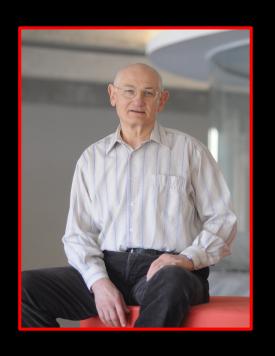
MOND: An alternative to particle dark matter Federico Lelli INAF - Arcetri Astrophysical Observatory



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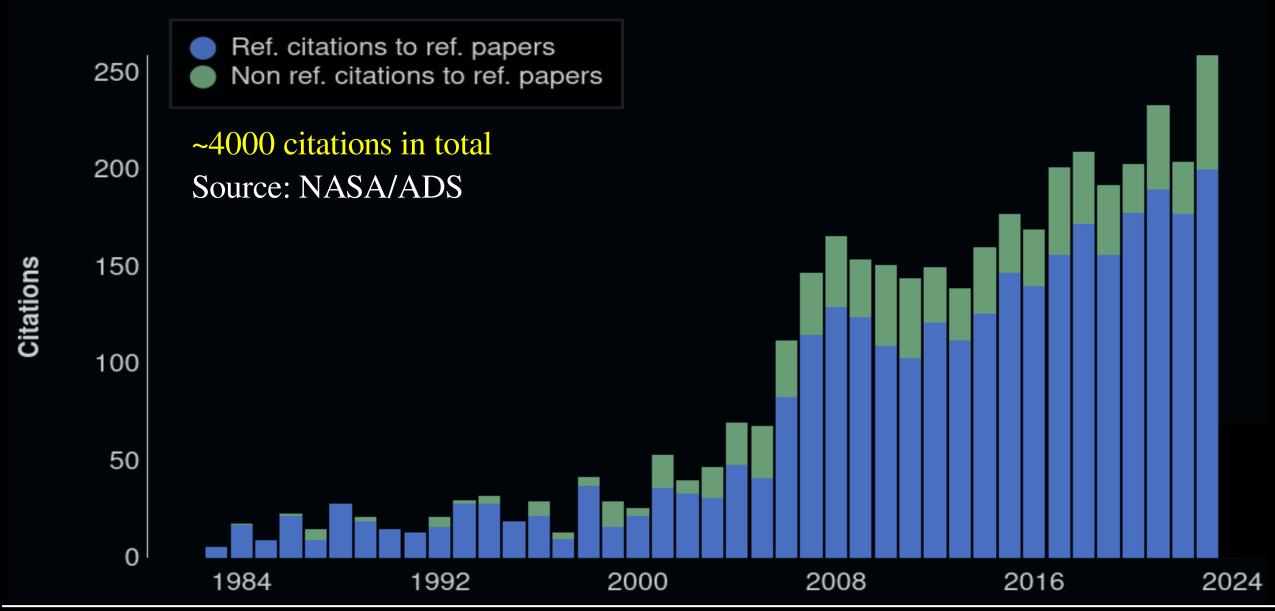


Proposed by Moderhai Milgrom (1983a, b, c)

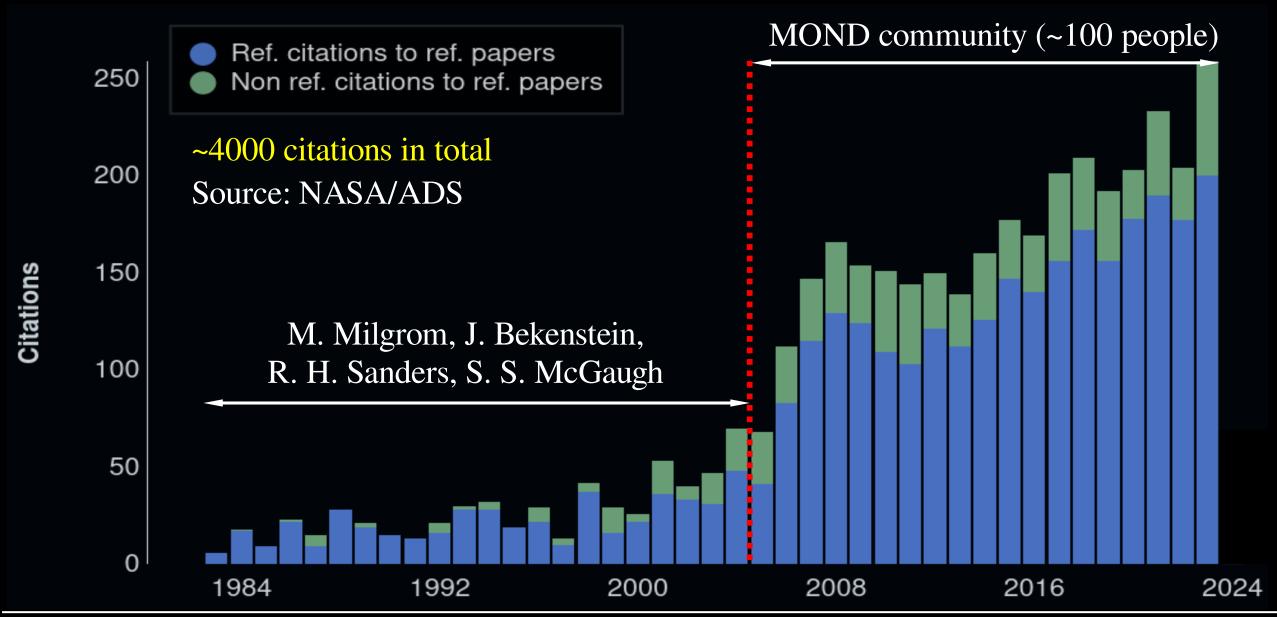
MOND still stands after more than 40 years.

Multifaceted paradigm including different theories at both the non-relativistic and relativistic levels.

Citations to the original MOND trilogy (Milgrom 1983a, b, c)



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similar role as c in Relativity and h in Quantum Mechanics

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 - $\vec{q}_N = -\vec{\nabla} \phi_N$ Newtonian gravitational field (from the Poisson's equation)

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$$V_{N}$$
 V_{N} V_{N}

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 Flat rotation curve at large R!

Or impose scale invariance (Milgrom 2009, ApJ): $(\vec{x}, t) \rightarrow (\lambda \vec{x}, \lambda t)$ V is invariant!

A MODIFICATION OF THE NEWTONIAN DYNAMICS: IMPLICATIONS FOR GALAXIES¹

M. MILGROM

Department of Physics, Weizmann Institute, Rehovot, Israel; and The Institute for Advanced Study Received 1982 February 4; accepted 1982 December 28

ABSTRACT

I use a modified form of the Newtonian dynamics (inertia and/or gravity) to describe the motion of bodies in the gravitational fields of galaxies, assuming that galaxies contain no hidden mass, with the following main results.

- 1) The Keplerian, circular velocity around a finite galaxy becomes independent of r at large radii, thus resulting in asymptotically flat velocity curves.
- 2 The asymptotic circular velocity (V_{∞}) is determined only by the total mass of the galaxy (M): $V_{\infty}^4 = a_0 GM$, where a_0 is an acceleration constant appearing in the modified dynamics. This relation is consistent with the observed Tully-Fisher relation if one uses a luminosity parameter which is proportional to the observable mass.
- (3) The discrepancy between the dynamically determined Oort density in the solar neighborhood and the density of observed matter disappears.
- 4. The rotation curve of a galaxy can remain flat down to very small radii, as observed, only if the galaxy's average surface density Σ falls in some narrow range of values which agrees with the Fish and Freeman laws. For smaller values of Σ , the velocity rises more slowly to the asymptotic value.
- 5. The value of the acceleration constant, a_0 , determined in a few independent ways is approximately $2 \times 10^{-8} (H_0/50 \text{ km s}^{-1} \text{ Mpc}^{-1})^2 \text{ cm s}^{-2}$, which is of the order of $CH_0 = 5 \times 10^{-8} (H_0/50 \text{ km s}^{-1} \text{ Mpc}^{-1}) \text{ cm s}^{-2}$.

The main predictions are:

- 1. Rotation curves calculated on the basis of the *observed* mass distribution and the modified dynamics should agree with the observed velocity curves.
 - 2. The $V_{\infty}^4 = a_0 GM$ relation should hold exactly.
- \bigcirc An analog of the Oort discrepancy should exist in all galaxies and become more severe with increasing r in a predictable way.

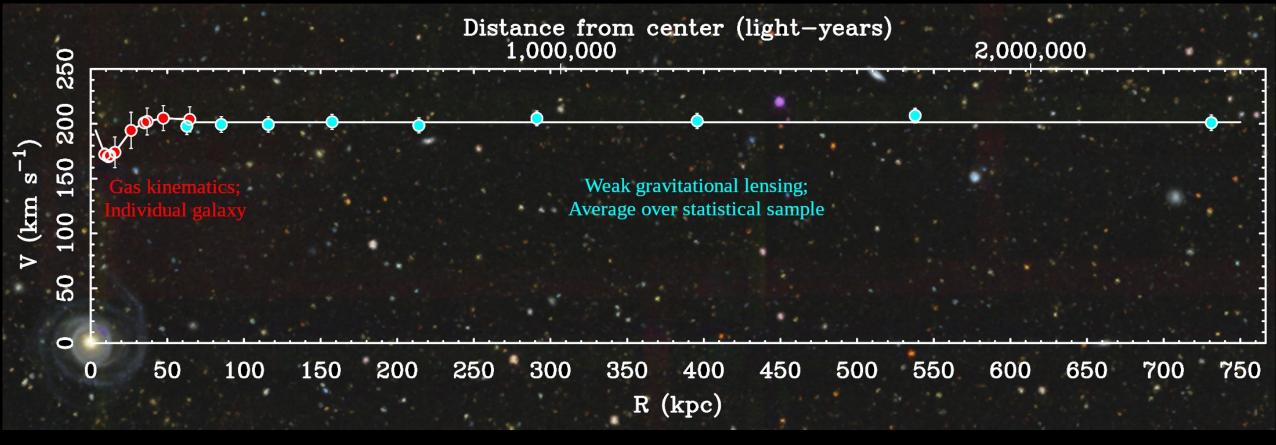


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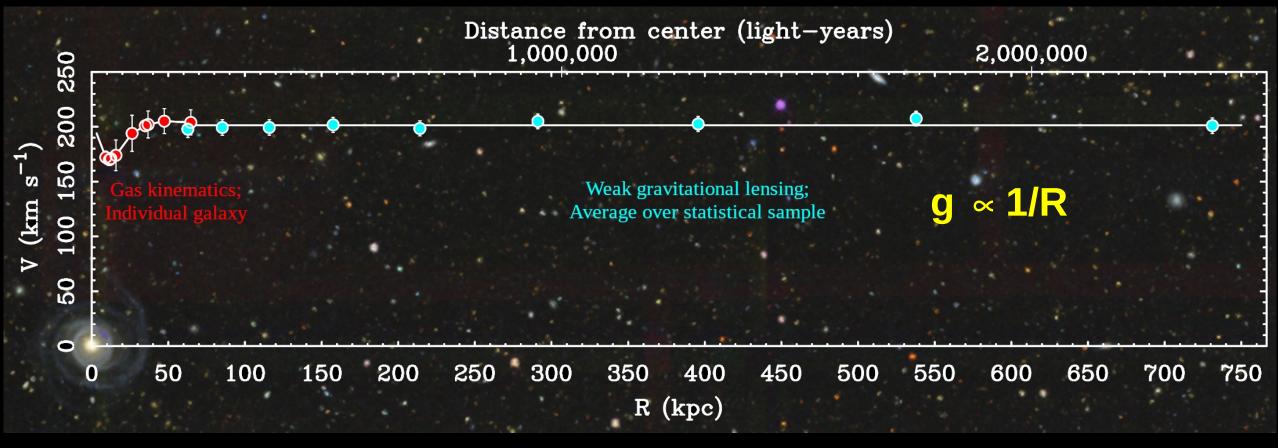
(1) $V_c = const$ at large R for isolated objects

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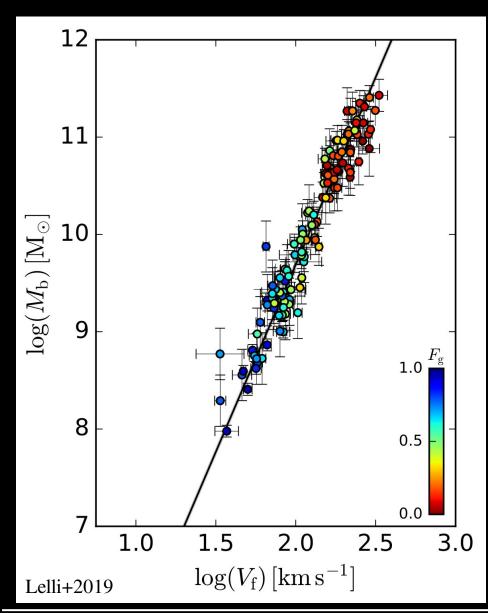
See Mistele+2024, ApJ

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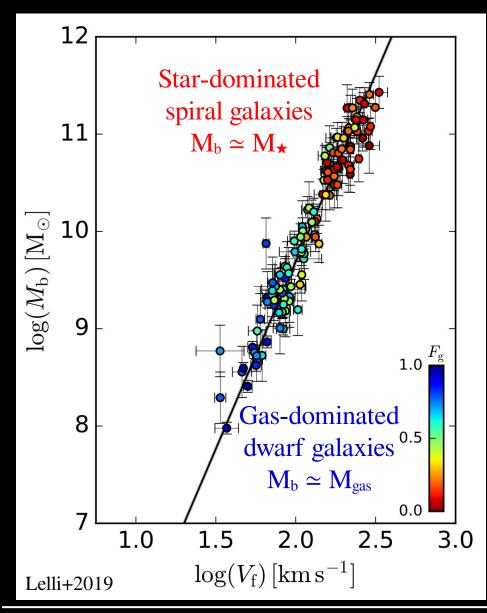
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(2) $V_f^4 = a_0 G M_b$ for isolated objects



Tully-Fisher relation (1977, A&A):

HI linewidth vs Luminosity for bright spirals

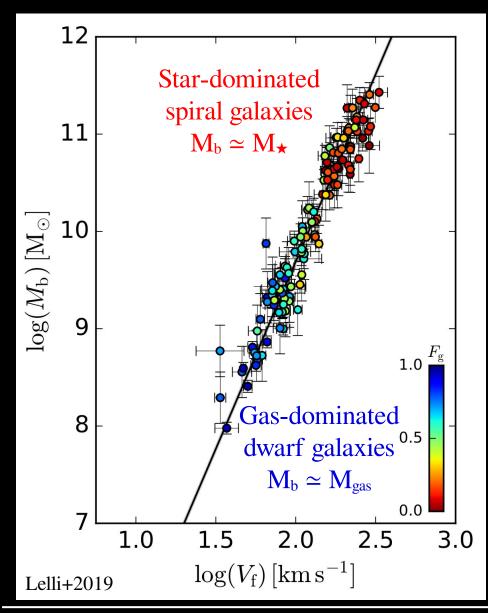


Tully-Fisher relation (1977, A&A):

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Four independent predictions in one equation:

(i) Key quantities are V_f and M_b (stars+gas) $\rightarrow OK$

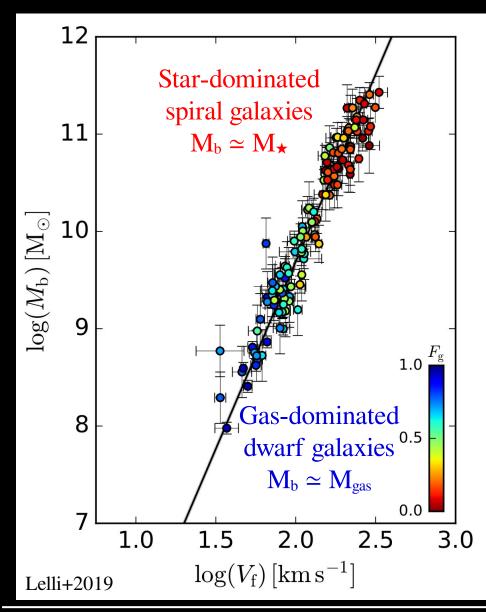


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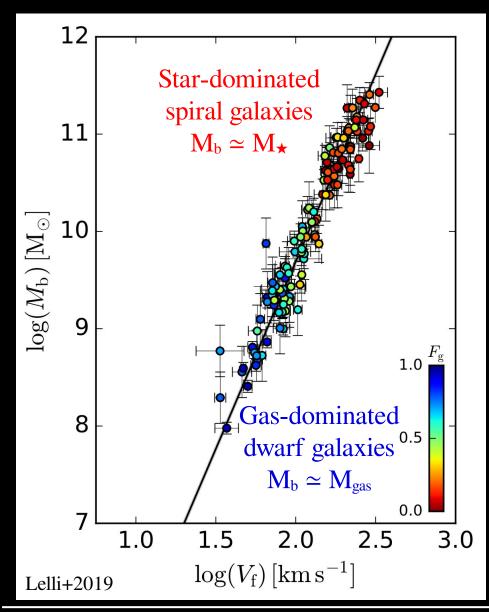


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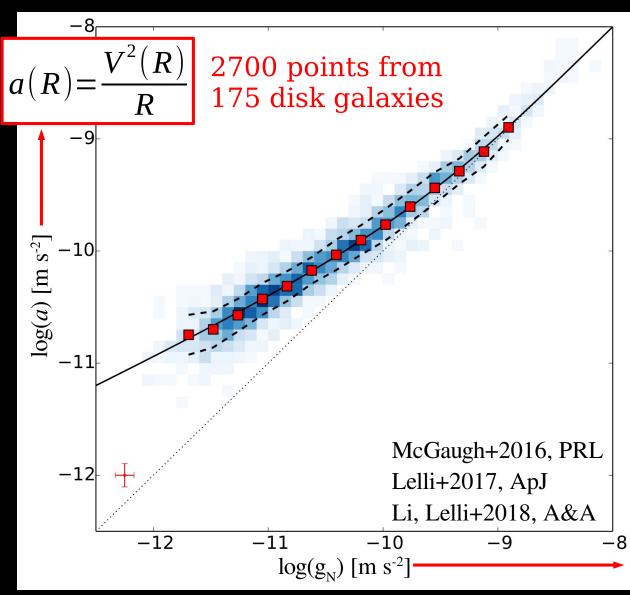


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- (iv) NO dependence on other properties (e.g., R_e , Σ_e), so NO intrinsic scatter along the relation \rightarrow OK

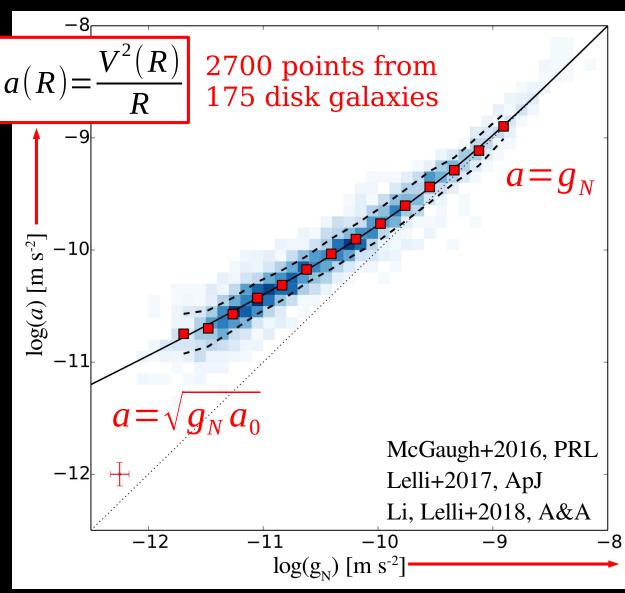


Radial Acceleration Relation (RAR)

(i) Fully empirical, independent of MOND

$$\nabla^2 \Phi_N(R,z) = 4\pi G \rho_b(R,z)$$

$$g_N(R,z=0) = -\nabla \Phi_N(R,z=0)$$

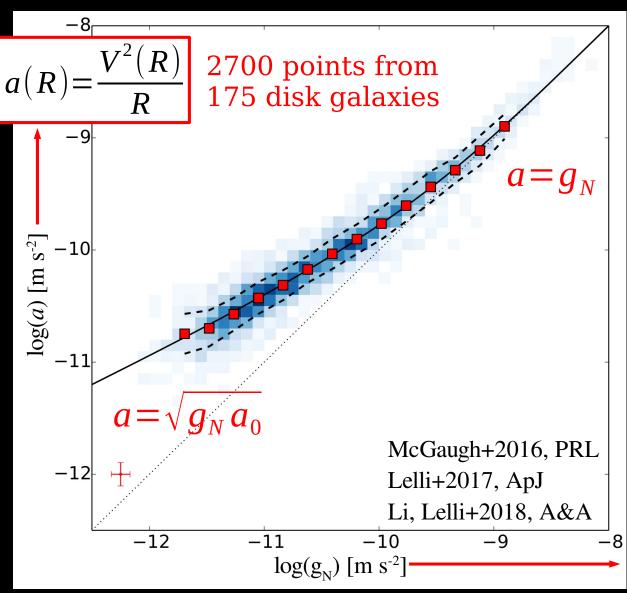


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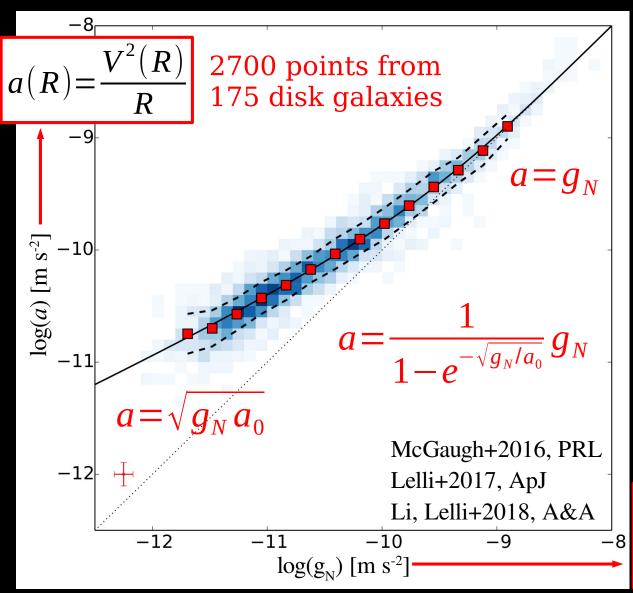


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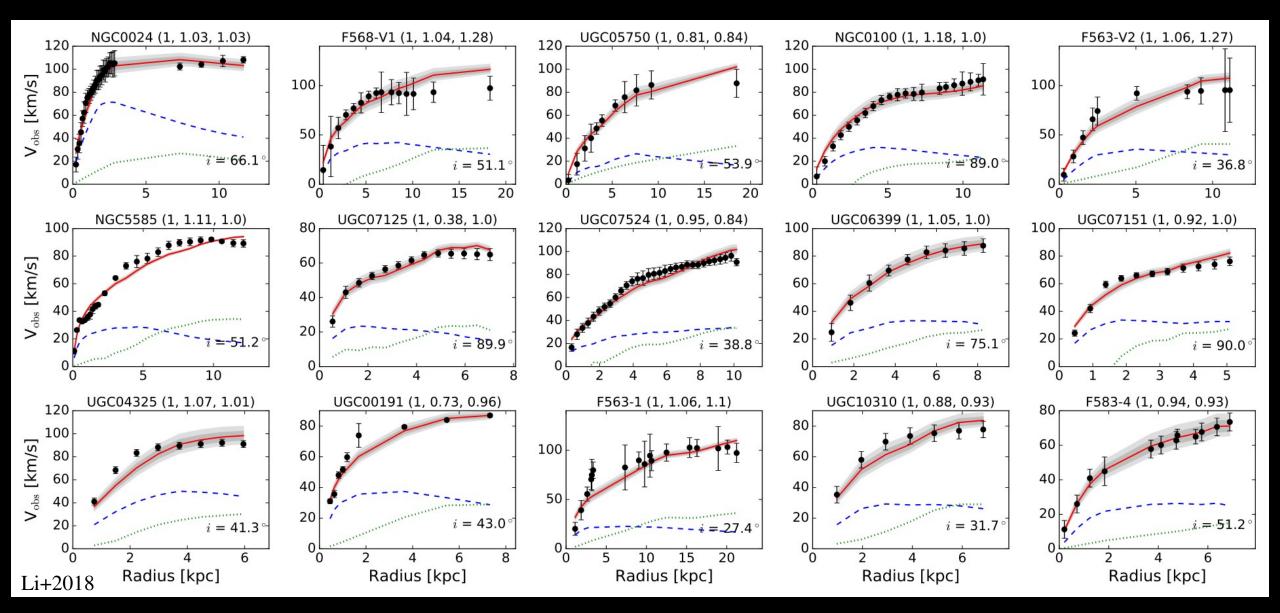
- (i) Fully empirical, independent of MOND
- (ii) MOND asymptotic limits \rightarrow OK
- (iii) No other dependencies $\rightarrow \overline{OK}$
- (iv) MOND interpolation function μ or ν

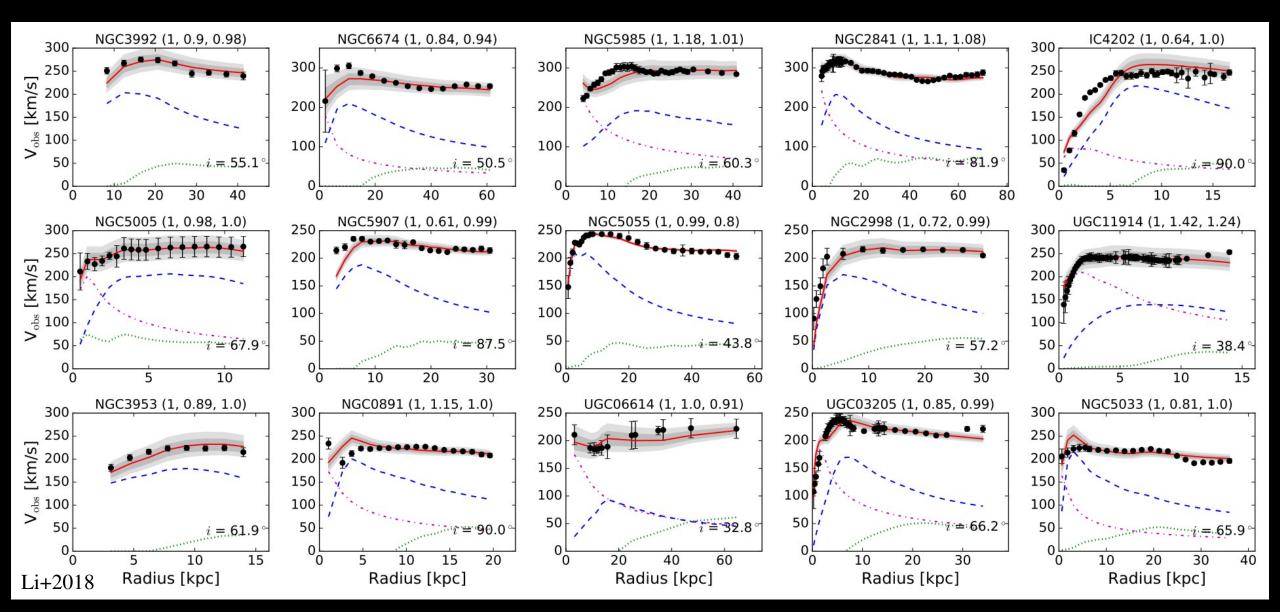
$$a\mu(\frac{a}{a_0})=g_N \iff a=\nu(\frac{g_N}{a_0})g_N$$

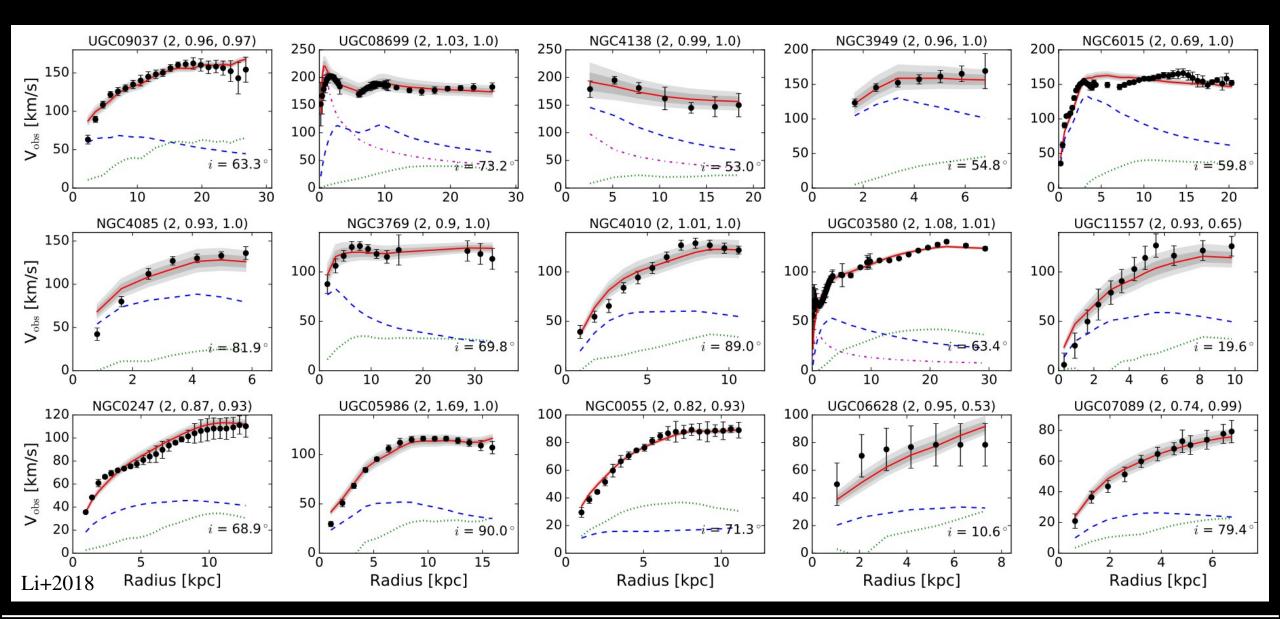
We can now assume $v(g_N/a_0)$ and predict rotation curves given ρ_b (within the errors)

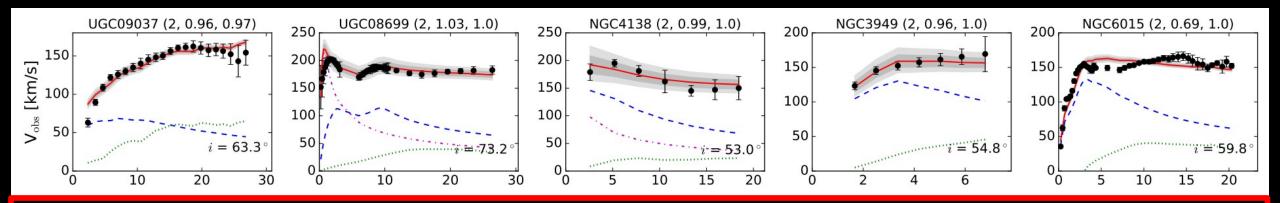
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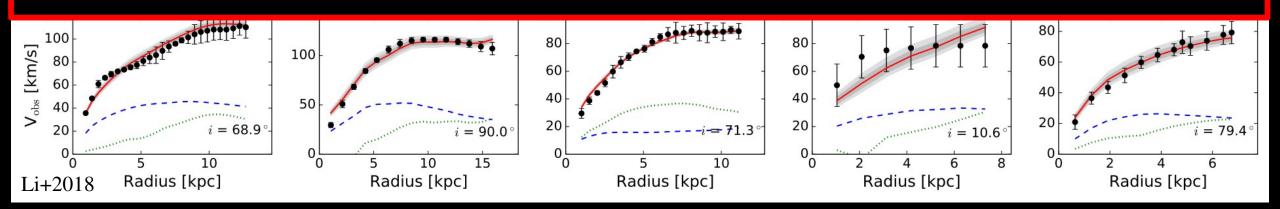








For the full sample of 175 galaxies, see Li, Lelli, McGaugh et al. (2018) or the SPARC database (http://astroweb.cwru.edu/SPARC/)



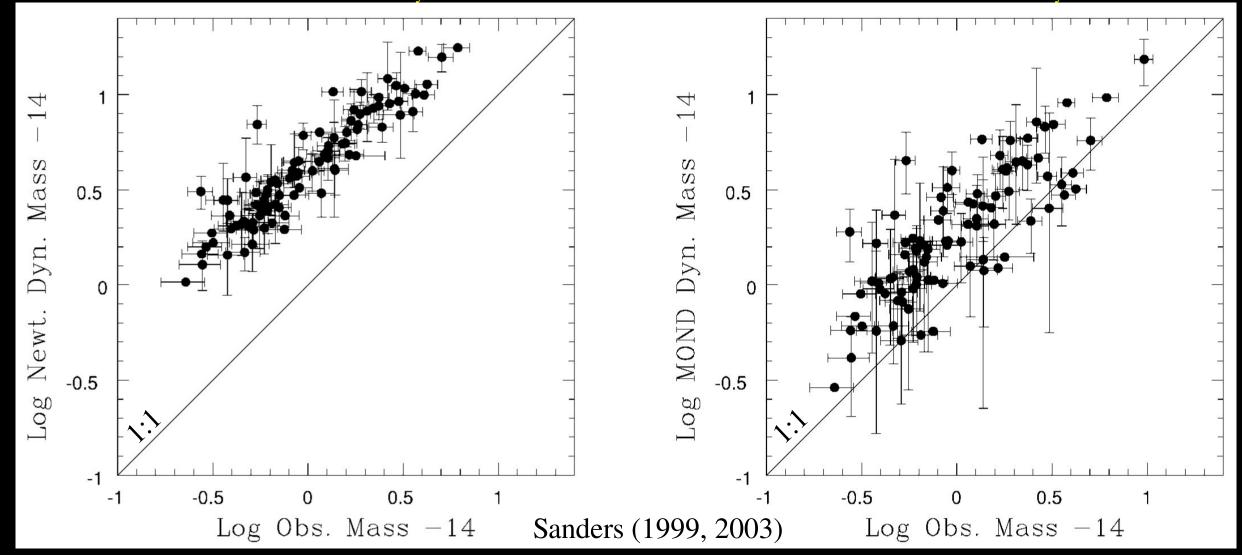
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Galaxy Clusters: Long-standing problem for MOND

Newtonian analysis: $M_{dvn}/M_{bar} \approx 5$

MOND analysis: $M_{dyn}/M_{bar} \simeq 2$



Nature of the missing mass?

• Dark baryons such as compact clouds of cold gas (Milgrom 2008)

For HI clouds $(T\sim10^4 \text{ K}) \rightarrow \text{M}_{cl} < 10^5 \text{ M}_{\odot} \text{ and } \text{R}_{cl} < 50 \text{ pc}$ (Kelleher & Lelli 2024)

Below current HI detection limits; SKA may detect them.

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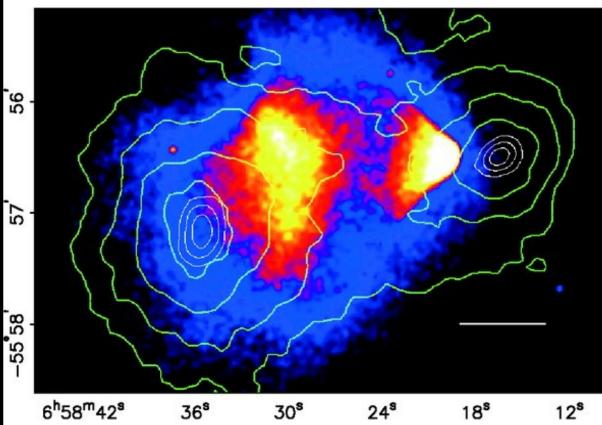
(Kelleher & Lelli 2024)

Below current HI detection limits; SKA may detect them.

• Sterile neutrinos of ~10 eV (Angus 2008, 2009) but CMB power spectrum may be problematic (Thomas+2016, Kopp+2018, Ilic+2021)

Bullet cluster in MOND: missing mass must be collisionless

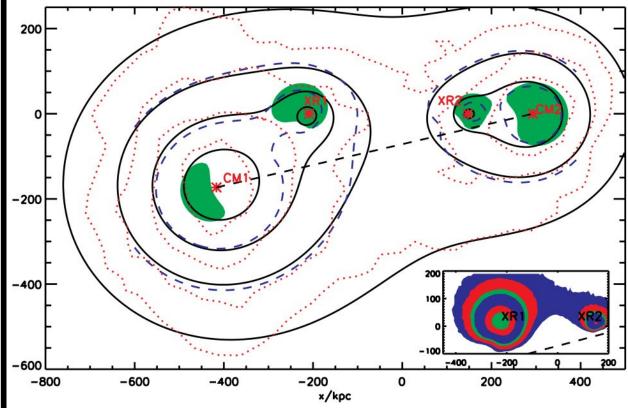




Green: Observed lensing map (total mass)

Blue/Red/Yellow: X-ray emission (hot gas)

MOND (Angus+2006, MNRAS; Angus+2007, ApJ)



Red: Observed lensing convergence map

Black: TeVeS model with 2eV neutrinos

Blue: total surface densities (baryons+υ)

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Cosmology with modified Newtonian dynamics (MOND)

R. H. Sanders

Kapteyn Astronomical Institute, Groningen, The Netherlands

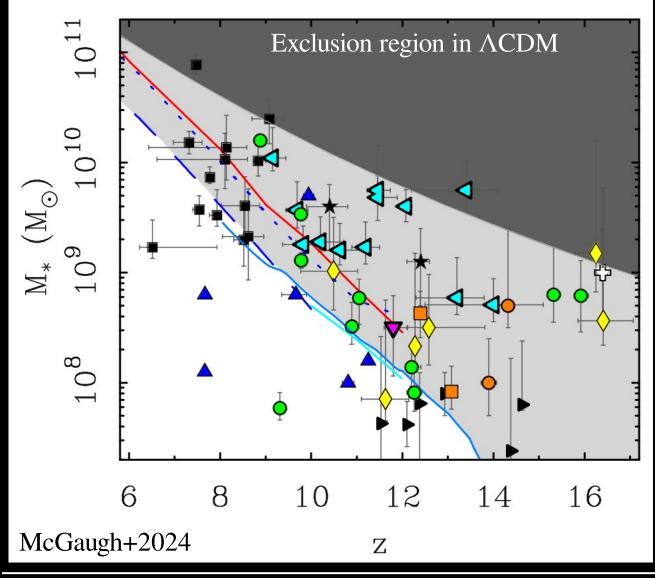
Accepted 1998 January 13. Received 1997 December 22; in original form 1997 October 17

ABSTRACT

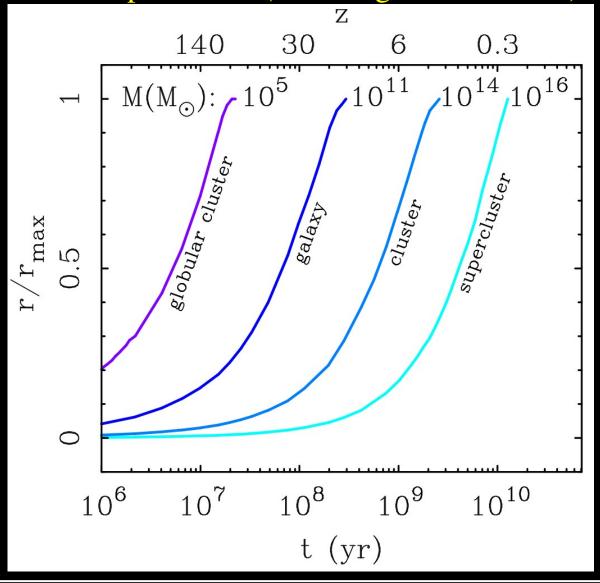
It is well known that the application of Newtonian dynamics to an expanding spherical region leads to the correct relativistic expression (the Friedmann equation) for the evolution of the cosmic scalefactor. Here, the cosmological implications of Milgrom's modified Newtonian dynamics (MOND) are considered by means of a similar procedure. Earlier work by Felten demonstrated that in a region dominated by modified dynamics the expansion cannot be uniform (separations cannot be expressed in terms of a scalefactor) and that any such region will eventually recollapse regardless of the initial expansion velocity and mean density. Here I show that, because of the acceleration threshold for the MOND phenomenology, a region dominated by MOND will have a finite size which, in the earlier Universe (z > 3), is smaller than the horizon scale. Therefore, uniform expansion and homogeneity on the horizon scale are consistent with MONDdominated non-uniform expansion and the development of inhomogeneities on smaller scales. In the radiation-dominated era, the amplitude of MOND-induced inhomogeneities is much smaller than that implied by observations of the cosmic background radiation, and the thermal and dynamical history of the Universe is identical to that of the standard big bang model. In particular, the standard results for primordial nucleosynthesis are retained. When matter first dominates the energy density of the Universe, the cosmology diverges from that of the standard model. Objects of galaxy mass are the first virialized objects to form (by z = 10), and larger structure develops rapidly. At present, the Universe would be inhomogeneous out to a substantial fraction of the Hubble radius.

Massive galaxies at z > 10 is the new normal!

Compilation of JWST observations



MOND prediction (following Sanders 1998)



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Modified Gravity (
$$\rightarrow \nabla^2 \Phi = 4\pi G \rho$$
)

Modified Inertia (\rightarrow F = ma)

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Non-relativistic Lagrangian theories:

AQUAL (Bekenstein & Milgrom 1984)

QUMOND (Milgrom 2010, 2023)

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Heuristic ideas:

Mach's principle for inertia? $a_0 \simeq c \Lambda^{1/2} \rightarrow \text{quantum vacuum? (Milgrom 1999)}$

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MOND-like DM models:

Bipolar DM (Blanchet+2008, 2009, 2015, 2017)

Superfluid DM (Berezhiani & Khoury 2015)

Baryon-interacting DM (Famaey+2018, +2020)

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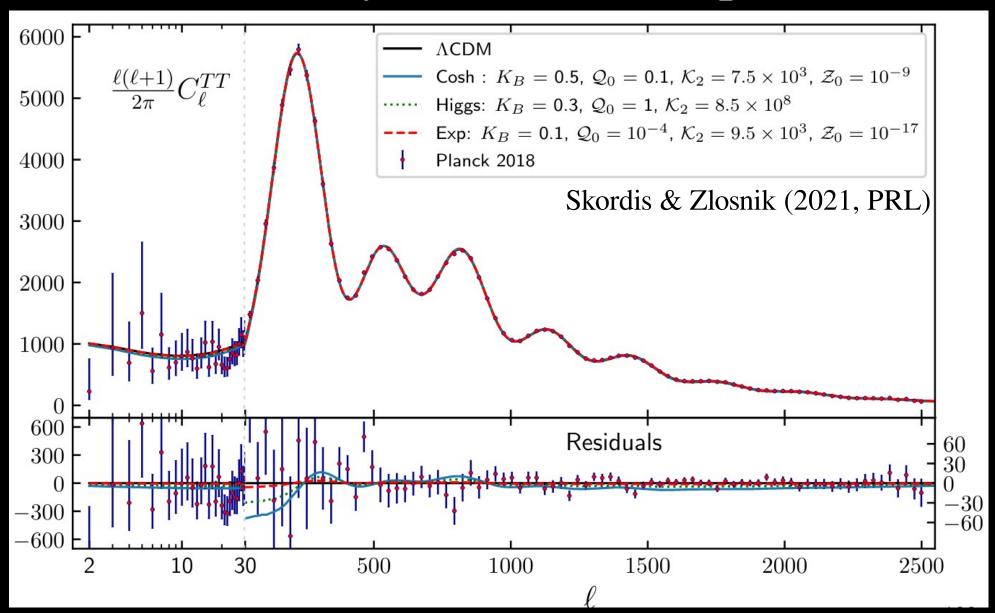
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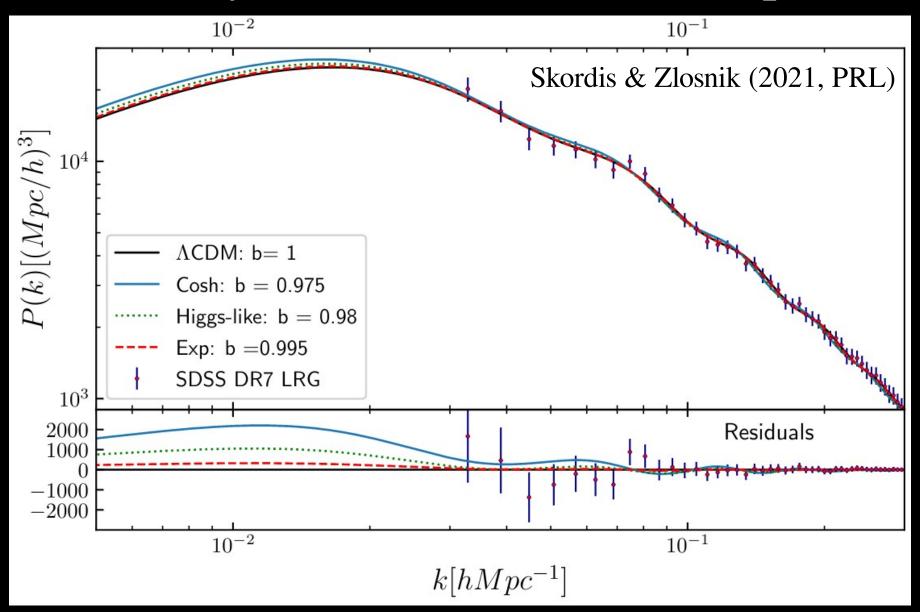
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AeST theory: CMB Power Spectrum



AeST theory: Linear Matter Power Spectrum



Why taking MOND seriously?



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• MOND made multiple *a-priori* predictions, later confirmed by observations. A-priori predictions are the core of the scientific method.

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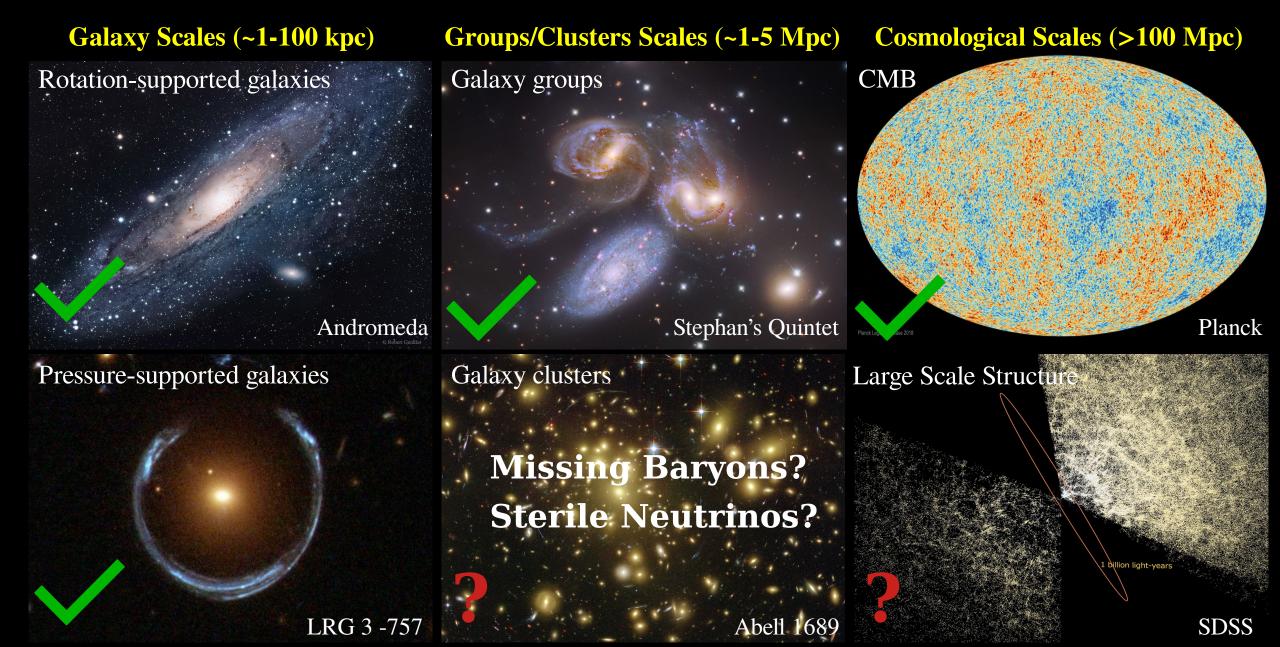
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- Even if MOND is wrong (or DM is detected), there must be a baryon DM coupling in galaxies acting in a MOND way. It tells us on the nature of DM.





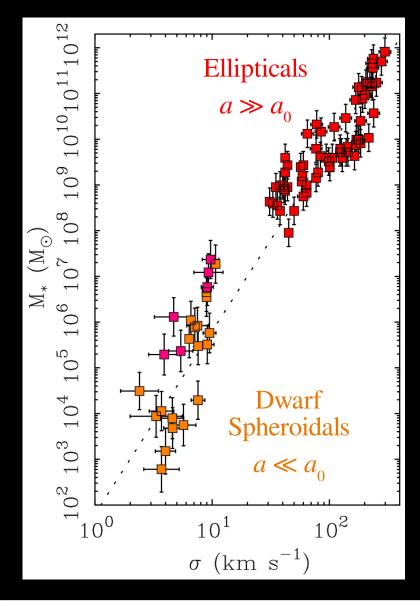
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- Even if MOND is wrong (or DM is detected), there must be a baryon DM coupling in galaxies acting in a MOND way. It tells us on the nature of DM.
- Over the past ~10 years, substantial progress in MOND research in both observations and theory (new tests, new ideas, new relativistic theories).

Summary: Status of MOND at Various Scales



More Slides

(4) $\sigma_{\rm V}^{4} \simeq a_0^{\rm G} \, {\rm M}_{\rm b}^{\rm f}$ for quasi-isothermal systems \rightarrow pressure-supported gals



Faber-Jackson relation (1976, ApJ) for ellipticals Three a-priori independent predictions in one equation:

- (i) Slope should be exactly $4 \rightarrow OK$
- (ii) Normalization is $a_0G \rightarrow OK$ with BTFR estimate!
- (iii) No dependence on other quantities IF $a \ll a_0 \rightarrow OK$

 σ_{v} is measured at R < R_e (containing half luminosity):

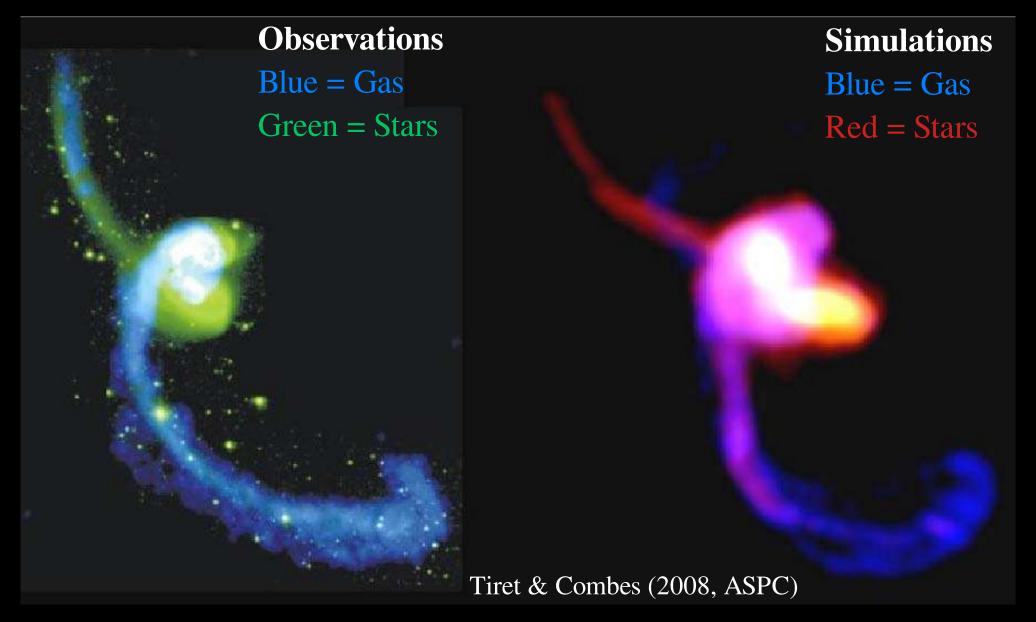
For dwarf spheroidals: $a \ll a_0$ at R<R_e \rightarrow MOND regime

For giant ellipticals: $a \gg a_0$ at R<R_e \rightarrow Newtonian regime

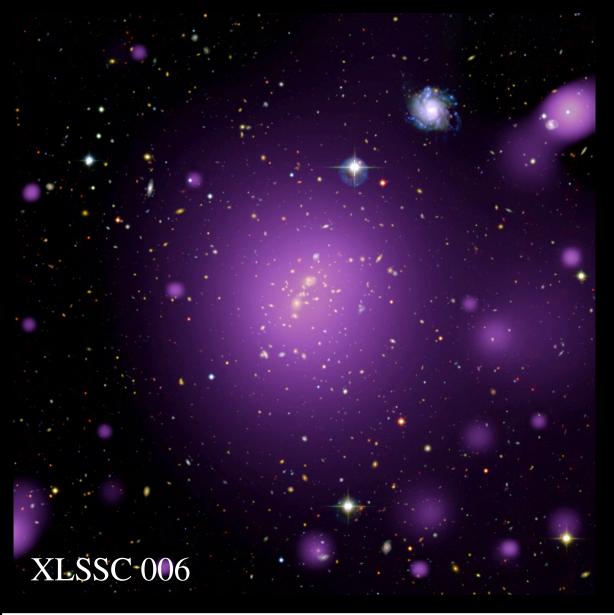
$$\frac{\sigma_V^2}{R} \simeq \frac{GM}{P^2} \longrightarrow M \simeq \sigma_V^2 R_e$$

Fundamental plane of ellipticals (Djorgovski & Davis 1987; Dressler 1987)

Interacting & Merging Galaxies: The Antennae



Galaxy Clusters: bound systems with ~100-1000 galaxies



Observed baryon budget:

~10% galaxies (optical & NIR)

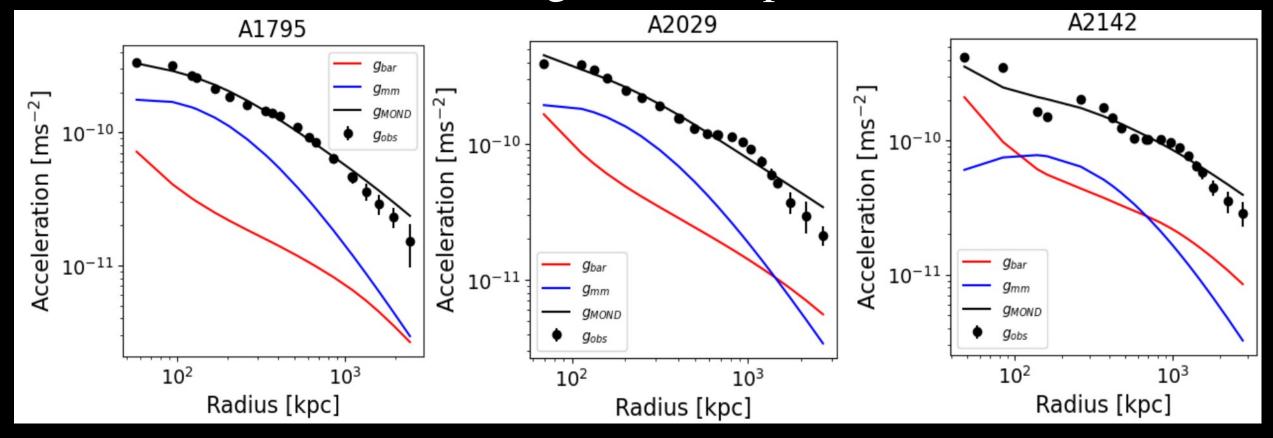
~90% hot ionized gas (X rays)

Sphere in Hydrostatic Equilibrium

$$g_{obs} = \frac{-1}{\rho_{gas}} \frac{\partial P_{gas}}{\partial r} \quad P_{gas} = \frac{k_B \rho_{gas} T_{gas}}{w m_p}$$

$$\Rightarrow g_{obs} = \frac{-k_B}{w \, m_p \, \rho_{gas}} \frac{\partial}{\partial r} \left| \rho_{gas} \, T_{gas} \right|$$

MOND fits with a missing mass component (Kelleher & Lelli 2024)



Density profile of missing mass:

$$\rho_{mm}(r) = \frac{\rho_0}{(1+r/r)^4}$$

Converged (physical) total mass:

$$\longrightarrow M_{mm,tot} = \int_{0}^{\infty} \rho(r) 4\pi r^{2} dr = \frac{4\pi}{3} \rho_{0} r_{s}^{3}$$

MOND – Cosmology Connection?

Two numerical coincidences (Milgrom 1983a, ApJ; Milgrom 1999, PhLA):

$$a_0 \simeq \frac{H_0 \cdot c}{2\pi}$$

 H_0 = Hubble constant \rightarrow maybe $a_0(t) \sim H(t)$?

$$a_0 \simeq \frac{c^2 \sqrt{\Lambda/3}}{2\pi}$$

 $a_0 \simeq \frac{c^2 \sqrt{\Lambda/3}}{2 \pi}$ $\Lambda = \text{Cosmological constant} \rightarrow \text{relation to Dark Energy?}$

IF this numerology has some deeper, fundamental meaning:

Either the state of the Universe at large enters in local dynamics, or

the same parameters enters both Cosmology (Λ) and local dynamics (a_0).

AQUAL = Aquadratic Lagrangian (Bekenstein & Milgrom 1984, ApJ)

$$S_{N} = \int dt \, L_{N} = \int dt \, d^{3}x \left| \rho \frac{V^{2}}{2} - \frac{\left| \overrightarrow{\nabla} \Phi \right|^{2}}{8\pi G} - \rho \Phi \right| \quad \text{Lagrangian is quadratic in } \nabla \Phi$$

$$-\frac{a_{0}^{2}}{8\pi G} F \left| \frac{\left| \overrightarrow{\nabla} \Phi \right|^{2}}{a_{0}^{2}} \right| \quad F(z) \to z \text{ for } z = |\nabla \Phi|^{2}/a_{0}^{2} \gg 1$$

$$F(z) \to z^{3/2} \text{ for } z = |\nabla \Phi|^{2}/a_{0}^{2} \ll 1$$

$$\frac{\delta S}{\delta \Phi} = 0 \rightarrow \nabla \cdot \left| \mu \left| \frac{|\vec{\nabla} \Phi|}{a_0} \right| \vec{\nabla} \Phi \right| = 4 \pi G \rho \quad \text{Modified Poisson's Equation}$$

$$\mu(x) = \frac{dF(z)}{dz}$$
 $z = x^2$ $F(z)$ provides the interpolation function $\mu = v^{-1}$

QUMOND = Quasi-Linear MOND (Milgrom 2010, MNRAS)

$$S_N = \int dt \, L_N = \int dt \, d^3x \left| \rho \frac{V^2}{2} - \frac{|\nabla \Phi|^2}{8\pi G} - \rho \Phi \right| \quad \text{Single gravitational potential } \Phi$$

$$\frac{-1}{8\pi G} \left| 2 \overrightarrow{\nabla} \Phi \cdot \overrightarrow{\nabla} \Phi_N - a_0^2 Q \left| \frac{\left| \overrightarrow{\nabla} \Phi_N \right|^2}{a_0^2} \right| \right|$$
 Two potentials: Φ and Φ_N !

QUMOND = Quasi-Linear MOND (Milgrom 2010, MNRAS)

$$S_N = \int dt \, L_N = \int dt \, d^3 x \left| \rho \frac{V^2}{2} - \frac{\left| \nabla \Phi \right|^2}{8\pi G} - \rho \Phi \right|$$
 Single gravitational potential Φ

$$\frac{-1}{8\pi G} \left[2 \overrightarrow{\nabla} \Phi \cdot \overrightarrow{\nabla} \Phi_N - a_0^2 Q \left| \frac{|\overrightarrow{\nabla} \Phi_N|^2}{a_0^2} \right| \right]$$
 Two potentials: Φ and Φ_N !

Principle of least Action varying Φ , Φ_N and $\overline{x} \to \text{set of 3 equations}$

$$\nabla^2 \Phi_N = 4 \pi G \rho$$
 — Standard, linear Poisson's equation for Φ_N

$$\nabla^2 \Phi = \overrightarrow{\nabla} \cdot \left[\mathbf{v} \left| |\overrightarrow{\nabla} \Phi_N| / a_0 \right| \overrightarrow{\nabla} \Phi_N \right] \longrightarrow \text{Non-linear step: get } \Phi \text{ from } \Phi_N \quad \mathbf{v}(\sqrt{x}) = \frac{d \, Q(x)}{x}$$

$$\vec{a} = -\vec{\nabla} \Phi$$
 — Acceleration/force set by second potential Φ

Relativistic MOND → add degrees of fredoom (new fields) to GR

- Tensor $g_{\mu\nu}$ \rightarrow Einstein's metric
- Scalar ϕ \rightarrow for the DM effect in the NR limit (Bekenstein & Milgrom 1984, ApJ)
- Vector A^µ → for gravitational lensing (Sanders 1997, ApJ; Bekenstein 2004, PRD)
- Free function(s) \rightarrow interpolation function(s) (non-fundamental effective theories?)

AeST: Aether Scalar-Tensor theory (Skordis & Zlosnik 2019, PRD; 2021, PRL; 2022, PRD)

- (1) Grav. Waves: $B^{\mu} = e^{-2\phi}A^{\mu}$ (timelike) \rightarrow theories $\{B_{\mu}, g_{\mu\nu}\} \rightarrow c_{GW} = c_{EM}$
- (2) Cosmology: k-essence term for $\phi \to \rho_{\phi} \propto a(t)^{-3}$ (like dust or CDM)

MOND as Modified Inertia (Milgrom 1994, 1999, 2022)

$$\overrightarrow{A}[\overrightarrow{x}(t); a_0] = -\overrightarrow{\nabla} \Phi_N$$
 \overline{A} is a functional of the full trajectory $\overline{x}(t)$
For $a \gg a_0$, $A \to a = d^2x/dt^2$ (Newton's 2nd law)

No full theory yet, but two general results:

(1) Imposing Newtonian and MOND limits + Galilei Invariance $\vec{x}(t) \rightarrow \vec{x}(t) + \vec{v_0}t$

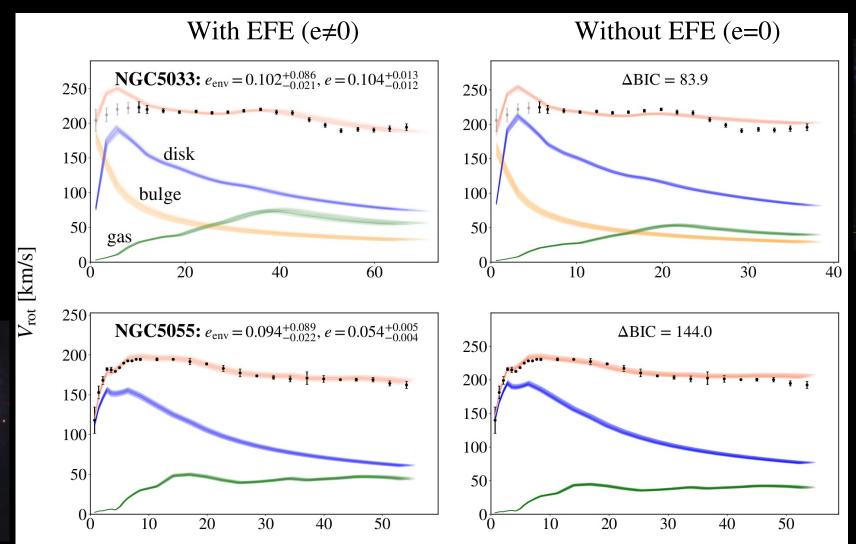
Theory is time non-local:
$$\vec{A}[\vec{x}(t), a_0] \neq F(\frac{d^i \vec{x}}{dt^i}; i=1, 2, ...N)$$

Accelerations at (\bar{x}, t) depend on the full orbital history!

(2) For purely circular orbits: $\vec{a}\mu(\frac{a}{a_0}) = \vec{g}_N$ holds exactly (RAR for disk galaxies)

The interpolation function is a derived concept valid for circular orbits!

EFE is weak: individual detections only in extreme cases

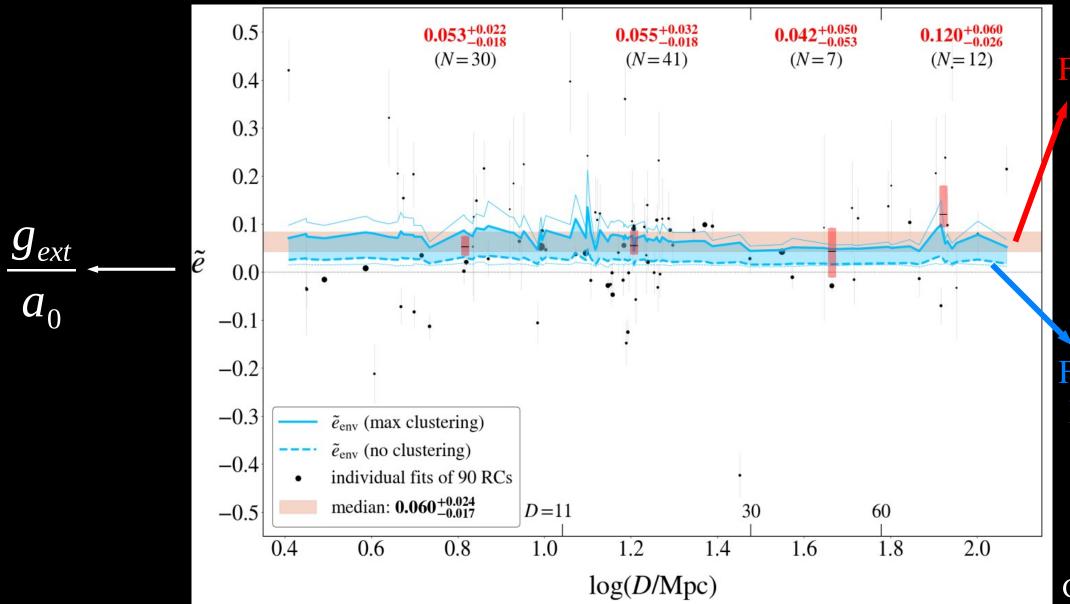


Chae+2020, 2021

NGC 5033

NGC 5055

Statistical approach: EFE>0 at >4\sigma and agrees with LSS



From Rotation

Curve Fits

From Baryon Large Scale Structure

Chae+2020, 2021

Galaxy Clusters on the Radial Acceleration Relation

