

Tidal deformability of black holes surrounded by thin accretion disks

Enrico Cannizzaro

Istituto Superior Tecnico, Lisbon

In collaboration with Valerio De Luca, Paolo Pani

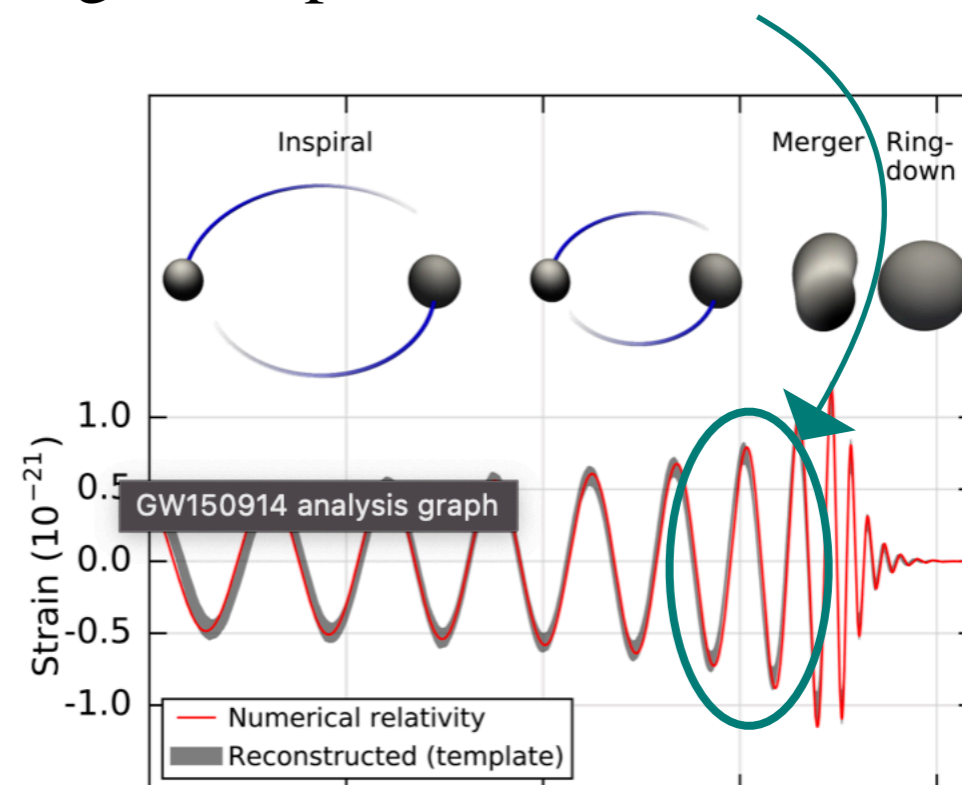
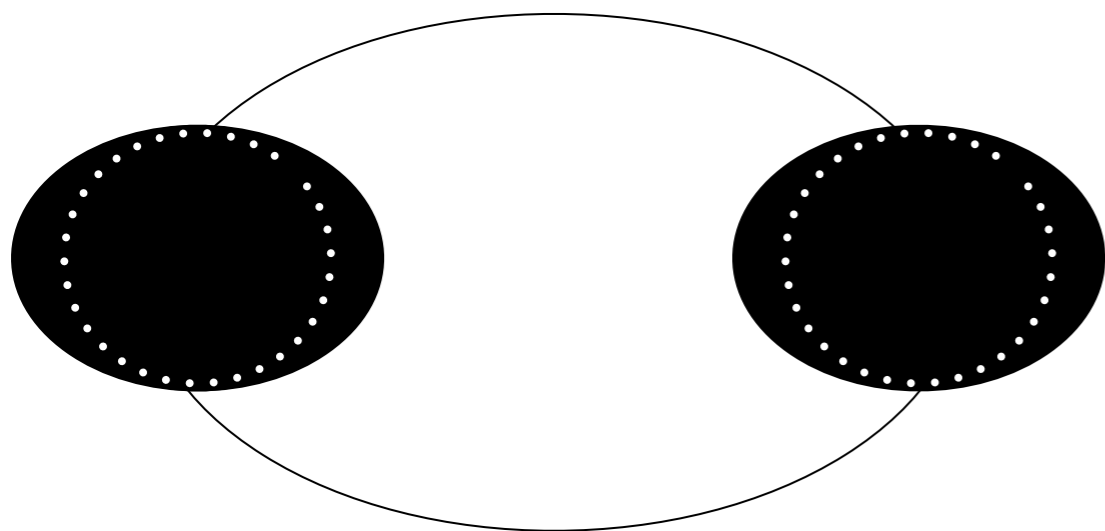


grit
gravitation in técnico



Tidal Deformability and Love Numbers

- Self-gravitating objects may be deformed in external tidal fields (e.g. companions in a binary system)
- Deformation is quantified using Tidal Love Numbers (TLNs), which describe the linear response of the body, akin to a “gravitational susceptibility”
- TLNs depend on the theory of gravity and the internal structure of the object, and they affect the orbital dynamics by leaving a footprint at 5PN order



Neutron Stars

Neutron stars can be tidally deformed, with a TLN dependent on their equation of state

$$\tilde{\Lambda} \propto O(EOS) \left(\frac{R_{NS}}{M_{NS}} \right)^5$$

Hinderer (2008)

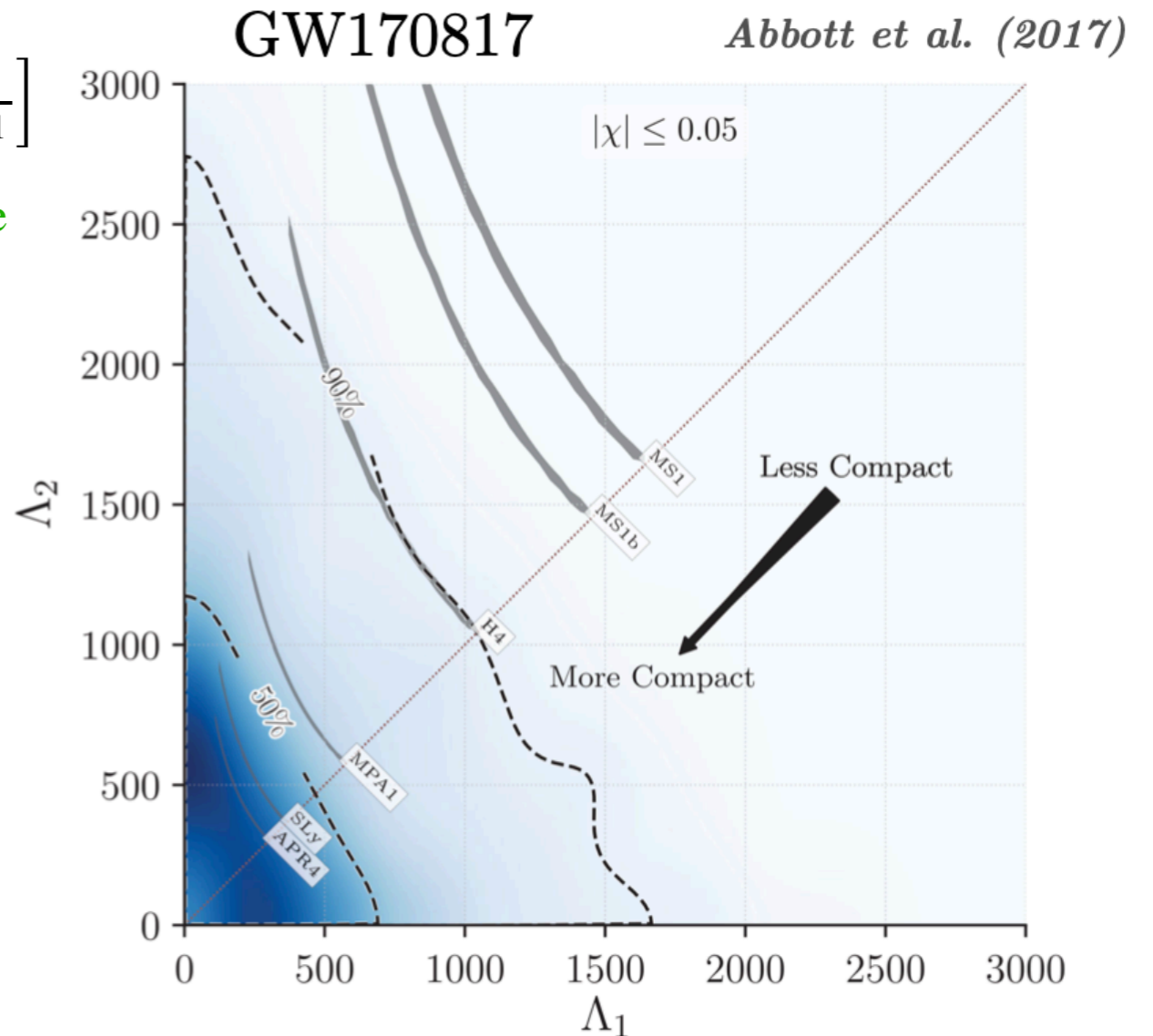
Gravitational potential of deformed body (Newtonian):

$$U_{tot} = -\frac{GM}{r} - \sum_{\ell=2}^{\infty} \sum_{m=-\ell}^{\ell} Y_{\ell m} \left[\underbrace{\frac{(\ell-2)!}{\ell!} \mathcal{E}_{\ell m} r^{\ell}}_{\text{Tidal field}} - \frac{(2\ell-1)!!}{\ell!} \frac{I_{\ell m}}{r^{\ell+1}} \right]_{\text{Body response}}$$

Linear response theory:

$$I_{\ell m}(\omega) = -\frac{(\ell-2)!}{(2\ell-1)!} \lambda_{\ell m}(\omega) r_h^{2\ell+1} \mathcal{E}_{\ell m}(\omega)$$

$$k_{\ell m} = \text{Re}(\lambda_{\ell m})$$



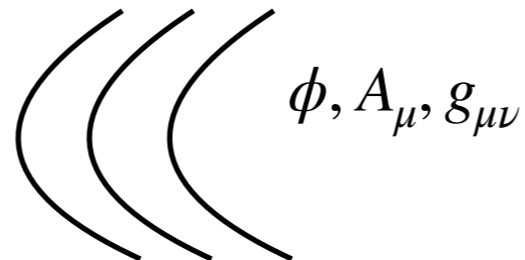
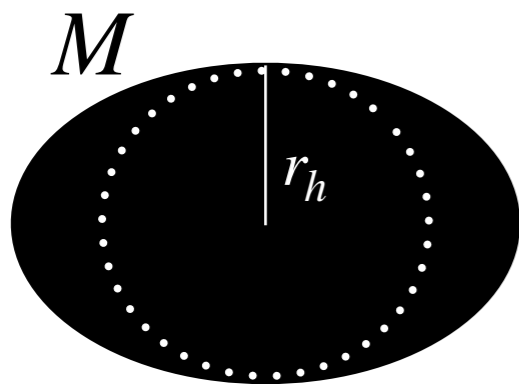
Neutron Stars

Neutron stars can be tidally deformed, with a TLN dependent on their equation of state

$$\tilde{\Lambda} \propto O(EOS) \left(\frac{R_{NS}}{M_{NS}} \right)^5$$

Hinderer (2008)

GR case: perturbation theory - massless fields in curved spacetime



Equation of motion: $\mathcal{O}_s \psi_s = 0$

Matching at the star surface: ψ_s regular at $r = r_h$

$$\psi_s \propto r^{\ell+1} \left[1 + k_s^{(\ell)} \left(\frac{r}{r_h} \right)^{-2\ell+1} \right]$$

Tidal deformability of black holes

In $D=4$, the static TLNs of asymptotically flat BHs in General Relativity vanish.

Fragile condition. $k_2 \neq 0$ in a plethora of different scenarios:



Beyond GR theories

Cardoso+ (2017,2018), De Luca+(2023), Barura+(2024)



**BH mimickers and
exotic compact objects**

*Pani+(2015), Cardoso+ (2017, 2019), Herdeiro+(2020),
Berti+(2024)*



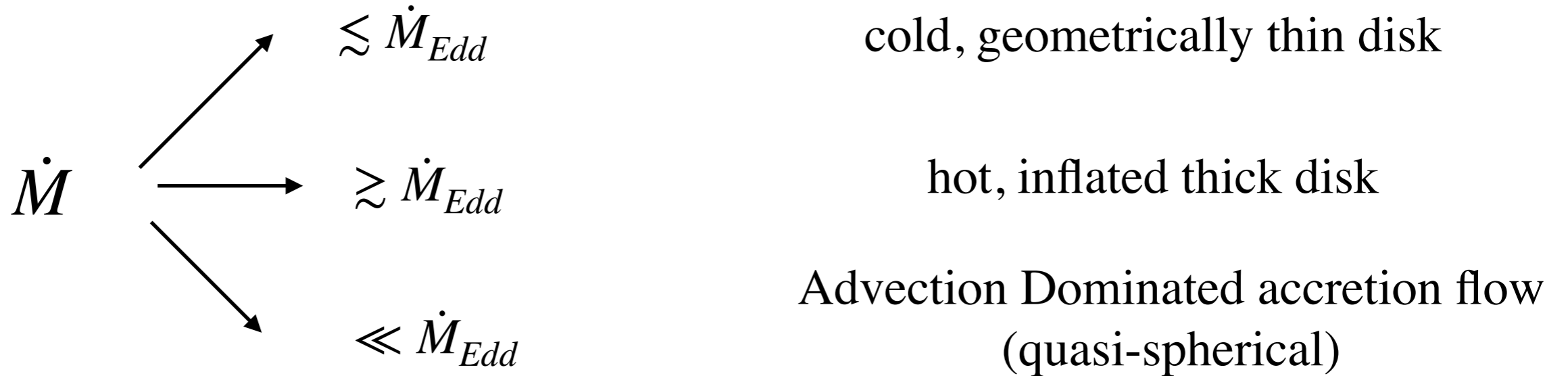
Higher dimensions

*Hui+(2021), Kol+(2012), Cardoso+(2019), Rodriguez+(2023),
Charalambous+(2023,2024), Ma+(2024)*

Implications in fundamental physics: Tidal tests of general relativity!

Why so vacuum?

Accretion disks are ubiquitous in astrophysics: Matter (dust, plasma) spirals around the BH due to gravity



- Can feature very high densities
($n_e \approx 10^{19} M_\odot / M \text{ cm}^{-3}$)
- Less compact than neutron stars (more deformable?)

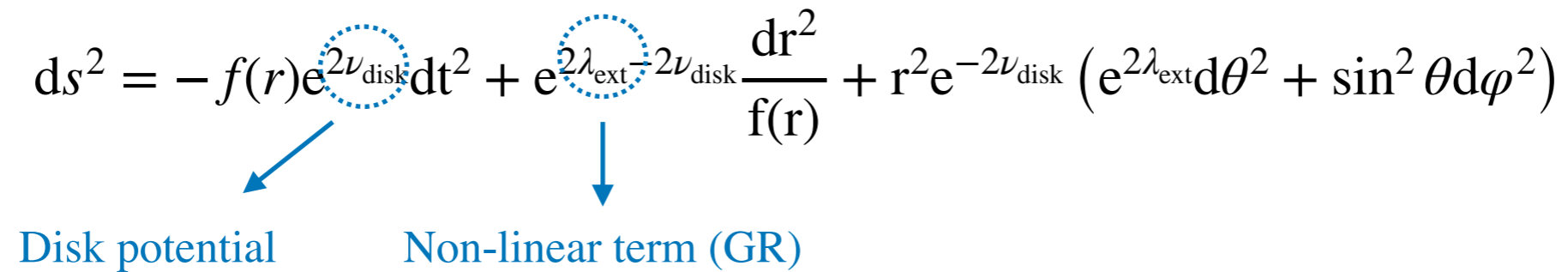


What is the deformability of a BH+disk system?

BH+thin disk geometry

“Non-linear superposition” of black hole and thin disk:

$$ds^2 = -f(r)e^{2\nu_{\text{disk}}}dt^2 + e^{2\lambda_{\text{ext}}-2\nu_{\text{disk}}}\frac{dr^2}{f(r)} + r^2e^{-2\nu_{\text{disk}}}(e^{2\lambda_{\text{ext}}}d\theta^2 + \sin^2\theta d\varphi^2)$$



Disk potential Non-linear term (GR)

- Fully relativistic axisymmetric and static metric
- The disk stretches from the horizon to infinity, and vanishes in the extremities
- Analytic expressions for the disk functions
- Density profile is astrophysically realistic (equatorial plane)

Model parameters of the disk:

$$\left\{ \begin{array}{ll} \epsilon = M_{\text{disk}}/M_{BH} & \text{mass of the disk compared to BH mass} \\ \tilde{b} & \text{location of the disk peak} \end{array} \right.$$

Tidal deformability of BHs with thin accretion disks

At leading order in ϵ :

Spin-0

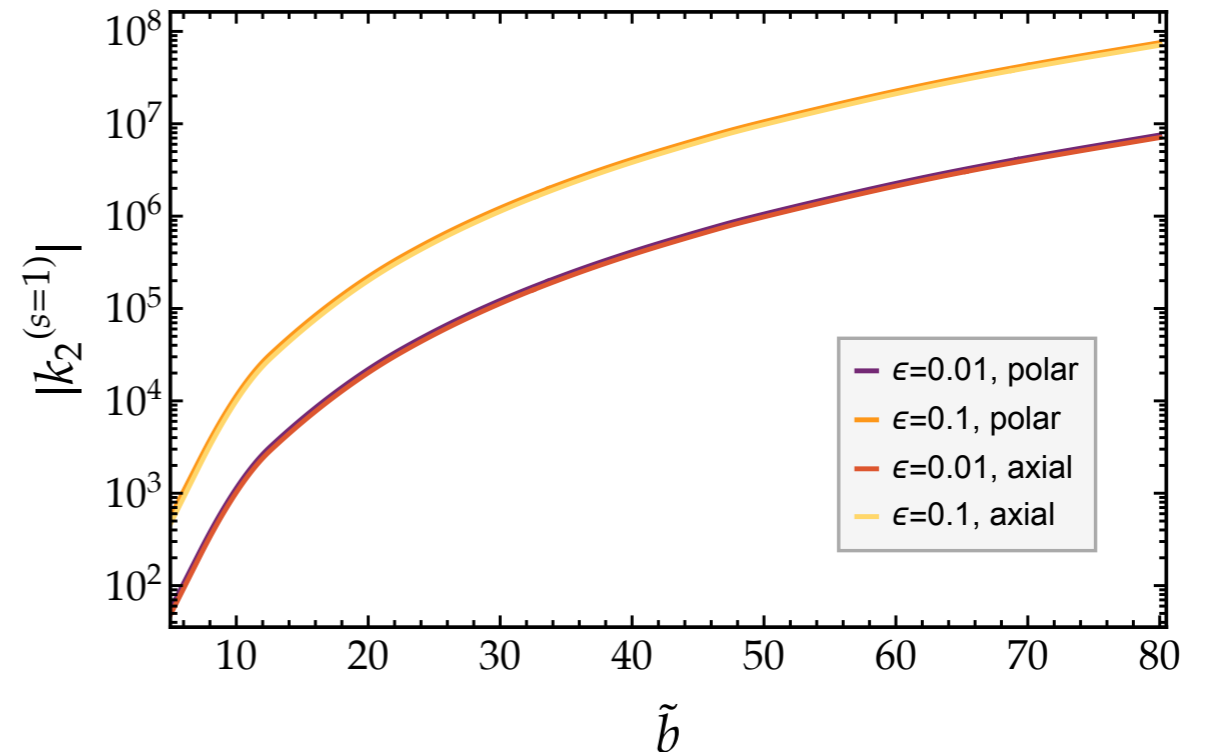
$$k_2 = \frac{\epsilon}{768} (432\tilde{b}^4 - 208\tilde{b}^2 + 9)$$

Typical values:

$$\epsilon \rightarrow [10^{-4}, 10^{-1}]$$

$$b \rightarrow [r_{\text{ISCO}}, 30M]$$

Spin-1



Scaling $k_2 \propto \epsilon \tilde{b}^4$ is in agreement with the scaling of TLNs of dressed BHs in other environments (boson clouds, thin shells, Einstein clusters) *De Luca+(2021), Duque+(2019,2021)*

Environmental Effects vs Modified GR

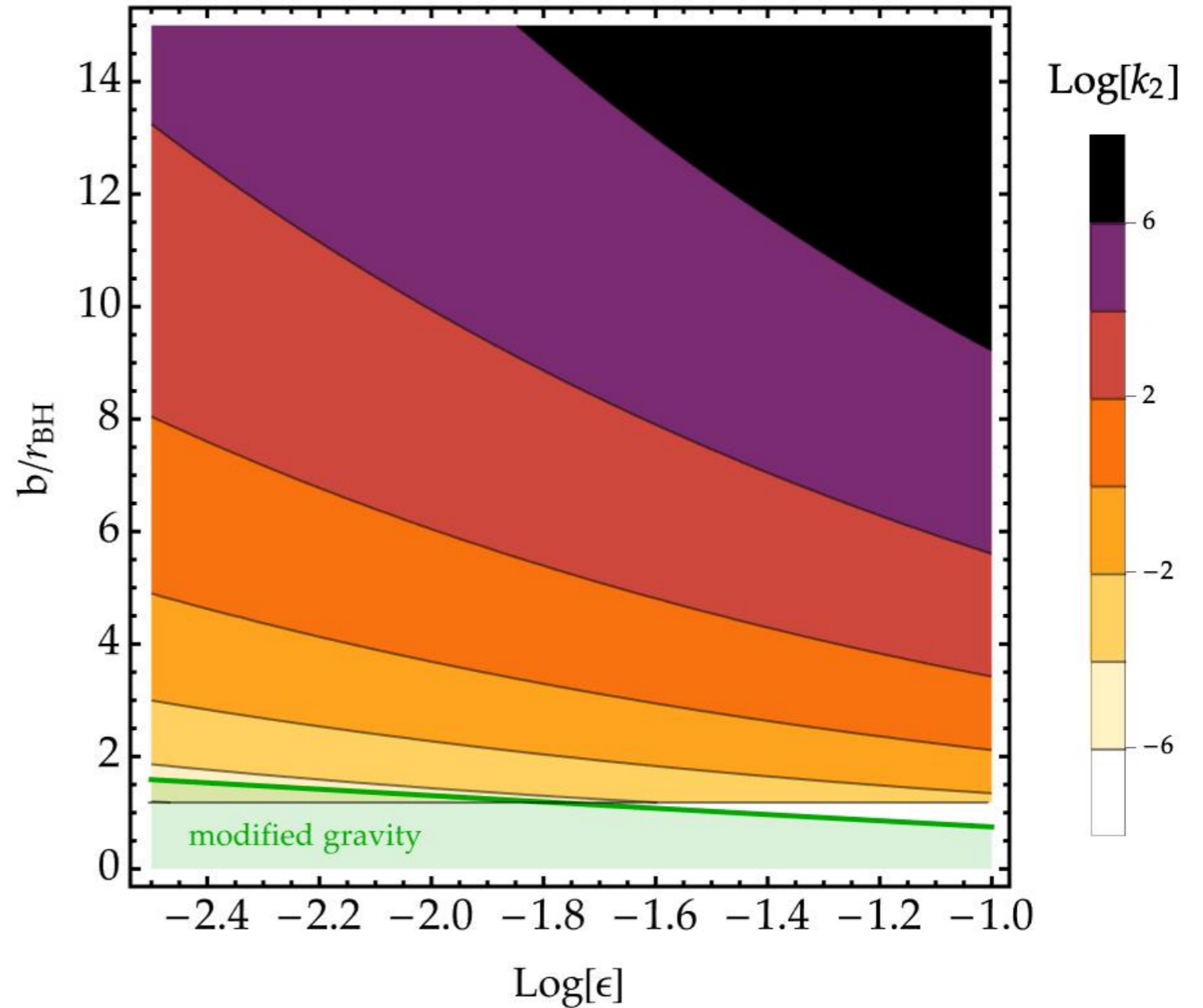
Lots of beyond GR theories feature
Non-vanishing Love Numbers



$$k_2^{\text{MG}} \propto \frac{\alpha_{\text{MG}}}{G\Lambda_{\text{MG}}^2} \left(\frac{1}{Gm_{\text{BH}}} \right)$$

Can the TLN of the BH+disk system
jeopardize tidal tests of GR?

Yes, whenever $\tilde{b} \gtrsim 5$, $\epsilon \gtrsim 10^{-2.5}$



Tidal disruption and frequency dependent TLNs

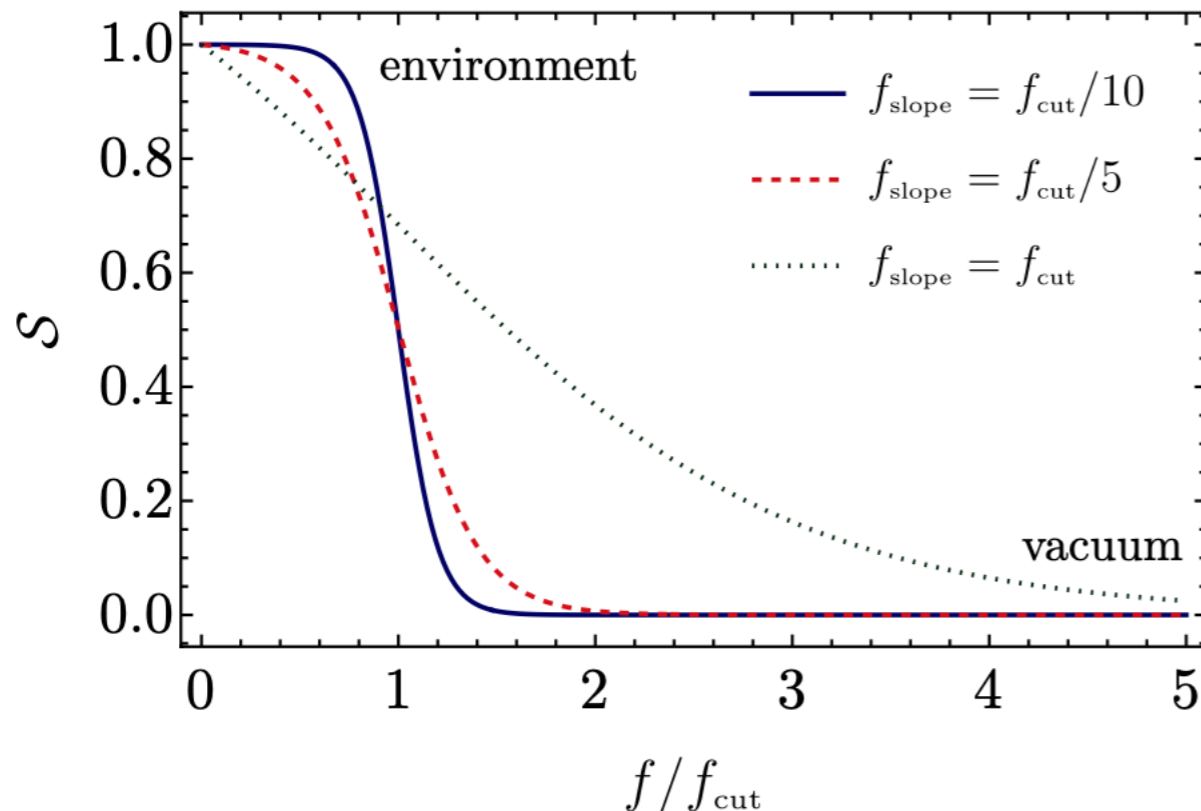
The accretion disks will get tidally disrupted at the Roche radius:

$$r_{Roche} = 2\gamma m_1(1 + \tilde{b}) \left[\frac{m_2}{m_1(1 + \epsilon)} \right]^{1/3} \quad \text{at a GW frequency} \quad f_{cut} = \frac{1}{2\sqrt{2}\pi\gamma^{3/2}m_1} \sqrt{\frac{(1 + \epsilon)(m_1 + m_2)}{(1 + \tilde{b})^3 m_2}}$$



Large \tilde{b} : Large TLN but small Roche frequency

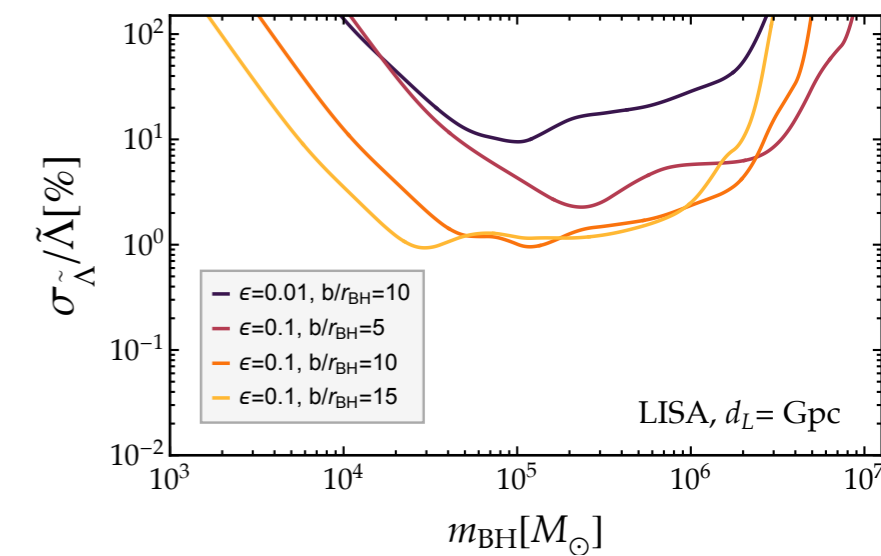
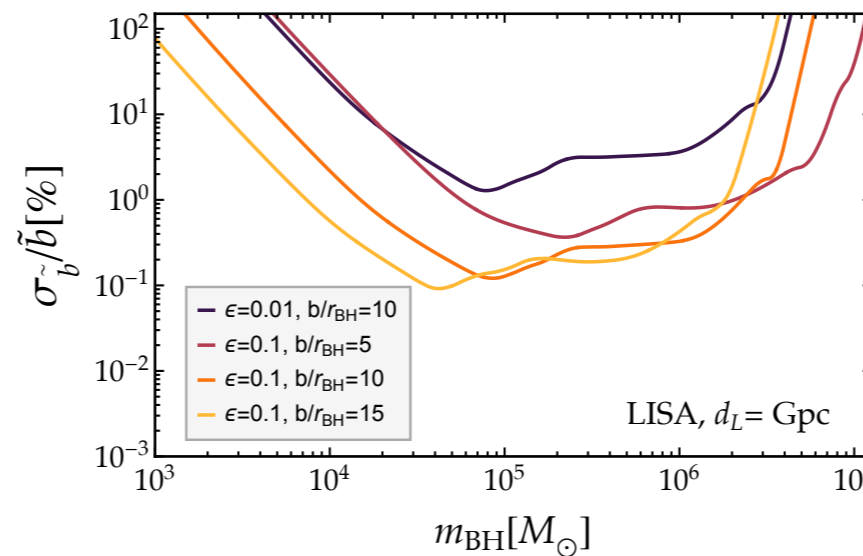
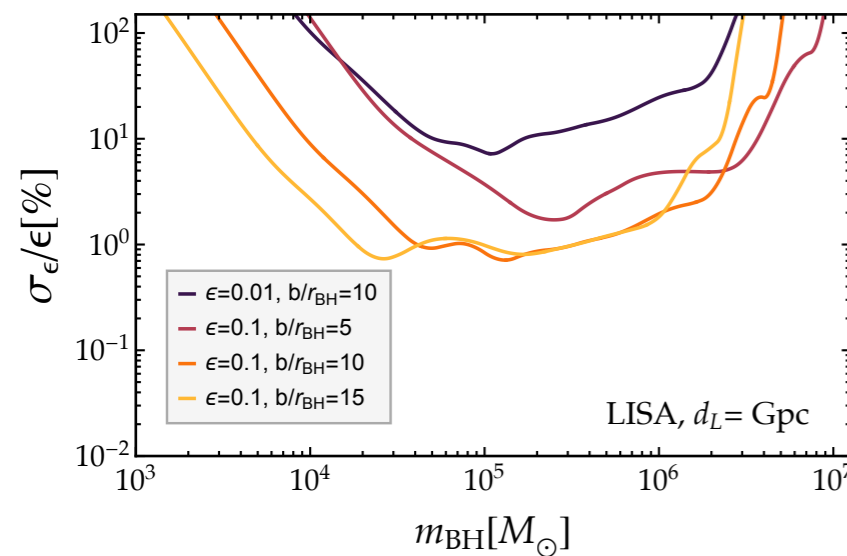
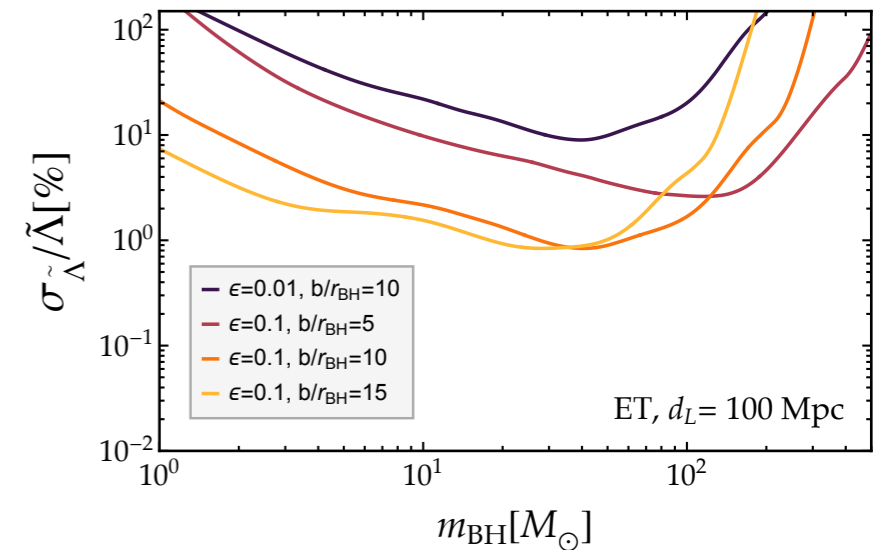
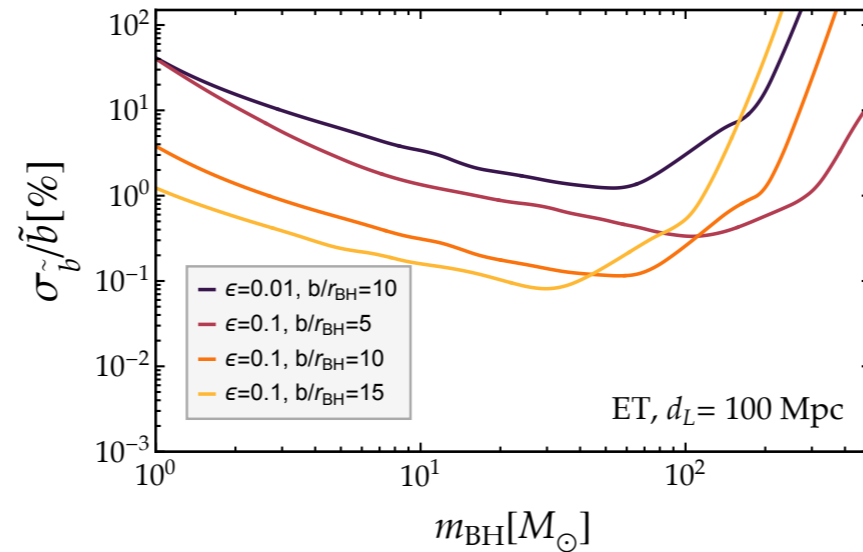
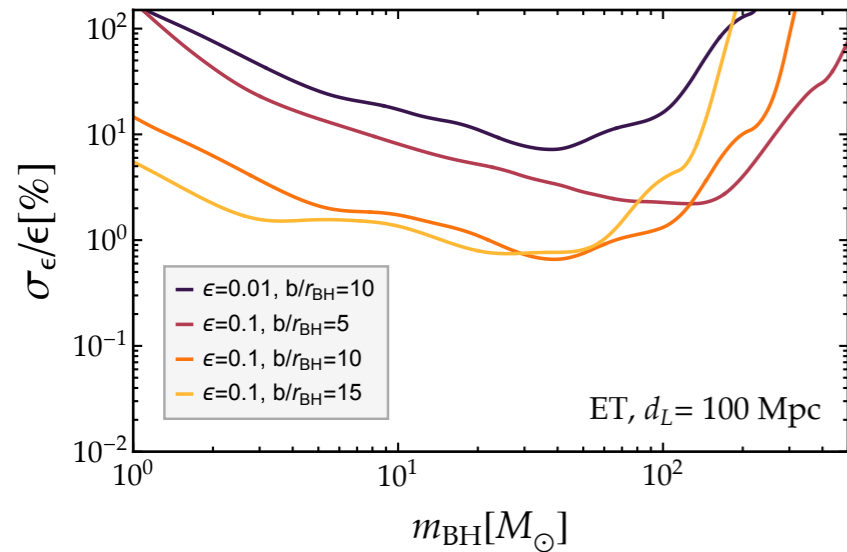
Proxy for frequency dependent tidal deformability



$$\tilde{\Lambda} \rightarrow \mathcal{S}(f) \cdot \tilde{\Lambda} = \frac{1 + e^{-f_{cut}/f_{slope}}}{1 + e^{(f-f_{cut})/f_{slope}}} \cdot \tilde{\Lambda}$$

Detectability at ET and LISA: Fisher matrix analysis

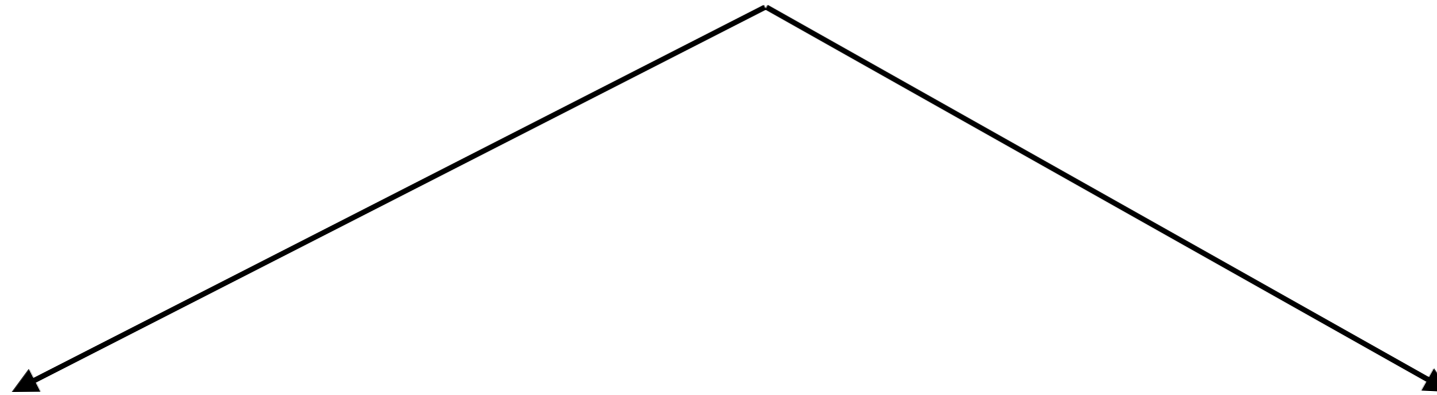
$$\tilde{h}(f) = C_{\Omega} \mathcal{A}_{\text{PNE}} e^{i\psi_{\text{PP}}(f) + i\psi_{\text{Tidal}}(f)}$$



All tidal parameters can be measured with high accuracy with next generation detectors!

Conclusions

- Presence of a thin disk around a BH induce a non-vanishing tidal deformability



Fundamental

Physics implications:

Such effect can easily jeopardize tidal tests of theories beyond GR and BH mimickers

Astrophysics

implications:

Disk parameters could be measured with high accuracy with LISA and third generation detectors

Thank you!