorem on triangle and box anomaly

$$F^{3\pi}(0,0,0) = \frac{1}{eF_{\pi}^2} F_{\pi\gamma\gamma}(0,0)$$

is valid up to two loops for leading logs beyond the soft-photon limit

## Summary

- neutral sector of low energy QCD: the important theoretical challenge of understanding EFT
- BSM: tantalizing theoretical connection with a different BSM sector!
- In this talk focused mainly on  $\pi^0 \to \gamma \gamma$  decay.
- summary: more theoretical understanding needed, and maybe more calculations needed

theory:  $(8.09 \pm 0.11) \text{ eV}$ PrimEx I+II:  $(7.80 \pm 0.12) \text{ eV}$ 

or equivalently  $\pi^0$  lifetime:

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theory: 8.04 \pm 0.11 \times 10^{-17} s
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Thank you for your attention!