



# Introduction to plasma acceleration

Gemma Costa

On behalf of the SPARC\_LAB collaboration



- Motivation to plasma based accelerators
- What is a plasma?
- Plasma based acceleration
- Plasma source
- Beam driven scheme
- What is a high power laser?
- Laser scheme
- Plasma acceleration @LNF: SPARC\_LAB
- EuPRAXIA and EuAPS projects

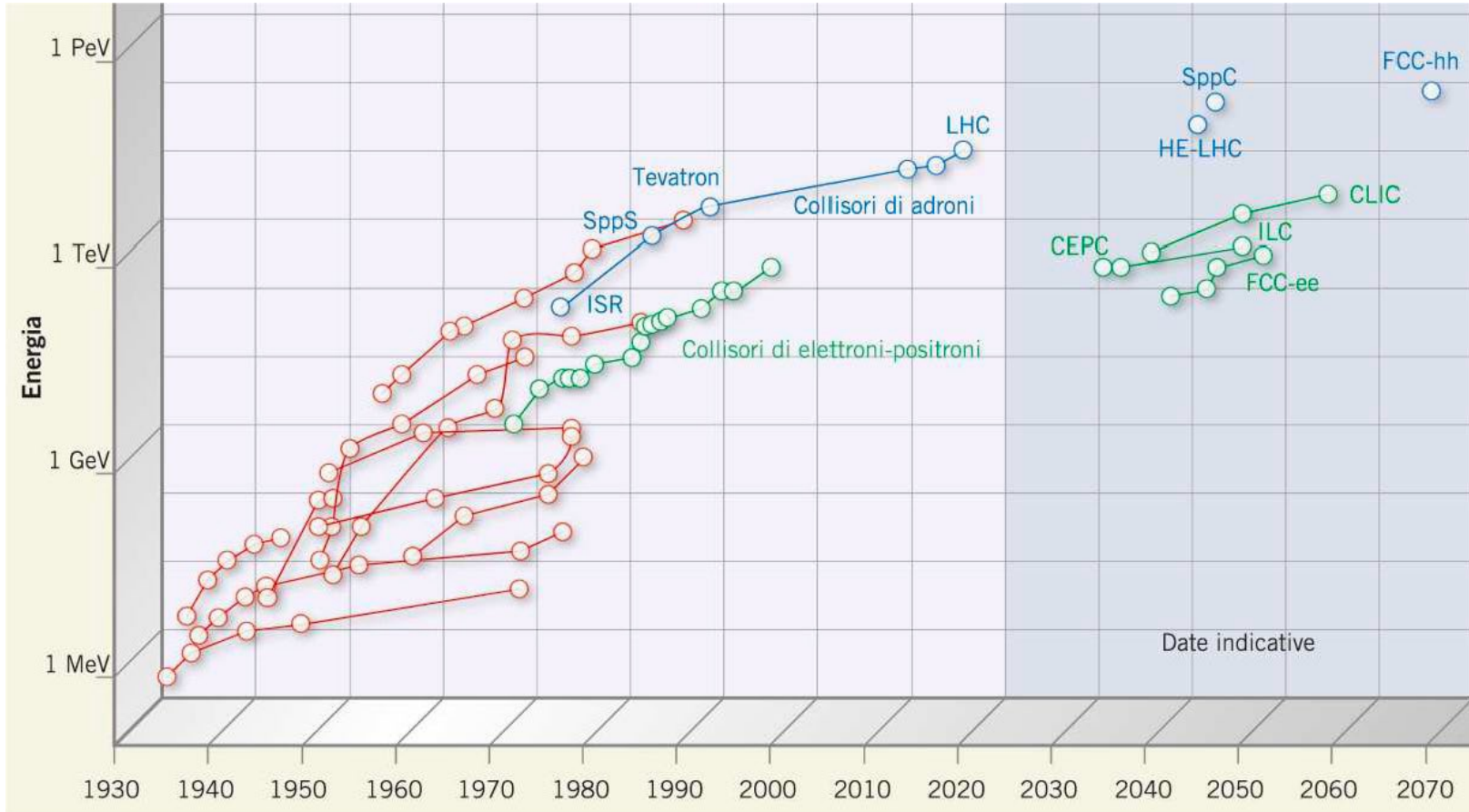
# Motivation to plasma based accelerators

Accelerators		1994	2014
<b>Industrial</b>		<b>&gt;4500</b>	<b>27 000</b>
	Electron accelerators >300 keV	1500	~5000
	Electron accelerators <300 keV	>1000	~8000
	Ion implanters and ion analysis	>2000	~12 000
	Neutron generators		~2000
<b>Science</b>		<b>~1000</b>	<b>~1200</b>
<b>Medicine</b>		<b>~4200</b>	<b>~14 000</b>
	Electron accelerators	~4000	~13 000
	Proton and ion accelerators	17	~59
	Production of radioisotopes	~200	~1100
<b>TOTAL</b>		<b>&gt;9700</b>	<b>42 000</b>

Chernyaev, A. P., and S. M. Varzar. "Particle accelerators in modern world." Physics of Atomic Nuclei 77.10 (2014): 1203-1215.

# Motivation to plasma based accelerators

The accelerating gradients have more or less remained constant over the past few decades, in the order of 20–50 MV/m

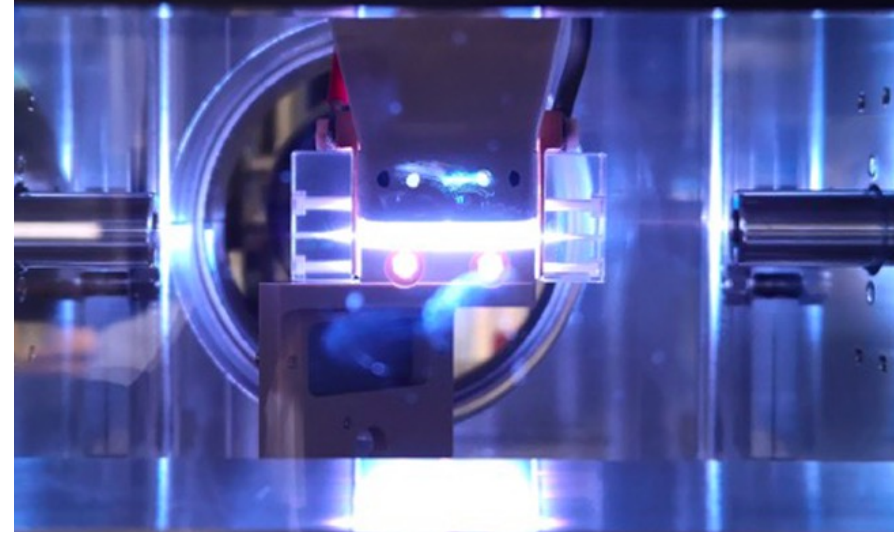


# Motivation to plasma based accelerators



Conventional RF structures:

$(20 - 40) \text{ MV/m}$   
range  $m - km$

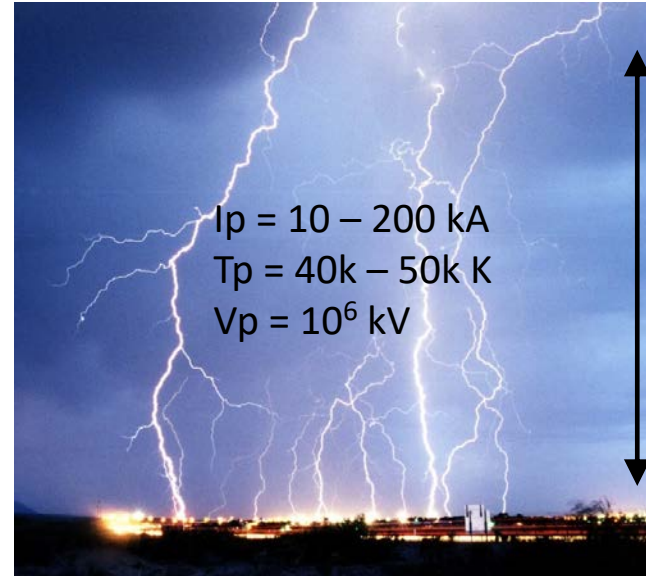
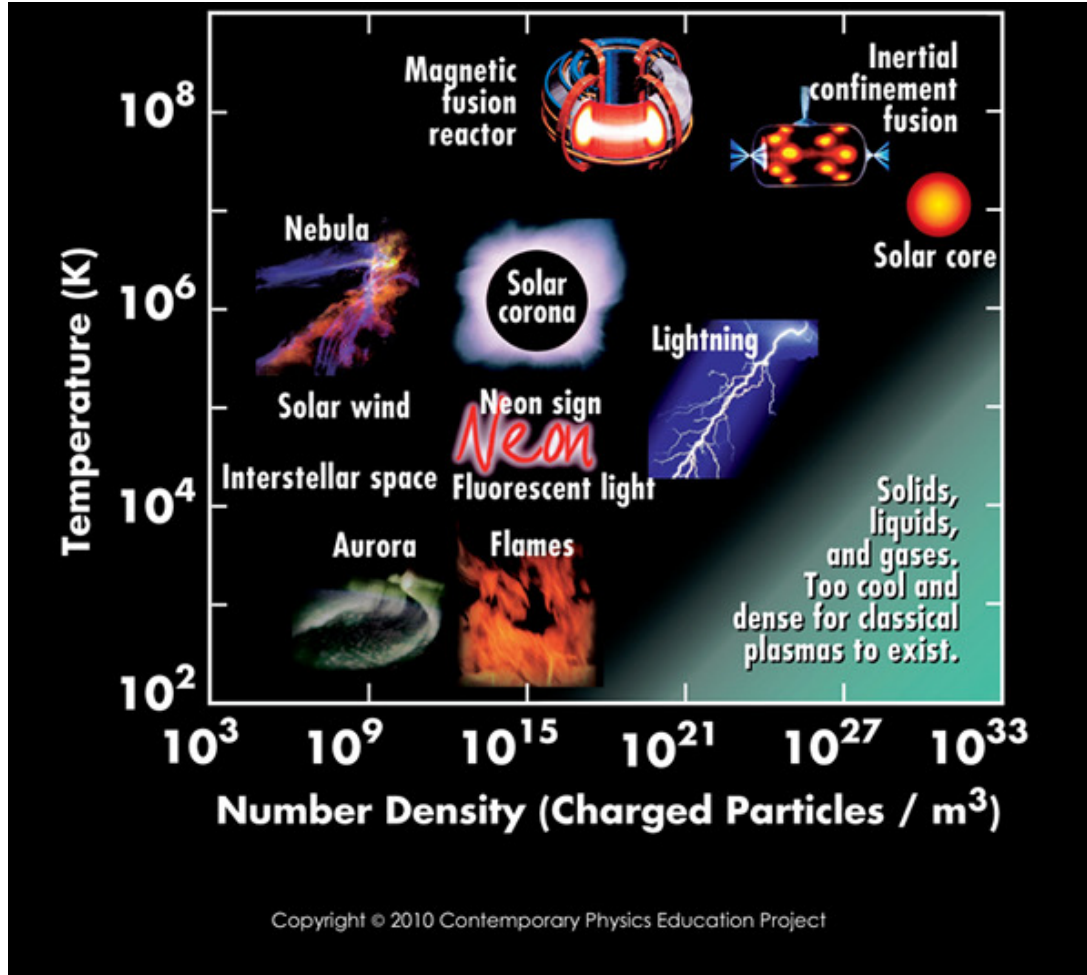


Plasma module:

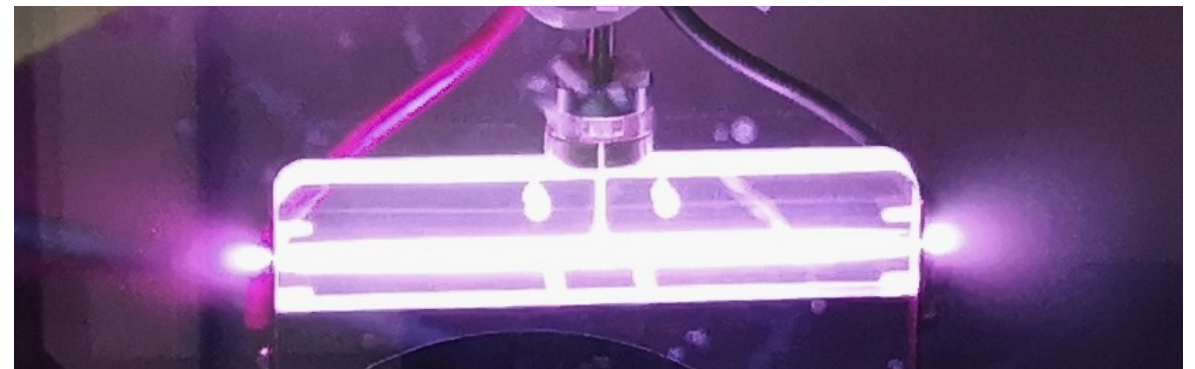
$E [\text{GV/m}] = 96 (n_e [\text{cm}^{-3}])^{1/2}$   
 $n_e = 10^{18} \text{ cm}^{-3} \Rightarrow 100 \text{ GV/m}$   
range  $mm - cm$

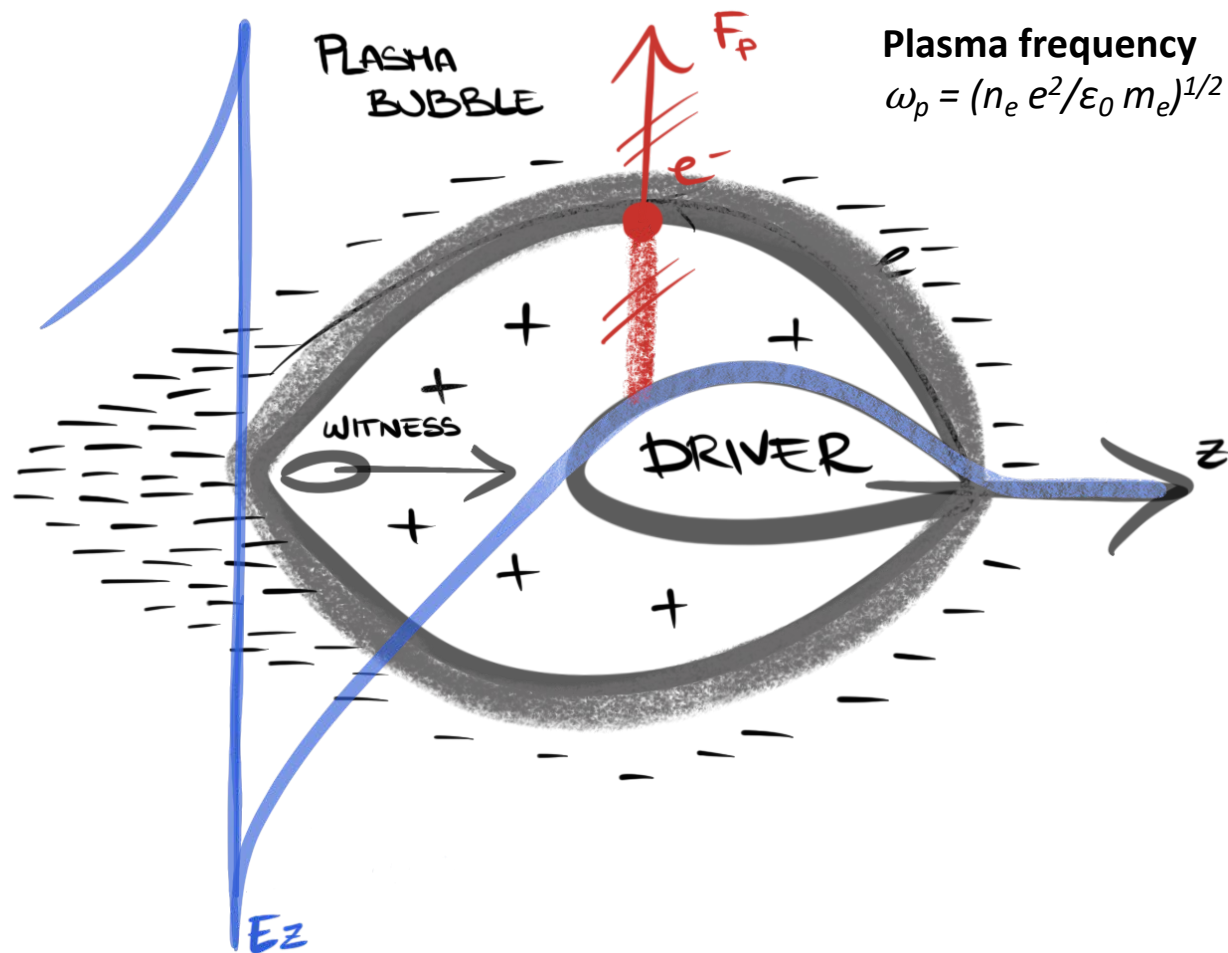
# What is a plasma?

- Quasi-neutral gas of charged particles showing collective behavior
- Local balance inside a plasma:  $n_e = Z n_i$



Plasma source for accelerators:  
 $I_p = 0.3 - 2 \text{ kA}$   
 $T_p = 20\text{k} - 100\text{k K}$   
 $V_p = 5 - 20 \text{ kV}$





Electron beam

or

Laser pulse

$$I = 10^{18} \text{ W/cm}^2 \text{ and } c\tau < \lambda_p/2$$

cold wave-breaking  
 Self-injection

VOLUME 43, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JULY 1979

**Laser Electron Accelerator**

T. Tajima and J. M. Dawson

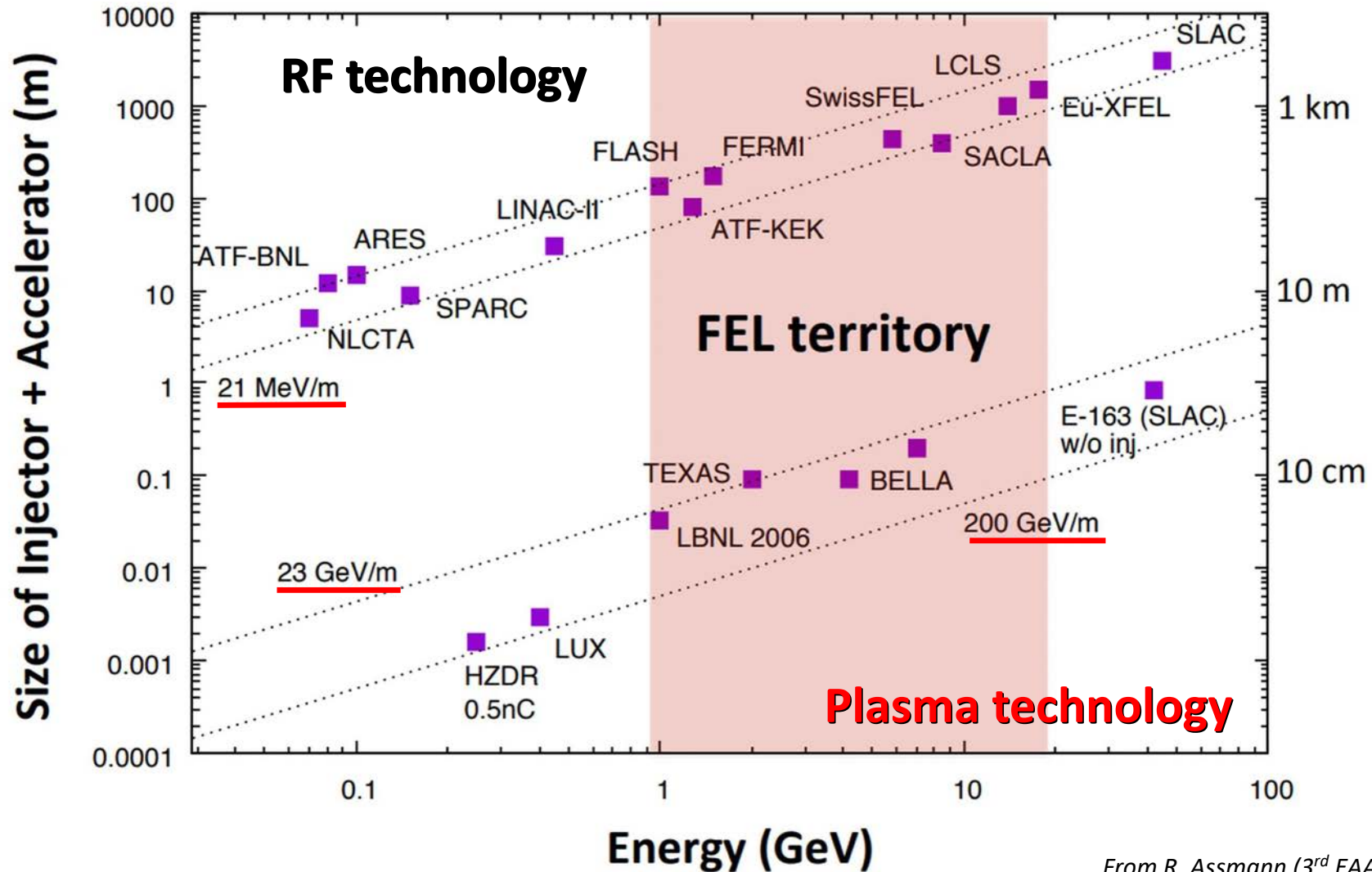
*Department of Physics, University of California, Los Angeles, California 90024*

(Received 9 March 1979)

An intense electromagnetic pulse can create a weak of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density  $10^{18}\text{W}/\text{cm}^2$  shone on plasmas of densities  $10^{18}\text{cm}^{-3}$  can yield gigaelectronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.

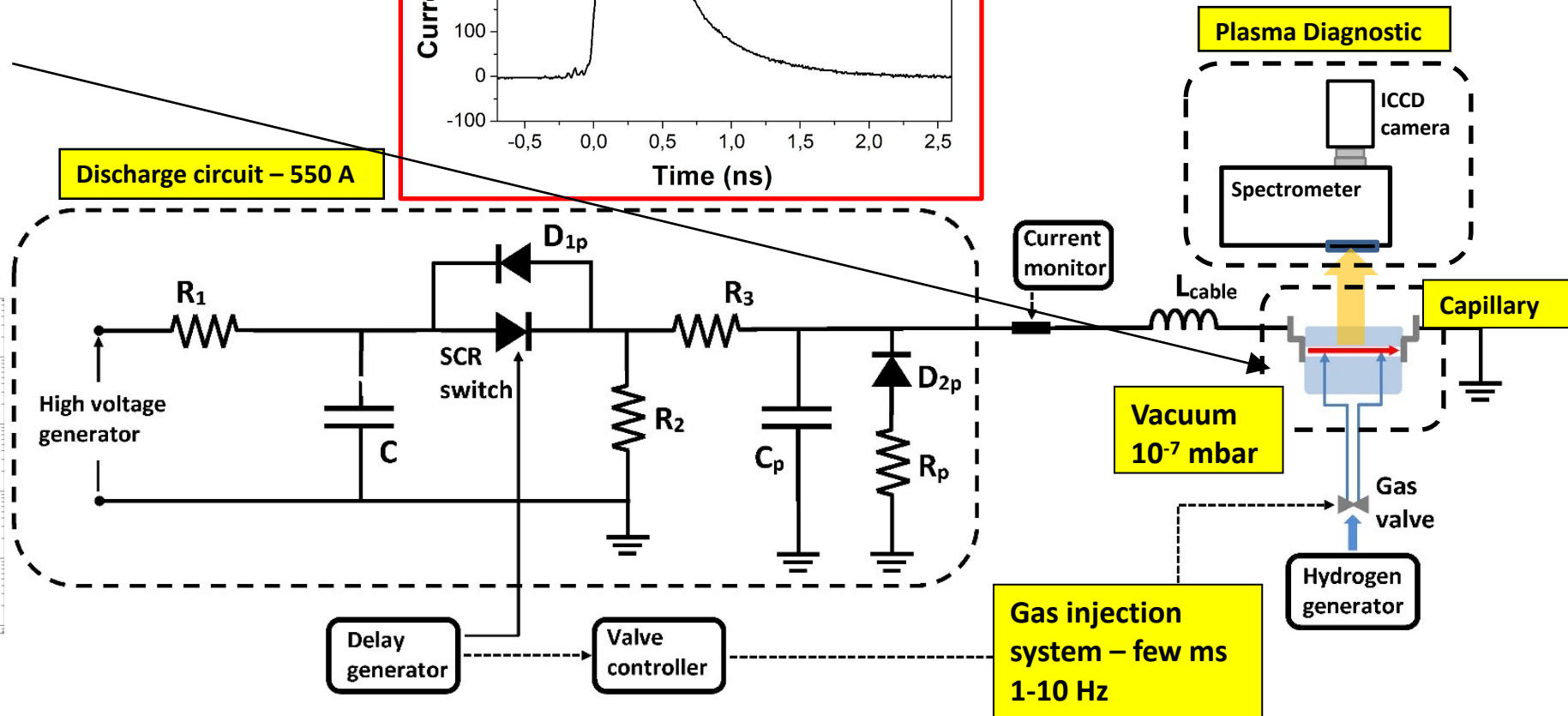
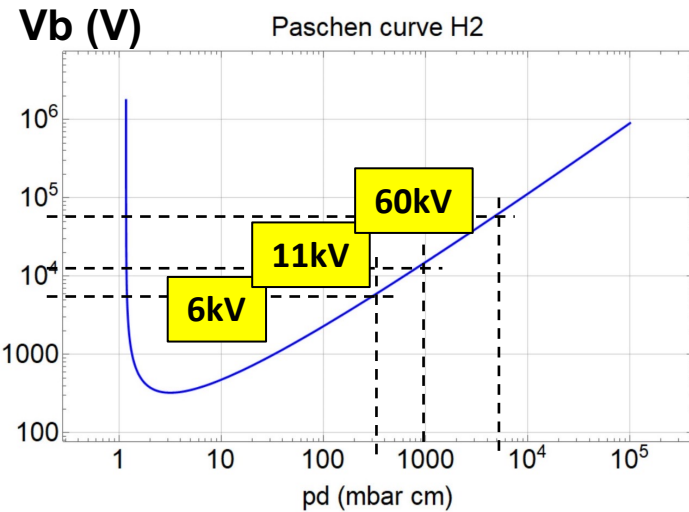
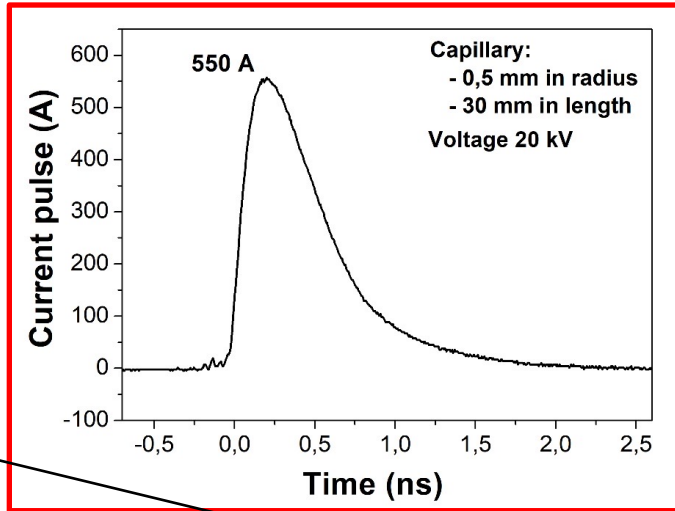
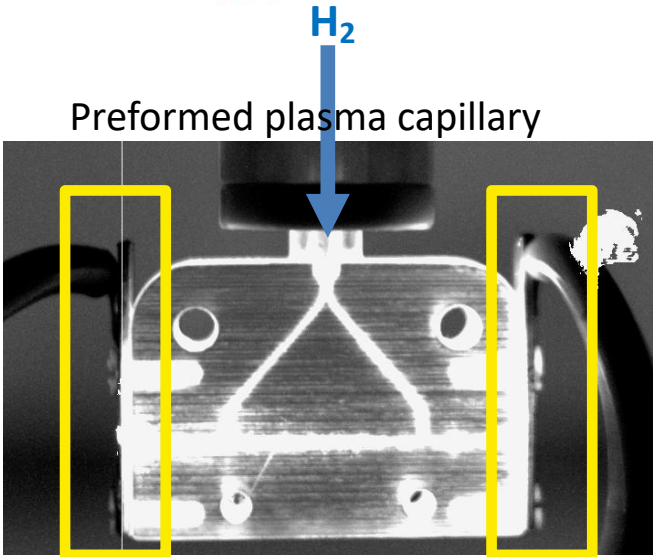


# Standard vs Plasma accelerators



From R. Assmann (3<sup>rd</sup> EAAC Workshop, 2017)

# Plasma source

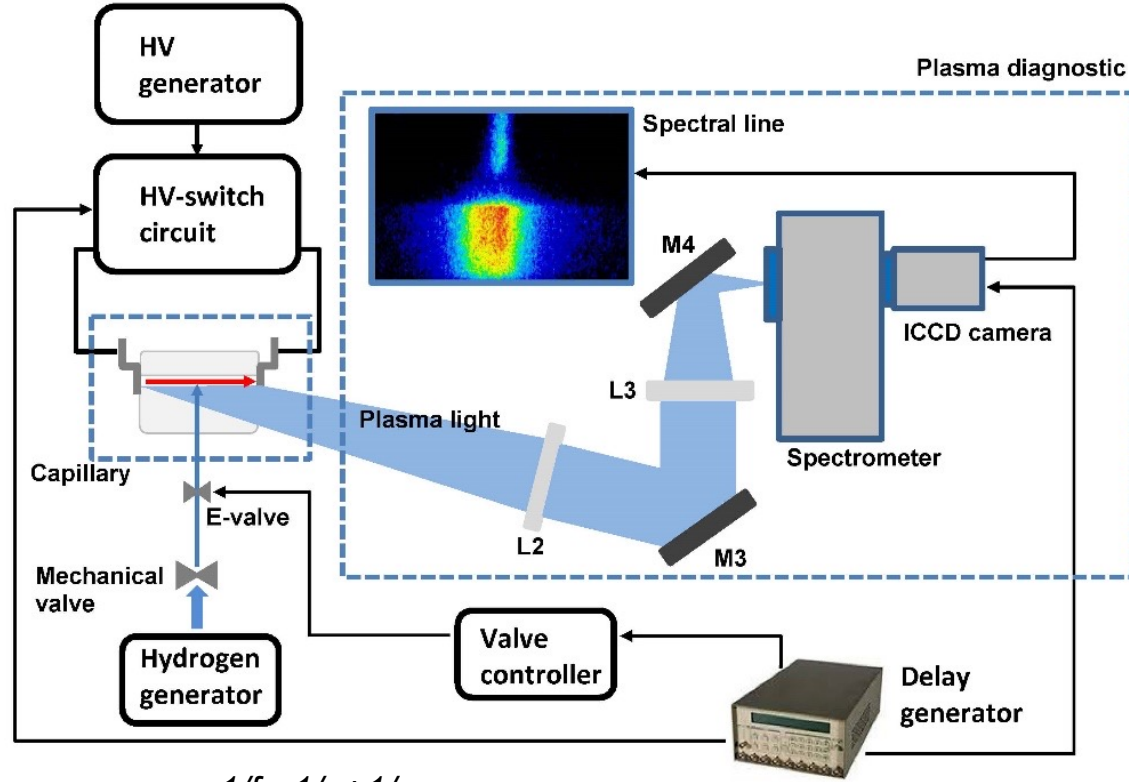
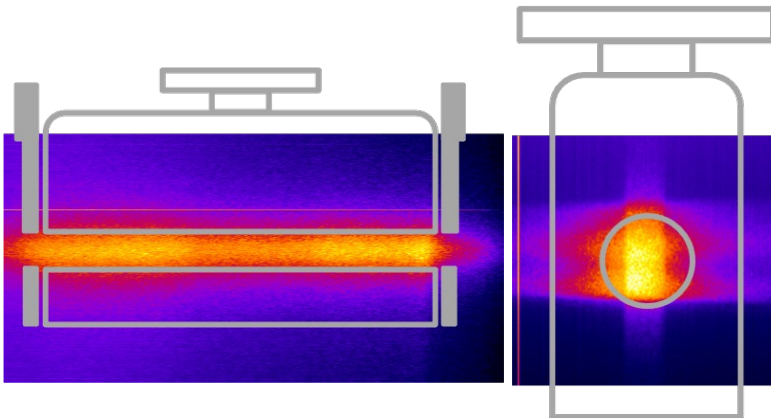


# Plasma source

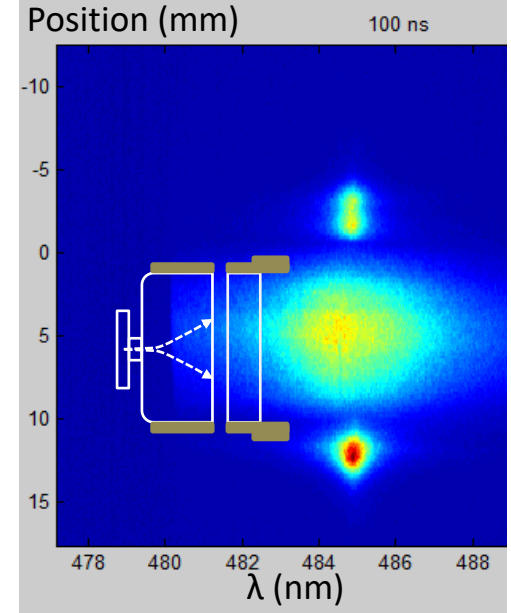
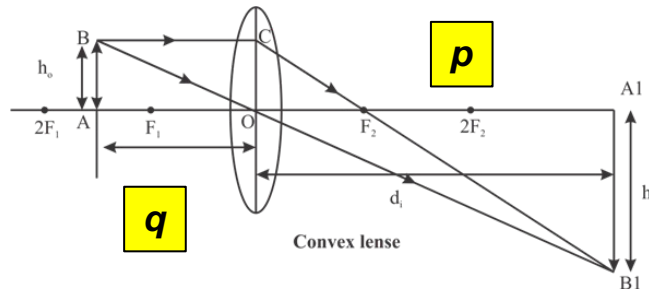
Stark Broadening:

$$n_e = 8.022 \times 10^{12} \left( \frac{\Delta\lambda}{\alpha} \right)^{\frac{3}{2}} \text{ cm}^{-3}$$

- $\Delta\lambda = \text{FWHM } H_{\alpha\beta}$
- $H_{\alpha} = 656 \text{ nm}, H_{\beta} = 486 \text{ nm}$
- $\alpha = \text{const.}$



$$1/f = 1/p + 1/q$$



# Beam driven plasma acceleration

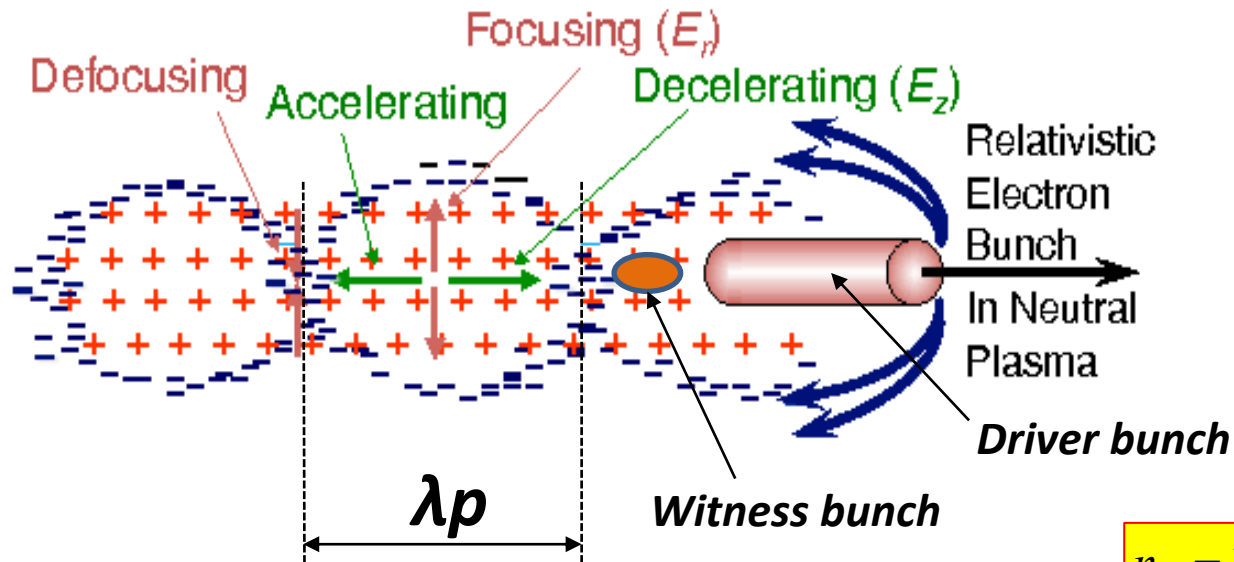
$$\omega_p = \sqrt{\frac{n_p e^2}{m_e \epsilon_0}}$$

$$\lambda_p = \frac{2\pi c}{\omega_p}$$

$$E_{cold,wb} [V/m] = \frac{m_e \omega_p c}{e} \approx 96 \sqrt{n_0 [cm^{-3}]}$$

$$\alpha = \frac{n_{bunch}}{n_p}$$

Linear regime (nb < np)  
 Weakly NL regime (nb = np)  
 Non-linear regime (nb >> np)



$$n_p = 10^{17} \text{ cm}^{-3}$$

$$\omega_p = 2 \times 10^{13} \text{ Hz}$$

$$E_0 = 30 \frac{GV}{m}$$

$$\lambda_p = 100 \text{ } \mu\text{m} \text{ (300 fs)}$$

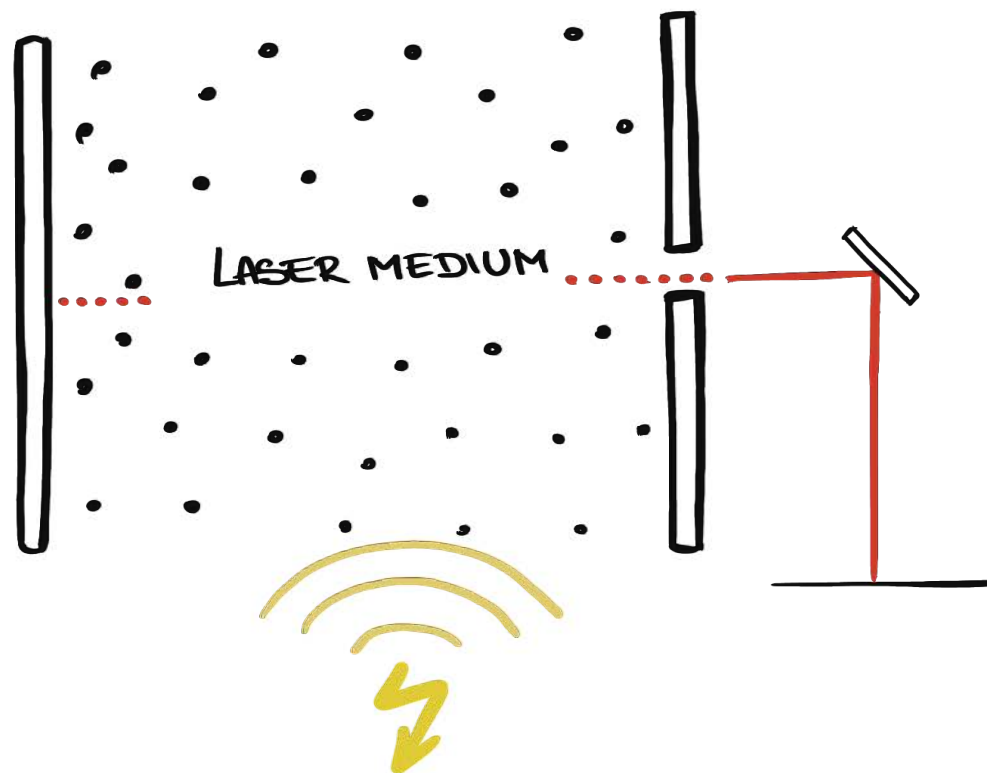
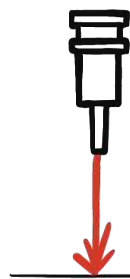
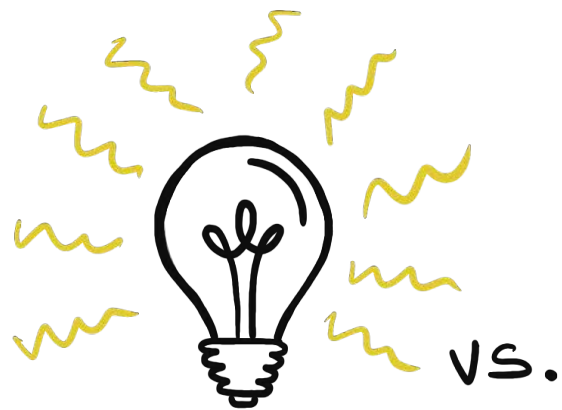
$$\sigma_{zD} \approx 25 \text{ } \mu\text{m} \text{ (75 fs)}$$

$$\sigma_{xD} \approx 2 \text{ } \mu\text{m}$$

How can we get an  
**ultra-short high-power**  
pulse?

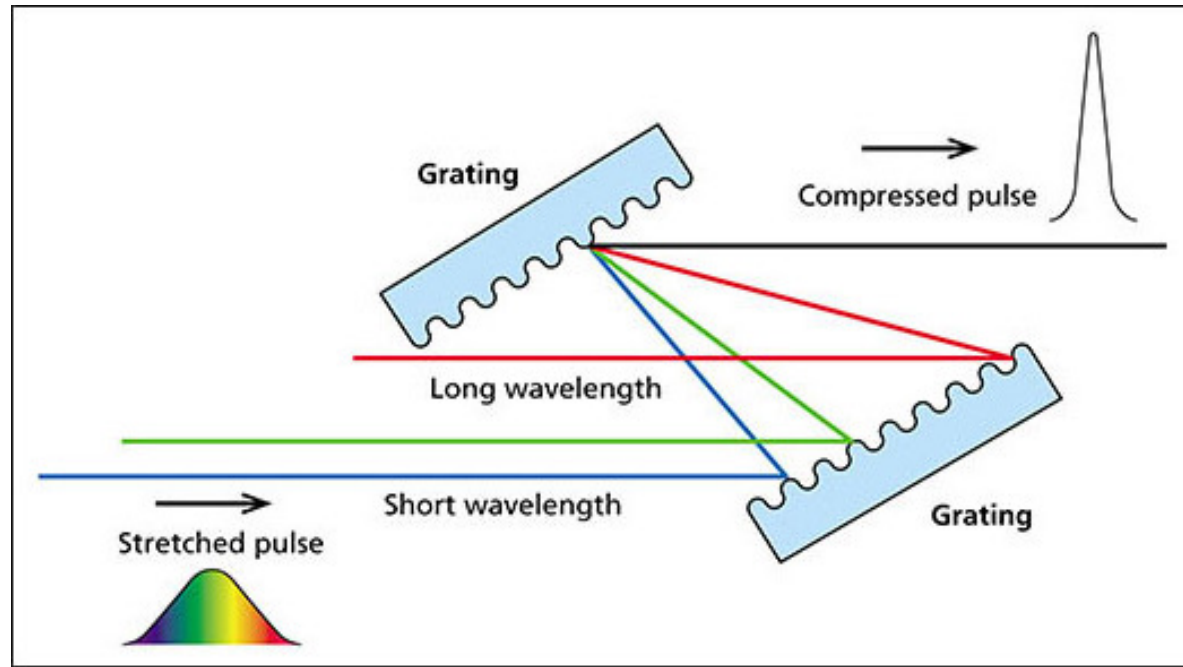


# What is a laser?



**L**IGHT  
**A**MPLIFICATION  
BY **S**TIMULATED  
**E**MISSION  
OF **R**ADIATION

# What is a high power laser?



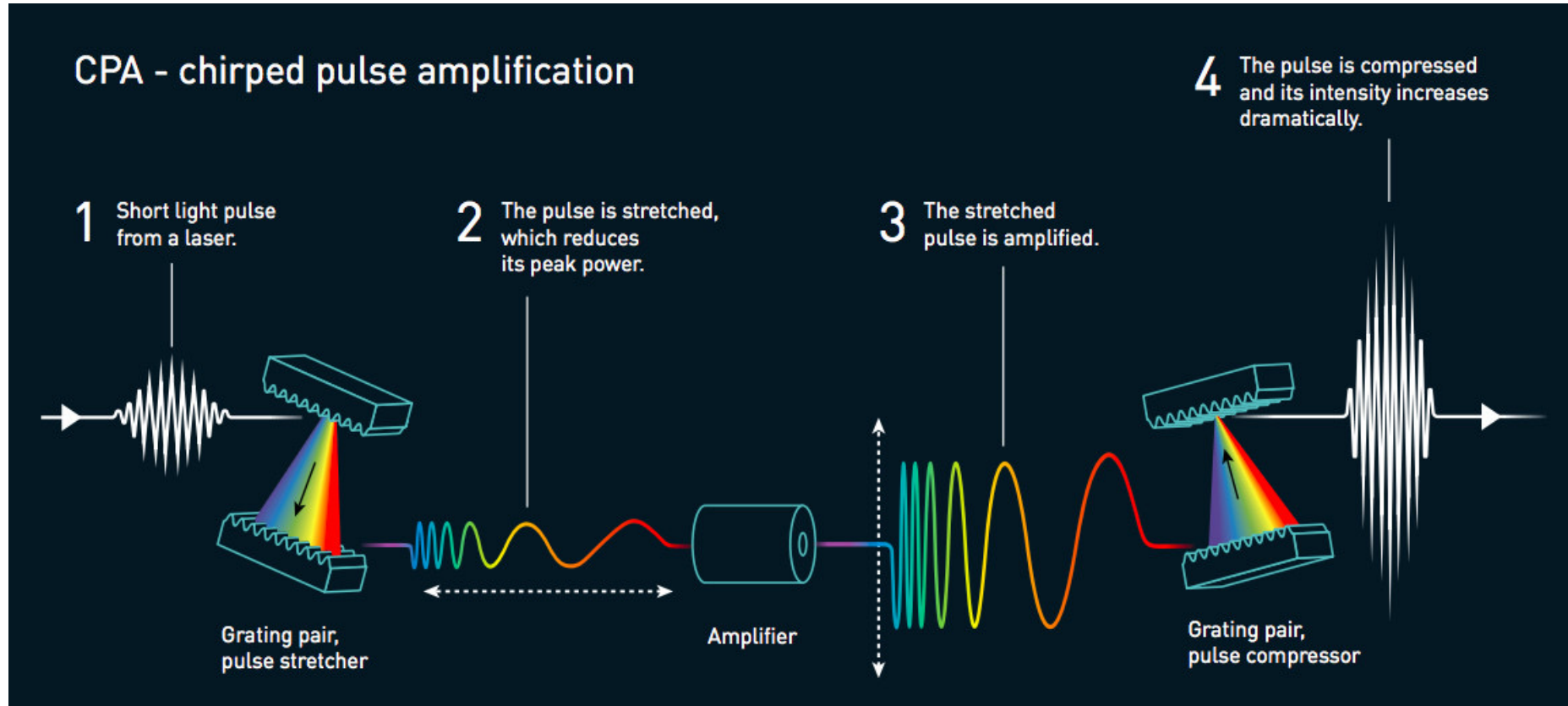
## Grating compressor

Red light → longer distance → closer to blue

Blue light → shorter distance → catches up

# What is a high power laser?

## Chirped Pulse Amplification

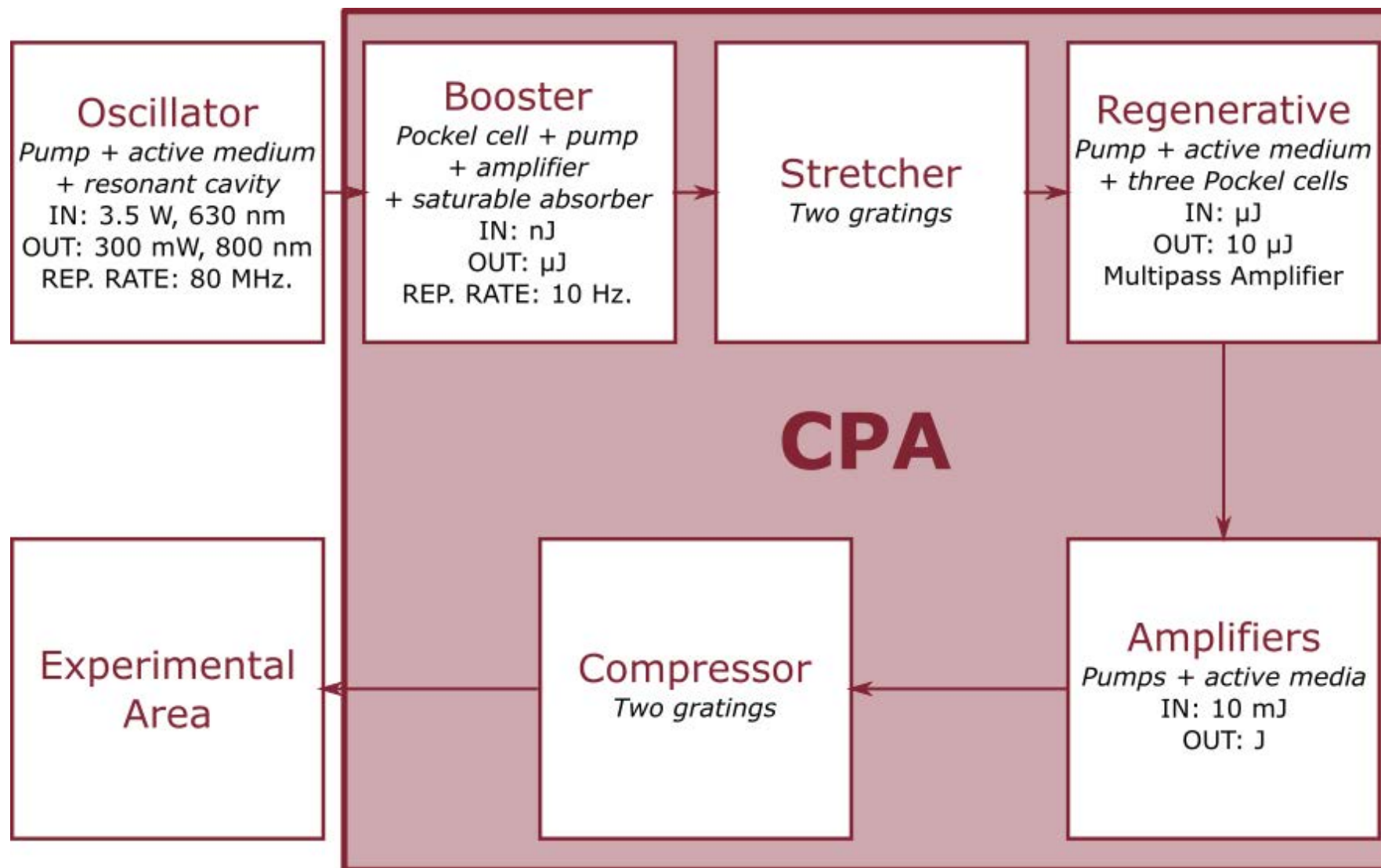


Physics NOBEL PRIZE 2018

**Short pulse → Stretch in time (chirp) → Amplify → Compress in time**

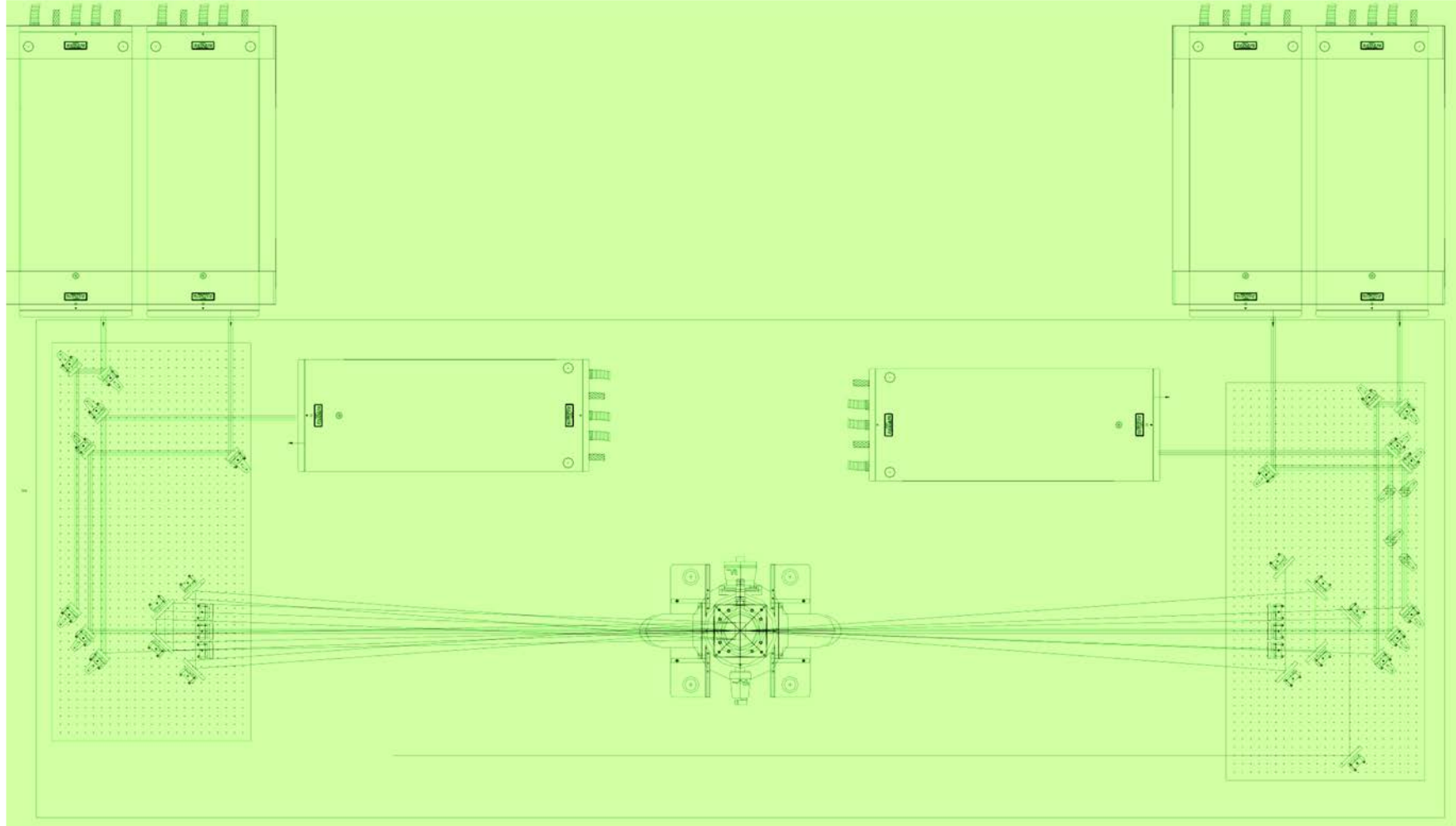


# What is a high power laser?



# What is a high power laser?

example of a Multipass Amplifier



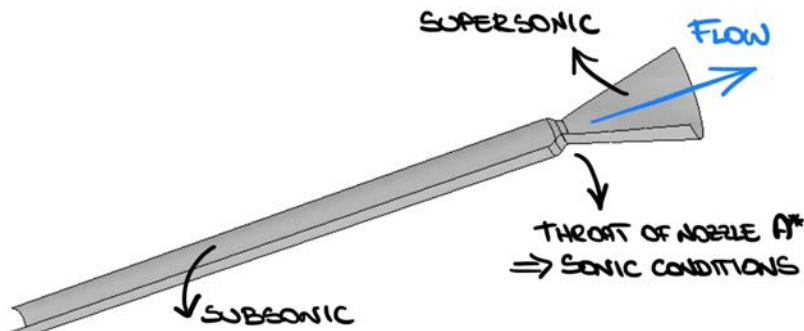
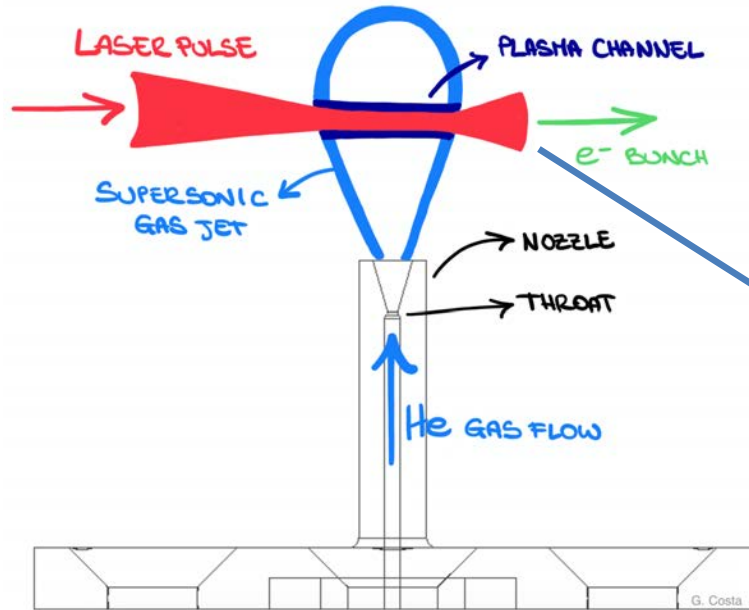
Normalized laser intensity  $a_0 = eE_0/(m_e\omega_p c)$

Electron motion in a laser field oscillation in the direction of the electric field

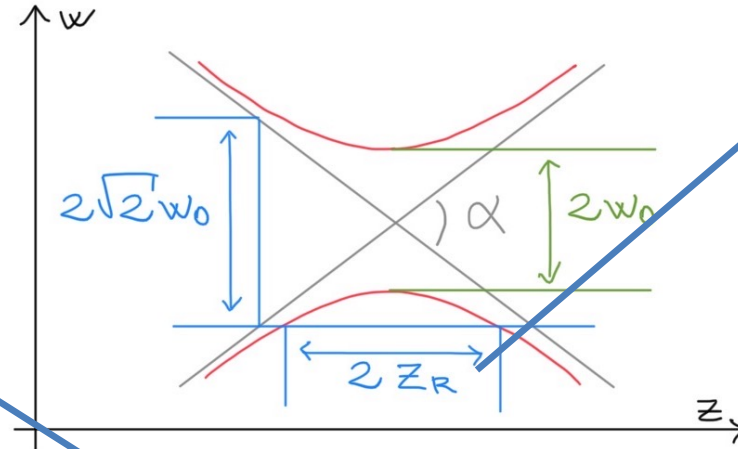
$$\text{velocity } \beta = -a_0 \sin(\omega_p t)$$

- $a_0 < 1$  linear regime  $\Rightarrow$  no wavebreaking
- $a_0 \approx 1$  non-linear regime  $\Rightarrow v_e \approx c \Rightarrow$  self-injection
- $a_0 \gg 1$  strongly non-linear regime  $\Rightarrow$  plasma bubble can focus and accelerate electrons to high energies

Gas jet flow within a **supersonic** nozzle



Along the propagation direction of a laser beam that is focused



**Rayleigh length  $z_R$**

distance from the laser waist  $w_0$  to the point where the cross-sectional area is doubled

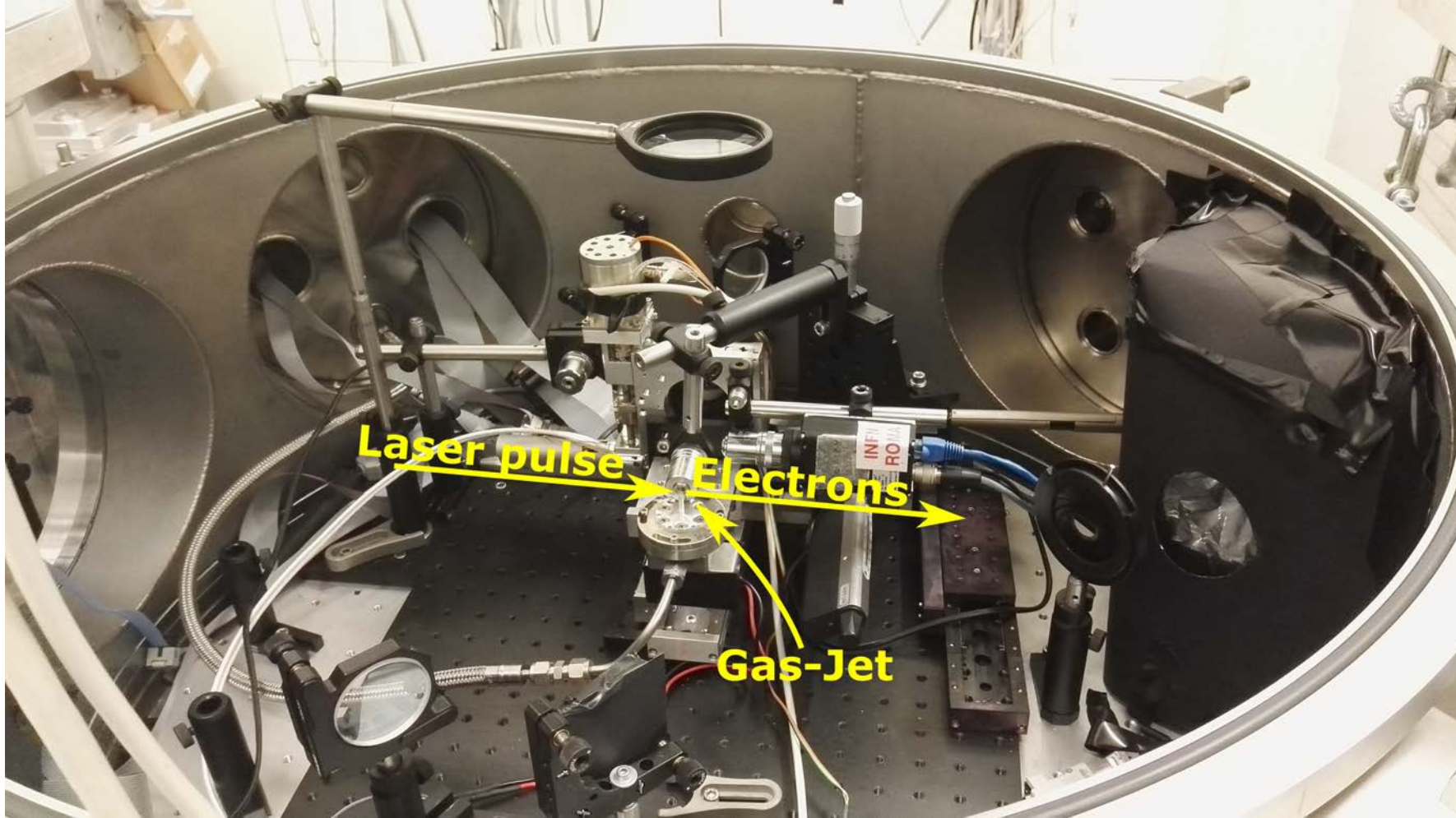
Laser with intensity transverse gradient

plasma refractive index acts as a convex lens

Laser **self-focusing** in plasma

increasing the pulse confinement length

## Laser interaction chamber



## Plasma density interferometric measurements

Mach Zehnder interferometer:

Path difference between probe and reference laser

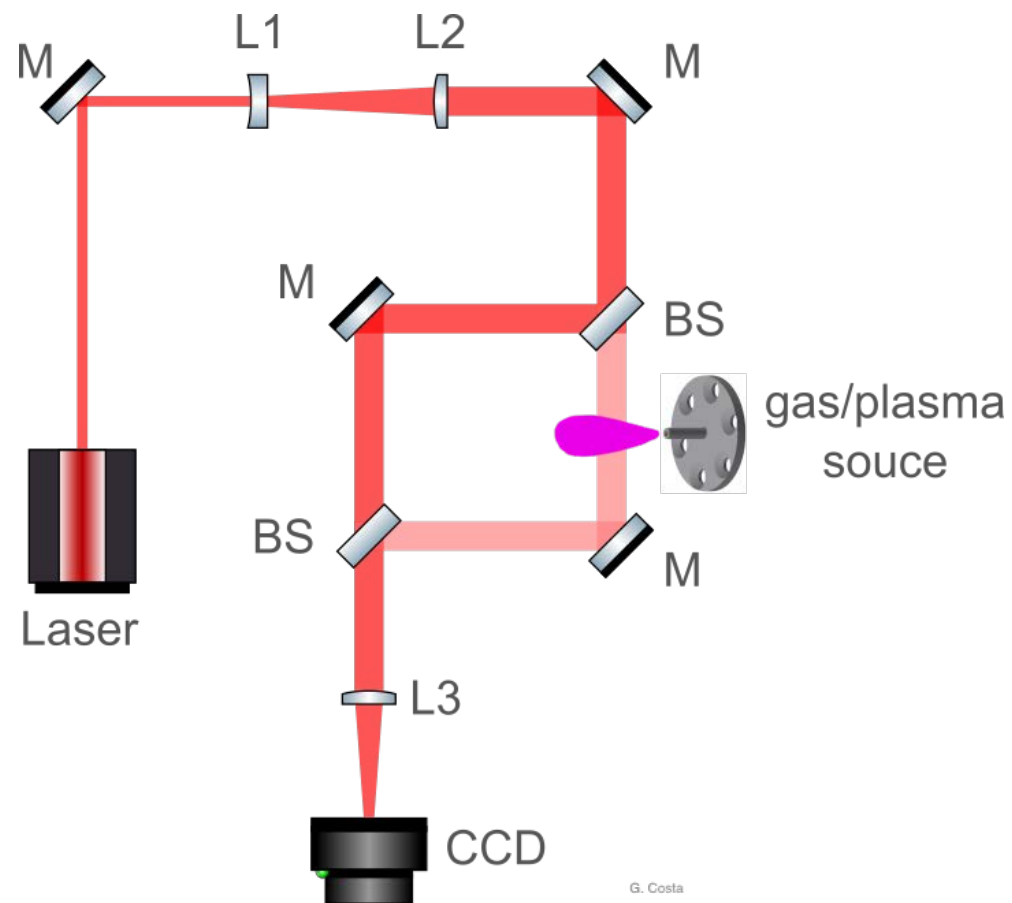


interference fringes

Plasma refractive index

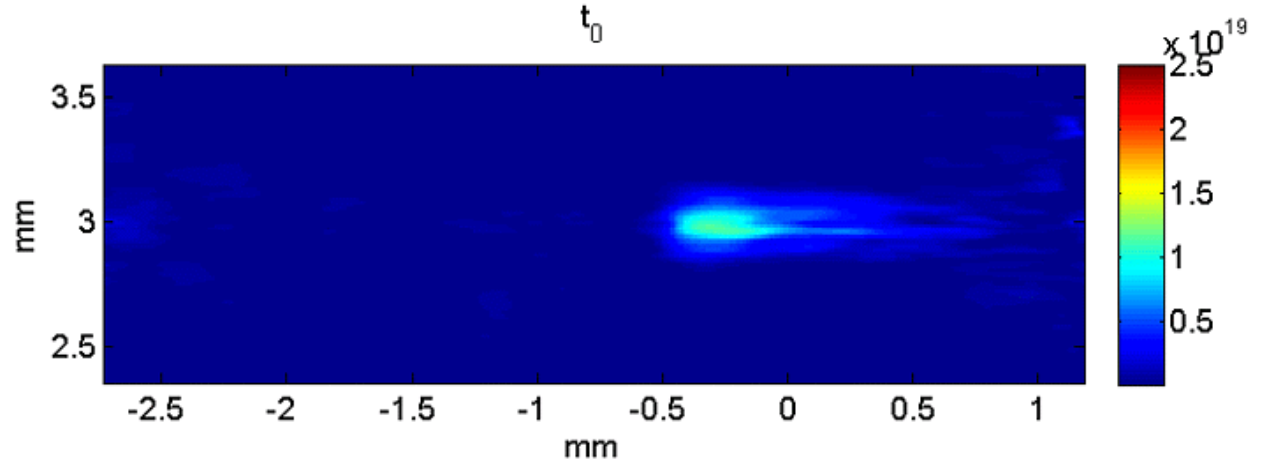
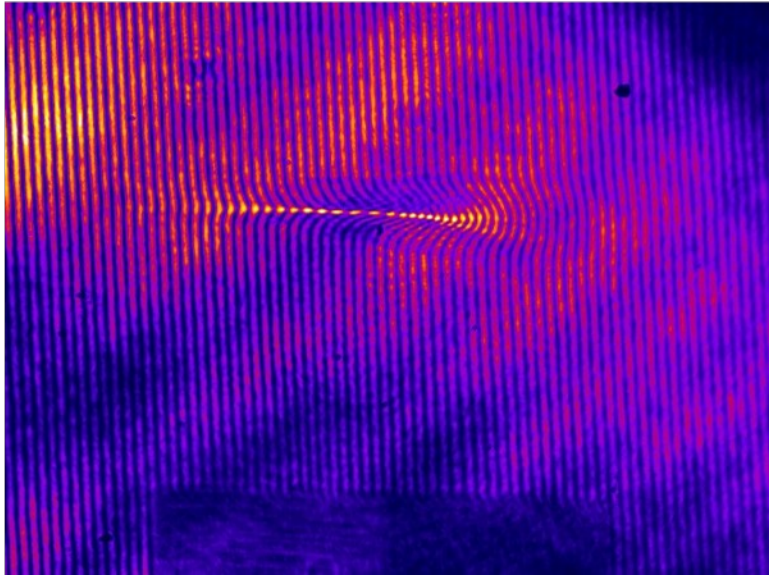
$$\eta = \left(1 - \frac{n_e}{n_c}\right)^{\frac{1}{2}}$$

- $n_e$  = e<sup>-</sup> plasma density
- $n_c$  = critical e<sup>-</sup> density

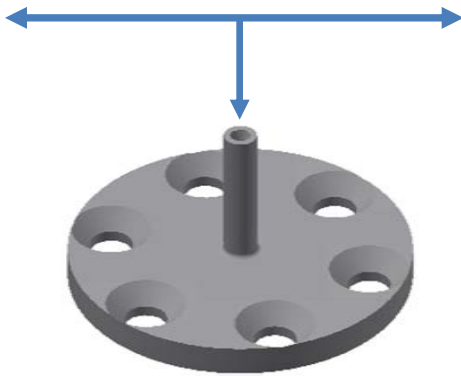


# Laser plasma acceleration

High power pulse: few J – 25 fs – 20 μm @focus → Laser intensity  $10^{19}$  W/cm<sup>2</sup>



typical densities  $\sim 10^{19}$  cm<sup>-3</sup> → self-injection!



Quality of produced and accelerated electron bunches:

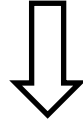
- divergence
- energy spread -> emittance

Limitations of this scheme:

- diffraction
- pump depletion
- dephasing

# Laser plasma acceleration

$e^-$  energy gain  $\propto$  Acceleration field  $\times$  Acceleration length  $\longrightarrow \pi Z_R$



Rayleigh length  $Z_R = \pi \frac{r_i^2}{\lambda_0}$   
 $\Rightarrow$  Diffraction

## optical guiding in plasma

- gaussian intensity laser pulse  $I(r) = I_0 e^{-\left(\frac{2r^2}{r_i^2}\right)}$

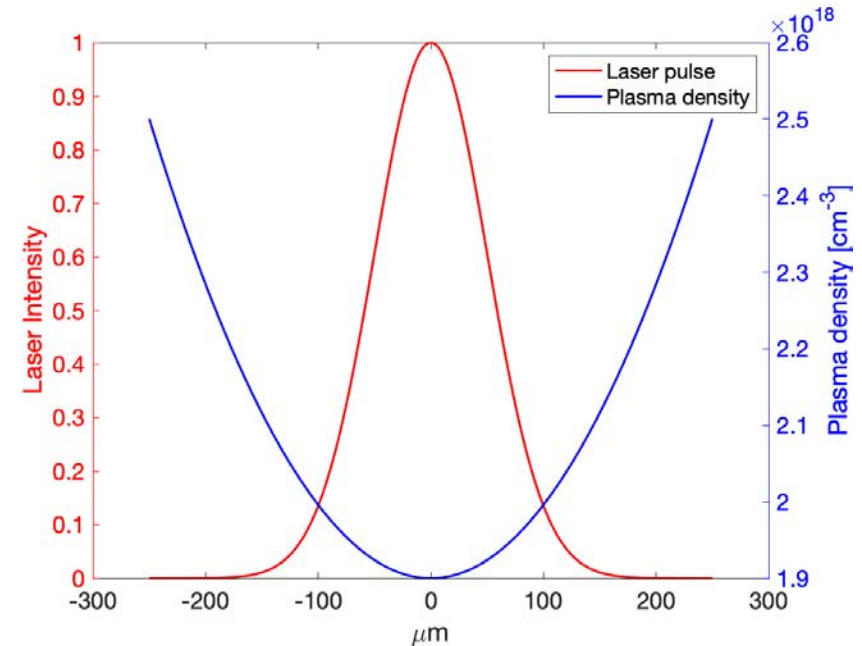
- parabolic density distribution plasma channel

$$n(r) = n_0 + \left(\frac{r}{r_m}\right)^2 n_d \quad \text{for } r \leq r_m$$

- if waist radius  $r_i = r_m = (\pi r_e n_d)^{-\frac{1}{2}}$ 
  - $\lambda_0$  = laser wavelength
  - $r_m$  = matching radius
  - $r_e = e^-$  radius
  - $n_0$  = plasma density on the axis
  - $n_d$  = density difference axis-walls

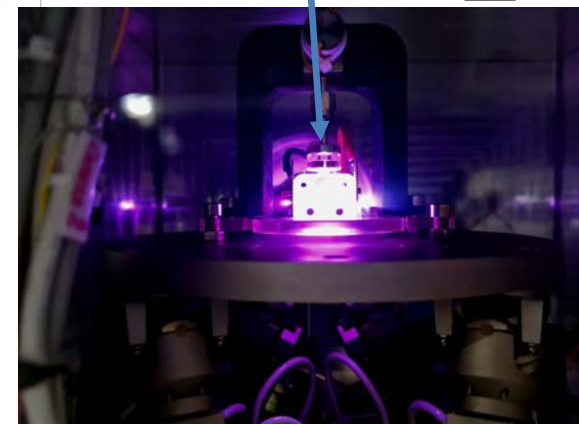
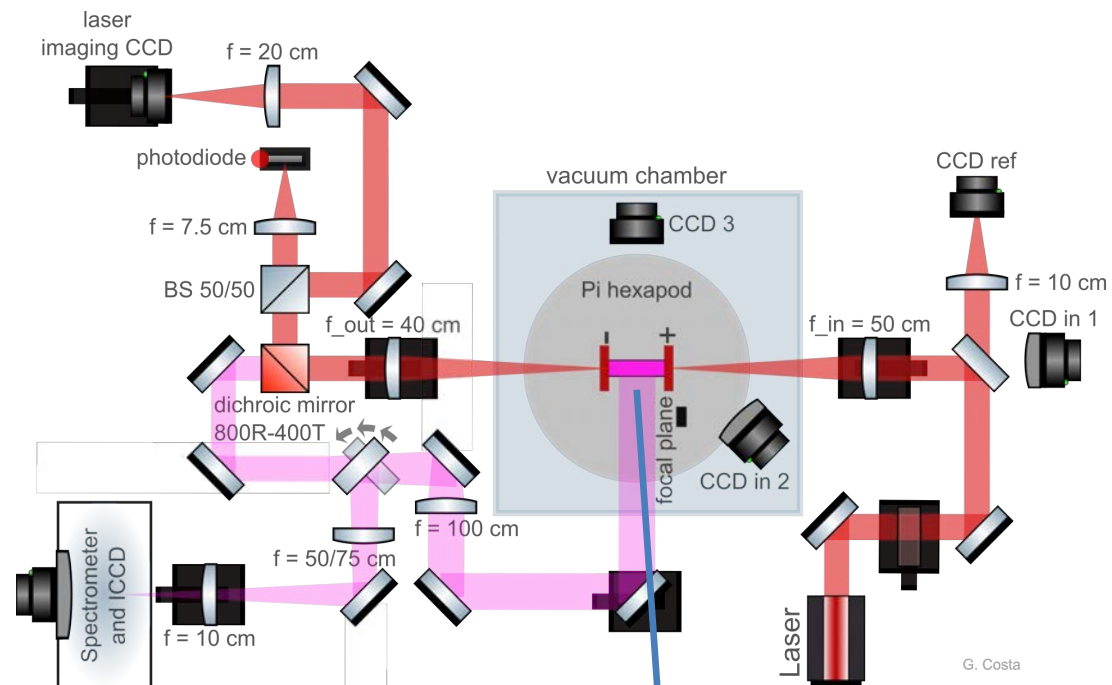
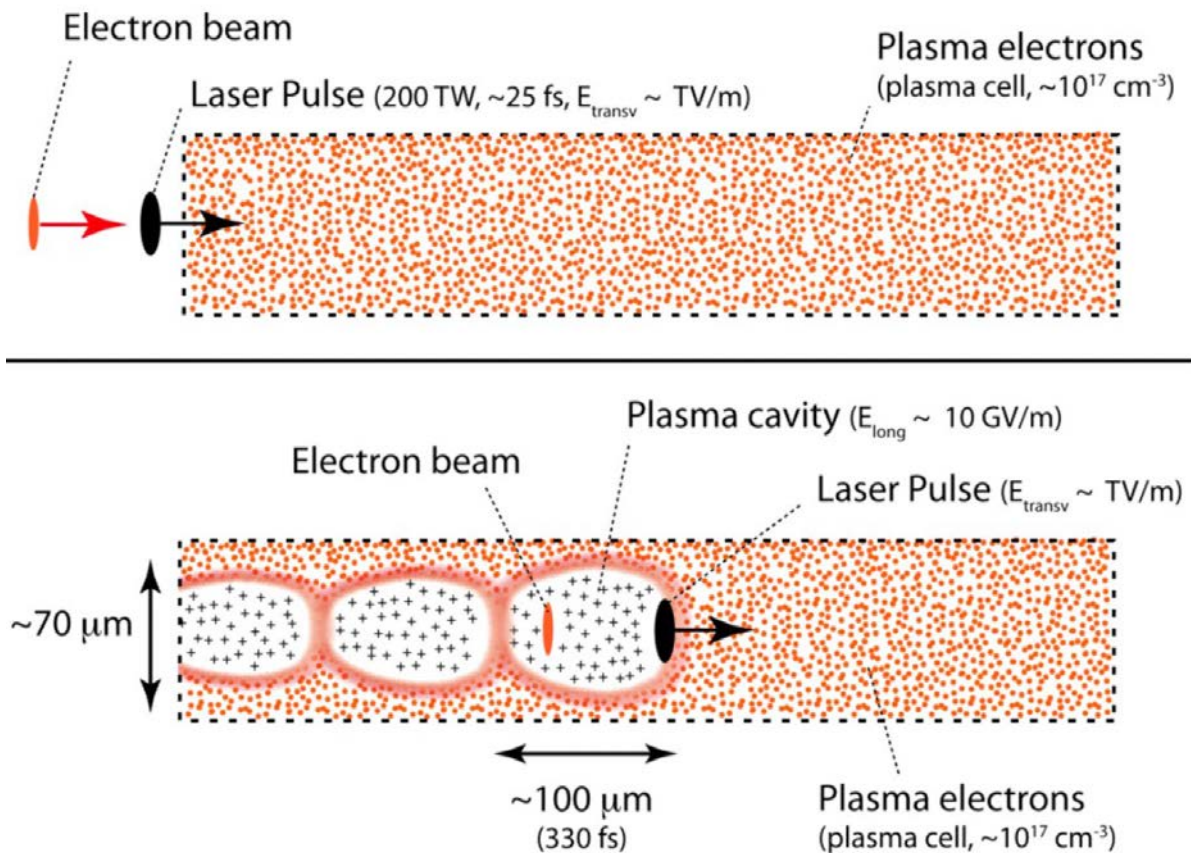
ex: for  $r_i = 60 \mu m$

- $n_0 \sim 2.5 \times 10^{18} \text{ cm}^{-3}$
- $n_d \sim 6 \times 10^{17} \text{ cm}^{-3}$



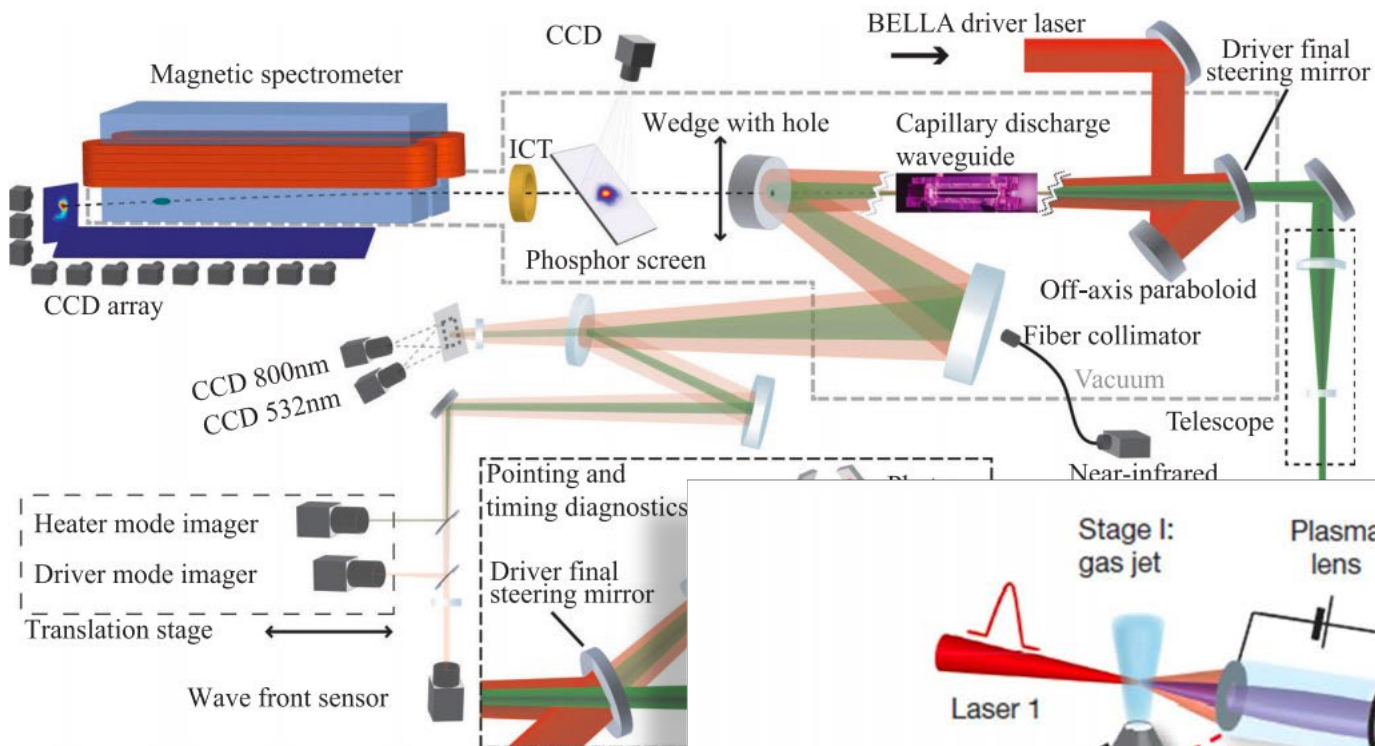


## Laser external injection scheme



G. Costa

# Laser plasma acceleration

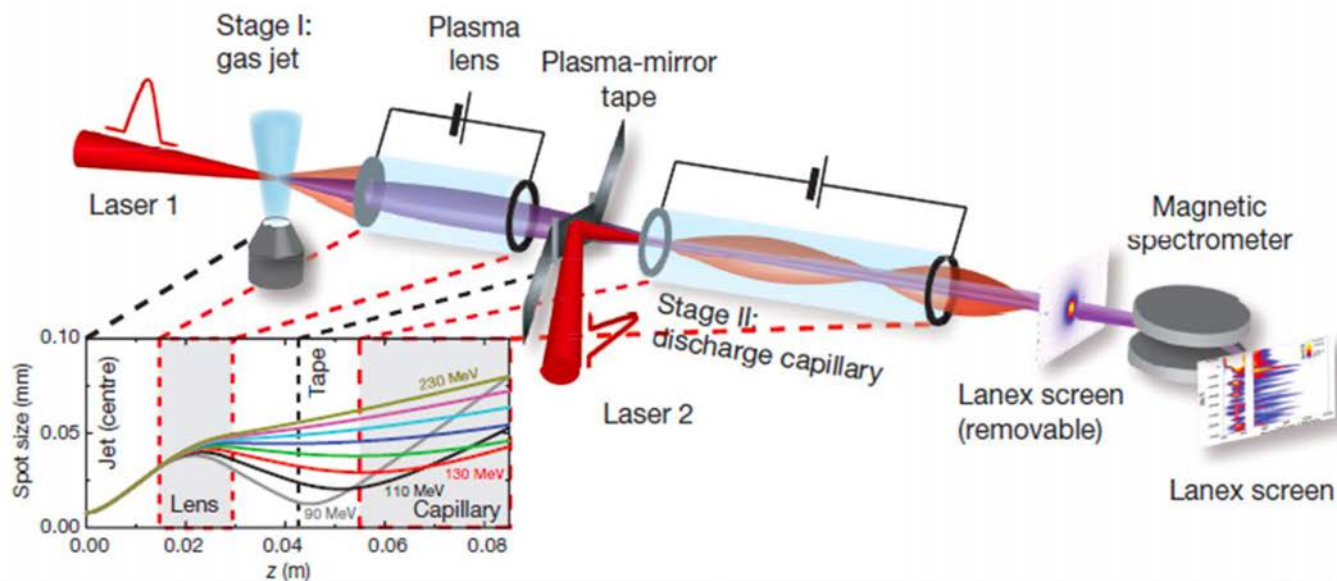


## 7.8 GeV in 20 cm

Gonsalves, A. J., et al. "Petawatt laser guiding and electron beam acceleration to 8 GeV in a laser-heated capillary discharge waveguide." *Physical review letters* 122.8 (2019): 084801.

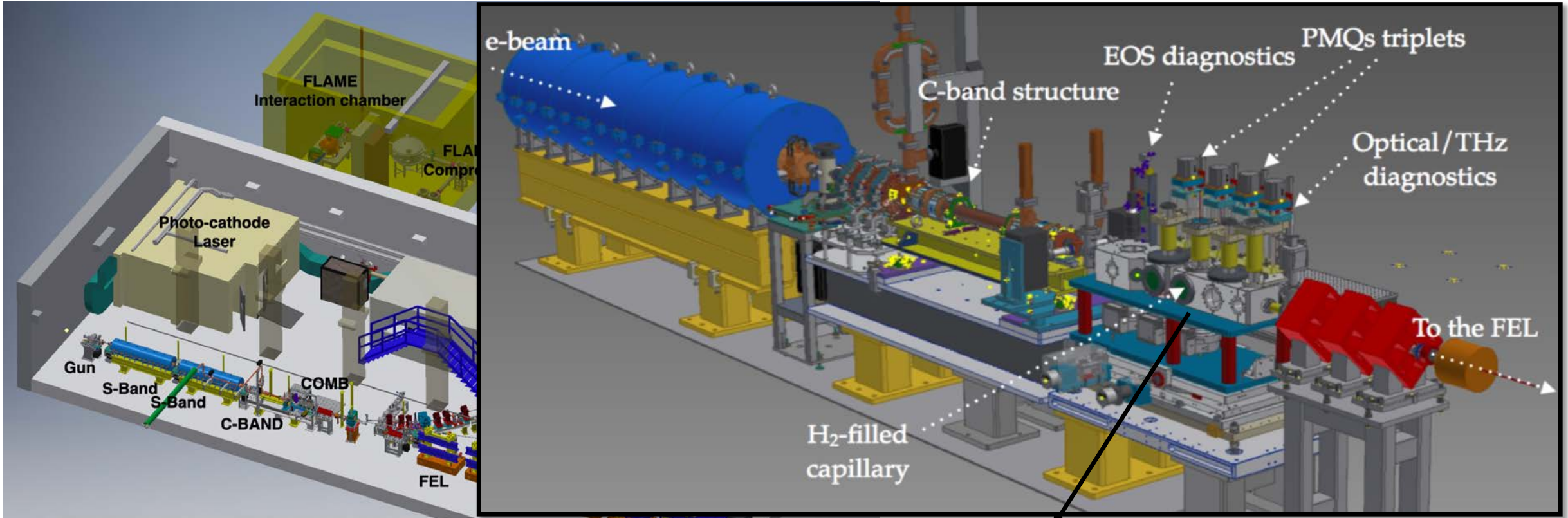
## Staging

*Nature* 530.7589 (2016): 190-193.

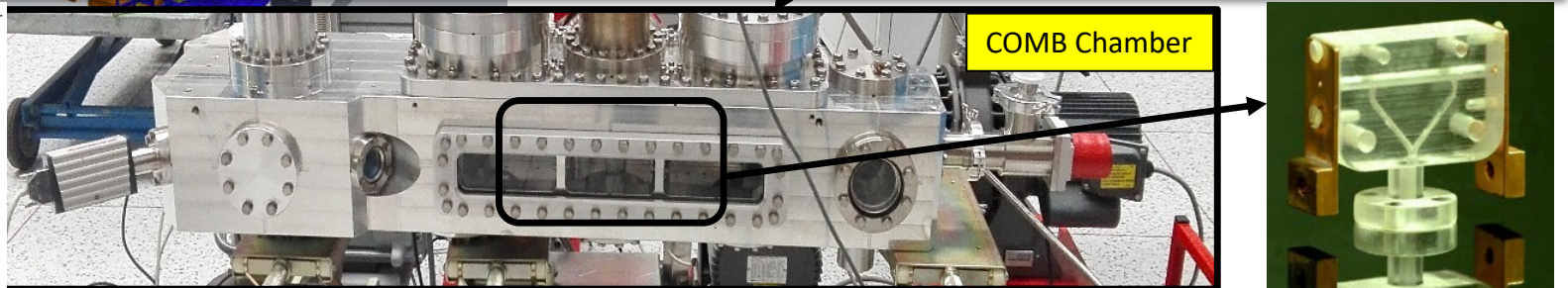


## Recap different schemes of plasma acceleration

- Laser only: *easiest* to implement, requires to tune the laser and the target, **but** difficult control over the whole process
- Electrons only: easier implementation than external injection, no need for independent synchronization system and driver guiding, **but** it depends heavily on the ability to properly tailor the driver(s) and witness phase spaces
- Laser + electrons: in principle has the best potentialities in term of e-beam brightness and energy, **but** it is the hardest to implement for laser guiding and synchronization issues



Energy	(30 – 180) MeV
Energy Spread	$\simeq 0.01\% - 1\%$
Charge	10 pC – 1 nC
Bunch length range (FWHM)	50 fs – 10 ps
Normalised Emittance	(1 – 3) mm mrad
Max Rep. Rate	10 Hz



## FLAME Laser

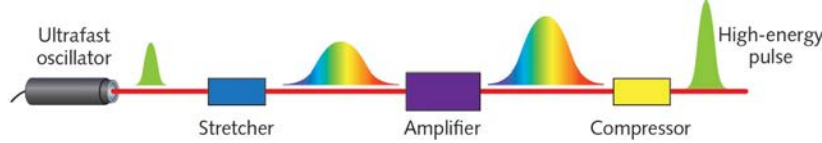
### MAIN PULSE

Peak power	250 TW
Energy on target	7 J
Rep. rate	10 Hz
Temporal length	25 fs
FW 1/e <sup>2</sup> @ focus	20 μm
Intensity	10 <sup>19</sup> W/cm <sup>2</sup>
Contrast-ratio	10 <sup>10</sup>

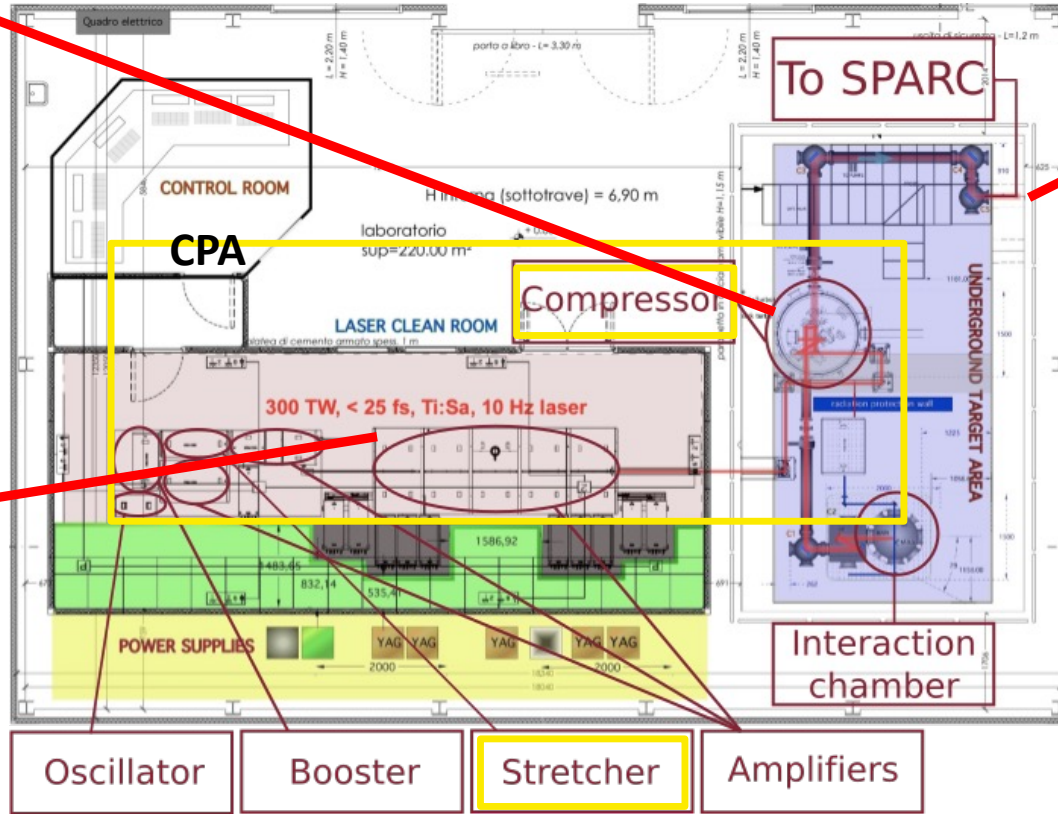
### PROBE PULSE

Energy on target	10 mJ
Temporal length	50 fs
FW 1/e <sup>2</sup> @ focus	120 μm
Intensity	10 <sup>16</sup> W/cm <sup>2</sup>

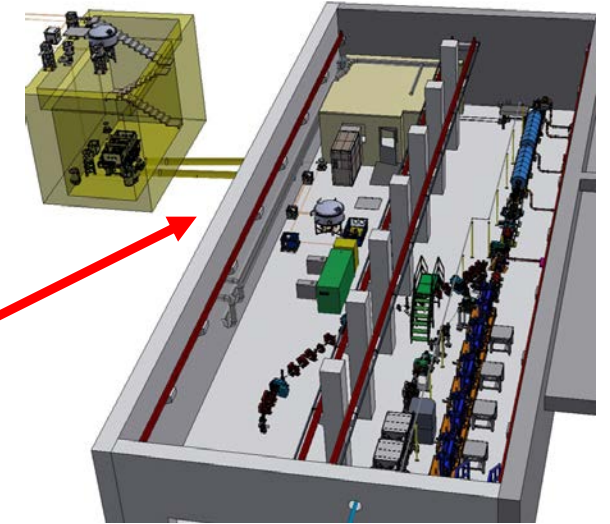
High-power Ti:Sa based laser, ultra-short CPA based pulse



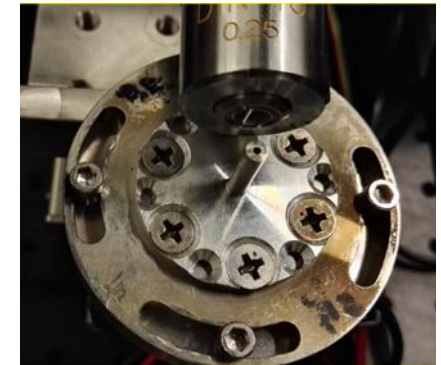
Strickland and Mourou 1985  
Nobel Prize in Physics 2018



80 MHz, 220 mW

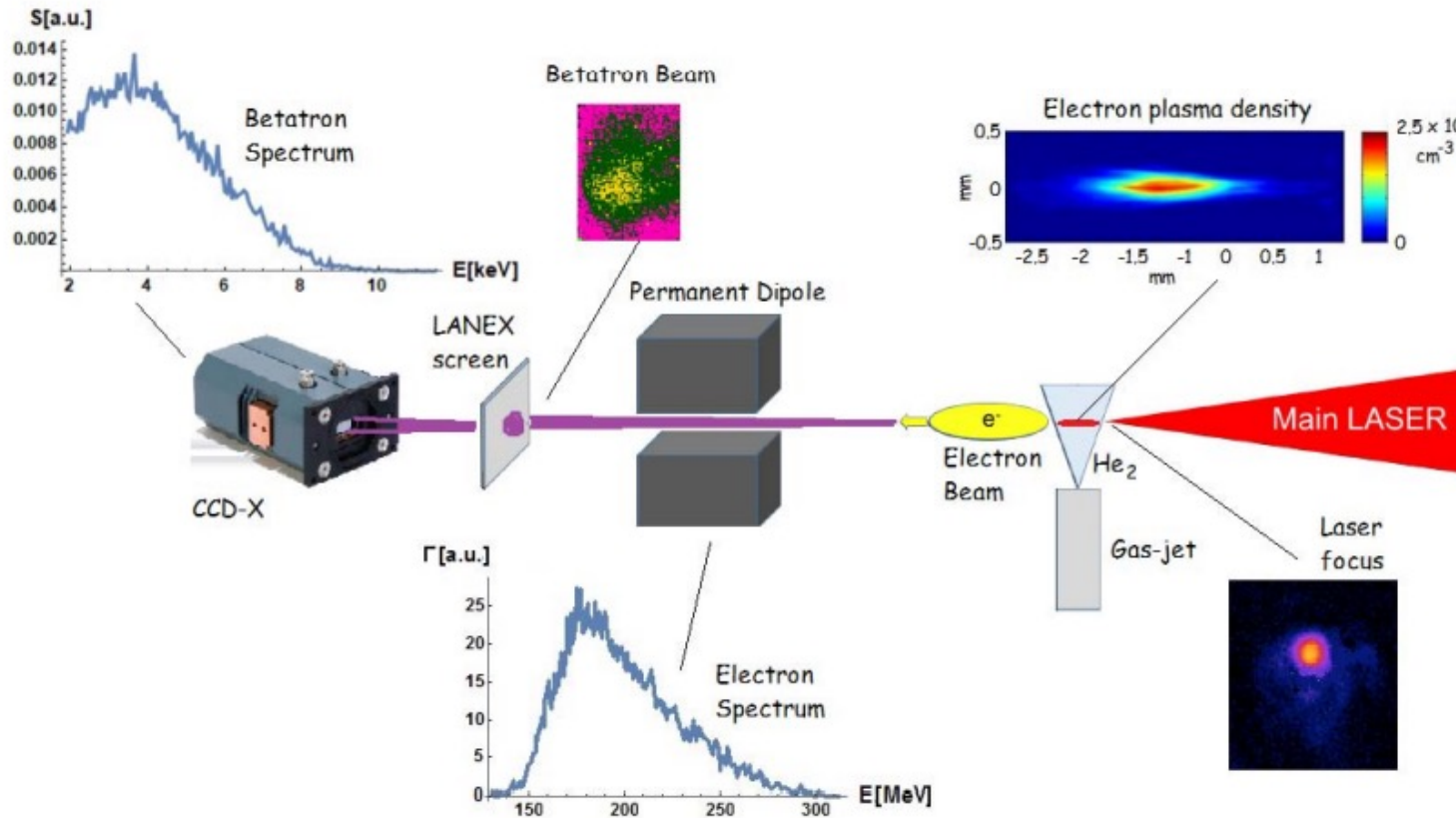


supersonic nozzle for self-injection experiments  
and X-rays radiation emission

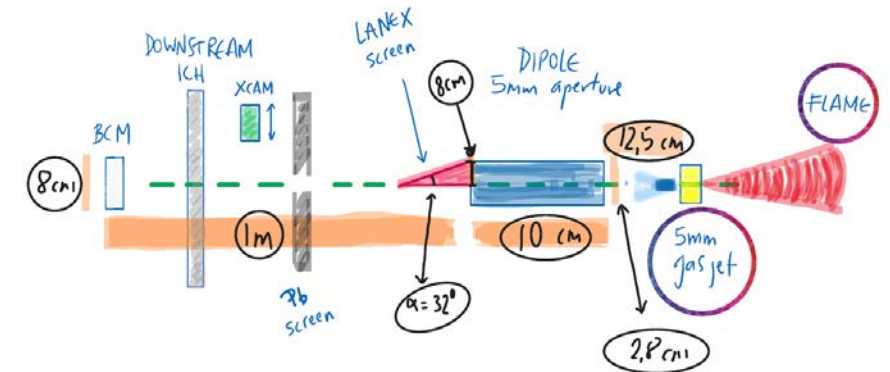


FG Bisesto, et al., doi:10.1016/j.nima.2018.02.027

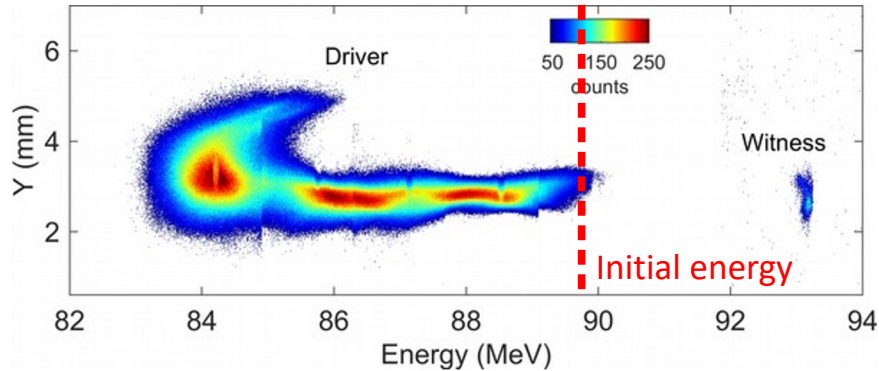
## FLAME experimental setup per plasma acceleration



- Laser energy 6 J
- Laser temporal length 30 fs FWHM
- Laser focal spot 18-20 μm FWHM

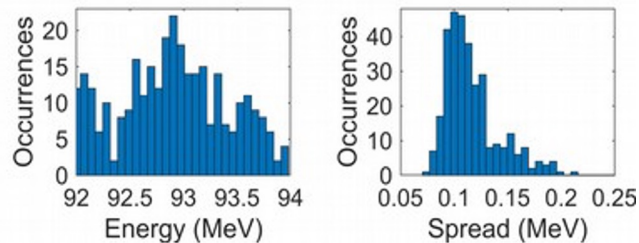
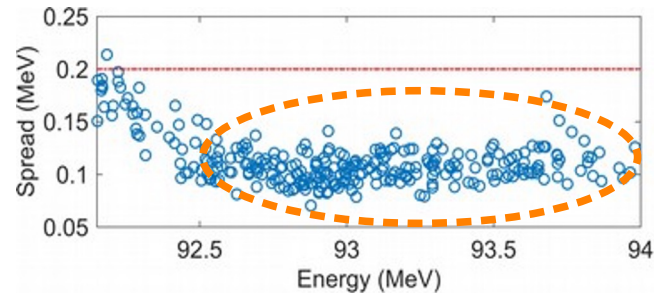


Courtesy M Galletti



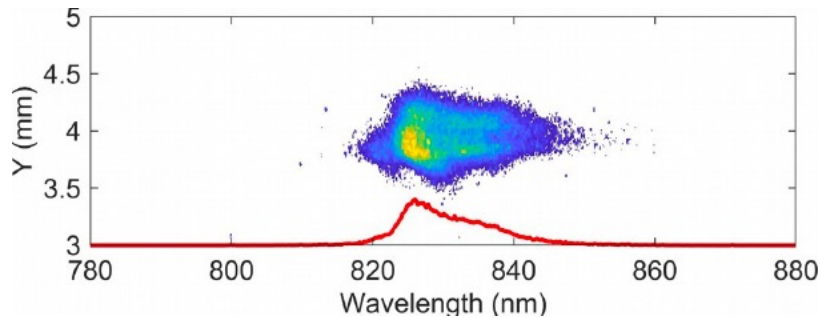
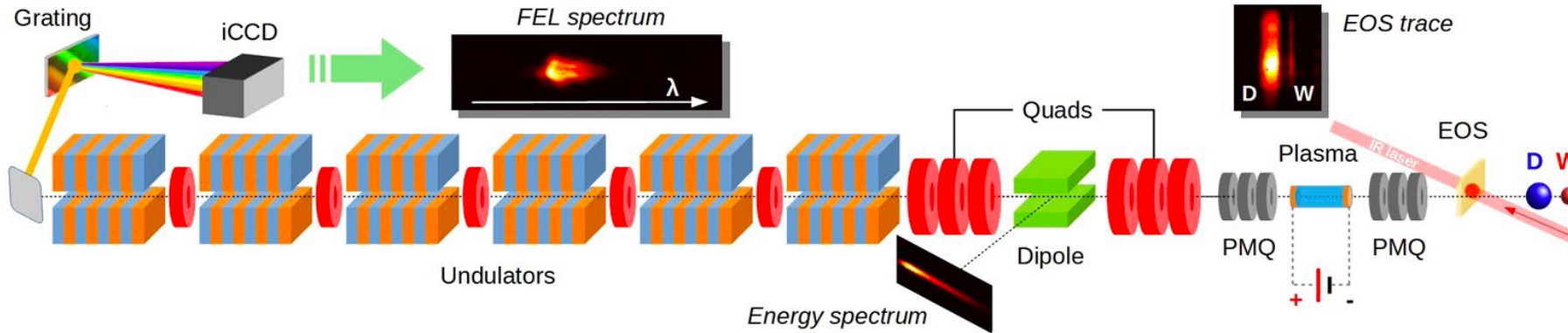
## Plasma acceleration results @SPARC\_LAB

- Two-bunches configuration produced directly at the cathode with laser-comb technique
- 200 pC driver (charge increased up to 350 pC) followed by witness bunch (20 pC)
- 4 MeV acceleration in 3 cm plasma with 200 pC driver
- Ultra-short durations (200 fs + 30 fs) obtained with velocity-bunching technique
- 133 MV/m accelerating gradient
- $2 \times 10^{15} \text{ cm}^{-3}$  plasma density
- Demonstration of projected energy spread compensation: spread from 0.2% to 0.12%



## Energy spread minimization in a beam-driven plasma wakefield accelerator

R. Pompili<sup>1</sup>✉, D. Alesini<sup>1</sup>, M. P. Anania<sup>1</sup>, M. Behtouei<sup>1</sup>, M. Bellaveglia<sup>1</sup>, A. Biagioni<sup>1</sup>, F. G. Bisesto<sup>1</sup>, M. Cesarini<sup>1,2</sup>, E. Chiadroni<sup>1</sup>, A. Cianchi<sup>3</sup>, G. Costa<sup>1</sup>, M. Croia<sup>1</sup>, A. Del Dotto<sup>1</sup>, D. Di Giovenale<sup>1</sup>, M. Diomedea<sup>1</sup>, F. Dipace<sup>1</sup>, M. Ferrario<sup>1</sup>, A. Giribono<sup>1</sup>, V. Lollo<sup>1</sup>, L. Magnisi<sup>1</sup>, M. Marongiu<sup>1</sup>, A. Mostacci<sup>1,2</sup>, L. Piersanti<sup>1</sup>, G. Di Pirro<sup>1</sup>, S. Romeo<sup>1</sup>, A. R. Rossi<sup>4</sup>, J. Scifo<sup>1</sup>, V. Shpakov<sup>1</sup>, C. Vaccarezza<sup>1</sup>, F. Villa<sup>1</sup> and A. Zigler<sup>1,5</sup>



## Proof of SASE and Seeded FEL driven by PWFA @SPARC\_LAB

- Proof-of-principle experiment to demonstrate high-quality PWFA acceleration able to drive a Free-Electron Laser
- Witness is completely characterized (energy, spread, X/Y emittance) allowing to match it into the undulators beamline
- Jitter is online monitored with Electro-Optical Sampling (EOS) diagnostics
- Imaging spectrometer with iCCD used to detect FEL radiation
- Spectrum of the SASE FEL radiation emitted at 830 nm

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**Free-electron lasing with compact beam-driven plasma wakefield accelerator**

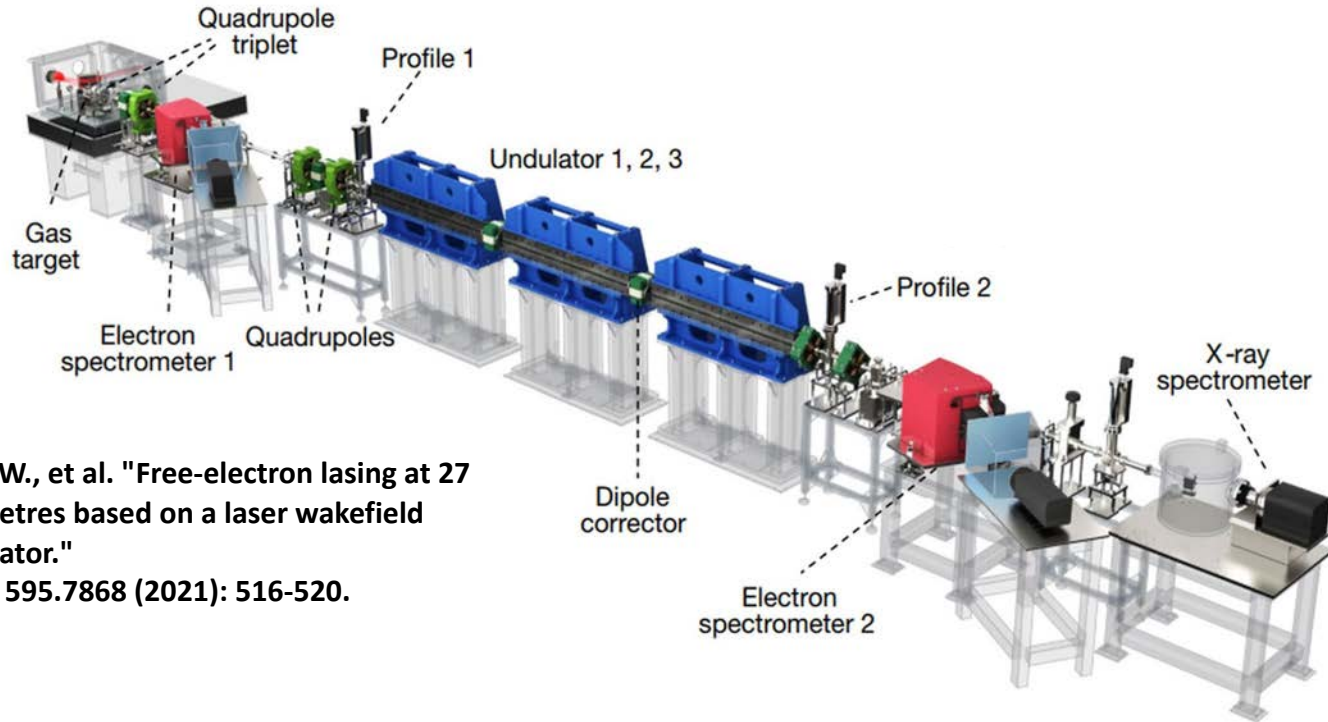
R. Pompili, D. Alesini, ... M. Ferrario + Show authors

Nature 605, 659–662 (2022) | Cite this article

In collaboration with

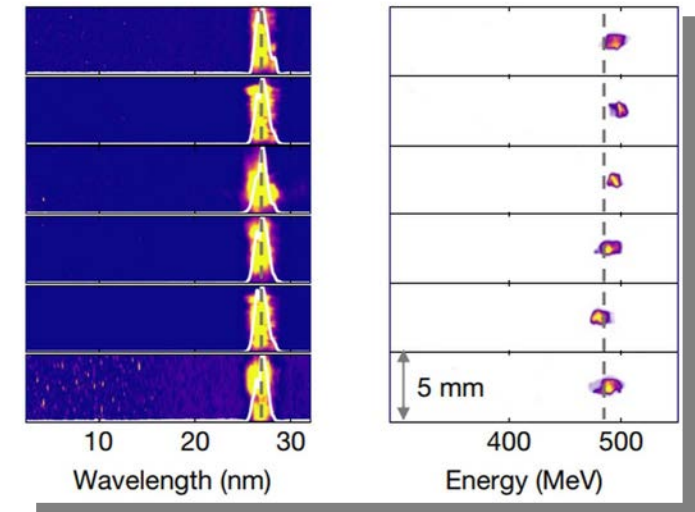
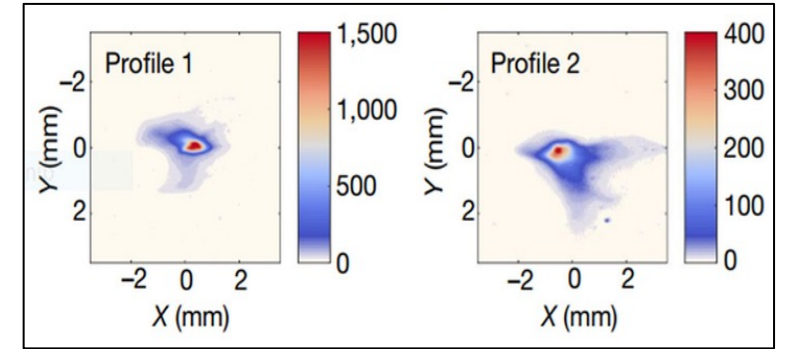






Wang, W., et al. "Free-electron lasing at 27 nanometres based on a laser wakefield accelerator." *Nature* 595.7868 (2021): 516-520.

- laser focused on a gas-jet
- Electron beam generated from a 200 TW ( $I \sim 4 \times 10^{18} \text{ W/cm}^2$ )
- ✓ Peak energy  $\sim 490 \text{ MeV}$ , 0.5% spread (measured), emittance  $0.5 \text{ \mu m}$  (estimated)
- ✓ Radiation energy from 0.5 to 150 nJ



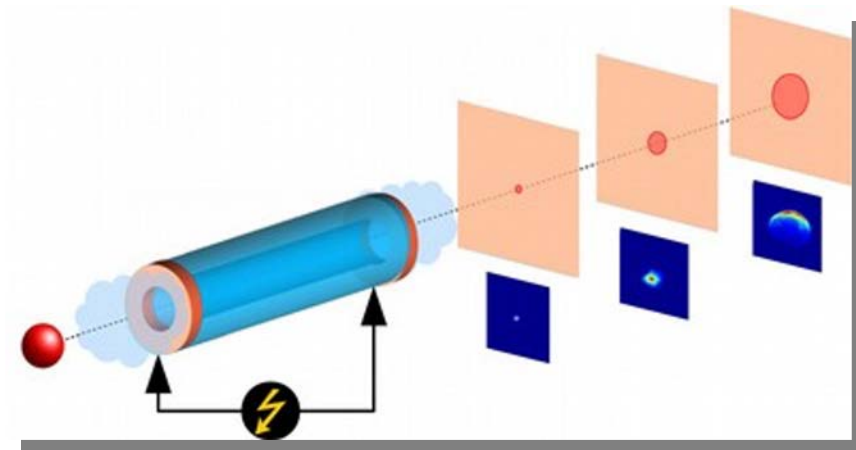
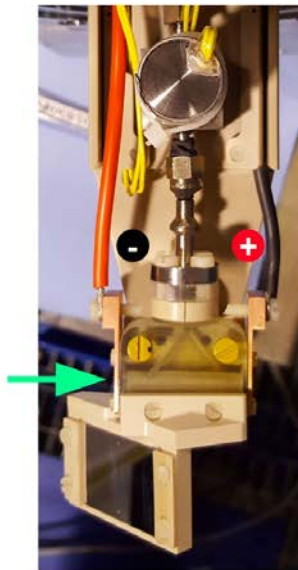
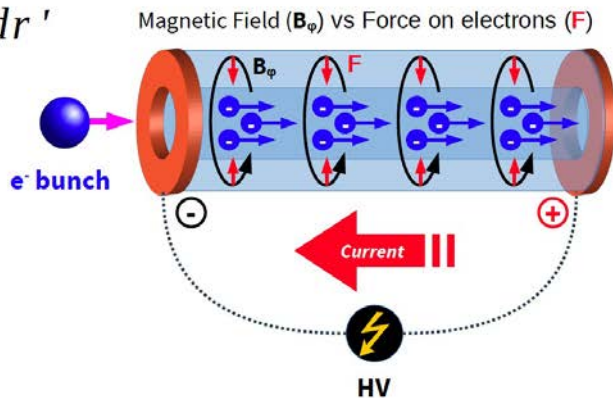
## Focusing with active-plasma lenses

Pompili, R., et al., Physical review letters 121.17 (2018): 174801.  
 Pompili, R., et al., Applied Physics Letters 110.10 (2017): 104101.

Focusing produced by electric discharge in plasma-filled capillary

*Magnetic field follows Ampere Law*

$$B_{\phi}(r) = \frac{\mu_0}{r} \int_0^r J(r') r' dr'$$



Weak chromaticity

*Like in quadrupoles  $\rightarrow K \sim 1/\gamma$*

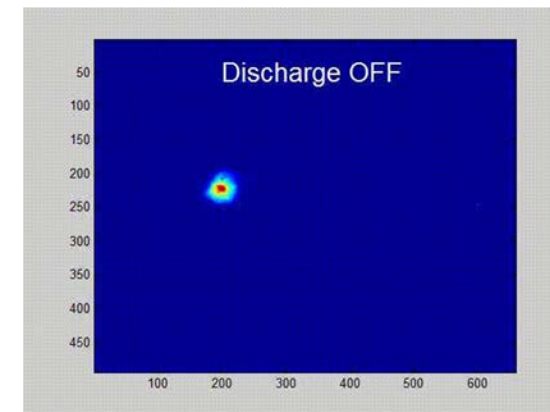
Radially symmetric

*Like in solenoids*

Compactness

*kT/m magnetic field  $\rightarrow$  much larger than strongest quadrupoles available (PMQ)*

Not sensitive to beam distribution



Courtesy R Pompili

EUROPEAN  
PLASMA RESEARCH  
ACCELERATOR WITH  
EXCELLENCE IN  
APPLICATIONS



- Frascati future facility
- Beam-driven plasma accelerator
- Europe most compact and most southern FEL
- The world most compact RF accelerator (X band with CERN)

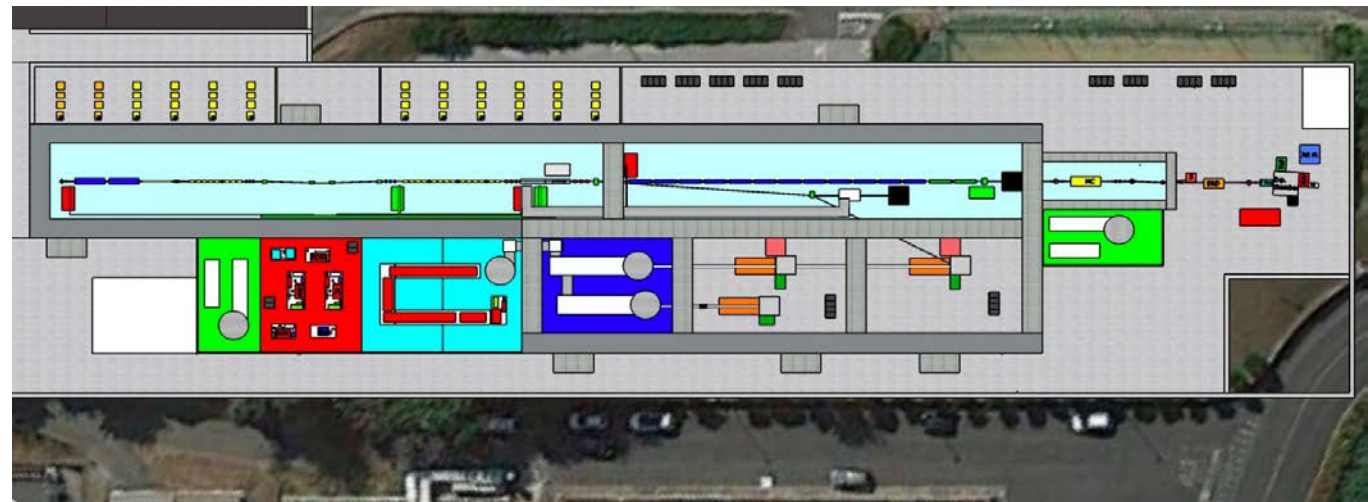
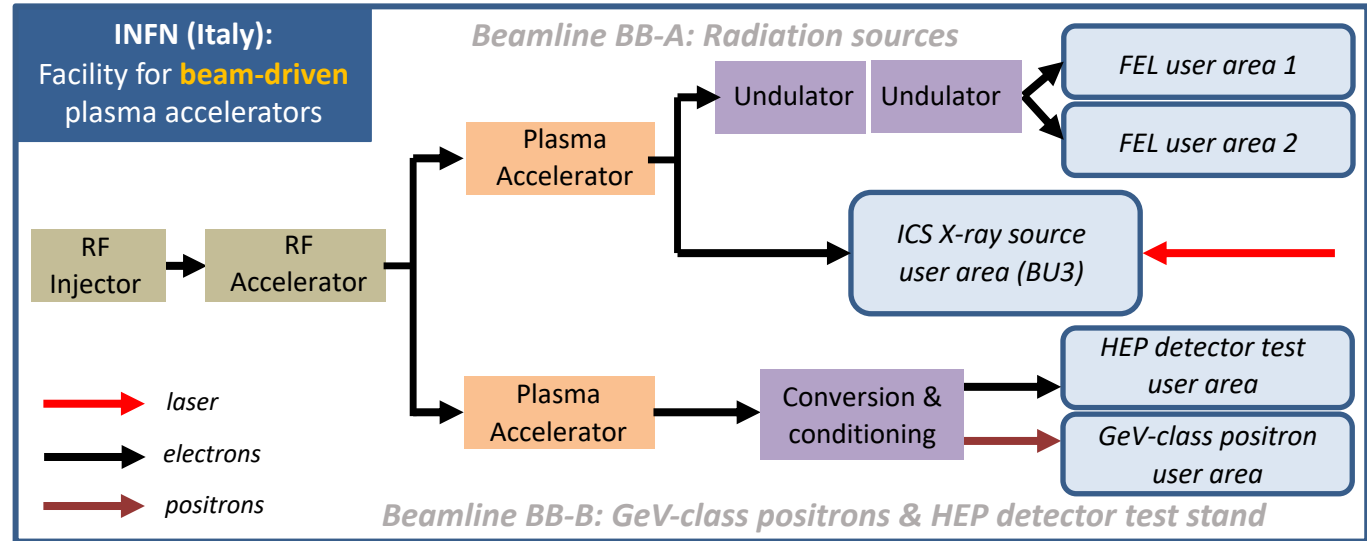
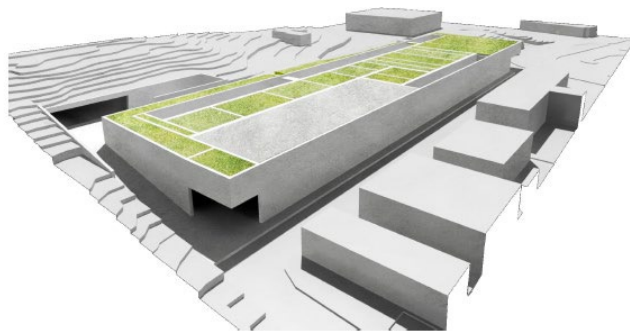
- Electron energy 1-5 GeV
- FEL user facility 1 GeV – 3 nm
- 500 MeV by RF Linac + 500 MeV by Plasma (LWFA or PWFA)
- 1 GeV by X-band RF Linac only



CDR-1, April 2018

## EuPRAXIA@SPARC\_LAB

Conceptual Design Report



## The PNRR EuAPS Project WP2

EuAPS will be the first brick of EuPRAXIA, a user facility based on the radiation emitted by electrons plasma accelerated

- The source will be hosted at LNF-INFN
- Several parts will be realized at CNR (Photon Diagnostics) and at Tor Vergata (User end Station)
- INFN-Mi will take care of simulation and data analysis
- Trieste University focuses on applications

Where	Target
INFN-Mi	Simulation & Data Analysis
LNF-INFN	Plasma source
LNF-INFN	Synchronization
CNR-Potenza	Photon Diagnostics
Tor Vergata	End user station
CNR- Montelibretti	Photon time diagnostics

Parameter	Value	unit
Electron beam Energy	100 - 500	MeV
Plasma Density	$10^{18} - 10^{19}$	$cm^{-3}$
Photon Critical Energy	1 - 10	keV
Number of Photons/pulse	$10^6 - 10^9$	
Repetition rate	1 - 5	Hz
Beam divergence	3 - 20	mrad

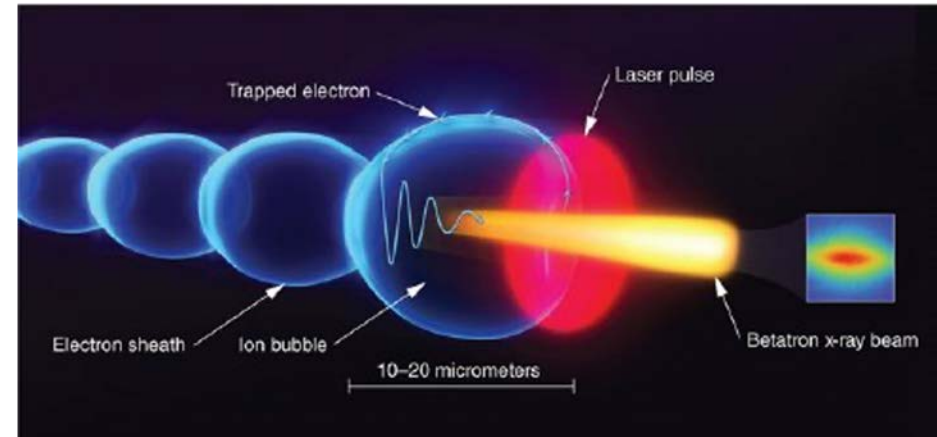
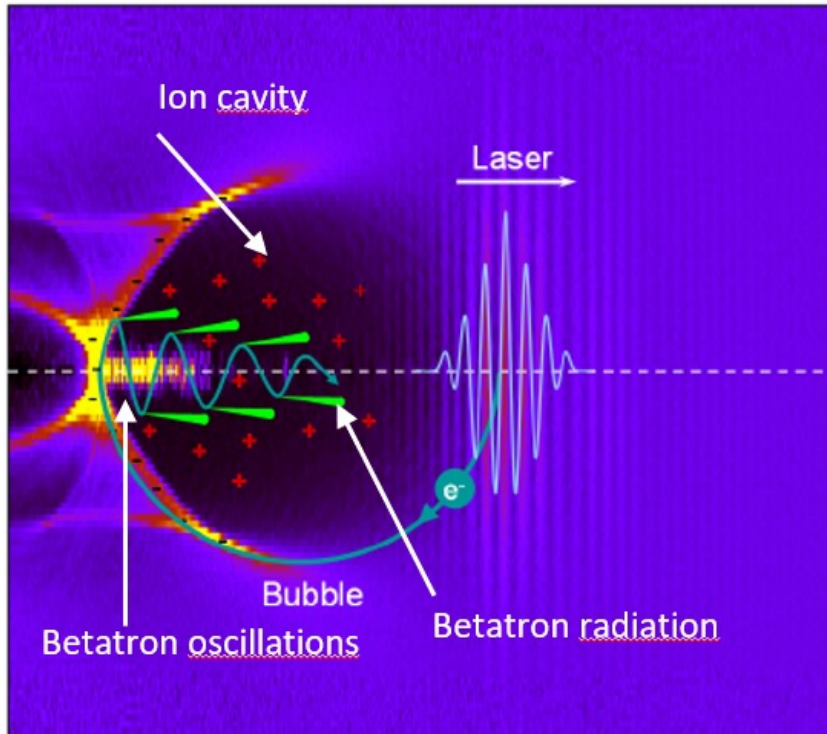
Courtesy A Cianchi

Betatron radiation emission

Longitudinal electric field → acceleration along the laser propagation axis

Radial electric field → oscillations around the reference trajectory

Emission of synchrotron-like radiation = **Betatron radiation**



- Energy from soft to hard X-rays
- High peak brilliance
- Spatially coherent
- Temporally incoherent
- Pulse ~ few fs

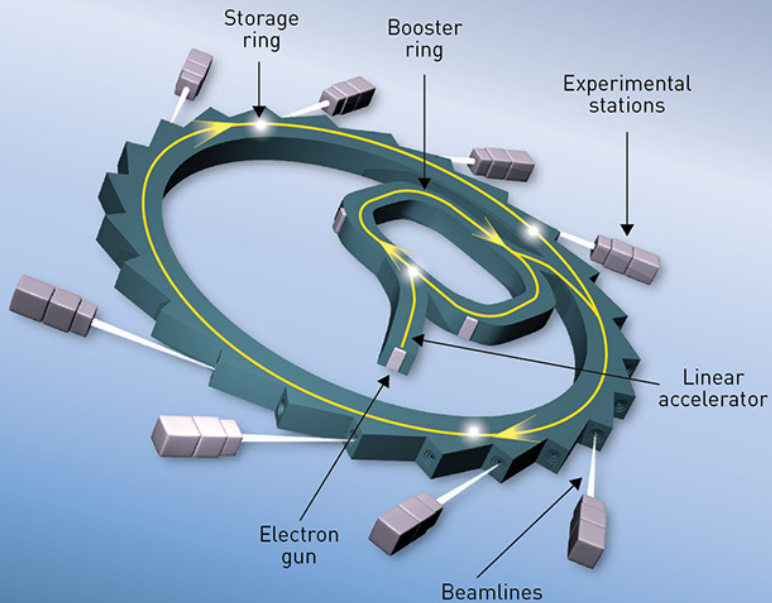
$$B \approx 10^{20} \frac{\text{Photons}}{s \cdot \text{mm}^2 \cdot \text{mrad}^2 \cdot 0.1\% \text{ BW}}$$

## Synchrotrons vs FELs

### Synchrotrons and X-ray FELs

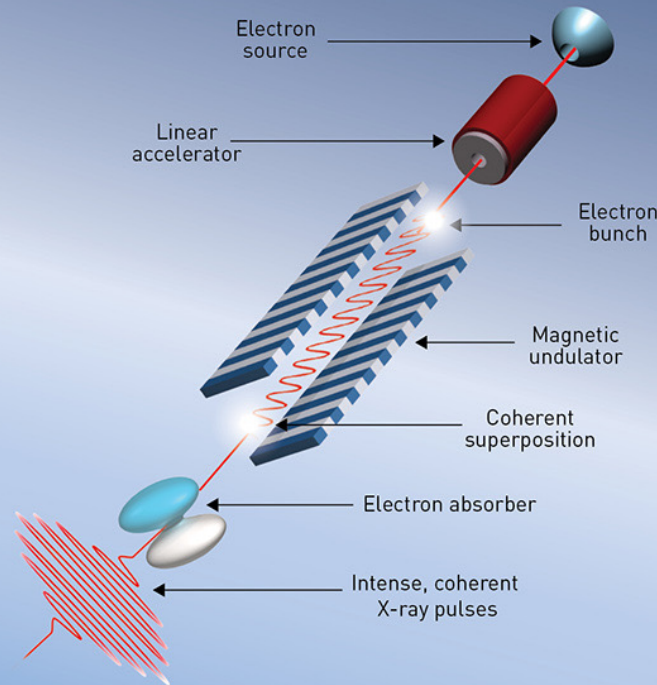
#### Synchrotron light source

Electrons, accelerated to near light speed in a linear accelerator and booster ring, whirl around in a larger storage ring, creating X-rays that feed beamlines for multiple experimental stations



#### X-ray free-electron laser (FEL)

In FELs, accelerated electron bunches are "wiggled" in a magnetic undulator, causing them to throw off coherent, bright and laser-like X-ray beams for experiments



### Betatron radiation:

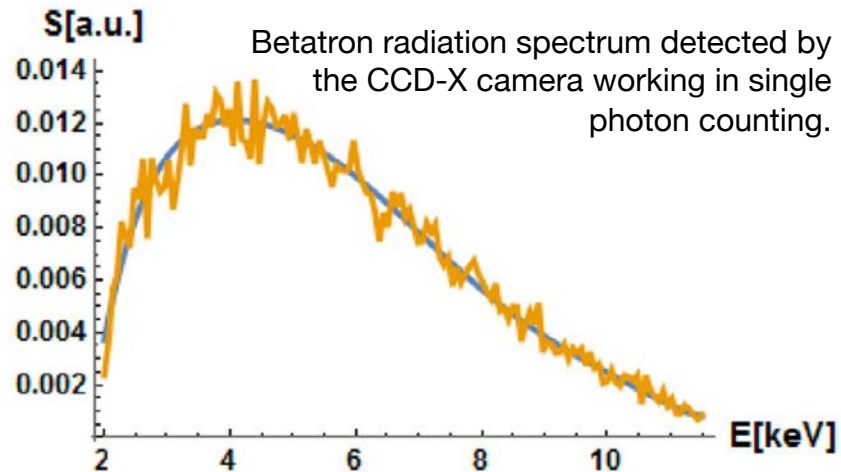
- Large bandwidth like a Synchrotrons
- Short pulse duration like a FEL

$$B = \frac{d^4 N}{dt d\Omega dS d\lambda / \lambda}$$

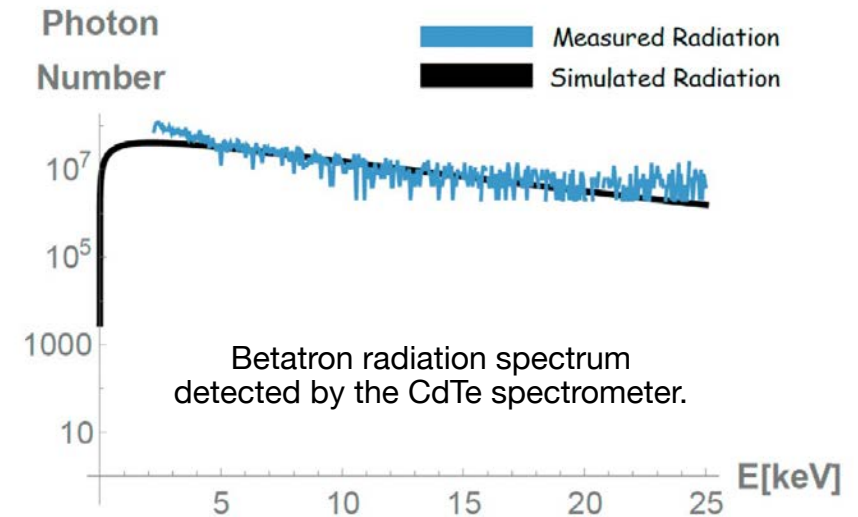
Ph/ (s mm<sup>2</sup> mrad<sup>2</sup> 0.1% of bandwidth)

Patricia Daukantas Synchrotron Light Sources for the 21st Century Optics & Photonics News Settembre 2021

## Betatron radiation test @FLAME



- Laser energy  $1 J$  - temporal length  $30 fs$  - focal spot rms  $5 \mu m$
- Plasma density  $10^{19} cm^{-3}$  - acceleration length  $1.1 mm$
- $e^-$  energy  $300 MeV$  - energy spread  $20\%$  - bunch charge  $5 pC$

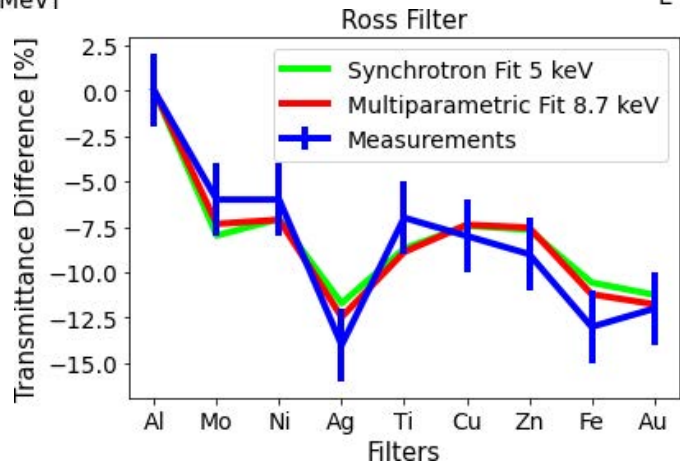
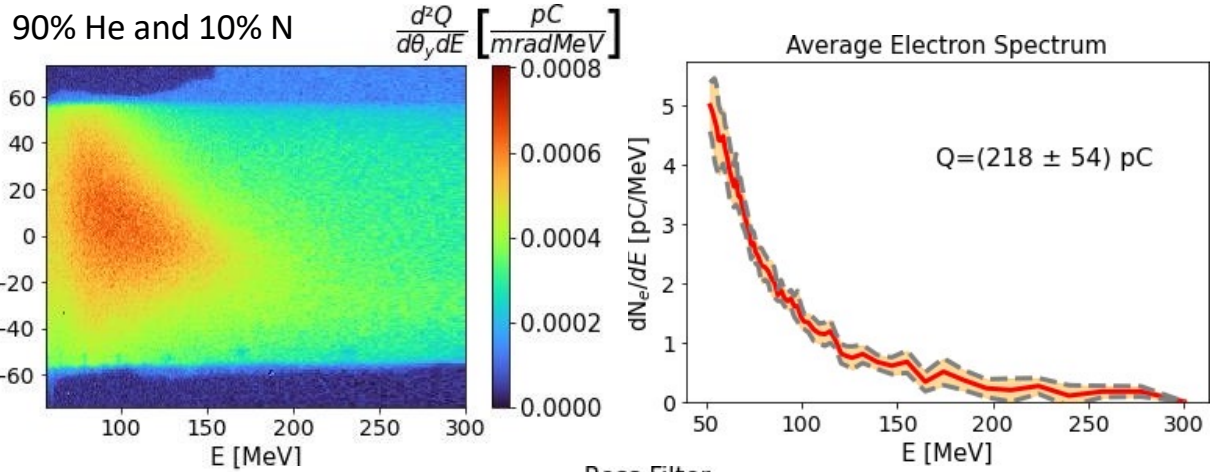


- Laser energy  $1.5 J$  - temporal length  $35 fs$  - focal spot rms  $5 \mu m$
- Plasma density  $6 \times 10^{18} cm^{-3}$  - acceleration length  $1 mm$
- $e^-$  energy  $200 MeV$  - energy spread  $30\%$  - bunch charge  $5 pC$

A. Curcio et al., First measurements of betatron radiation at FLAME laser facility, Nucl. Instr. Meth. B (2017), <http://dx.doi.org/10.1016/j.nimb.2017.03.106>



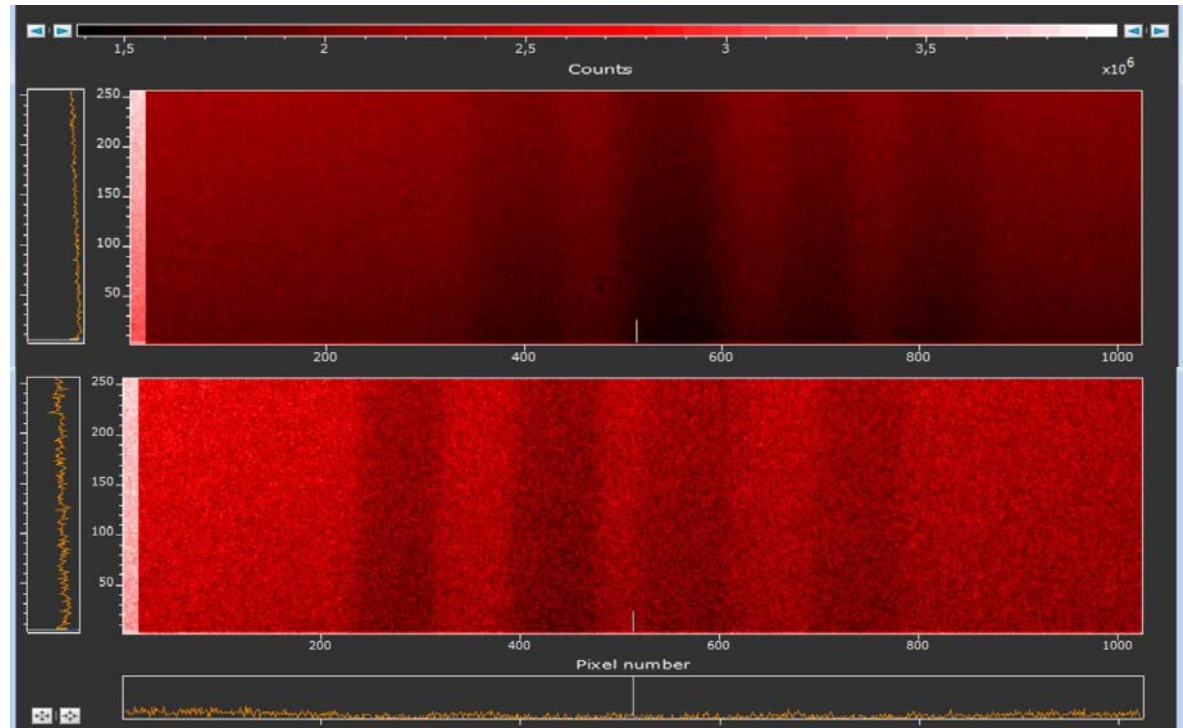
## Ongoing betatron radiation test @FLAME



Ross-filter:

Ti 15  $\mu\text{m}$  – Ag 33  $\mu\text{m}$  – Ni 7  $\mu\text{m}$  – Mo 4  $\mu\text{m}$

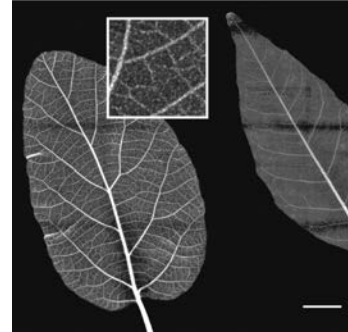
Au 6  $\mu\text{m}$  – Fe 25  $\mu\text{m}$  – Zn 10  $\mu\text{m}$  – Cu 8  $\mu\text{m}$



## Photon Science from betatron radiation

- Imaging of biological and cultural heritage samples:  
Exploits the brilliance and coherence of betatron radiation, requires small divergence and good focusing
- Static X-ray Spectroscopy and Ultra-fast X-ray spectroscopies:  
The second one requires timing between pump and probe pulses, exploits the fs pulse duration
- Wide angle scattering, diffraction:  
Requires monochromatic beams with high flux

### Imaging – The pilot experiment

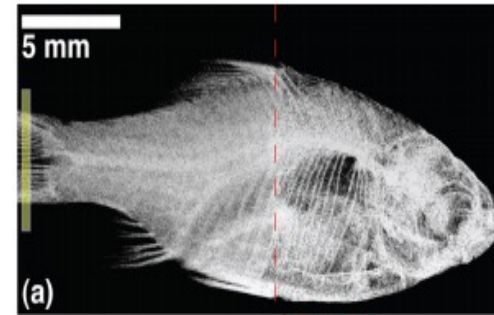


X-ray imaging of leaves (and wood) aiming at the (tens of) microns resolution.

Experiments performed with the broad radiation spectrum filtered by different materials to obtain difference maps emphasizing the presence of heavy metal contaminants → pollution control.

Reale et al. - MIDIX Soft X-rays microradiography

### Single shot phase contrast X-ray imaging:



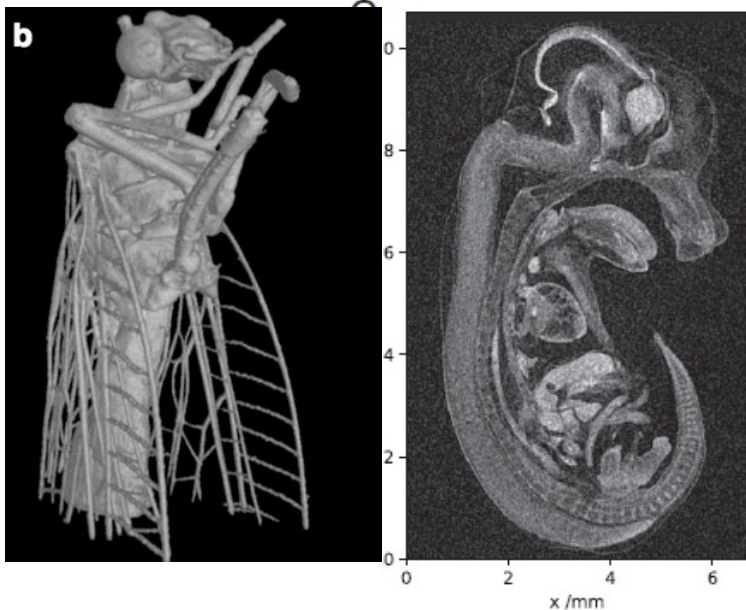
X-ray absorption contrast image of an orange tetra fish. The spectrum is synchrotron like with a critical energy  $E_c \sim 10$  keV. The phase contrast images are taken in a single shot 30 fs exposure.

Ref. Kneip, S., et al. "X-ray phase contrast imaging of biological specimens with femtosecond pulses of betatron radiation from a compact laser plasma wakefield accelerator." **APL** 99.9 (2011): 093701.

Courtesy F Stellato

## Imaging, Tomography and Phase Contrast

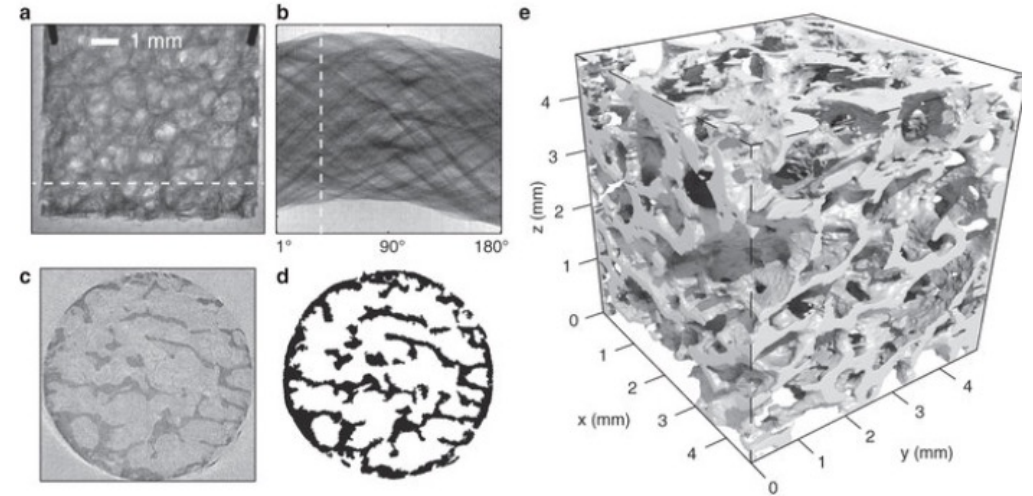
Betatron sources can fill the gap between synchrotrons and X-ray tubes for imaging and **Computer Tomography (CT)**



Betatron sources have a spatial coherence that allows performing Phase Contrast Imaging (PCI). In PCI, it is measured the difference in wavefront, while in traditional imaging, it is measured the difference in the X-ray absorption coefficient between different objects.

PCI provides better contrast than radiography, especially when dealing with biological samples.

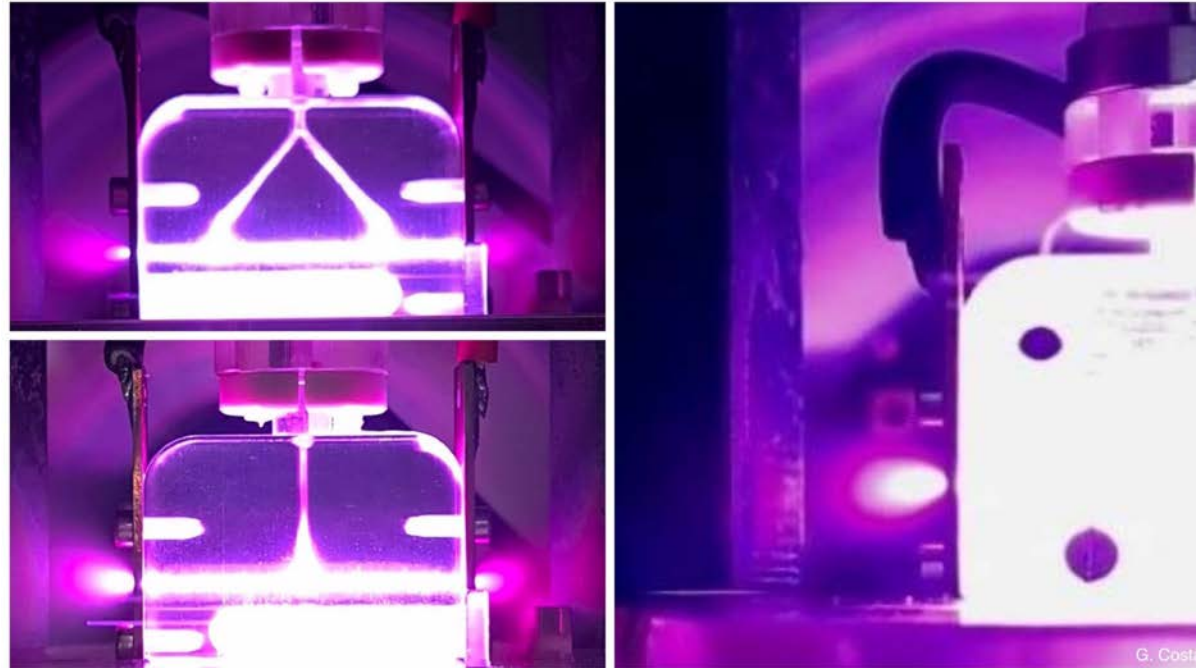
Guo *et al.* Scientific Reports 2019  
 Cole *et al.* PNAS 2018  
 Wenz *et al.* Nature communications 2015



**3D medical imaging:** Tomographic reconstruction of bone sample: (a) A raw image of the bone sample recorded on the xray camera. (c) Application of the inverse Radon transform to the sinogram in (b). (d) Pixels are classified as bone (black) or vacuum (white). (e) Stacking together 1300 slices generates a 3 D voxel map of the bone sample. An isosurface marking the detailed structure of the bone surface is constructed, rendered using a ray-tracing method.

Ref. Cole, J. M., et al. "Laser-wakefield accelerators as hard x-ray sources for 3D medical imaging of human bone." *Scientific reports* 5.1 (2015): 1-7.

Thank you for the attention!



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