OPTIMAL ANTI-FERROMAGNETS FOR LIGHT DARK MATTER DETECTION

Angelo Esposito







Istituto Nazionale di Fisica Nucleare

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- One of the strongest evidences for physics beyond the Standard Model
- However... huge possible mass range
 →
 detection techniques
 vary widely depending on the dark matter mass

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 $10^{6} \,\mathrm{eV}$ $10^{9} \,\mathrm{eV}$

 $10^{11} \, \mathrm{eV}$

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 $10^{-22} \, \text{eV}$



 $10^{-4} \, \text{eV}$

 $10^{40} \, \mathrm{eV}$















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- Need new materials and/or observables



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 To evade this we must look into inelastic processes, which relax the kinematics constraints

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Condensed matter	Atomic physics	Nuclear physics	
$q \ll 1/a_B$	$q \sim 1/a_B$	$q \lesssim \Lambda_{QCD}$	
10) ⁶ 1	$0^8 \qquad m_{\chi}$	(eV)



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_	d << 110B	d nor		1~	<u>Q</u> CD		>
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• Need theoretical tools that allow to solve or bypass these problems



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Typically, no more single particle final states
 signatures

[e.g., Trickle et al. - JHEP 2020, 1910.08092; Griffin et al. - PRD 2020, 1910.10716]



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• A growing field, with many interesting ideas at the edge between condensed matter and particle physics



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 - I. Solid crystals (GaAs, SiO, ...) -> multi-phonon

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[for a review, Kahn, Lin - Rept.Prog.Phys. 2022, 2108.03239]

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 Ways to detect few magnons have already been proposed (TES, MKID, quantum sensors)

[Trickle, Zhang, Zurek - PRL 2020, 1905.13744; Lachance-Quirion et al. - Science Advances 2017; Lachance-Quirion et al. - Science 2020]



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- Two benchmark models:

$$\mathcal{L}_{m.d.} \sim V_{\mu\nu} \,\bar{\chi} \sigma^{\mu\nu} \chi + V_{\mu} \,\bar{e} \gamma^{\mu} e$$
$$\mathcal{L}_{p.m.} \sim \phi \,\bar{\chi} \chi + \phi \,\bar{e} i \gamma^5 e$$

[e.g., Sigurdson et al. - PRD 2004, astro-ph/ 0406355; Chang, Weiner, Yavin - PRD 2010, 1007.4200]

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• At low energies:

$$\begin{aligned} \mathscr{L}_{m.d.} &\to \chi^{\dagger} \sigma^{i} \chi \left(\delta^{ij} - \nabla^{-2} \nabla^{i} \nabla^{j} \right) \, s_{i} \\ \mathscr{L}_{p.m.} &\to \chi^{\dagger} \chi \, \nabla^{-2} \, \nabla \cdot \mathbf{s} \end{aligned}$$



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[Trickle, Zhang, Zurek - PRL 2020, 1905.13744; Mitridate et al. - PRD 2020, 2005.10256; Chigus, Moroi, Nakayama - PRD 2020, 2001.10666; Trickle, Zhang, Zurek - PRD 2022, 2009.13534]



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Conservation of magnetization

only one magnon emitted

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$$\omega_{max} = E_{\chi} \frac{4 m_{\theta}/m_{\chi}}{(1 + m_{\theta}/m_{\chi})^2} \quad \text{with} \quad m_{\theta} \sim 1 \text{ MeV} \quad \longrightarrow \qquad \begin{array}{c} \text{inefficient for} \\ m_{\chi} \lesssim \mathcal{O}(\text{MeV}) \end{array}$$

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• Nickel-oxide has $v_{\theta} \sim v_{\chi} \rightarrow$ very efficient at absorbing dark matter energy [AE, Pavaskar - PRD (2023), 2210.13516]

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- In anti-ferromagnets you can always emit magnon and anti-magnon pairs and compensate for the charge
- Multi-magnon emission process evade the kinematical constraints and get down to $m_{\chi} \sim O(\text{keV})$



MAGNONS

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• At low energies/momenta magnons can be described by an EFT, invariant under the full symmetry group

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$$\mathbf{n}(x) = e^{i\left[\theta^{1}(x)J_{1} + \theta^{2}(x)J_{2}\right]} \cdot \hat{\mathbf{z}} \xrightarrow{SO(3)} R \cdot \mathbf{n}(x)$$
magnon fields



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$$= c_1 \left(\dot{\theta}^a \right)^2 - c_2 \left(\nabla_i \theta^a \right)^2 + \dots$$

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- Structure completely dictated by symmetry \rightarrow just need c_1
- This allows to bypass difficulties in the standard treatment (failure of the Holsten-Primakoff approach) [Dyson - Phys. Rev. 1956]

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• Use standard QFT methods to compute event rates

[AE, Pavaskar - PRD (2023), 2210.13516]



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Thank you for the attention!



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