## Manifesting hidden dynamics of a sub-component dark matter

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|--|---|---|
| Introduction   | Refe  | rence model   |
| n multi-component dark matter (DM) scenarios, less attention<br>been given to sub-dominant components of DM and the<br>corresponding impact on their detectability. They are often t<br>to be hidden from observations in direct and indirect detection<br>DM due to its small fraction.<br>However, a sub-dominant DM component can play a domina<br>in the dynamics of and detection of dark sector. The strategy<br>probe a sub-dominant DM component relies on its cosmolog<br>evolution which is sensitive to the interaction within a dark s<br>Here, we show a case when the dynamics within a dark sector<br>affect the detectability of a sub-component DM. | Image: Solution of the sector.For concreteness, we take scenario:Solution of the sector.1. Minimally two DM componentsImage: Solution of the sector.1. The light boosted DM graph with large kinetic energy   | a multi-component boosted DM<br>ponents: $\chi_0$ (heavy) and $\chi_1$ (light)<br>ent is the heavy DM $\chi_0$ which does not<br>s with the SM particles. Its relic<br>ed by the freeze-out of $\chi_0\chi_0 \leftrightarrow \chi_1\chi_1$<br>freeze-out mechanism.<br>M freezes out later lowering the<br>ared to $\chi_0$ .<br>$\chi_1$ can be produced via $\chi_0\chi_0 \rightarrow \chi_1\chi_1$<br>gy about $m_{\chi_0}$ .  |
|  | Results   |   |
| After the freeze-out of $\chi_0$ , the Boltzmann equation is<br>$\frac{dY_{\chi_1}}{dx} \simeq -\frac{\lambda_{\chi_1}(x)}{x} \left[ Y_{\chi_1}^2 - (Y_{\chi_1}^{eq}(x))^2 - Y_{ast.}^2(x) \right]$ $\chi_0 \chi_0 \to \chi_1 \chi_1$ $\chi_0 \chi_0 \to \chi_1 \chi_1$ $\chi_{\chi_i} = s \langle \sigma_i v_{rel} \rangle / H$ $\langle \sigma_0 v_{rel} \rangle = \langle \sigma v \rangle_{\chi_0 \chi_0 \to \chi_1 \chi_1}$ $\langle \sigma_1 v_{rel} \rangle = \langle \sigma v \rangle_{\chi_1 \chi_1 \to SM}$ $Y_{ast.}(x) = \sqrt{\frac{\langle \sigma_0 v_{rel} \rangle}{\langle \sigma_1 v_{rel} \rangle}} Y_{\chi_0}(x)$                                 | <ul> <li>If the ratio r<sub>1</sub> ≡ Ω<sub>χ1</sub>/Ω<sub>DM, tot</sub> = equilibrium with SM longer x<sub>f.o.,χ1</sub> &gt; 20 and hence Y<sup>eq</sup><sub>χ1</sub> </li> <li>Then the production rate of negligible compared to the or Y<sub>χ1</sub> compared to the standa</li> <li>So, for a fixed r<sub>1</sub> ≪ 1, the defective: ⟨σ<sub>1</sub>v<sub>rel</sub>⟩ is proportion p-wave dominant annihilation</li> <li>Then, the cosmo/astro observation observation observation of the standa</li> </ul>  | « 1, the light DM $\chi_1$ is in thermal<br>than the conventional WIMP case, i.e,<br>« $Y_{ast.}$ at the freeze-out of $\chi_1$ .<br>$\chi_1$ from $\chi_0\chi_0 \rightarrow \chi_1\chi_1$ is non-<br>depletion by $\chi_1\chi_1 \rightarrow$ SM, increasing<br>rd WIMP freeze-out case.<br>pletion $\chi_1\chi_1 \rightarrow$ SM should be<br>onal to $1/r_{1^2}$ , $1/r_{1^3}$ for the <i>s</i> -wave and<br>on process, respectively.<br>rvables of $\chi_1$ proportional to $n_{\chi_1}^2 \langle \sigma_1 v r_{1^2} \rangle$ |
| • Long after the freeze-out of $\chi_1\chi_1 \leftrightarrow$ SM, the energetically produced $\chi_1$ from $\chi_0\chi_0 \rightarrow \chi_1\chi_1$ can transfer its energy to relic $\chi_1$ via self-interactions, increasing its kinetic energy. Called the <i>self-heating mechanism</i> , which is important for <i>p</i> -wave dominant $\chi_1\chi_1 \rightarrow$ SM annihilation case.  | Y<br>the The bounds from the observery<br>or γ-rays in our galaxy, DM d<br>those constraining the warm<br>very strongly to the sub-com<br>$((\frac{10^{-1}}{U} + \frac{10^{-2}}{U})^{-1} + \frac{10^{-2}}{U} + 1$ | Pations of BBN, CMB, diffuse X-rays<br>irect detection experiments, and<br>dark matter (for $r_1 \ge 7\%$ ) apply<br>ponent DM $\chi_1$ .   |



## Discussion

- In multi-component DM scenarios, a sub-component DM ( $\chi_1$ ) can severely affect the cosmo/astro observables unlike conventional expectation; due to the strong current bounds the *p*-wave dominant  $\chi_1\chi_1 \rightarrow$  SM annihilation case is preferred.
- Self-heating from  $\chi_0\chi_0 \rightarrow \chi_1\chi_1$  followed by the self-interactions among  $\chi_1$  arises and changes the evolution of  $T_{\chi_1}$  even after its freeze-out. Then, the constraints for warm dark matter can be applied to the *p*-wave dominant  $\chi_1\chi_1 \rightarrow SM$  process for  $r_1 \ge 7\%$ .
- Direct detection constraints may change and complementary searches in accelerators can interplay with cosmo/astro observations to probe the detailed structure of dark sector.



 $10^{1}$ 

 $m_{\chi_1}$  [MeV]

 $m_{\chi_0} = 30 \, \text{MeV}$ 

10

 $10^{-5}$ 

 $10^{-6}$ 

 $10^{0}$ 



- $\langle \sigma_1 v_{rel} \rangle$ : *p*-wave dominant
- Reference model: singlet scalar  $\chi_1$  interacting with SM with a dark photon (A')

 $10^{1}$ 

 $m_{\chi_1}$  [MeV]

 $m_{\chi_0} = 30 \,\mathrm{MeV}$ 

 $\sigma/m = 0.1 \, {\rm cm}^2/$ 

 $10^{2}$ 

- Green: N<sub>eff</sub>, Pink: Warm DM
- A future discovery of  $\chi_1$  (blue star) can be inconsistent with the cosmo/astro bounds.

## References

 $10^{-5}$ 

 $10^{-6}$ 

 $10^{0}$ 

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