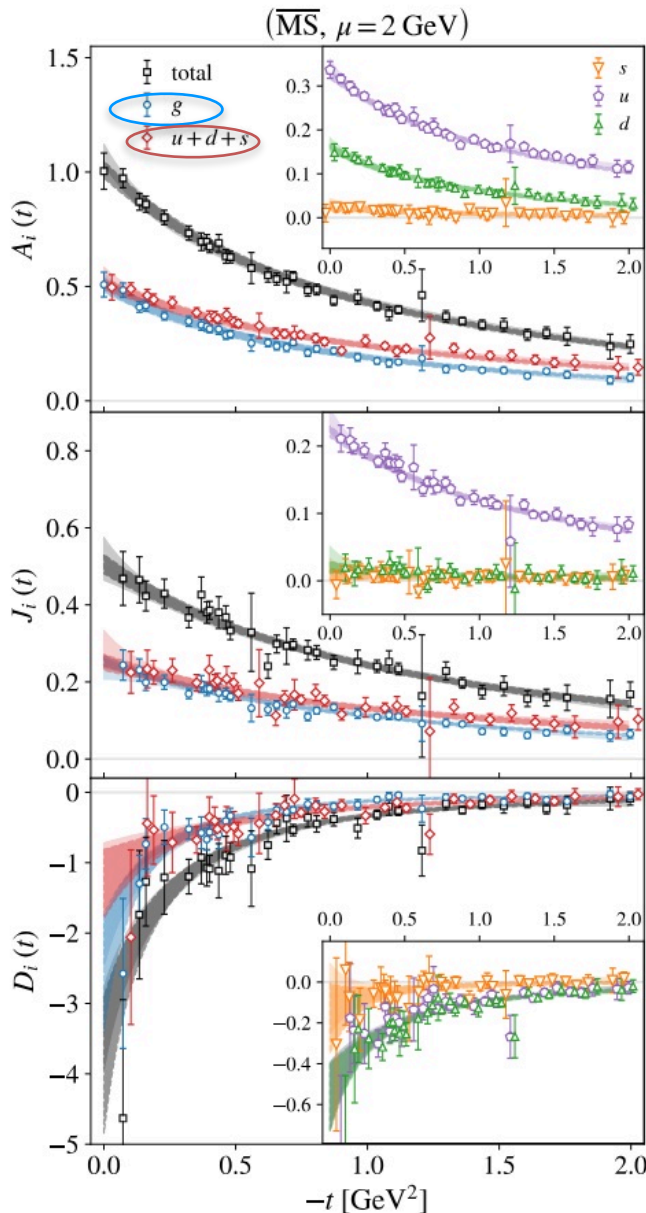


Threshold charmonium production at JLab22 - an access to gluon structure of the proton

Lubomir Pentchev

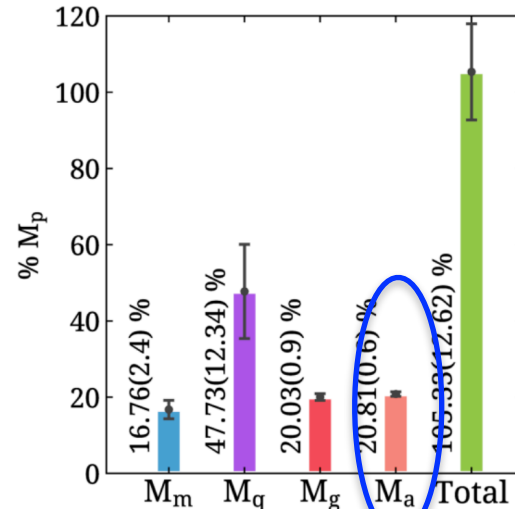


- Gluon contribution to the mechanical properties of the proton equally important as the quark one:

Lattice calculations of Gravitational Form Factors (GFFs) show similar contributions from **gluons (g)** and **quarks (u+d+s)**.

Hackett, Pefkou, Shanahan arxiv:2310.08484 (2023)

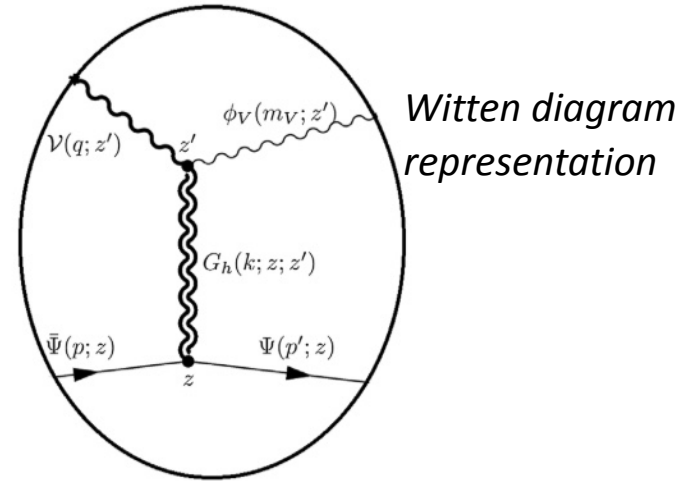
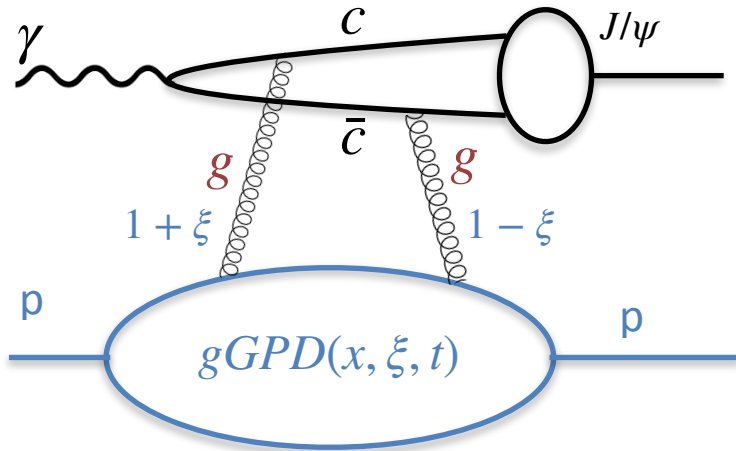
- Quark masses and kinetic energies of quarks and gluons are not enough to explain the mass of the proton: gluon condensate, or anomalous contribution to the mass of the proton is significant:



C. Alexandrou *et al.*, (ETMC), PRL 119, 142002 (2017)

C. Alexandrou *et al.*, (ETMC), PRL 116, 252001 (2016)

Threshold charmonium photoproduction - GPD and holographic approaches



$$\left(\frac{d\sigma}{dt}\right)_{\gamma p \rightarrow J/\psi p} = F(E_\gamma) \xi^{-4} [G_0(t) + \xi^2 G_2(t)] + \dots$$

$$\left(\frac{d\sigma}{dt}\right)_{\gamma p \rightarrow J/\psi p} = N(E_\gamma) [H_0(t) + \eta^2 H_2(t)] + \dots$$

$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t)\right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t)\right)^2$$

$$G_2(t) = 2\mathcal{A}_g^{(2)}(t)\mathcal{C}_g(t) + 2\frac{t}{4m^2}\mathcal{B}_g^{(2)}(t)\mathcal{C}_g(t) - \left(\mathcal{A}_g^{(2)}(t) + \mathcal{B}_g^{(2)}(t)\right)^2$$

$$\mathcal{A}_g^{(2)}(t) = A_g(t)$$

$$\mathcal{B}_g^{(2)}(t) = B_g(t)$$

$$\mathcal{C}_g(t) = 4C_g(t)$$

$$\xi = \frac{M_{J/\psi}^2 - t}{2(s - m^2) - M_{J/\psi}^2 + t}$$

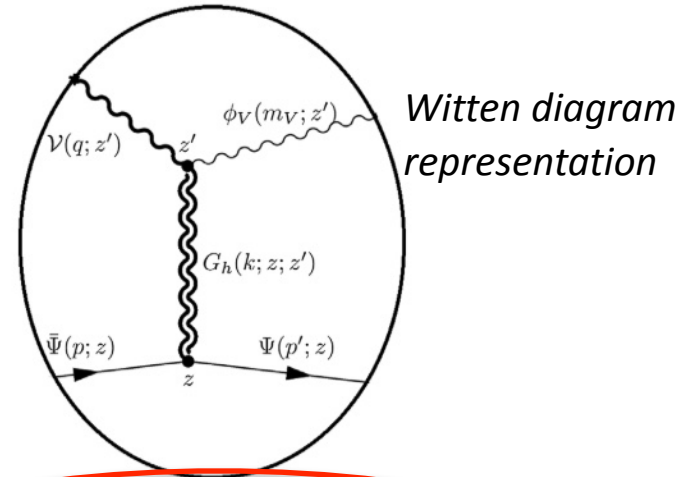
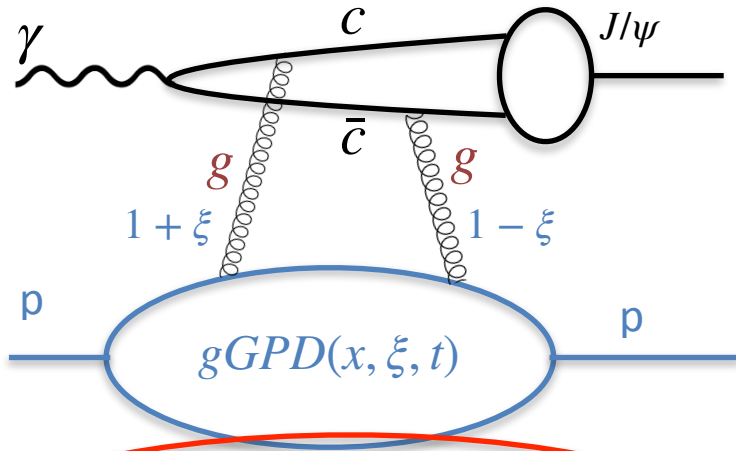
$$H_0(t) = A_g^2(t)$$

$$H_2(t) = 8A_g(t)C_g(t)$$

$$\eta = \frac{M_{J/\psi}^2}{2(s - m^2) - M_{J/\psi}^2 + t}$$

Holographic analysis by Mamo and Zahed PRD 106 (2022), PRD, PRD 101 (2020), Hatta and Yang PRD 98 (2018)

Threshold charmonium photoproduction - GPD and holographic approaches



$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = F(E_\gamma) \xi^{-4} [G_0(t) + \xi^2 G_2(t)] + \dots$$

$$\left(\frac{d\sigma}{dt} \right)_{\gamma p \rightarrow J/\psi p} = N(E_\gamma) [H_0(t) + \eta^2 H_2(t)] + \dots$$

$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t) \right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t) \right)^2$$

$$H_0(t) = A_g^2(t)$$

$$G_2(t) = 2\mathcal{A}_g^{(2)}(t)\mathcal{C}_g(t) + 2\frac{t}{4m^2}\mathcal{B}_g^{(2)}(t)\mathcal{C}_g(t) - \left(\mathcal{A}_g^{(2)}(t) + \mathcal{B}_g^{(2)}(t) \right)^2$$

$$H_2(t) = 8A_g(t)C_g(t)$$

$$\mathcal{A}_g^{(2)}(t) = A_g(t)$$

$$\mathcal{B}_g^{(2)}(t) = B_g(t)$$

$$\mathcal{C}_g(t) = 4C_g(t)$$

$$\xi = \frac{M_{J/\psi}^2 - t}{2(s - m^2) - M_{J/\psi}^2 + t}$$

$$\eta = \frac{M_{J/\psi}^2}{2(s - m^2) - M_{J/\psi}^2 + t}$$

leading moments for high ξ values

for large N_c and strong α_s

Holographic analysis by Mamo and Zahed PRD 106 (2022), PRD, PRD 101 (2020), Hatta and Yang PRD 98 (2018)

Form Factor Functions

$G_0(t)$ - GPD

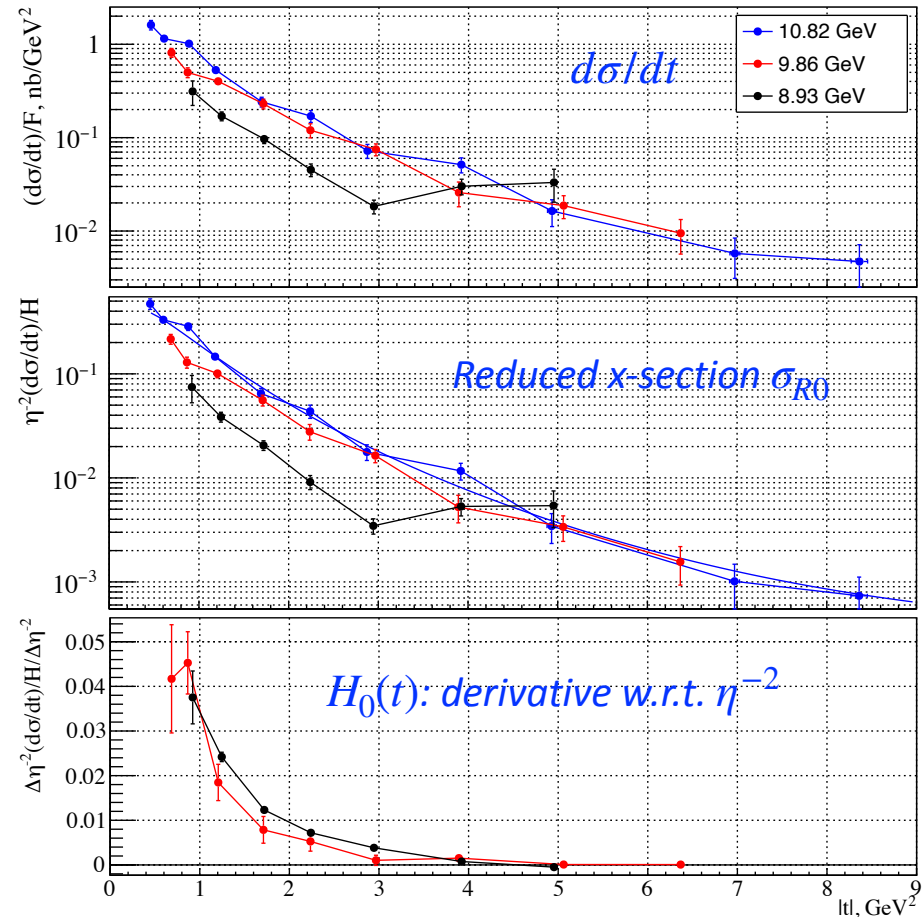
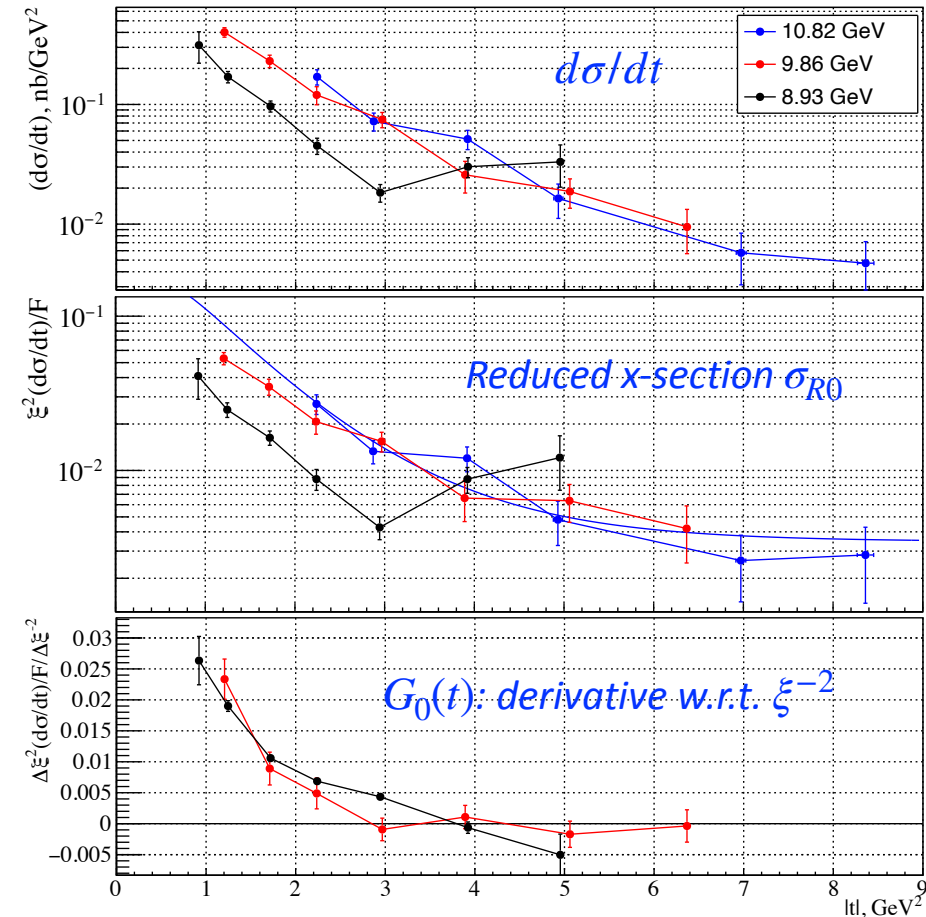
$$\sigma_{R0} = \frac{d\sigma}{dt} \frac{\xi^2}{F(E_\gamma)} = \xi^{-2} G_0(t) + G_2(t)$$

$$G_0(t) = \left[\sigma_{R0}(E_i, t) - \sigma_{R0}(E_j, t) \right] / \left[\xi^{-2}(E_i, t) - \xi^{-2}(E_j, t) \right]$$

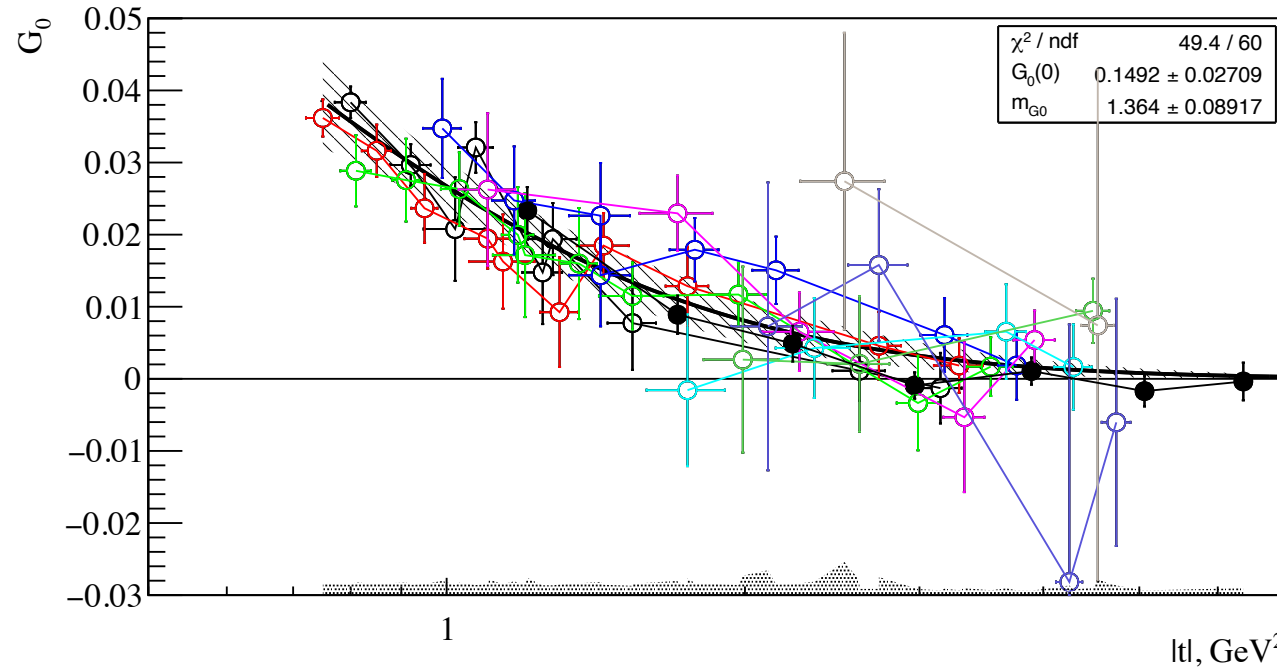
$H_0(t)$ - Holography

$$\sigma_{R0} = \frac{d\sigma}{dt} \frac{\eta^{-2}}{H(E_\gamma)} = \eta^{-2} H_0(t) + 4H_2(t)$$

$$H_0(t) = \left[\sigma_{R0}(E_i, t) - \sigma_{R0}(E_j, t) \right] / \left[\eta^{-2}(E_i, t) - \eta^{-2}(E_j, t) \right]$$



Gluon Form Factors - energy independence

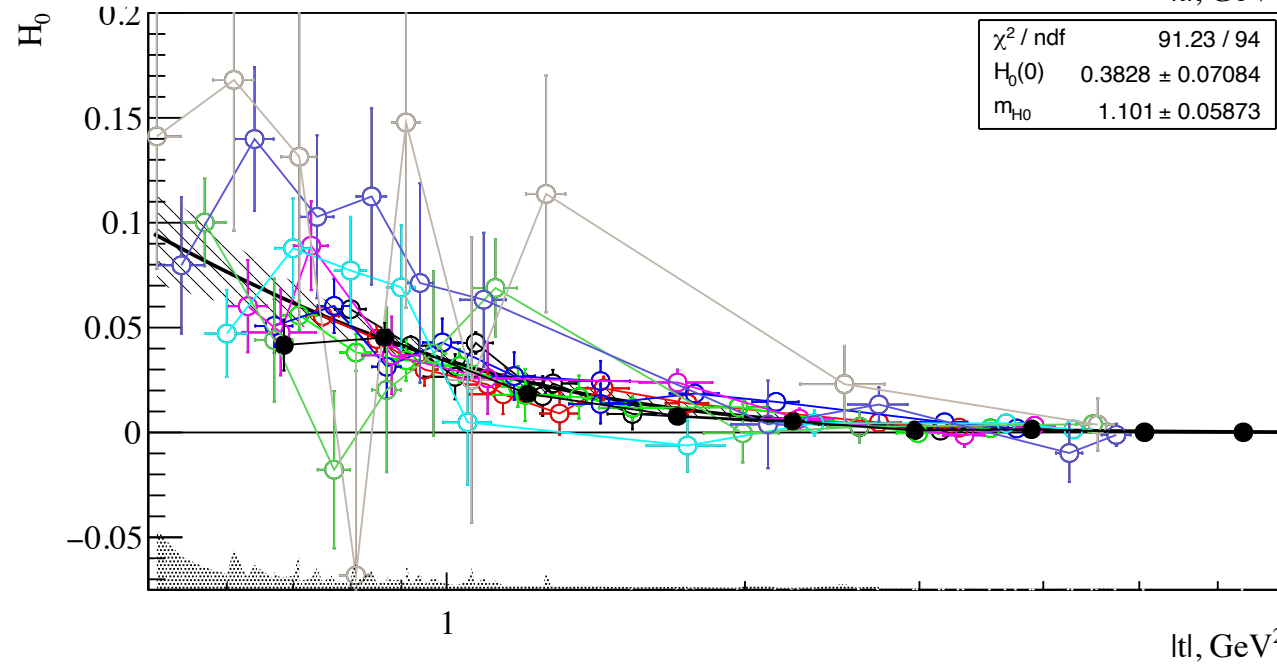


Energy independence of the $G(t)/H(t)$ functions (within errors) in agreement with the ξ/η -scaling

Fits with:

$$\frac{G_0(0)}{(1 - t/m_{G_0}^2)^4}$$

$$\frac{H_0(0)}{(1 - t/m_{H_0}^2)^4}$$



Using GlueX and J/ψ -007 data - different colors - different energies

Phys.Rev.C 108 (2023)

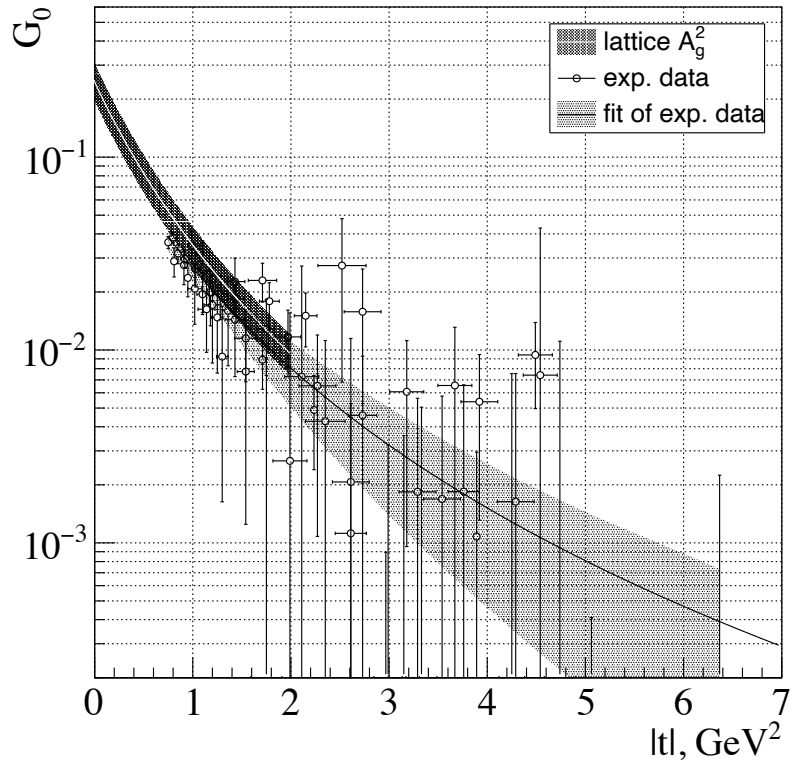
Nature 615 (2023)

LP and E.Chudakov

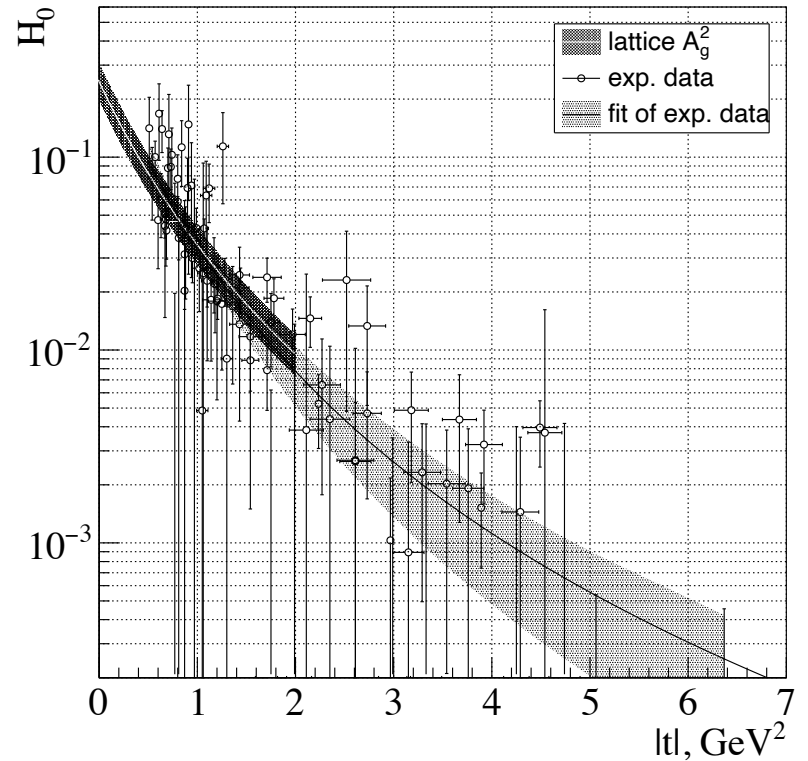
arXiv:2404.18776

Assuming leading-term approximation - data vs lattice

GPD



Holographic



$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t)\right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t)\right)^2 \quad \text{for high } \xi \text{ values}$$

$$\mathcal{A}_g^{(2)}(t) = A_g(t) \quad \text{leading-moment approximation}$$

$$\mathcal{B}_g^{(2)}(t) = B_g(t) \quad \text{approximation}$$

neglecting $B_g(t)$

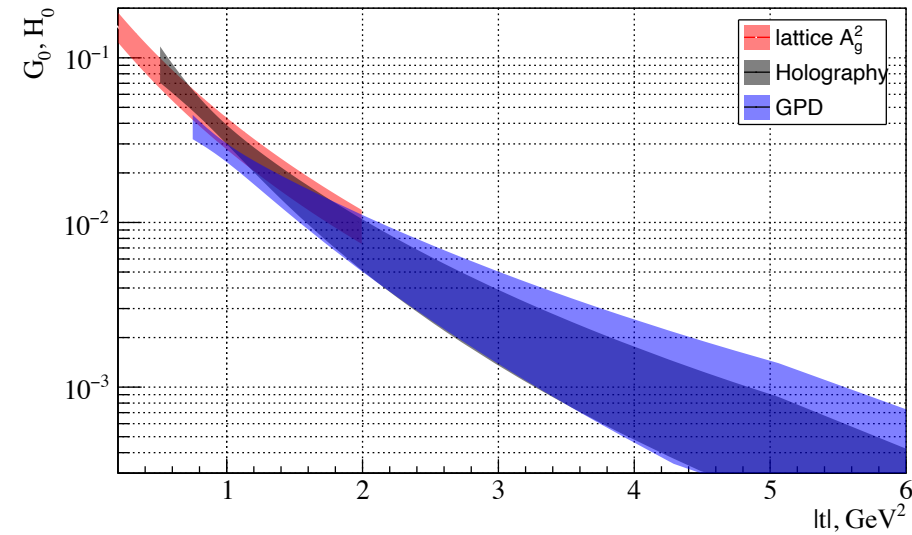
$$G_0(t) = A_g^2(t)$$

$$B_g(t) = 0$$

$$H_0(t) = A_g^2(t)$$

for large N_c and strong α_s

Gluon Gravitational Form Factors - summary

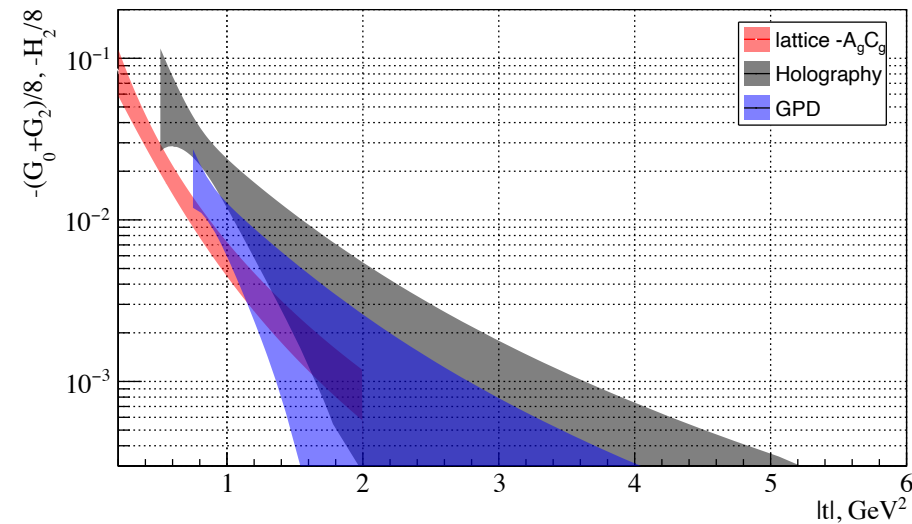


- Assuming $d\sigma/dt \sim \xi^{-4}[G_0(t) + \xi^2 G_2(t)]$ (GPD) or $d\sigma/dt \sim H_0(t) + \eta^2 H_2(t)$ (Holography), we extracted the **FF functions** kinematically from the data and found they are **indeed energy independent** (within the experimental errors)
- In leading-term approximation (and neglecting B_g):

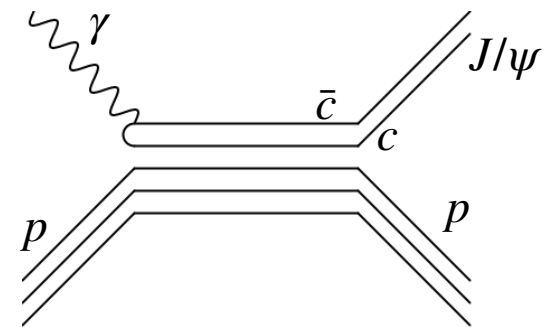
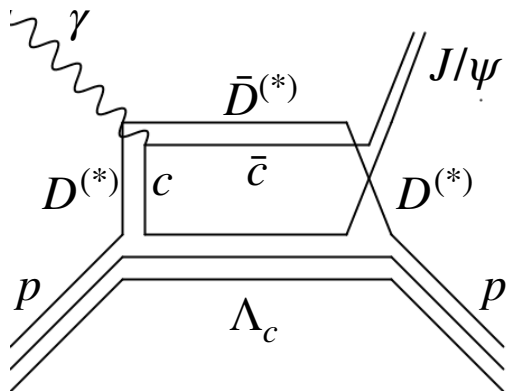
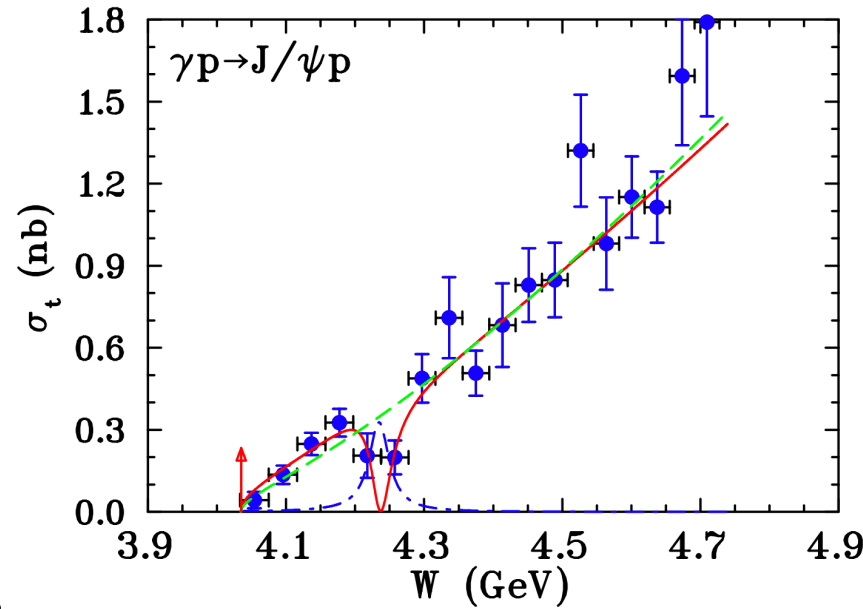
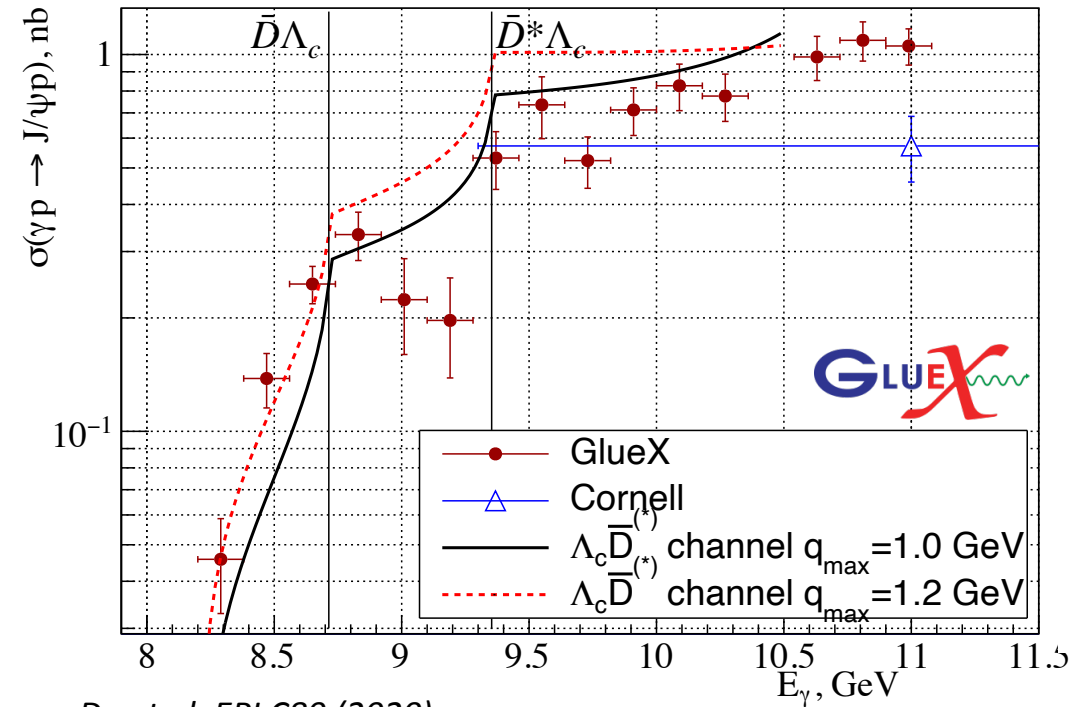
$$G_0(t) = H_0(t) = A_g^2(t) \text{ and}$$

$$G_0(t) + G_2(t) = H_2(t) = 8C_g(t)A_g(t)$$

- General agreement b/n extracted FFs using two diametric theories, each with specific corrections very different in nature (higher moments, $1/N_c, \dots$) and **agreement with lattice** - possible conclusion: **the corrections to the theories are not dominant**
- The above results should not be considered as a proof of the validity of the two theories, just that the data are consistent with them; in fact the observed ξ/η -scaling may come from more general considerations

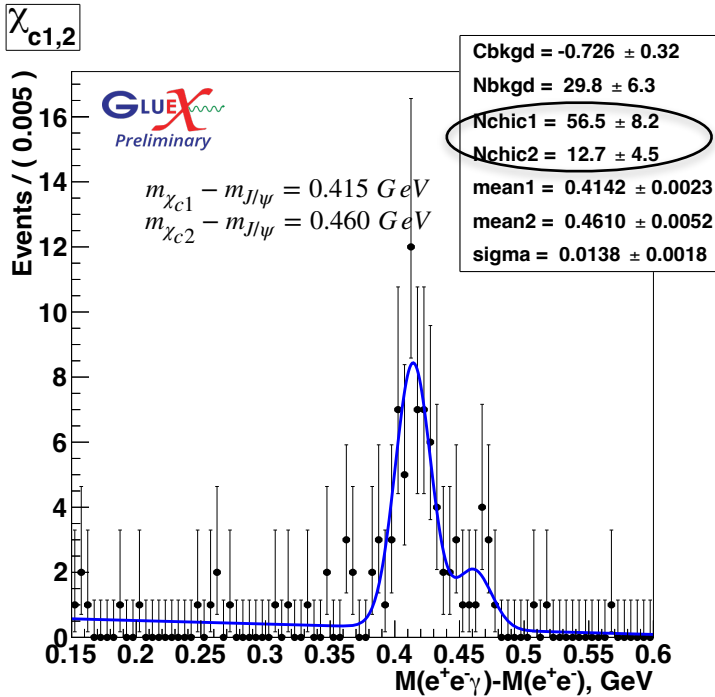


Other reaction mechanisms: open-charm, resonance



JPAC Phys.Rev.D 108 (2023)

Higher-mass charmonium states at threshold



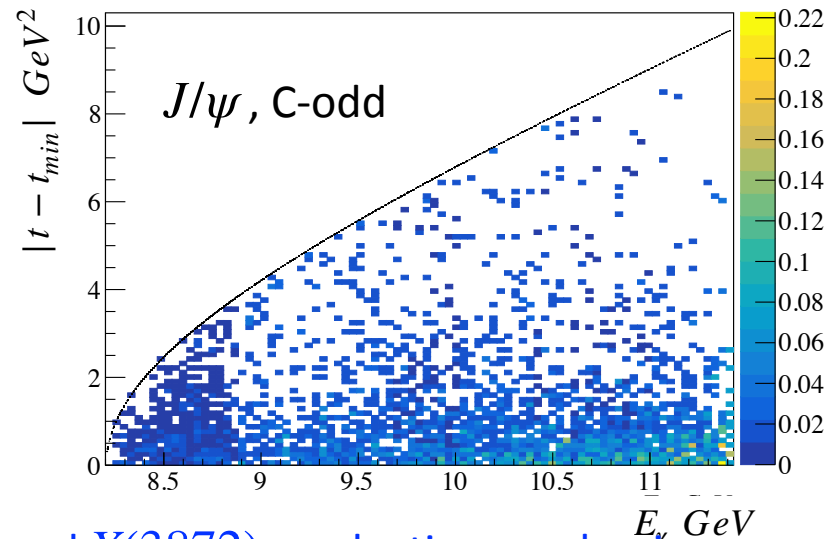
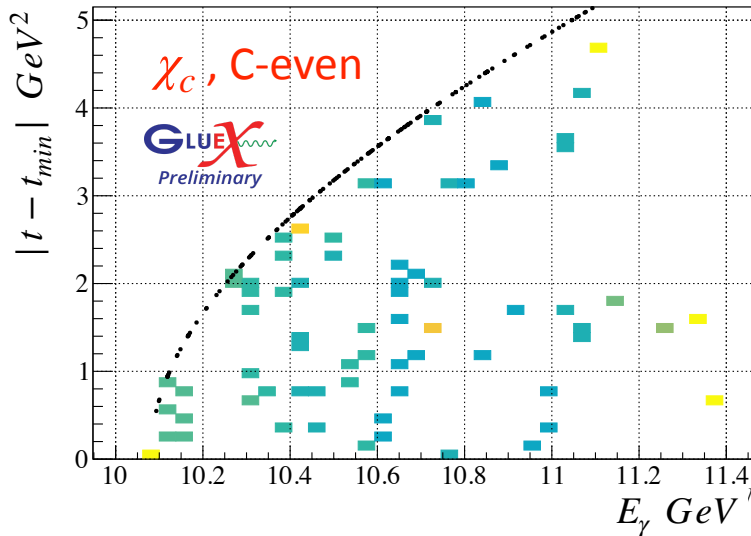
$$\gamma p \rightarrow \chi_c p \rightarrow (J/\psi \gamma) p \rightarrow (e^+ e^- \gamma) p$$

$\chi_{c1}(3511)$ and $\chi_{c2}(3556)$, 1^{++} and 2^{++} ,

$$E_\gamma^{thr} = 10.1 \text{ GeV}$$

- First ever evidence for photoproduction of C-even charmonium
- Studying χ_c states - complementary to J/ψ in understanding reaction mechanism near threshold

Dramatic difference: χ_c distribution in (E_γ, t) vs J/ψ

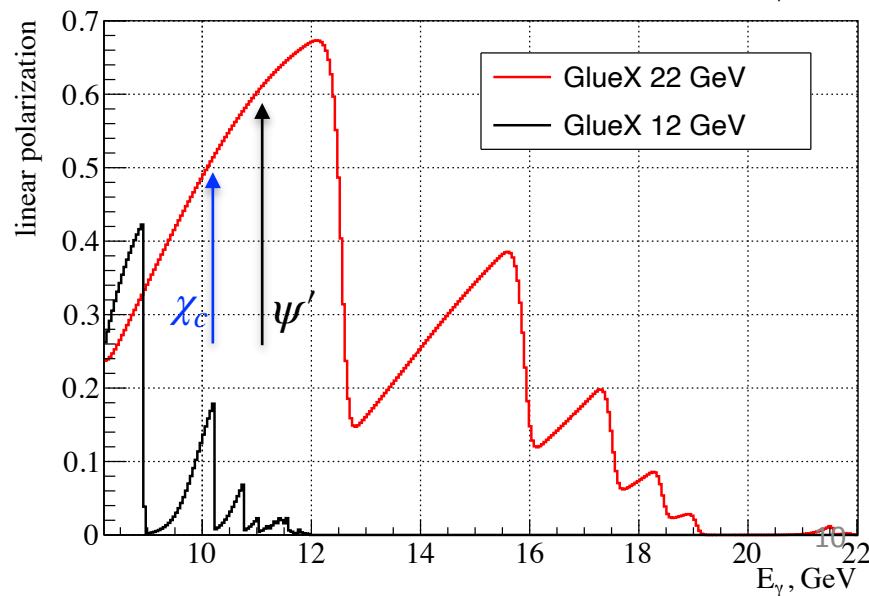
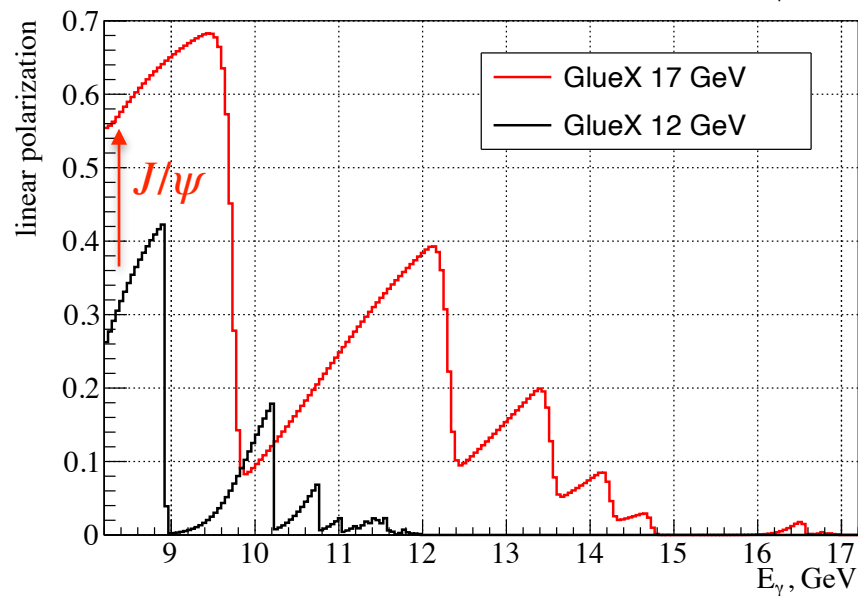
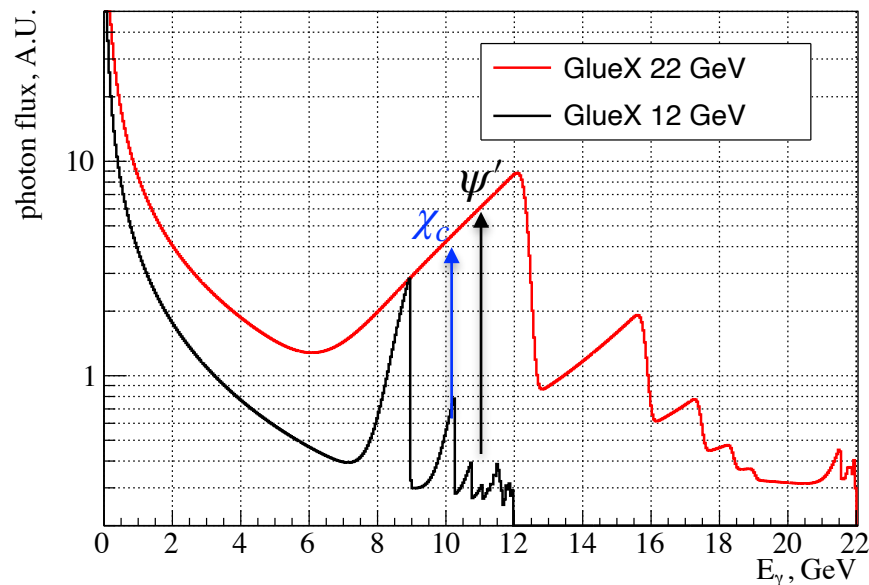
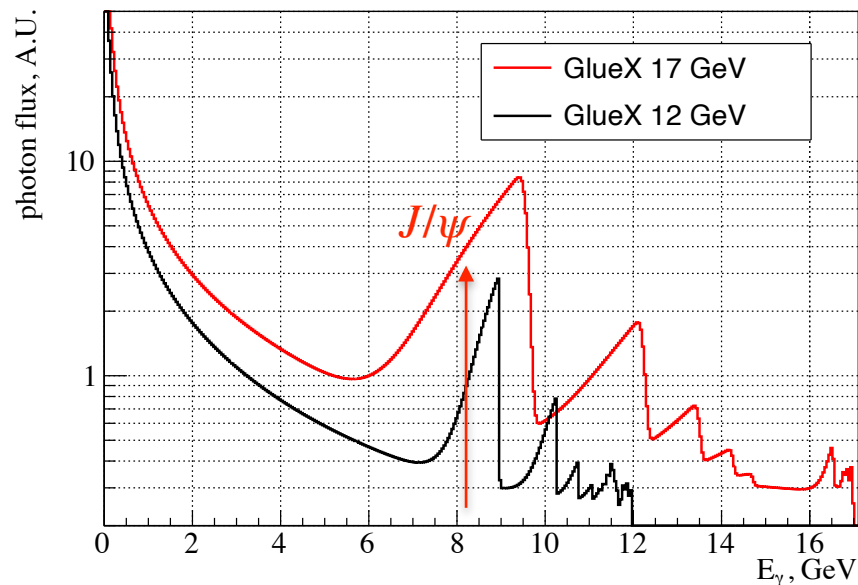


Studying χ_c can help to understand $X(3872)$ production mechanism

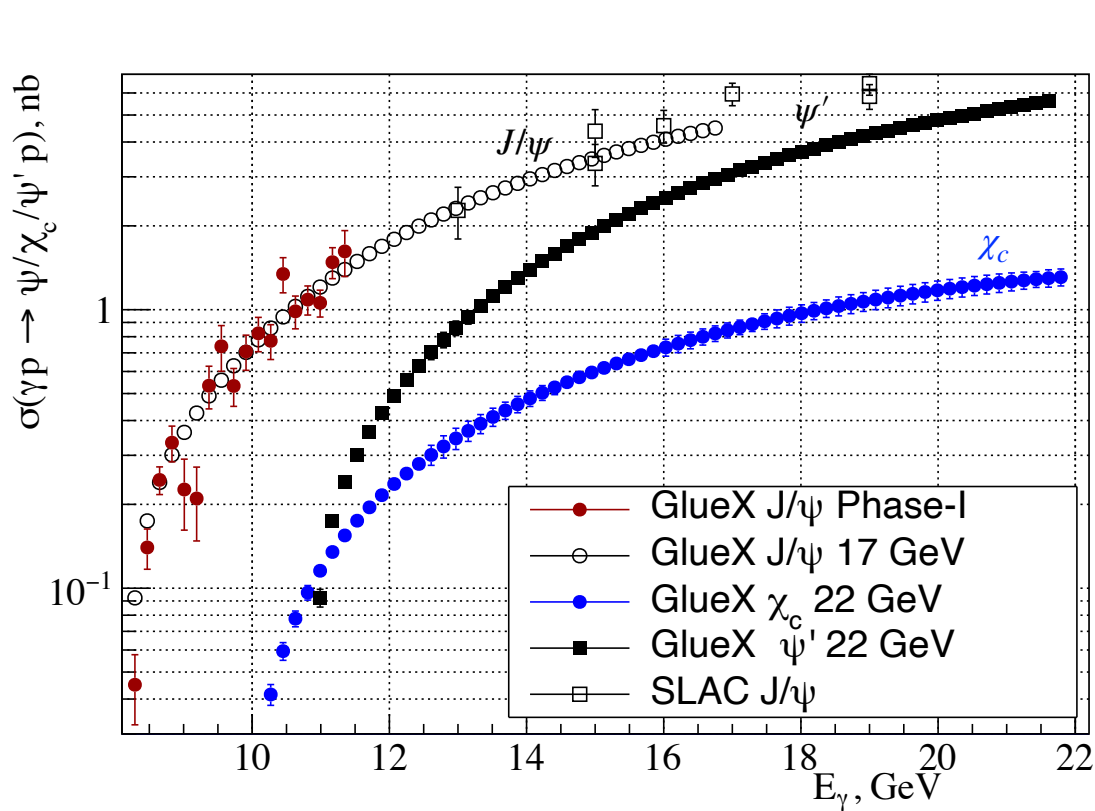
Threshold charmonium photoproduction at JLab22 with GlueX

GlueX uses polarized photon beam from coherent Bremsstrahlung

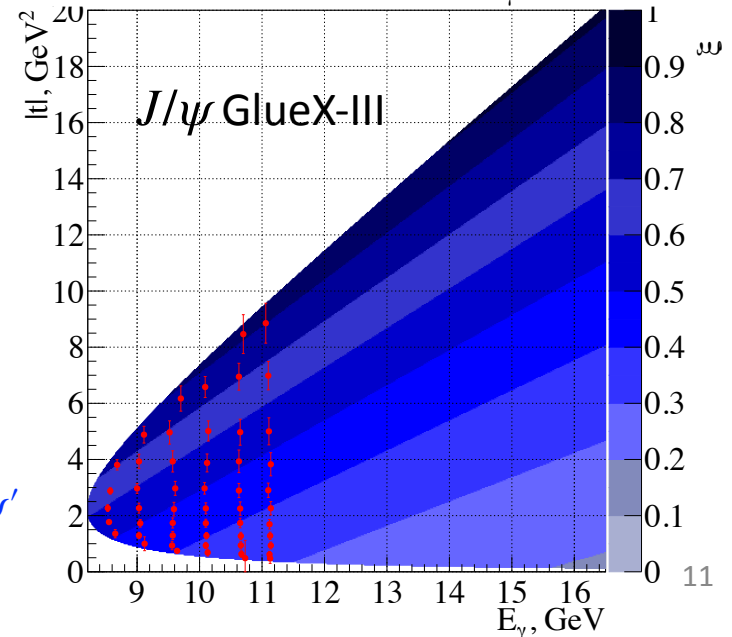
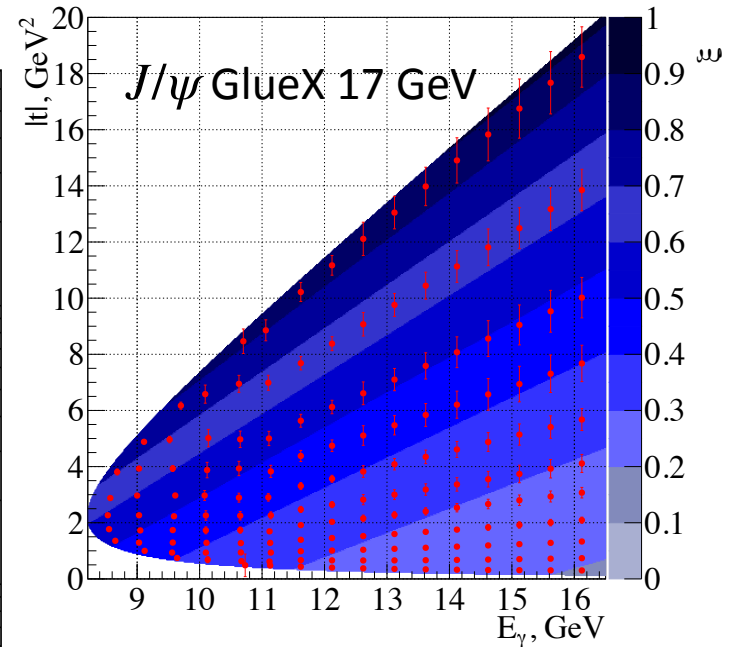
Taking advantage of increased end-point (electron beam energy):



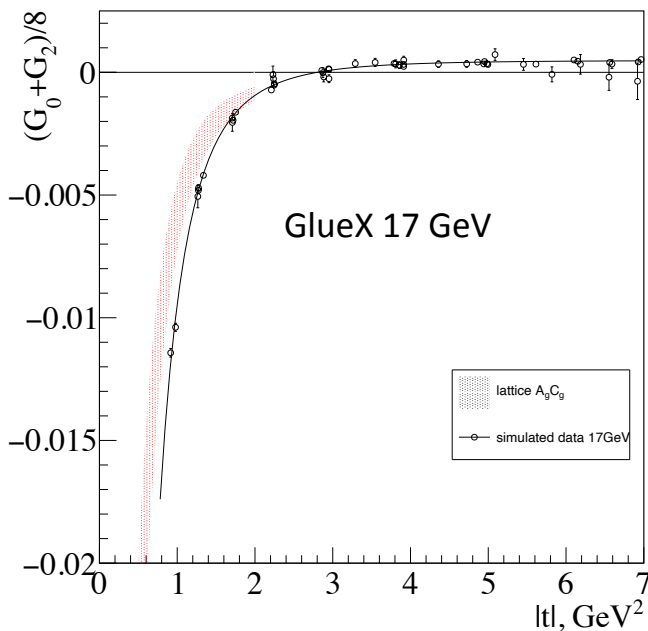
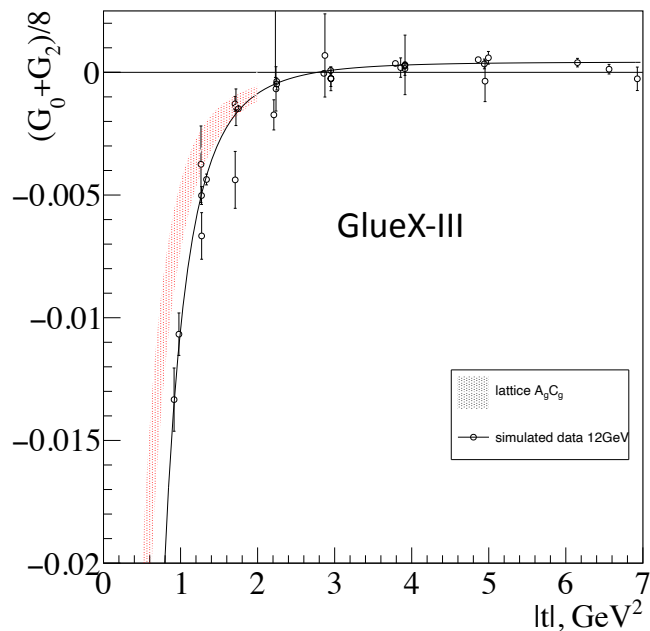
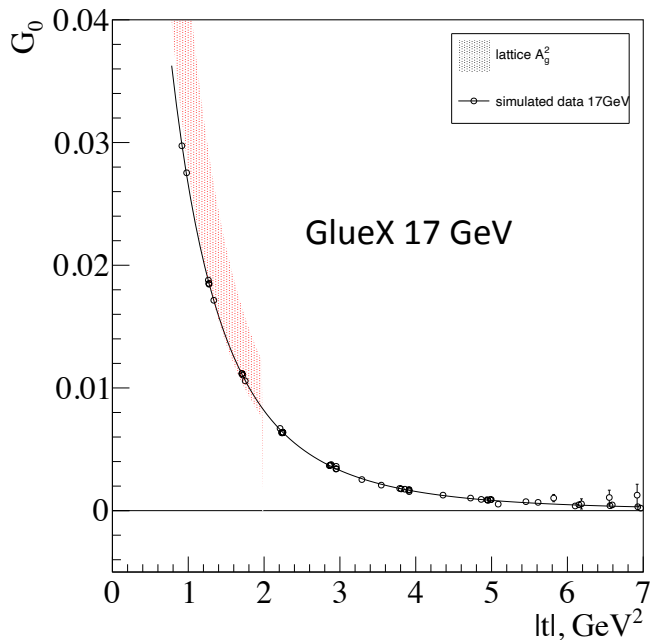
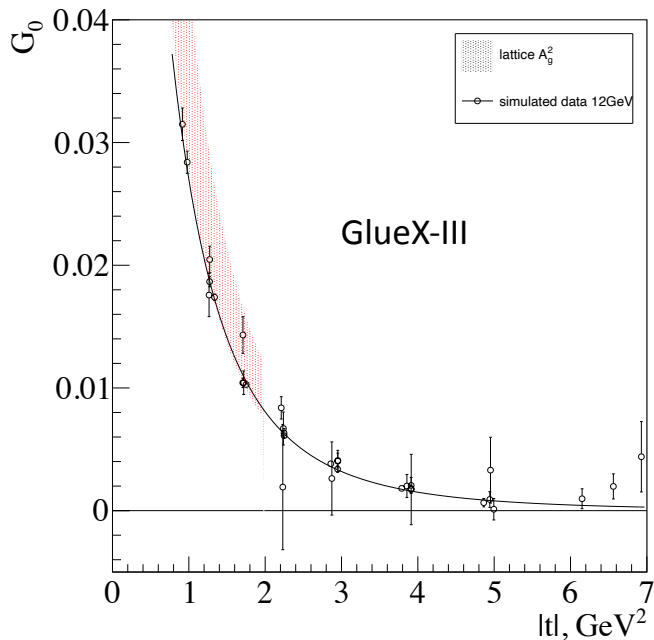
Threshold charmonium photoproduction at JLab22 with GlueX



- Absolute normalization and error estimation: scaling existing measurements (<12 GeV) by anticipated luminosity increase
- Energy dependence extrapolation: using fits of J/ψ data for J/ψ and ψ' , for χ_c based on VM exchange model (JPAC, *Phys.Rev.D* 102 (2020))
- Efficiencies estimated with MC
- All JLab22 projections for 100 PAC days: 80k J/ψ , 8k χ_c , 18k ψ'



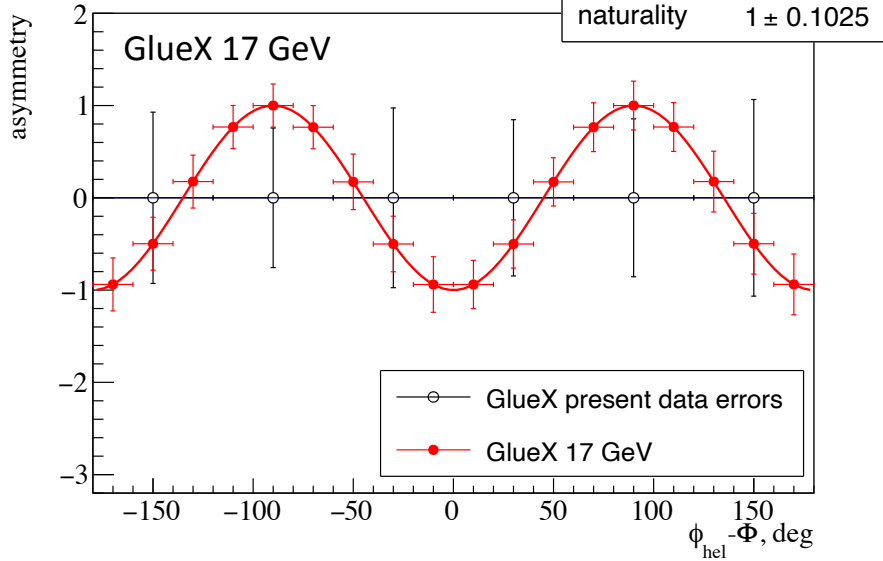
Threshold charmonium photoproduction at JLab22 with GlueX



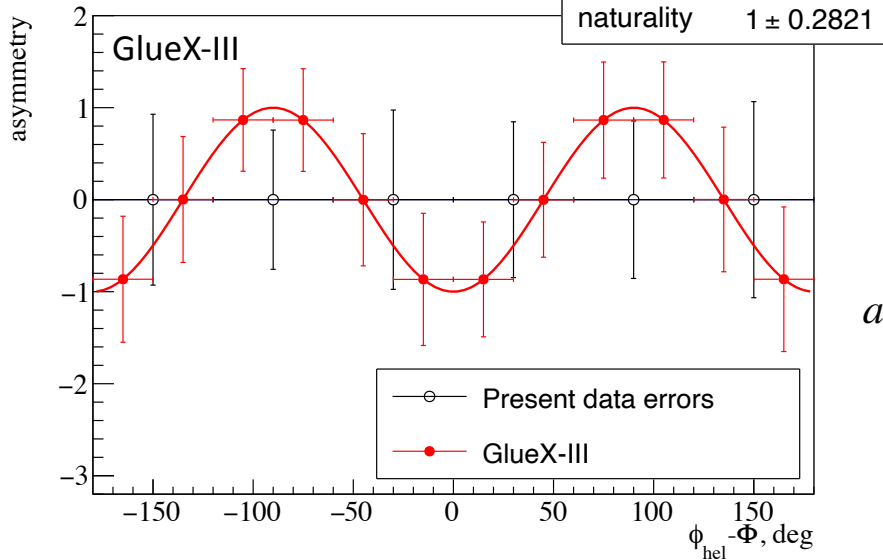
- Anticipated results for the extracted gluon Form Factors, assuming skewness scaling is valid
- Data randomized around a fit of the current data

Threshold charmonium photoproduction at JLab22 with GlueX

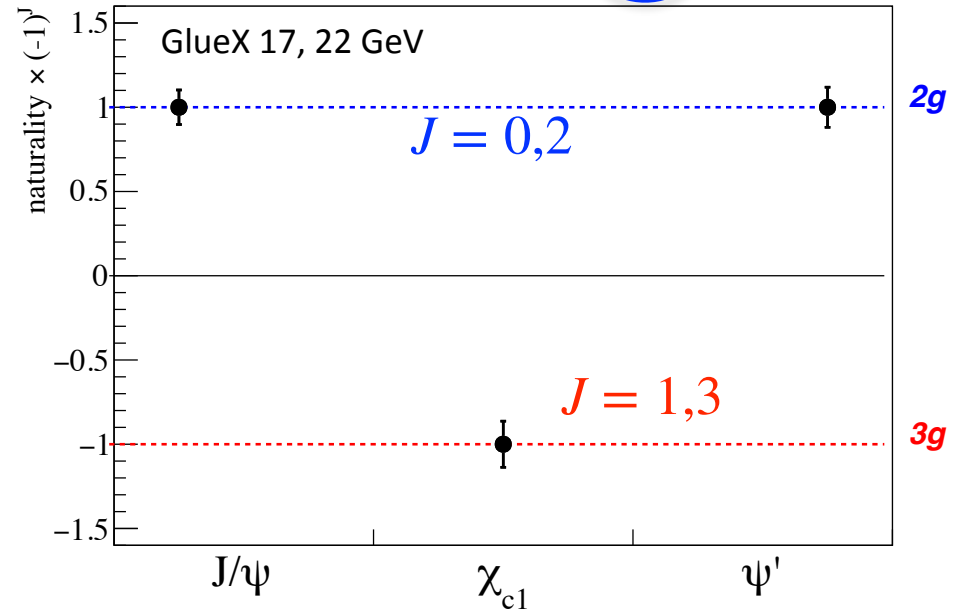
$\gamma p \rightarrow J/\psi p \rightarrow e^+e^- p$



$\gamma p \rightarrow J/\psi p \rightarrow e^+e^- p$



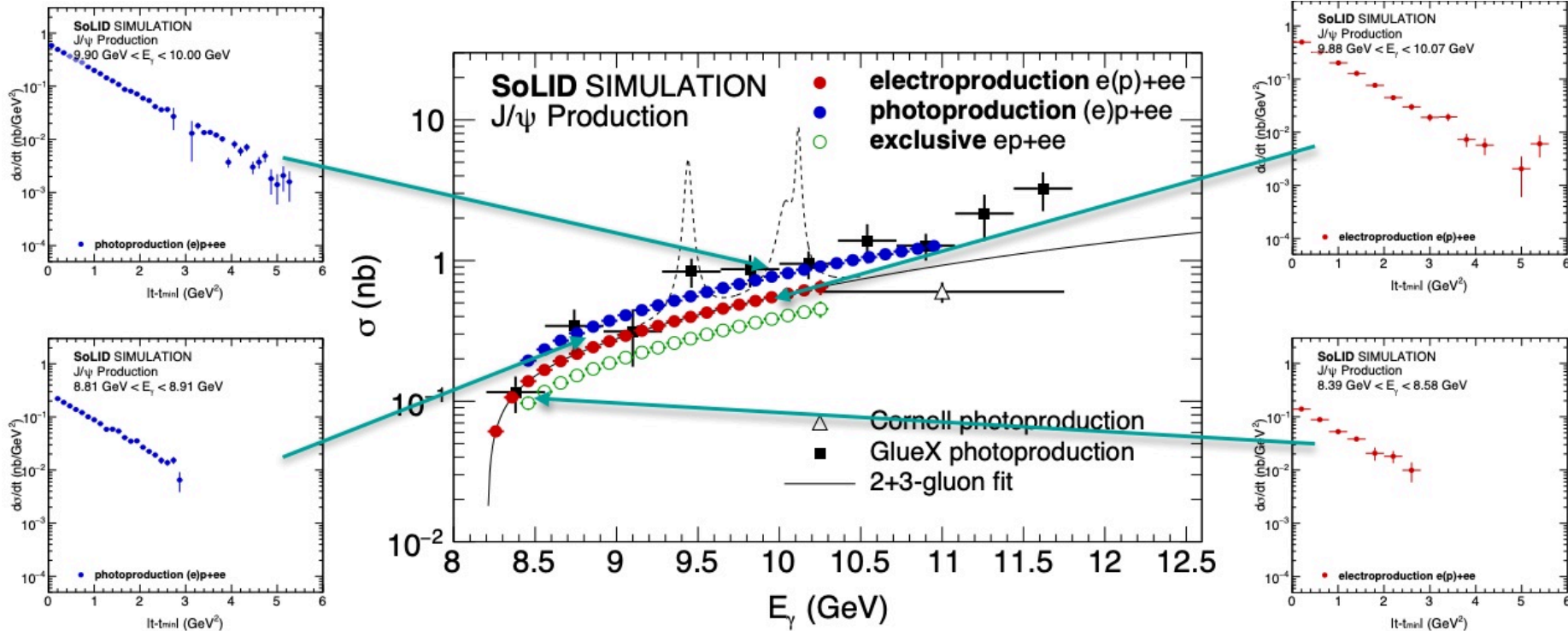
$$\text{naturality} \times (-1)^J = P$$



$$\begin{aligned} \text{asymmetry} &= \frac{2}{P_\gamma} \frac{Y_{J/\psi}(0) - Y_{J/\psi}(90)}{Y_{J/\psi}(0) + Y_{J/\psi}(90)} = \\ &= -(\rho_{1-1}^1 - \text{Im}\rho_{1-1}^2) \cos[2(\phi_{hel} - \Phi)] \end{aligned}$$

$$\text{naturality} = (-1)^J P$$

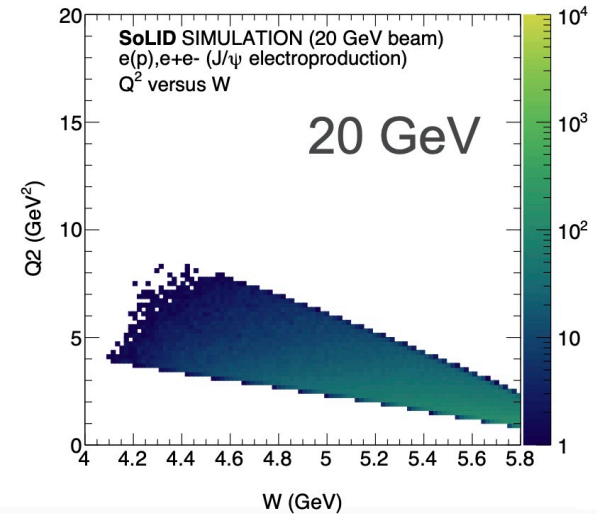
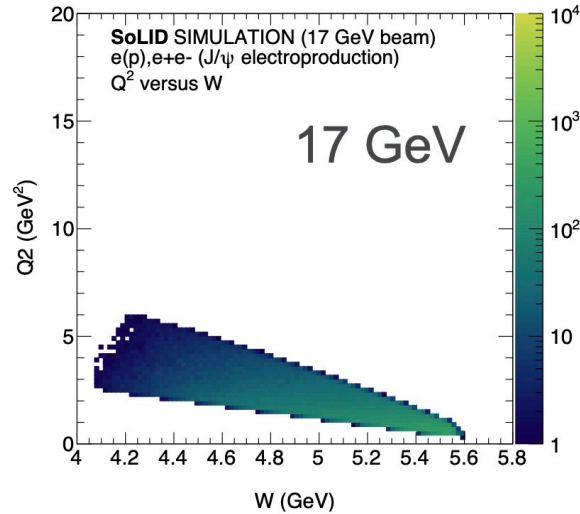
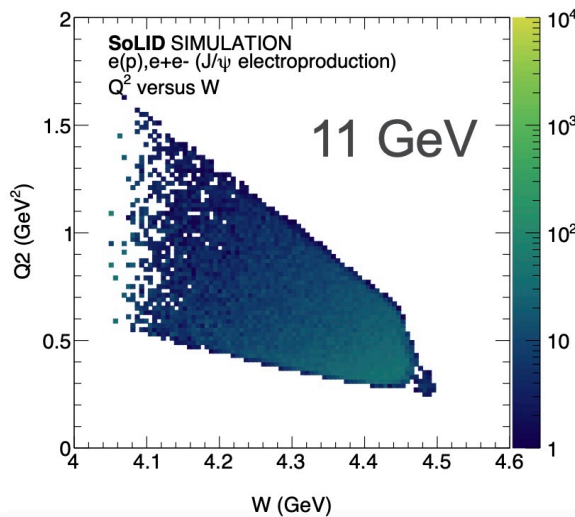
SoLID at JLab12



From S.Joosten talk at "J/ψ and Beyond" workshop

- SoLID at 12 GeV would give $\sim 1M J/\psi$'s per year in photoproduction covering practically full t -region
- ... and $\sim 20k J/\psi$'s in electro-production with $Q^2 < 1.5 \text{ GeV}^2$, which however is not enough to study mass dependence of the charmonium production, as $Q^2 \ll M_{J/\psi}^2$
 see Boussarie and Hatta, *Phys. Rev. D* 101 (2020)

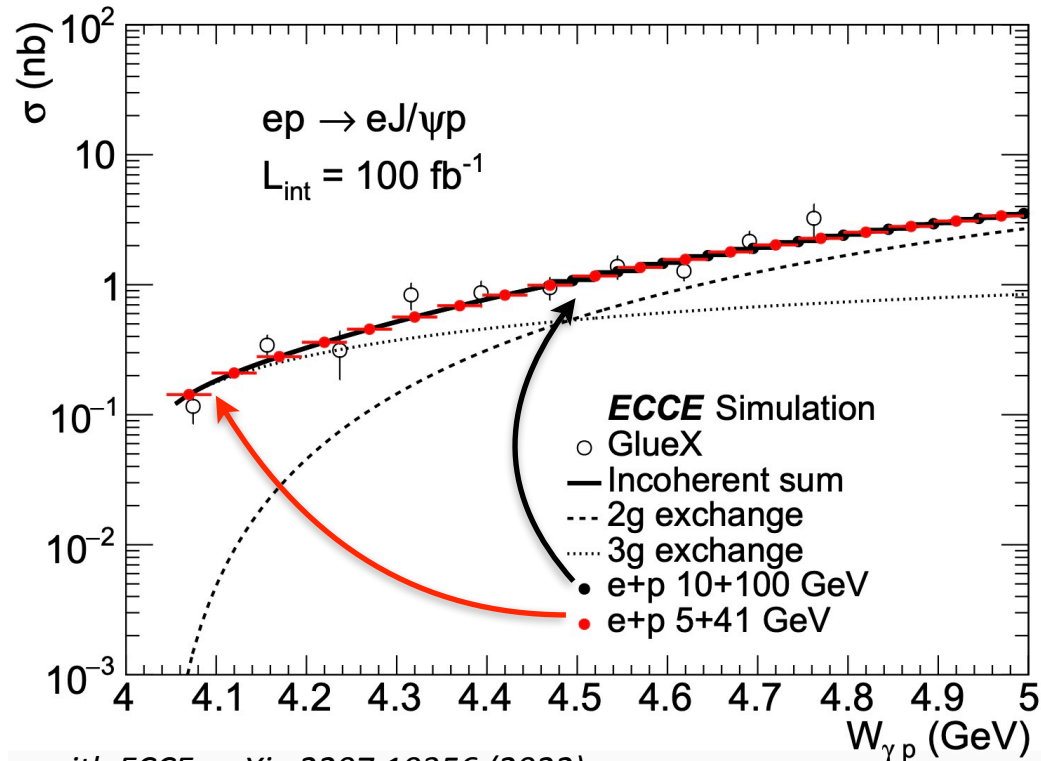
SoLID at JLAB22



From S.Joosten talk at “ J/ψ and Beyond” workshop

- However SoLID at 22 GeV in J/ψ electroproduction can reach $Q^2 \sim 8 \text{ GeV}^2 \approx M_{J/\psi}^2$, effectively increasing the mass of charmonium
- SoLID angular acceptance is in the range $7 - 28^\circ$, limiting in J/ψ photoproduction medium t -region for $E_\gamma > 12 \text{ GeV}$
- Good coverage for $\psi(2S)$ with acceptance limitations for $E_\gamma > 16 \text{ GeV}$ in photoproduction

J/ψ production close to threshold at EIC



Exclusive J/ψ detection with ECCE: [arXiv:2207.10356](https://arxiv.org/abs/2207.10356) (2022)

- 5+41 GeV beam configuration needed to reach threshold, however it has **very low luminosity** (4.4 fb^{-1} annually)
- Difficult to detect scattered electron with energy close to electron beam energy - **bad $W_{\gamma p}$ resolution**
- Using the detected J/ψ and proton to calculate $W_{\gamma p}$ also leads to a bad resolution (*Sylvester Joosten*)

Heavy quarkonium production close to threshold

	GlueX 22GeV	Solid 22GeV	EIC
J/ψ photoproduction	✓	✓ Acceptance limitations for $E_\gamma > 12$ GeV	?
J/ψ electroproduction	✗	✓ unique up to $Q^2 = 8$ GeV ²	?
χ_c photoproduction	✓ unique	✗	?
$\psi(2S)$ photoproduction	✓	✓ Acceptance limitations for $E_\gamma > 16$ GeV	?
$\psi(2S)$ electroproduction	✗	✓ up to $Q^2 = 1.5$ GeV ²	?
$J/\psi, \chi_c, \psi(2S)$ linear polarization	✓ unique	✗	✗
Υ	✗	✗	✓ unique

- GlueX has linear polarization and almost full acceptance for multi-particle final states (including photons) - unique in polarization measurements and χ_c states
- SoLID has very high luminosity, relatively wide acceptance, capable of reaching high Q^2 values in electroproduction (unique)
- EIC energies much above charmonia thresholds, detection is questionable, however very suitable for studying Υ production at threshold

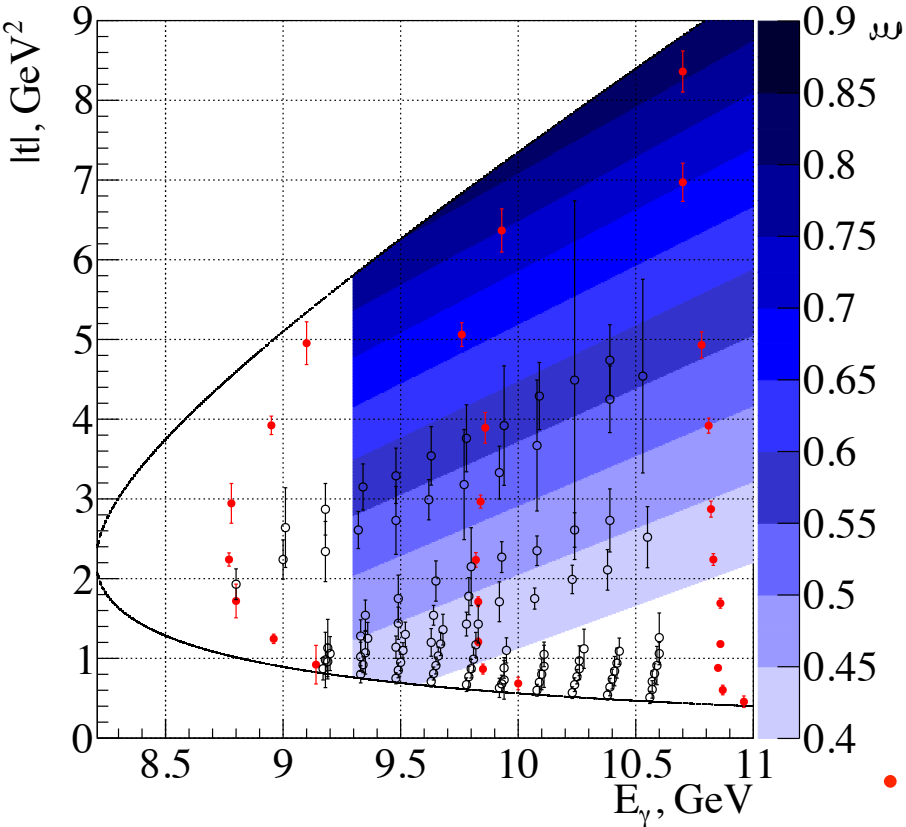
Conclusions

- It is remarkable to see the **gluon GFFs extracted from J/ψ data** using two diametric approaches (without any external constraints) **on the same scale with the lattice calculations**. However, more precise relation to the gluon properties of the proton requires comprehensive experimental and theoretical studies to better understand the reaction mechanism.
- **CEBAF energy upgrade adds three new dimensions** to these studies:
 - Threshold production of **higher-mass charmonium states** (with different quantum numbers) - $\psi(2S)$ (SoLID, GlueX), χ_c (GlueX)
 - Threshold charmonium **electroproduction at high Q^2** (SoLID)
 - **Polarization measurements with high FOM** (GlueX)
 - What about open-charm production, that is supposed to dominate with increasing the energy
- EIC would extend the charmonium studies at JLab to the bottomonium sector

Back up slides

Data used in gluon Form Factor extraction

GPD



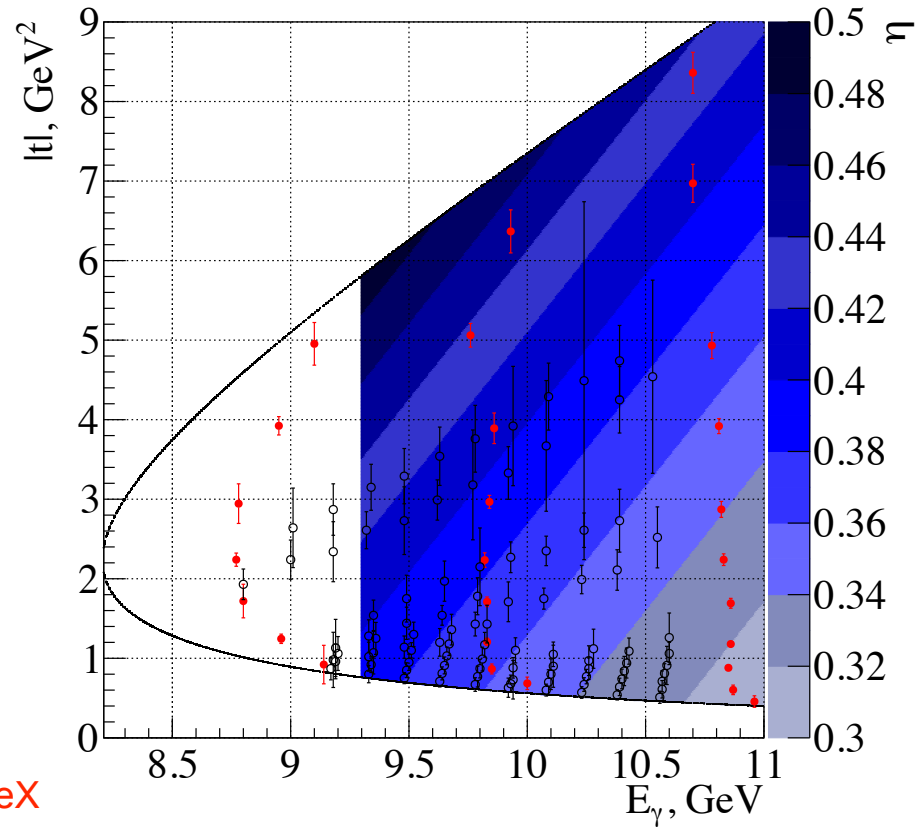
$\xi > 0.4$

$E_\gamma > 9.3 \text{ GeV}$ (away from $\bar{D}\Lambda_C$ and $\bar{D}^*\Lambda_C$ thresholds)

• GlueX

◦ J/ψ -007
Nature 615 (2023)

Holography



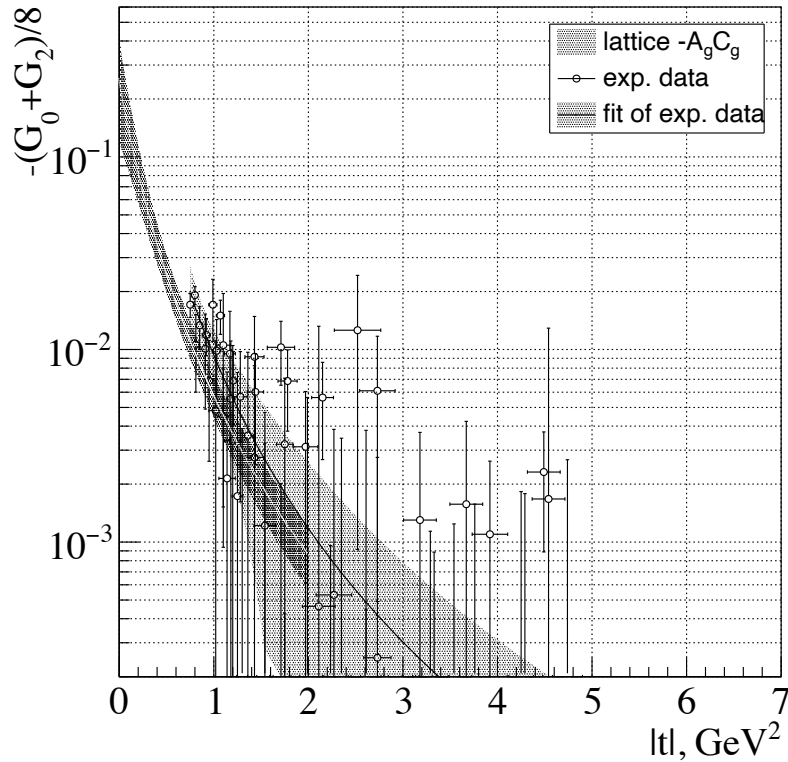
no constraints on η

$E_\gamma > 9.3 \text{ GeV}$ (away from $\bar{D}\Lambda_C$ and $\bar{D}^*\Lambda_C$ thresholds)

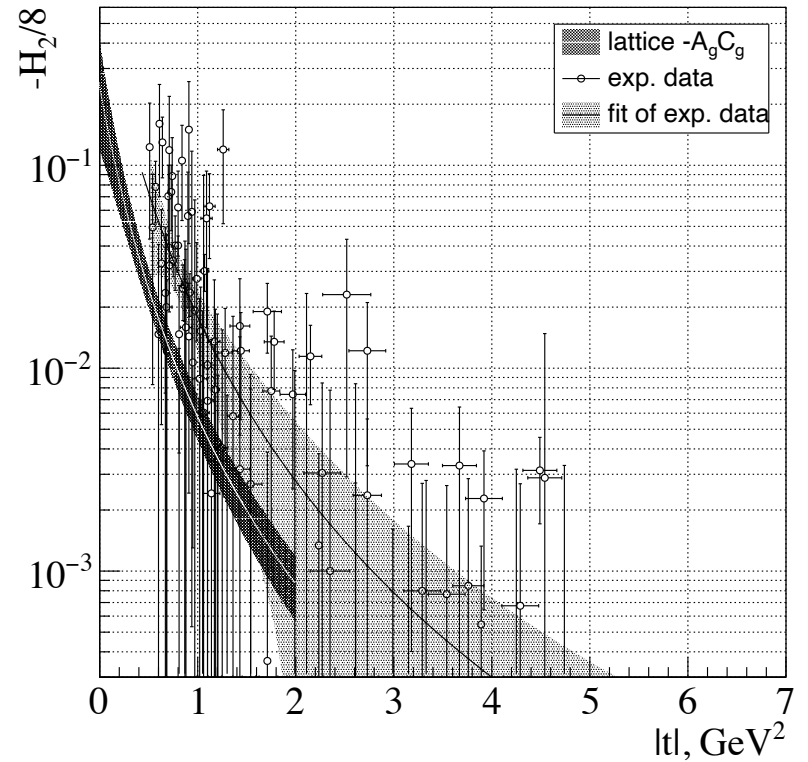
Error bars: relative errors in A.U. (not related to y-axis)

Assuming leading-term approximation - data vs lattice

GPD



Holographic



$$G_0(t) = \left(\mathcal{A}_g^{(2)}(t) \right)^2 - \frac{t}{4m^2} \left(\mathcal{B}_g^{(2)}(t) \right)^2$$

$$G_2(t) = 2\mathcal{A}_g^{(2)}(t)\mathcal{C}_g(t) + 2\frac{t}{4m^2}\mathcal{B}_g^{(2)}(t)\mathcal{C}_g(t) - \left(\mathcal{A}_g^{(2)}(t) + \mathcal{B}_g^{(2)}(t) \right)^2$$

$$\mathcal{A}_g^{(2)}(t) = A_g(t), \quad \mathcal{B}_g^{(2)}(t) = B_g(t),$$

$$\mathcal{C}_g(t) = 4C_g(t) \quad \text{leading-moment approximation}$$

neglecting $B_e(t)$

$$G_0(t) + G_2(t) = 8A_g(t)C_g(t)$$

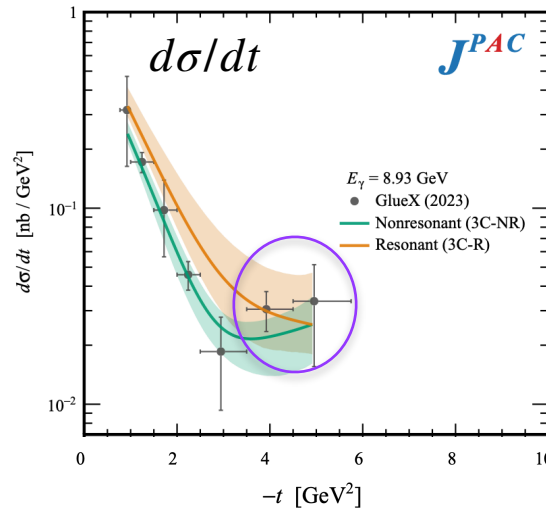
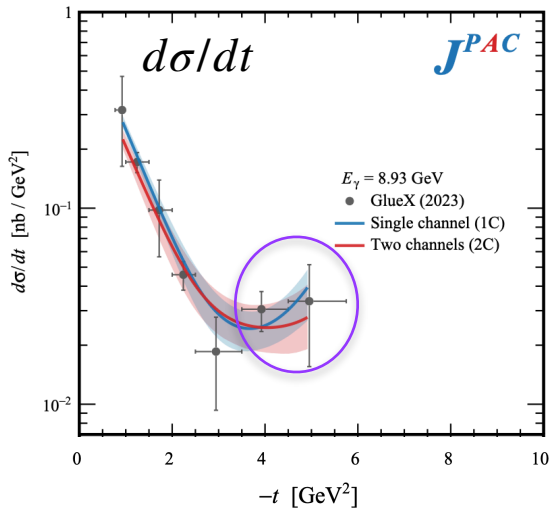
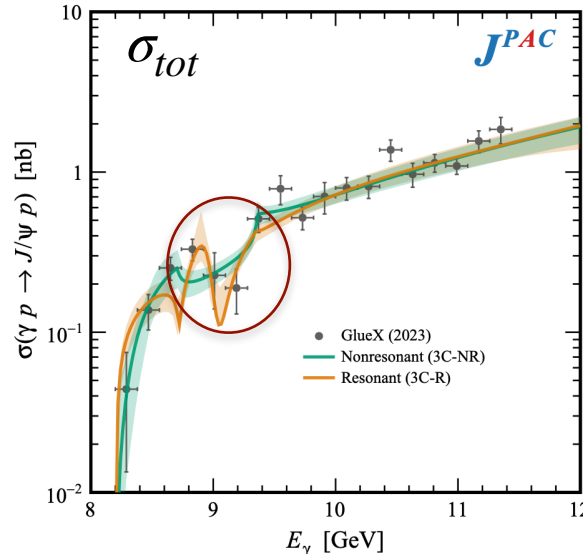
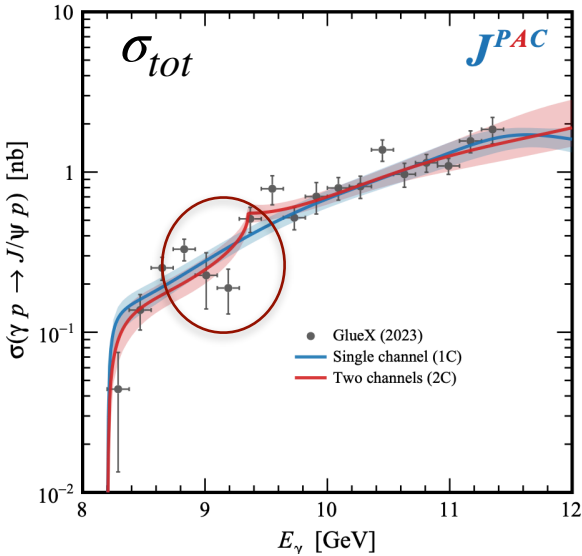
for high ξ values

$$B_g(t) = 0$$

$$H_2(t) = 8A_g(t)C_g(t)$$

for large N_c and strong α_s

Phenomenological approach: JPAC results



Phenomenological model based on s-channel PW expansion ($l \leq 3$):

- (1C) $J/\psi p$ interaction
- (2C) $J/\psi p$ and $\bar{D}^* \Lambda_C$
- (3C-NR) $J/\psi p$, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (non-resonant solution)
- (3C-NR) $J/\psi p$, $\bar{D} \Lambda_C$, $\bar{D}^* \Lambda_C$ (resonant solution)

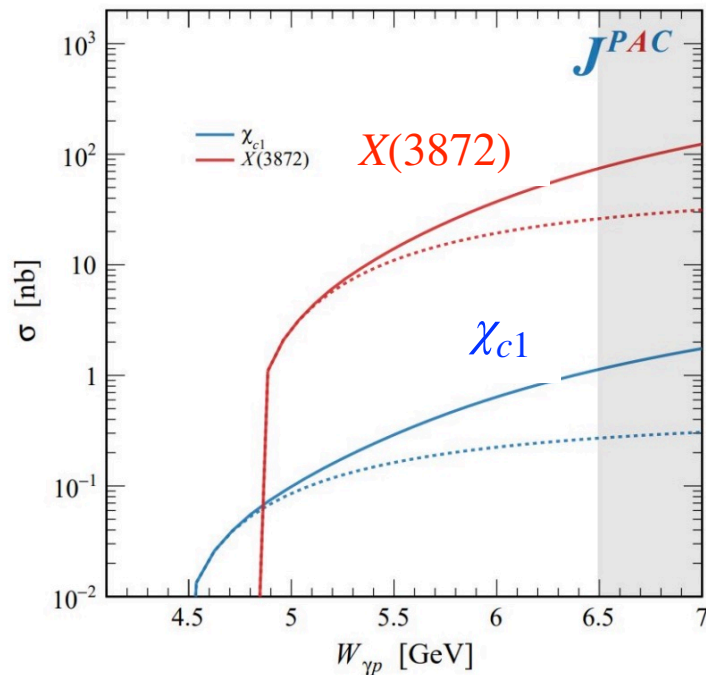
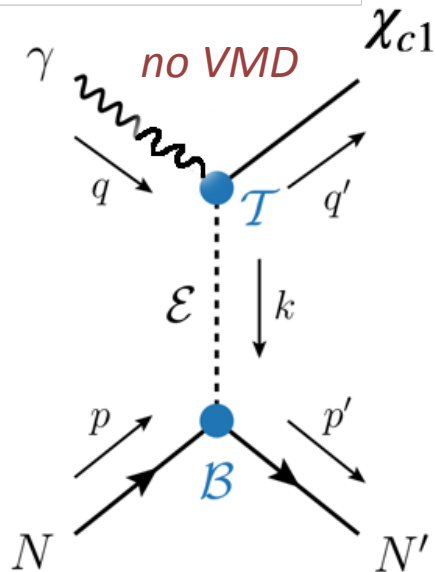
No stat. significant preference:

- 9 GeV structure requires sizable contribution from open charm
- Severe violation of VMD and factorization not excluded
- s-channel resonance not excluded
- t-enhancement indicates s-channel contribution: due to proximity to threshold or open-charm exchange

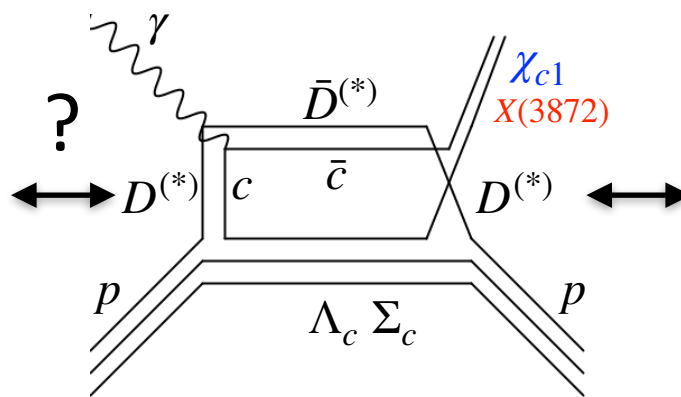
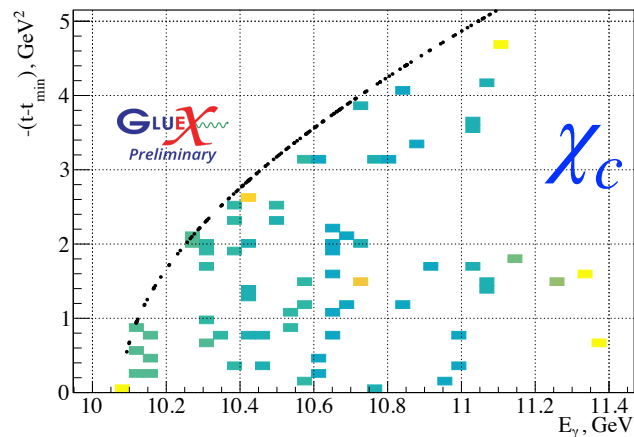
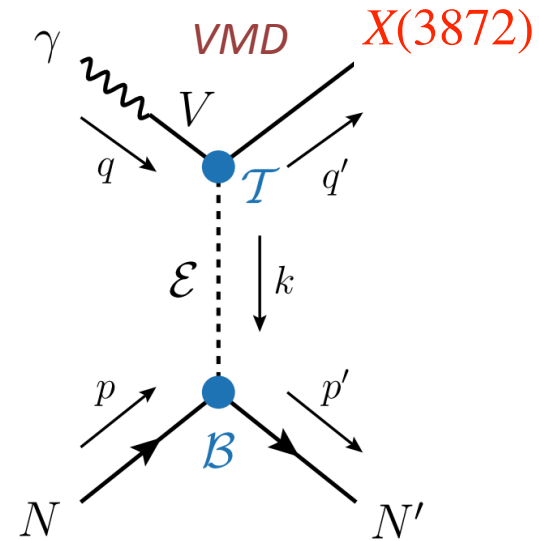
JPAC Phys.Rev.D 108 (2023)

Global fit of both Hall C & D $d\sigma/dt(t)$ and Hall D $\sigma_{tot}(E_\gamma)$

χ_c vs $X(3872)$ production



JPAC, PRD 102 (2020). t -channel production



$X(3872)$

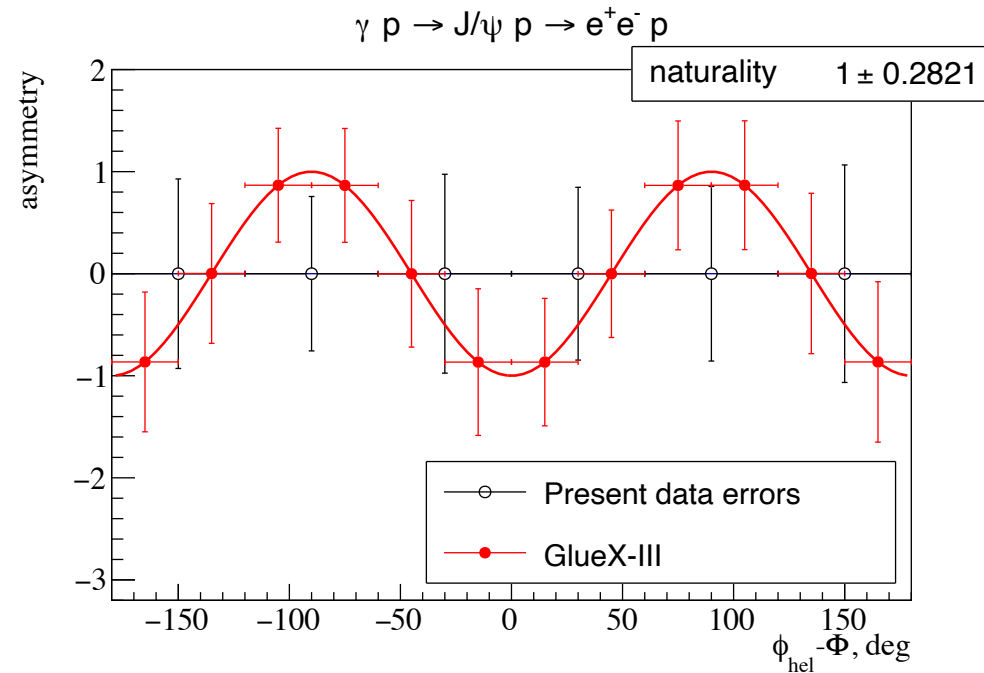
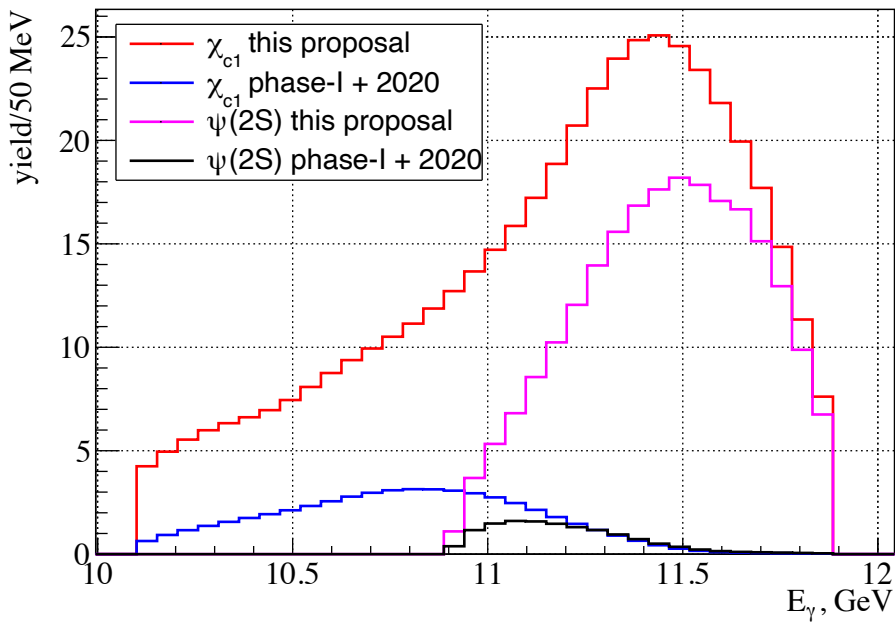
?

Recall discussions on Monday - Daniel's talk p.9

Studying χ_c can help to understand $X(3872)$ production mechanism

Prospect for charmonium threshold production with GlueX-III

Run Period	J/ψ	χ_{c1}	$\psi(2S)$
2016-2020 Phase I-II	3,960	55	12
2023-2025 Phase II (planned)	3,615	48	11
Phase III (proposal)	11,271	364	178
Projected Total	18,846	467	201



Prospect for charmonium threshold production with GlueX-III

